Controls on upper ocean salinity variability in the eastern subpolar North Atlantic during 1992-2017

Ali Hasan Siddiqui¹, Thomas W N Haine¹, An T Nguyen², and Martha Weaver Buckley³

¹Johns Hopkins University ²University of Texas-Austin ³George Mason University

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Abstract

The eastern subpolar North Atlantic upper ocean salinity undergoes decadal fluctuations. A large fresh anomaly event occurred during 2012–2016. Using the ECCO state estimate, we diagnose and compare mechanisms of this low salinity event with that of the 1990s fresh anomaly event. To avoid erroneous interpretations of physical mechanisms due to reference salinity values in the freshwater budget, we perform a salt mass content budget analysis of the eastern subpolar North Atlantic. It shows that the recent fresh anomaly occurs due to the circulation of anomalous salinity by mean currents entering the eastern subpolar basin from its western boundary via the North Atlantic Current. This is in contrast to the early 1990s, when the dominant mechanism governing the fresh anomaly was the transport of the mean salinity field by anomalous currents across the southern boundary of the subpolar North Atlantic.

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Ali H. Siddiqui¹, Thomas W. N. Haine¹, An T. Nguyen², Martha Buckley³

¹Department of Earth and Planetary Sciences, Johns Hopkins University, Baltimore, MD, USA ²Oden Institute for Computational Engineering and Sciences, University of Texas at Austin, Austin, TX, USA

³Atmospheric, Oceanic and Earth Sciences, George Mason University, Fairfax, VA, USA

Key Points:

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- Two fresh anomalies observed in the eastern subpolar North Atlantic upper ocean during 1992–2017 share similar spatial characteristics.
 Salt budget analysis shows the 2012–2016 fresh anomaly in the upper 200 m occurs due to transport of anomalous salinity by mean currents.
 In contrast, the fresh anomaly in the 1990s is due to anomalous circulation of the
 - In contrast, the fresh anomaly in the 1990s is due to anomalous circulation of the mean salinity field.

Corresponding author: Ali H. Siddiqui, asiddi24@jhu.edu

15 Abstract

The eastern subpolar North Atlantic upper ocean salinity undergoes decadal fluctuations. 16 A large fresh anomaly event occurred during 2012–2016. Using the ECCO state estimate, 17 we diagnose and compare mechanisms of this low salinity event with that of the 1990s 18 fresh anomaly event. To avoid erroneous interpretations of physical mechanisms due to 19 reference salinity values in the freshwater budget, we perform a salt mass content bud-20 get analysis of the eastern subpolar North Atlantic. It shows that the recent fresh anomaly 21 occurs due to the circulation of anomalous salinity by mean currents entering the east-22 ern subpolar basin from its western boundary via the North Atlantic Current. This is 23 in contrast to the early 1990s, when the dominant mechanism governing the fresh anomaly 24 was the transport of the mean salinity field by anomalous currents across the southern 25

²⁶ boundary of the subpolar North Atlantic.

27 Plain Language Summary

On decadal time scales, the eastern subpolar North Atlantic shifts between a salty 28 and fresh upper ocean. These changes are significant and need to be investigated because 29 of their impacts on global ocean circulation and regional ocean biogeochemistry. Between 30 2012 and 2016, there was a large event where this region became fresher than it had been 31 in over a century. A previous similar event occurred in the early 1990s, but with a smaller 32 magnitude. We use a numerical model of the ocean to figure out why these events oc-33 curred. Our study shows that there were two different mechanisms at play. The recent 34 event occurred because a lot of fresh water came in from the west by the mean currents. 35 The 1990s event occurred because ocean currents shifted and brought fresh water from 36 the south. 37

³⁸ 1 Introduction

Large scale low salinity events occur in the eastern subpolar North Atlantic Ocean (ESNA) on decadal time scales. Based on observations, the subpolar North Atlantic has been undergoing such decadal salinity changes since at least the early 20th century (Sundby & Drinkwater, 2007; Dickson et al., 1988; Dooley et al., 1984; R. Zhang & Vallis, 2006; Dickson et al., 1988; Belkin et al., 1998; Belkin, 2004). During the 1992–2017 period, there were two fresh anomaly events in the ESNA reaching maximum freshwater accumulation in 1995 and 2016 respectively.

We highlight previous discussion in the literature on mechanisms that control the salinity in the ESNA. Changes in the strength and size of the subpolar gyre (SPG) and local atmospheric forcing in the ESNA are proposed mechanisms involving processes local to the subpolar gyre (Fox et al., 2022; Holliday et al., 2020). Remote mechanisms to modulate the salinity in the ESNA on such time scales include advection of salt anomalies from the Arctic or the subtropics (J. Zhang et al., 2021; Yeager et al., 2012; Häkkinen et al., 2011; Thierry et al., 2008; Sundby & Drinkwater, 2007; Holliday, 2003).

The strength and size of the subpolar gyre has been hypothesized to play an im-53 portant role in setting the salinity variability in the ESNA (Holliday, 2003; Hátún et al., 54 2005; Häkkinen & Rhines, 2004; Sarafanov et al., 2008; Yeager et al., 2012; Häkkinen et 55 al., 2011; Thierry et al., 2008), especially in the context of the warming and salinifica-56 tion that occurred in the mid 1990s to 2000s. The expansion of the SPG reduces the con-57 tribution of salty subtropical waters to the ESNA, reducing the salinity; in 1994 and 2016 58 sea surface height (SSH) contours show an expanded subpolar gyre and fresh anomalies 59 in the ESNA (Fig. 1). In contrast, the contraction of the SPG allows more subtropical 60 waters into the ESNA, increasing salinity; in 2008 the subpolar gyre is contracted, as SSH 61 contours retreat westward, and the ESNA is saltier (Fig. 1). 62

Subsequent studies have refined diagnostics for studying the relationship between 63 the subpolar gyre strength and ESNA salinity; rather than considering SSH-based in-64 dices of the subpolar gyre strength ((e.g., Häkkinen & Rhines, 2004)), Tesdal et al. (2018) 65 analyze a density-based gyre index, which is a proxy for the baroclinic strength of the 66 gyre (Koul et al., 2020). Foukal and Lozier (2017, 2018) suggest that the salinity in the 67 ESNA is more strongly influenced by the intergy transport, which is modulated by the 68 Atlantic meridional overturning circulation (AMOC). Koul et al. (2020) perform Lagrangian 69 tracking experiments based on multiple definitions of the SPG strength during 1993–2016 70 and conclude that 64.8% virtual floats reaching ESNA originate from subtropical wa-71 ters. Contributions from subpolar-sourced waters increase five-fold during an expanded 72 state of the SPG, however. 73

Another method to diagnose mechanisms controlling temperature and salinity vari-74 ability is by performing budget calculations. To address the decadal SST variability in 75 the subpolar North Atlantic, Piecuch et al. (2017) calculated the heat budget for 46° -76 65° N and concluded that the warming in the late 1990s and subsequent cooling since 2008 77 is primarily driven by oceanic advective flux convergence. The anomalous flux conver-78 gence is dominated by anomalies across the southern boundary $(46^{\circ}N)$. Similar studies 79 by Oldenburg et al. (2018) and Tesdal and Haine (2020) reach the same conclusion on 80 the dominance of the southern boundary advection in setting subpolar North Atlantic 81 heat and freshwater variability. However, the role of the AMOC is unclear in this mech-82 anism. Similarly, Sanders et al. (2022) investigate the 2015 anomalous cooling in the east-83 ern and central subpolar region (defined over 50–20°W, 43–63°N) using a mixed layer 84 heat budget. They observe that surface heat loss initiates and drives the cooling, with 85 advection sustaining the anomaly in the region (as expected from Tesdal and Abernathey 86 (2021)). They also emphasize the role of vertical diffusion across the base of the mixed 87 layer in the re-emergence of the anomaly during summer of 2014. 88

⁸⁹ Bryden et al. (2020) observed that there has been a mean increase of 0.12 ± 0.04 ⁹⁰ Sv of freshwater flux (relative to a reference salinity of 35.17 psu) after 2010 compared ⁹¹ to before 2009. This increase is about 10% of the 2004–2009 average flux. They propose ⁹² that a freshening of 0.062 ± 0.013 Sv over the eastern subpolar gyre during 2014–16 rel-⁹³ ative to 2007–09 is primarily due to the reduction of the AMOC by 2.5 Sv after 2009.

So far, we have discussed how circulation variability within the Atlantic basin con-94 tributes to salinity anomalies in the ESNA. For the 2012–2017 freshening event, Holliday 95 et al. (2020) describe the primary mechanism of net freshwater gain in the upper 200 m 96 of the Iceland basin as the rerouting of Arctic-sourced Labrador Current water into the 97 northern branch of the North Atlantic Current (NAC; Reverdin et al., 2003). It is mod-98 ulated by changes in the SPG strength driven by changes in atmospheric forcing. Re-99 cently, Fox et al. (2022) highlighted that reduced surface heat loss led to an increase in 100 warmer (less dense) waters in the Labrador Sea. The transport of these less dense wa-101 ters from the upper ocean layers through the Labrador Current along with reduced vol-102 ume transport from the Gulf Stream drove the cooling and freshening in the eastern sub-103 polar region. 104

However, some studies suggest an important role for interactions between the sub-105 polar North Atlantic and the Arctic. J. Zhang et al. (2021) and Sundby and Drinkwa-106 ter (2007) attribute ESNA freshening events during 1983–1995 and 1947–2000 to the ex-107 port of freshwater buildup in the Arctic. They suggest that sea ice and liquid freshwa-108 ter anomalies travel via the Fram Strait and Davis Strait to the Labrador Sea and cir-109 culate around the eastern subpolar gyre in the North Atlantic Current. The proposed 110 111 mechanism of freshwater buildup in the 1990s is increased freshwater flux from the Davis Strait (Belkin, 2004), which entered the Labrador Sea and propagated around the east-112 ern subpolar gyre (Sundby & Drinkwater, 2007). 113



Figure 1: Annually-averaged subpolar North Atlantic upper-ocean salinity (0–200 m, colors) and sea-surface height (SSH; contours) averaged over one year preceding the salinity field. The SSH field is from the AVISO dataset and the salinity field is from the EN4 product. Following Chafik et al. (2019), the grey contours range from -0.8 m to 0.8 m with a spacing of 0.1 m and represent the mean dynamic topography (CNES-CLS2013 MDT). The red contours are -0.3, -0.2, -0.1 m and represent the three branches of the NAC. A Gaussian filter is used to smooth the SSH field with a scale of 1.25°. Modified from Fig. 3 of Weijer et al. (2022).

In this paper, we focus on what sets the upper ocean salinity in the ESNA on decadal time scales with emphasis on the two recent freshening events in the 1990s and 2010s using observations and modelling tools. We look at salt content anomaly budgets to explore oceanic mechanisms and further investigate the contribution of surface freshwater forcing in setting upper ocean salinity in the region.

In section 2 we discuss the methods used in this paper and investigate upper ocean salinity variability as seen in observations. In section 3, we compare and contrast the two fresh anomaly events observed during 1992–2017 using observations and ocean state estimates. We then diagnose the salinity variability using salt budget analysis for the ESNA and discuss potential mechanisms for the salinity variability.

124 2 Methods

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2.1 Datasets

The main tool for our analysis of upper ocean salinity is the ECCO (Estimating 126 the Circulation and Climate of the Ocean) version 4 release 4 (Forget et al., 2015) and 127 ASTE (The Arctic Subpolar Gyre sTate Estimate) release 1 (Nguyen et al., 2021) ocean 128 state estimates. The ECCOv4r4 state estimate (ECCO Consortium et al., n.d.) is a dynamically-129 consistent, data-constrained solution of the MITgcm model for the period 1992–2017. 130 This allows for the construction of realistic closed budgets of volume, heat and salt. The 131 horizontal resolution is 1°. ASTE_R1 is also a data-constrained ocean-sea ice model-data 132 output that covers the Arctic Ocean and Atlantic ocean with lateral boundaries at 47.5° N 133 in the North Pacific and 32.5°S in the South Atlantic for 2002–2017; the ASTE_R1 hor-134 izontal resolution is $1/3^{\circ}$. 135

We utilize a number of observational datasets to evaluate ECCOv4r4 and ASTE_R1. First, we compare to the EN4 hydrographic dataset (Good et al., 2013), which is an observational product compiled by the UK Met office. We use data from two hydrographic surveys to evaluate the performance of ECCOv4r4 for our analysis. The Extended EL-LET line (Holliday & Cunningham, 2013) is a section from Iceland to Scotland and the OVIDE line (Daniault et al., 2016) is a combination of sections from the southern tip

of Greenland to Portugal. These are the same sections used by Holliday et al. (2020). 142 We compare the salinity anomalies in the sections observed in the model data with the 143 ELLET and OVIDE sections. The data from the hydrographic survey carried out over 144 June–July, 2016 and May–June, 2016 for the OVIDE and ELLET line, respectively, is 145 compared with monthly-mean salinity anomaly from the model data. We consider the 146 upper 200 m annually-averaged ESNA salinity anomaly for 1992–2017 using the EN4 hy-147 drographic dataset (Good et al., 2013) along with ECCOv4r4 and ASTE_R1 datasets. 148 The upper ocean 200 m is used to remain consistent with previous work that investigates 149 upper ocean salinity in the ESNA (Holliday et al., 2020; Koul et al., 2020). 150

We use ECCOv4r4 to track the evolution of salinity anomalies for the two fresh anomaly events on monthly and also annual time scales. To investigate the mechanisms governing this evolution, we carry out salt budget analysis with ECCOv4r4 for the period 1992– 2017 for the entire subpolar North Atlantic (SPNA) and the ESNA. The salt budget analysis builds on the work by Buckley et al. (2014, 2015); Piecuch et al. (2017); Oldenburg et al. (2018); Tesdal and Haine (2020); Nguyen et al. (2021) to investigate heat and salinity variability in the subpolar North Atlantic.

The ESNA is defined as a box over $10-30^{\circ}$ W, $46-65^{\circ}$ N and the SPNA is defined as the North Atlantic between $45^{\circ}-65^{\circ}$ N. Fig. 2 shows the location of the ESNA.

160 3 Results

In this section, we evaluate the ECCOv4r4 model data. Next, we consider the spatial and temporal structure of the fresh anomaly events in the subpolar region. Then we investigate the salt budget results.

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3.1 Evaluation of ECCOv4r4

Here we present a comparison of ECCOv4r4 to ship-based hydrographic sections. 165 Fig. 2a shows the positions of the OVIDE and ELLET sections overlaid with sea level 166 height anomaly contours from ECCO. In summer of 2016, the upper 1000 m of the Rock-167 all Trough has the highest salinity in the ELLET section. The salinity decreases below 168 1000m, which is seen in both the observations and ECCO data (Fig. 3). ECCO overes-169 timates the salinity in the upper 1000 m, however. This is also true on the OVIDE sec-170 tion from the Reykjanes Ridge to the Iberian abyssal plain during summer of 2016. The 171 ECCO salinity in the Irminger Sea during summer of 2016 is realistic, however. ECCO 172 also captures the sub-surface salinity minimum over the Iberian abyssal plain. On the 173 ELLET section, ECCO overestimates the salinity in the 0–200 m Iceland Basin by around 174 0.06 psu, and underestimates it by 0.04 psu in the Rockall Trough. Similarly, on the OVIDE 175 section, ECCO overestimates the 0–200 m salinity in the ESNA control volume (Fig. 2) 176 by around 0.05 psu. The time series (Fig. 2) of 0-200 m depth-averaged salinity anoma-177 lies for the ESNA in ECCO accurately shows the fresh anomaly observed during 2012-178 2017 and 1992–1995 compared to EN4. The salinity difference between ECCO and EN4 179 is 0.027 ± 0.046 psu (from Fig. 2b, but using the absolute salinity values, not the anoma-180 lies). The biases in the ELLET and OVIDE sections quoted above are consistent with 181 this ECCO/EN4 difference, and are relatively small compared to the salinity fluctuations 182 over time seen in Fig. 2b. This builds confidence in the use of ECCO for our analysis. 183 The only noticeable disagreement between EN4 and ECCO seen in Fig. 2b is during 1995-184 1996, when the ESNA shows anomalous positive salinity anomalies in EN4, whereas ECCO 185 shows negative salinity anomalies. Spatial maps of 0–200 m salinity anomalies show a 186 large positive anomaly situated south of the Grand Banks in 1995 in both datasets (see 187 Supplemental Figs. 1–3). In EN4 this anomaly spreads throughout the ENSA in 1996, 188 but it does not spread so far east in ECCO. 189



Figure 2: (a) Subpolar North Atlantic (SPNA) sea level height anomaly contours (spacing of 0.04 m) in the ECCOv4r4 dataset averaged over 1992–2017. The Extended ELLET line (EEL) and OVIDE sections are also shown. (b) Upper ocean 200 m salinity anomaly time series for the eastern Subpolar North Atlantic (ESNA) from EN4 (red), ECCOv4r4 (black) and ASTE_R1 (grey) datasets. The ESNA is defined as 45–65°N, 10–30°W (purple box in (a)).

3.2 Spatial and Temporal Structure of Upper Ocean (S_{200}) Salinity Anomalies

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We first establish the occurrence of two fresh anomalies in the upper 200 m of the 192 ESNA using the ECCOv4r4, ASTE_R1, and EN4 datasets (Fig. 2). We observe a fresh 193 anomaly in the ESNA in the early 1990s, after which there is a prolonged period of salin-194 ification until 2011, and a reversal to freshening thereafter. The first fresh anomaly event 195 (F_1) is observed from 1992 until 1995 in EN4 and 1997 in ECCOv4r4. For the second 196 fresh anomaly event (F_2) , all three datasets show the Iceland basin salinity anomaly drop-197 ping below zero after 2012 until 2017. This is also reflected in the decadal cycle in the 198 upper ocean salinity maps for the ESNA region in the ECCO, ASTE_1 and EN4 datasets 199

(see Supplemental Fig. 1–3). We label the 1990s fresh event as F_1 and the 2010s fresh 200

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event as F_2 .

-2000

-3000

ECCO

0

200

 $\dot{400}$

600

Distance from first station [km]

(b)

800

1000

1200



Figure 3: Comparison of ECCOv4r4 salinity for (a) June–July 2016 with the OVIDE section and (b) May–June 2016 with the ELLET line. The ECCOv4r4 salinity sections are taken at the same times as the field observations. Vertical purple lines indicate parts of the sections inside the ESNA control volume defined in Fig. 2 (for the OVIDE section 700-2290 km, and for the ELLET section 0-1150 km, are within the ESNA control volume). Colorbar limits are different for the two section plots.

Spatial trends in the upper ocean (S_{200}) salinity are computed using EN4 and EC-202 COv4r4. During 2005–16, both products show a statistically significant freshening in the 203 ESNA at 95% confidence intervals using the student's t-distribution. ECCOv4r4 and EN4 204 disagree on trends in the Labrador Sea and the Grand Banks region, however (Fig. 4). 205

Next, we consider the annually averaged anomalies in the ECCO data and com-206 pare the two fresh anomaly events, F_1 and F_2 , which are shown in Fig. 5. As we trace 207 the freshwater event in the 1990s (F_1) we observe a fresh ESNA and saltier western sub-208 polar gyre (SPG) in 1992. During 1993 and 1994, a fresh anomaly is situated in the Labrador 209 Sea, and in the Iceland Basin in 1995. By 1996, the signal fades away from the ESNA. 210



Figure 4: Spatial distribution of linear trends in the upper-ocean salinity (0–200 m) over 2005–2016 using monthly mean fields of (a) EN4 and (b) ECCOv4r4. Dotted regions display insignificant trends calculated using the students-t test with a p-value of 0.05.

The 2010s event, has a similar fresh anomaly in the Labrador Sea in 2013, and also in the Iceland Basin in 2016. Note that in both events, there is a positive salinity anomaly

south of the Grand Banks region, preceding the maximum freshening in 1995 and 2016.

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Figure 5: Annually averaged anomaly maps for the upper 200 m salinity (S_{200}) during 1992–1994 and 2012–2014 in ECCOv4r4.

We investigate this further using the ECCO data by tracking salinity anomalies along 214 the western SPG and along the Gulf Stream with the Iceland Basin as a common ter-215 minus. We create a section following mean sea level anomaly contours around the sub-216 polar gyre which begins south of Denmark strait. The section follows the -0.6 m mean 217 sea level anomaly contour along the East Greenland Current around the southeast coast 218 of Greenland. The section continues along the West Greenland Current to the entrance 219 of Baffin Bay, where it retroflects and follows the Baffin Island Current, eventually reach-220 ing the Labrador Sea. At this point, the section follows the -0.8 m mean sea level anomaly 221 contour as it retroflects east of the Flemish Cap and follows the path of the northern branch 222 of the NAC to the Iceland basin. This section represents a potential subpolar pathway 223 for the anomaly propagation. 224

To investigate the salinity anomalies along a potential subtropical pathway, we create a section along the Gulf Stream which also terminates in the Iceland Basin. Both sections are inspired by the Lagrangian studies carried out by Burkholder and Lozier (2014);

Foukal and Lozier (2018); Koul et al. (2020); J. Zhang et al. (2021) where passive trac-

ers are tracked to the Iceland basin in a variety of experiments.



Figure 6: (a) Continuous sections along the western subpolar gyre boundary (purple; 56 stations) and along the Gulf Stream (red; 28 stations). The sections intersect at the NAC (station A1) and are merged from there to the Iceland Basin. (b) Hovmöller diagram of monthly salinity anomalies in the ECCO data along the sections. Vertical dashed lines are shown at Davis Strait (DS) and the Grand Banks (GB) region.

We first inspect the subtropical section. We find that there is a general tendency for there to be salinity anomalies of opposite sign along the Gulf Stream path and in the subpolar gyre, as noted by prior studies (Buckley et al., 2014; Joyce & Zhang, 2010; Sanchez-

Franks & Zhang, 2015; Yeager, 2015; Hátún et al., 2009; R. Zhang, 2008; Nye et al., 2011; 233 Yan et al., 2017, 2018). For both subpolar freshening events, there are positive salinity 234 anomalies along the Gulf Stream path. This positive salinity anomaly signal south of the 235 Grand Banks is also observed in the years preceding the fresh anomaly event in the an-236 nually averaged salinity anomaly maps (Fig. 5). Along the subpolar section, high fre-237 quency freshening/salinification events can be tracked from Denmark Strait along the 238 East Greenland Current and West Greenland Current. These represent seasonal fluctu-239 ations along the Greenland coast. Along the Labrador Current (from Davis Strait to the 240 Grand Banks), the characteristics of the freshening/salinification events changes with 241 lower-frequency variations than seen along the East and West Greenland Currents. From 242 the Grand Banks to the Iceland Basin, the freshening/salinification events are at even 243 lower frequencies. The relationship between salinity anomalies in the Icelandic Basin and 244 those in the East/West Greenland Current and Labrador Current is complex, with some 245 indications of signal propagation along the subpolar gye pathways, as suggested by Holliday 246 et al. (2020) and Fox et al. (2022). We can now explore the role of circulation changes 247 quantitatively by constructing a salt mass budget for the region. 248

3.3 Salt Budget Analysis

We construct a budget of the salt mass, an extensive quantity, which has a more accurate closure in the ECCO model output than a budget of salinity, an intensive quantity. The role of surface freshwater in the form of P-E+R (Precipitation-Evaporation+Runoff) is analyzed separately in a salinity budget calculation. We also avoid analyzing freshwater budgets due to ambiguities associated with reference salinity in such calculations (Schauer & Losch, 2019).

The salt conservation equation for the non-linear free surface in ECCOv4r4 is expressed in z^* coordinates (see equation (3) in Forget et al. (2015)). In z^* coordinates, sea surface height variations, η , are proportionally divided between ocean layers: $z^* = (z - \eta)/(H + \eta)$ (equation (1) in Piecuch (2017)), where z is the fixed vertical coordinate and H is the ocean depth.

We express the volume and time integrated salt content, M(t), for the control-volume V as

$$M(t) \equiv \underbrace{\rho_0 \int_{t^*=0}^{t^*=t} \int_V \frac{\partial(s^*S)}{\partial t} dV^* dt^*}_{\text{Salt Mass}} = \underbrace{\rho_0 \int_{t^*=0}^{t^*=t} \int_V -\nabla_{z^*} \cdot (s^*S\mathbf{v}_{res}) - \frac{\partial(Sw_{res})}{\partial z^*} dV^* dt^*}_{\text{Advection}} + \underbrace{\rho_0 \int_{t^*=0}^{t^*=t} \int_V s^*\mathcal{F}_S dV^* dt^*}_{\text{Forcing}} + \underbrace{\rho_0 \int_{t^*=0}^{t^*=t} \int_V s^*(D_{\sigma,S} + D_{\perp,S}) dV^* dt^*}_{\text{Diffusion}}.$$
(1)

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In this equation $s^{\star} = 1 + \eta/H$ is a scaling factor, $\nabla_{z^{\star}}$ indicates the gradient at con-267 stant z^{\star} , $(\mathbf{v}_{res}, w_{res})$ are the residual velocity fields defined as the sum of the Eulerian 268 and bolus velocities, \mathcal{F}_S is the forcing at the surface due to surface salt exchange due 269 to sea ice melting/formation and a redistribution of the surface flux in the vertical col-270 umn, and $D_{\sigma,S}$ and $D_{\perp,S}$ are diffusive processes parameterized along iso-neutral and ver-271 tical directions, respectively. We consider V to be the upper 200 m in the ESNA; we also 272 consider the upper 200 m for the whole SPNA (this assumption is tested below in sec-273 tion 3.4; see Appendix A for more details). Eq. (1) expresses that the total time inte-274 grated salt mass is balanced by the time integrated horizontal and vertical advective con-275 vergence of salt flux, diapycnal and isopycnal diffusion, and surface forcing. The four named 276 terms are each computed individually from the ECCO output, which allows us to test 277 the closure of (1). The ratio of the residual (left hand side minus right hand side) to the 278 salt mass (left hand side) is of $O(10^{-4})$. Details on how to close the salt budget in the 279 ECCO dataset are provided in Piecuch (2017). 280

We remove the mean seasonal cycle and a linear trend from all terms in the time 281 integrated salt budget. This gives us the salt content anomaly time series, which is shown 282 in Fig. 7 (labelled "salt anomaly"). We observe a negative salt content anomaly of the 283 SPNA during 1992–1997 and 2012–2017. For the ESNA, the salt content anomaly re-284 mains negative during 1992–1997 and 2012–2017. The salt mass anomaly increases and 285 reaches a maximum in 2004 for the entire SPNA, and in 2008 for the ESNA. We find that 286 both the diffusion and advection terms contribute to the salt content anomaly, and these 287 terms are strongly anti-correlated. The surface salt forcing (due to brine rejection) has 288 a negligible impact. 289

From 1995–2000 and 2004–2012, the change in the diffusion term is larger than the change in the advection term in the SPNA, and it drives the salt mass change. For the ESNA, this is seen during 1995–2000 and 2008–2011. We highlight the years in yellow/blue when the advection term increases/decreases rapidly in the two basins.



Figure 7: (a) Time and volume integrated salt anomaly budget for the upper 200 m of the (a) Subpolar North Atlantic and (b) the eastern subpolar North Atlantic (ESNA, see Fig. 2 for the definition of the ESNA). F_1 and F_2 are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

We further investigate the advection and diffusion terms separately. We decompose the advection term into time-average and anomaly terms, following Dong and Sutton (2002); Doney et al. (2007); Buckley et al. (2015); Piecuch et al. (2017) and Tesdal and Haine (2020). We write the advection term \mathcal{A} in (1), as $\mathcal{A} = \overline{\mathcal{A}} + \mathcal{A}'$, where the overbar denotes time averaging, i.e $\bar{\mathbf{v}} = \frac{1}{(t_f - t_i)} \int_{t_i}^{t_f} \mathbf{v} dt$ and the prime denotes departure from the time average. Initial and final time values in the averaging period are denoted by t_i and t_f . The averaging period is 1992–2017. This decomposes the advection term into variability produced by changes in the circulation $(\mathbf{v}'_{res}\bar{S})$, variability produced by changes in salinity $(\bar{\mathbf{v}}_{res}S')$, and that due to the co-variability of the circulation with the salinity $(\mathbf{v}'_{res}S' - \overline{\mathbf{v}'_{res}S'})$.

We now express $\mathbf{v}_{res} = \mathbf{v}_e + \mathbf{v}_b$, i.e., the total velocity is the sum of Eulerian (\mathbf{v}_e) and bolus (\mathbf{v}_b) velocities, and similarly for the vertical speeds (w_e, w_b) . Re-arranging the terms in Eq. (A2) gives:

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$$\mathcal{A}' = -\rho_0 \int^t \int_V \left\{ \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e S' + \frac{\partial(\bar{w}_e S')}{\partial z^\star} \right)}_{\bar{v}_e S'} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \mathbf{v}'_e \bar{S} + \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \mathbf{v}'_e \bar{S} + \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S$$

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$$\underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} (\mathbf{v}'_{e}S' - \overline{\mathbf{v}'_{e}S'}) + \frac{\partial (w'_{e}S' - \overline{w'_{e}S'})}{\partial z^{\star}} \end{pmatrix}}_{v'_{e}S' - \overline{v'_{e}S'}} \\ \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}S' + \frac{\partial (\overline{w}_{b}S')}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \mathbf{v}'_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot$$

$$\underbrace{\left(\nabla_{z^{\star}} \cdot s^{\star}(\mathbf{v}_{b}'S' - \overline{\mathbf{v}_{b}'S'}) + \frac{\partial(w_{b}'S' - \overline{w_{b}'S'})}{\partial z^{\star}}\right)}_{v_{b}'S' - \overline{v_{b}'S'}}\right\} dV^{\star} dt^{\star} + \epsilon. \quad (2)$$

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The sum of all the terms on the right hand side of Eq. (2) should equal the anomaly in the convergence of advective salt flux, which is an output in ECCOv4r4. When we compute the individual terms from monthly mean velocities and salinity (interpolated to the model velocity grid points), however, it introduces a numerical error, which we call ϵ . Volume integrated ratios of ϵ to each individual term are O(10⁻¹) (see Fig. 9).

The advection anomaly term is shown in Fig. 8. For the whole SPNA, during the 318 F_1 event, the variability in the advective convergence term is dominated by the anoma-319 lous Eulerian advection of mean salinity $(v'_e \bar{S})$. During 1992–1995, it is twice that of the 320 mean circulation of anomalous salinity $(\bar{v}_e S')$. For the F₂ event the mean bolus advec-321 tion of anomalous salinity $(\bar{v}_b S')$ tends to drive the freshening, as the Eulerian terms bal-322 ance each other out. Note that the advective convergences of anomalous circulation of 323 mean salinity and the mean circulation of anomalous salinity are anti-correlated during 324 both events. Anti-correlation of advective convergences due to salinity variations and due 325 to geostrophic velocity variations is expected when isobars and isohalines are aligned, 326 as shown for heat transport convergences by Buckley et al. (2015). 327

The decomposition of the anomalous advection term for the ESNA is different. The 328 F_1 negative salt anomaly is still dominated by the anomalous circulation of mean salin-329 ity $(v'_e S)$. In contrast, the driver for the F_2 anomaly is the mean circulation of anoma-330 lous salinity $(\bar{v}_e S')$. Also, unlike the SPNA, the anomalous circulation of mean salinity 331 is not anti-correlated to the mean circulation of anomalous salinity. This indicates that 332 either (1) the isobars and isotherms are not strongly aligned, which would occur if the 333 density field in the ESNA is only weakly dependent on salinity, or (2) there is a strong 334 contribution of agoestrophic transports to the advective convergences. For example, there 335 is no expected anticorrelation between salt transport convergences due to salinity vari-336 ations and those due to Ekman velocity variations. 337

To further investigate the $\bar{v}_e S'$ term, we look at the decomposed anomalous advection for the ESNA in terms of flux entering the control volume from the boundaries. We rewrite Eq. (A2), apply the Gauss-divergence theorem, and express the divergence of anomalous advection in the ESNA as a sum of salt fluxes across the four lateral boundaries and



Figure 8: Decomposition of the anomalous advection term \mathcal{A}' in the salt budget for the upper 200 m of the (a) Subpolar North Atlantic and (b) Eastern subpolar North Atlantic (ESNA, see Fig. 2 for the definition of the ESNA). The anomalous advection term is decomposed into contributions from: changes in the circulation $(v'_{res}\bar{S})$, changes in the salinity $(\bar{v}_{res}S')$, and changes due to the co-variability of the circulation with the salinity $(v'_{res}S' - v'_{res}S')$. This decomposition is made for both Eulerian and bolus velocity components. See Eq. (2) for details on the terms, and see also Fig. 7 for information on the full salt budget. F_1 and F_2 are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

across the 200 m vertical boundary: 342

 $\mathcal{A}' = -\rho_0 \int^t \int_B [\underbrace{\mathbf{s}^* \bar{\mathbf{v}}_{res} S' + \bar{w}_{res} S'}_{\bar{v}S'} + \underbrace{\mathbf{s}^* \mathbf{v}'_{res} \bar{S} + w'_{res} \bar{S}}_{\mathbf{v}'\bar{S}} + \underbrace{\mathbf{s}^* (\mathbf{v}'_{res} S' - \frac{v'\bar{S}}{\mathbf{v}'_{res} S'}) + (w'_{res} S' - \overline{w'_{res} S'})}_{v'S' - \overline{v'S'}} \cdot \hat{\mathbf{n}}] \, dB \, dt^*. \quad (3)$

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Here, B represents the four lateral surfaces, i.e., south, north, east and west, along with 347 the vertical boundary at 200 m. The $\bar{v}_e S'$ term across each face is also decomposed into 348 contributions from Eulerian and bolus components. These components are shown in Fig. 9. 349 We observe that during F_2 , there is an anomalous Eulerian salt flux $(\bar{v}S')$ entering the 350 western boundary. It is responsible for bringing saltier water as $\bar{v}S'$ along the NAC. This 351 is in contrast to the F_1 event where the advection is driven by the $v'\bar{S}$ term. A balance 352 between the Eulerian flux of $v'\bar{S}$ across the southern and western boundaries along with 353



Figure 9: Decomposition of the anomalous advection term \mathcal{A}' in the salt budget for the upper 200 m of the ESNA across the lateral and vertical boundaries. Shown here is the contribution from changes in the circulation with a mean salinity field $(v'\bar{S})$ and changes in salinity along mean flow $(\bar{v}S')$. This decomposition is made for both Eulerian and bolus velocity components. See Eq. (3) for details on the terms, and see also Fig. 7 for information on the full salt budget. F₁ and F₂ are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

the bolus flux across the 200 m horizon keep the ESNA fresh. This is the major difference between the two events.

The diffusion component of the salt mass anomaly budget is expressed as

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$$\mathcal{D}' = \rho_0 \int^t \int_V s^\star \underbrace{(D_{\sigma,S}}_{\mathcal{D}_H} + \underbrace{D_{\perp,S}}_{\mathcal{D}_V}) dV^\star dt^\star.$$
(4)

Fig. 10 shows the decomposition of the diffusion term into vertical and horizontal com-358 ponents. These terms are similar for both the ESNA and SPNA. The vertical diffusion 359 across 200 m dominates horizontal diffusion. Combined with the advection term, it is 360 the vertical diffusion that drives the salt mass anomaly during 1995–2000 to increase, 361 thereby reversing the fresh anomaly of the 1990s. It also drives the reversal in the salt 362 mass anomaly for 2004–2012. At this time, a decrease in the advection term brings the 363 salt mass anomaly below zero, initiating the F_2 event in both basins (see Fig. 7). The 364 jumps in the vertical diffusion term, which appear as staircases in the time series, are 365 associated with winter-time deep mixing events. The purple mixed layer depth timeseries 366 in Fig. 10 shows this association. 367

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3.4 Variation of Budget Terms With Depth

As the total advection and vertical diffusion terms are of the same magnitude, we investigate their sensitivity to varying the depth of the control volume from its nominal value of 200 m. Fig. 11 shows the change in the ESNA advective and diffusive con-



Figure 10: Decomposition of the diffusion term (\mathcal{D}') in the anomalous salt mass budget for the upper 200 m of the (a) Subpolar North Atlantic and (b) Eastern subpolar North Atlantic (ESNA, see Fig. 2 for the definition of the ESNA). Decomposition terms include horizontal (D_H) and vertical diffusion (D_V) , as in (4). The area-averaged mixed layer depth for each basin is shown in purple. F₁ and F₂ are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

vergence as the integration depth of the bottom of the control volume increases from 10 372 m to 5500 m. With increasing depth of integration, the contribution of the total diffu-373 sive flux in the salt tendency increases and then decreases, reaching a peak at 200 m, which 374 is the nominal value. For integration depths greater than 200 m, there is a trade-off (anti-375 correlation) between the contribution of the advective convergence term and the diffu-376 sive convergence term. For a control volume spanning the whole water column (integra-377 tion depth of 5500 m), the diffusion term is small and advection dominates. Fig. 11 con-378 textualizes the nominal value of 200 m used to study freshening in the ESNA by Holliday 379 et al. (2020); J. Zhang et al. (2021) and Fox et al. (2022). From a salt budget perspec-380 tive, the choice of 200 m maximizes the role of diffusion in modulating the salinity anoma-381 lies in the ESNA. 382

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3.5 Role of Precipitation, Evaporation, and Runoff

Performing a salt budget analysis for a control volume in the ocean does not account for changes in the salinity due to surface freshwater exchange, i.e., precipitation (P), evaporation (E), and runoff (R). To account for freshwater forcing from the atmosphere, ECCOv4r4 provides a diagnostic representing P-E fluxes and freshwater input from river runoff, R. We use the seawater volume budget and the salt budget to estimate the contribution of P-E+R in changing the salinity in the ESNA and SPNA. The volume conservation in ECCOv4r4 is (Forget et al. (2015); see their equation 3)

$$\frac{1}{H}\frac{\partial\eta}{\partial t} = -\nabla_{z^{\star}} \cdot s^{\star}\mathbf{v} - \frac{\partial w}{\partial z^{\star}} + s^{\star}\mathcal{F}.$$
(5)



Figure 11: Variation of the advection and diffusion term (horizontal diffusion D_H and vertical diffusion D_V with depth of the ESNA control volume used in the salt mass budget analysis. The 50 ECCO depth levels (from 10 m to 5000 m) are plotted. The black line denotes the 200 m isobath and the cyan and magenta lines denote 10 m and 5500 m depth levels respectively. F_1 and F_2 are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

Here, η is the sea surface height, $\mathbf{v} = (u_e, v_e, w_e)$ are the horizontal and vertical Eu-392 lerian velocity components. The other terms are the same as those in Eq. (A1). Eq. (5)393 expresses that the rate of change of the volume is a sum of surface freshwater forcing and 394 advective volume-flux divergence. Integration of Eq. (5) in space and time is used to cal-395 culate the volume-integrated anomaly in the total mass of the ESNA and SPNA (EC-396 COv4r4 makes the Boussinesq approximation, so seawater volume is proportional to sea-397 water mass). As Eq. (1), we express this as 398

$$M_{s}(t) \equiv \underbrace{\int_{V}^{t} \int_{V} \rho_{0} \, dV^{\star} \, dt}_{\text{Seawater Mass}} = \underbrace{\rho_{0} \int_{V}^{t} \int_{V} -\nabla_{z^{\star}} \cdot (s^{\star} \mathbf{v}) - \frac{\partial(w)}{\partial z^{\star}} \, dV^{\star} \, dt}_{\text{Advection}} + \underbrace{\rho_{0} \int_{V}^{t} \int_{V} s^{\star} \mathcal{F} \, dV^{\star} \, dt}_{\text{Forcing}}.$$
(6)

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Similar to the salt mass anomaly analysis, we remove the seasonality and long term trends 403 from the seawater mass time series $M_s(t)$. This gives us the time integrated seawater mass 404 anomaly, which is shown in Fig. 12 (labelled "seawater anomaly"). We find that the to-405 tal mass of the SPNA and the ESNA does not change significantly over 1992–2017. This 406 is due to a compensation between the anomalous P-E+R and the convergence of mass 407 over the basins. Josey and Marsh (2005) and Holliday et al. (2020) show that the ESNA 408 received anomalous positive P-E+R, which was primarily precipitation (P) during 1992– 409

1999 and 2012–2017. This is balanced by an increased mass flux exiting the basin dur-410

411

ing the same period.



Figure 12: (a) Seawater mass anomaly budget for the upper 200 m of the (a) Subpolar North Atlantic and (b) the eastern subpolar North Atlantic (ESNA, see Fig. 2 for the definition of the ESNA). F_1 and F_2 are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

We can now decompose the total contribution to changes in the average salinity 412 of the control volume using a combination of the salt mass and seawater mass budget. 413 We express the salt tendency in the salt conservation equation using 414

$$\frac{\partial (s^{\star}S)}{\partial t} = s^{\star} \frac{\partial S}{\partial t} + S \frac{\partial s^{\star}}{\partial t}$$

and 417

415 416

418 419

$$\frac{\partial s^{\star}}{\partial t} = \frac{1}{H} \frac{\partial \eta}{\partial t} \tag{8}$$

(7)

(from the definition of s^{\star}). Along with Eq. (5), we can rewrite the salt conservation equa-420 tion (A1) as an equation for the tendency of salinity (similar to equation (12) in Piecuch 421 (2017)), namely, 422

$$\underbrace{\frac{\partial S}{\partial t}}_{424} = \underbrace{\frac{1}{s^{\star}} \left[S \nabla_{z^{\star}} \cdot (s^{\star} \mathbf{v}) + S \frac{\partial w}{\partial z^{\star}} - \nabla_{z^{\star}} \cdot (s^{\star} S \mathbf{v}_{res}) - S \frac{\partial w_{res}}{\partial z^{\star}} \right] + D_s + \mathcal{F}_s}_{\Delta S_{Atm.}} - \underbrace{SF}_{\Delta S_{Atm.}}.$$
(9)

We integrate Eq. (9) in space and time, remove the seasonality and long term trends 425 from the salinity time series to get the changes in salinity over the SPNA and the ESNA. 426 This is shown in Fig. 13 (labelled " ΔS "). We sum the contributions of ocean advection, 427



Figure 13: Salinity contribution from the atmosphere and ocean for the upper 200 m of the (a) Subpolar North Atlantic and (b) the eastern subpolar North Atlantic (ESNA, see Fig. 2 for the definition of the ESNA). F_1 and F_2 are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

diffusion and surface salt exchange and call it ΔS_{Ocean} . The contribution of surface freshwater exchange is labelled $\Delta S_{Atm.}$.

For the SPNA, we find that the precipitation plays an important role during F_1 , but not during F_2 . After the F_1 event, the changes in salinity are due to ocean advection and diffusion. For the ESNA, precipitation is important for both the fresh anomaly events. Josey and Marsh (2005) and Holliday et al. (2020) conclude the same for the Iceland basin.

435 4 Discussion

Anomalous low salinity events in the upper 200 m eastern subpolar North Atlantic during the 1990s (F_1) and 2010s (F_2) appear to be qualitatively similar when observed as salinity anomaly signals. Using salt mass budget analysis in the ECCOv4r4 state estimate, we distinguish between the two events and conclude the following:

- 1. The fresh anomaly in the 1990s (F₁) occurs due to anomalous circulation of mean salinity $(v'\bar{S})$. This primarily occurs as the Eulerian component of $v'\bar{S}$ across the southern and western boundaries and the bolus component of the vertical flux of $v'\bar{S}$ across the 200 m horizon in the ESNA (Fig. 9).
- 2. The fresh anomaly during 2012–2017 (F₂) is due to the mean circulation of anomalous salinity ($\bar{v}S'$). This is entirely due to the Eulerian component of $\bar{v}S'$ across the western face of the Iceland Basin (Fig. 9).

447 448 449 3. Vertical diffusive flux across the 200 m depth boundary reverses the freshening after F_1 by increasing the salt mass content during 1995–2000. It also controls the salinity during the period leading up to the F_2 event (Fig. 10).

Therefore, although the two salinity anomalies are similar, they do not share the same mechanisms. The Great Salinity Anomaly (GSA) event (F_1) as described in Belkin (2004) propagates via the Labrador Sea and into the Iceland Basin in the mid 1990s. The fresh anomaly observed in the ESNA in the mid 2010s (F_2) has similar characteristics. The difference between the two is that F_2 is dominated by circulation of anomalous salinity along mean currents into the Iceland Basin and F_1 is dominated by anomalous circulation of mean salinity (Fig. 8).

We note that vertical diffusion has an important role to play in driving the neg-457 ative salinity anomalies to positive salinity anomalies. The 200 m isobath becomes sig-458 nificant in pushing salt flux into the upper layer of the ocean. As we increase the ver-459 tical depth of integration in the salt budget, the vertical exchange of salt via diffusion 460 decreases and the horizontal advective exchange increases. Sanders et al. (2022) conduct 461 a similar study to investigate the 2015 anomalous cooling in the eastern and central SPNA 462 (defined over 50–20°W, 43–63°N, which overlaps with, but extends further west than, 463 our definition of the ESNA, Fig. 2). They analyze the heat budget averaged over the mixed 464 layer using ECCOv4r3 and ECCOv4r4. They observe surface heat loss to initiate and 465 drive the cooling, with advection sustaining the anomaly in the region (as expected from 466 Tesdal and Abernathey (2021)). They also emphasize the role of vertical diffusion across 467 the base of the mixed layer in the re-emergence of the anomaly during summer of 2014. 468 The importance of horizontal advection and vertical diffusion for the temperature anoma-469 lies is similar to the importance of these processes for the salinity anomalies studied here 470 (Fig 9). 471

The surface freshwater fluxes (P-E+R) play a significant role in enhancing the fresh anomalies in the ESNA during the F₂ event. This was also noted by Holliday et al. (2020).

Holliday et al. (2020) state that the 2010s (F_2) event does not share the same char-474 acteristics of the 20th century GSAs. They argue that this event shows no freshening 475 in the Labrador Sea, unlike GSAs (Belkin (2004); Sundby and Drinkwater (2007)). How-476 ever, in the ECCOv4r3 freshwater budget for the Labrador Sea explored by Tesdal and 477 Haine (2020) there is an increased freshwater flux from the Labrador Sea via the Labrador 478 Current. This compensates an increased freshwater flux across the Davis Strait so that 479 little net salinity change occurs in the Labrador Sea over this period. Still, recent pa-480 pers argue that the 2010s (F_2) freshening event resembles a GSA. Specifically, Devana 481 et al. (2021) observe that salinity decrease from November 2015 to March 2017 in the 482 Iceland-Scotland Overflow Water is similar to the freshening observed in the 1990s. Also, 483 Biló et al. (2022) note that the salinity decrease during 2016–2019 in the Irminger Sea 484 $(0.04 \text{ psu year}^{-1})$, following the freshening in the Iceland basin, is among the highest ever 485 recorded. 486

Realistic, physically-consistent state estimates, such as ECCOv4r4 and ASTE_1 used 487 here, are valuable to diagnose mechanisms of large-scale inter-annual salinity and tem-488 perature fluctuations because closed volume, heat, and salt budgets can be constructed. 489 Apart from possible bias (section 3.1), such state estimates have some drawbacks, how-490 ever. State estimates do not include error estimates in the model output fields, for in-491 stance (Piecuch et al., 2017). Tesdal and Haine (2020) address this issue by using ± 2 492 standard deviations of the monthly lateral ECCO fluxes as a substitute for formal un-493 494 certainty estimates. Another limitation is that state estimates typically span a relatively short period (1992–2017 in the case of ECCOv4r4). Therefore, investigations of low-frequency 495 (decadal to centennial) salinity and temperature variability are not possible with these 496 products. 497

For future studies, it is important to investigate such low-frequency fluctuations 498 in the ESNA. Notably, the importance of the meridional overturning circulation, which 499 has not been addressed in this study, and atmospheric forcing needs further attention. 500 Coupled climate models are a promising resource for these studies because they also al-501 low construction of closed volume, heat, and salt budgets, but with much longer dura-502 tions. 503

5 Data Availability Statement 504

We use an open source python package, OceanSpy (https://oceanspy.readthedocs 505 .io; Almansi et al., 2019), to create and analyze synthetic hydrographic sections in the 506 model data. We also use the python package gcm-filters (https://gcm-filters.readthedocs 507 .io; Loose et al., 2022) to apply spatial Gaussian filters for smoothing the AVISO data. 508 The ECCO dataset is publicly available on the SciServer system (Medvedev et al., 2016) 509 and at https://podaac.jpl.nasa.gov/ECCO. ASTE_1 data is publicly available on https:// 510 arcticdata.io/catalog/portals/ASTE. The EN4 data are available at www.metoffice 511 .gov.uk/hadobs/en4/. Python scripts with Juptyer notebooks used for analyzing ECCO 512 budgets can be accessed at https://github.com/asiddi24/Siddiqui_et_al_JGR_Oceans 513 _2024 along with scripts used for generating all figures in the manuscript. 514

Appendix A Salt Budget equations 515

This appendix is useful in understanding the formulation of Eqs. (1, 2, 3). The z^* 516 coordinate is used in ECCOv4r4 to allow for exact tracer conservation (Campin et al., 517 2004) and to improve representation of flow over steep topography (Adcroft & Campin, 518 2004). Physically, z^* allows for variations in the non-linear free surface to be distributed 519 throughout the vertical water column. Using this coordinate, Forget et al. (2015) express 520 the salt conservation equation as 521

$$\frac{\partial(s^*S)}{\partial t} = -\nabla_{z^*} \cdot (s^* S \mathbf{v}_{res}) - \frac{\partial(S w_{res})}{\partial z^*} + s^* (\mathcal{F}_S + D_{\sigma,S} + D_{\perp,S}).$$
(A1)

Integrating Eq. (A1) over a spatial domain V and over time yields a time series for 523 the salt mass, M(t), expressed in Eq. (1). 524

For the anomalous advection term,
$$\mathcal{A}'$$
, consider

$$\mathbf{v}_{res} = \mathbf{\bar{v}}_{res} + \mathbf{v}'_{res}$$

527
$$w_{res} = \bar{w}_{res} + w'_{res}$$

528 $S = \bar{S} + S',$

$$S = S$$

This decomposition implies that

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$$\bar{\mathbf{v}}_{res}\bar{S} = \overline{\mathbf{v}_{res}S'} - \overline{\mathbf{v}'_{res}S'},$$

 $\bar{w}_{res}\bar{S} = \overline{w_{res}S} - \overline{w'_{res}S'}.$

The anomalous advection term is therefore: 534 535

$$\mathcal{A}' = -\rho_0 \int^t \int_V \nabla_{z^\star} \cdot \left[\left(s^\star (\bar{\mathbf{v}}_{res} S' + \mathbf{v'}_{res} \bar{S} + \mathbf{v'}_{res} S' - \overline{\mathbf{v'}_{res} S'} \right) \right] dV^\star dt^\star - \rho_0 \int^t \int_V \frac{\partial(\bar{w}_{res} S' + w'_{res} \bar{S} + w'_{res} S' - \overline{w'_{res} S'})}{\partial z^\star} dV^\star dt^\star.$$
(A2)

This is rearranged in Eq. (2). 539

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Controls on upper ocean salinity variability in the eastern subpolar North Atlantic during 1992–2017

Ali H. Siddiqui¹, Thomas W. N. Haine¹, An T. Nguyen², Martha Buckley³

¹Department of Earth and Planetary Sciences, Johns Hopkins University, Baltimore, MD, USA ²Oden Institute for Computational Engineering and Sciences, University of Texas at Austin, Austin, TX, USA

³Atmospheric, Oceanic and Earth Sciences, George Mason University, Fairfax, VA, USA

Key Points:

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- Two fresh anomalies observed in the eastern subpolar North Atlantic upper ocean during 1992–2017 share similar spatial characteristics.
 Salt budget analysis shows the 2012–2016 fresh anomaly in the upper 200 m occurs due to transport of anomalous salinity by mean currents.
 In contrast, the fresh anomaly in the 1990s is due to anomalous circulation of the
 - In contrast, the fresh anomaly in the 1990s is due to anomalous circulation of the mean salinity field.

Corresponding author: Ali H. Siddiqui, asiddi24@jhu.edu

15 Abstract

The eastern subpolar North Atlantic upper ocean salinity undergoes decadal fluctuations. 16 A large fresh anomaly event occurred during 2012–2016. Using the ECCO state estimate, 17 we diagnose and compare mechanisms of this low salinity event with that of the 1990s 18 fresh anomaly event. To avoid erroneous interpretations of physical mechanisms due to 19 reference salinity values in the freshwater budget, we perform a salt mass content bud-20 get analysis of the eastern subpolar North Atlantic. It shows that the recent fresh anomaly 21 occurs due to the circulation of anomalous salinity by mean currents entering the east-22 ern subpolar basin from its western boundary via the North Atlantic Current. This is 23 in contrast to the early 1990s, when the dominant mechanism governing the fresh anomaly 24 was the transport of the mean salinity field by anomalous currents across the southern 25

²⁶ boundary of the subpolar North Atlantic.

27 Plain Language Summary

On decadal time scales, the eastern subpolar North Atlantic shifts between a salty 28 and fresh upper ocean. These changes are significant and need to be investigated because 29 of their impacts on global ocean circulation and regional ocean biogeochemistry. Between 30 2012 and 2016, there was a large event where this region became fresher than it had been 31 in over a century. A previous similar event occurred in the early 1990s, but with a smaller 32 magnitude. We use a numerical model of the ocean to figure out why these events oc-33 curred. Our study shows that there were two different mechanisms at play. The recent 34 event occurred because a lot of fresh water came in from the west by the mean currents. 35 The 1990s event occurred because ocean currents shifted and brought fresh water from 36 the south. 37

³⁸ 1 Introduction

Large scale low salinity events occur in the eastern subpolar North Atlantic Ocean (ESNA) on decadal time scales. Based on observations, the subpolar North Atlantic has been undergoing such decadal salinity changes since at least the early 20th century (Sundby & Drinkwater, 2007; Dickson et al., 1988; Dooley et al., 1984; R. Zhang & Vallis, 2006; Dickson et al., 1988; Belkin et al., 1998; Belkin, 2004). During the 1992–2017 period, there were two fresh anomaly events in the ESNA reaching maximum freshwater accumulation in 1995 and 2016 respectively.

We highlight previous discussion in the literature on mechanisms that control the salinity in the ESNA. Changes in the strength and size of the subpolar gyre (SPG) and local atmospheric forcing in the ESNA are proposed mechanisms involving processes local to the subpolar gyre (Fox et al., 2022; Holliday et al., 2020). Remote mechanisms to modulate the salinity in the ESNA on such time scales include advection of salt anomalies from the Arctic or the subtropics (J. Zhang et al., 2021; Yeager et al., 2012; Häkkinen et al., 2011; Thierry et al., 2008; Sundby & Drinkwater, 2007; Holliday, 2003).

The strength and size of the subpolar gyre has been hypothesized to play an im-53 portant role in setting the salinity variability in the ESNA (Holliday, 2003; Hátún et al., 54 2005; Häkkinen & Rhines, 2004; Sarafanov et al., 2008; Yeager et al., 2012; Häkkinen et 55 al., 2011; Thierry et al., 2008), especially in the context of the warming and salinifica-56 tion that occurred in the mid 1990s to 2000s. The expansion of the SPG reduces the con-57 tribution of salty subtropical waters to the ESNA, reducing the salinity; in 1994 and 2016 58 sea surface height (SSH) contours show an expanded subpolar gyre and fresh anomalies 59 in the ESNA (Fig. 1). In contrast, the contraction of the SPG allows more subtropical 60 waters into the ESNA, increasing salinity; in 2008 the subpolar gyre is contracted, as SSH 61 contours retreat westward, and the ESNA is saltier (Fig. 1). 62

Subsequent studies have refined diagnostics for studying the relationship between 63 the subpolar gyre strength and ESNA salinity; rather than considering SSH-based in-64 dices of the subpolar gyre strength ((e.g., Häkkinen & Rhines, 2004)), Tesdal et al. (2018) 65 analyze a density-based gyre index, which is a proxy for the baroclinic strength of the 66 gyre (Koul et al., 2020). Foukal and Lozier (2017, 2018) suggest that the salinity in the 67 ESNA is more strongly influenced by the intergy transport, which is modulated by the 68 Atlantic meridional overturning circulation (AMOC). Koul et al. (2020) perform Lagrangian 69 tracking experiments based on multiple definitions of the SPG strength during 1993–2016 70 and conclude that 64.8% virtual floats reaching ESNA originate from subtropical wa-71 ters. Contributions from subpolar-sourced waters increase five-fold during an expanded 72 state of the SPG, however. 73

Another method to diagnose mechanisms controlling temperature and salinity vari-74 ability is by performing budget calculations. To address the decadal SST variability in 75 the subpolar North Atlantic, Piecuch et al. (2017) calculated the heat budget for 46° -76 65° N and concluded that the warming in the late 1990s and subsequent cooling since 2008 77 is primarily driven by oceanic advective flux convergence. The anomalous flux conver-78 gence is dominated by anomalies across the southern boundary $(46^{\circ}N)$. Similar studies 79 by Oldenburg et al. (2018) and Tesdal and Haine (2020) reach the same conclusion on 80 the dominance of the southern boundary advection in setting subpolar North Atlantic 81 heat and freshwater variability. However, the role of the AMOC is unclear in this mech-82 anism. Similarly, Sanders et al. (2022) investigate the 2015 anomalous cooling in the east-83 ern and central subpolar region (defined over 50–20°W, 43–63°N) using a mixed layer 84 heat budget. They observe that surface heat loss initiates and drives the cooling, with 85 advection sustaining the anomaly in the region (as expected from Tesdal and Abernathey 86 (2021)). They also emphasize the role of vertical diffusion across the base of the mixed 87 layer in the re-emergence of the anomaly during summer of 2014. 88

⁸⁹ Bryden et al. (2020) observed that there has been a mean increase of 0.12 ± 0.04 ⁹⁰ Sv of freshwater flux (relative to a reference salinity of 35.17 psu) after 2010 compared ⁹¹ to before 2009. This increase is about 10% of the 2004–2009 average flux. They propose ⁹² that a freshening of 0.062 ± 0.013 Sv over the eastern subpolar gyre during 2014–16 rel-⁹³ ative to 2007–09 is primarily due to the reduction of the AMOC by 2.5 Sv after 2009.

So far, we have discussed how circulation variability within the Atlantic basin con-94 tributes to salinity anomalies in the ESNA. For the 2012–2017 freshening event, Holliday 95 et al. (2020) describe the primary mechanism of net freshwater gain in the upper 200 m 96 of the Iceland basin as the rerouting of Arctic-sourced Labrador Current water into the 97 northern branch of the North Atlantic Current (NAC; Reverdin et al., 2003). It is mod-98 ulated by changes in the SPG strength driven by changes in atmospheric forcing. Re-99 cently, Fox et al. (2022) highlighted that reduced surface heat loss led to an increase in 100 warmer (less dense) waters in the Labrador Sea. The transport of these less dense wa-101 ters from the upper ocean layers through the Labrador Current along with reduced vol-102 ume transport from the Gulf Stream drove the cooling and freshening in the eastern sub-103 polar region. 104

However, some studies suggest an important role for interactions between the sub-105 polar North Atlantic and the Arctic. J. Zhang et al. (2021) and Sundby and Drinkwa-106 ter (2007) attribute ESNA freshening events during 1983–1995 and 1947–2000 to the ex-107 port of freshwater buildup in the Arctic. They suggest that sea ice and liquid freshwa-108 ter anomalies travel via the Fram Strait and Davis Strait to the Labrador Sea and cir-109 culate around the eastern subpolar gyre in the North Atlantic Current. The proposed 110 111 mechanism of freshwater buildup in the 1990s is increased freshwater flux from the Davis Strait (Belkin, 2004), which entered the Labrador Sea and propagated around the east-112 ern subpolar gyre (Sundby & Drinkwater, 2007). 113



Figure 1: Annually-averaged subpolar North Atlantic upper-ocean salinity (0–200 m, colors) and sea-surface height (SSH; contours) averaged over one year preceding the salinity field. The SSH field is from the AVISO dataset and the salinity field is from the EN4 product. Following Chafik et al. (2019), the grey contours range from -0.8 m to 0.8 m with a spacing of 0.1 m and represent the mean dynamic topography (CNES-CLS2013 MDT). The red contours are -0.3, -0.2, -0.1 m and represent the three branches of the NAC. A Gaussian filter is used to smooth the SSH field with a scale of 1.25°. Modified from Fig. 3 of Weijer et al. (2022).

In this paper, we focus on what sets the upper ocean salinity in the ESNA on decadal time scales with emphasis on the two recent freshening events in the 1990s and 2010s using observations and modelling tools. We look at salt content anomaly budgets to explore oceanic mechanisms and further investigate the contribution of surface freshwater forcing in setting upper ocean salinity in the region.

In section 2 we discuss the methods used in this paper and investigate upper ocean salinity variability as seen in observations. In section 3, we compare and contrast the two fresh anomaly events observed during 1992–2017 using observations and ocean state estimates. We then diagnose the salinity variability using salt budget analysis for the ESNA and discuss potential mechanisms for the salinity variability.

124 2 Methods

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2.1 Datasets

The main tool for our analysis of upper ocean salinity is the ECCO (Estimating 126 the Circulation and Climate of the Ocean) version 4 release 4 (Forget et al., 2015) and 127 ASTE (The Arctic Subpolar Gyre sTate Estimate) release 1 (Nguyen et al., 2021) ocean 128 state estimates. The ECCOv4r4 state estimate (ECCO Consortium et al., n.d.) is a dynamically-129 consistent, data-constrained solution of the MITgcm model for the period 1992–2017. 130 This allows for the construction of realistic closed budgets of volume, heat and salt. The 131 horizontal resolution is 1°. ASTE_R1 is also a data-constrained ocean-sea ice model-data 132 output that covers the Arctic Ocean and Atlantic ocean with lateral boundaries at 47.5° N 133 in the North Pacific and 32.5°S in the South Atlantic for 2002–2017; the ASTE_R1 hor-134 izontal resolution is $1/3^{\circ}$. 135

We utilize a number of observational datasets to evaluate ECCOv4r4 and ASTE_R1. First, we compare to the EN4 hydrographic dataset (Good et al., 2013), which is an observational product compiled by the UK Met office. We use data from two hydrographic surveys to evaluate the performance of ECCOv4r4 for our analysis. The Extended EL-LET line (Holliday & Cunningham, 2013) is a section from Iceland to Scotland and the OVIDE line (Daniault et al., 2016) is a combination of sections from the southern tip

of Greenland to Portugal. These are the same sections used by Holliday et al. (2020). 142 We compare the salinity anomalies in the sections observed in the model data with the 143 ELLET and OVIDE sections. The data from the hydrographic survey carried out over 144 June–July, 2016 and May–June, 2016 for the OVIDE and ELLET line, respectively, is 145 compared with monthly-mean salinity anomaly from the model data. We consider the 146 upper 200 m annually-averaged ESNA salinity anomaly for 1992–2017 using the EN4 hy-147 drographic dataset (Good et al., 2013) along with ECCOv4r4 and ASTE_R1 datasets. 148 The upper ocean 200 m is used to remain consistent with previous work that investigates 149 upper ocean salinity in the ESNA (Holliday et al., 2020; Koul et al., 2020). 150

We use ECCOv4r4 to track the evolution of salinity anomalies for the two fresh anomaly events on monthly and also annual time scales. To investigate the mechanisms governing this evolution, we carry out salt budget analysis with ECCOv4r4 for the period 1992– 2017 for the entire subpolar North Atlantic (SPNA) and the ESNA. The salt budget analysis builds on the work by Buckley et al. (2014, 2015); Piecuch et al. (2017); Oldenburg et al. (2018); Tesdal and Haine (2020); Nguyen et al. (2021) to investigate heat and salinity variability in the subpolar North Atlantic.

The ESNA is defined as a box over $10-30^{\circ}$ W, $46-65^{\circ}$ N and the SPNA is defined as the North Atlantic between $45^{\circ}-65^{\circ}$ N. Fig. 2 shows the location of the ESNA.

160 3 Results

In this section, we evaluate the ECCOv4r4 model data. Next, we consider the spatial and temporal structure of the fresh anomaly events in the subpolar region. Then we investigate the salt budget results.

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3.1 Evaluation of ECCOv4r4

Here we present a comparison of ECCOv4r4 to ship-based hydrographic sections. 165 Fig. 2a shows the positions of the OVIDE and ELLET sections overlaid with sea level 166 height anomaly contours from ECCO. In summer of 2016, the upper 1000 m of the Rock-167 all Trough has the highest salinity in the ELLET section. The salinity decreases below 168 1000m, which is seen in both the observations and ECCO data (Fig. 3). ECCO overes-169 timates the salinity in the upper 1000 m, however. This is also true on the OVIDE sec-170 tion from the Reykjanes Ridge to the Iberian abyssal plain during summer of 2016. The 171 ECCO salinity in the Irminger Sea during summer of 2016 is realistic, however. ECCO 172 also captures the sub-surface salinity minimum over the Iberian abyssal plain. On the 173 ELLET section, ECCO overestimates the salinity in the 0–200 m Iceland Basin by around 174 0.06 psu, and underestimates it by 0.04 psu in the Rockall Trough. Similarly, on the OVIDE 175 section, ECCO overestimates the 0–200 m salinity in the ESNA control volume (Fig. 2) 176 by around 0.05 psu. The time series (Fig. 2) of 0-200 m depth-averaged salinity anoma-177 lies for the ESNA in ECCO accurately shows the fresh anomaly observed during 2012-178 2017 and 1992–1995 compared to EN4. The salinity difference between ECCO and EN4 179 is 0.027 ± 0.046 psu (from Fig. 2b, but using the absolute salinity values, not the anoma-180 lies). The biases in the ELLET and OVIDE sections quoted above are consistent with 181 this ECCO/EN4 difference, and are relatively small compared to the salinity fluctuations 182 over time seen in Fig. 2b. This builds confidence in the use of ECCO for our analysis. 183 The only noticeable disagreement between EN4 and ECCO seen in Fig. 2b is during 1995-184 1996, when the ESNA shows anomalous positive salinity anomalies in EN4, whereas ECCO 185 shows negative salinity anomalies. Spatial maps of 0–200 m salinity anomalies show a 186 large positive anomaly situated south of the Grand Banks in 1995 in both datasets (see 187 Supplemental Figs. 1–3). In EN4 this anomaly spreads throughout the ENSA in 1996, 188 but it does not spread so far east in ECCO. 189



Figure 2: (a) Subpolar North Atlantic (SPNA) sea level height anomaly contours (spacing of 0.04 m) in the ECCOv4r4 dataset averaged over 1992–2017. The Extended ELLET line (EEL) and OVIDE sections are also shown. (b) Upper ocean 200 m salinity anomaly time series for the eastern Subpolar North Atlantic (ESNA) from EN4 (red), ECCOv4r4 (black) and ASTE_R1 (grey) datasets. The ESNA is defined as 45–65°N, 10–30°W (purple box in (a)).

3.2 Spatial and Temporal Structure of Upper Ocean (S_{200}) Salinity Anomalies

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We first establish the occurrence of two fresh anomalies in the upper 200 m of the 192 ESNA using the ECCOv4r4, ASTE_R1, and EN4 datasets (Fig. 2). We observe a fresh 193 anomaly in the ESNA in the early 1990s, after which there is a prolonged period of salin-194 ification until 2011, and a reversal to freshening thereafter. The first fresh anomaly event 195 (F_1) is observed from 1992 until 1995 in EN4 and 1997 in ECCOv4r4. For the second 196 fresh anomaly event (F_2) , all three datasets show the Iceland basin salinity anomaly drop-197 ping below zero after 2012 until 2017. This is also reflected in the decadal cycle in the 198 upper ocean salinity maps for the ESNA region in the ECCO, ASTE_1 and EN4 datasets 199

(see Supplemental Fig. 1–3). We label the 1990s fresh event as F_1 and the 2010s fresh 200

201

event as F_2 .

-2000

-3000

ECCO

0

200

 $\dot{400}$

600

Distance from first station [km]

(b)

800

1000

1200



Figure 3: Comparison of ECCOv4r4 salinity for (a) June–July 2016 with the OVIDE section and (b) May–June 2016 with the ELLET line. The ECCOv4r4 salinity sections are taken at the same times as the field observations. Vertical purple lines indicate parts of the sections inside the ESNA control volume defined in Fig. 2 (for the OVIDE section 700-2290 km, and for the ELLET section 0-1150 km, are within the ESNA control volume). Colorbar limits are different for the two section plots.

Spatial trends in the upper ocean (S_{200}) salinity are computed using EN4 and EC-202 COv4r4. During 2005–16, both products show a statistically significant freshening in the 203 ESNA at 95% confidence intervals using the student's t-distribution. ECCOv4r4 and EN4 204 disagree on trends in the Labrador Sea and the Grand Banks region, however (Fig. 4). 205

Next, we consider the annually averaged anomalies in the ECCO data and com-206 pare the two fresh anomaly events, F_1 and F_2 , which are shown in Fig. 5. As we trace 207 the freshwater event in the 1990s (F_1) we observe a fresh ESNA and saltier western sub-208 polar gyre (SPG) in 1992. During 1993 and 1994, a fresh anomaly is situated in the Labrador 209 Sea, and in the Iceland Basin in 1995. By 1996, the signal fades away from the ESNA. 210



Figure 4: Spatial distribution of linear trends in the upper-ocean salinity (0–200 m) over 2005–2016 using monthly mean fields of (a) EN4 and (b) ECCOv4r4. Dotted regions display insignificant trends calculated using the students-t test with a p-value of 0.05.

The 2010s event, has a similar fresh anomaly in the Labrador Sea in 2013, and also in the Iceland Basin in 2016. Note that in both events, there is a positive salinity anomaly

south of the Grand Banks region, preceding the maximum freshening in 1995 and 2016.

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Figure 5: Annually averaged anomaly maps for the upper 200 m salinity (S_{200}) during 1992–1994 and 2012–2014 in ECCOv4r4.

We investigate this further using the ECCO data by tracking salinity anomalies along 214 the western SPG and along the Gulf Stream with the Iceland Basin as a common ter-215 minus. We create a section following mean sea level anomaly contours around the sub-216 polar gyre which begins south of Denmark strait. The section follows the -0.6 m mean 217 sea level anomaly contour along the East Greenland Current around the southeast coast 218 of Greenland. The section continues along the West Greenland Current to the entrance 219 of Baffin Bay, where it retroflects and follows the Baffin Island Current, eventually reach-220 ing the Labrador Sea. At this point, the section follows the -0.8 m mean sea level anomaly 221 contour as it retroflects east of the Flemish Cap and follows the path of the northern branch 222 of the NAC to the Iceland basin. This section represents a potential subpolar pathway 223 for the anomaly propagation. 224

To investigate the salinity anomalies along a potential subtropical pathway, we create a section along the Gulf Stream which also terminates in the Iceland Basin. Both sections are inspired by the Lagrangian studies carried out by Burkholder and Lozier (2014);

Foukal and Lozier (2018); Koul et al. (2020); J. Zhang et al. (2021) where passive trac-

ers are tracked to the Iceland basin in a variety of experiments.



Figure 6: (a) Continuous sections along the western subpolar gyre boundary (purple; 56 stations) and along the Gulf Stream (red; 28 stations). The sections intersect at the NAC (station A1) and are merged from there to the Iceland Basin. (b) Hovmöller diagram of monthly salinity anomalies in the ECCO data along the sections. Vertical dashed lines are shown at Davis Strait (DS) and the Grand Banks (GB) region.

We first inspect the subtropical section. We find that there is a general tendency for there to be salinity anomalies of opposite sign along the Gulf Stream path and in the subpolar gyre, as noted by prior studies (Buckley et al., 2014; Joyce & Zhang, 2010; Sanchez-

Franks & Zhang, 2015; Yeager, 2015; Hátún et al., 2009; R. Zhang, 2008; Nye et al., 2011; 233 Yan et al., 2017, 2018). For both subpolar freshening events, there are positive salinity 234 anomalies along the Gulf Stream path. This positive salinity anomaly signal south of the 235 Grand Banks is also observed in the years preceding the fresh anomaly event in the an-236 nually averaged salinity anomaly maps (Fig. 5). Along the subpolar section, high fre-237 quency freshening/salinification events can be tracked from Denmark Strait along the 238 East Greenland Current and West Greenland Current. These represent seasonal fluctu-239 ations along the Greenland coast. Along the Labrador Current (from Davis Strait to the 240 Grand Banks), the characteristics of the freshening/salinification events changes with 241 lower-frequency variations than seen along the East and West Greenland Currents. From 242 the Grand Banks to the Iceland Basin, the freshening/salinification events are at even 243 lower frequencies. The relationship between salinity anomalies in the Icelandic Basin and 244 those in the East/West Greenland Current and Labrador Current is complex, with some 245 indications of signal propagation along the subpolar gye pathways, as suggested by Holliday 246 et al. (2020) and Fox et al. (2022). We can now explore the role of circulation changes 247 quantitatively by constructing a salt mass budget for the region. 248

3.3 Salt Budget Analysis

We construct a budget of the salt mass, an extensive quantity, which has a more accurate closure in the ECCO model output than a budget of salinity, an intensive quantity. The role of surface freshwater in the form of P-E+R (Precipitation-Evaporation+Runoff) is analyzed separately in a salinity budget calculation. We also avoid analyzing freshwater budgets due to ambiguities associated with reference salinity in such calculations (Schauer & Losch, 2019).

The salt conservation equation for the non-linear free surface in ECCOv4r4 is expressed in z^* coordinates (see equation (3) in Forget et al. (2015)). In z^* coordinates, sea surface height variations, η , are proportionally divided between ocean layers: $z^* = (z - \eta)/(H + \eta)$ (equation (1) in Piecuch (2017)), where z is the fixed vertical coordinate and H is the ocean depth.

We express the volume and time integrated salt content, M(t), for the control-volume V as

$$M(t) \equiv \underbrace{\rho_0 \int_{t^*=0}^{t^*=t} \int_V \frac{\partial(s^*S)}{\partial t} dV^* dt^*}_{\text{Salt Mass}} = \underbrace{\rho_0 \int_{t^*=0}^{t^*=t} \int_V -\nabla_{z^*} \cdot (s^*S\mathbf{v}_{res}) - \frac{\partial(Sw_{res})}{\partial z^*} dV^* dt^*}_{\text{Advection}} + \underbrace{\rho_0 \int_{t^*=0}^{t^*=t} \int_V s^*\mathcal{F}_S dV^* dt^*}_{\text{Forcing}} + \underbrace{\rho_0 \int_{t^*=0}^{t^*=t} \int_V s^*(D_{\sigma,S} + D_{\perp,S}) dV^* dt^*}_{\text{Diffusion}}.$$
(1)

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In this equation $s^{\star} = 1 + \eta/H$ is a scaling factor, $\nabla_{z^{\star}}$ indicates the gradient at con-267 stant z^{\star} , $(\mathbf{v}_{res}, w_{res})$ are the residual velocity fields defined as the sum of the Eulerian 268 and bolus velocities, \mathcal{F}_S is the forcing at the surface due to surface salt exchange due 269 to sea ice melting/formation and a redistribution of the surface flux in the vertical col-270 umn, and $D_{\sigma,S}$ and $D_{\perp,S}$ are diffusive processes parameterized along iso-neutral and ver-271 tical directions, respectively. We consider V to be the upper 200 m in the ESNA; we also 272 consider the upper 200 m for the whole SPNA (this assumption is tested below in sec-273 tion 3.4; see Appendix A for more details). Eq. (1) expresses that the total time inte-274 grated salt mass is balanced by the time integrated horizontal and vertical advective con-275 vergence of salt flux, diapycnal and isopycnal diffusion, and surface forcing. The four named 276 terms are each computed individually from the ECCO output, which allows us to test 277 the closure of (1). The ratio of the residual (left hand side minus right hand side) to the 278 salt mass (left hand side) is of $O(10^{-4})$. Details on how to close the salt budget in the 279 ECCO dataset are provided in Piecuch (2017). 280

We remove the mean seasonal cycle and a linear trend from all terms in the time 281 integrated salt budget. This gives us the salt content anomaly time series, which is shown 282 in Fig. 7 (labelled "salt anomaly"). We observe a negative salt content anomaly of the 283 SPNA during 1992–1997 and 2012–2017. For the ESNA, the salt content anomaly re-284 mains negative during 1992–1997 and 2012–2017. The salt mass anomaly increases and 285 reaches a maximum in 2004 for the entire SPNA, and in 2008 for the ESNA. We find that 286 both the diffusion and advection terms contribute to the salt content anomaly, and these 287 terms are strongly anti-correlated. The surface salt forcing (due to brine rejection) has 288 a negligible impact. 289

From 1995–2000 and 2004–2012, the change in the diffusion term is larger than the change in the advection term in the SPNA, and it drives the salt mass change. For the ESNA, this is seen during 1995–2000 and 2008–2011. We highlight the years in yellow/blue when the advection term increases/decreases rapidly in the two basins.



Figure 7: (a) Time and volume integrated salt anomaly budget for the upper 200 m of the (a) Subpolar North Atlantic and (b) the eastern subpolar North Atlantic (ESNA, see Fig. 2 for the definition of the ESNA). F_1 and F_2 are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

We further investigate the advection and diffusion terms separately. We decompose the advection term into time-average and anomaly terms, following Dong and Sutton (2002); Doney et al. (2007); Buckley et al. (2015); Piecuch et al. (2017) and Tesdal and Haine (2020). We write the advection term \mathcal{A} in (1), as $\mathcal{A} = \overline{\mathcal{A}} + \mathcal{A}'$, where the overbar denotes time averaging, i.e $\bar{\mathbf{v}} = \frac{1}{(t_f - t_i)} \int_{t_i}^{t_f} \mathbf{v} dt$ and the prime denotes departure from the time average. Initial and final time values in the averaging period are denoted by t_i and t_f . The averaging period is 1992–2017. This decomposes the advection term into variability produced by changes in the circulation $(\mathbf{v}'_{res}\bar{S})$, variability produced by changes in salinity $(\bar{\mathbf{v}}_{res}S')$, and that due to the co-variability of the circulation with the salinity $(\mathbf{v}'_{res}S' - \overline{\mathbf{v}'_{res}S'})$.

We now express $\mathbf{v}_{res} = \mathbf{v}_e + \mathbf{v}_b$, i.e., the total velocity is the sum of Eulerian (\mathbf{v}_e) and bolus (\mathbf{v}_b) velocities, and similarly for the vertical speeds (w_e, w_b) . Re-arranging the terms in Eq. (A2) gives:

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$$\mathcal{A}' = -\rho_0 \int^t \int_V \left\{ \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e S' + \frac{\partial(\bar{w}_e S')}{\partial z^\star} \right)}_{\bar{v}_e S'} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \mathbf{v}'_e \bar{S} + \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \mathbf{v}'_e \bar{S} + \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S}} + \underbrace{\left(\nabla_{z^\star} \cdot s^\star \bar{\mathbf{v}}_e - \frac{\partial(w'_e \bar{S})}{\partial z^\star} \right)}_{v'_e \bar{S$$

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$$\underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} (\mathbf{v}'_{e}S' - \overline{\mathbf{v}'_{e}S'}) + \frac{\partial (w'_{e}S' - \overline{w'_{e}S'})}{\partial z^{\star}} \end{pmatrix}}_{v'_{e}S' - \overline{v'_{e}S'}} \\ \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}S' + \frac{\partial (\overline{w}_{b}S')}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \mathbf{v}'_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} \end{pmatrix}}_{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S})}{\partial z^{\star}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \frac{\partial (w'_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}}_{b}\overline{S} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot s^{\star} \overline{\mathbf{v}} + \underbrace{\begin{pmatrix} \nabla_{z^{\star}} \cdot$$

$$\underbrace{\left(\nabla_{z^{\star}} \cdot s^{\star}(\mathbf{v}_{b}'S' - \overline{\mathbf{v}_{b}'S'}) + \frac{\partial(w_{b}'S' - \overline{w_{b}'S'})}{\partial z^{\star}}\right)}_{v_{b}'S' - \overline{v_{b}'S'}}\right\} dV^{\star} dt^{\star} + \epsilon. \quad (2)$$

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The sum of all the terms on the right hand side of Eq. (2) should equal the anomaly in the convergence of advective salt flux, which is an output in ECCOv4r4. When we compute the individual terms from monthly mean velocities and salinity (interpolated to the model velocity grid points), however, it introduces a numerical error, which we call ϵ . Volume integrated ratios of ϵ to each individual term are O(10⁻¹) (see Fig. 9).

The advection anomaly term is shown in Fig. 8. For the whole SPNA, during the 318 F_1 event, the variability in the advective convergence term is dominated by the anoma-319 lous Eulerian advection of mean salinity $(v'_e \bar{S})$. During 1992–1995, it is twice that of the 320 mean circulation of anomalous salinity $(\bar{v}_e S')$. For the F₂ event the mean bolus advec-321 tion of anomalous salinity $(\bar{v}_b S')$ tends to drive the freshening, as the Eulerian terms bal-322 ance each other out. Note that the advective convergences of anomalous circulation of 323 mean salinity and the mean circulation of anomalous salinity are anti-correlated during 324 both events. Anti-correlation of advective convergences due to salinity variations and due 325 to geostrophic velocity variations is expected when isobars and isohalines are aligned, 326 as shown for heat transport convergences by Buckley et al. (2015). 327

The decomposition of the anomalous advection term for the ESNA is different. The 328 F_1 negative salt anomaly is still dominated by the anomalous circulation of mean salin-329 ity $(v'_e S)$. In contrast, the driver for the F_2 anomaly is the mean circulation of anoma-330 lous salinity $(\bar{v}_e S')$. Also, unlike the SPNA, the anomalous circulation of mean salinity 331 is not anti-correlated to the mean circulation of anomalous salinity. This indicates that 332 either (1) the isobars and isotherms are not strongly aligned, which would occur if the 333 density field in the ESNA is only weakly dependent on salinity, or (2) there is a strong 334 contribution of agoestrophic transports to the advective convergences. For example, there 335 is no expected anticorrelation between salt transport convergences due to salinity vari-336 ations and those due to Ekman velocity variations. 337

To further investigate the $\bar{v}_e S'$ term, we look at the decomposed anomalous advection for the ESNA in terms of flux entering the control volume from the boundaries. We rewrite Eq. (A2), apply the Gauss-divergence theorem, and express the divergence of anomalous advection in the ESNA as a sum of salt fluxes across the four lateral boundaries and



Figure 8: Decomposition of the anomalous advection term \mathcal{A}' in the salt budget for the upper 200 m of the (a) Subpolar North Atlantic and (b) Eastern subpolar North Atlantic (ESNA, see Fig. 2 for the definition of the ESNA). The anomalous advection term is decomposed into contributions from: changes in the circulation $(v'_{res}\bar{S})$, changes in the salinity $(\bar{v}_{res}S')$, and changes due to the co-variability of the circulation with the salinity $(v'_{res}S' - v'_{res}S')$. This decomposition is made for both Eulerian and bolus velocity components. See Eq. (2) for details on the terms, and see also Fig. 7 for information on the full salt budget. F_1 and F_2 are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

across the 200 m vertical boundary: 342

 $\mathcal{A}' = -\rho_0 \int^t \int_B [\underbrace{\mathbf{s}^* \bar{\mathbf{v}}_{res} S' + \bar{w}_{res} S'}_{\bar{v}S'} + \underbrace{\mathbf{s}^* \mathbf{v}'_{res} \bar{S} + w'_{res} \bar{S}}_{\mathbf{v}'\bar{S}} + \underbrace{\mathbf{s}^* (\mathbf{v}'_{res} S' - \frac{v'\bar{S}}{\mathbf{v}'_{res} S'}) + (w'_{res} S' - \overline{w'_{res} S'})}_{v'S' - \overline{v'S'}} \cdot \hat{\mathbf{n}}] \, dB \, dt^*. \quad (3)$

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Here, B represents the four lateral surfaces, i.e., south, north, east and west, along with 347 the vertical boundary at 200 m. The $\bar{v}_e S'$ term across each face is also decomposed into 348 contributions from Eulerian and bolus components. These components are shown in Fig. 9. 349 We observe that during F_2 , there is an anomalous Eulerian salt flux $(\bar{v}S')$ entering the 350 western boundary. It is responsible for bringing saltier water as $\bar{v}S'$ along the NAC. This 351 is in contrast to the F_1 event where the advection is driven by the $v'\bar{S}$ term. A balance 352 between the Eulerian flux of $v'\bar{S}$ across the southern and western boundaries along with 353



Figure 9: Decomposition of the anomalous advection term \mathcal{A}' in the salt budget for the upper 200 m of the ESNA across the lateral and vertical boundaries. Shown here is the contribution from changes in the circulation with a mean salinity field $(v'\bar{S})$ and changes in salinity along mean flow $(\bar{v}S')$. This decomposition is made for both Eulerian and bolus velocity components. See Eq. (3) for details on the terms, and see also Fig. 7 for information on the full salt budget. F₁ and F₂ are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

the bolus flux across the 200 m horizon keep the ESNA fresh. This is the major difference between the two events.

The diffusion component of the salt mass anomaly budget is expressed as

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$$\mathcal{D}' = \rho_0 \int^t \int_V s^\star \underbrace{(D_{\sigma,S}}_{\mathcal{D}_H} + \underbrace{D_{\perp,S}}_{\mathcal{D}_V}) dV^\star dt^\star.$$
(4)

Fig. 10 shows the decomposition of the diffusion term into vertical and horizontal com-358 ponents. These terms are similar for both the ESNA and SPNA. The vertical diffusion 359 across 200 m dominates horizontal diffusion. Combined with the advection term, it is 360 the vertical diffusion that drives the salt mass anomaly during 1995–2000 to increase, 361 thereby reversing the fresh anomaly of the 1990s. It also drives the reversal in the salt 362 mass anomaly for 2004–2012. At this time, a decrease in the advection term brings the 363 salt mass anomaly below zero, initiating the F_2 event in both basins (see Fig. 7). The 364 jumps in the vertical diffusion term, which appear as staircases in the time series, are 365 associated with winter-time deep mixing events. The purple mixed layer depth timeseries 366 in Fig. 10 shows this association. 367

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3.4 Variation of Budget Terms With Depth

As the total advection and vertical diffusion terms are of the same magnitude, we investigate their sensitivity to varying the depth of the control volume from its nominal value of 200 m. Fig. 11 shows the change in the ESNA advective and diffusive con-



Figure 10: Decomposition of the diffusion term (\mathcal{D}') in the anomalous salt mass budget for the upper 200 m of the (a) Subpolar North Atlantic and (b) Eastern subpolar North Atlantic (ESNA, see Fig. 2 for the definition of the ESNA). Decomposition terms include horizontal (D_H) and vertical diffusion (D_V) , as in (4). The area-averaged mixed layer depth for each basin is shown in purple. F₁ and F₂ are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

vergence as the integration depth of the bottom of the control volume increases from 10 372 m to 5500 m. With increasing depth of integration, the contribution of the total diffu-373 sive flux in the salt tendency increases and then decreases, reaching a peak at 200 m, which 374 is the nominal value. For integration depths greater than 200 m, there is a trade-off (anti-375 correlation) between the contribution of the advective convergence term and the diffu-376 sive convergence term. For a control volume spanning the whole water column (integra-377 tion depth of 5500 m), the diffusion term is small and advection dominates. Fig. 11 con-378 textualizes the nominal value of 200 m used to study freshening in the ESNA by Holliday 379 et al. (2020); J. Zhang et al. (2021) and Fox et al. (2022). From a salt budget perspec-380 tive, the choice of 200 m maximizes the role of diffusion in modulating the salinity anoma-381 lies in the ESNA. 382

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3.5 Role of Precipitation, Evaporation, and Runoff

Performing a salt budget analysis for a control volume in the ocean does not account for changes in the salinity due to surface freshwater exchange, i.e., precipitation (P), evaporation (E), and runoff (R). To account for freshwater forcing from the atmosphere, ECCOv4r4 provides a diagnostic representing P-E fluxes and freshwater input from river runoff, R. We use the seawater volume budget and the salt budget to estimate the contribution of P-E+R in changing the salinity in the ESNA and SPNA. The volume conservation in ECCOv4r4 is (Forget et al. (2015); see their equation 3)

$$\frac{1}{H}\frac{\partial\eta}{\partial t} = -\nabla_{z^{\star}} \cdot s^{\star}\mathbf{v} - \frac{\partial w}{\partial z^{\star}} + s^{\star}\mathcal{F}.$$
(5)



Figure 11: Variation of the advection and diffusion term (horizontal diffusion D_H and vertical diffusion D_V with depth of the ESNA control volume used in the salt mass budget analysis. The 50 ECCO depth levels (from 10 m to 5000 m) are plotted. The black line denotes the 200 m isobath and the cyan and magenta lines denote 10 m and 5500 m depth levels respectively. F_1 and F_2 are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

Here, η is the sea surface height, $\mathbf{v} = (u_e, v_e, w_e)$ are the horizontal and vertical Eu-392 lerian velocity components. The other terms are the same as those in Eq. (A1). Eq. (5)393 expresses that the rate of change of the volume is a sum of surface freshwater forcing and 394 advective volume-flux divergence. Integration of Eq. (5) in space and time is used to cal-395 culate the volume-integrated anomaly in the total mass of the ESNA and SPNA (EC-396 COv4r4 makes the Boussinesq approximation, so seawater volume is proportional to sea-397 water mass). As Eq. (1), we express this as 398

$$M_{s}(t) \equiv \underbrace{\int_{V}^{t} \int_{V} \rho_{0} \, dV^{\star} \, dt}_{\text{Seawater Mass}} = \underbrace{\rho_{0} \int_{V}^{t} \int_{V} -\nabla_{z^{\star}} \cdot (s^{\star} \mathbf{v}) - \frac{\partial(w)}{\partial z^{\star}} \, dV^{\star} \, dt}_{\text{Advection}} + \underbrace{\rho_{0} \int_{V}^{t} \int_{V} s^{\star} \mathcal{F} \, dV^{\star} \, dt}_{\text{Forcing}}.$$
(6)

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Similar to the salt mass anomaly analysis, we remove the seasonality and long term trends 403 from the seawater mass time series $M_s(t)$. This gives us the time integrated seawater mass 404 anomaly, which is shown in Fig. 12 (labelled "seawater anomaly"). We find that the to-405 tal mass of the SPNA and the ESNA does not change significantly over 1992–2017. This 406 is due to a compensation between the anomalous P-E+R and the convergence of mass 407 over the basins. Josey and Marsh (2005) and Holliday et al. (2020) show that the ESNA 408 received anomalous positive P-E+R, which was primarily precipitation (P) during 1992– 409

1999 and 2012–2017. This is balanced by an increased mass flux exiting the basin dur-410

411

ing the same period.



Figure 12: (a) Seawater mass anomaly budget for the upper 200 m of the (a) Subpolar North Atlantic and (b) the eastern subpolar North Atlantic (ESNA, see Fig. 2 for the definition of the ESNA). F_1 and F_2 are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

We can now decompose the total contribution to changes in the average salinity 412 of the control volume using a combination of the salt mass and seawater mass budget. 413 We express the salt tendency in the salt conservation equation using 414

$$\frac{\partial (s^{\star}S)}{\partial t} = s^{\star} \frac{\partial S}{\partial t} + S \frac{\partial s^{\star}}{\partial t}$$

and 417

415 416

418 419

$$\frac{\partial s^{\star}}{\partial t} = \frac{1}{H} \frac{\partial \eta}{\partial t} \tag{8}$$

(7)

(from the definition of s^{\star}). Along with Eq. (5), we can rewrite the salt conservation equa-420 tion (A1) as an equation for the tendency of salinity (similar to equation (12) in Piecuch 421 (2017)), namely, 422

$$\underbrace{\frac{\partial S}{\partial t}}_{424} = \underbrace{\frac{1}{s^{\star}} \left[S \nabla_{z^{\star}} \cdot (s^{\star} \mathbf{v}) + S \frac{\partial w}{\partial z^{\star}} - \nabla_{z^{\star}} \cdot (s^{\star} S \mathbf{v}_{res}) - S \frac{\partial w_{res}}{\partial z^{\star}} \right] + D_s + \mathcal{F}_s}_{\Delta S_{Atm.}} - \underbrace{SF}_{\Delta S_{Atm.}}.$$
(9)

We integrate Eq. (9) in space and time, remove the seasonality and long term trends 425 from the salinity time series to get the changes in salinity over the SPNA and the ESNA. 426 This is shown in Fig. 13 (labelled " ΔS "). We sum the contributions of ocean advection, 427



Figure 13: Salinity contribution from the atmosphere and ocean for the upper 200 m of the (a) Subpolar North Atlantic and (b) the eastern subpolar North Atlantic (ESNA, see Fig. 2 for the definition of the ESNA). F_1 and F_2 are fresh anomaly events in the two basins during 1992–1997 and 2012–2017 respectively. Yellow/blue shading indicates periods of increased/decreased advection of salt mass.

diffusion and surface salt exchange and call it ΔS_{Ocean} . The contribution of surface freshwater exchange is labelled $\Delta S_{Atm.}$.

For the SPNA, we find that the precipitation plays an important role during F_1 , but not during F_2 . After the F_1 event, the changes in salinity are due to ocean advection and diffusion. For the ESNA, precipitation is important for both the fresh anomaly events. Josey and Marsh (2005) and Holliday et al. (2020) conclude the same for the Iceland basin.

435 4 Discussion

Anomalous low salinity events in the upper 200 m eastern subpolar North Atlantic during the 1990s (F_1) and 2010s (F_2) appear to be qualitatively similar when observed as salinity anomaly signals. Using salt mass budget analysis in the ECCOv4r4 state estimate, we distinguish between the two events and conclude the following:

- 1. The fresh anomaly in the 1990s (F₁) occurs due to anomalous circulation of mean salinity $(v'\bar{S})$. This primarily occurs as the Eulerian component of $v'\bar{S}$ across the southern and western boundaries and the bolus component of the vertical flux of $v'\bar{S}$ across the 200 m horizon in the ESNA (Fig. 9).
- 2. The fresh anomaly during 2012–2017 (F₂) is due to the mean circulation of anomalous salinity ($\bar{v}S'$). This is entirely due to the Eulerian component of $\bar{v}S'$ across the western face of the Iceland Basin (Fig. 9).

447 448 449 3. Vertical diffusive flux across the 200 m depth boundary reverses the freshening after F_1 by increasing the salt mass content during 1995–2000. It also controls the salinity during the period leading up to the F_2 event (Fig. 10).

Therefore, although the two salinity anomalies are similar, they do not share the same mechanisms. The Great Salinity Anomaly (GSA) event (F_1) as described in Belkin (2004) propagates via the Labrador Sea and into the Iceland Basin in the mid 1990s. The fresh anomaly observed in the ESNA in the mid 2010s (F_2) has similar characteristics. The difference between the two is that F_2 is dominated by circulation of anomalous salinity along mean currents into the Iceland Basin and F_1 is dominated by anomalous circulation of mean salinity (Fig. 8).

We note that vertical diffusion has an important role to play in driving the neg-457 ative salinity anomalies to positive salinity anomalies. The 200 m isobath becomes sig-458 nificant in pushing salt flux into the upper layer of the ocean. As we increase the ver-459 tical depth of integration in the salt budget, the vertical exchange of salt via diffusion 460 decreases and the horizontal advective exchange increases. Sanders et al. (2022) conduct 461 a similar study to investigate the 2015 anomalous cooling in the eastern and central SPNA 462 (defined over 50–20°W, 43–63°N, which overlaps with, but extends further west than, 463 our definition of the ESNA, Fig. 2). They analyze the heat budget averaged over the mixed 464 layer using ECCOv4r3 and ECCOv4r4. They observe surface heat loss to initiate and 465 drive the cooling, with advection sustaining the anomaly in the region (as expected from 466 Tesdal and Abernathey (2021)). They also emphasize the role of vertical diffusion across 467 the base of the mixed layer in the re-emergence of the anomaly during summer of 2014. 468 The importance of horizontal advection and vertical diffusion for the temperature anoma-469 lies is similar to the importance of these processes for the salinity anomalies studied here 470 (Fig 9). 471

The surface freshwater fluxes (P-E+R) play a significant role in enhancing the fresh anomalies in the ESNA during the F₂ event. This was also noted by Holliday et al. (2020).

Holliday et al. (2020) state that the 2010s (F_2) event does not share the same char-474 acteristics of the 20th century GSAs. They argue that this event shows no freshening 475 in the Labrador Sea, unlike GSAs (Belkin (2004); Sundby and Drinkwater (2007)). How-476 ever, in the ECCOv4r3 freshwater budget for the Labrador Sea explored by Tesdal and 477 Haine (2020) there is an increased freshwater flux from the Labrador Sea via the Labrador 478 Current. This compensates an increased freshwater flux across the Davis Strait so that 479 little net salinity change occurs in the Labrador Sea over this period. Still, recent pa-480 pers argue that the 2010s (F_2) freshening event resembles a GSA. Specifically, Devana 481 et al. (2021) observe that salinity decrease from November 2015 to March 2017 in the 482 Iceland-Scotland Overflow Water is similar to the freshening observed in the 1990s. Also, 483 Biló et al. (2022) note that the salinity decrease during 2016–2019 in the Irminger Sea 484 $(0.04 \text{ psu year}^{-1})$, following the freshening in the Iceland basin, is among the highest ever 485 recorded. 486

Realistic, physically-consistent state estimates, such as ECCOv4r4 and ASTE_1 used 487 here, are valuable to diagnose mechanisms of large-scale inter-annual salinity and tem-488 perature fluctuations because closed volume, heat, and salt budgets can be constructed. 489 Apart from possible bias (section 3.1), such state estimates have some drawbacks, how-490 ever. State estimates do not include error estimates in the model output fields, for in-491 stance (Piecuch et al., 2017). Tesdal and Haine (2020) address this issue by using ± 2 492 standard deviations of the monthly lateral ECCO fluxes as a substitute for formal un-493 494 certainty estimates. Another limitation is that state estimates typically span a relatively short period (1992–2017 in the case of ECCOv4r4). Therefore, investigations of low-frequency 495 (decadal to centennial) salinity and temperature variability are not possible with these 496 products. 497

For future studies, it is important to investigate such low-frequency fluctuations 498 in the ESNA. Notably, the importance of the meridional overturning circulation, which 499 has not been addressed in this study, and atmospheric forcing needs further attention. 500 Coupled climate models are a promising resource for these studies because they also al-501 low construction of closed volume, heat, and salt budgets, but with much longer dura-502 tions. 503

5 Data Availability Statement 504

We use an open source python package, OceanSpy (https://oceanspy.readthedocs 505 .io; Almansi et al., 2019), to create and analyze synthetic hydrographic sections in the 506 model data. We also use the python package gcm-filters (https://gcm-filters.readthedocs 507 .io; Loose et al., 2022) to apply spatial Gaussian filters for smoothing the AVISO data. 508 The ECCO dataset is publicly available on the SciServer system (Medvedev et al., 2016) 509 and at https://podaac.jpl.nasa.gov/ECCO. ASTE_1 data is publicly available on https:// 510 arcticdata.io/catalog/portals/ASTE. The EN4 data are available at www.metoffice 511 .gov.uk/hadobs/en4/. Python scripts with Juptyer notebooks used for analyzing ECCO 512 budgets can be accessed at https://github.com/asiddi24/Siddiqui_et_al_JGR_Oceans 513 _2024 along with scripts used for generating all figures in the manuscript. 514

Appendix A Salt Budget equations 515

This appendix is useful in understanding the formulation of Eqs. (1, 2, 3). The z^* 516 coordinate is used in ECCOv4r4 to allow for exact tracer conservation (Campin et al., 517 2004) and to improve representation of flow over steep topography (Adcroft & Campin, 518 2004). Physically, z^* allows for variations in the non-linear free surface to be distributed 519 throughout the vertical water column. Using this coordinate, Forget et al. (2015) express 520 the salt conservation equation as 521

$$\frac{\partial(s^*S)}{\partial t} = -\nabla_{z^*} \cdot (s^* S \mathbf{v}_{res}) - \frac{\partial(S w_{res})}{\partial z^*} + s^* (\mathcal{F}_S + D_{\sigma,S} + D_{\perp,S}).$$
(A1)

Integrating Eq. (A1) over a spatial domain V and over time yields a time series for 523 the salt mass, M(t), expressed in Eq. (1). 524

For the anomalous advection term,
$$\mathcal{A}'$$
, consider

$$\mathbf{v}_{res} = \mathbf{\bar{v}}_{res} + \mathbf{v}'_{res}$$

527
$$w_{res} = \bar{w}_{res} + w'_{res}$$

528 $S = \bar{S} + S',$

$$S = S$$

This decomposition implies that

531

$$\bar{\mathbf{v}}_{res}\bar{S} = \overline{\mathbf{v}_{res}S'} - \overline{\mathbf{v}'_{res}S'},$$

 $\bar{w}_{res}\bar{S} = \overline{w_{res}S} - \overline{w'_{res}S'}.$

The anomalous advection term is therefore: 534 535

$$\mathcal{A}' = -\rho_0 \int^t \int_V \nabla_{z^\star} \cdot \left[\left(s^\star (\bar{\mathbf{v}}_{res} S' + \mathbf{v'}_{res} \bar{S} + \mathbf{v'}_{res} S' - \overline{\mathbf{v'}_{res} S'} \right) \right] dV^\star dt^\star - \rho_0 \int^t \int_V \frac{\partial(\bar{w}_{res} S' + w'_{res} \bar{S} + w'_{res} S' - \overline{w'_{res} S'})}{\partial z^\star} dV^\star dt^\star.$$
(A2)

This is rearranged in Eq. (2). 539

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522

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Supporting Information for "Controls on upper ocean salinity variability in the eastern subpolar North Atlantic during 1992–2017"

Ali H. Siddiqui¹, Thomas W. N. Haine¹, An T. Nguyen², Martha Buckley³

¹Department of Earth and Planetary Sciences, Johns Hopkins University, Baltimore, MD, USA

²Oden Institute for Computational Engineering and Sciences, University of Texas at Austin, Austin, TX, USA

³Atmospheric, Oceanic and Earth Sciences, George Mason University, Fairfax, VA, USA

Contents of this file

1. Figures S1 to S3

January 5, 2024, 4:14am

Introduction Figures depicting the spatial variability of annually averaged 0-200 m salinity anomaly (S_{200}) anomalies using the ECCOv4r4, EN4 and ASTE_1 products are presented here.



Figure S1. Annually averaged 0–200 m salinity anomaly (S_{200}) using the ECCOv4r4 January 5, 2024, 4:13am dataset for 1992–2017.



Figure S2. Annually averaged 0–200 m salinity anomaly (S_{200}) using the EN4 dataset for 1990–2018.



Figure S3. Annually averaged 0–200 m salinity anomaly (S_{200}) using the ASTE_1 dataset for 2002–2017.

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