Advancements in Planetary Unstructured Equivalent Source Inversion and Current Circulation Modeling Technology for Earth's Magnetic Field

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Abstract

This study presents a novel approach to modeling the Earth's geomagnetic field, which originates from electric currents approximately 2,900 km beneath the surface, crucial for understanding planetary dynamics. We introduce a method for inverting a planetary-scale equivalent magnetization source and develop a 3-D equivalent electric current circulation model from this source, enhancing understanding of these deep currents. This research signifies the first use of unstructured tetrahedral magnetization inversion technology for planet-scale magnetic data interpretation and equivalent source model construction. Validated through a synthetic case study, the method is applied to the International Geomagnetic Reference Field (IGRF) and SWARM satellite datasets, comprising 35,768 magnetic vectors from two orbital altitudes. Employing various mesh configurations, we construct and compare detailed current source models from these datasets. The effectiveness of our equivalent current sources is confirmed by comparison with dynamo research findings, demonstrating significant advancements in geomagnetic field modeling, particularly in interpretability, and providing novel insights into Earth's magnetic phenomena.

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2	Circulation Modeling Technology for Earth's Magnetic Field				
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16	Key Points:				
17	• Introduction of a Novel Technology for Constructing Planetary Equivalent Magnetization				
18	Sources of Geomagnetic Field.				
19	• Inversion of 3-D Unstructured Tetrahedral Equivalent Sources in Cartesian Coordinates for				
20	Geomagnetic Modeling.				
21	• Enhancing Geomagnetic Field Analysis through 3-D Electric Current Circulation Modeling				
22	with an Equivalent Source.				

23 Abstract

24 This study presents a novel approach to modeling the Earth's geomagnetic field, which 25 originates from electric currents approximately 2,900 km beneath the surface, crucial for 26 understanding planetary dynamics. We introduce a method for inverting a planetary-scale 27 equivalent magnetization source and develop a 3-D equivalent electric current circulation model 28 from this source, enhancing understanding of these deep currents. This research signifies the first 29 use of unstructured tetrahedral magnetization inversion technology for planet-scale magnetic data interpretation and equivalent source model construction. Validated through a synthetic case 30 31 study, the method is applied to the International Geomagnetic Reference Field (IGRF) and 32 SWARM satellite datasets, comprising 35,768 magnetic vectors from two orbital altitudes. 33 Employing various mesh configurations, we construct and compare detailed current source 34 models from these datasets. The effectiveness of our equivalent current sources is confirmed by 35 comparison with dynamo research findings, demonstrating significant advancements in 36 geomagnetic field modeling, particularly in interpretability, and providing novel insights into 37 Earth's magnetic phenomena.

38

39 Plain Language Summary

The Earth's geomagnetic field plays a crucial role in shielding us from harmful solar winds and cosmic rays. This protective field arises from electric currents generated within the Earth's core, a dynamic realm composed of circulating liquid iron and nickel around a solid inner core. Measuring these deep-seated currents directly poses a significant challenge due to their inaccessibility. To date, scientists have effectively modeled the Earth's magnetic field using advanced techniques like spherical harmonic analysis and the concept of magnetic dipoles.

However, a major hurdle remains in accurately depicting the geomagnetic field using 46 magnetization volumes of arbitrary shapes and in unraveling the intricate patterns of the core's 47 current circulation. In this study, we unveil a novel method that enables the creation of a detailed, 48 49 three-dimensional model of the geomagnetic field's equivalent magnetization source comprising millions of unstructured tetrahedrons, provides a deeper and more nuanced insight into the 50 51 Earth's inner workings. Our comprehensive 3-D model closely aligns with the currents' 52 circulation patterns observed in dynamo research, affirming the validity and effectiveness of our 53 approach. This pioneering method marks a significant leap forward in our understanding of the 54 geomagnetic field, offering an invaluable tool for future research in Earth sciences.

55

56 **1. Introduction**

57 The Earth's magnetic field is generated by electric currents flowing deep within the core, approximately 2,900 km beneath the surface. Directly measuring these internal currents presents 58 59 formidable challenges due to the inaccessible and heterogeneous nature of the deep core media. 60 The geomagnetic field's secular variations, exemplified by phenomena such as the South Atlantic Anomaly (SAA), asymmetry in the main dipolar field, intense equatorial gyres at the core-mantle 61 62 boundary (CMB) (Jackson 2003), and significant static flux bundles observed in Canada and Siberia (Bloxham & Gubbins 1985), are all attributed to these fluid electric currents (Davies & 63 64 Constable 2017). Consequently, modeling the geomagnetic field to explore these internal electric 65 current sources has been a significant and extensively researched task over recent decades.

Spherical harmonic analysis (SHA) has been historically the primary method for 66 describing the geomagnetic field (Whaler & Gubbins, 1981; Alken et al., 2021). To gain a deeper 67 68 understanding of the geomagnetic field's source and to predict short-term secular variations, 69 researchers have introduced equivalent source methods, such as the magnetic dipole approach 70 (Dampney, 1969; Hansen & Miyazaki, 1984). Bhattacharyya and Chan (1977) suggested the use 71 of a set of magnetic dipoles placed on a surface as the equivalent source for interpolating 72 magnetic data across various surfaces. Initially introduced by Mayhew (1979) for crustal 73 geomagnetic field modeling, the Equivalent Source Dipole (ESD) method has since been widely 74 applied in diverse planetary magnetic field studies. For instance, Purucker et al. (2000) used the 75 ESD method to invert the radial magnetic field of Mars. Langlais et al. (2004) enhanced the ESD 76 technique to model the Martian lithospheric magnetic field's three components. Oliveira et al. 77 (2015) developed a time-dependent equivalent source dipole model for planetary magnetic fields 78 and their secular variations. Buffett (2015) investigated geomagnetic polarity transitions using an ESD model, and Langlais et al. (2019) introduced a new ESD model of the Martian crustal magnetic field using data from the MGS and MAVEN satellites. Most recently, Hood et al. (2021) employed the ESD technique to produce a comprehensive map of the lunar crustal magnetic field. These studies collectively demonstrate that the dipole source equivalent is a potent and efficient approach for addressing the complexities of the geomagnetic field and achieving well-fitting data.

85 Additional efforts have been directed towards improving source information in geomagnetic field models by using current loops as more realistic equivalent sources for aligning 86 87 with observations of the geomagnetic field (Zidarov & Petrova, 1974). Peddie (1979) created 88 eleven sets of current loops based on the Biot-Savart law, designed to fit the geomagnetic field. 89 He proposed that these loops could effectively be considered as dipoles with negligibly small 90 radii. Expanding on this concept, Demina and Farafonova (2016) investigated the possibilities 91 and errors involved in estimating the parameters of current loops of varying radii. Recently, 92 Rong et al. (2021) methodically inverted parameters of single circular loops to analyze the 93 geomagnetic dipole field. They utilized magnetic vectors sampled from Swarm satellites, 94 providing a comprehensive discussion on the equivalence between dipole and current loop sources in geomagnetic modeling. 95

While dipole and equivalent current loop sources have received extensive research attention over the past decades, there has also been a growing focus on the mesh-based volume equivalent source approach in small-scale regional magnetic surveys. Hansen and Miyazaki (1984) conceptualized a planar equivalent source composed of divided structured mesh elements. Li and Oldenburg (2010) innovated further by employing a sparse representation in the wavelet domain to construct a mesh-based equivalent susceptibilities source. Their work demonstrated 102 that this approach enhances the stability of data modeling in low-latitude regions. Pilkington and 103 Boulanger (2017) conducted a comparative analysis between source-based and field-based 104 methods, highlighting the potential of equivalent source methods for high-accuracy regional 105 magnetic field modeling. In an effort to minimize data prediction misfits, Li et al. (2019) 106 expanded the mesh-based source layer from a single to a double layer. More recently, Zuo et al. 107 (2020) showed that the accuracy of geomagnetic field continuation and transformation could be 108 further improved by approximating the actual source with a 3-D deep mesh and depth-weighted 109 functions, thereby enhancing the resemblance of the inverted model to the actual source.

110 While numerous methods employing volume equivalent sources exist, their application in 111 the study of planetary magnetic fields remains unexplored. A primary challenge arises from the 112 fact that these methods are predominantly tailored to model secondary fields, which are induced by Earth's geomagnetic field. In contrast, planetary geomagnetic fields, originating from internal 113 114 dynamo processes, or remanent magnetization, exhibit fundamentally different mechanics. 115 Additionally, fitting observational data within spherical or Cartesian coordinates using an 116 equivalent source proves more arduous for planetary fields than for regional magnetic field 117 modeling in structured meshes. Recent advancements demonstrate the efficacy of utilizing a 3D 118 unstructured grid for magnetic data inversion (Zuo et al., 2021). The approach shows significant 119 potential in accurately recovering sources with complex geometries and enhancing data fitting 120 precision. Building upon previous research in magnetization dipole equivalent sources and 121 volume equivalent source studies, we propose the development of a 3-D unstructured 122 magnetization equivalent source. This approach would involve employing a magnetization 123 forward modeling method within Cartesian coordinates, in conjunction with the development of 124 an equivalent current sources model, which aimed at enhancing the understanding of the

geomagnetic field. The potential and practical applicability of this method are corroborated through both synthetic and real data examples, drawing upon insights from previous dynamo studies.

128 2.1 Forward Modeling of Unstructured Magnetization Sources

129 Consider a static magnetic field generated by an unstructured magnetization vector model 130 $\mathbf{m} = [\mathbf{m}_x, \mathbf{m}_y, \mathbf{m}_z]$, under the assumption of no free currents. In accordance with Maxwell's 131 equations, this magnetic field **B** can be described by the Poisson equation, represented as Eq. 132 (1):

133
$$\begin{array}{c} -\nabla \cdot \left(\mu_0 (\nabla \phi - \mathbf{m})\right) = 0 \\ \mathbf{B} = \mu_0 \nabla \phi \end{array} \quad \text{in } \Omega \in \Re^3 \tag{1a}$$
(1b)

In this context, ϕ represents the potential field within the 3-D modeling domain Ω , as defined in Eq. (1a), and μ_0 is the permeability of free space. We utilized the third boundary condition, which describes the behavior of ϕ and its derivatives $\partial \phi / \partial \mathbf{n}$ at a specific distance rand direction $\hat{\mathbf{r}}$ from the source \mathbf{m} to the modeling boundary, as expressed in Eq. (2).

138
$$\frac{\partial \phi}{\partial \mathbf{n}} + \left(\frac{\hat{\mathbf{r}} \cdot \hat{\mathbf{n}}}{r}\right) \phi = 0 \qquad \text{on } \Gamma_s \tag{2}$$

139 Here, $\hat{\mathbf{n}}$ represents the outward normal vector of the boundary Γ_s . Subsequently, a weak 140 form of the governing equations is derived using Galerkin's method, as outlined in Eq. (3):

141
$$R = \int_{\Omega} \phi \nabla w d\Omega + \int_{\Omega} w \nabla \cdot \mathbf{M} d\Omega - \int_{\Gamma_{s}} w \frac{\partial \phi}{\partial \mathbf{n}} d\Gamma_{s}$$
(3)

142 In this framework, w denotes the basis function. We have developed a classical quadratic

finite element discretization scheme, as described by Jin (2002), within Cartesian coordinates. This scheme is employed to solve for the unknown potential field ϕ , details of which are elaborated in Text S1 and illustrated in Fig. S1. Subsequently, the magnetic field is estimated according to $\mathbf{B} = \mu_0 \nabla \phi$.

147 2.2 Equivalent Magnetization Model Inversion

An unstructured mesh Finite Element (FE) inversion scheme has been designed, building upon the forward modeling method. Taking into account the planetary spherical inverse model within Cartesian coordinates, we incorporate both the depth-weighted function and the regularization term into the inverse objective function. This is succinctly presented in Eq. (4).

152
minimze
$$\phi = \phi_d + \phi_m = \frac{1}{2} \left\| \mathbf{W}_d \left(\mathbf{F}(\mathbf{m}) - \mathbf{d}^{obs} \right) \right\|^2 + \frac{\beta}{2} \left\| \mathbf{W}_s \mathbf{W}_r (\mathbf{m} - \mathbf{m}_0) \right\|^2$$
(4)
where $\mathbf{W}_r^2 = \frac{1}{(R+h)^3}$

153 **m** represents the magnetization vector, which is derived by inverting the observational data, $\mathbf{d}^{obs} = [\mathbf{B}_x, \mathbf{B}_y, \mathbf{B}_z]$. The term \mathbf{W}_d corresponds to the measurement bias of the data. The 154 155 variable R denotes the radial distance from the discretized element to the sphere's surface, while 156 *h*h signifies the altitude of the observation. \mathbf{m}_0 is the reference model vector. The spatially dependent weighting function, denoted as W_r^2 , is calculated based on the volume of each 157 tetrahedron element in the mesh. \mathbf{W}_s is the smoothing matrix. The regularization coefficient, β , 158 159 can be estimated by balancing the data misfit error and the model dispersion. To solve the 160 objective function iteratively, a nonlinear conjugate gradient optimization method, commonly 161 applied in electromagnetic field inverse problems, is employed (referencing Newman & 162 Alumbaugh 2000). Furthermore, we have derived the gradient function of this objective function, 163 as shown in Eq. (5)

164
$$\frac{\partial \Phi}{\partial \mathbf{m}} = \mathbf{J}_{s}^{T} \mathbf{W}_{d}^{T} \mathbf{W}_{d} \left(F(\mathbf{m}) - \mathbf{d}^{obs} \right) + \beta \mathbf{W}_{s}^{T} \mathbf{W}_{s} \mathbf{W}_{r}^{T} \mathbf{W}_{r} \left(\mathbf{m} - \mathbf{m}_{0} \right)$$
(5)

165 where \mathbf{J}_s is the sensitive matrix, as Eq. (6) expressed.

166
$$\mathbf{J}_{s} = \begin{bmatrix} \partial \mathbf{B}_{x} / \partial \mathbf{m} \\ \partial \mathbf{B}_{y} / \partial \mathbf{m} \\ \partial \mathbf{B}_{z} / \partial \mathbf{m} \end{bmatrix} = \mu_{0} \begin{bmatrix} Q_{x} \\ Q_{y} \\ Q_{z} \end{bmatrix} \frac{\partial \phi}{\partial \mathbf{m}}$$
(6)

167 2.2 Estimation of the Equivalent Current Circulation Mode

The toroidal currents \mathbf{J}_{T} within the Earth's core are integral in generating the observed geomagnetic field **B**. In this study, we propose a methodology to estimate an equivalent current circulation model \mathbf{J}_{m} based on the inverted equivalent 3-D magnetization model **m** derived from **B**, as delineated in Eq. (7).

172

$$\mathbf{J}_{\mathbf{m}} = \nabla \times \mathbf{m} \tag{7}$$

173 It is a typical unique solving problem that can directly transform a magnetization model ${\bf m}$ to a corresponding current model ${\bf J}_{{\bf m}}$. In accordance with classical electromagnetic theory, 174 this processes adhere to Ampère's law (Text S3). This electric current model constructing scheme 175 176 avoid using the inaccessibility of physical properties for the deep, inhomogeneous core media. As for solving the magnetization model is a non-unique problem, the derived current model J_m 177 178 is also a non-unique equivalent magnetization electric current. We transfer the non-real existing 179 magnetization vectors \mathbf{m} to an equivalent current circulation model, which is an approximated 180 solution for conjecturing the actual electric currents in the deep core interior to facilitate and

181 enhance the understanding of the origination geomagnetic field.

Transforming a magnetization model \mathbf{m} into a corresponding current model $\mathbf{J}_{\mathbf{m}}$ 182 183 represents a typical unique solution problem, aligning with classical electromagnetic theory and 184 adhering to Ampère's law, as detailed in Text S3. This method of constructing an electric current 185 model circumvents the challenges posed by the inaccessibility of physical properties in the deep, 186 inhomogeneous core. Given the non-uniqueness of the magnetization model solution, the 187 resultant current model $\mathbf{J}_{\mathbf{m}}$ also embodies a non-unique equivalent of the magnetization electric 188 current. We convert the non-realistically existing magnetization vectors \mathbf{m} into an equivalent 189 current circulation model. This approach serves as an approximated solution, offering insights 190 into the actual electric currents deep within the core and thereby facilitating a better 191 understanding of the geomagnetic field's origin.

192 **3. Results**

3.1 Synthetic Example

194 To validate our inverse scheme, we use a synthetic model that consists of a highconductivity circular ring. This ring has an outer diameter of 2,000 kilometers and a width of 100 195 196 kilometers, and it carries a current with an intensity of 10 A/m², as depicted in Fig. 1a., designed 197 to simulate an equatorial toroidal current beneath the core-mantle boundary (CMB) within the 198 Earth's core. Employing Eq. (7), we derive the equivalent magnetization model **m** for this 199 current. As illustrated in Fig. 1a, the distribution of the equivalent magnetization vector model is 200 notably non-uniform (see subplot in Fig. 1a), with the highest magnetization vector magnitudes 201 (1.1E6 A/m) concentrated near the circular surface, decreasing rapidly towards the ring's center. 202 This magnetization model \mathbf{m} effectively replicates a magnetic field \mathbf{B} with an overall dipolefield characteristic, as shown in Fig. 1b. For our synthetic case, we use 35,768 magnetic observation vector data points, evenly distributed at altitudes of 450 km and 530 km, to simulate measurements from a dual-satellite orbit scenario. To enhance visual clarity, only half sections of these observation points are shown in Fig. 1a. The inversion mesh in our study comprises 2,337,505 unstructured tetrahedrons (illustrated in Fig. S1), ensuring a finely divided and uniform mesh for the core space."





210

Figure 1: Synthetic Forward and Inversion Example

a. Simulated Circular Current (indicated by green arrows) within the core region (represented by

a red sphere section) alongside the equivalent magnetization model (depicted with light blue

arrows) in the subplot. Observation sphere's surface at altitudes of 480 km and 530 km is shown,
with the outer section radius exaggerated for clarity. b. Simulated Magnetic Field visualized in 3D space. c. Section of the Inverted Magnetization Model, with the actual magnetization model
denoted by a gray circle. d. Recovered Magnetization Model, highlighting areas where values
exceed 2,000 A/m (gray circle). e. Comparative analysis of the recovered and actual current in
the core's Southern hemisphere. f. Detailed section of Recovered Magnetization Vectors, with an
enlarged detail in the subplot.

In our inversion process, we achieved a well-converged misfit of 0.68%, leading to the successful recovery of a 3D magnetization model, as depicted in Figs. 1c and 1d. From this magnetization model, we derived a current circulation model, which is presented in Fig. 1e. The magnetic field generated by the circular current displays a distinct downward dipole-like field characteristic (Fig. 1b), closely resembling the dipole component of the Earth's magnetic field in terms of strength and 3D orientation.

226 The inverted equivalent magnetization source \mathbf{m} is positioned at depths that correspond 227 accurately with the actual model, as illustrated in Figures 1c and 1d. Its maximum magnitude 228 reaches 3,171 A/m, noticeably smaller than that of the actual model. This discrepancy partly 229 arises because the volume of the recovered model \mathbf{m} is larger than that of the actual one. 230 Notably, the estimated current, positioned at a depth of 3,179 km, is shallower by 600 km 231 compared to the actual currents at 3,779 km depth, as illustrated in Fig. 1e. To understand the 232 origin of this bias, we examined a cross-section of the inverted magnetization vectors (Fig. 1f). A 233 comparison of the recovered magnetization vectors near the source region (subplot in Fig. 1f) 234 with the actual magnetization source vectors (subplot in Fig. 1a) reveals significant differences in 235 their orientations. The directions of the inverted magnetization vectors align with the primary

236 dipole field direction (downwards in the source region). This is consistent with classical inverse 237 procedures, where only magnetization sources that conform to the observed field direction can be 238 recovered. For example, the anomalous field generated by locally opposing magnetization 239 vectors (upwards) in the rotational magnetization source (Fig. 1a, subplot) is neutralized by the 240 intense dipole field and becomes indiscernible in the observational data, and the corresponding 241 upwards magnetization sources also remain unrecovered. Therefore, this bias originates from the 242 limited information contained in the observational data and cannot be rectified merely by 243 improving the inversion method.

244 **3.2 IGRF Example**

245 Magnetic vector data points, totaling 35,768, gathered at altitudes of 450 km and 530 km 246 above the Earth's surface, are sourced from the IGRF-13 dataset (Alken et al., 2021). These data 247 points follow the configuration depicted in Text S2 and are collected with the time parameter 248 'frozen' at January 1, 2023. In our analysis, we also take into account the electrical conductivity 249 of the inner core, incorporating this region into the inversion process. The inversion mesh 250 utilized is the same as the one used in the synthetic experiment (as shown in Fig. S1). Notably, 251 the inversion process achieves a commendable data prediction error level of 1% for the core's 252 equivalent magnetization model \mathbf{m} . Subsequently, the current circulation model is determined 253 based on this magnetization model, as illustrated in Figures 2a to 2f.



254

255 Figure 2: Inverted Electric Current Circulation Model Derived from the IGRF Dataset.

a. Horizontal Section of the Current Model within the Core Region. The large semi-gray sphere

is a visual aid representing the Earth's surface for locational context. b. Vertical Section of the

258 Current Model along 0° Longitude, showing the current distribution and flow at this meridian. c. 259 Vertical Section of the Current Model along 90° Longitude, illustrating the currents along this 260 longitudinal slice. d. Vectors Representation of the Current Model beneath the African 261 Hemisphere within the Core. A gray sphere, positioned 800 km below the core-mantle boundary 262 (CMB), is used to enhance the visibility of shallow intense currents. The sizes and lengths of the 263 vectors are indicative of the corresponding current intensities. e. Vectors Representation of the 264 Current Model beneath the Antarctic Hemisphere, depicting the current flow and intensity in this region. f. Vectors Representation of the Current Model beneath the Arctic Ocean Hemisphere. 265

As illustrated in the sectional view of Fig. 2a, intense currents are predominantly concentrated near the core-mantle boundary (CMB). The highest-intensity currents within the current circulation model are located in the equatorial region. Notably, beneath the South American continent, the equatorial gyre exhibits a discontinuity (indicated by red arrows) and extends into the deep core interior. Fig. 2b and 2c showcase the presence of weaker toroidal circular currents throughout the core space, which will be discussed in the following section.

Vector maps have been employed to analyze the circulation patterns of these intense 272 273 currents near the CMB surface (from the CMB up to a depth of 800 km). As depicted in Fig. 2d, 274 the most intense electric currents traverse the African hemisphere in a clockwise direction, 275 spanning from the eastern Indian Ocean to the Pacific Ocean. A significant retrograde current 276 branches off beneath the South Atlantic Ocean, moving eastward and forming a coherent gyre 277 structure. This gyre then connects with a high-latitude currents gyre observed in the Antarctic 278 region, as shown in Fig. 2e. In contrast, currents in the North Pole region, detailed in Fig. 2f, 279 cover a broader area but with lower intensities. Three relatively weaker current gyres are 280 identified in Siberia, Canada, and the Arctic Ocean, as indicated in Fig. 2f.

281 **3.2 SWARM Example**

282 We selected a total of 24,192 magnetic vector field data points from the Swarm Level 1b data product, collected via two distinct satellite orbits: Alpha, orbiting at an altitude of 475 km, 283 284 and Bravo, at 502 km above the Earth's surface (Dataset S1). These data were gathered during 285 August 22-27, 2022, specifically between local times 22:00 and 03:00, to minimize ionospheric contributions. The simultaneous inversion of two satellite vector datasets was chosen for two 286 287 primary reasons. First, the varying altitudes provide rich, gradient magnetometry-like 288 information. Second, using the same number of observation points at a single altitude would 289 significantly increase the number of divided cells in the model. Distributing observations across 290 two altitudes helps mitigate this issue by avoiding the introduction of unexpected divided cells.

To ensure the reliability of the data under varying magnetic conditions, we applied stringent filters: Kp < 2, |Dst| < 20 nT, and $F10.7 < 103 \ 10^{22} Wm^{-2}Hz^{-1}$. This filtration was crucial to maintain data quality and focus on periods with relatively stable magnetic conditions. Additionally, we restricted the data to intervals of 500 measurement points to achieve as close to a homogenous distribution as possible, as depicted in Fig. S3.



296

297 Figure 3: Inverted Electric Current Circulation Model Derived from the Swarm Dataset.

a. Horizontal Section of the Current Model within the Core Region, showcasing the layout and
flow of currents. b. Vertical Section of the Current Model along 0° Longitude, illustrating the

current distribution and dynamics at this meridian. c. Vertical Section of the Current Model along
90° Longitude. d. Vector Representation of the Current Model beneath the African Hemisphere
within the Core, showing the pattern and intensity of currents. e. Vector Representation of the
Current Model beneath the Antarctic Hemisphere, highlighting the current flows in this region. f.
Amplified Vector Representation of the Current Model beneath the Arctic Ocean Hemisphere.
The arrow lengths in each subplot are zoomed to emphasize the gyre structure, and they are not
to scale across all subplots. This amplification approach is consistent across IGRF experiments

The inversion results of the Swarm dataset are illustrated using the plotting template of Fig. 2 for comparative analysis. The sections depicting inverted currents in the *x*, *y*, and *z* directions within the core space are presented in Figs. 3a, 3b, and 3c, respectively. Figs. 3d, 3e, and 3f provide insights into the directions and intensities of currents beneath the CMB, extending 800 km deep below the equatorial, Antarctic, and Arctic Ocean regions. An animated version of this model is available as Movie S1, and for a comprehensive overview of the full set of Swarm current examples, refer to Fig. S4.

314 The Swarm model still exhibits an intense toroidal current structure, as shown in Fig. 3a. 315 Although the maximum current intensity of this Swarm section is 3.6E-3 A/m², which is less than 316 that of the corresponding IGRF model section (4.4E-3 A/m^2 , Fig. 2a), the distributions of the 317 equatorial current gyre in both the Swarm and IGRF model sections (Figs. 2a and 3a) are similar. 318 This includes the discontinuity of the equatorial gyre beneath the South American continent. 319 Similar patterns in the inverted current circulation are also observed in the other directional 320 sections (Figs. 3b and 3c) compared to the IGRF model (Figs. 2b and 2c). For the shallow 321 currents beneath the CMB (0-800km), the Swarm model displays a similar equatorial intense 322 gyre extending from the eastern Indian Ocean to the Pacific Ocean in a clockwise direction,

connecting to a retrograde current gyre beneath the South Pacific Ocean (Fig. 3d). The intense
current gyre located below the Antarctic hemisphere (Fig. 3e) and the weaker current gyres
observed in Siberia, Canada, and the Arctic Ocean (Fig. 3f) are consistent with those in the IGRF
inverted current model (Figs. 2e and 2f), exhibiting similar locations and flow patterns.

327 To further illuminate the 3-D structure of the inverted current model throughout the entire 328 depth of the Earth's core, we implemented a filtering procedure to remove all currents with 329 magnitudes less than 5.6E-4 A/m². This process allowed us to create a 3D current vector map 330 focused predominantly on the more intense current flows. Given that the current circulation 331 patterns of the IGRF model and the Swarm model are similar, only the Swarm model is depicted 332 for simplicity and clarity. To provide a comprehensive visualization of this complex structure, we 333 have divided the 3-D map into two vertical section segments: the Eurasia hemisphere and the 334 American hemisphere. These segments are displayed in both top-view and side-view 335 perspectives, as shown in Figs. 4a-d.



- 336
- 337

Figure 4: 3-D Structure of the Recovered Core Current Model.

338 a. Top View of the Eurasia Hemisphere Current Model: This profile displays the 0-180° 339 meridian at the North Pole. Gray sphere sections represent the CMB surface of the Southern 340 Hemisphere (inner surface of the sphere), with an overlay map indicating specific positions. b. 341 Side View of the Current Model depicted in Fig. 4a. The dynamics and spatial arrangement of 342 the currents are further detailed in Movie S2, the animated version of this figure. c. Top View of 343 the Current Model in the American Hemisphere: This profile shows the 180°-360° meridian at 344 the North Pole. Gray sphere sections again mark the CMB surface of the Southern Hemisphere (inner surface of the sphere). d. Side View of the Current Model depicted in Fig. 4c. For an 345 346 animated and dynamic representation of this model, refer to Movie S3.

347

The most intense currents in both hemispheres are identified as the equatorial gyre currents near the CMB, consistent with the *x*-direction model sections shown in Figs. 2a and 3a. In the deeper core, a mid-scale toroidal gyre branches off from the inner side of this intense equatorial gyre, as depicted in Fig. 4a. This mid-scale gyre primarily originates from the Eurasia hemisphere, extending down to the core's center. From the side view in Fig. 4b, it's evident that this inner gyre also extends vertically, spanning from the Northern Hemisphere's high-latitude regions to the South Pole's high latitudes.

355 The other part of the mid-scale toroidal gyre is illustrated in the sections for the other 356 hemisphere (Figs. 4c and 4d). As an integral part of the inner gyre circulation, these flows 357 continue in a clockwise direction, merging into the outer, more intense equatorial gyre (Fig. 4c). 358 From the side view in this hemisphere, depicted in Fig. 4d, the inner gyre currents extend to the 359 high-latitude region below the Antarctic continent, connecting with the middle-scale shallow 360 gyre in the region (Fig. 3e). In summary, alongside the prominent large-scale equatorial gyre, 361 there appears to be a tangentially oriented, cylindrical mid-scale toroidal gyre deep within the 362 core, contributing to a complex overall circulation pattern.

To evaluate the inverted equivalent magnetization source, we estimated the radial magnetic flux at the CMB based on this source, as shown in Fig. 5. The uncertainty of this map is approximately 0 to 6.62 A/m, inferred from the inverse error level of 1%.

366



Figure 5: Radial Magnetic Flux on the CMB Estimated Using the Swarm Dataset. The continent
 shapes provide orientation.

370 **4. Discussion**

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371 Although our proposed method aims to invert an equivalent magnetization with accurate 372 location and distribution to improve data modeling accuracy, and the details of current 373 circulation are thoroughly analyzed in the experimental section, it is important to clarify that the 374 primary objective of this study is not to invert the currents in the core. This is because volume 375 magnetization sources, such as those shown in Fig. 1f, are often complex and challenging to 376 interpret. Therefore, it is necessary to transform these into more comprehensible forms, like the 377 current circulation model, which offers an enhanced perspective compared to the conventional ambient magnetization volume, thereby facilitating the interpretation of geomagnetic field data. 378

Estimating the depth of the current source within the Earth's core inherently involves biases, as evidenced by our synthetic experiment. Although the equivalent magnetization source was accurately placed in terms of depth, the resultant current source showed a depth bias, being approximately 600 km shallower. This bias may arise from three key factors: Firstly, observational data inherently contain limited information, often missing crucial aspects of the 384 actual magnetization vector source, particularly where opposing magnetization is negated by the 385 dominant dipole field. Secondly, the inversion process for a magnetization source is 386 characteristically non-unique, making the derived current source an approximated, and 387 potentially divergent, representation of the actual source. Thirdly, our inverse method is not 388 tailored for precise 3-D current modeling; it is designed to provide an auxiliary, equivalent 389 source model to enhance understanding of the geomagnetic field, rather than to replicate the 390 actual source depth accurately. These factors collectively underscore the challenges in accurately 391 determining the depth of current sources in geomagnetic field modeling. Analogous to findings 392 from our synthetic experiment, it is plausible that the depths of current sources in the IGRF and 393 Swarm examples are shallower than the actual sources. This suggests that the true depth of the 394 actual source model might be approximately 600 km deeper than our estimates. Therefore, this 395 depth bias, or inherent uncertainty, must be carefully considered and accounted for in the 396 interpretation of any dataset.

397 To assess the similarity in distribution between the estimated models and the actual fluid 398 dynamics within the core's interior, we reference results from previous dynamo research. 399 Initially, we compare the radial magnetic flux derived in our study with findings from dynamo 400 area studies. Jackson (2003) analyzed MAGSAT satellite data, identifying intense flux spots or 401 bands in the equatorial regions on the CMB, and presented a radial flux map with flux intensity 402 ranging approximately between [-900 μ T, 900 μ T]. Korte & Holme (2010) constructed the 403 time-averaged field structure at the CMB over four different time scales. When comparing our 404 results (as illustrated in Fig. 5) with the corresponding radial field maps from these studies-405 specifically the 2000 BC map in Jackson (2003) and the average field for 1840 AD to 1950 AD 406 in Korte & Holme (2010)—a similarity in the distribution of fields can be observed. All maps 407 exhibit intense flux patches primarily in the equatorial and polar regions. The numerical range of 408 our map, [-500 μ T, 500 μ T] (Fig. 5a), aligns more closely with the results of Korte & Holme 409 (2010) ([-400 μ T, 400 μ T]) than with those of Jackson (2003). Furthermore, Jackson (2003) 410 hypothesized that these intense flux spots might be associated with equatorial fluid structures 411 near the CMB. Our current model supports this hypothesis, as shown in Figs. 2d and 3d, where 412 an intense equatorial current is observed beneath the region, dominating the circulation pattern.

413 Secondly, Stump and Pollack (1998) estimated a 2-D spherical current sheet model of the 414 Earth's inner core by inverting spherical harmonic coefficients, based on Maxwell's equations. 415 Their analysis resulted in a 2-D map depicting current flow directions, forming a complete global 416 circulation pattern. The surface circulation pattern they presented shows remarkable similarity to 417 our 3-D current circulation model on the CMB surface, as depicted in Figs. 2d-f and 3d-f. Both 418 models feature a dominant equatorial gyre linked to a retrograde current beneath the South 419 Atlantic Ocean, two weaker current gyres beneath Siberia and Canada, and a singular gyre in the 420 Antarctic region. Moreover, the direction and intensity trends of these current flows are 421 strikingly consistent across both models.

422 Although the current models derived from the IGRF and Swarm datasets display 423 remarkable similarities in their circulation patterns, there are notable differences between them. 424 We attribute these discrepancies to two primary factors. First, the IGRF-13's maximum spherical 425 harmonic degree expansion is 13. While this represents a significant improvement in capturing 426 smaller-scale internal signals compared to previous versions, the direct measurements of the 427 geomagnetic field by the Swarm satellites encompass a broader range of scale signals and 428 inferences. This, we believe, is the principal factor contributing to the observed differences 429 between the models. Second, the datasets differ in their collection timeframes. The IGRF dataset 430 was compiled with a 'frozen' time parameter set to January 1, 2023, whereas the Swarm dataset 431 consists of data filtered from measurements taken during August 22-27, 2022, between 22:00 and 432 03:00 local time. From this perspective, it is reasonable to expect some variation in models 433 derived from the same inversion method when the observational datasets differ in such respects.

434 The dynamo theory posits that both toroidal and poloidal currents coexist within the core. 435 However, it is the toroidal current that contributes to the external magnetic field, while the 436 toroidal field generated by the poloidal current is zero at CMB. So the actual current 3-D pattern 437 in the core cannot be completely recovered due to the limitations of external and long-distance 438 observation data. This study focuses on inverting the magnetic satellite-observed external field 439 and developing equivalent magnetization and current sources to modelling and gain a 440 comprehensive understanding of the geomagnetic field. Nevertheless, it is essential to clarify that 441 exploring the actual current flow patterns and conducting dynamo simulation research are 442 outside the scope of this study. Our forthcoming research will be concentrated on interpreting the 443 secular variations in conjunction with a time serial dataset inverted.

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534 **Open Research**

The magnetic dataset in Section 3.1 is extracted from the IGRF model (Alken et. al, 2021). The satellite magnetic observations are obtained from the European Space Agency (ESA) (Olsen et al, 2013). MATLAB (The Mathworks, 2021) is employed for plotting the 2-D figures. The unstructured meshes used in this study are generated using Gmsh (Geuzaine & Remacle, 2009). Paraview (Ayachit & Utkarsh, 2015) is utilized for visualizing the 3-D models.

541	Supplementary Materials
542	Supplementary Text
543	Figs. S1 to S3
544	Dataset S1
545	Movie S1 to S3
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547	Acknowledgements
548	
549	Competing interests
550	The authors have no conflicts of interest to declare.
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552	Additional information
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