

Response of Mixed Layer Depth Variability to Ocean Eddies and Atmospheric Noise in the Southern Ocean

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December 4, 2023

Ocean Eddies and Atmospheric Noise Drive Mixed Layer Depth Variability in the Southern Ocean

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Key Points:

- Simulations that account for ocean eddies show a significantly deeper average mixed layer depth (MLD) compared to those that do not account for eddy effects.
- In regions with strong eddy activities, reduced atmospheric noise in simulations results in higher MLD variability, driven more by increased ocean current variability than by reduced atmospheric influence.
- Atmospheric noise suppresses ocean's natural variability, particularly diminishing the ocean's inherent influence on MLD variations during ocean-atmosphere coupling.

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Abstract

We investigate the impact of atmospheric noise and model resolution on the relationship between oceanic currents, SST, and mixed layer depth (MLD) in the Southern Ocean, using global climate simulations and interactive ensemble experiments with the NCAR Community Climate System Model version 4.0, at both low (LR) and high spatial resolution (HR) in the ocean. Atmospheric noise is the variability from internal atmospheric dynamics, independent of low-frequency changes or anomalies in boundary conditions or atmospheric composition. The interactive ensemble coupling approach reduces atmospheric noise at the air-sea interface, enabling us to isolate its impact by comparing interactive ensemble simulations to directly coupled (control) runs. We assess the importance of ocean mesoscale currents by contrasting LR and HR simulations. The HR simulations that resolves ocean eddies shows deeper MLD compared to the non-eddy-resolving simulations, which is most likely due to excessive re-stratification by the parameterized eddies in the non-eddy-resolving simulations. In the HR simulations, reduced atmospheric noise amplifies mesoscale ocean currents in the interactive ensemble, which leads to increased SST and MLD variance in the Antarctic Circumpolar Current and Western Boundary Current regions. Furthermore, wind stress feedback interacting with ocean eddies modulates Ekman transport in eddy-resolving simulations, whereas in non-eddy-resolving simulations, Ekman transport is solely driven by atmospheric noise. This study addresses a gap in understanding the importance of oceanic intrinsic variability in driving MLD variability and demonstrating that atmospheric noise suppresses ocean's natural variability during atmosphere-ocean coupling.

Plain Language Summary

This paper investigates how ocean currents and atmospheric noise affect the Southern Ocean's mixed layer depth (MLD) variability. The mixed layer is the ocean's upper layer where temperature and salinity are relatively constant. The study uses computer simulations to analyze the impact of ocean currents and atmospheric noise on MLD variability. The results show that ocean currents significantly impact MLD variability, and the atmospheric noise suppresses the intrinsic oceanic variability in driving the MLD variations during the atmosphere-ocean coupling. The study also highlights the importance of eddy-resolving simulations in capturing the impact of ocean mesoscale currents on MLD variability. The findings of this study can help us better understand the complex interactions between the ocean and atmosphere in the Southern Ocean and improve our ability to predict future changes in the ocean and climate.

1 Introduction**1.1 Ocean Mixed Layer Variability and Ocean Eddies**

The ocean mixed layer plays a vital role in modulating exchanges of heat, fresh-water, and gases such as oxygen and carbon dioxide between the atmosphere and the ocean. Consequently, changes in mixed layer depth (MLD) can have profound implications for climate modeling, affecting exchange rates between the atmosphere, upper ocean, and deep ocean. The entrainment of carbon-rich, oxygen-poor waters into the mixed layer drives the interannual variability of air-sea oxygen and carbon dioxide fluxes in the Southern Ocean (Verdy et al., 2007). Additionally, MLD can modulate air-sea interaction by changing the effective heat capacity of the upper ocean: a shallower MLD results in lower heat capacity and increased sensitivity of sea surface temperature (SST) to surface heat flux. At the same time, a deeper MLD leads to higher heat capacity and reduced SST sensitivity (Tozuka & Cronin, 2014). Therefore, the contribution of surface heat fluxes to surface frontogenesis and frontolysis depends on their gradients and the distribution of MLD (Tozuka et al., 2018; Gao et al., 2023).

68 The Southern Ocean is an integral part of the global overturning circulation since
 69 the upwelling in the Southern Ocean is a vital branch of the circulation (Speer et al., 2000;
 70 Marshall & Speer, 2012). The formation of the deep winter mixed layer has been linked
 71 to the intermediate water masses, Subantarctic Mode Water (SAMW) and Antarctic In-
 72 termediate Water (AAIW), in the Southern Ocean (McCartney, 1977; Rintoul, 2002; Sallée
 73 et al., 2006; J. Holte & Talley, 2009; Lee et al., 2011; J. W. Holte et al., 2012). These
 74 intermediate water masses control the ventilation of the thermocline of the subtropical
 75 gyres in the Southern Hemisphere and contribute to the changes in heat, carbon, and
 76 productivity globally (Sloyan & Rintoul, 2001; Sarmiento et al., 2004; Sallée et al., 2012).
 77 Therefore, it is essential that we understand the MLD variability in the Southern Ocean.

78 1.2 Atmospheric Noise and Interactive Ensembles

79 Atmospheric noise is defined here as the variability due to internal atmospheric dy-
 80 namics independent of low-frequency variability or anomalies in boundary conditions (e.g.,
 81 SST anomalies, soil moisture, snow cover, sea-ice) or in atmospheric composition (i.e.,
 82 aerosols). For conceptual simplicity, the atmospheric forcing can be separated into sig-
 83 nal and noise components: (1) SST-driven atmospheric variability, which is referred to
 84 as "signal", and (2) the stochastic internal dynamics that is not directly driven by Sea
 85 Surface Temperature Anomalies, which is referred to as "atmospheric noise".

86 SST-driven atmospheric variability (the "signal") refers to atmospheric response
 87 to mesoscale SST anomalies (Small et al., 2008). For example, Kirtman et al. (2012) demon-
 88 strates that the correlation between the upward turbulent heat fluxes and SST anoma-
 89 lies in the Southern Ocean is positive in simulations that resolve mesoscale eddies, which
 90 implies that the atmosphere dampens SST variability. As a result, ocean eddy ampli-
 91 tude is enhanced when mesoscale SST-driven atmospheric processes are absent in the
 92 air-sea coupling (Kirtman et al., 2017). Gao et al. (2022) shows that mesoscale currents
 93 induce SST anomalies that are subsequently dampened by atmospheric heat fluxes. SST
 94 anomalies also impact near-surface wind, cloud properties, and rainfall in the Southern
 95 Ocean by affecting turbulence in the atmospheric boundary layer (Frenger et al., 2013).

96 The role of the oceanic signal can be isolated using the interactive ensemble cou-
 97 pling strategy, which reduces the atmospheric noise at the air-sea interface. Barsugli &
 98 Battisti (1998) provides a stochastically forced conceptual model which shows the inter-
 99 action between the atmosphere and ocean in midlatitudes amplifies variance within both
 100 systems and attenuates the energy exchange between them. The paper also mentions that
 101 a principal consequence of this air-sea thermal coupling is the mitigation of thermal damp-
 102 ing in midlatitude regions. Additionally, the article states that this study serves as a foun-
 103 dational framework for future explorations using the interactive ensemble design. The
 104 interactive ensemble approach introduced an ensemble of atmospheric models coupled
 105 to one ocean model to isolate the impact of internal atmospheric variability (Kirtman
 106 & Shukla, 2002). Many previous studies have proven its utility for quantifying how mesoscale
 107 air-sea coupling affects climate predictability (Wu & Kirtman, 2005; Lopez & Kirtman,
 108 2014; Kirtman et al., 2017; Bishop et al., 2017). For example, Kirtman et al. (2017) demon-
 109 strated that the ocean mesoscale activity increases model-estimated climate predictabil-
 110 ity by increasing the dependency of atmospheric internal dynamics on the SST-driven
 111 signal.

112 These studies suggest that neglecting mesoscale air-sea coupling can lead to inac-
 113 curate representation of ocean mesoscale variability, and potential biases in long-term
 114 climate modeling. This study will explore the importance of two main processes for mesoscale
 115 variability in MLD: 1) mesoscale oceanic currents, by comparing eddy-resolving to non-
 116 eddy-resolving ocean models; and 2) internal atmospheric noise at the air-sea interface,
 117 by using the interactive ensemble coupling technique.

Table 1. Global climate model configuration and experiments. Adapted from Kirtman et al. (2017)

Experiments	Ocean	Atmosphere	ensemble size
LRC	1°lat x 1°lon	0.5 °	1
LRIE	1 lat °x 1°lon	0.5 °	1 ocean, 10 atmosphere
HRC	0.1°lat x 0.1°lon	0.5 °	4
HRIE	0.1°lat x 0.1°lon	0.5 °	1 ocean, 10 atmosphere

2 Data and Methodology

2.1 Model Experiments and Data

The global climate simulations are based on the NCAR Community Climate System Model version 4.0 Gent et al. (2011). The atmospheric component is based on the Community Atmospheric Model version 4 and the oceanic part – the Parallel Ocean Program version 2 (Smith et al., 2010). Kirtman & Shukla (2002) and Kirtman et al. (2017) describe the interactive ensemble coupling strategy and details of the experiments used in this study, so here we only provide brief descriptions. We analyze four experiments: two control experiments and two interactive ensembles, at both low- and high spatial resolution in the ocean (Table 1). The standard low-resolution (LRC experiment is a 255-year present-day climate simulation, where the first 40 years are considered spin-up. The LRC experiments use a 1° atmospheric model coupled to ocean and sea ice models with a zonal resolution of 1.2° and meridional resolution that varies from 0.27° at the equator to the 0.54° in the mid-latitudes. The high-resolution (HR) experiments are based on a 4-member ensemble where each ensemble also uses a present-day forcing. The HR experiments use a 0.5° atmospheric model coupled to 0.1° ocean and sea ice models. We analyze the monthly data between 35 °S to 60 °S to include most of the Southern Ocean while avoiding the region with sea ice. We also selected 30-year-long data from year 121 to 150 in each experiment for comparison. We plot an altimetry-derived geographical position of the Subantarctic Front (SAF) and Polar Front (PF) in each figure (Park & Durand, 2019).

2.2 Interactive Ensemble Technique

The intention of the interactive ensemble coupling strategy is to suppress the atmospheric noise at the sea surface. Multiple realizations (ensemble members) of the atmospheric component are coupled to a single realization of the ocean component. The ensemble-mean fluxes of heat, momentum, and freshwater from atmospheric ensemble members are used to force the ocean component, while each atmospheric ensemble member has the same SST forcing produced by the ocean component. This coupling technique is to have the ensemble mean of the atmospheric models continuously interact with the ocean model as the coupled system evolves (Kirtman & Shukla, 2002).

To estimate the coupling strength and identify processes that drive the SST variability, Kirtman et al. (2005), Kirtman et al. (2017) and Zhang & Kirtman (2019) applied SST variance ratio test based on the Hasselmann (1976) hypothesis. Here we provide a brief description of this conceptual model following their examples.

In the control experiment, in which one atmospheric model is coupled to one ocean model, we assume that an oceanic variable O (such as SST) and an atmospheric variable A at time "n+1" are determined by their values at the previous time:

$$A^{n+1} = \alpha O^n + \mu A^n + N \quad (1)$$

$$O^{n+1} = \beta A^n + \delta O^n + P \quad (2)$$

155 where μ and δ represent the memory of the previous state, and α and β are the coupling
 156 coefficients that are bounded between 0 and 1. N and P stand for the uncoupled inter-
 157 nal noise in atmospheric and ocean components, respectively, which is assumed to be Gaus-
 158 sian and white. Ocean noise represents the effects of internal ocean dynamics, which in-
 159 clude mesoscale advection.

160 In the interactive ensemble, multiple atmospheric models (ensemble members) are
 161 coupled to one ocean model. Assuming M is the number of atmospheric models that are
 162 coupled with one ocean model, the equations 1 can be generalized into a set of equations
 163 representing the atmospheric interactive ensembles:

$$A_1^{n+1} = \alpha O^n + \gamma A_1^n + N_1, \quad (3)$$

$$A_2^{n+1} = \alpha O^n + \gamma A_2^n + N_2, \quad (4)$$

$$\dots \quad (5)$$

$$A_M^{n+1} = \alpha O^n + \gamma A_M^n + N_M, \quad (6)$$

$$O^{n+1} = \frac{\beta}{M} \sum_{k=1}^M A_k^n + \delta O^n + P, \quad (7)$$

164 where the atmospheric models are represented by $A_1^n, A_2^n, \dots, A_M^n$ with internal noise
 165 N_1, N_2, \dots, N_M .

166 The ratio between variance in the control and IE simulations can serve to quan-
 167 tify the impact of atmospheric noise and interpret the interactive ensemble. Taking O
 168 in the above theoretical model to be SST, the variance ratio can be diagnosed in terms
 169 of the coupling strength and the amplitude of atmospheric and oceanic noise forcing:

$$\frac{\text{Variance}(SST_{IE})}{\text{Variance}(SST_{CTRL})} = \frac{\beta^2 \sigma_N^2 / M + \sigma_P^2}{\beta^2 \sigma_N^2 + \sigma_P^2} \quad (8)$$

170 where σ_N^2 and σ_P^2 is the variance of the internal atmospheric and oceanic noise, re-
 171 spectively. Following the example of Kirtman et al. (2017) and Zhang & Kirtman (2019),
 172 who applied the variance ratio test solely to SST, we broaden the application of their
 173 theoretical model to include additional oceanic variables, such as MLD and current speed.
 174 The variance ratio test applies to terms like MLD and currents since MLD is implicitly
 175 coupled to the atmosphere and currents are directly coupled via the wind stress. For ex-
 176 ample, we can quantify the impact of atmospheric noise and ocean noise on the MLD
 177 variability, by analyzing the ratio of the MLD variance in the IE to that in the control
 178 experiment:

$$\frac{\text{Variance}(MLD_{IE})}{\text{Variance}(MLD_{CTRL})} = \frac{\beta^2 \sigma_N^2 / M + \sigma_P^2}{\beta^2 \sigma_N^2 + \sigma_P^2} \quad (9)$$

179 where σ_N^2 and σ_P^2 is the variance of internal atmospheric and oceanic noise, respectively.

180 For LRIE and HRIE, there are $M = 10$ atmosphere components coupled to 1 ocean
 181 component (Table 1). Suppose the SST or MLD variance ratio is between 0.1 and 1.0.
 182 In that case, the ocean noise (internal ocean dynamics), coupled feedback, non-linearity,

183 or a combination of these three elements can play a role, and the variability is only partially
 184 forced by the atmospheric noise. If the SST or MLD variance ratio exceeds 1.0,
 185 the reduction of atmospheric noise in the interactive ensemble enhances the oceanic vari-
 186 ance compared to the control experiment. We will conclude that, in this case, unstable
 187 coupling and non-linearity are essential, which means that a linear conceptual model can-
 188 not be used to explain the variability. In this case, the non-linear climate system can be
 189 chaotic in which the noise is "state-dependent", and the unstable coupled feedback af-
 190 fects the variability (Kirtman et al., 2017). Similarly, we can analyze the variance ra-
 191 tio of other oceanic variables, such as the surface current speed and Ekman transport.

192 **3 Results**

193 In this section, we first discuss the relationship between SST, MLD variability, and
 194 ocean currents by the eddy-resolving and non-eddy-resolving simulations. Hereafter, we
 195 use the terms "mesoscale ocean currents" and "ocean eddies" interchangeably, referring
 196 to ocean currents that occur at spatial scales of around 10 to 100 kilometers and tem-
 197 poral scales of days to months. We then explore the importance of atmospheric noise in
 198 SST and MLD variability by comparing the control simulations to the interactive ensem-
 199 ble.

200 **3.1 Importance of ocean currents in SST variability**

201 In high-resolution (HR) experiments, the variance of sea surface temperature (SST)
 202 is notably higher, particularly in Antarctic Circumpolar Current (ACC) regions and west-
 203 ern boundary current regions, such as the Agulhas Current areas, when compared to the
 204 results of low-resolution (LR) experiments. There are several reasons why the HR ex-
 205 periments produce significantly higher and more realistic SST variability compare to the
 206 LR experiments: 1) heat advection by ocean eddies creates SSTA in regions with strong
 207 oceanic currents, which demonstrates that resolving Southern Ocean eddies is critical
 208 for getting the SST variability right; 2) MLD variability modulates the relationship be-
 209 tween SSTA and the eddy advection of heat. The MLD in the HR experiment is deeper
 210 than in LRC and exhibits higher variance in general, which further enhances the impor-
 211 tance of oceanic advection in SST variability. These effects improve the difference in SSTA
 212 variance between LR and HR experiments. Besides, these effects possibly increase the
 213 HRIE/HRC variance ratio by enhancing the importance of ocean dynamics compared
 214 to atmospheric noise.

215 **3.2 Importance of ocean currents in MLD**

216 We first discuss the relationship between MLD and SST. The climatological annual
 217 cycle has not been removed from these values, and the Variations are heavily in-
 218 fluenced by seasonal variability. The variations of the MLD are negatively correlated with
 219 SST in most of the regions in the Southern Ocean (Fig.1 a-b), which means that cooler
 220 SST corresponds to deeper MLD. This is a relationship we can expect in seasonal vari-
 221 ations in MLD, which deepens in winter and shoals in summer. However, this simple one-
 222 dimensional relationship breaks down in regions with strong oceanic currents, such as
 223 the Antarctic Circumpolar Current (ACC) and Western Boundary Currents (WBC, Fig.1
 224 b). This breakdown is more pronounced in the HRC simulation than in LRC, largely
 225 because of fast time-mean oceanic currents (Fig.1) and stronger mesoscale variability (Fig.2
 226 e-f) in HRC. The surface Eddy Kinetic Energy (EKE) is calculated as $EKE = 1/2(U'^2 + V'^2)^{1/2}$,
 227 where U' and V' are the velocity departure from the 30-year-mean surface current
 228 speed. Note that most of the variability in velocities comes from mesoscale currents,
 229 although the seasonal anomalies in current speed will also contribute to EKE.

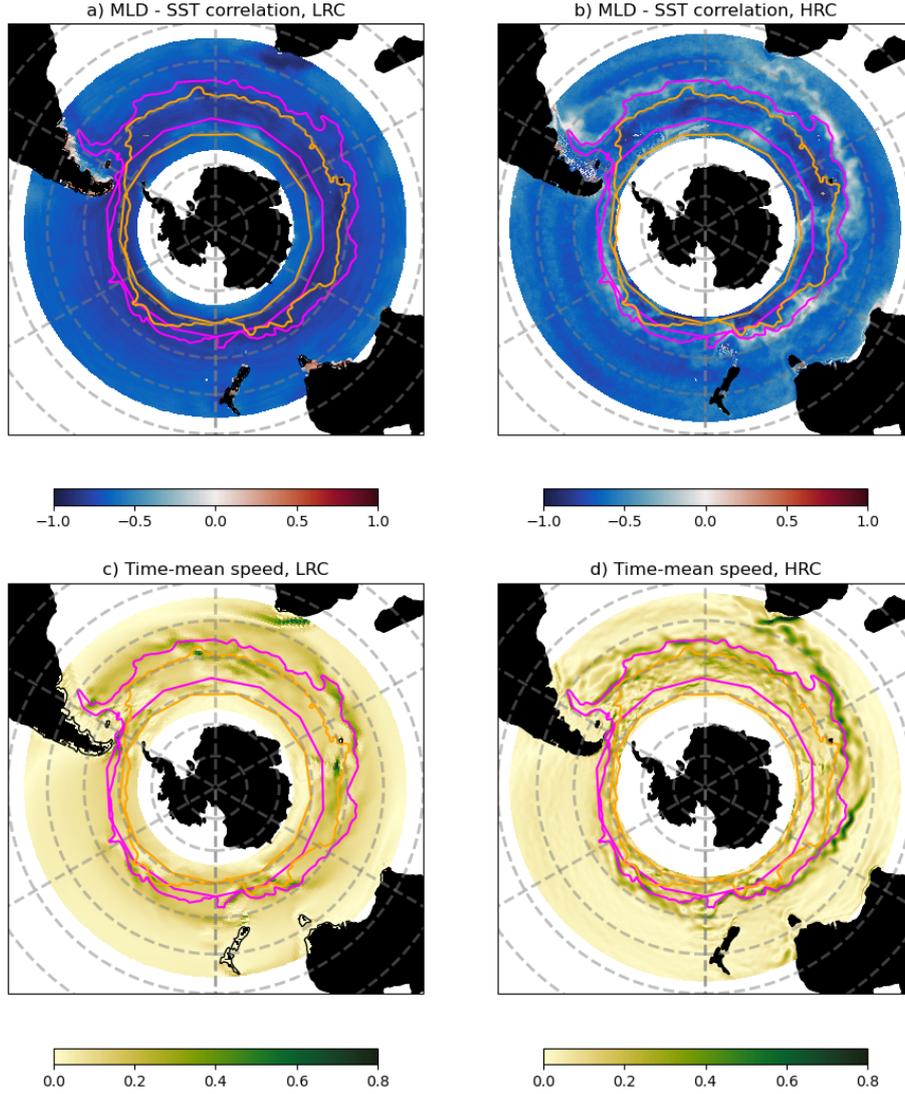


Figure 1. The correlation coefficient between the SST and MLD in a) LRC and b) HRC. The time-mean surface current speed in c) LRC and d) HRC, from year 121 to 150. Polar Front (PF, orange line) and Subantarctic Front (SAF, magenta line).

230 The weak relation between SST and MLD in the ACC region suggests that a one-
 231 dimensional atmospheric forced mixed layer model Kraus et al. (1967) does not apply
 232 in the regions of strong advection. This conclusion is in agreement with the findings of
 233 Gao et al. (2023) who concluded that the buoyancy advection shear by oceanic currents
 234 generally counteracts the atmospheric buoyancy forcing in driving the mixed layer vari-
 235 ability.

236 The time-mean MLD in HRC is significantly deeper than that in LRC in most South-
 237 ernn (Fig. 2). This is consistent with Lee et al. (2011), who discovered that the winter
 238 MLD in eddy-permitting ocean simulations aligns closely with observed data, while the
 239 winter MLD in coarse-resolution ocean models tends to be too shallow. The most sig-
 240 nificant disparities were identified within the Agulhas Current system, where a higher

241 surface heat loss over the Agulhas Return Current and a deeper mixed layer were ob-
 242 served in eddy-permitting simulations. In this study, we also find significant differences
 243 between the HRC and LRC experiments in ACC and WBC regions, including the Ag-
 244 ulhas Current system, part of the Brazil Current, and the Brazil–Malvinas Confluence
 245 (Fig.2a-b). In addition to the mean current strength, the difference between the HRC
 246 and LRC experiments is also dependent on EKE: the MLD in the HRC simulation is shall-
 247 lower in regions of higher EKE (WBC and ACC regions) and deeper in the other areas
 248 (Fig.2e-f).

249 It is unclear if we can explain the deepening of the MLD in HRC by the action of
 250 mesoscale eddies alone. On average, eddies are assumed to re-stratify the base of the mixed
 251 layer (Fox-Kemper et al., 2008; Fox-Kemper & Ferrari, 2008). At the same time, Gao
 252 et al. (2022) demonstrates that mesoscale buoyancy advection can also deepen the mixed
 253 layer, counteracting the atmospheric forcing. Additionally, mesoscale eddies are param-
 254 eterized with the Gent & McWilliams (1990) scheme (hereafter "GM") in the non-eddy-
 255 resolving LRC and LRIE experiments and the model used by Lee et al. (2011). There-
 256 fore, it is possible that the GM parameterization overestimates the re-stratifying effects
 257 of mesoscale buoyancy advection. It is also worth noting that, in most of the Southern
 258 Ocean, the SST in HRC is significantly warmer than that in LRC Kirtman et al. (2012)
 259 and cannot explain the deeper MLD in HRC. Therefore, it is sensible to assume the im-
 260 portance of ocean buoyancy advection in driving the MLD variability.

261 The reduction of atmospheric noise in LRIE relative to LRC mostly leads to deep-
 262 ening of MLD north of the SAF and to shoaling of MLD south of the SAF (Fig.2a). There-
 263 fore, the southerly slope of MLD is generally reduced due to atmospheric noise. The time-
 264 mean MLD in LRIE is, in contrast, shallower than in the HRC experiment in most re-
 265 gions of the Southern Ocean, except the Agulhas Current region, part of the Brazil Cur-
 266 rent, and the Brazil–Malvinas Confluence region (Fig.2b). Gao et al. (2023) concludes
 267 that the atmospheric forcing and mixing induce MLD variability, while the oceanic ad-
 268 vection of buoyancy essentially balances these atmospheric effects. With the reduction
 269 of atmospheric forcing, the strong buoyancy advection in ACC and WBC regions can
 270 become unbalanced and lead to the deepening of MLD.

271 **3.3 Importance of atmospheric noise in SST variability**

272 The response of SST variability to the reduction in the atmospheric noise in the
 273 interactive ensemble simulations differs between HR and LR simulations. Here SST anoma-
 274 lies (SSTAs) are defined as the departure from the monthly climatology. The absence
 275 of eddies in LRC leads to the lower SSTA variance (Fig.4) than in HRC, which is not
 276 surprising given weaker buoyancy advection in the LR case. Consistent with this result,
 277 the variance ratio (Fig.3) is also lower than in the HR simulations, which suggests that
 278 SSTA are primarily caused by the anomalies in the atmospheric forcing, which are sig-
 279 nificantly reduced in LRIE. The following analysis suggests that the importance of at-
 280 mospheric noise is overestimated in the LR simulations.

281 The noise reduction in LRIE relative to HRC enhances the SST variability in the
 282 ACC and WBC regions, where SSTA variance ratio exceeds 1.0 (Fig.3). This result is
 283 consistent with the findings in Kirtman et al. (2017). The results suggests that the SST
 284 variability in these regions is attributed to the intrinsic ocean dynamics, unstable cou-
 285 pled feedback and nonlinear dynamics. This is consistent with Gao et al. (2022) who found
 286 that the SSTA variability in the Southern Ocean in the regions with fast oceanic cur-
 287 rents is driven by the intrinsic ocean dynamics rather than the atmospheric forcing. In
 288 other regions, the variance ratio is mostly between 0.5 and 1.0, which suggests the SSTA
 289 variability is partially driven by atmospheric noise and partially by coupled feedback,
 290 ocean eddies or non-linearity, or a combination of the three. In addition, the SSTA vari-
 291 ance ratio can exceed 2.0 in the regions south of 60S, and these values may be exagger-

292 ated by biases in the Antarctic sea ice and excessive westerly winds in CCSM4 (Kirt-
 293 man et al., 2017). The enhanced SST variance in the interactive ensemble simulations
 294 is consistent with the increase in the upper-ocean currents, which is discussed in the next
 295 section.

296 3.4 Importance of atmospheric noise in MLD variability

297 The air-sea interaction over Southern Ocean eddies can produce substantial MLD
 298 variability, which is underestimated in the non-eddy-resolving ocean models (Fig.5). MLD
 299 anomalies (MLDAs) are defined as the departure from the monthly climatology. Gao et
 300 al. (2023) suggests that oceanic mesoscale currents compensate the atmosphere-induced
 301 variations in MLD: while the atmospheric forcing and oceanic vertical mixing induce the
 302 MLD variability, the oceanic advection of buoyancy counteracts these atmospheric ef-
 303 fects. The anomalies in atmospheric fluxes result either from intrinsic atmospheric vari-
 304 ability or from an SST-forced response. The analysis of this section will help to estimate
 305 the relative importance of these two processes because the IE reduces the former, inter-
 306 nal source for variability.

307 In both the HRIE and LRIE experiments, the MLDA variance is suppressed in most
 308 regions because of the reduced atmospheric noise (Fig.6). The changes are nevertheless
 309 dramatically different between the LR and HR simulations. In the LR experiments, the
 310 MLDA variance ratio is below 0.5 in most of the Southern Ocean (Fig.6a), which sug-
 311 gests that MLDA variability is partially forced by the atmosphere noise, and partially
 312 results from the coupled feedback, non-linearity and ocean noise. In contrast, in the HR
 313 experiments, the MLDA variance ratio is not only overall higher, but exceeds 1.0 in the
 314 ACC and WBC regions (Fig.6b).

315 In the ACC and WBC regions, the MLDA variance in HRIE is enhanced due to
 316 the reduced atmospheric noise. The SSTA variance is also enhanced there (Fig.3), how-
 317 ever, the correlation between the SST and MLD anomalies is low in these regions (Fig.1),
 318 which means that the increased SST variance cannot explain the increase in MLD vari-
 319 ance. We can, however, explain the increased MLDA variance with our findings in Gao
 320 et al. (2023). There, we concluded that while the atmospheric forcing and oceanic ver-
 321 tical mixing induce MLD variability, the oceanic advection of buoyancy counteracts these
 322 atmospheric effects. Results of Gao et al. (2023) further show that when mesoscale anoma-
 323 lies are removed from the surface fluxes of heat and momentum, the MLD variability can
 324 increase, and this effect is most pronounced in local winter. These conclusions suggest
 325 that when the atmospheric stochastic forcing is suppressed in HRIE, the oceanic advec-
 326 tion can become partially unbalanced and enhance MLD variance.

327 More variance ratio test on the surface current speed. In the LR experiment (Fig.7a),
 328 the surface current variance ratio is mostly between 0.1 and 1.0, which suggests a com-
 329 bination of the atmospheric noise and internal dynamics drives surface present variabil-
 330 ity. In the HR experiments, in contrast, the current speed variance ratio is larger than
 331 1.0 in most of the Southern Ocean (Fig.7b). This means the speed variance in HRIE in-
 332 creases due to the reduced atmospheric noise. The increase is especially pronounced in
 333 the Southern Indian and Atlantic sectors of the Southern Ocean, indicating the critical
 334 role of unstable coupling and nonlinearity. In contrast, the ratio is less than 1.0 in the
 335 Pacific sector (Fig.7). This pattern is similar to the MLDA variance ratio, which is also
 336 lower in the Pacific sector of the Southern Ocean (Fig.6). The similarity suggests the im-
 337 portance of ocean advection in MLD variance. The current speed and EKE are not as
 338 strong in the Southern Pacific sector (Fig.1 and Fig.2), and the ocean-atmosphere in-
 339 teractions have more considerable relative importance.

340 Compared to the SSTA variance ratio, HR experiments exhibit smaller regions where
 341 the MLDA variance ratio exceeds 1.0. This difference indicates that atmospheric noise
 342 plays a bigger role in driving MLD variability than SST variability. Atmospheric noise,

343 such as wind and air temperature variations, can significantly impact Mixed Layer Depth
 344 (MLD) more than Sea Surface Temperature (SST) due to the direct and immediate im-
 345 pact of wind stirring. This process mixes the ocean’s surface layer to varying depths, lead-
 346 ing to significant MLD variability. While SST is also affected by atmospheric conditions,
 347 the high heat capacity of the ocean means that a larger amount of heat exchange is re-
 348 quired to alter its temperature significantly. Therefore, atmospheric noise tends to drive
 349 more variability in MLD than SST.

350 This heightened variability reflects a non-linear and unstable coupling. Essentially,
 351 this means that the relationship between the different forces (ocean currents, atmospheric
 352 noise, etc.) is not straightforward. Instead, these forces interact in complex, unpredictable
 353 ways, which is in nature of ‘non-linearity.’ ‘Unstable coupling’ indicates that the rela-
 354 tionship between these elements isn’t stable or consistent and can change rapidly.

355 This variability could be due to the increased Mixed Layer Depth (MLD), induced
 356 by the rise in atmospheric noise and atmosphere-driven mixing (Fig.6). As the MLD in-
 357 creases, the inertia (or resistance to change) in the upper ocean also increases. This could
 358 then lead to a decrease in ocean current variability. The state-dependence of atmospheric
 359 noise and the chaotic nature of the system could offer possible explanations for these ob-
 360 served phenomena.

361 This difference is, however, reversed in the LR simulations. The atmospheric noise
 362 mainly drives the SST variability in the LR experiments whereas the MLDA variance
 363 is due to a combination of the atmospheric noise, oceanic dynamics, and coupled feed-
 364 backs. Weaker oceanic currents and air-sea feedback onto SSTA in the LR simulations
 365 explain this.

366 3.5 Ocean Eddies Modulating Ekman Transport

367 To get further insight into the amplification of ocean currents in the absence of at-
 368 mospheric noise, we investigate the response of Ekman transport to atmospheric forc-
 369 ing. The Ekman transport velocity (unit: m^2/s) in the u and v directions are calculated
 370 as below:

$$U_{Ekman} = \frac{\tau^y}{\rho_0 f}, \quad (10)$$

$$V_{Ekman} = \frac{-\tau^x}{\rho_0 f} \quad (11)$$

$$(12)$$

371 where τ^y and τ^x is the meridional and zonal wind stress at the sea surface, respec-
 372 tively. ρ_0 is the reference density ($1025 \text{ kg}/m^3$) and f is the Coriolis parameter. We next
 373 calculate the variance in the magnitude of the Ekman transport $\sqrt{U_{Ekman}^2 + V_{Ekman}^2}$
 374 and the variance ratio.

375 Fig.8 emphasizes the role of wind stress feedback and Southern Ocean eddies in shap-
 376 ing Ekman transport in our eddy-resolving experiments: the figure displays the influ-
 377 ence of eddies on Ekman transport, as illustrated through the contrast between LR and
 378 HR experiments. Importantly, the figure shows a decrease in Ekman variance in response
 379 to reduced atmospheric noise, a trend that is more significant in HRIE than in LRIE.
 380 Fig.8 shows the impact of mesoscale currents on Ekman transport by comparing the LR
 381 and HR experiments. In LR experiment, the variance ratio is below 0.1, which indicates
 382 atmospheric noise alone drives the variability of the Ekman current speed. In HR ex-
 383 periment, the variance ratio is between 0.1 and 0.5, which suggests the Ekman transport
 384 variance is partially forced by the atmosphere noise, and partially by coupled feedback

385 and internal variability. For example, the presence of ocean eddies can modify South-
 386 ern Ocean winds at the air-sea interface, which in turn alter the Ekman transport (Small
 387 et al., 2008; Frenger et al., 2013; Perlin et al., 2020). Mesoscale variability can also have
 388 an impact on wind stress through surface current speed correction: surface ocean cur-
 389 rents modulate turbulent air-sea exchanges by changing the velocity shear between the
 390 atmosphere and oceanic surface. Current speed correction to the wind stress acts as a
 391 “top drag” (Dewar & Flierl, 1987; Duhaut & Straub, 2006; Gaube et al., 2015), since the
 392 enhanced Ekman pumping leads to relaxation of the thermocline. Importantly, the sen-
 393 sitivity of Ekman currents to atmospheric noise cannot explain amplification of surface
 394 currents in the HRIE simulations.

395 4 Summary and Discussion

396 The objective of this study is explore the impact of atmospheric noise and model
 397 resolution on the relationship between oceanic currents, SST and MLD variability. This
 398 study utilizes the interactive ensemble coupling method, which reduces the atmospheric
 399 in both low- and high-resolution simulations. We analyzed the MLD variability and ex-
 400 amined a variance ratio between the IE and control experiments, which allowed us to
 401 estimate the relative importance of the atmospheric noise in the MLD variability. Based
 402 on the analysis in Gao et al. (2022) and Gao et al. (2023), we expect variance based on
 403 monthly means to be a good measure of mesoscale anomalies. The impact of the strong
 404 oceanic currents, fronts and eddies were further assessed by comparing the eddy-resolving
 405 (HR) experiments to non-eddy-resolving (LR) experiments.

406 The strong negative correlation between SST and MLD breaks down in regions with
 407 strong large-scale currents and mesoscale activity, namely within the ACC fronts and
 408 in the WBC regions. The time-mean MLD is also significantly deeper in the presence
 409 of eddies even though SST in HR is warmer than in LR (Kirtman et al., 2012). Simi-
 410 larly, Lee et al. (2011) found the winter MLD in coarse-resolution coupled ocean model
 411 is too shallow. The MLD difference also depends on EKE: the MLD in HR simulations
 412 is shallower in regions of higher EKE and deeper elsewhere. These results suggest a cru-
 413 cial role of mesoscale ocean currents in driving SST and MLD anomalies. Eddies are widely
 414 assumed to re-stratify the ocean, but our results indicate that the GM parameterization
 415 may overestimate the re-stratifying role of ocean eddies. Gao et al. (2023) further demon-
 416 strates that the effects of eddies on MLD are more complex, and the corresponding buoy-
 417 ancy advection can even de-stratify the ocean below the mixed layer and deepen the MLD.

418 The results further demonstrate that SST variability is mainly driven by oceanic
 419 processes rather than atmospheric noise in the ACC and WBC regions. This result is
 420 consistent with Gao et al. (2022): the SST variability is driven by intrinsic ocean dynam-
 421 ics instead of atmospheric forcing. In such “quiet” regions of the Southern Ocean as the
 422 Pacific sector, the role of the atmosphere is more significant, and the SST variability is
 423 jointly driven by atmospheric noise and oceanic internal dynamics. Significantly, the re-
 424 duction of the atmospheric noise in HRIE even enhances the SST variability in the ACC
 425 and WBC regions. However, it is still unclear why the SST variability is enhanced in these
 426 regions, and this topic deserves further investigation.

427 The atmospheric noise and the Southern Ocean eddies both control MLD variabil-
 428 ity. The differences between MLD variance ratio in LR and HR experiments demonstrate
 429 the importance of intrinsic ocean dynamics, especially in the ACC and WBC regions.
 430 Compared to the SST variability, however, the atmospheric forcing plays a more signif-
 431 icant role in driving MLD variability in the HR experiments than in the LR runs. This
 432 result is consistent with Gao et al. (2023): in the Southern Ocean, SST variability is mainly
 433 driven by the intrinsic oceanic dynamic, while the MLD variability is caused by both at-
 434 mospheric forcing and oceanic dynamics. Consistent with previous studies such as Zhang
 435 & Kirtman (2019), the atmospheric noise suppresses the upper-oceanic variability, and

436 the upper-ocean mesoscale variability in MLD, SST, and surface currents intensify in HRIE.
 437 Given a weak correlation between SST and MLD anomalies in these regions, it is nat-
 438 ural to assume that the increase in SST and MLD variance in the absence of atmospheric
 439 noise are both caused by stronger oceanic currents. Gao et al. (2023) concludes that while
 440 the atmospheric forcing and oceanic vertical mixing induce MLD variability, the oceanic
 441 advection of buoyancy counteracts these atmospheric effects. Consistent with this find-
 442 ing, when the atmospheric stochastic forcing is suppressed in HRIE, the oceanic advec-
 443 tion becomes unbalanced and can thus act to enhance MLD variance. Schneider et al.
 444 (2023) points out that ocean dynamics plays a minimal role in SO decadal variability
 445 in non-eddy-resolving models, which supports our conclusion that ocean mesoscale dy-
 446 namics have a large role in SO variability. It is still, however, unclear why the variance
 447 in oceanic currents increase in HRIE, and this question deserves further investigation.

448 The suppression of surface currents by atmospheric noise cannot be explained by
 449 changes in the Ekman currents alone. The variance in the Ekman transport is decreased
 450 in HRIE and LRIE simulations, but the reduction is different between the LR and HR
 451 simulations for several reasons. The Ekman transport is modulated by the wind stress
 452 feedback over ocean eddies (SST anomalies) in the HR experiments, while the Ekman
 453 transport is driven solely by atmospheric noise in the LR experiments. wind speed af-
 454 fected by SST anomalies, which created wind stress feedback to the ocean (Seo et al. (2016)).
 455 In the eddy-resolving experiments, the presence of ocean eddies can modify winds at the
 456 air-sea interface, which in turn alters the Ekman transport in the Southern Ocean (Small
 457 et al., 2008; Frenger et al., 2013; Perlin et al., 2020). The eddies also have an impact on
 458 the wind stress through surface current speed correction (Dewar & Flierl, 1987; O’Neill
 459 et al., 2003; Duhaut & Straub, 2006; Gaube et al., 2015). In other words, ocean eddies
 460 affect Ekman transport by modifying the wind forcing over the sea surface. Neglecting
 461 the eddy-wind coupling in non-eddy-resolving experiments may lead to imbalance in the
 462 zonal-mean steady-state circulation, since eddy-induced circulation compensates for the
 463 Ekman transport in the eddy-resolving models in the Southern Ocean (Abernathey et
 464 al., 2011; Marshall & Speer, 2012).

465 By examining the eddy-resolving and non-eddy-resolving experiments, we found
 466 that the air-sea interaction at mesoscale can cause significant differences in the variabil-
 467 ity of SST, MLD ocean currents, and Ekman transport. By comparing the interactive
 468 ensemble and control experiments, we conclude that non-eddy-resolving ocean models
 469 oftentimes overestimate the role of atmospheric noise and overlook the importance of ocean
 470 dynamics. In eddy-resolving ocean models, the ocean eddies, air-sea coupled feedback
 471 and non-linearity become more important to the mixed layer dynamics. We can further
 472 assess the eddy-induced effects on atmosphere-ocean coupling by using experiments with
 473 multiple ocean ensemble members coupled to one atmospheric component. Although run-
 474 ning 10s global ocean models simultaneously is an ambitious task, we believe this is pos-
 475 sible in the near future considering the fast development of computational power nowa-
 476 days.

477 5 Open Research

478 The numerical model and data are available upon request. The Python code and
 479 jupyter notebook used to produce the results of this study are shared through the GitHub
 480 repository at https://github.com/yugaophd/SO_CCSM4.

481 Acknowledgments

482 This study is supported by the National Science Foundation (NSF) Research, USA, Award
 483 No. 1559151 and 1849990 and the National Aeronautics and Space Administration (NASA)
 484 grant No. 80NSSC20K1136. We would also like to acknowledge the following grants: Na-
 485 tional Oceanic and Atmospheric Administration (NOAA) grant NA18OAR4310293, NA15OAR4320064,

486 NA20OAR4320472, and NSF grant No. OCE1419569, OCE1559151 and AGS2241538.
 487 Benjamin P. Kirman is the William R Middelthun Chair of Earth Sciences and is grate-
 488 ful for the associated support. We acknowledge the computing resources provided by the
 489 University of Miami’s Center of Computational Science and the high-performance com-
 490 puting support from Cheyenne provided by NCAR’s Computational and Information Sys-
 491 tems Laboratory, sponsored by NSF. We thank Dr. Michael Spall of Woods Hole Oceanog-
 492 raphic Institution and Dr. Lisa Beal of the University of Miami for their helpful com-
 493 ments and suggestions. We also thank the editors and reviewers for their thoughtful com-
 494 ments.

495 References

- 496 Abernathey, R., Marshall, J., & Ferreira, D. (2011). The dependence of south-
 497 ern ocean meridional overturning on wind stress. *Journal of Physical Oceanogra-*
 498 *phy*, *41*(12), 2261 - 2278. Retrieved from [https://journals.ametsoc.org/view/](https://journals.ametsoc.org/view/journals/phoc/41/12/jpo-d-11-023.1.xml)
 499 journals/phoc/41/12/jpo-d-11-023.1.xml doi: 10.1175/JPO-D-11-023.1
- 500 Barsugli, J. J., & Battisti, D. S. (1998, February). The Basic Effects of At-
 501 mosphere–Ocean Thermal Coupling on Midlatitude Variability. *Journal*
 502 *of the Atmospheric Sciences*, *55*(4), 477–493. Retrieved 2021-10-18, from
 503 [http://journals.ametsoc.org/doi/10.1175/1520-0469\(1998\)055<0477:](http://journals.ametsoc.org/doi/10.1175/1520-0469(1998)055<0477:TBEOAO>2.0.CO;2)
 504 [TBEOAO>2.0.CO;2](http://journals.ametsoc.org/doi/10.1175/1520-0469(1998)055(0477:TBEOAO)2.0.CO;2) doi: 10.1175/1520-0469(1998)055(0477:TBEOAO)2.0.CO;2
- 505 Bishop, S. P., Small, R. J., Bryan, F. O., & Tomas, R. A. (2017, October). Scale
 506 Dependence of Midlatitude Air–Sea Interaction. *Journal of Climate*, *30*(20),
 507 8207–8221. Retrieved 2020-06-19, from [https://journals.ametsoc.org/](https://journals.ametsoc.org/jcli/article/30/20/8207/33173/Scale-Dependence-of-Midlatitude-AirSea-Interaction)
 508 [jcli/article/30/20/8207/33173/Scale-Dependence-of-Midlatitude-AirSea](https://journals.ametsoc.org/jcli/article/30/20/8207/33173/Scale-Dependence-of-Midlatitude-AirSea-Interaction)
 509 [-Interaction](https://journals.ametsoc.org/jcli/article/30/20/8207/33173/Scale-Dependence-of-Midlatitude-AirSea-Interaction) doi: 10.1175/JCLI-D-17-0159.1
- 510 Dewar, W. K., & Flierl, G. R. (1987). Some effects of the wind on rings. *Journal of*
 511 *Physical Oceanography*, *17*(10), 1653 - 1667. Retrieved from [https://journals](https://journals.ametsoc.org/view/journals/phoc/17/10/1520-0485_1987_017_1653_seotwo_2_0_co_2.xml)
 512 [.ametsoc.org/view/journals/phoc/17/10/1520-0485_1987_017_1653_seotwo_2](https://journals.ametsoc.org/view/journals/phoc/17/10/1520-0485_1987_017_1653_seotwo_2_0_co_2.xml)
 513 [_0_co_2.xml](https://journals.ametsoc.org/view/journals/phoc/17/10/1520-0485_1987_017_1653_seotwo_2_0_co_2.xml) doi: 10.1175/1520-0485(1987)017<1653:SEOTWO>2.0.CO;2
- 514 Duhaut, T. H. A., & Straub, D. N. (2006). Wind stress dependence on ocean
 515 surface velocity: Implications for mechanical energy input to ocean circula-
 516 tion. *Journal of Physical Oceanography*, *36*(2), 202 - 211. Retrieved from
 517 <https://journals.ametsoc.org/view/journals/phoc/36/2/jpo2842.1.xml>
 518 doi: 10.1175/JPO2842.1
- 519 Fox-Kemper, B., & Ferrari, R. (2008). Parameterization of mixed layer eddies. part
 520 ii: Prognosis and impact. *Journal of Physical Oceanography*, *38*(6), 1166 - 1179.
 521 Retrieved from [https://journals.ametsoc.org/view/journals/phoc/38/6/](https://journals.ametsoc.org/view/journals/phoc/38/6/2007jpo3788.1.xml)
 522 [2007jpo3788.1.xml](https://journals.ametsoc.org/view/journals/phoc/38/6/2007jpo3788.1.xml) doi: 10.1175/2007JPO3788.1
- 523 Fox-Kemper, B., Ferrari, R., & Hallberg, R. (2008). Parameterization of mixed layer
 524 eddies. part i: Theory and diagnosis. *Journal of Physical Oceanography*, *38*(6),
 525 1145 - 1165. Retrieved from [https://journals.ametsoc.org/view/journals/](https://journals.ametsoc.org/view/journals/phoc/38/6/2007jpo3792.1.xml)
 526 [phoc/38/6/2007jpo3792.1.xml](https://journals.ametsoc.org/view/journals/phoc/38/6/2007jpo3792.1.xml) doi: 10.1175/2007JPO3792.1
- 527 Frenger, I., Gruber, N., Knutti, R., & Münnich, M. (2013, August). Imprint of
 528 Southern Ocean eddies on winds, clouds and rainfall. *Nature Geoscience*, *6*(8),
 529 608–612. Retrieved 2021-10-31, from [http://www.nature.com/articles/](http://www.nature.com/articles/ngeo1863)
 530 [ngeo1863](http://www.nature.com/articles/ngeo1863) doi: 10.1038/ngeo1863
- 531 Gao, Y., Kamenkovich, I., & Perlin, N. (2023). Origins of mesoscale mixed-layer
 532 depth variability in the southern ocean. *Ocean Science*, *19*(3), 615–627. Retrieved
 533 from <https://os.copernicus.org/articles/19/615/2023/> doi: 10.5194/os-19-
 534 -615-2023
- 535 Gao, Y., Kamenkovich, I., Perlin, N., & Kirtman, B. (2022). Oceanic advec-
 536 tion controls mesoscale mixed layer heat budget and air–sea heat exchange in
 537 the southern ocean. *Journal of Physical Oceanography*, *52*(4), 537 - 555. Re-

- 538 trieved from [https://journals.ametsoc.org/view/journals/phoc/52/4/](https://journals.ametsoc.org/view/journals/phoc/52/4/JPO-D-21-0063.1.xml)
539 JPO-D-21-0063.1.xml doi: <https://doi.org/10.1175/JPO-D-21-0063.1>
- 540 Gaube, P., Chelton, D. B., Samelson, R. M., Schlax, M. G., & O'Neill, L. W.
541 (2015). Satellite observations of mesoscale eddy-induced Ekman pumping.
542 *Journal of Physical Oceanography*, *45*(1), 104 - 132. Retrieved from [https://](https://journals.ametsoc.org/view/journals/phoc/45/1/jpo-d-14-0032.1.xml)
543 journals.ametsoc.org/view/journals/phoc/45/1/jpo-d-14-0032.1.xml doi:
544 10.1175/JPO-D-14-0032.1
- 545 Gent, P. R., Danabasoglu, G., Donner, L. J., Holland, M. M., Hunke, E. C., Jayne,
546 S. R., ... Zhang, M. (2011, October). The Community Climate System Model
547 Version 4. *Journal of Climate*, *24*(19), 4973–4991. Retrieved 2021-02-26,
548 from <https://journals.ametsoc.org/doi/10.1175/2011JCLI4083.1> doi:
549 10.1175/2011JCLI4083.1
- 550 Gent, P. R., & McWilliams, J. C. (1990). Isopycnal mixing in ocean circulation
551 models. *Journal of Physical Oceanography*, *20*(1), 150 - 155. doi: 10.1175/1520
552 -0485(1990)020<0150:IMIOCM>2.0.CO;2
- 553 Hasselmann, K. (1976, December). Stochastic climate models Part I. Theory. *Tellus*,
554 *28*(6), 473–485. Retrieved 2021-10-21, from [http://tellusa.net/index.php/](http://tellusa.net/index.php/tellusa/article/view/11316)
555 [tellusa/article/view/11316](http://tellusa.net/index.php/tellusa/article/view/11316) doi: 10.1111/j.2153-3490.1976.tb00696.x
- 556 Holte, J., & Talley, L. (2009, September). A New Algorithm for Finding Mixed
557 Layer Depths with Applications to Argo Data and Subantarctic Mode Water For-
558 mation. *Journal of Atmospheric and Oceanic Technology*, *26*(9), 1920–1939.
559 Retrieved 2021-10-08, from [http://journals.ametsoc.org/doi/10.1175/](http://journals.ametsoc.org/doi/10.1175/2009JTECH0543.1)
560 2009JTECH0543.1 doi: 10.1175/2009JTECH0543.1
- 561 Holte, J. W., Talley, L. D., Chereskin, T. K., & Sloyan, B. M. (2012, March).
562 The role of air-sea fluxes in Subantarctic Mode Water formation: SAMW For-
563 mation. *Journal of Geophysical Research: Oceans*, *117*(C3), n/a–n/a. Re-
564 trieved 2020-06-18, from <http://doi.wiley.com/10.1029/2011JC007798> doi:
565 10.1029/2011JC007798
- 566 Kirtman, B. P., Bitz, C., Bryan, F., Collins, W., Dennis, J., Hearn, N., ... Verten-
567 stein, M. (2012, September). Impact of ocean model resolution on CCSM
568 climate simulations. *Climate Dynamics*, *39*(6), 1303–1328. Retrieved 2021-
569 01-07, from <http://link.springer.com/10.1007/s00382-012-1500-3> doi:
570 10.1007/s00382-012-1500-3
- 571 Kirtman, B. P., Pegion, K., & Kinter, S. M. (2005, July). Internal Atmospheric
572 Dynamics and Tropical Indo-Pacific Climate Variability. *Journal of the At-*
573 *mospheric Sciences*, *62*(7), 2220–2233. Retrieved 2021-10-21, from [https://](https://journals.ametsoc.org/doi/10.1175/JAS3449.1)
574 journals.ametsoc.org/doi/10.1175/JAS3449.1 doi: 10.1175/JAS3449.1
- 575 Kirtman, B. P., Perlin, N., & Siqueira, L. (2017, December). Ocean eddies and
576 climate predictability. *Chaos: An Interdisciplinary Journal of Nonlinear Science*,
577 *27*(12), 126902. Retrieved 2021-08-09, from [http://aip.scitation.org/doi/10](http://aip.scitation.org/doi/10.1063/1.4990034)
578 .1063/1.4990034 doi: 10.1063/1.4990034
- 579 Kirtman, B. P., & Shukla, J. (2002, May). Interactive coupled ensemble: A new cou-
580 pling strategy for CGCMs: Interactive Coupled Ensemble. *Geophysical Research*
581 *Letters*, *29*(10), 5–1–5–4. Retrieved 2021-08-09, from [http://doi.wiley.com/10](http://doi.wiley.com/10.1029/2002GL014834)
582 .1029/2002GL014834 doi: 10.1029/2002GL014834
- 583 Kraus, E. B., Turner, J. S., & Hole, W. (1967). A one-dimensional model of the sea-
584 sonal thermocline. *Tellus*, *9*.
- 585 Lee, M.-M., Nurser, A. J. G., Stevens, I., & Sallée, J.-B. (2011, August). Subduc-
586 tion over the Southern Indian Ocean in a High-Resolution Atmosphere–Ocean
587 Coupled Model. *Journal of Climate*, *24*(15), 3830–3849. Retrieved 2021-10-
588 19, from <http://journals.ametsoc.org/doi/10.1175/2011JCLI3888.1> doi:
589 10.1175/2011JCLI3888.1
- 590 Lopez, H., & Kirtman, B. P. (2014, September). WWBs, ENSO predictabil-
591 ity, the apring barrier and extreme events: WWBs and ENSO Predictability.

- 592 *Journal of Geophysical Research: Atmospheres*, 119(17), 10,114–10,138. Re-
 593 triewed 2021-10-20, from <http://doi.wiley.com/10.1002/2014JD021908> doi:
 594 10.1002/2014JD021908
- 595 Marshall, J., & Speer, K. (2012, March). Closure of the meridional overturning
 596 circulation through Southern Ocean upwelling. *Nature Geoscience*, 5(3), 171–180.
 597 Retrieved 2021-11-02, from <http://www.nature.com/articles/ngeo1391> doi:
 598 10.1038/ngeo1391
- 599 McCartney, M. S. (1977). Subantarctic mode water. *A Voyage of Discovery, M. An-
 600 gel, Ed., Elsevier*, 103–119.
- 601 O'Neill, L. W., Chelton, D. B., & Esbensen, S. K. (2003). Observations of sst-
 602 induced perturbations of the wind stress field over the southern ocean on sea-
 603 sonal timescales. *Journal of Climate*, 16(14), 2340 - 2354. Retrieved from
 604 <https://journals.ametsoc.org/view/journals/clim/16/14/2780.1.xml>
 605 doi: 10.1175/2780.1
- 606 Park, Y.-H., & Durand, I. (2019, April). *Altimetry-driven Antarctic Circumpolar
 607 Current fronts*. SEANOE. Retrieved from <https://doi.org/10.17882/59800>
 608 doi: 10.17882/59800
- 609 Perlin, N., Kamenkovich, I., Gao, Y., & Kirtman, B. P. (2020, September). A
 610 study of mesoscale air–sea interaction in the Southern Ocean with a regional
 611 coupled model. *Ocean Modelling*, 153, 101660. Retrieved 2021-01-07, from
 612 <https://linkinghub.elsevier.com/retrieve/pii/S1463500320301621> doi:
 613 10.1016/j.ocemod.2020.101660
- 614 Rintoul, S. R. (2002). Ekman Transport Dominates Local Air–Sea Fluxes in Driving
 615 Variability of Subantarctic Mode Water. *Journal of Physical Oceanography*, 32,
 616 14.
- 617 Sallée, J.-B., Matear, R. J., Rintoul, S. R., & Lenton, A. (2012, August). Lo-
 618 calized subduction of anthropogenic carbon dioxide in the Southern Hemi-
 619 sphere oceans. *Nature Geoscience*, 5(8), 579–584. Retrieved 2021-11-08, from
 620 <http://www.nature.com/articles/ngeo1523> doi: 10.1038/ngeo1523
- 621 Sallée, J.-B., Wienders, N., Speer, K., & Morrow, R. (2006, December). Forma-
 622 tion of subantarctic mode water in the southeastern Indian Ocean. *Ocean Dynam-
 623 ics*, 56(5-6), 525–542. Retrieved 2021-01-07, from [http://link.springer.com/10](http://link.springer.com/10.1007/s10236-005-0054-x)
 624 [.1007/s10236-005-0054-x](http://link.springer.com/10.1007/s10236-005-0054-x) doi: 10.1007/s10236-005-0054-x
- 625 Sarmiento, J. L., Gruber, N., Brzezinski, M. A., & Dunne, J. P. (2004, Jan-
 626 uary). High-latitude controls of thermocline nutrients and low latitude bio-
 627 logical productivity. *Nature*, 427(6969), 56–60. Retrieved 2021-11-08, from
 628 <http://www.nature.com/articles/nature02127> doi: 10.1038/nature02127
- 629 Schneider, E. K., Kirtman, B. P., & Perlin, N. (2023). The role of atmospheric
 630 noise in decadal sst variability. *Journal of Climate*, 36(7), 2147 - 2166. Re-
 631 trieved from [https://journals.ametsoc.org/view/journals/clim/36/7/](https://journals.ametsoc.org/view/journals/clim/36/7/JCLI-D-22-0399.1.xml)
 632 [JCLI-D-22-0399.1.xml](https://journals.ametsoc.org/view/journals/clim/36/7/JCLI-D-22-0399.1.xml) doi: <https://doi.org/10.1175/JCLI-D-22-0399.1>
- 633 Seo, H., Miller, A. J., & Norris, J. R. (2016, February). Eddy–Wind Inter-
 634 action in the California Current System: Dynamics and Impacts. *Jour-
 635 nal of Physical Oceanography*, 46(2), 439–459. Retrieved 2020-06-08, from
 636 <http://journals.ametsoc.org/doi/10.1175/JPO-D-15-0086.1> doi:
 637 10.1175/JPO-D-15-0086.1
- 638 Sloyan, B. M., & Rintoul, S. R. (2001). Circulation, renewal, and modification of
 639 antarctic mode and intermediate water. *Journal of Physical Oceanography*, 31(4),
 640 1005 - 1030. Retrieved from [https://journals.ametsoc.org/view/journals/](https://journals.ametsoc.org/view/journals/phoc/31/4/1520-0485-2001.031_1005_cramoa.2.0.co.2.xml)
 641 [phoc/31/4/1520-0485-2001.031_1005_cramoa.2.0.co.2.xml](https://journals.ametsoc.org/view/journals/phoc/31/4/1520-0485-2001.031_1005_cramoa.2.0.co.2.xml) doi: 10.1175/1520-
 642 [0485\(2001\)031<1005:CRAMOA>2.0.CO;2](https://journals.ametsoc.org/view/journals/phoc/31/4/1520-0485-2001.031_1005_cramoa.2.0.co.2.xml)
- 643 Small, R., deSzoeko, S., Xie, S., O'Neill, L., Seo, H., Song, Q., ... Minobe, S.
 644 (2008, August). Air–sea interaction over ocean fronts and eddies. *Dynam-
 645 ics of Atmospheres and Oceans*, 45(3-4), 274–319. Retrieved 2020-06-17, from

- 646 <https://linkinghub.elsevier.com/retrieve/pii/S0377026508000341> doi:
647 10.1016/j.dynatmoce.2008.01.001
- 648 Smith, R., Jones, P., Briegleb, B., Bryan, F., Danabasoglu, G., Dennis, J., ... Yea-
649 ger, S. (2010). The Parallel Ocean Program (POP) Reference Manual. *Community*
650 *Climate System Model*, 141.
- 651 Speer, K., Rintoul, S. R., & Sloyan, B. (2000, December). The Diabatic Dea-
652 con Cell*. *Journal of Physical Oceanography*, 30(12), 3212–3222. Re-
653 trieved 2021-11-08, from [http://journals.ametsoc.org/doi/10.1175/1520-0485\(2000\)030<3212:TDDC>2.0.CO;2](http://journals.ametsoc.org/doi/10.1175/1520-0485(2000)030<3212:TDDC>2.0.CO;2) doi: 10.1175/1520-0485(2000)030<3212:TDDC>2.0.CO;2
- 654 Tozuka, T., & Cronin, M. F. (2014, April). Role of mixed layer depth in surface
655 frontogenesis: The Agulhas Return Current front: Tozuka and Cronin: Frontogen-
656 sis: Role of mixed layer depth. *Geophysical Research Letters*, 41(7), 2447–2453.
657 Retrieved 2020-06-05, from <http://doi.wiley.com/10.1002/2014GL059624> doi:
658 10.1002/2014GL059624
- 659 Tozuka, T., Ohishi, S., & Cronin, M. F. (2018, July). A metric for surface heat flux
660 effect on horizontal sea surface temperature gradients. *Climate Dynamics*, 51(1-
661 2), 547–561. Retrieved 2020-06-05, from <http://link.springer.com/10.1007/s00382-017-3940-2> doi: 10.1007/s00382-017-3940-2
- 662 Verdy, A., Dutkiewicz, S., Follows, M. J., Marshall, J., & Czaja, A. (2007,
663 June). Carbon dioxide and oxygen fluxes in the Southern Ocean: Mecha-
664 nisms of interannual variability: Southern Ocean Carbon and Oxygen Fluxes.
665 *Global Biogeochemical Cycles*, 21(2), N/A–N/A. Retrieved 2021-10-24, from
666 <http://doi.wiley.com/10.1029/2006GB002916> doi: 10.1029/2006GB002916
- 667 Wu, R., & Kirtman, B. P. (2005, August). Roles of Indian and Pacific Ocean
668 air–sea coupling in tropical atmospheric variability. *Climate Dynamics*, 25(2-3),
669 155–170. Retrieved 2021-10-18, from <http://link.springer.com/10.1007/s00382-005-0003-x> doi: 10.1007/s00382-005-0003-x
- 670 Zhang, W., & Kirtman, B. (2019, March). Estimates of Decadal Climate Pre-
671 dictability From an Interactive Ensemble Model. *Geophysical Research Letters*,
672 46(6), 3387–3397. Retrieved 2021-10-18, from <https://onlinelibrary.wiley.com/doi/10.1029/2018GL081307> doi: 10.1029/2018GL081307

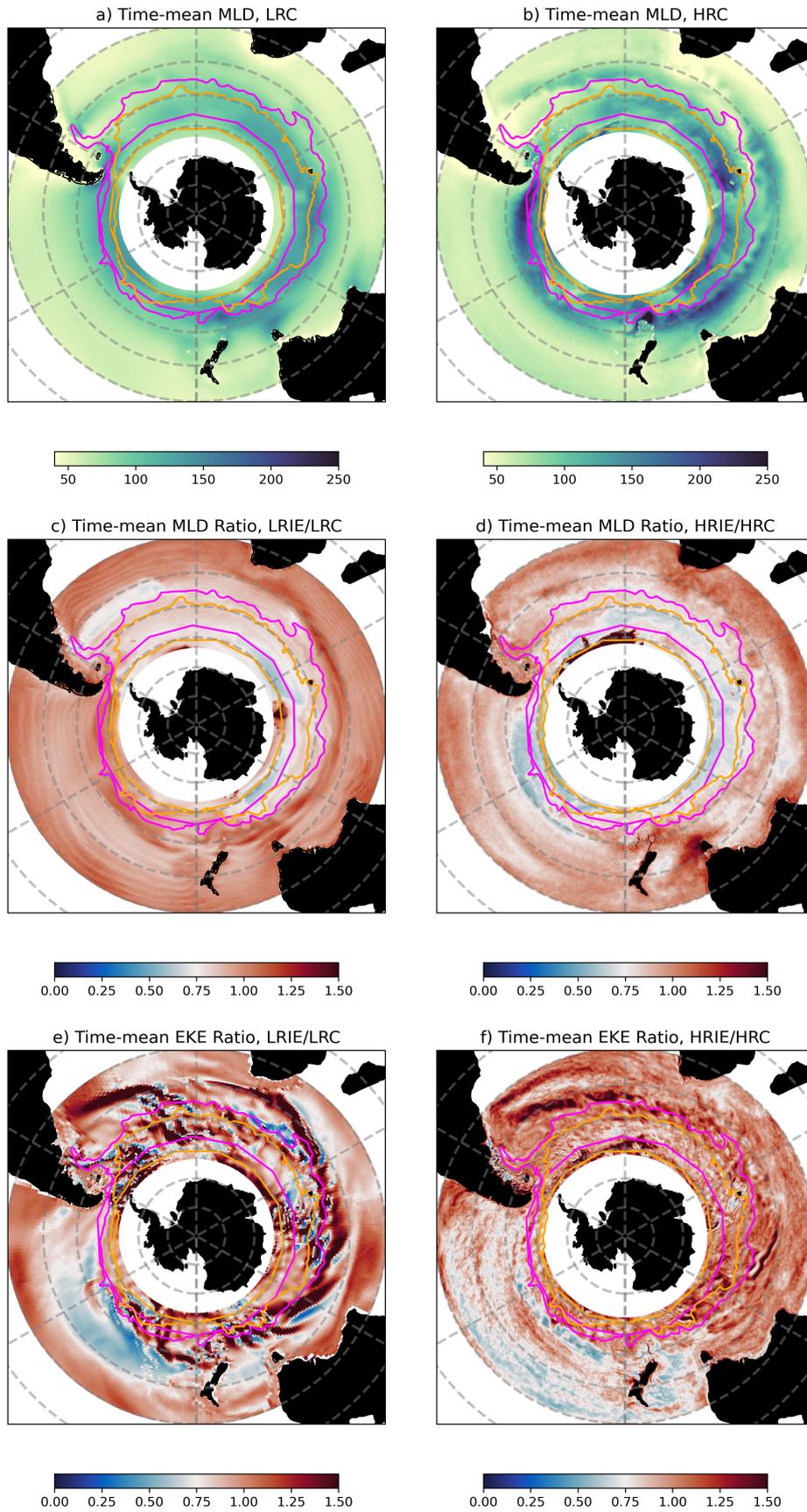


Figure 2. Time-mean MLD in a) LRC and b) HRC, from year 121 to 150. Polar Front (PF, orange line) and Subantarctic Front (SAF, magenta line).

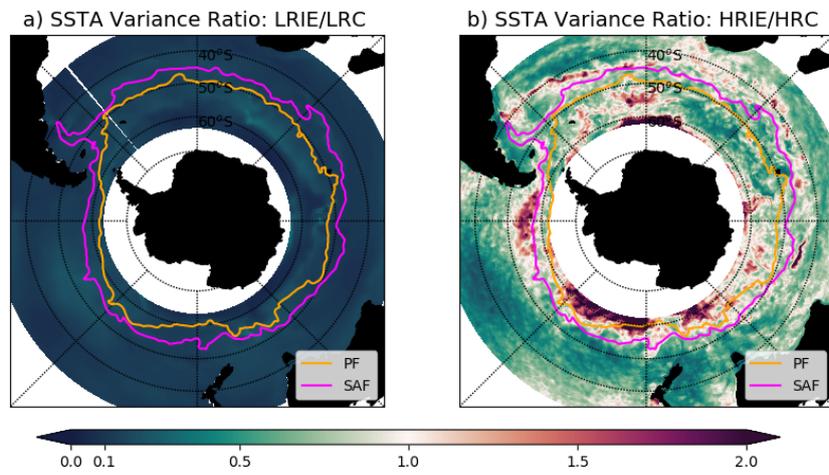


Figure 3. a) SSTA variance ratio of LRIE to LRC and b) SSTA variance ratio of HRIE to HRC. SSTA are the departures from the monthly SST climatology. Polar Front (PF, blue line) and Subantarctic Front (SAF, orange line).

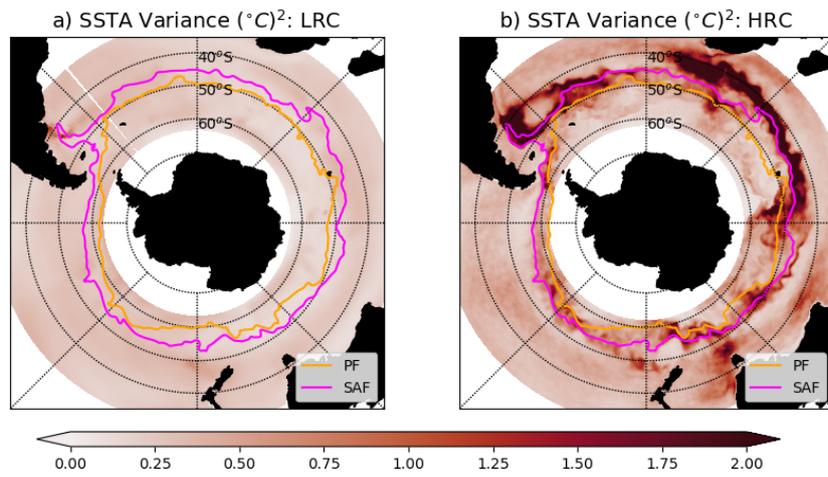


Figure 4. The variance of MLDA in a) LRC, b) HRC from year 121 to 150. SSTA anomalies (SSTAs) are defined as departures from the monthly climatology. Polar front (PF, orange line) and Subantarctic Front (SAF, magenta line).

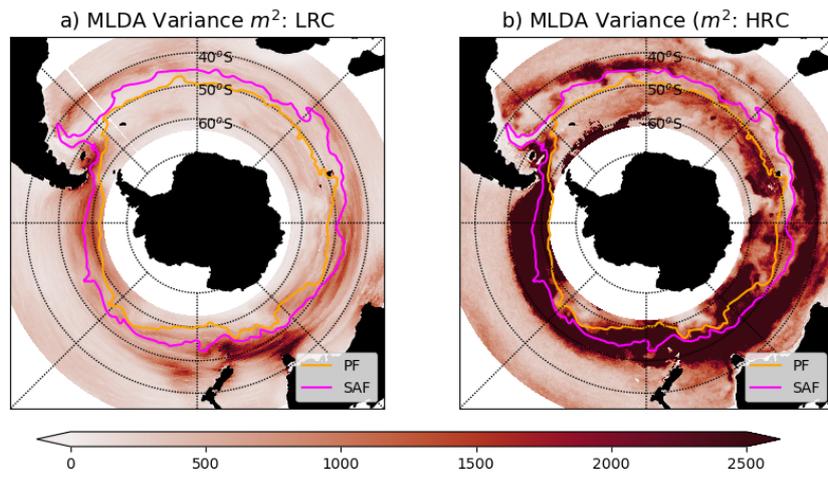


Figure 5. Variance of MLDA in a) LRC, b) HRC from year 121 to 150. MLD anomalies are defined as departures from the monthly climatology. Polar front (PF, orange line) and Subantarctic Front (SAF, magenta line).

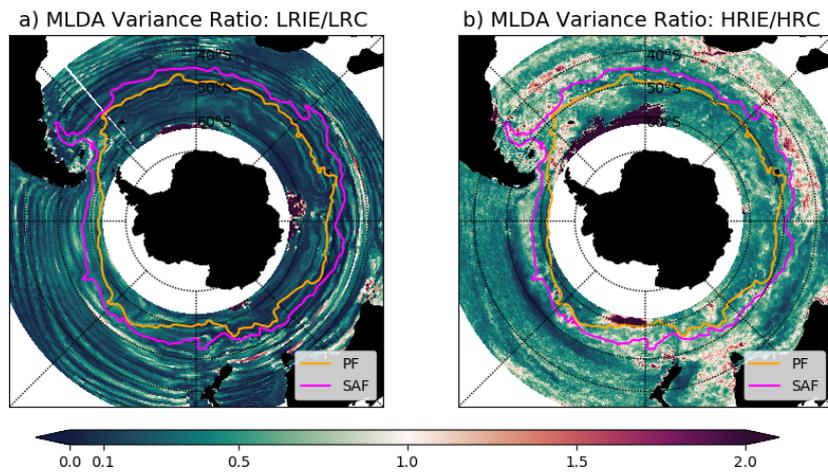


Figure 6. a) MLD variance ratio of LRIE to LRC and b) MLD variance ratio of HRIE to HRC. MLD anomalies are defined as departures from the monthly climatology. Polar Front (PF, blue line) and Subantarctic Front (SAF, orange line).

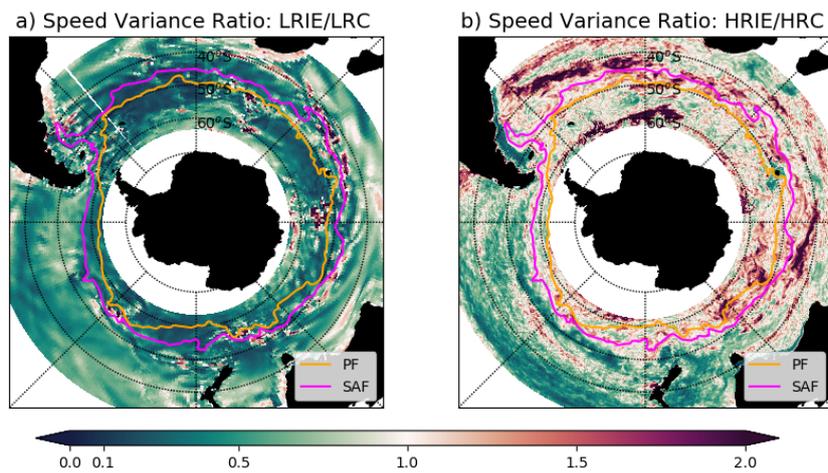


Figure 7. Surface current speed variance ratio in a) HRC and b) LRC, from year 121 to 150. Polar Front (PF, orange line) and Subantarctic Front (SAF, magenta line).

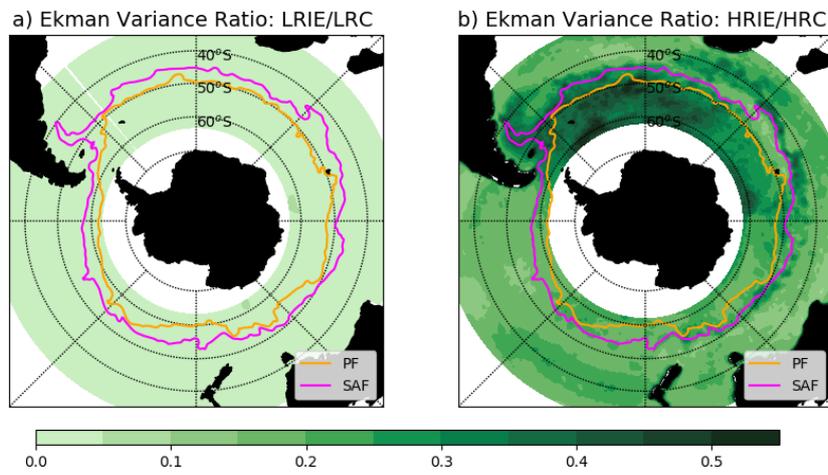


Figure 8. a) Ekman transport variance ratio of LRIE to LRC and b) Ekman transport variance ratio of HRIE to HRC. Polar Front (PF, blue line) and Subantarctic Front (SAF, orange line).