

# Blue Flashes as Counterparts to Narrow Bipolar Events: the Optical Signal of Shallow In-Cloud Discharges

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## Key Points:

- Blue luminous events are only detected in the 337 nm photometer and its filtered camera, with no signal in 777.4 nm optical band.
- Blue luminous events are found to be associated with the negative narrow bipolar events inside the thunderstorm.
- Narrow bipolar events have streamer-like structures and the numbers of streamer branches for them were evaluated.

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## Abstract

On August 7, 2019, eight Blue Luminous Events (BLUEs) above a thunderstorm over South China were detected by the Modular Multispectral Imaging Array (MMIA) of the Atmosphere-Space Interactions Monitor (ASIM), the Vaisala GLD360 global lightning network and the very low frequency (VLF)/low frequency (LF) sensor at Guangzhou. All events coincided with negative-polarity Narrow Bipolar Event (NBE) radio detections. The 337-nm photometer detected the strongest optical signal; in some events there was a weak signal in 180-230 nm but no signal was detected in 777.4 nm. The light-curves are consistent with sources spanning a few kilometers below the cloud tops and allow us to estimate the source extension and optical energy in the 337-nm band, corresponding to the energy emitted by fast breakdown involving around  $10^9$  streamer initiation events.

## Plain Language Summary

Installed on the International Space Station (ISS), the Atmosphere-Space Interactions Monitor (ASIM) is designed to observe Earth thunderstorms from space. Often it detects bursts of blue light emerging from active thunderclouds. These detections have been previously linked to radio signals named Narrow Bipolar Events (NBEs) that are routinely detected during a thunderstorm. Here we analyze the blue flashes from a storm that profusely produced negative-polarity NBEs. The optical signal can be understood as being produced by extended events close to the cloud top and we show that it is consistent with the barely understood electrical discharge process called fast breakdown that involves a huge number of thin, bright filaments called streamers.

## 1 Introduction

Blue Luminous Events (hereafter, BLUEs), such as blue jets and blue starters, are transient electrical phenomena that occur in active thunderstorms and are characterized by optical signals strongly dominated by the blue range of the spectrum and lasting from about one millisecond to hundreds of milliseconds. They were first reported by Wescott et al. (1995, 1996) in the framework of the Sprites 1994 aircraft campaign. After these initial reports, BLUEs have been observed and investigated mainly from space-based platforms: for example, from black-and-white camera of the Space Shuttle (Boeck et al., 1998), from the limb-pointing Imager of Sprites/Upper Atmospheric Lightning (ISUAL) onboard FORMOSAT-2 (Kuo et al., 2005; Chou et al., 2011, 2018), from off-the-shelf cameras on the International Space Station (ISS) (Chanrion et al., 2017) and, most recently, from the Modular Multispectral Imaging Array (MMIA) of the Atmosphere-Space Interactions Monitor (ASIM) onboard the ISS since April 2, 2018 (Soler et al., 2020; Neubert et

51 al., 2021). BLUES can also be observed from the ground, as reported by Edens (2011), but this is  
52 complicated due to Rayleigh scattering and extinction by intervening clouds.

53 As we define them here, BLUES comprise several phenomena. Blue jets normally emerge out  
54 of the cloud top and reach altitudes about 40 km - 50 km at speeds of the order of 100 km/s (Wescott  
55 et al., 2001). Blue starters terminate at lower altitudes, typically around 18 km but sometimes up  
56 to 25 km and advance with velocities 10 km/s to 100 km/s (Wescott et al., 1996); they are often  
57 described as initial phases of blue jets (Pasko, 2008) and hence their name. Other types of BLUES  
58 have also been reported, including small blue surface discharges (sometimes called *blue glimpses*)  
59 that appear to “dance” on the upper layer of the cloud at a rate of about 120 per minute (Chanrion  
60 et al., 2017) and the *gnomes* that emerge directly from the cloud top within  $\sim 1$  km, similar to blue  
61 starters, but with brighter and more uniform optical emission and much more compact shape (Lyons  
62 et al., 2003). *Giant jets*, which travel from cloud tops to the lower ionosphere, may also be classified  
63 as BLUES although they involve measurable emissions in the 777.4-nm band (van der Velde et al.,  
64 2019).

65 Each of these types of phenomena exhibits a different morphology but they likely share com-  
66 mon physical processes. The blue color indicates the presence of electron-impact excitation of  
67 molecular nitrogen (Pasko, 2008; Surkov & Hayakawa, 2020; F. Gordillo-Vázquez & Pérez-Invernón,  
68 2021) and the weak or absent atomic oxygen line at 777.4 nm, indicates that air does not reach high  
69 temperatures, typically associated with lightning leaders at ground level. This points to streamer  
70 coronas being the key component of BLUES, a conclusion supported by the close association be-  
71 tween BLUES and Narrow Bipolar Events (NBEs) (F. Liu et al., 2018; Soler et al., 2020), which  
72 are radio emissions also attributed to corona discharges in thunderclouds (Rison et al., 2016; Tilles  
73 et al., 2019; N. Liu et al., 2019). It is thus likely that both BLUES and NBEs are electromagnetic  
74 manifestations of large streamer coronas (or *fast breakdown*, a term coined by Rison et al. (2016)  
75 that we also adopt here).

76 Some distinctive features of each type of BLUE arise from their extension and their location  
77 inside the thundercloud. For example, Soler et al. (2020) analyzed a set of 10 BLUES associated with  
78 positive NBEs and at a considerable depth inside the cloud, presumably between the main negative  
79 and the upper positive charge region of the cloud. As these events are deeply buried in the cloud,  
80 the scattering by cloud droplets and ice crystals blurs their image as observed from above, resulting  
81 in a diffuse blob that can be identified with the blue glimpses reported by Chanrion et al. (2017).

82 Here we focus on events that are close to the cloud top, perhaps partially outside the cloud.  
83 This location suggest an origin between the upper negative region of the cloud and the positively  
84 charged screening layer, and this is supported by radio detections that associate these events with  
85 negative-polarity NBEs (Lyu et al., 2015; Wu et al., 2012; Smith et al., 1999). Because the emissions

86 come from close to the cloud top, optical radiation is less affected by scattering, leading to a more  
87 robust inference of source characteristics. This allows us to compare to radio observations of fast  
88 breakdown.

## 89 **2 Instruments and Observations**

90 The Modular Multispectral Imaging Array (MMIA) is a component of the Atmosphere-Space  
91 Interactions Monitor (ASIM), a mission launched on April 2, 2018 and installed on the International  
92 Space Station (Chanrion et al., 2019; Neubert et al., 2019). MMIA observes in ultraviolet and near-  
93 infrared wavelengths, points towards the nadir and contains three photometers and two cameras.  
94 The three photometers, with a temporal sampling rate of 100 kHz, include one in the UV band at  
95 180 nm to 230 nm, and two others sensitive to the same wavelengths as two installed cameras: in  
96 the near-UV at the strongest spectral line of the nitrogen second positive system (337 nm) and at  
97 the strongest lightning emission line, OI (777.4 nm). The spatial resolution of the cameras is around  
98 400 m  $\times$  400 m at the nadir point and their integration time is 83.3 ms.

99 On the evening of August 7 2019, above an intense localized thunderstorm over Southern  
100 China, there were eight blue luminous events (BLUES) simultaneously detected by MMIA, the  
101 ground-based Vaisala GLD360 global lightning network and the ground-based very low frequency  
102 (VLF)/low frequency (LF) sensor at Guangzhou (see Table 1 for further details). All of them were  
103 detected by MMIA's photometer and camera filtered at 337 nm; some events had a detectable signal  
104 in the 180 nm to 230 nm photometer but there was no signal in the 777.4 nm photometer and camera  
105 at the  $3\sigma$  confidence level. Depending on the event this implies that the 777.4-nm flux was at least  
106 between 50 and 300 times weaker than the 337-nm flux (see Figure S1 in supplemental material  
107 for more details). The rise times of the events in the 337 nm photometer are below 56  $\mu$ s, with the  
108 shortest of them being unresolved by the 10  $\mu$ s sampling time of the photometer. The peak bright-  
109 ness ranges from 20 to 140  $\mu$ W/m<sup>2</sup>, which is among the brightest signals detected by MMIA. The  
110 brightness and quick rise of the events indicate that they originate close to the cloud tops or perhaps  
111 slightly above them. Note however that below we show that most of the emissions were partially  
112 scattered by the cloud and that the photometer light curve is not indicative of the true source duration.

113 We sketch the context of the eight BLUES in Figure 1 which, in panels (a) and (b), plots the  
114 intra-cloud (IC) and cloud-to-ground (CG) flashes and the eight BLUES superimposed on the cloud  
115 top height (CTH) provided by the Fengyun-4A (FY-4A) satellite (Yang et al., 2017) for the time  
116 period from 13:04:00 to 13:07:00 UTC. During these three minutes, there were 522 lightning events  
117 with 240 CG and 282 IC flashes reported by GLD360 (see Figure 1). Two of the BLUE events (with  
118 ID 5 and 7) were missing from GLD360 so for all the BLUES we use the location provided by the  
119 lightning location systems (LLSs) in Guangzhou province (Chen et al., 2012).

120 The absolute timing uncertainty of MMIA is on the order of tens of milliseconds but we can  
121 correct the MMIA times to sub-millisecond accuracy by comparing flash times provided by GLD360  
122 to MMIA 777.4 nm-pulses. In our case we found that the systematic time shift with respect to the  
123 ground-based measurements experienced a time adjustment at around 13:06:07, the time corrections  
124 before and after the time adjustment are  $(-23.3 \pm 0.3)$  ms and  $(-6.2 \pm 0.5)$  ms, respectively (see  
125 Figure S2 in Supplemental material for further details). Note that the time shift  $-23.3$  ms is similar  
126 to the estimations for other thunderstorms such as the  $-28.7$  ms inferred by Soler et al. (2020) or the  
127  $-16.37$  ms from Neubert et al. (2021).

128 With this time correction we find that each of the eight BLUEs has a radio signal that, when  
129 back-propagated to the source, is within 0.7 ms of the optical peak. All VLF/LF waveforms of the  
130 BLUEs were unambiguously classified as negative NBEs (see Figure 2 for examples) measured by  
131 the vertical electric field antenna (frequency bandwidth 800 Hz to 400 kHz) located about 105 km  
132 away at Guangzhou station of Jianghuai Area Sferic Array (JASA) (Qin et al., 2015; F. Liu et al.,  
133 2018). The radio signals of the eight events, along with other positive NBEs at deeper locations in  
134 the same thunderstorm, are analyzed with more details in a complementary publication (F. Liu et al.,  
135 2021).

136 Figure 1 shows in panels (c) and (d) a composition of all camera images for the BLUE events  
137 (always from the 337 nm-filtered camera). To produce this picture we have added the projection  
138 of each of the eight MMIA images into the Earth surface according to coordinates introduced by  
139 the ASIM pipeline. The resulting locations differ noticeably from those provided by LLSs and the  
140 distribution is more spread out. We attribute this to uncertainties in the camera orientation. Note  
141 also that several of the images exhibit a sharp peak that appears to emerge from the middle of the  
142 diffuse blob: this is a blooming artifact of the CCD camera.

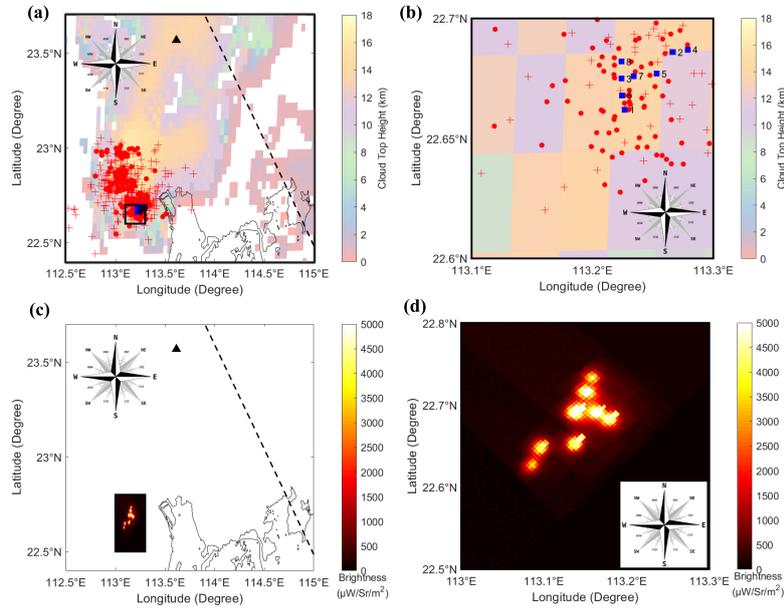
143 To understand better the relation between the BLUE emissions and their parent thunderstorm,  
144 we examined the progression of the cloud Top Blackbody Brightness temperature (TBB in  $K$ ) pro-  
145 vided by the Himawari-8 satellite (Bessho et al., 2016) with ten-minute resolution. Figure S3 in  
146 Supplemental material displays the TBB around the time of our detections. The BLUE events origi-  
147 nated from the boundary of a rapidly-evolving thunderstorm cell. This suggests that rapid turbulent  
148 mixing of the screening layer plays a role in the inception of fast breakdown (Krehbiel et al., 2008;  
149 Lyons et al., 2003) or the occurrence of groups of localized NBEs is associated with dynamically  
150 intense convection (Bandara et al., 2021). Note that the altitudes provided by FY-4A that we use  
151 here are likely underestimates since negative NBEs are usually associated with deep convection and  
152 detected in overshooting cloud tops (F. Liu et al., 2018; Wu et al., 2013).

**Table 1.** The eight BLUES simultaneously detected by MMIA, ground-based Vaisala GLD360 global lightning network and the ground-based VLF/LF sensor at Guangzhou. All the detection times have been corrected to the time with respect to the BLUES source locations.

ID	Date (Year/Month/Day)	MMIA time UTC(Source)	MMIA corrected time UTC(Source)	GLD360 time UTC(Source)	VLF/LF time UTC(Source)	Rise Time <sup>a</sup> ( $\mu$ s)	Time duration <sup>b</sup> ( $\mu$ s)	Peak Brightness ( $\mu$ W/m <sup>2</sup> )
1	2019/08/07	13:05:56.9362	13:05:56.9595	13:05:56.9594	13:05:56.9594	31.17	371.36	20.11
2	2019/08/07	13:05:58.6317	13:05:58.6550	13:05:58.6549	13:05:58.6549	8.04	196.72	142.18
3	2019/08/07	13:06:01.7568	13:06:01.7801	13:06:01.7800	13:06:01.7799	56.30	758.77	40.61
4	2019/08/07	13:06:09.5668	13:06:09.5730	13:06:09.5722	13:06:09.5723	8.71	399.46	97.88
5	2019/08/07	13:06:16.6329	13:06:16.6391	-	13:06:16.6384	13.07	912.33	44.51
6	2019/08/07	13:06:20.9670	13:06:20.9732	13:06:20.9726	13:06:20.9726	14.08	486.78	120.30
7	2019/08/07	13:06:30.4934	13:06:30.4996	-	13:06:30.4993	13.07	334.22	46.48
8	2019/08/07	13:06:31.6557	13:06:31.6619	13:06:31.6616	13:06:31.6615	13.40	237.87	39.96

<sup>a</sup> Rise time is calculated using the linear interpolation by taking the time for the amplitude of a photometer signal to rise from 10% to 90%. Note that the sampling time is 10  $\mu$ s so the rise is unresolved in several events.

<sup>b</sup> Time duration is calculated using the linear interpolation by the time interval for the amplitude of a photometer signal to rise from 10% and fall to 10%.



**Figure 1.** Distribution of the CG/IC lightning and eight BLUES along with Cloud Top Height (CTH) at the time period from 13:04:00 to 13:07:00 UTC (a) and the zoom of its black rectangular region (b) (Red dots: CG lightning detected by GLD360, Red crosses: IC lightning detected by GLD360 and Blue squares: BLUES detected by LLSs); eight BLUES images detected in the 337 nm filtered camera of MMIA (c) and the zoom (d). The ground-based VLF/LF sensor at Guangzhou is shown as black triangle. The footprints of ASIM are shown in black dashed line.

### 3 Light-scattering model

To better understand the MMIA observations we compare them now to a simple model where the light source is a thin, straight, uniformly bright segment and the cloud is homogeneous with a planar upper boundary. We neglect the intrinsic duration of the source, assuming that all light is emitted instantaneously.

Because photons can be scattered many times before they exit the cloud, an impulsive optical flash results in a temporally stretched light curve. To understand this curve we start with the expression for a point-like source buried in the cloud. Using the diffusion approximation for the propagation of photons inside the cloud proposed by Koshak et al. (1994), Soler et al. (2020) gave an analytical expression for this curve, which was derived in more detail by Luque et al. (2020). Adopting the normalization and the notation of the later, the photon flux exiting the cloud top reads

$$\Gamma(t) = NF(t) = \frac{Ne^{-t/\tau_A - \tau_D/t}}{\pi^{1/2}\tau_D} \left(\frac{t}{\tau_D}\right)^{-3/2}, \quad (1)$$

where  $F(t)$  is the flux per photon in the source,  $N$  is the total number of source photons,  $\tau_A$  is the mean absorption time of the photons inside the cloud and  $\tau_D = L^2/4D$  is, given a diffusion coefficient  $D$ , the characteristic time of diffusion for the distance  $L$  between the source and the cloud top. The derivation of these magnitudes from the microscopic properties of the cloud is given by Koshak et al. (1994) and reviewed by Luque et al. (2020). Equation (1) is valid for  $t > 0$ , where the time origin is the moment of light emission. For a distant observer, differences in light travel time from different points in the cloud are not significant so one can reinterpret the time in equation (1) with  $t = 0$  being the arrival time of an unscattered photon.

To obtain the light curve for an extended source that spans altitudes from the cloud top to a maximum depth  $L_0$  we integrate  $(N/L_0)F(t)dL$  from 0 to  $L_0$  (the factor  $N/L_0$  is the linear density of source photons, assumed uniform). The result is

$$\Gamma_L(t) = N\left(\frac{D}{t}\right)^{1/2} \left(1 - e^{-\tau_D/t}\right) e^{-t/\tau_A}, \quad \tau_D = L_0^2/4D. \quad (2)$$

Note that this expression disregards any part of the source above the cloud top. Some photons emitted outside the cloud propagate directly to the detector and others are back-scattered by the upper cloud surface after a small number of scattering events. These emissions have an effect only on a few data points in a photometer with a 10  $\mu$ s time resolution. We therefore do not account for these emissions which, although may be present, do not dominate the photometer light-curves.

For short times after the emission, equation (2) predicts a  $\sim t^{-1/2}$  dependence for the flux that is cut off by either photon absorption with a time-scale  $\tau_A$  or by the finite size of the source, with a time-scale  $\tau_D$  (see Figure S4 in supplemental material). The cutoff from the mean absorption time is likely negligible: Luque et al. (2020) estimates  $\tau_A \approx 2.3$  ms for clouds composed by water droplets

175 with an effective radius of  $20\ \mu\text{m}$  and a droplet density of  $10^8\ \text{m}^{-3}$  but this possibly overestimates the  
 176 absorption. The cloud tops are dominated by ice particles, which absorb radiation at  $337\ \text{nm}$  several  
 177 orders of magnitude less efficiently than water (Warren & Brandt, 2008). Besides, the available  
 178 estimates of the extinction coefficient (Peterson, 2020; Platt, 1997) also lead to absorption times  
 179 significantly longer than the duration of our events. Hence here we assume  $\tau_A \gg \tau_D$ .

180 As we show in Figure S4 in the supplemental material, most of the recorded BLUEs light-  
 181 curves have the shape predicted by equation (2). In the figure we plot a least-squares fit of the  
 182 observational data to the model with two parameters: an overall amplitude factor and the decay time  
 183  $\tau_D$ . To reduce the effect of the emissions from outside the cloud discussed above, we disregard the  
 184 data points at the peak of the light-curve. The good fit of most events indicate that indeed they  
 185 originate from sources that extend below the cloud top. Event 5 is the only one that does not show a  
 186 clear  $t^{-1/2}$  decay, possibly because there was a gap between the source and the cloud top or because  
 187 light emissions were inhomogeneous or long-lasting. In events 1 and 7 there is weak secondary  
 188 activity 1 ms to 2 ms after the main peak that distorts the estimate of the cutoff time  $\tau_D$ .

189 Leaving aside events where  $\tau_D$  was estimated poorly, this cutoff time ranges between 0.5 ms  
 190 and 1.6 ms. The smallest diffusion coefficient proposed by Soler et al. (2020),  $D = 3 \times 10^9\ \text{m}^2\text{s}^{-1}$   
 191 yields a range of lengths for the optical sources of  $L_0 = 2.4\ \text{km}$  to  $4.4\ \text{km}$ . However, the evaluated  
 192 results will be affected by the uncertainties that surround our modeling of the cloud composition.

193 Next we extend our model to include the propagation of the signal to the MMIA instruments,  
 194 accounting for Rayleigh scattering by the atmosphere and for the non-isotropic (approximately Lam-  
 195 bertian) emission pattern from the cloud tops. We use the radiative transfer Monte Carlo code Cloud-  
 196 Scat.jl (Luque et al., 2020) and run simulations of uniformly bright, straight vertical sources, with  
 197 the lengths  $L_0$  derived above, in a homogeneous cloud that spans altitudes from 7 km to the cloud top  
 198 height derived by the Fengyun-4A (FY-4A) satellite (listed in Table 2). The scattering parameters in  
 199 the cloud are those for a density of  $10^8\ \text{m}^{-3}$  spherical ice particles with  $20\ \mu\text{m}$  radius. The relative  
 200 positions between the source and the observer reproduce the conditions of each of the eight BLUEs  
 201 in our dataset.

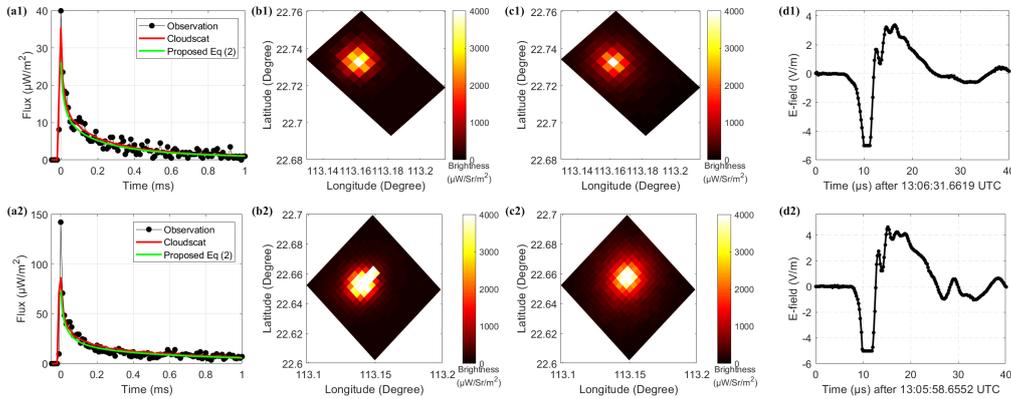
202 In Figure 2 we show the results of the Monte Carlo code comparing with 337-nm photometer  
 203 and camera observations for the event 2 and 8 (Additional comparisons can be found in Figure  
 204 S5-S12 in supplemental material). The photometer light curves calculated from CloudScat model  
 205 closely follow the analytical estimate of equation (2) and are a good fit to the observations. The  
 206 simulated camera images are also reasonably close to MMIA's records although they are slightly  
 207 more compact. This is a possible indication of a non-negligible source width on the order of the  
 208 camera resolution of about 400 m.

**Table 2.** Model-inferred properties of the eight BLUE events. We list the cloud top height measured by FY-4A at the event location, the best-fit cutoff time  $\tau_D$  (see equation (2)), the resulting source length  $L_0 = (4D\tau_D)^{1/2}$  with  $D = 3 \times 10^9 \text{ m}^2\text{s}^{-1}$ , the total optical energy in the 337-nm band of the second positive system of nitrogen and an estimation of the number of streamer branching events in the fast breakdown processes that we assume that originated the events.

ID	Cloud top height (km)	Cutoff time ( $\tau_D$ ) (ms)	Source length ( $L_0$ ) (km)	Optical energy at 337 nm (J)	Branching events
1	11.2	19.6 <sup>a</sup>	- <sup>a</sup>	$1.8 \times 10^4$	$1.4 \times 10^9$
2	11.6	1.6	4.4	$1.9 \times 10^4$	$1.4 \times 10^9$
3	14.0	1.2	3.8	$1.2 \times 10^4$	$9.1 \times 10^8$
4	13.1	0.5	2.4	$1.3 \times 10^4$	$9.8 \times 10^8$
5	11.6	8.6 <sup>b</sup>	- <sup>b</sup>	$2.4 \times 10^4$	$1.9 \times 10^9$
6	14.0	0.8	3.1	$1.8 \times 10^4$	$1.4 \times 10^9$
7	14.0	3.8 <sup>a</sup>	- <sup>a</sup>	$1.6 \times 10^4$	$1.2 \times 10^9$
8	14.0	0.5	2.5	$4.6 \times 10^3$	$3.6 \times 10^8$

<sup>a</sup> In events 1 and 7 there is secondary activity that distorted the estimation of the cutoff time  $\tau_D$  and the source length.

<sup>b</sup> Event 5 has a light-curve that cannot be explained by an impulsive, uniformly bright source.



**Figure 2.** Comparing the modeling results against the observation of MMIA for event 8 (a1 - d1) and event 2 (a2 - d2). The 337 nm photometer signals (a1, a2): MMIA observation (black dotted line), modeling results by using CloudScat model (red line) and the proposed equation (2) (green line). The images measured by 337 nm filtered camera of MMIA (b1, b2). The simulated images obtained from CloudScat model (c1, c2). The waveforms of the NBEs detected from the ground-based VLF/LF sensor at Guangzhou (d1, d2).

209

In the results presented here we always consider that the top of the source coincides with the top

210

of the cloud. As we discuss above, the effect of light emissions outside the cloud is too impulsive to

211 compare against the MMIA photometer and is possibly dominated by the intrinsic time-dependence  
 212 of the source. We performed additional Monte Carlo simulations that confirm that the photometer  
 213 light-curves are compatible with source tops within a few hundred meters of the cloud top, either  
 214 above or below it. The VLF/LF waveforms of negative NBEs for event 2 and 8 are also shown in  
 215 Figure 2.

216 The CloudScat.jl code outputs a photon flux at the observer’s location in units of photons per  
 217 unit time and unit surface that reach a detector for each photon in the source whereas the MMIA  
 218 photometers are calibrated in terms of power per unit surface (irradiance). The conversion factor is  
 219 the total energy of the event in the 337 band,  $E = Nhc/\lambda$ , where  $N$  is the total number of photons  
 220 emitted by the source,  $h$  is the Planck’s constant,  $c$  is the speed of light and  $\lambda = 337$  nm. By  
 221 comparing the results of our Monte Carlo code to the MMIA data we found the best-matching total  
 222 energy of each event. Because the events that we analyze are close to the cloud tops and thus barely  
 223 affected by in-cloud absorption, our estimates of  $E$  are weakly sensitive to our model assumptions  
 224 and thus provide a reasonably precise picture of the actual source emission intensity of the BLUE  
 225 events. The estimated energies are listed in Table 2.

226 N. Liu et al. (2019) analyzed radio spectra of NBEs and concluded that they can be understood  
 227 as systems of  $10^7$  to  $10^8$  streamers. In that analysis the key feature of a streamer is a current moment  
 228 that increases rapidly on a time scale of about one nanosecond, which is the timescale of streamer  
 229 initiation in numerical simulations. Here we also consider that a single nanosecond event may  
 230 produce more than one streamer, as is the case in a bifurcated tree. Denoting by  $b$  the mean number  
 231 of streamers emerging from an event, we have  $M = bK$ , where  $M$  is the total number of streamers  
 232 (unbifurcated branches) and  $K$  is the number of initiation events (most likely bifurcations from other  
 233 streamers). Then the total streamer length contained in one fast breakdown process is  $\ell bK$ , where  
 234  $\ell$  is the mean length between bifurcations (but see Nijdam et al. (2020) for a discussion of the  
 235 difficulties involved in precisely defining this quantity). If a streamer emits  $\eta$  photons per unit length  
 236 as it propagates, the total number of emitted photons is  $N = \eta\ell bK$ .

237 It follows from the scaling laws for streamers (Ebert et al., 2010) that  $\eta$  depends on air density  
 238 only through collisional quenching of the radiative states of  $N_2$  (the remaining factor is proportional  
 239 to  $\pi R^2 n_e$  where  $R$  is the streamer radius, which scales as the inverse of air density, and  $n_e$  is the  
 240 electron density, which scales as the square of air density). The numerical simulations by Malagón-  
 241 Romero & Luque (2019) predict a time-integrated photon yield of about  $2 \times 10^{18} \text{ m}^{-3}$  in a streamer  
 242 of radius  $R \approx 2$  mm. From the spectra presented by F. J. Gordillo-Vázquez et al. (2012), we infer  
 243 that about 30% of these emissions are inside the 5-nm window of the 337-nm MMIA photometer,  
 244 yielding  $\eta \approx (n_0/n) \times 2 \times 10^{12} \text{ m}^{-1}$ , where  $n_0/n$  is the ratio of air density at atmospheric-pressure,  
 245  $n_0$ , to the density at the altitude of interest,  $n$ . Here we take  $n_0/n \approx 6$  for an altitude of 14 km. Briels  
 246 et al. (2008) observed a ratio of branching length to streamer radius of about 20, so a radius of 2 mm

247 at atmospheric pressure translates into  $\ell \approx (n_0/n) \times 4 \text{ cm} = 24 \text{ cm}$ . We take the branching number  $b$   
 248 to be 2, although there are evidences that it may possibly be larger (Heijmans et al., 2013, 2015).

249 The estimation of the number of streamer branches for all the BLUEs are listed in Table 2.  
 250 Our results are 10-100 times above those derived from radio spectra by N. Liu et al. (2019). One  
 251 possible reason for this disagreement is the uncertainties in our assumed parameters. For example,  
 252 the estimated  $K$  is highly sensitive to the assumed streamer radius: had we chosen a radius of 5 mm  
 253 at atmospheric pressure, the estimation of  $K$  would be reduced by about a factor 15. It is also  
 254 possible that a large fraction of the optical signal in fast breakdown is emitted not close to streamer  
 255 heads but from long-lived glows, as is the case in sprites (Luque et al., 2016; Pérez-Invernón et al.,  
 256 2020).

#### 257 **4 Discussion and conclusions**

258 The eight BLUE events that we analyze in this paper expand and complete the picture of fast  
 259 breakdown as the source of both optical blue-dominated emissions and radio pulses detected as  
 260 NBEs in the VLF/LF bands or high-amplitude noise in VHF. All events were strongly detected in  
 261 the photometer and camera filtered at 337 nm; in some events there was a weak signal in 180-230 nm  
 262 but with no signal in 777.4 nm photometer and camera.

263 As in previous studies (Wescott et al., 1995, 1996; Chanrion et al., 2017), the BLUEs appeared  
 264 temporally isolated from either CG or IC flashes detected by the GLD360 network. However, all  
 265 the BLUEs coincide with NBEs observed by the ground-based VLF/LF sensor at Guangzhou. This  
 266 strengthens the connection between BLUEs and negative NBEs (Chou et al., 2018; F. Liu et al.,  
 267 2018) and further supports that NBEs originate from non-thermal, streamer processes (Rison et al.,  
 268 2016; Lyu et al., 2019; Tilles et al., 2019; Soler et al., 2020).

269 The rise times of the blue events in the 337 nm photometer are between  $10 \mu\text{s}$  to  $70 \mu\text{s}$  with  
 270 peak irradiance varying from  $20 \mu\text{Wm}^{-2}$  to  $140 \mu\text{Wm}^{-2}$ . The brightness and short rise times suggest  
 271 a source close or even slightly above the cloud tops and this is supported by our modeling results  
 272 based either on the diffusion approximation by Koshak et al. (1994) or on a Monte Carlo radiative  
 273 transfer code (Luque et al., 2020). Since all events are identified as negative NBEs, this is consistent  
 274 with previous studies that localize the initiation of most of negative NBEs between the upper positive  
 275 charge region and the screening charge region of the thunderstorm (Smith et al., 1999; Wu et al.,  
 276 2012; Lyu et al., 2015). The variation in the rise times between different events may be due to  
 277 differences in the intrinsic time dependence of the optical sources but this is equally well explained  
 278 by a finite distance to the cloud top or from non-uniformities of the optical sources below the cloud.

279 Our estimates of the total optical energy within the 337-nm band provide a new constraint for  
 280 models of fast breakdown. The present understanding of these events is still limited and it is difficult

281 to translate this energy into microscopical properties of fast breakdown. However our results confirm  
282 that fast breakdown involves more than  $10^7$  streamers, as inferred by N. Liu et al. (2019) and further  
283 analyzed by Cooray et al. (2020).

284 Future investigations should address the underlying physics of fast breakdown and its global  
285 significance, including its relation to lightning initiation. Data from the ASIM mission will likely  
286 play a decisive role in this research.

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