

Multi-layer Seismic Anisotropy Beneath Greenland

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Key Points:

- 299 new *KS shear wave splitting measurements made using broadband stations across Greenland
- Variations of fast direction with back-azimuth can be explained by two layers of anisotropy which are consistent across Greenland
- Lower layer consistent with asthenospheric shear; upper layer consistent with past lithospheric orogenic deformation

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Abstract

Seismic anisotropy provides insight into past episodes of lithospheric deformation and the orientations of strain in the underlying asthenosphere. The Greenland mantle has played host to a rich history of tectonic processes, including multiple orogenies and plume-lithosphere interactions. This study presents new measurements of SKS splitting that reveal strong variations in fast polarization direction with back-azimuth that are consistent across Greenland, including at stations where splitting measurements have not previously been reported. We compared observed fast polarization directions to the predictions of two-layer models with olivine-orthopyroxene anisotropy. The family of models which provides acceptable misfits at 95% confidence indicates an upper layer olivine a-axis azimuth of $226 \pm 2.9^\circ$ and a lower layer olivine a-axis azimuth of $124 \pm 2.7^\circ$ and non-zero axis dips are required. These models are consistent with asthenospheric anisotropy aligned approximately parallel to the direction of plate motion and lithospheric anisotropy due to Proterozoic and Archean orogenic fabrics.

Plain Language Summary

Measurements of seismic anisotropy (the direction-dependent variation of seismic wavespeed) provide useful information about the orientation of deformation in the Earth. We measured seismic anisotropy using shear waves refracted through the outer core and recorded by stations in Greenland. Due to new stations and data, this study includes more measurements of the effects of anisotropy than previously possible. We show that a model with two layers of anisotropy explains dominant patterns in the fast vibration direction of the shear waves as a function of the angle at which they approach each station. We suggest that the lower layer reflects deformation in the asthenospheric mantle induced by the motion of the plate above, and the shallow layer reflects coherent deformation in the continental lithosphere of Greenland due to its history of plate collisions.

1 Introduction

Nearly all of Greenland's surface geology is inaccessible because it is covered by the Greenland Ice Sheet. Therefore, geophysical investigations are especially important in furthering our understanding of Greenland's subglacial lithospheric structure. Greenland is a region of interest as its lithosphere contains cratonic material and records the his-

43 tory of Archean, Proterozoic, and Paleozoic orogenies, and also could provide insight into
44 the history of the Iceland plume (e.g Henriksen et al., 2009).

45 The majority of Greenland is Precambrian and has been modified by multiple tec-
46 tonic (orogenic and rifting) events (e.g Henriksen et al., 2009). Of particular note is the
47 Trans-Hudson Orogeny which was a widespread set of plate collisions that helped to build
48 Laurentia around 1.8 Ga (e.g St-Onge et al., 2009). Orogenic belts from this event can
49 be found across North America; in Greenland this includes the Rinkian and Nagssug-
50 toqidian belts that bound major crustal blocks (e.g Antonijevic & Lees, 2018; Dahl-Jensen
51 et al., 2003; Henriksen et al., 2009). During the Silurian, the continent-continent colli-
52 sion of Laurentia and Baltica developed the Eastern Greenland Caledonides, resulting
53 in complex thrust architecture along the Eastern coast (e.g Dawes, 2009; Higgins & Leslie,
54 2000). More recently, Greenland has been modified volcanically and thermally by the
55 passage of the Iceland plume underneath Greenland between 70-40 Ma (e.g Lawver &
56 Muller, 1994). However, studies differ regarding the exact path of the plume under Green-
57 land and its effects on the overlying lithosphere (Braun et al., 2007; Forsyth et al., 1986;
58 Lawver & Muller, 1994; Steffen et al., 2018; Steinberger et al., 2019). Regional-scale seis-
59 mic imaging has helped elucidate this tectonic history, constrain the temperature and
60 composition of the Greenland mantle, and interrogate plume-lithosphere interactions.
61 Surface-wave tomography (Darbyshire et al., 2018; Lebedev et al., 2018; Levshin et al.,
62 2001) has detected thick depleted cratonic mantle lithosphere and lithospheric structures
63 modified by multiple tectonic events. H-K stacking and synthetic modeling of receiver
64 functions have helped to constrain regional crustal thickness and composition (Dahl-Jensen
65 et al., 2003; Kumar et al., 2005, 2007). Images from body-wave tomography (Toyokuni
66 et al., 2020) show a NW-SE low-velocity anomaly within the mantle, coincident with heat
67 flow anomalies interpreted as evidence for plate movement over the Iceland plume. Sur-
68 face wave tomography (Levshin et al., 2017; Lebedev et al., 2018; Mordret, 2018; Pour-
69 point et al., 2018) has also been used to identify this thermal signature. In the mantle,
70 a common source of anisotropy is the lattice preferred orientation of minerals such as
71 olivine and orthopyroxene; in conditions where the mantle is relatively dry and/or low
72 stress, shear wave splitting fast polarization directions are thought to align approximately
73 parallel to the direction of horizontal flow (Karato et al., 2008; Long & Becker, 2010).
74 Shape preferred orientation of velocity heterogeneity can also cause anisotropy (Holtzman
75 et al., 2003).

76 Relatively few studies of seismic anisotropy exist for the mantle beneath Green-
77 land. Azimuthal anisotropy in global-scale images (Ekström, 2011; Schaeffer et al., 2016)
78 is difficult to interpret due to the coarse parameterization of these models, which pro-
79 vide only a small number of data points in Greenland. A regional-scale surface-wave study
80 (Darbyshire et al., 2018) shows weak anisotropy at shallow mantle depths with a NW-
81 SE fast direction beneath the central latitudes of Greenland and NE-SW fast directions
82 in the far north and south, but only provides constraints in the uppermost mantle. Pre-
83 vious shear wave splitting measurements are predominantly N-NE in southern Green-
84 land, and more variable elsewhere (e.g. Ucisik et al., 2008). A lateral gradient in anisotropy
85 near the southern coast of Greenland has also been measured with quasi-love waves (Servali
86 et al., 2020).

87 Shear-wave splitting arises when anisotropic media polarize shear wave particle mo-
88 tions that travel at different velocities, and the polarization direction of the fast shear
89 wave (Φ) and the time delay (δt) between the two split waves measured at the receiver
90 are commonly used to characterize the anisotropy. The presence of multiple layers of anisotropy,
91 with different a-axis azimuth, a-axis plunge, and/or strength result in back-azimuthal
92 variation of the measured splitting parameters. When detected, back-azimuthal varia-
93 tions of apparent splitting parameters are a useful tool for measuring the variation of
94 anisotropy with depth (e.g. Savage & Silver, 1993; Silver & Savage, 1994; Levin et al.,
95 1999). A range of approaches have been applied to this problem, including exploration
96 of the large model space using the neighbourhood algorithm (e.g. Wookey, 2012; Yuan
97 & Levin, 2014). Grid searches through model parameter space have also been used to
98 constrain a-axis azimuth, plunge, and anisotropy strength (Abt et al., 2010). Forward
99 modeling that parameterizes the anisotropy in each layer with a fast polarization direc-
100 tion and splitting time is also sometimes employed, reducing the parameter space (e.g.
101 Aragon et al., 2017; Hammond et al., 2014; Wookey, 2012) In a limited number of cases,
102 tomographic approaches have been applied to shear-wave splitting from local (Abt & Fis-
103 cher, 2008; Abt et al., 2009; Calixto et al., 2014) and teleseismic events (Long et al., 2008;
104 Mondal & Long, 2020).

105 Although prior studies have measured shear wave splitting in Greenland (Clement
106 et al., 1994; Helffrich et al., 1994; Vinnik et al., 1992; Ucisik et al., 2005, 2008), clear vari-
107 ations in splitting parameters with back-azimuth diagnostic of multiple layers of anisotropy
108 have not been resolved (e.g. Ucisik et al., 2008). In this study, we measure shear wave

109 splitting across Greenland using decades of new data and dozens of new stations, observe
110 systematic variations in fast polarization direction with back-azimuth, and model these
111 patterns with two-layer anisotropy.

112 **2 Data and Methods**

113 **2.1 Data**

114 We measured shear-wave splitting fast polarizations and delay times from *KS phases
115 using a new dataset collected from 27 stations deployed on the Greenland ice sheet and
116 coast, as well as stations on Ellesmere Island (ALE) and on Jan Mayen Island (JMIC)
117 (Table S1). Seismic data used in this analysis were acquired at broadband stations, which
118 were deployed for different periods of time (Table S1), ranging from four months (for sta-
119 tions part of temporary deployments) to nearly 30 years. Stations include those from net-
120 work codes DK (the Danish Seismological Network), GE (GEOFON), XF (GLISN), G
121 (GEOSCOPE), CN (Canadian National Seismograph Network) and II (the IRIS/USGS
122 Global Seismographic Network) (Table S1). Station spacing varies dramatically, from
123 more than 200 km on the ice sheet to less than 50 km on the coast, with stations mostly
124 distributed along the coast. We employed BH* channels sampled at 100 Hz. We selected
125 earthquakes of magnitude greater than 6.0, between epicentral distances of 90° and 130°
126 from each station.

127 **2.2 Measurement Methods**

128 To measure shear-wave splitting, we employed the SplitLab software (Wüstefeld
129 et al., 2008). We filtered waveforms between 0.01 and 0.1 Hz before manually inspect-
130 ing and windowing data around the *KS phase. We report results of the transverse-component
131 minimization method (Silver & Chan, 1991), although we only retained measurements
132 whose uncertainties overlapped those from eigenvalue minimization. We imposed sev-
133 eral other criteria to distinguish a measurement as high quality (e.g. Fig. 1): the *KS
134 phase is a clear arrival with a signal-to-noise ratio ≥ 2 on the Q component; *KS phases
135 (isolated or coincident) behave like a split *KS phase, i.e. the energy on the transverse
136 component decreases and elliptical particle motion becomes linear when splitting is re-
137 moved from the waveforms; the 95% error surfaces for the transverse-component min-
138 imization method and eigenvalue minimization method are close to an ellipse; the un-

Splitting Measurement

Event: **08-Jul-2008 (190) 09:13 -15.99N -71.75E 123km Mw=6.2**
 Station: **SUMG** Backazimuth: **211.9°** Distance: **91.08°** SNR_{sc}: **16.0**

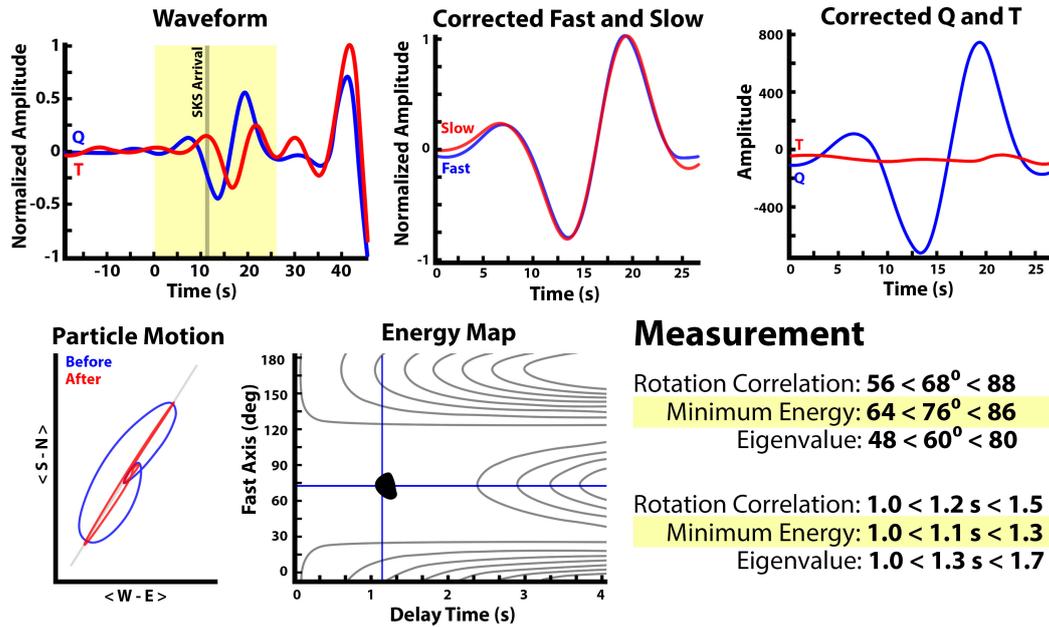


Figure 1. Example of a high quality splitting measurement obtained using SplitLab (Wüstefeld et al., 2008). Upper left: Q and T components of initial waveforms. An SKS phase is highlighted in the yellow window. Upper center: SKS phase waveform on fast and slow polarization components, shifted to remove the splitting lag time. Upper right: The shifted waveform components on the Q and T components. Lower left: Horizontal components of the SKS phase before (blue) and after (red) the splitting lag time was removed. Lower center: Surface of energy on the T component as a function of trial splitting fast direction and delay time. Splitting parameters within 95% confidence of the best-fitting values lie within the shaded contour. Lower right: fast direction and splitting time (with uncertainties) for each of the measurement methods.

139 certainty range in splitting time does not overlap zero nor does it exceed 4 seconds; the
 140 uncertainty in fast direction is less than $\pm 30^\circ$ for the transverse-component minimiza-
 141 tion method.

3 Results

3.1 Shear-wave Splitting Results

We measured a total of 299 high quality shear wave splitting measurements (Table S2). At many stations, there is significant variation in the measured fast directions with back-azimuth (Fig. 2, Fig. S1). In particular, the stations with the largest number of measurements (e.g. ALE, NEEM, SCO) all show clear variation in fast direction with back-azimuth. For example, at station NEEM (Fig. 3a), fast directions span 27° – 171° . At other stations, the fast directions are clustered around a single value (e.g. NE6, TULEG, KULLO, DAG, ISOG). However, at some stations the distribution of measurements with back-azimuth is insufficient to determine whether back-azimuthal variation in fast direction exists. In addition, there is little geographic coherence in mean fast direction between stations (Fig. 2, Fig. S1).

To examine whether measurements at individual stations can be fit by a single layer of anisotropy, we determined the single horizontal olivine a-axis orientation whose predicted shear-wave splitting fast directions produce the minimum summed circular misfit when compared to the observed fast directions at the station. Our measurements fall into three back-azimuth ranges of width 120° . If the maximum misfit (with respect to the best-fitting a-axis orientation) to a single observation in any cluster is greater than 30° , we deem the fast directions as not fit by a single layer. Stations are deemed ambiguous if the maximum circular misfit is less than 30° for all observed fast directions, but data do not exist in all three back-azimuth bins, which could result in under-sampling of the predictions of an underlying two-layer anisotropy pattern. Using this definition, only stations NE6 and ISOG are consistent with a single layer of anisotropy (Fig. 2).

In addition, when all fast polarizations are plotted together (Fig. 3b), their overall pattern of fast direction variation with back-azimuth is broadly consistent, including stations which can and cannot be fit by a single layer of anisotropy (Fig. 2, Fig. S1). This broad pattern of fast direction variation in back-azimuth (Fig. 3b) persists in regional sub-groups of stations, for example those north and south of 70° N (Fig. 3b).

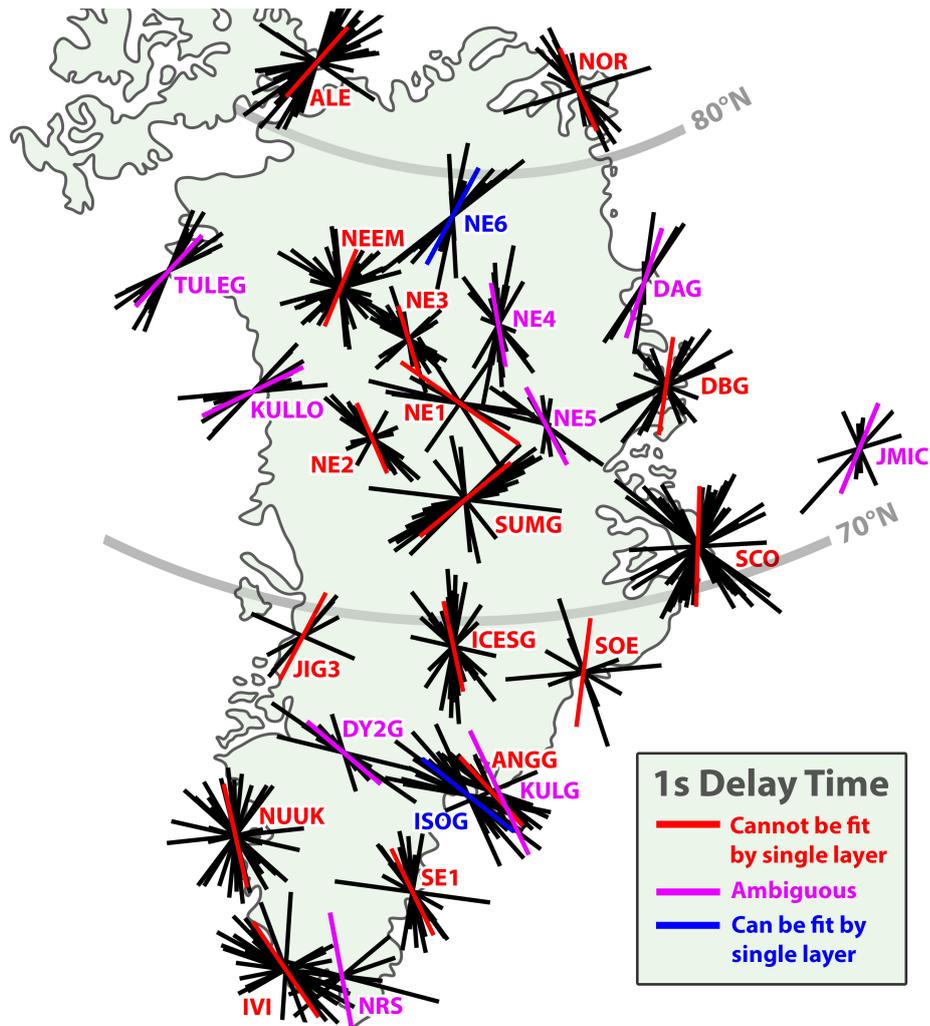


Figure 2. Map of individual splitting measurements (black lines) and average fast direction (colored lines) measured at each station. Color shows classification of whether the fast directions at the station can or cannot be fit by a single layer of anisotropy. The length of the lines corresponds to splitting time; the time scale appears in the legend.

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3.2 Modeling Two Layers of Anisotropy

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To constrain the variation of anisotropy with depth implied by the observed variation of shear-wave splitting fast directions with back-azimuth, we compare observed fast directions to the predictions of two-layer anisotropy models. We assume an isotropic crust of thickness 40 km (Darbyshire et al., 2018), an anisotropic mantle lithosphere between 40 km and 150 km, and an anisotropic asthenosphere between 150 and 300 km. The mantle lithosphere thickness is based on thermally defined lithospheric thickness values for

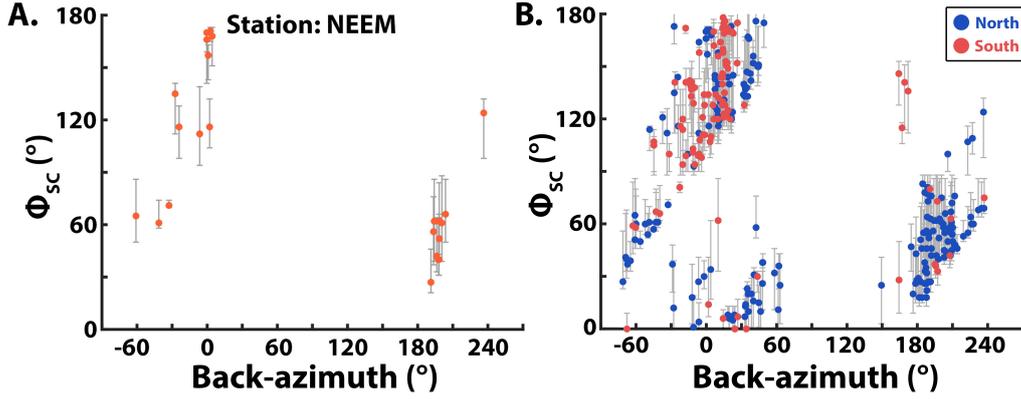


Figure 3. a) Fast directions versus back-azimuth at station NEEM. Significant variation in fast direction occurs which cannot be explained by a single layer of anisotropy. Plots for all other stations can be found in Fig. S1. b) All measurements of fast direction as a function of back-azimuth for our study area separated into northern ($< 70^\circ\text{N}$) and southern ($< 70^\circ\text{N}$) groups based on station latitude. $N=299$.

177 Greenland (Steinberger & Becker, 2018). To define the tensor of anisotropic elastic co-
 178 efficients, we assumed a mantle composed of 70% olivine and 30% orthopyroxene. The
 179 model has six free parameters: olivine a-axis azimuth (θ) and plunge (δ), and anisotropy
 180 strength (α) in each of the two layers. Anisotropy strength is defined as the percentage
 181 of total single crystal anisotropy. In other words, 100% anisotropy would be the elastic
 182 coefficients for pure olivine and orthopyroxene, aligned with respect to each other so that
 183 the a-axis of olivine is parallel to the c-axis of orthopyroxene, and the b-axis of olivine
 184 is parallel to the a-axis of orthopyroxene (Mainprice & Silver, 1993)

185 We predict shear-wave splitting parameters for each back-azimuth in the observed
 186 splitting dataset using the approximate particle motion perturbational method (Fischer
 187 et al., 2000); this approach has been shown to match results generated using pseudospec-
 188 tral synthetics (Hung & Forsyth, 1998). The code rotates and time-shifts an initial lin-
 189 ear wavelet of period 10 s using the Christoffel matrix for the anisotropy in the lower layer,
 190 and then rotates and time-shifts the resulting particle motion for the anisotropy in the
 191 upper layer. Shear-wave splitting parameters are then measured from the synthetic wave-
 192 form using the eigenvalue minimization method (Silver & Chan, 1991), which for the noise-
 193 free synthetics used in the modeling yields identical results to the transverse energy min-
 194 imization method that was applied to the data.

195 Because of the non-linear relationship of anisotropy model parameters to shear-wave
 196 splitting predictions, we employ a grid-search approach to determine the best-fitting model
 197 parameters. Due to the large number of model parameter combinations, we first com-
 198 pare the observed fast directions to predicted fast directions from a more coarsely sam-
 199 pled grid of parameters, and then implement a finer grid search around the minima re-
 200 solved from the coarser grid. In the coarse grid search, a-axis azimuth varies in incre-
 201 ments of 10° between 0° - 360° from north, a-axis plunge varies in increments of 10° be-
 202 tween 0° - 50° from horizontal, and the strength of anisotropy varies in increments of 10%
 203 between 0%-50%. In the finer grid search, we probe a-axis azimuths in increments of 2°
 204 in a range $\pm 20^\circ$ away from the best-fitting value from the coarse grid search, and probe
 205 dip and strength along the same spacing as in the coarse grid search.

206 Using the measurements of fast directions from all stations (Fig. 3b, Table S2), the
 207 coarse grid search yields a global minimum RMS misfit of 2.81 at $\theta_{deep} = 130^\circ$, $\delta_{deep} =$
 208 40° , $\alpha_{deep} = 50\%$; $\theta_{shallow} = 230^\circ$, $\delta_{shallow} = 30^\circ$, $\alpha_{shallow} = 40\%$ (Fig. 4). There
 209 are other local minima, but these do not minimize misfit. The finer grid search finds a
 210 better-fitting model with a misfit of approximately 2.26 (Fig. 4). The best-fitting pa-
 211 rameters for the finer grid search are $\theta_{deep} = 124^\circ$, $\delta_{deep} = 50^\circ$, $\alpha_{deep} = 50\%$; $\theta_{shallow} =$
 212 226° , $\delta_{shallow} = 20^\circ$, $\alpha_{shallow} = 40\%$. We use an F-test (Snedecor & Cochran, 1991)
 213 to determine the family of models that fit the observations within the 95% confidence
 214 limits of the best-fitting model.

215 From the fine grid search, we find a total of 39 models which satisfy the 95% con-
 216 fidence interval constraint, and adequately predict the large-scale variation of fast axis
 217 with back-azimuth (Fig. 5, left). As the width of the parameter histograms (Fig. 5, right)
 218 indicate, our grid search places robust constraints on the a-axis azimuths in the upper
 219 and lower layers. Acceptable a-axis azimuths in the upper layer vary from 222° to 232° ,
 220 and in the lower layer from 120° to 130° . A-axis plunge in the lower layer is at 50° , while
 221 acceptable values of a-axis plunge in the upper layer range from 10° to 30° . Among the
 222 parameters we probe in our grid search, the strength of anisotropy in each layer is the
 223 least well constrained. Unlike dip or layer a-axis orientation, it does not result in sharp
 224 discontinuities in the variation of fast axis with back-azimuth, and the model misfits are
 225 thus the least sensitive to it. Furthermore, the strength of anisotropy trades off with the
 226 thickness of each layer, as well as with the a-axis plunge (Abt et al., 2010). Nonetheless,
 227 the large values of strength highlight distinct and strong anisotropy in each layer.

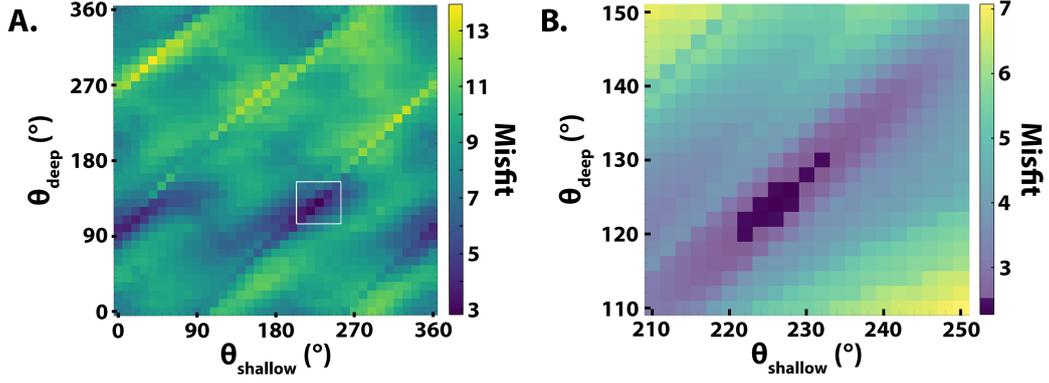


Figure 4. a): Misfit surface along a constant $\delta_{deep} = 40^\circ$, $\alpha_{deep} = 50\%$, $\delta_{shallow} = 30^\circ$, and $\alpha_{shallow} = 40\%$ showing variation in misfit as a function of θ_{deep} and $\theta_{shallow}$. The best fitting model from the coarse grid search is at $\theta_{deep} = 130^\circ$, $\delta_{deep} = 40^\circ$, $\alpha_{deep} = 50\%$; $\theta_{shallow} = 230^\circ$, $\delta_{shallow} = 30^\circ$, $\alpha_{shallow} = 40\%$ with a misfit of 2.8145. This model does not satisfy the F-test criterion corresponding to the finer grid search. The a-axis range probed in the finer grid search, which encloses the best-fitting model from the coarse grid search, is outlined in white. b): As in left, but for the finer grid search. This misfit surface is along constant $\delta_{deep} = 50^\circ$, $\alpha_{deep} = 50\%$, $\delta_{shallow} = 20^\circ$, and $\alpha_{shallow} = 40\%$. The darkened region in the right panel indicates parameters that lie within the 95% confidence F-test limits of the best-fitting model.

228 The delay times we measure exhibit significant scatter (Fig S2). Due to this, back-
 229 azimuthal variation in the delay times is not discernable and the distribution of delay
 230 times is fairly unimodal, centered on a mean of ≈ 1.7 s, albeit with a large standard de-
 231 viation of ≈ 0.5 s. As a result, we follow the convention used in many mantle-scale shear
 232 wave splitting studies (Aragon et al., 2017; Hammond et al., 2014; Dubé et al., 2020) and
 233 do not attempt to incorporate predictions of variations in delay times in our final grid
 234 searches for the best-fitting model parameters in each layer.

235 To investigate how the non-uniform back-azimuthal sampling of fast direction pat-
 236 terns affects the resolvability of model parameters, we conducted a series of tests on syn-
 237 thetic datasets that have the same back-azimuthal distribution as the observed fast di-
 238 rections. One test explores the case in which two-layer anisotropy has the same mean
 239 model parameters as the model that best fits the data, but the parameters are allowed
 240 to vary about those means following a Gaussian distribution with a standard deviation
 241 of 30° for a-axis azimuth, 10° for a-axis plunge, and 10% for anisotropy strength. This

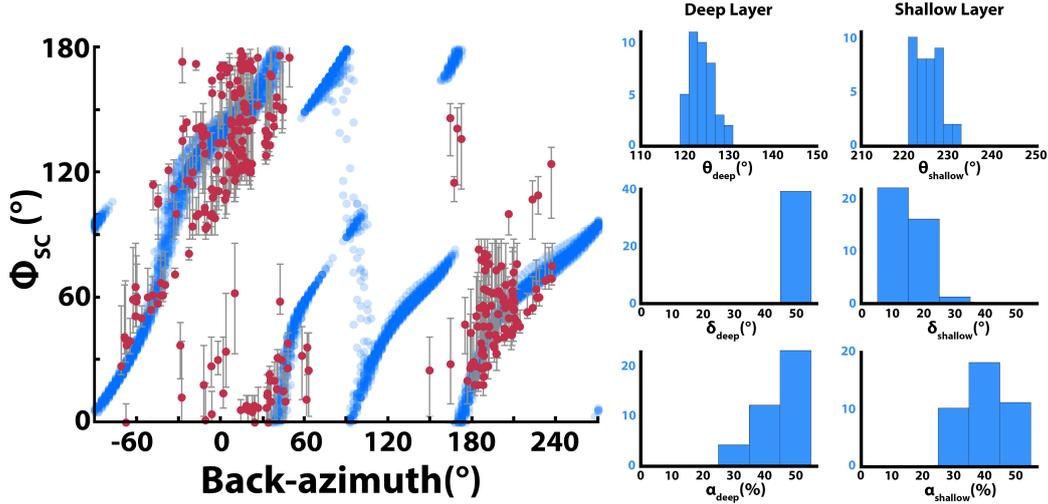


Figure 5. Left panel: suite of best-fitting model predictions of fast axis variation (blue) as a function of back-azimuth, obtained from the fine grid searches that satisfy the misfit criterion corresponding to the F-test at 95% confidence level, overlain on the fast-axis measurements from this study (red). Right panels: Histograms of the model parameters satisfying the misfit criterion.

242 case is intended to represent deformation as a function of depth which is similar across
 243 Greenland, but which varies laterally to a moderate degree. From this distribution of
 244 model parameters, a set of model parameters was drawn and fast directions were pre-
 245 dicted for each back-azimuth in the real dataset for 50 different draws of model param-
 246 eters. To generate the synthetic dataset, at every back-azimuth, we draw a value from
 247 one of the 50 different fast axis predictions. The model which best fits the synthetic dataset
 248 was then determined using the coarse grid of model parameter predictions. One of the
 249 100 iterations is shown in Fig. 6a. This process was repeated 100 times, and the result-
 250 ing distribution of best-fitting model parameters is shown in Fig. 6c-h (blue histograms)
 251 together with the input distribution of model parameters (pink histograms). The retrieved
 252 model parameters are broadly similar to the input distribution of model parameters, in
 253 particular for the upper and lower layer a-axis azimuths which are the best resolved model
 254 parameters. This result supports the argument that meaningful anisotropy parameters
 255 can be retrieved from fast direction data, even when the underlying model varies mod-
 256 erately.

257 In a second test, the possibility that sub-regions have simpler single-layer anisotropy
258 is added to the first scenario. In this test, 50% of the 50 model parameter sets come from
259 the Gaussian distributions about the best-fitting two-layer model (as in Fig. 6a) and 50%
260 are drawn from a distribution of one-layer anisotropy models. In the one-layer model dis-
261 tribution, horizontal a-axis azimuths have means of 50° and 115° , each with a standard
262 deviation of 15° . Again, the process is repeated 100 times, and the resulting best-fitting
263 model parameters are shown in Fig. 6i-n. Although the introduction of the one-layer mod-
264 els produces larger differences between the retrieved (blue) and input (pink) two-layer
265 model parameter distributions, the retrieved a-axis azimuths fall within the input dis-
266 tribution. This result indicates that retrieved two-layer a-axis azimuths can be obtained
267 not only when the underlying model varies moderately, but also when the regional dataset
268 also reflects sub-regions that contain one-layer anisotropy.

269 Additional synthetic tests are described in the supplement (Figs. S5-S7 and Text
270 S1). These tests include a scenario in which the fast directions at each back-azimuth are
271 randomly drawn from the total distribution of fast directions (Fig. S7). This test is equiv-
272 alent to assuming that each synthetic fast direction represents a localized region of one-
273 layer anisotropy, and that any apparent pattern of fast direction with back-azimuth is
274 coincidental. For 100 versions of this case, the resulting distribution of retrieved two-layer
275 models contain model parameters across the range of possible values, and a-axis azimuth
276 ranges for each layer are not well-constrained. These results fundamentally differ from
277 those produced by fitting the observed fast directions. We conclude that the observed
278 pattern of fast-direction versus back-azimuth in Greenland is not coincidental, and that
279 two-layer anisotropy (or at least depth-varying anisotropy) is required.

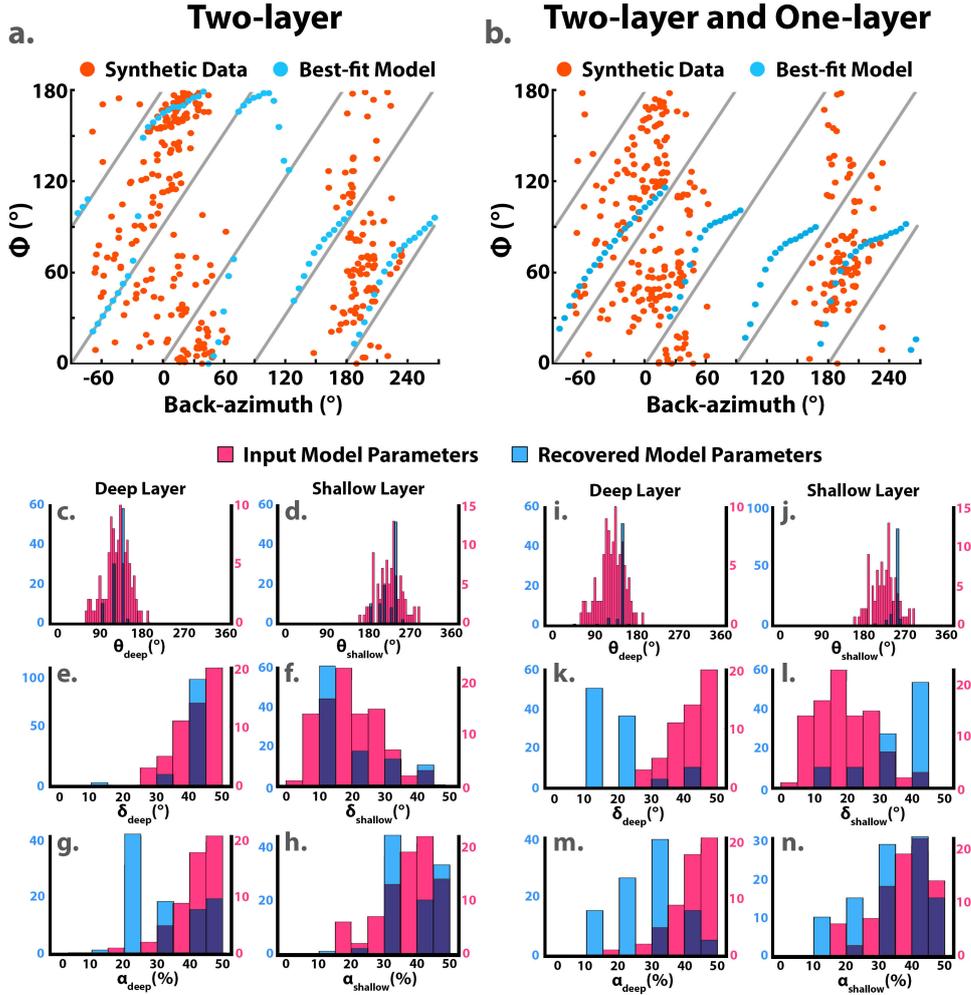


Figure 6. Synthetic test that explores how moderate random variation in two-layer model parameters impacts recovery of their distribution, without (a, c-h) and with (b, i-n) additional one-layer anisotropy. (a) One example of synthetic data (orange points) that are predictions of an input model where parameters are drawn from a Gaussian distributions where the mean is from the best-fitting model for the observed Greenland fast directions. Predictions of the retrieved model (blue) that best fits the synthetic data. (c-h) Distributions of model parameters recovered from fitting synthetic data (blue bars) for 100 cases of model parameters (pink bars) drawn from a Gaussian distribution where the mean is from the best-fitting model for the observations. (b) One example of synthetic data (orange points) that are predictions of a set of input models where 50% are drawn from a Gaussian distributions as described in (a) and 50% are drawn from one-layer models. (i-n) Distributions of model parameters recovered from fitting synthetic data (blue bars) for 100 cases of model parameters (pink bars) as described in (b).

4 Discussion

4.1 Comparison to Prior Studies of Anisotropy

Previous work in Greenland has suggested that significant differences in crustal azimuthal anisotropy exist between northern and southern Greenland (Darbyshire et al., 2018). However, we do not find any significant difference in the back-azimuthal pattern of fast directions between these regions (Fig. 3b). We also do not detect a difference in the pattern for stations on the Greenland ice sheet versus those on the coast. Although ice is an anisotropic mineral, the likely contribution to total splitting observed should be small, especially because there is likely not a coherent fabric throughout an entire column of ice within the ice sheet (e.g. Bentley, 1972; Harland et al., 2013; Smith et al., 2017; Thorsteinsson et al., 1997; Thorsteinsson, 2000).

The widespread coherence of back-azimuthal fast direction variation across the entirety of Greenland is a key feature of our results. Comparison of shear-wave splitting measurements from some previous studies in Greenland and the Canadian high arctic (Dubé et al., 2020; Helffrich et al., 1994) indicates broad agreement (Fig. 7) with the back-azimuthal dependence of the fast direction observed here. However, the fast direction distribution with back-azimuth from Ucisik et al. (2008) is less similar. Strong fast direction variations with back-azimuth were found in some studies from other regions of the Canadian shield (Fig. S3), but differences in these patterns relative to those in this study suggest regional variations in anisotropic parameters (Bastow et al., 2011; Darbyshire et al., 2015; Liddell et al., 2017; Snyder et al., 2013).

Among published models of azimuthal anisotropy, the regional models based on group velocity from Darbyshire et al. (2018) provide information at lateral scales most comparable to our results. At shallow mantle depths, Darbyshire et al. (2018) indicate NE-SW fast directions in the far north and south of Greenland, which are consistent with a-axis azimuths in the lithospheric layer of the best-fitting models found here, and NW-SE fast direction beneath the central latitudes of Greenland, which do not agree with our lithospheric parameters. However, the Darbyshire et al. (2018) group velocity results indicate weak anisotropy at mantle depths and represent constraints on only the shallow lithospheric mantle, leaving open the possibility that the two studies are compatible.

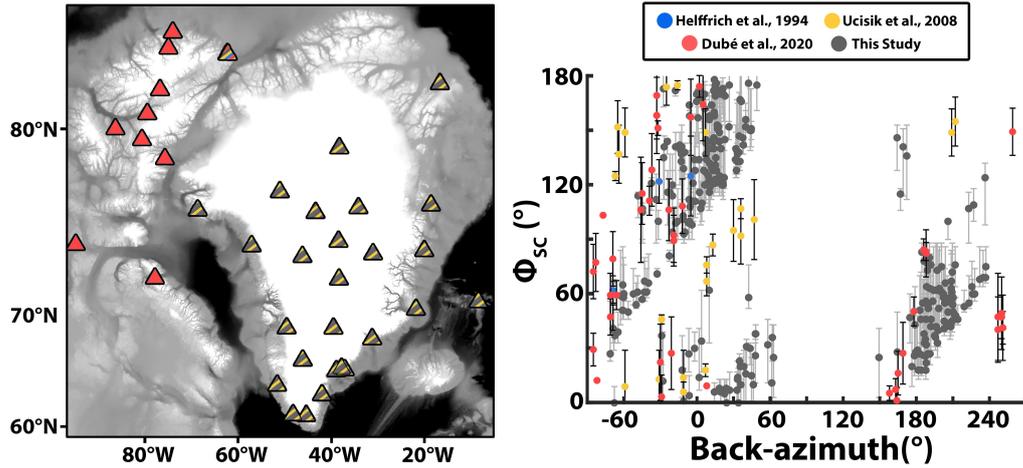


Figure 7. Left: Locations of stations outside of and in Greenland used in certain other studies. Colors are the same in the map and back-azimuth plot. Right: Comparison of our shear wave splitting fast polarizations (gray) with the results from other studies. $N_{other\ studies} = 70$.

311 4.2 Interpretation of Two-layer Anisotropy Models in Terms of Man- 312 tle Deformation

313 A simple, first-order prediction for shear in the asthenosphere is that it would be
314 driven by absolute plate motion, resulting in olivine a-axis azimuths that are parallel to
315 absolute plate motion. Olivine a-axes in the deeper layer of the two-layer anisotropy mod-
316 els that provide acceptable fits to the observed Greenland fast directions (120° to 130°)
317 are approximately aligned with absolute plate motion in Greenland assuming the no-net-
318 rotation reference frame plate motion model of Argus et al. (2011).

319 Consistency of asthenospheric a-axes azimuths and plate motion shear differs from
320 the conclusions of some previous studies of anisotropy (Darbyshire et al., 2015; Liddell
321 et al., 2017). These studies assert that asthenospheric anisotropy parallel to plate mo-
322 tion should not be expected, because the North American plate speed (~ 20 mm/yr) is
323 slower than what is required (~ 40 mm/yr) to develop basal drag fabric (Debayle & Ri-
324 card, 2013). However, the 40 mm/yr threshold (Debayle & Ricard, 2013) refers to whole-
325 plate alignment of fabric with plate motion. Indeed, Debayle and Ricard (2013) state
326 that for slow moving continental plates, the correlation between asthenospheric fabric
327 and plate motion is more complicated, but that agreement can persist over large scales,
328 citing central and eastern North America as an example. Thus, when considering a study

329 region which is smaller than an entire plate, slow plate velocities do not rule out astheno-
330 spheric fabric in agreement with plate motion, either from basal drag or secondary con-
331 vection (Debayle & Ricard, 2013).

332 We also compared the parameters of the deeper layer of anisotropy to models of
333 mantle flow that account for mantle temperature, buoyancy and viscosity, as well plate
334 motion boundary conditions. However, due to different boundary conditions and other
335 model assumptions, predictions for asthenospheric flow directions beneath Greenland dif-
336 fer between studies, and it is possible to find models which are broadly consistent with
337 the acceptable a-axes found here (Conrad & Behn, 2010; Marquart et al., 2007) or in-
338 consistent (e.g. Colli et al., 2018; Conrad & Behn, 2010; Marquart et al., 2007; Mihalfy
339 et al., 2008).

340 The agreement between our well-constrained deep layer a-axis fast azimuths and
341 the no-net-rotation plate motion directions from Argus et al. (2011) make a strong case
342 for asthenospheric anisotropy produced by shearing due to plate motion. However, the
343 fact that acceptable a-axis plunges are $\sim 50^\circ$ pose a complication for this model. Even
344 though a-axis plunge is less well-resolved than a-axis azimuth in the two-layer model-
345 ing, models with near-horizontal lower layer a-axes produce significantly worse fits to the
346 observed fast polarization directions. The apparent a-axis plunges suggest vertical flow
347 components, for example due to asthenosphere diverted beneath Greenland's thick cra-
348 tonic lithosphere or at the edges of a potential channel of thin lithosphere created by the
349 thermal signature of the Iceland hotspot (Fig. 8).

350 The relationship between the shallow layer a-axis azimuths inferred from the mod-
351 eling (222° to 232°) and lithospheric deformation fabrics is difficult to evaluate because
352 the Greenland ice sheet occludes much of the geologic evidence typically used to com-
353 pare lithospheric anisotropy fabrics with the deformation signatures of major tectonic
354 events. Nonetheless, inferred a-axis azimuths are consistent with deformation fabrics from
355 Proterozoic and Archean orogenic events in western and northern Greenland. The Trans-
356 Hudson Orogeny, which occurred 1.8 Ga, is responsible for Greenland's prominent Nagssug-
357 toqidian belt, although the direction of compression and shape of the tectonic bound-
358 ary is obscured by the ice sheet and has been interpolated many ways across Greenland
359 (Antonijevic & Lees, 2018; Dawes, 2009; Henriksen et al., 2009; Pourpoint et al., 2018).
360 Unobscured by the ice sheet, shear zones in Western Greenland closely associated with

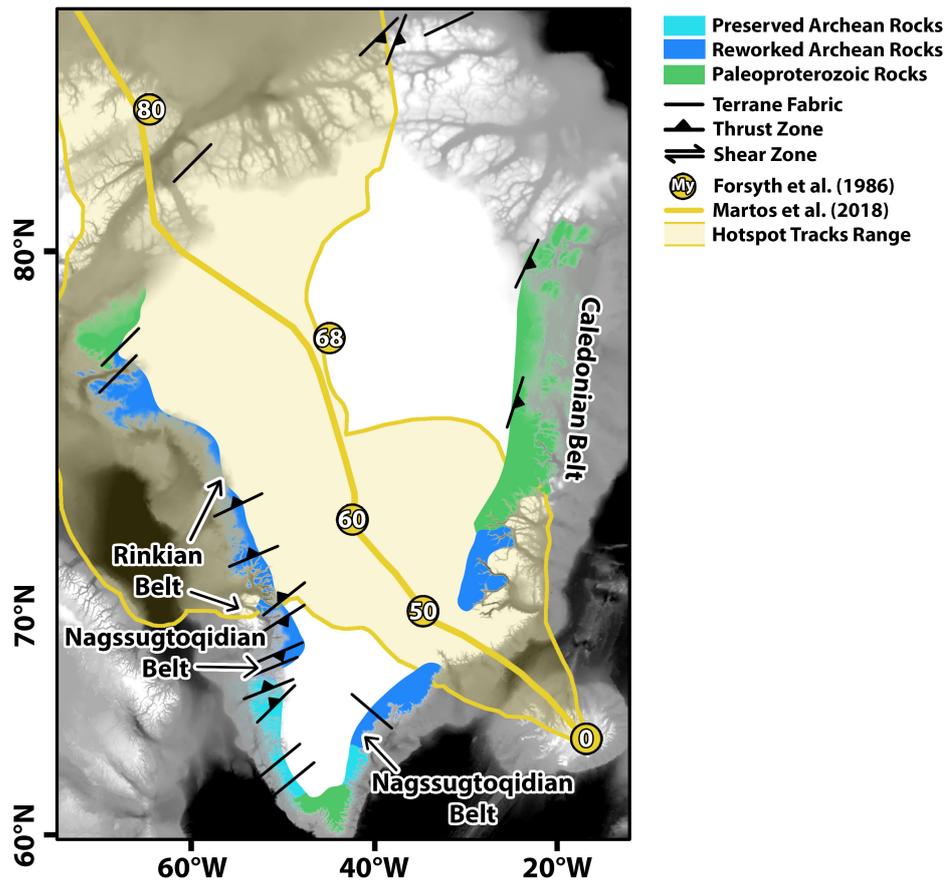


Figure 8. Map of relevant geology and tectonic features. Underlying gray field is ETOPO1 topography (Amante & Eakins, 2009) with basement ages, terrane fabrics, thrust zones, and orogenic belts from Henriksen et al. (2009); Dawes (2009); Higgins and Leslie (2000) Van Gool et al., (2002), Sanborn-Barrie et al. (2017), and Van Gosen and Piepjon (1999). The yellow-shaded region represents the extent of previously proposed hotspot tracks as compiled in Martos et al. (2018); the bold yellow line is their proposed hotspot track. Numbers show hotspot surface projections at different times from Forsyth et al. (1986).

361 the Nagssugtoqidian orogen trend ENE-WSW to NE-SW (van Gool et al., 2002; Bak et
 362 al., 1975), parallel to local thrust zones and older Archean terranes (van Gool et al., 2002;
 363 Henriksen et al., 2009; Korstgård et al., 1987). NW-directed thrusting has also been in-
 364 ferred in the Proterozoic Rinkian orogen to the north in western Greenland (Sanborn-
 365 Barrie et al., 2017; van Gool et al., 2002), and in thrust zones at Greenland's northern
 366 margin (Von Gosen & Piepjohn, 1999). These indicators of lithospheric deformation are
 367 consistent with acceptable a-axis azimuths from the two-layer anisotropy modeling, and

368 this agreement suggests that similarly-oriented deformation fabrics are also present be-
369 neath the ice sheet.

370 However, deformation indicators in eastern Greenland are less consistent with the
371 overall NE-SW acceptable lithospheric a-axis orientations. In southeastern Greenland,
372 van Gool et al. (2002) infer an ESE structural grain. At two stations within this zone
373 (ANGG and KULG), fast directions are predominantly NW-SE, raising the possibility
374 that lithospheric deformation in this zone is rotated from the shallow layer trend indi-
375 cated by the modeling of the complete set of Greenland stations. In addition, thrust fronts
376 in the Paleozoic Greenland Caledonides are oriented \sim N-S (Dawes, 2009). This incon-
377 consistency with the overall NE-SW oriented shallow layer a-axis orientation may reflect an
378 unresolved local variation in lithospheric a-axis alignment, or that lithospheric fabric as-
379 sociated with the Caledonian orogeny was limited in its depth extent, possibly due to
380 decoupling of the Laurentian retro-lithosphere (Hodges, 2016).

381 Local deformation associated with rifting has impacted both eastern and western
382 Greenland at differing scales. The orientation of extension associated with Labrador sea
383 rifting in the west is parallel to the shallow layer a-axis orientations, and strong crustal
384 anisotropy associated with mineral alignment during this process is resolved by Clement
385 et al. (1994) via shear-wave splitting. On the other hand, local rifting basins in East Green-
386 land show W-E and NW-SE extension (Henriksen et al., 2009), the latter being perpen-
387 dicular to the shallow a-axis orientations inferred here. This discrepancy may not be sig-
388 nificant if orogenic deformation over longer length-scales dominates lithospheric fabrics,
389 relative to more localized rifting events.

390 By comparing SKS and SKKS measurements, Dubé et al. (2020) show that anisotropy
391 in the lower mantle impacts measurements at station ALE. Lower mantle anisotropy has
392 been imaged below Iceland and shown to impact differential SKS-SKKS measurements
393 at Greenland stations (Wolf et al., 2019). We did not resolve consistent discrepancies be-
394 tween SKS and SKKS splitting in our dataset (Fig. S4). Nonetheless, further work should
395 be conducted to constrain the extent to which anisotropy from lower mantle or crustal
396 sources (e.g. Clement et al., 1994) may impact the shear-wave splitting measurements
397 in this study.

398 **5 Conclusions**

399 Using 299 new splitting measurements from stations across Greenland, we have found
400 a consistent pattern of fast direction variation in back-azimuth which indicates the pres-
401 ence of multi-layer anisotropy. We used grid searches to solve for two-layer models of anisotropy
402 that provide acceptable fits to the fast directions. Acceptable a-axis azimuths are 222°
403 to 232° in the shallow layer, and 120° to 130° in the deep layer.

404 The modeling results are consistent with an interpretation where anisotropy in the
405 deeper layer represents asthenospheric shearing due to plate motion in a no-net-rotation
406 reference frame. The upper layer is consistent with lithospheric anisotropy due to Pro-
407 terozoic and Archean orogenic events, as indicated by tectonic fabrics in western and north-
408 ern Greenland. The strong variations in back-azimuthal pattern of fast directions in Green-
409 land, combined with prior work in the Canadian high arctic, are consistent with coher-
410 ent lithospheric deformation from Proterozoic and Archean orogenesis on a broader scale
411 than previously appreciated.

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