

# Global structure of magnetotail reconnection revealed by mining space magnetometer data

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## Key Points:

- Global structure of magnetotail reconnection inferred from data mining matches its locations revealed by in-situ observations
- Reconstructed magnetotail reconnection structures include X- and O-lines, as well as magnetic nulls
- Reconstructed multiscale current sheet structure supports its formation mechanism by quasi-adiabatic ion motions

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**Abstract**

Reconnection in the magnetotail occurs along so-called X-lines, where magnetic field lines tear and detach from plasma on microscopic spatial scales (comparable to particle gyroradii). In 2017–2020 the Magnetospheric MultiScale (MMS) mission detected X-lines in the magnetotail enabling their investigation on local scales. However, the global structure and evolution of these X-lines, critical for understanding their formation and total energy conversion mechanisms, remained virtually unknown because of the intrinsically local nature of observations and the extreme sparsity of concurrent data. Here we show that mining a multi-mission archive of space magnetometer data collected over the last 25+ years and then fitting a magnetic field representation modeled using flexible basis-functions, faithfully reconstructs the global pattern of X-lines; 24 of the 26 modeled X-lines match ( $B_z = 0$  isocontours are within  $\sim 2$  Earth radii or  $R_E$ ) or nearly match ( $B_z = 2$  nT isocontours are within  $\sim 2R_E$ ) the locations of the MMS encountered reconnection sites.

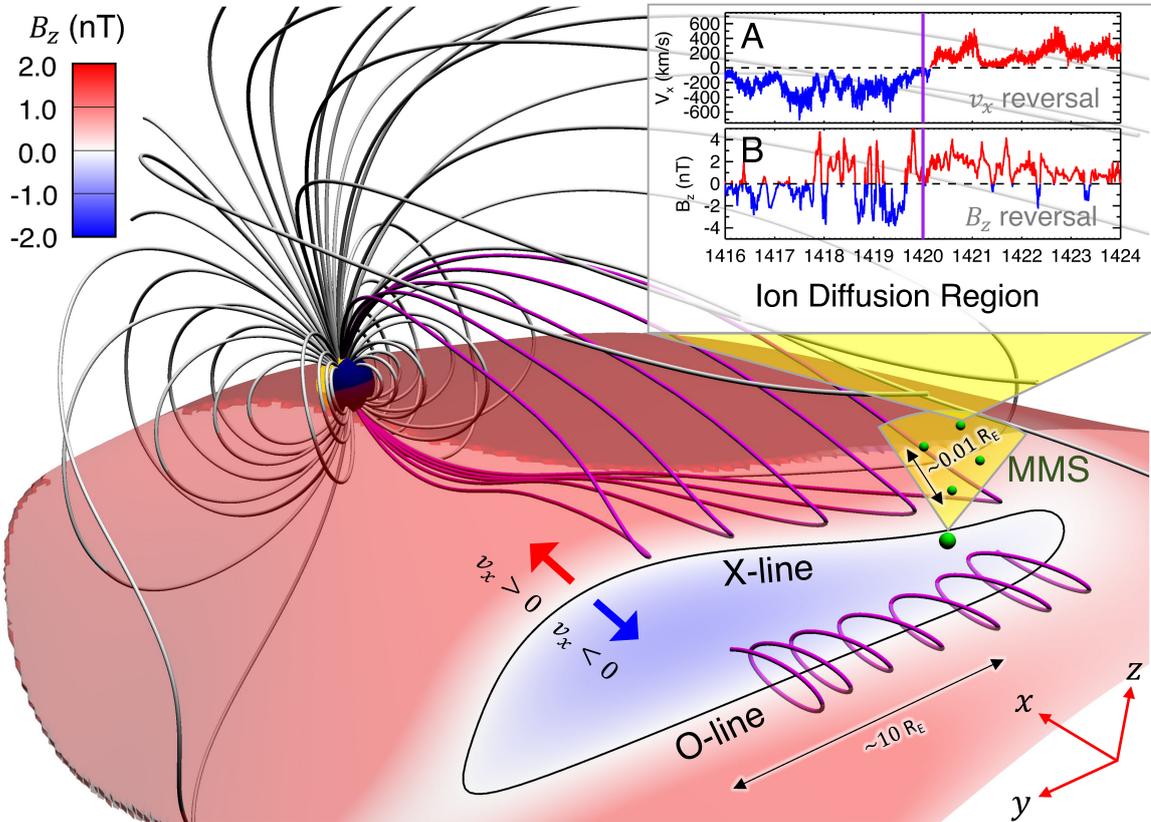
**Plain Language Summary**

Magnetic reconnection is a fundamental process in plasmas which couples microscopic scales ( $\sim$ electron to proton gyroradii) to explosive macroscopic phenomena many orders of magnitude larger, such as solar flares and geomagnetic storms/substorms. Reconnection forms along “X-lines”, rifts where oppositely directed magnetic field lines are forced together. In the Earth’s magnetosphere, reconnection has been observed by satellites at isolated locations; however, the large-scale structure of X-lines and their time evolution remains unknown because of the rarity and local nature of observations. Here, ground based measurements of geomagnetic activity and solar wind measurements are used to data-mine 25+ years of magnetometer data from 22 Earth-orbiting satellites, which are then utilized to reconstruct the global magnetic field associated with X-lines in Earth’s magnetosphere. We show that these reconstructions pinpoint the reconnection locations by verifying their consistency with direct spacecraft observations.

**1 Introduction**

X-lines are one of the most fundamental structures in magnetized plasmas, particularly in space, where they link global or even astronomical scale processes to those on the single particle orbit scale, thereby allowing those microscale processes to shape the universe (Ji et al., 2022). Dungey (1961) suggested that the interaction between Earth’s magnetic dipole and the solar wind causes reconnection of magnetic field lines on both the day and nightsides of Earth’s magnetosphere. The shape of these reconnecting field lines resembles the letter “X” and extends tens of Earth radii ( $R_E = 6,371.2$  kilometers) in the dawn-dusk direction thus forming X-lines (Figure 1). An X-line divides space into four sectors. In one pair of opposing sectors, the magnetic field and plasma converge towards the center of the X while in the other pair they are rapidly ejected from it. This reconnection process transforms energy stored in the magnetic field into particle kinetic and thermal energy, making it an efficient energy converter and particle accelerator (Ji et al., 2022). X-lines couple kinetic processes on proton and even electron gyroradius scales ( $\lesssim 0.01R_E$ ) (Torbert et al., 2018) to space weather phenomena on global scales: such as solar flares, coronal mass ejections, and magnetospheric storms and substorms ( $\sim 10R_E$ ) (Camporeale, 2019). This range of scales is so immense that its modeling has become one of the major challenges for nascent exascale computing (Ji et al., 2022).

While the microscale physics of reconnection in the magnetosphere has been studied in detail using recent multi-probe satellite missions (Angelopoulos et al., 2008; Burch et al., 2016; Burch et al., 2016; Torbert et al., 2018), its global structure is difficult to infer from data due to their paucity (rarity and locality): at any moment the huge volume of the magnetosphere ( $\gtrsim 10^5 R_E^3$ ) is probed by less than a dozen spacecraft (Sitnov et al., 2019). Understanding the global structure of reconnection is fundamental for determining substorm triggering mechanisms (Sitnov et al., 2019) and the total energy conversion during storms and substorms (Angelopoulos et al., 2013; Angelopoulos et al., 2020). Further, if X-line maps can be constructed from data, these maps could guide



**Figure 1.** 3D global picture of the magnetosphere and local MMS observations for 5 August 2020 (event Y in Table S2) in GSM coordinates. It shows that the data mining reconstructed X-line hits one of 26 ion diffusion region (IDR) encounters observed by the MMS mission during 2017–2020. It includes selected field lines and the color coded magnetic field distribution,  $B_z$ , sampled at the center of the tail current sheet taking into account deformation effects caused by the tilt angle of the Earth’s dipole axis as is detailed in the Supporting Information (SI). The  $B_z = 0$  isocontour is shown by the black line (the color table is saturated at  $|B_z| = 2$  nT to better reveal the isocontour). The inset shows key IDR parameters: (A) the proton bulk flow velocity component  $v_x$  and (B) the magnetic field  $B_z$ , from the MMS4 probe (the small green spheres show the MMS tetrahedral configuration) whose location is marked by the larger green sphere near the equatorial plane. The purple vertical line marks the reconstruction moment, 5 August 2020, 14:20 UT. The 3D visualizations are constructed using the VisIt visualization tool (Childs et al., 2012).

68 large-scale magnetohydrodynamic simulations of the magnetosphere by introducing a non-zero  
69 resistivity at their locations (Birn et al., 1996).

70 On the dayside, the X-line location can be readily estimated from the global geometry of  
71 the solar wind and Earth’s magnetic fields along with other well-defined physical parameters (Fuselier  
72 et al., 2011). In contrast, nightside reconnection is much less understood. Here, the solar wind-  
73 magnetosphere interaction stretches the dipole field lines in the antisunward direction forming  
74 the magnetotail while storing energy in the magnetic field. The release of this stored energy via  
75 reconnection is often unsteady and spontaneous. Observations of substorms (Russell & McPher-  
76 ron, 1973; Hones Jr., 1984; Baker et al., 1996; Angelopoulos et al., 2008, 2013) suggest that new  
77 X-lines form in the tail at distances of 10–30 $R_E$  and that this distance is controlled by the solar  
78 wind input (Nagai et al., 2005). However, despite decades of debate and being the target of ded-  
79 icated satellite missions (Nagai et al., 2005; Angelopoulos et al., 2008; Burch et al., 2016), the  
80 factors that determine the emergence, location, size, and shape of nightside X-lines remain a ma-  
81 jor mystery in heliophysics.

82 The recent four-probe Magnetospheric MultiScale (MMS) mission (Burch et al., 2016) en-  
83 abled microscopic analysis of magnetotail reconnection down to electron gyroradius scales (Torbert  
84 et al., 2018). During four years of MMS observations, 26 potential X-line encounters were found  
85 in the magnetotail (A. J. Rogers et al., 2019; A. Rogers et al., 2021), where explosive reconnect-  
86 ion causes substorms (Angelopoulos et al., 2008; Angelopoulos et al., 2020; Sitnov et al., 2019).  
87 They were detected in the form of Ion Diffusion Regions (IDRs) characterized by reversals of  
88 the North-South component of the magnetic field,  $B_z$ , and of the Sun-Earth component of the pro-  
89 ton bulk flow velocity,  $v_x$ , (Fig. 1 inset).

90 In this study, the global structure of magnetotail reconnection is derived from a large set  
91 of historic satellite magnetometer measurements using an advanced data mining approach. We  
92 show that our technique provides evidence justifying the global reconnection structure: the ob-  
93 tained contours delineating  $B_z$  reversals pass through most of the micro-scale IDRs discovered  
94 by MMS. We further discuss implications of the obtained magnetotail picture to the multiscale  
95 structure of its current sheet, and then describe its uncertainty and in-situ validation errors. Through-  
96 out this study, vector quantities are represented in the Geocentric Solar Magnetospheric System  
97 (GSM).

## 98 2 Data Mining Solution of the Data Paucity Problem

99 The key to solving the data paucity problem lies in the recurrent nature and repeatable pat-  
100 tern of storms and substorms. The storm recurrence time for medium intensity storms is approx-  
101 imately two weeks (Reyes et al., 2021), while it is 2–4 h for periodic substorm (Borovsky & Yaky-  
102 menko, 2017). This repeatability allows the magnetic field to be reconstructed not only from ob-  
103 servations at the moment of interest but also from records identified via mining the space mag-  
104 netometer archive by searching for other times when the magnetosphere was in a similar global  
105 state. The magnetospheric state is characterized using geomagnetic indices (metrics of magnetic  
106 activity derived from networks of ground magnetometers) and solar wind conditions. Specifi-  
107 cally, the magnetospheric state is defined using a 5-D state-space vector,  $\mathbf{G}(t) = (G_1, \dots, G_5)$ , formed  
108 from the geomagnetic storm index ( $SMRc$ ), substorm index ( $SML$ ), their time derivatives, and the  
109 solar wind electric field parameter ( $vB_z^{IMF}$ ; where  $v$  is the solar wind speed and  $B_z^{IMF}$  is the North-  
110 South component of the Interplanetary Magnetic Field, IMF). The  $SMR$  and  $SML$  ( $SMRc$  is a pressure-  
111 corrected  $SMR$  (Tsyganenko et al., 2021)) indices are provided by the SuperMag project (Gjerloev,  
112 2012) and represent variations of the ground-based magnetometer records from low/mid- and high-  
113 latitude stations respectively analogous to the  $Sym-H$  and  $AL$  indices used before (Sitnov et al.,  
114 2008; G. K. Stephens et al., 2019).  $G_{1-5}(t)$  are normalized by their standard deviations, smoothed  
115 over storm or substorm scales, and sampled at a 5-min cadence, as is detailed in (G. K. Stephens  
116 & Sitnov, 2021) and in the Supporting Information (SI). Including the time derivatives of these  
117 activity indices allows the data mining (DM) procedure to differentiate between storm and sub-  
118 storms phases as well as capturing memory effects of the magnetosphere as a dynamic system

119 (Sitnov et al., 2001). The space magnetometer archive contains data from 22 satellites (includ-  
 120 ing the four MMS probes) spanning the years 1995–2020 resulting in 8,649,672 magnetic field  
 121 measurements after being averaged over 5 and 15 min time windows as is further described in  
 122 Figure S1 and Table S1 of the SI.

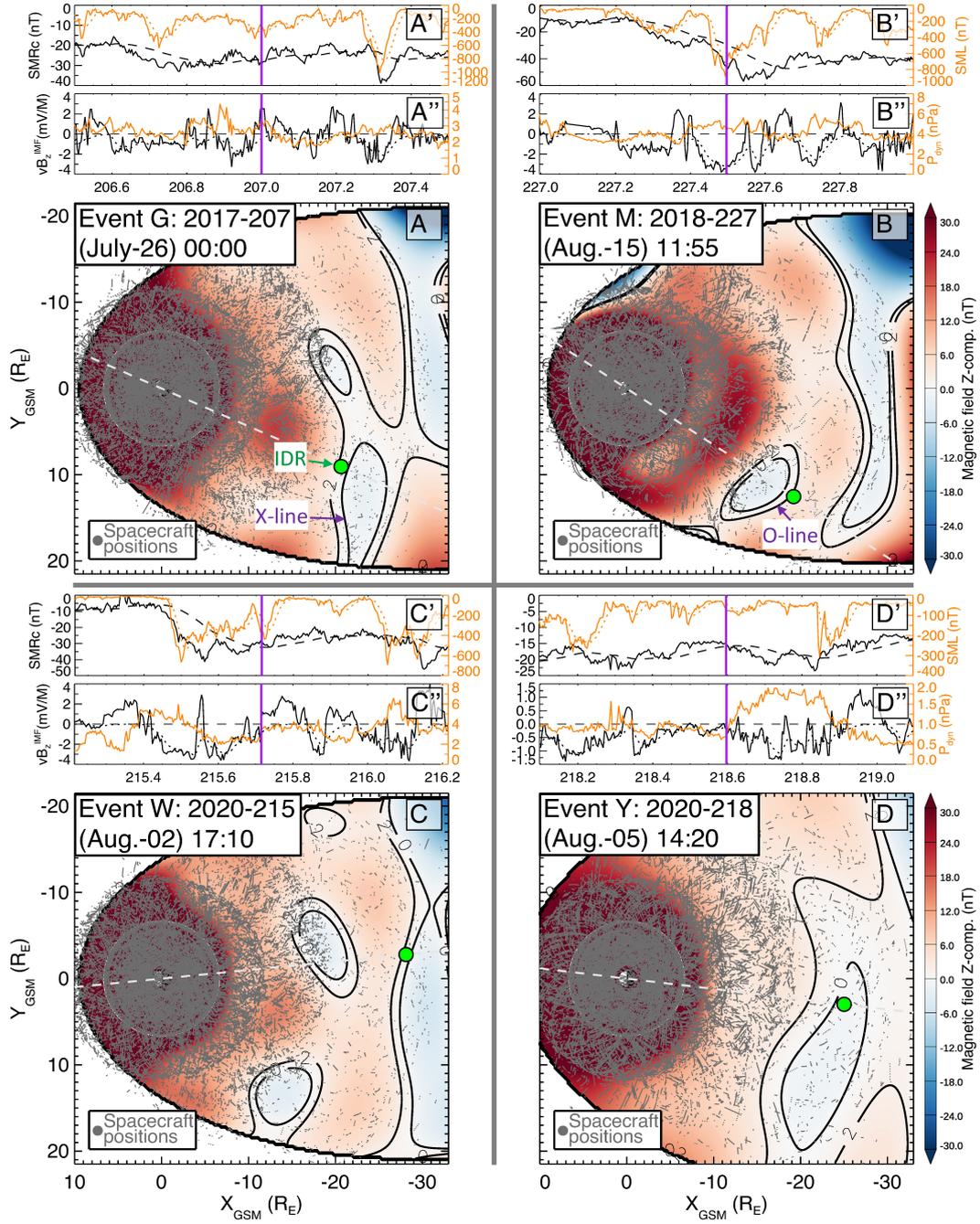
123 The DM algorithm employed is based on the k-nearest neighbor (kNN) classifier method  
 124 (Wettschereck et al., 1997; Sitnov et al., 2008). To illustrate the algorithm, assume the magnetic  
 125 field reconstruction,  $\mathbf{B}(t)$ , is sought for a query time  $t = t^{(q)}$ . This corresponds to a particular point  
 126 in the 5-D state-space,  $\mathbf{G}^{(q)} = \mathbf{G}(t^{(q)})$ . Surrounding this point will be other points,  $\mathbf{G}^{(i)}$ , in close  
 127 proximity to it; i.e., its nearest neighbors (NNs). Distances between points in state-space are com-  
 128 puted using the Euclidean metric:  $R_i = |\mathbf{G}^{(i)} - \mathbf{G}^{(q)}|$ . Time-adjacent NNs form intervals in time  
 129 that identify a subset of the magnetometer database used to fit the analytical formulation of the  
 130 magnetic field, yielding  $\mathbf{B}(t^{(q)})$ . The specific choice of the number of NNs to use in the recon-  
 131 struction,  $k_{NN}$ , is dictated by a balance between over- and under-fitting. G. K. Stephens and Sit-  
 132 nov (2021) found the optimal number to be  $k_{NN} = 32,000$  for tail reconstructions of substorms,  
 133 corresponding to  $\sim 1\%$  of the total database. The resulting subset is composed of a very small  
 134 number ( $\sim 1$ – $10$ ) of real (from the event of interest) but many ( $\sim 10^5$ ) virtual (from other events)  
 135 satellites.

136 The large number of virtual points enables new magnetic field architectures (Tsyganenko  
 137 & Sitnov, 2007; G. K. Stephens et al., 2019), which differ from classical empirical models with  
 138 custom-tailored modules (e.g., Tsyganenko & Sitnov, 2005) by utilizing regular basis function  
 139 expansions for the major magnetospheric current systems, to be used for the reconstructions. In  
 140 particular, all near-equatorial currents are approximated by two expansions representing general  
 141 current distributions of thick and thin current sheets with different thickness parameters  $D$  and  
 142  $D_{TCS}$ . The latter accounts for the formation of ion-scale thin current sheets (TCS) prior to sub-  
 143 storm onset (V. Sergeev et al., 2011), as is further detailed below. The independence of the cur-  
 144 rent sheet expansions is provided by the constraint  $D_{TCS} < D$ . To improve the reconstructions,  
 145 while fitting the magnetic field model with the NN subset, the spacecraft data were additionally  
 146 weighted: in the real space (to mitigate the inhomogeneity of their radial distribution (Tsyganenko  
 147 & Sitnov, 2007)) and in the state-space (to reduce the uncertainty and bias toward weaker activ-  
 148 ity regions (G. K. Stephens & Sitnov, 2021)) as it is further detailed in the SI.

149 The resulting DM reconstruction of the magnetic field during the early expansion phase  
 150 of the 5 August 2020 substorm (Figure 1) reveals the formation of an X-line at  $r \approx 23R_E$  in the  
 151 tail. This data-derived image of the X-line resembles sketches of solar flare arcades (e.g., Shiota  
 152 et al., 2005) but with a fundamental advantage that it is backed by a quantitative description. The  
 153 X-line appears on the dusk flank of the tail illustrated as the earthward part of the  $B_z = 0$  isocon-  
 154 tour in the equatorial plane (black line). It also corresponds to an earthward edge of a relatively  
 155 long ( $25R_E$ ) spiral structure, shown by the sample field lines that encircle the tailward part of the  
 156  $B_z = 0$  isocontour and form a magnetic O-line. The large green sphere in Figure 1 indicates the  
 157 location of the MMS satellites at this moment with the inset, which shows the observed  $B_z$  and  
 158  $v_x$ , demonstrating it was one of the fortuitous IDR encounters.

### 159 3 Ion Diffusion Regions

160 The main goal of the MMS mission (Burch et al., 2016) was the detection and investiga-  
 161 tion of reconnection regions in the magnetosphere and its boundary. That goal was relatively easy  
 162 to achieve at the magnetopause because of its regular structure (Fuselier et al., 2011) and in the  
 163 magnetosheath due to multiple reconnection sites in its turbulent plasma volume (Phan et al., 2018).  
 164 By contrast, only a handful of fortunate X-line encounters were detected/investigated in the mag-  
 165 netotail (Torbert et al., 2018; Chen et al., 2019). In this regard, the proposed DM reconstructions  
 166 offer an attractive opportunity to explore the dynamics of magnetotail topology on a global scale,  
 167 and its fidelity can be demonstrated by comparing our results with MMS observations. Magnetic  
 168 reconnection can be directly observed if and when a spacecraft fortuitously flies through an Ion  
 169 Diffusion Region (IDR), as shown in Figure 1. A recent systematic survey of MMS plasma and



**Figure 2.** IDRs and the equatorial magnetic field landscape. (A–D) Color-coded distributions of the equatorial magnetic field,  $B_z$ , with  $B_z = 0$  and 2 nT isocontours (black lines), big green dots pointing to the IDRs, and gray dots showing the spacecraft positions for the NN subsets used in the DM method for four IDR events, G, M, W and Y. Panels on top of each equatorial  $B_z$  distribution show the global context of the considered events in terms of (A'–D') the storm and substorm indices  $SMRc$  (black),  $SML$  (orange), and (A''–D'') the solar wind/IMF parameters  $vB_z^{IMF}$  (black) and  $P_{dyn}$  (orange) with the purple vertical line marking the event time.

170 field data in 2017 (A. J. Rogers et al., 2019) identified 12 such magnetotail IDRs, defined as cor-  
 171 related reversals of the proton bulk flow velocity,  $v_x$ , and the North-South magnetic field,  $B_z$ , as  
 172 shown in the Figure 1 inset, along with additional Hall magnetic and electric field signatures. That  
 173 analysis was later extended to 2018–2020 for a total of 26 IDR events (A. Rogers et al., 2021)  
 174 labeled here A–Z, “IDR alphabet”, listed in Table S2 in the SI.

175 Figure 1 shows that the DM reconstruction correctly identifies one of the IDR regions, namely  
 176 event Y (5 August 2020 14:20 UT), whose  $v_x$  and  $B_z$  reversals are shown in the inset. The pro-  
 177 jection of the magnetic field at the center of the tail current sheet into the equatorial plane is dis-  
 178 played in Figure 2D showing that the  $B_z = 0$  contour passes within  $\sim 1R_E$  of the IDR observed  
 179 by MMS. This success is remarkable given that only  $\sim 0.03\%$  of the measurements used to re-  
 180 construct the magnetic field were taken from this event, with the other 99.97% coming from other  
 181 similar events identified using the above described DM approach. Below we show that similar  
 182  $B_z = 0$  contours pass through most of the microscale IDRs discovered by MMS.

### 183 3.1 Reconstructed X- and O-lines in the Equatorial Plane

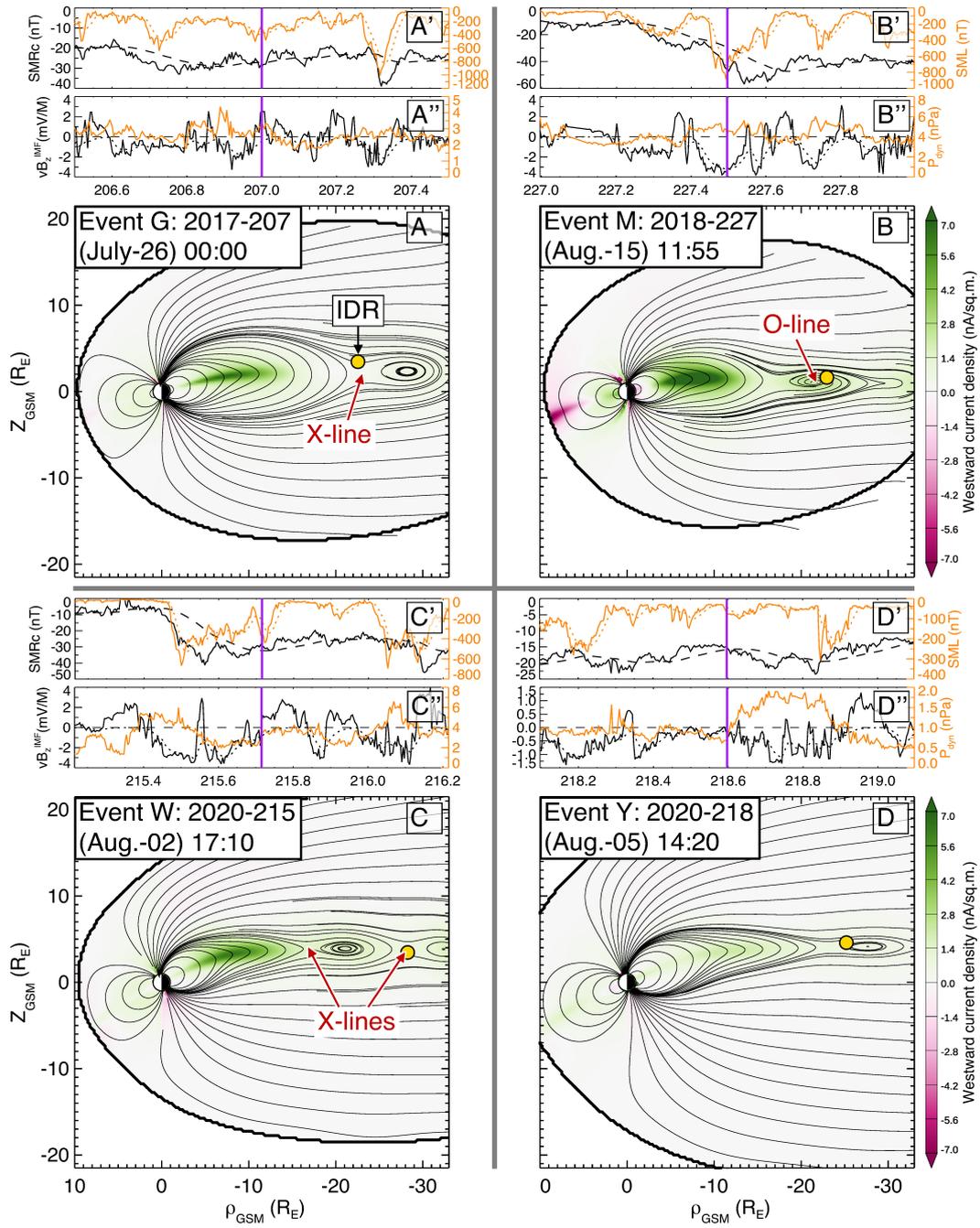
184 The reconstructions of 3 other events (G, M, W) presented in Figure 2 also show the  $B_z =$   
 185 0 contours pass within  $\sim 1R_E$  of the observed IDRs (the exact distances are provided in Table S2).  
 186 Closer examination shows that only events G, W, and Y are X-lines, whereas event M corresponds  
 187 to an O-line. Indeed, since the microscale formation of the MMS tetrahedron cannot determine  
 188 X-line motions using timing analysis, (e.g., Eastwood et al., 2010), or by framing the X-lines by  
 189 being tailward and earthward of them (Angelopoulos et al., 2008), it cannot distinguish whether  
 190 they are X- or O-lines.

191 The equatorial X-line reconstructions for the remaining 22 IDR events are provided in Figs. S3–  
 192 S8. They reveal that the DM approach hits 16 of the 26 IDRs (Figs. 2 and S3–S5), that is, the  $B_z =$   
 193 0 contours pass within  $\lesssim 2R_E$  of the IDRs, which includes 11 X-lines and 5 O-lines. Another  
 194 8 events are near hits, that is, the IDRs are located within  $\lesssim 2R_E$  from  $B_z = 2$  nT contours (Figs. S6  
 195 and S7). Only in 2 events (B, F) do the reconstructed  $B_z = 0$  contours miss their IDR targets (Fig-  
 196 ure S8); however, both events have a plausible explanation. Event B occurs during weak mag-  
 197 netospheric activity ( $SML \approx 0$ ) with effectively no solar wind/IMF input ( $vB_z^{IMF} > 0$ ) while event F  
 198 takes place during the middle of a several hours long gap in solar wind and IMF data (they are  
 199 interpolated in the reconstruction).

### 200 3.2 Reconnection Features in the Meridional Planes

201 The corresponding meridional slices through the planes containing the IDRs of the Fig-  
 202 ure 2 events (G, M, W, Y) are shown in Figure 3, illustrating the magnetic topology and distri-  
 203 butions of electric currents, while the remainder of the IDR alphabet (Figures S3–S8) is shown  
 204 in Figures S9–S14. The figures clarify that the observed  $B_z = 0$  contours indeed represent X-  
 205 and O-lines similar to the 3D magnetotail field geometry shown in Figure 1. They also confirm  
 206 the quasi-2-D nature of reconnection apparently imposed by the North-South symmetry of the  
 207 magnetotail (e.g., Tsyganenko & Fairfield, 2004) which is drastically different from the inher-  
 208 ently 3-D reconnection processes in the solar corona (Liu et al., 2016) and rapidly rotating plan-  
 209 ets (Griton et al., 2018).

210 These meridional distributions resemble empirical visualizations of reconnection in lab-  
 211 oratory plasmas, which became possible due to their large number of real probes (up to 200) and  
 212 additional symmetry constraints, such as the cylindrical symmetry imposed by the toroidal-shaped  
 213 flux cores in the PPPL Magnetic Reconnection Experiment (MRX) (Ji et al., 2022). Still, in con-  
 214 trast to MRX, magnetotail reconnection is only quasi 2-D due to the finite length of the X-line  
 215 forming a closed loop with the O-line, as well as the explicit 3-D effects, such as null-points (e.g.,  
 216 Greene, 1988; Ji et al., 2022). Null-points in the tail were indeed inferred from the four-probe  
 217 Cluster observations (Xiao et al., 2006). An example of the null-point pair seen in our DM re-  
 218 construction of event Y is presented in Figure S15. Additional deviations from the simple 2-D



**Figure 3.** IDRs against the meridional current and magnetic field distributions. (A–D) Color-coded distribution of the electric current perpendicular (westward positive) to the meridional plane, which contains the corresponding IDR (white dashed lines in Figure 2), for four events shown in Figure 2 with the similar format for global parameters (A'–D') and (A''–D'') on top of each distribution. The IDRs are shown here by big orange dots. Thin and thick lines show the magnetic field lines and the magnetospheric boundary (magnetopause).

219 picture could be due to a strong IMF  $B_y$  (e.g., Cowley, 1981) or North-South oscillations of the  
 220 tail current sheet that resemble a flapping flag (e.g., V. A. Sergeev et al., 2006; Sitnov et al., 2019).

221 Note that in the present reconstructions the original multiscale tail model (G. K. Stephens  
 222 et al., 2019) with the embedded TCS structure has been further generalized to verify the possi-  
 223 ble physical mechanisms of the TCS formation. It can be explained, (e.g., Sitnov et al., 2006),  
 224 by figure-eight like Speiser (1965) proton orbits. If this is the case, the parameter  $D_{TCS}$  of the mag-  
 225 netic field model should depend on the distance  $\rho$  from the Earth because the Speiser orbit size,  
 226  $\rho_{Si}$ , is inversely proportional to the magnetic field outside the sheet,  $B_L$ , which itself depends on  
 227  $\rho$  (Wang et al., 2004). To take this effect into account the magnetic field architecture was further  
 228 generalized using the approximation  $D_{TCS}(\rho) = [D_*^{-1} + \alpha \exp(-\beta\rho)]^{-1}$  with free parameters  
 229  $\alpha$ ,  $\beta$  and  $D_*$  to be inferred from data. The fitting details provided in Figure S2 suggest that the  
 230 scaling  $D_{TCS} \propto B_L^{-1} \propto \rho_{Si}$  does indeed take place, which supports the theoretical mechanism of  
 231 the TCS formation related to the Speiser orbits.

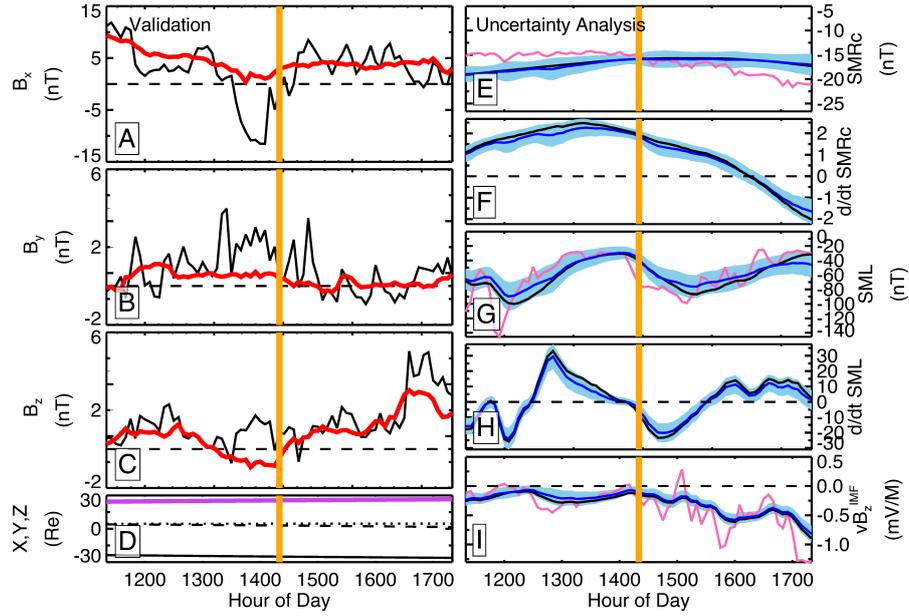
## 232 4 Validation and Uncertainty Quantification

233 An example of in-situ validation of this global reconstruction, rarely possible because of  
 234 the data paucity, is shown in Figs. 4A–4D for the MMS magnetic field observations of the tail  
 235 during event Y. It reveals relatively large deviations in the magnetic field components  $B_{x,y}$  par-  
 236 allel to the current plane (Figs. 4A, 4B). They are likely caused by the flapping North-South mo-  
 237 tions of the current sheet as a whole (V. A. Sergeev et al., 2006) that were found in MMS obser-  
 238 vations as well (Farrugia et al., 2021). These motions are spontaneous and may appear in differ-  
 239 ent phases of activity, so it is not surprising that they are not captured by the DM reconstructions.  
 240 At the same time, the  $B_z$  magnetic field is reproduced even better than it appears in observations  
 241 after 5-min averages (compare the black line in Figure 4C with the inset in Figure 1). Thus, hit-  
 242 ting 24 out of 26 IDRs, achieved in this study, shows (i) how to overcome the curse of data paucity  
 243 for in-situ data and (ii) presents solid evidence that not only validates our DM reconstructions,  
 244 but also helps understand the reconnection mechanisms and its consequences.

245 The fidelity of the present reconstructions can also be seen from the uncertainty analysis  
 246 presented in Figs. 4E–I. It compares 5 original binning parameters of the magnetosphere with  
 247 their means and standard deviations over the NN bins. The closeness of means to the original pa-  
 248 rameters  $G_{1-5}$  and small relative values of deviations suggest that the selected NNs closely fol-  
 249 low the magnetospheric dynamics, especially on substorm scales (Figs. 4G–4H).

## 250 5 Conclusions

251 This picture of the 2017–2020 MMS IDR alphabet suggests that, in spite of the extreme  
 252 paucity of in-situ observations, DM successfully reconstructs the overall structure of magneto-  
 253 tail X- and O-lines because they are strongly self-organized on the global scale. The X-lines vary  
 254 in length from 5 to  $40R_E$ , with the shorter ones forming inside of  $\sim 20R_E$  while the longer ones,  
 255  $\sim 40R_E$ , appear beyond  $25R_E$ . The concurrent appearance of such near-Earth and midtail X-lines  
 256 is consistent with the original conjectures regarding new X-line formation during substorms (Hones Jr.,  
 257 1984). It also explains the detection of X-lines as discrete points in radial distance in remote sens-  
 258 ing (Angelopoulos et al., 2013, Fig. 3C) as well as the stepwise retreat of magnetic reconnection  
 259 regions suggested by their auroral manifestations and confirmed by in-situ observations (Ieda et  
 260 al., 2016). The persistent formation of X-lines near  $30R_E$  has also been confirmed by the statisti-  
 261 cal analysis of the travelling compression regions (Imber et al., 2011). The success of our X-  
 262 line reconstruction indicates that year after year, the spatial/temporal patterns of storms and sub-  
 263 storms in the Earth’s magnetotail are highly recurrent and hence reproducible with historic data,  
 264 while magnetic reconnection controls the global state of the magnetosphere reflected in its ac-  
 265 tivity indices, their trends, and the solar wind energy input.



**Figure 4.** Validation and uncertainty analysis for event Y. (A)–(C) The 5-min averaged GSM magnetic field components (black lines) and their DM reconstructions (red lines). (D) MMS ephemeris (in GSM) X (solid line), Y (dashed line), Z (dash-dotted line) and the radial distance (pink line). (E)–(I) The storm/substorm state binning parameters  $\langle SMRc \rangle$ ,  $D\langle SMRc \rangle/Dt$ ,  $\langle SML \rangle$ ,  $D\langle SML \rangle/Dt$ , and  $\langle vB_z^{IMF} \rangle$  as described in the SI, shown by black lines as compared to their means over the NNs (blue lines). The light blue shading shows the standard deviations  $\pm 1\sigma$  of the NNs. Pink lines in Figs. 4E, 4G, and 4I show the original 5-min OMNI data for the parameters  $SMRc$  (pressure-corrected  $SMR$  (Tsyganenko et al., 2021)),  $SML$ , and  $vB_z^{IMF}$ . Yellow bars show the moment of the spatial reconstruction 5 August 2020, 14:20 UT shown in Figs. 1, 2D and 3D.

## 266 6 Open Research

267 The data used in the paper are archived on Zenodo (G. Stephens et al., 2022). For each of  
 268 the 26 IDR events, files are included that detail: time intervals identified using the nearest-neighbor  
 269 search and the resulting subset of magnetometer data and their associated weights, files contain-  
 270 ing the fit set of coefficients and parameters for the model, and the digital model output data that  
 271 were used in constructing the figures. The compiled magnetometer database used in this study  
 272 is available on the SPDF website (Korth et al., 2018). This study extended this database with the  
 273 addition of MMS magnetometer data which has also been included in the Zenodo archive. The  
 274 SMR and SML indices obtained from the SuperMAG web page are also included in the Zenodo  
 275 archive. The data describing the solar wind conditions were taken from the 5-min OMNI data  
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