## Red Line Diffuse-Like Aurora Driven by Time Domain Structures Associated with Braking Magnetotail Flow Bursts

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### Abstract

Magnetotail earthward-propagating fast plasma flows provide important pathways for magnetosphere-ionosphere coupling. This study reexamines a flow-related red-line diffuse-like aurora event previously reported by Liang et al., (2011), utilizing THEMIS and ground-based auroral observations from Poker Flat. We find that time domain structures (TDSs) within the flow bursts efficiently drive electron precipitation below a few keV, aligning with predominantly red-line auroral intensifications in this non-substorm event. The diffuse-like auroras sometimes coexisted with or potentially evolved from discrete forms. We forward model red-line diffuse-like auroras due to TDS-driven precipitation, employing the time-dependent TREx-ATM auroral transport code. The good correlation (0.77) between our modeled and observed red line emissions underscores that TDSs are a primary driver of the red-line diffuse-like auroras, though whistler-mode wave contributions are needed to fully explain the most intense red-line emissions.

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### **Key Points:**

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15	•	Predominantly red-line auroras are linked to flow bursts, TDSs, and ${<}1~{\rm keV}$ elec-
16		tron precipitation
17	•	For the first time, red-line diffuse-like auroras have been forward-modeled using
18		TREX-ATM with TDS inputs
19	•	A good correlation between forward-modeled and observed red-line emissions sug-
20		gests that TDSs are a major driver

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### 21 Abstract

Magnetotail earthward-propagating fast plasma flows provide important pathways 22 for magnetosphere-ionosphere coupling. This study reexamines a flow-related red-line 23 diffuse-like aurora event previously reported by Liang et al. (2011), utilizing THEMIS 24 and ground-based auroral observations from Poker Flat. We find that time domain struc-25 tures (TDSs) within the flow bursts efficiently drive electron precipitation below a few 26 keV, aligning with predominantly red-line auroral intensifications in this non-substorm 27 event. The diffuse-like auroras sometimes coexisted with or potentially evolved from dis-28 29 crete forms. We forward model red-line diffuse-like auroras due to TDS-driven precipitation, employing the time-dependent TREx-ATM auroral transport code. The good 30 correlation ( $\sim 0.77$ ) between our modeled and observed red line emissions underscores 31 that TDSs are a primary driver of the red-line diffuse-like auroras, though whistler-mode 32 wave contributions are needed to fully explain the most intense red-line emissions. 33

### <sup>34</sup> Plain Language Summary

Fast plasma flows in the magnetotail, traveling earthward at several hundred kilo-35 meters per second, transport energetic particles and magnetic flux into the inner mag-36 netosphere. Upon braking near Earth's high magnetic flux regions, they trigger plasma 37 instabilities and waves, leading to increased electric currents and particle precipitation 38 in the polar regions. This precipitation, depending on its driver, results in either diffuse-39 like auroras from electron pitch-angle scattering, or discrete auroras from field-aligned 40 electron acceleration and currents. Our case study highlights the important role of time-41 domain structures in diffuse-like aurora generation during flow braking. This reveals a 42 new aspect of magnetosphere-ionosphere coupling: the generation of diffuse-like auro-43 ras through electron scattering by time-domain structures in braking flow bursts. 44

### 45 1 Introduction

Since their discovery in the 1970's (Scarf et al., 1974; Gurnett et al., 1976), broadband electrostatic noise in the frequency range of tens of Hz up to several kHz have been
observed to be ubiquitous in various space plasmas (Ergun et al., 1998; Cattell et al.,
2005; Williams et al., 2006). These fluctuations, when examined in the time domain (Matsumoto
et al., 1994), manifest unique localized electrostatic solitary structures, comprising mostly
electron phase-space holes and double layers, which are collectively termed as time-domain
structures (TDSs) (Mozer et al., 2015).

Recently, there have been active discussions on the role of TDS-induced electron 53 scattering in generating diffuse auroras (Mozer et al., 2018; Nishimura et al., 2018; Vasko 54 et al., 2017, 2018; Shen et al., 2020, 2021). While Mozer et al. (2018) proposed that TDSs 55 contribute to electron precipitation into pulsating auroras, Nishimura et al. (2018) ar-56 gued that TDSs were more likely linked to discrete auroras or non-pulsating diffuse au-57 roras. Furthermore, a few studies showed that TDSs can efficiently scatter less than a 58 few keV electrons into the loss cone, thereby contributing to the diffuse auroral precip-59 itation (Vasko et al., 2017, 2018; Shen et al., 2020, 2021). Although individual conjunc-60 tion events have revealed moderate correlations between TDSs and diffuse auroras (see 61 Supporting Information in Shen et al., 2020), and their similar statistical distributions 62 in the nightside plasma sheet lend support (Newell et al., 2009; Malaspina et al., 2015), 63 direct evidence remains elusive. 64

Prior statistical studies revealed that TDSs are abundant in the plasma sheet within
most plasma injections and braking fast ion flows or bursty bulk flows (BBFs) (Gurnett
& Frank, 1977; Angelopoulos et al., 1992; Ergun et al., 2015; Malaspina et al., 2015). These
fast ion flows often manifest in the ionosphere as north-south auroral structures or "stream-

ers" (e.g., Lyons et al., 2012; Henderson, 2012; Nishimura et al., 2011). Auroral stream-69 ers are typically discrete in nature and are linked to upward field-aligned currents form-70 ing on the western edge of flow bursts (Nakamura et al., 2001; Nishimura et al., 2011). 71 However, some streamers may coexist with or evolve into diffuse-like aurora enhance-72 ments as they move equatorward and reach the inner edge of the plasma sheet (Henderson 73 et al., 1998). Within the braking flows, TDSs may contribute to electron precipitation 74 into structured diffuse-like and non-accelerated auroras (Sergeev et al., 2004; Shen et al., 75 2020).76

77 Low-energy ( $< \sim 1 \text{ keV}$ ) electron precipitation more effectively excites red-line (630) nm) auroras, whereas higher-energy (>1 keV) electrons predominantly induce green-line 78 (557.7 nm) and blue-line (427.8 nm) auroras. Thus, TDSs are likely linked to red-line-79 dominated auroras, especially when the plasma sheet state is prime for such low-energy 80 precipitation, like quiet times with low temperatures. In fact, to reliably identify TDS-81 driven diffuse auroras, non-substorm conditions are more suitable, as it is often challeng-82 ing to disentangle aggregated effects from different, concurrent drivers during substorms 83 that preclude direct linkage to flow bursts and TDSs (e.g., Sergeev et al., 2012; Shen et 84 al., 2023). 85

One potential example has been reported by Liang et al. (2011), where conjugate 86 THEMIS and ground-based optical and radar observations revealed flow-related auro-87 ral signatures, including high-latitude poleward boundary intensifications (PBIs), streamer-88 like discrete auroras, and red-line diffuse-like auroras that are associated with the dis-89 crete forms. In that study, THEMIS waveform data and potential TDS signatures have 90 not been investigated, leaving the impact of TDS-induced electron scattering on red-line 91 auroras unresolved. In this paper, we revisit this event, aiming to provide the first for-92 ward modeling and data-model comparison of red-line diffuse auroras driven by TDSs. 93

### <sup>94</sup> 2 Instrumentation and Models

We use the following data from THEMIS (Angelopoulos, 2008): electron and ion 95 fluxes and moments measured by the Electrostatic Analyzers (ESA) and the Solid State 96 Telecope (SST) instruments in the energy range of several eV up to 900 keV (McFadden 97 et al., 2008; Angelopoulos et al., 2008), DC vector magnetic fields at spin resolution ( $\sim$ 3 98 s) measured by the Fluxgate Magnetometers (FGM) (Auster et al., 2008), electric and 99 magnetic field wave power spectra within 1 Hz-4 kHz (FBK with 6 frequency bands and 100 FFP with 64 bands), and waveform data at 8,192 samples per second (sps, DC-coupled), 101 measured by the Electric Field Instrument (EFI), the search coil magnetometer (SCM), 102 and the Digital Fields Board (DFB) (Le Contel et al., 2008; Bonnell et al., 2008; Cully 103 et al., 2008). 104

The flow-related red-line auroral event occurred during 10:30–11:30 UT on March 105 3, 2009. Following Liang et al. (2011), we primarily use the ground-based Poker Flat merid-106 ian spectrograph (PFMSP) for auroral observations, recording every 15 seconds and cap-107 turing emissions in the red (630 nm), green (557.7 nm), blue (427.8 nm), and proton  $H_{\beta}$ 108 (486.1 nm) lines along the meridian. Complementarily, the Poker Flat digital all-sky cam-109 era (DASC) provided 2D auroral data to elucidate the longitudinal or magnetic local time 110 (MLT) extent of the red-line diffuse auroras. Note that the DASC does not provide ab-111 solute intensities of 630 nm emissions and it recorded images at 41-s cadence for this event. 112 We also use auroral observations from the THEMIS whitelight all-sky imager (ASI) at 113 Fort Yukon with 3-s cadence (Mende et al., 2008). 114

Although TDSs include different types of nonlinear electrostatic structures as noted by Mozer et al. (2015), electron phase space holes are the predominant type in the near-Earth plasma sheet (e.g., Ergun et al., 2015; Malaspina et al., 2018). Electron scattering by electron holes can be quantified using the quasi-linear approach initially devel<sup>119</sup> oped for an ensemble of plane waves (Kennel & Engelmann, 1966; Lyons, 1974). We use <sup>120</sup> the formulation of local pitch-angle diffusion coefficients  $D_{\alpha\alpha}$  derived in Vasko et al. (2018) <sup>121</sup> and refined by Shen et al. (2021). We bounce average the local diffusion rates using the <sup>122</sup> standard procedure (Lyons et al., 1972; Glauert & Horne, 2005):

$$\langle D_{\alpha_{eq}\alpha_{eq}}\rangle = v^{-1}\tau_B^{-1} \int D_{\alpha\alpha} \left(\partial \alpha_{eq}/\partial \alpha\right)^2 \left(\partial s/\partial \lambda\right) d\lambda/\cos\alpha,\tag{1}$$

where the integration is over the period of bounce motion  $\tau_B$ , and  $s = s(\lambda)$  is the length 123 of a field line. We use the relatively realistic T89 magnetic field model for bounce-averaging 124 (Tsyganenko, 1989; Ma et al., 2012). The latitudinal profile of  $\mathbf{B}$  near the equator is ad-125 justed to reduce stretching so as to align with local THEMIS magnetic field observations 126 (e.g., Ni et al., 2012). In addition, bounce-averaged pitch angle diffusion coefficients due 127 to transient whistler-mode waves are also calculated using the Full Diffusion Code (Ni 128 et al., 2008; Ma et al., 2020), including the Landau (n=0) and higher order  $(n \text{ up to } \pm 10)$ 129 cyclotron harmonic resonances. 130

We use the full precipitating electron distributions calculated from quasilinear dif-131 fusion theory to forward model TDS-driven red-line auroras. Because the varying timescale 132 of precipitation is typically shorter than the radiative timescale ( $\sim 110$  s) of 630 nm emis-133 sion, it is imperative to use a time-dependent auroral transport code, such as the TREx-134 ATM model (Liang et al., 2016). TREX-ATM is a time-dependent ionosphere aurora model, 135 adopting the two-stream electron transport code embedded in the GLOW model (Solomon 136 et al., 1988) and ambipolar diffusion to compute the electron transport in the atmosphere, 137 with additional capabilities to compute the impact ionization, secondary electron pro-138 duction, and impact excitation of neutrals. 139

### 140 **3 Results**

The event featuring bursty bulk flows and red-line auroras was observed from 11:40 141 to 11:20 UT on March 3, 2009. During this time, THEMIS-C spacecraft detected a se-142 ries of high-speed flow bursts (>400 km/s) in the Earth's plasma sheet, as shown in Fig-143 ure 1. Figures 1a–1c illustrate the magnetic fields threading all five THEMIS spacecraft, 144 modeled by T89, alongside the 2D ionospheric red-line auroras observed by Poker Flat 145 DASC. The red-line auroras spanned nearly  $45^{\circ}$  in magnetic longitude or 3 hrs in MLT, 146 primarily drifted eastward, occasionally moving equatorward, and exhibited both diffuse-147 like and discrete enhancements. 148

For this event complementary ground-based measurements, including THEMIS ASIs 149 and Poker Flat incoherent scatter radar, provided synergistic auroral and ionospheric 150 density observations as detailed by Liang et al. (2011). While not described again here, 151 the key findings include: (1) flow-related auroras featured intermittent poleward bound-152 ary intensifications (PBIs) and streamer-like green line emissions, occasionally coincid-153 ing with or transitioning to predominantly red-line diffuse-like auroras at lower latitudes; 154 (2) ionospheric density altitude profiles were consistent with the red-line emissions and 155 soft (<1 keV) electron precipitation; (3) southward-moving density patches corresponded 156 to earthward flow bursts and red-line auroral enhancements. It was noted that the ab-157 sence of substantial density enhancements above 200 km and the latitudinally-extended 158 red-line diffuse-like emissions were inconsistent with typical small-scale Alfvénic auro-159 ras and broadband acceleration (Chaston et al., 2003; Liang et al., 2019). However, for 160 some observed discrete auroras, such as the bright stripes around  $-100^{\circ}$  magnetic lon-161 gitude and discrete structures near the poleward boundary (likely linked to PBIs) in Fig-162 ure 1c, Alfvénic acceleration cannot be ruled out (Damiano et al., 2015; Tian et al., 2021; 163 Hull et al., 2022), though these aspects are not the main focus of this study. Auroral fea-164 tures will be further examined using multi-spectra PFMSP measurements. 165

Figures 1b–1c indicate that of the five THEMIS spacecraft, THB and THC were closest to midnight, with their footprints less than 1 hr *MLT* east of Poker Flat. THC



Figure 1. Overview of the flow-related red-line diffuse aurora event observed by THEMIS-C and Poker Flat all-sky camera (DASC). (a–b) Field line configuration threading the five THEMIS spacecraft in the GSM coordinates. Colored vectors represent ion flow velocities for each spacecraft. (c) Poker Flat DASC red-line auroras measured near 11:17 UT. The grids represent AACGM magnetic latitude and magnetic longitude, while the blue line marks the actual local midnight. The assumed emission height and approximate footprints of THEMIS-B (red box), THEMIS-C (green box), and Poker Flat have been mapped to ~200 km altitude using different Tsyganenko models of T89, T96, and T01. (d) THEMIS AE index. (e) Magnetic fields. (f) Ion flow velocities. (g-h) Electron energy spectrogram from 7 eV up to 900 keV. (i) Magnetic field FFP spectra in 10 Hz–2 kHz. (j) Electric field FBK spectra up to 2 kHz.



Figure 2. Zoomed-in view of the plasma waves within the two flow bursts near 10:58 UT (a–e) and 11:17 UT (f–j). Panels from top to bottom are: ion flow velocities, FFP magnetic field spectra, FBK electric field spectra, waveform measurements of high-frequency (>20 Hz and >50 Hz) electric fields in the field-aligned coordinates (**B** is in the z direction), and expanded views of milisecond-scale electron phase-space holes, also known as time-domain structures (TDSs).

detected all flow bursts exceeding 400 km/s within BBFs, whereas THB, positioned fur-168 ther earthward, registered only a weaker flow burst ( $\sim 200 \text{ km/s}$ ) that had waned dur-169 ing inward propagation. Figure 1d shows that the AE index remained below 80 nT through-170 out the event. Three strong flow bursts were recorded between 10:40 and 11:20 UT, each 171 lasting  $\sim 5$  minutes. The latter two, occurring closer to the equator, were accompanied 172 by  $|B_x|$  decreases, notable  $B_z$  dipolarizations (Figure 1e), and slightly energized electron 173 spectra (Figure 1h). Despite heightened flow activity, the plasma sheet temperature re-174 mained below 1 keV (Liang et al., 2011). Figures 1i–1j show that the flow bursts cor-175 responded to enhanced broadband electrostatic fluctuations (or TDSs) with frequencies 176 ranging from tens of Hz to several kHz. Additionally, sporadic electromagnetic wave power 177 enhancements, mostly below  $\sim 20$  Hz, indicate the presence of kinetic Alfvén waves (KAWs) 178 (e.g., Chaston et al., 2012; Malaspina et al., 2018; Shen et al., 2023). 179

Figure 2 provides a detailed view of wave spectra and waveform measurements as-180 sociated with two flow bursts near 10:58 UT and 11:17 UT. Throughout the flow inter-181 vals, there were continuous enhancements in broadband electrostatic fluctuations. Fig-182 ures 2d, 2e, 2i, and 2j clearly show that these broadband electrostatic fluctuations pre-183 dominantly consisted of electron phase-space holes, identified by a deficit of the electron 184 phase space density that exhibits bipolar parallel (in the z direction) and unipolar per-185 pendicular (in the x and y direction) electric fields, on the timescales of milisecond or 186 on the spatial scales of the local Debye length (Muschietti et al., 1999; Hutchinson, 2017; 187 Lotekar et al., 2020). In contrast, whistler-mode waves were more ephemeral, character-188 ized by sporadic, narrow-band magnetic power enhancements lasting less than one minute 189 within the flow intervals (Figures 2b and 2g). 190



Figure 3. (a, c) TDS- and whistler-driven bounce-averaged electron pitch-angle diffusion rates  $(\langle D_{\alpha_{eq}\alpha_{eq}}\rangle)$  as a function of energy and pitch angle. We use adjusted T89 model for bounce averaging. (b, d) Example energy flux distributions extrapolated to the loss cone (blue) and calculated within the loss cone (red) based on  $\langle D_{\alpha_{eq}\alpha_{eq}}\rangle$ . The gray-dashed curves indicate Maxwellian fittings of the precipitating distributions due to TDSs and whistlers, indicating characteristic energies of 1.5 keV and 1.8 keV near 11:15:18 UT.

Based on the observed characteristics of TDSs and whistler-mode waves, we pro-191 ceed to calculate bounce-averaged electron pitch-angle diffusion rates  $\langle D_{\alpha_{eq}\alpha_{eq}} \rangle$  and re-192 sulting electron distributions precipitating into the ionosphere (Shen et al., 2021; Ma et 193 al., 2012, 2020). TDSs can be generated locally at various latitudes and continuously along 194 the propagation of flow bursts and injections toward Earth (Ergun et al., 1998; Franz 195 et al., 2005; Mozer et al., 2015; Malaspina et al., 2015). We assume a latitudinal distri-196 bution of electron holes to be within  $\pm 25^{\circ}$  near the equator; increasing the latitude ex-197 tent to  $\pm 45^{\circ}$  enhances the diffusion rate near the loss cone by <20%. For transient whistler-198 mode waves, we determine the average wave spectra near 11:15 UT (Figure 2g) as hav-199 ing a Gaussian with a lower limit  $\omega_1=0.15\Omega_{ce}$ , a mean frequency  $\omega_2=0.3\Omega_{ce}$ , an upper 200 limit  $\omega_m = 0.45 \Omega_{ce}$ , a semi-bandwidth  $\delta \omega = 0.075 \Omega_{ce}$ , and a mean wave amplitude  $B_w = 50$ 201 pT. No whistler-mode waves were measured by waveform data in our event, thus we do 202 not have direct information on their wave normals. For the nightside equatorial plasma 203 sheet (Li et al., 2011; Agapitov et al., 2013; Meredith et al., 2021), we can assume par-204 allel whistlers confined within  $\pm 15^{\circ}$  latitude near the equator, having Gaussian wave nor-205 mals with a width  $\theta_w = 10^\circ$ , a minimum  $\theta_{min} = 0^\circ$ , and a maximum  $\theta_{max} = 30^\circ$ . We set 206 the background  $n_e \sim 0.6 \text{ cm}^{-3}$ ,  $f_{pe}/f_{ce} \sim 14.4$ , and  $L \sim 14.0$  based on THEMIS obser-207 vations. For bounce averaging, we use the T89 model with slight adjustments near the 208 equator to align with THEMIS-C local B observations. 209

Figures 3a and 3c present the calculated  $\langle D_{\alpha_{eq}\alpha_{eq}} \rangle$  as a function of energy and pitch 210 angle, resulting from electron scattering by TDSs and whistler-mode waves, respectively. 211 Near the loss cone ( $\sim 1^{\circ}$ ), TDSs efficiently induce pitch-angle scattering at energies be-212 low a few keV, with  $\langle D_{\alpha_{eq}\alpha_{eq}} \rangle$  on the order of  $10^{-4}$ – $10^{-3}$  s<sup>-1</sup>, assuming  $E_w \sim 1$  mV/m. 213 Due to large  $f_{pe}/f_{ce}$  for this event, whistler-mode waves can also drive efficient electron 214 scattering at energies as low as 200 eV, but the scattering rates exhibit a broad peak near 215 1 keV, extending up to 100 keV. This suggests that whistler-mode waves tend to pro-216 duce more energetic precipitation spectra compared with TDSs. 217

Example precipitating electron energy distributions are shown in Figures 3b and 3d. The differential energy flux (red pluses) within the loss cone can be estimated as  $x(E)J(E, \alpha_{LC})$ , where

$$x(E) = 2 \int_0^1 I_0(Z_0\tau)\tau d\tau / I_0(Z_0), \qquad (2)$$

<sup>221</sup> being the index of loss cone filling,  $J(E, \alpha_{LC})$  is the electron differential energy flux near <sup>222</sup> the loss cone (blue stars),  $I_0$  is the modified Bessel function with an argument  $Z_0 \simeq \alpha_{LC} / \sqrt{\langle D_{\alpha\alpha} \rangle_{LC} \cdot \tau_{loss}}$ <sup>223</sup> (Kennel & Petschek, 1966), and  $\tau_{loss}$  is assumed to be half of the bounce period.

Time-varying precipitating electron full distributions are input into the time-dependent 224 TREX-ATM code to forward model red-line auroras and compare their evolution with 225 that of the measured red-line diffuse-like auroras by PFMSP. We begin by examining PFMSP 226 auroral observations. Figures 4a–4d display approximately correlated ion flow bursts, TDS 227  $E_w$ , and enhanced auroral activities in the green line (577.7 nm), blue line (427.8 nm), 228 and red line (630 nm). To obtain wave amplitudes of TDSs from FBK spectra (Figure 2), 229 we integrate wave power in the frequency range of 20–900 Hz, as informed by TDS wavelet 230 analyses (Shen et al., 2021). The flow bursts (>400 km/s) were associated with inter-231 mittent poleward boundary intensifications (PBIs) near  $160^{\circ}$  elevation angle (or  $\sim 68^{\circ}$ 232 magnetic latitudes) in the green and blue line emissions, exhibiting distinct equatorward 233 movement. These features align with auroral streamers associated with earthward flow 234 channels in the plasma sheet (e.g., Lyons et al., 2012; Henderson, 2012; Nishimura et al., 235 2011). 236

However, the majority of the auroral enhancements were noted in the red line, pre-237 dominantly equatorward of the green-line enhancements. These enhanced red-line emis-238 sions, characterized by "patchiness" or short-lived activations, coincided with TDS am-239 plitude increases and exhibited discernible equatorward movement. As noted by Liang 240 et al. (2011), this movement was also in accord with the equatorward shift of soft elec-241 tron precipitation and enhanced density patches at altitudes above 150 km. Figure 4f 242 suggests that TDS-driven precipitation has characteristic energies mostly below 1 keV. 243 consistent with the overwhelmingly red-line emissions and soft electron precipitation. In 244 contrast, assuming the presence of whistler-mode waves throughout the flow bursts near 245 10:58 UT and 11:17 UT, these would produce precipitation with characteristic energies 246 exceeding 1 keV. 247

To compare with TREX-ATM model output, we need to identify PFMSP red-linediffuse-248 like aurora intensifications that were approximately conjugate with THEMIS observa-249 tions. First, we focus on the time interval between 10:45 and 11:30 UT, during which 250 the longitudinal separation between the PFMSP meridian and projected THEMIS-C foot-251 prints was less than 15°, or equivalently  $|\Delta MLT| < 1$  hr (Figure 4g). Second, we ad-252 just the PFMSP red-line data counts by subtracting a uniform background attributed 253 to ambient diffuse auroras, likely arising from non-TDS mechanisms like diffuse proton 254 precipitation (e.g., Lyons et al., 2015) as indicated by PFMSP  $H_{\beta}$ -line emissions. Third, 255 after applying Van Rhijin correction to data counts for aspect angle effects, we identify 256 an optimum range of elevation angles of 139°-146°, which probably corresponded to THEMIS 257 measurements. These optimal angles were determined by searching the maximum cor-258



Figure 4. Auroral features and TREx-ATM model results. (a) Ion flow velocity  $V_x$ . (b) Rootmean-square TDS electric field amplitudes  $|E_w|$  (20-900 Hz). (c-e) PFMSP 557.7-nm green-line, 427.8-nm blue-line, and 630-nm red-line intensity; dashed lines mark optimal elevation angle range for THEMIS conjunction. (f) Characteristic energies of electrons precipitated by TDS (black) and whistler-mode waves (blue). (g) Geographic longitudes of PFMSP and THEMIS-C footprints. (h) 630-nm intensity: PFMSP data (red) vs TDS-driven TREx-ATM model results (black, gray) based on THEMIS maximum and minimum fluxes around the equator (see Supporting Information). (i) TREx-ATM results with additional inputs from assumed, persistent whistler-mode waves; only maximum electron fluxes used. (j) Example ASI whitelight image showing diffuse-like auroras near PFMSP that likely evolved from discrete auroras. (k) Correlation coefficients for PFMSP and TREx-ATM modeled TDS-driven red-line auroras at various elevation angles.

relation coefficients between the modeled and observed red-line emissions across all elevation angles (Figure 4k).

The event exhibited coexisting auroral types due to various mechanisms. TDS-induced 261 soft electron precipitation leads to a broader altitude distribution of both 630 nm and 262 557.7 nm emissions, unlike the narrow altitude range associated with monoenergetic elec-263 tron acceleration in discrete arcs. In oblique optical observations, like those from PFMSP 264 and THEMIS ASI, this broader distribution may appear as more "diffusive" auroral forms 265 (see Figure 4j around the PFMSP meridian). Additional ASI auroral images are provided 266 in Supporting Information. The 630-nm emission, particularly sensitive to soft electron 267 precipitation, aligns closely with TDS effects, which we mainly focus on in the study. We 268 have carefully examined the observations from the Fort Yukon ASI (see Supplementary 269 Movie in Liang et al. (2011)) and excluded three periods when discrete arcs predomi-270 nately contribute to PFMSP red-line emissions. 271

Figure 4h shows a correlation coefficient of 0.77 between PFMSP observed and TREx-272 ATM forward-modeled red-line emissions, using inputs solely from TDS-driven precip-273 itation. The TREX-ATM model results, indicated by the black and gray lines, are based 274 on maximum and minimum flux levels measured by THEMIS near the equator during 275 10:40–11:20 UT. These in-situ measured fluxes were largely controlled by the spacecraft's 276 proximity to the true equator, where  $|B_{\tau}|$  is minimum (see Supporting Information). Fig-277 ure 4h reveals that the majority of the observed red-line diffuse-like auroral enhancements 278 during flow bursts can be explained by TDS-driven model results, albeit slightly under-279 estimated for the most intense emissions. Green bars highlight the time intervals of dis-280 crete emission peaks, unrelated to TDSs, near 11:00, 11:05, and 11:10 UT. Including these 281 intervals in our analysis lowers the correlation to 0.67. Figure 4i suggests that includ-282 ing potential whistler-mode wave effects along with TDSs in the model offers a better 283 explanation for the two bright red spots near 10:58 UT and 11:17 UT. This implies that 284 whistler-mode waves might also contribute to the most intense red-line diffuse emissions 285 during flow bursts. The compound effects of whistler-mode waves and field-aligned ac-286 celeration on energetic electron precipitation obscure TDS's role in 557.7-nm emissions. 287 Nonetheless, we have analyzed TDS-driven TREX-ATM 557.7-nm outputs and observed 288 a weaker correlation with PFMSP's green line ( $\sim 0.6$ ). 289

### <sup>290</sup> 4 Discussion

Plasma sheet flow bursts are typically linked to ionospheric auroral PBIs and equatorward-291 moving auroral streamers, as evidenced by white-light all-sky imager (ASI) observations, 292 often around substorm times (Donovan et al., 2008; Nishimura et al., 2011; Forsyth et 293 al., 2020; Yadav et al., 2022). In our non-substorm event, while green-line PBIs and forms 294 resembling equatorward-moving auroral streamers were noted, the most prominent fea-295 ture was the red-line auroras. These auroras extended towards lower latitudes and of-296 ten appeared more diffuse than discrete near PFMSP, likely due to wave scattering. Liang 297 et al. (2011) suggests that the predominance of red-line emissions was probably due to 298 a cooler inner plasma sheet ( $T_e < \sim 1 \text{ keV}$ ). Despite strong flows and potential associ-299 ation with hot electrons in the tail leading to streamer-inducing energetic precipitation 300 (Angelopoulos et al., 1992; Nishimura et al., 2020), there was an absence of significant 301 energetic injections intruding the inner plasma sheet around midnight (Figure 1g). This 302 could be attributed to differences in adiabatic motion or Alfvén layers between soft and 303 energetic electrons as they undergo convective, curvature and gradient drifts during in-304 ward propagation (e.g., Gabrielse et al., 2017). 305

While red-line emissions naturally accompany green-line emissions due to the  $O({}^{1}S)$ to  $O({}^{1}D)$  transition linked with 557.7 nm photo emission (Solomon et al., 1988), the dominance of red-line auroral emissions indicates a major low-energy precipitation population, distinct from the energetic precipitation responsible for green line excitation. This

low-energy precipitation was confirmed by ionospheric density altitude profiles measured 310 by PFISR (Liang et al., 2011). During the event period of our interest, ECH waves were 311 not observed, ruling them out as a cause for the auroral precipitation (Ni et al., 2012; 312 Zhang et al., 2015). Our investigations, encompassing wave spectra and waveform anal-313 yses, quasilinear calculations, and aurora forward modeling, point to electron scatter-314 ing by TDSs as the primary driver of the observed low-energy precipitation and red-line 315 diffuse-like auroras during this non-substorm event. Furthermore, during more active times 316 like substorms, TDS wave amplitudes tend to intensify (Ergun et al., 2015; Shen et al., 317 2021; Khazanov et al., 2021). It remains to be determined how much TDSs contribute 318 to the generation of red-line and green-line diffuse-like auroras during those active times. 319

### 320 5 Conclusion

This letter revisits a flow-related red-line aurora event initially reported by Liang 321 et al. (2011). Analyzing THEMIS spectra and waveform measurements, we have revealed 322 an abundance of time-domain structures (TDSs) associated with magnetotail flow bursts. 323 Applying quasilinear calculations, we have estimated TDS-driven electron distributions 324 precipitating into the ionosphere. Employing the time-dependent TREx-ATM auroral 325 transport code, we have forward-modeled red-line auroras due to TDS-driven precipi-326 tation. The good correlation of  $\sim 0.77$  between modeled and observed red-line emissions 327 suggests that TDSs likely cause the observed red-line diffuse-like auroras, which some-328 times coexisted with or likely evolved from discrete forms. However, to fully explain the 329 most intense red-line emissions, contributions from whistler-mode waves are also neces-330 sary. Our study suggests that wave-driven red-line diffuse-like auroras, associated with 331 braking flow bursts, offer a distinct yet complementary pathway of MI-coupling compared 332 with well-established flow-driven discrete auroral streamers, typically seen in the green 333 line and caused by field-aligned potential acceleration. 334

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### 343 Open Research

THEMIS data is available at http://themis.ssl.berkeley.edu/data/themis/. Poker Flat DASC and MSP can be accessed at http://optics.gi.alaska.edu/optics/ ?q=archive. Data access and processing was done using SPEDAS V4.1, see Angelopoulos et al. (2019).

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#### Red Line Diffuse-Like Aurora Driven by Time Domain 1 Structures Associated with Braking Magnetotail Flow 2 Bursts 3

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### **Key Points:**

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15	•	Predominantly red-line auroras are linked to flow bursts, TDSs, and ${<}1~{\rm keV}$ elec-
16		tron precipitation
17	•	For the first time, red-line diffuse-like auroras have been forward-modeled using
18		TREX-ATM with TDS inputs
19	•	A good correlation between forward-modeled and observed red-line emissions sug-
20		gests that TDSs are a major driver

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### 21 Abstract

Magnetotail earthward-propagating fast plasma flows provide important pathways 22 for magnetosphere-ionosphere coupling. This study reexamines a flow-related red-line 23 diffuse-like aurora event previously reported by Liang et al. (2011), utilizing THEMIS 24 and ground-based auroral observations from Poker Flat. We find that time domain struc-25 tures (TDSs) within the flow bursts efficiently drive electron precipitation below a few 26 keV, aligning with predominantly red-line auroral intensifications in this non-substorm 27 event. The diffuse-like auroras sometimes coexisted with or potentially evolved from dis-28 29 crete forms. We forward model red-line diffuse-like auroras due to TDS-driven precipitation, employing the time-dependent TREx-ATM auroral transport code. The good 30 correlation ( $\sim 0.77$ ) between our modeled and observed red line emissions underscores 31 that TDSs are a primary driver of the red-line diffuse-like auroras, though whistler-mode 32 wave contributions are needed to fully explain the most intense red-line emissions. 33

### <sup>34</sup> Plain Language Summary

Fast plasma flows in the magnetotail, traveling earthward at several hundred kilo-35 meters per second, transport energetic particles and magnetic flux into the inner mag-36 netosphere. Upon braking near Earth's high magnetic flux regions, they trigger plasma 37 instabilities and waves, leading to increased electric currents and particle precipitation 38 in the polar regions. This precipitation, depending on its driver, results in either diffuse-39 like auroras from electron pitch-angle scattering, or discrete auroras from field-aligned 40 electron acceleration and currents. Our case study highlights the important role of time-41 domain structures in diffuse-like aurora generation during flow braking. This reveals a 42 new aspect of magnetosphere-ionosphere coupling: the generation of diffuse-like auro-43 ras through electron scattering by time-domain structures in braking flow bursts. 44

### 45 1 Introduction

Since their discovery in the 1970's (Scarf et al., 1974; Gurnett et al., 1976), broadband electrostatic noise in the frequency range of tens of Hz up to several kHz have been
observed to be ubiquitous in various space plasmas (Ergun et al., 1998; Cattell et al.,
2005; Williams et al., 2006). These fluctuations, when examined in the time domain (Matsumoto
et al., 1994), manifest unique localized electrostatic solitary structures, comprising mostly
electron phase-space holes and double layers, which are collectively termed as time-domain
structures (TDSs) (Mozer et al., 2015).

Recently, there have been active discussions on the role of TDS-induced electron 53 scattering in generating diffuse auroras (Mozer et al., 2018; Nishimura et al., 2018; Vasko 54 et al., 2017, 2018; Shen et al., 2020, 2021). While Mozer et al. (2018) proposed that TDSs 55 contribute to electron precipitation into pulsating auroras, Nishimura et al. (2018) ar-56 gued that TDSs were more likely linked to discrete auroras or non-pulsating diffuse au-57 roras. Furthermore, a few studies showed that TDSs can efficiently scatter less than a 58 few keV electrons into the loss cone, thereby contributing to the diffuse auroral precip-59 itation (Vasko et al., 2017, 2018; Shen et al., 2020, 2021). Although individual conjunc-60 tion events have revealed moderate correlations between TDSs and diffuse auroras (see 61 Supporting Information in Shen et al., 2020), and their similar statistical distributions 62 in the nightside plasma sheet lend support (Newell et al., 2009; Malaspina et al., 2015), 63 direct evidence remains elusive. 64

Prior statistical studies revealed that TDSs are abundant in the plasma sheet within
most plasma injections and braking fast ion flows or bursty bulk flows (BBFs) (Gurnett
& Frank, 1977; Angelopoulos et al., 1992; Ergun et al., 2015; Malaspina et al., 2015). These
fast ion flows often manifest in the ionosphere as north-south auroral structures or "stream-

ers" (e.g., Lyons et al., 2012; Henderson, 2012; Nishimura et al., 2011). Auroral stream-69 ers are typically discrete in nature and are linked to upward field-aligned currents form-70 ing on the western edge of flow bursts (Nakamura et al., 2001; Nishimura et al., 2011). 71 However, some streamers may coexist with or evolve into diffuse-like aurora enhance-72 ments as they move equatorward and reach the inner edge of the plasma sheet (Henderson 73 et al., 1998). Within the braking flows, TDSs may contribute to electron precipitation 74 into structured diffuse-like and non-accelerated auroras (Sergeev et al., 2004; Shen et al., 75 2020).76

77 Low-energy ( $< \sim 1 \text{ keV}$ ) electron precipitation more effectively excites red-line (630) nm) auroras, whereas higher-energy (>1 keV) electrons predominantly induce green-line 78 (557.7 nm) and blue-line (427.8 nm) auroras. Thus, TDSs are likely linked to red-line-79 dominated auroras, especially when the plasma sheet state is prime for such low-energy 80 precipitation, like quiet times with low temperatures. In fact, to reliably identify TDS-81 driven diffuse auroras, non-substorm conditions are more suitable, as it is often challeng-82 ing to disentangle aggregated effects from different, concurrent drivers during substorms 83 that preclude direct linkage to flow bursts and TDSs (e.g., Sergeev et al., 2012; Shen et 84 al., 2023). 85

One potential example has been reported by Liang et al. (2011), where conjugate 86 THEMIS and ground-based optical and radar observations revealed flow-related auro-87 ral signatures, including high-latitude poleward boundary intensifications (PBIs), streamer-88 like discrete auroras, and red-line diffuse-like auroras that are associated with the dis-89 crete forms. In that study, THEMIS waveform data and potential TDS signatures have 90 not been investigated, leaving the impact of TDS-induced electron scattering on red-line 91 auroras unresolved. In this paper, we revisit this event, aiming to provide the first for-92 ward modeling and data-model comparison of red-line diffuse auroras driven by TDSs. 93

### <sup>94</sup> 2 Instrumentation and Models

We use the following data from THEMIS (Angelopoulos, 2008): electron and ion 95 fluxes and moments measured by the Electrostatic Analyzers (ESA) and the Solid State 96 Telecope (SST) instruments in the energy range of several eV up to 900 keV (McFadden 97 et al., 2008; Angelopoulos et al., 2008), DC vector magnetic fields at spin resolution ( $\sim$ 3 98 s) measured by the Fluxgate Magnetometers (FGM) (Auster et al., 2008), electric and 99 magnetic field wave power spectra within 1 Hz-4 kHz (FBK with 6 frequency bands and 100 FFP with 64 bands), and waveform data at 8,192 samples per second (sps, DC-coupled), 101 measured by the Electric Field Instrument (EFI), the search coil magnetometer (SCM), 102 and the Digital Fields Board (DFB) (Le Contel et al., 2008; Bonnell et al., 2008; Cully 103 et al., 2008). 104

The flow-related red-line auroral event occurred during 10:30–11:30 UT on March 105 3, 2009. Following Liang et al. (2011), we primarily use the ground-based Poker Flat merid-106 ian spectrograph (PFMSP) for auroral observations, recording every 15 seconds and cap-107 turing emissions in the red (630 nm), green (557.7 nm), blue (427.8 nm), and proton  $H_{\beta}$ 108 (486.1 nm) lines along the meridian. Complementarily, the Poker Flat digital all-sky cam-109 era (DASC) provided 2D auroral data to elucidate the longitudinal or magnetic local time 110 (MLT) extent of the red-line diffuse auroras. Note that the DASC does not provide ab-111 solute intensities of 630 nm emissions and it recorded images at 41-s cadence for this event. 112 We also use auroral observations from the THEMIS whitelight all-sky imager (ASI) at 113 Fort Yukon with 3-s cadence (Mende et al., 2008). 114

Although TDSs include different types of nonlinear electrostatic structures as noted by Mozer et al. (2015), electron phase space holes are the predominant type in the near-Earth plasma sheet (e.g., Ergun et al., 2015; Malaspina et al., 2018). Electron scattering by electron holes can be quantified using the quasi-linear approach initially devel<sup>119</sup> oped for an ensemble of plane waves (Kennel & Engelmann, 1966; Lyons, 1974). We use <sup>120</sup> the formulation of local pitch-angle diffusion coefficients  $D_{\alpha\alpha}$  derived in Vasko et al. (2018) <sup>121</sup> and refined by Shen et al. (2021). We bounce average the local diffusion rates using the <sup>122</sup> standard procedure (Lyons et al., 1972; Glauert & Horne, 2005):

$$\langle D_{\alpha_{eq}\alpha_{eq}}\rangle = v^{-1}\tau_B^{-1} \int D_{\alpha\alpha} \left(\partial \alpha_{eq}/\partial \alpha\right)^2 \left(\partial s/\partial \lambda\right) d\lambda/\cos\alpha,\tag{1}$$

where the integration is over the period of bounce motion  $\tau_B$ , and  $s = s(\lambda)$  is the length 123 of a field line. We use the relatively realistic T89 magnetic field model for bounce-averaging 124 (Tsyganenko, 1989; Ma et al., 2012). The latitudinal profile of  $\mathbf{B}$  near the equator is ad-125 justed to reduce stretching so as to align with local THEMIS magnetic field observations 126 (e.g., Ni et al., 2012). In addition, bounce-averaged pitch angle diffusion coefficients due 127 to transient whistler-mode waves are also calculated using the Full Diffusion Code (Ni 128 et al., 2008; Ma et al., 2020), including the Landau (n=0) and higher order  $(n \text{ up to } \pm 10)$ 129 cyclotron harmonic resonances. 130

We use the full precipitating electron distributions calculated from quasilinear dif-131 fusion theory to forward model TDS-driven red-line auroras. Because the varying timescale 132 of precipitation is typically shorter than the radiative timescale ( $\sim 110$  s) of 630 nm emis-133 sion, it is imperative to use a time-dependent auroral transport code, such as the TREx-134 ATM model (Liang et al., 2016). TREX-ATM is a time-dependent ionosphere aurora model, 135 adopting the two-stream electron transport code embedded in the GLOW model (Solomon 136 et al., 1988) and ambipolar diffusion to compute the electron transport in the atmosphere, 137 with additional capabilities to compute the impact ionization, secondary electron pro-138 duction, and impact excitation of neutrals. 139

### 140 **3 Results**

The event featuring bursty bulk flows and red-line auroras was observed from 11:40 141 to 11:20 UT on March 3, 2009. During this time, THEMIS-C spacecraft detected a se-142 ries of high-speed flow bursts (>400 km/s) in the Earth's plasma sheet, as shown in Fig-143 ure 1. Figures 1a–1c illustrate the magnetic fields threading all five THEMIS spacecraft, 144 modeled by T89, alongside the 2D ionospheric red-line auroras observed by Poker Flat 145 DASC. The red-line auroras spanned nearly  $45^{\circ}$  in magnetic longitude or 3 hrs in MLT, 146 primarily drifted eastward, occasionally moving equatorward, and exhibited both diffuse-147 like and discrete enhancements. 148

For this event complementary ground-based measurements, including THEMIS ASIs 149 and Poker Flat incoherent scatter radar, provided synergistic auroral and ionospheric 150 density observations as detailed by Liang et al. (2011). While not described again here, 151 the key findings include: (1) flow-related auroras featured intermittent poleward bound-152 ary intensifications (PBIs) and streamer-like green line emissions, occasionally coincid-153 ing with or transitioning to predominantly red-line diffuse-like auroras at lower latitudes; 154 (2) ionospheric density altitude profiles were consistent with the red-line emissions and 155 soft (<1 keV) electron precipitation; (3) southward-moving density patches corresponded 156 to earthward flow bursts and red-line auroral enhancements. It was noted that the ab-157 sence of substantial density enhancements above 200 km and the latitudinally-extended 158 red-line diffuse-like emissions were inconsistent with typical small-scale Alfvénic auro-159 ras and broadband acceleration (Chaston et al., 2003; Liang et al., 2019). However, for 160 some observed discrete auroras, such as the bright stripes around  $-100^{\circ}$  magnetic lon-161 gitude and discrete structures near the poleward boundary (likely linked to PBIs) in Fig-162 ure 1c, Alfvénic acceleration cannot be ruled out (Damiano et al., 2015; Tian et al., 2021; 163 Hull et al., 2022), though these aspects are not the main focus of this study. Auroral fea-164 tures will be further examined using multi-spectra PFMSP measurements. 165

Figures 1b–1c indicate that of the five THEMIS spacecraft, THB and THC were closest to midnight, with their footprints less than 1 hr *MLT* east of Poker Flat. THC



Figure 1. Overview of the flow-related red-line diffuse aurora event observed by THEMIS-C and Poker Flat all-sky camera (DASC). (a–b) Field line configuration threading the five THEMIS spacecraft in the GSM coordinates. Colored vectors represent ion flow velocities for each spacecraft. (c) Poker Flat DASC red-line auroras measured near 11:17 UT. The grids represent AACGM magnetic latitude and magnetic longitude, while the blue line marks the actual local midnight. The assumed emission height and approximate footprints of THEMIS-B (red box), THEMIS-C (green box), and Poker Flat have been mapped to ~200 km altitude using different Tsyganenko models of T89, T96, and T01. (d) THEMIS AE index. (e) Magnetic fields. (f) Ion flow velocities. (g-h) Electron energy spectrogram from 7 eV up to 900 keV. (i) Magnetic field FFP spectra in 10 Hz–2 kHz. (j) Electric field FBK spectra up to 2 kHz.



Figure 2. Zoomed-in view of the plasma waves within the two flow bursts near 10:58 UT (a–e) and 11:17 UT (f–j). Panels from top to bottom are: ion flow velocities, FFP magnetic field spectra, FBK electric field spectra, waveform measurements of high-frequency (>20 Hz and >50 Hz) electric fields in the field-aligned coordinates (**B** is in the z direction), and expanded views of milisecond-scale electron phase-space holes, also known as time-domain structures (TDSs).

detected all flow bursts exceeding 400 km/s within BBFs, whereas THB, positioned fur-168 ther earthward, registered only a weaker flow burst ( $\sim 200 \text{ km/s}$ ) that had waned dur-169 ing inward propagation. Figure 1d shows that the AE index remained below 80 nT through-170 out the event. Three strong flow bursts were recorded between 10:40 and 11:20 UT, each 171 lasting  $\sim 5$  minutes. The latter two, occurring closer to the equator, were accompanied 172 by  $|B_x|$  decreases, notable  $B_z$  dipolarizations (Figure 1e), and slightly energized electron 173 spectra (Figure 1h). Despite heightened flow activity, the plasma sheet temperature re-174 mained below 1 keV (Liang et al., 2011). Figures 1i–1j show that the flow bursts cor-175 responded to enhanced broadband electrostatic fluctuations (or TDSs) with frequencies 176 ranging from tens of Hz to several kHz. Additionally, sporadic electromagnetic wave power 177 enhancements, mostly below  $\sim 20$  Hz, indicate the presence of kinetic Alfvén waves (KAWs) 178 (e.g., Chaston et al., 2012; Malaspina et al., 2018; Shen et al., 2023). 179

Figure 2 provides a detailed view of wave spectra and waveform measurements as-180 sociated with two flow bursts near 10:58 UT and 11:17 UT. Throughout the flow inter-181 vals, there were continuous enhancements in broadband electrostatic fluctuations. Fig-182 ures 2d, 2e, 2i, and 2j clearly show that these broadband electrostatic fluctuations pre-183 dominantly consisted of electron phase-space holes, identified by a deficit of the electron 184 phase space density that exhibits bipolar parallel (in the z direction) and unipolar per-185 pendicular (in the x and y direction) electric fields, on the timescales of milisecond or 186 on the spatial scales of the local Debye length (Muschietti et al., 1999; Hutchinson, 2017; 187 Lotekar et al., 2020). In contrast, whistler-mode waves were more ephemeral, character-188 ized by sporadic, narrow-band magnetic power enhancements lasting less than one minute 189 within the flow intervals (Figures 2b and 2g). 190



Figure 3. (a, c) TDS- and whistler-driven bounce-averaged electron pitch-angle diffusion rates  $(\langle D_{\alpha_{eq}\alpha_{eq}}\rangle)$  as a function of energy and pitch angle. We use adjusted T89 model for bounce averaging. (b, d) Example energy flux distributions extrapolated to the loss cone (blue) and calculated within the loss cone (red) based on  $\langle D_{\alpha_{eq}\alpha_{eq}}\rangle$ . The gray-dashed curves indicate Maxwellian fittings of the precipitating distributions due to TDSs and whistlers, indicating characteristic energies of 1.5 keV and 1.8 keV near 11:15:18 UT.

Based on the observed characteristics of TDSs and whistler-mode waves, we pro-191 ceed to calculate bounce-averaged electron pitch-angle diffusion rates  $\langle D_{\alpha_{eq}\alpha_{eq}} \rangle$  and re-192 sulting electron distributions precipitating into the ionosphere (Shen et al., 2021; Ma et 193 al., 2012, 2020). TDSs can be generated locally at various latitudes and continuously along 194 the propagation of flow bursts and injections toward Earth (Ergun et al., 1998; Franz 195 et al., 2005; Mozer et al., 2015; Malaspina et al., 2015). We assume a latitudinal distri-196 bution of electron holes to be within  $\pm 25^{\circ}$  near the equator; increasing the latitude ex-197 tent to  $\pm 45^{\circ}$  enhances the diffusion rate near the loss cone by <20%. For transient whistler-198 mode waves, we determine the average wave spectra near 11:15 UT (Figure 2g) as hav-199 ing a Gaussian with a lower limit  $\omega_1=0.15\Omega_{ce}$ , a mean frequency  $\omega_2=0.3\Omega_{ce}$ , an upper 200 limit  $\omega_m = 0.45 \Omega_{ce}$ , a semi-bandwidth  $\delta \omega = 0.075 \Omega_{ce}$ , and a mean wave amplitude  $B_w = 50$ 201 pT. No whistler-mode waves were measured by waveform data in our event, thus we do 202 not have direct information on their wave normals. For the nightside equatorial plasma 203 sheet (Li et al., 2011; Agapitov et al., 2013; Meredith et al., 2021), we can assume par-204 allel whistlers confined within  $\pm 15^{\circ}$  latitude near the equator, having Gaussian wave nor-205 mals with a width  $\theta_w = 10^\circ$ , a minimum  $\theta_{min} = 0^\circ$ , and a maximum  $\theta_{max} = 30^\circ$ . We set 206 the background  $n_e \sim 0.6 \text{ cm}^{-3}$ ,  $f_{pe}/f_{ce} \sim 14.4$ , and  $L \sim 14.0$  based on THEMIS obser-207 vations. For bounce averaging, we use the T89 model with slight adjustments near the 208 equator to align with THEMIS-C local B observations. 209

Figures 3a and 3c present the calculated  $\langle D_{\alpha_{eq}\alpha_{eq}} \rangle$  as a function of energy and pitch 210 angle, resulting from electron scattering by TDSs and whistler-mode waves, respectively. 211 Near the loss cone ( $\sim 1^{\circ}$ ), TDSs efficiently induce pitch-angle scattering at energies be-212 low a few keV, with  $\langle D_{\alpha_{eq}\alpha_{eq}} \rangle$  on the order of  $10^{-4}$ – $10^{-3}$  s<sup>-1</sup>, assuming  $E_w \sim 1$  mV/m. 213 Due to large  $f_{pe}/f_{ce}$  for this event, whistler-mode waves can also drive efficient electron 214 scattering at energies as low as 200 eV, but the scattering rates exhibit a broad peak near 215 1 keV, extending up to 100 keV. This suggests that whistler-mode waves tend to pro-216 duce more energetic precipitation spectra compared with TDSs. 217

Example precipitating electron energy distributions are shown in Figures 3b and 3d. The differential energy flux (red pluses) within the loss cone can be estimated as  $x(E)J(E, \alpha_{LC})$ , where

$$x(E) = 2 \int_0^1 I_0(Z_0\tau)\tau d\tau / I_0(Z_0), \qquad (2)$$

<sup>221</sup> being the index of loss cone filling,  $J(E, \alpha_{LC})$  is the electron differential energy flux near <sup>222</sup> the loss cone (blue stars),  $I_0$  is the modified Bessel function with an argument  $Z_0 \simeq \alpha_{LC} / \sqrt{\langle D_{\alpha\alpha} \rangle_{LC} \cdot \tau_{loss}}$ <sup>223</sup> (Kennel & Petschek, 1966), and  $\tau_{loss}$  is assumed to be half of the bounce period.

Time-varying precipitating electron full distributions are input into the time-dependent 224 TREX-ATM code to forward model red-line auroras and compare their evolution with 225 that of the measured red-line diffuse-like auroras by PFMSP. We begin by examining PFMSP 226 auroral observations. Figures 4a–4d display approximately correlated ion flow bursts, TDS 227  $E_w$ , and enhanced auroral activities in the green line (577.7 nm), blue line (427.8 nm), 228 and red line (630 nm). To obtain wave amplitudes of TDSs from FBK spectra (Figure 2), 229 we integrate wave power in the frequency range of 20–900 Hz, as informed by TDS wavelet 230 analyses (Shen et al., 2021). The flow bursts (>400 km/s) were associated with inter-231 mittent poleward boundary intensifications (PBIs) near  $160^{\circ}$  elevation angle (or  $\sim 68^{\circ}$ 232 magnetic latitudes) in the green and blue line emissions, exhibiting distinct equatorward 233 movement. These features align with auroral streamers associated with earthward flow 234 channels in the plasma sheet (e.g., Lyons et al., 2012; Henderson, 2012; Nishimura et al., 235 2011). 236

However, the majority of the auroral enhancements were noted in the red line, pre-237 dominantly equatorward of the green-line enhancements. These enhanced red-line emis-238 sions, characterized by "patchiness" or short-lived activations, coincided with TDS am-239 plitude increases and exhibited discernible equatorward movement. As noted by Liang 240 et al. (2011), this movement was also in accord with the equatorward shift of soft elec-241 tron precipitation and enhanced density patches at altitudes above 150 km. Figure 4f 242 suggests that TDS-driven precipitation has characteristic energies mostly below 1 keV. 243 consistent with the overwhelmingly red-line emissions and soft electron precipitation. In 244 contrast, assuming the presence of whistler-mode waves throughout the flow bursts near 245 10:58 UT and 11:17 UT, these would produce precipitation with characteristic energies 246 exceeding 1 keV. 247

To compare with TREX-ATM model output, we need to identify PFMSP red-linediffuse-248 like aurora intensifications that were approximately conjugate with THEMIS observa-249 tions. First, we focus on the time interval between 10:45 and 11:30 UT, during which 250 the longitudinal separation between the PFMSP meridian and projected THEMIS-C foot-251 prints was less than 15°, or equivalently  $|\Delta MLT| < 1$  hr (Figure 4g). Second, we ad-252 just the PFMSP red-line data counts by subtracting a uniform background attributed 253 to ambient diffuse auroras, likely arising from non-TDS mechanisms like diffuse proton 254 precipitation (e.g., Lyons et al., 2015) as indicated by PFMSP  $H_{\beta}$ -line emissions. Third, 255 after applying Van Rhijin correction to data counts for aspect angle effects, we identify 256 an optimum range of elevation angles of 139°-146°, which probably corresponded to THEMIS 257 measurements. These optimal angles were determined by searching the maximum cor-258



Figure 4. Auroral features and TREx-ATM model results. (a) Ion flow velocity  $V_x$ . (b) Rootmean-square TDS electric field amplitudes  $|E_w|$  (20-900 Hz). (c-e) PFMSP 557.7-nm green-line, 427.8-nm blue-line, and 630-nm red-line intensity; dashed lines mark optimal elevation angle range for THEMIS conjunction. (f) Characteristic energies of electrons precipitated by TDS (black) and whistler-mode waves (blue). (g) Geographic longitudes of PFMSP and THEMIS-C footprints. (h) 630-nm intensity: PFMSP data (red) vs TDS-driven TREx-ATM model results (black, gray) based on THEMIS maximum and minimum fluxes around the equator (see Supporting Information). (i) TREx-ATM results with additional inputs from assumed, persistent whistler-mode waves; only maximum electron fluxes used. (j) Example ASI whitelight image showing diffuse-like auroras near PFMSP that likely evolved from discrete auroras. (k) Correlation coefficients for PFMSP and TREx-ATM modeled TDS-driven red-line auroras at various elevation angles.

relation coefficients between the modeled and observed red-line emissions across all elevation angles (Figure 4k).

The event exhibited coexisting auroral types due to various mechanisms. TDS-induced 261 soft electron precipitation leads to a broader altitude distribution of both 630 nm and 262 557.7 nm emissions, unlike the narrow altitude range associated with monoenergetic elec-263 tron acceleration in discrete arcs. In oblique optical observations, like those from PFMSP 264 and THEMIS ASI, this broader distribution may appear as more "diffusive" auroral forms 265 (see Figure 4j around the PFMSP meridian). Additional ASI auroral images are provided 266 in Supporting Information. The 630-nm emission, particularly sensitive to soft electron 267 precipitation, aligns closely with TDS effects, which we mainly focus on in the study. We 268 have carefully examined the observations from the Fort Yukon ASI (see Supplementary 269 Movie in Liang et al. (2011)) and excluded three periods when discrete arcs predomi-270 nately contribute to PFMSP red-line emissions. 271

Figure 4h shows a correlation coefficient of 0.77 between PFMSP observed and TREx-272 ATM forward-modeled red-line emissions, using inputs solely from TDS-driven precip-273 itation. The TREX-ATM model results, indicated by the black and gray lines, are based 274 on maximum and minimum flux levels measured by THEMIS near the equator during 275 10:40–11:20 UT. These in-situ measured fluxes were largely controlled by the spacecraft's 276 proximity to the true equator, where  $|B_{\tau}|$  is minimum (see Supporting Information). Fig-277 ure 4h reveals that the majority of the observed red-line diffuse-like auroral enhancements 278 during flow bursts can be explained by TDS-driven model results, albeit slightly under-279 estimated for the most intense emissions. Green bars highlight the time intervals of dis-280 crete emission peaks, unrelated to TDSs, near 11:00, 11:05, and 11:10 UT. Including these 281 intervals in our analysis lowers the correlation to 0.67. Figure 4i suggests that includ-282 ing potential whistler-mode wave effects along with TDSs in the model offers a better 283 explanation for the two bright red spots near 10:58 UT and 11:17 UT. This implies that 284 whistler-mode waves might also contribute to the most intense red-line diffuse emissions 285 during flow bursts. The compound effects of whistler-mode waves and field-aligned ac-286 celeration on energetic electron precipitation obscure TDS's role in 557.7-nm emissions. 287 Nonetheless, we have analyzed TDS-driven TREX-ATM 557.7-nm outputs and observed 288 a weaker correlation with PFMSP's green line ( $\sim 0.6$ ). 289

### <sup>290</sup> 4 Discussion

Plasma sheet flow bursts are typically linked to ionospheric auroral PBIs and equatorward-291 moving auroral streamers, as evidenced by white-light all-sky imager (ASI) observations, 292 often around substorm times (Donovan et al., 2008; Nishimura et al., 2011; Forsyth et 293 al., 2020; Yadav et al., 2022). In our non-substorm event, while green-line PBIs and forms 294 resembling equatorward-moving auroral streamers were noted, the most prominent fea-295 ture was the red-line auroras. These auroras extended towards lower latitudes and of-296 ten appeared more diffuse than discrete near PFMSP, likely due to wave scattering. Liang 297 et al. (2011) suggests that the predominance of red-line emissions was probably due to 298 a cooler inner plasma sheet ( $T_e < \sim 1 \text{ keV}$ ). Despite strong flows and potential associ-299 ation with hot electrons in the tail leading to streamer-inducing energetic precipitation 300 (Angelopoulos et al., 1992; Nishimura et al., 2020), there was an absence of significant 301 energetic injections intruding the inner plasma sheet around midnight (Figure 1g). This 302 could be attributed to differences in adiabatic motion or Alfvén layers between soft and 303 energetic electrons as they undergo convective, curvature and gradient drifts during in-304 ward propagation (e.g., Gabrielse et al., 2017). 305

While red-line emissions naturally accompany green-line emissions due to the  $O({}^{1}S)$ to  $O({}^{1}D)$  transition linked with 557.7 nm photo emission (Solomon et al., 1988), the dominance of red-line auroral emissions indicates a major low-energy precipitation population, distinct from the energetic precipitation responsible for green line excitation. This

low-energy precipitation was confirmed by ionospheric density altitude profiles measured 310 by PFISR (Liang et al., 2011). During the event period of our interest, ECH waves were 311 not observed, ruling them out as a cause for the auroral precipitation (Ni et al., 2012; 312 Zhang et al., 2015). Our investigations, encompassing wave spectra and waveform anal-313 yses, quasilinear calculations, and aurora forward modeling, point to electron scatter-314 ing by TDSs as the primary driver of the observed low-energy precipitation and red-line 315 diffuse-like auroras during this non-substorm event. Furthermore, during more active times 316 like substorms, TDS wave amplitudes tend to intensify (Ergun et al., 2015; Shen et al., 317 2021; Khazanov et al., 2021). It remains to be determined how much TDSs contribute 318 to the generation of red-line and green-line diffuse-like auroras during those active times. 319

### 320 5 Conclusion

This letter revisits a flow-related red-line aurora event initially reported by Liang 321 et al. (2011). Analyzing THEMIS spectra and waveform measurements, we have revealed 322 an abundance of time-domain structures (TDSs) associated with magnetotail flow bursts. 323 Applying quasilinear calculations, we have estimated TDS-driven electron distributions 324 precipitating into the ionosphere. Employing the time-dependent TREx-ATM auroral 325 transport code, we have forward-modeled red-line auroras due to TDS-driven precipi-326 tation. The good correlation of  $\sim 0.77$  between modeled and observed red-line emissions 327 suggests that TDSs likely cause the observed red-line diffuse-like auroras, which some-328 times coexisted with or likely evolved from discrete forms. However, to fully explain the 329 most intense red-line emissions, contributions from whistler-mode waves are also neces-330 sary. Our study suggests that wave-driven red-line diffuse-like auroras, associated with 331 braking flow bursts, offer a distinct yet complementary pathway of MI-coupling compared 332 with well-established flow-driven discrete auroral streamers, typically seen in the green 333 line and caused by field-aligned potential acceleration. 334

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### 343 Open Research

THEMIS data is available at http://themis.ssl.berkeley.edu/data/themis/. Poker Flat DASC and MSP can be accessed at http://optics.gi.alaska.edu/optics/ ?q=archive. Data access and processing was done using SPEDAS V4.1, see Angelopoulos et al. (2019).

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# Supporting Information for "Red Line Diffuse-Like Aurora Driven by Time Domain Structures Associated with Braking Magnetotail Flow Bursts"

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Introduction

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This Supporting Information includes Figure (S1) illustrating the variability of parallel electron energy fluxes (0°-22.5° in pitch angle) near the loss cone. These fluxes are closely related to the magnetic field  $B_x$  parameter, reflecting the spacecraft's proximity to the true magnetic equator. The observed flux variability primarily results from changing magnetic field configurations during bursty bulk flows (see Figure 1 in the main text). For accurate calculations of precipitating electron distributions corresponding to ground-based red-line diffuse auroral observations, it is necessary to use realistic equatorial fluxes measured when  $|B_x|$  is near its minimum at the equator.

Figure 1a in the main text shows that THEMIS-B and THEMIS-C were near the equator, while the other three spacecraft were farther away. Here Figure S1 demonstrates that electron energy fluxes measured by all spacecraft increase as  $B_x$  decreases. However, unlike the fluxes from THEMIS-B and THEMIS-C (panels b and c), which plateau, those from THEMIS-A (panel a), THEMIS-D (panel d), and THEMIS-E (panel e) do not. Based on this dependence on  $B_x$ , we select equatorial flux data solely from THEMIS-B and THEMIS-C during intervals when  $B_x$  is near its minimum, i.e., within -8 to 5 nT. The fast survey (peer) electron energy flux data were collected during the time intervals of 10:20–11:40 UT, whereas the burst mode (peeb) electron fluxes were measured during 10:53–11:22 UT. Because burst mode data provide better angular coverage thus more reliable parallel flux measurements, we use burst mode data to obtain realistic electron fluxes. In calculating precipitating electron distributions, we keep THEMIS-measured local energy distribution unchanged but proportionally adjust the maximum flux level to

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the maximum and minimum equatorial fluxes according to panel f of Figure S1. The flux-adjusted energy distributions are used to drive the TREx-ATM model runs.

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We have included one figure (Figure S2) that displays four characteristic ASI auroral images from Fort Yukon, showcasing discrete and diffuse auroras observed at 10:49:30, 10:59:18, 11:09:48, and 11:13:45 UT. These images are projected to an altitude of 110 km and presented in AACGM latitude and longitude coordinates, with the scanning path of the Poker Flat meridian spectrograph marked by a blue line. Note that the three time intervals of discrete auroras identified in the main text based on ASI imagery are 10:59:09–10:59:45 UT, 11:02:10–11:05:33 UT, and 11:08:30–11:11:54 UT.





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Figure S1. (a)THEMIS-A peer (reduced mode) parallel  $(0^{\circ}-22.5^{\circ})$  in pitch angle) electron energy fluxes versus  $B_x$ . (b-e) Same as (a) but for THEMIS-B, C, D and E, respectively. (f) THEMIS-B and THEMIS-C peeb parallel electron energy fluxes versus energy. Electron fluxes are included only when  $B_x$  is in the range of [-8, 5] nT, indicating near-equatorial measurements. The maximum and minimum flux levels are marked by the gray solid and dahsed lines.



**Figure S2.** Examples of white-light ASI images captured at Fort Yukon near the PFMSP meridian, following the format of Figure 4j in the main text: (a) Diffuse structures overlaying discrete auroras, observed at 10:49:30 UT. (b) Discrete aurora beside fading diffuse-like auroras, observed at 10:59:18 UT. (c) Discrete auroral arc, observed at 11:09:48 UT. (d) Diffuse aurora to the east of PFMSP, observed at 11:13:45 UT.

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