

Internal tide variability off Central California: multiple sources, seasonality, and eddying background

Tongxin Cai¹, Zhongxiang Zhao¹, Eric A D’Asaro¹, Jinbo Wang², and Lee-Lueng Fu³

¹University of Washington

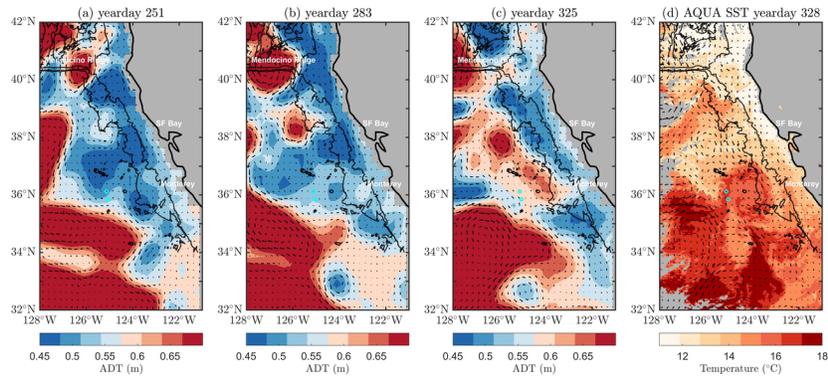
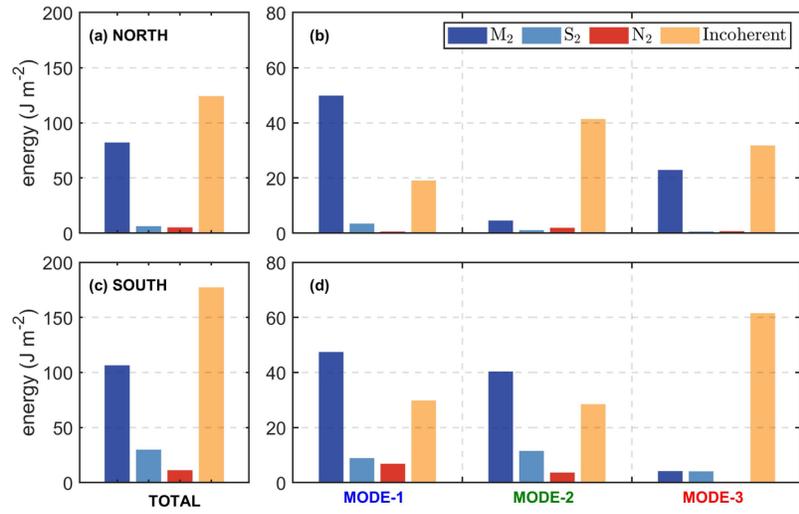
²Scripps Institution of Oceanography

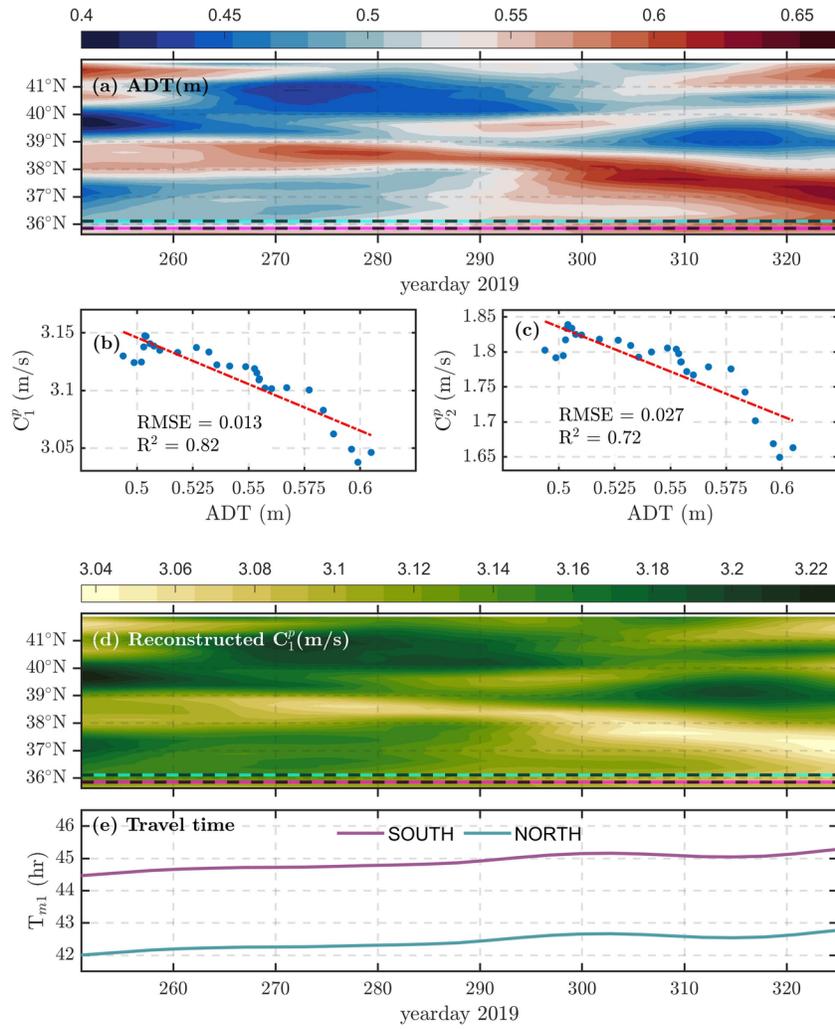
³Jet Propulsion Laboratory

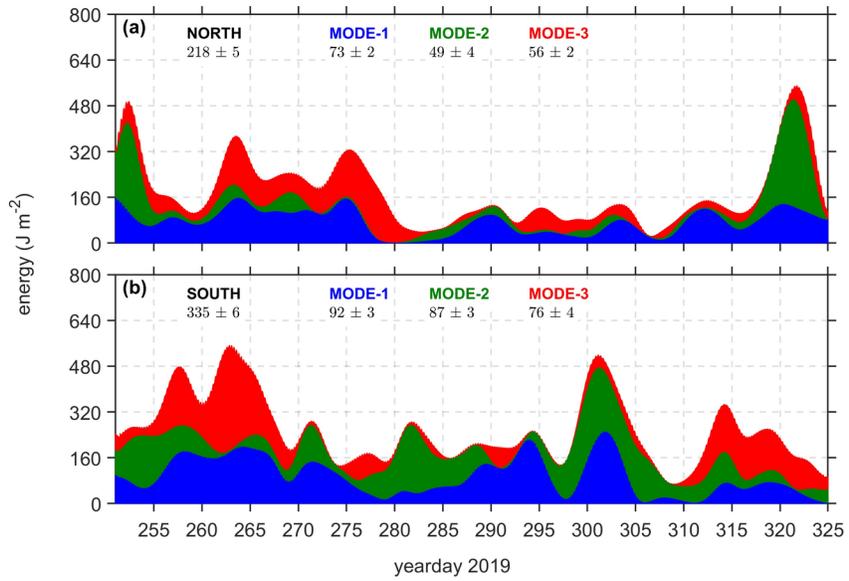
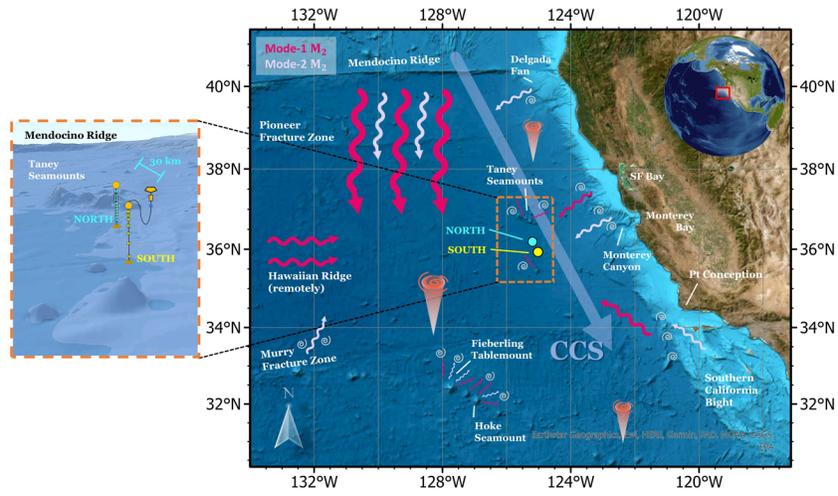
February 26, 2024

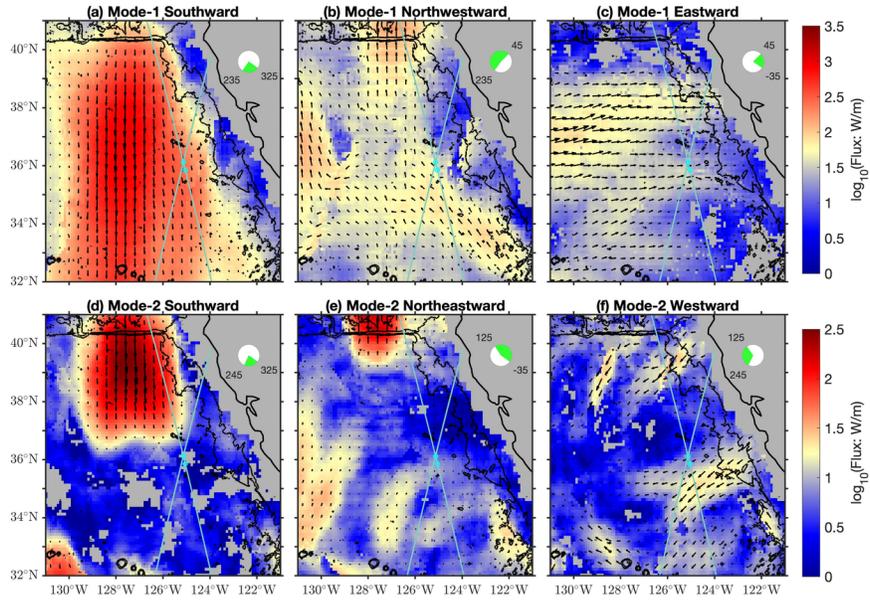
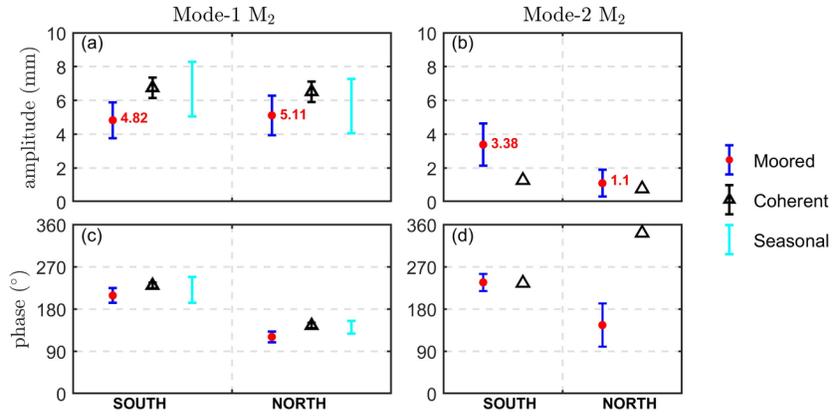
Abstract

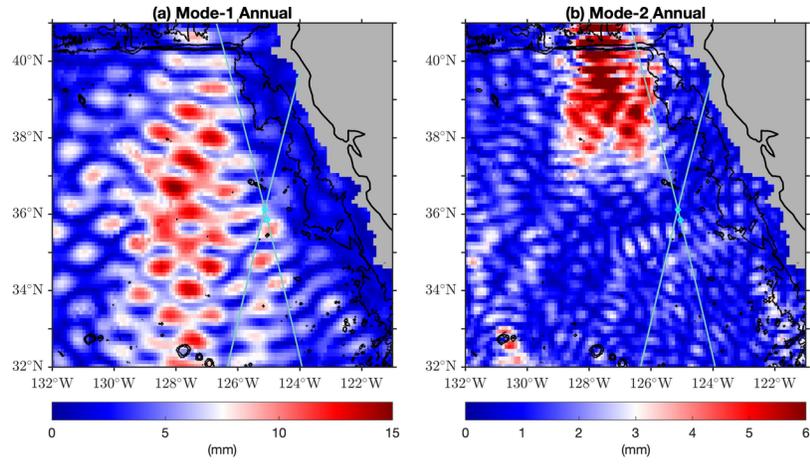
Two moorings deployed for 75 days in 2019 and long-term satellite altimetry data reveal a spatially complex and temporally variable internal tidal field at the SWOT Cal/Val site off central California due to the interference of multiple seasonally-variable sources. Coherent tides account for $\sim 45\%$ of the potential energy. The south mooring exhibits more energetic semidiurnal tides, while the north mooring displays stronger mode-1 M₂ with an amplitude of ~ 5.1 mm. These findings from in situ observations align with the analysis of 27-year altimetry data. The altimetry results indicate that the complex internal tidal field is attributed to multiple sources. Mode-1 tides primarily originate from the Mendocino Ridge and the 36.5° – 37.5° N California continental slope, while mode-2 tides are generated by local seamounts and Monterey Bay. The generation and propagation of these tides are influenced by mesoscale eddies and seasonal stratification. Seasonality is evident for mode-1 waves from three directions. Southward components from the Mendocino Ridge consistently play a dominant role (~ 268 MW) yearlong. We observed the strongest eastward waves during the fall and spring seasons, generated remotely from the Hawaiian Ridge. Westward waves from the 36.5° – 37.5° N California continental slope are weakest during summer, while those from the Southern California Bight are weakest during spring. The highest variability of energy flux is found in the westward waves ($\pm 22\%$), while the lowest is in the southward waves ($\pm 13\%$). These findings emphasize the importance of incorporating the seasonality and spatial variability of internal tides for the SWOT internal tidal correction.

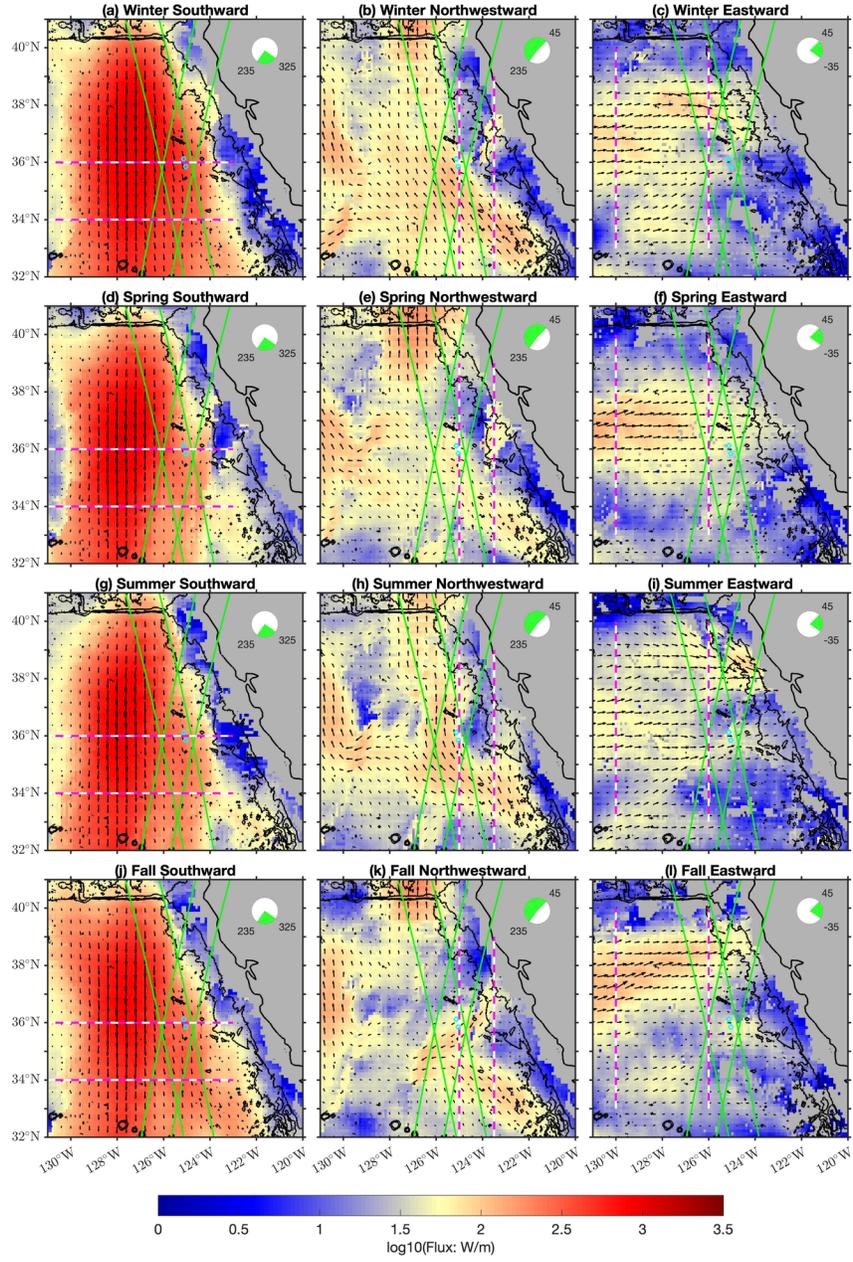


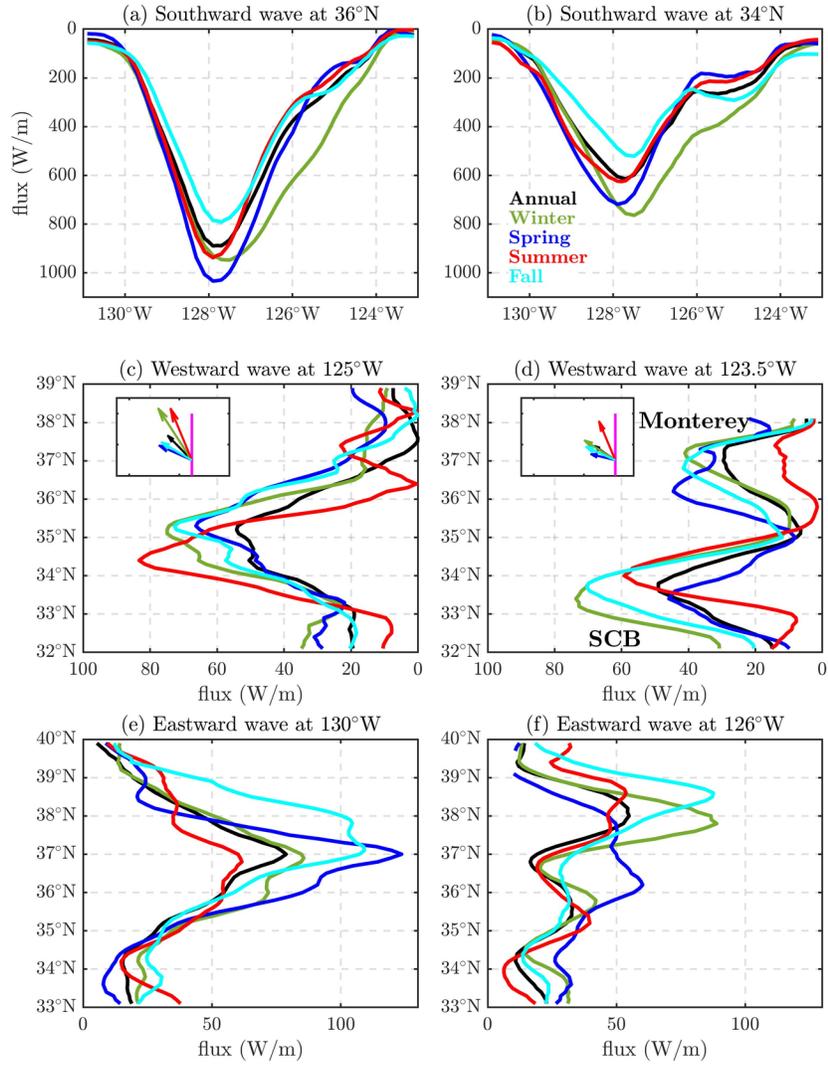


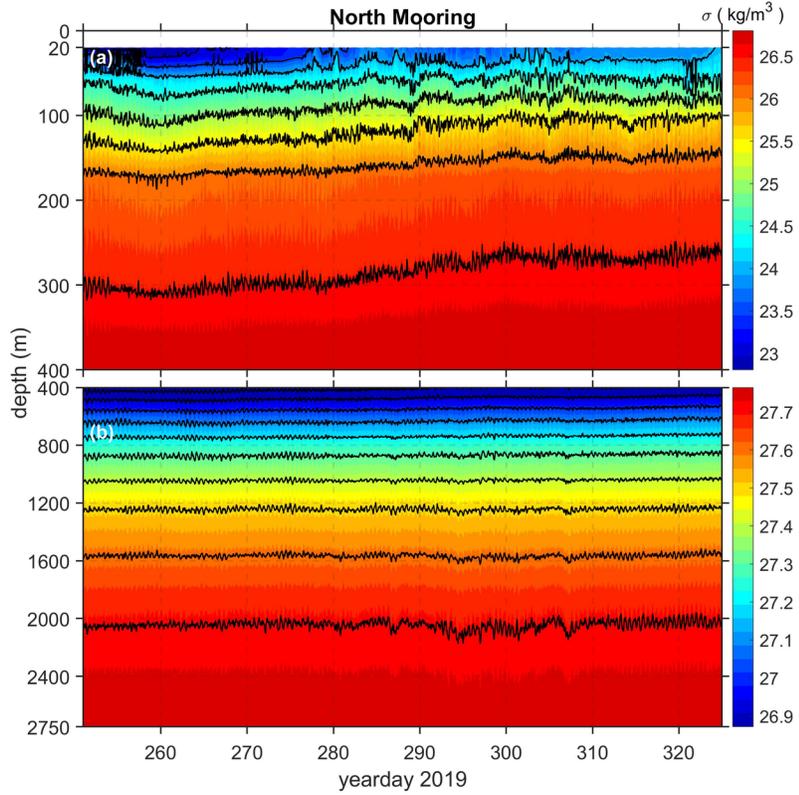
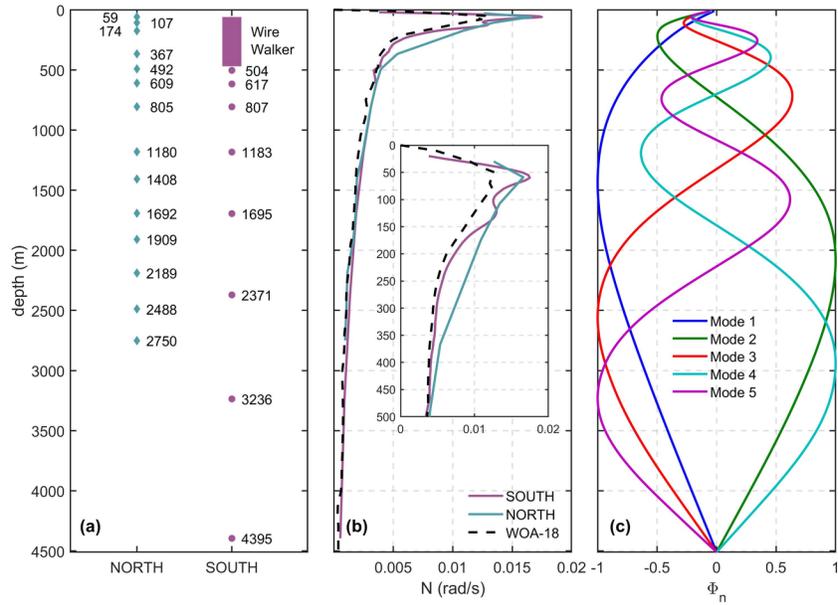


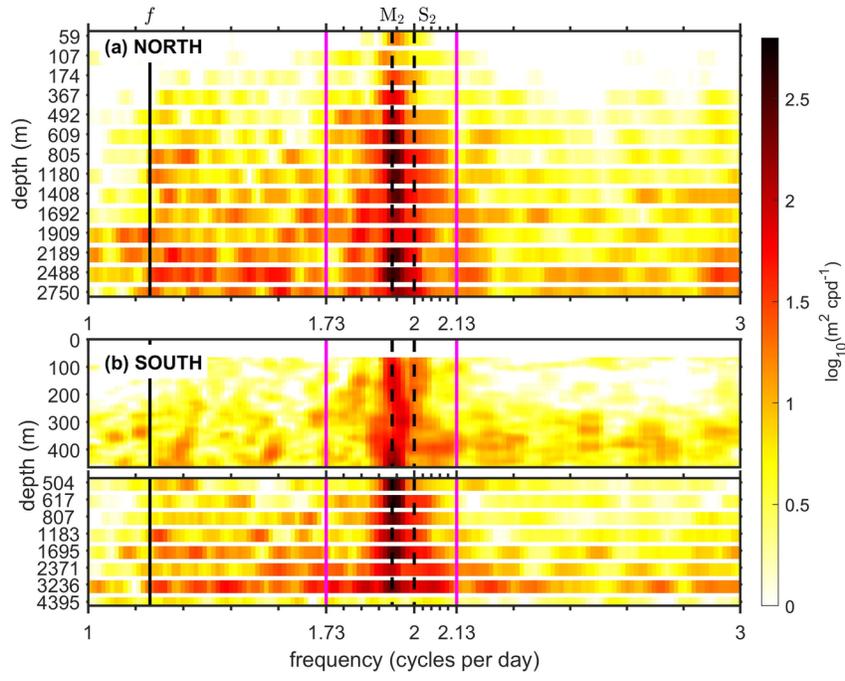












1 **Internal tide variability off Central California: multiple**
2 **sources, seasonality, and eddy background**

3 **Tongxin Cai^{1,2}, Zhongxiang Zhao^{1,2}, Eric D'Asaro^{1,2}, Jinbo Wang³, Lee-Lueng**
4 **Fu³**

5 ¹Applied Physics Laboratory, University of Washington, Seattle, Washington, USA.

6 ²School of Oceanography, University of Washington, Seattle, Washington, USA.

7 ³Jet Propulsion Laboratory, California Institute of Technology, Pasadena, California, USA.

8 **Key Points:**

- 9
- 10 • Temporal and spatial variations of semidiurnal internal tides are observed using
11 in situ moorings and satellite altimetry
 - 12 • Complex internal tide field is caused by multiple generation sources, seasonal strat-
13 ification, and mesoscale eddies
 - 14 • The three generation sources of M_2 internal tides in this region are subject to strong
but different seasonalities

Abstract

Two moorings deployed for 75 days in 2019 and long-term satellite altimetry data reveal a spatially complex and temporally variable internal tidal field at the SWOT Cal/Val site off central California due to the interference of multiple seasonally-variable sources. Coherent tides account for $\sim 45\%$ of the potential energy. The south mooring exhibits more energetic semidiurnal tides, while the north mooring displays stronger mode-1 M_2 with an amplitude of ~ 5.1 mm. These findings from in situ observations align with the analysis of 27-year altimetry data. The altimetry results indicate that the complex internal tidal field is attributed to multiple sources. Mode-1 tides primarily originate from the Mendocino Ridge and the $36.5\text{--}37.5^\circ\text{N}$ California continental slope, while mode-2 tides are generated by local seamounts and Monterey Bay. The generation and propagation of these tides are influenced by mesoscale eddies and seasonal stratification. Seasonality is evident for mode-1 waves from three directions. Southward components from the Mendocino Ridge consistently play a dominant role (~ 268 MW) yearlong. We observed the strongest eastward waves during the fall and spring seasons, generated remotely from the Hawaiian Ridge. Westward waves from the $36.5\text{--}37.5^\circ\text{N}$ California continental slope are weakest during summer, while those from the Southern California Bight are weakest during spring. The highest variability of energy flux is found in the westward waves ($\pm 22\%$), while the lowest is in the southward waves ($\pm 13\%$). These findings emphasize the importance of incorporating the seasonality and spatial variability of internal tides for the SWOT internal tidal correction.

Plain Language Summary

This study explores the variations of internal tides, which are waves at tidal frequencies beneath the ocean surface. They play a crucial role in deep-ocean mixing, ocean circulation, and the overall climate system by transporting nutrients, heat, and carbon within the ocean. Our research area is off central California. We use both in situ measurement and satellite observation to understand how internal tide change over time and space. Our discoveries suggest that five primary sources, changing ocean currents, and seasonal variations of internal tides, contribute to these tidal changes and create the complicated tidal field off central California.

1 Introduction

Investigating the internal tidal field off the U.S. west coast is like peeling an onion. Despite years of collaborative efforts within the research community, there are still many layers to uncover due to its complexity. The complexity, which manifests as temporal and spatial variations, is mainly related to the origins, pathways, and dissipation of internal tides. Previous observations and numerical simulation have identified the Mendocino Ridge (Alford, 2010), continental slope (G. S. Carter et al., 2005; M. Buijsman et al., 2012), local seamounts (Kunze & Toole, 1997), and the Hawaiian Ridge (Zhao, 2019), as the primary sources of these internal tides. After being generated, internal tides in the California Current System (CCS) are subject to the modulation of time-varying mesoscale eddies and background currents (Kurian et al., 2011), leading to temporal variations across various time scales. These combined influences contribute to the intricate nature of the internal tidal field off the U.S. west coast, which poses a significant challenge in unraveling the underlying dynamics. Research on internal tides holds significance for biological production and climate change because the fluctuations of heat, energy, nutrients, and other climatically significant tracers, such as carbon and greenhouse gases, within the ocean interior are influenced by internal tides and the resulting vertical mixing (Sharples et al., 2007; Melet et al., 2022). Here, we analyze the spatial and temporal variations of internal tides off central California using both 3-month moored data and 27-year satellite altimetry observation.

65 There are four primary sources of internal tides off the U.S. west coast. First, the
66 Mendocino Ridge contributes strong internal tides, primarily the M_2 constituent (Althaus
67 et al., 2003; Alford, 2010), which subsequently propagate in a north-south direction (Zhao
68 et al., 2019). Second, internal tides have been identified along the continental slope off
69 Washington State (Alford et al., 2012), Oregon State (Martini et al., 2011), and Cali-
70 fornia State. In California, Monterey Bay (G. S. Carter et al., 2005; Zhao et al., 2012;
71 Terker et al., 2014) and the South California Bight (M. Buijsman et al., 2012; Johnston
72 & Rudnick, 2015) have been focal points of research from both observational and numer-
73 ical perspectives. Additionally, local seamounts, such as Fieberling Tablemount (32.5°N ,
74 127.7°W) and Hoke Seamount (32.1°N , 126.9°W), play a role in internal tide generation
75 (Kunze & Toole, 1997; Zhao, 2018). More recently, satellite altimetry data (Zhao, 2019)
76 have provided evidence of another source of internal tides, demonstrating that far-field
77 internal tides originate remotely from the Hawaiian Ridge. The significance of these re-
78 motely generated tides to regional internal tidal field has been underscored through sim-
79 ulations (Siyabola et al., 2023).

80 Another factor that makes internal tides complicated off the U.S. west coast is their
81 temporal variability. The impact of time-varying stratification and background currents
82 on the generation and propagation of internal tides can occur over different time scales.
83 It can happen over a short period of a few days or on longer time scales such as seasons
84 and years. Interactions with mesoscale eddies and large-scale ocean circulations are sug-
85 gested to be one of the main drivers of the temporal and spatial variation of internal tides
86 (Zaron & Egbert, 2014; Kelly et al., 2016), leading to energy conversion, propagation speed
87 and direction changes, and phase variations of internal tides (Rainville & Pinkel, 2006;
88 Zilberman et al., 2011; Huang et al., 2018). Another influential factor is seasonal strat-
89 ification, which affects internal tides in terms of incoherence, propagation direction, am-
90 plitude, and energy flux (Zhao et al., 2012; Shriver et al., 2014; Ansong et al., 2017). How-
91 ever, seasonal variation of internal tides is not solely attributed to stratification, but also
92 to the seasonality of other ocean processes (Sasaki et al., 2014; Qiu et al., 2014; Zhao,
93 2021). The CCS region exhibits seasonal variations in eddy kinetic energy and mean cur-
94 rent patterns (Haney et al., 2001; Checkley Jr & Barth, 2009; Rudnick et al., 2017). In
95 addition, interference due to wave-wave interaction and the absence of comprehensive
96 4-dimensional observations hinder the way to dynamically link these main drivers to in-
97 ternal tide features (M. C. Buijsman et al., 2017), leading to incomplete understanding
98 of temporal variations of internal tides. Here, using the advanced wave decomposition
99 method (Zhao & Qiu, 2023), we focus on investigating the seasonal variations of inter-
100 nal tides in the presence of mesoscale eddies off California, where a complex internal tidal
101 field is seen from observations.

102 This study is also motivated by the availability of moored observation from the Sur-
103 face Water and Ocean Topography (SWOT) mission pre-launch campaign in 2019 (J. Wang
104 et al., 2022), and the latest advanced satellite altimetry model (Zhao, 2022) that is able
105 to derive the seasonality of internal tides. The global climatological seasonality of inter-
106 nal tides was successfully extracted by subsetting altimetry SSH data into four seasons,
107 leveraging a mapping method that incorporates two key techniques: plane wave anal-
108 ysis and spatial band-pass filtering (Zhao, 2021; Zhao & Qiu, 2023). Characterized by
109 its global coverage and minimal errors, the latest altimetry model enables a global as-
110 sessment of seasonal variations of internal tides while offering a meaningful comparison
111 with in situ observations and numerical simulations. Combining moored observations and
112 satellite altimetry offers a unique perspective on internal tide in a complex ocean envi-
113 ronment (Köhler et al., 2019; Löb et al., 2020). In our study region, where multisource
114 internal tide interference patterns are present (Rainville et al., 2010), a comprehensive
115 understanding of internal tides necessitates the complementary use of moored and al-
116 timetry data.

117 The goals of this study are as follows: (1) to reveal the temporal and spatial vari-
 118 ations of internal tides from moored observations; (2) to evaluate the performance and
 119 reliability of the 27-year-coherent altimetry model; (3) to elaborate the distinct charac-
 120 teristics of mode-1 and mode-2 internal tides in the CCS region, considering the influ-
 121 ence of multiple sources and eddying background; and (4) to explore the seasonality of
 122 mode-1 tides. Specifically, we examine the temporal and spatial variations in modal com-
 123 position and coherence of the semidiurnal internal tide. Our analysis primarily focuses
 124 on the mode-1 and mode-2 M_2 internal tides observed by moorings and compares them
 125 with 27-year satellite altimetry observations. Furthermore, we delve into the contribu-
 126 tion of waves from each direction and consider the influence of mesoscale eddies. Lastly,
 127 we explore the seasonality of mode-1 tides in each direction using the latest seasonal al-
 128 timetry models. Through this study, our aim is to provide a comprehensive understand-
 129 ing of the characteristics of semidiurnal internal tides in the CCS region.

130 This paper is organized as follows: Section 2 provides an overview of the data from
 131 the SWOT mission pre-launch campaign and processing methods. Section 3 introduces
 132 the satellite altimetry data and presents two key techniques utilized to extract the in-
 133 ternal tidal signal. Section 4 presents the findings from the moored observations and in-
 134 cludes a comparison with satellite observations discussed in Section 5. Additionally, Sec-
 135 tion 5 delves into the distinct generation and propagation characteristics of mode-1 and
 136 mode-2 M_2 tides. The seasonality of mode-1 tides in the CCS region is further explored
 137 in Section 6. Finally, in Section 7, we summarize the results and draw our conclusions.

138 2 SWOT Pre-Launch Field Campaign

139 2.1 Field Campaign

140 The SWOT Calibration/Validation (Cal/Val) pre-launch field campaign was car-
 141 ried out in a region located about 300 km west of Monterey, California, from Septem-
 142 ber 2019 until January 2020 (Figure 1). Three moorings were deployed along one Sentinel-
 143 3A satellite track in the CCS region. Two of these three moorings are studied in this work;
 144 the PMEL/WHOI mooring was at 125.13°W , 36.12°N (hereinafter as “the north moor-
 145 ing” based on latitudinal position) and the Scripps Institution of Oceanography (SIO)
 146 mooring was at 125.05°W , 35.85°N (hereinafter as “the south mooring”). Both moorings
 147 provide hydrographic temperature and salinity measurements, bottom pressure from bot-
 148 tom pressure recorders (BPRs). In addition, the surface buoy on the north mooring was
 149 equipped with a GPS sensor measuring the true Sea Surface Height.

150 This study uses data from salinity, temperature, and pressure instruments on the
 151 north mooring and the south mooring. The north mooring has 18 fixed CTD (Conduc-
 152 tivity, Temperature, and Depth) sensors located unevenly throughout the ocean column,
 153 measuring temperature, salinity, and conductivity with a sample interval of 1 minute.
 154 The south mooring has a Wirewalker profiler equipped with Sea-Bird Electronics SBE37-
 155 IM and RBR Concerto, which crawls up and down along the mooring wire from the sur-
 156 face to about 500 m. They provide temperature and salinity measurements of the wa-
 157 ter column with a vertical resolution of about 29 m and deliver one up or down profile
 158 every 18.6 minutes on average. Below 500-m depth, 8 fixed CTD are positioned unevenly
 159 towards the bottom with a sampling interval of 10 minutes. Figure 2a shows details of
 160 the instrument arrangements for both moorings. Further information on the data can
 161 be found in J. Wang et al. (2022).

162 2.2 Data Processing

163 The data obtained from 14 fixed CTDs at the north mooring are utilized after the
 164 quality control (see details in the Supporting Information). A consistent time period of
 165 75 days, from yearday 251 (9 September 2019) to yearday 325 (22 November 2019), is

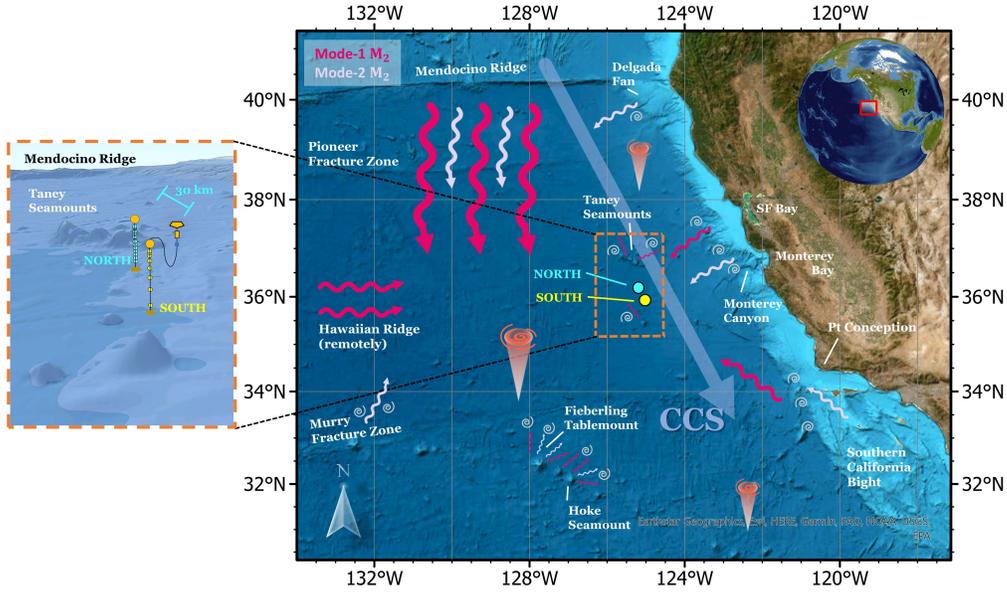


Figure 1: Study region in the California Current System (CCS). Bathymetry is mapped using data from GEBCO, with key topographic features labeled. Two moorings deployed during the SWOT pre-launch campaign 2019 are labeled in cyan (the north mooring) and yellow (the south mooring). Major sources of mode-1 (red) and mode-2 (purple) M_2 internal tides in this region are marked as curved arrows, based on the 27-year-coherent satellite altimetry model. The figure in the left of an orange box is the zoom-in view of the two moorings 30 km apart.

166 chosen for both moorings. In this paper, yearday 251 is 00:00 UTC on 9 September 2019.
 167 To facilitate comparison, the data from the upper 500 m at the south mooring are grid-
 168 ded onto a uniform 1-hr temporal and 5-m vertical grid using linear interpolation. The
 169 vertical displacement is calculated by determining the potential density anomaly (σ) from
 170 temperature and salinity data using the Gibbs Sea Water Function and the Thermody-
 171 namic Equation of Seawater 2010 software (McDougall & Barker, 2011). Figure 3 illus-
 172 trates the time series of potential density at the north mooring, while a similar figure
 173 for the south mooring can be found in the Supporting Information.

174 The ocean conditions can be obtained by calculating the buoyancy frequency (N^2)
 175 using a CTD-profile created from the *World Ocean Atlas 2018* (WOA18) (Zweng et al.,
 176 2019). The dashed line in Figure 2b represents the stratification profile obtained from
 177 a climatological analysis. This profile aligns with the measurements from the two moor-
 178 ings (represented by solid lines), with the exception of the upper 500 m (as shown in the
 179 close-up view). In order to preserve the seasonal (fall) ocean condition in the CTD-profile,
 180 we use data acquired from the Wirewalker Profiler in the south mooring for the upper
 181 500 m while data from WOA18 are employed for the remaining depth.

182 2.3 Vertical Displacement and its Frequency Spectra

183 Displacement of isopycnal η_σ is computed by potential density profiles via the re-
 184 lation

$$\eta_\sigma(z, t) = \frac{\sigma'(z, t)}{\frac{d\sigma}{dz}} \quad (1)$$

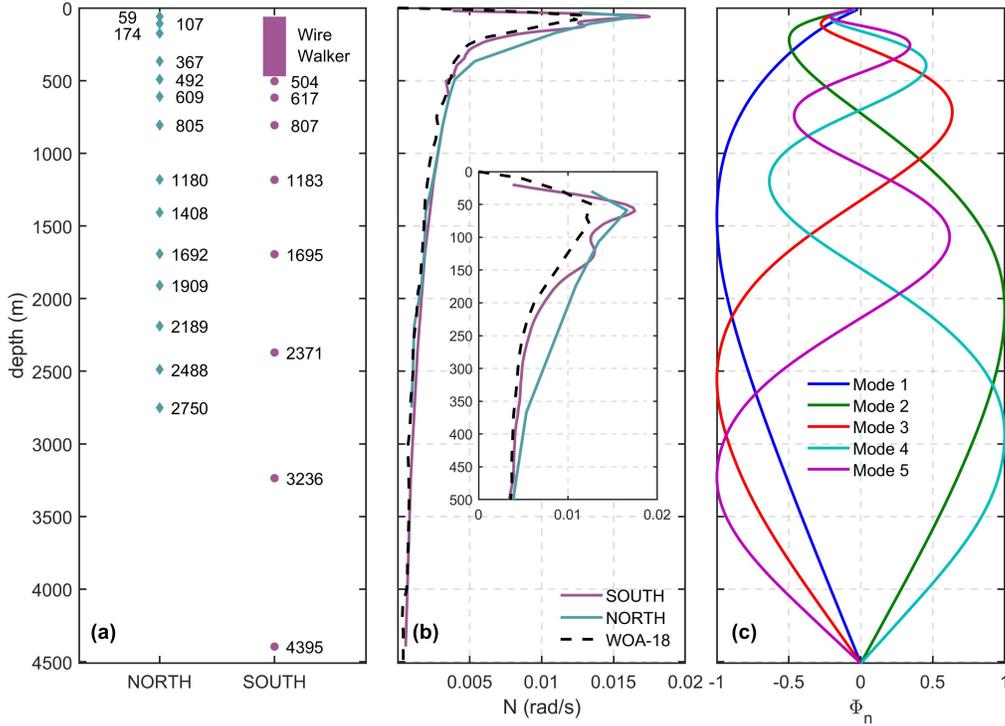


Figure 2: Mooring instrumentation and ocean stratification profiles. (a) The north mooring instruments (cyan diamonds as fixed CTDs) and the south mooring instruments (magenta asterisks as fixed CTDs and the box as the Wirewalker Profiler). Depth (m) of each fixed CTD is provided. (b) Brunt-Väisälä frequency N profiles (in rad/s). The solid lines are mooring measurements (cyan for the north mooring and magenta for the south mooring). The black dashed line is the WOA18 annual mean hydrographic data. The close-up view is of the upper 500 m. (c) Normalized vertical structure of the first five baroclinic modes (in colors) of internal tides for vertical displacement.

185 The gradient of potential density $\frac{d\sigma}{dz}$ is from CTD-profile and the perturbation $\sigma'(z, t) =$
 186 $\sigma(z, t) - \bar{\sigma}(z)$, where $\sigma(z, t)$ is the instantaneous density anomaly and $\bar{\sigma}(z)$ is the time
 187 mean of the potential density anomaly profile. We adjust the displacement by remov-
 188 ing the components of pressure variations arising from the mooring design (see details
 189 in the Supporting Information). The corrected data are consistent with those from J. Wang
 190 et al. (2022).

191 Figure 4 shows the spectrum of η_{tide} as a function of depth for the two moorings
 192 below the mixed layer in the upper ocean. The spectra are computed using a sine mul-
 193 titaper method (Thomson, 1982) with two sine tapers giving a degree of freedom (DOF)
 194 of 4. A smoothing process is applied to geometrically smooth the spectrum over 1/250
 195 of the total bandwidth. The resulting spectrum resolution is 0.0135 cycles per day. Strong
 196 semidiurnal tidal signals are apparent for both moorings. The nonlinearity of internal
 197 tides is indicated by the presence of overtones (i.e. M_4 in here) and some near-inertial
 198 motion (f) is also seen. The 95% and 50% confidence intervals of the spectrum can be
 199 referred to Supporting Information.

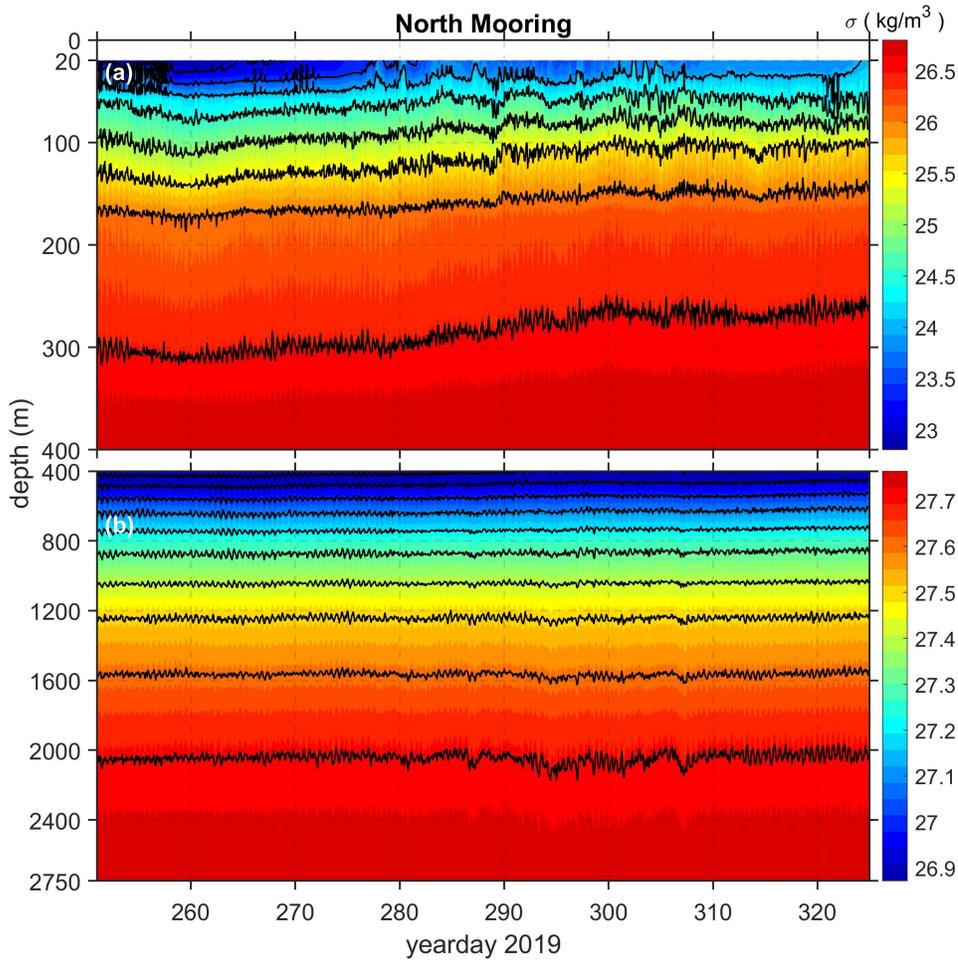


Figure 3: Hourly gridded potential density anomaly σ at the north mooring (a) at upper 400 m and (b) 400–2750 m. Black contour lines are isopycnals with a constant density value. Note that there are different ranges of colormap for (a) and (b).

2.4 Band-pass Filtering and Harmonic Analysis

The temporal variability of each tidal component is examined by band-pass filtering via fourth-order Butterworth and harmonic analysis. This passing band includes M_2 and S_2 tidal constituents and is referred to Zhao et al. (2010)’s set up, which is centered at the M_2 tidal frequency ($2.23 \times 10^{-5} \text{s}^{-1}$) with zero-phase response and quarter-power points at $2.01 \times 10^{-5} \text{s}^{-1}$ and $2.47 \times 10^{-5} \text{s}^{-1}$, i.e., 1.73–2.13 cpd. These frequency limits are wide enough to capture the majority of semidiurnal signals but narrow enough to separate them from other nontidal motions. The available data record is long enough to perform this filtering without being concerned with leakage or ringing, which are artifacts that can occur in the filtered signal due to the finite length of the data record and the characteristics of the filter.

The band passed semidiurnal signals are a combination of M_2 , S_2 , N_2 and incoherent constituents. The 75-day data record is long enough to separate M_2 , S_2 , and N_2 so

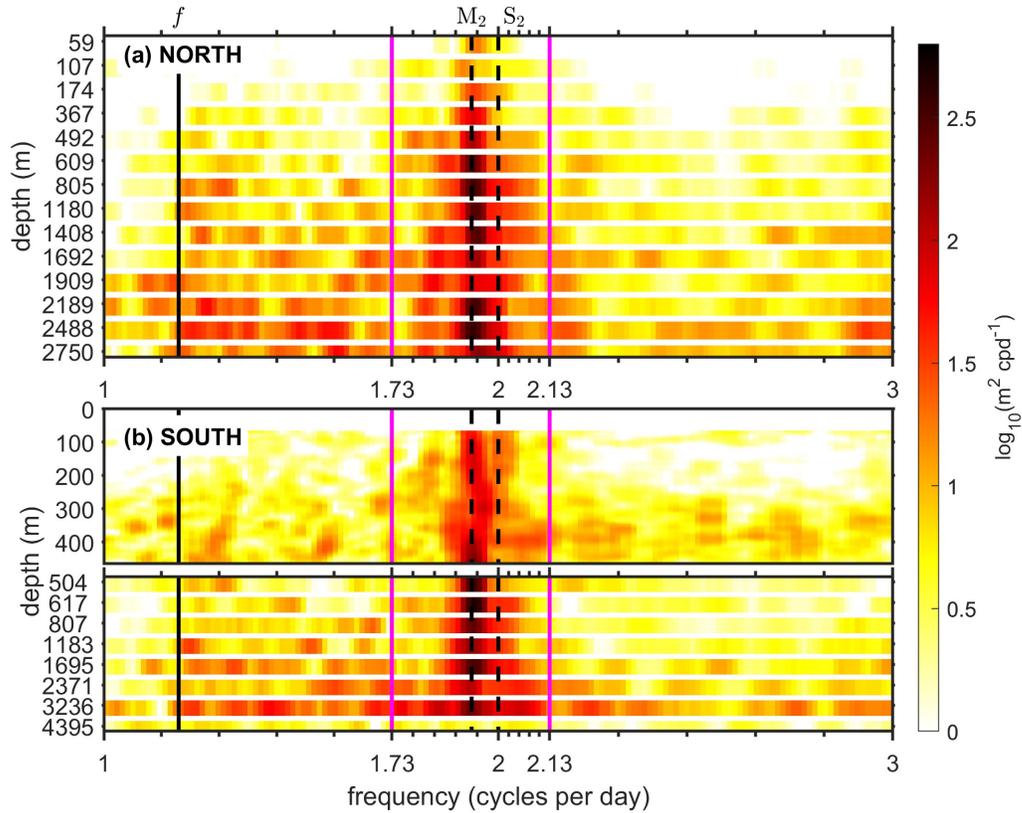


Figure 4: Spectrum of tidal displacement of (a) the north mooring and (b) the south mooring at depth of every CTD sensor. The spectra are calculated using a sine multitaper method giving a degree of freedom (DOF) of 2. A smoothing process is applied to geometrically smooth a spectrum, covering over 1/250 of the total bandwidth. Major frequency are labeled: M_2 and S_2 are as dashed black lines, inertial frequencies is as a solid black line, and band-pass limits are as solid magenta lines.

213 that they can be extracted by harmonic analysis (Pawlowicz et al., 2002). As coherent
 214 M_2 , S_2 , and N_2 signals dominate, K_2 constituent is neglected. The baroclinic vertical
 215 displacement is expressed as

$$\eta'_{semi} = \eta'_{M_2} + \eta'_{S_2} + \eta'_{N_2} + \eta'_{in} \quad (2)$$

216 where η'_{in} indicates the incoherent portion.

217 2.5 Modal Decomposition

218 Internal tides can be described by a superposition of discrete baroclinic modes that,
 219 for horizontally uniform $N(z)$ and no background shear, propagate as linear waves. There-
 220 fore, to analyze the semidiurnal vertical displacement within the chosen frequency band,
 221 displacement is projected onto these baroclinic modes. As described by Zhao et al. (2016),
 222 the baroclinic modes for vertical displacement, $\Phi(z)$, are calculated by the eigenvalue equa-
 223 tion (Wunsch, 1975; Munk, 1981),

$$\frac{d^2\Phi(z)}{dz^2} + \frac{N^2(z)}{c_n^2}\Phi(z) = 0 \quad (3)$$

224 $\Phi(0) = \Phi(-H) = 0$ are rigid-lid boundary conditions in location with depth H on a
 225 flat bottom. Subscript n is the vertical normal mode number and c_n is the eigenvalue
 226 velocity (Gill & Adrian, 1982). $N(z)$ is taken from the CTD-profile. The energy estimates
 227 can be severely limited by vertical gaps in the measurements, but it is possible to rep-
 228 resent internal tides by combining several distinct baroclinic modes (Nash et al., 2005;
 229 Zhao et al., 2012). The water column coverage is sufficient to compute the lowest five
 230 vertical modes, as shown in the Supporting Information.

231 After computing five-mode solutions for both moorings, the baroclinic displacement
 232 is expressed as

$$\eta'(z, t) = \sum_{n=1}^5 \eta'_n(t)\Phi_n(z) \quad (4)$$

233 where $\Phi_n(z)$ represents the vertical structure of the n th baroclinic mode and $\eta'_n(t)$ is the
 234 time-varying displacement of the n th baroclinic mode. At each time, $\eta'_n(t)$ is determined
 235 by least squares modal fitting.

236 Depth-integrated available potential energy (APE) is determined by the baroclinic
 237 displacement $\eta'(z, t)$

$$APE = \frac{1}{2}\rho_0 \int_{-H}^0 \langle N^2(z)\eta'^2(z, t) \rangle dz \quad (5)$$

238 with the unit of J/m^2 , where the angle brackets are the average over one tidal cycle, ρ_0
 239 is the vertically averaged water potential density, and $N(z)$ is the buoyancy frequency
 240 from the CTD-profile. Horizontal kinetic energies (HKE) and flux (F) are unavailable
 241 due to a lack of moored measurement of baroclinic current velocity $\mathbf{u}(z)$.

242 In order to compare with satellite altimetry, the sea surface height anomalies (SSHAs)
 243 are calculated with interior isopycnal displacement for each mode η'_n derived from above,
 244 which can be expressed as

$$SSHA_n = \kappa\eta'_n(t) \quad (6)$$

245 which κ is the conversion ratio depending on latitude, mode number, and frequency. $\kappa =$
 246 1.1×10^{-3} for M_2 mode-1 tide and $\kappa = 0.7 \times 10^{-3}$ for M_2 mode-2 tide in this site. For
 247 convenience, SSHAs are then converted from meters to millimeters.

248 3 Satellite Altimetry Model

249 Two kinds of satellite altimetry models are used in this study: the 27-year-coherent
 250 model and the climatologically seasonal model.

3.1 Satellite Altimetry data

Following the new mapping technique described in Zhao and Qiu (2023), the regional M_2 internal tidal field is mapped using 27 years (1993-2019) of satellite data from multiple altimetry missions. The sea surface height (SSH) data from seven exact-repeated satellite missions are combined into four data sets based on their orbital configurations, including 254 tracks from TPJ (TOPEX/Poseidon-Jason), 254 tracks from TPT (TOPEX/Poseidon-Jason tandem), 1002 tracks from ERS (*European Remote Sensing Satellite-2*), and 488 tracks from GFO (*Geosat Follow-On*). The merged data sets have denser ground tracks and higher spatial resolution compared to each individual mission with sparse tracks, enabling the development of an accurate internal tide model. Previous studies (Zhao, 2021; Zhao & Qiu, 2023) used the same data, except with a 25-year (1993-2017) altimetry record. Standard corrections are applied to all SSH measurements to address atmospheric effects, surface wave bias, and geophysical effects. The corrections for the ocean barotropic tide, polar tide, solid Earth tide, and loading tide are conducted using theoretical or empirical models. A high-pass filter with a cutoff wavelength of 2000 km is used for along-track filtering to remove mesoscale motions.

3.2 Mapping Procedure and Techniques

Two key techniques, plane wave analysis and 2D spatial filtering, are applied to the mode-1 and mode-2 M_2 mapping procedures. Instead of point-wise harmonic analysis, plane wave analysis (Zhao et al., 2016; Zhao, 2016) extracts internal tides by fitting plane waves using all altimetry measurements in one given fitting window that is 160 km in width. In overlapping fitting windows, least squares fitting is used to calculate the amplitudes a , phases ϕ , and propagation directions θ of the target internal tidal waves, following

$$\eta(x, y, z) = \sum_{m=1}^M a_m \cos(kx \cos(\theta_m) + ky \sin(\theta_m) - \omega t - \phi_m) \quad (7)$$

where ω and k are the frequency and wavenumber of M_2 , x and y are the local Cartesian coordinates, and t is the time. M is the number of internal waves extracted in each window via an iterative algorithm. Five waves are fitted for both mode-1 mode-2. Then M_2 internal tides are mapped at regular spatial grids.

2D spatial filtering aims to remove higher baroclinic modes and nontidal noise by employing a horizontal band-pass filter. The filter has a bandwidth of [0.8 1.25] times the regional mean wavelength, which is tested empirically with several values. For this method to work effectively, it is crucial that the variance of internal tides is mostly around the theoretical wavenumber (Zhao et al., 2019) and the bandwidth is as narrow as possible without eliminating the real signals. The wavelength (wavenumber) of M_2 internal tides depends on factors such as ocean depth, latitude, mode number, and ocean stratification. In Section 3.3, we will address the determination of this prerequisite parameter, with particular emphasis on accounting for seasonal variation.

The 27-year-coherent internal tide model is constructed following the mapping procedure described in Zhao and Qiu (2023), which involves three steps: (1) plane wave analysis to map internal tides at a 160 km \times 160 km window with 5 waves, (2) 2D spatial filtering to clean internal tides based on wavenumber, (3) multidirectional decomposition using plane wave analysis within the same window as step (1) to separate tidal waves by propagation directions. In the end, the internal tidal field is mapped on the grid of $0.1^\circ \times 0.1^\circ$ for mode-1 and $0.05^\circ \times 0.05^\circ$ for mode-2. This new mapping method significantly reduced model error and has been compared and assessed with an independent data set from *CryoSat-2*. The resultant tidal models exhibit minimal error, making it possible to resolve weak seasonal signals of internal tides from different propagating directions.

3.3 Seasonal Data Subsetting

The climatologically seasonal internal tide models are built with four seasonal subsets of altimetry data and WOA18 climatologies, following the method from Zhao (2021). The four seasonal subsets consist of January, February, and March for the winter model, April, May, and June for the spring model, July, August, and September for the summer model, and October, November, and December for the fall model. The seasonal models are developed following the same mapping procedure as the 27-year-coherent one, but with the respective data subset. Zhao (2021) employed this approach to study the seasonality of M_2 mode-1 internal tides.

To consider the seasonal variations from the altimetry models, the M_2 wavelength (wavenumber), one of the prerequisite parameters, is calculated for the four seasons using the ocean stratification profiles from the WOA18 climatological seasonal hydrography. At each $0.25^\circ \times 0.25^\circ$ grid point of the WOA18 data set, the vertical structure and wavelengths are determined by solving the Sturm-Liouville orthogonal equation (3) and $\lambda = \frac{c_p^n}{\omega}$. The largest mode-1 M_2 internal tides are our focus for seasonality analysis.

3.4 Energetics

The depth-integrated energy flux can be calculated from the satellite-derived SSHAs following

$$Flux = \frac{1}{2} a^2 F_n(\omega, H, f, N) \quad (8)$$

where a is the SSH amplitude. This equation (Zhao et al., 2016; Zhao, 2018) involves the transfer function F_n , which is the other prerequisite parameter dependent on the frequency ω , water depth H , local inertial frequency f , and stratification N . The transfer function is derived using the hydrographic profiles from the WOA18 data set. Since there are five waves at each grid point, the total values we discuss later are the scalar (energy) and vector (flux) sums of these waves.

4 Mooring Observations

In this section, we will present the observed time-varying internal tide energy of different modes and constituents at the two moorings. Our results indicate that (1) there are significant temporal and spatial variations of internal tides in the region; (2) the south mooring has a greater semidiurnal tidal energy, while the north mooring has a higher amplitude of M_2 mode-1 internal tide; (3) mode-1 tides covary at the two moorings, while mode-2 tides are weakly correlated; and (4) the deceleration of phase velocity may be associated with the formation of a warm-core anticyclone.

4.1 Time Series

To evaluate the temporal variations of internal tides at each mooring and their spatial disparities, we compute the vertical-integrated available potential energy (APE) from baroclinic displacement η in mode 1-3 using Equation (5). In addition, we calculate the time-mean total energy of the lowest-three modes by summing up the time-averaged energy in each mode, with a 95% confidence interval provided (Figure 5).

The time-averaged energy in each mode at the south mooring is higher than at the north mooring. At the north mooring (Figure 5a), the energy in the lowest-three modes is 218 ± 5 J/m². Contrary to the expected case described by de Lavergne et al. (2019), which suggests a strong decay of both energy and conversion rate with increasing mode number, we find that mode-3 tide (56 ± 2 J/m²) and mode-2 tide (49 ± 4 J/m²) are of similar magnitude. We acknowledge that there are uncertainties in estimating the modal contribution due to observations characterized by incomplete vertical spatial coverage.

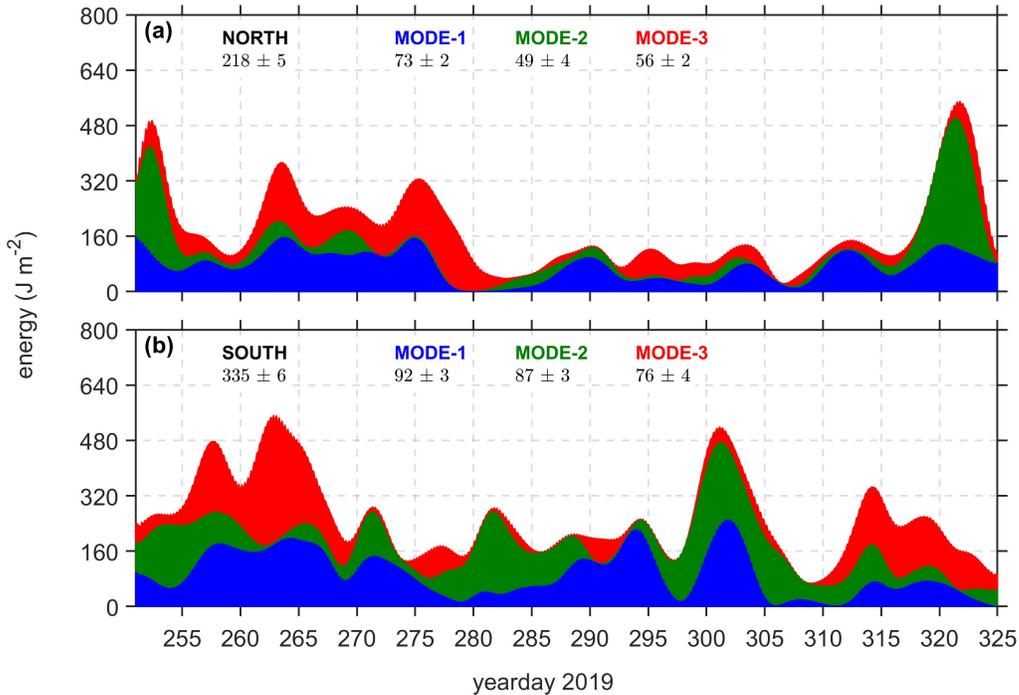


Figure 5: Time series of semidiurnal internal tide vertically integrated available potential energy (APE, J m^{-2}) in modes 1–3 (stacked colors) at (a) the north mooring and (b) the south mooring. The time-averaged energies in modes 1–3 and in total are given. The 95% confidence interval is listed behind each value. Temporal variations of semidiurnal internal tides are seen from both moorings.

344 Therefore, we focus only on mode-1 and mode-2 internal tides here (represented by blue
 345 and green colors in Figure 5). Overall, it is clear that the majority of the measured energy
 346 is contained in low-mode tides (i.e., mode 1-3 with 81%). At the south mooring (Fig-
 347 ure 5b), the energy of the total lowest-three modes is $335 \pm 6 \text{ J/m}^2$ and the energy de-
 348 creases as mode number increases. These variations of energy for dynamics over a sep-
 349 aration scale of $\mathcal{O}(30)$ km between the two moorings indicate a spatially complex inter-
 350 nal tidal field in this region.

351 At both moorings, the internal tides have significant temporal variations. At the
 352 north mooring, there are specific periods, such as those spanning yeardays 257-265 and
 353 yeardays 315-325, exhibit synchronized changes among different tidal modes. Conversely,
 354 during other periods like yeardays 266-272 and yeardays 277-285, tides in different modes
 355 manifest incoherent behavior, signifying a lack of consistent temporal alignment. Even
 356 when the changes in different tidal modes align, these temporal changes are not neces-
 357 sarily in proportion. For instance, despite mode-1 predominates over the whole period,
 358 mode-2 (green in Figure 5a) get excited during yearday 315-325, which could be attributed
 359 to fluctuations in the background currents and eddies. Substantial variations in the en-
 360 ergy time series are also evident at the south mooring. During certain periods, such as
 361 yearday 268-272 and yearday 295-305, there is consistency in how energy changes in dif-
 362 ferent modes. However, overall, energy variations in different modes often do not follow
 363 a coherent or synchronized pattern, indicating temporal incoherence. We did not see an
 364 obvious spring-neap cycle of semidiurnal tides from the time series of both moorings, which

365 is likely due to the extremely weak S_2 tide. According to satellite observations (see Section
 366 5), S_2 is associated with an SSH signal of ~ 2 mm, while M_2 signal is ~ 10 mm
 367 in this region.

368 Mode-1 tides covary at the two moorings while mode-2 tides are weakly correlated.
 369 Mode-1 tides, for example, weaken around yearday 280 and get stronger afterward for
 370 both moorings. In contrast, the peak of mode-2 tides from the north mooring at around
 371 yearday 322 is not seen from the south mooring. If the observed tides from these two
 372 moorings were only from the Mendocino Ridge in the north, we would not expect to see
 373 such significant spatial differences, especially for mode-2 tides. Therefore, we argue that
 374 these spatial variations are contributed by multiwave interference and different gener-
 375 ation sites for mode-1 and mode-2 tides. This hypothesis will be verified by the inter-
 376 nal tidal field from satellite observations in the next section.

377 4.2 Tidal Constituents

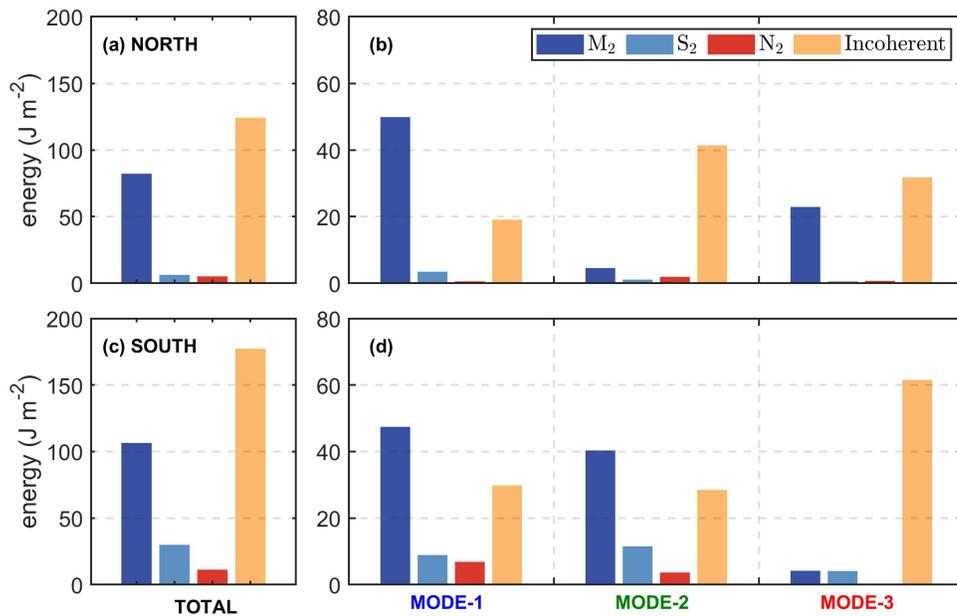


Figure 6: (a) Partition of energy by tidal constituents at the north mooring. Modal decomposition is applied to each tidal constituent and (b) shows the partition on the lowest-three modes (in the x axis). The same analysis for the south mooring is presented in (c) and (d).

378 We employ harmonic analysis to assess the energy of different semidiurnal tidal con-
 379 stituents, including M_2 , S_2 , and N_2 . The coherent and incoherent portions are defined
 380 in Section 2.4. The coherence of internal tides varies with different modes due to their
 381 unique vertical structure and propagation velocity (Rainville & Pinkel, 2006; Ponte &
 382 Klein, 2015). Therefore, investigating the coherence of internal tides mode-by-mode is
 383 necessary. To achieve this, we utilize modal decomposition techniques.

384 M_2 tides are dominant for mode-1 and mode-2 tides at both moorings. At the north
 385 mooring (Figure 6a and 6b), M_2 is dominant with 84 J/m^2 (38% of total semidiurnal

energy), while S_2 and N_2 are only 6 J/m^2 (3%) each. M_2 also has the highest partition of energy among all semidiurnal constituents for each mode. Similarly, at the south mooring (Figures 6c and 6d), M_2 has the greatest partition with 113 J/m^2 (33%), compared to 30 J/m^2 (9%) for S_2 and 12 J/m^2 (3%) for N_2 . Considering constituent partitions in each mode, M_2 is dominant in both mode-1 (48%) and mode-2 (49%). Although the total semidiurnal tide energy is higher at the south mooring, M_2 mode-1 energy is higher at the north mooring.

Both moorings exhibit a large incoherent portion (yellow columns in Figure 6). The incoherent portion (129 J/m^2 , 57% at the north mooring, 187 J/m^2 , 55% at the south mooring) is higher than any single constituent and exceeds the total amount of all coherent components. This large incoherent portion is probably caused by the influence of California currents and eddies, which decrease the coherent fraction of tidal energy by wave refraction (Rainville & Pinkel, 2006). Nontidal noise, such as that arising from the "swing" mooring configuration and the relatively short observation period (~ 3 months), could also contribute to the large incoherent portion. In particular, the incoherent part of mode-3 at the south mooring, which accounts for over 87% of the total energy in that mode, is likely unrealistic and could be the result of nontidal noise. The "real" incoherent portion of the internal tide is unreliable when the signal-to-noise ratio is low. Overall, the observed incoherent tides from both moorings are close to the globally-averaged 45% (Zaron & Ray, 2017) or 49% (Nelson et al., 2019) semidiurnal nonstationary variance fraction (SNVF).

In terms of M_2 tides, mode-1 (69%) dominates mode-2 (10%) tides at the north mooring while mode-1 (48%) and mode-2 (49%) tides are comparable at the south mooring. In addition, mode-1 tides have similar energy levels between the north mooring (50 J/m^2) and the south mooring (46 J/m^2). However, relatively strong mode-2 M_2 tides with 44 J/m^2 are observed at the south mooring, compared to 5 J/m^2 at the north mooring. These results support our hypothesis above that multiwave interference happens here and that mode-1 and mode-2 tides originate from different generation sites and consistent with the speculations by J. Wang et al. (2022).

4.3 Changing phase velocity

In our previous discussion, we suggested the potential contribution of mesoscale currents and eddies to the incoherent component of internal tides. Here, we will explore this statement in more detail by examining the phase velocity of internal tides. J. Wang et al. (2022) detected the development of a warm-core anticyclonic mesoscale eddy from the mooring array during the pre-launch campaign. The three moorings were within the meander on 8 September and on the edge of the formed eddy by the end of the deployment on 24 November (Figure 7). The formation of this eddy coincides with the different temporal variations of energy in different semidiurnal modes (Figure 5). The phase velocity of internal waves is dependent on the ocean stratification. To assess the impact of background currents and mesoscale eddies on the temporal variations of internal tides, we derive the time series of phase velocity c_p for mode-1 and mode-2 tides (see equation in the Supporting Information). The ocean stratification required for these phase velocity calculations is based on the CTD-profile derived from WOA18 and the Wirewalker Profiler. Following the methodology outlined by Kerry et al. (2016), we employ a 3-day averaging for the buoyancy profile. This specific duration is chosen because the background mesoscale field displays minimal variability over this time scale.

There is a good linear relationship between absolute dynamic tomography (ADT) and the phase speed at the mooring location from Figures 8b and 8c, with R^2 of 0.82 for mode-1 and 0.72 for mode-2. Assuming this relationship is consistently applicable in the surrounding region, we can reconstruct the phase speed c_p from ADT in other latitudes (Figure 8a). Mode-1 tides are mainly southbound from the Mendocino Ridge, ac-

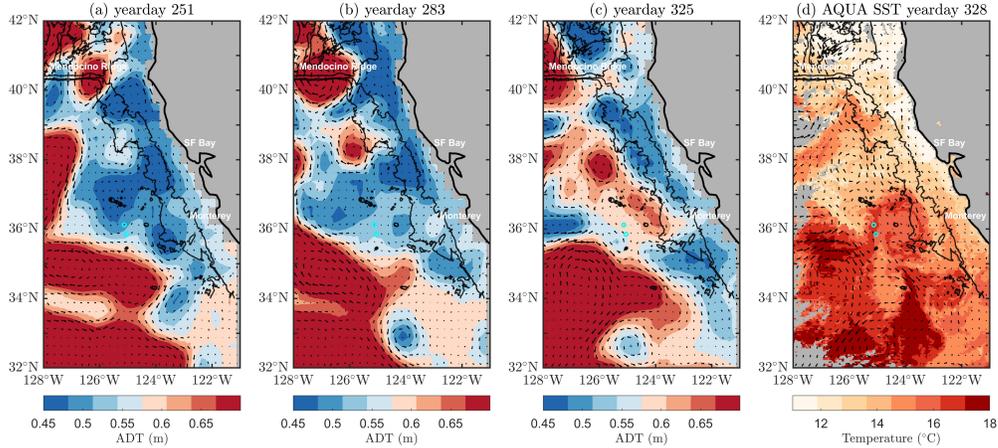


Figure 7: The absolute dynamic topography (ADT, color) and the surface geostrophic velocity anomaly (arrows) on (a) 8 September (yearday 251), (b) 10 October (yearday 283), and (c) 21 November (yearday 325), corresponding to the start, middle, and end of the pre-launch campaign in 2019. The ADT and surface geostrophic velocity field are from the Copernicus Marine Service. The cyan-colored dots mark the locations of the two moorings. A warm-core anticyclonic mesoscale eddy was formed close to the moorings. (d) The sea surface temperature (SST) after the formation of an eddy around November 24 (yearday 328), supporting the existence of an anticyclonic eddy. The SST data are MODIS Aqua Level 3 SST MID-IR Daily 4km product, from the Physical Oceanography Distributed Active Archive Center (PO.DAAC). Surface geostrophic velocity fields are provided for reference. Contours for the 3000-m and 3800-m isobath are shown.

437 cording to the altimetry data, which will be elaborated on in the next section. We thus
 438 reconstruct the time series of the phase velocity of mode-1 tides from the mooring and
 439 all the way up to 40.4°N (Figure 8d). Then we are able to derive the travel time of the
 440 wave propagating from the generation source (i.e., the Mendocinal Ridge) to both moor-
 441 ings by integrating the phase speed along latitude.

442 Mesoscale eddies are likely responsible for the increased travel time of mode-1 tides
 443 from their generation source to the mooring locations. The travel time (Figure 8e) of mode-
 444 1 M_2 tides to the north mooring shows a slight increase from 42.0 hours to 42.8 hours
 445 (2%). Similarly, the south mooring, located 30 km away, experiences waves with longer
 446 travel times by nearly an hour in yearday 325 after the eddy passed by. Mode-1 waves
 447 take from 44.4 hours to 45.3 hours (2%) to reach the south mooring. Hence, there is sim-
 448 ilar effect of mesoscale dynamics on mode-1 tides at both moorings. Although the re-
 449 sponse of mode-2 or higher mode tides to eddies may be stronger (Dunphy et al., 2017;
 450 Löb et al., 2020), the lack of comprehensive in-situ data and the effect of multiple sources
 451 of internal tides with equal contributions, make it challenging to provide a more quan-
 452 titative picture here. Additional research to detail the mechanism of wave-mesoscale in-
 453 teraction is needed. Ongoing researches involve both numerical simulations and theo-
 454 retical analyses, focusing on topics such as internal tide advection and refraction, enhanced
 455 dissipation of low-mode tides, and upscale energy transfer (Rainville & Pinkel, 2006; Sav-
 456 age et al., 2020; Y. Wang & Legg, 2023; Shakespeare, 2023). These studies inspire the
 457 design of future field programs that seek evidence for validation and potential adjust-
 458 ments for the parameterization and approximation in these theoretical and numerical
 459 models.

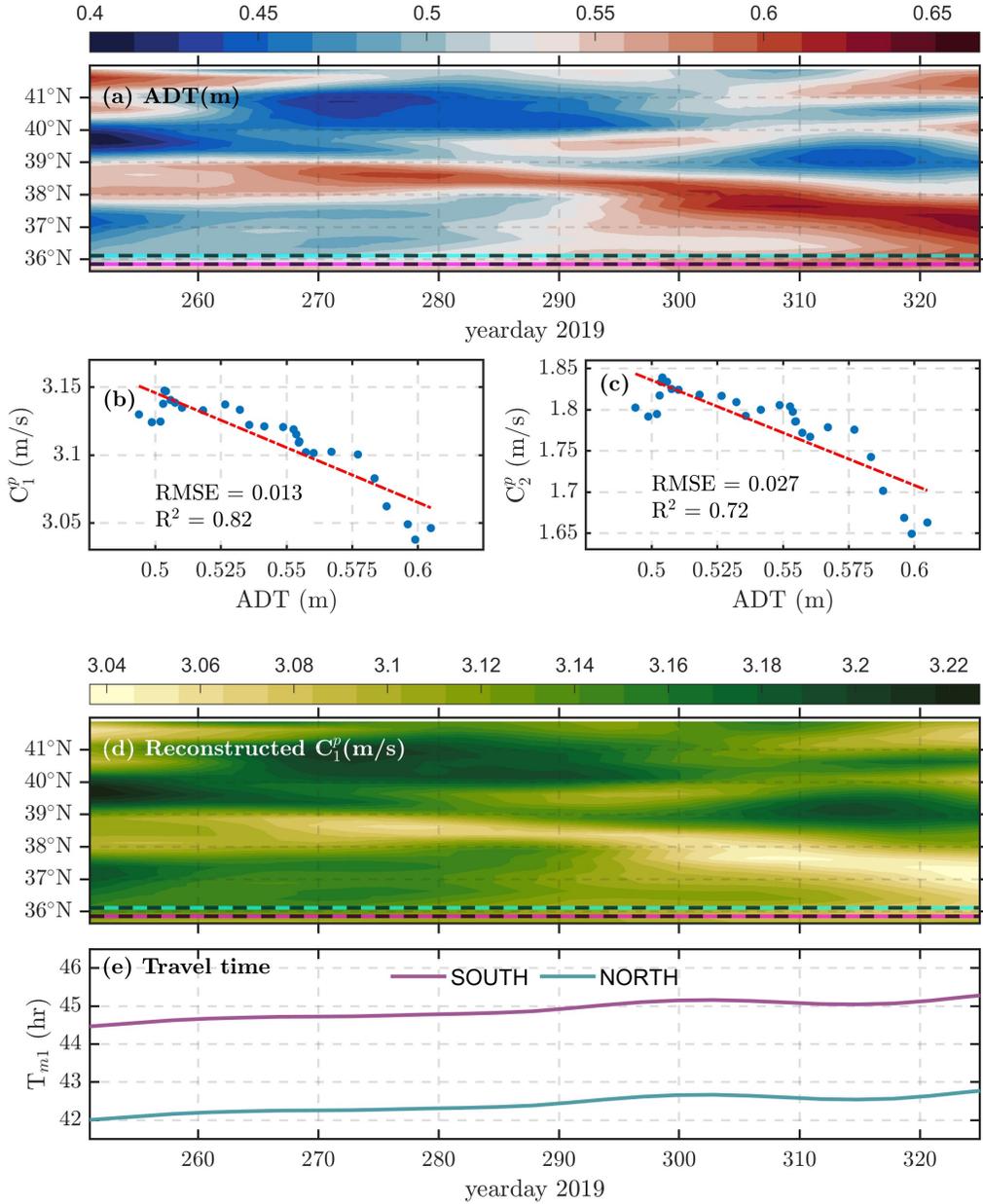


Figure 8: (a) Hovmöller diagram of the SLA at 125.1°W . Solid lines represent the latitudes of the north mooring (black-cyan striped lines) and the south mooring (black-magenta striped lines). (b) The correlation of absolute dynamic tomography (ADT) and phase velocity c_p of mode-1 M_2 tide. Three-day averaging is applied to the hourly buoyancy frequency profile. The blue dots are from moored observation. A linear fit is applied, and the fitted value is shown as a red dashed line. Root-mean-square error and R^2 are provided. (c) Same as (b) but for mode-2 M_2 tide. Using the linear relationship derived from moored observation, reconstructed phase velocity across latitude toward the sources of M_2 tides (the Mendocino Ridge at around 40.4°N) can be calculated from ADT (m). (d) The time series of the reconstructed phase velocity of mode-1 derived from ADT from (a). (e) The variability of the travel time (hr) of mode-1 tides over the record is estimated by integrating the phase velocity from the source (the Mendocino Ridge at 40.4°N) to the two moorings locations in the southward propagation direction.

5 Comparisons with Satellite Observations

The information obtained from in situ observations is insufficient for reconstructing the complete life cycle of internal tides. To better comprehend the internal tide in this region, we compare moored observations with internal tide models that are based on 27 years of satellite altimetry data. Our findings are as follows: (1) the amplitude and phase of both mode-1 and mode-2 M_2 internal tides extracted from the moorings are in good agreement with those obtained by satellite observations. (2) Despite the two moorings being only 30 km apart, there are spatial variations of M_2 internal tides due to interference from waves arriving from all directions. (3) We observe different features of mode-1 and mode-2 M_2 internal tides, resulting from distinct generation sites. Specifically, mode-1 tides mainly originate from the Mendocino Ridge and 36.5–37.5°N California continental slope, while mode-2 tides primarily come from local seamounts and Monterey Bay.

5.1 Altimetry Result

The SSHAs of mode-1 and mode-2 tides, derived from the 27-year-coherent M_2 altimetry model described in Section 3, reveal a complex internal tidal field in the studied region (Figure 9). This complexity is attributed to the presence of multiple sources for internal tides in the region, including the Hawaiian Ridge, the California continental slope, the Mendocino Ridge, and local generation over nearby seamounts. The superposition of multidirectional waves leads to the formation of standing-wave patterns. For mode-1 (Figure 9a), the predominant tidal waves propagate in north-south direction, originating from the Mendocino Ridge. Though the Mendocino Ridge is also a significant source for mode-2 tide (Figure 9b), the southward waves have a shorter excursion and do not reach the moorings location (cyan circles). Instead, the main sources of mode-2 at the two moorings are tidal beams originating from Monterey Bay and the Southern California Bight. However, the interference of multiple waves limits us to accurately determine the propagation direction of individual tidal beam and quantify its energy. Employing the multiwave decomposition approach is the key to overcoming the challenge and its effectiveness has been demonstrated in prior research (Zhao & Qiu, 2023).

The altimetry model offers a two-dimensional perspective on the generation and propagation of internal tides, which provides valuable context for interpreting the point-wise information obtained from mooring measurements. As such, the combination of these two data sets allows for a more comprehensive and nuanced analysis of the internal tidal field. For instance, the altimetry model can provide valuable insights into the spatial distribution of the mode-1 tidal beam and its relation to the mooring locations. This information can support the interpretation of the relatively small tidal amplitudes observed at the moorings, considering their proximity to the edge of the mode-1 tidal beam (Figure 9a). However, before delving into the detailed analysis, it is essential to establish the coherence and reliability of the two data sets to confidently utilize the altimetry model to shed light on the mooring observations.

5.2 Comparison with Moored Data

We compare the amplitude and phase of M_2 tides at the two mooring locations from moored and satellite altimetry observations. Figure 10 shows a high level of agreement, highlighting the precision and dependability of both data sources. Specifically, our analysis focuses on mode-1 and mode-2 M_2 signals due to their substantial energy content and their strong detectability through satellite observations.

The moored mode-1 M_2 tides (Figures 10a and 10b) exhibit an amplitude of 4.8 mm and a phase of 208 degrees at the south mooring. At the north mooring, the corresponding values are 5.1 mm and 121 degrees, both with a 95% confidence interval. Com-

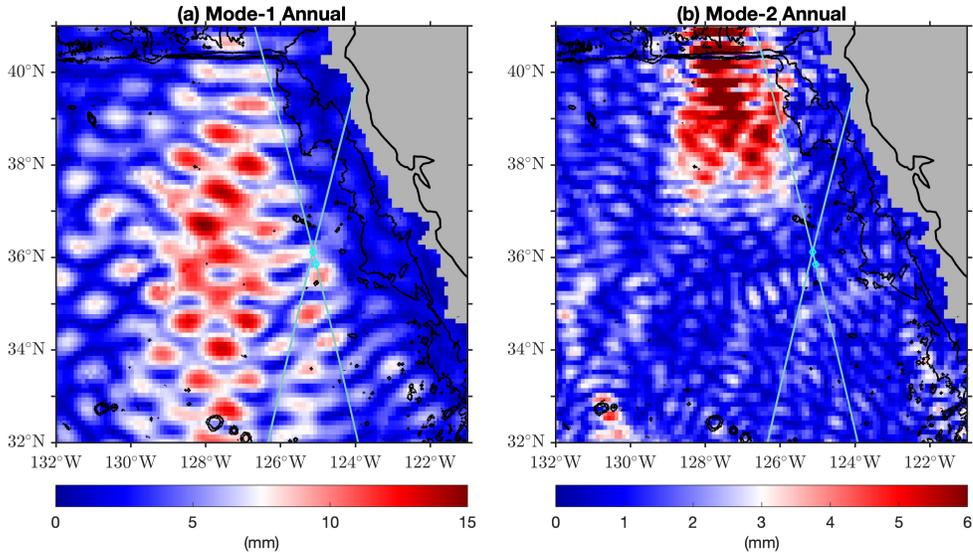


Figure 9: The SSHAs of M_2 (a) mode-1 and (b) mode-2 internal tides from the 27-year-coherent altimetry model. Note that different colorbar ranges are used for mode-1 and mode-2. Two cyan circles show the location of the two moorings from the SWOT pre-launch campaign. Two cyan lines crossing the north mooring are Sentinel-3A satellite tracks (S3A-140 and S3A-318). Contours for the 3000-m and 3800-m isobath are shown.

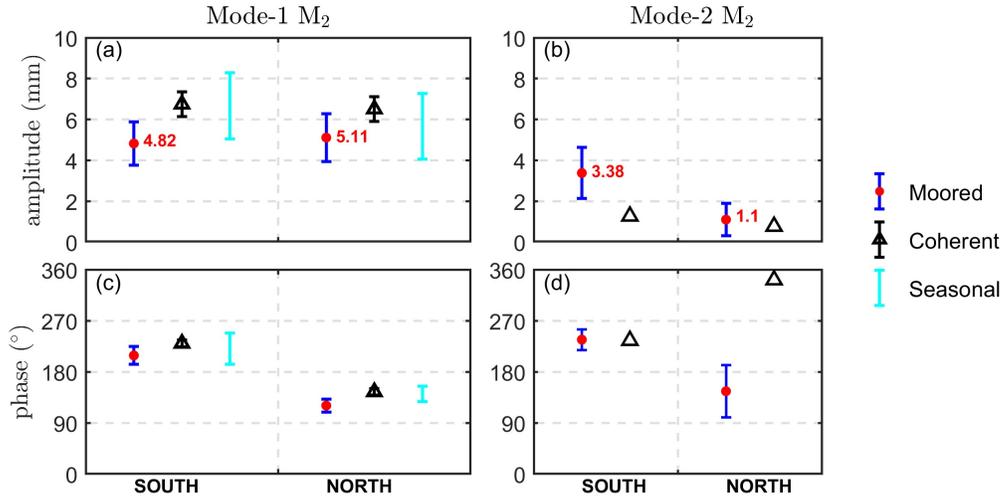


Figure 10: Comparison of moored and altimetry baroclinic SSHAs. (a) The amplitude (mm) and (c) phase ($^{\circ}$) of mode-1 M_2 SSHAs. The moored data are represented by red dots with a 95% confidence interval as blue bars. The amplitude is labeled explicitly. The black triangles depict results from the 27-year-coherent internal tide model. The black error bars for the 27-year-coherent model are ± 0.6 mm for amplitude and $\pm 6^{\circ}$ for phase. Four climatologically seasonal internal tide models span a range in cyan. The tidal features at the south mooring are plotted in the left column and those at the north mooring are in the right column. (b) The amplitude (mm) and (d) phase ($^{\circ}$) of mode-2 M_2 SSHAs.

paratively, the satellite altimetry models, depicted as black triangles for the 27-year-coherent model and cyan bars for the seasonal models, demonstrate a good agreement with the moored data at both locations, similar to the findings reported by Zhao et al. (2016). However, differences arise due to the disparity in record length, influencing the partitioning between incoherent and coherent signals. The altimetry observations used in the model span a much longer period (27 years) compared to the limited 3-month duration of the moored data. Extended observations enable the analysis to filter out the temporally variable component, resulting from interaction with other ocean dynamics and changing stratification, thus leading to bias-low result. Furthermore, the temporal variations of mode-1 tides, as discussed in Section 4, contributes to this imperfect correspondence. The altimetry measurements rarely capture the tidal variability associated with the advection and refraction caused by mesoscale eddies and currents. Nevertheless, this overall similarity emphasizes the accuracy and reliability of data obtained from moorings and satellite altimetry, taking into account the length of the recorded time series. The amplitude and phase of seasonal models cover a reasonably wide range.

The extraction of mode-2 M_2 tides poses greater challenges compared to mode-1 tides due to their relatively small amplitude and stronger seasonal variability. The moored mode-2 M_2 tides (Figures 10c and 10d) display an amplitude of 3.4 mm and a phase of 236 degrees at the south mooring. At the north mooring, the corresponding values are 1.1 mm and 146 degrees, both with a 95% confidence interval. The modest amplitude of mode-2 tides renders them more susceptible to noise. Furthermore, the combined effects of tidal interference and prominent seasonal variations contribute to the divergence between the results obtained from the two data sets. Despite these inherent difficulties and uncertainties, the moored and satellite findings regarding mode-2 M_2 tides exhibit consistency.

5.3 Generation and Propagation of Mode-1 and Mode-2 Tides

Consistent findings from both moorings and altimetry models reveal significant spatial variations of M_2 tides between the two moorings. In order to further investigate the altimetry results, we employ the 2D spatial filtering and plane wave analysis methods (Section 3). This method enables us to decompose the internal tidal field into different distinct propagation directions, providing a more detailed perspective on individual waves. Here, we decompose the 27-year-coherent altimetry results into three directions based on the dominant generation sites. Mode-1 tides are decomposed into southward waves (235° – 325°) from the Mendocino Ridge, eastward waves (-35° – 45°) from the Hawaiian Ridge, and northwestward waves (45° – 235°) from the local seamounts and continental slope. Mode-2 tides are decomposed into southward waves (245° – 325°) from the Mendocino Ridge, westward waves (125° – 245°) mostly from the continental slope, and northeastward waves (-35° – 125°) from the local seamounts. Through this decomposition, we are able to isolate and examine each wave, eliminating the interference caused by multiple waves (Zhao et al., 2019).

The mode-1 and mode-2 M_2 internal tides originate from different generation sites, based on the 27-year-coherent internal tide model. Mode-1 tides predominantly come from the Mendocino Ridge at 40.4°N , exhibiting a clear southward wave signal as depicted in Figure 11a. These waves propagate through both moorings, thereby explaining the relatively strong covariance observed in the moored data (Figure 5). These southward mode-1 waves are consistent with the previous in situ observation in this region (Alford, 2010; Musgrave et al., 2017). Interestingly, our analysis also indicates that local seamounts do not significantly contribute to the southward propagation of mode-1 tides, suggesting that these dominant and relatively larger mode-1 are not sensitive to minor topographic features. Internal tides from the California continental slope propagate northwestward (Figure 11b). Specifically, waves come from the Southern California Bight and the 36.5 – 37.5°N continental slope. In addition, two moorings are affected

561 by the eastward tidal waves from Hawaiian Ridge (Figure 11c). These remotely gener-
 562 ated waves, originating outside of this region, have been recognized as significant contri-
 563 butors to the internal tide energetics in previous studies (Ray & Zaron, 2016; Zhao
 564 et al., 2016; Siyanbola et al., 2023; Mazloff et al., 2020).

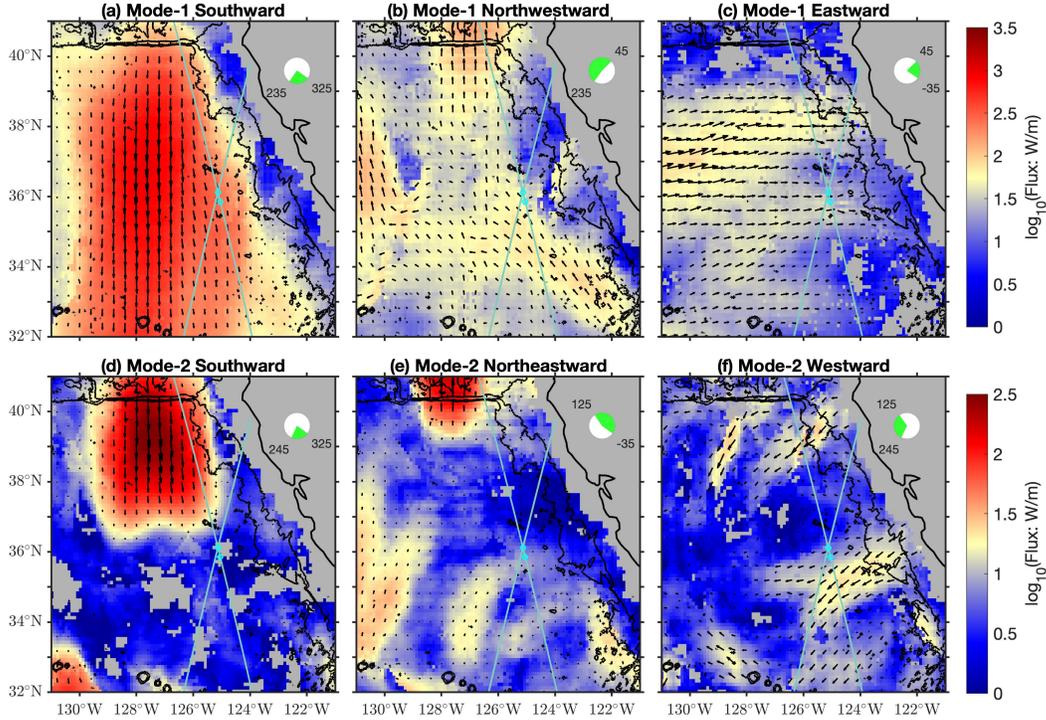


Figure 11: Fluxes of regional (a-c) mode-1 and (d-f) mode-2 M_2 internal tides from the 27-year-coherent model are shown in logarithmic scale. The internal tidal field has been decomposed into three components by propagation direction (directional range is shown as a green pie chart in the right upper corner). Colors and arrows indicate the magnitude and direction of internal tides, respectively. Note that different color bar ranges are used for different modes. Two cyan circles show the location of the two moorings from the SWOT pre-launch campaign. Two cyan lines crossing the north mooring are Sentinel-3A satellite tracks (S3A-140 and S3A-318). Contours for the 3000-m and 3800-m isobath are shown.

565 However, the behavior of mode-2 tides presents a different story. As illustrated in
 566 Figure 11d, the southward flux of mode-2 tides originating from the Mendocino Ridge
 567 (40.4°N) diminishes around 36.5°N . Consequently, unlike mode-1 tides, southward mode-
 568 2 tides have minimal impact on the mooring locations, likely due to dissipation or scatter-
 569 ing processes. These processes can cause the mode-2 tides to dissipate into incoher-
 570 ence or scatter into higher modes. This finding is consistent with the simulation obtained
 571 from MITgcm, indicating that the southward mode-2 tide propagates only a quarter of
 572 the distance covered by the mode-1 tide (Zhao et al., 2019). Instead, the mode-2 tides
 573 detected by the two moorings are northeastward waves (Figure 11e) generated by local
 574 seamounts such as Fieberling Seamount and Hoke Seamount (Kunze & Toole, 1997; Zhao,
 575 2018), and westward waves (Figure 11f) from the continental slope, including Monterey
 576 bay (G. Carter, 2010). Unlike mode-1, the remotely generated mode-2 tides from Hawai-
 577 ian Ridge dissipate along the way and barely reach this region, i.e., there is no sign of

578 eastward waves to the moorings location. The presence of multiple sources for mode-2
 579 tides, combined with the complex sea surface height (SSH) field resulting from tidal in-
 580 terference observed in satellite observations (Figure 9b), explains the weak correlation
 581 of mode-2 tides at the two moorings (Figure 5). Furthermore, it is worth noting that the
 582 position of the two moorings in the SWOT pre-launch campaign did not align with any
 583 mode-2 tidal beam (Figure 9b), resulting in a relatively attenuated signal compared to
 584 that of mode-1 tides.

585 Overall, this result highlights the complexity of the internal wave field in this re-
 586 gion and emphasizes the importance of utilizing advanced techniques, such as 2D spa-
 587 tial filtering and plane wave analysis, to directionally decompose and investigate indi-
 588 vidual wave characteristics. The observed diverse generation and propagation of M_2 mode-
 589 1 and mode-2 tides aligns with the findings obtained from the MITgcm simulation (Zhao
 590 et al., 2019).

591 6 Seasonal Variations

592 In the CCS region, the generation and propagation of internal tides are influenced
 593 by seasonal changes in stratification, background currents, and eddies (Zhao et al., 2012;
 594 Johnston & Rudnick, 2015). For example, in winter with weak stratification, tides prop-
 595 agate more slowly (Zhao, 2021). This weakened stratification is likely due to the cool-
 596 ing of the surface waters and weaker alongshore winds south of Cape Mendocino, which
 597 result in less restratification (Checkley Jr & Barth, 2009). The propagation speed of tidal
 598 waves during different seasons provides valuable information about ocean stratification
 599 and heat distribution.

600 To address this, we utilize the latest seasonal altimetry model (Section 3) to inves-
 601 tigate the seasonal variations of mode-1 M_2 internal tides. The same mapping procedure
 602 employed in the 27-year-coherent model is used, but with four seasonal subsets. We de-
 603 compose the waves into three propagation directions, maintaining the same range as in
 604 the 27-year-coherent model for comparison. Different seasonal models are analyzed by
 605 looking at the SSHAs and the magnitude and direction of energy flux in the CCS region.

606 6.1 Interference Patterns

607 The mode-1 internal tidal field associated with sea surface height anomalies (SSHAs)
 608 exhibits a complex pattern in all of the seasonal models (see Figure S8 in the Support-
 609 ing Information). The averaged Pearson correlation coefficient of SSHA between every
 610 two seasonal models is 0.84, indicating the statistical importance of the seasonality on
 611 the internal tidal field. For the SWOT mission (swaths in green in Figure S8), it serves
 612 as a compelling example of why it is crucial to account for the complexity and season-
 613 ality of internal tides when applying tidal correction. Due to the complex multi-wave in-
 614 terference, it is challenging to quantitatively analyze the seasonality in this region. There-
 615 fore, we employ multi-wave decomposition techniques for each seasonal model.

616 6.2 Generation Sites

617 We decompose the waves from all four seasonal models into three propagation di-
 618 rections and examine the energy flux (W/m) in each direction (Figure 12). To quantify
 619 the seasonal effects on the generation and propagation of mode-1 tides, for each direc-
 620 tion, we analyze data along two cross sections roughly perpendicular to the propagation
 621 shown by the striped lines.

622 The southward waves (Figures 12a, 12d, 12g, and 12j) originating from the Men-
 623 docino Ridge play a consistently dominant role throughout the year. We focus on two
 624 zonal cross sections at 34°N and 36°N (striped lines). The 36°N section represents the

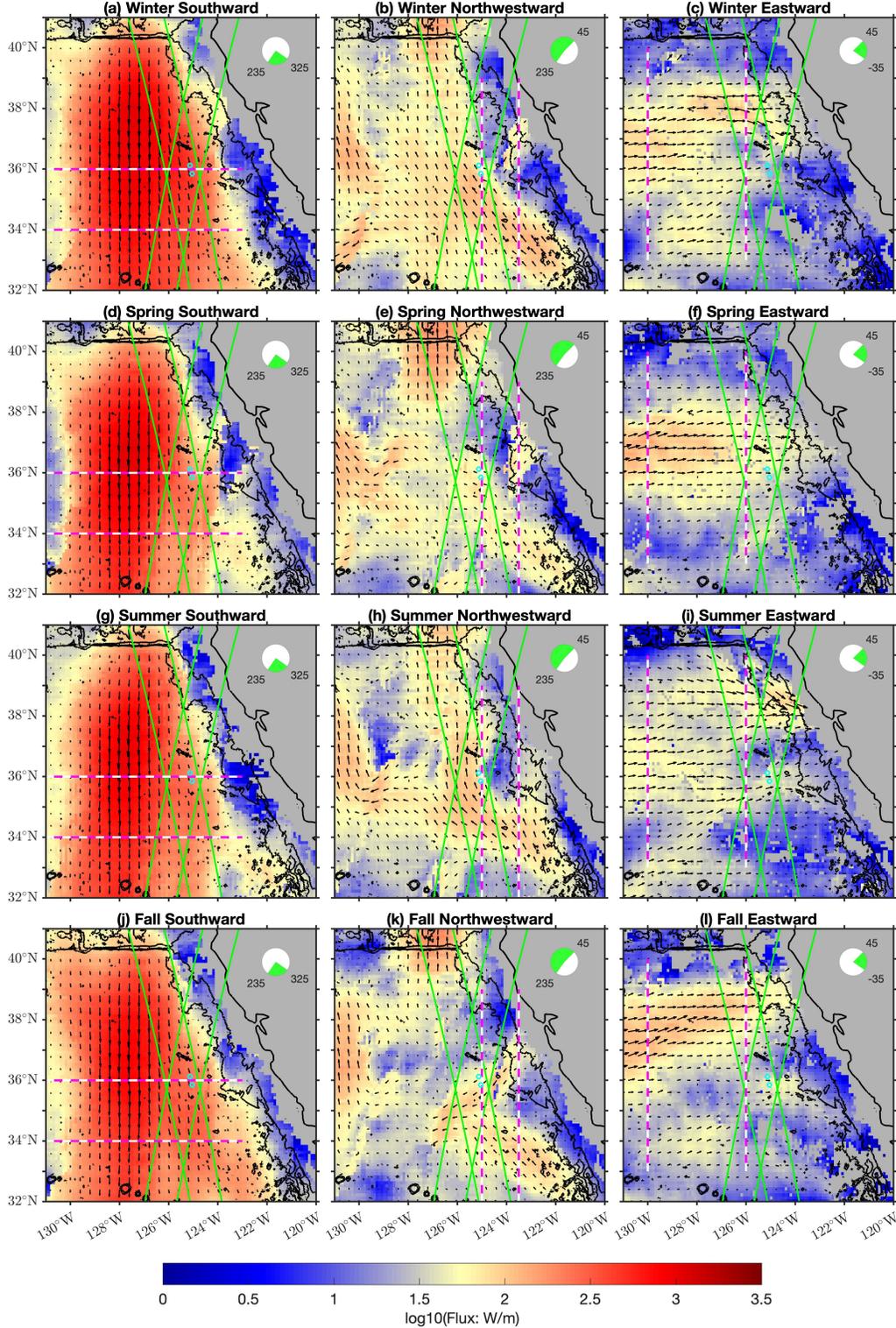


Figure 12: Fluxes of regional mode-1 M_2 internal tides from four climatologically seasonal model, (a–c) winter, (d–f) spring, (g–i) summer, and (j–l) fall, all of which are shown on a logarithmic scale. The internal tidal field has been decomposed into three components by propagation direction. Directional range is shown as a green pie chart. Colors and arrows indicate the magnitude and direction of internal tides, respectively. Two cyan circles show the location of the two moorings from the SWOT campaign. Green lines are the SWOT Cal/Val swath tracks. For each component, the two cross sections (striped lines) are given. The zonal cross sections are chosen at 34°N and 36°N for the southward waves. The meridional cross sections are chosen at 123.5°W and 125°W for the northwestward waves, and at 126°W and 130°W for the eastward waves.

energy peak of the southward waves, while the 34°N section represents the energy dissipation during propagation. For each section, we integrate the energy flux between 123°W and 131°W. The result will be discussed in the following section.

For the northwestward waves (Figures 12b, 12e, 12h, and 12k), we focus on the seasonal variations of the tidal beam from the Southern California Bight (SCB) and the tidal beam from the 36.5–37.5°N continental slope (hereinafter as “36.5–37.5°N”) in each seasonal model. The complex topography of islands, ridges, sills, deep basins, headlands, bays, and shelves in the SCB leads to an active internal wave field (Lerczak et al., 2003; M. Buijsman et al., 2012). The tidal beams from two sources are consistent with the MIT-gcm simulation (see fig. 8 from Zhao et al. (2019)). We select two meridional sections at 123.5°W and 125°W (striped lines) and integrate the energy flux between 32°N and 39°N. The two sections are chosen at the location before (123.5°W) and after (125°W) the waves from two sources merge. More quantitative analysis of the relative strengths of the two sources and their seasonality will be discussed in the next section.

The eastward waves (Figures 12c, 12f, 12i, and 12l), mainly generated from the Hawaiian Ridge, are evident in spring and fall. We quantify these seasonal variations by comparing the energy flux across two meridional sections at 126°W and 130°W (striped lines), both spanning between 33°N and 40°N. However, it is challenging to determine the main drivers of the seasonality of eastward waves. Factors such as background currents, eddies, and refraction of steep topography can alter the long-range waves generated from the Hawaiian Ridge after traveling 3,000 km (Dunphy & Lamb, 2014; Ponte & Klein, 2015). In addition, there are eastward tides possibly generated from or scattered by the local seamounts (e.g., the Spiess Seamounts Chain) and the fracture zone (e.g., the Murray Fracture Zone). This complexity of multiple sources contributes to the broad tidal beam, especially observed in the winter and summer models.

6.3 Cross Section Energy Flux

A cursory glance above indicates seasonal variations of internal tides from different directions. A more quantitative statement is obtained by looking at the energy flux through cross sections. By examining the distinct zonal (southward waves) or meridional (eastward and westward waves) variations of the cross-beam energy flux among four seasonal models and the 27-year-coherent model, we aim to gain a comprehensive understanding of the magnitude and direction of energy transfer. To facilitate this analysis, the cross-beam energy fluxes are averaged within 0.5-degree-wide sections and smoothed every 5 grid points along each cross section. Moreover, we will integrate and compare the cross-beam energy flux among the seasonal models and the 27-year-coherent model.

The analysis of southward waves (Figure 13a and b) reveals distinct energy flux patterns along the cross sections. At both the 34°N and 36°N sections, the highest flux peaks are observed around 128°W. The spring season (blue) exhibits the strongest flux peak at the 36°N section, while the winter season (green) shows the highest flux peak at the 34°N section. In contrast, the fall season (cyan) exhibits the weakest peaks in both sections, indicating an attenuation in southward tidal wave. The width of the tidal beam is approximately 400–500 m at the 36°N section, with the widest beams observed during the winter season in both sections. Particularly between 124°W and 128°W, the flux is exceptionally elevated during the winter season. As the internal tides propagate approximately 222 km to the 34°N section, an average of 20% of their energy flux dissipates, with the spring season experiencing the highest dissipation (22.5%). Notably, at the 34°N section, all seasons experience a flux reduction around 126°W, possibly due to refraction from steep topography. The 27-year-coherent model (black) generally represents the average of the four seasonal models. The cross-beam integrated energy flux is strongest during winter (321 MW and 260 MW) and weakest during fall (231 MW and 186 MW) at both sections.

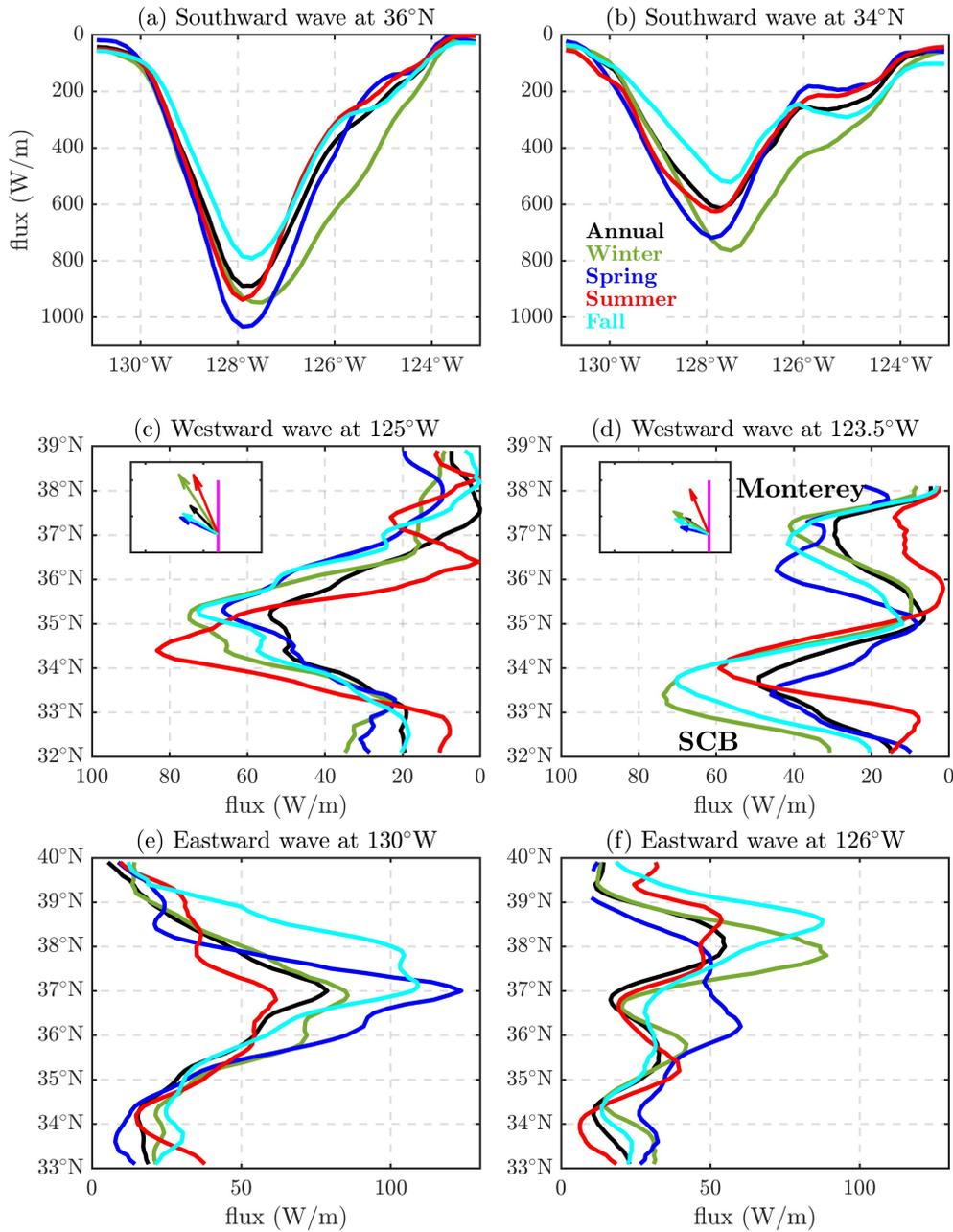


Figure 13: Seasonality of cross-beam energy flux from three directions. Southward energy flux is illustrated across (a) 36°N and (b) 34°N . Note that the flux is shown in the reversed direction of the y axis to align with southward waves. At both latitudes, each section spans between 123°W and 131°W . Westward energy flux is presented across (c) 125°W and (d) 123.5°W . Note that the flux is shown in the reversed direction of the x axis to align with westward waves. Each section spans between 32°N and 39°N . The cross-section summation of energy flux for five models is shown as vectors in the inset at the upper left corner. Eastward energy flux is showcased across (e) 130°W and (f) 126°W . Each section spans between 33°N and 40°N . The 27-year-coherent model is in black and the four seasonal models are in green for winter, blue for spring, red for summer, and cyan for fall.

676 For westward waves, both the relative strength of two sources (Figure 13d) and the
 677 total energy flux after their merge (Figure 13c) exhibit significant seasonal variations.
 678 At the 123.5°W section (Figure 13d), the presence of two distinct flux peaks signifies the
 679 different tidal waves originating from the SCB and the 36.5–37.5°N, consistent with Fig-
 680 ure 12. The flux south of 35°N represents the northwestward waves from the SCB, while
 681 the flux north of 35°N represents the southwestward waves from the 36.5–37.5°N. The
 682 strength of the northwestward tides from the SCB remains relatively consistent across
 683 all seasons, except for a 40% weakening during the spring season compared to the av-
 684 erage from the other three seasonal models. Similarly, the strength of the southwestward
 685 tides from the 36.5–37.5°N is relatively consistent across all seasons, except for a signif-
 686 icant decrease to 25% of the average energy flux during the summer season compared
 687 to other three seasonal models. This weakening of internal tides at the 36.5–37.5°N dur-
 688 ing summer leads to a variation in the direction of the integrated energy flux in the sum-
 689 mer model, where the flux is mainly determined by the northwestward waves from the
 690 SCB (see inset in Figure 13d). After the waves from the two sources merge at 125°W
 691 (Figure 13c), the meridional distributions of flux across different models are generally
 692 similar. However, there are some differences. In summer, the flux peak is shifted south-
 693 ward and observed at 34.5°N, while in the other three seasons, the peaks occur at 35.3°N.
 694 This shift is possibly due to weak generation from the 36.5–37.5°N. Additionally, dur-
 695 ing the summer season, a second peak at 37.3°N is observed, representing tides originated
 696 from the continental slope north of San Francisco Bay (e.g., Arena Canyon and Bodega
 697 Canyon). The integrated energy flux varies in magnitude and direction among the mod-
 698 els (insets in Figure 13c and d), indicating significant seasonality. The phase differences
 699 among the seasonal models lead to lower tidal energy flux in the 27-year-coherent model.

700 Turning to eastward waves (Figure 13e and f), we observe flux is intensified at 130°W
 701 in both the spring and fall models (Figure 13e), consistent with the distinct tidal beam
 702 observed in Figure 12f and 12l. This energy flux is primarily generated by remotely gen-
 703 erated waves originating from the Hawaiian Ridge. While the peaks at 37°N are higher
 704 during the spring season, the fall season exhibits the strongest integrated energy flux (41
 705 MW), mainly attributed to the relatively strong tidal flux between 37.5°N and 39°N. At
 706 the 126°W section, the flux peak is around 38°N. Particularly during the fall and win-
 707 ter seasons, the flux peaks are twice as strong as those from the 27-year-coherent and
 708 other seasonal models. Dissipation occurs in all seasons after tides propagate to 126°W.
 709 However, the energy flux redistribution observed in summer and the formation of the win-
 710 ter peak after propagating 400 km indicate the influence of local generation from nearby
 711 seamounts and refraction of the fracture zone. These factors contribute to the complex-
 712 ity of the eastward tidal wave dynamics in the region.

713 To summarize, southward waves from the Mendocino Ridge consistently play a dom-
 714 inant role throughout the year, with maximum amplitude in winter and the minimum
 715 in fall. However, during fall and spring, we observe the strongest eastward waves, gen-
 716 erated remotely from the Hawaiian Ridge. Westward waves from the 36.5–37.5°N con-
 717 tinental slope are weakest during summer while those from the Southern California Bight
 718 are weakest during spring. To quantify the seasonal variability for waves from each di-
 719 rection, we calculate the coefficient of variation of integrated energy flux in four seasons.
 720 The westward waves have the highest variability of flux with $\pm 22\%$, while the south-
 721 ward waves have the lowest variability with $\pm 13\%$.

722 As a simplified representation of the complex internal tidal field, this cross-section
 723 analysis could potentially underestimate the magnitude of the energy flux, as it only ac-
 724 counts for the portion that is orthogonal to the section. Also, the seasonal variations may
 725 be dependent on the definition of four seasons and corresponding ocean conditions. The
 726 definition of seasons and corresponding ocean conditions can vary depending on the re-
 727 search and the specific region of study. For the CCS region, some studies have utilized
 728 the alongshore wind direction as a criterion for defining seasons. In this approach, upwelling-

729 favorable conditions are characterized by equatorward winds, while poleward winds and
 730 storms indicate downwelling-favorable conditions (Checkley Jr & Barth, 2009; Dettinger,
 731 2011). This leads to a longer summer (June-September) and winter (December-February).
 732 Other factors, such as water temperature, energy sink from wind-current feedback (Delpech
 733 et al., 2023), and local atmospheric conditions, can also influence the seasonal variabil-
 734 ity of internal tides. The underlying mechanisms driving these variations warrant fur-
 735 ther investigation. Despite these considerations, the evident seasonality of internal tides
 736 in the region has significant implications for ocean mixing and circulation. The inclu-
 737 sion of seasonal variability in ocean models is crucial for capturing the dynamic nature
 738 of internal tides and their interactions with other oceanic processes. By incorporating
 739 seasonal variations, models can better represent the complex temporal dynamics of in-
 740 ternal tides, leading to improved predictions and understanding of oceanic phenomena.

741 7 Conclusions and Discussion

742 The study examines the temporal and spatial variations of semidiurnal internal tides
 743 off central California. This is achieved by utilizing both moored data from the SWOT
 744 pre-launch campaign in 2019 and internal tidal models from 27 years of altimetry. Pro-
 745 nounced semidiurnal internal tides are observed at both moorings. The south mooring
 746 exhibits stronger semidiurnal tidal energy, while the north mooring shows higher am-
 747 plitudes of the mode-1 M_2 internal tide. A warm anticyclone eddy during the measure-
 748 ments may have slowed the propagation speed of internal tides, leading to temporal vari-
 749 ability. Mode-1 tides from the two moorings are temporally correlated, whereas mode-
 750 2 tides are not. This discrepancy is likely caused by complex interference patterns re-
 751 sulting from waves originating from different directions.

752 The satellite models help explain the spatial variation of M_2 tides observed by the
 753 moorings and provides a comprehensive description of mode-1 and mode-2 tides in the
 754 region. The agreement between the moored and satellite results, in terms of both am-
 755 plitude and phase, supports the reliability of the satellite altimetry model. Different char-
 756 acteristics are observed between mode-1 and mode-2 M_2 tides, indicating distinct gen-
 757 eration sources. Mode-1 tides are primarily generated from the Mendocino Ridge and
 758 the 36.5–37.5°N California continental slope, while mode-2 tides originate mostly from
 759 local seamounts and Monterey Bay. Additionally, seasonal variations are observed in the
 760 generation and propagation of the regional mode-1 M_2 internal tides. The winter sea-
 761 son exhibits the strongest southward waves from the Mendocino Ridge and westward waves
 762 from the continental slope. In contrast, the fall season shows the strongest eastward waves,
 763 generated remotely from the Hawaiian Ridge, while exhibiting the weakest southward
 764 waves. Westward waves are weakest during the summer, possibly due to weak genera-
 765 tion from the continental slope, increased dissipation during propagation, or a combi-
 766 nation of both factors. Overall, the westward waves have the highest seasonal variabil-
 767 ity of tidal flux with $\pm 22\%$, while the southward waves have the lowest variability with
 768 $\pm 13\%$.

769 This analysis has limitations. The moorings have finite vertical resolution which
 770 limits the ability to accurately resolve the high modes (Nash et al., 2005). The analy-
 771 sis finds relatively weak internal tides compared to other regions such as the Hawaiian
 772 Ridge, the South China Sea, the Tasman Sea and the Mid-Atlantic Ridge (Alford et al.,
 773 2015; Zhao et al., 2016; Xu et al., 2016) which may introduce uncertainties due to the
 774 lower signal to noise. This is partially addressed by a sensitivity analysis, accurate cor-
 775 rection for the mooring motion, and robust statistical analysis with 95% confidence in-
 776 tervals. The relatively short mooring records may not be directly comparable to the 27-
 777 year average of satellite altimetry and the point mooring measurements may not be di-
 778 rectly comparable to the 160-km averaged satellite data. In particular, estimating the
 779 impact of mesoscale eddies on internal tides solely through short-term two-mooring mea-
 780 surements is challenging and these results are only suggestive, but offer some insights;

781 an array of moorings with a longer measurement period would be better (Huang et al.,
782 2018).

783 This study enhances our understanding of internal tide variability within the CCS
784 region, providing valuable insights for future research for SWOT and numerical mod-
785 eling endeavors (Arbic, 2022). For the SWOT tidal aliasing issue due to its long repeat
786 cycle, it is crucial to correct for unresolved internal tides before deriving and analyzing
787 submesoscale dynamics from the SWOT data, especially in regions where significant mode-
788 1 and mode-2 baroclinic tides exist (Qiu et al., 2018; Kelly et al., 2021; Carrere et al.,
789 2021). Our findings suggest that the incorporation of seasonal variability of internal tides
790 holds significant potential to improve the SWOT tidal correction. By quantifying the
791 contributions of internal tide and investigating its dynamics in this region, researchers
792 can fully explore the potential of observation-based data sets in studying various scales
793 and enhancing our understanding of air-ocean dynamics across different temporal and
794 spatial extents, ultimately impacting large-scale climate dynamics (Farrar et al., 2020).

795 8 Open Research

796 The SWOT pre-launch field campaign 2019-2020 data were downloaded from the
797 NASA Physical Oceanography Distributed Active Archive Center ([https://podaac.jpl
798 .nasa.gov/announcements/2022-06-09-SWOT-2019-2020-Prelaunch-Oceanography
799 -Field-Campaign-Dataset-Release](https://podaac.jpl.nasa.gov/announcements/2022-06-09-SWOT-2019-2020-Prelaunch-Oceanography-Field-Campaign-Dataset-Release)). The MODIS Aqua Level 3 SST product was down-
800 loaded from NASA Physical Oceanography Distributed Active Archive Center. The World
801 Ocean Atlas 2018 is produced and made available by the NOAA National Centers for
802 Environmental Information (<https://www.ncei.noaa.gov/products/world-ocean-atlas>).
803 The absolute dynamic topography and the surface geostrophic velocity anomaly data are
804 collected from the Copernicus Marine Service (<https://doi.org/10.48670/moi-00148>).
805 The bathymetry information is referred to in the General Bathymetric Chart of the Oceans
806 (GEBCO, <https://www.gebco.net/>). The satellite altimetry internal tide models will
807 be made public on the acceptance of this paper.

808 Acknowledgments

809 This work was supported by NSF projects OCE1947592, OCE2149028 and 80NSSC19K1007
810 P0007. We acknowledge the efforts of those involved in and supporting the public avail-
811 ability of the SWOT pre-launch campaign mooring data and the satellite altimeter data
812 used here. We express gratitude to the anonymous reviewers for their valuable sug-
813 gestions and comments.

814 References

- 815 Alford, M. H. (2010). Sustained, full-water-column observations of internal waves
816 and mixing near Mendocino Escarpment. *Journal of physical oceanography*,
817 *40*(12), 2643–2660.
- 818 Alford, M. H., Mickett, J. B., Zhang, S., MacCready, P., Zhao, Z., & Newton, J.
819 (2012). Internal waves on the washington continental shelf. *Oceanography*,
820 *25*(2), 66–79.
- 821 Alford, M. H., Peacock, T., MacKinnon, J. A., Nash, J. D., Buijsman, M. C., Centu-
822 rioni, L. R., . . . others (2015). The formation and fate of internal waves in the
823 South China Sea. *Nature*, *521*(7550), 65–69.
- 824 Althaus, A. M., Kunze, E., & Sanford, T. B. (2003). Internal tide radiation from
825 mendocino escarpment. *Journal of Physical Oceanography*, *33*(7), 1510–1527.
- 826 Ansong, J. K., Arbic, B. K., Alford, M. H., Buijsman, M. C., Shriver, J. F., Zhao,
827 Z., . . . others (2017). Semidiurnal internal tide energy fluxes and their variabil-
828 ity in a Global Ocean Model and moored observations. *Journal of Geophysical*

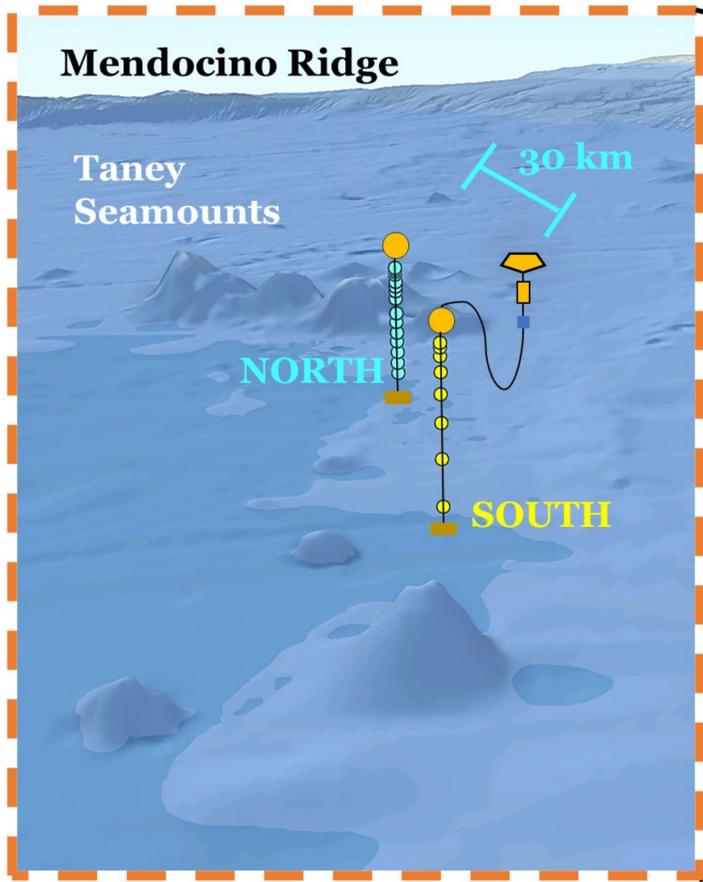
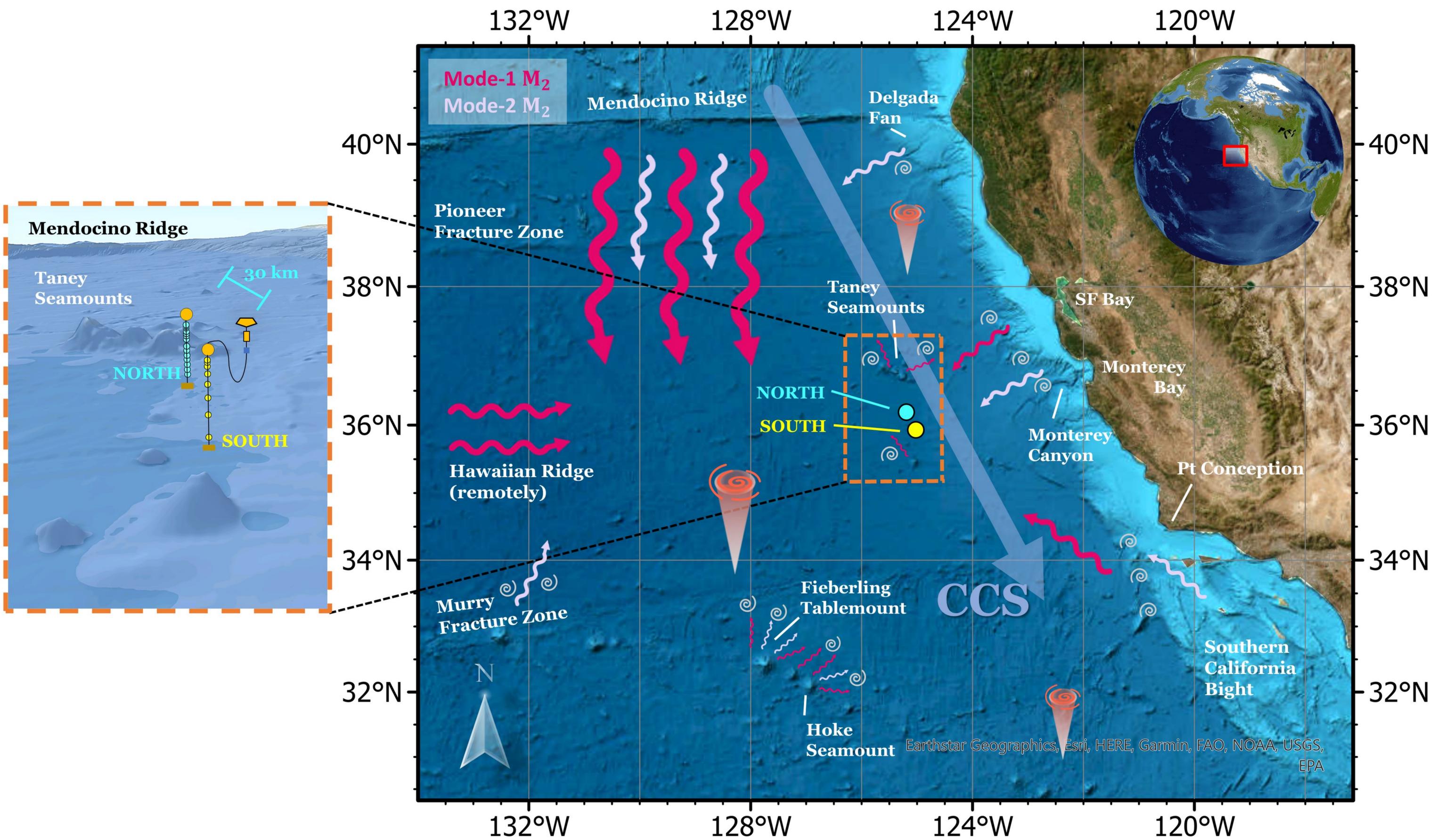
- 829 *Research: Oceans*, 122(3), 1882–1900.
- 830 Arbic, B. K. (2022). Incorporating tides and internal gravity waves within global
831 ocean general circulation models: A review. *Progress in Oceanography*,
832 102824.
- 833 Buijsman, M., Uchiyama, Y., McWilliams, J., & Hill-Lindsay, C. (2012). Modeling
834 semidiurnal internal tide variability in the Southern California Bight. *Journal*
835 *of Physical Oceanography*, 42(1), 62–77.
- 836 Buijsman, M. C., Arbic, B. K., Richman, J. G., Shriver, J. F., Wallcraft, A. J., &
837 Zamudio, L. (2017). Semidiurnal internal tide incoherence in the equatorial
838 pacific. *Journal of Geophysical Research: Oceans*, 122(7), 5286–5305.
- 839 Carrere, L., Arbic, B. K., Dushaw, B., Egbert, G., Erofeeva, S., Lyard, F., . . . oth-
840 ers (2021). Accuracy assessment of global internal-tide models using satellite
841 altimetry. *Ocean Science*, 17(1), 147–180.
- 842 Carter, G. (2010). Barotropic and baroclinic m 2 tides in the monterey bay region.
843 *Journal of Physical Oceanography*, 40(8), 1766–1783.
- 844 Carter, G. S., Gregg, M. C., & Lien, R.-C. (2005). Internal waves, solitary-like
845 waves, and mixing on the Monterey Bay shelf. *Continental Shelf Research*,
846 25(12-13), 1499–1520.
- 847 Checkley Jr, D. M., & Barth, J. A. (2009). Patterns and processes in the California
848 Current System. *Progress in Oceanography*, 83(1-4), 49–64.
- 849 de Lavergne, C., Falahat, S., Madec, G., Roquet, F., Nycander, J., & Vic, C. (2019).
850 Toward global maps of internal tide energy sinks. *Ocean Modelling*, 137, 52–
851 75.
- 852 Delpech, A., Barkan, R., Renault, L., McWilliams, J., Siyanbola, O. Q., Buijsman,
853 M. C., & Arbic, B. K. (2023). Wind-current feedback is an energy sink for
854 oceanic internal waves. *Scientific Reports*, 13(1), 5915.
- 855 Dettinger, M. (2011). Climate change, atmospheric rivers, and floods in California—a
856 multimodel analysis of storm frequency and magnitude changes. *JAWRA Jour-*
857 *nal of the American Water Resources Association*, 47(3), 514–523.
- 858 Dunphy, M., & Lamb, K. G. (2014). Focusing and vertical mode scattering of the
859 first mode internal tide by mesoscale eddy interaction. *Journal of Geophysical*
860 *Research: Oceans*, 119(1), 523–536.
- 861 Dunphy, M., Ponte, A. L., Klein, P., & Le Gentil, S. (2017). Low-mode internal
862 tide propagation in a turbulent eddy field. *Journal of Physical Oceanography*,
863 47(3), 649–665.
- 864 Farrar, J. T., D’Asaro, E., Rodriguez, E., Shcherbina, A., Czech, E., Matthias, P.,
865 . . . others (2020). S-MODE: The Sub-Mesoscale Ocean Dynamics Experi-
866 ment. In *Igarss 2020-2020 ieee international geoscience and remote sensing*
867 *symposium* (pp. 3533–3536).
- 868 Gill, A. E., & Adrian, E. (1982). *Atmosphere-ocean dynamics* (Vol. 30). Academic
869 press.
- 870 Haney, R. L., Hale, R. A., & Dietrich, D. E. (2001). Offshore propagation of eddy
871 kinetic energy in the california current. *Journal of Geophysical Research:*
872 *Oceans*, 106(C6), 11709–11717.
- 873 Huang, X., Wang, Z., Zhang, Z., Yang, Y., Zhou, C., Yang, Q., . . . Tian, J. (2018).
874 Role of mesoscale eddies in modulating the semidiurnal internal tide: Observa-
875 tion results in the northern South China Sea. *Journal of Physical Oceanogra-*
876 *phy*, 48(8), 1749–1770.
- 877 Johnston, T. S., & Rudnick, D. L. (2015). Trapped diurnal internal tides, propagat-
878 ing semidiurnal internal tides, and mixing estimates in the California Current
879 System from sustained glider observations, 2006–2012. *Deep Sea Research Part*
880 *II: Topical Studies in Oceanography*, 112, 61–78.
- 881 Kelly, S. M., Lermusiaux, P. F., Duda, T. F., & Haley, P. J. (2016). A coupled-mode
882 shallow-water model for tidal analysis: Internal tide reflection and refraction
883 by the Gulf Stream. *Journal of Physical Oceanography*, 46(12), 3661–3679.

- 884 Kelly, S. M., Waterhouse, A. F., & Savage, A. C. (2021). Global Dynamics of the
885 Stationary M_2 Mode-1 Internal Tide. *Geophysical Research Letters*, *48*(11),
886 e2020GL091692.
- 887 Kerry, C. G., Powell, B. S., & Carter, G. S. (2016). Quantifying the incoherent M_2
888 internal tide in the Philippine Sea. *Journal of Physical Oceanography*, *46*(8),
889 2483–2491.
- 890 Köhler, J., Walter, M., Mertens, C., Stiehler, J., Li, Z., Zhao, Z., . . . Rhein, M.
891 (2019). Energy flux observations in an internal tide beam in the eastern North
892 Atlantic. *Journal of Geophysical Research: Oceans*, *124*(8), 5747–5764.
- 893 Kunze, E., & Toole, J. M. (1997). Tidally driven vorticity, diurnal shear, and tur-
894 bulence atop Fieberling Seamount. *Journal of Physical Oceanography*, *27*(12),
895 2663–2693.
- 896 Kurian, J., Colas, F., Capet, X., McWilliams, J. C., & Chelton, D. B. (2011). Eddy
897 properties in the California Current System. *Journal of Geophysical Research:
898 Oceans*, *116*(C8).
- 899 Lerczak, J. A., Winant, C., & Hendershott, M. (2003). Observations of the semi-
900 diurnal internal tide on the southern California slope and shelf. *Journal of Geo-
901 physical Research: Oceans*, *108*(C3).
- 902 Löb, J., Köhler, J., Mertens, C., Walter, M., Li, Z., von Storch, J.-S., . . . Rhein, M.
903 (2020). Observations of the low-mode internal tide and its interaction with
904 mesoscale flow south of the Azores. *Journal of Geophysical Research: Oceans*,
905 *125*(11), e2019JC015879.
- 906 Martini, K. I., Alford, M. H., Kunze, E., Kelly, S. M., & Nash, J. D. (2011). Ob-
907 servations of internal tides on the oregon continental slope. *Journal of Physical
908 Oceanography*, *41*(9), 1772–1794.
- 909 Mazloff, M. R., Cornuelle, B., Gille, S. T., & Wang, J. (2020). The importance of
910 remote forcing for regional modeling of internal waves. *Journal of Geophysical
911 Research: Oceans*, *125*(2), e2019JC015623.
- 912 McDougall, T. J., & Barker, P. M. (2011). Getting started with TEOS-10 and the
913 Gibbs Seawater (GSW) oceanographic toolbox. *Scor/Iapso WG*, *127*(532), 1–
914 28.
- 915 Melet, A. V., Hallberg, R., & Marshall, D. P. (2022). The role of ocean mixing in
916 the climate system. In *Ocean mixing* (pp. 5–34). Elsevier.
- 917 Munk, W. (1981). Internal waves and small-scale processes. *Evolution of physical
918 oceanography*.
- 919 Musgrave, R., MacKinnon, J., Pinkel, R., Waterhouse, A., Nash, J., & Kelly, S.
920 (2017). The influence of subinertial internal tides on near-topographic turbu-
921 lence at the Mendocino Ridge: Observations and modeling. *Journal of Physical
922 Oceanography*, *47*(8), 2139–2154.
- 923 Nash, J. D., Alford, M. H., & Kunze, E. (2005). Estimating internal wave energy
924 fluxes in the ocean. *Journal of Atmospheric and Oceanic Technology*, *22*(10),
925 1551–1570.
- 926 Nelson, A. D., Arbic, B. K., Zaron, E. D., Savage, A. C., Richman, J. G., Buijsman,
927 M. C., & Shriver, J. F. (2019). Toward realistic nonstationarity of semidiurnal
928 baroclinic tides in a hydrodynamic model. *Journal of Geophysical Research:
929 Oceans*, *124*(9), 6632–6642.
- 930 Pawlowicz, R., Beardsley, B., & Lentz, S. (2002). Classical tidal harmonic anal-
931 ysis including error estimates in MATLAB using T_TIDE. *Computers & Geo-
932 sciences*, *28*(8), 929–937.
- 933 Ponte, A. L., & Klein, P. (2015). Incoherent signature of internal tides on sea level
934 in idealized numerical simulations. *Geophysical Research Letters*, *42*(5), 1520–
935 1526.
- 936 Qiu, B., Chen, S., Klein, P., Sasaki, H., & Sasai, Y. (2014). Seasonal mesoscale and
937 submesoscale eddy variability along the North Pacific Subtropical Countercur-
938 rent. *Journal of Physical Oceanography*, *44*(12), 3079–3098.

- 939 Qiu, B., Chen, S., Klein, P., Wang, J., Torres, H., Fu, L.-L., & Menemenlis, D.
 940 (2018). Seasonality in transition scale from balanced to unbalanced motions in
 941 the world ocean. *Journal of Physical Oceanography*, *48*(3), 591–605.
- 942 Rainville, L., Johnston, T. S., Carter, G. S., Merrifield, M. A., Pinkel, R., Worcester,
 943 P. F., & Dushaw, B. D. (2010). Interference pattern and propagation of the M
 944 2 internal tide south of the Hawaiian Ridge. *Journal of physical oceanography*,
 945 *40*(2), 311–325.
- 946 Rainville, L., & Pinkel, R. (2006). Propagation of low-mode internal waves through
 947 the ocean. *Journal of Physical Oceanography*, *36*(6), 1220–1236.
- 948 Ray, R. D., & Zaron, E. D. (2016). M₂ internal tides and their observed wavenum-
 949 ber spectra from satellite altimetry. *Journal of Physical Oceanography*, *46*(1),
 950 3–22.
- 951 Rudnick, D. L., Zaba, K. D., Todd, R. E., & Davis, R. E. (2017). A climatology of
 952 the california current system from a network of underwater gliders. *Progress in*
 953 *Oceanography*, *154*, 64–106.
- 954 Sasaki, H., Klein, P., Qiu, B., & Sasai, Y. (2014). Impact of oceanic-scale inter-
 955 actions on the seasonal modulation of ocean dynamics by the atmosphere. *Nature*
 956 *communications*, *5*(1), 5636.
- 957 Savage, A. C., Waterhouse, A. F., & Kelly, S. M. (2020). Internal tide nonstation-
 958 arity and wave–mesoscale interactions in the tasman sea. *Journal of Physical*
 959 *Oceanography*, *50*(10), 2931–2951.
- 960 Shakespeare, C. J. (2023). Eddy acceleration and decay driven by internal tides.
 961 *Journal of Physical Oceanography*.
- 962 Sharples, J., Tweddle, J. F., Mattias Green, J., Palmer, M. R., Kim, Y.-N., Hick-
 963 man, A. E., ... others (2007). Spring-neap modulation of internal tide mixing
 964 and vertical nitrate fluxes at a shelf edge in summer. *Limnology and Oceanog-*
 965 *raphy*, *52*(5), 1735–1747.
- 966 Shriver, J. F., Richman, J. G., & Arbic, B. K. (2014). How stationary are the inter-
 967 nal tides in a high-resolution global ocean circulation model? *Journal of Geo-*
 968 *physical Research: Oceans*, *119*(5), 2769–2787.
- 969 Siyanbola, O. Q., Buijsman, M. C., Delpech, A., Renault, L., Barkan, R., Shriver,
 970 J. F., ... McWilliams, J. C. (2023). Remote internal wave forcing of regional
 971 ocean simulations near the US West Coast. *Ocean Modelling*, *181*, 102154.
- 972 Terker, S. R., Girton, J. B., Kunze, E., Klymak, J. M., & Pinkel, R. (2014). Obser-
 973 vations of the internal tide on the California continental margin near Monterey
 974 Bay. *Continental Shelf Research*, *82*, 60–71.
- 975 Thomson, D. J. (1982). Spectrum estimation and harmonic analysis. *Proceedings of*
 976 *the IEEE*, *70*(9), 1055–1096.
- 977 Wang, J., Fu, L.-L., Haines, B., Lankhorst, M., Lucas, A. J., Farrar, J. T., ... oth-
 978 ers (2022). On the Development of SWOT In Situ Calibration/Validation for
 979 Short-Wavelength Ocean Topography. *Journal of Atmospheric and Oceanic*
 980 *Technology*, *39*(5), 595–617.
- 981 Wang, Y., & Legg, S. (2023). Enhanced dissipation of internal tides in a mesoscale
 982 baroclinic eddy. *Journal of Physical Oceanography*, *53*(10), 2293–2316.
- 983 Wunsch, C. (1975). Internal tides in the ocean. *Reviews of Geophysics*, *13*(1), 167–
 984 182.
- 985 Xu, Z., Liu, K., Yin, B., Zhao, Z., Wang, Y., & Li, Q. (2016). Long-range propaga-
 986 tion and associated variability of internal tides in the South China Sea. *Jour-*
 987 *nal of Geophysical Research: Oceans*, *121*(11), 8268–8286.
- 988 Zaron, E. D., & Egbert, G. D. (2014). Time-variable refraction of the internal tide
 989 at the Hawaiian Ridge. *Journal of physical oceanography*, *44*(2), 538–557.
- 990 Zaron, E. D., & Ray, R. D. (2017). Using an altimeter-derived internal tide model
 991 to remove tides from in situ data. *Geophysical Research Letters*, *44*(9), 4241–
 992 4245.
- 993 Zhao, Z. (2016). Using CryoSat-2 altimeter data to evaluate M₂ internal tides ob-

- 994 served from multisatellite altimetry. *Journal of Geophysical Research: Oceans*,
 995 *121*(7), 5164–5180.
- 996 Zhao, Z. (2018). The global mode-2 M_2 internal tide. *Journal of Geophysical Re-*
 997 *search: Oceans*, *123*(11), 7725–7746.
- 998 Zhao, Z. (2019). Mapping internal tides from satellite altimetry without blind direc-
 999 tions. *Journal of Geophysical Research: Oceans*, *124*(12), 8605–8625.
- 1000 Zhao, Z. (2021). Seasonal mode-1 M_2 internal tides from satellite altimetry. *Journal*
 1001 *of Physical Oceanography*, *51*(9), 3015–3035.
- 1002 Zhao, Z. (2022). Development of the yearly mode-1 M_2 internal tide model in 2019.
 1003 *Journal of Atmospheric and Oceanic Technology*, *39*(4), 463–478.
- 1004 Zhao, Z., Alford, M. H., Girton, J. B., Rainville, L., & Simmons, H. L. (2016).
 1005 Global observations of open-ocean mode-1 M_2 internal tides. *Journal of Physi-*
 1006 *cal Oceanography*, *46*(6), 1657–1684.
- 1007 Zhao, Z., Alford, M. H., Lien, R.-C., Gregg, M. C., & Carter, G. S. (2012). Internal
 1008 tides and mixing in a submarine canyon with time-varying stratification. *Jour-*
 1009 *nal of Physical Oceanography*, *42*(12), 2121–2142.
- 1010 Zhao, Z., Alford, M. H., MacKinnon, J. A., & Pinkel, R. (2010). Long-range prop-
 1011 agation of the semidiurnal internal tide from the Hawaiian Ridge. *Journal of*
 1012 *Physical Oceanography*, *40*(4), 713–736.
- 1013 Zhao, Z., & Qiu, B. (2023). Seasonal west-east seesaw of M_2 internal tides from the
 1014 Luzon Strait. *Journal of Geophysical Research: Oceans*, e2022JC019281.
- 1015 Zhao, Z., Wang, J., Menemenlis, D., Fu, L.-L., Chen, S., & Qiu, B. (2019). Decom-
 1016 position of the multimodal multidirectional M_2 internal tide field. *Journal of*
 1017 *Atmospheric and Oceanic Technology*, *36*(6), 1157–1173.
- 1018 Zilberman, N., Merrifield, M., Carter, G., Luther, D., Levine, M., & Boyd, T. J.
 1019 (2011). Incoherent nature of M_2 internal tides at the Hawaiian Ridge. *Journal*
 1020 *of physical oceanography*, *41*(11), 2021–2036.
- 1021 Zweng, M., Seidov, D., Boyer, T., Locarnini, M., Garcia, H., Mishonov, A., . . . oth-
 1022 ers (2019). World ocean atlas 2018, volume 2: Salinity.

Figure 1.



Earthstar Geographics, Esri, HERE, Garmin, FAO, NOAA, USGS, EPA

Figure 2.

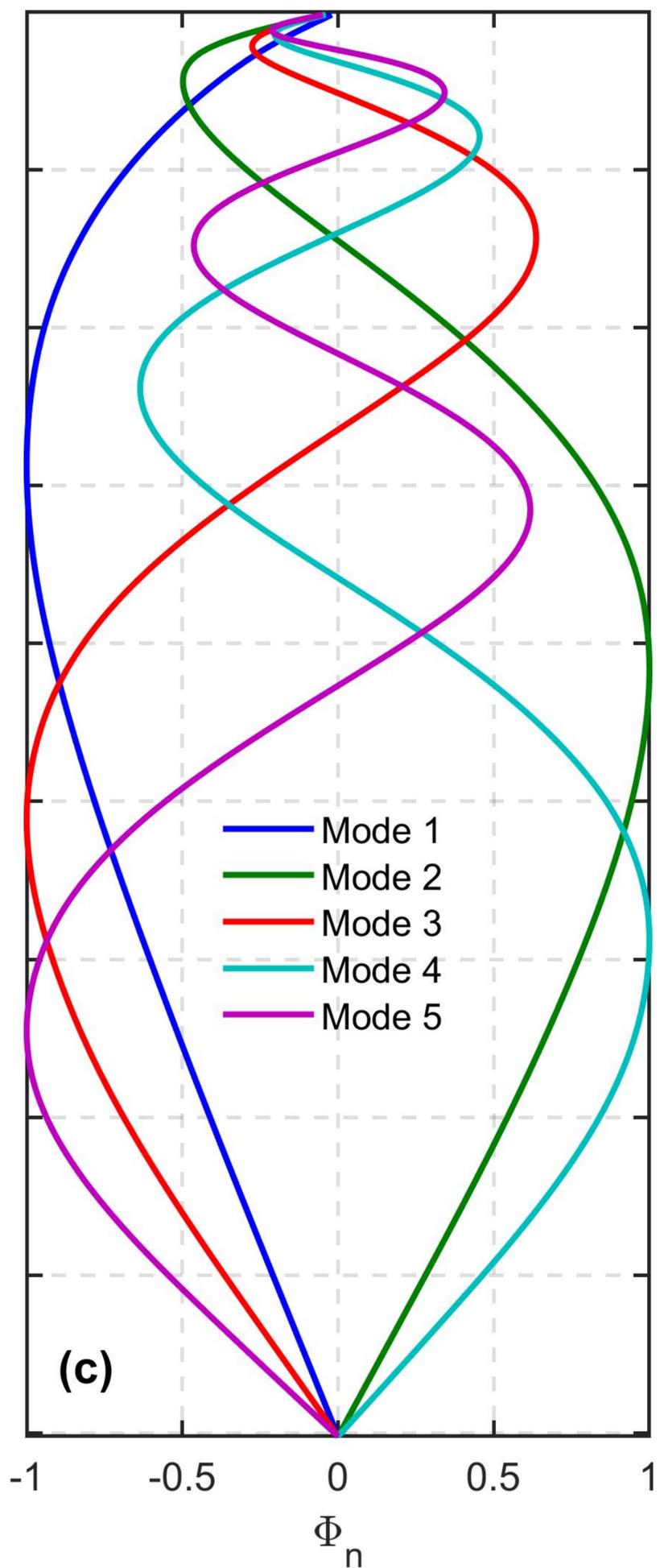
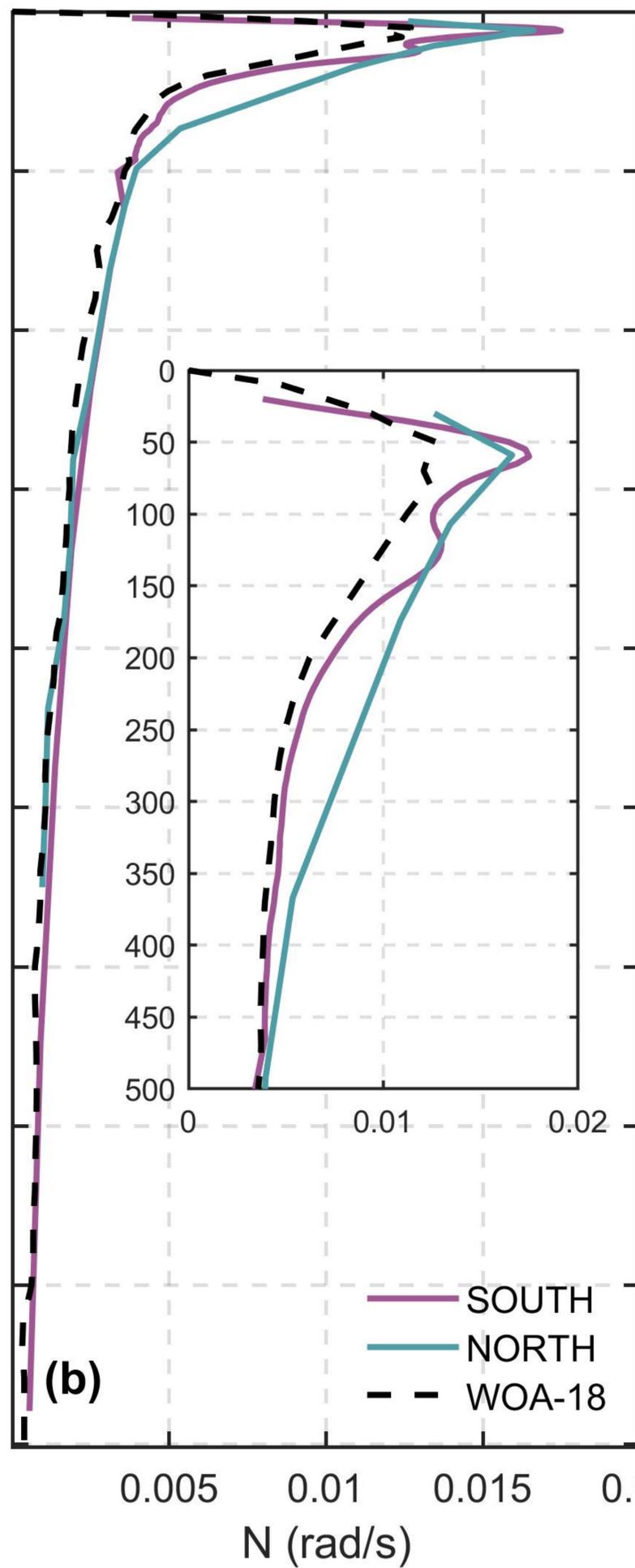
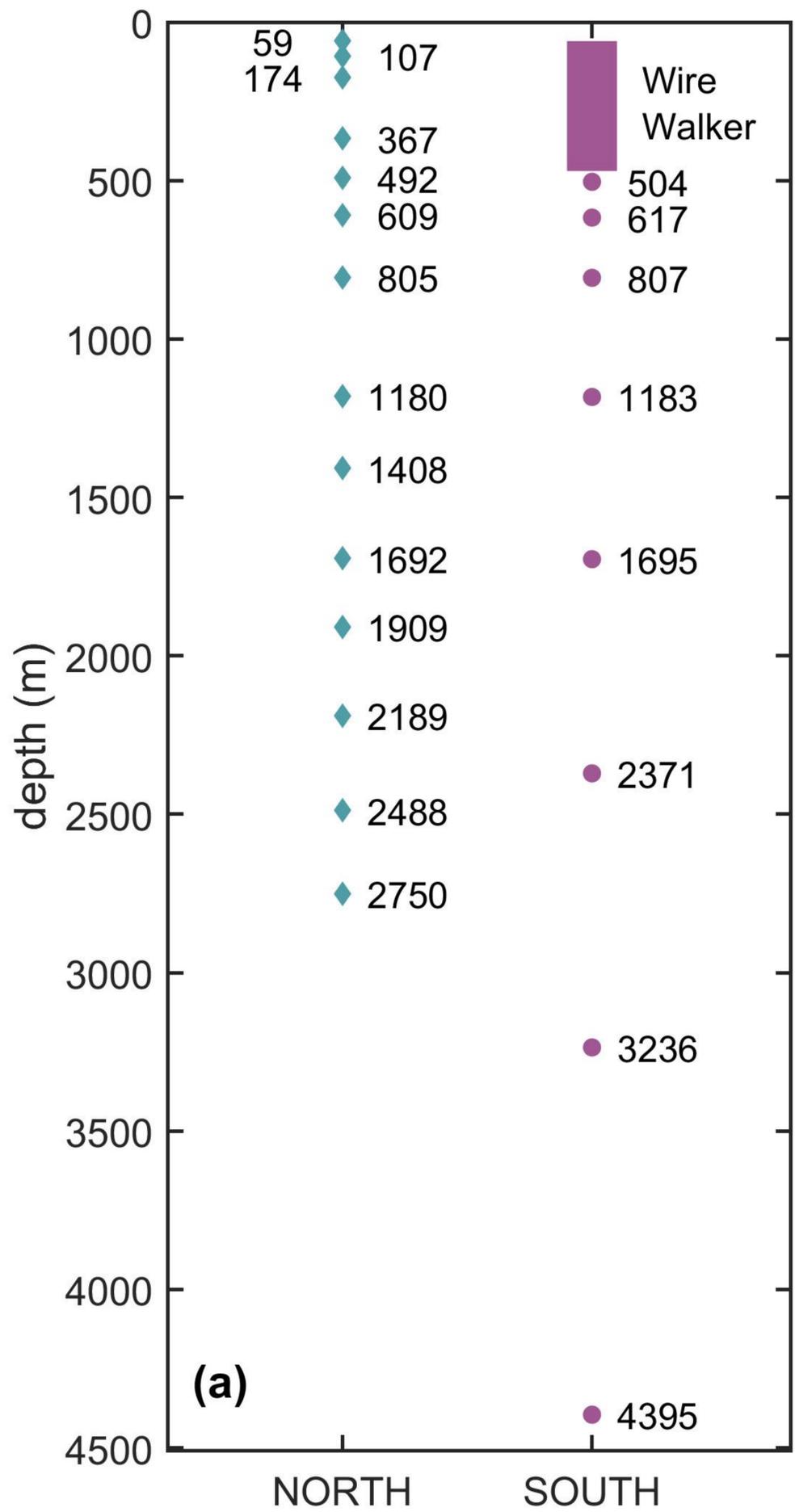


Figure 3.

North Mooring

σ (kg/m³)

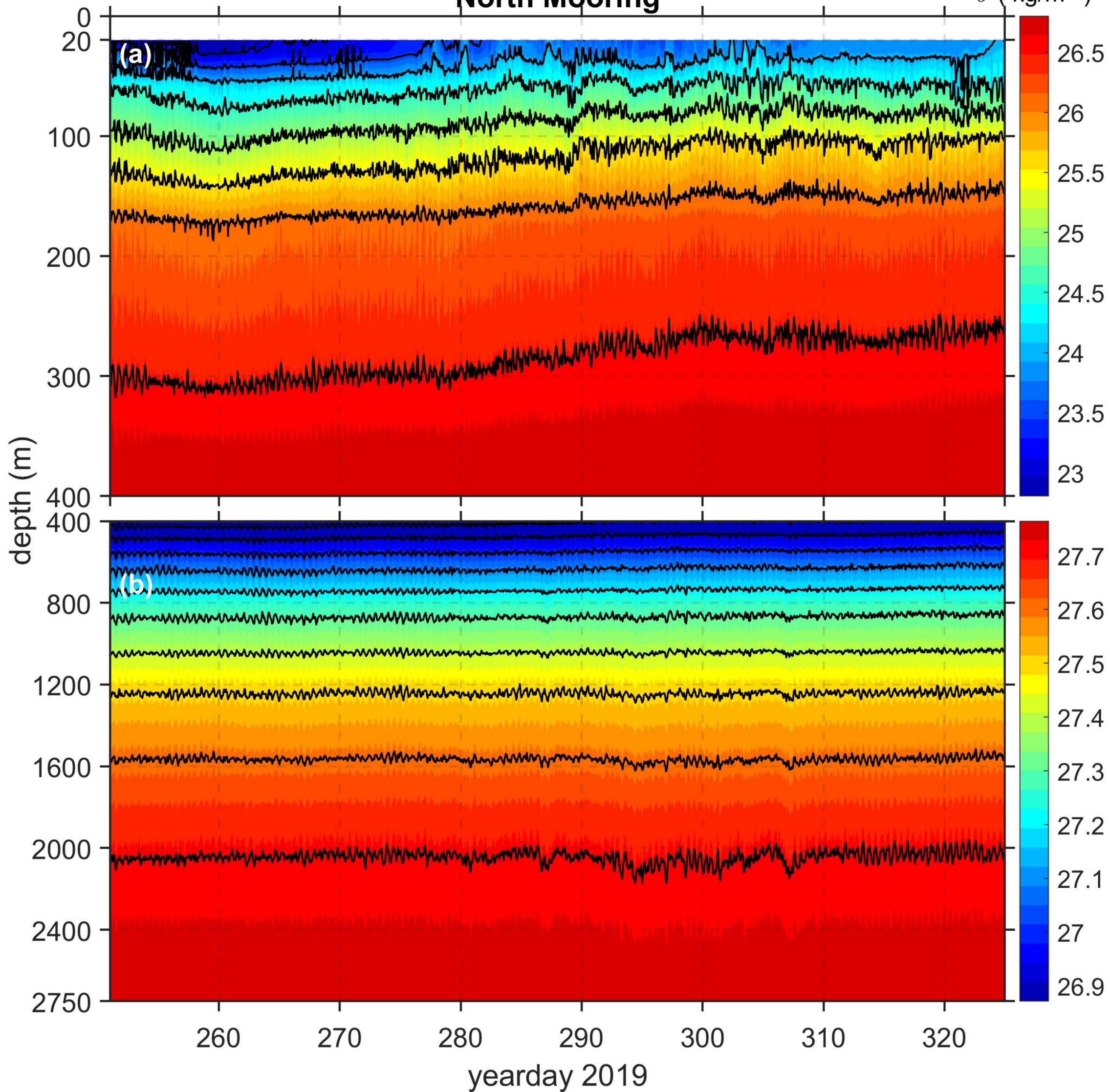


Figure 4.

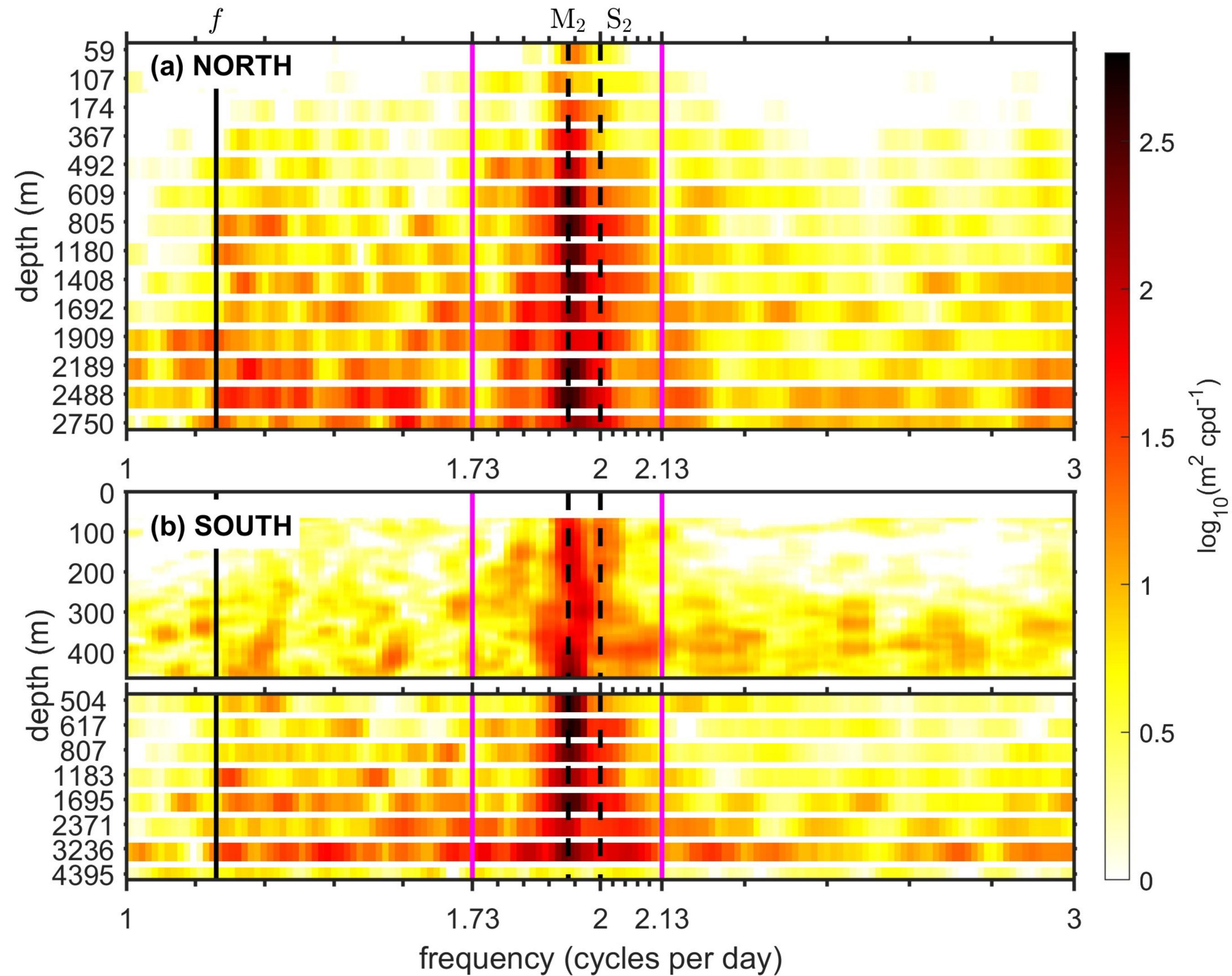


Figure 5.

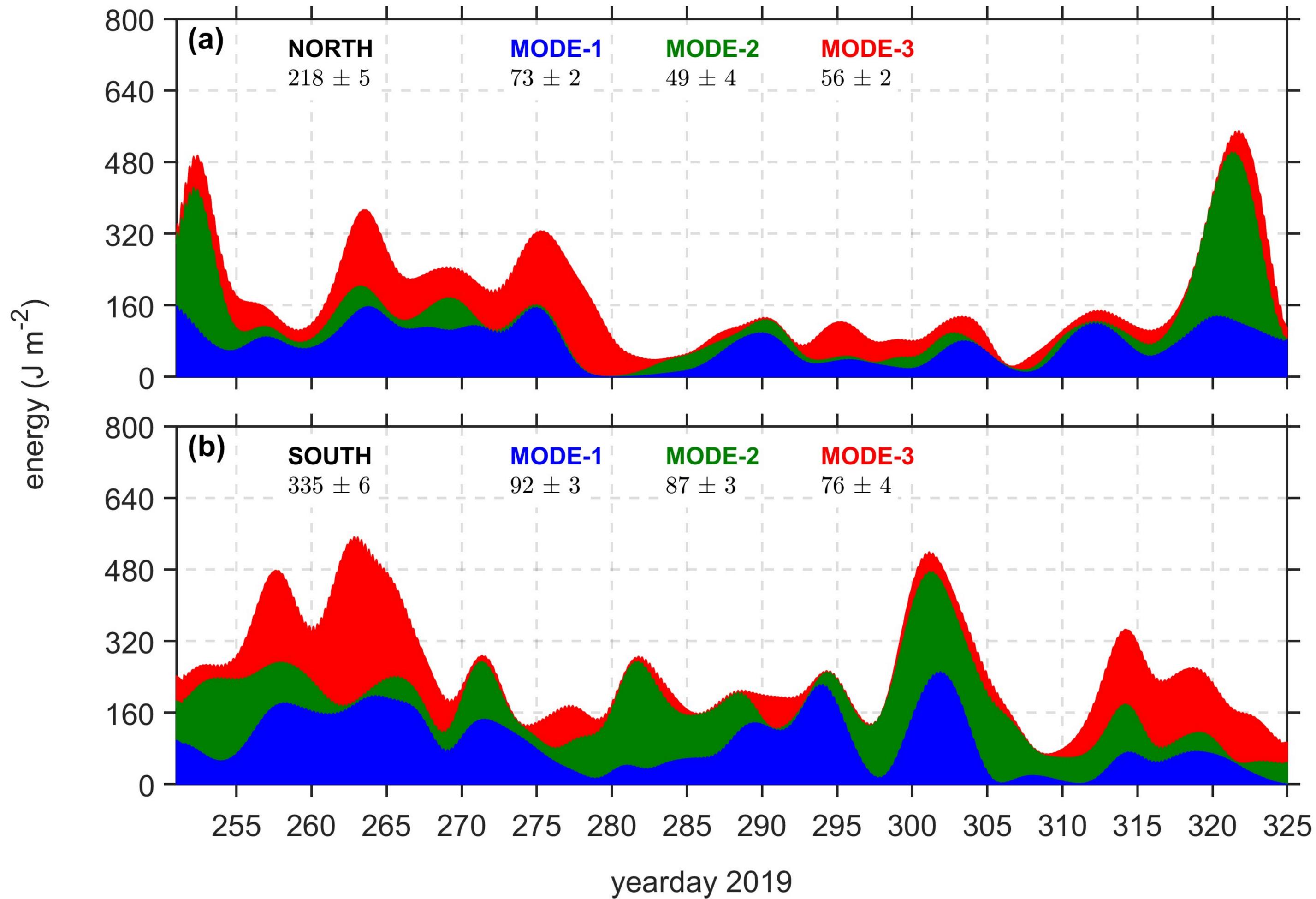


Figure 6.

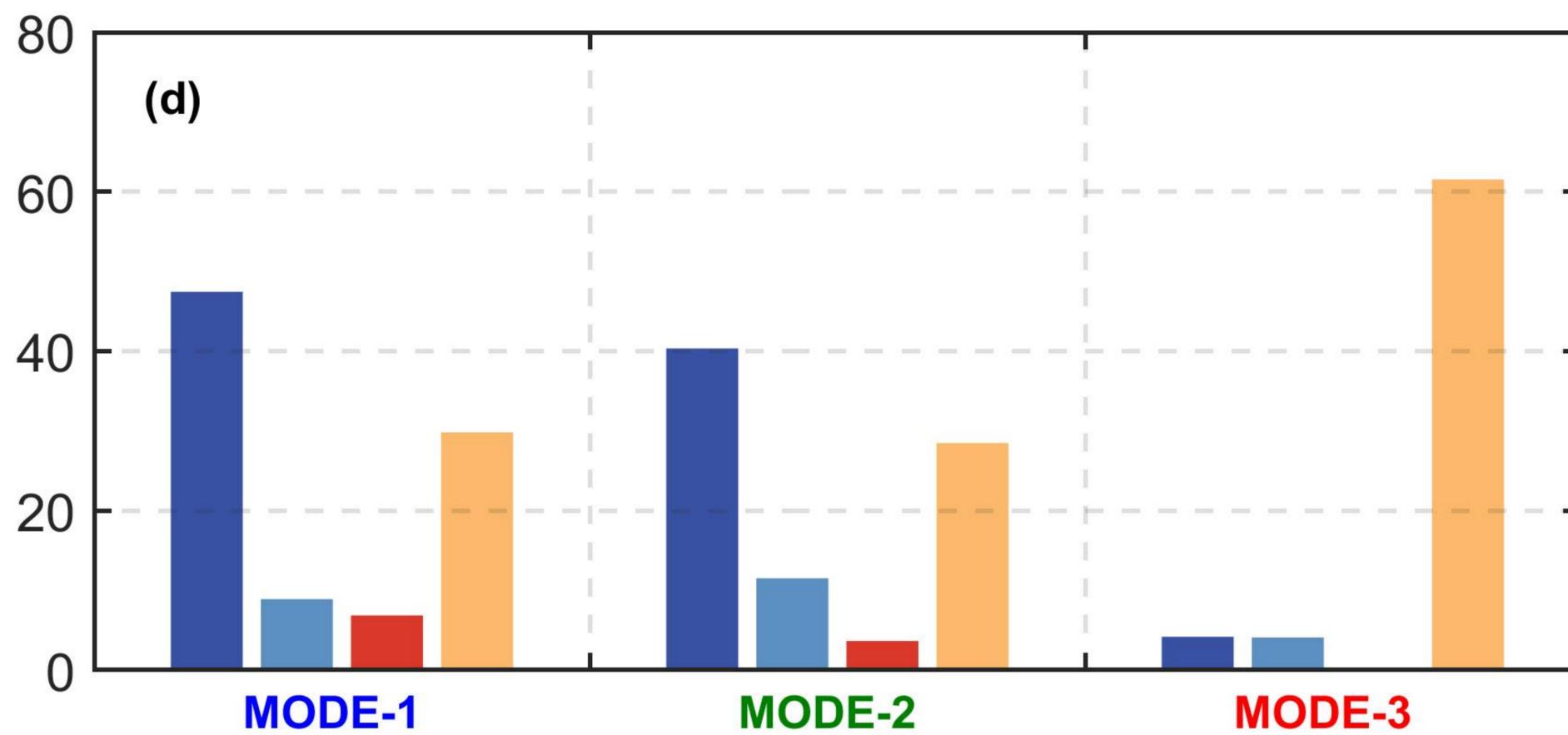
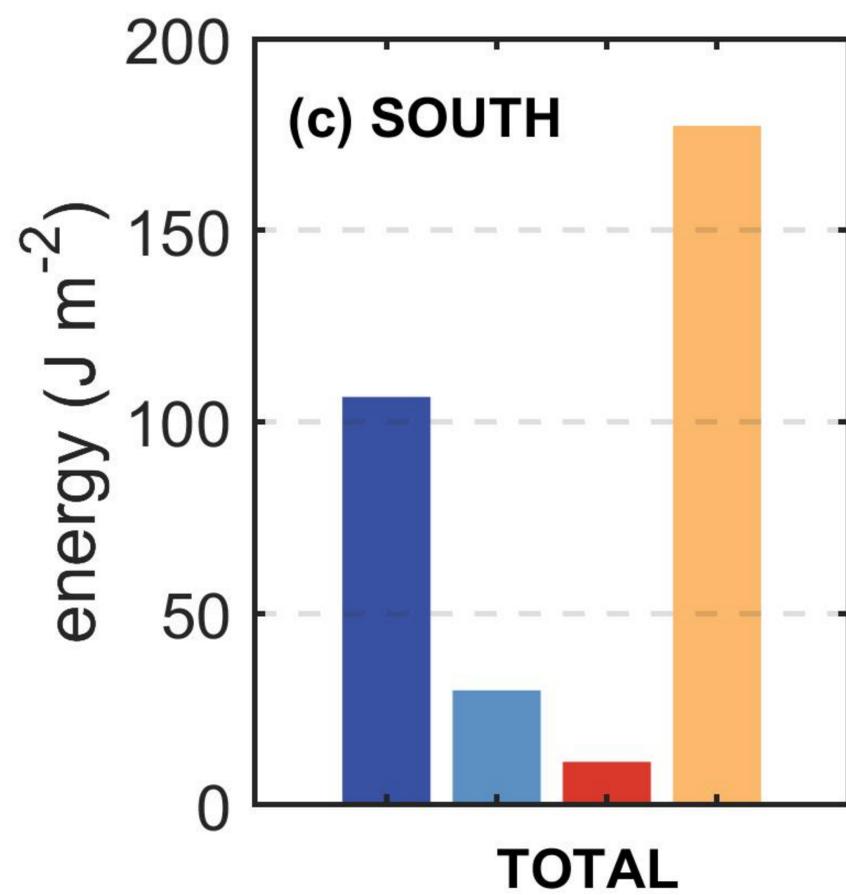
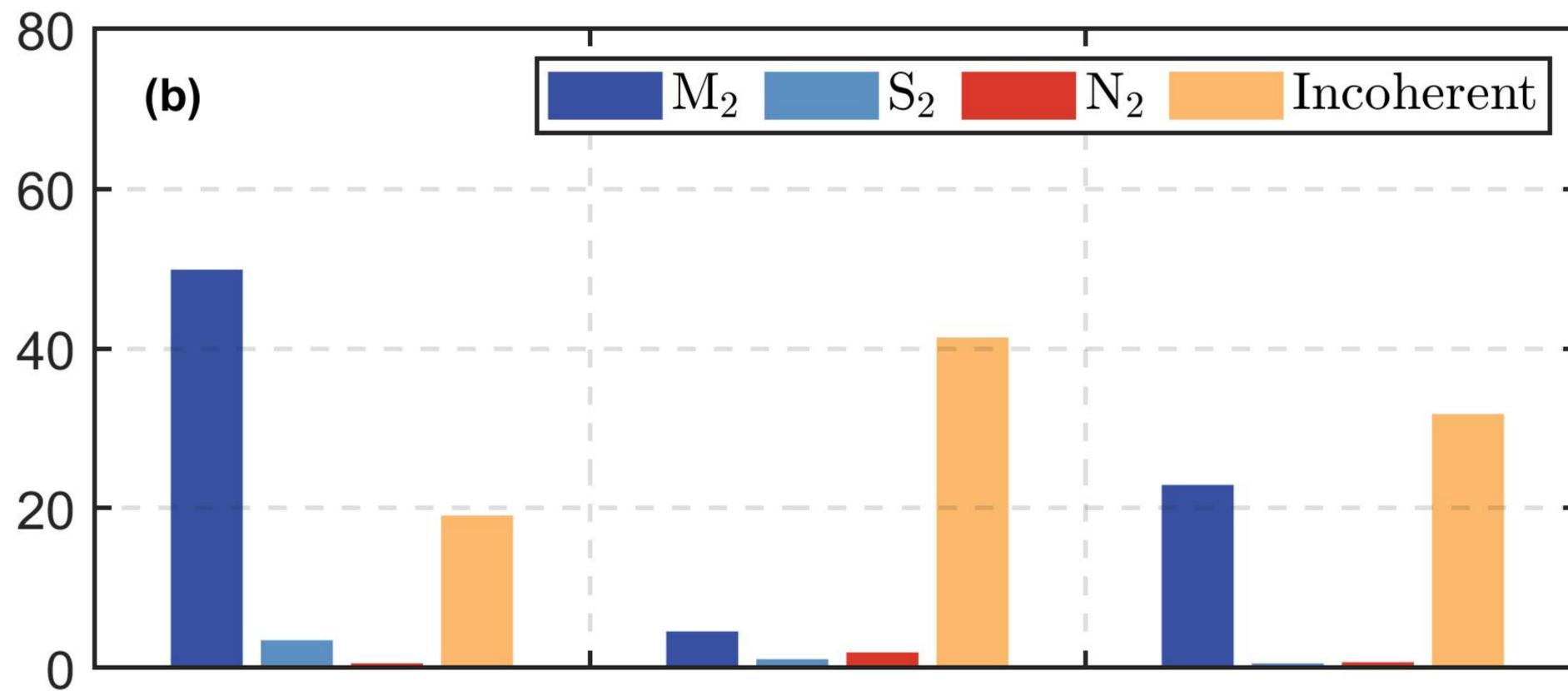
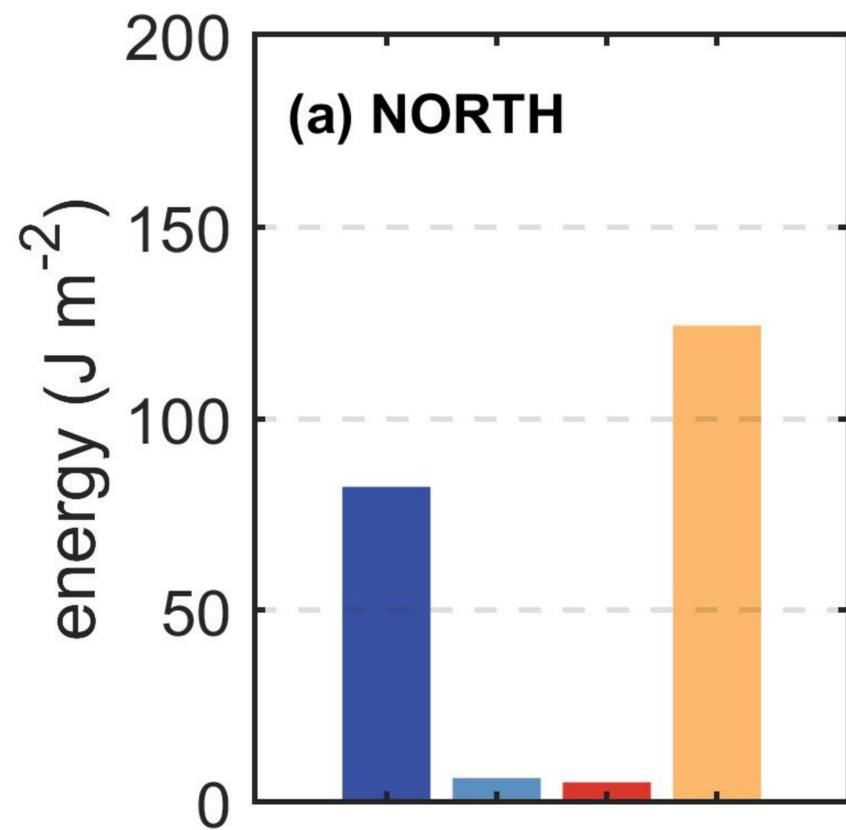
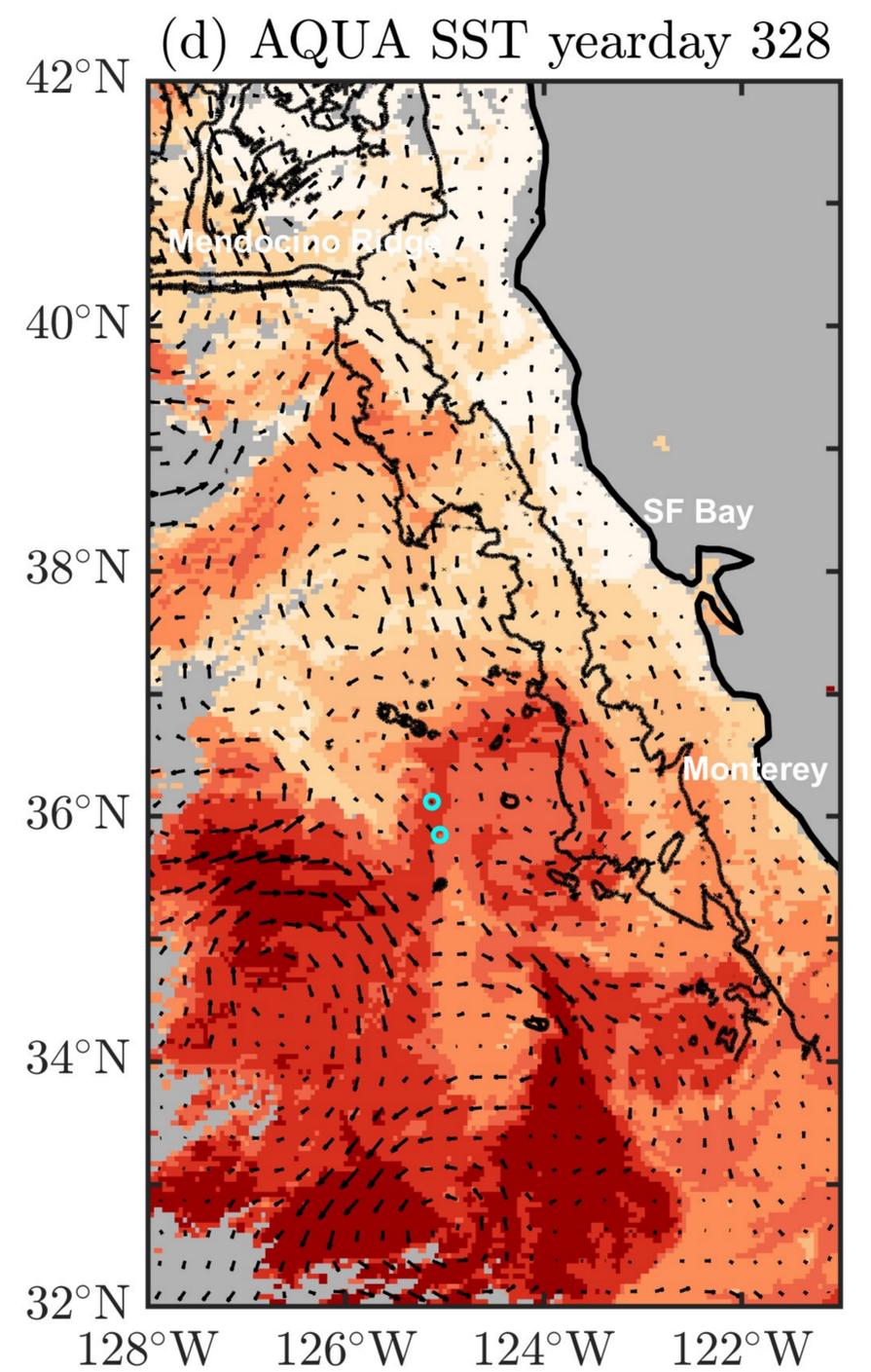
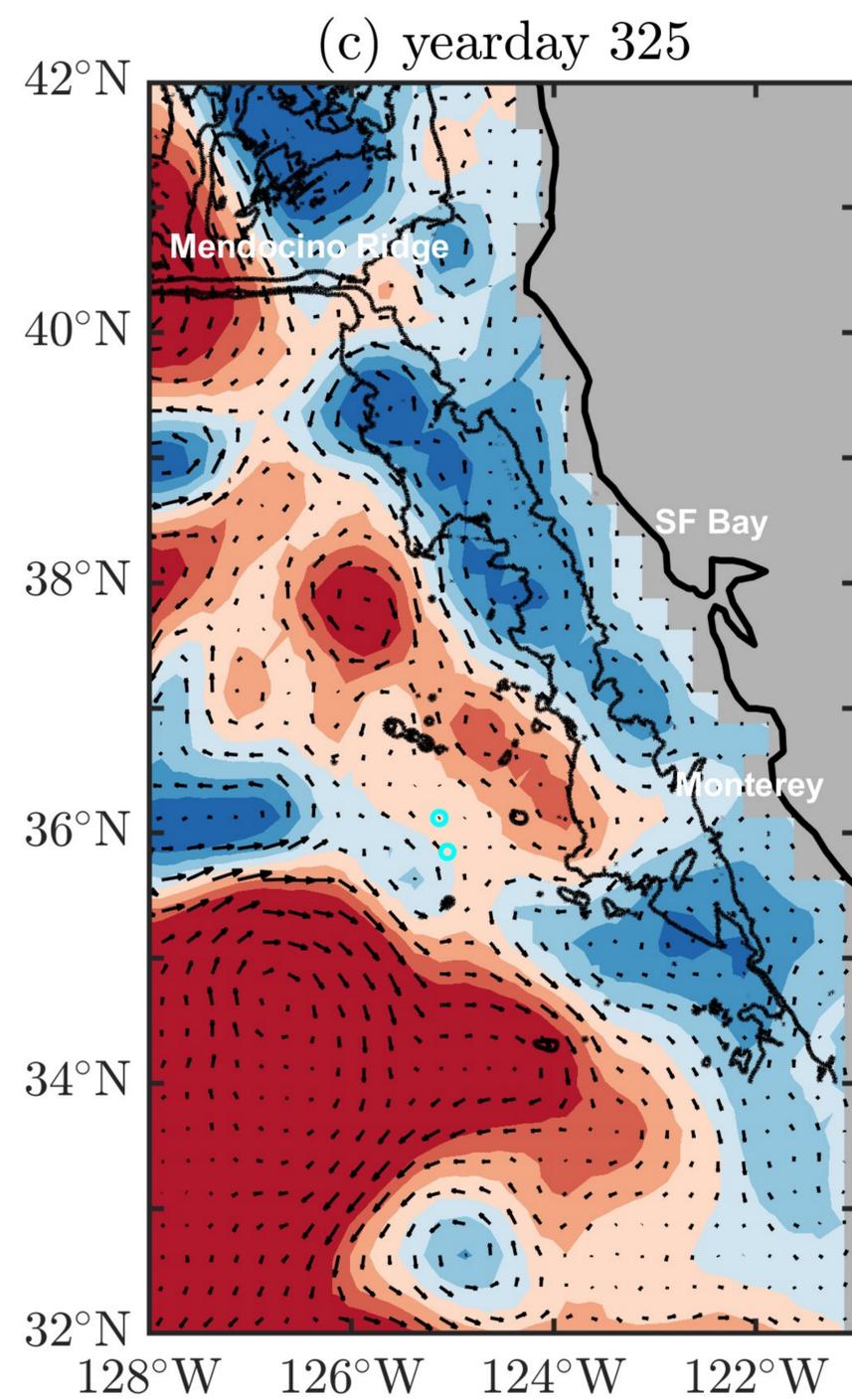
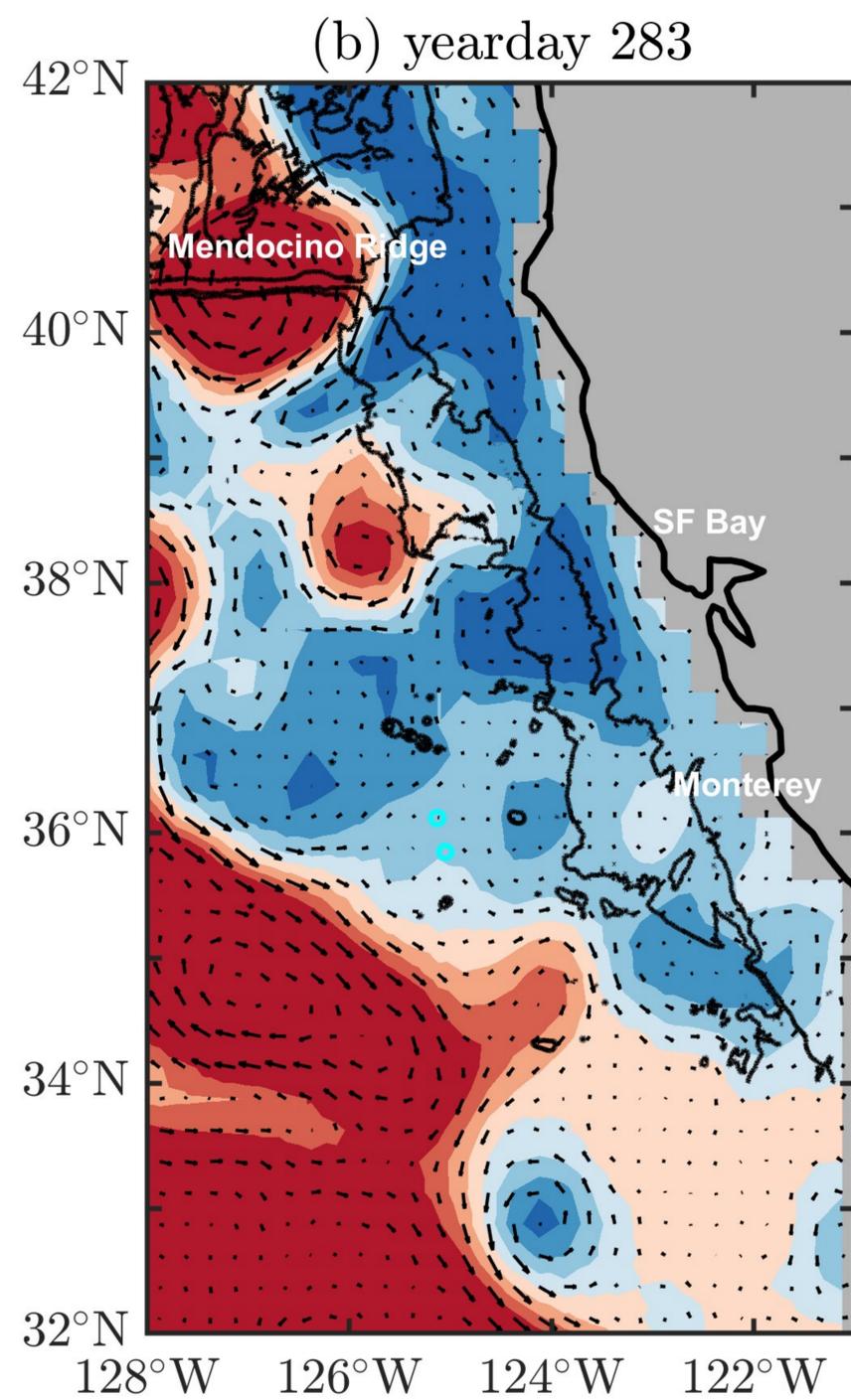
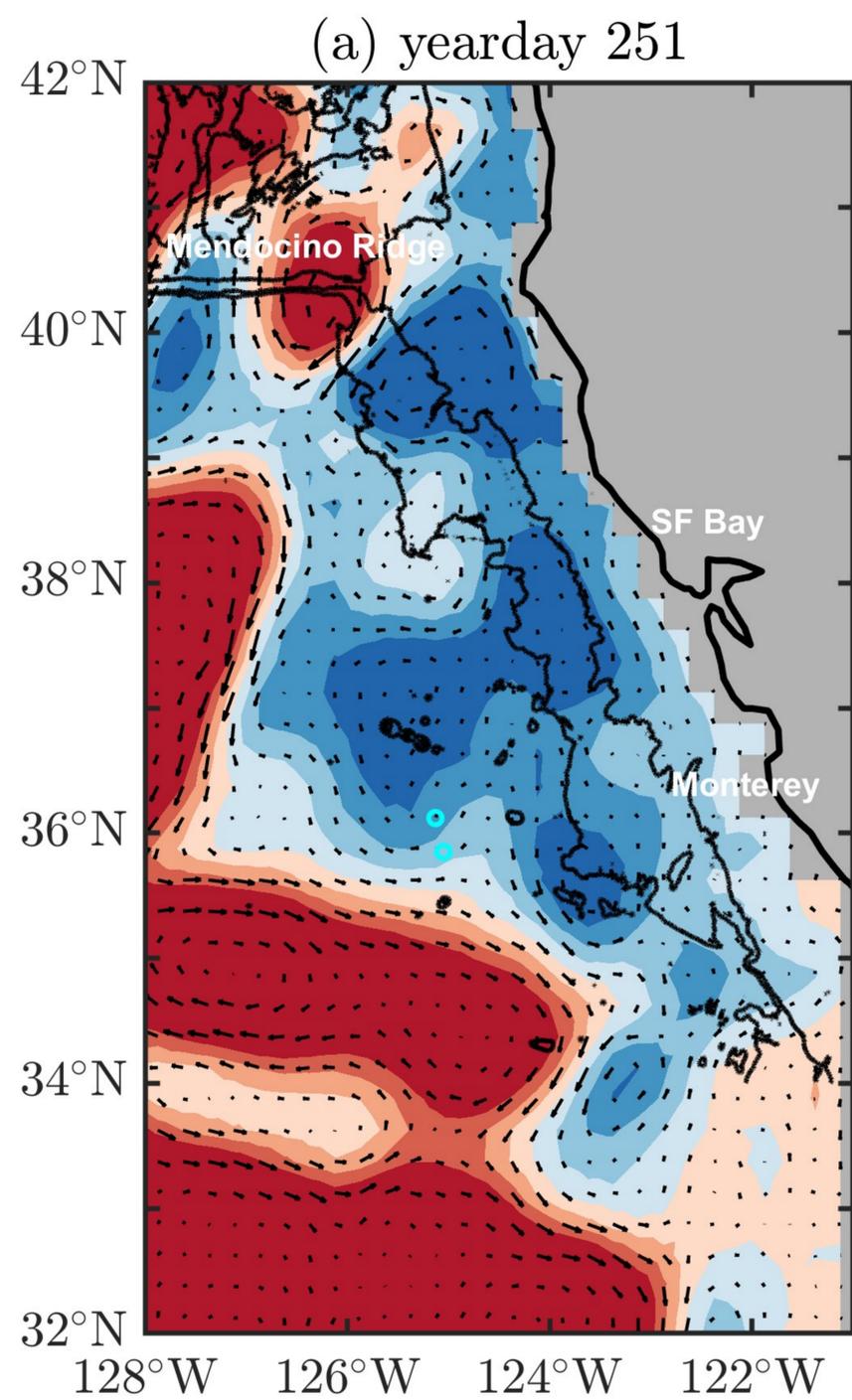


Figure 7.



0.45 0.5 0.55 0.6 0.65
ADT (m)

0.45 0.5 0.55 0.6 0.65
ADT (m)

0.45 0.5 0.55 0.6 0.65
ADT (m)

12 14 16 18
Temperature (°C)

Figure 8.

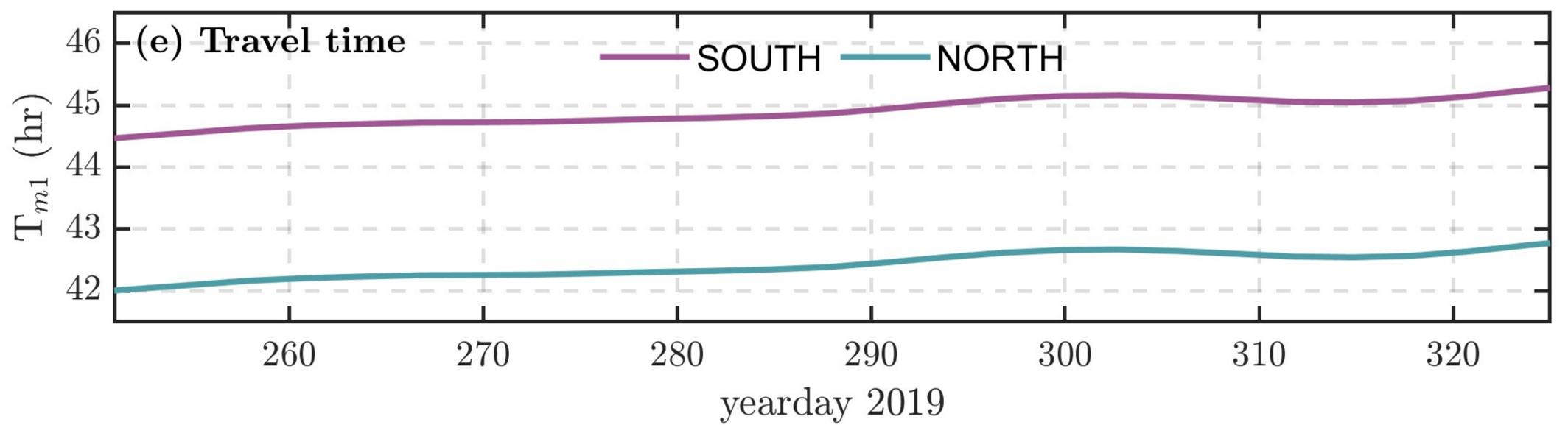
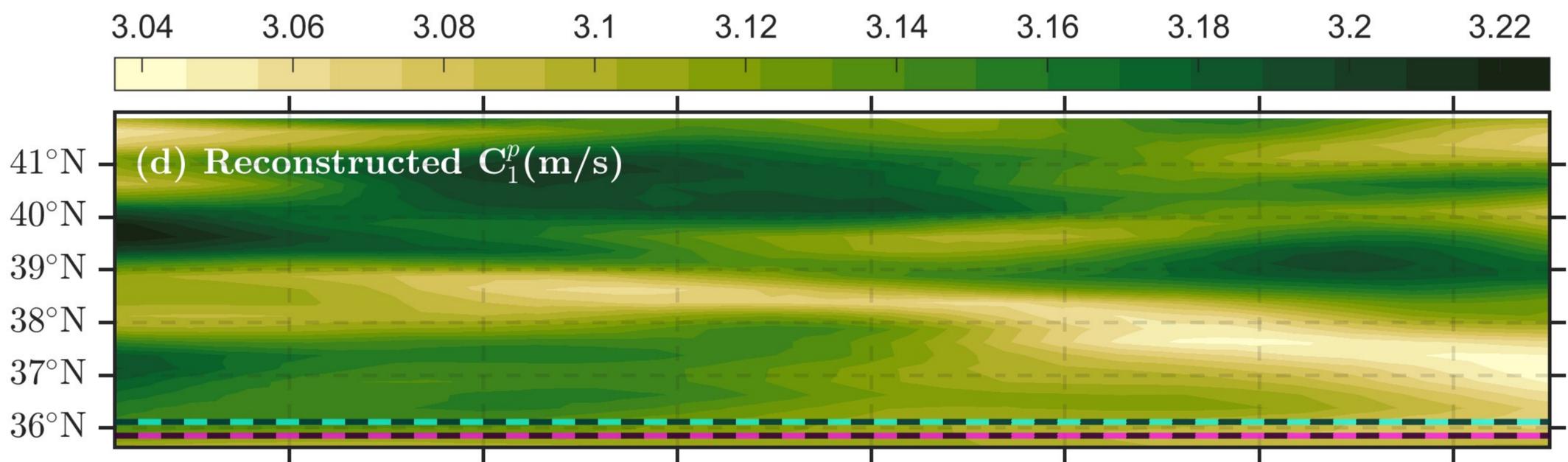
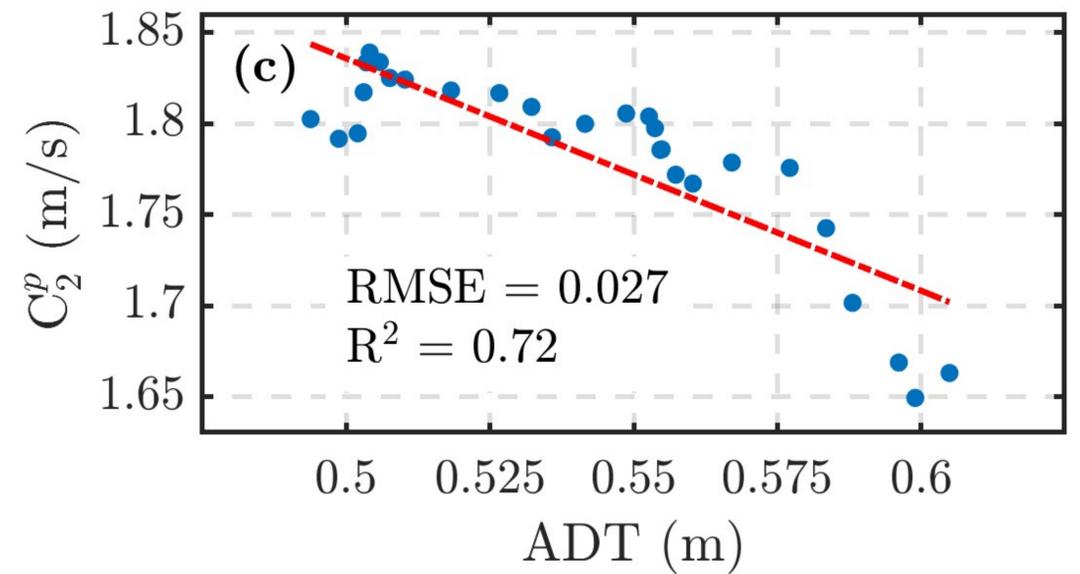
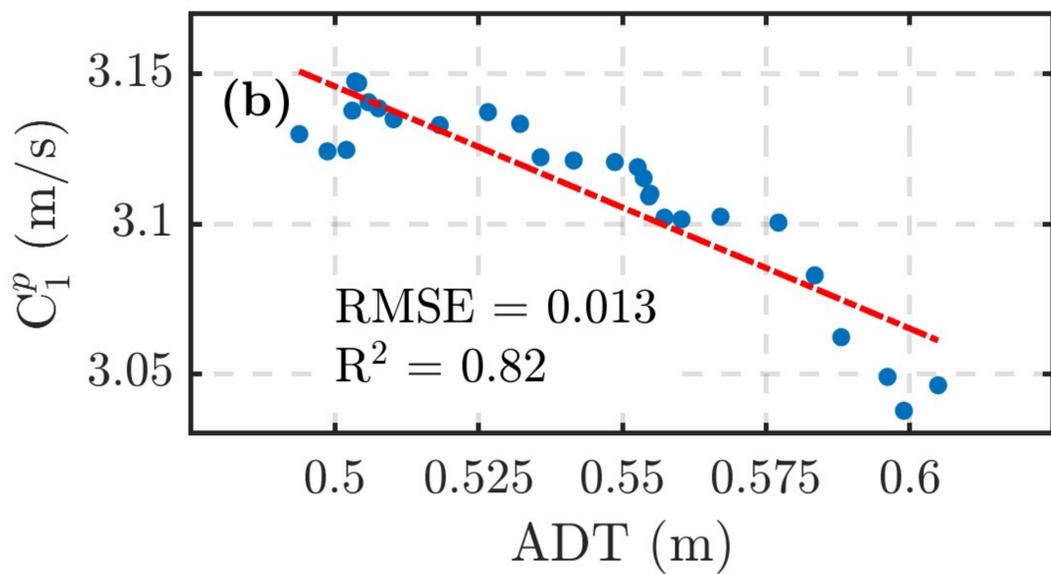
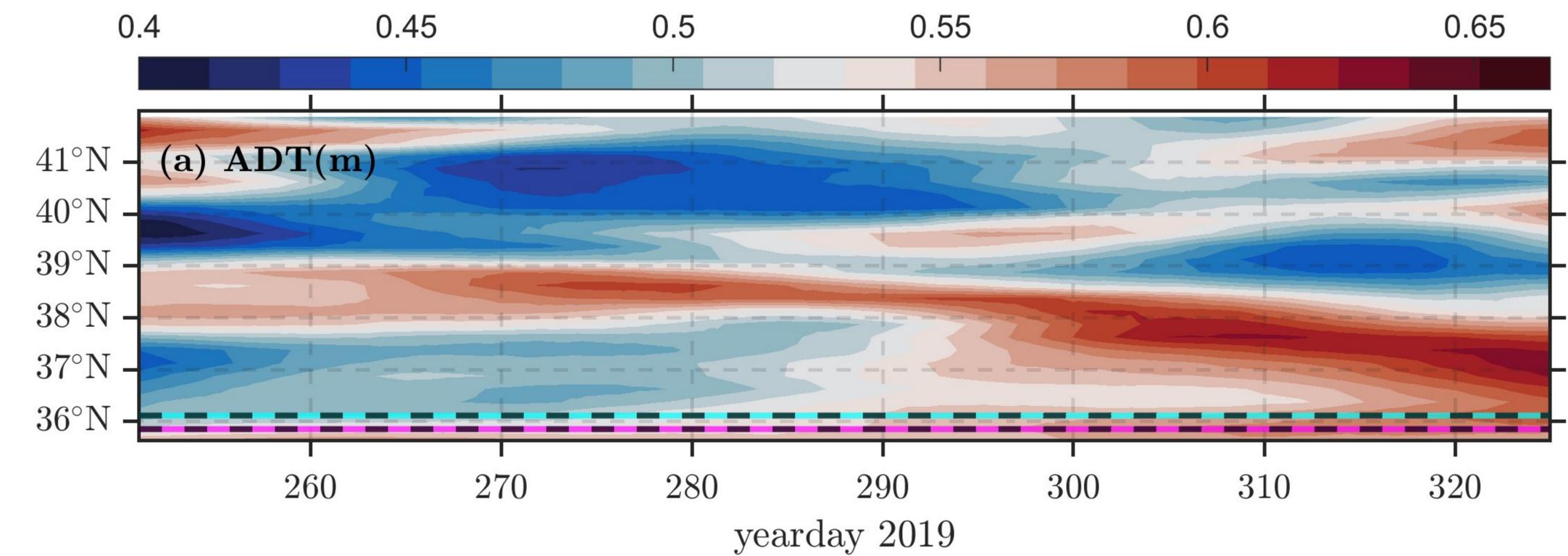
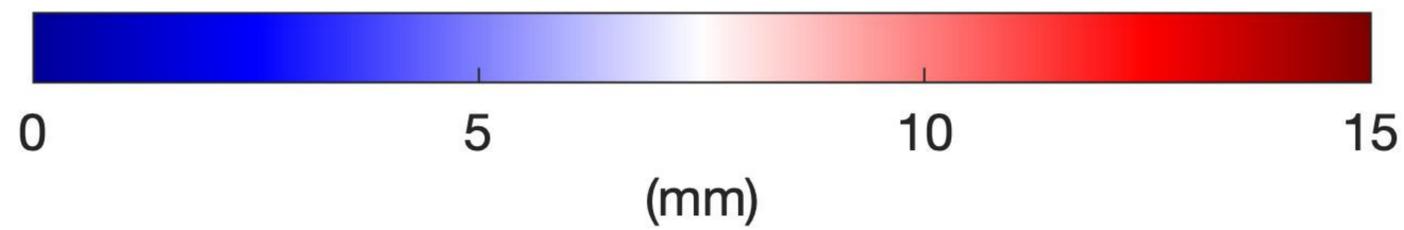
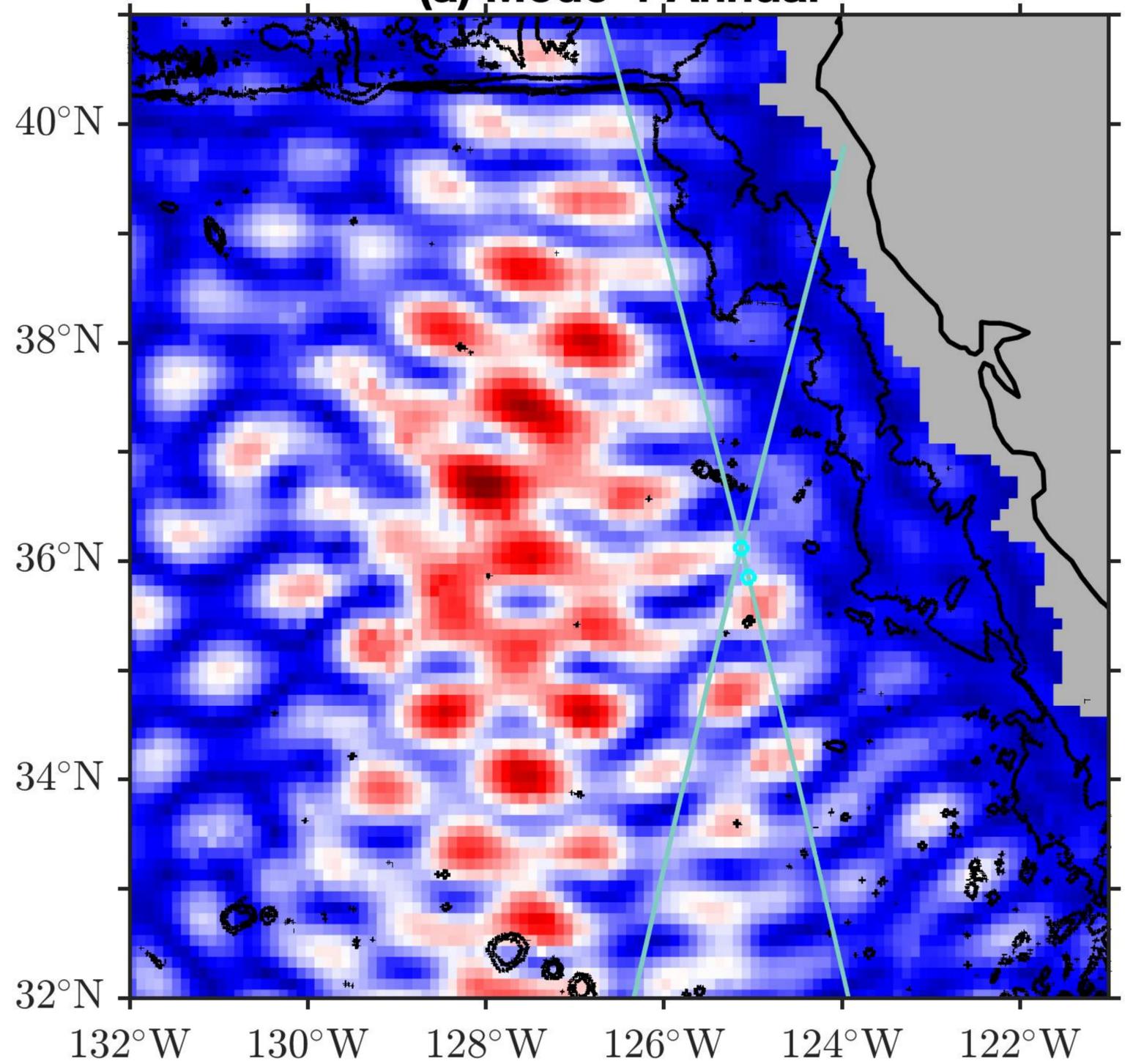


Figure 9.

(a) Mode-1 Annual



(b) Mode-2 Annual

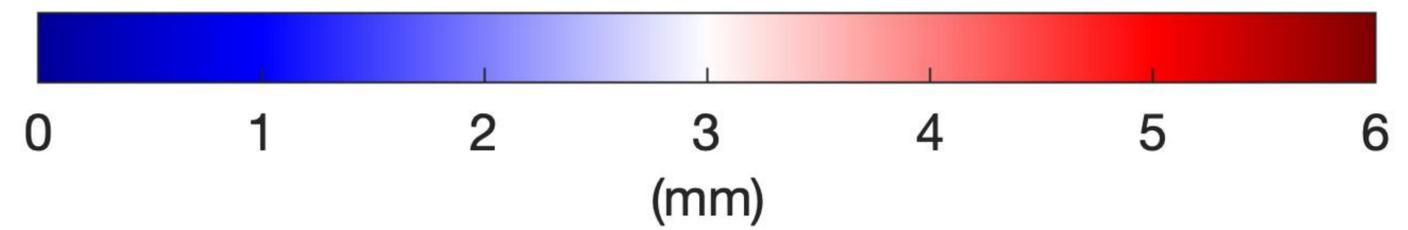
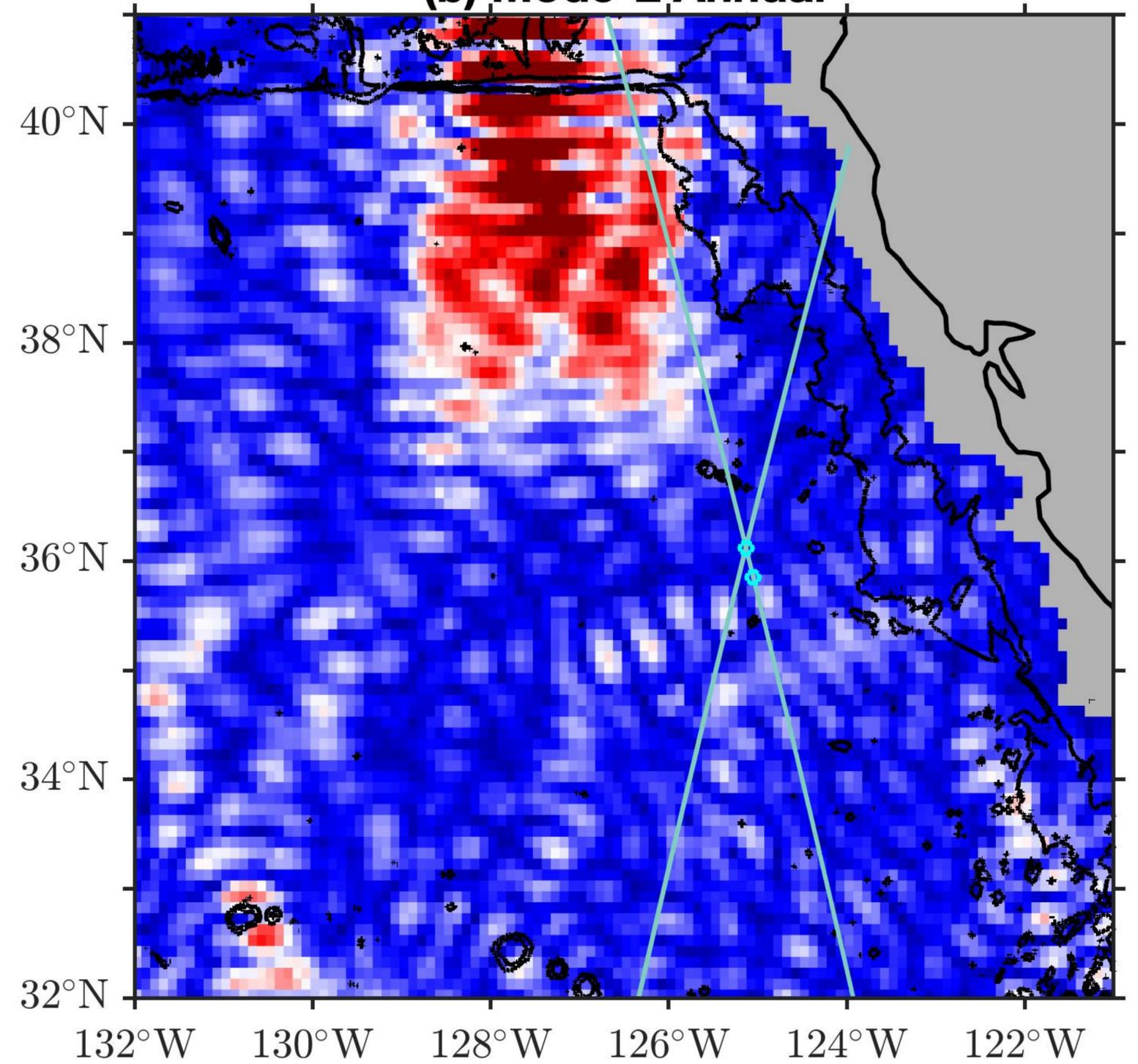


Figure 10.

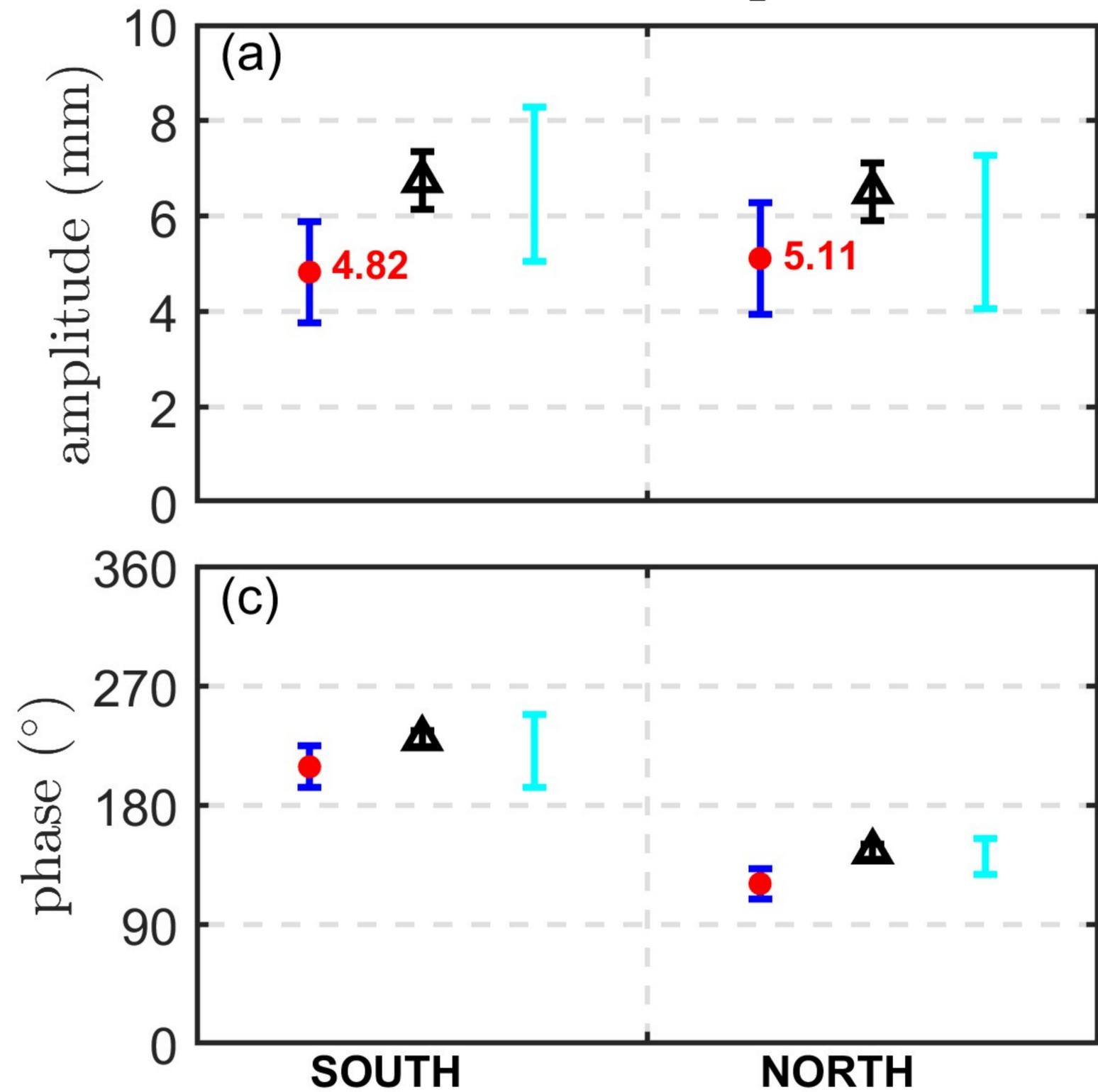
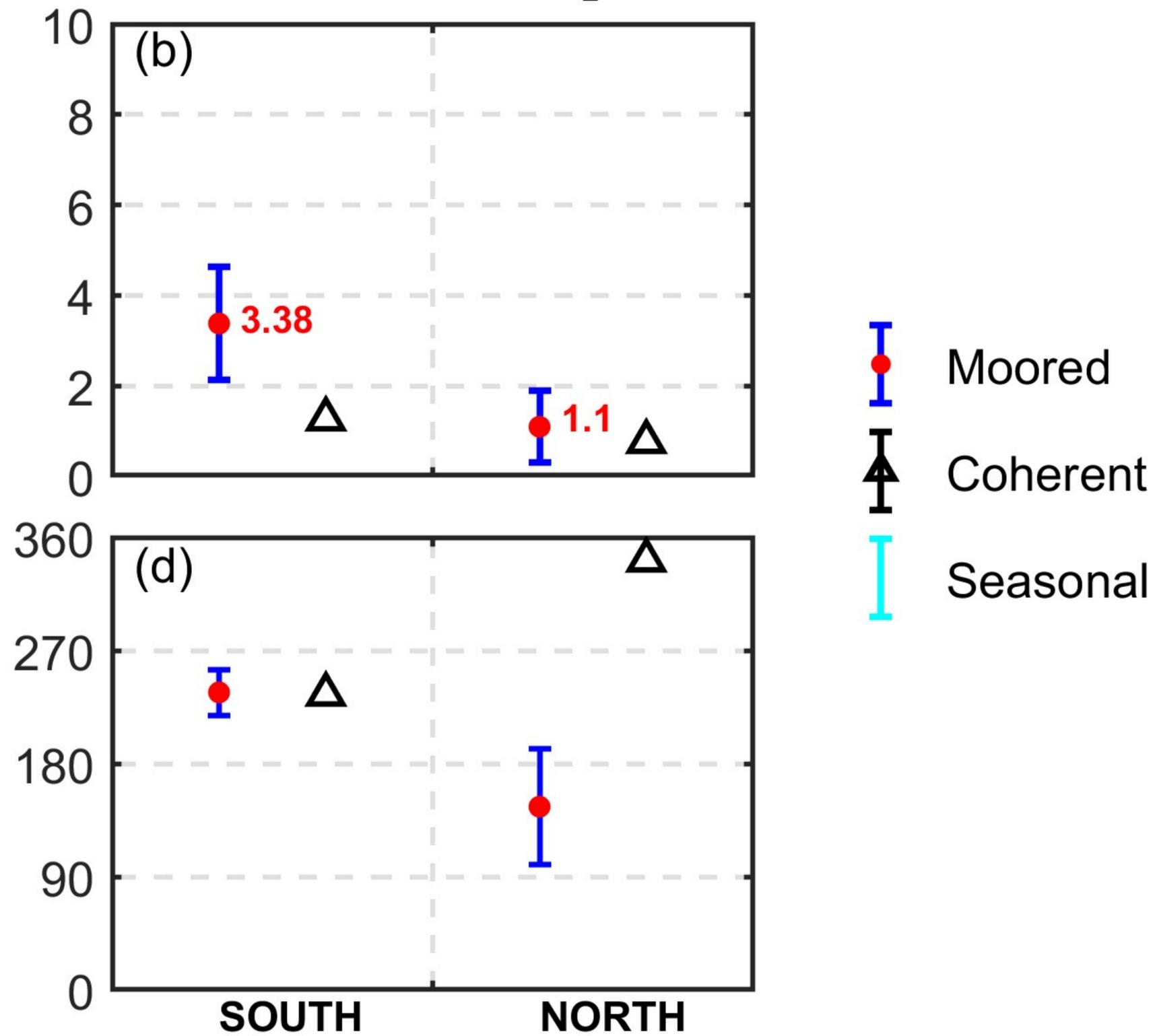
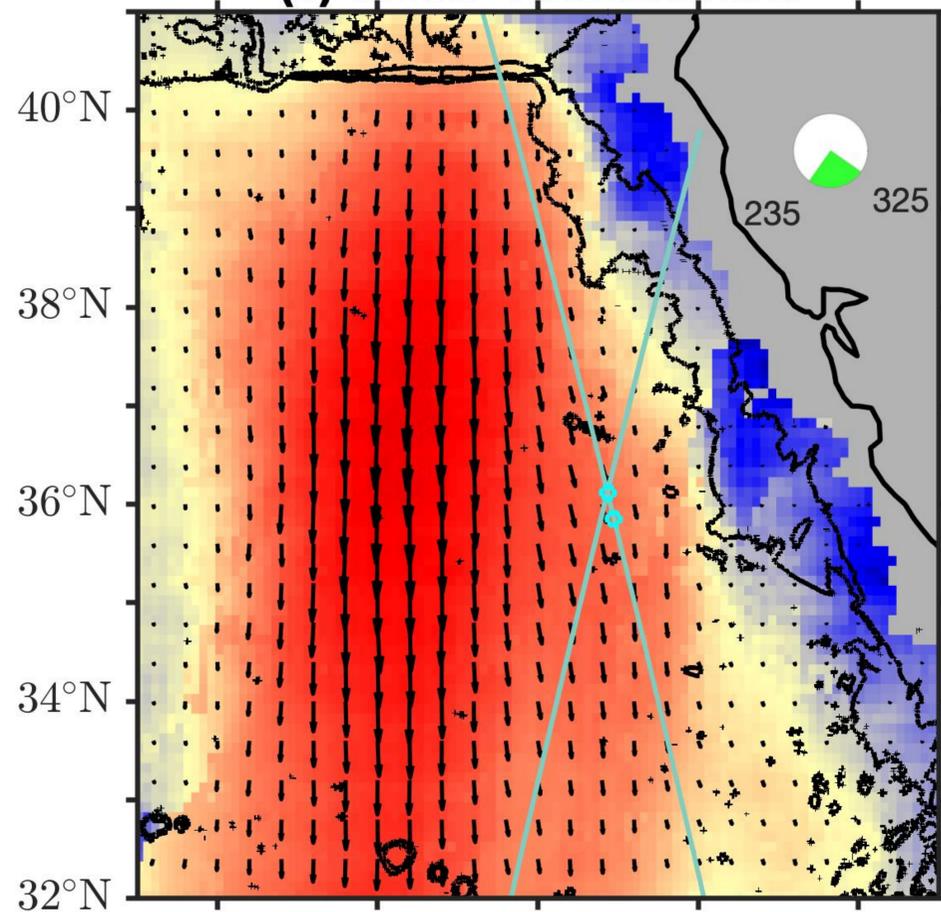
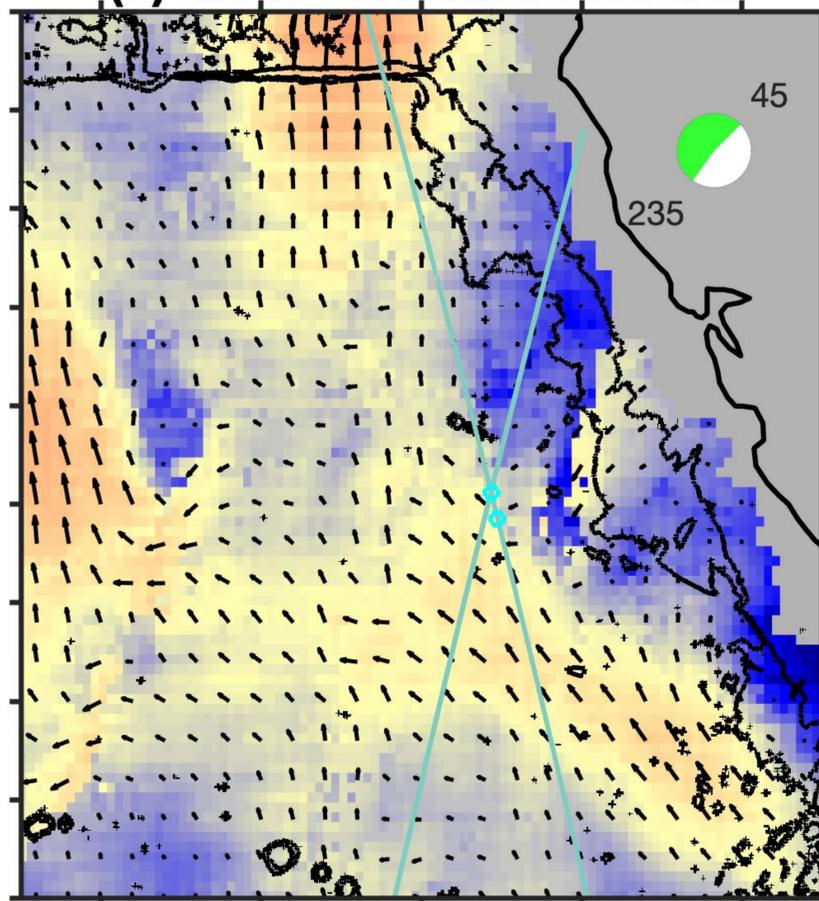
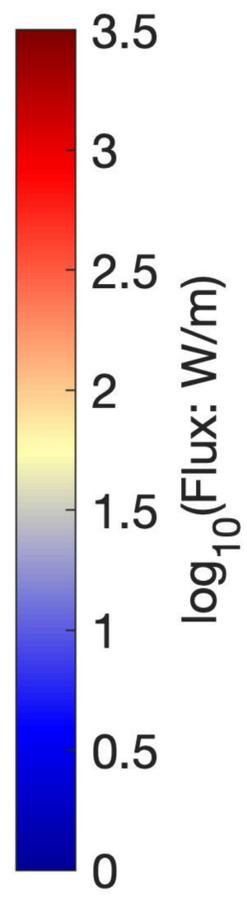
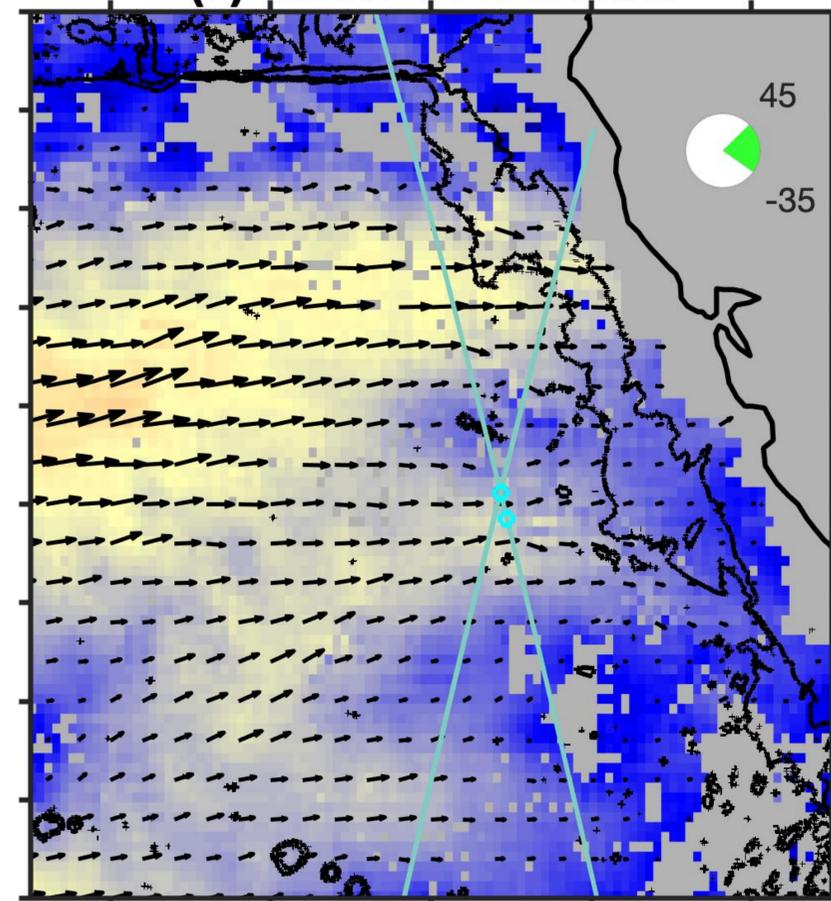
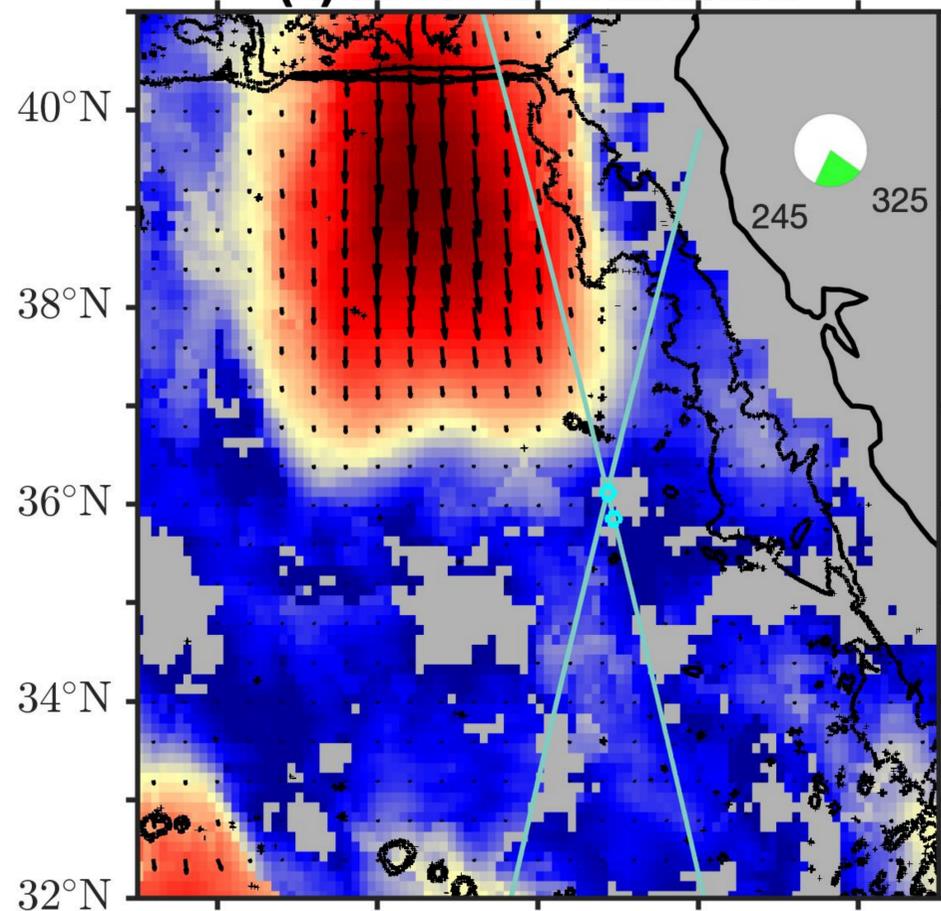
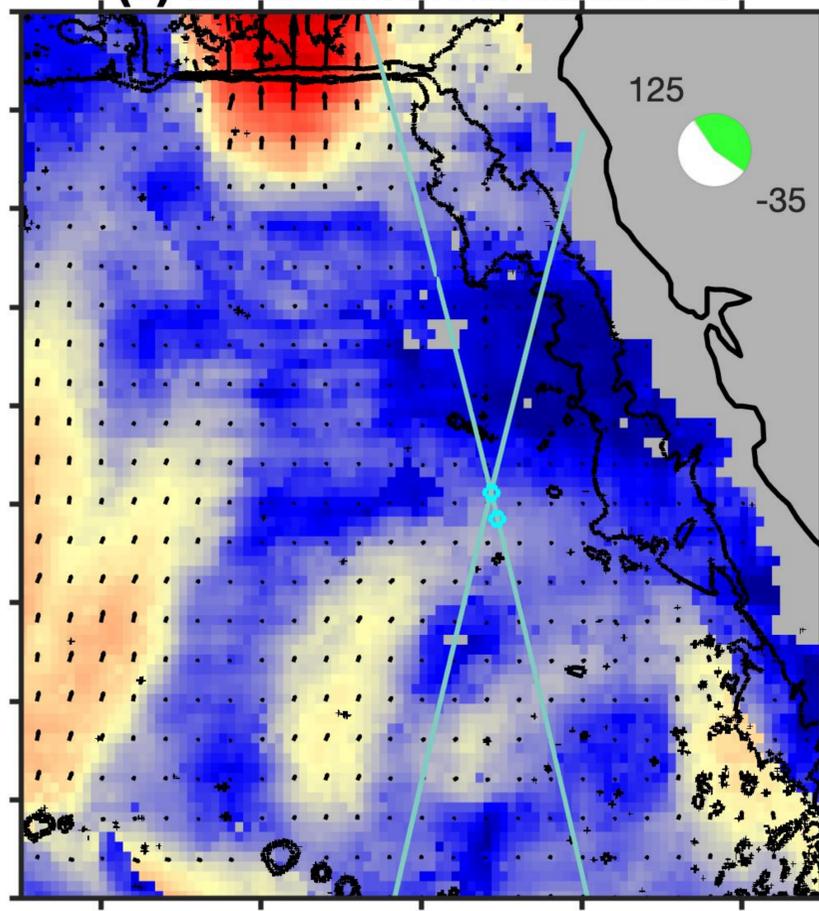
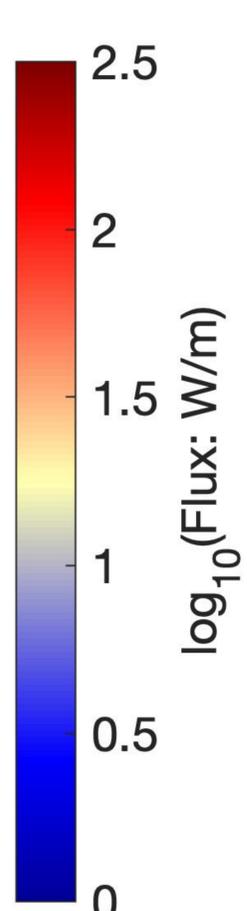
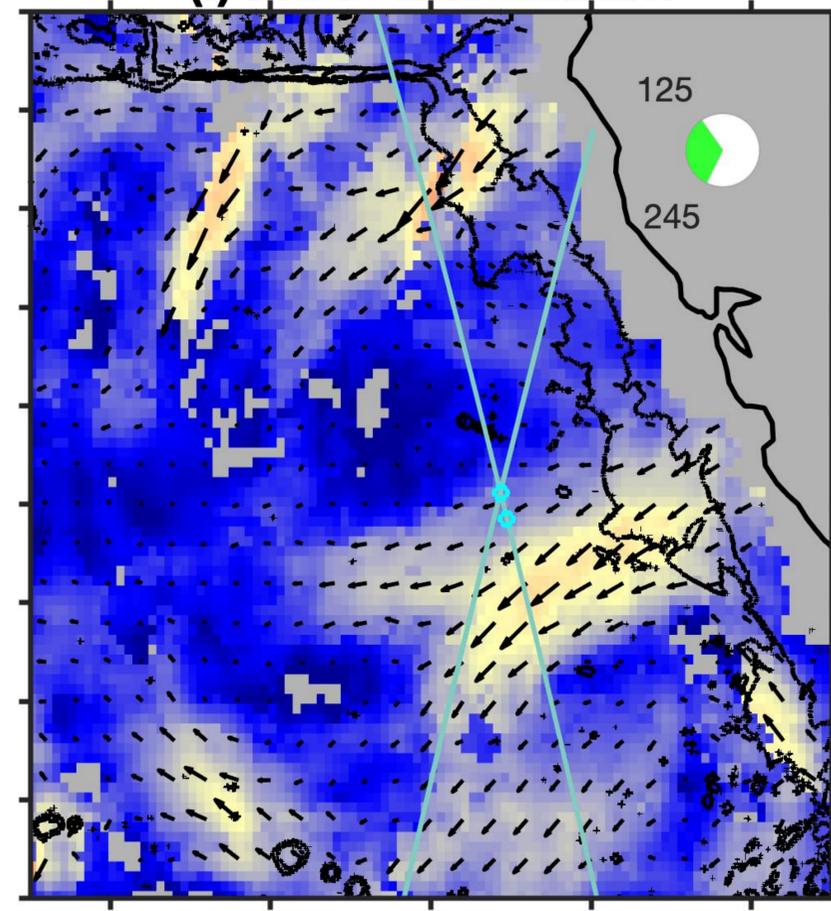
Mode-1 M_2 Mode-2 M_2 

Figure 11.

(a) Mode-1 Southward**(b) Mode-1 Northwestward****(c) Mode-1 Eastward****(d) Mode-2 Southward****(e) Mode-2 Northeastward****(f) Mode-2 Westward**

130°W 128°W 126°W 124°W 122°W

130°W 128°W 126°W 124°W 122°W

130°W 128°W 126°W 124°W 122°W

Figure 12.

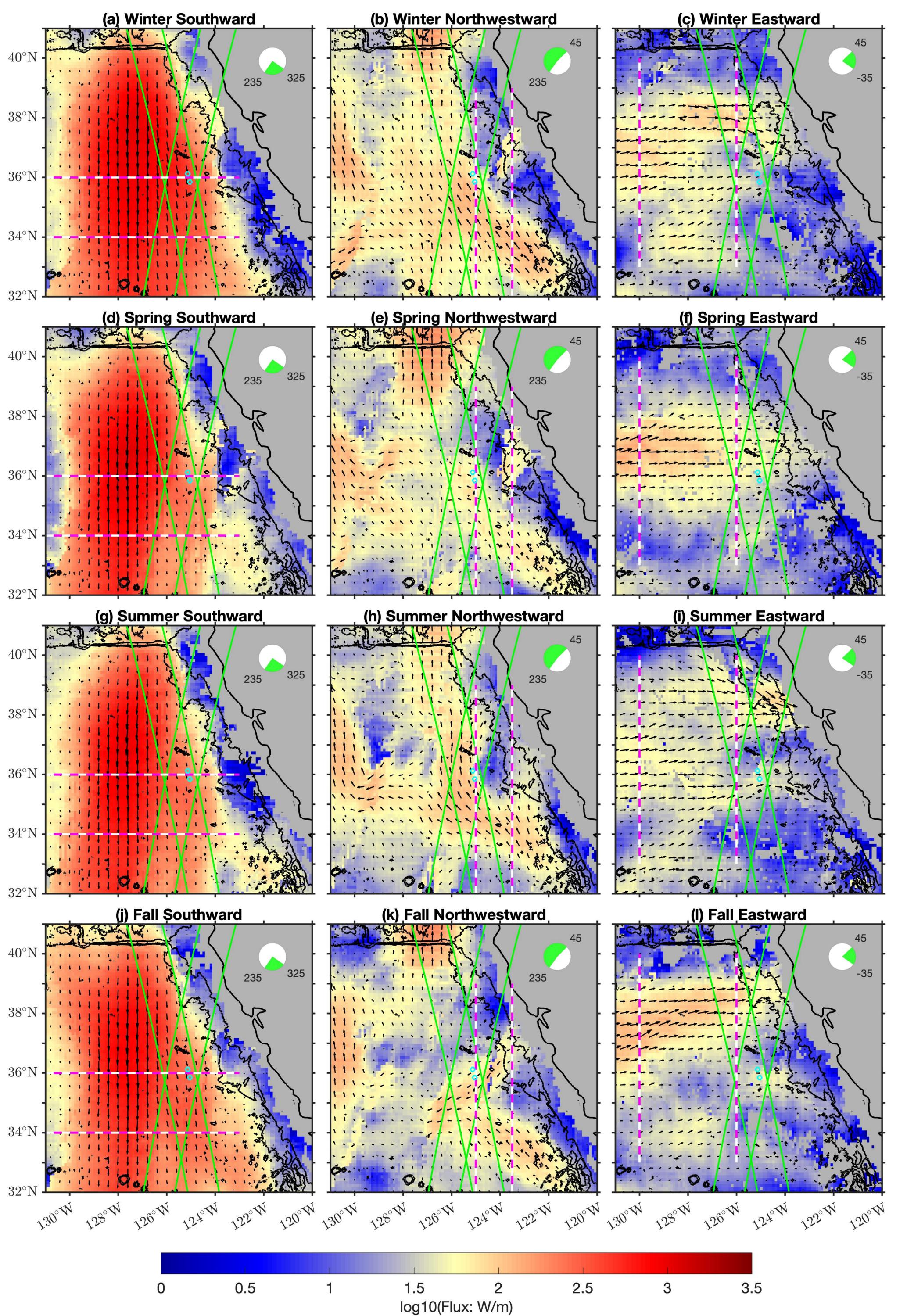
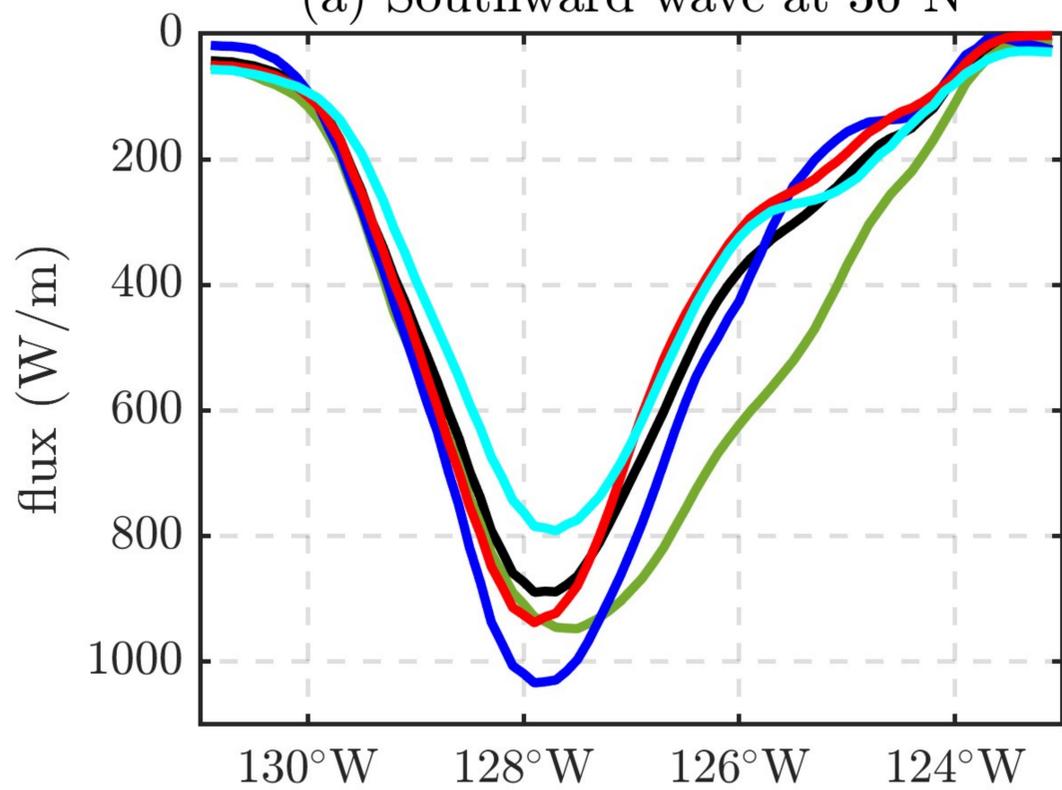
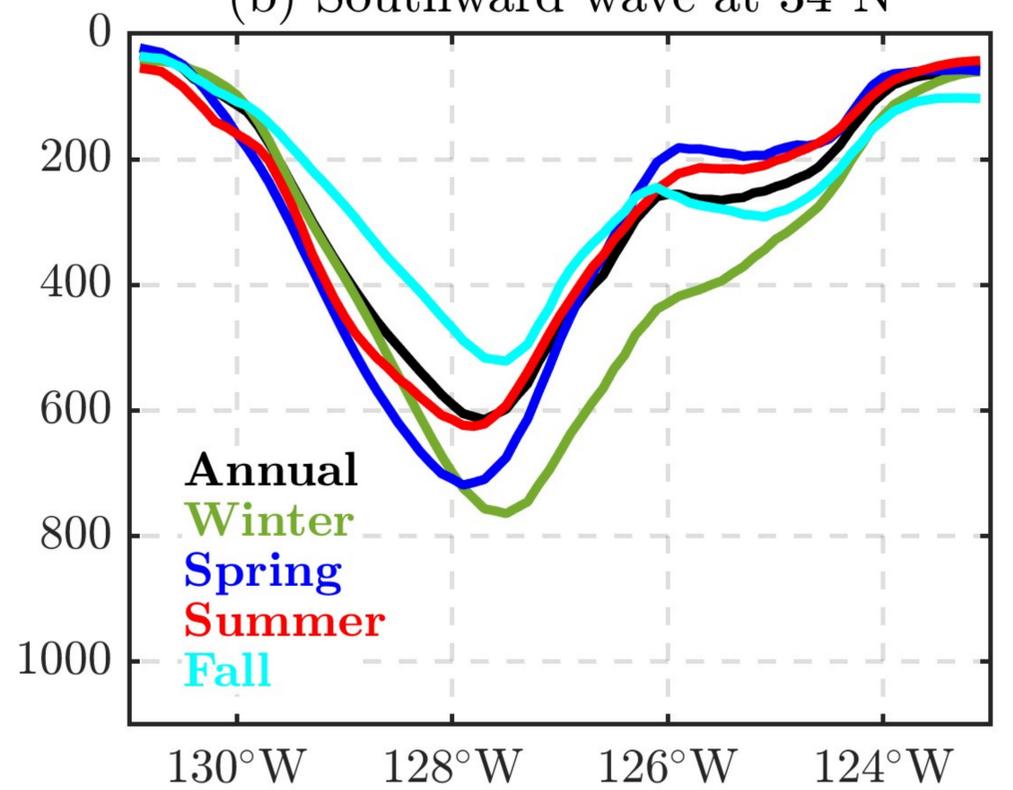


Figure 13.

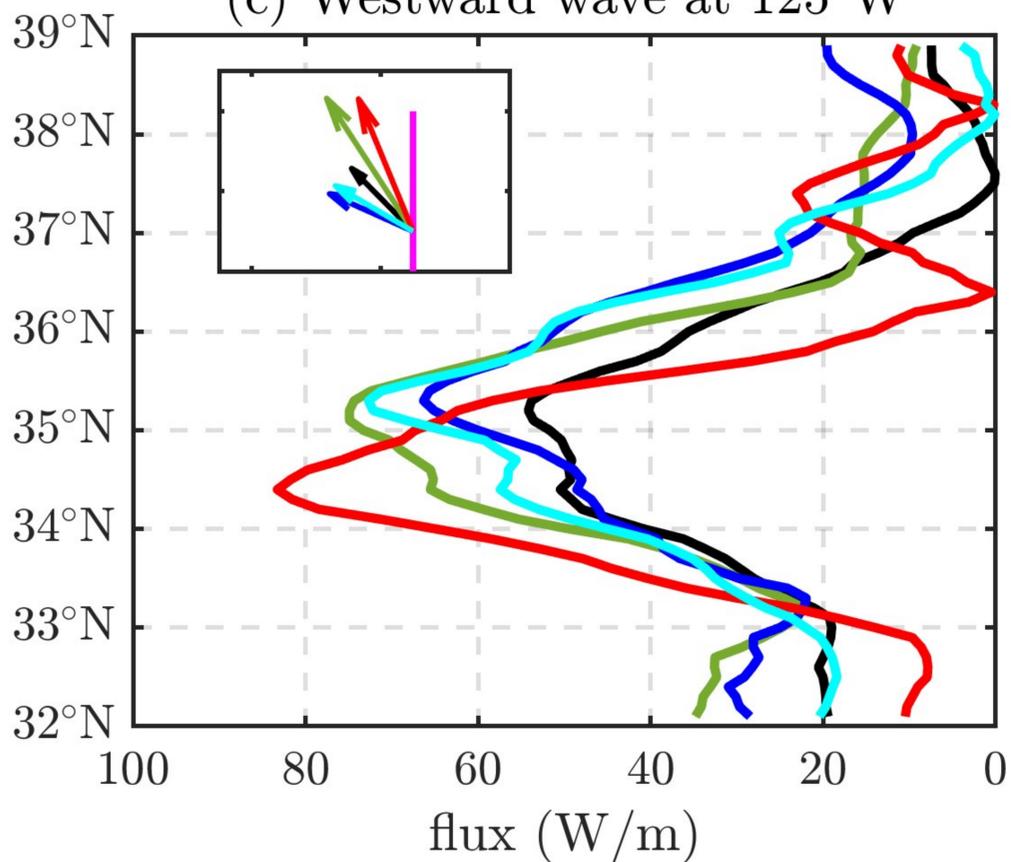
(a) Southward wave at 36°N



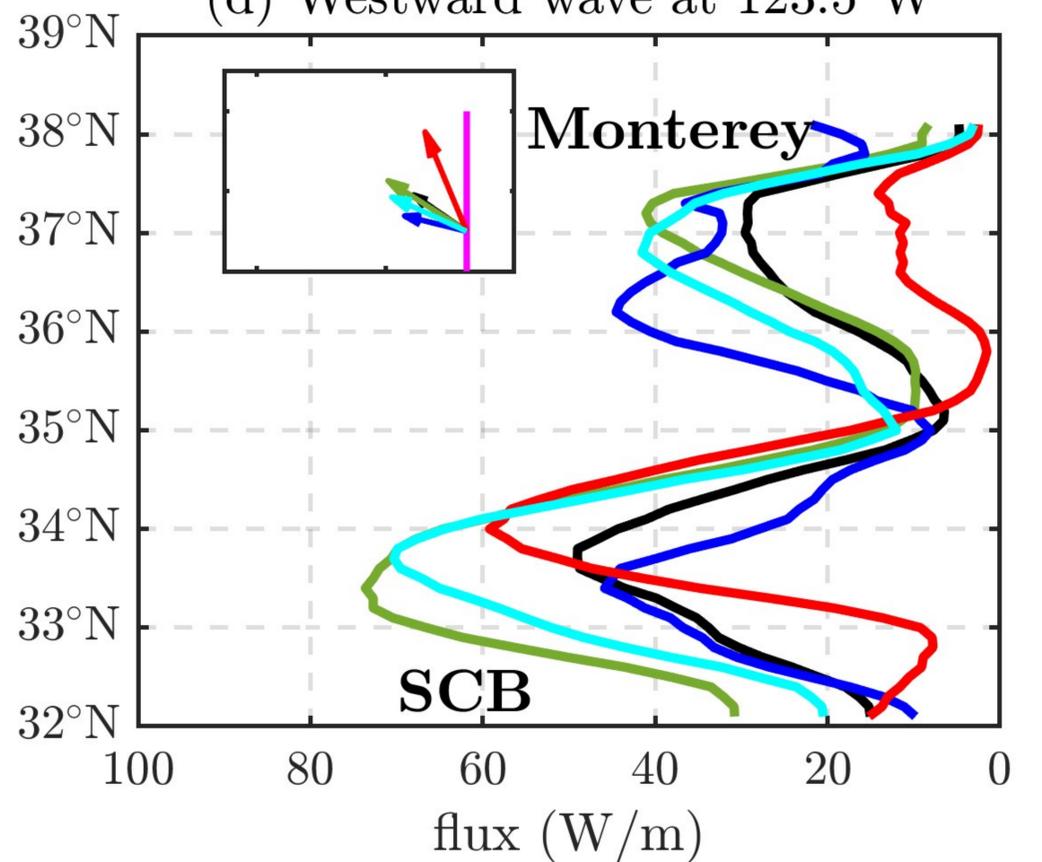
(b) Southward wave at 34°N



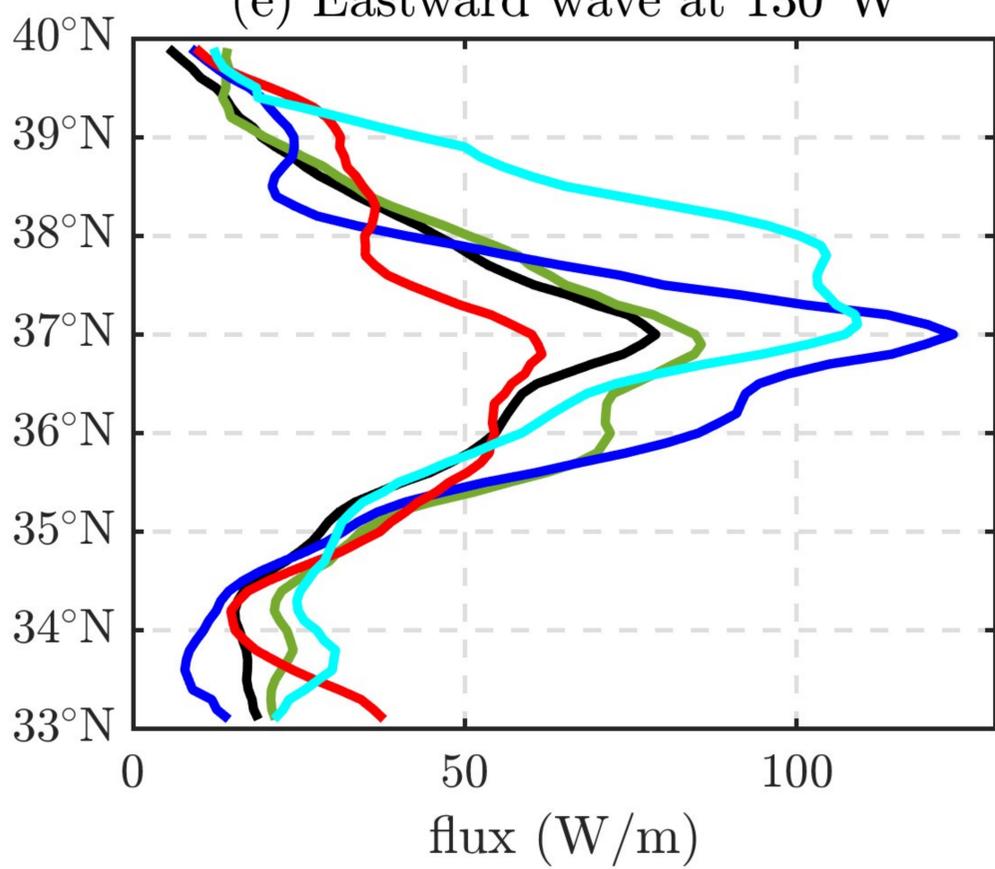
(c) Westward wave at 125°W



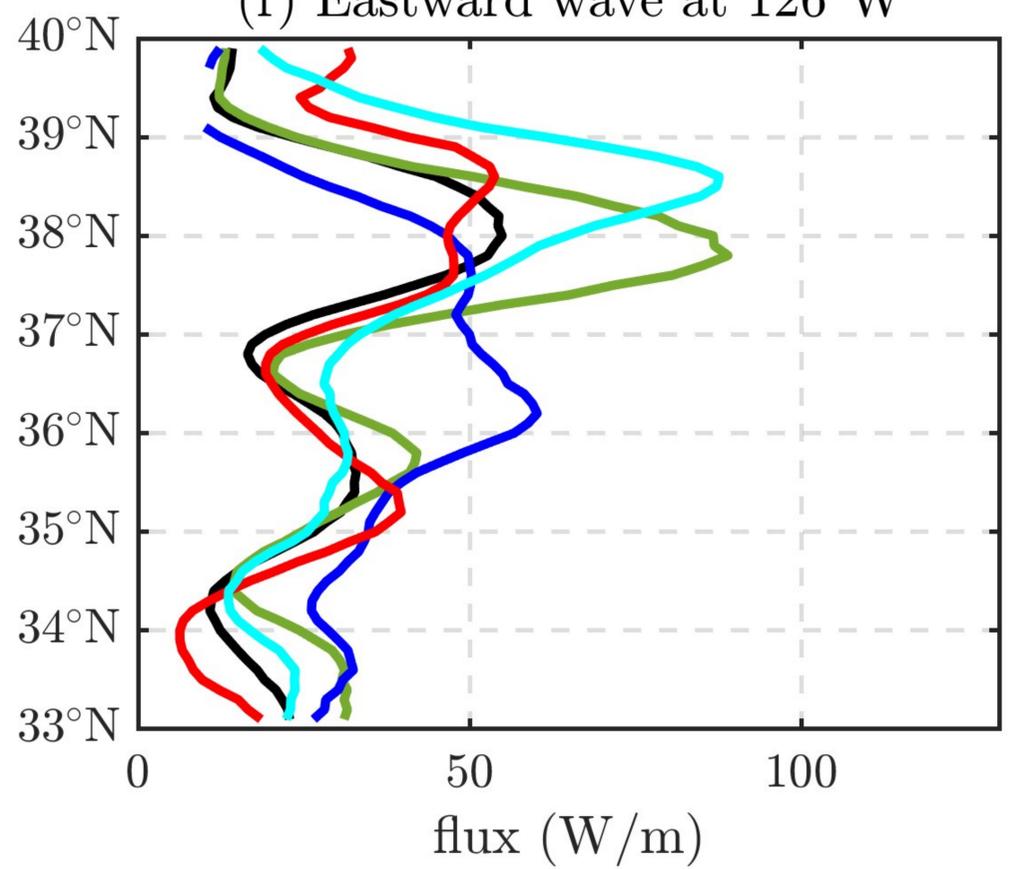
(d) Westward wave at 123.5°W



(e) Eastward wave at 130°W



(f) Eastward wave at 126°W



Supporting Information for “Time-Varying Internal Tides Revealed by Mooring Measurements in SWOT Cal/Val Pre-Launch Field Campaign 2019”

Tongxin Cai^{1,2}, Zhongxiang Zhao^{1,2}, Eric D’Asaro^{1,2}, Jinbo Wang³,

Lee-Lueng Fu³

¹Applied Physics Laboratory, University of Washington, Seattle, Washington, USA.

²School of Oceanography, University of Washington, Seattle, Washington, USA.

³Jet Propulsion Laboratory, California Institute of Technology, Pasadena, California, USA.

Contents of this file

1. Figures S1 to S8

Introduction

1. Data processing and equations

1.1. Quality control

Several steps of quality control are conducted. First, data below the surface mixed layer are selected, with a mixed layer depth of 60 m chosen for both moorings (Figure S1). Additionally, unrealistic extreme values or missing values are identified and removed. Specifically, the bottom 4510-m CTD data are excluded due to data corruption (Wang et

Corresponding author: T. Cai, School of Oceanography, University of Washington, 1503 NE Boat Street, Seattle, WA 98195-7940, USA (joycecai@uw.edu)

al., 2022), and the 261-m CTD data are disregarded due to a high level of noise observed in the frequency spectrum analysis (Figure S4d).

1.2. Buoyancy frequency

The buoyancy frequency is defined as

$$N^2(z) = -\frac{g}{\rho_0} \frac{d\sigma(z)}{dz} \quad (1)$$

1.3. Displacement correction

The pressure measurement taken at each CTD and the configuration of the mooring have revealed that the north mooring experienced a pull-down of approximately 300 m due to its “slack” design. As a result, the CTDs were not precisely fixed at the intended pressure level, resulting in slight vertical movements (Figure S3), especially in deeper waters (beyond 1000 m). Therefore, it is essential to adjust the vertical displacement at each depth by removing the component caused by pressure variations η_P , which we define as follows:

$$\eta_P(z, t) = \overline{P}(z, t) - P(z, t) \quad (2)$$

Here $\overline{P}(z, t)$ is the 10-day moving-average pressure at depth z . An example of η_P is shown in Figure S6b for the sensor at 2750 m from the north mooring. By taking account of the small vertical motion of CTDs, we have the vertical internal tide displacement η_{tide} as

$$\eta_{tide}(z, t) = \eta_\sigma(z, t) + \eta_P(z, t) \quad (3)$$

The data from the south mooring with taut design were less affected, but it is still crucial to apply the correction. The displacement correction at 2750 m for the north mooring is illustrated in Figure S6.

1.4. Phase velocity

In a nonrotating fluid, the eigenvalue velocity c_n is equal to phase velocity and group velocity. If under the influence of Earth's rotation Ω , the phase velocity c_p of each mode can be calculated based on dispersion relation following (Rainville & Pinkel, 2006; Zhao, 2021)

$$c_p^n = \frac{\omega}{\sqrt{\omega^2 - f^2}} c_n \quad (4)$$

where ω is the tidal frequency in this study and f is the inertial frequency. The phase velocity c_p^n of each mode varies with ocean stratification, as it is determined by the eigenvalue velocity c_n , which is a function of the buoyancy frequency $N(z)$ and depth H . The phase velocity at each time is then projected onto each mode by addressing a least squares problem.

2. Spectrum of vertical displacement

Prominent semidiurnal signals are observed across sensors in various depth below mixed layer depth at the north mooring (Figure S4). The significance of these tidal peaks is statistically confirmed within both the 95% (dim gray) and 50% (dark gray) confidence intervals (CI). To compute the spectra, a sine multitaper method was employed, utilizing a degree of freedom (DOF) of 4. Additionally, a geometric smoothing process was applied, spanning 1/250 of the total bandwidth, to enhance spectral coherence. At the south mooring (Figure S5), the measurements obtained from the fixed CTDs below 500 meters also exhibit dominant semidiurnal signals, characterized by notable peaks of the M_2 constituent and their statistical significance. At the sensor positioned at a depth of 4395 meters (Figure S5h), near the bottom (4516 m), the vertical displacement is primar-

ily influenced by turbulence induced by currents and/or waves within the weakly-stratified bottom boundary layer (Garrett, 2003; Wunsch et al., 2004; Kunze, 2017) .

3. Mode fitting number sensitivity analysis

Theoretical considerations of modal decomposition suggest that the number of modes employed for fitting does not significantly affect the obtained results due to the orthogonality of modes. However, practical challenges arise when performing on data sets characterized by vertical spatial gaps (Nash et al., 2005). These challenges are particularly pronounced for higher-mode signals due to their vertical structure and relatively weak magnitude, especially in scenarios where the available upper ocean data is sparse or lacks deep ocean observations (Zhao et al., 2010). Additionally, the computational burden associated with fitting a large number of modes is considerable. Consequently, determining the optimal number of modes for the decomposition process becomes imperative.

To evaluate the influence of incomplete water column coverage in mooring configurations in the campaign, we conducted a sensitivity analysis by varying the number of modes used for mode fitting. Specifically, we examined six distinct scenarios: fitting only mode 1, fitting mode 1-2, fitting mode 1-3, fitting mode 1-5, fitting mode 1-8, and fitting mode 1-10. The energy of the low-mode tide (mode 1-3) was compared across these scenarios, as depicted in Figure S7. Notably, the energy of the low-mode tide in both moorings converged when employing five or more modes for fitting. Considering the computational costs involved, it is evident that fitting the lowest five modes suffices for our analytical purposes, particularly when focusing on mode-1 and mode-2.

References

de Boyer Montégut, C., Madec, G., Fischer, A. S., Lazar, A., & Iudicone, D. (2004).

Mixed layer depth over the global ocean: An examination of profile data and a profile-based climatology. *Journal of Geophysical Research: Oceans*, *109*(C12).

Garrett, C. (2003). Internal tides and ocean mixing. *Science*, *301*(5641), 1858–1859.

Kunze, E. (2017). Internal-wave-driven mixing: Global geography and budgets. *Journal of Physical Oceanography*, *47*(6), 1325–1345.

Li, Q., & Fox-Kemper, B. (2017). Assessing the effects of Langmuir turbulence on the entrainment buoyancy flux in the ocean surface boundary layer. *Journal of Physical Oceanography*, *47*(12), 2863–2886.

Nash, J. D., Alford, M. H., & Kunze, E. (2005). Estimating internal wave energy fluxes in the ocean. *Journal of Atmospheric and Oceanic Technology*, *22*(10), 1551–1570.

Rainville, L., & Pinkel, R. (2006). Propagation of low-mode internal waves through the ocean. *Journal of Physical Oceanography*, *36*(6), 1220–1236.

Wang, J., Fu, L.-L., Haines, B., Lankhorst, M., Lucas, A. J., Farrar, J. T., ... others (2022). On the Development of SWOT In Situ Calibration/Validation for Short-Wavelength Ocean Topography. *Journal of Atmospheric and Oceanic Technology*, *39*(5), 595–617.

Wunsch, C., Ferrari, R., et al. (2004). Vertical mixing, energy, and the general circulation of the oceans. *Annual Review of Fluid Mechanics*, *36*(1), 281–314.

Zhao, Z. (2021). Seasonal mode-1 M_2 internal tides from satellite altimetry. *Journal of Physical Oceanography*, *51*(9), 3015–3035.

Zhao, Z., Alford, M. H., MacKinnon, J. A., & Pinkel, R. (2010). Long-range propagation of the semidiurnal internal tide from the Hawaiian Ridge. *Journal of Physical*

Oceanography, 40(4), 713–736.

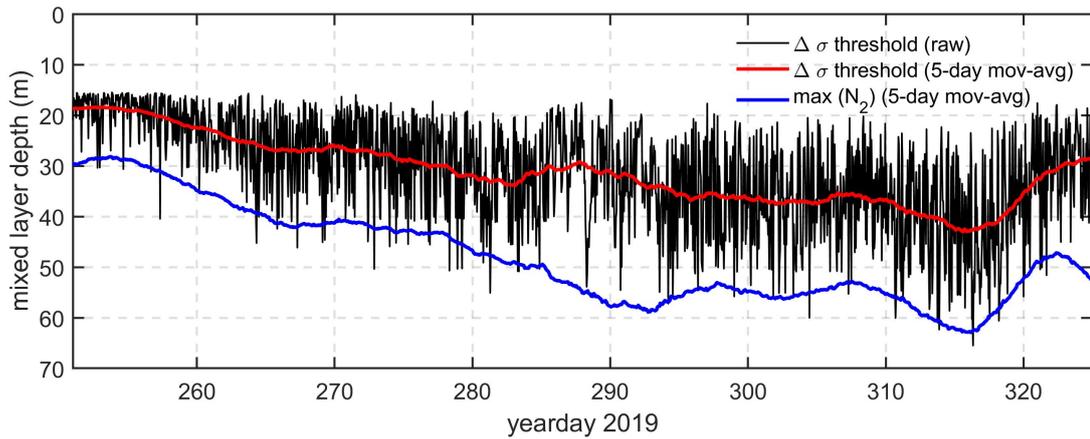


Figure S1. Mixed layer depth (MLD, in unit of m) at the south mooring. Two criteria are used: $\Delta\sigma$ criteria (de Boyer Montégut et al., 2004) and maximum buoyancy frequency (N^2) criteria (Li & Fox-Kemper, 2017). The threshold of $\Delta\sigma = 0.03kg/m^3$ and its temporal variation of MLD is plotted as a black line (raw data) and a red line (after 5-day moving averaging). The blue line represent the MLD using maximum buoyancy frequency criteria, also after 5-day moving average. Consistent deepening of the MLD is observed, starting from 25m and ending with 40m, with the maximum depth reaching 60m.

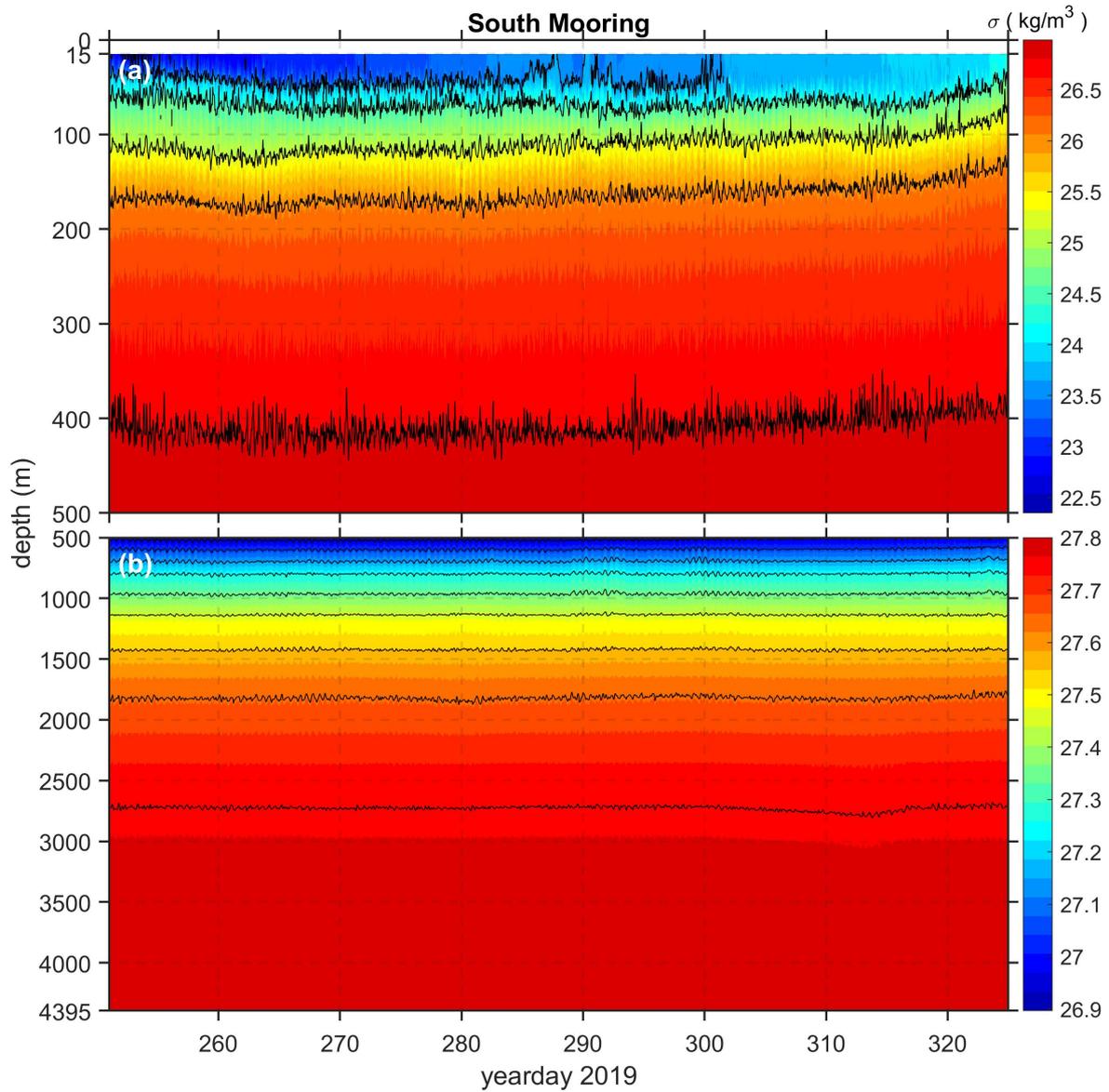


Figure S2. One-hour grided potential density σ (kg/m^3) at the south mooring (a) at upper 500 m from WireWalker Profiler and (b) 500 m - 4390 m from fixed CTDs. Colors indicate potential density σ (kg/m^3) with blue as lighter and red as denser. Black contour lines are isopycnals with constant density value. Note that there are different colorbar limits for (a) and (b).

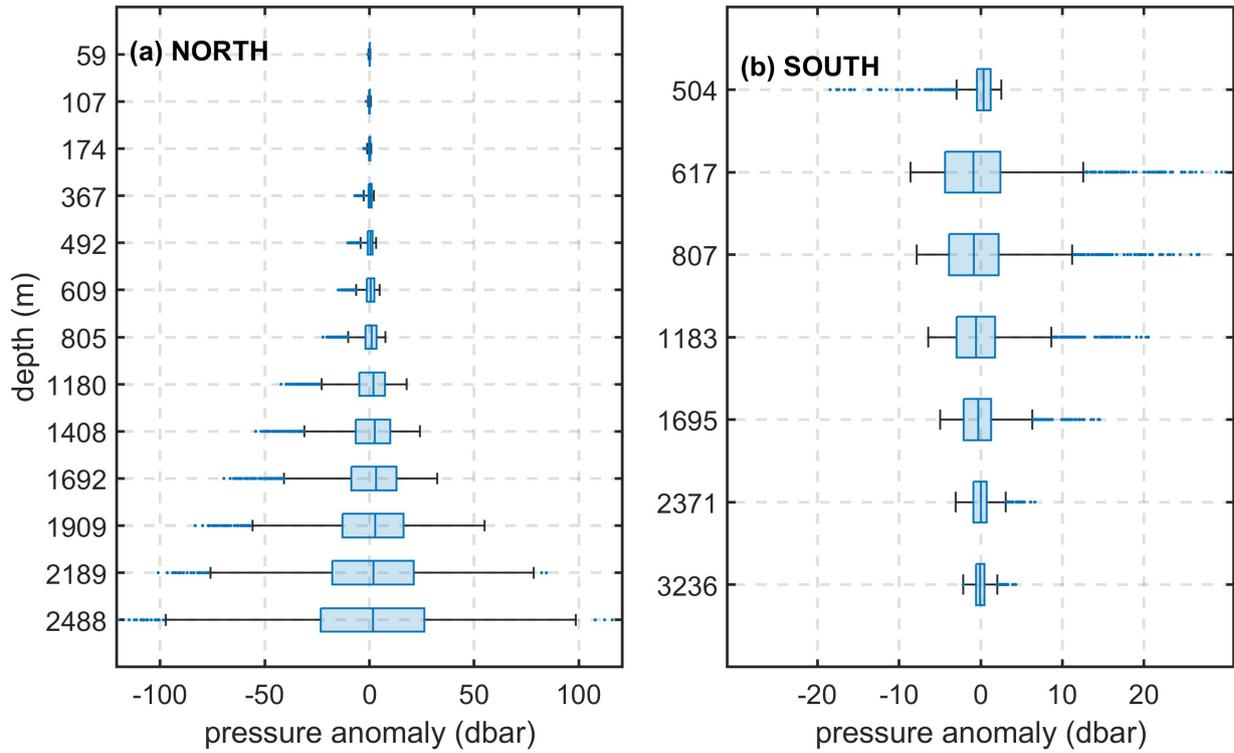


Figure S3. Box plot of pressure anomaly (dBar) of fixed CTDs from (a) the north mooring and (b) the south mooring. Due to the mooring configuration, there is large pressure variation from CTDs at the north mooring, especially in the deeper ocean. There is less effect on the south mooring.

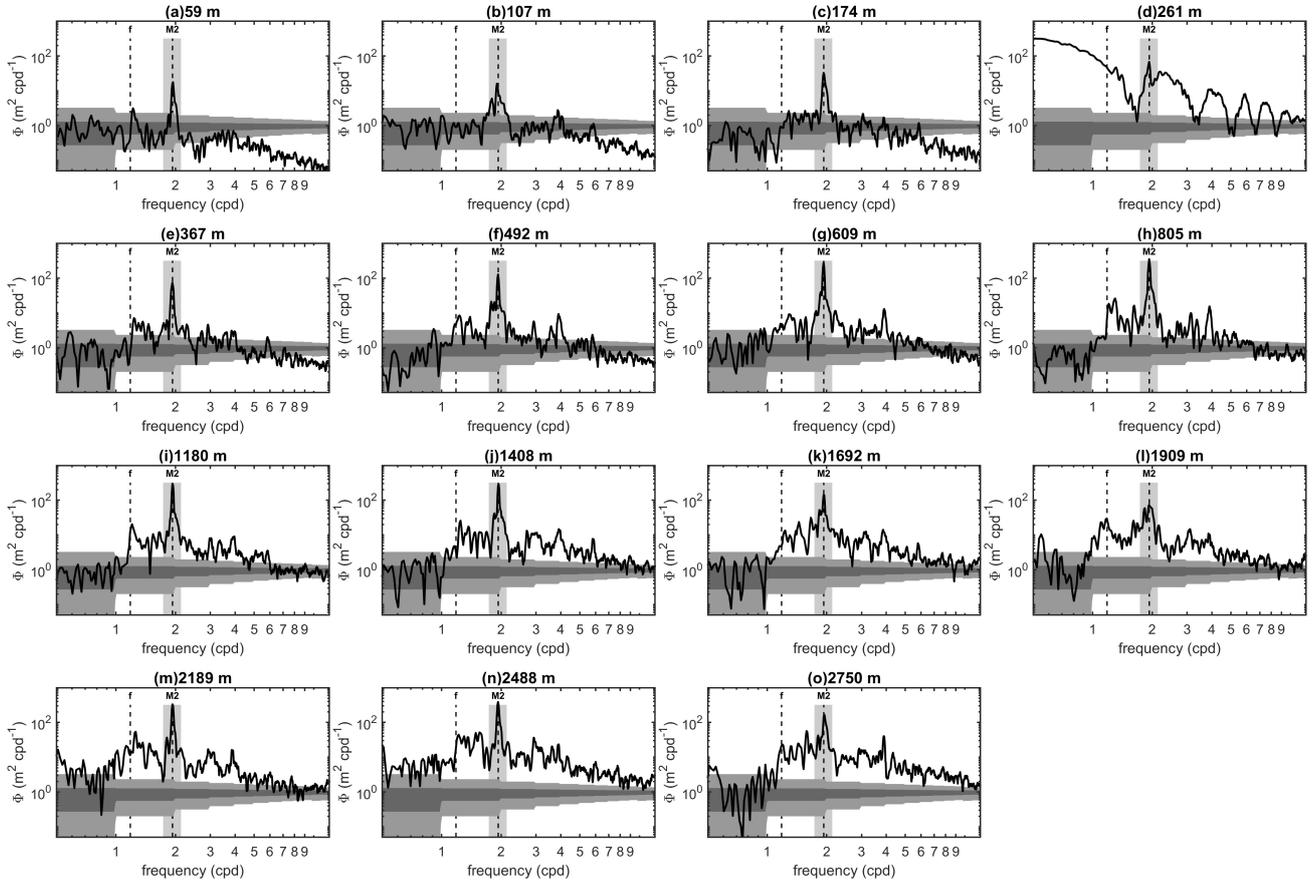


Figure S4. The spectrum of tidal displacement from every sensor at the north mooring. Dim gray are 95% Confident Interval (CI) and dark gray are 50% CI. The semidiurnal band used for filtering are shown in light gray. The two dashed lines indicate the Coriolis f and M_2 frequency. (d) The sensor at 261 m shows high level of noise and uncertainty. Therefore, it is disregard in the tidal analysis.

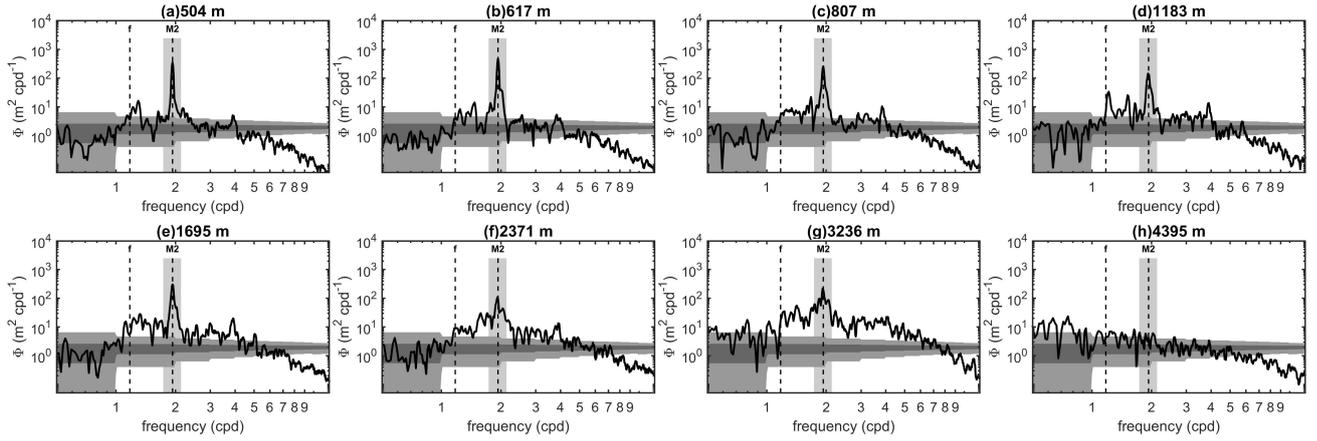


Figure S5. Same as Figure S4 but only fixed sensors at the south mooring below 500 m.

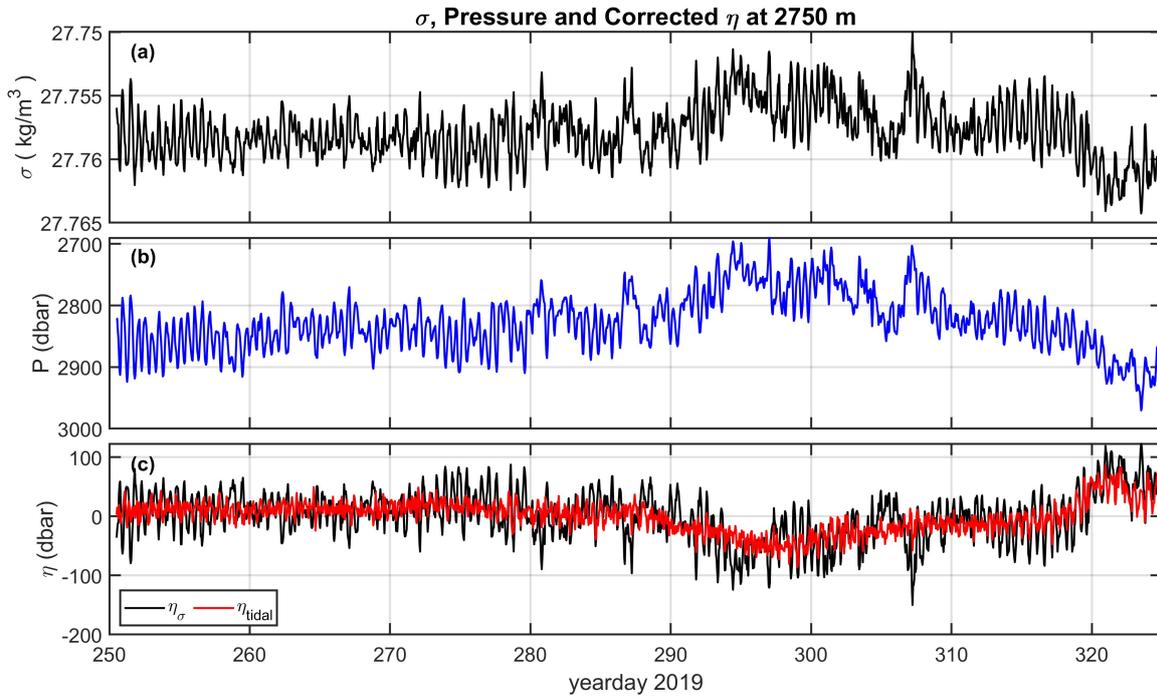


Figure S6. Time series of (a) the potential density anomaly σ , (b) pressure, and (c) the vertical displacement η of the north mooring at 2750 m. The black line is the total displacement η_σ measured, and the red line is the corrected displacement due to internal tide η_{tide} .

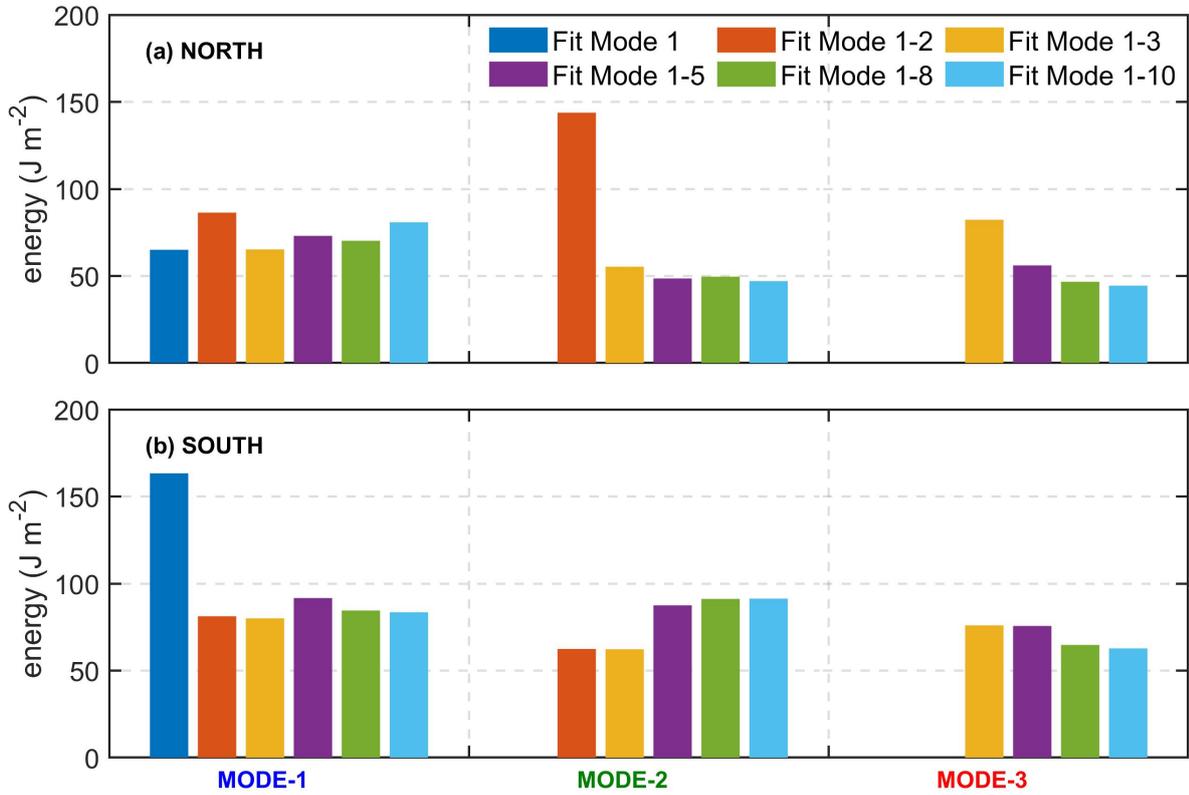


Figure S7. Energy of mode 1-3 (x axis) when mode fitting with different mode number at (a) the north mooring and (b) the south mooring. Six scenarios are examined and shown in different color bars; (blue) fitting only mode 1, (orange) fitting mode 1-2, (yellow) fitting mode 1-3, (purple) fitting mode 1-5, (green) fitting mode 1-8 and (blue) fitting mode 1-10.

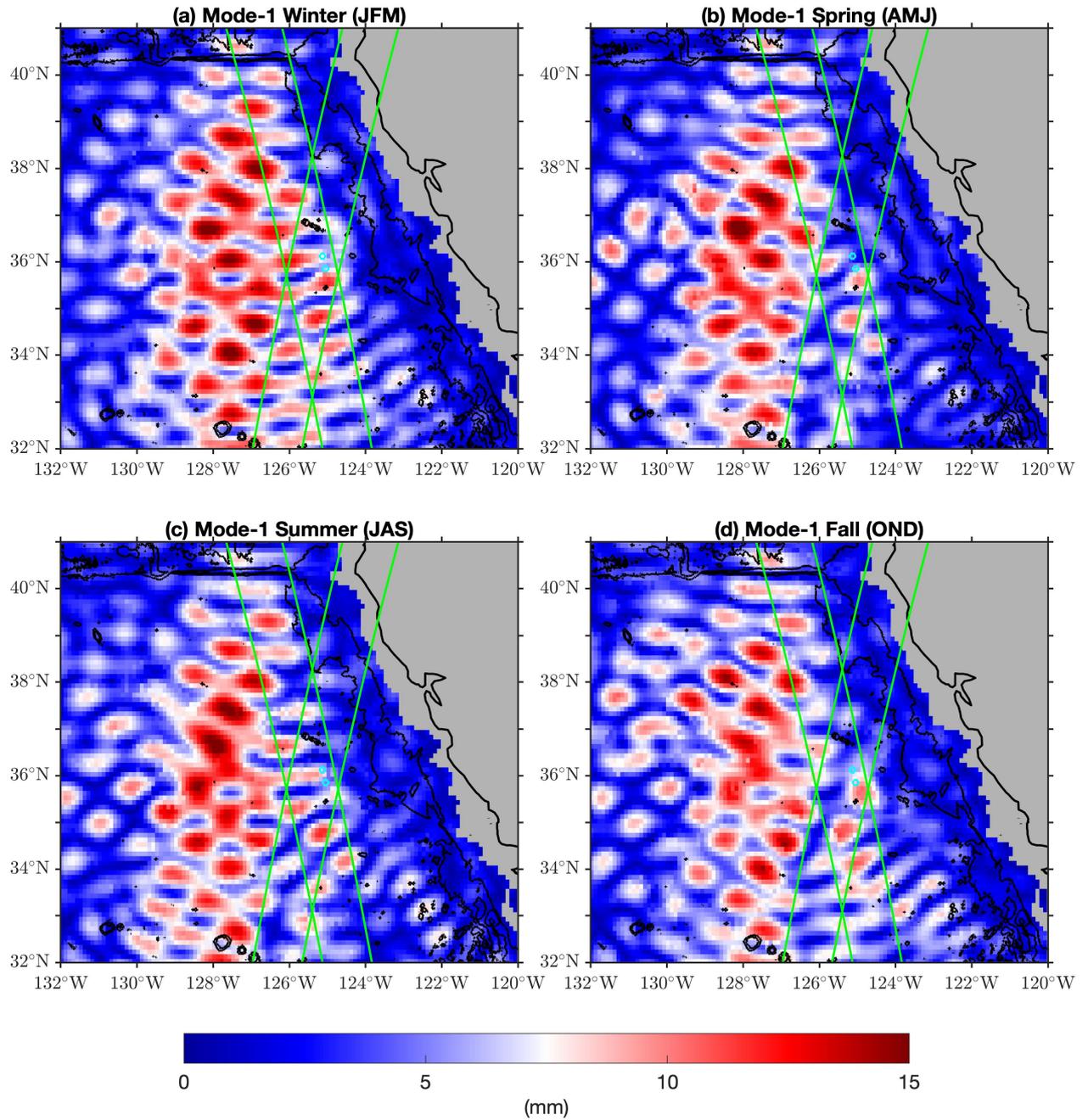


Figure S8. The SSHAs (mm) of mode-1 M_2 internal tides from four climatologically seasonal models. Each seasonal model consists of data from three months: (a) January, February, and March for winter, (b) April, May, and June for spring, (c) July, August, and September for summer, (d) October, November, and December for fall. Green lines are the SWOT Cal/Val swath tracks and cyan circles are the two moorings from the SWOT pre-launch campaign. Contours for the 3000-m and 3800-m isobath are shown.

January 6, 2024, 4:43am