

Consecutive Ruptures on a Complex Conjugate Fault System During the 2018 Gulf of Alaska Earthquake

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Introduction

This Supporting Information contains numerical tests for validation of the developed finite-fault inversion method (Text S1, Figures S1 to S4, and Table S1). Sensitivity of the

finite-fault inversion to assumptions of model fault depth and rupture velocity is shown in Text S2, and Figures S5 and S6. Comparison with the conventional smoothness constraints is also shown in Text S2 and Figures S7 and S8. The possibility of dummy imaging of reverberations is evaluated in Text S3 and Figure S9. Waveform fits for the main result and full snapshots of the rupture evolution are shown in Figures S10 and S11. Table S2 shows the near-fault velocity structure used for calculating Green's function. Table S3 provides the set of smoothness constraints adopted for the main result.

Text S1.

We perform the numerical tests to evaluate effects of the improved smoothness constraints and the horizontal non-rectangular model fault plane. To generate synthetic waveforms, orthogonal three faults were assumed (Figure S1a). Then, we assume the pure strike-slip rupture which spherically spread from a hypocenter at a depth of 30 km on the central fault, named F2, with a rupture velocity of 3.0 km/s (Figure S1a and b). The moment rate function of input model has peaks at 9 and 23 s (Figure S1c). We add a random Gaussian noise to the calculated Green's function, for which the standard deviation is 5% of maximum amplitude of each calculated Green's function. We also add a random Gaussian noise with zero mean and a standard deviation of $1.0 \mu m$ as background noise. We generate the synthetic waveforms at 78 stations used in the inversion of the 2018 Alaska earthquake (Figure 2c).

We compare results of four cases: (1) the rectangular model fault plane and the conventional smoothness constraints; (2) the rectangular model fault plane and the improved smoothness constraints; (3) the non-rectangular model fault plane and the conventional smoothness constraints, and (4) the non-rectangular model fault plane and the improved smoothness constraints.

We set the horizontal rectangular model fault plane with a width and length of 120 km to cover the input three faults (Figure S2a). The depth of the model fault plane is set to 30 km, which corresponds to the centroid depth of input source model. The spatial knot interval is set to 10 km. For the cases (3) and (4), we design the non-rectangular model fault plane based on the input three faults (Figure S2d). For all cases, the potency-rate density function at each knot is represented as a linear combination of B-spline functions over a duration of 30 s with an interval of 0.8 s and the rupture front velocity set at 7.0 km/s. We adopt the improved smoothness constraints at the cases (2) and (4) by referring to the input focal mechanism (Table S1).

In the case (1), the resultant moment rate function is smoother than the input one and has only one peak at 12 s, which is about 3 s later than the first peak of the input (Figure S2b). The normalized L2 norm, which represents the degree of misfit between the input and the resultant moment rate function (hereinafter called "the L2 norm" for simplicity), was 0.245. The snapshots show a wider potency-rate density distribution than the input, making it difficult to identify the fault geometry and interpret the source process (Figure S3a and b). Figure S4 shows the self-normalized potency-rate function for each basis component, obtained by taking a spatial integration of the potency-rate density function for each basis component. In the case (1), the potency-rate function of $M1$ component (Kikuchi & Kanamori, 1991), corresponding to the input slip

direction, is smoother than those of other components (Figure S4). This is because the conventional smoothness constraints work to excessively smooth out the dominant basis components.

In the case (2), the moment rate function yields two peaks at 10 and 23 s, which close to the input peaks (Figure S2c). The L2 norm is 0.074. The improved smoothness constraints remove the bias in the resultant potency-rate function of the *M1* component (Figure S4) and thus the spatiotemporal potency-rate density distribution of the case (2) is finer than that of the case (1) (Figure S3c). However, in the case (2), the image looks too blurry to resolve two independent ruptures of the input model due to insufficient spatial resolution (snapshot at 15 and 20 s in Figure S3c).

In the case (3), the moment rate function has two peaks at 11 and 23 s, which close to the input peaks (Figure S2e). The L2 norm is 0.135 and slightly larger than the case (2). The potency-rate density distribution of the case (3) at 15 and 20 s resolves two ruptures, which are not well resolved in the case (2) (15 and 20 s in Figure S3d). The spatial resolution of the inversion results is improved because the model space modification according to the input fault geometry is identical to implicitly introducing a *priori* constraint of the fault geometry (e.g., aftershock distribution). The model space reduction also contributes to reduce computational costs, which is useful for analyses of earthquakes having a vast source area, such as the 2018 Alaska earthquake.

In the case (4), the moment rate function reproduces the input in detail (Figure S2f), and the rupture evolution is fine enough to reproduce the input (Figure S3e). The L2 norm of the case (4) is 0.071, which is the minimum value among the four cases. Thus, we conclude from our numerical tests that the optimum strategy should be by using both the improved smoothness constraints and the horizontal non-rectangular model fault plane.

Text S2.

We evaluate the sensitivity of the inversion results by perturbing the model parameters. We perform the inversion analyses by changing the model fault plane depth to 33.6 ± 5 km. The obtained snapshots show the rupture pattern is insensitive to the model fault depths (Figure S5). We also check the inversion results by changing the assumption of maximum rupture velocity to 3 and 5 km/s. We resolve the similar rupture processes for the maximum rupture velocities at 5 and 7 km/s (Figure S6b and c). However, when assuming 3 km/s, the model does not clearly show the A2 rupture (Figure S6a). This is due to the limited model space that could artificially vanish the possible slip behavior beyond the designated rupture front.

We also perform the inversion analysis with the conventional smoothness constraints to evaluate the effect of the improved smoothness constraints for the real earthquake (the 2018 Gulf of Alaska earthquake). The inversion results with the conventional smoothness constraints show almost the similar source process of the results obtained by the improved smoothness constraints (Figure S7). However, the spatiotemporal rupture propagation of the conventional smoothness constraints is smoother than that of the improved ones by the excessive smoothing for the most

dominant $M1$ component for the earthquake (Figure S8), which provides the blurrier image, making it difficult to clearly resolve the multiple sub-events (Figure S7).

Text S3.

As shown in Figures 2 and S10, our finite-fault model sufficiently reproduces the complicated observed teleseismic P waves, resulting in showing the complex-multiple rupture episodes. On the other hand, pulses of the observed waveforms may include later arrivals due to structure complexities in the source region (e.g., Fan & Shearer, 2018; Yue et al., 2017). However, the theoretical Green's functions, assuming a 1D-layered structure model, are often poorly modeled for reverberations of dipping near-source bathymetry (Wiens 1987, 1989), which may induce artificial imaging of multiple-shock sequence. In principal, seismograms of relatively smaller earthquakes with a similar focal mechanism that occurred near the target earthquake can be regarded as an empirical Green's function (EGF) under an assumption that the moment rate function of that small earthquake is simple and short (Hartzell 1978; Dreger 1994). We here employ the EGFs instead of the theoretical Green's functions to evaluate whether multiple-shock sequence that we resolve is likely from the source effect or the reverberations. We deconvolve the EGFs from the observed waveforms of the 2018 Alaska earthquake for each station to remove the effects of the earth response including possible reverberations.

We select three events from the GCMT catalog (Dziewonski et al., 1981; Ekström et al., 2012) as the EGFs with the clear first P -phase motion and high signal-to-noise ratios (Figure S9a). The EGFs and the mainshock data are band-passed between 0.01 and 2 Hz and converted into ground velocities with a sampling interval of 0.1 s. We solve the least squares problem using the non-negative least squares algorithm of Lawson and Hanson (1974). We perform deconvolution for both a maximum source duration of 65 and 27 s to evaluate the validity of the sub-events resolved after 27 s for the mainshock (Figure S9b, c, d, and e).

The normalized moment rate functions obtained in the maximum length of 65 s show non-negligible moment release even after 27 s (Figure S9b, c, d and e). If the subevents after 27 s were artifacts caused by the reverberations of the initial rupture, the observed waveforms would be reproducible by convolving the moment rate function up to 27 s with the EGF. However, the synthesized waveforms obtained from the 27-s-moment-rate function fails to reproduce the several pulses of the observed waveforms, while the synthesized waveforms obtained from the 65-s-moment-rate function better fits the observed waveforms, suggesting that the subevents after 27 s should be necessary to explain the observed data (Figure S9b, c, d and e).

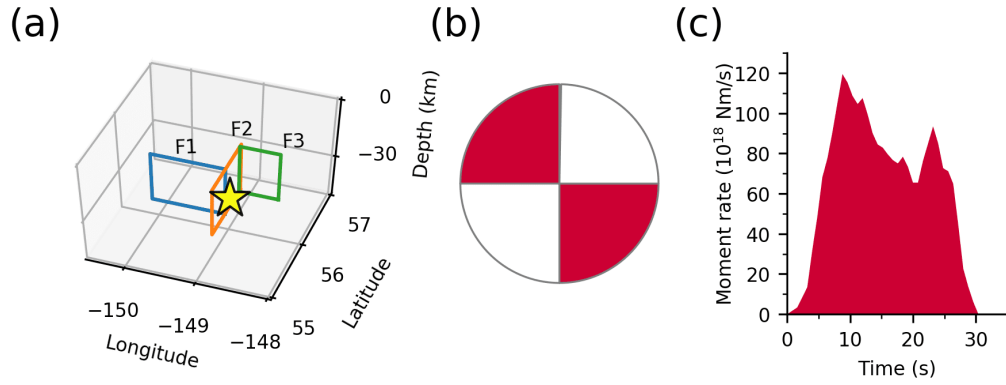


Figure S1. Input source model for the numerical tests. (a) Input fault geometry. The star indicates the location of the initial break. All three faults have dips of 90° and down-dip widths of 25 km. F1 (Blue) has a strike of 90° and a length of 65 km. F2 (Orange) has a strike of 0° and a length of 100 km. F3 (Green) has a strike of 90° and a length of 35 km. (b) Total focal mechanism of the input slip-rate. (c) Input moment rate function.

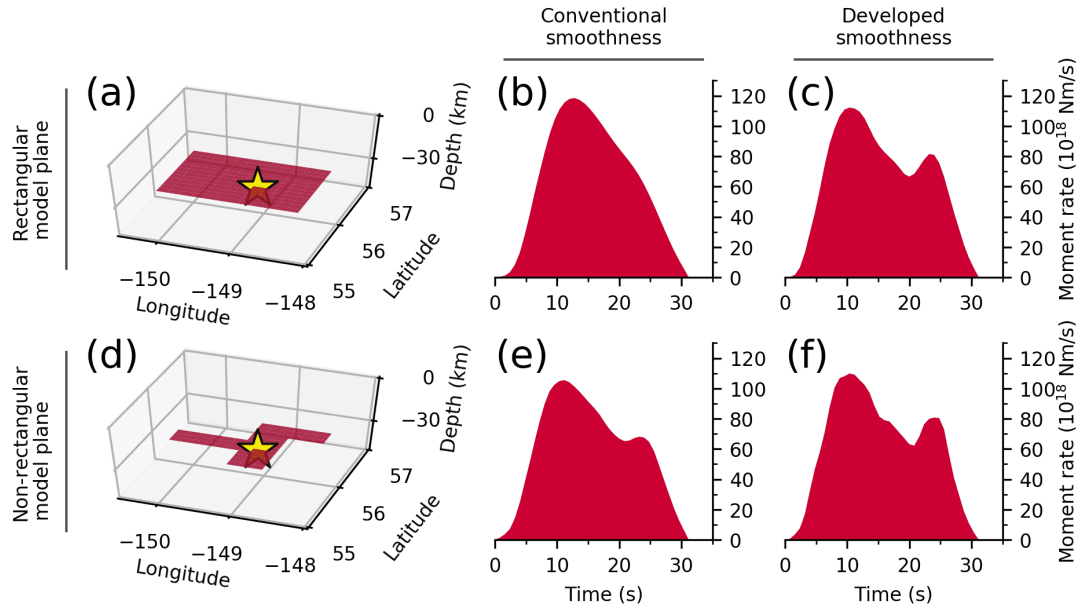
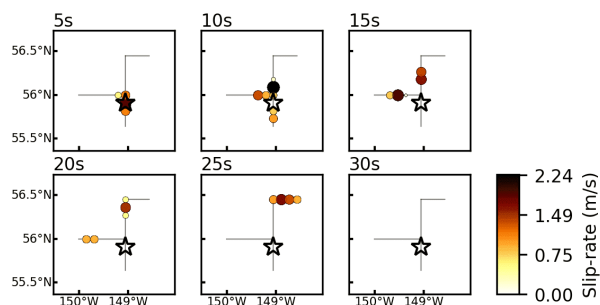


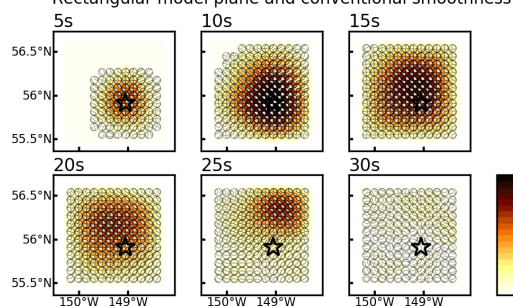
Figure S2. Assumed model fault planes and resultant moment rate functions for the numerical tests. (a) and (d) represent the rectangular model fault plane and the non-rectangular model fault plane, respectively. (b), (c), (e) and (f) show the moment rate functions obtained by the rectangular model fault plane and the conventional smoothness constraints (case 1), the rectangular model fault plane and the improved smoothness constraints (case 2), the non-rectangular model fault plane and the conventional smoothness constraints (case 3), and the non-rectangular model fault plane and the improved smoothness constraints (case 4), respectively.

(a) Input



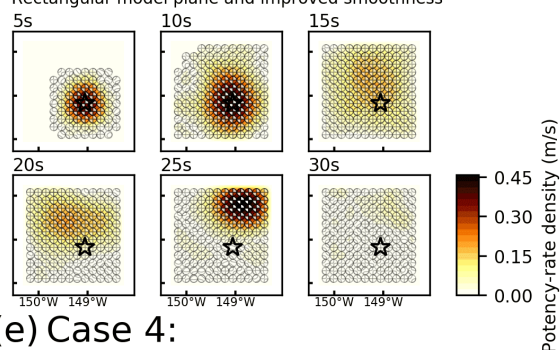
(b) Case 1:

Rectangular model plane and conventional smoothness



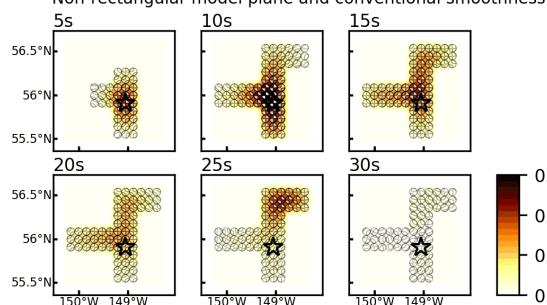
(c) Case 2:

Rectangular model plane and improved smoothness



(d) Case 3:

Non-rectangular model plane and conventional smoothness



(e) Case 4:

Non-rectangular model plane and improved smoothness

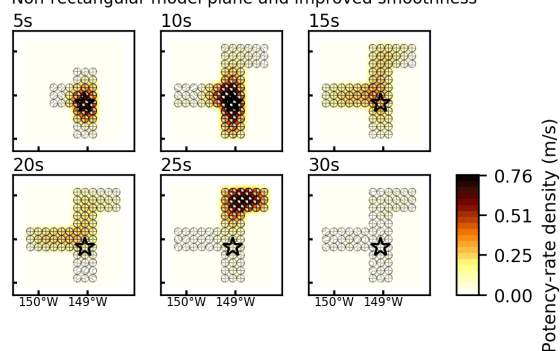


Figure S3. Snapshots of (a) input slip-rate and (b) to (e) resultant potency-rate density tensors for each numerical-test case every 5 s. The star denotes the initial breaking point. The dots in panel (a) denote the input source positions. Color of these dots represents the value of slip-rate. The gray line in panel (a) represents the input fault geometry.

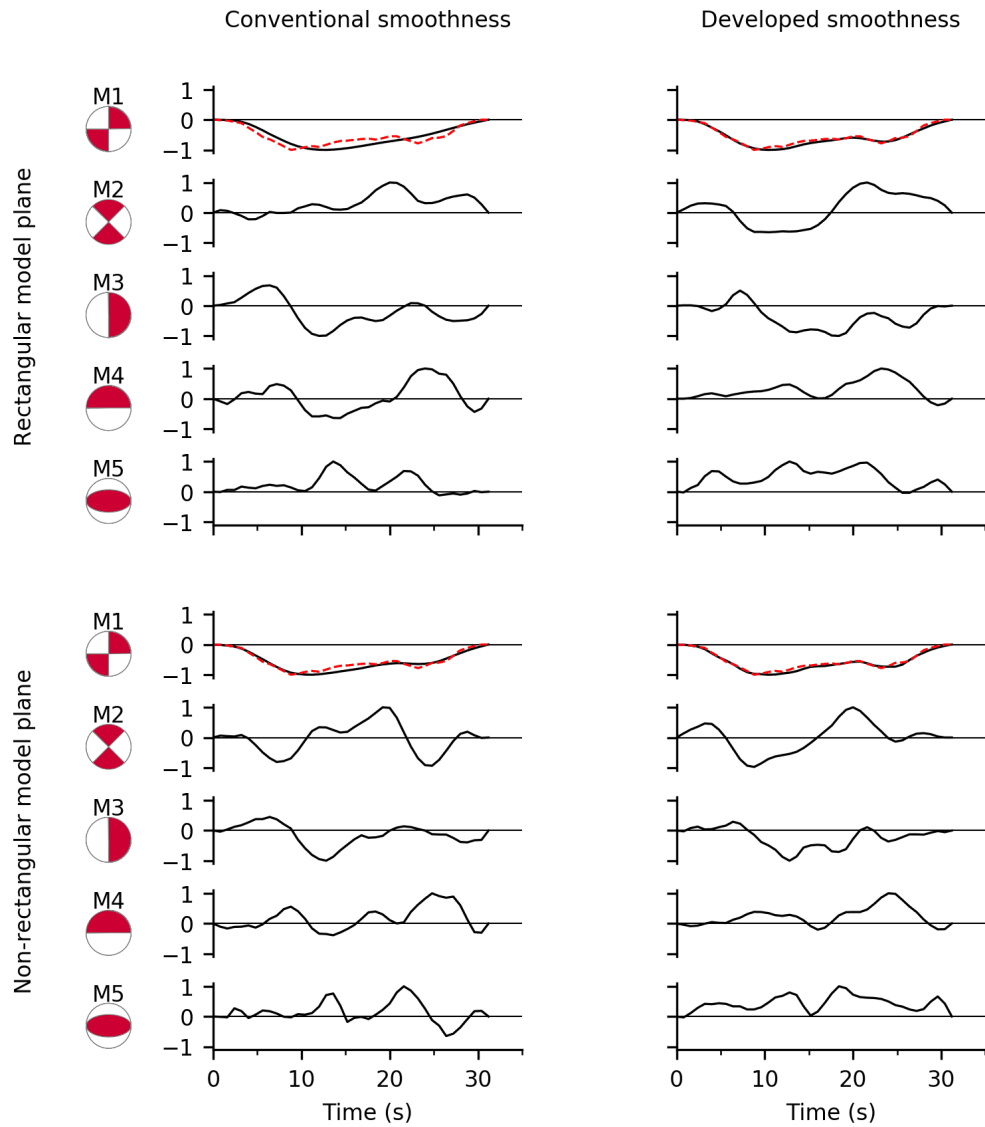


Figure S4. Comparison of the input slip-rate function (dashed line) and the potency-rate functions for each basis component, obtained by taking a spatial integration of the potency-rate density function (Kikuchi & Kanamori, 1991) (solid line). Each trace is self-normalized.

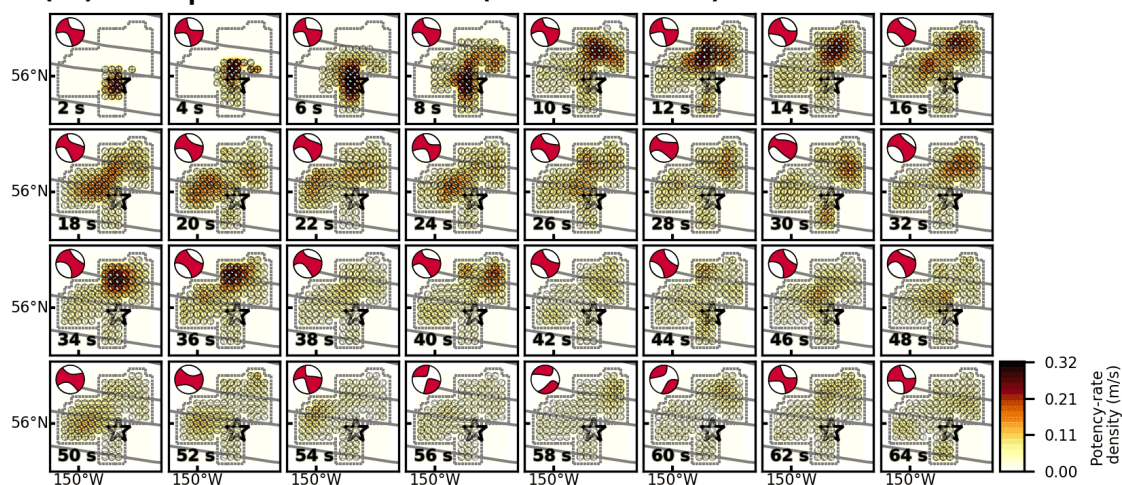
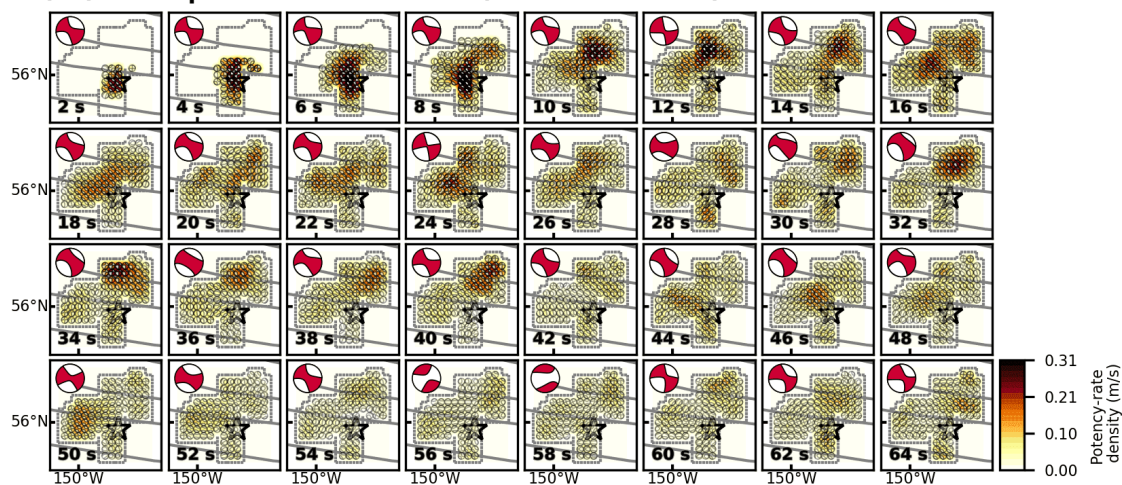
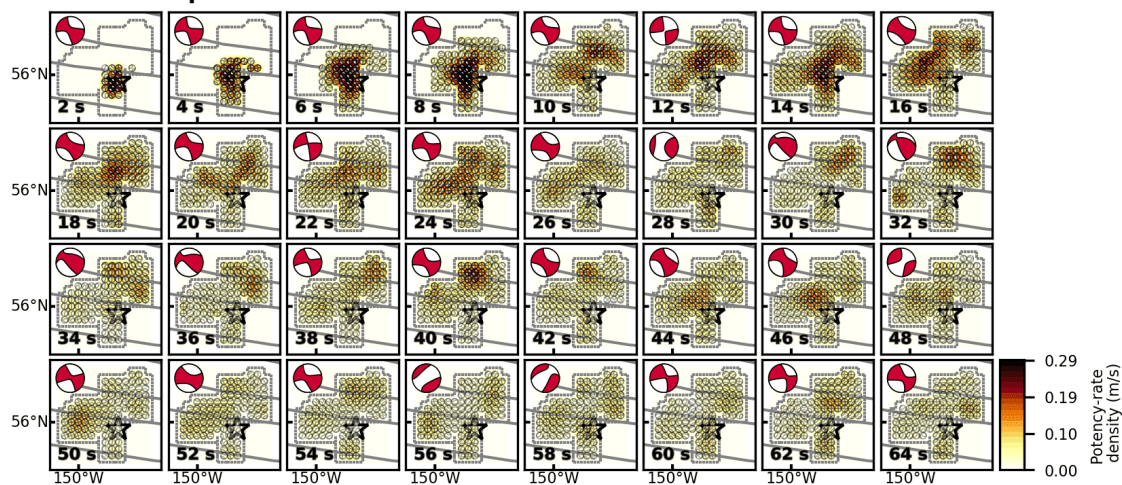
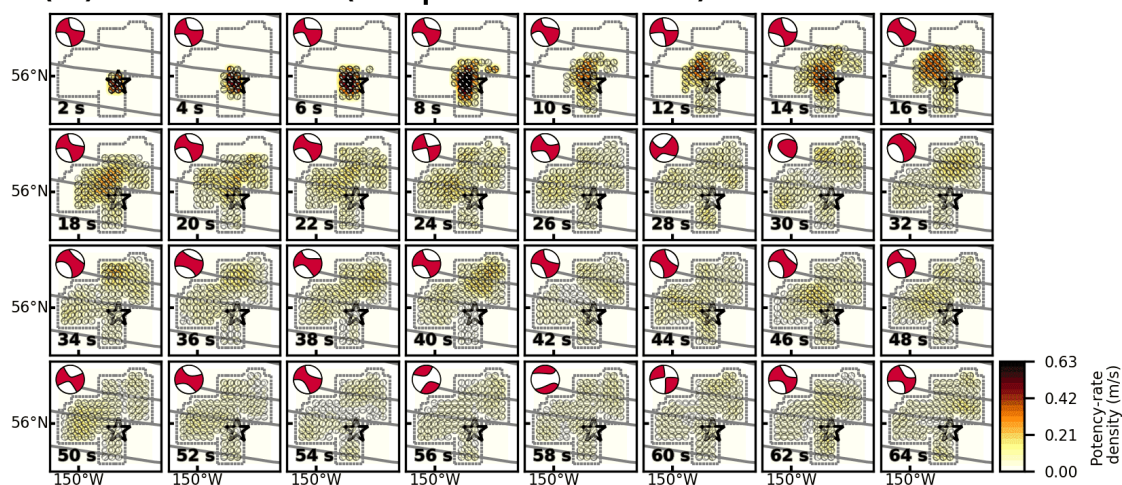
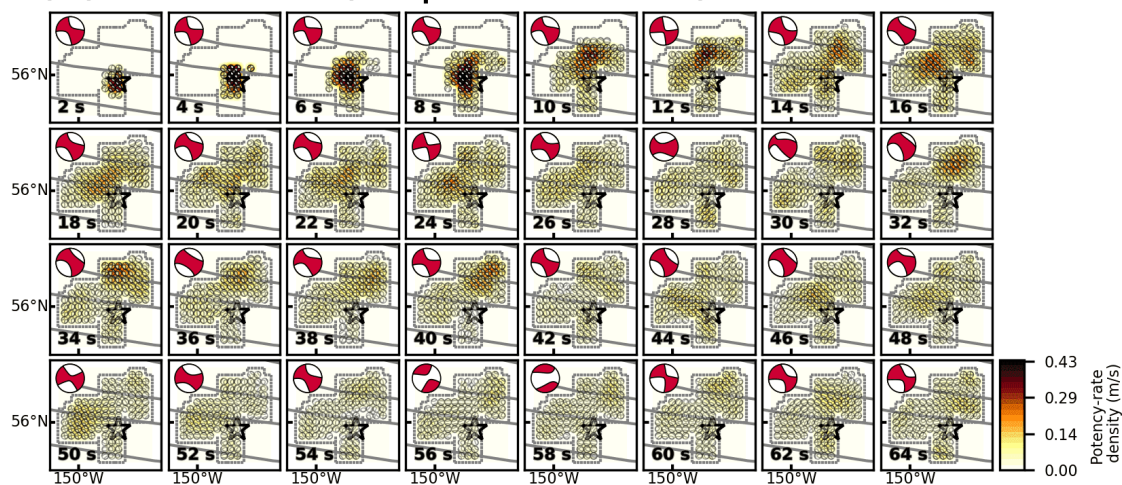
(a) Depth 28.6 km (V_r 7 km/s)(b) Depth 33.6 km (V_r 7 km/s)(c) Depth 38.6 km (V_r 7 km/s)

Figure S5. Summary of snapshots of the potency-rate density tensors for the different assumptions of model fault plane depth for the 2018 Alaska earthquake. The depth of each snapshot is (a) 28.6 km, (b) 33.6 km, and (c) 38.6 km. The rupture front velocity (V_r) is 7 km/s for all the snapshots. The corresponding time after onset for each snapshot is noted at the bottom-left of each panel. The dotted line shows the border of the assumed model fault plane. The star and solid lines indicate the epicenter (AEC) and the fracture zones (Matthews et al, 2011; Wessel et al., 2015), respectively. The large beachball in each panel indicates the corresponding total moment tensor at each time.

(a) Vr 3km/s (Depth 33.6 km)



(b) Vr 5km/s (Depth 33.6 km)



(c) Vr 7km/s (Depth 33.6 km)

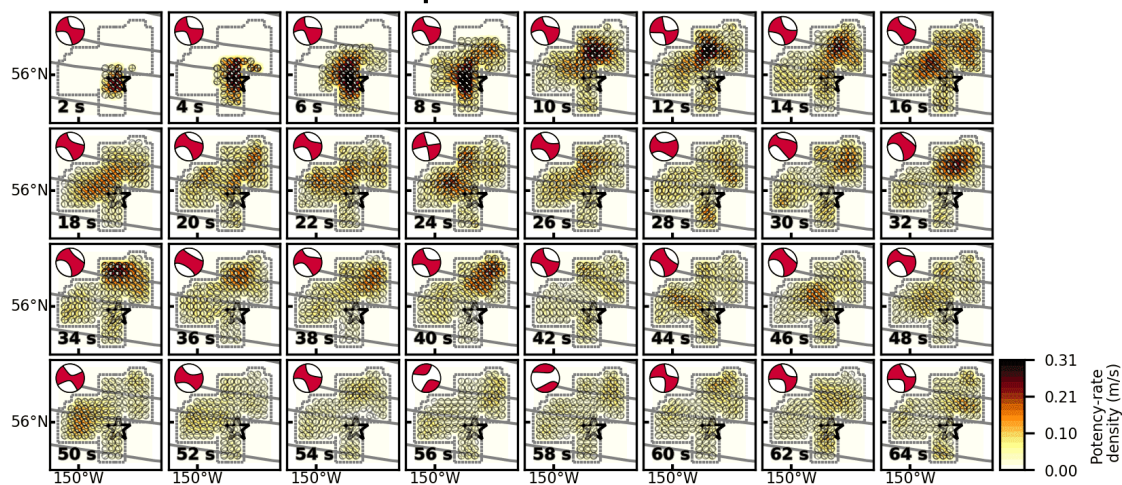


Figure S6. Summary of snapshots of the potency-rate density tensors for the different assumptions of rupture front velocity (V_r) for the 2018 Alaska earthquake. The rupture front velocity of each snapshot is (a) 3 km/s, (b) 5 km/s, and (c) 7 km/s. The model fault plane depth is 33.6 km for all the snapshots.

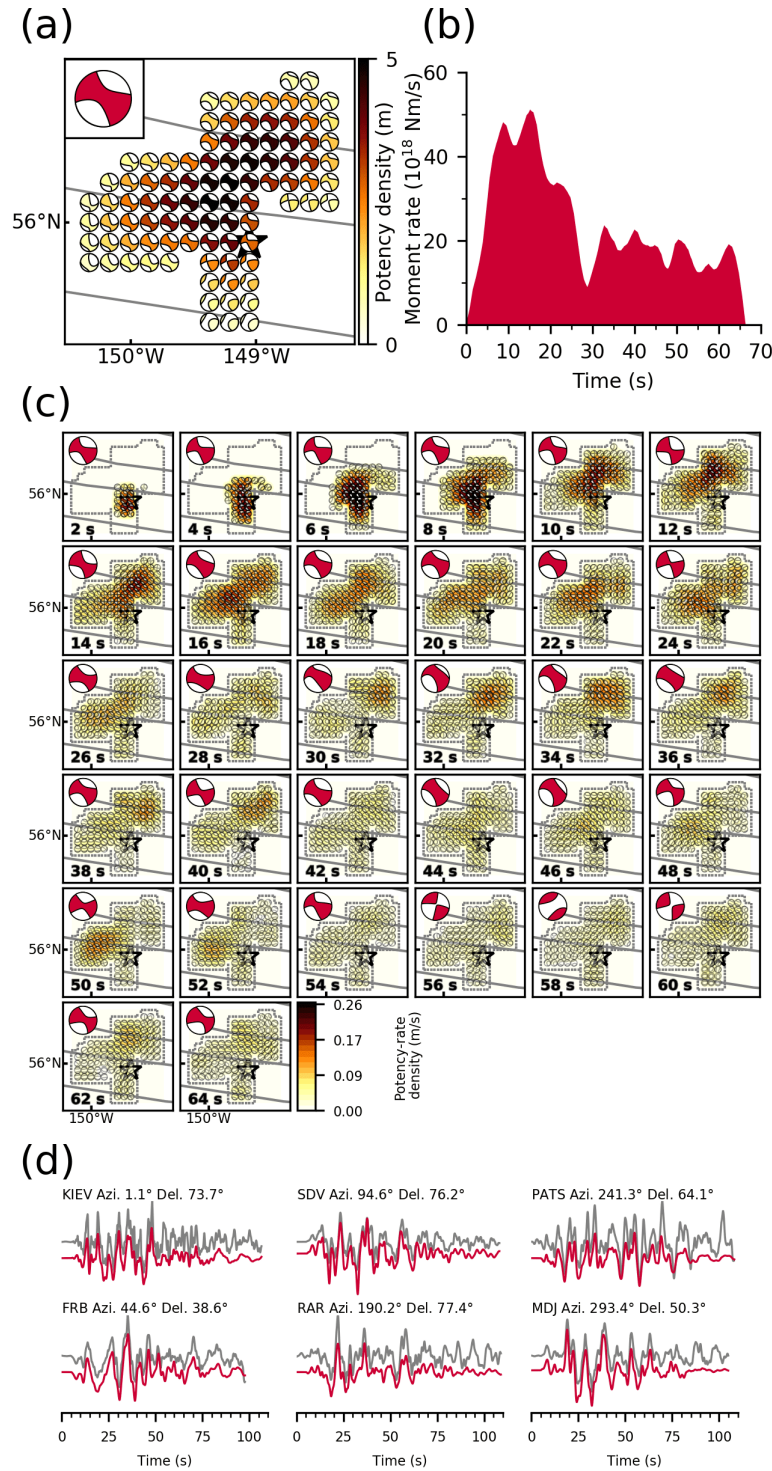


Figure S7. Summary of result obtained by the conventional smoothness constraints for the 2018 Alaska earthquake. (a) Map projection of the potency density tensor distribution on the assumed model fault plane. The star and solid lines indicate the epicenter (AEC) and the fracture zones (Matthews et al, 2011; Wessel et al., 2015), respectively. Inset is the total moment tensor. (b) The moment rate function. (c)

203 Snapshots of the potency-rate density tensors every 2 s. The dotted line shows the
204 border of the assumed model fault plane. The large beachball in each panel indicates the
205 corresponding total moment tensor at each time. (d) Comparison of observed
206 waveforms (gray) with synthetic waveforms (red) at the selected stations. The station
207 location is shown in Figure 2c.
208

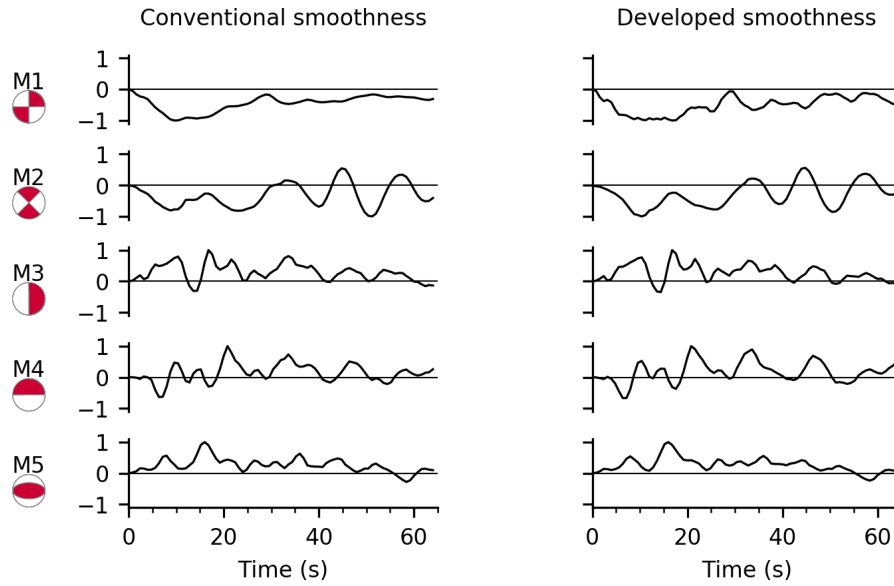
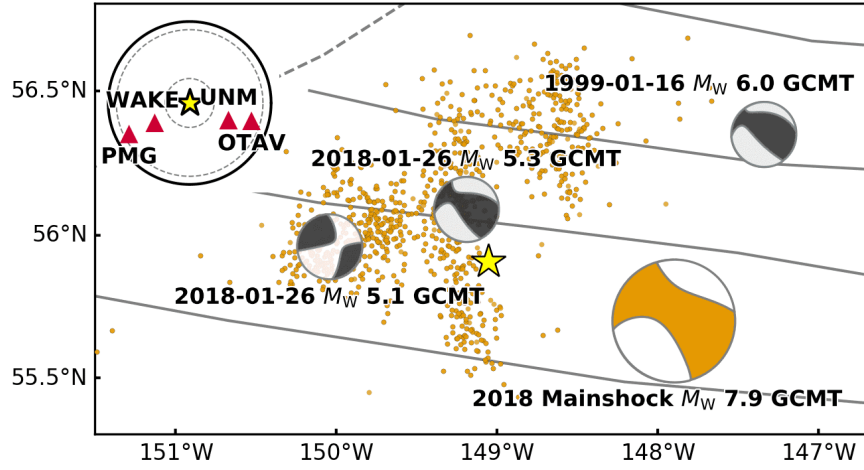
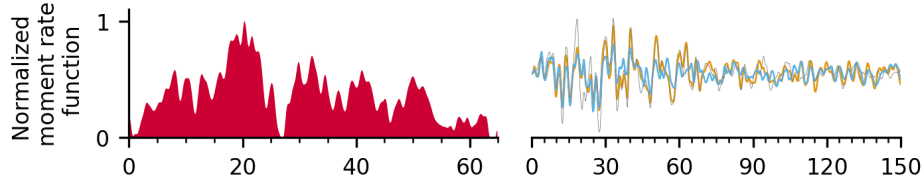


Figure S8. Comparison of the potency-rate functions of the 2018 Alaska earthquake obtained by the conventional (left column) and improved smoothness constraints (right column) for each basis double-couple component (Kikuchi & Kanamori, 1991). Each trace is self-normalized.

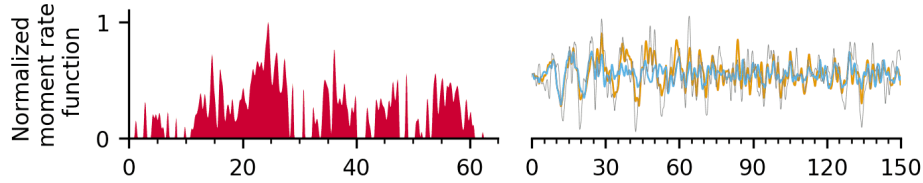
(a)



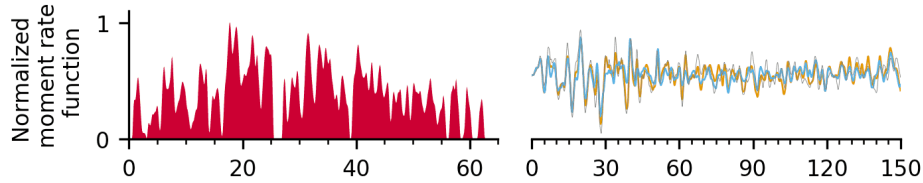
(b) OTAV Azi. 106° Del. 79°/EGF; $M_{5.3}$



(c) WAKE Azi. 240° Del. 49°/EGF; $M_{5.0}$



(d) UNM Azi. 113° Del. 52°/EGF; 1999M6.0



(e) PMG Azi. 242° Del. 83°/EGF; 1999M6.0

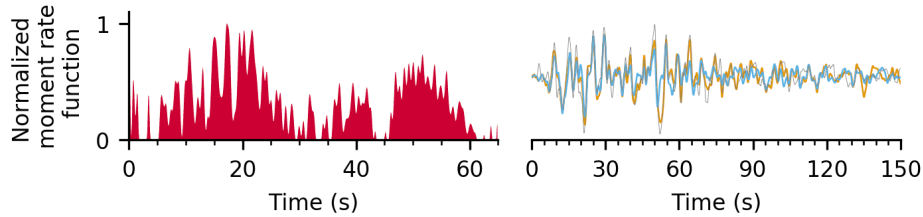


Figure S9. Summary of the EGF analysis. (a) Map projection of the GCMT solutions of the main shock (orange beachball) and events used as the EGFs (black beachballs). The star is the mainshock epicenter, and orange dots are aftershocks ($M \geq 3$) that occurred within one week of the mainshock; all epicentral locations are from AEC. Dashed and solid lines represent the plate boundaries (Bird, 2003) and the fracture zones (Matthews et al, 2011; Wessel et al., 2015), respectively. Inset is an azimuthal equidistant projection of the station distribution. (b) to (e) show the normalized moment rate function (left) and waveform fittings (right). Gray trace is the observed waveform. Also shown is the synthetic waveforms obtained by using 65-s-moment-rate function (orange) and 27-s-moment-rate function (blue). Each panel is labeled with the station name, azimuth (Azi.) and epicentral distance (Del.) from the mainshock, and the event name used as the EGF.

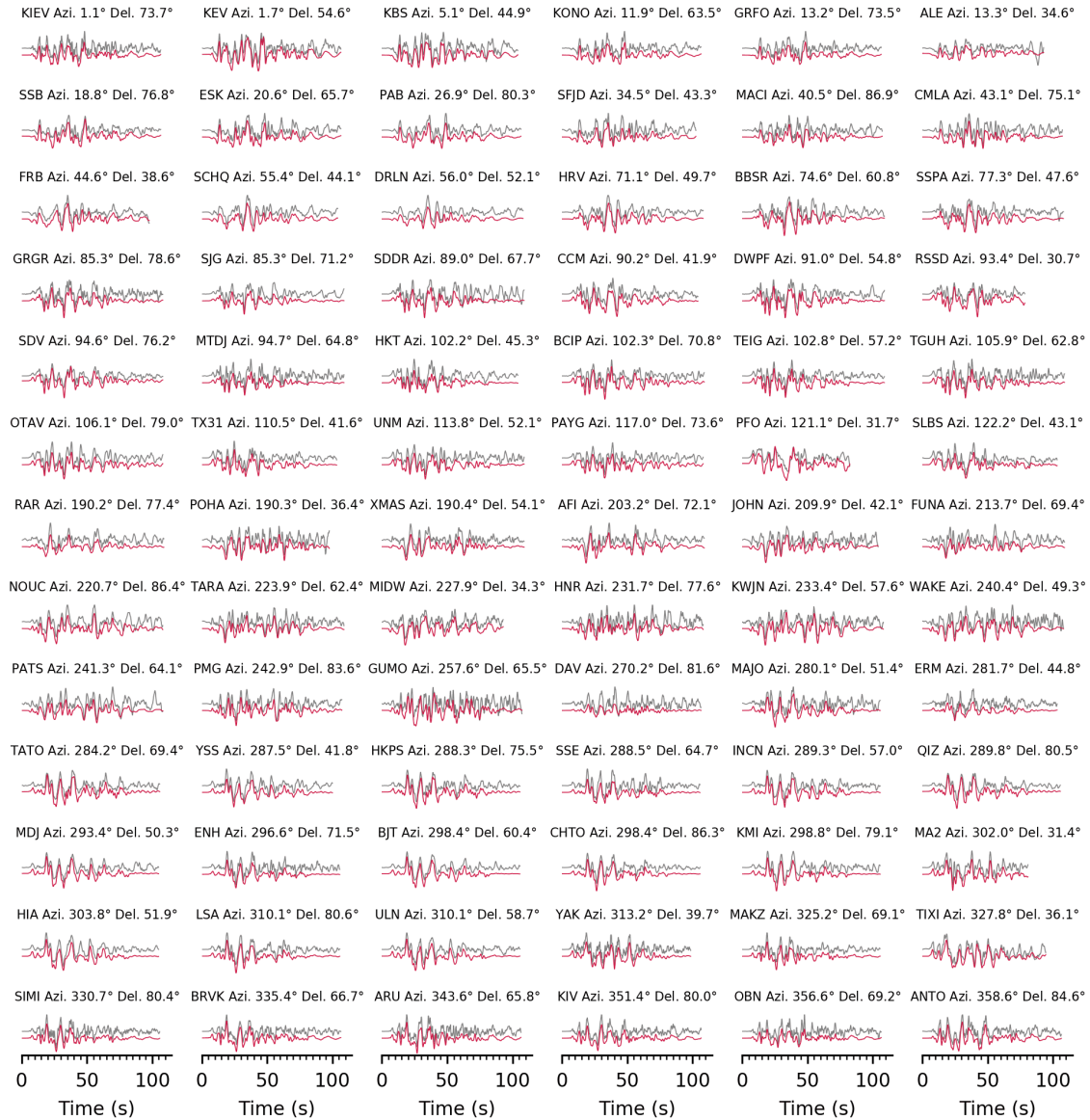


Figure S10. Comparison of observed waveforms (gray) with synthetic waveforms (red) for the main inversion results of the 2018 Alaska earthquake. Each panel is labeled with the station name, azimuth (Azi.), and epicentral distance (Del.) from the mainshock.

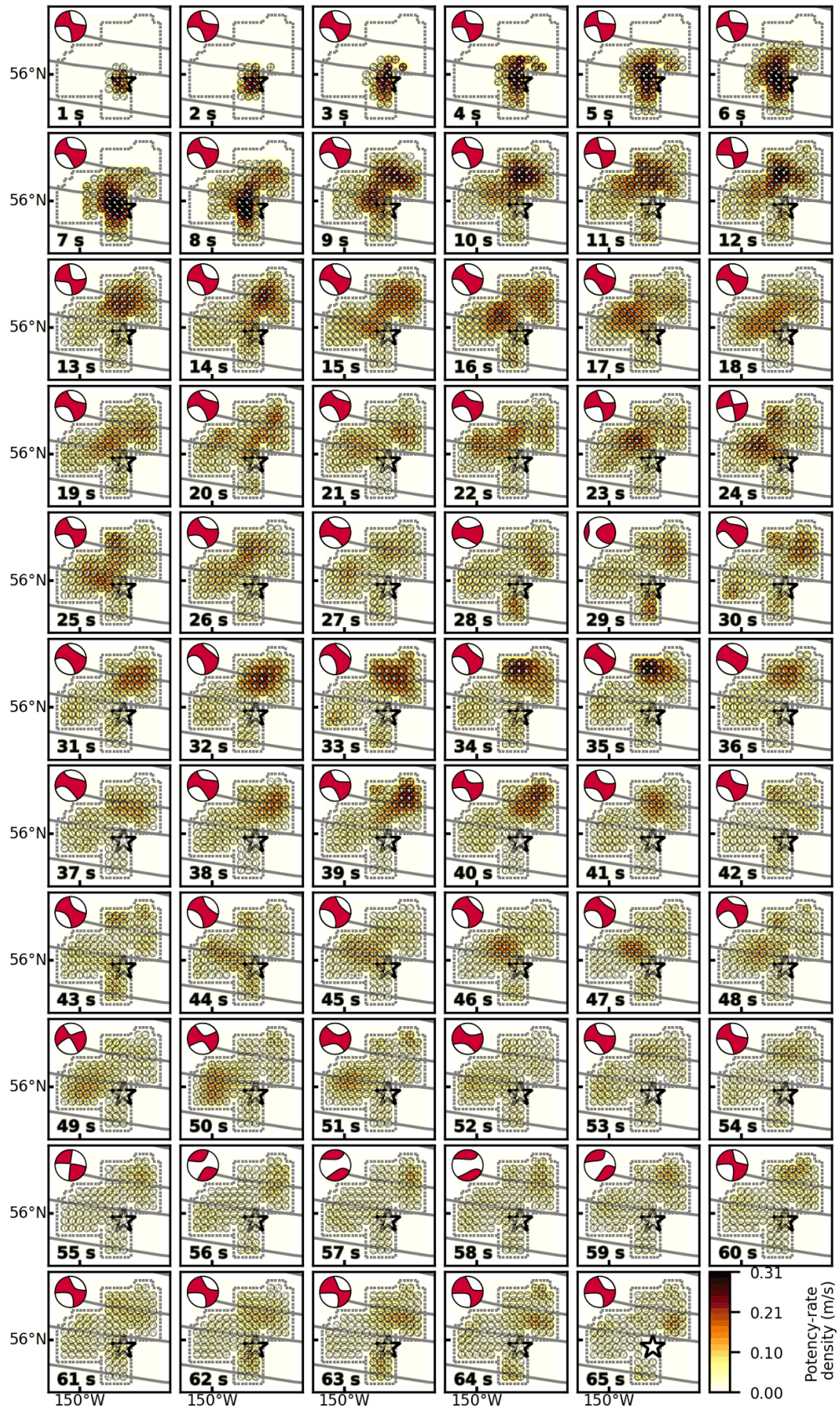


Figure S11. Snapshots of the potency-rate density tensors every 1 s for the 2018 Alaska earthquake. The dotted line shows the border of the assumed model fault plane. The star and solid lines indicate the epicenter (AEC) and the fracture zones (Matthews et al, 2011; Wessel et al., 2015), respectively. The large beachball in each panel indicates the corresponding total moment tensor at each time.

Table S1. Factors of the smoothness constraint of each potency component for the numerical tests. The number q represents **M1** to **M5** components defined by Kikuchi and Kanamori (1991). $|m_q|$ is the absolute value of the total potency derived from the input total moment tensor (Figure S1b). The scaling factor k was set so that $\min(k|m_q|) = 1$.

q	1	2	3	4	5
$ m_q $	1.0005	0.0000	0.0000	0.0000	0.0000
$k m_q $	10.0000	1.0000	1.0000	1.0000	1.0000

Table S2. CRUST1.0 structural velocity model (Laske et. al., 2013).

V_P (km/s)	V_S (km/s)	Density (10^3 kg/m ³)	Thickness (km)
1.50	0.00	1.02	4.30
1.85	0.41	1.87	0.39
5.00	2.70	2.55	0.66
6.50	3.70	2.85	1.47
7.10	4.05	3.05	4.53
8.08	4.49	3.33	0.00

Table S3. Factors of the smoothness constraint of each potency component for the analysis of the 2018 Alaska earthquake. $|m_q|$ is the absolute value of the total potency of each potency component derived from the GCMT solution (Figure 1). The scaling factor k was set so that $\min(k|m_q|) = 1$.

q	1	2	3	4	5
$ m_q $	0.7900	0.2500	0.3600	0.1900	0.2400
$k m_q $	4.1579	1.3158	1.8947	1.0000	1.2632

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