

1 **Classification of the Below-Cloud Mixing State Over**  
2 **the Southern Ocean Using In-Situ and**  
3 **Remotely-Sensed Measurements**

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7 **Key Points:**

- 8 • A new, correlation-based method determines the mixing state of the below-  
9 cloud layer with an accuracy of 76%.  
10 • The below-cloud layer is only ever poorly-mixed when winds are below 8 m  
11 s<sup>-1</sup> and the near-surface atmosphere is neutrally stable.  
12 • Sea spray particles are available to low-level cloud over the Southern Ocean  
13 more than 80% of the time throughout austral summer.

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**Abstract**

We demonstrate that the relationship between the abundance of particulate surface area observed at sea-level and measurements of backscattered light by a ceilometer can be used to classify the mixing state of the atmospheric layer beneath the lowest observed cloud, where the relationship is defined by the Spearman Rank correlation. The accuracy of this correlation-based method was compared to two methods of detecting boundary layer decoupling based on radiosonde measurements. An optimized version of the new methodology correctly determined the mixing state of the below-cloud layer for  $76 \pm 4\%$  of the radiosondes available for comparison. Further, it was more accurate than an alternative ground-based metric used to determine the below-cloud mixing state. For the majority of the time series in which the correlation analysis could be applied, the below-cloud boundary layer was well-mixed (54%), or else fog was present (27%), which indicated that aerosol particles observed at sea-level often have a direct pathway into low-cloud (81%). In the remaining analysis period, the near-surface atmospheric layer was stable and the atmospheric layer near the ocean surface was decoupled from the overlying cloud (19%). Forecasts from the Antarctic Mesoscale Prediction System also support our findings, showing that conditions that mix aerosol particles from the ocean surface to the lowest observed cloud occur 84% of the time over the open Southern Ocean. As a result, aerosol particles measured near sea-level are often tightly coupled to low-cloud formation over the Southern Ocean, highlighting the utility of shipborne aerosol observations in the region.

**Plain Language Summary**

Particles suspended in the atmosphere (aerosol) act as seeds for cloud droplet formation. The abundance of such particles directly influences the opacity of clouds, while their physical and chemical characteristics govern if and when those cloud droplets freeze. As a result, both the amount of solar radiation a cloud can reflect and the temperature of waters below are sensitive to the quantity and type of particles available to the cloud. We present a new methodology for understanding the conditions in which low-level clouds have direct access to the large and diverse reservoir of particles in the surface layer. We find that meteorological conditions which transfer particles from sea-level to low-level cloud are satisfied up to 81% of the time over the Southern Ocean. This suggests that the particles we observe near the surface almost always play a significant role in the formation of low-level cloud.

**1 Introduction**

The balance of incoming and outgoing shortwave radiation is mediated by the presence of clouds, and to a lesser extent, aerosol particles. While aerosol particles are significantly smaller than clouds, they are the seeds of cloud droplet formation (Pruppacher et al., 1998). The chemical and physical nature of these particles strongly determines both the ultimate phase of cloud droplets (e.g. liquid or ice) and the resulting size distribution of cloud droplets (Twomey, 1977). It therefore follows that errors in how these particles are represented within global climate models can cause significant climatological biases in the radiative balance. In particular, it has been observed that uncertainties in predicting cloud phase leads to substantial biases in the radiation balance within the cold sector of Southern Ocean cyclones (Bodas-Salcedo et al., 2014). While the abundant cyclones of the Southern Ocean occur solely as a function of favourable synoptic conditions (Irving et al., 2010), a global climate model's predictions of cloud phase in the cold sector of Southern Ocean cyclones (Vergara-Temprado et al., 2018), and in the wider Southern Ocean (Schuddeboom et al., 2019), is extremely sensitive to the properties of

64 particles in the underlying boundary layer. In addition, recent observations directly  
65 show that the large biases between modeled and observed outgoing shortwave ra-  
66 diation (Trenberth & Fasullo, 2010) are related to errors in how low-level cloud are  
67 represented within global climate models (Kuma et al., 2020). Improving our under-  
68 standing of the conditions in which particles can reach, and form, low-level cloud is  
69 therefore crucial to understanding the radiation balance over the Southern Ocean.

70 As wind speeds have increased over the Southern Ocean (Young et al., 2011;  
71 Hande et al., 2012a), there is significant interest in how naturally-produced particles  
72 impact cloud formation and the optical properties of the resultant clouds (McCoy  
73 et al., 2015), and whether this interaction represents a substantial climate feed-  
74 back (Korhonen et al., 2010). It is well-known that increasing the population of  
75 cloud condensation nuclei (CCN) directly increases the opacity of the overlying  
76 cloud (Twomey, 1977). Increases in wind speed over the open ocean will enhance  
77 the flux of sea spray particles (SSPs) from breaking waves (Hartery et al., 2020). In  
78 most regions of the Southern Ocean SSPs are the only local source of ice-nucleating  
79 particles (INPs) (DeMott et al., 2016), a region almost entirely devoid of such par-  
80 ticles (Bigg & Hopwood, 1963). While other, more potent, INPs like dust particles  
81 may be entrained into the boundary layer in specific seas (e.g. coastal seas near  
82 Patagonia), ice nucleating particles collected on Southern Ocean voyages have a  
83 much weaker surface activity than dust particles, which reflects the predominant  
84 abundance of sea spray (McCluskey et al., 2018). These particles can have a sub-  
85 stantial influence on the radiative and physical properties of the resulting cloud. Not  
86 only are ice clouds much less opaque (Hu et al., 2010), they are much more likely to  
87 precipitate (Borys et al., 2003). Thus, changes in the abundance of SSPs may have  
88 significant impacts on cloud radiative properties.

89 One of the challenges in unravelling aerosol–cloud interactions over the South-  
90 ern Ocean is that the region is frequently covered in cloud (80% of the time; Haynes  
91 et al. (2011)), which leads to difficulties in monitoring low-level clouds (McErlich et  
92 al., 2021) and the structure of boundary layer below (Hande et al., 2015). While in  
93 situ observational records of radiosondes from Macquarie Island provide rich data  
94 on the thermodynamic structure of the Southern Ocean boundary layer (Hande et  
95 al., 2012b), a lack of accompanying observations of CCN, INPs, and in situ micro-  
96 physical properties of low-level cloud leaves a gap in our understanding of how these  
97 particles interact with cloud over the Southern Ocean. Previous research, such as  
98 the dedicated ACE-1 (Russell et al., 1998), SOCEX (Boers et al., 1998), HIPPO  
99 (Wofsy, 2011) and more recently SOCRATES (McFarquhar et al., 2020) campaigns  
100 have used aircraft observations to bridge this knowledge gap. However, aircraft can  
101 only fly in a limited range of conditions, as the strong vertical wind shear and icing  
102 conditions present within boundary layer cloud poses a significant threat. By con-  
103 trast, ship-based measurements can be made in nearly all conditions. Here, we use  
104 measurements on the R/V *Tangaroa* during a voyage to the Ross Sea in the austral  
105 summer of 2018 to establish conditions in which particles near the surface are tur-  
106 bulently mixed to cloud base (Kremser et al., 2020). Establishing conditions when  
107 sea-level measurements are relevant to cloud will enable future research to better ex-  
108 ploit sea-level measurements in aerosol–cloud interaction studies, and adds value to  
109 the growing catalogue of near-surface measurements available from recent voyages.

## 110 2 Measurements

111 Over the course of a voyage between New Zealand and the Ross Sea, air was  
112 drawn from the mast of the R/V *Tangaroa* ( $\sim 20$  m above sea level “a.s.l.”) to a  
113 shipping container laboratory ( $\sim 2$  m a.s.l.) via 40 m of conductive hose. Within the  
114 laboratory, a passive cavity aerosol spectrometer probe (PCASP-100X; Droplet Mea-  
115 surement Technologies) and a differential mobility particle sizer (DMPS, TSI) mea-

116 sured the ambient concentration of particles suspended in the atmosphere (Kremser  
 117 et al., 2020). The PCASP measured the number concentration size spectra of par-  
 118 ticles suspended in the boundary layer in 30 size bins (0.1–3.0  $\mu\text{m}$ ) every minute.  
 119 The DMPS measured the number concentration size spectra in the size range 0.02–  
 120 0.3  $\mu\text{m}$  every 10 minutes. Following Modini et al. (2015) and Quinn et al. (2017), we  
 121 fit three lognormal size distributions to estimate the average diameter and number  
 122 concentration of Aitken, accumulation and coarse mode particulate. With rare ex-  
 123 ception, coarse mode particulate is almost entirely composed of sea spray particles  
 124 (SSPs) in the marine environment (Modini et al., 2015; Quinn et al., 2017); hence,  
 125 we will refer to the coarse mode as the SSP mode throughout the remainder of this  
 126 work. The PCASP was used exclusively to estimate the average size and abun-  
 127 dance of SSPs, while the DMPS was used for the Aitken and accumulation mode  
 128 particles. When data from the DMPS were not available, measurements from the  
 129 PCASP were used to constrain the abundance and size of accumulation mode parti-  
 130 cles. Further details on sampling set-up and analysis, including correction factors for  
 131 losses through the sampling line and methods for handling contamination from ship  
 132 exhaust, are described in Hartery et al. (2020) and Kremser et al. (2020). In par-  
 133 allel to the size-resolved particle concentration spectra generated by the SMPS and  
 134 PCASP, the total number of cloud condensation nuclei (CCN) was measured using  
 135 a CCN counter (CCNC-100; Droplet Measurement Technologies). The CCN counter  
 136 sampled from the same sampling conduit that drew ambient air to the PCASP and  
 137 DMPS. A measurement of the average number of ambient CCN was made twice an  
 138 hour at intervals of 0.1% supersaturation between 0.2–1.0%.

139 A ceilometer (CHM-15K; Lufft) transmitted pulses of laser light at a wave-  
 140 length of 1064 nm and recorded the total power of light scattered back to the  
 141 ceilometer per laser pulse from different levels of the atmosphere. Measurements  
 142 were recorded at a temporal resolution of one record per minute and a vertical reso-  
 143 lution of 15 m. For each record, the instrument also estimated the cloud base height,  
 144  $z_{CBH}$ . A raw quality control flag provided by the instrument was used to screen for  
 145 field-of-view contamination from fog or residual precipitation on the outer optical  
 146 window. A micro-rain radar (MRR-2; Metek) operated in proximity was also used to  
 147 detect and screen for precipitation events.

148 An Automated Weather Station (AWS) provided by New Zealand MetService  
 149 was positioned above the bridge of the R/V *Tangaroa* at 22.5 m. Relevant mea-  
 150 surements included ambient pressure, air temperature, relative humidity, long and  
 151 shortwave radiation fluxes, wind speed, and wind direction. Measurements from the  
 152 AWS were corrected to a height of 10 m according to the COARE 3.5 bulk-flux al-  
 153 gorithms (Edson et al., 2013) as detailed in Hartery et al. (2020). The bulk seawater  
 154 temperature was measured at a depth of 5.5 m below sea level with a thermistor  
 155 (SBE38; Sea-Bird Scientific). We also used the COARE 3.5 bulk-flux algorithms  
 156 (Edson et al., 2013) to calculate the sea skin temperature from the bulk tempera-  
 157 ture, accounting for long and shortwave fluxes (Edson et al., 2013).

158 Fifty-seven meteorological balloons were launched during the voyage. The ra-  
 159 diosondes (iMet-ABx; InterMet) recorded pressure, relative humidity, temperature  
 160 and wind speed. In quality control, two of the radiosondes were found to have a  
 161 faulty relative humidity sensor and one had more than one faulty sensor, leaving 54  
 162 useful profiles of the boundary layer. The radiosondes were launched approximately  
 163 twice daily once the ship was further south than 60° S.

164 Regional meteorological forecasts were downloaded from the Antarctic  
 165 Mesoscale Prediction System (AMPS). AMPS initializes a new forecast every twelve  
 166 hours, with subsequent output provided every three hours. AMPS provides forecasts  
 167 within several nested spatial grids. However, only forecasts for the outermost spa-  
 168 tial grid, “domain 1,” were used as it was the only grid which fully contained the

169 ship track. Domain 1 has a horizontal resolution of 24 km and is a  $544 \times 412$  grid  
 170 centred on  $90^\circ\text{S}$ . AMPS uses the Mellor-Yamada-Janjić (MYJ) scheme, a 2.5-level  
 171 closure model of turbulence, to predict the behavior of the planetary boundary layer  
 172 (PBL). AMPS calculates the height of the PBL to be the height at which the turbu-  
 173 lent kinetic energy falls below a pre-determined threshold (Janjic, 2001). To allow  
 174 for a brief model spin-up, only forecasts between 3–12 hours were used (Jolly et al.,  
 175 2016).

### 176 3 Methods

#### 177 3.1 Classification of the Below-Cloud Layer

178 The suspended particle cross-sectional surface area,  $A$ , was calculated from the  
 179 number concentration size spectra measured by the PCASP:

$$180 \quad A(t) = \int \frac{dn(t, D_p)}{d \log D_p} \pi \left( \frac{D_p}{2} \right)^2 d \log D_p \quad (1)$$

181 Where  $D_p$  is the particle diameter,  $n$  is the partial concentration of particles, and  
 182  $t$  is time. Note that as this is a correlation-based study, a more exact treatment of  
 183 the interaction of particulate with light which accounts for both Mie and Rayleigh  
 184 scattering (e.g. Bohren and Huffman (1983)) is not strictly necessary. In addition,  
 185 such calculations would necessitate *a priori* information about particle composition  
 186 and morphology which were not available for this study. The geometric surface area  
 187 is dominated by the sea spray and accumulation mode particles (97%, on average;  
 188 Fig. 1d), which the PCASP can readily measure.

189 To classify the below-cloud layer mixing state, we calculated rolling Spearman  
 190 Rank correlation coefficients centred on each hour of observation between the sea-  
 191 level concentration of aerosol surface area,  $A(t)$ , and the background-corrected total  
 192 power of backscattered light received by the ceilometer,  $P_c$ . The Spearman Rank  
 193 correlation coefficient was used as non-linearities related to the two-way transmis-  
 194 sion of light through an atmospheric layer are likely; however, the Pearson moment  
 195 correlation coefficient produced qualitatively similar results. Before calculating the  
 196 correlation coefficients, four quality control measures were implemented to ensure  
 197 that the calculated correlation coefficients would be meaningful. First, the obser-  
 198 vations were screened based on the ceilometer’s quality control flag and the ship  
 199 contamination flag described in Kremser et al. (2020). Second, only backscattered  
 200 light retrieved from heights below the 10<sup>th</sup>-percentile of CBH were studied. This  
 201 step ensured that correlations between backscattered light from a given altitude and  
 202 sea-level particulate surface area resulted from co-variations of the abundance of  
 203 aerosol particles at the surface and the selected height and not from variations in the  
 204 presence of cloud droplets. Next, we performed a signal-to-noise analysis, where the  
 205 signal-to-noise ratio (SNR) is defined as follows:

$$206 \quad \text{SNR} = \frac{P(t, z)}{P_{bg}(t)} \quad (2)$$

207  $P_{bg}$  is the ceilometer’s background signal, which the instrument measures at the end  
 208 of its laser pulse cycle,  $P$  is the raw laser power received by the instrument and  $z$   
 209 is the altitude a.s.l. from which the backscattered light was retrieved. We removed  
 210 any data points from profiles which had an SNR less than two. Following the SNR  
 211 analysis, the total power of backscattered light detected by the ceilometer,  $P(t, z)$ ,  
 212 was corrected for the background signal:

$$213 \quad P_c(t, z) = P(t, z) - P_{bg}(t) \quad (3)$$

214 Following the initial quality control, rolling correlation coefficients were calculated  
 215 between  $A(t)$  and  $P_c(t, z)$ . This was completed in a two-step process. First, a sub-set

216 of the time-series,  $T$ , was defined:

$$217 \quad T = \{t_{i-\Delta t/2}, t_{i-\Delta t/2+1}, \dots, t_{i+\Delta t/2}\} \quad (4)$$

218 where  $t_i$  is a specific time in the observation period and  $\Delta t$  defines the temporal  
 219 width of the sub-set around  $t_i$ . In this work, temporal widths between 1 and 20  
 220 hours were studied. As observations were recorded every minute, the sub-set  $T$   
 221 contained at least 60 data points and at most 1200. Spearman Rank Correlation  
 222 coefficients,  $r_s$ , were then calculated as follows:

$$223 \quad r_s(x, y) = \frac{\text{cov}(\text{rank}(x), \text{rank}(y))}{\sigma_{\text{rank}(x)}\sigma_{\text{rank}(y)}} \quad (5)$$

224 where  $\text{rank}(x)$  is a function which assigns an integer ranking to each value of a set  
 225  $x$ ;  $\text{cov}(x, y)$  is the covariance of two sets of data,  $x$  and  $y$ ; and  $\sigma_x$  is the standard  
 226 deviation of the set  $x$  (Spearman, 1904). Here,  $x$  and  $y$  are the sub-sets of  $A(t)$  and  
 227  $P_c(t, z)$  defined by  $T$ . Applying equations 4 & 5 to the entire time-series forms a  
 228 matrix,  $R_{\Delta t}(t, z)$ . A detailed justification of this range of time-scales is provided  
 229 in Section 4.3. Two additional post-processing procedures were implemented after  
 230 the correlations were calculated. If a subset,  $T$ , contained less than 20 valid data  
 231 points, then the correlation coefficient was labelled as not a number. For the remain-  
 232 ing data, a significance test was performed for each correlation value to ensure that  
 233 the value was significantly larger than zero ( $p < 0.05$ ). If the calculated correlation  
 234 coefficient failed the significance test, it was re-assigned a value of zero.

235 Once the fully quality-controlled correlation analysis had been completed, we  
 236 developed a simple metric to classify the mixing state of the atmospheric layer below  
 237 the lowest observed cloud. First, the average below-cloud correlation coefficient,  
 238  $\bar{r}_{bc}$ , was calculated. When  $\bar{r}_{bc} > 0$  ( $p < 0.05$ ), the below-cloud layer was classified:  
 239 “well-mixed”. In such cases, particles observed at sea-level were considered to be  
 240 well-mixed into the overlying cloud. However, if the average below-cloud correlation  
 241 coefficient didn’t exceed zero, then the surface layer below-cloud layer was classified:  
 242 “poorly-mixed”. In these cases, the atmospheric layer at the surface was assumed  
 243 to be decoupled from the overlying cloud. Finally, if there were insufficient data  
 244 points in any sub-set of the time series,  $T$ , then the correlation analysis was unable  
 245 to classify the mixing state of the below-cloud layer for that period.

### 246 3.2 Validation

247 To validate the proposed methodology and classification metric, we compared  
 248 results to four separate methods of determining the mixing state of the below-cloud  
 249 layer. The first two methods were variations on a conventional radiosonde analysis,  
 250 one was a surface-based method and the final method was a model-based method.

251 We compared the classification of the below-cloud mixing state according to  
 252 the correlation metric to two methods for detecting boundary layer decoupling based  
 253 on radiosonde profiles. The first method searched for maxima in the virtual poten-  
 254 tial temperature gradient ( $\partial\theta_v/\partial z^{-1}$ ) (Hande et al., 2012b). If a local maxima in  
 255 the virtual potential temperature gradient was detected and found to exceed 10 K  
 256  $\text{km}^{-1}$ , then the height at which this occurred was labelled as the main inversion,  
 257 or the boundary layer height. The method then searched for secondary maxima  
 258 larger than 5 K  $\text{km}^{-1}$  below the main inversion. If secondary inversions exist, then  
 259 the boundary layer is decoupled (Hande et al., 2012b). To be consistent with our  
 260 methodology, which can only classify the atmospheric layer below the lowest ob-  
 261 served cloud, the below-cloud layer was only labelled as decoupled if a secondary  
 262 inversion was located between the surface and the cloud.

263 A second method for detecting below-cloud decoupling was adapted from  
 264 Truong et al. (2020). In this method, a main inversion was only identified if a local

265 maximum in the virtual potential temperature gradient exceeded  $14 \text{ K km}^{-1}$ . To  
 266 detect decoupling, the decoupling parameter  $\mu$  was studied (Truong et al., 2020; Yin  
 267 & Albrecht, 2000). The decoupling parameter,  $\mu$ , is defined as follows:

$$268 \quad \mu = - \left( \frac{\partial \theta}{\partial z} - \frac{0.608 \theta}{1 + 0.608 r} \frac{\partial r}{\partial z} \right) \quad (6)$$

269 Where  $r$  is the water vapour mixing ratio and  $\theta$  is the potential temperature. Yin  
 270 and Albrecht (2000) devised  $\mu$  for their study of “transition layers” in the bound-  
 271 ary layer, as it is more sensitive to changes in the water vapour mixing ratio than  
 272 the vertical gradient of virtual potential temperature and is therefore more likely  
 273 to detect subtle boundary layer features like decoupling. Decoupling of the bound-  
 274 ary layer over the Southern Ocean was only detected when a value of  $\mu$  exceeded  
 275 2.5 times its average value throughout the boundary layer (Truong et al., 2020).  
 276 To be consistent with our method, we adapted this method to only classify the  
 277 below-cloud layer as decoupled if the threshold for  $\mu$  was exceeded in the below-  
 278 cloud layer (Truong et al., 2020). We used a simple optimization methodology to  
 279 determine which combination of time-scale,  $\Delta t$ , and correlation threshold,  $r_t$ , best  
 280 predicted the state of coupling between the surface and cloud layers as compared to  
 281 the reference methods (Hande et al., 2012b; Truong et al., 2020).

282 To provide a benchmark for our methodology, we compared the optimized  
 283 performance of the correlation-based method against another surface-based method-  
 284 ology for defining the mixing state of the below-cloud layer (Jones et al., 2011).  
 285 Briefly, if the difference in height between the observed cloud base height (CBH)  
 286 and lifted condensation level (LCL) exceeded 150 m, then the below-cloud layer was  
 287 considered to be decoupled from the cloud (and well-mixed otherwise). For these  
 288 calculations, the LCL was calculated from the AWS measurements, where the LCL  
 289 represents the height at which a cloud is expected to form based on a parcel of air  
 290 adiabatically ascending through a well-mixed boundary layer (Romps, 2017). Here,  
 291 we used the 1-hour averaged LCL and for consistency, the 10<sup>th</sup>-percentile of CBH  
 292 within each hour.

293 While the radiosonde profiles collected throughout the voyage provided a ro-  
 294 bust benchmark for the new methodology, radiosonde data were available at most  
 295 twice-a-day. To increase our confidence in the methodology, we compared its classifi-  
 296 cation of the below-cloud layer to near-surface measures of atmospheric stability. We  
 297 used two measures of near-surface atmospheric stability: the square of the Brunt-  
 298 Väisälä Frequency,  $N^2$ , and the 10-m wind speed. Values of  $N^2$  were calculated from  
 299 the AWS measurements and the COARE 3.5 bulk-flux algorithms:

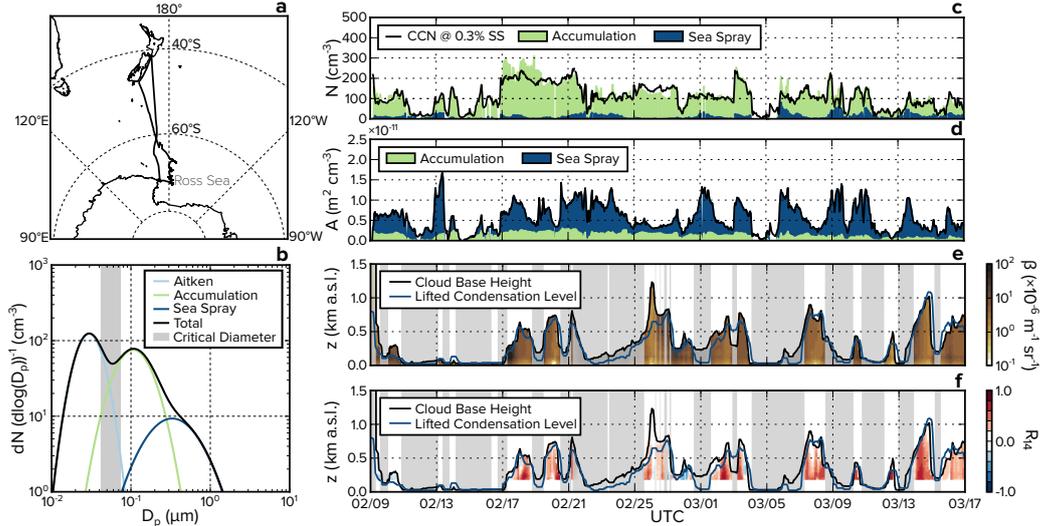
$$300 \quad N^2 = \frac{g}{\theta_v} \frac{\partial \theta_v}{\partial z} \quad (7)$$

302 where  $g$  is the gravitational acceleration,  $\theta_v$  is the virtual potential temperature,  
 303 and  $z$  the height above sea level. This stability analysis was combined with fore-  
 304 casts from AMPS to quantitatively define conditions in which aerosol–cloud coupling  
 305 was expected. As a coarse proxy for aerosol–cloud coupling, we investigated the  
 306 difference in the LCL and the predicted planetary boundary layer height (PBL) in  
 307 the AMPS forecasts. If the planetary boundary layer exceeded the lifted conden-  
 308 sation level, then aerosol particles measured at the ocean surface were considered  
 309 well-mixed to the minimum height where cloud could have occurred.

## 310 4 Results

### 311 4.1 Time Series Analysis

312 Throughout the voyage to and from the Ross Sea (voyage track shown in  
 313 Fig. 1a), the number–size distribution of particulate was predominantly trimodal as



**Figure 1.** (a) The track of the R/V *Tangaroa* during the Marine Environment and Ecosystem Voyage. (b) A typical size distribution for particles in the Southern Ocean. The expected range of cloud activation diameters for marine stratus is shown in grey. (c) The sea-level abundance of sea spray particles (SSPs; blue filled region) and accumulation mode particles (green filled region) is compared to the abundance of cloud condensation nuclei (CCN) at a supersaturation of 0.3% (black line). (d) The abundance of suspended surface area was calculated from the measured particle size distributions (Eq. 1). (e) A contour plot of the attenuated backscatter coefficient measured by the CHM-15K ceilometer. The lifted condensation level ( $z_{LCL}$ ) and cloud base height ( $z_{CBH}$ ) are also shown for reference. (f) Spearman Rank correlation coefficients between the sea-level abundance of particulate surface area and ceilometer backscatter are shown. Time periods when the ceilometer optical window was obscured, the cloud base was below 200 m, fog was present, or the aerosol sampling system was contaminated by ship exhaust are shaded.

314 seen in Fig. 1b. The representative number–size distribution shown in Fig. 1b was  
 315 constructed by taking voyage wide averages of the total number, width and median  
 316 size of the individual modes that were fit to the observations. The appearance of  
 317 these modes is consistent with previous observations in marine settings (Bates et  
 318 al., 1998; Quinn et al., 2017). A large majority of the particles in the smallest two  
 319 modes, the Aitken (30 nm,  $\sigma = 1.4$ ) and accumulation modes (100 nm,  $\sigma = 1.6$ ), are  
 320 thought to be produced as a single mode from homogeneous nucleation of volatile  
 321 sulfate species, with mode separation occurring as a result of cloud-processing  
 322 (Hoppel et al., 1986). These particles are nucleated in-situ from the condensation  
 323 of oxidized marine gasses and grow via self-coagulation and condensation. In con-  
 324 trast, sea spray particles (400 nm,  $\sigma = 2$ ) are directly generated from breaking ocean  
 325 waves, and tend to be much larger than particles in the Aitken and accumulation  
 326 mode (Prather et al., 2013). Note that size statistics presented in this section have  
 327 been corrected to a relative humidity of 80%. For sulfate and sea spray particles,  
 328 a particle at a relative humidity of 80% is approximately twice as large compared to  
 329 when it is dry (Gerber, 1985).

330 A representative size distribution of particles observed in the Southern Ocean  
 331 marine boundary layer at a relative humidity of 80% is shown in Fig. 1b. The  
 332 bifurcation of the Aitken and accumulation modes occurs when these particles  
 333 pass through non-precipitating cloud, since only the largest particles will be acti-

334 vated (Hoppel et al., 1986). Previous research has shown that the supersaturation  
 335 of water vapour within nascent marine stratus is relatively modest ( $<0.3\%$ ; Hegg  
 336 et al. (2009)). An estimation of the activation diameter based on a supersaturation  
 337 of  $0.3\%$ , and a range of particle hygroscopicity parameters is also shown in Fig. 1b.  
 338 The estimation of the range of activation diameter is based on the  $\kappa$ -Köhler model  
 339 for a range of expected hygroscopicity values (Petters & Kreidenweis, 2007). This  
 340 coincides well with the local minimum between the Aitken and accumulation mode,  
 341 supporting the cloud-processing hypothesis of Hoppel et al. (1986).

342 Fig. 1c displays the number of particles in both the accumulation and sea  
 343 spray modes, as these are the only particles relevant to cloud formation. This is  
 344 compared to the number concentration of CCN measured at a fixed supersaturation  
 345 of  $0.3\%$ . As expected, these two measurements are highly correlated. Across the  
 346 entire voyage, SSPs did not comprise a substantial fraction of CCN ( $14\%$ ). However,  
 347 in the latter half of the voyage we encountered several low pressure systems. These  
 348 cyclones were accompanied by high winds, resulting in substantial wave-breaking  
 349 and subsequent SSP generation in the region. This led to an enhanced relevance of  
 350 SSPs to the total CCN population ( $20\%$ ).

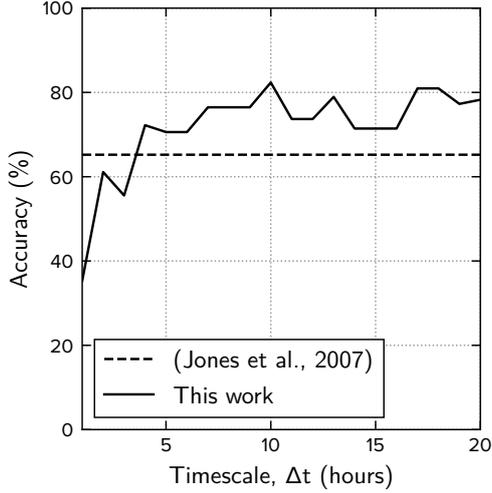
351 Fig. 1d shows the abundance of suspended particle surface area. Despite the  
 352 relatively low abundance of SSPs by number, the total amount of particulate surface  
 353 area is strongly dominated by variations in their abundance. In Fig. 1e, the time se-  
 354 ries of attenuated backscatter profiles measured by a coincident ceilometer is shown,  
 355 along with rolling averages of cloud base height and the lifted condensation level.  
 356 As demonstrated both empirically and theoretically, if the difference between cloud  
 357 base height and lifted condensation level is less than  $150$  m, the below-cloud layer  
 358 can be considered well-mixed (Jones et al., 2011). As a result, it is clear that there  
 359 was significant coupling between the surface layer and overlying cloud for much of  
 360 the time-series.

## 361 4.2 New Classification Methodology

362 We used the Spearman Rank correlation analysis between suspended particle  
 363 surface area at sea-level (Fig. 1d) and ceilometer backscatter from particles overhead  
 364 (Fig. 1e) to assess whether our measurements at the surface were representative of  
 365 the below-cloud population of CCN. Fig. 1f displays strong correlations between  
 366 these two quantities over time-scales of 14 hours when fog, precipitation, or con-  
 367 tamination from ship exhaust did not inhibit the analysis. This suggests that the  
 368 Southern Ocean boundary layer was consistently well-mixed throughout this mea-  
 369 surement campaign. We note that correlation coefficients could not be calculated  
 370 below  $200$  m, as these data typically failed the SNR analysis. While one would nor-  
 371 mally expect a large backscattered signal close to a lidar, and thus a high SNR, the  
 372 returning backscatter is not well-aligned with the FOV of the receiving optics in the  
 373 near-range, resulting in a low SNR. As an additional control, the significance of the  
 374 calculated correlation coefficients was assessed with a two-way t-test.

## 375 4.3 Comparison to Radiosonde Analysis

376 To validate the correlation analysis and establish the most accurate time-  
 377 scale for calculating correlation coefficients, we analyzed the 57 radiosonde profiles  
 378 recorded throughout the voyage. For each radiosonde, we used two gradient methods  
 379 to detect whether the surface layer was decoupled from the cloud layer (Hande et  
 380 al., 2012b). Out of the 57 radiosondes, three could not be used for analysis due to  
 381 faulty sensors, and 21 were launched when the cloud base height was below  $200$  m.  
 382 In such cases, there was insufficient ceilometer data to perform the correlation anal-  
 383 ysis, as the power of the returning backscatter was on the same order of magnitude



**Figure 2.** The accuracy with which the new correlation metric correctly classified the mixing state of the below-cloud layer for different correlation time-scales,  $\Delta t$ . The accuracy of the method was calculated in reference to two radiosonde analyses which classified the mixing state of the below-cloud layer based on thermodynamic gradients ( $N = 26$ ; non-precipitating conditions, no fog,  $CBH > 200$  m) (Hande et al., 2012b; Truong et al., 2020). The accuracy of the proposed method can be compared to the accuracy of another ground-based methodology of determining the state of below-cloud mixing (Jones et al., 2011).

384 as the instrument noise due to the FOV effects described earlier. In three additional  
 385 cases, the gradient methods did not detect a boundary layer. As a result, there were  
 386 only 30 radiosondes available for which the correlation analysis was valid. In these  
 387 remaining 30 cases, the radiosonde-based methods of classifying the below-cloud  
 388 layer differed only slightly. Overall, the below-cloud layer appeared well-mixed in  
 389 83% of profiles according to the criteria of Truong et al. (2020), and 90% of profiles  
 390 according to Hande et al. (2012b).

391 The classifications of the below-cloud mixing state defined by the radiosonde  
 392 analyses were used to define the optimal time-scale,  $\Delta t$ , and threshold,  $r_t$ , for the  
 393 correlation analysis. For consistency, only the profiles for which both radiosonde  
 394 methods agreed on the mixing state of the below-cloud layer were used as a refer-  
 395 ence when calculating the accuracy ( $N = 26$ ). In Fig. 2, the accuracy with which the  
 396 correlation analysis determined the mixing state of the below-cloud layer is shown  
 397 as a function of time-scale. Across all time-scales, the threshold for detecting a well-  
 398 mixed below-cloud layer was  $\overline{r_{bc}} > 0$ , where  $\overline{r_{bc}}$  is the average correlation coefficient  
 399 between sea-level and the 10<sup>th</sup>-percentile of cloud base height. As a benchmark, we  
 400 have also shown the accuracy of another ground-based method for determining the  
 401 below-cloud mixing state (Jones et al., 2011).

402 Fig. 2 demonstrates that the accuracy of the correlation-based method in-  
 403 creased from 35% to 76% as the time-scale increased, until time-scales of 7 hours  
 404 or longer were reached. Differences in accuracy at time-scales beyond 7 hours were  
 405 negligible considering the sample size ( $N = 26$ ). The increase in accuracy with in-  
 406 creasing time-scale is a direct result of increasing the number of samples in the  
 407 subset  $T$  (defined in Eq. 4) used for calculating the correlation coefficient. While  
 408 shorter time-scales are likely more representative of the time-scale of turbulence,

**Table 1.** This table summarizes statistics comparing measurements from radiosondes launched throughout the voyage and predictions from AMPS below 3 km (a.s.l.).

Statistic	P (hPa)	T (K)	$T_d$ (K)	U ( $\text{m s}^{-1}$ )
RMSD	2.0	1.2	3.2	2.7
Bias	0.3	–	-0.8	0.7
$R^2$	1	0.96	0.87	0.74

RMSD: Root Mean Squared Deviation

 $R^2$ : Pearson Correlation Coefficient

409 there is also a higher likelihood that the remaining noise in the ceilometer obser-  
410 vations will result in weaker correlations which fail the two-way t-test ( $p > 0.05$ ).  
411 Increasing the time-scale resulted in more consistent correlation coefficients across  
412 time-scales and more statistically significant results overall. The accuracy of this  
413 method also suggests that despite there being longer time-scale phenomena which  
414 could also correlate particulate surface area and backscatter (e.g. frontal systems,  
415 convective forcing at cloud top, precipitation, turbulent perturbations of relative  
416 humidity, air mass history, etc.), these phenomena are not likely to result in sub-  
417 stantial misclassification of the below-cloud layer. However, considering that such  
418 long time-scale phenomena do exist and may be more prevalent in other regions or  
419 observation periods, correlation coefficients calculated over time-scales beyond those  
420 presented here should be avoided as false positives and false negatives are likely to  
421 become more abundant.

422 Finally, we compared the accuracy of our new methodology to another method  
423 of remotely classifying the below-cloud mixing state (Jones et al., 2011). In this  
424 case, the referenced method was only 65% accurate at determining the mixing state  
425 of the below-cloud layer, whereas the proposed method was 76% accurate when  
426 correlation time-scales greater than 7 hours were considered. While the set of ra-  
427 diosondes for which we could compare both methods was quite limited ( $N = 26$ ),  
428 these results suggest that the proposed method more accurately classified the mixing  
429 state of the below-cloud layer than the referenced method ( $p < 0.05$ ).

430 Overall, the correlation analysis found that the below-cloud layer was well-  
431 mixed for 14% of the entire time series and poorly-mixed just 5%. Fog was found to  
432 occur 7% of the time, where fog was diagnosed when the relative humidity was mea-  
433 sured to be 100% and cloud base was less than 50 m. The remaining portion of the  
434 time series could not be analyzed (74%), as one or more of the following occurred:  
435 the ceilometer’s quality control flag was raised; the ship exhaust contaminated the  
436 aerosol sample; or, the cloud base was below 200 m but greater than 50 m, such  
437 that the entire profile of below-cloud backscatter failed the SNR analysis due to a  
438 lack of overlap between the FOV of the ceilometer’s optical system and the return-  
439 ing backscatter. While this may seem like a large loss of the time-series, if a given  
440 radiosonde was only representative of conditions for the hour of measurements in  
441 which it was operating, then the radiosonde analysis provided data for just 7% of  
442 the time series. With the proposed correlation analysis, we were able to classify the  
443 boundary layer for 26% of the time series, a marked improvement.

#### 444 4.4 Comparison to Stability Analysis

445 While the comparison to conventional radiosonde analyses provided evidence  
446 that the correlation analysis accurately classified the below-cloud mixing state, it  
447 still seemed prudent to evaluate the analysis against other metrics of atmospheric

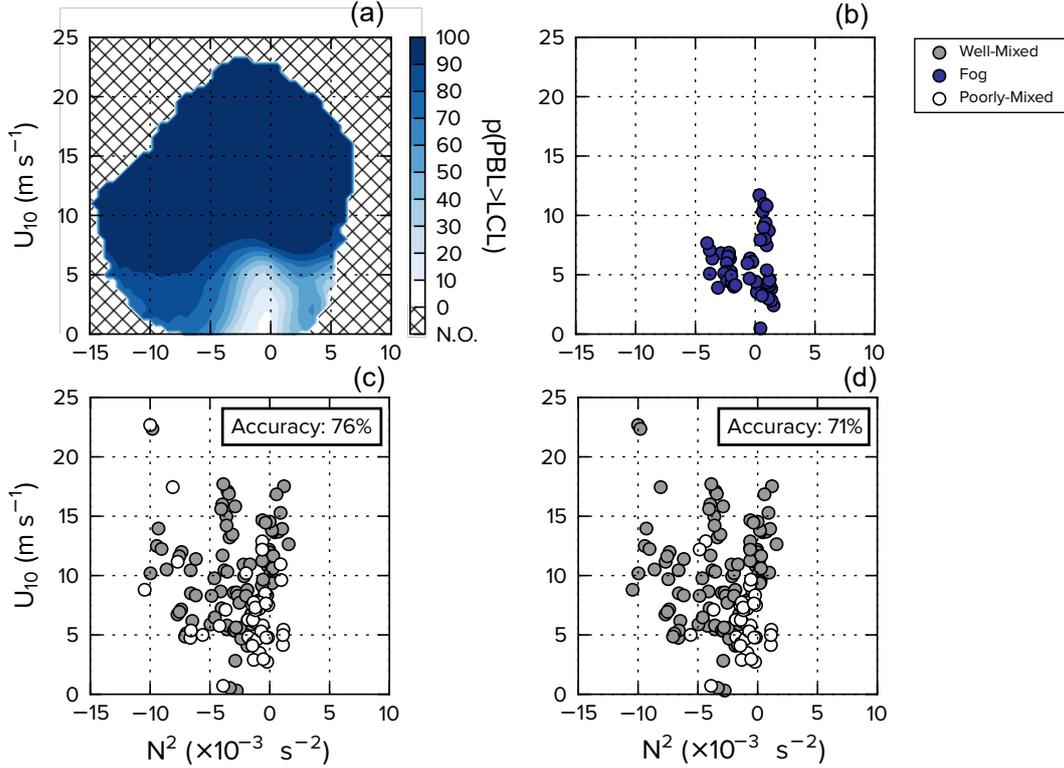
448 mixing. Here, we examine expected rates of occurrence for aerosol-cloud coupling  
 449 based on AMPS forecasts for the period of study.

450 First, the forecasts were compared to our observations from all available ra-  
 451 diosondes, except those with non-functioning RH sensors. Table 1 provides a sum-  
 452 mary of various comparison statistics between forecasts and measurements. The  
 453 statistics were only calculated below 3 km to restrict the comparison to relevant  
 454 planetary boundary layer (PBL) and lifted condensation level (LCL) heights. These  
 455 are presented in Table 1.

456 Overall, observed and measured values of the selected variables were rea-  
 457 sonably well correlated. However, there were minor biases worth mentioning. In  
 458 Table 1, statistically significant biases between modelled and measured values of  
 459 pressure, dew point temperature, and wind speed were observed ( $p < 0.001$ ). Within  
 460 AMPS, the height of the PBL is determined according to the turbulent kinetic en-  
 461 ergy profile (Janjic, 2001). This implies that the height of the PBL may have been  
 462 under-estimated by AMPS. The dew point temperature was also negatively biased  
 463 as a result of the over-abundance of water vapour in the AMPS boundary layer re-  
 464 lative to observations. This implies that the LCL was also under-predicted. Still,  
 465 considering the spatial and temporal scale of the AMPS forecasts, the agreement be-  
 466 tween model and measured values was quite good, and highly statistically significant  
 467 ( $p < 0.001$ ).

468 In Fig. 3a, the frequency of occurrence with which the depth of the planetary  
 469 boundary layer exceeded the lifted condensation level is shown, based on forecasts  
 470 from AMPS between 40 – 70 S (excluding areas less than 100 km from a coast).  
 471 The frequency of occurrence is shown as a function of two variables which are often  
 472 used to describe the stability of the near-surface atmospheric layer: the square of the  
 473 Brunt-Väisälä frequency,  $N^2$ , and the 10-m wind speed,  $U_{10}$ . The results in Fig. 3a  
 474 demonstrate that in near-neutral stability ( $N^2 \sim 0$ ) and weak winds ( $U_{10} \sim 0$ ), the  
 475 layer below the LCL was less-likely to be well-mixed, as the PBL was too shallow.  
 476 However, in all other cases, the boundary layer was likely well-mixed as the PBL  
 477 exceeded the LCL. Note that there was still a small percentage of the time when the  
 478 layer below the LCL was well-mixed despite the surface layer being near-neutrally  
 479 stable. While a well-mixed boundary layer would not be expected in such cases, the  
 480 near-surface layer is typically much shallower than the LCL, and is therefore not  
 481 always a perfect determinant of the mixing state of the entire layer below the LCL.  
 482 However, it is clear from Fig. 3a that in most other conditions, the layer below the  
 483 LCL is almost guaranteed to be well-mixed. Overall, The AMPS analysis in Fig. 3a  
 484 provides a general rule of thumb: if the 10 m wind speed exceeds  $8 \text{ m s}^{-1}$  then the  
 485 boundary layer will be well-mixed to the LCL, regardless of the near-surface stabil-  
 486 ity.

487 In Fig. 3c and d, the classification of the below-cloud layer according to the  
 488 correlation metric is shown for time-scales of 7 and 14 hours, respectively. We can  
 489 see that despite the accuracy with which the correlation analysis at 7 hour and 14  
 490 hour time-scales classified the mixing state of the boundary layer (Fig. 2), the cor-  
 491 relation metric calculated over a 14 hour time-scale provided a more qualitatively  
 492 consistent result with the AMPS analysis. In comparing Figs. 3c & d, it is clear  
 493 that the correlation metric calculated over a time-scale of 7 hours misclassified the  
 494 boundary layer more frequently, as a poorly mixed boundary layer is not expected  
 495 to occur at all if  $N^2 < -5 \times 10^{-3} \text{ s}^{-2}$  or  $U_{10} > 8 \text{ m s}^{-1}$  (Fig. 3a). Barring a few  
 496 exceptions, Fig. 3d shows that the correlation metric at a time-scale of 14 hours  
 497 typically only classified the below-cloud layer as decoupled only when the stability of  
 498 the near-surface layer was near-neutral and winds were less than  $8 \text{ m s}^{-1}$ , consistent  
 499 with the AMPS analysis. While only two time-scales are presented here, analysis at  
 500 all time-scales longer than 14 hours produced qualitatively similar results. Finally, in



**Figure 3.** (a) The frequency with which the height of the planetary boundary layer (PBL) predicted by AMPS was higher than the lifted condensation level (LCL) over the open Southern Ocean (40–70 S, >100 km from coastline) in February and March 2018 (N.O. = Conditions occurred less than frequently than 0.001%). (b) The occurrence of fog (CBH < 50 m, RH = 100%). (c) The classification of the below-cloud layer based on correlation coefficients calculated over 7-hour timescales (non-precipitating conditions; CBH > 200 m). The measure of accuracy is in reference to the radiosonde analyses (Fig. 2). (d) As in (c), but for a time-scale of 14-hours.

501 Fig. 3b, it is clear that fog tended to occur only in both near-neutral stability ( $N^2 >$   
 502  $-5 \times 10^{-3}$  s<sup>-2</sup>) and low winds, or stable conditions, consistent with advection fog.

503 One limitation of this analysis is that a cloud is not necessarily guaranteed to  
 504 occur at the LCL. As such, a direct quantitative comparison between Figs. 3a, c &  
 505 d is not possible, as cloud was always occurring in the subset of data we were able  
 506 to analyze but may not have been occurring in the AMPS forecasts. Still, we found  
 507 this figure to be a useful qualitative reference for our methodology. In addition, it  
 508 demonstrates that even though the fraction of the time-series available for analysis  
 509 via the correlation metric is low (26%;  $\sim 10$  days of observations), the conditions  
 510 encountered within this subset of the data are representative of the wide set of con-  
 511 ditions forecast by AMPS. As a result, statistics presented in the previous section  
 512 can be used to conclude that in non-precipitating conditions, the below-cloud layer  
 513 over the Ross Sea was likely well-mixed 54% of the time, poorly-mixed 19% of the  
 514 time, and contained fog 27% of the time.

## 5 Discussion

In this work, we were interested in understanding how often aerosol particles measured near the surface of the ocean were relevant to low cloud formation over the Southern Ocean. We proposed a new methodology, based on the correlation of particle surface area and ceilometer backscatter, which identified when aerosol particles observed at the surface were available to the lowest observed cloud. To validate the proposed methodology, we needed an accurate reference classification of the boundary layer against which we could compare our results. Here, we modified two radiosonde-based methodologies which determined the mixing state of the entire boundary layer based on gradients of thermodynamic variables (Hande et al., 2012b; Truong et al., 2020). Radiosonde-based methodologies were selected as the preferred reference methodology, as the observations were sensitive to fine thermodynamic changes in the boundary layer. These reference methodologies were modified to simply determine the mixing state of the atmosphere between the surface and the base of the lowest observed cloud. This allowed us to optimize the parameters of our proposed methodology (correlation time-scale, threshold of correlation strength) such that the predicted mixing state of the below-cloud layer best matched the referenced radiosonde methodologies. In the comparison (Fig. 2), the proposed correlation-based method correctly classified the mixing state of the below-cloud layer  $76 \pm 4\%$  of the time for correlation time-scales greater than 7 hours. The accuracy of our method was then compared to a more simple metric for classifying the mixing state of the below cloud layer, which was only accurate 65% of the time (Jones et al., 2011).

In a more qualitative comparison (Fig. 3), the classification of the below-cloud mixing state by the proposed methodology was also shown to be consistent with surface-based measurements of atmospheric stability and model predictions of turbulence in a wide range of conditions. The high accuracy of the new methodology’s predictions in comparison to radiosonde-based methods, in situ observations of near-surface atmospheric stability, and model forecasts of boundary layer turbulence gives us high confidence that the proposed method is accurate even when reference data is not available.

With the accuracy of our proposed methodology validated against multiple methods of determining the below-cloud mixing state, we can compare statistics to previous observations in the Southern Ocean. We find that while the below-cloud layer was often well-mixed, this was not always guaranteed. It is well-known that the marine boundary layer can stratify into a near-surface boundary layer and a sub-cloud layer (Garratt, 1994). In fact, radiosondes launched from Macquarie Island (54.62°S, 158.85°E) over the past two decades found that the boundary layer was well-mixed just 17.8% of the time (Hande et al., 2012b). In contrast, our time series analysis showed that in non-precipitating conditions, the below-cloud layer was well-mixed 54% of the time. This seems to be in stark contrast to the accuracy data presented in Fig. 2. However, the difference in frequency of occurrence comes primarily from a difference in the definition of decoupling. The method presented in this work was only designed to detect whether the boundary layer was well-mixed up to the lowest cloud. In contrast, the method used to analyze radiosondes launched from Macquarie Island was designed to detect decoupling throughout the entire boundary layer (Hande et al., 2012b). However, multi-layer clouds are frequently observed over the Southern Ocean (Hande et al., 2012b). In such settings, the inversion atop the lowest cloud will tend to decouple the atmospheric layer beneath the cloud from the rest of the boundary layer. Despite this decoupling, near-surface air is typically still well-mixed up to the lowest cloud, as cloud was often present in the atmospheric layer beneath the decoupling height (Hande et al., 2012b). As a result of this inconsistency, statistics retrieved from the radiosonde analysis at Macquarie

568 Island cannot be easily used to infer information about aerosol–cloud coupling and  
569 are not directly comparable to this study.

570 Overall, the results from the correlation analysis highlight that particles are al-  
571 most always available to the lowest cloud (Fig. 1f). The percentage of time in which  
572 aerosol–cloud coupling occurred within the valid section of our time series is simply  
573 the sum total of the rates of occurrence of fog and a well-mixed below-cloud layer:  
574 81%. Forecasts from AMPS tend to agree, as the layer of the atmosphere below the  
575 LCL was found to be well-mixed 84% of the time over the Southern Ocean through-  
576 out February and March 2018. Based on the good agreement between these methods  
577 of defining the below-cloud mixing state, we are confident in concluding that sea  
578 spray particles are available to low-level cloud over the Southern Ocean more than  
579 80% of the time in austral summer. As Kuma et al. (2020) noted, the ability to  
580 correctly predict the occurrence of low cloud is a critical necessity for improving  
581 the Southern Ocean shortwave radiation bias. The proposed method increases our  
582 understanding of these low clouds and the particles which help form them.

583 For instance, we found that the number of CCN at a supersaturation of 0.3%  
584 was consistent with the number of particles in the accumulation and sea spray mode.  
585 As a supersaturation of 0.3% is the expected water vapor supersaturation within  
586 marine stratocumulus (Hegg et al., 2009), this suggests that sea-level observations  
587 may provide a good constraint on the number of cloud droplets in a wide variety  
588 of conditions. We found that despite being readily-available to nascent clouds, sea  
589 spray particles were typically outnumbered by smaller, cloud-processed accumulation  
590 mode particles (Fig. 1c), consistent with previous studies (Quinn et al., 2017). How-  
591 ever, in addition to abundance, the ice-nucleating ability of particles is known to be  
592 a strong determinant of cloud phase and albedo: a climate model which determined  
593 the primary nucleation of ice within low-level clouds according to the abundance and  
594 type of boundary layer ice-nucleating particles found that predictions of cloud opac-  
595 ity were significantly more accurate in the cold sector of Southern Ocean cyclones  
596 relative to simpler glaciation schemes (Vergara-Temprado et al., 2018). Though less  
597 numerous than accumulation mode particles, sea spray particles are thought to be  
598 the only local source of ice-nucleating particles (DeMott et al., 2016) in a region  
599 that is often devoid of more potent ice nuclei (e.g. dust) (McCluskey et al., 2018).  
600 This study highlights that sea spray particles are available to many more cloud  
601 systems than just within the cold sector of cyclones. As a result, climate models  
602 which implement glaciation schemes that connect the primary nucleation of ice to  
603 the microphysical properties of aerosol particles will likely see more widespread im-  
604 provement to the Southern Ocean shortwave radiation bias. It also highlights that  
605 should models adopt more complex models of cloud glaciation, then they must also  
606 more carefully parameterize the flux of sea spray particles (Hartery et al., 2020).

607 The new method does not come without limitations, however. Depending on  
608 the FOV of the ceilometer’s optical receiver, the ability of the analysis to analyze  
609 below-cloud coupling in low cloud settings ( $CBH < 200$  m) can be severely impaired  
610 due to an incomplete overlap of the returning laser beam and the receiver’s FOV.  
611 In addition, though we have provided a reasonably comprehensive validation of the  
612 appropriate time-scale for calculating correlation coefficients and the threshold for  
613 classification of the below-cloud layer, there are potentially instances where the cor-  
614 relation analysis could trigger false positives and false negatives in other synoptic  
615 settings. These include, but are not limited to, frontal systems, convective forcing  
616 at cloud top, precipitation, turbulent perturbations of relative humidity, air mass  
617 history, etc. Still, given the accuracy of the methodology as quantitatively compared  
618 to the radiosonde analyses, and qualitatively to a forecast analysis, we are confident  
619 in the results presented as they pertain to this specific region and period of study.  
620 As an added benefit, the proposed method also uses instruments which function

621 nearly autonomously, with little need for oversight or on-site personnel. In contrast,  
 622 radiosonde programs require highly-trained personnel and can only be launched in a  
 623 limited set of meteorological conditions. It becomes exceedingly difficult to success-  
 624 fully launch a radiosonde once winds surpass  $15 \text{ m s}^{-1}$ , and potentially dangerous  
 625 when aboard a research vessel in unfavorable wave conditions. As a result, statis-  
 626 tics of boundary layer mixing collected from radiosonde programs are likely skewed  
 627 towards calm conditions. Finally, it is worth mentioning that there are likely un-  
 628 intended, negative environmental consequences of leaving irretrievable radiosonde  
 629 packages in the Southern Ocean.

## 630 6 Conclusions

631 In this work we presented a new technique for determining the state of bound-  
 632 ary layer mixing based on the value of the Spearman Rank correlation coefficient  
 633 calculated between sea-level observations of suspended particle surface area and  
 634 ceilometer backscatter. When data was available, these correlations were often high,  
 635 implying that particles measured at sea-level were well-mixed throughout the bound-  
 636 ary layer and were therefore readily-available to nascent, low-level cloud. From this  
 637 analysis, a simple metric was created to diagnose whether coupling occurred or not.  
 638 This revealed that in non-precipitating conditions the boundary layer was well-  
 639 mixed 54% of the time, contained fog 27% of the time, and was poorly-mixed just  
 640 19% of the time. This simple metric based on the correlation analysis was compared  
 641 to two conventional radiosonde analyses. The correlation-based metric accurately  
 642 classified the mixing state of the boundary layer 76% of the time when correlation  
 643 coefficients were calculated over periods longer than 7 hours. This is a noticeable im-  
 644 provement over the accuracy of a simpler ground-based method (65%). In addition,  
 645 the frequency of occurrence of below-cloud mixing estimated by the correlation-  
 646 based metric was qualitatively consistent with an analysis of mixing based on the  
 647 near-surface stability within regional forecasts. We estimate that aerosol will have  
 648 a direct pathway into low cloud either through a well-mixed below cloud layer or  
 649 surface-level fog, 81% of the time when clouds are present. Thus, in situ sea-level  
 650 observations of particulate offer substantial insight into cloud formation over the  
 651 Southern Ocean in a wide set of conditions.

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 665 at: [https://zenodo.org/record/4060237#.X\\_xcK-hKiMo](https://zenodo.org/record/4060237#.X_xcK-hKiMo). AMPS data used in  
 666 this study were downloaded from: [https://www.earthsystemgrid.org/project/](https://www.earthsystemgrid.org/project/amps.html)  
 667 [amps.html](https://www.earthsystemgrid.org/project/amps.html).

## 668 References

669 Bates, T. S., Huebert, B. J., Gras, J. L., Griffiths, F. B., & Durkee, P. A. (1998).

- 670 International global atmospheric chemistry (igac) project's first aerosol char-  
 671 acterization experiment (ace 1): Overview. *Journal of Geophysical Research:*  
 672 *Atmospheres*, 103(D13), 16297–16318.
- 673 Bigg, E. K., & Hopwood, S. C. (1963). Ice nuclei in the antarctic. *Journal of the*  
 674 *Atmospheric Sciences*, 20(3), 185–188. doi: 10.1175/1520-0469(1963)020<0185:  
 675 INITA>2.0.CO;2
- 676 Bodas-Salcedo, A., Williams, K. D., Ringer, M. A., Beau, I., Cole, J. N., Dufresne,  
 677 J.-L., ... Yokohata, T. (2014). Origins of the solar radiation biases over the  
 678 Southern Ocean in CFMIP2 models. *Journal of Climate*, 27(1), 41–56.
- 679 Boers, R., Jensen, J., & Krummel, P. (1998). Microphysical and short-wave ra-  
 680 diative structure of stratocumulus clouds over the southern ocean: Summer  
 681 results and seasonal differences. *Quarterly Journal of the Royal Meteorological*  
 682 *Society*, 124(545), 151–168.
- 683 Bohren, C. F., & Huffman, D. R. (1983). Absorption and scattering by a sphere. *Ab-*  
 684 *sorption and scattering of light by small particles*.
- 685 Borys, R. D., Lowenthal, D. H., Cohn, S. A., & Brown, W. O. J. (2003). Moun-  
 686 taintop and radar measurements of anthropogenic aerosol effects on snow  
 687 growth and snowfall rate. *Geophysical Research Letters*, 30(10). doi:  
 688 10.1029/2002GL016855
- 689 DeMott, P. J., Hill, T. C. J., McCluskey, C. S., Prather, K. A., Collins, D. B., Sulli-  
 690 van, R. C., ... Franc, G. D. (2016). Sea spray aerosol as a unique source of ice  
 691 nucleating particles. *Proceedings of the National Academy of Sciences*, 113(21),  
 692 5797–5803. doi: 10.1073/pnas.1514034112
- 693 Edson, J. B., Jampana, V., Weller, R. A., Bigorre, S. P., Plueddemann, A. J.,  
 694 Fairall, C. W., ... Hersbach, H. (2013). On the exchange of momentum over  
 695 the open ocean. *Journal of Physical Oceanography*, 43(8), 1589–1610. doi:  
 696 10.1175/JPO-D-12-0173.1
- 697 Garratt, J. (1994). Review: the atmospheric boundary layer. *Earth-Science Reviews*,  
 698 37(1), 89 - 134. Retrieved from [http://www.sciencedirect.com/science/  
 699 article/pii/0012825294900264](http://www.sciencedirect.com/science/article/pii/0012825294900264) doi: [https://doi.org/10.1016/0012-8252\(94\)  
 700 90026-4](https://doi.org/10.1016/0012-8252(94)90026-4)
- 701 Gerber, H. E. (1985). *Relative-humidity parameterization of the Navy Aerosol Model*  
 702 *(NAM)* (Tech. Rep.). Washington, DC: Naval Research Lab.
- 703 Hande, L., Siems, S., & Manton, M. (2012a). Observed trends in wind speed over  
 704 the Southern Ocean. *Geophysical Research Letters*, 39(11). doi: 10.1029/  
 705 2012GL051734
- 706 Hande, L., Siems, S., Manton, M., & Belusic, D. (2012b). Observations of wind  
 707 shear over the southern ocean. *Journal of Geophysical Research: Atmospheres*,  
 708 117(D12). doi: 10.1029/2012JD017488
- 709 Hande, L. B., Siems, S. T., Manton, M. J., & Lenschow, D. H. (2015). An eval-  
 710 uation of cosmic radio occultation data in the lower atmosphere over the  
 711 southern ocean. *Atmospheric Measurement Techniques*, 8(1), 97–107. Re-  
 712 trieved from <https://amt.copernicus.org/articles/8/97/2015/> doi:  
 713 10.5194/amt-8-97-2015
- 714 Hartery, S., Toohey, D., Revell, L., Sellegri, K., Kuma, P., Harvey, M., & McDon-  
 715 ald, A. (2020). Constraining the surface flux of sea spray particles from the  
 716 southern ocean. *Journal of Geophysical Research: Atmospheres*, 125(4). doi:  
 717 10.1029/2019JD032026
- 718 Haynes, J. M., Jakob, C., Rossow, W. B., Tselioudis, G., & Brown, J. (2011). Major  
 719 characteristics of southern ocean cloud regimes and their effects on the energy  
 720 budget. *Journal of Climate*, 24(19), 5061–5080. doi: 10.1175/2011JCLI4052.1
- 721 Hegg, D. A., Covert, D. S., Jonsson, H. H., & Woods, R. (2009). Differentiating  
 722 natural and anthropogenic cloud condensation nuclei in the california coastal  
 723 zone. *Tellus B: Chemical and Physical Meteorology*, 61(4), 669–676. doi:  
 724 10.1111/j.1600-0889.2009.00435.x

- 725 Hoppel, W. A., Frick, G. M., & Larson, R. E. (1986). Effect of nonprecipitating  
726 clouds on the aerosol size distribution in the marine boundary layer. *Geophysical*  
727 *Research Letters*, *13*(2), 125-128. doi: 10.1029/GL013i002p00125
- 728 Hu, Y., Rodier, S., Xu, K.-M., Sun, W., Huang, J., Lin, B., ... Josset, D. (2010). Oc-  
729 currence, liquid water content, and fraction of supercooled water clouds from  
730 combined caliop/iir/modis measurements. *Journal of Geophysical Research:*  
731 *Atmospheres*, *115*(D4).
- 732 Irving, D., Simmonds, I., & Keay, K. (2010). Mesoscale cyclone activity over the ice-  
733 free southern ocean: 1999–2008. *Journal of Climate*, *23*(20), 5404-5420. doi: 10  
734 .1175/2010JCLI3628.1
- 735 Janjic, Z. I. (2001). *Nonsingular implementation of the mellor-yamada level 2.5*  
736 *scheme in the ncep meso model* (Office Note No. 437). National Center for En-  
737 vironmental Prediction (NCEP).
- 738 Jolly, B., McDonald, A. J., Coggins, J. H. J., Zawar-Reza, P., Cassano, J., Laz-  
739 zara, M., ... Dale, E. (2016). A validation of the Antarctic mesoscale pre-  
740 diction system using self-organizing maps and high-density observations  
741 from SNOWWEB. *Monthly Weather Review*, *144*(9), 3181-3200. doi:  
742 10.1175/MWR-D-15-0447.1
- 743 Jones, C. R., Bretherton, C. S., & Leon, D. (2011). Coupled vs. decoupled boundary  
744 layers in vocals-rex. *Atmospheric Chemistry and Physics*, *11*(14), 7143–7153.  
745 Retrieved from <https://acp.copernicus.org/articles/11/7143/2011/>  
746 doi: 10.5194/acp-11-7143-2011
- 747 Korhonen, H., Carslaw, K. S., Forster, P. M., Mikkonen, S., Gordon, N. D., &  
748 Kokkola, H. (2010). Aerosol climate feedback due to decadal increases in  
749 southern hemisphere wind speeds. *Geophysical Research Letters*, *37*(2). doi:  
750 10.1029/2009GL041320
- 751 Kremser, S., Harvey, M., Kuma, P., Hartery, S., Saint-Macary, A., McGregor, J., ...  
752 Parsons, S. (2020). Southern ocean cloud and aerosol data: a compilation of  
753 measurements from the 2018 southern ocean ross sea marine ecosystems and  
754 environment voyage. *Earth System Science Data*(submitted).
- 755 Kuma, P., McDonald, A. J., Morgenstern, O., Alexander, S. P., Cassano, J. J., Gar-  
756 rett, S., ... Williams, J. (2020). Evaluation of southern ocean cloud in the  
757 hadgem3 general circulation model and merra-2 reanalysis using ship-based  
758 observations. *Atmospheric Chemistry and Physics*, *20*(11), 6607–6630. doi:  
759 10.5194/acp-20-6607-2020
- 760 McCluskey, C. S., Hill, T. C. J., Humphries, R. S., Rauker, A. M., Moreau, S.,  
761 Strutton, P. G., ... DeMott, P. J. (2018). Observations of ice nucleating  
762 particles over Southern Ocean waters. *Geophysical Research Letters*, *45*(21),  
763 11,989-11,997. doi: 10.1029/2018GL079981
- 764 McCoy, D. T., Burrows, S. M., Wood, R., Grosvenor, D. P., Elliott, S. M., Ma,  
765 P.-L., ... Hartmann, D. L. (2015). Natural aerosols explain seasonal and spa-  
766 tial patterns of Southern Ocean cloud albedo. *Science Advances*, *1*(6). doi:  
767 10.1126/sciadv.1500157
- 768 McErlich, C., McDonald, A., Schuddeboom, A., & Silber, I. (2021). Comparing  
769 satellite- and ground-based observations of cloud occurrence over high south-  
770 ern latitudes. *Journal of Geophysical Research: Atmospheres*, *126*(6). doi:  
771 <https://doi.org/10.1029/2020JD033607>
- 772 McFarquhar, G. M., Bretherton, C., Marchand, R., Protat, A., DeMott, P. J.,  
773 Alexander, S. P., ... McDonald, A. (2020). Observations of clouds, aerosols,  
774 precipitation, and surface radiation over the southern ocean: An overview of  
775 capricorn, marcus, micre and socrates. *Bulletin of the American Meteorological*  
776 *Society*, 1 - 92. doi: 10.1175/BAMS-D-20-0132.1
- 777 Modini, R. L., Frossard, A. A., Ahlm, L., Russell, L. M., Corrigan, C. E., Roberts,  
778 G. C., ... Leaitch, W. R. (2015). Primary marine aerosol-cloud interactions off  
779 the coast of California. *Journal of Geophysical Research: Atmospheres*, *120*(9),

- 780 4282-4303. doi: 10.1002/2014JD022963
- 781 Petters, M. D., & Kreidenweis, S. M. (2007). A single parameter representation  
782 of hygroscopic growth and cloud condensation nucleus activity. *Atmospheric*  
783 *Chemistry and Physics*, 7(8), 1961–1971. Retrieved from [http://www.atmos-](http://www.atmos-chem-phys.net/7/1961/2007/)  
784 [chem-phys.net/7/1961/2007/](http://www.atmos-chem-phys.net/7/1961/2007/) doi: 10.5194/acp-7-1961-2007
- 785 Prather, K. A., Bertram, T. H., Grassian, V. H., Deane, G. B., Stokes, M. D., De-  
786 Mott, P. J., ... Zhao, D. (2013). Bringing the ocean into the laboratory to  
787 probe the chemical complexity of sea spray aerosol. *Proceedings of the National*  
788 *Academy of Sciences*, 110(19), 7550–7555. doi: 10.1073/pnas.1300262110
- 789 Pruppacher, H. R., Klett, J. D., & Wang, P. K. (1998). *Microphysics of clouds and*  
790 *precipitation*. Taylor & Francis.
- 791 Quinn, P. K., Coffman, D. J., Johnson, J. E., Upchurch, L. M., & Bates, T. S.  
792 (2017). Small fraction of marine cloud condensation nuclei made up of sea  
793 spray aerosol. *Nature Geoscience*, 10, 674–679. doi: 10.1038/ngeo3003
- 794 Romps, D. M. (2017). Exact expression for the lifting condensation level. *Journal of*  
795 *the Atmospheric Sciences*, 74(12), 3891-3900. doi: 10.1175/JAS-D-17-0102.1
- 796 Russell, L. M., Lenschow, D. H., Laursen, K. K., Krummel, P. B., Siems, S. T.,  
797 Bandy, A. R., ... Bates, T. S. (1998). Bidirectional mixing in an ace 1 marine  
798 boundary layer overlain by a second turbulent layer. *Journal of Geophysical*  
799 *Research: Atmospheres*, 103(D13), 16411–16432.
- 800 Schuddeboom, A., Varma, V., McDonald, A. J., Morgenstern, O., Harvey, M., Par-  
801 sons, S., ... Furtado, K. (2019). Cluster-based evaluation of model compen-  
802 sating errors: A case study of cloud radiative effect in the Southern Ocean.  
803 *Geophysical Research Letters*, 46(6), 3446-3453. doi: 10.1029/2018GL081686
- 804 Spearman, C. (1904). The proof and measurement of association between two  
805 things. *The American Journal of Psychology*, 15(1), 72–101. Retrieved from  
806 <http://www.jstor.org/stable/1412159>
- 807 Trenberth, K. E., & Fasullo, J. T. (2010). Simulation of present-day and twenty-  
808 first-century energy budgets of the Southern Oceans. *Journal of Climate*,  
809 23(2), 440–454.
- 810 Truong, S. C. H., Huang, Y., Lang, F., Messmer, M., Simmonds, I., Siems, S. T., &  
811 Manton, M. J. (2020). A climatology of the marine atmospheric boundary  
812 layer over the southern ocean from four field campaigns during 2016–2018.  
813 *Journal of Geophysical Research: Atmospheres*, 125(20), e2020JD033214. doi:  
814 <https://doi.org/10.1029/2020JD033214>
- 815 Twomey, S. (1977). The influence of pollution on the shortwave albedo of  
816 clouds. *Journal of the Atmospheric Sciences*, 34(7), 1149-1152. doi:  
817 10.1175/1520-0469(1977)034<1149:TIOPOT>2.0.CO;2
- 818 Vergara-Temprado, J., Miltenberger, A. K., Furtado, K., Grosvenor, D. P., Ship-  
819 way, B. J., Hill, A. A., ... Carslaw, K. S. (2018). Strong control of southern  
820 ocean cloud reflectivity by ice-nucleating particles. *Proceedings of the National*  
821 *Academy of Sciences*, 115(11), 2687–2692. doi: 10.1073/pnas.1721627115
- 822 Wofsy, S. C. (2011). Hiaper pole-to-pole observations (hippo): fine-grained, global-  
823 scale measurements of climatically important atmospheric gases and aerosols.  
824 *Philosophical Transactions of the Royal Society A: Mathematical, Physical and*  
825 *Engineering Sciences*, 369(1943), 2073-2086. doi: 10.1098/rsta.2010.0313
- 826 Yin, B., & Albrecht, B. A. (2000). Spatial variability of atmospheric boundary layer  
827 structure over the eastern equatorial pacific. *Journal of Climate*, 13(9), 1574 -  
828 1592. doi: 10.1175/1520-0442(2000)013<1574:SVOABL>2.0.CO;2
- 829 Young, I. R., Zieger, S., & Babanin, A. V. (2011). Global trends in wind speed and  
830 wave height. *Science*, 332(6028), 451–455. doi: 10.1126/science.1197219