

1 **Estimating the occurrence of geomagnetic activity using the Hilbert-Huang**
2 **transform and extreme value theory**

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15 **Key Points**

- 16 • A Hilbert-Huang transform is applied to the geomagnetic aa data to identify
17 the solar cycle dependency in the data
- 18 • Extreme value theory is applied to the aa data separately in solar minimum
19 and maximum conditions
- 20 • March 1989 (Quebec) event is shown to overall be a 1-in-25 year event. But a
21 1-in-130 year event during solar minimum

22 **Abstract**

23 In this paper extreme value theory (EVT) has been used to estimate the return levels
24 for geomagnetic activity based on the aa index. The aa index is the longest,
25 continuously recorded, geomagnetic dataset (from 1868 – Present). This long, 150
26 year, dataset is an ideal candidate for extreme value analysis. However the data are
27 not independent and identically distributed as required for EVT since they are
28 impacted by the approximately 11 year solar cycle. The Hilbert-Huang Transform
29 has been used to identify the solar cycle component in the data and the data has
30 been split into solar maximum and minimum times. In these two regimes the
31 generalised extreme value distribution has been fit to the datasets. These have also
32 been combined for an estimate of the overall return times. The results suggest that
33 the largest event in the database (March 1989) is a one in 25 year event. However,
34 considering separate solar maximum and minimum times has a large impact on the
35 return times. During solar minimum conditions the return time of the March 1989
36 event is 130 years. This suggests that the occurrence of extreme space weather
37 events is conditionally dependent on where in the Solar Cycle we are.

38

39 **Introduction**

40 Geomagnetic storms are disturbances in the Earth's magnetosphere. They are
41 caused by changes in the solar wind which impact the magnetosphere. Two of the
42 main causes for these changes in the solar wind are coronal mass ejections (CMEs)
43 and high-speed solar wind streams (HSSs) [*Schwenn, 2007*]. CMEs usually have
44 large speeds (approximately five times faster than the background solar wind), high
45 energies and large magnetic field strengths [*Riley and Love, 2017*]. HSSs come from

46 solar coronal holes and the fast wind from these regions interacts with the slower
47 upstream wind which create co-rotating interaction regions (CIRs) [*Garion et al.*,
48 2018]. These regions have increased magnetic field strength and higher particle
49 density [*Schwenn, 2007*]. A lot of the largest space weather impacts are associated
50 with geomagnetic storms: geomagnetically induced currents, radio scintillation, solar
51 energetic particle events and enhanced fluxes of relativistic electrons [*Cannon,*
52 2013].

53

54 Indices are used to quantify the relative strength of geomagnetic events, these
55 include Dst ([*World Data Center for Geomagnetism Kyoto et al.*, 2015b]), Kp, Ap
56 ([*Bartels, 1957*]), AE, AO, AL, AU ([*World Data Center for Geomagnetism Kyoto et*
57 *al.*, 2015a]), am, as, an ([*Mayaud, 1980*]), aa and Aa ([*Mayaud, 1972*]). The indices
58 are calculated at different cadences from hourly to daily. Of these, the Dst index is
59 used for identifying and quantifying the severity of a geomagnetic storms [e.g. *Loewe*
60 *and Prohss, 1997*]. However the various indices are, for the most part, closely related
61 to each other.

62

63 Due to the impact of space weather events on human health and technology there is
64 interest in estimating the return time for the most extreme events. A common
65 reference point for extreme space weather events is the so called 'Carrington event'
66 [*Carrington, 1859*], which was one of the largest space weather events in the last
67 200 years [*Cliver and Svalgaard, 2004*]. A key question is what is the return time
68 (likelihood) of extreme space weather events. This is a difficult question to answer
69 satisfactorily as it requires investigating the tails of probability distributions, where

70 there is little data. However, this can be done rigorously using extreme value
71 theory/statistics (EVT).

72

73 EVT is mathematical rigorous and provides sensible measures of uncertainty, which
74 can be very large when there are few data points. Therefore one of the key
75 difficulties associated with using EVT is the need for sets of large independent large
76 samples. For example when looking at extreme temperatures (hot or cold) in
77 meteorology the annual maxima or minima are suitable time scales since, on the
78 whole, yearly temperatures are independent, whilst daily temperatures are not. In the
79 space weather domain it would be ideal to take solar cycle (11 year) minima or
80 maxima time series. Unfortunately there is not enough recorded data to have enough
81 data points remaining for effective analysis.

82

83 A number of authors have applied EVT to the space weather domain, including:
84 [*Elvidge and Angling, 2018; Koons, 2001; Meredith et al., 2015; Silbergleit, 1996;*
85 *1999; Siscoe, 1976; Thomson et al., 2011; Tsubouchi and Omura, 2007*]. Extreme
86 space weather events have also been estimated using techniques other than EVT
87 such as fitting power law distributions, log-normal distributions and generalized
88 Pareto distribution [*Chapman et al., 2020; Chapman et al., 2018; Riley, 2012; Riley*
89 *and Love, 2017*]. The overall goal of each of the papers is to try to quantify the
90 statistics of a particular measurable associated with an extreme space weather
91 event.

92

93 In terms of investigating geomagnetic activity: *Silbergleit* [1996] and *Tsubouchi and*
94 *Omura* [2007] used EVT to investigate extreme events in the DST index using 23
95 and 44 years of data respectively, *Koons* [2001] used the Ap index using 66 years of
96 data, *Siscoe* [1976] and *Silbergleit* [1999] variants of the aa index, using 91 and 124
97 years of data, and *Thomson et al.* [2011] the rate of change of the magnetic field
98 using 31 years. *Riley and Love* [2017] also estimated the probability of extreme DST
99 events using 60 years of data. Of those, the 91 and 124 year datasets of *Siscoe*
100 [1976] and *Silbergleit* [1999] are very useful for EVT since they capture the longest
101 time period (8 – 11 solar cycles). However, as well as the length of the datasets, the
102 number of data used is also crucial in reducing the uncertainty in the analysis.
103 *Siscoe* [1976] performed EVT only using the three largest events in each solar cycle
104 during the test period, resulting in 27 data points and *Silbergleit* [1999] used the
105 maximum value from each solar cycle, resulting in 12 data points. Whilst these
106 approaches break the data into suitable scale sizes, few data points remain which
107 means there are substantial uncertainties in the results.

108

109 In this paper the aa index (1868 – 2018; 150 years) is analysed using the annual
110 maximum values. This results in the largest temporal span of data for EVT in the
111 space weather domain. Recent work by *Chapman et al.* [2020] has used a linear
112 “mapping” between the top few percent of the aa index and the annual minimum Dst
113 index value to estimate the probability of extreme Dst values. This work takes
114 advantage of the extra data points from this aa-to-Dst mapped set.

115 Whilst geomagnetic storms tend to last between two and seven days, the annual
116 maxima has been chosen to increase the likelihood of having independent,

117 identically distributed data, a requirement for extreme value theory. Using, for
118 example, weekly rather than annual maxima can introduce further dependencies in
119 the dataset as individual events may, or may not, originate from the same solar
120 active region. This could result in the analysed data not being identically distributed.
121 However the disadvantage of using annual maxima is that the solar variability over
122 the ~11 year cycle remains embedded in the data. This temporal dependence is
123 accounted for by using the Hilbert-Huang transform [*Huang and Wu, 2008*] to split
124 the data into solar maximum and minimum times.

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126

127 **Extreme Value Theory**

128 EVT provides a sophisticated approach for estimating probability distribution
129 functions, and specifically for looking at the tail of such distributions. The method
130 avoids any starting assumption about the underlying distribution [*Coles, 2001*]. The
131 key result from EVT is the Fisher-Tippett-Gnedenko (FTG) theorem which states that
132 the maximum of an independent and identically distributed (iid) random variable
133 converges to one of only three possible distributions: the Gumbel distribution
134 [*Gumbel, 1935*], the Fréchet distribution [*Fréchet, 1927*], or the Weibull distribution
135 [*Weibull, 1951*], which can be grouped into the generalized extreme value
136 distribution.

137

138 Specifically, for a sequence of iid random variables X_1, X_2, \dots, X_n with common
139 distribution function F let $M_n = \max \{X_1, \dots, X_n\}$ and $w = \sup \{x: F(x) < 1\}$ then

140

141 $\Pr(M_n \leq x) = \Pr(X_1 \leq x, \dots, X_n \leq x) = F^n(x)$ (Equation 1)

142

143 Then as $n \rightarrow \infty$, $F^n(x) \rightarrow 0$ if $x < w$ and $F^n(x) \rightarrow 1$ otherwise, as such $M_n \rightarrow w$. To

144 avoid a degenerate distribution $F^n(x)$ is normalised. Assuming there is a non-

145 degenerate distribution G such that, for normalising constants $a_n > 0$ and b_n :

146

147 $\lim_{n \rightarrow \infty} F^n(a_n x + b_n) = G(x)$ (Equation 2)

148

149 where G is the generalized extreme value (GEV) distribution defined by

150

151 $G(x) = \exp \left\{ -1 \left[1 + \xi \left(\frac{x-\mu}{\sigma} \right) \right]^{-\frac{1}{\xi}} \right\}$ (Equation 3)

152

153 defined for $1 + \frac{\xi(x-\mu)}{\sigma} > 0$ and where μ is the location parameter, $\sigma > 0$ the scale

154 parameter and ξ the shape parameter [Coles, 2001]. For $\xi < 0$ the GEV reduces to

155 the Weibull distribution, for $\xi > 0$ the Fréchet distribution and in the limit $\xi \rightarrow 0$, $G(x)$

156 reduces to

157

158 $G(x) = \exp \left\{ - \exp \left(- \frac{x-\mu}{\sigma} \right) \right\}$, (Equation 4)

159

160 the Gumbel distribution.

161

162 The parameters of the GEV are usually estimated using a maximum log likelihood
163 method [Coles, 2001]. However, the requirement that the variables must be iid is
164 usually a barrier with using raw data directly and some form of pre-processing is
165 normally required.

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168 **Data**

169 The aa index is a global geomagnetic index which is based on the largest horizontal
170 deviation of the magnetic field measured in nT. It is based on data from two nearly
171 antipodal stations, one in the UK and another in Australia and has been continuously
172 recorded since 1868. Over the 150 years the stations where the data have been
173 recorded has changed. In order to maintain a constant value for the index the
174 weighting of the different stations have varied over time (Table 1).

175

176 *Table 1: Weighting factors for the stations used to compile the aa index [International*
177 *Service of Geomagnetic Indices, 2013].*

Northern hemisphere station (UK)			Southern hemisphere station (Australia)		
Time range	Station	Weighting factor	Time range	Station	Weighting factor
1868 - 1925	Greenwich	1.007	1868 - 1919	Melbourne	0.967
1926 - 1956	Abinger	0.934	1920 - 1979	Toolangi	1.033
1957 - Present	Hartland	1.059	1980 - Present	Canberra	1.084

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179

180 The main advantage of using the aa index for EVT is that it is the longest running
181 planetary index of geomagnetic activity. This long sample time helps in reducing the
182 uncertainties in the EVT extrapolation, however on this time scale the impact of the
183 solar cycles becomes apparent, which has been shown to have an impact on the
184 results of extreme value modelling [*Riley and Love, 2017*]. Figure 1 shows the time
185 series of the aa index in the top panel (each point is the annual maximum aa value)
186 and the bottom panel of the figure shows the periodogram created with a Hamming
187 windowing function on the data. The large peak in the periodogram corresponds to
188 10.7 years and is the solar cycle contribution to the data. To perform EVT on the
189 dataset this temporal dependence should be accounted for. In this work the Hilbert-
190 Huang transform (HHT) is used to identify solar maximum and minimum times, which
191 are assumed to each be iid, and EVT can be performed on each.

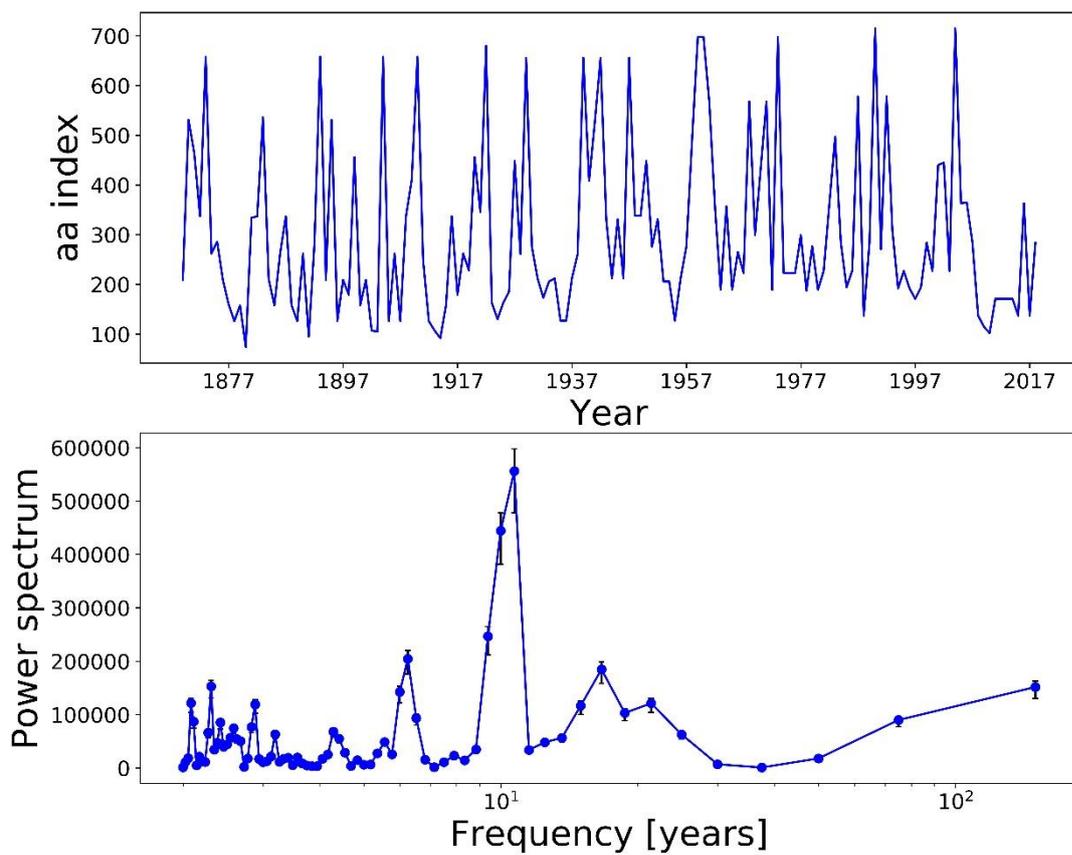
192

193 The HHT decomposes a time series into intrinsic mode functions (IMFs) and then
194 finds the instantaneous frequency of each IMF [*Norden E Huang and Wu, 2008*]. The
195 first step of the HHT is to use empirical mode decomposition (EMD) [*N. E. Huang et*
196 *al., 1998*]. Similar to the Fourier and Wavelet transform, EMD splits a signal into its
197 components, called intrinsic mode functions (IMFs). An IMF is a function in which the
198 number of extrema and zero-crossings differ by at most one and at each point the
199 mean value of the envelopes, defined by the local maxima and minima, is zero. The
200 sum of the IMFs reconstitute the original signal. Hilbert spectral analysis (HSA) can
201 then be used by applying the Hilbert transform to each IMF to find the instantaneous
202 frequency [*N. E. Huang et al., 1998*]. Unlike Fourier and Wavelet transforms HHT is

203 an algorithmic approach rather than theoretical. However the main advantage of the
204 HHT over Fourier and Wavelet is that it is suitable for nonlinear and non-stationary
205 data.

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208

209 Figure 1. Top panel shows the time series of the annual maxima aa index. Bottom
210 panel shows the periodogram of the series.

211

212

213 Applying the HHT to the aa index data (top panel of Figure 1) results in nine IMFs
214 shown in Figure 2. The top panel of the figure shows the original aa index values (in
215 red), then each of the IMFs are shown in blue. Using HSA the envelope of each IMF
216 has been found (shown in green) and the value of the instantaneous time period
217 (ITP) (one over the instantaneous frequency) is shown in the upper left of each IMF
218 plot. A single value for the ITP is found by fitting a linear polynomial through the data.
219 In each case the polynomial was of order zero (as expected) and the constant term
220 is shown. From the figure it can be seen that the third IMF has an ITP which
221 corresponds to the peak time period from Figure 1. This provides confidence that this
222 particular IMF corresponds to the sunspot cycle component in the aa index data.

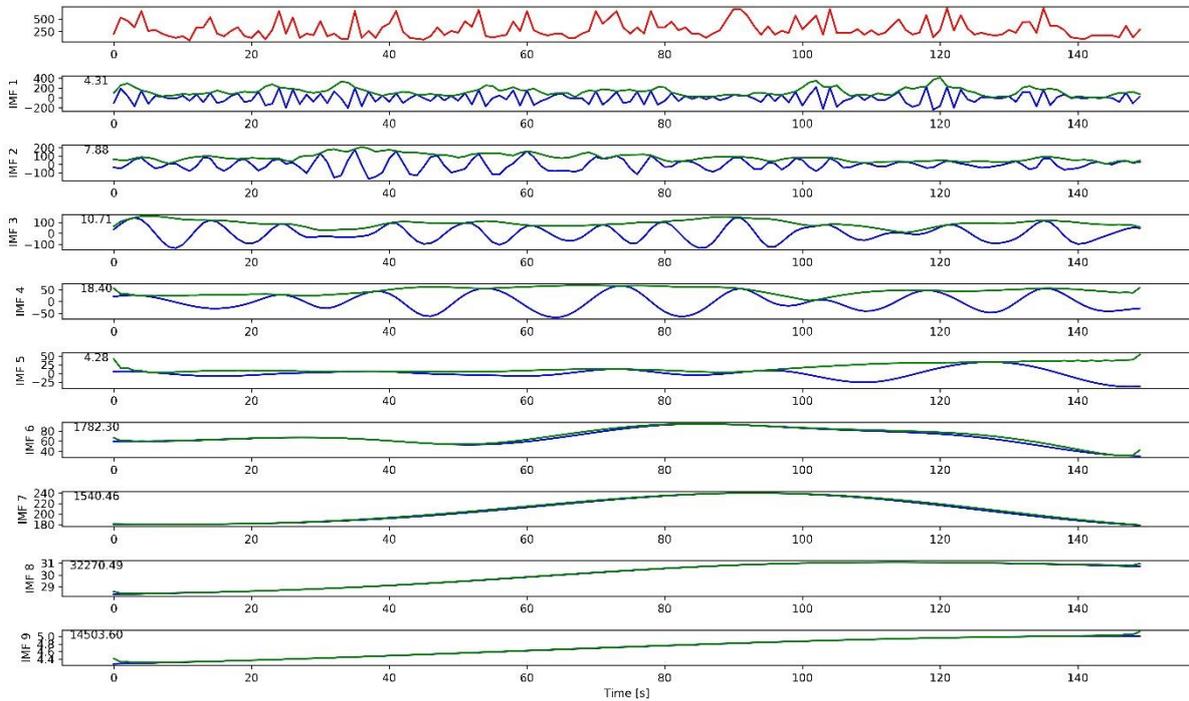
223

224 This particular IMF could be used in the fitting of the GEV as part of a temporally
225 varying location and scale parameters [Coles, 2001]. However this requires
226 propagation of the IMF forward in time. With only simplistic ways of interpolating the
227 IMF forward in time (it is hard to predict the next solar cycle) it makes finding the
228 return times difficult (a key attribute for extreme value modelling).

229

230 Instead the aa variables can be made to be iid by separating the solar maximum and
231 minimum conditions. In each of the the two cases it can be assumed that they are
232 distributed according to the same function. Rather than using a separate dataset to
233 try and determine when these times are, the third (solar cycle dependent) IMF can
234 be used. The IMF is centred at zero and positive values can be used to describe
235 solar maximum times, whilst negative values can be used as solar minimum times.
236 Comparing the IMF estimated solar maximum/minimum times to the sunspot cycle

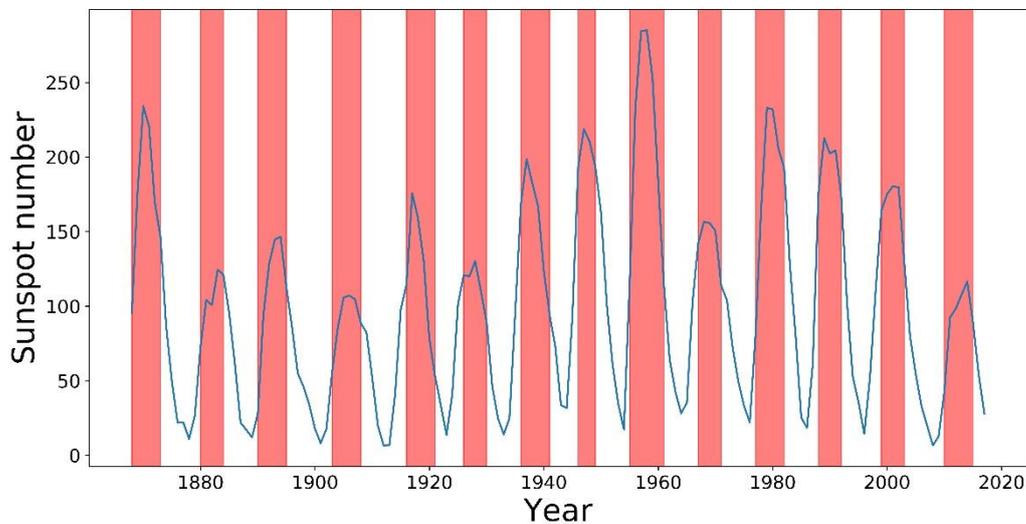
237 (as independent verification of this approach) shows that the estimated maximum
238 overlaps with the peak of the sunspot number as would be expected (Figure 3).



239

240 Figure 2. aa index decomposed into its 9 IMFs. The original aa index is shown in the
241 top panel (red) and each panel below shows an IMF (blue). For each IMF the
242 envelope has been found using the Hilbert transform (green) and the dominant
243 instantaneous frequency value is shown in the upper left of each IMF plot.

244



245

246 Figure 3. Annual maximum sunspot number [SILSO World Data Center, 2018]. Solar
247 maximum times (from the third IMF from the aa index) are shaded.

248

249 **Results**

250

251 Fitting the GEV to the solar maximum and minimum time series using least log
252 likelihood results in the estimated μ, σ, ξ (and standard errors) as shown in Table 2.
253 One method for verifying the quality of the fit of the GEV distribution is by looking at
254 the quantile plot which shows the pairs

255

$$256 \left\{ G^{-1} \left(\frac{i}{n+1} \right), x_i \right\}, \quad i \in \{1, \dots, n\} \quad (\text{Equation 5})$$

257

258 for the ordered annual aa values $\{x_1, x_2, \dots, x_n\}$ and where $G^{-1}(x)$ is the inverse of
259 Equation 3, given by [Coles, 2001]:

260

$$261 G^{-1} \left(\frac{i}{n+1} \right) = \mu - \frac{\sigma}{\xi} \left(1 - \left(-\log \left(\frac{i}{n+1} \right) \right)^{-\xi} \right). \quad (\text{Equation 6})$$

262

263 The quantile plot for solar maximum and minimum times is shown in Figure 4. The
264 plot being roughly linear is an indication of a good agreement between the model fit
265 and the empirical data.

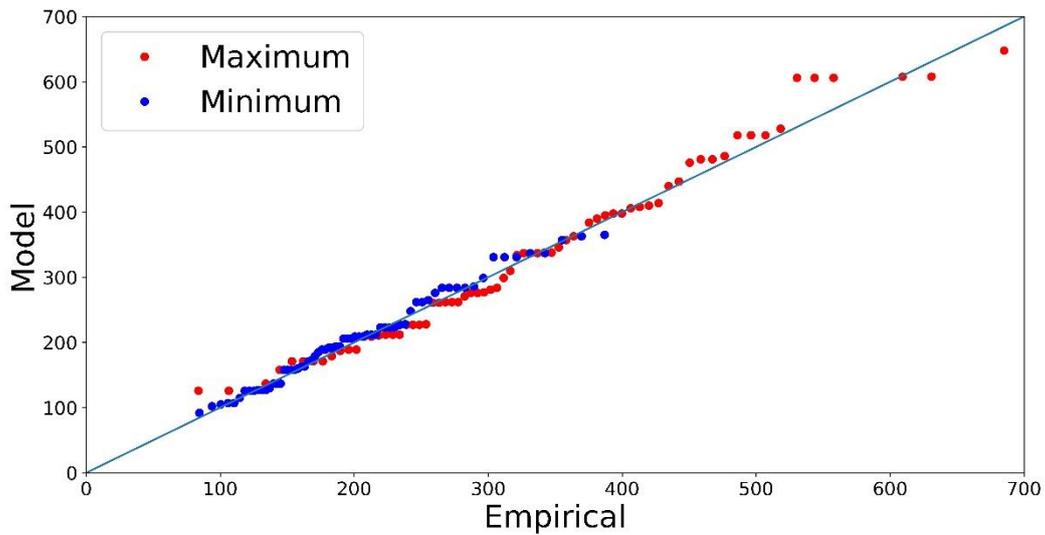
266

267 Table 2. The GEV fit parameters for the solar maximum and minimum conditions, the
 268 standard error is shown in parenthesis.

	Number data points	μ	σ	ξ
Maximum	78	279 (17.8)	130 (12.6)	-0.03 (0.12)
Minimum	72	174 (9.24)	68.8 (7.16)	0.15 (0.10)

269

270



271

272 Figure 4. Quantile plot of the GEV fit for solar maximum and minimum times.

273

274

275 Using the fitted GEV distributions the return time for any given event can be
 276 estimated for either solar maximum or minimum conditions using the values in Table
 277 2. However another useful return time would combine both solar maximum and
 278 minimum. The combined return time can be found by solving

279

$$G_{max}(z_p)G_{min}(z_p) = 1 - \frac{1}{p} \quad \text{(Equation 7)}$$

281

282 for z_p , the return level with return period $1/p$ (z_p is expected to be exceeded by the
 283 annual maximum aa value with probability p), and where G_{max}, G_{min} are the GEV
 284 distributions defined by the parameters in Table 2 [Coles, 2001]. The return levels for
 285 10, 50, 100, 500 and 1000 years are shown in Table 3, with the standard errors
 286 shown in parenthesis. It is interesting to note the difference in return times between
 287 solar maximum and minimum times. This suggests that the probability of an extreme
 288 event is conditionally based on what part of the solar cycle we are in. This is in
 289 agreement with the findings of *Riley and Love* [2017] who determined that, assuming
 290 a power law distribution, the probability of geomagnetic storm exceeding the
 291 Carrington event (in terms of Dst) was 1.4% during solar minimum conditions and
 292 28% for solar maximum conditions.

293

294 Table 3. Return levels for 10, 50, 100, 500 and 1000 year return periods for solar
 295 maximum and minimum conditions as well combined results. The standard error is
 296 shown in parenthesis.

	10-year	50-year	100-year	500-year	1000-year
Maximum	611 (50)	873 (146)	989 (208)	1271 (408)	1399 (519)
Minimum	365 (30)	565 (89)	670 (131)	969 (287)	1127 (386)
Combined	631 (60)	898 (155)	1019 (216)	1321 (429)	1463 (562)

297

298

300 These results have been compared to previous EVT work on the Ap index (an index
 301 similar to the aa index) undertaken by *Koons* [2001]. To compare the results, since
 302 the data are on different scales, they have been normalised by dividing through by
 303 the peak value of the March 1989 event [*Feynman and Hundhausen*, 1994]. This is
 304 the largest value in the aa index database (tied with the “Halloween” event of 2003)
 305 and the second largest in Ap (the largest is an event is from November 1960). *Koons*
 306 [2001] provides the fit parameters for the Gumbel distribution which was shown to
 307 have the best fit with the data. However no standard error was reported in the
 308 results. So for ease of comparison in this work the same 66 years of data was
 309 analysed and fit with the same Gumbel distribution as in *Koons* [2001] ($\mu = 99.1409$
 310 and $\sigma = 42.9416$). The normalised results for both the aa and Ap index data are
 311 shown in Table 4.

312

313 Table 4. Return levels for this work as well as previous EVT work on Ap [*Koons*,
 314 2001]. The values are normalised by dividing through by the index value of the event
 315 in March 1989.

	Normalised 10-year	Normalised 50-year	Normalised 100-year	Normalised 500-year	Normalised 1000-year
Aa (Combined)	0.88 (0.08)	1.26 (0.21)	1.43 (0.30)	1.85 (0.60)	2.05 (0.79)
Ap [<i>Koons</i> , 2001]	0.80 (0.08)	1.09 (0.17)	1.21 (0.22)	1.43 (0.34)	1.54 (0.42)

316

317

318 Comparing the EVT results between the aa and Ap index show that they are very
319 similar for 'short' return times (10 years) with values of 0.88 and 0.80 respectively.
320 These differences widen as the return periods get longer. This is to be expected as
321 in this work 150 years of aa data have been used for the estimates compared to 66
322 years of Ap data from *Koons* [2001]. The extra data points should provide better
323 estimates for the longer return periods, especially the 100-year return level.

324

325 **Conclusions**

326

327 Extreme value theory (EVT) has been used to estimate the return levels for
328 geomagnetic activity based on the aa index. The aa index is the longest,
329 continuously recorded, geomagnetic dataset. This long, 150 year, dataset is an ideal
330 candidate for extreme value analysis. Whilst the aa index is not the most commonly
331 used space weather index its close relationship with the more commonly used Dst
332 index [*Chapman et al., 2020*] implies that similar geoeffective impacts of extreme Dst
333 events would be felt during extreme aa events. However, the aa data are not
334 independent and identically distributed (iid) as required for EVT as they are impacted
335 by the approximately 11 year solar cycle. The Hilbert-Huang Transform has been
336 used to identify the solar cycle component in the data and the data have been split
337 into solar maximum and minimum times. In these two regimes the variables are
338 assumed to be iid, and the generalised extreme value (GEV) distribution has been fit
339 to the two datasets. These have also been combined for an estimate of the return
340 times.

341

342 The results suggest that the largest event in the database (March 1989 / October
343 2003) is a one in 25 year event (but with a standard error of 86 years). Whilst this
344 may seem counter-intuitive, since there are only two events of that size (aa of 715) in
345 the 150 year database, there are in total eight events where the aa index exceeds
346 650. Considering separate solar maximum and minimum times has a large impact on
347 the return time. During solar minimum the return time of the March 1989 event is 130
348 years (with a standard error of 145 years). This value seems reasonable since there
349 is one event with an aa > 650 during solar minimum in the 150 year database
350 (August 4th 1972; [Knipp et al., 2018]). However it is in contrast to the results of
351 *Riley and Love* [2017] who report the probability as ~1-in-700 years during solar
352 minimum (assuming a power law distribution). This demonstrates the uncertainty that
353 can arise when extrapolating extreme events by using different underlying
354 assumptions. Quantifying the impact of solar minimum is of particular importance in
355 quantifying the likelihood associated with extreme space weather events if a period
356 of extended solar minimum is entered. It has been estimated that there is a 15 - 20%
357 chance of returning to Maunder Minimum-like conditions within the next 40 years
358 [*Ineson et al., 2015; Lockwood, 2010*].

359

360 **Acknowledgements**

361 The aa data can be downloaded from British Geological Survey
362 (http://www.geomag.bgs.ac.uk/data_service/data/magnetic_indices/aaindex.html).

363

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