

1 **Variations of the Heppner-Maynard Boundaries on Northern**  
2 **Hemisphere Measured by SuperDARN During the Extremely Radial**  
3 **IMFs**

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15 **Key point**

- 16 ● The averaged HMB midnight latitude driven by radial IMF is measured.  
17 ● The seasonal variation of  $B_X$ -HMB correlation coefficients is different with  
18 the existence of single lobe reconnection.  
19 ● The radial IMF effects on the S-M-I coupling sometimes should not be  
20 ignored, especially for the anti-sunward periods during the wintertime.

21 **Abstract**

22 The Heppner-Maynard Boundary (HMB) represents the equatorward extent of the  
23 ionospheric convection pattern and can be used as a proxy for the low latitude of the  
24 auroral oval. We present a statistical study of the radial interplanetary magnetic field  
25 (IMF) effects on the HMB midnight latitude calculated from SuperDARN  
26 measurements between January 2002 and December 2017. We found the average  
27 values of HMB midnight latitude during both sunward and anti-sunward radial IMF  
28 are higher than  $65.5^\circ$ . There is a negative correlation between the magnitude of  $B_X$   
29 and HMB midnight latitude, although this effect is not obvious. Moreover, the  
30 seasonal variation of  $B_X$ -HMB correlation coefficients is different with the existence  
31 of single lobe reconnection. At the anti-sunward radial period, the correlation  
32 coefficient is up to 0.54 in wintertime. It would be caused by the enhanced lobe  
33 reconnection rates, which related to the special configuration between the solar-wind  
34 and magnetopause. This is the first long-term statistical study focused on HMB during  
35 radial IMF conditions in the context of solar wind-magnetosphere-ionosphere  
36 coupling. The results suggest that the effect of IMF  $B_X$  should not be ignored in the  
37 northern hemisphere wintertime especially during the anti-sunward radial IMF  
38 conditions.

39 **Keywords:** radial IMF, Heppner-Maynard Boundary, IMF  $B_X$ , S-M-I coupling

40 **Plain Language Summary**

41 The magnetic reconnection between solar wind and magnetosphere can drive a  
42 convection structure in the high-latitude ionosphere. Heppner and Maynard developed  
43 a method to calculate the equatorward boundary of the ionospheric convection, which  
44 is called Heppner-Maynard Boundary (HMB). HMB provides an import clue on solar  
45 wind-magnetosphere-ionosphere coupling. In this paper, we report the effects of the  
46 radial interplanetary magnetic field (IMF) on HMB in northern hemisphere. Our  
47 statistical results indicate that there is a negative and weak correlation between the  
48 strength of IMF  $B_X$  and HMB magnitude at midnight. The existence of reconnection  
49 happened on northern hemispheric high-latitude magnetosphere can affect the IMF  
50  $B_X$ -HMB correlation coefficients, and also makes seasonal variation of coefficients  
51 different. Moreover, the effect of IMF  $B_X$  should not be ignored in northern  
52 hemisphere winter time during the negative radial IMF conditions.

53 **1. Introduction**

54 The solar wind-magnetosphere-ionosphere (S-M-I) coupling controlled mainly by  
55 the magnetic reconnection occurring at the dayside magnetopause and in the  
56 magnetotail is a key issue in space physics. It is well-known that the interplanetary  
57 magnetic field (IMF)  $B_Z$  component, which plays a crucial role in the reconnection  
58 rates and energy transmission from the solar wind to the magnetosphere. For  
59 southward IMF  $B_Z$ , large energy input could occur on the dayside magnetopause due  
60 to the low latitude reconnection between the IMF and the geomagnetic field (Akasofu,  
61 1981; Dungey, 1961; Lu et al., 2013). During northward IMF, the energy and mass  
62 input is typically thought to be low (e.g., Lu et al., 2013). The amount of the open  
63 magnetic flux is modulated by the dynamic reconnection processes. The open-closed  
64 field line boundary (OCB) which is a surrounding region also called polar cap is the  
65 interface between geomagnetic field lines that are open to solar wind and closed to the  
66 opposite hemisphere (e.g., Lockwood, 1998; Rae et al., 2004; Kabin et al.,  
67 2004; Wang et al., 2014). The latitudinal location and movement of the OCB is of

68 importance as it reflects the balance of magnetic reconnection on the magnetopause  
 69 and in the magnetotail, and thus, reveals information about the total amount of open  
 70 magnetic flux in the magnetosphere (Milan et al., 2007; Lu et al., 2019).

71 A function can describe the rate of expanding-contracting polar cap paradigm  
 72 (ECPC) at a given time (Imber et al., 2013a, 2013b) :

$$73 \quad \frac{dF_{PC}}{dt} = \Phi_D - \Phi_N = \frac{d}{dt} \int_{PC} \mathbf{B} \cdot dS \quad (1)$$

74 where  $F_{PC}$  depicts the open magnetic flux in the polar cap,  $\Phi_D$  marks the dayside  
 75 reconnection rate,  $\Phi_N$  is the reconnection rate in the nightside, the ionospheric  
 76 magnetic field vector is represented by  $\mathbf{B}$ , and the integral is taken over the polar cap  
 77 area. This function suggests that the reconnection in the dayside makes the polar cap  
 78 expand and the reconnection in the magnetotail makes the opposite change.  
 79 Therefore, the dynamic balance between the reconnections occurring in dayside and  
 80 nightside is kept by a phenomenon known as steady magnetospheric convection. The  
 81 observations of 25 nightside reconnection events provide evidence for the dynamic  
 82 changes (Milan et al., 2007).

83 While direct measurement of  $\Phi_N$  is difficult, many studies focus on  $\Phi_D$ . Kan &  
 84 Lee (1979) gives an early function:

$$85 \quad \Phi = B_S V_S \left(\frac{\theta}{2}\right) l_0 \quad (2)$$

86 where  $\Phi$  represents the polar cap potential as an approximation of  $\Phi_D$ ,  $B_S$  marks the  
 87 IMF southward component,  $V_S$  depicts the plasma inflow speed,  $\theta$  is the IMF clock  
 88 angle, and  $l_0 = 7R_E$ , representing the effective length of the reconnection site. In other  
 89 studies,  $B_S$  can be replaced by  $B_T$  or  $B_{YZ}$  ( it is noticed that here  $B_T$  represents the  
 90 modulus of the  $B_Y$  and  $B_Z$  components rather than  $\mathbf{B}$  total.  $B_{YZ}^2 = B_Y^2 + B_Z^2$  )  
 91 (Borovsky E., 2008; Milan et al., 2012; Newell et al., 2007; Scurry & Russell, 1991;  
 92 Temerin & Li, 2006; Vasyliunas et al., 1982; Wygant et al., 1983). In these previous  
 93 studies, the contribution of IMF  $B_X$  was not taken into account in the S-M-I coupling.  
 94 Even in widely used MHD models, such as SWMF/BATS-R-US provided by  
 95 Community Coordinated Modeling Center (CMCC),  $B_X$  can only be set to 0 or a  
 96 constant (e.g., (Rae et al., 2004; Raeder et al., 2001; Wang et al., 2014; Lu et al.,  
 97 2019) to avoid the non-zero magnetic field in the outer boundary.

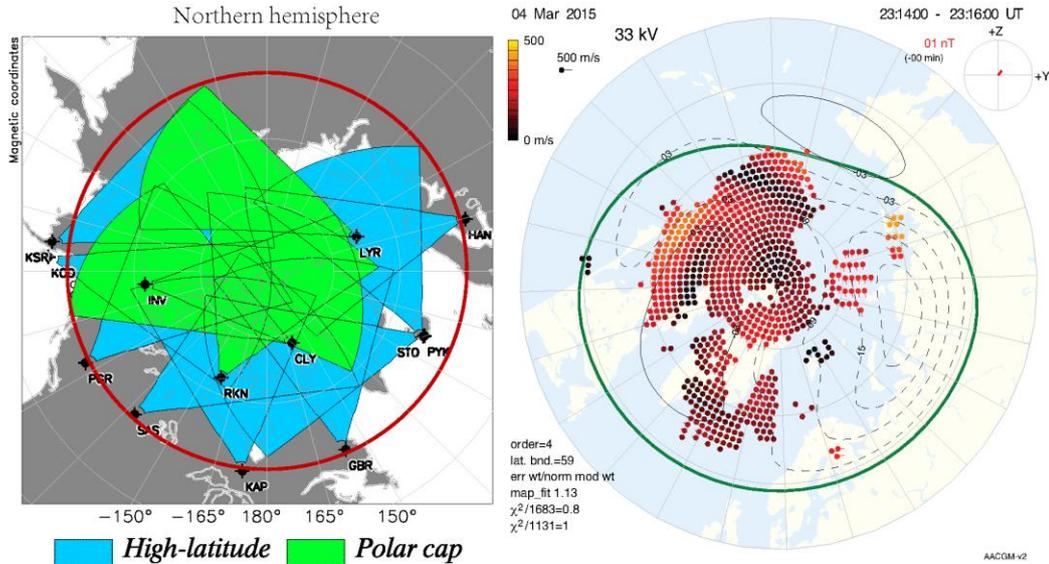
98 As mentioned above, previous works completely ignored the contribution of  $B_X$   
 99 on the S-M-I coupling process. However, Belenkaya (1998) found that the  
 100 ionospheric convection patterns changed with the angle of IMF in the X-Z plane,  
 101 which was depended on 4 different topological structures between solar wind and  
 102 magnetosphere. Other studies indicated that  $B_X$  can affect the asymmetry of  
 103 reconnection for large scale coupling system (e.g., Peng et al., 2010; Tang et al.,  
 104 2013). In addition,  $B_X$  will be transformed into the magnetosheath field  $B_Z$  with  
 105 different polarities in the southern and northern hemispheres during radial  
 106 IMF ( $|B_X|/|B_t| \geq 0.9$ ) (Pi et al., 2017). There is no doubt that this change in direction can  
 107 affect magnetic reconnection rates. A radial IMF, which is dominated by  $B_X$ , is a  
 108 special and stable period of IMF normally with low plasma density, low velocity, low  
 109 temperature, low dynamic pressure, and weak magnetic field (Pi et al., 2014). Radial  
 110 IMF periods account for an important duration of total time (~16% for purely radial  
 111 IMF periods and 10-15% for long periods larger than 4h) (Suvorova et al., 2010; Pi et  
 112 al., 2014). Here, we aim to find out the relationship between radial IMFs and large  
 113 scale S-M-I coupling, and when  $B_X$  effect has to be considered in the coupling.

114 The Heppner-Maynard Boundary (HMB) represents the equatorward extent of the  
 115 ionospheric convection pattern. The characteristic shape of the boundary is

116 determined by Dynamics Explorer 2 (DE2) electric field data (Heppner & Maynard,  
117 1987) and is pole centered. That means the size of the boundary can be represented by  
118 the latitude at HMB midnight. For a large amount of statistical studies based on good  
119 Super Dual Auroral Radar Network (SuperDARN) data coverage, the HMB midnight  
120 latitude could be used as a proxy for the latitude of the auroral oval in the northern  
121 hemisphere (Boakes et al., 2008; Imber et al., 2013b). HMB is calculated by  
122 combining ionospheric model data (excluding  $B_X$ ) with SuperDARN observation  
123 (including  $B_X$ ). In this way, the influence of  $B_X$  on auroral oval can be revealed by  
124 HMB indirectly. In this study, we focus on radial IMF events and study the  
125 relationship between HMB measured by SuperDARN and  $B_X$ . That is, in our events,  
126  $B_X$  dominates IMF with almost zero  $B_Y$  and  $B_Z$  components. This kind of parameter  
127 induces a very pure IMF  $B_X$  condition, thus the interferences that come from  $B_Y$  and  
128  $B_Z$  components are eliminated effectively.

## 129 **2. Data source and methodology**

130 In this work, IMF data is obtained from the OMNI database. Previous works have  
131 demonstrated that, following a change of the IMF, 20 min is long enough to produce  
132 a reconfiguration for the magnetosphere-ionosphere (M-I) coupling system (Murr &  
133 Hughes, 2007; Grocott & Milan, 2014). Here we also choose 20 mins as the  
134 minimum timescale in this study. Radial IMF durations of 20, 30, 40, 50, 60 and 90  
135 mins are selected as time bins, in each bin the radial IMF duration should be  $\geq 90\%$   
136 time of bin. This criterion means that, for 40 mins' time bin, the total radial time is  
137 36 mins at least. For SuperDARN data, the ionospheric echoes from polar cap and  
138 high-latitude SuperDARN radars are considered as the database. Thomas &  
139 Shepherd (2018) shown that under the stronger solar wind driving condition, the  
140 inclusion of mid-latitude radar data at the equatorward extent of the ionospheric  
141 convection can increase the measured cross-polar cap potential ( $\Phi$  in function 2) by  
142 40%, but the increase is very small for weak solar wind driving condition. To keep  
143 the same SuperDARN coverage in different years and to exclude the inaccuracy  
144 caused by extended convection patterns, the mid-latitude radars data are not  
145 included. The selected radars are shown in Fig. 1a. Shadow colored areas are the  
146 field of view of radars, blue represents high-latitude radar, green is polar cap radar.  
147 The red circle marks the geomagnetic latitude at  $60^\circ$ .



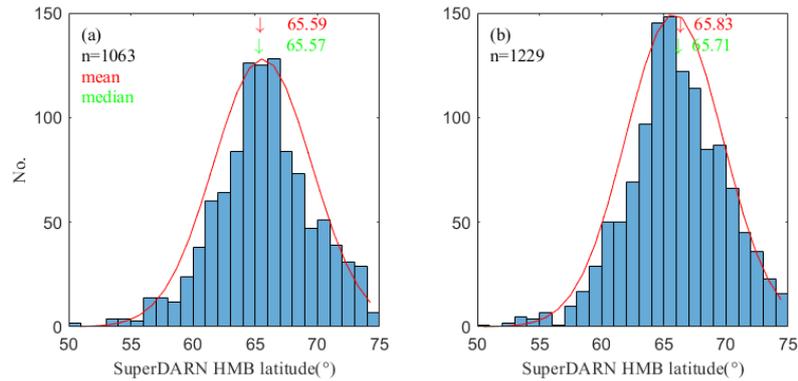
148 Fig. 1. SuperDARN radar coverages and an example map during the radial IMF. (a)  
 149 Radar coverages. (b) An example HMB map. The solid green line is the HMB for this  
 150 map, and the midnight latitude is  $59^\circ$ .

151 We use the Radar Software Toolkit 4.2 (RST4.2), a free authorized and  
 152 powerful software to process SuperDARN data (Barnes & Greenwald, 2005), to  
 153 calculate the HMB boundary from radars' rawacf data. Here we use a standard  
 154 criterion to determining the HMB (Imber et al., 2013a): the velocity threshold is set  
 155 for  $100\text{m/s}$ , and the number threshold of the effective radar backscatter points in a  
 156 convection map is 150. Thus, HMB can be obtained every 2 min.

157 A convection map on the northern hemisphere during a typical radial IMF  
 158 condition is shown in Fig. 1b. The convection map is calculated from the averaged  
 159 radar data assigned to a 2min scan period (2015-03-04 23:14UT-23:16UT). Magnetic  
 160 local noon locates to the top of the figure and dusk is to the left. The colored scatters  
 161 represent fitted ionospheric echoes from radar observations, with the color and the  
 162 vector direction corresponding to the velocity of the plasma flow given by the fourth-  
 163 order spherical harmonic fit. The dashed and solid black lines depict the electrostatic  
 164 potential, the contours of constant electrostatic potential also represent plasma flow  
 165 streamlines for the whole polar ionosphere. The right upper corner marks the IMF  
 166 condition, which shows noticeable small  $B_Y$  and  $B_Z$ . The solid green line is the HMB  
 167 for this map constraining the extent of the mapping, and its midnight latitude is  $59^\circ$ .  
 168 Based on the above criterion, we calculate all the HMB maps between 2002 and 2017,  
 169 and our database can be seen in the appendix.

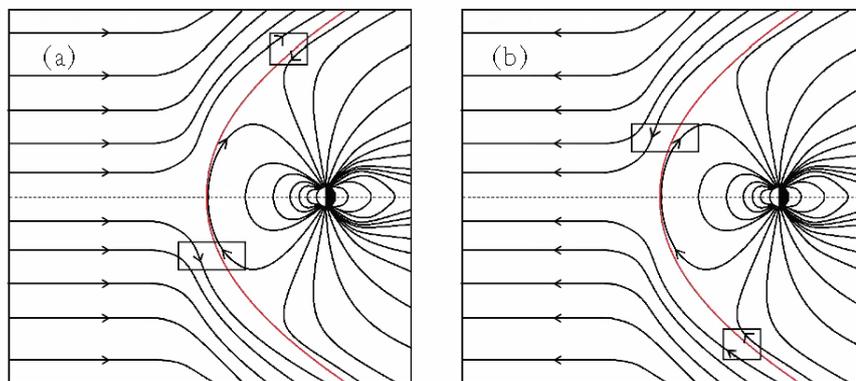
### 170 3.Statistical results and discussion

171 We average the HMB midnight latitudes for each radial IMF event for the  
 172 northern hemisphere and then get 2292 data in total. Fig. 2 is the distribution of the  
 173 HMB midnight latitudes. The x-axis is midnight latitude and the y-axis is the number  
 174 of events. n represents the number of total events, the red (green) numeric value is  
 175 mean (median) value for each panel. Fig. 2a shows HMB latitudes for anti-sunward  
 176  $B_X$  on the northern hemisphere. Event number, mean and median values of the HMB  
 177 at midnight are 1063, 65.59 and 65.57, respectively. Fig. 2b gives the HMB latitudes  
 178 in sunward  $B_X$ . The corresponding event number, mean and median are 1229, 65.83  
 179 and 65.71, respectively. The red line in each panel marks a Gauss Fitting Curve.



180 Fig. 2. The distribution of the HMB midnight latitudes. n is the number of total  
 181 events, the red/green numeric value is the mean/median for each panel. (a) HMB  
 182 latitudes in anti-sunward radial, (b) HMB latitudes in sunward radial.

183 It is shown that, under the radial IMF condition, the majority distributions of  
 184 HMB midnight latitude is around  $65.57^{\circ}$ - $65.83^{\circ}$ , which are apparently higher than the  
 185 average value of  $63.1^{\circ}$  for the substorm periods revealed by Imber et al. (2013b), and  
 186 also larger than the average of  $64^{\circ}$  for the time period from January 1996 to August  
 187 2012 as shown in Imber et al. (2013a). This result signifies that during radial IMFs,  
 188 ionospheric convection pattern is smaller. It suggests that the energy coupling  
 189 efficient inputting from the solar wind into the magnetosphere is evidently weak. This  
 190 is consistent with our understanding of the radial IMF, whose driving effect is weak  
 191 (compared with southward IMF). The HMB midnight latitude for anti-sunward radial  
 192 IMF is smaller than that in the sunward radial IMF. It suggests that the anti-sunward  
 193 radial IMF injects more energy into the ionosphere. The difference of topological  
 194 structure between anti-sunward and sunward radial IMF in large scale S-M-I coupling  
 195 can be seen in Fig3.



196 Fig. 3. Schematic illustration of the topological structure between IMF and  
 197 magnetosphere. The red line indicates the boundary of the magnetosphere on dayside,  
 198 the rectangles are possible locations of magnetic reconnection. (a) anti-sunward  
 199 radial, (b) sunward radial.

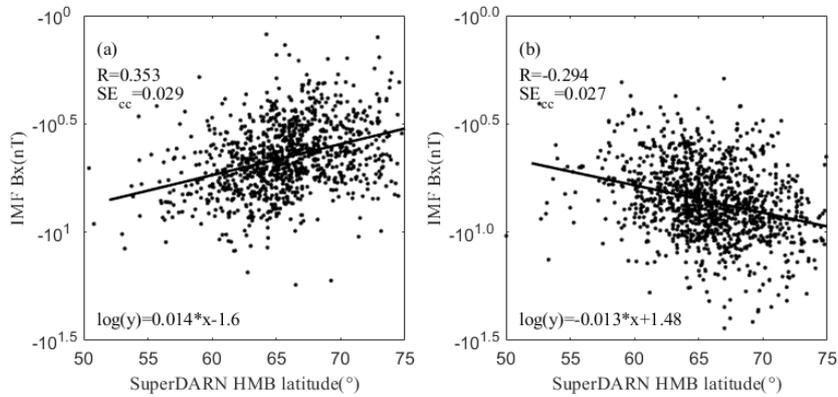
200 Fig. 3 is a schematic illustration in the GSM X-Z plane showing the topological  
 201 structure of radial events between the IMF and the magnetosphere. Fig. 3a (3b) shows  
 202 the case of an anti-sunward (sunward) radial case. The black lines represent the  
 203 magnetic field lines of the solar wind and the Earth. The red line marks the boundary  
 204 of the magnetosphere on the dayside. The rectangles label possible locations of  
 205 magnetic reconnection. A low-latitude reconnection and a single lobe reconnection  
 206 can be found during a radial IMF. It is easy to understand that low-latitude

207 reconnection contributes to both hemispheres while single lobe reconnection can only  
 208 affect one hemisphere. In an ideal situation (zero for  $B_Z$  and  $B_Y$ ), we assume that the  
 209 transformation rate of  $B_X$  to  $B_Z$  equal in northern and southern hemispheres.  
 210 Accordingly, we can infer that the low-latitude reconnection rates almost the same in  
 211 Fig. 3a and 3b.

### 212 3.1 Correlation analysis

213 In S-M-I coupling, the magnetic field in local region of reconnection is an  
 214 important parameter (e.g., Cassak & Shay, 2007). It is difficult to obtain the magnetic  
 215 field characters in the local magnetosheath for all the time. In addition, due to the  
 216 draping effect, IMF  $B_X$  component will be transformed into  $B_Z$  component in the  
 217 magnetosheath (Pi et al., 2017). Therefore, for statistical study, the magnetic field in  
 218 magnetosheath can be reflected by IMF  $B_X$  in radial IMF.

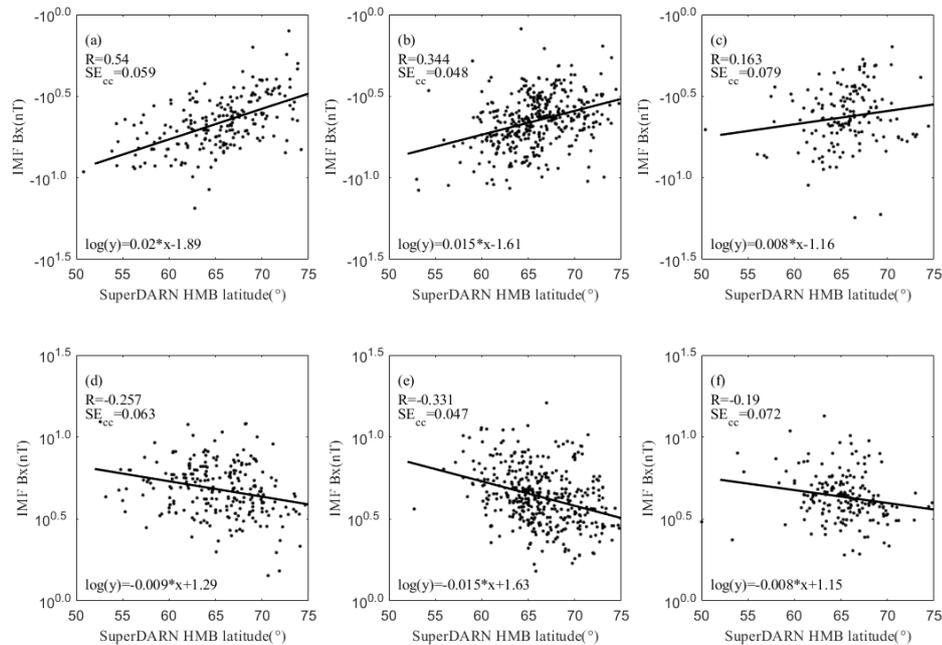
219 In this section, we discuss the correlation between HMB midnight latitude and  $B_X$   
 220 component during radial IMFs. Fig. 4 shows scatter plots of HMB midnight latitudes  
 221 and  $B_X$  magnitude. The x-axis is HMB midnight latitude, and the y-axis is the strength  
 222 of IMF  $B_X$ .  $R$  represents the Pearson correlation coefficient,  $SE_{cc}$  marks the standard  
 223 deviation of correlation coefficient ( $SE_{cc} = \sqrt{(1 - R^2)/(n - 2)}$ ,  $n$  means counts of  
 224 events). The black line represents the linear fitting line and its corresponding  
 225 expression is shown at the bottom of each panel. (e.g., in Fig. 4a left the HMB  
 226 midnight latitude ( $x$ ) and IMF  $B_X$  ( $y$ ) satisfy  $\log(y) = 0.014x - 1.6$ , and  $R$  ( $SE_{cc}$ ) is  
 227 0.353 (0.029)). All anti-sunward radial events are shown in the left panel of Fig.4a,  
 228 and the right figure gives all sunward radial events.



229 Fig. 4. The correlation between HMB midnight latitude and  $B_X$ . (a) all radial events in  
 230 anti-sunward cases. (b) all radial events in sunward cases.

231 In both statistical classifications, a negative correlation between the strength of  
 232  $B_X$  and HMB midnight latitude can be found. The absolute value of the correlation  
 233 coefficient with 0.353 in the anti-sunward radial IMF is larger than that in the  
 234 sunward radial IMF (0.294). Two coefficients both show weak correlation, it  
 235 suggested that  $B_X$  has relatively smaller contribution on the magnetic reconnection  
 236 process. In addition, the AE index is very closed in two cases (anti-sunward 142.2;  
 237 sunward 119.4), we can infer that difference in nightside reconnection rates is small.  
 238 Therefore, difference in correlation coefficients is mainly caused by the asymmetric  
 239 topological structure between the IMF and the magnetosphere (e.g., locations of  
 240 rectangles on Fig. 3).

242 Previous studies have shown that the subsolar X-line location (the region where  
 243 reconnection may occur) can shift poleward from the subsolar point, and the X-line  
 244 location has a seasonal dependence (e.g., Trattner et al., 2007; Hoshi et al., 2018).  
 245 This shift will undoubtedly affect the coupling process. In this section, we study the  
 246 effect of seasonal variations on the correlation between IMF  $B_X$  and HMB midnight  
 247 latitude.



248 Fig. 5. seasonal variations. Top row is plots for anti-sunward radial events, and  
 249 sunward radial events are shown in bottom row. The columns from left to right are  
 250 radial events in winter, equinoxes and summer.

251 Fig. 5 shows the seasonal variations. In Fig.5 the top row is plots for anti-sunward  
 252 radial events, and sunward radial events are shown in the bottom row. The columns  
 253 from left to right are radial events in winter (Dec/Jan), equinoxes (Mar/Apr/Sep/Oct)  
 254 and summer (Jun/Jul). The correlation varies significantly in different panels of Fig.5.  
 255 The correlation in anti-sunward radial events between  $B_X$  and HMB midnight  
 256 latitude is extremely high during wintertime, the coefficient is up to 0.54 (Fig. 5a). In this  
 257 case, the effect of IMF  $B_X$  on HMB is very large for a chaotic system. Statistical  
 258 results suggest that negative IMF  $B_X$  plays an important role in the S-M-I coupling  
 259 process during wintertime. Top row figures of Fig. 5 show that, with the northward  
 260 shift of the subsolar point, the correlation coefficients in anti-sunward radial events  
 261 become weaker from winter to summer. In the bottom figures of Fig. 5, for positive  
 262  $B_X$  events, the correlation coefficients become stronger with the equatorward moving  
 263 of subsolar point. However, the correlation between positive  $B_X$  cases is weak in all  
 264 seasons. As shown in right figures of Fig. 5, very weak correlation can be found in  
 265 both directions of  $B_X$  during summertime. In summer, solar radiation has a great  
 266 influence on the photoionization effect on the dayside ionosphere, which causes the  
 267 enhancement of the dayside ionospheric conductivity. It is expected that the dayside  
 268 R1 current will increase, accordingly, the magnetosphere shrinks on the dayside and  
 269 expands on the nightside (Ohtani et al., 2014). The energy transmission from the solar

270 wind to the Earth will be affected by the morphological structure change of  
271 magnetopause (Jing et al., 2014; Merkin et al., 2005; Ohtani et al., 2014; Raeder et al.,  
272 2001). That's the reason for a weaker correlation in summer.

273 For positive  $B_X$  events, the topological structure is similar to the southward IMF  
274 on the northern hemisphere. When the northern hemisphere is in summer (winter )  
275 and the north (south ) pole tilts toward the Sun, the reconnection location under finite  
276 dipole tilt shifts toward the winter (summer ) hemisphere (GSM coordinate) (Hoshi et  
277 al., 2018; Komar et al., 2014; Russell et al., 2003; Trenchi et al., 2008), and the  
278 reconnection rates can be reduced by this kind of X-line shift (Borovsky et al., 2008;  
279 Cassak & Shay, 2007). Furthermore, the  $SE_{cc}$  in each panel is small, which suggests  
280 that the correlation coefficients have enough credibility. Therefore, a negative  
281 correlation can be found between the correlation coefficient and the absolute value of  
282 subsolar point latitude in sunward radial.

283 It is more interesting in negative  $B_X$  events. It is not straightforward to explain  
284 why the anti-sunward radial events have such a high correlation in winter. We  
285 separate each negative  $B_X$  events in different seasons according to different directions  
286 of  $B_Z$  (positive, negative) and calculate the correlation coefficients respectively. The  
287 correlation coefficients for positive (negative)  $B_Z$  are listed as follows: winter 0.538  
288 (0.567); equinoxes 0.205 (0.471). We can see that the effect of  $B_Z$  is more evident in  
289 equinoxes. As mentioned in section 1, due to the draping effect, the IMF  $B_X$  will be  
290 transformed into the magnetosheath magnetic field  $B_Z$  with different polarities on the  
291 southern and northern hemispheres. In simple geometric theory, when the  
292 transformation line with the same (opposite) direction as the IMF  $B_Z$ , the influence of  
293 IMF  $B_X$  will be larger (smaller). It suggests that the transformation is mainly near the  
294 middle plane, and mainly affects low latitude reconnection. The correlation  
295 coefficient shows almost the same in winter, which seems to indicate that the main  
296 reason for such large correlation coefficient in negative  $B_X$  events during wintertime  
297 is not due to the low latitude reconnection.

298 As shown in Fig. 3, in anti-sunward radial condition, the northern hemisphere  
299 will be affected by both low-latitude and single lobe reconnection. We check the  
300 averaged AE index values (winter 149.7; equinoxes 144.6), which suggests that the  
301 nightside reconnection rate is almost identical inferred from AE. As mentioned above,  
302 we speculate that the dramatically high correlation in anti-sunward radial condition  
303 during wintertime is mainly due to the effect of lobe reconnection. There are some  
304 issues with lobe reconnection. The energy transport related to lobe reconnection is  
305 usually much less than the one associated with the low latitude magnetopause  
306 reconnection, the related region is also much smaller, usually limited to above  $80^\circ$   
307 MLAT on the dayside (06–18 MLT) in the ionospheric height (Reistad et al., 2019).  
308 With the enhancement of ionospheric conductivity, the lobe reconnection rate  
309 increases correspondingly (Paschmann et al., 2003; Reistad et al., 2019). This  
310 indicates that although single lobe reconnection cannot produce open magnetic lines  
311 (Imber et al., 2006), the special configuration between the magnetopause and solar  
312 wind during anti-sunward radial period seems to have a great impact on lobe  
313 reconnection rate during wintertime and inject much energy into the polar ionosphere.  
314 What occurred for anti-sunward radial in winter time on northern hemisphere is  
315 interesting and not confirmed. We will take a further study on this issue in future  
316 work.

#### 317 4. Summary and expectation

318 In this study, we select radial IMF events to study the correlation between IMF  
319  $B_X$  component and HMB, because the radial IMF can ideally ignore most of the other  
320 factors. The main findings in this work can be summarized as follows:

- 321 1. The solar wind is continuously stable and weak during radial IMF, and under  
322 this condition, HMB midnight latitude is higher than that during both long-  
323 term period (1997-2012) and high geomagnetic activity cases (Imber et al.,  
324 2013a, 2013b).
- 325 2.  $B_X$  has a small effect on the coupling process, and a negative correlation can  
326 be found between the strength of  $B_X$  and HMB midnight latitude. The  
327 correlation coefficient in the northern hemisphere is only 0.29~0.35.
- 328 3. During northern hemispheric wintertime, the correlation coefficient between  
329 the strength of  $B_X$  and HMB midnight latitude is up to 0.54 in anti-sunward  
330 radial events. In such situation, the effect of  $B_X$  in the coupling process is  
331 important and can't be ignored.
- 332 4. When the direction of the magnetic field line in magnetosheath is antiparallel  
333 to that of the geomagnetic field line in the low latitudes and parallel in the  
334 lobe region, the correlation coefficients is larger in equinoxes, but is smaller  
335 during the winter and summer. When the magnetic field direction in  
336 magnetosheath is anti-parallel with the geomagnetic field in low latitudes and  
337 also anti-parallel in the lobe region, a negative correlation between  
338 correlation coefficients and latitude of subsolar point can be found.

339 As a statistical study, we find an abnormal large correlation coefficient between  
340  $B_X$  and HMB in the winter northern hemisphere when  $B_X$  is negative. Although we  
341 present a possible explanation, further research is needed.

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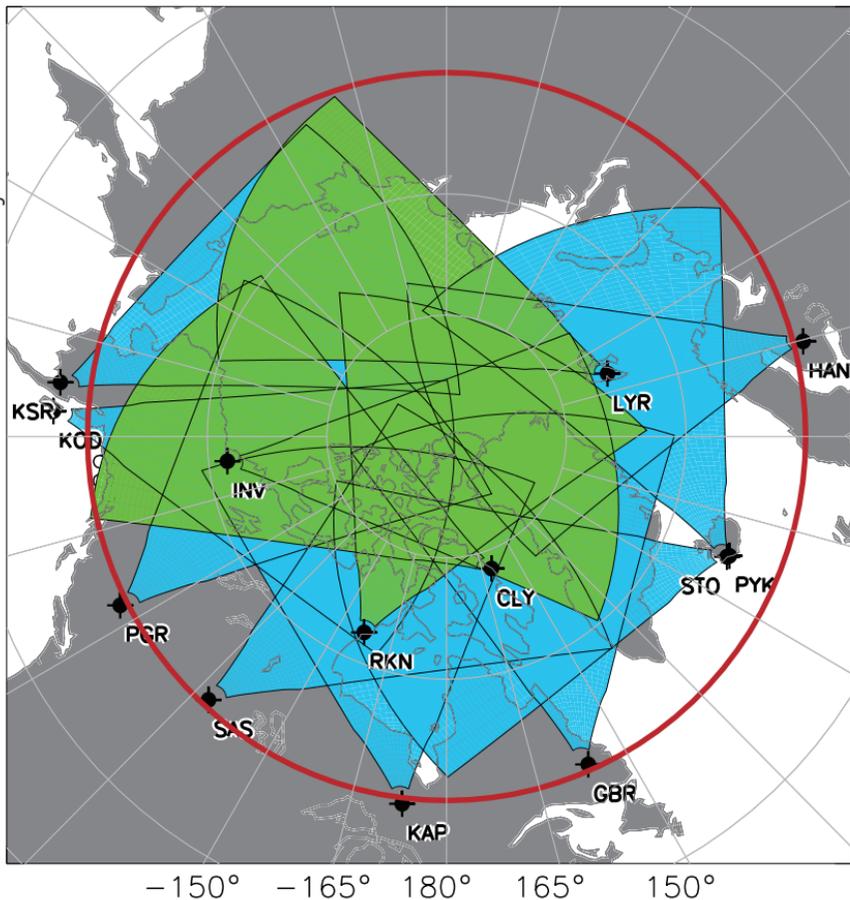
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Figure 1.

# Northern hemisphere

Magnetic coordinates

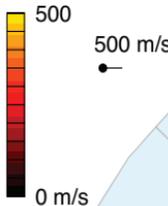


*High-latitude*     *Polar cap*

04 Mar 2015

33 kV

23:14:00 - 23:16:00 UT



order=4  
lat. bnd.=59  
err wt/norm mod wt  
map\_fit 1.13  
 $\chi^2/1683=0.8$   
 $\chi^2/1131=1$

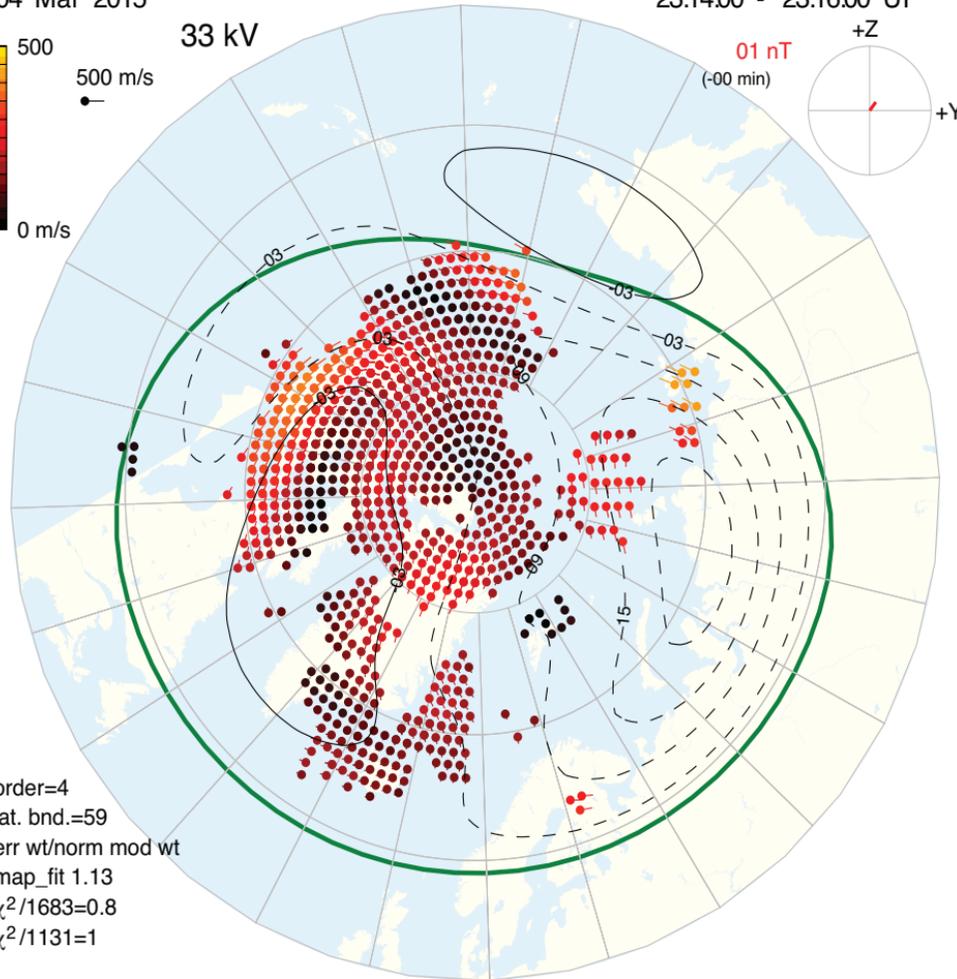


Figure 2.

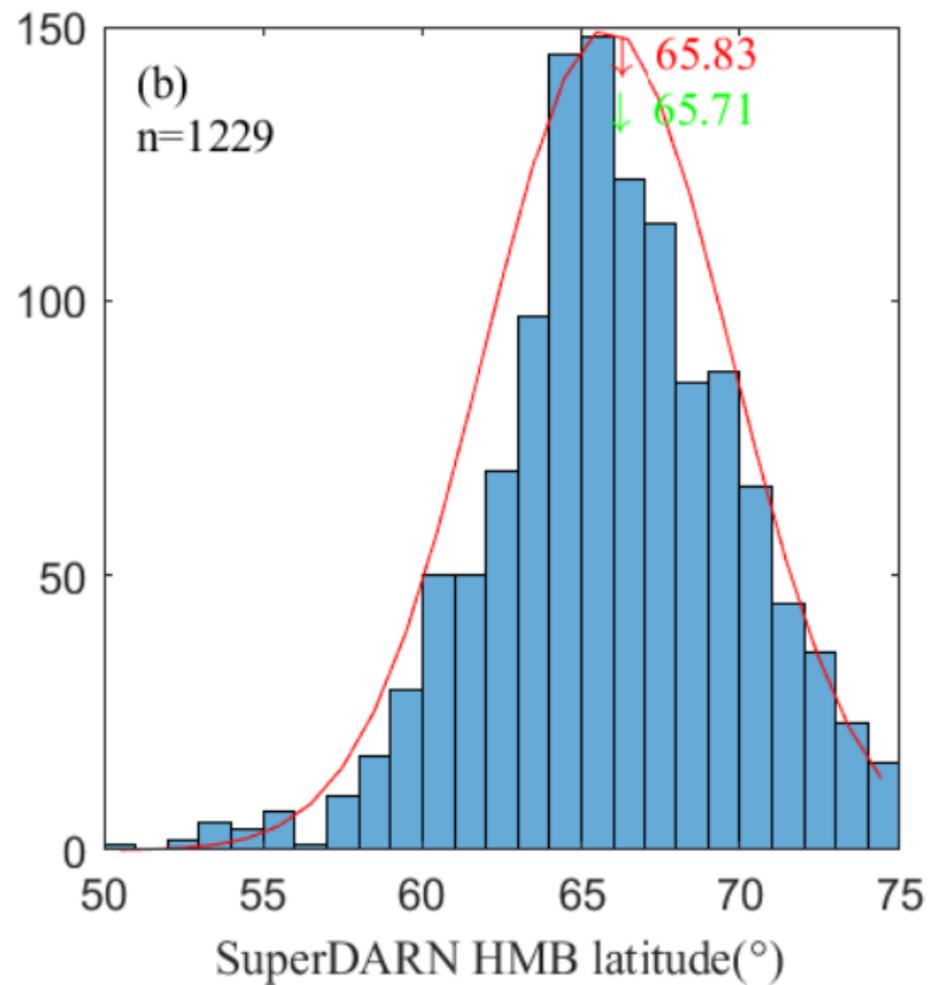
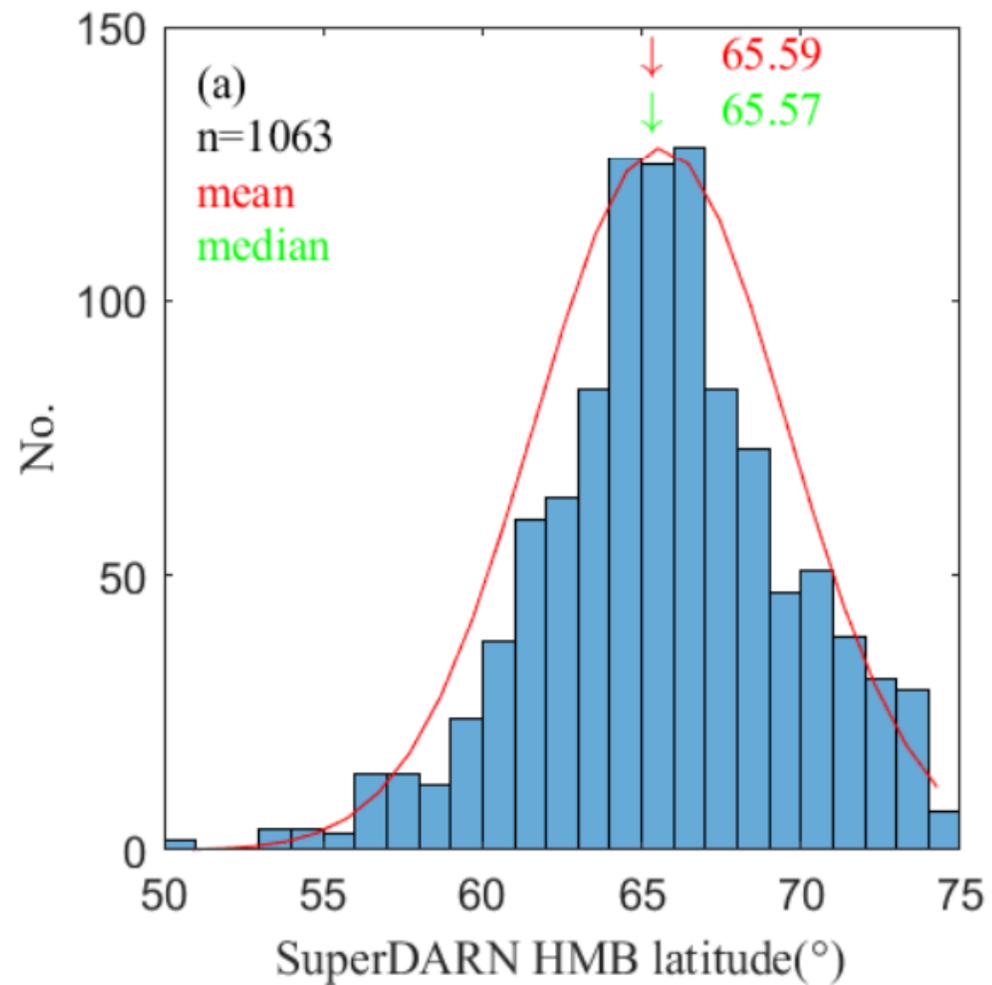
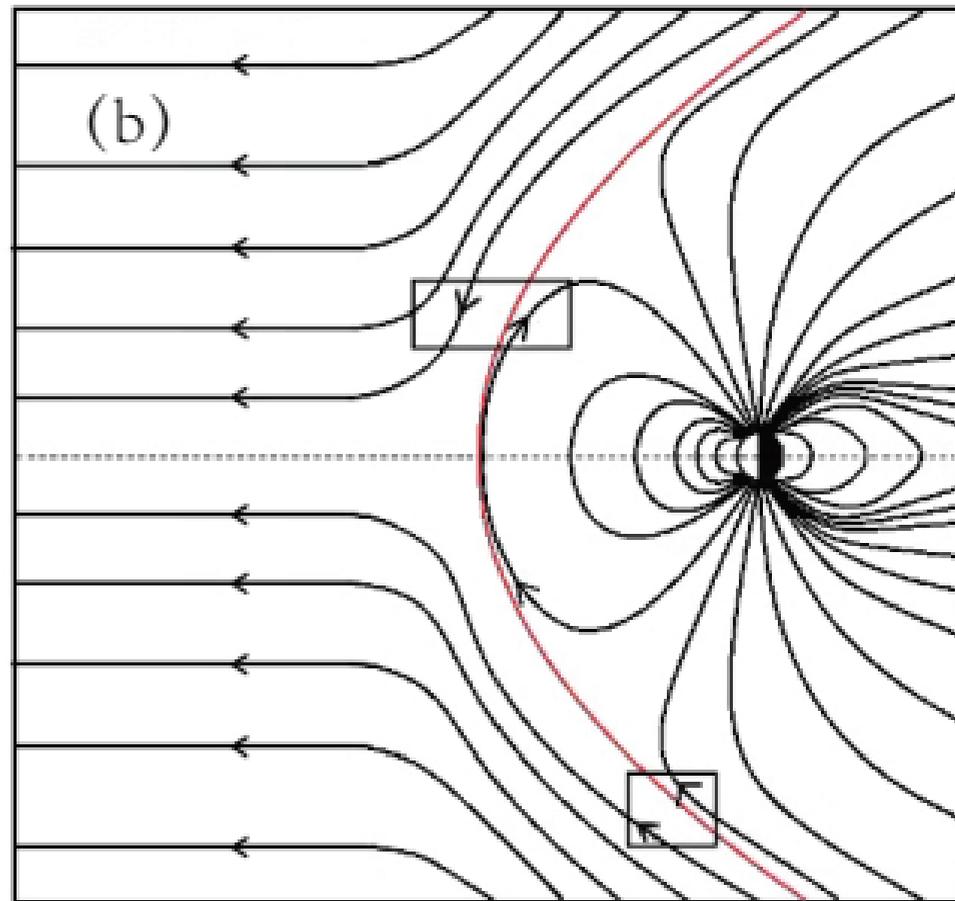
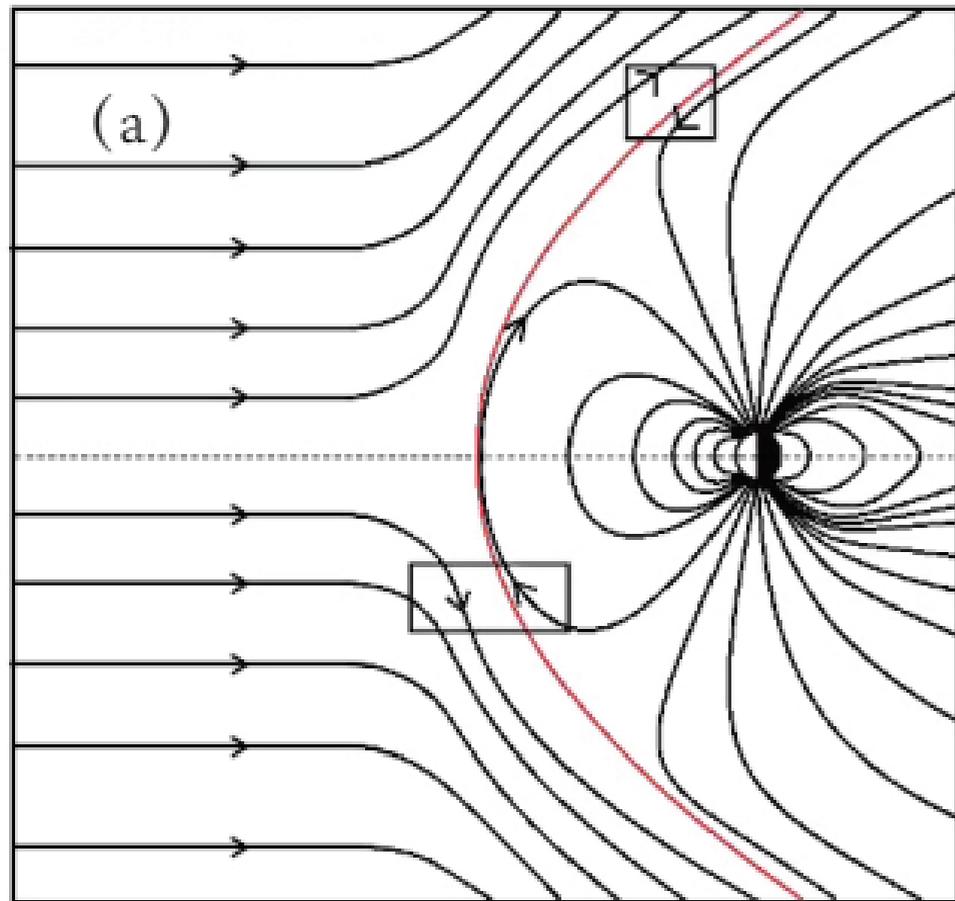


Figure 3.



**Figure 4.**

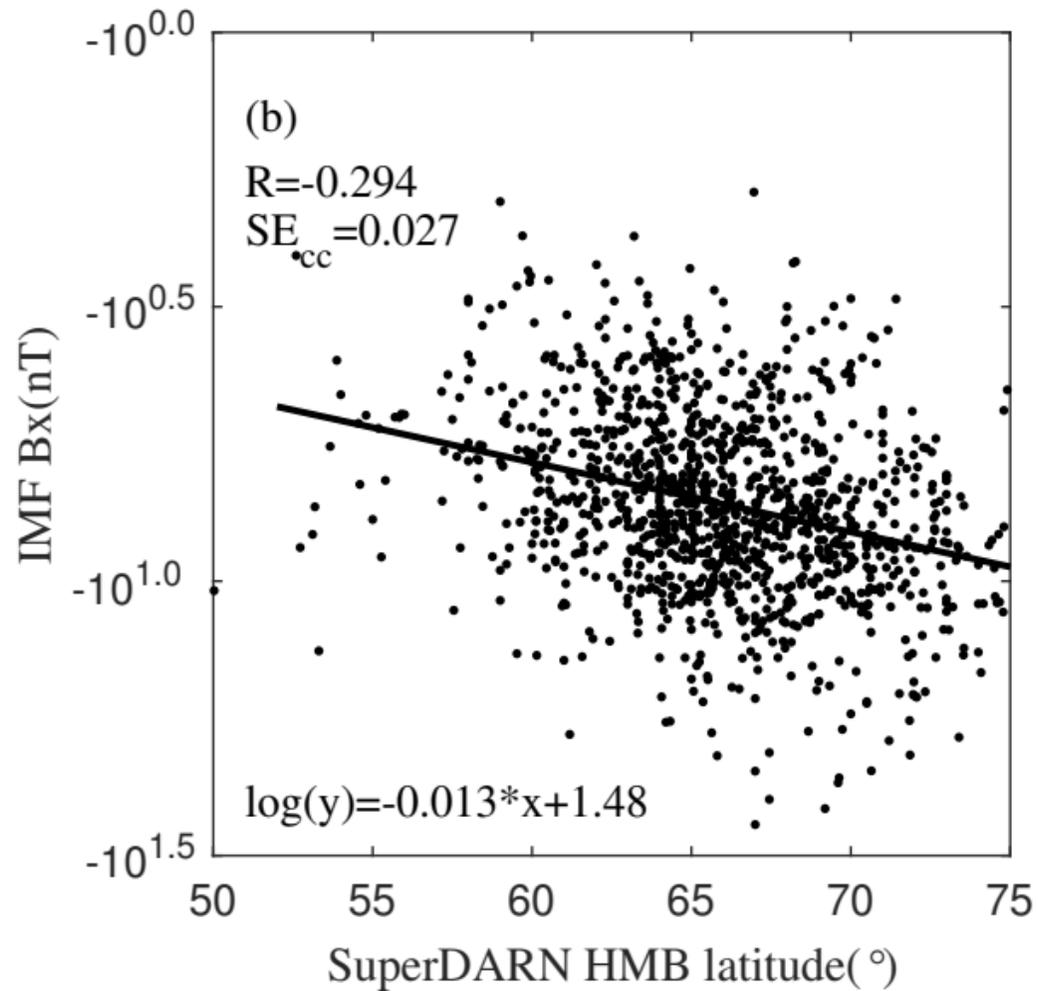
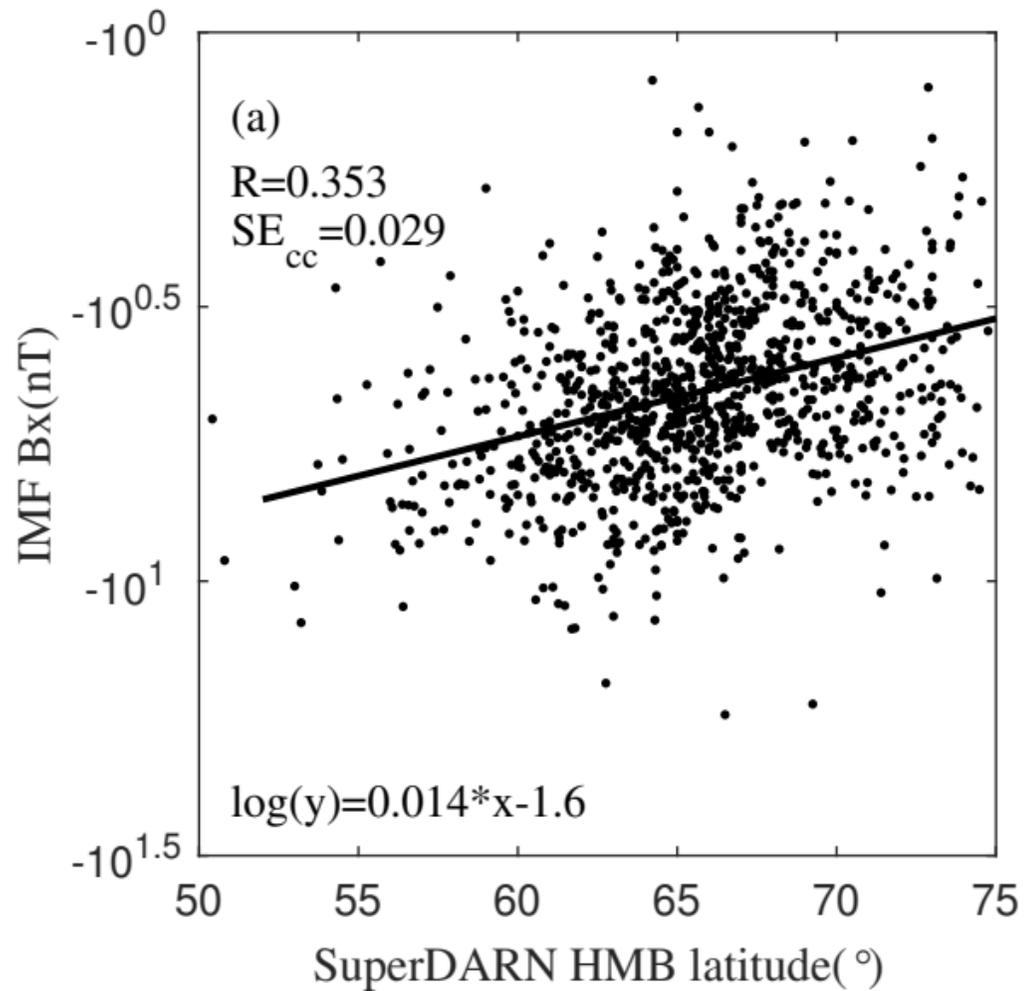


Figure 5.

