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Investigation of Ionospheric Small-Scale Plasma Structures associated with Particle Precipitation

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Key Points:

- Enhanced values of the GNSS-based 1-second amplitude scintillation index are observed at auroral intensity gradients.
- The IFLC is often elevated at the same time as the S4 index, confirming the diffractive nature of observed scintillation events.
- Enhanced ionospheric currents due to auroral particle precipitation contribute to structuring below the Fresnel scale.

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Abstract

We investigate the role of auroral particle precipitation in small-scale (below hundreds of meters) plasma structuring in the auroral ionosphere over the Arctic. To the scope, we together analyse data recorded by an Ionospheric Scintillation Monitor Receiver (ISMR) of Global Navigation Satellite System (GNSS) signals and by an All-Sky Camera located in Longyearbyen, Svalbard (Norway). We leverage on the raw GNSS samples provided at 50 Hz by the ISMR to evaluate amplitude and phase scintillation indices at 1 s time resolution and the Ionosphere-Free Linear Combination at 20 ms time resolution. The simultaneous use of the 1 s GNSS-based scintillation indices allows identifying the scale size of the irregularities involved in plasma structuring in the range of small (up to few hundreds of meters) and medium-scale size ranges (up to few kilometers) for GNSS frequencies and observational geometry. Additionally, they allow identifying the diffractive and refractive nature of the found fluctuations on the recorded GNSS signals. Six strong auroral events and their effects on plasma structuring are studied. Plasma structuring down to scales of hundreds of meters are seen when strong gradients in auroral emissions at 557.7 nm cross the line of sight between the GNSS satellite and receiver. Local magnetic field measurements confirm small-scale structuring processes coinciding with intensification of ionospheric currents. Since 557.7 nm emissions primarily originate from the ionospheric E-region, plasma instabilities from particle precipitation at E-region altitudes are considered to be responsible for the signatures of small-scale plasma structuring highlighted in the GNSS scintillation data.

1 Introduction

Particle precipitation into the E and F-region ionosphere can cause plasma irregularities and structuring on various scales (Moen et al., 2013; Zou et al., 2013; van der Meeren et al., 2015; De Franceschi et al., 2019; Jin et al., 2016; Dahlgren et al., 2017; Fæhn Follestad et al., 2020; Urbar et al., 2022; Alfonsi et al., 2022; Enengl et al., 2023). As radio waves propagate through the ionosphere they undergo diffraction and refraction (Keskinen & Ossakow, 1983; Huba et al., 1985; Kintner & Seyler, 1985). This results in rapid fluctuations of the phase and amplitude as observed by ground-based receivers (Hey et al., 1946; Kintner et al., 2007). Irregularities below the Fresnel's scale for Global Navigation Satellite System (GNSS) frequencies and observational geometry (order of few hundreds of meters) induce diffractive effects on both the amplitude and the phase of the signals recorded at ground. In fact, those small-scale irregularities act as a new wave source for the GNSS planar waves and cause interference at the receiver level, which is the primary source of scintillation (Basu et al., 1998; Wernik et al., 2003; Kintner et al., 2007). The recent literature highlighted that the fluctuations of this kind are the only one deemed to be called "scintillation", because of the stochastic nature of the induced fluctuations (McCaffrey & Jayachandran, 2019; Ghobadi et al., 2020). Besides scintillation, fluctuations on the GNSS phase measurements are induced by irregularities covering the full ranges of scales. These fluctuations are deterministic in nature, due to the dispersive nature of the ionospheric medium, and covers the full range of the ionospheric irregularities scale sizes (McCaffrey & Jayachandran, 2019). Most of the recent research about GNSS scintillation has been dedicated to efficiently evaluate the diffractive and refractive contribution in the recorded GNSS fluctuations, especially in the high-latitude ionosphere, where the highly dynamic plasma complicates such a task (McCaffrey & Jayachandran, 2019; Ghobadi et al., 2020; Spogli et al., 2021; Wang et al., 2018, 2022; Conroy et al., 2022). In this study, we investigate the role of the particle precipitation in small-scale plasma structuring in the E and F-region ionosphere and its effects on trans-ionospheric radio wave propagation over Svalbard (Norway).

One visual signature of ionosphere-magnetosphere coupling through Birkeland currents on Earth is the aurora. Using all-sky imagers with filters for the green and red auroral emission lines gives an indication of the location and altitude of the particle pre-

70 precipitation energy deposition. As the geomagnetic activity grows, particle precipitation
 71 leads to brighter aurora. Auroral emissions indicate particle precipitation into the E and
 72 F-region altitudes; green (557.7 nm) emissions correspond primarily to E-region altitudes,
 73 and the red (630.0 nm) emissions to F-region altitudes. The influx of energetic particles
 74 leads to increased plasma density and enhanced conductance in the ionosphere. The iono-
 75 spheric plasma is then subject to electric field effects, density irregularities and parti-
 76 cle streams. Instabilities associated with particle precipitation are kinetic instabilities,
 77 such as the two-stream instabilities or current-driven instabilities (Kropotkin, 2016). Neu-
 78 trals are also present in the E-region ionosphere, which provides conditions for the Farley-
 79 Buneman instability that arises from the difference between the electron and ion veloci-
 80 ty, caused by collisions of the ions with neutrals (Farley Jr., 1963; Buneman, 1963; Ro-
 81 gister & D'Angelo, 1970; Treumann, 1997). Given the impacts of particle streams into
 82 ionospheric plasma conditions, particle precipitation is proposed to contribute to iono-
 83 spheric small-scale plasma structuring (Greenwald et al., 2002).

84 Ionospheric plasma structuring can be indirectly observed by Ionospheric Scintil-
 85 lation Monitor Receivers (ISMRs). The above mentioned amplitude and phase fluctu-
 86 ations are usually studied through the phase (σ_ϕ) and amplitude (S4) scintillation in-
 87 dices. As anticipated above, phase scintillations are driven by ionospheric irregularities
 88 at small wave numbers and below the first Fresnel's radius for GNSS frequencies and ob-
 89 servational geometry. Phase scintillation below the first Fresnel radius is the result of
 90 interference between different phases and can be thought of as diffractive, while phase
 91 scintillation at small wave numbers is considered refractive and produced by fluctuations
 92 of plasma density integrated along the signal path (Kintner et al., 2007). However, the
 93 overall phase fluctuations measured at ground can be split into a refractive and diffrac-
 94 tive contribution. The refractive part of the signal is deterministic, as it can be related
 95 to the electron density and wave frequency and can be accounted for. The diffractive part
 96 is stochastic (Kintner et al., 2007; McCaffrey & Jayachandran, 2019). Amplitude scin-
 97 tillations are related to diffractive effects, as the refraction impacts the wave path and
 98 power but will not cause amplitude fluctuations (McCaffrey & Jayachandran, 2019). Re-
 99 fraction is observed for small and large scale structures, while diffraction occurs for struc-
 100 tures equal to or below the Fresnel's scale (Kintner et al., 2007; Zheng et al., 2022). The
 101 first order refractive contribution to GNSS phase fluctuations can be accounted for by
 102 calculating the Ionosphere-Free Linear Combination (IFLC) (Carrano et al., 2013). The
 103 IFLC is a combination of the carrier phases of two received waves at different frequen-
 104 cies able to account for the bulk of the refractive contribution induced by the ionosphere
 105 on the phase measurements (McCaffrey & Jayachandran, 2019). The use of the IFLC
 106 has been recently adopted to improve the phase detrending scheme adopted to evalu-
 107 ate the phase scintillation index. In fact, phase detrending is a delicate issue for ISMRs
 108 located in the high-latitude regions (Forte, 2005; Beach, 2006), for which the standard
 109 detrending scheme based on a 6th-order Butterworth high-pass filter with a cutoff fre-
 110 quency of 0.1 Hz is not effective in eliminating the bulk of the refractive contributions
 111 above the Fresnel's scale. To address this issue, the Fast Iterative Filtering technique (Cicone,
 112 2020; Cicone & Zhou, 2021) was recently used to determine a more suitable and adap-
 113 tive detrending scheme, with related cutoff frequency, to provide a phase scintillation in-
 114 dex in which the bulk of the refractive part is eliminated (Ghobadi et al., 2020; Spogli
 115 et al., 2021). Bearing this in mind, it depends on the purpose and whether one is inter-
 116 ested in the refractive contributions as a proxy for plasma structuring processes to de-
 117 cide on the choice of indices and on the related detrending scheme for GNSS phase mea-
 118 surements. Investigating the IFLC and scintillation indices along with the temporal and
 119 spatial evolution of particle precipitation, may improve our understanding of when/where
 120 small-scale structures are present, their role in plasma structuring processes and what
 121 their nature is.

122 Previous studies have already laid a pathway to understand the occurrence of iono-
 123 spheric irregularities in relation to particle precipitation and their impact on trans-ionospheric

124 radio wave propagation at high latitudes. The control of the interplanetary magnetic field
 125 (IMF) extends to small-scale irregularities of plasma density associated with the large-
 126 scale structures in the ionosphere. IMF Bz northward conditions showed moderate lev-
 127 els of amplitude and phase scintillation were observed with highly variable decorrelation
 128 times in a study by Basu et al. (1991). The relationship between auroral particle pre-
 129 cipitation, electric fields and ionospheric irregularities is a result of magnetosphere-ionosphere
 130 coupling. It confirms the dependence of ionospheric plasma instabilities on electric fields
 131 and precipitation-induced electron density gradients (Greenwald et al., 2002). Meso-scale
 132 electron density irregularities on the dayside auroral region are found pole-ward of the
 133 nominal cusp region (Basu et al., 1998). Irregularities on the pole-ward side of the au-
 134 rora are predominantly smaller than the Fresnel scale (observed as amplitude scintilla-
 135 tion) (Conroy et al., 2022). Plasma structuring as driven by particle precipitation is pro-
 136 posed to be driven at the boundaries of the auroral precipitation (such as Kelvin-Helmholtz
 137 and/or Farley-Buneman) down to E-region altitudes and to play a main role in plasma
 138 structuring on various scales Enengl et al. (2023). Pole-ward moving auroral forms can
 139 cause strong ionospheric irregularities, capable of causing more severe disturbances in
 140 the cusp ionosphere for navigation signals than polar cap patches (Oksavik et al., 2015).
 141 Loss of lock events consistently appeared consistently at edges of a auroral form named
 142 westward travelling surge in a study by Semeter et al. (2017). The authors concluded
 143 the E-region (near the oxygen emissions) was the source of the irregularities.

144 Moving auroral structures in the E-region occurred simultaneously with the Global
 145 Positioning System (GPS) signals showing the movement of the ionospheric regions caus-
 146 ing diffractive fading as observed by Smith et al. (2008). Refraction and diffraction oc-
 147 cur simultaneously during scintillation (Zheng et al., 2022). Zheng et al. (2022) observed
 148 that IFLC, S4 and phase indices showed consistent fluctuations for most scintillation events.
 149 However, they observed that IFLC and S4 did not always simultaneously correspond to
 150 scintillations, as IFLC was enhanced during the geomagnetic storm and S4 is not. They
 151 suggested that IFLC during the geomagnetic storm is caused by the increased high-frequency
 152 phase power, which is related to the enhanced density of small-scale irregularities dur-
 153 ing storm periods.

154 The following questions still remain open: i) how do spatial and temporal charac-
 155 teristics of small-scale structures relate to auroral emissions? ii) how does particle pre-
 156 cipitation contribute to Fresnel’s scale structuring? In this work, we study how small-
 157 scale plasma structuring is spatially and temporally related to particle precipitation. The
 158 effects of small-scale structuring on trans-ionospheric radio-wave propagation are used
 159 to investigate times of elevated scintillation indices and understand the plasma condi-
 160 tions and generation of small-scale structures. To quantify small-scale structuring pro-
 161 cesses, we use scintillation indices. For auroral precipitation, we measure auroral inten-
 162 sity for different emission lines.

163 2 Approach and Data Selection

164 The scintillation indices used in this study are the amplitude scintillation index (S4)
 165 and the phase scintillation index (σ_ϕ) (Fremouw et al., 1978). We further use the ionosphere-
 166 free linear combination (IFLC). The scintillation indices and IFLC are calculated start-
 167 ing using amplitude and phase, which in ISMRs are usually sampled at 50 to 100 Hz rate
 168 for each satellite in view. The considered ISMR is a Septentrio PolaRxS/PolaRx5s multi-
 169 frequency multi-constellation receiver (Bougard et al., 2011) operated by the Istituto Nazionale
 170 di Geofisica e Vulcanologia (INGV, Italy) and situated at Longyearbyen, Svalbard (ge-
 171 ographic coordinates: 78.17° N, 15.99° E). The receiver is operational since January 2019
 172 and the scintillation data provided are part of the scintillation data collection (Upper
 173 atmosphere physics and radiopropagation Working Group et al., 2020) available in the
 174 electronic Space Weather Upper atmosphere (eSWua, eswua.ingv.it) data portal. In this
 175 study, we concentrate on data from the GPS and GALILEO constellation. GPS satel-

lites are indicated by G followed by their PRN code (e.g. G01), and GALILEO satellites by E followed by their PRN code (e.g. E01). The cut-off angle to avoid multipath effects is set to 35° . Scintillation indices and IFLC are then derived according to equations presented below (Yeh & Liu, 1982; de Paula et al., 2021; Carrano et al., 2013).

The signal intensity I is calculated using the signal amplitude A :

$$I = A^2 = I_c^2 + Q_c^2, \quad (1)$$

with I_c and Q_c being the signal intensity in-phase and quadrature components (Briggs & Parkin, 1963; Yeh & Liu, 1982; Van Dierendonck A. J.; Klobuchar J.; Hua, 1993). I is then detrended using a 6th-order high-pass Butterworth filter, with a cutoff frequency of 0.1 Hz. The amplitude scintillation index S_4 is then given by:

$$S_4 = \sqrt{\frac{\langle I^2 \rangle - \langle I \rangle^2}{\langle I \rangle^2}}. \quad (2)$$

The phase was detrended with a 6th-order high-pass Butterworth filter, with a cutoff frequency of 0.1 Hz, and the accumulated phase was calculated by $2\pi\phi_c$, with ϕ_c being the carrier phase. Subsequently, the phase scintillation index σ_ϕ is calculated by:

$$\sigma_\phi = \sqrt{\langle \phi^2 \rangle - \langle \phi \rangle^2}, \quad (3)$$

as the standard deviation of the detrended measured phase (Yeh & Liu, 1982; Kintner et al., 2007).

ISMRs traditionally have a firmware able to provide the scintillation indices in quasi real-time with a 1 minute time resolution. For the purpose of our work, we decided to consider scintillation data at a better time resolution, i.e., 1 second. Such approach has been proven to be more effective in detecting small-scale irregularities triggered by particle precipitation (Materassi & Mitchell, 2007; Wang et al., 2021), which are otherwise almost absent when considered on the standard 1-minute basis (see, e.g., the climatological pictures provided in (Spogli et al., 2009; Alfonsi et al., 2011; Moen et al., 2013; De Franceschi et al., 2019)). According to what reported in (Forte & Radicella, 2002) and in their eqn. (6) and (7), scintillation indices at 1 s allow studying spatial scales of the irregularities that are below 100 m, i.e. below the Fresnel's scale, and able account to spatial variability of particle precipitation patterns.

When a GNSS signal propagates through the ionosphere, refraction occurs. The phase delay caused by this refraction is given as the refractive index ionospheric correction term R_{corr} :

$$R_{corr} = \frac{40.3}{f_c^2} \int n_e dL, \quad (4)$$

with f being the signal carrier frequency and the electron density n_e is integrated along the raypath dL (Kashcheyev et al., 2012; Zheng et al., 2022). For propagation paths, which are the same for different carrier frequencies, the R_{corr} ratio of two carrier waves is given by

$$\frac{R_{corr1}}{R_{corr2}} = \frac{f_1^2}{f_2^2}. \quad (5)$$

Carrier phase variations L_ϕ at a constant frequency are given by

$$L_\phi = r + \lambda N - R_{corr}, \quad (6)$$

with the signal wavelength λ , integer ambiguity N , and the geometric distance r . The ionosphere-free linear combination is then calculated by:

$$IFLC = \frac{f_1^2 \phi_{L1} - f_2^2 \phi_{L2}}{f_1^2 - f_2^2} = r + \frac{f_1^2}{f_1^2 - f_2^2} \lambda N - \frac{f_2^2}{f_1^2 - f_2^2} \lambda N, \quad (7)$$

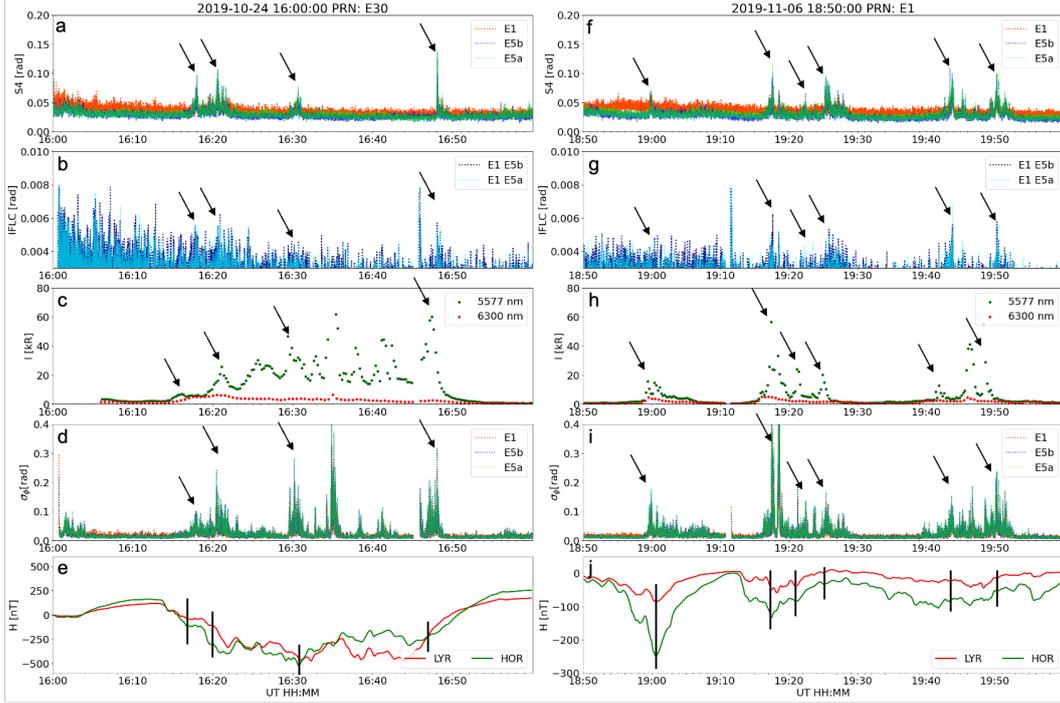
211 with the two carrier frequencies f_1 and f_2 and the corresponding carrier phases ϕ_{L1} and
 212 ϕ_{L2} . The ionospheric refractive index correction term R_{corr} is now removed and the IFLC
 213 is non-refractive (Cordes et al., 1986; McCaffrey & Jayachandran, 2019; Carrano et al.,
 214 2013; Zheng et al., 2022). The IFLC is then considered at the same rate of the I_c and
 215 Q_c samples, i.e., 50Hz.

216 In this work, the auroral emissions were captured by an all-sky camera (ASC) Keo
 217 Sentry 4ix Monochromatic Imagers from KEO Scientific, operated by the University of
 218 Oslo (Norway) situated in Longyearbyen Svalbard (geographic coordinates: 78.15° N,
 219 16.04° E). The imager is equipped with narrow band-pass filters to monitor 557.7 nm
 220 (green) and 630.0 nm (red) auroral emissions, recording images every 15–30 seconds. The
 221 green auroral emissions are projected to an altitude of 150 km (Partamies et al., 2022;
 222 Enengl et al., 2023). The red auroral images were projected to an altitude of 250 km.
 223 Observations of the brightest aurora may be influenced by the way the camera and is
 224 operated and calibrated. These observations are nevertheless a good measure for a com-
 225 parison between the auroral activity and the σ_ϕ indices.

226 In this article, we focus on the time period of 2019-2020 (during a solar minimum)
 227 the first season where the ISMR was operational. The availability of clear all-sky, au-
 228 roral images of intense particle precipitation, and simultaneous scintillation receiver data
 229 is the first filtering stage of our data set. Subsequently, we filter for a Kp index larger
 230 than 1.3 and solar wind speeds exceeding 400 km/s, to ensure ongoing geomagnetic ac-
 231 tivity. The geomagnetic and solar wind parameters were downloaded from the NASA/GSFC's
 232 OMNI data set through OMNIWeb (King & Papitashvili, 2005). Further, all events show
 233 moderate local deflections in the horizontal magnetic field H component (over 100 nT).
 234 The H component is recorded by a magnetometer network around Svalbard operated
 235 by the Tromsø Geophysical Observatory (Tanskanen, 2009). The signature of the enhance-
 236 ment of the westward electrojet and the substorm current wedge in superposition with
 237 eastward electrojet enhancements is imprinted as a decrease in the H component of the
 238 magnetic field at high latitudes (Akasofu, 1965; D'Onofrio et al., 2014). Six events ful-
 239 fil these criteria and are studied in detail.

240 3 Observations

241 The six strong particle precipitation events are shown in the context of scintilla-
 242 tion indices and magnetometer data indicating structuring processes and geomagnetic
 243 conditions. Figure 1 presents the two event types we encountered: long-lasting and short-
 244 lived precipitation events (Figures 1c and 1h, respectively). The data from 24th Octo-
 245 ber 2019 and 6th November 2019 show observations of increased S4 index (Figures 1a
 246 and 1f), indicating below Fresnel's scale plasma structuring. Peaks in S4 index values
 247 (shown by black arrows in Figures 1a and 1f) are recorded at: 16:18, 16:20, 16:31, 16:48
 248 UT on 24th October 2019 and 19:00, 19:18, 19:22, 19:26, 19:44, 19:50 UT on 6th Novem-
 249 ber 2019. To investigate the nature and cause of these elevated S4 observations, the IFLC,
 250 auroral intensity (I), σ_ϕ and H component are investigated. The peaks corresponding
 251 to elevated S4 indices are marked for all parameters in all panels by black arrows. The
 252 IFLC is shown in Figures 1b and 1g. Besides the elevated noise level at the beginning
 253 of both event windows, correspondence between S4 and IFLC is observed for the elevated
 254 S4 values. The IFLC indicates a diffractive contribution to scintillation. The auroral in-
 255 tensity at the piercing point of the satellite at the auroral emission plane is presented
 256 in Figures 1c and 1h. Here, we can observe the key difference between the two events:
 257 the 24th October 2019 event has an onset of auroral particle precipitation around 16:15
 258 and stays above 15 kR auroral intensity for the period after 16:20 UT to 16:48 UT. The
 259 6th November 2019 event shows multiple peaks in auroral intensity at around 19:00 UT,
 260 19:18, 19:21, 19:25, 19:42, 19:47. All of them are happening in a close minute window
 261 of the S4 index peaks, shortly before or after. The parameter σ_ϕ is shown in Figures 1d
 262 and 1i. The σ_ϕ index is a measure of small and large-scale structuring in the ionosphere.



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Figure 1. Small-scale plasma structures during particle precipitation. The data from 24th October 2019 and 6th November 2019 show observations of scintillation indices (S4 - panel a and f; σ_ϕ - panel d and i), the IFLC (panel b and g), magnetometer data (panel e and j), and auroral intensity (panel c and h). The peaks in S4 and IFLC mark strong intensity gradients (e.g. commence and end) of the elevated auroral intensity during strongly negative H, suggesting small-scale structuring during the commence and fading of particle precipitation. Most S4 peaks are accompanied by elevated σ_ϕ , indicating simultaneous large and small-scale structuring.

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Its diffractive contribution indicating small-scales and refractive contributions all-scale ranges. The elevations in the σ_ϕ index are roughly co-located with the elevated S4 index, however not all of the σ_ϕ variations can be associated with S4 index variations, e.g. on 24th October 2019 16:35. The H component is a measure for substorm activity, shown in Figures 1e and 1j. The H component decreases at the onset of auroral particle events. We find that the changes in the H-component measured by the LYR and HOR magnetometer (see red dashed line) are mostly coinciding with the S4 variations (indicated by black vertical lines).

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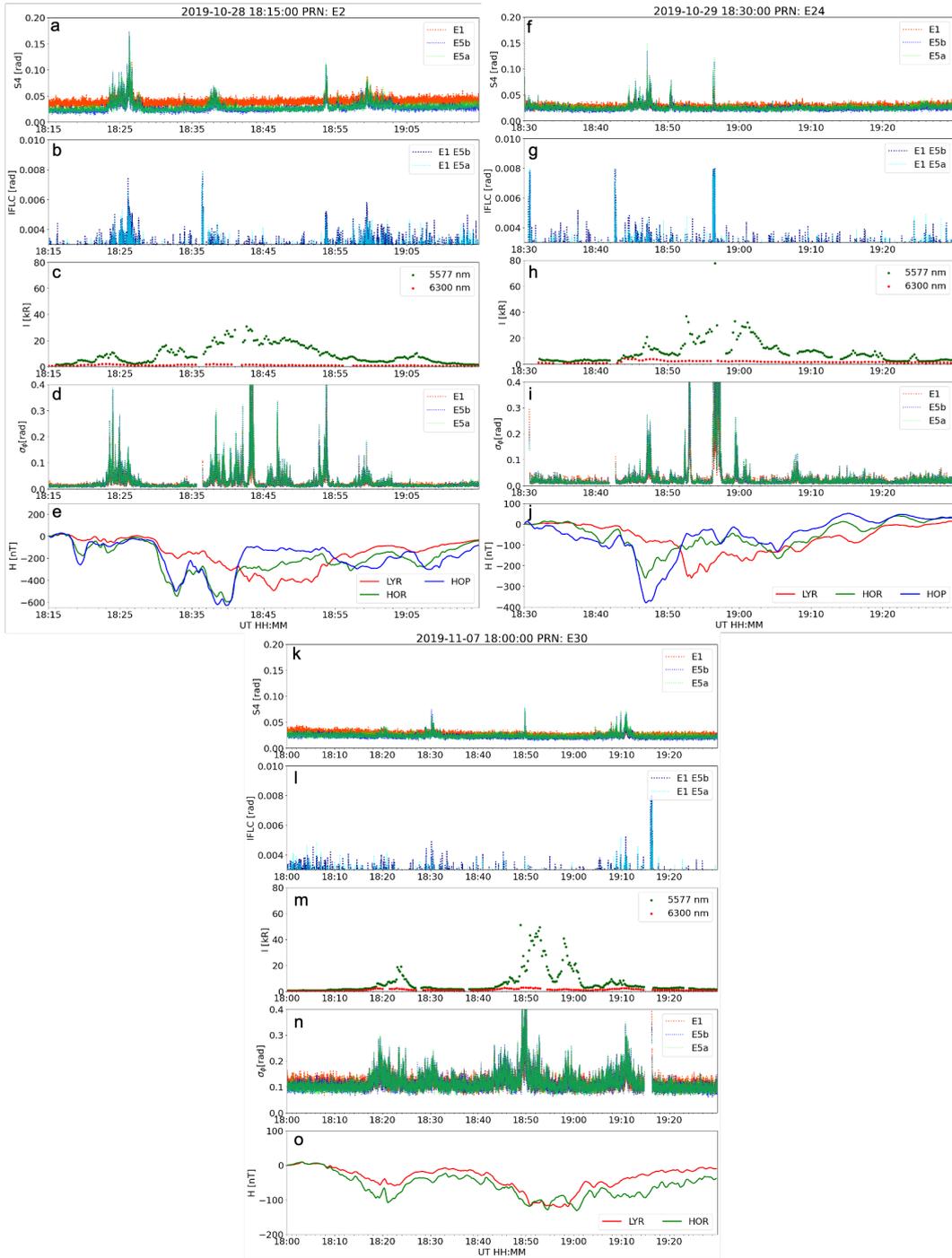
Additional three events are classified as long-lasting particle precipitation events as shown in the auroral intensity (Figures 2c, 2h and 2m). The events of 28th October 2019, 29th October 2019 and 7th November 2019 show the same features as the long-lasting event of 24th October 2019 in Figure 1. Comparing the S4 index shown in Figures 2a, 2f and 2k with the auroral intensity in Figures 2c, 2h and 2m shows that the elevated S4 marks the beginnings and ends of intervals of increased auroral emission intensity. This can be seen for the S4 peak at 18:28 and 19:00 UT on 28th October 2019 (2a) and in relation to the auroral intensity fading at 18:25 and 18:55 UT (Figure 2c). The IFLC (Figures 2b, 2g and 2l) and S4 index correlate for times of elevated S4 index, see e.g. 18:25, 18:35 and 18:55 UT on 28th October 2019. However, the IFLC can be noisy at times (Figure 2g at 18:30-18:45 UT on 29th October 2019). σ_ϕ variations (Figures 2d,

290 2i and 2n) are present for all occasions of elevated S4. The H component (Figures 2e,
291 2j and 2o) shows variations at aurora onset.

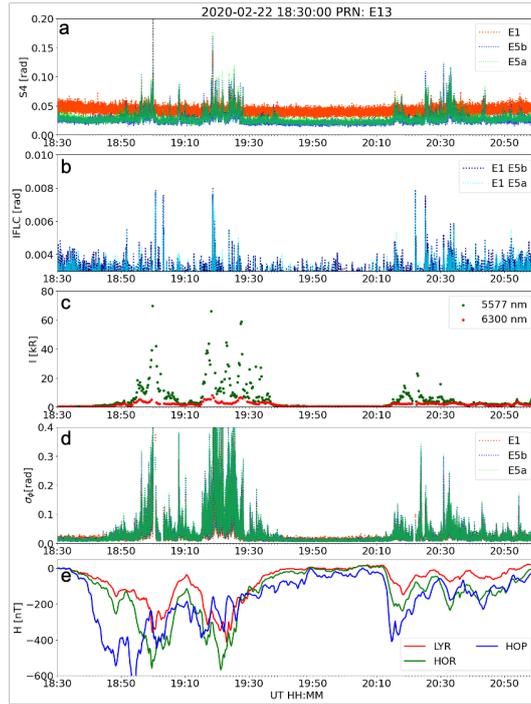
300 A second short-lived particle precipitation event is shown in Figure 3. Just as the
301 event of 6th November 2019 shown in Figure 1, it shares the characteristics of having
302 multiple short intense peaks in the auroral emission intensity, Figure 3 (Figure 3c), co-
303 located with the spikes in S4 (Figure 3a). While for long-lasting events, the S4 index is
304 elevated at intensity gradients, mainly marking the beginning and end of the auroral pre-
305 cipitation, for short-lasting events the precipitation can be intense and elevate the S4 in-
306 dex almost simultaneously. The H component (Figure 3e) repeats the pattern and shows
307 simultaneous decreases as the auroral intensity spikes. Most S4 peaks are accompanied
308 by elevated IFLC and σ_ϕ index.

317 To understand small-scale structuring in a spatial-temporal context, the piercing
318 points of the satellites are projected onto the auroral all-sky images, see Figures 4 and
319 5. At two occasions the satellite PRN E30 passes through a boundary of auroral precip-
320 itation (Figures 4a and 4c) into weaker emissions/background ionosphere (Figures 4b
321 and 4d). The auroral intensity at the piercing point of PRN E30 is peaking at 16:29:23
322 UT (Figure 4a) and 16:47:23 UT (Figure 4c), see also Figure 1c. The S4 index is peak-
323 ing at 16:30:38 and 16:48:08 UT, see Figures 4b, 4d and Figure 1a. This happens at the
324 same time as PRN E30 is crossing an auroral intensity gradient into weak auroral emis-
325 sions (Figure 4b) / or even no auroral emissions (Figure 4d). Small-scale structures are
326 observed at auroral intensity gradients.

327 Auroral particle precipitation crossing over piercing points of multiple satellites is
328 shown in Figure 5. At 19:17:23 UT PRN G24 is outside the auroral form (Figure 5a) and
329 shows no elevated S4 (Figure 5e). As the auroral form moves poleward, PRN G24 un-
330 dergoes an intensity gradient (Figures 5b and 5c). The S4 index onsets at about 19:18
331 UT after passing through the intensity gradient (Figures 5d and 5e). PRN E1 is within
332 strong auroral emissions at 19:17 UT, see Figure 5a. The S4 index is elevated latest at
333 19:17:38 UT (Figures 5b and 1f, and stays elevated at 19:17:53 UT (Figures 5c and 1f).
334 However at 19:18 when PRN E1 is at the center of high auroral intensity emissions (Fig-
335 ure 5d), the S4 index has returned to a local minimum before the next peak. PRN E9
336 does not pass through any auroral emissions in the investigated time window and does
337 not show any elevated S4 variations (Figures 5a-d and 5g). The S4 index of PRN E21
338 shows multiple spikes between 19:17-19:18 UT, Figures 5a-d, while it sits on a chang-
339 ing gradient of auroral intensity. Again, small-scale structures are observed at auroral
340 intensity gradients.

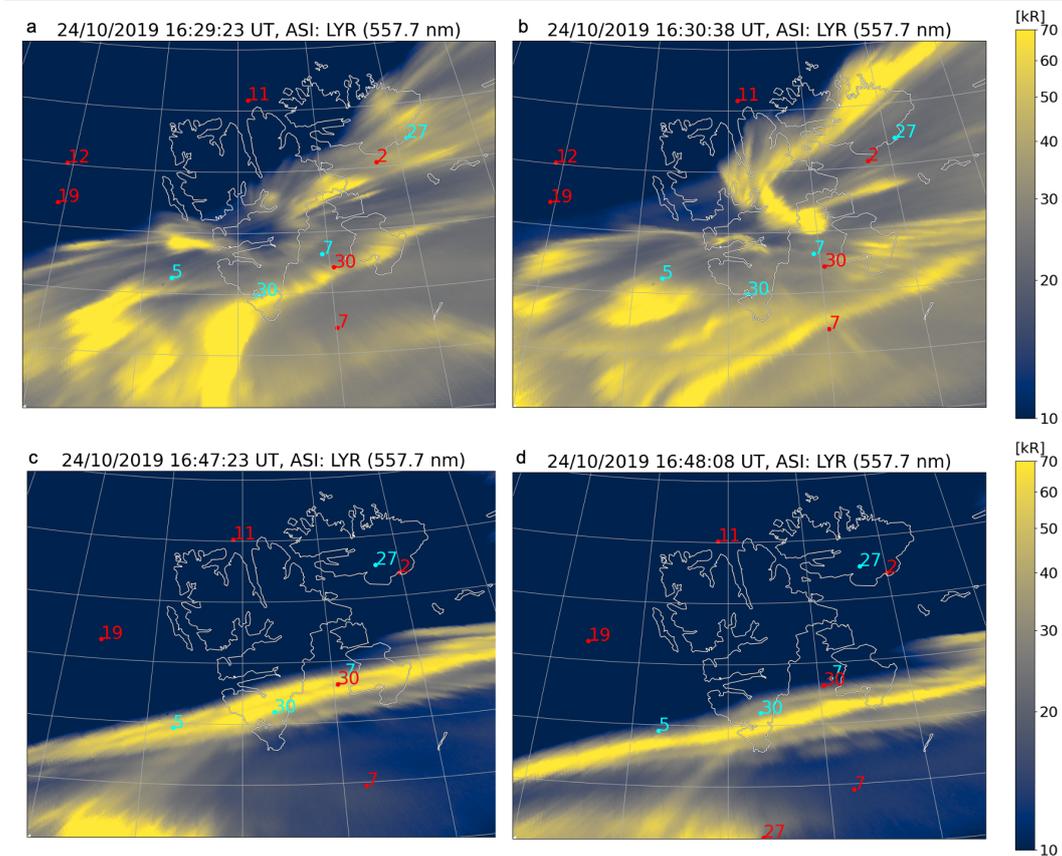


292 **Figure 2.** Small-scale plasma structures during long-lasting particle precipitation. The data
 293 from 28th October 2019, 29th October 2019 and 7th November 2019 show observations of scin-
 294 tillation indices (S_4 - panel a, f and k; σ_ϕ - panel d, i and n), the IFLC (panel b, g, and l), mag-
 295 netometer data (panel e, j and o) and auroral intensity (panel c, h and m). The peaks in S_4 and
 296 IFLC mark strong intensity gradients (e.g. commence and end) of the elevated auroral inten-
 297 sity during strongly negative H . This suggests small-scale structuring during the commence and
 298 fading of particle precipitation. Most S_4 peaks are accompanied by elevated σ_ϕ , indicating simul-
 299 taneous large and small-scale structuring.

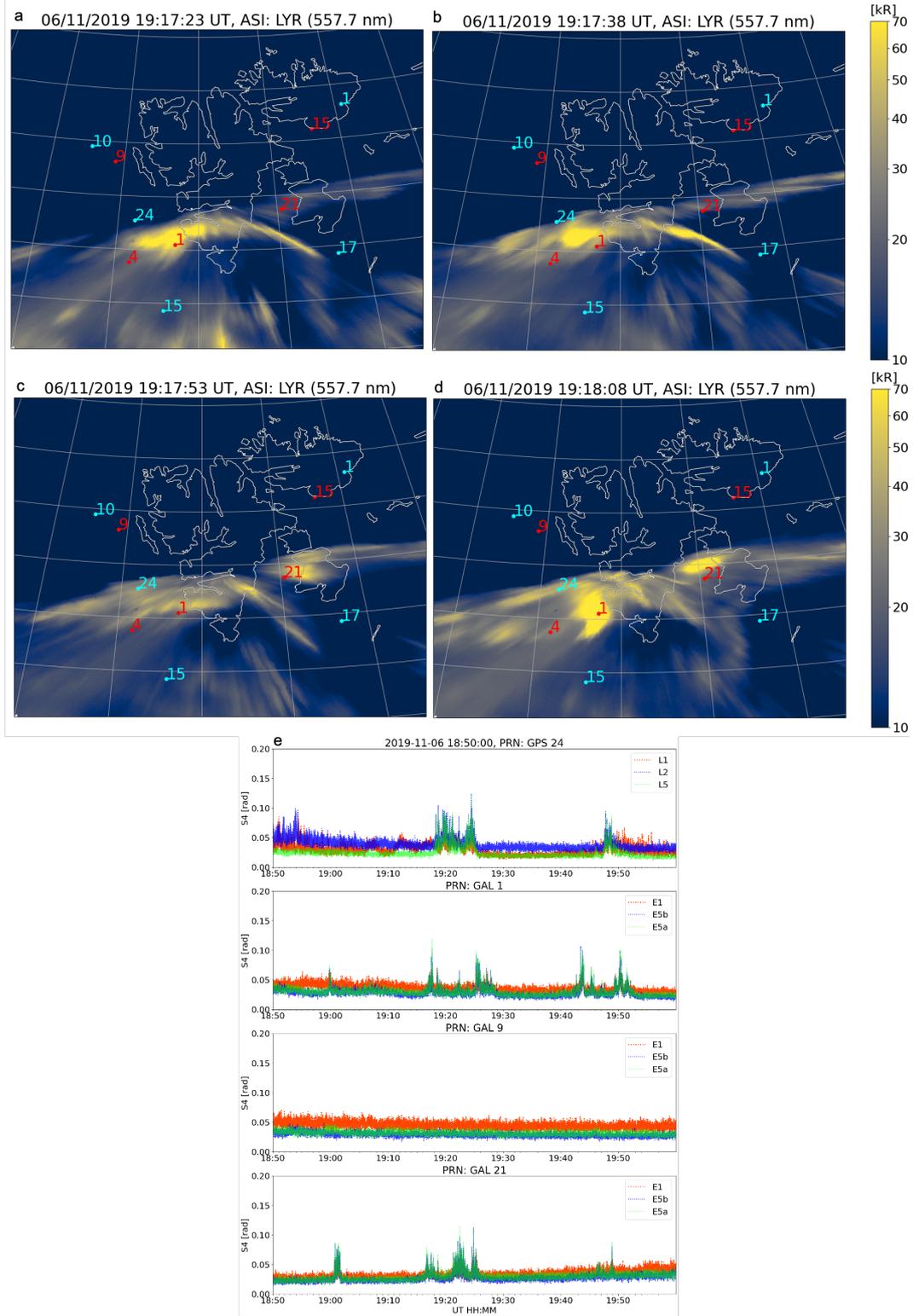


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310 **Figure 3.** Small-scale plasma structures during short-lived particle precipitation event. The
 311 data from 22th February 2020 show observations of scintillation indices (S_4 - panel a; σ_ϕ - panel
 312 d), the IFLC (panel b), magnetometer data (panel e) and auroral intensity (panel c). The peaks
 313 in S_4 and IFLC mark strong intensity gradients of the elevated auroral intensity during strongly
 314 negative H . This suggests small-scale structuring at large intensity gradients of particle pre-
 315 cipitation. Most S_4 peaks are accompanied by elevated σ_ϕ , indicating simultaneous large and
 316 small-scale structuring.



341 **Figure 4.** Spatial evolution of auroral particle precipitation and small-scale structuring at the
 342 satellite piercing point. The Figure shows all-sky images and GNSS positions relative to another
 343 projected onto 150 km altitude together with a map of Svalbard. The spatial evolution remains
 344 the same as for azimuth/elevation skyplots, as shown in the supplementary material (Figure Sky-
 345 plot.png for 24th October 2019 16:30:38). Intense yellow corresponds to strong auroral intensity.
 346 GPS satellite piercing points are shown in cyan, GALILEO satellite piercing points are shown in
 347 red. Two incidents on 24th October 2019 of satellite PRN E30 passing through strong auroral
 348 emissions with a strong consequent response in increased S4 index values indicating small-scale
 349 structuring. The auroral intensity is peaking at 16:29 and 16:47 UT as shown in panels a and c.
 350 And after that, we observe elevated S4 index at 16:30 and 16:48 UT at low auroral intensity, see
 351 panels b and d. Elevated S4 indices are observed at auroral intensity gradients.



352

353 **Figure 5.** Auroral particle precipitation over four different satellite piercing points (PRN GPS
 354 24, E1, E9 and E21). Panels a-d show GNSS positions and all-sky images projected to 150 km
 355 altitude together with a map of Svalbard. Intense yellow corresponds to strong auroral intensity.
 356 GPS satellite piercing points are shown in cyan, GALILEO satellite piercing points are shown in
 357 red. The S4 index variations from 18:50 to 20:00 UT are shown in panel e for the satellites PRN
 358 G24, E1, E9 and E21. Elevated S4 indices are observed for PRN G24, E1 and E21 coinciding
 359 with the particle precipitation gradient. PRN E9 does not pass through auroral precipitation and
 360 observes no small-scale plasma structuring. -12-

4 Discussion

We have analysed the response of the amplitude scintillation index S4, evaluated on a 1 s level, to auroral particle precipitation in the evening sector 16:00–21:00 UT (MLT: 19:00–24:00) over Svalbard. Six events from the 2019/2020 season were studied with data from the LYR all-sky and LYR GNSS receiver and show intense auroral emissions (above 30 kR) and solar wind speeds exceeding 400 km/s accompanied with moderate local deflections in the H component (over 100 nT) for $K_p > 1.3$. The events were sorted into two different types of precipitation: long-lasting (> 5 min) and short-lived (< 5 min).

In all observed events, the green auroral emissions peaked above 30 kR and were correlated with scintillation indices, whereas the red auroral emissions remained weak (below 10 kR). At times of elevated σ_ϕ and S4 values, the red auroral emissions are weak, whereas the green auroral emissions correspond to the σ_ϕ variations. This suggests that in particular E-region plasma processes and instabilities driven by auroral particle precipitation resulting in green emissions may lead to small- and large-scale plasma structuring processes that impact GNSS radio waves. In the following, we will focus on the green auroral emissions in relation to the derived scintillation indices.

For the long-lasting events, multiple peaks were observed in the S4 index indicating the beginning, end, and the intensity gradients. On 24th October 2019, the onset at 16:15 UT and fading at 16:47 UT of the particle precipitation are reflected in the S4 index increases, and so is the sharp auroral intensity increase at 16:29 UT, see again Figures 1a and 1c. The particle precipitation in the event on 28th October 2019 commences at 18:22 UT, a sharp increase in the auroral intensity is also noted at 18:36 UT and the intensity decrease is gradual after 18:53. All of these times are accompanied by S4 index increases, see again Figures 2a and 2c. The third long-lasting particle precipitation event on the 29th October 2019 shows a particle precipitation onset at 18:45 UT correlated with S4 index variations. During the event, the auroral intensity peaks around 18:57 UT and with it the S4 index, see again Figures 2f and 2h. No variations of the S4 index are found at the fading of the precipitation in this event. The event on 7th November 2019 shows an increase in the S4 index at its abrupt start at 18:50 UT, and the S4 index varies after 19:03 UT and at 19:10 UT as the intensity is fading. A shorter, just above 5-minute long event is shown at 18:20–18:30 UT with an elevated S4 index marking its beginning and end, see again Figures 2k and 2m.

The short-lived events show similar signatures of correlated precipitation and S4 index as the long-lasting events, however, as they are short the S4 index variations do not cease during the event. On the 6th November 2019, multiple peaks in the auroral intensity (19:00, 19:18, 19:22, 19:26, 19:44, 19:50 UT) are correlated to the peaks in the S4 index, see again Figures 1f and 1h. As the onset and fading happen in 1–5 minute intervals, the S4 index shows continuous variations in response. On 22nd February 2020, there were multiple short-lived auroral emission peaks that were close to one another between 19:10–19:40 UT when the satellite signal is subject to intensification of the S4 index above the noise level. The same can be observed for the other peaks (18:58, 20:30 UT) in the investigated time window, see again Figures 3a and 3c.

A correspondence between the IFLC and S4 index can be observed for most of the above listed elevated S4 index variations related to auroral particle precipitation, when there is low noise in the IFLC values (all events except 24th October 2019 at 16:00–16:20 UT), see Figures 1, 2 and 3- S4 and IFLC.

During all above listed elevated S4 index variations, we find elevated σ_ϕ values. This means that we find small-scale plasma structuring accompanied with larger scale plasma structuring. The σ_ϕ index magnitude, however, is not a measure to predict the S4 index magnitudes. Some σ_ϕ variations do not correspond to any S4 variations, such as on 24th October 2019 at 16:35 (Figure 1d) UT and 28th October 2019 at 18:40 UT (Fig-

412 ure 2d). The σ_ϕ index can also be elevated during long-lasting events, while the S4 index
 413 is specifically elevated at auroral intensity gradients.

414 Sharp decreases in the H component are found in relation to the S4 index increases.
 415 However, as the receiver to satellite line-of-sight is at a distance to the magnetometer
 416 stations, there is a delay between the observed H decreases and the S4 increases, see again
 417 Figures 1, 2 and 3 panels for the H component of the magnetic field. This indicates the
 418 ionospheric current system contributing to Fresnel's scale structuring.

419 Comparing the spatial evolution of the auroral emissions with the satellite pierc-
 420 ing points and elevated S4 index (Figures 4 and 5) confirms that elevated S4 indices are
 421 found near auroral emission gradients. The highest S4 values are located at the edges
 422 of auroral forms (when moving in/out of a form) or at local intensification or weaken-
 423 ing of auroral precipitation along the line of sight. In short, elevated S4 indices are ob-
 424 served at or after auroral intensity gradients along the line of sight from the satellite to
 425 the receiver.

426 In summary: (1) Elevated S4 values are associated with auroral intensity gradients
 427 as different levels of particle precipitation pass through the line of sight of the satellite
 428 to the receiver. (2) The IFLC often corresponds to increases of S4 index variations. (3)
 429 Intensification of ionospheric currents contributes to Fresnel-scale structuring.

430 In the following, we interpret the observations in terms of ionospheric plasma pro-
 431 cesses and the impacts on trans-ionospheric radio wave propagation.

432 During ionosphere-magnetosphere coupling, strong electric fields and currents can
 433 dissipate energy and excite plasma instabilities that create turbulence in the E-region
 434 ionosphere. The turbulence can induce nonlinear currents and strong electron heating,
 435 which can increase the global ionospheric conductance (Schlegel & St.-Maurice, 1981;
 436 Bahcivan, 2007; Dimant & Oppenheim, 2011a, 2011b). The E-region ionospheric plasma
 437 is subject to collisions with available neutrals, and energetic particle beams from the par-
 438 ticle precipitation. Particle precipitation into the ionospheric E-region leads to a widespread
 439 irregularity dissipation and a redistribution of energy and plays a main role in E-region
 440 large-scale ionospheric plasma structuring (Makarevich et al., 2021; Enengl et al., 2023).

441 Both, short-lived and long-lasting particle precipitation events, show the same re-
 442 sponse of S4 index variations to precipitation, namely elevated S4 values linked to au-
 443 roral intensity gradients. Auroral intensity gradients may indicate that processes at the
 444 auroral boundary, between background plasma and an energetic particle precipitation
 445 beam, contribute to small-scale plasma structuring in the E-region ionosphere. The Farley-
 446 Buneman instability (FBI) extracts its energy from the velocity difference of electrons
 447 streaming past ions, which collide with neutrals (Farley Jr., 1963; Buneman, 1963). The
 448 FBI may explain why the S4 index variations are present at the edge of strong particle
 449 precipitation gradients.

450 Dimant et al. (2021) states that during geomagnetic storms intense electron pre-
 451 cipitation frequently penetrates into regions of the strong electric fields in the E-region
 452 ionosphere. Without precipitation, strong electric fields drive E-region instabilities, which
 453 further lead to plasma turbulence and increased conductance. Electron precipitation, how-
 454 ever, dramatically raises the instability threshold, and can even largely suppress the in-
 455 stability inside the auroral regions. Our observations agree with Dimant et al. (2021),
 456 as we do not observe plasma structuring processes within auroral emissions, but we do
 457 observe plasma structuring at the boundaries of the particle precipitation signature.

458 E-region instabilities can cause electron heating, enhanced plasma particle trans-
 459 port and small-scale plasma turbulence that modifies large-scale ionospheric conductance
 460 and with it the entire dynamics of the near-Earth's plasma (Dimant et al., 2021). Our
 461 observations of co-located large-scale σ_ϕ index observations corresponding to small-scale

462 amplitude index variations confirm the hypothesis of E-region dynamics affected by par-
 463 ticle precipitation leading to large and small-scale (above and below Fresnel scale) plasma
 464 structuring.

465 Field-aligned currents couple the magnetosphere to the ionosphere. Particle pre-
 466 cipitation is a signature of this phenomenon and can directly drive ionospheric plasma
 467 by structured precipitation and result in elevated scintillation index values (Boström,
 468 1964; Carter et al., 2016; Xiong et al., 2020; Fæhn Follestad et al., 2020). In our obser-
 469 vations, the drops in the H component confirm geomagnetic perturbations at the onset
 470 of particle precipitation and elevated scintillation indices. It suggests that the measured
 471 currents in fact contribute to FBI driving and cause plasma structuring on various scales
 472 at the auroral boundary. We can conclude that field-aligned currents not only drive large-
 473 scale plasma structuring but also below Fresnel-scale structuring, as shown by the simul-
 474 taneous response of scintillation indices (σ_ϕ and S4 index variations) at the edge of field-
 475 aligned current signatures (the precipitation signatures and H component).

476 For the studied events, we find IFLC variations at the same time as elevated S4
 477 index values. Zheng et al. (2022) stated that although both, IFLC and S4, are indicat-
 478 ing diffractive variations, they however are not always correlated. This is in agreement
 479 with our study, as we find IFLC correlated with S4 for most of the S4 enhancements,
 480 but not all IFLC variations are reflected in the S4 index. This may be due to high noise
 481 levels in the IFLC. The IFLC can be used as a measure for diffractive effects as the re-
 482 fractive part is removed. High-frequency, refractive, variations in the GPS carrier phase
 483 can be wrongly classified as scintillation, where refractive variations are deterministic
 484 and diffractive variations are stochastic (McCaffrey & Jayachandran, 2019). Care must
 485 be taken at the choice of the cutoff frequency for phase detrending or a fast iterative fil-
 486 tering detrending scheme can be used to provide a more realistic determination of the
 487 phase scintillation index (Ghobadi et al., 2020). We therefore carefully studied S4 and
 488 IFLC to correctly classify diffractive variations. Our results show that diffractive effects
 489 impact the signal at the onset, fading or strong changes of particle precipitation. Since
 490 σ_ϕ variations are not always found simultaneously with S4 index variations, some σ_ϕ vari-
 491 ations can be of purely refractive nature caused by convected larger plasma structures
 492 during the particle precipitation. While at strong precipitation gradients, the σ_ϕ index
 493 also includes diffractive variations but seems otherwise to mainly include refractive sig-
 494 nal variations. Conroy et al. (2022) suggests that high latitude phase variations are mainly
 495 of refractive nature. As the IFLC and S4 show little to no variation during long-lasting
 496 precipitation events, they are observed as refractive and deterministic. Only at the au-
 497 roral intensity gradients (e.g. beginning and end) the scintillations are classified as diffrac-
 498 tive. For short-lasting precipitation, we find IFLC and S4 elevated, therefore they are
 499 fully classified as diffractive. In summary, we observe diffractive scintillation accompa-
 500 nished by amplitude scintillation during strong gradients (e.g. short-lived events, begin-
 501 ning and ending of long-lasting events) leading to stochastic effects.

502 5 Conclusions

503 This paper investigated the response of the S4 index, evaluated from GNSS signals
 504 on a 1 second basis. to intense auroral particle precipitation events over Svalbard. The
 505 S4 index is a proxy for small-scale plasma structuring, which we find to be associated
 506 with intense particle precipitation. Elevated S4 index values can be observed under var-
 507 ious conditions, even without strong auroral emissions. Here we specifically looked at
 508 strong auroral emissions in the evening sector during ongoing geomagnetic activity quan-
 509 tified by local deflections of the horizontal magnetic field component and combined with
 510 all-sky camera and scintillation receiver data at high latitudes (Longyearbyen, Svalbard).
 511 We identified 6 intense precipitation events to study the S4 response. Our results show
 512 that:

- 513 1. Clear increases in the S4 values above the noise floor level are observed at large
 514 auroral intensity gradients, as the bright aurora is moving in/out of the line of sight
 515 from the satellite to the ground receiver. The values of S4 indicate a weak scin-
 516 tillation activity, as S4 is never exceeding 0.2.
- 517 2. The IFLC is often elevated simultaneously with the S4 index, confirming the diffrac-
 518 tive nature of the events.
- 519 3. Significant increases in the ionospheric currents (auroral particle precipitation) con-
 520 tribute to plasma structuring below the Fresnel's scale.

521 Why small-scale plasma structuring processes in the E-region are observed specif-
 522 ically at auroral intensity gradients remains an open question. Further statistical stud-
 523 ies need to be carried out to investigate plasma instabilities like Farley-Buneman as a
 524 likely cause for small-scale structuring. Multi-instrument case studies, including in-situ
 525 measurements by rockets and satellites, are required to understand the combination of
 526 simultaneous large-scale and small-scale structuring processes driven by auroral parti-
 527 cle precipitation, their spatial/temporal evolution and their connection with the mag-
 528 netic field geometry linked to the IMF conditions.

529 Open Research Section

530 The scintillation data managed by the Istituto Nazionale di Geofisica e Vulcanolo-
 531 gia (INGV) are available at eSWua web portal (www.eswua.ingv.it), that is operated by
 532 the Upper Atmosphere Physics and Radio propagation group of INGV.

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