# Upper-mantle anisotropy in the southeastern margin of Tibetan Plateau revealed by fullwave SKS splitting intensity tomography

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#### Abstract

The southeastern margin of the Tibetan Plateau has undergone complex deformation since the Cenozoic, resulting in a high level of seismicity and seismic hazard. Knowledge about the seismic anisotropy provides important insight about the deformation mechanism and the regional seismotectonics beneath this tectonically active region. In this study, we conduct fullwave multi-scale tomography to investigate the seismic anisotropy in the southeastern margin of the Tibetan Plateau. Broadband records at 111 permanent stations in the region from 470 teleseismic events are used to obtain 5,216 high-quality SKS splitting intensity measurements, which are then inverted in conjunction with 3D sensitivity kernels to obtain an anisotropic model with multi-scale resolution. Resolution tests show that our dataset recovers anisotropy anomalies reasonably well on the scale of  $1^{\circ} \times 1^{\circ}$  horizontally and ~100 km vertically. Our result suggests that in the southeastern margin of the Tibetan Plateau the deformation in the lithosphere and asthenosphere are decoupled. The anisotropy in the lithosphere varies both laterally and vertically as a result of dynamic interactions of neighboring blocks as well as lithospheric reactivation. The anisotropy in the asthenosphere largely follows the direction of regional absolute plate motion. The SKS splittings observed at the surface are shown to be consistent with the vertical integral of our depth-dependent anisotropy model over lithospheric and asthenospheric depths.

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15	Key Points:
16	• A 3D shear-wave anisotropy model for the SE margin of the Tibetan Plateau is obtained by
17	fullwave SKS splitting intensity tomography
18	• Result shows a decoupling of lithosphere and asthenosphere deformations in the region, but a
19	coupling under Tibetan Plateau
20	• Lithospheric anisotropy has a complex pattern, whereas asthenospheric anisotropy follows the
21	APM

#### 23 Abstract

The southeastern margin of the Tibetan Plateau has undergone complex deformation since the Cenozoic, 24 25 resulting in a high level of seismicity and seismic hazard. Knowledge about the seismic anisotropy provides important insight about the deformation mechanism and the regional seismotectonics beneath this tectonically 26 active region. In this study, we conduct fullwave multi-scale tomography to investigate the seismic 27 anisotropy in the southeastern margin of the Tibetan Plateau. Broadband records at 111 permanent stations 28 in the region from 470 teleseismic events are used to obtain 5,216 high-quality SKS splitting intensity 29 measurements, which are then inverted in conjunction with 3D sensitivity kernels to obtain an anisotropic 30 model with multi-scale resolution. Resolution tests show that our dataset recovers anisotropy anomalies 31 reasonably well on the scale of 1° x 1° horizontally and ~100 km vertically. Our result suggests that in the 32 southeastern margin of the Tibetan Plateau the deformation in the lithosphere and asthenosphere are decoupled. 33 The anisotropy in the lithosphere varies both laterally and vertically as a result of dynamic interactions of 34 neighboring blocks as well as lithospheric reactivation. The anisotropy in the asthenosphere largely follows the 35 direction of regional absolute plate motion. The SKS splittings observed at the surface are shown to be 36 consistent with the vertical integral of our depth-dependent anisotropy model over lithospheric and 37 asthenospheric depths. 38

Keywords: seismic anisotropy; splitting intensity; finite-frequency; fullwave tomography; southeastern
Tibetan Plateau

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#### 42 Plain Language Summary

The southeastern margin of the Tibetan Plateau has undergone significant deformation since the Cenozoic due to the collision between the Eurasian and Indian plates in the south and interactions with the Yangtze Craton in the east. Knowledge about the upper mantle seismic anisotropy helps us understand the

deformation and dynamic evolution. In this study, we conduct fullwave multi-scale anisotropy 46 tomography for the southeastern margin of the Tibetan Plateau using 5,216 high-quality SKS splitting 47 intensity measurements obtained from the broadband records of 470 teleseismic events at 111 stations. Our 48 result shows a decoupling between the lithosphere and asthenosphere deformations in the southeastern margin 49 of the Tibetan Plateau. The anisotropy in the lithosphere varies both laterally and vertically as a result of the 50 dynamic interactions of neighboring blocks as well as lithospheric reactivation. The anisotropy in the 51 asthenosphere is largely parallel to the regional absolute plate motion, and the SKS splitting observed at the 52 surface is the result of vertical integration of the anisotropy effect through the lithosphere and asthenosphere. 53

#### 55 **1 Introduction**

The ongoing Indian-Eurasian continental collision has resulted in the largest plateau on Earth and caused widespread deformation in central and east Asia (Yin & Harrison, 2000; Kind et al., 2002). Despite numerous efforts, questions remain about the dynamics of the lithospheric deformation and asthenosphere flow beneath the Tibetan Plateau and the surrounding regions (Royden et al., 2008).

The southeastern (SE) margin of the Tibetan Plateau involves many active tectonic blocks (Figure 1), 60 including the Songpan-Ganzi Terrane (SGT), the Sichuan Basin (SCB), the Sichuan-Yunnan Rhombic 61 Block (SYRB), the Indo-China Block (ICB), the Qiangtang Block (QTB), and the Yangtze Craton (YTC). 62 The SGT is part of central Tibetan Plateau. Its eastern part is separated from the SCB and SYRB by the 63 Longmenshan Fault (LMSF) and Lijiang-Xiaojinhe Fault (LJ-XJHF), respectively, and bounded in the south 64 by the Jinshajiang Fault (JSJF) from the QTB. The convergence between the Indian and Eurasian plates 65 caused the SGT to expand eastward against the SCB during the Cenozoic (Yin & Harrison, 2000). GPS 66 observations show eastward crustal motion of the eastern SGT with the crustal strain rate decreasing 67 abruptly from ~20 mm/year (relative to the YTC reference frame) in the interior SGT to ~3-4 mm/year or 68 less in the vicinity of central and southern segments of the LMSF, indicating that the eastward expansion 69 of the SGT is apparently resisted by the SCB (Shen et al., 2005; Zhang, 2013) In addition, there is also an 70 71 apparent clockwise rotation of the GPS velocities around the Eastern Himalaya Syntax (EHS). Lowvelocity and high-conductivity anomalies in the mid-lower crust under the SGT revealed by geophysical 72 73 studies (Zhao et al., 2012; Bao et al., 2020) suggest the existence of mid-lower crustal flow. However, 74 these geophysical anomalies show strong lateral heterogeneity in eastern Tibet, implying a complex process of deformation in the region. The SCB and SYRB are both parts of the YTC (Zhang et al., 2013; 75 76 Li et al., 2021). The former forms the rigid and stable northwestern margin of the YTC, while the crust of 77 the latter is extruding southeastward along the Anninghe-Zemuhe Fault (ANH-ZMHF) and Xiaojiang Fault (XJF) in the east and the Red River Fault (RRF) in the southwest (Zhang et al., 2003). The Lancangjiang
Fault (LCJF) separates the narrow QTB in the east and the ICB in the west. Crustal movements are
predominantly characterized by a clockwise rotation around the Eastern Himalayan Syntaxis (EHS),
transforming the movement of the plateau material from eastward north of the syntaxis to southeastward
and southward further south (Wang & Shen, 2020).





Figure 1. Map of the tectonic environment of SE margin of Tibetan Plateau with seismic stations (greenfilled triangles) and epicenters (red open circles) of earthquakes of magnitude 5 and above from 2000 to

2022. The blue triangle marks the station WXT for which SKS waveforms, splitting intensities and 86 sensitivity kernels are shown in Figures 4, 5 and 7, respectively. Major active faults are shown by thick 87 black lines with abbreviated names in black, including LMSF: Longmenshan Fault; XSHF: Xianshuihe 88 Fault; JSJF: Jinshajiang Fault; LCJF: Lancangjiang Fault; LJ-XJHF: Lijiang-Xiaojinhe Fault; ANH-89 ZMHF: Anninghe-Zemuhe Fault; XJF: Xiaojiang Fault; and RRF: Red River Fault. Major active tectonic 90 blocks are indicated by abbreviated texts in blue, including SGT: Songpan-Ganzi Terrane; SCB: Sichuan 91 Basin; QTB: Qiangtang Block; SYRB: Sichuan-Yunnan Rhombic Block; YTC: Yangtze Craton; and ICB: 92 Indo-China Block. EHS stands for the Eastern Himalaya Syntax. Blue cross arrows show the surface strain 93 rate derived from GPS observation (Wang & Shen, 2020). Background color shows the topography. The 94 black box in the inset map indicates the location of the main figure. 95

Over the past two decades, a large number of seismic stations have been deployed in the SE margin 96 of the Tibetan Plateau. Waveforms recorded by the growing number of broadband stations have provided 97 crucial data for studying the structure and dynamics of the crust and upper mantle beneath the region, such 98 as variation in crustal thickness (Wang et al., 2017; Xu et al., 2020), widespread low-velocity anomalies in 99 mid-lower crust revealed by receiver function analysis (Hu et al., 2005; Xu et al., 2007; Zhang et al., 2009; 100 Wang et al., 2010), Lg-wave high-attenuation zones (Zhao et al., 2013; Wei & Zhao, 2019), joint inversion 101 of receiver function and surface wave dispersion (Liu et al., 2014), and body- and surface-wave 102 tomographies (Huang et al., 2002; Wang et al., 2003; Huang et al., 2009; Li et al., 2009; Yang et al., 2019; 103 Wei & Zhao, 2022). 104

Across the LMSF, the drastic change in elevation from ~5–6 km in the west to a few hundred meters in the east suggests significant variations in the lithospheric thickness, as illustrated in Figure 2, where the global model LITHO1.0 (Pasyanos et al., 2014) is shown north of 27°N. South of 27°N, LITHO1.0 is rather uniform, and the regional model of Yang et al. (2017) is depicted. In this region, the lithosphere has a thickness of less than 100 km in the northern and southern parts but ~150 km at the mid latitudes. The thickest lithosphere is more than 200 km beneath the SCB.



Figure 2. Lithosphere thickness map. The global model LITHO1.0 (Pasyanos et al., 2014) is shown for the region north of 27°N, whereas to the south the regional model of Yang et al. (2017) is depicted. Names of major faults and tectonic blocks are the same as in Figure 1.

115 Seismic anisotropy in the upper mantle is an important proxy for deformation. There have been 116 numerous studies devoted to crustal anisotropy beneath the SE Tibetan Plateau utilizing different methods, 117 such as Pms splitting (Sun et al., 2012; Cai et al., 2016; Han et al., 2020), anisotropic tomography of P 118 and Pn waves (Lei et al., 2014; Huang et al., 2018), surface wave anisotropic tomography (Yang et al.,

2010; Yao et al., 2010; Legendre et al., 2015; Zhang et al., 2023), and the splitting of shear waves (Shi et al., 119 2012). In the upper crust, the fast axis directions are mainly parallel to the strike of active faults (e.g., Yao 120 121 et al., 2010; Shi et al., 2012; Huang et al., 2018), whereas the anisotropic pattern in the lower crust is different. Huang et al. (2018) showed that the fast velocity direction deviates from the strikes of active faults 122 significantly using P-wave anisotropic tomography. Han et al. (2020) used Markov-chain Monte Carlo 123 124 inversion of receiver functions to isolate the effect of potential dipping interfaces. Their results showed that the fast axis directions in the lower crust are in good agreement with topography, implying that the 125 gravitational potential may be the driving force for the crustal deformation in southeastern Tibet. 126

XKS-wave splitting is routinely used to investigate the anisotropic structure in the upper mantle (Long 127 & Becker, 2010). Flesch et al. (2005) jointly analyzed GPS, surface geology and shear-wave splitting 128 measurements and argued for vertically coherent deformation in the crust and upper mantle in the Tibetan 129 Plateau but a decoupling in Yunnan Province in southwestern China. Lev et al. (2006) also suggested a 130 lithosphere-asthenosphere decoupling in Yunnan using shear-wave splitting observations, but they were not 131 132 able to determine the degree of coupling beneath the Tibetan Plateau. Based on a joint analysis of a larger dataset of shear-wave splitting measurements and GPS observations, Wang et al. (2008) argued for crust-133 mantle coupling in the Tibetan Plateau and the surrounding regions. 134

Substantial efforts have been made to develop a theoretical framework as well as practical strategies for inversion for the 3D distribution of anisotropy. A fullwave approach has been developed for the measurement of shear-wave splitting intensities and interpretation in terms of shear-wave azimuthal anisotropy parameters (Chevrot, 2000; Favier & Chevrot, 2003; Chevrot, 2006; Sieminski et al., 2008; Monteiller & Chevrot, 2011; Lin et al., 2014a), which has been applied to anisotropy tomographies for southern California (Monteiller & Chevrot, 2011; Lin et al., 2014b), the High Lava Plain (Mondal & Long, 2020), and the southeastern Tibetan Plateau (Huang & Chevrot, 2021). The depth variations of anisotropy obtained by these studies have shed new lights in understanding the sources of anisotropy and the associated mantledynamics.

In this study, we conduct a fullwave multiscale anisotropy tomography for the SE margin of Tibetan Plateau. We collect seismic records at regional permanent broadband stations from globally distributed earthquakes and obtain high-quality measurements of SKS splitting intensities. We then invert the splitting intensities using a wavelet-based parameterization of the 3D model to achieve a multi-scale resolution of the anisotropic structure. We also provide an interpretation of our anisotropic model for the SE margin of Tibetan Plateau in terms of regional geodynamics.

150 **2 Data and Methods** 

#### 151 **2.1 Waveform records**

We collect waveforms recorded by 111 permanent broadband stations (green triangles in Figure 1) deployed in the study region. To guarantee a wide range of azimuthal distribution, we select events of magnitude  $Mw \ge 5.5$  from 2009 to 2020, located in the epicentral distance range of 90°–130°. After quality control of the waveforms and removal of outliers of the data (see Section 2.2), a total of 470 events are used in the subsequent inversions. Figure 3 displays the event distribution.

#### 157 **2.2 Splitting intensity measurements**

Shear-wave splitting measurement is nowadays a routine procedure in the study of seismic anisotropy. Several previous studies have documented the measured SKS splitting parameters (fast directions and delay times) at stations in our study region (e.g., Chang et al., 2015; Yang et al., 2018; Liu et al., 2020; Huang & Chevrot, 2021; Li et al., 2021). In this study, we invert for the 3D anisotropy structure using the splitting intensity (SI) measurements obtained by computing the zero-lag cross-correlation between the transverse-component record and the time derivative of the radial-component record (Chevrot, 2000).



Figure 3. Distribution of 470 teleseismic events used for SKS splitting intensity inversion in this study.
Events of magnitudes Mw5.5 and greater in the epicentral distance range of 90°–130° during 2009–2020
are selected. The red box in the center indicates the study area.

### 168 For a given station, the SI of the SKS wave from the *i*-th event is defined as

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$$S_{i} = -2 \frac{\int_{t_{i1}}^{t_{i2}} \dot{u}_{i}^{R}(t) u_{i}^{T}(t) dt}{\int_{t_{i1}}^{t_{i2}} [\dot{u}_{i}^{R}(t)]^{2} dt},$$
 (1)

where  $[t_{i1}, t_{i2}]$  is the time window for the SKS wave, and  $u_i^R(t)$  and  $u_i^T(t)$  are the radial and transversecomponent records, respectively, from the *i*-th event. A dot above a variable indicates derivative with respect to time. The conventional SKS splitting parameters at the given station, namely the fast-direction  $S_i = \Delta t \sin 2(\theta - \theta_i^b),$ 

(2)

azimuth  $\theta$  and delay time  $\Delta t$ , are related to the SIs measured at the station from all events through a sinusoidal curve fitting (Chevrot, 2000; Lin et al., 2014a):

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176 where 
$$\theta_i^{b}$$
 is the back azimuth of the *i*-th event.

In this study, we obtain the SI measurements of SKS waves with the help of SplitRacer (Link et al., 177 178 2022), an efficient and automatic toolbox developed for the measurement and quality control of XKS splittings. An example of SplitRacer processing is shown in Figure 4. We use SplitRacer to determine the 179 SKS time window automatically, followed by a manual check on the quality of the SKS signals. Then, we 180 calculate the SIs using Eq. (1). The period band we use in this study is 8-50 s considering the dominant 181 periods of the SKS signals as well as minimizing the interference with neighboring phases. As shown in 182 Figure 4, the spectral powers of the radial- and transverse-component records are calculated by the short-183 time Fourier transform (Quatieri, 2006) and summed. Then, the dominant frequency band of the SKS 184 waveform can be identified (between the red dashed horizontal lines in Figure 4c). At each time, the 185 powers within the dominant frequency band are summed (Figure 4d), which defines the SKS window by 186 the two crossing points at 50% of the peak level (vertical green lines in Figure 4d). The final SKS window 187 is given by either expanding or shrinking the 50% energy window to a fixed 30-s window (red vertical 188 lines). After completing the quality check using SplitRacer, we obtain a total of 12,457 SKS wave SI 189 measurements. 190

Following Chevrot (2000), we estimate the uncertainty of each SI measurement using the following equation

$$\sigma_{i} = \sqrt{\frac{1}{N_{i}} \left\{ \sum_{j=1}^{N_{i}} \left[ u_{i}^{T}(t_{j}) \right]^{2} - \frac{S_{i}^{2}}{4} \sum_{j=1}^{N_{i}} \left[ \dot{u}_{i}^{R}(t_{j}) \right]^{2} \right\}},$$
(3)

where  $S_i$  is the *i*-th SI measurement, and  $N_i$  is the number of time samples used for the window to obtain the measurement.



Figure 4. An example of SKS window selection using SplitRacer. (a) and (b) show normalized radial and 197 transverse records, respectively, around the SKS arrival at station WXT (blue triangle in Figure 1) from 198 the 24 June 2019 earthquake in New Zealand with a back azimuth of 122°. The green lines mark the 199 theoretical SKS and SKKS arrival times corresponding to the event of interest. Red vertical lines mark 200 the start and end times of the final window used for SI measurement. (c) Summed spectral power of the 201 spectrogram of the radial and transverse records in (a) and (b). The red horizontal line marks the frequency 202 203 of maximum spectral power, and the red dashed lines show the frequency bounds of more than 80% of the maximum. The green vertical lines mark the time window in which the summed spectral power is 204 205 more than 50% of the maximum. (d) Summed spectral power over all frequencies. The green vertical lines mark the time window in which the summed spectral energy is more than 50% of the maximum. (e) 206 STA/LTA ratio of the radial-component record used as quality check. The blue vertical line denotes the 207 time when the STA/LTA ratio reaches its peak within the window determined in (d). The red horizontal 208 209 line is the acceptance threshold of 2.1 (Link et al., 2022).

210 We further clean our dataset of outliers based on two criteria: (1) At a given station, SI measurements

with uncertainties  $\sigma_i$  larger than 1.5 times the average uncertainty  $\bar{\sigma}$  for that station, i.e.  $\sigma_i > 1.5\bar{\sigma}$ , are removed; (2) all SI measurements at a given station are fit by a sinusoidal curve, and measurements that deviate from the sinusoidal curves by more than  $2\sigma_i$  are also removed. After removing the outliers, we retain a total of 5,216 high-quality SI measurements as our final dataset for subsequent anisotropy inversion. The standard deviation of the final dataset (the average of uncertainties of all retained data) is 0.074 s. Figure 5 shows the effect on the distribution of SI measurements before and after applying the above two criteria for station WXT, which reduces the number of SI data from 184 to 61.



Figure 5. Variations of measured SI with event back azimuth for station WXT (blue triangle in Figure 1) before (a) and after (b) removal of SI outliers. The vertical error bars show two standard deviations  $(2\sigma_i)$ of individual measurements. The dashed lines represent the sinusoidal curves that best fit the measurements. According to Eq. (2), the conventional SKS splitting parameters for this station are:  $\Delta t =$ 1.08 s and  $\theta = 109.14^{\circ}$ .

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Based on Eq. (2), we can obtain the conventional fast-direction azimuth  $\theta$  and delay time  $\Delta t$  at each station using all the SI measurements, and they are displayed in Figure 6. The spatial distribution of  $\theta$  has a similar pattern as seen in previous studies, i.e. a generally NW-SE oriented fast axis and a nearly uniform E-W alignment of the fast-axis directions south of ~26°N latitude.



Figure 6. Black line segments show the conventional SKS fast-direction azimuths and delay times at all station obtained from the SI measurements according to Eq. (2). Red line segments are fast-direction azimuths and delay times derived from Eq. (2) based on the model-predicted SIs obtained by integrating the sensitivity kernels (see Figure 7) with the anisotropy model in Figure 12 according to Eq. (4). Names of major faults and tectonic blocks are the same as in Figure 1.

#### 234 **2.3 Multiscale inversion for 3D anisotropic structure**

We implement the full-wave multiscale anisotropy tomography framework developed by Lin et al.

236 (2014b). Here we briefly describe our methodology. Interested readers may consult Lin et al. (2014b) for

a full description of the method.

Becker et al. (2006) pointed out that the upper-mantle anisotropy is mainly hexagonal with strongly

correlated parameters  $\varepsilon$  and  $\gamma$  quantifying the P- and S-wave anisotropy, respectively, and the ellipticity 239 parameter  $\delta$ . Furthermore, Zhao & Chevrot (2011) demonstrated that the SI of an SKS wave is mostly 240 sensitive to  $\gamma$ , and is about 10 times larger than to the isotropic heterogeneity. This allows us to simplify 241 242 the modeling of upper mantle anisotropy in two aspects: (1) it is sufficient to use a 1D reference model; and (2) the anisotropy can be described by only two parameters: the strength of anisotropy  $\gamma$  and the azimuth 243  $\phi_f$  of the symmetry axis of anisotropy. We assume that the symmetry axis is horizontal since the sub-244 vertically propagating SKS waves are insensitive to the dip angle of the symmetry axis of anisotropy 245 (Mondal & Long, 2019). Note that the anisotropy strength  $\gamma$  and azimuth  $\phi_f$  are spatially varying, and 246 the dependence of the SI on  $\phi_f$  is nonlinear. As documented in Favier & Chevrot (2003), we can 247 introduce two independent parameters,  $\gamma_c = \gamma \cos(2\phi_f)$  and  $\gamma_s = \gamma \sin(2\phi_f)$ , to enable a linear 248 relationship with the SI measurement 249

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$$S = \iiint \left[ K_{\gamma_c}^S(\mathbf{r})\gamma_c(\mathbf{r}) + K_{\gamma_s}^S(\mathbf{r})\gamma_s(\mathbf{r}) \right] d\mathbf{r}, \tag{4}$$

where  $K_{\gamma_c}^S$  and  $K_{\gamma_s}^S$  are the Fréchet sensitivity kernels of the splitting intensity S to  $\gamma_c$  and  $\gamma_s$ , respectively. 251 After obtaining  $\gamma_c$  and  $\gamma_s$  from a linear inversion, we can obtain the more familiar anisotropy parameters 252 by utilizing the relations:  $\phi_f(\mathbf{r}) = 0.5 \tan^{-1}[\gamma_s(\mathbf{r})/\gamma_c(\mathbf{r})]$  and  $\gamma(\mathbf{r}) = \sqrt{[\gamma_c(\mathbf{r})]^2 + [\gamma_s(\mathbf{r})]^2}$ . In this 253 study, we use PREM (Dziewonski & Anderson, 1981) as the reference model and compute the sensitivity 254 kernels using the normal-mode summation algorithm developed by Zhao & Chevrot (2011). Examples of 255 the Fréchet kernels for  $\gamma_c$  and  $\gamma_s$  are shown in Figure 7. Chevrot (2006) first noted the similarity between 256 the definitions of SI in Eq. (1) and finite-frequency traveltimes (e.g. Dahlen et al., 2000; Zhao et al., 2000), 257 which implies that the sensitivity kernels of SI to anisotropy parameters exhibit the distinct banana-doughnut 258 shapes typical in the sensitivities of finite-frequency traveltimes to isotropic velocities, as shown in Figure 7. 259



Figure 7. Examples of the sensitivity kernels of the SI to shear wave anisotropy parameters  $\gamma_c$  (a) and  $\gamma_s$ (b) shown in mapviews for the 146-km depth and in profiles (insets) along source-receiver path AA' for station WXT (blue triangle in Figure 1). The SKS wave is from the 24 June 2019 event in New Zealand with a back azimuth of 122°, and the waveforms are shown in Figure 4.

As in Lin et al. (2014b), we adopt a wavelet-based model parameterization (Chiao & Kuo, 2001) to obtain multi-scale resolutions in both sparsely- and densely-sampled regions of data coverage. The inverse problem can be expressed as:

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$$(\mathbf{G}\mathbf{W}^{-1})(\mathbf{W}\mathbf{m}) = \mathbf{d},\tag{5}$$

where **G** is the Gram matrix containing the Fréchet kernels, **W** is the 3D wavelet transform matrix, **m** is the vector containing the model parameters at spatial nodes, and **d** is the data vector of the SI measurements. In this study, we first parameterize the model by a 3D mesh with  $33 \times 33$  nodes horizontally and 17 nodes vertically, and apply the operator  $(\mathbf{W}^{-1})^T$  on each row of **G**. A damped least-squares solution to the inverse problem in Eq. (5) is solved by the LSQR algorithm (Paige & Saunders, 1982), with the damping factor  $\lambda$  selected empirically by a series of inversion experiments (see Section 2.4), and the final model can be obtained by an inverse wavelet transform. Readers may refer to Hung et al. (2011) for implementation details of the multi-scale parameterization. The wavelet approach achieves a higher spatial resolution where path coverage is better and a poorer resolution in less well-sampled regions, leading to results with objective and data-driven multi-scale resolution.

#### 279 **2.4 Resolution tests**

Careful selection of the damping factor  $\lambda$  and objectively evaluating the resolution for a given dataset in an inversion problem are paramount to interpreting the inversion results. Thus, it is important to characterize how reliable our anisotropic models are, and to limit our interpretations to robust features that are wellresolved and well-constrained by the observations.

We present three sets of resolution tests to illustrate the selection of the optimal damping factor and 284 285 the resulting resolution of our SI dataset. In each resolution test, an input model is constructed for an assumed spatial variation of anisotropy strength  $\gamma$  and azimuth  $\phi_f$ . Then, the spatial distributions of  $\gamma_c$ 286 and  $\gamma_s$  are obtained from  $\gamma$  and  $\phi_f$ , and synthetic SI data for the input model are calculated for all SKS 287 paths in our final dataset by integrating the products of the 3D sensitivity kernels  $K_{\gamma_c}^S$  and  $K_{\gamma_s}^S$  with  $\gamma_c$  and 288  $\gamma_s$  according to Eq. (4). Gaussian noise with the same standard deviation as our final inversion dataset 289 (0.074 s) is added to the synthetic SI data. We first perform a series of tests using checkerboards of two 290 different sizes (1° x 1° and 1.5° x 1.5°) with horizontally alternating azimuthal angles of fast axes  $\phi_f =$ 291 90° and  $\phi_f = 45^\circ$  but a fixed anisotropy strength  $\gamma = 4\%$ . Figure 8 shows the input and recovered models 292 for different values of the damping factor  $\lambda$  at a depth of 55 km for the 1° x 1° checkerboard test. The 293 effect of damping factor on the anisotropy pattern in the inversion result is obvious. The recovery results 294 suggest an optimal damping factor of  $\lambda = 4$ , and the checkerboard pattern is well-resolved in most parts 295 of the study area except in the northwest and southeast corners where stations are more sparse. Results for 296 the recovered models at different depths as well as for 1.5° x 1.5° checkerboards are presented in Figures 297

298 S1 and S2 in Supporting Information. Results of similar tests with alternating orthogonal azimuths of the

fast axis are presented Figures S3 and S4 in Supporting Information.



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Figure 8. Resolution tests for the azimuth of symmetry axis using 1° x 1° checkerboard. (Top-left) The input model has horizontally alternating azimuthal angles of fast axes  $\phi_f = 90^\circ$  and  $\phi_f = 45^\circ$  shown by both the color and the directions of the line segments, and a fixed anisotropy strength  $\gamma = 4\%$  represented by the lengths of the line segments. The two white dashed lines show the locations of profiles NS and N'S' in Figure 11. The other panels show the recovered models at a depth of 55 km for different damping factors  $\lambda$ . We choose  $\lambda = 4$  as the optimal value for the damping factor (top-right panel). Red triangles show locations of stations used.

In the second set of tests, we use input checkerboard models with a fixed azimuthal angle of symmetry axis  $\phi_f = 22.5^\circ$  but alternating perturbations of anisotropy strength of  $\delta \gamma = \pm 3\%$  relative to a background anisotropy strength of  $\gamma = 4\%$ . Figure 9 shows the recovery results at the depth of 55 km using different damping factors for the 1° x 1° input checkerboard model. Results for the recovered models at different depths as well as for 1.5° x 1.5° checkerboards are presented in Figures S5 and S6 in Supporting Information.



Figure 9. Resolution test using 1° x 1° checkerboard. (Top-left) The input model has horizontally alternating anisotropy strengths shown by the colors representing perturbations of  $\delta \gamma = \pm 3\%$  relative to a background anisotropy strength of  $\gamma = 4\%$ , and a fixed azimuthal angle of symmetry axis  $\phi_f = 22.5^\circ$ . The other panels show the recovered models at a depth of 55 km for different damping factors  $\lambda$ . We choose  $\lambda = 4$  as the optimal value for the damping factor (top-right panel).

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The tests in Figures 8 and 9 both show that the damping factor of  $\lambda = 4$  yields the best resolved checkerboard pattern. For convenience, resolution test results at different depths for damping factor  $\lambda =$ 4 are collected and displayed in Figure 10.



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Figure 10. Resolution test using 1° x 1° checkerboard with damping factor  $\lambda = 4$ . (Upper panels) 324 Recovered model at the depths of 99 km, 165 km, 208 km and 274 km for an input model having 325 horizontally alternating azimuthal angles of fast axes  $\phi_f = 90^\circ$  and  $\phi_f = 45^\circ$  with a fixed anisotropy 326 strength  $\gamma = 4\%$ . The input model and the recovered model at 55-km depth are shown in Figure 8. (Lower 327 panels) Recovered model at the depths of 99 km, 165 km, 208 km and 274 km for an input model with 328 329 horizontally alternating anisotropy strengths of  $\gamma = 1\%$  and  $\gamma = 7\%$  (perturbations of  $\delta \gamma = \pm 3\%$ relative to a background anisotropy strength of  $\gamma = 4\%$ ) with a fixed azimuthal angle of symmetry axis 330  $\phi_f = 22.5^{\circ}$ . The input model and the recovered model at 55-km depth are shown in Figure 9. 331

The third set of tests is intended to probe the depth resolution of our SI dataset with full-waveform anisotropy tomography. We use input models with a fixed anisotropy strength of  $\gamma = 4\%$  but several layers of alternating azimuthal angles of symmetry axes. Figure 11 shows the recovery results for a 4layer input model along a north-south vertical cross-section (cross-section NS in the top-left panel of Figure 8) for different damping factors.



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Figure 11. Resolution tests for an input model with 4 layers of alternating azimuthal angles of symmetry axes  $\phi_f = 45^\circ$  and  $\phi_f = 135^\circ$  but a fixed anisotropy strength of  $\gamma = 4\%$ . Shown here are the input model (top-left panel) and recovered models for different damping factors  $\lambda$  along the north-south cross-section NS through the middle of the study region (see top-left panel in Figure 8). The optimal damping is  $\lambda = 4$ (middle-left panel).

Recovered models along another profile as well as for a 2-layer input model are presented in Supporting Information (Figures S7-S10). These results also suggest an optimal damping factor of  $\lambda = 4$ . Based on all the resolution test results, we determine  $\lambda = 4$  to be the optimal damping factor for the SI data inversions in this study. It should be noted that the noise level in the SI measurements can affect the choice of the damping factor. The higher the noise level, the larger the damping factor (see the tests in Figures S11–S14).

The above resolution tests demonstrate that with the available distributions of seismic stations and teleseismic earthquakes, our SKS wave SI dataset can resolve reasonably well the shear-wave azimuthal anisotropy with a horizontal dimension of  $1^{\circ}$  x  $1^{\circ}$  and a vertical thickness of ~100 km in the main part of the study region. Near the western border and in the SE corner, the resolution is poor due to larger station spacing and fewer crossing SKS ray paths.

354 **3 Result** 

As stated in Section 2.2, our final dataset contains 5,216 SI measurements at 111 permanent seismic

stations from 470 events, which are used to invert for the 3D azimuthal anisotropy structure of the study
region. Figure 12 displays our inverted model at five representative depths. The azimuths of fast axis along
four vertical profiles are shown in Figure 13.



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Figure 12. Three-dimensional anisotropic model for the SE margin of Tibetan Plateau at 55 km, 99 km, 165 km, 208 km and 274 km depths. The anisotropy strength and the azimuth of the fast axes are shown by the background color and the black line segments, respectively. The lengths of the line segments are proportional to the anisotropy strength. The black and green arrows denote the absolute plate motion (APM) of the Eurasian and Indian plates, respectively, according to the model NNR-MORVEL56 (Argus et al., 2011). The white dashed lines show locations of the profiles shown in Figure 13. Names of major

faults and tectonic blocks are the same as in Figure 1. Blue triangles show locations of stations used. Regions in which the recovery errors are above 20% are mask out. The recovery error is defined by  $E = (|\phi_f - \tilde{\phi}_f|/90^\circ) \times 100\%$ , where  $\tilde{\phi}_f$  and  $\phi_f$  are the input and recovered azimuthal angles in the resolution tests in Figures 8 and 10.





the background color. Note that red ( $\sim 0^{\circ}$ ) and blue ( $\sim 180^{\circ}$ ) colors both indicate nearly N-S fast axis, whereas yellow color indicates nearly E-W fast axis. Moho and lithosphere-asthenosphere boundary (LAB) are depicted by thick black and white lines, respectively. Names of major faults and tectonic blocks are the same as in Figure 1.

As a first check of the inversion result, we calculate the model-predicted conventional SKS splitting time  $\Delta t$  and fast-direction azimuth  $\theta$  by integrating the anisotropy model (Figure 12) with sensitivity kernels (Figure 7) according to Eq. (4) and compare with observations. The comparison in Figure 6 shows excellent agreement between the model-predicted splitting parameters with observations at all stations.

We also carry out a recovery test for our inversion result, in which we use the model in Figure 12 as the input model and calculate the synthetic SI data, then invert them using the same model discretization and damping factor as in the inversion of the real data. The comparison in Figure 14 shows that the output model strongly resembles the input model at all depths. The synthetic inversion successfully reproduces all the major features of the input model.

386 4 Discussion

#### 387 **4.1 Comparison with previous results**

SKS splitting observation has been widely used to investigate the upper mantle anisotropy in the SE margin of 388 Tibetan Plateau (e.g., Lev et al., 2006; Flesch et al., 2005; Wang et al., 2008; Shi et al., 2012; Chang et al., 2015; 389 Yang et al., 2018; Liu et al., 2020; Huang & Chevrot, 2021; Li et al., 2021). In Figure 15, we compare our 390 splitting result with those of Chang et al. (2015) and Liu et al. (2020) whose SKS splitting data are available 391 392 from the Shear Wave Splitting Product Query website (http://ds.iris.edu/spud/swsmeasurement). At most stations, predictions by our model are consistent with both studies. Major differences occur at two stations at 393 the southern tip of SCB (highlighted by the blue ellipse in Figure 15), where anomalously large splitting times 394 395 can be seen in the results of Liu et al. (2020), shown by the red line segments. There are also discrepancies in both delay time and fast-axis azimuth at a few stations across the LJ-XJHF (highlighted by the red ellipses in 396

Figure 15) where Chang et al. (2015) reported large splitting times and nearly N-S fast axes, whereas our model predicts moderate-to-small splitting times and more E-W orientation of fast axes. Note that a strictly quantitative comparison of different studies is difficult due to differences in earthquake and window selections, data processing and frequency bands used in the measurement of SKS splitting parameters.



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Figure 14. Result of the recovery test. The recovered and input anisotropic models are compared at five depths marked at the bottom-right corner in each panel. Black and red line segments denote the directions of fast axes in the input and recovered models, respectively, and their lengths indicates the anisotropy strength. Names of major faults and tectonic blocks are the same as in Figure 1. Blue triangles show locations of stations used.



Figure 15. Comparison of SKS splitting times and fast axis directions predicted by our anisotropy model (Figure 12) with previous studies. Black line segments are results from the present study, whereas blue and red line segments show results from Chang et al. (2015) and Liu et al. (2020), respectively. Blue and red ellipses highlight stations where there are large discrepancies in splitting parameters between this study and Liu et al.

(2020) and Chang et al. (2015), respectively. Names of major faults and tectonic blocks are the same as in
Figure 1.

Depth variation of anisotropy can also be investigated by surface-wave tomography studies (e.g. 414 Montagner & Nataf, 1986; Debayle et al., 2005; Lin et al., 2011). Bao et al. (2020) and Zhang et al. (2023) 415 used Rayleigh-wave dispersion curves measured from ambient noise and teleseismic earthquakes to invert 416 for the azimuthal anisotropy in the SE margin of Tibetan Plateau. In general, both surface-wave results 417 show a lower anisotropy strength of 2-3% with smoother spatial variation patterns than our model 418 obtained from SKS splitting intensity. At the relatively shallow depths (~55 km) in the vicinity of the 419 LMSF and LJ-XJHF, all models exhibit fast axes nearly parallel to the fault zones. Surface-wave 420 tomography acquires depth resolution from different depth distributions of sensitivity of surface waves of 421 different frequencies, but its horizontal resolution is limited due to the horizontal averaging of horizontally 422 propagating surface waves. In addition, almost all surface-wave tomography studies use fundamental 423 424 Rayleigh waves, which also limits the depth resolution. On the other hand, our fullwave SI inversion uses nearly vertically propagating SKS waves, which acquires spatial resolution from the 3D sensitivity kernels 425 (Figure 7) with higher resolution horizontally than vertically. Therefore, discrepancies between models 426 obtained by surface-wave dispersion and SKS SI can be expected given the current station distribution 427 and available data. 428

#### 429 **4.2 Overall variation of shear-wave anisotropy**

Our anisotropy model exhibits an overall anisotropy strength  $\gamma$  of up to 6% in the lithosphere and asthenosphere. Both the strength and fast-axis show strong lateral and vertical variations (Figures 12 and 13). Within the lithosphere (depths of 55 km and 99 km), the anisotropy is strong in the vicinity of the LMSF and LJ-XJHF, the major fault zones in the region that form the boundary between the Tibetan Plateau and the YTC. At greater depths in the asthenosphere (e.g. 274 km), the entire region exhibits strong anisotropy with a largely NW-SE fast axis (predominantly green color in the lower parts in Figure 13), parallel to the direction of GPS observed surface velocity (Figure 1) and the APM (black arrows in
Figure 12), and turns to more E-W south of ~26°N. At intermediate depths (165 km and 208 km), the
anisotropy is highly variable, both in strength and direction of symmetry axis. Therefore, our model shows
an apparent decoupling of lithosphere and asthenosphere deformation in the SE margin of Tibetan Plateau,
consistent with previous studies (e.g. Flesch et al., 2005).

441 Most previous studies have reported two major features in the SKS splitting pattern in the region: an overall NW-SE fast axis and an E-W alignment of fast axes south of ~26°N latitude (e. g., Lev et al., 2006; Wang et 442 al., 2008; Chang et al., 2015; Huang & Chevrot, 2021), as shown in Figures 6 and 15. Our model shows that 443 this relatively simple splitting pattern observed at seismic stations on the surface is a manifestation of the vertical 444 averaging of complex variation of anisotropy over lithospheric and asthenospheric depths. For instance, along 445 the LMSF, the fast axes in the lithospheric depths (e.g. 55 km and 99 km) are mainly oriented NE-SW with 446 relatively high anisotropy strength, but change to NW-SE in the asthenosphere. As a result, we observe 447 relatively small SKS splitting times (< 0.5 s) with a NW-SE fast axis at the surface. On the other hand, the 448 449 moderate SKS splitting times (1 s or less) and E-W alignment of fast axes south of ~26°N are resulted primarily from a strong E-W azimuthal anisotropy (green and yellow colors south of 26°N and below the LAB in 450 451 Figures 13c and 13d) generated by an E-W asthenospheric flow due to the eastward subduction of the 452 Indian Plate under Myanmar (e.g. Yang et al., 2022). Regional variation of lithospheric and asthenospheric 453 contributions to anisotropy leads to small but discernable differences in SKS splitting observations. For example, under the SYRB, the anisotropy in the lithosphere between 24°N and 26°N is relatively strong 454 (55 km and 99 km depths in Figure 12) with different orientations of fast axes. The vertical integration of 455 456 lithospheric and asthenospheric anisotropy results in horizontally variable SKS splitting times and fast axes observed at seismic stations on the surface, as shown by the stations in SYRB in Figures 6 and 15. 457 In the rest of the study region south of 26°N, the asthenospheric contribution dominates, and the surface 458

459 SKS splitting parameters exhibit relatively uniform E-W fast axes.

#### 460 **4.3 The Songpan-Ganzi Terrane**

In our study region, the SGT is bounded by the JSJF in the southwest and LMSF and LJ-XJHF in the 461 southeast (Li et al., 2021). In the eastern margin of the SGT, the eastward expansion of the Tibetan Plateau 462 is blocked by the SCB, part of the YTC. Our model shows two distinct types of anisotropy. In the area far 463 away from the LMSF, the anisotropy is moderate to strong (3–6%) with NW-SE fast axes in both 464 lithosphere and asthenosphere (light blue and green colors under Tibetan Plateau in Figures 13b and 13c), 465 parallel to the GPS surface velocity (Figure 1), the regional APM direction (black arrows in Figure 12) 466 and the asthenospheric flow, suggesting a vertical coupling of deformation there, in agreement with Flesch 467 et al. (2005). The NW-SE anisotropy in the SGT is consistent with previous tomography results from 468 Rayleigh waves (Legendre et al., 2015; Bao et al., 2020) and Pn waves (Lei et al., 2014). However, in the 469 area near the LMSF, the fast axis in the lithosphere (e.g. 55 km and 99 km depths) becomes NE-SW, 470 parallel to the proposed lithospheric material flow as a result of the resistance of SCB to the eastward 471 472 expansion of the Tibetan Plateau at LMSF. Near the LMSF, lithospheric material of the SGT appears to diverge at ~102° longitude, going in opposite directions (NE and SW). In the asthenosphere, the fast axis 473 returns to NW-SE, consistent with the regional APM. The relatively sharp turn of the fast axis in the 474 475 lithosphere from SE-NW to NE-SW at the southern tip of SGT is a clear evidence for the redirection of the lithospheric flow from eastward to southward in the SE margin of Tibetan Plateau. 476

#### 477 **4.4 The Yangtze Craton**

The YTC is a Precambrian continental block which was accreted to the North China Craton in the Triassic. Our study region covers the western margin of the YTC, namely the SCB and SYRB, and our model shows that the anisotropy in the region varies significantly both horizontally and vertically, as a result of the complex geodynamic evolution of the YTC involving multiple rounds of lithospheric reactivation, such

as the Permian-Triassic Emeishan flood basalt eruption (e.g. Xu et al., 2001), as well as interactions with 482 surrounding tectonic blocks. In our model, the SCB exhibits a moderate anisotropy with a NE-SW orientation 483 484 near the LMSF at shallow depths (e.g. 55 km and 99 km depths in Figure 12 and yellow and red colors under LMSF in Figure 13b), presumably the effect of a dominant shearing in the vicinity of the LMSF by 485 the moving SGT lithosphere in NE and SW directions. Further east, E-W compression gradually takes 486 over the NE-SW shear at shallow depths, and the fast axis of anisotropy turns to more E-W (e.g. 55 km, 487 99 km and 165 km). This turn of anisotropy from NE-SW to more E-W away from the LMSF has also 488 been observed in surface-wave studies (e.g. Zhang et al., 2023). At greater depths in the lithosphere (208 489 km in Figure 12 and red and blue colors under SCB in Figure 13d), the anisotropy under the SCB becomes 490 nearly N-S, which is consistent with a frozen anisotropy in the lithosphere generated by the mantle flow 491 in the Cenozoic (Li et al., 2021) without being modified by the SGT-SCB block interaction at shallower 492 depths. At asthenospheric depths (274 km), the anisotropy is largely aligned with the regional APM. 493

The anisotropy under the SYRB appears to be continuous across the ANH-ZMHF and XJF, 494 495 suggesting that both fault systems are crustal boundaries where the SYRB is escaping in the SE direction. In the lithosphere (55 km and 99 km under SYRB and above 208 km depth east of ANH-ZMHF between 496 the SCB and 26°N), the anisotropy has a complex pattern of relatively small-scale horizontal variations of 497 498 fast axes, perhaps a result of multiple phases of lithospheric reactivations that have modified the previously frozen anisotropy. In the asthenosphere (> 100 km depth in SYRB and > 200 km between the SCB and 499 500 26°N), the anisotropy largely follows the direction of the regional APM or mantle flow. Around 208-km 501 depth, the azimuthal anisotropy south of the SCB is very weak, which may indicate a more complex local 502 mantle flow pattern due to the regionally predominant NW-SE mantle flow being disturbed by the root of 503 the YTC lithosphere. South of 26°N, the anisotropy under the YTC returns to largely parallel to the 504 direction of APM.

#### 505 **5 Conclusions**

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In this study, we have carried out fullwave multiscale tomography to obtain a 3D model for the shearwave anisotropy in the SE margin of Tibetan Plateau. A total of 5,216 high-quality SKS splitting intensities are obtained from the broadband records of 470 teleseismic events at 111 permanent stations after a series of quality control measures. In conjunction with the 3D sensitivity kernels and a waveletbased parameterization, the splitting intensities are inverted to achieve a data-driven multi-scale resolution to anisotropy structure in the upper mantle.

The vertical variation of the anisotropy in our result indicates that the lithospheric and asthenospheric deformations are decoupled in the SE margin of the Tibetan Plateau. On the other hand, the anisotropy appears to be vertically consistent under the Songpan-Ganzi Terrane, suggesting a coupling of the deformations in the lithosphere and asthenosphere in the interior of the Tibetan Plateau. It is worthwhile to point out that the vertical resolution is limited due to the nearly vertical propagation of the SKS waves.

As a result, the Moho discontinuity and crustal anisotropy are poorly resolved in our anisotropy model.

The strength of anisotropy in our model is spatially variable and can reach 6%, with strongest anisotropy in the asthenosphere presumably due to large-scale and relatively steady mantle flow, and in the lithosphere in the vicinity of the Longmenshan Fault and Lijiang-Xiaojinhe Fault, probably due to large shearing effect generated by the relative movement between the Songpan-Ganzi Terrane and the Sichuan Basin.

The azimuth of the fast axis of anisotropy in the asthenosphere largely follows the direction of regional absolute plate motion or mantle flow, i.e. mostly southeastward beneath the Songpan-Ganzi Terrane and the Yangtze Craton and nearly east-west south of 26°N latitude. In the lithosphere, however, the fast axis is highly variable. In the Sichuan Basin, the frozen anisotropy dominates in the deep lithosphere; whereas at shallower depths, the anisotropy is modified by the interaction with the Songpan-Ganzi Terrane into southeastward in the vicinity of the Longmenshan Fault and nearly east-west further east. In the Sichuan-Yunnan Rhombic Block and east of the Anninghe-Zemuhe Fault, the azimuth of the fast axis of anisotropy exhibits complex spatial pattern due to multiple phases of lithospheric reactivation. The vertical integration of the contributions from complex lithospheric anisotropy and relatively uniform asthenospheric anisotropy gives rise to the seemingly simple pattern of conventional SKS splitting parameters observed at the surface. Our 3D model of azimuthal anisotropy provides important new insights into the lithospheric and asthenospheric dynamics in the SE margin of Tibetan Plateau.

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543

#### 544 **Open Research**

The raw waveform data, the measured splitting intensities as well as the obtained anisotropic model can
be accessed at Lin, Y., & Zhao, L. (2024).

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#### 769 Figure Captions

Figure 1. Map of the tectonic environment of SE margin of Tibetan Plateau with seismic stations (green-770 771 filled triangles) and epicenters (red open circles) of earthquakes of magnitude 5 and above from 2000 to 2022. The blue triangle marks the station WXT for which SKS waveforms, splitting intensities and 772 sensitivity kernels are shown in Figures 4, 5 and 7, respectively. Major active faults are shown by thick 773 black lines with abbreviated names in black, including LMSF: Longmenshan Fault; XSHF: Xianshuihe 774 Fault; JSJF: Jinshajiang Fault; LCJF: Lancangjiang Fault; LJ-XJHF: Lijiang-Xiaojinhe Fault; ANH-775 ZMHF: Anninghe-Zemuhe Fault; XJF: Xiaojiang Fault; and RRF: Red River Fault. Major active tectonic 776 blocks are indicated by abbreviated texts in blue, including SGT: Songpan-Ganzi Terrane; SCB: Sichuan 777 Basin; QTB: Qiangtang Block; SYRB: Sichuan-Yunnan Rhombic Block; YTC: Yangtze Craton; and ICB: 778 Indo-China Block. EHS stands for the Eastern Himalaya Syntax. Blue cross arrows show the surface strain 779 rate derived from GPS observation (Wang & Shen, 2020). Background color shows the topography. The 780 781 black box in the inset map indicates the location of the main figure.

**Figure 2.** Lithosphere thickness map. The global model LITHO1.0 (Pasyanos et al., 2014) is shown for the region north of 27°N, whereas to the south the regional model of Yang et al. (2017) is depicted. Names of major faults and tectonic blocks are the same as in Figure 1.

Figure 3. Distribution of 470 teleseismic events used for SKS splitting intensity inversion in this study.
Events of magnitudes Mw5.5 and greater in the epicentral distance range of 90°–130° during 2009–2020
are selected. The red box in the center indicates the study area.

Figure 4. An example of SKS window selection using SplitRacer. (a) and (b) show normalized radial and 788 789 transverse records, respectively, around the SKS arrival at station WXT (blue triangle in Figure 1) from 790 the 24 June 2019 earthquake in New Zealand with a back azimuth of 122°. The green lines mark the theoretical SKS and SKKS arrival times corresponding to the event of interest. Red vertical lines mark 791 792 the start and end times of the final window used for SI measurement. (c) Summed spectral power of the spectrogram of the radial and transverse records in (a) and (b). The red horizontal line marks the frequency 793 794 of maximum spectral power, and the red dashed lines show the frequency bounds of more than 80% of the maximum. The green vertical lines mark the time window in which the summed spectral power is 795 796 more than 50% of the maximum. (d) Summed spectral power over all frequencies. The green vertical lines 797 mark the time window in which the summed spectral energy is more than 50% of the maximum. (e) 798 STA/LTA ratio of the radial-component record used as quality check. The blue vertical line denotes the time when the STA/LTA ratio reaches its peak within the window determined in (d). The red horizontal 799

- line is the acceptance threshold of 2.1 (Link et al., 2022).
- **Figure 5**. Variations of measured SI with event back azimuth for station WXT (blue triangle in Figure 1)

before (a) and after (b) removal of SI outliers. The vertical error bars show two standard deviations  $(2\sigma_i)$ of individual measurements. The dashed lines represent the sinusoidal curves that best fit the measurements. According to Eq. (2), the conventional SKS splitting parameters for this station are:  $\Delta t =$ 1.08 s and  $\theta = 109.14^{\circ}$ .

- **Figure 6**. Black line segments show the conventional SKS fast-direction azimuths and delay times at all station obtained from the SI measurements according to Eq. (2). Red line segments are fast-direction azimuths and delay times derived from Eq. (2) based on the model-predicted SIs obtained by integrating the sensitivity kernels (see Figure 7) with the anisotropy model in Figure 12 according to Eq. (4). Names of major faults and tectonic blocks are the same as in Figure 1.
- Figure 7. Examples of the sensitivity kernels of the SI to shear wave anisotropy parameters  $\gamma_c$  (a) and  $\gamma_s$ (b) shown in mapviews for the 146-km depth and in profiles (insets) along source-receiver path AA' for station WXT (blue triangle in Figure 1). The SKS wave is from the 24 June 2019 event in New Zealand with a back azimuth of 122°, and the waveforms are shown in Figure 4.
- Figure 8. Resolution tests for the azimuth of symmetry axis using 1° x 1° checkerboard. (Top-left) The input model has horizontally alternating azimuthal angles of fast axes  $\phi_f = 90^\circ$  and  $\phi_f = 45^\circ$  shown by both the color and the directions of the line segments, and a fixed anisotropy strength  $\gamma = 4\%$  represented by the lengths of the line segments. The two white dashed lines show the locations of profiles NS and N'S' in Figure 11. The other panels show the recovered models at a depth of 55 km for different damping factors  $\lambda$ . We choose  $\lambda = 4$  as the optimal value for the damping factor (top-right panel). Red triangles show locations of stations used.
- Figure 9. Resolution test using 1° x 1° checkerboard. (Top-left) The input model has horizontally alternating anisotropy strengths shown by the colors representing perturbations of  $\delta \gamma = \pm 3\%$  relative to a background anisotropy strength of  $\gamma = 4\%$ , and a fixed azimuthal angle of symmetry axis  $\phi_f = 22.5^\circ$ . The other panels show the recovered models at a depth of 55 km for different damping factors  $\lambda$ . We choose  $\lambda = 4$  as the optimal value for the damping factor (top-right panel).
- Figure 10. Resolution test using 1° x 1° checkerboard with damping factor  $\lambda = 4$ . (Upper panels) Recovered model at the depths of 99 km, 165 km, 208 km and 274 km for an input model having horizontally alternating azimuthal angles of fast axes  $\phi_f = 90^\circ$  and  $\phi_f = 45^\circ$  with a fixed anisotropy strength  $\gamma = 4\%$ . The input model and the recovered model at 55-km depth are shown in Figure 8. (Lower

- panels) Recovered model at the depths of 99 km, 165 km, 208 km and 274 km for an input model with horizontally alternating anisotropy strengths of  $\gamma = 1\%$  and  $\gamma = 7\%$  (perturbations of  $\delta\gamma = \pm 3\%$ relative to a background anisotropy strength of  $\gamma = 4\%$ ) with a fixed azimuthal angle of symmetry axis  $\phi_f = 22.5^\circ$ . The input model and the recovered model at 55-km depth are shown in Figure 9.
- Figure 11. Resolution tests for an input model with 4 layers of alternating azimuthal angles of symmetry axes  $\phi_f = 45^\circ$  and  $\phi_f = 135^\circ$  but a fixed anisotropy strength of  $\gamma = 4\%$ . Shown here are the input model (top-left panel) and recovered models for different damping factors  $\lambda$  along the north-south cross-section NS through the middle of the study region (see top-left panel in Figure 8). The optimal damping is  $\lambda = 4$ (middle-left panel).
- Figure 12. Three-dimensional anisotropic model for the SE margin of Tibetan Plateau at 55 km, 99 km, 840 165 km, 208 km and 274 km depths. The anisotropy strength and the azimuth of the fast axes are shown 841 by the background color and the black line segments, respectively. The lengths of the line segments are 842 proportional to the anisotropy strength. The black and green arrows denote the absolute plate motion 843 (APM) of the Eurasian and Indian plates, respectively, according to the model NNR-MORVEL56 (Argus 844 et al., 2011). The white dashed lines show locations of the profiles shown in Figure 13. Names of major 845 faults and tectonic blocks are the same as in Figure 1. Blue triangles show locations of stations used. 846 847 Regions in which the recovery errors are above 20% are mask out. The recovery error is defined by E = $(|\phi_f - \tilde{\phi}_f|/90^\circ) \times 100\%$ , where  $\tilde{\phi}_f$  and  $\phi_f$  are the input and recovered azimuthal angles in the 848 849 resolution tests in Figures 8 and 10.
- Figure 13. Vertical profiles of the azimuths of fast axis in our 3D anisotropy model along (a) 26°N latitude, (b) 32°N latitude, (c) 100°E longitude and (d) 104°E longitude. The azimuths of the fast axis are shown by the background color. Note that red (~0°) and blue (~180°) colors both indicate nearly N-S fast axis, whereas yellow color indicates nearly E-W fast axis. Moho and lithosphere-asthenosphere boundary (LAB) are depicted by thick black and white lines, respectively. Names of major faults and tectonic blocks are the same as in Figure 1.
- Figure 14. Result of the recovery test. The recovered and input anisotropic models are compared at five depths marked at the bottom-right corner in each panel. Black and red line segments denote the directions of fast axes in the input and recovered models, respectively, and their lengths indicates the anisotropy strength. Names of major faults and tectonic blocks are the same as in Figure 1. Blue triangles show locations of stations used.
- Figure 15. Comparison of SKS splitting times and fast axis directions predicted by our anisotropy model

- 862 (Figure 12) with previous studies. Black line segments are results from the present study, whereas blue and red
- line segments show results from Chang et al. (2015) and Liu et al. (2020), respectively. Blue and red ellipses
- <sup>864</sup> highlight stations where there are large discrepancies in splitting parameters between this study and Liu et al.
- 865 (2020) and Chang et al. (2015), respectively. Names of major faults and tectonic blocks are the same as in
- Figure 1.

1	Manuscript submitted to Journal of Geophysical Research: Solid Earth
2	Supporting Information for
3	Upper-mantle anisotropy in the southeastern margin of Tibetan Plateau revealed by
4	fullwave SKS splitting intensity tomography
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11	
12	This Supporting Information contains 14 Supplementary Figures S1-S14.
13	Figure S1: Resolution test results at different depths for 45° variation in the azimuth of
14	symmetry axis using 1° x 1° checkerboard and different damping factors.
15	Figure S2: Resolution test results at different depths for 45° variation in the azimuth of
16	symmetry axis using 1.5° x 1.5° checkerboard and different damping factors.
17	Figure S3: Resolution test results at different depths for 90° variation in the azimuth of
18	symmetry axis using 1° x 1° checkerboard and different damping factors.
19	Figure S4: Resolution test results at different depths for 90° variation in the azimuth of
20	symmetry axis using 1.5° x 1.5° checkerboard and different damping factors.
21	Figure S5: Resolution test results at different depths for anisotropy strength using 1° x 1°
22	checkerboard and different damping factors.
23	Figure S6: Resolution test results at different depths for anisotropy strength using 1.5° x
24	1.5° checkerboard and different damping factors.
25	Figure S7: Resolution tests for the azimuth of symmetry axis using a 2-layer input model
26	and different damping factors.
27	Figure S8: Resolution tests for anisotropy strength using a 2-layer input model and
28	different damping factors.
29	Figure S9: Resolution tests for the azimuth of symmetry axis using a 4-layer input model
30	and different damping factors.
31	Figure S10: Resolution tests for anisotropy strength using a 4-layer input model and
32	different damping factors.

- 33 Figure S11. Resolution tests result at 55-km depth for the azimuth of symmetry axis using
- 34 1° x 1° checkerboard with a substantially increased standard deviation (0.3 s) of the
- 35 Gaussian noise.
- 36 **Figure S12**. Result at different depth for the same resolution test as Figure S11.
- 37 Figure S13. Resolution tests result at 55-km depth for anisotropy strength using 1° x 1°
- 38 checkerboard with a substantially increased standard deviation (0.3 s) of the Gaussian noise.
- **Figure S14**. Result at different depth for the same resolution test as Figure S13.
- 40

## 41 Supplementary Figures



Figure S1. Resolution tests for the azimuth of symmetry axis using 1° x 1° checkerboard. The input model (top-left panel of Figure 8) has horizontally alternating azimuthal angles of fast axes  $\phi_f = 90^\circ$  and  $\phi_f = 45^\circ$  shown by both the color and the directions of the line segments, and a fixed anisotropy strength  $\gamma = 4\%$  represented by the lengths of the line segments. Shown here are recovered models using different damping factors (left to right) at different depths (top to bottom).





**Figure S2.** Resolution tests for the azimuth of symmetry axis using  $1.5^{\circ} \times 1.5^{\circ}$ checkerboard. The input model (top-left panel) has horizontally alternating azimuthal angles of fast axes  $\phi_f = 90^{\circ}$  and  $\phi_f = 45^{\circ}$  shown by both the color and the directions of the line segments, and a fixed anisotropy strength  $\gamma = 4\%$  represented by the lengths of the line segments. The rest of the panels show recovered models using different damping factors (left to right) at different depths (top to bottom).





57 **Figure S3**. Resolution tests for the azimuth of symmetry axis using 1° x 1° checkerboard. 58 The input model (top-left panel) has horizontally alternating azimuthal angles of fast axes 59  $\phi_f = 45^\circ$  and  $\phi_f = 135^\circ$  shown by both the color and the directions of the line segments, 60 and a fixed anisotropy strength  $\gamma = 4\%$  represented by the lengths of the line segments. 61 Shown here are recovered models using different damping factors (left to right) at different 62 depths (top to bottom).



64 **Figure S4**. Resolution tests for the azimuth of symmetry axis using  $1.5^{\circ} \times 1.5^{\circ}$ 65 checkerboard. The input model (top-left panel) has horizontally alternating azimuthal 66 angles of fast axes  $\phi_f = 45^{\circ}$  and  $\phi_f = 135^{\circ}$  shown by both the color and the directions of 67 the line segments, and a fixed anisotropy strength  $\gamma = 4\%$  represented by the lengths of 68 the line segments. The rest of the panels show recovered models using different damping 69 factors (left to right) at different depths (top to bottom).



Figure S5. Resolution tests for anisotropy strength using 1° x 1° checkerboard. The input model (top-left panel of Figure 8) has horizontally alternating anisotropy strengths shown by the colors representing perturbations of  $\delta \gamma = \pm 0.03$  relative to a background anisotropy strength of  $\gamma = 0.04$  and a fixed azimuthal angle of symmetry axis  $\phi_f = 22.5^\circ$ . Shown here are recovered models using different damping factors (left to right) at different depths (top to bottom).





Figure S6. Resolution tests for anisotropy strength using  $1.5^{\circ} \times 1.5^{\circ}$  checkerboard. The input model (top-left panel) has horizontally alternating anisotropy strengths shown by the colors representing perturbations of  $\delta \gamma = \pm 0.03$  relative to a background anisotropy strength of  $\gamma = 0.04$  and a fixed azimuthal angle of symmetry axis  $\phi_f = 22.5^{\circ}$ . The rest of the panels show recovered models using different damping factors (left to right) at different depths (top to bottom).



Figure S7. Resolution tests for an input model with 2 layers of different azimuthal angles of symmetry axes  $\phi_f = 45^\circ$  and  $\phi_f = 135^\circ$  but a fixed anisotropy strength of  $\gamma = 4\%$ .

88 Shown here are the input models (top two panels) along the NS and N'S' cross-sections 89 (see top-left panel in Figure 8 for the locations of the cross-sections) and recovered models

90 for different damping factors  $\lambda$  (lower panels).



Figure S8. Resolution tests for an input model with 2 layers of different anisotropy strengths shown by the colors representing perturbations of  $\delta \gamma = \pm 0.03$  relative to a background anisotropy strength of  $\gamma = 0.04$  and a fixed azimuthal angle of symmetry axis  $\phi_f = 22.5^\circ$ . Shown here are the input models (top two panels) along the NS and N'S' crosssections (see top-left panel in Figure 8 for the locations of the cross-sections) and recovered models for different damping factors  $\lambda$  (lower panels).



**Figure S9.** Resolution tests for an input model with 4 layers of alternating azimuthal angles of symmetry axes  $\phi_f = 45^\circ$  and  $\phi_f = 135^\circ$  but a fixed anisotropy strength of  $\gamma = 4\%$ . Shown here are the input model (top panel) along the N'S' cross-section (see top-left panel in Figure 8 for the location of the cross-section) and recovered models for different damping factors  $\lambda$  (lower panels).



104

105 **Figure S10**. Resolution tests for an input model with 4 layers of alternating anisotropy 106 strengths shown by the colors representing perturbations of  $\delta \gamma = \pm 0.03$  relative to a 107 background anisotropy strength of  $\gamma = 0.04$  and a fixed azimuthal angle of symmetry axis 108  $\phi_f = 22.5^\circ$ . Shown here are the input models (top two panels) along the NS and N'S' cross-109 sections (see top-left panel in Figure 8 for the locations of the cross-sections) and recovered 110 models for different damping factors  $\lambda$  (lower panels).





Figure S11. Resolution tests for the azimuth of symmetry axis using 1° x 1° checkerboard with substantially increased standard deviation (0.3 s) of the Gaussian noise. (Top-left) The input model has horizontally alternating azimuthal angles of fast axes  $\phi_f = 90^\circ$  and  $\phi_f =$ 45° shown by both the color and the directions of the line segments, and a fixed anisotropy

strength  $\gamma = 4\%$  represented by the lengths of the line segments. The other panels show

118 the recovered models at a depth of 55 km for different damping factors  $\lambda$ .





Figure S12. Recovered models for the same resolution test in Figure S11 using for different
 damping factors (left to right) at different depths (top to bottom).





Figure S13. Resolution test for anisotropy strength using 1° x 1° checkerboard with substantially increased standard deviation (0.3 s) of the Gaussian noise. (Top-left) The input model has horizontally alternating anisotropy strengths shown by the colors representing perturbations of  $\delta \gamma = \pm 3\%$  relative to a background anisotropy strength of  $\gamma = 4\%$ , and a fixed azimuthal angle of symmetry axis  $\phi_f = 22.5^\circ$ . The other panels show the recovered models at a depth of 55 km for different damping factors  $\lambda$ .





Figure S14. Recovered models for the same resolution test in Figure S13 using for different
 damping factors (left to right) at different depths (top to bottom).