

# Opposite phase changes of precipitation annual cycle over land and ocean under global warming

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18    **Keypoints:**

- 19    1. Precipitation annual cycle shows a phase delay over land but a phase advance over ocean under  
20    global warming;
- 21    2. The land delay is due to enhanced effective atmospheric heat capacity and the ocean advance is  
22    linked to land-ocean rainfall shift in summer;
- 23    3. The enhanced and reduced surface temperature annual cycle over ocean versus land contribute  
24    to the rainfall shift via energetic constraint.

## 25    **Abstract**

26    The annual cycle of precipitation is a fundamental aspect of the global water cycle. Climate  
27    warming induces amplitude enhancement and phase delay in the zonal-mean tropical precipitation.  
28    Here, we report a land-ocean contrast in the phase response of precipitation annual cycle, with a  
29    delay over land and an advance over ocean as climate warms. Although two-thirds of the Earth's  
30    surface are covered by ocean, land dominates the zonal-mean phase delay, attributable to an  
31    increase in the effective atmospheric heat capacity. The phase advance over ocean is associated  
32    with a precipitation shift from land to ocean during the peak rainy season. This shift is well  
33    constrained by the energetic and related to a land-ocean contrast in the amplitude change of surface  
34    temperature annual cycle: seasonally different wind changes enhance this amplitude over ocean,  
35    while increased effective atmospheric heat capacity and surface cooling feedback reduce the  
36    amplitude over land.

## 37 **Plain language summary**

38 The seasonal monsoon rainfall provides the water resource for ~40% of the world population, but  
39 it has been a longstanding challenge to predict when monsoon rain will arrive, especially at the  
40 regional scale. By separating land from ocean, we revealed that the previous finding of seasonal  
41 delay of zonal-mean rainfall under warming mainly occurs over land, where most human activities  
42 take place, while ocean shows a phase advance. This contrasting phase behavior between land and  
43 ocean is shown to be rooted in the first-principle physical constraints related to energy and climate  
44 feedbacks, thus confidence is assigned to the differential phase response between the tropical land  
45 and ocean under climate warming.

## 1. Introduction

Over the annual cycle, the most prominent hydrological feature is the advance and retreat of tropical rainfall following the movement of monsoons and inter-tropical convergence zone (ITCZ). The annual cycle of tropical rainfall is ultimately driven by solar insolation, but its exact evolution is determined by a myriad of factors, including land-ocean distribution, land surface properties and ocean heat fluxes. Global warming is poised to substantially delay the cyclic progression of tropical rainfall, especially over the monsoonal regions (Biasutti and Sobel 2009; Seth et al. 2011 2013; Pascale et al. 2016; Song et al. 2018a) and enhance its amplitude (Chou et al. 2011; Chou and Lan 2012; Huang et al. 2013). The enhanced tropical precipitation annual cycle mainly occurs over ocean and is due to the increase of moisture under global warming as a “wet-get-wetter” response (Held and Soden 2006; Chou et al. 2009). The seasonal delay of tropical rainfall will induce seasonally-dependent subtropical high responses under warming (Song et al. 2018a, b).

In the past decade, the atmospheric energetic framework (e.g., Kang et al. 2008; Frierson et al. 2013; Schneider et al. 2014; Boos and Korty 2016; Biasutti et al. 2018) has been developed to relate the ITCZ shift to changes in the net energy input and/or the effective atmospheric heat capacity (Song et al. 2018). The causes of the delay are only beginning to be unveiled through the lens of this energetic perspective (Biasutti and Sobel 2009; Song et al. 2018a). Both the seasonal delay of sea surface temperature (SST) and the annual-mean SST warming can lead to a seasonal delay of tropical precipitation (Biasutti and Sobel 2009; Song et al. 2018; Dwyer et al. 2014). The former is attributed to the high-latitude sea ice melting, which boosts the effective heat capacity of the system by exposing the thick ocean mixed layer to the atmosphere (Dwyer et al. 2012; Donohoe and Battisti 2013). The latter can be explained by the increase of the effective atmospheric heat

capacity mainly due to the increase in atmospheric moisture with warming, resulting in a more sluggish response of atmospheric heat transport (AHT) and precipitation to the seasonal solar forcing (Song et al. 2018a; Cronin and Emanuel 2013).

Mechanisms that have been used to explain the phase change in the annual cycle of precipitation do not distinguish between the response over land and ocean. As land and ocean differ in many ways, differentiating their phase responses is a natural next step towards understanding regional water cycle in a warming climate. In this study, we find that under global warming, the phase of tropical precipitation is delayed over land but advanced over ocean. Interestingly, despite covering only one-third of the Earth's surface, it is land, not ocean, that dominates the phase changes of zonal-mean precipitation annual cycle. The opposite phase changes between land and ocean will have profound implications for both terrestrial/marine ecosystems and human activities.

## **2. Data and methods**

### **2.1 Model simulation data**

Monthly mean data of multiple variables from 37 Coupled Model Intercomparison Project Phase 5 (CMIP5; Taylor et al. 2012) models and daily precipitation data from 40 members of CESM1 Large Ensemble (LENS) Project (Kay et al. 2015) are used in this study. The statistics of the present-day climate for 1962-2005 are from the historical (HIST) simulations while those for the future climate of 2056-2099 are from the RCP8.5 (RCP85) scenario simulations.

To examine the role of wind-evaporation-SST (WES) feedback on the phase changes of precipitation annual cycle, we also analyze two sets of experiments: one with the same fully

coupled CESM1 as used for the CESM LENS Project and the other with a partially coupled configuration of CESM1 (Liu et al. 2018). Each set consists of a pair of 100-year long simulations with one under the condition of pre-industrial CO<sub>2</sub> (1xCO<sub>2</sub>) and the other with CO<sub>2</sub> quadrupling (4xCO<sub>2</sub>). In the partially coupled set, the wind speed from the pre-industrial simulation is prescribed through the bulk formula that is used to compute the evaporation over ocean in both the 1xCO<sub>2</sub> and 4xCO<sub>2</sub> runs, so that the part of the climate change response due to WES feedback is disabled. Comparing the response to 4xCO<sub>2</sub> forcing between the fully coupled and partially coupled sets allows the role of the WES feedback to be isolated. See Liu et al. (2018) for more details about the design of the experiments.

## 2.2 Phase estimated from the first Fourier harmonic

To estimate the phase of the annual cycle, Fourier transformation is performed to fit the related monthly time series to a Fourier harmonic (i.e., a sinusoidal function with time) with an angular frequency of  $2\pi/12\text{mon}$ . As the first Fourier harmonic explains more than 90% of variance in most models, we only use it to determine the phase of the annual cycle.

## 2.3 Atmospheric energetic framework

According to the atmospheric energy equation (Neelin and Held 1987), the divergence of AHT ( $\nabla \cdot AHT$ ) is equal to the difference between the net energy input to the atmosphere  $F_{net}$  and the column moist static energy (MSE) tendency  $\frac{\partial \langle h \rangle}{\partial t}$ :

$$\nabla \cdot AHT = F_{net} - \frac{\partial \langle h \rangle}{\partial t}, \quad (1)$$

where the angle bracket represents vertical integration between the surface and the top of the model,  $F_{net}$  includes sensible heat flux, latent heat flux, net longwave radiation and net shortwave radiation,  $h$  is the MSE. Following Boos and Korty (2016), we also isolate the divergent

component of the energy flux by solving the inverse Laplacian of energy flux potential  $\chi$ , which is defined as

$$\nabla^2 \chi = \nabla \cdot AHT \quad (2)$$

### 3. Results

Figure 1 shows the future change in the annual cycle of tropical precipitation over land and ocean in both hemispheres. Climatologically, when the peak solar forcing moves across the equator from one hemisphere to another, tropical precipitation increases rapidly in the warmed hemisphere, but the timing of precipitation peak differs between land and ocean, with the former leading the latter by about a month due to its lower heat capacity (solid lines in Fig. 1a-d). Under global warming, there is an intriguing contrast in the future changes between land and ocean (bars in Fig. 1a-d): over land, precipitation anomalies peak much later than the climatological peak, while the opposite is true over ocean, indicating a phase delay in the precipitation annual cycle over land but a phase advance over ocean. The larger amplitude changes during the transition seasons over land than ocean accentuate the dominance of the land response in the zonal-mean response.

The phase changes over land versus ocean (0-40°N/S) are further quantified. Most CMIP5 models (Fig. 1e) and all LENS members (Fig. 1f) exhibit a seasonal delay over land, with median values of 2 to 4.5 days, and a seasonal advance over ocean, with a similar range of median values. For the Northern Hemispheric (NH) land and ocean, 34 and 33 out of 37 CMIP5 models show the phase delay and advance, respectively. Similarly for the Southern Hemispheric (SH) land and ocean, 36 and 29 out of 37 CMIP5 models show the phase delay and advance, respectively.



To decipher the contrasting precipitation phase changes between land and ocean, we conduct Empirical Orthogonal Function (EOF) analysis on the annual cycle of precipitation and identify two dominant modes (Fig. 2; Wang and Ding 2008): one characterized by an inter-hemispheric contrast (EOF1; accounting for 67.9% of variance) with a peak during local summer and the other featuring a distinct land-sea contrast (EOF2; accounting for 18.6% of variance) with a peak during local spring. Precipitation center tends to be well collocated with the local maxima of column MSE (contours in Figs. 2a-b). The phase change in both PCs is in quadrature to the phase of the corresponding climatological PCs, indicating a phase delay (Fig. 2c-d). The phase delay in the EOF1 has been attributed to both the phase delay in the annual cycle of SST and the annual-mean SST warming as explained above. The similar phase delays in PC1 and PC2 are mainly the result of EOF analysis that requires the EOF modes to be orthogonal to each other. The seasonal delay in both PC1 and PC2 is quite robust among the CMIP5 models and LENS simulations (Fig. 2e), with 36 out of 37 CMIP5 models and 39 out of 40 LENS members showing a delay. Another robust aspect in the response of the two EOF modes under warming is the enhancement of the amplitude in all CMIP5 models and LENS members, with a median increase of ~10% (Fig. 2f), consistent with the enhanced annual cycle of the global monsoon rainfall (Kitoh et al. 2013; Lee and Wang 2014).

Were the spatial pattern of the leading two EOF modes unchanged, the delay in their phases would imply a phase delay in the precipitation annual cycle over both land and ocean. But it turns out not to be the case (Fig. 1). As the two leading EOF modes of precipitation annual cycle explain more than 85% of the total seasonal variance, we construct the precipitation annual cycle over land and ocean ( $P_L$  and  $P_O$ , respectively) based on the two modes as follows:

$$P_L = A_{1,L}PC1 + A_{2,L}PC2 \quad (3)$$

$$P_O = A_{1,O}PC1 + A_{2,O}PC2 \quad (4)$$

Here,  $PC1$  and  $PC2$  are the corresponding normalized principle components,  $A_{1,L}$  and  $A_{2,L}$  are the amplitudes of EOF1 and EOF2 averaged over land, respectively;  $A_{1,O}$  and  $A_{2,O}$  are the same but over ocean.

Given the orthogonality of PCs, we can express the  $PC1$  and  $PC2$  as two sinusoidal functions with  $PC2$  leading  $PC1$  by  $\pi/2$ . Then, the annual cycle of  $P_L$  and  $P_O$  can be written as:

$$P_L = \sqrt{A_{1,L}^2 + A_{2,L}^2} \sin(t - \phi_1 - \phi_L) \quad (5)$$

$$P_O = \sqrt{A_{1,O}^2 + A_{2,O}^2} \sin(t - \phi_1 - \phi_O), \quad (6)$$

Here,  $t$  is time,  $\phi_1$  denotes the phase of  $PC1$ . It is found that the reconstructed precipitation annual cycle and the actual one is almost identical. This way, the phases of  $P_L$  and  $P_O$  are  $\phi_1 + \phi_L$  and  $\phi_1 + \phi_O$ , respectively.  $\phi_L = -\arcsin\left(\frac{1}{\sqrt{\left(\frac{A_{1,L}}{A_{2,L}}\right)^2 + 1}}\right)$  and  $\phi_O = \arcsin\left(\frac{1}{\sqrt{\left(\frac{A_{1,O}}{A_{2,O}}\right)^2 + 1}}\right)$  represent extra

phase associated with the relative contribution of EOF1 versus EOF2 to the land and ocean annual cycle, respectively. Under warming, if the amplitude of EOF1 increases more than EOF2 over land,  $\frac{A_{1,L}}{A_{2,L}}$  becomes larger and the resultant change in  $\phi_L$  would imply a delay. The opposite is true over ocean (note the opposite signs between  $\phi_L$  and  $\phi_O$ ). The actual land-ocean phase change differences (Fig. 1e-f) can be well explained by their amplitude-based estimates (i.e.,  $\phi_L - \phi_O$ ), with correlation of 0.97 and 0.89 in NH and SH, respectively (Supplementary Fig. 1), lending support to the decomposition framework here. Hence, the land-sea difference in phase changes of

precipitation annual cycles can also be casted in terms of the spatial pattern changes in EOF1 and EOF2 (i.e.,  $\frac{A_{1,L}}{A_{2,L}}$  and  $\frac{A_{1,O}}{A_{2,O}}$ ).

Figure 3a shows a weak reduction in  $\frac{A_{1,L}}{A_{2,L}}$ , but a marked increase in  $\frac{A_{1,O}}{A_{2,O}}$  in both hemispheres ( $\frac{A_{1,O}}{A_{2,O}}$  increases by 38% and 29% for NH and SH, respectively) under warming, with the former opposing slightly the delay in  $\phi_1$  by an amount  $\phi_L$ , while the latter overwhelming the positive change in  $\phi_1$ , resulting in a net phase advance in  $P_O$  (i.e.,  $\phi_O + \phi_1 < 0$ ) shown in Fig. 1. Further examination indicates that the increase in  $\frac{A_{1,O}}{A_{2,O}}$  is dominated by the increase in  $A_{1,O}$  (Fig. 3b), with  $A_{2,O}$  showing negligible changes (not shown). The change of EOF1 precipitation pattern (Fig. 3c) features large increases over broad tropical oceans, while the change is only marginal over tropical land, indicating a shift of precipitation from land to ocean during local summer when EOF1 peaks. This shift is robust in terms of the land-minus-ocean precipitation change among CMIP5 models (Supplementary Fig. 2a).

What tilts the land-ocean precipitation balance towards ocean in EOF1? As the “wet-get-wetter” response doesn’t hold well over land due to the limited moisture supply (Fasullo 2012; Chadwick et al. 2013; Byrne and O’Gorman 2015; Donat et al. 2016), it may not be surprising that the land precipitation change is subdued relative to the ocean response. Here we explore other possible mechanisms via the energetic perspective. Viewing from the energetic constraint on precipitation, the land-ocean precipitation shift in EOF1 should be associated with an AHT from ocean to land, and indeed this is evidenced by the direct calculation of the change in AHT divergence contrast between land and ocean (Fig. 3c; Supplementary Fig. 2). This divergent flow of MSE appears to be from the warmer ocean to the colder land (Fig. 3c&d). According to the convective quasi-equilibrium argument suitable for the tropical convective regions (Emanuel

1995; Shekar and Boos 2016), the column MSE is approximately in equilibrium with the sub-cloud entropy, which is in turn regulated by the surface temperature (TS) when the surface relative humidity is high. With the diffusive nature of energy transport (Boos and Korty 2016), AHT should point from higher MSE to lower MSE regions and hence from the relatively warmer TS to the colder TS regions. The correlation between the land-sea contrast of AHT divergence and TS changes among CMIP5 models is 0.81 and 0.63 in NH and SH, respectively. Even over land, the corresponding correlation is 0.70 and 0.54, still statistically significant at 1% level. Previous studies (e.g., Hurley and Boos 2013; Roderick et al. 2014) suggested that over land, TS is mainly driven by precipitation when evaporation is limited. If so, the TS change would be negatively correlated with the change in precipitation and AHT divergence. Here, the significantly positive correlation between TS and AHT divergence changes over land implies the decreased TS over land does play a role in the suppressed precipitation in summer. As such, the land-sea precipitation shift during summer is linked to the opposite amplitude changes in the TS annual cycle over land and ocean, which is the focus of our investigation next.

Over ocean, a robust climate change response is the strengthening of surface wind in the winter subtropics and the weakening in the summer subtropics (Sobel and Camargo; contours in Fig. 3d). As wind speed is identified to be important in the tropical SST warming pattern formation (Xie et al. 2010), the contrasting wind response induces relative cooling in the winter subtropics and warming in the summer subtropics via WES feedback (Xie and Philander 1994; Lu and Zhao 2012), leading to the enhanced SST annual cycle (Fig. 3d). Thus, more rainfall occurs over ocean during summer, resulting in a phase advance over ocean. To further support our hypothesis, the CESM1 experiments with active WES reproduce the opposite phase changes of precipitation annual cycle over land and ocean (red crosses in Fig. 1f). When the WES feedback is disabled in

the 4xCO<sub>2</sub> experiments, the phase advance over ocean is completely nullified in NH and reduced by more than half in SH, whereas the land phase delay mostly persists (blue crosses in Fig. 1f).

In contrast to the increased amplitude of TS over ocean, the amplitude of TS is reduced over land, most notably in the monsoon regions (Fig. 3d). The governing mechanism for the annual cycle of temperature over land can be captured by an atmosphere-land surface interaction model forced by a periodic solar insolation:

$$C_A \frac{dT_A}{dt} = \lambda(T_S - T_A) - B T_A \quad (7)$$

$$C_S \frac{dT_S}{dt} = -\lambda(T_S - T_A) + R, \quad (8)$$

where  $T_A$  is the mass-weighted temperature of the column atmosphere, which is dominated by the near surface air temperature,  $T_S$  is the land TS,  $C_A$  and  $C_S$  are the effective heat capacities of the atmosphere and surface, respectively. The total surface-atmosphere energy exchange including longwave radiation, sensible and latent heat fluxes is parameterized to be linearly proportional to  $T_S - T_A$  with an exchange coefficient  $\lambda$  (as similarly treated in Barsugli and Battisti (1998) and Zhou and Xie (2018)).  $B$  is the bulk feedback parameter accounting for all the radiative processes at the top of atmosphere, including water vapor, cloud, lapse rate and Planck feedbacks. Since the adjustment of  $T_A$  through the top of the atmosphere is slower than that at the surface,  $B$  is generally smaller than  $\lambda$  (Barsugli and Battisti 1998). The system is forced at the surface by a periodic solar radiation  $R$  with a frequency of yr<sup>-1</sup>, hence shortwave absorption by the atmosphere is neglected.

Over tropical land,  $C_S$  is estimated to be one order smaller than  $C_A$ . Thus, adding Eqs. (7) and (8) together and ignoring the small term associated with  $C_S$  reduce the system to a single variable system:

$$C_A \frac{dT_A}{dt} = -BT_A + R, \quad (9)$$

from which the amplitude relation between  $T_A$  and  $R$  can be derived:

$$|T_A| = \frac{|R|}{\sqrt{B^2 + \omega^2 C_A^2}}, \quad (10)$$

where  $\omega = 2\pi yr^{-1}$ . Under the condition  $\omega C_A \ll \lambda$ , which holds well for the modern climate condition, the amplitude relation between  $T_A$  and  $T_S$  can be written as:

$$|T_S| \approx (1 + \frac{B}{\lambda})|T_A|, \quad (11)$$

Thus, for a given solar insolation, the amplitude of atmospheric temperature decreases with the increase in the effective atmospheric heat capacity and bulk feedback parameter  $B$ . The same may also be said of the amplitude of  $T_S$  considering the linear relationship between  $T_S$  and  $T_A$  from Eq. (11).

Following Schwartz (2007), the effective atmospheric heat capacity is defined as  $C_A = \frac{\partial \langle h \rangle / \partial t}{\partial \langle T \rangle / \partial t}$ , with  $\langle h \rangle$  and  $\langle T \rangle$  representing the vertically-integrated MSE and temperature. This definition of  $C_A$  allows to estimate it by regressing the monthly tendency of multi-model ensemble mean  $\langle h \rangle$  against that of  $\langle T \rangle$ . Based on this definition,  $C_A$  can be scaled as  $C_A \approx c_p + \frac{L_v \langle q \rangle}{\langle T \rangle}$ , where  $\langle q \rangle$  is the vertically integrated water vapor. As  $\langle q \rangle$  increases nonlinearly with temperature following the Clausius-Clapeyron relation,  $C_A$  increases as the mean climate warms. Song et al. (2018a) has already shown that the annual cycle of  $\partial \langle h \rangle / \partial t$  is robustly enhanced mainly due to the contribution from  $\partial \langle L_v q \rangle / \partial t$ , which scales linearly with global warming. Here, we provide further evidences for the increase of effective atmospheric heat capacity (Supplementary Fig. 3). In the current climate,  $C_A$  is 2162.10 and 2704.01 J kg<sup>-1</sup> K<sup>-1</sup> for NH and

SH, respectively. Under global warming,  $C_A$  is increased to 2393.10 and 3271.39 J kg<sup>-1</sup> K<sup>-1</sup>, with an increase rate of 10.7% and 21.0% for NH and SH, respectively. The enhancement of  $C_A$  is ubiquitous across all CMIP5 models examined (Supplementary Fig. 4). As no robust change in  $B$  is found over tropical land under warming in CMIP5 models, the robust increase of the effective atmospheric heat capacity acts to reduce the amplitude of the  $T_A$  and  $T_S$  annual cycle over land, consistent with Cronin and Emanuel (2013) based on radiative-convective equilibrium simulations.

Since a good proportionality holds between  $T_A$  and  $T_S$  with small phase differences for the system above, we can further simplify Eq. (8) into a surface energy balance model for  $T_S$ :

$$C_S \frac{dT_S}{dt} = -\beta T_S + R, \quad (12)$$

where  $\beta$  is a parameter to measure the surface cooling feedback, encompassing processes of net longwave radiative cooling and turbulent heat fluxes. A similar surface energy model has been used to understand the temperature annual cycle in an extraterrestrial planet (Mitchell et al. 2014), where only longwave radiation is considered in the damping term. The amplitude of  $T_S$  governed by Eq. (12) is determined by

$$|T_S| = \frac{|R|}{\sqrt{\beta^2 + \omega^2 C_S^2}}, \quad (13)$$

and is expected to decrease with an increase of the feedback parameter  $\beta$ . We estimate  $\beta$  by regressing the total surface energy flux excluding the shortwave fluxes on the annual cycle of TS for both current and future climates. The resultant  $\beta$  is generally positive over land in the current climate (Supplementary Fig. 5a), vindicating its physical meaning as a damping parameter. In the tropical monsoonal regions,  $\beta$  indeed increases across the majority of CMIP5 models under the

RCP85 scenario and geographically, it coincides with regions of decreased TS annual cycle (Fig. 3d and Supplementary Fig. 5b). Averaged over the land monsoonal regions, the change of  $\beta$  is significantly (at 1% level) correlated with that of summer-vs-winter TS difference across the CMIP5 models, with correlations reaching -0.68 in NH and -0.46 in SH (or -0.81 with the outlier models excluded) (Supplementary Fig. 6). Taken together, both the increase of  $C_A$  and  $\beta$  work in tandem to dampen the TS amplitude over land.

#### 4. Conclusions

This study discovers intriguing opposite phase changes in the precipitation annual cycle between land and ocean under climate warming: a delay over land, but an advance over ocean. Although land only covers one-third of the Earth's surface, it dominates the phase delay of zonal-mean precipitation annual cycle previously found (Biasutti and Sobel 2009; Song et al. 2018a). The delay over land is mainly determined by the increased effective atmospheric heat capacity under warming, the same mechanism responsible for the zonal-mean precipitation delay (Song et al. 2018a). The phase advance over ocean is related to a shift of precipitation center towards ocean in the local summer, which overwhelms the phase delay resulting from the increased effective atmospheric heat capacity. The shift of precipitation from land to ocean during the peak rainy season corresponds well to the pattern of AHT divergence change, epitomizing the energetic constraint on the tropical precipitation change even at the regional scale. The AHT divergence change under climate warming is closely related to the opposite changes in the amplitude of TS annual cycle between land and ocean. The strengthened surface wind in winter cools SST and the weakened one in summer warms SST, enhancing the amplitude of SST annual cycle. Over land, on the other hand, both the increased effective atmospheric heat capacity and surface cooling



305 feedback over monsoonal regions act to dampen the amplitude of TS annual cycle. The  
306 combination of enhanced TS amplitude over ocean and reduced TS amplitude over land rebalances  
307 the atmospheric column energy in favor of ocean precipitation. As a result, the center of  
308 precipitation shifts towards ocean, which is further manifested as a differential phase change  
309 between land and ocean.

310         Considering the earlier transition from dry to wet seasons over land than ocean in the  
311 climatological annual cycle, the phase delay over land and phase advance over ocean under  
312 warming will make the two more in sync in a warmer future. The delayed onset of rainfall as well  
313 as the dampened annual cycle of temperature over land would have profound economic and  
314 societal impacts, especially for regions like Asia, Africa and Latin America, where the local rainfed  
315 agriculture and economy are under the sway of monsoonal rain.

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322 [data.org/sim/gcm\\_monthly/AR5/Reference-Archive.html](http://www.ipcc-data.org/sim/gcm_monthly/AR5/Reference-Archive.html)) and NCAR CESM1 group, which  
323 makes the large-ensemble (LENS) experiments available  
324 (<http://www.cesm.ucar.edu/projects/community-projects/LENS/>).

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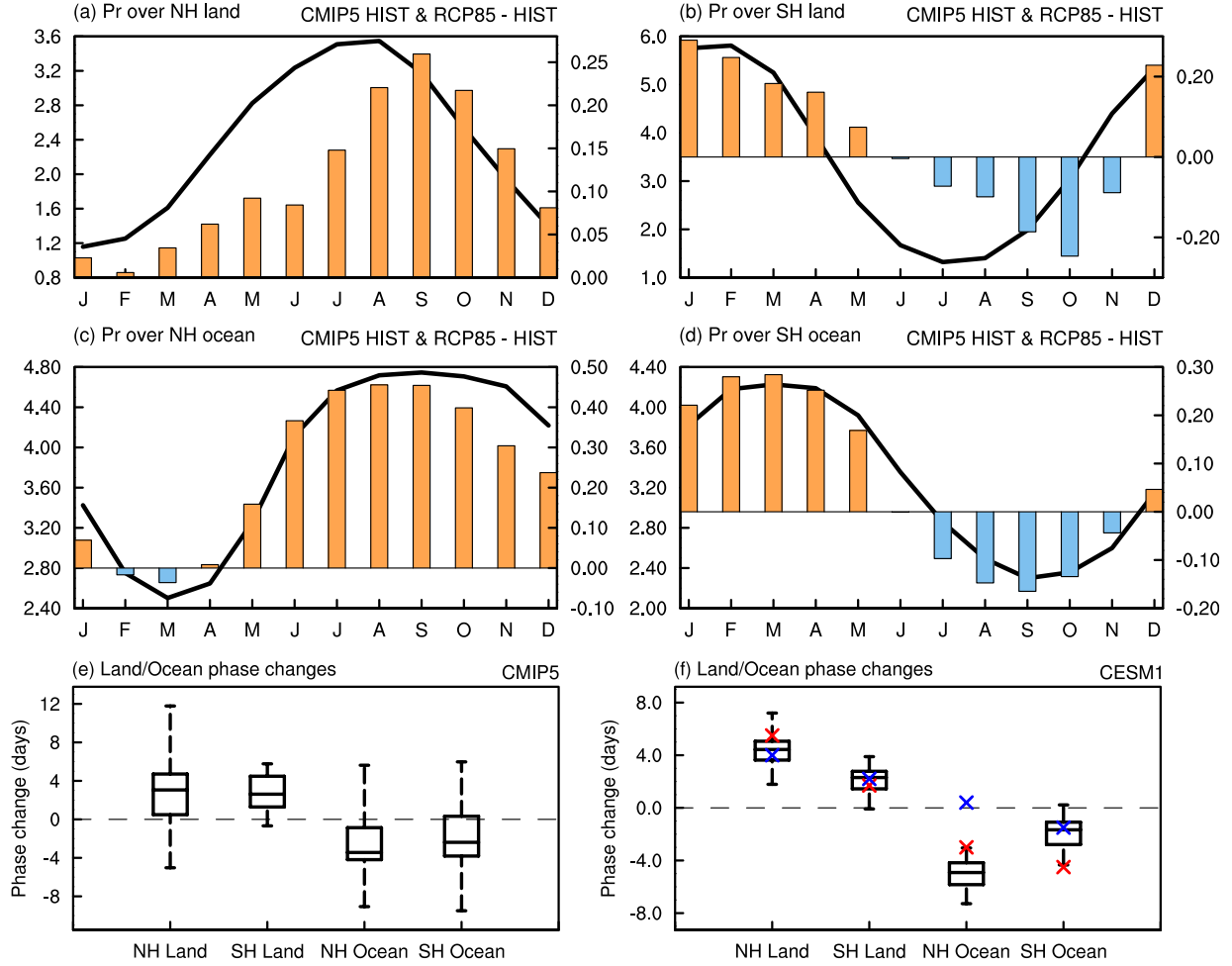
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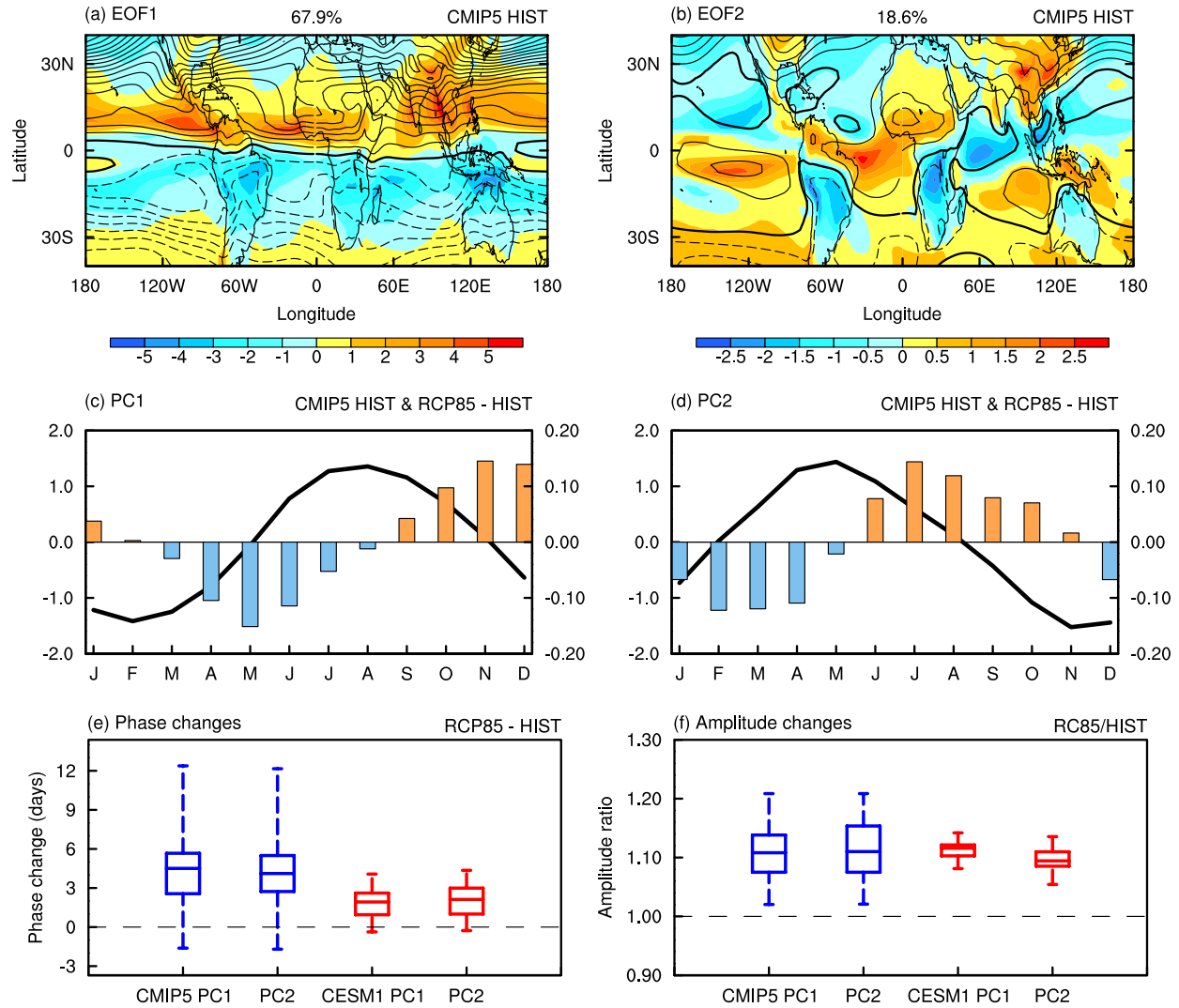
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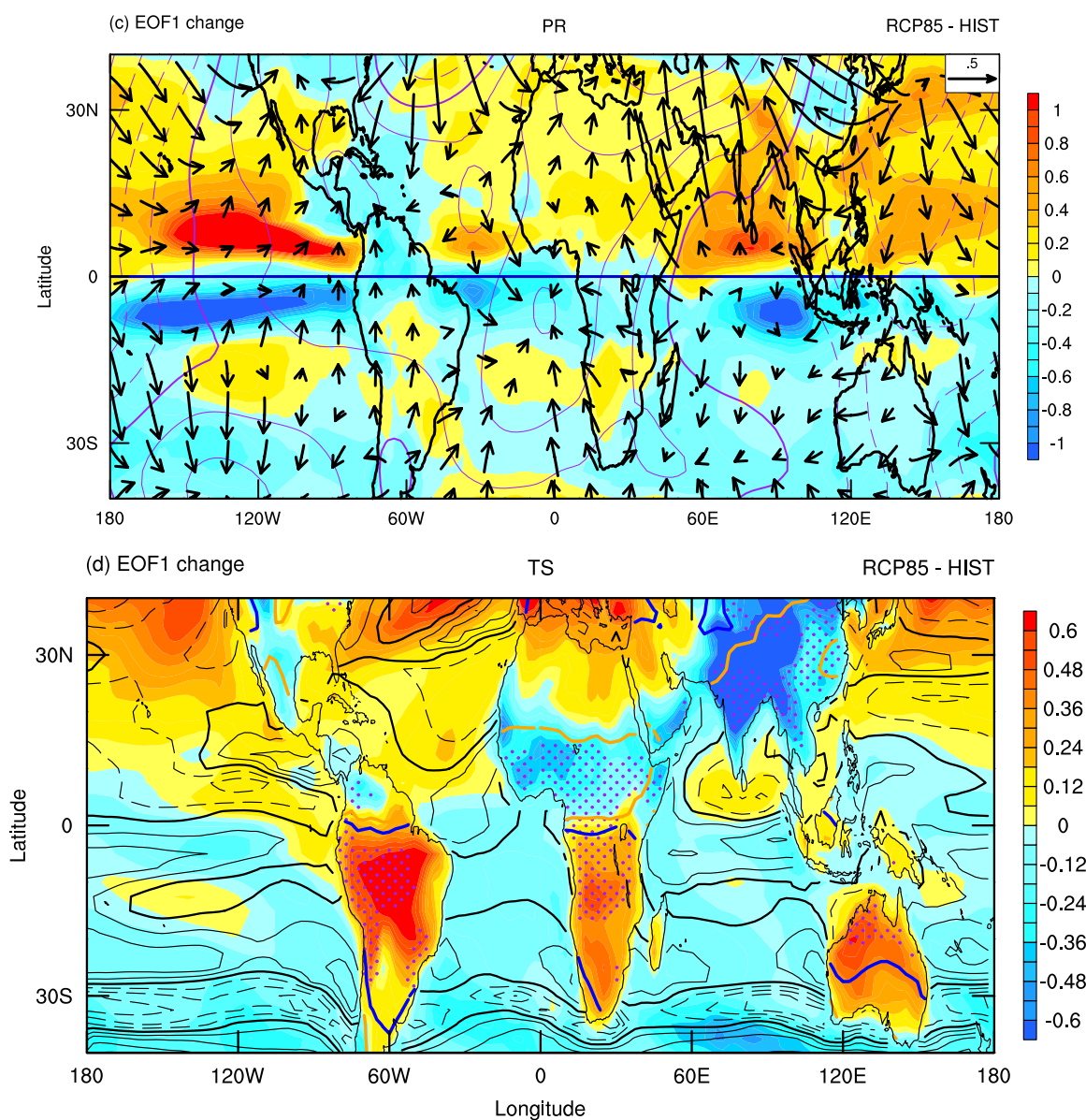
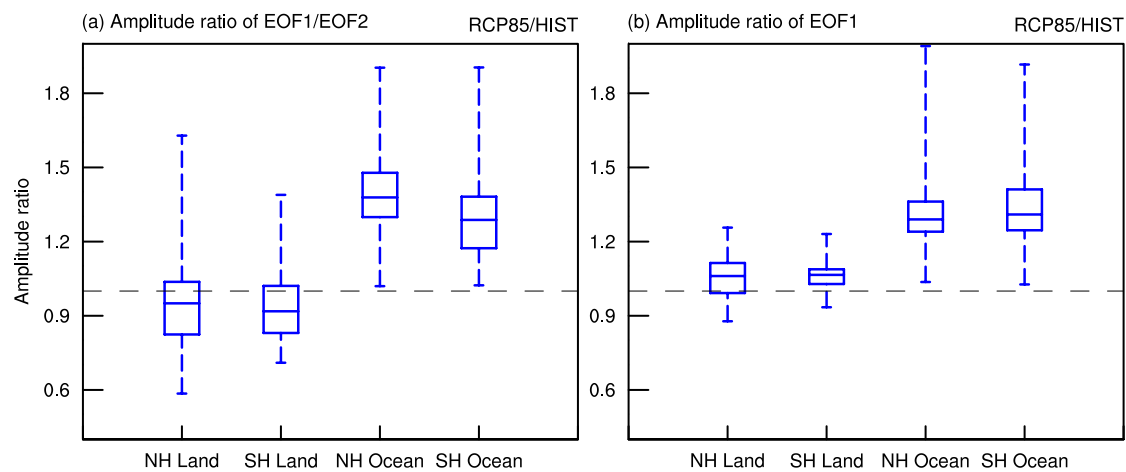
**Figure 1** The annual cycle of precipitation (unit: mm day<sup>-1</sup>) in HIST runs (black line) and its future changes between RCP85 and HIST runs over (a) North Hemispheric (NH) land (0°-40°N), (b) South Hemispheric (SH) land (0°-40°S), (c) NH ocean (0°-40°N) and (d) SH land (0°-40°S). Box-plots for the future phase changes in the precipitation annual cycle over land and ocean from (e) CMIP5 monthly data and (f) CESM1 LENS daily data. The lines in each box represent the 25th percentile, median, and the 75th percentile, and the whiskers represent the minimum and maximum of the multi-model (or LENS) ensemble. In (f), the red and blue crosses indicate the phase changes between the 4xCO<sub>2</sub> and 1xCO<sub>2</sub> experiments of CESM1 with and without WES feedback, respectively.





**Figure 2** The spatial patterns of (a) EOF1 and (b) EOF2 in HIST runs and the corresponding normalized principle components (c) PC1 and (d) PC2 in HIST (black line) and their corresponding future changes between RCP85 and HIST. In a (b), overlaid as contours is the vertically-integrated MSE (unit:  $10^7 \text{ J kg}^{-1}$ ) regressed upon the normalized PC1 (PC2) in HIST runs. The bold line indicates zero and solid (dashed) lines represent positive (negative) values with an interval of  $0.5 \times 10^7 \text{ J kg}^{-1}$ . Box-plots for (e) the phase changes and (f) the amplitude ratio between RCP85 and HIST for the two EOF modes from CMIP5 (blue) and CESM1 LENS (red).

435 The lines in each box represent the 25th percentile, median, and the 75th percentile, and the  
436 whiskers represent the minimum and maximum of the multi-model (or LENS) ensemble.



438 **Figure 3** Box-plots comparing the ratio of the amplitude of (a) EOF1 over land  $A_{1,L}$  divided by  
 439 EOF2 over land  $A_{2,L}$  and EOF1 over ocean  $A_{1,O}$  divided by EOF2 over ocean  $A_{2,O}$  and (b) EOF1  
 440 over land  $A_{1,L}$  and EOF1 over ocean  $A_{1,O}$  between the RCP85 and HIST runs (i.e., RCP85 divided  
 441 by HIST) among CMIP5 models. Future change of (c) precipitation (shaded; unit: mm day<sup>-1</sup>),  
 442 energy flux potential (contour; unit: W) and divergent energy flow (vectors; unit: 10<sup>7</sup> W m<sup>-1</sup>) and  
 443 (d) TS (unit: K) regressed on the PC1 of precipitation annual cycle between the RCP85 and HIST  
 444 runs. In c, the bold line is zero and the solid (dashed) contours denote positive (negative) values  
 445 with the contour interval of 0.4x10<sup>13</sup> W. The line patterns over ocean in d have the same meaning  
 446 as c except that it is the change in the 1000 hPa wind speed congruent with EOF1 and the interval  
 447 is 0.2 m s<sup>-1</sup>. The orange and blue contours over land in d indicate the 0.5 and -0.5 mm day<sup>-1</sup> contours  
 448 of precipitation EOF1 shown in Fig. 2a, respectively, demarcating the boundaries of the land  
 449 monsoonal regions. The purple dots over land indicate regions of increased surface cooling  
 450 feedback parameter  $\beta$  under global warming. The horizontal lines in a and b represent the  
 451 minimum, 25th percentile, median, 75th percentile, and the whiskers are the maximum and  
 452 minimum values of the multi-model ensemble.