

The Combined Influence of Lower Band Chorus and ULF waves on Radiation Belt Electron Fluxes at Individual L-shells

Laura E. Simms¹, Mark J. Engebretson¹, Craig J. Rodger², Stavros Dimitrakoudis³, I.R. Mann³, and Peter J. Chi⁴

¹ Augsburg University, Minneapolis, MN, USA

² University of Otago, Dunedin, New Zealand

³ University of Alberta, Edmonton, Canada

⁴ UCLA, Los Angeles, CA, USA

Corresponding author: Laura E. Simms (simmsl@augsborg.edu)

Key Points:

- VLF and ULF waves are associated with same day electron (>1.5 MeV) decreases but with flux enhancement 1-2 days later.
- ULF is nearly always more influential than chorus, but electron enhancement is most effective when both wave types are present.
- Both VLF and ULF are more influential at L 4-6 than at lower altitudes.

Abstract

We investigate the timing and relative influence of VLF in the chorus frequency range observed by the DEMETER spacecraft and ULF wave activity from ground stations on daily changes in electron flux (0.23 to over 2.9 MeV) observed by the HEO-3 spacecraft. At each L shell, we use multiple regression to investigate the effects of each wave type and each daily lag independent of the others. We find that reduction and enhancement of electrons occur at different time scales. Chorus power spectral density and ULF wave power are associated with immediate electron decreases on the same day but with flux enhancement 1-2 days later. ULF is nearly always more influential than chorus on both increases and decreases of flux, although chorus is often a significant factor. There was virtually no difference in correlations of ULF Pc3, Pc4, or Pc5 with electron flux. A synergistic interaction between chorus and ULF waves means that enhancement is most effective when both waves are present, pointing to a two-step process where local acceleration by chorus waves first energizes electrons which are then brought to even higher energies by inward radial diffusion due to ULF waves. However, decreases in flux due to these waves act additively. Chorus and ULF waves combined are most effective at describing changes in electron flux at >1.5 MeV. At lower L (2-3), correlations between ULF and VLF (likely hiss) with electron flux were low. The most successful models, over $L=4-6$, explained up to 47.1% of the variation in the data.

1. Introduction

Both chorus and ULF Pc5 waves are thought to influence electron flux levels, increasing them in some situations and decreasing them in others. Observations of single storms show that interactions with chorus can scatter electrons into the loss cone (Clilverd et al., 2016; Horne and Thorne, 2003; Shprits et al., 2007, 2008; Thorne et al., 2005) while ULF Pc5 waves can lead to loss through a combination of outward radial diffusion and magnetopause shadowing (Katsavrias et al., 2015; Kellerman and Shprits, 2012; Loto'aniu et al., 2010; Mann et al., 2012; Ozeke et al., 2020; Turner et al., 2012). However, chorus may also result in flux enhancement through local acceleration (Horne et al., 2005; Katsavrias et al., 2015; Reeves, 2013; Shprits et al., 2008; Summers et al., 1998), while ULF Pc5 waves accelerate electrons through inward radial diffusion (e.g., Hao et al., 2019; Katsavrias et al., 2019; Mann et al., 2004) and direct interaction (Claudepierre et al., 2013; Hao et al., 2014; Zong et al., 2009; Zong et al., 2017).

Both wave types often occur simultaneously before or during changes in electron flux levels (Li et al., 2005). It requires further investigation to determine whether they contribute equally to electron flux levels. Katsavrias et al. (2015) found that over $L=3-5$, the electron population above 300 MeV/G increased when both chorus (inferred from the ratio of precipitating to trapped electron fluxes using POES data in low Earth orbit) and ULF wave power were high, but depleted when only ULF wave activity was high. They interpreted this to mean that chorus was the dominant driver of electron acceleration, but ULF waves, via outward diffusion, resulted in depletion. In a different pair of storms, inferred chorus appeared to accelerate electrons to relativistic energies while ULF waves further accelerated this

population to ultrarelativistic energies through inward radial diffusion (Katsavrias et al., 2019). Using results from superposed epoch analysis, flux enhancements were associated with VLF activity (inferred from microbursts) around $L=4.5$, but at L shells above that, the flux association was stronger with ULF activity (O'Brien et al., 2003). Analyzing effects simultaneously using multiple regression, ULF wave power at geosynchronous orbit was found to be of somewhat more influence on electron flux levels (>1.5 MeV) than satellite-observed chorus power spectral density (Simms et al., 2018ab).

There is, however, still a need to investigate these simultaneous wave effects by L shell, over multiple days, using analyses which study the effect of each factor independent of the others. We use partial correlation analysis to study the influence of each wave type over several days. From these analyses, we are able to determine both when the waves are most influential, and to separate the positive and negative effects that occur on different days. Using multiple regression, we study the separate effects of each wave type on each day. We also explore the nonlinear action of waves on electron flux by including quadratic terms, and whether ULF and VLF waves act both additively and synergistically by including a multiplicative term (Neter et al., 1985).

In this study, we use electron flux data (HEO-3 satellite) gathered at each L shell ($L=1.75-9$). We use satellite-observed VLF power spectral density (PSD) in the lower band chorus range from the DEMETER satellite, binned by L -shell over $L=1.75-6$. These waves are assumed to be predominately chorus above $L=3$ and predominately hiss below this (Bortnik et al., 2008; Carpenter and Park, 1973). This differs from previous studies of VLF correlation by L shell with electron flux as the VLF measure is neither a proxy from microburst observation nor ground-based VLF data which does not as accurately reflect VLF wave activity occurring in space (Simms et al., 2019). We compare the effects of ULF Pc3, Pc4, and Pc5 data from ground-based stations, Pc5 from $L>3$.

2. Data

Data covered the time period 11 Aug 2004-27 July 2007. Electron flux data were obtained from the HEO-3 satellite which lies in a roughly 12 hour highly elliptical orbit with good data coverage in the L range from 2 to 9 (Fennell et al., 2004; Fennell and Roeder, 2008). The orbit is such that it cycles through all MLT. We use the E2 telescope channel (>0.23 MeV) and the omnidirectional sensors E4-E6 channels (>0.6 , >1.5 , and >2.9 MeV) (all sampling number of electrons/second) with observations binned by L shell (IGRF) with, for example, $L=2$ including $L=2.0-2.99$. The E2 telescope temperature rose with altitude at $L \geq 5$ which increased noise levels. This may mean that results in the >0.23 MeV channel at higher L are less reliable. We use only data collected over the northern hemisphere (high altitude) where dwell time at each L shell is longer. This reduced variability in the data as pitch angles differ between northern and southern hemispheres. We did, however, retain both "even" and "odd" orbits, which sample different MLT, to obtain a more representative average of the electron population. Note that L is correlated with latitude, with lower L shells ($L \leq 3$) sampled from near the equator. As high energy protons can contaminate the electron channels, we removed days on which solar proton events were occurring, as well as one day following (<ftp://ftp.swpc.noaa.gov/pub/indices/SPE.txt>). However, there may also be proton contamination outside solar proton events. To correct for this, we removed

data falling far above the cloud of data points when proton vs. electron counts are plotted (O'Brien, 2012). This removed more lower L shell observations ($L < 4$). For this reason, results from lower L shells are not as robust.

We obtained VLF power spectral density ($\log_{10} [\mu V^2 / m^2 / Hz]$) from dayside, northern hemisphere passes (LT 10:30) from the Instrument Champ Electrique (ICE) on the DEMETER satellite (Berthelier et al., 2006). There was good data coverage over $L=1-6$ (IGRF). We limit the VLF data to the lower band chorus range (0.1-0.5 fce). We expect this dataset is likely dominated by lower band chorus activity, but acknowledge that other VLF wave activity could be present. Therefore, more generally, this dataset contains whistler mode waves. In particular, this VLF wave band at lower L shells (within the plasmasphere) is dominated by hiss (Bortnik et al., 2008). Chorus waves predominate above the plasmopause ($L \sim 3$) and hiss within the plasmasphere below this (Carpenter and Park, 1973). We use only the dayside pass of the satellite because this range of VLF is found over a broader range of latitudes on dayside rather than nightside and is not as influenced by geomagnetic activity (Tsurutani & Smith, 1977; Li et al., 2009, Thorne et al., 2010). This may represent only a sample of overall global activity as satellites can only sample one small area of the magnetosphere at a time. In particular, it may result in a certain L shell being more heavily sampled at a particular local time.

ULF Pc3, Pc4, and Pc5 wave power (nT^2/Hz) was obtained from individual McMAC ground stations at $L=2.5$ (Bennington or BENN) and $L=3.4$ (Glyndon or GLYN) and CARISMA ground stations at $L=4.06$ (Pinawa or PINA), $L=5.15$ (Island Lake or ISLL), $L=6.15$ (Gillam or GILL), and $L=7.44$ (Fort Churchill or FCHU). Wave power was calculated using a Fourier transform with a 1 h window. This allows a good discernment of waves down to the minimum (Pc5) frequency of 1.67 mHz. There were strong correlations between the 3 ULF wave bands at all the ground stations, and each band correlated extremely similarly with electron flux. In order to compare most easily to the ULF index, we use Pc5 in most of our analyses, however, given the high correlations between bands, each of Pc3, Pc4, and Pc5 are a nearly identical proxy for the other two. Daily averaged ULF Pc5 wave power was obtained from a ground-based ULF Pc5 index covering local times 0500 – 1500 in the Pc5 range (2-7 mHz) obtained from magnetometers stationed at 60-70° N CGM (Corrected GeoMagnetic) latitude (Kozyreva et al., 2007).

The log values of lower band chorus PSD, ULF power, and flux data were all daily averaged. Predictor variables (VLF and ULF wave activity) were lagged at 0 (same day; Lag 0), 1 (previous day; Lag 1), 2, and 3 days as most of the correlation between these waves and electron flux occurs within this time frame (Mann et al., 2004; Simms et al., 2018a). For the correlation and regression analyses, daily change in electron flux was calculated by subtracting the daily average of the previous day from the current day's average. All single correlations reported are standard Pearson correlation coefficients. Partial correlation between two factors fixes other variables at a given level to control for their effects (Neter et al., 1985).

The flux observations show serial autocorrelation ($p < 0.0001$ using Durbin-Watson test, see Neter et al., 1985), as each day is correlated with the previous day. To correct for this, we include previous day's flux as a predictor in the regression models. The result of this is that the regressions are essentially AR1 (autoregressive at one time step, represented by the previous day's flux term) differenced (as the observations are daily change) models (Hyndman and Athanasopoulos, 2018; Simms et al., 2019). The addition of the AR1 term reduces the autocorrelation enough that we can have confidence in the p-values of the statistical tests. Nonlinear effects of waves on electron flux are explored by including quadratic terms in the regression analyses, while the synergistic combined action of ULF and VLF waves is tested by including a multiplicative term (Neter et al., 1985).

Statistical analyses were performed in MATLAB.

3. Results

Daily average electron fluxes peak at $L=4$, with lower values observed in the slot region ($L=2$) and beyond geosynchronous orbit ($L>6$) (Figure 1). In contrast, \log_{10} lower band chorus PSD (from DEMETER satellite) increases steadily up to $L=6$ (Figure 2).

The power of all three ULF bands from the ground stations (Pc3, Pc4, and Pc5) rise over $L=2-6$, with very low power seen at $L2$ (Figure 3). ULF Pc4 has the highest power, but there are such high correlations between these three bands (0.9432 - 0.9989) that any band is an almost exact proxy for the other two.

3.1 Correlations between Chorus PSD and Electron Flux

Neither hiss nor chorus PSD on the same day (Lag 0) is strongly associated with electron flux difference at >1.5 MeV with correlations ranging from -0.02 to 0.07 (Table 1). When chorus is measured the day before (Lag 1), there are stronger correlations at $L>3$, but all are below 0.40 (Figure 4; Table 1).

Table 1. Correlation of same day (Lag 0) and previous day (Lag 1) hiss ($L<4$) and chorus ($L>3$) PSD with >1.5 MeV electron flux daily change.

L shell	Same day (Lag 0)	Day previous (Lag 1)
$L=2-2.99$	0.04	0.03
$L=3-3.99$	0.07	0.15
$L=4-4.99$	-0.02	0.37
$L=5-5.99$	-0.02	0.35
$L=6-6.99$	-0.05	0.30

Both VLF in the chorus band and flux differences show little spread below $L=3$. Values are tightly clustered at low VLF PSD (not much VLF is seen at these L shells) and around 0 change in electrons. This may reduce the likelihood of seeing much influence of the VLF whistler mode waves on changes in electron flux at $L=2-2.99$.

We perform partial correlations between VLF waves at each L shell (2-6) over 0-3 days with daily flux difference at the four energy levels (Figure 5, $r < 0.10$ lie within the gray band). At $L=2-2.99$ (hiss), nearly all correlation magnitudes were less than 0.1. Over $L=4-5$, chorus measured on the same day is associated with a drop in flux at the 3 higher energies. Chorus measured one day before (Lag 1) is most associated with an increase in flux at all 4 energies over $L=4-6$. However, none of the correlations exceed 0.4. These low correlations suggest that the chorus is not the only driver of electron changes. However, we note that the partial correlation analysis differentiates the Lag 0 effect more clearly than the simple correlation analysis of Table 1. Over $L=4-5$, at >1.5 MeV, electron reductions at Lag 0 due to chorus are stronger ($L4$: $r = -0.32$, $L5$: $r = -0.21$ in the partial correlations). The Lag 1 partial correlation with flux enhancement is more similar ($L4$: $r = 0.38$, $L5$: $r = 0.35$ in the >1.5 MeV partial correlations) to that found with simple correlation.

3.2 Correlations between ULF Power and Electron Flux

Despite the obvious differences in power, extremely high correlations (0.9432 - 0.9989) between the three ULF bands (Pc3, Pc4, and Pc5) suggest that any one could serve as an excellent proxy for the others. Simple correlations between each ULF band (Lag 1) and >1.5 MeV electron flux at L=5 are identical (Figure 6). Partial correlations over 0-3 days previous between ULF in each band with the daily differences in the 4 electron flux energies at L = 5 show this to be the case (Figure 7). The pattern of correlation was virtually the same between each ULF band and electron flux. These correlations were similarly indistinguishable at the other L shells. We continue our analysis using ULF Pc5 data.

ULF power shows a similar pattern to that of the effect of chorus (or hiss at the lower L shells) on >1.5 MeV flux: Lag 1 correlations tend to be positive and of greater magnitude than Lag 0 correlations. Both Lag 0 and Lag 1 correlations are greatest in magnitude over L=4-7 (Table 2; Figure 8). ULF Pc5 at L=5 is the most highly correlated ($r=0.34$).

Table 2. Correlation of same day (Lag 0) and previous day (Lag 1) ULF Pc5 power with >1.5 MeV electron flux daily change.

L shell	a. Same day (Lag 0)	b. Day previous (Lag 1)
L=2 (BENN)	0.05	0.06
L=3 (GLYN)	0.06	0.19
L=4 (PINA)	-0.19	0.33
L=5 (ISLL)	-0.19	0.34
L=6 (GILL)	-0.20	0.30
L=7 (FCHU)	-0.20	0.18

Partial correlations of ULF power at each L shell (1-7) with the flux differences over 0-3 day lags show a similar pattern to that of VLF waves (Figure 9). At L=2, all correlations were low (close to or less than $|0.1|$). At L=3, the most important ULF correlations are at Lag 1, but these are all ≤ 0.25 . Over L=4-7, the peak in positive association occurs with previous day's ULF, while the strongest negative correlations are with Lag 0 ULF. Using partial correlation analysis, we can more clearly see the electron decreases at Lag 0 associated with ULF waves than the individual correlations of Table 2 would suggest. The middle range of flux energies (>0.6 MeV and >1.5 MeV) show the highest response (r as low as -0.4 at Lag 0 and as high as 0.4 at Lag 1). ULF power from 2 days previous is less associated with flux changes and that from 3 days prior is below $|0.10|$.

3.3 The ULF Index Is a Reasonable Proxy for ULF Pc5 Power at Each L Shell

If ULF power measurements from each L shell are not available, the ULF index can be a practical replacement. The correlation of the index with ULF Pc5 power at ground stations at each L shell is reasonably high (0.74 – 0.85, Table 3a). The correlation of daily change of electron flux at each L shell with the Lag 1 index is also very similar to that with each ground station, despite not being centered on a particular L shell (compare Table 2b with Table 3b).

Table 3. At each L shell, the correlation of the ULF Pc5 index a. with ULF Pc5 power at individual CARISMA ground stations and b. with next day's >1.5 MeV electron flux daily change.

L shell	a. ULF Index with ground station	b. Electron flux with Lag 1 ULF Index
L=2	0.74	0.05
L=3	0.83	0.19
L=4	0.85	0.38
L=5	0.80	0.33
L=6	0.78	0.24
L=7	0.76	0.13

3.4 Combined Effects of Chorus and ULF Pc5

A key question is whether the two most important factors, chorus (at $L > 3$) and ULF Pc5, each still correlate with flux changes when the other factor is accounted for. To determine this, we perform multiple regressions including both wave types, producing standardized regression coefficients so as to compare relative strengths of predictors on a common scale. As the highest partial correlations were on Lag 0-2 (same day to two days previous) we include all these measures over $L=3-6$, the L shells in which both chorus and ULF power have good data coverage and where the individual correlations with flux were highest. Previous day's flux (Lag 1 daily flux change) is included as a predictor in the model to correct for serial autocorrelation (leading to an essentially AR1 model).

Over $L=4-6$, ULF Pc5 consistently shows a stronger influence than chorus over Lag 0-1, with a negative influence at Lag 0 and positive influence at Lag 1 (Figure 10; red line indicates statistically significant coefficients). It is notable that the influence of ULF Pc5 Lag 2 increases from low to high flux energy levels, while that at Lag 1 decreases. This indicates that it takes up to several days for processes driven by these waves to accelerate electrons up to the higher energies. At the 3 highest energies, ULF Pc5 influence at each time step is strongest at $L=4$ (decreasing through $L=6$) and at >0.6 MeV flux (decreasing up to >2.9 MeV). At $L=3$, Lag 1 ULF is also much more important than chorus. However, there is more effect of chorus at Lag 0 at the higher 2 energies, and the ULF correlation is positive at the 2 lower energies.

This combined variable set (lags 0-2 of both chorus and ULF and previous day's flux) is most strongly associated with flux differences at >1.5 MeV. R^2 (percent of variability in the data explained by the regression model) peaks at 47.1% at $L=5$ in the >1.5 MeV energy range. This corresponds roughly to a correlation of 0.69 (the square root of R^2). However, the variability explained is not much lower at the other 3 energies. The rough correlation estimate is similar at >0.6 and >2.9 MeV ($r = 0.66$ and 0.65), but somewhat lower at >0.23 MeV (0.58). At each L shell and energy, the correlation of flux with chorus and ULF Pc5 combined is higher than in the simple correlations. This indicates that both factors and all lags are needed for a fuller description of flux changes.

3.5 Nonlinear Effects of Chorus and ULF Pc5 on Flux

However, these linear, additive models may not completely capture the combined effects of chorus and ULF waves on electron flux. The effect of each factor may be neither linear nor independent of the other wave type. For each of Lag 0 and 1, we include squared terms to describe nonlinear effects, and a multiplicative interaction term to describe possible synergistic effects between the chorus and ULF wave effects. Regression surfaces show the predicted combined response of flux difference to both wave types with the marginal effects shown by the slope of the surface at each edge. The marginal effect is the effect of one predictor variable when the other predictor is held constant at 0.

When waves are measured on the same day as flux change (Lag 0) chorus and ULF waves appear to act independently. There is no strong multiplicative interaction visible in the surface plots (Figure 11). Over L 4-5, at the three higher energies, the linear response to ULF power is negative with an increasingly negative response at higher ULF power (the nonlinear term). In contrast, at L=3, ULF is associated with flux increases, particularly at the lower energies. (This can also be seen in the regression coefficients of Figure 10.) The response to VLF waves is linear and positive over L=4-6.

However, a synergistic action is seen between wave types at Lag 1 at the three highest energies (Figure 12). Flux responds more positively to the highest chorus or ULF levels when the other wave type is also at a high level. This can be seen on the surface plots in the unexpected rise in flux difference in the farthest corner (black arrows in 12b mark some of the strongest examples of this). However, as much of the chorus and ULF individual effects may be explained by their action within this multiplicative interaction term, the remaining individual marginal effects may appear negative (blue and orange arrows of 12b). This is most notable with chorus and suggests that much of the positive effect of chorus is the result of joint action with ULF waves rather than from its own individual, additive effects. Increased electron flux, therefore, may be the result of a combination of processes driven by ULF and chorus waves. Flux enhancement is less likely to occur when either wave type occurs alone. Although the nonlinear component of chorus waves still tends to result in flux decreases, ULF waves show some positive, nonlinear effects, with the highest ULF power resulting in higher flux increases (green arrow of 12b).

4. Discussion

In this study, we find that chorus and ULF waves are most explanatory of relativistic flux changes over L=4-6, with their strongest influence seen on the >1.5 MeV electrons. Reductions and enhancements of flux due to these waves occur at different time scales, with decreases occurring near simultaneously with increased wave activity, while enhancement, particularly of higher energy electrons due to ULF waves, occurs over a longer time scale.

When measured on the same day as flux (Lag 0), chorus (above L=3) and ULF waves predict electron decreases (at >0.6 MeV and above). There was little association between flux changes at higher energies and Lag 0 hiss (below L=4), although we note a weak positive correlation at L=3 in the >0.23 flux band. We assume these decreases may represent electron loss. This short timescale has been previously noted for VLF waves: loss due to scattering by chorus or hiss occurs at < 1 day (Summers et al., 2007) and within the plasmasphere (Jaynes et al., 2014).

However, when measured on the previous day (Lag 1), these waves are associated with flux enhancements on the following day, consistent with previous observations of chorus (at L>3) and hiss (below L=3) (Summers et al., 2007), as well as ULF influences on electron flux around 1 MeV (Elkington

et al., 1999). Peak correlations between ULF waves and electron flux have been noted at a lag of 2-3 days (Mann et al., 2004). Our use of partial correlation analysis has refined the peak of the enhancement correlation to a one day lag by removing the loss influences that occur more immediately.

Over $L=4-6$, ULF Pc5 wave activity from two days previous (Lag 2) is also influential at higher L shells and/or at higher electron energies. At Lag 2 (waves measured two days before flux), above $L=3$, the influence of ULF Pc5 increases from low to high electron energy levels, while that of the Lag 1 response drops off. Enhancement processes due to ULF waves, therefore, appear to energize electrons in stages, bringing them from lower to the highest energies over a period of days, a response that has been noted previously to geomagnetic indices and solar wind speed (Rodger et al., 2010). This trend was seen for chorus influences as well over $L=4-5$, but it is much less marked. It is also notable that the Lag 2 ULF influence increases at higher L shells, while the influence of VLF waves drops off. The implications of these trends are that inward radial diffusion, driven by ULF waves, takes some time (> 1 day) to bring electrons to the higher energies, that radial diffusion is of more importance at $L>3$ with increasing influence at higher L shells, and that local acceleration due to VLF waves is only an important factor at $L<5$.

We note that L shell and latitude are correlated in the HEO-3 dataset, resulting in a correlation between pitch angle and L shell. As wave action may depend on pitch angle, this may be at least part of the reason behind the correlational differences seen between flux changes and waves at different L shells (Shprits et al., 2008).

Previously, it was suggested that electron flux peaks are due to both VLF/ELF and ULF acceleration equally near $L=4.5$, but to ULF acceleration alone at and above geosynchronous orbit (O'Brien et al., 2003). However, we have here found that the association of ULF waves with flux changes is stronger than that with VLF waves at all L shells studied (at Lag 1). Although we note that the chorus (VLF) effect is strongest relative to the ULF correlation at $L=4$, its standardized regression coefficient is still always less than the corresponding ULF coefficient. This difference in relative effect of VLF vs. ULF between studies may result from our use of satellite-observed VLF data which is a more accurate depiction of chorus rather than the microburst proxy data used in the previous study. However, there are other differences. We do not analyze flux changes solely following storms, choosing instead to measure daily changes; we use multiple regression to statistically determine the simultaneous effects of VLF and ULF; and the levels of geomagnetic activity between our time period of study (2004-2007) and that of the previous paper (1996-2001) are dissimilar. All these differences may contribute to finding a stronger relative enhancement effect of ULF waves over $L=3-6$.

Above $L=3$, electron decreases (at Lag 0) are more strongly associated with ULF Pc5 wave activity than with chorus. This has been previously observed in several storms (Katsavrias et al., 2015), but we show this with a statistical analysis here. We also show that there is some contribution to electron reductions from chorus, presumably due to scattering of electrons into the loss cone. The strongest effect of chorus on electron decreases occurs at the highest electron energy (>2.9 MeV), with the influence dropping at each higher L shell.

4.1 Nonlinear influences of ULF and VLF wave activity

While at Lag 1 the linear regression model shows ULF Pc5 wave power is more strongly associated with flux changes than chorus, the addition of nonlinear (quadratic) and interaction (multiplicative) terms describes a more nuanced relationship. Flux enhancement is more likely when both chorus and ULF waves are high. Their action is more effective in combination than when alone. This was found previously for electrons at geosynchronous orbit (Simms et al., 2018b) but we now confirm this finding for lower L shells. Not all associations between flux change and wave activity are linear when the multiplicative interaction term is included in the model, nor are they all positive. These marginal (individual) negative responses in the nonlinear models result from much of the flux response being tied up in the multiplicative interaction term. They are what is left over after the more significant, and positive, interactive effect is accounted for. This argues that the processes associated with chorus (local acceleration) and those associated with ULF waves (inward radial diffusion) are both necessary to produce flux enhancements and do not act independently. This is consistent with the proposed two-step process where local acceleration by chorus waves first energizes electrons which are then brought to even higher energies by inward radial diffusion (Jaynes et al., 2015; Katsavrias et al., 2019; Zhao et al., 2019).

The same cannot be said for flux decreases on the same day (Lag 0 analyses; L 4-6). In this case, chorus and ULF waves act independently. There is no strong multiplicative interaction visible in the surface plots. Those processes associated with loss due to chorus (which scatters electrons into the loss cone) and ULF waves (via outward radial diffusion) are able to operate independently. There is some nonlinearity to the flux response to ULF waves with a stronger decrease in flux at high ULF in the lowest energies and less decrease in flux at high ULF in the higher energies.

5. Summary

- a. We use multiple regression to control for the effects of other variables and to determine the timescale of electron decreases vs. enhancement. By holding other factors constant in this way, we see that the individual effects on electron reductions and enhancements due to chorus and ULF waves are often stronger than it would appear from individual correlations.
- b. Over L=4-6 (4.0-6.99), both chorus and ULF Pc5 drive immediate electron decreases and delayed enhancement.
- c. ULF waves consistently show a stronger influence on electron enhancement than do chorus waves. ULF power is also often more associated with electron reductions than chorus.
- d. There is a synergistic interaction between chorus and ULF wave activity on electron enhancement. This means that their combined effect is stronger than would be

384 expected. This points to a two-step process of electron acceleration where local
385 acceleration by chorus waves energizes electrons which are subsequently brought to
386 even higher energies by inward radial diffusion.

- 387
- 388 e. However, chorus and ULF waves drive electron decreases additively. In other words, the
389 actions of the two wave types act independently.
- 390
- 391 f. Contributions of ULF Pc5 and what are likely hiss to electron decreases and
392 enhancement are lower at $L < 4$ than at $L \geq 4$, and minimal at $L = 2$.

393

394

395 Acknowledgements

396 We thank T.P. O'Brien for suggesting this study, advice, and for providing electron data from the HEO-3
397 satellite (data available at: <http://virbo.org/HEO>). We thank R. Gamble for preparing and J.-J. Berthelier
398 for providing DEMETER ICE data, available at the CDPP (Centre de données de la Physique des Plasmas)
399 website: <https://cdpp-archive.cnes.fr/>. CARISMA data are available at
400 <https://www.carisma.ca/carisma-data-repository>. McMAC data can be accessed through the CDAWeb
401 at <https://cdaweb.sci.gsfc.nasa.gov/index.html/>. The ULF Pc5 index is available at
402 http://ulf.gcras.ru/plot_ulf.html. The suggestions of an anonymous reviewer greatly improved this
403 manuscript.

404 Work at Augsburg University was supported by NSF grant AGS-1651263.

405

406 Literature Cited

407

408 Berthelier, J. J., Godefroy, M., Leblanc, F., Malingre, M., Menvielle, M., Lagoutte, D., et al. (2006). ICE,
409 the electric field experiment on DEMETER. *Planetary and Space Science*, 54(5), 456-471,
410 <https://doi.org/10.1016/j.pss.2005.10.016>

411

412 Blake, J. B., D. N. Baker, N. Turner, K. W. Ogilvie, and R. P. Lepping (1997), Correlation of changes in the
413 outer-zone relativistic-electron population with upstream solar wind and magnetic field measurements,
414 *Geophysical Research Letters*, 24, 927, doi.org/10.1029/97GL00859

415

416 Carpenter, D.L. and C. G. Park (1973), On what ionospheric workers should know about the
417 plasmopause-plasmasphere, *Reviews of Geophysics*, 11, 133,
418 <https://doi.org/10.1029/RG011i001p00133>

419

420 Claudepierre, S. G., I. R. Mann, K. Takahashi, J. F. Fennell, M. K. Hudson, J. B. Blake, J. L. Roeder, J. H.
421 Clemmons, H. E. Spence, G. D. Reeves, D. N. Baker, H. O. Funsten, R. H. W. Friedel, M. G. Henderson,

C. A. Kletzing, W. S. Kurth, R. J. MacDowall, C. W. Smith, and J. R. Wygant (2013), Van Allen Probes observation of localized drift resonance between poloidal mode ultra-low frequency waves and 60 keV electrons, *Geophysical Research Letters*, 40, 4491–4497, doi:10.1002/grl.50901.

Clilverd, M A, C J Rodger, M Andersson, P T Verronen, and A Seppälä (2016), Linkages between the radiation belts, polar atmosphere and climate: electron precipitation through wave particle interactions, in *Waves, particles and storms in geospace*, edited by Georgios Balasis, Ioannis A. Daglis, and Ian R. Mann, Chapter 14, 355-376, Oxford University Press, <http://dx.doi.org/10.1093/acprof:oso/9780198705246.003.0015>, ISBN: 9780198705246

Elkington, S.R., M.K. Hudson, A. A. Chan (1999), Acceleration of relativistic electrons via drift-resonant interaction with toroidal-mode Pc-5 ULF oscillations, *Geophysical Research Letters*, VOL. 26, NO. 21, 3273-3276, doi.org/10.1029/1999GL003659

Fennell J.F., J. B. Blake, R. Friedel, and S. Kanekal (2004), The Energetic Electron Response to Magnetic Storms: HEO Satellite Observations, Aerospace Report No. TR-2004(8570)-5

Fennell, J. F. and J. Roeder (2008), Storm time phase space density radial profiles of energetic electrons for small and large K values: SCATHA results, *Journal of Atmospheric and Solar-Terrestrial Physics* 70(14):1760-1773, doi: 10.1016/j.jastp.2008.03.014

Hao, Y. X., Q.-G. Zong, Y. F. Wang, X.-Z. Zhou, Hui Zhang, S. Y. Fu, Z. Y. Pu, H. E. Spence, J. B. Blake, J. Bonnell, J. R. Wygant, and C. A. Kletzing (2014), Interactions of energetic electrons with ULF waves triggered by interplanetary shock: Van Allen Probes observations in the magnetotail, *Journal of Geophysical Research Space Physics*, 119, 8262–8273, doi:10.1002/2014JA020023.

Hao, Y. X., Zong, Q.-G., Zhou, X.-Z., Rankin, R., Chen, X. R., Liu, Y., et al. (2019). Global-scale ULF waves associated with SSC accelerate magnetospheric ultrarelativistic electrons. *Journal of Geophysical Research: Space Physics*, 124, 1525–1538. <https://doi.org/10.1029/2018JA026134>

Horne, R. B., and R. M. Thorne (2003), Relativistic electron acceleration and precipitation during resonant interactions with whistler-mode chorus, *Geophysical Research Letters*, 30(10), 1527, doi:10.1029/2003GL016973

Horne, R. B., R. M. Thorne, S. A. Glauert, J. M. Albert, N. P. Meredith, and R. R. Anderson (2005), Timescale for radiation belt electron acceleration by whistler mode chorus waves, *Journal of Geophysical Research*, 110, A03225, doi:10.1029/2004JA010811

Hyndman, R.J., & Athanasopoulos, G. (2018) *Forecasting: principles and practice*, 2nd edition, OTexts: Melbourne, Australia. [OTexts.com/fpp2](https://www.otexts.com/fpp2)

Jaynes, A.N., X. Li, Q.G. Schiller, L.W. Blum, W. Tu, D.L. Turner, B. Ni, J. Bortnik, D.N. Baker, S.G. Kanekal, J.B. Blake and J. Wygant (2014), Evolution of relativistic outer belt electrons during an extended quiescent period, *Journal of Geophysical Research*, 119, 9558–9566, doi:10.1002/2014JA020125.

Jaynes, A. N., et al. (2015), Source and seed populations for relativistic electrons: Their roles in radiation belt changes, *Journal of Geophysical Research Space Physics*, 120, 7240–7254, doi:10.1002/2015JA021234.

Katsavrias, C., I. A. Daglis, W. Li, S. Dimitrakoudis, M. Georgiou, D. L. Turner, and C. Papadimitriou, (2015), Combined effects of concurrent Pc5 and chorus waves on relativistic electron dynamics, *Annales Geophysicae*, 33, pp 1173–1181, doi:10.5194/angeo-33-1173-2015.

Katsavrias, C., Sandberg, I., Li, W., Podladchikova, O., Daglis, I. A., Papadimitriou, C., et al. (2019). Highly relativistic electron flux enhancement during the weak geomagnetic storm of April–May 2017. *Journal of Geophysical Research: Space Physics*, 124, 4402–4413. <https://doi.org/10.1029/2019JA026743>

Kellerman, A. C., and Y. Y. Shprits (2012), On the influence of solar wind conditions on the outer-electron radiation belt, *Journal of Geophysical Research*, 117, A05217, doi:10.1029/2011JA017253.

Kozyreva, O., V. Pilipenko, M. J. Engebretson, K. Yumoto, J. Watermann, and N. Romanova (2007), In search of a new ULF Pc5 wave index: Comparison of Pc5 power with dynamics of geostationary relativistic electrons, *Planetary Space Science*, 55, 755–769.

Li, L., J. Cao, and G. Zhou (2005), Combined acceleration of electrons by whistler-mode and compressional ULF turbulences near the geosynchronous orbit, *Journal of Geophysical Research*, 110, A03203, doi:10.1029/2004JA010628.

Li, W., R. M. Thorne, V. Angelopoulos, J. Bortnik, C. M. Cully, B. Ni, O. LeContel, A. Roux, U. Auster, and W. Magnes (2009), Global distribution of whistler-mode chorus waves observed on the THEMIS spacecraft, *Geophysical Research Letters*, 36, L09104, doi:10.1029/2009GL037595.

Loto'aniu, T., Thorne, R., Fraser, Brian, & Summers, D. (2006). Estimating relativistic electron pitch angle scattering rates using properties of the electromagnetic ion cyclotron wave spectrum. *Journal of Geophysical Research*. 111. 10.1029/2005JA011452

Loto'aniu, T. M., H. J. Singer, C. L. Waters, V. Angelopoulos, I. R. Mann, S. R. Elkington, and J. W. Bonnell (2010), Relativistic electron loss due to ultralow frequency waves and enhanced outward radial diffusion, *Journal of Geophysical Research*, 115, A12245, doi:10.1029/2010JA015755

Mann, I.R., T.P. O'Brien, D.K. Milling (2004), Correlations between ULF wave power, solar wind speed, and relativistic electron in the magnetosphere: solar cycle dependence, *Journal of Atmospheric and Solar-Terrestrial Physics* 66 (2004) 187– 198, doi:10.1016/j.jastp.2003.10.002

Mann, I. R., Murphy, K. R., Ozeke, L. G., Rae, I., Milling, D. K., and Kale, A. (2012), The role of ultralow frequency waves in radiation belt dynamics, *Geoph. Monog. Series*, 199, 69–92, doi:10.1029/2012GM001349

Neter, J., Wasserman, W., & Kutner, M. H. (1985). *Applied linear statistical models*, (p. 1127). Homewood, Ill: Richard D. Irwin, Inc.

O'Brien, T. P., K. R. Lorentzen, I. R. Mann, N. P. Meredith, J. B. Blake, J. F. Fennell, M. D. Looper, D. K. Milling, and R. R. Anderson (2003), Energization of relativistic electrons in the presence of ULF power and MeV microbursts: Evidence for dual ULF and VLF acceleration, *Journal of Geophysical Research*, 108(A8), 1329, doi:10.1029/2002JA009784

O'Brien, T.P. (2012), Data Cleaning Guidelines for AE-9/AP-9 Datasets, Aerospace Report No. TOR-2012(1237)-4

Ozeke, L. G., Mann, I. R., Dufresne, S. K. Y., Olifer, L., Morley, S. K., Claudepierre, S. G., et al. (2020). Rapid outer radiation belt flux dropouts and fast acceleration during the March 2015 and 2013 storms: The role of ULF wave transport from a dynamic outer boundary. *Journal of Geophysical Research: Space Physics*, 125, e2019JA027179. <https://doi.org/10.1029/2019JA027179>

Reeves, G. D., et al. (2013), Electron acceleration in the heart of the Van Allen radiation belts, *Science*, 341, 999–994, doi:10.1126/science.1237743

Rodger, C. J., Clilverd, M. A., Green, J. C., and Lam, M. M. (2010), Use of POES SEM-2 observations to examine radiation belt dynamics and energetic electron precipitation into the atmosphere, *J. Geophys. Res.*, 115, A04202, doi:10.1029/2008JA014023

Shprits, Y.Y., Subbotin, D.A., Meredith, D.P., Elkington, S.R. (2008), Review of modeling of losses and sources of relativistic electrons in the outer radiation belt II: Local acceleration and loss, *Journal of Atmospheric and Solar-Terrestrial Physics*, 70:1694–1713, doi:10.1016/j.jastp.2008.06.014

Shprits, Y. Y., Thorne, R. M., Friedel, R., Reeves, G. D., Fennell, J., Baker, D. N., and Kanekal, S. G. (2006), Outward radial diffusion driven by losses at magnetopause, *Journal of Geophysical Research*, 111, A11214, doi:10.1029/2006JA011657

Shprits, Y. Y., Meredith, N. P., and Thorne, R. M. (2007), Parameterization of radiation belt electron loss timescales due to interactions with chorus waves, *Geophysical Research Letters*, 34, L11110, doi:10.1029/2006GL029050

Shprits, Y. Y., A. Kellerman, N. Aseev, A. Y. Drozdov, and I. Michaelis (2017), Multi-MeV electron loss in the heart of the radiation belts, *Geophysical Research Letters*, 44, 1204–1209, doi:10.1002/2016GL072258.

Simms, L., Engebretson, M., Clilverd, M., Rodger, C., Lessard, M., Gjerloev, J., & Reeves, G. (2018a). A distributed lag autoregressive model of geostationary relativistic electron fluxes: Comparing the influences of waves, seed and source electrons, and solar wind inputs. *Journal of Geophysical Research: Space Physics*, 123. <https://doi.org/10.1029/2017JA025002>

Simms, L. E., Engebretson, M. J., Clilverd, M. A., Rodger, C. J., & Reeves, G. D. (2018b). Nonlinear and synergistic effects of ULF Pc5, VLF chorus, and EMIC waves on relativistic electron flux at geosynchronous orbit. *Journal of Geophysical Research: Space Physics*, 123. <https://doi.org/10.1029/2017JA025003>

Simms, L. E., Engebretson, M. J., Clilverd, M. A., & Rodger, C. J. (2019). Ground-based observations of VLF waves as a proxy for satellite observations: Development of models including the influence of solar illumination and geomagnetic disturbance levels. *Journal of Geophysical Research: Space Physics*, 124, 2682–2696. <https://doi.org/10.1029/2018JA026407>

Summers, D., R.M. Thorne, and F. Xiao (1998), Relativistic theory of wave-particle resonant diffusion with application to electron acceleration in the magnetosphere, *Journal of Geophysical Research*, 103, 20,487–20,500, <https://doi.org/10.1029/98JA01740>

Summers, D., B. Ni, and N. P. Meredith (2007), Timescales for radiation belt electron acceleration and loss due to resonant wave-particle interactions: 2. Evaluation for VLF chorus, ELF hiss, and electromagnetic ion cyclotron waves, *Journal of Geophysical Research*, 112, A04207, doi:10.1029/2006JA011993.

Tan, L. C., X. Shao, A. S. Sharma, and S. F. Fung (2011), Relativistic electron acceleration by compressional-mode ULF waves: Evidence from correlated Cluster, Los Alamos National Laboratory spacecraft, and ground-based magnetometer measurements, *Journal of Geophysical Research*, 116, A07226, doi:10.1029/2010JA016226.

Thorne, R.M. (2010), Radiation belt dynamics: The importance of wave-particle interactions, *Geophysical Research Letters*, 37, L22107, doi:10.1029/2010GL044990.

Thorne, R. M., T. P. O'Brien, Y. Y. Shprits, D. Summers, and R. B. Horne (2005), Timescale for MeV electron microburst loss during geomagnetic storms, *Journal of Geophysical Research*, 110, A09202, doi:10.1029/2004JA010882.

Tsurutani, B. T. and Edward J. Smith (1977), Two types of magnetospheric ELF chorus and their substorm dependences, *Journal of Geophysical Research*, 82, 5112–5128, doi: 10.1029/JA082i032p05112.

Turner, D. L., Shprits, Y., Hartinger, M., and Angelopoulos, V. (2012), Explaining sudden losses of outer radiation belt electrons during geomagnetic storms, *Nature Physics Lett.*, 8, 208–212, doi:10.1038/NPHYS2185

Turner, D. L., et al. (2017), Investigating the source of near-relativistic and relativistic electrons in Earth's inner radiation belt, *Journal of Geophysical Research Space Physics*, 122, 695–710, doi:10.1002/2016JA023600.

Usanova, M. E., et al. (2014), Effect of EMIC waves on relativistic and ultrarelativistic electron populations: Ground-based and Van Allen Probes observations, *Geophysical Research Letters*, 41, 1375–1381, doi:10.1002/2013GL059024.

Zhao, H., Baker, D. N., Li, X., Malaspina, D. M., Jaynes, A. N., & Kanekal, S. G. (2019). On the acceleration mechanism of ultrarelativistic electrons in the center of the outer radiation belt: A statistical study. *Journal of Geophysical Research: Space Physics*, 124, <https://doi.org/10.1029/2019JA027111>

Zong, Q.-G., X.-Z. Zhou, Y. F. Wang, X. Li, P. Song, D. N. Baker, T. A. Fritz, P. W. Daly, M. Dunlop, and A. Pedersen (2009), Energetic electron response to ULF waves induced by interplanetary shocks in the outer radiation belt, *Journal of Geophysical Research*, 114, A10204, doi:10.1029/2009JA014393.

Zong, Q. R. Rankin, X. Zhou (2017), The interaction of ultra-low-frequency pc3-5 waves with charged particles in Earth's magnetosphere. *Reviews of Modern Plasma Physics*, 1(1), 10, doi: 10.1007/s41614-017-0011-4

Figure Captions

Figure 1. Mean daily electron fluxes at 4 energy ranges (>0.23 , >0.6 , >1.5 , and >2.9 MeV) over $L=2-9$ (northern hemisphere).

Figure 2. Mean daily lower band chorus PSD over $L=2-6$ (northern hemisphere).

Figure 3. Mean daily ULF power (Pc3, Pc4, and Pc5) at McMAC stations BENN ($L=2.5$), GLYN ($L=3.4$), and CARISMA stations PINA ($L=4.06$), ISLL ($L=5.15$), GILL ($L=6.15$), and FCHU ($L=7.44$).

Figure 4. Scatterplots of mean daily lower band chorus \log_{10} PSD on previous day (Lag 1) vs. daily change in >1.5 MeV \log_{10} electron flux channel over $L=2-6$ (a-e). The correlation coefficient (r) is given for each L .

Figure 5. Partial correlations between mean daily lower band chorus \log_{10} PSD lagged over 0-3 days and daily change in 4 electron \log_{10} flux energy ranges (>0.23 , >0.6 , >1.5 , and >2.9 MeV) over $L=2-7$ (a-f) (northern hemisphere). Correlations <0.10 lie within the gray area.

Figure 6. Scatterplots of ULF power (Lag 1) and >1.5 MeV electron flux at $L=5$. a. ULF Pc3, b. ULF Pc4, c. ULF Pc5. All ULF bands show nearly identical correlations with flux.

Figure 7. Partial correlations at $L=5$ between mean daily ULF a. Pc3, b. Pc4, and c. Pc5 over 0-3 days previous and daily change in 4 electron \log_{10} flux energy ranges (>0.23 , >0.6 , >1.5 , and >2.9 MeV).

Figure 8. Scatterplots of mean daily ULF Pc5 \log_{10} power on previous day (Lag 1) vs. daily change in >1.5 MeV \log_{10} electron flux channel over $L=2-7$ (a-f). The correlation coefficient (r) is given for each L .

Figure 9. Partial correlations between mean daily ULF Pc5 lagged over 0-3 days and daily change in 4 electron \log_{10} flux energy ranges (>0.23 , >0.6 , >1.5 , and >2.9 MeV) over $L=2-7$ (a-f). Correlations <0.10 lie within the gray area.

Figure 10. Standardized regression coefficients for each of $L=3-6$ predicting daily change in 4 electron flux energy ranges (a. >0.23 , b. >0.6 , c. >1.5 , and d. >2.9 MeV) using lower band chorus PSD and ULF Pc5 power averaged over the same day (Lag 0; white), the previous day (Lag 1; light blue), and 2 days previous (Lag 2; dark blue) (northern hemisphere). Red lines denote statistically significant coefficients. Percent of variation in flux difference explained by the model is given at the top of each panel. This corresponds to the correlation (r) given within the panel. Lag 1 electron flux is added to each analysis to correct for serial autocorrelation. ULF at GLYN for $L=3$, PINA at $L=4$, ISLL at $L=5$, and GILL at $L=6$.

Figure 11. Predicting change in daily flux using linear and quadratic terms of Lag 0 \log_{10} (ULF Pc5 power) and chorus \log_{10} (PSD) as well as their multiplicative interaction term over $L=3-6$ (a-d) at 4 flux energy

644 ranges (>0.23 , >0.6 , >1.5 , and >2.9 MeV). Lag 1 electron flux is added to each analysis to correct for
645 serial autocorrelation. ULF at GLYN for L=3, PINA at L=4, ISLL at L=5, and GILL at L=6.

646 Figure 12. Predicting change in daily flux using linear and quadratic terms of Lag 1 \log_{10} (ULF Pc5 power)
647 and chorus \log_{10} (PSD) as well as their multiplicative interaction term over L=4-6 (a-d) at 4 flux energy
648 ranges (>0.23 , >0.6 , >1.5 , and >2.9 MeV). Black arrows indicate examples of strong multiplicative
649 interaction; blue arrows are examples of ULF effect becoming negative at high ULF power when the ULF-
650 VLF multiplicative interaction is strong; green arrows indicate nonlinear increased effect at high ULF
651 power. Lag 1 electron flux is added to each analysis to correct for serial autocorrelation. ULF at GLYN for
652 L=3, PINA at L=4, ISLL at L=5, and GILL at L=6.

Figure 1.

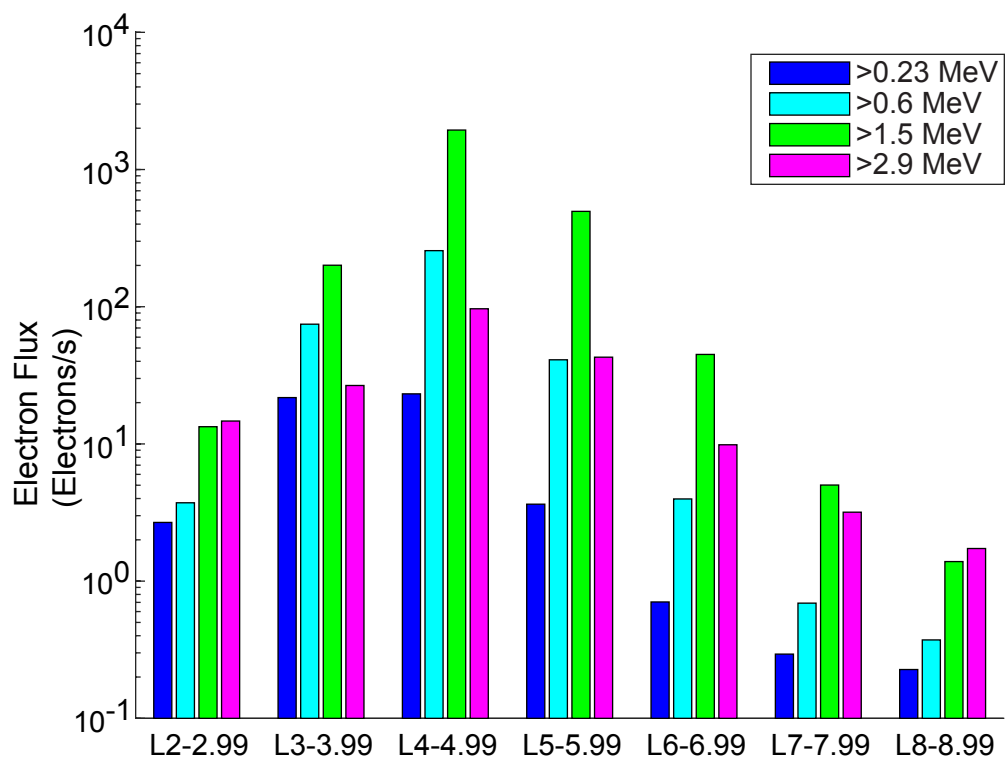


Figure 2.

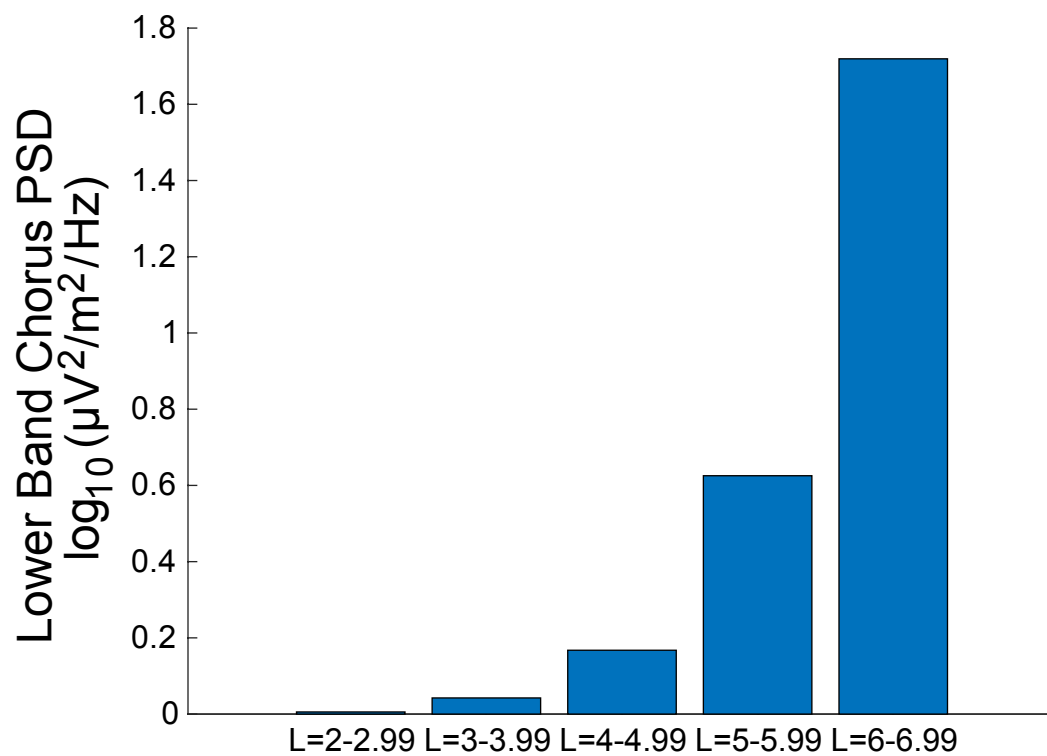


Figure 3.

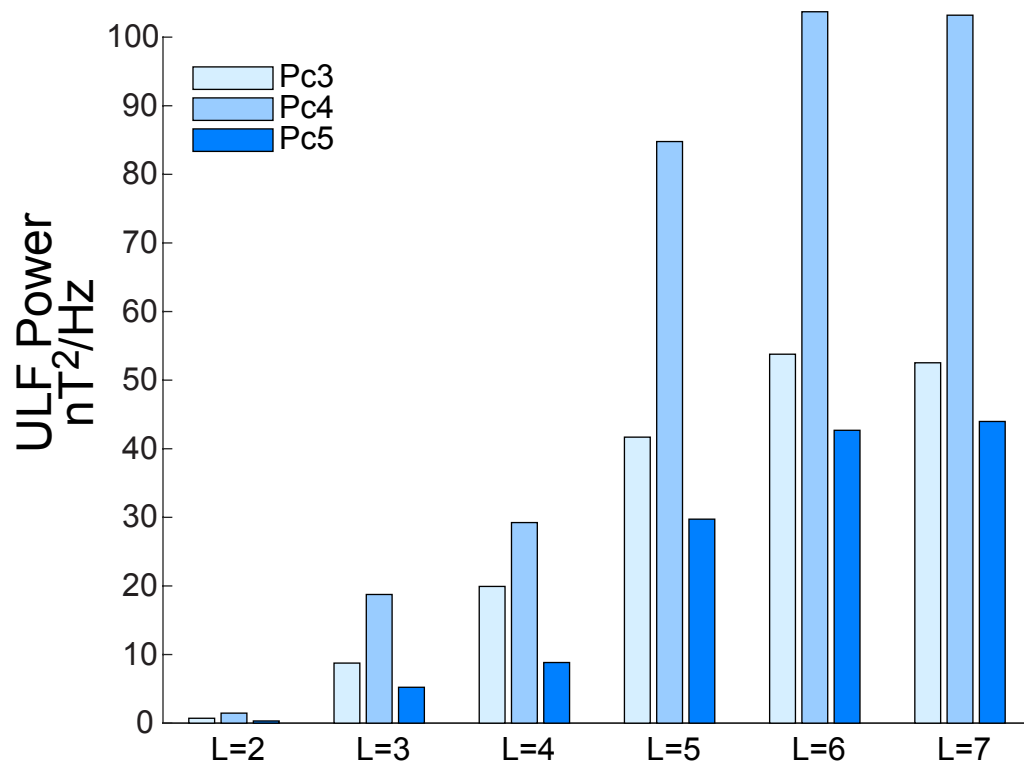


Figure 4.

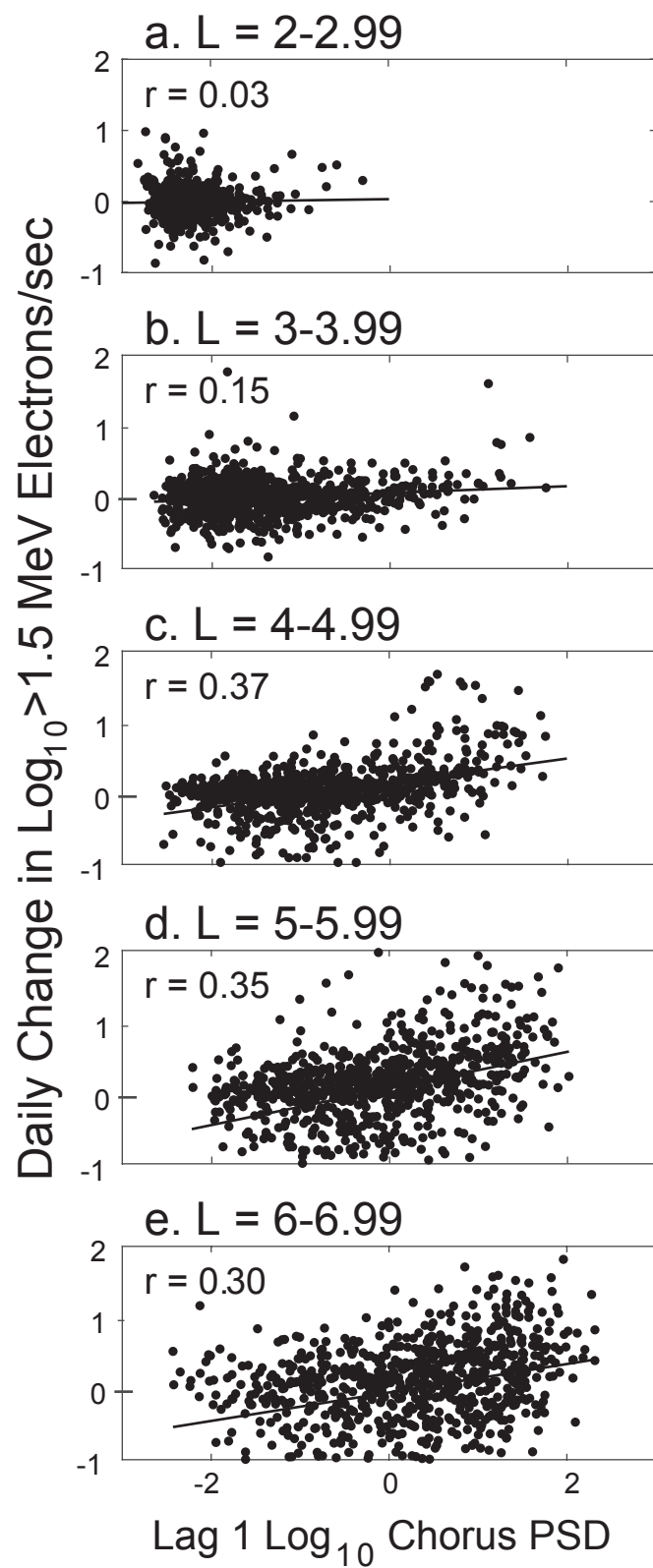


Figure 5.

Partial Correlations of Lower Band Chorus PSD
with Daily Change in Electrons/sec

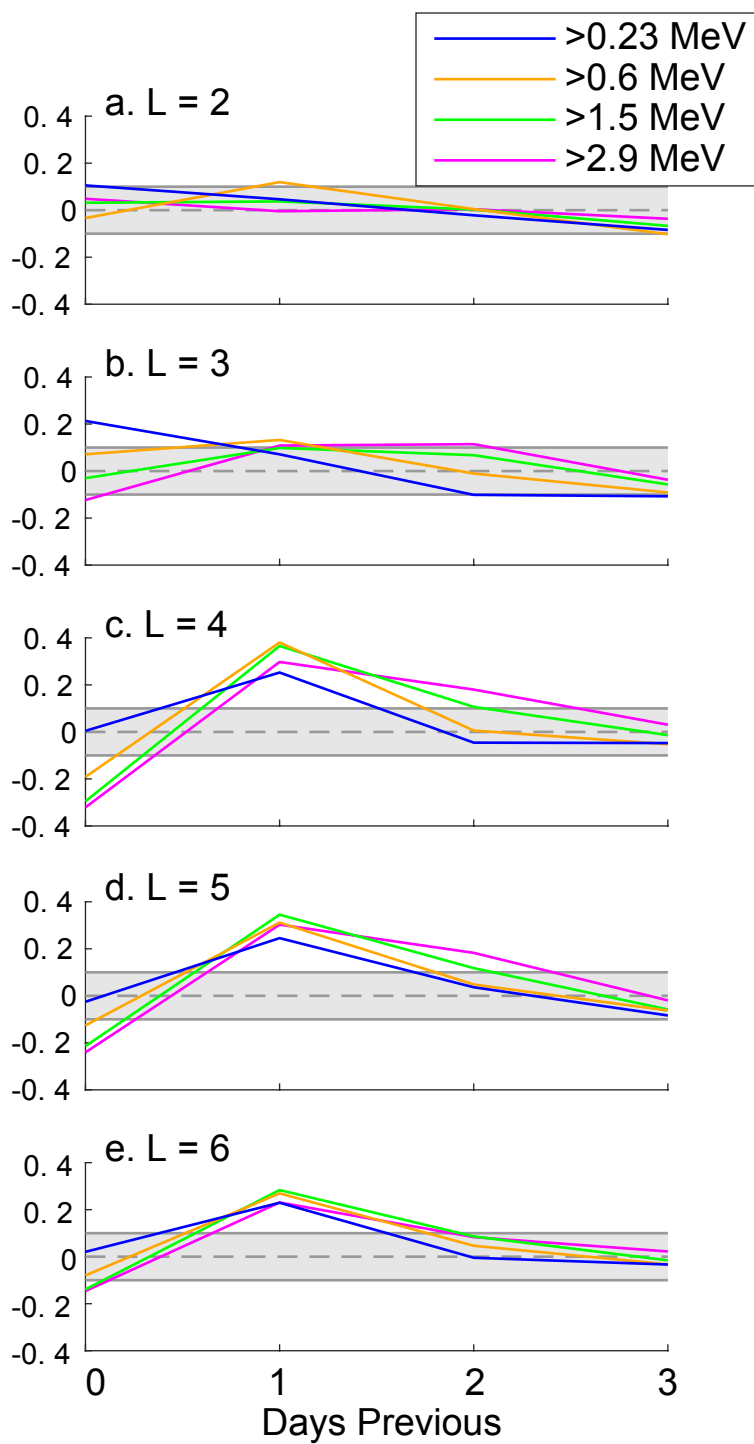
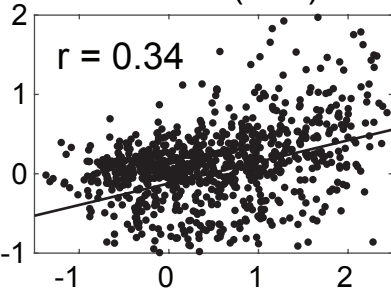


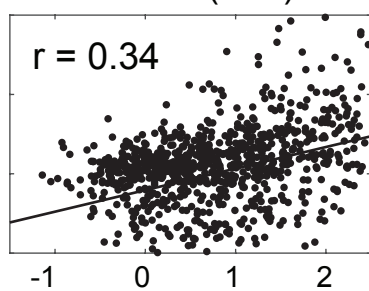
Figure 6.

Daily Change Log_{10}
>1.5 MeV Electrons/sec

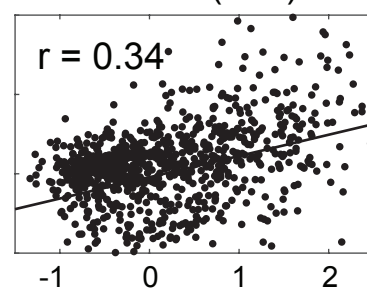
a. ULF Pc3 (L=5)



b. ULF Pc4 (L=5)



c. ULF Pc5 (L=5)



Lag 1 Log_{10} ULF Power

Figure 7.

Partial Correlation of ULF Power (L=5)
with Daily Change in Electrons/sec

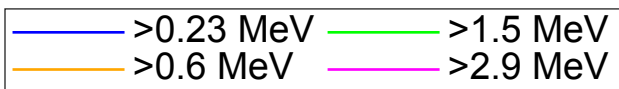
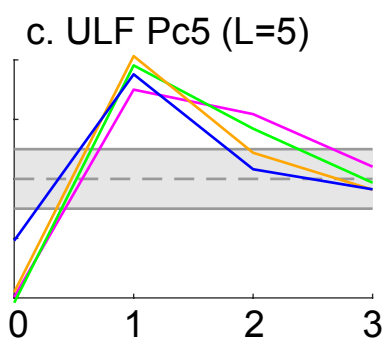
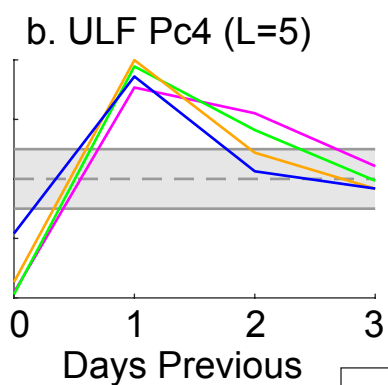
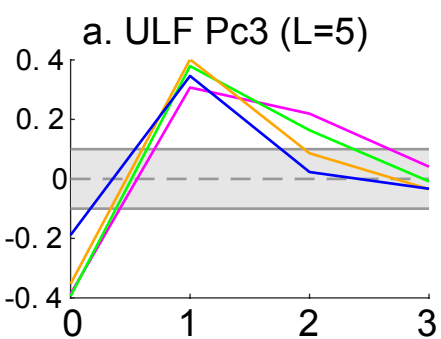


Figure 8.

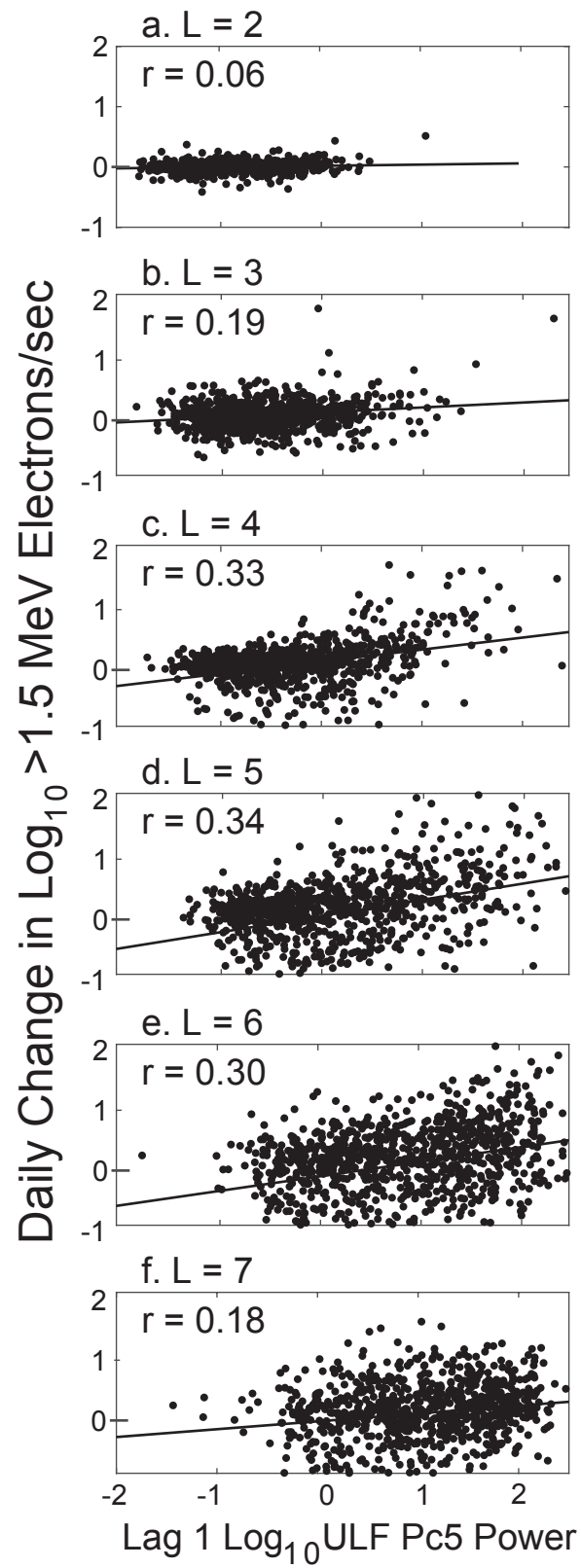


Figure 9.

Partial Correlations of ULF Pc5 Power
with Daily Change in Electrons/sec

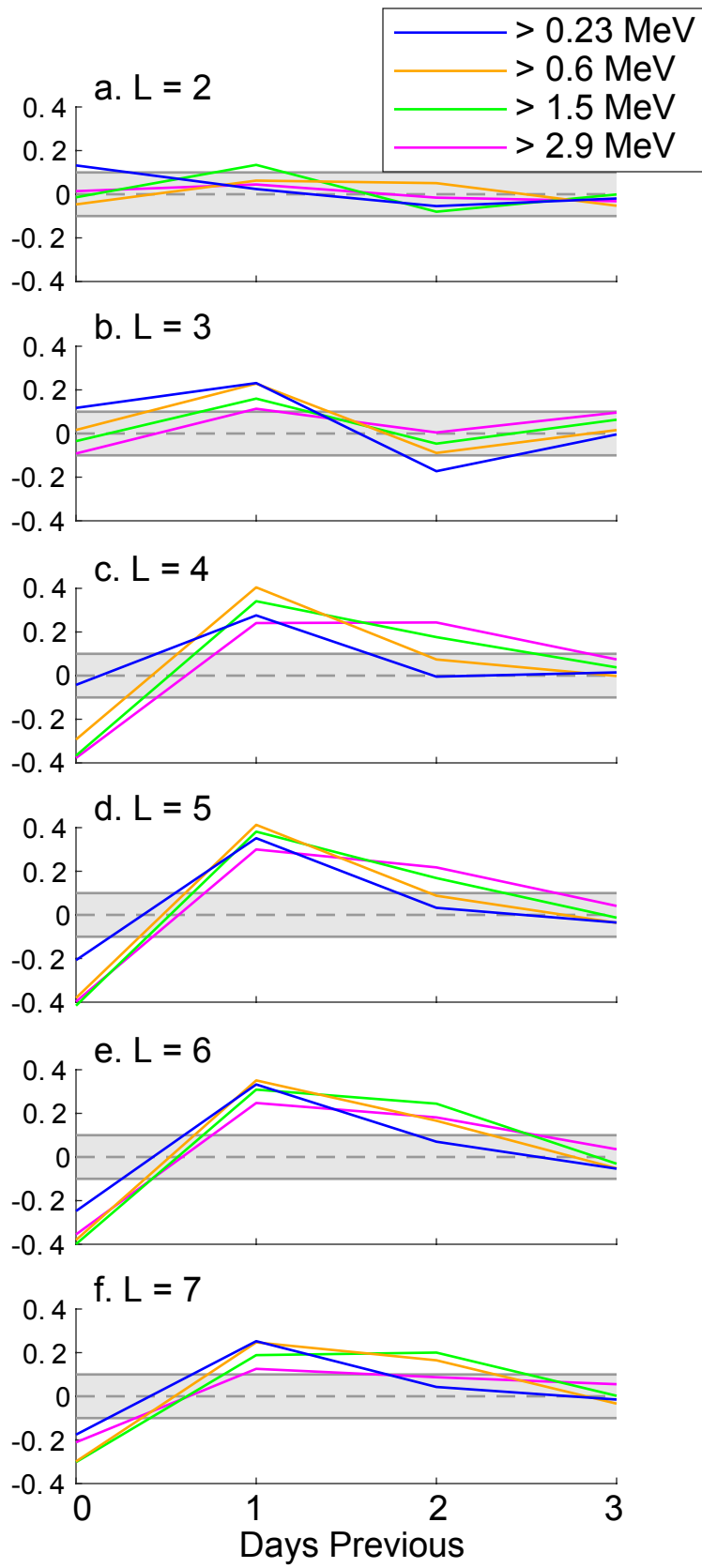


Figure 10.

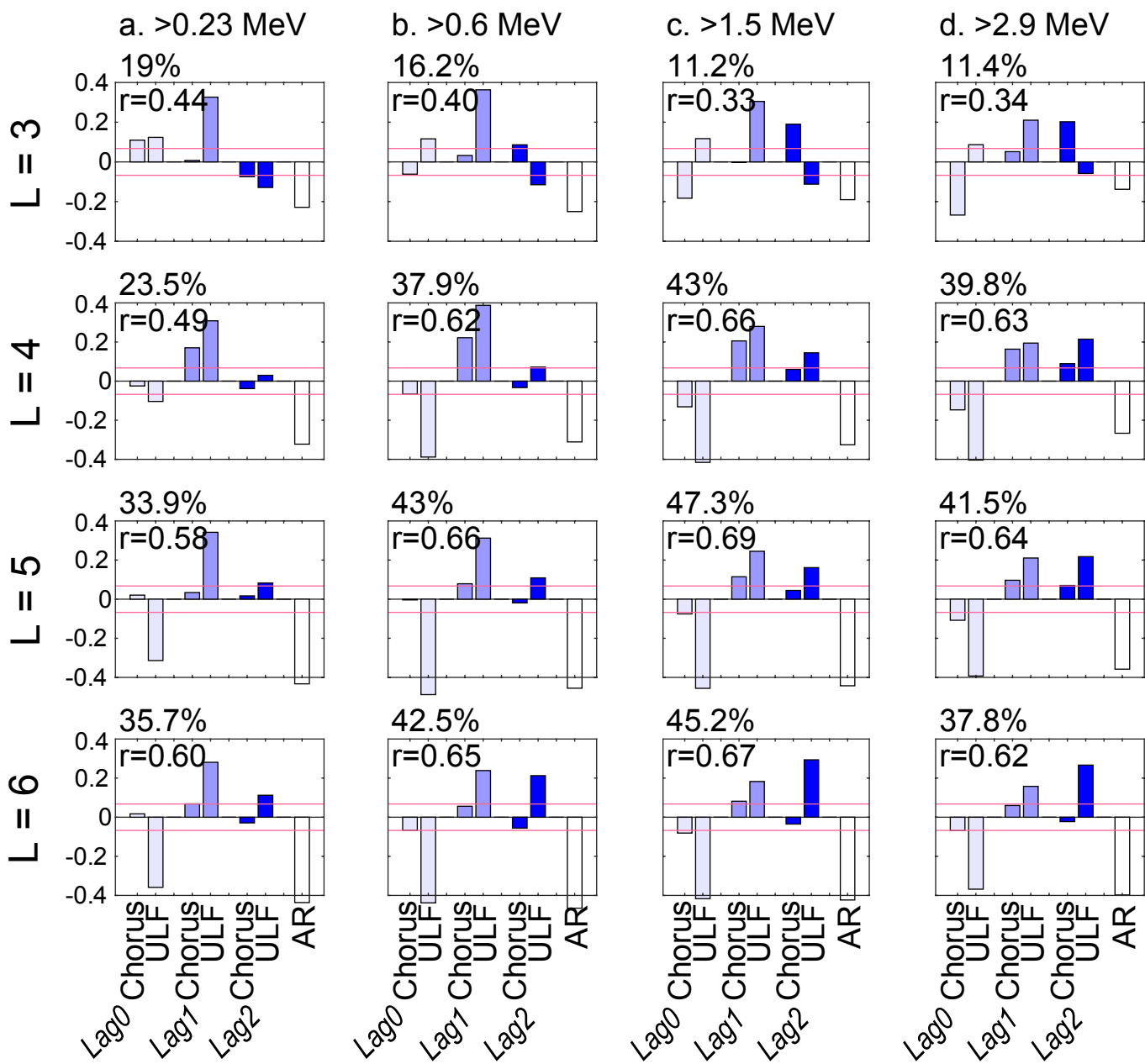


Figure 11.

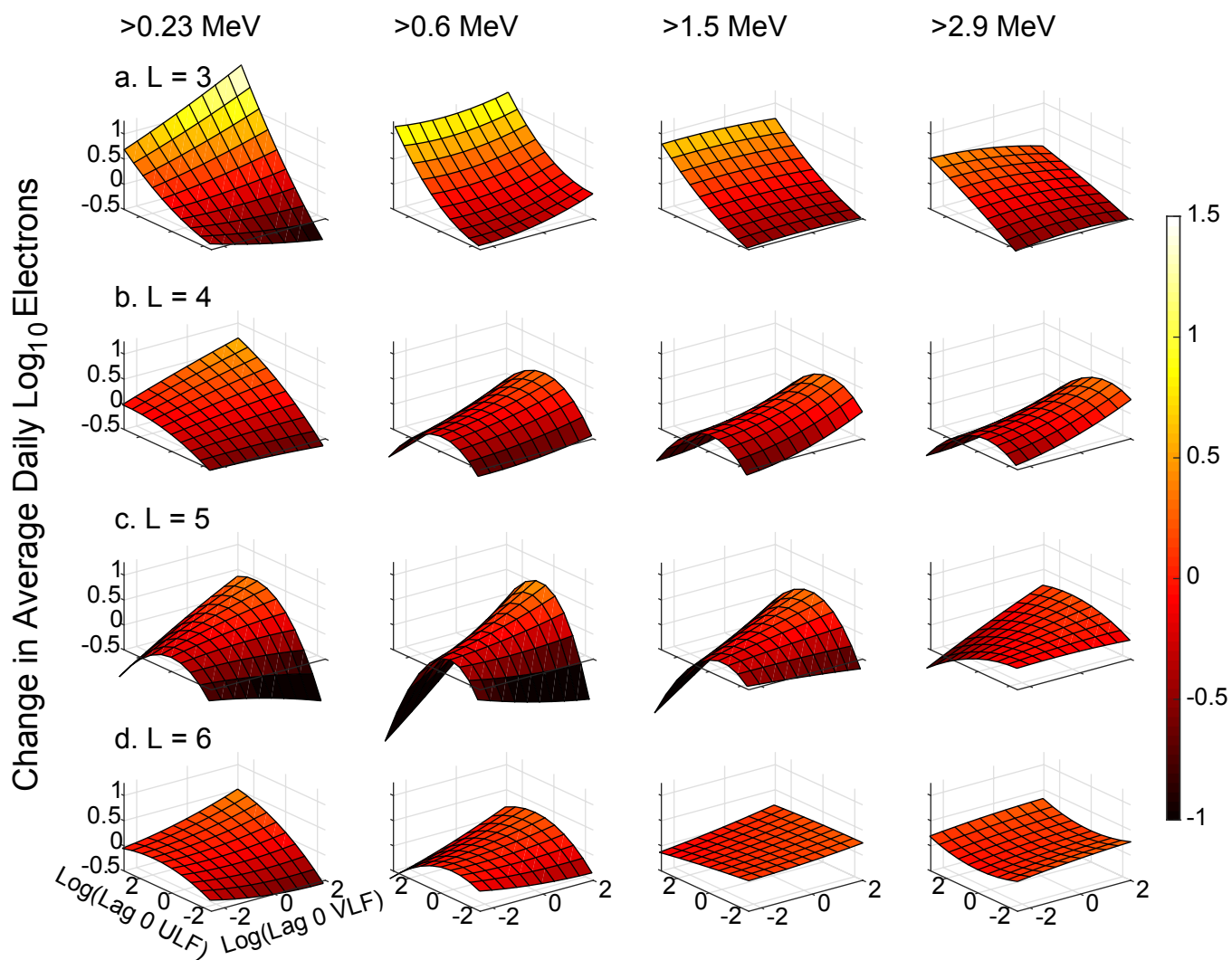


Figure 12.

