

Abstract

A basin-wide significant shallow bias is found in the southern tropical Pacific thermocline in an ensemble of models from the coupled model intercomparison project phases 5 and 6. In contrast to observations, where the southern thermocline is far deeper than its northern counterpart, models have a hemispherically symmetric tropical thermocline. The shallow thermocline bias is closely related to the well known double intertropical convergence zone (ITCZ) bias, as shown by ensemble partitioning. The physical thermocline (i.e., depth of maximal vertical thermal gradient) is more strongly linked to the double ITCZ bias than the commonly used 20°C isotherm thermocline proxy. A shallow thermocline bias is further found to be associated with wider separation of double ITCZ peaks, stronger southern precipitation, a stronger cold tongue, and a spurious south equatorial counter current. Climatic implications and feedback mechanisms between the biases are discussed.

Plain Language Summary

In this study we find that in global climate model simulations the depth of the tropical south Pacific thermocline is shallower than in reality. The thermocline is the layer that separates the warm upper ocean from its cold depths. When the thermocline is shallower, the capacity of the upper ocean to absorb heat from the atmosphere is reduced. This means that the error in the modeled depth of the thermocline could alter the modeled heat balance between the ocean and atmosphere, as well as the heat imbalance between the northern and southern hemispheres. We further find that a well known error of excessive rain in the tropical south Pacific, which has been a major problem in models for the past three decades and remains unsolved, is worse in models that have this thermocline simulation problem. By drawing connections between the two model inaccuracies, progress can be made in understanding their sources. Correcting the thermocline simulation error and the errors associated with it is critical to the reliability of global scale model predictions.

1 Introduction

The thermocline is the layer in the ocean of maximal temperature decline. It separates the upper ocean, where air-sea interactions occur, from the abyssal ocean, affording the two regions distinct circulations, thermodynamic properties and characteristic timescales (Knauss & Garfield, 2016). Impacts of thermocline depth on the climate system are diverse and fundamental. In particular, it influences sea surface temperature (SST) by setting the thermal foot-

41 print of Ekman divergence driven upwelling. Such is the case, for example, along the equa-
42 torial Pacific and Atlantic cold tongues, where a strip of cold surface water extends from the
43 east due to zonal sloping of the thermocline (Sverdrup et al., 1942). Thermocline depth also
44 regulates ocean heat storage, ocean energy transport and, by extension, the global energy bud-
45 get (Boccaletti et al., 2004; Vialard et al., 2001; Schott et al., 2004). Since thermocline depth
46 marks the lower boundary of the ocean layer most susceptible to mechanical and radiative forc-
47 ing, its representation in climate models is paramount (Harper, 2000; Burls et al., 2017; Thomas
48 & Fedorov, 2017).

49 Models participating in phases 3 and 5 of the Coupled Model Intercomparison Project
50 (CMIP) are known to systematically produce a too shallow equatorial thermocline (e.g., Zheng
51 et al., 2012; Li & Xie, 2012, 2014; Castaño-Tierno et al., 2018). However, the meridional struc-
52 ture of the tropical thermocline and its relation to the intertropical convergence zone (ITCZ,
53 a zonal band near the equator where tropical rain is concentrated) has not been fully inves-
54 tigated in coupled General Circulation Models (GCMs). Tropical Pacific meridional thermo-
55 cline structure is particularly interesting in light of the well known double ITCZ bias (DIB),
56 characterized by excessive precipitation south of the equator primarily in the Pacific (Fig. 1a-
57 c; Mechoso et al., 1995; Lin, 2007; Adam, Schneider, & Brient, 2018). A DIB is often ac-
58 companied by a cold-tongue bias, where the modeled cold tongue is too cold and westward
59 extended (Li & Xie, 2014). Coupled ocean-atmosphere dynamics have been shown to be at
60 the root of the DIB, but the source of the bias remains unknown.

61 In the tropics, a warm well-mixed surface layer typically sits above a strong thermocline
62 year-round (Knauss & Garfield, 2016). A strong or stable thermocline (i.e., a layer of large
63 vertical thermal gradient, $\partial_z \theta$) forms a barrier that weakens mixing between the upper and abyssal
64 layers and is often shallower and thinner (Vallis, 2017). Thermocline depth may be shoaled
65 by modulation of the diffusivity of heat downward and advection of cold water upwards be-
66 tween the abyssal and upper layers, but also by surface warming or abyssal cooling through
67 increased stratification (Vallis, 2017). Wind driven Ekman transport is pivotal to thermocline
68 structure as well: subtropical current convergence causes subduction and thermocline deep-
69 ening; current divergence at the equator, along eastern coastlines and near the ITCZ, causes
70 upwelling and thermocline shoaling (Luyten et al., 1983). In the limit of geostrophy, the ther-
71 mocline shoals to the left (right) of ocean currents when facing downstream in the northern
72 (southern) hemisphere (Knauss & Garfield, 2016). A shallow thermocline is expected to be

73 more reactive to surface processes such as wind stress and radiative changes on seasonal and
74 inter-annual timescales.

75 In this paper we study the tropical Pacific thermocline in an ensemble of CMIP phases
76 5 and 6 (CMIP5/6) models. Comparison with reanalysis data reveals a significant shallow bias
77 in the southern tropical Pacific thermocline. Section 2 describes the data and diagnostics. In
78 section 3 we characterize the shallow thermocline bias. We then examine tropical influences
79 of the bias in section 4 and end in section 5 with a discussion of possible sources and impli-
80 cations, as related to the double ITCZ and cold tongue problems.

81 **2 Data and Methods**

82 **2.1 Data**

83 For oceanic reference data, we used the European Center for Medium-range Weather Fore-
84 casts Ocean Reanalysis/Analysis System phases 4 and 5 (ORAS4/5) (Zuo et al., 2019), the Global
85 Ocean Data Assimilation System (GODAS) (Behringer & Xue, 2004) and the Simple Ocean
86 Data Assimilation version 3.12.2 (SODA3) reanalysis (Carton et al., 2018). Precipitation data
87 was taken from the National Oceanic and Atmospheric Administration’s (NOAA) Global Pre-
88 cipitation Climatology Project (GPCP), version 2.2 (Adler et al., 2003) and from the Climate
89 Prediction Center (CPC) Merged Analysis of Precipitation (CMAP) (P. Xie & Arkin, 1996).
90 SST data was taken from version 3b of the Extended Reconstructed SST (ERSST.v3b) dataset,
91 provided by NOAA’s National Climatic Data Center (Smith et al., 2008). For model data, we
92 analyzed 37 CMIP5 and 36 CMIP6 models (see Tables T1/2), based on availability. We used
93 monthly data from the first realization of historical simulations (coupled GCMs driven by pre-
94 scribed atmospheric compositions). All analyses were performed on monthly climatology com-
95 puted between 1979–2005.

96 **2.2 Diagnostics**

97 *Thermocline Depth Estimation*

98 Due to sparse ocean measurements and low vertical ocean model resolution, thermocline
99 depth is commonly approximated as an isotherm (Fiedler, 2010). In the tropical Pacific, the
100 20°C isotherm (Z_{20}) has been used widely in observational and model analyses (e.g., Zheng
101 et al., 2012; Li & Xie, 2012)). However, important distinctions exist between the depth of the
102 physical thermocline (i.e., depth of maximal $\partial_z \theta$, denoted here as Z_{tc}) and isothermal defi-

103 nitions, especially in their sensitivity to tropical mechanical forcing. Specifically, Z20 tends
 104 to be deeper, flatter and less sensitive to seasonal warming (Castaño-Tierno et al., 2018). There-
 105 fore, in the present study we look at both Ztc and Z20.

Z20 is computed through linear interpolation of ocean potential temperature (Castaño-Tierno et al., 2018). Ztc is calculated using

$$Z_{tc} = \frac{\int_{z_t}^{z_b} (\partial_z \theta)^n z dz}{\int_{z_t}^{z_b} (\partial_z \theta)^n dz} \quad (1)$$

106 where subscripts t and b denote top and bottom depths, respectively, and $\partial_z \theta$ is the vertical
 107 gradient of the potential temperature calculated between each two consecutive levels. This re-
 108 duces grid dependence, where n acts as a smoothing parameter, such that for $n = 1$, Ztc is the
 109 depth of the centroid of $\partial_z \theta(z)$, and for $n \rightarrow \infty$, Ztc is the depth of the exact half-grid local
 110 maximum with no smoothing (Adam, Grise, et al., 2018). We took z_t to be the average of the
 111 upper two levels (~ 10 meters in most datasets), z_b as the average of the two levels closest
 112 to 400 meters in each dataset, and a smoothing factor of $n=15$. Results are not sensitive to ei-
 113 ther parameter.

114 *Tropical Asymmetry Index*

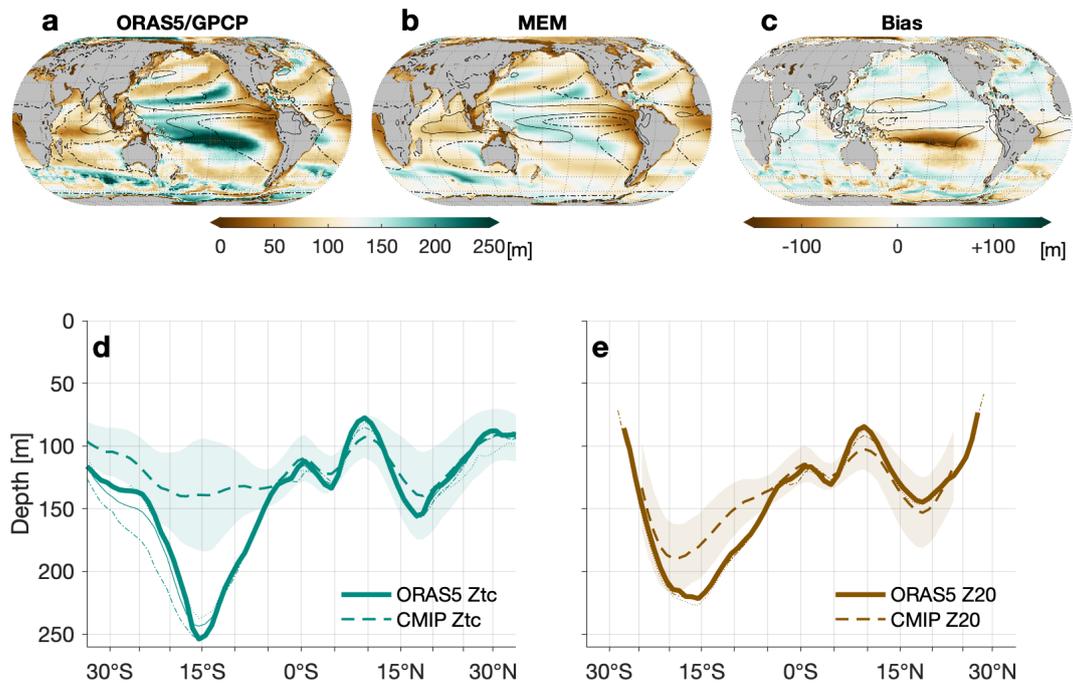
115 The tropical Pacific asymmetry index A_f (Hwang & Frierson, 2013) of a field f zon-
 116 ally averaged within the Pacific ($160\text{--}260^\circ$), is defined as the (area-weighted) north-tropical
 117 average ($0\text{--}20^\circ\text{N}$) minus the south-tropical average ($0\text{--}20^\circ\text{S}$), normalized by the tropical av-
 118 erage ($20^\circ\text{S}\text{--}20^\circ\text{N}$):

$$A_f = \frac{\bar{f}_{0-20N} - \bar{f}_{20S-0}}{\bar{f}_{20S-20N}}. \quad (2)$$

119 A_f is positive if f is greater north of the equator than south, negative if f is greater south of
 120 the equator than north, and zero if the distribution is symmetric about the equator.

121 **3 Thermocline Biases**

129 Thermocline depth (Ztc) is shown in Fig. 1a for the ORAS5 reanalysis. Tropical Pacific
 130 deepening is the most prominent feature, found between 10° and 20° off the equator, angled
 131 relative to the equator such that maximal depth is closer to the equator in the west. Substan-
 132 tial hemispheric asymmetry also stands out, with the southern deepening greater and wider.
 133 Mean Ztc in a combined ensemble of CMIP5 and CMIP6 models (MEM) is shown in Fig. 1b.



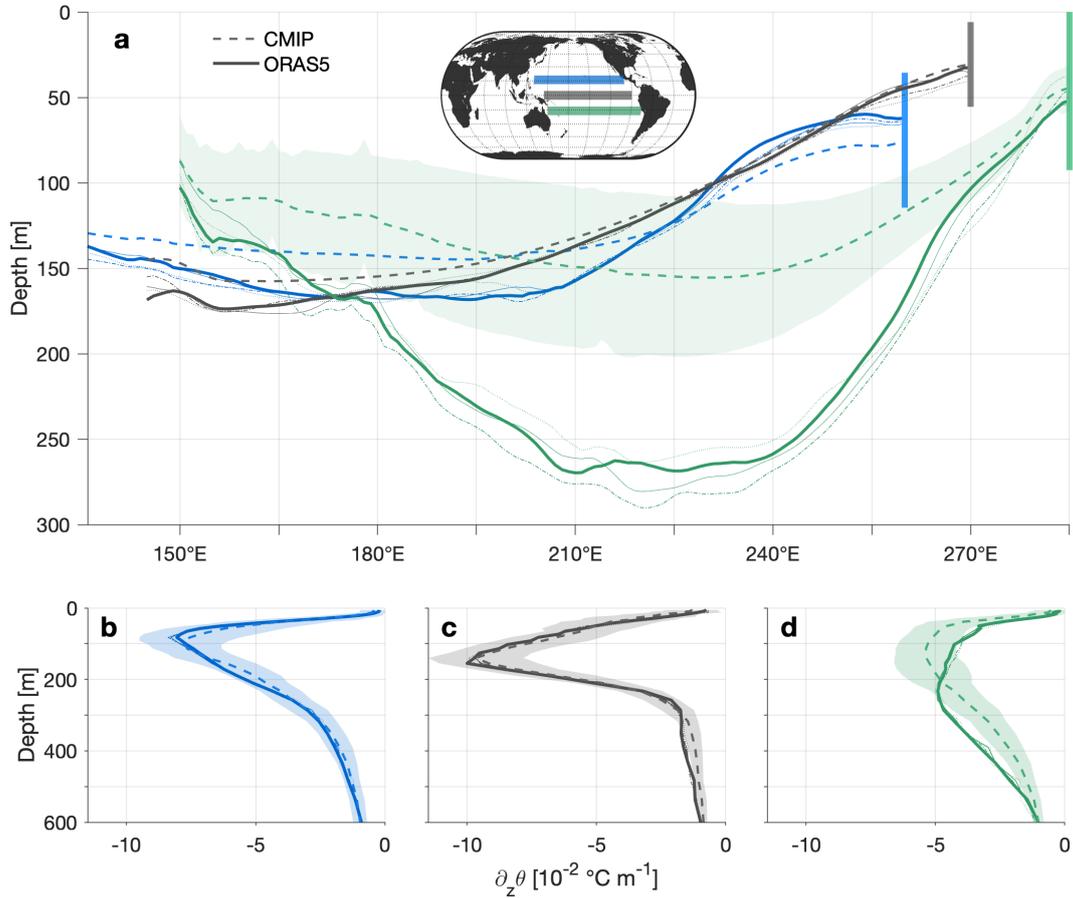
122 **Figure 1.** Annual mean thermocline depth (Z_{tc} , color contours) and precipitation (line contours) in (a)
 123 observations (ORAS5 and GPCP), (b) CMIP model ensemble mean (MEM), and (c) MEM bias. Solid (dash-
 124 dotted) precipitation contours indicate values greater (smaller) than 3 mm day^{-1} in panels (a,b) and 1.5 mm
 125 day^{-1} in panel (c). (d) Thermocline depth and (e) 20°C isotherm depth (Z_{20}) in the Pacific ($160\text{-}260^\circ\text{E}$) in
 126 ORAS5 (thick solid), ORAS4 (thin solid), GODAS (dotted), SODA (dash-dotted) and MEM (dashed) with
 127 one standard deviation of the ensemble (shaded). All data is averaged from 1979–2005. To avoid seasonal
 128 biasing, Z_{20} is truncated poleward of $\sim 25^\circ$ due to (hemispheric) summer surfacing.

134 The subtropical Pacific deepening in models is severely less pronounced and less hemispher-
 135 ically asymmetric, producing a mild bias in the northern hemisphere (NH) but an acute bias
 136 in the southern hemisphere (SH) (Fig. 1c). A similar southern bias of slightly lesser intensity
 137 is found in the south western Atlantic. We hereafter focus our analysis on the tropical Pacific,
 138 but note that the key elements of the shallow thermocline bias exist in the Atlantic, with unique
 139 features owing to the dynamic differences between the basins, which are beyond the scope of
 140 this work.

141 Pacific zonal mean Z_{tc} and Z_{20} are shown in Figs. 1d and 1e, respectively. The observed
 142 thermocline is shallow in the subtropics, reaches maximal depth around 15° off the equator
 143 and shoals from there equatorward. The exception to this pattern is the NH shoaling around
 144 9°N , just poleward of the ITCZ, caused by wind-driven divergent meridional currents. Hemi-
 145 spheric asymmetry is clearly observed, with over 100 meters differentiating the deepest south-
 146 ern point from the deepest northern point, leading to a tropical Pacific asymmetry index of $A_{Z_{tc}} =$
 147 -0.32 . Z_{20} exhibits all of these features as well, with a similar asymmetry index of $A_{Z_{20}} =$
 148 -0.31 . Equatorward of 10° Z_{20} deviates minimally from Z_{tc} . However, poleward of this, Z_{20}
 149 is generally shallower than Z_{tc} and diverges from Z_{tc} completely as it surfaces poleward of
 150 $\sim 25\text{--}35^\circ$ (depending on the season).

151 Biases in the NH are generally insignificant, perhaps with the exception of a mild deep-
 152 ening bias in the vicinity of the ITCZ (Zhu et al., 2021). In contrast with observations, the ther-
 153 mocline in models is significantly more flat and hemispherically symmetric, with a mean mod-
 154 eled $A_{Z_{tc}}$ value of -0.11 . A robust bias is seen between $5^\circ\text{--}30^\circ\text{S}$, which reaches over 100 me-
 155 ters at the deepest point of the observed thermocline. Inter-model spread is also maximal at
 156 this point, reaching a standard deviation of 50 meters. A similar but weaker bias is found in
 157 Z_{20} , with model $A_{Z_{20}} = -0.15$. The biased region is also where Z_{20} and Z_{tc} diverge most,
 158 but in opposite directions in observation and models. Since Z_{20} is deeper than Z_{tc} in mod-
 159 els, but shallower than Z_{tc} in observations, temperature at Z_{tc} is distinctly warmer in mod-
 160 els than in observations as a result of the shallow bias. We note that the bias in Z_{tc} is not sea-
 161 sonal (Fig. S1a); SH Z_{20} seasonal variance is very weak in observations, and even more so
 162 in models (Fig. S1b).

169 To further understand the spatial structure of the thermocline bias, Fig. 2a shows merid-
 170 ional averages over equatorial and sub-equatorial sections ($5^\circ\text{N}\text{--}5^\circ\text{S}$, and $10^\circ\text{--}20^\circ$ in each hemi-
 171 sphere). The southern bias is pan-Pacific, and in proportion to thermocline depth, such that



163 **Figure 2.** (a) Thermocline depth meridionally averaged in northern (10-20°N, 130-260°E, blue), equato-
 164 rial (5°S-5°N, 145-270°E, grey) and southern (10-20°S, 150-285°E, green) sections, as marked in colors on
 165 the inset map. Vertical lines indicate the maximal span of ± 1 standard deviation of model spread within the
 166 averaging region for each section. (b,c,d) Vertical potential temperature gradient, averaged meridionally and
 167 zonally over the sections shown in (a). The mixed layer at the top is imperceptible because of low vertical
 168 resolution. Datasets are denoted as in Fig. 1(d,e).

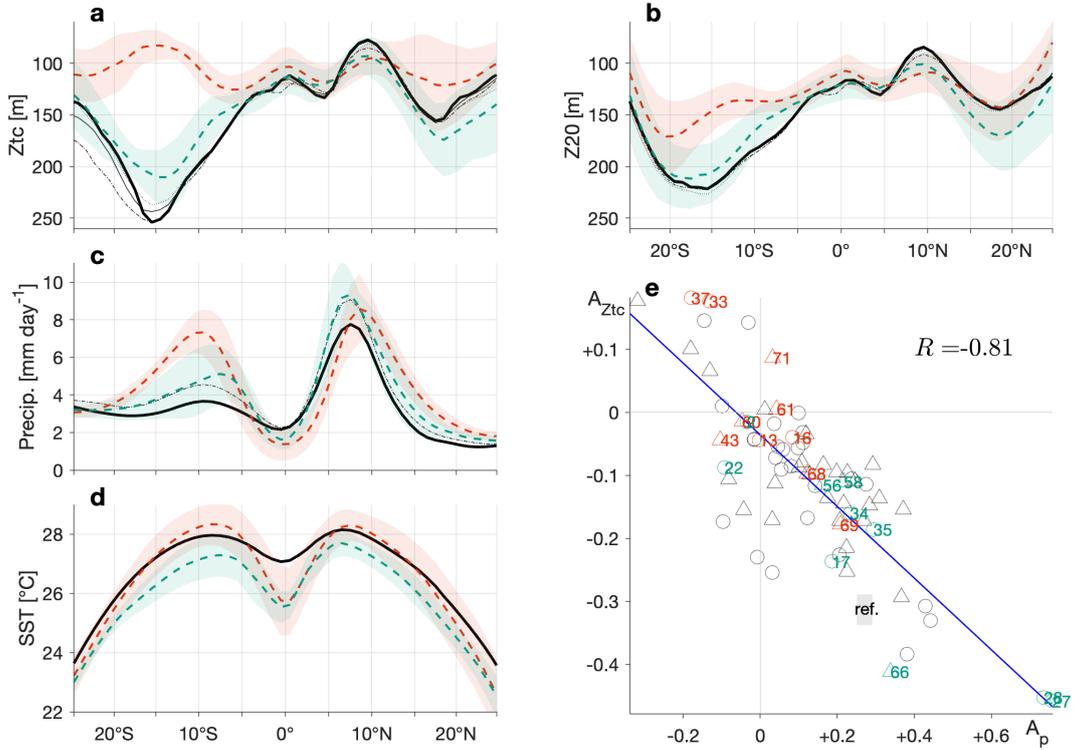
172 modeled zonal thermocline sloping is too weak. The NH and equatorial sections also have weak
 173 modelled zonal gradients, though both the observed gradient and the bias in it are smaller. As
 174 a result, zonal thermocline sloping is equivalent in all three sections in models. A similar but
 175 weaker bias is seen in Z20 in the southern section, which is further extended to the west rel-
 176 ative to Ztc, since Z20 does not shoal in the west as Ztc does (Fig. S2a; Wang et al., 2000).

177 Figures 2b–d show vertical profiles of the vertical potential temperate gradient for the
 178 equatorial and sub-equatorial sections. The thermocline in the equatorial and NH sections is
 179 on average stronger and narrower than that of the SH, in models and observations alike. In
 180 the SH section, in addition to the shallow bias, the modeled thermocline is narrower than ob-
 181 served. This is true in section mean (Fig. 2d), as well as locally within it (Fig. S2b-g). No clear
 182 bias is seen in the strength of the thermocline. This is surprising, as one might expect a shal-
 183 lower and narrower thermocline to be less diffuse (Vallis, 2017). It is curious that while the
 184 modeled thermocline is more symmetric about the equator than observed, its strength is as asym-
 185 metric as in observations. This may point to surface rather than internal ocean processes at
 186 play. More significant positive biases in thermocline strength are found closer to the equator
 187 where the thermocline shoals (Fig. S3).

188 **4 Inter-Model variance**

196 To characterize the climate state associated with the shallow thermocline bias, in Fig. 3
 197 we partition the ensemble, identifying the 10 models with the largest thermocline bias (shal-
 198 lowest thermocline) and smallest bias (deepest thermocline) in the 10-20°S section within the
 199 Pacific (160-260°E), where model spread is largest (Fig. 1d). The results are not sensitive to
 200 this particular range, nor to the number of models or to specific ensemble members, and nearly
 201 equal representation of CMIP5 and CMIP6 is found in both end member groups (see Tables
 202 T1/2). We find that the shallow thermocline models have a ridged southern thermocline, while
 203 the deep thermocline models have the observed southern dip, albeit not quite as pronounced
 204 (Fig. 3a). Accordingly, the ridged models have very weak asymmetry, while in dipped mod-
 205 els $A_{Ztc} = -0.22$, still significantly lower than observations. The ridged thermocline models
 206 have a shallower thermocline across the tropical Pacific; however, a large and robust bias is
 207 only found in the SH.

208 Precipitation in the ridged models has a stronger southern ITCZ bias, wider separation
 209 of precipitation peaks, and near zero asymmetry (Fig. 3c). In contrast, dipped models have a



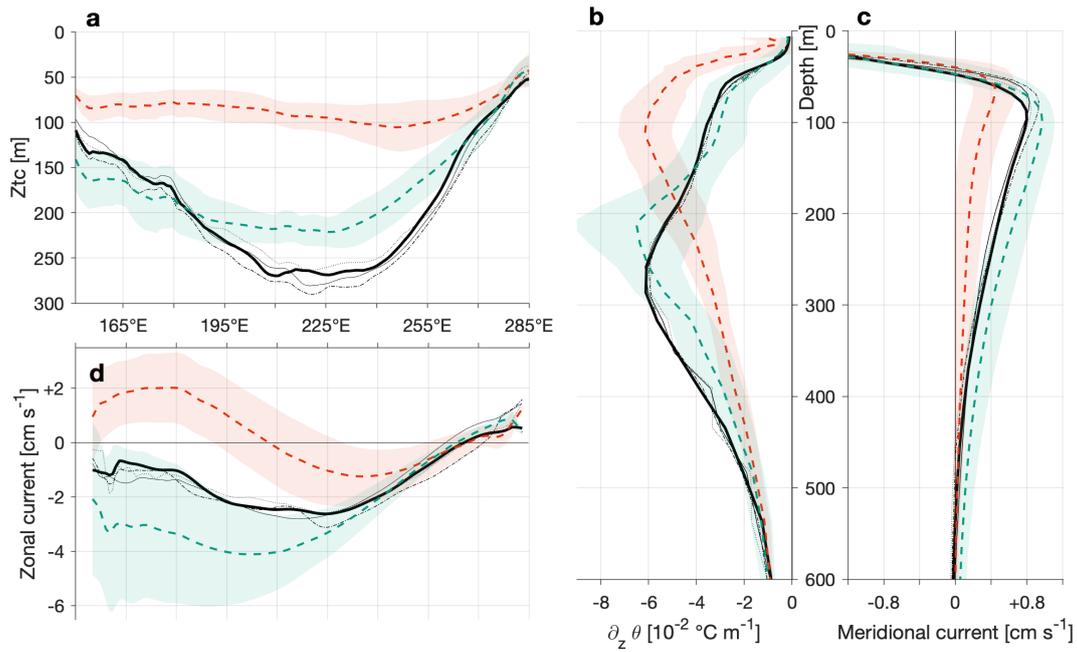
189 **Figure 3.** (a) Thermocline depth (Z_{tc}), (b) 20°C isotherm depth (Z_{20}), (c) precipitation, and (d) sea surface
 190 temperature for observations (black; a,b: see Fig. 1(d,e) for notation of datasets; c: GPCP (solid), CMAP
 191 (dash-dotted); e: ERSST), and mean of "ridged" models (red), and of "dipped" models (turquoise) (see text
 192 for partitioning), with ± 1 inter-model standard deviation (shaded). (e) Relation of precipitation (A_p) and ther-
 193 mocline ($A_{Z_{tc}}$) asymmetry indices: grey rectangle gives observed range; circles and triangles denote CMIP5
 194 and CMIP6 models, respectively, colors correspond to the model groups in (a-d) and numbers correspond to
 195 Tables T1/2. All quantities are averaged in the Pacific sector (160-260°E).

210 precipitation asymmetry index of 0.28, similar to the mean observed value (0.27). A close re-
211 lation is found between asymmetry of thermocline depth and of precipitation (Fig. 3e), with
212 correlation for the full ensemble of models $R = -0.81$, stronger than that found between pre-
213 cipitation and Z20 asymmetry ($R = -0.75$) or precipitation and SST asymmetry ($R = 0.76$).
214 Of 73 models in the CMIP ensemble, only four have greater thermocline asymmetry than ob-
215 servations.

216 Dipped models have a tropical-wide cold SST bias (Fig. 3d), consistent with a gener-
217 ally deeper thermocline. The equatorial cold tongue region has a cold bias in both ensemble
218 members; however, the off-equatorial meridional SST gradient is much larger in ridged mod-
219 els than in dipped ones or in observations. The resultant stronger cold tongue is consistent with
220 a shallow equatorial thermocline bias (Zheng et al., 2012). Also in agreement with these re-
221 sults, off-equatorial meridional SST gradients have been shown to be as important to surface
222 convergence and ITCZ formation as SST itself (Back & Bretherton, 2009). Moreover, the ridged
223 models are asymmetrically warmer just south of the equator, creating an asymmetry bias in
224 SST gradients that promotes south-equatorial deep convection and a DIB (Zhou et al., 2020).

225 Ridging is not seen to the same extent in Z20 (Fig. 3b). Instead, Ztc ridging is expressed
226 in Z20 as a flattening in place of deepening. Nonetheless, lack of asymmetry is nearly as great
227 in Z20 of the ridged thermocline, with $A_{Z20} = -0.10$. Z20 has no southern bias in the dipped
228 group, though its hemispheric asymmetry is reduced relative to observations ($A_{Z20} = -0.18$).
229 The temperature at the thermocline is therefore unbiased in the dipped group, whereas the south-
230 ern thermocline is too warm in the ridged group. Lack of ridging in Z20 in the south may in-
231 dicate a fundamental difference relative to the northern ridging (see discussion).

235 Shown in Fig. 4a, a strong separation between the groups emerges in SH thermocline
236 zonal sloping, which is nearly as large in the dipped group as in observations. In contrast, a
237 zonal slope is practically absent in the ridged group, to the point that thermocline asymme-
238 try becomes positive westward of 230°E (compare to Fig. S4a). But even the dipped group
239 exhibits only a mild degree of western shoaling, possibly related to a westward extended cold
240 tongue and resultant warm pool cooling. Vertically, the thermocline is thicker and deeper in
241 observations than in either group, with the ridged group most narrow (Fig. 4b). Thermocline
242 strength does not have a clear bias among models, further questioning the role of stability (see
243 text related to Fig. 2d).



232 **Figure 4.** (a) Thermocline depth (Z_{tc}) (10-20°S). (b) Vertical potential temperature gradient and (c)
 233 Meridional current (10-20°S, 200-250°E). (d) Zonal current (8-13°S, upper 600 meter vertical mean).
 234 Datasets denoted as in Fig. 3(a,b).

244 Meridional current velocity in the region of 10-20°S and 200-250°E, is shown in Fig. 4c.
 245 Near surface poleward currents (outside the plotted range) are larger in the ridged models than
 246 in the dipped models or observations. The current changes direction from southward to north-
 247 ward around a depth of 40 meters in both groups of models and in reanalyses. The lower branch
 248 equatorward current in the ridged group is significantly weaker than the dipped group or ob-
 249 servations, consistent (through geostrophy) with the weaker thermocline deepening slope to
 250 its left (west) in Fig. 4a. This weakening is also consistent with current velocities becoming
 251 much smaller below the thermocline: the northward return current reaches half of its maxi-
 252 mal magnitude at a depth of ~ 140 meters in the ridged group, ~ 270 meters in the dipped group
 253 and ~ 230 meters in observations. The meridional velocity bias may therefore be interpreted
 254 as a shrinking of the lower branch of the southern subtropical cell (Schott et al., 2004), which
 255 in turn may indicate altered ocean energy transport (Held, 2001).

256 In observations, northward displacement of the ITCZ and resultant southerly cross-equatorial
 257 wind drive the eastward north equatorial counter current, situated between the ITCZ and the
 258 equator (Vallis, 2017); but the SH and equatorial zonal surface currents are westward. In con-
 259 trast with observations, ridged models have an eastward south equatorial counter current (SECC)

260 in the western Pacific between 8-13°S (Fig. 4d). Shoaling of the thermocline south of the SECC
 261 is expected, and thus so is a ridged thermocline. Consistent with Zhang & Song (2010), who
 262 found a similar spurious SECC in the NCAR CCSM3 model, this bias can be tied to a south-
 263 ward meridional wind stress bias (Fig. S5), associated with a southern ITCZ, through Ekman
 264 coupling.

265 **5 Discussion and Conclusions**

266 It is well known that the double ITCZ bias (DIB) has accompanying SST, surface wind
 267 and cloud radiative biases (Lin, 2007; Li & Xie, 2014; Zhang et al., 2019). However, to un-
 268 derstand the fully coupled nature of these biases requires an understanding of their three-dimensional
 269 oceanic expression. Here we show that a significant southern shallow thermocline bias is a promi-
 270 nent feature of the tropical Pacific problem, and is equally present in both CMIP5 and CMIP6
 271 models. The bias is manifested in both the physical thermocline (Z_{tc} , the depth of maximal
 272 vertical temperature gradient) and the commonly-used 20°C isotherm, but is significantly more
 273 pronounced in Z_{tc} . The bias is spatially coincident with precipitation biases, as evident in Fig. 1c,
 274 and its severity increases with the severity of the DIB (Fig. 3). Since water mass within the
 275 equatorial Pacific thermocline largely originates from the SH (Goodman et al., 2005), the south-
 276 ern Pacific shallow thermocline bias is a likely culprit of equatorial biases as well.

277 The key properties of the thermocline bias are:

- 278 1. A shallow bias between 5-30°S (Fig. 1d) and 160-260°E (Fig. 2a), up to 100 meters
 279 in ensemble mean.
- 280 2. Weak inter-hemispheric thermocline asymmetry ($A_{Z_{tc}} = -0.11$ in models compared to
 281 -0.32 in observations).
- 282 3. Thermocline and precipitation tropical asymmetry indices are strongly correlated ($R =$
 283 -0.81 , Fig. 3e).

284 Features unique to the most severely biased models are the:

- 285 1. Ridged thermocline south of the equator (Fig. 3a),
- 286 2. Lack of zonal thermocline sloping (Fig. 4a),
- 287 3. Wider and stronger double ITCZ (Fig. 3c),
- 288 4. Warmer southern off-equatorial SST (Fig. 3d),
- 289 5. Stronger cold tongue (Fig. 3d),

- 290 6. Reduced meridional return current (Fig. 4c), and
291 7. Spurious south equatorial counter current (SECC, Fig. 4d).

292 Extensive work done related to the DIB over two and a half decades since it was first
293 recognized by Mechoso et al. (1995), has treated it nearly exclusively in the annual mean. How-
294 ever, as Mechoso et al. (1995) and subsequent works pointed to, (e.g., De Szoeke & Xie, 2008;
295 Li & Xie, 2014; Adam, Schneider, & Brient, 2018) this bias has distinct seasonal aspects, which
296 peak in boreal spring. Thus, mechanisms that might explain the bias must account for its sea-
297 sonality. In particular, mixed layer depth, and to a large extent tropical thermocline depth, which
298 determine the thermal inertia of tropical waters on seasonal timescales, can modulate the sea-
299 sonal amplitude of ITCZ migrations. Indeed, slab ocean experiments showed that ITCZ sea-
300 sonal migration off the equator has a larger amplitude for shallower mixed layer depths (Dono-
301 hoe et al., 2014). Moreover, to be convinced of the importance of thermal inertia to ITCZ mi-
302 grations one need not look any further than the zonal variation in the extent to which the ITCZ
303 travels meridionally on seasonal time scales, in direct relation to the zonal sloping of the Pa-
304 cific thermocline. A smaller mass of water active in seasonal radiative heating and cooling will
305 result in low thermal inertia, leading to (1) SST more easily increasing above the threshold
306 required for deep convection, and (2) the latitude of maximal SST migrating further poleward.
307 In this sense, a shallow thermocline is as much a condition for deep convection as it is a re-
308 sult of it (Knauss & Garfield, 2016). This mechanism offers a seasonal link between the DIB
309 and the shallow thermocline bias, as it suggests that the summer hemisphere thermocline depth
310 is the relevant variable to setting seasonal ITCZ migration.

311 The SECC, found here to be associated with thermocline ridging, has been suggested
312 by Zhang & Song (2010) to be part of a positive feedback: the erroneous eastward current car-
313 ries warm water from the western warm pool eastward, heating southeastern SST and promot-
314 ing south-equatorial deep convection, a southern ITCZ and southward cross equatorial wind,
315 which further strengthens the SECC through Ekman coupling. This mechanism is consistent
316 with the SECC (Fig. 4d) and thermocline ridging (Fig. 4d) found here, concentrated in the west-
317 ern Pacific (Compare S6a and S6c). Interestingly, the spurious southern Ztc ridging is absent
318 in Z20 (see text related to Fig. 3b). Local shoaling of Z20 and Ztc can result from positive
319 wind stress curl, which leads to surface current divergence. Conversely, surface warming can
320 lead to opposite trends in Ztc and Z20: deepening of Z20 because of overall warming of the
321 upper ocean, but shoaling of Ztc, if the top part of the thermocline layer heats more than its

322 bottom part (Yang & Wang, 2009). This is the case, for example, in the western Pacific warm
323 pool, where Ztc shoals significantly but Z20 does not (compare Fig. 2a to Fig. S2a). In the
324 case of southern off-equatorial ridging, eastern Pacific Ztc shoals but does not ridge (Fig. S6a,b);
325 but both western Pacific Ztc and Z20 exhibit ridging (Fig. S6c,d). This may indicate that the
326 ridging biases originate from erroneous positive wind stress curl in the western Pacific, and
327 translate to eastern biases through the resultant SECC, which carries heat to the southeastern
328 Pacific.

329 Thermocline structure is influenced not only by local atmospheric forcing, but by inde-
330 pendent resolved and unresolved ocean processes. Representation of turbulent mixing in ocean
331 models in particular remains a major challenge (Guilyardi et al., 2009) and is critical to ther-
332 mocline properties, their relation to SST and the positive feedback between wind strength, SST
333 and the position of the ITCZ that maintains hemispheric asymmetry (S.-P. Xie & Philander,
334 1994). It follows that the thermocline bias might develop regardless of atmospheric biases. The
335 evidence presented in this paper begs the questions: What part of the bias is set by atmospheric
336 forcing? And to what extent does thermocline depth determine the tropical precipitation dis-
337 tribution?

338 One limitation of our results is the isolated treatment of the thermocline, without address-
339 ing salinity gradients and the resulting, possibly divergent, pycnocline. Though salinity gra-
340 dients are mild in most of the tropics, heavy precipitation areas, such as the western warm pool
341 can cause extreme surface freshening and a strong halocline, shallower than the thermocline.
342 This results in a shallow pycnocline above the thermocline and a barrier layer between the two
343 that further prevents mixing and entrainment from below (Sprintall & Tomczak, 1992; Breugem
344 et al., 2008). Salinity can thus affect both the thermocline itself and the importance we attribute
345 to it. Since we find significant precipitation biases in the thermocline bias region, and resul-
346 tant salinity biases are known to exist (Zhi et al., 2019), salinity may be important, and should
347 be addressed in future work. With that said, the thermal profile of the ocean column has im-
348 portant ramifications in and of itself, for mean ocean heat transport when combined with ve-
349 locity fields, and for the local heat budget, just to name two.

350 To conclude, a clear connection is found here between the shallow thermocline bias in
351 the southern Pacific and the well-known double ITCZ bias. Possible positive feedbacks among
352 these include:

- 353 1. A shallow thermocline reduces thermal inertia, causing the ITCZ to more easily develop
 354 south of the equator, which creates wind convergence and Ekman current divergence
 355 at the ITCZ, further shoaling the thermocline (Donohoe et al., 2014).
- 356 2. Wind anomalies associated with the southward position of the ITCZ in models lead to
 357 a spurious SECC and shoaling south of it. This, in turn, causes eastward heat advec-
 358 tion which reinforces the southward shift of the ITCZ and maintains a SECC and shoaled
 359 thermocline (Zhang & Song, 2010)

360 Ramifications of this bias to the fidelity of model predictions are potentially extensive.
 361 The dynamic and thermodynamic feedbacks mentioned here therefore merit further exploration.

362 **Acknowledgments**

363 The research was supported by the Israeli Science Foundation grant 1185/17. We thank Ita-
 364 mar Yacoby, Moran Erez and Eilat Elbaum for their helpful comments along the process. All
 365 of the data used in the analysis presented here is publicly available. We thank the climate mod-
 366 eling groups for producing and making available their model output, the Earth System Grid
 367 Federation (ESGF) for archiving the data and providing access, and the multiple funding agen-
 368 cies who support CMIP and ESGF. All CMIP data analyzed here are available from the ESGF
 369 at <https://esgf-node.llnl.gov/projects/esgf-llnl>.

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