Sensitivity of the Shallow-to-Deep Convective Transition to Moisture and Wind Shear in the Amazon

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Abstract

Deep convection is the primary influence on weather and climate in tropical regions. However, understanding and simulating the shallow-to-deep (STD) convective transition has long been challenging. Here, we conduct high-resolution numerical simulations to assess the environmental controls on the evolution of isolated convection in the Amazon during the wet season. Observations and large-scale forcing derived through the constrained variational analysis approach for the GoAmazon2014/5 experiments are used in the simulations and model validation. The model consistently reproduces the GOAmazon observations for precipitation, moisture, and surface fluxes of radiation, latent and sensible heat. Through sensitivity experiments, we examine the relative importance of moisture and vertical wind shear in controlling the STD convective transition. Reducing the pre-convective humidity within the lower 1.5 km significantly suppresses vertical development and lowers the ice water path. Additionally, the maximum precipitation rate decreases almost quadratically with column water vapor. Conversely, a reduction of column water vapor above 1.5 km by a factor of two or more is necessary to produce a comparable decrease in ice water path or precipitation. Moderate low-level wind shear facilitates the STD transition, leading to an earlier peak of ice water compared to stronger wind shear or its absence. Although upper-level wind shear negatively influences high cloud formation, its role in controlling the STD transition is relatively smaller than that of low-level wind shear. Our results help quantify the role of moisture and wind shear on the STD transition, but also suggest that dynamic factors may exert a more pronounced influence.

Sensitivity of the Shallow-to-Deep Convective 1 Transition to Moisture and Wind Shear in the Amazon 2

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Key Points: 10

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11	•	SAM-LSM consistently reproduces the $GoAmazon2014/5$ observations for precip-
12		itation, moisture, and surface fluxes during the wet season.
13	•	Daytime convection shows a noticeable sensitivity to pre-convective low-level hu-
14		midity and a weaker response to free troposphere humidity.
15	•	Vertical wind shear has a lesser influence than humidity on the shallow-to-deep
16		convective transition.

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17 Abstract

Deep convection is the primary influence on weather and climate in tropical regions. 18 However, understanding and simulating the shallow-to-deep (STD) convective transition 19 has long been challenging. Here, we conduct high-resolution numerical simulations to 20 assess the environmental controls on the evolution of isolated convection in the Amazon 21 during the wet season. Observations and large-scale forcing derived through the constrained 22 variational analysis approach for the GoAmazon2014/5 experiments are used in the sim-23 ulations and model validation. The model consistently reproduces the GOAmazon ob-24 25 servations for precipitation, moisture, and surface fluxes of radiation, latent and sensible heat. Through sensitivity experiments, we examine the relative importance of mois-26 ture and vertical wind shear in controlling the STD convective transition. Reducing the 27 pre-convective humidity within the lower 1.5 km significantly suppresses vertical devel-28 opment and lowers the ice water path. Additionally, the maximum precipitation rate de-29 creases almost quadratically with column water vapor. Conversely, a reduction of col-30 umn water vapor above 1.5 km by a factor of two or more is necessary to produce a com-31 parable decrease in ice water path or precipitation. Moderate low-level wind shear fa-32 cilitates the STD transition, leading to an earlier peak of ice water compared to stronger 33 wind shear or its absence. Although upper-level wind shear negatively influences high 34 cloud formation, its role in controlling the STD transition is relatively smaller than that 35 of low-level wind shear. Our results help quantify the role of moisture and wind shear 36 on the STD transition, but also suggest that dynamic factors may exert a more pronounced 37 influence. 38

³⁹ Plain Language Summary

The Amazon rainforest plays a vital role in the Earth's climate system. However, 40 it is not entirely understood how environmental conditions control the evolution from 41 fair weather conditions to severe thunderstorms in regions of the deep Tropics. We ad-42 dress this problem utilizing numerical simulations that capture the interactions between 43 the forest, atmosphere, and clouds. Atmospheric modeling data developed for the GoA-44 mazon2014/5 experiment are used to initialize our Amazon-based simulations. The model 45 consistently reproduces the Amazon environment throughout the period of our simula-46 tions, which covers December 2014. Additionally, we contrast the model results between 47 the control simulation and experiments in which the moisture or wind is modified to eval-48 uate their relative importance to cloud development and precipitation. Lower tropospheric 49 moisture is critical to cloud growth. The amount of moisture in the air above 1.5 km has 50 a minor contribution to cloud development and precipitation. Low-level wind of mod-51 erate strength facilitates cloud development during the afternoon. The upper-level wind 52 negatively affects the ice formation in high clouds. These results help strengthen our knowl-53 edge of tropical convection, critical for improving numerical model performance. 54

55 1 Introduction

Deep convection dominates the weather and climate in the tropics. Nevertheless, 56 comprehending and simulating the convective processes is a formidable challenge due to 57 the wide range of spatial and temporal scales involved (Mapes et al., 2009; Moncrieff et 58 al., 2012; Zhang et al., 2013). Shallow cumulus convection, a small-scale phenomenon 59 lasting tens of minutes and covering spatial scales of the order of a few kilometers, of-60 tentimes evolves into deep convective clouds covering tens of kilometers within typical 61 time scales of 2 to 4 hours (Wu et al., 2009; Hohenegger & Stevens, 2013; Adams et al., 62 2013; Henkes et al., 2021; Powell, 2022). Moreover, deep convection frequently becomes 63 organized and experiences upscale growth into mesoscale convective systems (MCSs) with 64 lifetimes spanning hours to a day and ranging in horizontal scale from 100 km to 1,000 65 km (Houze Jr, 2004). Likewise, land-atmosphere interactions and complex physical pro-66

cesses ranging from cloud microphysics to the generation of gravity waves are intrinsically tied to deep convection (Silva Dias et al., 2002; Mapes et al., 2006; Mapes & Neale,
2011; Jewtoukoff et al., 2013; Gupta et al., 2023).

General circulation models (GCMs) rely on parameterizations of convective pro-70 cesses and typically struggle to reproduce the shallow-to-deep (STD) convective tran-71 sition over continental regions (Betts, 2002; Betts & Jakob, 2002; Bechtold et al., 2004; 72 Grabowski et al., 2006). Their simulated precipitation peaks much earlier than observed 73 (Lin et al., 2000; Betts, 2002; Collier & Bowman, 2004; Dai & Trenberth, 2004), which 74 75 is an important source of bias and uncertainty in GCMs to this day (Sherwood et al., 2014; Stevens & Bony, 2013; Itterly et al., 2018; Maher et al., 2018; Freitas et al., 2020, 76 2024). To circumvent the inherent challenges posed by convective parameterizations, cloud-77 resolving models (CRMs), which explicitly resolve the up- and downdrafts in clouds, have 78 been used to study convective processes over continental and oceanic regions. For ex-79 ample, M. Khairoutdinov and Randall (2006) conducted the first high-resolution numer-80 ical simulations to investigate the STD transition over the Amazon. Their findings high-81 lighted the importance of cold pools in forcing the development of deep convection, while 82 the impact of vertical wind shear and free tropospheric preconditioning were relatively 83 minor. As part of the EUROCS (EUROpean Cloud Systems study), Derbyshire et al. 84 (2004) evaluated the sensitivity of cumulus convection to free tropospheric humidity. Un-85 like M. Khairoutdinov and Randall (2006), they observed intense deep precipitating con-86 vection in moister scenarios, whereas only shallow convection was evident in the driest 87 scenario. Waite and Khouider (2010) conducted idealized numerical simulations over the 88 tropical Atlantic Ocean. Their study emphasized the importance of congestus precon-89 ditioning, which reduces the impact of entrainment on cloud buoyancy, ultimately lead-90 ing to the STD transition. In contrast, Hohenegger and Stevens (2013) showed that the 91 transition from congestus to deep convective clouds occurs on shorter time scales than 92 required for congestus clouds to moisten the atmosphere sufficiently. This implies that 93 dynamic factors play a more substantial role in driving convection. While CRM stud-94 ies offer valuable insights into physical convective processes, they still require validation 95 through high-resolution observations, which have typically been lacking in tropical rain-96 forests. 97

In the Amazon, important, but often limited field campaigns, have explored dif-98 ferent aspects of tropical convection. Adams et al. (2015) established the Amazon Dense 99 GNSS Meteorological Network, a one-year campaign to observe the interaction between 100 water vapor fields and deep convection. Adams et al. (2013) also utilized GNSS/GPS 101 data from a long-term single site (July 2008 to December 2011) in Manaus, Brazil to eval-102 uate the water vapor convergence associated with the STD transition and found a weak 103 and quasi-linear convergence timescale of approximately 8 hours, followed by a robust 104 and non-linear convergence timescale of approximately 4 hours during the STD transi-105 tion. Later, Adams et al. (2017) employed this dense network data to investigate how 106 vapor fields evolve spatially during during the STD transition. Their results were con-107 sistent with the 4-hour STD timescale and the spatial evolution was reflective of the wa-108 ter vapor convergence posited in the single site study. More recently, the Green Ocean 109 Amazon (GOAmazon) 2014/5 Experiment (Martin et al., 2016, 2017) was carried out 110 from 2014 to 2015 in the central Amazon, providing the most comprehensive set of ob-111 servations of clouds and aerosols in the Amazon to date. Analyzing this dataset during 112 the dry season (June-September), Ghate and Kollias (2016) noted an excess of water va-113 por above 2 km during the early morning when contrasting locally-driven precipitating 114 days and nonprecipitating days. Conversely, Zhuang et al. (2017) and Tian et al. (2021) 115 observed that deep convective days exhibit relatively higher moisture extending from the 116 surface to mid-levels in all seasons. Schiro et al. (2016) showed a robust correlation be-117 tween total column water vapor and precipitation in both the central Amazon and the 118 tropical western Pacific. Furthermore, Schiro and Neelin (2019) demonstrated a strong 119 connection between the initiation and likelihood of daytime precipitation and the bound-120

ary layer and lower free troposphere moisture content. Previous studies do not completely
agree on the relative importance of vertical wind shear. For example, while Zhuang et
al. (2017) indicated that more intense low-level and deep-layer bulk wind shears facilitate the STD transition during the dry season (June-September), Chakraborty et al. (2018)
suggested that a more intense low-level shear could inhibit deep convection during the
transition season (August-November), especially if it increases the entrainment of dry
air.

In this paper, we conduct high-resolution model simulations to assess the role of 128 moisture and vertical wind shear in controlling the STD convective transition in the Ama-129 zon. First, we focus on model validation employing GoAmazon data for the period of 130 December 2014. Then, we conduct idealized sensitivity experiments in which either mois-131 ture or large-scale wind are modified at different atmospheric levels to assess their rel-132 ative importance in the development of deep convection. The paper is structured as fol-133 lows: Section 2 shows the study area. Section 3 describes the material and methods. Sec-134 tion 4 covers the model validation. Sensitivity experiments for moisture and wind shear 135 are conducted in section 5. A discussion of the results is given in section 6. Section 7 con-136 tains the conclusions. 137

138 2 Study Region

The Amazon Basin is bordered by significant altitudes (Figure 1a), primarily in 139 the western region, where some peaks in the Andes Mountains rise well over 6,000 me-140 ters in elevation above sea level. However, the simulations are conducted over the GoA-141 mazon2014/5 campaign region in the central Amazon, where the topography can be ad-142 equately considered as an extensive plain with minimal variations (< 130 m in our do-143 main of interest, section 3.2). During the experiment, most of the observations were taken 144 from the T3 site, located 70 km downwind of Manaus, in Manacapuru (3.21°S, 60.60°W), 145 a site characterized by a pasture surrounded by forest and close to the intersection of 146 the Solimões River and Negro River (Figure 1b). 147



Figure 1. (a) Land topography and ocean depth (NOAA National Centers for Environmental Information, 2022) around the Amazon. (b) Land cover (Friedl et al., 2010) around the GoAmazon2014/5 sites. The dashed circle with a radius of 202 km centered at the T1 site (in Manaus) corresponds to the S-band radar domain. The dotted circle with a radius of 110 km shows the domain of the large-scale forcing developed for the GoAmazon2014/5 Experiment (Tang et al., 2016). We also indicate the Amazon, Solimões, and Negro Rivers on the map. Land cover is from 2014, based on the Moderate Resolution Imaging Spectroradiometer (MODIS) - International Geosphere-Biosphere Programme (IGBP) land cover classification system. The Amazon Basin contour is provided by Mayorga et al. (2012).

¹⁴⁸ **3** Material and Methods

¹⁴⁹ **3.1 Data**

For model validation, we use campaign observations of precipitation, moisture, ra-150 diation, and surface latent and sensible heat fluxes. Precipitation is based on the SIPAM 151 S-band radar measurements (Schumacher & Funk, 2018), which we average over the do-152 main of the control runs. Sensible heat flux (H) and latent heat flux (LE) are from the 153 Quality Controlled Eddy Correlation (QCECOR) Flux Measurement (ARM, 2014b). Sur-154 face radiation fluxes are from the Sky Radiation Radiometers (SKYRAD) and Ground 155 Radiation Radiometers (GNDRAD) (ARM, 2013). Column water vapor (CWV) is cal-156 culated from the balloon-borne sounding system (SONDE), which provides the vertical 157 profiles of thermodynamic conditions 4 times per day during the period of this study, 158 at 02, 08, 14, and 20 LST (ARM, 2014a). 159

Large-scale atmospheric fields of water vapor mixing ratio, temperature, wind, and 160 moisture and temperature tendencies are based on 3-hour Constrained Variational Anal-161 ysis Data (VARANAL). This assimilation product was developed using atmospheric fields 162 from the European Centre for Medium-Range Weather Forecasts (ECMWF) ERA-Interim 163 reanalysis (Dee et al., 2011), which were mainly constrained by the SIPAM S-band radar 164 precipitation rate and ARM surface fluxes through the column heat and moisture bud-165 get analysis (Tang et al., 2016). The VARANAL data represent an average over the anal-166 ysis domain centered at T1 site, covering a radius of 110 km (Figure 1b). 167

The Moderate Resolution Imaging Spectroradiometer (MODIS) provides data prod-168 ucts of land cover type (MCD12Q1 Version 6) and leaf area index (MCD15A2H Version 169 6.1) (Friedl et al., 2010). Specifically, we use the land cover product based on the Inter-170 national Geosphere-Biosphere Programme (IGBP) land cover classification system. Silt, 171 clay, and sand content in the soil are based on in-situ measurements of the soil type "Terra 172 Firme" (Terra-firma) described in Table 1 on Schaefer et al. (2017). Soil temperature 173 and wetness are based on the NASA Global Land Data Assimilation System (GLDAS) 174 Noah Land Surface Model L4 3 hourly 0.25 x 0.25 degree V2.1 (Rodell et al., 2004). 175

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3.2 Model Configuration

We employ the numerical model System for Atmospheric Modeling (SAM), ver-177 sion 6.11.8, which solves the anelastic equations of motion and uses liquid water static 178 energy, total nonprecipitating, and precipitating water as thermodynamic prognostic vari-179 ables (M. F. Khairoutdinov & Randall, 2003). The equations are solved using lateral pe-180 riodic boundary conditions. A prognostic turbulent kinetic energy 1.5-order closure scheme 181 is used to parameterize subgrid-scale effects. Different microphysics parameterizations 182 are available, including the single-moment (Morrison, 2003), double-moment (Morrison 183 et al., 2005), and Predicted Particle Properties (P3) (Morrison & Milbrandt, 2015) schemes. 184 The radiative heating can be prescribed or calculated by choosing a radiation scheme, 185 either the Community Atmosphere Model (CAM3) (Collins et al., 2006) or the Rapid 186 Radiative Transfer Model (RRTM) (Mlawer et al., 1997) schemes. Surface fluxes can be 187 prescribed or calculated using Monin-Obukhov similarity theory or a simplified Land Sur-188 face Model (LSM) (Lee & Khairoutdinov, 2015), which is only compatible with the CAM3 189 radiation scheme for the current SAM-LSM version. 190

The baseline configuration for the simulations analyzed in this paper considers a domain of $200 \times 200 \times 27$ km³. This choice was made primarily to reasonably accommodate MCSs, given that they typically span about 100 km (Houze Jr, 2004). The horizontal resolution is 500 m, and the vertical resolution varies: it starts at a minimum of 50 m below 1.5 km and increases to 300 m in the upper troposphere. From there, it gradually stretches up to 500 m at the model's upper boundary, which reaches 27 km, resulting in 128 vertical levels. The temporal resolution is 5 seconds, and instantaneous model fields and statistics are output every 30 minutes. The control simulation uses the P3 microphysics scheme. The CAM3 radiation scheme is called every 150 seconds. Surface fluxes are calculated through the LSM (see Section 3.3).

The large-scale forcing is based on the VARANAL dataset for the period of December 2014 in the central Amazon. Winds were nudged with a 2-hour timescale throughout the simulation. The water vapor mixing ratio was nudged only during the spin-up, considered as the period from 1-5 December 2014, with a timescale of 6 hours.

For the purpose of model validation, we conducted additional simulations where the only modificatio was the choice of the microphysics scheme: single-moment, doublemoment, or P3 schemes. We also assessed model sensitivity to resolution and domain size by performing additional simulations at 250 m resolution or with a 400×400×27 km³ domain.

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3.3 Land Surface Model Configuration

The simplified Land Surface Model uses a minimalist set of parameters to solve the 211 transport of heat, moisture, and radiation in the soil and vegetation and calculate the 212 transfer of momentum between the surface and the atmosphere (Lee & Khairoutdinov, 213 2015). It adequately idealizes the land-atmosphere interactions, which fundamentally in-214 fluence convection over the Amazon forest (Silva Dias et al., 2002; Betts, 2002). To de-215 fine the characteristics of the surface and the vegetation for the LSM, we use the MODIS 216 data of land cover type during 2014 and mean leaf area index (LAI) for December 2014 217 over a domain of 200×200 km² centered at T3 site (see Figure S1). These are associated 218 with the period and area of our simulations. The surface LAI was set to $4.4 \text{ m}^2 \text{ m}^{-2}$, 219 which corresponds to the domain average in satellite observations. In addition, we choose 220 to cover the surface uniformly with evergreen broadleaf forests, which cover 83% of the 221 area in observations. However, based on several tests conducted to optimize the LSM 222 parameters, we modify the default near-infrared visible albedo for vegetation from 0.20 223 to 0.30, the root length from 150 cm to 200 cm, and the displacement height factor from 224 0.68 to 0.65. The corresponding displacement height is $0.65 \times 20 \text{ m} = 13 \text{ m}$, where 20 225 m is the default value of the height of the canopy. These modifications improve the agree-226 ment between the observations and the simulated surface radiation, latent and sensible 227 heat fluxes. 228

The soil is simulated using 11 layers from the surface down to a depth of 400 cm. 229 Clay and sand contents for each layer are based on in-situ measurements on Terra-firma 230 forests (section 3.1). The initial conditions for soil temperature and soil wetness are based 231 on GLDAS Noah, which provides information on 4 layers: 0-10, 10-40, 40-100, and 100-232 200 cm. The LSM soil layers close to the surface, which experience greater diurnal cy-233 cle variation, are interpolated using the nearest neighbor method. The deeper soil lay-234 ers are interpolated (and extrapolated) linearly. The initial profile of soil temperature 235 and wetness is shown in Figure S2. 236

3.4 Cloud Regime Days

For the sensitivity experiments, we select a set of deep convective (Deep) days from 238 the control simulation (P3 scheme, with a horizontal resolution of 500 m and a domain 239 size of $200 \times 200 \times 27 \text{ km}^3$) and perturb the sounding or the large-scale forcing imposed. 240 For the Deep selection, we require that the domain average of total ice presents a dis-241 tinct deepening during the afternoon, characteristic of the STD convective transition. 242 The chosen Deep days are December 17th, 21st, 23rd, and 26th. For comparison, we also 243 select a set of shallow cumulus (ShCu) days from the control run. We identify four days 244 with negligible ice content and minimum surface precipitation: December 9th, 13th, 27th, 245

and 28th. Figure S3 shows the profile of cloud liquid and total ice for our selection of
 cloud regime days.

248 4 Model Validation

A comparison between simulated and observed CWV, precipitation rate, LE, and H is shown in Figure 2. Results are shown for simulations using the single-moment, doublemoment, and P3 microphysics schemes, in addition to the higher-resolution (P3/250m) and larger domain (P3/400km) runs.



Figure 2. Comparison of modeling results (solid colors), large-scale forcing (dashed gray), and observations (dotted) for (a) Precipitation rate (mm/hr), (b) Column water vapor (cm), (c) Latent heat flux (W/m^2) , and (d) Sensible heat flux (W/m^2) averaged over the model domain.

Simulated and observed CVW values agree very well in the first week, likely a re-253 sult of the nudging imposed during the spin-up period, after which some differences be-254 gin to appear. CWV values are generally higher for the P3 scheme and lower for the single-255 moment run. The P3 cases exhibit the strongest correlation (Pearson) with the obser-256 vations: 0.78 for P3/250m and 0.75 for both P3 and P3/400km (see the Taylor diagrams 257 in Figure S4). Conversely, the single-moment scheme shows a weaker correlation with 258 the observed CWV (0.56), while the double-moment scheme correlation is close to that 259 of the P3 scheme (0.70). For the standard deviation of CWV, the model values range 260 from 0.24 mm (P3/250 m) to 0.28 mm (Single-Moment), while the observations indicate 261 a value of 0.34 mm. Despite this difference, these statistics suggest that the model can 262

reproduce observed moisture content reasonably well for at least one month without resorting to any water vapor nudging.

The different simulations closely reproduce the observed surface precipitation rate, 265 with correlations ranging from 0.76 (P3/250m) to 0.79 (P3/400 km). The simulations ex-266 hibit better agreement for lower precipitation rates, while they tend to underestimate 267 the most intense precipitation events, which are associated with MCSs. Moreover, the 268 model precipitation did not show significant sensitivity to the microphysics, spatial res-269 olution, or domain size. We hypothesize that the model's underestimation of intense sur-270 271 face precipitation could potentially be attributed to the periodic boundary conditions. These conditions might prevent the advection of MCSs that could have developed in ar-272 eas outside the domain. Nevertheless, our validation results remain satisfactory, partic-273 ularly considering our primary focus is locally-driven STD convective transitions. 274

Observed surface fluxes are reproduced reasonably well in the model runs. LE cor-275 relations with observations vary from 0.81 (double-moment) to 0.84 (P3, P3/250m, and 276 P3/400km), while the H correlations range from 0.78 (double-moment) to 0.80 (P3, P3/250m, 277 and P3/400km). The model only slightly overestimates the standard deviation of the ob-278 served mean LE, with the difference between model runs and observations being less than 279 2 W m^{-2} . However, it should be noted that the ECOR flux measurement system pro-280 vides local measurements of surface fluxes in a grassland region (T3 site, see Figure 1b), 281 while the model provides an average for an area of $200 \times 200 \text{ km}^2$ (or $400 \times 400 \text{ km}^2$ for 282 P3/400km), entirely covered by evergreen broadleaf forest. These differences make the 283 qualitative agreement between model simulations and observations all the more remark-284 able. 285

To evaluate the surface radiation budget, Figure 3 compares modeled and observed 286 surface shortwave and longwave fluxes, including both downward and upward compo-287 nents. There is high-frequency variability in the observations that is not present in the 288 model, likely because its values correspond to horizontal averages over the domain, whereas 289 observational values are taken at the T3 site. Nevertheless, the model reproduces the 290 observations satisfactorily for downward/upward surface shortwave and upward longwave 291 fluxes (correlation ranges 0.82-0.86, Figure S5). In the case of downward longwave fluxes, 292 the correlation is weaker, ranging from 0.57 (single and double-moment) to 0.62 (P3, P3/250m, 293 and P3/400km), although these values are reasonable. 294

²⁹⁵ Overall, our simulations with different microphysics schemes compared reasonably ²⁹⁶ well with the observations considered in our validation analysis. The exception was the ²⁹⁷ column water vapor, where the P3 scheme showed a stronger correlation with the ob-²⁹⁸ servations. While neither the higher resolution (P3/250m) nor the larger domain size (P3/400km) ²⁹⁹ simulations demonstrated significant improvements over the P3 case, they significantly ³⁰⁰ increased computational costs. This motivated our choice of the P3 scheme with 500 m ³⁰¹ horizontal resolution and a $200 \times 200 \text{ km}^2$ domain size as the control run configuration ³⁰² which underlies the results presented below.

5 Sensitivity Experiments

To evaluate the role of moisture and vertical wind shear in the STD convective transition, a series of sensitivity experiments are carried out. First, we perturb the water vapor profile at low levels and the free troposphere to investigate the importance of lowand mid-level preconditioning. For vertical wind shear, we modify the structure of the low or upper-level jets to evaluate the relative importance of wind shear at different levels. The results in this section are associated with the mean composites for the four Deep or ShCu (section 3.4) simulated days.

Figure 4 shows the composite of cloud liquid (r_l) , total ice (r_i) , and rainwater (r_r) mixing ratios for the Deep and ShCu days averaged over the model domain. In addition



Figure 3. Similar to Figure 2, but for (a) downward shortwave, (b) upward shortwave, (c) downward longwave, and (d) upward longwave flux at the surface (W/m^2) averaged over the model domain.

to liquid water, we include the convective boundary layer (CBL) height (magenta dashed 313 line), defined as the height at which the buoyancy flux reaches its first local minimum. 314 The figure shows that the CBL height follows closely cloud base and the lifting conden-315 sation level, with values reaching a maximum of 1.30 km at 14:45 LST on Deep days and 316 1.41 km at 16:15 on ShCu days. Both regimes exhibit a peak in r_l associated with shal-317 low convection, below 3 km between 10-14 LST. Additionally, Deep days show two peaks 318 in r_i . The first occurs between 12-14 LST at upper levels (> 8 km), associated with deep 319 convection driven by surface heating (Martin et al., 2016; Tang et al., 2016; Zhuang et 320 al., 2017; Tian et al., 2021). The second peak occurs a few hours later, between 16-18 321 LST, associated with the late afternoon STD convection transition triggered by these 322 land-atmosphere interactions. 323

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5.1 Low-Level Moisture Experiment

For the low-level moisture experiment, a moisture perturbation within the lower 1.5 km of the domain is introduced. We introduce this perturbation by multiplying the water vapor mixing ratio by a constant factor. To ensure smoothness of the vertical profile, we linearly decrease the factor to 0 between altitudes of 1.25 km and 1.75 km. For each of the Deep days selected, the model is restarted from the control run at 02 LST and the perturbation was applied. Finally, the factors are selected such that the CWV for the perturbed profile at 02 LST drops by 1, 2, 3, and 4 mm.



Figure 4. Composites for the Deep (top) and ShCu (bottom) regime days showing the diurnal cycle of domain-averaged (a,d): Cloud liquid water (colormap) and convective boundary layer (magenta dashed line), (b,e): Total ice, and (c,f): Rain content. The Deep selection includes the 17th, 21st, 23rd, and 26th of December 2014, while the ShCu selection comprises the days of the 9th, 13th, 27th, and 28th (section 3.4).

Figure 5 depicts the 30-minute average water vapor profiles at 02:15 LST, 08:15 332 LST, and 14:15 LST, along with the daytime CWV for both moisture sensitivity exper-333 iments (low-level and free troposphere, with the latter described in the next section). The 334 low-level dry perturbations (continuous lines) diminish from nighttime to afternoon due 335 to latent heat flux and moisture tendencies. Above the perturbed region, the mixing ra-336 tio values are remarkably similar, indicating minimal vertical mixing. At 14 LST, be-337 fore the late afternoon STD convective transition, the lower tropospheric water vapor 338 for the case where CWV drops by $3 \text{ mm} (BL_{3mm})$ is similar to those for the ShCu days. 339 In terms of CWV, the experiments BL_{3mm} and BL_{4mm} demonstrate lower CWV values 340 compared to ShCu days in the early morning. However, while the ShCu composite re-341 mains relatively stable throughout the diurnal cycle, the Deep composite exhibits wa-342 ter vapor convergence, leading to higher CWV values than the ShCu days for all exper-343 iments in the afternoon. 344

Figure 6 shows the magnitude (colors) and relative (contours) difference between 345 experiments and control case for cloud liquid, total ice, and rain domain-averaged mix-346 ing ratios. In addition, the liquid water path, ice water path, and surface precipitation 347 (lines) are also included. Cloud liquid water is reduced up to 75% near the cloud base, 348 with a more extensive impact observed for drier scenarios from 10 to 12 LST. Above 3 349 km, the amount of liquid water experiences a significant increase of up to 100% during 350 the afternoon in drier scenarios, reflecting a greater presence of warm clouds. Addition-351 ally, the peak in liquid water in drier cases occurs later, shifting from 11:45 LST in the 352 control case to 13:15 LST for experiments BL_{3mm} and BL_{4mm} . 353

Ice water content shows a significant sensitivity to low-level dry perturbations. For example, the control case exhibits an ice water path maximum of 110.1 g m⁻², declining to 67.2 g m⁻², 56.2 g m⁻², 40.0 g m⁻², and reaching a minimum of 9.5 g m⁻² with a decrease in CWV by 1, 2, 3, and 4 mm at 02 LST, respectively. The relative differences compared to the control case are 33.5%, 44.4%, 60.4%, and 90.6%, respectively, indicating a non-linear decrease in the ice water path as a function of the change in CWV.



Figure 5. Moisture perturbation. Specific humidity (g/kg) profile at (a,e) 02:15 LST, (b,f) 08:15 LST, (c,g) 14:15 LST, and time series of (d,h) column water vapor (mm) for low-level (upper panels, solid colors) and free troposphere (lower panels, solid colors) moisture experiments. The Deep composite (control) is the dashed black line, and the ShCu composite is the dashed magenta line.

For the rain content, a decrease in drier scenarios is observed from early morning 360 to early afternoon. While the reduction remains insignificant when reducing CWV by 361 1 mm, contours of 50% and 75% emerge in the drier cases. In terms of precipitation rate, 362 the control case shows a peak of 0.86 mm hr^{-1} , declining to 0.78, 0.68, 0.55, and 0.42363 mm hr^{-1} with a decrease in CWV by 1, 2, 3, and 4 mm at 02 LST, respectively. The 364 decrease in the maximum precipitation rate resulting from changes in CWV can be ef-365 fectively modeled by a quadratic function (Figure S6). Note that despite the observed 366 sensitivity of ice content to low-level dry perturbations, where ice content becomes neg-367 ligible in the driest scenario, the model still produces significant amounts of warm pre-368 cipitation for all experiments. 369

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5.2 Free Troposphere Moisture Experiment

The experiments conducted in the free troposphere are analogous to the low-level moisture experiments, differing only in that the perturbation is applied above 1.5 km. We select multiplicative factors such that the CWV for the perturbed profile at 02 LST drops by 2, 4, 6, and 8 mm, respectively. These changes in CWV correspond to double what is applied in the previous section, and our choice is motivated by the weaker sensitivity of free troposphere humidity to convection, as we will show below.

The dry perturbations in the free troposphere diminish throughout the day, similarly to the low-level perturbations (Figure 5). The drier scenarios also exhibit slightly drier conditions at lower levels by early afternoon compared to the control run. The free troposphere experiments start with a lower CWV value than that of shallow days. However, the case FT_{2mm} already presents higher CWV values than the shallow composite in the early morning, while FT_{4mm} shows higher values than shallow days only in the



Figure 6. Low-level moisture experiment. The composites show the diurnal cycle of domainaveraged anomalies for (a,d,g,j) Cloud liquid water (g/kg), (b,e,h,k) Total ice (g/kg), and (c,f,i,l)Rain (g/kg) mixing ratios. Each row corresponds to a different decrease in CWV, (a-c) 1 mm, (d-f) 2 mm, (g-i) 3 mm, and (j-l) 4 mm. The colors indicate the absolute difference between each experiment and the control, while the contours show relative differences of 50% (dotted), 75% (dashed), and 100% (solid). Liquid and ice water paths (right axis, g/m^2) are presented along with cloud liquid and total ice, respectively, and surface precipitation rate (right axis, mm hr⁻¹) is shown alongside rainwater. The solid line represents the experiment, while the dashed line represents the control runs conducted during the Deep days.

afternoon (around 14 LST). For the drier cases, FT_{6mm} and FT_{8mm} , their CWV remains lower than that of the shallow days throughout the simulation.

Figure 7 presents the results for the free-troposphere moisture experiments and control runs during the Deep days. While applying a dry perturbation above 1.5 km leads to a reduction in cloud liquid water and liquid water path throughout the troposphere, this impact is relatively minor compared to what is observed in the low-level experiments.

³⁸⁹ Cloud ice water exhibits a greater sensitivity to the free troposphere perturbations. ³⁹⁰ However, the impact is still relatively minor compared to the perturbations at low lev-³⁹¹ els. For instance, when the perturbation in the free troposphere and at low levels leads ³⁹² to a 2 mm drop in CWV, the maximum ice water path is 99.1 g m⁻² and 56.2 g m⁻², ³⁹³ respectively. Similarly, with a 4 mm drop in CWV, the maximum ice water path is 60.9 ³⁹⁴ g m⁻² and 9.5 g m⁻² for the perturbations in the free troposphere and at low levels, re-³⁹⁵ spectively. Moreover, the driest free troposphere case (FT_{8mm}) still exhibits significant



Figure 7. Same as Figure 6, but for the free troposphere moisture experiment instead, where the decreases in CWV were 2, 4, 6, and 8 mm.

ice water path values, with a maximum of 45.0 g m⁻² at 17:45 LST, being even greater than the value observed for the experiment where CWV is reduced by 3 mm at low levels (BL_{3mm}, 40.0 g m⁻²).

Rain content also shows a reasonable sensitivity to the free troposphere perturba-399 tions, although it is minor compared to low-level perturbations. For example, when the 400 perturbation in the free troposphere and at low levels leads to a 2 mm drop in CWV, 401 the peak of precipitation is 0.82 mm hr^{-1} and 0.68 mm hr^{-1} , respectively. Similarly, with 402 a 4 mm drop in CWV, the peak for the free troposphere experiment remains the same 403 $(0.82 \text{ mm hr}^{-1})$ while the low-level perturbation shows 0.42 mm hr⁻¹. Finally, the dri-404 est free troposphere case (FT_{8mm}) shows a maximum of 0.64 mm hr⁻¹, which better re-405 lates with experiment BL_{2mm} (0.68 mm hr⁻¹). 406

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5.3 Wind Jet Experiment

To assess the impact of wind shear on the deepening of convective clouds in the Amazon, we perform sensitivity experiments in which the low- or high-level jets are modified. Each jet's intensity, amplitude, and altitude are changed separately. The modified wind profiles are used to force the model, with nudging applied with a timescale of 2 hours throughout the simulation.

In order to have better control of the shape of the wind profiles and easily generate sensitivity tests, we approximate the imposed wind using an analytical formula. More

specifically, considering the shape of the wind speed, we model this quantity as the su-415 perposition of two Gaussian functions, each representing a jet. The average wind speed 416 is fitted to this function, and the fitting parameters are given in Table S1. The wind di-417 rection was fitted to a piece-wise linear function. (Figure 8b). The wind direction is con-418 stant in the bottom ~ 2 km. It veers clockwise at a constant rate of about 14° km⁻¹ from 419 2 to 15 km, and counterclockwise at a rate of -28° km⁻¹ from 15 to 20 km, and it is con-420 stant above 20 km (not shown). The wind speed and direction imposed for control runs 421 are shown in Figure 8a-b (solid black line). 422



Figure 8. Wind profiles for the jet experiments control run showing: (a) large-scale wind speed and (b) wind direction, as measured (blue) and as idealized (black). Sensitivity experiments perturbed the wind speed profile of either the (c) low-level or (d) upper-level wind jets by intensifying (blue), widening (orange), shifting (green), or removing the jet (red).

For each jet, we conduct 4 experiments in which we (1) increase the jet strength, (2) increase the width, (3) shift the peak position, and (4) remove the jet entirely while keeping the wind direction constant in all cases. The modified wind profiles are shown in Figures 8c-d. For each experiment, the model is restarted from the control run at 02 LST of each one of the four Deep days selected, and the modified wind profiles are applied to force the model with a 2-h nudging timescale.

Figure 9 presents the results for the low-level jet experiments. While the jet po-429 sition significantly affects cloud liquid water, the jet width has a negligible impact. The 430 higher position possibly enhances the low-level updrafts, leading to an increase of r_l above 431 2 km, particularly in the late afternoon (around 16 LST), where r_l can increase by as 432 much as 100%. Removing the low-level jet also exhibits a similar impact on cloud liq-433 uid water, although the anomalies are smaller than those associated with the higher jet. 434 The stronger jet only slightly impacts the cloud liquid water. Although the jet influences 435 the cloud water profile, the integrated liquid water path is similar in all experiments and 436 does not significantly differ from the control. 437

The cloud ice content is more significantly affected by the low-level jet. The con-438 trol experiment exhibits peaks in ice water path at 12:45 LST and 16:15 LST. The stronger 439 jet simulation shows positive anomalies before 12 LST and after 16 LST, resulting in an 440 increase in r_i of up to 100%, with negative anomalies observed in between. For the wider 441 jet, a positive anomaly dominates throughout the diurnal cycle, especially between 12-442 17 LST, when r_i increases up to 100% above 12 km. However, there is only a modest 443 increment in ice water path. The higher jet shows a decrease ranging from 50-75% in r_i 444 around 12 LST and an increase afterwards. There is a delay of the convective activity, 445 with a suppression at 12 LST and an increase around 16 LST. When the low-level jet 446 is removed, the ice water path is substantially reduced, both at 12 and 16 LST (up to 447 75%), with the maximum being reached only around 17:45. 448



Figure 9. Low-level jet experiment. The composites show the diurnal cycle of domainaveraged anomalies for (a,d,g,j) Cloud liquid water (g/kg), (b,e,h,k) Total ice (g/kg), and (c,f,i,l) Rain (g/kg) mixing ratios. Each row corresponds to a different low-level jet perturbation, (a-c) strength, (d-f) width, (g-i) position, and (j-l) removed entirely. The colors indicate the absolute difference between each experiment and the control, while the contours show relative differences of 50% (dotted), 75% (dashed), and 100% (solid). Liquid and ice water paths (right axis, g/m²) are presented along with cloud liquid and total ice, respectively, and surface precipitation rate (right axis, mm hr⁻¹) is shown alongside rainwater. The solid line represents the experiment, while the dashed line represents the control runs conducted during the Deep days.

Rain content follows the changes in total ice mixing ratio, with negative (positive) 449 anomalies where ice decreases (increases). However, changes in r_r are less significant, with 450 smaller areas showing changes greater than 50%. While the surface precipitation shows 451 minimal impact in the stronger and wider jet experiments, it exhibits notable differences 452 for the other two cases. For the higher jet, precipitation decreases around noon and in-453 creases in the late afternoon, following the changes in ice water path. In the absence of 454 a jet, precipitation is particularly reduced between 12-15 LST, with the noon peak be-455 ing roughly 40% lower than that observed in the control runs. 456

Figure 10 presents the results for the upper-level jet experiments. There is a striking contrast with the low-level jet results. The upper-level jet affects convection only in the upper troposphere, with negligible impacts on liquid water. Overall, the experiments show an alternating increasing and decreasing ice content pattern. This is related to a delay in convection, which can be more easily noticed on the ice water path curves. Both



Figure 10. Same as Figure 9, but for the upper-level jet experiment instead.

the noon and afternoon peaks are displaced to later times, and the afternoon peak also 462 gets more intense. The exception is the wider jet experiment, where the noon peak slightly 463 increases and is not delayed. The changes are more significant after 17 LST when all ex-464 periments exhibit an increase of up to 100% in ice water. In the case of the removed jet, 465 the afternoon peak occurs 30 minutes later (16:45 LST). For the stronger and wider jet 466 experiments, the peak is delayed by 1 hour (17:15 LST), and for the jet with a relatively 467 lower position, the peak occurs 1.5 hours later (17:45 LST). Nonetheless, there is no sig-468 nificant change in the rain content and surface precipitation. 469

470 6 Discussion

While several studies in the literature have employed CRMs to simulate convec-471 tive properties in both continental (M. F. Khairoutdinov & Randall, 2003; M. Khairout-472 dinov & Randall, 2006; Henderson & Pincus, 2009; Cecchini et al., 2022) and ocean (Blossey 473 et al., 2007; M. F. Khairoutdinov et al., 2009; Liu et al., 2015; Blossey et al., 2021) re-474 gions, we conduct simulations specifically for the central Amazon by coupling a CRM 475 with a LSM, thus explicitly representing biosphere-atmosphere feedbacks, which plays 476 a crucial role on convection over the Amazon tropical forest (Silva Dias et al., 2002). Vilà-477 Guerau de Arellano et al. (2020) utilized a large-eddy model combined with a different 478 LSM to investigate the diurnal cycle of energy, moisture, and carbon dioxide from clear 479 to cloudy conditions. Nevertheless, their study relied only upon a specific day during the 480 Amazon dry season and did not address deep convection. Although the recent work by 481 Gonçalves et al. (2022) also utilized a CRM coupled with an LSM to simulate convec-482

tion in the central Amazon, they did not evaluate the evolution of moisture content and 483 surface latent and sensible heat fluxes. This restricts the validation of their simulations 484 in reproducing these crucial convective properties, which we have addressed in our study. 485 Specifically, we evaluate the model performance using single-moment, double-moment, and P3 microphysics schemes. Only the CWV exhibits a noticeable sensitivity to the mi-487 crophysics, where the P3 scheme shows the strongest correlation (Pearson, 0.8) with ob-488 servations. Varying the horizontal resolution (from 500 to 250 m) and domain size (from 489 $200 \times 200 \times 27$ to $400 \times 400 \times 27$ km³) reveals minimal sensitivity to the model results. For 490 the surface fluxes, the P3 microphysics scheme indicates a correlation of about 0.8 for 491 both latent heat and sensible heat fluxes. Overall, our results demonstrate that our sim-492 ulations satisfactorily reproduce the convective properties in the central Amazon. 493

Our sensitivity experiments indicate that the humidity in the early morning at lower 494 levels plays a crucial role in the late afternoon STD convective transition in the Ama-495 zon. By reducing water vapor in the lowest 1.5 km, the diurnal peak of the ice water path 496 substantially drops with changes in CWV, whereas a change of 2 mm causes a reduc-497 tion of ice in the range of 50-75%, and a change of 4 mm results in negligible ice during the simulation. Moreover, the maximum precipitation rate demonstrates an approx-499 imately quadratic decrease with variations in the low-level CWV. Schiro et al. (2016) ex-500 amined the relationship between precipitation and CWV by calculating the condition-501 ally averaged precipitation to CWV using local observations of both variables at corre-502 sponding times. They observed that the probability and intensity of precipitation can 503 be roughly characterized by an exponential function of CWV magnitude. It is essential 504 to highlight that our approach differs from that of Schiro et al. (2016). We used the domain-505 average for maximum afternoon precipitation, while CWV precedes the diurnal cycle. 506 This procedural difference somewhat justifies why we observed a quadratic relationship 507 instead of an exponential one, as there is no contradiction between these results. Fur-508 thermore, our findings indicated that achieving a comparable reduction in ice water path 509 or precipitation, as observed in the low-level experiments, requires reducing the column 510 water vapor by a factor of two or more in the free troposphere. 511

While M. Khairoutdinov and Randall (2006) conducted experiments that differed 512 from those designed in this study and were limited to a single idealized case during TRMM-513 LBA on February 23, 1999, they similarly indicated that free troposphere precondition-514 ing plays a minor role in convection in the Amazon. Based on GoAmazon2014/5 obser-515 vations, Ghate and Kollias (2016) noted that locally-driven precipitating days during the 516 dry season show an early morning water vapor excess above the boundary layer while 517 Zhuang et al. (2017); Tian et al. (2021) found that deep convective days exhibit a moister 518 environment extending from the surface to higher levels, regardless of the season. Schiro 519 and Neelin (2019) showed that the onset and probability of the STD transition are closely 520 linked to both lower-free-tropospheric moisture (700–900 hPa) and boundary layer mois-521 ture. Conversely, MSC likelihood rises with higher lower-free-tropospheric humidity, while 522 the relationship with boundary layer moisture is less distinct. The relative importance 523 of moisture to convection can also vary based on the regions being studied. Focusing on 524 the Tropical Western Pacific region on Nauru Island, Holloway and Neelin (2009) found 525 a strong correlation between observed precipitation and moisture variability in the free 526 troposphere, with limited variability in the boundary layer. Additionally, Bretherton et 527 al. (2004) also highlighted the importance of free-tropospheric humidity to convection 528 over the Tropical Oceans. 529

Vertical wind shear primarily impacts the peak timing of ice water in our simulations. Furthermore, our findings indicate that convection is enhanced during the afternoon when the low-level wind is idealized using a jet of larger width, moderate strength, and with a relatively higher peak position from around 2 to 4 km. Conversely, the upperlevel wind has a minor influence on convective intensity. M. Khairoutdinov and Randall (2006) designed experiments employing an idealized large-scale wind forcing and a free

wind shear environment. Similar to our results, the STD transition was not prevented 536 by removing vertical wind shear. Cecchini et al. (2022) also conducted numerical exper-537 iments to quantify the impact of vertical wind shear in the central Amazon, specifically 538 targeting shallow cumulus convection during a typical day in the dry season. By intro-539 ducing incremental changes in the large-scale wind speed across the entire vertical do-540 main, the authors observed a weakening of convective intensity, suggesting that verti-541 cal wind shear prevents the STD convective transition. Here, we have identified that the 542 vertical level of wind shear significantly influences its impact on convection. Moreover, 543 while stronger vertical wind shear suppresses the initial phase of convection, the STD 544 convective transition still occurs, albeit with a delay ranging from a few minutes to an 545 hour. 546

In contrast to prior observational studies, our modeling results offer quantitative 547 insights into the role of vertical wind shear in Amazonian convection. For example, Zhuang 548 et al. (2017) observed that ShCu days are linked to stronger mid-level wind shear dur-549 ing the wet season. We observe that when the upper-level jet is shifted from around 12 550 to 8 km, which is related to mid-level wind shear, the trigger for the STD transition is 551 only delayed by about 1.5 hours. While Chakraborty et al. (2018) indicated that more 552 intense low-level shear is associated with shallow convection during the transition sea-553 son, we observe that stronger low-level shear has little influence on cloud liquid water 554 and might only provoke a delay in the late afternoon STD transition. 555

Thus, our findings suggest that a moderate shear environment might more efficiently 556 separate downdrafts and updrafts within the cloud while concurrently organizing the con-557 vergence of low-level water vapor within the cloud layers. A wider jet leads to a smoother 558 and more gradual shift in wind shear, extending from the surface to higher altitudes, thereby 559 also organizing the water vapor convergence from below the cloud base to higher levels. 560 Given the dependence of water vapor convergence on low-level humidity, convection demon-561 strates heightened sensitivity to boundary layer humidity. Meanwhile, the upper-level 562 jet primarily impacts the extensive cloud anvil, exerting a relatively minor influence on 563 ice content above 8 km. Finally, these results collectively suggest that dynamic factors 564 may exert a more pronounced influence on convection in the Amazon. 565

566 7 Conclusions

While numerous observational studies have explored the environmental controls on 567 convection in the Amazon (Itterly et al., 2016; Ghate & Kollias, 2016; Zhuang et al., 2017; 568 Schiro et al., 2016, 2018; Chakraborty et al., 2018; Tian et al., 2021; Giangrande et al., 569 2023), we have specifically addressed this problem through high-resolution idealized sim-570 ulations. We employ the System for Atmospheric Modeling (SAM) model coupled with 571 a LSM to perform simulations for the Amazon region in December 2014. The model is 572 forced with the large-scale fields from the variational analysis, and the observations from 573 the GoAmazon2014/5 experiment are used to validate the model results. The LSM de-574 fault input parameters are modified according to in-situ and satellite observations over 575 the Amazon region, and fine-tuning tests focused on improving the model agreement with 576 the observations. The simulations consistently reproduce the observations for precipi-577 tation, column water vapor (CWV), surface latent and sensible heat fluxes, and surface 578 radiation fluxes. Sensitivity tests demonstrate that simulations conducted using a single-579 moment microphysics scheme drifted towards a drier state, while simulations with the 580 P3 microphysics scheme more closely reproduce the observed water budget. For a more 581 detailed validation of the LSM, having more comprehensive observations of the soil prop-582 erties (e.g., temperature and wetness down to 4 m) would be necessary. 583

In light of recent observational studies addressing the shallow-to-deep (STD) convective transition (Ghate & Kollias, 2016; Zhuang et al., 2017; Tian et al., 2021), our study has the advantage of conducting idealized sensitivity experiments in which only one en-

vironmental control—moisture or vertical wind shear at low or high levels—is modified 587 at a time. This approach efficiently isolates their influence in controlling convection. The 588 pre-convective humidity at low levels had the greatest impact on convection. The diur-589 nal peak in the ice water path robustly decays with changes in CWV within the lower 1.5 km. To have a comparable impact on the diurnal cycle of convection, it is necessary 591 to reduce free tropospheric CWV by approximately twice the amount in the lower lev-592 els. Vertical wind shear mainly affects the ice water peak timing. A wider low-level jet 593 of moderate strength possibly facilitates the STD convective transition by organizing low-594 level water vapor convergence and potentially separating downdrafts and updrafts within 595 the cloud. The upper-level wind shear has a minor influence over convection in the Ama-596 zon. 597

While our results provide quantitative information on the role of moisture and wind 598 shear in convection, we suggest that sensitivity experiments be conducted using differ-599 ent cloud-resolving models. For instance, SAM uses periodic lateral conditions, artifi-600 cially impacting the numerical results. Using multiple models can aid in evaluating the 601 robustness of the findings and identifying potential model biases. Although our sensitivity experiments identified that the maximum afternoon precipitation rate decreases 603 roughly quadratically with changes in pre-convective CWV, particularly in the low-level 604 experiment, it is important to note that this relationship was derived from only four val-605 ues of moisture perturbation. This limitation restricts the significance of the findings. 606 The robustness of this conclusion should be further addressed, with a particular focus 607 on understanding the associated mechanism for this relationship. Finally, we also rec-608 ommend that future studies conduct specific experiments to investigate the role of wa-609 ter vapor convergence and the effects of large-scale wind direction on deep convection. 610

611 8 Data Availability

The GoAmazon2014/5 observations are publicly available at https://www.arm.gov/ 612 research/campaigns/amf2014goamazon. The large-scale forcing data based on the vari-613 ational analysis for the GoAmazon2014/5 experiment is available at the ARM Archive: 614 http://iop.archive.arm.gov/arm-iop/Oeval-data/xie/scm-forcing/iop_at_mao/. 615 The Moderate Resolution Imaging Spectroradiometer (MODIS) data for land cover and 616 leaf area index can be downloaded through the Application for Extracting and Explor-617 ing Analysis Ready Samples ($A\rho\rho EEARS$, https://appeears.earthdatacloud.nasa 618 .gov/). The Global Land Data Assimilation System (GLDAS) data for soil tempera-619 ture and soil wetness are available at https://disc.gsfc.nasa.gov/datasets/GLDAS 620 _NOAH025_3H_2.1/summary. For further assistance concerning the model input files and 621 the necessary modifications in the SAM's source code to perform the moisture pertur-622 bations as described in section 5.1 and section 5.2, please refer to the author. 623

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Sensitivity of the Shallow-to-Deep Convective 1 Transition to Moisture and Wind Shear in the Amazon 2

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Key Points: 10

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11	•	SAM-LSM consistently reproduces the $GoAmazon2014/5$ observations for precip-
12		itation, moisture, and surface fluxes during the wet season.
13	•	Daytime convection shows a noticeable sensitivity to pre-convective low-level hu-
14		midity and a weaker response to free troposphere humidity.
15	•	Vertical wind shear has a lesser influence than humidity on the shallow-to-deep
16		convective transition.

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17 Abstract

Deep convection is the primary influence on weather and climate in tropical regions. 18 However, understanding and simulating the shallow-to-deep (STD) convective transition 19 has long been challenging. Here, we conduct high-resolution numerical simulations to 20 assess the environmental controls on the evolution of isolated convection in the Amazon 21 during the wet season. Observations and large-scale forcing derived through the constrained 22 variational analysis approach for the GoAmazon2014/5 experiments are used in the sim-23 ulations and model validation. The model consistently reproduces the GOAmazon ob-24 25 servations for precipitation, moisture, and surface fluxes of radiation, latent and sensible heat. Through sensitivity experiments, we examine the relative importance of mois-26 ture and vertical wind shear in controlling the STD convective transition. Reducing the 27 pre-convective humidity within the lower 1.5 km significantly suppresses vertical devel-28 opment and lowers the ice water path. Additionally, the maximum precipitation rate de-29 creases almost quadratically with column water vapor. Conversely, a reduction of col-30 umn water vapor above 1.5 km by a factor of two or more is necessary to produce a com-31 parable decrease in ice water path or precipitation. Moderate low-level wind shear fa-32 cilitates the STD transition, leading to an earlier peak of ice water compared to stronger 33 wind shear or its absence. Although upper-level wind shear negatively influences high 34 cloud formation, its role in controlling the STD transition is relatively smaller than that 35 of low-level wind shear. Our results help quantify the role of moisture and wind shear 36 on the STD transition, but also suggest that dynamic factors may exert a more pronounced 37 influence. 38

³⁹ Plain Language Summary

The Amazon rainforest plays a vital role in the Earth's climate system. However, 40 it is not entirely understood how environmental conditions control the evolution from 41 fair weather conditions to severe thunderstorms in regions of the deep Tropics. We ad-42 dress this problem utilizing numerical simulations that capture the interactions between 43 the forest, atmosphere, and clouds. Atmospheric modeling data developed for the GoA-44 mazon2014/5 experiment are used to initialize our Amazon-based simulations. The model 45 consistently reproduces the Amazon environment throughout the period of our simula-46 tions, which covers December 2014. Additionally, we contrast the model results between 47 the control simulation and experiments in which the moisture or wind is modified to eval-48 uate their relative importance to cloud development and precipitation. Lower tropospheric 49 moisture is critical to cloud growth. The amount of moisture in the air above 1.5 km has 50 a minor contribution to cloud development and precipitation. Low-level wind of mod-51 erate strength facilitates cloud development during the afternoon. The upper-level wind 52 negatively affects the ice formation in high clouds. These results help strengthen our knowl-53 edge of tropical convection, critical for improving numerical model performance. 54

55 1 Introduction

Deep convection dominates the weather and climate in the tropics. Nevertheless, 56 comprehending and simulating the convective processes is a formidable challenge due to 57 the wide range of spatial and temporal scales involved (Mapes et al., 2009; Moncrieff et 58 al., 2012; Zhang et al., 2013). Shallow cumulus convection, a small-scale phenomenon 59 lasting tens of minutes and covering spatial scales of the order of a few kilometers, of-60 tentimes evolves into deep convective clouds covering tens of kilometers within typical 61 time scales of 2 to 4 hours (Wu et al., 2009; Hohenegger & Stevens, 2013; Adams et al., 62 2013; Henkes et al., 2021; Powell, 2022). Moreover, deep convection frequently becomes 63 organized and experiences upscale growth into mesoscale convective systems (MCSs) with 64 lifetimes spanning hours to a day and ranging in horizontal scale from 100 km to 1,000 65 km (Houze Jr, 2004). Likewise, land-atmosphere interactions and complex physical pro-66

cesses ranging from cloud microphysics to the generation of gravity waves are intrinsically tied to deep convection (Silva Dias et al., 2002; Mapes et al., 2006; Mapes & Neale,
2011; Jewtoukoff et al., 2013; Gupta et al., 2023).

General circulation models (GCMs) rely on parameterizations of convective pro-70 cesses and typically struggle to reproduce the shallow-to-deep (STD) convective tran-71 sition over continental regions (Betts, 2002; Betts & Jakob, 2002; Bechtold et al., 2004; 72 Grabowski et al., 2006). Their simulated precipitation peaks much earlier than observed 73 (Lin et al., 2000; Betts, 2002; Collier & Bowman, 2004; Dai & Trenberth, 2004), which 74 75 is an important source of bias and uncertainty in GCMs to this day (Sherwood et al., 2014; Stevens & Bony, 2013; Itterly et al., 2018; Maher et al., 2018; Freitas et al., 2020, 76 2024). To circumvent the inherent challenges posed by convective parameterizations, cloud-77 resolving models (CRMs), which explicitly resolve the up- and downdrafts in clouds, have 78 been used to study convective processes over continental and oceanic regions. For ex-79 ample, M. Khairoutdinov and Randall (2006) conducted the first high-resolution numer-80 ical simulations to investigate the STD transition over the Amazon. Their findings high-81 lighted the importance of cold pools in forcing the development of deep convection, while 82 the impact of vertical wind shear and free tropospheric preconditioning were relatively 83 minor. As part of the EUROCS (EUROpean Cloud Systems study), Derbyshire et al. 84 (2004) evaluated the sensitivity of cumulus convection to free tropospheric humidity. Un-85 like M. Khairoutdinov and Randall (2006), they observed intense deep precipitating con-86 vection in moister scenarios, whereas only shallow convection was evident in the driest 87 scenario. Waite and Khouider (2010) conducted idealized numerical simulations over the 88 tropical Atlantic Ocean. Their study emphasized the importance of congestus precon-89 ditioning, which reduces the impact of entrainment on cloud buoyancy, ultimately lead-90 ing to the STD transition. In contrast, Hohenegger and Stevens (2013) showed that the 91 transition from congestus to deep convective clouds occurs on shorter time scales than 92 required for congestus clouds to moisten the atmosphere sufficiently. This implies that 93 dynamic factors play a more substantial role in driving convection. While CRM stud-94 ies offer valuable insights into physical convective processes, they still require validation 95 through high-resolution observations, which have typically been lacking in tropical rain-96 forests. 97

In the Amazon, important, but often limited field campaigns, have explored dif-98 ferent aspects of tropical convection. Adams et al. (2015) established the Amazon Dense 99 GNSS Meteorological Network, a one-year campaign to observe the interaction between 100 water vapor fields and deep convection. Adams et al. (2013) also utilized GNSS/GPS 101 data from a long-term single site (July 2008 to December 2011) in Manaus, Brazil to eval-102 uate the water vapor convergence associated with the STD transition and found a weak 103 and quasi-linear convergence timescale of approximately 8 hours, followed by a robust 104 and non-linear convergence timescale of approximately 4 hours during the STD transi-105 tion. Later, Adams et al. (2017) employed this dense network data to investigate how 106 vapor fields evolve spatially during during the STD transition. Their results were con-107 sistent with the 4-hour STD timescale and the spatial evolution was reflective of the wa-108 ter vapor convergence posited in the single site study. More recently, the Green Ocean 109 Amazon (GOAmazon) 2014/5 Experiment (Martin et al., 2016, 2017) was carried out 110 from 2014 to 2015 in the central Amazon, providing the most comprehensive set of ob-111 servations of clouds and aerosols in the Amazon to date. Analyzing this dataset during 112 the dry season (June-September), Ghate and Kollias (2016) noted an excess of water va-113 por above 2 km during the early morning when contrasting locally-driven precipitating 114 days and nonprecipitating days. Conversely, Zhuang et al. (2017) and Tian et al. (2021) 115 observed that deep convective days exhibit relatively higher moisture extending from the 116 surface to mid-levels in all seasons. Schiro et al. (2016) showed a robust correlation be-117 tween total column water vapor and precipitation in both the central Amazon and the 118 tropical western Pacific. Furthermore, Schiro and Neelin (2019) demonstrated a strong 119 connection between the initiation and likelihood of daytime precipitation and the bound-120

ary layer and lower free troposphere moisture content. Previous studies do not completely
agree on the relative importance of vertical wind shear. For example, while Zhuang et
al. (2017) indicated that more intense low-level and deep-layer bulk wind shears facilitate the STD transition during the dry season (June-September), Chakraborty et al. (2018)
suggested that a more intense low-level shear could inhibit deep convection during the
transition season (August-November), especially if it increases the entrainment of dry
air.

In this paper, we conduct high-resolution model simulations to assess the role of 128 moisture and vertical wind shear in controlling the STD convective transition in the Ama-129 zon. First, we focus on model validation employing GoAmazon data for the period of 130 December 2014. Then, we conduct idealized sensitivity experiments in which either mois-131 ture or large-scale wind are modified at different atmospheric levels to assess their rel-132 ative importance in the development of deep convection. The paper is structured as fol-133 lows: Section 2 shows the study area. Section 3 describes the material and methods. Sec-134 tion 4 covers the model validation. Sensitivity experiments for moisture and wind shear 135 are conducted in section 5. A discussion of the results is given in section 6. Section 7 con-136 tains the conclusions. 137

138 2 Study Region

The Amazon Basin is bordered by significant altitudes (Figure 1a), primarily in 139 the western region, where some peaks in the Andes Mountains rise well over 6,000 me-140 ters in elevation above sea level. However, the simulations are conducted over the GoA-141 mazon2014/5 campaign region in the central Amazon, where the topography can be ad-142 equately considered as an extensive plain with minimal variations (< 130 m in our do-143 main of interest, section 3.2). During the experiment, most of the observations were taken 144 from the T3 site, located 70 km downwind of Manaus, in Manacapuru (3.21°S, 60.60°W), 145 a site characterized by a pasture surrounded by forest and close to the intersection of 146 the Solimões River and Negro River (Figure 1b). 147



Figure 1. (a) Land topography and ocean depth (NOAA National Centers for Environmental Information, 2022) around the Amazon. (b) Land cover (Friedl et al., 2010) around the GoAmazon2014/5 sites. The dashed circle with a radius of 202 km centered at the T1 site (in Manaus) corresponds to the S-band radar domain. The dotted circle with a radius of 110 km shows the domain of the large-scale forcing developed for the GoAmazon2014/5 Experiment (Tang et al., 2016). We also indicate the Amazon, Solimões, and Negro Rivers on the map. Land cover is from 2014, based on the Moderate Resolution Imaging Spectroradiometer (MODIS) - International Geosphere-Biosphere Programme (IGBP) land cover classification system. The Amazon Basin contour is provided by Mayorga et al. (2012).

¹⁴⁸ **3** Material and Methods

¹⁴⁹ **3.1 Data**

For model validation, we use campaign observations of precipitation, moisture, ra-150 diation, and surface latent and sensible heat fluxes. Precipitation is based on the SIPAM 151 S-band radar measurements (Schumacher & Funk, 2018), which we average over the do-152 main of the control runs. Sensible heat flux (H) and latent heat flux (LE) are from the 153 Quality Controlled Eddy Correlation (QCECOR) Flux Measurement (ARM, 2014b). Sur-154 face radiation fluxes are from the Sky Radiation Radiometers (SKYRAD) and Ground 155 Radiation Radiometers (GNDRAD) (ARM, 2013). Column water vapor (CWV) is cal-156 culated from the balloon-borne sounding system (SONDE), which provides the vertical 157 profiles of thermodynamic conditions 4 times per day during the period of this study, 158 at 02, 08, 14, and 20 LST (ARM, 2014a). 159

Large-scale atmospheric fields of water vapor mixing ratio, temperature, wind, and 160 moisture and temperature tendencies are based on 3-hour Constrained Variational Anal-161 ysis Data (VARANAL). This assimilation product was developed using atmospheric fields 162 from the European Centre for Medium-Range Weather Forecasts (ECMWF) ERA-Interim 163 reanalysis (Dee et al., 2011), which were mainly constrained by the SIPAM S-band radar 164 precipitation rate and ARM surface fluxes through the column heat and moisture bud-165 get analysis (Tang et al., 2016). The VARANAL data represent an average over the anal-166 ysis domain centered at T1 site, covering a radius of 110 km (Figure 1b). 167

The Moderate Resolution Imaging Spectroradiometer (MODIS) provides data prod-168 ucts of land cover type (MCD12Q1 Version 6) and leaf area index (MCD15A2H Version 169 6.1) (Friedl et al., 2010). Specifically, we use the land cover product based on the Inter-170 national Geosphere-Biosphere Programme (IGBP) land cover classification system. Silt, 171 clay, and sand content in the soil are based on in-situ measurements of the soil type "Terra 172 Firme" (Terra-firma) described in Table 1 on Schaefer et al. (2017). Soil temperature 173 and wetness are based on the NASA Global Land Data Assimilation System (GLDAS) 174 Noah Land Surface Model L4 3 hourly 0.25 x 0.25 degree V2.1 (Rodell et al., 2004). 175

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3.2 Model Configuration

We employ the numerical model System for Atmospheric Modeling (SAM), ver-177 sion 6.11.8, which solves the anelastic equations of motion and uses liquid water static 178 energy, total nonprecipitating, and precipitating water as thermodynamic prognostic vari-179 ables (M. F. Khairoutdinov & Randall, 2003). The equations are solved using lateral pe-180 riodic boundary conditions. A prognostic turbulent kinetic energy 1.5-order closure scheme 181 is used to parameterize subgrid-scale effects. Different microphysics parameterizations 182 are available, including the single-moment (Morrison, 2003), double-moment (Morrison 183 et al., 2005), and Predicted Particle Properties (P3) (Morrison & Milbrandt, 2015) schemes. 184 The radiative heating can be prescribed or calculated by choosing a radiation scheme, 185 either the Community Atmosphere Model (CAM3) (Collins et al., 2006) or the Rapid 186 Radiative Transfer Model (RRTM) (Mlawer et al., 1997) schemes. Surface fluxes can be 187 prescribed or calculated using Monin-Obukhov similarity theory or a simplified Land Sur-188 face Model (LSM) (Lee & Khairoutdinov, 2015), which is only compatible with the CAM3 189 radiation scheme for the current SAM-LSM version. 190

The baseline configuration for the simulations analyzed in this paper considers a domain of $200 \times 200 \times 27$ km³. This choice was made primarily to reasonably accommodate MCSs, given that they typically span about 100 km (Houze Jr, 2004). The horizontal resolution is 500 m, and the vertical resolution varies: it starts at a minimum of 50 m below 1.5 km and increases to 300 m in the upper troposphere. From there, it gradually stretches up to 500 m at the model's upper boundary, which reaches 27 km, resulting in 128 vertical levels. The temporal resolution is 5 seconds, and instantaneous model fields and statistics are output every 30 minutes. The control simulation uses the P3 microphysics scheme. The CAM3 radiation scheme is called every 150 seconds. Surface fluxes are calculated through the LSM (see Section 3.3).

The large-scale forcing is based on the VARANAL dataset for the period of December 2014 in the central Amazon. Winds were nudged with a 2-hour timescale throughout the simulation. The water vapor mixing ratio was nudged only during the spin-up, considered as the period from 1-5 December 2014, with a timescale of 6 hours.

For the purpose of model validation, we conducted additional simulations where the only modificatio was the choice of the microphysics scheme: single-moment, doublemoment, or P3 schemes. We also assessed model sensitivity to resolution and domain size by performing additional simulations at 250 m resolution or with a 400×400×27 km³ domain.

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3.3 Land Surface Model Configuration

The simplified Land Surface Model uses a minimalist set of parameters to solve the 211 transport of heat, moisture, and radiation in the soil and vegetation and calculate the 212 transfer of momentum between the surface and the atmosphere (Lee & Khairoutdinov, 213 2015). It adequately idealizes the land-atmosphere interactions, which fundamentally in-214 fluence convection over the Amazon forest (Silva Dias et al., 2002; Betts, 2002). To de-215 fine the characteristics of the surface and the vegetation for the LSM, we use the MODIS 216 data of land cover type during 2014 and mean leaf area index (LAI) for December 2014 217 over a domain of 200×200 km² centered at T3 site (see Figure S1). These are associated 218 with the period and area of our simulations. The surface LAI was set to $4.4 \text{ m}^2 \text{ m}^{-2}$, 219 which corresponds to the domain average in satellite observations. In addition, we choose 220 to cover the surface uniformly with evergreen broadleaf forests, which cover 83% of the 221 area in observations. However, based on several tests conducted to optimize the LSM 222 parameters, we modify the default near-infrared visible albedo for vegetation from 0.20 223 to 0.30, the root length from 150 cm to 200 cm, and the displacement height factor from 224 0.68 to 0.65. The corresponding displacement height is $0.65 \times 20 \text{ m} = 13 \text{ m}$, where 20 225 m is the default value of the height of the canopy. These modifications improve the agree-226 ment between the observations and the simulated surface radiation, latent and sensible 227 heat fluxes. 228

The soil is simulated using 11 layers from the surface down to a depth of 400 cm. 229 Clay and sand contents for each layer are based on in-situ measurements on Terra-firma 230 forests (section 3.1). The initial conditions for soil temperature and soil wetness are based 231 on GLDAS Noah, which provides information on 4 layers: 0-10, 10-40, 40-100, and 100-232 200 cm. The LSM soil layers close to the surface, which experience greater diurnal cy-233 cle variation, are interpolated using the nearest neighbor method. The deeper soil lay-234 ers are interpolated (and extrapolated) linearly. The initial profile of soil temperature 235 and wetness is shown in Figure S2. 236

3.4 Cloud Regime Days

For the sensitivity experiments, we select a set of deep convective (Deep) days from 238 the control simulation (P3 scheme, with a horizontal resolution of 500 m and a domain 239 size of $200 \times 200 \times 27 \text{ km}^3$) and perturb the sounding or the large-scale forcing imposed. 240 For the Deep selection, we require that the domain average of total ice presents a dis-241 tinct deepening during the afternoon, characteristic of the STD convective transition. 242 The chosen Deep days are December 17th, 21st, 23rd, and 26th. For comparison, we also 243 select a set of shallow cumulus (ShCu) days from the control run. We identify four days 244 with negligible ice content and minimum surface precipitation: December 9th, 13th, 27th, 245

and 28th. Figure S3 shows the profile of cloud liquid and total ice for our selection of
 cloud regime days.

248 4 Model Validation

A comparison between simulated and observed CWV, precipitation rate, LE, and H is shown in Figure 2. Results are shown for simulations using the single-moment, doublemoment, and P3 microphysics schemes, in addition to the higher-resolution (P3/250m) and larger domain (P3/400km) runs.



Figure 2. Comparison of modeling results (solid colors), large-scale forcing (dashed gray), and observations (dotted) for (a) Precipitation rate (mm/hr), (b) Column water vapor (cm), (c) Latent heat flux (W/m^2) , and (d) Sensible heat flux (W/m^2) averaged over the model domain.

Simulated and observed CVW values agree very well in the first week, likely a re-253 sult of the nudging imposed during the spin-up period, after which some differences be-254 gin to appear. CWV values are generally higher for the P3 scheme and lower for the single-255 moment run. The P3 cases exhibit the strongest correlation (Pearson) with the obser-256 vations: 0.78 for P3/250m and 0.75 for both P3 and P3/400km (see the Taylor diagrams 257 in Figure S4). Conversely, the single-moment scheme shows a weaker correlation with 258 the observed CWV (0.56), while the double-moment scheme correlation is close to that 259 of the P3 scheme (0.70). For the standard deviation of CWV, the model values range 260 from 0.24 mm (P3/250 m) to 0.28 mm (Single-Moment), while the observations indicate 261 a value of 0.34 mm. Despite this difference, these statistics suggest that the model can 262

reproduce observed moisture content reasonably well for at least one month without resorting to any water vapor nudging.

The different simulations closely reproduce the observed surface precipitation rate, 265 with correlations ranging from 0.76 (P3/250m) to 0.79 (P3/400 km). The simulations ex-266 hibit better agreement for lower precipitation rates, while they tend to underestimate 267 the most intense precipitation events, which are associated with MCSs. Moreover, the 268 model precipitation did not show significant sensitivity to the microphysics, spatial res-269 olution, or domain size. We hypothesize that the model's underestimation of intense sur-270 271 face precipitation could potentially be attributed to the periodic boundary conditions. These conditions might prevent the advection of MCSs that could have developed in ar-272 eas outside the domain. Nevertheless, our validation results remain satisfactory, partic-273 ularly considering our primary focus is locally-driven STD convective transitions. 274

Observed surface fluxes are reproduced reasonably well in the model runs. LE cor-275 relations with observations vary from 0.81 (double-moment) to 0.84 (P3, P3/250m, and 276 P3/400km), while the H correlations range from 0.78 (double-moment) to 0.80 (P3, P3/250m, 277 and P3/400km). The model only slightly overestimates the standard deviation of the ob-278 served mean LE, with the difference between model runs and observations being less than 279 2 W m^{-2} . However, it should be noted that the ECOR flux measurement system pro-280 vides local measurements of surface fluxes in a grassland region (T3 site, see Figure 1b), 281 while the model provides an average for an area of $200 \times 200 \text{ km}^2$ (or $400 \times 400 \text{ km}^2$ for 282 P3/400km), entirely covered by evergreen broadleaf forest. These differences make the 283 qualitative agreement between model simulations and observations all the more remark-284 able. 285

To evaluate the surface radiation budget, Figure 3 compares modeled and observed 286 surface shortwave and longwave fluxes, including both downward and upward compo-287 nents. There is high-frequency variability in the observations that is not present in the 288 model, likely because its values correspond to horizontal averages over the domain, whereas 289 observational values are taken at the T3 site. Nevertheless, the model reproduces the 290 observations satisfactorily for downward/upward surface shortwave and upward longwave 291 fluxes (correlation ranges 0.82-0.86, Figure S5). In the case of downward longwave fluxes, 292 the correlation is weaker, ranging from 0.57 (single and double-moment) to 0.62 (P3, P3/250m, 293 and P3/400km), although these values are reasonable. 294

²⁹⁵ Overall, our simulations with different microphysics schemes compared reasonably ²⁹⁶ well with the observations considered in our validation analysis. The exception was the ²⁹⁷ column water vapor, where the P3 scheme showed a stronger correlation with the ob-²⁹⁸ servations. While neither the higher resolution (P3/250m) nor the larger domain size (P3/400km) ²⁹⁹ simulations demonstrated significant improvements over the P3 case, they significantly ³⁰⁰ increased computational costs. This motivated our choice of the P3 scheme with 500 m ³⁰¹ horizontal resolution and a $200 \times 200 \text{ km}^2$ domain size as the control run configuration ³⁰² which underlies the results presented below.

5 Sensitivity Experiments

To evaluate the role of moisture and vertical wind shear in the STD convective transition, a series of sensitivity experiments are carried out. First, we perturb the water vapor profile at low levels and the free troposphere to investigate the importance of lowand mid-level preconditioning. For vertical wind shear, we modify the structure of the low or upper-level jets to evaluate the relative importance of wind shear at different levels. The results in this section are associated with the mean composites for the four Deep or ShCu (section 3.4) simulated days.

Figure 4 shows the composite of cloud liquid (r_l) , total ice (r_i) , and rainwater (r_r) mixing ratios for the Deep and ShCu days averaged over the model domain. In addition



Figure 3. Similar to Figure 2, but for (a) downward shortwave, (b) upward shortwave, (c) downward longwave, and (d) upward longwave flux at the surface (W/m^2) averaged over the model domain.

to liquid water, we include the convective boundary layer (CBL) height (magenta dashed 313 line), defined as the height at which the buoyancy flux reaches its first local minimum. 314 The figure shows that the CBL height follows closely cloud base and the lifting conden-315 sation level, with values reaching a maximum of 1.30 km at 14:45 LST on Deep days and 316 1.41 km at 16:15 on ShCu days. Both regimes exhibit a peak in r_l associated with shal-317 low convection, below 3 km between 10-14 LST. Additionally, Deep days show two peaks 318 in r_i . The first occurs between 12-14 LST at upper levels (> 8 km), associated with deep 319 convection driven by surface heating (Martin et al., 2016; Tang et al., 2016; Zhuang et 320 al., 2017; Tian et al., 2021). The second peak occurs a few hours later, between 16-18 321 LST, associated with the late afternoon STD convection transition triggered by these 322 land-atmosphere interactions. 323

324

5.1 Low-Level Moisture Experiment

For the low-level moisture experiment, a moisture perturbation within the lower 1.5 km of the domain is introduced. We introduce this perturbation by multiplying the water vapor mixing ratio by a constant factor. To ensure smoothness of the vertical profile, we linearly decrease the factor to 0 between altitudes of 1.25 km and 1.75 km. For each of the Deep days selected, the model is restarted from the control run at 02 LST and the perturbation was applied. Finally, the factors are selected such that the CWV for the perturbed profile at 02 LST drops by 1, 2, 3, and 4 mm.



Figure 4. Composites for the Deep (top) and ShCu (bottom) regime days showing the diurnal cycle of domain-averaged (a,d): Cloud liquid water (colormap) and convective boundary layer (magenta dashed line), (b,e): Total ice, and (c,f): Rain content. The Deep selection includes the 17th, 21st, 23rd, and 26th of December 2014, while the ShCu selection comprises the days of the 9th, 13th, 27th, and 28th (section 3.4).

Figure 5 depicts the 30-minute average water vapor profiles at 02:15 LST, 08:15 332 LST, and 14:15 LST, along with the daytime CWV for both moisture sensitivity exper-333 iments (low-level and free troposphere, with the latter described in the next section). The 334 low-level dry perturbations (continuous lines) diminish from nighttime to afternoon due 335 to latent heat flux and moisture tendencies. Above the perturbed region, the mixing ra-336 tio values are remarkably similar, indicating minimal vertical mixing. At 14 LST, be-337 fore the late afternoon STD convective transition, the lower tropospheric water vapor 338 for the case where CWV drops by $3 \text{ mm} (BL_{3mm})$ is similar to those for the ShCu days. 339 In terms of CWV, the experiments BL_{3mm} and BL_{4mm} demonstrate lower CWV values 340 compared to ShCu days in the early morning. However, while the ShCu composite re-341 mains relatively stable throughout the diurnal cycle, the Deep composite exhibits wa-342 ter vapor convergence, leading to higher CWV values than the ShCu days for all exper-343 iments in the afternoon. 344

Figure 6 shows the magnitude (colors) and relative (contours) difference between 345 experiments and control case for cloud liquid, total ice, and rain domain-averaged mix-346 ing ratios. In addition, the liquid water path, ice water path, and surface precipitation 347 (lines) are also included. Cloud liquid water is reduced up to 75% near the cloud base, 348 with a more extensive impact observed for drier scenarios from 10 to 12 LST. Above 3 349 km, the amount of liquid water experiences a significant increase of up to 100% during 350 the afternoon in drier scenarios, reflecting a greater presence of warm clouds. Addition-351 ally, the peak in liquid water in drier cases occurs later, shifting from 11:45 LST in the 352 control case to 13:15 LST for experiments BL_{3mm} and BL_{4mm} . 353

Ice water content shows a significant sensitivity to low-level dry perturbations. For example, the control case exhibits an ice water path maximum of 110.1 g m⁻², declining to 67.2 g m⁻², 56.2 g m⁻², 40.0 g m⁻², and reaching a minimum of 9.5 g m⁻² with a decrease in CWV by 1, 2, 3, and 4 mm at 02 LST, respectively. The relative differences compared to the control case are 33.5%, 44.4%, 60.4%, and 90.6%, respectively, indicating a non-linear decrease in the ice water path as a function of the change in CWV.



Figure 5. Moisture perturbation. Specific humidity (g/kg) profile at (a,e) 02:15 LST, (b,f) 08:15 LST, (c,g) 14:15 LST, and time series of (d,h) column water vapor (mm) for low-level (upper panels, solid colors) and free troposphere (lower panels, solid colors) moisture experiments. The Deep composite (control) is the dashed black line, and the ShCu composite is the dashed magenta line.

For the rain content, a decrease in drier scenarios is observed from early morning 360 to early afternoon. While the reduction remains insignificant when reducing CWV by 361 1 mm, contours of 50% and 75% emerge in the drier cases. In terms of precipitation rate, 362 the control case shows a peak of $0.86 \text{ mm } \text{hr}^{-1}$, declining to 0.78, 0.68, 0.55, and 0.42363 mm hr^{-1} with a decrease in CWV by 1, 2, 3, and 4 mm at 02 LST, respectively. The 364 decrease in the maximum precipitation rate resulting from changes in CWV can be ef-365 fectively modeled by a quadratic function (Figure S6). Note that despite the observed 366 sensitivity of ice content to low-level dry perturbations, where ice content becomes neg-367 ligible in the driest scenario, the model still produces significant amounts of warm pre-368 cipitation for all experiments. 369

370

5.2 Free Troposphere Moisture Experiment

The experiments conducted in the free troposphere are analogous to the low-level moisture experiments, differing only in that the perturbation is applied above 1.5 km. We select multiplicative factors such that the CWV for the perturbed profile at 02 LST drops by 2, 4, 6, and 8 mm, respectively. These changes in CWV correspond to double what is applied in the previous section, and our choice is motivated by the weaker sensitivity of free troposphere humidity to convection, as we will show below.

The dry perturbations in the free troposphere diminish throughout the day, similarly to the low-level perturbations (Figure 5). The drier scenarios also exhibit slightly drier conditions at lower levels by early afternoon compared to the control run. The free troposphere experiments start with a lower CWV value than that of shallow days. However, the case FT_{2mm} already presents higher CWV values than the shallow composite in the early morning, while FT_{4mm} shows higher values than shallow days only in the



Figure 6. Low-level moisture experiment. The composites show the diurnal cycle of domainaveraged anomalies for (a,d,g,j) Cloud liquid water (g/kg), (b,e,h,k) Total ice (g/kg), and (c,f,i,l)Rain (g/kg) mixing ratios. Each row corresponds to a different decrease in CWV, (a-c) 1 mm, (d-f) 2 mm, (g-i) 3 mm, and (j-l) 4 mm. The colors indicate the absolute difference between each experiment and the control, while the contours show relative differences of 50% (dotted), 75% (dashed), and 100% (solid). Liquid and ice water paths (right axis, g/m^2) are presented along with cloud liquid and total ice, respectively, and surface precipitation rate (right axis, mm hr⁻¹) is shown alongside rainwater. The solid line represents the experiment, while the dashed line represents the control runs conducted during the Deep days.

afternoon (around 14 LST). For the drier cases, FT_{6mm} and FT_{8mm} , their CWV remains lower than that of the shallow days throughout the simulation.

Figure 7 presents the results for the free-troposphere moisture experiments and control runs during the Deep days. While applying a dry perturbation above 1.5 km leads to a reduction in cloud liquid water and liquid water path throughout the troposphere, this impact is relatively minor compared to what is observed in the low-level experiments.

³⁸⁹ Cloud ice water exhibits a greater sensitivity to the free troposphere perturbations. ³⁹⁰ However, the impact is still relatively minor compared to the perturbations at low lev-³⁹¹ els. For instance, when the perturbation in the free troposphere and at low levels leads ³⁹² to a 2 mm drop in CWV, the maximum ice water path is 99.1 g m⁻² and 56.2 g m⁻², ³⁹³ respectively. Similarly, with a 4 mm drop in CWV, the maximum ice water path is 60.9 ³⁹⁴ g m⁻² and 9.5 g m⁻² for the perturbations in the free troposphere and at low levels, re-³⁹⁵ spectively. Moreover, the driest free troposphere case (FT_{8mm}) still exhibits significant



Figure 7. Same as Figure 6, but for the free troposphere moisture experiment instead, where the decreases in CWV were 2, 4, 6, and 8 mm.

ice water path values, with a maximum of 45.0 g m⁻² at 17:45 LST, being even greater than the value observed for the experiment where CWV is reduced by 3 mm at low levels (BL_{3mm}, 40.0 g m⁻²).

Rain content also shows a reasonable sensitivity to the free troposphere perturba-399 tions, although it is minor compared to low-level perturbations. For example, when the 400 perturbation in the free troposphere and at low levels leads to a 2 mm drop in CWV, 401 the peak of precipitation is 0.82 mm hr^{-1} and 0.68 mm hr^{-1} , respectively. Similarly, with 402 a 4 mm drop in CWV, the peak for the free troposphere experiment remains the same 403 $(0.82 \text{ mm hr}^{-1})$ while the low-level perturbation shows 0.42 mm hr⁻¹. Finally, the dri-404 est free troposphere case (FT_{8mm}) shows a maximum of 0.64 mm hr⁻¹, which better re-405 lates with experiment BL_{2mm} (0.68 mm hr⁻¹). 406

407

5.3 Wind Jet Experiment

To assess the impact of wind shear on the deepening of convective clouds in the Amazon, we perform sensitivity experiments in which the low- or high-level jets are modified. Each jet's intensity, amplitude, and altitude are changed separately. The modified wind profiles are used to force the model, with nudging applied with a timescale of 2 hours throughout the simulation.

In order to have better control of the shape of the wind profiles and easily generate sensitivity tests, we approximate the imposed wind using an analytical formula. More

specifically, considering the shape of the wind speed, we model this quantity as the su-415 perposition of two Gaussian functions, each representing a jet. The average wind speed 416 is fitted to this function, and the fitting parameters are given in Table S1. The wind di-417 rection was fitted to a piece-wise linear function. (Figure 8b). The wind direction is con-418 stant in the bottom ~ 2 km. It veers clockwise at a constant rate of about 14° km⁻¹ from 419 2 to 15 km, and counterclockwise at a rate of -28° km⁻¹ from 15 to 20 km, and it is con-420 stant above 20 km (not shown). The wind speed and direction imposed for control runs 421 are shown in Figure 8a-b (solid black line). 422



Figure 8. Wind profiles for the jet experiments control run showing: (a) large-scale wind speed and (b) wind direction, as measured (blue) and as idealized (black). Sensitivity experiments perturbed the wind speed profile of either the (c) low-level or (d) upper-level wind jets by intensifying (blue), widening (orange), shifting (green), or removing the jet (red).

For each jet, we conduct 4 experiments in which we (1) increase the jet strength, (2) increase the width, (3) shift the peak position, and (4) remove the jet entirely while keeping the wind direction constant in all cases. The modified wind profiles are shown in Figures 8c-d. For each experiment, the model is restarted from the control run at 02 LST of each one of the four Deep days selected, and the modified wind profiles are applied to force the model with a 2-h nudging timescale.

Figure 9 presents the results for the low-level jet experiments. While the jet po-429 sition significantly affects cloud liquid water, the jet width has a negligible impact. The 430 higher position possibly enhances the low-level updrafts, leading to an increase of r_l above 431 2 km, particularly in the late afternoon (around 16 LST), where r_l can increase by as 432 much as 100%. Removing the low-level jet also exhibits a similar impact on cloud liq-433 uid water, although the anomalies are smaller than those associated with the higher jet. 434 The stronger jet only slightly impacts the cloud liquid water. Although the jet influences 435 the cloud water profile, the integrated liquid water path is similar in all experiments and 436 does not significantly differ from the control. 437

The cloud ice content is more significantly affected by the low-level jet. The con-438 trol experiment exhibits peaks in ice water path at 12:45 LST and 16:15 LST. The stronger 439 jet simulation shows positive anomalies before 12 LST and after 16 LST, resulting in an 440 increase in r_i of up to 100%, with negative anomalies observed in between. For the wider 441 jet, a positive anomaly dominates throughout the diurnal cycle, especially between 12-442 17 LST, when r_i increases up to 100% above 12 km. However, there is only a modest 443 increment in ice water path. The higher jet shows a decrease ranging from 50-75% in r_i 444 around 12 LST and an increase afterwards. There is a delay of the convective activity, 445 with a suppression at 12 LST and an increase around 16 LST. When the low-level jet 446 is removed, the ice water path is substantially reduced, both at 12 and 16 LST (up to 447 75%), with the maximum being reached only around 17:45. 448



Figure 9. Low-level jet experiment. The composites show the diurnal cycle of domainaveraged anomalies for (a,d,g,j) Cloud liquid water (g/kg), (b,e,h,k) Total ice (g/kg), and (c,f,i,l) Rain (g/kg) mixing ratios. Each row corresponds to a different low-level jet perturbation, (a-c) strength, (d-f) width, (g-i) position, and (j-l) removed entirely. The colors indicate the absolute difference between each experiment and the control, while the contours show relative differences of 50% (dotted), 75% (dashed), and 100% (solid). Liquid and ice water paths (right axis, g/m²) are presented along with cloud liquid and total ice, respectively, and surface precipitation rate (right axis, mm hr⁻¹) is shown alongside rainwater. The solid line represents the experiment, while the dashed line represents the control runs conducted during the Deep days.

Rain content follows the changes in total ice mixing ratio, with negative (positive) 449 anomalies where ice decreases (increases). However, changes in r_r are less significant, with 450 smaller areas showing changes greater than 50%. While the surface precipitation shows 451 minimal impact in the stronger and wider jet experiments, it exhibits notable differences 452 for the other two cases. For the higher jet, precipitation decreases around noon and in-453 creases in the late afternoon, following the changes in ice water path. In the absence of 454 a jet, precipitation is particularly reduced between 12-15 LST, with the noon peak be-455 ing roughly 40% lower than that observed in the control runs. 456

Figure 10 presents the results for the upper-level jet experiments. There is a striking contrast with the low-level jet results. The upper-level jet affects convection only in the upper troposphere, with negligible impacts on liquid water. Overall, the experiments show an alternating increasing and decreasing ice content pattern. This is related to a delay in convection, which can be more easily noticed on the ice water path curves. Both



Figure 10. Same as Figure 9, but for the upper-level jet experiment instead.

the noon and afternoon peaks are displaced to later times, and the afternoon peak also 462 gets more intense. The exception is the wider jet experiment, where the noon peak slightly 463 increases and is not delayed. The changes are more significant after 17 LST when all ex-464 periments exhibit an increase of up to 100% in ice water. In the case of the removed jet, 465 the afternoon peak occurs 30 minutes later (16:45 LST). For the stronger and wider jet 466 experiments, the peak is delayed by 1 hour (17:15 LST), and for the jet with a relatively 467 lower position, the peak occurs 1.5 hours later (17:45 LST). Nonetheless, there is no sig-468 nificant change in the rain content and surface precipitation. 469

470 6 Discussion

While several studies in the literature have employed CRMs to simulate convec-471 tive properties in both continental (M. F. Khairoutdinov & Randall, 2003; M. Khairout-472 dinov & Randall, 2006; Henderson & Pincus, 2009; Cecchini et al., 2022) and ocean (Blossey 473 et al., 2007; M. F. Khairoutdinov et al., 2009; Liu et al., 2015; Blossey et al., 2021) re-474 gions, we conduct simulations specifically for the central Amazon by coupling a CRM 475 with a LSM, thus explicitly representing biosphere-atmosphere feedbacks, which plays 476 a crucial role on convection over the Amazon tropical forest (Silva Dias et al., 2002). Vilà-477 Guerau de Arellano et al. (2020) utilized a large-eddy model combined with a different 478 LSM to investigate the diurnal cycle of energy, moisture, and carbon dioxide from clear 479 to cloudy conditions. Nevertheless, their study relied only upon a specific day during the 480 Amazon dry season and did not address deep convection. Although the recent work by 481 Gonçalves et al. (2022) also utilized a CRM coupled with an LSM to simulate convec-482

tion in the central Amazon, they did not evaluate the evolution of moisture content and 483 surface latent and sensible heat fluxes. This restricts the validation of their simulations 484 in reproducing these crucial convective properties, which we have addressed in our study. 485 Specifically, we evaluate the model performance using single-moment, double-moment, and P3 microphysics schemes. Only the CWV exhibits a noticeable sensitivity to the mi-487 crophysics, where the P3 scheme shows the strongest correlation (Pearson, 0.8) with ob-488 servations. Varying the horizontal resolution (from 500 to 250 m) and domain size (from 489 $200 \times 200 \times 27$ to $400 \times 400 \times 27$ km³) reveals minimal sensitivity to the model results. For 490 the surface fluxes, the P3 microphysics scheme indicates a correlation of about 0.8 for 491 both latent heat and sensible heat fluxes. Overall, our results demonstrate that our sim-492 ulations satisfactorily reproduce the convective properties in the central Amazon. 493

Our sensitivity experiments indicate that the humidity in the early morning at lower 494 levels plays a crucial role in the late afternoon STD convective transition in the Ama-495 zon. By reducing water vapor in the lowest 1.5 km, the diurnal peak of the ice water path 496 substantially drops with changes in CWV, whereas a change of 2 mm causes a reduc-497 tion of ice in the range of 50-75%, and a change of 4 mm results in negligible ice during the simulation. Moreover, the maximum precipitation rate demonstrates an approx-499 imately quadratic decrease with variations in the low-level CWV. Schiro et al. (2016) ex-500 amined the relationship between precipitation and CWV by calculating the condition-501 ally averaged precipitation to CWV using local observations of both variables at corre-502 sponding times. They observed that the probability and intensity of precipitation can 503 be roughly characterized by an exponential function of CWV magnitude. It is essential 504 to highlight that our approach differs from that of Schiro et al. (2016). We used the domain-505 average for maximum afternoon precipitation, while CWV precedes the diurnal cycle. 506 This procedural difference somewhat justifies why we observed a quadratic relationship 507 instead of an exponential one, as there is no contradiction between these results. Fur-508 thermore, our findings indicated that achieving a comparable reduction in ice water path 509 or precipitation, as observed in the low-level experiments, requires reducing the column 510 water vapor by a factor of two or more in the free troposphere. 511

While M. Khairoutdinov and Randall (2006) conducted experiments that differed 512 from those designed in this study and were limited to a single idealized case during TRMM-513 LBA on February 23, 1999, they similarly indicated that free troposphere precondition-514 ing plays a minor role in convection in the Amazon. Based on GoAmazon2014/5 obser-515 vations, Ghate and Kollias (2016) noted that locally-driven precipitating days during the 516 dry season show an early morning water vapor excess above the boundary layer while 517 Zhuang et al. (2017); Tian et al. (2021) found that deep convective days exhibit a moister 518 environment extending from the surface to higher levels, regardless of the season. Schiro 519 and Neelin (2019) showed that the onset and probability of the STD transition are closely 520 linked to both lower-free-tropospheric moisture (700–900 hPa) and boundary layer mois-521 ture. Conversely, MSC likelihood rises with higher lower-free-tropospheric humidity, while 522 the relationship with boundary layer moisture is less distinct. The relative importance 523 of moisture to convection can also vary based on the regions being studied. Focusing on 524 the Tropical Western Pacific region on Nauru Island, Holloway and Neelin (2009) found 525 a strong correlation between observed precipitation and moisture variability in the free 526 troposphere, with limited variability in the boundary layer. Additionally, Bretherton et 527 al. (2004) also highlighted the importance of free-tropospheric humidity to convection 528 over the Tropical Oceans. 529

Vertical wind shear primarily impacts the peak timing of ice water in our simulations. Furthermore, our findings indicate that convection is enhanced during the afternoon when the low-level wind is idealized using a jet of larger width, moderate strength, and with a relatively higher peak position from around 2 to 4 km. Conversely, the upperlevel wind has a minor influence on convective intensity. M. Khairoutdinov and Randall (2006) designed experiments employing an idealized large-scale wind forcing and a free

wind shear environment. Similar to our results, the STD transition was not prevented 536 by removing vertical wind shear. Cecchini et al. (2022) also conducted numerical exper-537 iments to quantify the impact of vertical wind shear in the central Amazon, specifically 538 targeting shallow cumulus convection during a typical day in the dry season. By intro-539 ducing incremental changes in the large-scale wind speed across the entire vertical do-540 main, the authors observed a weakening of convective intensity, suggesting that verti-541 cal wind shear prevents the STD convective transition. Here, we have identified that the 542 vertical level of wind shear significantly influences its impact on convection. Moreover, 543 while stronger vertical wind shear suppresses the initial phase of convection, the STD 544 convective transition still occurs, albeit with a delay ranging from a few minutes to an 545 hour. 546

In contrast to prior observational studies, our modeling results offer quantitative 547 insights into the role of vertical wind shear in Amazonian convection. For example, Zhuang 548 et al. (2017) observed that ShCu days are linked to stronger mid-level wind shear dur-549 ing the wet season. We observe that when the upper-level jet is shifted from around 12 550 to 8 km, which is related to mid-level wind shear, the trigger for the STD transition is 551 only delayed by about 1.5 hours. While Chakraborty et al. (2018) indicated that more 552 intense low-level shear is associated with shallow convection during the transition sea-553 son, we observe that stronger low-level shear has little influence on cloud liquid water 554 and might only provoke a delay in the late afternoon STD transition. 555

Thus, our findings suggest that a moderate shear environment might more efficiently 556 separate downdrafts and updrafts within the cloud while concurrently organizing the con-557 vergence of low-level water vapor within the cloud layers. A wider jet leads to a smoother 558 and more gradual shift in wind shear, extending from the surface to higher altitudes, thereby 559 also organizing the water vapor convergence from below the cloud base to higher levels. 560 Given the dependence of water vapor convergence on low-level humidity, convection demon-561 strates heightened sensitivity to boundary layer humidity. Meanwhile, the upper-level 562 jet primarily impacts the extensive cloud anvil, exerting a relatively minor influence on 563 ice content above 8 km. Finally, these results collectively suggest that dynamic factors 564 may exert a more pronounced influence on convection in the Amazon. 565

566 7 Conclusions

While numerous observational studies have explored the environmental controls on 567 convection in the Amazon (Itterly et al., 2016; Ghate & Kollias, 2016; Zhuang et al., 2017; 568 Schiro et al., 2016, 2018; Chakraborty et al., 2018; Tian et al., 2021; Giangrande et al., 569 2023), we have specifically addressed this problem through high-resolution idealized sim-570 ulations. We employ the System for Atmospheric Modeling (SAM) model coupled with 571 a LSM to perform simulations for the Amazon region in December 2014. The model is 572 forced with the large-scale fields from the variational analysis, and the observations from 573 the GoAmazon2014/5 experiment are used to validate the model results. The LSM de-574 fault input parameters are modified according to in-situ and satellite observations over 575 the Amazon region, and fine-tuning tests focused on improving the model agreement with 576 the observations. The simulations consistently reproduce the observations for precipi-577 tation, column water vapor (CWV), surface latent and sensible heat fluxes, and surface 578 radiation fluxes. Sensitivity tests demonstrate that simulations conducted using a single-579 moment microphysics scheme drifted towards a drier state, while simulations with the 580 P3 microphysics scheme more closely reproduce the observed water budget. For a more 581 detailed validation of the LSM, having more comprehensive observations of the soil prop-582 erties (e.g., temperature and wetness down to 4 m) would be necessary. 583

In light of recent observational studies addressing the shallow-to-deep (STD) convective transition (Ghate & Kollias, 2016; Zhuang et al., 2017; Tian et al., 2021), our study has the advantage of conducting idealized sensitivity experiments in which only one en-

vironmental control—moisture or vertical wind shear at low or high levels—is modified 587 at a time. This approach efficiently isolates their influence in controlling convection. The 588 pre-convective humidity at low levels had the greatest impact on convection. The diur-589 nal peak in the ice water path robustly decays with changes in CWV within the lower 1.5 km. To have a comparable impact on the diurnal cycle of convection, it is necessary 591 to reduce free tropospheric CWV by approximately twice the amount in the lower lev-592 els. Vertical wind shear mainly affects the ice water peak timing. A wider low-level jet 593 of moderate strength possibly facilitates the STD convective transition by organizing low-594 level water vapor convergence and potentially separating downdrafts and updrafts within 595 the cloud. The upper-level wind shear has a minor influence over convection in the Ama-596 zon. 597

While our results provide quantitative information on the role of moisture and wind 598 shear in convection, we suggest that sensitivity experiments be conducted using differ-599 ent cloud-resolving models. For instance, SAM uses periodic lateral conditions, artifi-600 cially impacting the numerical results. Using multiple models can aid in evaluating the 601 robustness of the findings and identifying potential model biases. Although our sensitivity experiments identified that the maximum afternoon precipitation rate decreases 603 roughly quadratically with changes in pre-convective CWV, particularly in the low-level 604 experiment, it is important to note that this relationship was derived from only four val-605 ues of moisture perturbation. This limitation restricts the significance of the findings. 606 The robustness of this conclusion should be further addressed, with a particular focus 607 on understanding the associated mechanism for this relationship. Finally, we also rec-608 ommend that future studies conduct specific experiments to investigate the role of wa-609 ter vapor convergence and the effects of large-scale wind direction on deep convection. 610

611 8 Data Availability

The GoAmazon2014/5 observations are publicly available at https://www.arm.gov/ 612 research/campaigns/amf2014goamazon. The large-scale forcing data based on the vari-613 ational analysis for the GoAmazon2014/5 experiment is available at the ARM Archive: 614 http://iop.archive.arm.gov/arm-iop/Oeval-data/xie/scm-forcing/iop_at_mao/. 615 The Moderate Resolution Imaging Spectroradiometer (MODIS) data for land cover and 616 leaf area index can be downloaded through the Application for Extracting and Explor-617 ing Analysis Ready Samples ($A\rho\rho EEARS$, https://appeears.earthdatacloud.nasa 618 .gov/). The Global Land Data Assimilation System (GLDAS) data for soil tempera-619 ture and soil wetness are available at https://disc.gsfc.nasa.gov/datasets/GLDAS 620 _NOAH025_3H_2.1/summary. For further assistance concerning the model input files and 621 the necessary modifications in the SAM's source code to perform the moisture pertur-622 bations as described in section 5.1 and section 5.2, please refer to the author. 623

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Supporting Information for "Sensitivity of the Shallow-to-Deep Convective Transition to Moisture and Wind Shear in the Amazon"

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Closed Shrublands

Nate

ciduous Broadleaf Forests Evergreen Broadleaf Forests

Evergreen Needleleaf Forests

40

80

120

x (km)

2

0

200

160

80 -

40

0 | 0

40

limões Rive

120

x (km)

80

160

200

Figure S1. (a) Land cover type and (b) LAI on SAM's coordinate. The 200x200 km² domain is centered at the T3 site (3.21°S, 60.60°W). Land cover is from 2014, and LAI is based on the average for December 2014. We also indicate in (a) the Solimões River and Negro River.



Figure S2. (a) Soil temperature and (b) soil wetness initial condition. GLDAS Noah data for 1 December 2014 at 00 UTC.



Figure S3. Cloud regime days. The first row (a,c,e,g) shows the cloud liquid, and the second row (b,d,f,h) shows the total ice mixing ratio profile for the selected Deep days. The third row (i,k,m,o) shows the cloud liquid, and the fourth row (j,l,n,p) shows the total ice mixing ratio profile for the selected ShCu days.



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Taylor Diagrams. (a) Precipitation rate. (b) Column water vapor. (c) Latent heat Figure S4. flux. (d) Sensible heat flux. The statistics correspond to the standard deviation of the mean and Pearson correlation.

0.8 1.2 Standard deviation (W/m²)

1.6

0.4

3.0 4.5 Standard deviation (W/m²)

1.5

6.0

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Figure S5. Taylor Diagrams. (a) Surface downward shortwave flux. (b) Surface upward shortwave flux. (c) Surface downward longwave flux. (d) Surface upward longwave flux. The statistics correspond to the standard deviation of the mean and Pearson correlation.



Figure S6. Maximum precipitation as a function of the change in CWV is depicted for the (a) Low-level and (b) free troposphere moisture experiments. The blue solid line with square markers represents the model results, while the orange dashed line shows the corresponding quadratic fit.

Table S1. Function parameters utilized in the idealization of the large-scale horizontal wind of the form: $v_{spd}(z) = v_0 + a_1 \exp\left(-\frac{(z-\overline{z}_1)^2}{2\sigma_1^2}\right) + a_2 \exp\left(-\frac{(z-\overline{z}_2)^2}{2\sigma_2^2}\right)$. For the control wind, $v_0 = 0.5$ m s⁻¹, $a_1 = 8.4$ m s⁻¹, $\overline{z}_1 = 2.1$ km, $\sigma_1 = 1.6$ km, $a_2 = 5.4$ m s⁻¹, $\overline{z}_2 = 11.75$ km, and $\sigma_2 = 2.5$ km. For the experiments (low- or high-level jets), the change in jet properties is achieved by varying only one control wind parameter at a time, as follows:

Experiment	Jet property	v_0	a_1	\overline{z}_1	σ_1	a_2	\overline{z}_2	σ_2
Both	Control	0.5	8.4	2.1	1.6	5.4	11.75	2.5
Low-level	$\operatorname{strength}$		13					
Low-level	removed		0					
Low-level	position			4.5				
Low-level	width				2.7			
Upper-level	$\operatorname{strength}$					10		
Upper-level	removed					0		
Upper-level	position						9	
Upper-level	width							3.5

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