Nearshore flow dynamics over shore-oblique bathymetric features during storm wave conditions

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Abstract

Shore-oblique bathymetric features occur around the world and have been statistically correlated with enhanced shoreline retreat on sandy beaches. However, the physical mechanisms that explain a causal relationship are not well understood. In this study, radar remote sensing observations and results from a phase-resolved numerical model explore how complex morphology alters nearshore hydrodynamics. Observations at selected times during high-energy storm events as well as a suite of idealized simulations indicate that shore-oblique features induce strong spatial variations in the water surface elevation and wave breaking patterns. Re-emergent offshore flows and longshore current accelerations occur near the apex of the oblique nearshore features. The results suggest that complex bathymetric morphology exerts a powerful control on nearshore hydrodynamics and increases the potential for enhanced cross-shore and alongshore sediment transport, thus contributing to localized erosional zones.

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Key Points:

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9	•	Radar and numerical modeling indicate variability in surf zone width and wave
10		height
11	•	Offshore-directed flows occur in a region with shore-oblique sand bars and erosional
12		hot spots
13	•	Morphologic influence and directional variability are quantified with idealized model
14		simulations

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15 Abstract

Shore-oblique bathymetric features occur around the world and have been statistically 16 correlated with enhanced shoreline retreat on sandy beaches. However, the physical mech-17 anisms that explain a causal relationship are not well understood. In this study, radar 18 remote sensing observations and results from a phase-resolved numerical model explore 19 how complex morphology alters nearshore hydrodynamics. Observations at selected times 20 during high-energy storm events as well as a suite of idealized simulations indicate that 21 shore-oblique features induce strong spatial variations in the water surface elevation and 22 wave breaking patterns. Re-emergent offshore flows and longshore current accelerations 23 occur near the apex of the oblique nearshore features. The results suggest that complex 24 bathymetric morphology exerts a powerful control on nearshore hydrodynamics and in-25 creases the potential for enhanced cross-shore and alongshore sediment transport, thus 26 contributing to localized erosional zones. 27

²⁸ Plain Language Summary

Near the shoreline, underwater topography is affected by sea level, waves, currents, 29 tides, and geological characteristics. Typically, sandbars are oriented parallel to the coast-30 line, but shore-oblique sandbars have also been identified around the world. In many in-31 stances, these oblique features have been correlated with zones of erosion, however the 32 explanation for this statistical relationship is not fully understood. In this study, remote 33 sensing observations and modeling results explore how complex underwater topographies 34 alter coastal wave energy and flow patterns. Selected stormy periods were observed in 35 addition to a set of idealized simulations. Results indicate that shore-oblique features 36 cause localized changes to wave heights, wave breaking, and current speeds. These al-37 terations contribute to the presence of fast offshore-directed rip currents and accelera-38 tions in shore-parallel currents, which would enhance the potential for sand to be trans-39 ported away from these zones. The findings suggest that the changes to the flow field 40 induced by the oblique features contribute to zones of high rates of erosion. 41

42 **1** Introduction

The coastal nearshore region, defined as the transition zone between the shoreline 43 and the inner shelf, is a dynamic zone shaped by the interplay of many physical processes 44 occurring at various spatial and temporal scales. The level and distribution of incom-45 ing wave energy exerts a powerful control on nearshore forces and alters currents, sed-46 iment transport, and morphology (Wright & Short, 1984; Svendsen, 1984; Gallagher et 47 al., 1998). When propagating at oblique incident angles to the bathymetric contours, sur-48 face waves transfer momentum and drive longshore currents, moving material along the 49 coastline (Longuet-Higgins, 1970; Guza et al., 1986). These currents meander and change 50 velocity in response to bathymetric variability (Garnier et al., 2013). At moderate in-51 cident angles, gradients in wave-driven setup of the mean water level can cause local along-52 shore flows to converge and flow offshore as rip-currents (Castelle et al., 2016; Moulton 53 et al., 2017). 54

These strong offshore-directed rip currents can scour channels in nearshore sand-55 bars and alter patterns in wave dissipation, causing the local mean water level to slope 56 towards the rip channel (Haller et al., 2002). This feedback reinforces offshore flows and 57 can erode the beach profile (Komar & McDougal, 1988). For weak offshore flows, wave 58 energy dissipation can be reduced in rip-channels, leading to higher wave heights that 59 break close to the shoreline (Haller et al., 2002) and increased vulnerability to erosion 60 (Holman & Sallenger, 1993). In some cases, stable patterns in the underlying bathymetry 61 (e.g., submarine canyons) produce morphologically-driven recurring offshore flows by redi-62 recting wave energy (Long & Ozkan-Haller, 2005; Magne et al., 2007; O'Dea et al., 2021). 63 For example, dramatic variations in refraction and wave energy occur across the Scripps 64

Canyon in southern California, with increased local wave heights at the head of the canyon
 (Magne et al., 2007).

Morphologically-driven processes are often shaped by the underlying geology (also 67 referred to as antecedent or framework geology). The geologic framework, generated over 68 very long time scales, is characterized by deposits or bedrock underlying modern sed-69 iments to a depth of approximately 10 m (Browder & McNinch, 2006). Along the east 70 coast of North America, anomalous nearshore bathymetric features have been identified 71 and connected with geological history: shore-normal or shore-oblique troughs and bars 72 73 occur in New York (Schwab et al., 2000), New Jersey (Snedden et al., 1994), Virginia (Colman et al., 1990; Browder & McNinch, 2006), North Carolina (McNinch, 2004), Florida 74 (Houser et al., 2008; Barrett & Houser, 2012), northeastern New Zealand (Green et al., 75 2004), Prince Edward Island (Wernette & Houser, 2022) among many other locations 76 and are commonly related to paleo-river channels. These features alter alongshore pat-77 terns in wave breaking (Safak et al., 2017), modify the distribution of wave energy (Mulligan 78 et al., 2019a), and have been hypothesized to act as conduits for offshore sediment trans-79 port (Thieler et al., 1995) and enhanced alongshore advection (Gutierrez et al., 2005). 80 However, the physical mechanisms linking framework geology to alterations in hydro-81 dynamics and morphodynamics across the shoreface are not well understood. 82

In this study, we investigate the influence of a complex bathymetric framework on 83 nearshore flow dynamics during storm conditions at Kitty Hawk, North Carolina, USA. 84 We hypothesize that shore-oblique bathymetric features alter mean nearshore currents 85 and induce localized zones of higher wave height and wave setup on the shoreface, which 86 may affect sediment transport patterns and the morphology of this nearshore region. Both 87 real and idealized storm events are simulated with a phase-resolving numerical model 88 for a wide range of incident wave angles. The event-based simulations focus on time pe-89 riods with large waves corresponding to Hurricane Jose and Hurricane Maria in 2017, 90 that were also observed with an X-band radar (XBR) remote sensing system. The ide-91 alized simulations isolate the effects of the bathymetry and incident wave angle on the 92 nearshore hydrodynamics. The novel combination of remote sensing and numerical re-93 sults are synthesized to provide new insight on wave-driven flows over complex geologically-94 controlled bathymetric features. 95

⁹⁶ 1.1 Regional setting

Kitty Hawk (KH) is located on the Outer Banks of North Carolina (Fig. 1a). The 97 Outer Banks are a series of long barrier islands segmented by inlets that formed in a mi-98 crotidal, wave-dominated environment (Hayes, 1979), and are subject to storms (e.g., 99 Tropical Cyclones and Nor'Easters) and high rates of sea level rise (Sallenger et al., 2012; 100 Kemp et al., 2017). These barrier islands have been shaped into cuspate forelands that 101 divide large estuaries (Albemarle and Pamlico Sounds) from the Atlantic Ocean. Cur-102 rently, few inlets connect the sounds to the ocean and the barrier islands receive a very 103 limited supply of sediment from riverine sources (Culver et al., 2007; Mulligan et al., 2019b). 104 Located 5 km west of the modern transgressive shoreline at KH, a progradational beach 105 ridge complex formed 3 ka - 2 ka before present during a period of rapid Holocene sea 106 level rise and abundant sediment availability (Mallinson et al., 2008). The present-day 107 coastal morphology and dynamics are influenced by the complex underlying regional ge-108 ologic framework. 109

¹¹⁰ 2 Shore-oblique features

The segment of coast is dominated by shore-parallel depth contours except near KH, where the nearshore region includes a series of bathymetric undulations with bars and troughs as shown in Fig. 1b. These bathymetric variations are situated in mean water depths ranging from 2-12 m and are oriented offshore at an average angle of 42° from

the shoreline (Schupp et al., 2006). The morphologic features are referred to as shore-115 oblique features (SOFs) and, within the region considered in the present study, include 116 a northern (SOF-N) and southern (SOF-S) trough shown in Fig. 1b. Seismic imaging 117 studies have associated these bathymetric anomalies to underlying Pleistocene paleo-channels, 118 specifically the paleo-Roanoke River (Boss et al., 2002; McNinch, 2004; Browder & Mc-119 Ninch, 2006). Cycles of eustatic sea-level changes that occurred concurrently with glacial 120 episodes during the Pleistocene epoch dissected numerous fluvial channels into the un-121 derlying Quaternary strata (Boss et al., 2002). During the ensuing Holocene transgres-122 sion, these channels were drowned and infilled with muds, peats, sands, and gravels (Riggs 123 et al., 1992, 1995; Schwartz & Birkemeier, 2004). 124

The SOFs are relatively stable but can migrate approximately 250 m alongshore 125 in response to individual storm events (McNinch & Miselis, 2012). Schupp et al. (2006) 126 found that all movement was confined to zones shoreward of the -9 m bathymetric con-127 tour. They also slowly migrate downdrift on a decadal timescale. Between 2004 and 2017, 128 SOF-N migrated south by approximately 600 m and deepened by 0.5 m (Szczyrba et al., 129 2023a). Regionally, these SOFs have been correlated with areas of high shoreline vari-130 ability and, on longer time scales, high long-term shoreline change rates. McNinch (2004) 131 visually identified a correlation between the SOFs and high shoreline change rates and 132 Schupp et al. (2006) quantified this correlation on a regional scale as statistically signif-133 icant. However, the nearshore hydrodynamic processes that physically link the SOFs with 134 erosion have not been well defined. 135

3 Observations

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3.1 Bathymetry and wave data

In June, 2017, bathymetric profiles were collected to support the local beach nourishment monitoring plan (APTIM, 2017). The nearshore survey lines were collected by a survey vessel and spaced approximately 150 m meters apart, with sounding points sampled every 7-8 m (APTIM, 2017). The survey data were interpolated onto a 5 m x 5 m regular grid in ArcGIS 10.7.1, encompassing two shore-oblique bathymetric features (SOF-N and SOF-S). This bathymetric grid, shown in Fig. 1c, is used as input to the numerical model.

Wave observations used as boundary conditions in the numerical model were sourced 145 from the US Army Corps of Engineers Field Research Facility (FRF) at Duck, located 146 14 km north of KH. Wave spectra were observed by a 1 MHz Nortek Acoustic Wave and 147 Current (AWAC) profiler, located offshore in a nominal depth of 11 m. Previous stud-148 ies have shown that the bulk wave statistics including significant wave height (H_s) , mean 149 direction (θ_m) , and peak period (T_p) are highly correlated at the 11 m depth over a dis-150 tance of 33 km between the FRF and Jennette's Pier, located 20 km south of KH (Mulligan 151 et al., 2019a). Wave data were obtained for key times during major wave events in Septem-152 ber 2017. These correspond to a time during Hurricane Jose (denoted "J"; 19-September-153 2017 20:00:00 UTC), and two times during Hurricane Maria (denoted "M1" and "M2"; 154 27-September-2017 09:00:00 UTC and 28-September-2017 02:00:00 UTC) indicated in 155 Fig. 2 and Table 1. Directional energy-density wave spectra were used as offshore bound-156 ary conditions for the numerical model (Fig. 3). The directional spectra were calculated 157 from the first four Fourier coefficients, using the maximum entropy method (Lygre & Krogstad, 158 1986) with a directional resolution of 1° , a frequency resolution of 0.0075 Hz, and a fre-159 quency range of 0.0400-0.4975 Hz. The spectra were rotated 28° clockwise of true north 160 to match the orientation of the rotated bathymetric grid before they were input into the 161 model. 162

The times corresponding to J, M1, and M2 were selected because they have similar mean total water levels, and a range of energy levels and incident wave angles (Ta-

ble 1 and Fig. 3). Hurricane Jose at time J was a moderate wave height $(H_s = 2.8 \text{ m})$ 165 event with a positive incident angle (meaning north of shore-perpendicular) and mod-166 erate directional spreading (Fig. 3a, d, g). Hurricane Maria at time M1 was a high-energy 167 event, with an H_s of 4.0 m, long period waves, low directional spreading and an oblique 168 positive incident angle (Fig. 3b, e, h). As Hurricane Maria propagated along the coast-169 line, the dominant wind and wave direction shifted and the conditions at time M2, which 170 occurred 17 hours after M1, had moderate wave heights $(H_s = 2.6 \text{ m})$ and periods dur-171 ing the passage of the storm. Time M2 was characterized by high directional spreading 172 from the high-frequency oblique wind-waves forced by strong offshore winds (Fig. 3c, f, 173 i). 174

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3.2 X-band radar collection and processing

The Radar Inlet Observation System (RIOS) is a fully automated non-Doppler XBR 176 remote sensing system housed within a mobile trailer for rapid deployment (McNinch 177 et al., 2012; Humberston et al., 2019). RIOS emits microwave radio signals with a hor-178 izontal beam width of 1.2° and a vertical beam width of 25° and records the subsequent 179 backscatter. Bragg resonance is the primary mode of XBR backscatter, wherein radar 180 pulses interact with short capillary waves (1-2 cm wavelength, λ) on the water surface. 181 Backscatter increases when waves steepen, modulating the short-gravity wave field, and 182 become roughened as they approach the shoreline and break (Catalán et al., 2014). Thus, 183 RIOS provides a digital reconstruction of the incoming nearshore wave field as a series 184 of sequential spikes in the XBR backscatter return intensity. XBR can resolve the spatio-185 temporal evolution of incoming nearshore waves and the observations compare well to 186 both numerical models and in-situ data (Szczyrba et al., 2023b). Additional RIOS-system 187 details are described by Humberston et al. (2019) and McNinch et al. (2012). 188

RIOS was deployed in the study site at the location shown in Fig. 1b (36.066408°. 189 -75.690222°) in September, 2017. Nearshore sea surface conditions were observed for 14 190 minutes per hour between September 19-23 and September 27-30, capturing J, M1, and 191 M2 time periods. The XBR antenna completed a rotation every 1.67 s, resulting in a sam-192 pling frequency of 0.60 Hz. Raw backscatter data were transformed from polar coordi-193 nates onto a 5 m x 5 m Cartesian grid (Humberston et al., 2019) and smoothed in the 194 alongshore direction by 100 m using a multidimensional image averaging filter (imfilter). 195 The resulting grid was rotated 28° clockwise of true north to match the orientation of 196 the numerical modeling framework. Overall, the XBR footprint spanned 3 km alongshore 197 and 1 km offshore. The time-dependant surf zone width was calculated from 14 minute 198 time-averaged XBR backscatter data. Time-averaged XBR and optical data highlight 199 zones of high dissipation and are often used to determine the spatial patterns of wave 200 breaking (Lippmann et al., 1993; Brodie et al., 2018). Following the method described 201 by O'Dea et al. (2021), the offshore surf zone edge was identified as the maximum cross-202 shore gradient of the XBR backscatter intensity, calculated at each alongshore location. 203 Outliers were removed and the resulting surf zone edge lines were smoothed using a mov-204 ing mean filter with a window size of 100 m. 205

Methods from O'Dea et al. (2021) were also applied to search for and identify mor-206 phologically driven offshore-directed currents in the study region. This approach ana-207 lyzes backscatter intensity offshore of the surf zone where few waves are breaking. Off-208 shore flows escaping beyond the breaker line can be identified in the XBR data as alter-209 ations in surface roughness. Offshore directed rip currents enhance sea surface rough-210 ness, incite short-scale wave breaking (i.e. microbreaking), and elevate XBR-measured 211 backscatter intensity (Lyzenga, 1991; Plant et al., 2010; Haller et al., 2014; O'Dea et al., 212 2021). This effect has been confirmed in a study utilizing both XBR and GPS-equipped 213 floating drifters (Takewaka & Yamakawa, 2011) and with a cross-shore array of current 214 meters (Haller et al., 2014). To identify these zones of wave-current interaction, an along-215 shore transect of backscatter intensity was extracted at a location 100 m offshore of the 216

²¹⁷ surf zone edge during periods with waves with $H_s \geq 2$ m. The mean backscatter inten-²¹⁸ sity was removed from each transect to highlight zone of anomalously high return inten-²¹⁹ sities. The intensity anomaly transect data were then averaged together to identify per-²²⁰ sistent XBR backscatter anomalies in the 9-day observation period. The XBR anoma-²²¹ lies for J, M1, and M2 were also analyzed.

XBR observations were used to determine representative wave phase speeds. Pixel 222 intensity time series were input into the bathymetric-inversion algorithm cBathy (Holman 223 et al., 2013b; Holman & Bergsma, 2021). cBathy is typically used to estimate bathymetry 224 225 from optical imagery, often from Argus camera systems (e.g., Oades et al. (2023)), but can also process XBR data (Honegger et al., 2019). Resulting XBR-derived bathymet-226 ric surfaces have been applied in numerical modeling studies (O'Dea et al., 2021). The 227 algorithm estimates the wave numbers (k) of dominant wave frequencies (f) and uses 228 the linear dispersion equation to relate these k-f pairs to mean water depth. The k-f 229 pairs are extracted within mapped tiles by calculating the Fourier phase cross-spectra 230 between each pixel per tile. In the present study, tile spacings were set to be 10 m in the 231 cross-shore direction and 25 m in the alongshore direction. The algorithm produces a 232 matrix of wave numbers for each the most coherent 4 frequencies wherein the individ-233 ual matrices contain a range of frequencies that represent the most dominant frequency 234 calculated per pixel. 235

The most coherent k-f matrix was extracted to create wave speed maps of repre-236 sentative peak wave conditions for each event. cBathy's accuracy declines significantly 237 in very shallow water with h < -2 m (Honegger et al., 2019), therefore these regions were 238 excluded from further analyses. Cross-shore transects were extracted to support com-239 parisons with the modeling data with a 20 m moving mean applied to eliminate noise. 240 Using the peak frequency calculated from cBathy, the method of Streßer et al. (2022) 241 was adapted to convert XBR-measured wave phase speeds to representative wave height 242 estimates (see Eq. 12 in Streßer et al. (2022)). Their wave-by-wave approach applies a 243 physically derived scaling that relates changes in breaking phase speed to wave height 244 245 via the amplitude dispersion relation:

$$H_p = \frac{C_p^2}{g(\frac{1}{\gamma} + \alpha_{ad})} \tag{1}$$

This equation was applied using the celerity of the peak frequency (C_p) , the default breaking index (γ) value of 0.78 (Streßer et al., 2022; Larson & Kraus, 1994), and the recommended calibration coefficient (α_{ad}) value of 0.5 (Streßer et al., 2022) to estimate a representative wave height based on peak frequency values (H_p) . This method is valid within the surf zone, since it was developed based on the modified nonlinear shallow water phase speed for waves in very shallow water (i.e., $h:\lambda < 1/10$) from Hedges (1976).

²⁵³ 4 Numerical model

The phase-resolving 3D numerical model Simulating WAves till SHore (SWASH) 254 applies the nonlinear shallow water horizontal momentum equations and the nonhydro-255 static vertical momentum equation using a finite difference scheme (Zijlema et al., 2011; 256 Zijlema, 2020) to simulate the water surface and velocity fluctuations. SWASH was ap-257 plied to simulate several times (J, M1, M2) in the study period that correspond to radar 258 data collection times. In addition to the selected events, an idealized suite of 21 simu-259 lations were also performed to isolate key variables and specifically explore how the in-260 cident wave conditions affect the nearshore hydrodynamics. 261

4.1 Model setup

SWASH has been used to simulate wave-driven currents over variable bathymetry, 263 such as over submerged reefs (da Silva et al., 2023). Previous studies have also applied 264 SWASH to model regions near the present study area (Gomes et al., 2016; Mulligan et 265 al., 2019a; Szczyrba et al., 2023b) and similar numerical input parameters were applied 266 in this study. A 5 m x 5 m regular structured grid, extending 1000 m offshore and 3000 267 m alongshore, was constructed. A 300 m extension was added to the offshore edge of the 268 bathymetric grid (from x = 1000-1300 m) to create a uniform offshore region and pre-269 270 vent numerical anomalies from being introduced due to depth variations at the east boundary (without the extension, the bathymetric variations at the model boundary causes 271 numerical instabilities that are not realistic). The extension was created by linearly in-272 terpolating the offshore bathymetric edge to a constant depth of h = -12 m over a cross-273 shore distance of 100 m and then extending this constant depth over the remaining 200 274 m, following previous studies (O'Dea et al., 2021). The bathymetry was also extended 275 by 500 m in each alongshore direction by interpolating the sides to a uniform contour 276 configuration in order to enable the application of periodic alongshore boundary condi-277 tions in the numerical model. The bathymetry data was rotated 28° clockwise of true 278 north which oriented the mean shoreline parallel to the y-axis of the grid (Fig. 1c). To 279 resolve depth-dependent dynamics, three bathymetry-following vertical layers with equal 280 thickness were included. This vertical resolution is sufficient to resolve wave frequency-281 dispersion (Zijlema & Stelling, 2005; Smit et al., 2013). To confirm this, a sensitivity test 282 using 9 vertical layers was conducted to explore the flow structure in and around the SOFs 283 in further detail, and found no substantial differences from the 3-layer runs. 284

The Sommerfeld radiation condition was applied at the onshore boundary to ap-285 propriately limit wave reflection and the offshore boundary was weakly reflective. The 286 alongshore boundary was periodic, meaning that energy exiting one alongshore bound-287 ary re-entered the domain at the opposite alongshore boundary. The first 15 minutes of 288 modeling time were disregarded as model spin-up and the remaining 45 minutes were 289 analyzed. The wavemaker generated random wave time series that statistically matched 290 the input wave parameters with a cycle interval of 2700 seconds (45 minutes). Accord-291 ing to Zijlema et al. (2011), the cycle time should span 100-300 peak wave periods to pro-292 vide accurate statistical results and the selected cycle time spans over 200 wave periods 293 of the longest peak period simulated (M1). The model performed calculations with an 294 initial time step of 0.05 s that is automatically adjusted throughout the simulation ac-295 cording to the Courant number, which can reach a maximum of 0.5. Total water levels 296 were adjusted for the correct tidal stage and storm surge present during the simulated 297 events. Event times during hurricane wave events J (Jose), and M1, M2 (Maria) were 298 forced at the offshore boundary using the directional energy-density spectra for each cor-299 responding time (Fig. 3g-i). 300

4.2 Model output processing

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To enable direct comparison with the XBR cBathy products, spectral processing 302 extracted peak wave frequency parameters per grid cell from the J, M1, and M2 simu-303 lations. The fast Fourier transform was calculated at each grid cell with 256 samples and 304 a Hanning window with 50% overlap. The peak wave frequency per grid cell was also 305 extracted and used to calculate k and λ with the dispersion relation (Fenton & McKee, 306 1990; Dean & Dalrymple, 1991). The representative C_p was then calculated and input 307 into Eq. 1. This approach emulates the XBR processing steps. H_s was calculated across 308 the computational domain as four times the standard deviation of the water surface el-309 evation output (Raubenheimer et al., 1996). The spatial gradient of the water surface 310 elevation data highlighted wave crests as locations with steepest slopes. Vectors perpen-311 dicular to these crests at each grid cell were then time-averaged to provide local estimates 312 of mean wave angles across the study region (Szczyrba et al., 2023b). A dimensionless 313

wave breaking parameter (Q_b) was calculated from the binary wave breaking locations output from the model (Gomes et al., 2016). This parameter represents the relative intensity of breaking independent of the model time step and spatial resolution. These breaking zones were compared to the surf zone widths observed by XBR.

318 5 Results

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5.1 Event flow dynamics

Regions of persistently high XBR backscatter outside of the surf zone were iden-320 tified for J, M1, and M2 as well as the mean during all periods with $H_s \geq 2$ m observed 321 by XBR (Fig. 4). The strongest XBR anomalies occur at the southern edge of SOF-N 322 and the southern edge of SOF-S, located at y ~ 1600 m and y ~ 250 m (Fig. 4b-d). Each 323 event, as well as the study-period mean, indicate anomalously high backscatter signals 324 at these locations. Offshore flows enhance XBR backscatter intensity where wave-current 325 interactions roughen the sea surface. Even during lower energy periods, high backscat-326 ter indicative of offshore flow extends outside of the surf zone. For example, as shown 327 in Fig. 4d, at a time when $H_s = 2.3$ m, high backscatter return intensity (RI) is evident 328 at y ~ 1600 m and y ~ 250 m. The signals along the southern edges of both SOFs are 329 strongest during M2, an event with negatively incident (southerly) waves. Spikes in XBR 330 backscatter also occur along both the northern and southern edges of SOF-S during M2, 331 and in the channel trough backscatter measures abruptly weaken (Fig. 4b, c). In the mean 332 observations, this U-shaped pattern in the XBR anomaly can be observed in both SOFs 333 and is much wider in SOF-S than SOF-N. This U-shape indicates increased backscat-334 ter along both edges of each SOF, with the southern edge inducing stronger backscat-335 ter. M1, an event with large waves approaching the shoreline at a highly oblique, pos-336 itive (northerly) angle, only exhibits backscatter anomalies along the southern edges of 337 the SOFs, and no anomaly is detected near the northern edges. The strongest XBR anoma-338 lies align approximately with regions of higher local shoreline change rates, i.e., erosional 339 hot spots (List et al., 2006), with a slight offset (Fig. 4a, b). Each hot spot is offset north 340 of the southern edge of SOF-N and SOF-S, where the highest XBR anomalies occur. 341

The observed and simulated extent of the surf zone are shown in Fig. 5. Event M1 342 was the highest energy event with the widest surf zone (Fig. 5a). Mean surf zone widths 343 are 223 m, 281 m, and 269 m for J, M1, and M2, respectively, however the extents vary 344 considerably in the alongshore direction. The surf zone of M1 is widest north of SOF-345 N, beginning offshore at $x \sim 340$ m. The surf zone of M2 is widest south of SOF-N, oc-346 curring at $x \sim 320$ m. During all events, surf zones narrow along the shoreward apex of 347 SOF-N and widen on either side, coincident with the feature's edges. SOF-N modulates 348 surf zone widths more than SOF-S. The offshore extent of the surf zone identified by XBR 349 (Fig. 5a) was compared to the breaking parameter (Q_b) calculated from the numerical 350 model (Fig. 5b). The XBR-sensed surf zone edge aligns with the offshore edge of the most 351 intense breaking locations in the SWASH simulations and displays similar lateral vari-352 ability in break-point locations. During M1, Q_b is most intense at x ~ 200 m, y ~ 2,000, 353 which is located along the northern edge of SOF-N (Fig. 5b). 354

Spatial patterns in C_p shown in Fig. 6 are similar in XBR observations (Fig. 6a) 355 and SWASH results (Fig. 6b), wherein C_p remains high within the shore-oblique troughs 356 as the waves approach the shoreline, coinciding with the bathymetric contours. Peak fre-357 quency waves enter the domain at 9-10 m/s and slow to 6-7 m/s along the 4 m bathy-358 metric line. Across a variety of cross-shore profiles, XBR and SWASH estimates agree 359 well (Fig. 6c-e). Mean C_p errors are 0.26, 0.26, and -0.04 m/s for J, M1, and M2, respec-360 tively, however errors vary spatially. In the cross-shore, error is highest shoreward of x 361 \sim 335 m where the model slightly underestimates C_p by 0.47 m/s when compared to XBR 362 observations. In the alongshore, the model underestimates C_p north of y ~ 1,750 m and 363

overestimates south of y \sim 1,500 m. Across all periods simulated, mean error is highest between y \sim 925 - 1,225 m with a maximum of -0.53 m/s.

Using Eq. 1, spatial distributions of C_p were converted into representative wave heights (H_p) to further compare XBR and model results as shown in Fig. 7. Patterns in H_p compare well between XBR and the model. Wave heights vary laterally across the study domain, with higher wave heights reaching closer to the shoreline in the lee of the troughs of the SOFs. Wave dissipation begins along the 4 m contour lines. Root mean square errors (RMSE) are 0.76 m, 0.72 m, and 0.65 m across the entire domain and 1.35 m, 0.61 m, and 0.57 m within the surf zone for J, M1, and M2, respectively.

Fig. 8 displays the spatial patterns in relative wave height H_s/h (Fig. 8a) and wave 373 frequency (Fig. 8b) alongside the alongshore distributions of the normalized maximum 374 H_s (Fig. 8b) and the water surface elevation gradient (Fig. 8c) The distribution of mod-375 eled H_s varies across space and by event (Fig. 8a, b). Relative wave heights are elevated 376 north and south of SOF-N around y ~ 2000 and 1500 m, and in these zones high rel-377 ative wave heights extend to the 6 m contour (Fig. 8a). H_s values increase near the SOFs 378 and decrease in the region between SOF-N and SOF-S (Fig. 8b). During all events, H_s 379 increases along the edges of the SOFs, following the bathymetric contours, and decreases 380 along the central axes (Fig. 8b). The alongshore differences in H_s at a particular cross-381 shore location varies between 0.2-0.4 m. The alongshore variability in wave heights also 382 induces strong variations in the water surface elevation gradient (Fig. 8c). At SOF-N, 383 the minimum alongshore gradient occurs at y = 1665 m and then sharply increases on 384 either side of that location across all events. The mean wave frequency is also elevated 385 along the southern edges of SOF-N and SOF-S, coincident with the abrupt alongshore 386 variations in wave setup (Fig. 8c), indicative of wave-current interactions. The wave fre-387 quency changes displayed in Fig. 8d appear similar to the XBR backscatter patterns dis-388 played in Fig. 4d. 389

Simulated wave angle and velocity results near SOF-N during J, M1, and M2 are 390 shown in 9. The SOFs modulate patterns of wave refraction during each event (Fig. 9a-391 c). M1 waves maintain positive approach angles across the domain even near the shore-392 line but are moderated closer towards shore-normal in the lee of SOF-N. J and M2 waves 393 refract into positive and negative directions, creating zones of convergence and divergence 394 around SOF-N. In both events, divergence occurs along the northern edge of SOF-N while 395 convergence occurs along the southern edge. A similar, but less severe effect, is observed around SOF-S. The features cause asymmetric wave refraction, whereby wave entering 397 at moderate angles refract away from the trough and toward the edges of the SOFs (Fig. 398 9a-c), increasing wave heights along the edges and sheltering the zone immediately in 399 the lee of the trough (Fig. 8b). 400

Across all three selected times, cross-shore velocities are directed offshore down the 401 oblique central axis of SOF-N (Fig. 9d-f), coinciding with the locations of wave diver-402 gence (Fig. 9a-c). These flows decelerate and broaden as they move offshore. Cross-shore 403 velocities route back onshore at the southern edge outside of the surf zone and this shoreward-404 directed patch occurs 50 m farther offshore during M1 than during J and M2. North of 405 SOF-N, cross-shore velocities are also directed onshore during J and M2, but this effect 406 is weak during M1. The offshore and onshore-directed velocities thus form asymmetric 407 circulation cells, with the stronger cell occurring along SOF-N's southern edge. The lo-408 cation of the maximum and minimum cross-shore velocities remain fairly stable across 409 all three periods simulated. Alongshore velocities accelerate and meander significantly 410 offshore near SOF-N in J and M2 (Fig. 9g-i). The acceleration in both events occurs near 411 412 $x \sim 250$ m, $y \sim 1750$ m. During M1, the alongshore velocities are very strong and meander slightly in response to the SOF contours. The longshore current reverses direc-413 tion during M2 at x \sim 300 m, y \sim 1950 m, and the nexus of the reversal coincides with 414 a location of high cross-shore velocity. 415

During J and M2, the mean flow is characterized by two circulation cells between 416 (1) $x \sim 0.500$ m, $y \sim 2000-2500$ m and (2) $x \sim 0.500$ m, $y \sim 1500-2000$ m (Fig. 9j-l). 417 The southern circulation cell is stronger than the northern cell. While water moves off-418 shore down the axis of SOF-N in J and M2, the mean current merely meanders in re-419 sponse SOF in M1 since the strong inertia of the longshore current overpowers poten-420 tial circulation cells. The offshore-directed flows in J and M2 maintain high velocities 421 offshore of the surf zone, with fast flows identifiable up to 2 surf zone widths away from 422 shoreline. Offshore of SOF-N at x = 600 m, the maximum velocities occur in similar lo-423 cations (U = 0.57 m/s, 0.90 m/s, 0.54 m/s at y = 1820 m, 1515 m, 1500 m for J, M1, 424 M2, respectively) and the minimum velocities also occur in similar locations (U = 0.45425 m/a, 0.79 m/s, 0.40 m/s at y = 2275 m, 2065 m, 2255 m for J, M1, M2, respectively). 426 The alongshore velocity component plays a larger role than the cross-shore velocity com-427 ponent in controlling the overall mean current. 428

429

5.2 Idealized effects of incident wave conditions

A second suite of idealized simulations isolated the effects of various wave conditions including incident θ_p and H_s on nearshore hydrodynamics. These simulations were forced by a JONSWAP spectra with a T_p of 12 s and zero directional spreading. The H_s values varied between 1.0, 3.0, and 5.0 m while θ_p varied between -15 to 15° in intervals of 5°. It should be noted that the ideal simulation with an H_s of 5.0 m and θ_p of 15° failed because large waves at highly oblique angles are challenging to simulate (Baker et al., 2021). Thus, results for that simulations pertain to an H_s of 4.9 m at θ_p of 15°.

Model results of Q_b across a variety of incident wave approach angles are shown 437 in Fig. 10. Incident waves that approach from $+15^{\circ}$ and transition to -15° emulate a com-438 mon pathway of tropical cyclones in this region and Q_b varies depending on the incom-439 ing θ_p (Fig. 10). Q_b is dispersed relatively evenly alongshore between the 6 m depth con-440 tour and the shoreline when incident waves approach at $+15^{\circ}$. However, Q_b begins fur-441 ther offshore along the edges of SOF-N and closer to the shoreline at the shoreward apex 442 of the SOF. Between $+10^{\circ}$ - 0° , Q_b begins to concentrate along the northern edge of SOF-443 N. The breaking region is most localized when incident waves approach at -5° , focused 444 at the shoreward apex of SOF-N (near x = 150 m, y = 1750 m) as well as SOF-S (near 445 x = 250 m, y = 500 m). North of SOF-N most breaking occurs at the shoreline. At -446 10° , waves break again across the northern sandbar. Finally, at -15° , breaking is more 447 widely distributed across the axes of both SOF-N and SOF-S. 448

Simulated U across a variety of incident wave approach angles are shown in Fig. 449 11. When waves approach the shoreline at highly oblique angles, a strong longshore cur-450 rent develops with minor offshore movement (Fig. 11a-b, f-g). The longshore currents 451 accelerate at the apex of SOF-N, near y = 1900 m, and SOF-S, near y = 500 m, and the 452 flow also meanders slightly offshore in these zones. With moderate incident θ_p , several 453 offshore circulation cells develop. During moderately positive incident angles, offshore 454 flows route down the northern edges of the SOFs (Fig. 11b, c) while during moderately 455 negative incident angles, they flow down the southern edges of the SOFs (Fig. 11e, f). 456 During periods of negative wave incidents (Fig. 11e-g), a localized zone of flow acceler-457 ation away from the shoreline develops at the location where the shoreline angle changes 458 orientation, between y = 1400 - 1600 m. 459

Across all simulations, weak offshore flows are consistently present along either the northern or southern edges of SOF-N and SOF-S and these offshore currents exit the computational domain, 1 km away from the shoreline (Fig. 11). However, at moderate incident angles, the offshore flows strengthen, coalesce, and develop complex pathways. Near the shoreline, a series of alongshore feeder currents converge and route offshore across a wide zone around SOF-N (Fig. 11c-e). Beginning at an incident angle of $+5^{\circ}$, four distinct zones of offshore-routed flows develop at y = 250 m, y = 1100 m, y = 1600 m, and y = 2500 m (Fig. 11c). These flows strengthen as incident angles shift towards -5° and begin to converge into two dominant offshore routing zones (near y = 750 m and 1800 m) which are continually present as wave angles continue to become more oblique at - 10° , identifiable at y = 450 m and y = 1750 m (Fig. 11f).

471 6 Discussion

472 6.1 Model error assessment

492

Several past studies have evaluated a similar model configuration at the FRF and 473 found that model results of H_s and energy-density spectra compare well with in-situ sen-474 sors located across the nearshore region (Gomes et al., 2016) as well as spatially-continuous 475 XBR estimates of wave angles and wave breaking patterns (Szczyrba et al., 2023b). In the absence of in-situ sensors situated within the study site, the accuracy of the J, M1, 477 and M2 simulations is assessed through a comparison to the XBR observations of surf 478 zone width, celerity, and wave height. The alongshore variability in surf zone width mea-479 sured by XBR is well represented by patterns of Q_b simulated in the model (Fig. 5b). 480 Mean celerity errors are low, ranging from -0.04 to 0.26 m/s. Errors are highest shore-481 ward of x = 400 m, where high rates of bathymetric change occur (Fig. 6). 482

Estimates of H_p are also in good agreement and both depict similar alongshore vari-483 ability in response to the complex bathymetric configuration (Fig. 7). The best avail-484 able bathymetric data were surveyed in June of 2017. Subsequent changes in the nearshore 485 morphology, particularly in the surf and swash zones, would lead to differences in wave 486 parameters estimated from XBR and SWASH and contribute to the identified errors. Past 487 studies have indicated that the SOFs in the region remain in fixed locations even after 488 energetic wave events (McNinch, 2004; Browder & McNinch, 2006; Schupp et al., 2006). 489 therefore it is expected that high errors due to alterations in the bathymetry would be 490 confined to zones closest to the shoreline. 491

6.2 Identification of rip currents

Anomalously high XBR backscatter signals seaward of the surf zone can occur where 493 incoming waves interact with surface current convergences, such as rip currents and other 494 offshore flows (Takewaka & Yamakawa, 2011; Haller et al., 2014; O'Dea et al., 2021). Backscat-495 ter anomalies indicative of strong offshore currents are visible in nearly every collection 496 hour with incident $H_s \ge 2$ m (Fig. 4d). These offshore flows emerge in two primary lo-497 cations, between y = 200 - 400 m and between y = 1500 - 2000 m. Flow structures re-498 main visible for several hours across a variety of incident angles, water levels, wave pe-499 riods, and directional spreading conditions. The visible plumes also extend 1-3 surf-zone 500 widths beyond the XBR-identified surf zone edge (Fig. 4), supporting the findings of Kumar 501 et al. (2021) and O'Dea et al. (2021). 502

Because the observed currents re-emerge in similar positions, they are likely mor-503 phologically controlled (Short, 2007). These flow structures are more obscured when wind 504 levels in the region reach above 10 m/s (Fig. 4c) because XBR images become saturated 505 during periods of high water surface roughness (Haller et al., 2014). Another limitation 506 of XBR is the requisite for depth-limited wave breaking. Small waves below ~ 1 m are 507 not large enough to be resolved by XBR, although this threshold is site-specific (McNinch, 508 2007). However, rip currents are also more likely to be generated as wave energy increases 509 (MacMahan et al., 2005), thus XBR remains a useful tool to identify them. 510

The observed reoccurring offshore flows coincide with the locations of SOF-N and SOF-S and are also represented in simulations of J, M1, and M2 (Fig. 8d and Fig. 9jl). Because the maximum current velocities outside of the surf zone at x = 600 m reach values between 0.54 - 0.90 m/s during these events, these flows would pose a risk to human safety and thus can be considered morphologically-controlled rip currents (Leatherman & Fletemeyer, 2011). Similar re-emergent flows are also represented in the idealized simulations with H_s values of 1 m, 3 m, and 5 m (Fig. 11). Even during low energy conditions ($H_s = 1$ m), narrow bands of weak offshore flow appear in a similar location shoreward of SOF-S (flows originating near y = 500 m) and SOF-N (flows originating near y = 1500 m).

During higher energy conditions $(H_s = 3 \text{ m and } 5 \text{ m})$, the origin of the offshore flow 521 within the surf zone is more variable, however the flows preferentially route along either 522 523 the northern or southern edges of SOF-N and SOF-S. When waves approach at a positive angle, depth-averaged currents are directed onshore along the southern edge (up-524 stream side of the SOF) and are offshore-directed along the northern edge (downstream 525 side). The opposite circulation pattern emerges when waves approach at a negative an-526 gle, with offshore flows directed along the southern edge. This patterns is similar to ob-527 servations of a dredged rip channel by Moulton et al. (2017), wherein a commonly ob-528 served circulation pattern included meandering longshore currents with offshore-directed 529 flows occurring on the downstream rip channel wall. 530

Other studies have linked complex bathymetry to rip current generation. Along the 531 northern shore of Prince Edward Island, Canada, Wernette and Houser (2022) proposed 532 a relationship between paleo-river channels identified by ground-penetrating radar and 533 rip currents observed by aerial imagery. At the present study site, results indicate that 534 the SOFs act as a pseudo-transverse bar and rip morphology, which is associated with 535 a common class of rip currents (Holman et al., 2006; Short, 2007; Turner et al., 2007). 536 However, the SOFs remain fixed in place with little change, as opposed to ephemeral sand-537 bar features that migrate shoreward after a storm reset (Houser et al., 2020). This pro-538 cess is akin to the shoreface-connected ridges on Fire Island, New York that have also 539 been shown to modify wave refraction and produce alongshore variable wave breaking 540 patterns (Safak et al., 2017). 541

542

6.3 Rip current generation

Rip currents are driven by gradients in wave momentum (Longuet-Higgins & Stew-543 art, 1964) resulting from alongshore variations in wave breaking (Lippmann & Holman, 544 1989). These variations can be caused by hydrodynamic forcing, such as intersecting wave 545 trains, standing edge waves, and wave spreading (Bowen, 1969; Suanda & Feddersen, 2015; 546 Kumar & Feddersen, 2017; Moulton et al., 2023), or artificial structures, such as piers 547 and groynes (Pattiaratchi et al., 2009). However, they are more often induced by along-548 shore variations in surf zone morphology (Bowen, 1969; Dalrymple, 1978). While the SOFs 549 at this site are not morphologically pronounced within the surf zone, they are persistent 550 nearshore structures that visibly impact the configuration of the bathymetry in areas of 551 depth-limited breaking during higher energy wave events (e.g., the 4 m isobath, Fig. 1). 552 Thus, the SOFs induce substantial alongshore variability in wave setup (Fig. 8c). 553

Nearshore wave heights respond to changes in nearshore bathymetry and the present 554 results indicate that larger waves reach closer to the shoreline along the axis of the SOFs 555 (Fig. 7), supporting the conclusions of Mulligan et al. (2019a) and observations by Sonu 556 (1972) at other sites. Waves refract across the obliquely variable bathymetry (Fig. 9a-557 c), creating zones of wave energy convergence and alongshore variations in H_s (Fig. 8a, 558 b). As a result, alongshore variations in wave breaking and surf zone widths also occur 559 (Fig. 5), which can drive rip currents (Bowen, 1969; Dalrymple, 1978). Higher setup is 560 generated near the shoreward side edges of the SOFs due to variability in wave break-561 ing (Fig. 8c), leading to a pressure gradient that forces water to flow away from regions 562 of intense breaking towards zones of lower setup. When these alongshore flows converge, 563 they route offshore (Fig. 9), interact with the incoming wave field, and cause localized 564 zones of higher frequency waves (Fig. 8d) and rougher water (Fig. 4b). These results 565

support the findings of experimental rip channel studies (Haller et al., 2002; Moulton et al., 2017).

568

6.4 Influence of wave climate

The mean incident wave conditions near KH throughout 2017 included a signifi-569 cant wave height of 0.87 m, mean incident angle of 17°, and 32° of directional spread-570 ing. However, Nor'easters and tropical cyclones generate a variety of high-energy inci-571 dent wave conditions in this region. When waves approach the SOFs at angles close to 572 the orientation of the axis (negative incident angles, in our coordinate system), wave break-573 ing occurs in concentrated local patches (Fig. 10). When approaching from the oppo-574 site direction (positive incident wave angles), wave breaking is more widely distributed 575 along the 4 m isobath. These findings support Safak et al. (2017), who concluded that 576 when waves approach shoreface-connected ridges at angles similar to the angle of the ridge 577 crests, the features focused the most wave energy. Strong offshore flows (i.e., rip currents) 578 are more likely to occur when wave angles approach the site between -10° to 10° (Fig. 579 11). Rip currents are more prevalent during conditions of moderate incident wave an-580 gles (Engle, 2002; MacMahan et al., 2005; Dusek & Seim, 2013a) because highly oblique 581 angles induce strong longshore currents that inhibit offshore flow (Kumar et al., 2011). 582 During positive, oblique incident conditions (15°) , the longshore current accelerates where 583 the contours of the SOFs pinch towards the shoreline (Fig. 11a) whereas during nega-584 tive, oblique conditions (-15°) , this acceleration occurs south of the SOFs. 585

Numerous weak offshore flows occur during periods of low energy, however these 586 offshore flows re-emerge in similar positions across a variety of incident wave angles. As 587 H_s increases, these offshore flows coalesce into larger, stronger offshore flows resembling 588 rip currents (Fig. 11 that preferentially route down either the northern or southern edges 589 of the SOFs when incident wave angles enter the domain between -10° to 10° . Fig. 12 590 explores the idealized effect of incident H_s and θ on maximum U (Fig. 12a, c-e) and vari-591 ability of U (Fig. 12b, f-h) at the mean XBR-measured surf zone location as well as two 592 zones located 250 m and 500 m offshore. Within the surf zone, maximum U values are 593 mostly found at the shoreward axis of SOF-N, located at approximately y = 1650 -594 2000m (Fig. 12a). Additionally, U generally increases with increasingly oblique waves, 595 owing to the generation of a strong longshore current (Fig. 12c). The variability of U596 is higher for negatively incident waves, but is highest for waves of moderate incident an-597 gles (Fig. 12f). Both mean U and the variability of U decrease offshore and increase with 598 larger H_s . At all cross-shore locations, the variability of U decreases with increasingly 599 oblique incident waves (Fig. 12f-h). The locations of maximum U variability are con-600 centrated along the southern edges of both SOF-N and SOF-S (Fig. 12b). 601

The idealized simulations do not represent realistic field conditions, but are intended 602 to isolate the response to different incident wave conditions. This set of simulations in-603 cluded JONSWAP spectra with zero directional spreading. Several studies have suggested 604 that rip currents are more pervasive during narrow-banded incident conditions (Dusek 605 et al., 2011; Dusek & Seim, 2013b), therefore the idealized results can not be used to pre-606 dict rip current activity. However, XBR images indicate that rip currents at this site are 607 indeed generated throughout a range of incident angles and directional spreading val-608 ues (Fig. 4). The purpose of the idealized results is to explore the sensitivity and response 609 of the SOFs to isolated variables (i.e. θ , H_s). It is also well known that rip current in-610 tensity varies in response to the total water level (e.g., tidal stage) and that strong rip 611 currents can occur at low tide even during periods of low energy conditions due to en-612 hanced wave breaking over surf zone sandbars (Brander & Short, 2001; MacMahan et 613 al., 2005; Voulgaris et al., 2011). Overall, bathymetrically controlled rip currents tend 614 to be strongest and most prevalent during periods of low water levels, high H_s , moder-615 ate incident *theta*, and low spreading (Haller et al., 2002; MacMahan et al., 2010; Moul-616 ton et al., 2017). 617

6.5 Implications for hot spots of erosion

Prior studies have demonstrated that, at the regional scale (i.e., ~ 40 km along-619 shore length). SOFs in the Outer Banks correlate significantly with erosional hot spots 620 (McNinch, 2004; Schupp et al., 2006). The physical mechanisms contributing to the causal 621 relationship have previously not been determined. The results of this research suggest 622 that strong cross-shore fluxes (i.e., rip currents) are concentrated near the SOFs during 623 periods with moderate incident wave angles. Between 2012 and 2022, 72 beach rescues 624 specifically related to rip currents occurred within the study site (Kitty Hawk Ocean Res-625 cue, personal communication, September 28, 2023). Rip currents and surf zone eddies 626 are the primary mechanisms for cross-shore sediment transport (MacMahan et al., 2006; 627 Dalrymple et al., 2011; Castelle et al., 2016) and Splinter and Palmsten (2012) found that 628 areas of higher dune erosion occur where rip currents are present directly offshore. Dur-629 ing periods of both high wave energy and moderate wave angles, strong offshore-directed 630 flows may pronounce offshore transport, leading to enhanced erosion and undulations 631 in the shape of the shoreline, similar to the results found in Fire Island, New York by 632 Safak et al. (2017). During conditions with more oblique incident wave angles, longshore 633 currents accelerate near the SOFs due to alongshore variability in bathymetry, which would 634 also contribute to higher rates of sediment transport. Therefore, across a range of inci-635 dent wave conditions and energy levels, a combination of longshore current accelerations 636 and far-reaching offshore flows could contribute to the severity of beach erosion near the 637 SOFs at this site. 638

Bathymetric surveys indicate that between 2004 and 2017 the trough of the 4 m 639 contour within SOF-N and SOF-S deepened and moved south by 300 - 400 m (Szczyrba 640 et al., 2023a). This might explain the southerly offset between the origin of the morpho-641 logically controlled rip currents and the long-term shoreline change rates (Fig. 4). We 642 hypothesize that bathymetric hysteresis causes the erosional hot spot to lag behind the 643 southerly movement of the SOFs and, therefore, in the future the hot spot will migrate 644 south. When comparing the long-term shoreline change rate data released in 2004, 2013, 645 and 2020, there is some evidence that the hot spot is intensifying and expanding south-646 wards as the SOF troughs have deepened and moved south, although these data are cu-647 mulative long-term averages that do not highlight year to year changes. Nevertheless, 648 according to these data, between 2004 and 2020 the rate of erosion at SOF-N increased 649 from -0.36 m/yr to -0.70 m/yr and the SOF-S from -0.79 m/yr to -1.00 m/yr. The ar-650 eas affected have also expanded south by 100 m (North Carolina Division Of Coastal Man-651 agement, 2021). 652

However, this hypothesis is complicated by the beach nourishment projects that 653 occur in this region nearly every five years. A monitoring report conducted one year af-654 ter the conclusion of the 2017 beach nourishment project in KH observed that that up 655 to 60% of the sand volume loss occurred adjacent to the SOFs, although the authors cited 656 challenges in calculating sediment volumes around the complex bathymetric structures 657 (APTIM, 2019). The report concluded that future monitoring should investigate the im-658 pact of these SOFs on nourishment performance. The results of this study suggest that 659 the SOFs act as a conduit for enhanced cross-shore exchange of sediment because, dur-660 ing periods of moderate wave heights and angles, enhanced flow velocities are directed 661 offshore down the SOF edges and persist beyond the surf zone edge. During periods with 662 strong longshore currents, the flow accelerates and meanders near the SOFs, which would 663 further enhance sediment transport away from the SOF and contributing to the erosional 664 hot spot. These findings emphasize the importance of both cross- and longshore trans-665 port on nearshore morphological evolution, supporting the conclusions of Thieler et al. 666 (1995) and Gutierrez et al. (2005) at other sites. 667

The storm-induced wave forcing addressed in this research also affects shoreline change on an inter-annual to decadal time scale (Splinter et al., 2014). As climate change alters ocean temperatures, the intensity and duration of the Atlantic hurricane season may increase (Knutson et al., 2010; Walsh et al., 2016). This elevates the likelihood of sequential high-energy wave events (i.e., storm clusters) that inhibit beach recovery (Coco et
al., 2014) and result in higher rates of shoreline recession (Dodet et al., 2019). Geologically inherited complex bathymetry produces variability in beach response and recovery, whereby influencing the pattern of barrier island transgression over time (Houser,
2012). With a more intense wave climate, complex nearshore bathymetric features will
continue to drive spatially variable circulation patterns and affect erosional hot spots.

778 7 Summary and Conclusions

The influence of complex bathymetry on nearshore flow dynamics in an area with 679 erosional hot spots was explored with a combination of remotely-sensed data and a nu-680 merical model. An X-Band Radar (XBR) system was deployed at the study site in Kitty 681 Hawk, North Carolina, zfor nine days at the end of September, 2017. This period coin-682 cided with the passage of Hurricane Jose and Hurricane Maria, and these events were 683 also simulated with a phase-resolving numerical model. In total, 24 simulations were per-684 formed: 3 selected times during storm events (J, M1, M2) that were also observed by 685 XBR and 21 idealized simulations that explored a range of incident wave conditions. The 686 event simulations were validated against XBR observations of wave breaking patterns, 687 celerity, and representative wave height estimates. 688

The XBR observations were used to identify two zones of persistently high sea-surface 689 backscatter during energetic periods $(H_s \ge 2 \text{ m})$, indicative of strong offshore flows. Across 690 a wide variety of incident wave conditions, the offshore flows re-emerged near shore oblique 691 features (SOFs) and are concluded to be bathymetrically controlled rip currents. The 692 wave-current interactions associated with these offshore flows roughens the sea surface. 693 enabling them to be easily identified with XBR, and also increases the wave frequency 694 in localized zones. The location of the offshore currents are also linked to the incident 695 wave direction and the bathymetry. When waves approach the shoreline at a positive 696 angle, the offshore currents flow along the northern SOF edges while when wave approach 697 at a negative angle, the flows route along the southern edges. The rip currents result from 698 wave breaking across the alongshore-varying bathymetry near the SOFs, supported by 699 both XBR and numerical modeling data. Numerical simulations indicate that relative 700 wave heights and wave setup also vary across the cross-section of the SOFs. When the 701 wave field is oblique, a strong alongshore current represses offshore flows, and this cur-702 rent accelerates and meanders where the SOF contours pinch towards the shoreline. 703

Idealized simulations indicate that the complex bathymetry exerts a strong con-704 trol over nearshore wave heights, refraction patterns, and mean currents. The XBR ob-705 servations were used to identify morphologically-driven persistent offshore-directed flows 706 just south of long-term erosional hot spots. Future studies, using other numerical mod-707 els that simulate sediment movement and bed elevation change, could directly model the 708 morphodynamic evolution of zones with complex bathymetry to evaluate impacts on ero-709 sional hot spot intensity and evolution. Persistent offshore flows and alongshore flow ac-710 celeration may exacerbate erosion and enhance the shoreline undulations at this site. This 711 underscores the importance of both cross-shore and alongshore sediment transport on 712 the nearshore morphology and shoreline evolution on sandy beaches with nearshore bathy-713 metric features. 714

715 8 Data Availability Statement

All model inputs, radar data, processed model output data, and figure generation
 MATLAB codes used in this study are hosted on the Borealis data repository at Queen's
 University titled "Nearshore Flow Dynamics – SWASH and XBand Radar", via DOI:
 https://borealisdata.ca/privateurl.xhtml?token=fa6593ed-6ea6-437c-8734-4711be8950c8

- ⁷²⁰ with data use license agreement CC-BY 4.0. This is a private link for reviewers only, and
- the public link will be included upon manuscript acceptance.

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- versity, and Dr. Spicer Bak as well as Dr. Kate Brodie from the U.S. Army Engineer Re search and Development Center.
- 732 Acronyms
- 733 **AWAC** Acoustic Wave and Currents profiler
- ⁷³⁴ **FRF** Field Research Facility
- 735 **KH** Kitty Hawk
- 736 **RIOS** Radar Inlet Observation System
- ⁷³⁷ **R.I.** Return intensity of X-band radar
- 738 **RMSE** Root mean square error
- ⁷³⁹ **SOF** Shore-oblique feature
- 740 **SOF-N** Northern shore-oblique feature
- ⁷⁴¹ **SOF-S** Southern shore-oblique feature
- 742 **SWASH** Simulating WAves till SHore
- 743 **TKE** Turbulent kinetic energy
- 744 XBR X-band radar

745 Notation

- α_{ad} Calibration coefficient, Eq. 12 from Streßer et al. (2022)
- C_{p} Celerity of peak frequency waves [m/s]
- $_{748}$ **D**₅₀ Median grain size [mm]
- η_{49} η Water surface elevation [m]
- f Wave frequency [Hz]
- 751 γ Breaking index (0.78)
- $_{752}$ **g** Gravitational acceleration $[m/s^2]$
- $_{753}$ **h** Water depth [m]
- H_p Height of peak frequency waves [m]
- 755 H_s Significant wave height [m]
- 756 k Wave number [1/m]
- 757 λ Wavelength [m]
- ⁷⁵⁸ $\boldsymbol{\theta}_{m}$ Mean wave angle [°]
- $au_{59} au$ Shear $[\mathrm{N}/m^2]$
- 760 T_p Peak period [s]
- $_{761}$ U Depth-averaged mean velocity [m/s]
- $_{^{762}}$ U_{*cr} Critical velocity [m/s]
- 763 $oldsymbol{Q}_{b}$ Breaking parameter

Table 1. Wave conditions during the hindcast simulated event (J, M1, M2) including significant wave height (H_s) , peak period (T_p) , total water level (η) , mean wave angle relative to perpendicular of the study site's mean shoreline (θ_m) , and peak directional spreading (σ_p) .

Label	UTC	H_s [m]	T_p [s]	$\eta~[{\rm m}]$	θ_m [°]	σ_p [°]
J	19-Sep-2017 20:00:00	2.8	10.0	0.9	3.9	36.3
M1	27-Sep-2017 09:00:00	4.0	13.3	0.7	16.9	23.2
M2	28-Sep-2017 02:00:00	2.6	10.3	0.8	-37.1	58.5

Measured at the FRF's 11 m AWAC.



Figure 1. (a) Study region location within the Outer Banks, (b) June 2017 bathymetry measured offshore of Kitty Hawk, North Carolina where the red star represents the XBR location and labels N and S identify the locations of SOF-N and SOF-S, respectively, and (c) rotated bathymetric grid used for numerical modeling.



Figure 2. Wave conditions throughout the study period: (a) significant wave height (H_s) , (b) peak period (T_p) , and (c) mean wave angle relative to perpendicular of the mean shoreline (θ_m) . Wave angles are described as going towards, meaning positive values indicate an approach angle north of shore-perpendicular and negative values indicate an approach angle south of shore-perpendicular. Vertical dashed lines indicate the simulated hours during Hurricanes Jose and Maria (J, M1, M2) and shaded regions are periods observed by XBR.



Figure 3. Wave spectra for time periods (left column) J, (middle column) M1, and (right column) M2: (a, b, c) energy-density, (d, e, f) directional distribution of wave angles, and (g, h, i) directional energy-density spectra used to force simulations at the offshore boundary. Red lines in (g, h, i) are the mean direction per frequency calculated from Kuik et al. (1988). Wave angles are described as going towards, meaning positive values indicate an approach angle north of shore-perpendicular and negative values indicate an approach angle south of shore-perpendicular.



Figure 4. Alongshore variability in: (a) the long-term average shoreline change rate over 1940-2016 with shaded regions indicating the troughs of SOF-N and SOF-s;(b) XBR backscatter intensity anomalies for J (blue line), M1 (black line), M2 (orange line), and the mean during the observation period (thick black line); (c) time-averaged XBR backscatter return intensity (R.I.) during event M2 with the white lines indicating the corresponding XBR estimated offshore edge of the surf zone; and (d) time-averaged XBR backscatter during a moderate energy event ($H_s = 2.3$ m) on September 20, 2017 at 04:13 UTC.



Figure 5. Surf zone detection from XBR observations and model: (a) XBR measurements of the cross-shore location of the offshore surf zone edge during event J (blue line), M1 (black line), and M2 (orange line) with shaded regions indicating the SOF troughs; (b) XBR-derived M1 surf zone edge (black line) over the simulated fraction of breaking waves (Q_b) for M1.



Figure 6. XBR observations and model results of wave celerity at the peak frequency (C_p) during M2: (a) XBR observations in areas deeper than 2 m, (b) SWASH results, and (c, d, e) cross-shore transects of modeling results (solid lines) and XBR observations (dashed lines) at alongshore locations (c) y = 2250 m, (d) y = 1500 m, and (e) y = 750 m. White dashed lines in (a) and (b) indicate the cross-shore transect locations.



Figure 7. Estimates of the peak wave height (H_p) for event M2 from (a) XBR observations and (b) modelling results. Solid black lines in (a) and (b) represent the XBR-measured surf zone edge of M2.



Figure 8. Wave height estimates from numerical simulations: (a) the spatial distribution of relative wave heights (H_s/h) during M1, (b) maximum H_s divided by mean H_s at each alongshore location for J (blue line), M1 (black line), and M2 (orange line), (c) maximum alongshore gradient in setup, and (d) saturated frequency map during M1. Gray shaded regions in (b) and (c) indicate the approximate bounds of the shore-oblique troughs.



Figure 9. Flow dynamics around the northern shore-oblique feature between y = 1500 m -2500 m for time periods (left column) J, (middle column) M1, and (right column) M2. Flow dynamics displayed include: (a, b, c) wave angle (θ) relative to shore-perpendicular, (d, e, f) cross-shore velocity (u), (g, h, i) alongshore velocity (v), and (j, k, l) overall depth averaged current velocity (U).



Figure 10. Fraction of breaking waves (Q_b) calculated from $H_s = 3$ m idealized simulations with various incident peak wave angles (θ_p) : (a) 15°, (b) 10°, (c) 5°, (d) 0°, (e) -5°, (f) -10°, and (g) -15°.



Figure 11. Depth-averaged current velocities (U) calculated from $H_s = 3$ m idealized simulations with various incident peak wave angles (θ_p): (a) 15°, (b) 10°, (c) 5°, (d) 0°, (e) -5°, (f) -10°, and (g) -15°.



Figure 12. Locations of maximum U (a) and locations of maximum variability of U (b) during a variety of incident wave angles and significant wave heights (H_s) , where circles represent an incoming H_s of 1 m, triangles represent H_s of 3 m, and stars represent H_s of 5 m. Quantified influence of incoming wave angle (θ) and H_s on maximum U (c-e) and the variability of U (f-h) in the surf, intermediate, and offshore zones. Cross-shore locations of the surf, intermediate, and offshore zones are depicted in (a) and (b) as white dashed lines.

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Nearshore flow dynamics over shore-oblique bathymetric features during storm wave conditions

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Key Points:

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9	•	Radar and numerical modeling indicate variability in surf zone width and wave
10		height
11	•	Offshore-directed flows occur in a region with shore-oblique sand bars and erosional
12		hot spots
13	•	Morphologic influence and directional variability are quantified with idealized model
14		simulations

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15 Abstract

Shore-oblique bathymetric features occur around the world and have been statistically 16 correlated with enhanced shoreline retreat on sandy beaches. However, the physical mech-17 anisms that explain a causal relationship are not well understood. In this study, radar 18 remote sensing observations and results from a phase-resolved numerical model explore 19 how complex morphology alters nearshore hydrodynamics. Observations at selected times 20 during high-energy storm events as well as a suite of idealized simulations indicate that 21 shore-oblique features induce strong spatial variations in the water surface elevation and 22 wave breaking patterns. Re-emergent offshore flows and longshore current accelerations 23 occur near the apex of the oblique nearshore features. The results suggest that complex 24 bathymetric morphology exerts a powerful control on nearshore hydrodynamics and in-25 creases the potential for enhanced cross-shore and alongshore sediment transport, thus 26 contributing to localized erosional zones. 27

²⁸ Plain Language Summary

Near the shoreline, underwater topography is affected by sea level, waves, currents, 29 tides, and geological characteristics. Typically, sandbars are oriented parallel to the coast-30 line, but shore-oblique sandbars have also been identified around the world. In many in-31 stances, these oblique features have been correlated with zones of erosion, however the 32 explanation for this statistical relationship is not fully understood. In this study, remote 33 sensing observations and modeling results explore how complex underwater topographies 34 alter coastal wave energy and flow patterns. Selected stormy periods were observed in 35 addition to a set of idealized simulations. Results indicate that shore-oblique features 36 cause localized changes to wave heights, wave breaking, and current speeds. These al-37 terations contribute to the presence of fast offshore-directed rip currents and accelera-38 tions in shore-parallel currents, which would enhance the potential for sand to be trans-39 ported away from these zones. The findings suggest that the changes to the flow field 40 induced by the oblique features contribute to zones of high rates of erosion. 41

42 **1** Introduction

The coastal nearshore region, defined as the transition zone between the shoreline 43 and the inner shelf, is a dynamic zone shaped by the interplay of many physical processes 44 occurring at various spatial and temporal scales. The level and distribution of incom-45 ing wave energy exerts a powerful control on nearshore forces and alters currents, sed-46 iment transport, and morphology (Wright & Short, 1984; Svendsen, 1984; Gallagher et 47 al., 1998). When propagating at oblique incident angles to the bathymetric contours, sur-48 face waves transfer momentum and drive longshore currents, moving material along the 49 coastline (Longuet-Higgins, 1970; Guza et al., 1986). These currents meander and change 50 velocity in response to bathymetric variability (Garnier et al., 2013). At moderate in-51 cident angles, gradients in wave-driven setup of the mean water level can cause local along-52 shore flows to converge and flow offshore as rip-currents (Castelle et al., 2016; Moulton 53 et al., 2017). 54

These strong offshore-directed rip currents can scour channels in nearshore sand-55 bars and alter patterns in wave dissipation, causing the local mean water level to slope 56 towards the rip channel (Haller et al., 2002). This feedback reinforces offshore flows and 57 can erode the beach profile (Komar & McDougal, 1988). For weak offshore flows, wave 58 energy dissipation can be reduced in rip-channels, leading to higher wave heights that 59 break close to the shoreline (Haller et al., 2002) and increased vulnerability to erosion 60 (Holman & Sallenger, 1993). In some cases, stable patterns in the underlying bathymetry 61 (e.g., submarine canyons) produce morphologically-driven recurring offshore flows by redi-62 recting wave energy (Long & Ozkan-Haller, 2005; Magne et al., 2007; O'Dea et al., 2021). 63 For example, dramatic variations in refraction and wave energy occur across the Scripps 64

Canyon in southern California, with increased local wave heights at the head of the canyon
 (Magne et al., 2007).

Morphologically-driven processes are often shaped by the underlying geology (also 67 referred to as antecedent or framework geology). The geologic framework, generated over 68 very long time scales, is characterized by deposits or bedrock underlying modern sed-69 iments to a depth of approximately 10 m (Browder & McNinch, 2006). Along the east 70 coast of North America, anomalous nearshore bathymetric features have been identified 71 and connected with geological history: shore-normal or shore-oblique troughs and bars 72 73 occur in New York (Schwab et al., 2000), New Jersey (Snedden et al., 1994), Virginia (Colman et al., 1990; Browder & McNinch, 2006), North Carolina (McNinch, 2004), Florida 74 (Houser et al., 2008; Barrett & Houser, 2012), northeastern New Zealand (Green et al., 75 2004), Prince Edward Island (Wernette & Houser, 2022) among many other locations 76 and are commonly related to paleo-river channels. These features alter alongshore pat-77 terns in wave breaking (Safak et al., 2017), modify the distribution of wave energy (Mulligan 78 et al., 2019a), and have been hypothesized to act as conduits for offshore sediment trans-79 port (Thieler et al., 1995) and enhanced alongshore advection (Gutierrez et al., 2005). 80 However, the physical mechanisms linking framework geology to alterations in hydro-81 dynamics and morphodynamics across the shoreface are not well understood. 82

In this study, we investigate the influence of a complex bathymetric framework on 83 nearshore flow dynamics during storm conditions at Kitty Hawk, North Carolina, USA. 84 We hypothesize that shore-oblique bathymetric features alter mean nearshore currents 85 and induce localized zones of higher wave height and wave setup on the shoreface, which 86 may affect sediment transport patterns and the morphology of this nearshore region. Both 87 real and idealized storm events are simulated with a phase-resolving numerical model 88 for a wide range of incident wave angles. The event-based simulations focus on time pe-89 riods with large waves corresponding to Hurricane Jose and Hurricane Maria in 2017, 90 that were also observed with an X-band radar (XBR) remote sensing system. The ide-91 alized simulations isolate the effects of the bathymetry and incident wave angle on the 92 nearshore hydrodynamics. The novel combination of remote sensing and numerical re-93 sults are synthesized to provide new insight on wave-driven flows over complex geologically-94 controlled bathymetric features. 95

⁹⁶ 1.1 Regional setting

Kitty Hawk (KH) is located on the Outer Banks of North Carolina (Fig. 1a). The 97 Outer Banks are a series of long barrier islands segmented by inlets that formed in a mi-98 crotidal, wave-dominated environment (Hayes, 1979), and are subject to storms (e.g., 99 Tropical Cyclones and Nor'Easters) and high rates of sea level rise (Sallenger et al., 2012; 100 Kemp et al., 2017). These barrier islands have been shaped into cuspate forelands that 101 divide large estuaries (Albemarle and Pamlico Sounds) from the Atlantic Ocean. Cur-102 rently, few inlets connect the sounds to the ocean and the barrier islands receive a very 103 limited supply of sediment from riverine sources (Culver et al., 2007; Mulligan et al., 2019b). 104 Located 5 km west of the modern transgressive shoreline at KH, a progradational beach 105 ridge complex formed 3 ka - 2 ka before present during a period of rapid Holocene sea 106 level rise and abundant sediment availability (Mallinson et al., 2008). The present-day 107 coastal morphology and dynamics are influenced by the complex underlying regional ge-108 ologic framework. 109

¹¹⁰ 2 Shore-oblique features

The segment of coast is dominated by shore-parallel depth contours except near KH, where the nearshore region includes a series of bathymetric undulations with bars and troughs as shown in Fig. 1b. These bathymetric variations are situated in mean water depths ranging from 2-12 m and are oriented offshore at an average angle of 42° from

the shoreline (Schupp et al., 2006). The morphologic features are referred to as shore-115 oblique features (SOFs) and, within the region considered in the present study, include 116 a northern (SOF-N) and southern (SOF-S) trough shown in Fig. 1b. Seismic imaging 117 studies have associated these bathymetric anomalies to underlying Pleistocene paleo-channels, 118 specifically the paleo-Roanoke River (Boss et al., 2002; McNinch, 2004; Browder & Mc-119 Ninch, 2006). Cycles of eustatic sea-level changes that occurred concurrently with glacial 120 episodes during the Pleistocene epoch dissected numerous fluvial channels into the un-121 derlying Quaternary strata (Boss et al., 2002). During the ensuing Holocene transgres-122 sion, these channels were drowned and infilled with muds, peats, sands, and gravels (Riggs 123 et al., 1992, 1995; Schwartz & Birkemeier, 2004). 124

The SOFs are relatively stable but can migrate approximately 250 m alongshore 125 in response to individual storm events (McNinch & Miselis, 2012). Schupp et al. (2006) 126 found that all movement was confined to zones shoreward of the -9 m bathymetric con-127 tour. They also slowly migrate downdrift on a decadal timescale. Between 2004 and 2017, 128 SOF-N migrated south by approximately 600 m and deepened by 0.5 m (Szczyrba et al., 129 2023a). Regionally, these SOFs have been correlated with areas of high shoreline vari-130 ability and, on longer time scales, high long-term shoreline change rates. McNinch (2004) 131 visually identified a correlation between the SOFs and high shoreline change rates and 132 Schupp et al. (2006) quantified this correlation on a regional scale as statistically signif-133 icant. However, the nearshore hydrodynamic processes that physically link the SOFs with 134 erosion have not been well defined. 135

3 Observations

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3.1 Bathymetry and wave data

In June, 2017, bathymetric profiles were collected to support the local beach nourishment monitoring plan (APTIM, 2017). The nearshore survey lines were collected by a survey vessel and spaced approximately 150 m meters apart, with sounding points sampled every 7-8 m (APTIM, 2017). The survey data were interpolated onto a 5 m x 5 m regular grid in ArcGIS 10.7.1, encompassing two shore-oblique bathymetric features (SOF-N and SOF-S). This bathymetric grid, shown in Fig. 1c, is used as input to the numerical model.

Wave observations used as boundary conditions in the numerical model were sourced 145 from the US Army Corps of Engineers Field Research Facility (FRF) at Duck, located 146 14 km north of KH. Wave spectra were observed by a 1 MHz Nortek Acoustic Wave and 147 Current (AWAC) profiler, located offshore in a nominal depth of 11 m. Previous stud-148 ies have shown that the bulk wave statistics including significant wave height (H_s) , mean 149 direction (θ_m) , and peak period (T_p) are highly correlated at the 11 m depth over a dis-150 tance of 33 km between the FRF and Jennette's Pier, located 20 km south of KH (Mulligan 151 et al., 2019a). Wave data were obtained for key times during major wave events in Septem-152 ber 2017. These correspond to a time during Hurricane Jose (denoted "J"; 19-September-153 2017 20:00:00 UTC), and two times during Hurricane Maria (denoted "M1" and "M2"; 154 27-September-2017 09:00:00 UTC and 28-September-2017 02:00:00 UTC) indicated in 155 Fig. 2 and Table 1. Directional energy-density wave spectra were used as offshore bound-156 ary conditions for the numerical model (Fig. 3). The directional spectra were calculated 157 from the first four Fourier coefficients, using the maximum entropy method (Lygre & Krogstad, 158 1986) with a directional resolution of 1° , a frequency resolution of 0.0075 Hz, and a fre-159 quency range of 0.0400-0.4975 Hz. The spectra were rotated 28° clockwise of true north 160 to match the orientation of the rotated bathymetric grid before they were input into the 161 model. 162

The times corresponding to J, M1, and M2 were selected because they have similar mean total water levels, and a range of energy levels and incident wave angles (Ta-

ble 1 and Fig. 3). Hurricane Jose at time J was a moderate wave height $(H_s = 2.8 \text{ m})$ 165 event with a positive incident angle (meaning north of shore-perpendicular) and mod-166 erate directional spreading (Fig. 3a, d, g). Hurricane Maria at time M1 was a high-energy 167 event, with an H_s of 4.0 m, long period waves, low directional spreading and an oblique 168 positive incident angle (Fig. 3b, e, h). As Hurricane Maria propagated along the coast-169 line, the dominant wind and wave direction shifted and the conditions at time M2, which 170 occurred 17 hours after M1, had moderate wave heights $(H_s = 2.6 \text{ m})$ and periods dur-171 ing the passage of the storm. Time M2 was characterized by high directional spreading 172 from the high-frequency oblique wind-waves forced by strong offshore winds (Fig. 3c, f, 173 i). 174

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3.2 X-band radar collection and processing

The Radar Inlet Observation System (RIOS) is a fully automated non-Doppler XBR 176 remote sensing system housed within a mobile trailer for rapid deployment (McNinch 177 et al., 2012; Humberston et al., 2019). RIOS emits microwave radio signals with a hor-178 izontal beam width of 1.2° and a vertical beam width of 25° and records the subsequent 179 backscatter. Bragg resonance is the primary mode of XBR backscatter, wherein radar 180 pulses interact with short capillary waves (1-2 cm wavelength, λ) on the water surface. 181 Backscatter increases when waves steepen, modulating the short-gravity wave field, and 182 become roughened as they approach the shoreline and break (Catalán et al., 2014). Thus, 183 RIOS provides a digital reconstruction of the incoming nearshore wave field as a series 184 of sequential spikes in the XBR backscatter return intensity. XBR can resolve the spatio-185 temporal evolution of incoming nearshore waves and the observations compare well to 186 both numerical models and in-situ data (Szczyrba et al., 2023b). Additional RIOS-system 187 details are described by Humberston et al. (2019) and McNinch et al. (2012). 188

RIOS was deployed in the study site at the location shown in Fig. 1b (36.066408°. 189 -75.690222°) in September, 2017. Nearshore sea surface conditions were observed for 14 190 minutes per hour between September 19-23 and September 27-30, capturing J, M1, and 191 M2 time periods. The XBR antenna completed a rotation every 1.67 s, resulting in a sam-192 pling frequency of 0.60 Hz. Raw backscatter data were transformed from polar coordi-193 nates onto a 5 m x 5 m Cartesian grid (Humberston et al., 2019) and smoothed in the 194 alongshore direction by 100 m using a multidimensional image averaging filter (imfilter). 195 The resulting grid was rotated 28° clockwise of true north to match the orientation of 196 the numerical modeling framework. Overall, the XBR footprint spanned 3 km alongshore 197 and 1 km offshore. The time-dependant surf zone width was calculated from 14 minute 198 time-averaged XBR backscatter data. Time-averaged XBR and optical data highlight 199 zones of high dissipation and are often used to determine the spatial patterns of wave 200 breaking (Lippmann et al., 1993; Brodie et al., 2018). Following the method described 201 by O'Dea et al. (2021), the offshore surf zone edge was identified as the maximum cross-202 shore gradient of the XBR backscatter intensity, calculated at each alongshore location. 203 Outliers were removed and the resulting surf zone edge lines were smoothed using a mov-204 ing mean filter with a window size of 100 m. 205

Methods from O'Dea et al. (2021) were also applied to search for and identify mor-206 phologically driven offshore-directed currents in the study region. This approach ana-207 lyzes backscatter intensity offshore of the surf zone where few waves are breaking. Off-208 shore flows escaping beyond the breaker line can be identified in the XBR data as alter-209 ations in surface roughness. Offshore directed rip currents enhance sea surface rough-210 ness, incite short-scale wave breaking (i.e. microbreaking), and elevate XBR-measured 211 backscatter intensity (Lyzenga, 1991; Plant et al., 2010; Haller et al., 2014; O'Dea et al., 212 2021). This effect has been confirmed in a study utilizing both XBR and GPS-equipped 213 floating drifters (Takewaka & Yamakawa, 2011) and with a cross-shore array of current 214 meters (Haller et al., 2014). To identify these zones of wave-current interaction, an along-215 shore transect of backscatter intensity was extracted at a location 100 m offshore of the 216

²¹⁷ surf zone edge during periods with waves with $H_s \geq 2$ m. The mean backscatter inten-²¹⁸ sity was removed from each transect to highlight zone of anomalously high return inten-²¹⁹ sities. The intensity anomaly transect data were then averaged together to identify per-²²⁰ sistent XBR backscatter anomalies in the 9-day observation period. The XBR anoma-²²¹ lies for J, M1, and M2 were also analyzed.

XBR observations were used to determine representative wave phase speeds. Pixel 222 intensity time series were input into the bathymetric-inversion algorithm cBathy (Holman 223 et al., 2013b; Holman & Bergsma, 2021). cBathy is typically used to estimate bathymetry 224 225 from optical imagery, often from Argus camera systems (e.g., Oades et al. (2023)), but can also process XBR data (Honegger et al., 2019). Resulting XBR-derived bathymet-226 ric surfaces have been applied in numerical modeling studies (O'Dea et al., 2021). The 227 algorithm estimates the wave numbers (k) of dominant wave frequencies (f) and uses 228 the linear dispersion equation to relate these k-f pairs to mean water depth. The k-f 229 pairs are extracted within mapped tiles by calculating the Fourier phase cross-spectra 230 between each pixel per tile. In the present study, tile spacings were set to be 10 m in the 231 cross-shore direction and 25 m in the alongshore direction. The algorithm produces a 232 matrix of wave numbers for each the most coherent 4 frequencies wherein the individ-233 ual matrices contain a range of frequencies that represent the most dominant frequency 234 calculated per pixel. 235

The most coherent k-f matrix was extracted to create wave speed maps of repre-236 sentative peak wave conditions for each event. cBathy's accuracy declines significantly 237 in very shallow water with h < -2 m (Honegger et al., 2019), therefore these regions were 238 excluded from further analyses. Cross-shore transects were extracted to support com-239 parisons with the modeling data with a 20 m moving mean applied to eliminate noise. 240 Using the peak frequency calculated from cBathy, the method of Streßer et al. (2022) 241 was adapted to convert XBR-measured wave phase speeds to representative wave height 242 estimates (see Eq. 12 in Streßer et al. (2022)). Their wave-by-wave approach applies a 243 physically derived scaling that relates changes in breaking phase speed to wave height 244 245 via the amplitude dispersion relation:

$$H_p = \frac{C_p^2}{g(\frac{1}{\gamma} + \alpha_{ad})} \tag{1}$$

This equation was applied using the celerity of the peak frequency (C_p) , the default breaking index (γ) value of 0.78 (Streßer et al., 2022; Larson & Kraus, 1994), and the recommended calibration coefficient (α_{ad}) value of 0.5 (Streßer et al., 2022) to estimate a representative wave height based on peak frequency values (H_p) . This method is valid within the surf zone, since it was developed based on the modified nonlinear shallow water phase speed for waves in very shallow water (i.e., $h:\lambda < 1/10$) from Hedges (1976).

²⁵³ 4 Numerical model

The phase-resolving 3D numerical model Simulating WAves till SHore (SWASH) 254 applies the nonlinear shallow water horizontal momentum equations and the nonhydro-255 static vertical momentum equation using a finite difference scheme (Zijlema et al., 2011; 256 Zijlema, 2020) to simulate the water surface and velocity fluctuations. SWASH was ap-257 plied to simulate several times (J, M1, M2) in the study period that correspond to radar 258 data collection times. In addition to the selected events, an idealized suite of 21 simu-259 lations were also performed to isolate key variables and specifically explore how the in-260 cident wave conditions affect the nearshore hydrodynamics. 261

4.1 Model setup

SWASH has been used to simulate wave-driven currents over variable bathymetry, 263 such as over submerged reefs (da Silva et al., 2023). Previous studies have also applied 264 SWASH to model regions near the present study area (Gomes et al., 2016; Mulligan et 265 al., 2019a; Szczyrba et al., 2023b) and similar numerical input parameters were applied 266 in this study. A 5 m x 5 m regular structured grid, extending 1000 m offshore and 3000 267 m alongshore, was constructed. A 300 m extension was added to the offshore edge of the 268 bathymetric grid (from x = 1000-1300 m) to create a uniform offshore region and pre-269 270 vent numerical anomalies from being introduced due to depth variations at the east boundary (without the extension, the bathymetric variations at the model boundary causes 271 numerical instabilities that are not realistic). The extension was created by linearly in-272 terpolating the offshore bathymetric edge to a constant depth of h = -12 m over a cross-273 shore distance of 100 m and then extending this constant depth over the remaining 200 274 m, following previous studies (O'Dea et al., 2021). The bathymetry was also extended 275 by 500 m in each alongshore direction by interpolating the sides to a uniform contour 276 configuration in order to enable the application of periodic alongshore boundary condi-277 tions in the numerical model. The bathymetry data was rotated 28° clockwise of true 278 north which oriented the mean shoreline parallel to the y-axis of the grid (Fig. 1c). To 279 resolve depth-dependent dynamics, three bathymetry-following vertical layers with equal 280 thickness were included. This vertical resolution is sufficient to resolve wave frequency-281 dispersion (Zijlema & Stelling, 2005; Smit et al., 2013). To confirm this, a sensitivity test 282 using 9 vertical layers was conducted to explore the flow structure in and around the SOFs 283 in further detail, and found no substantial differences from the 3-layer runs. 284

The Sommerfeld radiation condition was applied at the onshore boundary to ap-285 propriately limit wave reflection and the offshore boundary was weakly reflective. The 286 alongshore boundary was periodic, meaning that energy exiting one alongshore bound-287 ary re-entered the domain at the opposite alongshore boundary. The first 15 minutes of 288 modeling time were disregarded as model spin-up and the remaining 45 minutes were 289 analyzed. The wavemaker generated random wave time series that statistically matched 290 the input wave parameters with a cycle interval of 2700 seconds (45 minutes). Accord-291 ing to Zijlema et al. (2011), the cycle time should span 100-300 peak wave periods to pro-292 vide accurate statistical results and the selected cycle time spans over 200 wave periods 293 of the longest peak period simulated (M1). The model performed calculations with an 294 initial time step of 0.05 s that is automatically adjusted throughout the simulation ac-295 cording to the Courant number, which can reach a maximum of 0.5. Total water levels 296 were adjusted for the correct tidal stage and storm surge present during the simulated 297 events. Event times during hurricane wave events J (Jose), and M1, M2 (Maria) were 298 forced at the offshore boundary using the directional energy-density spectra for each cor-299 responding time (Fig. 3g-i). 300

4.2 Model output processing

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To enable direct comparison with the XBR cBathy products, spectral processing 302 extracted peak wave frequency parameters per grid cell from the J, M1, and M2 simu-303 lations. The fast Fourier transform was calculated at each grid cell with 256 samples and 304 a Hanning window with 50% overlap. The peak wave frequency per grid cell was also 305 extracted and used to calculate k and λ with the dispersion relation (Fenton & McKee, 306 1990; Dean & Dalrymple, 1991). The representative C_p was then calculated and input 307 into Eq. 1. This approach emulates the XBR processing steps. H_s was calculated across 308 the computational domain as four times the standard deviation of the water surface el-309 evation output (Raubenheimer et al., 1996). The spatial gradient of the water surface 310 elevation data highlighted wave crests as locations with steepest slopes. Vectors perpen-311 dicular to these crests at each grid cell were then time-averaged to provide local estimates 312 of mean wave angles across the study region (Szczyrba et al., 2023b). A dimensionless 313

wave breaking parameter (Q_b) was calculated from the binary wave breaking locations output from the model (Gomes et al., 2016). This parameter represents the relative intensity of breaking independent of the model time step and spatial resolution. These breaking zones were compared to the surf zone widths observed by XBR.

318 5 Results

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5.1 Event flow dynamics

Regions of persistently high XBR backscatter outside of the surf zone were iden-320 tified for J, M1, and M2 as well as the mean during all periods with $H_s \geq 2$ m observed 321 by XBR (Fig. 4). The strongest XBR anomalies occur at the southern edge of SOF-N 322 and the southern edge of SOF-S, located at y ~ 1600 m and y ~ 250 m (Fig. 4b-d). Each 323 event, as well as the study-period mean, indicate anomalously high backscatter signals 324 at these locations. Offshore flows enhance XBR backscatter intensity where wave-current 325 interactions roughen the sea surface. Even during lower energy periods, high backscat-326 ter indicative of offshore flow extends outside of the surf zone. For example, as shown 327 in Fig. 4d, at a time when $H_s = 2.3$ m, high backscatter return intensity (RI) is evident 328 at y ~ 1600 m and y ~ 250 m. The signals along the southern edges of both SOFs are 329 strongest during M2, an event with negatively incident (southerly) waves. Spikes in XBR 330 backscatter also occur along both the northern and southern edges of SOF-S during M2, 331 and in the channel trough backscatter measures abruptly weaken (Fig. 4b, c). In the mean 332 observations, this U-shaped pattern in the XBR anomaly can be observed in both SOFs 333 and is much wider in SOF-S than SOF-N. This U-shape indicates increased backscat-334 ter along both edges of each SOF, with the southern edge inducing stronger backscat-335 ter. M1, an event with large waves approaching the shoreline at a highly oblique, pos-336 itive (northerly) angle, only exhibits backscatter anomalies along the southern edges of 337 the SOFs, and no anomaly is detected near the northern edges. The strongest XBR anoma-338 lies align approximately with regions of higher local shoreline change rates, i.e., erosional 339 hot spots (List et al., 2006), with a slight offset (Fig. 4a, b). Each hot spot is offset north 340 of the southern edge of SOF-N and SOF-S, where the highest XBR anomalies occur. 341

The observed and simulated extent of the surf zone are shown in Fig. 5. Event M1 342 was the highest energy event with the widest surf zone (Fig. 5a). Mean surf zone widths 343 are 223 m, 281 m, and 269 m for J, M1, and M2, respectively, however the extents vary 344 considerably in the alongshore direction. The surf zone of M1 is widest north of SOF-345 N, beginning offshore at x \sim 340 m. The surf zone of M2 is widest south of SOF-N, oc-346 curring at $x \sim 320$ m. During all events, surf zones narrow along the shoreward apex of 347 SOF-N and widen on either side, coincident with the feature's edges. SOF-N modulates 348 surf zone widths more than SOF-S. The offshore extent of the surf zone identified by XBR 349 (Fig. 5a) was compared to the breaking parameter (Q_b) calculated from the numerical 350 model (Fig. 5b). The XBR-sensed surf zone edge aligns with the offshore edge of the most 351 intense breaking locations in the SWASH simulations and displays similar lateral vari-352 ability in break-point locations. During M1, Q_b is most intense at x ~ 200 m, y ~ 2,000, 353 which is located along the northern edge of SOF-N (Fig. 5b). 354

Spatial patterns in C_p shown in Fig. 6 are similar in XBR observations (Fig. 6a) 355 and SWASH results (Fig. 6b), wherein C_p remains high within the shore-oblique troughs 356 as the waves approach the shoreline, coinciding with the bathymetric contours. Peak fre-357 quency waves enter the domain at 9-10 m/s and slow to 6-7 m/s along the 4 m bathy-358 metric line. Across a variety of cross-shore profiles, XBR and SWASH estimates agree 359 well (Fig. 6c-e). Mean C_p errors are 0.26, 0.26, and -0.04 m/s for J, M1, and M2, respec-360 tively, however errors vary spatially. In the cross-shore, error is highest shoreward of x 361 \sim 335 m where the model slightly underestimates C_p by 0.47 m/s when compared to XBR 362 observations. In the alongshore, the model underestimates C_p north of y ~ 1,750 m and 363

overestimates south of y \sim 1,500 m. Across all periods simulated, mean error is highest between y \sim 925 - 1,225 m with a maximum of -0.53 m/s.

Using Eq. 1, spatial distributions of C_p were converted into representative wave heights (H_p) to further compare XBR and model results as shown in Fig. 7. Patterns in H_p compare well between XBR and the model. Wave heights vary laterally across the study domain, with higher wave heights reaching closer to the shoreline in the lee of the troughs of the SOFs. Wave dissipation begins along the 4 m contour lines. Root mean square errors (RMSE) are 0.76 m, 0.72 m, and 0.65 m across the entire domain and 1.35 m, 0.61 m, and 0.57 m within the surf zone for J, M1, and M2, respectively.

Fig. 8 displays the spatial patterns in relative wave height H_s/h (Fig. 8a) and wave 373 frequency (Fig. 8b) alongside the alongshore distributions of the normalized maximum 374 H_s (Fig. 8b) and the water surface elevation gradient (Fig. 8c) The distribution of mod-375 eled H_s varies across space and by event (Fig. 8a, b). Relative wave heights are elevated 376 north and south of SOF-N around y ~ 2000 and 1500 m, and in these zones high rel-377 ative wave heights extend to the 6 m contour (Fig. 8a). H_s values increase near the SOFs 378 and decrease in the region between SOF-N and SOF-S (Fig. 8b). During all events, H_s 379 increases along the edges of the SOFs, following the bathymetric contours, and decreases 380 along the central axes (Fig. 8b). The alongshore differences in H_s at a particular cross-381 shore location varies between 0.2-0.4 m. The alongshore variability in wave heights also 382 induces strong variations in the water surface elevation gradient (Fig. 8c). At SOF-N, 383 the minimum alongshore gradient occurs at y = 1665 m and then sharply increases on 384 either side of that location across all events. The mean wave frequency is also elevated 385 along the southern edges of SOF-N and SOF-S, coincident with the abrupt alongshore 386 variations in wave setup (Fig. 8c), indicative of wave-current interactions. The wave fre-387 quency changes displayed in Fig. 8d appear similar to the XBR backscatter patterns dis-388 played in Fig. 4d. 389

Simulated wave angle and velocity results near SOF-N during J, M1, and M2 are 390 shown in 9. The SOFs modulate patterns of wave refraction during each event (Fig. 9a-391 c). M1 waves maintain positive approach angles across the domain even near the shore-392 line but are moderated closer towards shore-normal in the lee of SOF-N. J and M2 waves 393 refract into positive and negative directions, creating zones of convergence and divergence 394 around SOF-N. In both events, divergence occurs along the northern edge of SOF-N while 395 convergence occurs along the southern edge. A similar, but less severe effect, is observed around SOF-S. The features cause asymmetric wave refraction, whereby wave entering 397 at moderate angles refract away from the trough and toward the edges of the SOFs (Fig. 398 9a-c), increasing wave heights along the edges and sheltering the zone immediately in 399 the lee of the trough (Fig. 8b). 400

Across all three selected times, cross-shore velocities are directed offshore down the 401 oblique central axis of SOF-N (Fig. 9d-f), coinciding with the locations of wave diver-402 gence (Fig. 9a-c). These flows decelerate and broaden as they move offshore. Cross-shore 403 velocities route back onshore at the southern edge outside of the surf zone and this shoreward-404 directed patch occurs 50 m farther offshore during M1 than during J and M2. North of 405 SOF-N, cross-shore velocities are also directed onshore during J and M2, but this effect 406 is weak during M1. The offshore and onshore-directed velocities thus form asymmetric 407 circulation cells, with the stronger cell occurring along SOF-N's southern edge. The lo-408 cation of the maximum and minimum cross-shore velocities remain fairly stable across 409 all three periods simulated. Alongshore velocities accelerate and meander significantly 410 offshore near SOF-N in J and M2 (Fig. 9g-i). The acceleration in both events occurs near 411 412 $x \sim 250$ m, $y \sim 1750$ m. During M1, the alongshore velocities are very strong and meander slightly in response to the SOF contours. The longshore current reverses direc-413 tion during M2 at x \sim 300 m, y \sim 1950 m, and the nexus of the reversal coincides with 414 a location of high cross-shore velocity. 415

During J and M2, the mean flow is characterized by two circulation cells between 416 (1) $x \sim 0.500$ m, $y \sim 2000-2500$ m and (2) $x \sim 0.500$ m, $y \sim 1500-2000$ m (Fig. 9j-l). 417 The southern circulation cell is stronger than the northern cell. While water moves off-418 shore down the axis of SOF-N in J and M2, the mean current merely meanders in re-419 sponse SOF in M1 since the strong inertia of the longshore current overpowers poten-420 tial circulation cells. The offshore-directed flows in J and M2 maintain high velocities 421 offshore of the surf zone, with fast flows identifiable up to 2 surf zone widths away from 422 shoreline. Offshore of SOF-N at x = 600 m, the maximum velocities occur in similar lo-423 cations (U = 0.57 m/s, 0.90 m/s, 0.54 m/s at y = 1820 m, 1515 m, 1500 m for J, M1, 424 M2, respectively) and the minimum velocities also occur in similar locations (U = 0.45425 m/a, 0.79 m/s, 0.40 m/s at y = 2275 m, 2065 m, 2255 m for J, M1, M2, respectively). 426 The alongshore velocity component plays a larger role than the cross-shore velocity com-427 ponent in controlling the overall mean current. 428

429

5.2 Idealized effects of incident wave conditions

A second suite of idealized simulations isolated the effects of various wave conditions including incident θ_p and H_s on nearshore hydrodynamics. These simulations were forced by a JONSWAP spectra with a T_p of 12 s and zero directional spreading. The H_s values varied between 1.0, 3.0, and 5.0 m while θ_p varied between -15 to 15° in intervals of 5°. It should be noted that the ideal simulation with an H_s of 5.0 m and θ_p of 15° failed because large waves at highly oblique angles are challenging to simulate (Baker et al., 2021). Thus, results for that simulations pertain to an H_s of 4.9 m at θ_p of 15°.

Model results of Q_b across a variety of incident wave approach angles are shown 437 in Fig. 10. Incident waves that approach from $+15^{\circ}$ and transition to -15° emulate a com-438 mon pathway of tropical cyclones in this region and Q_b varies depending on the incom-439 ing θ_p (Fig. 10). Q_b is dispersed relatively evenly alongshore between the 6 m depth con-440 tour and the shoreline when incident waves approach at $+15^{\circ}$. However, Q_b begins fur-441 ther offshore along the edges of SOF-N and closer to the shoreline at the shoreward apex 442 of the SOF. Between $+10^{\circ}$ - 0° , Q_b begins to concentrate along the northern edge of SOF-443 N. The breaking region is most localized when incident waves approach at -5° , focused 444 at the shoreward apex of SOF-N (near x = 150 m, y = 1750 m) as well as SOF-S (near 445 x = 250 m, y = 500 m). North of SOF-N most breaking occurs at the shoreline. At -446 10° , waves break again across the northern sandbar. Finally, at -15° , breaking is more 447 widely distributed across the axes of both SOF-N and SOF-S. 448

Simulated U across a variety of incident wave approach angles are shown in Fig. 449 11. When waves approach the shoreline at highly oblique angles, a strong longshore cur-450 rent develops with minor offshore movement (Fig. 11a-b, f-g). The longshore currents 451 accelerate at the apex of SOF-N, near y = 1900 m, and SOF-S, near y = 500 m, and the 452 flow also meanders slightly offshore in these zones. With moderate incident θ_p , several 453 offshore circulation cells develop. During moderately positive incident angles, offshore 454 flows route down the northern edges of the SOFs (Fig. 11b, c) while during moderately 455 negative incident angles, they flow down the southern edges of the SOFs (Fig. 11e, f). 456 During periods of negative wave incidents (Fig. 11e-g), a localized zone of flow acceler-457 ation away from the shoreline develops at the location where the shoreline angle changes 458 orientation, between y = 1400 - 1600 m. 459

Across all simulations, weak offshore flows are consistently present along either the northern or southern edges of SOF-N and SOF-S and these offshore currents exit the computational domain, 1 km away from the shoreline (Fig. 11). However, at moderate incident angles, the offshore flows strengthen, coalesce, and develop complex pathways. Near the shoreline, a series of alongshore feeder currents converge and route offshore across a wide zone around SOF-N (Fig. 11c-e). Beginning at an incident angle of $+5^{\circ}$, four distinct zones of offshore-routed flows develop at y = 250 m, y = 1100 m, y = 1600 m, and y = 2500 m (Fig. 11c). These flows strengthen as incident angles shift towards -5° and begin to converge into two dominant offshore routing zones (near y = 750 m and 1800 m) which are continually present as wave angles continue to become more oblique at - 10° , identifiable at y = 450 m and y = 1750 m (Fig. 11f).

471 6 Discussion

472 6.1 Model error assessment

492

Several past studies have evaluated a similar model configuration at the FRF and 473 found that model results of H_s and energy-density spectra compare well with in-situ sen-474 sors located across the nearshore region (Gomes et al., 2016) as well as spatially-continuous 475 XBR estimates of wave angles and wave breaking patterns (Szczyrba et al., 2023b). In the absence of in-situ sensors situated within the study site, the accuracy of the J, M1, 477 and M2 simulations is assessed through a comparison to the XBR observations of surf 478 zone width, celerity, and wave height. The alongshore variability in surf zone width mea-479 sured by XBR is well represented by patterns of Q_b simulated in the model (Fig. 5b). 480 Mean celerity errors are low, ranging from -0.04 to 0.26 m/s. Errors are highest shore-481 ward of x = 400 m, where high rates of bathymetric change occur (Fig. 6). 482

Estimates of H_p are also in good agreement and both depict similar alongshore vari-483 ability in response to the complex bathymetric configuration (Fig. 7). The best avail-484 able bathymetric data were surveyed in June of 2017. Subsequent changes in the nearshore 485 morphology, particularly in the surf and swash zones, would lead to differences in wave 486 parameters estimated from XBR and SWASH and contribute to the identified errors. Past 487 studies have indicated that the SOFs in the region remain in fixed locations even after 488 energetic wave events (McNinch, 2004; Browder & McNinch, 2006; Schupp et al., 2006). 489 therefore it is expected that high errors due to alterations in the bathymetry would be 490 confined to zones closest to the shoreline. 491

6.2 Identification of rip currents

Anomalously high XBR backscatter signals seaward of the surf zone can occur where 493 incoming waves interact with surface current convergences, such as rip currents and other 494 offshore flows (Takewaka & Yamakawa, 2011; Haller et al., 2014; O'Dea et al., 2021). Backscat-495 ter anomalies indicative of strong offshore currents are visible in nearly every collection 496 hour with incident $H_s \geq 2$ m (Fig. 4d). These offshore flows emerge in two primary lo-497 cations, between y = 200 - 400 m and between y = 1500 - 2000 m. Flow structures re-498 main visible for several hours across a variety of incident angles, water levels, wave pe-499 riods, and directional spreading conditions. The visible plumes also extend 1-3 surf-zone 500 widths beyond the XBR-identified surf zone edge (Fig. 4), supporting the findings of Kumar 501 et al. (2021) and O'Dea et al. (2021). 502

Because the observed currents re-emerge in similar positions, they are likely mor-503 phologically controlled (Short, 2007). These flow structures are more obscured when wind 504 levels in the region reach above 10 m/s (Fig. 4c) because XBR images become saturated 505 during periods of high water surface roughness (Haller et al., 2014). Another limitation 506 of XBR is the requisite for depth-limited wave breaking. Small waves below ~ 1 m are 507 not large enough to be resolved by XBR, although this threshold is site-specific (McNinch, 508 2007). However, rip currents are also more likely to be generated as wave energy increases 509 (MacMahan et al., 2005), thus XBR remains a useful tool to identify them. 510

The observed reoccurring offshore flows coincide with the locations of SOF-N and SOF-S and are also represented in simulations of J, M1, and M2 (Fig. 8d and Fig. 9jl). Because the maximum current velocities outside of the surf zone at x = 600 m reach values between 0.54 - 0.90 m/s during these events, these flows would pose a risk to human safety and thus can be considered morphologically-controlled rip currents (Leatherman & Fletemeyer, 2011). Similar re-emergent flows are also represented in the idealized simulations with H_s values of 1 m, 3 m, and 5 m (Fig. 11). Even during low energy conditions ($H_s = 1$ m), narrow bands of weak offshore flow appear in a similar location shoreward of SOF-S (flows originating near y = 500 m) and SOF-N (flows originating near y = 1500 m).

During higher energy conditions $(H_s = 3 \text{ m and } 5 \text{ m})$, the origin of the offshore flow 521 within the surf zone is more variable, however the flows preferentially route along either 522 523 the northern or southern edges of SOF-N and SOF-S. When waves approach at a positive angle, depth-averaged currents are directed onshore along the southern edge (up-524 stream side of the SOF) and are offshore-directed along the northern edge (downstream 525 side). The opposite circulation pattern emerges when waves approach at a negative an-526 gle, with offshore flows directed along the southern edge. This patterns is similar to ob-527 servations of a dredged rip channel by Moulton et al. (2017), wherein a commonly ob-528 served circulation pattern included meandering longshore currents with offshore-directed 529 flows occurring on the downstream rip channel wall. 530

Other studies have linked complex bathymetry to rip current generation. Along the 531 northern shore of Prince Edward Island, Canada, Wernette and Houser (2022) proposed 532 a relationship between paleo-river channels identified by ground-penetrating radar and 533 rip currents observed by aerial imagery. At the present study site, results indicate that 534 the SOFs act as a pseudo-transverse bar and rip morphology, which is associated with 535 a common class of rip currents (Holman et al., 2006; Short, 2007; Turner et al., 2007). 536 However, the SOFs remain fixed in place with little change, as opposed to ephemeral sand-537 bar features that migrate shoreward after a storm reset (Houser et al., 2020). This pro-538 cess is akin to the shoreface-connected ridges on Fire Island, New York that have also 539 been shown to modify wave refraction and produce alongshore variable wave breaking 540 patterns (Safak et al., 2017). 541

542

6.3 Rip current generation

Rip currents are driven by gradients in wave momentum (Longuet-Higgins & Stew-543 art, 1964) resulting from alongshore variations in wave breaking (Lippmann & Holman, 544 1989). These variations can be caused by hydrodynamic forcing, such as intersecting wave 545 trains, standing edge waves, and wave spreading (Bowen, 1969; Suanda & Feddersen, 2015; 546 Kumar & Feddersen, 2017; Moulton et al., 2023), or artificial structures, such as piers 547 and groynes (Pattiaratchi et al., 2009). However, they are more often induced by along-548 shore variations in surf zone morphology (Bowen, 1969; Dalrymple, 1978). While the SOFs 549 at this site are not morphologically pronounced within the surf zone, they are persistent 550 nearshore structures that visibly impact the configuration of the bathymetry in areas of 551 depth-limited breaking during higher energy wave events (e.g., the 4 m isobath, Fig. 1). 552 Thus, the SOFs induce substantial alongshore variability in wave setup (Fig. 8c). 553

Nearshore wave heights respond to changes in nearshore bathymetry and the present 554 results indicate that larger waves reach closer to the shoreline along the axis of the SOFs 555 (Fig. 7), supporting the conclusions of Mulligan et al. (2019a) and observations by Sonu 556 (1972) at other sites. Waves refract across the obliquely variable bathymetry (Fig. 9a-557 c), creating zones of wave energy convergence and alongshore variations in H_s (Fig. 8a, 558 b). As a result, alongshore variations in wave breaking and surf zone widths also occur 559 (Fig. 5), which can drive rip currents (Bowen, 1969; Dalrymple, 1978). Higher setup is 560 generated near the shoreward side edges of the SOFs due to variability in wave break-561 ing (Fig. 8c), leading to a pressure gradient that forces water to flow away from regions 562 of intense breaking towards zones of lower setup. When these alongshore flows converge, 563 they route offshore (Fig. 9), interact with the incoming wave field, and cause localized 564 zones of higher frequency waves (Fig. 8d) and rougher water (Fig. 4b). These results 565

support the findings of experimental rip channel studies (Haller et al., 2002; Moulton et al., 2017).

568

6.4 Influence of wave climate

The mean incident wave conditions near KH throughout 2017 included a signifi-569 cant wave height of 0.87 m, mean incident angle of 17°, and 32° of directional spread-570 ing. However, Nor'easters and tropical cyclones generate a variety of high-energy inci-571 dent wave conditions in this region. When waves approach the SOFs at angles close to 572 the orientation of the axis (negative incident angles, in our coordinate system), wave break-573 ing occurs in concentrated local patches (Fig. 10). When approaching from the oppo-574 site direction (positive incident wave angles), wave breaking is more widely distributed 575 along the 4 m isobath. These findings support Safak et al. (2017), who concluded that 576 when waves approach shoreface-connected ridges at angles similar to the angle of the ridge 577 crests, the features focused the most wave energy. Strong offshore flows (i.e., rip currents) 578 are more likely to occur when wave angles approach the site between -10° to 10° (Fig. 579 11). Rip currents are more prevalent during conditions of moderate incident wave an-580 gles (Engle, 2002; MacMahan et al., 2005; Dusek & Seim, 2013a) because highly oblique 581 angles induce strong longshore currents that inhibit offshore flow (Kumar et al., 2011). 582 During positive, oblique incident conditions (15°) , the longshore current accelerates where 583 the contours of the SOFs pinch towards the shoreline (Fig. 11a) whereas during nega-584 tive, oblique conditions (-15°) , this acceleration occurs south of the SOFs. 585

Numerous weak offshore flows occur during periods of low energy, however these 586 offshore flows re-emerge in similar positions across a variety of incident wave angles. As 587 H_s increases, these offshore flows coalesce into larger, stronger offshore flows resembling 588 rip currents (Fig. 11 that preferentially route down either the northern or southern edges 589 of the SOFs when incident wave angles enter the domain between -10° to 10° . Fig. 12 590 explores the idealized effect of incident H_s and θ on maximum U (Fig. 12a, c-e) and vari-591 ability of U (Fig. 12b, f-h) at the mean XBR-measured surf zone location as well as two 592 zones located 250 m and 500 m offshore. Within the surf zone, maximum U values are 593 mostly found at the shoreward axis of SOF-N, located at approximately y = 1650 -594 2000m (Fig. 12a). Additionally, U generally increases with increasingly oblique waves, 595 owing to the generation of a strong longshore current (Fig. 12c). The variability of U596 is higher for negatively incident waves, but is highest for waves of moderate incident an-597 gles (Fig. 12f). Both mean U and the variability of U decrease offshore and increase with 598 larger H_s . At all cross-shore locations, the variability of U decreases with increasingly 599 oblique incident waves (Fig. 12f-h). The locations of maximum U variability are con-600 centrated along the southern edges of both SOF-N and SOF-S (Fig. 12b). 601

The idealized simulations do not represent realistic field conditions, but are intended 602 to isolate the response to different incident wave conditions. This set of simulations in-603 cluded JONSWAP spectra with zero directional spreading. Several studies have suggested 604 that rip currents are more pervasive during narrow-banded incident conditions (Dusek 605 et al., 2011; Dusek & Seim, 2013b), therefore the idealized results can not be used to pre-606 dict rip current activity. However, XBR images indicate that rip currents at this site are 607 indeed generated throughout a range of incident angles and directional spreading val-608 ues (Fig. 4). The purpose of the idealized results is to explore the sensitivity and response 609 of the SOFs to isolated variables (i.e. θ , H_s). It is also well known that rip current in-610 tensity varies in response to the total water level (e.g., tidal stage) and that strong rip 611 currents can occur at low tide even during periods of low energy conditions due to en-612 hanced wave breaking over surf zone sandbars (Brander & Short, 2001; MacMahan et 613 al., 2005; Voulgaris et al., 2011). Overall, bathymetrically controlled rip currents tend 614 to be strongest and most prevalent during periods of low water levels, high H_s , moder-615 ate incident *theta*, and low spreading (Haller et al., 2002; MacMahan et al., 2010; Moul-616 ton et al., 2017). 617

6.5 Implications for hot spots of erosion

Prior studies have demonstrated that, at the regional scale (i.e., ~ 40 km along-619 shore length). SOFs in the Outer Banks correlate significantly with erosional hot spots 620 (McNinch, 2004; Schupp et al., 2006). The physical mechanisms contributing to the causal 621 relationship have previously not been determined. The results of this research suggest 622 that strong cross-shore fluxes (i.e., rip currents) are concentrated near the SOFs during 623 periods with moderate incident wave angles. Between 2012 and 2022, 72 beach rescues 624 specifically related to rip currents occurred within the study site (Kitty Hawk Ocean Res-625 cue, personal communication, September 28, 2023). Rip currents and surf zone eddies 626 are the primary mechanisms for cross-shore sediment transport (MacMahan et al., 2006; 627 Dalrymple et al., 2011; Castelle et al., 2016) and Splinter and Palmsten (2012) found that 628 areas of higher dune erosion occur where rip currents are present directly offshore. Dur-629 ing periods of both high wave energy and moderate wave angles, strong offshore-directed 630 flows may pronounce offshore transport, leading to enhanced erosion and undulations 631 in the shape of the shoreline, similar to the results found in Fire Island, New York by 632 Safak et al. (2017). During conditions with more oblique incident wave angles, longshore 633 currents accelerate near the SOFs due to alongshore variability in bathymetry, which would 634 also contribute to higher rates of sediment transport. Therefore, across a range of inci-635 dent wave conditions and energy levels, a combination of longshore current accelerations 636 and far-reaching offshore flows could contribute to the severity of beach erosion near the 637 SOFs at this site. 638

Bathymetric surveys indicate that between 2004 and 2017 the trough of the 4 m 639 contour within SOF-N and SOF-S deepened and moved south by 300 - 400 m (Szczyrba 640 et al., 2023a). This might explain the southerly offset between the origin of the morpho-641 logically controlled rip currents and the long-term shoreline change rates (Fig. 4). We 642 hypothesize that bathymetric hysteresis causes the erosional hot spot to lag behind the 643 southerly movement of the SOFs and, therefore, in the future the hot spot will migrate 644 south. When comparing the long-term shoreline change rate data released in 2004, 2013, 645 and 2020, there is some evidence that the hot spot is intensifying and expanding south-646 wards as the SOF troughs have deepened and moved south, although these data are cu-647 mulative long-term averages that do not highlight year to year changes. Nevertheless, 648 according to these data, between 2004 and 2020 the rate of erosion at SOF-N increased 649 from -0.36 m/yr to -0.70 m/yr and the SOF-S from -0.79 m/yr to -1.00 m/yr. The ar-650 eas affected have also expanded south by 100 m (North Carolina Division Of Coastal Man-651 agement, 2021). 652

However, this hypothesis is complicated by the beach nourishment projects that 653 occur in this region nearly every five years. A monitoring report conducted one year af-654 ter the conclusion of the 2017 beach nourishment project in KH observed that that up 655 to 60% of the sand volume loss occurred adjacent to the SOFs, although the authors cited 656 challenges in calculating sediment volumes around the complex bathymetric structures 657 (APTIM, 2019). The report concluded that future monitoring should investigate the im-658 pact of these SOFs on nourishment performance. The results of this study suggest that 659 the SOFs act as a conduit for enhanced cross-shore exchange of sediment because, dur-660 ing periods of moderate wave heights and angles, enhanced flow velocities are directed 661 offshore down the SOF edges and persist beyond the surf zone edge. During periods with 662 strong longshore currents, the flow accelerates and meanders near the SOFs, which would 663 further enhance sediment transport away from the SOF and contributing to the erosional 664 hot spot. These findings emphasize the importance of both cross- and longshore trans-665 port on nearshore morphological evolution, supporting the conclusions of Thieler et al. 666 (1995) and Gutierrez et al. (2005) at other sites. 667

The storm-induced wave forcing addressed in this research also affects shoreline change on an inter-annual to decadal time scale (Splinter et al., 2014). As climate change alters ocean temperatures, the intensity and duration of the Atlantic hurricane season may increase (Knutson et al., 2010; Walsh et al., 2016). This elevates the likelihood of sequential high-energy wave events (i.e., storm clusters) that inhibit beach recovery (Coco et
al., 2014) and result in higher rates of shoreline recession (Dodet et al., 2019). Geologically inherited complex bathymetry produces variability in beach response and recovery, whereby influencing the pattern of barrier island transgression over time (Houser,
2012). With a more intense wave climate, complex nearshore bathymetric features will
continue to drive spatially variable circulation patterns and affect erosional hot spots.

778 7 Summary and Conclusions

The influence of complex bathymetry on nearshore flow dynamics in an area with 679 erosional hot spots was explored with a combination of remotely-sensed data and a nu-680 merical model. An X-Band Radar (XBR) system was deployed at the study site in Kitty 681 Hawk, North Carolina, zfor nine days at the end of September, 2017. This period coin-682 cided with the passage of Hurricane Jose and Hurricane Maria, and these events were 683 also simulated with a phase-resolving numerical model. In total, 24 simulations were per-684 formed: 3 selected times during storm events (J, M1, M2) that were also observed by 685 XBR and 21 idealized simulations that explored a range of incident wave conditions. The 686 event simulations were validated against XBR observations of wave breaking patterns, 687 celerity, and representative wave height estimates. 688

The XBR observations were used to identify two zones of persistently high sea-surface 689 backscatter during energetic periods $(H_s \ge 2 \text{ m})$, indicative of strong offshore flows. Across 690 a wide variety of incident wave conditions, the offshore flows re-emerged near shore oblique 691 features (SOFs) and are concluded to be bathymetrically controlled rip currents. The 692 wave-current interactions associated with these offshore flows roughens the sea surface. 693 enabling them to be easily identified with XBR, and also increases the wave frequency 694 in localized zones. The location of the offshore currents are also linked to the incident 695 wave direction and the bathymetry. When waves approach the shoreline at a positive 696 angle, the offshore currents flow along the northern SOF edges while when wave approach 697 at a negative angle, the flows route along the southern edges. The rip currents result from 698 wave breaking across the alongshore-varying bathymetry near the SOFs, supported by 699 both XBR and numerical modeling data. Numerical simulations indicate that relative 700 wave heights and wave setup also vary across the cross-section of the SOFs. When the 701 wave field is oblique, a strong alongshore current represses offshore flows, and this cur-702 rent accelerates and meanders where the SOF contours pinch towards the shoreline. 703

Idealized simulations indicate that the complex bathymetry exerts a strong con-704 trol over nearshore wave heights, refraction patterns, and mean currents. The XBR ob-705 servations were used to identify morphologically-driven persistent offshore-directed flows 706 just south of long-term erosional hot spots. Future studies, using other numerical mod-707 els that simulate sediment movement and bed elevation change, could directly model the 708 morphodynamic evolution of zones with complex bathymetry to evaluate impacts on ero-709 sional hot spot intensity and evolution. Persistent offshore flows and alongshore flow ac-710 celeration may exacerbate erosion and enhance the shoreline undulations at this site. This 711 underscores the importance of both cross-shore and alongshore sediment transport on 712 the nearshore morphology and shoreline evolution on sandy beaches with nearshore bathy-713 metric features. 714

715 8 Data Availability Statement

All model inputs, radar data, processed model output data, and figure generation
 MATLAB codes used in this study are hosted on the Borealis data repository at Queen's
 University titled "Nearshore Flow Dynamics – SWASH and XBand Radar", via DOI:
 https://borealisdata.ca/privateurl.xhtml?token=fa6593ed-6ea6-437c-8734-4711be8950c8

- ⁷²⁰ with data use license agreement CC-BY 4.0. This is a private link for reviewers only, and
- the public link will be included upon manuscript acceptance.

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- versity, and Dr. Spicer Bak as well as Dr. Kate Brodie from the U.S. Army Engineer Re search and Development Center.
- 732 Acronyms
- 733 **AWAC** Acoustic Wave and Currents profiler
- ⁷³⁴ **FRF** Field Research Facility
- 735 **KH** Kitty Hawk
- 736 **RIOS** Radar Inlet Observation System
- ⁷³⁷ **R.I.** Return intensity of X-band radar
- 738 **RMSE** Root mean square error
- ⁷³⁹ **SOF** Shore-oblique feature
- 740 **SOF-N** Northern shore-oblique feature
- ⁷⁴¹ **SOF-S** Southern shore-oblique feature
- 742 **SWASH** Simulating WAves till SHore
- 743 **TKE** Turbulent kinetic energy
- 744 XBR X-band radar

745 Notation

- α_{ad} Calibration coefficient, Eq. 12 from Streßer et al. (2022)
- C_{p} Celerity of peak frequency waves [m/s]
- $_{748}$ **D**₅₀ Median grain size [mm]
- η_{49} η Water surface elevation [m]
- f Wave frequency [Hz]
- 751 γ Breaking index (0.78)
- $_{752}$ **g** Gravitational acceleration $[m/s^2]$
- $_{753}$ **h** Water depth [m]
- H_p Height of peak frequency waves [m]
- 755 H_s Significant wave height [m]
- 756 k Wave number [1/m]
- 757 λ Wavelength [m]
- ⁷⁵⁸ $\boldsymbol{\theta}_{m}$ Mean wave angle [°]
- $au_{59} au$ Shear $[\mathrm{N}/m^2]$
- 760 T_p Peak period [s]
- $_{761}$ U Depth-averaged mean velocity [m/s]
- $_{^{762}}$ U_{*cr} Critical velocity [m/s]
- 763 $oldsymbol{Q}_{b}$ Breaking parameter

Table 1. Wave conditions during the hindcast simulated event (J, M1, M2) including significant wave height (H_s) , peak period (T_p) , total water level (η) , mean wave angle relative to perpendicular of the study site's mean shoreline (θ_m) , and peak directional spreading (σ_p) .

Label	UTC	H_s [m]	T_p [s]	$\eta~[{\rm m}]$	θ_m [°]	σ_p [°]
J	19-Sep-2017 20:00:00	2.8	10.0	0.9	3.9	36.3
M1	27-Sep-2017 09:00:00	4.0	13.3	0.7	16.9	23.2
M2	28-Sep-2017 02:00:00	2.6	10.3	0.8	-37.1	58.5

Measured at the FRF's 11 m AWAC.



Figure 1. (a) Study region location within the Outer Banks, (b) June 2017 bathymetry measured offshore of Kitty Hawk, North Carolina where the red star represents the XBR location and labels N and S identify the locations of SOF-N and SOF-S, respectively, and (c) rotated bathymetric grid used for numerical modeling.



Figure 2. Wave conditions throughout the study period: (a) significant wave height (H_s) , (b) peak period (T_p) , and (c) mean wave angle relative to perpendicular of the mean shoreline (θ_m) . Wave angles are described as going towards, meaning positive values indicate an approach angle north of shore-perpendicular and negative values indicate an approach angle south of shore-perpendicular. Vertical dashed lines indicate the simulated hours during Hurricanes Jose and Maria (J, M1, M2) and shaded regions are periods observed by XBR.



Figure 3. Wave spectra for time periods (left column) J, (middle column) M1, and (right column) M2: (a, b, c) energy-density, (d, e, f) directional distribution of wave angles, and (g, h, i) directional energy-density spectra used to force simulations at the offshore boundary. Red lines in (g, h, i) are the mean direction per frequency calculated from Kuik et al. (1988). Wave angles are described as going towards, meaning positive values indicate an approach angle north of shore-perpendicular and negative values indicate an approach angle south of shore-perpendicular.



Figure 4. Alongshore variability in: (a) the long-term average shoreline change rate over 1940-2016 with shaded regions indicating the troughs of SOF-N and SOF-s;(b) XBR backscatter intensity anomalies for J (blue line), M1 (black line), M2 (orange line), and the mean during the observation period (thick black line); (c) time-averaged XBR backscatter return intensity (R.I.) during event M2 with the white lines indicating the corresponding XBR estimated offshore edge of the surf zone; and (d) time-averaged XBR backscatter during a moderate energy event ($H_s = 2.3$ m) on September 20, 2017 at 04:13 UTC.



Figure 5. Surf zone detection from XBR observations and model: (a) XBR measurements of the cross-shore location of the offshore surf zone edge during event J (blue line), M1 (black line), and M2 (orange line) with shaded regions indicating the SOF troughs; (b) XBR-derived M1 surf zone edge (black line) over the simulated fraction of breaking waves (Q_b) for M1.



Figure 6. XBR observations and model results of wave celerity at the peak frequency (C_p) during M2: (a) XBR observations in areas deeper than 2 m, (b) SWASH results, and (c, d, e) cross-shore transects of modeling results (solid lines) and XBR observations (dashed lines) at alongshore locations (c) y = 2250 m, (d) y = 1500 m, and (e) y = 750 m. White dashed lines in (a) and (b) indicate the cross-shore transect locations.



Figure 7. Estimates of the peak wave height (H_p) for event M2 from (a) XBR observations and (b) modelling results. Solid black lines in (a) and (b) represent the XBR-measured surf zone edge of M2.



Figure 8. Wave height estimates from numerical simulations: (a) the spatial distribution of relative wave heights (H_s/h) during M1, (b) maximum H_s divided by mean H_s at each alongshore location for J (blue line), M1 (black line), and M2 (orange line), (c) maximum alongshore gradient in setup, and (d) saturated frequency map during M1. Gray shaded regions in (b) and (c) indicate the approximate bounds of the shore-oblique troughs.



Figure 9. Flow dynamics around the northern shore-oblique feature between y = 1500 m -2500 m for time periods (left column) J, (middle column) M1, and (right column) M2. Flow dynamics displayed include: (a, b, c) wave angle (θ) relative to shore-perpendicular, (d, e, f) cross-shore velocity (u), (g, h, i) alongshore velocity (v), and (j, k, l) overall depth averaged current velocity (U).



Figure 10. Fraction of breaking waves (Q_b) calculated from $H_s = 3$ m idealized simulations with various incident peak wave angles (θ_p) : (a) 15°, (b) 10°, (c) 5°, (d) 0°, (e) -5°, (f) -10°, and (g) -15°.



Figure 11. Depth-averaged current velocities (U) calculated from $H_s = 3$ m idealized simulations with various incident peak wave angles (θ_p): (a) 15°, (b) 10°, (c) 5°, (d) 0°, (e) -5°, (f) -10°, and (g) -15°.



Figure 12. Locations of maximum U (a) and locations of maximum variability of U (b) during a variety of incident wave angles and significant wave heights (H_s) , where circles represent an incoming H_s of 1 m, triangles represent H_s of 3 m, and stars represent H_s of 5 m. Quantified influence of incoming wave angle (θ) and H_s on maximum U (c-e) and the variability of U (f-h) in the surf, intermediate, and offshore zones. Cross-shore locations of the surf, intermediate, and offshore zones are depicted in (a) and (b) as white dashed lines.

764 **References**

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