



## Abstract

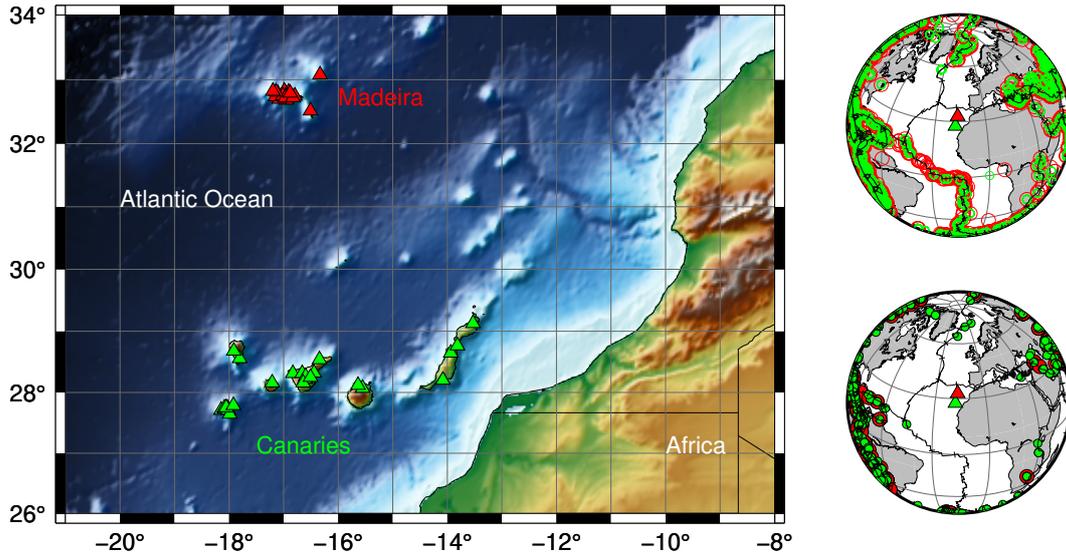
The Canary and Madeira Islands are two distinct hotspots in the Central-East Atlantic that are close to each other. Their volcanism is generally attributed to underlying mantle plumes, but the detailed structure of these plumes is still not well understood. The thermal and compositional structure of the plume introduces complexities in the phase transitions of the mantle, which impact the depth and magnitude of seismic discontinuities. We use 1268 high-quality receiver functions from stations located at the two hotspots to detect P-to-s converted phases through a common-conversion point stacking approach and conduct a detailed analysis of mantle seismic discontinuities. The results show that both hotspots are characterized by a thin mantle transition zone (MTZ), with sharp 410 and 660 discontinuities at depths of 429-420 km and 647-664 km, beneath the Canaries and Madeira respectively. The results indicate that the Canary plume crosses the MTZ, whereas the Madeira plume mainly influences the upper portion of the MTZ. Furthermore, we find reliable detections of a sharp X discontinuity beneath the Canaries at 289 km. Its presence suggests the accumulation of silica-rich recycled eclogite at these depths. We also use the amplitudes of P410s and PXs to derive velocity jumps at corresponding discontinuities. Based on these measurements, we estimate that the basalt proportion is 60-80%, with accumulation being more significant in the Canaries than in Madeira. The MTZ thickness, the presence of the X discontinuity, and the high basalt proportion provide compelling evidence for a deep-rooted thermochemical plume beneath the study area.

## Plain Language Summary

In this study, we examine the Canary and Madeira Islands, two separate groups of islands in the middle of the Atlantic Ocean, not far from each other. While it is generally accepted that volcanic activity in these regions is driven by underlying mantle plumes, the precise characteristics of these plumes remain relatively unknown. Recent tomography studies have revealed distinct structures of the plumes, with the Canary plume extending vertically throughout the upper mantle, while the Madeira plume is visible up to a depth of around 300 km, beyond which the resolution of the images deteriorates. To enhance our understanding, we use receiver functions, derived from the analysis of seismic waveforms, particularly the conversion of compressional waves to shear waves at seismic discontinuities within the Earth. Receiver functions offer a higher resolution compared to tomographic images, primarily due to their ability to directly measure seismic wave conversions at specific depths, reflecting where mantle minerals undergo changes in their crystal structure. This analysis provides valuable insights into the variations in both temperature and mineral chemistry in the mantle, as both factors significantly influence the precise depths of crystal rearrangements. Our research provides substantial evidence supporting the existence of a deep-seated thermochemical plume beneath the study area. However, our findings suggest that beneath Madeira, the plume may be less vigorous or presently disconnected from its deep source, which is consistent with the information obtained from tomographic images.

## 1 Introduction

The Canary and Madeira archipelagos are part of two intraplate volcanic provinces in the Central-East Atlantic Ocean close to the African coast (Figure 1). Both provinces consist of linear chains of volcanic islands and seamounts that display a general SW-NE progression of increasing volcanism age, aligning with the movement of the African plate. They are both marked by relatively low-magnitude seismic activity, which on the Canaries is mainly associated with volcanic processes (Carracedo & Troll, 2021). Although still under debate, the volcanism of both archipelagos is attributed to the interaction be-



**Figure 1.** Map of the Canary and Madeira archipelagos displaying station locations (left), with red and green triangles representing stations in Madeira and the Canaries, respectively. On the right, the global distribution of all teleseismic earthquakes recorded in the seismic stations (top) and the ones used in the final analysis (bottom) after rigorous quality controls, with red triangles marking the positions of the Madeira archipelago and green for the Canary archipelago. Events follow the same color coding as the triangles, with green indicating events analyzed for the Canaries and red for those related to Madeira.

71 tween the African Plate and underlying mantle plumes (Mata et al., 1998; Geldmacher  
72 et al., 2001, 2005; Civiero et al., 2021; Negredo et al., 2022).

73 The most recent travel-time tomography models image slow anomalies under both  
74 archipelagos, attributed to mantle upwellings (Civiero et al., 2018; Civiero, Custódio, et  
75 al., 2019). In a recent comprehensive analysis employing tomography and paleo-reconstruction  
76 models, Civiero et al. (2021) proposed that both the Madeira and Canary hotspots are  
77 fed by plume-like structures overlying the seismically slow Central-East Atlantic Anomaly  
78 (CEAA). This anomaly emerges from depths below 700 kilometres and maintains a con-  
79 nection with the African LLSVP (Large low-shear-velocity provinces). However, the two  
80 hotspot systems seem to be in different stages of development. According to this model,  
81 while the Canary plume-like structure is rooted in the CEAA, tomography results hint  
82 that the small Madeira plumelet (100 km width) appears detached from it. The plumelet  
83 extends from the surface down to at least 300 km, below which the resolution degrades.  
84 The presence of mantle upwellings beneath the Canary and Madeira archipelagos is also  
85 supported by new shear-wave splitting measurements (Schlaphorst et al., 2022). How-  
86 ever, in the Canary Province many of the predictions of plume theory (Morgan, 1972)  
87 are not met, e.g., volcano ages do not follow a consistent linear age–distance relation-  
88 ship, with coeval volcanism occurring across several hundreds of kilometres (Geldmacher  
89 et al., 2001, 2005). Insights from numerical modelling of mantle flow highlight the po-  
90 tential role of edge-driven convection in deflecting the mantle plume beneath the Canary  
91 archipelago, providing a plausible explanation for the complex age progression of the mag-  
92 matism and widespread volcanism (Manjón-Cabeza Córdoba & Ballmer, 2022).

93 Studying the seismic discontinuities beneath both hotspots can provide new insights  
94 into the mantle thermal and compositional properties of these regions, thus allowing the  
95 assessment of the origin of these hotspots. Seismic discontinuities mark discernible changes

96 in the density, shear and compressional wave speeds of the mantle as depth increases.  
 97 The occurrence, depth and seismic visibility of these changes are dependent on the tem-  
 98 perature and composition variability of the mantle. The upper mantle has predominantly  
 99 peridotite composition, where the dominant mineral is olivine, usually corresponding to  
 100 more than 50% volume, accompanied by orthopyroxene and, in most fertile compositions,  
 101 by clinopyroxene and an aluminous phase that changes with pressure (Ziberna et al., 2013,  
 102 e.g.) and references therein. In this dominant peridotite composition are dispersed eclogite  
 103 domains resulting from the recycling, via subduction, into the mantle of altered oceanic  
 104 crust (e.g. Allègre & Turcotte, 1986; Helffrich & Wood, 2001). Such predominant olivine-  
 105 rich parageneses change in depth as a result of several phase transitions (PT) that ex-  
 106 plain the observations of sharp seismic discontinuities at the base of the upper mantle:  
 107 at a depth of 410 km the olivine transforms to wadsleyite (ol→wa) with a positive Clapey-  
 108 ron slope, at 520 km the wadsleyite transforms to ringwoodite (wa→ri) with a positive  
 109 Clapeyron slope and at 660 km ringwoodite dissociates into bridgmanite ((Mg, Fe)-perovskite)  
 110 and magnesiowustite (ri→br+mw) with a negative Clapeyron slope (e.g. Helffrich, 2000;  
 111 Helffrich & Wood, 2001). The corresponding seismic discontinuities are named 410, 520  
 112 and 660, after the approximate depth at which the PT occur. The region of the man-  
 113 tle bounded by the 410 and 660 is known as the MTZ, an acronym of mantle tran-  
 114 sition zone.

115 The thermal and compositional structure of the mantle introduces complexities in  
 116 the PT. Temperature anomalies in the mantle move the PT to different pressures (depths)  
 117 according to their respective Clapeyron slopes (Bina & Helffrich, 1994). Due to the op-  
 118 posite sign in the Clapeyron slopes of the olivine-related PT responsible for the 410 and  
 119 660, the depth changes of these discontinuities are anti-correlated in response to a ther-  
 120 mal anomaly. Therefore, while the 410 becomes shallower in colder regions and deeper  
 121 in hotter ones, the depth changes of the 660 are opposite. Consequently, the MTZ is ex-  
 122 pected to be thicker when crossed by subducted slabs and thinner in high temperature  
 123 regions (Collier & Helffrich, 2001; Helffrich, 2000; Lawrence & Shearer, 2006; Vidale &  
 124 Benz, 1992). This expected behaviour of the MTZ has been used as an indication of ther-  
 125 mal anomalies, akin to a mantle thermometer (e.g. Helffrich, 2000).

126 The presence of other mineral phase transformations and variations in composi-  
 127 tion may introduce further complexities in the structure of the MTZ and in the man-  
 128 tle above. For example, at extremely high temperatures, a PT in majorite garnet occurs  
 129 at 660-700 km which is thought to become dominant (Weidner & Wang, 1998; Hirose,  
 130 2002; Xu et al., 2008). The majorite garnet PT has a positive Clapeyron slope, which  
 131 contrasts with the positive slope of the ringwoodite PT. As a result, in a hot mantle, the  
 132 660 would occur at greater depths. This could lead to a positively correlated depressed  
 133 topography on the 410 and 660 discontinuities with a minor overall impact on the thick-  
 134 ness of the MTZ (Deuss et al., 2006; Deuss, 2007). Furthermore, this PT takes place across  
 135 a broad depth interval, significantly broadening the 660 discontinuity.

136 Mantle upwellings rooted deep in the mantle not only exhibit high temperatures  
 137 but also different composition from the surrounding mantle. As part of mantle convec-  
 138 tion, cold, ancient oceanic crust is transformed into eclogite in subduction zones, trans-  
 139 ported to the lower mantle and then is brought back up into the upper mantle once en-  
 140 trained in hot mantle plumes. Because of their unique chemical signature, these plumes  
 141 are known as thermochemical plumes rather than purely thermal plumes (see Koppers  
 142 et al. (2021) for a review). The presence of a consistent seismic discontinuity approxi-  
 143 mately 300 km beneath various hotspots is considered strong evidence of the thermo-  
 144 chemical nature of deep-seated plumes (Schmerr, 2015; Kemp et al., 2019; Pugh et al.,  
 145 2021). This discontinuity named the 300 or X discontinuity, is attributed to the coesite-  
 146 to-stishovite (co→st) phase transition in regions of the mantle enriched in eclogite (Williams  
 147 & Revenaugh, 2005), as expected for a thermochemical plume.

Receiver functions and precursor studies have examined the MTZ beneath the Canary Islands at a regional scale, while those from Madeira are contributions from global studies. In the Canaries, studies indicate a thinner MTZ, possibly due to hot mantle material (Deuss, 2007; Martinez-Arevalo et al., 2013; Saki et al., 2015). For Madeira, studies show a moderately thinner MTZ, supporting the plume theory (Houser et al., 2008; Lawrence & Shearer, 2008; Deuss, 2009). However, a detailed analysis of MTZ discontinuities is lacking in both regions, especially for Madeira.

In this study, we investigate the fine structure of the upper mantle and MTZ beneath both archipelagos using receiver functions. In particular, we estimate the time, amplitude and frequency content of P-to-s converted phases from which we derived the depth, magnitude and width of velocity jumps of mantle discontinuities. Through a comprehensive interpretation of our findings in the context of mineral physics and geodynamics, we inferred the thermochemical nature of the Canary mantle plume. Beneath Madeira, our findings confirm the observations from Civiero et al. (2021) suggesting that the plume appears to be disconnected from the lower mantle and primarily affects the upper portion of the MTZ.

## 2 Data

For the Canary archipelago, we use data from 23 stations of the Spanish Digital Seismic Network (operated by the *Instituto Geográfico Nacional*) and one station from the Global Seismograph Network. These stations were active in varying periods, ranging from 2008 to 2017. For Madeira, we use data from 16 stations from the temporary DOCTAR network that was active during 2011/2012 (Matos et al., 2015), as well as 2 permanent stations of the Portuguese National Seismic Network (Instituto Português do Mar e da Atmosfera, I.P., 2006). See Figure 1 for station deployment and Table S1 in the supplementary material for the names and geographic coordinates of stations.

We selected the waveforms of 1241 events for the Canary archipelago and 1268 events for Madeira, of Mw between 5.5 and 7 and from epicentral distances between 30° and 95°. After discarding incomplete ZNE triplets, we obtain 13578 three-component waveforms for the Canary archipelago and 6481 for Madeira.

## 3 Method

To investigate the structure of the upper mantle, we employ receiver functions (RFs) and common conversion point stacking techniques.

### 3.1 Receiver functions

We use teleseismic *Pds* phases, i.e. the phases resulting from the direct *P* wave converting to an *S* wave at a seismic discontinuity at a depth *d*. Because of their polarization and their almost vertical incidence, *Pds* conversions are recorded on the radial (*R*) component of teleseismic earthquakes. Due to their arrival during the coda of the *P* phase has an abundance of other phases (e.g., multiple reflections and scattered waves), as well as the weak amplitude of the *Pds* phases, direct detection in individual seismograms is difficult. However, converted phases are expected to be coherent with the waveform of the main *P* arrival for conversion at discontinuities which are thinner than one-half of the *P*-wavelength (Richards, 1972; Paulssen, 1988; Bostock, 1999). Therefore, it is possible to extract them by waveform similarity. This can be done using deconvolution of the *P* phase in the vertical component (*P<sub>Z</sub>*) from *R*, which is known as the receiver function (RF) technique (Phinney, 1964; Vinnik, 1977; Langston, 1979; Ammon, 1991). The deconvolution process eliminates the source component and isolates the P-to-s conversion peaks from discontinuities beneath the station.

195 We compute RFs using the processing methods outlined in Bonatto et al. (2015,  
 196 2020). The processing workflow is illustrated in Figure 2a and described in detail in Sup-  
 197 plementary Section 1. After the processing and the quality controls, we obtain 948 RFs  
 198 for the Canaries and 320 RFs for the Madeira archipelago in five frequency bands: 0.02-  
 199 0.12 Hz, 0.02-0.2 Hz, 0.02-0.32 Hz, 0.02-0.48 Hz, 0.02-0.64 Hz. Table S1 summarises the  
 200 number of teleseismic records per station and the corresponding final number of good-  
 201 quality RFs.

### 202 3.2 Stacking of receiver functions with moveout correction

203 During the stacking process, we apply a linear moveout correction, also known as  
 204 a slowness stack, to account for the time delay between the arrival of seismic waves at  
 205 different distances from the source. This correction ensures that the seismic signals are  
 206 accurately aligned and stacked, resulting in improved signal detection and identification.

207 To perform the stacking process with a linear moveout correction, we use two tech-  
 208 niques: the phase weighted stack (PWS) (Schimmel & Paulssen, 1997) and the linear stack  
 209 (LS). The PWS is a non-linear stacking technique that incorporates the coherence of the  
 210 instantaneous phase as a weight in the linear stack, enhancing the signal-to-noise ratio  
 211 (SNR) and improving the accuracy of signal detection.

212 The stacking process with moveout correction, accounting for both stacking tech-  
 213 niques, is applied as follows:

$$S(s_j, t) = \frac{1}{N} \sum_{i=1}^N RF_i(t + s_j(\Delta_i - \Delta_{REF}))c_i \quad (1)$$

214 where  $s_j$  refers to a specific relative slowness with respect to the  $P$  phase,  $t$  is relative  
 215 time with respect to the  $P$  phase,  $\Delta_i$  is the epicentral distance for  $RF_i$ ,  $\Delta_{REF}$  is the ref-  
 216 erence epicentral distance, and  $N$  is the total number of stacked RFs. The reference dis-  
 217 tance,  $\Delta_{REF}$ , of  $75^\circ$  is chosen for all the slowness stacks to account for the distribution  
 218 of epicentral distances which are mainly concentrated between  $60^\circ$  and  $90^\circ$ .

219 The weight  $c_i$  in equation 1 is the phase stack and it accounts for the instantaneous  
 220 phase similarity:

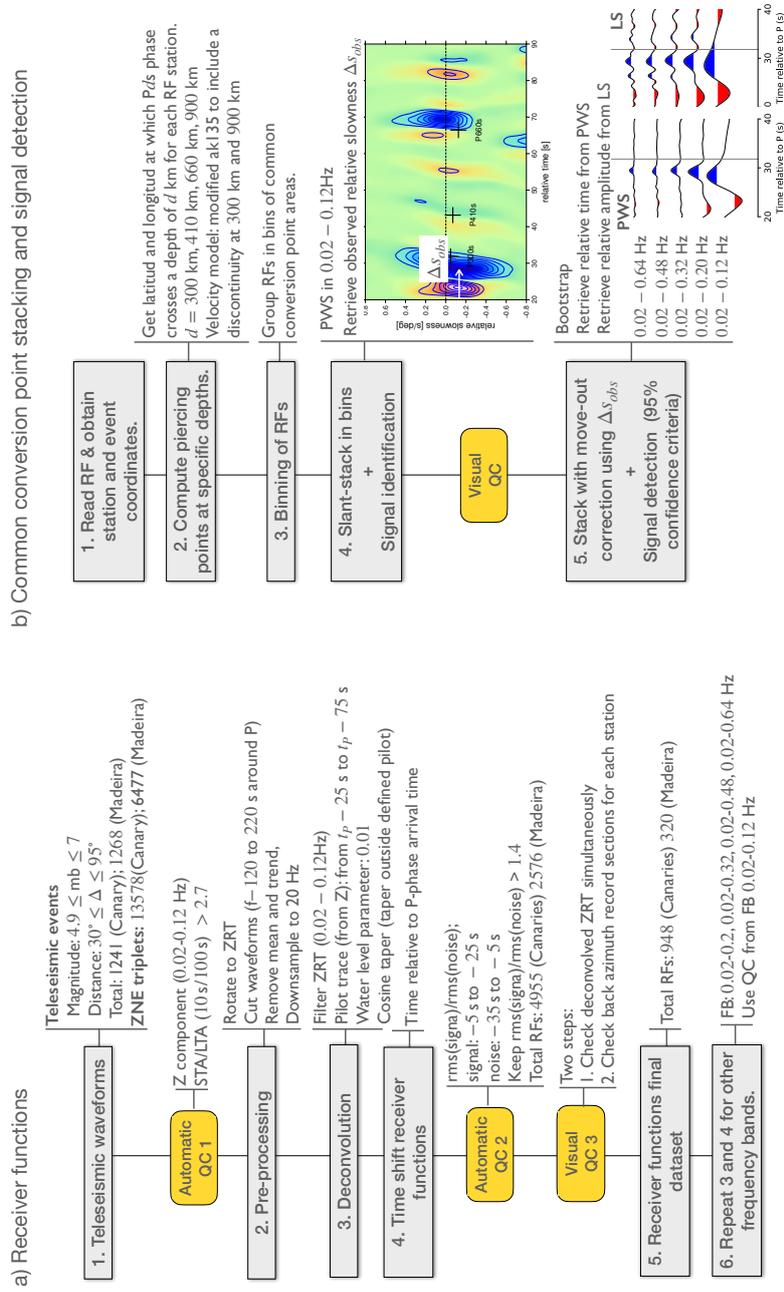
$$c_i = \left| \frac{1}{N} \sum_{k=1}^N e^{i\phi_k(t+s_j(\Delta_i-\Delta_{REF}))} \right|^\nu \quad (2)$$

221  $\phi_k$  is the instantaneous phase for the  $k$ -th RF, and the parameter  $\nu$  controls the weight-  
 222 ing of phase similarity and dissimilarity. The move-out corrected LS is defined with  $\nu =$   
 223  $0$  ( $c_i = 1$ ), while for the PWS we choose a value of  $\nu = 2$ .

#### 224 3.2.1 Identification of converted phases in global stacks

225 The average fine structure of the upper mantle in the study area is obtained by stack-  
 226 ing all the RFs that passed the quality controls in the lowest frequency band (0.02–0.12  
 227 Hz). This global stack is useful for identifying converted phases from major discontinu-  
 228 ities and from minor reflectors that are prevalent throughout the area.

229 Figures 3 a and b display relative time-distance stacks of RFs for the Canaries and  
 230 Madeira datasets, respectively, with the  $P$  phase set as time zero. Each trace at a given  
 231 epicentral distance is computed from the stack (PWS) of RFs with epicentral distances  
 232 within a  $5^\circ$  interval and uses a fixed relative slowness parameter of  $-0.1$  s/ $^\circ$  to perform  
 233 the move-out correction within that distance interval. Since converted phases have a steeper



**Figure 2.** Processing workflow of a) Receiver functions and b) Common conversion point stacking of RFs.

234 angle of incidence than the direct  $P$  phase,  $Pds$  phases are always expected to align along  
 235 a negative slope line (i.e., negative relative slowness), as shown with the continuous black  
 236 lines. Conversely, multiples have a shallower angle of incidence and a corresponding posi-  
 237 tive slope (i.e., positive relative slowness), as shown with the dashed lines. Slowness stacks,  
 238 also referred to as vespagrams or relative time-slowness stacks, are a useful tool for un-  
 239 ambiguously identifying  $Pds$  phases. Figure 3 c and e show the slowness stacks for the  
 240 Canary dataset and Figure 3 d and f show the slowness stack for the Madeira dataset.  
 241 Here, all RFs are stacked applying a linear move-out correction for a range of slowness  
 242 values and the result is plotted in the relative time-slowness domain, where the refer-  
 243 ence at zero is the  $P$  phase. Conversions from discontinuities in the upper mantle are  
 244 expected with negative relative slowness, while multiples are expected with positive rel-  
 245 ative slowness.

### 246 **3.2.2 Common conversion point stacking and signal detection**

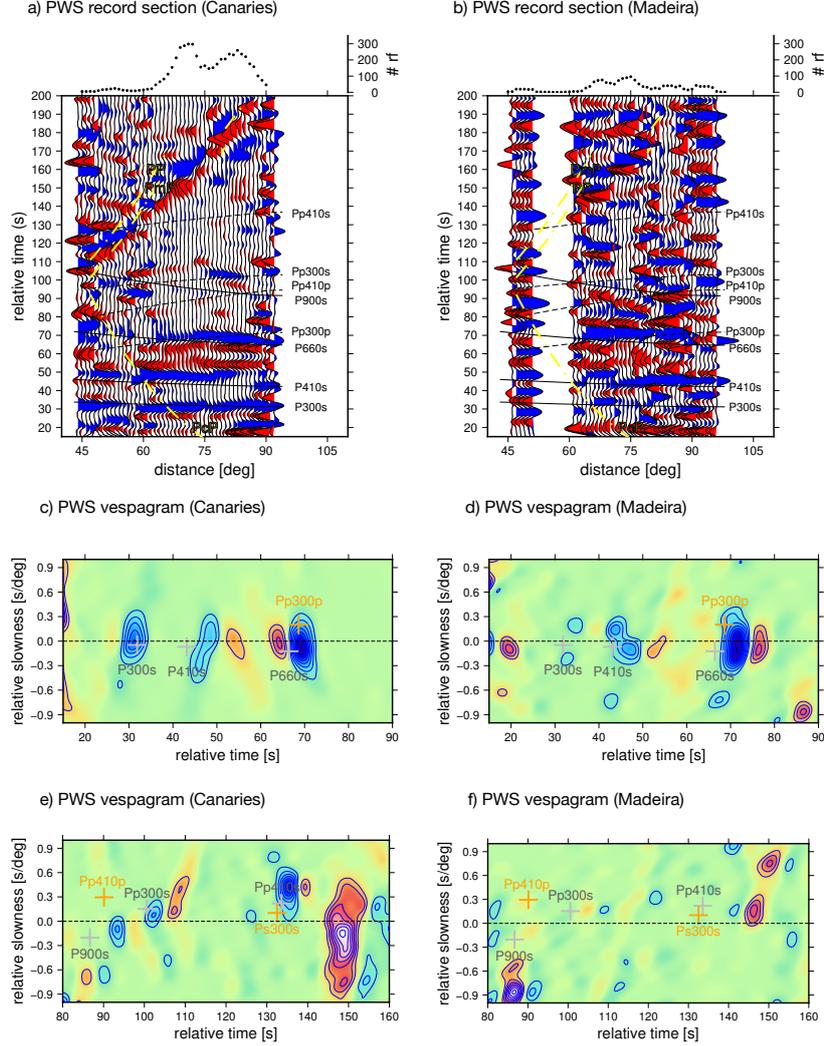
247 The lateral variations of the detected discontinuities are investigated using a common-  
 248 conversion-point (CCP) stacking approach as in Bonatto et al. (2020). The processing  
 249 workflow for the CCP stacking and signal detection is illustrated in Figure 2b. We first  
 250 retrieve station and event coordinates for each RF and compute piercing points of con-  
 251 verted phases  $P300s$ ,  $P410s$  and  $P660s$  at their corresponding conversion depths: 300  
 252 km, 410 km, and 660 km. The CCP stacking is performed for these three conversion depths,  
 253 which correspond to the reference depth of our target discontinuities. For each conver-  
 254 sion depth, we stack the RFs with conversion point coordinates in the same circular area  
 255 or bin. To ensure reliable measurements, we set the bin radius for each conversion depth  
 256 to be approximately the same as the size of the first Fresnel Zone at an intermediate fre-  
 257 quency of 0.2 Hz:  $r_{P300s, P410s} = 0.5^\circ$  and  $r_{P660s} = 0.75^\circ$ . The centres of these bins  
 258 are equally spaced every  $0.7^\circ$  to enable an overlap of the sampled areas. If the number  
 259 of RFs in the bin is less than 25, we increase the bin radius to  $r_2 = 1.5r_1$ , or to  $r_3 =$   
 260  $2r_1$  if numbers remain low. For each bin and conversion depth, we first stack the RFs  
 261 and perform a visual quality control on individual slowness stacks in the lowest frequency  
 262 band 0.02 – 0.12 Hz. We assign a quality level, ranging from 1 to 3, to each stack de-  
 263 pending on the following criteria:

- 264 1: The signal identification is very clear, without other signals near the target phase.
- 265 2: Other coherent signals near the target phase, but the signal is still clear.
- 266 3: Multiple coherently stacked signals of similar amplitude near the target signal or  
 267 non-detection.

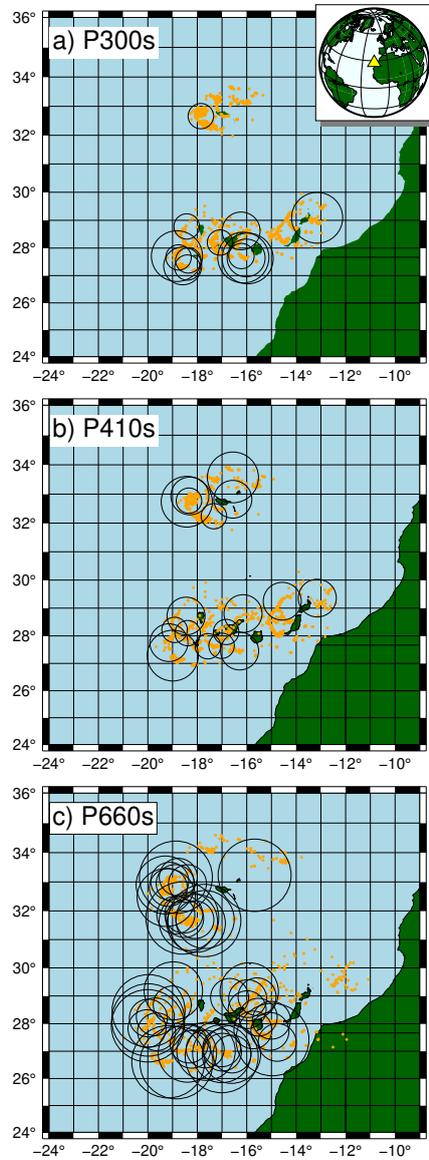
268 Representative examples of stacks classified as quality 1, 2, and 3 for  $P300s$  are shown  
 269 in Figs. S1 to S11 of Supplementary Section 2. Stacks with quality 1 and 2 are consid-  
 270 ered as good quality stacks and stacks of quality 3 are not considered in the analysis.

271 We use the good quality stacks to estimate the relative slowness at which each phase  
 272 is detected. Then, we compute the stacking with move-out correction using this slow-  
 273 ness value; i.e., fixed  $s_j$  in Eq. 1. Mean and standard deviation of stacks are calculated  
 274 using a bootstrap resampling algorithm with replacement (i.e., a waveform may be cho-  
 275 sen multiple times). Subsequently, we measure the relative travel time values from both  
 276 stacks (PWS and LS) and the amplitude of converted phases from LS. We only include  
 277 signals in the analysis that meet the 95% confidence criteria, which requires the ampli-  
 278 tude of a detected signal to be at least twice its amplitude standard error.

279 Finally, we retain only those detected signals from good-quality stacks with more  
 280 than 25 stacked RFs. These signals should also show negative slowness and positive am-  
 281 plitude, as expected for the target-converted phases. The CCP bins with robust detec-  
 282 tions of converted phases are marked on the maps in Figure 4.



**Figure 3.** Record sections for (a) the 948 RFs from the Canary Island dataset and (b) the 320 RFs from the Madeira dataset. Each trace in the record sections results from stacking (PWS) the RFs with epicentral distances within a  $5^\circ$  interval around a given epicentral distance in the horizontal axis. Blue and red colours represent positive and negative amplitudes. Black lines show the theoretical travel time curves for a slightly modified ak135 reference model that includes imperceptible density increase at 300 km and 900 km. Dashed yellow lines represent travel time curves for potentially interfering phases (PcP, PmP, PP). The number of RFs is indicated at the top of the record section. (c)-(f) Stacking (PWS) in the relative time-slowness domain for the Canary Islands and Madeira dataset in different time intervals. Crosses mark the expected time and slowness of different phases for positive (gray) and negative polarities (orange) in the modified ak135 reference model.



**Figure 4.** Location of common conversion point bins with robust detections (with PWS and LS) of the converted phases: a) P300s, b) P410s, and c) P660s. Circles correspond to CCP bins with signals that satisfy the 95% confidence criteria. The size of the circles are proportional to the radius of the CCP bin; from the smallest to the largest, the radii of the circular sectors are:  $0.5^\circ$ ,  $0.75^\circ$ ,  $1.0^\circ$ ,  $1.25^\circ$ ,  $1.5^\circ$ . Orange dots indicate the piercing points of corresponding phases.

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### 3.3 Time-to-depth conversion

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To convert relative times to depth, we use tomography-derived velocities. Velocity anomalies are extracted from two tomography models: the global model MIT08 (Li et al., 2008) and the regional model IBEM-P18 (for P waves)/IBEM-S19 (for S waves) (Civiero et al., 2018; Civiero, Custódio, et al., 2019). While the regional model has a higher resolution, it does not cover the entire area of interest. However, we include it in the analysis to assess the consistency of the results in the overlapping region with both models. Since MIT08 only accounts for P-wave anomalies, we approximate corresponding S-wave anomalies to be 1.3 times P-wave anomalies at the surface, increasing linearly to 3 times at the base of the mantle, following Ritsema and Van Heijst (2002). To account for possible resolution overestimation given by noiseless synthetic tests (Li et al., 2008; Civiero, Armitage, et al., 2019), we assume that the amplitudes of anomalies recovered by the inversion could be about 50% than those of unknown real anomalies. Therefore, we compute the corrections using MIT08 and IBEM-P18/S19 anomalies multiplied by a factor of 2.

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For each tomography model, we compute time corrections for each CCP bin and converted phase. We use TauP toolkit (Crotwell et al., 1999) to create a modified *ak135* velocity model for each bin which includes the anomalies. We use this new velocity model to compute the theoretical relative travel times of converted phases (for a reference distance of  $75^\circ$ ). Then, we use these relative travel times to compute time corrections with respect to *ak135*, i.e., the difference between travel times obtained with *ak135* and with the modified *ak135*. Finally, we add the correction to the observed relative times of detected signals and the corrected time values are converted to depth using the reference velocity model *ak135* for a reference distance of  $75^\circ$ . Although the tomographic models used to compute the corrections have large uncertainties, it should be noted that the overall effect of the corrections is to locate the discontinuities closer to their actual depth.

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### 3.4 Temperature variations in the MTZ

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The thickness of the mantle transition zone (MTZ) at each location is calculated by subtracting the depths of the 410 and 660 discontinuities ( $z_{660} - z_{410}$ ). Variations in the MTZ thickness can be attributed to temperature variations, as thermal anomalies affect the depths of the 410 and 660 discontinuities differently.

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The conversion of anomalous MTZ thickness to temperature is accomplished through the use of Clapeyron slopes associated with the olivine phase transitions (refer to Table S2). Clapeyron slopes ( $dP/dT$  or  $\gamma$ ) can be written as

$$\left(\frac{dP}{dT}\right)_H \approx \frac{dP}{dz} \frac{\delta z_H}{\delta T} \quad (3)$$

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where  $H$  stands for 410 or 660,  $dz_H$  corresponds to  $dz_{410} = 410 - z_{410}$  or  $dz_{660} = 660 - z_{660}$  and  $dP/dz$  is the upper-mantle pressure increase with depth which is about  $100/3$  MPa  $\text{km}^{-1}$  in *PREM* (Dziewonski & Anderson, 1981).

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As proposed by Helffrich (2000) (2000), we adopt the assumption that the deflection of discontinuities is due to vertically consistent temperature changes. Therefore, we can convert the difference in MTZ thickness to temperature variation using the equation below

$$\delta T = \left[ \left(\frac{dP}{dT}\right)_{660} - \left(\frac{dP}{dT}\right)_{410} \right]^{-1} \frac{dP}{dz} \delta MTZ \quad (4)$$

324 where  $\delta MTZ = MTZ_{thickness} - 250$  km or  $\delta MTZ = MTZ_{thickness} - 242$  km, if we  
 325 consider the reference value for the MTZ thickness from the global RFs study of Lawrence  
 326 and Shearer (2006).

### 327 **3.5 Frequency and amplitude analysis of converted phases**

328 Studying the frequency visibility and the amplitude of seismic phases provides im-  
 329 portant constraints on the sharpness and velocity contrast of the velocity gradients as-  
 330 sociated with mantle discontinuities. (Shearer & Flanagan, 1999; van der Meijde et al.,  
 331 2003; Juliá, 2007; Bonatto et al., 2020; Pugh et al., 2021). These parameters and the dis-  
 332 continuity depth are directly linked to the underlying mechanisms that give rise to the  
 333 seismic boundaries and play an important role in interpreting our observations.

#### 334 **3.5.1 Sharpness of discontinuities**

335 Seismic discontinuities are generally approximated as linear velocity gradients, which  
 336 act as a low pass filter to converted waves. In particular, the amplitudes of the converted  
 337 waves at the discontinuity are significant when the gradient transition interval,  $\Delta z$ , is  
 338 smaller than half of the wavelength,  $\lambda_P$ , of the incident  $P$ -wave (Bostock, 1999; Paulssen,  
 339 1988). This is

$$\Delta z < \frac{\lambda_P}{2} = \frac{v_P}{2f} \quad (5)$$

340 where  $v_p$  is the  $P$ -wave velocity at the corresponding discontinuity depth. Therefore, to  
 341 constrain the thickness of discontinuities, we need to find the highest frequency  $f^{max}$   
 342 for which the amplitude of converted phases is visible. We estimate this frequency value by  
 343 analyzing the amplitudes of converted phases in the CCP linear stacks across the dif-  
 344 ferent frequency bands that we use to compute the RFs. We use the following criteria:  
 345  $f^{max}$  is the frequency above which the amplitude drops below 60-70% of the amplitude  
 346 at the lowest frequency. Finally, the upper bound for the thickness of the discontinuity  
 347 is obtained by substituting  $f$  with  $f^{max}$  in Eq. 5. This rule of thumb has been used be-  
 348 fore to constrain the thickness of mantle discontinuities (e.g. Paulssen, 1988; van der Mei-  
 349 jde et al., 2003; Jenkins et al., 2017; Bonatto et al., 2020).

#### 350 **3.5.2 Velocity contrast**

351 Reflection and transmission coefficients describe how the energy of a seismic wave  
 352 is partitioned when it encounters a boundary between two elastic media. They are ob-  
 353 tained as the amplitude ratio of the incoming wave and the reflected or transmitted waves.  
 354 The coefficients depend on several factors, including the velocities and densities of the  
 355 seismic waves on both sides of the interface (the velocity and density contrasts), and the  
 356 angle of incidence of the incoming wave. The elementary formulas to compute the co-  
 357 efficients can be found in equations 5.39 and 5.40 of Aki and Richards (2002).

358 To estimate the velocity contrast across each detected discontinuity, we manually  
 359 compare the theoretical coefficients of converted phases, assuming different velocity con-  
 360 trasts, with the estimated amplitudes of the  $Pds$  phases in our stacked RFs. This com-  
 361 parison allows us to obtain a rough assessment of the velocity jump. To minimize atten-  
 362 uation effects due to the finite thickness of discontinuities, we perform this analysis for  
 363 RFs in the lower frequency band: 0.02-0.2 Hz. In this approach, we assume that  $Pds$  phases  
 364 are primarily sensitive to S-wave velocity contrasts, as demonstrated by Juliá Juliá (2007).

## 4 Results

We investigate the upper mantle and MTZ structure using 948 high-quality RFs around the Canaries and 320 around Madeira. We first create time-distance stacks and vespagrams for the entire dataset of each archipelago to identify converted phases of possible discontinuities prevalent throughout the area.

The number of events observed at stations in the Canary archipelago that pass the quality controls are higher across almost the entire distance range, resulting in stronger signals of the major converted phases in the time-distance stacks (Figures 3 a, b). The vespagrams confirm that we obtain clear converted phases from the globally detected 410 and 660 discontinuities beneath both volcanic provinces (Figures 3 c, d). However, a clear converted phase at a reflector near 300 km depth is only observed beneath the Canaries. Based on the robust detections of the 410 and 660, we have confidence that this measurement of the X discontinuity is not an anomaly caused by data coverage. In contrast to the Canary archipelago, no strong multiples (Pp300s and Pp410s) from the reflectors at the X and 410 discontinuities can be observed around the Madeira archipelago (Figures 3 e, f). Likewise, a converted phase at a discontinuity of approximately 800-900 km depth can be observed around the Canaries but is not present around Madeira. This phenomenon will be subject to further investigation in a future study.

### 4.1 Topography of discontinuities and MTZ thickness

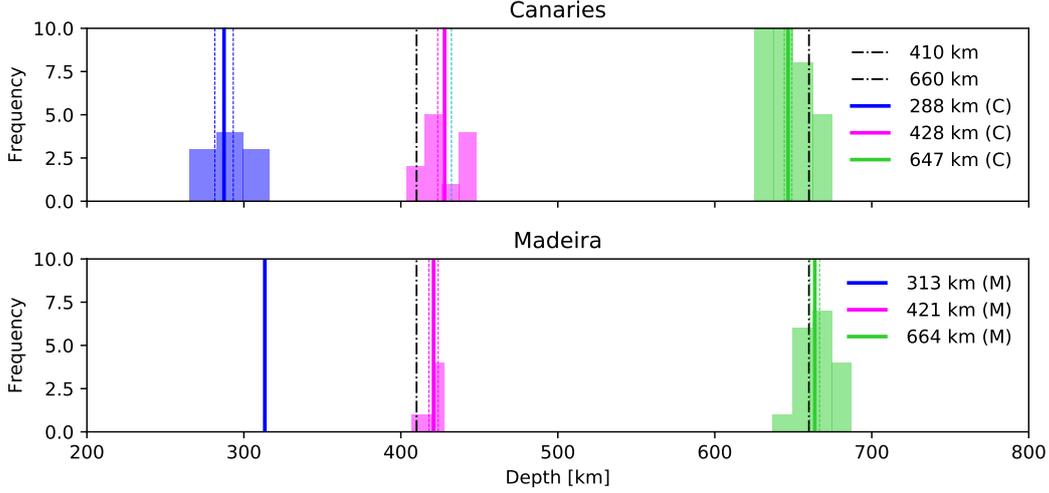
In order to investigate the spatial distribution of the converted phases, we apply the CCP stacking method. Then, we use the relative times of the converted phases to estimate the depth of the discontinuities from tomography-derived velocities with the global model MIT08 (Li et al., 2008).

Our analysis indicates that the X discontinuity is situated at a mean depth of  $288 \pm 6$  km ( $\bar{x} \pm 1\sigma_{\bar{x}}$ ) beneath the Canaries (Figure 5), with a range spanning from 251 km to 316 km. It is prominent around the western and central part of the Canaries with depths shallower than 300 km towards the south and deeper than 300 km towards the north (Figure 6a), likely stemming from a larger azimuthal coverage from the west and south. For Madeira, we have only retained one reliable detection of this discontinuity in the west, with a depth of 313 km.

On average, the 410 is located deeper beneath the Canaries (Figure 5) with a mean depth of  $428 \pm 4$  km and a range of 403 km to 448 km, in comparison to a mean depth of  $421 \pm 3$  and a range of 407 km to 428 km beneath Madeira. Whereas in the Canaries it is located deeper towards the western part of the archipelago with values larger than 410 km southwest of the islands of La Palma, Tenerife and Gran Canaria, and smaller values to the northeast, a lateral trend is not as evident around the Madeira archipelago. This difference can be attributed to Madeira's smaller geographical extent compared to the Canaries, as well as the fact that the stations collectively cover a significantly smaller area. We observe, larger values in the northwest and a shallower value towards the south, making this pattern less pronounced.

In contrast, the 660 is shallower beneath the Canary archipelago with a mean depth of  $647 \pm 2$  km and ranges of 625 km to 675 km, in comparison to a mean depth of  $664 \pm 3$  km and a range of 637 km to 687 km beneath Madeira. Based on the piercing point distribution, nearly all results can be found towards the west of Madeira and around the western and central part of the Canaries (Figure 6c). In Madeira, the 660 is detected at larger depths towards the west, whereas it is shallower towards the east.

The general trend of a deeper 410 and a shallower 660 in the Canaries results in a thinner MTZ there at  $216 \pm 2$  km compared to  $237 \pm 2$  km around Madeira (Figures 6 d). These MTZ thickness values are derived from depth estimates of the 410 and 660 dis-



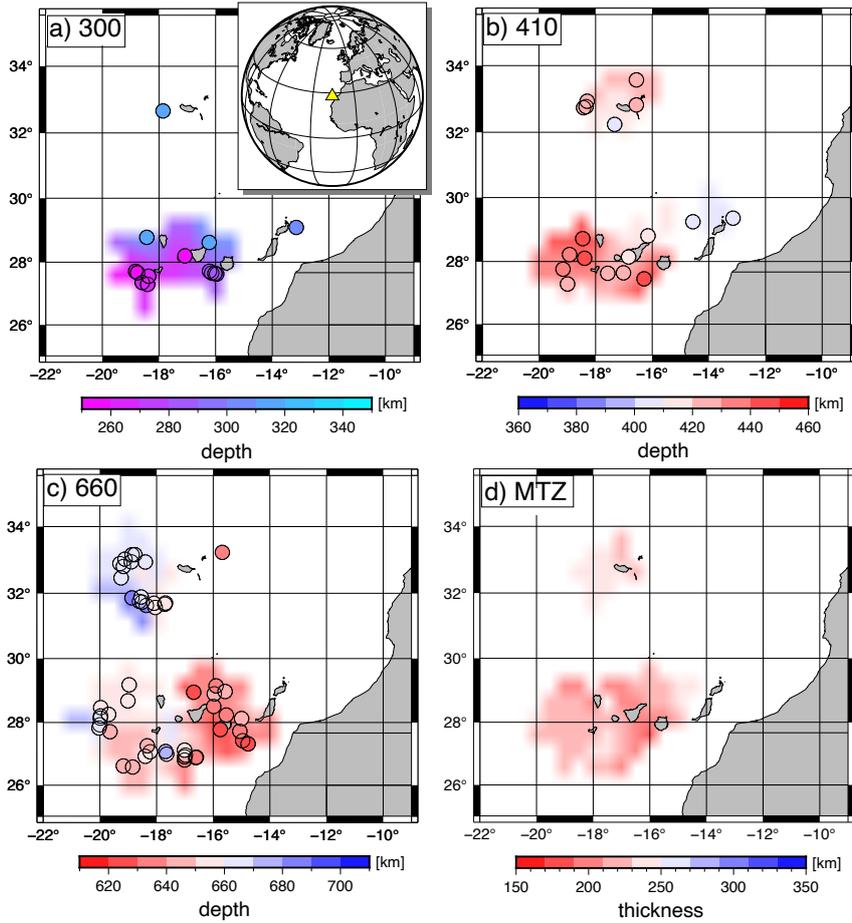
**Figure 5.** Histograms of estimated depth values, after time corrections, for the X, 410, and 660 discontinuities beneath the Canaries and Madeira. Vertical solid lines (blue, pink, green and black) correspond to mean depth values for both archipelagos (given in the legend, where M stands for Madeira and C for the Canaries), dashed lines correspond to standard errors of mean values and dash-dotted lines are reference depth values in a modified ak135 model that incorporates a discontinuity at 300 km.

414 continuities, constrained to regions where observations for both discontinuities are avail-  
 415 able.

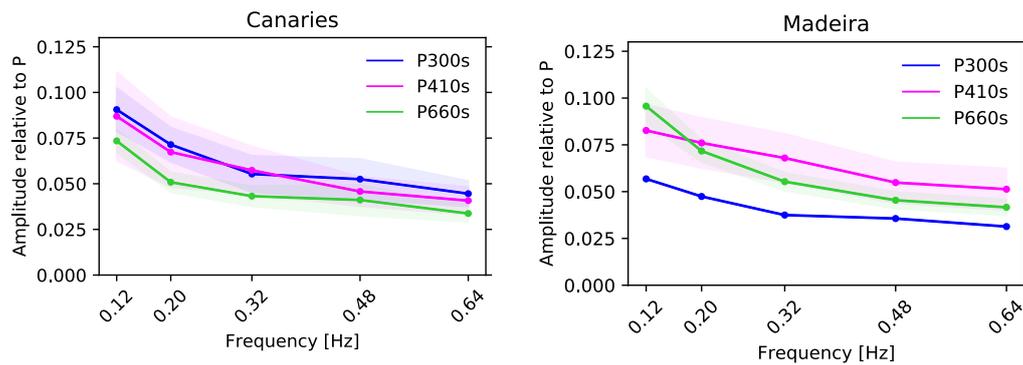
## 416 4.2 Sharpness and velocity contrast of detected discontinuities

417 Figure 7 shows the average amplitudes of the  $P300s$ ,  $P410s$ , and  $P660s$  phases relative  
 418 to  $P$  across five frequency bands for both archipelagos. As expected for linear ve-  
 419 locity gradients, the amplitude of all converted phases decreases with higher upper-frequency  
 420 boundaries. We apply the criteria defined in section 3.5 to constrain the thickness of the  
 421 discontinuities. In the Canary archipelago,  $P300s$ ,  $P410s$ , and  $P660s$  exhibit an ampli-  
 422 tude reduction of  $\sim 30\text{-}40\%$  (w.r.t the lowest frequency band) at 0.32 Hz. Therefore, we  
 423 consider  $f^{max} = 0.32$  Hz for all converted phases, indicating that the detected discon-  
 424 tinuities are sharp. Using Eq. 5 and corresponding values of  $v_p$  in ak135, we obtain  $\Delta z_{300,Can} <$   
 425  $13$  km,  $\Delta z_{410,Can} < 14$  km, and  $\Delta z_{660,Can} < 15$  km. In Madeira,  $P300s$  and  $P410s$   
 426 exhibit an amplitude drop of  $\sim 30\text{-}40\%$  at 0.48 Hz, leading to sharper discontinuities:  $\Delta z_{300,Mad} <$   
 427  $9$  km,  $\Delta z_{410,Mad} < 9$  km. For  $P660s$ , the maximum frequency for which significant en-  
 428 ergy is converted is 0.32 Hz, which corresponds to  $\Delta z_{660,Mad} < 15$  km.

429 In both archipelagos the relative amplitude of  $P300s$ ,  $P410s$ , and  $P660s$  relative  
 430 to  $P$  is roughly 0.04 to 0.09 (Figure 7). As previously noted, we observed that as the upper-  
 431 frequency boundaries increase, the amplitudes of all signals decrease. However, whereas  
 432 the amplitude of  $P300s$  is close to the amplitude of  $P410s$  across all frequency bands  
 433 in the Canaries ( $0.09 \pm 0.01$  at 0.02-0.12 Hz to  $0.04 \pm 0.01$  at 0.02-0.64 Hz), it is the weak-  
 434 est around the Madeira archipelago ( $0.06 \pm 0.02$  at 0.02-0.12 Hz to  $0.03 \pm 0.01$  at 0.02-0.64  
 435 Hz). This discrepancy is further emphasized by the fact that only one robust measure-  
 436 ment of this discontinuity could be obtained there. Our observed relative amplitudes at  
 437 0.02-0.20 Hz for the Canary and Madeira archipelagos are displayed in Table 1. The ta-



**Figure 6.** Maps showing discontinuity depth (a, b, c) and MTZ thickness surfaces (d), derived from depth estimates of the 410 and 660 discontinuities, constrained to regions where observations for both discontinuities are available. Circles with thicker lines correspond to robust detections in the PWS and LS, i.e., signals that satisfy the 95% confidence criteria. The ones with a thinner line correspond to robust detection only in the LS (the signal is visible in the PWS but it does not satisfy the 95% confidence criteria).



**Figure 7.** Change of relative amplitude of converted phases for different higher cut-off frequencies for the Canary (left) and Madeira (right) archipelagos. Values are given as a fraction of the initial *P*-phase amplitude. The shaded areas show uncertainties. Note, that *P*300s in Madeira consists of only one measurement.

	phase	observed amplitude at 0.2 Hz	theoretical coefficients	$v_S$ jump [%]	relative $v_S$ jump 410/300	410/660
<i>Madeira</i>						
	P300s	0.047	$0.047 \pm 0.001$	5.1	1.6	1.1
	P410s	$0.076 \pm 0.013$	$0.076 \pm 0.002$	8.1		
	P660s	$0.072 \pm 0.005$	$0.072 \pm 0.002$	7.6		
<i>Canaries</i>						
	P300s	$0.071 \pm 0.009$	$0.071 \pm 0.002$	7.6	0.96	1.3
	P410s	$0.067 \pm 0.019$	$0.067 \pm 0.002$	7.3		
	P660s	$0.051 \pm 0.005$	$0.051 \pm 0.002$	5.5		

**Table 1.** Observed amplitudes and theoretical coefficients used to constrain the velocity jumps ( $v_S$  jump) at target discontinuities. The theoretical coefficient is computed by averaging the transmission coefficients across a range of epicentral distances from  $60^\circ$  to  $90^\circ$ .

438 ble also shows the estimated velocity jumps and their corresponding theoretical trans-  
439 mission coefficients.

## 440 5 Discussion

### 441 5.1 Mantle transition zone in the Canary and Madeira archipelagos

442 In Supplementary Section 3, we examine the time corrections and the depth cor-  
443 relations between the topography of the 410 and 660 discontinuities for both tomographic  
444 models, IBEM-P18/-S19 and MIT08, to assess the accuracy of our time-to-depth con-  
445 versions. Based on this analysis, we conclude that MIT-P08 likely places the 410 and 660  
446 discontinuities slightly deeper than their true depths. Conversely, the MTZ thickness re-  
447 mains unaffected by the mantle structure above it. The MTZ thickness is primarily in-  
448 fluenced by the internal structure of the MTZ, characterized by smaller anomalies and  
449 consequently smaller errors. This is evident in Figure S14, showing similar time correc-  
450 tions for the *P410s* and *P660s*.

#### 451 5.1.1 Previous studies

452 A thinned MTZ has been found beneath the Canary and Madeira archipelagos in  
453 previous studies using RFs and SS/PP precursors. The effect is more pronounced ( $\sim 15$ –  
454  $20$  km thinner than the global average) beneath the Canaries (Li et al., 2003; Deuss, 2007;  
455 Gu et al., 2009; Martinez-Arevalo et al., 2013; Saki et al., 2015, e.g.), but still evident  
456 beneath Madeira (Houser et al., 2008; Lawrence & Shearer, 2008; Deuss, 2009, e.g.). There  
457 are different interpretations for the thinning of the MTZ. Whereas some studies explain  
458 those anomalies through the presence of hot mantle material, thus supporting the the-  
459 ory of existing plumes (Gu et al., 2009, e.g.), other studies find that the thinning is sim-  
460 ilar to the widespread thinning of the MTZ beneath oceans observed from SS precur-  
461 sor studies; this thinning is not solely restricted to plumes (Li et al., 2003; Martinez-Arevalo  
462 et al., 2013, e.g.).

463 The study of Civiero et al. (2021), based on results from high-resolution seismic  
464 tomography (Civiero et al., 2018; Civiero, Custódio, et al., 2019), shear-wave splitting  
465 (Schlaphorst et al., 2022) and gravity (Sandwell et al., 2014) along with plate reconstruc-  
466 tion showed that the Madeira and Canary volcanism is fed by plume-like structures over-  
467 lying the seismically slow Central-East Atlantic Anomaly (CEAA) developing as a dome-  
468 like instability beneath  $700$  km depth in connection to the African LLSVP. However, the

469 two hotspot systems seem to be in different stages of development. While the Canary  
 470 plume-like structure is still rooted in the CEEA, tomography results suggest that the  
 471 Madeira plumelet is presently detached from it, extending from the surface down to at  
 472 least 300 km depth (below which the resolution degrades), which was interpreted as a  
 473 result of the Madeira plumelet being probably at a dying stage.

### 474 **5.1.2 Our study**

475 Our results support the presence of a hot rising upwelling beneath the Canary Is-  
 476 lands that extends through the MTZ. We find that the MTZ is thinner in the Canary  
 477 archipelago (Figure 6). This is attributed to a depressed 410 and an uplifted 660, which  
 478 is consistent with the MTZ being crossed by a hot mantle upwelling. The negative cor-  
 479 relation observed between the depths of the 410 and 660 discontinuities beneath the Ca-  
 480 nary Islands confirms that the polymorphs olivine and ringwoodite are the dominant min-  
 481 erals influencing the topography of both discontinuities.

482 In Madeira, the MTZ is also thinner, primarily due to a depressed 410 and a stan-  
 483 dard 660, suggesting that the thermal anomaly predominantly affects the 410 disconti-  
 484 nuity. This is compatible with a low-velocity anomaly elongating downward to the top  
 485 of the MTZ, i.e. extending more than 100 km deeper than the bulk of the anomaly im-  
 486 aged in the model by Civiero et al. (2021). The hypothesis of a possible deeper exten-  
 487 sion of the Madeira plumelet was admitted by the authors, on the basis of considerations  
 488 about the size of the anomaly and the coarser model resolution below this area. Our re-  
 489 sults provide complementary information that reinforces the hypothesis that the plumelet  
 490 extends deeper than 300 km, but does not fully cross the MTZ and is detached from the  
 491 CEEA.

492 By utilizing the Clapeyron slope of the olivine and ringwoodite PT, it becomes pos-  
 493 sible to estimate the total temperature variations across the MTZ beneath both archipela-  
 494 gos. If we substitute  $\gamma_{660} = -2.0$  MPa/K,  $\gamma_{410} = +4.00$  MPa/K and a reference MTZ  
 495 thickness of 242 km from Lawrence and Shearer (2006) global RFs study in Eq. 4, the  
 496 average MTZ thickness beneath the Canary archipelago ( $216 \pm 2$  km) translates to a  
 497 temperature difference of  $+144 \pm 11$  K. Similarly, the average MTZ thickness beneath  
 498 the Madeira archipelago ( $237 \pm 2$  km) translates to a temperature difference of  $+33 \pm$   
 499  $11$  K. These estimates provide insights into the thermal structure of the mantle anoma-  
 500 lies beneath these volcanic regions. Excess temperatures within the mantle beneath these  
 501 volcanic islands were independently confirmed using the chemical composition of the olivines,  
 502 the liquidus phase of the alkaline magmas erupting on them (Putirka, 2008).

## 503 **5.2 Evidences of thermochemical plume beneath the Canaries**

504 The detection of *PXs* (also named *P300s* throughout the article) serves as com-  
 505 pelling evidence for the existence of chemical heterogeneities within the plume material.

### 506 **5.2.1 The X discontinuity**

507 The X discontinuity has been observed in a variety of tectonic environments: sub-  
 508 duction zones (Zhang & Lay, 1993; Schmerr, 2015), stable continents (Wajeman, 1988;  
 509 Pugh et al., 2023), several hotspots (Courtier et al., 2007; Bagley & Revenaugh, 2008;  
 510 Schmerr, 2015; Kemp et al., 2019; Pugh et al., 2021, 2023), and ridges (Schmerr, 2015).  
 511 However, it seems not to be omnipresent throughout the planet (e.g. Schmerr et al., 2013;  
 512 Williams & Revenaugh, 2005). This suggests that such discontinuity is probably linked  
 513 to some form of mantle heterogeneity(ies) promoting a specific PT. These observations  
 514 over a broad range of possible mantle temperatures have led to several explanations for  
 515 the X (summarized in Table S3). Out of the different mechanisms that have been pro-  
 516 posed to explain the X discontinuity, the silica phase transition from coesite to stishovite

517 has been suggested as the most plausible explanation for the X in hot mantle upwellings,  
 518 such as the Canary archipelago, where it was previously observed (Pugh et al., 2021).  
 519 The presence of this discontinuity implies the presence of polymorphic silica-rich phases  
 520 beneath the study region at these depths. Within a plume setting, this material might  
 521 correspond to recycled eclogite, more silica-rich than the dominant peridotite, aligning  
 522 with thermochemical plume models (Kemp et al., 2019; Pugh et al., 2021, 2023).

523 The presence of eclogite beneath hotspots may be explained by bidirectional mass  
 524 transfer in the dynamic Earth’s interior, in agreement with a general model first proposed  
 525 by Hofmann and White (1982). The basaltic oceanic crust is recycled in the mantle through  
 526 subduction, leading to the production of eclogitic rocks, denser than the dominant ul-  
 527 trabasic compositions (peridotites in the upper mantle) (Niu, 2018). This contributes  
 528 to the generation of chemical heterogeneities that can eventually accumulate close to the  
 529 base of the lower mantle (slab graveyards; e.g., van der Meer et al., 2018; Jones et al.,  
 530 2021) explaining partially the origin of LLSVPs (Koppers et al., 2021). These LLVPs,  
 531 which are located beneath Africa and the Pacific (e.g. French & Romanowicz, 2015) have  
 532 been considered to act as plume nurseries for major deep-mantle-sourced hotspots like  
 533 the Canaries one (Koppers et al., 2021). These LLVPs-rooted upwellings carry recycled  
 534 oceanic crust (containing eclogite-type materials) to the upper levels of the mantle (Hofmann  
 535 & White, 1982; Tackley, 2000; Koppers et al., 2021). Owing to uranium enrichment of  
 536 oceanic crust during sea-water alteration and lead depletion during subduction-induced  
 537 dehydration, high  $^{206}\text{Pb}/^{204}\text{Pb}$  is considered the best geochemical proxy for the pres-  
 538 ence in the mantle of ancient recycled oceanic crust, which is considered the cause for  
 539 the HIMU mantle component (e.g. White, 2015). Lavas from both Canary (Gurenko et  
 540 al., 2006, 2009; Day et al., 2010) and Madeira (e.g. Mata et al., 1998; Geldmacher & Ho-  
 541 hornle, 2000; Gurenko et al., 2013) show evidence for the contribution of the HIMU man-  
 542 tle component, translated in moderately high  $^{206}\text{Pb}/^{204}\text{Pb}$  (up to 20.270 and up to 19.989,  
 543 respectively). These isotopic signatures independently support the geophysical evidence  
 544 for the presence of eclogites in the mantle upwellings associated with those volcanic provinces.  
 545 Recent numerical modelling indicates that mantle plumes can entrain about 20% basalt  
 546 component as recycled eclogite, which after leads to the formation of an eclogitic reser-  
 547 voir at depths of around 400 to 300 km (Ballmer et al., 2013; Dannberg & Sobolev, 2015).  
 548 Since an eclogite proportion larger than 40% is required to seismically observe the X dis-  
 549 continuity (Schmerr, 2015; Kemp et al., 2019; Pugh et al., 2021), it has been suggested  
 550 that with time the accumulation of eclogite above the MTZ would increase to the lev-  
 551 els required to detect the *PXs* phase (Pugh et al., 2021).

552 The absence of a consistent *PXs* phase beneath Madeira suggests that, if a basalt  
 553 reservoir is present, it likely contains less than 40% basalt. Conversely, localised detec-  
 554 tions of the *PXs* phase beneath Madeira and the Canaries can be interpreted as indi-  
 555 cations of localised accumulations of basalt.

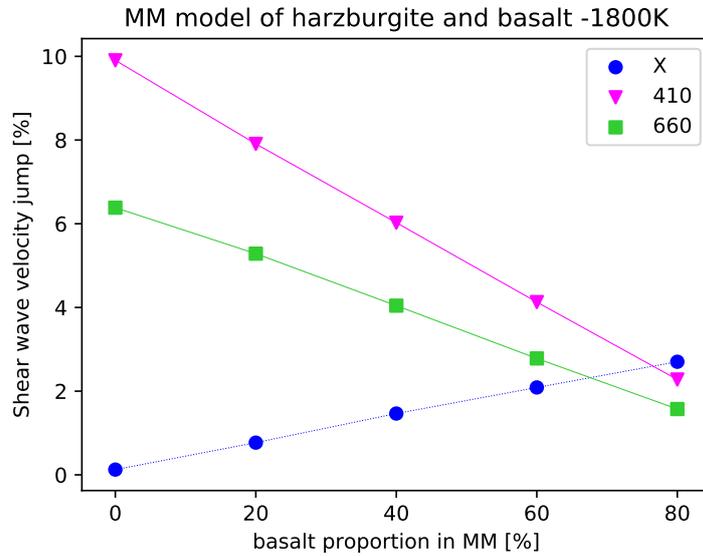
556 Furthermore, it is plausible to hypothesise that the accumulation of the basalt reser-  
 557 voir beneath the Canary Islands is facilitated by the long-lasting and continuous hotspot  
 558 feeding up to the present day. In contrast, if Madeira is in a declining phase, the accu-  
 559 mulation of this reservoir may have been comparatively less efficient and/or may have  
 560 already diminished after the downward retreat of the denser eclogite material, after the  
 561 decrease of the plume positive buoyancy.

### 562 **5.2.2 Compositional constraints from amplitude analysis**

563 A mechanical mixture (MM) of basalt (MORB) and harzburgite, with basalt in-  
 564 creasing with depth, accurately replicates the seismic data for realistic thermal struc-  
 565 tures of the upper mantle (Cammarano et al., 2009). First principle calculations indi-  
 566 cate that the depth variations of the MTZ discontinuities are primarily influenced by tem-  
 567 perature, whereas the magnitude of the corresponding velocity jumps is mainly affected

568 by composition (e.g. Xu et al., 2008). Applying a simplified theoretical analysis to con-  
 569 strain the basalt proportion above the MTZ, we utilize the velocity contrast obtained  
 570 in the amplitude analysis of the detected phases and the velocity contrast in a MM model  
 571 of the mantle.

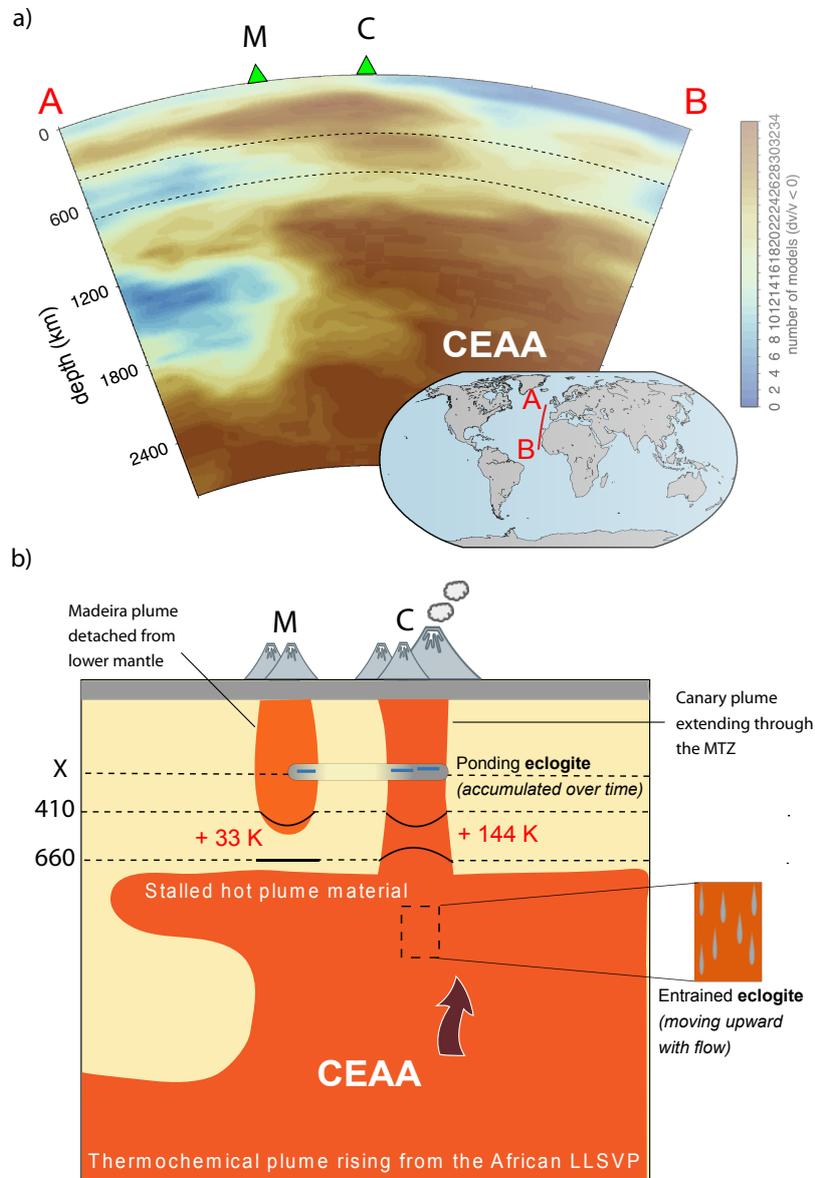
572 Theoretical velocity contrasts for the  $co \rightarrow st$ ,  $ol \rightarrow wa$  and  $ri \rightarrow br + mw$  phases are com-  
 573 puted from the velocity models of Xu et al. (2008) for a potential temperature of 1800  
 574 K and basalt proportions ranging from 0% to 80%. Notably, the 410 discontinuity in these  
 575 models aligns with the estimated depth beneath the Canaries, where a consistent  $PXs$   
 576 phase is detected. The predicted shear wave velocity jumps at depths of 300 km, 410 km,  
 577 and 660 km, corresponding to the mentioned PT, are depicted in Figure 8 and summa-  
 578 rized in the first four columns of Table 2. The analysis reveals that as the proportion  
 579 of basalt increases, the velocity jumps for the 410 and 660 decrease, whereas the trend  
 580 is opposite for the X discontinuity.



**Figure 8.** Predicted shear wave velocity jumps (in %) at depths of 300 km, 410 km, and 660 km, observed in a mechanical mixture (MM) of harzburgite and basalt with varying basalt proportions ( $f_b$ ), obtained using the formula  $V_{Sjump} = (v_2 - v_1)/2(v_1 + v_2)$ , where  $v_1$  and  $v_2$  represent the velocities of the upper and lower layers, respectively.

581 Our estimated velocity jumps (column four in Table 1) are obtained from the am-  
 582 plitude analysis of detected converted phases. These phases can be difficult to detect on  
 583 individual RFs as arrivals are often relatively weak and close to the noise level. Occa-  
 584 sionally, some phases show anomalously large amplitudes, in which case they can be iden-  
 585 tified individually.

586 Generally, these signals are most reliably detected when a large number of RFs from  
 587 small CCP areas are stacked but even then converted phases can be weak, thus making  
 588 it challenging to robustly identify them. For example, for  $P300s$ , we only obtained 18  
 589 robust detections out of 34 good-quality CCP stacks. Moreover, the amplitude of con-  
 590 verted phases in stacked RFs exhibits variability due to factors such as noise, focusing,  
 591 and defocusing by small-scale topography variations. Consequently, constraints on ab-  
 592 solute velocity jumps based on these amplitudes may not be accurate. Therefore, rather  
 593 than interpreting the absolute values of velocity jumps, we focus on the relative values



**Figure 9.** a) Vote cross-section created with the SubMachine tool (<http://www.earth.ox.ac.uk/smachine/>) using 34 P- and S-wave tomographic models and illustrating the most robust long-wavelength low-velocity anomalies in the mantle beneath Canaries (C) and Madeira (M). b) Cartoon illustrating the Canary Plume's extension through the upper mantle and the detached Madeira upwelling, incorporating constraints from our analysis (X discontinuity, eclogite above the MTZ, topography of MTZ discontinuities and temperature of MTZ).

basalt [%] $f_b$	$v_s$ jump [%]			relative $v_s$ jump	
	300	410	660	410/300	410/660
0	0.18	9.94	7.37	-	1.34
20	0.82	7.97	5.96	9.72	1.34
40	1.46	6.04	4.55	4.14	1.33
60	2.08	4.15	3.11	<b>1.99</b>	1.33
80	2.70	2.25	1.65	<b>0.83</b>	1.35

**Table 2.** Predicted shear wave velocity jumps at depths of 300 km, 410 km, and 660 km, observed in a mechanical mixture of harzburgite and basalt with varying proportions of basalt,  $f_b$  and a potential temperature of 1800 K. The last two columns display the relative velocity jumps between 410 and 300, and 410 and 660 as the proportion of basalt increases.

594 of the estimated velocity jumps (last two columns of Table 1). The corresponding the-  
 595 oretical values (relative velocity jumps w.r.t the 410) for varying proportions of basalt  
 596 are displayed in the last two columns of Table 2. This theoretical analysis reveals that  
 597 the relative velocity jump between the 410 and 660 remains relatively constant for vary-  
 598 ing fractions of basalt. However, the relative velocity jump between the 410 and X de-  
 599 creases as the proportion of basalt increases. In other words, the strength of the X dis-  
 600 continuity becomes more prominent for higher values of  $f_b$ , while the opposite trend is  
 601 observed for the 410 and 660 discontinuities.

602 Our estimations for the relative velocity jumps between the 410 and X are 0.96 and  
 603 1.6 (Table 1), for the Canary and Madeira archipelagos, respectively. These values cor-  
 604 respond to a proportion of basalt between 60% and 80% (Table 2). Based on the reli-  
 605 ability of the X discontinuity in the Canaries dataset, we conclude that the accumula-  
 606 tion of basalt is more significant there. For Madeira, the accumulation of basalt appears  
 607 to be very localized.

## 608 6 Conclusions

609 Using P RFs, we conduct a systematic regional study to map out mantle seismic  
 610 discontinuities beneath the Madeira and Canary archipelagos. We identify converted phases  
 611 beneath both archipelagos for transition zone discontinuities, namely the 410 and 660.  
 612 In the Canary Islands, the 410 is deeper ( $428 \pm 4$  km) and the 660 is shallower ( $647 \pm$   
 613  $2$  km). The MTZ is thinner ( $216 \pm 2$  km) compared to the global average suggesting the  
 614 presence of a hot mantle upwelling characterised by excess temperatures that we esti-  
 615 mate of  $+144 \pm 11$  K with respect to the surrounding mantle. Furthermore, the domi-  
 616 nant minerals influencing the topography of the 410 and 660 are, respectively, olivine and  
 617 its polymorph ringwoodite. The discontinuity sharpness analysis provides additional sup-  
 618 port to this conclusion, as it reveals clear and well-defined 660 discontinuity features, which  
 619 is in contrast to a broad 660 that would arise due to garnet being the dominant mineral.  
 620 Beneath Madeira, the 660 discontinuity is close to the reference value and the 410 dis-  
 621 continuity is found at greater depths ( $421 \pm 3$  km). Here the MTZ is cooler than be-  
 622 low the Canaries, with estimated excess temperatures of only  $+33 \pm 11$  K with respect  
 623 to the surrounding mantle. Furthermore, beneath the Canaries, we have reliable detec-  
 624 tions of the X discontinuity, attributed to the coesite-to-stishovite PT in regions of the  
 625 upper mantle enriched in eclogite (as expected for thermochemical plumes). The X dis-  
 626 continuity is in the range of 251-316 km depth in a broad area beneath the Canaries, but  
 627 it is observed at only one isolated location beneath Madeira at 313 km. We estimate that  
 628 the basalt proportion in the plume, between 300-400 km, is 60-80%, with accumulation  
 629 being more significant in the Canaries than in Madeira. The MTZ thickness, the pres-

630 ence of the X discontinuity, and the high basalt fraction provide compelling evidence for  
 631 a deep-seated thermochemical plume beneath the study area, which feeds an active up-  
 632 welling that presently crosses the MTZ beneath the Canaries. Beneath Madeira, where  
 633 only the upper portion of the MTZ seems to be affected and to a minor extent, the up-  
 634 welling could be less vigorous or presently disconnected from its deep source, consistently  
 635 with tomographic images. Our findings are summarized in the conceptual model of Fig-  
 636 ure 9.

## 637 Open Research Section

638 Data from IPMA permanent stations (PM) are available from IPMA at [http://](http://ceida.ipma.pt)  
 639 [ceida.ipma.pt](http://ceida.ipma.pt). Data from DOCTAR experiment (Y7) are available on request from  
 640 the GFZ at <https://dataservices.gfz-potsdam.de/portal/>. Data from the Global  
 641 Seismograph Network (IU) are available from the IRIS Data Management Center (IRIS-  
 642 DMC): <http://service.iris.edu/fdsnws/dataselect/1/>. Data from the Spanish Dig-  
 643 ital Seismic Network (ES) are available on request from the Institut Cartogràfic i Geològic  
 644 de Catalunya (ICGC): <http://ws.icgc.cat/fdsnws/dataselect/1/>.  
 645 DOI information or FDSN network information, where available, are:  
 646 IU: <https://doi.org/10.7914/SN/IU>  
 647 ES: <https://10.7914/SN/ES>  
 648 Y7: [https://www.fdsn.org/networks/detail/Y7\\_2011/](https://www.fdsn.org/networks/detail/Y7_2011/)

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