

On the variability of total electron content over Europe during the 2009 and 2019 Northern Hemisphere SSWs

T. A. Siddiqui¹, Y. Yamazaki², C. Stolle^{2,3}, A. Maute⁴, J. Laštovička⁵, I. K. Edemskiy⁵, Z. Mošna⁵

¹Leibniz-Institute of Atmospheric Physics at the , Kühlungsborn, Germany

²GFZ German Research Centre for Geosciences, Potsdam, Germany

³Faculty of Science, University of Potsdam, Potsdam, Germany

⁴High Altitude Observatory, National Center for Atmospheric Research, Boulder, Colorado, USA

⁵Institute of Atmospheric Physics, Czech Academy of Sciences,, Prague, Czech Republic

Key Points:

- The large-scale dynamical variabilities associated with 2009 and 2019 SSWs were simulated using WACCM-X
- Numerical experiments are carried out with TIE-GCM driven by WACCM-X to isolate out the geomagnetic and lower atmospheric effects during both SSWs
- Simulation results show that the TEC variability over Europe was caused predominantly by geomagnetic forcing for the 2019 SSW

Corresponding author: T. A. Siddiqui, siddiqui@iap-kborn.de

17 **Abstract**

18 The nature of the variability of the Total Electron Content (TEC) over Europe is inves-
 19 tigated during the 2009 and 2019 Northern Hemisphere (NH) SSW events in this study.
 20 As the TEC variability is driven by geomagnetic and lower atmospheric forcing mech-
 21 anisms, we investigate the dominant drivers and their respective contributions to TEC
 22 changes during both SSW events. We simulate the SSWs using the Whole Atmosphere
 23 Community Climate Model eXtended version (WACCM-X) and compare the semidiur-
 24 nal solar and lunar tidal variabilities in the mesosphere-lower thermosphere (MLT) re-
 25 gion. Further, in order to assess the mechanisms responsible for the TEC variability dur-
 26 ing both SSWs, we run numerical experiments using the National Center for Atmospheric
 27 Research (NCAR) Thermosphere-Ionosphere Electrodynamics General Circulation Model
 28 (TIE-GCM). We constrain the TIE-GCM lower boundary with the WACCM-X fields and
 29 carry out simulations both with and without geomagnetic forcing for each of the SSWs.
 30 The TIE-GCM simulations allow us to isolate the geomagnetic and lower atmospheric
 31 forcing effects on the TEC. We find that there was a major enhancement in daytime TEC
 32 over Europe during the 2019 SSW event, which was predominantly geomagnetically forced
 33 ($\sim 80\%$), while for the 2009 SSW, the major variability in TEC was accounted for by lower
 34 atmospheric forcing.

35 **1 Introduction**

36 The axial tilt of the Earth creates a strong temperature gradient between the pole
 37 and mid-latitudes due to lack of solar heating over the polar regions in wintertime. This
 38 temperature gradient along with the Earth's coriolis force results in the formation of strato-
 39 spheric polar vortex (SPV) that appear over high-latitudes every winter before break-
 40 ing down in the summer when the polar regions start becoming warmer (e.g., Polvani
 41 et al., 2013). The SPV manifests itself in the form of planetary-scale westerly (flowing
 42 from west to east) winds that encircle the pole at mid-to higher latitudes. Due to inter-
 43 action with the planetary waves (PWs) that originate from the troposphere, the SPV
 44 experiences large intra- and inter annual variability. The PWs can propagate vertically
 45 into the stratosphere when the stratospheric winds are westerly (Charney & Drazin, 1961)
 46 and as the PWs break in the stratosphere (McIntyre & Palmer, 1983), they deposit their
 47 momentum, which leads to the deceleration of the polar vortex. Due to larger topographic
 48 and land-sea contrasts, the PW activity is higher in the Northern Hemisphere (NH) as

49 compared to the Southern Hemisphere (SH) (e.g., van Loon et al., 1973), which results
50 in the NH stratospheric polar vortex being weaker and much more variable than the one
51 in SH (Vaugh & Randel, 1999). Certain extreme cases of PW activity can lead to a break-
52 down of the polar vortex and a reversal of the westerly flow resulting in magnificent tran-
53 sient meteorological events known as sudden stratospheric warmings (SSWs) (e.g., Scher-
54 hag, 1952; Matsuno, 1971).

55 SSWs are large-scale events that result in an increase in the stratospheric polar tem-
56 perature usually by several tens of degrees within a few days, which reverses the merid-
57 ional temperature gradient (i.e., polar temperatures exceed those at mid-latitudes), and
58 concurrently, a deceleration of the westerly zonal mean zonal wind (ZMW) (e.g., An-
59 dreds et al., 1987). In the literature, SSWs are classified as major and minor warmings
60 depending on the extent of increase in polar stratospheric temperature and ZMW re-
61 versal. According to the definition of World Meteorological Organization (WMO), an
62 SSW event is said to be “major” if the westerly zonal mean flow, poleward of 60° at 10
63 hPa, reverses to easterly along with the reversal of the meridional temperature gradi-
64 ent. A “minor” SSW event occurs when the polar stratospheric temperature increases
65 by at least 25 K within a week or faster. Major SSWs are common in the NH and oc-
66 cur with a frequency of 0.6/year (e.g., Charlton & Polvani, 2007; Butler et al., 2015) whereas
67 SSWs in SH are extremely rare. Along with minor SSWs in August 2010 (Eswaraiah et
68 al., 2017) and September 2019 (Yamazaki et al., 2020), only one major warming in Septem-
69 ber 2002 (e.g. Baldwin, 2003; Allen et al., 2003) has ever been recorded in the SH.

70 The breaking of PWs during the SSWs and its associated effects are not only lim-
71 ited to polar stratosphere, but are rather witnessed across different latitudes and alti-
72 tudes (Pedatella et al., 2018). A mean meridional circulation is induced in the strato-
73 sphere as a result of PWs breaking (Haynes et al., 1991), which leads to an upwelling
74 at equatorial latitudes and results in adiabatic cooling over these regions (Fritz & Soules,
75 1970). Accompanying the warming in the stratosphere is the cooling in the mesosphere
76 at polar latitudes (e.g., Labitzke, 1972; Liu & Roble, 2002) and warming in the meso-
77 sphere at equatorial latitudes (e.g., Garcia, 1987; Chandran & Collins, 2014). The NH
78 SSW related effects are also witnessed in the SH in the form of mesospheric warming and
79 a decrease in the occurrence of polar mesospheric clouds through inter-hemispheric cou-
80 pling mechanisms (e.g., Karlsson et al., 2009; Körnich & Becker, 2010). The associated

81 effects of NH SSW in the thermosphere have also been reported to result in warming at
82 mid- and high latitudes (e.g., Goncharenko & Zhang, 2008; Funke et al., 2010).

83 In the last decade, evidence of the SSW impact in the ionosphere, mostly facilitated
84 by the extremely quiet solar and geomagnetic activity levels in the last solar cycle, have
85 been reported in numerous studies following the seminal works by Goncharenko and Zhang
86 (2008) and Chau et al. (2009). Observations have revealed a consistent enhanced semi-
87 diurnal perturbation pattern at low-latitudes in ionospheric vertical plasma drifts (e.g.,
88 Chau et al., 2009), electron densities (e.g., Lin et al., 2013) and equatorial electrojet (e.g.,
89 Vineeth et al., 2009; Fejer et al., 2010; Yamazaki et al., 2012) in response to SSWs. These
90 perturbations have been linked to the changes in migrating semidiurnal solar (SW2) and
91 lunar (M2) atmospheric tides during SSWs. The SW2 and M2 tides are generated in the
92 lower atmospheric regions and are able to propagate upward to the dynamo-region heights
93 where they influence the generation of electric fields in the ionosphere (Baker et al., 1953).
94 The SW2 enhancement during SSWs is thought to be due to the changes in tidal prop-
95 agation conditions (Jin et al., 2012), non-linear interaction with the stationary PWs (Liu
96 et al., 2010) and changes in the stratospheric ozone distribution (Goncharenko et al., 2012;
97 Siddiqui et al., 2019). The cause of M2 amplification is suggested to be a result of back-
98 ground zonal mean zonal wind changes, which shifts the secondary (Pekeris) resonance
99 peak of the atmosphere towards the period of M2 tide (Forbes & Zhang, 2012). It has
100 also been found during SSWs that the relative enhancement of M2 tide in the mesosphere-
101 lower thermosphere (MLT) and ionosphere is larger than that of SW2 (Pedatella, Liu,
102 Richmond, Maute, & Fang, 2012) and the comparably smaller amplitudes of M2 can even
103 exceed those of SW2 (e.g., Chau et al., 2015; Siddiqui et al., 2018).

104 A large number of studies that have reported the impact of SSWs on the ionosphere
105 have focused on the variability at low-latitudes (e.g., Chau et al., 2012; Yiğit & Medvedev,
106 2015, and references therein). In particular, most of these studies have focused on the
107 January 2009 major NH SSW (e.g., Manney et al., 2009), which occurred in the begin-
108 ning of solar cycle 24 under extremely quiet solar and geomagnetic conditions, and dis-
109 cussed the ionospheric impacts over different longitudinal sectors (e.g., Chau et al., 2010;
110 Goncharenko et al., 2010; Fejer et al., 2010; Xiong et al., 2013; Patra et al., 2014; Ya-
111 dav et al., 2017; Liu et al., 2019) during this prolonged SSW event. Towards the end of
112 solar cycle 24, another major SSW event was recorded in the NH under similar quiet so-
113 lar and geomagnetic activity conditions in the final weeks of December 2018 and in the

114 beginning of January 2019. As the occurrence of SSWs under such favorable conditions
115 is seldom, this event provides us further opportunities to investigate the ionospheric im-
116 pacts of SSWs. Compared to the investigation of SSW related ionospheric variabilities
117 at equatorial and low-latitudes, the mid-latitude ionosphere variability has not yet been
118 thoroughly investigated. Although the evidence of mid-latitude ionosphere variability
119 during SSWs have been reported in some studies (e.g., Goncharenko et al., 2013; Polyakova
120 et al., 2014; Chen et al., 2016), the Total Electron Content (TEC) data over the Euro-
121 pean region, which hosts a dense global navigation satellite system (GNSS) networks,
122 have not been much exploited. In this study, we compare the observed TEC over Eu-
123 rope during the 2018/2019 and 2008/2009 SSWs and investigate the dominant mecha-
124 nisms behind the variabilities using simulations.

125 The structure of this paper is as follows. In Section 2, the descriptions of the mod-
126 els and the experiments are provided followed by the information about the data sets
127 that are used in this study. In Section 4, we present our results followed by discussion
128 in Section 5. The summary and conclusions from this work are presented at the end.

129 **2 Model descriptions and experiment settings**

130 **2.1 WACCM-X**

131 We use WACCM-X version 2.0 (Liu et al., 2018), a configuration of the NCAR Com-
132 munity Earth System Model (CESM 2.0; Hurrell et al., 2013) to perform the model sim-
133 ulations in this study. WACCM-X extends from the surface to the upper thermosphere
134 with its top boundary, depending on the solar and geomagnetic activity, lying between
135 500 and 700 km. The vertical resolution of WACCM-X above the stratosphere is one fourth
136 of a scale height, and the horizontal resolution is $1.9^\circ \times 2.5^\circ$ in latitude and longitude,
137 respectively. WACCM-X is built upon the Whole Atmosphere Community Climate Model
138 (WACCM) (Marsh et al., 2013) and Community Atmosphere Model version 4 (Neale et
139 al., 2013). The new version of WACCM-X has a coupled ionosphere and incorporates
140 self-consistent low-mid-latitude ionospheric electrodynamics adapted from the Thermosphere-
141 Ionosphere Electrodynamics General Circulation Model (TIE-GCM). At high-latitudes,
142 WACCM-X uses an empirical electric potential pattern (Heelis et al., 1982), which is pa-
143 rameterized by the 3-hour geomagnetic Kp index and an auroral precipitation oval based

144 on the formulation described by Roble and Ridley (1987). More details about the phys-
145 ical processes included in WACCM-X 2.0 can be found in Liu et al. (2018).

146 WACCM-X provides a comprehensive tool to study the entire atmosphere-ionosphere
147 system. The impact of lower atmospheric forcing on the upper atmospheric variability
148 can be studied using WACCM-X during specific time periods by constraining the tro-
149 pospheric and stratospheric dynamics to meteorological reanalysis fields. In the present
150 study, we use the specified dynamics (SD) set up in WACCM-X to simulate the SSWs
151 by constraining the winds and temperatures from 0-50 km towards the National Aero-
152 nautics and Space Administration Modern Era Retrospective Analysis for Research and
153 Applications (MERRA) Version 2 (Gelaro et al., 2017). Using the approach described
154 in Kunz et al. (2011), the WACCM-X model fields are constrained to the MERRA2 me-
155 teorological fields at every model time step (i.e., 5 min). The default SD-WACCM-X set
156 up does not include forcing from the M2 tide but we implement it in our simulations based
157 on the method described by Pedatella et al. (2012) because the M2 forcing becomes an
158 important source of MLT and ionospheric variability during SSWs. Hourly outputs of
159 winds, temperature and geopotential height are obtained for the 2009 and 2019 SSWs
160 through the SD-WACCM-X runs in this study.

161 2.2 TIE-GCM

162 TIE-GCM is a three dimensional, self-consistent numerical model of the coupled
163 thermosphere-ionosphere system that has been developed at the High Altitude Obser-
164 vatory at the National Center for Atmospheric Research (NCAR). The model spans from
165 ~ 97 km to about 450-600 km depending on the solar cycle activity. In this study, we use
166 TIE-GCM version 2.0 with a horizontal resolution of 2.5° by 2.5° in geographic longi-
167 tude and latitude and a vertical resolution of 0.25 times the scale height. The input pa-
168 rameters for TIE-GCM include the solar XUV, EUV and FUV spectral fluxes that are
169 defined by the EUVAC model (Richards et al., 1994) using the F10.7 index. The global
170 electric potential due to the wind dynamo is solved by the self-consistent TIE-GCM iono-
171 spheric electrodynamic at low- and mid-latitudes. At high-latitudes, however, the elec-
172 tric potential is prescribed through empirical convection electric field patterns using the
173 Heelis (Heelis et al., 1982) or Weimer (Weimer, 2005) models. TIE-GCM also uses an
174 analytical auroral model to account for high-latitude auroral particle precipitation in the

175 default set up. The upper part of WACCM-X is based on TIE-GCM and it uses this same
176 auroral model.

177 The effect of lower atmospheric tidal forcing can be specified in TIE-GCM using
178 the tidal perturbations at its lower boundary. The amplitudes and phases of upward prop-
179 agating atmospheric tides specified at the model lower boundary in the default setup are
180 based on the Global Scale Wave Model (GSWM) (Hagan et al., 1999). The default TIE-
181 GCM lower boundary (LB) assumes constant neutral temperature ($T=181$ K), geopo-
182 tential height ($Z=96.37$ km) and zero horizontal winds (e.g., Maute, 2017). For a more
183 realistic LB conditions, TIE-GCM includes the option to specify hourly inputs at its LB
184 from any other source. In this study, we specify the hourly WACCM-X outputs for the
185 2009 and 2019 SSWs at the TIE-GCM LB and carry out two simulations for each of the
186 SSWs in order to examine the effects of lower atmospheric forcing on the thermosphere-
187 ionosphere system. In the first simulation setup (hereafter referred to as S1), the TIE-
188 GCM forced by WACCM-X is run in its default setting and the obtained day-to-day iono-
189 spheric variability from this run includes the effects of both geomagnetic and lower at-
190 mospheric forcings. In the second simulation setup (hereafter referred to as S2), we turn-
191 off the geomagnetic forcing and carry out a similar run for both the SSWs. The day-to-
192 day ionospheric variability resulting from the second run arises solely from lower atmo-
193 spheric forcing during both SSWs. The geomagnetic forcing is turned off from the TIE-
194 GCM runs by reducing the hemispheric power from 18 GW in the first setup to 0.1 GW
195 and the cross-polar cap potential from 30 kV in the first setup to 0.1 kV. Additionally,
196 the Heelis convection model and the analytical auroral model have also been turned off
197 in the second setup to remove the magnetospheric energy input. The experiment setups
198 used in this study have been summarized in Table 1. We used TIE-GCM instead of WACCM-
199 X to carry out these two simulations because of the former being computationally less
200 expensive than the latter.

201 **3 Data Sets**

202 For this study we use the GPS TEC data from the MIT Haystack Observatory's
203 Madrigal database (Rideout & Coster, 2006), which incorporates the data from over 2000
204 GPS receivers worldwide. The processed TEC data from the MIT Automated Process-
205 ing of GPS (MAPGPS) software provides estimates of TEC over 1° by 1° (latitude by

206 longitude) bins with a temporal resolution of 5 minutes. The unit of GPS TEC data is
 207 TECu, where 1 TECu is defined as 10^{16} electrons/m².

208 Hourly values of solar flux ($F_{10.7}$) (Tapping, 2013) have been downloaded from the
 209 Space Physics Data facility of the Goddard Space Flight Center through the OMNIWeb
 210 data interface to plot the levels of solar activity during the 2009 and 2019 SSWs. The
 211 3-hourly Kp indices are downloaded from the website of German Research Centre for Geo-
 212 sciences (GFZ), Potsdam to monitor the geomagnetic activity levels during the two SSW
 213 events. We also use the version 5 temperature data from the Microwave Limb Sounder
 214 (MLS) onboard the NASA's Earth Observing System (EOS) Aura satellite (Waters et
 215 al., 2006) to compare and validate the temperature obtained from WACCM-X simula-
 216 tions.

217 4 Results

218 4.1 Zonal mean and PW variability during 2009 and 2019 SSWs

219 In Figure 1, the zonal mean temperatures averaged between 70° and 80° N dur-
 220 ing January - February 2009 (Figures 1a and 1c) and December 2018 - January 2019 (Fig-
 221 ures 1b and 1d) from Aura Microwave Limb Sounder (MLS) observations (top panels)
 222 and SD-WACCM-X simulations (bottom panels) are presented. The vertical white lines
 223 in the figure mark the day of polar vortex weakening (PVW). As an alternative to the
 224 classical definition of SSW provided by the WMO, PVW has been used to correlate the
 225 tidal enhancements in the MLT and ionosphere with the magnitude of the reversal of strato-
 226 spheric zonal mean zonal wind (ZMZW) (e.g., Zhang & Forbes, 2014; Chau et al., 2015;
 227 Siddiqui et al., 2015). A PVW day is identified by locating the earliest and most extreme
 228 reversal of ZMZW at 70°N and 48 km altitude (1 hPa) that occurs simultaneously with
 229 the increase in zonal mean temperature at North Pole and 40 km altitude (3 hPa) be-
 230 tween December and February.

231 It can be seen that there is a reasonable agreement between the observed and sim-
 232 ulated zonal mean temperatures during the considered time intervals. With the onset
 233 of the SSW, the warm stratopause in each of the SSW events starts to descend from its
 234 climatological position near 60 km (0.2 hPa) toward lower altitudes resulting in warm-
 235 ing at these heights (e.g., Labitzke, 1981) before breaking down completely. The stratopause
 236 then reappears at higher altitude before slowly returning to its original location. This

237 altitudinal shift and reemergence of the stratopause is called an elevated stratopause event
 238 and is associated with major SSW events (e.g., Siskind et al., 2007; Manney et al., 2008).
 239 In WACCM-X simulations the elevated stratopause events are being reproduced for both
 240 the SSWs but their appearances tend to occur slightly earlier in comparison with the Aura
 241 MLS observations. The PVW date for the 2009 SSW event occurs on day 2 (Jan 23) and
 242 for the 2019 SSW event on day -3 (28 Dec). The warming of the stratosphere is accom-
 243 panied by cooling in the mesosphere in both these events. In case of 2009 SSW, meso-
 244 spheric cooling can be observed to start around day 20 above 0.1 hPa in Figure 1a and
 245 in the case of 2019 SSW close to day -8 in Figure 1b. The model simulations are able
 246 to reproduce these features to a large extent and the overall qualitative agreement in tem-
 247 perature is good but some discrepancies can also be seen especially at altitudes where
 248 SD-WACCM-X stops being constrained to MERRA2, i.e. above 50 km (~ 1 hPa). The
 249 quantitative differences between SD-WACCM-X and Aura MLS observations start to be-
 250 come more clear above this altitude and is more pronounced in the MLT region. These
 251 temperature differences could be related to gravity wave forcing, which is parameterized
 252 in SD-WACCM-X and may be contributing to the discrepancy between the observations
 253 and model simulations (e.g., Smith, 2012).

254 The left panels of Figure 2 presents the ZMZW at 60°N from SD-WACCM-X for
 255 the time intervals that include the 2009 and 2019 SSW events while the right panels show
 256 the Kp index and daily solar flux (s.f.u) conditions. The dotted black lines show the PVW
 257 days in all the panels. For the 2009 SSW, the ZMZW in the stratosphere between 1-10
 258 hPa changes from eastward to westward direction starting around day 20 in Figure 2a.
 259 At 10 hPa, the ZMZW remains in the reversed direction until day 54 while at 1 hPa, where
 260 the PVW is defined, the ZMZW reaches a peak reversal on day 23 with a value of about
 261 -52 m/s. For the 2019 SSW, the reversal of the ZMZW between 1-10 hPa is seen towards
 262 the end of December around day -8 in Figure 2c. The westward direction of ZMZW at
 263 10 hPa remains till day 21 while at 1 hPa, the ZMZW is found to reach a peak rever-
 264 sal on day -3 with a value of -15 m/s. From the comparison of the ZMZW in Figures 2a
 265 and 2c, it can be clearly seen that the 2009 SSW event was stronger in terms of ZMZW
 266 reversal and more prolonged than the 2019 SSW event.

267 From the Kp indices and solar flux values in Figures 2b and 2d, it can be inferred
 268 that both the 2009 and 2019 SSWs were recorded under periods of low solar and geo-
 269 magnetic activities. The solar flux levels during the 2009 SSW and 2019 SSWs remained

270 below 75 s.f.u for both the events. The geomagnetic activity during the onset and peak
 271 phase of 2009 SSW hovered mostly around $K_p \leq 2+$ with an exception on day 19 when
 272 K_p values reached 4o during the 3-hourly intervals. The 2019 SSW event showed higher
 273 geomagnetic activity levels as compared to the 2009 SSW event with brief periods of spike
 274 in K_p values that reached up to 4+ on day -3 (Dec 28) and 5o on day 5 but overall the
 275 period during the 2019 SSW remained geomagnetically quiet.

276 Figure 3 presents the variability of PWs with wave number 1 (PW1) and 2 (PW2)
 277 at 10 hPa in temperature from SD-WACCM-X for the 2009 and 2019 SSWs. In this fig-
 278 ure, the colorbar scales of PW1 are chosen to be twice as large in magnitude as compared
 279 to those for PW2. In Figures 3a and 3c, the amplitudes of PW1 and PW2, respectively,
 280 are presented for the 2009 SSW event. The enhancement of PW1 is seen in NH high-
 281 latitudes particularly around 60°N with peaks on days 6 and 22 while in case of PW2,
 282 the enhancement begins in the second week of January and peaks on day 19. Based on
 283 the enhanced amplitudes of PW2 in Figure 3c, we find that our results are consistent
 284 with earlier studies (e.g., Manney et al., 2009) that have classified the 2009 SSW as PW2
 285 forced split SSW event. The amplitudes of PW1 and PW2 for the 2009 SSW are also
 286 similar to the results shown by Pedatella et al. (2014) using four different whole atmo-
 287 sphere models.

288 The amplitude of PW1 and PW2 in temperature are presented in Figures 3b and
 289 3d, respectively, for the 2019 SSW event. The PW1 amplitudes show enhancement pole-
 290 ward of 60°N during December with maxima on day -7. The peak PW1 amplitudes for
 291 this event are almost twice as large as compared to that of peak PW1 for the 2009 SSW.
 292 The PW2 amplitudes are smaller in magnitude as compared to that of PW1 for this event
 293 but enhancement in PW2 is also seen poleward of 60°N towards the end of December
 294 with maxima centered on day -9. The 2019 SSW has been classified in literature as nei-
 295 ther a typical displaced nor a typical split SSW event (Rao et al., 2019) but rather a mixed
 296 type event, which was initially a displaced SSW and later became a split SSW. From Fig-
 297 ure 3b, it can be inferred by the dominance of PW1 amplitudes that the 2019 SSW must
 298 have started as a displaced SSW event.

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4.2 Migrating semidiurnal tides during 2009 and 2019 SSWs

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It is well established now that the primary reason for the variability in the MLT and ionosphere during SSWs is due to the modulation of atmospheric tides. In particular, the variability of SW2 and M2 have been found to be the most significant based on modeling and observational studies (e.g., Vineeth et al., 2009; Chau et al., 2009; Fejer et al., 2010; Pedatella & Forbes, 2010; Goncharenko et al., 2012; Yamazaki et al., 2012; Forbes & Zhang, 2012; Lin et al., 2013; Pedatella et al., 2014; He et al., 2017; Siddiqui et al., 2018; Hibbins et al., 2019). Atmospheric tides refer to planetary-scale oscillations of the atmosphere that are mainly excited by gravitational forces of the moon and by thermal forcing from the sun (e.g., Lindzen & Chapman, 1969; Forbes & Garrett, 1978; Forbes, 1982). These oscillations have periods and sub-periods of a solar or a lunar day. The solar tides form the dominating component of the tidal oscillations and are predominantly thermally forced. The excitation mechanisms include daily periodic absorption of solar energy by tropospheric water vapour and stratospheric ozone (e.g., Forbes & Wu, 2006; Zhang et al., 2010) while the relatively smaller lunar tides are mainly forced due to the lunar gravitational effects on the Earth's atmosphere. In this section, we investigate the variability of the migrating semidiurnal solar (SW2) and semidiurnal lunar (M2) tides during the two SSW events.

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The hourly outputs of neutral temperature from SD-WACCM-X simulations are used to extract the components of the solar and lunar tides by performing a least squares fit with a moving window of 15 days at each latitude. The following equation based on Pedatella et al. (2012) has been used to make the fit:

$$\sum_{n=0}^3 \sum_{s=n-5}^{n+5} A_{n,s} \cos(n\Omega t + s\lambda - \phi_{n,s}) + \sum_{s=-3}^3 L_s \cos(2\tau + (s-2)\lambda - \Phi_{n,s}) \quad (1)$$

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where $\Omega = \frac{2\pi}{24}$ hour⁻¹, t is universal time in hours, n denotes the harmonics of a solar day, λ is the longitude, s is the zonal wave number, $A_{n,s}$ and $\phi_{n,s}$ are the amplitude and phase of the respective solar tidal components. For a wave propagating westward $s > 0$, while for a wave propagating eastward $s < 0$. L_s and $\Phi_{n,s}$ represent the amplitude and the phase of the semidiurnal lunar tide, respectively.

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Figure 4 presents the amplitudes and phases of the SW2 (top panels) and M2 (bottom panels) tides in neutral temperature at 1×10^{-4} hPa (~ 110 km altitude) from SD-

328 WACCM-X simulations between January and February 2009. The vertical white lines
 329 mark the day of PVW. In Figure 4a, the SW2 tidal variability shows enhanced ampli-
 330 tudes of up to 21 K at Southern Hemisphere (SH) mid-latitudes on day 15 and then again
 331 amplitudes of over 25 K around day 35. In the NH, the SW2 enhancements are only seen
 332 after day 25 but enhanced amplitudes of up to 24 K are found from simulations. Follow-
 333 ing the PVW day, there is a sudden decrease in the amplitude of SW2 at SH mid-latitudes.
 334 This feature of SW2 was also reported in a study by Pedatella et al. (2014) where they
 335 compared the temporal variability of SW2 during the 2009 SSW event using four dif-
 336 ferent whole atmosphere models. Figure 4c presents the amplitude of M2 tide where we
 337 notice its enhancement in both hemispheres a few days after the PVW day. Another en-
 338 hancement of M2 is seen in the NH after day 40. The M2 variability from SD-WACCM-
 339 X during this SSW event is consistent with the results of Zhang and Forbes (2014), in
 340 which similar observations of M2 have been reported from neutral temperature measure-
 341 ments at 110 km from the Sounding of the Atmosphere using Broadband Emission Ra-
 342 diometry (SABER) instrument onboard the Thermosphere Ionosphere Mesosphere En-
 343 ergetics Dynamics (TIMED) satellite. The SW2 and M2 phases in UT are presented in
 344 Figures 4b and 4d, respectively. At NH low- and mid-latitudes, the SW2 phase shows
 345 a noticeable change in phase before and after the peak PVW day. At SH mid-latitudes,
 346 there is a change in SW2 phase of up to 1 h coinciding with the weakening of SW2. The
 347 decrease of SW2 phase by a few hours in the ionosphere was also reported by Pedatella
 348 et al. (2014) and Lin et al. (2013) during the period where the SW2 amplitude weakened.
 349 The phase of the M2 tide shows more discernible phase shifts at low- and mid-latitudes
 350 around the PVW, which becomes relatively more stable after day 40.

351 Figure 5 shows the amplitudes and phases of SW2 and M2 tides in neutral tem-
 352 perature at 1×10^{-4} hPa (~ 110 km altitude) between December 2018 and January 2019.
 353 Similar to the SW2 variability during the 2009 SSW event, the SW2 amplitude in Fig-
 354 ure 5a shows enhancement at SH mid-latitudes on either side of the PVW day with a
 355 reduction of SW2 amplitude in between. The SW2 amplitude reaches a peak of around
 356 17 K on day -12 and around 20 K on day 8 in the SH. In the NH, the SW2 amplifica-
 357 tion starts close to the PVW day and it reaches a value of around 18 K on day 9. In Fig-
 358 ure 5c, the first M2 enhancement in SH with peak amplitudes of around 6.5 K happens
 359 a few days after the PVW day. The second M2 enhancement in the SH with peak am-
 360 plitude of around 7.5 K is seen on day 11. Following Chau et al. (2015) and Conte et al.

(2017), we also performed tidal analyses using a 21-day running window to further reduce any artifacts or ambiguity between the determination of SW2 and M2 tides but we still found the amplification of M2 to be similar to as shown in Figure 5c using a 15-day window. We find that the M2 amplitude during the 2019 SSW is far smaller as compared to that during the 2009 SSW event. One reason for the difference in amplification could be related to the timing of the SSW event relative to the phase of the moon. Pedatella and Liu (2013) have shown from simulation results that the lunar tidal response in the ionosphere is dependent upon the phase of the moon relative to the timing of the SSW event. Further, Fejer et al. (2010) have found that the lunar effects in the ionosphere during SSWs amplify close to the new or the full moon days. The new moon day for the 2009 SSW occurred 3 days after the PVW day while for the 2019 it was recorded 6 days before the PVW. This difference in the timing of the new moon days relative to the PVW days may also be contributing to the reduced amplification of M2 during the 2019 SSW. The SW2 and M2 phases are presented in Figures 5b and 5d, respectively. At low-latitudes in NH, the SW2 phase shows a decrease of 1-2 hours around the PVW day. The decrease in the phase is also seen in the SH around the PVW day similar to the reduction seen during the 2009 SSW event. From the M2 phase plot it is found that the M2 phase shift is more noticeable in the SH as compared to the NH around the PVW day.

Figure 6 presents the M2 tidal amplitudes for the 2009 and 2019 SSW events obtained using the V2.0 temperature measurements from the SABER instrument onboard the TIMED satellite. We employ the least-squares fitting method mentioned in Zhang and Forbes (2014) to determine the M2 amplitudes from zonally averaged SABER temperature residuals. As the period of the M2 tide from the frame of the TIMED satellite is 11.85 days (Forbes et al., 2013), we use a 12-day moving window and fit only the M2 tide to the daily zonally averaged temperature residuals. Figures 6a and 6b show the amplitude of the M2 tide for the 2009 and 2019 SSWs, respectively. The vertical white lines show the days of PVW. The enhancement of M2 following the PVW can be seen at low- and mid- latitudes for the 2009 SSW event in Figure 6b. This plot of M2 amplitude is similar to the one shown in Zhang and Forbes (2014) (see Figure 1). The M2 tides in neutral temperature from SD-WACCM-X simulations for the 2009 SSW, shown in Figure 4c, match very well with the M2 from SABER temperature observations in terms of the timing of the M2 enhancement. In the NH, the variability of the M2 amplitudes from SD-WACCM-X simulations is slightly more consistent with those obtained from

394 SABER temperature observations as the M2 peaks are located at similar latitudes. In
 395 the SH, the M2 peaks from SABER observations on day 27 are slightly more equator-
 396 ward as compared to those from SD-WACCM-X simulations. The M2 amplitudes for the
 397 2019 SSW presented in Figure 6b shows that the level of M2 enhancement for this SSW
 398 was clearly lower than for the 2009 SSW. The M2 enhancements can be seen to occur
 399 towards the end of December 2018 between days -7 and 0 poleward of 40°S and around
 400 day 22 above 20°N. The enhancement in the SH towards the end of December is also cap-
 401 tured in M2 from SD-WACCM-X simulations in Figure 5c. But the M2 amplitudes from
 402 SD-WACCM-X do not exactly reproduce the observations for the 2019 SSW event as the
 403 M2 peak around day 22 in the NH is not seen in the simulations. There is similarity be-
 404 tween the observed and simulated M2 amplitudes in the sense that the enhancements
 405 were particularly weaker for the 2019 SSW as compared to the 2009 SSW event, which
 406 is confirmed from both the simulated and observed M2 amplitudes.

407 **4.3 GPS TEC variability over Europe during 2009 and 2019 SSWs**

408 In this section of the paper, we first analyze the variability of TEC over Europe
 409 from GPS observations during the 2019 SSW. The readers may note that there has been
 410 a change in the order of presentation of the SSWs in this section. The GPS TEC obser-
 411 vations over Europe during the 2009 SSW is presented later for the purpose of compar-
 412 ison. Following Goncharenko et al. (2010), we first define the mean state of the quiet-
 413 time TEC during the 2019 SSW to investigate the TEC variability over Europe associ-
 414 ated with this event. We select the quiet period prior to the SSW onset between 11 and
 415 20 December, 2018 with $Kp \leq 3$ on most days and solar flux levels below 70 s.f.u. to es-
 416 timate the mean TEC values. The perturbations in TEC during the 2019 SSW, ΔTEC ,
 417 is then calculated by using,

$$\Delta TEC = TEC - TEC_{mean} \quad (2)$$

418 In Figure 7, we present the ΔTEC (in TECu) observed between December 26, 2018
 419 and January 6, 2019 at 12 UT over Europe. The most notable feature in this figure is
 420 the large positive ΔTEC that is recorded on December 28 over the whole region in gen-
 421 eral and over the South-West part of Europe in particular. Apart from this day, ΔTEC
 422 values also show notable enhancements over Europe the following day on December 29

423 and again on January 5 but with lower magnitudes. A depletion in ΔTEC across South-
 424 ern Europe is also observed on January 3. We take the averaged value of TEC over the
 425 region shown in Figure 7 for each UT and present the temporal evolution of averaged
 426 TEC (Figure 8a) and ΔTEC (Figure 8b) across Europe as a function of local time. We
 427 use latitudinal and longitudinal bands between $35^\circ\text{-}60^\circ\text{N}$ and $15^\circ\text{W-}30^\circ\text{E}$, respectively,
 428 to calculate these values. The dotted black lines mark the PVW day. The diurnal vari-
 429 ation of TEC is evident in Figure 8a with the TEC values increasing gradually from morn-
 430 ing to afternoon hours and then after attaining a maximum during afternoon decreas-
 431 ing gradually post sunset. The day-to-day variability of TEC, which is subject to solar,
 432 geomagnetic and lower atmospheric forcing, can also be seen in Figure 8a. The spike in
 433 both averaged TEC and ΔTEC values on days -3 (Dec 28) and -2 (Dec 29) between 10-
 434 12 UT is also clear in this figure. From the Kp values shown for this time interval in Fig-
 435 ure 2d, we notice that there was an increase in Kp index on December 28 with the max-
 436 imum values on this day reaching up to 4+ and not decreasing below 3o. It is also to
 437 be noted that the timing of TEC enhancement is coinciding with the PVW day for the
 438 2019 SSW event. As the influence of SSW events on the variability of TEC is now rel-
 439 atively well known (e.g., Goncharenko et al., 2010; Yue et al., 2010; Sumod et al., 2012;
 440 Goncharenko et al., 2013; Polyakova et al., 2014; Vieira et al., 2017), there is a motiva-
 441 tion to investigate whether the TEC and ΔTEC spikes were linked to increased atmo-
 442 spheric forcing or to geomagnetic forcing. In the next section, we investigate the dom-
 443 inant forcing mechanism that is causing this TEC variability over Europe in more de-
 444 tail with the help of TIE-GCM simulations but first we present the TEC and ΔTEC vari-
 445 ability over Europe for the 2009 SSW event in Figure 9 for comparison with those from
 446 the 2019 SSW.

447 We first select the quiet period prior to the SSW onset between 3 and 12 January,
 448 2009 with $K_p \leq 3$ on most days and solar flux levels below 70 s.f.u. to estimate the quiet-
 449 time mean TEC values. These exact dates have been used by Goncharenko et al. (2010)
 450 in an earlier study to characterize the mean behavior of the quiet-time ionosphere dur-
 451 ing the 2009 SSW. ΔTEC is then calculated using equation 2 as before. Along with the
 452 diurnal and day-to-day variability of TEC, we see major TEC enhancements in Figure
 453 9a between 10 and 12 UT on days 41 and 46. In Figure 9b, ΔTEC shows the perturba-
 454 tions of TEC from the quiet-time mean values and the large TEC perturbations partic-
 455 ularly after day 40 is more clearly noticeable. The spikes in ΔTEC on days 41 and 46

456 between 10-12 UT are 3-5 TECu greater than the mean TEC values at these hours. We
 457 have limited the presentation of TEC data to day 50 to minimize the effects of seasonal
 458 transition as daytime TEC values start to depart from the calculated mean values. We
 459 notice from Figure 2c that the Kp values on day 41 remained below 1 and on day 46 be-
 460 low 3. Further, the solar flux values ranged between 65-68 s.f.u. on these days. Based
 461 on the levels of geomagnetic activities on these days, it is fair to assume that this TEC
 462 driver solely cannot explain the large TEC perturbations that are witnessed in Figure
 463 9b. From the results of Goncharenko et al. (2012), it is known that the TEC variances
 464 (computed as a departure from the mean daytime values) at low-latitudes in the Amer-
 465 ican sector can be up to 5 times larger than those obtained from the International Ref-
 466 erence Ionosphere (IRI) model for almost a month after the 2009 SSW event. The large
 467 TEC variances in their study was attributed to the modified tidal forcing associated with
 468 the 2009 SSW. In the next section, we explore using TIE-GCM simulations whether the
 469 enhancement in the averaged TEC plot in Figure 9a is also linked to the increased lower
 470 atmospheric forcing during the 2009 SSW event.

471 5 Discussions

472 5.1 Variability of simulated TEC over Europe during 2009 and 2019 SSWs

473 Based on the GPS observations, it seems that the TEC variability over Europe dur-
 474 ing the 2019 SSW may also be affected by moderately enhanced geomagnetic activity
 475 levels as seen by the increase in Kp values. As the TEC variability is influenced by both
 476 geomagnetic and lower atmospheric forcing, it is imperative to separate the effects of these
 477 two processes and assess their individual contributions towards the variability of TEC.
 478 To separate the influence of geomagnetic and lower atmospheric forcing on the TEC, we
 479 carry out two simulations using TIE-GCM. For the 2019 SSW, the TIE-GCM model is
 480 run with and without geomagnetic forcing to isolate the mechanisms behind the TEC
 481 variability during these events. In Figure 10, the TEC over Europe, which is derived from
 482 TIE-GCM simulations, is presented for the 2019 SSW. Figure 10a shows the average TEC
 483 over Europe for the TIE-GCM run where both the geomagnetic and lower atmospheric
 484 forcings have been switched on (S1). We note that the TEC derived from TIE-GCM sim-
 485 ulations are able to reproduce the TEC spikes on days -3 (Dec 28) and 5 (Jan 5) and
 486 are qualitatively similar in comparison to the averaged GPS TEC in Figure 8a. The mod-
 487 eled and observed TEC may only be compared in a qualitative sense owing to the up-

488 per boundary limit of TIE-GCM, which extends to only about 750 km in altitude, hin-
 489 ders a quantitative comparison with GPS TEC observations. The modeled TEC results
 490 demonstrate that the TIE-GCM simulation includes the various forcing mechanisms that
 491 are responsible for the observed TEC variability and can be used to filter the TEC vari-
 492 ability associated with SSW. The averaged TEC enhancements over Europe reach val-
 493 ues of about 8.5 TECu at 12 UT on day -3 and 6.5 TECu at 13 UT on day 5. In com-
 494 parison to the quiet-time seasonal TEC levels, the TEC spikes represent an increase of
 495 more than 100% at 12 UT on day -3 and more than 50% at 13 UT on day 5, which is
 496 due to the combined effects of geomagnetic and lower atmospheric forcing.

497 In Figure 10b, the averaged TEC over Europe is presented for the TIE-GCM run
 498 where the geomagnetic forcing have been switched off and only lower atmospheric forc-
 499 ing remains active (S2). It can be noticed that there is an apparent reduction of the av-
 500 eraged TEC values after turning off geomagnetic forcing and the peak TEC values have
 501 been reduced to less than 5 TECu. The major TEC enhancements in this plot can be
 502 seen on days -3 and 4 with peak TEC values reaching 5 and 4.5 TECu, respectively, be-
 503 tween 10-14 UT. In Figure 10c, the difference of TEC values from the two simulations
 504 ($TEC_{diff} = TEC_{S1} - TEC_{S2}$) is presented. The filled contour lines are plotted when ab-
 505 solute value of TEC_{diff} exceeds 1 TECu. It can be clearly noted that the major differ-
 506 ence in TEC is seen on days -3 and 5 when TEC_{diff} values between 12-14 UT reach 3.6
 507 and 2.3 TECu, respectively. Through this plot, the contribution to TEC variability solely
 508 due to geomagnetic forcing can be assessed. It can also be easily inferred that the ma-
 509 jor processes behind the averaged TEC spikes on days -3 and 5 between 10-14 UT are
 510 related to geomagnetic forcing. For a quantitative breakdown of the contribution of ge-
 511 omagnetic and lower atmospheric forcing to TEC variability, we first calculate the sea-
 512 sonal TEC levels ($TEC_{seasonal}$) from the quiet-days between 11 and 20 December in TEC_{S2} .
 513 The contribution of geomagnetic forcing, in percentage, to the TEC variability can then
 514 be calculated by using the relation,

$$\frac{TEC_{diff} \times 100\%}{(TEC_{S1} - TEC_{seasonal})} \quad (3)$$

515 The contribution of geomagnetic forcing to the TEC variability in percentage is plot-
 516 ted through the dashed black and blue open contour lines at 40 and 80% levels, respec-
 517 tively.

518 For the TEC enhancements on day -3 (Dec 28) at 12 UT, the geomagnetic contri-
 519 bution comes out to be 78% and the remaining 22% is the contribution due to lower at-
 520 mospheric forcing. In case of the TEC enhancements on day 5 (Jan 5) at 12 UT, the ge-
 521 omagnetic and lower atmospheric contributions comes out to be about 82% and 18%,
 522 respectively.

523 The individual contributions of the geomagnetic and lower atmospheric forcing on
 524 the TEC variability over Europe is also assessed for the 2009 SSW event. We present
 525 the averaged TEC over Europe derived from TIE-GCM simulations for the 2009 SSW
 526 in Figure 11. The averaged TEC derived from the TIE-GCM run with S1 and S2 setups
 527 are presented in Figure 11a and 11b, respectively. A qualitative comparison with the GPS
 528 TEC observations in Figure 9a suggests that the primary features of the averaged TEC
 529 variability have been consistent in the simulations. The comparatively lower levels of av-
 530 eraged TEC before day 30 between 10-15 UT and the moderately enhanced averaged TEC
 531 levels after this day has been correctly reproduced in the simulation. The spike in the
 532 average TEC on day 45 (February 14), seen more pronounced in Figure 9a, has also been
 533 reproduced in the simulations but it is delayed by an hour in comparison with the ob-
 534 servations. The spike in TIE-GCM derived TEC reaches 5.3 TECu on day 45 at 13 UT.
 535 Compared to the quiet-time seasonal variations at this UT, which is calculated using the
 536 TEC values between day 3 and 12, the increase in TEC comes out to be 1.2 TECu on
 537 day 45. This represents an increase in TEC values by about 30% from seasonal varia-
 538 tions. In comparison with the 2019 SSW, we notice that the major source of averaged
 539 TEC variability for the 2009 SSW comes due to the lower atmospheric forcing. This point
 540 becomes even more clear through Figure 11c, which shows the difference of the averaged
 541 TEC values from the two simulations. The filled contour lines in this figure are again
 542 shown for values greater than 1 TECu and the dashed open contour lines mark the con-
 543 tribution of the geomagnetic forcing to the TEC variability in percentage. We notice that
 544 unlike the 2019 SSW, the TEC difference plot for the 2009 SSW points to reduced con-
 545 tribution from geomagnetic forcing to the average TEC variability over Europe during
 546 this SSW.

547 **5.2 Possible reasons for the observed TEC variability**

548 Most of the studies that have documented the ionospheric effects of SSWs, espe-
 549 cially during the 2009 SSW, have focused on the variability at equatorial and low-latitudes

(e.g., Chau et al., 2012; Yiğit & Medvedev, 2015, and references therein). While it is now accepted that the mechanisms causing the variability at these latitudes are driven by the changes in the vertically propagating semidiurnal solar and lunar tides, the mechanisms responsible for the mid-latitude ionospheric variability during SSWs are not as well understood. Simulation results by Pedatella and Maute (2015) have shown that the variability of the mid-latitude ionospheric F-region peak height (hmF2) during SSWs is predominantly driven by the field-aligned neutral winds, which is modulated by the M2 tidal enhancements. Yue et al. (2010) observed the global ionospheric response using Constellation Observing System for Meteorology, Ionosphere, and Climate (COSMIC) satellites during the 2009 SSW and suggested the changes in the neutral wind and composition due to direct propagation of tides as another mechanism for the ionospheric mid-latitude variability during SSWs. It is also known from observations and modeling studies that the influence of SSWs at mid- and high-latitudes ionosphere is generally smaller as compared to that at low-latitudes. Oyama et al. (2014) used the FORMOSAT-3/ COSMIC peak ionospheric electron density (NmF2) data and found that changes in mid-latitude NmF2 to be only between 20-30% during the 2009 SSW, which was comparably much lower than the changes at low-latitudes. Their observation results were found to be consistent with the simulations shown by Pedatella and Maute (2015). For the 2009 and 2019 SSW event, our results also show similar numbers as the TEC variability over Europe increases by \sim 20-30% with respect to seasonal variations due to SSW associated effects.

During the 2019 SSW event, along with the increase in Kp values to 4+ on December 28, the meridional component of the interplanetary magnetic field (IMF) in Geocentric Solar Magnetospheric (GSM) coordinate system, B_z turned southward and reached up to -7.5 nT, the Auroral Electrojet (AE) index reached to levels >500 nT. The symmetric disturbance field in H (SYM-H) index declined from 26 nT on December 27 to -30 nT on December 28. Based on the statistical measures, the geomagnetic activity parameters resemble the conditions of a weak geomagnetic storm (e.g., Yokoyama & Kamide, 1997). The sudden surge in TEC observations over Europe on December 28 between 10-12 UT as seen in Figure 7 could be resulting due to a result of the positive storm effect mechanism. During geomagnetic storms, the relative increase in the ionospheric plasma with respect to quiet-time conditions is referred to as positive storm effect or positive ionospheric storms. (e.g., Matsushita, 1959; Goncharenko et al., 2007; Astafyeva et al., 2016). The positive storm effect can arise due to the change in the direction of the merid-

583 ional thermospheric winds from poleward to equatorward, which results in the transport
584 of plasma along the magnetic field lines to altitudes where the recombination rates are
585 lower (e.g., Jones & Rishbeth, 1971). During daytime, this results in an increase in the
586 F-region plasma densities. Another mechanism that plays an important role in the in-
587 crease of plasma densities during geomagnetic storms is the penetration of high-latitude
588 convection electric fields into the low-latitude ionosphere. This phenomenon more com-
589 monly known as prompt penetration electric fields (PPEF) (e.g., Sastri, 1988; Abdu et
590 al., 1995) can increase the plasma densities in the equatorial ionization anomaly (EIA)
591 and shift the EIA crests poleward, in some cases, by upto 15 degrees in latitude (e.g.,
592 Astafyeva et al., 2020). The poleward shift of the EIA may also result in the observed
593 TEC surge over Europe. During geomagnetically active times, ionospheric plasma den-
594 sities is driven by a combination of these drivers and the interplay of these processes could
595 be responsible for the observed spike in TEC seen across Europe in Figure 7.

596 It is well known that the dynamics of the thermosphere-ionosphere system is greatly
597 modified due to increased geomagnetic forcing processes at high-latitudes (e.g., Prólss,
598 1995). Pedatella (2016) performed a numerical simulation using TIE-GCM that show-
599 cased the ionospheric variability in response to a SSW event that was contrived to oc-
600 cur simultaneously along with a major geomagnetic storm. The results from their study
601 showed that the TEC changes due to a geomagnetic storm are significantly ($\sim 50-100\%$)
602 different when the effects of SSW were included in the simulation. Pedatella (2016) con-
603 cluded that the changes in the thermospheric composition due to SSWs (Korenkov et
604 al., 2012) may influence the geomagnetic storm related composition changes in the ther-
605 mosphere. The TEC variability over Europe seen in Figure 7 could also be influenced
606 by the changes in thermospheric composition due to both SSW and geomagnetic related
607 effects. However, more research is needed to understand the role and contribution of ther-
608 mospheric composition changes due to different drivers that lead to ionospheric variabil-
609 ities.

610 Similar to the 2019 SSW event, another SSW event in January 2012 was accom-
611 panied by a moderate geomagnetic storm and the TEC disturbances during this period
612 were studied at low- and mid-latitudes in the Brazilian sector by Vieira et al. (2017). Sim-
613 ilar to our observations in Figure 7, a spike in daytime TEC values, was seen in their re-
614 sults following the occurrence of a moderate geomagnetic storm during this SSW event.
615 To separate the potential contribution of geomagnetic and lower atmospheric drivers that

616 are responsible for the TEC variability, Vieira et al. (2017) used a regional empirical TEC
617 model. Another study by Liu et al. (2019) analyzed the TEC response to the 2018 SSW
618 at low- and mid-latitude stations in China and attempted to separate the effects of iono-
619 spheric drivers on TEC by correlating the TEC perturbations with F10.7 solar flux, Ap
620 index and the solar wind speed. Spikes in TEC on certain days during this event was
621 attributed to the increased geomagnetic activity in their study. The use of simulations
622 to study the ionospheric effects solely due to lower atmospheric forcing have been im-
623 plemented in some earlier studies. Pedatella and Maute (2015) simulated the ionospheric
624 effects during the 2013 SSW event only due to lower atmospheric forcing by running the
625 TIME-GCM simulations under constant solar and geomagnetic activity levels. Yamazaki
626 et al. (2014) used a similar method by running the TIE-GCM under constant solar and
627 magnetospheric energy inputs to study the day-to-day variability of the equatorial elec-
628 trojet due to lower atmospheric forcing. Our method builds upon the techniques that
629 are presented in Pedatella and Maute (2015) and Yamazaki et al. (2014) and provides
630 a tool to separate the geomagnetic and lower atmospheric forcing effects on the ionospheric
631 variability. The occurrence of an SSW event during quiet-time ionospheric conditions
632 is a rarity and there have been many SSWs that occur under active geomagnetic con-
633 ditions, which complicates the separation of SSW driven ionospheric variability.

634 **6 Conclusions**

635 The variability in the mid-latitude TEC over Europe was investigated during the
636 2019 SSW in the present study using GPS TEC observations and TIE-GCM simulations.
637 The main feature of the TEC response during this SSW was a dramatic spike in the day-
638 time TEC that lasted for a couple of days. The geomagnetic activity indices suggest that
639 the 2019 SSW period was also accompanied by weak geomagnetic storm like conditions,
640 which coincided with the spike in TEC values. As the TEC variability is influenced by
641 both geomagnetic and lower atmospheric forcings, we used TIE-GCM simulations to in-
642 vestigate the contributions of each of the individual mechanisms towards the TEC en-
643 hancement. To quantify the isolated influence of either geomagnetic or lower atmospheric
644 forcing on TEC, we first force the TIE-GCM lower boundary with the output from WACCM-
645 X simulations performed over the 2019 SSW period, and then conduct two numerical sim-
646 ulations. The first TIE-GCM simulation includes both geomagnetic and lower atmospheric
647 forcing while in the second simulation the geomagnetic forcing has been turned off. We

ascertain the individual contributions of geomagnetic and lower atmospheric forcing towards the sudden TEC enhancement that is witnessed during the 2019 SSW through these two simulations. We further compare the TEC variability during 2019 SSW over Europe with that of the TEC variability during the 2009 SSW. Based on the results we summarize our findings as follows:

1. It is found that the spike in TEC over Europe during the 2019 SSW is forced dominantly due to increase in geomagnetic activity which accounts for about $\sim 80\%$ of the TEC variability while the remaining variability is accounted for by the lower atmospheric forcing.
2. In contrast, the variability of TEC over Europe during the 2009 SSW event was up to 30% in comparison to seasonal variations and was predominantly due to lower atmospheric forcing.
3. The mid-latitude TEC variability associated with lower atmospheric forcing during the 2019 SSW lies between 20-30% relative to seasonal TEC values, which is similar to the levels reported by previous observation and modeling based studies that documented the mid-latitude ionospheric variability during the 2009 SSW.

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1053 **Table 1.** TIE-GCM simulations: (left to right) experiment setup and the representation of
 1054 geomagnetic and lower atmospheric forcings in the model runs.

Experiment setup	Geomagnetic Forcing	Lower atmospheric forcing
S1	On	On
S2	Off	On

1055 **Figure 1.** Daily zonal mean temperature (K) averaged between 70° and 80°N as a function
1056 of pressure is presented from (a) Aura MLS observations (c) SD-WACCM-X simulations for the
1057 2009 SSW. The same is presented in Figures 1b and 1d, except for the 2019 SSW. The vertical
1058 dashed black lines mark the day of PVW for the corresponding SSWs.

1059 **Figure 2.** Daily zonal mean zonal wind (ZMZW) (m/s) at 60°N as a function of pressure is
1060 presented from SD-WACCM-X simulations for (a) 2009 SSW (c) 2019 SSW. The Kp index and
1061 solar flux levels are presented for the 2009 SSW in Figure 2b and for the 2019 SSW in Figure 2d.
1062 The vertical dashed black lines mark the day of PVW for the corresponding SSWs.

1063 **Figure 3.** Planetary wave 1 amplitude of temperature at 10 hPa (~ 30 km) from SD-
1064 WACCM-X simulations for (a) 2009 SSW (b) 2019 SSW. The same is presented for planetary
1065 wave 2 for (c) 2009 SSW and (d) 2019 SSW. The vertical white dashed lines mark the day of
1066 PVW for the corresponding SSWs.

1067 **Figure 4.** SW2 tidal amplitude (a) and phase (b) in neutral temperature at 1×10^{-4} hPa
1068 (~ 110 km) from SD-WACCM-X for the 2009 SSW event. The same is presented for M2 tidal
1069 amplitude (c) and phase (d). The vertical white dashed lines mark the day of PVW for the
1070 corresponding SSWs.

1072 **Figure 6.** Amplitude of M2 tide at 110 km from SABER V2.0 temperature measurements
1073 for (a) 2009 SSW (b) 2019 SSW. The vertical white dashed lines mark the day of PVW for the
1074 corresponding SSWs.

1075 **Figure 7.** Daily TEC perturbations (ΔTEC) over Europe from GPS TEC observations be-
1076 tween December 26, 2018 and January 6, 2019 at 12 UT.

1077 **Figure 8.** Daily averaged (a) TEC and (b) TEC perturbations (Δ TEC) over Europe from
1078 GPS TEC observations as a function of universal time for the 2019 SSW. The contour lines in (b)
1079 are only plotted when absolute Δ TEC exceeds 1 TECu. The vertical black dashed lines mark the
1080 day of PVW.

1081

Figure 9. Same as Figure 8 except for the 2009 SSW.

1082 **Figure 10.** TIE-GCM derived daily averaged TEC over Europe for simulation setup (a) S1
1083 and (b) S2, as a function of universal time for the 2019 SSW. The difference of (a) and (b) is
1084 plotted in (c). The filled contour lines in (c) are only plotted when absolute TEC difference
1085 exceeds 1 TECu. The dashed black and blue open contour lines mark the contribution of geo-
1086 magnetic forcing to the TEC variability at 40 and 80% levels, respectively. The vertical black
1087 dashed lines mark the day of PVW.

1088

Figure 11. Same as Figure 10 except for the 2009 SSW.

Figure 1.

Zonal mean temperature (K), 70-80 N

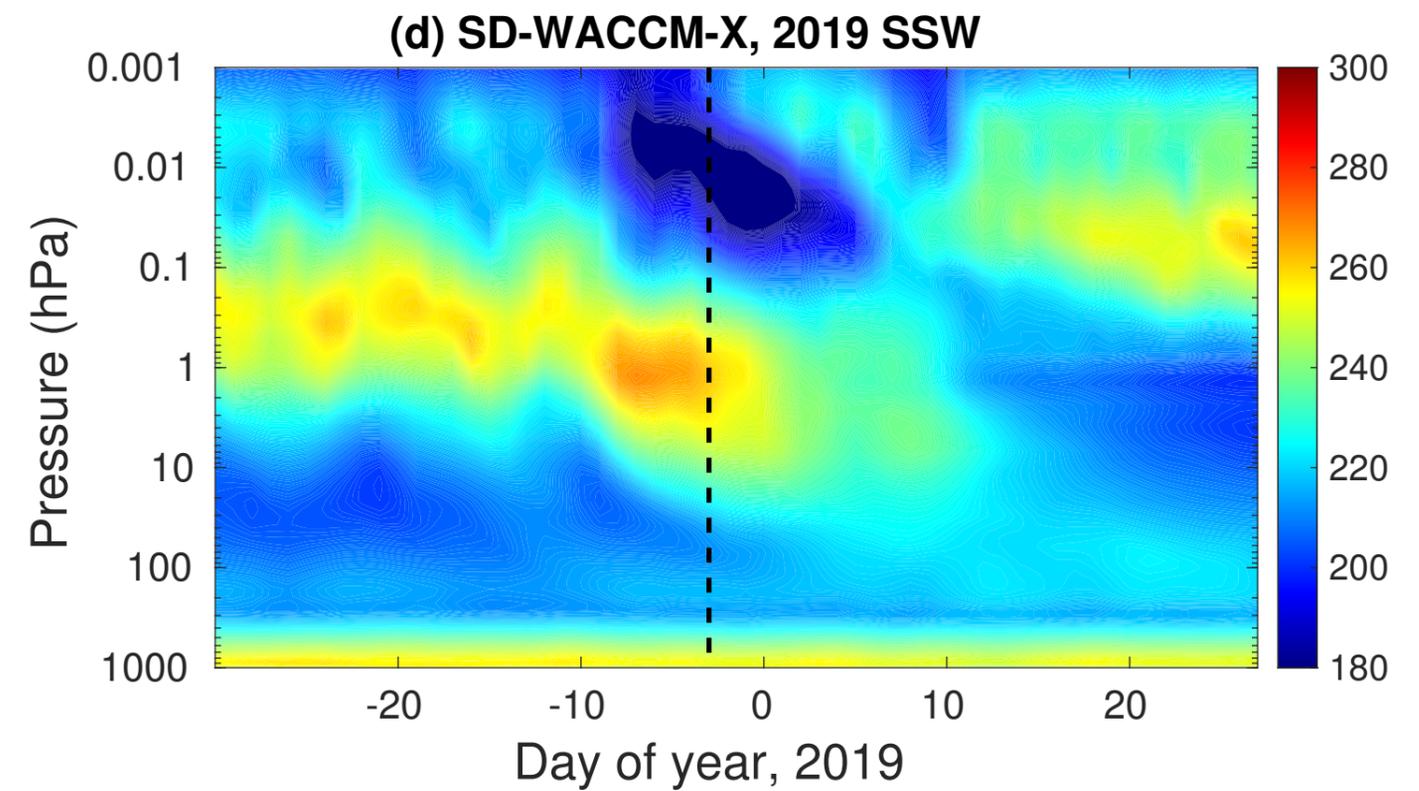
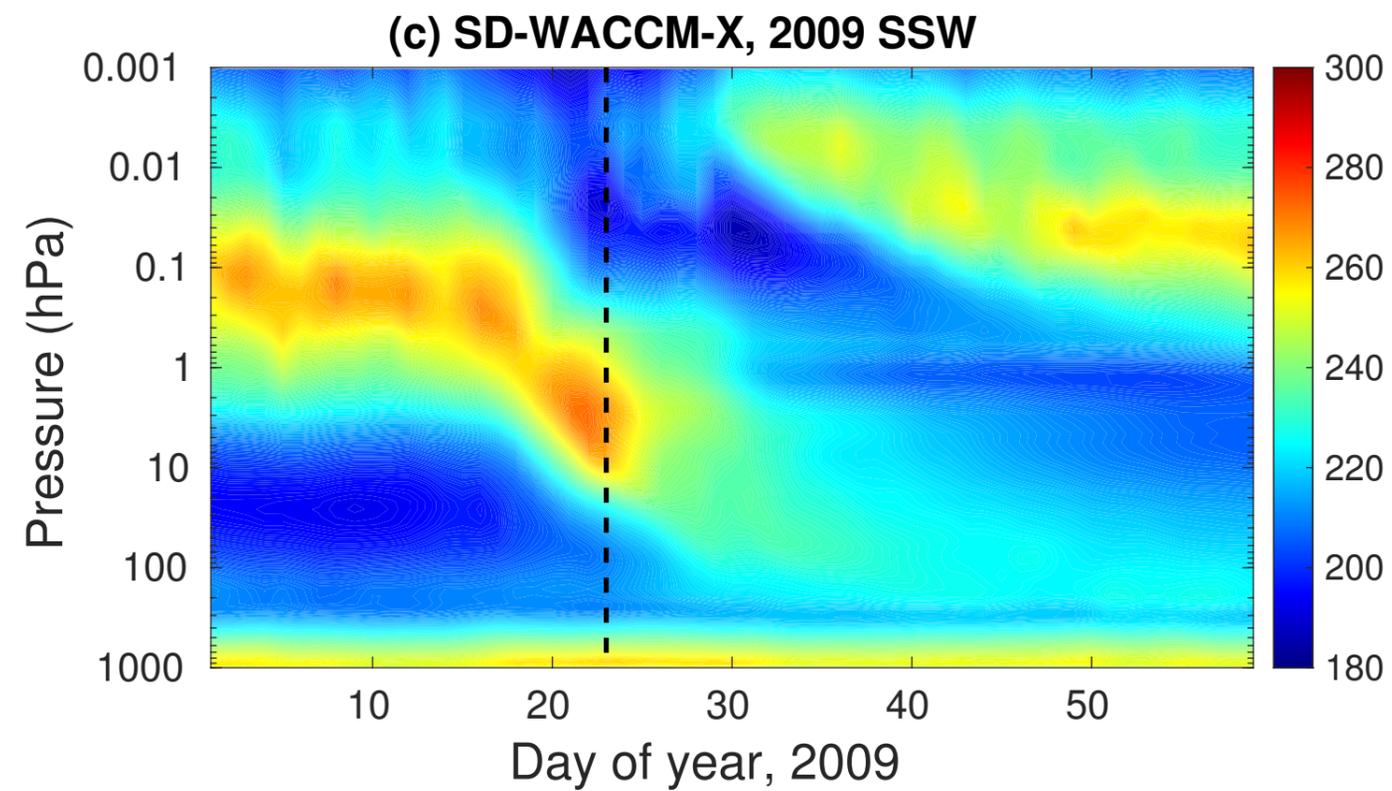
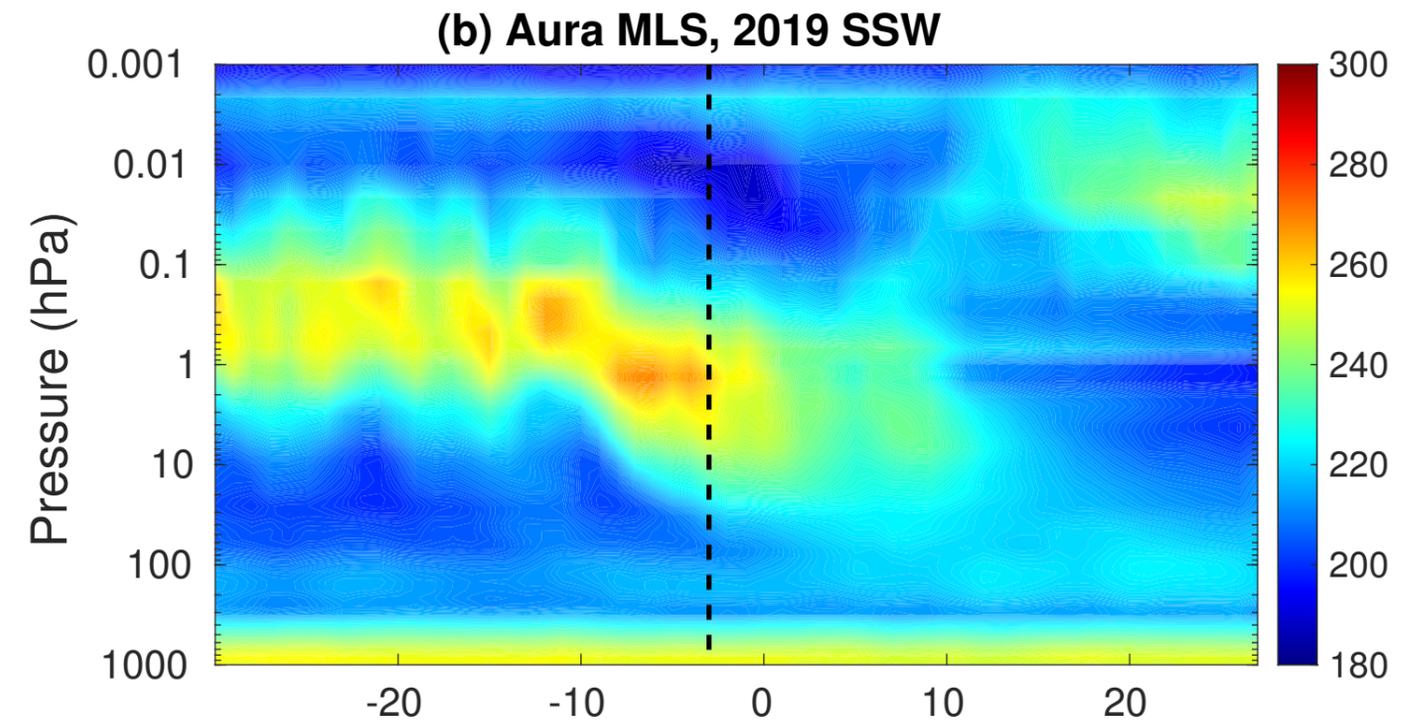
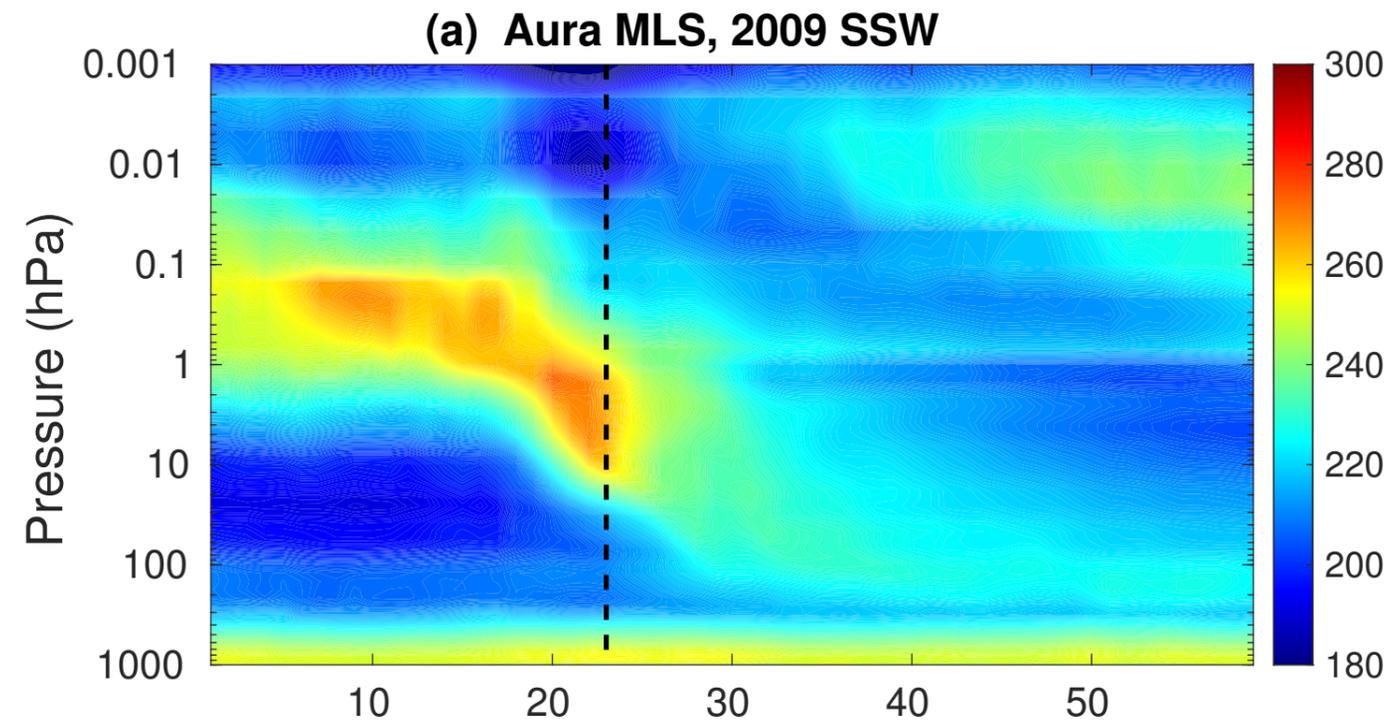


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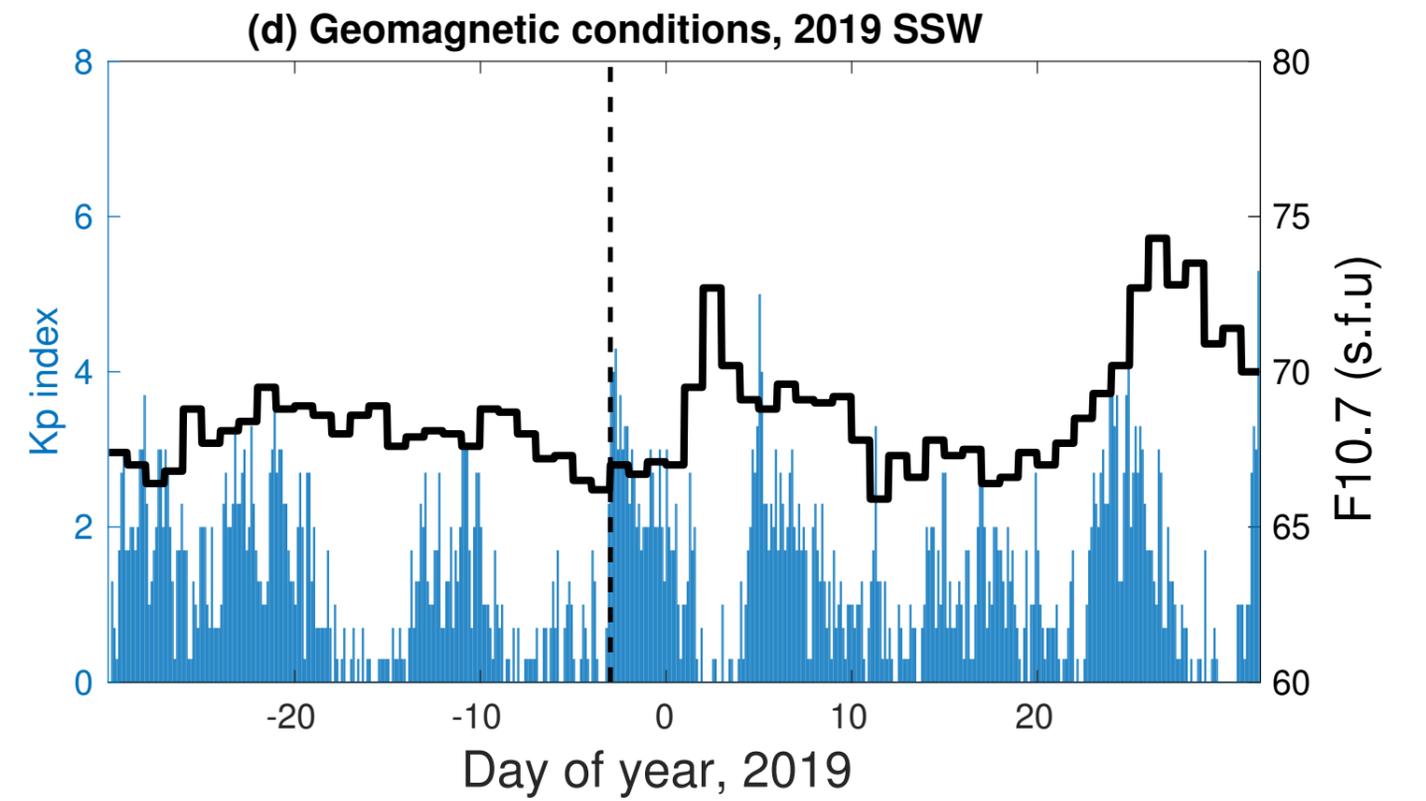
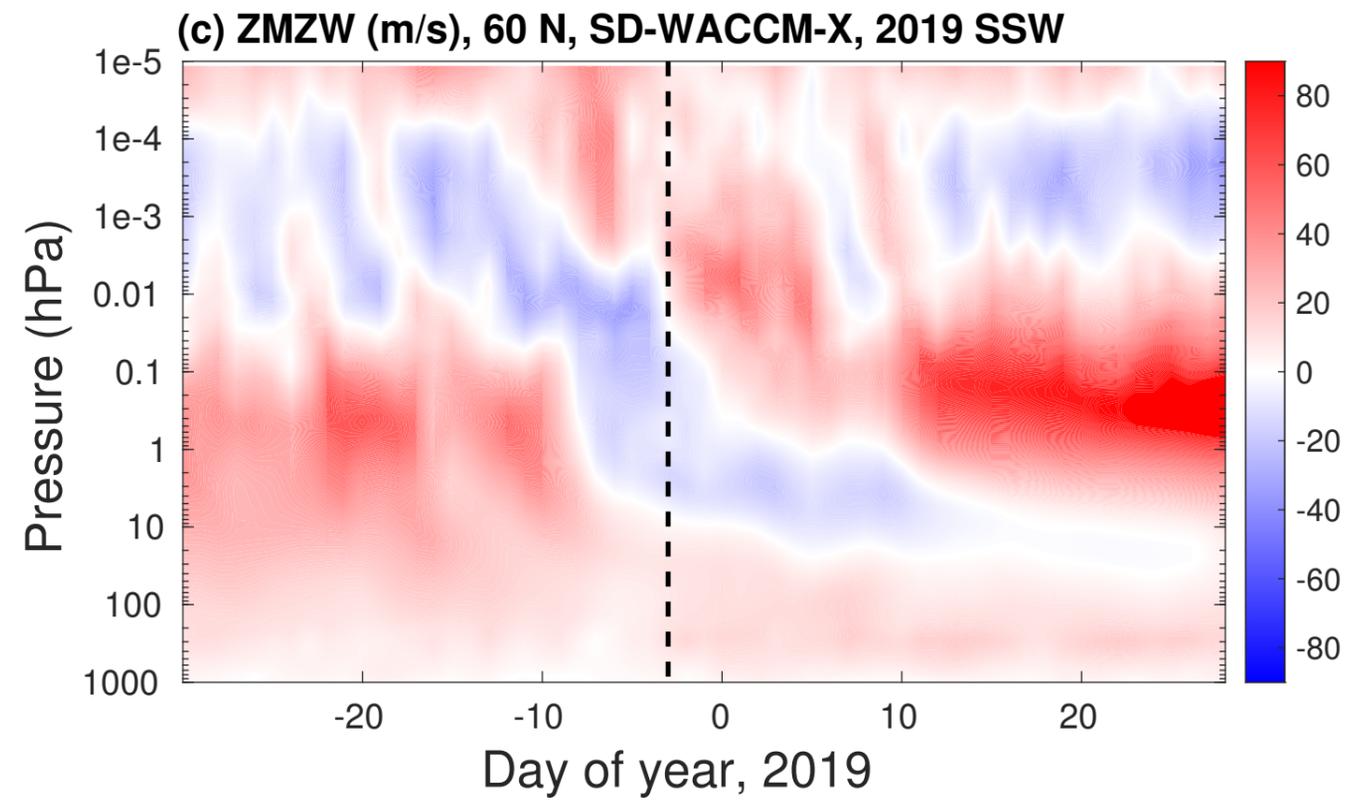
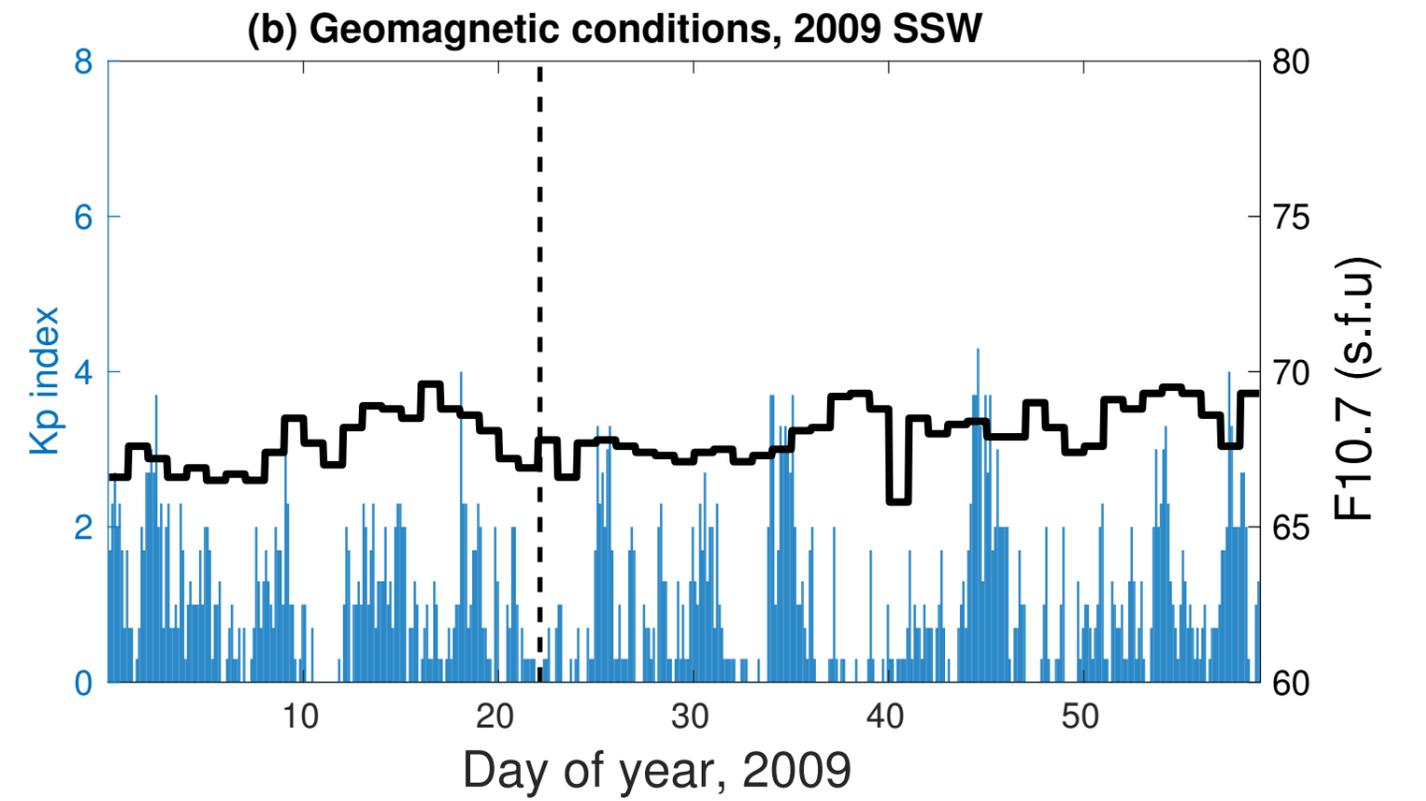
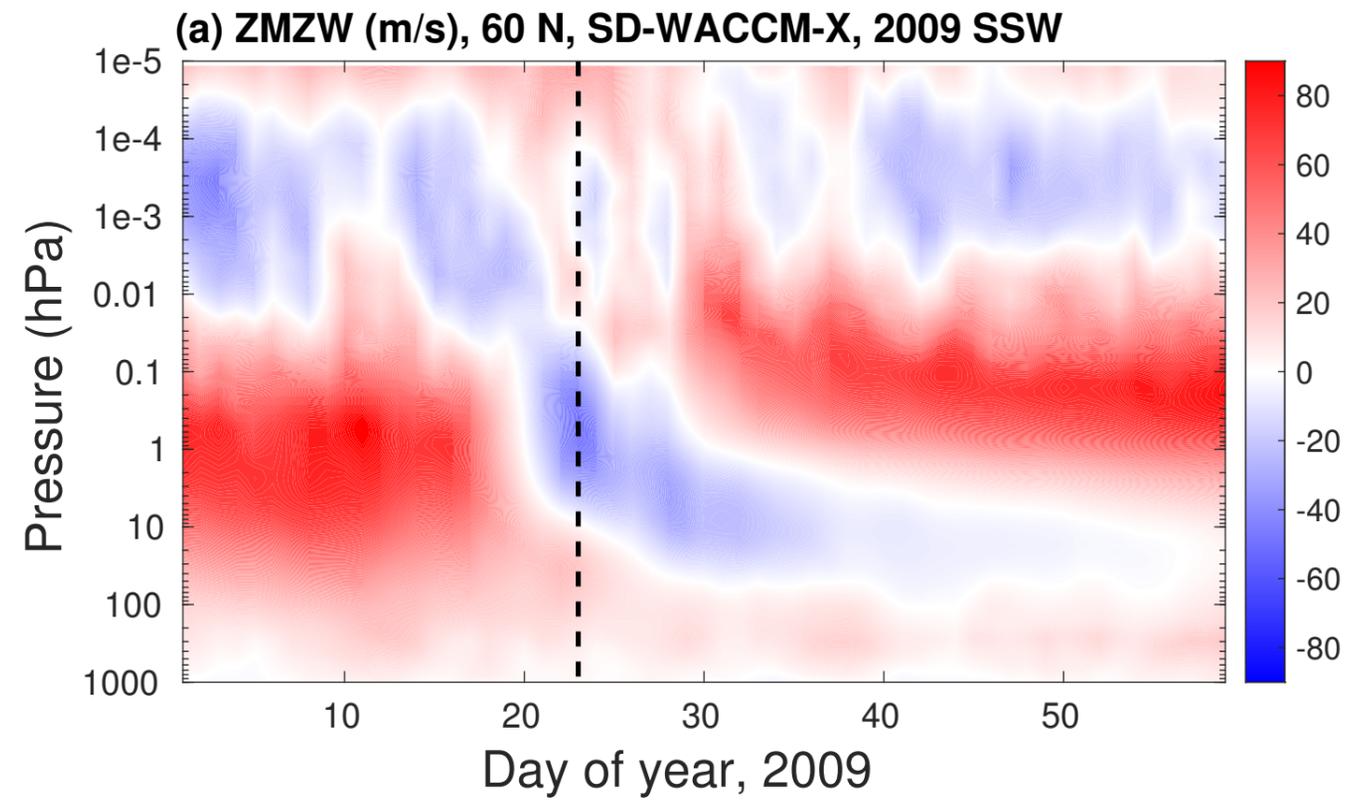
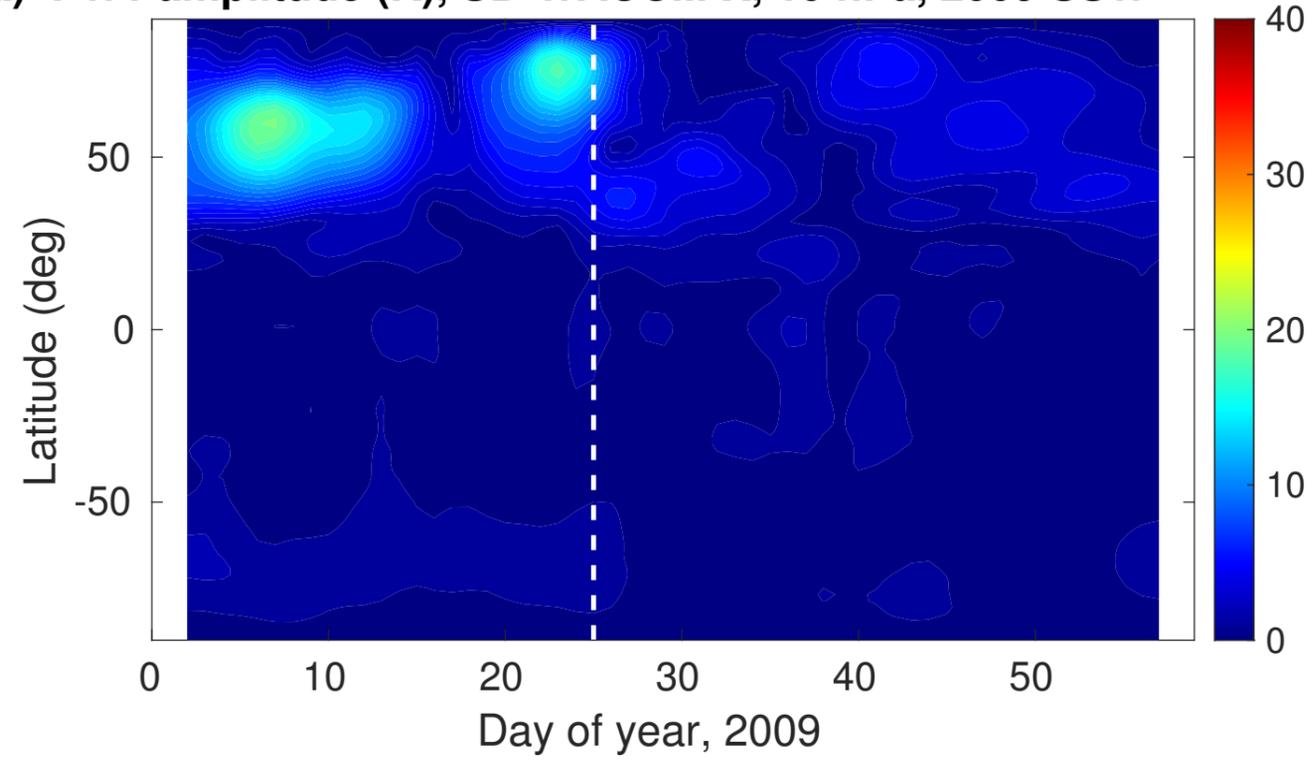
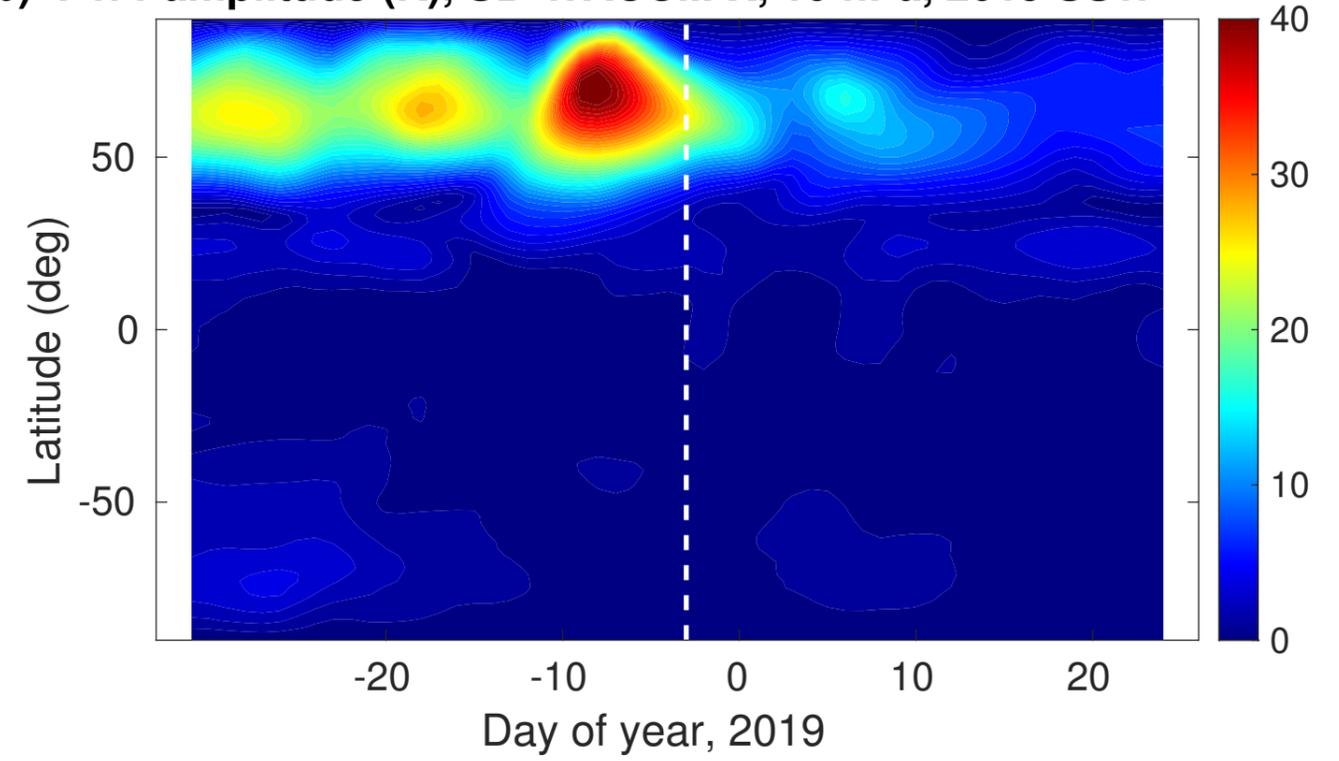


Figure 3.

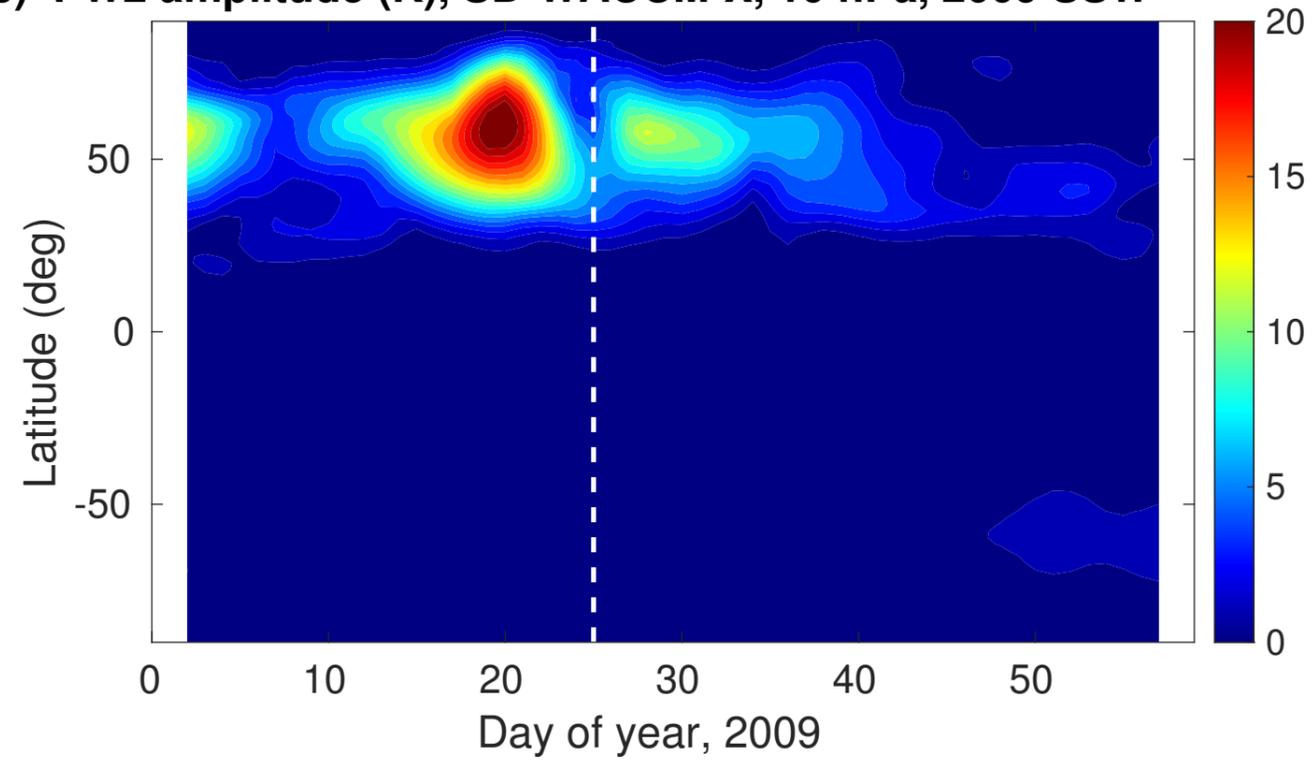
(a) PW1 amplitude (K), SD-WACCM-X, 10 hPa, 2009 SSW



(b) PW1 amplitude (K), SD-WACCM-X, 10 hPa, 2019 SSW



(c) PW2 amplitude (K), SD-WACCM-X, 10 hPa, 2009 SSW



(d) PW2 amplitude (K), SD-WACCM-X, 10 hPa, 2019 SSW

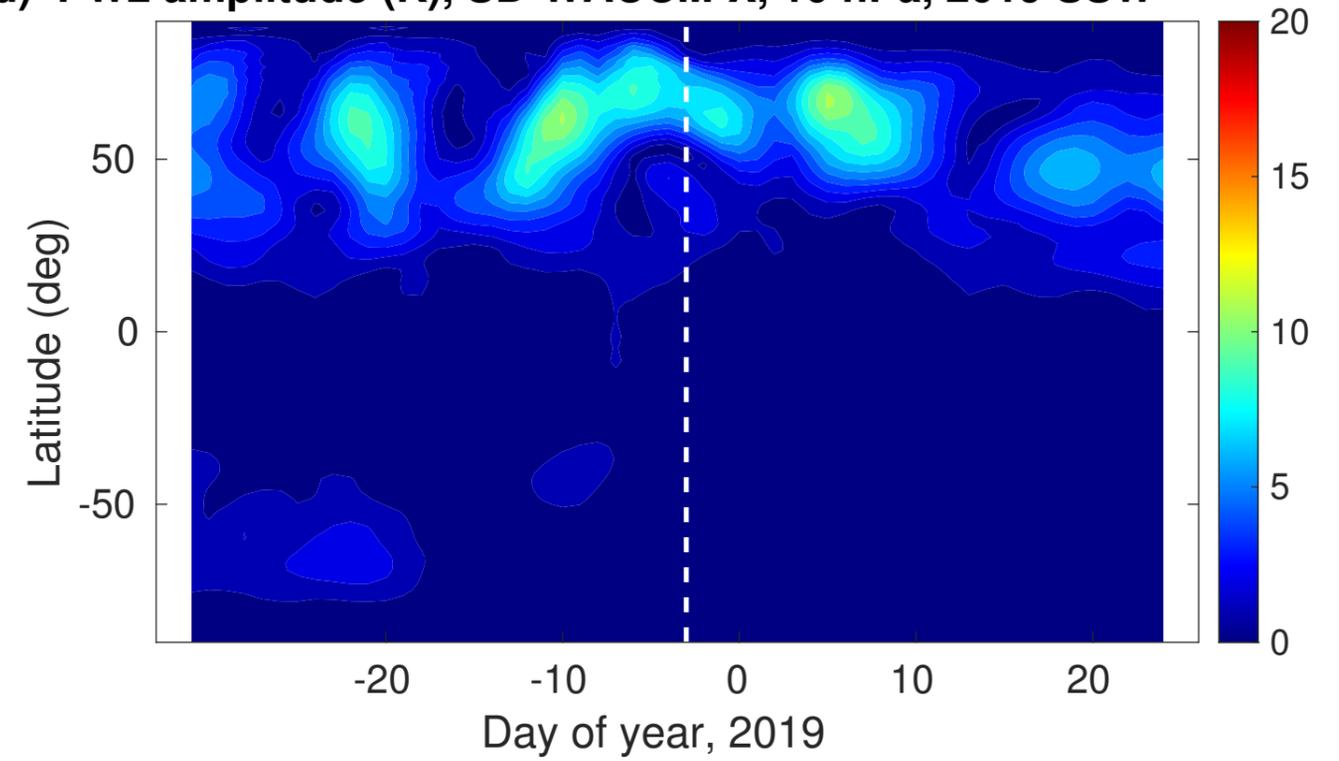
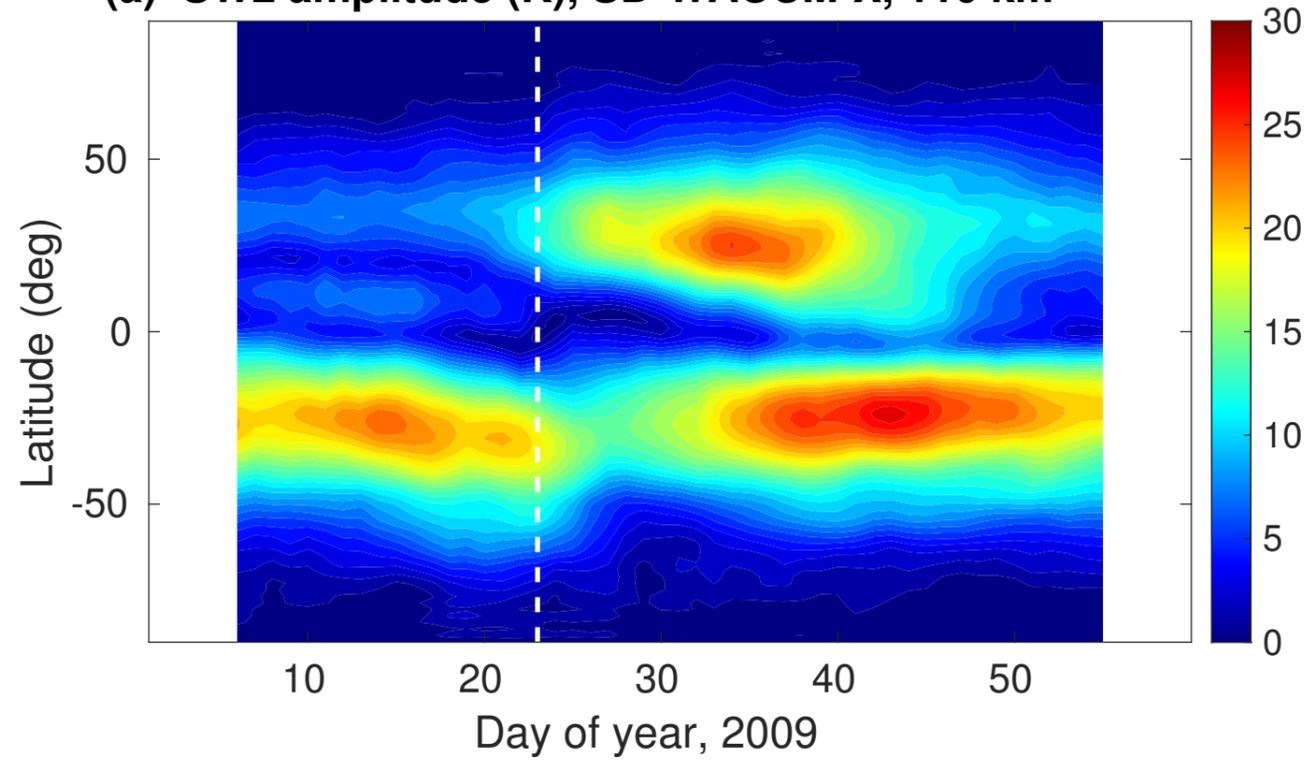
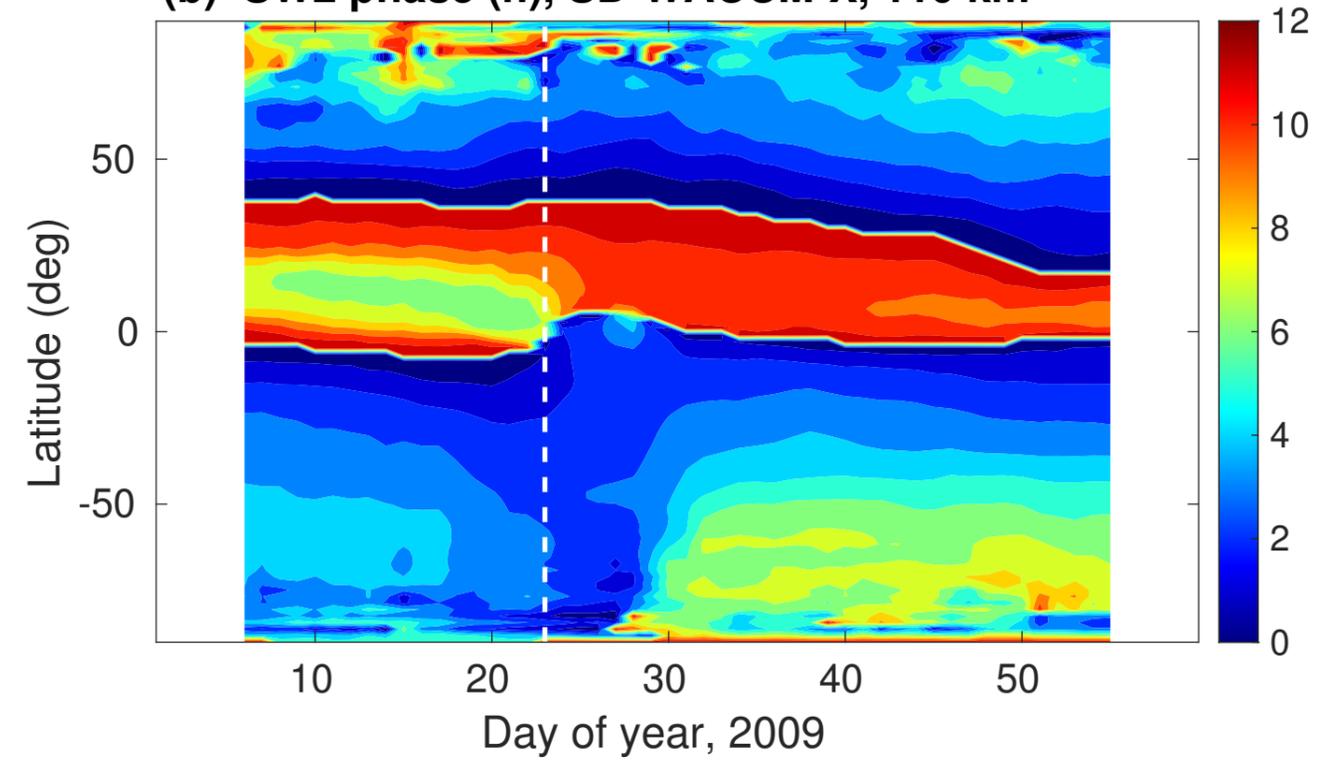


Figure 4.

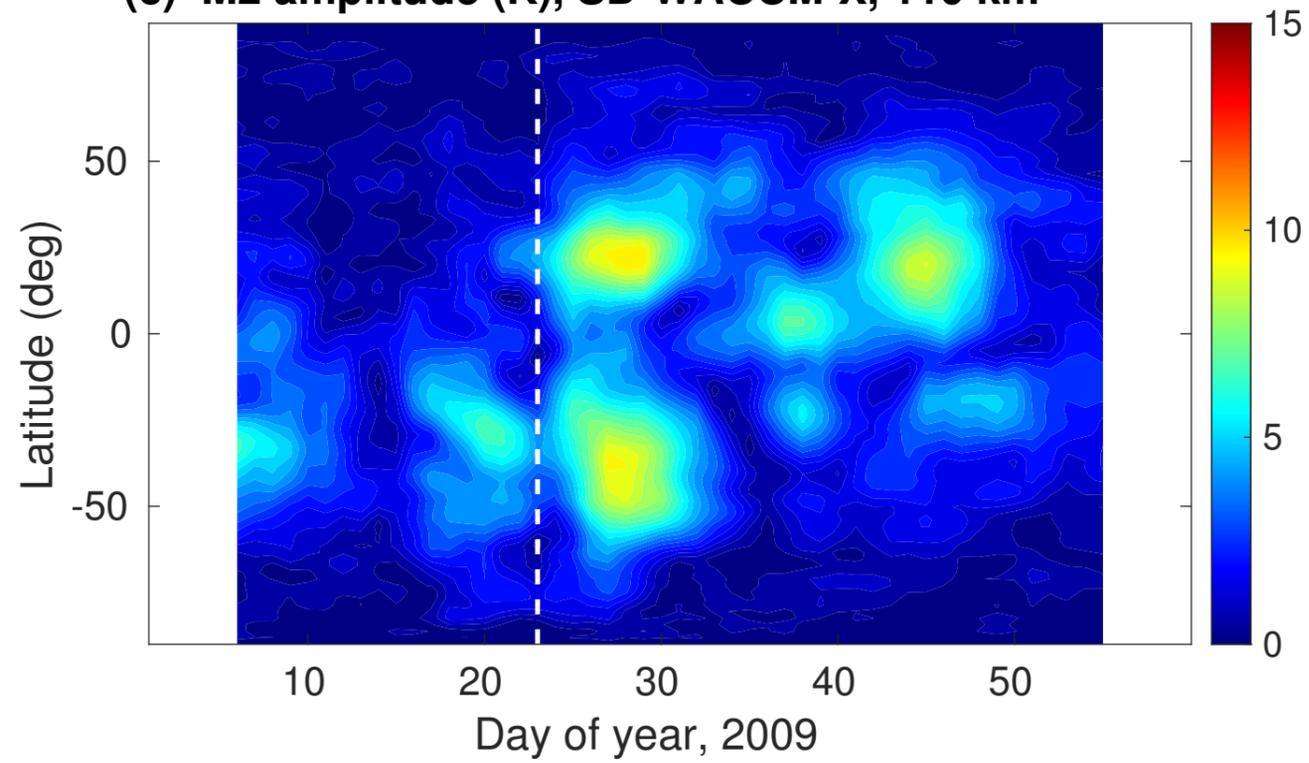
(a) SW2 amplitude (K), SD-WACCM-X, 110 km



(b) SW2 phase (h), SD-WACCM-X, 110 km



(c) M2 amplitude (K), SD-WACCM-X, 110 km



(d) M2 phase (h), SD-WACCM-X, 110 km

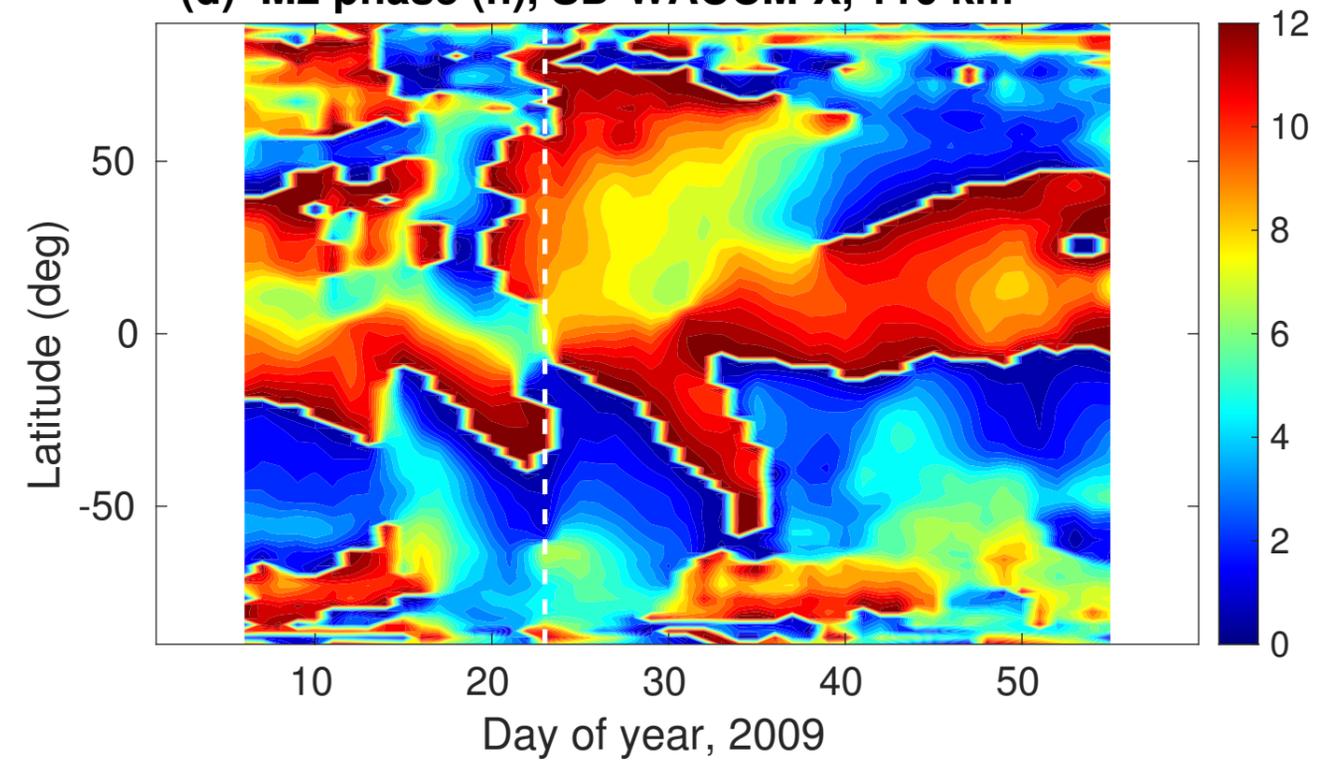
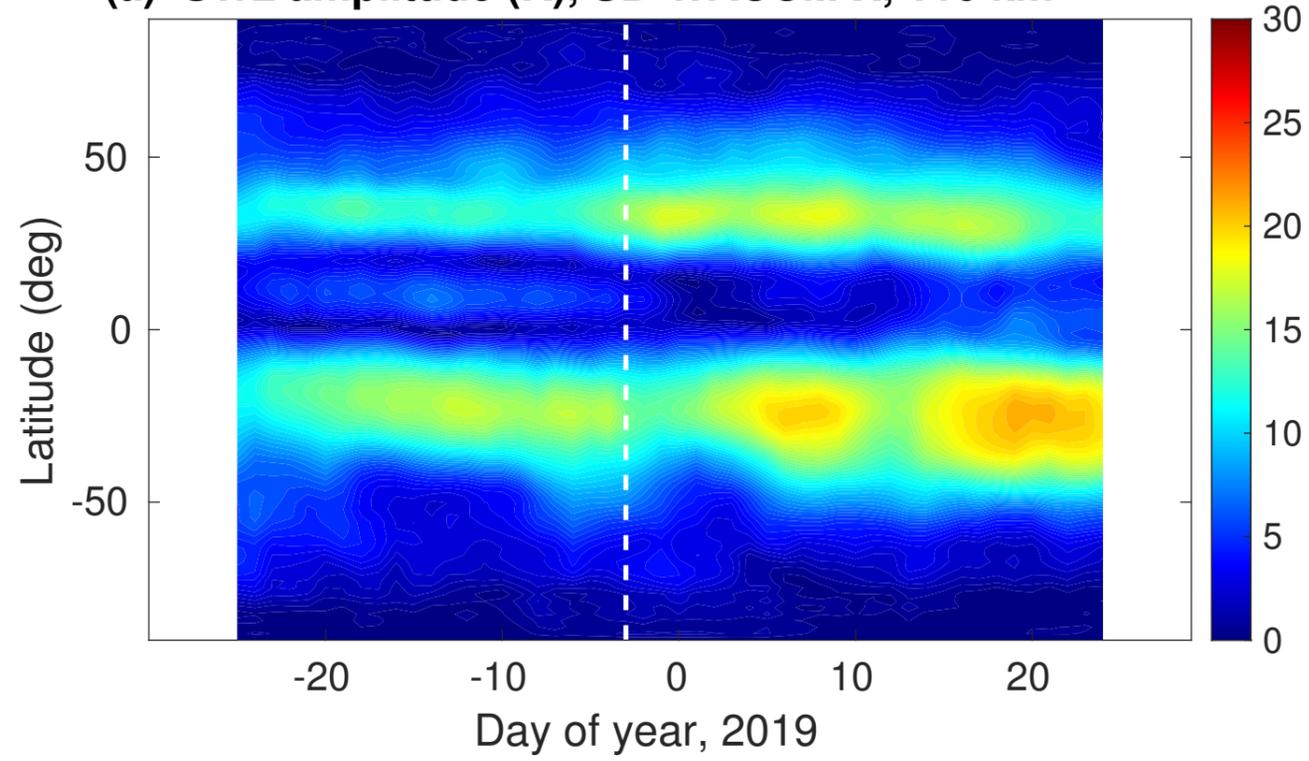
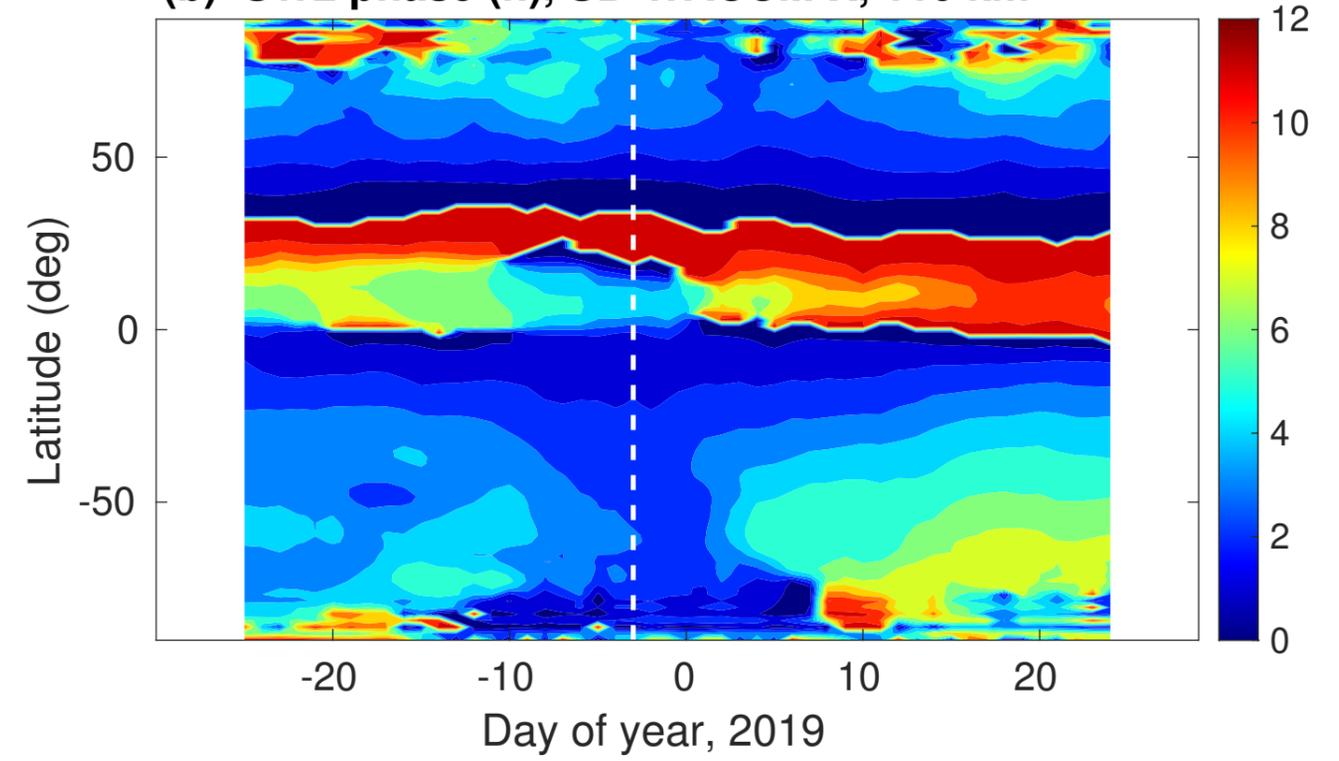


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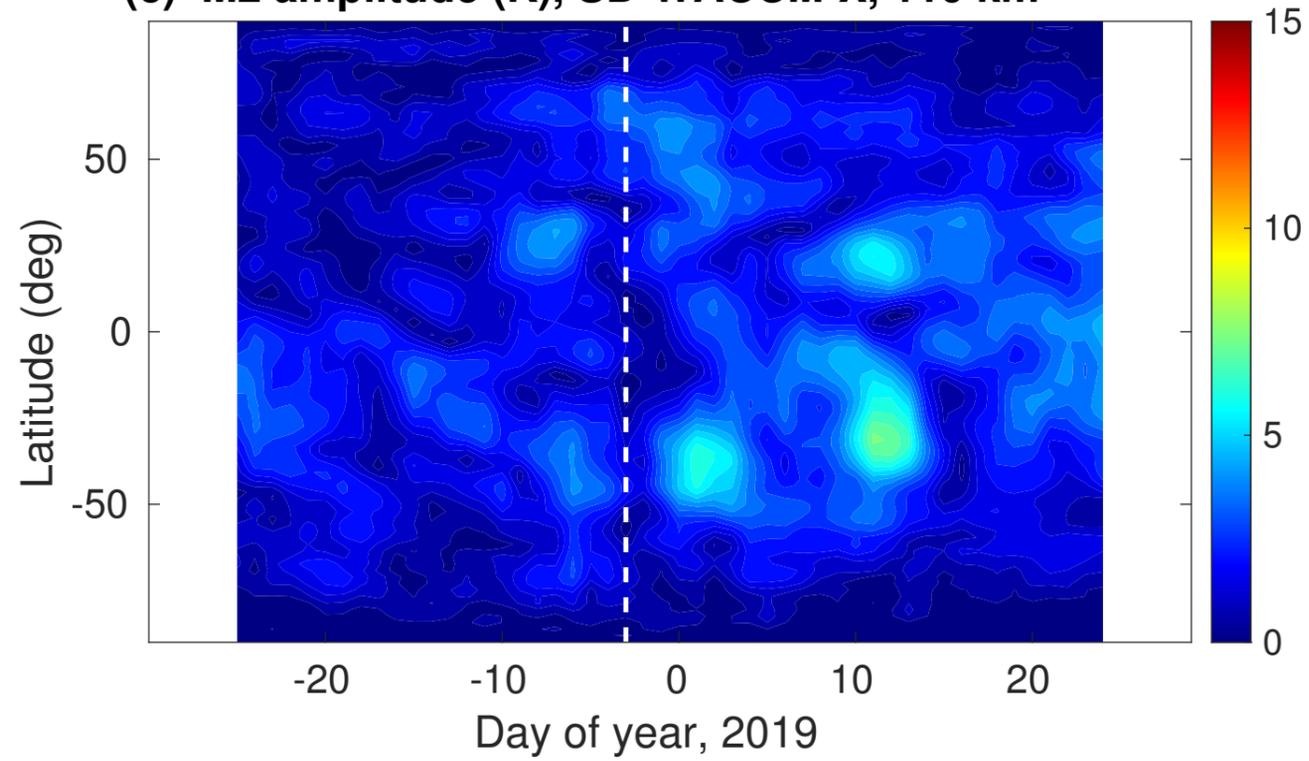
(a) SW2 amplitude (K), SD-WACCM-X, 110 km



(b) SW2 phase (h), SD-WACCM-X, 110 km



(c) M2 amplitude (K), SD-WACCM-X, 110 km



(d) M2 phase (h), SD-WACCM-X, 110 km

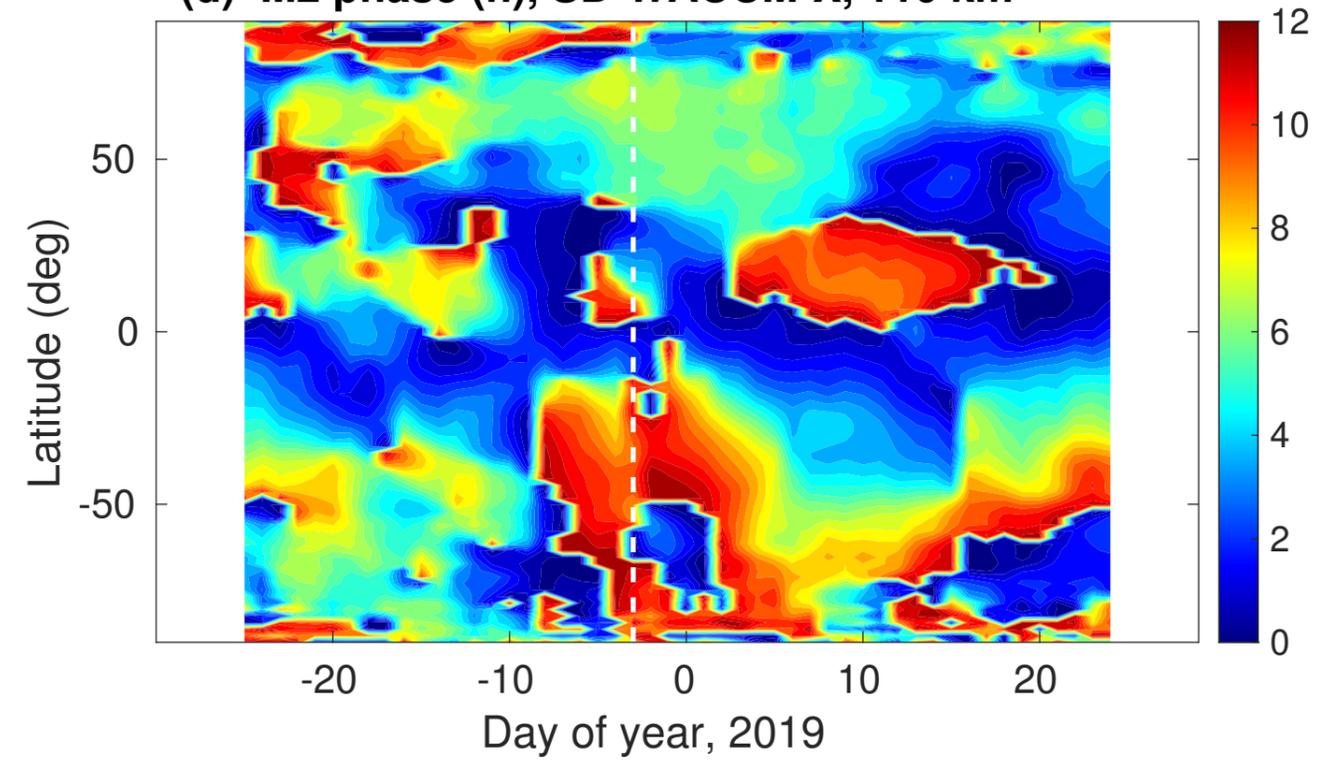
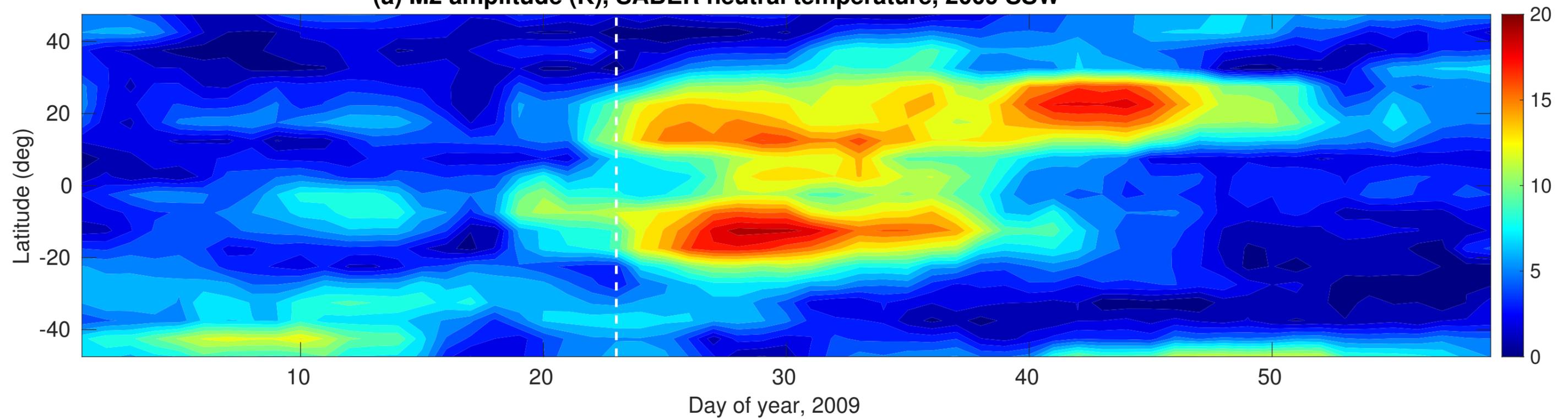


Figure 6.

(a) M2 amplitude (K), SABER neutral temperature, 2009 SSW



(b) M2 amplitude (K), SABER neutral temperature, 2019 SSW

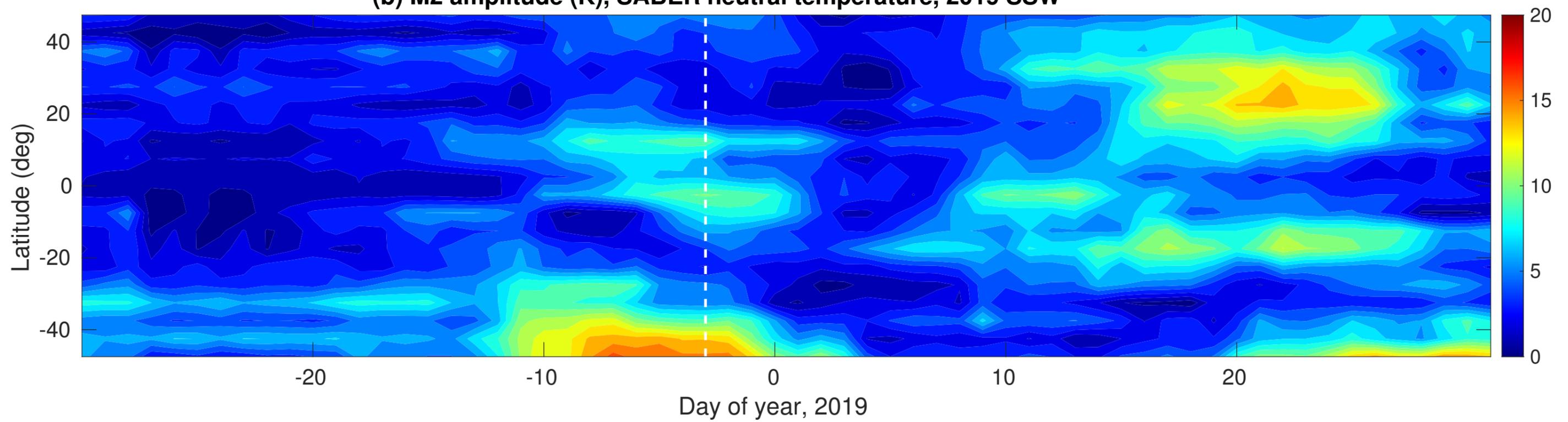


Figure 7.

Δ TEC over Europe at 12 UT

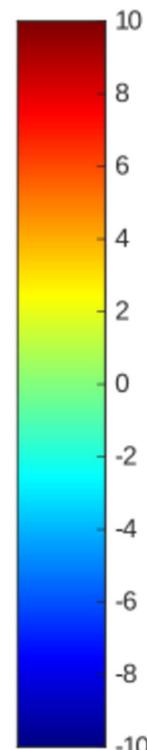
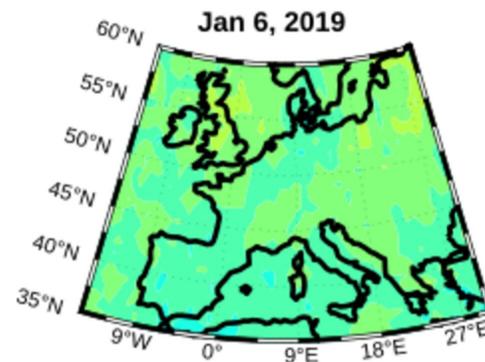
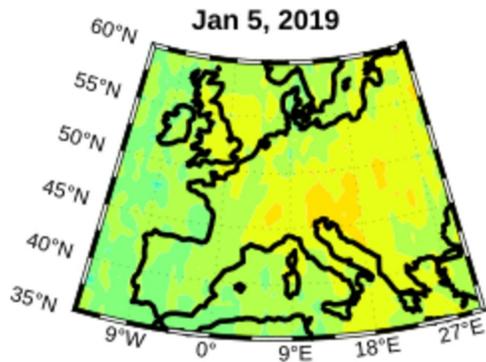
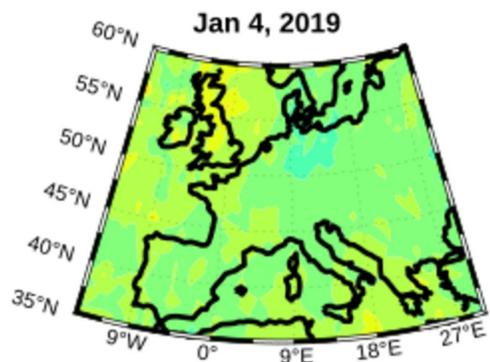
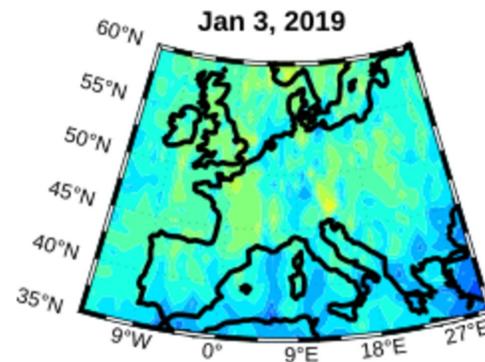
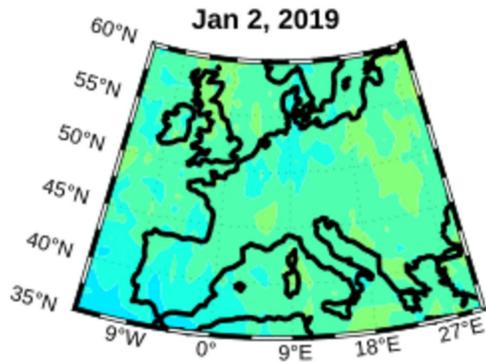
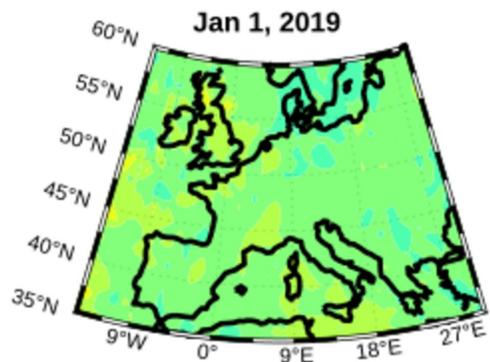
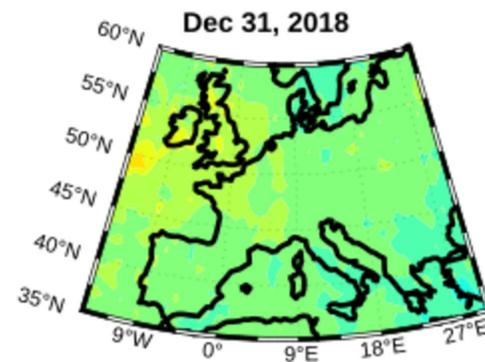
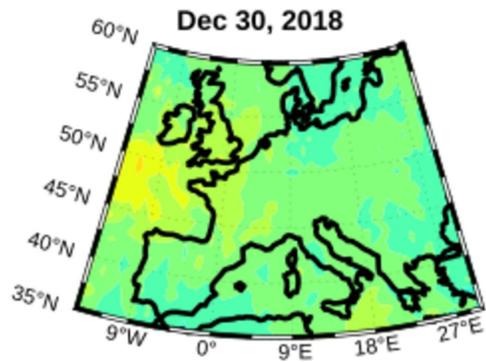
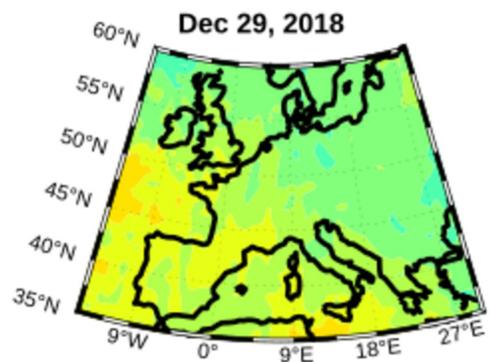
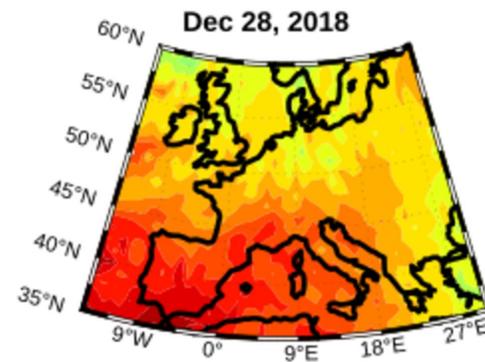
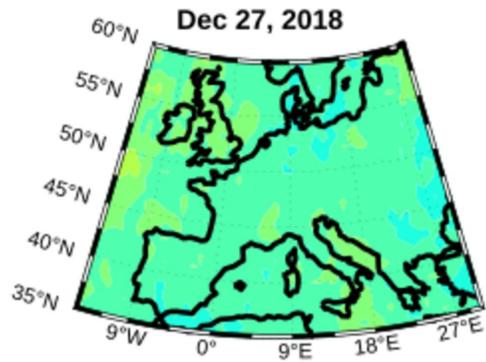
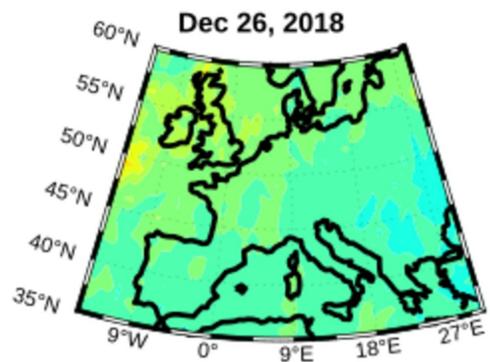
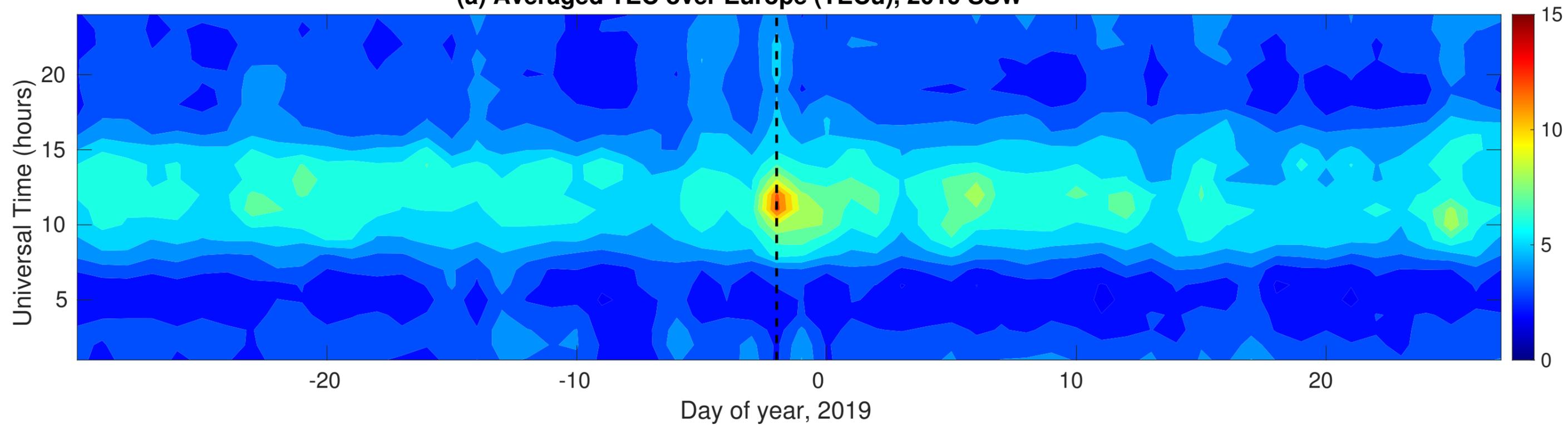


Figure 8.

(a) Averaged TEC over Europe (TECu), 2019 SSW



(b) Δ TEC over Europe (TECu), 2019 SSW

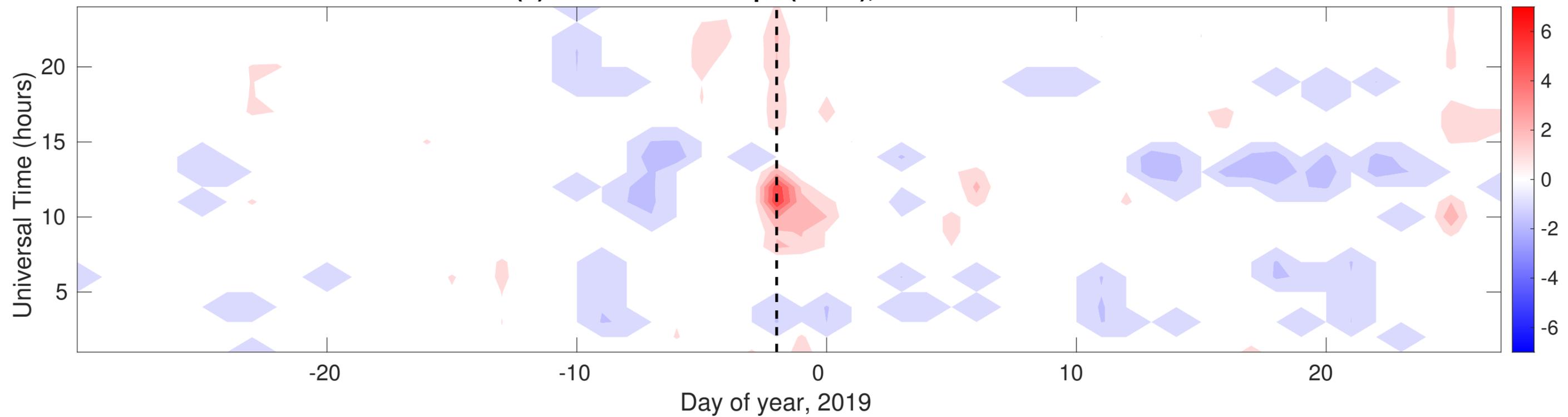
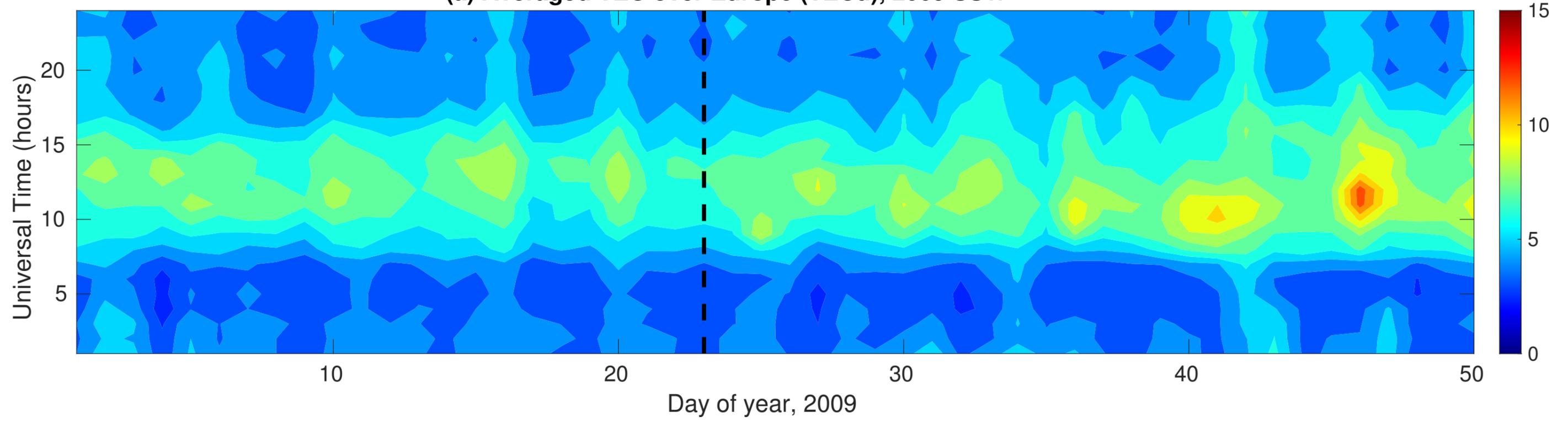


Figure 9.

(a) Averaged TEC over Europe (TECu), 2009 SSW



(b) Δ TEC over Europe (TECu), 2009 SSW

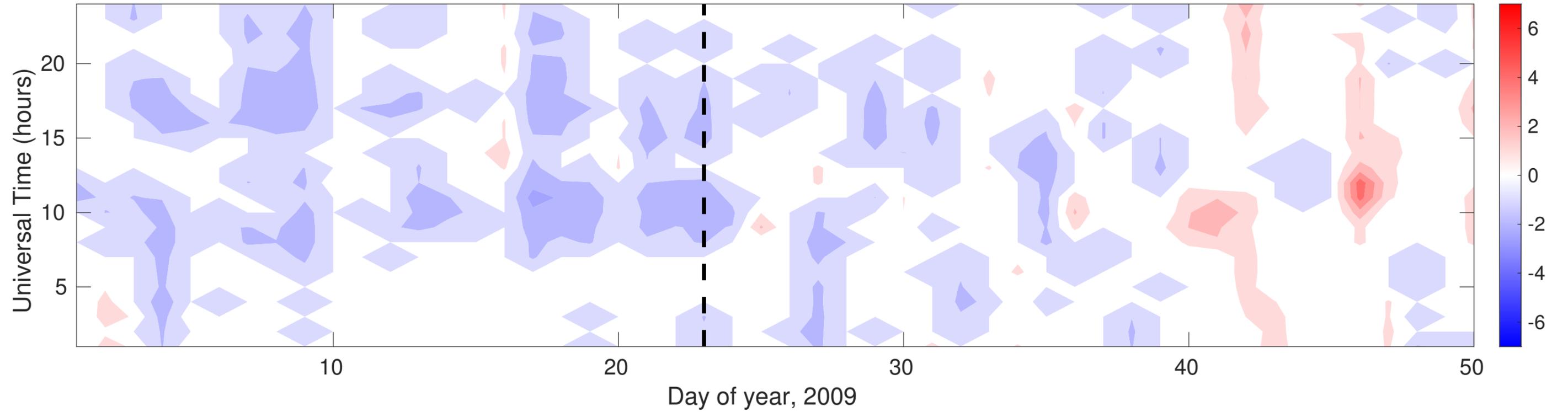
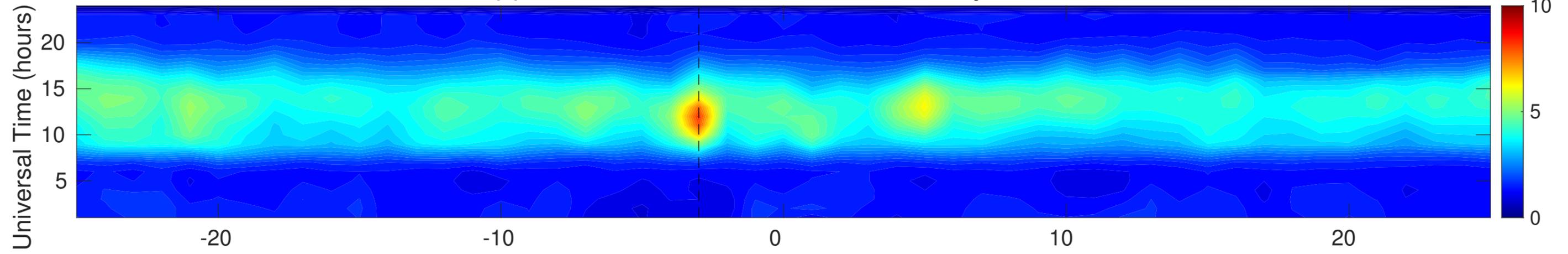
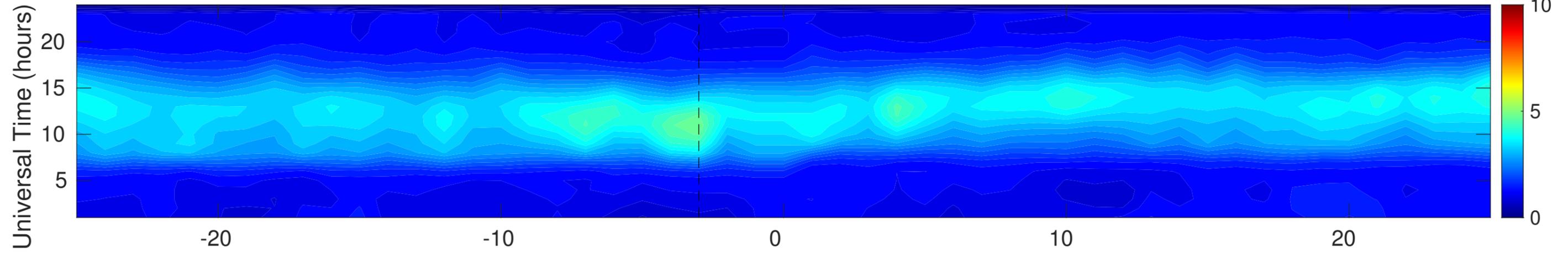


Figure 10.

(a) TEC from TIE-GCM, 2019 SSW, S1 setup



(b) TEC from TIE-GCM, 2019 SSW, S2 setup



(c) TEC from TIE-GCM, 2019 SSW, S1-S2

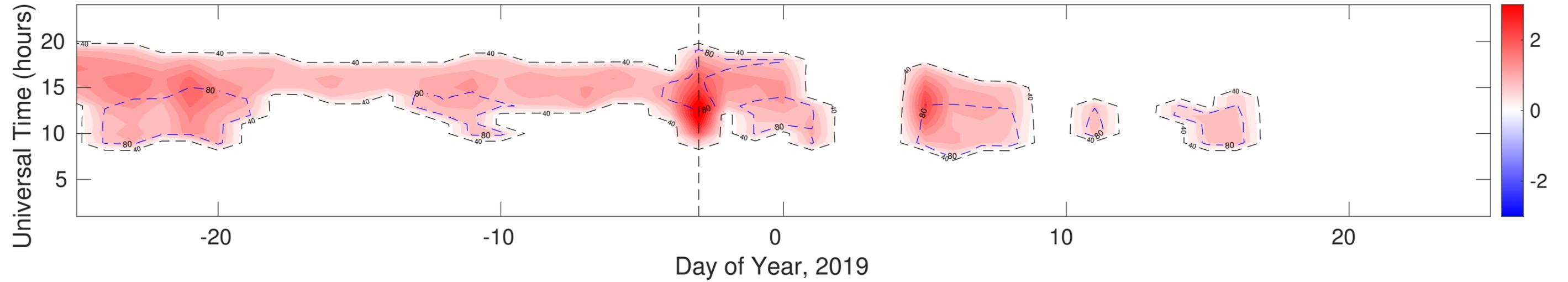
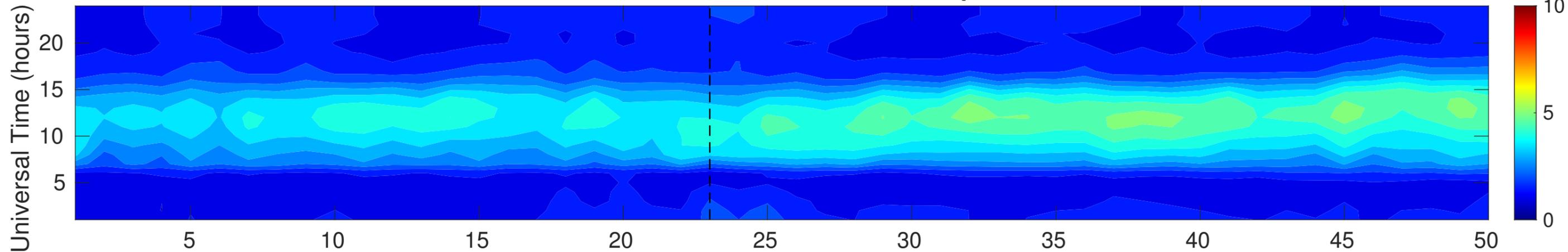
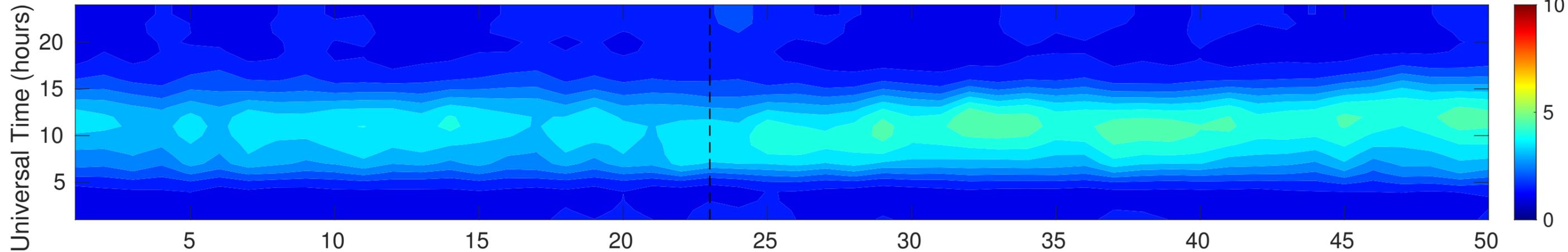


Figure 11.

TEC from TIE-GCM, 2009 SSW, S1 setup



TEC from TIE-GCM, 2009 SSW, S2 setup



(c) TEC from TIE-GCM, 2009 SSW, S1-S2

