

1 **Climatology of High-frequency Gravity Waves**
2 **Observed by an Airglow Imager at Andes Lidar**
3 **Observatory**

4 **Bing Cao¹, Alan Z. Liu²**

5 ¹Institute of Geophysics and Planetary Physics, Scripps Institution of Oceanography, University of
6 California San Diego, La Jolla, CA, USA

7 ²Center for Space and Atmospheric Research and Department of Physical Sciences, Embry-Riddle
8 Aeronautical University, Daytona Beach, FL, USA

9 **Key Points:**

- 10 • A large number of high-frequency gravity waves were observed near mesopause
11 region over the Andes by an airglow imager.
- 12 • Preferential propagation direction of the waves shows seasonal dependence, pole-
13 ward in austral summer and equator-ward in austral winter.
- 14 • Convective activities are a likely wave source, playing an important role in shap-
15 ing the observed wave directionality.

Abstract

The long-term climatology of high-frequency quasi-monochromatic gravity waves is presented using multi-year airglow images observed at Andes Lidar Observatory (ALO, 30.3°S, 70.7°W) in northern Chile. A large number of high-frequency gravity waves were retrieved from OH airglow images. The distribution of primary wave parameters including horizontal wavelength, vertical wavelength, intrinsic wave speed, and intrinsic wave period are obtained and are in the ranges of 20–30 km, 15–25 km, 50–100 ms⁻¹, and 5–10 min, respectively. The waves tend to propagate against the local background winds and show clear seasonal variations. In austral winter (May–Aug), the observed wave occurrence frequency is higher and preferential wave propagation is equator-ward. In austral summer (Nov–Feb), the wave occurrence frequency is lower and the waves mostly propagate pole-ward. Critical-layer filtering plays an important role in determining the preferential propagation direction in certain months, especially for waves with a small observed phase speed (less than typical background winds). The wave occurrence and preferential propagation direction are shown to be related to the locations of convection activities nearby and their relative distance to ALO. However, other possible wave sources such as secondary wave generation and possible ducted propagation cannot be ruled out. The estimated momentum fluxes have typical values of a few to 10 m²s⁻² and show seasonal variations with a clear anti-correlation with local background wind directions.

1 Introduction

Airglow refers to the emissions of photons in the Earth atmosphere via chemiluminescence processes, that mainly result from the reactions among species such as atomic oxygen, atomic nitrogen, and hydroxyl radicals (Khomich et al., 2008). Several of these emissions originate within the Mesosphere and Lower Thermosphere (MLT) region (altitude range around 80–100 km) as thin luminous layers with typical thickness of 6–10 km (Full Width at Half Maximum, or FWHM). Historically, the first airglow emissions to be investigated were the green ionized oxygen (OI) line (557.7 nm) with peak altitude at ~96 km and the yellow Na line (589.2 nm) with peak altitude at ~90 km. But the brightest source of airglow is the hydroxyl (OH) Meinel bands emission (peak altitude at ~87 km) which radiates over a broad spectral range (0.7–4.0 μm) primarily in the near-infrared band. Many studies have revealed that these airglow emissions are very useful tracers to retrieve the atmospheric properties and study the dynamical processes such

48 as instabilities, ripples, small scale gravity waves, as well as larger scale atmospheric waves
49 such as tides and planetary waves (e.g., Medeiros et al., 2007; T. Li et al., 2009; Cao et
50 al., 2016; J. Li et al., 2017).

51 The atmospheric flow in the MLT region is dominated by abundant atmospheric
52 waves, of which gravity waves are an important type with large varieties in wave char-
53 acteristics and potential sources. High-frequency atmospheric gravity waves carry sig-
54 nificant amount of momentum from lower atmosphere. The dissipation and breaking of
55 these waves have large impacts to the circulation through momentum deposition to the
56 background flow. Airglow imaging systems are most sensitive to this part of gravity waves
57 spectrum because of the high horizontal and temporal resolution (J. Hecht et al., 2001;
58 Ejiri et al., 2003; J. H. Hecht et al., 2004). Gravity wave information can be inferred from
59 the wave induced emission intensity fluctuations detected by such imaging systems. These
60 gravity waves are revealed with typical horizontal wavelengths of 20 to 100 km, intrin-
61 sic wave periods of 5 to 10 min, and horizontal phase speeds between 30 to 100 ms^{-1} (Taylor,
62 1997; Ejiri et al., 2003; Z. Li et al., 2011). The momentum flux estimated from airglow
63 emission perturbation has an average magnitude of 5–10 m^2s^{-2} (J. Tang, Kamalabadi,
64 et al., 2005; Y. Tang et al., 2014). Studies based on airglow observations suggest that
65 the wave propagation in the mid-latitudes often shows an annual variation: pole-ward
66 in summer and equator-ward in winter. Several mechanisms such as critical layer filter-
67 ing (Taylor et al., 1993), ducted wave propagation (Walterscheid et al., 1999), variations
68 of the location of wave sources (Nakamura et al., 2003) and Doppler-shift by the back-
69 ground winds (Z. Li et al., 2011) were proposed to explain the directionality of wave propa-
70 gation. In general, these mechanisms all play some roles in affecting the wave propa-
71 gation directions but their relative importance varies with seasons and geographic loca-
72 tions.

73 Andes Lidar Observatory (ALO) is located at Cerro Pachón (30.3°S, 70.7°W) on
74 the west side of Andes ridge, which is generally aligned in the north-south direction and
75 extends several thousands kilometers in South America. The elevation of ALO is 2530
76 m compared to the ridge peak altitude of 4500–5000 m. Many satellite observations have
77 revealed the existence of gravity wave hot spots in the stratosphere over southern An-
78 des (Hoffmann et al., 2013; Hindley et al., 2015). It is believed that the major wave sources
79 are subtropical deep convection in low and mid-latitude, and orographic sources at lat-
80 itudes of 40°S to 70°S during austral winter time (Jiang et al., 2004). Whether these ac-

81 tive gravity waves in stratosphere reach higher MLT region before they break remains
82 an unanswered question. Airglow imaging systems, together with other observation in-
83 struments, have been utilized to depict a complete picture of gravity wave propagation
84 from stratosphere to mesosphere (Bossert et al., 2014; Fritts et al., 2019; Taylor et al.,
85 2019). By comparing the results with previous deployment of airglow imager at other
86 locations, the similarities as well as differences of wave characteristics, preferential prop-
87 agation directions, and momentum flux may reflect the generality and specialty of wave
88 sources and background winds at different locations.

89 In this study, we present an application of the long-term dataset of mesospheric
90 airglow and wind observations that were acquired in the Andes. The dataset is used to
91 study the distribution of the intrinsic gravity wave parameters, their dominant propa-
92 gation directions and possible controlling mechanisms, as well as variation of momen-
93 tum flux and its relationship with background wind. The study is organized as follow-
94 ing: section 2 briefly describes the instrumentation from which the data were retrieved
95 and section 3 describes the dataset and methodology. Sections 4 and 5 present the main
96 results and discussion. The summary and conclusions are presented in section 6.

97 **2 Instrumentation**

98 **2.1 Airglow Imager**

99 An all-sky airglow imager is equipped with a cooled charge-coupled device (CCD)
100 and a fish-eye lens to collect the airglow emissions from all the sky. One or several nar-
101 row width bandpass filters are used to distinguish the emissions of different spectrum
102 range from different altitude ranges (Taylor et al., 1995). The airglow imager operated
103 at ALO was equipped with two filters to capture OH and OI emissions alternately at night
104 during the low moon period throughout the year. The integration times for the OH and
105 OI images are 1 min and 1.5 min, respectively. For OH Meinel band emission, the band-
106 width of the filter is 750–930 nm with a notch at 865 nm to exclude the molecular oxy-
107 gen emission. The airglow emissions were collected by a 1024×1024 CCD array and
108 then binned to a 512×512 array to increase signal-to-noise ratio. When the field-of-view
109 is limited within $\pm 45^\circ$ zenith angle, the airglow images cover an area of about 200×200
110 km^2 with a resolution better than 1 km/pixel if projected to OH airglow altitude at ~ 87
111 km.

112 2.2 Meteor Radar

113 The ALO meteor radar uses a SKiYMET radar system (Franke et al., 2005) op-
114 erating at 40.92 MHz. There are two major components of the radar. The transmitter
115 is a three-element Yagi antenna directed toward the zenith with a transmitted power of
116 approximately 170 W from a 13.3 μm pulse length, 6 kW peak envelop power and 466
117 μm inter-pulse period. The meteor trails were illuminated by the radiated energy. The
118 receiver is comprised of five three-element Yagi antenna oriented along two orthogonal
119 baselines and they sampled every 13.3 μm , resulting in 2 km range resolution. The backscat-
120 tered signals from meteor trail are received by different antennas at different arrival an-
121 gle and timing. Then, the interferometry method was performed to determine the po-
122 sition of meteor trail in the sky. Wind velocities were retrieved from the continuous track-
123 ing of trail positions and Doppler shifts (Hocking et al., 2001) with the assumption that
124 the horizontal wind field is almost uniform and stationary within the spatio-temporal
125 window and the vertical wind is negligible. The meteor radar provides continuous hourly-
126 averaged horizontal winds between 80 and 100 km (Franke et al., 2005). The winds around
127 the OH airglow layer were calculated through Gaussian-weighted averaging centered at
128 87 km with a window of 5 km.

129 3 Data and Methodology

130 A narrow-band sodium wind/temperature lidar, an all-sky airglow imager, a pho-
131 tometer, and a meteor radar have been deployed to ALO since September 19, 2009. The
132 airglow imager captured only OH images before Aug 2011, OH and OI images alternately
133 after that. The possible influence of the different image timing on wave extraction is dis-
134 cussed in details in Supporting Information. The meteor radar had some technical is-
135 sues in mid 2014, thus no more wind data were observed afterwards. Therefore, we only
136 processed the airglow data when the meteor radar wind data were available (2009 to 2014).
137 The number of hours when OH airglow images were obtained are summarized in Table
138 1. There are about 300–600 hours of data in each calendar month accumulated in 6 years.
139 The amount of data enables a robust analysis of seasonal variations of gravity waves.

140 Before airglow images can be used for wave extraction, there are several pre-processing
 141 procedures that need to be implemented. Firstly, all the stars present on the images need
 142 be removed. Secondly, images need to be unwrapped to correct the spatial distortions
 143 due to fish-eye lens and emission intensity variation due to van Rhijn effect. Thirdly, the
 144 Milky Way over ALO in southern hemisphere is present and close to zenith most of the
 145 time and is much brighter than the airglow emission within the imager observational band-
 146 width. Therefore, an additional procedure of removing the Milky Way (Z. Li et al., 2014)
 147 is necessary and applied before gravity waves can be identified.

148 High frequency, quasi-monochromatic gravity waves are identified from the images
 149 using a series of procedures described in detail in J. Tang, Franke, et al. (2005) and J. Tang,
 150 Kamalabadi, et al. (2005) and briefed here. Three consecutive images (I_1, I_2, I_3) were
 151 used to form two consecutive time-differenced (TD) images ($TD_1 = I_2 - I_1, TD_2 =$
 152 $I_3 - I_2$) for spectral analysis. Horizontal wave parameters including wavelength, observed
 153 phase speed, propagation direction and relative airglow perturbation amplitude (I'_{OH}/\bar{I}_{OH})
 154 were derived from each set of two TD images. Intrinsic phase speed and intrinsic frequency
 155 are derived with background winds provided by meteor radar. Vertical wavelength is cal-
 156 culated using a simplified dispersion relationship (equation 24 of Fritts and Alexander
 157 (2003)) with buoyancy frequency near the OH airglow layer derived using temperature
 158 from NRLMSISE-00 empirical model (Picone et al., 2002). The relative airglow inten-
 159 sity amplitude is calculated by dividing the perturbation intensity I'_{OH} by the average
 160 intensity \bar{I}_{OH} of the star-free and de-trended images after excluding the dark current
 161 and background emission, which is assumed to be 30% of total emission intensity (Swenson
 162 & Mende, 1994). The gravity wave momentum flux was derived based on the intrinsic
 163 wave parameters and relative temperature amplitude, converted from I'_{OH}/\bar{I}_{OH} using
 164 the airglow model described in Liu and Swenson (2003). The total gravity wave momen-
 165 tum flux is calculated using the following equation:

$$F_m = \frac{k}{m} \frac{g^2}{N^2} \left\langle \left(\frac{T'}{\bar{T}} \right)^2 \right\rangle = \frac{k}{m} \frac{g^2}{N^2 C^2} \left\langle \left(\frac{I'_{OH}}{\bar{I}_{OH}} \right)^2 \right\rangle (\text{m}^2 \text{s}^{-2}), \quad (1)$$

166 of which k, m are the horizontal and vertical wavenumber, N^2 is the squared buoyancy
 167 frequency and C is the cancellation factor, which is a function dependent on wave in-
 168 trinsic parameters, especially vertical wavelengths (Liu & Swenson, 2003; Hickey & Yu,
 169 2005). For each set of images (two TD images or three raw images), there can be zero
 170 to multiple gravity waves identified and counted within this period of triple image in-

Table 1. Number of accumulated hours of OH airglow images captured in each month at ALO.

Year	Jan	Feb	Mar	Apr	May	Jun	Jul	Aug	Sep	Oct	Nov	Dec	Total
2009									61.5	49.8	141.9	284.0	537.2
2010	176.7	183.7	213.5	321.5	194.9	128.8	212.7	270.7	120.2			23.5	1846.4
2011	117.7	5.1	30.5	133.7	251.2	238.1	258.7	187.9	161.8	120.4	160.6	210.8	1876.7
2012	66.2	13.3	112.2	13.7	6.4	11.3	233.1	6.5	120.2	35.6	54.5	32.9	705.7
2013		59.9	121.3	107.4		92.0	74.8	27.8	40.9	112.4	63.9	9.0	709.4
2014	120.4	90.5	61.6	90.4	53.0		42.12	80.9					539.0
Total	481.0	352.3	539.1	666.7	505.5	470.3	821.5	573.8	504.6	318.2	420.9	560.2	6214.2

171 tegration time. There exist some gravity wave events lasting longer time and showing
172 up in multiple sets of images. In our analysis, one persistent wave event will be counted
173 as multiple waves that are retrieved from different sets of TD images. Therefore, the statis-
174 tics based on wave counts represent the overall duration of gravity waves, not numbers
175 of coherent gravity wave events. However, the statistics of gravity wave events are also
176 analyzed.

177 There are a few steps in the data processing that need extra attention when dis-
178 cussing results in the following sections. Firstly, the TD method acts as a high-pass fil-
179 ter and excludes stationary and slower wave features such as mountain waves. There-
180 fore, the analysis here only include high frequency gravity waves. The influence of TD
181 method is discussed in details in Support Information. Secondly, in order to find intrin-
182 sic wave parameters, the Doppler shift correction is applied after the observed (ground-
183 based) wave parameters are calculated. This is different from the conventional method
184 where the raw images were shifted opposite to the background wind before the wave pa-
185 rameters were estimated (Z. Li et al., 2011). This is to avoid any possible image dete-
186 rioration in shifting images. Thirdly, some pixels on the imager CCD were broken af-
187 ter Nov 2012 and a black band about 20 km wide showed up in the side of airglow im-
188 ages. The bad pixels were cropped which makes the images used for wave extraction slightly
189 smaller than previously used. This brings little difference in extracted wave parameters
190 since the size of remaining images is still much larger than typical wavelength.

191 **4 Results**

192 **4.1 Wave Parameters Statistics**

193 Figure 1 demonstrates the histograms (frequency) for typical gravity wave param-
194 eters, including horizontal wavelength, vertical wavelength, observed phase speed, intrin-
195 sic phase speed, intrinsic period and wave amplitude. The bin sizes for them are 2.5 km,
196 2.5 km, 5 ms^{-1} , 5 ms^{-1} , 1 min and 0.1%, respectively. The normalized frequencies are
197 divided by the bin width to make the histograms akin to probability density functions.
198 In order to evaluate the robustness of the histogram, Bootstrapping method is used to
199 estimate the 95% confidence intervals for each frequency. In this study, the number of
200 identified waves are more than 60000, the histograms are robust as indicated by small
201 statistical uncertainties. The horizontal wavelengths of most waves are less than 100 km

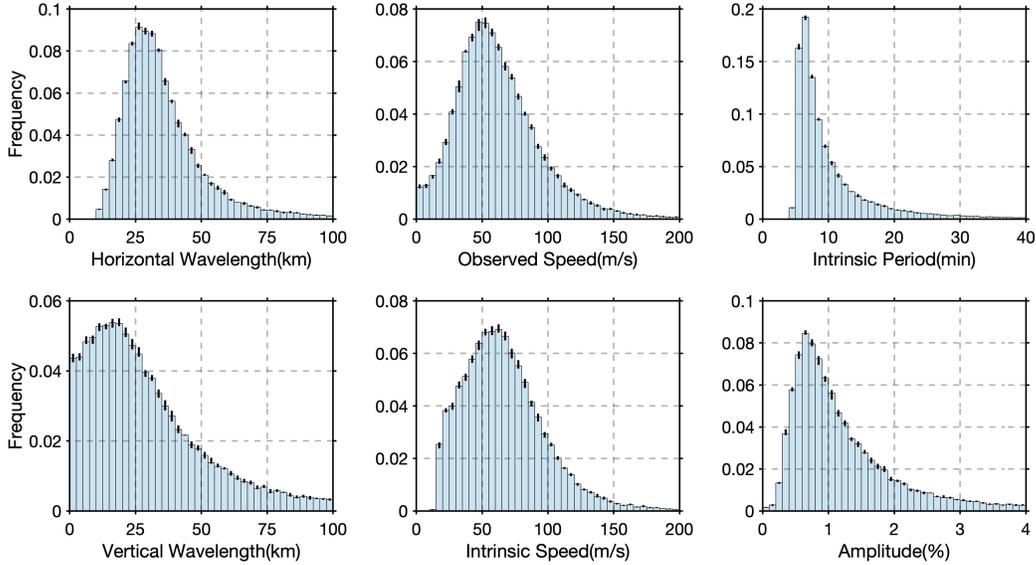


Figure 1. Histograms of gravity wave parameters (from top to bottom, left to right), horizontal wavelength, vertical wavelength, observed phase speed, intrinsic phase speed, period and relative intensity. Small vertical solid lines on top of each bar indicate the 95% confidence interval for each frequency.

202 with peaks near 20–30 km. The vertical wavelengths are mostly larger than 10 km and
 203 with peaks near 15–25 km range. These wavelengths are similar to those found in Maui
 204 (Z. Li et al., 2011) and other sites (Taylor et al., 1993; Nakamura et al., 1999; J. H. Hecht
 205 et al., 2004; Dou et al., 2010). Due to the cancellation effects of wave perturbations in
 206 airglow layer (Liu & Swenson, 2003), waves with vertical wavelength smaller than the
 207 thickness of airglow layer will be greatly attenuated in airglow images. As indicated by
 208 the calculation of vertical wavelength, most of the waves (84%) identified from airglow
 209 images are freely-propagating waves (with positive vertical wavenumber) of which most
 210 of waves (81%) has vertical wavelength larger than 10 km. The calculation of vertical
 211 wavelength requires background temperature which is retrieved from an empirical model
 212 instead of realistic observations. So the distribution of vertical wavelength is treated as
 213 reliable only in climatological and statistical perspective. When daily variations are con-
 214 sidered, discrepancies are expected. Those waves with very short vertical wavelength (≤ 10
 215 km) might be due to inaccuracy of the model data. The observed (ground-based) hor-
 216 izontal phase speeds peak near 45–55 ms^{-1} , while intrinsic horizontal phase speeds peak
 217 near 60–70 ms^{-1} , which indicates waves mostly propagate against background winds.

218 For the wave intrinsic period, the short-period (high-frequency) waves dominate with pe-
 219 riod mostly less than 10 min, with peak near 5–6 min. Due to the fact that most grav-
 220 ity waves propagate against the background wind, the waves are Doppler-shifted to higher
 221 intrinsic frequency and large vertical wavelength, which makes high-frequency waves more
 222 likely to be observed in airglow images. The wave induced emission intensities are less
 223 than 2–3% and peak near 0.5–0.6%.

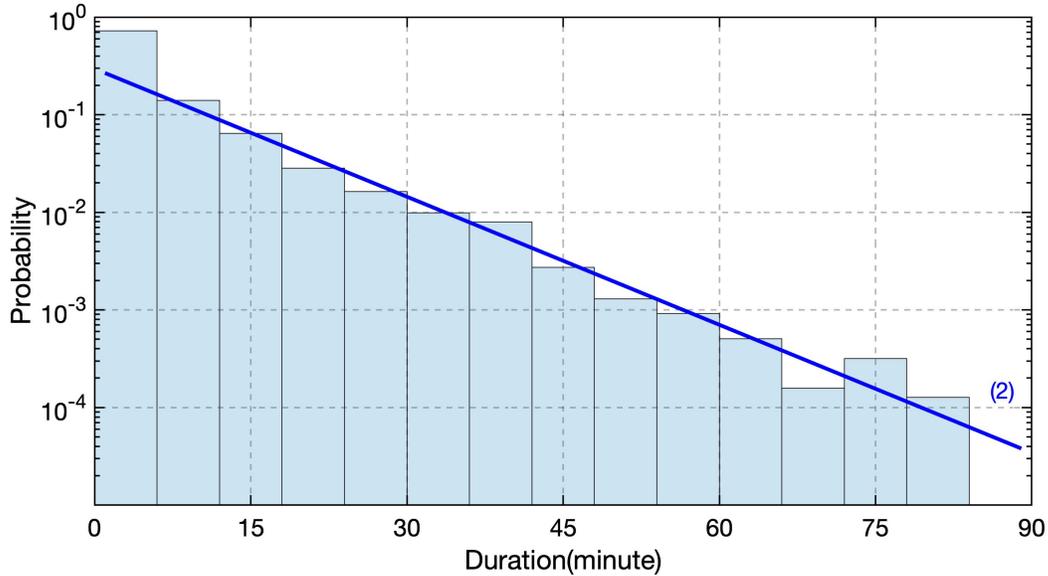


Figure 2. Histogram of the wave event duration. Thick straight lines are from least square fitting. The numbers in the parenthesis at the tails of histograms are the number of wave events at corresponding probabilities.

224 There exists coherent and persistent ‘wave events’ lasting longer than the minimum
 225 duration of a ‘wave’, which is the time of a set of three consecutive images. Hereafter,
 226 a ‘wave event’ refers to a coherent gravity wave composed of several consecutive ‘waves’,
 227 which refers to the wave identified from a set of three images. Complete wave events were
 228 distinguished by identifying consecutive waves with similar parameters, including prop-
 229 agation direction, wavelength and period. Horizontal propagation azimuth and wavelength/wavenumber
 230 were chosen as the primary criteria because they are directly retrieved from 2-D airglow
 231 images. After some tentative tests, 15° and 0.001 km^{-1} are chosen as threshold values.
 232 After the wave event detection was implemented, most waves were identified as part of
 233 a persistent wave event, the remaining waves that do not belong to any wave events were
 234 treated as isolated and associated with the minimum duration.

As shown in Figure 2, the probability density function of wave event duration mostly follows an exponential distribution, i.e. a straight line in semi-log coordinate. The wave events associated with the minimum duration fall in the first bin. The longest duration identified from the data is about 80 min. But it is very rare with only 2 wave events identified in more than 6 years. In order to obtain the mathematical function of the probability distribution, a least-square fitting is applied on the histograms based on following formula: $y = \frac{1}{\tau_0} \exp\left(-\frac{x}{\tau_0}\right)$ of which τ_0 and τ_0^2 are the mean and variance for exponential distributions. The fitting was done in semi-log coordinates that a straight line was fitted to find out the slope ($-1/\tau_0$). Finally, τ_0 is determined as 9.22 min with a 95% confidence interval of 8.28–10.16 min. Theoretically, the mean duration is projected to be 9.22 min for all wave events. However, the actual mean duration of all waves events, including those events with minimum duration, are calculated to be 7.6 min. This is due to the probability of the minimum duration has some derivation from the exponential distribution. Multiple factors such as possible wave breaking, wave packet traversing the imager field-of-view, and wave source characters could contribute to this observed distribution of wave duration. With limited information especially about the background atmosphere status, it is hard to deduce the possible mechanisms that would result in this distribution. Further modeling studies are needed to investigate it in depth.

4.2 Propagation Direction

The distribution of wave characteristics such as wavelength, period and phase speed does not vary much with seasons. However, the preferential wave propagation directions shows clear seasonal dependence. The distribution of wave propagation and corresponding background wind directions are shown by the histogram in polar coordinate in Figure 3. The histograms are organized by calendar month, four rows are austral summer, fall, winter and spring. There are about 2000–5000 waves identified in each calendar month. Overall, gravity waves tend to propagate against background wind especially during summer and winter time. In summer time (Dec to Feb), the dominant wave propagation direction is mostly southward/polar-ward while the background wind is northward. In winter time (Jun to Aug), the dominant wave propagation direction is northward/equatorward while the winds are southward or southeastward. In spring, the preferential directions show a tendency of transition from northward to southward. Opposite transition

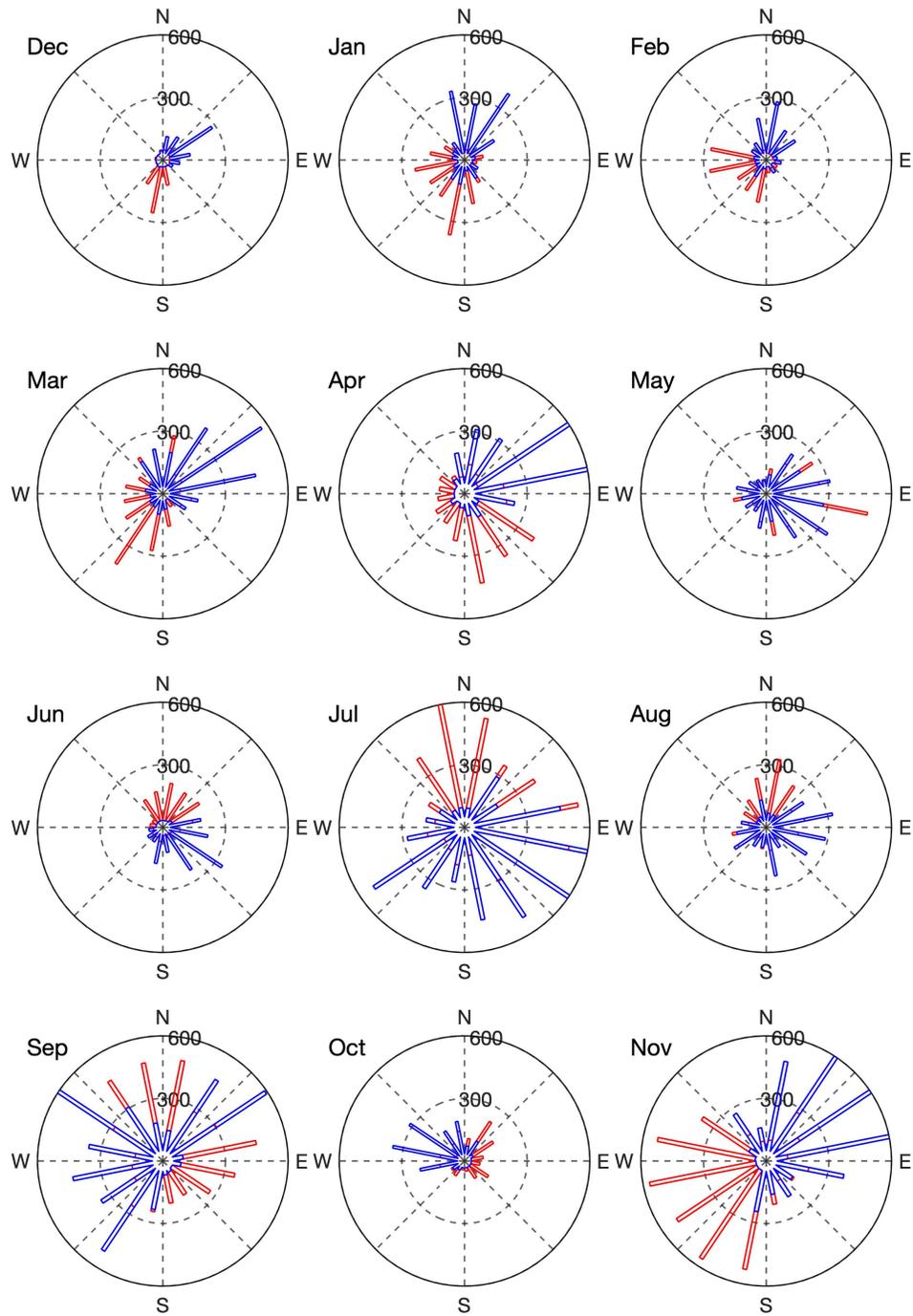


Figure 3. Histograms of (red) wave propagation direction and (blue) background wind direction in each calendar month at a 22.5° azimuth angle bin. The numbers (300, 600) at different radii are the number of waves.

266 can be found in fall. The preferential propagation direction are contributed by multi-
 267 ple factors including potential wave source locations and background wind filtering.

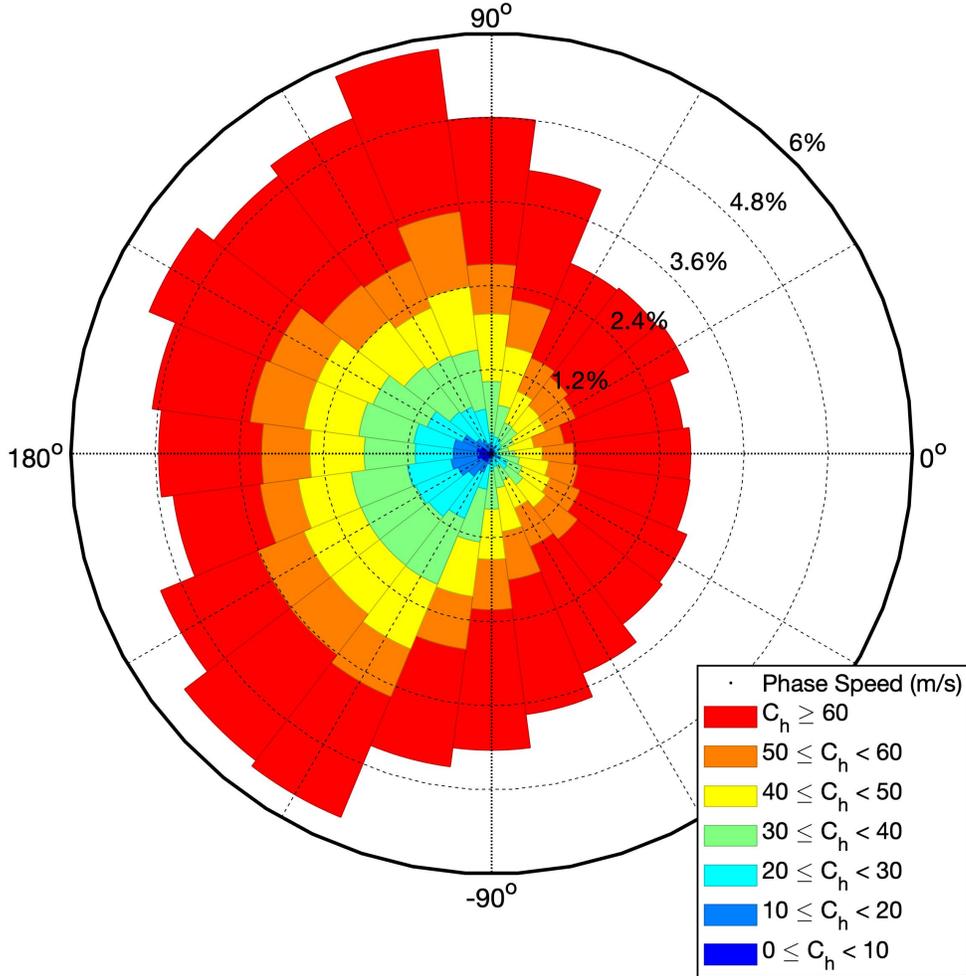


Figure 4. Histogram of azimuth differences between gravity wave propagation and background wind directions for the waves of different observed phase speed. The numbers at different radii are percentage of waves.

268 In order to evaluate the relationship between propagation and background wind
 269 directions of each individual wave, the azimuth angle differences between wave propa-
 270 gation and background wind directions are calculated for waves with different phase speed.
 271 As shown in Figure 4, the azimuth angle differences are mostly toward the hemisphere
 272 of 180°. However, the distribution have some dependence on the phase speed. For waves
 273 with observed phase speed less than 20 ms^{-1} , it is prominent that the azimuth angle dif-

274 ferences are highly clustered around 180° . This means those waves mostly propagate against
 275 the winds which is an indicator of critical layer filtering of waves propagating along the
 276 winds if any. The distribution around 180° becomes less concentrated for larger observed
 277 phase speeds. For waves with phase speed between 20 and 40 ms^{-1} , they mostly prop-
 278 agate toward opposite direction with background wind but with a boarder range. For
 279 those faster waves with phase speed larger than 50 ms^{-1} , their propagation shows lit-
 280 tle dependence on background wind and can propagate at any directions with respect
 281 to background wind. The monthly mean horizontal winds in the OH airglow layer are
 282 around $30\text{--}40 \text{ ms}^{-1}$. Background winds would be able to filter out waves with observed
 283 phase speed similar or smaller than wind speed. However, background winds tend to ex-
 284 ert less influence on those faster waves through critical layers filtering. Besides the ef-
 285 fects of critical layer, waves propagate along the background winds are Doppler-shifted
 286 to smaller vertical wavelength thus larger shear may occur to make waves more easily
 287 to break down due to instability. For faster waves, there could be other factors contribut-
 288 ing to the preferential propagation direction.

289 4.3 Background Filtering

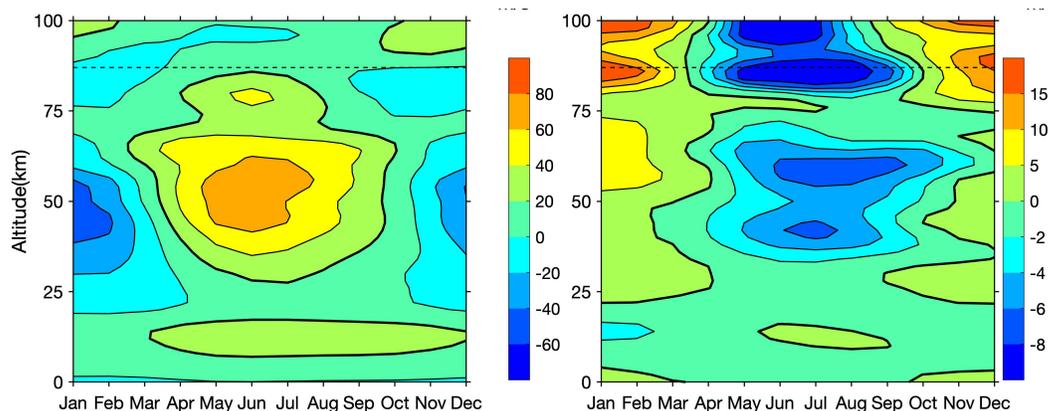


Figure 5. Monthly mean (left) zonal and (right) meridional winds retrieved from HWM-14 model, averaged between 00:00 to 06:00 UT at location of ALO, horizontal dashed lines indicate the mean altitude of OH airglow layer.

290 Background atmosphere where the waves propagate through plays an important
 291 role in controlling the prevailing propagation direction through critical-layer filtering. When
 292 gravity waves reach a layer where wave observed phase speed equals background wind

293 speed, waves will be absorbed or filtered. The Doppler-shifted or intrinsic frequency $\hat{\omega}$
 294 can be related to observed frequency ω by

$$\hat{\omega} = \omega \left(1 - \frac{\bar{u} \cos \phi + \bar{v} \sin \phi}{c} \right), \quad (2)$$

295 of which the term $\bar{u} \cos \phi + \bar{v} \sin \phi$ is the background wind (\bar{u}, \bar{v}) projected to wave prop-
 296 agation direction. ‘Blocking diagram’ (Taylor et al., 1993; Medeiros et al., 2003) is in-
 297 troduced to demonstrate the ‘forbidden zone’ of gravity waves, i.e., the range of phase
 298 speed c and propagation azimuth angle ϕ of waves that would be filtered out in certain
 299 background wind profiles where $\hat{\omega} \leq 0$ is satisfied.

300 Currently, there is no complete observations of atmospheric winds from source level
 301 to airglow layer near ALO. We turn to the model winds retrieved from Horizontal Wind
 302 Model-14 (HWM-14) (Drob et al., 2015), which reasonably reproduces climatological winds.
 303 Figure 5 shows the monthly mean zonal and meridional winds at ALO. Only winds be-
 304 tween 00:00 and 06:00 UT are selected to match the timing of airglow images at night.
 305 At ALO, zonal winds in the stratosphere are eastward in austral winter and westward
 306 in summer, with largest magnitudes exceeding $\pm 60 \text{ ms}^{-1}$. Meridional winds magnitudes
 307 are much smaller and are mostly polar-ward but equator-ward in summer above 50 km.
 308 In Figure 6, ‘blocking diagrams’ were plotted for each month using the monthly aver-
 309 aged wind profiles from HWM-14 at ALO. They represent the effects of critical layer fil-
 310 tering on gravity waves accumulated in the altitude range from 15 km that is above most
 311 convective activities to 87 km that is the peak altitude of OH airglow. The observed phase
 312 speed and propagation direction of all waves are demonstrated by scattered dots. The
 313 ‘forbidden zones’ of gravity waves predicted by critical layer filtering theory are mostly
 314 along west and east directions due to much larger amplitudes of zonal wind component
 315 especially in stratosphere. As shown in Figure 6, a lot of waves can be found in the pre-
 316 dicted ‘forbidden zones’ in some months. This might be due to the discrepancies between
 317 modeled and realistic winds. However, areas around certain smaller phase speeds and
 318 directions show up as hollows in the scattered plots in multiple months such as May, Jun,
 319 Oct and Nov. The absence of these waves indicates the effects of critical layer filtering,
 320 as they are filtered out by the realistic background winds that are not reflected in an em-
 321 pirical wind model.

322 Here, critical layers filtering predicted by HWM-14 model can not explain the wave
 323 propagation direction well. The monthly mean winds retrieved from HWM-14 cannot

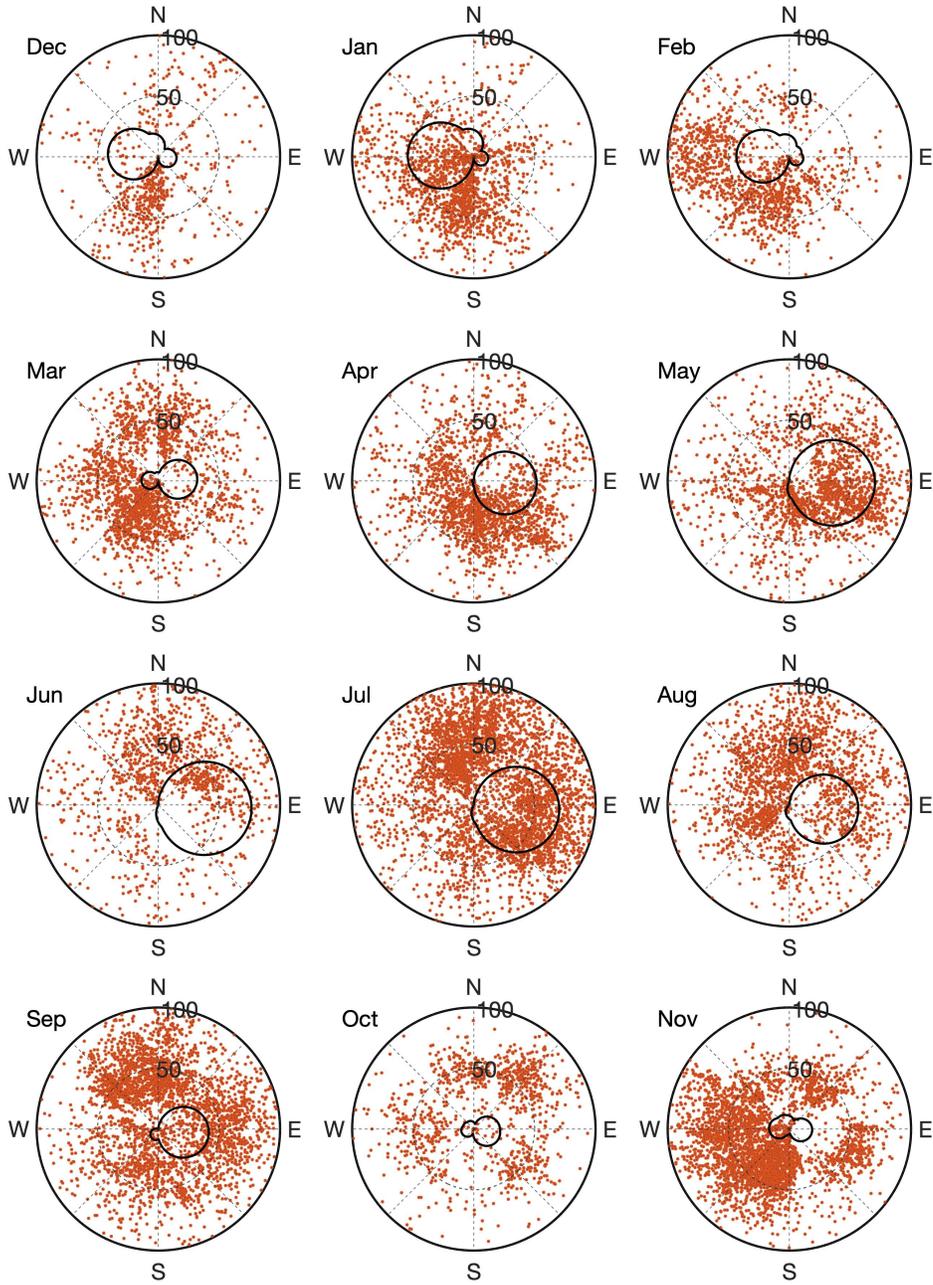


Figure 6. Scatter plots of observed phase speed ($0\text{--}100\text{ ms}^{-1}$). One dot represents an identified wave with certain phase speed and propagation direction. Small amount of waves with phase speed larger than 100 ms^{-1} are not included here. Area inside the solid black lines are the ‘forbidden zone’ predicted by critical layer filtering theory.

324 capture the short-period variation of the real winds such as tidal influences, day-to-day
325 variability and any waves that have period longer than gravity waves that are observed
326 by airglow imager. Time-varying background winds reduce the effects of critical layer
327 filtering because a lot of waves have less time to interact with varying winds and/or changes
328 of ground phase speed that can be critically filtered (Heale & Snively, 2015). This is es-
329 pecially true for the waves observed by airglow imagers that are mostly high-frequency,
330 with periods less than 15 min.

331 **4.4 Convective Wave Sources**

332 Multiple hot-spots of gravity waves have been revealed by many previous studies
333 using satellite observations and models (Jiang et al., 2004; Geller et al., 2013; Hoffmann
334 et al., 2013, 2016) over the South America and southeast Pacific. Convection and oro-
335 graphic sources were found to be two most likely ones around this region. Vadas, Tay-
336 lor, et al. (2009) used ray-tracing to locate the potential wave sources and found out the
337 convection is likely the sources of mesospheric gravity waves observed by an OH airglow
338 imager in Brazil.

339 Outgoing Longwave Radiation (OLR) is a measure of the amount of energy emit-
340 ted to space from earth's surface, including oceans and atmosphere. OLR values are of-
341 ten used as a good proxy for convection in tropical and subtropical regions. In general,
342 smaller values indicate stronger convective activities because they are associated with
343 high cloud tops with lower temperature. Interpolated monthly mean OLR data (Liebmann
344 & Smith, 1996) was acquired from Physical Sciences Laboratory of NOAA. The data has
345 a 2.5° by 2.5° spatial resolution and global coverage. In Figure 7, the OLR intensities
346 of each month are averaged between 2009 and 2014, and shown by colors with a reversed
347 color-scale. The histogram of propagation direction of gravity waves is also shown on the
348 map. In each calendar month, the occurrence frequency of gravity waves is quantified
349 as the ratio of the number of identified waves to the number of images, which is a proxy
350 of the relative likelihood of occurrence of gravity waves in each months. As shown in Fig-
351 ure 7, the occurrence frequencies are high over winter and early spring (Jun to Oct) and
352 low over summer and fall (Feb to May). Regarding of the wave propagation direction,
353 the occurrence frequency is generally higher when convection is identified at closer dis-
354 tance, especially within 250 km.

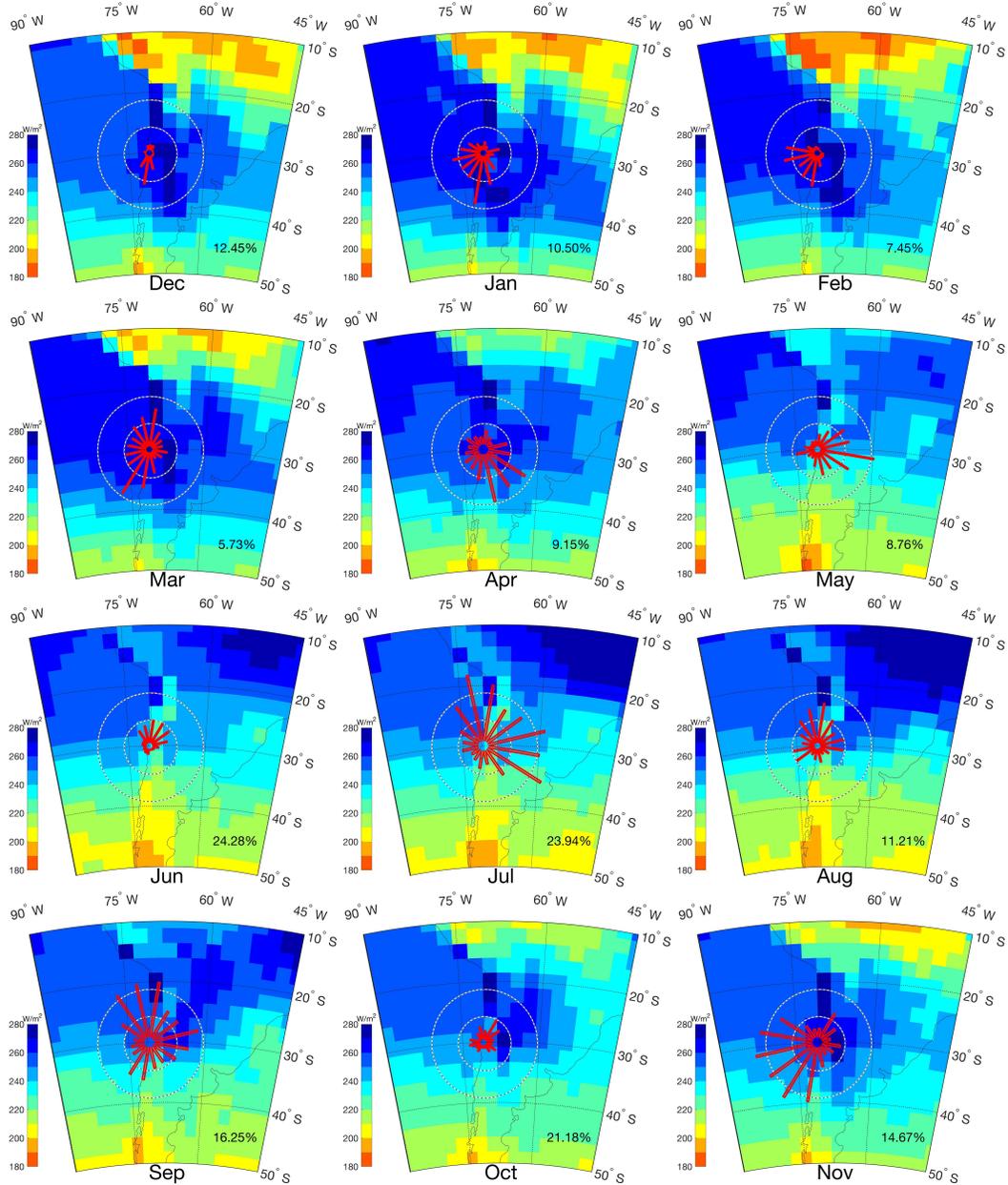


Figure 7. Monthly mean OLR (color shading) overlapped with histograms (in red) on the map showing the coastline of South America. The polar histograms show the propagation direction of gravity waves, same as red in Figure 3. The radius of the two circles centered at ALO represent 250 and 500 waves for histogram. The numbers in the lower right corner of each panel are the wave occurrence frequencies, see text for definition. The color-scale of OLR values is reverse so warmer(red) colors indicate stronger convective activities.

355 On the continent of South America, there are a few notable areas with strong con-
 356 vection on the north and south side, including Amazon Basin in the tropics and La Palate
 357 Basin in the subtropics ($\sim 30^\circ\text{S}$). They provide a large amount of moisture and energy
 358 for deep convection and precipitation (Insel et al., 2010; Romatschke & Houze, 2010).
 359 ALO is located on the west side of Andes which has an average elevation of 4–5 km and
 360 blocks warm moist air from the east. The convective activities indicated by OLR inten-
 361 sity show clear seasonal variations and high correlations with the wave propagation di-
 362 rection. From late spring to early fall (Nov to Mar), strong convective activities show
 363 up in the Amazon Basin and expand to a large area. Some of these convections, espe-
 364 cially in summer, are close to ALO within several hundred km. The observed waves have
 365 a clear preference of southwestward propagation but with lower occurrence frequencies.
 366 In winter (Jun to Oct), the closest and strongest convective source is over the Pacific Ocean
 367 to the southwest of ALO and coast area to the south of ALO, during which the wave prop-
 368 agation is clearly northeastward or northward. The convection is much closer to ALO
 369 during this time, the occurrence frequencies are highest in a year. From spring to early
 370 fall (Sep to Apr), there is also a strong and localized convective source over La Palate
 371 Basin to the east and southeast of ALO. This feature is not evident in OLR, but was demon-
 372 strated by precipitation and lighting data (Rasmussen et al., 2014, 2016). The wave prop-
 373 agation shows a preference of westward or northwestward in some spring and summer
 374 months (Sep, Nov, Jan, Feb, Mar), consistent with this wave source.

375 4.5 Momentum Fluxes

376 Figure 8 shows the monthly mean zonal and meridional gravity wave momentum
 377 fluxes ($\langle u'w' \rangle$ and $\langle v'w' \rangle$) with zonal and meridional background winds averaged over
 378 22:00–06:00 UT in the OH airglow layer. Overall, the zonal and meridional momentum
 379 fluxes have the magnitudes of several m^2s^{-2} with meridional component slightly larger
 380 than zonal one. The mean momentum flux magnitudes are very small especially consid-
 381 ering the low density in middle atmosphere, each individual wave might not exert large
 382 influence on the background. However, a large amount of these waves carrying little mo-
 383 mentum flux still show evident effects on the background. Both momentum flux com-
 384 ponents tend to toward the opposite direction of background winds. Zonal momentum
 385 flux is mostly westward and zonal wind is mostly eastward. There are some intra-seasonal
 386 variations in zonal momentum flux and wind. The opposite directionality between merid-

387 ional momentum flux and wind is more distinct. Meridional momentum flux shows a clear
 388 annual oscillation with northward maximum near austral winter time and southward max-
 389 imum in summer. Gravity wave momentum fluxes at mesopause altitude are affected by
 390 both wave sources in the lower atmosphere and critical layer filtering by the mean flow
 391 between the sources and mesopause (Z. Li et al., 2011). The change of momentum flux
 392 is related to the variation of the location of primary wave sources, which are mostly lo-
 393 cate at east and northeast of ALO in summer and south in winter. The momentum flux
 394 was estimated based on equation 1 using extra information from empirical models, it is
 395 also counted as reliable in a climatological perspective.

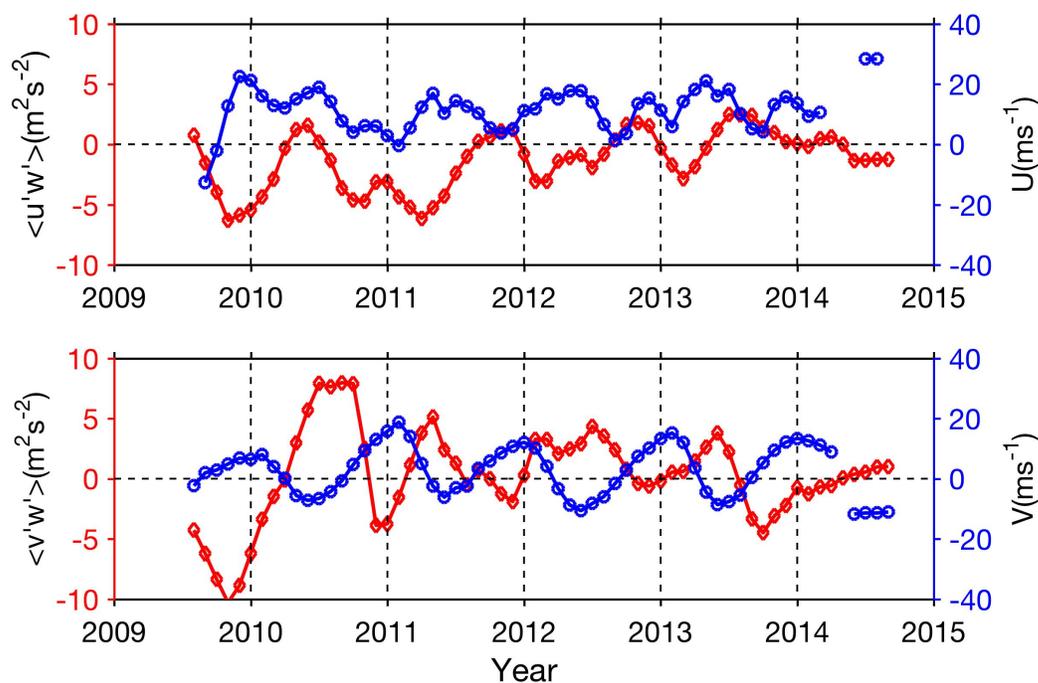


Figure 8. Monthly mean (top) zonal and (bottom) meridional (left axis, red) momentum flux and (right axis, blue) wind from 2009 to 2014. The zonal and meridional winds are averaged between 22:00 to 06:00 UT.

396 5 Discussions

397 In the wave extraction method, a set of three consecutive images are used to ob-
 398 tain two TD images for co-spectral analysis. For the airglow imager at ALO, only OH
 399 airglow images were captured with a 1-min integration time before 25 Aug 2011 and OH
 400 and OI images were captured alternately with 1-min and 1.5-min integration times af-

401 terwards. A gravity wave event has to last 3 min and 6 min to be identified in these two
402 different configurations (see Figures S1(a) and S1(b)), respectively. For the long term
403 climatology study, in order to minimize the discrepancies due to different integration times,
404 a trade-off made in the data processing is to skip every other image for the time period
405 when only OH airglow images were captured (see Figure S1(c)). The detailed discus-
406 sions regarding the influences of the TD method and different integration times are pre-
407 sented in the Supporting Information.

408 In this study, we focus on convection as a primary candidate of wave sources. Clear
409 correlations are revealed between convective activities and observed waves characteris-
410 tics including wave occurrence frequency and propagation direction. The distance be-
411 tween possible wave source area and ALO where the waves were observed is an impor-
412 tant factor. The high-frequency gravity waves tend to propagate upward in a more steep
413 path and thus likely to have a nearby source located within 100–200 km range (Vadas,
414 Yue, et al., 2009). As shown in Figure 7, some intense convective activities in summer
415 time are more than 1000 km away from ALO such as the Amazon Basin in the north and
416 northeast and the area at the east coast of South America. Simulation studies have shown
417 that long-range propagation of gravity waves in MLT region is possible through ducted
418 propagation (J. H. Hecht et al., 2001; Snively & Pasko, 2008; Snively et al., 2013; Heale
419 et al., 2014). However, airglow images retrieved from a single layer could not distinguish
420 whether these waves are ducted. Some of the waves propagating southward and south-
421 westward in late spring and summer (Nov to Mar) are possibly ducted considering the
422 long distance between waves sources and ALO. However, there also exists other wave sources
423 nearby beyond the convection, such as secondary wave generation.

424 Southern Andes has been reported in many studies as a hot-spot of orographically-
425 generated gravity waves. Satellite observations and modeling reveal highest wave occur-
426 rence at mid-fall to mid-spring (Apr-Oct) (Hoffmann et al., 2013; Alexander et al., 2015;
427 Hoffmann et al., 2016). The hot-spots concentrate around the west coast of South Amer-
428 ica along the ridge of Andes, extending from 30°S to the tip at 60°S. It is found that the
429 wave activities are closely correlated with lower-level zonal flow over topography around
430 winter time. Large amount of northward propagating waves are observed in this study
431 with highest occurrence frequencies around the same period over ALO. These waves orig-
432 inate from southern area where the core of the hot-spot is located. Even though these
433 mesospheric high-frequency gravity waves observed by the airglow imager are not directly

434 generated by orographic sources, there also exist the possibility of secondary wave gen-
435 eration (Vadas et al., 2003; Bossert et al., 2017) due to nonlinear interaction or wave dis-
436 sipation. Orographic gravity waves have near-zero ground phase speed. They are absorbed
437 and dissipate near the zero-wind layer. As revealed by the monthly mean horizontal winds
438 in Figure 5, there exists zero-wind layers in stratosphere beneath the OH airglow layer
439 between Apr and Sep. Even though the climatological model winds do not capture the
440 short term variations, the background atmosphere is still favorable for the breaking of
441 mountain waves and resulting secondary wave generation in those months.

442 The discussion of wave-background interaction below the airglow layer altitude is
443 limited because of the use of climatological model winds, which do not duplicate the re-
444 alistic winds and fully explain the observed results. Further studies utilizing more re-
445 alistic reanalysis data would be beneficial to evaluate the critical layer filtering, and pro-
446 vide proper background conditions for a ray-tracing modeling study to locate the source
447 area and identify possible ducted propagation. Any waves that cannot be traced back
448 all the way to the troposphere might be accounted to aforementioned mechanisms.

449 **6 Summary and Conclusions**

450 The long-term dataset from 2009 to 2014, retrieved by an all-sky airglow imager
451 at ALO, is used to investigate the characteristics of high-frequency quasi-monochromatic
452 gravity waves. The typical horizontal wavelengths are around 20–40 km and ground-based
453 horizontal phase speeds are between 40 and 70 ms^{-1} . The intrinsic periods of gravity
454 waves cluster around 4–10 min. However, those wave parameters are jointly limited by
455 the ‘observation filter’ effort of the airglow imager, and the images processing method.
456 And some parameters such as vertical wavelength and momentum flux are estimated us-
457 ing information of empirical models, they are reliable only in statistical and climatolog-
458 ical perspective. The observed gravity waves tend to propagate against the local back-
459 ground wind in most months and also show strong seasonal dependence in the prefer-
460 ential propagation direction. The duration of coherent ‘wave events’ is found to follow
461 an exponential distribution, with a mean duration about 7–9 min. It is not yet fully un-
462 derstood the mechanism that leads to such a distribution. The mean wave momentum
463 flux estimated from airglow data has a much smaller magnitude of several m^2s^{-2} com-
464 pared to those distinct waves investigated in case studies. However, these waves asso-
465 ciated little momentum flux contribute significantly to alter background collectively be-

466 cause of their much higher probability of presence (Cao & Liu, 2016). The wave momen-
467 tum flux tends to be toward opposite direction of background winds in airglow layer, es-
468 pecially in meridional direction. These results are consistent with previous studies based
469 on airglow images from other mid-latitude sites such as Fort Collins, CO (20°N) (Y. Tang
470 et al., 2014), Maui, HI (20°N) (Z. Li et al., 2011), Shigaraki, Japan (35°N) (Nakamura
471 et al., 1999) and Urbana, IL (40°N) (J. H. Hecht et al., 2001). In addition to the sim-
472 ilarities, the presented wave characteristics especially the preferential propagation direc-
473 tion and occurrence frequency show high correlation with localized environment. The
474 new results add information of high-frequency gravity waves in the mid-latitude of South-
475 ern Hemisphere that is beneficial to the understanding of gravity waves in global scale.

476 It has been suggested that source locations where the waves are generated and back-
477 ground wind where the waves propagate through, cooperatively determine the observed
478 wave characteristics in MLT region. ALO is located at a place near or within the zone
479 of influence of several remarkable convection sources. During the austral summer, the
480 convection over Amazon Basin is dramatically strong and expands over a vast area. Those
481 waves with southwestward propagation direction could originate from there and might
482 be associated with ducted long range propagation. Even the stratospheric zonal wind
483 are mostly westward in this season, the wave sources overwhelm the background wind
484 filtering effect in determining the directionality of wave propagation direction. In win-
485 ter time, the closer convection is over the Pacific Ocean or coast area to the south of ALO,
486 this could mostly explain the northeastward and northward preferential propagation di-
487 rection. Critical-layer filtering predicted by model winds could not explain the propa-
488 gation direction preference well in most months. However, some hollow zones exist in
489 ‘block diagram’ that indicate the efforts of filtering of the waves with slower velocities.
490 The opposite direction of gravity waves and local background wind also indicates the fil-
491 tering effects of critical layer on slower waves. The results of this study do not show that
492 the anisotropy of propagation direction was entirely due to wave filtering by stratospheric
493 winds (Taylor et al., 1993; Medeiros et al., 2003) as the background winds exert effects
494 mainly on slower waves. The locations of wave sources and where they are observed play
495 a more important role in shaping the prevailing wave propagation. In this study, the re-
496 lationship between the observed waves and potential sources is described mostly qual-
497 itatively. The relationship between the strength of the convective activities and wave oc-
498 currence frequency, wave amplitude can be described with some dependence on distance.

499 In other words, it could be possible to quantify the influential area of certain convective
 500 activities. This would provide some insight regarding a simplified assumption in the grav-
 501 ity wave parameterization that the horizontal propagation of waves are neglected.

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 510 from <https://zenodo.org/record/4267147>. The OLR data is provided by the NOAA/OAR/ESRL
 511 PSL, Boulder, Colorado, USA from their website at <https://psl.noaa.gov/>. The NRLMSISE-
 512 00 and HWM-14 models data are generated by functions embedded in Matlab Aerospace
 513 Toolbox (<https://www.mathworks.com/products/aerospace-toolbox.html>).

514 **References**

- 515 Alexander, P., de la Torre, A., Schmidt, T., Llamedo, P., & Hierro, R. (2015).
 516 Limb sounders tracking topographic gravity wave activity from the strato-
 517 sphere to the ionosphere around midlatitude Andes. *Journal of Geophysi-
 518 cal Research: Space Physics*, *120*(10), 9014-9022. Retrieved from [https://](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2015JA021409)
 519 agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2015JA021409 doi:
 520 [10.1002/2015JA021409](https://doi.org/10.1002/2015JA021409)
- 521 Bossert, K., Fritts, D. C., Pautet, P. D., Taylor, M. J., Williams, B. P., & Pendle-
 522 ton, W. R. (2014). Investigation of a mesospheric gravity wave ducting
 523 event using coordinated sodium lidar and Mesospheric Temperature Mapper
 524 measurements at ALOMAR, Norway (69°N). *J. Geophys. Res.*, *119*(16),
 525 9765-9778. doi: [10.1002/2014JD021460](https://doi.org/10.1002/2014JD021460)
- 526 Bossert, K., Kruse, C. G., Heale, C. J., Fritts, D. C., Williams, B. P., Snively,
 527 J. B., ... Taylor, M. J. (2017). Secondary gravity wave generation over
 528 New Zealand during the DEEPWAVE campaign. *Journal of Geophysi-
 529 cal Research: Atmospheres*, *122*(15), 7834-7850. Retrieved from <https://>

- 530 agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2016JD026079 doi:
531 10.1002/2016JD026079
- 532 Cao, B., Heale, C. J., Guo, Y., Liu, A. Z., & Snively, J. B. (2016). Observation
533 and modeling of gravity wave propagation through reflection and critical layers
534 above Andes Lidar Observatory at Cerro Pachón, Chile. *Journal of Geo-*
535 *physical Research: Atmospheres*, *121*(21), 12,737–12,750. Retrieved from
536 <http://dx.doi.org/10.1002/2016JD025173> doi: 10.1002/2016JD025173
- 537 Cao, B., & Liu, A. Z. (2016). Intermittency of gravity wave momentum flux in the
538 mesopause region observed with an all-sky airglow imager. *Journal of Geophys-*
539 *ical Research: Atmospheres*, *121*(2), 650–663. Retrieved from [http://dx.doi](http://dx.doi.org/10.1002/2015JD023802)
540 [.org/10.1002/2015JD023802](http://dx.doi.org/10.1002/2015JD023802) doi: 10.1002/2015JD023802
- 541 Dou, X., Li, T., Tang, Y., Yue, J., Nakamura, T., Xue, X., ... She, C.-Y. (2010).
542 Variability of gravity wave occurrence frequency and propagation direction
543 in the upper mesosphere observed by the OH imager in Northern Colorado.
544 *Journal of Atmospheric and Solar-Terrestrial Physics*, *72*(5–6), 457–462.
545 Retrieved from [http://www.sciencedirect.com/science/article/pii/](http://www.sciencedirect.com/science/article/pii/S1364682610000131)
546 [S1364682610000131](http://www.sciencedirect.com/science/article/pii/S1364682610000131) doi: <http://dx.doi.org/10.1016/j.jastp.2010.01.002>
- 547 Drob, D. P., Emmert, J. T., Meriwether, J. W., Makela, J. J., Doornbos, E., Conde,
548 M., ... Klenzing, J. H. (2015). An update to the Horizontal Wind Model
549 (HWM): The quiet time thermosphere. *Earth and Space Science*, *2*(7), 301–
550 319. Retrieved from [https://agupubs.onlinelibrary.wiley.com/doi/abs/](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2014EA000089)
551 [10.1002/2014EA000089](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2014EA000089) doi: <https://doi.org/10.1002/2014EA000089>
- 552 Ejiri, M. K., Shiokawa, K., Ogawa, T., Igarashi, K., Nakamura, T., & Tsuda,
553 T. (2003). Statistical study of short-period gravity waves in OH and OI
554 nightglow images at two separated sites. *J. Geophys. Res.*, *108*(D21). doi:
555 doi:10.1029/2002JD002795
- 556 Franke, S. J., Chu, X., Liu, A. Z., & Hocking, W. K. (2005). Comparison of meteor
557 radar and Na Doppler lidar measurements of winds in the mesopause region
558 above Maui, Hawaii. *J. Geophys. Res.*, *110*(D9). doi: 10.1029/2003JD004486
- 559 Fritts, D. C., & Alexander, M. J. (2003). Gravity wave dynamics and effects in the
560 middle atmosphere. *Rev. Geophys.*, *41*(1). doi: 10.1029/2001RG000106
- 561 Fritts, D. C., Wang, L., Taylor, M. J., Pautet, P.-D., Criddle, N. R., Kaifler, B.,
562 ... Liley, B. (2019). Large-Amplitude Mountain Waves in the Meso-

- 563 sphere Observed on 21 June 2014 During DEEPWAVE: 2. Nonlinear Dy-
 564 namics, Wave Breaking, and Instabilities. *Journal of Geophysical Re-*
 565 *search: Atmospheres*, 124(17-18), 10006-10032. Retrieved from [https://](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2019JD030899)
 566 agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2019JD030899 doi:
 567 <https://doi.org/10.1029/2019JD030899>
- 568 Geller, M. A., Alexander, M. J., Love, P. T., Bacmeister, J., Ern, M., Hertzog, A.,
 569 ... Zhou, T. (2013). A comparison between gravity wave momentum fluxes
 570 in observations and climate models. *Journal of Climate*, 26(17), 6383–6405.
 571 Retrieved from <http://dx.doi.org/10.1175/JCLI-D-12-00545.1> doi:
 572 [10.1175/JCLI-D-12-00545.1](http://dx.doi.org/10.1175/JCLI-D-12-00545.1)
- 573 Heale, C. J., & Snively, J. B. (2015). Gravity wave propagation through a verti-
 574 cally and horizontally inhomogeneous background wind. *Journal of Geophys-*
 575 *ical Research: Atmospheres*, 120(12), 5931–5950. Retrieved from [http://dx](http://dx.doi.org/10.1002/2015JD023505)
 576 [.doi.org/10.1002/2015JD023505](http://dx.doi.org/10.1002/2015JD023505) doi: 10.1002/2015JD023505
- 577 Heale, C. J., Snively, J. B., & Hickey, M. P. (2014). Numerical simulation of the
 578 long-range propagation of gravity wave packets at high latitudes. *Journal of*
 579 *Geophysical Research: Atmospheres*, 119(19), 11,116–11,134. Retrieved from
 580 <http://dx.doi.org/10.1002/2014JD022099> doi: 10.1002/2014JD022099
- 581 Hecht, J., Walterscheid, R. L., & Vincent, R. A. (2001). Airglow observations of dy-
 582 namical (wind shear-induced) instabilities over Adelaide, Australia, associated
 583 with atmospheric gravity waves. *J. Geophys. Res.*, 106(D22), 28189-28197.
- 584 Hecht, J. H., Kovalam, S., May, P. T., Mills, G., Vincent, R. A., Walterscheid, R. L.,
 585 & Woithe, J. (2004). Airglow imager observations of atmospheric gravity waves
 586 at Alice Springs and Adelaide, Australia during the Darwin Area Wave Experi-
 587 ment (DAWEX). *J. Geophys. Res.*, 109(D20). doi: 10.1029/2004JD004697
- 588 Hecht, J. H., Walterscheid, R. L., Hickey, M. P., & Franke, S. J. (2001). Climatology
 589 and modeling of quasi-monochromatic atmospheric gravity waves observed over
 590 Urbana Illinois. *J. Geophys. Res.*, 106(D6), 5181–5195.
- 591 Hickey, M. P., & Yu, Y. (2005). A full-wave investigation of the use of a 'cancell-
 592 ation factor' in gravity wave–OH airglow interaction studies. *J. Geophys. Res.*,
 593 110. doi: 10.1029/2003JA010372
- 594 Hindley, N. P., Wright, C. J., Smith, N. D., & Mitchell, N. J. (2015). The southern
 595 stratospheric gravity wave hot spot: individual waves and their momentum

- 596 fluxes measured by COSMIC GPS-RO. *Atmospheric Chemistry and Physics*,
 597 15(14), 7797–7818. Retrieved from [http://www.atmos-chem-phys.net/15/](http://www.atmos-chem-phys.net/15/7797/2015/)
 598 7797/2015/ doi: 10.5194/acp-15-7797-2015
- 599 Hocking, W. K., Fuller, B., & Vandeppeer, B. (2001). Real-time determination of
 600 meteor-related parameters utilizing modern digital technology. *J. Atmos. Sol.-*
 601 *Terr. Phys.*, 63(2-3), 155-169.
- 602 Hoffmann, L., Grimsdell, A. W., & Alexander, M. J. (2016). Stratospheric gravity
 603 waves at Southern Hemisphere orographic hotspots: 2003–2014 AIRS/Aqua
 604 observations. *Atmospheric Chemistry and Physics*, 16(14), 9381–9397. Re-
 605 trieved from <https://www.atmos-chem-phys.net/16/9381/2016/> doi:
 606 10.5194/acp-16-9381-2016
- 607 Hoffmann, L., Xue, X., & Alexander, M. J. (2013). A global view of strato-
 608 spheric gravity wave hotspots located with Atmospheric Infrared Sounder
 609 observations. *Journal of Geophysical Research: Atmospheres*, 118(2), 416–
 610 434. Retrieved from <http://dx.doi.org/10.1029/2012JD018658> doi:
 611 10.1029/2012JD018658
- 612 Insel, N., Poulsen, C. J., & Ehlers, T. A. (2010). Influence of the Andes Mountains
 613 on South American moisture transport, convection, and precipitation. *Climate*
 614 *Dynamics*, 35(7), 1477–1492. Retrieved from [http://dx.doi.org/10.1007/](http://dx.doi.org/10.1007/s00382-009-0637-1)
 615 [s00382-009-0637-1](http://dx.doi.org/10.1007/s00382-009-0637-1) doi: 10.1007/s00382-009-0637-1
- 616 Jiang, J. H., Wang, B., Goya, K., Hocke, K., Eckermann, S. D., Ma, J., ... Read,
 617 W. G. (2004). Geographical distribution and interseasonal variability of tropi-
 618 cal deep convection: UARS MLS observations and analyses. *J. Geophys. Res.*,
 619 209(D03111). doi: doi:10.1029/2003JD003756
- 620 Khomich, V. Y., Semenov, A. I., & Shefov, N. N. (2008). *Airglow as an indicator of*
 621 *upper atmospheric structure and dynamics*. Springer.
- 622 Li, J., Li, T., Dou, X., Fang, X., Cao, B., She, C.-Y., ... Thorsen, D. (2017). Char-
 623 acteristics of ripple structures revealed in OH airglow images. *Journal of*
 624 *Geophysical Research: Space Physics*. Retrieved from [http://dx.doi.org/](http://dx.doi.org/10.1002/2016JA023538)
 625 [10.1002/2016JA023538](http://dx.doi.org/10.1002/2016JA023538) doi: 10.1002/2016JA023538
- 626 Li, T., She, C. Y., Liu, H.-L., Yue, J., Nakamura, T., Krueger, D. A., ... Wang,
 627 S. (2009). Observation of local tidal variability and instability, along
 628 with dissipation of diurnal tidal harmonics in the mesopause region over

- 629 Fort Collins, Colorado (41 °N, 105 °W). *J. Geophys. Res.*, *114*, D06106,
 630 doi:10.1029/2008JD011089. Retrieved from [http://www.agu.org/pubs/
 631 crossref/2009/2008JD011089.shtml](http://www.agu.org/pubs/crossref/2009/2008JD011089.shtml)
- 632 Li, Z., Liu, A. Z., Lu, X., Swenson, G. R., & Franke, S. J. (2011). Gravity
 633 wave characteristics from OH airglow imager over Maui. *J. Geophys. Res.*,
 634 *116*(D22). doi: 10.1029/2011JD015870
- 635 Li, Z., Liu, A. Z., & Sivjee, G. G. (2014). Removing Milky Way from airglow im-
 636 ages using principal component analysis. *J. Atmos. Sol.-Terr. Phys.*, *110-111*,
 637 50-57. doi: 10.1016/j.jastp.2014.01.016
- 638 Liebmann, B., & Smith, C. A. (1996). Description of a complete (interpolated)
 639 outgoing longwave radiation dataset. *Bulletin of the American Meteorological
 640 Society*, *77*(6), 1275–1277. Retrieved from [http://www.jstor.org/stable/
 641 26233278](http://www.jstor.org/stable/26233278)
- 642 Liu, A. Z., & Swenson, G. R. (2003). A modeling study of O₂ and OH airglow per-
 643 turbations induced by atmospheric gravity waves. *J. Geophys. Res.*, *108*(D4).
 644 doi: 10.1029/2002JD002474
- 645 Medeiros, A. F., Takahashi, H., Buriti, R. A., Fachine, J., Wrasse, C. M., & Gobbi,
 646 D. (2007). MLT gravity wave climatology in the South America equatorial
 647 region observed by airglow imager. *Annales Geophysicae*, *25*(2), 399–406. Re-
 648 trieved from <https://angeo.copernicus.org/articles/25/399/2007/> doi:
 649 10.5194/angeo-25-399-2007
- 650 Medeiros, A. F., Taylor, M. J., Takahashi, H., Batista, P. P., & Gobbi, D. (2003).
 651 An investigation of gravity wave activity in the low-latitude upper mesosphere:
 652 Propagation direction and wind filtering. *J. Geophys. Res.*, *108*(D14), 4411,
 653 doi:10.1029/2002JD002593.
- 654 Nakamura, T., Aono, T., Tsuda, T., Admiranto, A. G., Achmad, E., & Suranto.
 655 (2003). Mesospheric gravity waves over a tropical convective region ob-
 656 served by OH airglow imaging in Indonesia. *Geophysical Research Letters*,
 657 *30*(17). Retrieved from <http://dx.doi.org/10.1029/2003GL017619> doi:
 658 10.1029/2003GL017619
- 659 Nakamura, T., Higashikawa, A., Tsuda, T., & Matsushita, Y. (1999). Seasonal
 660 variations of gravity wave structures in OH airglow with a CCD imager
 661 at Shigaraki. *Earth, Planets and Space*, *51*(7), 897–906. Retrieved from

- 662 <http://dx.doi.org/10.1186/BF03353248> doi: 10.1186/BF03353248
- 663 Picone, J. M., Hedin, A. E., & Drob, D. (2002). NRLMSISE-00 empirical model of
664 the atmosphere: Statistical comparisons and scientific issues. *J. Geophys. Res.*,
665 *107*(A12), SIA 15-1–SIA 15-16. doi: 10.1029/2002JA009430
- 666 Rasmussen, K. L., Chaplin, M. M., Zuluaga, M. D., & Houze, R. A. (2016). Con-
667 tribution of Extreme Convective Storms to Rainfall in South America. *Jour-
668 nal of Hydrometeorology*, *17*(1), 353-367. Retrieved from [https://doi.org/10
669 .1175/JHM-D-15-0067.1](https://doi.org/10.1175/JHM-D-15-0067.1) doi: 10.1175/JHM-D-15-0067.1
- 670 Rasmussen, K. L., Zuluaga, M. D., & Houze Jr., R. A. (2014). Severe convection and
671 lightning in subtropical South America. *Geophysical Research Letters*, *41*(20),
672 7359-7366. Retrieved from [https://agupubs.onlinelibrary.wiley.com/
673 doi/abs/10.1002/2014GL061767](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2014GL061767) doi: 10.1002/2014GL061767
- 674 Romatschke, U., & Houze, R. A. (2010). Extreme summer convection in South
675 America. *Journal of Climate*, *23*(14), 3761-3791. Retrieved from [http://dx
676 .doi.org/10.1175/2010JCLI3465.1](http://dx.doi.org/10.1175/2010JCLI3465.1) doi: 10.1175/2010JCLI3465.1
- 677 Snively, J. B., Nielsen, K., Hickey, M. P., Heale, C. J., Taylor, M. J., & Moffat-
678 Griffin, T. (2013). Numerical and statistical evidence for long-range ducted
679 gravity wave propagation over Halley, Antarctica. *Geophysical Research Let-
680 ters*, *40*(18), 4813–4817. Retrieved from [http://dx.doi.org/10.1002/
681 grl.50926](http://dx.doi.org/10.1002/grl.50926) doi: 10.1002/grl.50926
- 682 Snively, J. B., & Pasko, V. P. (2008). Excitation of ducted gravity waves in the
683 lower thermosphere by tropospheric sources. *Journal of Geophysical Research:
684 Space Physics*, *113*(A6). Retrieved from [http://dx.doi.org/10.1029/
685 2007JA012693](http://dx.doi.org/10.1029/2007JA012693) doi: 10.1029/2007JA012693
- 686 Swenson, G. R., & Mende, S. B. (1994). OH emission and gravity waves (including
687 a breaking wave) in all-sky imagery from Bear Lake, UT. *Geophysical Research
688 Letters*, *21*(20), 2239–2242. Retrieved from [http://dx.doi.org/10.1029/
689 94GL02112](http://dx.doi.org/10.1029/94GL02112) doi: 10.1029/94GL02112
- 690 Tang, J., Franke, S. J., Kamalabadi, F., & Swenson, G. R. (2005). Motion extraction
691 of atmospheric waves from spectroscopic imaging. *Geosci. Rem. Sens. Lett.*,
692 *2*(2), 229- 232.
- 693 Tang, J., Kamalabadi, F., Franke, S. J., Liu, A. Z., & Swenson, G. R. (2005). Es-
694 timation of gravity wave momentum flux using spectroscopic imaging. *IEEE*

- 695 *Transactions on Geoscience and Remote Sensing*, *43*(1), 103-109.
- 696 Tang, Y., Dou, X., Li, T., Nakamura, T., Xue, X., Huang, C., ... Avery, S. (2014).
697 Gravity wave characteristics in the mesopause region revealed from OH airglow
698 imager observations over Northern Colorado. *Journal of Geophysical Research:*
699 *Space Physics*, *119*(1), 630–645. doi: 10.1002/2013JA018955
- 700 Taylor, M. J. (1997). A review of advances in imaging techniques for measuring
701 short period gravity waves in the mesosphere and lower thermosphere. *Adv.*
702 *Space Res.*, *19*(4), 667-676.
- 703 Taylor, M. J., Gu, Y. Y., Tao, X., & Gardner, C. S. (1995). An investiga-
704 tion of intrinsic gravity wave signatures using coordinated lidar and night-
705 glow image measurements. *Geophys. Res. Lett.*, *22*(20), 2853-2856. doi:
706 10.1029/95GL02949
- 707 Taylor, M. J., Pautet, P.-D., Fritts, D. C., Kaifler, B., Smith, S. M., Zhao, Y., ...
708 Russell III, J. M. (2019). Large-Amplitude Mountain Waves in the Mesosphere
709 Observed on 21 June 2014 During DEEPWAVE: 1. Wave Development, Scales,
710 Momentum Fluxes, and Environmental Sensitivity. *Journal of Geophysical*
711 *Research: Atmospheres*, *124*(19), 10364-10384. Retrieved from [https://](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2019JD030932)
712 agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2019JD030932 doi:
713 <https://doi.org/10.1029/2019JD030932>
- 714 Taylor, M. J., Ryan, E. H., Tuan, T. F., & Edwards, R. (1993). Evidence of pref-
715 erential directions for gravity wave propagation due to wind filtering in the
716 middle atmosphere. *Journal of Geophysical Research: Space Physics*, *98*(A4),
717 6047–6057. doi: 10.1029/92JA02604
- 718 Vadas, S. L., Fritts, D. C., & Alexander, M. J. (2003). Mechanism for the genera-
719 tion of secondary waves in wave breaking regions. *Journal of the Atmospheric*
720 *Sciences*, *60*(1), 194–214. doi: 10.1175/1520-0469(2003)060<0194:MFTGOS>2.0
721 .CO;2
- 722 Vadas, S. L., Taylor, M. J., Pautet, P.-D., Stamus, P. A., Fritts, D. C., Liu, H.-L.,
723 ... Takahashi, H. (2009). Convection: the likely source of the medium-scale
724 gravity waves observed in the OH airglow layer near Brasilia, Brazil, dur-
725 ing the SpreadFEx campaign. *Annales Geophysicae*, *27*(1), 231–259. doi:
726 10.5194/angeo-27-231-2009
- 727 Vadas, S. L., Yue, J., She, C.-Y., Stamus, P. A., & Liu, A. Z. (2009). A model

728 study of the effects of winds on concentric rings of gravity waves from a con-
729 vective plume near Fort Collins on 11 May 2004. *J. Geophys. Res.*, *114*. doi:
730 10.1029/2008JD010753

731 Walterscheid, R. L., Hecht, J. H., Vincent, R. A., Reid, I. M., Woithe, J., & Hickey,
732 M. P. (1999). Analysis and interpretation of airglow and radar observations of
733 quasi-monochromatic gravity waves in the upper mesosphere and lower ther-
734 mosphere over Adelaide, Australia (35°S, 138°E). *J. Atmos. Sol.-Terr. Phys.*,
735 *61*, 461-478.