

Coupled Aqua and Ridge Planets in the Community Earth System Model

Xiaoning Wu¹, Kevin A. Reed¹, Christopher L. P. Wolfe¹, Gustavo M. Marques², Scott D. Bachman², Frank O. Bryan²

¹School of Marine and Atmospheric Sciences, Stony Brook University, Stony Brook, NY

²Climate and Global Dynamics Laboratory, National Center for Atmospheric Research, Boulder, CO

Key Points:

- Two baseline examples of fully coupled CESM with idealized ocean geometry, Aqua and Ridge, are presented
- With sufficient resolution, coupled Aqua has a global cold belt of equatorial upwelling and corresponding “reverse Hadley” cells
- Ridge’s zonal asymmetry is crucial for making its circulations more Earth-like compared to Aqua

Corresponding author: Xiaoning Wu, xiaoning.wu.1@stonybrook.edu

Abstract

Idealized models can reveal insights into Earth’s climate system by reducing its complexities. However, their potential is undermined by the scarcity of fully coupled idealized models with components comparable to contemporary, comprehensive Earth System Models. To fill this gap, we compare and contrast the climates of two idealized planets which build on the Simpler Models initiative of the Community Earth System Model (CESM). Using the fully coupled CESM, the Aqua configuration is ocean-covered except for two polar land caps, and the Ridge configuration has an additional pole-to-pole grid-cell-wide continent. Contrary to most sea surface temperature profiles assumed for atmosphere-only aquaplanet experiments with the thermal maximum on the equator, the coupled Aqua configuration is characterized by a global cold belt of wind-driven equatorial upwelling, analogous to the eastern Pacific cold tongue. The presence of the meridional boundary on Ridge introduces zonal asymmetry in thermal and circulation features, similar to the contrast between western and eastern Pacific. This zonal asymmetry leads to a distinct climate state from Aqua, cooled by $\sim 2^{\circ}\text{C}$ via the radiative feedback of clouds and water vapor. The meridional boundary of Ridge is also crucial for producing a more Earth-like climate state compared to Aqua, including features of atmospheric and ocean circulation, the seasonal cycle of the Intertropical Convergence Zone, and the meridional heat transport. The mean climates of these two basic configurations provide a baseline for exploring other idealized ocean geometries, and their application for investigating various features and scale interactions in the coupled climate system.

Plain Language Summary

Simplified climate models can improve our understanding of the Earth’s climate system by stripping down its complexities. Previous simplified climate models — with idealized ocean shapes — have laid great groundwork, but their coarse resolution and overly reduced model components are hard to relate to contemporary models for international climate assessments. We fill this gap by presenting two simplified climate models with components and resolution similar to that of state-of-the-art Earth system models. Aqua is ocean-covered except for two polar land caps, and Ridge has an additional pole-to-pole strip continent. Ridge’s ocean, like the Pacific, has a western pool and eastern cold tongue upwelled from below. On Aqua, without continents blocking the east-west direction, the equatorial upwelling extends globally, forming a cold belt. This results in a warmer global climate on Aqua than Ridge, as clouds over Ridge’s warm pool reflect away more solar radiation than Aqua’s cold and dry equatorial region. Ridge’s strip continent in the north-south direction makes its climate more Earth-like than Aqua, including circulation and poleward transport of energy. The capacity of the Aqua and Ridge planets enables the application to problems of scientific interest and societal impacts, such as El Niño and hurricanes.

1 Introduction

Idealized models are illuminating tools for understanding Earth’s climate system (Held, 2005; Maher et al., 2019). By reducing the complexities of the coupled climate system in terms of boundary conditions or model physics, idealized models have helped advance the scientific understanding of various aspects and scales of the climate system (e.g., Manabe & Bryan, 1969; Ferreira et al., 2010; Wolfe & Cessi, 2010; Abernathey et al., 2013; Voigt & Shaw, 2015; Chavas et al., 2017; Brunetti et al., 2019), as well as the evaluation and development of climate model components (Chang et al., 2001; Reed & Jablonowski, 2012; Bachman & Fox-Kemper, 2013; Herrington & Reed, 2017). The availability of idealized models, embedded within a hierarchy of complexity leading up to state-of-the-art, comprehensive Earth System Models used for climate projection and assess-

63 ments (Eyring et al., 2016), can serve as a valuable resource for climate research and ed-
64 ucation (Jeevanjee et al., 2017; Polvani et al., 2017; Schultz et al., 2017).

65 Focusing on the atmosphere-ocean system, ocean-covered representations of Earth
66 (commonly referred to as aquaplanets) have been widely used for either the atmospheric
67 or ocean component at various degree of complexity, but fully coupled configurations are
68 relatively scarce. For the atmospheric component, there is a rich history of application
69 for aquaplanets (Neale & Hoskins, 2000; Blackburn et al., 2013), with either prescribed
70 sea surface temperature (e.g., Medeiros et al., 2016) or slab ocean configurations (e.g.,
71 Donohoe et al., 2014; Benedict et al., 2017) as the simplified lower boundary condition,
72 forgoing ocean dynamics. Example topics of study using aquaplanet configurations in-
73 clude the hemispheric asymmetry in tropical rainfall (Frierson et al., 2013), the length
74 scale of extratropical storm tracks (Kaspi & Schneider, 2011), and the effect of off-equatorial
75 thermal forcing on tropical cyclone activity (Ballinger et al., 2015). For the ocean com-
76 ponent forced by a prescribed atmosphere, idealized ocean basins are used for understand-
77 ing the overturning circulation (Wolfe & Cessi, 2010; Jones & Cessi, 2016; Cessi & Jones,
78 2017; Ferrari et al., 2017) and factors affecting salinity (Jones & Cessi, 2017, 2018). For
79 global and coupled configurations, earlier works (Smith et al., 2006; Farneti & Vallis, 2009)
80 have explored the global climates of selected ocean geometries. Other notable examples
81 using coupled aquaplanets include a hierarchy of idealized ocean geometries (Marshall
82 et al., 2007; Enderton & Marshall, 2009; Ferreira et al., 2010). These simplified designs
83 demonstrate remarkable resemblance to the observed Earth climate on the planetary scale,
84 including the meridional heat transport (Czaja & Marshall, 2006; Marshall et al., 2007;
85 Enderton & Marshall, 2009) and ocean salinity contrast (Ferreira et al., 2010; Nilsson
86 et al., 2013). However, these configurations, oriented towards the global-scale ocean cir-
87 culation with extremely simplified atmospheres (e.g., Molteni, 2003) at $\sim 3^\circ$ horizontal
88 resolution or coarser, do not aim to address important atmospheric processes that de-
89 pend on higher horizontal and vertical resolution, or more complete model physics (Ballinger
90 et al., 2015; Herrington & Reed, 2017).

91 In summary, a gap is present in the hierarchy between previously available ideal-
92 ized models and comprehensive Earth System Models. Specifically, there is currently no
93 coupled idealized model available with comprehensive model physics equivalent to those
94 used for the Coupled Model Intercomparison Project (CMIP; Eyring et al., 2016) for both
95 the atmospheric and ocean components. This lack of availability undermines the appli-
96 cation of idealized modeling to process-level understanding of CMIP-class models where
97 atmosphere-ocean coupling plays a key role, and precludes the full investigation of scale
98 interactions of scientific and societal interest at the atmosphere-ocean interface (e.g., Scoc-
99 cimarro et al., 2017; Carranza et al., 2018; Li & Sriver, 2018).

100 To fill this gap, by building on the Simpler Models initiative (Polvani et al., 2017,
101 <http://www.cesm.ucar.edu/models/simpler-models>) of the Community Earth System Model
102 (CESM; Hurrell et al., 2013; Danabasoglu et al., 2020), we have developed two fully cou-
103 pled baseline configurations with idealized ocean geometry. The new development brings
104 unique, CMIP-relevant modeling capabilities into the idealized framework. In this study,
105 we present the mean climates of the two configurations and discuss the contrast between
106 them. The first one, Aqua, is ocean-covered except for minimal polar land caps; the sec-
107 ond one, Ridge, has a single meridional boundary. Comparing and contrasting with pre-
108 vious idealized studies, these two configurations demonstrate the role of ocean geome-
109 try in the coupled climate state, including impacts on meridional heat transport. The
110 assessment of these two basic configurations provides a baseline for exploring additional
111 forms of idealized ocean geometries, and their application to the study of various fea-
112 tures and scale interactions in the coupled climate system.

113 This paper is organized as follows. Section 2 describes the details of model config-
114 uration, and the simulation data under analysis. Section 3 presents the mean climates
115 of the CESM Aqua and Ridge planets from the perspectives of the energy budget, the

116 large-scale circulation, and the meridional heat transport. Finally, Section 4 discusses
 117 the results in the context of previously documented models, and the outlooks for future
 118 work.

119 2 Data and Methods

120 The idealized configurations are developed in the framework of CESM (Hurrell et
 121 al., 2013; Danabasoglu et al., 2020), a state-of-the-art, community modeling tool. With
 122 numerous options for configuration and a vibrant user community, CESM provides the
 123 capacity to produce simulations for international climate assessments (Eyring et al., 2016),
 124 as well as reduced-complexity options for fundamental investigations and continued model
 125 component development (Polvani et al., 2017). We expand on currently available options
 126 of atmosphere-only or slab ocean aquaplanets (Medeiros et al., 2016; Benedict et al., 2017),
 127 and introduce fully coupled configurations with dynamical ocean.

128 Two types of idealized ocean geometries are configured, as shown in Fig. 1. For Aqua,
 129 the planet is ocean-covered except for two polar continents that reach down to 80°N/S.
 130 The presence of the polar continents, occupying minimal area, is required by the ocean
 131 grid. For Ridge (Smith et al., 2006; Enderton & Marshall, 2009), a single grid-cell-wide
 132 strip of pole-to-pole continent is added as a meridional boundary for the ocean basin.
 133 All land has zero orography.

134 The atmospheric component is the Community Atmosphere Model version 4 (CAM4;
 135 Neale et al., 2010). The choice of model version is made to balance complexity and com-
 136 putational cost. The finite-volume dynamical core, based on a regular latitude-longitude
 137 grid, is built upon a 2D shallow water approach (Lin & Rood, 1996, 1997) and mass-conservative
 138 in flux-form. The parameterization schemes include deep convection (Zhang & McFar-
 139 lane, 1995), shallow moist convection (Hack, 1994), dry boundary layer turbulence (Holtslag
 140 & Boville, 1993), and cloud physics, radiation, etc. further described in Neale et al. (2010).
 141 The horizontal resolution is nominally 1°, resulting in grid spacing of ~110 km in the
 142 tropical regions. In the vertical direction, the model is divided into 26 layers in a hybrid
 143 sigma-pressure coordinate system, with finer spacing near model bottom and top (~3
 144 hPa). Settings for the solar constant, dry mass, greenhouse gas concentrations, ozone
 145 distribution, and aerosols are adapted from the Aqua-Planet Experiment (Neale & Hoskins,
 146 2000).

147 The ocean component is the Modular Ocean Model version 6 (MOM6; Adcroft et
 148 al., 2019), the latest update to replace the previous CESM ocean component starting with
 149 CESM3. One advantage of MOM6 is the versatile specification of vertical layers via the
 150 use of the Arbitrary-Lagrangian-Eulerian algorithm (Hirt et al., 1974; Bleck, 1978). The
 151 horizontal resolution is nominally 2°, with equatorial refinement to 1°. The ocean max-
 152 imum depth is 4000 m, divided into 57 vertical layers, with thickness decreasing from
 153 ~250 m at the bottom to 2.5 m near the ocean surface. The effects of mesoscale eddies
 154 are parameterized by activating two schemes in the tracer equation. The first scheme
 155 follows the ideas of Gent et al. (1995), where available potential energy is removed from
 156 the large scale by flattening isopycnals. A constant thickness diffusivity of 2000 m² s⁻¹
 157 is used without any vertical structure. The associated eddy-induced transport is applied
 158 as a bolus velocity. To avoid the problems associated with layer thickness diffusion de-
 159 scribed by Holloway (1997), this scheme is implemented as an interface height diffusion.
 160 Following Solomon (1971) and Redi (1982), the second scheme represents the diffusive
 161 mixing of tracers along neutral surfaces, which is implemented using a finite-volume general-
 162 coordinate methodology. Again, a constant along-isopycnal tracer diffusivity of 2000 m²
 163 s⁻¹ is used. The K-Profile vertical mixing Parameterization (KPP; Large et al., 1994)
 164 is applied via the Community ocean Vertical Mixing (CVMix; Griffies et al., 2015) frame-
 165 work. The diapycnal diffusivity is 2×10^{-5} m²s⁻¹, the Laplacian horizontal viscosity
 166 is 1×10^4 m²s⁻¹, and the coefficient for quadratic bottom drag is 0.005. To provide to-

167 pographic form drag for balancing the momentum input from the atmosphere, we pre-
 168 scribe zonally and hemispherically symmetric bottom topography in analytical, sinusoidal
 169 form (see Fig. S1). This is particularly important for the Aqua case, and we use the same
 170 bottom topography in both cases for consistency. The bottom topography has the max-
 171 imum height of 500 m in the vertical, with horizontal length scale of 1000 km in the merid-
 172 ional direction, and 45° in the zonal direction. These length scales are chosen to avoid
 173 subgrid-scale signals.

174 The sea ice component is the Community Ice CodE version 5 (CICE5; Bailey et
 175 al., 2018), with full thermodynamics and dynamics. Since the quasi-equilibrium climate
 176 states of both configurations are too warm for sea ice formation (see Table 1), the sea
 177 ice component is present but not active for the period under analysis in the present study.
 178 As needed by the minimal presence of land, the land component is the Community Land
 179 Model version 5 (CLM5; Lawrence et al., 2019) at the same horizontal resolution as the
 180 atmospheric component. For the polar and ridge continents, the land surface type is set
 181 to wetland, which behaves most similarly to a slab ocean in comparison with other land
 182 surface types. Precipitation over land, a small amount, is returned to the ocean by ad-
 183 justing the water balance in the MOM6 component. The coupling is handled by the Com-
 184 mon Infrastructure for Modeling the Earth (CIME; <http://github.com/ESMCI/cime>, see
 185 description in Danabasoglu et al. (2020)). The coupling frequency for all components is
 186 hourly, based on the spatial resolution of model components.

187 For both Aqua and Ridge configurations, the diurnal cycle is retained, and an ide-
 188 alized seasonal cycle is imposed by setting the orbital obliquity to 23.3° . Model initial-
 189 ization is zonally symmetric for all components (atmosphere, ocean, land, and sea ice).
 190 On the National Science Foundation (NSF)-supported Cheyenne supercomputer housed
 191 at the National Center for Atmospheric Research (NCAR), the model achieves through-
 192 put of ~ 80 simulated years per wall-clock day, while archiving annually averaged out-
 193 put for the ocean and monthly averaged output for all other components. By Year 400
 194 of the 500-year integration, although the deep ocean is still drifting, the top-of-atmosphere
 195 (TOA) radiative balance has adjusted close to equilibrium for both configurations (im-
 196 balance $\sim O(0.1) \text{ Wm}^{-2}$, see Table 1 and Fig. S2). We discuss the climate state of Year
 197 401–500 in the following section, using monthly averaged output for the atmosphere and
 198 annually averaged output for the ocean.

199 3 Results

200 Fig. 1 illustrates the state of the coupled Aqua and Ridge planets, with snapshots
 201 of their oceans and atmospheres in boreal summer. Both planets are warm and ice-free.
 202 For the zonally symmetric Aqua, the sea surface temperature (SST; Fig. 1a) shows a global
 203 cold belt of equatorial upwelling that persists through the seasonal cycle (see animation
 204 in supplement). A common feature of coupled Aqua configurations with dynamical oceans
 205 (Smith et al., 2006; Marshall et al., 2007; Farneti & Vallis, 2009), this local SST min-
 206 imum on the equator is markedly different from the typical SST patterns used for atmosphere-
 207 only Aqua-Planet Experiments (Neale & Hoskins, 2000). For Ridge (Fig. 1b), the pres-
 208 ence of the meridional boundary leads to the formation of a western warm pool, limit-
 209 ing the global equatorial upwelling of Aqua to eastern upwelling in the cold tongue. Anal-
 210 ogous to the Pacific, besides the local equatorial upwelling, the equatorward eastern bound-
 211 ary current also contributes to the cold tongue via advection (Wyrski, 1981; Kessler, 2006).
 212 These SST patterns, in turn, influence the characteristics of their atmospheres. Both plan-
 213 ets exhibit a rich variety of synoptic systems, including extratropical storms and trop-
 214 ical cyclone-like vortices (Fig. 1c–d). For Aqua (Fig. 1c), on either side of the cold
 215 and dry equator, the atmosphere is remarkably rich in moisture even in the winter hemisphere.
 216 This is associated with Aqua’s unique circulation patterns in the seasonal cycle, discussed
 217 later in Section 3.2. For Ridge (Fig. 1d), the winter hemisphere is noticeably drier com-
 218 pared to its summer hemisphere, especially around the cold tongue and the eastern bound-

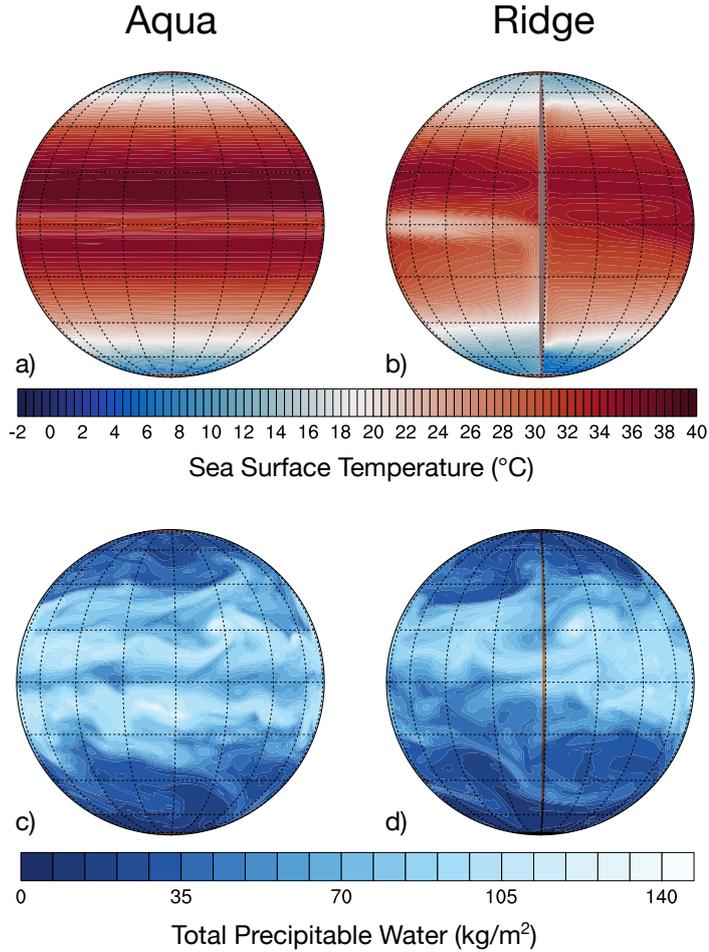


Figure 1. Illustration of the Aqua and Ridge planets. The polar land caps and the ridge continent are marked in brown. (a–b) SST ($^{\circ}\text{C}$) for August (100-yr climatology), showing the global cold belt of equatorial upwelling on Aqua, and the eastern and western boundary currents on Ridge (see animation of the seasonal cycle in supplement); (c–d) Instantaneous snapshots of total precipitable water (kgm^{-2}) from boreal summer, displaying various synoptic systems.

219 any current. The presence of the western warm pool is reflected in the rich reservoir of
 220 atmospheric moisture in the region.

221 The contrast in these thermodynamic and dynamic features, with an emphasis on
 222 the zonal asymmetry of Ridge, is further detailed in Fig. 2 with the 100-year climatol-
 223 ogy. For Aqua (Fig. 2, left column), the equatorial atmosphere is uniformly associated
 224 with subsidence (Fig. 2a), as a result of local SST minimum in the equatorial region. Driven
 225 by mild easterly wind stress (Fig. 2c), the equatorial belt of upwelling (Fig. 2e) produces
 226 a shallow thermocline in the ocean underneath (Fig. 2i). In contrast, Ridge (Fig. 2, right
 227 column) produces many Pacific-like features: a Walker-like circulation (Fig. 2b) devel-
 228 ops, with convection over the moist western warm pool, and subsidence over the dry east-
 229 ern cold tongue; the convergence of zonal wind stress around 120°E (Fig. 2d) marks the
 230 location of the warmest equatorial SST (Fig. 2f), producing a zonal SST gradient of $\sim 8^{\circ}\text{C}$
 231 averaged over 5°N – 5°S (Fig. 2h), contrary to Aqua’s zonal uniformity (Fig. 2g). Corre-

Table 1. Statistics of the global mean, annually averaged over Year 401-500. Global mean ocean salinity is a constant value of 34.969 psu for both planets, due to the absence of sea ice.

| | Unit | Aqua | | Ridge | |
|--------------------------------------|---------------------|----------|--------|----------|--------|
| | | Avg. | Stdev. | Avg. | Stdev. |
| Surface temperature | °C | 27.466 | 0.104 | 25.503 | 0.071 |
| Surface pressure | hPa | 1016.580 | 0.067 | 1015.690 | 0.040 |
| Total cloud fraction | fraction | 0.444 | 0.002 | 0.472 | 0.004 |
| Cloud radiative forcing | Wm ⁻² | -23.166 | 0.250 | -25.857 | 0.315 |
| Total precipitable water | kgm ⁻² | 58.070 | 0.694 | 49.194 | 0.396 |
| Precipitation rate | mmday ⁻¹ | 4.384 | 0.020 | 4.182 | 0.014 |
| Net shortwave (TOA) | Wm ⁻² | 261.507 | 0.223 | 257.822 | 0.345 |
| Net longwave (TOA) | Wm ⁻² | 261.129 | 0.286 | 258.091 | 0.236 |
| Net shortwave (ocean surface) | Wm ⁻² | 183.856 | 0.318 | 181.852 | 0.390 |
| Net longwave (ocean surface) | Wm ⁻² | -44.443 | 0.378 | -48.341 | 0.252 |
| Downwelling longwave (ocean surface) | Wm ⁻² | 424.391 | 1.057 | 408.145 | 0.644 |
| Latent (ocean surface) | Wm ⁻² | -129.284 | 0.586 | -123.159 | 0.387 |
| Sensible (ocean surface) | Wm ⁻² | -9.683 | 0.107 | -10.578 | 0.092 |
| Ocean potential temperature | °C | 8.566 | 0.015 | 7.553 | 0.026 |

spondingly, the equatorial thermocline (Fig. 2j) deepens from the eastern end: the 18°C isotherm deepens all the way to ~300 m at the western boundary, whereas the 28°C isotherm reaches maximum depth in the middle of the ocean basin before shoaling again in the west. Note that with active atmosphere, the western warm pool is established at a distance away from the western boundary (~ 1/3 of the basin width), as opposed to immediately against the western boundary in an ocean-only model forced by prescribed, zonally uniform wind.

The fundamental role of the meridional ocean boundary in determining the global climate, as suggested by Figs. 1 and 2, are further analyzed in the subsections below. Contrasting the climates of Aqua and Ridge, we explore the following aspects: the global energy budget, the large-scale circulation with seasonality, and the resulting meridional heat transport.

3.1 Global Energy Budget and Balance

The differences between the global mean climates of Aqua and Ridge is presented in Table 1, which summarizes the statistics of global mean values concerning the energy budget and the water cycle over the annually averaged 100-year period under analysis. In virtually all aspects, the differences between the global mean state of Aqua and Ridge are well beyond the range of their respective interannual variability, as measured by the standard deviation of the global mean.

The warmth of the climate states — with ~27°C global mean surface temperature for Aqua — are comparable to Smith et al. (2006), although greater contrast between Aqua and Ridge is presented here. Aqua is ~2°C warmer in global mean surface temperature and ~1°C warmer in global mean ocean potential temperature compared to Ridge (Table 1). In the energy budget, this corresponds to greater net shortwave heating at TOA, as well as at the ocean surface. The radiative forcing of clouds plays a large role in the cooling of Ridge relative to Aqua: the prominent cloud radiative cooling in the tropics, due to the presence of the western warm pool on Ridge with its convective activities, is reflected in the global mean.

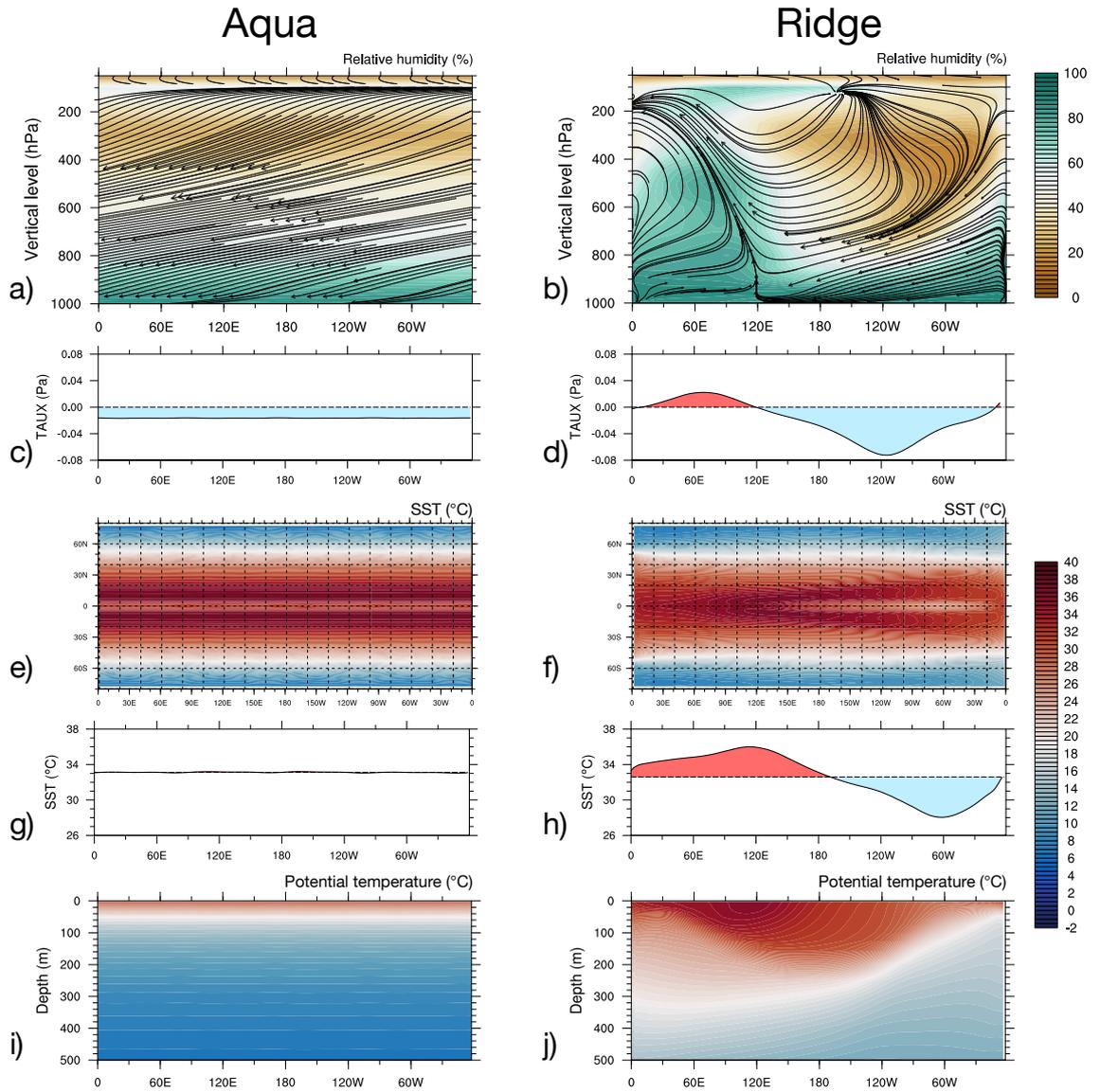


Figure 2. Zonal features in the tropics, 100-yr climatology: (a–b) Zonal circulation in the atmosphere with Walker-like feature on Ridge, seen in relative humidity (colored shading), and streamline of zonal and vertical velocity (solid arrows). Vertical velocity is scaled by a factor of 50 for visualization; (c–d) Zonal gradient of zonal wind stress (Pa), the dashed horizontal line marking zero; (e–f) SST ($^{\circ}\text{C}$); (g–h) Zonal gradient of SST ($^{\circ}\text{C}$), the dashed horizontal line marking the zonal average value; (i–j) Equatorial thermocline, as seen in potential temperature ($^{\circ}\text{C}$). All panels except for (e–f) are averaged 5°N – 5°S .

260 The meridional structure of the energy budget is further detailed in Fig. 3. In the
 261 zonal average of the TOA radiative budget (Fig. 3a–d), both Aqua and Ridge qualita-
 262 tively resemble Earth observations (e.g., Stephens et al., 2015). The extent of the trop-
 263 ics is essentially identical for both planets, with poleward limits at 37.2°N/S as defined
 264 by TOA radiative surplus. In the zonal average, the net tropical heating of Aqua is greater
 265 relative to Ridge at both TOA and the ocean surface. At TOA, Aqua receives more short-
 266 wave (Fig. 3a–b) and integrated net surplus heating (Fig. 3c–d) than Ridge. Over the
 267 ocean surface (Fig. 3e–f), the heating of Aqua relative to Ridge in the deep tropics is mostly
 268 due to greater net shortwave and lesser latent heat loss over the equatorial cold belt (see
 269 Fig. S3). Specifically, the presence of the western warm pool on Ridge (Fig. 2, right col-
 270 umn) reduces surface shortwave flux via cloud forcing, and enhances latent heat loss of
 271 the ocean by greater evaporation associated with its warmer temperature. These effects
 272 are analogous to observed surface heat fluxes in the Pacific (e.g., Grist & Josey, 2003),
 273 where the Eastern Pacific cold tongue is a region of greater ocean heating than the rest
 274 of tropical Pacific. In this sense, these heating effects are expanded to the entire equa-
 275 torial cold belt on Aqua, contributing to its warmer climate.

276 The warmer climate of Aqua reinforces a more intense water cycle than Ridge. In
 277 the global average (Table 1), Aqua’s intensified water cycle relative to Ridge is reflected
 278 in its slightly higher surface pressure due to water vapor pressure, higher total precip-
 279 itable water by 18%, and higher precipitation rate by 4.8% (Table 1). The percentage
 280 of precipitation increase on Aqua relative to Ridge is consistent with the latent heating
 281 of their atmospheres, at a lesser fractional increase than for total precipitable water, as
 282 discussed by Pendergrass and Hartmann (2014). Aqua’s fractional increase of precipi-
 283 tation with regard to global mean surface temperature is also in line with those reported
 284 from CMIP5 warming experiments (Collins et al., 2013). On Aqua, the higher amount
 285 of water vapor – a greenhouse gas – helps to maintain its warm state, as shown in the
 286 dramatic warming by downwelling longwave compared to Ridge (Table 1). Furthermore,
 287 the meridional structures of some relevant fields are shown in Fig. 4, and the zonally av-
 288 eraged vertical structures of moisture and salinity are shown in Fig. 5. In the zonal av-
 289 erage, both planets have two Intertropical Convergence Zones (ITCZs), with Aqua hav-
 290 ing higher peaks in precipitation (Fig. 4b) and moisture (Fig. 5a–b) than Ridge. The re-
 291 sulting patterns of freshwater forcing (Fig. 4h) correspond to near-surface salinity of the
 292 ocean (Fig. 5c–d). It is worth noting that ”double ITCZs” are a common feature of atmosphere-
 293 only aquaplanets with prescribed equatorial thermal maximum (Blackburn et al., 2013;
 294 Medeiros et al., 2016), and the coupled SST patterns of Aqua and Ridge (Fig. 4a) are
 295 perhaps even more conducive to such structures.

296 As defined by the TOA radiative budget in Fig. 3, the boundary of the tropics and
 297 the descending branch of the Hadley cell (see Fig. 8a–b and later discussion) coincides
 298 with many dynamical features in the zonal average (Fig. 4): the peaks in surface pres-
 299 sure (Fig. 4c), the switching of direction of zonal wind stress (Fig. 4e) and peaks in wind
 300 stress curl (Fig. 4f), and the deepening of the mixed layer depth towards higher latitudes
 301 (Fig. 4g). In Fig. 4g, the zonal asymmetry in Ridge’s tropical thermocline (Fig. 2) is re-
 302 sponsible for deeper mixed layer depth in the deep tropics than Aqua. These contrasts
 303 in the circulation pattern are further discussed in the next subsection.

304 3.2 Large-Scale Circulation

305 For both the atmosphere and the ocean, Fig. 6 shows features of the horizontal cir-
 306 culation, while Fig. 7 shows the vertical structures of the zonally averaged zonal flows.

307 For the atmosphere, the impact of ocean geometry is mediated by SST. In the sur-
 308 face pressure field (Fig. 6a–b), compared to Aqua’s zonally uniform belt of subtropical
 309 high, Ridge has more defined centers of subtropical highs over its eastern boundary cur-
 310 rents (see Fig. 6b). In the vertical structure, the contrast between the zonally averaged

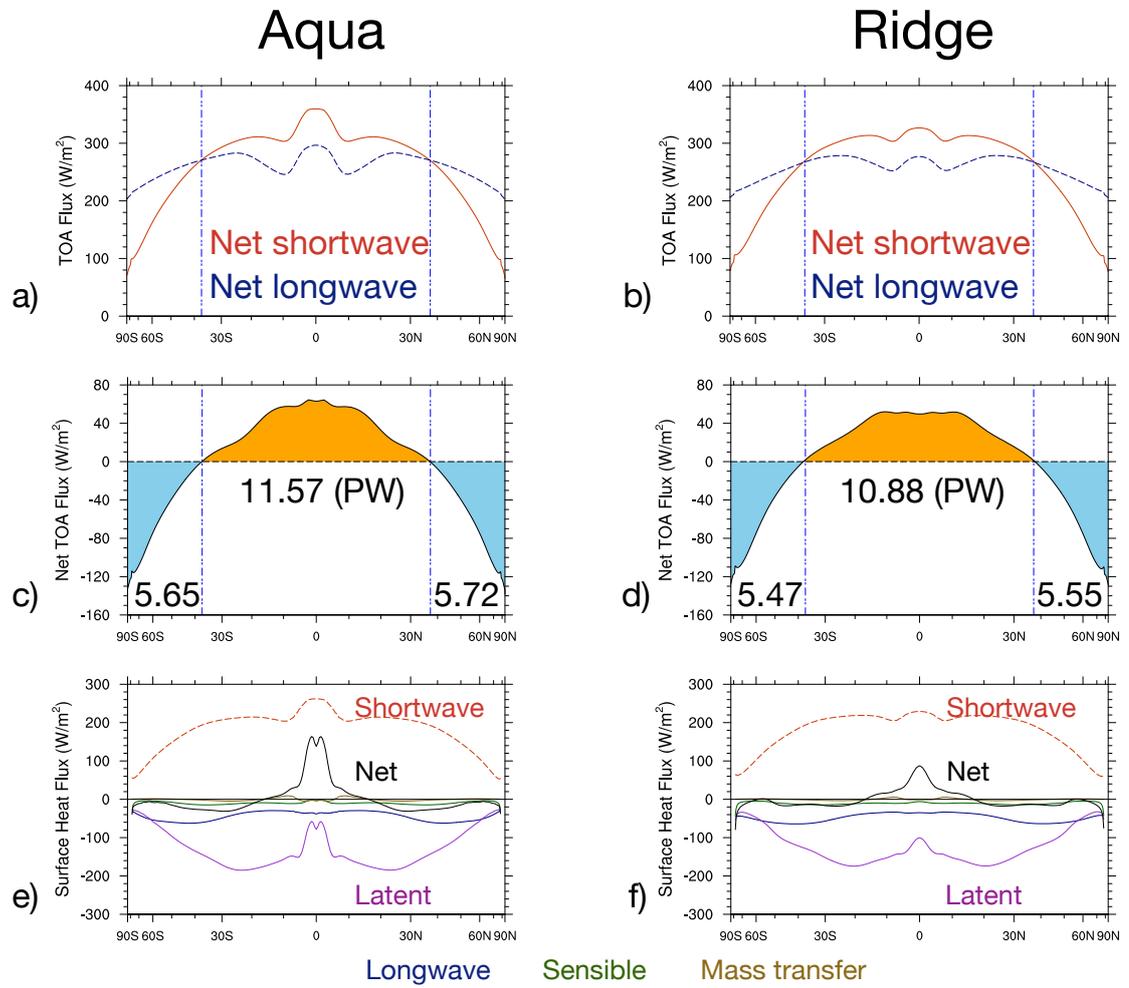


Figure 3. Zonally averaged energy budgets, 100-yr climatology: (a-b) Top-of-atmosphere (TOA) fluxes (Wm^{-2}); (c-d) Net TOA flux (Wm^{-2}) derived from (a-b), labeled with the integrated total amount of tropical surplus (shaded in orange) and extratropical deficit (shaded in blue), in petawatt (PW); (e-f) Ocean surface heat fluxes (Wm^{-2}). The x-axis is scaled by $\sin(\text{lat})$ to reflect the proportion of surface area, with minor tick marks at 10° intervals.

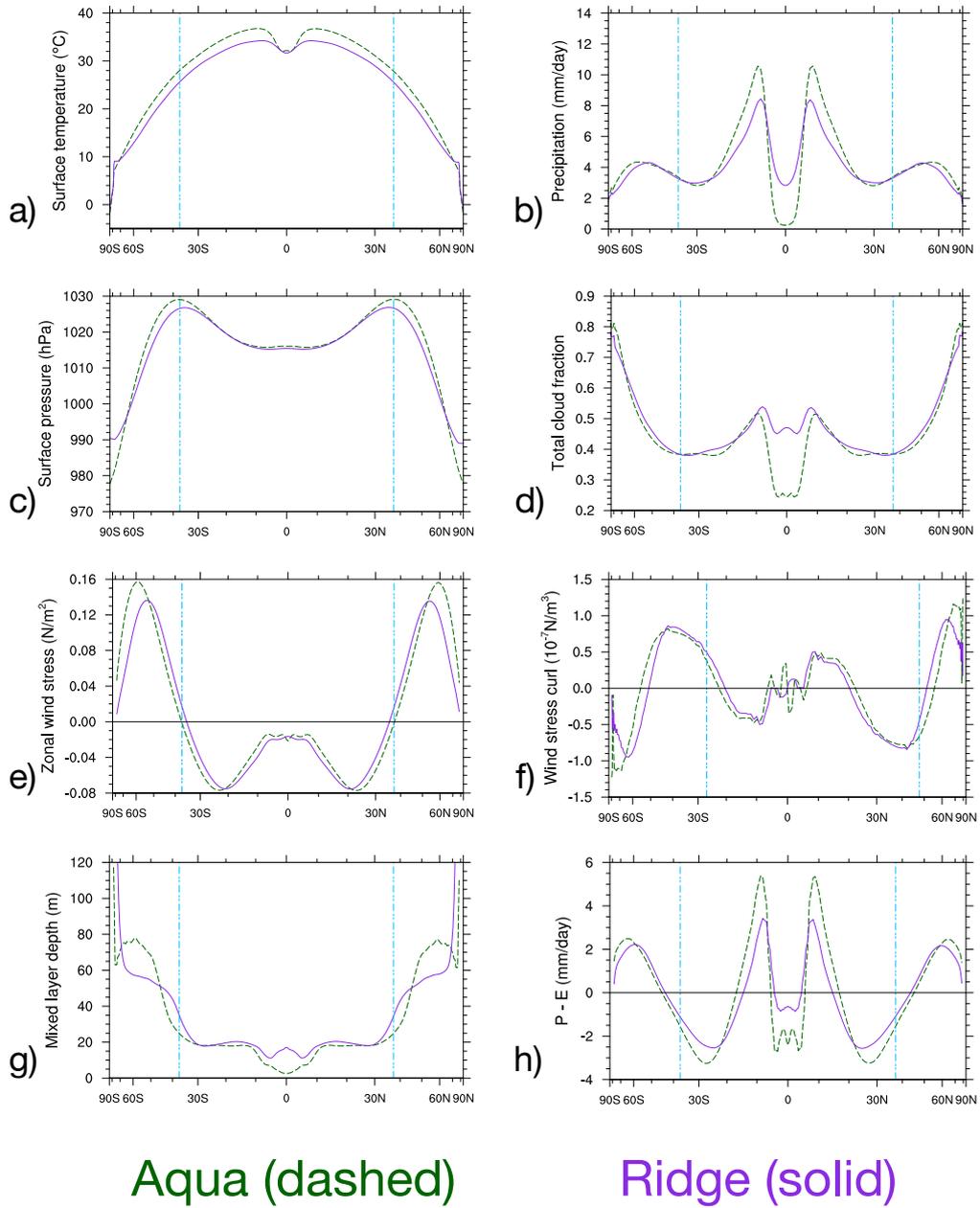


Figure 4. Zonal average profiles, 100-yr climatology: (a) Surface temperature ($^{\circ}\text{C}$); (b) Precipitation rate (mmday^{-1}); (c) Surface pressure (hPa); (d) Total cloud fraction (fraction); (e) Zonal wind stress (Nm^{-2}); (f) Curl of zonal wind stress (10^{-7}Nm^{-3}); (g) Ocean mixed layer depth (m); (h) Precipitation minus evaporation (mmday^{-1}). The vertical blue lines mark the extent of the tropics, as defined by TOA radiative budget (see Fig. 3).

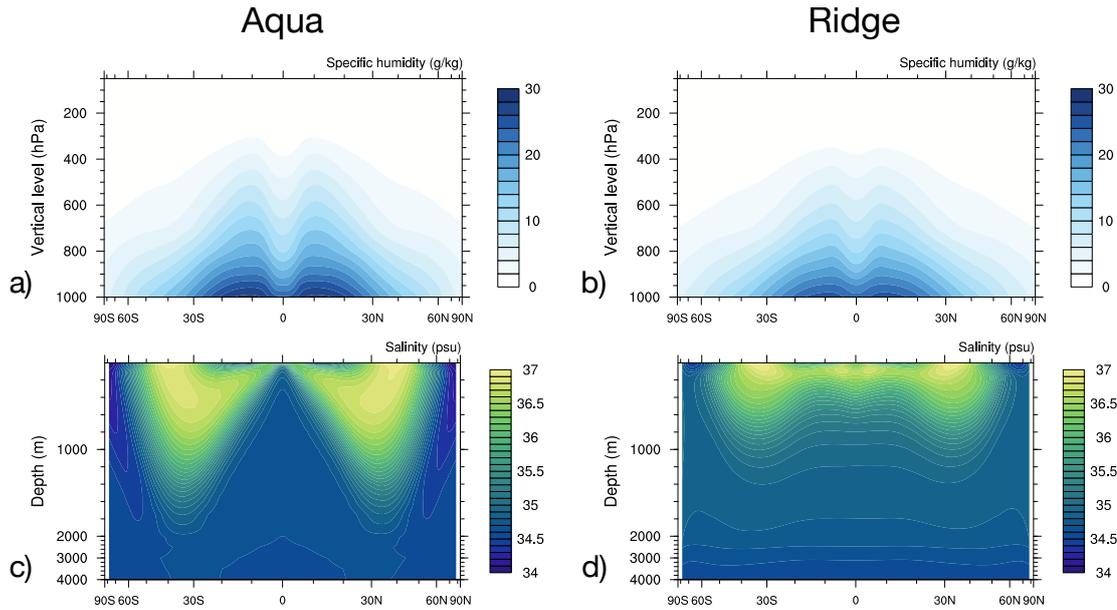


Figure 5. Zonally averaged vertical sections, 100-yr climatology: (a–b) Specific humidity (gkg^{-1}); (c–d) Salinity (psu). For the ocean (c–d), the depths below 2000 m are linearly shrunk as labeled.

311 zonal wind of Aqua and Ridge (Fig. 7a–b) is influenced by their surface temperature gra-
 312 dients (Fig. 4a) through the thermal wind relationship (cf. Enderton & Marshall, 2009).
 313 Due to enhanced ocean heat transport to the extratropics via western boundary currents
 314 (cf. Enderton & Marshall, 2009; Vallis & Farneti, 2009), the meridional gradient of Ridge’s
 315 surface temperature is flattened relative to Aqua (Fig. 4a). Consequently, the greater
 316 surface temperature gradient of Aqua results in greater vertical wind shear, stronger west-
 317 ward flows accumulating upward over the equator, and stronger subtropical and polar
 318 jets in the upper levels (Fig. 7a–b).

319 For the ocean, the defining horizontal circulations – zonal for Aqua and gyral for
 320 Ridge – are shown in Fig. 6c–d. On zonally unbounded Aqua, the rapid zonal flows re-
 321 sult in ~ 1800 Sv of globally integrated net zonal transport. On bounded Ridge, the gy-
 322 ral flows $\sim O(100)$ Sv arise from Sverdrup dynamics, corresponding to the meridional
 323 distribution of surface wind stress (Fig. 4e–f). These gyres suggest analogues of the Pa-
 324 cific’s equatorial counter-currents and the western and eastern boundary systems. Fig. 7c–
 325 d presents the zonally averaged vertical structure of these currents in the zonal direc-
 326 tion. On Aqua (Fig. 7c), the direction of the zonal currents corresponds to the surface
 327 wind stress (Fig. 4c), with velocity dampening towards zero deeper down. Near the sur-
 328 face, the maximum velocity of the westward current reaches 2.29 ms^{-1} . Ridge, in con-
 329 trast, shows richer structure particularly in the tropics, with the presence of equatorial
 330 under- and counter-currents (Fig. 7d). The depth of the equatorial undercurrent at ~ 200
 331 m is consistent with the depth of the equatorial thermocline (Fig. 2i). These features are
 332 absent on Aqua, which cannot maintain zonal pressure gradients in its interior. Ridge’s
 333 maximum velocity, in the near-surface equatorial westward current, is 0.44 m^{-1} , about
 334 an order of magnitude lower than Aqua’s. The effect of Ridge’s meridional boundary is
 335 also seen in the meridional overturning circulation of both the atmosphere and the ocean

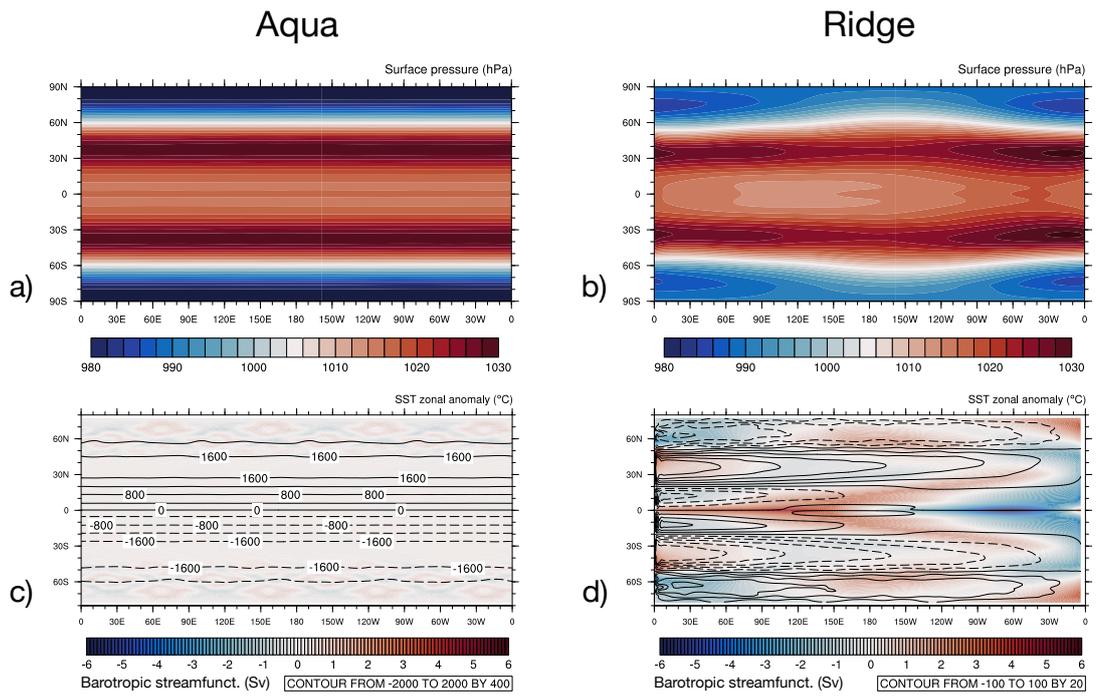


Figure 6. Plan views of 100-yr climatology: (a–b) Atmosphere: surface pressure (hPa); (c–d) Ocean: barotropic streamfunction (Sv, contour lines; solid is positive/clockwise, dashed is negative/counterclockwise), overlaid on the zonal anomaly of SST (°C, color shading). Note the difference in contouring intervals for the streamfunction (c–d). The pattern of Aqua’s SST zonal anomaly (panel c), barely visible, reflects the imprints of bottom topography.

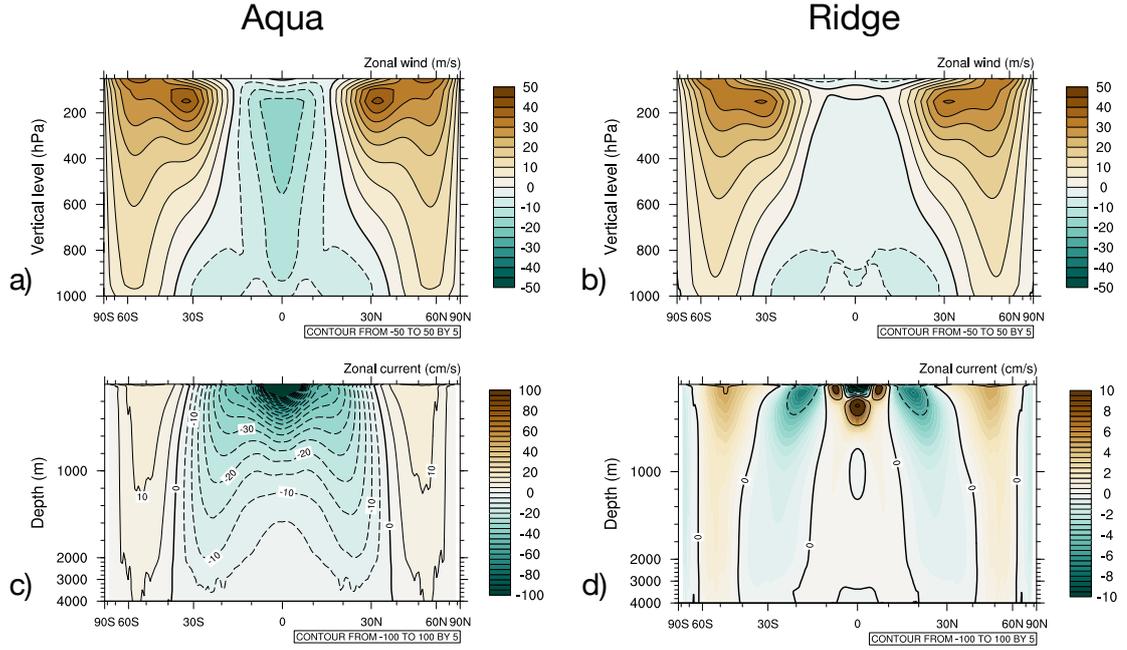


Figure 7. Zonally averaged vertical sections, 100-yr climatology: (a–b) Atmosphere: zonal wind (ms^{-1}); (c–d) Ocean: zonal current (cms^{-1}). Note that the color scale of panel (c) is an order of magnitude greater than panel (d).

336 (Fig. 8). For the atmosphere, in addition to the more familiar-looking overturning cells,
 337 the equatorial cold belt on Aqua leads to the formation of “reverse Hadley” cells in the
 338 deep tropics. On Ridge, this pattern is largely suppressed due to the western warm pool
 339 (Fig. 2, right column) that reduces the meridional gradient around the equatorial SST
 340 minimum in the zonal average (Fig. 4a). For the ocean, Aqua’s residual overturning broadly
 341 follows the isopycnals (Marshall & Radko, 2003; Wolfe & Cessi, 2011), forming deep sub-
 342 tropical cells (Fig. 8c). Alternatively, the residual overturning can be interpreted as the
 343 combination of the Eulerian mean and eddy components (see Fig. S4), where the com-
 344 pensating effect between the two components at high latitudes is analogous to the van-
 345 ishing Deacon cell in the Southern Ocean (cf. Smith et al., 2006; Marshall et al., 2007).
 346 For Ridge, the presence of zonal pressure gradient largely reduces the depth of the sub-
 347 tropical overturning cells. Under the influence of polar convection, the mid-depth (~ 1000
 348 m), diapycnal overturning cells in the midlatitudes are maintained by the balance be-
 349 tween cooling via upwelling and diffusive heating (W. H. Munk, 1966; W. Munk & Wun-
 350 sch, 1998).

351 An intriguing consequence of the “reverse Hadley” cells is observed in the season-
 352 ality of the ITCZs. The effect is most clearly seen in boreal or austral summer. Using
 353 boreal summer (June, July, and August) as an example, Fig. 9 shows the zonally aver-
 354 aged profiles of surface temperature and precipitation, and Fig. 10 shows the correspond-
 355 ing meridional overturning circulation in the atmosphere. For Aqua, despite higher SST
 356 in the summer hemisphere (Fig. 9a), its peak precipitation is in the winter hemisphere
 357 (Fig. 9b). This is a consequence of the persistence of Aqua’s equatorial cold belt throu-
 358 ghout the seasonal cycle, which creates a stand-alone “reverse Hadley” cell in the winter
 359 hemisphere (Fig. 10a) over the local SST minimum (Fig. 9a). The ascending branch of
 360 this overturning cell, at $\sim 10^\circ$ in the winter hemisphere, creates a narrow but extreme

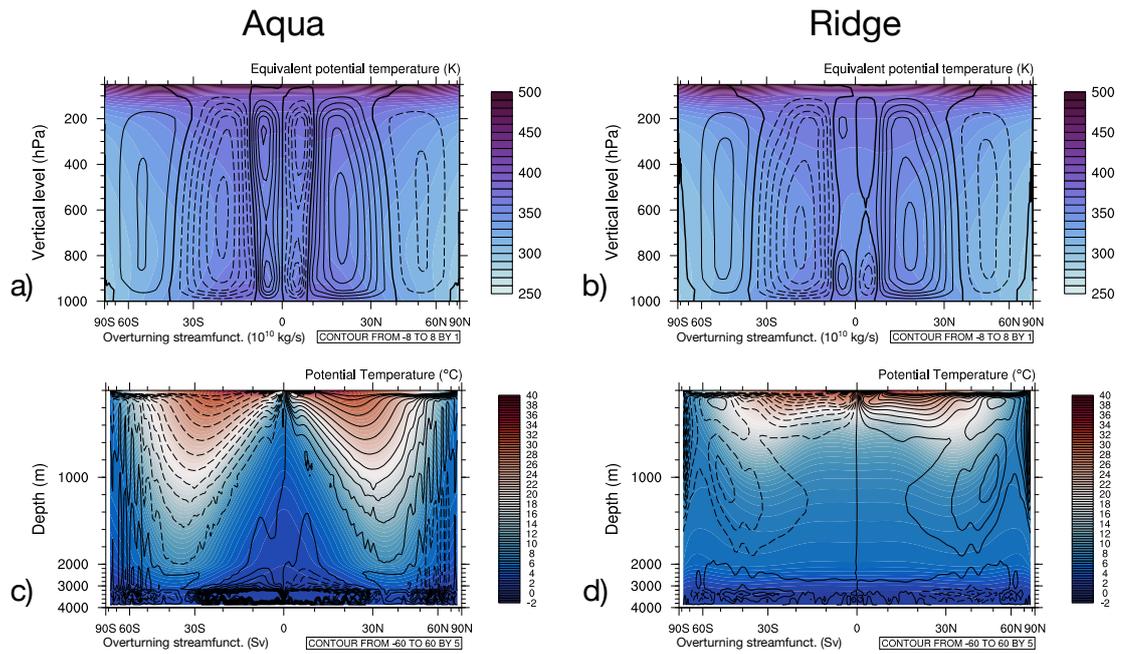


Figure 8. Zonally averaged vertical sections, 100-yr climatology: (a–b) Atmosphere: Eulerian meridional overturning streamfunction (10^{10} kgs^{-1} , contour lines; solid is positive/clockwise, dashed is negative/counterclockwise), overlaid on equivalent potential temperature (K, colored shading); (c–d) Ocean: residual overturning streamfunction (Sv, contour lines; solid is positive/clockwise, dashed is negative/counterclockwise), overlaid on potential temperature ($^{\circ}\text{C}$, colored shading).

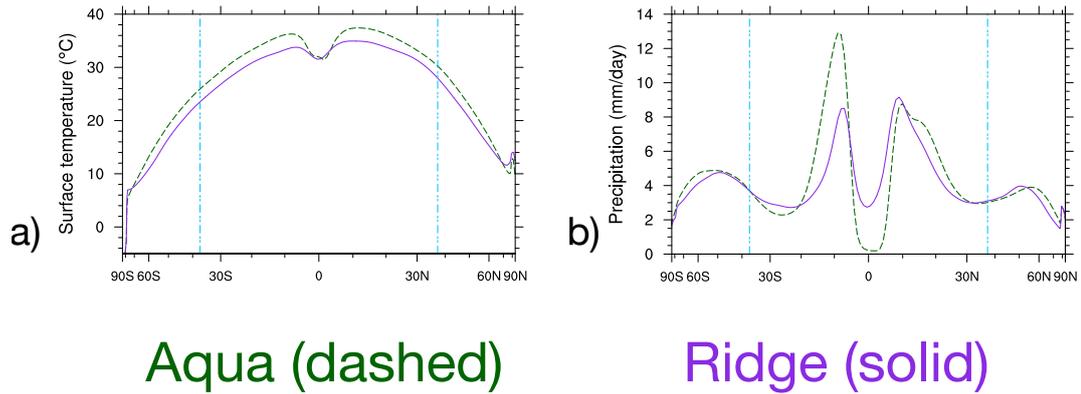


Figure 9. As Fig. 4(a–b), but for boreal summer (June, July, and August)

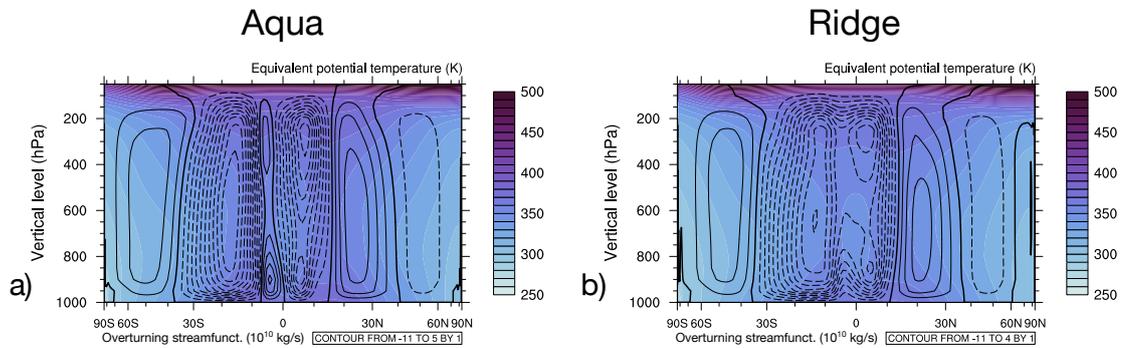


Figure 10. As Fig. 8(a–b), but for boreal summer (June, July, and August)

361 band of maximum precipitation (Fig. 9b) exceeding that of the summer hemisphere. In
 362 this regard, Ridge behaves more Earth-like: in the winter or summer season, although
 363 the presence of the eastern cold tongue manifests itself by affecting the magnitude of the
 364 cross-equatorial Hadley-like overturning (Fig. 10b), dynamically its convection-inducing
 365 effects in the winter flank are much reduced compared to Aqua’s case. For this reason,
 366 the zonally averaged maximum precipitation of Ridge remains in the summer hemisphere
 367 (Fig. 9b). As suggested by the somewhat counter-intuitive seasonal distribution of Aqua’s
 368 atmospheric moisture in Fig. 1c, this is yet another subtle aspect of how ocean geom-
 369 etry governs the state of the coupled climate, including the interaction between the large-
 370 scale circulation and the water cycle.

371 3.3 Meridional Heat Transport

372 The energy budget and circulation patterns of the two planets, as described by the
 373 previous subsections, drive the meridional heat transport (Fig. 11). For both planets,
 374 the total meridional heat transport peaks close to 6 PW in each hemisphere, at the bounds
 375 of the tropics. Ocean heat transport, dominating in the deep tropics, peaks at close to
 376 4 PW for Aqua, and around 2 PW for Ridge. Atmospheric heat transport, dominating

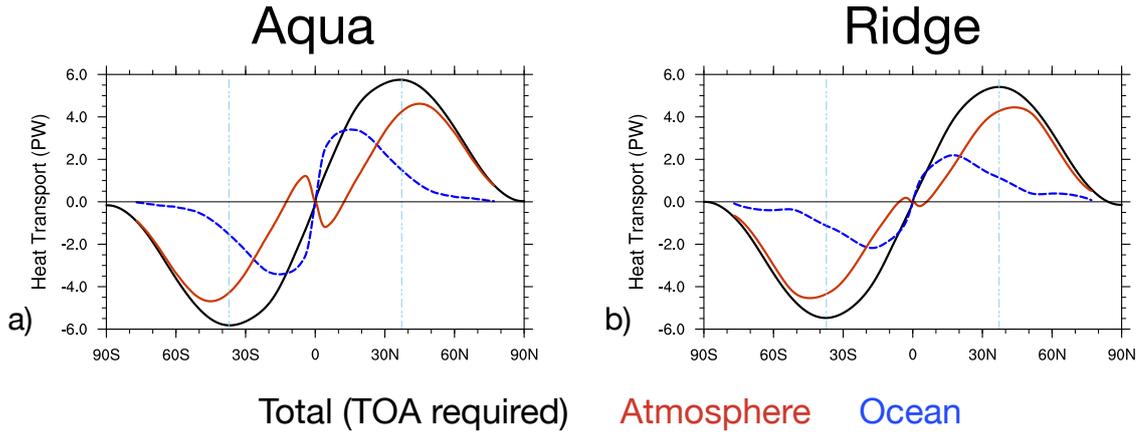


Figure 11. Meridional heat transport, 100-yr climatology: (a) Aqua; (b) Ridge. The vertical blue lines mark the extent of the tropics, as in Fig. 3(a-d) and Fig. 4.

377 in the extratropics, peaks around 50°N/S for both planets. Overall, as discussed by Enderton
 378 and Marshall (2009), the qualitative features and partition between the atmosphere and
 379 the ocean resemble Earth observations (Fasullo & Trenberth, 2008), with Ridge show-
 380 ing greater degrees of realism. Here we discuss the differences between Aqua and Ridge
 381 from energetic and dynamic perspectives (Armour et al., 2019).

382 Energetically, the TOA tropical surplus (Fig. 3c-d) requires greater amounts of to-
 383 tal meridional heat transport for Aqua than Ridge. Likewise, the excessive net heating
 384 of Aqua’s tropical ocean (Fig. 3e) results in greater amounts of ocean heat transport out
 385 of the tropics than Ridge (Fig. 3f). In particular, over the equatorial region, since the
 386 net heating at ocean surface exceeds that of TOA, it is implied that the atmosphere must
 387 compensate by transporting energy equatorward for those regions (Fig. 11).

388 Dynamically, these requirements are fulfilled by the meridional overturning circula-
 389 tion in both fluids (Fig. 8). As detailed in Czaja and Marshall (2006), the meridional
 390 heat transport by either fluid can be viewed as decomposed into two factors: the mag-
 391 nitude of the meridional overturning, and the energy contrast between the poleward and
 392 equatorward branches, as measured by moist static energy for the atmosphere and po-
 393 tential temperature for the ocean. For the atmosphere, the equatorward heat transport
 394 over 10°N/S is delivered by the “reverse Hadley” cells (Fig. 8a-b), which transport higher
 395 amounts of moist static energy in their equatorward upper branches than their poleward
 396 lower branches, at greater magnitude of overturning on Aqua than Ridge. It is worth
 397 noting that the Eulerian mean overturning in Fig. 8a-b only reflects the atmospheric heat
 398 transport by the mean flow, which dominates in the tropics, but gives way to the eddy
 399 component at higher latitudes (cf. Enderton & Marshall, 2009). For the ocean, the en-
 400 ergetically required ocean heat transport is accomplished by the residual overturning (Fig. 8c-
 401 d), where the thermal contrast between the poleward upper branch and equatorward lower
 402 branch is greater on Aqua than Ridge, as the equatorward branch of Aqua’s residual over-
 403 turning reaches near the bottom. On Ridge (Fig. 11b), the “kinks” in ocean heat trans-
 404 port, or local maximum at $\sim 20^{\circ}\text{N/S}$ and local minimum at $\sim 50^{\circ}\text{N/S}$, reflect the bound-
 405 ary of the gyres (Fig. 6d), absent on Aqua.

406 Overall, this comparison highlights the influence of the meridional boundary on merid-
 407 ional heat transport and its partition, via both the energetic requirements and the dy-
 408 namics (Czaja & Marshall, 2006; Enderton & Marshall, 2009). Particularly, in light of

409 having better resolved “reverse Hadley” circulation than earlier investigations (Czaja &
 410 Marshall, 2006; Smith et al., 2006; Farneti & Vallis, 2009) and the corresponding equa-
 411 torward heat transport by the atmosphere, we note the role of Ridge’s meridional bound-
 412 ary in shaping a more Earth-like pattern of meridional heat transport.

413 **4 Conclusions and Discussion**

414 In this study, we introduce the first two examples of fully coupled, idealized mod-
 415 els developed in the CESM Simpler Models framework. Building upon previous ideal-
 416 ized studies using aquaplanets at various degrees of complexity and atmosphere-ocean
 417 coupling, our work explores the coupled climate controlled by ocean geometry, represented
 418 by a meridional boundary present on Ridge and absent on Aqua. By using contempo-
 419 rary atmospheric and ocean model components at resolutions comparable to comprehen-
 420 sive Earth System Models, we aim to apply these idealized models to future studies of
 421 various features in the coupled climate system.

422 Contrasting the mean climates of the CESM Aqua and Ridge planets, the main con-
 423 clusions are summarized as follows:

- 424 1. With sufficient horizontal and vertical resolution, Aqua manifests a global cold belt
 425 of equatorial upwelling, while Ridge develops zonal contrast between its western
 426 warm pool and eastern cold tongue due to boundary dynamics;
- 427 2. Energetically, Aqua’s cold belt results in a climate state $\sim 2^\circ\text{C}$ warmer than Ridge
 428 on global average, due to the effects of tropical clouds and water vapor;
- 429 3. Dynamically, the meridional boundary of Ridge — with the resulting zonal asym-
 430 metry — is crucial for producing a climate system with more Earth-like features
 431 compared to Aqua, including atmospheric and ocean circulation, the seasonality
 432 of ITCZ, and the meridional heat transport.

433 In general, the CESM Aqua and Ridge planets present a number of qualitative fea-
 434 tures similar to those discussed by previous works (Smith et al., 2006; Enderton & Mar-
 435 shall, 2009; Farneti & Vallis, 2009), including the large-scale circulation and meridional
 436 heat transport. We discuss the following aspects of distinction from previous models:

- 437 1. **The climate contrast between Aqua and Ridge and the role of ocean
 438 geometry in planetary albedo.** Contrary to Enderton and Marshall (2009) where
 439 Aqua — with its sea ice — has a colder climate than the ice-free Ridge, in the present
 440 study ice-free CESM Aqua is warmer than Ridge. As discussed in Section 3.1, this
 441 is attributed to the tropical distribution of clouds, which largely dominates the
 442 planetary albedo in the absence of ice. Compared to ice-present climate states of
 443 Enderton and Marshall (2009), the contrast between ice-free Aqua and Ridge sug-
 444 gests a fundamentally different role of the ocean’s meridional boundary on the global
 445 climate: instead of reducing planetary albedo by the melting of sea ice via the ocean’s
 446 western boundary dynamics, the strip continent in CESM Ridge enhances plan-
 447 etary albedo through tropical clouds, via the formation of the western warm pool
 448 and atmospheric convection over it. While the quantitative effect likely depends
 449 on configurations of the atmospheric model including resolution, parameterization
 450 and other aspects affecting the representation of clouds (e.g. apparently minimized
 451 contrast in Smith et al., 2006), the qualitative contrast with Aqua has implica-
 452 tions for the investigation of ice-free warm states in Earth’s history or future. From
 453 a practical standpoint, for alternative applications of the current CESM Aqua and
 454 Ridge models, colder climate states with sea ice — when desired for certain in-
 455 vestigations — can potentially be achieved by parameter tuning in the atmospheric
 456 component.

- 457 **2. Aqua’s equatorial cold belt and the resulting “reverse Hadley” circu-**
 458 **lation.** In CESM Aqua, the atmospheric “reverse Hadley” cells over the equato-
 459 rial belt of upwelling are more distinctively represented than earlier models (Smith
 460 et al., 2006; Marshall et al., 2007; Farneti & Vallis, 2009), providing stronger con-
 461 trast against the corresponding Ridge configuration. While the coupled tropical
 462 dynamics of wind-driven equatorial upwelling and the corresponding atmospheric
 463 “reverse Hadley” cells are relatively straightforward, the representation of these
 464 features largely depends on the horizontal resolution of both model components
 465 for resolving the oceanic belt of upwelling and the narrow (less than 10° in the merid-
 466 ional extent) atmospheric cells. By using model components and resolution com-
 467 parable to that of CMIP, the assessment of these features — in contrast to Ridge
 468 or additional forms of ocean geometry — will have direct relevance to eastern Pa-
 469 cific upwelling and the corresponding regional meridional cells (e.g. Sun et al., 2019)
 470 in realistic, coupled Earth configurations.
- 471 **3. The location and intensity of Ridge’s western warm pool, and the as-**
 472 **sociated Walker circulation.** Compared to earlier Ridge models with a warm
 473 pool closer to the western boundary and a relatively weak zonal SST gradient (Smith
 474 et al., 2006; Enderton & Marshall, 2009; Farneti & Vallis, 2009), CESM Ridge has
 475 a climatological warm pool farther east (distance from the western boundary $\sim 1/3$
 476 of the basin width), and a zonal SST gradient comparable to the Pacific. Besides
 477 ocean dynamics, the roles of cloud forcing and wind stress in the formation of the
 478 warm pool are broadly consistent with some of earlier idealized studies on the West-
 479 ern Pacific (Clement et al., 2005; Watanabe, 2008a, 2008b). The question of con-
 480 trolling factors and mechanisms for the location and intensity of the warm pool
 481 can be a topic of further investigation in this coupled, idealized framework.

482 Furthermore, preliminary analysis on the sub-seasonal to interannual variability
 483 of Aqua and Ridge reveals promising features, including MJO- and ENSO-like modes
 484 on Ridge. These modes of tropical variability, in different forms for Aqua and Ridge with
 485 relevance to the interpretation of realistic Earth configurations, will be addressed in fu-
 486 ture work.

487 To conclude, the climate states of CESM Aqua and Ridge configurations showcase
 488 the capacity of the idealized coupled models to represent relatively well-understood dy-
 489 namics, while further enabling more detailed investigation of the coupled climate sys-
 490 tem. The newly available capacities — including aspects of cloud radiative effects, con-
 491 vection, and circulation — are due to increased resolution and more complete physics
 492 of CMIP-class components. By using CESM components, the close relationship between
 493 these idealized configurations and comprehensive, realistic Earth configurations fills a
 494 long-standing gap in the idealized modeling hierarchy. This addition to the hierarchy opens
 495 up new potential for the investigation of coupled atmosphere-ocean processes, as well as
 496 serving as test beds for model evaluation and development. The Aqua and Ridge con-
 497 figurations presented here are expected to be available in the next major release of CESM
 498 as part of the Simpler Models suite, potentially with the software for creating additional,
 499 customized ocean geometries. Building on the two baseline configurations of Aqua and
 500 Ridge, increasingly complex ocean geometries may be explored (e.g. Ferreira et al., 2010).
 501 Furthermore, with increased atmospheric and/or ocean resolution, the CESM idealized
 502 coupled models can provide insights into an even wider range of features and scale in-
 503 teractions of scientific and societal interests in the coupled climate system.

504 **Acknowledgments**

505 The simulation outputs under analysis are available on CISL’s Globally Accessible Data
 506 Environment. We thank the following collaborators for their help with simplified climate
 507 models (in alphabetical order): Alper Altuntas, Kyle Armour, David Bailey, Jim Bene-
 508 dict, Pedro Di Nezio, Erik Kluzek, Keith Lindsay, Brian Medeiros, Sarah Ragen, Mathew

509 Rothstein, and Andrew Shao. We also thank Anna-Lena Deppenmeier for helpful dis-
 510 cussion on the manuscript. Wu was supported by National Science Foundation (NSF)
 511 grant AGS1648629, the Advanced Study Program of NCAR, and the Junior Researcher
 512 Award of the Institute for Advanced Computational Science at Stony Brook University.
 513 Reed was supported by NSF grants AGS1648629 and AGS1830729. The National Cen-
 514 ter for Atmospheric Research (NCAR) is sponsored by the NSF under Cooperative Agree-
 515 ment 1852977. We acknowledge computing and data storage resources, including the Cheyenne
 516 supercomputer (doi:10.5065/D6RX99HX), provided by the Computational and Infor-
 517 mation Systems Laboratory (CISL) at NCAR.

518 References

- 519 Abernathey, R., Ferreira, D., & Klocker, A. (2013). Diagnostics of isopycnal mixing
 520 in a circumpolar channel. *Ocean Modelling*, *72*, 1–16.
- 521 Adcroft, A., Anderson, W., Balaji, V., Blanton, C., Bushuk, M., Dufour, C. O., ...
 522 others (2019). The GFDL global ocean and sea ice model OM4. 0: Model
 523 description and simulation features. *Journal of Advances in Modeling Earth*
 524 *Systems*, *11*(10), 3167–3211.
- 525 Armour, K. C., Siler, N., Donohoe, A., & Roe, G. H. (2019). Meridional atmospheric
 526 heat transport constrained by energetics and mediated by large-scale diffusion.
 527 *Journal of Climate*, *32*(12), 3655–3680.
- 528 Bachman, S., & Fox-Kemper, B. (2013). Eddy parameterization challenge suite I:
 529 Eady spindown. *Ocean Modelling*, *64*, 12–28.
- 530 Bailey, D., DuVivier, A., Holland, M., Hunke, E., Lipscomb, B., Briegleb, B., ...
 531 Schramm, J. (2018). *CESM CICE5 users guide* (Tech. Rep.). Tech. rep.
- 532 Ballinger, A. P., Merlis, T. M., Held, I. M., & Zhao, M. (2015). The sensitivity
 533 of tropical cyclone activity to off-equatorial thermal forcing in aquaplanet
 534 simulations [Journal Article]. *Journal of the Atmospheric Sciences*, *72*(6),
 535 2286–2302.
- 536 Benedict, J. J., Medeiros, B., Clement, A. C., & Pendergrass, A. G. (2017). Sen-
 537 sitivities of the hydrologic cycle to model physics, grid resolution, and ocean
 538 type in the aquaplanet Community Atmosphere Model. *Journal of Advances in*
 539 *Modeling Earth Systems*, *9*(2), 1307–1324.
- 540 Blackburn, M., Williamson, D. L., Nakajima, K., Ohfuchi, W., Takahashi, Y. O.,
 541 Hayashi, Y.-Y., ... others (2013). The aqua-planet experiment (APE): Control
 542 SST simulation. *Journal of the Meteorological Society of Japan. Ser. II*, *91*,
 543 17–56.
- 544 Bleck, R. (1978). On the use of hybrid vertical coordinates in numerical weather pre-
 545 diction models. *Monthly weather review*, *106*(9), 1233–1244.
- 546 Brunetti, M., Kasparian, J., & V erard, C. (2019). Co-existing climate attractors in a
 547 coupled aquaplanet. *Climate Dynamics*, *53*(9-10), 6293–6308.
- 548 Carranza, M. M., Gille, S. T., Franks, P. J., Johnson, K. S., Pinkel, R., & Girton,
 549 J. B. (2018). When mixed layers are not mixed: Storm-driven mixing and
 550 bio-optical vertical gradients in mixed layers of the southern ocean. *Journal of*
 551 *Geophysical Research: Oceans*, *123*(10), 7264–7289.
- 552 Cessi, P., & Jones, C. S. (2017). Warm-route versus cold-route interbasin exchange
 553 in the meridional overturning circulation. *J. Phys. Oceanogr.*, *47*(8), 1981–
 554 1997.
- 555 Chang, K.-I., Ghil, M., Ide, K., & Lai, C.-C. A. (2001). Transition to aperiodic
 556 variability in a wind-driven double-gyre circulation model. *Journal of physical*
 557 *oceanography*, *31*(5), 1260–1286.
- 558 Chavas, D. R., Reed, K. A., & Knaff, J. A. (2017). Physical understanding of the
 559 tropical cyclone wind-pressure relationship [Journal Article]. *Nature communi-*
 560 *cations*, *8*(1), 1360.

- 561 Clement, A. C., Seager, R., & Murtugudde, R. (2005). Why are there tropical warm
562 pools? *Journal of climate*, *18*(24), 5294–5311.
- 563 Collins, M., Knutti, R., Arblaster, J., Dufresne, J.-L., Fichefet, T., Friedlingstein, P.,
564 ... others (2013). Long-term climate change: projections, commitments and
565 irreversibility. In *Climate change 2013-the physical science basis: Contribution*
566 *of working group I to the fifth assessment report of the intergovernmental panel*
567 *on climate change* (pp. 1029–1136). Cambridge University Press.
- 568 Czaja, A., & Marshall, J. (2006). The partitioning of poleward heat transport be-
569 tween the atmosphere and ocean [Journal Article]. *Journal of the atmospheric*
570 *sciences*, *63*(5), 1498–1511.
- 571 Danabasoglu, G., Lamarque, J.-F., Bacmeister, J., Bailey, D., DuVivier, A., Ed-
572 wards, J., ... others (2020). The Community Earth System Model ver-
573 sion 2 (CESM2). *Journal of Advances in Modeling Earth Systems*, *12*(2),
574 e2019MS001916.
- 575 Donohoe, A., Frierson, D. M., & Battisti, D. S. (2014). The effect of ocean mixed
576 layer depth on climate in slab ocean aquaplanet experiments [Journal Article].
577 *Climate dynamics*, *43*(3-4), 1041–1055.
- 578 Enderton, D., & Marshall, J. (2009). Explorations of atmosphere–ocean–ice climates
579 on an aquaplanet and their meridional energy transports. *Journal of the Atmo-*
580 *spheric Sciences*, *66*(6), 1593–1611.
- 581 Eyring, V., Bony, S., Meehl, G. A., Senior, C. A., Stevens, B., Stouffer, R. J., &
582 Taylor, K. E. (2016). Overview of the Coupled Model Intercomparison Project
583 Phase 6 (CMIP6) experimental design and organization. *Geoscientific Model*
584 *Development (Online)*, *9*(LLNL-JRNL-736881).
- 585 Farneti, R., & Vallis, G. (2009). An Intermediate Complexity Climate Model (IC-
586 CMp1) based on the GFDL flexible modelling system. *Geoscientific Model De-*
587 *velopment*, *2*(2), 73.
- 588 Fasullo, J. T., & Trenberth, K. E. (2008). The annual cycle of the energy budget.
589 part II: Meridional structures and poleward transports. *Journal of Climate*,
590 *21*(10), 2313–2325.
- 591 Ferrari, R., Nadeau, L.-P., Marshall, D. P., Allison, L. C., & Johnson, H. L. (2017).
592 A model of the ocean overturning circulation with two closed basins and a
593 reentrant channel. *J. Phys. Oceanogr.*, *47*(12), 2887–2906.
- 594 Ferreira, D., Marshall, J., & Campin, J.-M. (2010). Localization of deep water for-
595 mation: Role of atmospheric moisture transport and geometrical constraints on
596 ocean circulation [Journal Article]. *Journal of Climate*, *23*(6), 1456–1476.
- 597 Frierson, D. M., Hwang, Y.-T., Fučkar, N. S., Seager, R., Kang, S. M., Donohoe,
598 A., ... Battisti, D. S. (2013). Contribution of ocean overturning circulation
599 to tropical rainfall peak in the northern hemisphere [Journal Article]. *Nature*
600 *Geoscience*, *6*(11), 940–944.
- 601 Gent, P. R., Willebrand, J., McDougall, T. J., & McWilliams, J. C. (1995). Parame-
602 terizing eddy-induced tracer transports in ocean circulation models. *Journal of*
603 *Physical Oceanography*, *25*(4), 463–474.
- 604 Griffies, S. M., Levy, M., Adcroft, A. J., Danabasoglu, G., Hallberg, R. W., Jacob-
605 sen, D., ... Ringler, T. (2015). Theory and numerics of the Community Ocean
606 Vertical Mixing (CVMIX) project. *Tech. Rep.*.
- 607 Grist, J. P., & Josey, S. A. (2003). Inverse analysis adjustment of the SOC air–sea
608 flux climatology using ocean heat transport constraints. *Journal of Climate*,
609 *16*(20), 3274–3295.
- 610 Hack, J. J. (1994). Parameterization of moist convection in the National Center for
611 Atmospheric Research community climate model (CCM2). *Journal of Geophys-*
612 *ical Research: Atmospheres*, *99*(D3), 5551–5568.
- 613 Held, I. M. (2005). The gap between simulation and understanding in climate mod-
614 eling. *Bulletin of the American Meteorological Society*, *86*(11), 1609–1614.
- 615 Herrington, A. R., & Reed, K. A. (2017). An explanation for the sensitivity of the

- 616 mean state of the community atmosphere model to horizontal resolution on
 617 aquaplanets. *Journal of Climate*, 30(13), 4781–4797.
- 618 Hirt, C. W., Amsden, A. A., & Cook, J. (1974). An arbitrary Lagrangian-Eulerian
 619 computing method for all flow speeds. *Journal of computational physics*,
 620 14(3), 227–253.
- 621 Holloway, G. (1997). Eddy transport of thickness and momentum in layer and level
 622 models. *Journal of physical oceanography*, 27(6), 1153–1157.
- 623 Holtslag, A., & Boville, B. (1993). Local versus nonlocal boundary-layer diffusion in
 624 a global climate model. *Journal of Climate*, 6(10), 1825–1842.
- 625 Hurrell, J. W., Holland, M. M., Gent, P. R., Ghan, S., Kay, J. E., Kushner, P. J.,
 626 ... others (2013). The Community Earth System Model: a framework for
 627 collaborative research. *Bulletin of the American Meteorological Society*, 94(9),
 628 1339–1360.
- 629 Jeevanjee, N., Hassanzadeh, P., Hill, S., & Sheshadri, A. (2017). A perspective
 630 on climate model hierarchies. *Journal of Advances in Modeling Earth Systems*,
 631 9(4), 1760–1771.
- 632 Jones, C. S., & Cessi, P. (2016). Interbasin transport of the meridional overturning
 633 circulation. *J. Phys. Oceanogr.*, 46(4), 1157–1169.
- 634 Jones, C. S., & Cessi, P. (2017). Size matters: Another reason why the Atlantic is
 635 saltier than the Pacific. *J. Phys. Oceanogr.*, 47(11), 2843–2859.
- 636 Jones, C. S., & Cessi, P. (2018). Components of upper-ocean salt transport by the
 637 gyres and the meridional overturning circulation. *J. Phys. Oceanogr.*, 48(10),
 638 2445–2456.
- 639 Kaspi, Y., & Schneider, T. (2011). Downstream self-destruction of storm tracks
 640 [Journal Article]. *Journal of the Atmospheric Sciences*, 68(10), 2459–2464.
- 641 Kessler, W. S. (2006). The circulation of the eastern tropical Pacific: A review.
 642 *Progress in Oceanography*, 69(2-4), 181–217.
- 643 Large, W. G., McWilliams, J. C., & Doney, S. C. (1994). Oceanic vertical mixing: A
 644 review and a model with a nonlocal boundary layer parameterization. *Reviews*
 645 *of Geophysics*, 32(4), 363–403.
- 646 Lawrence, D. M., Fisher, R. A., Koven, C. D., Oleson, K. W., Swenson, S. C., Bo-
 647 nan, G., ... others (2019). The community land model version 5: Description
 648 of new features, benchmarking, and impact of forcing uncertainty. *Journal of*
 649 *Advances in Modeling Earth Systems*.
- 650 Li, H., & Sriviver, R. L. (2018). Impact of tropical cyclones on the global ocean:
 651 Results from multi-decadal global ocean simulations isolating tropical cyclone
 652 forcing [Journal Article]. *Journal of Climate*(2018).
- 653 Lin, S.-J., & Rood, R. B. (1996). Multidimensional flux-form semi-Lagrangian trans-
 654 port schemes. *Monthly Weather Review*, 124(9), 2046–2070.
- 655 Lin, S.-J., & Rood, R. B. (1997). An explicit flux-form semi-Lagrangian shallow-
 656 water model on the sphere. *Quarterly Journal of the Royal Meteorological Soci-*
 657 *ety*, 123(544), 2477–2498.
- 658 Maher, P., Gerber, E. P., Medeiros, B., Merlis, T. M., Sherwood, S., Sheshadri, A.,
 659 ... Zurita-Gotor, P. (2019). Model hierarchies for understanding atmospheric
 660 circulation. *Reviews of Geophysics*, 57(2), 250–280.
- 661 Manabe, S., & Bryan, K. (1969). Climate calculations with a combined ocean-
 662 atmosphere model. *Journal of the Atmospheric Sciences*, 26(4), 786–789.
- 663 Marshall, J., Ferreira, D., Campin, J.-M., & Enderton, D. (2007). Mean climate and
 664 variability of the atmosphere and ocean on an aquaplanet [Journal Article].
 665 *Journal of the Atmospheric Sciences*, 64(12), 4270–4286.
- 666 Marshall, J., & Radko, T. (2003). Residual-mean solutions for the antarctic cir-
 667 cumpolar current and its associated overturning circulation. *Journal of Physi-*
 668 *cal Oceanography*, 33(11), 2341–2354.
- 669 Medeiros, B., Williamson, D. L., & Olson, J. G. (2016). Reference aquaplanet cli-
 670 mate in the Community Atmosphere Model, Version 5. *Journal of Advances in*

- 671 *Modeling Earth Systems*, 8(1), 406–424.
- 672 Molteni, F. (2003). Atmospheric simulations using a GCM with simplified phys-
673 ical parametrizations. I: Model climatology and variability in multi-decadal
674 experiments. *Climate Dynamics*, 20(2-3), 175–191.
- 675 Munk, W., & Wunsch, C. (1998). Abyssal recipes II: Energetics of tidal and wind
676 mixing. *Deep-sea research. Part I, Oceanographic research papers*, 45(12),
677 1977–2010.
- 678 Munk, W. H. (1966). Abyssal recipes. *Deep Sea Research and Oceanographic Ab-*
679 *stracts*, 13(4), 707–730.
- 680 Neale, R. B., Chen, C.-C., Gettelman, A., Lauritzen, P. H., Park, S., Williamson,
681 D. L., ... others (2010). Description of the NCAR Community Atmosphere
682 Model (CAM 4.0). *NCAR Tech. Note NCAR/TN-486+ STR*, 1(1), 1–12.
- 683 Neale, R. B., & Hoskins, B. J. (2000). A standard test for AGCMs including their
684 physical parametrizations: I: The proposal. *Atmospheric Science Letters*, 1(2),
685 101–107.
- 686 Nilsson, J., Langen, P. L., Ferreira, D., & Marshall, J. (2013). Ocean basin geometry
687 and the salinification of the atlantic ocean. *Journal of Climate*, 26(16), 6163–
688 6184.
- 689 Pendergrass, A. G., & Hartmann, D. L. (2014). The atmospheric energy constraint
690 on global-mean precipitation change. *Journal of climate*, 27(2), 757–768.
- 691 Polvani, L., Clement, A., Medeiros, B., Benedict, J., & Simpson, I. (2017). When
692 less is more: Opening the door to simpler climate models, *Eos*, 98. *Eos, Trans-*
693 *actions American Geophysical Union*, 99(3), 15–16.
- 694 Redi, M. H. (1982). Oceanic isopycnal mixing by coordinate rotation. *Journal of*
695 *Physical Oceanography*, 12(10), 1154–1158.
- 696 Reed, K. A., & Jablonowski, C. (2012). Idealized tropical cyclone simulations of in-
697 termediate complexity: A test case for AGCMs. *Journal of Advances in Model-*
698 *ing Earth Systems*, 4(2).
- 699 Schultz, D. M., Fairman Jr, J. G., Anderson, S., & Gardner, S. (2017). Build your
700 own earth: A web-based tool for exploring climate model output in teaching
701 and research. *Bulletin of the American Meteorological Society*, 98(8), 1617–
702 1623.
- 703 Scoccimarro, E., Fogli, P. G., Reed, K. A., Gualdi, S., Masina, S., & Navarra, A.
704 (2017). Tropical cyclone interaction with the ocean: The role of high-frequency
705 (subdaily) coupled processes [Journal Article]. *Journal of Climate*, 30(1),
706 145–162.
- 707 Smith, R. S., Dubois, C., & Marotzke, J. (2006). Global climate and ocean circula-
708 tion on an aquaplanet ocean–atmosphere general circulation model. *Journal of*
709 *climate*, 19(18), 4719–4737.
- 710 Solomon, H. (1971). On the representation of isentropic mixing in ocean circulation
711 models. *Journal of Physical Oceanography*, 1(3), 233–234.
- 712 Stephens, G. L., O’Brien, D., Webster, P. J., Pilewski, P., Kato, S., & Li, J.-l.
713 (2015). The albedo of earth. *Reviews of geophysics*, 53(1), 141–163.
- 714 Sun, Y., Li, L. Z., Ramstein, G., Zhou, T., Tan, N., Kageyama, M., & Wang, S.
715 (2019). Regional meridional cells governing the interannual variability of the
716 hadley circulation in boreal winter. *Climate dynamics*, 52(1-2), 831–853.
- 717 Vallis, G. K., & Farneti, R. (2009). Meridional energy transport in the coupled
718 atmosphere–ocean system: Scaling and numerical experiments. *Quarterly Jour-*
719 *nal of the Royal Meteorological Society: A journal of the atmospheric sciences,*
720 *applied meteorology and physical oceanography*, 135(644), 1643–1660.
- 721 Voigt, A., & Shaw, T. A. (2015). Circulation response to warming shaped by radi-
722 ative changes of clouds and water vapour. *Nature Geoscience*, 8(2), 102–106.
- 723 Watanabe, M. (2008a). Two regimes of the equatorial warm pool. part I: A simple
724 tropical climate model. *Journal of climate*, 21(14), 3533–3544.

- 725 Watanabe, M. (2008b). Two regimes of the equatorial warm pool. part II: Hybrid
726 coupled GCM experiments. *Journal of climate*, *21*(14), 3545–3560.
- 727 Wolfe, C. L., & Cessi, P. (2010). What sets the strength of the middepth stratifica-
728 tion and overturning circulation in eddying ocean models? *Journal of Physical*
729 *Oceanography*, *40*(7), 1520–1538.
- 730 Wolfe, C. L., & Cessi, P. (2011). The adiabatic pole-to-pole overturning circulation.
731 *Journal of Physical Oceanography*, *41*(9), 1795–1810.
- 732 Wyrski, K. (1981). An estimate of equatorial upwelling in the pacific. *Journal of*
733 *Physical Oceanography*, *11*(9), 1205–1214.
- 734 Zhang, G. J., & McFarlane, N. A. (1995). Sensitivity of climate simulations to
735 the parameterization of cumulus convection in the Canadian Climate Centre
736 general circulation model. *Atmosphere-ocean*, *33*(3), 407–446.