# Responses of field-aligned currents and equatorial electrojet to sudden decrease of solar wind dynamic pressure during the March 2023 geomagnetic storm

Guan Le<sup>1</sup>, Guiping Liu<sup>1</sup>, Endawoke Yizengaw<sup>2</sup>, Chin-Chun Wu<sup>3</sup>, Yihua Zheng<sup>1</sup>, Sarah Kimberly Vines<sup>4</sup>, and Natalia Buzulukova<sup>5</sup>

<sup>1</sup>NASA Goddard Space Flight Center <sup>2</sup>The Aerospace Corporation <sup>3</sup>Naval Research Laboratory <sup>4</sup>Johns Hopkins University Applied Physics Laboratory <sup>5</sup>UMCP/NASA GSFC

April 26, 2024

#### Abstract

We present a study of the magnetosphere-ionosphere coupling during the 23 March 2023 magnetic storm, focusing on the effect of the drastic decrease of the solar wind dynamic pressure occurred during the main phase. Our observations show that the negative pressure pulse had significant impact to the magnetosphere-ionosphere system. It weakened large-scale field-aligned currents and paused the progression of the storm main phase for  $\tilde{}$  3 hrs. Due to the sudden decrease of the plasma convection after the negative pressure pulse, the low-latitude ionosphere was over-shielded and experienced a brief period of westward penetration electric field, which reversed the direction of the equatorial electrojet. The counter electrojet was observed both in space and on the ground. A transient, localized enhancement of downward field-aligned current was observed near dawn, consistent with the mechanism for transmitting MHD disturbances from magnetosphere to the ionosphere after the negative pressure pulse.

1	Solar wind-magnetosphere-ionosphere coupling and its impact on equatorial
2	ionospheric electrodynamics during the 23 March 2023 geomagnetic storm:
3	Effect of sudden decrease of solar wind dynamic pressure
4	
5	Guan Le <sup>1</sup> , Guiping Liu <sup>1</sup> , and Endawoke Yizengaw <sup>2</sup> , Chin-Chun Wu <sup>3</sup> , Yihua Zheng <sup>1</sup> , Sarah Vines <sup>4,5</sup> ,
6	Natalia Buzulukova <sup>1,6</sup>
7	
8	<sup>1</sup> NASA Goddard Space Flight Center, Greenbelt, MD (guan.le@nasa.gov)
9	<sup>2</sup> The Aerospace Corporation, El Segundo, CA
10	<sup>3</sup> Naval Research Laboratory, Washington, DC
11	<sup>4</sup> Applied Physics Laboratory, Laurel, MD
12	<sup>5</sup> Now at Southwest Research Institute, San Antonio, TX
13	<sup>6</sup> University of Maryland, College Park, MD
14	
15	Corresponding author: Guan Le (guan.le@nasa.gov)
16	
17 18	Key Points:
19 20 21 22	1. Direct evidence of prompt penetration of electric field in the equatorial ionosphere caused by negative solar wind pressure pulse
23 24 25	2. Transient counter electrojet caused by westward penetration electric field after the arrival of negative pressure pulse
26 27 28 29	<ol> <li>Significant decrease of global large-scale FACs and transient enhancement of localized FAC in response to negative pressure pulse</li> </ol>

#### 30 Abstract

We present a study of the magnetosphere-ionosphere coupling during the 23 March 2023 magnetic 31 32 storm, focusing on the effect of the drastic decrease of the solar wind dynamic pressure occurred 33 during the main phase. Our observations show that the negative pressure pulse had significant impact to the magnetosphere-ionosphere system. It weakened large-scale field-aligned currents and 34 35 paused the progression of the storm main phase for  $\sim 3$  hrs. Due to the sudden decrease of the 36 plasma convection after the negative pressure pulse, the low-latitude ionosphere was over-shielded and experienced a brief period of westward penetration electric field, which reversed the direction 37 of the equatorial electrojet. The counter electrojet was observed both in space and on the ground. A 38 39 transient, localized enhancement of downward field-aligned current was observed near dawn, 40 consistent with the mechanism for transmitting MHD disturbances from magnetosphere to the 41 ionosphere after the negative pressure pulse.

42

#### 43 Plain Language Summary

44 The solar wind is a continuous stream of charged particles blowing from the Sun. The Earth's 45 magnetic field forms a protective shield around our planet, called the magnetosphere, which deflects most of the solar wind particles away from the Earth. Disturbances in the solar wind can interact 46 47 with the magnetosphere and impact the Earth's upper atmosphere (ionosphere). The interaction 48 creates electric fields forcing charged particles to move in the magnetosphere, which creates electric 49 currents flowing along the magnetic field lines connecting to the high-latitude ionosphere and drives the movement of charged particles there. The low-latitude ionosphere is generally shielded from 50 51 these electric fields. Sudden changes in the solar wind can break such balance, leading to the 52 electric field penetration to low latitudes. We examined how the magnetosphere and ionosphere 53 interacted during the 23 March 2023 geomagnetic storm, focusing on what happened when the solar 54 wind dynamic pressure suddenly decreased. We found the pressure drop caused a sudden decrease of the high-latitude electric field, resulting in a brief period of overshielding and the electric field in 55 the equatorial ionosphere reversed its direction. This changed the direction of the equatorial 56 electrojet, a major electric current in the ionosphere at the magnetic equator. 57

58 1 Introduction

59

In steady-state conditions, the low-latitude ionosphere is shielded from the high-latitude convection
electric field due to the partial ring current-associated region-2 (R2) field-aligned currents (FACs)
which act to oppose the electric field associated with region-1 (R1) FACs (e.g., Southwood, 1977).
However, it can be directly coupled to the magnetospheric disturbances through prompt penetration
of the convection electric field during active times (Nishida, 1968; Jaggi and Wolf, 1973; Fejer et
al., 1979).

66

67 The equatorial electrojet (EEJ), an intense band of eastward electric current flowing along the dayside magnetic equator in the E-region ionosphere (~110 km altitude), is driven by an eastward 68 69 zonal electric field from plasma-neutral collisional interactions known as the E-region wind dynamo 70 (Richmond, 1973; Heelis, 2004). The intensity and polarity of the EEJ respond directly to the 71 perturbations of the zonal electric field. Variations of the EEJ often serve as an indicator for the 72 equatorial zonal electric field perturbations, which can be caused by either neutral wind changes 73 from lower atmosphere forcing or prompt penetration electric fields (PPEFs) from enhanced 74 magnetosphere-ionosphere (M-I) coupling. Many studies have used EEJ variations to probe the 75 presence of PPEFs that are attributed to interplanetary magnetic field (IMF) variations (e.g., 76 Yizengaw et al., 2011, 2016) or solar wind dynamic pressure pulses (e.g., Nilam et al., 2020, 2023). 77 Understanding the sources and the process of PPEFs continues to be a subject of ongoing 78 investigation (Kelly et al., 2003; Fejer et al., 2024).

79

80 This paper reports the observations of the M-I coupling and its effect on the equatorial ionosphere 81 in response to a sudden decrease of the solar wind dynamic pressure during the main phase of the 82 23 March 2023 geomagnetic storm. Figure 1 shows 1-min resolution OMNI data for the IMF and solar wind parameters along with ground-based SYM-H index for 23-25 March 2023. This large 83 storm (minimum Dst  $\sim$  -170 nT, Kp  $\sim$  7) was associated with the passage of an interplanetary 84 85 coronal mass ejection (ICME), triggered by the southward IMF in both the sheath and the ICME 86 regions. A drastic density decrease was observed at the boundary crossing from the sheath to the ICME by the WIND spacecraft. As a result, a significant negative solar wind pressure pulse hit the 87

88 Earth's magnetosphere during the main phase of the storm (1440 UT, marked by the red dashed line



89 in Figure 1). The solar wind density as well as the dynamic pressure decreased by a factor of  $\sim 10$ .

Figure 1. The 1-min resolution OMNI data with IMF/solar wind parameters (top 7 panels) and SYM-H index
(bottom panel) for 23 March 2023 magnetic storm. The negative pressure pulse during the main phase of the
storm is marked by the red dashed line.

94

90

95

96 We examine how FACs at high latitudes and the EEJ at the equator responded to the negative 97 pressure pulse using both space and ground-based magnetic field data. In the following sections, we 98 first present evidence for a transient PPEF associated with the pressure pulse from the ground based 99 EEJ observations. Then we examine the response of large-scale FACs globally by AMPERE and 100 locally by Swarm satellites. We also analyze the EEJ observations in space by Swarm, which 101 provide additional evidence for the transient PPEF associated with the pressure pulse. Finally, we 102 discuss the dynamic processes involving solar wind pressure pulse interacting with the 103 magnetosphere and coupling into the polar ionosphere, that allow us to understand the behaviors of 104 the equatorial ionosphere. 105

## 106 2 Observations

#### 108 2.1 Ground-based Observations of the EEJ

109

110 The EEJ signals can be obtained from a pair of ground magnetometer stations located near the 111 magnetic equator on the same meridian, one directly under the EEJ at the equator (within  $\pm 3.5^{\circ}$ ) and the other just off the EEJ region ( $6^{\circ}-9^{\circ}$  from the magnetic equator) (Anderson et al., 2004; 112 113 Yizengaw et al., 2014). The EEJ signals can only be detected by the station at the magnetic equator 114 because the EEJ current is confined in a narrow latitudinal band (within  $\pm 3^{\circ}$ ). But both stations are 115 expected to record the same magnetic field variations from other large-scale current sources, such as 116 the solar quiet (Sq) currents, the ring current, and the magnetopause current. The EEJ signals are 117 extracted from the difference of the H-components between the two stations. In this study, we used two pairs of geomagnetic observatories at two meridians (~80°W and ~50°W). One pair is located 118 119 at Jicamarca (JICA, 11.95°S/76.87°W GEO, MLat = 0.6°N) and Piura (PIUR, 5.2°S/80.6°W GEO, Mlat = 6.9°N) in Peru. The other pair is located at Tatuoca, Brazil (TTB, 1.21°S/48.5°W GEO) and 120 Kourou, French Guvana (KOU, 5.21°N/52.7°W GEO). TTB and KOU are well located under and 121 122 far enough from the EEJ, respectively. They are within the region of South Atlantic Anomaly with rapid northward moving of the magnetic equator, and the magnetic equator passed the TTB in 123 124 March 2013 (Morschhauser et al., 2017).

125

Figure 2 shows the magnetic field observations from the 2 pairs of ground observatories on 23 March 2023 with three subpanels for each pair, from top to bottom, showing the H-component with the background removed ( $\delta H$ ) off the magnetic equator, at the magnetic equator, and the EEJ signal ( $\delta H_{EEJ}$ , the differences between  $\delta H$  at the geomagnetic equator and off the equator), respectively. The horizontal bar in the 3<sup>rd</sup> subpanel indicates dayside hours (6-18 LT) at the equator station. The red dashed line indicates the time of the negative pressure pulse (1440 UT) in Figure 1. The local

time (LT) of the pressure pulses at the two equator stations are also noted in Figure 2.

133

134 The eastward zonal electric field from the wind dynamo drives the eastward EEJ, producing a

positive magnetic field perturbation ( $\delta H_{EEJ} > 0$ ) in the dayside. This is generally the case in Figure 2

except for a brief period immediately following the negative pressure pulse. There was a transient

137 negative impulse of the H-component at all the stations, consisting of a sharp decrease (~6 min) and

a relatively gradual (~ 1 hour) return, apparently due to the sudden decrease of the magnetopause
current and expansion of the magnetosphere in response to the negative pressure pulse (Araki and
Nagano, 1988). However, the transient negative impulse at the equator station is much stronger than
its off-equator counterpart, and the EEJ signature reversed its sign showing a transient counter
electrojet flowing westward. This observation indicates the negative pressure pulse set up a
transient westward electric field (~ 1 hour) in the equatorial ionosphere.



145

146 Figure 2. Ground-based observations of the H-component from 2 pairs of ground observatories on 23 March

147 2023, JICA-PIUR and TTB-KOU, respectively. The red dashed line marks the negative pressure pulse in Figure

- 148 1. The black horizontal bars indicate the daytime (06-16 LT) at the equator stations.
- 149
- 150

#### 151 2.2 AMPERE Observations of Large-scale FACs

153 AMPERE observations of large-scale FACs are derived from global measurements of magnetic 154 field perturbations from the Iridium constellation of more than 70 near-polar orbiting satellites 155 [Anderson et al., 2000]. It collects 10-min data to generate one global pattern of large-scale FAC 156 distributions and provides a continuous monitor of the state of the global M-I system. (AMPERE 157 data will unlikely reveal transient and localized variations due to the limitation of spatial and 158 temporal resolution.) Figure 3 shows the AMPERE observations of the total field-aligned currents 159 flowing into and out of the ionosphere on 23 March 2023 (Figure S1 provides the magnetic field 160 perturbations and global FAC maps). The total upward current out of one hemisphere is calculated 161 by integrating all the upward current density over the entire area above 40° latitude, and likewise for the total downward current. Again, the red dashed vertical line corresponds to the negative pressure 162 163 pulse in Figure 1.

164

Starting from  $\sim 07$  UT, the total FACs gradually intensified as the storm progressed with the SYM-H index became more negative, representing an increasing active magnetosphere as FACs facilitate the electromagnetic energy input from the magnetosphere into the ionosphere. There is a brief period ( $\sim 1$  hr) of total current drop starting at  $\sim 13$  UT, apparently associated with the northward excursion of the IMF Bz component (Figure 1) which turned off the dayside reconnection and reduced the magnetospheric convection temporally.

171

172 Figure 3 shows the total currents responded to the negative pressure in two stages. The total currents

dropped sharply at ~1440 UT due to the sudden sunward motion of the magnetopause and

174 expansion of the magnetosphere. The sudden reduction of the magnetopause current also caused a

step decrease of the SYM-H index (Figure 1). Then the total currents continued to decrease

176 gradually. The decreasing trend of the SYM-H index has flattened out within the storm main phase,

indicating the pause of the ring current development (Figure 1). This is expected as IMF Bz

178 fluctuated around zero and the expanded magnetosphere adjusted to the new state of reduced

179 geomagnetic activity level. At ~1630 UT, the IMF Bz gradually turned southward, which

terminated the decreasing trend of the total currents. At  $\sim 18$  UT, both the total currents (Figure 3)

and the SYM-H index (Figure 1) showed that the magnetospheric activities began to intensify

182 rapidly with the prolonged steady southward IMF in the ICME. In summary of the AMPERE

183 observations, large-scale FACs were significantly weakened by the negative pressure pulse.



Figure 3. AMPERE Observations of the total amount of upward and downward FACs in northern and southernhemisphere, respectively.

**191 2.3** Swarm Observations of FACs and EEJ

192

193 Swarm is a three-satellite mission in a high-inclination (87.5°) low-Earth orbit, which provides 194 vector magnetic field data for frequent in situ measurements of FACs at high latitudes (Lühr et al., 195 2015) and scale magnetic field strength for the EEJ in the equatorial region (Alken et al., 2015). 196 Among the three satellites, A and C form a pair flying side by side at the same altitude (~460km) 197 with a longitudinal separation of 1.4°. Swarm B has slightly higher altitude (~530km) and its orbital 198 plane slowly drifts apart from those of Swarm A/C. In this study, we used two official Swarm level-199 2 data products: (1) the vector magnetic field residuals  $\delta B$  for the study of FACs, and (2) the height-200 integrated latitudinal profile of eastward EEJ current. The EEJ current profile is estimated from the 201 Swarm scalar magnetic field measurements by isolating the EEJ signal from the many other 202 geomagnetic sources and then fitting the EEJ signal with a line current model (Alken et al., 2015). 203 The EEJ current peak at the magnetic equator provides a good estimate of the EEJ strength. 204 Figure 4 presents an overview of the Swarm observations. Figure 4a shows the spacecraft orbits for 205 206 the polar cap pass near 1440 UT, the intervals marked by the red bars in Figures 4b/4c. Figures 4b and 4c contain 4 hours of Swarm vector magnetic field residuals  $\delta B$  in solar magnetic (SM) 207 208 coordinate system centered at 1440 UT (red dashed line) for Swarm A and B, respectively. Swarm 209 C data are nearly the same as Swarm A (not shown). During this interval, Swarm made 5 passes of 210 the polar cap, denoted by N (S) for the northern (southern) hemisphere, and 3 crossings of the 211 dayside magnetic equator marked by MEq and the blue dashed lines. The perturbations in  $\delta B$  are 212 the signals of FACs, occurring at auroral latitudes on both sides of the magnetic pole. The 213 latitudinal profiles of the estimated EEJ current at the dayside magnetic equator crossings are 214 presented in Figures 4d-4f for Swarm A and 4g-4i for Swarm B. The positive current is for eastward EEJ. 215

216

Both Swarm A and B were in the dayside morning sector over the northern polar cap at the time of
the negative pressure pulse (red dots in Figure 4a). In Figure 4a, the tick marks on each trajectory
are separated by 10 min. The red arrows indicate the directions of the spacecraft motion. Swarm A
was moving from nightside to dayside and Swarm B from dayside to nightside with ~ 2 hr local
time separation of the orbital planes.

223 In Figures 4b&4c, the FACs observed before the negative pressure pulse were generally stronger 224 than those after at Swarm, in agreement with the AMPERE observations. The only exception is that 225 the FAC signal was significantly enhanced to  $\sim 2000$  nT in magnitude shortly after the negative 226 pressure pulse at Swarm A (highlighted in yellow in Figure 4b) at ~7 LT (Figure 4a). The magnetic 227 field perturbations were mainly in the -x direction (anti-sunward), which is the signature of a pair of 228 FACs flowing downward at higher and upward at lower latitudes, respectively. The enhanced FAC 229 pair had the same polarity of the regular R1/R2 FACs in the dawn sector. The enhanced  $\delta Bx$ 230 magnitude was mainly due to the much-enhanced dawnward FAC at higher latitudes since the 231 gradient (i.e., time rate of change) of  $\delta Bx$  was significantly higher at the poleward edge. The FACs 232 observed by Swarm B at nearly the same time (yellow-highlighted interval in Figure 4c) but at ~11 233 LT (Figure 4a) did not show the same feature, neither did the subsequent FACs in the pre-midnight 234 sector. When Swarm A returned to the same region in next orbit about 90 min later (~ 1615 UT), the FACs have returned to the weakened state. These observations indicate the much-enhanced 235 236 downward FAC is a localized (near dawn) and transient (duration < 90 min) phenomenon in 237 response to the sudden decrease of the solar wind dynamic pressure. The AMPERE observations 238 did not capture such a localized transient response.

239

We now examine the EEJ profiles. As Swarm B is much closer to the local noon at the dayside 240 241 equator, the EEJ signal is expected to be much stronger at Swarm B than Swarm A. Before the 242 negative pressure pulse, the EEJ profile is not well defined at Swarm A (1323 UT, Figure 4d), 243 mostly likely due to a very weak EEJ in early morning. But closer to the local noon, Swarm B 244 detected the typical eastward EEJ profile at 1252 UT (Figure 4e) and 1426 UT (Figure 4f). Then 245 about 17 min after the negative pressure pulse, Swarm A observed a well-defined westward EEJ, or 246 counter electrojet (Figure 4e). The observed counter electrojet appeared to be a transient 247 phenomenon. The EEJ returned to nominal eastward direction in the next two profiles, 1601 UT at 248 Swarm B (Figure 4i) and 1632 UT at Swarm A (Figure 4f). These observations are in agreement 249 with the ground-based EEJ currents in Figure 2. 250

- 251
- 252



Figure 4. Swarm A and B observations of FACs and the EEJ: (a) Spacecraft trajectories near the negative
pressure pulse; (b-c) the vector magnetic field residuals; (d-i) the latitudinal profiles of the EEJ around the
magnetic equator.

- 261 **3** Discussion
- 262

263 We first summarize the observations presented above.

- The solar wind dynamic pressure decreased suddenly and significantly at the boundary of
   the ICME that caused the 23 March 2023 magnetic storm. The negative pressure pulse
   arrived at the Earth at 1440 UT during the main phase of the storm and the IMF Bz
   fluctuated between northward and southward (Figure 1).
- The total large-scale FAC currents flowing into and out of the ionosphere decreased
   significantly soon after the arrival of the negative pressure pulse based (Figure 3). The
   overall geomagnetic activity level in the magnetosphere was weakened for more than 3 hrs,
   which paused the progression of the storm main phase. The activity level picked up again
   only after the IMF Bz turned strongly southward for an extended period during the passage
   of the ICME.
- Swarm A observed a significant enhancement of the downward FAC at the poleward edge
   of the FAC region near dawn shortly after the negative pressure pulse, which appeared to be
   localized and transient (Figure 4). Nearly simultaneous Swarm B observations closer to the
   local noon showed weakened FACs, consistent with the AMPERE observations.
- A transient counter electrojet was observed both in space by Swarm A (Figure 4) and on the ground (Figure 2) within minutes after the arrival of the negative pressure pulse. The counter electrojet lasted for ~ 1 hr and then returned to its regular eastward direction. The observed transient reversal of the EEJ to the westward direction suggests that the equatorial ionosphere experienced a brief period of a westward electric field after the negative pressure pulse.
- 284

These observations demonstrate the profound impact to the M-I system by the negative pressure pulse. The observed counter electrojet clearly shows that a transient westward electric field associated with the negative pressure pulse penetrated to the equatorial ionosphere from overshielding (Hori et al., 2012; Fujita et al., 2012). The penetration electric field was much stronger in magnitude than the background eastward electric field from the wind dynamo so that the overall zonal electric field was reversed. Our observations indicated there was a sudden decrease of the dawn-to-dusk (eastward) convection electric field as evident by the sudden decrease of the total

292 FAC currents flowing into and out of the polar ionosphere immediately after the negative pressure

- 293 pulse (Figure 3). The total FACs then gradually decrease with a time scale of hours. However,
- 294 SYM-H, the ring current index, was flatten out in the same period, indicating the ring current did
- not immediately respond to the weakened convection electric field (Figure 1). The delayed response
- of the ring current reflects the time scale for the M-I system to gradually adjust to the expanded
- state of the magnetosphere with decreased level of plasma convection (Earle and Kelley, 1987).
- 298 Thus, there was a short period when the low-latitude ionosphere was over-shielded and experienced
- a dusk-to-dawn (westward) electric field. Based on the duration of the counter-electrojet in the
- 300 ground-based observations (Figure 2), the response of the ring current-R2 FAC system was delayed
- 301 for  $\sim 6$  min, and it took  $\sim 1$  hr for the M-I system to gradually adjust itself to the decreased plasma
- 302 convection level and the low-latitude ionosphere to return to be fully shielded.
- 303

304 To understand the transient responses and localized enhancement of FACs, it is necessary to review 305 the current understanding of the underlying physical process. The M-I system responds to a sudden 306 pressure pulse in two phases, including a preliminary impulse (PI) and a two-stage main impulse 307 (MI) (e.g., Tamao, 1964a&b; Araki, 1977; Araki and Allen, 1982). The PI is due to the propagation and conversion of a compressional wave front launched from the magnetopause when the 308 309 magnetosphere is suddenly compressed or expanded. The PI is transient by nature because its driver 310 is the interaction between the pressure pulse and the magnetopause, which disappears in minutes 311 after the impulse front propagates away from the dayside.

312

313 Although more previous studies focused on sudden pressure increases than decreases, the basic 314 physics is the same. Based on Tamao's (1964a&b) pioneer work, Araki (1994) proposed a M-I 315 coupling PI model to explain the global observations after geomagnetic sudden commencements. 316 As illustrated in their Figure 12, the magnetopause moves inward and the dawn-to-dusk 317 magnetopause current increases when the solar wind dynamic pressure suddenly increases. A 318 compressional MHD wave is excited on the magnetopause, which propagates into the equatorial 319 magnetosphere. The solar wind-magnetosphere interaction as a dynamo generates an enhanced 320 dusk-to-down electric field at the magnetopause ( $J \cdot E < 0$ ). A dusk-to-dawn electric field and 321 associated inertia electric current are induced inside the magnetosphere. The extra magnetopause 322 current and the inertia current would form a counterclockwise current loop. The compressional

323 wave will be converted into the transverse Alfven wave due to the nonuniformity of the

magnetosphere (Tamao, 1964b; Southwood and Kivelson, 1990). When the compressional wave

front reaches the region where the Alfven speed has a largest spatial gradient, converted Alfven

326 waves are generated and propagate along the field lines with associated FACs. A pair of FACs will

327 be a part of the current loop, downward in the dusk side and upward in the dawn side. This process

328 happens in time scale of minutes. So, the pair of FACs exists transiently at lower latitudes than the

329 regular R1 currents with opposite polarity. A quantitative detail of the PI process is provided in the

330 MHD simulations by Fujita et al. (2003a&b, 2005), and the source region of the MHD wave mode

331 conversion for the generation of the transient FACs was found to be in the region of 6 < L < 7

**332** (Fujita et al., 2003a).

333

334 In the case of negative pressure pulses, the observations by Araki (1988) and simulations by Fujita 335 et al. (2004, 2012) showed that the magnetospheric and ionospheric signatures mostly mirror those 336 in pressure pulses. The negative pressure pulse causes the expansion of the magnetosphere and a 337 decrease of the magnetopause current. The PI is associated with a dawn-to-dusk transient dynamo 338 electric field at the magnetopause and induced electric field in the magnetosphere. The equatorial 339 current loop would be clockwise to effectively reduce the magnetic field strength in the 340 magnetosphere, and the pair of transient FACs would be downward in the dawnside and upward in 341 the duskside, in the same polarity of the regular R1 currents. The transient and localized 342 enhancement of the downward FAC observed by Swarm A near dawn (Figure 4) matches the 343 predicted polarity of the FACs. However, our observations differ in an important aspect from the 344 model prediction. The transient, localized FAC enhancement was observed at the poleward edge of the FAC region, implying the source region was near the magnetopause, as in the earliest work of 345 346 Tamao (1964a). Further theorical and numerical investigation is still needed to understand the 347 source region of the transient FACs during the PI. In addition, understanding the role of the ring 348 current/R2 FAC system to the undershielding/overshielding and its restoration is particularly needed 349 in future simulations.

350

351 4 Conclusions

353	A drastic decrease of the solar wind dynamic pressure occurred during main phase of the 23 March
354	2023 geomagnetic storm in association with the boundary between the ICME and its sheath. Our
355	observations show that the negative pressure pulse had significant impact to the M-I system. It
356	weakened the overall geomagnetic activities and plasma convection and paused the progression of
357	the storm main phase for $\sim$ 3 hrs. Due to the sudden decrease of the dawn-to-dusk convection
358	electric field, there was a transient period when the low-latitude ionosphere was over-shielded and
359	experienced a brief period of dusk-to-dawn (westward) penetration electric field. The transient
360	westward penetration electric field reversed the direction of the equatorial electrojet, and the
361	counter electrojet was observed both in space and on the ground. The response of the ring current-
362	R2 FAC system was delayed for ~6 min, and it took ~ 1 hr for the M-I system to adjust itself to the
363	decreased plasma convection level until the low-latitude ionosphere was fully shielded again.
364	Although the overall large-scale FACs were weakened by the negative pressure pulse, a transient,
365	localized enhancement of downward FAC was observed near dawn, consistent with the mechanism
366	for transmitting MHD disturbances in the M-I coupling after the negative pressure pulse. But the
367	latitudinal location of the localized FAC enhancement differed from the model prediction, which
368	calls further investigation of the MI coupling in response to the pressure pulse.
369	
370	
371	Acknowledgements
372	
373	GL thanks Lan Jian for helpful discussion. EY was partially supported by the AFOSR (FA9550-20-
374	1-0119) and NSF (AGS-1848730) grants. NB was partially supported by NASA Internal Scientist
375	Funding Model on Mesoscale Dynamics.
376	
377	

# 379 Data Availability Statement

- 380
- 381 The OMNI data are available at <u>https://omniweb.gsfc.nasa.gov</u>. The JICA and PIUR magnetometer
- data are available at <u>https://zenodo.org/records/10823058</u>. The KOU and TTB magnetometer data
- are available at INTERMAGNET (<u>www.intermagnet.org</u>). The AMPERE data are available at
- 384 <u>https://ampere.jhuapl.edu</u>. The Swarm data are accessible at
- 385 <u>https://earth.esa.int/eogateway/missions/swarm/data.</u>

386

387

389	References
390	
391	Alken, P., S. Maus, A. Chulliat, P. Vigneron, O. Sirol, and G. Hulot (2015), Swarm equatorial
392	electric field chain: First results, Geophys. Res. Lett., 42, 673-680.
393	https://doi.org/10.1002/2014GL062658
394	
395	Anderson, D., A. Anghel, J. Chau, and O. Veliz (2004), Daytime vertical $E \times B$ drift velocities
396	inferred from ground-based magnetometer observations at low latitudes, Space Weather, 2, S11001.
397	https://doi.org/10.1029/2004SW000095
398	
399	Araki, T. (1977), Global structure of geomagnetic sudden commencements, Planet. Space Sci., 25,
400	373. https://doi.org/10.1016/0032-0633(77)90053-8
401	
402	Araki, T. (1994). A physical model of the geomagnetic sudden commencement. Geophysical
403	Monograph-American Geophysical Union, 81, 183-183. https://doi.org/10.1029/GM081p0183
404	
405	T. Araki, J. H. Allen (1982), Latitudinal reversal of polarization of the geomagnetic sudden
406	commencement, J. Geophys. Res., 87, 5207-5216. <u>https://doi.org/10.1029/JA087iA07p05207</u>
407	
408	Araki, T. and H. Nagano (1988), Geomagnetic response to sudden expansions of the
409	magnetosphere, J. Geophys. Res., 93, 3983–3988. <u>https://doi.org/10.1029/JA093iA05p03983</u>
410	
411	Earle, G. D., and M. C. Kelley (1987), Spectral studies of the sources of ionospheric electric fields,
412	J. Geophys. Res., 92(A1), 213–224. https://doi.org.10.1029/JA092iA01p00213
413	
414	Fejer, B., C. Gonzales, D. Farley, M. Kelley, and R. Woodman (1979), Equatorial electric fields
415	during magnetically disturbed conditions 1. The effect of the interplanetary magnetic field, J.
416	Geophys. Res., 84(A10), 5797-5802. https://doi.org/10.1029/JA084iA10p05797
417	

- 418 Fejer BG, Laranja SR and Condor P (2024), Multi-process driven unusually large equatorial
- 419 perturbation electric fields during the April 2023 geomagnetic storm. Front. Astron. Space Sci.
- 420 11:1351735. <u>https://doi.org/10.3389/fspas.2024.1351735</u>
- 421
- 422 Fujita, S., T. Tanaka, T. Kikuchi, K. Fujimoto, K. Hosokawa, and M. Itonaga (2003a), A numerical
- 423 simulation of the geomagnetic sudden commencement: 1. Generation of the field-aligned current
- 424 associated with the preliminary impulse, J. Geophys. Res., 108(A12), 1416.
- 425 <u>https://doi.org/10.1029/2002JA009407</u>
- 426
- 427 Fujita, S., T. Tanaka, T. Kikuchi, K. Fujimoto, and M. Itonaga (2003b), A numerical simulation of
- 428 the geomagnetic sudden commencement: 2. Plasma processes in the main impulse, J. Geophys.
- 429 Res., 108(A12), 1417. <u>https://doi.org/10.1029/2002JA009763</u>
- 430
- Fujita, S., T. Tanaka, T. Kikuchi, and S. Tsunomura (2004), A numerical simulation of a negative
  sudden impulse, Earth Planets Space, 56, 463–472. <u>https://doi.org/10.1186/BF03352499</u>
- 433
- 434 Fujita, S., T. Tanaka, and T. Motoba (2005), A numerical simulation of the geomagnetic sudden
- 435 commencement: 3. A sudden commencement in the magnetosphere-ionosphere compound system,
- 436 J. Geophys. Res., 110, A11203. <u>https://doi.org/10.1029/2005JA011055</u>
- 437
- 438 Fujita, S., H. Yamagishi, K. T. Murata, M. Den, and T. Tanaka (2012), A numerical simulation of a
- 439 negative solar wind impulse: Revisited, J. Geophys. Res., 117, A09219.
- 440 <u>https://doi.org/10.1029/2012JA017526</u>
- 441
- 442 Heelis, R. A. (2004). Electrodynamics in the low and middle latitude ionosphere: A tutorial. *Journal*
- 443 of Atmospheric and Solar-Terrestrial Physics, 66(10), 825-
- 444 838. <u>https://doi.org/10.1016/j.jastp.2004.01.034</u>
- 445
- 446 Hori, T., A. Shinbori, N. Nishitani, T. Kikuchi, S. Fujita, T. Nagatsuma, O. Troshichev, K. Yumoto,
- 447 A. Moiseyev, and K. Seki (2012), Evolution of negative SI-induced ionospheric flows observed by

448	SuperDARN King Salmon HF radar, J. Geophys. Res., 117, A12223,
449	https://doi.org/10.1029/2012JA018093
450	
451	Jaggi, R. K., and R. A. Wolf (1973), Self-consistent calculation of the motion of a sheet of ions in
452	the magnetosphere, J. Geophys. Res.; Space Physics, 78(16), 2852-2866.
453	https://doi.org/10.1029/JA078i016p02852
454	
455	Kelley, M. C., J. J. Makela, J. L. Chau, and M. J. Nicolls (2003), Penetration of the solar wind
456	electric field into the magnetosphere/ionosphere system, Geophys. Res. Lett., 30, 1158.
457	https://doi.org/10.1029/2002GL016321
458	
459	Morschhauser, A., Brando Soares, G., Haseloff, J., Bronkalla, O., Protásio, J., Pinheiro, K., and
460	Matzka, J. (2017): The magnetic observatory on Tatuoca, Belém, Brazil: history and recent
461	developments, Geosci. Instrum. Method. Data Syst., 6, 367-376. https://doi.org/10.5194/gi-6-367-
462	2017
463	
464	Nishida, A. (1968), Coherence of geomagnetic DP 2 fluctuations with interplanetary magnetic
465	variations, J. Geophys. Res.: Space Physics, 73, 5549-5559.
466	https://doi.org/10.1029/JA073i017p05549
467	
468	Richmond, A. D. (1973). Equatorial electrojet-1. Development of a model including winds and
469	instabilities. Journal of Atmospheric and Terrestrial Physics, 35(6), 1083-
470	1103. <u>https://doi.org/10.1016/0021-9169(73)90007-x</u>
471	
472	Southwood, D. J. (1977), The role of hot plasma in magnetospheric convection, J. Geophys. Res.,
473	82(35), 5512-5520. https://doi.org/10.1029/JA082i035p05512
474	
475	Southwood, D. J., and M. G. Kivelson (1990), The magnetohydrodynamic response of the
476	magnetospheric cavity to changes in solar wind pressure, J. Geophys. Res., 95, 2301-2309.
477	https://doi.org/10.1029/JA095iA03p02301
478	

479	Tamao, T. (1964a), The structure of three-dimensional hydromagnetic waves in a uniform cold
480	plasma, J. Geomagn. Geoelectr., 16, 89–114, 1964a. https://doi.org/10.5636/jgg.16.89
481	
482	Tamao, T. (1964b), A hydromagnetic interpretation of geomagnetic SSC*, Rep. Ionos. Space Res.
483	Jpn., 18, 16–31.
484	
485	Yizengaw, E., M. B. Moldwin, A. Mebrahtu, B. Damtie, E. Zesta, C. E. Valladares, and P. Doherty
486	(2011), Comparison of storm time equatorial ionospheric electrodynamics in the African and
487	American sectors, Journal of Atmospheric and Solar-Terrestrial Physics, 73, 156-163.
488	https://doi.org/10.1016/j.jastp.2010.08.008
489	
490	Yizengaw, E., M. B. Moldwin, E. Zesta, C. M. Biouele, B. Damtie, A. Mebrahtu, B. Rabiu, C. F.
491	Valladares, and R. Stoneback (2014), The longitudinal variability of equatorial electrojet and
492	vertical drift velocity in the African and American sector, Ann. Geophys., 32, 231-238.
493	https://doi.org/10.5194/angeo-32-231-2014
494	

- 495 Yizengaw, E., Moldwin, M. B., Zesta, E., Magoun, M., Pradipta, R., Biouele, C. M., et al. (2016).
- Response of the equatorial ionosphere to the geomagnetic DP 2 current system. Geophysical 496
- 497 Research Letters, 43(14), 7364–7372. https://doi.org/10.1002/2016g1070090

1	Solar wind-magnetosphere-ionosphere coupling and its impact on equatorial
2	ionospheric electrodynamics during the 23 March 2023 geomagnetic storm:
3	Effect of sudden decrease of solar wind dynamic pressure
4	
5	Guan Le <sup>1</sup> , Guiping Liu <sup>1</sup> , and Endawoke Yizengaw <sup>2</sup> , Chin-Chun Wu <sup>3</sup> , Yihua Zheng <sup>1</sup> , Sarah Vines <sup>4,5</sup> ,
6	Natalia Buzulukova <sup>1,6</sup>
7	
8	<sup>1</sup> NASA Goddard Space Flight Center, Greenbelt, MD (guan.le@nasa.gov)
9	<sup>2</sup> The Aerospace Corporation, El Segundo, CA
10	<sup>3</sup> Naval Research Laboratory, Washington, DC
11	<sup>4</sup> Applied Physics Laboratory, Laurel, MD
12	<sup>5</sup> Now at Southwest Research Institute, San Antonio, TX
13	<sup>6</sup> University of Maryland, College Park, MD
14	
15	Corresponding author: Guan Le (guan.le@nasa.gov)
16	
17 18	Key Points:
19 20 21 22	1. Direct evidence of prompt penetration of electric field in the equatorial ionosphere caused by negative solar wind pressure pulse
23 24 25	2. Transient counter electrojet caused by westward penetration electric field after the arrival of negative pressure pulse
26 27 28 29	<ol> <li>Significant decrease of global large-scale FACs and transient enhancement of localized FAC in response to negative pressure pulse</li> </ol>

#### 30 Abstract

We present a study of the magnetosphere-ionosphere coupling during the 23 March 2023 magnetic 31 32 storm, focusing on the effect of the drastic decrease of the solar wind dynamic pressure occurred 33 during the main phase. Our observations show that the negative pressure pulse had significant impact to the magnetosphere-ionosphere system. It weakened large-scale field-aligned currents and 34 35 paused the progression of the storm main phase for  $\sim 3$  hrs. Due to the sudden decrease of the 36 plasma convection after the negative pressure pulse, the low-latitude ionosphere was over-shielded and experienced a brief period of westward penetration electric field, which reversed the direction 37 of the equatorial electrojet. The counter electrojet was observed both in space and on the ground. A 38 39 transient, localized enhancement of downward field-aligned current was observed near dawn, 40 consistent with the mechanism for transmitting MHD disturbances from magnetosphere to the 41 ionosphere after the negative pressure pulse.

42

#### 43 Plain Language Summary

44 The solar wind is a continuous stream of charged particles blowing from the Sun. The Earth's 45 magnetic field forms a protective shield around our planet, called the magnetosphere, which deflects most of the solar wind particles away from the Earth. Disturbances in the solar wind can interact 46 47 with the magnetosphere and impact the Earth's upper atmosphere (ionosphere). The interaction 48 creates electric fields forcing charged particles to move in the magnetosphere, which creates electric 49 currents flowing along the magnetic field lines connecting to the high-latitude ionosphere and drives the movement of charged particles there. The low-latitude ionosphere is generally shielded from 50 51 these electric fields. Sudden changes in the solar wind can break such balance, leading to the 52 electric field penetration to low latitudes. We examined how the magnetosphere and ionosphere 53 interacted during the 23 March 2023 geomagnetic storm, focusing on what happened when the solar 54 wind dynamic pressure suddenly decreased. We found the pressure drop caused a sudden decrease of the high-latitude electric field, resulting in a brief period of overshielding and the electric field in 55 the equatorial ionosphere reversed its direction. This changed the direction of the equatorial 56 electrojet, a major electric current in the ionosphere at the magnetic equator. 57

58 1 Introduction

59

In steady-state conditions, the low-latitude ionosphere is shielded from the high-latitude convection
electric field due to the partial ring current-associated region-2 (R2) field-aligned currents (FACs)
which act to oppose the electric field associated with region-1 (R1) FACs (e.g., Southwood, 1977).
However, it can be directly coupled to the magnetospheric disturbances through prompt penetration
of the convection electric field during active times (Nishida, 1968; Jaggi and Wolf, 1973; Fejer et
al., 1979).

66

67 The equatorial electrojet (EEJ), an intense band of eastward electric current flowing along the dayside magnetic equator in the E-region ionosphere (~110 km altitude), is driven by an eastward 68 69 zonal electric field from plasma-neutral collisional interactions known as the E-region wind dynamo 70 (Richmond, 1973; Heelis, 2004). The intensity and polarity of the EEJ respond directly to the 71 perturbations of the zonal electric field. Variations of the EEJ often serve as an indicator for the 72 equatorial zonal electric field perturbations, which can be caused by either neutral wind changes 73 from lower atmosphere forcing or prompt penetration electric fields (PPEFs) from enhanced 74 magnetosphere-ionosphere (M-I) coupling. Many studies have used EEJ variations to probe the 75 presence of PPEFs that are attributed to interplanetary magnetic field (IMF) variations (e.g., 76 Yizengaw et al., 2011, 2016) or solar wind dynamic pressure pulses (e.g., Nilam et al., 2020, 2023). 77 Understanding the sources and the process of PPEFs continues to be a subject of ongoing 78 investigation (Kelly et al., 2003; Fejer et al., 2024).

79

80 This paper reports the observations of the M-I coupling and its effect on the equatorial ionosphere 81 in response to a sudden decrease of the solar wind dynamic pressure during the main phase of the 82 23 March 2023 geomagnetic storm. Figure 1 shows 1-min resolution OMNI data for the IMF and solar wind parameters along with ground-based SYM-H index for 23-25 March 2023. This large 83 storm (minimum Dst  $\sim$  -170 nT, Kp  $\sim$  7) was associated with the passage of an interplanetary 84 85 coronal mass ejection (ICME), triggered by the southward IMF in both the sheath and the ICME 86 regions. A drastic density decrease was observed at the boundary crossing from the sheath to the ICME by the WIND spacecraft. As a result, a significant negative solar wind pressure pulse hit the 87

88 Earth's magnetosphere during the main phase of the storm (1440 UT, marked by the red dashed line



89 in Figure 1). The solar wind density as well as the dynamic pressure decreased by a factor of  $\sim 10$ .

Figure 1. The 1-min resolution OMNI data with IMF/solar wind parameters (top 7 panels) and SYM-H index
(bottom panel) for 23 March 2023 magnetic storm. The negative pressure pulse during the main phase of the
storm is marked by the red dashed line.

94

90

95

96 We examine how FACs at high latitudes and the EEJ at the equator responded to the negative 97 pressure pulse using both space and ground-based magnetic field data. In the following sections, we 98 first present evidence for a transient PPEF associated with the pressure pulse from the ground based 99 EEJ observations. Then we examine the response of large-scale FACs globally by AMPERE and 100 locally by Swarm satellites. We also analyze the EEJ observations in space by Swarm, which 101 provide additional evidence for the transient PPEF associated with the pressure pulse. Finally, we 102 discuss the dynamic processes involving solar wind pressure pulse interacting with the 103 magnetosphere and coupling into the polar ionosphere, that allow us to understand the behaviors of 104 the equatorial ionosphere. 105

## 106 2 Observations

#### 108 2.1 Ground-based Observations of the EEJ

109

110 The EEJ signals can be obtained from a pair of ground magnetometer stations located near the 111 magnetic equator on the same meridian, one directly under the EEJ at the equator (within  $\pm 3.5^{\circ}$ ) and the other just off the EEJ region ( $6^{\circ}-9^{\circ}$  from the magnetic equator) (Anderson et al., 2004; 112 113 Yizengaw et al., 2014). The EEJ signals can only be detected by the station at the magnetic equator 114 because the EEJ current is confined in a narrow latitudinal band (within  $\pm 3^{\circ}$ ). But both stations are 115 expected to record the same magnetic field variations from other large-scale current sources, such as 116 the solar quiet (Sq) currents, the ring current, and the magnetopause current. The EEJ signals are 117 extracted from the difference of the H-components between the two stations. In this study, we used two pairs of geomagnetic observatories at two meridians (~80°W and ~50°W). One pair is located 118 119 at Jicamarca (JICA, 11.95°S/76.87°W GEO, MLat = 0.6°N) and Piura (PIUR, 5.2°S/80.6°W GEO, Mlat = 6.9°N) in Peru. The other pair is located at Tatuoca, Brazil (TTB, 1.21°S/48.5°W GEO) and 120 Kourou, French Guvana (KOU, 5.21°N/52.7°W GEO). TTB and KOU are well located under and 121 122 far enough from the EEJ, respectively. They are within the region of South Atlantic Anomaly with rapid northward moving of the magnetic equator, and the magnetic equator passed the TTB in 123 124 March 2013 (Morschhauser et al., 2017).

125

Figure 2 shows the magnetic field observations from the 2 pairs of ground observatories on 23 March 2023 with three subpanels for each pair, from top to bottom, showing the H-component with the background removed ( $\delta H$ ) off the magnetic equator, at the magnetic equator, and the EEJ signal ( $\delta H_{EEJ}$ , the differences between  $\delta H$  at the geomagnetic equator and off the equator), respectively. The horizontal bar in the 3<sup>rd</sup> subpanel indicates dayside hours (6-18 LT) at the equator station. The red dashed line indicates the time of the negative pressure pulse (1440 UT) in Figure 1. The local

time (LT) of the pressure pulses at the two equator stations are also noted in Figure 2.

133

134 The eastward zonal electric field from the wind dynamo drives the eastward EEJ, producing a

positive magnetic field perturbation ( $\delta H_{EEJ} > 0$ ) in the dayside. This is generally the case in Figure 2

except for a brief period immediately following the negative pressure pulse. There was a transient

137 negative impulse of the H-component at all the stations, consisting of a sharp decrease (~6 min) and

a relatively gradual (~ 1 hour) return, apparently due to the sudden decrease of the magnetopause
current and expansion of the magnetosphere in response to the negative pressure pulse (Araki and
Nagano, 1988). However, the transient negative impulse at the equator station is much stronger than
its off-equator counterpart, and the EEJ signature reversed its sign showing a transient counter
electrojet flowing westward. This observation indicates the negative pressure pulse set up a
transient westward electric field (~ 1 hour) in the equatorial ionosphere.



145

146 Figure 2. Ground-based observations of the H-component from 2 pairs of ground observatories on 23 March

147 2023, JICA-PIUR and TTB-KOU, respectively. The red dashed line marks the negative pressure pulse in Figure

- 148 1. The black horizontal bars indicate the daytime (06-16 LT) at the equator stations.
- 149
- 150

#### 151 2.2 AMPERE Observations of Large-scale FACs

153 AMPERE observations of large-scale FACs are derived from global measurements of magnetic 154 field perturbations from the Iridium constellation of more than 70 near-polar orbiting satellites 155 [Anderson et al., 2000]. It collects 10-min data to generate one global pattern of large-scale FAC 156 distributions and provides a continuous monitor of the state of the global M-I system. (AMPERE 157 data will unlikely reveal transient and localized variations due to the limitation of spatial and 158 temporal resolution.) Figure 3 shows the AMPERE observations of the total field-aligned currents 159 flowing into and out of the ionosphere on 23 March 2023 (Figure S1 provides the magnetic field 160 perturbations and global FAC maps). The total upward current out of one hemisphere is calculated 161 by integrating all the upward current density over the entire area above 40° latitude, and likewise for the total downward current. Again, the red dashed vertical line corresponds to the negative pressure 162 163 pulse in Figure 1.

164

Starting from  $\sim 07$  UT, the total FACs gradually intensified as the storm progressed with the SYM-H index became more negative, representing an increasing active magnetosphere as FACs facilitate the electromagnetic energy input from the magnetosphere into the ionosphere. There is a brief period ( $\sim 1$  hr) of total current drop starting at  $\sim 13$  UT, apparently associated with the northward excursion of the IMF Bz component (Figure 1) which turned off the dayside reconnection and reduced the magnetospheric convection temporally.

171

172 Figure 3 shows the total currents responded to the negative pressure in two stages. The total currents

dropped sharply at ~1440 UT due to the sudden sunward motion of the magnetopause and

174 expansion of the magnetosphere. The sudden reduction of the magnetopause current also caused a

step decrease of the SYM-H index (Figure 1). Then the total currents continued to decrease

176 gradually. The decreasing trend of the SYM-H index has flattened out within the storm main phase,

indicating the pause of the ring current development (Figure 1). This is expected as IMF Bz

178 fluctuated around zero and the expanded magnetosphere adjusted to the new state of reduced

179 geomagnetic activity level. At ~1630 UT, the IMF Bz gradually turned southward, which

terminated the decreasing trend of the total currents. At  $\sim 18$  UT, both the total currents (Figure 3)

and the SYM-H index (Figure 1) showed that the magnetospheric activities began to intensify

182 rapidly with the prolonged steady southward IMF in the ICME. In summary of the AMPERE

183 observations, large-scale FACs were significantly weakened by the negative pressure pulse.



Figure 3. AMPERE Observations of the total amount of upward and downward FACs in northern and southernhemisphere, respectively.

**191 2.3** Swarm Observations of FACs and EEJ

192

193 Swarm is a three-satellite mission in a high-inclination (87.5°) low-Earth orbit, which provides 194 vector magnetic field data for frequent in situ measurements of FACs at high latitudes (Lühr et al., 195 2015) and scale magnetic field strength for the EEJ in the equatorial region (Alken et al., 2015). 196 Among the three satellites, A and C form a pair flying side by side at the same altitude (~460km) 197 with a longitudinal separation of 1.4°. Swarm B has slightly higher altitude (~530km) and its orbital 198 plane slowly drifts apart from those of Swarm A/C. In this study, we used two official Swarm level-199 2 data products: (1) the vector magnetic field residuals  $\delta B$  for the study of FACs, and (2) the height-200 integrated latitudinal profile of eastward EEJ current. The EEJ current profile is estimated from the 201 Swarm scalar magnetic field measurements by isolating the EEJ signal from the many other 202 geomagnetic sources and then fitting the EEJ signal with a line current model (Alken et al., 2015). 203 The EEJ current peak at the magnetic equator provides a good estimate of the EEJ strength. 204 Figure 4 presents an overview of the Swarm observations. Figure 4a shows the spacecraft orbits for 205 206 the polar cap pass near 1440 UT, the intervals marked by the red bars in Figures 4b/4c. Figures 4b and 4c contain 4 hours of Swarm vector magnetic field residuals  $\delta B$  in solar magnetic (SM) 207 208 coordinate system centered at 1440 UT (red dashed line) for Swarm A and B, respectively. Swarm 209 C data are nearly the same as Swarm A (not shown). During this interval, Swarm made 5 passes of 210 the polar cap, denoted by N (S) for the northern (southern) hemisphere, and 3 crossings of the 211 dayside magnetic equator marked by MEq and the blue dashed lines. The perturbations in  $\delta B$  are 212 the signals of FACs, occurring at auroral latitudes on both sides of the magnetic pole. The 213 latitudinal profiles of the estimated EEJ current at the dayside magnetic equator crossings are 214 presented in Figures 4d-4f for Swarm A and 4g-4i for Swarm B. The positive current is for eastward EEJ. 215

216

Both Swarm A and B were in the dayside morning sector over the northern polar cap at the time of
the negative pressure pulse (red dots in Figure 4a). In Figure 4a, the tick marks on each trajectory
are separated by 10 min. The red arrows indicate the directions of the spacecraft motion. Swarm A
was moving from nightside to dayside and Swarm B from dayside to nightside with ~ 2 hr local
time separation of the orbital planes.

223 In Figures 4b&4c, the FACs observed before the negative pressure pulse were generally stronger 224 than those after at Swarm, in agreement with the AMPERE observations. The only exception is that 225 the FAC signal was significantly enhanced to  $\sim 2000$  nT in magnitude shortly after the negative 226 pressure pulse at Swarm A (highlighted in yellow in Figure 4b) at ~7 LT (Figure 4a). The magnetic 227 field perturbations were mainly in the -x direction (anti-sunward), which is the signature of a pair of 228 FACs flowing downward at higher and upward at lower latitudes, respectively. The enhanced FAC 229 pair had the same polarity of the regular R1/R2 FACs in the dawn sector. The enhanced  $\delta Bx$ 230 magnitude was mainly due to the much-enhanced dawnward FAC at higher latitudes since the 231 gradient (i.e., time rate of change) of  $\delta Bx$  was significantly higher at the poleward edge. The FACs 232 observed by Swarm B at nearly the same time (yellow-highlighted interval in Figure 4c) but at ~11 233 LT (Figure 4a) did not show the same feature, neither did the subsequent FACs in the pre-midnight 234 sector. When Swarm A returned to the same region in next orbit about 90 min later (~ 1615 UT), the FACs have returned to the weakened state. These observations indicate the much-enhanced 235 236 downward FAC is a localized (near dawn) and transient (duration < 90 min) phenomenon in 237 response to the sudden decrease of the solar wind dynamic pressure. The AMPERE observations 238 did not capture such a localized transient response.

239

We now examine the EEJ profiles. As Swarm B is much closer to the local noon at the dayside 240 241 equator, the EEJ signal is expected to be much stronger at Swarm B than Swarm A. Before the 242 negative pressure pulse, the EEJ profile is not well defined at Swarm A (1323 UT, Figure 4d), 243 mostly likely due to a very weak EEJ in early morning. But closer to the local noon, Swarm B 244 detected the typical eastward EEJ profile at 1252 UT (Figure 4e) and 1426 UT (Figure 4f). Then 245 about 17 min after the negative pressure pulse, Swarm A observed a well-defined westward EEJ, or 246 counter electrojet (Figure 4e). The observed counter electrojet appeared to be a transient 247 phenomenon. The EEJ returned to nominal eastward direction in the next two profiles, 1601 UT at 248 Swarm B (Figure 4i) and 1632 UT at Swarm A (Figure 4f). These observations are in agreement 249 with the ground-based EEJ currents in Figure 2. 250

- 251
- 252



Figure 4. Swarm A and B observations of FACs and the EEJ: (a) Spacecraft trajectories near the negative
pressure pulse; (b-c) the vector magnetic field residuals; (d-i) the latitudinal profiles of the EEJ around the
magnetic equator.

- 261 **3** Discussion
- 262

263 We first summarize the observations presented above.

- The solar wind dynamic pressure decreased suddenly and significantly at the boundary of
   the ICME that caused the 23 March 2023 magnetic storm. The negative pressure pulse
   arrived at the Earth at 1440 UT during the main phase of the storm and the IMF Bz
   fluctuated between northward and southward (Figure 1).
- The total large-scale FAC currents flowing into and out of the ionosphere decreased
   significantly soon after the arrival of the negative pressure pulse based (Figure 3). The
   overall geomagnetic activity level in the magnetosphere was weakened for more than 3 hrs,
   which paused the progression of the storm main phase. The activity level picked up again
   only after the IMF Bz turned strongly southward for an extended period during the passage
   of the ICME.
- Swarm A observed a significant enhancement of the downward FAC at the poleward edge
   of the FAC region near dawn shortly after the negative pressure pulse, which appeared to be
   localized and transient (Figure 4). Nearly simultaneous Swarm B observations closer to the
   local noon showed weakened FACs, consistent with the AMPERE observations.
- A transient counter electrojet was observed both in space by Swarm A (Figure 4) and on the ground (Figure 2) within minutes after the arrival of the negative pressure pulse. The counter electrojet lasted for ~ 1 hr and then returned to its regular eastward direction. The observed transient reversal of the EEJ to the westward direction suggests that the equatorial ionosphere experienced a brief period of a westward electric field after the negative pressure pulse.
- 284

These observations demonstrate the profound impact to the M-I system by the negative pressure pulse. The observed counter electrojet clearly shows that a transient westward electric field associated with the negative pressure pulse penetrated to the equatorial ionosphere from overshielding (Hori et al., 2012; Fujita et al., 2012). The penetration electric field was much stronger in magnitude than the background eastward electric field from the wind dynamo so that the overall zonal electric field was reversed. Our observations indicated there was a sudden decrease of the dawn-to-dusk (eastward) convection electric field as evident by the sudden decrease of the total

292 FAC currents flowing into and out of the polar ionosphere immediately after the negative pressure

- 293 pulse (Figure 3). The total FACs then gradually decrease with a time scale of hours. However,
- 294 SYM-H, the ring current index, was flatten out in the same period, indicating the ring current did
- not immediately respond to the weakened convection electric field (Figure 1). The delayed response
- of the ring current reflects the time scale for the M-I system to gradually adjust to the expanded
- state of the magnetosphere with decreased level of plasma convection (Earle and Kelley, 1987).
- 298 Thus, there was a short period when the low-latitude ionosphere was over-shielded and experienced
- a dusk-to-dawn (westward) electric field. Based on the duration of the counter-electrojet in the
- 300 ground-based observations (Figure 2), the response of the ring current-R2 FAC system was delayed
- 301 for  $\sim 6$  min, and it took  $\sim 1$  hr for the M-I system to gradually adjust itself to the decreased plasma
- 302 convection level and the low-latitude ionosphere to return to be fully shielded.
- 303

304 To understand the transient responses and localized enhancement of FACs, it is necessary to review 305 the current understanding of the underlying physical process. The M-I system responds to a sudden 306 pressure pulse in two phases, including a preliminary impulse (PI) and a two-stage main impulse 307 (MI) (e.g., Tamao, 1964a&b; Araki, 1977; Araki and Allen, 1982). The PI is due to the propagation and conversion of a compressional wave front launched from the magnetopause when the 308 309 magnetosphere is suddenly compressed or expanded. The PI is transient by nature because its driver 310 is the interaction between the pressure pulse and the magnetopause, which disappears in minutes 311 after the impulse front propagates away from the dayside.

312

313 Although more previous studies focused on sudden pressure increases than decreases, the basic 314 physics is the same. Based on Tamao's (1964a&b) pioneer work, Araki (1994) proposed a M-I 315 coupling PI model to explain the global observations after geomagnetic sudden commencements. 316 As illustrated in their Figure 12, the magnetopause moves inward and the dawn-to-dusk 317 magnetopause current increases when the solar wind dynamic pressure suddenly increases. A 318 compressional MHD wave is excited on the magnetopause, which propagates into the equatorial 319 magnetosphere. The solar wind-magnetosphere interaction as a dynamo generates an enhanced 320 dusk-to-down electric field at the magnetopause ( $J \cdot E < 0$ ). A dusk-to-dawn electric field and 321 associated inertia electric current are induced inside the magnetosphere. The extra magnetopause 322 current and the inertia current would form a counterclockwise current loop. The compressional

323 wave will be converted into the transverse Alfven wave due to the nonuniformity of the

magnetosphere (Tamao, 1964b; Southwood and Kivelson, 1990). When the compressional wave

front reaches the region where the Alfven speed has a largest spatial gradient, converted Alfven

326 waves are generated and propagate along the field lines with associated FACs. A pair of FACs will

327 be a part of the current loop, downward in the dusk side and upward in the dawn side. This process

328 happens in time scale of minutes. So, the pair of FACs exists transiently at lower latitudes than the

329 regular R1 currents with opposite polarity. A quantitative detail of the PI process is provided in the

330 MHD simulations by Fujita et al. (2003a&b, 2005), and the source region of the MHD wave mode

331 conversion for the generation of the transient FACs was found to be in the region of 6 < L < 7

**332** (Fujita et al., 2003a).

333

334 In the case of negative pressure pulses, the observations by Araki (1988) and simulations by Fujita 335 et al. (2004, 2012) showed that the magnetospheric and ionospheric signatures mostly mirror those 336 in pressure pulses. The negative pressure pulse causes the expansion of the magnetosphere and a 337 decrease of the magnetopause current. The PI is associated with a dawn-to-dusk transient dynamo 338 electric field at the magnetopause and induced electric field in the magnetosphere. The equatorial 339 current loop would be clockwise to effectively reduce the magnetic field strength in the 340 magnetosphere, and the pair of transient FACs would be downward in the dawnside and upward in 341 the duskside, in the same polarity of the regular R1 currents. The transient and localized 342 enhancement of the downward FAC observed by Swarm A near dawn (Figure 4) matches the 343 predicted polarity of the FACs. However, our observations differ in an important aspect from the 344 model prediction. The transient, localized FAC enhancement was observed at the poleward edge of the FAC region, implying the source region was near the magnetopause, as in the earliest work of 345 346 Tamao (1964a). Further theorical and numerical investigation is still needed to understand the 347 source region of the transient FACs during the PI. In addition, understanding the role of the ring 348 current/R2 FAC system to the undershielding/overshielding and its restoration is particularly needed 349 in future simulations.

350

351 4 Conclusions

353	A drastic decrease of the solar wind dynamic pressure occurred during main phase of the 23 March
354	2023 geomagnetic storm in association with the boundary between the ICME and its sheath. Our
355	observations show that the negative pressure pulse had significant impact to the M-I system. It
356	weakened the overall geomagnetic activities and plasma convection and paused the progression of
357	the storm main phase for $\sim$ 3 hrs. Due to the sudden decrease of the dawn-to-dusk convection
358	electric field, there was a transient period when the low-latitude ionosphere was over-shielded and
359	experienced a brief period of dusk-to-dawn (westward) penetration electric field. The transient
360	westward penetration electric field reversed the direction of the equatorial electrojet, and the
361	counter electrojet was observed both in space and on the ground. The response of the ring current-
362	R2 FAC system was delayed for ~6 min, and it took ~ 1 hr for the M-I system to adjust itself to the
363	decreased plasma convection level until the low-latitude ionosphere was fully shielded again.
364	Although the overall large-scale FACs were weakened by the negative pressure pulse, a transient,
365	localized enhancement of downward FAC was observed near dawn, consistent with the mechanism
366	for transmitting MHD disturbances in the M-I coupling after the negative pressure pulse. But the
367	latitudinal location of the localized FAC enhancement differed from the model prediction, which
368	calls further investigation of the MI coupling in response to the pressure pulse.
369	
370	
371	Acknowledgements
372	
373	GL thanks Lan Jian for helpful discussion. EY was partially supported by the AFOSR (FA9550-20-
374	1-0119) and NSF (AGS-1848730) grants. NB was partially supported by NASA Internal Scientist
375	Funding Model on Mesoscale Dynamics.
376	
377	

# 379 Data Availability Statement

- 380
- 381 The OMNI data are available at <u>https://omniweb.gsfc.nasa.gov</u>. The JICA and PIUR magnetometer
- data are available at <u>https://zenodo.org/records/10823058</u>. The KOU and TTB magnetometer data
- are available at INTERMAGNET (<u>www.intermagnet.org</u>). The AMPERE data are available at
- 384 <u>https://ampere.jhuapl.edu</u>. The Swarm data are accessible at
- 385 <u>https://earth.esa.int/eogateway/missions/swarm/data.</u>

386

387

389	References
390	
391	Alken, P., S. Maus, A. Chulliat, P. Vigneron, O. Sirol, and G. Hulot (2015), Swarm equatorial
392	electric field chain: First results, Geophys. Res. Lett., 42, 673-680.
393	https://doi.org/10.1002/2014GL062658
394	
395	Anderson, D., A. Anghel, J. Chau, and O. Veliz (2004), Daytime vertical $E \times B$ drift velocities
396	inferred from ground-based magnetometer observations at low latitudes, Space Weather, 2, S11001.
397	https://doi.org/10.1029/2004SW000095
398	
399	Araki, T. (1977), Global structure of geomagnetic sudden commencements, Planet. Space Sci., 25,
400	373. https://doi.org/10.1016/0032-0633(77)90053-8
401	
402	Araki, T. (1994). A physical model of the geomagnetic sudden commencement. Geophysical
403	Monograph-American Geophysical Union, 81, 183-183. https://doi.org/10.1029/GM081p0183
404	
405	T. Araki, J. H. Allen (1982), Latitudinal reversal of polarization of the geomagnetic sudden
406	commencement, J. Geophys. Res., 87, 5207-5216. <u>https://doi.org/10.1029/JA087iA07p05207</u>
407	
408	Araki, T. and H. Nagano (1988), Geomagnetic response to sudden expansions of the
409	magnetosphere, J. Geophys. Res., 93, 3983–3988. <u>https://doi.org/10.1029/JA093iA05p03983</u>
410	
411	Earle, G. D., and M. C. Kelley (1987), Spectral studies of the sources of ionospheric electric fields,
412	J. Geophys. Res., 92(A1), 213–224. https://doi.org.10.1029/JA092iA01p00213
413	
414	Fejer, B., C. Gonzales, D. Farley, M. Kelley, and R. Woodman (1979), Equatorial electric fields
415	during magnetically disturbed conditions 1. The effect of the interplanetary magnetic field, J.
416	Geophys. Res., 84(A10), 5797-5802. https://doi.org/10.1029/JA084iA10p05797
417	

- 418 Fejer BG, Laranja SR and Condor P (2024), Multi-process driven unusually large equatorial
- 419 perturbation electric fields during the April 2023 geomagnetic storm. Front. Astron. Space Sci.
- 420 11:1351735. <u>https://doi.org/10.3389/fspas.2024.1351735</u>
- 421
- 422 Fujita, S., T. Tanaka, T. Kikuchi, K. Fujimoto, K. Hosokawa, and M. Itonaga (2003a), A numerical
- 423 simulation of the geomagnetic sudden commencement: 1. Generation of the field-aligned current
- 424 associated with the preliminary impulse, J. Geophys. Res., 108(A12), 1416.
- 425 <u>https://doi.org/10.1029/2002JA009407</u>
- 426
- 427 Fujita, S., T. Tanaka, T. Kikuchi, K. Fujimoto, and M. Itonaga (2003b), A numerical simulation of
- 428 the geomagnetic sudden commencement: 2. Plasma processes in the main impulse, J. Geophys.
- 429 Res., 108(A12), 1417. <u>https://doi.org/10.1029/2002JA009763</u>
- 430
- Fujita, S., T. Tanaka, T. Kikuchi, and S. Tsunomura (2004), A numerical simulation of a negative
  sudden impulse, Earth Planets Space, 56, 463–472. <u>https://doi.org/10.1186/BF03352499</u>
- 433
- 434 Fujita, S., T. Tanaka, and T. Motoba (2005), A numerical simulation of the geomagnetic sudden
- 435 commencement: 3. A sudden commencement in the magnetosphere-ionosphere compound system,
- 436 J. Geophys. Res., 110, A11203. <u>https://doi.org/10.1029/2005JA011055</u>
- 437
- 438 Fujita, S., H. Yamagishi, K. T. Murata, M. Den, and T. Tanaka (2012), A numerical simulation of a
- 439 negative solar wind impulse: Revisited, J. Geophys. Res., 117, A09219.
- 440 <u>https://doi.org/10.1029/2012JA017526</u>
- 441
- 442 Heelis, R. A. (2004). Electrodynamics in the low and middle latitude ionosphere: A tutorial. *Journal*
- 443 of Atmospheric and Solar-Terrestrial Physics, 66(10), 825-
- 444 838. <u>https://doi.org/10.1016/j.jastp.2004.01.034</u>
- 445
- 446 Hori, T., A. Shinbori, N. Nishitani, T. Kikuchi, S. Fujita, T. Nagatsuma, O. Troshichev, K. Yumoto,
- 447 A. Moiseyev, and K. Seki (2012), Evolution of negative SI-induced ionospheric flows observed by

448	SuperDARN King Salmon HF radar, J. Geophys. Res., 117, A12223,
449	https://doi.org/10.1029/2012JA018093
450	
451	Jaggi, R. K., and R. A. Wolf (1973), Self-consistent calculation of the motion of a sheet of ions in
452	the magnetosphere, J. Geophys. Res.; Space Physics, 78(16), 2852-2866.
453	https://doi.org/10.1029/JA078i016p02852
454	
455	Kelley, M. C., J. J. Makela, J. L. Chau, and M. J. Nicolls (2003), Penetration of the solar wind
456	electric field into the magnetosphere/ionosphere system, Geophys. Res. Lett., 30, 1158.
457	https://doi.org/10.1029/2002GL016321
458	
459	Morschhauser, A., Brando Soares, G., Haseloff, J., Bronkalla, O., Protásio, J., Pinheiro, K., and
460	Matzka, J. (2017): The magnetic observatory on Tatuoca, Belém, Brazil: history and recent
461	developments, Geosci. Instrum. Method. Data Syst., 6, 367-376. https://doi.org/10.5194/gi-6-367-
462	2017
463	
464	Nishida, A. (1968), Coherence of geomagnetic DP 2 fluctuations with interplanetary magnetic
465	variations, J. Geophys. Res.: Space Physics, 73, 5549-5559.
466	https://doi.org/10.1029/JA073i017p05549
467	
468	Richmond, A. D. (1973). Equatorial electrojet-1. Development of a model including winds and
469	instabilities. Journal of Atmospheric and Terrestrial Physics, 35(6), 1083-
470	1103. <u>https://doi.org/10.1016/0021-9169(73)90007-x</u>
471	
472	Southwood, D. J. (1977), The role of hot plasma in magnetospheric convection, J. Geophys. Res.,
473	82(35), 5512-5520. https://doi.org/10.1029/JA082i035p05512
474	
475	Southwood, D. J., and M. G. Kivelson (1990), The magnetohydrodynamic response of the
476	magnetospheric cavity to changes in solar wind pressure, J. Geophys. Res., 95, 2301-2309.
477	https://doi.org/10.1029/JA095iA03p02301
478	

479	Tamao, T. (1964a), The structure of three-dimensional hydromagnetic waves in a uniform cold
480	plasma, J. Geomagn. Geoelectr., 16, 89–114, 1964a. https://doi.org/10.5636/jgg.16.89
481	
482	Tamao, T. (1964b), A hydromagnetic interpretation of geomagnetic SSC*, Rep. Ionos. Space Res.
483	Jpn., 18, 16–31.
484	
485	Yizengaw, E., M. B. Moldwin, A. Mebrahtu, B. Damtie, E. Zesta, C. E. Valladares, and P. Doherty
486	(2011), Comparison of storm time equatorial ionospheric electrodynamics in the African and
487	American sectors, Journal of Atmospheric and Solar-Terrestrial Physics, 73, 156-163.
488	https://doi.org/10.1016/j.jastp.2010.08.008
489	
490	Yizengaw, E., M. B. Moldwin, E. Zesta, C. M. Biouele, B. Damtie, A. Mebrahtu, B. Rabiu, C. F.
491	Valladares, and R. Stoneback (2014), The longitudinal variability of equatorial electrojet and
492	vertical drift velocity in the African and American sector, Ann. Geophys., 32, 231-238.
493	https://doi.org/10.5194/angeo-32-231-2014
494	

- 495 Yizengaw, E., Moldwin, M. B., Zesta, E., Magoun, M., Pradipta, R., Biouele, C. M., et al. (2016).
- Response of the equatorial ionosphere to the geomagnetic DP 2 current system. Geophysical 496
- 497 Research Letters, 43(14), 7364–7372. https://doi.org/10.1002/2016g1070090



#### Journal of Geophysical Research Space Physics

#### Supporting Information for

# Solar wind-magnetosphere-ionosphere coupling and its impact on equatorial ionospheric electrodynamics during the 23 March 2023 geomagnetic storm: Effect of sudden decrease of solar wind dynamic pressure

Guan Le<sup>1</sup>, Guiping Liu<sup>1</sup>, and Endawoke Yizengaw<sup>2</sup>, Chin-Chun Wu<sup>3</sup>, Yihua Zheng<sup>1</sup>, Sarah Vines<sup>4,5</sup>, Natalia Buzulukova<sup>1,6</sup>

<sup>1</sup>NASA Goddard Space Flight Center, Greenbelt, MD (guan.le@nasa.gov)
 <sup>2</sup>The Aerospace Corporation, El Segundo, CA
 <sup>3</sup>Naval Research Laboratory, Washington, DC
 <sup>4</sup>Applied Physics Laboratory, Laurel, MD
 <sup>5</sup>Now at Southwest Research Institute, San Antonio, TX
 <sup>6</sup>University of Maryland, College Park, MD

**Contents of this file** Figures S1

#### Introduction

This supporting information provides the AMPERE observations of global magnetic field perturbations and large-scale field-aligned current maps before and after the negative pressure pulse at 1440 UT on 23 March 2023.



**Figure S1**. AMPERE global maps of observed magnetic field perturbations (left), fitted magnetic field perturbations (middle), and derived large-scale field-aligned currents patterns (right): (a) Before the negative pressure pulse 1430-1440 UT; (b) After the negative pressure pulse 1440-1450 UT.