

# Mineral-melt partitioning of redox-sensitive elements

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## Abstract

Elements with variable valence state (i.e. redox-sensitive) often show contrasting mineral/melt partition coefficients as a function of oxygen fugacity ( $fO_2$ ) in magmatic systems. This is because trace-element incorporation into crystal lattices depends on the charge, size, and crystal-field stabilization energy of atoms, all of which differ greatly between oxidized and reduced species of the same element. This has two critical implications: (1) petrologic/geochemical modelling of partitioning behavior of redox-sensitive trace-elements in magmatic systems requires some knowledge of their oxidation state, and (2) the oxidation state of magmatic systems may be inferred from partitioning relations of redox-sensitive trace elements preserved in mineral and melt phases of rapidly cooled magmas. The advantage of this oxybarometric approach is that mineral/melt partitioning relations are not sensitive to late stage degassing, charge-transfer on quenching, or surficial alteration. In this chapter we discuss the theoretical treatment of experimental mineral/melt partitioning data of redox-sensitive trace elements, and review aspects concerning the partitioning behavior of well-known redox-sensitive elements, including transition metals (Ti, V, Cr, Fe), rare earth elements (Ce, Eu), U, and siderophile elements (Mo, W, Re, and platinum group elements) under planetary magmatic  $fO_2$  conditions.

## 1. Introduction

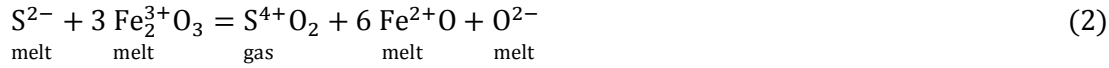
Although it is possible to determine valence-state ratios of redox-sensitive trace elements in natural and synthetic silicate glasses using techniques such X-ray absorption spectroscopy (e.g. Sutton and Newville, 2014), it may be challenging to relate these ratios to the oxygen fugacity ( $fO_2$ ) of a silicate melt. This is because many elements exchange electrons with Fe, the most abundant redox-sensitive element in most magmas, upon quenching. For example, a mid-ocean ridge basaltic (MORB) melt at 1400 °C equilibrated at  $fO_2$  equivalent to FMQ-1.7 (i.e. 1.7 log units below the  $fO_2$  defined by the fayalite-magnetite-quartz buffer at the same temperature) has  $Cr^{2+}/\Sigma Cr \sim 0.5$  (where  $\Sigma Cr = Cr^{2+} + Cr^{3+}$ ), but upon quenching to a glass all the  $Cr^{2+}$  is converted to  $Cr^{3+}$  by reaction with  $Fe^{3+}$  (Berry et al., 2003):



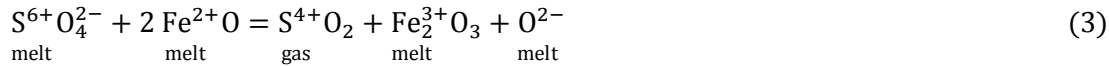
Similar reactions with Fe have been documented for Eu and Ce (Burnham and Berry, 2014; Burnham et al., 2015), but because they require only the transfer of electrons, they are probably unavoidable. Consequently, valence-state ratios of trace elements in natural Fe-bearing glasses may not be representative of their high-temperature melt chemistry.

Provided that these elements are only present at low abundances relative to Fe, the exchange minimally affects the  $Fe^{3+}/\Sigma Fe$  (where  $\Sigma Fe = Fe^{2+} + Fe^{3+}$ ) of the magma. However, this ratio can still be affected by degassing of volatile species (e.g. Burgisser and Scaillet, 2007;

Mathez, 1984; Sato, 1978) as well as surficial oxidation (e.g. Rhodes and Vollinger, 2005). Sulfur degassing appears to be particularly important in the Earth because of the strong capacity of S to either reduce or oxidize Fe depending on the initial oxidation state of the magma (e.g. de Moor et al., 2013; Moussallam et al., 2016). For instance, the  $\text{Fe}^{3+}/\Sigma\text{Fe}$  of reduced magmas, where  $\text{S}^{2-}$  (sulfide) is the dominant melt species, may be further reduced upon degassing of  $\text{SO}_2$ :



whereas the  $\text{Fe}^{3+}/\Sigma\text{Fe}$  of oxidized magmas, where  $\text{S}^{6+}$  (sulfate) is the dominant melt species, may be further oxidized upon degassing of  $\text{SO}_2$ :



While potentially significant in their own right, such effects obscure the primary magmatic redox characteristics. In contrast, high-temperature distributions of trace elements between minerals and melts (i.e. partitioning relations) are readily preserved in rocks that have cooled at moderate or rapid rates, and hence allow the petrologist to ‘see’ melt chemistry that more accurately reflects the properties of the source of a magma. Furthermore, analysis of trace elements in minerals and glasses is readily achieved using benchtop equipment, unlike some spectroscopic techniques such as XANES that require access to the handful of suitably-equipped synchrotron beamlines worldwide. Hence, mineral/melt partitioning allows more convenient insights into the redox states of magmas.

To derive viable oxybarometers from mineral/melt partitioning, experimental data and a thermodynamic framework are required. Experimental equipment, such as gas-mixing tube furnaces, allows access to  $\sim 25$  log units of  $f\text{O}_2$  at 1400 °C; from 10 log units below the iron-wüstite buffer (IW-10) to six log units above the hematite-magnetite buffer (HM+6). For several redox-sensitive elements this is sufficient to encompass the entire transition from one valence state to another (where the redox reaction involves the exchange of 1 electron, such as  $\text{Cr}^{2+}$  to  $\text{Cr}^{3+}$ , the transition occurs over 16 log units, whereas for elements whose redox transition involves the exchange of 8 electrons, such as  $\text{S}^{2-}$  to  $\text{S}^{6+}$ , the transition occurs over 2 log units). This allows the partitioning behavior of the end-members to be determined. Nevertheless, in several cases it is not feasible to do this. For instance, where one valence state is highly volatile (e.g.  $\text{Cr}^{6+}$  or  $\text{Mo}^{6+}$ ), or where an intermediate valence state always exists alongside a higher or lower valence state (e.g.  $\text{V}^{3+}$ ,  $\text{V}^{4+}$  and  $\text{U}^{5+}$ ), or when extreme redox conditions necessary to constrain the partitioning behavior of a valence state alone cannot be achieved in the laboratory. In such cases, it may be necessary to model the expected behavior of a valence state.

Lattice strain theory has become one of the fundamental concepts for understanding trace element partitioning into minerals (Blundy and Wood, 1994; Wood and Blundy, 2003). The essence of the theory is that substitution of a trace element  $i$  onto a cation site in a crystal lattice is energetically neutral if the substituent ion has the same ionic radius ( $r_i = r_0$ ) and charge as the ion it replaces, resulting in a partition coefficient  $D_0$ . Where its ionic charge and/or radius differs, electrostatic work is done in repelling or attracting the adjacent anions (usually oxygen), and this work results in the trace element being less compatible than the one with optimal charge and radius, i.e.  $D_i < D_0$ . The dependence of  $\log D_i$  on  $r_i$  is approximately parabolic;  $D_0$  likewise has a quadratic dependence on atomic number,  $Z$  (Fig. 1). These two controlling variables are the reason that the various valence states of an element can have striking differences in their partitioning behavior. Consider the removal of an electron from  $\text{Eu}^{2+}$ :



which not only increases the charge of the cation, but also (because of its higher effective nuclear charge) cause a contraction of the electronic orbitals, resulting in a decrease in ionic radius. This same principle applies to other major and trace elements capable of more than one valence state. As charge and ionic radius are principal controls on partitioning, it is almost inevitable that different valence states of an element will behave like different elements. Routine analytical techniques do not measure the speciation of redox-sensitive elements, which are usually reported as total concentrations without regard for the important differences between their valence states.

There are however challenges in the application of lattice strain theory to predict partition coefficients. These come mainly from crystal field effects and from the multitude of substitution mechanisms that are possible for a single cation in mineral sites.

Furthermore, there are important crystal-chemical implications for the fact that different valence states, such as  $\text{Cr}^{2+}$  and  $\text{Cr}^{3+}$ , behave like different elements. Crystals must have an overall neutral charge, hence when  $\text{Cr}^{2+}$  substitutes for  $\text{Mg}^{2+}$  in olivine, this criterion is observed, with the simple substitution  $\text{Mg}^{2+} \leftrightarrow \text{Cr}^{2+}$  with the endmember  $\text{Cr}^{2+}$ -olivine forming according to the reaction:



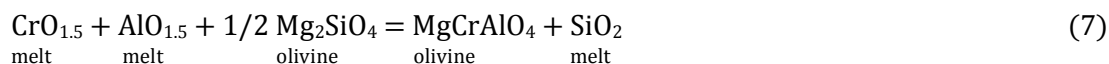
which, because forsterite ( $\text{Mg}_2\text{SiO}_4$ ) saturation requires  $2\text{MgO} + \text{SiO}_2$  and hence a fixed relationship between the activities of  $\text{SiO}_2$  and  $\text{MgO}$ , can equivalently be written:



However, in order to substitute  $\text{Cr}^{3+}$  into the olivine structure, accommodation must be made for the difference in charge. One possibility involves the creation of vacancies [vac] on the octahedral site:  $2 \text{Mg}^{2+} \leftrightarrow 4/3 \text{Cr}^{3+} + 2/3[\text{vac}]$ . The reaction for formation of the corresponding endmember component is:



Another possibility involves coupled substitution with Al and Si on the tetrahedral site:  $\text{Mg}^{2+} + \text{Si}^{4+} \leftrightarrow \text{Cr}^{3+} + \text{Al}^{3+}$ . The reaction for formation of the corresponding endmember component is:



Yet a third is tied to the incorporation of water:



It can be seen from Equations 5-8, that these different substitutions each have their own dependence on the activity of  $\text{SiO}_2$ , and some additionally depend on the activities of  $\text{Al}_2\text{O}_3$  and  $\text{H}_2\text{O}$ . Furthermore, other components in the melt (e.g.  $\text{CaO}$ ) affect the activity coefficients of  $\text{CrO}$  and  $\text{CrO}_{1.5}$ , and accordingly the transition between  $\text{Cr}^{2+}$  and  $\text{Cr}^{3+}$  and their partition coefficients are also sensitive to melt composition (Berry et al., 2006). Consequently, the examples of partition coefficients quoted in this review should be taken as illustrative, rather than definitive.

In this chapter we summarize current knowledge on the effect of oxygen fugacity on the partitioning relations of redox-sensitive elements with relevant application to magmatic processes. We start by outlining some of the theoretical background behind the treatment of experimental partitioning data of redox-sensitive trace elements. This is followed by detailed sections on mineral/melt partitioning of transition metals (Ti, V, Cr, Fe), rare earth elements (Ce, Eu), uranium (U), and siderophile elements (Mo, W, Re, and platinum group elements).

## 2. Theoretical background

The bulk partition coefficient for a redox-variable element M ( $D_M$ ) is given by the sum of the contributions from each valence state (a+, b+, c+), which is equivalent to the average of the partition coefficients for each valence state ( $D_{M^{a+}}$ ,  $D_{M^{b+}}$ ,  $D_{M^{c+}}$ ) weighted by their relative abundances:

$$D_M = D_{M^{a+}} \left( \frac{M^{a+}}{\sum M} \right) + D_{M^{b+}} \left( \frac{M^{b+}}{\sum M} \right) + D_{M^{c+}} \left( \frac{M^{c+}}{\sum M} \right) + \dots \quad (9)$$

where  $\sum M = M^{a+} + M^{b+} + M^{c+} + \dots$ . Any thermodynamic treatment of mineral/melt partitioning of redox-sensitive elements, however, requires that both the homogeneous redox equilibria between the components in the melt phase and the heterogeneous equilibria of the mineral-melt partitioning reaction are considered.

### 2.1. Homogeneous equilibria

The general reaction describing the oxidation of a redox-variable element M in a melt is:



where  $x$  is the effective valence state of the metal cation M, and  $n$  is the number of electrons involved in the redox process. The equilibrium constant ( $K$ ) for this homogeneous reaction is given by:

$$K_{10} = \frac{X_{M^{(x+n)+}O_{(x+n)/2}} \gamma_{M^{(x+n)+}O_{(x+n)/2}}}{X_{M^{x+}O_{x/2}} \gamma_{M^{x+}O_{x/2}} (fO_2)^{n/4}} \quad (11)$$

where  $X$  is mole fraction and  $\gamma$  is the activity coefficient of each component relative to the standard states implicit in the definition of the Gibb's free energy of the reaction at the temperature of interest. Assuming that, for the low concentrations of trace elements, the activity coefficients are constant for a given composition (Henry's Law), the expression may be simplified to:

$$K'_{10} = \left( \frac{[M^{(x+n)+}O_{(x+n)/2}]}{[M^{x+}O_{x/2}]} \right) (fO_2)^{-n/4} \quad (12)$$

where  $K'_{10} = K_{10} / (\gamma_{M^{(x+n)+}O_{(x+n)/2}} / \gamma_{M^{x+}O_{x/2}})$  and the values in square brackets represent concentrations converted from mole fractions (note that the conversion factor from mole fractions to concentration cancels out between the numerator and denominator in the quotient). Taking the logarithm of this equation and re-arranging gives:

$$\log \left( \frac{[M^{(x+n)+}O_{(x+n)/2}]}{[M^{x+}O_{x/2}]} \right) = \frac{n}{4} \log fO_2 + \log K'_{10} \quad (13)$$

In the case where only two valence states are possible in the melt, by defining  $\sum M = M^{x+} + M^{(x+n)+}$ , one obtains:

$$\frac{M^{x+}}{\sum M} = \frac{1}{1 + 10^{(n/4 \log fO_2 + \log K'_{10})}} \quad (14a)$$

and

$$\frac{M^{(x+n)+}}{\sum M} = \frac{1}{1 + 10^{-(n/4 \log fO_2 + \log K'_{10})}} \quad (14b)$$

From this equation, it can be seen that  $M^{x+} / \sum M$  takes a sigmoidal form, constrained to vary from 0 to 1, with  $\log fO_2$ . The slope is defined by  $n/4$  (i.e. the number of electrons in the redox reaction), and the position in  $\log fO_2$  space over which the transition from one valence state to the other occurs is determined by  $\log K'_{10}$  (i.e. the simplified equilibrium constant for the homogeneous redox reaction). For elements such as U and V, which have three and four possible valence states in the melt,  $\sum M \neq M^{x+} + M^{(x+n)+}$ ; instead, more complicated formulations are required. In the case of Cr, which can occur as  $Cr^{2+}$ ,  $Cr^{3+}$  and  $Cr^{6+}$ , there is essentially no coexistence of  $Cr^{2+}$  and  $Cr^{6+}$ , and hence Equations 14a,b are valid approximations.

The above treatment groups all  $M^{x+}$  together. However, from a melt structural viewpoint, it has been argued that  $Fe^{3+}$  and  $Eu^{3+}$  occur both as dissociated cations and in complexes such as  $FeO_2^-$  and  $EuO_2^-$  (Fraser, 1975; Ottonello et al., 2001). Although it would be possible to write separate reactions for oxidation of  $Fe^{2+}$  to  $Fe^{3+}$  and  $FeO_2^-$ , the resulting equations would have an identical mathematical form to the simple model we present here. Our activity coefficients, and the use of a modified equilibrium constant  $\log K'_{10}$ , accommodate the terms that would be necessary to account for the speciation of the participating ions. The use of a more refined structural model may, however, prove advantageous in describing the dependence of  $\log K'_{10}$  on melt composition. This is addressed in greater detail in Moretti (this volume).

## 2.2. Heterogeneous equilibria

In the mineral (min) phase, the concentration of M will be given by the sum of all possible valence states, each with its own partitioning reaction:

$$M^{x+}_{\text{mineral}} = [M^{x+}O_{x/2}]_{\text{melt}} + \sum [N^{x+}O_{x/2}]_{\text{melt}} \quad (15)$$

where  $N^{x+}O_{x/2}$  are the ‘stoichiometric’ control needed to form the phase component in which  $M^{x+}O_{x/2}$  occurs (O’Neill and Eggins, 2002). The equilibrium constant for this heterogeneous reaction is given by:

$$K_{15} = \frac{X_{M^{x+}}^{\min} \gamma_{M^{x+}O}^{\min}}{X_{M^{x+}O_{x/2}}^{\text{melt}} \gamma_{M^{x+}O_{x/2}}^{\text{melt}} \prod a_{N^{x+}O_{x/2}}^{\text{melt}}} \quad (16)$$

where,  $\prod a_{N^{x+}O_{x/2}}$  is the product of the activities of the components  $N^{x+}O_{x/2}$  that make up the ‘stoichiometric control’ in the melt. The mineral/melt partition coefficient for each valence state of M is given by:

$$D_{M^{x+}} = \frac{\gamma_{M^{x+}O}^{\min}}{K_{15} \left( \gamma_{M^{x+}O_{x/2}}^{\text{melt}} \prod a_{N^{x+}O_{x/2}}^{\text{melt}} c \right)} \quad (17)$$

where  $c$  is a constant to convert mole fractions to concentrations by weight. From this equation, it becomes clear that the mineral/melt partitioning of M will depend on the activities of the major element oxides in the system, even if they mix ideally.

### 2.3. Mineral/melt partitioning as a function of $fO_2$

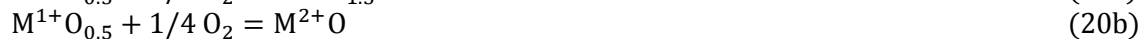
The mineral/melt partition coefficient of a redox-sensitive element M at constant pressure, temperature, and major element composition can be defined (alternatively to Equation 9) as the sum of the concentrations of all valence states in the mineral phase, divided by the sum of the concentration of all valence states in the melt phase:

$$D_M = \frac{\sum [M_{\min}^{x+}]}{\sum [M_{\text{melt}}^{x+}]} \quad (18)$$

Assuming that M occurs in the valence states 1+, 2+, 3+..., this equation can be re-arranged to:

$$D_M = \frac{\frac{[M_{\min}^{1+}]}{[M_{\text{melt}}^{1+}]} + \frac{[M_{\min}^{2+}]}{[M_{\text{melt}}^{1+}]} + \frac{[M_{\min}^{3+}]}{[M_{\text{melt}}^{1+}]} \dots}{1 + \frac{[M_{\text{melt}}^{2+}]}{[M_{\text{melt}}^{1+}]} + \frac{[M_{\text{melt}}^{3+}]}{[M_{\text{melt}}^{1+}]} \dots} \quad (19)$$

Homogeneous redox reactions can be written relating the different valence states to either the more reduced or oxidized end-members:



...

and the ratios of the concentrations of these related to  $K'$  and  $fO_2$  in the following way:

$$[M_{\text{melt}}^{3+}]/[M_{\text{melt}}^{1+}] = (fO_2)^{0.5} K'_{20a} \quad (21a)$$

$$[M_{\text{melt}}^{2+}]/[M_{\text{melt}}^{1+}] = (fO_2)^{0.25} K'_{20b} \quad (21b)$$

Substituting Equation 21 into Equation 19 and re-arranging gives:

$$D_M = \frac{D^{1+} + [D^{2+}(fO_2)^{0.25} K'_{20b}] + [D^{3+}(fO_2)^{0.5} K'_{20a}] + \dots}{1 + [(fO_2)^{0.25} K'_{20b}] + [(fO_2)^{0.5} K'_{20a}] + \dots} \quad (22)$$

For convenience, this expression can also be written in the following form:

$$D_M = \frac{D^{1+} + (D^{2+} 10^{0.25 \log fO_2 + \log K'_{20b}}) + (D^{3+} 10^{0.5 \log fO_2 + \log K'_{20a}}) + \dots}{1 + (10^{0.25 \log fO_2 + \log K'_{20b}}) + (10^{0.5 \log fO_2 + \log K'_{20a}}) + \dots} \quad (23)$$

which then allows for the log  $fO_2$  term to be substituted more readily by oxygen fugacity terms relative to buffers in the customary way (e.g.  $\Delta FMQ$ ,  $\Delta IW$ , etc.) as a form of minimizing the effect of temperature and pressure, if non-isobaric/isothermal data are considered.

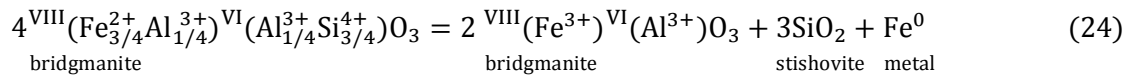
### 3. Transition metals (Fe, Cr, Ti, V)

A large number of transition metals (those elements whose atoms have incompletely filled  $3d$  orbitals) occur in more than one valence state under magmatic  $fO_2$  conditions (Fig. 2). Amongst the redox-sensitive first-row transition metals, Fe is a major constituent of the solar system while others like Ti, V, and Cr occur in minor-to-trace concentrations (Palme and O'Neill, 2014). Mineral/melt partition relations of transition metals are harder to constrain in comparison with other trace elements as their behavior frequently does not obey Henry's Law. Furthermore, their electronic configuration results in strong crystal-field effects (Burns, 1993) and hence partitioning behavior is often anomalous in the context of lattice strain theory.

#### 3.1. Iron (Fe)

As the only redox-sensitive cation that is a major element component in the solar system, iron has a fundamental role in magmatic redox processes. Arguably,  $fO_2$  in the solid Earth is entirely controlled by the partitioning behavior of iron, which occurs as  $Fe^{3+}$  (ferric iron),  $Fe^{2+}$  (ferrous iron) as well as  $Fe^0$  (metallic iron). The crystal-chemical controls on Fe partitioning are essential to understanding  $fO_2$  throughout the Earth's mantle. For example, in the dominant upper mantle phase olivine ( $Mg_2SiO_4$ ), iron occurs almost exclusively in the divalent form substituting for Mg in octahedral sites (i.e.  $Fe^{3+}/\Sigma Fe \sim 0$ ; O'Neill et al. 1993). This means that, in a system closed to oxygen, a reduced basaltic melt that crystallizes olivine will become oxidized (i.e. its  $Fe^{3+}/\Sigma Fe$  will increase). Conversely, an oxidized melt that crystallizes magnetite ( $Fe^{2+}Fe_2^{3+}O_4$ , with  $Fe^{3+}/\Sigma Fe \sim 0.66$ ) will become reduced (i.e. its  $Fe^{3+}/\Sigma Fe$  will decrease). Other ferromagnesian minerals in the upper mantle (e.g. pyroxenes, amphiboles, garnets) will also contribute to this mass balance.

Another fundamental example that illustrates the significance of crystal-chemical controls on the partitioning behavior of Fe comes from the most abundant mineral phase in the Earth's mantle, bridgmanite (Fe,Mg)(Si,Al)O<sub>3</sub>. Bridgmanite can contain up to 20 atomic % Fe, a substantial proportion of which is trivalent (Fe<sup>3+</sup>/ΣFe up to 0.6; McCammon, 1997; Frost et al., 2004). This preference of bridgmanite for ferric iron occurs because the coupled substitution of Fe<sup>3+</sup> on to the VIII-fold coordinated Mg site, charge balanced by Al substitution on to the VI-fold coordinated Si site, is energetically favorable (Kesson et al., 1995). Therefore, Fe<sup>3+</sup> occurs in experimentally grown Al-bearing bridgmanite even at very reducing conditions, with fO<sub>2</sub> maintained low via a disproportionation reaction that form metallic iron (e.g. Shim et al., 2017):



These types of disproportionation reactions are powerful in explaining the paradoxical evidence suggesting high Fe<sup>3+</sup> contents in the very reduced lower mantle (McCammon, 2005).

In other abundant ferromagnesian minerals such as pyroxenes, the incorporation of Fe<sup>2+</sup> and Fe<sup>3+</sup> is also strongly dictated by stoichiometric controls. The experimental data of McCanta et al. (2004), for instance, highlight significant steric constraints on the substitution of Fe<sup>3+</sup> for Fe<sup>2+</sup> in pigeonite (low Ca clinopyroxene) compared to augite (high Ca clinopyroxene). The limited incorporation of Fe<sup>3+</sup> into pigeonite is such that even under highly oxidizing conditions only a few percent of the total Fe will be trivalent. In contrast, alkali pyroxenes like aegirine (NaFe<sup>3+</sup>Si<sub>2</sub>O<sub>6</sub>) have a strong affinity for Fe<sup>3+</sup> but will not precipitate, even under oxidized conditions, unless there is sufficient Na available in the melt. As there is complete solid solution between augite and aegirine, Fe<sup>3+</sup>/Fe<sup>2+</sup> of a pyroxene has a complex dependence on melt composition when Fe is a major element, and lies beyond the scope of this discussion. A related issue arises with the partitioning of trace elements into magnetite, which is strongly fO<sub>2</sub>-dependent. However, this is mostly due to the stoichiometric control by the Fe<sup>3+</sup>/Fe<sup>2+</sup> ratio of the melt rather than any variation in oxidation state of the trace elements themselves. For example, Ti becomes less compatible from FMQ to FMQ+4, well outside the range at which any Ti<sup>3+</sup> is present, and this behavior is also exhibited by redox-invariant trace elements such as Zr and Hf (Sievwright et al., 2017). Hence, care should be taken when evaluating redox-related trends in Fe-rich systems to understand the factors contributing to variations in partition coefficients. The relationships described in this chapter are only applicable to trace components of melts and minerals.

Based on experiments with simplified synthetic compositions doped with small concentrations of Fe as a function of fO<sub>2</sub>, Mallmann and O'Neill (2009) showed that  $D_{\text{Fe}^{2+}} \gg D_{\text{Fe}^{3+}}$  in olivine and orthopyroxene, but  $D_{\text{Fe}^{2+}} < D_{\text{Fe}^{3+}}$  in clinopyroxene and spinel (Fig. 3a). Amphibole/melt data from King et al. (2000) indicates that the partitioning of both Fe<sup>3+</sup> and Fe<sup>2+</sup> are both close to unity, with  $D_{\text{Fe}^{3+}}$  about a factor of two higher than  $D_{\text{Fe}^{2+}}$  (though the range of fO<sub>2</sub> in this study, from FMQ-6 to FMQ+3, is not sufficient to constrain either Fe<sup>2+</sup> or Fe<sup>3+</sup>, so the differences in  $D$ s between the two species are likely larger). The data of Martel et al. (1999) and Krawczynski et al. (2012), which examined fO<sub>2</sub> over an even more limited range alongside the effects of H<sub>2</sub>O,  $T$  and  $P$ , show higher  $D_{\text{Fe}}/D_{\text{Mg}}$  at more oxidizing conditions, supporting a higher compatibility for Fe<sup>3+</sup> in igneous amphiboles. It should be noted in passing that amphibole Fe<sup>3+</sup>/ΣFe is readily altered by a redox reaction  $\text{Fe}^{2+} + \text{OH}^- = \text{Fe}^{3+} + \text{O}^{2-} + \frac{1}{2} \text{H}_2$ , and hence may not preserve magmatic values (Popp et al., 1995).

In non-ferromagnesian minerals such as feldspars, iron behaves as a typical trace element. In plagioclase feldspar for instance (Fig. 3a), Fe<sup>3+</sup> ( $r^{\text{IV}} = 0.49 \text{ \AA}$ )<sup>1</sup> is significantly more

<sup>1</sup> Ionic radii reported in this study were taken from Shannon (1976).



compatible than  $\text{Fe}^{2+}$  (Phinney, 1992; Wilke and Behrens, 1999) as it directly substitutes for  $\text{Al}^{3+}$  ( $r^{\text{IV}} = 0.39 \text{ \AA}$ ) in tetrahedral sites, as opposed to  $\text{Fe}^{2+}$  which is significantly smaller ( $r^{\text{VIII}} = 0.78 \text{ \AA}$ ) than the larger VIII-fold coordinated A-site ( $r_0 \sim 1.2 \text{ \AA}$ ) (Longhi et al., 1976). In some oxidized peralkaline lavas, alkali feldspars with up to 2.8 wt.%  $\text{Fe}_2\text{O}_3$  have been reported (Mann et al. 2006). Lundgaard and Tegner (2004) found a strong dependence of  $D_{\text{Fe}^{2+}}$  and  $D_{\text{Fe}^{3+}}$  on melt composition, particularly  $\text{SiO}_2$ , but not on plagioclase composition.

### 3.2. Chromium (Cr)

Chromium can exist as  $\text{Cr}^{2+}$ ,  $\text{Cr}^{3+}$  and  $\text{Cr}^{6+}$  in magmas (e.g. Berry and O'Neill, 2004; Schreiber and Haskin, 1976), but direct detection of  $\text{Cr}^{2+}$  in quenched basaltic melts is hindered by an electron exchange with Fe on quenching (Reaction 1; Berry et al., 2003). Mineral-melt partitioning relations, however, have long indicated the presence of  $\text{Cr}^{2+}$  in silicate melts over typical magmatic conditions (e.g. Barnes, 1986; Hanson and Jones, 1998; Postovetov and Roeder, 2000). The divalent and trivalent species of Cr overwhelmingly dominate over typical magmatic  $f\text{O}_2$ s (Fig. 2), with hexavalent Cr only becoming stable at very oxidizing conditions (usually  $> \text{FMQ}+4$ ; Berry and O'Neill, 2004). Because such highly oxidizing conditions are uncommon in magmatic systems, and  $\text{Cr}^{6+}$  appears to be highly incompatible in high-temperature mineral phases (Mallmann and O'Neill, 2009), the presence of small amounts of  $\text{Cr}^{6+}$  in magmas should have no significant effect on the bulk partitioning of Cr (although it should be noted that the transition from  $\text{Cr}^{3+}$  to  $\text{Cr}^{6+}$  is expected to occur over a relatively narrow range of  $f\text{O}_2$  given that the redox reaction involves three electrons). High-temperature experiments under the highly oxidizing conditions necessary to constrain the behavior of  $\text{Cr}^{6+}$  are limited (hexavalent Cr appears to be volatile at high temperature; Wijbrans et al., 2015), so the focus of our discussion will be on  $\text{Cr}^{2+}$  and  $\text{Cr}^{3+}$ .

Chromium is a major constituent of many rock-forming minerals (Burns and Burns, 1975), normally making end-member components where  $\text{Cr}^{3+}$  substitutes for  $\text{Al}^{3+}$  (and  $\text{Fe}^{3+}$ ). Some of these substitutions are complete solid solutions, most notably between spinel ( $\text{MgAl}_2^{3+}\text{O}_4$ ) and magnesiochromite ( $\text{MgCr}_2^{3+}\text{O}_4$ ), and between grossular ( $\text{Ca}_3\text{Al}_2^{3+}\text{Si}_3\text{O}_{12}$ ) and uvarovite ( $\text{Ca}_3\text{Cr}_2^{3+}\text{Si}_3\text{O}_{12}$ ). In most cases,  $\text{Cr}^{3+}$  only occurs in octahedral coordination, as a result of its large crystal-field stabilization energy (CFSE; Burns, 1993). The ionic radius of  $\text{Cr}^{3+}$  is also significantly larger than  $\text{Al}^{3+}$  (and hence  $\text{Si}^{4+}$ ), making it difficult for tetrahedral substitution. Incorporation of  $\text{Cr}^{2+}$  is unfavored due to its electronic configuration, which in octahedral coordination induces a destabilizing crystal-field energy and Jahn-teller distortions (Burns, 1975). Hence,  $\text{Cr}^{2+}$  can substitute for  $\text{Mg}^{2+}$  but only to a limited extent, and endmember  $\text{Cr}^{2+}$  silicates do not occur in nature (Li et al., 1995).

The importance of Cr to mantle phase equilibria (e.g. Liu and O'Neill, 2004), the link between chromite and precious metals (e.g. platinum group elements, PGE) in layered intrusions (e.g. Barnes, 1986), and its utility as an oxybarometer in magmatic systems (e.g. Postovetov and Roeder, 2000) has prompted numerous studies aimed at determining Cr solubility in silicate melts (e.g. Roeder and Reynolds, 1991) and its partitioning between associated phases. Here we focus on minerals where Cr substitutes as a trace component, as this allows us to constrain the effect of  $f\text{O}_2$  on the partitioning behavior of both  $\text{Cr}^{2+}$  and  $\text{Cr}^{3+}$ .

Amongst minerals where Cr substitutes as trace element, olivine has the largest available experimental partitioning dataset as a function of  $f\text{O}_2$ . In olivine,  $\text{Cr}^{2+}$  ( $r^{\text{VI}} = 0.80 \text{ \AA}$ ) is incorporated by direct replacement of  $\text{Mg}^{2+}$  ( $r^{\text{VI}} = 0.72 \text{ \AA}$ ) in octahedral sites without significant lattice distortion (Reaction 5; Jollands et al., 2018). Substitution of  $\text{Cr}^{3+}$  ( $r^{\text{VI}} = 0.615 \text{ \AA}$ ) into olivine octahedral sites, on the other hand, creates charge imbalance that can be accounted for either by octahedral vacancies (Reaction 6; Hanson and Jones, 1998; Jollands et al., 2018; Papike

et al., 2005), or by coupled substitution with  $\text{Al}^{3+}$  and  $\text{Si}^{4+}$  in tetrahedral sites (Reaction 7; Jollands et al., 2018).

Olivine/melt partitioning studies from Schreiber and Haskin (1976) and Hanson and Jones (1998) in the  $\text{CaO-MgO-Al}_2\text{O}_3\text{-SiO}_2$  (CMAS) show a systematic dependence of  $D_{\text{Cr}^{3+}}$  on melt composition (Fig. 4), though the dependence is not straightforward, suggesting that both substitution mechanisms described above likely operate. Hanson and Jones (1998), however, were able to parameterize  $D_{\text{Cr}^{3+}}$  between olivine and melt using the universal melt descriptor  $\text{NBO}/T$  ( $D_{\text{Cr}^{3+}} = -0.39 \text{ NBO}/T + 1.29$ ). For  $\text{Cr}^{2+}$ , however, partitioning between olivine and melt appears to be independent of melt composition, but varies slightly with temperature (Hanson and Jones, 1998). More often than not,  $D_{\text{Cr}^{2+}}$  is nearly equal to  $D_{\text{Cr}^{3+}}$  between olivine and melt (including for basaltic compositions), so that effectively no change in bulk Cr partitioning is detected as a function of  $f\text{O}_2$  (Mallmann and O'Neill, 2009; Hanson and Jones, 1998; Fig. 3b). Direct Cr valence measurements in olivine by XANES have confirmed these findings (Bell et al. 2014).

Mallmann and O'Neill (2009) found that  $\text{Cr}^{2+}$  is similarly partitioned amongst olivine, ortho- and clinopyroxene, but  $\text{Cr}^{3+}$  is significantly more compatible in pyroxenes, particularly clinopyroxene (Fig. 3b). This is not surprising given that  $\text{Cr}^{3+}$  ( $r^{\text{VI}} = 0.615 \text{ \AA}$ ) has an ionic radius that is similar to the ideal strain-free octahedral M1 site of clinopyroxenes for trivalent cations ( $r_{0^{3+},\text{M1}} \sim 0.668 \text{ \AA}$ ; Hill et al., 2000).  $\text{Cr}^{2+}$  ( $r^{\text{VI}} = 0.80 \text{ \AA}$ ), on the other hand, is significantly larger than the ideal M1 site of clinopyroxenes for divalent cations ( $r_{0^{2+},\text{M1}} \sim 0.69 \text{ \AA}$ ; Hill et al., 2000). Mallmann and O'Neill (2009) noted a significant difference between the partitioning of  $\text{Cr}^{3+}$  between clinopyroxene and melt obtained for their compositions V1 and V7 (which are identical except for the Cr-doping levels), suggesting that  $\text{Cr}^{3+}$  does not behave according to Henry's Law in the concentration range of these experiments ( $> 1.5 \text{ wt.}\%$   $\text{Cr}_2\text{O}_3$ ), probably because of the strong coupling of  $\text{Cr}^{3+}$  with  $\text{Al}^{3+}$  in pyroxenes (Klemme and O'Neill, 2000). Chromium is noticeably more compatible in augite compared to pigeonite and enstatite, given the availability of cations ( $\text{Al}^{3+}$  and  $\text{Na}^+$ ) for charge balancing the substitution of  $\text{Cr}^{3+}$  for Mg in octahedral sites (Karner et al., 2007; Mallmann and O'Neill, 2009). Interestingly,  $\text{Cr}^{2+}$  appears to be dominant species entering the structure of lower mantle phases such as ferropericlase and bridgmanite (Eeckhout et al., 2007). Neither  $\text{Cr}^{3+}$  nor  $\text{Cr}^{2+}$  appears to be easily incorporated into plagioclase (Aigner-Torres et al., 2007) as  $\text{Cr}^{3+}$  is too large for the tetrahedral site, and  $\text{Cr}^{2+}$  too small for the large distorted VIII-fold coordinated site.

Finally, we note that Cr is a well-known outlier in lattice-strain parabolas due to large crystal-field effects, so the prediction of  $\text{Cr}^{2+}$  and  $\text{Cr}^{3+}$  partitioning using lattice strain models should be done with great caution.

### 3.3. Titanium (Ti)

Titanium primarily occurs as  $\text{Ti}^{4+}$  over most of the range of  $f\text{O}_2$ s relevant to planetary scientists, but spectroscopic studies indicate that under very reducing conditions (i.e. below FMQ-4) a substantial fraction of Ti occurs as  $\text{Ti}^{3+}$  (Schreiber, 1977). The transition from  $\text{Ti}^{3+}$  to  $\text{Ti}^{4+}$  appears to be strongly dependent on melt composition (Borisov, 2012; Fig. 2). Hence, although the partitioning of Ti is a poor oxybarometer for terrestrial magmatism, it is markedly more relevant to understanding the geochemical evolution of the early solar nebula and reduced bodies such as the Moon (Fig. 2).

Hibonite ( $\text{CaAl}_{12}\text{O}_{19}$ ) occurs in the calcium aluminum-rich inclusions (CAIs) of chondritic meteorites. It is one of the earliest phases to condense from the solar nebula, and can contain up to 8 wt%  $\text{TiO}_2$  (Allen et al., 1978). The blue color of many hibonite grains has been shown to be due to the presence of  $\text{Ti}^{3+}$  (Burns and Burns, 1984; Sutton et al., 2017),

incorporated as a  $\text{CaTi}_2^{3+}\text{Al}_{10}\text{O}_{19}$  component, but  $\text{Ti}^{4+}$  can also be accommodated in hibonite via a coupled substitution with Mg, as a  $\text{CaMg}_2\text{Ti}_2^{4+}\text{Al}_8\text{O}_{19}$  component (Doyle et al., 2014). The coupled nature of the substitution of  $\text{Ti}^{4+}$  means that the availability of Mg exerts a strong control on  $\text{Ti}^{3+}/\Sigma\text{Ti}$  (where  $\Sigma\text{Ti} = \text{Ti}^{3+} + \text{Ti}^{4+}$ ) in hibonite; since it is often unclear which phases are in equilibrium in CAIs, this oxybarometer may have limited utility (Berry et al., 2017).

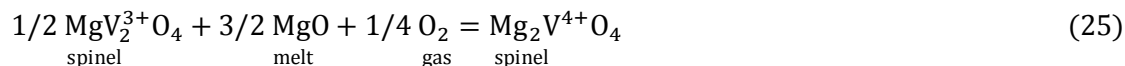
$\text{Ti}^{3+}$  is more compatible than  $\text{Ti}^{4+}$  in the majority of phases, including olivine, diopside, orthopyroxene and anorthite (Mallmann and O'Neill, 2009; Peters et al., 1995; Fig. 3c). Ti K-edge XANES measurements of both synthetic and natural lunar basalt samples have revealed that  $\text{Ti}^{3+}$  is enriched in ilmenite, armalcolite and clinopyroxene at  $f\text{O}_2$  of about FMQ-4.5 (Leitzke et al., 2018; Simon and Sutton, 2018). However, as it is most compatible in clinopyroxene, essentially all research has focused on this phase. Indeed, a  $\text{Ti}^{3+}$ -pyroxene (grossmanite,  $\text{CaTi}^{3+}\text{AlSiO}_6$ ) has been reported from the Allende meteorite, and  $\text{NaTi}^{3+}\text{Si}_2\text{O}_6$  has been synthesized experimentally (Ohashi et al., 1982; Ma and Rossman, 2009). Although quantitative oxybarometry is hampered by multiple sources of uncertainty, including the equilibrium mineral assemblage (as with hibonite), the mixing relations of Ti-bearing pyroxenes and the temperature of formation, decreasing  $\text{Ti}^{3+}/\Sigma\text{Ti}$  from core to rim of some CAIs suggests increasingly oxidizing conditions in the solar nebula over a period of up to 300 kyr, perhaps by as much as 6 – 7 log  $f\text{O}_2$  units (Young et al., 2012). However, in some CAIs an oxidation trend, followed by a more reduced rim, has been observed. This is attributed to incompatible  $\text{Ti}^{4+}$  concentrating in the melt phase, which was isolated from the solar nebula by a thick melilite mantle, followed by late-stage ingress of nebular gases to reduce the residual melt prior to crystallization of the rims (Papike et al., 2016).  $\text{Ti}^{3+}$  has also been observed in lunar pyroxenes (Burns et al., 1972), but is thought to be a minor component even on such a reduced planetary body (Papike et al., 2005). Terrestrial magmas generally form at far more oxidizing conditions and Ti is expected to occur exclusively as  $\text{Ti}^{4+}$  (Fig. 2). Pyroxenes in the native iron-bearing lavas of Disko Island, Greenland, have been suggested to contain some  $\text{Ti}^{3+}$  (Pedersen, 1981), but this is speculative and may not be compatible with the rather high Fe contents.

### 3.4. Vanadium (V)

Vanadium is the redox-sensitive trace element that occurs in the largest number of valence states at typical magmatic  $f\text{O}_2$ s (Fig. 2), namely  $\text{V}^{2+}$ ,  $\text{V}^{3+}$ ,  $\text{V}^{4+}$  and  $\text{V}^{5+}$  (Borisov et al., 1987; Schreiber et al., 1987; Sutton et al., 2005). Given the plethora of valence states with theoretically contrasting geochemical behaviors, significant effort has been made to characterize the effect of  $f\text{O}_2$  and other variables such as pressure, temperature, mineral and melt composition on the partitioning of V in silicate/oxide phases relevant to mantle melting and magmatic differentiation (e.g. Arató and Audétat, 2017; Canil 1997; 1999; 2002; Canil and Fedortchouk, 2000; 2001; Karner et al., 2008; Laubier et al., 2014; Mallmann and O'Neill, 2009; 2013; Papike et al., 2015; Righter et al. 2006a,b; 2011; Shearer et al., 2006; Sievwright et al., 2017; Toplis and Corgne 2002; Wijbrans et al., 2015). The end goal of many of these studies is to calibrate the sensitivity of V partitioning to  $f\text{O}_2$  as an oxybarometer.

Under terrestrial magmatic  $f\text{O}_2$  conditions (Fig. 2),  $\text{V}^{3+}$  and  $\text{V}^{4+}$  are the dominant valence states, and spinel the major mineral group host. In pure magnetite ( $\text{Fe}^{2+}\text{Fe}_2^{3+}\text{O}_4$ ), where the stoichiometry fixes  $\text{Fe}^{3+}/\text{Fe}^{2+}$  at 2, incorporation of vanadium occurs via an  $f\text{O}_2$ -independent electron-exchange reaction:  $\text{Fe}^{3+} + \text{V}^{3+} = \text{Fe}^{2+} + \text{V}^{4+}$  (O'Neill and Navrotsky, 1984). However, experimental studies of spinel group minerals have invariably shown strong dependence of vanadium partitioning with  $f\text{O}_2$  both in Fe-free and Fe-bearing systems (e.g. Arató and Audétat, 2017; Canil et al., 1999; Mallmann and O'Neill, 2009; Papike et al., 2015; Righter et al., 2006b; Sievwright et al., 2017; Toplis and Corgne, 2002). Importantly, because Fe is an essential component of magnetite, the partitioning of non-redox variable elements depends on  $f\text{O}_2$

(Sievwright et al., 2017); this is not the case for Fe-poor spinels. For example, the ratio  $V^{3+}/V^{4+}$  of pure spinel ( $MgAl_2O_4$ ) is controlled by the reaction:



It is therefore unsurprising that the bulk partitioning of V between spinel group minerals and silicate melt is a strong function of  $fO_2$  as well as composition (mineral and melt). Experimental data shows that  $D_{V^{3+}} \gg D_{V^{4+}} > D_{V^{5+}}$  between spinel and silicate melt (Fig. 3d), and that the bulk partitioning of vanadium for Cr-spinels is higher than that for Al-spinels (Canil, 2002; Mallmann and O'Neill, 2009). The negative correlation between Al and  $V^{3+}$  in experimentally grown spinels provides strong evidence that the preferred mode of vanadium substitution into spinel is by direct replacement of Al in octahedral sites (Balan et al., 2006; Canil, 2002; Righter et al., 2006b). However, the details of the various possible substitution mechanisms are complex, given the structure of spinels and the multivalent nature of their cations (e.g. Papike et al., 2015). Melt composition also appears to play a significant role, as demonstrated by V partitioning experiments between magnetite and silicic melts (Arató and Audétat, 2017).

Owing to its high charge and small ionic radius,  $V^{5+}$  is highly incompatible in all rock-forming minerals. When sufficiently oxidized conditions are achieved, the partitioning of  $V^{5+}$  can be obtained, though this value is typically  $< 0.01$  (e.g. Mallmann and O'Neill, 2009).  $V^{4+}$  has partition coefficients that fall between  $V^{3+}$  and  $V^{5+}$  (Fig. 3d) and therefore is much harder to constrain. The ionic radius of  $V^{4+}$  is similar to that of  $Ti^{4+}$  in both octahedral and tetrahedral coordination ( $r^VI = 0.58$  and  $0.605$  Å;  $r^IV = 0.46$  and  $0.42$  Å; for  $V^{4+}$  and  $Ti^{4+}$  respectively), and both have similar electronic configurations and thus no significant difference in CFSE, so an approximation can be obtained using lattice strain theory. In connection with that, we note that  $D_{V^{4+}}$  for pyroxene and olivine obtained by Mallmann and O'Neill (2009) are about an order of magnitude higher than what is inferred from lattice strain parabolas of other tetravalent cations. For some minerals, constraining the partitioning of  $V^{3+}$  may also be difficult when it shows compatibility between  $V^{2+}$  and  $V^{4+}$  (e.g. olivine or orthopyroxene). For clinopyroxene though,  $V^{3+}$  is clearly more compatible than both  $V^{2+}$  and  $V^{4+}$ , and its partitioning readily determined (Fig. 3d). Partition coefficients for  $V^{2+}$  can be obtained if sufficiently reduced experimental data is available, such as those reported by Mallmann and O'Neill (2009), but extrapolation from lattice strain parabolas should be avoided since  $V^{2+}$  has the same electronic configuration as  $Cr^{3+}$  and therefore large CFSE in octahedral coordination (Burns, 1993).

For mineral phases that only crystallize at high pressure from natural magmas, partitioning relations are harder to constrain as a function of  $fO_2$  since this intensive parameter needs to be buffered externally using double-capsule techniques or internally using low-solubility solid metal-oxide. Righter et al. (2011) extended the  $fO_2$  conditions obtained by Mallmann and O'Neill (2007, 2009) to constrain the  $D_{V^{3+}}$  in garnet at  $\sim 10$ . The compatibility order of vanadium valences in garnet is  $D_{V^{3+}} \gg D_{V^{4+}} > D_{V^{5+}}$ , with values comparable to clinopyroxene (Mallmann and O'Neill, 2009). XANES studies by Righter et al. (2011), however, suggest that small amounts of  $V^{2+}$  are incorporated into garnet, even under relatively oxidizing conditions.

Because of the complexities in developing a thermodynamic partitioning model that involves four different oxidation states, all of which likely varying differently with compositional parameters, Mallmann and O'Neill (2013) developed an empirical calibration of V partitioning between olivine and melt over a restricted range of  $fO_2$ s (from FMQ-4 to FMQ+4), where  $\log D_V$  vary more or less linearly with  $\log fO_2$ , accounting for the effects of olivine and melt composition. Similarly, Arató and Audétat (2017) developed an empirical calibration for magnetite. It would be useful to develop similar models for other common magmatic minerals, but this requires more extensive experimental databases since substitution of different V

valence states in minerals like pyroxene and garnet are more complex than in olivine and magnetite.

#### 4. Rare earths (Ce and Eu)

The rare earth elements (REE), including Y, have favorable characteristics for examining redox processes. Apart from Ce and Eu, the REE exist exclusively in the trivalent state and their ionic radii vary smoothly from La to Lu (with Y somewhere between Ho and Er). This provides excellent coverage for fitting a lattice strain parabola to the REE<sup>3+</sup>; anomalies relative to this parabola are readily quantifiable, and often visually striking. Zircon is notable because its REE patterns commonly feature both Ce and Eu anomalies (Fig. 5). They are typically quantified as Ce/Ce\* and Eu/Eu\*, where Ce\* and Eu\* are the values that would be predicted by interpolation or extrapolation. The simplest is the geometric mean of La and Pr (i.e., a linear interpolation on a logarithmic axis):

$$\text{Ce}^* = (\text{La}_N \times \text{Pr}_N)^{0.5} \quad (26)$$

where 'N' denotes normalization to CI chondrite, whole rock or matrix glass. La is often below the limit of detection for analyses of zircon, and so extrapolation from Pr-Nd-Sm may be used instead. Alternatively, polynomial fits to the REE pattern may be used (e.g. O'Neill, 2016; Burnham and Berry, 2017). There are systematic deviations between the methods, so values taken from different sources may need recalculation before being compared.

Cerium exists as both Ce<sup>3+</sup> and Ce<sup>4+</sup> in magmas. However, at typical magmatic fO<sub>2</sub> the proportion of Ce<sup>4+</sup> is small (Fig. 2), with Ce<sup>4+</sup>/ΣCe ~ 0.001 (Burnham and Berry, 2012; 2014). At such low levels, it can be seen from Fig. 6a that where Ce<sup>4+</sup> is less compatible than Ce<sup>3+</sup> there is no perceptible effect on the bulk partitioning of Ce: even in the most extreme case, i.e.  $D_{\text{Ce}^{4+}} = 0$ ,  $D_{\text{Ce}}$  would drop by 0.1 %, an amount that is far less than the achievable analytical accuracy. By a similar line of logic, it can be calculated that in order for a mineral in a terrestrial magma to develop a noticeable Ce anomaly, e.g. Ce/Ce\* = 1.1,  $D_{\text{Ce}^{4+}}$  must be greater than  $D_{\text{Ce}^{3+}}$  by a factor of 100. For minerals with a steep or strongly curved REE patterns, or larger analytical uncertainties, even a 10 % excess of Ce might not be a convincing anomaly.

Few minerals have a strong affinity for Ce<sup>4+</sup>. The most notable, zircon, will be addressed in detail below. This is a favorable substitution on account of the identical 4+ charge of Zr, and a similarly large ionic radius ( $r^{\text{VIII}} = 0.84$  and  $0.97$  Å for Zr and Ce<sup>4+</sup> respectively). Positive Ce anomalies have been reported from other Zr minerals such as baddeleyite (Schärer et al., 2011), wadeite (Jaques, 2016), elpidite and armstrongite (Gerdes et al., 2017). Thomson et al. (2016) reported both positive and negative Ce anomalies for inclusions in diamonds, but the precise interpretation of these grains is complicated by their polymineralic nature and a lack of experimental partitioning data.

Zircon is the commonest of these, and has attracted the most experimental work (Burnham and Berry, 2012; Trail et al., 2011; 2012). The most recent calibration of the Ce-in-zircon oxybarometer uses lattice strain theory to infer  $D_{\text{Ce}^{3+}}$  and  $D_{\text{Ce}^{4+}}$ , which allow Ce<sup>4+</sup>/ΣCe in the melt to be calculated and the corresponding fO<sub>2</sub> deduced (Smythe and Brenan, 2016). This approach avoids the challenge of growing analyzable well-equilibrated zircon crystals, but there is presently a discrepancy between the two available models for Ce speciation in silicate melts (Burnham and Berry, 2014; Smythe and Brenan, 2015) and hence further work is required to develop a robust version of this oxybarometer.

For other minerals, there has been no systematic study of the effect of fO<sub>2</sub>, but this can be predicted by a combination of lattice strain modelling (and for minerals with more than one

cation site, whose 4+ parabolas are underconstrained, by making the approximation  $D_{\text{Ce}^{4+}} = D_{\text{U}^{4+}}$  and  $\text{Ce}^{4+}/\Sigma\text{Ce}$  estimated from the optical basicity of the melt composition (Burnham and Berry, 2014). It can be seen that  $\text{Ce}^{4+}$  is less compatible than  $\text{Ce}^{3+}$  in Zr-free minerals and that, as explained above,  $D_{\text{Ce}}$  does not change appreciably in the  $f\text{O}_2$  range of natural magmas.

Although the foregoing discussion considered mineral/melt partitioning, it is not always possible to consider mineral compositions in terms of mineral/melt partition coefficients. In particular, plutonic rocks do not retain any melt fraction, but even volcanic rocks often contain a significant cargo of crystals that are substantially older and may not be in equilibrium with their present host (e.g. Murphy et al., 1998). In the case of Ce, it is possible to use the bulk rock or even CI chondrite as a normalization factor, because magma compositions with Ce anomalies are uncommon. It is often stated that subduction zone magmas can have  $\text{Ce}/\text{Ce}^* < 1$ , indicating the incorporation of Ce-depleted sediments. As highlighted by O'Neill (2016), however, many reported examples of Ce anomalies in lavas are poorly constrained. Relatively small negative Ce anomalies,  $\text{Ce}/\text{Ce}^* \sim 0.9$ , are known in a variety of subduction settings (e.g. Hastie et al., 2009; Woodhead, 1989). As zircon Ce anomalies are typically one or two orders of magnitude larger, failure to account for a pre-existing Ce anomaly in the magma results in an underestimate of  $\log f\text{O}_2$  of  $< 0.2$ . The same is not true for Eu.

Europium can exist as  $\text{Eu}^{2+}$  and  $\text{Eu}^{3+}$  in magmas. The transition occurs most rapidly below the FMQ buffer, a range over which  $\text{Fe}^{3+}/\Sigma\text{Fe}$  changes more slowly, and hence it is potentially a useful oxybarometer for reduced magmas (Fig. 2). The crystal chemistry of  $\text{Eu}^{2+}$  ( $r^{\text{VI}} = 1.17 \text{ \AA}$ ) can be thought of as similar to  $\text{Sr}^{2+}$  ( $r^{\text{VI}} = 1.18 \text{ \AA}$ ), and indeed there is even a  $\text{Eu}^{2+}$  feldspar (Kimata, 1988). In contrast,  $\text{Eu}^{3+}$  behaves as a typical middle REE. Accordingly,  $D_{\text{Eu}^{2+}} > D_{\text{Eu}^{3+}}$  for plagioclase and alkali feldspars and åkermanitic melilite (Kuehner et al., 1989), but for most other minerals  $D_{\text{Eu}^{3+}} > D_{\text{Eu}^{2+}}$  (Fig. 6b). Plagioclase feldspar and clinopyroxene are by far the most-studied in terms of Eu partitioning as a function of  $f\text{O}_2$  and the following discussion will be restricted to these, although phases such as titanite, with  $D_{\text{Eu}^{3+}}/D_{\text{Eu}^{2+}} \approx 2500$  may be important to consider for some igneous systems (Loader et al., 2017).

There are various models that allow plagioclase/melt  $D_{\text{Eu}^{2+}}$  and  $D_{\text{Eu}^{3+}}$  to be predicted using lattice strain theory (Aigner-Torres et al., 2007; Dohmen and Blundy, 2014). Corresponding models for clinopyroxene and other minerals have not been developed, and even in the case of plagioclase, Eu oxybarometry will be more precise where measured values of partition coefficients for proxy elements are used for  $D_{\text{Eu}^{2+}}$  and  $D_{\text{Eu}^{3+}}$ . Rearranging Equation 9, we can approximate  $\text{Eu}^{3+}/\Sigma\text{Eu}$  by:

$$\text{Eu}^{3+}/\Sigma\text{Eu} = (D_{\text{Eu}} - D_{\text{Sr}})/(D_{\text{MREE}} - D_{\text{Sr}}) \quad (27)$$

where  $D_{\text{Eu}}$  is the observed Eu partition coefficient and  $D_{\text{MREE}}$  is an estimate of  $D_{\text{Eu}^{3+}}$ , i.e.  $D_{\text{Sm}}$ ,  $D_{\text{Gd}}$ , or preferably  $(D_{\text{Sm}} \times D_{\text{Gd}})^{0.5}$ . For values of  $\text{Eu}^{3+}/\Sigma\text{Eu}$  in the range 0.1 to 0.9, the  $f\text{O}_2$  can be calculated by rearranging the general equation from Burnham et al. (2015):

$$\log f\text{O}_2 = 40.4 - 25640/T - 56.8\Lambda - 4 \log (\Sigma\text{Eu}/\text{Eu}^{3+} - 1) \quad (28)$$

where  $T$  is temperature in Kelvin, and  $\Lambda$  is optical basicity, which can be calculated as:

$$\Lambda = \frac{\frac{\text{SiO}_2}{62.6} + \frac{\text{TiO}_2}{36.3} + \frac{\text{Al}_2\text{O}_3}{56.6} + \frac{\text{FeO}}{55.8} + \frac{\text{MgO}}{51.7} + \frac{\text{CaO}}{56.1} + \frac{\text{Na}_2\text{O}}{53.2} + \frac{\text{K}_2\text{O}}{67.3}}{\frac{\text{SiO}_2}{30.04} + \frac{\text{TiO}_2}{39.9} + \frac{\text{Al}_2\text{O}_3}{34} + \frac{\text{FeO}}{55.8} + \frac{\text{MgO}}{40.3} + \frac{\text{CaO}}{56.1} + \frac{\text{Na}_2\text{O}}{62} + \frac{\text{K}_2\text{O}}{94.2}} \quad (29)$$

with all oxide components in wt. %, and no distinction drawn between  $\text{Fe}^{2+}$  and  $\text{Fe}^{3+}$ . The effect of  $\text{H}_2\text{O}$  was not included in the optical basicity model of Duffy (1993) and hence is not considered here. At natural concentrations,  $\text{MnO}$  and  $\text{P}_2\text{O}_5$  have negligible influence (e.g. 0.5 wt.%  $\text{P}_2\text{O}_5$  would change  $\log K'$  for  $\text{Eu}^{2+}$ - $\text{Eu}^{3+}$  by  $\sim 0.02$ ).

With values of  $\text{Eu}^{3+}/\Sigma\text{Eu}$  outside the range 0.1 to 0.9, the reliability of this type of oxybarometer is compromised by the shallow slope of  $M^{x+}/\Sigma M$  against  $f\text{O}_2$ . An error of  $\pm 0.03$  propagates to more than a log unit of  $f\text{O}_2$ . A key challenge in using Eu partitioning to determine  $f\text{O}_2$  is the strong dependence of  $\log K'$  on melt composition (Aigner-Torres et al., 2007; Burnham and Berry, 2015). The relationship given in Equation 28 was derived for alkali-free compositions, though experimental data for Na- and K-bearing compositions in the literature are generally consistent with this relationship. Moreover, although it lies outside the scope of this review, we note that in hydrothermal fluids the  $\text{Eu}^{2+}$ - $\text{Eu}^{3+}$  equilibrium is 12 times more sensitive to pH than to  $\log f\text{O}_2$ , and hence this variable must also be considered in low-temperature systems (Brugger et al., 2008).

## 5. Uranium (U)

Although it has been known for many years that U can exist as  $\text{U}^{4+}$ ,  $\text{U}^{5+}$  and  $\text{U}^{6+}$  in silicate melts (e.g. Calas 1979; Schreiber 1983), this fact has often been rather under-appreciated in the geochemical literature. The radioactivity of U generates heat as well as a series of decay products culminating in Pb, and hence understanding the geochemistry of U is particularly important. Interpretation of U-series disequilibrium depends on accurate partition coefficients for U and Th (as well as other elements in the series). Because of potential differences in redox between mid-oceanic ridges and subduction zone settings, for example, it is essential to consider the role of  $f\text{O}_2$  in controlling the partitioning of U relative to Th.

Partitioning of U as a function of  $f\text{O}_2$  has only been studied for a handful of minerals (Fig. 7). Despite the importance of garnet in influencing U-series disequilibrium (Beattie, 1993), little is known about the how its partitioning of U might be affected by  $f\text{O}_2$ . For all minerals studied to date, it appears that  $D_{\text{U}^{4+}} > D_{\text{U}^{5+}} > D_{\text{U}^{6+}}$ . As with the intermediate oxidation states of V, the partitioning of  $\text{U}^{5+}$  is difficult to constrain without independent spectroscopic measurements of  $\text{U}^{5+}/\Sigma\text{U}$  in the same melt composition. An additional complication arises because  $\text{U}^{6+}$  is volatile and is gradually lost from 1 atm experiments, leading to uncertainty on  $D_{\text{U}^{6+}}$  because diffusion is unlikely to occur fast enough to keep crystals in equilibrium with the changing melt composition (Burnham and Berry, 2012).

Although there are insufficient data to develop a general expression for  $\text{U}^{n+}/\Sigma\text{U}$  as a function of melt composition, temperature, pressure and  $f\text{O}_2$ , it is known that higher oxidation states are favored by lower temperatures and melts with higher optical basicities (Schreiber, 1983; Halse, 2014). Notably,  $\text{U}^{5+}$  appears to have a reduced stability field with increasing pressure (Halse, 2014): at higher pressure it appears that  $\text{U}^{4+}$  oxidizes directly to  $\text{U}^{6+}$ , leading to a sharper change in partition coefficients as a function of  $f\text{O}_2$  (Mallmann et al., 2016).

U-Pb dating of zircon is widely regarded as one of the most reliable techniques for determining rock ages. For zircons younger than a few million years old, it becomes important to consider the contribution of  $^{230}\text{Th}$  to the ingrowth of  $^{206}\text{Pb}$ , because minerals for which U is less compatible than Th will develop a  $^{206}\text{Pb}$  excess, and vice versa. Correcting for these consequences of U-series disequilibrium necessitates estimates of  $D_{\text{Th}}/D_{\text{U}}$ . U is more compatible than Th under reducing conditions, but much less compatible at high  $f\text{O}_2$ : the variation in terrestrial magmatic redox conditions equates to a  $\sim 100$  kyr range of possible age corrections (McLean et al., 2011).

## 6. Siderophile elements (Mo, W, Re, Pt group elements)

Siderophile elements (i.e. those with strong tendency to alloy with Fe) are important tracers of planetary differentiation owing to their propensity of fractionating into planetary cores. Because of this, a great deal of effort has been made to investigate the behavior of these elements during core formation and to determine their valence state in silicate melts.

### 6.1. Molybdenum (Mo), tungsten (W) and rhenium (Re)

Mo, W and Re are known to occur in 4+ and 6+ valence states at  $fO_2$ s relevant to planetary magmatism (Cottrell et al., 2009; Ertel et al., 2001; Holzheid et al., 1994; O'Neill & Eggins, 2002; O'Neill et al., 2008; Righter et al., 2016; Wade et al., 2012; 2013). While the hexavalent state of Mo, W and Re is dominant over most planetary conditions (Fig. 2), the mineral/melt partitioning behavior of these elements is strongly controlled by their tetravalent state (Fig. 8).

A common feature of the partitioning behavior of Mo, W and Re is that their 4+ valence state is significantly more compatible in silicates and oxides (e.g. spinel) than their 6+ valence state, the latter of which being almost perfectly incompatible in crystalline phases. The relative compatibility in silicates and oxides tends to follow the order  $D_{\text{cpx}} > D_{\text{opx}} > D_{\text{olv}} > D_{\text{spl}} \sim D_{\text{plg}}$ , independently of valence state (Fonseca et al., 2014; Leitzke et al., 2016; 2017; Mallmann and O'Neill, 2007; Wijbrans et al., 2015). For all three elements, the difference between the mineral/melt partition coefficient of their reduced and oxidized species spans three to five orders of magnitude (Fig. 8). The difference in partition coefficients between the valence states is most pronounced for clinopyroxene, with  $D_{6+} \sim 10^{-5}$  and  $D_{4+} \sim 3$ -6. The compatibility of tetravalent Re, Mo and W in clinopyroxene results from the nearly ideal size of these cations ( $r^{\text{VI}} = 0.66 \text{ \AA}$  for  $W^{4+}$ ,  $0.65 \text{ \AA}$  for  $Mo^{4+}$  and  $0.63 \text{ \AA}$  for  $Re^{4+}$ ) relative to the octahedral M1 site ( $r_0^{4+,M1} \sim 0.66 \text{ \AA}$ ; Hill et al., 2000), which makes the substitution into clinopyroxene lattice energetically favorable. While the mineral/melt partitioning of  $Re^{4+}$  and  $Mo^{4+}$  were determined directly by experimental data at reducing conditions (Mallmann and O'Neill, 2007; Leitzke et al., 2017), data for  $W^{4+}$  was constrained by Fonseca et al. (2014) using a lattice strain fit to Ti, Hf, Zr and Re. This is because the extremely reducing conditions necessary to reduce all  $W^{6+}$  to  $W^{4+}$  cannot be achieved in the laboratory. The important conclusion of these studies is that the partitioning behavior of these elements in magmatic systems will be completely different depending on  $fO_2$ , with significant implications for comparing lunar vs. terrestrial, or oceanic vs. arc magmatism.

The effect of composition on the mineral/melt partitioning of Mo, W and Re in pyroxenes has been addressed in some studies. Hill et al. (2000) and Righter and Shearer (2003) found a negative correlation between  $D_{\text{Mo,W}}$  and the Ca-Tschermak (CaTs) component in clinopyroxene. However, the opposite trend was found by Leitzke et al. (2017), who observed that the clinopyroxenes with the highest CaTs reported by Hill et al (2000) were formed from melt compositions with the highest CaO. Leitzke et al. (2017) speculated that, because the activity coefficients of  $MoO_3$  and  $MoO_2$  in silicate melts decrease with increasing CaO content (O'Neill and Eggins, 2002), the negative correlation found by Hill et al. (2000) could be due to the high CaO content of the silicate melt rather than pyroxene composition. Leitzke et al. (2016) investigated the effect of  $TiO_2$  in the silicate melt, observing a positive correlation between the  $TiO_2$  content of the melt on the partitioning of  $Mo^{4+}$  and  $W^{4+}$ , but little to no effect on the partitioning of  $Mo^{6+}$  and  $W^{6+}$  (perhaps because these cations are very incompatible). Conversely, partitioning coefficients for Hf and Zr decrease with  $TiO_2$  content of the melt. ( $D_{\text{Ti}}$  remained constant, as expected for Henry's Law behavior). Leitzke et al. (2016) argued that the introduction of  $CaTi^{4+}Al_2O_6$  into clinopyroxene with increasing melt  $TiO_2$  raises elastic strain in



the M1 site of clinopyroxene. The increased elastic strain (which results in a tighter parabola) leads to an increase in the partitioning values of cations whose size are close to ideal for that site (e.g.  $W^{4+}$ ), and a decrease in the partitioning values of those cations whose size are either smaller or larger than the ideal for that site (e.g.  $Zr^{4+}$  and  $Hf^{4+}$ ).

## 6.2. Platinum Group Elements (Ru, Rh, Pd, Os, Ir, Pt)

There have been few studies investigating the redox-dependent partitioning behavior of the Platinum Group Elements (PGE) during magmatism, with most focusing on the role of base metal sulfides and, to a lesser extent, oxides. The scarcity of partitioning data arises from the low solubility of PGE in silicate melts (typically  $< a few \mu g/g$ ), especially under reducing conditions (O'Neill et al., 1995), and from the almost unavoidable presence of sub-micrometer metallic nuggets that plague the analyses of PGEs in synthetic silicate glasses and often lead to erroneous solubility values (Ertel et al., 2008). The solubility of a PGE in a silicate melt in equilibrium with metal can give an indication of its valence state(s), as it will dissolve according to Equation 10 (with  $x = 0$  for the metal). However, despite significant effort in the last 20-30 years, PGE solubility results are still controversial. This leads to ambiguous interpretations regarding the PGE speciation as a function of  $fO_2$ . For example, solubility data for Ru are used to indicate that it either dissolves as  $Ru^{3+}$  (Borisov and Nachtweyh, 1998), or as both  $Ru^{3+}$  and  $Ru^{4+}$  (Laurenz et al., 2013) in silicate melts over planetary magmatic  $fO_2$ s. Similarly, Os appears to dissolve as  $Os^{3+}$  under most conditions, though the possibility of  $Os^{4+}$  being stable at more oxidizing conditions cannot be ruled out (Borisov and Walker, 2000; Fortenfant et al., 2006). Data for Ir are also ambiguous, indicating speciation as either  $Ir^{3+}$ , or combinations of  $Ir^{1+}$  -  $Ir^{3+}$  or  $Ir^{2+}$  -  $Ir^{3+}$  (Brenan et al., 2005; Borisov and Palme, 1995; Fonseca et al., 2011; O'Neill, 2015). Pt appears to occur exclusively as  $Pt^{2+}$ , except perhaps at very oxidizing conditions ( $FMQ > 5$ ) where  $Pt^{4+}$  may also be stable (Ertel et al., 1999). Of all PGEs, unambiguous evidence of redox-sensitivity over magmatic  $fO_2$ s only exists for palladium and rhodium, both of which occur as monovalent and divalent cations under typical magmatic  $fO_2$ s (e.g. Laurenz et al., 2010; Sommer, 2014; Fig. 2). Reports of other valence states for Pd and Rh in the literature arise from interpretation of micro(nano)nugget contaminated data (e.g. Borisov et al., 1994; Ertel et al., 1999).

Of the common magmatic minerals, spinel is the only mineral capable of accommodating significant amounts of PGEs, particularly Ru, Ir and Rh (e.g. Park et al., 2017). Wijbrans et al., (2015) reported a marked increase in  $D_{Rh}$  around FMQ in both Fe-free and Fe-rich spinels, indicating higher compatibility of  $Rh^{3+}$  relative to  $Rh^{2+}$ . Brenan et al. (2012) observed significant variations in the partitioning of several PGEs with  $fO_2$  for Fe- and Cr-bearing spinels. However, as noted above, the partitioning of trace elements into magnetite is controlled by the  $Fe^{3+}$  and  $Fe^{2+}$  content of the melt, and it is often difficult to deconvolve this effect, and that of drastically changing site occupancies in the spinel from that of valence state changes in the trace element of interest. Nevertheless, the data in the Fe-free system are consistent with a change from  $Rh^+$  to  $Rh^{2+}$ , and also with olivine/melt partitioning data that shows decreasing  $D_{Rh}$  with increasing  $fO_2$  (Brenan et al., 2003). Of the very few spinel/melt partitioning data available, Pd appears to be incompatible in spinel (Brenan et al., 2012).

## 7. Concluding remarks

In this chapter we have shown that the partitioning behavior of redox-sensitive trace elements will almost inevitably change with  $fO_2$  in magmatic systems. This is because the incorporation of cations into crystal lattices (at constant pressure, temperature and composition) depends on their charge, size, and crystal-field stabilization energies, all of which

likely vary substantially between the oxidized and reduced species of the same element. This dependence on  $fO_2$  means that petrologic/geochemical modelling involving elements such as Ti, Cr, V, Ce, Eu, U, Mo, Re, W, Pd, Rh (and probably others) must account for the oxidation state of the magma. As the only major-component redox-sensitive elements in many magmatic systems, the partitioning behavior of Fe is more likely to control rather than to reflect  $fO_2$ . Given the large range (10-15 order of magnitude; Fig. 2) of magmatic  $fO_2$  in planetary systems, significant differences in the partitioning behavior of individual elements ought to be expected. For instance, several studies have shown that certain trace elements behave in completely different ways during lunar compared to terrestrial magmatism, or during magmatism in mid-ocean ridge compared to subduction zone settings.

Nevertheless, if the partitioning of redox-sensitive trace elements is determined experimentally as a function of  $fO_2$  (and the effects of other parameters such as pressure, temperature and composition accounted for), then partitioning relations preserved in natural samples (e.g. glass and phenocryst in a volcanic rock) may provide useful ways of estimating the oxidation state of magmatic systems. This approach has the greatest advantage of being insensitive to late-stage degassing, charge-transfer (electron-exchange) reactions, or surficial alteration, all of which may potentially alter other magmatic redox proxies (e.g.  $Fe^{3+}/\Sigma Fe$ ).

The preceding discussion demonstrates the importance of a good understanding of redox in modelling or interpreting trace element concentrations in melts and minerals. However, it should also be apparent that further work is required to bring our ability to model the partitioning of redox-variable elements. First, there are several elements for which very little information is available: most notably the PGEs, which have been investigated in surprisingly few minerals and whose redox transitions are poorly constrained. Copper was excluded from this review because of a paucity of data, but the data of Liu et al. (2014, 2015) indicate that its partitioning changes significantly from FMQ+1 to FMQ+5, a range that is highly relevant to arc environments where essentially all the world's economic Cu deposits are located. Likewise, experimental partitioning data by Mallmann and O'Neill (2009) indicate a potential change of phosphorus valence states from 5+ to 3+ under conditions below FMQ-4, but the results could also be explained by volatility loss from the melt after crystallization. And, under ultra-reducing conditions ( $\sim$  FMQ-7), where Nb and Ta begin to show siderophile tendencies, there is evidence that they may take on lower valence states in silicate melts (Cartier et al., 2014).

Furthermore, the dependence of these redox transitions and partition coefficients on temperature and melt composition is complex, and only partially understood for the majority of the redox couples addressed here. As seen in Figure 4, melt composition affects the partitioning of different oxidation states of the same element in differing ways, an effect that is additional to the influence of stoichiometric controls on trace element partitioning. Similarly, melt compositional effects can shift the redox transition by up to 6 log units of  $fO_2$  (Fig. 2). Because of these complications, extreme caution should be taken when attempting to study partitioning of redox-variable elements outside their calibrated ranges.

Many factors are important in selecting a mineral/melt partitioning oxybarometer. First, the anticipated redox state of the system will give an indication of which redox couples are relevant, but a large contrast between the compatibilities of the valence states of the element can sufficiently outweigh this consideration (for example, zircon will develop a Ce anomaly at conditions that are many orders of magnitude more reducing than the midpoint of the  $Ce^{3+}$ - $Ce^{4+}$  equilibrium). The figures presented in this chapter may guide the geoscientist in determining whether the partition coefficient is likely to be sensitive to  $fO_2$  for her or his samples. As emphasized above, the availability of appropriate experimental constraints is important, as are more general petrological considerations, for instance resistance to alteration processes such as

diffusive re-equilibration. In the ideal case, multiple elements and minerals would be combined to check for internal consistency.

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## Figure captions

**Figure 1. (a)** Schematic illustration of lattice strain parameters  $D_0$ ,  $r_0$  and  $E$  for the mineral/melt partition coefficient ( $D$ ) of an element  $i$  as modelled by the Brice equation (Blundy and Wood, 1994);  $D_i = D_0 \exp([-4\pi EN_A(r_0 - r_i)^2 - 1/3(r_0 - r_i)^3]/RT)$ , and **(b, c, d)** measured changes in these parameters with cation valence state (i.e. charge) as determined experimentally for  $\text{CaSiO}_3$ /melt partitioning (Law et al., 2000).

**Figure 2.** Relative oxygen fugacity conditions (in log units relative to FMQ, the fayalite-magnetite-quartz buffer; values from O'Neill, 1987) for several planetary geochemical reservoirs (Frost and McCammon, 2008; Grossman et al. 2008; Wadhwa, 2008), and valence state transition ranges for the elements discussed in this chapter. See text for references. Also plotted for reference the IW (Fe-FeO), NNO (Ni-NiO), and HM (hematite-magnetite) buffers.

**Figure 3:** Mineral/melt partitioning of transition metals as a function of oxygen fugacity (in log units relative to FMQ, the fayalite-magnetite-quartz buffer; values from O'Neill, 1987) for selected minerals. Abbreviations: amph - amphibole, cpx - clinopyroxene, opx - orthopyroxene, olv - olivine, plag - plagioclase feldspar, spl - spinel, zrc - zircon. All curves shown in the main diagram are best fits of the experimental data to an equation with the same form as Equation 22. The fraction (from 0 to 1) of the different valence states of each element are plotted on top of the diagram. These were calculated using Equation 13 using available data. **(a)** Fe mineral/melt partitioning data from King et al., 2000 (amph), Mallmann and O'Neill, 2009 (spl,

olv, opx, cpx) and Phinney, 1992 (plag). **(b)** Cr mineral/melt partitioning data from Mallmann and O'Neill, 2009 (olv, opx, cpx). **(c)** Ti mineral/melt partitioning data from Mallmann and O'Neill, 2009 (olv, opx, cpx), Peters et al. 1995 (plag) and Burnham and Berry, 2012 (zrc). **(d)** (d) V mineral/melt partitioning data from Mallmann and O'Neill, 2009 (olv, opx, cpx, spl).

**Figure 4.** Comparison of Cr olivine( $\text{Fo}_{100}$ )/melt partitioning coefficients obtained at 1 bar and 1320 °C as a function of oxygen fugacity (in log units relative to FMQ, the fayalite-magnetite-quartz buffer; values from O'Neill, 1987) for three Fe-free compositions in the  $\text{CaO-MgO-Al}_2\text{O}_3\text{-SiO}_2$  system: FAS1, FAD1 and FAD3 (Hanson and Jones, 1998). The results illustrate a mild effect of melt composition on the partitioning of  $\text{Cr}^{3+}$  yet no significant effect on the partitioning of  $\text{Cr}^{2+}$ . The fraction (from 0 to 1) of the different valence states are plotted on top of the diagram.

**Figure 5.** Zircon/melt partition coefficients for REE elements determined experimentally at 1 bar as a function of oxygen fugacity by Burnham and Berry (2012). The colored segments highlight systematic changes in the partitioning of Ce and Eu with oxygen fugacity (labelled relative to the FMQ buffer).

**Figure 6.** Mineral/melt partitioning of redox-sensitive rare earth elements (REE), Ce **(a)** and Eu **(b)** as a function of oxygen fugacity (in log units relative to FMQ, the fayalite-magnetite-quartz buffer; values from O'Neill, 1987) for selected minerals. Abbreviations: bdd - baddeleyite, cpx - clinopyroxene (aluminous diopside), olv - olivine, opx - orthopyroxene, plag - plagioclase feldspar, ttn - titanite, zrc - zircon. Where data were not obtained as a function of  $f\text{O}_2$ , partition coefficients for  $\text{Ce}^{3+}$ ,  $\text{Ce}^{4+}$ ,  $\text{Eu}^{2+}$  and  $\text{Eu}^{3+}$  (or estimates thereof) were used in conjunction with the  $\text{Ce}^{4+}/\Sigma\text{Ce}$  and  $\text{Eu}^{3+}/\Sigma\text{Eu}$  models of Burnham and Berry (2014) and Burnham et al. (2015). Data from Klemme and Meyer, 2003 (bdd), Leitzke et al., 2017 (cpx-Ce, plag-Ce, olv), Sun et al., 1974 (cpx-Eu, plag-Eu, opx-Eu), van Kan Parker et al., 2011 (opx-Ce), Bachmann et al., 2005 (ttn), and Burnham and Berry, 2012 (zrc). See Fig. 3 for more details.

**Figure 7.** Mineral/melt partitioning of U as a function of oxygen fugacity (in log units relative to FMQ, the fayalite-magnetite-quartz buffer; values from O'Neill, 1987) for selected minerals. Abbreviations: cpx - clinopyroxene (diopside), olv - olivine, opx - orthopyroxene, zrc - zircon. Data from Fonseca et al., 2014 (opx, cpx, olv), and Burnham and Berry, 2012 (zrc). See Fig. 3 for more details.

**Figure 8.** Mineral/melt partitioning of siderophile elements (Mo, W and Re) as a function of oxygen fugacity (in log units relative to FMQ, the fayalite-magnetite-quartz buffer; values from O'Neill, 1987) for selected minerals. Abbreviations: cpx - clinopyroxene, opx - orthopyroxene, olv - olivine, spl - spinel, grt - garnet. Data from Leitzke et al., 2017 (Mo: cpx, opx, olv), Wijbrans et al., 2015 (Mo: spl), Fonseca et al., 2014 (W: opx, cpx, olv), and Mallmann and O'Neill, 2007 (Re: cpx, opx, spl, grt, ol). See Fig. 3 for more details.

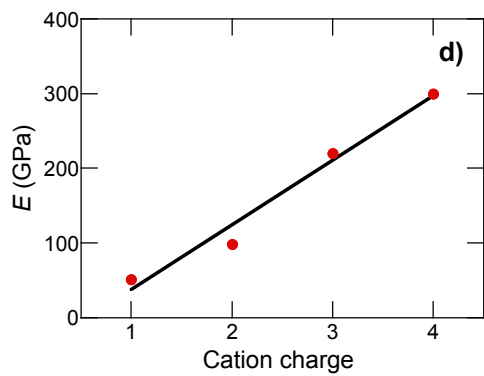
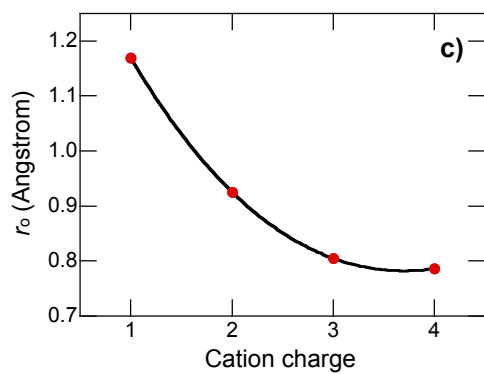
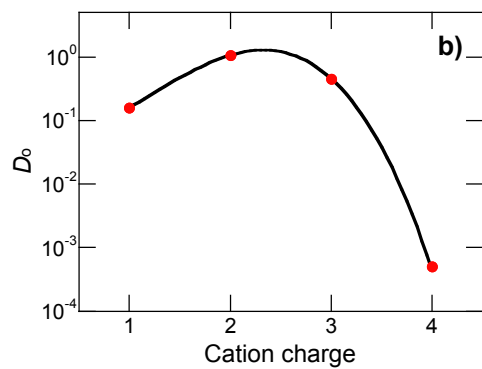
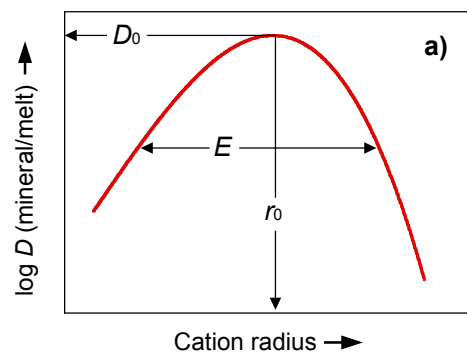


Figure 2

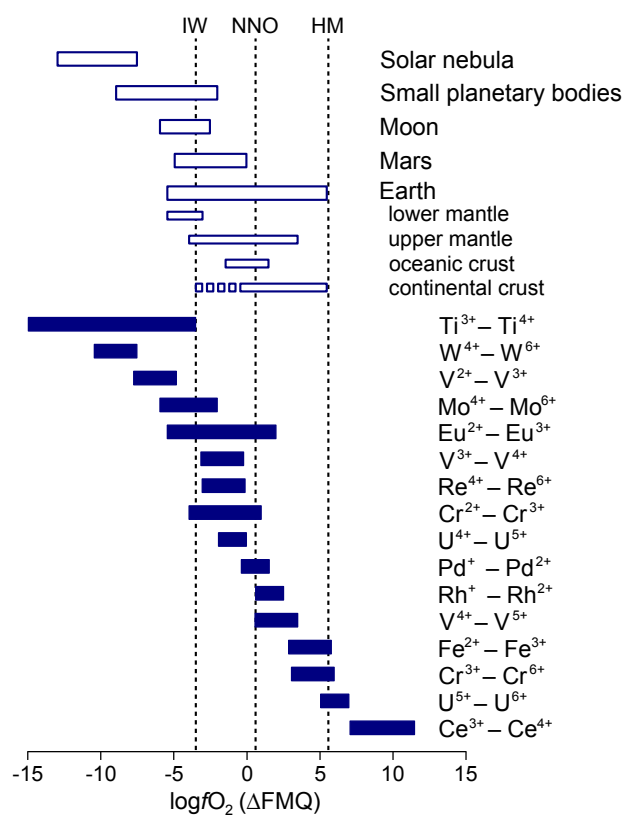


Figure 3

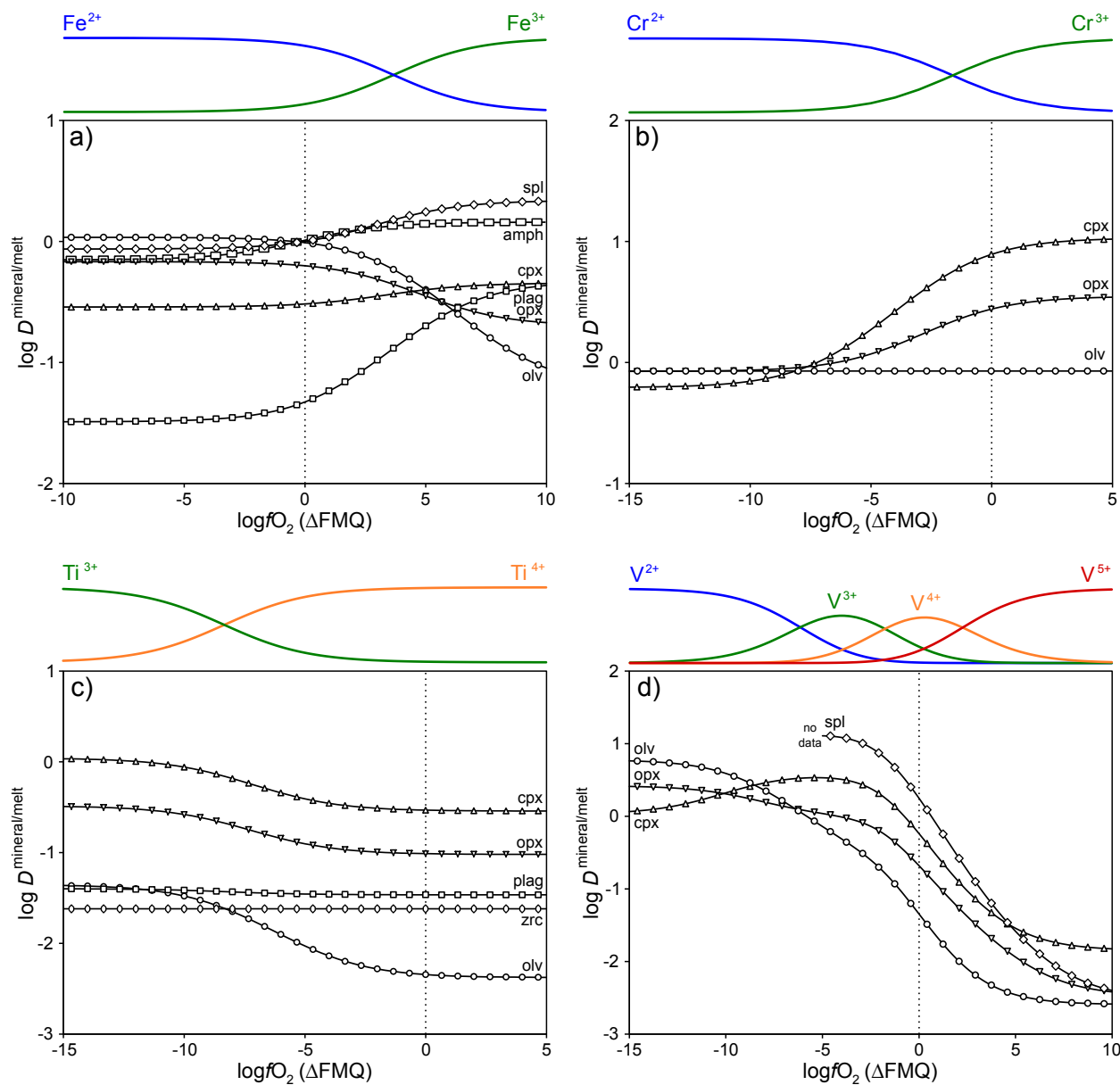




Figure 4

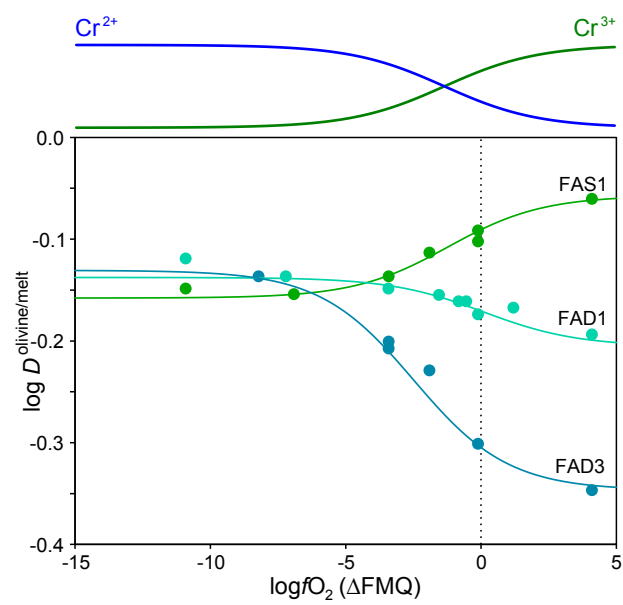


Figure 5

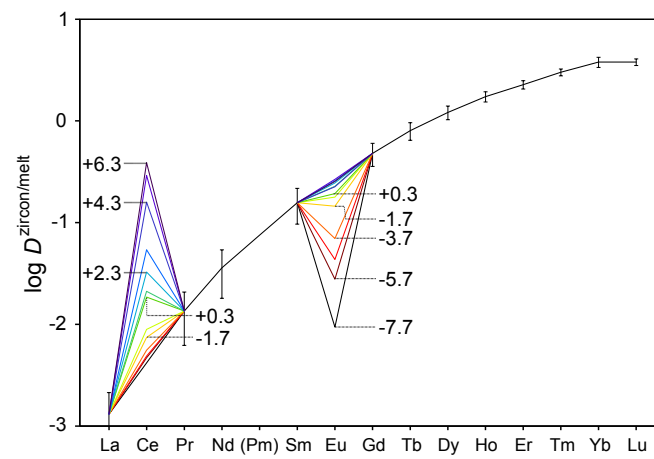


Figure 6

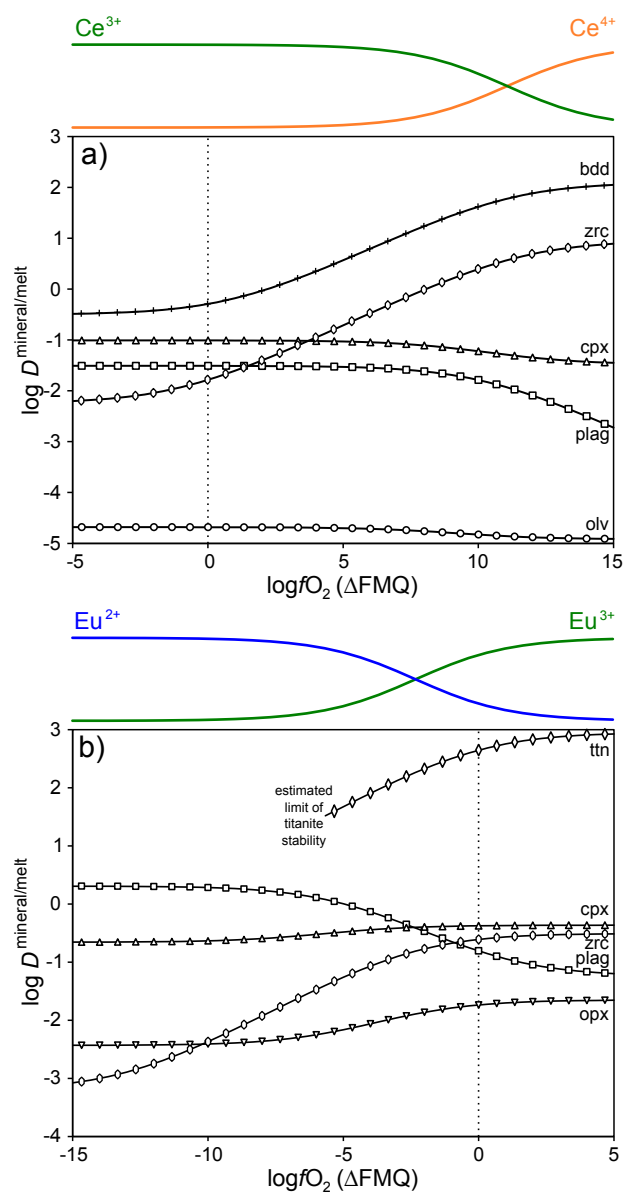


Figure 7

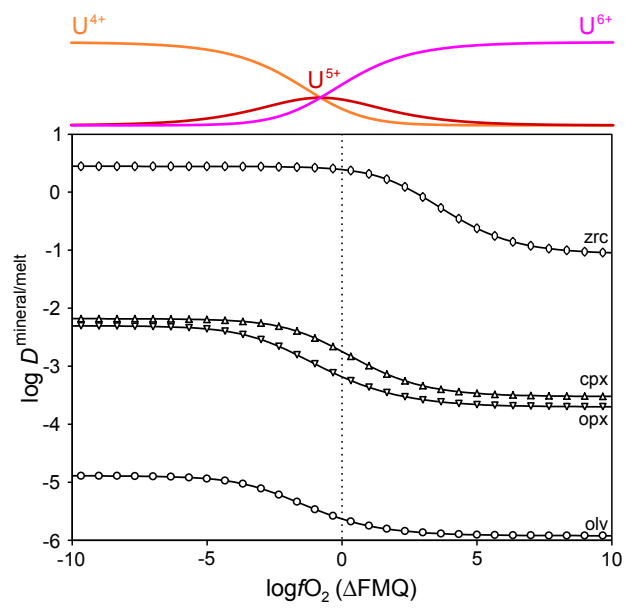


Figure 8

