# Evolution of the Climate Forcing During the Two Years after the Hunga Tonga-Hunga Ha'apai Eruption

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18								
19	Key Points							
20	• The Jan. 15, 2022, Hunga eruption increased aerosols and H <sub>2</sub> O in the southern							
21	hemisphere stratosphere and then dispersed throughout 2022/3.							
22	• Stratospheric water vapor, ozone, temperature, and aerosol optical depth							
23	contribute to the change in downward radiative fluxes.							
24	• Hunga produced a global change in tropopause downward radiative flux of -0.17							
25	and $\pm 0.07 \text{ W/m}^2$ over the two-year period.							
26								
27								

#### 28 Abstract

29

30 We calculate the climate forcing for the two years after the January 15, 2022, Hunga 31 Tonga-Hunga Ha'apai (Hunga) eruption. We use satellite observations of stratospheric

32 aerosols, trace gases and temperatures to compute the tropopause radiative flux changes

33 relative to climatology. Overall, the net downward radiative flux decreased compared to

- 34 climatology. Although the Hunga stratospheric water vapor anomaly increases the
- 35 downward infrared radiative flux, the solar flux reduction due to Hunga aerosol shroud
- 36 dominates the net flux over most of the two-year period. Decreases in temperature 37 produced by the Hunga stratospheric circulation changes contributes to the decrease in
- 38 downward flux; however, the Hunga induced decrease in ozone increases the net short-
- 39 wave downward flux creating small sub-tropical net flux increase in late 2022.
- 40 Coincident with the aerosols settling out, the water vapor anomaly disperses, and
- 41 circulation changes disappear so that the contrasting forcings all decrease together. By
- 42 the end of 2023, most of the Hunga induced radiative forcing changes have disappeared.
- 43 There is some disagreement in the satellite stratospheric aerosol optical depth (SAOD)
- 44 which we view as a measure of the uncertainty; however, SAOD uncertainty does not
- 45 alter our conclusion that, overall, aerosols dominate the radiative flux changes followed by temperature and ozone.
- 46

#### 47 48 **Plain Language Summary**

- 49 The Hunga Tonga-Hunga Ha'apai (Hunga) submarine volcanic eruption on January
- 50 15, 2022, produced aerosol and water vapor plumes in the stratosphere. These
- 51 plumes have persisted mostly in the Southern Hemisphere throughout 2022 and
- 52 into 2023. Enhanced tropospheric warming due to the added stratospheric water
- 53 vapor is offset by the larger stratospheric aerosol attenuation of solar radiation.
- 54 Hunga induced circulation changes that reduce ozone stratospheric ozone and 55 lower temperatures also play a role in the net forcing. The change in the radiative
- 56 flux could result in a very slight 2022/3 cooling in Southern Hemisphere. The Hunga
- 57 climate forcing has decreased to near zero by the end of 2023.
- 58

#### 59 **Index Terms**

- 60 0340 Middle atmosphere dynamics
- 61 0341 Middle atmosphere: constituent transport and chemistry
- 62
- 63 0370 Volcanic effects
- 64
- 65

#### 66 1. Introduction

- 67 The eruption of the Hunga Tonga-Hunga Ha'apai (Hunga) (20.54°S, 175.38°W)
- 68 submarine volcano on Jan. 15, 2022, sent material to the mesosphere (Proud et al., 2022;
- 69 Carr et al., 2022). Microwave Limb Sounder (MLS) measurements (Millán et al., 2022,

70 hereafter M22; Santee et al., 2023) and balloon sondes measurements (Vomel et al. 2022) 71 showed that a significant amount of water vapor was injected by the eruption into the 72 tropical Southern Hemisphere (SH) mid-stratosphere. Hunga also injected at least 0.5 Tg 73 of SO<sub>2</sub> into the stratosphere (Carn et al., 2022) although this amount may have been as 74 much as 1.5 Tg (Sellitto et al., 2024). The SO<sub>2</sub> oxidation forms a sulfate aerosol layer that 75 was detected by the Ozone Mapping and Profile Suite limb profiler (OMPS) (Taha et al., 76 2022) shortly after the eruption. The MLS estimated Hunga water injection was 77 unprecedented, up to 146 Tg or ~10% increase in the total stratospheric water vapor prior 78 to the eruption (M22). The water vapor and aerosol plumes from the HT eruption have 79 persisted in the SH throughout 2022 (Schoeberl et al., 2023a, b, hereafter S23a,b). The 80 presence of water vapor led to a stratospheric cooling of ~ 4 K in March and April 81 (Schoeberl et al., 2022, hereafter S22) due to the increased outgoing IR radiation. This 82 cooling produces a secondary circulation (Coy et al., 2023) that produced temperature 83 and ozone anomaly (Wang et al., 2023) in mid 2022.

84

The volcanically generated abundance of stratospheric aerosols causes a reduction in
solar radiative forcing and, if large enough, a decrease in tropospheric temperatures
(Aurby et al., 2021; Stenchikov, 2016; Hansen et al., 2002) which has been observed
(Fujiwara et al., 2020, Crutzen, 2006, Robock, 2000). Volcanic aerosols can persist in the
stratosphere for years and even self-loft through solar heating when the aerosols are
mixed with ash (Khaykin et al., 2022).

91

92 Changes in stratospheric water vapor can also cause changes in climate forcing (Forster 93 and Shine, 1999). Solomon et al. (2010) showed that the 10-year lower tropical 94 stratospheric decrease of  $\sim 0.4$  ppmv H<sub>2</sub>O in the tropical between 2000 and  $\sim 2005$  would 95 reduce tropospheric forcing by  $\sim 0.098 \text{ W/m}^2$  or about  $\sim 0.245 \text{ W/m}^2/\text{ppmv}$ . The water 96 vapor radiative forcing results from changes in the thermal IR emission and solar flux 97 absorption. The solar flux absorption by water vapor is generally smaller than the thermal 98 emission. Extending the Solomon et al. (2010) study, Dessler et al. (2013) determined 99 calculated a water vapor climate feedback parameter of 0.27 W/m<sup>2</sup>/ppmv. Banerjee et al. 100 (2019) analyzing CMIP5 models computed the stratospheric water vapor component of 101 the climate feedback parameter to be 0.14  $W/m^2/K$  for 4xCO<sub>2</sub>. Li and Newman (2020) 102 using the Goddard Earth Observing System Chemistry-Climate model computed a similar 103  $4xCO_2$  stratospheric water vapor feedback value of 0.11 W/m<sup>2</sup>/K.

104

105 Given the sensitivity of the climate to stratospheric water vapor, it is logical to assume 106 that Hunga might have a significant climate impact. Jenkins et al. (2022) used a 107 parameterized climate-response model to investigate the climate impact of the Hunga 108 water vapor plume. They neglected the impact of aerosols and only considered the 109 radiative forcing due to the water vapor injection and computed a  $0.12 \text{ W/m}^2$  increase in 110 tropospheric radiative forcing. However, Sellitto et al. (2022) and Zhu et al. (2023) 111 roughly estimated that the aerosol plume would produce a peak solar forcing reduction of 112 ~1.7-1.8 W/m<sup>2</sup>, exceeding the estimated of H<sub>2</sub>O forcing. S23b provided a more accurate 113 estimate confirming that the aerosols overwhelmed the water vapor flux increase during 114 the first year following the eruption.

114 the first year foll 115 In this study we extend the S23b computation of the radiative forcing into the second
year following the Hunga eruption. Our basic approach is the same as S23b, but in
addition we break out the various radiative forcing components in more detail. As before
we use the OMPS measurements of stratospheric aerosol extinction (Taha et al., 2022) to
compute the stratospheric aerosol optical depth (SAOD), but we also compare NASA
OMPS SAOD to the Stratospheric Aerosol and Gas Experiment III on the international

- space station (SAGE III/ISS) measurements of SAOD, and OMPS data processed by U.
- of Saskatchewan (USask) algorithm. The USask OMPS data are processed using a
   tomographic retrieval scheme (Bourassa et al., 2023) that is different from the NASA
- 124 tomographic retrieval scheme (Bourassa et al., 2023) that is different from the NA 125 algorithm. The tomographic retrieval has the advantage of correcting OMPS
- measurements for distortion around the edges of aerosol and cloud anomalies (see
- 127 Gorkavyi et al., 2021). We refer to these data as USask OMPS.
- 128
- 129 To estimate the trace gas radiative forcing we use the AER rapid radiative transfer model
- 130 (RRTM, Mlawer et al.,1997) to compute the changes in shortwave and longwave fluxes
- 131 at the tropopause. Our approach is to use the prior 10-year climatology (2012-2021) of
- 132 MLS trace gases and temperatures and then swap in the changes observed by MLS in the
- 133 2022-2023 period to compute the relative change in radiative forcing for each
- 134 component. This allows us to quantify the relative importance of various processes
- 135 contributing to the overall radiative impact. We focus on the downward longwave and 136 shortwave flux changes at the tropopause relative to a 10-year climatology. In general,
- 137 longwave IR radiation from the mid-stratosphere will be absorbed in the cold upper
- 138 troposphere whereas shortwave radiation will penetrate to the surface. The tropospheric
- 139 climatic response to these flux changes is beyond the scope of this study. Our goal is to 140 determine the net flux changes at the tropopause as the Hunga plume evolves.
- 141

# 142 **2. Observational Data Sets**

143

We use Microwave Limb Sounder (MLS) V5 for temperature and trace gas observations.
The data quality for the Hunga anomaly is detailed in M22 and MLS data is described in
Livesey et al. (2021). Other trace gases changes are described in Santee et al. (2023).
We restrict our constituent analysis to below 35 km well above the maximum Hunga
water vapor anomaly (~25 km). In addition, the climate forcing due water vapor
emissions above 30 km is negligible (Solomon et al., 2010). The daily MLS data sets are
averaged onto a 5°x10° latitude-longitude grid.

- 151
- We use NASA OMPS level-2 V2.1 aerosol extinction data (Taha et al.,2021) which
- provides aerosol retrievals up to 40 km. Although the extinction measurements by
   OMPS V2.1 are generally consistent with those made by SAGE III/ISS (Taha et al.,
- 155 2021), as shown in Gorkavyi et al. (2021), the NASA OMPS algorithm may overestimate
- 156 the aerosol extinction at the edges and below eruption clouds because of the limb viewing
- 157 geometry. Bourassa et al. (2023) developed a tomographic retrieval scheme that corrects
- 158 for the OMPS viewing geometry problems, and the USask OMPS retrieved Hunga
- extinction levels are roughly a factor of two smaller than NASA OMPS for the first four
- 160 months after the eruption. As Bourassa et al (2023) noted the aerosol distribution is

becomes more zonal after the first four months, the aerosol edges are disappearing andthe differences between the two extinction estimates is becomes smaller.

163

For both NASA OMPS and USask OMPS, we integrate the 745 nm extinction from the 1 km above the tropopause to 35 km to obtain SAOD. We use extinction measurements at

166 745 nm since this wavelength has good sensitivity to small particles and is less

167 contaminated by Rayleigh scattering than shorter wavelengths (Taha et al., 2021). The

168 tropopause information comes from the Modern-Era Retrospective analysis for

169 Research and Applications, Version 2 (MERRA2, see Gelaro et al., 2017). We start the

170 integration above the tropopause to eliminate the extinction by thin clouds near the

tropopause. Daily data are interpolated onto a 2° latitude zonal mean daily grid; we use a

- 172 10 day box-car smoother to reduce measurement noise.
- 173

174 We use the S23b algorithm to convert SAOD to solar flux reduction; the algorithm uses 175 the 550nm SAOD. To convert the SAOD at the 745 nm wavelength to 550nm we use the 176 Ångström exponent from SAGE III/ISS (Cisewski et al., 2014) calculated using 177 extinction coefficients at 550 nm and 756 nm. We use the SAGE Ångström exponent 178 instead of one derived from OMPS, because the OMPS Ångström exponent appears 179 inconsistent with the SAGE Ångström exponent likely due to limitation of the shorter 180 wavelength retrievals in the SH and lower altitudes, and, to some extent, the algorithm's 181 particle size assumptions.

182

183 **3. Analysis of Hunga Climate Impact** 

184

# 185 **3.1 Changes in constituent distributions and temperatures following Hunga** 186 eruption.

186 187

188 To interpret the changes in downward radiative fluxes, we need to assess how the 189 constituent distribution and temperatures change following the eruption relative to 190 climatology. Some of these changes are part of year-to-year variability in the 191 stratosphere (e.g. the quasi-biennial oscillation, QBO), whereas others are induced by the 192 Hunga water vapor and aerosol anomalies. Figure 1 shows the equatorial time series of 193 aerosols, water vapor, ozone, and temperature. Overlaid on each figure is the equatorial 194 zero wind line, showing the descent of the westerly phase of the QBO starting in April 195 2022, and the easterly phase starting in April 2023.

196

Fig. 1a shows that aerosols enter the tropical region shortly after the Hunga eruption on
Jan. 15, 2022. The aerosol concentration gradient follows the zero-wind line downward
as the meridional circulation associated with the QBO pushes aerosols southward
(Schoeberl et al., 2023a). In contrast, Figure 1b shows that the water vapor anomaly
moves steadily upward as part of the Brewer-Dobson (BD) circulation. The usual taperecorder signal is also evident in the figure with ascending smaller water vapor anomalies
in August 2022, May 2023, and August 2023.

204

Associated with the QBO westerly descent Fig. 1c shows an ozone increase moving with the zero-wind line. This ozone increase is also associated with the QBO secondary 207 circulation (Plumb and Bell, 1982) which advects higher concentration of ozone

downward and creates a warm temperature anomaly seen in Fig. 1d. The reverse occursfor the descending QBO easterly phase in 2023.

210

211 Figure 2 shows the time series as in Fig. 1 but at  $40^{\circ}$ S. Aerosols and water vapor arrive at 212 this latitude mostly after May 2022. This latitude is too far from the equatorial QBO to be 213 influenced by its secondary circulation. However, the water vapor anomaly is strongly 214 correlated with a decrease in ozone and temperature starting in May 2022 and ending in 215 December 2022. Wang et al. (2023) shows that this anomaly is the result of a weakening 216 of the descending branch of the BD circulation due to in situ radiative cooling associated 217 with Hunga water vapor. The descending branch transports ozone from higher altitudes 218 into the middle stratosphere and adiabatically warms the region. As the BD circulation 219 weakens, both an ozone and temperature anomaly form. Later, the water vapor anomaly 220 disperses, and as the summer SH BD circulation weakens, the anomaly fades.

221

222 Figure 3 shows OMPS 745 nm aerosol extinction coefficient at 20 km, as well as water 223 vapor at 25 km along with changes in ozone and temperature relative to the 10-year MLS 224 climatology at 25 km. This figure provides a third perspective on constituent changes. 225 The aerosol and water vapor anomalies stay isolated in the SH except for some initial 226 transport into the Northern Hemisphere (NH) shortly after the eruption (S23a). The 227 tropical temperature decrease in Feb.-April 2022 is due to radiative cooling by water 228 vapor (Schoeberl et al., 2022). The changes in ozone at the equator – the increase in 229 May-September 2022 and decrease in the same months in 2023 are associated with the 230 QBO circulation moving downward through this altitude region.

231

In the SH extra-tropics, the March 2022 temperature (Fig. 3c) decrease is due to water vapor cooling (S22), but the later temperature and ozone decrease further south is the result of the weakening BD circulation mentioned above (Wang et al., 2023). Under normal conditions, BD circulation adiabatically heats the extra-tropics and advects ozone into the lower stratosphere. The weaker BD circulation thus causes a temperature and ozone decrease.

238

# 239 **3.2 Aerosol Direct Forcing**

240

241 3.2.1 Aerosol measurements.

242

The S23b parameterization scheme is used to compute the direct solar forcing. This
scheme, as do the schemes shown in Table 1, uses the SAOD at 550 nm derived from
extinction measurements at 745 nm and converted to 550 nm using the SAGE Ångström
exponent. Figure 4 shows time series of the NASA OMPS 745 nm SAOD (4a), the
USask 745 nm SAOD (4b), and SAGE III/ISS 756 nm SAOD (4c). The SAGE
measurements are interpolated to the OMPS regular grid, but we show the SAGE
measurement points to show where the interpolation is filling in missing data. Figure 4d

compares the OMPS SAOD measurements converted to 550nm at 20°S and the SAGE

250 compares the OMPS SAOD measurements converted to 550nm at 20°S and the
 251 550 nm SAOD measurements interpolated to the OMPS grid.

- 253 Figure 4 shows the range of SAOD values with USask nearly a factor of two smaller than 254 NASA OMPS after the eruption. above, the main difference between the two OMPS 255 SAOD values is that USask corrects for the effects of inhomogeneity along the line of 256 sight (Bourassa et al., 2023), although differences can also be caused by the difference 257 between size distribution assumptions built into the two algorithms. However, from the 258 SAGE measurements, it is apparent that USask is low biased while the NASA is high 259 biased. Figure 4d shows that after August 2022, NASA OMPS comes into agreement 260 with SAGE and all three SAOD estimates converge in mid-2023. (The slower rise in 261 SAGE SAOD after the eruption; this is due to the SAGE measurement pattern which 262 missed the initial eruption latitude.) We view the differing SAOD estimates as a measure 263 of the uncertainty. In our forcing estimates below, we will show results with both NASA OMPS SAOD and USask SAOD – these tend to bracket the SAGE estimates. 264
- 265
- 266 The evolution of the SAOD reflects the evolution of the aerosol distribution shown in
- 267 Fig. 3a (also Taha et al., 2022). Both SAOD distributions show an initial high value 268 between 30°S and the equator until May-June when the SAOD shifts south. This shift is 269 also apparent aerosols at 20 km (Fig. 3a). The southward shift is due to increased 270 seasonal eddy transport. In April-July 2023, NASA OMPS shows a new anomaly in 271 SAOD. This anomaly is less evident in the USask product but is still present. The source 272 of this anomaly is unknown but may be a movement of Hunga aerosols toward mid-273 latitudes with the formation of the Antarctic polar vortex. No SH volcanic eruptions 274 occurred during this period. In any event, the exact source of this anomaly is uncertain. 275 By the end of 2023 the SAOD anomalies have largely disappeared.
- 276

277 3.2.2 Direct forcing parameterization278

- A variety of parameterizations have been used to convert global averaged SAOD into
- 280 global average solar direct forcing change ( $\Delta A$ ) as shown in Table 1. These

281 parameterizations take the form  $\Delta A = -R$  SAOD (550nm).

- 282
- 283

Table 1 Parameterization for SAOD Solar Forcing						
Reference	R					
Hansen et al. (2002)	21					
Yu and Huang (2023)	29.5 – Clear 15.7 – All sky					
Yu et al. (2023)	23					
S23b linear-log fit	19.5 SAOD < 0.015 5.58+1.26 log <sub>e</sub> (SAOD) SAOD > 0.015					

- 285 The linear parameterizations tend to underestimate the forcing of mid-sized eruptions
- 286 (e.g. El Chichón, SAOD of ~0.05;) and overestimate very large eruptions (e.g. Pinatubo,
- 287 SAOD of 0.2, Pitarai et al., 2006). The Yu and Huang (2023) parameterization was

- 288 developed from MERRA2 tropospheric AOD estimates during a non-volcanic period. In 289 their nomenclature, 'Clear' assumes no cloud reflectivity whereas 'All sky' includes 290 climatological cloud distributions. Most estimates of direct solar forcing changes do not 291 include cloud effects because the solar flux impact of clouds is considered a separate 292 uncertainty (Hansen et al., 2002). In this study, we will use the Yu and Huang (2023) 293 algorithm to estimate the impact of clouds on the Hunga direct forcing as was done in 294 S23b. We multiply  $\Delta A$  by the cosine of minimum solar zenith angle as a function of day 295 to approximate the change in solar forcing. Shortwave forcing computed by the RRTM
- includes the zenith angle variations.
- 297

Figure 5 shows the estimated reduction in solar forcing using NASA OMPS and USask
OMPS. The southward movement of the aerosol distribution in May 2022 is reflected in
the forcing shift. We also see a slight increase in forcing due to the April-July 2023
aerosol anomaly.

**3.3 Flux changes due to trace gases and temperatures** 

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# 303

304

In S23b we only considered the water vapor shortwave and longwave IR downward flux
changes at the tropopause relative to a 5-year MLS water vapor climatology. Here we
include the changes due to temperature, ozone, and water vapor relative to a 10-year
MLS climatology. The climatology averages out the QBO induced changes in
temperature as well as any year-to-year stratospheric variability that would normally
occur. This means that some of our computed radiative flux anomalies may be due to
processes not associated with Hunga trace gas anomalies (e.g. the QBO).

312

313 Our approach is to take the 10-year climatology of temperature and trace gases, then 314 insert the one of the 2022/3 anomalies fields and compare the changes in tropopause 315 downward fluxes to the climatology. For example, we insert the observed 2022/3 water 316 vapor anomaly into the climatology and compare the altered downward fluxes to the 317 downward flux climatology. This works well for isolating the effects of the Hunga water 318 vapor since seasonal variations in stratospheric water vapor is normally ~10% and Hunga 319 anomalies were up to 5 times large than the climatology. For ozone and temperature this 320 approach is more problematic since large natural changes occur that may be unrelated to 321 Hunga (e.g. the tropical QBO (Baldwin et al., 2001) and extra-tropical stratospheric 322 warmings (Veenus et al., 2023; Tao et al., 2015)).

- 323
- 324 3.3.1 Water vapor
- 325

The Hunga stratospheric water vapor anomaly will increase the tropopause downward long-wave IR flux and reduce the short-wave flux (Solomon et al., 2010). Figure 3a shows the zonal mean water vapor anomaly at 25 km, and Fig. 6 shows the corresponding changes in downward flux at the tropopause. The longwave increase is shown in Fig.6a and the shortwave decrease is shown in Fig. 6b. The change in H<sub>2</sub>O shortwave flux combines with the aerosols to reduce the direct solar forcing. The increased longwave flux, on the other hand, is absorbed in the upper troposphere where the temperatures are

333 significantly colder than the emitting region. Figure 6b shows that the increased long

334 wave flux is on the order of ~  $0.6 \text{ W/m}^2$  which is smaller than the aerosol reduction of 335 direct solar forcing (Fig. 5).

336

337 3.3.2 Ozone

338

339 Figure 3b shows the changes in ozone at 25 km and Fig. 7 shows the changes in the 340 downward flux. The changes in ozone are mostly driven by changes in circulation either 341 natural or by the secondary circulation produced by the *in situ* radiative cooling 342 associated with the water vapor anomaly (Santee et al., 2023; Wang et al., 2023). Smaller

343 ozone change may also be due to altered chemical processes (Wilmouth et al., 2023).

344

345 As discussed in Section 3.1, ozone changes at the equator (Figs. 1c, 3b) are due to the 346 descending westerly and then easterly QBO phases. In the SH extra-tropics, a decrease in 347 ozone occurs from March 2022 through October 2022 due to the relative weakening of 348 the downward branch of the downward BD circulation by water vapor radiative cooling 349 (Coy et al., 2022; Wang et al., 2023).

350

351 Fig. 7 shows that increases in ozone cause a decrease in shortwave tropopause flux and 352 an increase in longwave tropopause flux. The reverse is true for decreases in ozone. As 353 expected, Figure 7a shows a small increase in downward longwave IR flux associated 354 with the QBO driven ozone enhancement. A larger broader decrease in IR flux is 355 associated with the southern extra-tropical ozone decrease. The longwave flux changes 356 are small relative to the changes due to water vapor (Fig. 6b). In contrast, Fig. 7b shows 357 the shortwave flux changes are larger than the water vapor shortwave flux changes. The 358 equatorial short-wave ozone flux changes are also mostly associated with the QBO with a 359 relatively large flux decrease at the equator and a smaller flux increase in the southern 360 extratropical latitudes associated with the decrease of ozone. The shortwave ozone flux 361 changes are on the scale of the aerosol flux changes (Fig. 5).

362

363 3.3.3 Flux Changes due to Temperature

364

365 The changes in stratospheric temperature also alter the downward longwave fluxes by 366 radiatively important trace gases even though the gas concentrations are not significantly 367 altered by the eruption (e.g.  $N_2O$  see Santee et al. (2023), Fig. 1). To estimate the 368

tropopause downward flux changes due to temperature changes, we use the

369 climatological trace gas concentration and swap in the 2022/3 temperatures. Figs. 3a

- 370 shows the 25 km temperature differences from climatology.
- 371

372 Fig. 8 shows the temperature induced change in downward longwave fluxes. The 373 shortwave flux is not directly affected by the temperature changes and is not shown. The 374 flux change mirrors the temperature anomalies shown in Fig. 3a especially the impact of 375 the extra-tropical SH cooling from March 2022 - November 2022. As mentioned in 376 Section 3.1, this temperature anomaly is due to the weaker BD circulation and produces a 377 significant decrease in downward long-wave flux. In general, the changes in stratospheric

378 temperature are as large or larger than the downward long-wave flux than trace gas

379 anomalies.

### 380

#### 381 **3.4 Combined Fluxes** 382

383 3.4.1 Total flux changes

384

385 Figure 9 shows the total estimated tropopause flux changes following the Hunga 386 eruptions. We combine the temperature and trace gas shortwave and longwave downward 387 fluxes with the aerosol direct forcing changes. As with the aerosol direct forcing, short 388 wave forcing by ozone and water vapor is weighted by the solar zenith angle. 389 Figure 9a shows the zonal mean flux vs time using the NASA OMPS SAOD, while Fig. 390 9b shows the zonal mean fluxes using the USask OMPS SAOD, and parts 9c-e show the 391 component fluxes at the equator, 20°S and 40°S associated with Fig. 9a. The component 392 picture shows aerosols dominating the forcing with changes in temperature and ozone 393 (short wave) contributing next. The long-wave water vapor heating, the focus of Jenkins 394 et al. (2022), appears to be one of the least important components of the total flux after 395 the first few months.

396

397 Figure 9 shows that there is net SH cooling through most of the two-year period with 398 either NASA or USask SAOD. The exception is near 20° S in Fig. 9b,d where the fluxes 399 are slightly positive from August – November 2022. This is the period where the aerosol 400 distribution shifts southward and the increase in ozone short-wave flux exceeds the 401 aerosol decrease. Aerosols and short-wave flux variations in ozone and long wave flux 402 variations in temperature dominate the total flux changes. Recall that these flux changes 403 are relative to the 10-year climatology. Thus, some of these changes are natural (e.g. 404 QBO) and some generated by the Hunga anomalous circulation (Wang et. al., 2023).

405

406 Fig. 4a shows that the Hunga SAOD anomaly persists into the beginning of 2023 then 407 reaches a small second peak at higher latitudes in May 2023. This second peak shows up 408 in the forcing (Fig. 9a,e). The combination of aerosol direct forcing, temperature changes 409 and ozone recovery lead to net decrease in downward flux in 2023. By the end of 2023 410 the aerosol forcing has dwindled to near zero. Fluctuations in the 2023 forcing 411 components, aside from aerosols, appear to be mostly due to year-to-year variability.

412

413 3.4.2 All Sky Albedo

414

415 Yu and Huang (2023) developed a cloud correction for SAOD direct forcing (all sky 416 albedo) which essentially includes cloud reflection of solar radiation before it can reach 417 the surface. Most papers computing aerosol impact of volcanic emissions or fires do not 418 include cloud effects in computing the direct forcing. In our computation, cloud 419 reflectivity is applied to all the short-wave fluxes, thus all sky albedo reduces the 420 shortwave ozone and water vapor forcing as well as the aerosol direct effect. Fig. 10a 421 shows the impact of all sky albedo on the total forcing using NASA SAOD, and Fig 10b 422 shows the results using USask SAOD. These figures should be compared to Fig. 9a,b, 423 respectively. The USask SAOD case shows the smallest net forcing as expected.

424

425 3.4.3 Global and Hemispheric Forcing

# 426

- 427 To estimate the global forcing, we integrate the downward flux from  $\pm 60^{\circ}$ . The
- 428 hemispheric flux is computed from  $60^{\circ}$  S to the equator and from the equator to  $60^{\circ}$  N. In
- 429 Fig. 11 we show the two extreme cases NASA OMPS SAOD clear skies and USask
- 430 OMPS with all sky albedo to provide an estimate of the uncertainty in the forcing. The
- 431 two cases are shown in Fig. 9a and Fig. 10b. The 2022/2023 peak and average radiative
- 432 forcing is summarized in Table 2.
- 433

Our results show that most of the period there is global net cooling, except for the JanFeb. 2022 period right after the eruption, when the water vapor forcing peaked, and

- before most of the aerosol shroud has formed. A second region of very slightly positive
  forcing occurs 10°-40°S June-December 2022 where decreases in ozone generate an
  increase in tropopause shortwave flux.
- 439

The Jan-Feb. 2022 global flux increases occurs when a NH warming– not connected with Hunga (Fig. 3a) - exceeds the SH cooling. We also note that from Figs. 9,10 and Table 2, the aerosol reduction in direct forcing is largest in the SH where to which the aerosols are confined through most of the post eruption period (Fig. 5). We conclude that the Hunga peak global forcing is  $-0.475 \pm 0.145$  W/m<sup>2</sup>. In contrast, the Pinatubo global forcing was ~ -3.5 W/m<sup>2</sup> (Pitari et al., 2016), about 6-12 times larger than Hunga.

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Table 2. Forcing amounts in $W/m^2$										
Forcing	Peak	Peak SH	Peak NH	Average	Average	Average				
roreing	I buit	I can oll	I built I thi	Trenage	Trenage	Trenage				
	Global			Global	SH	NH				
					~					
NASA OMPS	-0.59	-0.75	-0.47	-0.24	-0.43	-0.05				
Clear Sky										
USask OMPS	-0.3	-0.55	-0.39	-0.1	-0.21	0.01				
All Sky										
Average	$-0.47 \pm$	$-0.65 \pm$	-0.43 ±	-0.17 ±	-0.32 ±	-0.025 $\pm$				
	0.14	0.10	0.04	0.07	0.11	0.02				

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# 451 **4.0 Summary and Conclusions**

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We have extended the S23b estimate of the Hunga radiative forcing through 2023. We have also added the changes in the long-wave and short-wave ozone and long-wave temperature radiative fluxes to our estimates. We note that there are differences between the SAOD estimates from NASA OMPS retrievals (Taha et al., 2021) and USask tomographic retrievals (Bourassa et al., 2023) post-Hunga. Both retrievals show contrasting hias compared to SACE III/ISS SAOD measurements (Fig. 4d) in the first

458 contrasting bias compared to SAGE III/ISS SAOD measurements (Fig. 4d) in the first

half of 2022, before the aerosol distribution has become zonal. After mid 2022 the

460 NASA and USask algorithms are in better agreement. We account for the SAOD

differences by performing radiative forcing estimates for both NASA and USask SAOD

- 462 retrievals as a measure of the forcing uncertainty. We also account for tropospheric cloud
- albedo using the Yu and Huang (2023) parameterization as was done is S23b.
- 464

Our earlier conclusion (S23b) that the 2022/2023 global Hunga impact is a reduction in 465 466 tropopause flux remains valid even with the aerosol uncertainty and ozone/temperature 467 effects included. As we previously found, the tropopause flux reduction is largely due to 468 the aerosol shroud which is mostly confined to the SH. Lower stratospheric SH Hunga 469 induced temperature changes (Wang et al., 2023; Santee et al., 2023) reinforce the 470 reduction downward radiative flux. Stratospheric ozone decreases produce an increase in 471 the shortwave flux, and this is an important contributor to the total flux. By the end of 472 2023 aerosols have settled out, the water vapor anomaly has largely dispersed in the 473 lower stratosphere, and the net forcing between  $\pm 60^{\circ}$  has dissipated.

474

The Hunga eruption cooled the climate, but the amount of cooling is so small it will be

476 difficult to extract the signal from tropospheric meteorological observations. The

477 secondary circulation induced by stratospheric water vapor cooling altered the

478 stratospheric temperature and ozone distribution which significantly contributed to

- 479 Hunga changes in the climate forcing.
- 480

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# 489 **Open Research**

- 490 The RTM used to estimate H<sub>2</sub>O IR cooling rates is from Atmospheric and
- 491 Environmental Research (RTE+RRTMGP) and can be freely downloaded at
- 492 <u>http://rtweb.aer.com/rrtm\_frame.html</u>.
- 493 OMPS data, Taha et al. (2021), is available at
- 494 https://disc.gsfc.nasa.gov/datasets/OMPS NPP LP L2 AER DAILY 2/summary,
- 495 DOI: https://doi.org/<u>10.5067/CX2B9NW6FI27</u> The algorithm is documented in
- Taha et al. (2021). Data are public with unrestricted access (registration required).
- 497 The OMPS USask data is available at https://doi.org/10.5281/zenodo.7293121
- 498 Aura MLS Level 2 data, Livesey et al. (2021) JPL D-33509 Rev. C, is available at
- 499 <u>https://disc.gsfc.nasa.gov/datasets?page=1&keywords=AURA%20MLS</u>
- 500 The temperature data is available at
- 501 https://acdisc.gesdisc.eosdis.nasa.gov/data/Aura MLS Level2/ML2T.004/
- 502 The V5 water vapor data is available at
- 503 https://acdisc.gesdisc.eosdis.nasa.gov/data/Aura MLS Level2/ML2H20.005/
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#### Fig. 1

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Figure 1 Equatorial cross section of aerosols and trace gases. Part a 745 nm extinction coefficient  $(km^{-1}) \pm 1^{\circ}$  of the equator. Part b, water vapor, Part c  $\Delta$  ozone (difference from 10-year climatology). Part d,  $\Delta$  temperature all  $\pm 2.5^{\circ}$ of the equator. Month letters starting 2022 are shown in the figure. Months are show as first letters and monthly regions are divided by white lines. The thick white line is the equatorial zero wind line from MERRA2 assimilation. Color bars indicate the scale.



701 <sup>Fig. 2</sup>

702 Figure 2. Same as Fig. 1 at 40° S



Fig. 3

Figure 3 Constituent zonal mean fields. Part a, aerosol extinction coefficient at 20 km. Parts b-d at 25 km. Part b, water vapor, Part c  $\Delta$  ozone, Part c  $\Delta$  temperature. Thin white horizontal line indicates the equator. Thick black line contour is zero in parts d and d.











Figure 4 Zonal mean SAOD time series. Part a, NASA OMPS, Part b, USask OMPS, both at 745 nm. Part c, SAGE
III/ISS SAOD at 756nm interpolated onto the OMPS grid. Part d shows the 20°S SAOD converted to 550nm using the
SAGE Ångström exponent from Parts a, b, and c.







#### 716 717 Figure 6 Part a, change in long-wave IR flux from climatology due to water vapor. Part b, change in short-wave flux. Months indicated as in Fig. 1.



723 *Figure 8 As in Fig. 7a but for change in temperature.* 



Figure 9 Part a. Net forcing including trace gas fluxes and aerosol direct forcing using NASA OMPS SAOD. Part b is similar to (a) but for USask SAOD. Parts c, d, e show the various flux components at the equator (Part b), 20°S (Part c) and 40°S (Part d) using NASA OMPS SAOD.









Fig. 11

Figure 10 Variations of global (black), SH (blue) and NH (cyan) downward fluxes. Part a, clear sky, NASA OMPS SAOD (see Fig. 9a). Part b, all sky USask SAOD (Fig. 10b).