Simulating Postfire Debris Flow Runout Using Morphodynamic Models and Stochastic Surrogates

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Abstract

Fire affects soil and vegetation, which in turn can promote the initiation and growth of runoff-generated debris flows in steep watersheds. Postfire hazard assessments often focus on identifying the most likely watersheds to produce debris flows, quantifying rainfall intensity-duration thresholds for debris flow initiation, and estimating the volume of potential debris flows. This work seeks to expand on such analyses and forecast downstream debris flow runout and peak flow depth. Here, we report on a high fidelity computational framework that enables debris flow simulation over two watersheds and the downstream alluvial fan, although at significant computational cost. We also develop a Gaussian Process surrogate model, allowing for rapid prediction of simulator outputs for untested scenarios. We utilize this framework to explore model sensitivity to rainfall intensity and sediment availability as well as parameters associated with saturated hydraulic conductivity, hydraulic roughness, grain size, and sediment entrainment. Simulation results are most sensitive to peak rainfall intensity and hydraulic roughness. We further use this approach to examine variations in debris flow inundation patterns at different stages of postfire recovery. Sensitivity analysis indicates that constraints on temporal changes in hydraulic roughness, saturated hydraulic conductivity, and grain size following fire would be particularly beneficial for forecasting debris flow runout throughout the postfire recovery period. The emulator methodology presented here also provides a means to compute the probability of a debris flow inundating a specific downstream region, consequent to a forecast or design rainstorm. This workflow could be employed in scenario-based planning for postfire hazard mitigation.

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9 Key Points:

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We simulate postfire debris-flow runout by accelerating and building surrogates for a morphodynamic model of runoff and sediment transport. With this acceleration, we can rapidly explore how postfire recovery influences debris flow runout as a function of time since fire. Surrogate-based postfire hazard analyses offer rapid assessments of downstream debris flow inundation due to forecast or design rainstorms.

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16 Abstract

Fire affects soil and vegetation, which in turn can promote the initiation and growth of 17 runoff-generated debris flows in steep watersheds. Postfire hazard assessments often fo-18 cus on identifying the most likely watersheds to produce debris flows, quantifying rain-19 fall intensity-duration thresholds for debris flow initiation, and estimating the volume 20 of potential debris flows. This work seeks to expand on such analyses and forecast down-21 stream debris flow runout and peak flow depth. Here, we report on a high fidelity com-22 putational framework that enables debris flow simulation over two watersheds and the 23 downstream alluvial fan, although at significant computational cost. We also develop a 24 Gaussian Process surrogate model, allowing for rapid prediction of simulator outputs for 25 untested scenarios. We utilize this framework to explore model sensitivity to rainfall in-26 tensity and sediment availability as well as parameters associated with saturated hydraulic 27 conductivity, hydraulic roughness, grain size, and sediment entrainment. Simulation re-28 sults are most sensitive to peak rainfall intensity and hydraulic roughness. We further 29 use this approach to examine variations in debris flow inundation patterns at different 30 stages of postfire recovery. Sensitivity analysis indicates that constraints on temporal 31 changes in hydraulic roughness, saturated hydraulic conductivity, and grain size follow-32 ing fire would be particularly beneficial for forecasting debris flow runout throughout the 33 postfire recovery period. The emulator methodology presented here also provides a means 34 to compute the probability of a debris flow inundating a specific downstream region, con-35 sequent to a forecast or design rainstorm. This workflow could be employed in scenario-36 based planning for postfire hazard mitigation. 37

³⁸ Plain Language Summary

Fire on steep hillslopes increases the potential for debris flows, or rapidly moving 39 mixtures of water, soil, ash, and rock, that can develop during intense rainfall. Debris 40 flows threaten communities situated downstream of steep, burned areas. Burn severity 41 and hillslope steepness give some indication of the potential for debris flow initiation in 42 response to a particular rainstorm, but these factors alone do not indicate which down-43 stream areas might be inundated by a flow. To investigate areas impacted by potential 44 debris flows, we utilize a computational model that represents the physical processes of 45 debris flow initiation and runout. Such process-based models are computationally inten-46 sive, which has limited their use in rapid hazard assessments. Here, we implement a de-47

⁴⁸ bris flow initiation and runout model in a high-performance computing framework. Even
⁴⁹ so, a typical debris flow simulation takes several hours to complete on a supercomputer.
⁵⁰ Thus, we also build statistical models of these physical models, so-called emulators. Emulators can rapidly approximate the debris flow simulation and thus offer a mechanism
⁵² to investigate outcomes of debris flow models over a wide range of scenarios. We investigate how debris flow inundation footprints vary with rainfall intensity and the impact
⁵⁴ of landscape recovery on debris flows.

55 1 Introduction

Fire alters soil and vegetation, leading to increases in runoff and erosion (Shakesby 56 & Doerr, 2006; Moody et al., 2013). In extreme cases, particularly when steep water-57 sheds burn at moderate or high severity, rapid entrainment of sediment into runoff can 58 produce debris flows (Kean et al., 2011; Gabet & Bookter, 2008; Esposito et al., 2023; 59 Conedera et al., 2003; Nyman et al., 2011; Diakakis et al., 2023). Postfire debris flows 60 generated by runoff are most common in the first year following fire, when fire-driven 61 reductions in soil infiltration capacity, rainfall interception, and hydraulic roughness are 62 most extreme (DeGraff et al., 2015; Hoch et al., 2021; Thomas et al., 2021; Esposito et 63 al., 2023; Graber et al., 2023). Due to the complex interactions among runoff, sediment 64 transport, and debris-flow initiation and runout following fire, mathematical models that 65 couple these processes have the potential to inform our understanding of postfire debris-66 flow hazards and how they change through time as landscapes recover (McGuire et al., 67 2021). Yet the exploration of postfire debris flow processes through application of mor-68 phodynamic models for runoff and sediment transport is often limited by the high di-69 mensionality, poor constraints on parameters, and substantial computation time of the 70 models. Quantification of uncertainties associated with sources and model parameters, 71 and incorporation of those uncertainties in probabilistic predictions of hazard, often re-72 quire use of simulation ensembles with hundreds or thousands of members, greatly in-73 creasing computational costs (Bayarri et al., 2015). In this work, we accelerate a recently 74 developed morphodynamic model of runoff and sediment transport (McGuire et al., 2017), 75 and pair model runs with stochastic surrogates for high-dimensional output (Gu & Berger, 76 2016) as a strategy for simulating postfire debris-flow runout. This acceleration also en-77 ables us to rapidly explore how temporal changes in fire-affected model input parame-78 ters influence debris flow runout. 79

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The processes leading to the initiation and growth of runoff-generated debris flows 80 following fire involve the generation of spatially distributed overland flow and its sub-81 sequent interaction with sediment on hillslopes and in channels (Santi et al., 2008; Sta-82 ley et al., 2014; McGuire et al., 2017; Guilinger et al., 2020). This presents a contrast 83 to debris flows that mobilize from shallow landslides, which initiate when infiltration pro-84 motes increases in pore-water pressure that causes a discrete mass of soil to become un-85 stable on a hillslope (Iverson et al., 1997). The source of sediment for postfire runoff-86 generated debris flows may come from a combination of processes, including widespread, 87 shallow erosion on hillslopes in response to raindrop-driven sediment transport and un-88 confined sheet flow, rill erosion on hillslopes in areas of concentrated flow, and channel 89 scour (Santi et al., 2008; Staley et al., 2014; McGuire et al., 2017; Tang et al., 2019). All 90 three processes are more efficient at eroding sediment following fire as a result of decreases 91 in ground cover and increases in runoff (Robichaud et al., 2016), particularly rill and chan-92 nel erosion processes where overland flow does the work to entrain and transport sed-93 iment (Sheridan et al., 2007; Wagenbrenner et al., 2010). In areas of unconfined, shal-94 low flow, raindrops facilitate sediment detachment and transport in combination with 95 runoff (Kinnell, 2005). Raindrop-driven sediment transport on hillslopes increases fol-96 lowing fire due to removal of the vegetation canopy, litter, and duff, that tend to shield 97 the soil surface from raindrop impact in unburned settings. 98

Models designed to simulate runoff-generated debris flows from initiation to depo-99 sition must therefore account for spatially distributed runoff and sediment transport as 100 well as changes in flow behavior resulting from spatial and temporal variations in sed-101 iment concentration. Fully developed debris flows are characterized by volumetric sed-102 iment concentrations in excess of 40-50%, though they initiate from runoff with initially 103 negligible sediment concentration. Due, in part, to the relatively flashy hydrologic re-104 sponse of watersheds burned by fire, postfire runoff-generated debris flows initiate in re-105 sponse to short-duration bursts of high intensity rainfall (Kean et al., 2011). Rainfall in-106 tensity averaged over a 15-minute time period, I_{15} , is correlated well with runoff mag-107 nitude at the outlet of small, recently burned watersheds and threshold values of I_{15} have 108 proven to be reasonable predictors for debris flow initiation in the western USA (Kean 109 et al., 2011; Staley et al., 2017). Rainstorms that contain multiple, distinct bursts of high 110 intensity rainfall, such as where I_{15} exceeds the debris flow threshold at multiple times 111 in a single event, may lead to multiple pulses of debris flow activity (Kean et al., 2011). 112

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One benefit of morphodynamic models that are capable of simulating the debris-flow life-113 cycle from initiation to runout is their ability to directly account for the impacts of tem-114 porally varying rainfall intensity on debris flow processes, including runout and inunda-115 tion extent. In contrast, models designed only to simulate debris flow runout processes 116 (i.e. neglecting runoff and sediment transport), can be employed by defining an inflow 117 hydrograph above the anticipated runout zone based on the rainfall hydrograph and an 118 estimated debris flow volume, or by allowing a pile of sediment and water to flow down-119 stream from a pre-defined initiation zone (Barnhart et al., 2021; Gorr et al., 2022; Gib-120 son et al., 2022). As a result, employing morphodynamics models to estimate debris flow 121 runout avoids introducing epistemic uncertainty associated with specifying a volume of 122 material associated with an inflow hydrograph. 123

McGuire et al. (2017) developed a model that accounts for infiltration, runoff, sed-124 iment transport, and changes to flow resistance driven by sediment concentration in or-125 der to simulate postfire debris flow initiation and growth. In this model, rainfall drives 126 sediment entrainment and transport processes that naturally lead to debris flow initi-127 ation when hydrogeomorphic conditions give rise to flows with sufficiently high sediment 128 concentrations. Since rainfall, runoff, and erosion processes are related to model param-129 eters known to change following fire, such as saturated hydraulic conductivity (Perkins 130 et al., 2022; Ebel et al., 2022; Thomas et al., 2021), hydraulic roughness (Stoof et al., 131 2015), and vegetation cover (Stoof et al., 2012), it is possible to use this framework to 132 explore how postfire recovery affects debris-flow initiation, growth, and runout. The model 133 proposed by McGuire et al. (2017) contains a number of parameters that are challeng-134 ing to constrain and is computationally intensive, especially when simulating debris flow 135 initiation and runout processes over large areas. Thus far, model applications have been 136 limited to examining debris-flow initiation processes in small headwater basins where prior 137 work and intensive field monitoring helped constrain parameter values (McGuire et al., 138 2021, 2017; Tang et al., 2019). Here, we employ adaptive mesh refinement and parallel 139 computations from the Titan framework (Patra et al., 2005; Dalbey et al., 2008) to make 140 model application more tractable over larger spatial domains, specifically with the goal 141 of simulating the entirety of the postfire debris-flow lifecycle from runoff generation and 142 debris flow growth to runout. In addition, we employ statistical emulators, which enable 143 fast approximations to solutions of the model equations, in order to explore the influ-144 ence of model parameters on debris-flow runout extent. 145

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Gaussian process emulators (GPs) are a powerful class of statistical surrogates that 146 enable rapid approximation and uncertainty quantification of computationally intensive 147 first principles (conservation laws) based models or, simulators (Currin et al., 1988; Sacks, 148 Welch, et al., 1989; Welch et al., 1992). In the context of postfire debris-flows, GPs of-149 fer a mechanism to quickly explore spatial patterns in peak flow depth and area inun-150 dated for various parameter settings. GPs also allow for uncertainty quantification via 151 Monte Carlo (MC) simulations of flow inundation. Along the way, GPs offer an approx-152 imate sensitivity analysis of physical parameters' effects on debris-flow inundation. Par-153 allel partial emulation (PPE) (Gu & Berger, 2016) extends the GP methodology to vector-154 valued outputs – in our case, flow depth at each map point (pixel). Some recent stud-155 ies use PPE for flow model sensitivity analysis and calibration (Zhao et al., 2021; Zhao 156 & Kowalski, 2022). In this work, we apply PPE to explore the effects of rainfall inten-157 sity and sediment availability as well as parameters associated with saturated hydraulic 158 conductivity, hydraulic roughness, grain size, and sediment entrainment on peak debris 159 flow depth and inundation extent. 160

¹⁶¹ 2 Study Area

The Thomas Fire ignited in December 2017 and burned more than 1100 km^2 , in-162 cluding a series of steep watersheds in the Santa Ynez Mountains above the community 163 of Montecito, USA (Figure 1a). On 9 January 2018, widespread rainfall developed over 164 the burned area in association with an atmospheric river (Oakley et al., 2018). A nar-165 row cold frontal rainband (NCFR), a relatively small-scale feature characterized by a band 166 of intense precipitation that forms along a cold front, moved over burned watersheds above 167 Montecito and produced a short-duration burst of intense rainfall. Peak 15-minute rain-168 fall intensities, I_{15} , in this area ranged from approximately 78 mm/h to 105 mm/h (Kean 169 et al., 2019). During this time period, rainfall intensities greatly exceeded the infiltra-170 tion capacity of the soil, leading to infiltration-excess overland flow that generated rills 171 on steep hillslopes (Alessio et al., 2021). The combination of intense hillslope erosion and 172 channel incision led to runoff-generated debris flows that traveled across the populated 173 alluvial fan (Kean et al., 2019; Morell et al., 2021; Alessio et al., 2021). The debris flows 174 that initiated in six watersheds above Montecito mobilized more than $630,000 \text{ m}^3$ of sed-175 iment and led to 23 fatalities, 408 damaged structures, and more than \$1 billion in dam-176 age (Kean et al., 2019; Lancaster et al., 2021). In this study, we focus on modeling de-177

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Figure 1. (a) The Thomas Fire ignited on 4 December 2017 and burned over 1140 km² in southern California, USA. (b) The western portion of the burned area included a series of steep watersheds (black rectangle) above the community of Montecito, located near Santa Barbara. (c) Our study focuses on two watersheds, Oak Creek (0.45 km²) and San Ysidro Creek (7.6 km²), located upstream of Montecito.

bris flow initiation, growth, and runout for two of these watersheds, Oak Creek and San 178 Ysidro Creek. These watersheds are well suited for our study since high resolution to-179 pographic and rainfall data are available and we can leverage data collection following 180 the fire (Kean et al., 2019) and previous work in the Transverse Ranges of southern Cal-181 ifornia (McGuire et al., 2016, 2021; Tang et al., 2019; McGuire et al., 2017) to constrain 182 model parameter ranges. In addition, the model domain, which includes the watersheds 183 and downstream alluvial fan, is large enough to make physics-based simulations com-184 putationally challenging (Figure 1c). 185

Roughly 85% of San Ysidro Creek, which has a total drainage area of 7.6 km², burned at moderate or high severity (Figure 1c). Approximately 49% of Oak Creek, which is substantially smaller at 0.45 km², burned at moderate or high severity. Both the San Ysidro

and Oak Creek watersheds are steep, with median slopes of 28 and 21 degrees, respec-

¹⁹⁰ tively. The debris flow that initiated in San Ysidro Creek mobilized a total volume of

 $_{191}$ 297,000 m³ and inundated an area of 905,000 m² while the smaller Oak Creek watershed

 $_{192}$ produced a debris flow that moved 10,000 m³ of sediment and inundated 102,000 m²

(Kean et al., 2019) (Figure 1c). The grain size distribution in the debris flows was bi-

¹⁹⁴ modal, consisting of a sandy matrix that suspended boulders with a large axis greater

than several meters (Kean et al., 2019).

¹⁹⁶ 3 Methodology

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3.1 Simulating runoff-generated debris flows with Titan2D

Steep watersheds recently burned by fire often experience greater amounts of runoff 198 and increased rates of sediment transport. Factors affecting rates of sediment transport, 199 and also the initiation and growth of runoff-generated debris flows, include rainfall in-200 tensity and duration, vegetation cover, soil infiltration capacity, and sediment charac-201 teristics (e.g. grain size, erodibility). The model developed by McGuire et al. (2017) rep-202 resents rainfall, infiltration, fluid flow, and sediment entrainment and deposition processes, 203 which makes it a useful framework for simulating runoff-generated debris flows in steep 204 terrain (McGuire et al., 2017; Tang et al., 2020). We provide a brief overview of the gov-205 erning equations, which we solve within the Titan2D framework (Patra et al., 2005; Simakov 206 et al., 2019). 207

The Titan2D code employs an adaptive mesh, finite volume scheme to solve hy-208 perbolic PDEs describing shallow-water like mass flows over digital elevation models of 209 real topography. Titan2D (Patra et al., 2005) was originally developed to solve the depth 210 averaged shallow-water mass flow equations by Savage and Hutter (1989). Titan was mod-211 ernized and restructured in 2019 (Simakov et al., 2019) to optimize storage and access 212 for parallel adaptive mesh refinement, and to facilitate the usage of new material mod-213 els. Using Titan2D to solve the model equations proposed by McGuire et al. (2017) there-214 fore offers several advantages, especially for simulating debris flows over spatial scales 215 of more than a few square kilometers. In particular, Titan2D uses partitioned hash-tables 216 for better memory allocation structures, allowing it to compute over a large domain em-217 ploying adaptive mesh refinement (AMR) and unrefinement, with computational efficiency. 218

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It is also well suited for parallel computing using MPI and OpenMP and dynamic load balancing for exploiting multiprocessor computing. The computational efficiencies attained make it tractable to simulate end-to-end flows from initiation to downstream inundation.

The equations representing the motion of fluid and sediment can be written as a set of depth averaged conservation laws. We have

$$\frac{\partial \mathbf{U}}{\partial t} + \frac{\partial \mathbf{F}}{\partial x} + \frac{\partial \mathbf{G}}{\partial y} = \mathbf{S_0} + \mathbf{S_1} + \mathbf{S_2},\tag{1}$$

226 where

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$$\mathbf{U} = \left\{ h \quad uh \quad vh \quad c_1h \quad \cdots \quad c_kh \right\}^T, \tag{2}$$

$$\mathbf{F} = \left\{ hu \quad hu^2 + \frac{1}{2}g_z h^2 \quad huv \quad huc_1 \quad \cdots \quad huc_k \right\}_{-}^{T}, \tag{3}$$

$$\mathbf{G} = \left\{ hv \quad huv \quad hv^2 + \frac{1}{2}g_z h^2 \quad hvc_1 \quad \cdots \quad hvc_k \right\}^T, \tag{4}$$

and where h, u, v, and c_i are flow depth, velocity along x-axis, velocity along y- axis, 230 and sediment concentration of particle size class i. Components of gravitational accel-231 eration in the x, y, and z directions are given by g_x , g_y , and g_z , respectively, and k de-232 notes the number of particle size classes. $\mathbf{S_0}, \mathbf{S_1}$ and $\mathbf{S_2}$ are source terms. $\mathbf{S_0}$ denotes 233 the contributions of mass sources and sinks associated with the effective rainfall rate, P_{eff} , 234 and the soil infiltration capacity, I, as well as momentum sources and sinks arising from 235 variations in topographic elevation, and spatial variations in sediment concentration and 236 debris flow resistance terms, S_x and S_y . Specifically, S_0 is given as 237

$$\mathbf{S_0} = \begin{cases} P_{eff} - I + \frac{\partial z}{\partial t} \\ -g_x h + \gamma_x - \psi S_x \\ -g_y h + \gamma_y - \psi S_y \\ 0 \\ \vdots \\ 0 \end{cases} \right\}.$$
(5)

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The debris flow resistance terms are scaled by ψ , which increases linearly from 0 to 1 as the volumetric sediment concentration increases from 0.2 to 0.4. This scaling factor gradually increases the importance of the debris flow resistance terms as volumetric sediment concentration approaches levels that are consistent with a transition from with flood flow to debris flow. The terms γ_x and γ_y account for the effects of spatially variable sediment

concentration and are given by

$$\gamma_x = \frac{-(\rho_s - \rho_w)g_z h^2}{2\rho_f} \frac{\partial c}{\partial x} \tag{6}$$

and

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$$\gamma_y = \frac{-(\rho_s - \rho_w)g_z h^2}{2\rho_f} \frac{\partial c}{\partial y} \tag{7}$$

Here, c denotes volumetric sediment concentration, $\rho_w = 1000 \text{ kg m}^{-3}$ the density of

water, $\rho_s = 2600 \text{ kg m}^{-3}$ the density of sediment, and $\rho_f = c\rho_f + (1-c)\rho_w$ the den-

 $_{241}$ sity of the flow. S_1 accounts for flow resistance using a depth-dependent Manning's for-

242 mulation, and is given as

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$$\mathbf{S_1} = \begin{cases} 0 \\ g_z \eta^2 h u \sqrt{h u^2 + h v^2} / h^{7/3} \\ g_z \eta^2 h v \sqrt{h u^2 + h v^2} / h^{7/3} \\ 0 \\ \vdots \\ 0 \\ 0 \\ \end{cases},$$
(8)

where η is the Manning friction coefficient. The friction coefficient varies with flow depth according to

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$$\eta = \begin{cases} \eta_0 (h/h_c)^{-\epsilon} & h \le h_c \\ \eta_0 & h > h_c \end{cases},$$
(9)

where η_0 is the hydraulic roughness coefficient, h_c is a critical flow depth and ϵ is a phenomenological exponent. Soil infiltration capacity, I, is represented by the Green-Ampt model where

$$I = k_s \frac{Z_f + h_f + h}{Z_f},\tag{10}$$

with k_s denoting saturated hydraulic conductivity, h_f the wetting front potential, $Z_f = V/(\theta_s - \theta_i)$ the depth of the wetting front, V the cumulative infiltrated depth, θ_s the volumetric soil moisture content at saturation, and θ_i the initial volumetric soil moisture content. The source term S_2 accounts for sediment entrainment and deposition processes, which are represented using the framework proposed (Hairsine & Rose, 1992a, ²⁵⁶ 1992b). In particular,

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$$\mathbf{S_2} = \begin{cases} 0 \\ 0 \\ e_1 + e_{r1} + r_1 + r_{r1} - d_1 \\ \vdots \\ e_k + e_{rk} + r_k + r_{rk} - d_k \end{cases},$$
(11)

where e_k and e_{rk} are sediment detachment and re-detachment rates due to raindrop im-258 pact for sediment particles in size class k, r_k and r_{rk} are rates of entrainment and re-259 entrainment due to runoff, and d_k is the effective deposition rate. The model differen-260 tiates between original soil, which has not yet been entrained and transported during the 261 modeled rainstorm, and deposited sediment, which has been detached and subsequently 262 deposited. Detachment rates for entraining original sediment and re-entraining deposited 263 sediment are computed differently. Sediment in the deposited layer may also fail en-masse 264 (McGuire et al., 2017). Rates of sediment entrainment and re-entrainment by runoff are 265 given by 266

$$r_k = (1 - H)p_k \frac{F(\Omega - \Omega_{cr})}{J},$$
(12)

268 and

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$$_{rk} = H \frac{m_k}{m_t} \frac{F(\Omega - \Omega_{cr})}{\frac{\rho_s - \rho_f}{\rho_s} gh}.$$
(13)

Here, m_k is the deposited sediment mass per unit area for sediment in size class k, m_t 270 is the total mass of deposited sediment per unit area, $H = \min(m_t/m_t^*, 1)$ accounts 271 for the degree to which deposited sediment shields the underlying bed from erosion, m_t^* 272 is the mass of deposited sediment needed to completely shield original sediment from ero-273 sion, ρ_f is the density of the flow, ρ_s is the density of sediment, F denotes the fraction 274 of stream power effective in sediment entrainment, $\Omega = \rho_f g S_f \sqrt{uh^2 + vh^2}$ is stream 275 power, and $S_f = \eta^2 (uh^2 + vh^2)h^{-10/3}$ is the friction slope. In this work, we consider 276 a single particle size class characterized by a representative particle diameter, δ . 277

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3.2 Rainfall and model parameters

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A digital elevation model (DEM) of the study area is input to the Titan2D simulation. Here, we use a 1 m DEM derived from post-event airborne lidar. Elevations and slopes at locations required by the computational mesh are obtained using a 9 point $(3 \times$ 3) finite difference stencil to interpolate on the DEM grid reducing the effects of artifacts and noise in the data (Patra et al., 2005). Errors in the DEM could be treated as
an uncertainty that is propagated through the simulation and subsequent analysis (Stefanescu,
Bursik, & Patra, 2012; Stefanescu, Bursik, Cordoba, et al., 2012), but we did not consider this in the present study.

Runoff and debris flows initiated in the study area in response to a short duration, high intensity burst of rainfall in the early morning hours of 9 January 2018 (Kean et al., 2019). All simulations use 1-minute rainfall intensity data derived from the KTYD rain gauge for a 20-minute time period that spans this short temporal window when rainfall intensity rapidly increased and debris flows initiated (Figure 2). The gauge is maintained by the Santa Barbara County Flood Control District and located approximately 5 km west of the San Ysidro Creek watershed.



Figure 2. Rainfall hyetograph derived from the KTYD rain gauge, located roughly 5 km west of the San Ysidro Creek watershed.

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Simulations were designed to explore the extent to which inundated area and peak 294 flow depths on the alluvial fan were influenced by rainfall intensity as well as several pa-295 rameters that may play critical roles in debris-flow initiation and growth. We explored 296 the impact of different rainfall intensities by multiplying the 1-minute rainfall intensity 297 time series by a rainfall intensity factor (RI_{fac}) that varied from 0.5 to 1.5. We also var-298 ied the representative particle diameter, δ , from 0.05–0.125 mm, the fraction of stream 299 power effective in entrainment, F, from 0.001 - 0.006, the hydraulic roughness coeffi-300 cient, n_0 , from 0.03 – 0.2, and saturated hydraulic conductivity, k_s , from 5 – 25 mm 301

h⁻¹. We further enforced a maximum soil thickness, r_{max} , that varied from 0.25–1.5 m to explore the role of sediment availability. All other parameters were fixed (Table 1, Table 2). We used a Latin hypercube sampling strategy to generate 64 different parameter sets from the ranges specified above (McKay et al., 1979). Figure 3 shows the maximum flow depth for three of the 64 training simulations chosen to demonstrate typical results with low, moderate, and long runout extents. Each of these simulations took several hours to complete on an HPC cluster using up to 16 cores on an intel Xeon Gold 6226R processor.



Figure 3. Maximum flow depth of three simulations resulting in (a) short runout (b) moderate runout and (c) long runout onto the alluvial fan. For contrast, the maximum colorbar limit is set to 2.5 m although maximum flow depth does exceed 2.5 m in some locations.

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We chose to focus on exploring the effects of rainfall intensity, δ , F, n_0 , k_s , and soil 310 thickness (sediment availability) since they control different aspects of the debris flow 311 initiation and growth process and, aside from rainfall intensity, they may all be strongly 312 affected by fire in our study area. Peak rainfall intensity over sub-hourly durations, par-313 ticularly the 15-minute duration, is correlated well with runoff in recently burned wa-314 tersheds in southern California (Kean et al., 2011). Peak 15-minute rainfall intensity is 315 also used in empirical models designed to predict postfire debris-flow likelihood and vol-316 ume in the western USA (Staley et al., 2017; Gartner et al., 2014). We therefore expect 317 that variations in rainfall intensity during the relatively short (< 0.5 hours) portion of 318 the rainstorm that we are modeling will influence debris flow processes. 319

We expect the representative grain size, δ , to be relatively small in areas of con-320 centrated flow immediately following fire in our study area given the propensity for post-321 fire dry ravel to transport hillslope sediment to channels and valley bottoms (Florsheim 322 et al., 1991; Lamb et al., 2011). Both δ and the amount of sediment available for trans-323 port, which we vary by enforcing a maximum soil thickness (r_{max}) throughout the model 324 domain, may vary as a function of time since fire as sediment is exported from postfire 325 rainstorms (Tang et al., 2019). Similarly, Liu et al. (2021) found that k_s and the Man-326 ning coefficient were lowest during rainstorms in the first year following a high severity 327 fire in the San Gabriel Mountains, southern California, and increased by factors of roughly 328 3-4 over the next 4 years. Immediately following fire in southern California, values for 329 the Manning coefficient and saturated hydraulic conductivity may be as low as 0.025-330 $0.07 \text{ sm}^{-1/3}$ and $1-6 \text{ mm h}^{-1}$, respectively (Rengers et al., 2016; Tang et al., 2019; 331 Liu et al., 2021). Kean et al. (2019) used post-event, point scale measurements with a 332 tension infiltrometer to estimate the geometric mean of saturated hydraulic conductiv-333 ity at 20 mm h^{-1} in the days following the Montecito debris flows. The effective frac-334 tion of stream power, F, may be expected to increase immediately following fire due to 335 reductions in roughness associated with ground cover and vegetation. Past studies sug-336 gest values of $F \approx 0.005$ perform reasonably well in steep, recently burned watersheds 337 in southern California (McGuire et al., 2017; Tang et al., 2019). 338

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3.3 Emulating debris flows

Statistical emulators are effectively probabilistic models of computationally inten-340 sive physical model systems or *simulators*. That is, statistical emulators relate a set of 341 user-defined inputs, often physical parameter specifications, to simulator output. Gaus-342 sian process emulators (GPs) are a popular class of surrogates for approximating and 343 quantifying uncertainties in simulators as they (almost) interpolate computer model out-344 put (Sacks, Schiller, & Welch, 1989; Sacks, Welch, et al., 1989; Santner et al., 2003; Ras-345 mussen & Williams, 2006). Further, the variance of the associated GP offers a quick mech-346 anism to assess the uncertainty of using the emulator in place of the simulator for model 347 prediction at untested inputs. Thus GP emulators offer a rapid and quantifiable mech-348 anism to approximate output from physical process models that are computationally in-349 tensive to exercise. The parallel partial emulator (PPE) (Gu & Berger, 2016) extends 350 this surrogate model to vector valued output. 351

Inputs to GP emulators are user defined. They are typically influential parameters, which show up within the governing dynamics, the forcing terms, or boundary conditions, as opposed to independent variables in the physical model. For the model described in Sec 3.1, we choose p = 6 parameters to define our input vector, namely those described in Sec 3.2 and given by $\mathbf{q} = [k_s, r_{max}, \eta_0, F, \delta, RI_{fac}]$. The rainfall intensity factor is a scaling of the true rainfall time series that triggered the debris flows in our study area.

In initial explorations, the ratio of pore fluid pressure to total basal normal stress, 359 λ , was also considered, but found to not substantially influence spatial patterns in peak 360 flow depth and runout extent on the fan. We will discuss the relationship between GP 361 emulators and sensitivity analysis further in Sec 4. The output under consideration, \mathbf{y} , 362 is the maximum (over time) flow depth at each of s = 1.4M map points. The main ob-363 jective of the emulator is to predict the output of the Titan2D model at an untested sce-364 nario, \mathbf{q}^* , given a relatively modest set of N training or design runs \mathbf{q}^{D} and each of their 365 corresponding inundation depth outputs, \mathbf{y}_j^{D} , $j = 1, \dots, N$. In this work, we take N =366 64 training runs and each output vector, $\mathbf{y}_{i}^{\mathrm{D}}$, is a 1.4M element vector recording the peak 367 flow depth at each map point. Collecting these outputs together, we have $Y^{\rm D}$, as a 64× 368 1.4M matrix of training run outputs. The 64 training run inputs are chosen by a Latin 369 hypercube design (McKay et al., 1979; Santner et al., 2003) covering the ranges of in-370 puts listed in Table 1. All other parameters are fixed (Table 2). To fit the emulator, these 371 parameter ranges are normalized to a unit hypercube. 372

Given the training data, $\{\mathbf{q}^{\mathrm{D}}, Y^{\mathrm{D}}\}$, to approximate the inundation resulting from an untested scenario, \mathbf{q}^* , we use the predictive mean of the PPE given by

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 $\tilde{\mathbf{y}}(\mathbf{q}^*) = \mathbf{h}^{\mathrm{T}}(\mathbf{q}^*)B + \mathbf{r}^{\mathrm{T}}(\mathbf{q}^*)R^{-1}(Y^{\mathrm{D}} - H^{\mathrm{D}}B),$ (14)

where R is an $N \times N$ (64×64 in this work) matrix of correlations between pairs of de-376 sign inputs, $\mathbf{r}(\mathbf{q}^*)$ is an $N \times 1$ vector of correlations between the untested input, \mathbf{q}^* , and 377 each of the input scenarios in the design, \mathbf{q}^{D} . Further, $\mathbf{h}(\mathbf{q})$ is a $l \times 1$ vector of regres-378 sion variables, often taken to be constant or linear in \mathbf{q} (i.e., l = 1 for constant case used 379 in this work and l = p + 1 for the linear case), and H^{D} is and $N \times l$ matrix where the 380 j^{th} row are the regression variables evaluated at the j^{th} design point, $\mathbf{h}^{\mathrm{T}}(\mathbf{q}_{j}^{\mathrm{D}})$. The ma-381 trix B is a $l \times s$ matrix of regression coefficients. Here, each of the s = 1.4M outputs 382 has its own set of regression coefficients, but a shared correlation structure. We use a 383

Model parameter	Min value	Max value	Range parameter
k_s : Saturated hydraulic conductivity (mm	5	20	3.5
h^{-1})			
r_{max} : Maximum soil thickness (m)	0.25	1.5	9.9
η_0 : hydraulic roughness coefficient (s m ^{-1/3})	0.03	0.2	0.15
F: Fraction of stream power effective in sedi-	0.6×10^{-3}	1.0×10^{-3}	9.9
ment detachment			
δ : Effective grain size (mm)	0.05	0.125	5.8
$RI_{\rm fac}$: Rainfall intensity factor	0.5	1.5	0.38

Table 1. Parameter ranges with units and GP range parameters (unit-less) for the six parameters that varied among the N=64 debris flow simulations.

Matérn 5/2 correlation function (Stein, 1999). For two scenarios, e.g., two input points 384 $\mathbf{q}_i = (x_{i1}, \ldots, x_{ip})^{\mathrm{T}}$ and $\mathbf{q}_j = (x_{j1}, \ldots, x_{jp})^{\mathrm{T}}$, the standardized distance and correla-385 tion between these input scenarios are given by 386

$$d_{k} = \left(\frac{|x_{ik} - x_{jk}|^{2}}{\theta_{k}^{2}}\right)^{1/2}, \quad k = 1, \dots, p$$

$$c(\mathbf{q}_{i}, \mathbf{q}_{j}) = \prod_{k=1}^{p} \left(1 + \sqrt{5}d_{k} + \frac{5}{3}d_{k}^{2}\right) \exp\left(-\sqrt{5}d_{k}\right), \quad (15)$$

389

respectively. The predictive variance for each output dimension (pixel) of the PPE is given 390 by 301

³⁹²
$$\mathbf{v}_{j}(\mathbf{q}^{*}) = \sigma_{j}^{2} \left(1 - \mathbf{r}^{\mathrm{T}}(\mathbf{q}^{*})R^{-1}\mathbf{r}(\mathbf{q}^{*}) + \left(\mathbf{h}(\mathbf{q}^{*}) - (H^{\mathrm{D}})^{\mathrm{T}}R^{-1}\mathbf{r}(\mathbf{q}^{*})\right)^{\mathrm{T}} \right)^{\mathrm{T}}$$

$$\times \left((H^{\mathrm{D}})^{\mathrm{T}} R^{-1} H^{\mathrm{D}} \right)^{-1} \left(\mathbf{h}(\mathbf{q}^{*}) - (H^{\mathrm{D}})^{\mathrm{T}} R^{-1} \mathbf{r}(\mathbf{q}^{*}) \right) \right), \tag{16}$$

where σ_j^2 , (j = 1, ..., s) is the scalar variance corresponding to each pixel's output. "Fit-395 ting" a PPE amounts to estimating the regression parameters in B, the scalar variances 396 at of each output, σ_j^2 , and the range parameters $\{\theta_k : k = 1, \ldots, p\}$. To do so, we use 397 the RobustGaSP package (Gu et al., 2018, 2019). On a laptop, fitting a PPE to 1.4M398 pixels of output with N = 64 training runs takes roughly 10 minutes. 399

GP emulators have been applied to Titan2D-based volcanic debris flows (Bayarri 400 et al., 2009, 2015; Spiller et al., 2014; Rutarindwa et al., 2019) and recently to other Titan2D-401 based debris flows (Zhao et al., 2021; Zhao & Kowalski, 2022). In each of these studies, 402

Symbol	Definition	Value	Unit
a_0	Detachability of original soil	1000	$\rm kg\ m^{-2}\ s^{-1}$
a_{d0}	Detachability of deposited sediment	2000	$\rm kg\ m^{-2}\ s^{-1}$
m_{t0}^*	Deposited sediment needed to shield original soil	2.7	${\rm kg}~{\rm m}^{-2}$
J	Specific energy of entrainment	15.125	$\mathrm{m}^2~\mathrm{s}^{-2}$
C	Effective cohesion	200	Pa
ϕ_{bed}	Basal friction angle	32	\deg
λ	Ratio of pore fluid pressure to total normal stress	0.8	-
C_v	Fraction vegetation cover	0	-
h_f	Wetting front poential	1	mm
$ heta_i$	Initial volumetric soil moisture	0.1	-
θ_s	Volumetric soil moisture at saturation	0.39	-
ϵ	Exponent in friction model	0.33	-
h_c	Critical depth in friction model	3	mm

Table 2. Model parameters using the same notation as (McGuire et al., 2017).

source terms (particularly debris mass or flux) were specified via ad-hoc parameteriza tions which are less appropriate for postfire, runoff-generated debris flows.

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3.4 Numerical Experiments

Evaluation of the GP emulator's mean quickly allows one to explore any output quantity of interest over the parameter space. Here we take the output quantity of interest, \mathbf{y} , to be the maximum debris-flow depth at all locations. Additionally, the variance of the GP emulator accounts for the uncertainty introduced by evaluating the GP mean, $\tilde{\mathbf{y}}$, instead of the debris-flow process model. We can break our exploration of numerical experiments into three groups.

First, we perform leave-one-out experiments as a test of the PPE performance. This experiment amounts to excluding one simulation at a time, fitting a GP to the 63 remaining simulations, and then comparing the GP predicted inundation of the left-out scenario to actual simulated inundation for that scenario. This is repeated for each of the N =64 simulations. Second, we explore the relative importance of different model parameters using the GP's range parameters. The range parameters are positive numbers indicating the influence of each model parameter on the model response – the smaller the range parameter, the more influence the corresponding model parameter has on the debris flow model (i.e. maximum flow depth). As such, these range parameters act as an effective sensitivity analysis.

Lastly, we employ the emulator to isolate and explore how flow extent/depths are 423 driven by (1) changes in rainfall intensity and (2) changes in saturated hydraulic con-424 ductivity and hydraulic roughness that occur as the landscape recovers. We focus, in par-425 ticular, on exploring the effects of postfire changes in saturated hydraulic conductivity 426 and hydraulic roughness since Liu et al. (2021) provide guidance for parameterizing these 427 effects in the nearby San Gabriel Mountains. Since the GP emulator enables rapid for-428 ward uncertainty quantification, we demonstrate how it can be used to accelerate a Monte 429 Carlo probability of inundation calculation for two cases, namely when the observed storm 430 occurs 2 months and 14 months postfire. 431

432 4 Results

433

4.1 Emulator performance

To test the performance of the GP emulator for approximating the Titan2D sim-434 ulations, we examine leave-one-out predictions. Of note, the range parameter estimates 435 were very stable. We found the coefficients of variation for each of the six values to be 436 between 0.01-0.10 indicating that the relative influence of any input to the GP was not 437 swayed strongly by any single flow simulation. For illustrative purposes, we focus on two 438 points of interest – one located in a channel and one on the adjacent fan surface where 439 flow is relatively unconfined. In Figure 4 for both cases we sort the simulations by their 440 left-out flow depths, \mathbf{y} , and predict each with a credible interval centered at the mean 441 of the GP, \tilde{y} . We find root mean squared errors from these leave-one-out experiments 442 of 0.12m and 0.17m for the locations on the fan and on the channel, respectively. Fur-443 ther, we see that 89% (fan location) and 94% (channel location) of the simulated depths 444 fall within their predictive credible intervals. These numbers are slightly below the an-445 ticipated 95%, but this is likely due to the relatively small training set, and in the case 446

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of the fan location, the fact that roughly half of the simulations resulted in no inundation, which is challenging for GP emulation (Spiller et al., 2023).



Figure 4. Leave-one-out experiments for (a) a location on the fan and (b) a location in the channel (see Figure 7 for details). In each panel, indices are sorted based on the simulated max flow depth (red stars). GP predictive means for these scenarios are plotted in black while the 95% credible intervals are plotted as vertical blue bars.

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4.2 Sensitivity analysis

A crucial step to fitting a GP is estimating the range parameters. Smaller range parameters indicate that the corresponding model parameter has more influence on the debris flow model output of interest (i.e. maximum flow depth). From table 1, one can see that the debris flow model is most sensitive to rainfall intensity and the hydraulic roughness coefficient; it is moderately sensitive to saturated hydraulic conductivity and effective grain size; and is relatively insensitive to the maximum soil thickness and to the fraction of stream power effective in sediment detachment.

Scaling the rainfall intensity time series has a substantial effect on inundation extent (Figure 5). As the rainfall leading to the flows after the Thomas fire were quite intense, it is not surprising to see significant runout even when $RI_{fac} = 0.5$, though the extent of inundation is diminished relative to cases with more intense rainfall, namely $RI_{fac} = 1$ and $RI_{fac} = 1.5$. More intense rainfall leads to both increased water runoff and sediment entrainment, leading to greater flow volumes and increases in peak flow depth and area inundated.



Figure 5. Maximum flow depth for 50% ($RI_{fac}=0.5$), 100% ($RI_{fac}=1$), and 150% ($RI_{fac}=1.5$) scaling of the rainfall time series. Other parameters are set at their nominal value, except for hydraulic roughness and saturated hydraulic conductivity which were set to the minimal values in table 1, consistent with values anticipated immediately after the fire. For contrast, the maximum colorbar limit is set to 2.5 m although in the channels towards the north, maximum flow depth exceeds 2.5m.

4.3 Effects of postfire recovery on runout

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Liu et al. (2021) developed parametric best-fit curves to model the change in sat-465 urated hydraulic conductivity and hydraulic roughness as a function of time following 466 fire in the nearby San Gabriel Mountains. Using these relationships, and setting others 467 to the center of their respective ranges, we use the GP emulator to explore the effects 468 of temporal changes in the hydraulic roughness coefficient and saturated hydraulic con-469 ductivity. Both peak flow depth and area inundated in response to the observed rain-470 storm would decrease substantially over the first six months following fire (Figure 6). For 471 example, US Highway 101, which runs perpendicular to the direction of flow near the 472 distal portion of the fan, would only be inundated when the rainstorm occurs within the 473 first 3 months following fire. If the observed rainstorm were to have occurred 12 months 474 following the fire, the simulated inundation area would be limited to channels near the 475 fan apex. 476

We can also explore the effects of rainfall intensity and temporal changes in hydraulic roughness and saturated hydraulic conductivity following fire by examining flow depth

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Figure 6. Maximum inundation for several values of the Manning coefficient and saturated hydraulic conductivity along with the corresponding time from Figures 7b and 7c of Liu et al. (2021), respectively. All of the other parameters are set to the center value of their range. For contrast, the max colorbar limit is set to 2.5 m although in the channels towards the north, max flow depth exceeds 2.5m.

at distinct points of interest. Again, we consider the same two points for illustrative pur-479 poses, one located in a channel and one on the adjacent fan surface where flow is rela-480 tively unconfined (Figure 7). For a given time since fire, peak flow depths are greater 481 in the channel relative to on the fan surface, as expected. Peak flow depth decreases grad-482 ually over the first several months at the point on the fan before dropping to near zero 483 after approximately six months. Peak flow depths decrease over the first year following 484 fire in the channel location from roughly 2.5 m to 1.5 m. Visualizing peak flow depths 485 as a function of time since fire and rainfall intensity can be helpful for assessing tempo-486 ral shifts in the magnitude of rainfall associated with potential debris-flow impacts at 487 different locations. For example, even a rainstorm characterized by $RI_{fac} = 1.5$ would 488 not result in peak flow depths greater than 20 cm after approximately 0.6 years follow-489 ing fire. 490

We further use the emulator to produce probabilistic maps of inundation at different times following fire (Figure 8). Differences in the spatial patterns of inundation



Figure 7. Exploration at one location on the fan (top row) and one location in the channel (bottom row). Panels in column (a) indicate locations of all emulated flow depths (black) and those being explored in detail (red). Panels in column (b) show peak flow depth as a function of time since fire. The hydraulic roughness coefficient and saturated hydraulic conductivity are parameterized as a function of time since (Liu et al., 2021) while all other parameters set at their central values. (Note that the vertical scales are different; the maximum flow depth on the fan is roughly 1.0 m, and that in the channel is roughly 2.5 m.) Panels in column (c) show peak flow depth versus rainfall intensity with the hydraulic roughness coefficient and saturated hydraulic conductivity set to their respectively minimum values (i.e. as would be expected prior to any recovery) and with all other parameters set at their central values. Panels in column (d) contain color maps for maximum flow depth at these two locations varying all combinations of rainfall intensity and time (via Manning coefficient and saturated hydraulic conductivity). The white contours indicate the values of time and rainfall intensity leading to an inundation of 10 cm or more.

- ⁴⁹³ likelihood are apparent between scenarios where the storm occurs 2 months following the
- fire versus 14 months. We identify a location as being inundated if peak flow depth ex-
- $_{495}$ ceeds > 0.1m. All parameters were set to their central values except for saturated hy-
- ⁴⁹⁶ draulic conductivity and the hydraulic roughness coefficient. The latter parameter is sam-
- ⁴⁹⁷ pled from the distribution suggested by Liu et al. (2021) while the former is set to 2 and
- ⁴⁹⁸ 14 month values, respectively, estimated from the same study. The probability MC cal-

culation was carried out with 100 samples.



Figure 8. Probability of inundation maps, where a location is considered inundated if the maximum flow depth exceeds 10 cm. To calculate this probability, all parameters except the Manning coefficient and saturated hydraulic conductivity are set to their central values with the latter set to values corresponding to the 2 month and 14 month estimates from Liu et al. (2021), respectively.

499

500 5 Discussion

Fire impacts on soil and vegetation properties that affect the initiation and growth of runoff-generated debris flows are most extreme in the first few months following fire (DeGraff et al., 2015; Thomas et al., 2021). Potentially rapid changes in hydrologic conditions following fire limit the time window for gathering data needed to constrain pa-

rameters for postfire runoff and erosion models, including the model used here. Aside 505 from rainfall intensity, which will not be affected by the fire, we found that hydraulic rough-506 ness, the representative grain size, and saturated hydraulic conductivity played the most 507 important roles in controlling debris flow inundation. Additional model testing across 508 fire-prone regions in different geologic and climate settings is needed to assess model per-509 formance and determine the extent to which results related to parameter sensitivity are 510 generalizable. Nonetheless, this result provides observational targets that can help fo-511 cus future efforts to collect perishable postfire data. 512

We hypothesize that hydraulic roughness plays an important role in controlling in-513 undated area and peak flow depths because of its influence on both modeled sediment 514 detachment rates and flow resistance. Saturated hydraulic conductivity will influence the 515 rate at which sediment is detached by overland flow since it exerts a strong control on 516 the magnitude of infiltration-excess overland flow that often dominates in postfire set-517 tings. Increased rates of sediment detachment lead to increases in flow volume, which 518 in turn acts to increase runout and inundation potential (Barnhart et al., 2021). Grain 519 size similarly influences flow volume since a larger grain size will encourage more rapid 520 deposition of sediment. 521

Our evaluation of parameter sensitivity indicates that constraints on postfire val-522 ues for hydraulic roughness, saturated hydraulic conductivity, and the grain size distri-523 bution of sediment entrained in debris flows would be particularly beneficial for improv-524 ing estimates of debris flow runout estimates. Burn severity is likely to play a substan-525 tial role in a fire's effect on these variables (Moody et al., 2015; McGuire & Youberg, 2020). 526 In addition, attempts to capture changes in debris flow runout as a function of time since 527 fire would benefit from methods to parameterize temporal changes in hydraulic rough-528 ness, saturated hydraulic conductivity, and the grain size distribution of sediment en-529 trained in debris flows. Fire-driven reductions in hydraulic roughness are commonly cited 530 as a cause for increased runoff and erosion (McGuire & Youberg, 2020; Stoof et al., 2015), 531 but there are few constraints on the temporal changes in hydraulic roughness following 532 fire, which may be facilitated by changes in vegetation cover and/or grain roughness. Par-533 ticularly in southern California (Doehring, 1968; Florsheim et al., 1991; DiBiase & Lamb, 534 2020) and other tectonically active regions in the western USA (Roering & Gerber, 2005), 535 fire can promote substantial increases in dry ravel activity on hillslopes that may reduce 536 hydraulic roughness by increasing the availability of fine sediment in channels. Hydraulic 537

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roughness may then increase over time as dry ravel deposits are progressively eroded dur-538 ing postfire rainstorms (Tang et al., 2019). Temporal changes in debris flow sediment 539 source locations (Guilinger et al., 2020) and coarsening of particle size distributions due 540 to preferential erosion of fines would also influence the effective grain size in the model. 541 In practice, it is not clear how to quantitatively connect this single grain size parame-542 ter to the particle size distribution of hillslope or channel sediment, especially when flows 543 contain boulders. Postfire changes in saturated hydraulic conductivity can be inferred 544 from calibration of hydrologic models (Liu et al., 2021), rainfall simulator experiments 545 at the small plot scale (Robichaud et al., 2016), and point scale measurements (Ebel, 2020; 546 Ebel et al., 2022; Perkins et al., 2022). While some general patterns have been observed 547 between time since fire and values of saturated hydraulic conductivity, there is substan-548 tial site-to-site variability (Ebel & Martin, 2017). The level of uncertainty in influential 549 model input parameters and how they change over time highlights the need for proba-550 bilistic assessments of debris flow runout, which emulators can help to achieve by facil-551 itating rapid exploration of large parameter spaces. 552

Rainfall is a necessary driver for debris flow initiation and the model was also sen-553 sitive to rainfall intensity, specifically a rainfall intensity factor which we used to scale 554 the rainfall intensity time series. This finding is consistent with observations that post-555 fire basin-scale sediment yields (Pak & Lee, 2008) and debris flow volume (Gartner et 556 al., 2014) increase with rainfall intensity averaged over durations of 60 minutes or less. 557 Short duration (sub-hourly) bursts of high intensity rainfall are effective at generating 558 infiltration-excess overland flow that can trigger debris flows in recently burned steep-559 lands (Kean et al., 2011; Nyman et al., 2011; Esposito et al., 2023). Emulators may be 560 useful for generating probabilistic maps of debris flow inundation in response to design 561 storms with different rainfall intensities or examining changes at particular points of in-562 terest. In cases where there are specific values at risk downstream of a burned area, rapid 563 exploration of debris flow characteristics (i.e. peak flow depth) as a function of rainfall 564 intensity could help define impact-based rainfall thresholds that could be used for plan-565 ning and warning purposes. In other words, one could take advantage of the emulator's 566 computational efficiency to determine, not only the rainfall intensity required to initi-567 ate a debris flow but also the rainfall intensity required to produce a debris flow that would 568 impact a prescribed area of interest with some prescribed depth of flow. 569

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The computational cost of many physically-based debris flow models is a limita-570 tion in applications that are time sensitive, such as rapid postfire hazard assessments. 571 Postfire debris flows in the western USA, such as those that occurred near Montecito, 572 may occur before the fire has been officially contained and within weeks or months of 573 fire ignition. The emulator methodology presented here provides one avenue for mini-574 mizing computation times, since an initial suite of simulations can be used to train the 575 emulator which can later be applied with substantially less computational effort to gen-576 erate a probabilistic hazard map for a specific scenario. An emulator may even be trained 577 prior to a fire. Analogous approaches have been employed in related applications (Rutarindwa 578 et al., 2019; Spiller et al., 2020). Within the context of postfire hazards, an emulator could 579 be used to assess debris-flow runout and inundation downstream of a burned area in re-580 sponse to a design or forecast rainstorm. Atmospheric model ensembles, for example, can 581 provide estimates of peak 15-minute rainfall intensity over watersheds of interest that 582 could be used to constrain a distribution of rainfall intensity factors (Oakley et al., 2023). 583

584 6 Conclusions

We develop a physics based high fidelity computationally expensive morphodynamic 585 model and cost effective surrogates based on Gaussian process models of postfire debris 586 flows. We employ the Gaussian Process surrogate model, or emulator, to approximate 587 peak debris flow depth from a physics-based morphodynamic model, Titan2D. The em-588 ulator is able to approximate the peak flow depth with a mean squared error that is gen-589 erally in the range of 0.1-0.2 m when using a modest training data set built from 64 590 Titan2D simulations. The range parameters associated with the emulator provide a met-591 ric for the relative importance of input parameters, which provides guidance for those 592 that are most important to constrain for forward modeling of debris flow runout. We find 593 that peak flow depths are most sensitive to changes in hydraulic roughness and a rain-594 fall intensity factor and are moderately sensitive to saturated hydraulic conductivity and 595 effective grain size. We highlight the emulator's ability to provide rapid estimates of peak 596 flow depth for parameter combinations that were not part of the training data set by gen-597 erating probabilistic maps of inundation as a function of time since fire. Inundation like-598 lihood changes substantially over the first year following the fire, driven by temporal vari-599 ations in hydraulic roughness and saturated hydraulic conductivity. Emulator-based anal-600 yses can facilitate rapid Monte Carlo calculations of inundation probability, making them 601

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a promising option for rapid postfire hazard assessments and scenario planning beforea fire starts.

⁶⁰⁴ Data Availability Statement

The debris flow model under consideration in this paper is from (McGuire et al., 2017) and it is accelerated by implementation in the Titan2D platform (Patra et al., 2005; Simakov et al., 2019). Parametric models of the Manning coefficient and saturated hydraulic conductivity versus time are available from (Liu et al., 2021) as are validated samples of those same parameters for debris flows 2 and 14 months post fire. Packages to implement the parallel partial emulator (Gu & Berger, 2016) are available in (Gu et al., 2019).

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Simulating Postfire Debris Flow Runout Using Morphodynamic Models and Stochastic Surrogates

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9 Key Points:

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We simulate postfire debris-flow runout by accelerating and building surrogates for a morphodynamic model of runoff and sediment transport. With this acceleration, we can rapidly explore how postfire recovery influences debris flow runout as a function of time since fire. Surrogate-based postfire hazard analyses offer rapid assessments of downstream debris flow inundation due to forecast or design rainstorms.

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16 Abstract

Fire affects soil and vegetation, which in turn can promote the initiation and growth of 17 runoff-generated debris flows in steep watersheds. Postfire hazard assessments often fo-18 cus on identifying the most likely watersheds to produce debris flows, quantifying rain-19 fall intensity-duration thresholds for debris flow initiation, and estimating the volume 20 of potential debris flows. This work seeks to expand on such analyses and forecast down-21 stream debris flow runout and peak flow depth. Here, we report on a high fidelity com-22 putational framework that enables debris flow simulation over two watersheds and the 23 downstream alluvial fan, although at significant computational cost. We also develop a 24 Gaussian Process surrogate model, allowing for rapid prediction of simulator outputs for 25 untested scenarios. We utilize this framework to explore model sensitivity to rainfall in-26 tensity and sediment availability as well as parameters associated with saturated hydraulic 27 conductivity, hydraulic roughness, grain size, and sediment entrainment. Simulation re-28 sults are most sensitive to peak rainfall intensity and hydraulic roughness. We further 29 use this approach to examine variations in debris flow inundation patterns at different 30 stages of postfire recovery. Sensitivity analysis indicates that constraints on temporal 31 changes in hydraulic roughness, saturated hydraulic conductivity, and grain size follow-32 ing fire would be particularly beneficial for forecasting debris flow runout throughout the 33 postfire recovery period. The emulator methodology presented here also provides a means 34 to compute the probability of a debris flow inundating a specific downstream region, con-35 sequent to a forecast or design rainstorm. This workflow could be employed in scenario-36 based planning for postfire hazard mitigation. 37

³⁸ Plain Language Summary

Fire on steep hillslopes increases the potential for debris flows, or rapidly moving 39 mixtures of water, soil, ash, and rock, that can develop during intense rainfall. Debris 40 flows threaten communities situated downstream of steep, burned areas. Burn severity 41 and hillslope steepness give some indication of the potential for debris flow initiation in 42 response to a particular rainstorm, but these factors alone do not indicate which down-43 stream areas might be inundated by a flow. To investigate areas impacted by potential 44 debris flows, we utilize a computational model that represents the physical processes of 45 debris flow initiation and runout. Such process-based models are computationally inten-46 sive, which has limited their use in rapid hazard assessments. Here, we implement a de-47

⁴⁸ bris flow initiation and runout model in a high-performance computing framework. Even
⁴⁹ so, a typical debris flow simulation takes several hours to complete on a supercomputer.
⁵⁰ Thus, we also build statistical models of these physical models, so-called emulators. Emulators can rapidly approximate the debris flow simulation and thus offer a mechanism
⁵² to investigate outcomes of debris flow models over a wide range of scenarios. We investigate how debris flow inundation footprints vary with rainfall intensity and the impact
⁵⁴ of landscape recovery on debris flows.

55 1 Introduction

Fire alters soil and vegetation, leading to increases in runoff and erosion (Shakesby 56 & Doerr, 2006; Moody et al., 2013). In extreme cases, particularly when steep water-57 sheds burn at moderate or high severity, rapid entrainment of sediment into runoff can 58 produce debris flows (Kean et al., 2011; Gabet & Bookter, 2008; Esposito et al., 2023; 59 Conedera et al., 2003; Nyman et al., 2011; Diakakis et al., 2023). Postfire debris flows 60 generated by runoff are most common in the first year following fire, when fire-driven 61 reductions in soil infiltration capacity, rainfall interception, and hydraulic roughness are 62 most extreme (DeGraff et al., 2015; Hoch et al., 2021; Thomas et al., 2021; Esposito et 63 al., 2023; Graber et al., 2023). Due to the complex interactions among runoff, sediment 64 transport, and debris-flow initiation and runout following fire, mathematical models that 65 couple these processes have the potential to inform our understanding of postfire debris-66 flow hazards and how they change through time as landscapes recover (McGuire et al., 67 2021). Yet the exploration of postfire debris flow processes through application of mor-68 phodynamic models for runoff and sediment transport is often limited by the high di-69 mensionality, poor constraints on parameters, and substantial computation time of the 70 models. Quantification of uncertainties associated with sources and model parameters, 71 and incorporation of those uncertainties in probabilistic predictions of hazard, often re-72 quire use of simulation ensembles with hundreds or thousands of members, greatly in-73 creasing computational costs (Bayarri et al., 2015). In this work, we accelerate a recently 74 developed morphodynamic model of runoff and sediment transport (McGuire et al., 2017), 75 and pair model runs with stochastic surrogates for high-dimensional output (Gu & Berger, 76 2016) as a strategy for simulating postfire debris-flow runout. This acceleration also en-77 ables us to rapidly explore how temporal changes in fire-affected model input parame-78 ters influence debris flow runout. 79

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The processes leading to the initiation and growth of runoff-generated debris flows 80 following fire involve the generation of spatially distributed overland flow and its sub-81 sequent interaction with sediment on hillslopes and in channels (Santi et al., 2008; Sta-82 ley et al., 2014; McGuire et al., 2017; Guilinger et al., 2020). This presents a contrast 83 to debris flows that mobilize from shallow landslides, which initiate when infiltration pro-84 motes increases in pore-water pressure that causes a discrete mass of soil to become un-85 stable on a hillslope (Iverson et al., 1997). The source of sediment for postfire runoff-86 generated debris flows may come from a combination of processes, including widespread, 87 shallow erosion on hillslopes in response to raindrop-driven sediment transport and un-88 confined sheet flow, rill erosion on hillslopes in areas of concentrated flow, and channel 89 scour (Santi et al., 2008; Staley et al., 2014; McGuire et al., 2017; Tang et al., 2019). All 90 three processes are more efficient at eroding sediment following fire as a result of decreases 91 in ground cover and increases in runoff (Robichaud et al., 2016), particularly rill and chan-92 nel erosion processes where overland flow does the work to entrain and transport sed-93 iment (Sheridan et al., 2007; Wagenbrenner et al., 2010). In areas of unconfined, shal-94 low flow, raindrops facilitate sediment detachment and transport in combination with 95 runoff (Kinnell, 2005). Raindrop-driven sediment transport on hillslopes increases fol-96 lowing fire due to removal of the vegetation canopy, litter, and duff, that tend to shield 97 the soil surface from raindrop impact in unburned settings. 98

Models designed to simulate runoff-generated debris flows from initiation to depo-99 sition must therefore account for spatially distributed runoff and sediment transport as 100 well as changes in flow behavior resulting from spatial and temporal variations in sed-101 iment concentration. Fully developed debris flows are characterized by volumetric sed-102 iment concentrations in excess of 40-50%, though they initiate from runoff with initially 103 negligible sediment concentration. Due, in part, to the relatively flashy hydrologic re-104 sponse of watersheds burned by fire, postfire runoff-generated debris flows initiate in re-105 sponse to short-duration bursts of high intensity rainfall (Kean et al., 2011). Rainfall in-106 tensity averaged over a 15-minute time period, I_{15} , is correlated well with runoff mag-107 nitude at the outlet of small, recently burned watersheds and threshold values of I_{15} have 108 proven to be reasonable predictors for debris flow initiation in the western USA (Kean 109 et al., 2011; Staley et al., 2017). Rainstorms that contain multiple, distinct bursts of high 110 intensity rainfall, such as where I_{15} exceeds the debris flow threshold at multiple times 111 in a single event, may lead to multiple pulses of debris flow activity (Kean et al., 2011). 112

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One benefit of morphodynamic models that are capable of simulating the debris-flow life-113 cycle from initiation to runout is their ability to directly account for the impacts of tem-114 porally varying rainfall intensity on debris flow processes, including runout and inunda-115 tion extent. In contrast, models designed only to simulate debris flow runout processes 116 (i.e. neglecting runoff and sediment transport), can be employed by defining an inflow 117 hydrograph above the anticipated runout zone based on the rainfall hydrograph and an 118 estimated debris flow volume, or by allowing a pile of sediment and water to flow down-119 stream from a pre-defined initiation zone (Barnhart et al., 2021; Gorr et al., 2022; Gib-120 son et al., 2022). As a result, employing morphodynamics models to estimate debris flow 121 runout avoids introducing epistemic uncertainty associated with specifying a volume of 122 material associated with an inflow hydrograph. 123

McGuire et al. (2017) developed a model that accounts for infiltration, runoff, sed-124 iment transport, and changes to flow resistance driven by sediment concentration in or-125 der to simulate postfire debris flow initiation and growth. In this model, rainfall drives 126 sediment entrainment and transport processes that naturally lead to debris flow initi-127 ation when hydrogeomorphic conditions give rise to flows with sufficiently high sediment 128 concentrations. Since rainfall, runoff, and erosion processes are related to model param-129 eters known to change following fire, such as saturated hydraulic conductivity (Perkins 130 et al., 2022; Ebel et al., 2022; Thomas et al., 2021), hydraulic roughness (Stoof et al., 131 2015), and vegetation cover (Stoof et al., 2012), it is possible to use this framework to 132 explore how postfire recovery affects debris-flow initiation, growth, and runout. The model 133 proposed by McGuire et al. (2017) contains a number of parameters that are challeng-134 ing to constrain and is computationally intensive, especially when simulating debris flow 135 initiation and runout processes over large areas. Thus far, model applications have been 136 limited to examining debris-flow initiation processes in small headwater basins where prior 137 work and intensive field monitoring helped constrain parameter values (McGuire et al., 138 2021, 2017; Tang et al., 2019). Here, we employ adaptive mesh refinement and parallel 139 computations from the Titan framework (Patra et al., 2005; Dalbey et al., 2008) to make 140 model application more tractable over larger spatial domains, specifically with the goal 141 of simulating the entirety of the postfire debris-flow lifecycle from runoff generation and 142 debris flow growth to runout. In addition, we employ statistical emulators, which enable 143 fast approximations to solutions of the model equations, in order to explore the influ-144 ence of model parameters on debris-flow runout extent. 145

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Gaussian process emulators (GPs) are a powerful class of statistical surrogates that 146 enable rapid approximation and uncertainty quantification of computationally intensive 147 first principles (conservation laws) based models or, simulators (Currin et al., 1988; Sacks, 148 Welch, et al., 1989; Welch et al., 1992). In the context of postfire debris-flows, GPs of-149 fer a mechanism to quickly explore spatial patterns in peak flow depth and area inun-150 dated for various parameter settings. GPs also allow for uncertainty quantification via 151 Monte Carlo (MC) simulations of flow inundation. Along the way, GPs offer an approx-152 imate sensitivity analysis of physical parameters' effects on debris-flow inundation. Par-153 allel partial emulation (PPE) (Gu & Berger, 2016) extends the GP methodology to vector-154 valued outputs – in our case, flow depth at each map point (pixel). Some recent stud-155 ies use PPE for flow model sensitivity analysis and calibration (Zhao et al., 2021; Zhao 156 & Kowalski, 2022). In this work, we apply PPE to explore the effects of rainfall inten-157 sity and sediment availability as well as parameters associated with saturated hydraulic 158 conductivity, hydraulic roughness, grain size, and sediment entrainment on peak debris 159 flow depth and inundation extent. 160

¹⁶¹ 2 Study Area

The Thomas Fire ignited in December 2017 and burned more than 1100 km^2 , in-162 cluding a series of steep watersheds in the Santa Ynez Mountains above the community 163 of Montecito, USA (Figure 1a). On 9 January 2018, widespread rainfall developed over 164 the burned area in association with an atmospheric river (Oakley et al., 2018). A nar-165 row cold frontal rainband (NCFR), a relatively small-scale feature characterized by a band 166 of intense precipitation that forms along a cold front, moved over burned watersheds above 167 Montecito and produced a short-duration burst of intense rainfall. Peak 15-minute rain-168 fall intensities, I_{15} , in this area ranged from approximately 78 mm/h to 105 mm/h (Kean 169 et al., 2019). During this time period, rainfall intensities greatly exceeded the infiltra-170 tion capacity of the soil, leading to infiltration-excess overland flow that generated rills 171 on steep hillslopes (Alessio et al., 2021). The combination of intense hillslope erosion and 172 channel incision led to runoff-generated debris flows that traveled across the populated 173 alluvial fan (Kean et al., 2019; Morell et al., 2021; Alessio et al., 2021). The debris flows 174 that initiated in six watersheds above Montecito mobilized more than $630,000 \text{ m}^3$ of sed-175 iment and led to 23 fatalities, 408 damaged structures, and more than \$1 billion in dam-176 age (Kean et al., 2019; Lancaster et al., 2021). In this study, we focus on modeling de-177

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Figure 1. (a) The Thomas Fire ignited on 4 December 2017 and burned over 1140 km² in southern California, USA. (b) The western portion of the burned area included a series of steep watersheds (black rectangle) above the community of Montecito, located near Santa Barbara. (c) Our study focuses on two watersheds, Oak Creek (0.45 km²) and San Ysidro Creek (7.6 km²), located upstream of Montecito.

bris flow initiation, growth, and runout for two of these watersheds, Oak Creek and San 178 Ysidro Creek. These watersheds are well suited for our study since high resolution to-179 pographic and rainfall data are available and we can leverage data collection following 180 the fire (Kean et al., 2019) and previous work in the Transverse Ranges of southern Cal-181 ifornia (McGuire et al., 2016, 2021; Tang et al., 2019; McGuire et al., 2017) to constrain 182 model parameter ranges. In addition, the model domain, which includes the watersheds 183 and downstream alluvial fan, is large enough to make physics-based simulations com-184 putationally challenging (Figure 1c). 185

Roughly 85% of San Ysidro Creek, which has a total drainage area of 7.6 km², burned at moderate or high severity (Figure 1c). Approximately 49% of Oak Creek, which is substantially smaller at 0.45 km², burned at moderate or high severity. Both the San Ysidro

and Oak Creek watersheds are steep, with median slopes of 28 and 21 degrees, respec-

¹⁹⁰ tively. The debris flow that initiated in San Ysidro Creek mobilized a total volume of

 $_{191}$ 297,000 m³ and inundated an area of 905,000 m² while the smaller Oak Creek watershed

 $_{192}$ produced a debris flow that moved 10,000 m³ of sediment and inundated 102,000 m²

(Kean et al., 2019) (Figure 1c). The grain size distribution in the debris flows was bi-

¹⁹⁴ modal, consisting of a sandy matrix that suspended boulders with a large axis greater

than several meters (Kean et al., 2019).

¹⁹⁶ 3 Methodology

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3.1 Simulating runoff-generated debris flows with Titan2D

Steep watersheds recently burned by fire often experience greater amounts of runoff 198 and increased rates of sediment transport. Factors affecting rates of sediment transport, 199 and also the initiation and growth of runoff-generated debris flows, include rainfall in-200 tensity and duration, vegetation cover, soil infiltration capacity, and sediment charac-201 teristics (e.g. grain size, erodibility). The model developed by McGuire et al. (2017) rep-202 resents rainfall, infiltration, fluid flow, and sediment entrainment and deposition processes, 203 which makes it a useful framework for simulating runoff-generated debris flows in steep 204 terrain (McGuire et al., 2017; Tang et al., 2020). We provide a brief overview of the gov-205 erning equations, which we solve within the Titan2D framework (Patra et al., 2005; Simakov 206 et al., 2019). 207

The Titan2D code employs an adaptive mesh, finite volume scheme to solve hy-208 perbolic PDEs describing shallow-water like mass flows over digital elevation models of 209 real topography. Titan2D (Patra et al., 2005) was originally developed to solve the depth 210 averaged shallow-water mass flow equations by Savage and Hutter (1989). Titan was mod-211 ernized and restructured in 2019 (Simakov et al., 2019) to optimize storage and access 212 for parallel adaptive mesh refinement, and to facilitate the usage of new material mod-213 els. Using Titan2D to solve the model equations proposed by McGuire et al. (2017) there-214 fore offers several advantages, especially for simulating debris flows over spatial scales 215 of more than a few square kilometers. In particular, Titan2D uses partitioned hash-tables 216 for better memory allocation structures, allowing it to compute over a large domain em-217 ploying adaptive mesh refinement (AMR) and unrefinement, with computational efficiency. 218

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It is also well suited for parallel computing using MPI and OpenMP and dynamic load balancing for exploiting multiprocessor computing. The computational efficiencies attained make it tractable to simulate end-to-end flows from initiation to downstream inundation.

The equations representing the motion of fluid and sediment can be written as a set of depth averaged conservation laws. We have

$$\frac{\partial \mathbf{U}}{\partial t} + \frac{\partial \mathbf{F}}{\partial x} + \frac{\partial \mathbf{G}}{\partial y} = \mathbf{S_0} + \mathbf{S_1} + \mathbf{S_2},\tag{1}$$

226 where

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$$\mathbf{U} = \left\{ h \quad uh \quad vh \quad c_1h \quad \cdots \quad c_kh \right\}^T, \tag{2}$$

$$\mathbf{F} = \left\{ hu \quad hu^2 + \frac{1}{2}g_z h^2 \quad huv \quad huc_1 \quad \cdots \quad huc_k \right\}_{-}^{T}, \tag{3}$$

$$\mathbf{G} = \left\{ hv \quad huv \quad hv^2 + \frac{1}{2}g_z h^2 \quad hvc_1 \quad \cdots \quad hvc_k \right\}^T, \tag{4}$$

and where h, u, v, and c_i are flow depth, velocity along x-axis, velocity along y- axis, 230 and sediment concentration of particle size class i. Components of gravitational accel-231 eration in the x, y, and z directions are given by g_x , g_y , and g_z , respectively, and k de-232 notes the number of particle size classes. $\mathbf{S_0}, \mathbf{S_1}$ and $\mathbf{S_2}$ are source terms. $\mathbf{S_0}$ denotes 233 the contributions of mass sources and sinks associated with the effective rainfall rate, P_{eff} , 234 and the soil infiltration capacity, I, as well as momentum sources and sinks arising from 235 variations in topographic elevation, and spatial variations in sediment concentration and 236 debris flow resistance terms, S_x and S_y . Specifically, S_0 is given as 237

$$\mathbf{S_0} = \begin{cases} P_{eff} - I + \frac{\partial z}{\partial t} \\ -g_x h + \gamma_x - \psi S_x \\ -g_y h + \gamma_y - \psi S_y \\ 0 \\ \vdots \\ 0 \end{cases} \right\}.$$
(5)

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The debris flow resistance terms are scaled by ψ , which increases linearly from 0 to 1 as the volumetric sediment concentration increases from 0.2 to 0.4. This scaling factor gradually increases the importance of the debris flow resistance terms as volumetric sediment concentration approaches levels that are consistent with a transition from with flood flow to debris flow. The terms γ_x and γ_y account for the effects of spatially variable sediment

concentration and are given by

$$\gamma_x = \frac{-(\rho_s - \rho_w)g_z h^2}{2\rho_f} \frac{\partial c}{\partial x} \tag{6}$$

and

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$$\gamma_y = \frac{-(\rho_s - \rho_w)g_z h^2}{2\rho_f} \frac{\partial c}{\partial y} \tag{7}$$

Here, c denotes volumetric sediment concentration, $\rho_w = 1000 \text{ kg m}^{-3}$ the density of

water, $\rho_s = 2600 \text{ kg m}^{-3}$ the density of sediment, and $\rho_f = c\rho_f + (1-c)\rho_w$ the den-

 $_{241}$ sity of the flow. S_1 accounts for flow resistance using a depth-dependent Manning's for-

242 mulation, and is given as

243
$$\mathbf{S_1} = \begin{cases} 0 \\ g_z \eta^2 h u \sqrt{h u^2 + h v^2} / h^{7/3} \\ g_z \eta^2 h v \sqrt{h u^2 + h v^2} / h^{7/3} \\ 0 \\ \vdots \\ 0 \\ 0 \\ \end{cases},$$
(8)

where η is the Manning friction coefficient. The friction coefficient varies with flow depth according to

246
$$\eta = \begin{cases} \eta_0 (h/h_c)^{-\epsilon} & h \le h_c \\ \eta_0 & h > h_c \end{cases},$$
(9)

where η_0 is the hydraulic roughness coefficient, h_c is a critical flow depth and ϵ is a phenomenological exponent. Soil infiltration capacity, I, is represented by the Green-Ampt model where

$$I = k_s \frac{Z_f + h_f + h}{Z_f},\tag{10}$$

with k_s denoting saturated hydraulic conductivity, h_f the wetting front potential, $Z_f = V/(\theta_s - \theta_i)$ the depth of the wetting front, V the cumulative infiltrated depth, θ_s the volumetric soil moisture content at saturation, and θ_i the initial volumetric soil moisture content. The source term S_2 accounts for sediment entrainment and deposition processes, which are represented using the framework proposed (Hairsine & Rose, 1992a, ²⁵⁶ 1992b). In particular,

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$$\mathbf{S_2} = \begin{cases} 0 \\ 0 \\ e_1 + e_{r1} + r_1 + r_{r1} - d_1 \\ \vdots \\ e_k + e_{rk} + r_k + r_{rk} - d_k \end{cases},$$
(11)

where e_k and e_{rk} are sediment detachment and re-detachment rates due to raindrop im-258 pact for sediment particles in size class k, r_k and r_{rk} are rates of entrainment and re-259 entrainment due to runoff, and d_k is the effective deposition rate. The model differen-260 tiates between original soil, which has not yet been entrained and transported during the 261 modeled rainstorm, and deposited sediment, which has been detached and subsequently 262 deposited. Detachment rates for entraining original sediment and re-entraining deposited 263 sediment are computed differently. Sediment in the deposited layer may also fail en-masse 264 (McGuire et al., 2017). Rates of sediment entrainment and re-entrainment by runoff are 265 given by 266

$$r_k = (1 - H)p_k \frac{F(\Omega - \Omega_{cr})}{J},$$
(12)

268 and

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$$_{rk} = H \frac{m_k}{m_t} \frac{F(\Omega - \Omega_{cr})}{\frac{\rho_s - \rho_f}{\rho_s} gh}.$$
(13)

Here, m_k is the deposited sediment mass per unit area for sediment in size class k, m_t 270 is the total mass of deposited sediment per unit area, $H = \min(m_t/m_t^*, 1)$ accounts 271 for the degree to which deposited sediment shields the underlying bed from erosion, m_t^* 272 is the mass of deposited sediment needed to completely shield original sediment from ero-273 sion, ρ_f is the density of the flow, ρ_s is the density of sediment, F denotes the fraction 274 of stream power effective in sediment entrainment, $\Omega = \rho_f g S_f \sqrt{uh^2 + vh^2}$ is stream 275 power, and $S_f = \eta^2 (uh^2 + vh^2)h^{-10/3}$ is the friction slope. In this work, we consider 276 a single particle size class characterized by a representative particle diameter, δ . 277

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3.2 Rainfall and model parameters

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A digital elevation model (DEM) of the study area is input to the Titan2D simulation. Here, we use a 1 m DEM derived from post-event airborne lidar. Elevations and slopes at locations required by the computational mesh are obtained using a 9 point $(3 \times$ 3) finite difference stencil to interpolate on the DEM grid reducing the effects of artifacts and noise in the data (Patra et al., 2005). Errors in the DEM could be treated as
an uncertainty that is propagated through the simulation and subsequent analysis (Stefanescu,
Bursik, & Patra, 2012; Stefanescu, Bursik, Cordoba, et al., 2012), but we did not consider this in the present study.

Runoff and debris flows initiated in the study area in response to a short duration, high intensity burst of rainfall in the early morning hours of 9 January 2018 (Kean et al., 2019). All simulations use 1-minute rainfall intensity data derived from the KTYD rain gauge for a 20-minute time period that spans this short temporal window when rainfall intensity rapidly increased and debris flows initiated (Figure 2). The gauge is maintained by the Santa Barbara County Flood Control District and located approximately 5 km west of the San Ysidro Creek watershed.



Figure 2. Rainfall hyetograph derived from the KTYD rain gauge, located roughly 5 km west of the San Ysidro Creek watershed.

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Simulations were designed to explore the extent to which inundated area and peak 294 flow depths on the alluvial fan were influenced by rainfall intensity as well as several pa-295 rameters that may play critical roles in debris-flow initiation and growth. We explored 296 the impact of different rainfall intensities by multiplying the 1-minute rainfall intensity 297 time series by a rainfall intensity factor (RI_{fac}) that varied from 0.5 to 1.5. We also var-298 ied the representative particle diameter, δ , from 0.05–0.125 mm, the fraction of stream 299 power effective in entrainment, F, from 0.001 - 0.006, the hydraulic roughness coeffi-300 cient, n_0 , from 0.03 – 0.2, and saturated hydraulic conductivity, k_s , from 5 – 25 mm 301

h⁻¹. We further enforced a maximum soil thickness, r_{max} , that varied from 0.25–1.5 m to explore the role of sediment availability. All other parameters were fixed (Table 1, Table 2). We used a Latin hypercube sampling strategy to generate 64 different parameter sets from the ranges specified above (McKay et al., 1979). Figure 3 shows the maximum flow depth for three of the 64 training simulations chosen to demonstrate typical results with low, moderate, and long runout extents. Each of these simulations took several hours to complete on an HPC cluster using up to 16 cores on an intel Xeon Gold 6226R processor.



Figure 3. Maximum flow depth of three simulations resulting in (a) short runout (b) moderate runout and (c) long runout onto the alluvial fan. For contrast, the maximum colorbar limit is set to 2.5 m although maximum flow depth does exceed 2.5 m in some locations.

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We chose to focus on exploring the effects of rainfall intensity, δ , F, n_0 , k_s , and soil 310 thickness (sediment availability) since they control different aspects of the debris flow 311 initiation and growth process and, aside from rainfall intensity, they may all be strongly 312 affected by fire in our study area. Peak rainfall intensity over sub-hourly durations, par-313 ticularly the 15-minute duration, is correlated well with runoff in recently burned wa-314 tersheds in southern California (Kean et al., 2011). Peak 15-minute rainfall intensity is 315 also used in empirical models designed to predict postfire debris-flow likelihood and vol-316 ume in the western USA (Staley et al., 2017; Gartner et al., 2014). We therefore expect 317 that variations in rainfall intensity during the relatively short (< 0.5 hours) portion of 318 the rainstorm that we are modeling will influence debris flow processes. 319

We expect the representative grain size, δ , to be relatively small in areas of con-320 centrated flow immediately following fire in our study area given the propensity for post-321 fire dry ravel to transport hillslope sediment to channels and valley bottoms (Florsheim 322 et al., 1991; Lamb et al., 2011). Both δ and the amount of sediment available for trans-323 port, which we vary by enforcing a maximum soil thickness (r_{max}) throughout the model 324 domain, may vary as a function of time since fire as sediment is exported from postfire 325 rainstorms (Tang et al., 2019). Similarly, Liu et al. (2021) found that k_s and the Man-326 ning coefficient were lowest during rainstorms in the first year following a high severity 327 fire in the San Gabriel Mountains, southern California, and increased by factors of roughly 328 3-4 over the next 4 years. Immediately following fire in southern California, values for 329 the Manning coefficient and saturated hydraulic conductivity may be as low as 0.025-330 $0.07 \text{ sm}^{-1/3}$ and $1-6 \text{ mm h}^{-1}$, respectively (Rengers et al., 2016; Tang et al., 2019; 331 Liu et al., 2021). Kean et al. (2019) used post-event, point scale measurements with a 332 tension infiltrometer to estimate the geometric mean of saturated hydraulic conductiv-333 ity at 20 mm h^{-1} in the days following the Montecito debris flows. The effective frac-334 tion of stream power, F, may be expected to increase immediately following fire due to 335 reductions in roughness associated with ground cover and vegetation. Past studies sug-336 gest values of $F \approx 0.005$ perform reasonably well in steep, recently burned watersheds 337 in southern California (McGuire et al., 2017; Tang et al., 2019). 338

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3.3 Emulating debris flows

Statistical emulators are effectively probabilistic models of computationally inten-340 sive physical model systems or *simulators*. That is, statistical emulators relate a set of 341 user-defined inputs, often physical parameter specifications, to simulator output. Gaus-342 sian process emulators (GPs) are a popular class of surrogates for approximating and 343 quantifying uncertainties in simulators as they (almost) interpolate computer model out-344 put (Sacks, Schiller, & Welch, 1989; Sacks, Welch, et al., 1989; Santner et al., 2003; Ras-345 mussen & Williams, 2006). Further, the variance of the associated GP offers a quick mech-346 anism to assess the uncertainty of using the emulator in place of the simulator for model 347 prediction at untested inputs. Thus GP emulators offer a rapid and quantifiable mech-348 anism to approximate output from physical process models that are computationally in-349 tensive to exercise. The parallel partial emulator (PPE) (Gu & Berger, 2016) extends 350 this surrogate model to vector valued output. 351

Inputs to GP emulators are user defined. They are typically influential parameters, which show up within the governing dynamics, the forcing terms, or boundary conditions, as opposed to independent variables in the physical model. For the model described in Sec 3.1, we choose p = 6 parameters to define our input vector, namely those described in Sec 3.2 and given by $\mathbf{q} = [k_s, r_{max}, \eta_0, F, \delta, RI_{fac}]$. The rainfall intensity factor is a scaling of the true rainfall time series that triggered the debris flows in our study area.

In initial explorations, the ratio of pore fluid pressure to total basal normal stress, 359 λ , was also considered, but found to not substantially influence spatial patterns in peak 360 flow depth and runout extent on the fan. We will discuss the relationship between GP 361 emulators and sensitivity analysis further in Sec 4. The output under consideration, \mathbf{y} , 362 is the maximum (over time) flow depth at each of s = 1.4M map points. The main ob-363 jective of the emulator is to predict the output of the Titan2D model at an untested sce-364 nario, \mathbf{q}^* , given a relatively modest set of N training or design runs \mathbf{q}^{D} and each of their 365 corresponding inundation depth outputs, \mathbf{y}_j^{D} , $j = 1, \dots, N$. In this work, we take N =366 64 training runs and each output vector, $\mathbf{y}_{i}^{\mathrm{D}}$, is a 1.4M element vector recording the peak 367 flow depth at each map point. Collecting these outputs together, we have $Y^{\rm D}$, as a 64× 368 1.4M matrix of training run outputs. The 64 training run inputs are chosen by a Latin 369 hypercube design (McKay et al., 1979; Santner et al., 2003) covering the ranges of in-370 puts listed in Table 1. All other parameters are fixed (Table 2). To fit the emulator, these 371 parameter ranges are normalized to a unit hypercube. 372

Given the training data, $\{\mathbf{q}^{\mathrm{D}}, Y^{\mathrm{D}}\}$, to approximate the inundation resulting from an untested scenario, \mathbf{q}^* , we use the predictive mean of the PPE given by

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 $\tilde{\mathbf{y}}(\mathbf{q}^*) = \mathbf{h}^{\mathrm{T}}(\mathbf{q}^*)B + \mathbf{r}^{\mathrm{T}}(\mathbf{q}^*)R^{-1}(Y^{\mathrm{D}} - H^{\mathrm{D}}B),$ (14)

where R is an $N \times N$ (64×64 in this work) matrix of correlations between pairs of de-376 sign inputs, $\mathbf{r}(\mathbf{q}^*)$ is an $N \times 1$ vector of correlations between the untested input, \mathbf{q}^* , and 377 each of the input scenarios in the design, \mathbf{q}^{D} . Further, $\mathbf{h}(\mathbf{q})$ is a $l \times 1$ vector of regres-378 sion variables, often taken to be constant or linear in \mathbf{q} (i.e., l = 1 for constant case used 379 in this work and l = p + 1 for the linear case), and H^{D} is and $N \times l$ matrix where the 380 j^{th} row are the regression variables evaluated at the j^{th} design point, $\mathbf{h}^{\mathrm{T}}(\mathbf{q}_{j}^{\mathrm{D}})$. The ma-381 trix B is a $l \times s$ matrix of regression coefficients. Here, each of the s = 1.4M outputs 382 has its own set of regression coefficients, but a shared correlation structure. We use a 383

Model parameter	Min value	Max value	Range parameter
k_s : Saturated hydraulic conductivity (mm	5	20	3.5
h^{-1})			
r_{max} : Maximum soil thickness (m)	0.25	1.5	9.9
η_0 : hydraulic roughness coefficient (s m ^{-1/3})	0.03	0.2	0.15
F: Fraction of stream power effective in sedi-	0.6×10^{-3}	1.0×10^{-3}	9.9
ment detachment			
δ : Effective grain size (mm)	0.05	0.125	5.8
$RI_{\rm fac}$: Rainfall intensity factor	0.5	1.5	0.38

Table 1. Parameter ranges with units and GP range parameters (unit-less) for the six parameters that varied among the N=64 debris flow simulations.

Matérn 5/2 correlation function (Stein, 1999). For two scenarios, e.g., two input points 384 $\mathbf{q}_i = (x_{i1}, \ldots, x_{ip})^{\mathrm{T}}$ and $\mathbf{q}_j = (x_{j1}, \ldots, x_{jp})^{\mathrm{T}}$, the standardized distance and correla-385 tion between these input scenarios are given by 386

$$d_{k} = \left(\frac{|x_{ik} - x_{jk}|^{2}}{\theta_{k}^{2}}\right)^{1/2}, \quad k = 1, \dots, p$$

$$c(\mathbf{q}_{i}, \mathbf{q}_{j}) = \prod_{k=1}^{p} \left(1 + \sqrt{5}d_{k} + \frac{5}{3}d_{k}^{2}\right) \exp\left(-\sqrt{5}d_{k}\right), \quad (15)$$

389

respectively. The predictive variance for each output dimension (pixel) of the PPE is given 390 by 301

³⁹²
$$\mathbf{v}_{j}(\mathbf{q}^{*}) = \sigma_{j}^{2} \left(1 - \mathbf{r}^{\mathrm{T}}(\mathbf{q}^{*})R^{-1}\mathbf{r}(\mathbf{q}^{*}) + \left(\mathbf{h}(\mathbf{q}^{*}) - (H^{\mathrm{D}})^{\mathrm{T}}R^{-1}\mathbf{r}(\mathbf{q}^{*})\right)^{\mathrm{T}} \right)^{\mathrm{T}}$$

$$\times \left((H^{\mathrm{D}})^{\mathrm{T}} R^{-1} H^{\mathrm{D}} \right)^{-1} \left(\mathbf{h}(\mathbf{q}^{*}) - (H^{\mathrm{D}})^{\mathrm{T}} R^{-1} \mathbf{r}(\mathbf{q}^{*}) \right) \right), \tag{16}$$

where σ_j^2 , (j = 1, ..., s) is the scalar variance corresponding to each pixel's output. "Fit-395 ting" a PPE amounts to estimating the regression parameters in B, the scalar variances 396 at of each output, σ_j^2 , and the range parameters $\{\theta_k : k = 1, \ldots, p\}$. To do so, we use 397 the RobustGaSP package (Gu et al., 2018, 2019). On a laptop, fitting a PPE to 1.4M398 pixels of output with N = 64 training runs takes roughly 10 minutes. 399

GP emulators have been applied to Titan2D-based volcanic debris flows (Bayarri 400 et al., 2009, 2015; Spiller et al., 2014; Rutarindwa et al., 2019) and recently to other Titan2D-401 based debris flows (Zhao et al., 2021; Zhao & Kowalski, 2022). In each of these studies, 402

Symbol	Definition	Value	Unit
a_0	Detachability of original soil	1000	$\rm kg\ m^{-2}\ s^{-1}$
a_{d0}	Detachability of deposited sediment	2000	$\rm kg\ m^{-2}\ s^{-1}$
m_{t0}^*	Deposited sediment needed to shield original soil	2.7	${\rm kg}~{\rm m}^{-2}$
J	Specific energy of entrainment	15.125	$\mathrm{m}^2~\mathrm{s}^{-2}$
C	Effective cohesion	200	Pa
ϕ_{bed}	Basal friction angle	32	\deg
λ	Ratio of pore fluid pressure to total normal stress	0.8	-
C_v	Fraction vegetation cover	0	-
h_f	Wetting front poential	1	mm
$ heta_i$	Initial volumetric soil moisture	0.1	-
θ_s	Volumetric soil moisture at saturation	0.39	-
ϵ	Exponent in friction model	0.33	-
h_c	Critical depth in friction model	3	mm

Table 2. Model parameters using the same notation as (McGuire et al., 2017).

source terms (particularly debris mass or flux) were specified via ad-hoc parameteriza tions which are less appropriate for postfire, runoff-generated debris flows.

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3.4 Numerical Experiments

Evaluation of the GP emulator's mean quickly allows one to explore any output quantity of interest over the parameter space. Here we take the output quantity of interest, \mathbf{y} , to be the maximum debris-flow depth at all locations. Additionally, the variance of the GP emulator accounts for the uncertainty introduced by evaluating the GP mean, $\tilde{\mathbf{y}}$, instead of the debris-flow process model. We can break our exploration of numerical experiments into three groups.

First, we perform leave-one-out experiments as a test of the PPE performance. This experiment amounts to excluding one simulation at a time, fitting a GP to the 63 remaining simulations, and then comparing the GP predicted inundation of the left-out scenario to actual simulated inundation for that scenario. This is repeated for each of the N =64 simulations. Second, we explore the relative importance of different model parameters using the GP's range parameters. The range parameters are positive numbers indicating the influence of each model parameter on the model response – the smaller the range parameter, the more influence the corresponding model parameter has on the debris flow model (i.e. maximum flow depth). As such, these range parameters act as an effective sensitivity analysis.

Lastly, we employ the emulator to isolate and explore how flow extent/depths are 423 driven by (1) changes in rainfall intensity and (2) changes in saturated hydraulic con-424 ductivity and hydraulic roughness that occur as the landscape recovers. We focus, in par-425 ticular, on exploring the effects of postfire changes in saturated hydraulic conductivity 426 and hydraulic roughness since Liu et al. (2021) provide guidance for parameterizing these 427 effects in the nearby San Gabriel Mountains. Since the GP emulator enables rapid for-428 ward uncertainty quantification, we demonstrate how it can be used to accelerate a Monte 429 Carlo probability of inundation calculation for two cases, namely when the observed storm 430 occurs 2 months and 14 months postfire. 431

432 4 Results

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4.1 Emulator performance

To test the performance of the GP emulator for approximating the Titan2D sim-434 ulations, we examine leave-one-out predictions. Of note, the range parameter estimates 435 were very stable. We found the coefficients of variation for each of the six values to be 436 between 0.01-0.10 indicating that the relative influence of any input to the GP was not 437 swayed strongly by any single flow simulation. For illustrative purposes, we focus on two 438 points of interest – one located in a channel and one on the adjacent fan surface where 439 flow is relatively unconfined. In Figure 4 for both cases we sort the simulations by their 440 left-out flow depths, \mathbf{y} , and predict each with a credible interval centered at the mean 441 of the GP, \tilde{y} . We find root mean squared errors from these leave-one-out experiments 442 of 0.12m and 0.17m for the locations on the fan and on the channel, respectively. Fur-443 ther, we see that 89% (fan location) and 94% (channel location) of the simulated depths 444 fall within their predictive credible intervals. These numbers are slightly below the an-445 ticipated 95%, but this is likely due to the relatively small training set, and in the case 446

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of the fan location, the fact that roughly half of the simulations resulted in no inundation, which is challenging for GP emulation (Spiller et al., 2023).



Figure 4. Leave-one-out experiments for (a) a location on the fan and (b) a location in the channel (see Figure 7 for details). In each panel, indices are sorted based on the simulated max flow depth (red stars). GP predictive means for these scenarios are plotted in black while the 95% credible intervals are plotted as vertical blue bars.

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4.2 Sensitivity analysis

A crucial step to fitting a GP is estimating the range parameters. Smaller range parameters indicate that the corresponding model parameter has more influence on the debris flow model output of interest (i.e. maximum flow depth). From table 1, one can see that the debris flow model is most sensitive to rainfall intensity and the hydraulic roughness coefficient; it is moderately sensitive to saturated hydraulic conductivity and effective grain size; and is relatively insensitive to the maximum soil thickness and to the fraction of stream power effective in sediment detachment.

Scaling the rainfall intensity time series has a substantial effect on inundation extent (Figure 5). As the rainfall leading to the flows after the Thomas fire were quite intense, it is not surprising to see significant runout even when $RI_{fac} = 0.5$, though the extent of inundation is diminished relative to cases with more intense rainfall, namely $RI_{fac} = 1$ and $RI_{fac} = 1.5$. More intense rainfall leads to both increased water runoff and sediment entrainment, leading to greater flow volumes and increases in peak flow depth and area inundated.



Figure 5. Maximum flow depth for 50% ($RI_{fac}=0.5$), 100% ($RI_{fac}=1$), and 150% ($RI_{fac}=1.5$) scaling of the rainfall time series. Other parameters are set at their nominal value, except for hydraulic roughness and saturated hydraulic conductivity which were set to the minimal values in table 1, consistent with values anticipated immediately after the fire. For contrast, the maximum colorbar limit is set to 2.5 m although in the channels towards the north, maximum flow depth exceeds 2.5m.

4.3 Effects of postfire recovery on runout

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Liu et al. (2021) developed parametric best-fit curves to model the change in sat-465 urated hydraulic conductivity and hydraulic roughness as a function of time following 466 fire in the nearby San Gabriel Mountains. Using these relationships, and setting others 467 to the center of their respective ranges, we use the GP emulator to explore the effects 468 of temporal changes in the hydraulic roughness coefficient and saturated hydraulic con-469 ductivity. Both peak flow depth and area inundated in response to the observed rain-470 storm would decrease substantially over the first six months following fire (Figure 6). For 471 example, US Highway 101, which runs perpendicular to the direction of flow near the 472 distal portion of the fan, would only be inundated when the rainstorm occurs within the 473 first 3 months following fire. If the observed rainstorm were to have occurred 12 months 474 following the fire, the simulated inundation area would be limited to channels near the 475 fan apex. 476

We can also explore the effects of rainfall intensity and temporal changes in hydraulic roughness and saturated hydraulic conductivity following fire by examining flow depth

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Figure 6. Maximum inundation for several values of the Manning coefficient and saturated hydraulic conductivity along with the corresponding time from Figures 7b and 7c of Liu et al. (2021), respectively. All of the other parameters are set to the center value of their range. For contrast, the max colorbar limit is set to 2.5 m although in the channels towards the north, max flow depth exceeds 2.5m.

at distinct points of interest. Again, we consider the same two points for illustrative pur-479 poses, one located in a channel and one on the adjacent fan surface where flow is rela-480 tively unconfined (Figure 7). For a given time since fire, peak flow depths are greater 481 in the channel relative to on the fan surface, as expected. Peak flow depth decreases grad-482 ually over the first several months at the point on the fan before dropping to near zero 483 after approximately six months. Peak flow depths decrease over the first year following 484 fire in the channel location from roughly 2.5 m to 1.5 m. Visualizing peak flow depths 485 as a function of time since fire and rainfall intensity can be helpful for assessing tempo-486 ral shifts in the magnitude of rainfall associated with potential debris-flow impacts at 487 different locations. For example, even a rainstorm characterized by $RI_{fac} = 1.5$ would 488 not result in peak flow depths greater than 20 cm after approximately 0.6 years follow-489 ing fire. 490

We further use the emulator to produce probabilistic maps of inundation at different times following fire (Figure 8). Differences in the spatial patterns of inundation



Figure 7. Exploration at one location on the fan (top row) and one location in the channel (bottom row). Panels in column (a) indicate locations of all emulated flow depths (black) and those being explored in detail (red). Panels in column (b) show peak flow depth as a function of time since fire. The hydraulic roughness coefficient and saturated hydraulic conductivity are parameterized as a function of time since (Liu et al., 2021) while all other parameters set at their central values. (Note that the vertical scales are different; the maximum flow depth on the fan is roughly 1.0 m, and that in the channel is roughly 2.5 m.) Panels in column (c) show peak flow depth versus rainfall intensity with the hydraulic roughness coefficient and saturated hydraulic conductivity set to their respectively minimum values (i.e. as would be expected prior to any recovery) and with all other parameters set at their central values. Panels in column (d) contain color maps for maximum flow depth at these two locations varying all combinations of rainfall intensity and time (via Manning coefficient and saturated hydraulic conductivity). The white contours indicate the values of time and rainfall intensity leading to an inundation of 10 cm or more.

- ⁴⁹³ likelihood are apparent between scenarios where the storm occurs 2 months following the
- fire versus 14 months. We identify a location as being inundated if peak flow depth ex-
- $_{495}$ ceeds > 0.1m. All parameters were set to their central values except for saturated hy-
- ⁴⁹⁶ draulic conductivity and the hydraulic roughness coefficient. The latter parameter is sam-
- ⁴⁹⁷ pled from the distribution suggested by Liu et al. (2021) while the former is set to 2 and
- ⁴⁹⁸ 14 month values, respectively, estimated from the same study. The probability MC cal-

culation was carried out with 100 samples.



Figure 8. Probability of inundation maps, where a location is considered inundated if the maximum flow depth exceeds 10 cm. To calculate this probability, all parameters except the Manning coefficient and saturated hydraulic conductivity are set to their central values with the latter set to values corresponding to the 2 month and 14 month estimates from Liu et al. (2021), respectively.

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500 5 Discussion

Fire impacts on soil and vegetation properties that affect the initiation and growth of runoff-generated debris flows are most extreme in the first few months following fire (DeGraff et al., 2015; Thomas et al., 2021). Potentially rapid changes in hydrologic conditions following fire limit the time window for gathering data needed to constrain pa-

rameters for postfire runoff and erosion models, including the model used here. Aside 505 from rainfall intensity, which will not be affected by the fire, we found that hydraulic rough-506 ness, the representative grain size, and saturated hydraulic conductivity played the most 507 important roles in controlling debris flow inundation. Additional model testing across 508 fire-prone regions in different geologic and climate settings is needed to assess model per-509 formance and determine the extent to which results related to parameter sensitivity are 510 generalizable. Nonetheless, this result provides observational targets that can help fo-511 cus future efforts to collect perishable postfire data. 512

We hypothesize that hydraulic roughness plays an important role in controlling in-513 undated area and peak flow depths because of its influence on both modeled sediment 514 detachment rates and flow resistance. Saturated hydraulic conductivity will influence the 515 rate at which sediment is detached by overland flow since it exerts a strong control on 516 the magnitude of infiltration-excess overland flow that often dominates in postfire set-517 tings. Increased rates of sediment detachment lead to increases in flow volume, which 518 in turn acts to increase runout and inundation potential (Barnhart et al., 2021). Grain 519 size similarly influences flow volume since a larger grain size will encourage more rapid 520 deposition of sediment. 521

Our evaluation of parameter sensitivity indicates that constraints on postfire val-522 ues for hydraulic roughness, saturated hydraulic conductivity, and the grain size distri-523 bution of sediment entrained in debris flows would be particularly beneficial for improv-524 ing estimates of debris flow runout estimates. Burn severity is likely to play a substan-525 tial role in a fire's effect on these variables (Moody et al., 2015; McGuire & Youberg, 2020). 526 In addition, attempts to capture changes in debris flow runout as a function of time since 527 fire would benefit from methods to parameterize temporal changes in hydraulic rough-528 ness, saturated hydraulic conductivity, and the grain size distribution of sediment en-529 trained in debris flows. Fire-driven reductions in hydraulic roughness are commonly cited 530 as a cause for increased runoff and erosion (McGuire & Youberg, 2020; Stoof et al., 2015), 531 but there are few constraints on the temporal changes in hydraulic roughness following 532 fire, which may be facilitated by changes in vegetation cover and/or grain roughness. Par-533 ticularly in southern California (Doehring, 1968; Florsheim et al., 1991; DiBiase & Lamb, 534 2020) and other tectonically active regions in the western USA (Roering & Gerber, 2005), 535 fire can promote substantial increases in dry ravel activity on hillslopes that may reduce 536 hydraulic roughness by increasing the availability of fine sediment in channels. Hydraulic 537

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roughness may then increase over time as dry ravel deposits are progressively eroded dur-538 ing postfire rainstorms (Tang et al., 2019). Temporal changes in debris flow sediment 539 source locations (Guilinger et al., 2020) and coarsening of particle size distributions due 540 to preferential erosion of fines would also influence the effective grain size in the model. 541 In practice, it is not clear how to quantitatively connect this single grain size parame-542 ter to the particle size distribution of hillslope or channel sediment, especially when flows 543 contain boulders. Postfire changes in saturated hydraulic conductivity can be inferred 544 from calibration of hydrologic models (Liu et al., 2021), rainfall simulator experiments 545 at the small plot scale (Robichaud et al., 2016), and point scale measurements (Ebel, 2020; 546 Ebel et al., 2022; Perkins et al., 2022). While some general patterns have been observed 547 between time since fire and values of saturated hydraulic conductivity, there is substan-548 tial site-to-site variability (Ebel & Martin, 2017). The level of uncertainty in influential 549 model input parameters and how they change over time highlights the need for proba-550 bilistic assessments of debris flow runout, which emulators can help to achieve by facil-551 itating rapid exploration of large parameter spaces. 552

Rainfall is a necessary driver for debris flow initiation and the model was also sen-553 sitive to rainfall intensity, specifically a rainfall intensity factor which we used to scale 554 the rainfall intensity time series. This finding is consistent with observations that post-555 fire basin-scale sediment yields (Pak & Lee, 2008) and debris flow volume (Gartner et 556 al., 2014) increase with rainfall intensity averaged over durations of 60 minutes or less. 557 Short duration (sub-hourly) bursts of high intensity rainfall are effective at generating 558 infiltration-excess overland flow that can trigger debris flows in recently burned steep-559 lands (Kean et al., 2011; Nyman et al., 2011; Esposito et al., 2023). Emulators may be 560 useful for generating probabilistic maps of debris flow inundation in response to design 561 storms with different rainfall intensities or examining changes at particular points of in-562 terest. In cases where there are specific values at risk downstream of a burned area, rapid 563 exploration of debris flow characteristics (i.e. peak flow depth) as a function of rainfall 564 intensity could help define impact-based rainfall thresholds that could be used for plan-565 ning and warning purposes. In other words, one could take advantage of the emulator's 566 computational efficiency to determine, not only the rainfall intensity required to initi-567 ate a debris flow but also the rainfall intensity required to produce a debris flow that would 568 impact a prescribed area of interest with some prescribed depth of flow. 569

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The computational cost of many physically-based debris flow models is a limita-570 tion in applications that are time sensitive, such as rapid postfire hazard assessments. 571 Postfire debris flows in the western USA, such as those that occurred near Montecito, 572 may occur before the fire has been officially contained and within weeks or months of 573 fire ignition. The emulator methodology presented here provides one avenue for mini-574 mizing computation times, since an initial suite of simulations can be used to train the 575 emulator which can later be applied with substantially less computational effort to gen-576 erate a probabilistic hazard map for a specific scenario. An emulator may even be trained 577 prior to a fire. Analogous approaches have been employed in related applications (Rutarindwa 578 et al., 2019; Spiller et al., 2020). Within the context of postfire hazards, an emulator could 579 be used to assess debris-flow runout and inundation downstream of a burned area in re-580 sponse to a design or forecast rainstorm. Atmospheric model ensembles, for example, can 581 provide estimates of peak 15-minute rainfall intensity over watersheds of interest that 582 could be used to constrain a distribution of rainfall intensity factors (Oakley et al., 2023). 583

584 6 Conclusions

We develop a physics based high fidelity computationally expensive morphodynamic 585 model and cost effective surrogates based on Gaussian process models of postfire debris 586 flows. We employ the Gaussian Process surrogate model, or emulator, to approximate 587 peak debris flow depth from a physics-based morphodynamic model, Titan2D. The em-588 ulator is able to approximate the peak flow depth with a mean squared error that is gen-589 erally in the range of 0.1-0.2 m when using a modest training data set built from 64 590 Titan2D simulations. The range parameters associated with the emulator provide a met-591 ric for the relative importance of input parameters, which provides guidance for those 592 that are most important to constrain for forward modeling of debris flow runout. We find 593 that peak flow depths are most sensitive to changes in hydraulic roughness and a rain-594 fall intensity factor and are moderately sensitive to saturated hydraulic conductivity and 595 effective grain size. We highlight the emulator's ability to provide rapid estimates of peak 596 flow depth for parameter combinations that were not part of the training data set by gen-597 erating probabilistic maps of inundation as a function of time since fire. Inundation like-598 lihood changes substantially over the first year following the fire, driven by temporal vari-599 ations in hydraulic roughness and saturated hydraulic conductivity. Emulator-based anal-600 yses can facilitate rapid Monte Carlo calculations of inundation probability, making them 601

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a promising option for rapid postfire hazard assessments and scenario planning beforea fire starts.

⁶⁰⁴ Data Availability Statement

The debris flow model under consideration in this paper is from (McGuire et al., 2017) and it is accelerated by implementation in the Titan2D platform (Patra et al., 2005; Simakov et al., 2019). Parametric models of the Manning coefficient and saturated hydraulic conductivity versus time are available from (Liu et al., 2021) as are validated samples of those same parameters for debris flows 2 and 14 months post fire. Packages to implement the parallel partial emulator (Gu & Berger, 2016) are available in (Gu et al., 2019).

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