

# Strong super-rotation during the 2018 Martian Global Dust Storm

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## Key Points:

- Super-rotation in the martian atmosphere was greatly enhanced during the Mars Year 34 Global Dust Storm
- Strengthened tropical winds were instrumental in transporting dust across the globe during the initial stages of the Global Dust Storm
- Tracking of dust plumes using data assimilation yielded estimates of wind speeds in the tropics of 10-15 m/s below 10 km

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## Abstract

We report on the strong enhancement of tropical winds in the martian atmosphere during the Global Dust Storm of Martian Year 34, as revealed by an assimilation of temperature and dust observations into a martian global climate model. We demonstrate that global super-rotation almost doubled at the peak of the dust storm period, as compared to Mars Year 33, which did not have a Global Dust Storm. The enhanced westerly jet played a key role in the initial stages of the dust storm, transporting dust across the planet and between different lifting centers. By tracking the progression of dust plumes, we derive estimates of wind speeds in the lower tropical atmosphere of 10-15 m/s below 10 km. The tropical region was dominated by diurnal and semi-diurnal tides, with a significant amplification of the gravest Kelvin mode during the onset of the storm.

## Plain Language Summary

Dust plays a major role in driving the behaviour of the atmosphere of Mars. During a Global Dust Storm, winds lift and transport dust throughout the atmosphere; in turn, dust affects wind direction and strength by heating and cooling the surrounding air. Using a technique that allowed us to combine satellite observations with simulations of the martian atmosphere, we demonstrate that winds at tropical latitudes were greatly strengthened during the 2018 martian Global Dust Storm as a result of heating by dust. We show that these winds played a crucial role in transporting dust across the globe during the initial stages of the storm. By tracking the locations of the observed dust plumes, we were able to estimate wind speeds in the lower atmosphere of Mars.

## 1 Introduction

A planetary atmosphere is said to be in a state of *super-rotation* if the total axial angular momentum of the atmosphere is greater than the angular momentum of an atmosphere in a state of pure solid body rotation (Read, 1986). Super-rotation usually takes the form of a prograde equatorial zonal jet, and it has been shown that such a state must be maintained by the equatorward transport of angular momentum by non-axisymmetric eddy motions (Hide, 1969). Identifying the exact mechanisms that initiate and maintain super-rotating jets remains a significant challenge in planetary atmospheric research, and is crucial for understanding the atmospheric dynamics of slow-rotating terrestrial planets such as Venus and Titan (Read & Lebonnois, 2018). The martian atmosphere is also a candidate for super-rotation, as first identified by Lewis and Read (2003).

Global Dust Storms (GDS) are planetary-scale dust lifting events in the martian atmosphere that occur every few martian years, and are the most dramatic of all the dust-driven atmospheric phenomena on the planet. During a GDS, large parts of the planet are enveloped in a shroud of dust that can persist for several months at a time, drastically affecting the atmospheric state. Such storms have all been observed to occur during the high dust loading season ( $L_S=180^\circ-360^\circ$ ), when the planet's orbit brings it closer to the Sun. The most recent GDS occurred in Mars Year (MY) 34 (using the naming convention of Clancy et al. (2000)). The storm initiated shortly after the northern hemisphere autumn equinox at  $L_S=187^\circ$ , and its evolution was monitored by multiple orbital, surface and Earth-based instruments (Guzewich et al., 2019; Kass et al., 2019; Smith, 2019; Sánchez-Lavega et al., 2019; Felici et al., 2019; Jain et al., 2020; Aoki et al., 2019; Stcherbinine et al., 2020).

The formation mechanisms of a GDS are still not well understood; in particular, it has not been established why a GDS forms out of regional storms in some years, but not others. Posited hypotheses include the re-distribution of dust between finite reservoirs between years (Mulholland et al., 2013; Newman & Richardson, 2015), and a weak

64 coupling between orbital and rotational angular momentum on Mars (Shirley, 2017; New-  
65 man et al., 2019).

66 Lewis and Read (2003) have shown that dust-driven heating can excite a super-  
67 rotating jet in the martian atmosphere. In dusty conditions (such as during a GDS), the  
68 enhanced absorption of solar and thermal radiation amplifies the thermal tides, which  
69 are normal modes of the atmospheric response to diurnal solar forcing. The westward  
70 propagation of these tides away from their tropical source regions then enables an east-  
71 ward super-rotating jet to form.

72 One might expect that a strengthening of the equatorial jet by the presence of dust  
73 would also affect dust transport, at least in the tropical regions where the jet is strongest.  
74 Transport of dust by the jet determines the locations of further dust-driven heating; in  
75 this way, there is an intimate link between dust-driven heating and the strength of the  
76 equatorial jet which may have implications for GDS evolution.

77 In this paper, we investigate dust-driven super-rotation of the martian atmosphere  
78 during the MY 34 GDS, using a data assimilation scheme (Lewis et al., 2007) to assim-  
79 ilate temperature and dust data from the Mars Climate Sounder (MCS) and Atmospheric  
80 Chemistry Suite (ACS) instruments into a Mars Global Climate Model (MGCM). We  
81 show that the atmosphere was strongly super-rotating during the dust storm period, driven  
82 by excited thermal tides. The enhanced equatorial jet is shown to play a significant role  
83 in transporting dust eastwards during the initial phases of the storm, and we estimate  
84 the strength of zonal winds by tracking the eastward migration of dust plumes.

85 In Section 2, we describe the model and assimilation scheme that we used for this  
86 study, and we briefly describe the observational datasets utilised in the assimilation. We  
87 also describe some super-rotation metrics that will be used later in the study. Results  
88 are presented in Section 3, and we discuss our findings in Section 4.

## 89 2 Data and Methods

### 90 2.1 Observation data

91 MCS is a passive 9-channel limb-scanning radiometer aboard the Mars Reconnaissance  
92 Orbiter (MRO) (McCleese et al., 2007). Measurements are taken at approximately  
93 3 am and 3 pm local time along a Sun-synchronous orbit, across a wavelength range of  
94 0.3–45  $\mu\text{m}$ . Radiance profiles were generated from the surface to around 80 km with an  
95 intrinsic vertical resolution of 5 km, enabling the retrieval of profiles of temperature, dust  
96 and water ice (Kleinböhl et al., 2009, 2017). For this assimilation study we made use of  
97 the temperature observations and derived column dust values from MCS (Montabone  
98 et al., 2015), which were available for the duration of MYs 33 and 34.

99 ACS (Korablev et al., 2018) is an array of three infrared spectrometers aboard the  
100 ExoMars Trace Gas Orbiter (TGO), that together provide spectral coverage over a wave-  
101 length range of 0.7–17  $\mu\text{m}$ . The near-infrared (NIR) channel uses an echelle grating with  
102 an Acousto-Optical Tunable filter with a resolving power of 25,000, to retrieve atmospheric  
103 density and temperature profiles based on the 1.57  $\mu\text{m}$  CO<sub>2</sub> band (Fedorova et al., 2020).  
104 The orbit of TGO allows for solar occultations that cross the terminator over a range  
105 of local times, a key feature of these datasets. ACS-NIR temperature profiles were in-  
106 cluded in our assimilation of MY 34 over the period  $L_S=163^\circ-360^\circ$ .

### 107 2.2 Global Climate Model and data assimilation scheme

108 The martian atmosphere is modelled using the UK version of the Laboratoire de  
109 Météorologique (LMD) MGCM (Forget et al., 1999). The LMD model contains numer-  
110 ous parameterizations of physical processes on Mars, including schemes for CO<sub>2</sub> conden-

111 sation and sublimation, radiative transfer, photochemistry, orographic and gravity wave  
 112 drag, cloud microphysics, convection and the boundary layer (Forget et al., 1999; Lefèvre  
 113 et al., 2004; Colaïtis et al., 2013; Navarro et al., 2014). Dust transport is via an inter-  
 114 active two-moment scheme that describes the particle size distribution using two trac-  
 115 ers, the mass mixing ratio and the number density (Madeleine et al., 2011; Spiga et al.,  
 116 2013).

117 The UK-specific variant of the model combines the physics schemes of the LMD  
 118 model with a spectral dynamical core (Hoskins & Simmons, 1975), a vertical finite dif-  
 119 ference scheme that conserves energy and angular momentum (Simmons & Burridge, 1981),  
 120 and a semi-Lagrangian tracer advection scheme (Newman et al., 2002). The model was  
 121 run at T42 resolution (triangular truncation at horizontal wavenumber 42), which cor-  
 122 responds to a physical grid resolution of  $3.75^\circ$ . 55 unevenly-spaced, terrain-following ver-  
 123 tical levels were used, with the model top at approximately 100 km.

124 The data assimilation scheme is an implementation of the Analysis Correction (AC)  
 125 scheme, originally developed for terrestrial applications (Lorenc et al., 1991) and later  
 126 modified and re-tuned for the martian context (Lewis et al., 1996, 2007). The scheme  
 127 employs a form of successive corrections, with analysis increments interlaced between dy-  
 128 namical timesteps. It has been successfully employed in assimilating profiles of temper-  
 129 ature and column dust opacity from the Thermal Emission Spectrometer (TES) (Lewis  
 130 & Barker, 2005; Montabone et al., 2005; Lewis et al., 2007) and Mars Climate Sounder  
 131 (MCS) (Steele et al., 2014; Holmes et al., 2019, 2020; Streeter et al., 2020) instruments.

132 We performed model runs for two consecutive years, MY 33 and MY 34, assimi-  
 133 lating spacecraft data into the model where available. Our primary focus is on the dust  
 134 storm period of MY 34,  $L_S=180^\circ-240^\circ$ . The MY 33 assimilation was used for compar-  
 135 ison as it was a year that did not have a GDS.

### 136 2.3 Super-rotation metrics

Super-rotation can be quantified by defining a global super-rotation index  $S$  (Lewis  
& Read, 2003) as

$$S = \frac{\iiint \rho u a \cos \phi dV}{\iiint \rho \Omega a^2 \cos^2 \phi dV} \times 100\%, \quad (1)$$

137 where  $\rho$  is density,  $u$  is the zonal wind velocity, and  $a$ ,  $\phi$  and  $\Omega$  are the planetary radius,  
 138 latitude and rotation rate respectively. The volume element is  $dV = a^2 \cos \phi d\lambda d\phi dz$ ,  
 139 where  $\lambda$  and  $z$  are east longitude and geometric height respectively. Each integral is per-  
 140 formed over the whole atmospheric volume. Positive values of  $S$  would indicate that the  
 141 atmosphere as a whole had more angular momentum than if it was in pure solid body  
 142 rotation.

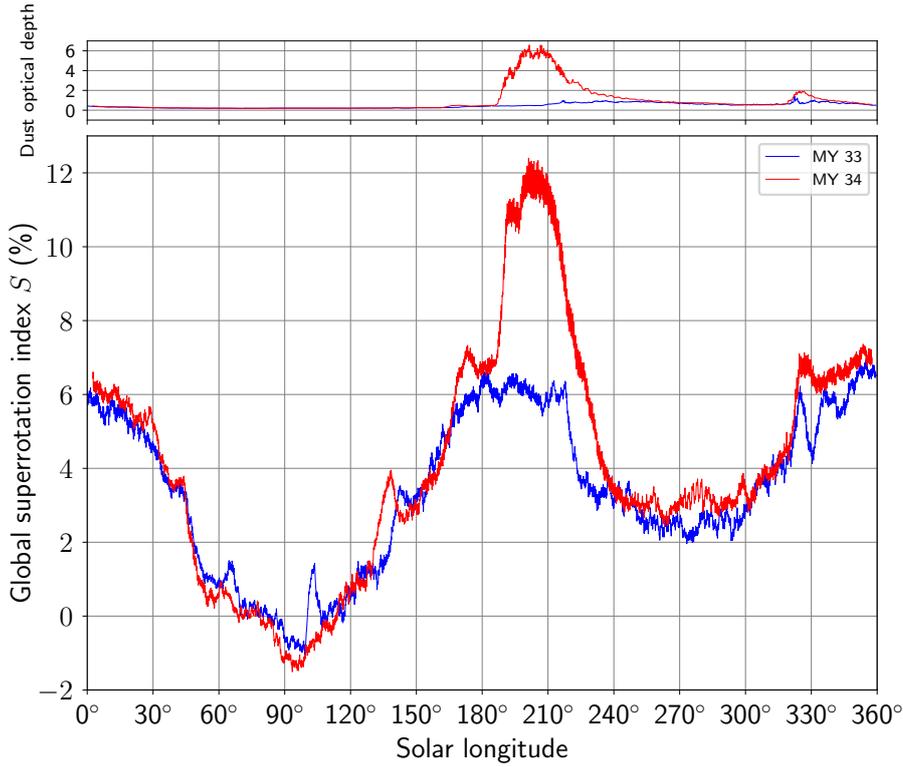
Alternatively, super-rotation can be diagnosed at a given latitude and height by  
defining a local super-rotation index  $s$  as

$$s = [\bar{m}/\Omega a^2 - 1] \times 100\%. \quad (2)$$

Here  $\bar{m}$  is the zonally averaged value of axial angular momentum  $m$ , which is defined as

$$m = a \cos \phi [\Omega a \cos \phi + u]. \quad (3)$$

143  $s$  is a local measure of super-rotation, and is essentially a comparison between the  
 144 zonal-mean angular momentum of a fluid parcel at a given latitude and height, against  
 145 the angular momentum of an equal-mass parcel at rest at the equator. Positive values  
 146 of  $s$  represent super-rotation.



**Figure 1.** Values of the global super-rotation metric  $S$  during MY 33 (blue) and MY 34 (red). Positive values indicate global super-rotation. The upper panel shows the corresponding column dust optical depth values over the equator.

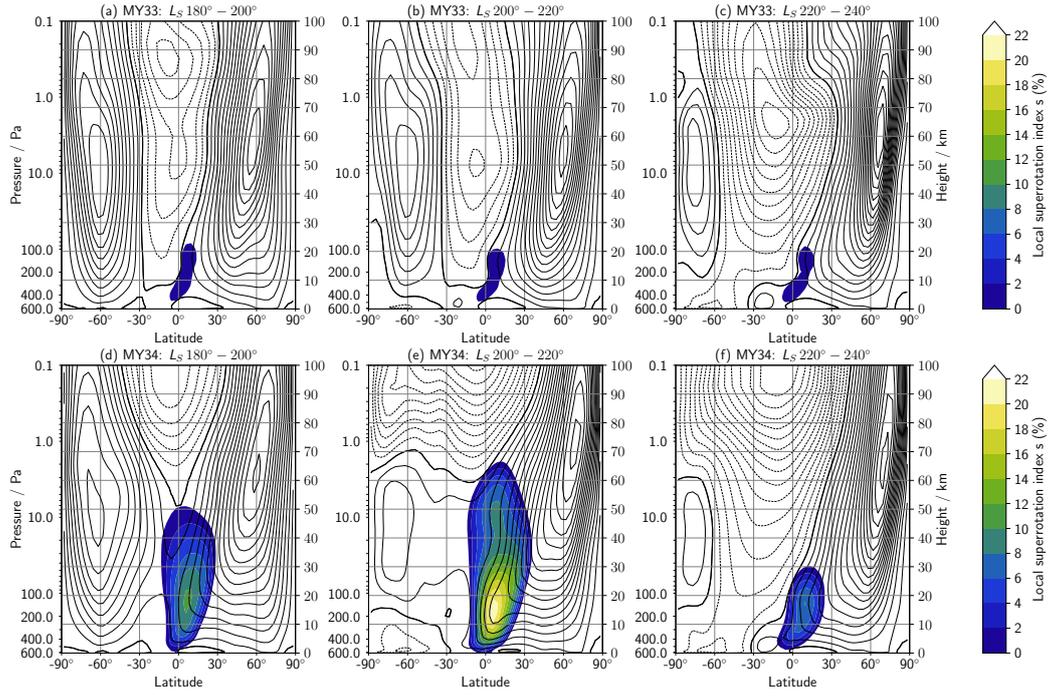
### 3 Results

#### 3.1 Atmospheric super-rotation

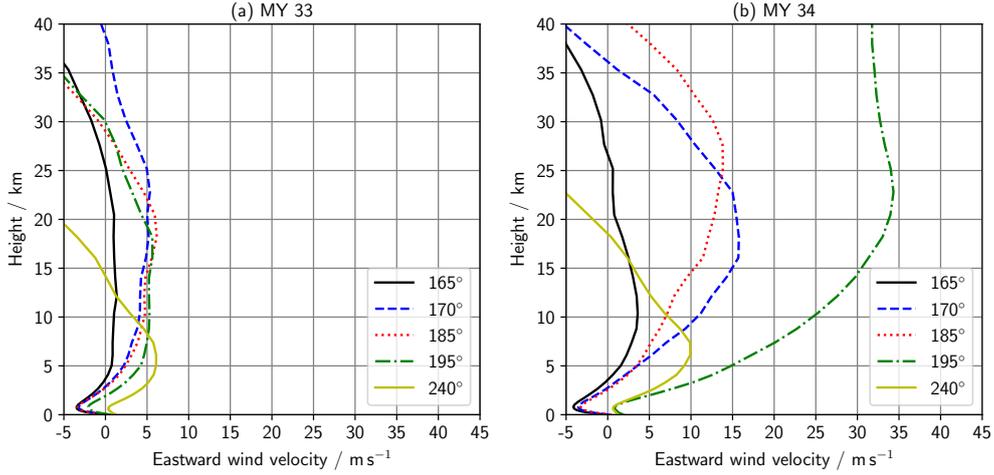
Fig. 1 shows the variation of the global super-rotation metric  $S$  through MY 33-34. A major feature in Fig. 1 is the large peak in  $S$  centered at  $L_S=200^\circ$  in MY 34, which clearly corresponds to the period of the MY 34 GDS. Global super-rotation in MY 34 first diverges significantly from MY 33 values at  $L_S=165^\circ$ , when the dust opacity in the tropics first begins to increase. After an initial peak to  $S = 7.2\%$  at  $L_S=173^\circ$ , global super-rotation increases sharply upon storm onset at  $L_S=187^\circ$ , reaching a maximum of  $S = 12.6\%$  at  $L_S=201^\circ$ . The peak value is twice as large as the corresponding MY 33 value of  $S = 6.2\%$ , indicating a significant increase in atmospheric angular momentum. After the peak,  $S$  decays and returns to background values at  $L_S=240^\circ$ .

Outside of the global dust storm period, the values of  $S$  are broadly similar in both years. Global super-rotation variations have a semi-annual structure, with broad peaks occurring during equinoxes and troughs occurring during the solstices, due to changes in the Hadley cell structure (Lewis & Read, 2003).  $S$  remains largely positive throughout both years of study, except during Northern summer solstice. The average values of  $S$  are 3.4% and 3.9% for MY 33 and MY 34 respectively.

Fig. 2 depicts values of zonal-mean zonal winds and the local super-rotation index across latitude and height, during the MY 34 GDS period in both years. The data have



**Figure 2.** Contours of zonal-mean zonal wind (black) and local super-rotation index (filled, positive values only) during the MY 34 GDS period for MY 33 (top) and MY 34 (bottom). Data in each column is averaged over a solar longitude block of width  $20^\circ$ . Zonal wind contours have been drawn at  $10 \text{ m s}^{-1}$  intervals, with eastward (westward) winds denoted by solid (dashed) lines. The bold solid contour is the zero wind line.



**Figure 3.** Profiles of zonal-mean zonal wind in the tropics during the MY 34 GDS. Zonal-mean zonal wind has been averaged across the (30° S - 30° N) latitude band, and is displayed at five different solar longitudes for MY 33 (left) and MY 34 (right) respectively.

166 been averaged in time over 20° periods of solar longitude. These broadly cover the ini-  
 167 tiation, peak and decay phases of the storm.

168 In all periods, super-rotation is much stronger in MY 34 than in MY 33. In the time  
 169 period of storm initiation (Figs. 2(a), 2(d)), the atmosphere is in an equinoctial state,  
 170 with extratropical westerly jets in both hemispheres and equatorial westerlies. This is  
 171 reflected by the positive local super-rotation values in the tropics, which a peak of  $s \approx$   
 172  $12\%$  at 20 km in MY 34 compared to  $s \approx 1\%$  in MY 33.

173 Super-rotation is strongest in the period  $L_S = 200^\circ - 220^\circ$  (Fig. 2(e)), which cor-  
 174 responds to the peak of the dust storm in MY 34. The super-rotating jet continues to  
 175 dominate the tropical band, and extends up to 60 km. There is a strong peak in local  
 176 super-rotation in the tropical atmosphere at 20 km, reflecting a significant acceleration  
 177 of tropical winds by the presence of dust.

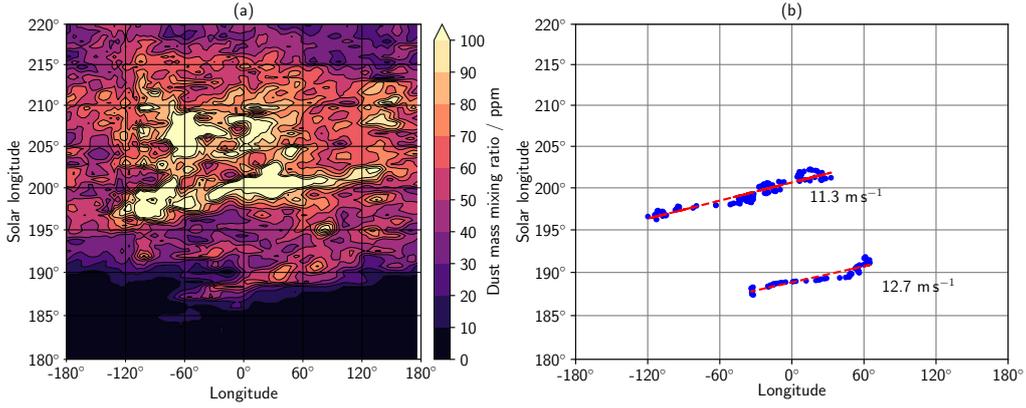
178 In the decay phase of the storm (Fig. 2(f)), easterly winds dominate the tropical  
 179 regions above 20 km. Super-rotation is greatly reduced in both intensity and extent, and  
 180 is confined to the lower atmosphere below 30 km.

### 181 3.2 Effect of super-rotation on winds and dust transport

182 Fig. 3 shows zonal-mean zonal wind profiles at four solar longitudes during the ini-  
 183 tiation phase of the GDS ( $L_S = 165^\circ, 170^\circ, 185^\circ, 195^\circ$ ), and at one solar longitude dur-  
 184 ing the decay phase of the GDS ( $L_S = 240^\circ$ ) for MY 33 (left) and MY 34 (right). Zonal  
 185 winds have been averaged over the tropical latitude band (30° S–30° N).

186 The tropical zonal wind profile does not change substantially between  $L_S = 165^\circ -$   
 187  $195^\circ$  of MY 33 (Fig. 3(a)). Easterly winds in the lowermost atmosphere are overlain by  
 188 westerly winds that peak at 20–25 km with velocities between 5–7 m s<sup>-1</sup>.

189 In contrast, enhanced super-rotation has a significant impact on the zonal wind pro-  
 190 files in MY 34. Whilst the wind profile for  $L_S = 165^\circ$  in MY 34 is similar to the winds in  
 191 MY 33 discussed above, by  $L_S = 170^\circ$  zonal wind speeds have increased at almost all heights.  
 192 The wind profile weakens slightly at  $L_S = 185^\circ$ , before strengthening substantially to reach



**Figure 4.** Eastward transport of dust by tropical winds during the MY 34 GDS. (a) Dust mass mixing ratio on the 290 Pa pressure level ( $\approx 8$  km above the surface). Values have been averaged around the equator ( $10^\circ$  S– $10^\circ$  N), and are in units of parts per million (ppm). (b) Estimates of zonal wind speed during the initiation phase of the GDS, derived by tracking the mean longitude of the 65 ppm (Chryse plume) and 100 ppm (Tharsis plume) contours.

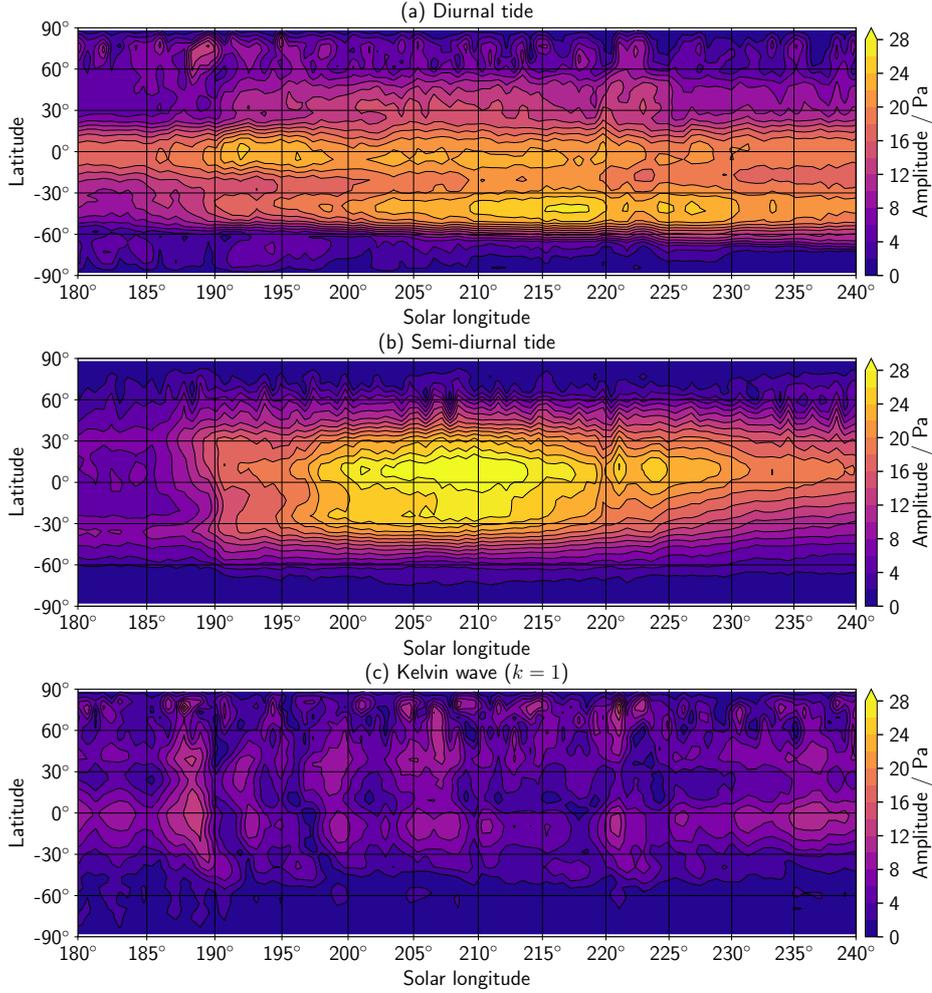
193 peak values of  $35 \text{ m s}^{-1}$  at  $L_S=195^\circ$ , with eastward winds penetrating all the way down  
 194 to the lowest model level (5 m). In the decay phase of the GDS period, wind profiles from  
 195 MY 34 start to converge with profiles from MY 33.

196 The enhanced winds in the tropics played a significant role in the transport and  
 197 distribution of dust during the initial phase of the GDS. Fig. 4(a) shows the evolution  
 198 of the dust mass mixing ratio, averaged around the equator ( $10^\circ$  S– $10^\circ$  N) and plotted  
 199 on the 290 Pa pressure level ( $\approx 8$  km above the surface). The first signs of enhanced dust  
 200 presence in the atmosphere appear in Chryse Planitia ( $40^\circ$  W,  $L_S=187^\circ$ ), from which a  
 201 coherent band of dust can be seen moving eastward from this initial dust excitation re-  
 202 gion. These winds rapidly transport the dust around the planet, eventually reaching the  
 203 Tharsis region where a large lifting event is initiated ( $90^\circ$  W,  $L_S=197^\circ$ ). Continued east-  
 204 ward dust transport results in dust being advected back to the original dust lifting re-  
 205 gion in Chryse.

206 We obtained indirect estimates of wind speeds on Mars by tracking the eastward  
 207 propagation rate of the dust plumes. At each model time point and vertical level, a di-  
 208 urnal moving average of the dust mass mixing ratio was calculated, and then averaged  
 209 across the tropical latitude band. Plots of the 65 ppm and 100 ppm contour lines were  
 210 constructed across longitude and height, and the mean longitudes of these contour lev-  
 211 els were used to estimate the plume locations. The analysis revealed two westward-propagating  
 212 plumes, shown in Fig. 4(b), with dust propagation speeds of  $12.7 \text{ m s}^{-1}$  (Chryse plume)  
 213 and  $11.3 \text{ m s}^{-1}$  (Tharsis plume). These results are within the range of model wind speeds  
 214 in the lower atmosphere below 10 km (Fig. 3(b)). We also note that these speeds would  
 215 not be consistent with model winds below 10 km in the absence of strong dust-driven forc-  
 216 ing (Fig. 3(a)), as these winds do not exceed  $5 \text{ m s}^{-1}$ .

### 217 3.3 Tidal excitation during the MY 34 GDS

218 We performed a Fourier analysis of surface pressure during the GDS period of MY  
 219 34. The daily pressure field for each sol was decomposed into the sum of Fourier modes  
 220 in wavenumber-frequency space.



**Figure 5.** Latitude-time variations of the largest surface pressure modes during the MY 34 GDS: (a) diurnal tide, (b) semi-diurnal tide, and (c)  $k = 1$  Kelvin wave.

221 The modes with the largest amplitudes over the GDS period were the Sun-synchronous  
 222 diurnal and semi-diurnal tides, which are excited by the diurnal radiative heating and  
 223 cooling of the atmosphere by solar and thermal radiation. During dusty periods, the  
 224 increased thermal forcing of the atmosphere results in the enhancement of these modes.

225 In addition to the tides, the wavenumber-1 Kelvin wave was also found to have sig-  
 226 nificant amplitude. The Kelvin wave is a resonant mode that is excited by the interac-  
 227 tion of the diurnal solar heating and the wavenumber-2 variations in topography and ther-  
 228 mal inertia on Mars (Zurek, 1976; Wilson & Hamilton, 1996).

229 Fig. 5 shows the latitudinal variation of these three largest wave modes over the GDS  
 230 period for MY 34. Amplitudes were generally strongest in the tropical region, and de-  
 231 creased toward the poles.

232 The diurnal tide was the largest mode during the initiation phase, with an ampli-  
 233 tude of 16 Pa (Fig. 5(a)). In comparison, the amplitudes of the other modes were quite  
 234 weak, with the amplitude of the semi-diurnal tide not exceeding 4 Pa. Upon the onset  
 235 of dust lifting in the Chryse Planitia region at  $L_S = 187^\circ$ , the amplitudes of all three  
 236 modes began to increase. Both the diurnal and semi-diurnal tides experienced a rela-

237 tively monotonic increase in amplitude with increased dust loading. In contrast, the Kelvin  
 238 wave response was more varied, increasing dramatically to a peak of 14 Pa at  $L_S=188^\circ$ ,  
 239 before decaying rapidly back to background levels of between 2–8 Pa.

240 As the storm intensified, the amplitude of both sun-synchronous tides increased.  
 241 Growth of the semi-diurnal tide was more significant due to the essentially linear rela-  
 242 tionship between the semi-diurnal tide and total dust loading as demonstrated by Lewis  
 243 and Barker (2005). It eventually surpassed the amplitude of the diurnal tide at the peak  
 244 of the dust storm, with a value of almost 30 Pa compared to 22 Pa. In the decay phase  
 245 of the storm, the amplitudes of both modes gradually reduced, with the amplitude of  
 246 the semi-diurnal tide decreasing faster than the amplitude of the diurnal tide.

## 247 4 Discussion

248 The tropical westerly jet was strongly enhanced during the MY 34 GDS. This had  
 249 a very large impact on the value of the global super-rotation index  $S$ , which doubled in  
 250 the dust storm period from 6.2% (MY 33) to 12.6% (MY 34). The jet dominated the  
 251 tropical atmosphere during the peak of the storm, with local super-rotation extending  
 252 up to 60 km. These changes in tropical wind behaviour had significant consequences for  
 253 atmospheric tracer transport.

254 Montabone et al. (2020) reported a value of  $S = 16\%$  during the peak of the MY  
 255 34 GDS, compared to a pre-dust storm value of 5%. Our results agree qualitatively, and  
 256 indicate a large increase in super-rotation during the dust storm period. Quantitative  
 257 differences are likely to be due to the different dynamical cores (spectral versus grid-point)  
 258 and vertical resolutions (50 versus 29) used by the different models (UK versus LMD).

259 Transport of dust in the initial phases of the storm was largely controlled by the  
 260 strength of the westerly tropical jet, which determined the timescale over which dust lofted  
 261 in Chryse Planitia was transported eastward towards Tharsis. The initial increase in the  
 262 strength of global super-rotation occurred between  $L_S=163^\circ$ – $170^\circ$ , when dust lifted in  
 263 the southern hemisphere began encroaching into the tropical band. The enhancement  
 264 of winds occurred quite some time before GDS onset at  $L_S=187^\circ$ , and might be an im-  
 265 portant precursor condition for equinoctial GDS development.

266 The jet-driven eastward transport of dust back into the Chryse basin from Thar-  
 267 sis during the GDS resulted in a significant rejuvenation of dust activity in that region.  
 268 This suggests that dust activity in Chryse Planitia may have decayed - and the dust storm  
 269 not have become global in extent - had the enhanced westerly equatorial jet not been  
 270 strong enough to transport dust from Tharsis into the Chryse basin.

271 The propagation of dust plumes in the assimilation was used to estimate wind speeds  
 272 in the lower martian atmosphere. The plume locations were constrained by satellite-derived  
 273 observations of the dust distribution, wherever available; at other times, model advec-  
 274 tion of the dust fields provided self-consistency and improved predictive capabilities at  
 275 times where observations were unavailable. The self-consistency of these results gives us  
 276 reason to expect winds on the order of  $10$ - $15 \text{ m s}^{-1}$  below 10 km during equinoctial global  
 277 dust storms.

278 Our results provide strong evidence that a strong super-rotating jet forms in the  
 279 tropics during an equinoctial GDS. This jet has a major role to play in determining the evo-  
 280 lution pathway of a GDS, and we have provided estimates of wind speeds in the lowest  
 281 10 km of the tropical atmosphere that are consistent with our claims. Future work in this  
 282 area could explore the dynamics of the super-rotating jet in different GDS years. In ad-  
 283 dition, the role of the Kelvin wave in the initial phase of the GDS would benefit from  
 284 further clarification.

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