

Vertically resolved analysis of the Madden-Julian Oscillation highlights the role of convective transport of moist static energy

Da Yang¹, Lin Yao¹, and Walter Hannah²

¹Department of the Geophysical Sciences, University of Chicago

²Atmospheric, Earth and Energy Division, Lawrence Livermore National Laboratory

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Abstract

We simulate the Madden-Julian oscillation (MJO) over an aquaplanet with uniform surface temperature using the multiscale modeling framework (MMF) configuration of the Energy Exascale Earth System Model (E3SM-MMF). The model produces MJO-like features that have a similar spatial structure and propagation behavior to the observed MJO. To explore the processes involved in the propagation and maintenance of these MJO-like features, we perform a vertically resolved moist static energy (MSE) analysis for the MJO. Unlike the column-integrated MSE analysis, our method emphasizes the local production of MSE variance and quantifies how individual physical processes amplify and propagate the MJO's characteristic vertical structure. We find that radiation, convection, and boundary layer processes all contribute to maintaining the MJO, balanced by the large-scale MSE transport. Furthermore, large-scale dynamics, convection, and boundary layer processes all contribute to the propagation of the MJO, while radiation slows the propagation. Additionally, we perform mechanism-denial experiments to examine the role of radiation and associated feedbacks in simulating the MJO. We find that the MJO can still self-emerge and maintain its characteristic structures without radiative feedbacks. This study highlights the role of convective MSE transport in the MJO dynamics, which was overlooked in the column-integrated MSE analysis.

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3 **moist static energy**

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5 ¹Department of the Geophysical Sciences, University of Chicago, Chicago, IL, USA

6 ²Atmospheric, Earth and Energy Division, Lawrence Livermore National Laboratory, Livermore, CA, USA

7 **Key Points:**

- 8 • We have successfully simulated the Madden-Julian Oscillation (MJO) using the E3SM-
9 MMF over an aquaplanet with uniform surface temperature.
10 • Vertically resolved analyses of moist static energy highlight the role of convection in
11 the maintenance and propagation of the MJO.
12 • Mechanism-denial experiments show that radiative feedbacks are not essential to sim-
13 ulate the MJO.

Corresponding author: Da Yang, dayang@uchicago.edu

14 **Abstract**

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 16 surface temperature using the multiscale modeling framework (MMF) configuration of the
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 19 To explore the processes involved in the propagation and maintenance of these MJO-like
 20 features, we perform a vertically resolved moist static energy (MSE) analysis for the MJO
 21 (Yao et al., 2022). Unlike the column-integrated MSE analysis, our method emphasizes the
 22 local production of MSE variance and quantifies how individual physical processes amplify
 23 and propagate the MJO’s characteristic vertical structure. We find that radiation, convec-
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 26 layer processes all contribute to the propagation of the MJO, while radiation slows the
 27 propagation. Additionally, we perform mechanism-denial experiments to examine the role
 28 of radiation and associated feedbacks in simulating the MJO. We find that the MJO can
 29 still self-emerge and maintain its characteristic structures without radiative feedbacks. This
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 31 overlooked in the column-integrated MSE analysis.

32 **Plain Language Summary**

33 We conduct simulations of the Madden-Julian oscillation (MJO) using a computer
 34 model that can explicitly simulate deep convective clouds. The simulated MJO behaves
 35 similarly to what has been observed in the real world in terms of its spatial structure and
 36 propagation. We then delve into the detailed mechanisms behind the MJO, using a method
 37 that analyzes how energy and moisture move vertically through the atmosphere, rather than
 38 just averaging these properties across the whole atmosphere. This novel analysis shows that
 39 radiation, convection, turbulence in the atmospheric boundary layer, and large-scale at-
 40 mospheric flows all play roles in sustaining the MJO and affect its eastward propagation.
 41 Interestingly, the MJO can still develop and maintain its unique features without the in-
 42 fluence of radiation, indicating other processes are also key. This research underscores the
 43 importance of understanding the vertical transport of energy and moisture by convective
 44 storms in studying the MJO, an aspect previously underappreciated in some simpler models
 45 and diagnoses.

46 **1 Introduction**

47 The Madden-Julian Oscillation (MJO) is a month-long, planetary-scale rainfall pattern
 48 in the tropical atmosphere (C. Zhang, 2005). It often initiates in the Indian Ocean and
 49 then propagates eastward at about 5 m s^{-1} . This propagation speed is about one third of
 50 convectively coupled equatorial Kelvin waves speed and is an order of magnitude smaller
 51 than the dry gravity wave speed in the tropical atmosphere. What provides the energy
 52 to maintain the planetary-scale circulation and rainfall pattern of the MJO? Why does it
 53 propagate eastward? Although the MJO was first discovered in the 1960s, there is still no
 54 consensus on the above questions (C. Zhang et al., 2020; Majda & Stechmann, 2009; Adames
 55 & Kim, 2016; Yang & Ingersoll, 2013, 2014; Wang et al., 2016). This lack of understanding
 56 impeded the progress in simulating the MJO in general circulation models (GCMs) (e.g.,
 57 Wang et al., 2018).

58 A popular method to study the MJO is to diagnose its moist static energy (MSE)
 59 budget (e.g., Andersen & Kuang, 2012; Pritchard & Yang, 2016; Arnold & Randall, 2015).
 60 We define the MSE as $h = c_p T + Lq + gz$, where c_p represents the specific heat capacity of the
 61 air at constant pressure, T represents temperature, L represents latent heat of condensation,

62 q represents specific humidity, g represents gravity acceleration, and z represents altitude.
 63 The MSE budget equation is given by

$$\partial_t h' + \nabla_h \cdot (\bar{u}h)' + \partial_p(\omega h)' = Q', \quad (1)$$

64 where $(\cdot)'$ represents MJO associated quantities, \bar{u} represents horizontal velocity, ω represents
 65 pressure velocity, and Q represents sources and sinks of MSE, including radiation,
 66 convection, boundary-layer turbulence and other sub-grid scale processes. In particular,
 67 Andersen and Kuang (2012) performed vertical integral of this budget, examining the main-
 68 tenance and eastward propagation mechanisms of MJO-associated MSE anomalies. This
 69 analysis framework focuses on horizontal variance of vertically integrated MSE and implic-
 70 itly assumes that MJO's vertical structure is not fundamental to its dynamics. For example,
 71 although Q might have complex vertical structures, it reduces to boundary contributions
 72 after the vertical integral. Let's consider the tendency generated by sub-grid scale vertical
 73 MSE transport Q_c :

$$\int_{p_s}^0 Q_c \frac{dp}{g} = - \int_{p_s}^0 \partial_p F_c \frac{dp}{g} = \frac{F_c|_{p_s} - F_c|_{p=0}}{g} = \frac{F_c|_{p_s}}{g}. \quad (2)$$

74 Here Q_c can include the effects of convection and boundary layer turbulence, p represents
 75 pressure, p_s represents surface pressure, g represents gravity acceleration, F_c represents
 76 MSE fluxes, $F_c|_{p_s}$ and $F_c|_{p=0}$ represent the convective MSE fluxes at the surface and top
 77 of the atmosphere, respectively. At the upper boundary, convective MSE flux is 0. Then,
 78 the vertically integrated contribution due to subgrid-scale vertical MSE transport becomes
 79 equivalent to the surface-flux contribution in this framework. This approach simplifies the
 80 diagnostic process but also revises the conceptual picture. It may appear that proper surface
 81 MSE fluxes would dictate successful MJO simulations regardless of the vertical distribution
 82 of convective MSE transport. This could be misleading.

83 To complement the vertically integrated analysis, we present a vertically resolved MSE
 84 analysis to study the MJO. This framework was first developed to study convective self-
 85 aggregation (Yao et al., 2022; Yao & Yang, 2023) and was subsequently applied to study
 86 tropical cyclones (B. Zhang et al., 2022). The vertically resolved analysis respects the
 87 characteristic vertical structure of the MJO and highlights the importance of convective
 88 MSE transport to the maintenance and eastward propagation of the MJO. We will present
 89 our methods in Section 2, simulation and analysis results in Section 3, and conclusion and
 90 discussion in Section 4.

91 2 The vertically resolved MSE analysis

92 Our diagnostic framework follows Yao and Yang (2023). The underlying assumption of
 93 the analysis is that the MJO has a characteristic vertical structure that is fundamental to the
 94 dynamics of the MJO. Then a physical process that has a positive pattern correlation with
 95 MJO-associated MSE anomaly $h'(x, y, p)$ increases MJO-associated MSE anomaly, thereby
 96 contributing to maintain the MJO. This is analogous to the idea that heating the warm
 97 part of the atmosphere increases available potential energy (Lorenz, 1954; Yang, 2018a).
 98 We project Equation (1) onto $h'(x, y, p)$ and get the contribution that each term makes to
 99 the maintenance of the MJO:

$$\int_{p_T}^{p_s} \frac{dp}{g} \left[\frac{1}{2} \partial_t (\overline{h'^2}) + \overline{h' \nabla_h \cdot (\bar{u}h)'} + \overline{h' \partial_p(\omega h)'} \right] = \int_{p_T}^{p_s} \overline{h' Q'} \frac{dp}{g}, \quad (3)$$

100 where $\overline{(\cdot)}$ represents horizontal average. Then we normalize Equation (3) by the total MSE
 101 variance \mathcal{A} and get the MSE variance budget equation in the unit of growth rate:

$$\text{contribution to growth} = \frac{\text{Equation (3)}}{\mathcal{A}}, \quad (4)$$

102 where

$$\mathcal{A} = \int_{p_T}^{p_S} \frac{h'^2 dp}{g}. \quad (5)$$

103 If the MJO has reached its maintenance stage (*i.e.* statistical equilibrium), $\partial_t h'$ no
 104 longer changes the overall amplitude of the MJO but instead describes the propagation of
 105 the MJO. Therefore, to assess the contribution of each term to this propagation, we project
 106 Equation (1) onto $\partial_t h'$:

$$\text{contribution to propagation} = \frac{\int_{p_T}^{p_S} \overline{\partial_t h' \cdot S} dp / g}{\mathcal{B}}, \quad (6)$$

107 where S represents a given term in Equation (1), and

$$\mathcal{B} = \int_{p_T}^{p_S} \frac{\overline{(\partial_t h')^2} dp}{g}. \quad (7)$$

108 In contrast to the vertically integrated analysis, our approach first calculates the spatial
 109 MSE variance and then performs the vertical integral. This subtle change in the operation
 110 order allows us to objectively diagnose if the vertical distribution of MSE fluxes, e.g., via
 111 convection, makes a significant contribution to the MJO's maintenance and propagation.

112 To the best of our knowledge, there are two major studies that have presented analysis
 113 results explicitly resolving the vertical dimension. Chikira (2014) noticed the limitations of
 114 the vertically integrated framework and performed a detailed budget analysis of the specific
 115 humidity anomalies associated with MJO, in a spirit similar to our study. However, the
 116 author did not quantify the contribution of each process to the development, maintenance,
 117 and propagation of the MJO. Wolding et al. (2016) assumed a weak horizontal temperature
 118 gradient and developed a vertically resolved analysis method for the MJO. That framework
 119 may work well in the free troposphere but does not apply to the boundary layer, where a
 120 substantial horizontal temperature gradient can be sustained.

121 3 Methods

122 3.1 Model Description

123 E3SM was originally forked from the NCAR CESM (Hurrell et al., 2013), but all model
 124 components have undergone significant development since then (Golaz et al., 2019; Xie et
 125 al., 2018). The dynamical core uses a spectral element method on a cubed-sphere geometry
 126 (Ronchi et al., 1996; Taylor et al., 2007). Physics calculations, including the embedded
 127 cloud-resolving models (CRMs) in E3SM-MMF, are performed on a finite volume grid that
 128 is slightly coarser than the dynamics grid, but more closely matches the effective resolution
 129 of the dynamics (Hannah et al., 2021).

130 The multi-scale modelling framework (MMF) configuration of E3SM (E3SM-MMF)
 131 was originally adapted from the super-parameterized CAM (SP-CAM; Khairoutdinov et
 132 al., 2005). E3SM-MMF has also undergone significant development, but still reproduces
 133 the general behavior of its predecessor (Hannah et al., 2020). The embedded CRM in
 134 E3SM-MMF is adapted from the System for Atmospheric Modeling (SAM) (Khairoutdinov
 135 & Randall, 2003), but rewritten in C++ using the performance portability library of Yet

136 Another Kernel Launcher (YAKL) ¹ to facilitate GPU hard acceleration. Microphysical
 137 processes are parameterized with a single moment scheme, and sub-grid scale turbulent
 138 fluxes within the CRM are parameterized using a diagnostic Smagorinsky-type closure.
 139 There is an additional boundary layer scheme outside of the CRM based on Holtslag and
 140 Boville (1993). This allows surface momentum fluxes to be mixed through the boundary
 141 layer prior to calling the global dynamics, which reduces a problematic near-surface wind
 142 bias. Aerosol and ozone concentrations are prescribed with present-day values.

143 E3SM-MMF uses a 60 layer vertical grid with 50 levels in the embedded CRM. The
 144 embedded CRM in E3SM-MMF uses a two-dimensional domain with 64 CRM columns in a
 145 north-south orientation and 2 km horizontal grid spacing. The global physics time step is set
 146 at 20 minutes with a CRM time step of 10 seconds. The CRM variance transport scheme of
 147 Hannah and Pressel (2022) is enabled to reduce grid-scale noise caused by variance trapping
 148 in the CRM (Hannah et al., 2022).

149 3.2 Model Simulations

150 The model was configured for radiative-convective equilibrium (RCE) according to the
 151 RCE model intercomparison project protocol (Wing et al., 2018). This includes globally
 152 homogeneous surface temperature of 300K and globally homogeneous downward shortwave
 153 radiation. Additionally, rotation was enabled to create an equatorial wave guide. This
 154 global RCE setup is a further simplification from aquaplanet simulations with meridional
 155 surface temperature gradient and avoids interference from middle latitude weather systems
 156 (e.g., Hu et al., 2008). The simulations were run for 9 years using 128 nodes of the NERSC
 157 Perlmutter. The standard global cube-sphere grid was used with 30x30 spectral elements
 158 per cube face and 2x2 FV physics cells per element (i.e. ne30pg2), with a physics grid
 159 spacing of 150 km.

160 To explore the role of cloud-radiative feedbacks in the MJO-like phenomena that emerge
 161 in these simulations we use two methods for spatially homogenizing the radiative tendencies.
 162 In the first method, abbreviated as “HomoRad”, we allow radiation to be calculated in every
 163 column at each global model physics step, but before tendencies can be applied to the state,
 164 we calculate a global average at each model level that is then applied to each column. This
 165 method allows the global equilibrium to adjust in time. In the second method, abbreviated
 166 as “FixedRad”, we use fixed profiles of longwave and shortwave radiative heating tendencies
 167 calculated as the global and temporal averaged profiles from the control run, which are then
 168 applied to the state instead of calling the radiation scheme.

169 4 Results

170 E3SM-MMF can successfully simulate the MJO over an aquaplanet with a uniform SST.
 171 Figure 1a plots anomalous outgoing longwave radiation (OLR) of our control simulation.
 172 There are small-scale, short-lived waves that propagate both eastward and westward. In
 173 addition, there are wave envelopes that span about half of the equatorial circumference and
 174 can last longer than 50 days. These large-scale signals are the MJOs and propagate eastward
 175 at about 9 m/s. We then perform a 2D Fourier transform of the OLR anomaly and plot its
 176 power spectrum (Fig. 2a), where the MJO stands out as the most dominant intraseasonal
 177 variability as in observations.

178 We construct the MJO composites following Ma and Kuang (2011). To detect con-
 179 vective centers, we first average the daily OLR data across the latitudinal range of 10°S to
 180 10°N. Then, we filter the meridionally averaged OLR within the MJO’s spectral window,
 181 targeting zonal wavenumbers 1 to 9 and periods from 20 to 100 days (Fig. 2a), using the

¹ see <https://github.com/mrnorman/YAKL> for more information

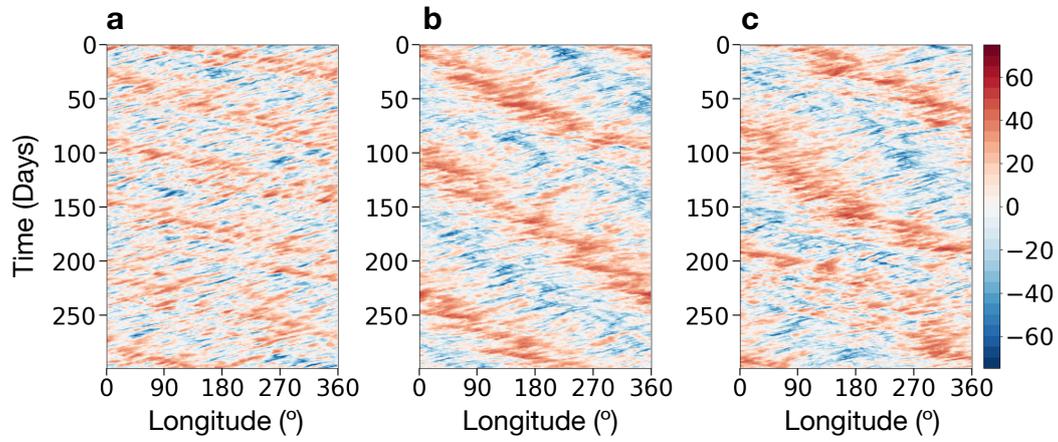


Figure 1. Hovmöller diagrams of OLR anomalies in (a) the control simulation, (b) the HomoRad simulation, and (c) the FixedRad simulation.

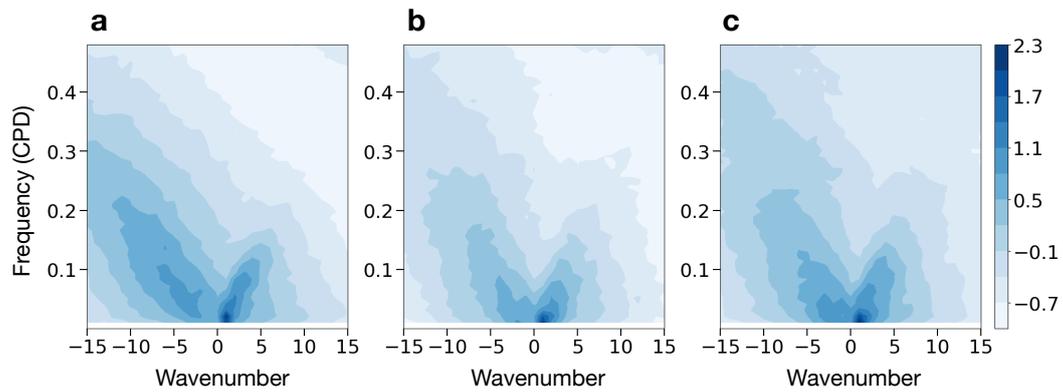


Figure 2. The power spectra of OLR in (a) the control simulation, (b) the HomoRad simulation, and (c) the FixedRad simulation.

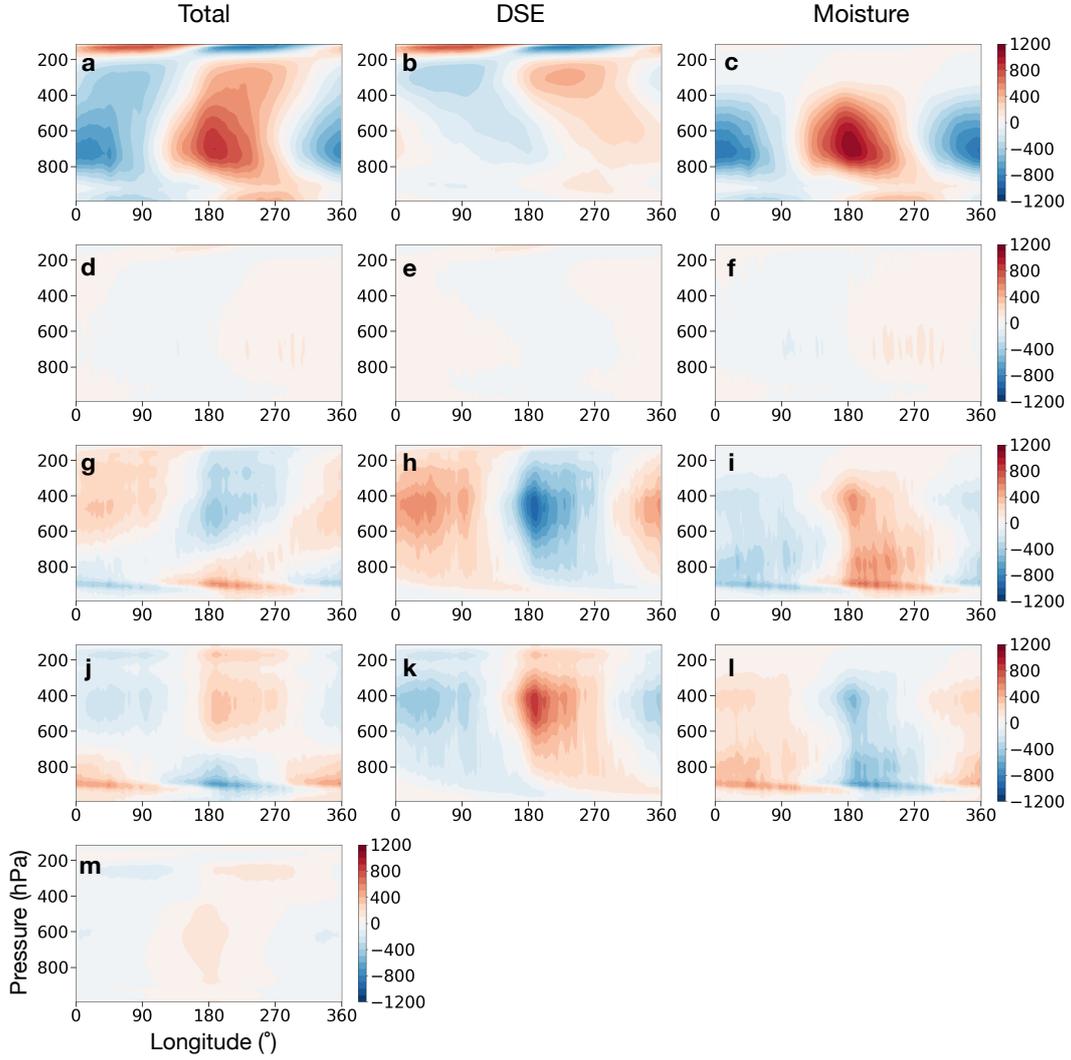


Figure 3. Vertical structures of the MSE budget of the MJO composite. The first column presents terms in Equation (1). Then we decompose these terms into components associated with dry static energy and specific humidity and plot them in the second and third columns. The first row represents MJO’s MSE anomalies (units: J/kg). The second row represents total MSE tendency $\partial_t \text{MSE}$ (units: J/kg/day). The third row represents the effect of large-scale dynamics in transporting and redistributing MSE. The fourth row represents the effect of CRM and boundary layer schemes in transporting and redistributing MSE. The last row represents the effect of radiation. Although convection does not change column-integrated MSE, it does redistribute MSE within a given column and thus changes local MSE.

182 method of Wheeler and Kiladis (1999). Thereafter, the lowest OLR value is identified as
 183 the convective center on each day. In creating the MJO composites, we calculate anomalies
 184 by subtracting zonal averages. Subsequently, we rearrange the data on each day to ensure
 185 that each convective center is strategically placed at 180° , which is the center longitude of
 186 the maps. Finally, we average the rearranged data over time to produce the composites.
 187 For 3D variables in the MSE budget equation, we further average the variables across the
 188 same latitudinal span (10°S to 10°N) to illustrate the characteristic vertical structures.

189 Figure S1a plots the composite OLR and 200-hPa geopotential anomalies of the MJO.
 190 The horizontal structure of the simulated MJO is similar to the observations and model
 191 simulations (Andersen & Kuang, 2012; Arnold & Randall, 2015; Straub & Haertel, 2005).
 192 Both the OLR and geopotential anomalies are symmetric about the equator and with peaks
 193 right at the equator. Negative OLR anomalies correspond to the convective regions of the
 194 MJO, spanning over 150° longitudes and 30° latitudes. The geopotential anomalies show
 195 distinct quadrupole structures, which is often interpreted as a Gill-type response (Gill, 1980)
 196 to a pair of heating and cooling anomalies on the equator, associated with MJO's enhanced
 197 and suppressed convection, respectively.

198 We apply the same compositing technique to analyze the MJO-associated MSE anom-
 199 lies (Fig. 3a). In this process, we rearrange the data so that the convective center of the
 200 MJO is at 180° longitude. There are large-scale ascending motions near the convective
 201 center and descending motions elsewhere. The composite MSE anomalies show a peak in
 202 the lower troposphere, around 650 hPa. Above this level, the MSE anomalies tilt eastward;
 203 below this level, the MSE anomalies tilt westward. Such a distinct vertical structure looks
 204 similar to what we observe in the real tropical atmosphere and may result from a combina-
 205 tion of the first, second, and potentially higher vertical modes (Andersen & Kuang, 2012;
 206 Haertel et al., 2008; Straub & Haertel, 2005). We further decompose the MSE anomalies
 207 into the dry static energy component ($\text{DSE} = c_p T + gz$) and the moisture component (Lq).
 208 The DSE component (Fig. 3b) is generally weaker than the moisture component (Fig. 3c),
 209 except for the upper troposphere. That may result from the weak buoyancy gradient nature
 210 of the tropical atmosphere. There, the Rossby number is large, and the Froude number is
 211 small. That leads to a small horizontal buoyancy gradient (Charney, 1963; Yang & Seidel,
 212 2020; Yang, 2018b; Seidel & Yang, 2020; Yang et al., 2022). Therefore, the vertical structure
 213 of the MSE anomalies mainly follows that of the moisture component.

214 The rest of Figure 3 detail composites of individual terms from the MSE budget equation
 215 (Equation 1), using a uniform color scale across all panels to ease the comparison of their
 216 magnitudes. Figure 3d-f shows the MSE tendency $\partial_t h'$ and its DSE ($\partial_t \text{DSE}$) and moisture
 217 ($\partial_t Lq$) components. The magnitude is weak, and the signal leads the MSE anomaly by
 218 a quarter cycle. This result indicates that the MJO would maintain its amplitude while
 219 propagating eastward. Thus $\partial_t h'$ does not project strongly onto the MSE anomaly but
 220 instead is responsible for the eastward propagation of the MJO.

221 Figure 3g shows the large-scale MSE convergence $-\nabla_{3D} \cdot (\bar{u}h)'$, where $(\bar{u}h)'$ represents
 222 MSE flux due to large-scale circulations. Near the convective center, $-\nabla_{3D} \cdot (\bar{u}h)'$ is positive
 223 in the lower troposphere, suggesting an up-gradient MSE transport; $-\nabla_{3D} \cdot (\bar{u}h)'$ becomes
 224 negative in the upper troposphere due to large-scale winds exporting MSE to the subsiding
 225 regions. To understand this rich vertical structure, we decompose $-\nabla_{3D} \cdot (\bar{u}h)'$ into DSE
 226 and moisture components, both of which have simpler vertical structures (Fig. 3h-i). Large-
 227 scale ascending motions are present over the convective regions, and they adiabatically cool
 228 (Fig. 3h) and moisten (Fig. 3i) the atmosphere throughout the depth of the troposphere,
 229 except for the near-surface levels to the west of the convective center. Meanwhile, large-
 230 scale descending motions are present and adiabatically warm and dry the atmosphere over
 231 clear-sky regions. Therefore, the competing effects of the DSE and moisture components
 232 lead to the complex vertical structure in Fig. 3g. Adiabatic cooling (and warming) proves
 233 more significant in the upper troposphere, whereas moistening (and drying) predominantly
 234 affects the lower troposphere.

235 In E3SM-MMF, the vertical MSE transports by convection and boundary layer tur-
 236 bulence are treated separately, but their effects are similar—redistributing MSE without
 237 changing the column-integrated MSE. Therefore, we consider them as an integral part and
 238 present their MSE transports together (Fig. 3j). As shown in Eqn. (2), convergence of
 239 convective MSE flux yields a positive local MSE tendency, and divergence of convective
 240 MSE flux yields a negative local MSE tendency. The overall effect is to stabilize the air
 241 column by transporting high MSE air from the boundary layer to the free troposphere. This
 242 transport of MSE is associated with convective heating that increases DSE (Fig. 3k) and
 243 condensation that decreases specific humidity (Fig. 3l). Their net effect reduces (increases)
 244 MSE in the lower (upper) troposphere over the convective regions.

245 There appears to be a significant compensation between the convective MSE transport
 246 and large-scale MSE convergence when we compare the third and fourth rows of Fig. 3.
 247 This compensation was discussed in detail by Emanuel et al. (1994). Large-scale flows
 248 converge high-MSE air in the boundary layer toward the convective center. Then, small-
 249 scale turbulence and convection transport the high-MSE air to the free troposphere, where
 250 the large-scale circulation exports it to surrounding areas. As we will discuss, the net
 251 effect of the large-scale and the small-scale parameterized dynamics tends to stabilize the
 252 circulation.

253 To further understand the collective effects of convection and large-scale circulations,
 254 we add the third and fourth rows of Fig. 3 together and get Fig. S2. We reproduce
 255 MJO’s MSE anomaly and its total tendency ($\partial_t \text{MSE}$) in Fig. S2 a & b for convenience.
 256 Fig. S2 c shows the MSE tendency due to convection and circulations, which has a spatial
 257 pattern very similar to that in Fig. S2 b, confirming that convection and circulations are
 258 responsible for MJO’s eastward propagation. We then decompose Fig. S2 c into DSE and
 259 moisture components shown in Fig. S2 d & e, respectively. Their contribution to the DSE
 260 tendency is largely out of phase with the MSE anomaly in Fig. S2 a, suggesting a damping
 261 or stabilizing effect. Their contribution to the moisture tendency leads the MSE anomaly
 262 by a quarter cycle and is in phase with $\partial_t \text{MSE}$, suggesting that the moisture component is
 263 mainly responsible for the propagation.

264 To put this result into context, we assume that the MJO’s MSE anomaly h' has the
 265 form of a propagating wave:

$$h' = \hat{h}e^{ikx + \sigma t},$$

266 where \hat{h} represents MJO’s amplitude, k and σ represent zonal wavenumber and growth rate
 267 respectively. Our diagnosis suggests that convection and circulation’s contribution to the
 268 DSE tendency is associated with the real part of σ , and their contribution to moisture is
 269 associated with the imaginary part.

270 Radiative heating anomalies appear largely in phase with the MSE anomalies, suggest-
 271 ing a positive contribution to maintaining the MJO (Fig. 3m). Over the convective regions,
 272 there is anomalous water vapor and cloud cover leading to positive radiative heating anom-
 273 alies that further support upward motions and convection. This forms a positive feedback
 274 loop between water vapor, clouds, and radiation. Over the dry subsiding regions, radiation
 275 effectively cools the atmosphere and induces subsidence. The amplitude of radiative heating
 276 anomalies is weak but consistent with pressure levels over the convective region. This par-
 277 ticular vertical structure makes radiative heating anomaly project strongly onto the MJO’s
 278 MSE anomaly.

279 We conduct a quantitative analysis of each term’s role in sustaining the MJO by plotting
 280 Equation (4) in Figure 4a. Radiation is identified as the primary factor in the MJO’s
 281 maintenance, aligning with previous findings from column-integrated analyses (Andersen
 282 & Kuang, 2012; Arnold & Randall, 2015; Pritchard & Yang, 2016). Additionally, the net
 283 influence of convection is significant, but its effects on DSE and moisture are opposite. For
 284 example, over the convective areas, drying effects through condensation and precipitation
 285 decrease MSE, while convective heating increases MSE. Large-scale flows tend to reduce the

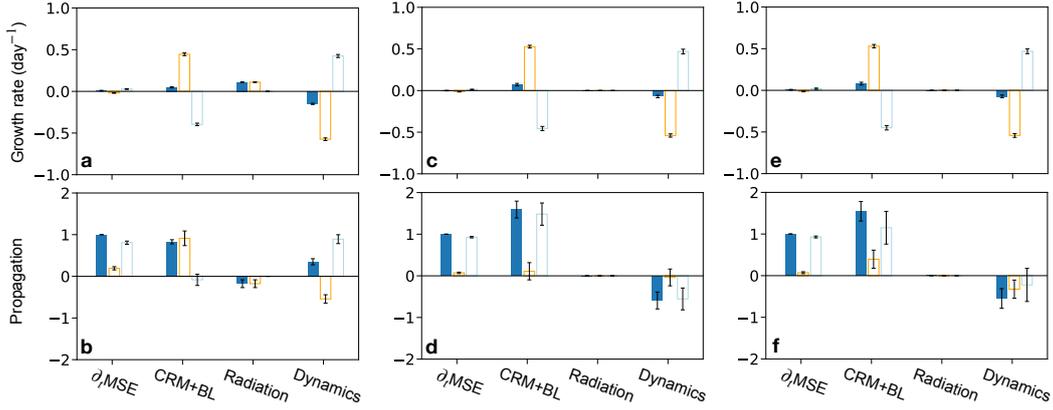


Figure 4. The MSE budget for the control simulation (first column), HomoRad simulation (second column), and FixedRad simulation (third column). The first and second rows show individual terms’ contribution to MJO’s maintenance and propagation, respectively. The solid blue bar represents MSE; the open orange and blue bars represent the corresponding DSE and moisture components, respectively. In the second and third columns, the radiative heating rate is horizontally uniform, so it does not contribute to maintaining and propagating MSE anomalies.

286 MJO amplitude mainly through adiabatically cooling the area with positive MSE anomalies.
 287 Again, the effects of convection and large-scale flows compensate significantly in maintaining
 288 the MJO.

289 We then assess each term’s contribution to MJO’s eastward propagation by plotting
 290 Equation (6) in Fig. 4b. Convection is primarily responsible for the eastward propagation
 291 of the MJO, predominantly through its heating, not its associated drying effect due to
 292 condensation. Radiation slightly hinders the eastward progression. On the other hand,
 293 large-scale dynamics support the eastward propagation, where adiabatic cooling retards the
 294 eastward propagation, and moistening effects favor the eastward propagation.

295 **5 Mechanism-denial experiments**

296 Although radiation is a major factor in maintaining the MJO, our diagnostic results
 297 show that convection’s contribution is of the same order of magnitude (Fig. 4a). This
 298 result motivates mechanism-denial experiments to examine if interactive radiation or ra-
 299 diative feedbacks are essential to the MJO. We perform two experiments, in which we
 300 horizontally homogenize the radiative heating rate and prescribe a uniform radiative heat-
 301 ing rate, respectively (Section 3c). Both methods can effectively decouple radiation from
 302 MJO’s thermodynamic and circulation patterns, switching off a positive feedback. However,
 303 E3SM-MMF can still simulate MJOs without the radiative feedback (Figs. 1, 2), and they
 304 show horizontal structures similar to those in the control simulation (Fig. S1). Although the
 305 simulation results disagree with the global 20-km resolution simulations by Khairoutdinov
 306 and Emanuel (2018), our results agree well with superparameterized GCM results in Arnold
 307 and Randall (2015) and Grabowski (2003). Through our vertically resolved MSE analyses,
 308 we find that convective MSE transport is sufficient to maintain the MJO (Fig. 4c, d).

309 It would be intuitive to assume that the MJO becomes weaker in the mechanism-denial
 310 experiments. However, the MJO’s OLR anomalies seem to become stronger, while wind
 311 anomalies (e.g., Rossby wave gyres) become weaker (Figs. 1, S1). This result may suggest

312 that the maintenance of MJO involves nonlinear processes, and understanding this counter-
313 intuitive result may require a finite-amplitude theory that goes beyond linear analysis.

314 6 Conclusion and Discussion

315 This paper presents MJO simulations over an aquaplanet with a uniform surface tem-
316 perature using E3SM-MMF. The simulated MJOs have similar spatial structures and prop-
317 agation behavior to observations. We consider that the MJO has a characteristic vertical
318 structure that is fundamental to its dynamics. Therefore, to understand the propagation
319 and maintenance of the MJO, we perform a vertically resolved MSE analysis for the MJO
320 (Yao et al., 2022). Our method quantifies how individual physical processes amplify and
321 propagate the characteristic vertical structure of the MJO. Our analyses show that both
322 radiation and convection (CRM + boundary layer) contribute to maintaining the MJO,
323 balanced by the large-scale dynamic transport of MSE. Furthermore, convection is primar-
324 ily responsible for the eastward propagation of the MJO, while radiation may slightly retard
325 the propagation. The diagnostic results seem to suggest that the MJO can still develop and
326 maintain even without interactive radiation and associated feedbacks. This hypothesis is
327 then confirmed by mechanism-denial experiments.

328 Although convection does not change the column-integrated MSE, convection can in-
329 deed change local MSE. Our diagnostic analysis and mechanism-denial experiments highlight
330 the role of convection in both MJO’s maintenance and propagation. This result may appear
331 to agree with a school of studies proposing that convection drives MJO (Yang & Ingersoll,
332 2013, 2014; Wang & Chen, 1989; Wang & Rui, 1990; Majda & Stechmann, 2009). However,
333 in the column integrated analysis, such vertical MSE transport reduces to boundary contri-
334 butions, and thus, the effects would be overlooked by design (see Eq. 2). It’s important to
335 note that several moisture-mode models for the MJO (e.g., Adames & Kim, 2016; Ahmed,
336 2021) predominantly use vertically integrated equations for temperature, moisture, or MSE.
337 This approach may inadvertently fail to capture the complete effects of convection on the
338 MJO’s maintenance and propagation, suggesting a potential area for improvements in these
339 models.

340 Why does super-parameterization of convection lead to better MJO simulations? What
341 does traditional convection parameterization lack in simulating the MJO? These questions
342 have puzzled our field for more than a decade. Our analysis method can help diagnose
343 GCM simulations and address the above questions. We plan to perform E3SM simulations
344 using the same model setup and conduct our vertically resolved MSE analysis for the sim-
345 ulated MJO. We will compare the results from E3SM and E3SM-MMF simulations with a
346 focus on convective MSE transport. It is likely that traditional convection parameteriza-
347 tions cannot efficiently transport boundary layer high-MSE air to the free troposphere and
348 properly distribute it, so convection’s contributions to MJO’s maintenance and propagation
349 are underestimated.

350 7 Open Research

351 The E3SM project, code, simulation configurations, model output, and tools to work
352 with the output are described on its website (<https://e3sm.org>). Instructions on how to
353 get started with running E3SM are available on the website (<https://e3sm.org/model/running-e3sm/e3sm-quick-start>). All code for E3SM may be accessed on the GitHub
354 repository (<https://github.com/E3SM-Project/E3SM>). The raw output data from E3SM-
355 MMF used in this study are archived in the National Energy Research Scientific Comput-
356 ing Center (NERSC). The specific branch used to conduct the simulations can be found at
357 <https://github.com/E3SM-Project/E3SM/tree/whannah/mmf/rce-with-rotation> and is
358 also archived at <https://doi.org/10.5281/zenodo.10989362>. The analysis code and
359 a condensed version of the data needed to reproduce our results are also archived at
360 <https://doi.org/10.5281/zenodo.10998360>.
361

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Supporting Information for ”Vertically resolved analysis of the Madden-Julian Oscillation highlights the role of convective transport of moist static energy”

Da Yang¹, Lin Yao¹, and Walter Hannah²

¹Department of the Geophysical Sciences, University of Chicago, Chicago, IL, USA

²Atmospheric, Earth and Energy Division, Lawrence Livermore National Laboratory, Livermore, CA, USA

Contents of this file

1. Text S1 to S2
2. Figures S1 to S2

Text S1. Horizontal Structures of the MJO composite

Figure S1 shows that the MJO composites in the control and the mechanism-denial experiments share similar horizontal structures.

Text S2. Vertical Structures of the MJO composite

Figure S2 shows the vertical structures of the MJO composite in the control experiment. The results complement Figure 3 by showing the combined contribution of convection

(CRM), boundary layer, and large-scale dynamics to the tendencies of DSE and moisture components.

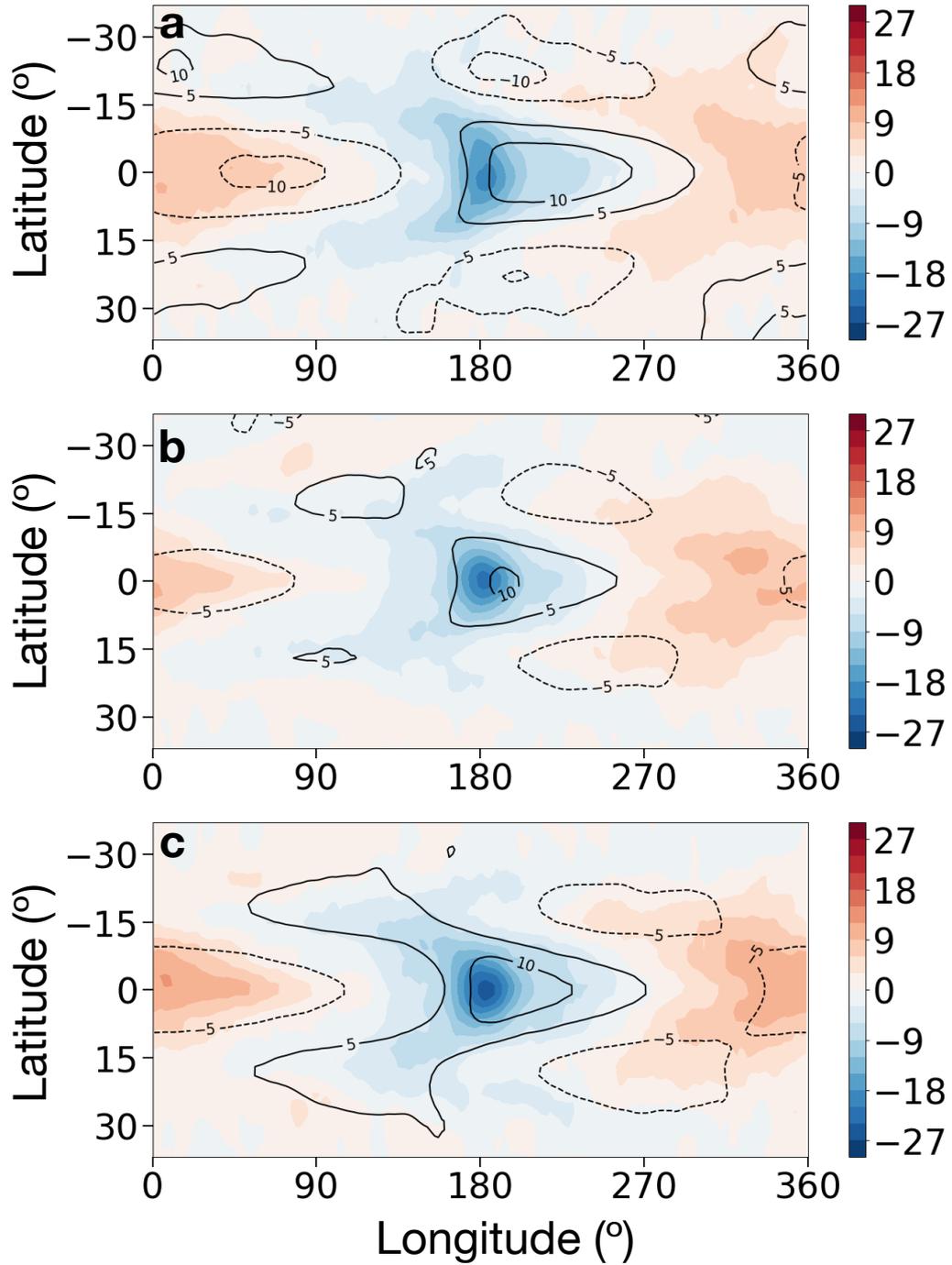


Figure S1. Horizontal structure of the MJO composite. (a) The control simulation. (b) The HomoRad simulation. (c) The FixedRad simulation. Color shading represents OLR anomalies, and contours represent geopotential anomalies associated with the MJO.

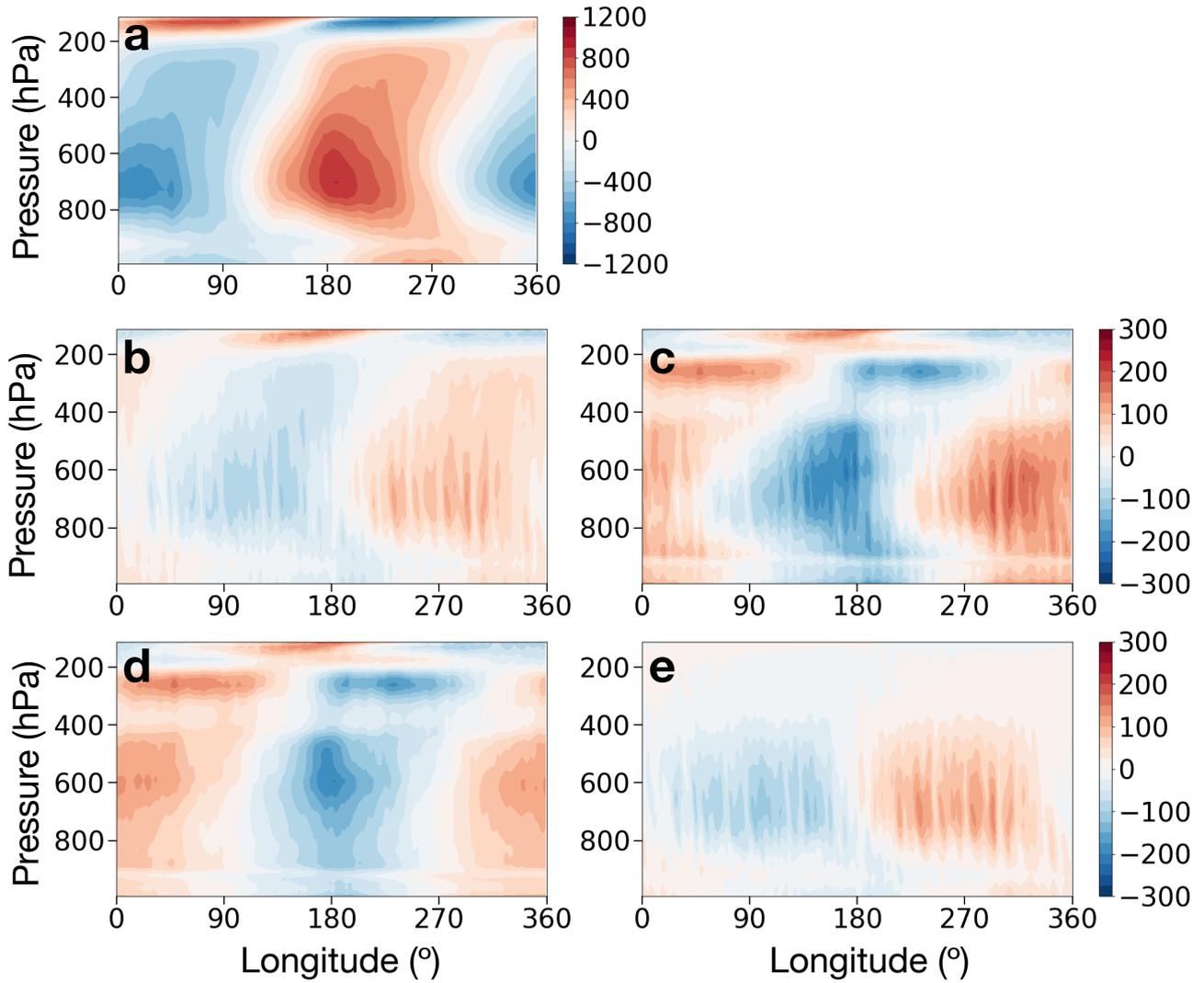


Figure S2. Vertical structures of the MSE budget of the MJO composite. (a) The MJO's MSE anomaly (J/kg). (b) Total MSE tendency $\partial_t \text{MSE}$ (J/kg/day). (c) Total MSE tendency due to convection (CRM), boundary layer, and large-scale dynamics (J/kg/day). (d) and (e) show the DSE and moisture components of (c), respectively (J/kg/day).