# An Improved Bio-Physical Parameterization for Radiant Heating in the Surface Ocean

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#### Abstract

Solar heating of the upper ocean is a primary energy input to the ocean-atmosphere system, and the vertical heating profile is modified by the concentration of phytoplankton in the water, with consequences for sea surface temperature and upper ocean dynamics. Despite the development of increasingly complex modeling approaches for radiative transfer in the atmosphere and upper ocean, the simple parameterizations of radiant heating used in most ocean models are plagued by errors and inconsistencies. There remains a need for a parameterization that is reliable in the upper meters and contains an explicitly spectral dependence on the concentration of biogenic material, while maintaining the computational simplicity of the parameterizations currently in use. In this work, we assemble simple, observationally-validated physical modeling tools for the key controls on ocean radiant heating, and simplify them into a parameterization that fulfills this need. We then use observations from 64 spectroradiometer depth casts across 6 cruises, 13 surface hyperspectral radiometer deployments, and 2 UAV flights to probe the accuracy and uncertainty associated with the new parameterization. We conclude with a case study using the new parameterization to demonstrate the impact of chlorophyll concentration on the structure of diurnal warm layers, an investigation that was not possible to conduct accurately using previous parameterizations. The parameterization presented in this work equips researchers to better model global patterns of sea surface temperature, diurnal warming, and mixed-layer depths, without a prohibitive increase in complexity.

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10	Key Points:				
11 12	• Using empirical physical modeling tools, we develop a bio-physical radiant heating parameterization that remedies a gap in the literature.				
13 14	• In-situ observations of the underwater and surface solar spectrum and the surface albedo are used to probe the parameterization's accuracy.				
15 16	• The parameterization is applied in a first-of-its-kind case study of the sensitivity of diurnal warm layers to chlorophyll concentration.				

### 17 Abstract

Solar heating of the upper ocean is a primary energy input to the ocean-atmosphere system, and 18 the vertical heating profile is modified by the concentration of phytoplankton in the water, with 19 consequences for sea surface temperature and upper ocean dynamics. Despite the development 20 of increasingly complex modeling approaches for radiative transfer in the atmosphere and upper 21 22 ocean, the simple parameterizations of radiant heating used in most ocean models are plagued by errors and inconsistencies. There remains a need for a parameterization that is reliable in the 23 upper meters and contains an explicitly spectral dependence on the concentration of biogenic 24 material, while maintaining the computational simplicity of the parameterizations currently in 25 use. In this work, we assemble simple, observationally-validated physical modeling tools for the 26 key controls on ocean radiant heating, and simplify them into a parameterization that fulfills this 27 28 need. We then use observations from 64 spectroradiometer depth casts across 6 cruises, 13 surface hyperspectral radiometer deployments, and 2 UAV flights to probe the accuracy and 29 uncertainty associated with the new parameterization. We conclude with a case study using the 30 new parameterization to demonstrate the impact of chlorophyll concentration on the structure of 31 diurnal warm layers, an investigation that was not possible to conduct accurately using previous 32 parameterizations. The parameterization presented in this work equips researchers to better 33 model global patterns of sea surface temperature, diurnal warming, and mixed-layer depths, 34 35 without a prohibitive increase in complexity.

#### 36 Plain Language Summary

37 Most of the sunlight that hits our planet is absorbed by the ocean as heat, and the vertical distribution of this heat influences the circulation patterns of the atmosphere and upper ocean. 38 Phytoplankton in the ocean also absorb sunlight, preventing it from heating deeper water. It is 39 important, therefore, for ocean models to have a realistic representation of the solar heating 40 profile as a function of the phytoplankton concentration. However, there are issues with current 41 parameterizations of solar heating in ocean models. To remedy this, we use intuitive physical 42 modeling tools to develop a new parameterization that balances simplicity and accuracy and can 43 be applied across all types of ocean modeling. We compare our parameterization to several types 44 45 of direct measurements to better understand its accuracy and limitations. Finally, we conclude with a case study demonstrating the value of the new parameterization in an application that was 46 previously not possible. 47

## 48 **1 Introduction**

Sunlight is the primary energy input to our planet, and with over 70% of the earth's surface 49 covered by water, solar heating of the ocean plays a fundamental role in the behavior of the 50 coupled ocean-atmosphere system at scales large and small. Though it is often conceptualized as 51 a surface flux to the well-mixed upper layer of the ocean, some wavelengths of sunlight can 52 penetrate quite deep before being completely absorbed as heat, and any light that penetrates 53 deeper than the mixed layer will reinforce the stratification of the underlying pycnocline, 54 effectively decoupling it from direct influence on the sea surface temperature (SST). The vertical 55 distribution of solar heating is sensitive to the amount of absorptive material in the water, and in 56 the open ocean where phytoplankton and their derivatives are the primary absorbing constituents 57 ("case I waters"), the spectral absorption has been empirically shown to follow a robust power-58 law relationship to the chlorophyll concentration (Morel, 1988; Morel et al., 2007; Morel & 59 60 Maritorena, 2001). Consequently, spatiotemporal variability of the phytoplankton system in the

61 surface ocean exerts a significant influence on SST, mixed-layer heat budgets, and upper ocean

dynamics (Gildor et al., 2003; Iudicone et al., 2008; Kahru et al., 1993; Kantha & Clayson, 1994;
Murtugudde et al., 2002; Ohlmann et al., 1998; Sathyendranath et al., 1991; Simpson & Dickey,

- Murtugudde et al., 2002; Ohlmann et al., 1998; Sathyendranath et al., 1991; Simpson & Dickey
   1981; Zaneveld et al., 1981). The formation and evolution of diurnal warm layers (DWLs),
- 65 which can span broad areas of the tropical oceans, is especially sensitive to variability in the
- vertical distribution of solar heating (Bellenger & Duvel, 2009; Dickey & Simpson, 1983). At
- 67 larger scales, phytoplankton-induced SST variability may modify the strength of the Hadley and
- 68 Walker circulations (Gnanadesikan & Anderson, 2009), steer tropical cyclones away from the
- 69 equator (Gnanadesikan et al., 2010), slow the meridional overturning circulation (Sweeney et al.,
- 2005), and perhaps even play a role in Arctic Amplification (Hill, 2008). Meanwhile tantalizing
- evidence has been mounting for a possible coupling between phytoplankton and both the
- 72 Madden-Julian Oscillation (MJO) and El Nino-Southern Oscillation (ENSO), in which wind
- <sup>73</sup> bursts upwell nutrients and stimulate phytoplankton blooms, which subsequently trap solar rediction near the surface, increasing SST and thereby atmospheric convection during the
- radiation near the surface, increasing SST and thereby atmospheric convection during thequiescent phases of these oscillatory climate modes (Jin et al., 2013; Jochum et al., 2010;
- <sup>75</sup> quiescent phases of these oscillatory enhance modes (Jin et al., 2013, Joenum et al., 2016,
   <sup>76</sup> Lengaigne et al., 2007; Marzeion et al., 2005; Siegel et al., 1995; Strutton & Chavez, 2004).

The coupling between ocean physics and biology is an important piece of the climate system that has yet to be fully untangled. In one study coupling an ocean biogeochemical model to an

79 ocean general circulation model, Manizza et al. (2005) found that modification of light

80 penetration by phytoplankton biomass amplified the seasonal cycles of temperature, mixed layer

- depth, and sea ice cover by 10%, and these physical changes caused increases in phytoplankton biomass, a positive feedback that further amplified the physical perturbations. Another study by
- biomass, a positive feedback that further amplified the physical perturbations. Another study by
   Wetzel et al. (2006) added a marine biogeochemical model to a fully coupled atmosphere-ocean
- climate model, and in addition to observing similar amplifications in the magnitude of the
- seasonal cycle, the timing of the seasons became more realistic, with spring starting two weeks
- 86 earlier. These coupled experiments rely on formulations for the physics of solar absorption that
- are expressed as a function of chlorophyll concentration, as this is a primary output of
- biogeochemical models. Modern remote-sensing tools enable the retrieval of attenuation

coefficient estimates at specific wavelengths from satellites, and this has led some investigators to suggest that the direct observations of the optical properties should be used in physical

- modeling (e.g. Liu et al., 2020). While there are good reasons for this approach and undoubtedly
- research contexts in which it is the smart choice, it introduces significant complexity and largely
- precludes the possibility of coupling the physics to the biology through the shared variable of

<sup>94</sup> chlorophyll concentration. Thus we emphasize the continued need for a simple, chlorophyll-

95 dependent parameterization of radiant heating in the modern oceanographer's toolkit.

Despite an abundance of evidence as to the importance of (primarily biogenic) in-water absorbing constituents to the mixed-layer heat budget and SST, the extensive literature exploring

this bio-physical interaction has been riddled with errors, due at least in part to the inter-

disciplinary nature of bio-physical reseach. We present here a brief history of the development of

100 expressions for radiant heating in the surface ocean, and a discussion of their mis-application in

101 various studies.

## 102 **1.1 Foundational Work**

Early research represented the underwater irradiance profile as a single exponential (e.g. Denman, 1973). While this is physically valid for a given wavelength, it is problematic when

applied to the broadband solar irradiance due to spectral variability in the exponential decay 105 rates, which span many orders of magnitude across the wavelengths in the solar spectrum. Kraus 106 (1972) is generally credited as the first to suggest using a sum of two exponentials. Paulson & 107 Simpson (1977) fit a double-exponential form to Jerlov (1968, 1976)'s data for different optical 108 water types, and this became the standard parameterization used in many mixed-layer models 109 even to this day (e.g. Price et al., 1986). While the two exponentials are nominally split 110 somewhere around the edge of the visible domain, the split is not explicitly spectral and the 111 exponentials are simply tuned to best fit the available data. Zaneveld & Spinrad (1980) employed 112 an alternative fitting approach to the same underlying Jerlov data, using an arc-tangent curve to 113 better represent the rapid extinction of irradiance in the upper 10 meters. While these 114 parameterizations were an important step along the journey, the approach of choosing between 115 qualitative optical water types is outdated, thus we suggest that parameterizations based on 116 Jerlov data should not be in use today. Several other parameterizations were also developed in 117 the early 1980s to pursue specific physical research questions. Paulson & Simpson (1981) fit 9 118 explicitly spectral bands originally calculated by Schmidt (1908) and catalogued in Defant 119 (1961) for the purpose of calculating solar absorption across the ocean skin – this 120 121 parameterization was employed in early iterations of the COARE cool-skin correction (Fairall, Bradley, Godfrey, et al., 1996), but was shown to be inaccurate and modified for newer versions 122 to better fit available data (Wick et al., 2005). Given that it is built on a very small sample of 123 124 laboratory data acquired at the turn of last century using measurement techniques that have since been improved upon many times over, it too should not be used. Finally, Soloviev (1982) 125 performed a three-exponential fit to a combination of Jerlov data and a near-surface underwater 126 irradiance profile provided by Ivanoff (1977) to estimate a parameterization that remains in use 127 in modern warm layer studies (Fairall, Bradley, Godfrey, et al., 1996). 128

As the importance of biological agents in modifying visible-wavelength attenuation became 129 evident, chlorophyll-dependent spectral parameterizations were developed for the visible domain 130 as an improvement over Jerlov's qualitative water types (e.g. Siegel & Dickey, 1987; Smith & 131 Baker, 1978). This culminated in the work of Morel (1988), who combined 176 spectral 132 downwelling irradiance measurements from 12 cruises and fitted a power-law relationship to 133 predict spectral attenuation from chlorophyll concentration in 5nm bins across the 400-700nm 134 visible domain. Morel & Antoine (1994) then expanded the relationship to cover the full solar 135 spectrum for use in radiant heating studies, and this has served as the basis for most subsequent 136 chlorophyll-dependent parameterizations. However, the underlying work of Morel (1988) has 137 since been updated several times to incorporate more field data and better measurements of the 138 water's contribution to the attenuation (Morel et al., 2007; Morel & Maritorena, 2001). 139

With the development of increasingly powerful radiative transfer models came the 140 opportunity to simulate a wide variety of conditions and develop parameterizations from model 141 outputs. Ohlmann & Siegel (2000) used the HYDROLIGHT model to generate a full suite of 142 irradiance profiles for a wide range of chlorophyll concentrations, solar zenith angles, and cloud 143 indices, and then fit their results to a 4-exponential form without explicit spectral partitions. In a 144 follow-up paper, Ohlmann (2003) provides a simplification of the model dependent only on 145 chlorophyll concentration for use in climate models, at the cost of accuracy in the upper meters. 146 And finally, Lee et al. (2005) tried to improve upon this approach by developing an explicitly 147 spectral 2-band parameterization using the HYDROLIGHT model, this time as a function of 148 satellite-observable properties (absorption and backscatter at 490nm). 149

#### 150 **1.2 Issues in the Literature**

Application of the parameterizations discussed in Section 1.1 to various modeling efforts has been subject to a number of errors. These errors do not necessarily call into question the fundamental findings of the studies, but they have been propagated through multiple publications to the point that they are deeply rooted in the literature, introducing unnecessary inaccuracy into our research on bio-optical phenomena.

One pervasive issue infecting decades of work is the misinterpretation of the two exponential 156 bands presented in Paulson & Simpson (1977). These exponential functions were developed by 157 fitting underwater profiles of broadband irradiance in different depth segments, with one 158 exponential fit to the irradiance profile below 10m, and the other fit to the irradiance profile in 159 the top 6m. The fractions of incident irradiance assigned to the two bands just reflect the best fit 160 to the data, not a division of the solar spectrum at a particular wavelength. Yet Manizza et al., 161 (2005) mistook the partitioning between Paulson & Simpson (1977)'s two exponential bands to 162 be an explicitly spectral split located at 700nm, consequently assigning 42% of the incident 163 irradiance to the visible (400:700nm) waveband in their spectral model. There is an abundance of 164 observational evidence that the visible fraction of incident irradiance is approximately 45-47% 165 under clear skies (Britton & Dodd, 1976; Frouin et al., 1989; Goldberg & Klein, 1977; Howell et 166 al., 1983; Ivanoff, 1977; Kirk, 1994; Kvifte et al., 1983; McCree, 1966; Moon, 1940; 167 Papaioannou et al., 1993; Rao, 1984; Rodskjer, 1983; Strutton & Chavez, 2004; Weiss & 168 Norman, 1985; Yocum et al., 1964), and increases up to  $\sim 60\%$  with increasing cloud cover 169 170 (Blackburn & Proctor, 1983; Frouin et al., 1989; Nann & Riordan, 1991; Siegel et al., 1999; Tanre et al., 1979, and see Section 4 for further observational validation). It is fortuitous that the 171 42% partition presented in Paulson & Simpson (1977) for Jerlov Type I waters is reasonably 172 close to the spectral estimate of 45-47% (had a different decision been made as to the qualitative 173 water type, the partitioning would have been considerably less accurate). Manizza et al. (2005) 174 then made the unsubstantiated assumption that the visible-wavelength band could be further split 175 176 at 530nm with 50% of the visible irradiance assigned to each band – as it turns out, the modeling tools we introduce in Section 2 suggest that a 45/55% split would be more appropriate. 177

178 This unnecessarily erroneous setup has since been propagated into major earth system models, including the American GFDL-ESM4.1 (Stock et al., 2020), Australian ACCESS-OM2 179 (Kiss et al., 2020), and Chinese BCC-CSM2 (Wu et al., 2021). It also appears regularly in 180 bespoke modeling efforts to elicit key features of bio-physical interactions (e.g. Gnanadesikan & 181 Anderson, 2009; Holmes et al., 2019; Kim et al., 2015; Lengaigne et al., 2007; Lim et al., 2019; 182 Moeller et al., 2019; Twelves et al., 2021). While we anticipate that the broad conclusions 183 184 arrived at in these studies are unlikely to be significantly altered by this error, there is simply no reason for it to remain an issue in ocean modeling going forward, and the parameterization 185 developed in this work offers a straightforward replacement. 186

In addition to the misuse of Paulson & Simpson (1977) for spectral partitioning, there are 187 other cases in the literature of the wrong parameterization being selected for the task at hand. For 188 example, Prytherch et al. (2013) updated the Diurnal Warm Layer (DWL) model of Fairall et al. 189 (1996) to use the chlorophyll-dependent parameterization of Ohlmann (2003) rather than the 190 default profile from Soloviev (1982). Yet Ohlmann (2003) explicitly states that "the 191 192 parameterization is not valid for depths shallower than 2m," making it a terrible choice for DWL modeling. Unsurprisingly, Prytherch et al. (2013) report that accurate representation of solar 193 absorption still appears to be a primary confounding factor in modeling DWLs. 194

195 Similarly, Gentemann et al. (2009) set out to improve upon the DWL model of Fairall et al.

196 (1996), and one of the key changes they chose to make was using the 9-band spectral

197 parameterization of Paulson & Simpson (1981), the reason given being that it displays higher

absorption near the surface compared to that of Soloviev (1982). As discussed in Section 1.1, this

parameterization is based on a small dataset acquired using outdated laboratory practices >100

years ago, and there is no justifiable reason for its continued use when far better studies have

been conducted in the interim. In the end, Gentemann et al. (2009) had to multiply the solar

- absorption profile by 1.2 in order to get their model to agree well with observations. The
   parameterization developed in this work provides a much better option for future DWL modeling
- 203 parameterization d 204 efforts.

## 205 **1.3 Purpose of This Work**

Thus far we have discussed how the simple parameterizations of upper ocean radiant heating used in most modern models are plagued by errors and inconsistencies. Indeed, there is no parameterization in use that fulfills these criteria:

- 209 1) computationally efficient enough for large models
- 210 2) accurate within the upper few meters of the ocean
- 3) accounts for the effect of biogenic substances
- 4) explicitly spectral in the photosynthetically active (400:700nm) domain

In this work, we develop a parameterization that fulfills the above criteria. We begin by 213 assembling simple, observationally-validated physical modeling tools for the key controls on 214 215 upper ocean radiant heating, creating an intuitive spectral model. We simplify the spectral model to a useable parameterization that can be directly implemented in existing ocean models, 216 comparing our results to prior parameterizations. We then use observational datasets to inform 217 our understanding of the uncertainties associated with this parameterization, and conclude with a 218 case study demonstrating the impact of chlorophyll concentration on the formation and evolution 219 of diurnal warm layers. 220

## 221 2 The Full Spectral Model

222 The purpose of this section is to develop a spectral model for ocean radiant heating that is conceptually and computationally accessible, and requires a minimum of input parameters. 223 224 While sophisticated radiative transfer modeling tools (see e.g. Mobley, 1994; Ricchiazzi et al., 1998 as a starting point) might enable a more complete description of the process, their 225 complexity presents an unnecessary impediment to our ultimate objective of a simple and 226 227 reliable radiant heating parameterization. By assembling pre-existing, observationally validated models for the solar spectrum at the ocean surface, the broadband albedo, and the spectral 228 underwater attenuation, we create a radiant heating model that is both accurate and accessible, 229 230 needing as inputs only the downwelling broadband irradiance, the chlorophyll concentration, and the time and location on the earth. The spectral model, summarized graphically in Figure 1 and 231 232 as a flow-chart in Figure 2, is both a step towards a useful parameterization, and a valuable tool to help build intuition for the interaction between sunlight and the atmosphere-ocean system. 233



Figure 1: Graphical depiction of spectral model components. (a) Clear-sky spectra from Diffey (2015) at 3 latitudes for an example date & time, differentiated by line style. For the 45N latitude, the modification of the spectrum by the cloud index model of Siegel (1999) is shown with color, demonstrating the preferential transmission of shorter wavelengths through clouds. (b) Broadband albedo from Payne (1972) as a function of atmospheric transmissivity and solar elevation, shown on a logarithmic color scale. (c) Diffuse attenuation coefficients at each wavelength in the solar spectrum, with chlorophyll-dependence in the UV and visible from Morel et al. (2007), and infrared coefficients from Bertie & Lan (1996).

#### 235 **2.1 Clear-Sky Solar Irradiance Spectrum**

To predict the solar irradiance spectrum at the ocean surface under clear skies, we 236 employ the model of Diffey (2015), who deliberately set out to "develop a model that is as 237 simple as it can be commensurate with delivering results of adequate accuracy." By starting from 238 a reference spectrum and accounting for solar angle, direct and diffuse radiation, Rayleigh 239 scattering, aerosol scattering and absorption, and ozone absorption, they demonstrate good 240 agreement with both observations and radiative transfer modeling using nothing more complex 241 242 than an Excel spreadsheet with time and position as inputs. The Excel sheet was carefully translated into Python for use in this study. 243

#### 244 **2.2 Spectral Influence of Clouds**

Having predicted the solar irradiance under a clear sky, we calculate the "Cloud Index" *CI* of the sky in terms of the ratio between observed and predicted broadband irradiances:

$$CI = 1 - \frac{SW_{observed}}{SW_{clear-sky}}$$

where  $SW_{clear-sky}$  comes from integrating the clear-sky spectrum predicted in Section 2.1. We

then employ the empirical model of Siegel et al. (1999) to relate our broadband Cloud Index to a

249 wavelength-specific cloud index, scaling each wavelength in the clear-sky spectrum by a

different amount to accurately reflect the changes in spectral composition caused by clouds
 (namely, preferential transmission of shorter wavelengths). This leaves us with a final estimate

for the solar spectrum at the ocean surface, which integrates to  $SW_{observed}$ .

#### 253 **2.3 Albedo**

The albedo of the sea surface does vary with wavelength, with shorter wavelengths 254 tending to have a slightly higher albedo than red and infrared (Ohlmann et al., 2000, see their 255 Figure 12). However, the magnitude of spectral variability and the absolute magnitude of the 256 albedo are both so small that the spectral effects are second- or even third-order from a radiant 257 258 heating perspective, and thus we do not account for them in this model. Rather, we use the empirical model of Payne (1972), which predicts a broadband albedo as a function of the sun 259 angle and the transmissivity of the atmosphere (equivalent to one minus the Cloud Index). This 260 lookup table was developed from four months of observations at the mouth of Buzzards Bay, 261 Massachussets, and validated against radiative transfer modeling by Ohlmann et al. (2000). The 262 albedo may also display a small wind speed dependence at low sun angles (Katsaros et al., 1985; 263 Payne, 1972) but low sun angles generally correspond to low absolute irradiances and are 264 therefore of minor concern in the context of radiant heating studies outside of the high latitudes. 265

#### 266 **2.4 Underwater Attenuation**

To account for the effect of biogenic substances in the water on the attenuation of UV and visible light, we use the bio-optical parameterization originally published by Morel (1988) and most recently updated in Morel et al. (2007). The parameterization rests on the assumption that at each wavelength, the spectral attenuation coefficient  $K_d$  can be decomposed into two additive parts:

$$K_d = K_w + K_{bio}$$

where  $K_w$  is the attenuation due to the water itself, and  $K_{bio}$  is the combined attenuation due to all biogenic substances in the water (including algal cells, detritus, colored dissolved organic matter, and other associated nonalgal organisms). The "bio-optical assumption" states that the total attenuating effect of all biogenic substances in the water co-varies with chlorophyll in a consistent way; the Morel publications have demonstrated that this assumption is reasonable in open ocean (or "case I") waters, with  $K_{bio}$  following a power-law relationship to chlorophyll concentration:

278 concentration:

$$K_{bio}(Chl,\lambda) = \chi(\lambda)[Chl]^{e(\lambda)}$$

where  $\chi(\lambda)$  and  $e(\lambda)$  are empirical parameters determined by a linear fit to [*Chl*] vs.  $K_{bio}$ observations in natural-log space at each wavelength  $\lambda$ . Morel et al. (2007) only published their parameterization down to 350nm, so we follow the approach of Morel & Antoine (1994) in extrapolating the parameterization between 300:350nm. For  $\lambda >$ 700nm, the attenuation due to the water itself is so strong that biogenic substances have negligible influence. We therefore use the "gold standard" laboratory data from Bertie & Lan (1996) for attenuation coefficients in the

infrared.



*Figure 2: Flowchart demonstrating how the modeling tools depicted in Figure 1 and described in Section 2 can be combined for use as a spectral radiant heating model.* 

#### 287 **3 The Final Parameterization**

Despite its simplicity, the spectral model derived in Section 2 is far too complex to be reasonably integrated into most ocean modeling. From large-scale coupled climate models to 1-D mixed-layer models, most ocean models represent the underwater irradiance profile as the sum of just a few exponentials. Our purpose in this section is therefore to simplify our spectral model to just a few wavelength bands while balancing simplicity and accuracy. To do this, we first simplify the underwater attenuation coefficients, and then determine how to partition the incident irradiance between the resulting spectral bands.

#### 295 **3.1 Simplifying the Model to a Useful Parameterization**

296 In order to simplify the parameters of the bio-optical fits and preserve the chlorophylldependence of the model in the UV and visible, we took logarithmic averages of each parameter 297  $(K_w, \gamma, \text{ and } e)$  within a specified wavelength band. Within the UV band from 300:400nm, we 298 improved the agreement with the spectral model by weighting the average using a typical 299 irradiance spectrum, which helps account for how little light there is in the 300:350nm range 300 compared to the 350:400nm range. For the Visible (PAR) wavelengths from 400:700nm, we 301 sought to divide the spectrum into the minimum number of bands necessary to accurately 302 reproduce the fully spectral output. While just two visible-wavelength bands have been 303 employed in the past (e.g. Manizza et al., 2005), we found that two bands introduced too much 304 vertical variability in the profiles compared to the full spectral model regardless of the location 305 of the split, while three bands (with splits at 510nm and 600nm) reproduce the full spectral 306 model almost perfectly. Parameters for the four resulting wavelength bands are given in Table 1. 307

The infrared wavelengths present a different challenge. While there is no chlorophylldependence that must be preserved through the simplification, the attenuation coefficients span so many orders of magnitude between 700:2500nm that the resulting irradiance profile in the upper meter of the ocean decreases much more rapidly than a simple exponential. In trying to fit the profile as a sum of exponentials, we found that we needed >10 exponential bands to reasonably approximate the profile, which presents a good deal of computational complexity compared to the parameterizations we intend to replace. Rather than accept a poor fit with fewer exponentials, we follow the pioneering example of Zaneveld & Spinrad (1980) in using an

arctangent curve to approximate the rapid decrease near the surface. We fit modeled profiles

317 generated from a wide variety of input conditions to the form:

$$e^{-C_1 z} [1 - C_2 \arctan(C_3 + C_4 z)]$$

and found the fit constants  $C_{1-4}$  (given in Table 1) that best reproduce the spectral model output.

Having defined five wavelength bands and produced reliable descriptions of each band's 319 decay with depth, we must determine what fraction of the incident broadband irradiance should 320 be assigned to each band. In order to do so, we initialized the spectral model with a range of 321 input conditions that span the parameter spaces for sun angle, cloud index, and chlorophyll, and 322 found the constant partitions between the five simplified bands that minimized the absolute 323 difference between the spectral model outputs and the five-band parameterization. The resulting 324 partitions are 6% UV, 49% infrared, and 45% visible (subdivided into 17% blue, 14% yellow, 325 and 14% red). These correspond to a cloud index of  $\sim 0.2$  in the spectral model. 326

Finally, the albedo is generally such a small percentage of the incident irradiance that 327 radiant heating parameterizations have historically treated it as a constant. As can be seen in 328 Figure 1b, the Payne (1972) model yields an albedo of ~0.055 across the majority of relevant 329 conditions, with slightly lower albedos under very clear skies, and significantly higher albedos in 330 conditions of very low sun angles. Because low sun angles correspond to low absolute 331 irradiances, neglecting the variability in albedo yields very small absolute errors. We therefore 332 follow the established literature (e.g. Fairall, Bradley, Rogers, et al., 1996) in setting a constant 333 albedo of 0.055 for our parameterization. The final parameterization is summarized in Table 1. 334

335

Waveband [nm]	Partition (F)	<b>Transmission Profile</b> (multiply by 0.945*SW)	Paramete	rs
300:400	$\begin{array}{c c} \textbf{:400} \\ \textbf{()} \\ ($	$Fe^{-K_d z}$	$K_w = 0$	0.0218
(UV)		where $K_d = K_w + \chi [Chl]^e$	$\chi = 0.1758$	
			<i>e</i> = 0	.6541
400:510	0.17	$Fe^{-K_d z}$	$K_w = 0$	0.0119
(Blue)		where $K_d = K_w + \chi [Chl]^e$	$\chi = 0.1048$	
			e = 0.6330	
510:600	0.14	$Fe^{-K_d z}$	$K_w = 0$	0.0665
(Yellow)		where $K_d = K_w + \chi [Chl]^e$	$\chi = 0.0582$	
			e = 0.5342	
600:700	0.14	$Fe^{-K_d z}$	$K_w = 0$	0.3608
(Red)	0.11	where $K_d = K_w + \chi [Chl]^e$	$\chi = 0.0585$	
			e = 0.4723	
700:2500	0.49	$Fe^{-C_1z}(1-C_2\arctan(C_3+C_4z))$	$C_1 = 1.87$	$C_2 = 0.47$
(IR)			$C_3 = 0.66$	$C_4 = 30$

336 Table 1: The Final Radiant Heating Parameterization.

#### 337 **3.2 Comparison to Existing Parameterizations**

Figure 3 provides a revealing comparison between the new parameterization and several 338 existing parameterizations, in terms of both the irradiance and heating profiles, in both linear and 339 logarithmic depth-space. The parameterization aligns well with the more complex formulation of 340 Ohlmann & Siegel (2000), particularly at high and moderate chlorophyll concentrations. The 341 342 divergence at the lowest chlorophyll concentration is expected, due to our implementation of a more recent iteration of the bio-optical relationship, which is based on the significantly lower 343 (and more accurate) values for the absorption of pure water from Pope & Fry (1997) as 344 implemented in Morel et al. (2007), rather than the Smith & Baker (1981) pure-water spectrum 345 implemented in Morel (1988) and used by Ohlmann & Siegel (2000). In fact, further revisions to 346 the pure-seawater absorption (Lee et al., 2015; Mason et al., 2016) and scattering (Zhang et al., 347 2009) spectra have been published since Morel et al. (2007), with even lower values in the blue 348 wavelengths, but the changes are small enough that they do not yield appreciable differences 349 when integrated for radiant heating calculations. Among the classical parameterizations without 350 Chlorophyll-dependence, we observe the best agreement with Soloviev (1982), particularly near 351 the surface. This reveals why efforts to improve DWL modeling have struggled when switching 352 away from Soloviev (1982), as other options (such as the Paulson & Simpson (1977) & (1981) 353 parameterizations also plotted in Figure 3) perform markedly worse in the upper meters. Indeed, 354 any parameterization built on Paulson & Simpson (1977), including Manizza et al. (2005), will 355 likely perform equally poorly near the surface. 356



**Comparison of Parameterizations** 

Figure 3: Comparison of the parameterization developed in this work to several parameterizations in the literature. The left panels show profiles of irradiance with both linear (top) and logarithmic (bottom) depth scales, while the right panels show the resulting profiles of heating. Ohlmann & Siegel 2000 is plotted in dotted lines colored by chlorophyll concentration, for comparison with our new parameterization plotted in solid lines colored by chlorophyll concentration. The classical parameterizations that do not include chlorophyll dependence are plotted in grey, distinguished by line style.

#### 358 4 Uncertainty & Observations

The simplest approach to estimating the uncertainty associated with our parameterization is to compare it to the outputs of the full spectral model, initialized across a broad range of input conditions. The differences – plotted against depth in Figure 4 – portray an envelope of potential uncertainty, and the color gradient reveals a pronounced Cloud Index dependence, with positive differences at low Cloud Indices and negative differences at high Cloud Indices. Line styles in

- Figure 4 represent different input chlorophyll concentrations; at a given Cloud Index, lower
- chlorophyll concentrations generally correspond to higher absolute differences at depth, simply
- because more light is able to penetrate deeper. However, the parameterization and model agree
- well at a Cloud Index of 0.2, and changes in chlorophyll concentration at this Cloud Index yield no change in the difference between model and parameterization. This implies that the
- chlorophyll-dependence is well captured by the parameterization, while the biggest factor that is
- 370 neglected in the parameterization is the spectral influence of clouds. The comparisons in Figure 4
- would suggest an uncertainty estimate of  $\pm 30 \text{ W/m}^2$  in the upper meters. However, although the
- spectral model is built on observationally-validated tools, we must be cautious in interpreting it
- as an absolutely true representation of real-world behavior. We therefore turn to several types of
- in-situ observations to shed further light on the validity of both the parameterization and the
- 375 spectral model.



Figure 4: Comparison of the simplified parameterization to the full spectral model in the upper 50m across the parameter space of input conditions. The colors reveal that the differences are largely due to the inclusion/exclusion of cloud effects. The lower panels

show the upper 10m split out by wavelength region, revealing the visible region to be the dominant source of variability.

#### 377 4.1 Profiling Multi-Spectral Radiometer Observations

We can test the visible-wavelength portions of both the full model and simplified 378 parameterization using a dataset of 140 profiling spectroradiometer depth casts (Satlantic Profiler 379 II equipped with 2 OCI 507 and 2 OCR 507 radiometers for a total of 14 wavelength channels 380 spanning 380-705nm) with coincident pyranometer and chlorophyll measurements, collected on 381 2 cruises in the tropical Atlantic (September 2015 and August 2016 onboard R/V Meteor), 3 382 cruises in the Gulf of Mexico (June 2015, August 2018 and July 2019 onboard R/V Endeavour), 383 and one cruise in the Tropical Pacific (November 2019 onboard R/V Falkor). The casts were 384 carefully quality-controlled, and only included in the final dataset if the standard deviation of 385 pyranometer measurements during the cast was less than 50  $W/m^2$  (because a single pyranometer 386 measurement and resulting Cloud Index value must be assigned to each cast), resulting in a final 387 dataset of 64 casts. Each spectrally integrated visible irradiance profile observation was then 388 compared to the profiles predicted by both the model and the parameterization. Figure 5 shows 389 the results of the comparison, with the surprising result that the simplified parameterization 390 appears to perform better than the full model at reproducing these observations, particularly at 391 shallow depths. The spread in differences suggests the parameterization is reliable to within 392 about  $\pm 25 \text{ W/m}^2$  in the visible wavelengths, which seem to be the main spectral region of 393 uncertainty based on Figure 4. Given that Cloud Index dependence is the primary difference 394 395 between parameterization and spectral model, the superior performance of the parameterization suggests the Cloud Index dependence of the spectral model may be too strong. We will further 396 investigate this possibility using a dataset of surface hyperspectral radiometer observations. 397





Figure 5: Comparison of integrated visible-wavelength irradiance profiles from 64 multi-spectral depth casts across 6 cruises to both the spectral model (blue) and simplified parameterization (orange). The map shows the locations of the cruises, and

the histograms show the distribution of model input parameters demonstrating the range of conditions sampled.

#### 399 **4.2 Surface Hyperspectral Radiometer Observations**

While onboard R/V Falkor in the Tropical Pacific near Fiji in November 2019, a floating 400 hyperspectral radiometer (Seabird HyperOCR) was deployed behind the ship for 30 minutes near 401 solar noon on 13 separate days to capture the downwelling spectrum in the 350-800nm range at 402 3.3nm resolution. These observations can be used to test the model predictions for partitioning of 403 the three visible wavelength bands as a function of Cloud Index. Because the pyranometer (Kipp 404 & Zonen CMP22) mounted on the ship's mast - which is used to calculate the Cloud Index -405 was several hundred meters away from the floating radiometer, there was a temporal lag in sky 406 conditions that is particularly evident on days with variable cloudiness, which are the most 407 important days for filling out the Cloud Index parameter space. We therefore aligned the 408 radiometer observations with the appropriate Cloud Index estimates by performing a lag 409 correlation between the pyranometer and the (spectrally integrated) radiometer, with the resulting 410 lags ranging between 19 and 72 seconds depending on the wind speed and direction. Figure 6 411 presents a comparison of observations and model output as a function of Cloud Index – note the 412 y-axes are given as fractions of the visible irradiance rather than total solar irradiance, as the 413 observations only span this range completely. Although the model generally predicts close to the 414 right partitions, the observations display minimal variation in the Cloud Index range of 0.0-0.2, 415 supporting the hypothesis that the model is over-sensitive to Cloud Index, particularly in 416 conditions of relatively few clouds.

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Figure 6: Partitioning of the three visible-wavelength bands defined in our parameterization as a function of Cloud Index. Floating hyperspectral radiometer observations are plotted in light grey, and binned by cloud index in blue. Model outputs are plotted in red.

#### 4.3 UAV Albedo Observations 419

While a constant albedo was chosen for simplicity in the final parameterization, the 420

Payne (1972) model displays albedos as high as 50% or more in low sun-angle conditions. For 421

high-latitude applications, therefore, a variable albedo will likely be required. For this reason, we 422

wish to interrogate the accuracy of the Payne (1972) model to determine its suitability for future 423

use. Two Uncrewed Aerial Vehicle (UAV) flights were performed during the same 2019 R/V 424

Falkor cruise with a payload of matched up- and down-looking Hukseflux SR-03 Pyranometers. 425 The downwelling irradiance measurements were corrected for platform motion following the 426 procedure outlined in Reineman et al. (2013; see also Equation 5.1 in Bannehr & Glover, 1991) 427 with a first-order Butterworth low-pass filter cutoff at 1/6 Hz. After correction, the dataset was 428 subset to reject all points with roll or pitch values more than 1 degree from neutral, and then 429 composited into 20-minute averages. Figure 7 shows the albedo observations from both flights 430 compared to the Payne (1972) model (with model input parameters plotted in the lower panels). 431 The observations and model output evolve similarly in time, but the observations are 432 systematically lower than the model output by about 15% on average. This is unsurprising given 433 that the Payne (1972) model is built on observations at the mouth of Buzzard's Bay, 434 435 Massachussets, a turbid coastal environment that does not provide a generalizable analog to open ocean conditions. However, it is encouraging that the temporal evolution of the model follows 436 the observations reasonably well, suggesting that the input parameters of solar elevation and 437 atmospheric transmissivity have been properly identified as the primary controls on albedo 438 variability. More observations of this kind are needed in open-ocean and high-latitude 439 environments to build a more robust version of this simple and potentially very powerful 440 modeling approach. Meanwhile, for studies employing a constant albedo, the limited 441 observations we have thus far suggest that  $\sim 0.045$  (15% lower than the current value) might be a 442 better choice; however, more observations are needed before we advocate for implementing such 443 a change. 444



Figure 7: Albedo as measured directly during two UAV flights (black), compared to output of the Payne (1972) empirical model (red), with the input parameters for the model plotted in the lower panels.

## 446 **5 Case Study: Sensitivity of Diurnal Warm Layers to Chlorophyll Concentration**

The accuracy of our new parameterization near the surface presents an unprecedented 447 opportunity to investigate the sensitivity of DWL formation to chlorophyll concentration. When 448 the COARE3.5 model is forced with shipboard observations from the 2019 R/V Falkor cruise 449 (referenced in previous sections), it predicts a DWL nearly every day. We therefore modified the 450 COARE3.5 model to use the new parameterization, and ran multiple iterations of the model fed 451 with different chlorophyll concentrations, with all other forcing taken from the shipboard 452 observations and consistent between runs. The results, shown in Figure 8, demonstrate the 453 impact that variations in chlorophyll can have on the magnitude and depth of DWLs. Under the 454 same forcing conditions, the range of chlorophyll concentrations likely to be found in the open 455 ocean can modify the strength of the DWL temperature gradient by several tenths of a degree, 456 with higher chlorophyll leading to warmer layers, but significant variability between days (due 457 primarily to the sensitivity of the DWL to wind forcing). Similarly, differences in chlorophyll 458 can lead to changes in the depth of the warm layer of several meters, with higher chlorophyll 459 corresponding to shallower layers. This is a generally intuitive result that more absorbing 460 material in the water should lead to warmer and shallower DWLs as more radiant heating is 461 trapped closer to the surface. These high-chlorophyll warm layers simultaneously have stronger 462 static stability and larger air-sea temperature differences, with the consequence that more of their 463 heat will be returned to the atmosphere via turbulent fluxes. This provides a clear mechanism by 464 which the chlorophyll concentration could modulate the strength of atmospheric convection. 465



Figure 8: Results of modifying the COARE3.5 diurnal warm layer model to use our new parameterization, and varying the input chlorophyll concentration while keeping all other inputs constant. Black lines show the default COARE3.5 model's predictions of warm layer magnitude and depth for the 2019 research cruise onboard R/V Falkor. Colored lines show the differences from the

black lines when the COARE3.5 model is run with the new parameterization at a variety of realistic chlorophyll concentrations.

### 467 6 Conclusion

We have developed the first radiant heating parameterization to offer both chlorophyll 468 dependence and accuracy in the upper meter of the ocean while maintaining a computational 469 simplicity reasonably comparable to its predecessors. Because it explicitly calculates the profile 470 of photosynthetically available radiation (visible light), it can be used in biological and coupled 471 bio-physical modeling as well. Using the new parameterization, we demonstrate the extent to 472 which chlorophyll concentration can affect the depth and magnitude of DWLs, providing 473 mechanistic insight into the possible interactions between phytoplankton and atmospheric 474 convection. Moving forwards, we expect the parameterization - and the discussion of historical 475 inaccuracies provided herein – to enable a needed step forward in accurate bio-physical 476 modeling of upper ocean stratification, dynamics, and mixed-layer depths. We suggest the use of 477 this parameterization in all future multi-level ocean modeling. 478

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# 486 **Open Research**

487 Data used in this study (Witte et al., 2024) is archived for public access on Columbia Academic

488 Commons (<u>https://doi.org/10.7916/wmdm-vm51</u>). The code used to produce the figures in this

study is available for public access on Github (<u>https://github.com/Zappa-Lab/Bio-Physical-</u>

- 490 <u>Radiant-Heating-Parameterization</u>).
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