Longwave Radiative Feedback Due to Stratiform and Anvil Clouds

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Abstract

Studies have implicated the importance of longwave (LW) cloud-radiative forcing (CRF) in facilitating or accelerating the upscale development of tropical moist convection. While different cloud types are known to have distinct CRF, their individual roles in driving upscale development through radiative feedback is largely unexplored. We hypothesize that CRF from stratiform regions will have the greatest effect on upscale tropical convection. We test this hypothesis by analyzing output from convection-permitting ensemble Weather Research and Forecasting (WRF) model simulations of tropical cyclone formation. Using a novel column-by-column cloud classification scheme introduced herein, we use this model output to identify the relative contribution of five cloud types (shallow, congestus, and deep convection; and stratiform and anvil clouds) to the direct LW radiative forcing and the upscale development of convection via LW moist static energy variance. Results indicate that stratiform and anvil regions contribute dominantly to the domain averages of these variables.

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Key Points:

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6	•	A novel column-by-column cloud microphysical classification scheme was devel-
7		oped for application with numerical model output
8	•	Radiative feedback due to stratiform and anvil clouds is a leading driver of trop-
9		ical convective upscale development
10	•	The local radiative forcing by deep convective regions is similar in magnitude to
11		stratiform but its impact is limited by its smaller area

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12 Abstract

Studies have implicated the importance of longwave (LW) cloud-radiative forcing (CRF) 13 in facilitating or accelerating the upscale development of tropical moist convection. While 14 different cloud types are known to have distinct CRF, their individual roles in driving 15 upscale development through radiative feedback is largely unexplored. We hypothesize 16 that CRF from stratiform regions will have the greatest effect on upscale tropical con-17 vection. We test this hypothesis by analyzing output from convection-permitting ensem-18 ble Weather Research and Forecasting (WRF) model simulations of tropical cyclone for-19 mation. Using a novel column-by-column cloud classification scheme introduced herein, 20 we use this model output to identify the relative contribution of five cloud types (shal-21 low, congestus, and deep convection; and stratiform and anvil clouds) to the direct LW 22 radiative forcing and the upscale development of convection via LW moist static energy 23 variance. Results indicate that stratiform and anvil regions contribute dominantly to the 24 domain averages of these variables. 25

²⁶ Plain Language Summary

Infrared or longwave radiation and its interaction with clouds is important in the 27 formation of tropical storms. Given the different shapes and distributions of distinct cloud 28 types, we hypothesize that they interact with longwave radiation differently, and there-29 fore exert different impacts on the organization of tropical convection. This issue has largely 30 31 been unexplored. To address this gap, we tested our hypothesis by analyzing numerical model simulations of the formation of two tropical cyclones. Further, we developed 32 a novel cloud classification scheme based on cloud properties that identifies five distinct 33 cloud types. Our results indicate that light-raining regions, such as stratiform and anvil, 34 contribute dominantly to the domain's longwave cloud-radiative heating and the moist-35 ening of convective regions. This is due to both these cloud types' strong greenhouse trap-36 ping effect and their extensive areal coverage, which spreads this effect over large regions 37 of a developing storm. 38

³⁹ 1 Introduction

Over the past few decades, we have gained a better understanding of cloud feed-40 backs and their importance in our climate on a global and regional scale (Zelinka et al., 41 2020). For example, we now recognize that cloud-radiative forcing (CRF) is an essen-42 tial maintenance mechanism of the Madden–Julian Oscillation (MJO; Adames & Kim, 43 2016; Ciesielski et al., 2017; Najarian & Sakaeda, 2023). We have further learned of CRF's 44 part in accelerating tropical cyclone (TC) genesis (e.g., J. H. Ruppert et al., 2020; Wu 45 et al., 2021). However, there remain uncertainties regarding how cloud feedbacks and 46 CRF affect the smaller-scale dynamics of moist convection (Bony et al., 2015). 47

Studies have begun to highlight the role of CRF in the dynamics of tropical con-48 vection (e.g., Bretherton et al., 2005; Wing et al., 2016; J. H. Ruppert et al., 2020), for 49 example, through the study of self-aggregation. Self-aggregation, the spontaneous ini-50 tiation and clustering of convection, develops in idealized model frameworks that are in 51 radiative-convective equilibrium (RCE), an approximation for the real tropical atmosphere 52 (Manabe & Strickler, 1964; Bretherton et al., 2005). A budget of moist static energy (MSE) 53 variance identifies multiple pathways for promoting convective upscale growth, which shows 54 that it is the longwave (LW) cloud effect that dominates the maintenance of a mature 55 cluster (Muller & Held, 2012; Wing & Cronin, 2016; Wing et al., 2017). The inclusion 56 of rotation provides an idealized analogue for TC development, wherein self-aggregation 57 takes the form of tropical cyclogenesis (Bretherton et al., 2005; Davis, 2015; Wing et al., 58 2016). Like in non-rotating frameworks, the cloud-LW radiative feedback considerably 59 aids the development of self-aggregation, but differing from non-rotating frameworks, sur-60 face flux feedback is also important to aggregation and its maintenance (Wing et al., 2016). 61

Nonetheless, prior to the development of strong surface winds (i.e., genesis of a surface
 vortex), the cloud-LW radiative feedback dominates aggregation, which takes the form
 of upscale convective development (J. H. Ruppert et al., 2020).

Recent studies have applied the concepts from RCE model frameworks into a real-65 world context and demonstrate that LW CRF indeed accelerates TC development (J. H. Rup-66 pert et al., 2020; Wu et al., 2021). One supported hypothesis for how this works is that 67 LW CRF promotes upward motion in moist regions (Bu et al., 2014; J. H. Ruppert et 68 al., 2020) and, thus, creates a thermally direct circulation that increases moisture in those 69 70 moist areas (Bretherton et al., 2005; Needham & Randall, 2021a, 2021b). Still, there is limited understanding of how the interaction between clouds, radiative heating, and convective-71 scale motions manifests in this feedback, which constitutes an important knowledge gap 72 that we seek to address here. 73

Tropical convection is composed of and closely linked to distinct cloud types, in-74 cluding shallow cumuli, congestus, deep cumulonimbi, stratiform, and anvil clouds (Johnson 75 et al., 1999). These clouds are distinct components of mesoscale convective systems (MCSs), 76 each with unique dynamic behavior, distribution, and composition (Houze Jr., 2004). To 77 our knowledge, no study to date has identified the unique role of specific cloud types on 78 CRF and their resulting influence on convective upscale growth, which is the focus of 79 our study. We specifically address the question of how different cloud types in organized 80 convective systems uniquely promote upscale development through LW radiative forc-81 ing. 82

Given the widespread, blanketing, and long-lived nature of both anvil clouds and 83 stratiform precipitation systems (Webster & Stephens, 1980; Houze, 1997; Schumacher 84 & Houze, 2003; Ahmed & Schumacher, 2015), we specifically hypothesize that these 85 cloud systems are the most important for promoting convective upscale de-86 velopment through LW CRF. We examine this hypothesis by determining the con-87 tributions of five different cloud types to LW CRF and the LW MSE variance source term 88 using a novel classification scheme that we apply to convection-permitting ensemble Weather 89 Research and Forecasting (WRF) model simulations of tropical organized deep convec-90 tion. We leverage two TC development events to do so, Super Typhoon Haiyan (2013) 91 and Hurricane Maria (2017), though we emphasize the early simulation periods prior to 92 any intense TCs. While we highlight our analysis of Haiyan, the results from Maria (Sup-03 porting Information; SI) support our hypothesis as well. Our results can guide future observational study of CRF. Furthermore, with CRF tied to a primary source of numer-95 ical model uncertainty (Morrison et al., 2020; Zelinka et al., 2020), this work may ulti-96 mately help identify new pathways to improve the numerical model prediction of weather 97 and climate. 98

99 2 WRF Simulations

To quantify the role of different precipitating cloud types on CRF and their asso-100 ciated impact on tropical convection, we simulate the TC development cases for Haiyan 101 and Maria through numerical model simulations. These storms were chosen because they 102 developed in a typical environment for tropical cyclogenesis, including weak vertical wind 103 shear and high sea surface temperatures (SSTs; J. H. Ruppert et al., 2020), and so rep-104 resent a larger population of TC cases. To support the notion that our results are gen-105 eralizable, i.e., not specific to post-TC-genesis conditions, we include in SI (Figures S1-106 S2) our results with only the first 24 hours of the simulation for comparison. The sim-107 ulations consist of a 10-member ensemble using the Advanced Research Weather Research 108 and Forecasting model (WRF-ARW, version 4.3.1; Skamarock et al., 2021). We use a 109 nested domain with the outer domain's initial and boundary conditions initiated from 110 the first 10 ensemble members of the NOAA-NCEP Global Ensemble Forecast System 111 (2015) retrieved from the NOAA National Centers for Environmental Information. The 112

model is run with 55 stretched vertical levels with a model top at 10 hPa and a two-nest 113 approach. The inner nest is 3-km grid spacing and approximately $3,600 \times 2,200$ km in scale. 114 The simulations are integrated from 0000 UTC 1–0000 UTC 5 Nov 2013 for Haiyan and 115 0000 UTC 14–0000 UTC 18 Sept 2017. The microphysics is represented by the two-moment 116 scheme of Thompson and Eidhammer (2014). The simulations have clouds interacting 117 with radiation as in nature using the SW and LW radiation schemes from the Rapid Ra-118 diative Transfer Model for GCMs (RRTMG) (Iacono et al., 2008), which is fully coupled 119 to the microphysics scheme. The other physics settings are as in J. H. Ruppert et al. (2020). 120 Our results for Maria are shown in the SI (Figures S3-S4). For all analysis, we exclude 121 the first 12 timesteps as "spin-up" time and 80 points from domain lateral boundaries. 122

¹²³ 3 Cloud Classification

To investigate our research question, we require a classification algorithm that can 124 accurately identify a range of cloud types. Traditional precipitation classifications rely 125 on low-level reflectivity and its gradients to identify stratiform and convective precip-126 itation (e.g., Steiner et al., 1995; Biggerstaff & Listemaa, 2000; Powell et al., 2016). As 127 we seek to more comprehensively capture three-dimensional cloud coverage, however, we 128 develop a column-by-column scheme that leverages full-column model hydrometeor mass 129 information, similar to Sui et al. (2007). Since it is column-by-column, the scheme can 130 also be effectively implemented during runtime in parallelized model frameworks, which 131 is a strength that we will exploit in a forthcoming study. While the scheme is developed 132 ad hoc for our purposes and is expected to be sensitive to model microphysical scheme 133 choice, its simplicity should make future implementations of the scheme straightforward, 134 subject to adjustments as necessary. 135

To build confidence in our algorithm, we compare the spatial distributions and mean 136 vertical motion profiles using our classification to that of the traditional reflectivity-based 137 scheme of Rogers (2010), which was developed for application to model output and as-138 signs cloud type based on reflectivity at 0.4 and 3 km elevation. Although we make this 139 comparison to a reflectivity-based classification, we neither expect nor desire our results 140 to perfectly match, since the motivation behind each algorithm is different. A scheme 141 based on low-level reflectivity will likely underestimate or incorrectly classify stratiform 142 and anvil regions considering reflectivity's sensitivity to large rain drops and high rain 143 rates. Since cloud-radiation interaction is not limited to strongly precipitating clouds, 144 we have designed our algorithm with the goal of capturing this broader population of 145 radiatively interactive clouds. 146

Most classification schemes are limited to three precipitation categories: convec-147 tive, stratiform, and a third category dependent on the algorithm. The scheme by Rogers 148 (2010) includes any il as its third category and the Sui et al. (2007) classification contains 149 a mixed category between convective and stratiform. We use the model hydrometeor in-150 formation in our scheme to further separate the categories, which include deep convec-151 tive, shallow convective, congestus, stratiform, anvil, and non-precipitating. In summary, 152 we seek to develop an approach that leverages model microphysical information, captures 153 the bulk convective and stratiform behavior as validated using well-established paradigms 154 of vertical motion, and includes additional classification sub-types to capture their dis-155 tinct radiative forcing signatures. 156

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3.1 Description and Development

Our classification is summarized in Figure 1. The first step of our classification determines if a cell contains cloud. This decision is determined via a total water path (TWP; the sum of rain, cloud, graupel, snow, and ice column-integrated mixing ratios) threshold of 0.1 mm. We found a TWP threshold to be a necessary cutoff to exclude many grid columns identified as containing spurious (i.e., small magnitude) hydrometeor amounts



Figure 1. Flow chart summarizing the categorization process for our precipitation classification algorithm.

associated with negligible rainfall and radiative forcing. We compared this TWP threshold to Rogers (2010) to confirm we were not cutting out a large population of cloud (Figure S5). Adjustments to the TWP threshold primarily effects shallow convective and anvil domain fractions, but does not strongly alter the radiative forcing statistics (not shown).

Next, bulk convective and stratiform categories are separated by a cloud ratio (CR) 167 threshold, as in Sui et al. (2007). The CR is the ratio of ice water path (IWP) to liq-168 uid water path (LWP). IWP is the sum of column-integrated graupel, snow, and cloud 169 ice mixing ratios and LWP is the sum of column-integrated rain and cloud water mix-170 ing ratios. Columns with a CR < 2 are considered convective and columns with CR \geq 171 2 are considered stratiform, assuming stratiform regions will be dominated by ice hydrom-172 eteors. We again compared the CR threshold to the Rogers (2010) scheme and confirmed 173 our threshold falls between its identified convective and stratiform populations (Figure 174 S6). Convective regions are further divided between deep convective, congestus, and shal-175 low, as follows. Grid cells are marked shallow if the column-integrated rain mixing ra-176 tio falls below a threshold of 0.1 mm considering that congestus and deep convective re-177 gions would have higher rain rates (Johnson et al., 1996). Deep convective is separated 178 from congestus by a graupel mixing ratio threshold of 10^{-4} mm, with deep convective 179 regions exceeding this threshold, on the basis that congestus have limited vertical ex-180 tent beyond the 0°C level (Johnson et al., 1999) and hence limited glaciation. Stratiform 181 is separated from anvil where columns exceed a rain mixing ratio of 0.01 mm, account-182 ing for stratiform having more precipitating liquid water content than anvil clouds (Houze, 183 1997; Houze Jr., 2004). While we lack a means for comparing these sub-classifications 184 with the traditional algorithm, we assess their averaged vertical motion profiles against 185 well-established vertical structures in convective and stratiform precipitation (Steiner 186 et al., 1995; Houze Jr., 2004) in the following subsection. 187



Figure 2. Maps comparing a) the new classification scheme, the b) traditional reflectivity classification, c) the LW ACRE, and d) 1000 hPa rain water mixing ratio. See color bars for algorithm classifications. All panels show the first ensemble member of Haiyan at 36 hours.

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3.2 Comparison and Validation

We first present horizontal maps (Figure 2) comparing the new algorithm and the 189 Rogers (2010) classification. The new classification produces the following cloud frac-190 tions for the domain shown (Figure 2): deep convective as 2.76%, congestus as 1.39%, 191 shallow convective as 4.04%, stratiform as 16.42%, anvil as 13.99%, and non-precipitating 192 as 61.4%. The reflectivity approach produces the following cloud fractions: convective 193 is 3.53%, stratiform is 10.09%, anvil is 6.98%, and non-precipitating is 79.4%. The new 194 classification marks more grid cells within the domain as cloud while also increasing the 195 number of points identified as anvil and stratiform compared to the reflectivity approach. 196 The increased count for these cloud types indicates that our algorithm is more sensitive 197 to cloudy columns with lower rain rates, which is an expected and intended result, given 198 our objectives. Additionally, we see more stratiform and anvil regions enveloping the deep 199

convection regions, the latter of which are often located on edges with the reflectivity based algorithm. Our algorithm also identifies shallow convection, which is incorporated
 with the general convective category when using reflectivity.

We next present the vertical motion (w) profiles averaged for each cloud type in 203 our classification alongside the adapted Rogers (2010) scheme in Figure 3a-c. All three 204 convective w-profiles have an expected bottom-heavy profile, with deep convective max-205 imizing around 550 hPa, congestus maximizing at 850 hPa, and shallow convective max-206 imizing at 900 hPa. The profile averaged across all three convective types is consistent 207 with the shape of reflectivity-based classification profile, albeit with a smaller magnitude 208 and broader peak. These differences are consistent with the capture of more weakly pre-209 cipitating columns of convection in the new algorithm. Further, for both stratiform and 210 anvil the new algorithm produces a "top-heavy" profile, with upper-level rising and low-211 level sinking motion, with anvil having a smaller magnitude than stratiform. In compar-212 ison with the reflectivity-based approach, the profiles are similar in shape but with slight 213 differences in magnitudes and with our classification having a 50-100 hPa higher inflec-214 tion point. The overall merit of this new classification scheme is supported by the con-215 sistencies between these w-profiles in the two algorithms and, more broadly, with well-216 established paradigms documented in the literature (Steiner et al., 1995; Houze, 1997; 217 Houze Jr., 2004). 218

We conclude that the new classification algorithm accurately identifies precipita-219 tion types for this model output. The scheme is less computationally bulky than tradi-220 tional reflectivity-based schemes as it is based on microphysics thresholds and is a column-221 by-column approach. This allows for cloud classification within the framework of mod-222 els without the need of neighboring cell information, which is computationally cumber-223 some in highly parallelized frameworks commonly used for convective-scale modeling. 224 However, this approach does have some weaknesses. By having the cloud classification 225 based on microphysics thresholds, the scheme inherently relies on specific hydrometeor 226 behavior and treatment. Different microphysics schemes have vastly different treatments 227 of hydrometeors (Morrison et al., 2020) and would hence require modification to apply 228 to other microphysics schemes. This caveat extends to reflectivity-based approaches, how-229 ever, since model-based reflectivity relies upon microphysical assumptions. Additionally, 230 the new algorithm can only identify one cloud type per column. So, if there are layers 231 of different cloud types present, only the most prominent type will be identified. Despite 232 these limitations, we deem our algorithm suitable for our science question. 233

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4.1 ACRE and CRF

We next seek to quantify each cloud mode's contribution to LW CRF and ACRE. 236 Using our column-by-column based classification algorithm, we calculate the domain-237 averaged and class-averaged LW ACRE. The *domain* average measures each cloud types' 238 contribution to the total LW ACRE in the domain, while the *class* average calculates 239 the mean LW ACRE averaged only over the cells of that type. Stratiform and anvil modes 240 contribute the most to the domain-averaged LW ACRE, with averages around 10 W m^{-2} 241 (Figure 4a). Of the convective types, deep convective has the greatest contribution to 242 the LW ACRE (2.5 W m^{-2}), with congestus and shallow convective points providing 243 the smallest contributions. The large contribution to the domain average by stratiform 244 and any ill modes is partly due to the larger area coverage of these cloud types (Figure 245 2a,b). 246

4 Longwave Radiative Features of Different Cloud Types

But area coverage is not the only reason stratiform and anvil regions have the greatest domain-averaged LW ACRE. When averaging ACRE by class, stratiform and anvil regions retain the highest ACRE value (Figure 4b). They are almost an order of mag-



Figure 3. Averaged profiles of vertical motion (w) (a-c) and CRF (d-f) in convective categories (left column), stratiform and anvil (middle column), and the total convective and stratiform categories (right column). In (a-c), the new classification is presented in solid lines and the reflectivity approach appear as dashed. In (d,e), the dashed black lines represent the averages across all shown categories within the respective panel. Plots include values from all 10 members and 85 timesteps of Haiyan.



Figure 4. Boxplots of domain-averaged (left) and class-averaged (right) LW ACRE (top) and $\hat{h}'NETLW'$ (bottom) by precipitation type. White circles indicate the mean value. Black diamonds represent outliers. Domain averages represent class-averaged values normalized by total grid cell count, and class-averaged values are normalized only by category cell count. Plots include values from all 10 members and 85 timesteps of Haiyan.

²⁵⁰ nitude greater than that of the congestus and shallow convective types. Deep convec-²⁵¹ tive has the greatest LW ACRE value of the three convective modes, with an average ²⁵² value of 85 W m⁻². This value is comparable in magnitude to that of stratiform and ²⁵³ anvil regions, suggesting comparable radiative forcing by these categories within a given ²⁵⁴ column. The combination of stratiform and anvil's large class-averaged LW ACRE with ²⁵⁵ their larger area coverage explains their much larger contribution to the domain-averaged ²⁵⁶ LW ACRE.

We next present vertical profiles of LW CRF to aid interpretation of these results 257 (Figure 3d-f). Cloud types with larger ACRE values exhibit a deep layer of positive CRF. 258 Stratiform and anvil modes are similar in CRF shape and magnitude, which is notewor-259 thy given their much different vertical motion magnitude (Figure 3b). Namely, the ra-260 diative forcing per unit vertical mass flux within a given layer is much larger for anvil 261 than stratiform. CRF in these cloud modes is positive from 200 hPa to the surface and 262 with maxima around 5 K day⁻¹ at 300 hPa, with strong cloud-top cooling above 200 hPa. 263 Of the convective categories, deep convective has the deepest layer of positive CRF of 264 2 K day⁻¹ from 200 hPa to the surface, with a small layer near 0 K day⁻¹ between 400 265 and 500 hPa. This may be due to layers of detrained cloud in association with the $0^{\circ}C$ 266 stable layer (Johnson et al., 1996). Like stratiform and anvil, deep convective has a strong 267 signature of cloud-top cooling above 200 hPa. Congestus and shallow convective modes 268 have maxima in the lower troposphere with cooling above due to their low cloud-top height. 269 Above 800 hPa, their CRF hovers around 0 K day⁻¹. The modest heating values in the 270 upper troposphere in these categories are likely a result of the algorithm only identify-271 ing one cloud type for each column, with these columns potentially including thin anvil 272 clouds. Otherwise, these results are consistent with expectations of cloud depth based 273 on mean vertical motion (Figure 3 a-c). 274

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4.2 LW MSE Variance Source Term

The MSE budget is a tool that allows us to assess and quantify upscale develop-276 ment and intensification of convection. This tool has been used to assess the influence 277 of radiative feedback in relation to both self-aggregation (Wing & Emanuel, 2014) and 278 TC genesis (J. H. Ruppert et al., 2020; Wu et al., 2021). Our interest in this budget is 279 primarily in the LW-source term as it has been shown to be the dominant term for main-280 taining and accelerating convective upscale development (at least, prior to TC genesis) 281 (Wing & Emanuel, 2014; Wing & Cronin, 2016; J. H. Ruppert et al., 2020). The LW MSE 282 variance term (h'NetLW') is the correlation between the anomaly of the density-weighted 283 vertical integral of MSE (h') and the anomalous column LW radiative flux convergence 284 (NetLW'), where anomalies are calculated as the deviation from the domain average. More 285 details on the calculation of this term can be found in Wing and Emanuel (2014) and 286 J. H. Ruppert et al. (2020). 287

When we average h'NetLW' by cloud type, we once again see that stratiform and 288 anvil regions contribute the most to the domain-averaged LW MSE variance source term. 289 Stratiform dominates, with an average of about 0.035 day^{-1} (Figure 4c). The stratiform 290 regions also have the highest class-averaged LW MSE variance of all the cloud types (Fig-291 ure 4d). Congestus and shallow convective points have the lowest class-averaged LW MSE 292 variance, as may be anticipated from LW ACRE (Figure 4 a-b). Surprisingly, deep con-293 vective regions have the second highest LW MSE variance, almost matching that of strat-294 iform. Anvil's class average follows deep convective with a value of 0.25 day^{-1} compared 295 to stratiform and deep connective's averages of about 0.4 day⁻¹. Like the LW ACRE, 296 the stratiform and anvil regions have much greater areal coverage than convective re-297 gions (Figure 2a), which explains the smaller domain-averaged value for convection. 298

The high column influence of stratiform and anvil clouds in terms of both LW ACRE and \hat{h}' NetLW', combined with their extensive areal coverage, indicates their unique importance for supporting upscale convective development in the tropics by amplifying MSE
 variance via radiative forcing. Although deep convective regions also have high values
 of LW MSE variance in a given column, the small regional coverage of this cloud type
 likely limits its area-averaged impact, which is how radiative forcing is linked to the ten dency of MSE variance (Wing & Emanuel, 2014). These findings emphasize the impor tant influences of stratiform and anvil regions on tropical convective organization through
 their radiative forcing, which have not been previously examined in this manner.

5 Summary and Conclusions

In this study, we have investigated the role of CRF of five different cloud types and 309 their ability to aid the organization of tropical convection via the analysis of convection-310 permitting WRF simulations conducted in the context of TC development. To accom-311 plish this, we developed a novel column-by-column cloud classification algorithm based 312 on microphysical thresholds. Our classification holds several advantages over low-level 313 reflectivity-based approaches, such as it is computationally efficient and can be run within 314 the framework of a numerical model, it is sensitive to cloud (including non-precipitating) 315 throughout the column, and it identifies five cloud modes (instead of two or three): shal-316 low, congestus, and deep convective; and stratiform and anvil. However, the disadvan-317 tages to this algorithm includes its likely sensitivity to different microphysics schemes 318 based on its threshold approach and that it can only identify one cloud type per column. 319 Despite these disadvantages, this approach to cloud classification allows for more ques-320 tions to be answered on the influence of cloud type, including our question of how dif-321 ferent cloud types in organized convective systems promote upscale development though 322 LW radiative forcing. 323

We hypothesized that stratiform and anvil regions would support convective or-324 ganization more than other categories through LW ACRE and the LW MSE variance 325 source term. We found that stratiform and anvil contributed the most to the domain-326 averaged ACRE and had greater class-averaged ACRE than that of the other types, in-327 dicating their important contribution to the direct LW radiative forcing. For the LW MSE 328 variance source term, stratiform and anvil again contributed the most to the domain av-329 erage. However, the class averages revealed that deep convective was on par with strat-330 iform regions, resulting in those two classes having the highest class averages. Anvil was 331 third highest, followed by shallow and congestus, which were much weaker. While the 332 class-averaged LW MSE variance source term indicates that the localized forcing by deep 333 convective, stratiform, and anvil clouds is comparable, anvil and stratiform clouds dom-334 inate in supporting convective upscale development owing to their much greater area cov-335 erage. Although we do not fully answer the question of how different cloud types in or-336 ganized convective systems uniquely promote upscale development through LW radia-337 tive forcing, we do provide support of our hypothesis and shed new light on the specific 338 cloud types most important to convective upscale via LW cloud feedback. Future work 339 will focus on the mechanisms CRF works through to promote organization within trop-340 ical convection. 341

³⁴² Open Research Section

The code needed to recreate the WRF simulations described in this study is published at Zenodo (J. Ruppert & Zhang, 2024). The code for the precipitation classification algorithm (Luschen & Ruppert, 2024b) and the analysis (Luschen & Ruppert, 2024a) presented here are available on Zenodo as well.

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³⁵⁵ contributed to this research.

356 References

- Adames, Á. F., & Kim, D. (2016, mar). The MJO as a Dispersive, Convectively Coupled Moisture Wave: Theory and Observations. *Journal of the Atmospheric Sciences*, 73(3), 913–941. doi: 10.1175/JAS-D-15-0170.1
- Ahmed, F., & Schumacher, C. (2015). Convective and stratiform components of the precipitation-moisture relationship. *Geophysical Research Letters*, 42(23), 10,453–10,462. doi: 10.1002/2015GL066957
- Biggerstaff, M. I., & Listemaa, S. A. (2000, December). An Improved Scheme for Convective/Stratiform Echo Classification Using Radar Reflectivity. Journal of Applied Meteorology and Climatology, 39(12), 2129–2150. doi: 10.1175/1520
 -0450(2001)040(2129:AISFCS)2.0.CO;2
- Bony, S., Stevens, B., Frierson, D. M. W., Jakob, C., Kageyama, M., Pincus, R., ...
 Webb, M. J. (2015, April). Clouds, circulation and climate sensitivity. *Nature Geoscience*, 8(4), 261–268. doi: 10.1038/ngeo2398
- 370Bretherton, C. S., Blossey, P. N., & Khairoutdinov, M. (2005, December).An371Energy-Balance Analysis of Deep Convective Self-Aggregation above Uni-372form SST. Journal of the Atmospheric Sciences, 62(12), 4273–4292.37310.1175/JAS3614.1
- Bu, Y. P., Fovell, R. G., & Corbosiero, K. L. (2014, May). Influence of
 Cloud–Radiative Forcing on Tropical Cyclone Structure. Journal of the At mospheric Sciences, 71(5), 1644–1662. doi: 10.1175/JAS-D-13-0265.1
- Ciesielski, P. E., Johnson, R. H., Jiang, X., Zhang, Y., & Xie, S. (2017). Relationships between radiation, clouds, and convection during DYNAMO. Journal of Geophysical Research: Atmospheres, 122(5), 2529–2548. doi: 10.1002/2016JD025965
- Davis, C. A. (2015, September). The Formation of Moist Vortices and Tropical
 Cyclones in Idealized Simulations. *Journal of the Atmospheric Sciences*, 72(9),
 3499–3516. doi: 10.1175/JAS-D-15-0027.1
- Houze, R. A. (1997, oct). Stratiform Precipitation in Regions of Convection: A Meteorological Paradox? Bulletin of the American Meteorological Society, 78(10), 2179–2196. doi: 10.1175/1520-0477(1997)078(2179:SPIROC)2.0.CO;2
- Houze Jr., R. A. (2004). Mesoscale convective systems. Reviews of Geophysics, 42(4). doi: 10.1029/2004RG000150
- Iacono, M. J., Delamere, J. S., Mlawer, E. J., Shephard, M. W., Clough, S. A., &
 Collins, W. D. (2008). Radiative forcing by long-lived greenhouse gases: Cal culations with the AER radiative transfer models. Journal of Geophysical
 Research: Atmospheres, 113(D13). doi: 10.1029/2008JD009944
- Johnson, R. H., Ciesielski, P. E., & Hart, K. A. (1996, jul). Tropical Inversions near
 the 0°C Level. Journal of the Atmospheric Sciences, 53(13), 1838–1855. doi:
 10.1175/1520-0469(1996)053(1838:TINTL)2.0.CO;2
- Johnson, R. H., Rickenbach, T. M., Rutledge, S. A., Ciesielski, P. E., & Schubert,
 W. H. (1999, aug). Trimodal Characteristics of Tropical Convection. Journal of Climate, 12(8), 2397–2418. doi: 10.1175/1520-0442(1999)012(2397: TCOTC)2.0.CO;2

400	Luschen, E., & Ruppert, J. (2024a, 2). Analysis Code for Geophysical Research Let-
401	ter paper. doi: 10.5281 /zenodo. 10645132
402	Luschen, E., & Ruppert, J. (2024b, 2). Column-based Precipitation Classification Al-
403	gorithm. doi: $10.5281/zenodo.10611873$
404	Manabe, S., & Strickler, R. F. (1964, July). Thermal Equilibrium of the Atmosphere
405	with a Convective Adjustment. Journal of the Atmospheric Sciences, $21(4)$,
406	361–385. doi: $10.1175/1520-0469(1964)021(0361:TEOTAW)2.0.CO;2$
407	Morrison, H., van Lier-Walqui, M., Fridlind, A. M., Grabowski, W. W., Harrington,
408	J. Y., Hoose, C., Xue, L. (2020). Confronting the Challenge of Modeling
409	Cloud and Precipitation Microphysics. Journal of Advances in Modeling Earth
410	Systems, 12(8), e2019MS001689. doi: 10.1029/2019MS001689
411	Muller, C. J., & Held, I. M. (2012, August). Detailed Investigation of the Self-
412	Aggregation of Convection in Cloud-Resolving Simulations. Journal of the At-
413	mospheric Sciences, 09(8), 2551-2505. doi: 10.1175/JAS-D-11-0257.1
414	Dadiating Enging During DVNAMO (AMIE – Journal of Coophysical Bassarch
415	Aduative Forcing During D1NAMO/AMIE. Journal of Geophysical Research:
416	Nachham M. P. & Pandall D. A. (2021a) Linking Atmospheric Cloud Padia
417	tive Effects and Tropical Precipitation Coophysical Research Letters 18(14)
410	e2021GL094004 doi: 10.1029/2021GL094004
420	Needham M B & Bandall D A (2021b) Riehl and Malkus Revisited: The Role
421	of Cloud Radiative Effects. Journal of Geophysical Research: Atmospheres.
422	<i>126</i> (16), e2021JD035019. doi: 10.1029/2021JD035019
423	NOAA-NCEP Global Ensemble Forecast System. (2015). NCEP GFS 0.25 De-
424	gree Global Forecast Grids Historical Archive. UCAR/NCAR - Research Data
425	Archive. doi: 10.5065/D65D8PWK
426	Powell, S. W., Houze, R. A., & Brodzik, S. R. (2016, March). Rainfall-Type Catego-
427	rization of Radar Echoes Using Polar Coordinate Reflectivity Data. Journal of
428	Atmospheric and Oceanic Technology, 33(3), 523–538. doi: 10.1175/JTECH-D
429	-15-0135.1
430	Rogers, R. (2010, January). Convective-Scale Structure and Evolution during a
431	High-Resolution Simulation of Tropical Cyclone Rapid Intensification. Journal
432	of the Atmospheric Sciences, 67(1), 44–70. doi: 10.1175/2009JAS3122.1
433	Ruppert, J., & Zhang, Y. (2024, 2). Ensemble WRF Simulations of Typhoon Haiyan
434	and Hurricane Maria. doi: 10.5281/zenodo.10572959
435	Ruppert, J. H., Wing, A. A., Tang, X., & Duran, E. L. (2020, November). The crit-
436	Ical role of cloud-infrared radiation feedback in tropical cyclone development. $P_{\text{reaccodings}}$ of the National Academy of Sciences $117/(45)$ 27884 27802 doi:
437	$10\ 1073\ /\text{pms}\ 9013584117$
438	Schumacher C k Houze R A (2003 jun) Stratiform Bain in the Tron-
439	ics as Seen by the TRMM Precipitation Badar* Journal of Climate
441	16(11), $1739-1756$. Retrieved from http://journals.ametsoc.org/
442	doi/10.1175/1520-0442(2003)016%3C1739:SRITTA%3E2.0.CO;2 doi:
443	10.1175/1520-0442(2003)016(1739:SRITTA)2.0.CO;2
444	Skamarock, W. C., Klemp, J. B., Dudhia, J. B., Gill, D. O., Barker, D. M., Duda,
445	M. G., Powers, J. G. (2021). A Description of the Advanced Re-
446	search WRF Model Version 4.3. NCAR Technical Note(July), 1–165. doi:
447	10.5065/1dfh-6p97
448	Steiner, M., Houze, R. A., & Yuter, S. E. (1995, September). Climatological Char-
449	acterization of Three-Dimensional Storm Structure from Operational Radar
450	and Rain Gauge Data. Journal of Applied Meteorology and Climatology, $34(9)$,
451	1978–2007. doi: $10.1175/1520-0450(1995)034(1978:CCOTDS)2.0.CO;2$
452	Sui, CH., Tsay, CT., & Li, X. (2007). Convective–stratiform rainfall separation by
453	cloud content. Journal of Geophysical Research: Atmospheres, 112(D14). doi:
454	10.1029/2006JD008082

455	Thompson, G., & Eidhammer, T. (2014, October). A Study of Aerosol Impacts on
456	Clouds and Precipitation Development in a Large Winter Cyclone. Journal of
457	the Atmospheric Sciences, 71(10), 3636–3658. doi: 10.1175/JAS-D-13-0305.1
458	Webster, P. J., & Stephens, G. L. (1980, jul). Tropical Upper-Tropospheric Ex-
459	tended Clouds: Inferences from Winter MONEX. Journal of the Atmospheric
460	Sciences, 37(7), 1521-1541. Retrieved from http://journals.ametsoc.org/
461	doi/abs/10.1175/1520-0469-37.7.1521 doi: 10.1175/1520-0469-37.7.1521
462	Wing, A. A., Camargo, S. J., & Sobel, A. H. (2016, July). Role of Radia-
463	tive–Convective Feedbacks in Spontaneous Tropical Cyclogenesis in Idealized
464	Numerical Simulations. Journal of the Atmospheric Sciences, 73(7), 2633–
465	2642. doi: 10.1175/JAS-D-15-0380.1
466	Wing, A. A., & Cronin, T. W. (2016). Self-aggregation of convection in long channel
467	geometry. Quarterly Journal of the Royal Meteorological Society, 142(694), 1-
	15 doi: 10.1002/ai.2628
468	$15. \ dol. \ 10.1002/\text{Q}.2020$
468 469	Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-
468 469 470	Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self- Aggregation in Numerical Simulations: A Review. Surveys in Geophysics,
468 469 470 471	Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self- Aggregation in Numerical Simulations: A Review. <i>Surveys in Geophysics</i> , 38(6). doi: 10.1007/s10712-017-9408-4
468 469 470 471 472	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling
468 469 470 471 472 473	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simula-
468 469 470 471 472 473 474	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi:
468 469 470 471 472 473 474 475	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269
468 469 470 471 472 473 474 475 476	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269 Wu, SN., Soden, B. J., & Nolan, D. S. (2021). Examining the Role of Cloud
468 469 470 471 472 473 474 475 476 477	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269 Wu, SN., Soden, B. J., & Nolan, D. S. (2021). Examining the Role of Cloud Radiative Interactions in Tropical Cyclone Development Using Satellite Mea-
468 469 470 471 472 473 474 475 475 477 478	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269 Wu, SN., Soden, B. J., & Nolan, D. S. (2021). Examining the Role of Cloud Radiative Interactions in Tropical Cyclone Development Using Satellite Measurements and WRF Simulations. Geophysical Research Letters, 48(15),
468 469 470 471 472 473 474 475 476 477 478 479	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269 Wu, SN., Soden, B. J., & Nolan, D. S. (2021). Examining the Role of Cloud Radiative Interactions in Tropical Cyclone Development Using Satellite Measurements and WRF Simulations. Geophysical Research Letters, 48(15), e2021GL093259. doi: 10.1029/2021GL093259
468 469 470 471 472 473 474 475 476 477 478 479 480	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269 Wu, SN., Soden, B. J., & Nolan, D. S. (2021). Examining the Role of Cloud Radiative Interactions in Tropical Cyclone Development Using Satellite Measurements and WRF Simulations. Geophysical Research Letters, 48(15), e2021GL093259. doi: 10.1029/2021GL093259 Zelinka, M. D., Myers, T. A., McCoy, D. T., Po-Chedley, S., Caldwell, P. M.,
468 469 470 471 472 473 474 475 476 477 478 478 478 478 480 481	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269 Wu, SN., Soden, B. J., & Nolan, D. S. (2021). Examining the Role of Cloud Radiative Interactions in Tropical Cyclone Development Using Satellite Measurements and WRF Simulations. Geophysical Research Letters, 48(15), e2021GL093259. doi: 10.1029/2021GL093259 Zelinka, M. D., Myers, T. A., McCoy, D. T., Po-Chedley, S., Caldwell, P. M., Ceppi, P., Taylor, K. E. (2020). Causes of higher climate sensitivity in

483 https://doi.org/10.1029/2019GL085782

Longwave Radiative Feedback Due to Stratiform and Anvil Clouds

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Key Points:

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6	•	A novel column-by-column cloud microphysical classification scheme was devel-
7		oped for application with numerical model output
8	•	Radiative feedback due to stratiform and anvil clouds is a leading driver of trop-
9		ical convective upscale development
10	•	The local radiative forcing by deep convective regions is similar in magnitude to
11		stratiform but its impact is limited by its smaller area

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12 Abstract

Studies have implicated the importance of longwave (LW) cloud-radiative forcing (CRF) 13 in facilitating or accelerating the upscale development of tropical moist convection. While 14 different cloud types are known to have distinct CRF, their individual roles in driving 15 upscale development through radiative feedback is largely unexplored. We hypothesize 16 that CRF from stratiform regions will have the greatest effect on upscale tropical con-17 vection. We test this hypothesis by analyzing output from convection-permitting ensem-18 ble Weather Research and Forecasting (WRF) model simulations of tropical cyclone for-19 mation. Using a novel column-by-column cloud classification scheme introduced herein, 20 we use this model output to identify the relative contribution of five cloud types (shal-21 low, congestus, and deep convection; and stratiform and anvil clouds) to the direct LW 22 radiative forcing and the upscale development of convection via LW moist static energy 23 variance. Results indicate that stratiform and anvil regions contribute dominantly to the 24 domain averages of these variables. 25

²⁶ Plain Language Summary

Infrared or longwave radiation and its interaction with clouds is important in the 27 formation of tropical storms. Given the different shapes and distributions of distinct cloud 28 types, we hypothesize that they interact with longwave radiation differently, and there-29 fore exert different impacts on the organization of tropical convection. This issue has largely 30 31 been unexplored. To address this gap, we tested our hypothesis by analyzing numerical model simulations of the formation of two tropical cyclones. Further, we developed 32 a novel cloud classification scheme based on cloud properties that identifies five distinct 33 cloud types. Our results indicate that light-raining regions, such as stratiform and anvil, 34 contribute dominantly to the domain's longwave cloud-radiative heating and the moist-35 ening of convective regions. This is due to both these cloud types' strong greenhouse trap-36 ping effect and their extensive areal coverage, which spreads this effect over large regions 37 of a developing storm. 38

³⁹ 1 Introduction

Over the past few decades, we have gained a better understanding of cloud feed-40 backs and their importance in our climate on a global and regional scale (Zelinka et al., 41 2020). For example, we now recognize that cloud-radiative forcing (CRF) is an essen-42 tial maintenance mechanism of the Madden–Julian Oscillation (MJO; Adames & Kim, 43 2016; Ciesielski et al., 2017; Najarian & Sakaeda, 2023). We have further learned of CRF's 44 part in accelerating tropical cyclone (TC) genesis (e.g., J. H. Ruppert et al., 2020; Wu 45 et al., 2021). However, there remain uncertainties regarding how cloud feedbacks and 46 CRF affect the smaller-scale dynamics of moist convection (Bony et al., 2015). 47

Studies have begun to highlight the role of CRF in the dynamics of tropical con-48 vection (e.g., Bretherton et al., 2005; Wing et al., 2016; J. H. Ruppert et al., 2020), for 49 example, through the study of self-aggregation. Self-aggregation, the spontaneous ini-50 tiation and clustering of convection, develops in idealized model frameworks that are in 51 radiative-convective equilibrium (RCE), an approximation for the real tropical atmosphere 52 (Manabe & Strickler, 1964; Bretherton et al., 2005). A budget of moist static energy (MSE) 53 variance identifies multiple pathways for promoting convective upscale growth, which shows 54 that it is the longwave (LW) cloud effect that dominates the maintenance of a mature 55 cluster (Muller & Held, 2012; Wing & Cronin, 2016; Wing et al., 2017). The inclusion 56 of rotation provides an idealized analogue for TC development, wherein self-aggregation 57 takes the form of tropical cyclogenesis (Bretherton et al., 2005; Davis, 2015; Wing et al., 58 2016). Like in non-rotating frameworks, the cloud-LW radiative feedback considerably 59 aids the development of self-aggregation, but differing from non-rotating frameworks, sur-60 face flux feedback is also important to aggregation and its maintenance (Wing et al., 2016). 61

Nonetheless, prior to the development of strong surface winds (i.e., genesis of a surface
 vortex), the cloud-LW radiative feedback dominates aggregation, which takes the form
 of upscale convective development (J. H. Ruppert et al., 2020).

Recent studies have applied the concepts from RCE model frameworks into a real-65 world context and demonstrate that LW CRF indeed accelerates TC development (J. H. Rup-66 pert et al., 2020; Wu et al., 2021). One supported hypothesis for how this works is that 67 LW CRF promotes upward motion in moist regions (Bu et al., 2014; J. H. Ruppert et 68 al., 2020) and, thus, creates a thermally direct circulation that increases moisture in those 69 70 moist areas (Bretherton et al., 2005; Needham & Randall, 2021a, 2021b). Still, there is limited understanding of how the interaction between clouds, radiative heating, and convective-71 scale motions manifests in this feedback, which constitutes an important knowledge gap 72 that we seek to address here. 73

Tropical convection is composed of and closely linked to distinct cloud types, in-74 cluding shallow cumuli, congestus, deep cumulonimbi, stratiform, and anvil clouds (Johnson 75 et al., 1999). These clouds are distinct components of mesoscale convective systems (MCSs), 76 each with unique dynamic behavior, distribution, and composition (Houze Jr., 2004). To 77 our knowledge, no study to date has identified the unique role of specific cloud types on 78 CRF and their resulting influence on convective upscale growth, which is the focus of 79 our study. We specifically address the question of how different cloud types in organized 80 convective systems uniquely promote upscale development through LW radiative forc-81 ing. 82

Given the widespread, blanketing, and long-lived nature of both anvil clouds and 83 stratiform precipitation systems (Webster & Stephens, 1980; Houze, 1997; Schumacher 84 & Houze, 2003; Ahmed & Schumacher, 2015), we specifically hypothesize that these 85 cloud systems are the most important for promoting convective upscale de-86 velopment through LW CRF. We examine this hypothesis by determining the con-87 tributions of five different cloud types to LW CRF and the LW MSE variance source term 88 using a novel classification scheme that we apply to convection-permitting ensemble Weather 89 Research and Forecasting (WRF) model simulations of tropical organized deep convec-90 tion. We leverage two TC development events to do so, Super Typhoon Haiyan (2013) 91 and Hurricane Maria (2017), though we emphasize the early simulation periods prior to 92 any intense TCs. While we highlight our analysis of Haiyan, the results from Maria (Sup-03 porting Information; SI) support our hypothesis as well. Our results can guide future observational study of CRF. Furthermore, with CRF tied to a primary source of numer-95 ical model uncertainty (Morrison et al., 2020; Zelinka et al., 2020), this work may ulti-96 mately help identify new pathways to improve the numerical model prediction of weather 97 and climate. 98

99 2 WRF Simulations

To quantify the role of different precipitating cloud types on CRF and their asso-100 ciated impact on tropical convection, we simulate the TC development cases for Haiyan 101 and Maria through numerical model simulations. These storms were chosen because they 102 developed in a typical environment for tropical cyclogenesis, including weak vertical wind 103 shear and high sea surface temperatures (SSTs; J. H. Ruppert et al., 2020), and so rep-104 resent a larger population of TC cases. To support the notion that our results are gen-105 eralizable, i.e., not specific to post-TC-genesis conditions, we include in SI (Figures S1-106 S2) our results with only the first 24 hours of the simulation for comparison. The sim-107 ulations consist of a 10-member ensemble using the Advanced Research Weather Research 108 and Forecasting model (WRF-ARW, version 4.3.1; Skamarock et al., 2021). We use a 109 nested domain with the outer domain's initial and boundary conditions initiated from 110 the first 10 ensemble members of the NOAA-NCEP Global Ensemble Forecast System 111 (2015) retrieved from the NOAA National Centers for Environmental Information. The 112

model is run with 55 stretched vertical levels with a model top at 10 hPa and a two-nest 113 approach. The inner nest is 3-km grid spacing and approximately $3,600 \times 2,200$ km in scale. 114 The simulations are integrated from 0000 UTC 1–0000 UTC 5 Nov 2013 for Haiyan and 115 0000 UTC 14–0000 UTC 18 Sept 2017. The microphysics is represented by the two-moment 116 scheme of Thompson and Eidhammer (2014). The simulations have clouds interacting 117 with radiation as in nature using the SW and LW radiation schemes from the Rapid Ra-118 diative Transfer Model for GCMs (RRTMG) (Iacono et al., 2008), which is fully coupled 119 to the microphysics scheme. The other physics settings are as in J. H. Ruppert et al. (2020). 120 Our results for Maria are shown in the SI (Figures S3-S4). For all analysis, we exclude 121 the first 12 timesteps as "spin-up" time and 80 points from domain lateral boundaries. 122

123 **3** Cloud Classification

To investigate our research question, we require a classification algorithm that can 124 accurately identify a range of cloud types. Traditional precipitation classifications rely 125 on low-level reflectivity and its gradients to identify stratiform and convective precip-126 itation (e.g., Steiner et al., 1995; Biggerstaff & Listemaa, 2000; Powell et al., 2016). As 127 we seek to more comprehensively capture three-dimensional cloud coverage, however, we 128 develop a column-by-column scheme that leverages full-column model hydrometeor mass 129 information, similar to Sui et al. (2007). Since it is column-by-column, the scheme can 130 also be effectively implemented during runtime in parallelized model frameworks, which 131 is a strength that we will exploit in a forthcoming study. While the scheme is developed 132 ad hoc for our purposes and is expected to be sensitive to model microphysical scheme 133 choice, its simplicity should make future implementations of the scheme straightforward, 134 subject to adjustments as necessary. 135

To build confidence in our algorithm, we compare the spatial distributions and mean 136 vertical motion profiles using our classification to that of the traditional reflectivity-based 137 scheme of Rogers (2010), which was developed for application to model output and as-138 signs cloud type based on reflectivity at 0.4 and 3 km elevation. Although we make this 139 comparison to a reflectivity-based classification, we neither expect nor desire our results 140 to perfectly match, since the motivation behind each algorithm is different. A scheme 141 based on low-level reflectivity will likely underestimate or incorrectly classify stratiform 142 and anvil regions considering reflectivity's sensitivity to large rain drops and high rain 143 rates. Since cloud-radiation interaction is not limited to strongly precipitating clouds, 144 we have designed our algorithm with the goal of capturing this broader population of 145 radiatively interactive clouds. 146

Most classification schemes are limited to three precipitation categories: convec-147 tive, stratiform, and a third category dependent on the algorithm. The scheme by Rogers 148 (2010) includes any il as its third category and the Sui et al. (2007) classification contains 149 a mixed category between convective and stratiform. We use the model hydrometeor in-150 formation in our scheme to further separate the categories, which include deep convec-151 tive, shallow convective, congestus, stratiform, anvil, and non-precipitating. In summary, 152 we seek to develop an approach that leverages model microphysical information, captures 153 the bulk convective and stratiform behavior as validated using well-established paradigms 154 of vertical motion, and includes additional classification sub-types to capture their dis-155 tinct radiative forcing signatures. 156

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3.1 Description and Development

Our classification is summarized in Figure 1. The first step of our classification determines if a cell contains cloud. This decision is determined via a total water path (TWP; the sum of rain, cloud, graupel, snow, and ice column-integrated mixing ratios) threshold of 0.1 mm. We found a TWP threshold to be a necessary cutoff to exclude many grid columns identified as containing spurious (i.e., small magnitude) hydrometeor amounts



Figure 1. Flow chart summarizing the categorization process for our precipitation classification algorithm.

associated with negligible rainfall and radiative forcing. We compared this TWP threshold to Rogers (2010) to confirm we were not cutting out a large population of cloud (Figure S5). Adjustments to the TWP threshold primarily effects shallow convective and anvil domain fractions, but does not strongly alter the radiative forcing statistics (not shown).

Next, bulk convective and stratiform categories are separated by a cloud ratio (CR) 167 threshold, as in Sui et al. (2007). The CR is the ratio of ice water path (IWP) to liq-168 uid water path (LWP). IWP is the sum of column-integrated graupel, snow, and cloud 169 ice mixing ratios and LWP is the sum of column-integrated rain and cloud water mix-170 ing ratios. Columns with a CR < 2 are considered convective and columns with CR \geq 171 2 are considered stratiform, assuming stratiform regions will be dominated by ice hydrom-172 eteors. We again compared the CR threshold to the Rogers (2010) scheme and confirmed 173 our threshold falls between its identified convective and stratiform populations (Figure 174 S6). Convective regions are further divided between deep convective, congestus, and shal-175 low, as follows. Grid cells are marked shallow if the column-integrated rain mixing ra-176 tio falls below a threshold of 0.1 mm considering that congestus and deep convective re-177 gions would have higher rain rates (Johnson et al., 1996). Deep convective is separated 178 from congestus by a graupel mixing ratio threshold of 10^{-4} mm, with deep convective 179 regions exceeding this threshold, on the basis that congestus have limited vertical ex-180 tent beyond the 0°C level (Johnson et al., 1999) and hence limited glaciation. Stratiform 181 is separated from anvil where columns exceed a rain mixing ratio of 0.01 mm, account-182 ing for stratiform having more precipitating liquid water content than anvil clouds (Houze, 183 1997; Houze Jr., 2004). While we lack a means for comparing these sub-classifications 184 with the traditional algorithm, we assess their averaged vertical motion profiles against 185 well-established vertical structures in convective and stratiform precipitation (Steiner 186 et al., 1995; Houze Jr., 2004) in the following subsection. 187



Figure 2. Maps comparing a) the new classification scheme, the b) traditional reflectivity classification, c) the LW ACRE, and d) 1000 hPa rain water mixing ratio. See color bars for algorithm classifications. All panels show the first ensemble member of Haiyan at 36 hours.

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3.2 Comparison and Validation

We first present horizontal maps (Figure 2) comparing the new algorithm and the 189 Rogers (2010) classification. The new classification produces the following cloud frac-190 tions for the domain shown (Figure 2): deep convective as 2.76%, congestus as 1.39%, 191 shallow convective as 4.04%, stratiform as 16.42%, anvil as 13.99%, and non-precipitating 192 as 61.4%. The reflectivity approach produces the following cloud fractions: convective 193 is 3.53%, stratiform is 10.09%, anvil is 6.98%, and non-precipitating is 79.4%. The new 194 classification marks more grid cells within the domain as cloud while also increasing the 195 number of points identified as anvil and stratiform compared to the reflectivity approach. 196 The increased count for these cloud types indicates that our algorithm is more sensitive 197 to cloudy columns with lower rain rates, which is an expected and intended result, given 198 our objectives. Additionally, we see more stratiform and anvil regions enveloping the deep 199

convection regions, the latter of which are often located on edges with the reflectivity based algorithm. Our algorithm also identifies shallow convection, which is incorporated
 with the general convective category when using reflectivity.

We next present the vertical motion (w) profiles averaged for each cloud type in 203 our classification alongside the adapted Rogers (2010) scheme in Figure 3a-c. All three 204 convective w-profiles have an expected bottom-heavy profile, with deep convective max-205 imizing around 550 hPa, congestus maximizing at 850 hPa, and shallow convective max-206 imizing at 900 hPa. The profile averaged across all three convective types is consistent 207 with the shape of reflectivity-based classification profile, albeit with a smaller magnitude 208 and broader peak. These differences are consistent with the capture of more weakly pre-209 cipitating columns of convection in the new algorithm. Further, for both stratiform and 210 anvil the new algorithm produces a "top-heavy" profile, with upper-level rising and low-211 level sinking motion, with anvil having a smaller magnitude than stratiform. In compar-212 ison with the reflectivity-based approach, the profiles are similar in shape but with slight 213 differences in magnitudes and with our classification having a 50-100 hPa higher inflec-214 tion point. The overall merit of this new classification scheme is supported by the con-215 sistencies between these w-profiles in the two algorithms and, more broadly, with well-216 established paradigms documented in the literature (Steiner et al., 1995; Houze, 1997; 217 Houze Jr., 2004). 218

We conclude that the new classification algorithm accurately identifies precipita-219 tion types for this model output. The scheme is less computationally bulky than tradi-220 tional reflectivity-based schemes as it is based on microphysics thresholds and is a column-221 by-column approach. This allows for cloud classification within the framework of mod-222 els without the need of neighboring cell information, which is computationally cumber-223 some in highly parallelized frameworks commonly used for convective-scale modeling. 224 However, this approach does have some weaknesses. By having the cloud classification 225 based on microphysics thresholds, the scheme inherently relies on specific hydrometeor 226 behavior and treatment. Different microphysics schemes have vastly different treatments 227 of hydrometeors (Morrison et al., 2020) and would hence require modification to apply 228 to other microphysics schemes. This caveat extends to reflectivity-based approaches, how-229 ever, since model-based reflectivity relies upon microphysical assumptions. Additionally, 230 the new algorithm can only identify one cloud type per column. So, if there are layers 231 of different cloud types present, only the most prominent type will be identified. Despite 232 these limitations, we deem our algorithm suitable for our science question. 233

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4.1 ACRE and CRF

We next seek to quantify each cloud mode's contribution to LW CRF and ACRE. 236 Using our column-by-column based classification algorithm, we calculate the domain-237 averaged and class-averaged LW ACRE. The *domain* average measures each cloud types' 238 contribution to the total LW ACRE in the domain, while the *class* average calculates 239 the mean LW ACRE averaged only over the cells of that type. Stratiform and anvil modes 240 contribute the most to the domain-averaged LW ACRE, with averages around 10 W m^{-2} 241 (Figure 4a). Of the convective types, deep convective has the greatest contribution to 242 the LW ACRE (2.5 W m^{-2}), with congestus and shallow convective points providing 243 the smallest contributions. The large contribution to the domain average by stratiform 244 and any ill modes is partly due to the larger area coverage of these cloud types (Figure 245 2a,b). 246

4 Longwave Radiative Features of Different Cloud Types

But area coverage is not the only reason stratiform and anvil regions have the greatest domain-averaged LW ACRE. When averaging ACRE by class, stratiform and anvil regions retain the highest ACRE value (Figure 4b). They are almost an order of mag-



Figure 3. Averaged profiles of vertical motion (w) (a-c) and CRF (d-f) in convective categories (left column), stratiform and anvil (middle column), and the total convective and stratiform categories (right column). In (a-c), the new classification is presented in solid lines and the reflectivity approach appear as dashed. In (d,e), the dashed black lines represent the averages across all shown categories within the respective panel. Plots include values from all 10 members and 85 timesteps of Haiyan.



Figure 4. Boxplots of domain-averaged (left) and class-averaged (right) LW ACRE (top) and $\hat{h}'NETLW'$ (bottom) by precipitation type. White circles indicate the mean value. Black diamonds represent outliers. Domain averages represent class-averaged values normalized by total grid cell count, and class-averaged values are normalized only by category cell count. Plots include values from all 10 members and 85 timesteps of Haiyan.

²⁵⁰ nitude greater than that of the congestus and shallow convective types. Deep convec-²⁵¹ tive has the greatest LW ACRE value of the three convective modes, with an average ²⁵² value of 85 W m⁻². This value is comparable in magnitude to that of stratiform and ²⁵³ anvil regions, suggesting comparable radiative forcing by these categories within a given ²⁵⁴ column. The combination of stratiform and anvil's large class-averaged LW ACRE with ²⁵⁵ their larger area coverage explains their much larger contribution to the domain-averaged ²⁵⁶ LW ACRE.

We next present vertical profiles of LW CRF to aid interpretation of these results 257 (Figure 3d-f). Cloud types with larger ACRE values exhibit a deep layer of positive CRF. 258 Stratiform and anvil modes are similar in CRF shape and magnitude, which is notewor-259 thy given their much different vertical motion magnitude (Figure 3b). Namely, the ra-260 diative forcing per unit vertical mass flux within a given layer is much larger for anvil 261 than stratiform. CRF in these cloud modes is positive from 200 hPa to the surface and 262 with maxima around 5 K day⁻¹ at 300 hPa, with strong cloud-top cooling above 200 hPa. 263 Of the convective categories, deep convective has the deepest layer of positive CRF of 264 2 K day⁻¹ from 200 hPa to the surface, with a small layer near 0 K day⁻¹ between 400 265 and 500 hPa. This may be due to layers of detrained cloud in association with the $0^{\circ}C$ 266 stable layer (Johnson et al., 1996). Like stratiform and anvil, deep convective has a strong 267 signature of cloud-top cooling above 200 hPa. Congestus and shallow convective modes 268 have maxima in the lower troposphere with cooling above due to their low cloud-top height. 269 Above 800 hPa, their CRF hovers around 0 K day⁻¹. The modest heating values in the 270 upper troposphere in these categories are likely a result of the algorithm only identify-271 ing one cloud type for each column, with these columns potentially including thin anvil 272 clouds. Otherwise, these results are consistent with expectations of cloud depth based 273 on mean vertical motion (Figure 3 a-c). 274

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4.2 LW MSE Variance Source Term

The MSE budget is a tool that allows us to assess and quantify upscale develop-276 ment and intensification of convection. This tool has been used to assess the influence 277 of radiative feedback in relation to both self-aggregation (Wing & Emanuel, 2014) and 278 TC genesis (J. H. Ruppert et al., 2020; Wu et al., 2021). Our interest in this budget is 279 primarily in the LW-source term as it has been shown to be the dominant term for main-280 taining and accelerating convective upscale development (at least, prior to TC genesis) 281 (Wing & Emanuel, 2014; Wing & Cronin, 2016; J. H. Ruppert et al., 2020). The LW MSE 282 variance term (h'NetLW') is the correlation between the anomaly of the density-weighted 283 vertical integral of MSE (h') and the anomalous column LW radiative flux convergence 284 (NetLW'), where anomalies are calculated as the deviation from the domain average. More 285 details on the calculation of this term can be found in Wing and Emanuel (2014) and 286 J. H. Ruppert et al. (2020). 287

When we average h'NetLW' by cloud type, we once again see that stratiform and 288 anvil regions contribute the most to the domain-averaged LW MSE variance source term. 289 Stratiform dominates, with an average of about 0.035 day^{-1} (Figure 4c). The stratiform 290 regions also have the highest class-averaged LW MSE variance of all the cloud types (Fig-291 ure 4d). Congestus and shallow convective points have the lowest class-averaged LW MSE 292 variance, as may be anticipated from LW ACRE (Figure 4 a-b). Surprisingly, deep con-293 vective regions have the second highest LW MSE variance, almost matching that of strat-294 iform. Anvil's class average follows deep convective with a value of 0.25 day^{-1} compared 295 to stratiform and deep connective's averages of about 0.4 day⁻¹. Like the LW ACRE, 296 the stratiform and anvil regions have much greater areal coverage than convective re-297 gions (Figure 2a), which explains the smaller domain-averaged value for convection. 298

The high column influence of stratiform and anvil clouds in terms of both LW ACRE and \hat{h}' NetLW', combined with their extensive areal coverage, indicates their unique importance for supporting upscale convective development in the tropics by amplifying MSE
 variance via radiative forcing. Although deep convective regions also have high values
 of LW MSE variance in a given column, the small regional coverage of this cloud type
 likely limits its area-averaged impact, which is how radiative forcing is linked to the ten dency of MSE variance (Wing & Emanuel, 2014). These findings emphasize the impor tant influences of stratiform and anvil regions on tropical convective organization through
 their radiative forcing, which have not been previously examined in this manner.

5 Summary and Conclusions

In this study, we have investigated the role of CRF of five different cloud types and 309 their ability to aid the organization of tropical convection via the analysis of convection-310 permitting WRF simulations conducted in the context of TC development. To accom-311 plish this, we developed a novel column-by-column cloud classification algorithm based 312 on microphysical thresholds. Our classification holds several advantages over low-level 313 reflectivity-based approaches, such as it is computationally efficient and can be run within 314 the framework of a numerical model, it is sensitive to cloud (including non-precipitating) 315 throughout the column, and it identifies five cloud modes (instead of two or three): shal-316 low, congestus, and deep convective; and stratiform and anvil. However, the disadvan-317 tages to this algorithm includes its likely sensitivity to different microphysics schemes 318 based on its threshold approach and that it can only identify one cloud type per column. 319 Despite these disadvantages, this approach to cloud classification allows for more ques-320 tions to be answered on the influence of cloud type, including our question of how dif-321 ferent cloud types in organized convective systems promote upscale development though 322 LW radiative forcing. 323

We hypothesized that stratiform and anvil regions would support convective or-324 ganization more than other categories through LW ACRE and the LW MSE variance 325 source term. We found that stratiform and anvil contributed the most to the domain-326 averaged ACRE and had greater class-averaged ACRE than that of the other types, in-327 dicating their important contribution to the direct LW radiative forcing. For the LW MSE 328 variance source term, stratiform and anvil again contributed the most to the domain av-329 erage. However, the class averages revealed that deep convective was on par with strat-330 iform regions, resulting in those two classes having the highest class averages. Anvil was 331 third highest, followed by shallow and congestus, which were much weaker. While the 332 class-averaged LW MSE variance source term indicates that the localized forcing by deep 333 convective, stratiform, and anvil clouds is comparable, anvil and stratiform clouds dom-334 inate in supporting convective upscale development owing to their much greater area cov-335 erage. Although we do not fully answer the question of how different cloud types in or-336 ganized convective systems uniquely promote upscale development through LW radia-337 tive forcing, we do provide support of our hypothesis and shed new light on the specific 338 cloud types most important to convective upscale via LW cloud feedback. Future work 339 will focus on the mechanisms CRF works through to promote organization within trop-340 ical convection. 341

³⁴² Open Research Section

The code needed to recreate the WRF simulations described in this study is published at Zenodo (J. Ruppert & Zhang, 2024). The code for the precipitation classification algorithm (Luschen & Ruppert, 2024b) and the analysis (Luschen & Ruppert, 2024a) presented here are available on Zenodo as well.

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³⁵⁵ contributed to this research.

356 References

- Adames, Á. F., & Kim, D. (2016, mar). The MJO as a Dispersive, Convectively Coupled Moisture Wave: Theory and Observations. *Journal of the Atmospheric Sciences*, 73(3), 913–941. doi: 10.1175/JAS-D-15-0170.1
- Ahmed, F., & Schumacher, C. (2015). Convective and stratiform components of the precipitation-moisture relationship. *Geophysical Research Letters*, 42(23), 10,453–10,462. doi: 10.1002/2015GL066957
- Biggerstaff, M. I., & Listemaa, S. A. (2000, December). An Improved Scheme for
 Convective/Stratiform Echo Classification Using Radar Reflectivity. Journal
 of Applied Meteorology and Climatology, 39(12), 2129–2150. doi: 10.1175/1520
 -0450(2001)040(2129:AISFCS)2.0.CO;2
- Bony, S., Stevens, B., Frierson, D. M. W., Jakob, C., Kageyama, M., Pincus, R., ...
 Webb, M. J. (2015, April). Clouds, circulation and climate sensitivity. *Nature Geoscience*, 8(4), 261–268. doi: 10.1038/ngeo2398
- 370Bretherton, C. S., Blossey, P. N., & Khairoutdinov, M. (2005, December).An371Energy-Balance Analysis of Deep Convective Self-Aggregation above Uni-372form SST. Journal of the Atmospheric Sciences, 62(12), 4273–4292.37310.1175/JAS3614.1
- Bu, Y. P., Fovell, R. G., & Corbosiero, K. L. (2014, May). Influence of
 Cloud–Radiative Forcing on Tropical Cyclone Structure. Journal of the At mospheric Sciences, 71(5), 1644–1662. doi: 10.1175/JAS-D-13-0265.1
- Ciesielski, P. E., Johnson, R. H., Jiang, X., Zhang, Y., & Xie, S. (2017). Relationships between radiation, clouds, and convection during DYNAMO. Journal of Geophysical Research: Atmospheres, 122(5), 2529–2548. doi: 10.1002/2016JD025965
- Davis, C. A. (2015, September). The Formation of Moist Vortices and Tropical
 Cyclones in Idealized Simulations. *Journal of the Atmospheric Sciences*, 72(9),
 3499–3516. doi: 10.1175/JAS-D-15-0027.1
- Houze, R. A. (1997, oct). Stratiform Precipitation in Regions of Convection: A Meteorological Paradox? Bulletin of the American Meteorological Society, 78(10), 2179–2196. doi: 10.1175/1520-0477(1997)078(2179:SPIROC)2.0.CO;2
- Houze Jr., R. A. (2004). Mesoscale convective systems. Reviews of Geophysics, 42(4). doi: 10.1029/2004RG000150
- Iacono, M. J., Delamere, J. S., Mlawer, E. J., Shephard, M. W., Clough, S. A., &
 Collins, W. D. (2008). Radiative forcing by long-lived greenhouse gases: Cal culations with the AER radiative transfer models. Journal of Geophysical
 Research: Atmospheres, 113(D13). doi: 10.1029/2008JD009944
- Johnson, R. H., Ciesielski, P. E., & Hart, K. A. (1996, jul). Tropical Inversions near
 the 0°C Level. Journal of the Atmospheric Sciences, 53(13), 1838–1855. doi:
 10.1175/1520-0469(1996)053(1838:TINTL)2.0.CO;2
- Johnson, R. H., Rickenbach, T. M., Rutledge, S. A., Ciesielski, P. E., & Schubert,
 W. H. (1999, aug). Trimodal Characteristics of Tropical Convection. Journal of Climate, 12(8), 2397–2418. doi: 10.1175/1520-0442(1999)012(2397: TCOTC)2.0.CO;2

400	Luschen, E., & Ruppert, J. (2024a, 2). Analysis Code for Geophysical Research Let-
401	ter paper. doi: 10.5281 /zenodo. 10645132
402	Luschen, E., & Ruppert, J. (2024b, 2). Column-based Precipitation Classification Al-
403	gorithm. doi: $10.5281/zenodo.10611873$
404	Manabe, S., & Strickler, R. F. (1964, July). Thermal Equilibrium of the Atmosphere
405	with a Convective Adjustment. Journal of the Atmospheric Sciences, $21(4)$,
406	361–385. doi: $10.1175/1520-0469(1964)021(0361:TEOTAW)2.0.CO;2$
407	Morrison, H., van Lier-Walqui, M., Fridlind, A. M., Grabowski, W. W., Harrington,
408	J. Y., Hoose, C., Xue, L. (2020). Confronting the Challenge of Modeling
409	Cloud and Precipitation Microphysics. Journal of Advances in Modeling Earth
410	Systems, 12(8), e2019MS001689. doi: 10.1029/2019MS001689
411	Muller, C. J., & Held, I. M. (2012, August). Detailed Investigation of the Self-
412	Aggregation of Convection in Cloud-Resolving Simulations. Journal of the At-
413	mospheric Sciences, 09(8), 2551-2505. doi: 10.1175/JAS-D-11-0257.1
414	Dadiating Enging During DVNAMO (AMIE – Journal of Coophysical Bassarch
415	Aduative Forcing During D1NAMO/AMIE. Journal of Geophysical Research:
416	Nachham M. P. & Pandall D. A. (2021a) Linking Atmospheric Cloud Padia
417	tive Effects and Tropical Precipitation Coophysical Research Letters 18(14)
410	e2021GL094004 doi: 10.1029/2021GL094004
420	Needham M B & Bandall D A (2021b) Riehl and Malkus Revisited: The Role
421	of Cloud Radiative Effects. Journal of Geophysical Research: Atmospheres.
422	<i>126</i> (16), e2021JD035019. doi: 10.1029/2021JD035019
423	NOAA-NCEP Global Ensemble Forecast System. (2015). NCEP GFS 0.25 De-
424	gree Global Forecast Grids Historical Archive. UCAR/NCAR - Research Data
425	Archive. doi: 10.5065/D65D8PWK
426	Powell, S. W., Houze, R. A., & Brodzik, S. R. (2016, March). Rainfall-Type Catego-
427	rization of Radar Echoes Using Polar Coordinate Reflectivity Data. Journal of
428	Atmospheric and Oceanic Technology, 33(3), 523–538. doi: 10.1175/JTECH-D
429	-15-0135.1
430	Rogers, R. (2010, January). Convective-Scale Structure and Evolution during a
431	High-Resolution Simulation of Tropical Cyclone Rapid Intensification. Journal
432	of the Atmospheric Sciences, 67(1), 44–70. doi: 10.1175/2009JAS3122.1
433	Ruppert, J., & Zhang, Y. (2024, 2). Ensemble WRF Simulations of Typhoon Haiyan
434	and Hurricane Maria. doi: 10.5281/zenodo.10572959
435	Ruppert, J. H., Wing, A. A., Tang, X., & Duran, E. L. (2020, November). The crit-
436	Ical role of cloud-infrared radiation feedback in tropical cyclone development. $P_{\text{reaccodings}}$ of the National Academy of Sciences $117/(45)$ 27884 27802 doi:
437	$10\ 1073\ /\text{pms}\ 9013584117$
438	Schumacher C k Houze R A (2003 jun) Stratiform Bain in the Tron-
439	ics as Seen by the TRMM Precipitation Badar* Journal of Climate
441	16(11), $1739-1756$. Retrieved from http://journals.ametsoc.org/
442	doi/10.1175/1520-0442(2003)016%3C1739:SRITTA%3E2.0.CO;2 doi:
443	10.1175/1520-0442(2003)016(1739:SRITTA)2.0.CO;2
444	Skamarock, W. C., Klemp, J. B., Dudhia, J. B., Gill, D. O., Barker, D. M., Duda,
445	M. G., Powers, J. G. (2021). A Description of the Advanced Re-
446	search WRF Model Version 4.3. NCAR Technical Note(July), 1–165. doi:
447	10.5065/1dfh-6p97
448	Steiner, M., Houze, R. A., & Yuter, S. E. (1995, September). Climatological Char-
449	acterization of Three-Dimensional Storm Structure from Operational Radar
450	and Rain Gauge Data. Journal of Applied Meteorology and Climatology, $34(9)$,
451	1978–2007. doi: $10.1175/1520-0450(1995)034(1978:CCOTDS)2.0.CO;2$
452	Sui, CH., Tsay, CT., & Li, X. (2007). Convective–stratiform rainfall separation by
453	cloud content. Journal of Geophysical Research: Atmospheres, 112(D14). doi:
454	10.1029/2006JD008082

455	Thompson, G., & Eidhammer, T. (2014, October). A Study of Aerosol Impacts on
456	Clouds and Precipitation Development in a Large Winter Cyclone. Journal of
457	the Atmospheric Sciences, 71(10), 3636–3658. doi: 10.1175/JAS-D-13-0305.1
458	Webster, P. J., & Stephens, G. L. (1980, jul). Tropical Upper-Tropospheric Ex-
459	tended Clouds: Inferences from Winter MONEX. Journal of the Atmospheric
460	Sciences, 37(7), 1521-1541. Retrieved from http://journals.ametsoc.org/
461	doi/abs/10.1175/1520-0469-37.7.1521 doi: 10.1175/1520-0469-37.7.1521
462	Wing, A. A., Camargo, S. J., & Sobel, A. H. (2016, July). Role of Radia-
463	tive–Convective Feedbacks in Spontaneous Tropical Cyclogenesis in Idealized
464	Numerical Simulations. Journal of the Atmospheric Sciences, 73(7), 2633–
465	2642. doi: 10.1175/JAS-D-15-0380.1
466	Wing, A. A., & Cronin, T. W. (2016). Self-aggregation of convection in long channel
467	geometry. Quarterly Journal of the Royal Meteorological Society, 142(694), 1-
	15 doi: 10.1002/ai.2628
468	$15. \ dol. \ 10.1002/\text{Q}.2020$
468 469	Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-
468 469 470	Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self- Aggregation in Numerical Simulations: A Review. Surveys in Geophysics,
468 469 470 471	Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self- Aggregation in Numerical Simulations: A Review. <i>Surveys in Geophysics</i> , 38(6). doi: 10.1007/s10712-017-9408-4
468 469 470 471 472	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling
468 469 470 471 472 473	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simula-
468 469 470 471 472 473 474	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi:
468 469 470 471 472 473 474 475	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269
468 469 470 471 472 473 474 475 476	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269 Wu, SN., Soden, B. J., & Nolan, D. S. (2021). Examining the Role of Cloud
468 469 470 471 472 473 474 475 476 477	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269 Wu, SN., Soden, B. J., & Nolan, D. S. (2021). Examining the Role of Cloud Radiative Interactions in Tropical Cyclone Development Using Satellite Mea-
468 469 470 471 472 473 474 475 475 477 478	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269 Wu, SN., Soden, B. J., & Nolan, D. S. (2021). Examining the Role of Cloud Radiative Interactions in Tropical Cyclone Development Using Satellite Measurements and WRF Simulations. Geophysical Research Letters, 48(15),
468 469 470 471 472 473 474 475 476 477 478 479	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269 Wu, SN., Soden, B. J., & Nolan, D. S. (2021). Examining the Role of Cloud Radiative Interactions in Tropical Cyclone Development Using Satellite Measurements and WRF Simulations. Geophysical Research Letters, 48(15), e2021GL093259. doi: 10.1029/2021GL093259
468 469 470 471 472 473 474 475 476 477 478 479 480	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269 Wu, SN., Soden, B. J., & Nolan, D. S. (2021). Examining the Role of Cloud Radiative Interactions in Tropical Cyclone Development Using Satellite Measurements and WRF Simulations. Geophysical Research Letters, 48(15), e2021GL093259. doi: 10.1029/2021GL093259 Zelinka, M. D., Myers, T. A., McCoy, D. T., Po-Chedley, S., Caldwell, P. M.,
468 469 470 471 472 473 474 475 476 477 478 478 478 478 480 481	 Wing, A. A., Emanuel, K., Holloway, C. E., & Muller, C. (2017). Convective Self-Aggregation in Numerical Simulations: A Review. Surveys in Geophysics, 38(6). doi: 10.1007/s10712-017-9408-4 Wing, A. A., & Emanuel, K. A. (2014). Physical mechanisms controlling self-aggregation of convection in idealized numerical modeling simulations. Journal of Advances in Modeling Earth Systems, 6(1), 59–74. doi: 10.1002/2013MS000269 Wu, SN., Soden, B. J., & Nolan, D. S. (2021). Examining the Role of Cloud Radiative Interactions in Tropical Cyclone Development Using Satellite Measurements and WRF Simulations. Geophysical Research Letters, 48(15), e2021GL093259. doi: 10.1029/2021GL093259 Zelinka, M. D., Myers, T. A., McCoy, D. T., Po-Chedley, S., Caldwell, P. M., Ceppi, P., Taylor, K. E. (2020). Causes of higher climate sensitivity in

483 https://doi.org/10.1029/2019GL085782

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Supporting Information for

Longwave Radiative Feedback Due to Stratiform and Anvil Clouds

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Introduction

The supporting figures included are to primarily show that Maria produces the same qualitative results as Haiyan shown within the main paper. We also include analysis for a 24-hour period (12-36 hours) at the beginning of the Haiyan simulation to demonstrate that our results are generalizable beyond the post-TC-genesis. Finally, we present histograms of the Rogers (2010) classification to support our development of the new classification scheme.



Figure S1. Same as Figure 3 but for the first 24 hours after "spin-up."



Figure S2. Same as Figure 4 but for the first 24 hours after "spin-up."



Figure S3. Same as Figure 3 but for Hurricane Maria.



Figure S4. Same as Figure 4 but for Hurricane Maria.



Figure S5. Total water path (TWP) frequency for the Rogers (2010) classification. The convective mode is in blue, the stratiform mode is in red, and the anvil mode is in black. The vertical black line indicates the TWP threshold for the new classification. Figure shows the first member of Haiyan at the timestep 36 hours.



Figure S6. Same as Figure S5 but for cloud ratio (CR).