

Quantifying the Direct Radiative Effect of Stratospheric Aerosols Using Radiative Kernels

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Abstract

To facilitate the quantification of the stratospheric aerosol radiative effect, this study generates a set of aerosol direct radiative effect (ADRE) kernels based on MERRA-2 reanalysis data. These radiative kernels measure the sensitivities of ADRE to perturbations in scattering and absorbing aerosol optical depth (AOD), respectively. Both broadband and band-by-band radiative kernels are developed to account for the wavelength dependency of ADRE. The broadband kernels are then emulated by a multivariate regression model, which predicts the kernel values from a handful of predictors, including the top-of-atmosphere (TOA) insolation, TOA reflectance, and stratospheric AOD. These kernels offer an efficient and versatile way to assess the ADRE of stratospheric aerosols. The ADREs of the 2022 Hunga volcano eruption and the 2020 Australia wildfire are estimated from the kernels and validated against radiative transfer model-calculated results. The Hunga eruption induced a global mean cooling forcing of -0.46 W/m^2 throughout 2022, while the Australia wildfire caused a warming forcing of $+0.28 \text{ W/m}^2$ from January to August. The kernel estimation can capture over 90% of the ADRE variance with relative error within 10%, in these assessments. The results demonstrate the spectral dependencies of stratospheric ADRE and highlight the distinct radiative sensitivity of stratospheric aerosols, which differs significantly from that of tropospheric aerosols.

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Key Points:

- A global dataset of radiative sensitivity kernels is developed to quantify stratospheric aerosol direct radiative effect (ADRE).
- An analytical model is developed to emulate the kernel values from a handful of predictor variables.
- The stratospheric aerosol kernels capture the spatiotemporally varying ADRE values of volcanic eruptions and wildfire events well.

16 **Abstract**

17 To facilitate the quantification of the stratospheric aerosol radiative effect, this study generates a
18 set of aerosol direct radiative effect (ADRE) kernels based on MERRA-2 reanalysis data. These
19 radiative kernels measure the sensitivities of ADRE to perturbations in scattering and absorbing
20 aerosol optical depth (AOD), respectively. Both broadband and band-by-band radiative kernels
21 are developed to account for the wavelength dependency of ADRE. The broadband kernels are
22 then emulated by a multivariate regression model, which predicts the kernel values from a
23 handful of predictors, including the top-of-atmosphere (TOA) insolation, TOA reflectance, and
24 stratospheric AOD. These kernels offer an efficient and versatile way to assess the ADRE of
25 stratospheric aerosols. The ADREs of the 2022 Hunga volcano eruption and the 2020 Australia
26 wildfire are estimated from the kernels and validated against radiative transfer model-calculated
27 results. The Hunga eruption induced a global mean cooling forcing of -0.46 W/m^2 throughout
28 2022, while the Australia wildfire caused a warming forcing of $+0.28 \text{ W/m}^2$ from January to
29 August. The kernel estimation can capture over 90% of the ADRE variance with relative error
30 within 10%, in these assessments. The results demonstrate the spectral dependencies of
31 stratospheric ADRE and highlight the distinct radiative sensitivity of stratospheric aerosols,
32 which differs significantly from that of tropospheric aerosols.

33 **Plain Language Summary**

34 Stratospheric aerosols influence the Earth's energy balance by scattering and absorbing solar
35 radiation, making it crucial to accurately measure their radiative impact. However, quantifying
36 the aerosol radiative impact is computationally expensive if using radiative transfer models. In
37 this work, we develop a set of aerosol radiative kernels, which can provide a flexible and
38 efficient means for calculating the radiative effects of stratospheric aerosols. The kernels have
39 been demonstrated to effectively quantify the radiative impacts of stratospheric aerosols resulting
40 from wildfire and volcanic eruption events.

42 **1 Introduction**

43 Stratospheric aerosols influence the Earth's radiative energy budget and have profound
44 climate impacts (Kremser et al., 2016). The largest contributor to stratospheric aerosols is
45 volcanic eruptions, which can inject a mixture of sulfur dioxide, sulfuric acid, and water directly
46 into the stratosphere, where they transform into stratospheric aerosols (Martinsson et al., 2019).
47 By increasing the reflection of solar radiation, those volcanic aerosols exert a negative radiative
48 forcing at the top-of-the-atmosphere (TOA), which can lead to pronounced surface cooling and
49 changes in atmospheric circulation and water cycle (Robock, 2000; Grinsted et al., 2007; Wu et
50 al., 2023; Günther et al., 2024). Apart from volcanic eruptions, wildfires-induced
51 pyrocumulonimbus (PyroCb) events can also transport a significant amount of carbonaceous
52 aerosols into the lower stratosphere (Fromm et al., 2010; Ohneiser et al., 2020; Liu et al., 2022;
53 Damany-Pearce et al., 2022). Observation and model studies suggest that the absorptivity of
54 biomass-burning aerosols can warm the stratosphere, deplete the stratospheric ozone, and modify
55 vertical dynamics and horizontal dispersion (Damany-Pearce et al., 2022; Ohneiser et al., 2020,
56 2023).

57 Although the importance of stratospheric aerosols is well recognized, the quantification
58 of their radiative effect has not been an easy task, as it requires the consideration of multiple
59 factors, including aerosol types, height and size distributions, as well as the environmental
60 factors at their locations (Weisenstein et al., 2015; MacMartin et al., 2017; Q.-R. Yu et al., 2019;
61 Visioni et al., 2020; P. Yu et al., 2023; Q. Yu et al., 2024). The straightforward and most

62 accurate way to quantify the aerosol direct radiative effect (ADRE) is the Partial Radiative
63 Perturbation (PRP) method, which requires running a radiative transfer model and differencing
64 the modeled radiative fluxes with and without aerosol perturbations, although this quantification
65 method is computationally expensive. Many studies used alternative approaches to estimate
66 ADRE, for example, by using an analytical relationship between the aerosol optical depth (AOD)
67 and the radiative effect. Hansen et al. (2005) estimated a radiative sensitivity of -22 W/m^2 per
68 unit AOD change, based on the simulation of the Pinatubo eruption case using a global climate
69 model. P. Yu et al. (2023) reported a similar scaling relation for stratospheric aerosols also based
70 on modeling experiments. (Schoeberl et al., 2023, 2024a) applied the radiative sensitivity kernels
71 of Q. Yu & Huang (2023b) to evaluate the climate impacts of the 2022 Hunga volcano eruption.
72 However, these kernels were derived based on the aerosol perturbations in the whole
73 atmospheric column, which is dominated by tropospheric, as opposed to stratospheric aerosols.

74 To the best of our knowledge, a global dataset of radiative sensitivity kernels specifically
75 developed for assessing the ADRE of stratospheric aerosols is still lacking. The existing global
76 aerosol kernels, including those of Q. Yu & Huang, (2023b) and Thorsen et al. (2020), were
77 developed with a focus on tropospheric aerosols, whose radiative sensitivity, as shown later in
78 this paper, differ markedly from stratospheric aerosols. A recent study by Gao et al. (2023) tested
79 the kernel quantification of the ADRE of tropopause aerosols, although the development was
80 limited to the East Asia region. A global kernel dataset, which can facilitate an efficient yet
81 accurate quantification of the spatiotemporally varying radiative impacts of stratospheric
82 aerosols, is expected to have a broad spectrum of applications. This is especially relevant given
83 the frequent occurrence of wildfires (Damany-Pearce et al., 2022), recent volcanic eruptions
84 (Taha et al., 2022), and the increasing discussions about stratospheric aerosol geoengineering
85 (Visioni et al., 2020).

86 It is well recognized that the aerosol optical properties, radiative transfer, and the
87 resulting aerosol radiative effects, all have a strong spectral dependence. For example, the
88 spectral dependence of AOD is often approximated using the Angstrom relationship (Ångström,
89 1929), although the Angstrom exponent (AE) may vary with wavelength (Schuster et al., 2006)
90 and height (Chen et al., 2020). Incorporating spectrally measured aerosol optical properties can
91 reduce uncertainty in the ADRE quantification (Chauvigné et al., 2021). Thorsen et al. (2020)
92 found that distinguishing column-integrated aerosol optical properties in the mid-visible and
93 near-infrared wavelengths can help constrain ADRE, pointing to the potential benefits of
94 developing band-by-band kernels. In addition, spectral kernels may take advantage of the
95 spectral AOD information, which is available from many state-of-the-art climate models as well
96 as satellite and ground-based measurements. Therefore, in addition to a set of broadband
97 stratospheric aerosol kernels, we also aim to produce an accompanying set of spectrally
98 decomposed, band-by-band kernels, to facilitate the use of spectral information in the ADRE
99 quantification.

100 Observational and modeling studies have shown that ADRE sensitivity is strongly
101 influenced by environmental conditions such as clouds, relative humidity, and surface albedo
102 (McComiskey et al., 2008; Loeb et al., 2019; Schoeberl et al., 2023; Q. Yu & Huang, 2023a,
103 2023b). However, the primary environmental factors affecting stratospheric ADRE sensitivity
104 and their underlying physics remain to be elucidated. Another objective of our study is to
105 investigate this environmental dependence. Integrating a physical model with statistical analyses,
106 we experiment with sorting the global aerosol kernels, which are conventionally computed on
107 geographic grids (latitude, longitude, and calendar month), based on the geophysical variables

108 that govern the kernel values according to radiative transfer physics. We aim to establish an
 109 analytical equation to capture the spatiotemporal variations of the kernel values. Such an
 110 analytical relation can be considered a physical (as opposed to geographical) kernel dataset and
 111 can be used for the ADRE quantification under arbitrary situations regardless of the geographic
 112 location, which potentially makes the kernels suitable for broader applications.

113 In summary, in this study, we aim to develop a set of radiative sensitivity kernels that are
 114 specifically designed for quantifying the stratospheric ADRE. The kernels developed here
 115 include both broadband and spectral band-by-band TOA flux kernels provided on conventional
 116 latitude-longitude-month grids, as well as physically sorted broadband kernels whose values are
 117 determined from analytical equations. The structure of this paper is as follows. Section 2 details
 118 the methods used to calculate both broadband and band-by-band kernels. Section 3 describes the
 119 development of physically sorted kernels. These aerosol kernels constitute a versatile means to
 120 quantify the stratospheric ADRE. We demonstrate the use and performance of these kernels by
 121 applying them to two cases: the 2022 Hunga volcanic eruption (Bourassa et al., 2023; Kloss et al.,
 122 2022; Taha et al., 2022) and the 2020 Australia wildfire (Ohneiser et al., 2020; Damany-Pearce
 123 et al., 2022; Sellitto et al., 2022) in Section 4. A summary is provided in Section 5.

124 **2 Data and Methods**

125 2.1 Stratospheric Aerosol Direct Radiative Effect

126 The stratospheric ADRE is calculated as the difference in net radiative fluxes at TOA
 127 with and without stratospheric aerosols:

$$128 \quad \text{Stratos ADRE} = F^{net}(\text{all aerosols}) - F^{net}(\text{no stratos aerosols}) \quad (1)$$

129 where $F^{net} = F^{\downarrow} - F^{\uparrow}$, with the downward flux F^{\downarrow} being positive.

130 In this study, we focus on the shortwave stratospheric ADRE at the TOA under the all-
 131 sky condition. This is because the longwave ADRE is orders of magnitude smaller (Reddy et al.,
 132 2005; Heald et al., 2014; Balmes & Fu, 2021), and aerosol scattering is often neglected in the
 133 longwave schemes of radiative transfer models (Mlawer et al., 1997, 2016), despite stratospheric
 134 aerosols being primarily scattering particles. However, the method described here can also be
 135 used to calculate aerosol kernels in the longwave spectrum, at the surface, or for atmospheric
 136 heating rate. Radiative fluxes are computed using the Rapid Radiative Transfer Model (RRTMG)
 137 (Mlawer et al., 1997, 2016). The required inputs for these calculations are obtained from the
 138 Modern-Era Retrospective Analysis for Research and Applications, Version 2 (MERRA-2)
 139 dataset (Gelaro et al., 2017). We use instantaneous atmospheric and cloud profiles, including air
 140 temperature and pressure, surface temperature, surface albedo, water vapor, ozone, specific
 141 humidity, cloud fraction, and the mass fraction of cloud liquid and ice water. The tropopause is
 142 defined according to the criterion of the World Meteorological Organization (WMO, 1957) as
 143 the lowest level where the temperature lapse rate decreases to 2 K/km or less, and the average
 144 lapse rate from this level to any level within the next 2 km does not exceed 2 K/km.

145 Aerosol optical properties are calculated based on the MERRA-2's instantaneous aerosol
 146 mixing ratio profiles, which include 72 layers. MERRA-2 provides 15 externally mixed aerosol
 147 tracers: hydrophobic and hydrophilic black and organic carbon, sulfate, dust (five size bins), and

148 sea salt (five size bins) (Randles et al., 2017). Aerosol optical properties vary with relative
149 humidity to account for hygroscopic growth. For computational efficiency, the 3-hourly
150 MERRA-2 inputs are resampled into a $2.5^\circ \times 2.5^\circ$ grid box. Our goal is to replicate the aerosol
151 radiative transfer calculations from the MERRA-2 dataset and isolate the impact of stratospheric
152 aerosols to study the stratospheric ADRE. Validations of aerosol optical property inputs and total
153 ADRE calculations against MERRA-2 diagnostic aerosol and radiation products are provided in
154 the Supporting Information (Figures S1-S3).

155 2.2 Computation of Stratospheric Aerosol Kernels

156 Aerosol radiative kernels ($\frac{\partial(ADRE)}{\partial x}$) are the partial derivative of the ADRE to an aerosol-
157 related property x such as AOD and single scattering albedo. These kernels represent how
158 ADRE responds to atmospheric aerosol perturbations. By multiplying the radiative kernels with
159 the changes in x , we can approximate the resulting change in ADRE, which provides a
160 convenient means for estimating the radiative impact of aerosols.

161 In this study, we develop kernels for both stratospheric scattering aerosol optical depth
162 (AOD_{scat}) and absorbing aerosol optical depth (AOD_{abs}). For each type of kernel, radiative
163 transfer calculations are performed twice: one with background aerosols and one with
164 perturbations in the stratospheric aerosols. The sizes of the perturbation are 0.1 for AOD_{scat} and
165 0.01 for AOD_{abs} at 550 nm. The perturbation magnitude differs between stratospheric AOD_{scat}
166 and AOD_{abs} due to the smaller background stratospheric AOD_{abs} compared to AOD_{scat} . We use
167 absolute perturbation values instead of relative ones (such as 1%) to minimize noise from
168 numerical errors caused by very small background AOD values. We have verified that the
169 radiative flux changes respond linearly to the AOD perturbations within the typical magnitudes
170 of stratospheric aerosol perturbations (ΔAOD_{scat} ranging from 10^{-3} to 1 and ΔAOD_{abs} ranging
171 from 10^{-4} to 1, respectively). The sum of the AOD_{scat} and AOD_{abs} effects can also be linearly
172 added to determine the total stratospheric ADRE. Sensitivity tests have also been conducted to
173 determine the impacts of perturbation height on the aerosol kernels. Results indicate minimal
174 difference between perturbing a single layer at random altitudes versus the entire stratosphere.
175 Therefore, for our perturbation runs, we assume a conserved vertical profile shape of
176 stratospheric aerosols. Details about sensitivity tests of linear scaling, linear additivity, and
177 height dependency of stratospheric aerosol kernels are provided in the Supporting Information
178 (Figures S4-S6).

179 The perturbation computations produce both broadband and band-by-band stratospheric
180 aerosol kernels. To account for the diurnal cycle, the 3-hourly kernels are averaged into monthly
181 mean values. These aerosol kernels are computed for an El Niño-Southern Oscillation (ENSO)
182 neutral year, 2022. The impact of interannual variability on aerosol kernels is small, as
183 demonstrated by the comparisons of monthly mean kernels between 2020 and 2022. The R-
184 squared values and Root Mean Squared Errors (RMSE) between the monthly mean AOD_{scat}
185 kernels in those two years are 92% and 3.62, respectively, while for the AOD_{abs} kernels, they are
186 97% and 46.90, respectively. Detailed comparisons are provided in the Supporting Information
187 (Figure S7), showing consistency in both spatial distributions and global mean values.
188

189 2.2.1 Broadband Aerosol Kernels

190 Given that solar energy peaks in the mid-visible bands and that aerosol optical properties
 191 are commonly observed in this range, we use the 550 nm AOD as the perturbation variable. The
 192 stratospheric AOD_{scat} and AOD_{abs} kernels are defined as follows:

$$193 \frac{\partial(ADRE)}{\partial(AOD_{scat})} = \frac{F^{net}(AOD_{scat}+\Delta AOD_{scat}, AOD_{abs}, SSA', g') - F^{net}(AOD_{scat}, AOD_{abs}, SSA, g)}{\Delta AOD_{scat}^{550}} \quad (2)$$

$$194 \frac{\partial(ADRE)}{\partial(AOD_{abs})} = \frac{F^{net}(AOD_{scat}, AOD_{abs}+\Delta AOD_{abs}, SSA', g') - F^{net}(AOD_{scat}, AOD_{abs}, SSA, g)}{\Delta AOD_{abs}^{550}} \quad (3)$$

195 In the unperturbed runs, the background aerosol profiles of AOD, single scattering albedo
 196 (SSA), and asymmetry factor (g) are taken from reconstructed MERRA-2 aerosol optical
 197 property profiles. In the perturbation runs, an aerosol layer representing the stratospheric aerosol
 198 perturbations is added to the background aerosol profile. This added aerosol layer has the
 199 scattering or absorbing AOD values at 550 nm of 0.1 and 0.01, respectively, and the incremental
 200 AOD values (ΔAOD_{scat} and ΔAOD_{abs}) at other wavelengths are prescribed according to the
 201 Angstrom relationship (with the AE being 1). For the scattering AOD perturbation, the SSA and
 202 g values of this added layer are assumed to be 1 and 0.7. The g value is based on annual and
 203 global mean asymmetry factor values reported by Ayash et al. (2008) as well as the background
 204 upper troposphere and lower stratosphere aerosol configurations in Sellitto et al. (2022).
 205 Weighted averaging is used to calculate the values of these aerosol properties in the perturbation
 206 runs. For AOD_{scat} perturbation runs,

$$207 SSA' = \frac{SSA^{550} AOD^{550} + 1 * \Delta AOD_{scat}^{550}}{AOD^{550} + \Delta AOD_{scat}^{550}} \quad (4)$$

$$208 g' = \frac{g^{550} SSA^{550} AOD^{550} + 0.7 * 1 * \Delta AOD_{scat}^{550}}{SSA^{550} AOD^{550} + 1 * \Delta AOD_{scat}^{550}} \quad (5)$$

209 For AOD_{abs} perturbation runs, the SSA and g values are

$$210 SSA' = \frac{SSA^{550} AOD^{550} + 0 * \Delta AOD_{abs}^{550}}{AOD^{550} + \Delta AOD_{abs}^{550}} \quad (6)$$

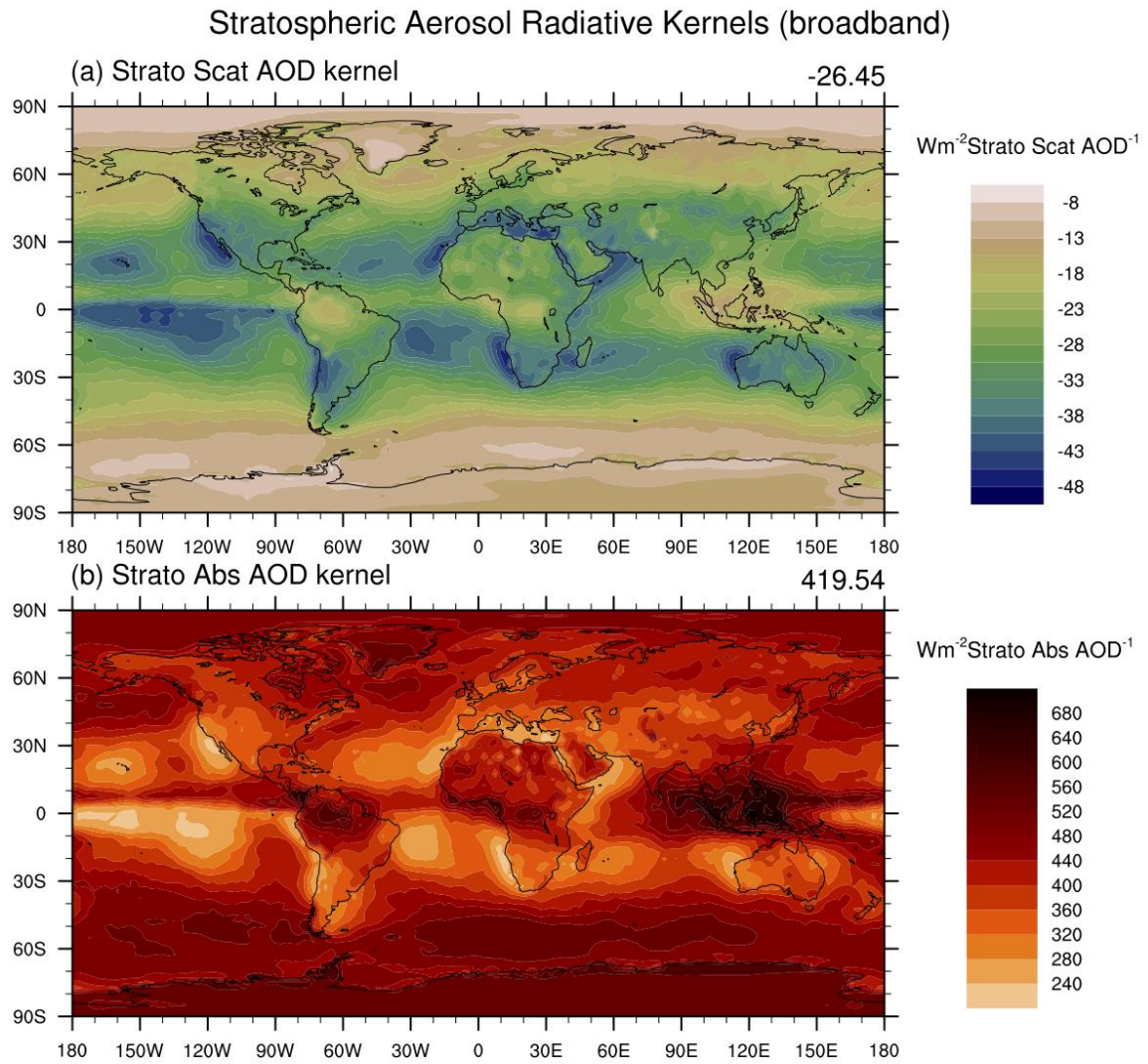
$$211 g' = g \quad (7)$$

212 To use the aerosol kernels derived here to calculate ADRE, users simply need to obtain
 213 stratospheric ΔAOD_{scat}^{550} and ΔAOD_{abs}^{550} values appropriate to the case of interest, and then
 214 multiply these with broadband kernel values.

$$215 \Delta ADRE = \frac{\partial(ADRE)}{\partial(AOD_{scat})} \cdot \Delta AOD_{scat}^{550} + \frac{\partial(ADRE)}{\partial(AOD_{abs})} \cdot \Delta AOD_{abs}^{550} \quad (8)$$

216 Figure 1 shows the global distribution of annual mean stratospheric AOD_{scat} and AOD_{abs}
 217 kernels, in the units of W/m^2 per unit change in stratospheric AOD. Both AOD_{scat} and AOD_{abs}
 218 kernels exhibit strong atmosphere dependencies. In cloudy regions (e.g., the Intertropical
 219 Convergence Zone, tropical eastern Atlantic, northwest Pacific Ocean, and Southern Ocean), the

220 sensitivity of stratospheric ADRE to stratospheric AOD_{scat} is relatively lower due to the presence
221 of underlying clouds, while the sensitivity to AOD_{abs} is relatively higher, compared to other
222 regions. This is because in the case of the scattering effect, clouds already brighten the
223 atmosphere and make the TOA radiation less sensitive to scattering aerosols and in the case of
224 the absorbing effect, clouds increase the solar radiation reflected into the stratosphere, thereby
225 amplifying the absorption by the stratospheric aerosols. Similar patterns are observed over the
226 polar and desert regions with high surface albedo. Because of their scattering or absorbing nature,
227 AOD_{scat} kernels are always negative, while AOD_{abs} kernels are always positive. In terms of
228 global means, a 0.1 increase in stratospheric AOD_{scat}^{550} results in a -2.65 W/m² cooling, while a 0.1
229 increase in AOD_{abs}^{550} results in a $+41.95$ W/m² warming at the TOA. Note that these sensitivity
230 values are larger than those reported by Q. Yu & Huang (2023b), particularly for absorbing
231 aerosols. This is because the kernels developed in this study focus exclusively on stratospheric
232 aerosols. These aerosols interact with a larger proportion of photons that have not been
233 attenuated by clouds or tropospheric absorbers. Additionally, underlying clouds enhance the
234 brightness of the troposphere, which further intensify the sensitivity of stratospheric ADRE to
235 AOD_{abs} .



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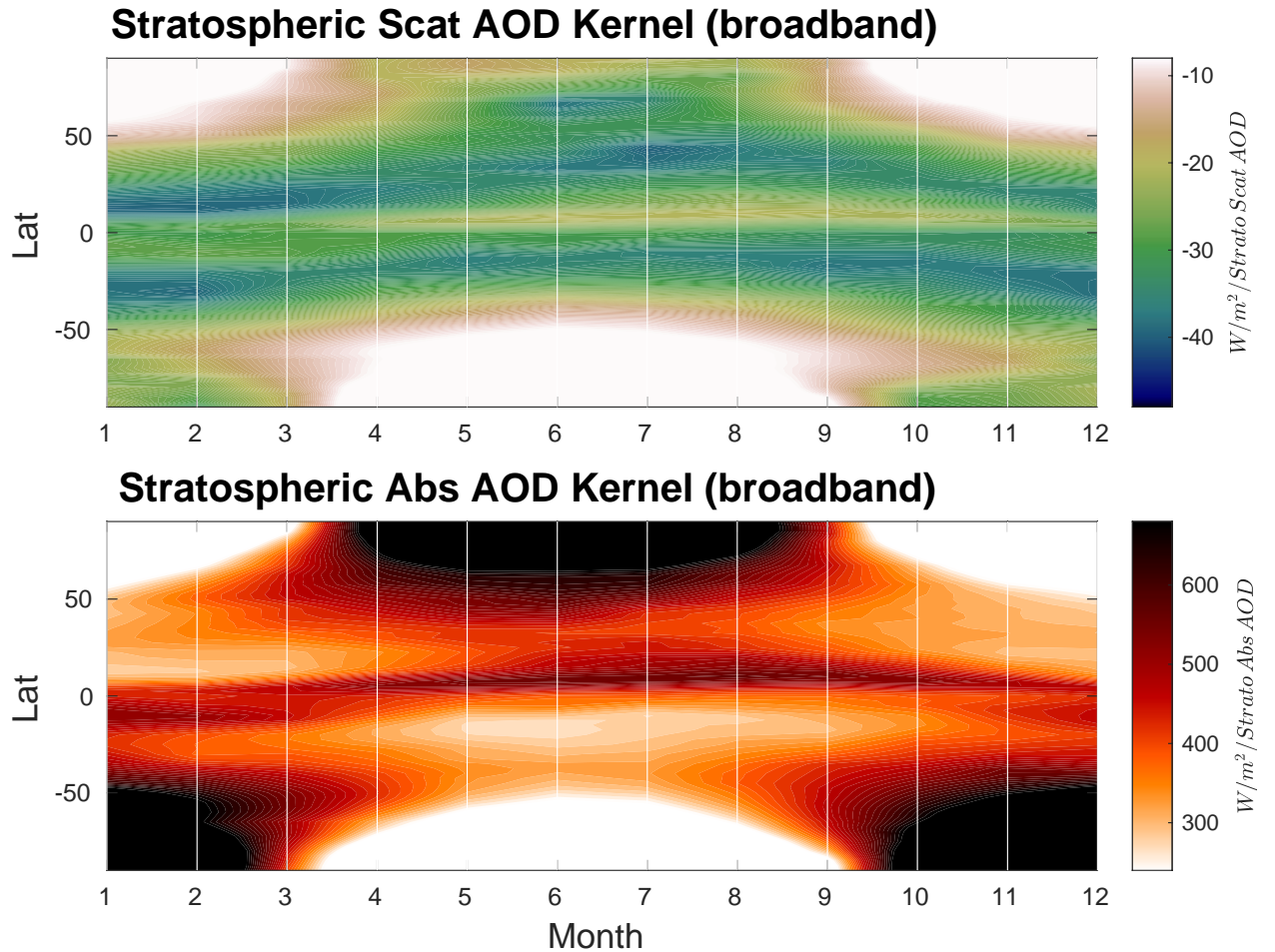
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Figure 1. Spatial distributions of annual mean broadband aerosol kernels (a) for stratospheric AOD_{scat} and (b) for stratospheric AOD_{abs} . The global mean and annual mean values are indicated in the upper right corner of each subplot. Kernels are shown in units of watts per square meter per unit change in stratospheric AOD at 550 nm.



241

242 Figure 2. Temporal variations of zonal mean broadband stratospheric aerosol kernels (a)
 243 for stratospheric AOD_{scat} and (b) for stratospheric AOD_{abs} .

244 Apart from the spatial inhomogeneity, stratospheric aerosol kernels also display strong
 245 temporal variations. Figure 2 displays the temporal variations in zonal mean stratospheric
 246 broadband AOD_{scat} and AOD_{abs} kernels. The pronounced latitudinal differences in aerosol
 247 kernels reflect patterns of solar insolation. In tropical regions, the sensitivity of ADRE to
 248 stratospheric aerosols remains high throughout the year, while polar regions show notable
 249 seasonal variations.

250

2.2.2 Band-by-band Aerosol Kernels

251 While broadband aerosol kernels are convenient to use, they rely on assumptions about
 252 the wavelength dependency of aerosol optical properties, which may not always be accurate. To
 253 facilitate a more flexible and accurate ADRE quantification, we leverage the band configuration
 254 of the RRTMG model to calculate a set of band-by-band stratospheric aerosol kernels. The
 255 RRTMG shortwave bands, detailed in Table 1, cover a spectrum from $0.2 \mu\text{m}$ to $12.2 \mu\text{m}$ across
 256 14 bands.

257

Table 1. RRTMG shortwave bands.

SW band	Wavenumber $\nu [cm^{-1}]$	Wavelength $\lambda [nm]$	AOD wavelength [nm]
Band 29	820-2600	12195- 3846	7082.2
Band 16	2600-3250	3846-3077	3444.7
Band 17	3250-4000	3077-2500	2777
Band 18	4000-4650	2500-2151	2320.2
Band 19	4650-5150	2151-1942	2044.2
Band 20	5150-6150	1942-1626	1778.4
Band 21	6150-7700	1626-1299	1455.2
Band 22	7700-8050	1299-1242	1270
Band 23	8050-12850	1242-778	944.3
Band 24	12850-16000	778-625	693.5
Band 25	16000-22650	625-442	527.1
Band 26	22650-29000	442-345	399.8
Band 27	29000-38000	345-263	329.1
Band 28	38000-50000	263-200	229.8

258

The stratospheric aerosol band-by-band kernels for AOD_{scat} and AOD_{abs} are expressed as:

$$259 \quad \frac{\partial ADRE^i}{\partial AOD_{scat}^i} = \frac{F^{net}(AOD_{scat}^i + \Delta AOD_{scat}^i, AOD_{abs}^i, SSA^i, g^i) - F^{net}(AOD_{scat}^i, AOD_{abs}^i, SSA^i, g^i)}{\Delta AOD_{scat}^i} \quad (9)$$

$$260 \quad \frac{\partial ADRE^i}{\partial AOD_{abs}^i} = \frac{F^{net}(AOD_{scat}^i, AOD_{abs}^i + \Delta AOD_{abs}^i, SSA', g') - F^{net}(AOD_{scat}^i, AOD_{abs}^i, SSA', g')}{\Delta AOD_{abs}^i} \quad (10)$$

261 In the equations above, i represents the i^{th} band in RRTMG. ΔAOD_{scat}^i and ΔAOD_{abs}^i are the
 262 added AOD perturbation at the i^{th} band, which vary with wavelength according to the Angstrom
 263 relation in our calculation.

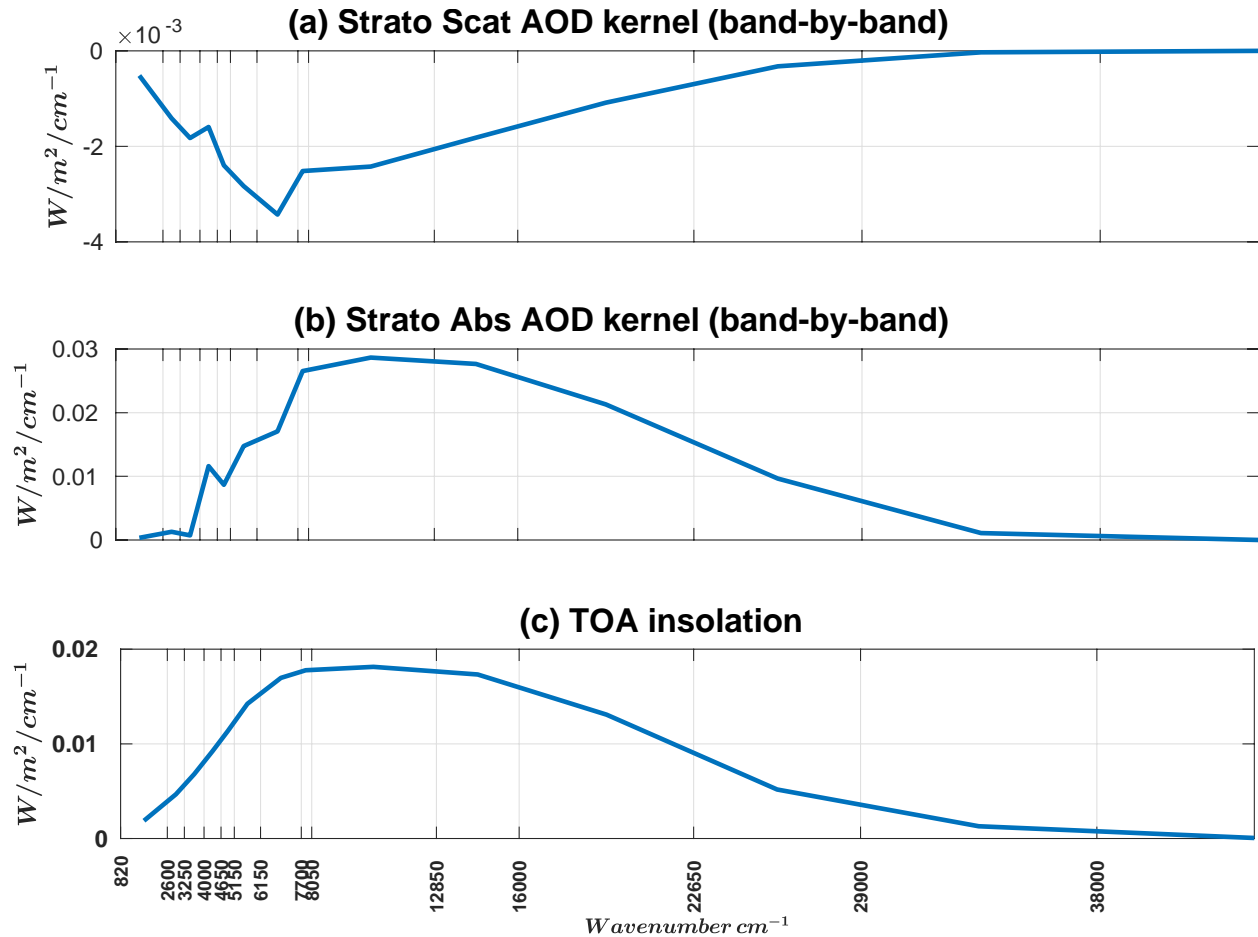
$$264 \quad \Delta AOD_{scat}^i = 0.1 * \left(\frac{\text{wavelength}^i}{550} \right)^{-1} \quad (11)$$

$$265 \quad \Delta AOD_{abs}^i = 0.01 * \left(\frac{\text{wavelength}^i}{550} \right)^{-1} \quad (12)$$

266 Note that for each band, perturbed AOD is calculated at the central wavelength following
 267 RRTMG configuration as listed in Table 1. The SSA and g calculations in the perturbation runs
 268 are similar to those in the broadband kernel calculation.

269 To use the band-by-band kernels, users need to obtain the ΔAOD_{scat}^i and ΔAOD_{abs}^i for
 270 each band, multiply them by band-by-band kernels, and sum over the 14 bands.

$$271 \quad \Delta ADRE = \sum_{i=16}^{29} \left(\frac{\partial ADRE^i}{\partial AOD_{scat}^i} \cdot \Delta AOD_{scat}^i \right) + \sum_{i=16}^{29} \left(\frac{\partial ADRE^i}{\partial AOD_{abs}^i} \cdot \Delta AOD_{abs}^i \right) \quad (13)$$



272

273 Figure 3. Global mean annual mean stratospheric aerosol band-by-band kernels for (a) AOD_{scat}
 274 and (b) AOD_{abs} . For demonstration purposes, kernels are normalized by the corresponding
 275 bandwidth. The normalized kernel unit is watts per meter squared per unit change in the
 276 respective stratospheric AOD per wavenumber. (c) Normalized spectral solar radiation.

277 Figure 3 presents the global mean band-by-band stratospheric AOD_{scat} and AOD_{abs}
 278 kernels. For comparison purposes, the spectral kernels are normalized by the bandwidth. The
 279 results indicate that the spectral signatures of the band-by-band aerosol kernels are primarily
 280 dominated by the strength of incoming solar radiation. The aerosol radiative sensitivity peaks
 281 from the near-ultraviolet band ($\sim 22650\text{ cm}^{-1}$) to the near-infrared band ($\sim 8080\text{ cm}^{-1}$), which
 282 corresponds to band 23 to 25 (442 nm-1242 nm) in RRTMG as indicated in Table 1. Accurately
 283 determining aerosol optical properties in these bands can help constrain the ADRE without
 284 needing the aerosol information across the full spectrum. Most aerosol retrieval products provide
 285 optical properties at a few discrete wavelengths ranging from near-ultraviolet to near-infrared.
 286 For example, the AEROSOL ROBOTIC NETWORK (AERONET) provides AOD products at 340, 380,
 287 440, 500, 675, 870, and 1020 nm (Giles et al., 2019). By interpolating observed AOD values at
 288 the central AOD wavelengths in the RRTMG configuration for relevant bands and assuming
 289 spectral dependence of optical properties for the remaining bands, users can calculate
 290 stratospheric ADRE more accurately than using broadband kernels. In the following section, we
 291 will use the spectral AOD observations to compare the ADRE values computed from the
 292 broadband and band-by-band stratospheric aerosol kernels.

293 2.3 OMPS Aerosol Data and Quality Control

294 To quantify the stratospheric ADRE, we utilize the aerosol extinction coefficient profiles
295 from the OMPS-LP Level 2 daily product. The Ozone Mapping and Profiler Suite (OMPS)
296 measures limb scattering of sunlight at tangent altitudes from ground level up to approximately
297 100 km with a vertical resolution of 1km (Flynn et al., 2006). The aerosol product from OMPS
298 has been widely used to study the stratospheric ADRE (Damany-Pearce et al., 2022; Bourassa et
299 al., 2023; Schoeberl et al., 2023, 2024). This study uses aerosol extinction coefficient retrievals
300 along the center slit (aligned with the orbital track) of the OMPS-LP. The retrieved extinction
301 profiles extend up to 40km, and quality control procedures are applied before the analysis
302 following Damany-Pearce et al. (2022). Only data with ResidualFlag = 0, SingleScatteringAngle
303 $\leq 145^\circ$, and SwathLevelQualityFlags with bits 0, 1, and 7 = 0 are considered valid. The
304 tropopause definition is consistent with that used in the kernel calculation. We integrate the
305 extinction coefficient throughout the stratosphere to calculate the stratospheric AOD. To
306 facilitate kernel application, we average the AOD data onto the same $2.5^\circ \times 2.5^\circ$ latitude-
307 longitude grid.

308 For using the broadband aerosol kernels, we choose the 869 nm extinction coefficient
309 from OMPS and scale it to 550 nm, assuming an AE value of 1. This AE value is chosen because
310 it represents the background stratospheric aerosol conditions and the specific conditions of the
311 Hunga aerosols, and has been applied in other similar studies (Schoeberl et al., 2023; Sellitto et
312 al., 2024). The 869 nm wavelength is chosen over other channels closer to 550 nm because
313 OMPS aerosol products have performance issues at shorter wavelengths in the southern
314 hemisphere (Taha et al., 2021).

315 For the band-by-band kernel application, we utilize extinction coefficients measured at
316 510 nm, 600 nm, 675 nm, 745 nm, 869 nm, and 997 nm, and interpolate extinction values to
317 527.1 nm, 693.5 nm, and 944.3 nm using measurements from the nearest wavelengths as
318 required by the aerosol kernels. For the remaining bands, we scale the extinction coefficient from
319 869 nm to the corresponding central AOD wavelength, assuming an AE of 1. In the following
320 section, we use the OMPS spectral AODs as an example to demonstrate the usage of our kernels.

321 Our goal is to estimate the changes in stratospheric ADRE (Δ ADRE) due to the 2022
322 Hunga volcanic eruption and the 2020 Australia wildfires using our aerosol kernels. We consider
323 the MERRA-2 stratospheric AOD as the background aerosol states because no eruptive
324 volcanoes are included in MERRA-2 after 2010 (Randal et al., 2016). Therefore, the
325 stratospheric AOD anomaly is calculated by subtracting the background stratospheric AOD
326 values given by MERRA-2 from OMPS stratospheric AOD. For the kernel application, the AOD
327 values in Equations (8) and (13) are the differences between OMPS and MERRA-2 stratospheric
328 AOD. To validate the performance of our aerosol kernels, we use the same AOD anomalies as
329 input to the RRMTG model to calculate the "truth" values of stratospheric ADRE for comparison.

330 **3 Physically Sorted Aerosol Kernels**

331 As shown in the previous section (e.g., Figures 1 and 2), there are strong spatial and
332 temporal variabilities in the kernel values. It is thus important to understand how the aerosol
333 properties and environmental variables (e.g., surface albedo and clouds) interact with each other

334 to influence the radiative sensitivity. To address this question, we follow a widely used
335 conceptual model of ADRE (Chlek & Coakley Jr, 1974; Haywood & Shine, 1995) to identify the
336 key factors and their expressions to use in an analytical model to predict the kernel values. We
337 then determine the coefficient values statistically using a multivariable regression method
338 following (Q. Yu & Huang, 2023a, 2023b). Different from the geographically gridded kernels
339 presented in the previous section, the physically sorted kernels developed here are not
340 constrained by space and time, allowing one to more flexibly estimate the stratospheric ADRE.

341 3.1 Physical Model

342 We follow the formulation of Haywood & Shine (1995), but consider the stratospheric
343 aerosols as a scattering layer and represent the troposphere-surface system as a whole with a
344 reflectance parameter at the tropopause. The all-sky stratospheric ADRE at the TOA can be
345 expressed as follows:

$$346 \quad ADRE = -ST_{at}^2 \beta \omega \tau \sec \theta \frac{(1-R_s)^2 - \frac{R_s(1-\omega)}{\beta} \left[\frac{2-\tau \sec \theta}{\omega} - \tau \sec \theta (2\beta-1) \right]}{1-R_s \beta \omega \tau \sec \theta} \quad (14)$$

347 The environment-related variables are solar insolation (S), atmospheric transmittance (T_{at}) above
 348 the aerosol layer, the solar zenith angle (θ), and tropopause reflectance (R_s). The aerosol-related
 349 variables are the aerosol backscattering ratio (β), aerosol single scattering albedo (ω), and
 350 aerosol optical depth (τ). The stratospheric ADRE is further expanded as:

$$351 \quad ADRE = -ST_{at}^2 \beta \omega \tau \sec \theta (1 + R_s \tau \sec \theta \beta \omega) \left\{ (1 - R_s)^2 - \frac{R_s(1-\omega)}{\beta} \left[\frac{2-\tau \sec \theta}{\omega} - \tau \sec \theta (2\beta - 1) \right] \right\}$$

352 (15)

353 The sensitivity of stratospheric ADRE to τ is

$$354 \quad \frac{\partial ADRE}{\partial \tau} =$$

$$355 \quad -ST_{at}^2 \beta \omega (1 - R_s)^2 (1 + 2\beta \omega R_s \tau \sec \theta) + ST_{at}^2 \beta \omega (1 + 2\beta \omega R_s \tau \sec \theta) \frac{R_s(1-\omega)}{\beta} \left[\frac{2-\tau \sec \theta}{\omega} - \right.$$

$$356 \quad \left. \tau \sec \theta (2\beta - 1) \right] + ST_{at}^2 \beta \omega (\tau \sec \theta + \beta \omega R_s \tau \sec \theta^2) \left[-\frac{R_s}{\beta} \frac{1-\omega}{\omega} - 2\beta + 1 \right] \quad (16)$$

357 Neglecting higher-order terms, Equation (16) is approximated as

$$358 \quad \frac{\partial ADRE}{\partial \tau} = -ST_{at}^2 [\beta \omega + R_s (2\beta \omega + 2 - 3\omega) - R_s^2 \beta \omega + R_s \tau \sec \theta (-2 + 3\omega - 2\omega \beta - \omega^2) +$$

$$359 \quad \beta \omega \tau \sec \theta] \quad (17)$$

360 This equation suggests that stratospheric aerosol kernels are influenced by these terms: R_s , $R_s \tau$,
 361 τ , and R_s^2 . The combination terms arise from the coupling effects between the stratospheric
 362 aerosol layer and the underlying troposphere-surface system. In the following section, we will
 363 use these terms as predictors to reproduce the spatiotemporally varying stratospheric aerosol
 364 kernels. The goal is to capture the physical processes governing ADRE sensitivity, which should
 365 be independent from geographic locations.

366 3.2 Statistical Model

367 Regression models have been a useful tool in predicting radiative forcing and capturing
 368 nonlinear radiative interactions in many studies (Huang et al., 2016; Datsieris et al., 2022; Q. Yu
 369 & Huang, 2023b, 2023a). In this work, we built a multi-variable regression model to represent
 370 the annual mean global stratospheric aerosol kernels following Q. Yu & Huang (2023b, 2023a).
 371 The model is expressed as:

$$372 \quad \frac{Y(i,j) - \bar{Y}}{\bar{Y}} = \sum_{k=1}^n A_k \frac{X_k(i,j) - \bar{X}_k}{\bar{X}_k} \quad (18)$$

373 Here, X are predictors (e.g., R_s , $R_s \tau$) at latitude i and longitude j . Y is either the broadband
 374 aerosol kernels for stratospheric AOD_{scat} or AOD_{abs} . A_k is the regression coefficient and n is the
 375 number of predictors. Note that the global field of Y is predicted by one uniform set of A_k values.
 376 Both predictors and predictands are normalized by their global mean values, denoted by a bar.

377 Following the physical model derived above, we select R_s , $R_s\tau$, τ , and R_s^2 as predictors.
 378 As TOA reflectance (R) is more easily obtained, we use it as a proxy for the tropopause
 379 reflectance. To accurately represent global aerosol kernels using as few predictors as possible,
 380 we have tested the performance of all possible combinations of predictors (listed in Supporting
 381 Information Table 1&2). Results suggest that the four predictors are sufficient to capture almost
 382 all main features of stratospheric aerosol kernels.

383 The physically sorted broadband aerosol kernels for stratospheric AOD_{scat} is given by

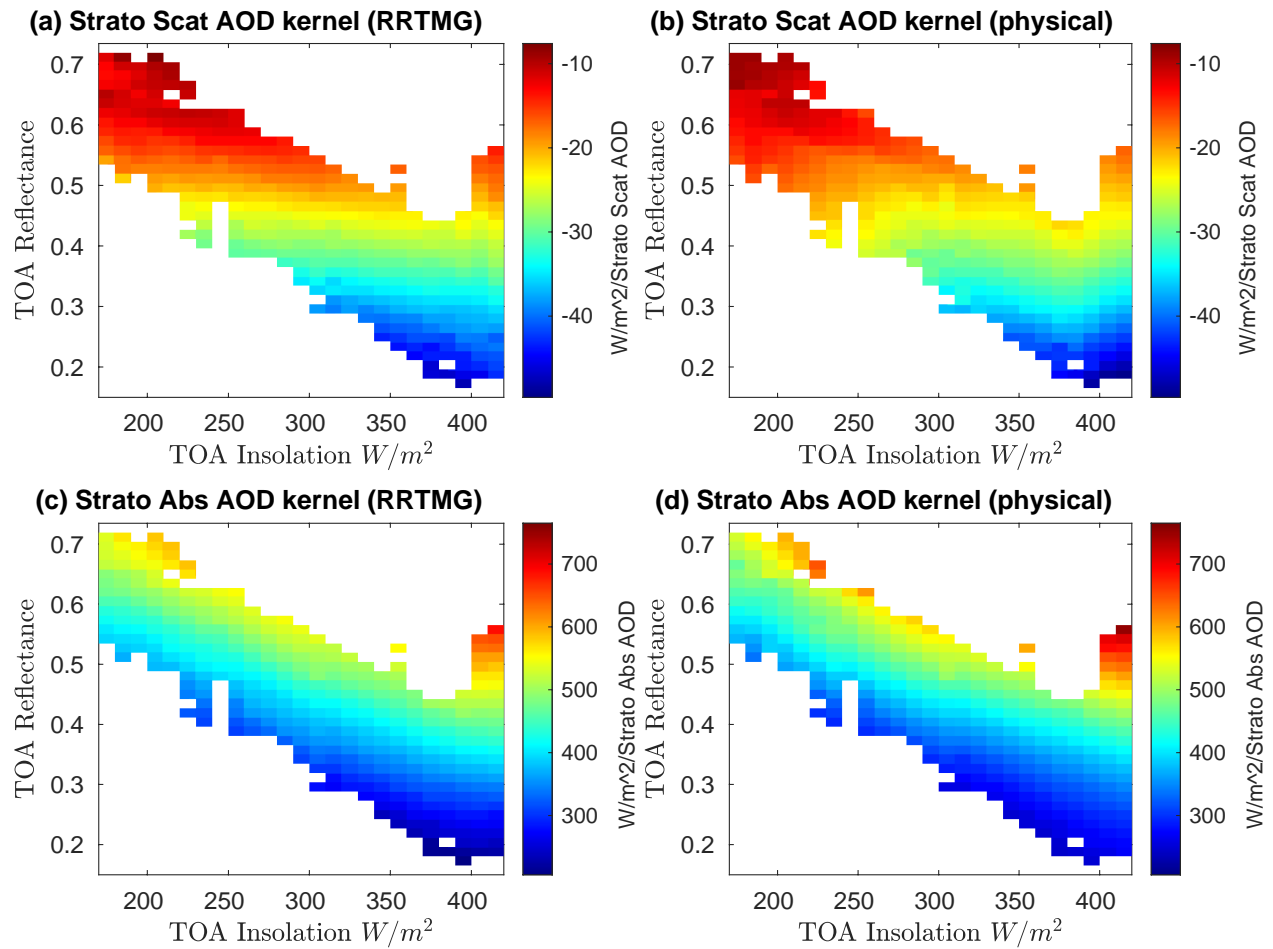
$$384 \frac{\frac{\partial(ADRE)}{\partial(AOD_{scat})} - (-0.076)}{S} = -2.264 \cdot \frac{R-0.413}{0.413} + 0.753 \frac{R^2-0.184}{0.184} + 0.671 \frac{\tau-0.002}{0.002} - 0.3186 \frac{R\tau-0.001}{0.001} \quad (19)$$

385 The physically sorted broadband aerosol kernels for stratospheric AOD_{abs} is given by

$$386 \frac{\frac{\partial(ADRE)}{\partial(AOD_{abs})} - 1.323}{S} = -0.313 \cdot \frac{R-0.413}{0.413} + 0.696 \frac{R^2-0.184}{0.184} - 0.175 \frac{\tau-0.002}{0.002} + 0.258 \frac{R\tau-0.001}{0.001} \quad (20)$$

387 The comparison of statistically fitted broadband aerosol kernels for stratospheric AOD_{scat} and
 388 AOD_{abs} against benchmark RRTMG calculations is shown in the Supporting Information
 389 (Figures S8). Results suggest that more than 94% of the spatial variance in aerosol kernels is
 390 captured by the regression model, indicating its effectiveness in predicting the variability of
 391 aerosol kernels.

392 Figure 4 displays the impact of environmental variables (TOA insolation S and
 393 reflectance R) on the distributions of annual mean global aerosol kernels. Generally speaking, an
 394 increase in solar insolation results in a larger magnitude of aerosol kernels, while a more
 395 reflective underlying "surface" (due to clouds or Earth's surface) leads to a less cooling or more
 396 warming impact on net TOA fluxes. The physically sorted aerosol kernels can well capture their
 397 sensitivity to those environmental variables. More importantly, they can estimate stratospheric
 398 ADRE sensitivity in idealized conditions where actual observations are lacking.



399

400 Figure 4. Distributions of broadband stratospheric AOD_{scat} and AOD_{abs} kernels as a function of
 401 TOA reflectance and TOA insolation. Left column: RRTMG-calculated stratospheric aerosol
 402 kernels; Right column: the physically sorted aerosol kernels predicted by the regression model.

403 4 Stratospheric Aerosol Kernel Applications

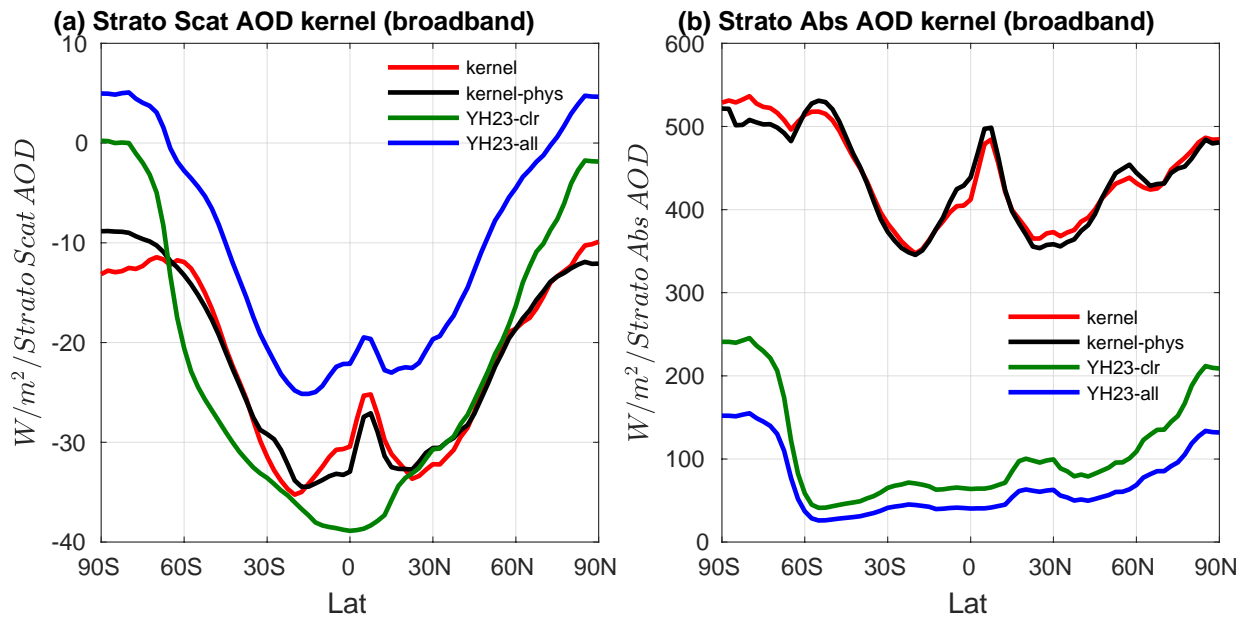
404 Multiplying aerosol radiative kernels by changes in stratospheric AOD from specific
 405 events (e.g., volcanic eruptions) provides estimates of the corresponding ADRE. In this section,
 406 we examine the radiative effects of the volcanic ash plume from the 2022 Hunga eruption and
 407 the biomass-burning aerosols from the 2020 Australia wildfires to demonstrate the application of
 408 stratospheric aerosol kernels. We compare the results of broadband, band-by-band, and
 409 physically sorted kernels.

410 4.1 Aerosol Radiative Kernel Comparisons

411 Given the kernels here are developed specifically for stratospheric aerosols, it is of
 412 interest to compare them with other kernels not designed this way. Besides the simple scaling
 413 relations given in the literature (e.g., Hansen et al. 2005; P. Yu et al. 2021), Q. Yu & Huang,
 414 (2023b, denoted as YH23 from here on) derived a set of global ADRE sensitivity kernels mainly
 415 for tropospheric aerosols and validated against the independent results of Thorsen et al. (2020).

416 Schoeberl et al. (2023, 2024) used the YH23 kernels to estimate the radiative impact of the
 417 Hunga eruption. We include the YH23 kernels for comparison in the following.

418 In Figure 5, we compare the zonal mean AOD_{scat} and AOD_{abs} sensitivity in YH23 with
 419 the broadband stratospheric aerosol kernels calculated by RRTMG and the statistical regression
 420 model. Results show that the aerosol kernels display significant latitudinal differences. For all-
 421 sky stratospheric AOD_{scat} kernels, the magnitude peaks in the subtropical regions because the
 422 relative brightness of aerosols is reduced above the tropical cloudy regions. The physically sorted
 423 kernels closely match the RRTMG results, indicating a good performance of the physical sorting
 424 method. Interestingly, the clear-sky, as opposed to the all-sky, AOD_{scat} kernels given by YH23
 425 render more similar magnitudes to the all-sky stratospheric AOD_{scat} kernels developed here,
 426 especially in the mid-latitudes. This is because the stratospheric aerosols are located above
 427 tropospheric clouds, which suppress the radiative sensitivity to tropospheric AOD perturbations
 428 but do not strongly affect the radiative effect of stratospheric aerosols. Compared to the AOD_{abs}
 429 kernels in YH23, the stratospheric AOD_{abs} kernels developed here are much larger due to the
 430 enhanced ADRE sensitivity to AOD_{abs} above bright underlying clouds.



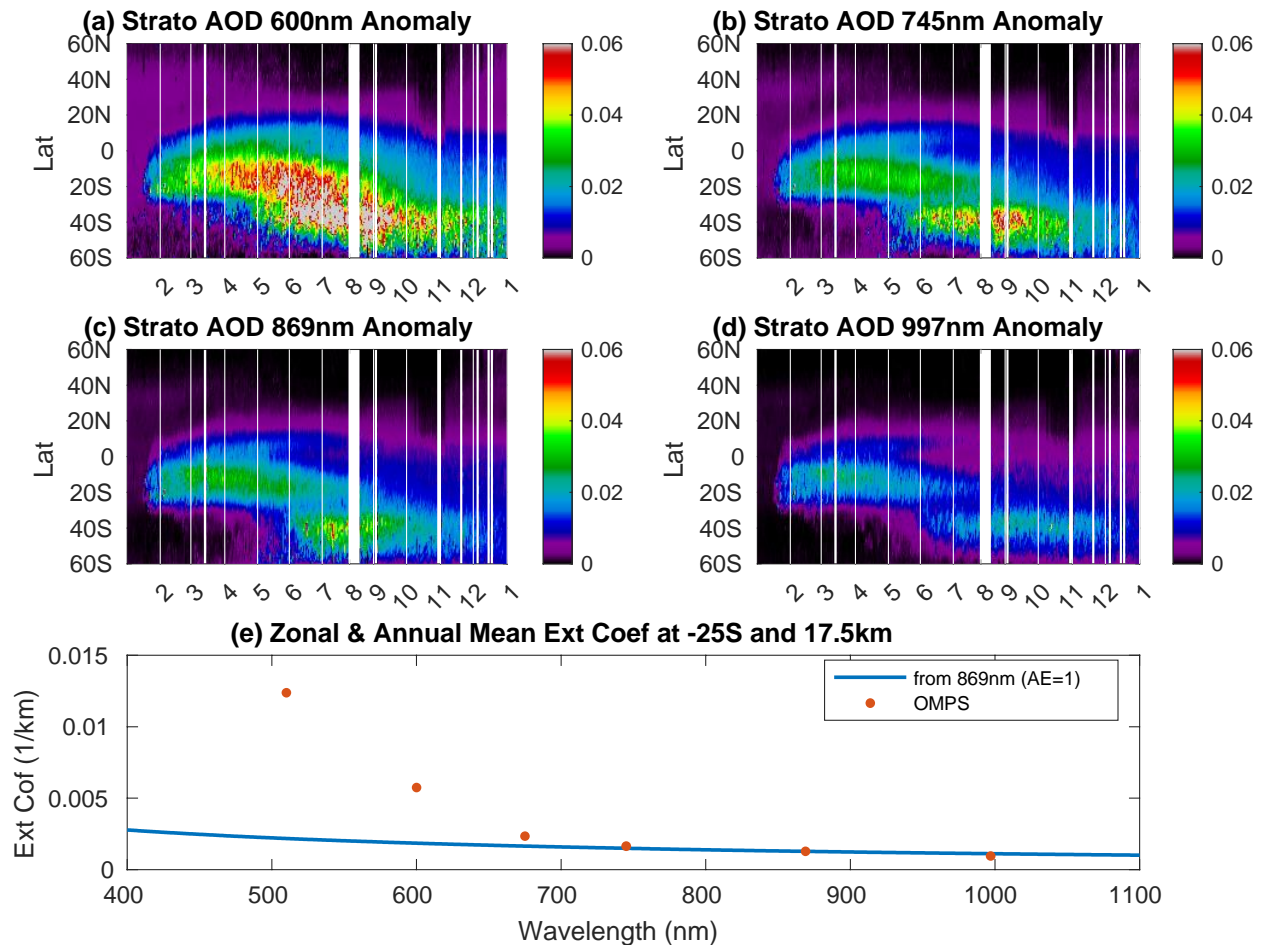
431
 432 Figure 5. Annual mean and zonal mean broadband stratospheric (a) AOD_{scat} and (b) AOD_{abs}
 433 radiative kernels. YH23-clr and YH23-all represent the clear-sky and all-sky scattering AOD
 434 radiative sensitivity quantified in Q. Yu & Huang (2023b) for tropospheric aerosols. Kernel and
 435 Kernel-phys indicate the broadband kernels calculated from RRTMG and emulated by a
 436 regression model, respectively.

437 4.2 2022 Hunga Volcanic Eruption

438 On January 15, 2022, the Hunga Tonga volcano (20.57°S , 175.38°W) erupted violently,
 439 releasing sulfur compounds and other aerosols into the atmosphere (Kloss et al., 2022; Taha et
 440 al., 2022; Schoeberl et al., 2023, 2024). To assess the corresponding ADRE, we first calculate
 441 the stratospheric AOD anomaly following Section 2.3.

442 Figure 6a-6d shows the evolution of the zonal mean stratospheric AOD anomaly relative
443 to the background at different wavelengths throughout 2022. Although the Hunga eruption
444 occurred in late January, the OMPS product showed little AOD signal initially because the
445 extinction retrieval becomes unreliable in the presence of clouds and optically thick aerosol
446 plumes (Taha et al., 2021). Over time, the aerosol plume descended to the lower stratosphere and
447 dispersed horizontally. Within four months after the eruption, the aerosols primarily remained in
448 tropical latitudes with some northward spread. This led to an initial AOD peak in the tropical
449 regions due to the immediate formation and accumulation of aerosols, as reported by other
450 studies (Schoeberl et al., 2023; Taha et al., 2022). As the southern hemisphere approached
451 winter, a meridional circulation developed between the tropics and subtropics to maintain the
452 thermal wind balance, known as the QBO direct, meridional, or secondary circulation (Strahan et
453 al., 2015). This circulation transported stratospheric aerosols into the mid-latitudes. Meanwhile,
454 the polar vortex acted as a barrier, causing the accumulated aerosols in the subtropics to create a
455 second AOD peak during July-September. The double peak features shown here were also
456 reported in other observations and model simulations (Wang et al., 2023; Schoeberl et al., 2024).
457 Gaps in the data are caused by spacecraft anomalies or failures to meet the data screening
458 criteria.

459 Figure 6e also suggests that the zonal average stratospheric AOD anomaly varies
460 significantly with wavelength, indicating that assuming a simple Angstrom exponent cannot fully
461 represent the wavelength dependency of AOD. Therefore, it is important to incorporate the band-
462 by-band kernels with AOD observations to accurately calculate the stratospheric ADRE.



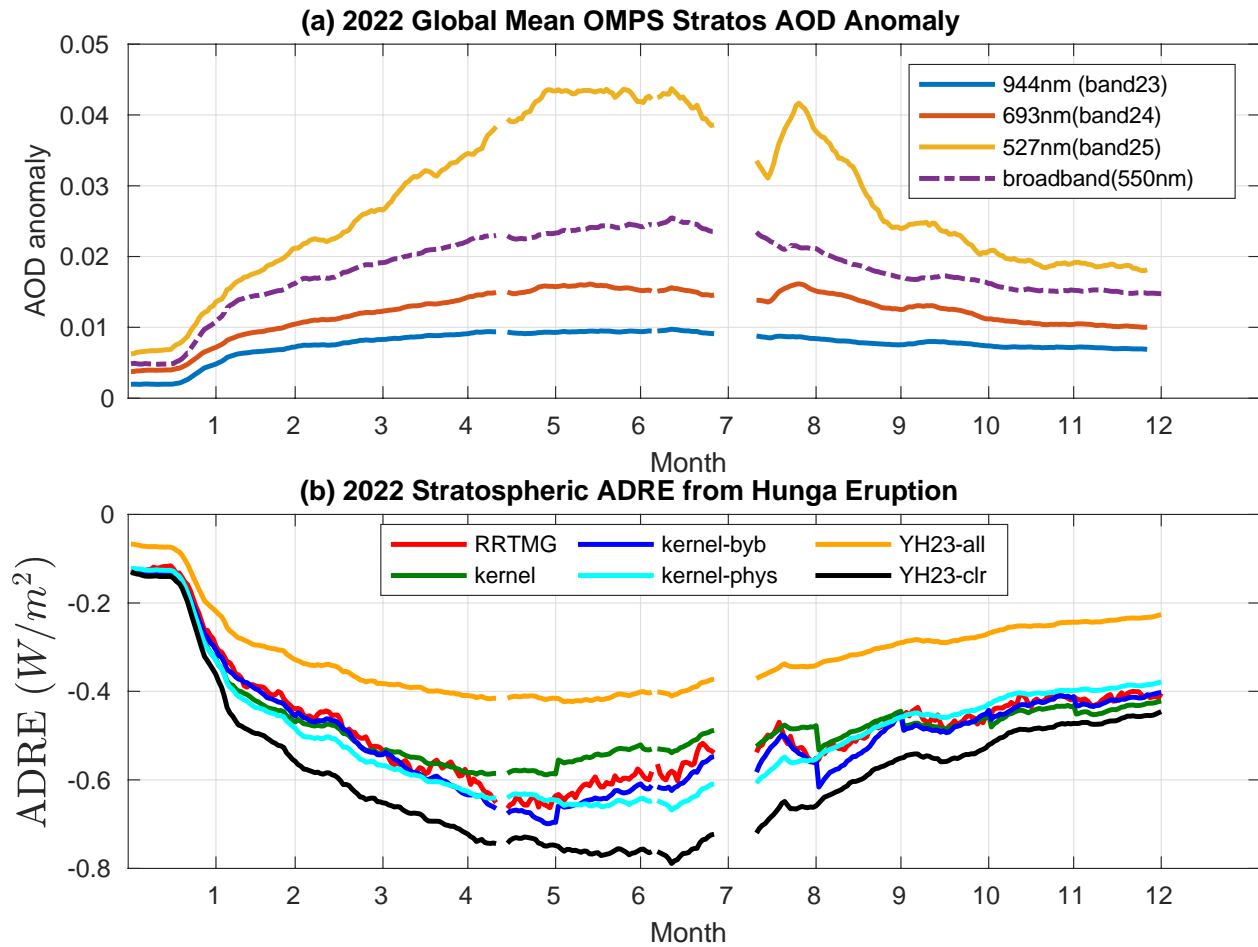
463

464 Figure 6. Latitude-time plots of the zonal mean stratospheric AOD anomaly at (a) 600 nm, (b)
 465 745 nm, (c) 869 nm, and (d) 997 nm from OMPS-LP in 2022, with the x-axis in (a)-(d)
 466 representing corresponding months. (e) Zonal and annual mean aerosol extinction coefficient at -
 467 25S and 17.5 km. The red dots represent OMPS observations, while the blue line shows the
 468 wavelength dependency assuming an AE of 1.

469 In the ADRE calculation, we assume the AOD anomaly with $SSA = 1$ because
 470 observations suggest that the absorbing particles in the volcanic ashes are of small amounts and
 471 do not significantly impact the radiative properties (Kloss et al., 2022). We also assume the
 472 stratospheric AOD anomalies from OMPS at a discrete set of wavelengths represent the
 473 observational truth. Figure 7a displays the global mean stratospheric AOD anomaly as calculated
 474 in Section 2.3 throughout 2022. We have listed the spectral AOD anomaly at RRTMG mid-
 475 visible bands (bands 23-25). These values are interpolated from the nearby wavelengths from
 476 OMPS. For the broadband AOD anomaly, we calculate the AOD using OMPS 869 nm, assuming
 477 an AE of 1. This way, we can estimate the relative errors of the broadband kernel method when
 478 there are observation uncertainties in AE. Results indicate distinct features in the spectra AOD,
 479 suggesting a peak in global mean AOD values around June.

480 We further calculated the stratospheric ADRE using both the stratospheric kernels
 481 developed here and the kernels from YH23. The YH23 kernels, although based on total column
 482 aerosols, have been used in stratospheric ADRE quantifications in Schoeberl et al. (2023, 2024).

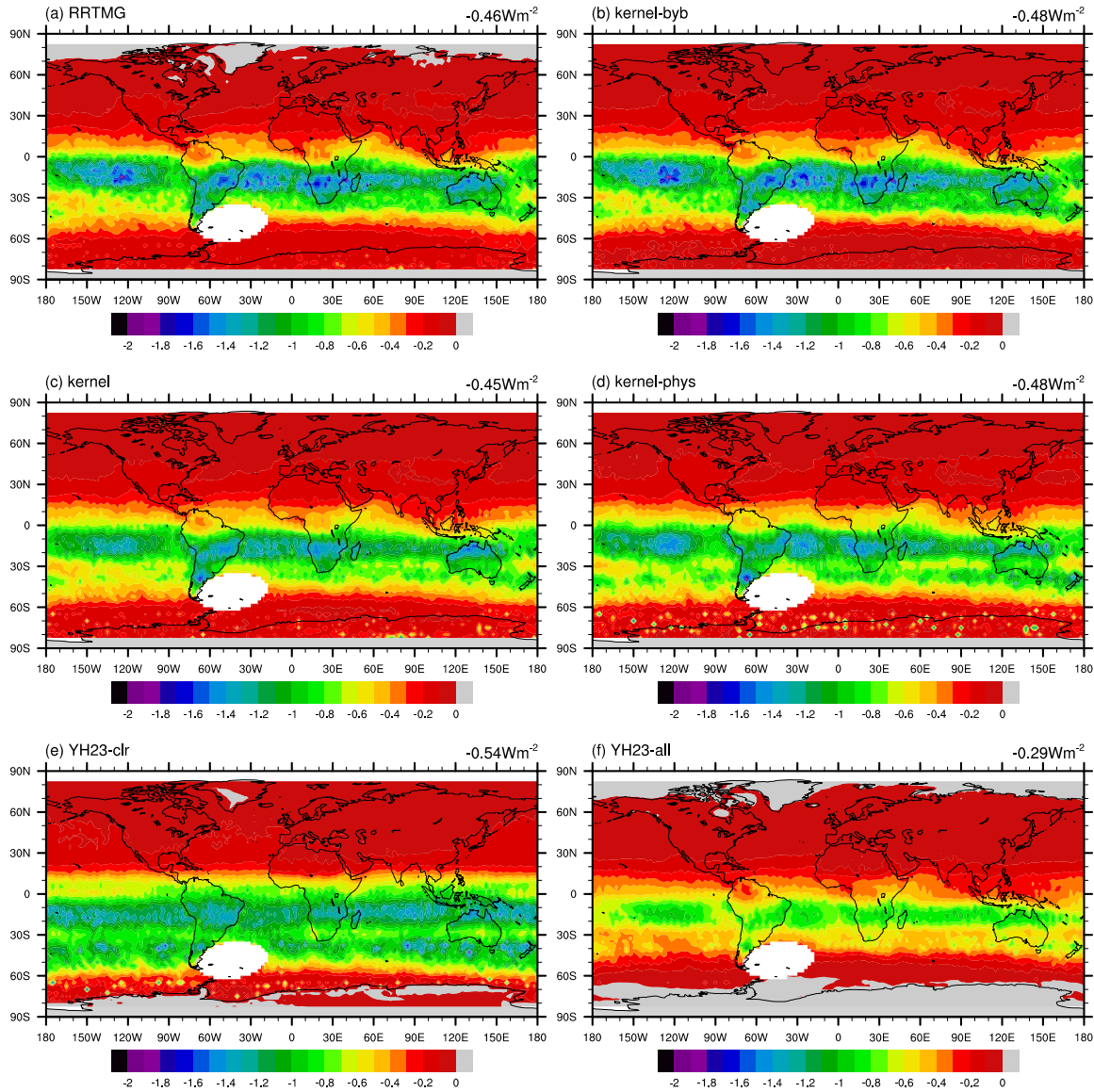
483 By including YH23 kernels in the comparison, we show the discrepancies that would be caused
484 by kernels not specifically made for stratospheric aerosols. Figure 7b shows the stratospheric
485 ADRE from the Hunga Eruption in 2022. For comparison, the RRTMG-calculated results based
486 on the band-by-band AOD inputs are indicated by the red line. In general, the ADRE peaks with
487 AOD near June, and using the band-by-band aerosol kernels can quantify it most accurately. The
488 performance of broadband and physically sorted stratospheric aerosol kernels is slightly worse
489 than that of the band-by-band kernels, as they fail to capture the wavelength dependency
490 information and the Angstrom exponent assumption may be inadequate. In terms of global mean
491 values, using the YH23 clear-sky kernel overestimates the cooling effect of the Hunga eruption,
492 while using the YH23 all-sky kernel underestimates it.



493

494 Figure 7. Time series of the global mean (a) stratospheric AOD anomaly from OMPS-LP
 495 following the Hunga Eruption in 2022 and (b) stratospheric ADRE from Hunga Eruption in 2022.
 496 YH23-clr and YH23-all represent the clear-sky and all-sky scattering AOD radiative sensitivity
 497 quantified in Q. Yu & Huang (2023b), respectively. Kernel, kernel-phys, and kernel-byb indicate
 498 the broadband kernels calculated from RRTMG, broadband kernels from the regression model,
 499 and the band-by-band kernels, respectively.

Stratospheric ADRE from Hunga Eruption



500

501 Figure 8. Annual mean stratospheric ADRE from the Hunga eruption in 2022, with global mean
 502 values indicated in the top right of each subplot. (a) RRTMG benchmark calculations; (b) Band-
 503 by-band kernel quantifications; (c) Broadband kernel quantifications; (d) Physically sorted kernel
 504 quantifications; (e) YH23 clear-sky kernel quantifications; (f) YH23 all-sky kernel
 505 quantifications. Global mean values are shown in the top right of each subplot.

506 Apart from the time evolution, we also compare the spatial patterns of stratospheric
 507 ADRE using different kernel schemes. Figure 8 displays the annual mean stratospheric ADRE
 508 from the Hunga eruption calculated from RRTMG as well as the kernels developed in this work.
 509 Results show that the volcanic eruption caused a uniform cooling in the southern hemisphere's
 510 tropical and subtropical regions due to the dispersion of aerosols described before. In terms of
 511 global mean ADRE, the Hunga eruption induced a cooling of -0.46 W/m^2 . All stratospheric

512 kernels developed in this work can reproduce the spatial features of ADRE relatively well, with
 513 the band-by-band kernels performing the best. Although the YH23 clear-sky scheme can
 514 approximately reproduce the global mean stratospheric ADRE values, it fails to capture the
 515 spatial patterns, especially over the cloudy regions.

516 Table 2 listed the R^2 and RMSE values comparing the ADRE induced by the Hunga
 517 eruption, calculated using different kernel schemes and the RRTMG model. Globally, the band-
 518 by-band, broadband, and physically sorted aerosol kernels capture 98.89%, 93.83%, and 94.33%
 519 of the variance in RRTMG-calculated ADRE, with RMSEs less than 0.04 W/m² (approximately
 520 8.7% relative to the global mean values). Using YH23 kernels results in RMSEs greater than
 521 0.11 W/m², which is 23.91% relative to the global mean.

522 Table 2. Performance of stratospheric kernels calculated in this study and kernels from YH23 in
 523 quantifying the ADRE of the 2022 Hunga volcanic eruption and 2020 Australia wildfire. R^2
 524 represents the coefficient of determination, and RMSE is the Root Mean Squared Error. Relative
 525 errors are calculated by dividing the RMSE by the global mean values. Broadband and
 526 physically sorted kernels are used under the assumption of AE being 1.

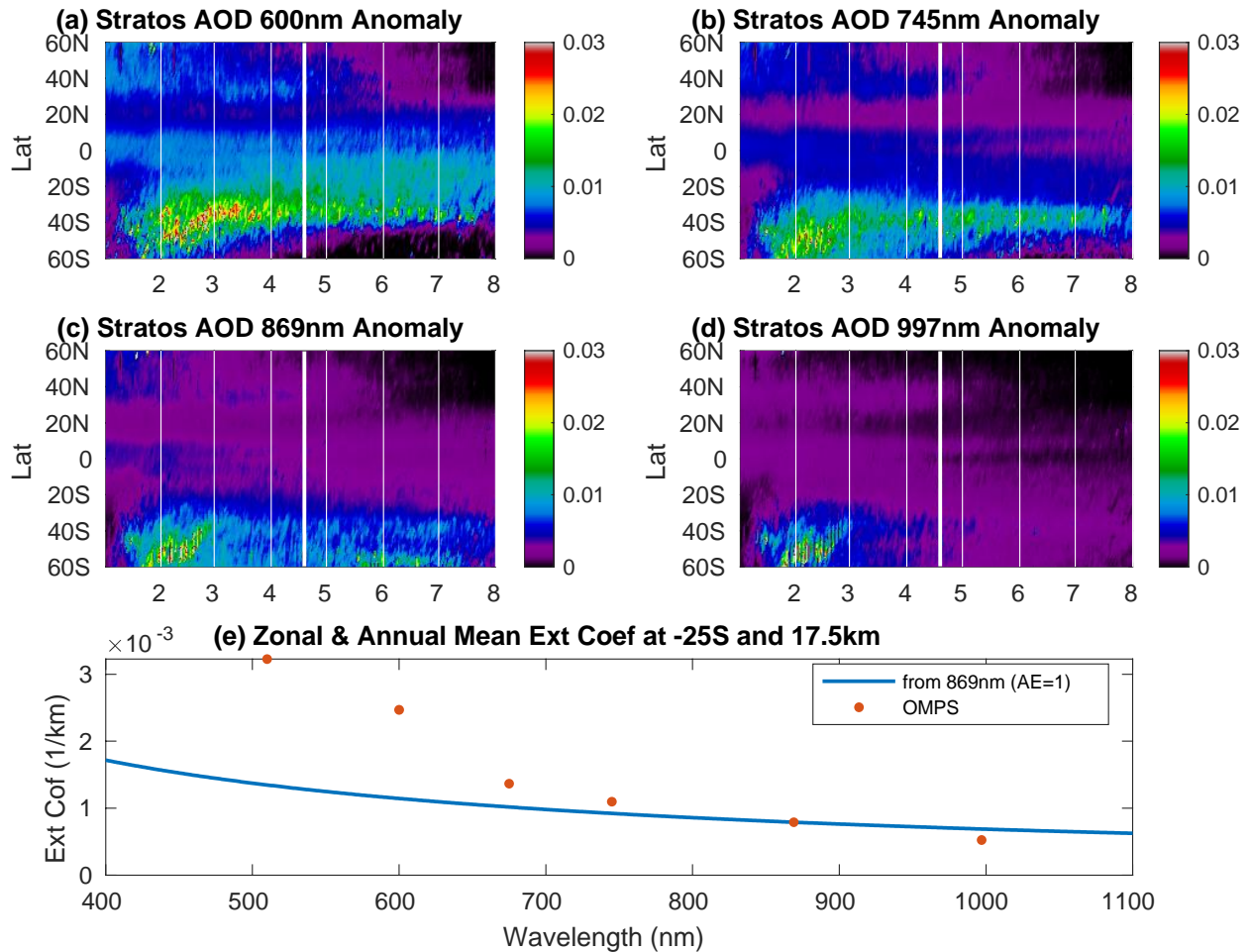
	2022 Hunga volcanic eruption			2020 Australia wildfire		
	R^2	RMSE (W/m ²)	Relative errors	R^2	RMSE(W/m ²)	Relative errors
Band-by-band kernels	98.89%	0.02	4.35%	99.02%	0.01	3.57%
Broadband kernels	93.83%	0.04	8.70%	94.75%	0.04	14.29%
Physically sorted kernels	94.33%	0.04	8.70%	83.19%	0.08	28.57%
YH23-clr	95.11%	0.11	23.91%	83.59%	0.37	132.14%
YH23-all	91.88%	0.17	36.96%	51.32%	0.31	110.71%

527

528 4.3 2020 Australia Wildfire

529 In late December 2019, massive bushfires occurred in southeastern Australia and lifted a
 530 considerable amount of smoke into the stratosphere via pyrocumulonimbus clouds. Unlike the

531 volcanic eruption case, we apply both the AOD_{scat} and AOD_{abs} kernels to study the stratospheric
 532 ADRE of the black-carbon-containing smoke particles as they are partly absorbing.



533

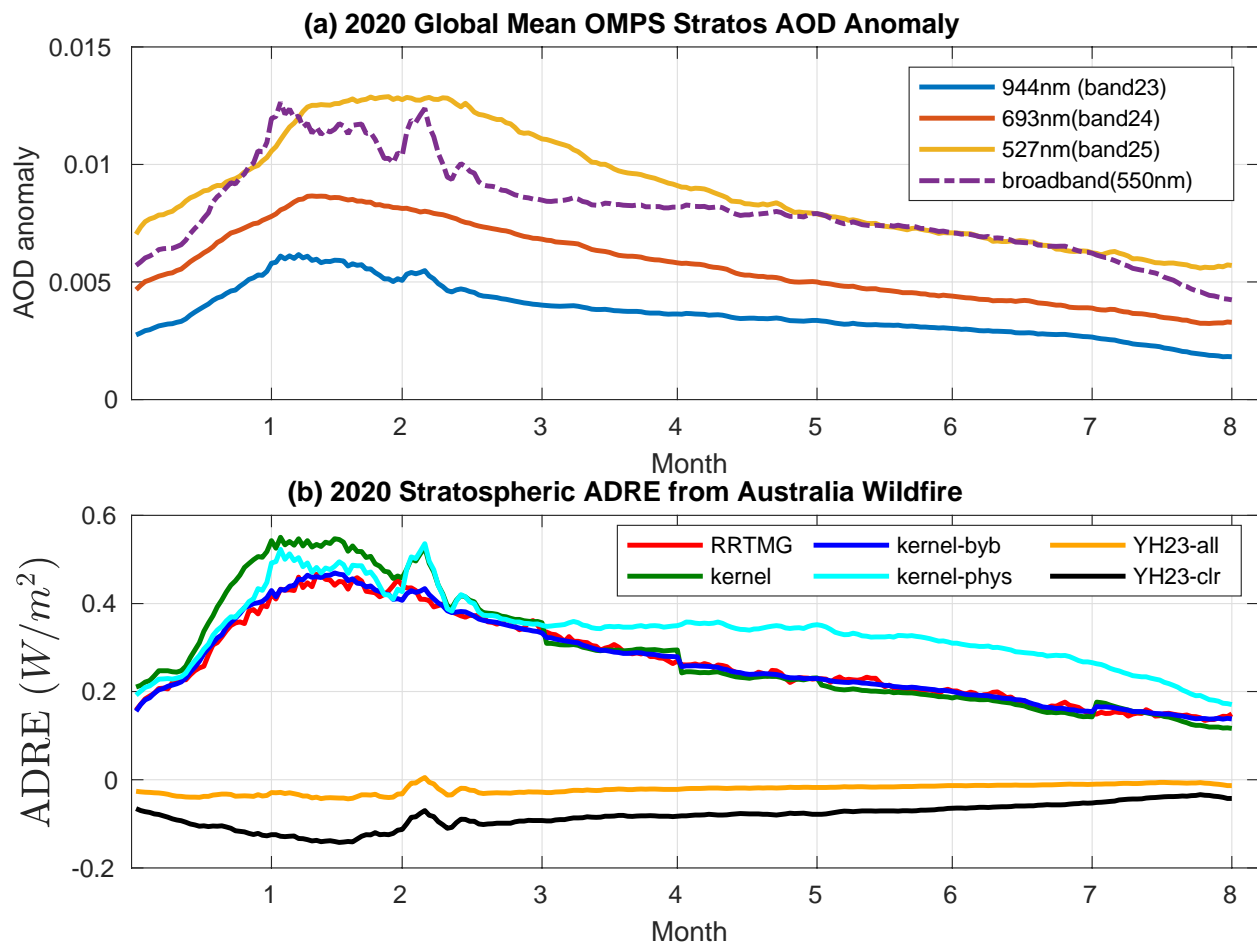
534 Figure 9. Same as Figure 6, but for the year 2020.

535 Figure 9 shows the zonal mean stratospheric AOD anomaly at different wavelengths,
 536 over 8 months starting in January 2020. After several periods of intense fires in early January
 537 2020, the stratospheric AOD reached a maximum in early February and decayed afterward. The
 538 delay in reaching the AOD peaks might be due to the subsequent self-lofting of upper
 539 tropospheric aerosols, caused by buoyancy changes from the aerosols absorbing solar radiation
 540 (Ohneiser et al., 2020). After being lifted, aerosols spread equatorward and dilute significantly,
 541 leading to a decrease in the stratospheric AOD anomaly which lasts around eight months. Same
 542 as Figure 6e, Figure 9e also suggests that the AOD wavelength dependency relationship is
 543 complex, and assuming a certain Angstrom exponent may not be sufficient to represent the band-
 544 by-band AOD.

545 To investigate the performance of kernels in quantifying wildfire-related events, we
 546 assume the stratospheric aerosol anomaly consists of aged biomass burning aerosols with a
 547 single scattering albedo at 550 nm of 0.86 as suggested by recent studies (Damany-Pearce et al.,
 548 2022; Ohneiser et al., 2020). Figure 10a shows the global mean stratospheric AOD anomaly in
 549 the mid-visible bands in 2020. The interpolated AOD at RRTMG bands shows distinct

550 differences compared to the broadband AOD, especially in February. Broadband AOD is
551 calculated from scaling 869 nm to 550 nm assuming an AE of 1. Stratospheric ADRE is further
552 calculated with known AOD and SSA. Figure 10b shows ADRE calculated from the RRTMG as
553 well as kernel methods. Results show that using YH23 kernel schemes significantly
554 underestimates the warming effect of stratospheric aerosols. This is because the stratospheric
555 AOD_{abs} kernels are nearly twice as large as the YH-clearsky AOD_{abs} kernels, whereas the
556 AOD_{scat} kernels show similar magnitudes compared with the YH-clearsky values (Figure 5b and
557 Figure S9 in the Supporting Information). These results further emphasize the need to use
558 kernels specifically designed for stratospheric aerosols to accurately quantify ADRE. As a
559 comparison, using the stratospheric aerosol kernels captures wildfire-induced ADRE relatively
560 well. The agreements in ADRE calculations indicate that our stratospheric aerosol kernel dataset
561 is applicable for quantifying ADRE regardless of aerosol types (either scattering or absorbing).
562 For the Australia wildfire, the carbonaceous aerosols led to a peak global mean warming of over
563 $+0.4 \text{ W/m}^2$ in mid-February 2020.

564 Figure 11 shows the comparison of the spatial distributions of ADRE calculated by
 565 different kernel schemes. Similar to the Hunga volcanic eruption case, the band-by-band kernel
 566 quantifications align most closely with the RRTMG-calculated results, with the R^2 being 99.02%
 567 (Table 2). Both the broadband and the physically sorted kernels slightly overestimate the aerosol
 568 warming, especially in the high-latitude regions. The slightly lower performance of broadband
 569 and physically sorted kernels in capturing the spatial patterns of ADRE might come from the
 570 bias in AOD_{550} , as using a uniform AE value for AOD wavelength conversion across the globe
 571 may not be representative. This aligns with other studies indicating that the Angstrom exponent
 572 can vary significantly (Malinina et al., 2019), particularly during wildfire events when the
 573 inclusion of organics can complicate the particle size distribution interpretation (Bourassa et al.,
 574 2019). Overall, the stratospheric kernels can capture more than 83.19% of the variance in the
 575 Australia wildfire-induced ADRE, with the RMSEs less than 0.08 W/m^2 (i.e. 28.57% relative to
 576 the global mean). Both Figure 11 and Table 2 suggest that YH23 kernel schemes are not suitable
 577 for quantifying the wildfire-dominated stratospheric ADRE because of its significant
 578 underestimation of AOD_{abs} kernels.



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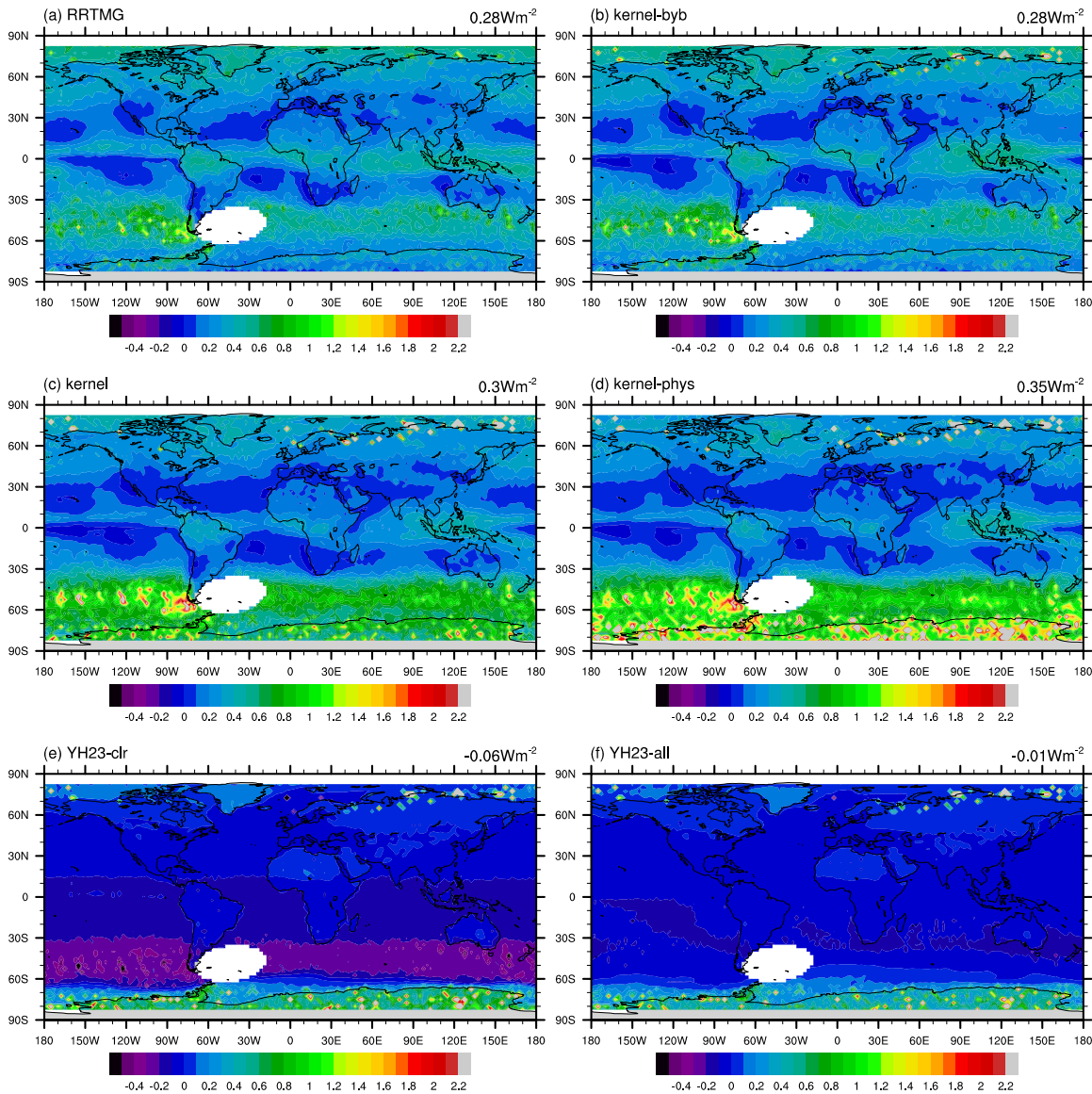
580

Figure 10. Same as Figure 7 but for the 2020 Australia wildfire.

581

582

Stratospheric ADRE from Australia Wildfire



583

584

Figure 11. Same as Figure 8 but for the 2020 Australia wildfire.

585 **5 Conclusions**

586 This paper provides, for the first time, a comprehensive set of radiative kernels for
 587 stratospheric aerosols. The kernels are derived for the scattering and absorbing aerosol optical
 588 depth, respectively, based on partial radiative perturbation (PRP) computations using one year of
 589 3-hourly MERRA-2 data. We analyzed the spatial variability of broadband aerosol kernels,
 590 demonstrating that they can be emulated as a joint function of solar insolation, TOA reflectance,
 591 and stratospheric aerosol optical depth. The developed aerosol radiative kernels provide a
 592 versatile tool for assessing the stratospheric ADRE of different aerosol types.

593 Stratospheric aerosol kernels exhibit significant spatial, temporal, and spectral variability
594 (Figures 1-3). Validation tests have been done to evaluate the aerosol height dependency, linear
595 scaling, and linear additivity of the kernels (Figures S4-S6). On a global scale, a 0.1 increase in
596 stratospheric AOD_{scat}^{550} leads to a cooling effect of -2.65 W/m^2 at the TOA, while a similar
597 increase in AOD_{abs}^{550} results in a warming effect of $+41.95 \text{ W/m}^2$ (Figure 1). The magnitude of
598 stratospheric aerosol kernels is greater than that for tropospheric aerosols (e.g., Thorsen et al.,
599 2020; Q. Yu & Huang, 2023b), particularly for absorbing aerosols. This is due to the higher
600 placement of aerosols, which interact with radiation less attenuated by clouds or tropospheric
601 absorbers. Additionally, underlying clouds enhance the brightness of dark surfaces, thus
602 amplifying the sensitivity of stratospheric ADRE to absorbing aerosols. Band-by-band aerosol
603 kernels were calculated for the 14 bands in RRTMG SW, with spectral signatures indicating
604 peak sensitivity from the near ultraviolet to the near-infrared (bands 23 to 25, 442 nm-1242 nm)
605 (Figure 3). Using discrete AOD observations at these wavelengths allows for a more accurate
606 constraint on ADRE.

607 From the single-layer aerosol analytical model, we identified that broadband aerosol
608 kernels are related to TOA insolation, tropospheric reflectance, and stratospheric aerosol optical
609 depth. We proposed a physically sorted set of aerosol kernels using a multivariate regression
610 model, which can effectively reproduce the RRTMG-calculated broadband kernels (Figure 4 &
611 S8). These physically sorted kernels are independent from geophysical location and can provide
612 first-order estimations of stratospheric ADRE using satellite measurements.

613
614 The kernels were applied to calculate the ADRE for two stratospheric aerosol injection
615 events: the 2022 Hunga volcanic eruption and the 2020 Australia wildfire. There is overall good
616 agreement between the RRTMG-calculated results and those obtained using the kernels (Figures
617 7-8, Figures 10-11, Table 2). Band-by-band kernels perform best by constraining the wavelength
618 dependency of AOD. Using band-by-band kernels can reproduce 99% of the ADRE variance
619 with relative errors of less than 4%. Using other stratospheric kernels can capture more than 90%
620 of the variance with relative errors of less than 10% (Table 2), despite the uncertainty in AE. The
621 stratospheric ADRE from the 2022 Hunga eruption peaked six months after the event, inducing a
622 global mean cooling of -0.46 W/m^2 (Figures 7-8). For the 2020 Australia wildfire, the
623 stratospheric ADRE peaks one month after the event and results in a global mean warming of
624 $+0.28 \text{ W/m}^2$ from January to August (Figures 10-11).

625 To accurately calculate stratospheric ADRE, users are recommended to use the band-by-
626 band kernels when reliable spectral AOD data is available. If such information is unavailable, the
627 broadband aerosol kernels can be used alternatively, although the results should be used with
628 caution as the broadband kernels are calculated based on an assumed Angstrom Exponent of 1.
629 The physically sorted kernels have the advantage of not being restricted to specific geographical
630 locations. With climate change, the aerosol-related and environmental conditions at a location
631 may change. In such cases, physically sorted kernels may have an advantage for the ADRE
632 quantification. Considering that there may be rapid stratospheric temperature adjustments in
633 response to the instantaneous perturbations of the aerosols, the kernels developed in this work
634 can be extended to include the radiative effects of such adjustments in future work, to provide an
635 estimation of the effective (adjusted) radiative effect of the stratospheric aerosols.

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639 research.

640 **Data Availability Statement**

641 The Modern-Era Retrospective Analysis for Research and Applications, Version 2 (MERRA-2)
642 data used to calculate stratospheric aerosol kernels are provided by NASA Global Modeling and
643 Assimilation Office (Randles et al., 2017). The aerosol mixing ratio, assimilated meteorological
644 fields, radiation and aerosol diagnostics data are available at Global Modeling and Assimilation
645 Office via <https://doi.org/10.5067/LTVB4GPCOTK2>,
646 <https://doi.org/10.5067/WWQSQ8IVFW8>, <https://doi.org/10.5067/Q9QMY5PBNV1T> and
647 <https://doi.org/10.5067/KLICLTZ8EM9D> (Gelaro et al., 2017). The dataset of stratospheric
648 aerosol direct radiative effect kernels (monthly mean broadband, band-by-band, and the annual
649 mean physically sorted ones), along with the scripts and data to reproduce the findings in this
650 paper, are available on Mendeley Data at <https://data.mendeley.com/datasets/t87tfnk2xd/1>
651

652 **References**

- 653 Ångström, A. (1929). On the atmospheric transmission of sun radiation and on dust in the air.
654 *Geografiska Annaler*, *11*(2), 156–166.
- 655 Ayash, T., Gong, S. L., Jia, C. Q., Huang, P., Zhao, T. L., & Lavoue, D. (2008). Global modeling
656 of multicomponent aerosol species: Aerosol optical parameters. *Journal of Geophysical*
657 *Research: Atmospheres*, *113*(D12). <https://doi.org/10.1029/2007JD008968>
- 658 Balmes, K. A., & Fu, Q. (2021). All-Sky Aerosol Direct Radiative Effects at the ARM SGP Site.
659 *Journal of Geophysical Research: Atmospheres*, *126*(17), e2021JD034933.
660 <https://doi.org/10.1029/2021JD034933>
- 661 Bourassa, A. E., Rieger, L. A., Zawada, D. J., Khaykin, S., Thomason, L. W., & Degenstein, D.
662 A. (2019). Satellite Limb Observations of Unprecedented Forest Fire Aerosol in the
663 Stratosphere. *Journal of Geophysical Research: Atmospheres*, *124*(16), 9510–9519.
664 <https://doi.org/10.1029/2019JD030607>

- 665 Bourassa, A. E., Zawada, D. J., Rieger, L. A., Warnock, T. W., Toohey, M., & Degenstein, D. A.
666 (2023). Tomographic Retrievals of Hunga Tonga-Hunga Ha'apai Volcanic Aerosol.
667 *Geophysical Research Letters*, *50*(3), e2022GL101978.
668 <https://doi.org/10.1029/2022GL101978>
- 669 Chauvigné, A., Waquet, F., Auriol, F., Blarel, L., Delegove, C., Dubovik, O., Flamant, C.,
670 Gaetani, M., Goloub, P., Loisil, R., Mallet, M., Nicolas, J.-M., Parol, F., Peers, F., Torres,
671 B., & Formenti, P. (2021). Aerosol above-cloud direct radiative effect and properties in
672 the Namibian region during the AEROSOL, RadiatiON, and CLOUDS in southern Africa
673 (AEROCLO-sA) field campaign – Multi-Viewing, Multi-Channel, Multi-Polarization
674 (3MI) airborne simulator and sun photometer measurements. *Atmospheric Chemistry and*
675 *Physics*, *21*(10), 8233–8253. <https://doi.org/10.5194/acp-21-8233-2021>
- 676 Chen, Z., Bhartia, P. K., Torres, O., Jaross, G., Loughman, R., DeLand, M., Colarco, P.,
677 Damadeo, R., & Taha, G. (2020). Evaluation of the OMPS/LP stratospheric aerosol
678 extinction product using SAGE III/ISS observations. *Atmospheric Measurement*
679 *Techniques*, *13*(6), 3471–3485. <https://doi.org/10.5194/amt-13-3471-2020>
- 680 Chlek, P., & Coakley Jr, J. A. (1974). Aerosols and climate. *Science*, *183*(4120), 75–77.
- 681 Damany-Pearce, L., Johnson, B., Wells, A., Osborne, M., Allan, J., Belcher, C., Jones, A., &
682 Haywood, J. (2022). Australian wildfires cause the largest stratospheric warming since
683 Pinatubo and extends the lifetime of the Antarctic ozone hole. *Scientific Reports*, *12*(1),
684 12665. <https://doi.org/10.1038/s41598-022-15794-3>
- 685 Datseris, G., Blanco, J., Hadas, O., Bony, S., Caballero, R., Kaspi, Y., & Stevens, B. (2022).
686 Minimal Recipes for Global Cloudiness. *Geophysical Research Letters*, *49*(20),
687 e2022GL099678. <https://doi.org/10.1029/2022GL099678>

- 688 Flynn, L. E., Seftor, C. J., Larsen, J. C., & Xu, P. (2006). The ozone mapping and profiler suite.
689 In *Earth Science Satellite Remote Sensing: Vol. 1: Science and Instruments* (pp. 279–
690 296). Springer.
- 691 Fromm, M., Lindsey, D. T., Servranckx, R., Yue, G., Trickl, T., Sica, R., Doucet, P., & Godin-
692 Beekmann, S. (2010). The untold story of pyrocumulonimbus. *Bulletin of the American*
693 *Meteorological Society*, *91*(9), 1193–1210.
- 694 Gao, J., Huang, Y., Peng, Y., & Wright, J. S. (2023). Aerosol Effects on Clear-Sky Shortwave
695 Heating in the Asian Monsoon Tropopause Layer. *Journal of Geophysical Research:*
696 *Atmospheres*, *128*(4), e2022JD036956. <https://doi.org/10.1029/2022JD036956>
- 697 Gelaro, R., McCarty, W., Suárez, M. J., Todling, R., Molod, A., Takacs, L., Randles, C. A.,
698 Darmenov, A., Bosilovich, M. G., Reichle, R., Wargan, K., Coy, L., Cullather, R., Draper,
699 C., Akella, S., Buchard, V., Conaty, A., Silva, A. M. da, Gu, W., ... Zhao, B. (2017). The
700 Modern-Era Retrospective Analysis for Research and Applications, Version 2 (MERRA-
701 2). *Journal of Climate*, *30*(14), 5419–5454. <https://doi.org/10.1175/JCLI-D-16-0758.1>
702 [dataset]
- 703 Giles, D. M., Sinyuk, A., Sorokin, M. G., Schafer, J. S., Smirnov, A., Slutsker, I., Eck, T. F.,
704 Holben, B. N., Lewis, J. R., Campbell, J. R., Welton, E. J., Korkin, S. V., & Lyapustin, A.
705 I. (2019). Advancements in the Aerosol Robotic Network (AERONET) Version 3
706 database – automated near-real-time quality control algorithm with improved cloud
707 screening for Sun photometer aerosol optical depth (AOD) measurements. *Atmospheric*
708 *Measurement Techniques*, *12*(1), 169–209. <https://doi.org/10.5194/amt-12-169-2019>

- 709 Grinsted, A., Moore, J. C., & Jevrejeva, S. (2007). Observational evidence for volcanic impact
710 on sea level and the global water cycle. *Proceedings of the National Academy of Sciences*,
711 *104*(50), 19730–19734. <https://doi.org/10.1073/pnas.0705825104>
- 712 Günther, M., Schmidt, H., Timmreck, C., & Toohey, M. (2024). Why does stratospheric aerosol
713 forcing strongly cool the warm pool? *Atmospheric Chemistry and Physics*, *24*(12), 7203–
714 7225. <https://doi.org/10.5194/acp-24-7203-2024>
- 715 Hansen, J. (2005). Efficacy of climate forcings. *Journal of Geophysical Research*, *110*(D18),
716 D18104. <https://doi.org/10.1029/2005JD005776>
- 717 Haywood, J. M., & Shine, K. P. (1995). The effect of anthropogenic sulfate and soot aerosol on
718 the clear sky planetary radiation budget. *Geophysical Research Letters*, *22*(5), 603–606.
719 <https://doi.org/10.1029/95GL00075>
- 720 Heald, C. L., Ridley, D. A., Kroll, J. H., Barrett, S. R. H., Cady-Pereira, K. E., Alvarado, M. J.,
721 & Holmes, C. D. (2014). Contrasting the direct radiative effect and direct radiative
722 forcing of aerosols. *Atmospheric Chemistry and Physics*, *14*(11), 5513–5527.
723 <https://doi.org/10.5194/acp-14-5513-2014>
- 724 Huang, Y., Tan, X., & Xia, Y. (2016). Inhomogeneous radiative forcing of homogeneous
725 greenhouse gases: Inhomogeneous Forcing of Homogeneous Gas. *Journal of*
726 *Geophysical Research: Atmospheres*, *121*(6), 2780–2789.
727 <https://doi.org/10.1002/2015JD024569>
- 728 Kloss, C., Sellitto, P., Renard, J.-B., Baron, A., Bègue, N., Legras, B., Berthet, G., Briaud, E.,
729 Carboni, E., Duchamp, C., Duflot, V., Jacquet, P., Marquestaut, N., Metzger, J.-M.,
730 Payen, G., Ranaivombola, M., Roberts, T., Siddans, R., & Jégou, F. (2022). Aerosol
731 Characterization of the Stratospheric Plume From the Volcanic Eruption at Hunga Tonga

- 732 15 January 2022. *Geophysical Research Letters*, 49(16), e2022GL099394.
733 <https://doi.org/10.1029/2022GL099394>
- 734 Kremser, S., Thomason, L. W., von Hobe, M., Hermann, M., Deshler, T., Timmreck, C., Toohey,
735 M., Stenke, A., Schwarz, J. P., Weigel, R., Fueglistaler, S., Prata, F. J., Vernier, J.-P.,
736 Schlager, H., Barnes, J. E., Antuña-Marrero, J.-C., Fairlie, D., Palm, M., Mahieu, E., ...
737 Meland, B. (2016). Stratospheric aerosol—Observations, processes, and impact on
738 climate. *Reviews of Geophysics*, 54(2), 278–335. <https://doi.org/10.1002/2015RG000511>
- 739 Liu, C.-C., Portmann, R. W., Liu, S., Rosenlof, K. H., Peng, Y., & Yu, P. (2022). Significant
740 Effective Radiative Forcing of Stratospheric Wildfire Smoke. *Geophysical Research*
741 *Letters*, 49(17), e2022GL100175. <https://doi.org/10.1029/2022GL100175>
- 742 Loeb, N. G., Wang, H., Rose, F. G., Kato, S., Smith, W. L., & Sun-Mack, S. (2019).
743 Decomposing Shortwave Top-of-Atmosphere and Surface Radiative Flux Variations in
744 Terms of Surface and Atmospheric Contributions. *Journal of Climate*, 32(16), 5003–
745 5019. <https://doi.org/10.1175/JCLI-D-18-0826.1>
- 746 MacMartin, D. G., Kravitz, B., Tilmes, S., Richter, J. H., Mills, M. J., Lamarque, J.-F., Tribbia, J.
747 J., & Vitt, F. (2017). The Climate Response to Stratospheric Aerosol Geoengineering Can
748 Be Tailored Using Multiple Injection Locations. *Journal of Geophysical Research:*
749 *Atmospheres*, 122(23), 12,574–12,590. <https://doi.org/10.1002/2017JD026868>
- 750 Malinina, E., Rozanov, A., Rieger, L., Bourassa, A., Bovensmann, H., Burrows, J. P., &
751 Degenstein, D. (2019). Stratospheric aerosol characteristics from space-borne
752 observations: Extinction coefficient and Ångström exponent. *Atmospheric Measurement*
753 *Techniques*, 12(7), 3485–3502. <https://doi.org/10.5194/amt-12-3485-2019>

- 754 Martinsson, B. G., Friberg, J., Sandvik, O. S., Hermann, M., van Velthoven, P. F. J., & Zahn, A.
755 (2019). Formation and composition of the UTLS aerosol. *Npj Climate and Atmospheric*
756 *Science*, 2(1), 1–6. <https://doi.org/10.1038/s41612-019-0097-1>
- 757 McComiskey, A., Schwartz, S. E., Schmid, B., Guan, H., Lewis, E. R., Ricchiazzi, P., & Ogren,
758 J. A. (2008). Direct aerosol forcing: Calculation from observables and sensitivities to
759 inputs. *Journal of Geophysical Research*, 113(D9), D09202.
760 <https://doi.org/10.1029/2007JD009170>
- 761 Mlawer, E. J., Iacono, M. J., Pincus, R., Barker, H. W., Oreopoulos, L., & Mitchell, D. L. (2016).
762 Contributions of the ARM Program to Radiative Transfer Modeling for Climate and
763 Weather Applications. *Meteorological Monographs*, 57(1), 15.1-15.19.
764 <https://doi.org/10.1175/AMSMONOGRAPHS-D-15-0041.1>
- 765 Mlawer, E. J., Taubman, S. J., Brown, P. D., Iacono, M. J., & Clough, S. A. (1997). Radiative
766 transfer for inhomogeneous atmospheres: RRTM, a validated correlated-k model for the
767 longwave. *Journal of Geophysical Research: Atmospheres*, 102(D14), 16663–16682.
768 <https://doi.org/10.1029/97JD00237>
- 769 Ohneiser, K., Ansmann, A., Baars, H., Seifert, P., Barja, B., Jimenez, C., Radenz, M., Teisseire,
770 A., Floutsi, A., Haarig, M., Foth, A., Chudnovsky, A., Engelmann, R., Zamorano, F.,
771 Bühl, J., & Wandinger, U. (2020). Smoke of extreme Australian bushfires observed in the
772 stratosphere over Punta Arenas, Chile, in January 2020: Optical thickness, lidar ratios,
773 and depolarization ratios at 355 and 532 nm. *Atmospheric Chemistry and Physics*,
774 20(13), 8003–8015. <https://doi.org/10.5194/acp-20-8003-2020>
- 775 Ohneiser, K., Ansmann, A., Witthuhn, J., Deneke, H., Chudnovsky, A., Walter, G., & Senf, F.
776 (2023). Self-lofting of wildfire smoke in the troposphere and stratosphere: Simulations

- 777 and space lidar observations. *Atmospheric Chemistry and Physics*, 23(4), 2901–2925.
778 <https://doi.org/10.5194/acp-23-2901-2023>
- 779 Randles, C. A., A. M. da Silva, V. Buchard, A. Darmenov, P. R. Colarco, V. Aquila, H. Bian, E.
780 P. Nowottnick, X. Pan, A. Smirnov, H. Yu, and R. Govindaraju, 2016. The MERRA-2
781 Aerosol Assimilation. *NASA Technical Report Series on Global Modeling and Data*
782 *Assimilation, NASA/TM-2016-104606*, Vol. 45, 143 pp. [Available online at
783 <https://gmao.gsfc.nasa.gov/pubs/docs/Randles887.pdf>]
- 784
785 Randles, C. A., Silva, A. M. da, Buchard, V., Colarco, P. R., Darmenov, A., Govindaraju, R.,
786 Smirnov, A., Holben, B., Ferrare, R., Hair, J., Shinozuka, Y., & Flynn, C. J. (2017). The
787 MERRA-2 Aerosol Reanalysis, 1980 Onward. Part I: System Description and Data
788 Assimilation Evaluation. *Journal of Climate*, 30(17), 6823–6850.
789 <https://doi.org/10.1175/JCLI-D-16-0609.1> [dataset]
- 790 Reddy, M. S., Boucher, O., Balkanski, Y., & Schulz, M. (2005). Aerosol optical depths and
791 direct radiative perturbations by species and source type. *Geophysical Research Letters*,
792 32(12). <https://doi.org/10.1029/2004GL021743>
- 793 Robock, A. (2000). Volcanic eruptions and climate. *Reviews of Geophysics*, 38(2), 191–219.
- 794 Schoeberl, M. R., Wang, Y., Taha, G., Zawada, D. J., Ueyama, R., & Dessler, A. (2024a).
795 Evolution of the Climate Forcing During the Two Years After the Hunga Tonga-Hunga
796 Ha’apai Eruption. *Journal of Geophysical Research: Atmospheres*, 129(14),
797 e2024JD041296. <https://doi.org/10.1029/2024JD041296>
- 798 Schoeberl, M. R., Wang, Y., Taha, G., Zawada, D. J., Ueyama, R., & Dessler, A. E. (2024b).
799 *Evolution of the Climate Forcing During the Two Years after the Hunga Tonga-Hunga*
800 *Ha’apai Eruption*. <https://doi.org/10.22541/essoar.171288896.63010190/v1>

- 801 Schoeberl, M. R., Wang, Y., Ueyama, R., Dessler, A., Taha, G., & Yu, W. (2023). The
802 Estimated Climate Impact of the Hunga Tonga-Hunga Ha'apai Eruption Plume.
803 *Geophysical Research Letters*, *50*(18), e2023GL104634.
804 <https://doi.org/10.1029/2023GL104634>
- 805 Schuster, G. L., Dubovik, O., & Holben, B. N. (2006). Angstrom exponent and bimodal aerosol
806 size distributions. *Journal of Geophysical Research: Atmospheres*, *111*(D7).
807 <https://doi.org/10.1029/2005JD006328>
- 808 Sellitto, P., Belhadji, R., Kloss, C., & Legras, B. (2022). Radiative impacts of the Australian
809 bushfires 2019–2020 – Part 1: Large-scale radiative forcing. *Atmospheric Chemistry and*
810 *Physics*, *22*(14), 9299–9311. <https://doi.org/10.5194/acp-22-9299-2022>
- 811 Sellitto, P., Belhadji, R., Legras, B., Podglajen, A., & Duchamp, C. (2024). *The optical*
812 *properties of stratospheric aerosol layer perturbation of the Hunga volcano eruption of*
813 *January 15th, 2022*. <https://doi.org/10.5194/egusphere-2024-1433>
- 814 Strahan, S. E., Oman, L. D., Douglass, A. R., & Coy, L. (2015). Modulation of Antarctic vortex
815 composition by the quasi-biennial oscillation. *Geophysical Research Letters*, *42*(10),
816 4216–4223. <https://doi.org/10.1002/2015GL063759>
- 817 Taha, G., Loughman, R., Colarco, P. R., Zhu, T., Thomason, L. W., & Jaross, G. (2022).
818 Tracking the 2022 Hunga Tonga-Hunga Ha'apai Aerosol Cloud in the Upper and Middle
819 Stratosphere Using Space-Based Observations. *Geophysical Research Letters*, *49*(19),
820 e2022GL100091. <https://doi.org/10.1029/2022GL100091>
- 821 Taha, G., Loughman, R., Zhu, T., Thomason, L., Kar, J., Rieger, L., & Bourassa, A. (2021).
822 OMPS LP Version 2.0 multi-wavelength aerosol extinction coefficient retrieval algorithm.

- 823 *Atmospheric Measurement Techniques*, 14(2), 1015–1036. <https://doi.org/10.5194/amt->
824 14-1015-2021
- 825 Thorsen, T. J., Ferrare, R. A., Kato, S., & Winker, D. M. (2020). Aerosol Direct Radiative Effect
826 Sensitivity Analysis. *Journal of Climate*, 33(14), 6119–6139.
827 <https://doi.org/10.1175/JCLI-D-19-0669.1>
- 828 Visioni, D., MacMartin, D. G., Kravitz, B., Lee, W., Simpson, I. R., & Richter, J. H. (2020).
829 Reduced Poleward Transport Due to Stratospheric Heating Under Stratospheric Aerosols
830 Geoengineering. *Geophysical Research Letters*, 47(17), e2020GL089470.
831 <https://doi.org/10.1029/2020GL089470>
- 832 Wang, X., Randel, W., Zhu, Y., Tilmes, S., Starr, J., Yu, W., Garcia, R., Toon, O. B., Park, M.,
833 Kinnison, D., Zhang, J., Bourassa, A., Rieger, L., Warnock, T., & Li, J. (2023).
834 Stratospheric Climate Anomalies and Ozone Loss Caused by the Hunga Tonga-Hunga
835 Ha’apai Volcanic Eruption. *Journal of Geophysical Research: Atmospheres*, 128(22),
836 e2023JD039480. <https://doi.org/10.1029/2023JD039480>
- 837 Weisenstein, D. K., Keith, D. W., & Dykema, J. A. (2015). Solar geoengineering using solid
838 aerosol in the stratosphere. *Atmospheric Chemistry and Physics*, 15(20), 11835–11859.
839 <https://doi.org/10.5194/acp-15-11835-2015>
- 840 Wu, C.-H., Lee, S.-Y., Tsai, I.-C., Shiu, C.-J., & Chen, Y.-Y. (2023). Volcanic contribution to
841 the 1990s North Pacific climate shift in winter. *Scientific Reports*, 13(1), 5672.
842 <https://doi.org/10.1038/s41598-023-32956-z>
- 843 Yu, P., Portmann, R. W., Peng, Y., Liu, C.-C., Zhu, Y., Asher, E., Bai, Z., Lu, Y., Bian, J., Mills,
844 M., Schmidt, A., Rosenlof, K. H., & Toon, O. B. (2023). Radiative Forcing From the

- 845 2014–2022 Volcanic and Wildfire Injections. *Geophysical Research Letters*, 50(13),
846 e2023GL103791. <https://doi.org/10.1029/2023GL103791>
- 847 Yu, Qiurun (2024), “Quantifying the Direct Radiative Effect of Stratospheric Aerosols Using
848 Radiative Kernels”, Mendeley Data, V1, doi: 10.17632/t87tfnk2xd.1 [dataset]
- 849
850 Yu, Q., & Huang, Y. (2023a). A Dissection of the Inter-Model Spread of the Aerosol Direct
851 Radiative Effect in CMIP6 Models. *Geophysical Research Letters*, 50(21),
852 e2023GL105112. <https://doi.org/10.1029/2023GL105112>
- 853 Yu, Q., & Huang, Y. (2023b). Distributions and Trends of the Aerosol Direct Radiative Effect in
854 the 21st Century: Aerosol and Environmental Contributions. *Journal of Geophysical
855 Research: Atmospheres*, 128(4). <https://doi.org/10.1029/2022JD037716>
- 856 Yu, Q., Jervis, D., & Huang, Y. (2024). Accounting for the effect of aerosols in GHGSat
857 methane retrieval. *Atmospheric Measurement Techniques*, 17(11), 3347–3366.
858 <https://doi.org/10.5194/amt-17-3347-2024>
- 859 Yu, Q.-R., Zhang, F., Li, J., & Zhang, J. (2019). Analysis of sea-salt aerosol size distributions in
860 radiative transfer. *Journal of Aerosol Science*, 129, 71–86.
861 <https://doi.org/10.1016/j.jaerosci.2018.11.014>
862



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Supporting Information for

Quantifying the Direct Radiative Effect of Stratospheric Aerosols Using Radiative Kernels

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Contents of this file

Figures S1 to S9

Tables S1 to S2

Jan 2020 AOD Comparison (40N, 115.625E)

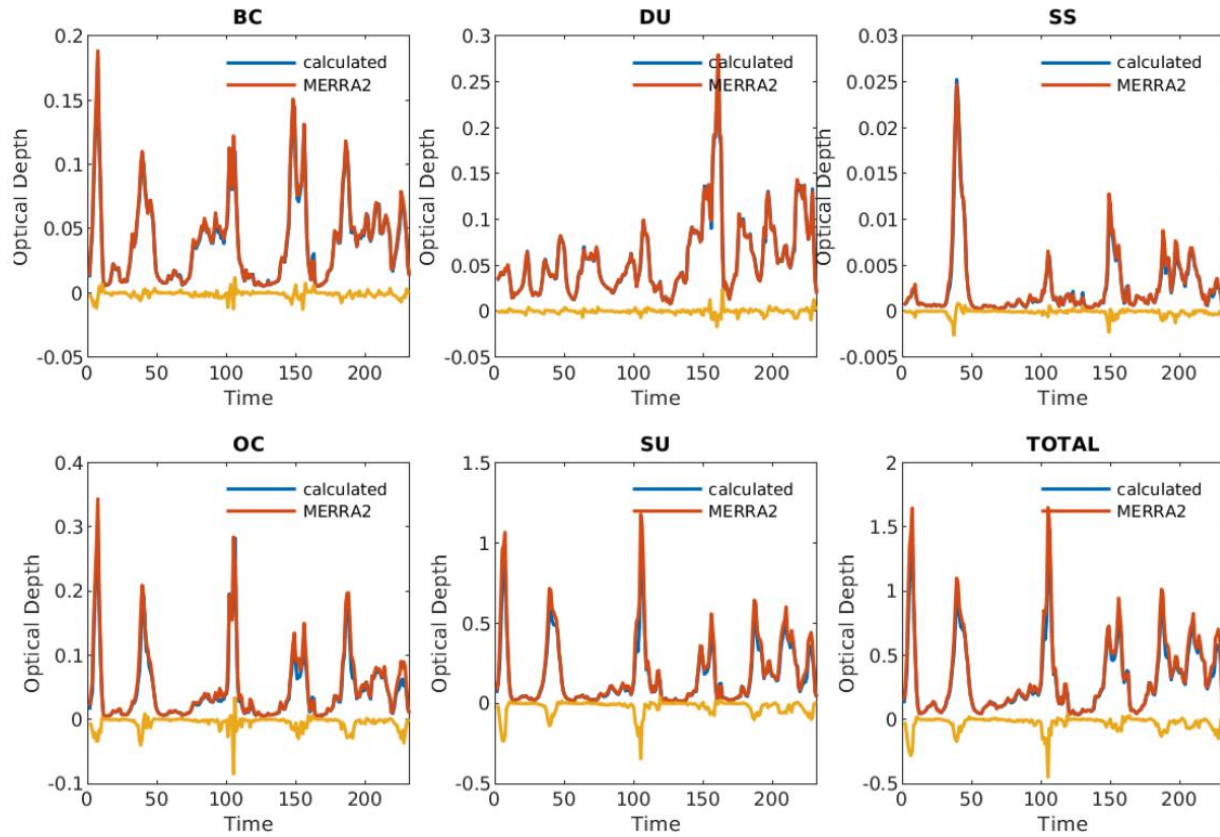


Figure S1. Validation of total column aerosol optical depth (AOD) reconstructed from MERRA-2 aerosol mixing ratio data for black carbon (BC), dust (DU), sea salt (SS), organic carbon (OC), and sulfate (SU) aerosols. The validation is conducted for Beijing, China, in January 2020.

Jan 2020 AOD Comparison (40N, 115.625E)

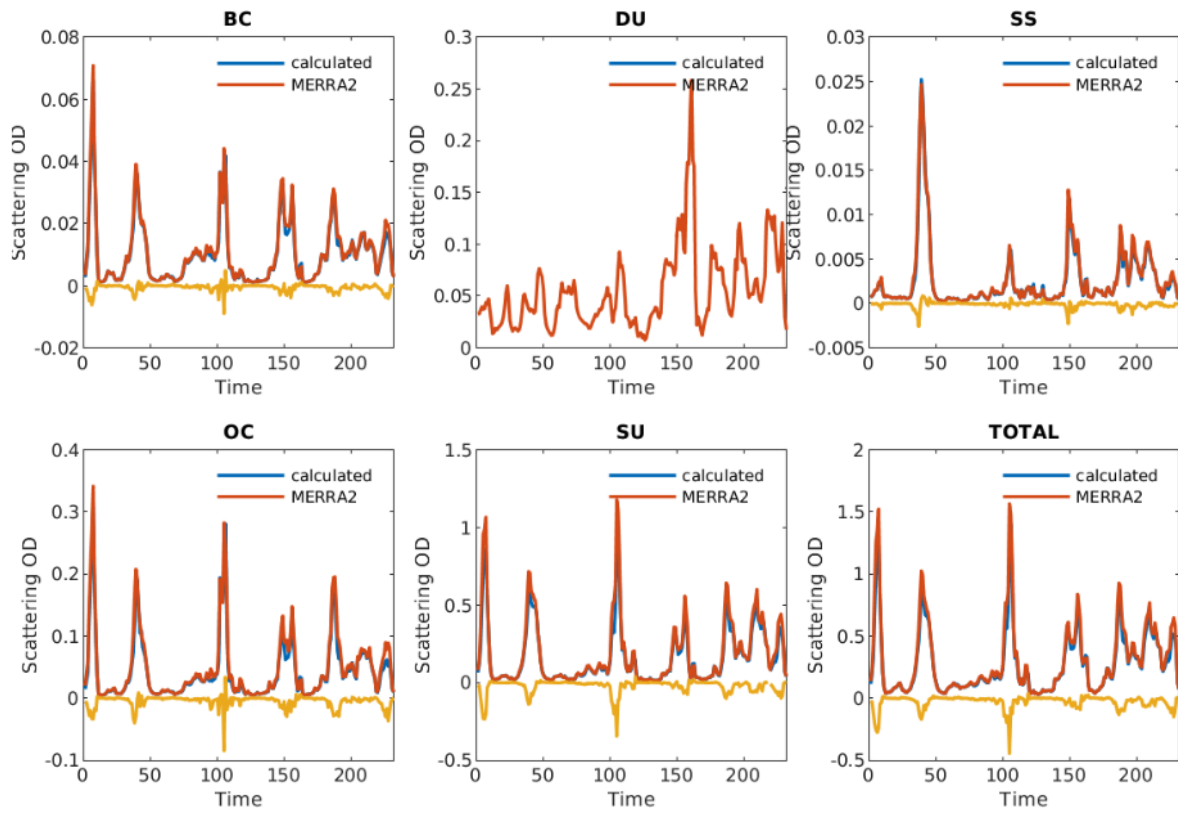


Figure S2. The same as Figure S1 but for total column scattering aerosol optical depth.

Allsky TOA ADRE (2020.0101)

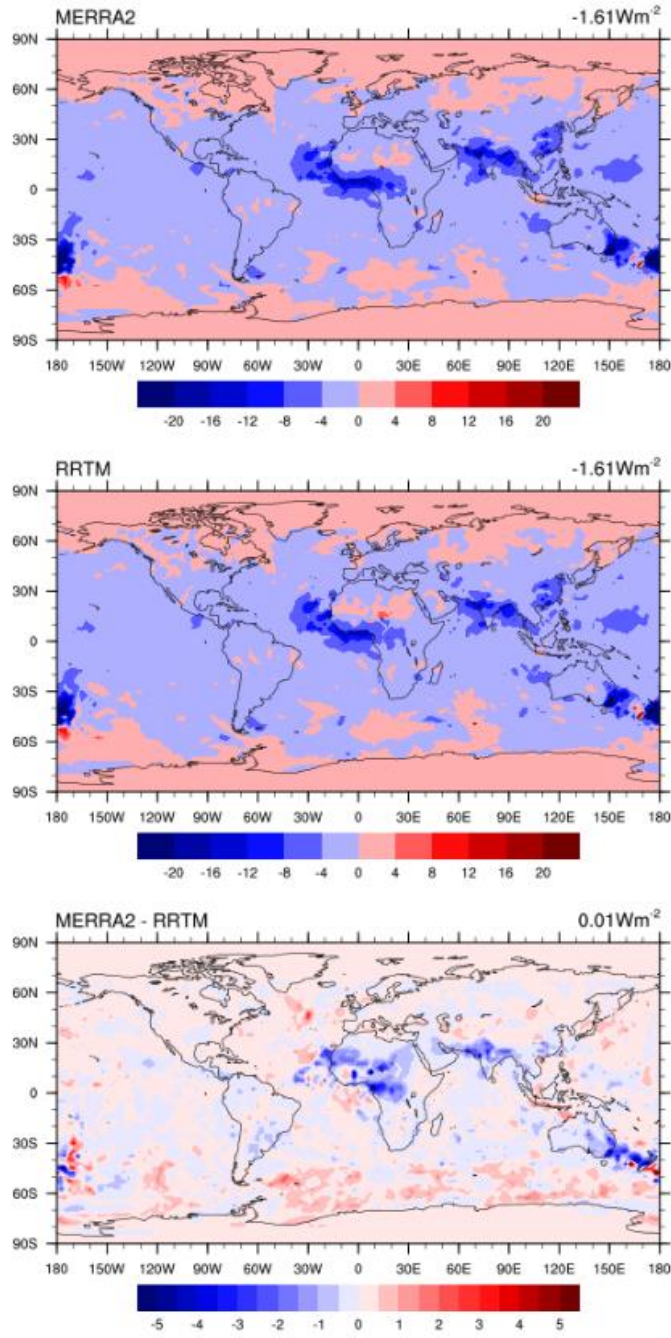


Figure S3. Validation of all-sky total column ADRE calculations using RRMTG for January 1st, 2020. Upper: ADRE from MERRA2; Middle: RRTMG-calculated ADRE; Bottom: Bias.

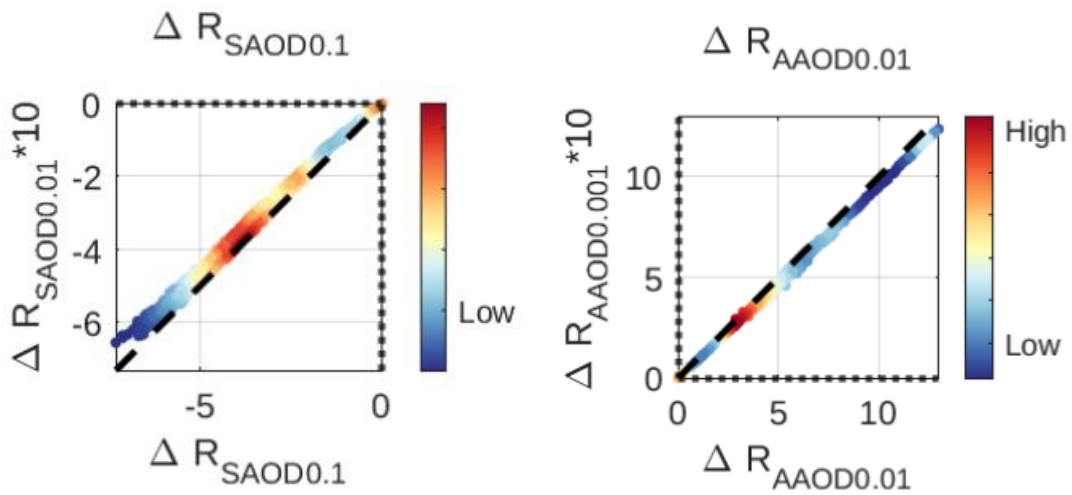


Figure S4. Linear scaling test for broadband Top-Of-Atmosphere (TOA) flux changes (ΔR) in response to perturbations in stratospheric scattering and absorbing AOD. Aerosols are placed at the 1st layer above tropopause. The scattering AOD perturbations are 0.1 and 0.01, while the absorbing AOD perturbations are 0.01 and 0.001, respectively.

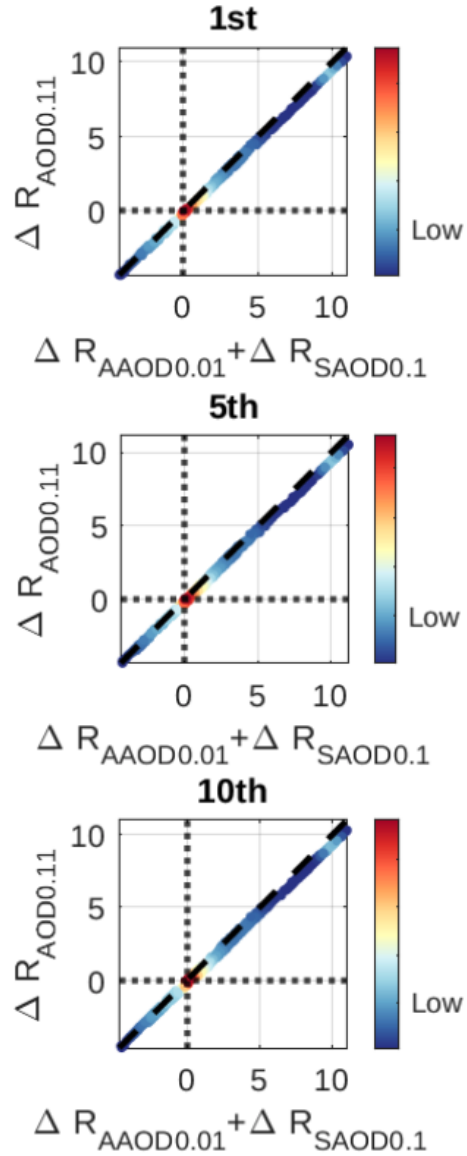


Figure S5. linear additivity test for broadband TOA flux changes (ΔR) in response to perturbations in both stratospheric scattering and absorbing AOD. Aerosols are positioned at the 1st, 5th, and 10th layer above the tropopause, respectively. The perturbations are set to 0.1 for scattering AOD and 0.01 for absorbing AOD. The summed ΔR for scattering and absorbing AOD perturbations shows good agreement with the results from total AOD perturbations.

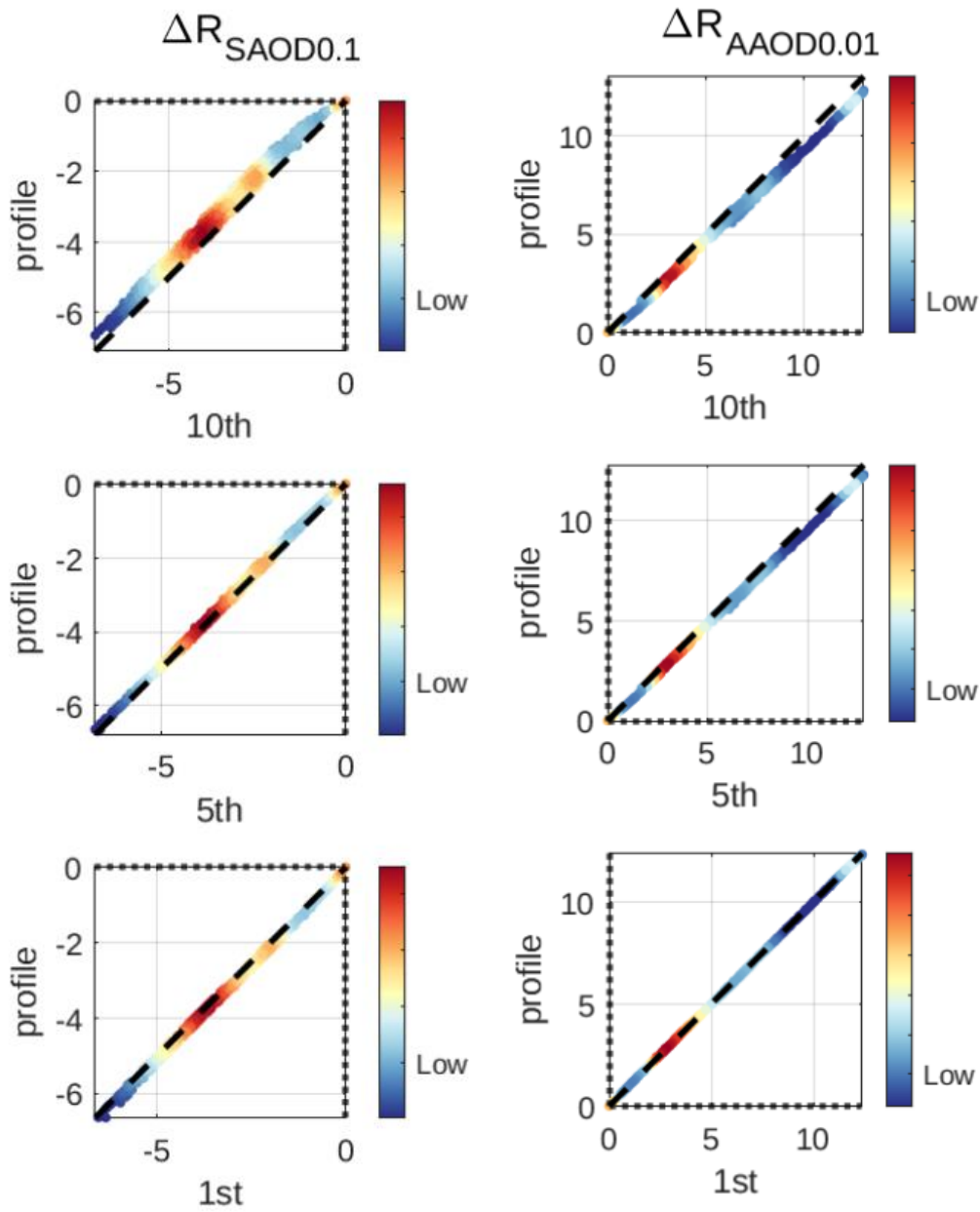


Figure S6. height dependency test for broadband TOA flux changes (ΔR) in response to aerosol perturbation layer height. The ADRE results from perturbing AOD at a single radom layer (e.g., 1st, 5th, 10th above the tropopause) are similar to those obtained from perturbing the entire stratospheric aerosol profiles.

Stratospheric Aerosol Radiative Kernels (broadband)

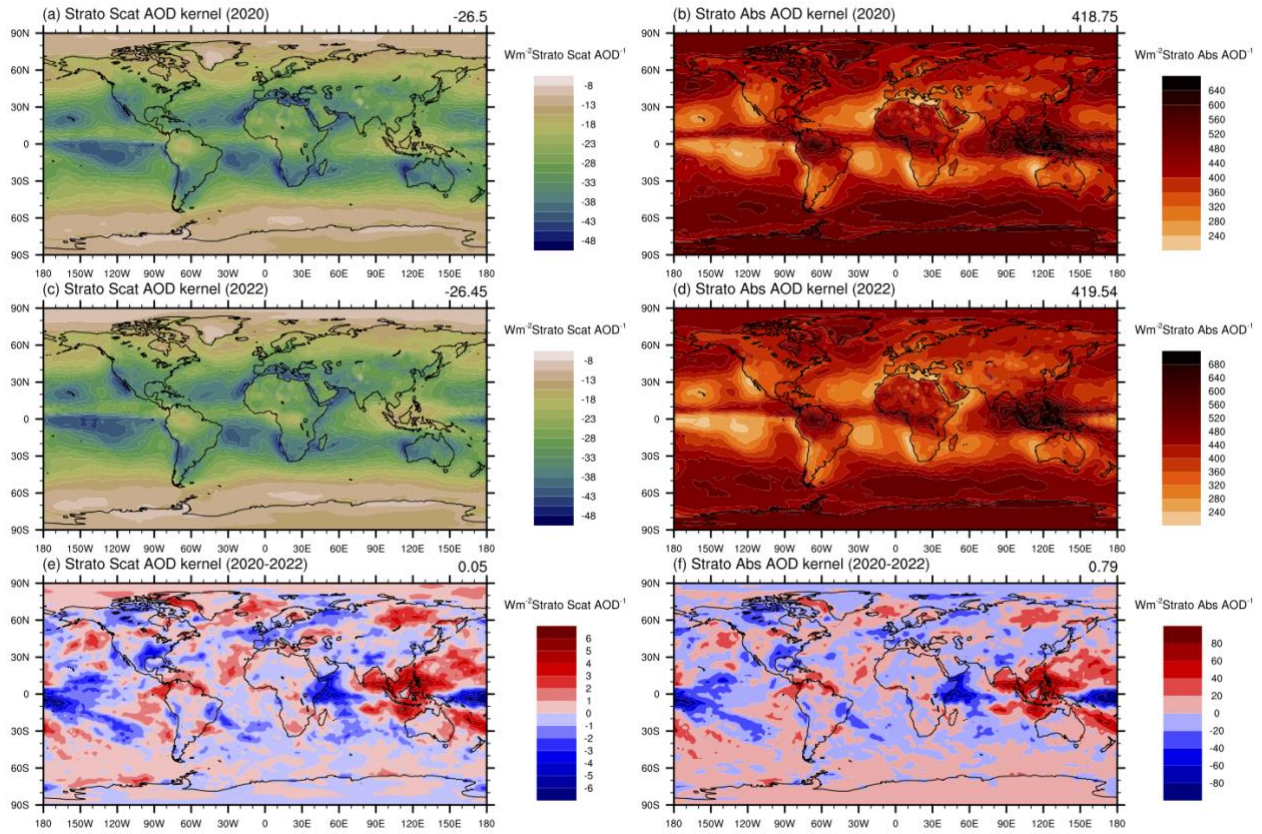


Figure S7. Comparisons between the annual mean stratospheric AOD_{scat} and AOD_{abs} kernels for the years 2020 and 2022. First row: 2020; Middle row: 2022; Bottom row: differences (2020-2022).

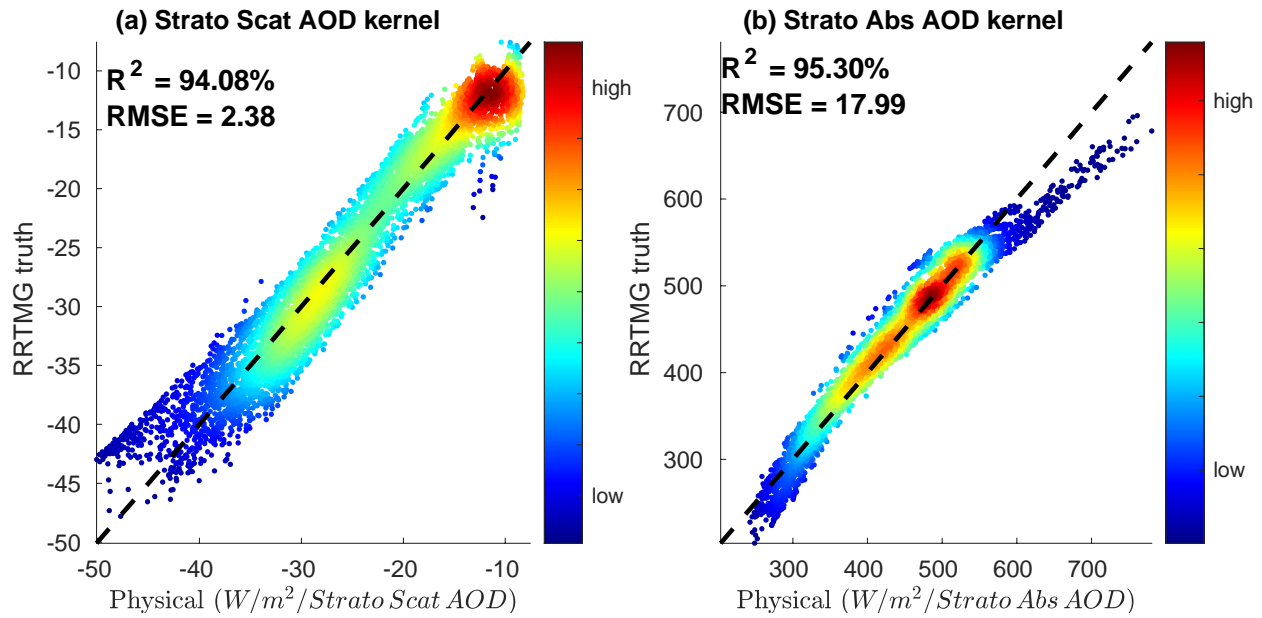


Figure S8. Validations of the physically sorted broadband aerosol kernels for stratospheric AOD_{scat} and AOD_{abs} against benchmark RRTMG calculations.

ADRE Kernel Comparison

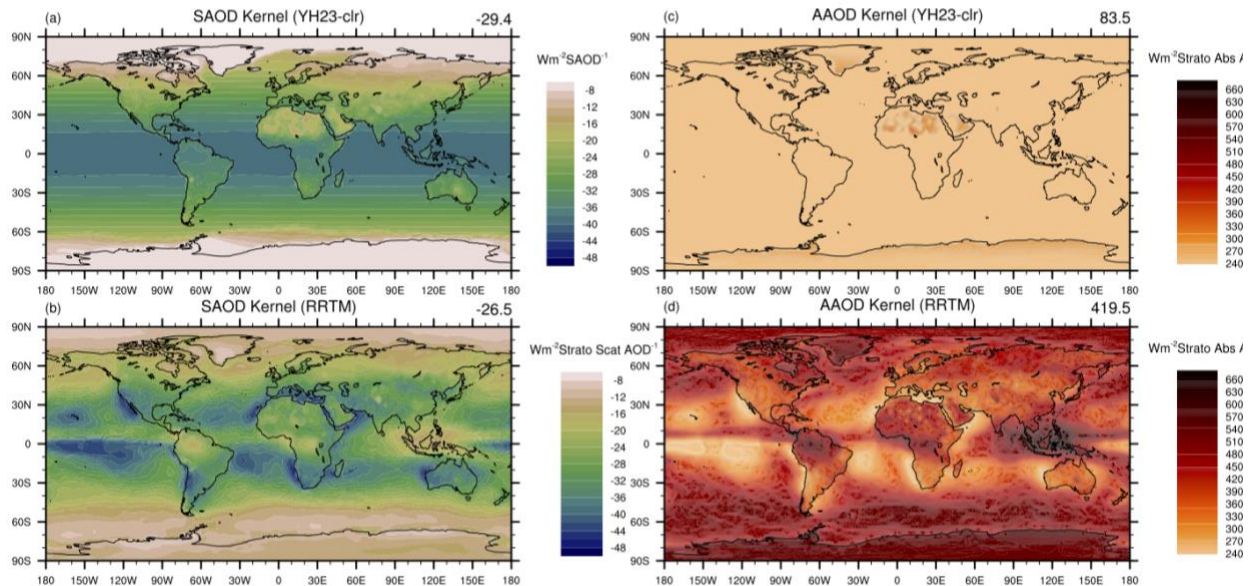


Figure S9. Comparisons between the stratospheric ADRE kernels developed in this work and the YH23 clear-sky kernels for total column aerosols. Left column: kernels for AOD_{scat} ; Right column: kernels for AOD_{abs}

Table 1 Evaluation of predictor performance for SAOD kernels normalized by insolation ($\frac{1}{S} \frac{\partial ADRE}{\partial SAOD}$). R is the TOA reflectance and τ is the stratospheric aerosol optical depth. R² represents the coefficient of determination, and RMSE is the Root Mean Squared Error.

1 predictor	R_s	R_s^2	τ	$R_s \tau$	-	-
R ²	86.04%	84.21%	69.82%	73.18%	-	-
RMSE	3.40	3.63	4.75	4.51	-	-
2 predictors	$(R_s \tau \& (R_s^2))$	$(R_s \tau \& \tau)$	$(R_s \tau \& (R_s))$	$(R_s^2) \& \tau$	$(R_s) \& \tau$	$(R_s^2) \& (R_s)$
R ²	92.50%	83.87%	92.50%	91.72%	93.91%	87.30%
RMSE	2.27	3.55	2.27	2.44	2.02	3.15
3 predictors	$(R_s^2) \& (R_s) \& \tau$	$(R_s^2) \& (R_s) \& (R_s \tau)$	$(R_s^2) \& \tau \& (R_s \tau)$			
R ²	94.01%	93.55%	92.50%	-	-	-
RMSE	rmse=2.99	rmse=3.06	rmse=3.26	-	-	-
4 predictors	$\tau \& R_s^2 \& (R_s) \& (R_s \tau)$	-	-	-	-	-
R ²	94.08%	-	-	-	-	-
RMSE	2.38	-	-	-	-	-

Table 2 Evaluations of predictor performances for AAOD kernels normalized by insolation ($\frac{1}{S} \frac{\partial ADRE}{\partial AAOD}$). R is the TOA reflectance and τ is the stratospheric aerosol optical depth.

1 predictor	R_s	R_s^2	τ	$R_s \tau$	-	-
R ²	74.03%	87.49%	1.97%	50.86%	-	-
RMSE	45.45	29.54	83.21	52.46	-	-
2 predictors	$(R_s \tau) \& (R_s^2)$	$(R_s \tau) \& \tau$	$(R_s \tau) \& (R_s)$	$(R_s^2) \& \tau$	$(R_s) \& \tau$	$(R_s^2) \& (R_s)$
R ²	90.43%	86.32%	82.59%	88.64%	76.37%	89.89%
RMSE	23.96	27.55	32.37	26.59	39.53	25.57
3 predictors	$(R_s^2) \& (R_s) \& \tau$	$(R_s^2) \& (R_s) \& (R_s \tau)$	$(R_s^2) \& \tau \& (R_s \tau)$			
R ²	90.43%	86.32%	82.59%	-	-	-
RMSE	23.96	27.55	32.37	-	-	-
4 predictors	$\tau \& R_s^2 \& (R_s) \& (R_s \tau)$	-	-	-	-	-
R ²	95.30%	-	-	-	-	-
RMSE	17.99	-	-	-	-	-

