

Assessment of the Madden-Julian Oscillation in CMIP6 Models based on Moisture Mode Theory

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Abstract

The moist processes of the Madden-Julian Oscillation (MJO) in the Coupled Model Intercomparison Project Phase 6 models are assessed using moisture mode theory-based diagnostics over the Indian Ocean (10°S-10°N, 75°E-100°E). Results show that no model can capture all the moisture mode properties relative to the reanalysis. Most models satisfy weak temperature gradient balance but have unrealistically fast MJO propagation and a lower moisture-precipitation correlation. Models that satisfy the most moisture mode criteria reliably simulate the background moist static energy (MSE) and low-level zonal winds compared to models that satisfy the least amount of criteria. The MSE budget associated with the MJO is also well-represented in the good models rather than in the poor models. Our results show that capturing the MJO's moisture mode properties over the Indian Ocean is associated with a more realistic representation of the MJO simulation and thus can be employed to diagnose MJO performance.

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2 **Models based on Moisture Mode Theory**

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5 **Key Points:**

- 6 • MJO simulation skill in 25 CMIP6 models is assessed using moisture mode the-
7 ory.
8 • No model can realistically reproduce all the moisture mode properties of the MJO
9 over the Indian Ocean.
10 • Models that best capture the MJO's moisture mode features exhibit more real-
11 istic mean states and MJO structure.

Abstract

The moist processes of the Madden-Julian Oscillation (MJO) in the Coupled Model Intercomparison Project Phase 6 models are assessed using moisture mode theory-based diagnostics over the Indian Ocean (10°S-10°N, 75°E-100°E). Results show that no model can capture all the moisture mode properties relative to the reanalysis. Most models satisfy weak temperature gradient balance but have unrealistically fast MJO propagation and a lower moisture-precipitation correlation. Models that satisfy the most moisture mode criteria reliably simulate the background moist static energy (MSE) and low-level zonal winds compared to models that satisfy the least amount of criteria. The MSE budget associated with the MJO is also well-represented in the good models rather than in the poor models. Our results show that capturing the MJO's moisture mode properties over the Indian Ocean is associated with a more realistic representation of the MJO simulation and thus can be employed to diagnose MJO performance.

Plain Language Summary

The Madden-Julian Oscillation is the most important tropical phenomenon that drives weather at the intraseasonal time scale. Although the MJO has been analyzed for the past decades, its simulation in climate models can still be improved. Previous studies have emphasized that the MJO evolution is tightly modulated by moisture fluctuations and posited the moisture mode theory to explain its behavior. Here, we show that no climate model can realistically reproduce the moist thermodynamics of the MJO, particularly its sensitivity to humidity anomalies. A few models can simulate some basic MJO features.

1 Introduction

The Madden-Julian Oscillation (MJO; Madden & Julian, 1971, 1972) is a planetary-scale envelope of convection that is coupled with the circulation and moisture (Raymond & Fuchs, 2009; Sobel & Maloney, 2013; Á. F. Adames & Kim, 2016, among others). This convective envelope often initiates over the Indian Ocean (IO) and propagates eastward at about 3 to 5 m s⁻¹ (C. Zhang & Ling, 2017; Rushley et al., 2022). The MJO affects weather and climate phenomena around the globe through its teleconnections, including Asia and Australian rainfall events (Chang et al., 2021; Bagtasa, 2020; Cowan et al., 2022; Dao et al., 2023), tropical cyclone genesis (J.-M. Chen et al., 2018; Rahul et al., 2022), El Niño Southern Oscillation and Atlantic Niño (Hendon et al., 2007; S.-K. Lee et al., 2023), as well as heatwaves and the frequency of tornadoes and hailstorms in the Northern America (Y.-Y. Lee & Grotjahn, 2019; Miller et al., 2022). In part due to these impacts, many studies in the last decades have tried to better understand the MJO through a combination of observation, theory, and modeling experiments (e.g., Raymond & Fuchs, 2009; Maloney et al., 2010; Sobel & Maloney, 2012; Á. F. Adames & Kim, 2016; Wang et al., 2016, and references therein).

Numerous studies have observed that the growth of MJO convection is associated with feedbacks that increase moisture anomalies (Sobel et al., 2014; Del Genio & Chen, 2015; B. Zhang et al., 2019). Moreover, the eastward propagation of MJO is predominantly governed by horizontal and vertical moisture advection (Kiranmayi & Maloney, 2011; Kim, Kug, & Sobel, 2014; K.-C. Tseng et al., 2015; Á. F. Adames & Wallace, 2015; Hung & Sui, 2018). These features have led to a view of MJO that has become a basis of moisture mode theory. Moisture mode theory posits that the MJO is tightly modulated and organized by moisture fluctuations, while temperature anomalies play a minor role because of weak temperature gradient (WTG) balance (Emanuel et al., 1994; Raymond & Fuchs, 2009; Sobel et al., 2014; Á. F. Adames & Kim, 2016; Á. F. Adames, 2017; Ahmed et al., 2021; A. F. Adames & Maloney, 2021; Mayta & Adames Corraliza, 2023, among others). The processes that lead to the moisture fluctuations also lead to

the evolution of moisture mode. According to these conditions, Mayta et al. (2022) proposed a series of moisture mode criteria to analyze the different tropical waves. Further, by using these criteria, Mayta and Adames Corraliza (2023) found that MJO behaves as a moisture mode only over the IO region. Outside this region, temperature fluctuations are as influential as moisture anomalies in MJO’s thermodynamics because a faster propagation of MJO prevents WTG balance.

Although our understanding of the MJO has significantly improved, accurate representation of MJO variability remains a major challenge in global climate models (GCMs; Kim et al., 2009; Ahn et al., 2017, 2020). It is well-documented that the failure of models to simulate the MJO is largely a result of inadequate treatment of deep cumulus convection, particularly its insufficient sensitivity to free tropospheric water vapor (e.g., Maloney & Hartmann, 2001; M.-I. Lee et al., 2003; Holloway et al., 2013; Kim, Lee, et al., 2014). Models in which convection is sensitive to water vapor fluctuations produce regions of precipitation that persist at the intraseasonal timescale, hence producing MJO activity. From this, models that have a strong coupling of precipitation with low-level wind field and moisture can simulate more realistic MJO convection (Holloway et al., 2013; Ahn et al., 2017). Furthermore, a strong horizontal gradient of mean state moisture can drive robust MJO propagation (Jiang, 2017; Ahn et al., 2020). All of these features are consistent with the MJO being at least partially explained as a moisture mode.

Based on these previous results, we hypothesize that the moisture mode properties of the MJO are essential for its realistic simulation. To this end, we seek to examine the MJO simulation in the Coupled Model Intercomparison Project Phase 6 (CMIP6) models based on the moisture mode framework (Ahmed et al., 2021; Mayta et al., 2022; Mayta & Adames, 2023). Specifically, we seek to answer the following questions:

- Q1: Can the global climate models reproduce the MJO moisture mode properties over the Indian Ocean?
- Q2: If a model can capture the moisture mode behaviors, does it mean that the model has better skills in the MJO simulation than others?

The structure of this research is as follows. Section 2 describes the datasets and methods. Section 3 diagnoses MJO simulation by the moisture mode theory. In section 4, we compare the good and poor simulations in the moisture mode behaviors against the observations. Major findings are summarized in section 5.

2 Data Description, Processing, and Diagnostics

2.1 Data Sources

25 CMIP6 models (Eyring et al., 2016) are adopted to evaluate MJO simulation in the 20-year (1995-2014) historical scenario (Table 1). We primarily use r1i1p1f1 ensemble member for most models, except for EC-Earth3 (r3i1p1f1), HadGEM3-GC31-LL (r1i1p1f3), HadGEM3-GC31-MM (r3i1p1f3), and UKESM1-0-LL (r1i1p1f2) based on their available data. Models that cannot provide all radiative fluxes to compute net radiation within the atmosphere are marked with asterisks (*).

Observation and reanalysis data are used as a reference for model simulations. We use the moisture, precipitation, temperature, horizontal winds, vertical velocity, geopotential height, radiation, and surface fluxes from the fifth generation of the European Centre for Medium-Range Weather Forecasts (ECMWF) reanalysis (ERA5; Hersbach et al., 2019). The outgoing longwave radiation (OLR) from NOAA Physical Sciences Laboratory (Liebmann & Smith, 1996) is used to calculate the MJO index. The precipitation from Tropical Rainfall Measuring Mission (TRMM; Kummerow et al., 2000) product is applied to compute the realistic wave responses by the space-time power spectra. All data

Models	Institutions	Lat × Lon
ACCESS-CM2*	Commonwealth Scientific and Industrial Research Organisation, Australian Research Council Centre of Excellence for Climate System Science	144 × 192
AWI-ESM-1-1-LR	Alfred Wegener Institute, Helmholtz Centre for Polar and Marine Research	96 × 192
BCC-ESM1*	Beijing Climate Center	64 × 128
CESM2		192 × 288
CESM2-FV2	National Center for Atmospheric Research, Climate and Global Dynamics Laboratory	96 × 144
CESM2-WACCM		192 × 288
CESM2-WACCM-FV2		96 × 144
EC-Earth3*	Rosby Center, Swedish Meteorological and Hydrological Institute	256 × 512
FGOALS-g3*	Chinese Academy of Sciences	80 × 180
GFDL-CM4	National Oceanic and Atmospheric Administration, Geophysical Fluid Dynamics Laboratory	90 × 144
HadGEM3-GC31-LL	Met Office Hadley Centre	144 × 192
HadGEM3-GC31-MM		324 × 432
IITM-ESM*	Centre for Climate Change Research, Indian Institute of Tropical Meteorology Pune	94 × 192
INM-CM4-8*	Institute for Numerical Mathematics, Russian Academy of Science	120 × 180
INM-CM5-0*		120 × 180
IPSL-CM6A-LR	Institut Pierre Simon Laplace	143 × 144
IPSL-CM6A-LR-INCA		143 × 144
KACE-1-0-G*	National Institute of Meteorological Sciences/Korea Meteorological Administration Japan Agency for Marine-Earth Science and Technology	144 × 192
MIROC6*	Atmosphere and Ocean Research Institute National Institute for Environmental Studies / RIKEN Center for Computational Science	128 × 256
MPI-ESM-1-2-HAM		96 × 192
MPI-ESM1-2-HR	Max Planck Institute for Meteorology	192 × 384
MPI-ESM1-2-LR		96 × 192
MRI-ESM2-0	Meteorological Research Institute	160 × 320
TaiESM1*	Research Center for Environmental Changes, Academia Sinica	192 × 288
UKESM1-0-LL	Met Office Hadley Centre	144 × 192

Table 1. List of 25 CMIP6 models used in this study, including their research centers and horizontal resolutions. The models with an asterisk (*) mean lack complete radiation fluxes for calculating net radiation heating within the atmosphere.

110 are interpolated into a uniform horizontal resolution of 2.5° longitude \times 2.5° latitude.
 111 We discuss the MJO activity during the extended boreal winter (November to April) when
 112 the MJO is more active (Q. Zhang et al., 2019; X. Li et al., 2020).

113 2.2 Filtering, EOF Analysis, and Regressions

114 The dominant mode of the MJO convection is derived through the empirical or-
 115 thogonal function (EOF) analysis of 20-96-days bandpass filtered OLR over the equa-
 116 torial belt (15°S - 15°N). The first EOF mode (EOF1) has the largest amplitude around
 117 90°E , corresponding to the MJO convection (not shown). The field variables are regressed
 118 onto the first principal component (PC1) time series to obtain a composite of the MJO
 119 evolution from -30 to 30 days, following the same process as previous studies (e.g., Á. F. Adames
 120 et al., 2021; Mayta et al., 2021; Mayta & Adames Corraliza, 2023). These perturbations
 121 are then scaled to one standard deviation of PC1. To make the strongest MJO convec-
 122 tion occur near 90°E at lag 0 day in all data, we refer to the basis function approach (J. Lee
 123 et al., 2019; Orbe et al., 2020, among others) to project simulated OLR anomalies onto
 124 the observed EOF1 and hence obtain the PC1 time series of each model.

125 2.3 Diagnostic Criteria

126 In order to evaluate the moist thermodynamics of simulated MJO, we apply the
 127 moisture mode criteria over the Indian Ocean (IO; 10°S - 10°N , 75°E - 100°E) region, where
 128 the MJO shows characteristics of a moisture mode (Mayta & Adames Corraliza, 2023).
 129 The criteria are (Ahmed et al., 2021; Mayta et al., 2022):

- 130 1. Wave must exhibit a large moisture signature that is highly correlated with the pre-
 131 cipitation anomalies

132 To be considered a moisture mode, the MJO's precipitation anomalies P' should
 133 be sensitive to the column water vapor $\langle q \rangle'$ variations. In other words, it must ex-
 134 hibit a high coherence between $\langle q \rangle'$ and P' ($R_{P,q}$) as follows,

$$\langle q \rangle' \propto P' \quad (1)$$

135 where $\langle \cdot \rangle \equiv \frac{1}{g} \int_{100}^{1000} (\cdot) dp$ is vertical integration from 1000 to 100 hPa, and primes
 136 ($'$) represent the regressed anomalies. The correlation should be higher than 0.9,
 137 indicating that the moisture fluctuations significantly modulate the precipitation
 138 evolution, at least 81% of the variance. The slope of $\langle q \rangle'$ to P' is the convective
 139 moisture adjustment time scale ($\tau_c \equiv \frac{\langle q \rangle'}{P'}$), defined as the time to remove col-
 140 umn moisture through the rainfall (Betts, 1986). The τ_c of MJO convection should
 141 be about 1 day over the IO region (Mayta & Adames Corraliza, 2023).

142 2. *The system must be in the weak temperature gradient (WTG) balance*

143 Under WTG approximation, the vertical advection of dry static energy $\langle \omega \partial_p s \rangle'$
 144 must exhibit a balance with apparent heat source $\langle Q_1 \rangle'$, expressed as:

$$\langle \omega \partial_p s \rangle' \simeq \langle Q_1 \rangle' \quad (2)$$

145 where $s = C_p T + gz$ is dry static energy (DSE). To satisfy this second criterion,
 146 the slope of $\langle Q_1 \rangle'$ to $\langle \omega \partial_p s \rangle'$ should be close to 1 in linear least-squares fitting, and
 147 their correlation must also be higher than 0.9.

148 3. *Moisture must govern the evolution of moist static energy*

149 If the MJO is a moisture mode, the column water vapor must be the main con-
 150 tributor to its moist static energy (MSE, m), giving the following relation,

$$\langle m \rangle' = \langle s \rangle' + \langle L_v q \rangle' \approx \langle L_v q \rangle' \quad (3)$$

151 To guarantee the approximation in Eq. (3), a slope of $\langle L_v q \rangle'$ to $\langle m \rangle'$ must be \sim
 152 1 in linear least-squares fitting ($S_{q,m}$), with a high coherence between both vari-
 153 ables (> 0.9).

154 4. N_{mode}

155 The dimensionless N_{mode} parameter is also adopted to quantify the relative im-
 156 portance of column water vapor versus temperature in the evolution of MSE (Á. F. Adames
 157 et al., 2019). N_{mode} can be defined as in Á. F. Adames et al. (2019) and Mayta
 158 et al. (2022) as follows,

$$N_{mode} \simeq \frac{c_p^2 \tau}{c^2 \tau_c} \quad (4)$$

159 where $c = 50 \text{ m s}^{-1}$ is the phase speed of a first baroclinic free gravity wave, c_p
 160 is the phase speed of MJO over the warm pool Indian Ocean, and τ is the char-
 161 acteristic temporal scale of MJO (i.e., ~ 37 days in the ERA5). The c_p is estimated
 162 by using the Radon Transform method (Radon, 1917; Mayta et al., 2023), which
 163 is described in the Supplementary Information (SI). The MJO can be classified
 164 as a moisture mode when $N_{mode} \ll 1$ (i.e., $\log_{10} N_{mode} < -0.5$).

165 3 Moist Thermodynamic Diagnostics of MJO Simulation

166 As in Mayta et al. (2022), the moisture mode criteria are applied to the reanaly-
 167 sis and models by constructing scatterplots. The results are summarized in Figure 1.

168 Reanalysis, as recently documented in Mayta and Adames Corraliza (2023), shows
 169 a high correlation between $\langle q \rangle'$ and P' over the IO region ($R_{P,q} = 0.95$), whereas the
 170 climate models depict an average value of 0.88 ± 0.05 . Among them, HadGEM3-GC31-
 171 LL, KACE-1-0-G, and TaiESM1 models have the highest correlation ($R_{P,q} = 0.94$). The
 172 remaining 15 models underestimate $R_{P,q}$ (< 0.9 ; black values). The τ_c of ERA5 is about

	$R_{P,q}$	τ_c	$S_{q,m}$	$\log_{10} N_{mode}$	
ERA5	0.95	1.02	0.98	-0.69	
ACCESS-CM2	0.88	1.10	0.96	-0.06	
AWI-ESM-1-1-LR	0.92	0.86	0.86	-0.97	×
BCC-ESM1	0.86	0.83	0.89	-0.64	
CESM2	0.85	1.15	0.87	-0.60	
CESM2-FV2	0.92	1.11	0.89	-0.61	○
CESM2-WACCM	0.86	1.05	0.92	-0.57	
CESM2-WACCM-FV2	0.88	1.15	0.89	-0.39	
EC-Earth3	0.89	1.28	1.00	-0.65	○
FGOALS-g3	0.87	1.08	0.87	-0.69	
GFDL-CM4	0.90	0.93	0.90	-0.45	○
HadGEM3-GC31-LL	0.94	1.19	0.97	-0.25	
HadGEM3-GC31-MM	0.92	1.15	0.94	-0.05	
IITM-ESM	0.93	0.99	0.82	-0.98	
INM-CM4-8	0.76	1.08	1.02	-0.02	×
INM-CM5-0	0.80	0.95	1.01	0.00	
IPSL-CM6A-LR	0.84	1.10	0.95	-0.68	
IPSL-CM6A-LR-INCA	0.79	0.98	0.88	-1.03	×
KACE-1-0-G	0.94	1.27	0.95	0.28	
MIROC6	0.88	1.20	0.92	-0.75	○
MPI-ESM-1-2-HAM	0.88	0.60	0.83	-0.94	×
MPI-ESM1-2-HR	0.91	0.83	0.87	-0.94	
MPI-ESM1-2-LR	0.88	0.68	0.85	-0.94	
MRI-ESM2-0	0.88	1.30	1.00	-0.20	
TaiESM1	0.94	1.37	1.05	-0.49	
UKESM1-0-LL	0.92	1.27	0.94	-0.27	

Figure 1. The values of criteria $R_{P,q}$, τ_c , $S_{q,m}$, and $\log_{10} N_{mode}$ from the ERA5 and 25 CMIP6 models. Numbers in blue represent values that satisfy the moisture mode criteria: (1) $R_{P,q} > 0.9$, and τ_c within ± 0.5 standard deviations relative to the ERA5 (0.93-1.12 days); (2) $S_{q,m} \sim 1$ (0.9-1.05); and (3) $\log_{10} N_{mode}$ ranges from -0.8 to -0.5. For model selection, the green boxes indicate model values within ± 1.5 standard deviations from the reanalysis, while the orange boxes represent the $\log_{10} N_{mode}$ within the range of -0.8 to -0.3. The relatively good and poor models are marked by the green circles and red crosses, respectively.

173 1.02 day. Only 10 models are within ± 0.5 standard deviation (SD) relative to the reanal-
 174 ysis (0.93-1.12 days). For WTG approximation (not shown), the slope of $\langle Q_1 \rangle'$ versus
 175 $\langle \omega \partial_p s \rangle'$ in reanalysis is 0.99. The values of the 25 models range from 0.98 to 1.06, and
 176 the multi-model mean is 1.01 ± 0.02 . The correlation between $\langle Q_1 \rangle'$ and $\langle \omega \partial_p s \rangle'$ is higher
 177 than 0.99 in the ERA5 and all models included. It suggests that these simulations largely
 178 satisfy the WTG balance over the IO region, so this criterion is not shown in Figure 1.
 179 $S_{q,m}$ in ERA5 is approximately 0.98 (Fig. 1). The mean of the 25 models is $\sim 0.92 \pm$
 180 0.06 , with most models showing values ranging from 0.89 to 1.02 (within 1.5 SD rela-
 181 tive to ERA5). However, the $S_{q,m}$ in the 11 models are lower than 0.9, particularly for
 182 the IITM-ESM and MPI-ESM-1-2-HAM models (< 0.85). The relatively low $S_{q,m}$ val-
 183 ues indicate that the contribution of $\langle s \rangle'$ to $\langle m \rangle'$ is more significant. All models and re-
 184 analysis have a high correlation coefficient (> 0.98) between moisture and MSE anomalies
 185 (not shown).

186 The $\log_{10} N_{mode}$ value of ERA5 is approximately -0.69 ($N_{mode} \sim 0.2$), indicating
 187 that MJO exhibits moisture mode behavior over the IO region, in agreement with Mayta
 188 and Adames Corraliza (2023). Eight models depict a $\log_{10} N_{mode} > -0.3$ ($N_{mode} > 0.5$),
 189 implying that their wave behavior is far from the moisture mode regime. It is worth not-
 190 ing that some models show almost good results for the first three moisture mode crite-
 191 ria but have $\log_{10} N_{mode} > -0.25$ (e.g., ACCESS-CM2, HadGEM3-GC31-MM, and KACE-
 192 1-0-G). Á. F. Adames et al. (2019) and Á. F. Adames (2022) performed a scale analy-
 193 sis and demonstrated that N_{mode} is largely determined by the phase speed of the wave.
 194 These models, as expected, simulate a faster MJO phase speed ($c_p > 8 \text{ ms}^{-1}$) than ERA5
 195 (Fig. S1). A high sensitivity of N_{mode} to c_p was found in these 25 CMIP6 models (fur-
 196 ther discussion in SI). On the other hand, the $\log_{10} N_{mode} < -0.8$ ($N_{mode} < 0.16$; e.g.,
 197 AWI-ESM-1-1-LR, IITM-ESM, IPSL-CM6A-LR-INCA, MPI-ESM-1-2-HAM, MPI-ESM1-
 198 2-HR, and MPI-ESM1-2-LR) are models with near stationary MJO-like behavior ($c_p <$
 199 2.2 ms^{-1}).

200 According to the moisture mode criteria for the MJO's behavior (e.g., $R_{P,q} > 0.9$,
 201 $\tau_c \sim 1$ day, $S_{q,m} \sim 1$, and that $\log_{10} N_{mode}$ ranges within -0.8 to -0.5; marked by the
 202 blue values in Fig. 1), no model accurately captures all the MJO's moisture mode prop-
 203 erties as previously observed. However, some models still have reasonable values close
 204 to the observations but with a slightly long τ_c or low $S_{q,m}$.

205 4 Comparison between observations, Good and Poor Models

206 In this section, we further discuss whether these models have better skills associ-
 207 ated with the MJO simulation than others if they can approximately capture the mois-
 208 ture mode behavior. To this end, we consider the relaxed criteria ranges within ± 1.5 SD
 209 relative to the reanalysis for $R_{P,q}$, τ_c , and $S_{q,m}$ (green boxes in Fig. 1). $R_{P,q}$ and τ_c must
 210 be considered as one criterion. $\log_{10} N_{mode}$ should be -0.8 to -0.3 (orange boxes) because
 211 $\log_{10} N_{mode} < -0.3$ represents a higher contribution from the moisture fluctuation than
 212 the temperature fluctuation ($N_{mode} < 0.5$). Based on these conditions, four relatively
 213 good models (RGMs; CESM2-FV2, EC-Earth3, GFDL-CM4, and MIROC6) and four
 214 relatively poor models (RPMs; AWI-ESM-1-1-LR, INM-CM4-8, IPSL-CM6A-LR-INCA,
 215 and MPI-ESM-1-2-HAM) are selected. We do not adopt more than one model from the
 216 same research institution to avoid the multi-model means that are dominated by sim-
 217 ilar simulations. If the models are from the same research center, the model with the best
 218 (worst) performance was selected in the relatively good (poor) model groups. For instance,
 219 in the poor model group, we selected MPI-ESM-1-2-HAM rather than MPI-ESM1-2-LR
 220 because the former has a lower values in τ_c and $S_{q,m}$ than the latter.

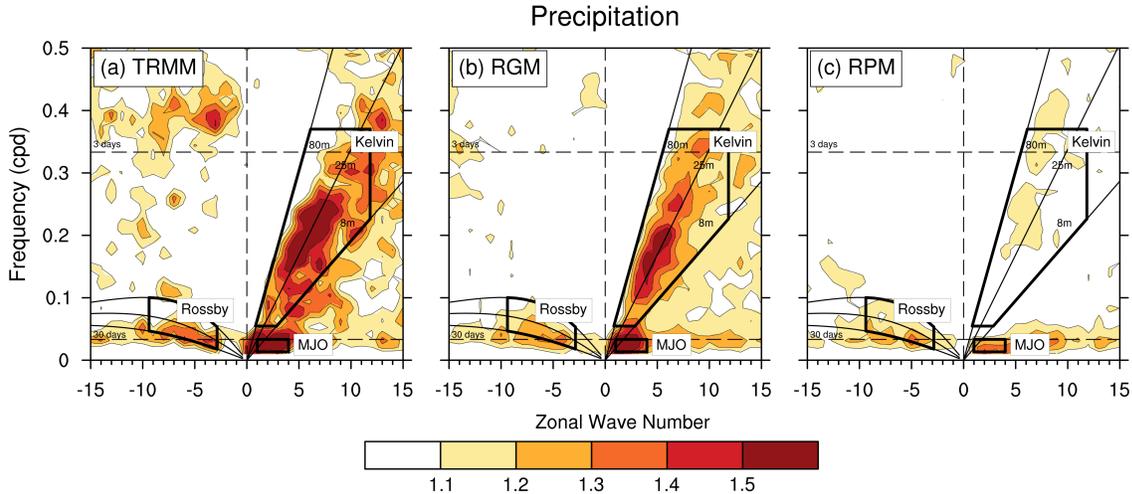


Figure 2. Space-time spectrum of the precipitation averaged between 10°S and 10°N for (a) TRMM, (b) the RGMs ensemble, and (c) the RPMs ensemble. The solid dispersion curves correspond to 8 m, 25 m, and 80 m equivalent depths. Color shading interval is 0.1.

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4.1 Space-Time Spectrum

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First, we computed space-time power spectra, making use of the fast Fourier transform (FFT). The calculation procedure is similar to those used by previous studies (e.g., Wheeler & Kiladis, 1999; Rushley et al., 2019; Y. Li et al., 2022, among others). We used precipitation from TRMM, RGMs, and RPMs as an input for the calculation. The results of TRMM and ERA5 are similar (not shown), although the ERA5 reanalysis (1995-2014) has a longer period than TRMM observation (1998-2014).

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Figure 2 shows the symmetric power spectra of precipitation in the frequency-wavenumber domain. For the MJO band (wavenumber $k = 1 - 4$, and period of 30 – 90 days), the TRMM and ensemble good models show strong spectra (power > 1.5), whereas it is relatively weak in the poor model group (power < 1.4). While precipitation exhibits a strong Kelvin wave signal in the observation and RGMs, such a signal is largely weak in the RPMs. Overall, the good model group can capture better wave signals and intensities than the poor model group.

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4.2 Mean State

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In previous studies, advection of the mean MSE has been found to be critical for MJO simulation (e.g., Jiang, 2017; Ahn et al., 2020; Ren et al., 2021, and others). On the other hand, the absence of mean low-level westerly winds in the western Pacific can lead to a non-propagating MJO in the model simulation (Inness & Slingo, 2003). Thus, it is worthwhile to compare the mean-state column-integrated MSE and 850-hPa zonal winds between the reanalysis, RGM, and RPM (Fig. 3). The pattern correlation (Cor) and root mean square deviation (RMSD) between the individual model group and the reanalysis are also shown in the upper right corner of each panel in Figure 3.

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For ERA5, the relatively high column MSE is concentrated over the Indo-Pacific warm pool and decreases with higher latitude (Fig. 3a). The models simulate a similar but underestimated column MSE distribution compared with the reanalysis. The RGM has a higher column MSE over the equatorial warm pool relative to the RPM, especially in the western Pacific with the MSE extreme. This leads to stronger background zonal

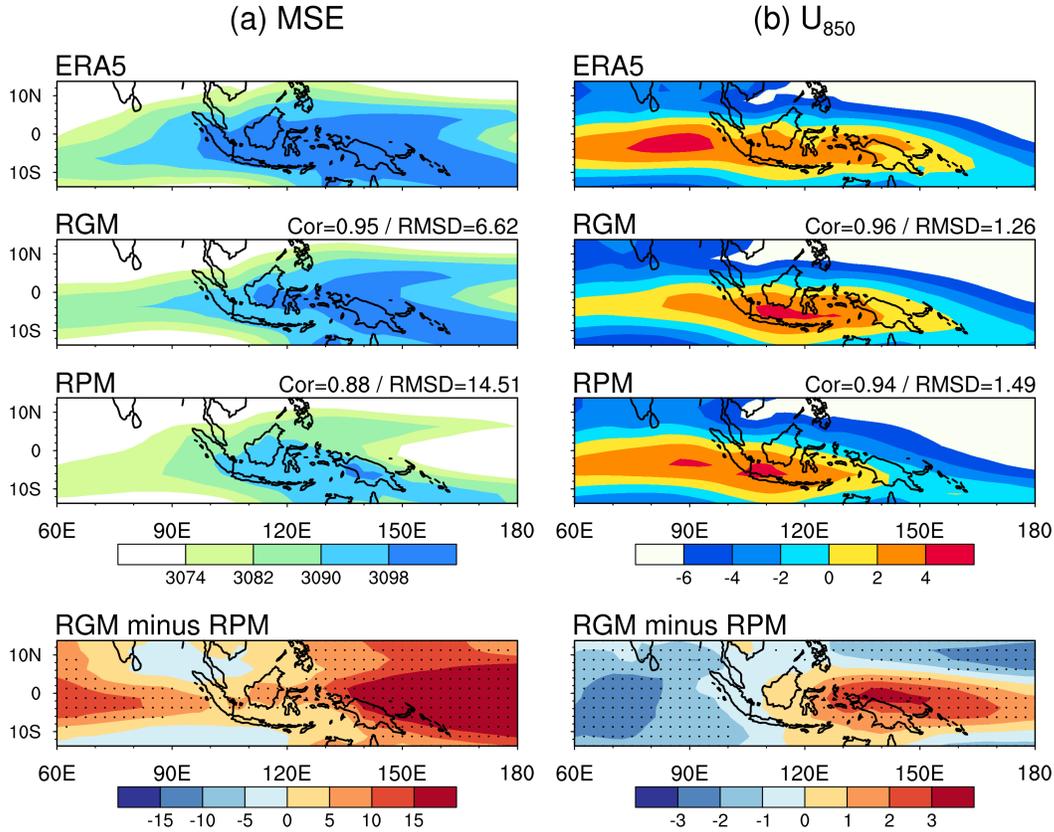


Figure 3. Spatial pattern of mean-state (a) column-integrated MSE (10^6 Jm^{-2}) and (b) 850-hPa zonal wind (ms^{-1}) for the boreal winter, derived from (top) the ERA5, (middle top) ensemble good model, (middle bottom) ensemble poor model group, and (bottom) the difference between RGM and RPM (RGM minus RPM). The gray dots in the bottom panels indicate statistical significance at the 95% confidence level. The pattern correlation (Cor) / root-mean-square deviation (RMSD) between the model group and reanalysis is presented at the top-right corners, respectively.

249 and meridional gradients of MSE in the good model group than in the poor model group.
 250 The observed westerlies cover the tropical warm pool from 60°E to 165°E , while the max-
 251 imum wind speed occurs in the IO region (Fig. 3b). For the model simulation, the peak
 252 of zonal winds appears near the Maritime Continent, resulting in weaker westerlies over
 253 the IO region than reanalysis. In addition, the good model group simulates weaker (stronger)
 254 westerlies than the poor model group in the Indian Ocean (western Pacific) region. The
 255 westerlies can extend toward 160°E in the RGM; however, they are replaced by the strong
 256 easterly winds at 140°E in the RPM, especially for AWI-ESM-1-1-LR, INM-CM4-8 and
 257 IPSL-CM6A-LR-INCA (not shown), where the absence of background westerlies might
 258 partially explains why their MJO convection can not propagate across the Maritime Con-
 259 tinent (Fig. S1).

260 4.3 Moist Static Energy Budget Analysis

261 The MSE budget is widely used to investigate the moist energy recharging and dis-
 262 charging associated with the MJO evolution (Inoue & Back, 2015; Ren et al., 2021; W.-
 263 L. Tseng et al., 2022, and others), taking the following form:

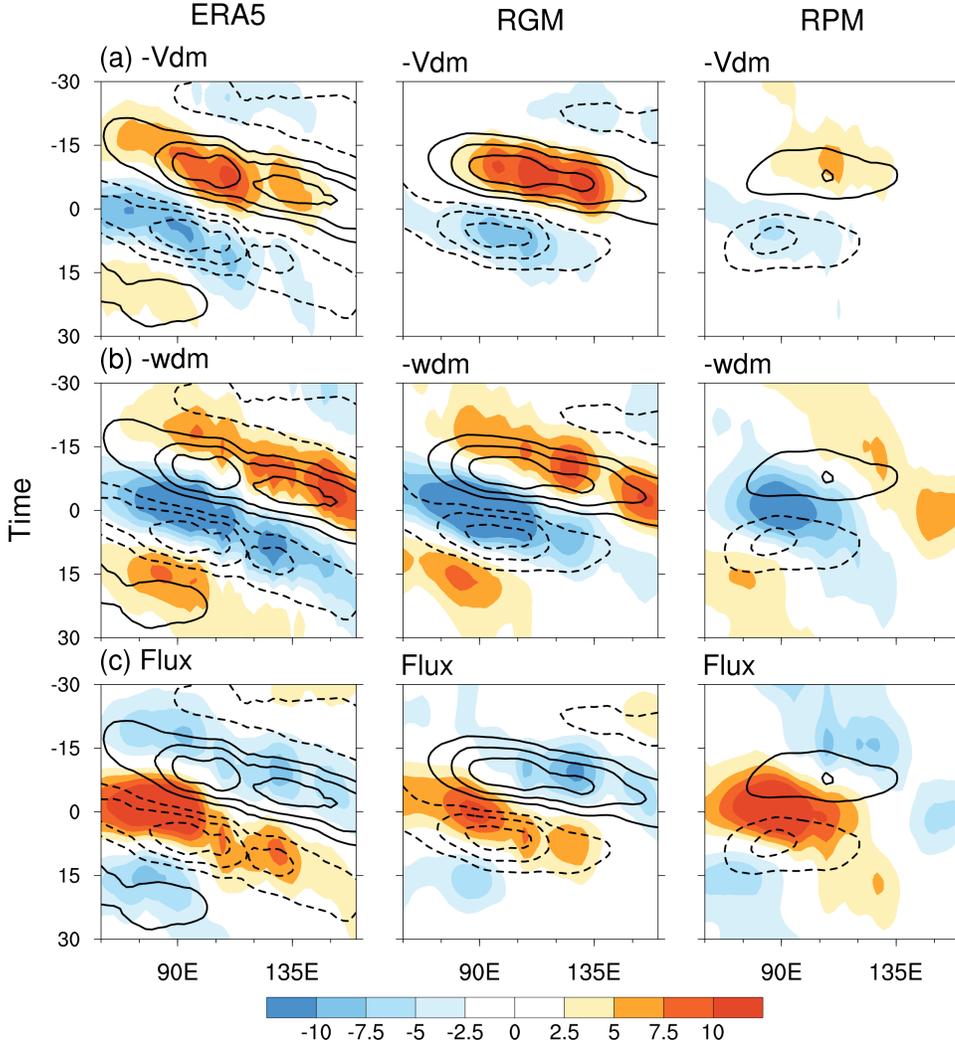


Figure 4. Hovmöller diagram of regressed MSE budget terms (shading) and $\langle \partial_t m \rangle'$ (contour) averaged between 10°S and 10°N from -30 to 30 days for the ERA5 (left panel), RGM (middle panel) and RPM (right panel). The individual terms are (a) $\langle -\mathbf{v} \cdot \nabla m \rangle'$, (b) $\langle -\omega \partial_p m \rangle'$, and (c) Flux' . The contour and shading intervals are 2 and 2.5 W m^{-2} , respectively.

$$\langle \partial_t m \rangle' = \langle -\mathbf{v} \cdot \nabla m \rangle' + \langle -\omega \partial_p m \rangle' + \text{Flux}' \quad (5)$$

264 where the left-side term in Eq. (5) is the MSE tendency. The first and second terms on
 265 the right-side represent the horizontal and vertical MSE advection, respectively. The other
 266 term of Eq. (5) is the flux term ($\text{Flux}' = \langle Q_r \rangle' + L_v E' + SH'$) that includes the col-
 267 umn radiative flux $\langle Q_r \rangle'$, surface latent heat flux $L_v E'$, and surface sensible heat flux
 268 SH' .

269 Figure 4 shows the Hovmöller diagram of the regressed MSE budget terms in Eq.
 270 5 for the ERA5 and model groups. The MSE budget terms display an eastward propa-
 271 gation in the reanalysis (Fig. 4a). $\langle -\mathbf{v} \cdot \nabla m \rangle'$ varies in phase with $\langle \partial_t m \rangle'$. RGM sim-
 272 ulations reproduce the eastward MJO convection with strong $\langle -\mathbf{v} \cdot \nabla m \rangle'$ and $\langle \partial_t m \rangle'$.
 273 The RPM exhibits weak and nearly non-propagating convection. In Figure 4b, $\langle -\omega \partial_p m \rangle'$

274 leads $\langle \partial_t m \rangle'$ in the reanalysis and the models. Compared with ERA5, $\langle -\omega \partial_p m \rangle'$ has a
 275 stronger drying effect (< 0) in the RGM and an underestimated amplitude in the RPM.
 276 The observed $Flux'$ exhibits a lagged evolution with $\langle \partial_t m \rangle'$ (Fig. 4c). The simulated
 277 $Flux'$ of RGM is weaker than the reanalysis, whereas the positive $Flux'$ in the RPM can
 278 not propagate into the western Pacific.

279 5 Summary and Conclusions

280 In this study, we applied the moisture mode theory-based diagnosis (Ahmed et al.,
 281 2021; Mayta et al., 2022; Mayta & Adames Corraliza, 2023) to assess the moisture mode
 282 properties of MJO over the Indian Ocean (10°S - 10°N , 75°E - 100°E) in the 25 CMIP6 mod-
 283 els. The following are answers to the two questions based on the results in Sections 3 to
 284 4:

285 Q1: Can the global climate models reproduce the MJO moisture mode properties over
 286 the Indian Ocean?

287 Our results demonstrate that none of the models used in this study could reliably
 288 reproduce all moist thermodynamic properties of the MJO as observed in Figure 1: (i)
 289 Few models showed a high correlation (greater than 0.9) between moisture and precip-
 290 itation anomalies and exhibited convective adjustment time scale (τ_c) that aligned with
 291 the reanalysis; (ii) All models can satisfy the criteria for weak temperature gradient (WTG)
 292 balance; (iii) Nevertheless, 11 models still exhibited an unrealistically high contribution
 293 from temperature fluctuations to the MSE anomalies; and (iv) limited number of mod-
 294 els showed values of N_{mode} that are close to those of the reanalysis data. High values of
 295 N_{mode} ($\gg 0.5$) or low N_{mode} ($\ll 0.16$) imply that many models showed unrealistically
 296 fast or nearly non-propagating MJO convection, respectively.

297 Q2: If a model can capture the moisture mode behaviors, does it mean that the model
 298 has better skills in the MJO simulation than others?

299 While no model fully captures the behavior of the MJO, there is a subset that per-
 300 forms reasonably well. These good models (e.g., CESM2-FV2, EC-Earth3, GFDL-CM4,
 301 and MIROC6) were selected based on their acceptable performance in the moisture mode
 302 criteria (Fig. 1). They also show a stronger wave response of the MJO signals compared
 303 to the relatively poor models (Fig. 2). The good model group realistically simulates the
 304 mean-state column MSE and low-level zonal winds (Fig. 3). MSE budget associated with
 305 the MJO is better represented in the good models rather than in the poor models group,
 306 especially in the MSE advection terms (Fig. 4). Our results indicate that models accu-
 307 rately capturing the moisture mode behavior of the MJO over the Indian Ocean demon-
 308 strate improved simulation of the MJO.

309 Most of these “acceptable” good models also depicted good performance in the MJO
 310 metrics proposed by previous studies (Ahn et al., 2020; Orbe et al., 2020; G. Chen et al.,
 311 2022; Y. Li et al., 2022). A robust MJO propagation is correlated with a more humid
 312 mean state with stronger horizontal moisture gradients, as well as more robust MJO wind
 313 anomalies (Ahn et al., 2020; Y. Li et al., 2022). This consistency makes sense since many
 314 previous studies have obtained these results under the a-priori assumption that the MJO
 315 behaves as a moisture mode. In other words, many previous studies implicitly assume
 316 that the moisture mode criteria are always satisfied, and that good MJO models are those
 317 that best simulate the processes that lead to the destabilization and propagation of mois-
 318 ture modes. These include having stronger horizontal moisture gradients that lead to
 319 more robust propagation via horizontal moisture advection, convection that is more sen-
 320 sitive to moisture variations, and a small effective gross moist stability (e.g., Benedict
 321 et al., 2014; Ahn et al., 2017, 2020).

Our study extends upon previous by showing that the expected moisture mode behavior only exists in models that more robustly simulate the MJO. Thus, the relatively good MJO models do not just simulate the processes that lead to the destabilization and propagation of moisture modes, they are also the models that best simulate the moisture mode behavior of the MJO over the Indian Ocean. Poorer models not only have weaker or non-propagating MJO-like variability, but this variability is inconsistent with moisture mode behavior. Thus, simulating an MJO that behaves as a moisture mode over the Indian may be synonymous with simulating a realistic MJO, and the four criteria used here appear to be useful diagnostic tools for evaluating MJO simulation performance.

In spite of these findings, we cannot say whether simulating the moisture mode behavior is what causes the models to perform better. It may be related to more realistic convection representation, or a combination of other factors. More work is needed to better understand the causality.

6 Open Research

We downloaded the CMIP6 model simulation outputs from the Lawrence Livermore National Laboratory (<https://esgf-node.llnl.gov/search/cmip6>). The interpolated OLR data was obtained from the NOAA (<https://psl.noaa.gov/data/gridded/data.interp.OLR.html>). The precipitation from Tropical Rainfall Measuring Mission (3B42) dataset was downloaded from the NASA (<https://disc.gsfc.nasa.gov/>). The reanalysis data was available at ECMWF (ERA5; <https://doi.org/10.24381/cds.adbb2d47>).

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Figure 1.

	$R_{P,q}$	τ_c	$S_{q,m}$	$\log_{10} N_{mode}$	
ERA5	0.95	1.02	0.98	-0.69	
ACCESS-CM2	0.88	1.10	0.96	-0.06	
AWI-ESM-1-1-LR	0.92	0.86	0.86	-0.97	✗
BCC-ESM1	0.86	0.83	0.89	-0.64	
CESM2	0.85	1.15	0.87	-0.60	
CESM2-FV2	0.92	1.11	0.89	-0.61	○
CESM2-WACCM	0.86	1.05	0.92	-0.57	
CESM2-WACCM-FV2	0.88	1.15	0.89	-0.39	
EC-Earth3	0.89	1.28	1.00	-0.65	○
FGOALS-g3	0.87	1.08	0.87	-0.69	
GFDL-CM4	0.90	0.93	0.90	-0.45	○
HadGEM3-GC31-LL	0.94	1.19	0.97	-0.25	
HadGEM3-GC31-MM	0.92	1.15	0.94	-0.05	
IITM-ESM	0.93	0.99	0.82	-0.98	
INM-CM4-8	0.76	1.08	1.02	-0.02	✗
INM-CM5-0	0.80	0.95	1.01	0.00	
IPSL-CM6A-LR	0.84	1.10	0.95	-0.68	
IPSL-CM6A-LR-INCA	0.79	0.98	0.88	-1.03	✗
KACE-1-0-G	0.94	1.27	0.95	0.28	
MIROC6	0.88	1.20	0.92	-0.75	○
MPI-ESM-1-2-HAM	0.88	0.60	0.83	-0.94	✗
MPI-ESM1-2-HR	0.91	0.83	0.87	-0.94	
MPI-ESM1-2-LR	0.88	0.68	0.85	-0.94	
MRI-ESM2-0	0.88	1.30	1.00	-0.20	
TaiESM1	0.94	1.37	1.05	-0.49	
UKESM1-0-LL	0.92	1.27	0.94	-0.27	

Figure 2.

Precipitation

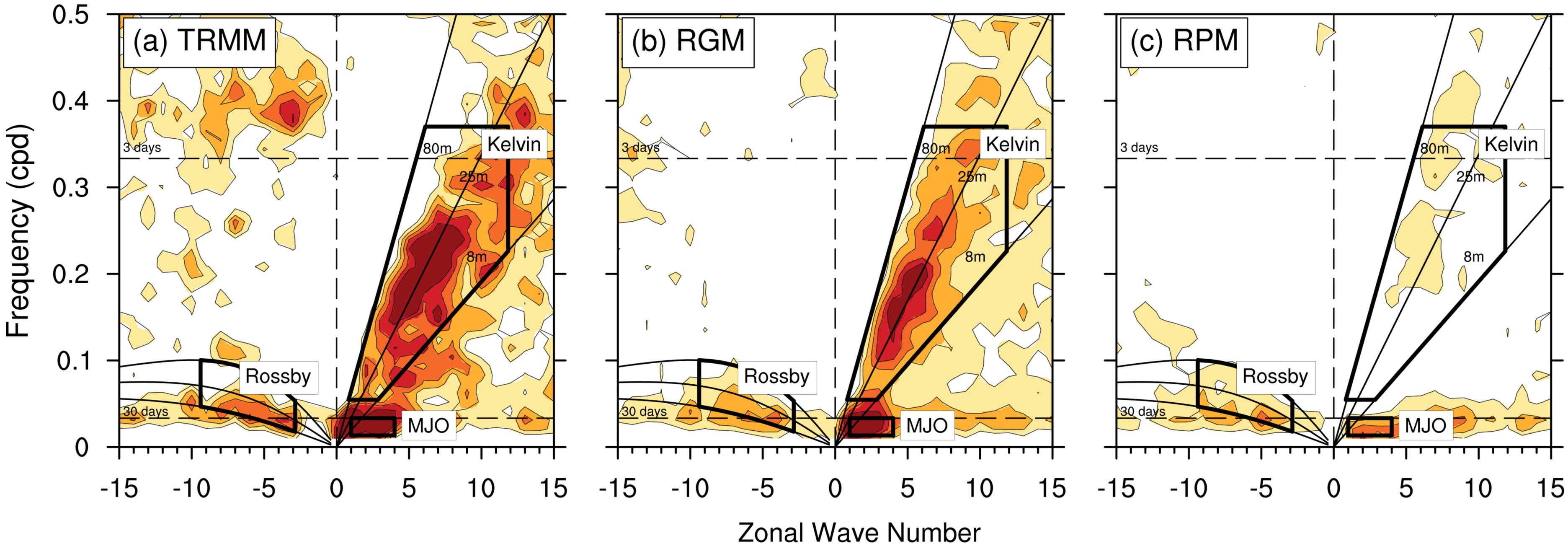
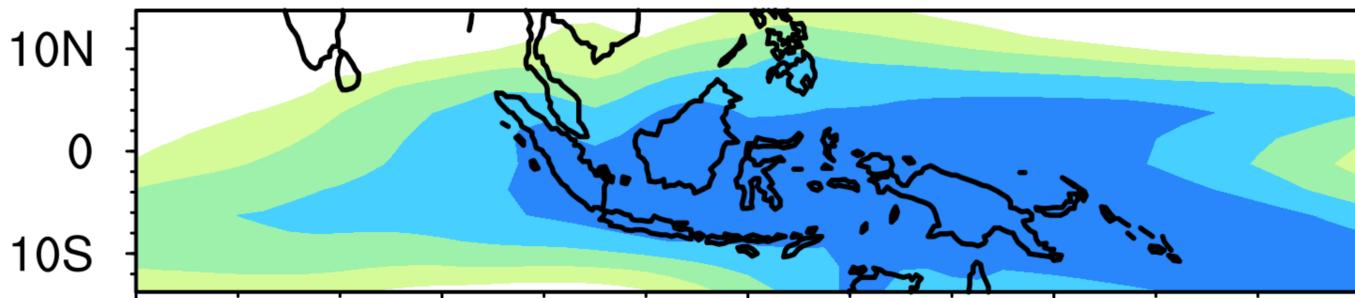


Figure 3.

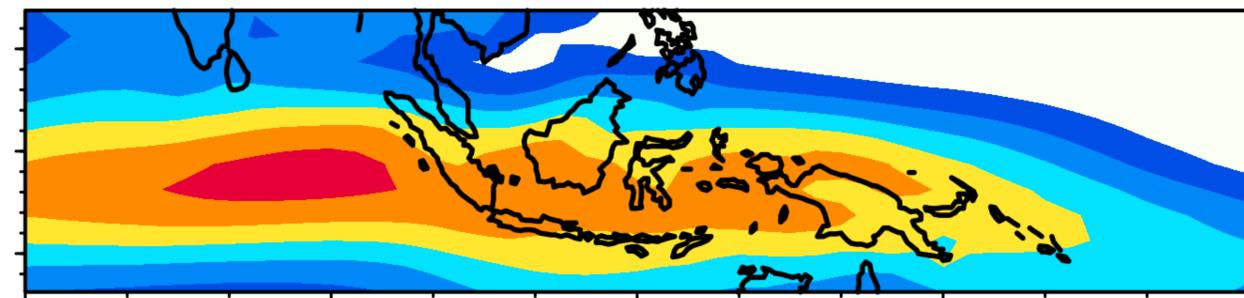
(a) MSE

(b) U_{850}

ERA5

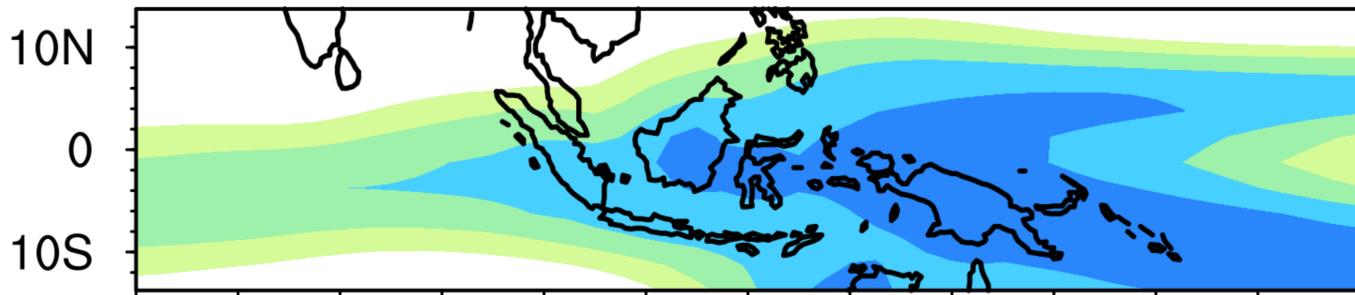


ERA5



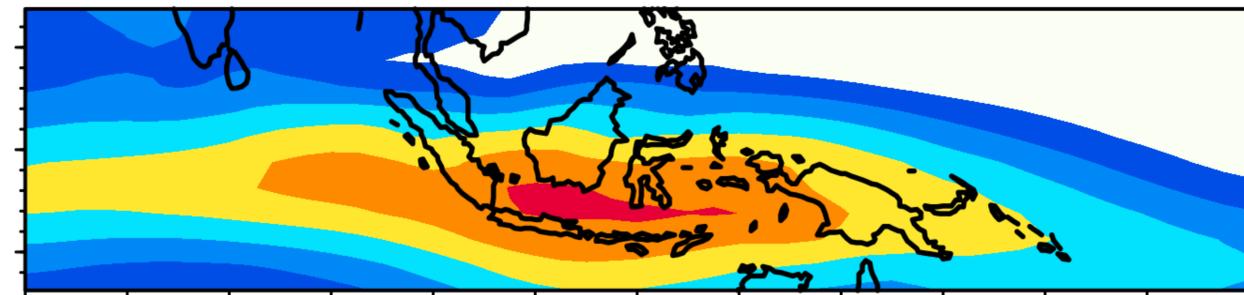
RGM

Cor=0.95 / RMSD=6.62



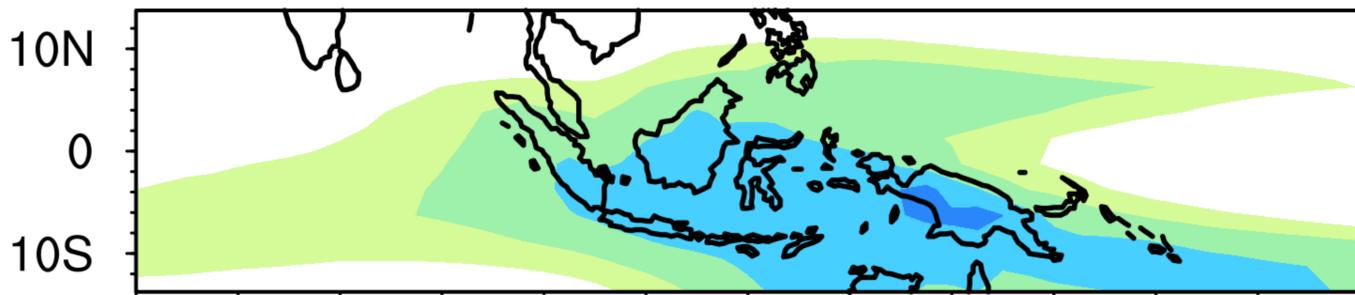
RGM

Cor=0.96 / RMSD=1.26



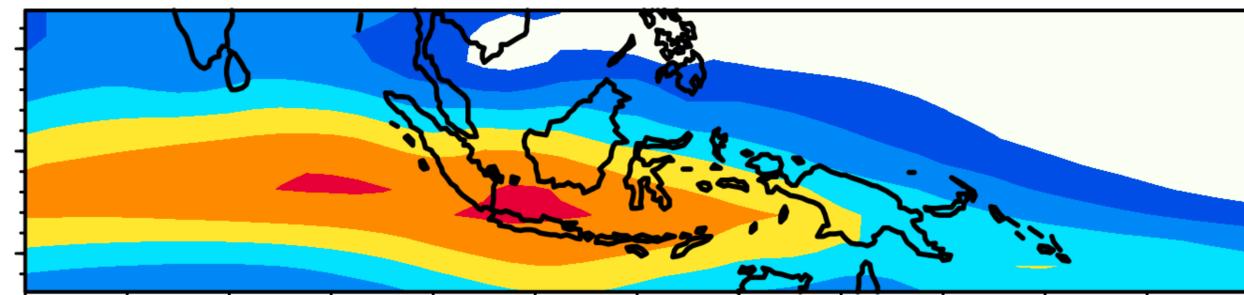
RPM

Cor=0.88 / RMSD=14.51

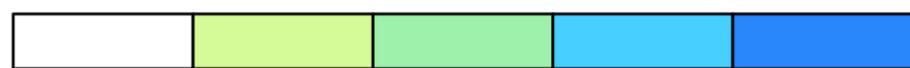


RPM

Cor=0.94 / RMSD=1.49



60E 90E 120E 150E 180



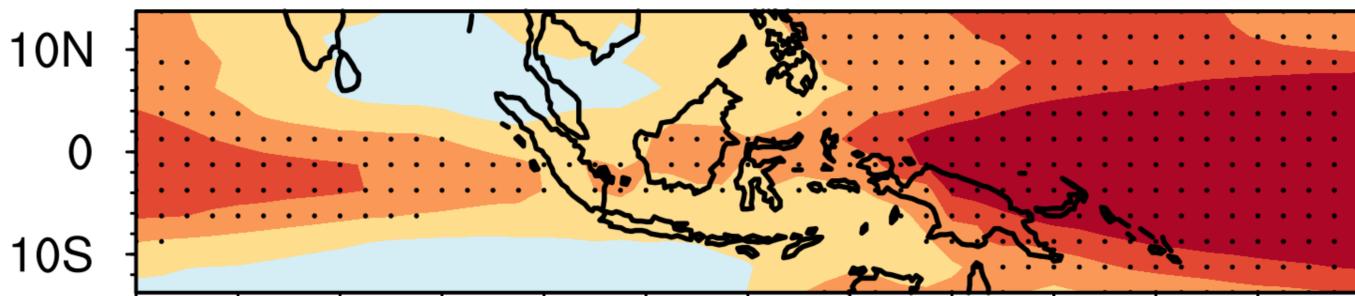
3074 3082 3090 3098

60E 90E 120E 150E 180

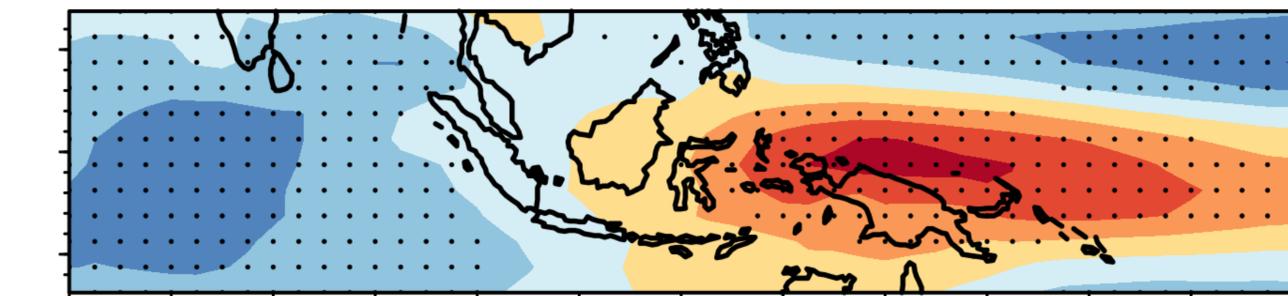


-6 -4 -2 0 2 4

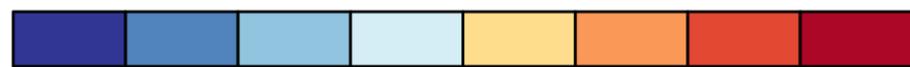
RGM minus RPM



RGM minus RPM

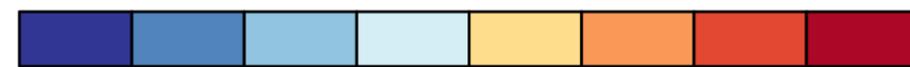


60E 90E 120E 150E 180



-15 -10 -5 0 5 10 15

60E 90E 120E 150E 180



-3 -2 -1 0 1 2 3

Figure 4.

ERA5

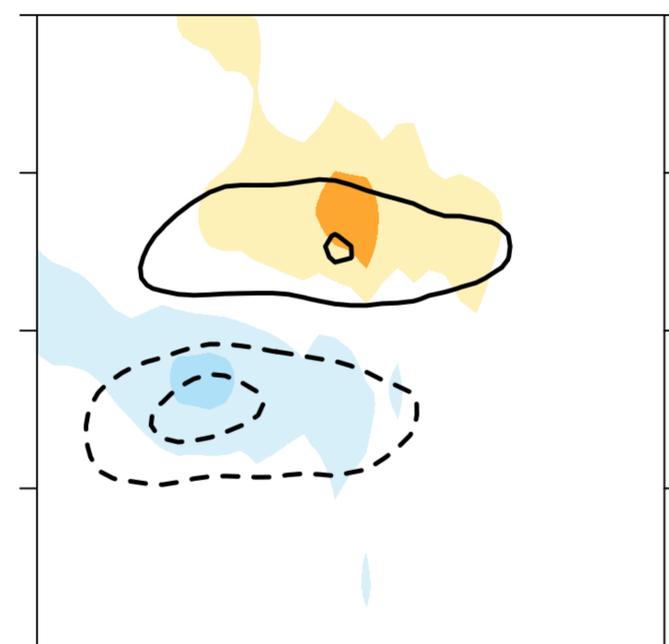
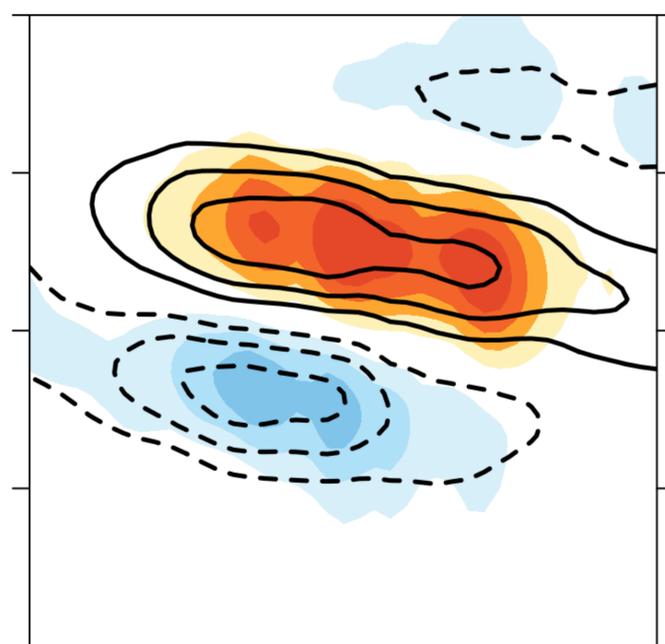
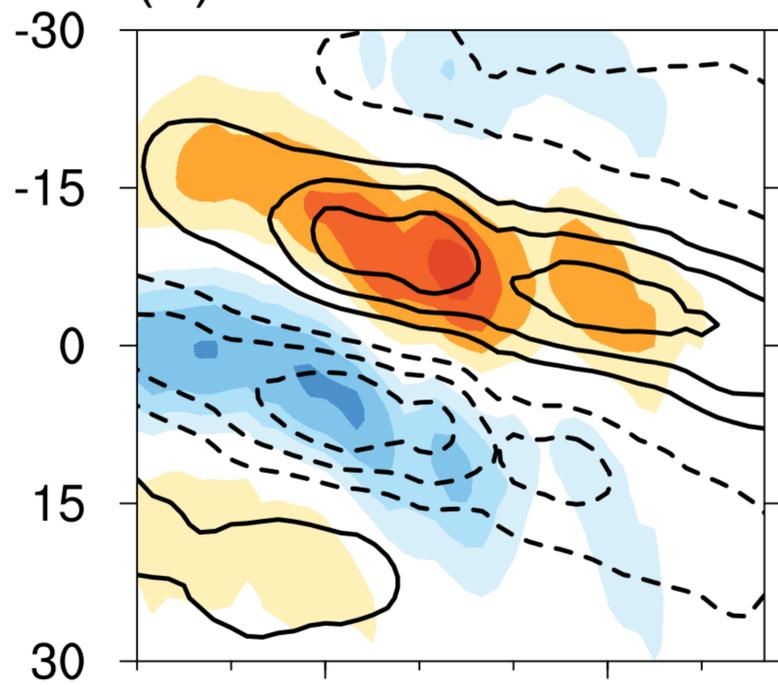
RGM

RPM

(a) -Vdm

-Vdm

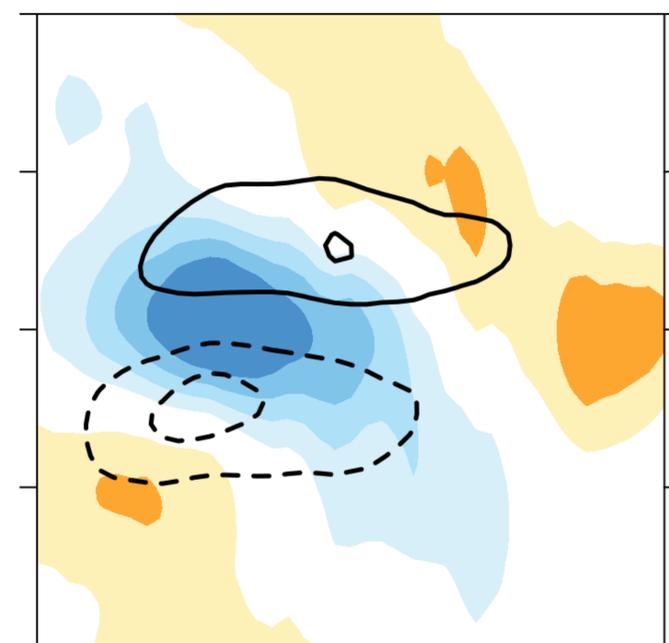
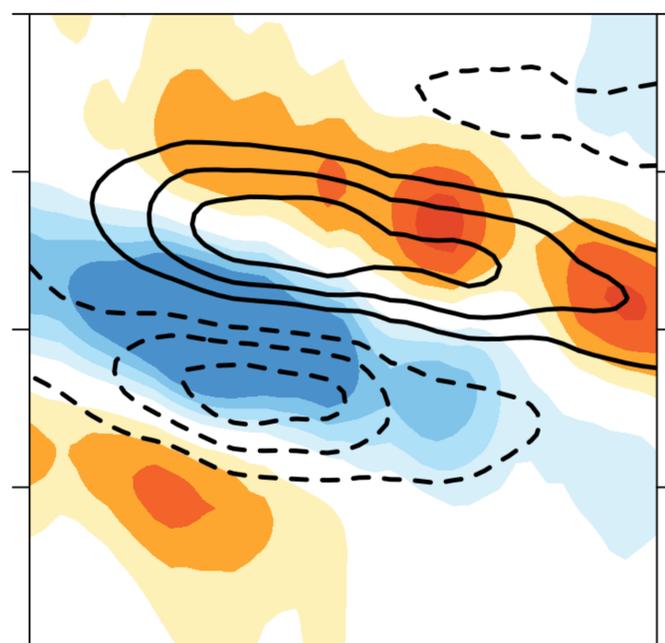
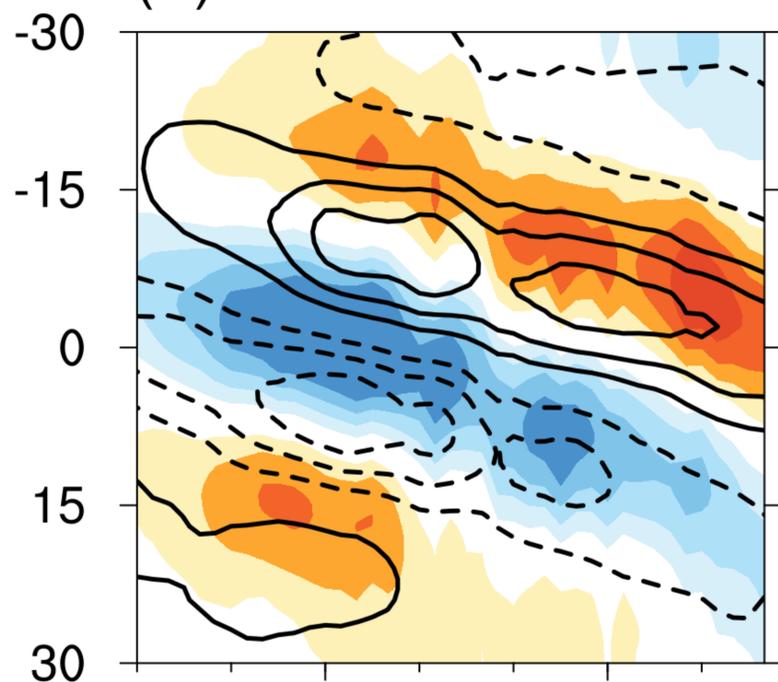
-Vdm



(b) -wdm

-wdm

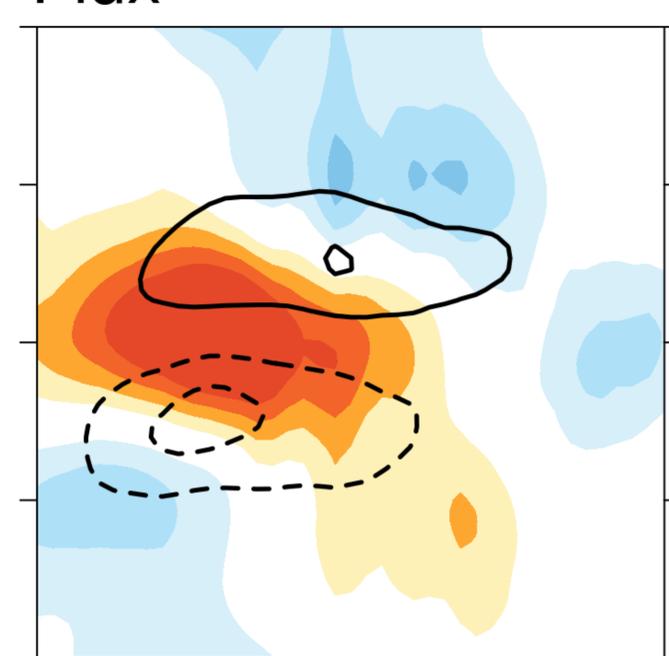
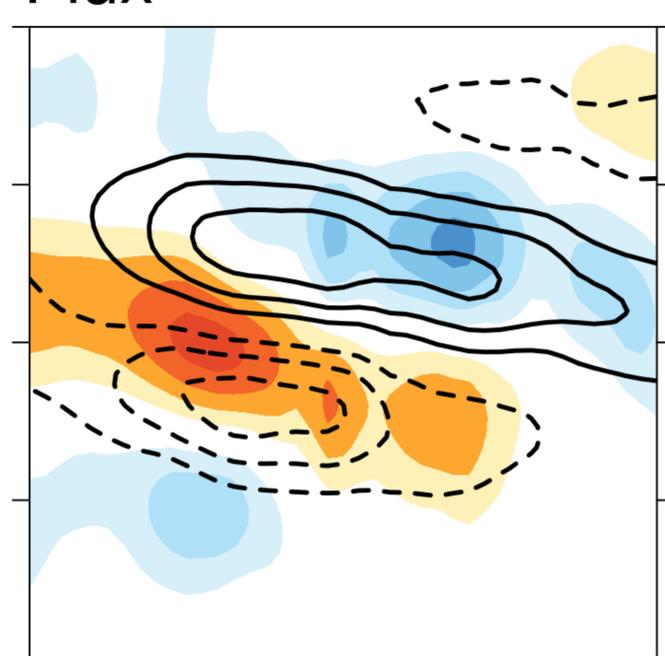
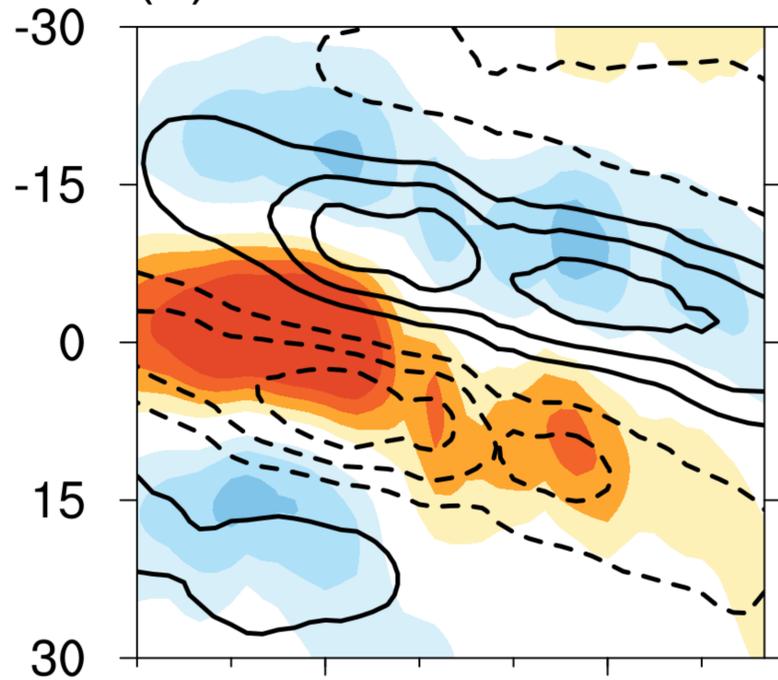
-wdm



(c) Flux

Flux

Flux



90E

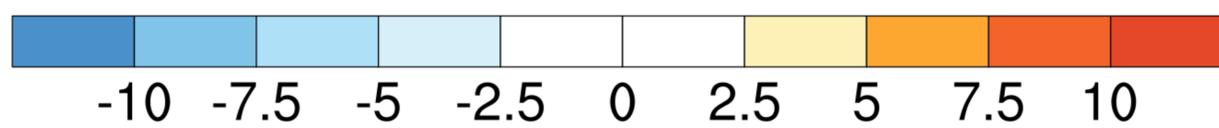
135E

90E

135E

90E

135E



Supplementary Information for “Assessment of the Madden-Julian Oscillation in CMIP6 Models based on Moisture Mode Theory”

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Contents of this file

1. MJO Phase speed: Radon Transform method
2. Parameters of N_{mode}
3. Figures S1, S2, S3, and S4

1. MJO Phase speed: Radon Transform method

Following the previous studies (Yang et al., 2007; Mayta et al., 2021, 2023), the Radon Transform method (Radon, 1917) is adopted to objectively evaluate MJO phase speed based on the longitude-time plot (as shown in Fig. S1). The convective envelopes are projected onto a plate that has an angle θ (0° to 180°) relative to the x -axis (Fig. S2a), and the largest total amplitude $\sum_i P^2(x'_i, \theta)$ will exist in θ_{max} (Fig. S2b). The phase speed can be estimated as

$$c_p = \frac{2\pi a \cos \psi}{360^\circ} \tan(\theta_{max}) \frac{\Delta x}{\Delta t} \quad (1)$$

where a is the earth's radius, and $\frac{2\pi a \cos \psi}{360^\circ}$ is the length of unit degree at latitude ψ . The Δx and Δt are the temporal and spatial ($^\circ$) resolutions of a data grid, respectively. Since the MJO propagation will accelerate when it crosses the Maritime Continent (Rushley et al., 2022), we only consider the range from 60°E to 120°E . For example, the c_p of ERA5 is 3.77 m s^{-1} obtained from $\theta_{max} = 49.5^\circ$ (Fig. S2b). The averaged c_p is $5.3 \pm 3.1 \text{ m s}^{-1}$ for 25 CMIP6 models, and 10 models simulate the $c_p > 5 \text{ m s}^{-1}$ (refer to Fig. S1). These results suggest that the MJO convection has a faster phase speed in model simulation than in the real world, consistent with the previous study of the MJO propagation in CMIP5 models (Ahn et al., 2017).

2. Parameters of N_{mode}

Here, we analyze whether the N_{mode} in the models is also sensitive to its phase speed. The equation of N_{mode} can be rewritten by common logarithm (\log_{10})

$$\log_{10} N_{mode} \approx 2 \log_{10} c_p + \log_{10} \tau + (-\log_{10} \tau_c) + (-2 \log_{10} c) \quad (2)$$

Since the c is consistent ($= 50 \text{ ms}^{-1}$), we neglect the last term. Figure S4 illustrates the scatterplots of the leading three terms in Eq. 2, including the slope (S) and correlation (R) to $\log_{10} N_{mode}$. For the 25 models, MJO phase speed has the highest relation with N_{mode} ($S = 1.37, R = 0.97$; Fig. S3a). Large and small N_{mode} values correspond to relatively fast and slow propagation, respectively. The temporal scale of wave shows a relatively weak effect and a negative correlation to N_{mode} ($S = -0.24, R = -0.79$; Fig. S3b). The contribution of $-\log_{10} \tau$ to $\log_{10} N_{mode}$ is much smaller than other terms ($S = -0.13, R = -0.54$; Fig. S3c). A high sensitivity of N_{mode} to c_p is also found in these 25 CMIP6 models, consistent with the observation in previous studies (e.g., Adames et al., 2019; Adames, 2022).

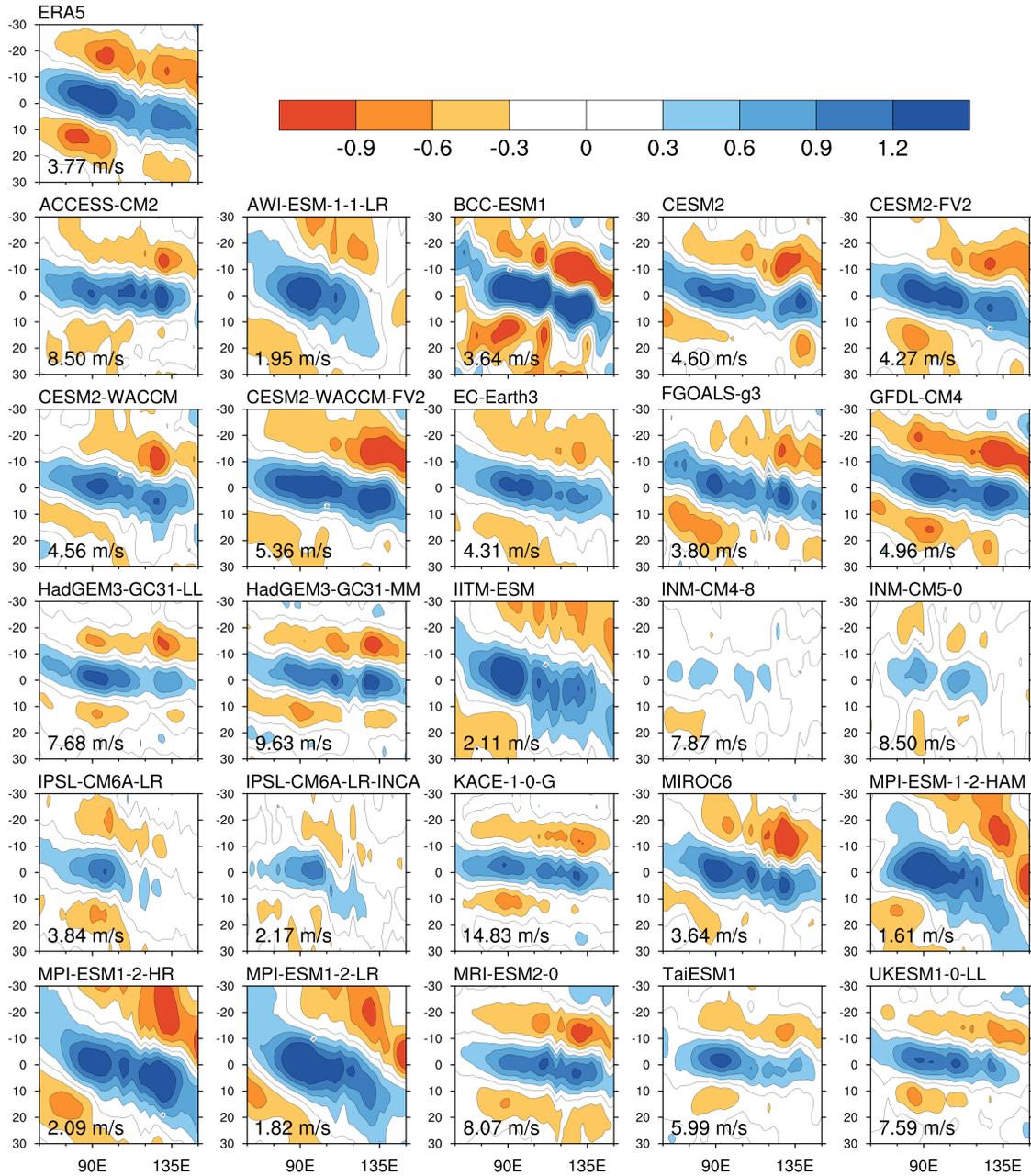


Figure S1. Hovmöller diagram of regressed precipitation averaged between 10°S and 10°N from -30 to 30 days for the ERA5 and 25 CMIP6 models. Values in the bottom-left corner represent the MJO phase speed over 60°E-120°E calculated by using the Radon Transform. The contour interval is 0.3 mm day⁻¹.

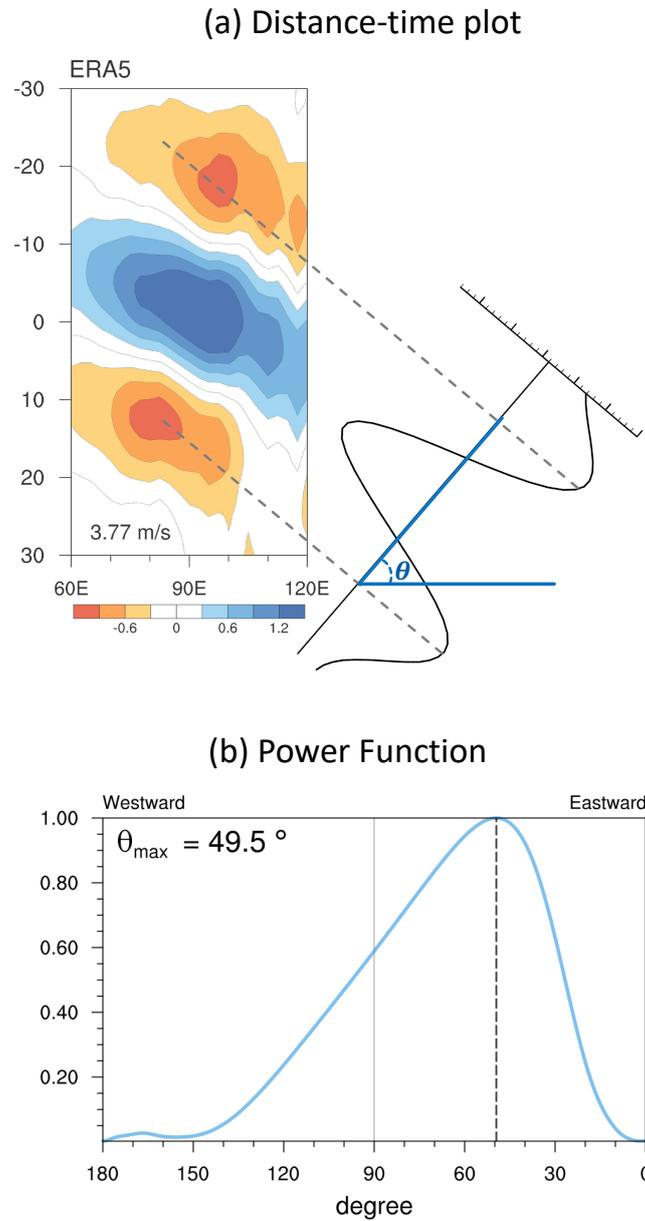


Figure S2. Schematic depicting the process of the Radon Transform method to estimate MJO phase speed in ERA5. (a) Hovmöller diagram projects onto a plate that has an angle θ with the x -axis. (b) The θ that has a maximum total absolute amplitude $\sum_i P^2(x'_i, \theta)$ is defined as θ_{\max} .

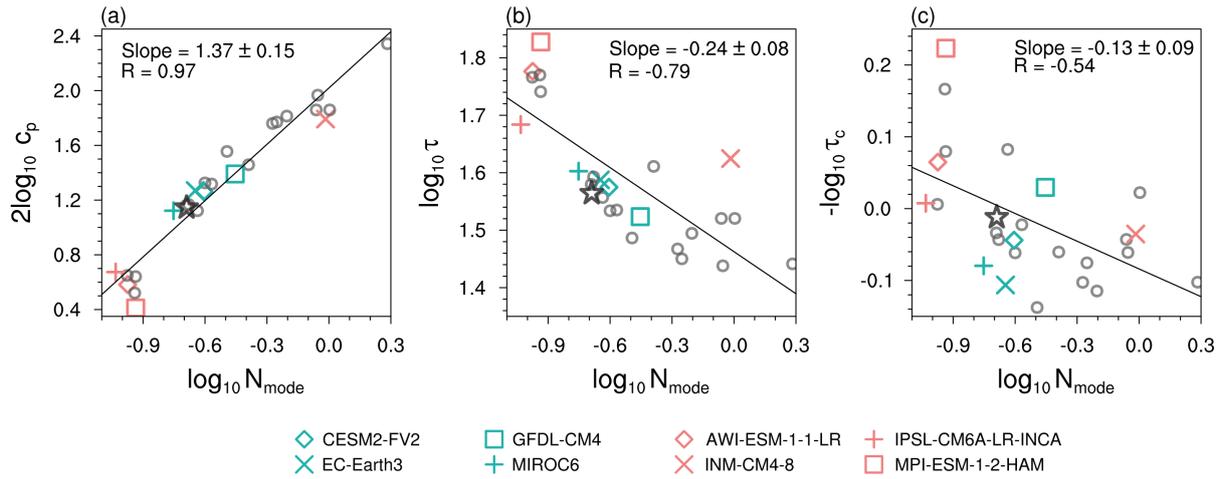


Figure S3. Scatterplot of (a) $2\log_{10} c_p$ (b) $\log_{10} \tau$, and (c) $-\log_{10} \tau_c$ versus $\log_{10} N_{mode}$ for ERA5 (black star) and 25 CMIP6 models. The green and pink markers symbolize four good and four poor models, respectively. Solid lines are calculated by the linear least-squares fitting. The values in the top represent the slope of linear fits and correlation coefficient for 25 model results. All correlation coefficients pass the 95% confidence interval (p -value < 0.05).

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