

1           **Geospace Concussion: Global reversal of ionospheric**  
2           **vertical plasma drift in response to a sudden**  
3           **commencement**

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17           **Key Points:**

- 18           • Dayside ionospheric plasma undergoes a transient motion from downward to up-  
19           ward during a sudden commencement (SC)
- 20           • Both observations and simulations show that the reversed vertical drift is a global  
21           response of the ionosphere to the SC
- 22           • The transient response is caused by a reversal of induced zonal electric field dur-  
23           ing the SC

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## Abstract

An interplanetary shock can abruptly compress the magnetosphere, excite magnetospheric waves and field-aligned currents, and cause a ground magnetic response known as a sudden commencement (SC). However, the transient ( $< \sim 1$  min) response of the ionosphere-thermosphere system during an SC has been little studied due to limited temporal resolution in previous investigations. Here, we report observations of a global reversal of ionospheric vertical plasma motion during an SC on 24 October 2011 using  $\sim 6$  s resolution SuperDARN ground scatter data. The dayside ionosphere suddenly moved downward during the magnetospheric compression due to the SC, lasting for only  $\sim 1$  min before moving upward. By contrast, the post-midnight ionosphere briefly moved upward then moved downward during the SC. Simulations with a coupled geospace model suggest that the reversed  $\vec{E} \times \vec{B}$  vertical drift is caused by a global reversal of ionospheric zonal electric field induced by magnetospheric compression during the SC.

## Plain Language Summary

It is well-known that a shock wave can suddenly compress objects they directly interact with. In this study, we report a special case in the geospace environment in which an interplanetary shock produced a concussion-like response in the ionosphere that was tens of thousands of kilometers away from the location where the shock first impacted. The ionized part of the atmosphere, or the ionosphere, was remotely connected to the magnetosphere - the region of geospace dominated by the Earth's magnetic field - via electric currents. When the magnetosphere was abruptly compressed after the shock arrival, a pair of electric currents flowing along the geomagnetic field lines was generated in the dayside mid-latitudes. The newly generated currents flipped the dayside ionospheric electric field from eastward to westward, leading to a downward motion of dayside ionospheric charged particles. Within one minute, the vertical motion and zonal electric field flipped again to the direction before the compression due to the generation of another pair of electric currents with an opposite sense to the first pair. This study depicts a global picture of the transient ionospheric response using multi-point high-resolution measurements and simulations with a state-of-the-art fully coupled geospace model.

## 1 Introduction

An interplanetary (IP) shock or a discontinuity in the solar wind can cause a sudden commencement (SC) in ground magnetic perturbations. It is also sometimes called sudden impulse or sudden storm commencement if followed by a geomagnetic storm (Joselyn & Tsurutani, 1990), hereinafter referred to as SC in this paper. Araki (1994) proposed a physical model to characterize the ground magnetic response during an SC. In the magnetosphere, two pairs of field-aligned currents (FACs) with opposite sense and induced electric fields are formed after the compression of the magnetosphere by the IP shock. The FACs and associated ionospheric currents produce a two-pulse signature in ground magnetic perturbations - a preliminary impulse (PI) followed by a main impulse (MI). The enhanced magnetopause current after the impinging of an IP shock produces a stepwise increase in the magnetic horizontal component, known as the SC.

SC impacts on the coupled magnetosphere-ionosphere-thermosphere (M-I-T) system have been extensively studied, including but not limited to the prompt acceleration of radiation belt electrons by the induced electric fields and subsequent ultra-low frequency (ULF) waves, enhanced ionospheric electron/ion temperature, F-region plasma uplift and frictional heating, and the generation of geomagnetically induced currents (e.g., Zong et al., 2009; Hudson et al., 2017; Zou et al., 2017; Belakhovsky et al., 2017; Kappenman, 2003). Global dayside ionosphere uplifting has long been reported to follow the SC due to the enhanced eastward electric fields on the dayside from penetrating interplanetary electric fields (e.g., Mannucci et al., 2005). However, less attention has been paid to the

74 ionospheric downward drift associated with the short-lived westward electric fields pre-  
 75 ceding the eastward electric fields. Early work since the 1960s reported frequency shifts  
 76 of high frequency (HF) Doppler sounders associated with SCs, called SCF (e.g., Davies  
 77 et al., 1962; Kanellakos & Villard, 1962; Huang et al., 1973). A model was proposed by  
 78 Huang (1976) to explain the HF Doppler effects of SCs and attributed the frequency shifts  
 79 to the vertical motions of the charged particles in the ionosphere forced by two oppos-  
 80 ing electric fields. According to HF Doppler sounder observations, SCF(+/-) is charac-  
 81 terized by a sharp positive frequency deviation spike followed by a prolonged negative  
 82 frequency deviation, and usually appears in the daytime and evening sectors (06-21 LT)  
 83 while SCF(-/+) is characterized by a negative frequency deviation followed by a positive  
 84 one, and occurs in the nighttime sector (21-06 LT). Previous reports of the positive pre-  
 85 liminary frequency deviations of SCF (i.e., the ionospheric downward motion) found they  
 86 are mostly constrained to low latitudes and not important due to small amplitudes and  
 87 a short duration (Kikuchi et al., 1985; Kikuchi, 1986).

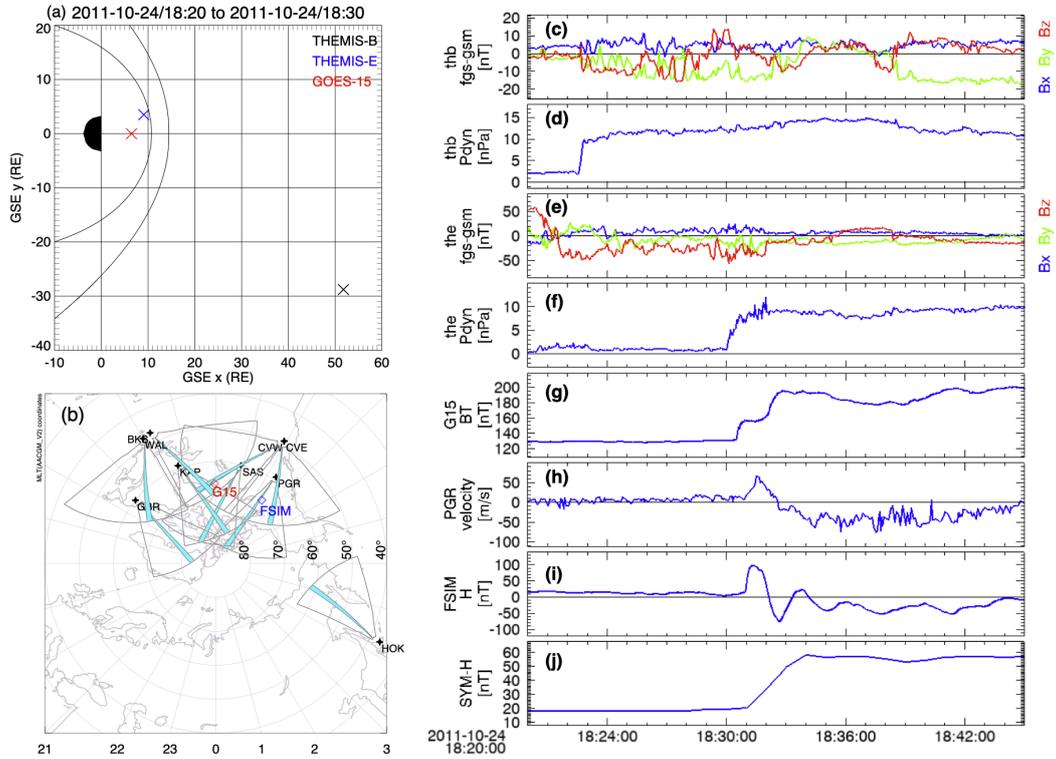
88 To understand the magnetospheric and ionospheric responses to SCs, many numer-  
 89 ical studies have also been conducted (e.g., Fujita et al., 2003a, 2003b; Kim et al., 2009;  
 90 Yu & Ridley, 2011; Zou et al., 2017; Ozturk et al., 2018; Fujita, 2019). However, most  
 91 previous SC simulations either ignored the processes occurring within one minute after  
 92 the SC or could not resolve such short time scale due to limited time resolution. For in-  
 93 stance, Kim et al. (2009) resolved MI-related vortex with global MHD simulations but  
 94 could not confirm PI-related vortex with 1-min resolution simulations. Zou et al. (2017)  
 95 investigated ionospheric SC effects with PFISR observations and global MHD simula-  
 96 tions and focused on the upward ion motion and plasma density and temperature vari-  
 97 ations. The transient impacts of SCs on the I-T system are still not well understood due  
 98 to lack of self-consistent M-I-T two-way coupled models and observations with high-temporal  
 99 resolution ( $< 1$  min). In particular, the ionospheric vertical drift related to SCs on global  
 100 scales and at high temporal resolution ( $< 1$  min) has not yet been well investigated or  
 101 understood. The main purpose of this study is to investigate the effects of SCs on the  
 102 I-T system and their temporal evolution using sub-minute, high cadence observations  
 103 and fully coupled whole geospace modeling.

## 104 **2 Observations and simulation results**

### 105 **2.1 Data Sets and Models**

106 Space and ground-based data sets and numerical simulations are used to investi-  
 107 gate geospace responses to an SC event on 24 October 2011 with a focus on the I-T ef-  
 108 fects. The data sets include two Time History of Events and Macroscale Interactions dur-  
 109 ing Substorms (THEMIS, Angelopoulos, 2009) spacecraft with THEMIS B located up-  
 110 stream in the solar wind and THEMIS E located inside the magnetosheath just before  
 111 the SC, the Geostationary Operational Environment Satellite (GOES, Singer et al., 1996)  
 112 15 satellite located inside the magnetosphere, and multiple ground magnetometers and  
 113 Super Dual Auroral Radar Network (SuperDARN) coherent scatter radars (Chisham et  
 114 al., 2007; Nishitani et al., 2019). The locations of the three spacecraft are shown in Geo-  
 115 centric Solar Ecliptic (GSE) coordinates in Figure 1a. Figure 1b shows the locations of  
 116 the ionospheric footprint of GOES 15 (red diamond), the Fort Simpson (FSIM) ground  
 117 magnetometer (blue diamond), and SuperDARN radar fields of view in altitude-adjusted  
 118 corrected geomagnetic (AACGM) coordinates (Shepherd, 2014).

119 The Multiscale Atmosphere-Geospace Environment (MAGE) model is a newly de-  
 120 veloped geospace model that is designed to study mesoscale processes in the coupled geospace  
 121 system. It consists of the Grid Agnostic MHD for Extended Research Applications (GAM-  
 122 ERA) global MHD model of the magnetosphere (B. Zhang et al., 2019; Sorathia et al.,  
 123 2020), the Rice Convection Model (RCM) model of the ring current (Toffoletto et al.,  
 124 2003), Thermosphere Ionosphere Electrodynamics General Circulation Model (TIEGCM)



**Figure 1.** Left: locations of (a) THEMIS B (black) and E (blue) spacecraft and GOES 15 (red) satellite in the X-Y plane in GSE coordinates from 18:20 UT to 18:30 UT on 24 October 2011; (b) locations of the ionospheric footprint of GOES 15 (red), the FSIM ground magnetometer (blue), and SuperDARN radar fields of view and THEMIS mode camping beams (cyan) in AACGM coordinates at 18:32 UT. Right: space and ground observations from 18:20 UT to 18:45 UT of (c-d) interplanetary magnetic field components and solar wind dynamic pressure from THEMIS B spacecraft measurements; (e-f) magnetic field components and dynamic pressure from THEMIS E spacecraft measurements; (g) total magnetic field from the GOES 15 satellite; (h) Doppler velocity measurements from the SuperDARN Prince George radar (beam 12 and gate 11); (i) detrended horizontal magnetic field from the FSIM ground magnetometer; (j) SYM-H index.

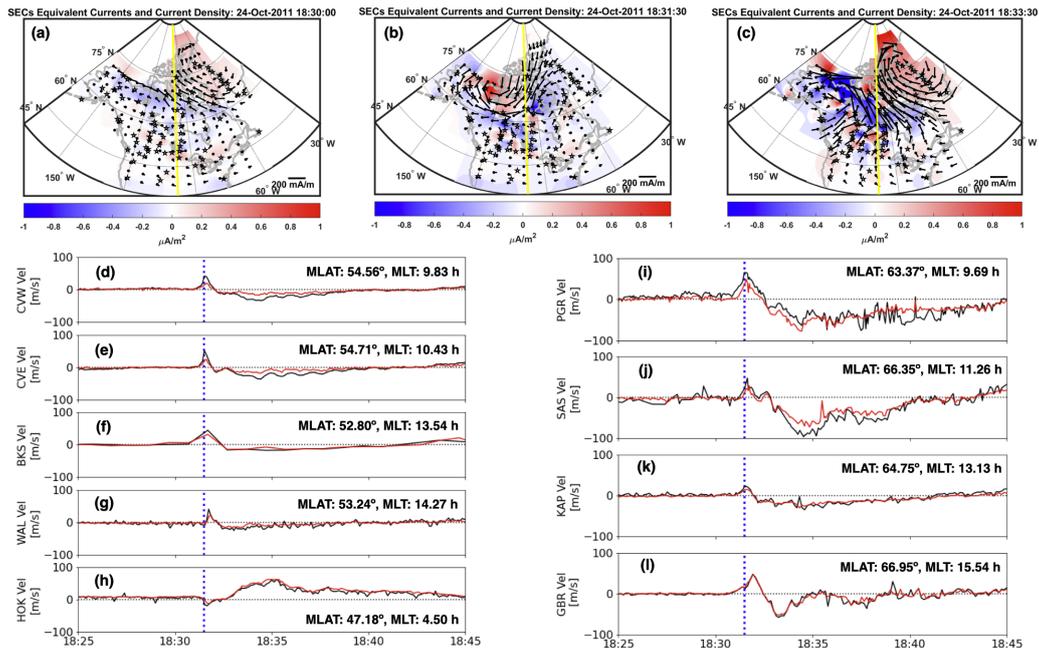
of the upper atmosphere (Richmond et al., 1992), and the RE-developed Magnetosphere-Ionosphere Coupler/Solver (REMIX) (Merkin & Lyon, 2010). Details about the model configuration used in this study can be found in Pham et al. (2022) and Lin et al. (2021).

## 2.2 Observations

Shown in Figure 1(right) are observations during the SC event on 24 October 2011. An IP shock was observed by THEMIS B at 18:22:30 UT with interplanetary magnetic field variations (Figure 1c) and a sharp solar wind dynamic pressure enhancement from about 2 nPa to 10 nPa (Figure 1d). THEMIS E spacecraft, initially inside the magnetosheath, observed gradually enhanced plasma pressure at  $\sim$ 18:30:00 UT and then crossed the bow shock to enter the solar wind at  $\sim$ 18:32:00 UT in response to the compression of the magnetosphere by the IP shock. The GOES 15 satellite detected enhanced magnetic field strength at 18:30:30 UT due to the compression by the IP shock (Figure 1g). A transient ( $\sim$  1 min) positive Doppler shift followed by a relatively long lasting negative Doppler shift was observed by multiple SuperDARN radars on the dayside with one example shown in Figure 1h from the Prince George (PGR) radar. A positive PI followed by a negative MI was observed by the FSIM ground magnetometer at 9.7 h magnetic local time (MLT) (Figure 1i), which is consistent with an upward FAC followed by a downward FAC after the SC in the morning sector as described by the Araki model. The SC signature with an enhancement in the SYM-H index (Figure 1j) occurred at 18:31 UT. Note SYM-H only has 1 min resolution. In addition to geomagnetic perturbations, this event was also reported by Shi et al. (2022) and Hartinger et al. (2020) to cause intense geoelectric field perturbations (1.67 V/km at 18:31:41 UT) over Minnesota in the United States.

The Spherical Elementary Current Systems (SECs; Weygand, 2009a, 2009b; Weygand et al., 2011) technique is applied to data from the widely spaced ground magnetometer arrays in North America and Western Greenland to obtain the equivalent ionospheric currents for this SC event. SECs equivalent currents (black vectors) and vertical current density (red-blue color map) are shown in Figure 2 top panels during (a) pre-SC at 18:30:00 UT, (b) PI phase at 18:31:30 UT, and (c) MI phase at 18:33:30 UT. The vertical yellow lines indicate local noon. The SECs equivalent currents from 18:25:00 UT to 18:31:00 UT look similar to those shown in Figure 2a with an anti-clockwise ionospheric current vortex and an upward vertical current (red) in the postnoon sector above  $60^\circ$  geographic latitude and an azimuthally extended downward current (blue) at  $65^\circ$ - $70^\circ$  geographic latitude. The PI related vertical currents first appeared at 18:31:00 UT at lower latitudes and moved poleward with an upward current (red) in the morning and a downward current (blue) in the afternoon sector at 18:31:30 UT as shown in Figure 2b. Figure 2c presents the follow up MI related vertical currents that are opposite to those in the PI phase, but is similar to the pre-SC currents (Figure 2a) with much stronger intensity and well defined current vortexes. These results are consistent with two pairs of FACs with opposite sense generated during the PI and MI phases from the physical model of SC in Araki (1994). An animation showing the evolution of the PI and MI related SECs currents at 10 s cadence can be found in the Supporting Information.

Ground backscatter echoes from SuperDARN coherent scatter radars are used to monitor ionospheric vertical drifts as shown in the bottom panels of Figure 2. Ground scatter echoes are typically formed during the daytime due to the high vertical gradient in the refractive index. The transmitted signal bends toward the ground and is reflected from surface roughness and returns to the radar following the same path. SuperDARN ground backscatter is sensitive to vertical ionospheric motions (Ponomarenko et al., 2003; Menk et al., 2003), and can be used to measure the vertical motion of the ionospheric layers through sunrise and sunset and also the vertical plasma motion associated with traveling ionospheric disturbances (e.g., Milan et al., 2013). In this paper, for the first time, this technique is used to study ionospheric vertical drifts associated with an



**Figure 2.** Top: Equivalent ionospheric currents (black vectors) and current density (red-blue color map with amplitude and sign given in the color bar at the bottom) at (a) 18:30:00 UT during pre-SC period, (b) 18:31:30 UT during the PI phase, (c) 18:33:30 UT during the MI phase. The vertical yellow line indicates local noon. Bottom: Doppler velocity from multiple SuperDARN radars (d-h) at middle latitudes and (i-l) high latitudes from 18:25:00 UT to 18:45:00 UT on 24 October 2011.

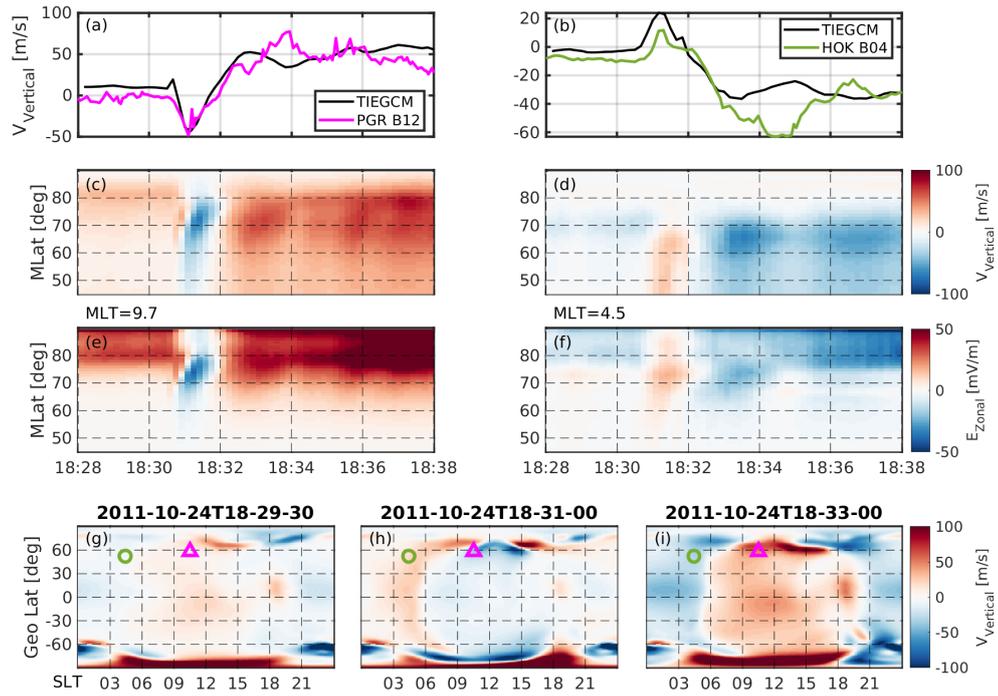
177 SC. This is made possible due to radars operating in a mode called THEMIS mode which  
 178 includes a camping beam; one that is revisited repeatedly during a typical scan. The THEMIS  
 179 mode is capable of sampling the camping beam (color coded in cyan in Figure 1b) ev-  
 180 ery  $\sim 6$  s and therefore capturing transient variations of  $< 1$  min associated with the SC.

181 The bottom panels in Figure 2 show Doppler velocity variations in ground scatter-  
 182 ter from multiple SuperDARN radars at middle to high latitudes. Black traces indicate  
 183 Doppler velocity obtained from a specific range-gate cell with the largest preliminary im-  
 184 pulse observed from the camping beam except for the BKS radar in Figure 2f which shows  
 185 observations from one normal beam 18 with a temporal resolution of 1 min. The median  
 186 velocity across multiple range gate cells from the selected beams at each recording time  
 187 was calculated and shown as red traces. The MLAT/MLT location of the ionospheric  
 188 reflection point of ground scatter at a specified range-gate cell is calculated assuming an  
 189 altitude of 250 km (Bristow et al., 1994) and shown on the right of each panel. A tran-  
 190 sient (1-2 min) positive Doppler shift followed by longer lasting ( $\sim 7$  min) negative Doppler  
 191 shift was observed by multiple SuperDARN radars on the dayside. Blue vertical dotted  
 192 lines indicate the time at 18:31:30 UT when the PGR radar first observed the peak of  
 193 the positive impulse. By contrast, observations from the Hokkaido East (HOK) radar  
 194 located post-midnight at  $\sim 4.5$  h MLT show the opposite Doppler velocity impulses (Fig-  
 195 ure 2h), that is, a transient negative Doppler shift followed by longer lasting positive one.  
 196 This is consistent with the HF Doppler sounder observations of SCF (+-) on the day-  
 197 side and SCF(-+) in the nighttime sector (21-06 LT). The positive (negative) Doppler  
 198 velocity from SuperDARN ground scatter indicates a downward (upward) plasma motion  
 199 which might be driven by a westward (eastward) electric field associated with the  
 200 SC. Note that the BKS radar only shows a clear positive impulse from Beams 18 to 22  
 201 (B18 are shown in Figure 2f). Due to a lower temporal resolution (1 min), the positive  
 202 impulse only consists of 1-2 data points which makes it difficult to be connected with  
 203 any physical phenomenon without the context provided by other high temporal resolu-  
 204 tion radar observations.

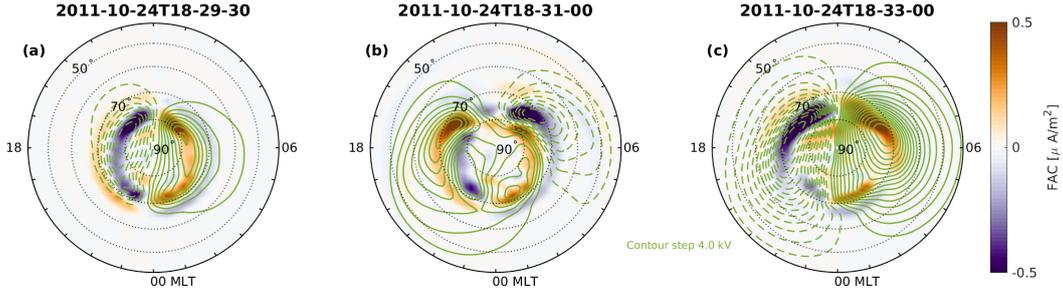
### 205 2.3 MAGE Simulations

206 In this study, we used the solar wind parameters measured by THEMIS B space-  
 207 craft to drive the MAGE model. As shown in Figure 1, THEMIS B was located upstream  
 208 of the bow shock, which provided closer to real-time information on the IP shock before  
 209 it arrived at the Earth with higher temporal resolution, compared to OMNI data. The  
 210 transient reversal of vertical plasma drifts shown in Figure 2 was reproduced by the MAGE  
 211 model. Figures 3a and 3b show the vertical plasma drift sampled from TIEGCM results  
 212 at two SuperDARN radar measurement locations, beam 12 of PGR (dayside near 9.7 h  
 213 MLT) and beam 4 of HOK (nightside at 4.5 h MLT), respectively. In this study, TIEGCM  
 214 has a time step of 5s and is output every 10 s. The observational data are shown with  
 215 the magenta and green curves for the two radars and the simulation results are shown  
 216 in black. Note that the SuperDARN Doppler velocity, which is positive for downward,  
 217 were transformed into the vertical direction by flipping the sign to directly compare with  
 218 TIEGCM outputs in Figures 3a-3b. The vertical drifts sampled at PGR turned down-  
 219 ward at 18:31 UT during the SC with a maximum speed of  $\sim 50$  m/s and became pos-  
 220 itive (upward) after 18:32 UT. The sampled vertical drifts at HOK showed a transient  
 221 upward motion of  $\sim 20$  m/s during the SC before turning downward after the SC. Note  
 222 that for this plot the SuperDARN measurements were shifted back in time by 30 s in  
 223 order to match the maximum downward drifts simulated by the model. The deviation  
 224 is likely due to uncertainty in the timing of solar wind parameters used to drive the model.

225 The similarity in measurements at multiple SuperDARN radars distributed widely  
 226 in local time and latitude suggests that the vertical plasma drift reversal is a global ef-  
 227 fect. Figures 3c and 3d show the keograms of vertical plasma drift sampled by the PGR  
 228 (9.7 h MLT) and HOK (4.5 h MLT) radars, respectively. The prenoon vertical drift was



**Figure 3.** (a-b) Median vertical plasma drifts measured by the SuperDARN PGR radar (magenta) and HOK radar (green), and TIEGCM samplings at the same locations (black). (c-f) Keograms of vertical plasma drifts and zonal electric field at 9.7 h MLT and 4.5 h MLT. (g-i) Vertical plasma drifts sampled at a mean altitude of 255 km from TIEGCM. The magenta triangle and green circle stand for the locations of the PGR and HOK radar observations, respectively.



**Figure 4.** MAGE-simulated northern ionospheric FACs (purple-orange color map) and convection responses to the IP shock. Positive currents (orange) are downward. The green contours show the convection potential separated by every 4.0 kV. Solid curves show positive potential.

229 downward over a broad range of latitudes for about one minute from 18:30:50 UT to 18:31:50  
 230 UT during the PI phase, while before and after the PI, it was upward at all latitudes.  
 231 By contrast, in the post-midnight sector, the vertical drift was downward at high lat-  
 232 itudes but it reversed to upward at middle and low latitude during the PI phase. This  
 233 ionospheric plasma motion is well described by the  $\vec{E} \times \vec{B}$  drift. Figures 3e and 3f show  
 234 the keograms of zonal electric fields at the same two MLTs. During the PI phase, the  
 235 zonal electric field was westward in the prenoon sector and eastward in the post-midnight  
 236 sector, which is consistent with the vertical drift response.

237 In Figures 3(g-i), we use the simulation results to depict the global picture of the  
 238 vertical plasma drift (zonal electric field) response during the SC. The vertical drifts were  
 239 sampled from TIEGCM outputs at a mean altitude of 255 km at the same three UTs  
 240 shown in Figures 2 (a-c) with the 30 s shift taken into account. The locations of PGR  
 241 and HOK measurements are denoted with a magenta triangle and a green circle, respec-  
 242 tively. Before the SC at 18:29:30 UT, the vertical plasma drifts were a few tens of m/s  
 243 upward on the dayside ( $6 < SLT < 18$ ) and slightly downward on the nightside. Dur-  
 244 ing the PI phase at 18:31:00 UT, however, the vertical drifts globally reversed to down-  
 245 ward on the dayside and upward on the nightside. The dayside downward plasma drift  
 246 reached a few tens of m/s at middle and low latitudes but exceeded 100 m/s at auro-  
 247 ral latitudes. The globally reversed vertical drifts only lasted for about one minute be-  
 248 fore they were reversed again, at 18:32:00 UT. The dayside upward drifts after the SC,  
 249 e.g., at 18:33:00 UT, were much stronger than those before the SC.

250 To understand the cause of the transient reversal of ionospheric vertical plasma drifts  
 251 and zonal electric fields, we analyzed the electrodynamic ionospheric response during the  
 252 SC with MAGE simulation results. Figure 4 shows the northern ionospheric FACs (purple-  
 253 orange color map) and convection potential (green contour) at the same three times as  
 254 shown in Figures 3(g-i). Before the shock arrival (Figure 4a), the ionosphere showed a  
 255 typical pattern of a pair of Region-1 FACs poleward of a pair of Region-2 FACs and two  
 256 cell convection with the dawnside (duskside) at positive (negative) potentials. As the shock  
 257 front arrived and propagated across the Earth (Figure 4b), a pair of Region-2 sense FACs  
 258 was generated at dayside mid-latitudes that propagated poleward. The postnoon down-  
 259 ward FAC and prenoon upward FAC requires a westward electric field for current clo-  
 260 sure, which drives a downward plasma motion on the dayside. The two-cell convection  
 261 before the SC was overtaken by a reversed two-cell convection gradually moving from  
 262 dayside to nightside. After the shock front completely passed over the Earth (Figure 4c),  
 263 the dayside ionospheric electric field reversed to eastward and the convection returned  
 264 to the regular two-cell pattern. An animation showing the MAGE-simulated evolution  
 265 of FACs and ionospheric convection pattern from 18:25:00 UT to 18:45:00 UT is provided

266 in the supporting information. The evolution of two pairs of FACs and their poleward  
 267 propagation from MAGE simulations are consistent with those from the SECs measure-  
 268 ments in Figures 2(a-c).

### 269 3 Discussion and Summary

270 Although the geospace response to an IP shock on scales of more than several min-  
 271 utes has been well investigated, the transient vertical ionospheric motion at sub-minute  
 272 resolution is scrutinized for the first time with high temporal resolution observations and  
 273 a coupled geospace model, the MAGE simulations. The THEMIS E spacecraft inside the  
 274 magnetosheath observed the arrival of the IP shock at  $\sim 18:30:00$  UT (Figures 1e-1f). Af-  
 275 ter about 30 s, the GOES-15 satellite located near noon detected the compression of the  
 276 magnetosphere (Figure 1g). On the ground, the PI- and MI-related FACs were clearly  
 277 seen following the SC as shown in Figures 2b-2c. MAGE simulations reproduced the geospace  
 278 responses including the westward and then eastward induced electric fields in the day-  
 279 side magnetosphere-ionosphere, two pairs of FACs with opposite sense, and ionospheric  
 280 convection reconfiguration associated with the newly formed FACs after the SC as shown  
 281 in Figure 4. Multiple-point radar measurements and the MAGE simulation results re-  
 282 veal that the concussion is a global response of the ionosphere to the IP shock. Figure  
 283 2 shows that transient downward motion was detected by all SuperDARN radars on the  
 284 dayside, while the HOK radar on the nightside detected upward ionosphere motion. Al-  
 285 though this study focuses on SuperDARN measurements, the transient vertical ionospheric  
 286 motion was also detected by other facilities. As shown in Figure S1, the Communica-  
 287 tions/Navigation Outage Forecasting System (CNOFS) satellite (de La Beaujardière et  
 288 al., 2004) detected a transient downward ion drift velocity of up to 30 m/s from 18:31:00  
 289 UT to 18:32:12 UT at 6.6 h MLT near the magnetic equator, similar to those reported  
 290 by R. Zhang et al. (2022) with a focus on ULF waves during SCs using CNOFS satel-  
 291 lite observations at a single location. MAGE simulation results reveal that the vertical  
 292 plasma drift was globally temporarily reversed on both the dayside and nightside (Fig-  
 293 ures 3g-3i). The PGR radar observations show the downward motion on the dayside reached  
 294  $\sim 70$  m/s at high latitudes (Figure 2i).

295 We suggest that the downward and subsequent upward ionospheric plasma verti-  
 296 cal drifts on the dayside were mainly driven by induced electric fields through  $\vec{E} \times \vec{B}$   
 297 with a transient westward electric field followed by a long-lasting eastward electric field  
 298 in the dayside magnetosphere and ionosphere. This evolution is identified in the Super-  
 299 DARN observations and reproduced by the MAGE simulations (Figures 3a and 3b). Al-  
 300 ternatively, the positive Doppler velocity in the SuperDARN ground scatter measure-  
 301 ments could result from changes in ionospheric refractive index and ray reflection height.  
 302 For instance, solar flares can instantaneously enhance the ionospheric electron density  
 303 and lower the F-region reflection height, causing the so-called Doppler flash (e.g., Kikuchi  
 304 et al., 1986; Chakraborty et al., 2018, 2021). However, this mechanism likely only plays  
 305 a minor role in the current study. We examined the vertical electron density profiles and  
 306 time series of F2 peak height (HmF2) at the locations of the PGR and HOK measure-  
 307 ments (Figures S2-S3) and found the relative variation of electron density was only 2%  
 308 and HmF2 variation was only  $\sim 2$  km during the PI phase, inadequate to cause the Doppler  
 309 shift measured by the radars. Nevertheless, it is possible that in some regions where shock  
 310 aurora are generated associated with the SC (Liu et al., 2015; Zhou et al., 2017), changes  
 311 in electron density might play a role. A further examination of the I-T effects during SCs  
 312 (e.g., shock aurora and electron temperature variations) is deferred to a future study us-  
 313 ing events when observations of these parameters are available (e.g., incoherent scatter  
 314 radar measurements from the 17 March 2015 storm).

315 To summarize, high-temporal resolution observations and the MAGE model sim-  
 316 ulation are used to investigate the effects of an SC on the geospace system, particularly  
 317 on the I-T system during the PI phase. We report for the first time using SuperDARN

ground scatter observations that the ionosphere undergoes a globally downward motion on the dayside and upward motion in the post midnight sector over 1 min during the PI phase, before it was gradually up lifted by an eastward electric field on the dayside during the longer-lasting MI phase. The high cadence outputs from the coupled geospace model of MAGE reveals for the first time that the ionospheric vertical motion related to SC is a global phenomenon with a larger impact than previously expected. This study advances our understanding of the effects of SCs in several ways:

1. Most previous studies focused on the dayside uplifting of the ionosphere due to limited temporal resolution while this study found that a transient downward drift ( $< 1$  min) precedes the ionosphere uplifting on the dayside following the SC.
2. This study utilized high temporal resolution ( $\sim 6$  s) ground scatter signatures in SuperDARN data to estimate ionospheric vertical drifts associated with an SC, whereas other SuperDARN observations using 1 min resolution data focused on ionospheric convection reconfiguration and radar backscatter echo responses associated with SCs (e.g., Coco et al., 2005; Kane & Makarevich, 2010; Boudouridis et al., 2011; Hori et al., 2012). Simultaneous observations from multiple SuperDARN radars provide direct evidence of the existence of the SC-related transient vertical drift in the ionosphere over a larger scale and with larger amplitudes than previously thought (e.g., Kikuchi et al., 1985; Kikuchi, 1986).
3. The coupled geospace model MAGE simulations with high temporal resolution revealed for the first time that the transient ion vertical drift associated with an SC is a global phenomenon (changes seen from the dayside to the nightside, and from the polar region to the equatorial region), whereas most previous MHD simulations concentrated on processes above 1 min time scale.

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## Data Availability Statement

The SECs are located at <http://vmo.igpp.ucla.edu/data1/SECS/>. The SYM-H index used in this paper was provided by the WDC for Geomagnetism, Kyoto (<http://wdc.kugi.kyoto-u.ac.jp/wdc/Sec3.html>). Access to SuperDARN data can be found at <http://vt.superdarn.org/tiki-index.php?page=Data+Access>. Data from the THEMIS mission can be found at <http://themis.ssl.berkeley.edu/data/themis/>. The GOES magnetic field data can be found at <https://satdat.ngdc.noaa.gov/sem/goes/data/full1/>. The MAGE simulation data are saved at this data repository: <https://doi.org/10.5065/xj5m-8t12>.

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