

1 **Short-scale variations in high-resolution**
2 **crystal-preferred orientation data in an alpine ice core**
3 **– do we need a new statistical approach?**

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12 **Key Points:**

- 13 • Crystal-preferred orientation data were collected in high-resolution in intervals of
14 up to 1 m length in a cold Alpine ice core.
- 15 • A strong anisotropy and high spatial variability on the cm-scale is found and chal-
16 lenges established practices in fabric sampling.
- 17 • Strain localisation is considered as main driver of short-scale variations in crystal-
18 orientation fabric in ice.

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Abstract

We analysed crystal-preferred orientation of c-axis and microstructure data from the Alpine ice core KCC at an unprecedented resolution and coverage of any Alpine ice core. We find that an anisotropic single-maximum fabric develops as early as 25 m depth in firn under vertical compression and strengthens under simple shear conditions towards the bedrock at 72 m depth. The analysis of continuously measured intervals with subsequent thin section samples from several depths of the ice core reveals a high spatial variability in the crystal orientation and crystal size on the 10 cm-scale as well as within a few centimeters. We quantify the variability and investigate the possible causes and links to other microstructural properties. Our findings support the hypothesis that the observed variability is a consequence of strain localisation on small spatial scales with influence on fabric and microstructure. From a methodological perspective, the results of this study lead us to challenge whether single thin sections from ice cores provide representative parameters for their depth to be used to infer the fabric development in a glacier on the large scale. Previously proposed uncertainty estimates for fabric and grain size parameters do not capture the observed variability. This might therefore demand a new scale-sensitive statistical approach.

1 Introduction

One of the foremost goals of contemporary glaciological research is to advance our understanding of the internal deformation of glaciers and ice sheets to better constrain the dynamic components of modelled projections of their development under changing climate conditions. To this end, the physical properties of polycrystalline glacier ice have been studied on a number of deep ice cores from the polar regions for more than 40 years, as reviewed in Faria, Weikusat, and Azuma (2014a). The analysis of crystal-preferred orientation (CPO, fabric) and microstructure (μS) on ice core samples provides evidence to draw conclusions on dominant deformation and recrystallisation (RX) regimes at different depths in the ice column (Weikusat, Jansen, et al., 2017), and may also be of importance for the interpretation of climate records (Faria, Freitag, & Kipfstuhl, 2010; Kennedy, Pettit, & DiPrinzio, 2013). To clarify the terminology in this study, *microstructure* includes the shape and size of grains, grain boundaries (GBs), bubbles and other spatial features in ice, while *fabric* is only used for crystallographic orientation of crystals.

The inherent plastic anisotropy of individual ice crystals (Duval, Ashby, & Anderman, 1983) leads to their alignment in preferred orientations under the local stress conditions along the flow path (Azuma & Higashi, 1985). Dynamic recrystallisation occurs as a consequence of the accumulated strain energy (Llorens et al., 2017; Schulson & Duval, 2009), and the presence of pore space and impurities contribute to the complex development of fabric and microstructure (Azuma, Miyakoshi, Yokoyama, & Takata, 2012; Durand, Weiss, et al., 2006; Eichler et al., 2017; Steinbach et al., 2016). As a result of these inter-dependent micro-scale processes, the viscous behaviour of the ice bulk is changing, which, in turn, can lead to an enhancing feedback on the deformation (Azuma & Goto-Azuma, 1996). To establish a thorough understanding of how dynamic processes on the micro-scale, recorded as a temporal snapshot in ice-core data, are connected with the observations on the macro-scale, remains a major challenge in the study of ice-physical properties (Eisen, Hamann, Kipfstuhl, Steinhage, & Wilhelms, 2007).

While it has repeatedly been demonstrated that it is essential to include the crystal anisotropy into ice flow models (Graham, Morlighem, Warner, & Treverrow, 2018; Martin, Gudmundsson, Pritchard, & Gagliardini, 2009; Pettit, Thorsteinsson, Jacobson, & Waddington, 2007; Seddik, Greve, Placidi, Hamann, & Gagliardini, 2008), it is difficult to obtain the necessary amount of data from an ice core which is representative for the CPO development in the vertical dimension, while the sampling in the horizontal dimension is, in any case, limited to the cross-section of the core. In fact, it is currently not constrained how much ice-core data can be considered as representative. Due to the time-consuming nature of fabric studies, sampling for this purpose is typically in

72 the order of one sample per tens of metres of ice core. Sampling is thus at a very low
 73 resolution with much less than a few per cent of the core analysed. (appendix A, Table A.1).
 74 Despite this, the results are assumed to reflect the general fabric evolution in the ice col-
 75 umn. However, some studies report the observation of small-scale features like fine-grained
 76 layers in the microstructure and sudden spatial changes between distinct fabric patterns
 77 on a resolution of about 15 m (DiPrinzio et al., 2005; Fitzpatrick et al., 2014; Jansen et
 78 al., 2016; Montagnat et al., 2014). It stands to question if these structures are common
 79 to most ice cores, potentially (1) distorting the perceived representation from low-resolution
 80 ice-core fabric studies, and (2) disregarding information on deformation processes on smaller
 81 scales.

82 Only a few studies are focused on the fabric of non-polar (e.g. alpine or sub-polar)
 83 glacier ice, with most of them being older than 30 years and/or from temperate glaciers
 84 (Table 1; see also the review paper from Hudleston, 2015). While the measurement tech-
 85 niques were certainly less sophisticated in the past, and these thus provided statistically
 86 less representative data sets, it was reported that samples from nearby locations (within
 87 a few meters) showed distinctly different fabric (Kamb, 1959; Rigsby, 1951) and bands
 88 with small crystals and single-maximum fabrics were observed in an ice core (Vallon, Pe-
 89 tit, & Fabre, 1976). Non-polar glaciers offer the opportunity to study glaciological pro-
 90 cesses that can be assumed to follow the same physical laws as in polar regions, on a smaller
 91 spatial scale and a more constrained temporal scale, providing upper limits for the du-
 92 ration of microstructural processes. The ice temperature is generally higher compared
 93 to the non-basal parts of polar ice sheets and often similar to the temperatures encoun-
 94 tered at the base of polar ice sheets, approaching the pressure-melting point, where tempera-
 95 ture-driven microstructural processes may be enhanced (Faria, Weikusat, & Azuma, 2014b).
 96 They can thus serve as in-situ or natural laboratories with deformation conditions that
 97 are not feasible in laboratory experiments and several orders closer to those in polar ice
 98 sheets. However, the majority of non-polar ice cores from the last 30 years were anal-
 99 ysed with a climatological focus (e.g. Bohleber, Wagenbach, Schner, & Bhm, 2013; Pre-
 100 unkert, Wagenbach, Legrand, & Vincent, 2000; Schwikowski, Brütsch, Gäggeler, & Schot-
 101 terer, 1999; Thompson et al., 2006). Only recently, some low-resolution fabric data from
 102 an ice core drilled in 1992 on the Tibetan Plateau (Thompson et al., 1995) were reported
 103 and compared to polar fabric data (Y. Li, Kipfstuhl, & Huang, 2017).

104 The cold high-Alpine glacier Colle Gnifetti (CG) has long been the subject of ex-
 105 tensive glaciological research, with a focus on its archive of midlatitude climate (Wagen-
 106 bach, Bohleber, & Preunkert, 2012). In the present study we provide the first compre-
 107 hensive fabric analysis on this Alpine glacier, obtained from a recent ice core. For the
 108 first time, we systematically analyse the variability of fabric on a scale of centimeters to
 109 decimeters within continuously analysed depth intervals.

110 While the crystallographic evolution is the main focus of the study, other physi-
 111 cal properties and microstructural parameters (grain size, density, bubbles) are consid-
 112 ered for the benefit of discussing the context of physical processes that could influence
 113 the fabric development at Colle Gnifetti. Following the findings of our study, we discuss
 114 the significance of defining a representative scale length for fabric studies.

115 2 Study Site

116 The high-Alpine glacier Colle Gnifetti, Monte Rosa Massif, on the border between
 117 Switzerland and Italy, is characterized by a high spatio-temporal variability in net snow
 118 accumulation which is primarily a result of insolation-driven snow consolidation coun-
 119 teracting the effect of wind erosion (Wagenbach, Mnnich, Schotterer, & Oeschger, 1988).
 120 On average, net snow accumulation is mainly made up by precipitation during the warm
 121 seasons, with a minor but highly variable contribution by winter snow. The average net
 122 accumulation rate at KCC was determined around 22 cm water equivalent per year (Bohle-
 123 ber et al., 2018).

Table 1. Early and more recent crystallographic studies on non-polar or subpolar glaciers. We would like to mention current projects involving crystal-orientation measurements on temperate and polythermal ice on Rhone Glacier, Switzerland (Hellmann et al., 2018) and Storglaciären, Sweden (Monz & Hudleston, 2018), and Jarvis Glacier, Alaska (Gerbi et al., 2018). *Section* refers to a thin section sample.

| Authors | Region | Type & Temperature | Fabric resolution; sample description |
|---|---|---|---|
| Rigsby (1951) | Emmons glacier, Mt. Rainier, Washington, USA | Mountain and valley glacier, temperate | 11 surface locations, stagnant and fast flowing ice; ~40 sections with 1725 crystals |
| Meier, Rigsby, and Sharp (1954) | Saskatchewan Glacier, Alberta, Canada | Valley glacier, temperate | 8 surface locations |
| Rigsby (1960) | Emmons, Saskatchewan and Malaspina Glacier, Alaska, USA; | 3 temperate, 2 polar sites | surface locations, 8000 crystals |
| | Moltke Glacier and Nunatarsuak Ice Ramp, Thule, Greenland | | |
| Kamb (1959), Allen, Kamb, Meier, and Sharp (1960) | Lower Blue Glacier, Washington, USA | Valley glacier, ablation zone, temperate | 9 surface locations, 144 sections |
| Taylor (1963) | Burroughs Glacier, Alaska | Ablation zone, temperate | 2 surface locations |
| Higashi (1967) | Mendenhall Glacier, Alaska | Temperate | 30 samples |
| Jonsson (1970) | Isfallsglaciären (and Storglaciären), Sweden | Small valley glacier, ablation zone, polythermal | 9 locations at glacier base (tunnel), oriented ice blocks, 81 sections |
| Hooke (1973) | Barnes ice cap, Baffin Island, Canada | Ice cap margin, cold (-10°C) | 4 shallow ice cores and surface locations; 18 sampling sites in intervals of $\sim 5\text{--}10\text{ m}$ |
| Vallon et al. (1976) | Vallée Blanche, Mt. Blanc, France | Valley glacier, accumulation zone, 3550 m a.s.l., temperate | Ice core (30 m firn and 150 m ice); ca. 50 sections in intervals of $\sim 10\text{ m}$ |
| Hudleston (1980) | Barnes ice cap, Baffin Island, Canada | Plateau glacier, 2400–2900 m a.s.l., temperate | 8 ice cores (along-flow transect), 5 in the ablation zone, 3 in the accumulation zone without reaching bedrock, 170 m in total; 348 sections with 3000 crystals |
| Tison and Hubbard (2000) | Glacier de Tsanfleuron, Switzerland | | Ice core, 309 m to bedrock; sampled $\sim 1\%$ in intervals of $\sim 10\text{ m}$ |
| Y. Li et al. (2017) | Guliya ice cap, Tibetan Plateau | Ice cap, 6200 m a.s.l., cold (-2°C) | |
| This study | Colle Gnifetti, Switzerland/Italy | Mountain glacier, 4484 m a.s.l., cold (below -10°C) | 92 sections from 72 m ice core, 12% of glacier thickness sampled, ~ 67000 crystals |

The most prominent studies showing fabric diagrams and giving microstructural quantities that could be found (and were accessible) in peer-reviewed literature in English are summarized here. We do not claim that this table is complete.

The ice core KCC was drilled to bedrock in 2013 (N 45° 55.736, E 7° 52.576, 4484 m a.s.l., Bohleber et al., 2018) in about 100 m distance to the ice core KCI, drilled in 2005 (Bohleber et al., 2013), on a flank towards an ice cliff. KCC is 72 m long and the firn-ice transition was found to be at a depth of about 36 m. A borehole temperature between -13.6°C in 13 m depth and -12.4°C at the bedrock was measured in 2014 (Hoelzle, Darms, Lthi, & Suter, 2011, no change since 2011, pers. comm. M. Hoelzle, University of Fribourg, 2014), implying that the ice is frozen to bedrock. Radiocarbon dating revealed a basal ice age of ~ 4000 years BP from the ice core sections at bedrock (Hoffmann et al., 2017). Therefore, we can assume that ice from different depths of the core has experienced quite a different deformation history despite the ice core’s short length. Surface velocities at Colle Gnifetti were last measured from 2014 to 2016 and were of the order of 1 m a^{-1} close to the KCC borehole (Licciulli, 2018).

3 Laboratory Measurements and Processing

The ice core was stored at -18°C during transport and at -30°C before and after processing. Vertical thin sections ($33 \times 100\text{ mm}^2$, ca. $300\ \mu\text{m}$ thick) of the ice core were prepared with a standard microtoming procedure in 13 intervals between 25 m and 72 m depth. Over each interval they provide continuous records of up to 110 cm length in depth. *Continuous* refers to measurements of several adjacent sections in an ice core segment or bag. The azimuthal orientation of the sections may vary due to deliberate rotation of core segments during processing, which was noted and considered in the analysis, or accidental loss of orientation information during core logging. Crystallographic c-axis orientations were measured by means of polarized light microscopy (e.g. Heilbronner & Barrett, 2014; Peterzell, Kohlmann, Wilson, Seiler, & Gleadow, 2009; Wilson, Russell-Head, & Sim, 2003) with an automatic fabric analyser from Russell-Head Instruments (FA, model G50). Additionally, microstructural maps (Kipfstuhl et al., 2006) were acquired with a Large Area Scanning Microscope (LASM, Schäfer+Kirchhoff, Binder, Garbe, Wagenbach, Kipfstuhl, & Freitag, 2013) from the same section surface to obtain supporting information on pore space and grain boundaries. The KCC fabric data consist of 92 vertical thin sections and samples 12% of the entire ice core.

For each pixel of the thin section the measurement provides the orientation of the crystallographic c-axis by two angles, azimuth ϑ in the interval $[0, 2\pi]$ and colatitude φ in the interval $[0, \pi/2]$, with respect to the (nearly) vertical ice-core axis. The c-axis is expressed as a vector \vec{c} with unit length:

$$\vec{c}(\vartheta, \varphi) = (\sin(\varphi) \cos(\vartheta), \sin(\varphi) \sin(\vartheta), \cos(\varphi)). \quad (1)$$

The obtained FA data are processed using the open-source software **cAxes** developed at AWI (Eichler, 2013) for automatic pixel-wise image analysis, excluding pixels with quality below 50% (Peterzell et al., 2009) and setting a lower grain size threshold of $A_{\min} = 0.2\text{ mm}^2$ (constrained by the section thickness $> 200\ \mu\text{m}$).

Eigenvalues λ_i ($i = 1, 2, 3$; $\lambda_1 \leq \lambda_2 \leq \lambda_3$; $\sum \lambda_i = 1$) of the second-order orientation tensor $a_{ij}^{(2)}$ are calculated from the c-axis distribution of a thin section sample (Adam, 1989; Durand, Gagliardini, et al., 2006; Wallbrecher, 1986) and are weighted with the grains’ cross-section area A as proposed and discussed by Gagliardini, Durand, and Wang (2004). These are standard parameters to classify the type and strength of the crystallographic-preferred orientation and can be associated with different deformation regimes (e.g. Weikusat, Jansen, et al., 2017). We can calculate a statistical uncertainty for the eigenvalues σ_λ depending on the sample size of the crystal ensemble N_g following Durand, Gagliardini, et al. (2006):

$$\sigma_\lambda = (-1.64\lambda^2 + 1.86\lambda - 0.14)N_g^{-1/2}. \quad (2)$$

We remark that this equation will not give a positive uncertainty estimate below $\lambda = 0.081$. The uncertainty σ_A , which accounts for sectioning and population effects when calculating the mean grain size \bar{A} as determined from the number of pixels per grain in

175 a 2-D section, is estimated by

$$176 \quad \sigma_A = 2\bar{A}(0.02 + 0.44 N_g^{-1/2}). \quad (3)$$

177 However, this uncertainty is based on the assumption of regular grain features during
 178 normal grain growth (Anderson, Grest, & Srolovitz, 1989; Durand, Gagliardini, et al.,
 179 2006) which can no longer be assumed to be the dominant process as soon as RX pro-
 180 cesses set in (Faria et al., 2014b). The measurement uncertainty associated with the FA
 181 instrument was shown to be small enough to be neglected compared to the sampling un-
 182 certainty (Montagnat et al., 2012).

183 For the centimeter-scale analysis along the measured core segments we apply a slid-
 184 ing window frame of 2 cm height and 2 mm step size for the calculation of second-order
 185 orientation tensor eigenvalues and mean grain size. The median number of grains per
 186 window with area $33 \times 20 \text{ mm}^2$ is 143 with an interquartile range of 90 – 217.

187 Additionally, the high-resolution density profile for the upper 50 m of the ice core
 188 was determined by evaluating 2D projected X-ray computer tomography images (XCT,
 189 Freitag, Wilhelms, & Kipfstuhl, 2004). Visual stratigraphy images (Svensson et al., 2005)
 190 were recorded over the full length and qualitatively analysed with focus on the visibil-
 191 ity of well-defined layers and layer inclination. The number density and size of closed-
 192 off bubbles in the ice was determined from the LASM images using the open-source soft-
 193 ware *ImageJ*. The records of dust-related impurities and stable water isotopes presented
 194 in Bohleber et al. (2018) were available for comparison with the crystal data (more ex-
 195 tensively discussed in Kerch, 2016).

196 4 Results

197 4.1 Physical Properties Profile

198 We present the comprehensive data set from a classical thin section-scale fabric anal-
 199 ysis, but from continuously measured intervals instead of just one thin section sample
 200 per interval. Figure 1 shows the depth profile of area-weighted second-order orientation
 201 tensor eigenvalues for KCC, the mean and maximum grain size per section and the high-
 202 resolution density from XCT. The uncertainties are calculated with Eq. 2 and 3. Ad-
 203 ditionally, the standard deviation of the grain size distribution is indicated towards higher
 204 grain sizes. Schmidt diagrams for individual thin sections are shown in appendix B, Fig. B.1.

205 The data reveal how a distinct anisotropy of c-axes evolves with depth. A weakly
 206 anisotropic fabric can already be observed in the firn. The crystal anisotropy quickly in-
 207 creases across the firn-ice transition down to 53 m, developing a broad single-maximum
 208 fabric. In the lower 20 m of the core a small portion of grains are forming a weak gir-
 209 dle pattern in combination with narrowing of the strong single maximum, with a min-
 210 imum spherical aperture (cone angle) of 9° at the bedrock, and increasingly inclined to
 211 the vertical, with a maximum inclination angle of the c-axis eigenvector at bedrock of
 212 27° .

213 Remarkable is the high variability of eigenvalues within continuously measured inter-
 214 vals, at all depths of the ice core, with the exception of the deepest 70 cm long inter-
 215 val close to bedrock. This variability exceeds the estimated population-dependent un-
 216 certainty and does not show a discernible trend. In the following, the c-axis eigenvalue
 217 λ_{3e} is considered as a single parameter to further investigate the small-scale fabric vari-
 218 ability, as single-maximum fabric is recognized as the dominant fabric type. Table 2 lists
 219 the average c-axis eigenvalue and difference between extreme values for each interval. The
 220 difference exceeds 0.1 in most intervals.

221 The mean grain size \bar{A} lies between 1 mm^2 and 18 mm^2 calculated from 155–1707
 222 grains per section, with 75% of the entire grain population smaller than 8 mm^2 . Note
 223 that for most sections 75% of all grains are smaller than the calculated mean grain size.
 224 The grain size tends to increase from the firn towards a depth of approximately 48 m and
 225 subsequently shows a sharp decrease in the interval near the bedrock. However, the vari-

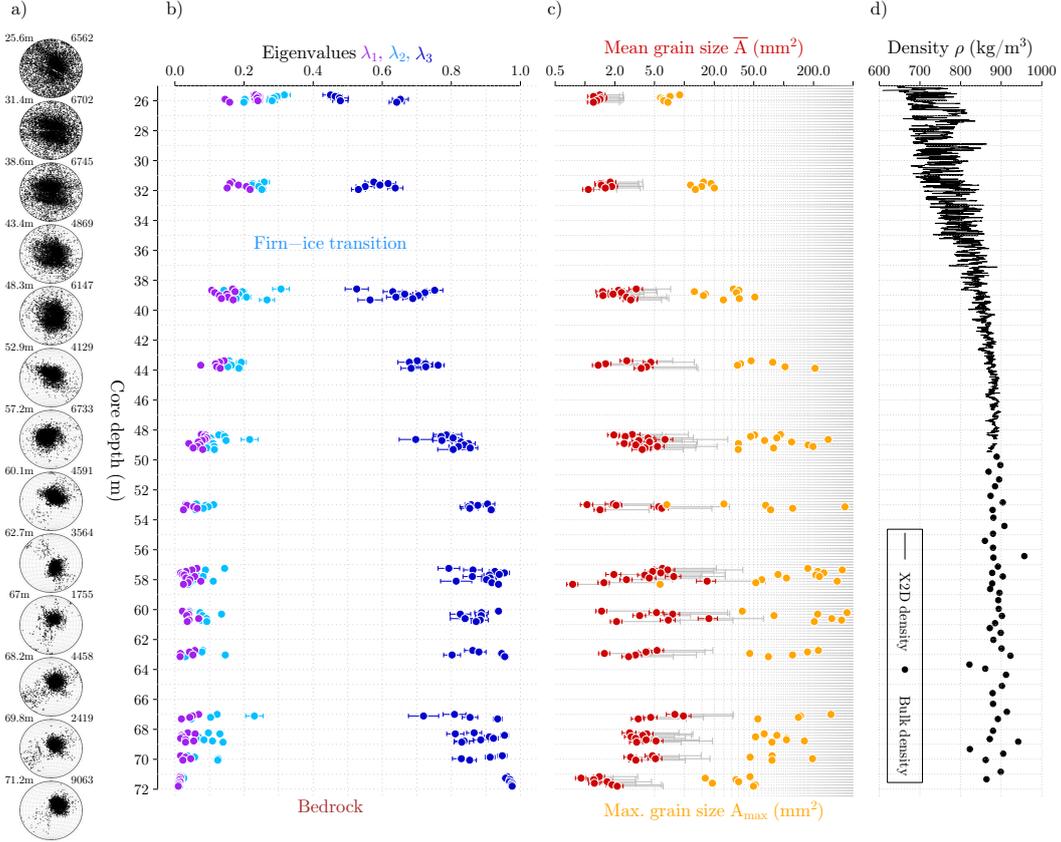


Figure 1. Evolution of a) c-axis distributions, b) c-axis orientation tensor eigenvalues (area-weighted), c) grain sizes and d) density with depth. Uncertainties $2\sigma_\lambda$ (blue) and $2\sigma_A$ (red) are calculated with Eq. 2 and 3. The standard deviation of the mean grain size (red) is additionally shown as grey whiskers (on a logarithmic scale). The maximum grain size is shown in orange. C-axis distributions are presented per continuously measured interval (Schmidt diagrams with the centre corresponding to the ice core axis). Each dot represents a single grain as identified from the thin section data. The top depth and number of grains are indicated on the top left/right of each pole figure. The azimuthal orientation of the Schmidt diagrams varies but could not be further constrained.

226 ability in mean grain size within an interval increases, in conjunction with a strong in-
 227 crease in the maximum grain size A_{\max} up to 4 cm^2 in 48–68 m depth. The calculated
 228 uncertainty associated with the sectioning and population effect for calculating the mean
 229 grain size is much lower than the standard deviation.

230 4.2 Visual Stratigraphy

231 We identify six depth zones in which the characterisation of the visual stratigraphy
 232 changes (Table 3, Fig. B.3), considering the contrast between and the thickness of
 233 alternating lighter and darker bands in the visual stratigraphy linescan images. For most
 234 of the core a stratigraphic order can be observed, which, however, cannot be discerned
 235 in the images in the bottom 5 m. The layer inclination varies strongly in the firn and in-
 236 creases towards bedrock in the ice with maximum values of 30° . As we do not know the
 237 orientation of the ice core (and thus that of the image plane) with respect to the geo-
 238 graphical coordinates and flow direction, layer inclination values can only be minimum

Table 2. Mean c-axis eigenvalues $\bar{\lambda}_3$ with standard deviation and variability for each interval (indicated by top depth and length) of continuous measurements, calculated on the thin section (Fig. 1) and cm-scale (Fig. 3).

| Top depth (m) | Length (m) | Mean c-axis eigenvalue $\bar{\lambda}_3$ (thin section-scale) | Variability $\lambda_{3,\max} - \lambda_{3,\min}$ (thin section-scale) | Mean c-axis eigenvalue $\bar{\lambda}_3$ (cm-scale) | Variability $\lambda_{3,\max} - \lambda_{3,\min}$ (cm-scale) |
|---------------|------------|---|--|---|--|
| 25.61 | 0.59 | 0.53 ± 0.09 | 0.2 | 0.54 ± 0.12 | 0.48 |
| 31.43 | 0.6 | 0.58 ± 0.04 | 0.11 | 0.6 ± 0.06 | 0.31 |
| 38.58 | 0.79 | 0.66 ± 0.07 | 0.23 | 0.67 ± 0.08 | 0.4 |
| 43.37 | 0.6 | 0.71 ± 0.03 | 0.08 | 0.72 ± 0.06 | 0.25 |
| 48.3 | 1.04 | 0.8 ± 0.04 | 0.16 | 0.81 ± 0.07 | 0.34 |
| 52.94 | 0.49 | 0.88 ± 0.03 | 0.07 | 0.87 ± 0.04 | 0.19 |
| 57.24 | 1.14 | 0.89 ± 0.05 | 0.16 | 0.9 ± 0.05 | 0.21 |
| 60.1 | 0.8 | 0.88 ± 0.04 | 0.11 | 0.9 ± 0.05 | 0.21 |
| 62.73 | 0.51 | 0.89 ± 0.06 | 0.15 | 0.9 ± 0.07 | 0.3 |
| 67 | 0.4 | 0.83 ± 0.09 | 0.21 | 0.85 ± 0.10 | 0.45 |
| 68.23 | 0.72 | 0.88 ± 0.05 | 0.14 | 0.89 ± 0.06 | 0.25 |
| 69.76 | 0.4 | 0.89 ± 0.05 | 0.12 | 0.89 ± 0.05 | 0.17 |
| 71.16 | 0.71 | 0.97 ± 0.01 | 0.02 | 0.96 ± 0.01 | 0.05 |

estimates. Also, a possible borehole inclination during drilling might have an unquantifiable effect on the layer inclination. However, borehole inclination measurements from 2016 revealed inclination angles below 5° between 10 and 60 m depth and $5-10^\circ$ down to bedrock (Licciulli, 2018). This is not sufficient to account for the observed layer inclination in the visual stratigraphy record and has probably mostly developed over 3 years since the drilling.

About 20 melt layers (thickness $\mathcal{O}(\text{cm})$) were found in the firn part of the ice core, as is to be expected at this drilling site (Alean, Haeberli, & Schdler, 1983), often with coarse grains in the vicinity of the melt features.

On visual inspection of the thin section microstructure images, we assess how often we observe a distinct change in grain size in layers, regardless of the mean grain size in these layers. In more than a third of all sections we find such grain size layer transitions. Fine-grained layers (19 observations, Fig. 2) occur most frequently at depths between the firn-ice transition and 60 m. An estimated 17% of all thin section samples contain thick layers of large grains (average maximum grain size: 2.5 cm^2) with some being part of the same large-grain layer extending over more than one section, between a depth of 48 and 69 m. Below 62 m grains of similar size are often clustered within sections of heterogeneous size distribution. Almost a third of all sections exhibit a heterogeneous grain size distribution without a clear layered structure. Only one third of all sections have a nearly homogeneous grain size distribution, mostly in the firn and very close to bedrock.

4.3 High-resolution Physical Properties

Following the stratigraphic observations we acknowledge the need to study fabric and microstructure on a smaller-than-section scale. We show the results from the high-resolution centimeter-scale analysis for all measured ice core intervals in Fig. 3. The c-axis eigenvalue λ_3 is given, together with the mean grain size \bar{A} , grain and bubble number, and the particle dust record from continuous flow analysis (Bohleber et al., 2018). A depth interval is shown in detail in Fig. 4, including the density ρ and bubble area A_b .

Table 3. Observations from visual stratigraphy line scan images (Fig. B.3). Several zones are identified by changes in the qualitative description; changes in inclination occur in slightly different depth zones. Inclination is counted from the horizontal perpendicular to the ice core axis.

| Zone | Depth | Qualitative description | Inclination |
|------|-----------|---|--|
| I | 0 – 14 m | bands (thinner than 25 cm) of weak contrast | 0 – 43 m: |
| II | 14 – 38 m | well-defined transitions between lighter/darker bands with decreasing thickness, bands (< 10 cm) disappear at the firn-ice-transition | 5 – 15° in 59% of images, varying strongly |
| III | 38 – 55 m | more uniform appearance, thickness of weakly discernible alternating bands exceeds those of zone II | 43 – 51 m: no or little inclination |
| IV | 55 – 64 m | clearer banding without well-defined transitions | 51 – 63 m: 10 – 15° |
| V | 64 – 71 m | layers ~ cm with good contrast, but increasingly irregular appearance | 63 m – bedrock: |
| VI | 71 – 72 m | uniform appearance | 15 – 20°, max. 30° |

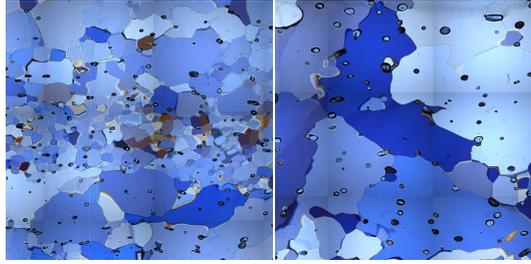


Figure 2. FA image details ($2.5 \times 2.5 \text{ cm}^2$). Left: Example of a fine-grained layer. Right: Example of a coarse-grained layer.

267 The high-resolution analysis reveals an extreme variability in the fabric and mi-
 268 crostructure parameters on the sub-decimeter scale with gradients in eigenvalue of up
 269 to 0.2 over a few centimeters. The uncertainty is small compared to the variations. The
 270 difference between extreme values within an interval is roughly twice as high as when
 271 comparing per-section eigenvalues (Table 2). Below 53 m depth the eigenvalue appears
 272 to spike towards lower values while it is high on average (Figure 3), matching the ob-
 273 servation of layers with larger and less oriented grains within a more strongly oriented
 274 matrix. Fine-grained layers cannot be resolved well as they are often thinner than the
 275 sliding computation window. The variability in the c-axis eigenvalue is accompanied by
 276 an equally variable mean grain size which is changing often in accordance with the eigen-
 277 value, with a significant (Spearman’s) rank correlation coefficient (r_s) of $-0.33 < r_s <$
 278 -0.79 for 8 out of 13 intervals (Appendix C: correlation computation). The anticor-
 279 relation between the mean grain size and the number of bubbles N_b is strong in all inter-
 280 vals ($-0.38 < r_s < -0.91$). In 10 out of 13 intervals we obtain significant correlation
 281 coefficients between the mean grain size and the dust content ($-0.51 < r_s < -0.9$).
 282 The close-up from 43 m depth (Fig. 4) shows that for a higher density ρ there are more
 283 but smaller bubbles ($0.56 < r_s < 0.87$ above 44 m depth). Furthermore, the density
 284 at this depth changes in accordance with the dust content.

5 Discussion

5.1 Depth-evolution of Crystal Anisotropy

The presented data of crystal-preferred orientation from a non-temperate, midlatitude, high-altitude glacier shows clearly that the development of an anisotropic CPO is not confined to the larger scale of polar ice sheets. In contrast, already in the firn, at one third of the glacier depth, we observe an anisotropic CPO pattern, as has been hypothesized by Diez et al. (2014) following a seismic velocity analysis at the site. The dominant fabric type is a single-maximum, which, in the upper part of the glacier, results from vertical compression in the firn and across the firn-ice transition, and may be enhanced by simple shear due to the surface slope at the flank of CG. In the lower third of the ice core, the inclined single maximum is consistent with simple shear (Azuma, 1994; Llorens et al., 2017; Qi et al., 2019, in press), as we can assume that the glacier is frozen to bedrock. This assessment is further supported by a recent ice-flow model for Colle Gnifetti (Licciulli, 2018), yielding high vertical strain rates $\dot{\epsilon}_{zz} < 1.9 \cdot 10^{-9} \text{ s}^{-1}$ and an elevated horizontal strain rate $\dot{\epsilon}_{zy}$ in the upper 10 m of the firn, and a steadily increasing horizontal strain rate from 30 m depth to bedrock ($\dot{\epsilon}_{zy} < 2.9 \cdot 10^{-10} \text{ s}^{-1}$). The faint girdle pattern we see in the Schmidt diagrams of the last 10 m above bedrock is indicative of a component of diverging flow (A. J. Gow & Meese, 2007). This suggests a more complex flow pattern in the bottom part of the glacier than the surface flowline implies at the KCC drilling site (Bohleber et al., 2018). It runs towards the eastern ice cliff, possibly with a component of extensional flow towards the saddle. However, the girdle is not as clearly recognized in single sections of individual intervals, supporting our approach of assessing continuously measured intervals to obtain a better representation of the crystal anisotropy for a given depth. When examined individually, some thin section samples exhibit specific structures, which are, in fact, contributing to the interval girdle. Figure 5a shows steeply inclined bands of crystals with higher colatitude (darker red) that deviate from the surrounding grain matrix. We interpret these to be tilted-lattice bands, which, for polar ice cores, have been shown to indicate small-scale folding (Jansen et al., 2016). In Fig. 5b a number of grains distributed over the section are constituting a second maximum of c-axes which is oriented approximately 90° from the dominant single maximum towards the horizontal. It is currently not resolved why these grains deviate from the matrix but a second maximum is usually interpreted as a consequence of recrystallisation under simple shear conditions (Alley, 1992; Qi et al., 2019, in press).

By including anisotropic deformation and the most relevant recrystallisation mechanisms in a microstructural model, Llorens et al. (2016) find that the grain size is mainly controlled by the recrystallisation processes, accompanied by secondary effects like grain dissection (Steinbach et al., 2017). Thus, the observed grain size variability with concentrations of very large or small grains is indicative of ongoing recrystallisation in this shallow mountain glacier. Exemplary microstructure images are shown in appendix B, Fig. B.2, illustrating the occurrence of various grain and subgrain topological features such as bulging GBs and intricate subgrain boundary structures. These are typically observed in polar firn and ice as well as experimentally deformed ice under comparably low stress undergoing dynamic recrystallisation (Hamann, Weikusat, Azuma, & Kipfstuhl, 2007; Kipfstuhl et al., 2009; Weikusat, Kuiper, Pennock, Kipfstuhl, & Drury, 2017), but the discussion of further microstructural parameters is beyond the intended scope of this paper.

5.2 Short-scale Variability

Our analyses on the sub-section scale show that layers of distinct grain size and layers of varying fabric strength can be observed at many depths. The variations are much stronger than previously known from low-resolution studies which could imply effects on the integrity of stratigraphic records (Faria et al., 2010) and also be highly relevant for the interpretation of geophysical data to map intraglacial structures (Diez et al., 2015; Eisen et al., 2007; Fujita, Maeno, & Matsuoka, 2006; Hofstede, Diez, Eisen, Jansen, &

338 Kristoffersen, 2012; Horgan et al., 2008; Kerch, Diez, Weikusat, & Eisen, 2018; Matsuoka
339 et al., 2003). The comparison of high-resolution data series reveals a good correlation
340 between the presented microstructural parameters, but it is very difficult to establish the
341 order of causal dependencies; feedbacks between coupled processes are very probable.

342 There is evidence on a large spatio-temporal scale in ice cores, that the microstruc-
343 ture could be governed by the impurity content (Durand, Weiss, et al., 2006). However,
344 the process identified to be the main driver for this relation (Zener pinning) is under dis-
345 cussion (Eichler et al., 2017). This raises the question whether the same mechanisms could
346 also play a role on the cm-scale? Our results strongly suggest that this is the case as par-
347 ticle content and mean grain size show a strong anticorrelation. While the ice from Colle
348 Gnifetti is relatively pure for a midlatitude glacier due to the high altitude, its dust con-
349 tent exceeds the dust load in Holocene ice from polar cores by a factor of 10 (Wagen-
350 bach & Geis, 1989), increasing the chance of observing a connection between microstruc-
351 ture and particle content. Recently developed techniques such as LA-ICP-MS (Della Lunga,
352 Müller, Rasmussen, Svensson, & Vallelonga, 2017; Reinhardt et al., 2001; Spaulding et
353 al., 2017) could provide spatially resolved data on particle distribution with respect to
354 microstructural features to study the interaction processes in detail.

355 Another parameter which is highly variable from the deposition processes on the
356 glacier surface is the density. Extremes in density are melt layers that are observed not
357 only in alpine glaciers but also in polar ice sheets (Schaller et al., 2016). We cannot di-
358 rectly connect outstanding layers in microstructure or fabric with original melt layers.
359 However, even without melt layers, initial density fluctuations in the snow, which could
360 be linked to the concentration of calcium in polar cores (Hrhold et al., 2012), can be ex-
361 pected to evolve into a sequence of layers in which bubbles have different characteris-
362 tics, e.g. many small bubbles or fewer large bubbles. Steinbach et al. (2016) show in their
363 microstructure model how bubbles can initiate strain localisation. In our data we see a
364 clear connection between the bubble number density and the mean grain size, which feeds
365 back into the idea that grain size variations are connected to density variations (Roes-
366 siger, Bons, & Faria, 2014). As mentioned before, Llorens et al. (2016) show that recryst-
367 tallisation processes are most evident in grain evolution. An existing grain size variabil-
368 ity would be enhanced by this.

369 But how does all this affect the fabric evolution and what are the consequences for
370 the bulk deformation? The small-scale variability in the CPO mostly correlates weakly
371 with the microstructure variability in general, with local exceptions of strong correla-
372 tion, especially where there is a strong contrast in microstructure parameters. Thus, the
373 role of the microstructure variations for the fabric variations appears to be that of ini-
374 tial perturbations which subsequently lead to local enhancement of strain rate. This has
375 been hypothesized before (Montagnat et al., 2014) but has not been presented with much
376 evidence due to a lack of continuous data. As a result of locally enhanced deformation
377 small-scale shear zones (Hudleston, 2015) might develop. This would, in turn, allow for
378 thick layers of ice deforming at a lower rate to be dominated by recrystallisation. How-
379 ever, we do not have information on the lateral scale to draw conclusions on the effect
380 on the bulk deformation. We certainly recommend to collect high-resolution borehole
381 deformation data to be able to connect fabric information and in-situ deformation on
382 this particular length scale, e.g. by means of ultrasonic profiling (Gusmeroli, Pettit, Kennedy,
383 & Ritz, 2012), even if they can only provide a reduced two-dimensional representation
384 of CPO. Our results are in line with Svensson, Baadsager, Persson, Hvidberg, and Siggaard-
385 Andersen (2003) who conducted a case study on the seasonal variability in ice crystal
386 properties on the cm-scale and observe small-grain layers corresponding to the spring
387 season but cannot find a seasonal variability of the *c*-axis orientations. Additionally, we
388 find no evidence of a connection between the thickness of fabric layers with the thick-
389 ness of annual layers as determined from impurity records (Bohleber et al., 2018).

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5.3 Do We Need a New Scale-sensitive Statistical Approach?

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From a methodological standpoint, our study provides evidence that, a priori, one thin section sample does not provide sufficient statistics to draw reliable conclusions for a depth interval of several tens of meters. An apparent fabric variability may simply be the result of statistical undersampling. The statistical uncertainty as proposed by Durand, Gagliardini, et al. (2006) does not appear to be an appropriate uncertainty measure for our data if this uncertainty is supposed to account for the variation in adjacent samples. Consequently, the task at hand is to establish which sample length scale is in fact representative for a given interval-of-interest. The same holds for the sampling interval, i.e. the distance between samples where no measurements are conducted. This sampling interval and the sample length make it mandatory to readdress the estimation of an appropriate statistical error based on such representative scales.

Furthermore, our results demonstrate that the choice of a representative length scale for CPO data is also dependent on the research question, considering the activity of several contributing and coupled processes on different length scales. More specific, the evaluation of fabric data needs to be adjusted to the resolution of complementary methods applied to the ice core or glacier, e.g. radar and seismics as well as continuous logging techniques, to advance our understanding of glaciological processes across different methodological approaches and spatial scales.

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6 Conclusions

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This study confirms the observation of a comparable development of crystal anisotropy in cold alpine glaciers and in polar glaciers or ice sheets, despite a factor of 35 (compared with the NEEM ice core) difference in thickness. Additionally, we observe and describe a strong variability in the microstructural parameters and crystal orientation during a high-resolution analysis of continuous measurement intervals. Thus, we can assume that this variability should also develop in polar glaciers, which emphasizes the usefulness of non-polar glacier environments for the study of glacial processes in general. We consider different causes for the variability, recognising that different causes may dominate on different length scales. Acknowledging the interconnectivity of the analysed parameters, we conclude that the high variability in CPO may develop as a result of deformation localisation (shear zones) on different length scales, following initial differences in density and microstructural parameters which, in turn, are influenced by the impurity concentration and the seasonal cycle of environmental conditions. The possible consequences of such shear zones on bulk deformation (e.g. folds) and the interpretation of stratigraphic parameters are still under investigation (Bons et al., 2016; Montagnat et al., 2013; Ran et al., 2018).

Our results lead us to question the current single-section mode of CPO evaluation and call for a new statistical approach based on a higher and scale-specific data coverage, and/or an adjusted uncertainty estimate. Furthermore, our findings reinforce the demand for high-resolution and high-coverage proxy methods (e.g. ultrasonic sounding in the borehole and the laboratory (Mikesell, van Wijk, Otheim, Marshall, & Kurbatov, 2017; Vaughan, van Wijk, Prior, & Bowman, 2016), high-resolution dielectric profiling (Wilhelms, 2005), polarimetric radar (Drews et al., 2012; J. Li et al., 2018) and seismics (Kerch, Diez, et al., 2018)) to be applied in order to obtain a more complete understanding of the vertical and lateral crystal anisotropy structure and its role for large-scale glacier deformation.

Table A.1. Resolution of fabric studies on polar ice cores. *Section* refers to a typically 10 cm long thin section sample; *continuous* refers to measurements of several adjacent sections in an ice core segment or bag.

| Authors | Ice core | Length [m] | Fabric resolution; intervals between samples | estimated % of total length |
|--|-------------------|------------|--|-----------------------------|
| A. Gow et al. (1997) | GISP2 | 3054 | ca. 500 sections | 1.7 % |
| Thorsteinsson, Kipfstuhl, and Miller (1997) | GRIP | 3029 | ca. 100 sections; mostly 25–55 m intervals | 0.3 % |
| Wang et al. (2002) | NGRIP | 2930 | 142 sections; 5–66 m intervals | 0.5 % |
| Montagnat et al. (2014), Eichler, Weikusat, and Kipfstuhl (2013) | NEEM | 2540 | ca. 700 sections; 10 m intervals, partially continuous (bags of 55 cm) | 3 % |
| Treverrow, Jun, and Jacka (2016) | Dome Summit South | 1196 | 185 sections; 5–6 m intervals | 0.9 % |
| Azuma et al. (1999) | Dome Fuji | 2500 | ca. 240 sections; 20 m interval, partially continuous | 1 % |
| A. J. Gow and Meese (2007), DiPrinzio et al. (2005) | Siple Dome | 1004 | ca. 100 sections; 20 m intervals | 1 % |
| Wang, Kipfstuhl, Azuma, Thorsteinsson, and Miller (2003), Durand et al. (2007, 2009) | EPICA Dome C | 3260 | ca. 200 sections; 11–50 m intervals | 0.7 % |
| Weikusat, Jansen, et al. (2017) | EPICA DML | 2774 | 210 sections; 50 m intervals, partially continuous | 0.8 % |
| Montagnat et al. (2012) | Talos Dome | 1620 | 10–20 m intervals | 1 % |
| Fitzpatrick et al. (2014) | WAIS | 3405 | 20 m intervals | 0.5 % |

436 A Collection of major polar fabric studies from the last 20 years

437 B Complementary figures

438 C Correlation computation for cm-scale data series

439 As the data are in general not normally distributed we compute correlation coef-
 440 ficients between various data sets available in this study by considering the ranks of the
 441 bivariate data (Spearman correlation coefficient)

$$442 \quad r_s = 1 - \frac{6 \sum_{i=1}^n d_i^2}{n(n^2 - 1)} \quad (\text{C.1})$$

443 with the sample size of paired values n , the rank difference d and $-1 < r_s < 1$. The
 444 significance of the obtained correlation coefficient is tested. The data of this study can
 445 be regarded as a time series for which each data point is not necessarily independent of

446 its neighbours (Mudelsee, 2003, "serial dependence"). The autocorrelation of a data se-
 447 ries of length n for a lag τ is used to evaluate the non-randomness of the data and pro-
 448 vide the distance or time lag within which subsequent data points cannot be considered
 449 as independent. Considering τ and the short length of data sequences computed from
 450 continuous fabric data we derive an effective sample size, i.e. smaller than the original
 451 sample size, from which a correlation with another data set might be calculated under
 452 the assumption of randomness of the data points. Only eigenvalues calculated from suc-
 453 cessive, but not overlapping, frames are used to derive correlation coefficients, which lim-
 454 its the sample size to 50 data points per meter. By choosing a different frame set that
 455 is shifted by $2n$ mm (with $n = [1,9]$) and repeated calculation we obtain several values
 456 for the correlation coefficient for the same continuous interval, providing a robust esti-
 457 mate of correlation on the 2 cm scale. All bivariate data sets are downsampled to the lower
 458 resolution of the two variables.

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 461 cans) are published in the open-access database PANGAEA® (Kerch, Eisen, et al., 2018,
 462 doi:10.1594/PANGAEA.887838) and also available upon request.

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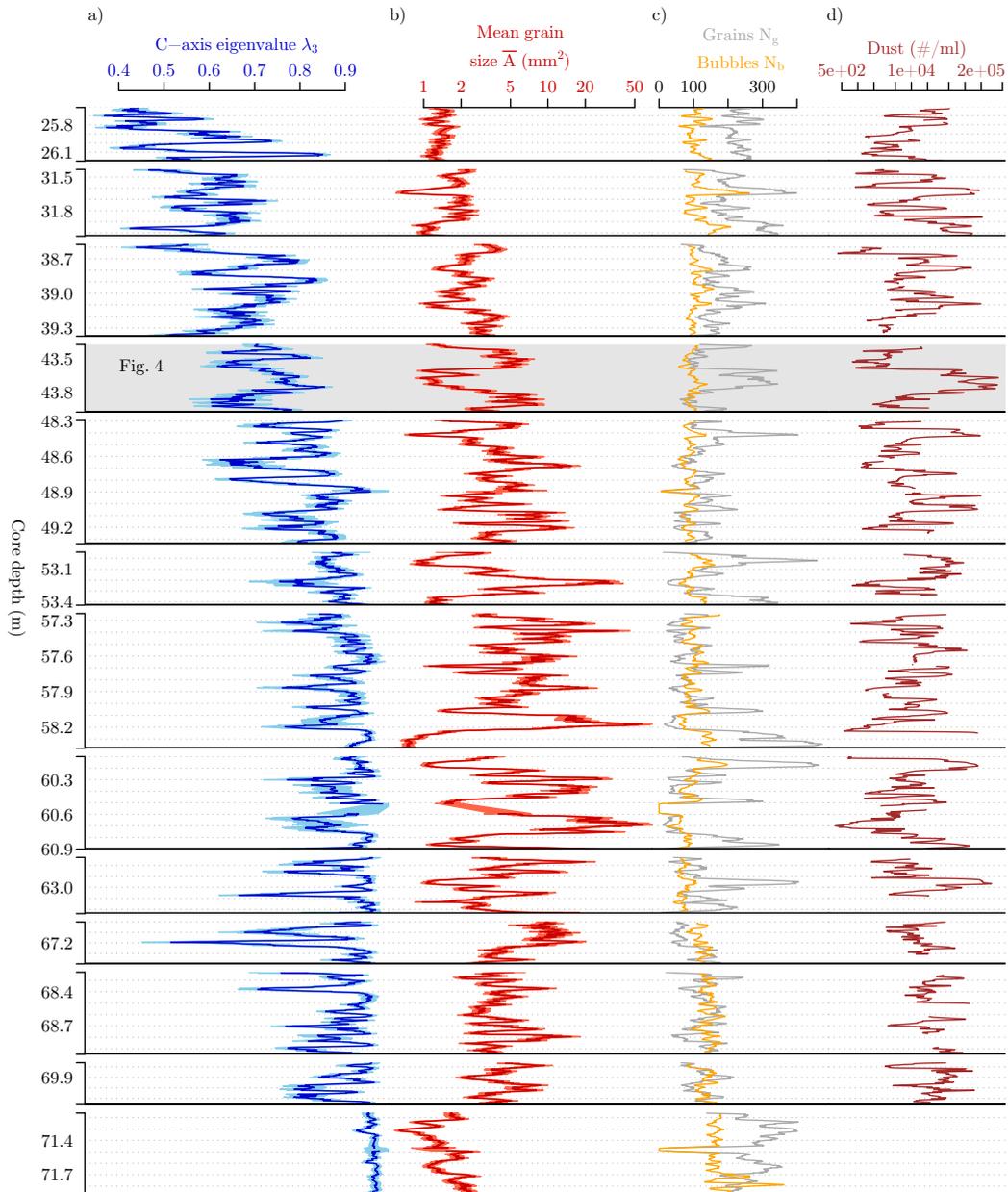


Figure 3. Physical properties parameters derived from cm-scale fabric analysis with 2 cm sliding window. The depth axis is discontinuous; horizontal black lines indicate axis breaks. a) Eigenvalue and b) mean grain size (logarithmic scale) are shown with uncertainty following Eq. 2 and 3. c) The grain number N_g is inversely related to the mean grain size and compared to the number of bubbles N_b . d) The insoluble particle concentration (referred to as "dust") from continuous flow analysis is additionally shown (logarithmic scale) for comparison. The shaded section is shown in detail in Fig. 4.

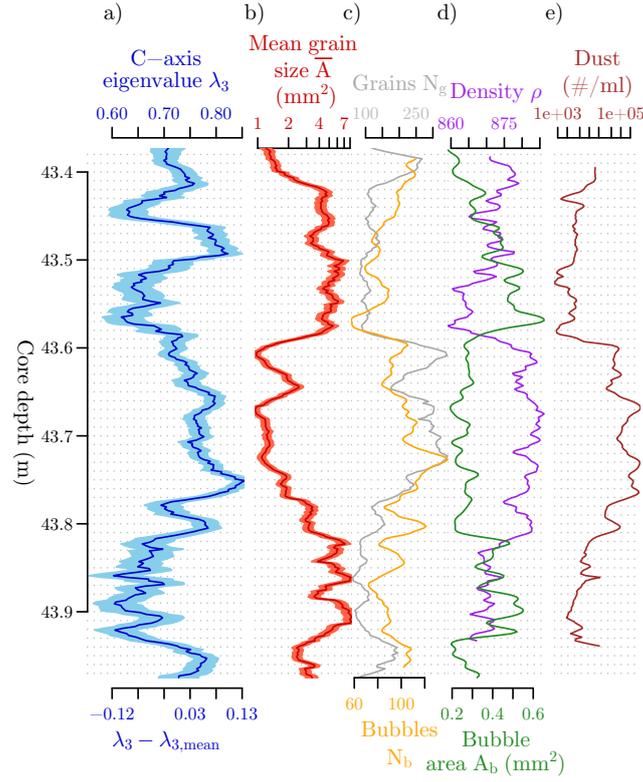


Figure 4. Excerpt from Fig. 3 showing a medium depth interval below the firn-ice transition. a) – c) as in Fig. 3. d) Density and the mean bubble area A_b are additionally shown. e) as d) in Fig. 3.

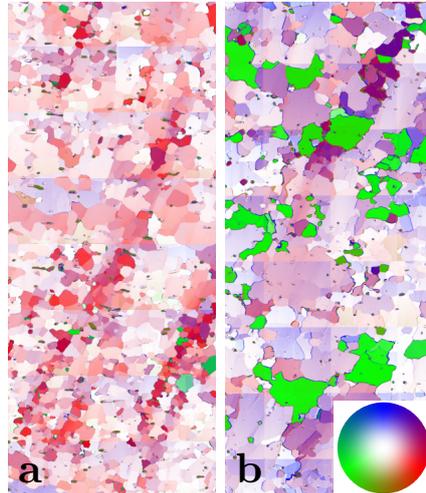


Figure 5. FA image details ($3.1 \times 7.5 \text{ cm}^2$). Left: Tilted-lattice bands (darker red) in 60.8 m depth. Right: Unconnected crystals (green) forming a second maximum inclined to the main single maximum (purple) in 70.1 m depth. The colour code indicates the c-axis orientation: the white center corresponds to a vertical orientation, parallel to the ice core axis.

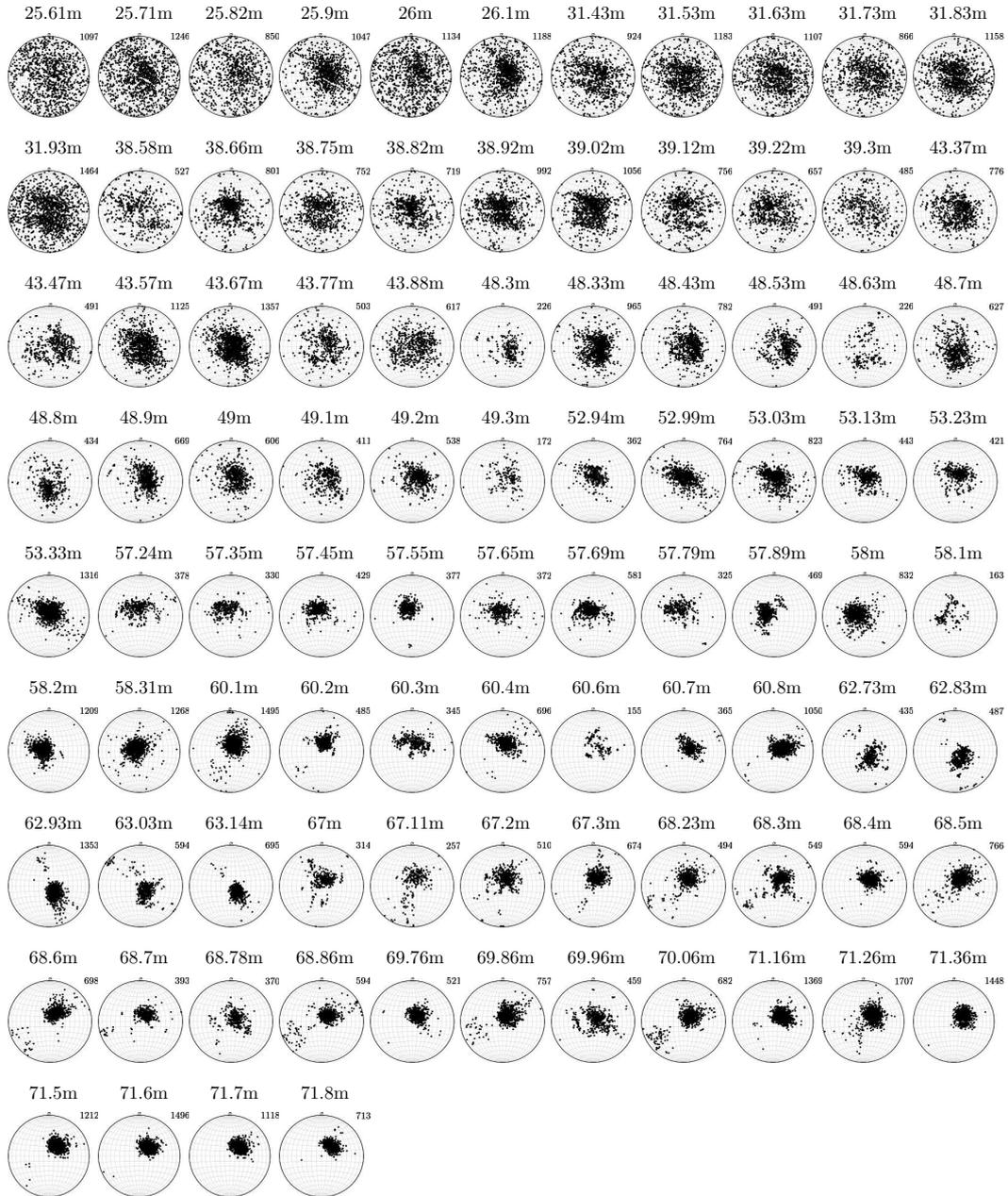


Figure B.1. Schmidt diagrams for all KCC thin sections. The number of grains is indicated on the top right of each polefigure.

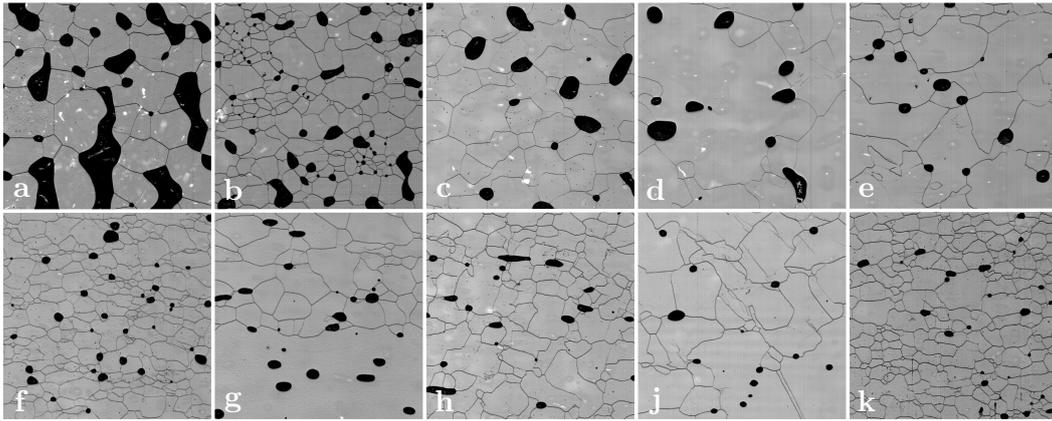


Figure B.2. Exemplary microstructure (LASM) image details (1 cm^2) from different depths of the ice core KCC. a) 25.9 m: firn. b) 31.63 m: deep firn with small-grain clusters in regions with many small bubbles. c) 38.66 m: bubble cluster just below firn-ice transition. d) Zigzagging subgrain boundaries. e) 49 m: Protruding grains and bubble alignment. f) 53.03 m: small-grain section. g) 57.35 m: small-to-large-grain transition; note that only within the large grain bubbles are not located on grain boundaries. h) 62.93 m: angular grains and inclined/elongated bubbles. j) 67.1 m: irregular grains and parallel GBs. k) 71.26 m: small-grain distribution close to bedrock.

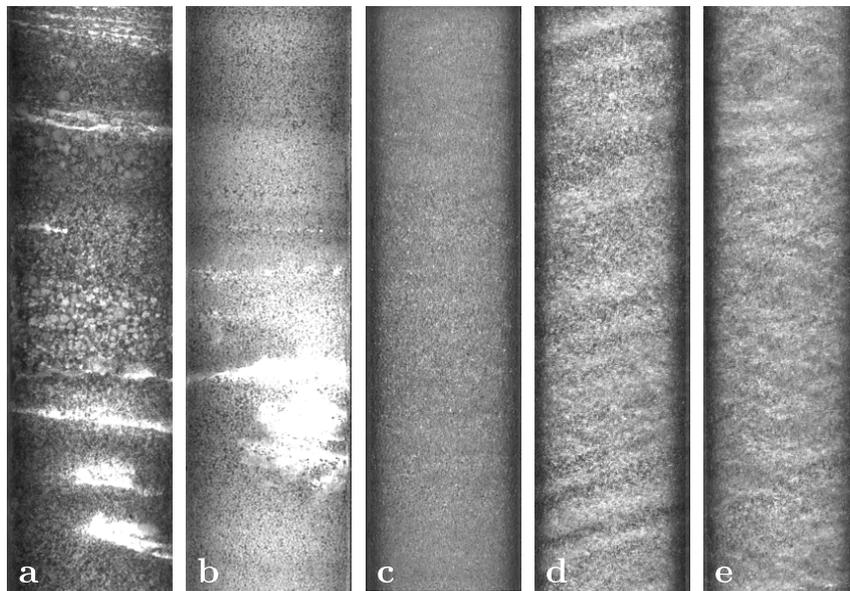


Figure B.3. Exemplary visual stratigraphy linescan images (approx. $8 \times 31 \text{ cm}^2$) from different depths of the ice core KCC. a) 6 m depth: thin melt layers, ice lenses and coarse grains in firn. b) 24 m depth: banding and large ice lense, presumably from percolating melt in deep firn. c) 47 m depth: barely visible transitions and no discernible inclination. d) 65 m depth: inclined layers. e) 70 m depth: Layers appear very irregular.