

Cloud Process Coupling and Time Integration in the E3SM Atmosphere Model

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Key Points:

- Changing time step has as large an effect on E3SMv1's climate as doubling horizontal resolution.
- Only some of the whole model's time step sensitivity can be attributed to individual processes.
- Coupling between processes is the most important source of model time step sensitivity.

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14 **Abstract**

15 In this study, we find significant sensitivity to the choice of time step for the Energy Ex-
 16 ascale Earth System Model’s atmospheric component, leading to large decreases in the
 17 magnitude of cloud forcing when the time step is reduced to 10 seconds. Reducing the
 18 time step size for the microphysics increases precipitation, leading to a drying of the at-
 19 mosphere and an increase in surface evaporation. This effect is amplified when the mi-
 20 crophysics is substepped together with other cloud physics processes. Coupling the model’s
 21 dynamics and physics more frequently reduces cloud fraction at lower altitudes, while
 22 producing more cloud liquid at higher altitudes. Reducing the deep convection time step
 23 also reduces low cloud mass and cloud fraction. Together, these results suggest that cloud
 24 physics in a global circulation model can depend strongly on time step, and in partic-
 25 ular on the frequency with which cloud-related processes are coupled with each other and
 26 with the model dynamics.

27 **Plain Language Summary**

28 Computer simulations of the Earth’s atmosphere take the state of the atmosphere
 29 at one point in time, then predict the state of the atmosphere a short interval of time
 30 into the future. The length of this time interval is known as the “time step”. By doing
 31 this repeatedly, models can produce a simulated history of the atmosphere for years or
 32 even centuries. A smaller time step size requires more computer time, but should ide-
 33 ally lead to more accurate results. In this study we reduce the time step for the atmo-
 34 sphere in the Energy Exascale Earth System Model from half an hour to ten seconds.
 35 The simulation with a smaller time step has more rain, which removes water from the
 36 atmosphere and reduces the fraction of the Earth’s surface that is covered by clouds. We
 37 also experiment with changing the time step for only some parts of the model and not
 38 others. We find that the effects of the time step size on the model are related mostly to
 39 the frequency of coupling between processes rather than the time step used for any in-
 40 dividual processes.

41 **1 Introduction**

42 Time integration strategies for general circulation models (GCMs) have grown more
 43 complex, both due to an increase in the number of separate processes within a model
 44 (allowing individual parameterizations to adopt different time integration strategies from
 45 the “host” model), and due to changes in process coupling, such as modifications that
 46 allow different processes to be run concurrently (Balaji et al., 2016; Donahue & Cald-
 47 well, 2020). This increase in complexity means that there is often no single time step for
 48 the model, but rather a set of interrelated time steps controlling the rate at which var-
 49 ious calculations are performed and allowed to interact. This increased complexity can
 50 be daunting, but it also provides the opportunity to conduct more detailed experiments
 51 regarding the effect of temporal resolution on GCMs. Past research has established that
 52 certain cloud processes can be disproportionately responsible for time integration error
 53 in GCMs (Wan et al., 2015). When the time steps used for particular processes can eas-
 54 ily be adjusted independently from one another, it becomes possible to study the time
 55 step sensitivity of these processes with minimal changes to the model code.

56 Prior research using the Community Atmosphere Model, versions 3 and 4, (CAM3/4),
 57 as well as its predecessor, the Community Climate Model, version 3, suggests that time
 58 step size for a GCM has significant effects on precipitation in an aquaplanet simulation,
 59 particularly in the intertropical convergence zone (ITCZ) (Williamson & Olson, 2003;
 60 Williamson, 2008; Mishra et al., 2008). Decreasing the time step increases total precip-
 61 itation in the ITCZ and the frequency of extreme precipitation events. The partition of
 62 precipitation between the large-scale and convective parameterizations is also affected,
 63 with an increase in the large-scale precipitation being responsible for the aforementioned

64 effects. For CAM3, the increase in total precipitation was found to be dependent on an
 65 increase in evaporation at the surface, which in turn was due to an increase in wind speed
 66 and a decrease in specific humidity near the surface (Mishra et al., 2008). Further re-
 67 search using CAM3 for real-planet simulations confirmed these results, and showed that
 68 the increased evaporation, in addition to producing increased precipitable water (and
 69 precipitation), also produced a larger cloud fraction at low altitude and an increased mag-
 70 nitude of radiative cloud forcing (Mishra & Shanay, 2011). Williamson (2013), using CAM4,
 71 found that the time step of the convective schemes, and in particular their rate of cou-
 72 pling with other processes, controlled the repartitioning of precipitation between large-
 73 scale and convective processes.

74 Yu and Pritchard (2015) experimented with changes in the global model time step
 75 for a superparameterized version of CAM3 (SPCAM3), without changing its cloud res-
 76 solving model’s time step. Decreasing the CAM time step increased overall precipita-
 77 tion in SPCAM3, as well as the frequency of heavy precipitation events, which also hap-
 78 pened for CAM3. However, reducing the time step size in SPCAM3 *decreased* precip-
 79 itable water, decreased both liquid and ice water path, and ultimately decreased the mag-
 80 nitude of radiative cloud forcings. This was hypothesized to be due to changes in con-
 81 vective organization, producing an increase in precipitation efficiency and effectively dry-
 82 ing out the atmosphere.

83 Although both CAM3 and SPCAM3 experience similar increases in surface evap-
 84 oration and precipitation (which must match in the long run, to balance the water bud-
 85 get), the proposed mechanisms driving these changes are different. In effect, the increased
 86 evaporation in CAM3 ”pushes” more precipitable water into the atmosphere, eventually
 87 forcing an increase in condensation and precipitation to remove this water, while the in-
 88 creased precipitation efficiency in SPCAM3 ”pulls” water out of the atmosphere, dry-
 89 ing it out and encouraging evaporation to replace the lost water.

90 This paper analyzes time step sensitivity in the Energy Earth System Model (E3SM),
 91 which shares very little of its physics with CAM3, and almost none with SPCAM3. Nev-
 92 ertheless, E3SM is descended from CAM3 and uses the same general strategies for cou-
 93 pling the physics parameterizations, dynamics, and surface components. We have recently
 94 examined the effect of time step size on a specific parameterization used by E3SM, the
 95 Morrison-Gettelman microphysics version 2 (MG2). We found that the total precipita-
 96 tion was not sensitive to changes to the MG2 time step alone, though we did see a mild
 97 increase in the ratio of stratiform to convective precipitation (Santos et al., 2020). While
 98 the details of the precipitation physics and the vertical distribution of rain mass changed
 99 significantly at small time steps, the effect on total precipitation reaching the surface was
 100 mild.

101 2 Model Description

102 E3SM version 1 (E3SMv1) is an earth system model developed by the U.S. Depart-
 103 ment of Energy, focusing on three main research topics: (1) the water cycle, (2) the cryosphere,
 104 and (3) biogeochemistry (Golaz et al., 2019). This study focuses on E3SMv1’s atmosphere
 105 model, EAMv1 (Rasch et al., 2019; Xie et al., 2018). For a run at ~ 100 km atmospheric
 106 resolution (1°), the standard time step for coupling between the physics parameteriza-
 107 tions, dynamical core, and surface parameterizations is 30 minutes (1800 seconds), with
 108 the various processes and dynamical core typically coupled using a “sequential split” method.
 109 This means that for each time step, each parameterization accepts a state that has al-
 110 ready been updated by applying the effects of previous parameterizations, and so at large
 111 time steps, the model physics depends on the order in which these updates are applied
 112 (Donahue & Caldwell, 2018). (This should be irrelevant for sufficiently small time steps,
 113 assuming that the model converges in this limit to the true solution of its spatially dis-
 114 cretized motivating equations.)

Component Name	Component Function(s)	Time Step Size(s)
Spectral element dynamics (HOMME)	Dynamics Tracer Advection	Vertical remapping: 900s Dynamics/Advection: 300s Hyperviscosity subcycle: 100s
Cloud Layers Unified By Binormals (CLUBB)	Turbulence Shallow Convection Stratiform Clouds	300s (looped with MG2)
Zhang-McFarlane scheme (ZM)	Deep Convection	1800s
Four-mode Modal Aerosol Module (MAM4)	Aerosols	1800s
Morrison-Gottelman scheme version 2 (MG2)	Stratiform Microphysics	300s (looped with CLUBB)
Linearized ozone chemistry (LINOZ2)	Ozone chemistry	1800s
Rapid Radiative Transfer Model for GCMs (RRTMG)	Radiative transfer	1800s/3600s (tendencies are only recalculated once per hour, but applied every time step)
Gravity wave scheme	Gravity wave propagation and breaking	1800s

Table 1. EAMv1 parameterizations and corresponding time steps for a $\Delta x \approx 100$ km run using default settings

115 Table 1 shows a summary of the main parameterizations of EAMv1, as well as the
116 relevant time steps for these schemes when default options are used. This is by neces-
117 sity a broad overview, since each of these parameterizations is a complex piece of soft-
118 ware in its own right. Note also that the CLUBB and MG2 parameterizations are sub-
119 stepped within a single loop, so each "sees" the updates from the other during each of
120 their smaller 300 second time steps.

121 The climate of EAMv1 using default settings is extensively documented, for instance
122 in Golaz et al. (2019), Rasch et al. (2019), and Xie et al. (2018). This work will focus
123 on the climatological differences between various modified runs using shorter time steps
124 and a control EAMv1 run at default settings.

125 This approach allows us to broadly discuss the magnitude and nature of time inte-
126 gration error in E3SM. However, we note that the model was developed and tuned for
127 much longer time steps than we use here, and therefore the tuning parameters are likely
128 set so as to partially cancel this time integration error. Thus, while the ALL10 run de-
129 scribed below should have a dramatically lower *time integration* error, we do not assert
130 that it is "better" than the default configuration as a production model, and in fact it
131 would require significant retuning (at a minimum) to be usable at all.

132 3 Methodology

133 We ran E3SM at a ~ 100 km atmospheric resolution (ne30_ne30 grid) with standard
 134 E3SMv1 tuning and prescribed sea surface temperature. These runs were performed using
 135 a maintenance version of E3SM 1.1 (hash 25c94366) with pre-industrial forcings (compset
 136 F1850C5AV1C-04P2) unless otherwise specified, and all had prescribed sea surface tem-
 137 perature (SST) and sea ice extent.

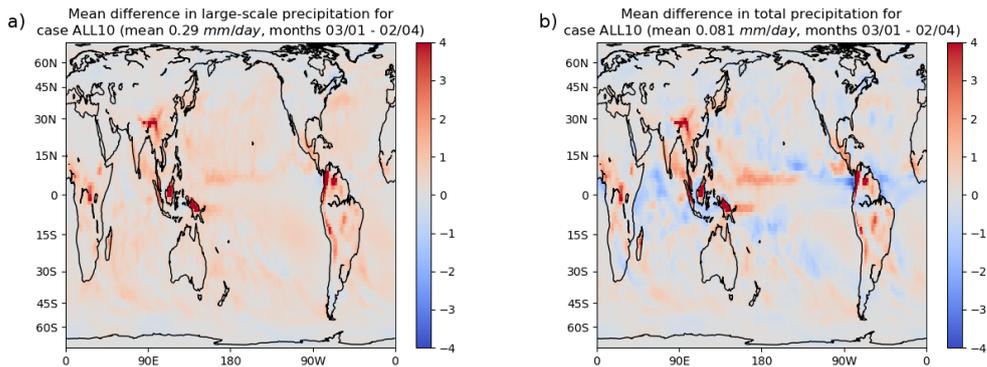
138 The control run (CTRL) is an out-of-the-box run using default settings. We com-
 139 pare this to a new run (ALL10) that changed the atmosphere’s dynamics-physics-surface
 140 coupling time step, known as “dtime” in the model settings, to 10 s. This time step is
 141 chosen to be as small as we could reasonably afford, within the constraints of the com-
 142 puting power we had available for this study. CTRL and ALL10 simulations are both
 143 3 years in length. Results based on years 1-2 were unchanged after adding year 3, which
 144 gives us confidence that 3 years is long enough to draw robust conclusions.

145 The dtime setting is often thought of as specifying the entire atmosphere’s time
 146 step, but there are three ways in which this is not quite true. First, the radiation param-
 147 eterization uses a fixed time step of once per hour, regardless of dtime. While the ALL10
 148 run does not modify the radiation time step, our tests with shorter runs show that the
 149 model is not especially sensitive to this time step. Second, the dynamics and cloud physics
 150 contain some substepping by default, though none run at a time step as small as 10 s.
 151 In the ALL10 run, we disable these forms of substepping, so that all major dynamics and
 152 physics processes aside from radiation run at the same small time step. Third, because
 153 the dtime setting also governs the rate at which EAMv1 exchanges information with the
 154 surface components, changing it forces a change in the land and sea ice components, which
 155 must run at a 10 s time step as well.

156 To investigate which processes within EAMv1 were most responsible for its over-
 157 all time step sensitivity, we configured a series of runs using built-in options to substep
 158 individual parameterizations at a 10 s time step. These runs include DYN10, which sub-
 159 stepped EAMv1’s spectral element dynamical core, CLUBB10, which substepped the CLUBB
 160 cloud physics parameterization, and MICRO10, which substepped the MG2 stratiform
 161 microphysics. Typically CLUBB and MG2 are substepped together within a single loop,
 162 so we also produced a CLUBBMICRO10 run that used this capability to run both schemes
 163 at 10 s. Because this run showed clear differences from the control, we produced a CLUBB-
 164 MICRO60 run where CLUBB and MG2 were substepped together at 60 s. Separately,
 165 we also produced a CLUBB10MICRO10 run, where the individual time steps for CLUBB
 166 and MG2 were reduced, but the two schemes were only coupled at the default rate of
 167 once per 300 s. To verify our prior belief that the model is less sensitive to the radiation
 168 time step size than to other physics time steps, we produced the ALLRAD10 run, which
 169 is identical to the ALL10 run except that the radiation is also run at a 10 second time
 170 step. Finally, we produced the ALL300 and ALL60 runs, which change dtime to 300 s
 171 and 60 s, respectively. ALL300 is useful as an additional control, because it decreases the
 172 dynamics-physics coupling time step, but does not decrease the CLUBB or MG2 time
 173 steps, nor does it decrease the dynamics time steps (except the remapping for the ver-
 174 tically Lagrangian advection scheme, which normally runs every 900 seconds). A sum-
 175 mary of all these runs can be found in Table 2.

176 In EAMv1, the most important parameterization that does *not* have this built-in
 177 substepping capability is the ZM deep convection, which always runs at the time step
 178 given by dtime. Substepping this parameterization individually is challenging and is de-
 179 scribed further in section 4.3.

Name	Substepped processes	Substep size	Run length
CTRL	None	N/A	38 mo
ALL10	Dynamics-physics coupling	10 s	38 mo
ALL60		60 s	30 d
ALL300		300 s	30 d
DYN10	Dynamics and tracer advection	10 s	30 d
MICRO10	MG2 microphysics	10 s	30 d
CLUBB10	CLUBB unified cloud parameterization	10 s	30 d
CLUBB10MICRO10	CLUBB and MG2	10 s	30 d
CLUBBMICRO10	CLUBB+MG2 combined loop	10 s	30 d
CLUBBMICRO60		60 s	
ALLRAD10	Dynamics-physics coupling and radiation	10 s	30 d

Table 2. Runs performed using the default aerosol scheme.**Figure 1.** Differences in a) large-scale precipitation and b) total precipitation between CTRL and ALL10 runs.

4 Results

4.1 Effects of Decreasing the Physics Time Step

Consistent with the literature on CAM, we find a 0.08 mm/d increase in global-mean total precipitation with 10 s time steps, mostly stemming from low latitudes. While convective precipitation decreases, large-scale precipitation increases by $\sim 60\%$ in the tropics (defined here as latitudes from 30S to 30N), especially over land (Figure 1a). The overall effect is an increase in tropical precipitation over the Pacific warm pool, and a near-doubling of precipitation in parts of Borneo, New Guinea, and Colombia (Figure 1b).

Evaporation must increase to fuel the precipitation increases. Over the oceans, this is due to slightly higher wind speed and lower near-surface relative humidity. The average 10 m wind speed in the tropics increases from 5.56 m/s to 5.70 m/s, and occurs mainly

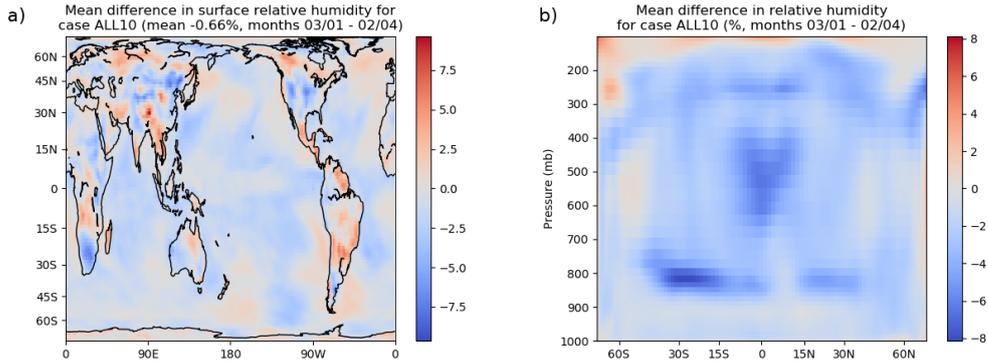


Figure 2. Differences in a) relative humidity and b) zonal mean relative humidity between CTRL and ALL10 runs.

191 in the northern Indian ocean and central Pacific, while latent heat flux increases in the
 192 same areas, by about 4 W/m^2 (not shown). This is consistent with the CAM3 literature,
 193 suggesting that an increase in wind speed contributes an increase in evaporation for short
 194 time steps.

195 Unlike in CAM3, the relative humidity decreases throughout the troposphere in the
 196 ALL10 run, as shown in Figure 2. The decreases in the 800-850 hPa layer correspond
 197 to a significant reduction in low cloud mass (Figure 3). This suggests that the increase
 198 in precipitation is caused primarily by increased precipitation efficiency. At the same time,
 199 we see an increase in cloud liquid above the boundary layer, particularly at high latitudes.
 200 In the ALL10 run, the cloud fraction not only decreases at lower levels, where less liq-
 201 uid cloud is present, but also at higher levels, where the average cloud mass mixing ra-
 202 tios are similar to or greater than their values in CTRL (Figure 4).

203 Consistent with the decrease in overall cloud mass and fraction, the ALL10 run shows
 204 substantially reduced radiative cloud forcing compared with CTRL, as can be seen in
 205 Figure 5. The effect on shortwave cloud forcing is especially large, the global mean be-
 206 ing reduced from -43.0 W/m^2 to -37.5 W/m^2 .

207 Spatial-pattern differences between CTRL and ALL10 are summarized by a Tay-
 208 lor diagram shown in Figure 6. We find that the effect of reducing the model time step
 209 to 10 seconds (black symbols) is comparable to the effect of doubling the model's hor-
 210 izontal grid spacing (red symbols), a natural standard for comparison. Historically, spa-
 211 tial resolution changes have been perceived as a major model change while accompany-
 212 ing time step changes have been taken for granted; Figure 6 illustrates that this view-
 213 point has shortcomings. The variables most affected by the time step are related to pre-
 214 cipitation, with the large-scale precipitation showing the most difference.

215 4.2 Effects of Changing the Physics Substepping

216 We did not have the computational resources to run all of our substepped config-
 217 urations for multiple years, but many of the effects of substepping are quite large and
 218 can easily be distinguished with only a few days of data. Since all runs started with the
 219 same initial conditions, they followed a similar trajectory for the first 15 days before be-
 220 ginning to diverge, so we focused on comparing the runs during this time period. (In Fig-
 221 ure 12, we have included data from 30 days, which shows an example of how, in these
 222 different runs, the global means of cloud-related variables become significantly less cor-
 223 related towards the end of the month.)

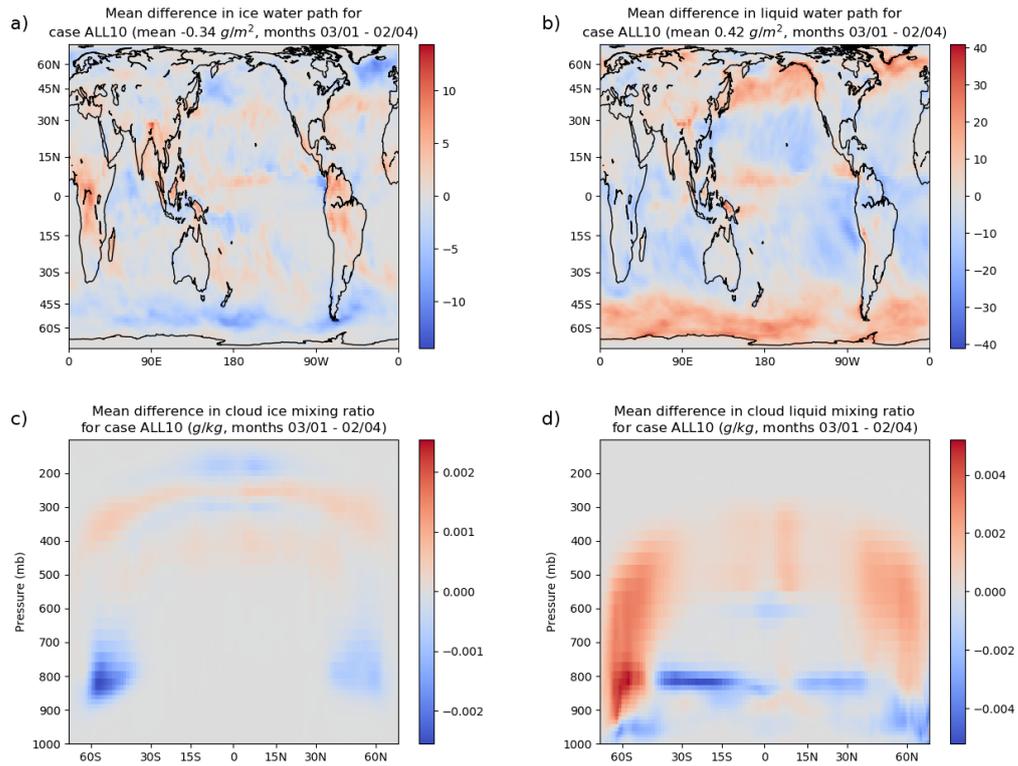


Figure 3. Differences in mass of cloud ice and liquid water between CTRL and ALL10 runs, measured by a) ice water path, b) liquid water path, c) zonal mean cloud ice mixing ratio, and d) zonal mean cloud liquid mixing ratio.

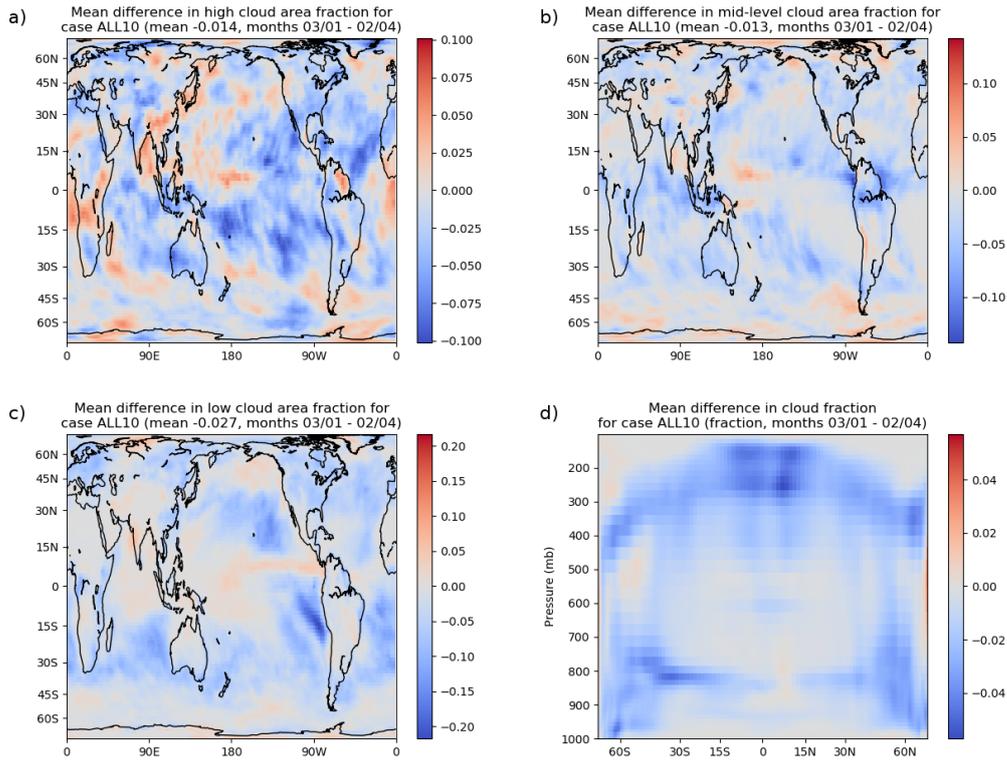


Figure 4. Differences in cloud fraction between CTRL and ALL10 runs: a) high cloud fraction (defined as the vertical integral above 400 mb), b) mid-level cloud fraction (vertical integral over the range 400-700 mb), c) low cloud fraction (vertical integral below 700 mb), and d) zonal mean cloud fraction at each level.

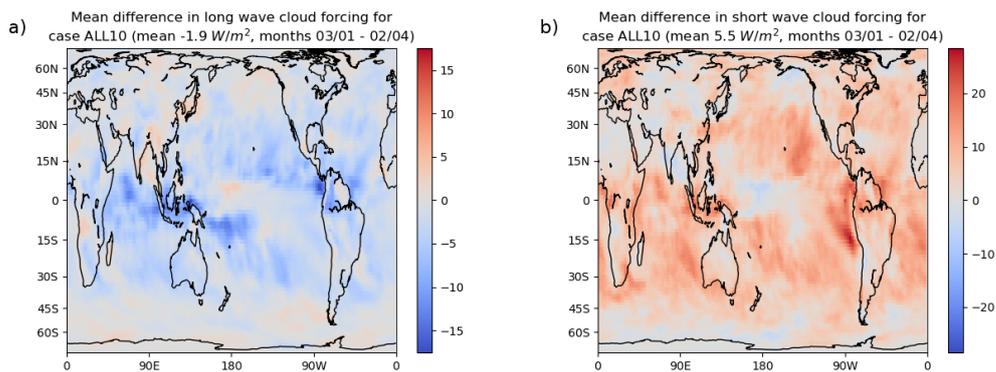


Figure 5. Differences in a) longwave cloud forcing and b) shortwave cloud forcing between CTRL and ALL10 runs.

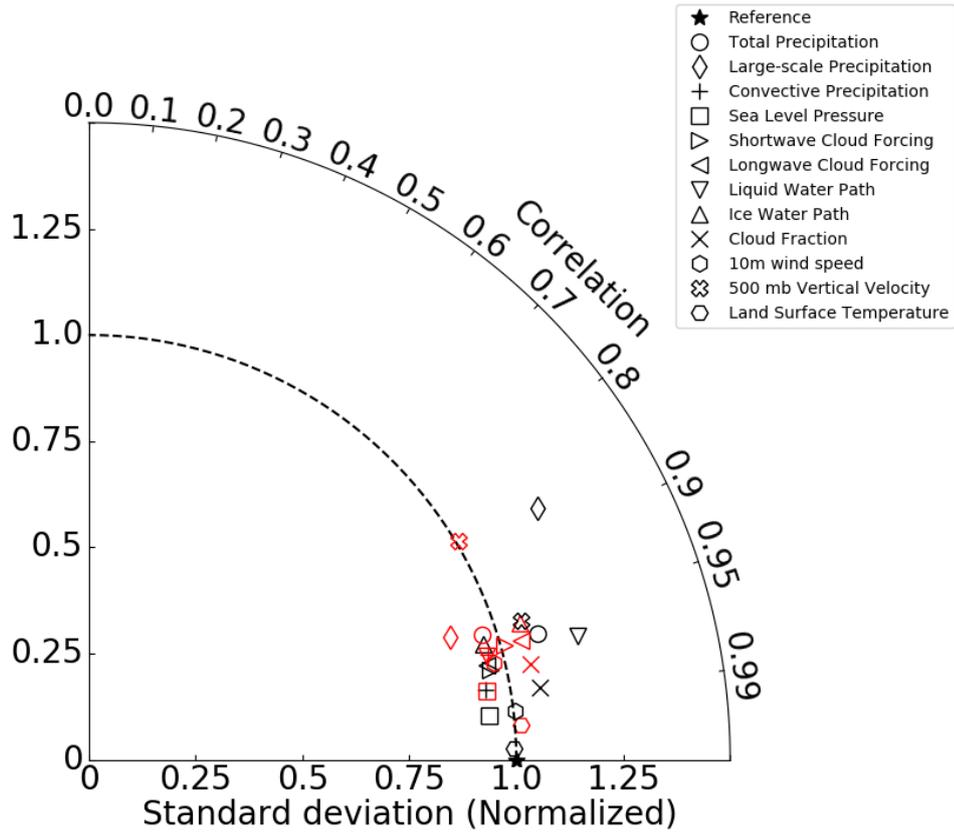


Figure 6. Taylor diagram comparing results from the CTRL and ALL10 runs (black), and comparing results from CTRL to a run with default settings using the ne16 grid (red), which has a grid spacing of $\sim 1.9^\circ$. This diagram uses values averaged over a three year time period starting March of the first simulated year for each run; only spatial variability is accounted for.

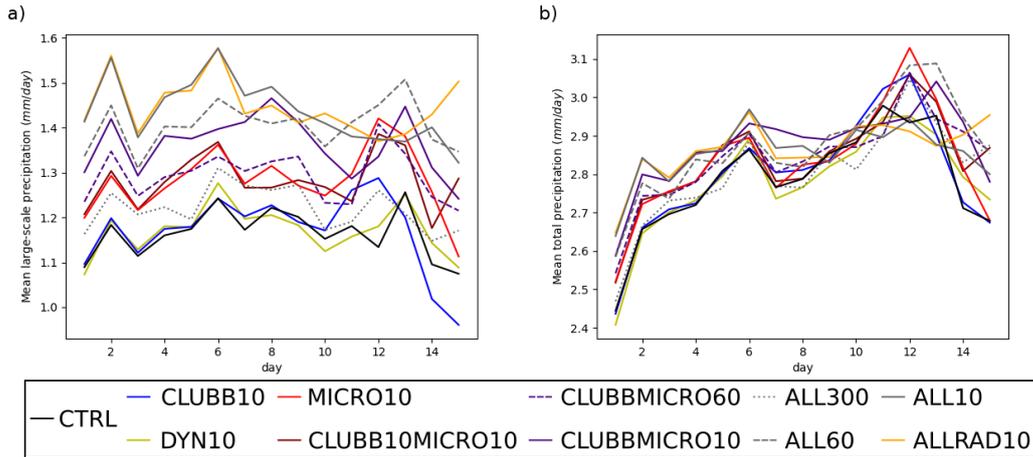


Figure 7. Daily global means of a) large-scale precipitation only and b) total precipitation for substepped runs.

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We first examine the changes in precipitation across these runs, shown in Figure 7. We can categorize our runs into five main categories based on large-scale precipitation:

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1. The lowest average large-scale precipitation rates are found in CTRL, CLUBB10, and DYN10, suggesting that the cloud physics is not sensitive to the CLUBB or dynamics time steps.
2. A slightly higher large-scale precipitation rate is found in the ALL300 run, which has a reduced time step for the ZM deep convection and an increased dynamics-physics coupling frequency, but does not change the CLUBB or MG2 time steps. This run shows a mild repartitioning of precipitation from the convective to large-scale category, perhaps consistent with the mechanisms described in Williamson (2013).
3. The MICRO10 and CLUBB10MICRO10 runs show signs of increased precipitation efficiency overall, increasing both large-scale *and* total precipitation.
4. One of the largest large-scale precipitation rates comes from CLUBB10MICRO10, which couples CLUBB to MG2 every 10 seconds, suggesting that an increase in CLUBB-MG2 coupling frequency has an additional effect beyond that from simply substepping MG2 more frequently. CLUBB10MICRO60 may see a similar effect, though it is not as clearly distinguished from the MICRO10 run.
5. The ALL10 and ALLRAD10 runs decrease both the dynamics-physics coupling time step and the time step used for all physics parameterizations, and these show the largest changes in precipitation.

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While the increase in large-scale precipitation seen in the MICRO10 run is substantial, the spatial pattern is quite different from the ALL10 case, as seen in Figure 8. In particular, the patterns observed in the ALL10 case, such as the increase in large-scale precipitation over land, only appear when both CLUBB and MG2 are substepped together, as in the CLUBB10MICRO10 case. We have also plotted the CLUBB10MICRO10 case, to show that this difference is not due simply to the reduction in the CLUBB time step alone.

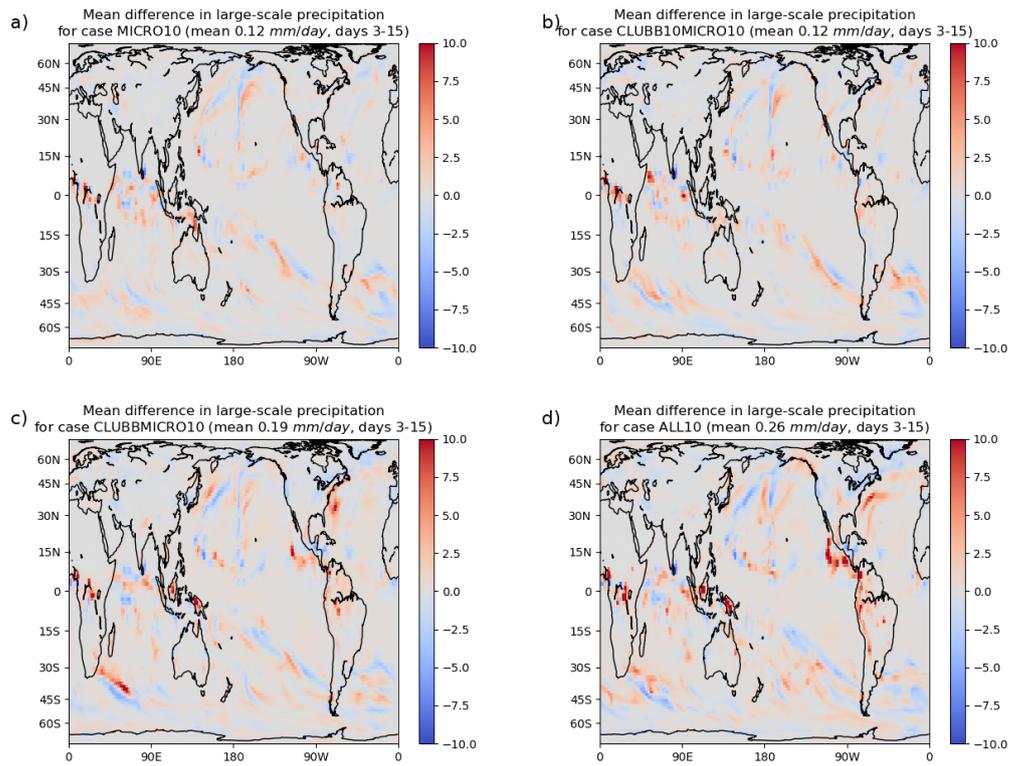


Figure 8. Differences in large-scale precipitation versus CTRL for a) MICRO10, b) CLUBB10MICRO10, c) CLUBBMICRO10, and d) ALL10.

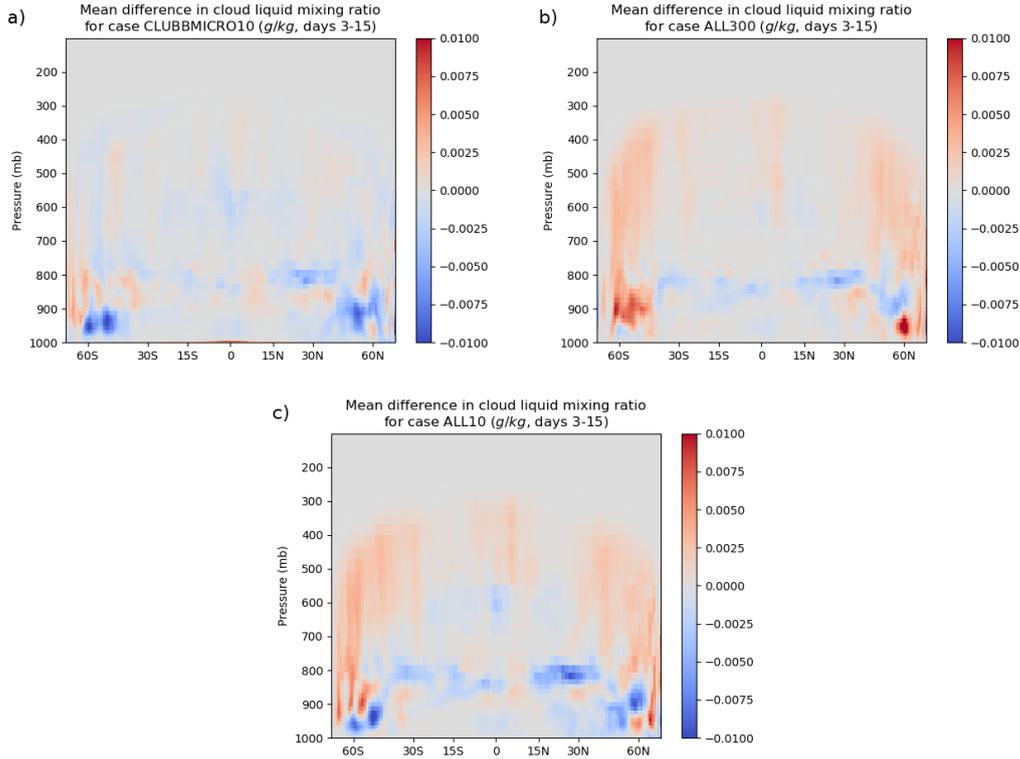


Figure 9. Differences in zonal mean cloud liquid mixing ratio versus CTRL for a) ALL300, b) CLUBBMICRO10, and c) ALL10.

253 Together, these results suggest that the main time steps affecting the precipitation
 254 are the MG2 microphysics time step and the coupling time step between the CLUBB
 255 and MG2 schemes. The change in the CLUBB and MG2 combined time step therefore
 256 explains most of the precipitation change noted earlier between the CTRL and ALL10
 257 runs. Either the dynamics-physics coupling time step or the ZM deep convection time
 258 step could also be affecting the precipitation between the convective and large-scale pro-
 259 cesses, since both these time steps are changed in the ALL10 and ALL300 runs. We will
 260 investigate this further in Section 4.3. The changes in total precipitation are mostly ap-
 261 parent for the first week of the run, after which the runs with default MG2 time step see
 262 a significant increase in convective precipitation (not shown), leading to no systematic
 263 difference between simulations after this point.

264 The particular pattern of decreased relative humidity found in the ALL10 run also
 265 appears in the CLUBBMICRO10 run, but not the MICRO10 run (not shown). However,
 266 as shown in Figure 9, the cloud liquid in the CLUBBMICRO10 run only matches the
 267 ALL10 run at low altitudes, while the ALL300 run is much more effective at matching
 268 ALL10 at higher altitudes, suggesting that the dynamics-physics coupling time step is
 269 responsible for these increases in cloud liquid. The CLUBBMICRO10 also consistently
 270 produces clouds in the lowest level of the atmosphere over land, particularly in South
 271 America and Southeast Asia, which are not seen in any other run. This may be an ar-
 272 tifact resulting from CLUBB and MG2 running at a time step much smaller than the
 273 atmosphere-land coupling interval.

274 The overall differences in ice and liquid water path are shown in Figure 10. We see
 275 that, in the tropics, the runs with reduced MG2 time step reduce the liquid water path

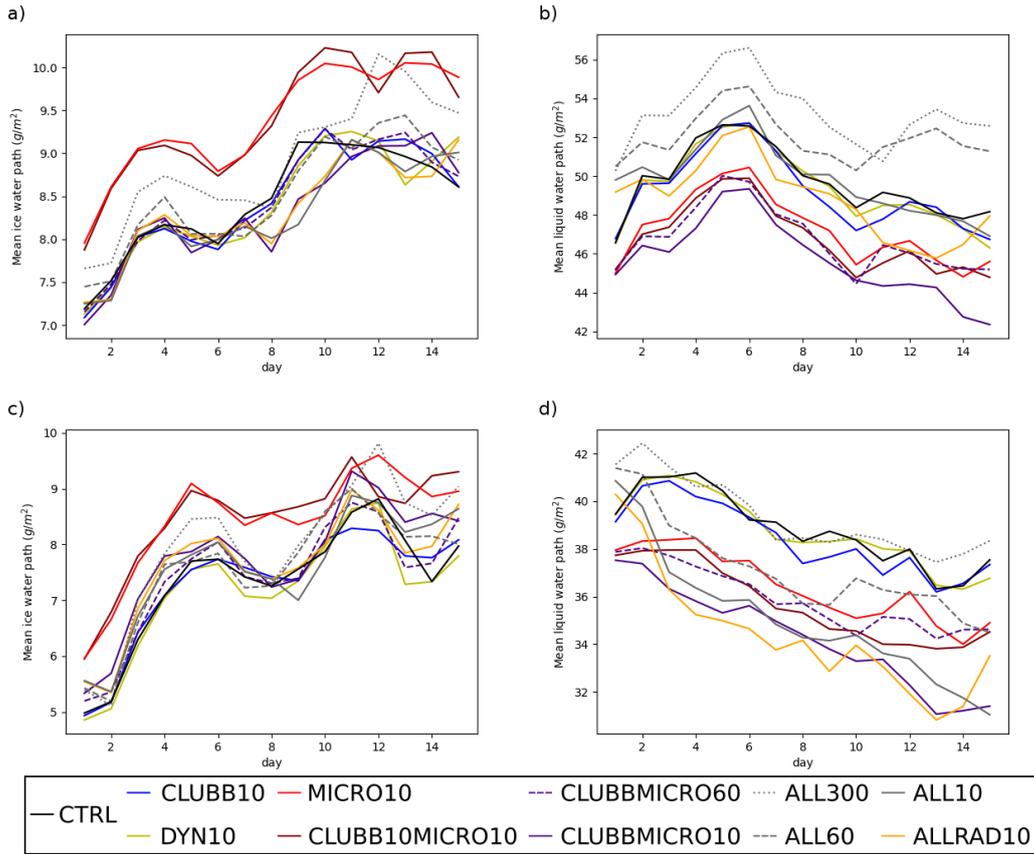


Figure 10. Daily means for a,c) ice water path and b,d) liquid water path for substepped runs. Plots a-b) show global means, while c-d) show means over low latitude grid points (30S–30N).

276 in a way that is similar to the ALL10 run, while the ALL300 run increases the ice and
 277 liquid water path everywhere. We also note that if MG2 is substepped independently
 278 from CLUBB, the ice water path increases significantly, but this does not occur when
 279 MG2 and CLUBB are substepped together.

280 We noted earlier that the ALL10 run caused a reduction in cloud fraction through-
 281 out most of the atmosphere, especially in the low cloud fraction. As shown in Figure 11,
 282 this effect seems to have different causes, depending on which level of the atmosphere
 283 is examined. Reductions in low cloud are primarily due to the reduction in the dynamics-
 284 physics coupling substep, but the CLUBB and MG2 combined time step has a greater
 285 effect on the cloud fraction above 700 mb.

286 Finally, we turn to the changes in radiative cloud forcing between runs. The mag-
 287 nitudes of both shortwave and longwave cloud forcing are reduced in the ALL10, ALL-
 288 RAD10, CLUBBMICRO10, and CLUBBMICRO60 runs, likely due to the significant de-
 289 creases in cloud fraction and liquid water path found in the tropics. The MICRO10 and
 290 CLUBB10MICRO10 runs, on the other hand, have a much larger ice water path, lead-
 291 ing to an increase in longwave cloud forcing. The ALL300 run has both a reduced low
 292 cloud fraction and an increase in liquid and ice water path, leading to a decrease in short-
 293 wave cloud forcing and no net change in longwave cloud forcing.

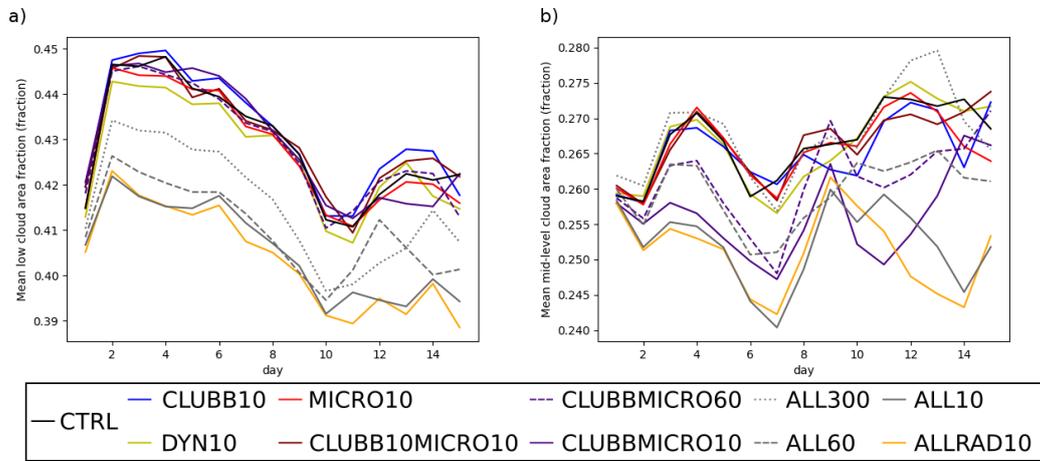


Figure 11. Daily global means of a) low cloud fraction (> 700 mb), and b) mid-level cloud fraction (400-700 mb) for substepped runs.

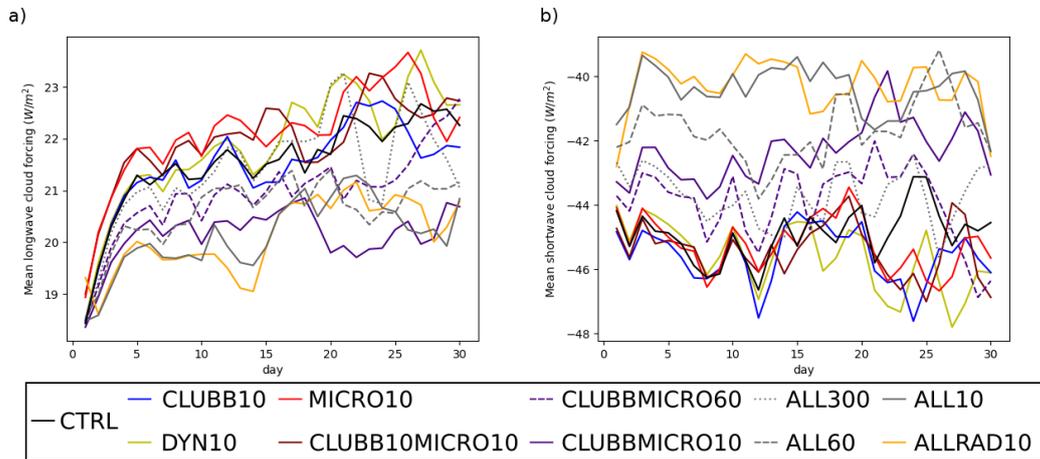


Figure 12. Daily global means of a) longwave cloud forcing and b) shortwave cloud forcing for substepped runs.

294 We notice that most variables take a few days for the differences between runs to
 295 fully develop, but the effect of a change in the MG2 time step strongly affects large-scale
 296 precipitation and ice water path within the first day. We suspect that most effects of a
 297 decreased time step require a certain degree of "spin up" in order for runs starting with
 298 the same initial condition to become more distinct. We hypothesize that the more in-
 299 stantaneous changes are primarily due to the direct effects of a decreased time step on
 300 microphysical process rates, which can respond directly and dramatically to changes in
 301 time step (Santos et al., 2020).

302 4.3 Substepping the ZM Deep Convection Scheme

303 So far, we have been unable to distinguish between the effect of substepping the
 304 ZM deep convection scheme and the effect of reducing the dynamics-physics coupling time
 305 step. In order to explore the effect of ZM substepping on results, we produced a simple
 306 set of code modifications to EAMv1 to allow this scheme to be substepped on its own.
 307 Most of these modifications are straightforward, since the main effect of ZM is simply
 308 to modify the state of the atmosphere for the next parameterization in the sequentially
 309 split physics. Two of ZM's outputs are precipitation process rates that are used by the
 310 modal aerosol scheme to calculate the total precipitation produced/evaporated by the
 311 deep convection. These rates are averaged over the whole time step in our modifications.

312 With this modified version of EAMv1, we were able to run the deep convection at
 313 a somewhat lower time step size, down to 300 s. However, the model becomes unstable
 314 if the ZM deep convection is a smaller time step (60 seconds or less) while the rest of the
 315 model runs at a default time step. Specifically, the wet deposition routines in the modal
 316 aerosols behave inappropriately, causing an unphysical increase in aerosol mass due to
 317 excessive water uptake, which in turn causes the aerosol optical depth to increase expo-
 318 nentially until the model crashes. Even in runs that did not crash, this behavior was present
 319 and had a significant impact on model physics. We were therefore unable to investigate
 320 the effect of ZM substepping further using the model's default configuration.

321 Fortunately, we do have an alternative, which is to run the model with prescribed
 322 aerosols, an ability commonly used for single column runs (Lebassi-Habtezion & Cald-
 323 well, 2015). This required switching to a configuration where prescribed aerosol data was
 324 available, so we used a year 2000 compset, FC5AV1C-04P2. As a result, these results
 325 cannot be directly compared to our previous runs, though we used the same spatial grid,
 326 and the physics of this compset is similar to our previous runs, aside from initial/boundary
 327 conditions. We reproduced the CTRL, ALL10, and CLUBBMICRO10 runs using pre-
 328 scribed aerosols, and further produced runs that substep ZM by itself, as well as runs
 329 that substep ZM and CLUBB/MG2 separately, and finally a run that substeps all three
 330 parameterizations. These simulations are summarized in Table 3.

331 First, we note that the CLD10PA run, which substeps the ZM deep convection along
 332 with CLUBB and MG2 in a single loop, has the same effect on the partitioning of pre-
 333 cipitation as seen in the ALL10PA run, being even closer to those results than the CLUBB-
 334 MICRO10PA run was. This difference can be seen in Figure 13, and suggests that the
 335 reduced ZM-CLUBB-MG2 coupling time step was responsible for the increased large-
 336 scale precipitation seen in the ALL300 and ALL10 runs in Figure 7. Figure 13 also shows
 337 that substepping ZM *by itself* has no effect on precipitation, since the ZM10PA run is
 338 similar to CTRLPA and the CLUBBMICRO10ZM10PA run is similar to CLUBBMICRO10PA.
 339 Williamson (2013) shows that the ZM scheme is not very active when coupled to other
 340 parameterizations at small time step sizes, when using a typical value of the convective
 341 relaxation time-scale (on the order of 1 hour, which is also the value in our experiments).
 342 We see the same shift from deep convection towards stratiform precipitation at short time
 343 steps.

Name	Substepped processes	Substep size	Run length
CTRLPA	None	N/A	30 d
ALL10PA	Dynamics-physics coupling	10 s	30 d
CLUBBMICRO10PA	CLUBB+MG2 combined loop	10 s	30 d
ZM10PA	ZM deep convection	10 s	30 d
CLUBBMICRO10ZM10PA	CLUBB+MG2 combined loop and ZM deep convection	10 s	30 d
CLD10PA	CLUBB+MG2+ZM combined loop	10 s	30 d

Table 3. Runs performed using prescribed aerosols

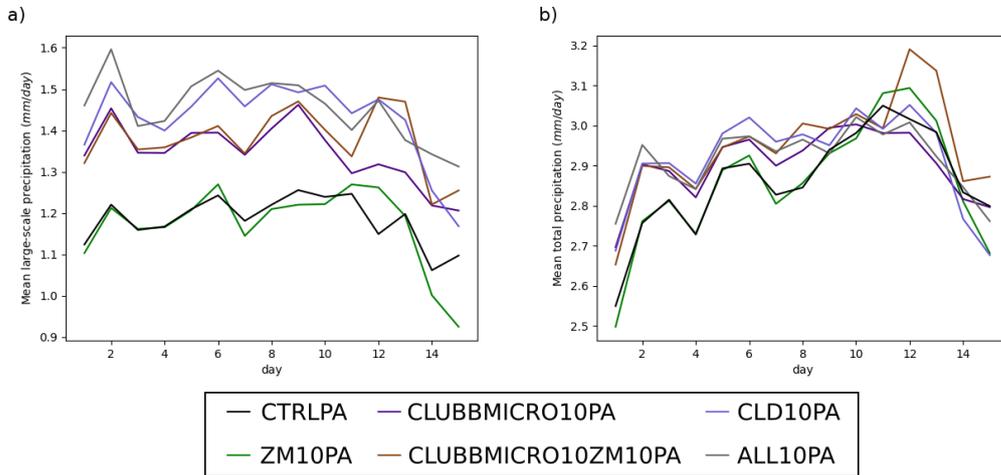


Figure 13. Daily global means of a) large-scale precipitation only and b) total precipitation for prescribed aerosol substepped runs.

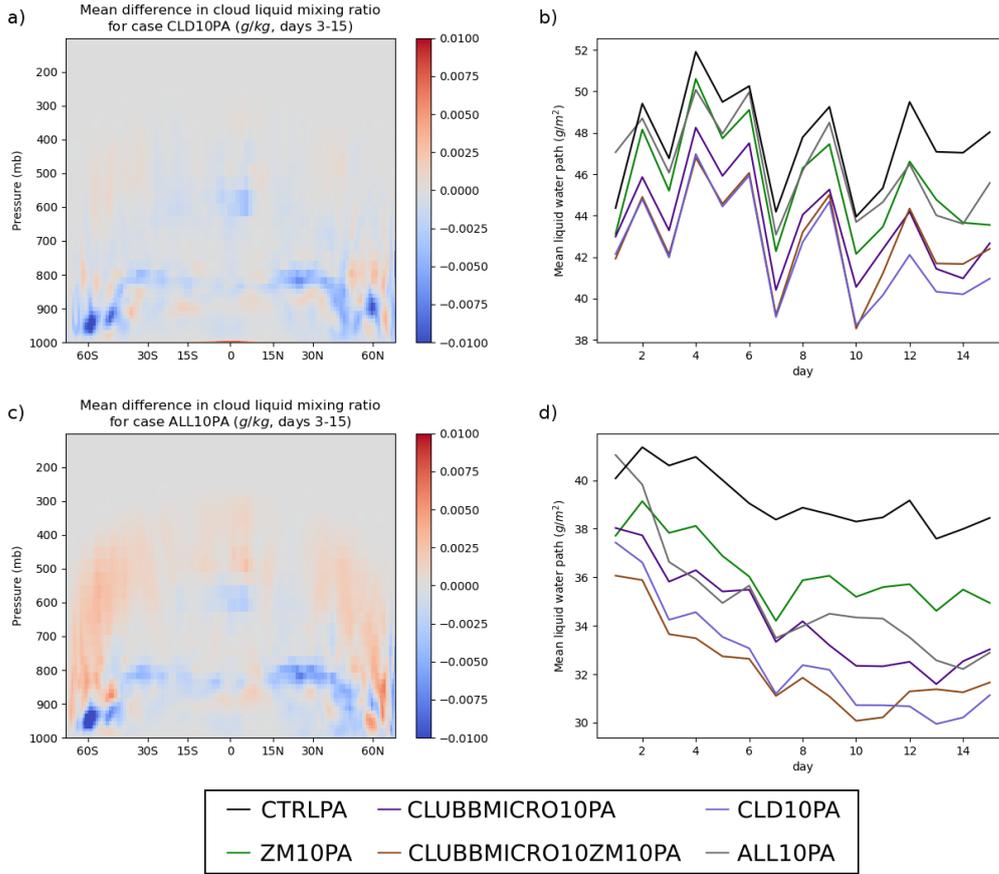


Figure 14. Left: Differences in zonal mean cloud liquid mixing ratio versus CTRLPA for a) CLD10PA and c) ALL10PA. Right: Daily liquid water path for prescribed aerosol substepped runs using b) a global mean, and d) a mean over low latitude grid points (30S–30N).

344 Second, we note that deep convection substepping, like substepping of the other
 345 parameterizations, causes a decrease in low cloud liquid mass, and contributes to the over-
 346 all pattern seen in the ALL10PA run. This is seen in Figure 14, where the distribution
 347 of liquid water below 750 mb agrees quite well between the CLD10PA run and the ALL10PA
 348 run. However, the increase in cloud liquid above this level is still absent from the CLD10PA
 349 run, implying that that increase requires more frequent dynamics-physics coupling to oc-
 350 cur. This means that the CLD10PA ”overshoots” the ALL10PA run in the tropics, hav-
 351 ing an even lower liquid water path. Unlike the effect of ZM substepping on precipita-
 352 tion, the effect on cloud liquid does not rely on coupling with CLUBB and MG2, since
 353 the CLUBBMICRO10ZM10PA run (not shown) and the CLD10PA run have fairly simi-
 354 lar liquid water path.

355 Other than this, ZM substepping accounts for very little of the differences between
 356 the CTRLPA and ALL10PA runs. There are small decreases in cloud fraction (not shown),
 357 but these are much weaker than the effect of increased dynamics-physics coupling fre-
 358 quency below 700 mb, and weaker than the effect of a smaller CLUBB and MG2 time
 359 step above 700 mb. The effect on shortwave cloud forcing, as seen in Figure 15, is thus
 360 also relatively small compared with the effect of changing the CLUBB and MG2 time
 361 step. The effect of ZM substepping on longwave cloud forcing may appear to be more
 362 significant, since the CLUBBMICRO10ZM10PA run has almost the same longwave cloud

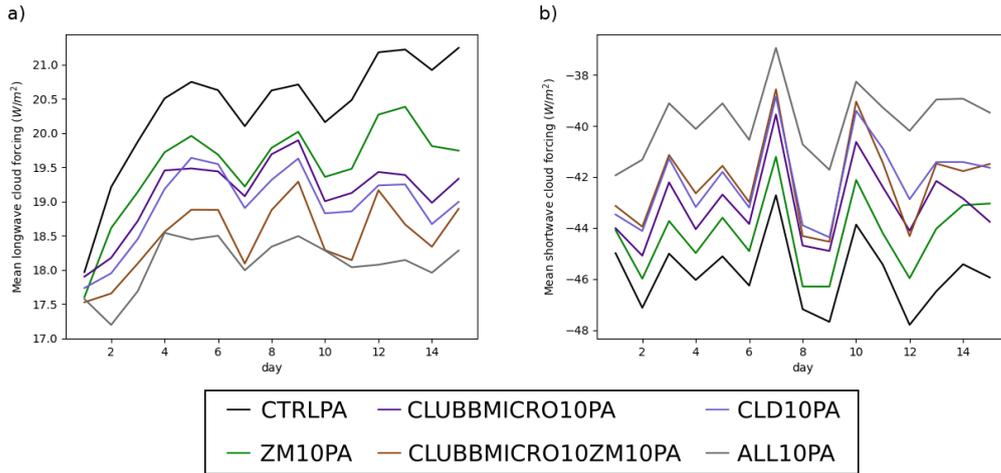


Figure 15. Daily global means of a) longwave cloud forcing and b) shortwave cloud forcing for prescribed aerosol substepped runs.

363 forcing as the ALL10PA run. However, the mechanism here is completely different; all
 364 runs where ZM is substepped, including CLUBBMICRO10ZM10PA, have a dramatically
 365 lowered ice water path (not shown). ALL10PA and CLUBBMICRO10PA, on the other
 366 hand, have an ice water path similar to the control, and so the decrease in longwave cloud
 367 forcing is instead attributable to decreased cloud fraction in these cases.

368 5 Conclusions

369 EAMv1 at its default 1800 second time step produces very different results from
 370 the same model at a 10 second time step, indicating that the release implementation should
 371 not be viewed as calculating the “time-resolved” solution to the system of equations that
 372 defines the model physics. The amount and regional distribution of precipitation, and
 373 especially the radiatively-important partitioning between large-scale and convective pre-
 374 cipitation, shows particularly strong sensitivity to the time step size. The cloud radi-
 375 ative forcings also differ by several watts per square meter, indicating that a reduction
 376 in the time step would require, at a minimum, significant retuning of the model to pro-
 377 duce reasonable results.

378 By experimenting with substepping of model components, we have been able to dis-
 379 tinguish three main model time steps that account for most of the changes seen between
 380 the 1800s and 10s versions of the model. We briefly summarize these findings here.

381 First, a reduction of the combined CLUBB and MG2 time step causes the follow-
 382 ing changes:

- 383 1. An increase in total precipitation, leading to a reduction in humidity and a reduc-
 384 tion in cloud liquid mass below 750 mb.
- 385 2. A reduction in the ratio of convective to large-scale precipitation.
- 386 3. Regional changes in precipitation, most notably including a massive increase in
 387 average precipitation on the maritime continent and in South America.
- 388 4. A reduction in cloud fraction above 700 mb, causing a large reduction in the mag-
 389 nitudes of both shortwave and longwave radiative cloud forcing.

390 The first two of these changes can be produced by changing the time step for MG2
 391 alone, but substepping CLUBB and MG2 together produces a significantly larger effect.

392 Second, a reduction of the dynamics-physics coupling time step (a.k.a. *dtime*) causes
 393 the following changes:

- 394 1. An increase in cloud mass in the upper troposphere, especially in the midlatitudes.
- 395 2. A substantial decrease in cloud fraction below 700 mb, causing further large de-
 396 creases in radiative cloud forcing, especially for shortwave radiation.

397 Third, a reduction of the ZM time step causes the following changes:

- 398 1. A further reduction in cloud liquid mass below 750 mb, though this corresponds
 399 to a reduction in net condensation minus evaporation, not an increase in precip-
 400 itation.
- 401 2. A small decrease in cloud fraction everywhere, causing further small decreases in
 402 radiative forcings.
- 403 3. A further reduction in the ratio of convective to large-scale precipitation. (How-
 404 ever, this only occurs when ZM is coupled more frequently with CLUBB and MG2,
 405 which in the original code can only be done by reducing *dtime*.)

406 These observations indicate that even when individual parameterizations seem to
 407 be well resolved in time, a low coupling frequency between parameterizations can still
 408 be a significant source of model error. This may be an underappreciated issue, since de-
 409 velopers of new parameterizations tend to focus on the time step of their own particu-
 410 lar parameterization rather than the frequency with which it is coupled to other parts
 411 of a model. For our case, while the MG2 and ZM parameterizations are each mildly sen-
 412 sitive to changes in time step, the majority of the time step sensitivity in E3SM is ac-
 413 tually due to sensitivity to the coupling frequencies between different processes.

414 In EAMv1, most calculations in the dynamics and many physics parameterizations
 415 are already running at a five minute time step or less, even for lower resolution runs. Donahue
 416 and Caldwell (2020) found that for a 1° simulation, halving the dynamics-physics cou-
 417 pling frequency only increased model cost by 20%. Our ALL300 run had one-sixth the
 418 *dtime* of the CTRL run, but only required 66% more core-hours per simulated year. This
 419 raises the question: why not use a shorter time step for the dynamics-physics coupling
 420 for all runs? In the short term, this would likely require some significant retuning of ex-
 421 isting models. In the long run, however, our results suggest that significant improvements
 422 could be attained by reducing the dynamics-physics coupling time step for future model
 423 development.

424 While simply reducing the model time step might be a reasonable way of dealing
 425 with E3SM's sensitivity to the dynamics-physics coupling frequency, the sensitivity to
 426 CLUBB+MG2 coupling is more difficult to address. CLUBB and MG2 together account
 427 for a large share of the model cost, so reducing their time steps by a factor of 30 could
 428 be around an order of magnitude slower than the default configuration. Since most mod-
 429 elers will not be able to accept such a large increase in computational cost, we can sug-
 430 gest a few other ways of working around this cost in future model development:

- 431 1. Reduce the cost of simulating the most expensive physical processes, e.g. by switch-
 432 ing to simpler implementations of these processes. For instance, a model that uses
 433 a less accurate approximation for some process, but is able to run at a higher tem-
 434 poral or spatial resolution, may end up being more accurate than a model that
 435 uses a more accurate set of equations.
- 436 2. Use alternative (e.g. higher-order) time integration schemes to lower the time in-
 437 tegration error at moderate time step sizes.

438 3. Redesign the physics to separate out processes that have a shorter or longer time
 439 scale. In the case of E3SM, this would mean refactoring CLUBB and MG2 (or any
 440 future set of schemes that cause similar issues), in order to isolate the parts of those
 441 parameterizations that are most responsible for the time step sensitivity of the over-
 442 all model. If this subset of physical processes can be calculated much more cheaply
 443 than the total cost of CLUBB and MG2, it could then be handled with a more
 444 accurate time integration scheme without incurring an excessive cost.

445 In practice, errors coming from a coarse temporal resolution are often handled by
 446 simply tuning GCMs so that they match observations when run with longer time step
 447 sizes, just as models are tuned for a particular horizontal and vertical resolution. While
 448 this is an effective approach for many studies to match current climate observations, if
 449 a model relies heavily on tuning to cancel large numerical errors, it is unlikely to have
 450 the correct sensitivity to forcing changes. This is especially a concern for studies that
 451 use a model to simulate conditions very different from those originally used to tune that
 452 model (e.g. for paleoclimate).

453 We recommend that GCM developers continue to study time step sensitivity by
 454 running experiments with full model physics. Even when the effect of time step on each
 455 individual parameterization is well known, the effect of process coupling can affect model
 456 behavior in unpredictable ways. Most users cannot afford to use sub-minute time steps
 457 for the entire model, and therefore are likely using a model that is not achieving its own
 458 converged small-time-step behavior. Therefore it is important to understand the limi-
 459 tations and biases present in workhorse models that have been tuned for coarse time step
 460 sizes.

461 **Acknowledgments**

462 The data used in this paper will be made available using Argonne National Laboratory's
 463 Petrel service (data still being uploaded, URL to be added). Until the upload is com-
 464 plete, the data can be found in an xz-compressed tar file at:

465 <https://drive.google.com/file/d/1yENgyLnaUNBGhEMMx0u6vbiVk51cqDMC/view>

466 This tar file contains data files from all runs used in this paper, labeled using the
 467 following case names:

- 468 1. CTRL - timestep_ctrl
- 469 2. ALL10 - timestep_all_10s
- 470 3. ALL60 - timestep_all_60s
- 471 4. ALL300 - timestep_all_300s
- 472 5. DYN10 - timestep_dyn_10s
- 473 6. MICRO10 - timestep_MG2_10s
- 474 7. CLUBB10 - timestep_CLUBB_10s
- 475 8. CLUBB10MICRO10 - timestep_CLUBB_10s_MG2_10s
- 476 9. CLUBBMICRO10 - timestep_CLUBB_MG2_10s
- 477 10. CLUBBMICRO60 - timestep_CLUBB_MG2_60s
- 478 11. ALLRAD10 - timestep_all_rad_10s
- 479 12. CTRLPA - timestep_presaer_ctrl
- 480 13. ALL10PA - timestep_presaer_all_10s
- 481 14. CLUBBMICRO10PA - timestep_presaer_CLUBB_MG2_10s
- 482 15. ZM10PA - timestep_presaer_ZM_10s
- 483 16. CLUBBMICRO10ZM10PA - timestep_presaer_CLUBB_MG2_10s_ZM_10s
- 484 17. CLD10PA - timestep_presaer_cld_10s

485 The ne16 grid data used for Figure 6 is labeled with the case name timestep_ctrl_ne16.
486 Data files are placed into the following directory:

- 487 1. timestep_monthly_avgs: Monthly average data sets for the CTRL and ALL10 runs.
488 Data from the ne16 grid run is also in this directory, both on the original grid (h0
489 files) and regridded to ne30 resolution.
- 490 2. timestep_monthly_avgs_lat_lon: Monthly data regridded to a rectangular latitude-
491 longitude grid with 1.4-degree resolution (used for plotting purposes).
- 492 3. timestep_daily_avgs: Daily average data sets for all runs (excepting the ne16 grid
493 run) for the first 30 days of simulations.
- 494 4. timestep_daily_avgs_lat_lon: Daily data regridded to a rectangular latitude-longitude
495 grid with 1.4-degree resolution (used for plotting purposes).

496 The code used to process the data in this study is available at:

497 https://github.com/quanttheory/E3SMTimestepStudy/tree/Santos_JAMES_2020

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