

The effect of increased water vapor from the Hunga Tonga-Hunga Ha’apai eruption on the Antarctic ozone hole

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Abstract

The eruption of the Hunga Tonga-Hunga Ha’apai volcano on 15 January 2022 was one of the most explosive eruptions of the last decades. The unprecedented amount of water vapor injected into the stratosphere increased the stratospheric water vapor burden by about 10%. Using model runs from the ATLAS chemistry and transport model and Microwave Limb Sounder (MLS) satellite observations, we show that while 20-40% more water vapor than usual was entrained into the Antarctic polar vortex as it formed (e.g., typical values of 4.6 ppm at 21.5 km increased to 6.7 ppm), the direct effect of the increased water vapor on Antarctic ozone depletion was minor. This is caused by the very low temperatures in the vortex, which limit water vapor to the saturation pressure and tend to reset any anomalies in water vapor by dehydration before they can have an effect on ozone loss.

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Plain Language Summary

The eruption of the Hunga Tonga-Hunga Ha’apai volcano on 15 January 2022 was one of the most explosive eruptions of the last decades. An unprecedented amount of water vapor was injected into the stratosphere, increasing the total stratospheric water vapor mass by about 10%. Using model runs and satellite observations, we show that while the dissipation of the plume increased water vapor in the Antarctic in 2023 by 20–40% at the beginning of the ozone hole season compared to earlier years, the effect of the increased water vapor on the Antarctic ozone hole was minor. This is caused by the very low temperatures in the vortex, which limit water vapor due to condensation and tend to reset any anomalies in water vapor before they can have an effect on ozone loss and sedimentation of condensed particles.

1 Introduction

The eruption of the Hunga Tonga-Hunga Ha’apai volcano on 15 January 2022 was one of the most violent eruptions of the last decades. It reached a volcanic explosivity index (VEI) of 5, and the plume reached an altitude of more than 50 km (e.g. Carr et al., 2022; Millán et al., 2022; Proud et al., 2022; Schoeberl et al., 2022). An unprecedented amount of water vapor was injected into the stratosphere, increasing the stratospheric water vapor burden by about 10% or 150 Tg (Millán et al., 2022; Schoeberl et al., 2022; Vömel et al., 2022). While the increased water vapor was not able to penetrate into the 2022 Antarctic vortex (Manney et al., 2023), the dissipating plume increased water vapor observed by MLS in the developing Antarctic vortex in 2023 by about 20% to 40% compared to earlier years (in the pressure range of 56.2–17.7 hPa or approximate altitude range 20–28 km). Figure 1 (a) shows vortex-averaged profiles of MLS measurements of water vapor (version 5, (Livesey et al., 2022)) in the developing Antarctic vortex for all years of the MLS record (2005–2023) on 20 May. Here and in the following, the vortex is defined as the volume inside the -36 PVU contour of modified potential vorticity (which scales potential vorticity to have a similar range of values throughout the stratosphere) calculated from a reference level of $\theta_0 = 475$ K (Lait, 1994). Typical profiles at the end of May in the years before 2023 are very similar. In comparison, the profile of 2023 shows increased values throughout a large vertical range. For example, values at 21.5 km increased from an average 4.6 ppm in earlier years to 6.7 ppm in 2023.

It was speculated that the increased water vapor from the eruption could lead to increased ozone depletion in the Antarctic ozone hole in 2023 (e.g. Millán et al., 2022; Manney et al., 2023). Water vapor can influence polar ozone depletion mainly by its effect on polar stratospheric clouds (PSCs), e.g. by changes in formation temperature thresholds, particle size distribution, or dehydration and denitrification. However, MLS measurements in 2023 show that ozone values are not exceptional and are well within the range of earlier years. This can be seen in Figure 1 (h), which shows vortex-averaged pro-

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76 We perform runs of the ATLAS chemistry and transport model (Wohlmann & Rex,
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85 reanalysis (provided on a $1.125^\circ \times 1.125^\circ$ horizontal grid, 3 h temporal resolution, and
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87 The model was run from 1 April 2023 to 1 October 2023. Chemical species are initial-
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96 The value was empirically adjusted to fit the water vapor measurements in earlier stud-
97 ies.

98 3 Results

99 Figure 2 (a) shows the time evolution of vortex-averaged water vapor in the ref-
100 erence run, the sensitivity run, and in the MLS data at 475 K potential temperature. The
101 reference run and MLS show excellent agreement over the complete time period. Wa-
102 ter vapor starts with values of about 6–7 ppm in May, but quickly decreases to values
103 of about 3–4 ppm between mid-May and the start of July. This decrease is caused by
104 dehydration, which limits water vapor to the saturation pressure through condensation
105 and sedimentation in the very cold polar vortex. The decrease is also visible in the MLS
106 profiles in Figures 1 (c, e, g), except at the highest altitudes where water vapor does not
107 exceed the saturation limit. Figures 1 (c, e, g) also show that water vapor is well within
108 the range of previous years by the beginning of July.

109 Figure 2 (a) shows that the sensitivity run starts at values below 5 ppm that do
110 not agree well with the measurements (as expected). However, the sensitivity run quickly
111 converges to the reference run by early July. This is because in both cases the temper-
112 atures are sufficiently low that water vapor abundances are mostly above the saturation
113 limit. Once the saturation limit is reached, water vapor in both runs is equalized.

114 Figure 2 (b) shows the corresponding time evolution of vortex-averaged ozone at
115 475 K. The difference between the reference run and the sensitivity run is very small through-
116 out the whole time period. This is because heterogeneous ozone loss is only significant

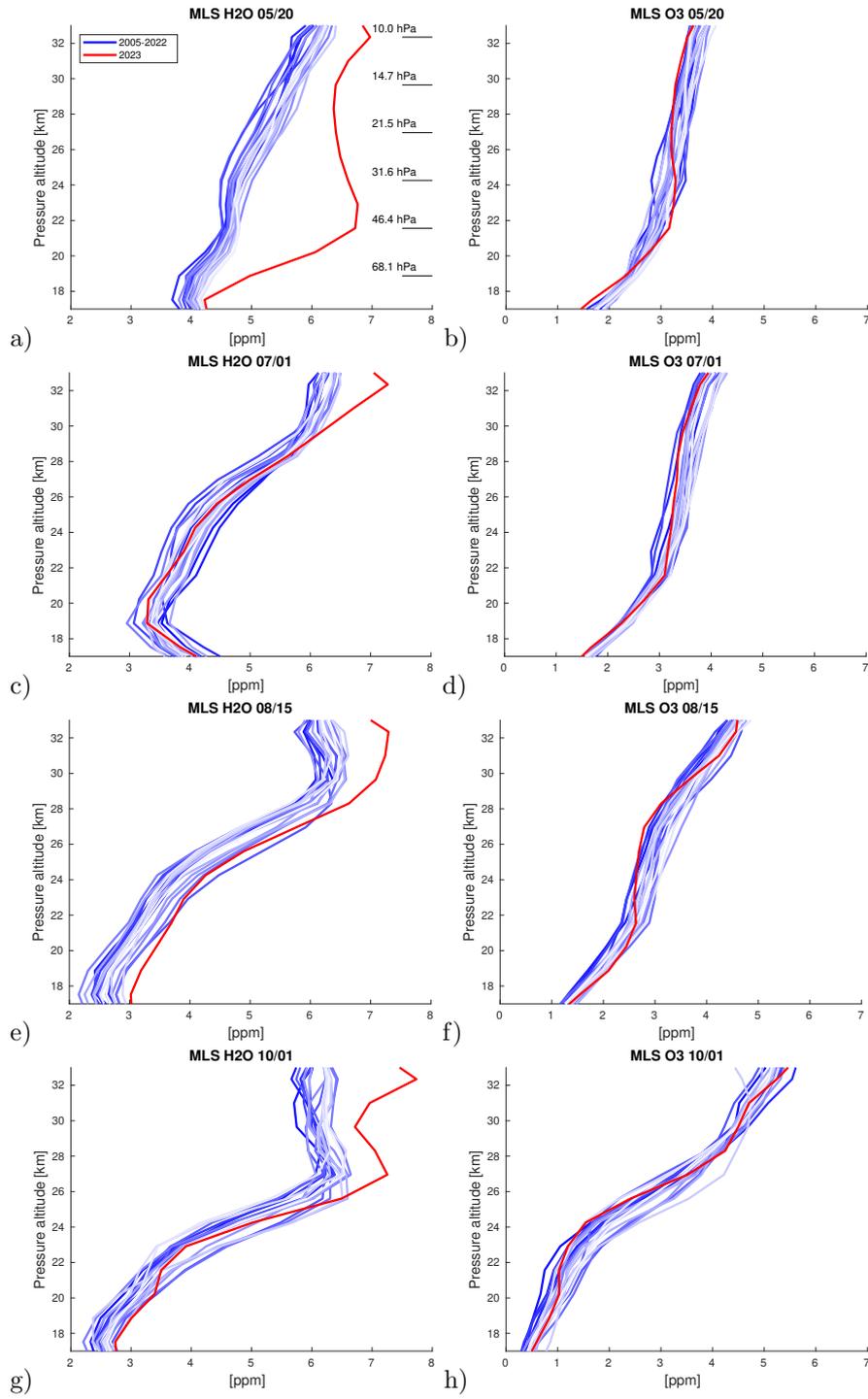


Figure 1. Left: Vortex-averaged water vapor profiles observed by MLS on 20 May (a), 1 July (c), 15 August (e) and 1 October (g) for all years of the MLS data record. 2005–2022 in shades of blue, 2023 is highlighted in red. Right: Same for ozone (b, d, f, h). Profiles are averaged over all MLS measurements of the given day inside the -36 PVU contour of modified potential vorticity. Every other MLS pressure level is indicated in panel (a).

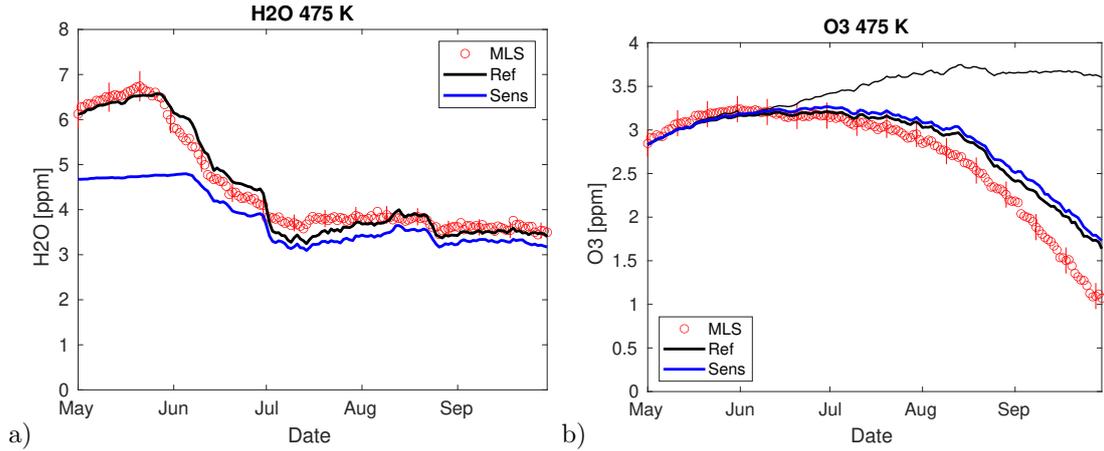


Figure 2. Left (a): Vortex-averaged water vapor at 475 K potential temperature observed by MLS in 2023 (red dots) and modelled by ATLAS (blue and black lines). The black line shows the reference run initialized with MLS measurements from 2023, while the blue line shows the sensitivity run initialized with MLS water vapor data from 2022 (i.e., without the effect of Hunga Tonga on water vapor). Right (b): Same for ozone. The thin black line shows a passive ozone tracer initialized on 1 June. The difference between the thin black line and the other lines quantifies the amount of ozone loss.

117 within the sunlit portion of the vortex from July through September. The thin black line
 118 shows a passive ozone tracer initialized on 1 June. The difference between the passive
 119 ozone tracer and the other lines quantifies the amount of ozone loss. There is only a short
 120 time period before July (cf. Figure 2 (a)) when differences in water vapor between the
 121 runs can have a direct effect on differences in ozone loss.

122 The difference between the modelled ozone and the passive ozone tracer underestimates
 123 ozone loss by about 0.7 ppm (30 %) compared to the difference between the pas-
 124 sive ozone tracer and MLS ozone. Since ATLAS runs for other Antarctic winters agree
 125 better with MLS (e.g., Wohltmann et al., 2021), this discrepancy might hint that other
 126 effects from the Hunga Tonga eruption are at work. However, as outlined before, it can-
 127 not be the direct effect of water vapor.

128 The results at other potential temperature levels lead to similar conclusions, and
 129 the effect of the increased water vapor remains small throughout the ozone column. Fig-
 130 ure 3 shows the chemical ozone loss modelled by ATLAS for the partial column from the
 131 lower model boundary at 157 hPa to 28.6 hPa. Ozone loss was determined by subtract-
 132 ing the passive ozone tracer initialized on 1 June from the modelled ozone values. Since
 133 the value of the passive ozone tracer is not known for air masses that entered through
 134 the upper model boundary after 1 June and descended in the vortex, the column is re-
 135 stricted to 28.6 hPa. The figure shows that the effect of the increased water vapor on
 136 the column loss in the model is small.

137 4 Discussion and Summary

138 There are several ways in which water vapor might change ozone loss in addition
 139 to dehydration:

- 140 • The threshold temperatures for the formation of all PSC types (supercooled ternary
 141 solution (STS), nitric acid trihydrate (NAT) and ice) are increased by the increased

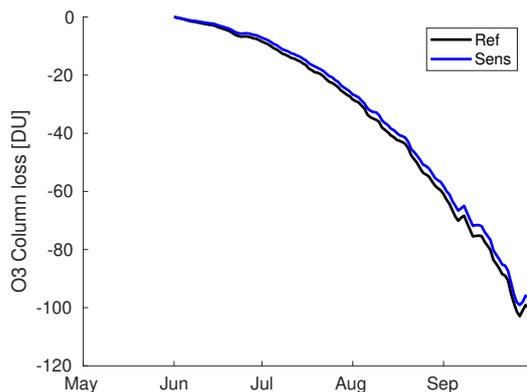


Figure 3. Vortex-averaged chemical loss of ozone modelled by ATLAS for the partial column from the lower model boundary at 157 hPa to 28.6 hPa. The black line shows the reference run initialized with MLS measurements from 2023, while the blue line shows the sensitivity run initialized with MLS water vapor data from 2022 (i.e., without the effect of Hunga Tonga on water vapor).

142 water vapor. For example, a change in water vapor from 5 ppm to 6 ppm increases
 143 the threshold temperature for ice by 1.1 K at 50 hPa (Marti & Mauersberger, 1993),
 144 and for NAT by 0.8 K (for a HNO_3 mixing ratio of 10 ppb) (D. Hanson & Mauers-
 145 berger, 1988). STS has no defined threshold temperature for formation but shows
 146 a gradual increase of reaction rates and droplet volume with lower temperatures.
 147 However, an increase in water vapor from 5 ppm to 6 ppm has to a good approx-
 148 imation the same effect on the STS reactivity (time needed for the $\text{HCl} + \text{ClONO}_2$
 149 reaction to completely deplete one of the reaction partners) as a temperature change
 150 of 1 K (D. R. Hanson & Ravishankara, 1994; Carslaw et al., 1995) (for 2 ppb HCl,
 151 1 ppb ClONO_2 , 0.15 ppb H_2SO_4 , number density of droplets 10 cm^{-3}).

152 While these changes increase the volume inside the vortex where the formation
 153 of PSCs is possible, this can only have an effect before dehydration sets in. The
 154 effects of changes in threshold temperatures and STS reactivity are included in
 155 ATLAS. Figure 2 (b) shows that they have no large effect on ozone depletion.

156 • Changes in water vapor have an effect on the particle size distribution of all cloud
 157 types. There is more water vapor available above the saturation limit, and one might
 158 expect larger particle sizes and larger surface area densities. However, particle for-
 159 mation and growth is a complex process, and this might not be straightforward.
 160 As in many chemistry and transport models, PSCs are treated in a somewhat sim-
 161 plified manner in ATLAS (Tritscher et al., 2021). For NAT and ice clouds, a con-
 162 stant number density is assumed, and a uniform particle size is then calculated
 163 from the HNO_3 and H_2O available above the saturation pressure. That is, the uni-
 164 form particle size will simply increase in our model. For STS, a constant number
 165 density and a log-normal distribution is assumed that is scaled with the total liq-
 166 uid volume.

167 Figure 2 (b) shows that changes in size distribution have only a small effect on
 168 ozone depletion in ATLAS. However, particle growth and formation might be more
 169 complex in reality and the effect on ozone depletion may be larger. Changes in
 170 size distribution may also have an effect after dehydration sets in (i.e., differences
 171 in the size distribution of NAT, ice and STS PSCs caused by increased water va-
 172 por in May and June might be persistent or lead to changes in the size distribu-
 173 tion in later months).

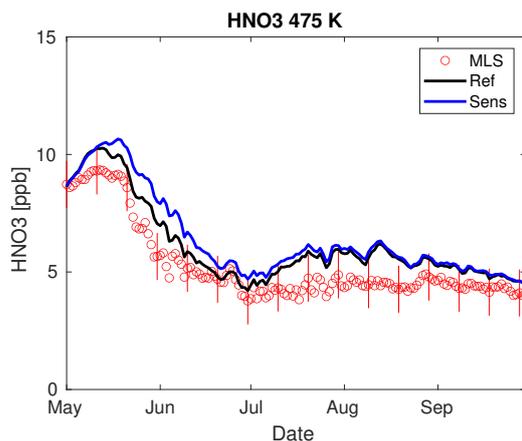


Figure 4. Vortex-averaged HNO_3 (gas phase) at 475 K potential temperature observed by MLS in 2023 (red dots) and modelled by ATLAS (blue and black lines). The black line shows the reference run initialized with MLS measurements from 2023, while the blue line shows the sensitivity run initialized with MLS water vapor data from 2022 (i.e., without the effect of Hunga Tonga on water vapor).

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- Since larger particles have a greater fall velocity, dehydration might also be faster in 2023. However, this is not considered in the simple ATLAS dehydration parameterization. The good agreement between ATLAS and MLS suggests that dehydration might be a fast enough process so that the speed of the process does not matter for the amount of ozone depletion.
 - Denitrification might also be affected by differences in the formation of the NAT particles. Denitrification is treated in a more sophisticated manner than dehydration in ATLAS and incorporates the nucleation, growth, sedimentation, and evaporation of individual particles with the DLAPSE model (Davies et al., 2005). Figure 4 shows that the changes in gas-phase HNO_3 (which are mainly caused by denitrification) are small between the reference run and the sensitivity run.

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In summary, we have shown that while the Hunga Tonga-Hunga Ha’apai eruption increased water vapor in the emerging Southern Hemisphere stratospheric polar vortex by 20–40% in 2023, ozone values at the end of September in the Antarctic ozone hole were in no way exceptional and were well within in the range of earlier years. ATLAS model runs indicate that the direct effect of the increased water vapor on Antarctic ozone depletion was minor. The reason for this are the very low temperatures in the vortex, which limit water vapor to the saturation pressure and tend to reset any anomalies in water vapor through dehydration before they can have an effect on ozone loss. However, the ATLAS runs underestimate observed ozone loss in 2023 by about 30% (in contrast to good agreement with observations in similar studies for other winters). Further studies would be needed to understand the causes of this discrepancy and whether it may be related to other effects from the Hunga Tonga-Hunga Ha’apai eruption in addition to the increased water vapor and dehydration.

198 Open Research Section

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ATLAS source code is available from the repository at <https://gitlab.awi.de/iwohltmann/atlas-julia>. MLS data are available at <https://disc.gsfc.nasa.gov/datasets?page=1&keywords=AURA%20MLS>. ECMWF ERA5 data are available at Hersbach et al. (2017), doi:10.24381/cds.143582cf.

202 **Acknowledgments**

203 ECMWF ERA5 and ERA5T data (Hersbach et al., 2017) was downloaded from
 204 the Copernicus Climate Change Service (C3S) (2023). The results contain modified Coper-
 205 nicus Climate Change Service information 2023. Neither the European Commission nor
 206 ECMWF is responsible for any use that may be made of the Copernicus information or
 207 data it contains.

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100 erence run, the sensitivity run, and in the MLS data at 475 K potential temperature. The
101 reference run and MLS show excellent agreement over the complete time period. Wa-
102 ter vapor starts with values of about 6–7 ppm in May, but quickly decreases to values
103 of about 3–4 ppm between mid-May and the start of July. This decrease is caused by
104 dehydration, which limits water vapor to the saturation pressure through condensation
105 and sedimentation in the very cold polar vortex. The decrease is also visible in the MLS
106 profiles in Figures 1 (c, e, g), except at the highest altitudes where water vapor does not
107 exceed the saturation limit. Figures 1 (c, e, g) also show that water vapor is well within
108 the range of previous years by the beginning of July.

109 Figure 2 (a) shows that the sensitivity run starts at values below 5 ppm that do
110 not agree well with the measurements (as expected). However, the sensitivity run quickly
111 converges to the reference run by early July. This is because in both cases the temper-
112 atures are sufficiently low that water vapor abundances are mostly above the saturation
113 limit. Once the saturation limit is reached, water vapor in both runs is equalized.

114 Figure 2 (b) shows the corresponding time evolution of vortex-averaged ozone at
115 475 K. The difference between the reference run and the sensitivity run is very small through-
116 out the whole time period. This is because heterogeneous ozone loss is only significant

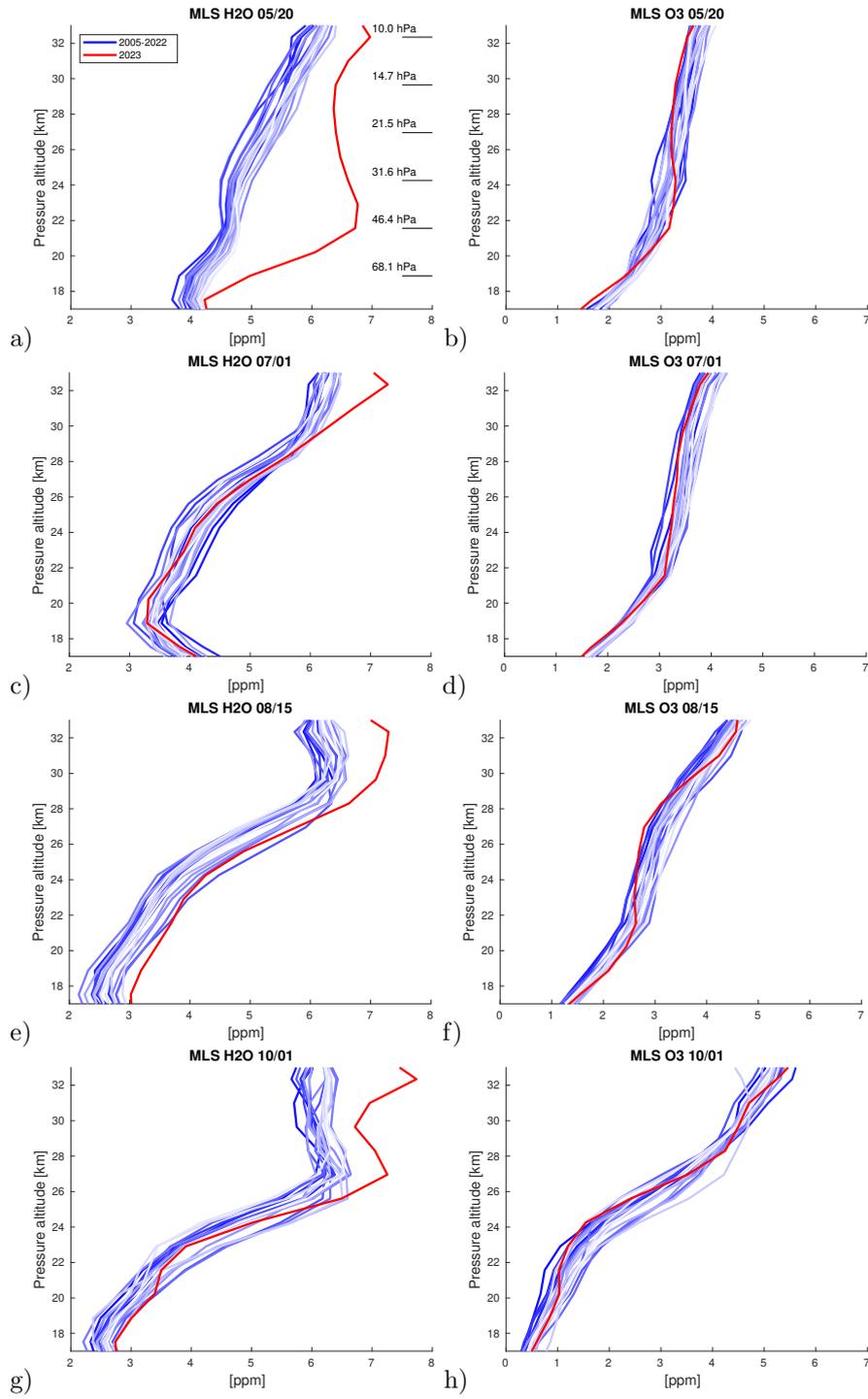


Figure 1. Left: Vortex-averaged water vapor profiles observed by MLS on 20 May (a), 1 July (c), 15 August (e) and 1 October (g) for all years of the MLS data record. 2005–2022 in shades of blue, 2023 is highlighted in red. Right: Same for ozone (b, d, f, h). Profiles are averaged over all MLS measurements of the given day inside the -36 PVU contour of modified potential vorticity. Every other MLS pressure level is indicated in panel (a).

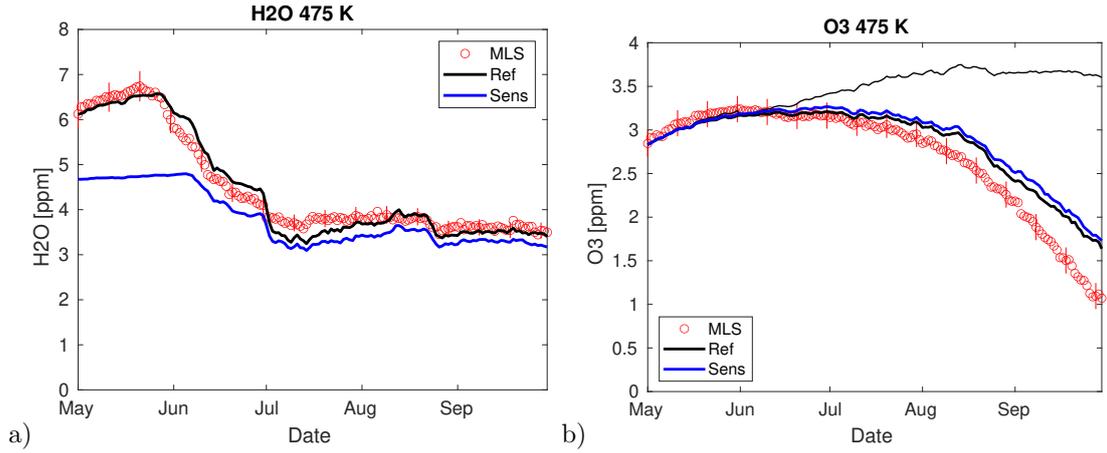


Figure 2. Left (a): Vortex-averaged water vapor at 475 K potential temperature observed by MLS in 2023 (red dots) and modelled by ATLAS (blue and black lines). The black line shows the reference run initialized with MLS measurements from 2023, while the blue line shows the sensitivity run initialized with MLS water vapor data from 2022 (i.e., without the effect of Hunga Tonga on water vapor). Right (b): Same for ozone. The thin black line shows a passive ozone tracer initialized on 1 June. The difference between the thin black line and the other lines quantifies the amount of ozone loss.

117 within the sunlit portion of the vortex from July through September. The thin black line
 118 shows a passive ozone tracer initialized on 1 June. The difference between the passive
 119 ozone tracer and the other lines quantifies the amount of ozone loss. There is only a short
 120 time period before July (cf. Figure 2 (a)) when differences in water vapor between the
 121 runs can have a direct effect on differences in ozone loss.

122 The difference between the modelled ozone and the passive ozone tracer underestimates
 123 ozone loss by about 0.7 ppm (30 %) compared to the difference between the pas-
 124 sive ozone tracer and MLS ozone. Since ATLAS runs for other Antarctic winters agree
 125 better with MLS (e.g., Wohltmann et al., 2021), this discrepancy might hint that other
 126 effects from the Hunga Tonga eruption are at work. However, as outlined before, it can-
 127 not be the direct effect of water vapor.

128 The results at other potential temperature levels lead to similar conclusions, and
 129 the effect of the increased water vapor remains small throughout the ozone column. Fig-
 130 ure 3 shows the chemical ozone loss modelled by ATLAS for the partial column from the
 131 lower model boundary at 157 hPa to 28.6 hPa. Ozone loss was determined by subtract-
 132 ing the passive ozone tracer initialized on 1 June from the modelled ozone values. Since
 133 the value of the passive ozone tracer is not known for air masses that entered through
 134 the upper model boundary after 1 June and descended in the vortex, the column is re-
 135 stricted to 28.6 hPa. The figure shows that the effect of the increased water vapor on
 136 the column loss in the model is small.

137 4 Discussion and Summary

138 There are several ways in which water vapor might change ozone loss in addition
 139 to dehydration:

- 140 • The threshold temperatures for the formation of all PSC types (supercooled ternary
 141 solution (STS), nitric acid trihydrate (NAT) and ice) are increased by the increased

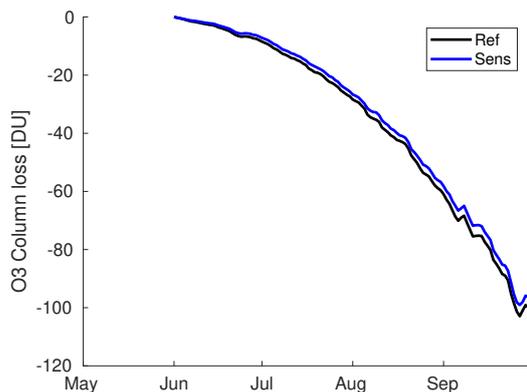


Figure 3. Vortex-averaged chemical loss of ozone modelled by ATLAS for the partial column from the lower model boundary at 157 hPa to 28.6 hPa. The black line shows the reference run initialized with MLS measurements from 2023, while the blue line shows the sensitivity run initialized with MLS water vapor data from 2022 (i.e., without the effect of Hunga Tonga on water vapor).

142 water vapor. For example, a change in water vapor from 5 ppm to 6 ppm increases
 143 the threshold temperature for ice by 1.1 K at 50 hPa (Marti & Mauersberger, 1993),
 144 and for NAT by 0.8 K (for a HNO_3 mixing ratio of 10 ppb) (D. Hanson & Mauers-
 145 berger, 1988). STS has no defined threshold temperature for formation but shows
 146 a gradual increase of reaction rates and droplet volume with lower temperatures.
 147 However, an increase in water vapor from 5 ppm to 6 ppm has to a good approx-
 148 imation the same effect on the STS reactivity (time needed for the $\text{HCl} + \text{ClONO}_2$
 149 reaction to completely deplete one of the reaction partners) as a temperature change
 150 of 1 K (D. R. Hanson & Ravishankara, 1994; Carslaw et al., 1995) (for 2 ppb HCl,
 151 1 ppb ClONO_2 , 0.15 ppb H_2SO_4 , number density of droplets 10 cm^{-3}).

152 While these changes increase the volume inside the vortex where the formation
 153 of PSCs is possible, this can only have an effect before dehydration sets in. The
 154 effects of changes in threshold temperatures and STS reactivity are included in
 155 ATLAS. Figure 2 (b) shows that they have no large effect on ozone depletion.

156 • Changes in water vapor have an effect on the particle size distribution of all cloud
 157 types. There is more water vapor available above the saturation limit, and one might
 158 expect larger particle sizes and larger surface area densities. However, particle for-
 159 mation and growth is a complex process, and this might not be straightforward.
 160 As in many chemistry and transport models, PSCs are treated in a somewhat sim-
 161 plified manner in ATLAS (Tritscher et al., 2021). For NAT and ice clouds, a con-
 162 stant number density is assumed, and a uniform particle size is then calculated
 163 from the HNO_3 and H_2O available above the saturation pressure. That is, the uni-
 164 form particle size will simply increase in our model. For STS, a constant number
 165 density and a log-normal distribution is assumed that is scaled with the total liq-
 166 uid volume.

167 Figure 2 (b) shows that changes in size distribution have only a small effect on
 168 ozone depletion in ATLAS. However, particle growth and formation might be more
 169 complex in reality and the effect on ozone depletion may be larger. Changes in
 170 size distribution may also have an effect after dehydration sets in (i.e., differences
 171 in the size distribution of NAT, ice and STS PSCs caused by increased water va-
 172 por in May and June might be persistent or lead to changes in the size distribu-
 173 tion in later months).

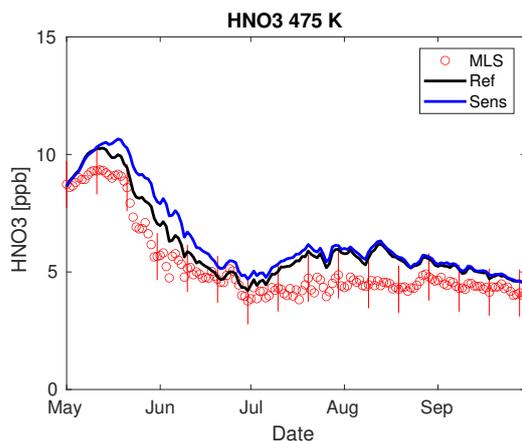


Figure 4. Vortex-averaged HNO_3 (gas phase) at 475 K potential temperature observed by MLS in 2023 (red dots) and modelled by ATLAS (blue and black lines). The black line shows the reference run initialized with MLS measurements from 2023, while the blue line shows the sensitivity run initialized with MLS water vapor data from 2022 (i.e., without the effect of Hunga Tonga on water vapor).

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- Since larger particles have a greater fall velocity, dehydration might also be faster in 2023. However, this is not considered in the simple ATLAS dehydration parameterization. The good agreement between ATLAS and MLS suggests that dehydration might be a fast enough process so that the speed of the process does not matter for the amount of ozone depletion.
 - Denitrification might also be affected by differences in the formation of the NAT particles. Denitrification is treated in a more sophisticated manner than dehydration in ATLAS and incorporates the nucleation, growth, sedimentation, and evaporation of individual particles with the DLAPSE model (Davies et al., 2005). Figure 4 shows that the changes in gas-phase HNO_3 (which are mainly caused by denitrification) are small between the reference run and the sensitivity run.

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In summary, we have shown that while the Hunga Tonga-Hunga Ha’apai eruption increased water vapor in the emerging Southern Hemisphere stratospheric polar vortex by 20–40% in 2023, ozone values at the end of September in the Antarctic ozone hole were in no way exceptional and were well within in the range of earlier years. ATLAS model runs indicate that the direct effect of the increased water vapor on Antarctic ozone depletion was minor. The reason for this are the very low temperatures in the vortex, which limit water vapor to the saturation pressure and tend to reset any anomalies in water vapor through dehydration before they can have an effect on ozone loss. However, the ATLAS runs underestimate observed ozone loss in 2023 by about 30% (in contrast to good agreement with observations in similar studies for other winters). Further studies would be needed to understand the causes of this discrepancy and whether it may be related to other effects from the Hunga Tonga-Hunga Ha’apai eruption in addition to the increased water vapor and dehydration.

198 Open Research Section

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ATLAS source code is available from the repository at <https://gitlab.awi.de/iwohltmann/atlas-julia>. MLS data are available at <https://disc.gsfc.nasa.gov/datasets?page=1&keywords=AURA%20MLS>. ECMWF ERA5 data are available at Hersbach et al. (2017), doi:10.24381/cds.143582cf.

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