

# Effects of balanced motions and unbalanced internal waves on steric height in the mid-latitude ocean

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## Abstract

The baroclinic component of the sea surface height, referred to as steric height, is governed by geostrophically balanced motions and unbalanced internal waves, and thus is an essential indicator of ocean interior dynamics. Using yearlong measurements from a mooring array, we assess the distribution of upper-ocean steric height across frequencies and spatial scales of O(1-20 km) in the northeast Atlantic. Temporal decomposition indicates that the two largest contributors to steric height variance are large-scale atmospheric forcing (32.8%) and mesoscale eddies (34.1%), followed by submesoscale motions (15.2%), semidiurnal internal tides (8%), super-tidal variability (6.1%) and near-inertial motions (3.8%). Structure function diagnostics further reveal the seasonality and scale dependence of steric height variance. In winter, steric height is dominated by balanced motions across all resolved scales, whereas in summer, unbalanced internal waves become the leading-order contributor to steric height at scales of a few kilometers.

1        **Effects of balanced motions and unbalanced internal**  
2        **waves on steric height in the mid-latitude ocean**

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8        **Key Points:**

- 9        • The distribution of steric height variance across frequencies and spatial scales of O(1-  
10        20 km) is revealed by yearlong mooring measurements.
- 11        • Balanced motions dominate the upper-ocean steric height variance, and account for  
12        ~83% of the total variance.
- 13        • Internal waves become increasingly important in summer, and are able to dominate  
14        over balanced motions at spatial scales of a few kilometers.

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**Abstract**

The baroclinic component of the sea surface height, referred to as steric height, is governed by geostrophically balanced motions and unbalanced internal waves, and thus is an essential indicator of ocean interior dynamics. Using yearlong measurements from a mooring array, we assess the distribution of upper-ocean steric height across frequencies and spatial scales of  $O(1-20\text{ km})$  in the northeast Atlantic. Temporal decomposition indicates that the two largest contributors to steric height variance are large-scale atmospheric forcing (32.8%) and mesoscale eddies (34.1%), followed by submesoscale motions (15.2%), semidiurnal internal tides (8%), super-tidal variability (6.1%) and near-inertial motions (3.8%). Structure function diagnostics further reveal the seasonality and scale dependence of steric height variance. In winter, steric height is dominated by balanced motions across all resolved scales, whereas in summer, unbalanced internal waves become the leading-order contributor to steric height at scales of a few kilometers.

**Plain Language Summary**

Steric height is the sea surface height component associated with changes in water-column density, and is typically contributed by ocean dynamic processes across a wide range of scales, from the large-scale ocean circulation to the small-scale wave motion. To investigate ocean dynamics using sea surface height data measured by satellites, it is crucial to comprehend the various constituents of steric height variability. In this study, the effects of balanced motions (e.g., eddies and ocean fronts) and unbalanced wave motions (e.g., internal waves) on steric height are quantified based on yearlong moored observations at a mid-latitude ocean site of the northeast Atlantic. Overall, balanced motions account for  $\sim 83\%$  of the upper-ocean steric height variance. Steric height variance also show notable seasonal variations and scale dependence. At spatial scales of  $O(10\text{ km})$ , the steric height is predominately determined by balanced motions throughout the year. By contrast, at spatial scales of  $O(1\text{ km})$ , unbalanced wave motions are the major contributor to steric height in summer whereas balanced motions still dominate in winter. This study provides a quantitative assessment of the effects of balanced motions and unbalanced wave motions on steric height, and provide insights for the exploration of next-generation high-resolution altimetry data.

## 1 Introduction

Sea surface height measurements over the global oceans are now routinely derived from satellite altimetry, and have greatly advanced our understanding of ocean dynamics over the last 30 years (Chelton et al., 2011). However, the conventional nadir radar altimeters can only resolve large-scale and mesoscale variability with horizontal resolutions of  $O(100\text{ km})$  (Ballarotta et al., 2019). Xu and Fu (2011, 2012) examined along-track sea surface height measurements to diagnose the dynamical regimes of the global oceans, and pointed out the necessity of higher resolution satellite altimeter for a more accurate detection of the dynamic characteristics globally. The next-generation wide-swath altimetry missions, such as the Surface Water and Ocean Topography (SWOT) altimeter mission which has been launched in December 2022 and the Chinese “Guanlan” mission which is in the early designing stage, are expected to provide, for the first time, 2D sea level maps globally at spatial scales down to the submesoscale (Chen et al., 2019; Fu & Ubelmann, 2014). SWOT spatial resolution is predicted to be of about 15-50 km depending on the local sea state and measurement noise (Morrow et al., 2019; Wang et al., 2019). One of the scientific challenges associated with increasing spatial resolution is that high-frequency internal waves with spatial scales comparable to submesoscale motions can also be observed, and will behave as “noises” for isolating submesoscale motions given the limited temporal revisit of altimetric satellites (e.g., tens of days). Therefore, it is important to quantitatively assess the respective contributions of submesoscale motions and internal waves to sea surface height.

Submesoscale processes are characterised by spatial scales of  $O(0.1\text{-}10\text{ km})$  and temporal scales of several hours to several days (McWilliams, 2016; Thomas et al., 2008; Callies et al., 2020). Recent studies revealed that submesoscale processes modulate the equilibrium state of the upper ocean through a bi-directional kinetic energy cascade (Ferrari & Wunsch, 2009; McWilliams, 2017; Qiu et al., 2017). Further, submesoscale motions are particularly effective at inducing intense vertical velocities in the upper ocean, which may exceed mesoscale processes by one order of magnitude (Klein et al., 2009; Lévy et al., 2018; Su et al., 2018). Given their importance in oceanic energy cascade and vertical tracer transport, the observation and understanding of submesoscales have become key scientific targets of the SWOT mission and the oceanographic community in general.

Internal waves are propagating disturbances in stably stratified fluids, with gravity acting as the restoring force (Gerkema & Zimmerman, 2008; Sutherland, 2010). The internal wave field can be divided into three components: near-inertial waves, internal tides and the internal-wave continuum. Near-inertial waves typically arise from strong variable winds over the ocean surface, and are expected to have minor signatures on sea surface height (Munk & Phillips, 1968; Fu, 1981). In contrast, internal tides and the internal wave continuum are more substantial contributors to the sea surface height field (Callies & Wu, 2019; Chereskin et al., 2019). Müller et al. (2015) demonstrated that high-resolution global ocean models with tidal and atmospheric forcing are beginning to resolve the internal-wave continuum. Savage et al. (2017) analyzed the components of sea surface height and their spatial distributions based on the outputs of global ocean models with different resolutions, and suggested that the internal wave signal will be high-frequency noise for future high-resolution altimeters and its non-stationary component will be difficult to predict. Based on a high-resolution MITgcm global model, Torres et al. (2018) analyzed contributions of balanced and unbalanced motions in kinetic energy and sea surface height, and found that the relative contributions of two classes of motions to various surface fields are complex and dependent on multiple factors, such as seasonal, geographical and the distribution of low and high eddy kinetic energy regions.

Steric height is the baroclinic component of sea surface height, and is mainly determined by dynamical processes such as balanced (sub)mesoscale motions and unbalanced high-frequency internal waves. A quantitative assessment of how these multiscale processes contribute to the steric height is the foundation for exploring the ocean dynamics from satellite sea surface height measurements (Baker-Yeboah et al., 2009; Gill & Niller, 1973).

98 However, observational studies on steric height variability down to the submesoscales are  
 99 still rare. Very recently, Miao et al. (2021) used temporal filtering to quantitatively assess  
 100 the influence of multiscale dynamic processes on steric height based on temperature/salinity  
 101 time series observed by a single mooring in the South China Sea, and showed that the rel-  
 102 ative contribution of submesoscale motions (7.2%) is smaller than those of diurnal (8.5%)  
 103 and semidiurnal (20.2%) internal tides. They also illustrated that steric height associated  
 104 with mesoscales and submesoscales are stronger in winter than summer but the opposite  
 105 occurs for tidal and super-tidal motions.

106 Similar to the findings reported in Miao et al. (2021), here we will show that there is  
 107 a seasonality in steric height of balanced (sub)mesoscale motions and unbalanced internal  
 108 waves in a mid-latitude ocean. In addition to that, we also examine the scale dependence of  
 109 steric height in the spatial domain based on structure function approaches. We show that  
 110 unbalanced internal waves only dominate over balanced (sub)mesoscale motions in summer  
 111 at spatial scales of a few kilometers in the study region.

## 112 2 Observations and Methods

### 113 2.1 Mooring data

114 The data used in this study were primarily collected from nine bottom-anchored subsur-  
 115 face moorings deployed over the Porcupine Abyssal Plain (48.63-48.75°N, 16.09-16.27°W)  
 116 site in the northeastern Atlantic Ocean for the period September 2012 - September 2013  
 117 (Figure 1a), as part of the OSMOSIS (Ocean Surface Mixing, Ocean Submesoscale Inter-  
 118 action Study) experiment (Buckingham et al., 2016; Yu et al., 2019; Erickson et al., 2020).  
 119 The OSMOSIS mooring site is in an abyssal plain of depth close to 4800 m, and is analogous  
 120 to many open ocean regions far away from western boundaries of ocean basins and from  
 121 complex topography. Nine moorings were arranged in two concentric quadrilaterals with  
 122 side lengths of  $\sim 13$  km (outer cluster) and 1-2 km (inner cluster) around a centrally located  
 123 mooring (Figure 1b), which can concurrently capture mesoscale and submesoscale signals to  
 124 a large extent. The mooring sensors comprised a series of Seabird MicroCAT conductivity-  
 125 temperature-depth (CTD) sensors and Nortek Aquadopp current meters at different depths,  
 126 spanning the approximate depth interval 30-530 m. The present study predominately uses  
 127 data from the CTDs.

128 The mooring measurements captured most of the pycnocline plus part of the ocean  
 129 interior throughout the year, and most of the mixed layer during winter months. Here, the  
 130 seasons are defined as follows: winter (December to April) and summer (June to August).  
 131 Temperature, salinity and pressure measurements were recorded with intervals every 5 min-  
 132 utes. We linearly interpolate measurements of temperature and salinity onto surfaces of  
 133 constant depth at 10 m intervals between depths of 50 m and 520 m for each mooring, and  
 134 onto uniform 10 minutes intervals between 5 September 2012 and 5 September 2013. Further  
 135 information regarding the OSMOSIS moorings, such as the detailed distribution of moored  
 136 instruments and associated observational uncertainties, can be found in the work by Yu et  
 137 al. (2019) or Naveira Garabato et al. (2022). Furthermore, the mooring measurements were  
 138 complemented by hydrographic observations acquired by two ocean gliders that navigated  
 139 in a bow-tie pattern across the mooring array for the entire sampling period (Damerell et  
 140 al., 2016; Thompson et al., 2016).

### 141 2.2 Steric height calculation

142 The steric height  $\xi$  is calculated as the integral of specific volume anomaly  $\delta$  from the  
 143 reference pressure  $p_{ref}$  to the pressure  $p$ , and is given by

$$144 \quad \xi = \frac{1}{g} \int_{p_{ref}}^p \delta dp, \quad (1)$$

145 where  $\delta = \alpha(S, T, p) - \alpha_0$ ,  $\alpha_0 = \alpha(35, 0, p)$  is the specific volume  $\alpha$  at local pressure with  
 146  $S = 35$  psu and  $T = 0^\circ\text{C}$ ,  $g$  is the gravitational acceleration. In this study, the reference  
 147 pressure  $p_{ref}$  is chosen to be 480 db, above which the nine OSMOSIS moorings all have  
 148 CTD measurements.

### 149 2.3 Band pass filtering

150 One way to isolate contributions of each dynamical process to steric height is by band  
 151 pass filtering applied in the frequency domain. Here, the frequency bands of interest are  
 152 defined as large-scale forcing (1/90–1/30 cpd), mesoscale (1/30–1/5 cpd), submesoscale  
 153 (1/5–1 cpd, following Naveira Garabato et al. (2022)), near-inertial (0.9–1.1 $f$ ), semidiurnal  
 154 (1.9–2.1 cpd) and supertidal band (>2.1 cpd), where  $f = 2\Omega \sin \phi$  is the inertial frequency  
 155 (with  $\Omega$  as the Earth’s angular velocity and  $\phi$  as latitude). A forth-order Butterworth filter  
 156 with a cutoff of these frequency bands are respectively applied on steric height time series  
 157 at 80 db, and we then compute the root mean square of each frequency band.

### 158 2.4 Second-order structure functions

159 We employ the second-order structure function approach to examine steric height vari-  
 160 ance across spatial scales. The second-order structure function for steric height between a  
 161 given location  $\mathbf{x}$  and another location separated from  $\mathbf{x}$  by the distance  $\mathbf{r}$  is defined as

$$162 D_\xi(\mathbf{r}, \mathbf{x}) = \overline{[\xi(\mathbf{x}) - \xi(\mathbf{x} + \mathbf{r})]^2}, \quad (2)$$

163 where the overbar denotes temporal means.

164 If a given homogeneous, isotropic turbulence spectrum (of energy or tracer variance)  
 165 has power-law behavior over a range of wavenumber between the energy injection and dis-  
 166 sipation scales, then a related scaling law for the structure function is expected (Webb,  
 167 1964). Following McCaffrey et al. (2015), the relationship between the spectral slope  $\lambda$  and  
 168 structure function slope  $\alpha$  is

$$169 \lambda = -\alpha - 1. \quad (3)$$

170 Based on nine distinct mooring sites and their separations, there are 36 combinations  
 171 of mooring pairs, covering spatial scales from 18.7 km down to 1.3 km. We also utilize a  
 172 low-pass filter (with a cutoff period at the local inertial period  $\sim 16$  hours) to decompose  
 173 steric height into sub-inertial and super-inertial components,  $\xi(z, t) = \xi(z, t)_{sub} + \xi(z, t)_{sup}$ .  
 174 For each component, we repeat the estimates of the second-order structure function.

### 175 2.5 Frequency-resolved structure functions

176 The frequency-solved structure function approach was developed by Callies et al. (2020),  
 177 and this approach allows for the assessment of steric height variance as a function of temporal  
 178 and spatial scales. To apply this approach, we first temporally detrend steric height time  
 179 series at nine moorings, then a Hann window is applied to each time series to minimize  
 180 spectral leakage. Subsequently, we calculate their temporal Fourier transform, and lastly  
 181 compute spatial second-order structure functions estimates. Specifically, at each mooring  
 182 location  $\mathbf{x}$ , the temporal Fourier transform of the steric height  $\xi$  is given by:

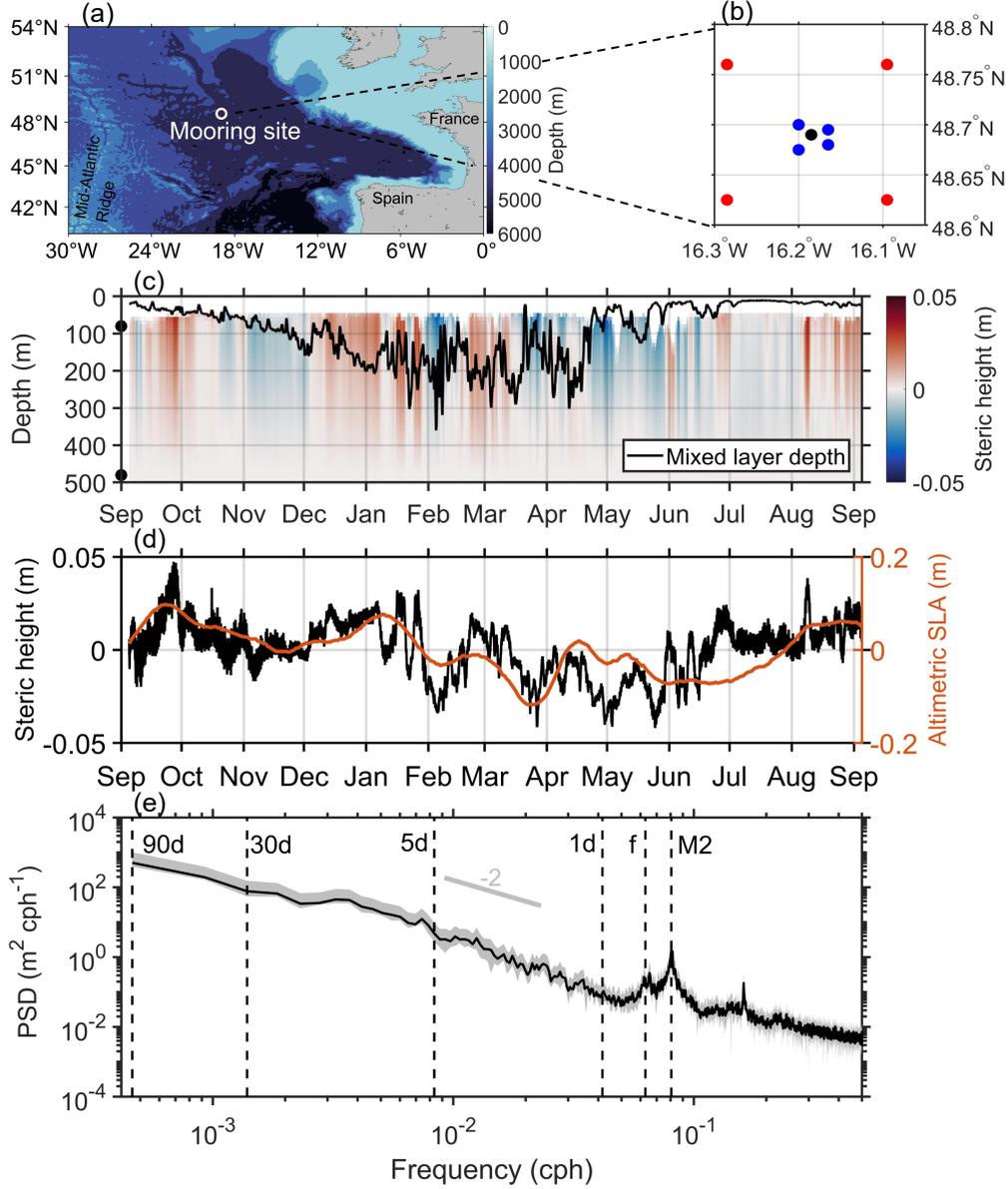
$$183 \tilde{\xi}(\mathbf{x}, \omega) = \int_{-\infty}^{+\infty} \xi(\mathbf{x}, t) e^{-i\omega t} dt. \quad (4)$$

184 For mooring pairs at locations  $\mathbf{x}$  and  $\mathbf{x} + \mathbf{r}$ , the frequency-resolved structure function  
 185 can be calculated by:

$$186 \Delta^\xi(\mathbf{r}, \omega) = \frac{1}{2} \langle \left| \tilde{\xi}(\mathbf{x} + \mathbf{r}, \omega) - \tilde{\xi}(\mathbf{x}, \omega) \right|^2 \rangle, \quad (5)$$

187  
188

where  $\omega$  is frequency,  $\Delta^\xi(\mathbf{r}, \omega)$  is the frequency-resolved second-order structure function as a function of  $\mathbf{r}$  and  $\omega$ , the angle brackets denote the average over nearby spatial distances.



**Figure 1.** (a) OSMOSIS study region in the northeast Atlantic, with bathymetry shown in the colormap on the right. The white circle denotes the location of the OSMOSIS mooring array. (b) Locations of inner (blue circles), outer (red circles) and center (black circle) moorings. (c) Depth-resolved time series of steric height for the central mooring. The two black dots on the y axis respectively indicate the depths corresponding to 80 db and 480 db, and the black line represents the mixed layer depth. (d) Time series of steric height at 80 db from the central mooring (black) and satellite altimetry sea level anomaly (red) from September 2012 to September 2013. (e) Power spectral densities of the steric height. Gray shading indicates 95% confidence intervals, and slope of  $-2$  is shown for reference. The vertical dashed lines indicate the 90 days, 30 days, 5 days, 1 day, near-inertial frequency (1/16 cph) and semidiurnal frequency (1/12.42 cph), respectively.

### 3 Results

#### 3.1 Temporal decomposition of steric height

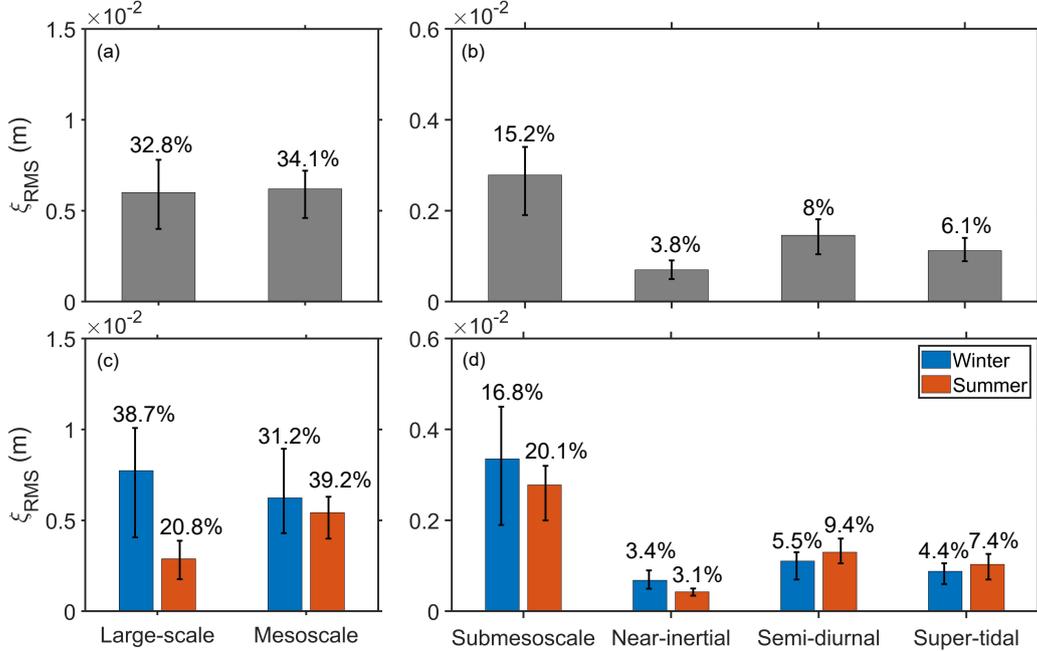
The time evolution of steric height over the annual cycle is shown in Figure 1c, contextualized with the glider-based mixed layer depth. The mixed layer depth is calculated from coincident glider data using a threshold value of potential density increase ( $\Delta\rho = 0.03 \text{ kg m}^{-3}$ ) from a near-surface value at 10 m (Damerell et al., 2016). Steric height is characteristically intensified near the surface and ranges from  $-0.05$  to  $0.05$  m. A closer look at steric height time series at 80 db reveals more details of its temporal variability (Figure 1d). Seasonal heating/cooling, which directly affects near-surface density, is one of the main factors determining the annual cycle of steric height in this region. This large-scale atmospheric forcing leads to predominantly positive values of steric height during non-winter months (such as September-November and July-September) and negative values during winter months (such as January-April; also see Figure S1). Steric height also displays profound high-frequency signals, especially in non-winter months when mixed layer is shallow and vertical stratification is strong. The wave signature in the steric height field is further confirmed by the frequency spectral analysis in Figure 1e, in which the most noticeable feature is the presence of a spectral peak at the semidiurnal tidal frequency that can only be accounted for by internal tides. In contrast, low-frequency eddy-like signals are more evident in winter months when mixed layer is deep and vertical stratification is considerably reduced.

Mooring-derived steric height shows a reasonably consistent trend with the sea level anomaly (SLA) observed by altimetry, although the former is smaller by a factor of 3-4 than the latter. Note that substantial corrections, such as dynamic atmospheric correction and ocean tide correction, have already been applied in the altimetry data, and thus the SLA is expected to be primarily contributed by steric height. The discrepancy in the mooring-based steric height and altimetry SLA likely stems from the missing contributions of steric height in the deep ocean (i.e., below 480 db) by the moorings, as in the comparison between seal-based upper-ocean steric height and altimetry SLA in the Southern Ocean (Siegelman et al., 2020). Additionally, Wang et al. (2018) compared different reference levels for the calculation of steric height to the full-depth steric height based on the output of a high-resolution model, and found the upper-ocean dominance of the steric height component on the sea surface height (accounting for  $>70\%$  of full-depth steric height for a reference level of 580 m in their case).

Frequency band integrated steric height variance are compared in Figure 2. Over the whole year, largest contributions to steric height variability are associated with large-scale atmospheric forcing and mesoscale eddies, constituting 32.8% and 34.1% of the total steric height variance, respectively. Submesoscale motions are the third contributor to the steric height, and account for about 15.2% of the total variance. By contrast, three unbalanced internal wave components (near-inertial, semidiurnal and super-tidal waves) all contribute less than submesoscale motions, making up 3.8%, 8% and 6.1%, respectively. Note that the steric height variance in those internal wave bands may, at least in part, arise from Doppler shifting of submesoscale motions by mesoscale eddies and/or the presence of background spectrum, which would lead to an overestimate of the relative contributions of high-frequency steric height variance.

We next examine the seasonality of steric height across frequencies. The large-scale atmospheric forcing, mesoscale eddies, submesoscale motions and near-inertial waves all produce higher root-mean-square values of steric height in winter compared to summer. The annual-mean surface heat flux is approximately  $-45 \text{ W m}^{-2}$ , indicating a stronger winter cooling compared to summer heating (Yu et al., 2019). This asymmetry between heating and cooling may explain the seasonal difference of steric height variance at lowest frequencies. The seasonality of steric height at mesoscale and submesoscale bands are also expected, because mesoscale eddies and the associated submesoscale motions are found to be more active in winter than in summer (Buckingham et al., 2016; Thompson et al., 2016).

241 Furthermore, the enhancement of near-inertial steric height variance in winter conforms to  
 242 expectations from the seasonality of near-inertial energy generated by surface winds (Yu et  
 243 al., 2022). However, an opposite seasonality is seen at semidiurnal and super-tidal bands,  
 244 where steric height variance is much stronger in summer than in winter. This is mainly  
 245 caused by stronger summer vertical stratification in the upper ocean (Callies et al., 2020;  
 246 Rocha, Gille, et al., 2016).



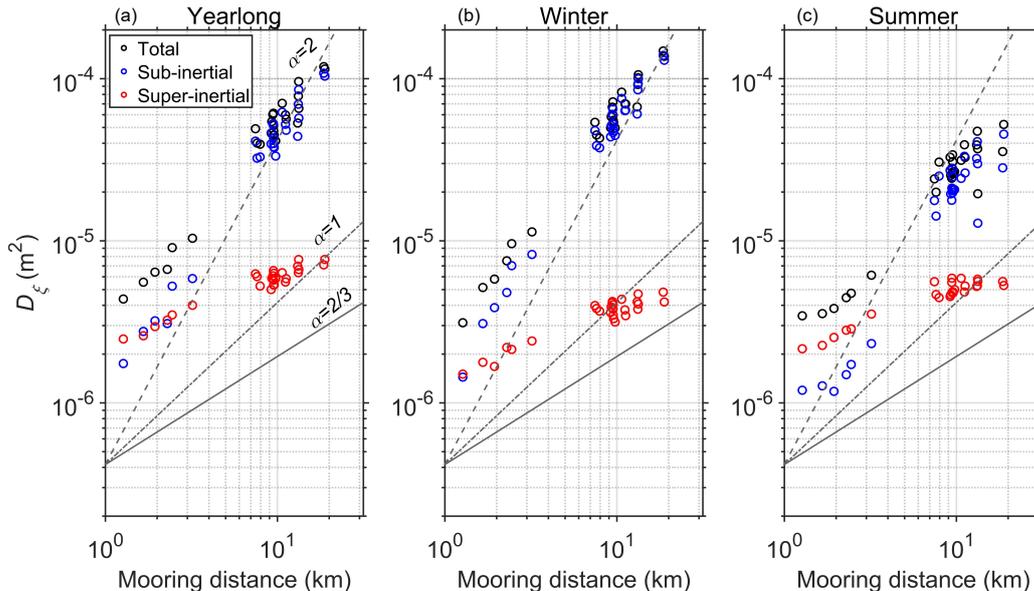
**Figure 2.** (a-b) Annual-mean and (c-d) seasonal-mean of the root-mean-square steric height at 80 db and for different frequency bands. Fraction of each component to the total steric height is marked. Black vertical bars illustrate the 95% confidence intervals estimated using a bootstrap approach.

### 247 3.2 Steric height variance across spatial scales

248 To assess the distribution of steric height variance across spatial scales, we compute  
 249 second-order structure functions of steric height at 80 db for yearlong, winter and summer,  
 250 respectively (Figure 3). Total steric height variance (black circles) exhibits a seasonality  
 251 across spatial scales of  $O(1-20 \text{ km})$ , elevated in winter and reduced in summer. The slope  
 252 for steric height structure functions in winter is comparable to 2, corresponding to a spectral  
 253 slope of  $k^{-3}$ , where  $k$  is the horizontal wavenumber. Note that the spectral slope of kinetic  
 254 energy estimated from the same observations data is between  $k^{-1.6}$  and  $k^{-2}$  (Erickson et al.,  
 255 2020), which indicates a sea surface height spectral slope of  $k^{-3.6}$  and  $k^{-4}$ , slightly steeper  
 256 than the slope  $k^{-3}$  found here. By contrast, the slope of steric height structure functions in  
 257 summer becomes flatter, indicating that steric height variance are reduced at larger scales  
 258 but enhanced at smaller scales.

259 The yearlong record also allows us to compare the sub-inertial (blue circles) and super-  
 260 inertial (red circles) steric height structure functions. In winter, the steric height variance is  
 261 predominately determined by its sub-inertial component, which also follows closely a spectral  
 262 slope of  $k^{-3}$ . The super-inertial component, however, contributes little to the total steric  
 263 height variance with a structure function slope near  $2/3$ , corresponding to a spectral slope

264 of  $k^{-5/3}$ . This slope is close to  $k^{-2}$  for sea level variance predicted by the GM spectrum  
 265 (Callies & Wu, 2019). Sub-inertial steric height structure functions at around 10 km are  
 266 one order of magnitude larger than the super-inertial ones, but their differences decrease  
 267 towards smaller scales. Notably, the sub-inertial and super-inertial steric height structure  
 268 functions are comparable in magnitude at smallest scales. However, a somewhat different  
 269 picture is seen in summer. While the sub-inertial steric height variance still dominates over  
 270 the super-inertial variance at spatial scales of order 10 km, the super-inertial component  
 271 is considerably larger than the sub-inertial component at spatial scales of O(1 km). This  
 272 suggests that high-frequency internal waves dominate over balanced (sub)mesoscale motions  
 273 at scales of O(1 km) in summer.

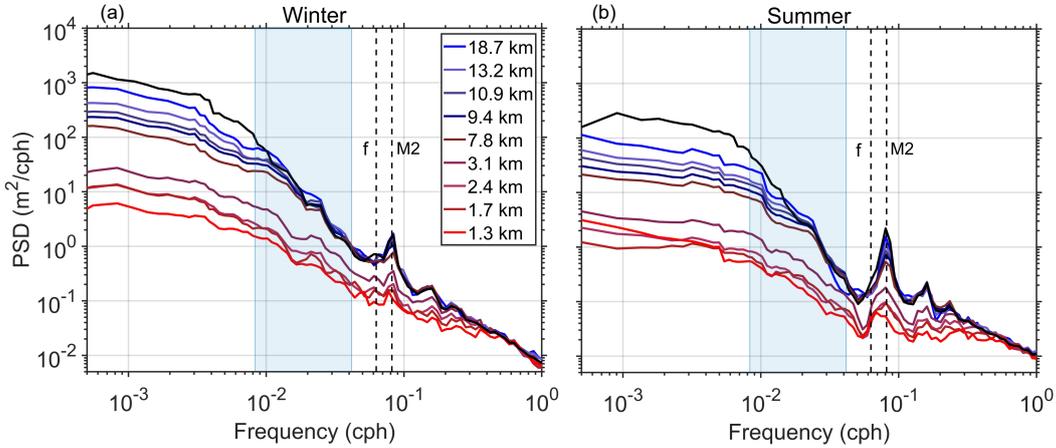


**Figure 3.** The second-order structure functions of steric height at 80 db during (a) the whole year, (b) winter and (c) summer. The black, blue and red circles denotes all, sub-inertial (>16 hours) and super-inertial (<16 hours) motions, respectively. Black lines provide a reference for 2, 1 and 2/3 power-law slopes.

274 The frequency-resolved structure functions of steric height provide useful insights into  
 275 the link between temporal and spatial scales resolved by the mooring array (Figure 4).  
 276 At lowest frequencies, the frequency spectrum (black curve) is considerably larger than  
 277 the frequency-resolved structure functions (colored curves) at all separation distances for  
 278 both winter and summer. This indicates that the steric height variance at low frequen-  
 279 cies is primarily controlled by spatial scales larger than the largest mooring sampled scale  
 280  $2r=37.4$  km. Within the submesoscale band, the frequency spectrum is notably larger in  
 281 winter than in summer, and the frequency-resolved structure functions for the larger sep-  
 282 arations ( $r=7.8$ – $18.7$  km) converge to the frequency spectrum. At highest frequencies, the  
 283 frequency-resolved structure functions converges to the frequency spectrum at all separa-  
 284 tions, indicating that at these periods, the steric height variance becomes decorrelated across  
 285 all mooring pairs.

286 Comparing our work to Callies et al. (2020), who carried out frequency-resolved struc-  
 287 ture functions of horizontal velocity, there are two notable distinctions. Firstly, at the  
 288 near-inertial frequency, the frequency-resolved structure functions of steric height vari-  
 289 ance do not exhibit peaks, primarily due to the negligible signature of near-inertial motion on the

290 steric height variance. Secondly, at semidiurnal frequency, the frequency-resolved structure  
 291 functions of steric height at large separations (e.g.,  $r=7.8-18.7$  km) are found to more closely  
 292 align with the frequency spectrum, unlike the significant gap showed in velocity frequency-  
 293 resolved structure functions. We speculate this difference is partially attributed to the  
 294 presence of large-scale barotropic tide signals in the velocity measurements while absence  
 295 in steric height estimates. To confirm this hypothesis, we further compared the frequency-  
 296 resolved structure functions and frequency spectrum by incorporating the barotropic semi-  
 297 diurnal tide sea level component into the steric height time series, and found a discernible  
 298 increase of spectral peak at the semidiurnal frequency but not for structure functions (Figure  
 299 S2). Another reason may be that low-mode baroclinic tides, which are featured with larger  
 300 spatial scales compared to high-mode components, are not fully captured due to a shallow  
 301 reference level of 480 db. Sea level for low-mode semidiurnal tides predicted by High Resolu-  
 302 tion Empirical Tides (HRET; Zaron et al., 2022) yield a root-mean-square value of 0.28 cm  
 303 at the mooring site, nearly a factor of 2 larger than 0.15 cm estimated from the semi-diurnal  
 304 band of mooring-derived steric height time series. This suggests that a fair amount of steric  
 305 height variability at semi-tidal frequencies, most likely associated with large-scale low-mode  
 306 components, is being missed due to the limited vertical extent of mooring observations.



**Figure 4.** Frequency-resolved structure functions of steric height at 80 db for (a) winter and (b) summer. The black lines show the frequency spectrum and the colored lines show the frequency-resolved structure functions for the separations  $r$  given in the legend. The light blue area represents the submesoscale range, corresponding to 1-5 days. The two vertical lines denote the near-inertial frequency (1/16 cph) and semidiurnal frequency (1/12.42 cph), respectively.

## 307 4 Conclusion

308 In this study, we provide observational insights, for the first time, into seasonal to sub-  
 309 mesoscale steric height variability down to scales of 1 km at a typical mid-latitude ocean site  
 310 of northeast Atlantic. Throughout the year, large-scale atmospheric forcing and mesoscale  
 311 eddies are the two largest contributors to steric height, with their respective relative contri-  
 312 butions of 32.8% and 34.1%, followed by submesoscale motions (15.2%), semidiurnal (8%),  
 313 supertidal (6.1%) and near-inertial (3.8%) waves. Another important finding is that contri-  
 314 butions of balanced motions and unbalanced internal waves to steric height display a strong  
 315 seasonal cycle and scale dependence. Low-frequency balanced motions largely dominate  
 316 the upper-ocean steric height variance in winter and at scales of  $O(10)$  km. In contrast,  
 317 high-frequency unbalanced internal waves become increasingly important in summer, and  
 318 are able to dominate over low-frequency balanced motions at scales of  $O(1)$  km).

319 The study by Miao et al. (2021) found some similar results to our findings regarding  
320 the variations in steric height between seasons caused by multiscale dynamic processes. One  
321 notable difference is the relative contribution of semidiurnal tides on steric height was found  
322 to be significantly greater than that of submesoscale processes in the South China Sea,  
323 regardless of the season. By contrast, our analysis found that the relative contribution of  
324 semidiurnal tides to the steric height is lower than that of submesoscale motions and that  
325 reported in Miao et al. (2021). It should be emphasized that the South China Sea is well  
326 known for its strong internal tides radiated from the Luzon Strait (Alford et al., 2015),  
327 while the OSMOSIS observational area is representative of mid-latitude open ocean regions  
328 characterised by moderate eddy kinetic energy and weak mean flow and internal waves.

329 Our results add observational assessment on the steric height variance in the subme-  
330 soscale range, which is crucial to interpret future altimetric high-resolution sea surface height  
331 maps. Apart from steric height, the sea surface height also includes contributions from the  
332 bottom pressure anomaly (mainly caused by barotropic tides in the open ocean) and the  
333 atmospheric pressure loading (i.e., the inverted barometer effect). A key outcome of the  
334 satellite-derived sea surface height is to estimate surface velocity via geostrophy. To this end,  
335 one needs to filter out the signals associated with bottom pressure anomaly and atmospheric  
336 pressure loading, and then subtract unbalanced internal wave signals. The barotropic tides  
337 and inverted barometer effect are large-scale signals in the sea surface height, and can be  
338 largely removed by spatial filtering. Such corrections have been widely applied to present  
339 altimetry measurements. It has been recently suggested that sea surface height variance  
340 spectra associated with balanced flow drop off steeply with wavelength and internal tides  
341 are likely the main factor affecting the accuracy to infer submesoscale balanced flow at  
342 the smallest scales resolved by SWOT (Callies & Wu, 2019; de Marez et al., 2023). For  
343 the SWOT mission, spatial filtering is likely the practical approach to further mitigate the  
344 effects of fast unbalanced variability given its long repeat sampling cycle (21 days).

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### 352 **Data Availability Statement**

353 All OSMOSIS mooring data are freely available, and are archived at the British Oceanographic  
354 Data Centre. Moored observations can be obtained from [https://www.bodc.ac](https://www.bodc.ac.uk/projects/data_management/uk/osmosis/)  
355 [.uk/projects/data\\_management/uk/osmosis/](https://www.bodc.ac.uk/projects/data_management/uk/osmosis/). Altimeter data were obtained from [https://](https://data.marine.copernicus.eu/product/)  
356 [data.marine.copernicus.eu/product/](https://data.marine.copernicus.eu/product/).

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1        **Effects of balanced motions and unbalanced internal**  
2        **waves on steric height in the mid-latitude ocean**

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8        **Key Points:**

- 9        • The distribution of steric height variance across frequencies and spatial scales of O(1-  
10        20 km) is revealed by yearlong mooring measurements.
- 11        • Balanced motions dominate the upper-ocean steric height variance, and account for  
12        ~83% of the total variance.
- 13        • Internal waves become increasingly important in summer, and are able to dominate  
14        over balanced motions at spatial scales of a few kilometers.

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**Abstract**

The baroclinic component of the sea surface height, referred to as steric height, is governed by geostrophically balanced motions and unbalanced internal waves, and thus is an essential indicator of ocean interior dynamics. Using yearlong measurements from a mooring array, we assess the distribution of upper-ocean steric height across frequencies and spatial scales of  $O(1-20\text{ km})$  in the northeast Atlantic. Temporal decomposition indicates that the two largest contributors to steric height variance are large-scale atmospheric forcing (32.8%) and mesoscale eddies (34.1%), followed by submesoscale motions (15.2%), semidiurnal internal tides (8%), super-tidal variability (6.1%) and near-inertial motions (3.8%). Structure function diagnostics further reveal the seasonality and scale dependence of steric height variance. In winter, steric height is dominated by balanced motions across all resolved scales, whereas in summer, unbalanced internal waves become the leading-order contributor to steric height at scales of a few kilometers.

**Plain Language Summary**

Steric height is the sea surface height component associated with changes in water-column density, and is typically contributed by ocean dynamic processes across a wide range of scales, from the large-scale ocean circulation to the small-scale wave motion. To investigate ocean dynamics using sea surface height data measured by satellites, it is crucial to comprehend the various constituents of steric height variability. In this study, the effects of balanced motions (e.g., eddies and ocean fronts) and unbalanced wave motions (e.g., internal waves) on steric height are quantified based on yearlong moored observations at a mid-latitude ocean site of the northeast Atlantic. Overall, balanced motions account for  $\sim 83\%$  of the upper-ocean steric height variance. Steric height variance also show notable seasonal variations and scale dependence. At spatial scales of  $O(10\text{ km})$ , the steric height is predominately determined by balanced motions throughout the year. By contrast, at spatial scales of  $O(1\text{ km})$ , unbalanced wave motions are the major contributor to steric height in summer whereas balanced motions still dominate in winter. This study provides a quantitative assessment of the effects of balanced motions and unbalanced wave motions on steric height, and provide insights for the exploration of next-generation high-resolution altimetry data.

## 1 Introduction

Sea surface height measurements over the global oceans are now routinely derived from satellite altimetry, and have greatly advanced our understanding of ocean dynamics over the last 30 years (Chelton et al., 2011). However, the conventional nadir radar altimeters can only resolve large-scale and mesoscale variability with horizontal resolutions of  $O(100\text{ km})$  (Ballarotta et al., 2019). Xu and Fu (2011, 2012) examined along-track sea surface height measurements to diagnose the dynamical regimes of the global oceans, and pointed out the necessity of higher resolution satellite altimeter for a more accurate detection of the dynamic characteristics globally. The next-generation wide-swath altimetry missions, such as the Surface Water and Ocean Topography (SWOT) altimeter mission which has been launched in December 2022 and the Chinese “Guanlan” mission which is in the early designing stage, are expected to provide, for the first time, 2D sea level maps globally at spatial scales down to the submesoscale (Chen et al., 2019; Fu & Ubelmann, 2014). SWOT spatial resolution is predicted to be of about 15-50 km depending on the local sea state and measurement noise (Morrow et al., 2019; Wang et al., 2019). One of the scientific challenges associated with increasing spatial resolution is that high-frequency internal waves with spatial scales comparable to submesoscale motions can also be observed, and will behave as “noises” for isolating submesoscale motions given the limited temporal revisit of altimetric satellites (e.g., tens of days). Therefore, it is important to quantitatively assess the respective contributions of submesoscale motions and internal waves to sea surface height.

Submesoscale processes are characterised by spatial scales of  $O(0.1\text{-}10\text{ km})$  and temporal scales of several hours to several days (McWilliams, 2016; Thomas et al., 2008; Callies et al., 2020). Recent studies revealed that submesoscale processes modulate the equilibrium state of the upper ocean through a bi-directional kinetic energy cascade (Ferrari & Wunsch, 2009; McWilliams, 2017; Qiu et al., 2017). Further, submesoscale motions are particularly effective at inducing intense vertical velocities in the upper ocean, which may exceed mesoscale processes by one order of magnitude (Klein et al., 2009; Lévy et al., 2018; Su et al., 2018). Given their importance in oceanic energy cascade and vertical tracer transport, the observation and understanding of submesoscales have become key scientific targets of the SWOT mission and the oceanographic community in general.

Internal waves are propagating disturbances in stably stratified fluids, with gravity acting as the restoring force (Gerkema & Zimmerman, 2008; Sutherland, 2010). The internal wave field can be divided into three components: near-inertial waves, internal tides and the internal-wave continuum. Near-inertial waves typically arise from strong variable winds over the ocean surface, and are expected to have minor signatures on sea surface height (Munk & Phillips, 1968; Fu, 1981). In contrast, internal tides and the internal wave continuum are more substantial contributors to the sea surface height field (Callies & Wu, 2019; Chereskin et al., 2019). Müller et al. (2015) demonstrated that high-resolution global ocean models with tidal and atmospheric forcing are beginning to resolve the internal-wave continuum. Savage et al. (2017) analyzed the components of sea surface height and their spatial distributions based on the outputs of global ocean models with different resolutions, and suggested that the internal wave signal will be high-frequency noise for future high-resolution altimeters and its non-stationary component will be difficult to predict. Based on a high-resolution MITgcm global model, Torres et al. (2018) analyzed contributions of balanced and unbalanced motions in kinetic energy and sea surface height, and found that the relative contributions of two classes of motions to various surface fields are complex and dependent on multiple factors, such as seasonal, geographical and the distribution of low and high eddy kinetic energy regions.

Steric height is the baroclinic component of sea surface height, and is mainly determined by dynamical processes such as balanced (sub)mesoscale motions and unbalanced high-frequency internal waves. A quantitative assessment of how these multiscale processes contribute to the steric height is the foundation for exploring the ocean dynamics from satellite sea surface height measurements (Baker-Yeboah et al., 2009; Gill & Niller, 1973).

98 However, observational studies on steric height variability down to the submesoscales are  
 99 still rare. Very recently, Miao et al. (2021) used temporal filtering to quantitatively assess  
 100 the influence of multiscale dynamic processes on steric height based on temperature/salinity  
 101 time series observed by a single mooring in the South China Sea, and showed that the rel-  
 102 ative contribution of submesoscale motions (7.2%) is smaller than those of diurnal (8.5%)  
 103 and semidiurnal (20.2%) internal tides. They also illustrated that steric height associated  
 104 with mesoscales and submesoscales are stronger in winter than summer but the opposite  
 105 occurs for tidal and super-tidal motions.

106 Similar to the findings reported in Miao et al. (2021), here we will show that there is  
 107 a seasonality in steric height of balanced (sub)mesoscale motions and unbalanced internal  
 108 waves in a mid-latitude ocean. In addition to that, we also examine the scale dependence of  
 109 steric height in the spatial domain based on structure function approaches. We show that  
 110 unbalanced internal waves only dominate over balanced (sub)mesoscale motions in summer  
 111 at spatial scales of a few kilometers in the study region.

## 112 2 Observations and Methods

### 113 2.1 Mooring data

114 The data used in this study were primarily collected from nine bottom-anchored subsur-  
 115 face moorings deployed over the Porcupine Abyssal Plain (48.63-48.75°N, 16.09-16.27°W)  
 116 site in the northeastern Atlantic Ocean for the period September 2012 - September 2013  
 117 (Figure 1a), as part of the OSMOSIS (Ocean Surface Mixing, Ocean Submesoscale Inter-  
 118 action Study) experiment (Buckingham et al., 2016; Yu et al., 2019; Erickson et al., 2020).  
 119 The OSMOSIS mooring site is in an abyssal plain of depth close to 4800 m, and is analogous  
 120 to many open ocean regions far away from western boundaries of ocean basins and from  
 121 complex topography. Nine moorings were arranged in two concentric quadrilaterals with  
 122 side lengths of  $\sim 13$  km (outer cluster) and 1-2 km (inner cluster) around a centrally located  
 123 mooring (Figure 1b), which can concurrently capture mesoscale and submesoscale signals to  
 124 a large extent. The mooring sensors comprised a series of Seabird MicroCAT conductivity-  
 125 temperature-depth (CTD) sensors and Nortek Aquadopp current meters at different depths,  
 126 spanning the approximate depth interval 30-530 m. The present study predominately uses  
 127 data from the CTDs.

128 The mooring measurements captured most of the pycnocline plus part of the ocean  
 129 interior throughout the year, and most of the mixed layer during winter months. Here, the  
 130 seasons are defined as follows: winter (December to April) and summer (June to August).  
 131 Temperature, salinity and pressure measurements were recorded with intervals every 5 min-  
 132 utes. We linearly interpolate measurements of temperature and salinity onto surfaces of  
 133 constant depth at 10 m intervals between depths of 50 m and 520 m for each mooring, and  
 134 onto uniform 10 minutes intervals between 5 September 2012 and 5 September 2013. Further  
 135 information regarding the OSMOSIS moorings, such as the detailed distribution of moored  
 136 instruments and associated observational uncertainties, can be found in the work by Yu et  
 137 al. (2019) or Naveira Garabato et al. (2022). Furthermore, the mooring measurements were  
 138 complemented by hydrographic observations acquired by two ocean gliders that navigated  
 139 in a bow-tie pattern across the mooring array for the entire sampling period (Damerell et  
 140 al., 2016; Thompson et al., 2016).

### 141 2.2 Steric height calculation

142 The steric height  $\xi$  is calculated as the integral of specific volume anomaly  $\delta$  from the  
 143 reference pressure  $p_{ref}$  to the pressure  $p$ , and is given by

$$144 \quad \xi = \frac{1}{g} \int_{p_{ref}}^p \delta dp, \quad (1)$$

145 where  $\delta = \alpha(S, T, p) - \alpha_0$ ,  $\alpha_0 = \alpha(35, 0, p)$  is the specific volume  $\alpha$  at local pressure with  
 146  $S = 35$  psu and  $T = 0^\circ\text{C}$ ,  $g$  is the gravitational acceleration. In this study, the reference  
 147 pressure  $p_{ref}$  is chosen to be 480 db, above which the nine OSMOSIS moorings all have  
 148 CTD measurements.

### 149 2.3 Band pass filtering

150 One way to isolate contributions of each dynamical process to steric height is by band  
 151 pass filtering applied in the frequency domain. Here, the frequency bands of interest are  
 152 defined as large-scale forcing (1/90–1/30 cpd), mesoscale (1/30–1/5 cpd), submesoscale  
 153 (1/5–1 cpd, following Naveira Garabato et al. (2022)), near-inertial (0.9–1.1 $f$ ), semidiurnal  
 154 (1.9–2.1 cpd) and supertidal band (>2.1 cpd), where  $f = 2\Omega \sin \phi$  is the inertial frequency  
 155 (with  $\Omega$  as the Earth’s angular velocity and  $\phi$  as latitude). A forth-order Butterworth filter  
 156 with a cutoff of these frequency bands are respectively applied on steric height time series  
 157 at 80 db, and we then compute the root mean square of each frequency band.

### 158 2.4 Second-order structure functions

159 We employ the second-order structure function approach to examine steric height vari-  
 160 ance across spatial scales. The second-order structure function for steric height between a  
 161 given location  $\mathbf{x}$  and another location separated from  $\mathbf{x}$  by the distance  $\mathbf{r}$  is defined as

$$162 D_\xi(\mathbf{r}, \mathbf{x}) = \overline{[\xi(\mathbf{x}) - \xi(\mathbf{x} + \mathbf{r})]^2}, \quad (2)$$

163 where the overbar denotes temporal means.

164 If a given homogeneous, isotropic turbulence spectrum (of energy or tracer variance)  
 165 has power-law behavior over a range of wavenumber between the energy injection and dis-  
 166 sipation scales, then a related scaling law for the structure function is expected (Webb,  
 167 1964). Following McCaffrey et al. (2015), the relationship between the spectral slope  $\lambda$  and  
 168 structure function slope  $\alpha$  is

$$169 \lambda = -\alpha - 1. \quad (3)$$

170 Based on nine distinct mooring sites and their separations, there are 36 combinations  
 171 of mooring pairs, covering spatial scales from 18.7 km down to 1.3 km. We also utilize a  
 172 low-pass filter (with a cutoff period at the local inertial period  $\sim 16$  hours) to decompose  
 173 steric height into sub-inertial and super-inertial components,  $\xi(z, t) = \xi(z, t)_{sub} + \xi(z, t)_{sup}$ .  
 174 For each component, we repeat the estimates of the second-order structure function.

### 175 2.5 Frequency-resolved structure functions

176 The frequency-solved structure function approach was developed by Callies et al. (2020),  
 177 and this approach allows for the assessment of steric height variance as a function of temporal  
 178 and spatial scales. To apply this approach, we first temporally detrend steric height time  
 179 series at nine moorings, then a Hann window is applied to each time series to minimize  
 180 spectral leakage. Subsequently, we calculate their temporal Fourier transform, and lastly  
 181 compute spatial second-order structure functions estimates. Specifically, at each mooring  
 182 location  $\mathbf{x}$ , the temporal Fourier transform of the steric height  $\xi$  is given by:

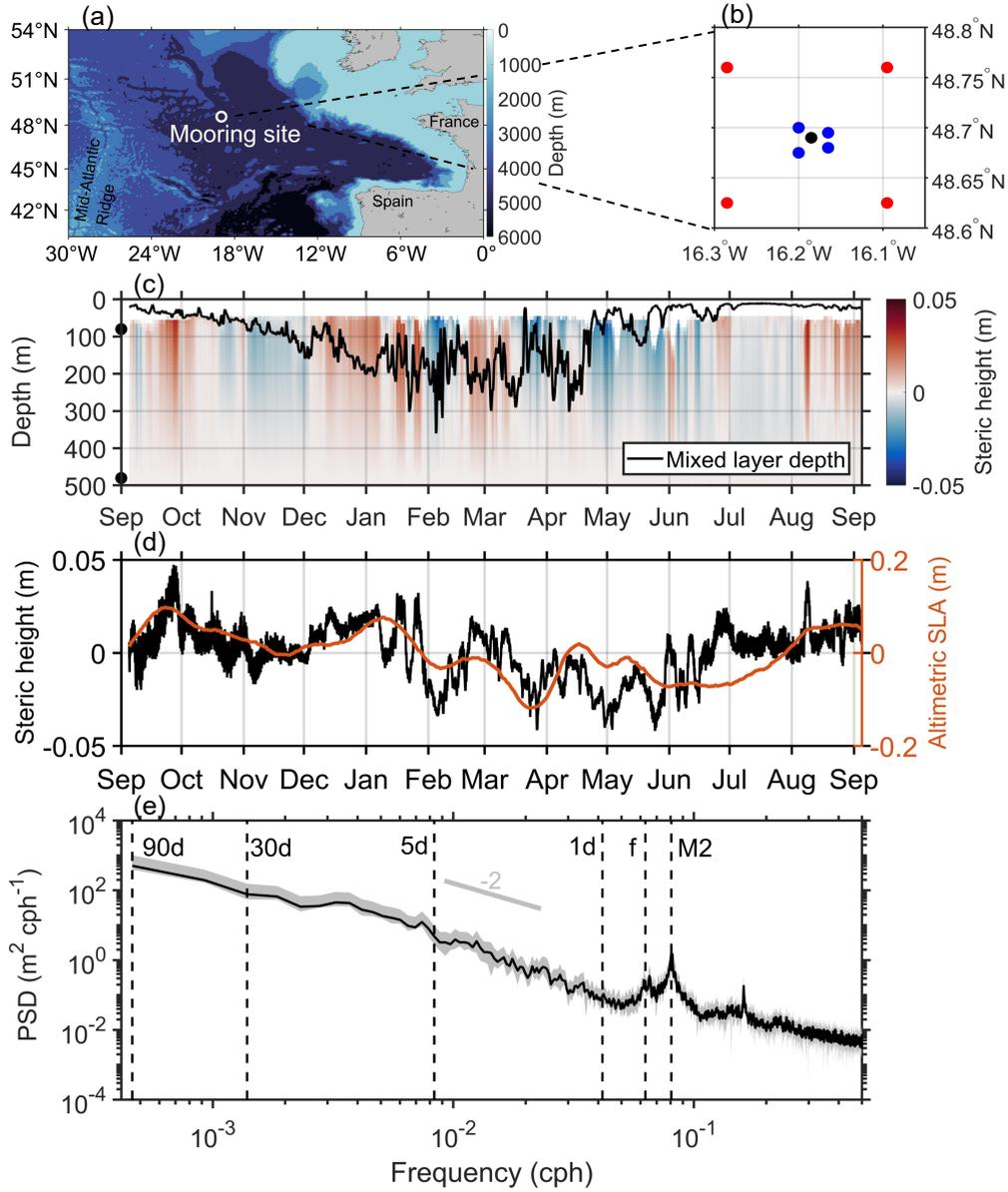
$$183 \tilde{\xi}(\mathbf{x}, \omega) = \int_{-\infty}^{+\infty} \xi(\mathbf{x}, t) e^{-i\omega t} dt. \quad (4)$$

184 For mooring pairs at locations  $\mathbf{x}$  and  $\mathbf{x} + \mathbf{r}$ , the frequency-resolved structure function  
 185 can be calculated by:

$$186 \Delta^\xi(\mathbf{r}, \omega) = \frac{1}{2} \langle \left| \tilde{\xi}(\mathbf{x} + \mathbf{r}, \omega) - \tilde{\xi}(\mathbf{x}, \omega) \right|^2 \rangle, \quad (5)$$

187  
188

where  $\omega$  is frequency,  $\Delta^\xi(\mathbf{r}, \omega)$  is the frequency-resolved second-order structure function as a function of  $\mathbf{r}$  and  $\omega$ , the angle brackets denote the average over nearby spatial distances.



**Figure 1.** (a) OSMOSIS study region in the northeast Atlantic, with bathymetry shown in the colormap on the right. The white circle denotes the location of the OSMOSIS mooring array. (b) Locations of inner (blue circles), outer (red circles) and center (black circle) moorings. (c) Depth-resolved time series of steric height for the central mooring. The two black dots on the y axis respectively indicate the depths corresponding to 80 db and 480 db, and the black line represents the mixed layer depth. (d) Time series of steric height at 80 db from the central mooring (black) and satellite altimetry sea level anomaly (red) from September 2012 to September 2013. (e) Power spectral densities of the steric height. Gray shading indicates 95% confidence intervals, and slope of  $-2$  is shown for reference. The vertical dashed lines indicate the 90 days, 30 days, 5 days, 1 day, near-inertial frequency (1/16 cph) and semidiurnal frequency (1/12.42 cph), respectively.

### 3 Results

#### 3.1 Temporal decomposition of steric height

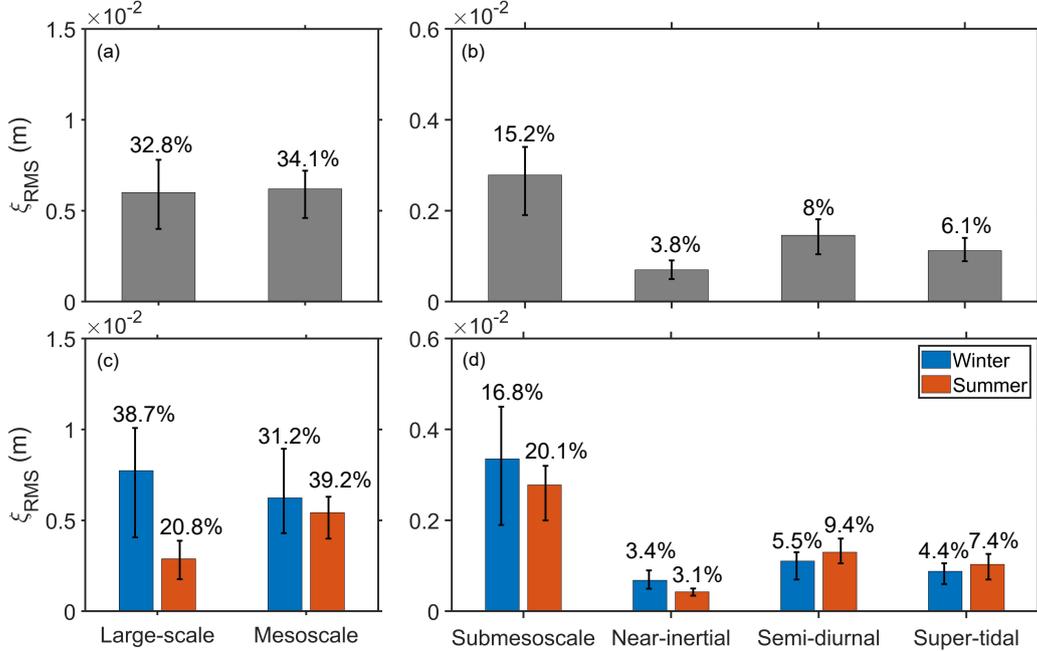
The time evolution of steric height over the annual cycle is shown in Figure 1c, contextualized with the glider-based mixed layer depth. The mixed layer depth is calculated from coincident glider data using a threshold value of potential density increase ( $\Delta\rho = 0.03 \text{ kg m}^{-3}$ ) from a near-surface value at 10 m (Damerell et al., 2016). Steric height is characteristically intensified near the surface and ranges from  $-0.05$  to  $0.05$  m. A closer look at steric height time series at 80 db reveals more details of its temporal variability (Figure 1d). Seasonal heating/cooling, which directly affects near-surface density, is one of the main factors determining the annual cycle of steric height in this region. This large-scale atmospheric forcing leads to predominantly positive values of steric height during non-winter months (such as September-November and July-September) and negative values during winter months (such as January-April; also see Figure S1). Steric height also displays profound high-frequency signals, especially in non-winter months when mixed layer is shallow and vertical stratification is strong. The wave signature in the steric height field is further confirmed by the frequency spectral analysis in Figure 1e, in which the most noticeable feature is the presence of a spectral peak at the semidiurnal tidal frequency that can only be accounted for by internal tides. In contrast, low-frequency eddy-like signals are more evident in winter months when mixed layer is deep and vertical stratification is considerably reduced.

Mooring-derived steric height shows a reasonably consistent trend with the sea level anomaly (SLA) observed by altimetry, although the former is smaller by a factor of 3-4 than the latter. Note that substantial corrections, such as dynamic atmospheric correction and ocean tide correction, have already been applied in the altimetry data, and thus the SLA is expected to be primarily contributed by steric height. The discrepancy in the mooring-based steric height and altimetry SLA likely stems from the missing contributions of steric height in the deep ocean (i.e., below 480 db) by the moorings, as in the comparison between seal-based upper-ocean steric height and altimetry SLA in the Southern Ocean (Siegelman et al., 2020). Additionally, Wang et al. (2018) compared different reference levels for the calculation of steric height to the full-depth steric height based on the output of a high-resolution model, and found the upper-ocean dominance of the steric height component on the sea surface height (accounting for  $>70\%$  of full-depth steric height for a reference level of 580 m in their case).

Frequency band integrated steric height variance are compared in Figure 2. Over the whole year, largest contributions to steric height variability are associated with large-scale atmospheric forcing and mesoscale eddies, constituting 32.8% and 34.1% of the total steric height variance, respectively. Submesoscale motions are the third contributor to the steric height, and account for about 15.2% of the total variance. By contrast, three unbalanced internal wave components (near-inertial, semidiurnal and super-tidal waves) all contribute less than submesoscale motions, making up 3.8%, 8% and 6.1%, respectively. Note that the steric height variance in those internal wave bands may, at least in part, arise from Doppler shifting of submesoscale motions by mesoscale eddies and/or the presence of background spectrum, which would lead to an overestimate of the relative contributions of high-frequency steric height variance.

We next examine the seasonality of steric height across frequencies. The large-scale atmospheric forcing, mesoscale eddies, submesoscale motions and near-inertial waves all produce higher root-mean-square values of steric height in winter compared to summer. The annual-mean surface heat flux is approximately  $-45 \text{ W m}^{-2}$ , indicating a stronger winter cooling compared to summer heating (Yu et al., 2019). This asymmetry between heating and cooling may explain the seasonal difference of steric height variance at lowest frequencies. The seasonality of steric height at mesoscale and submesoscale bands are also expected, because mesoscale eddies and the associated submesoscale motions are found to be more active in winter than in summer (Buckingham et al., 2016; Thompson et al., 2016).

241 Furthermore, the enhancement of near-inertial steric height variance in winter conforms to  
 242 expectations from the seasonality of near-inertial energy generated by surface winds (Yu et  
 243 al., 2022). However, an opposite seasonality is seen at semidiurnal and super-tidal bands,  
 244 where steric height variance is much stronger in summer than in winter. This is mainly  
 245 caused by stronger summer vertical stratification in the upper ocean (Callies et al., 2020;  
 246 Rocha, Gille, et al., 2016).



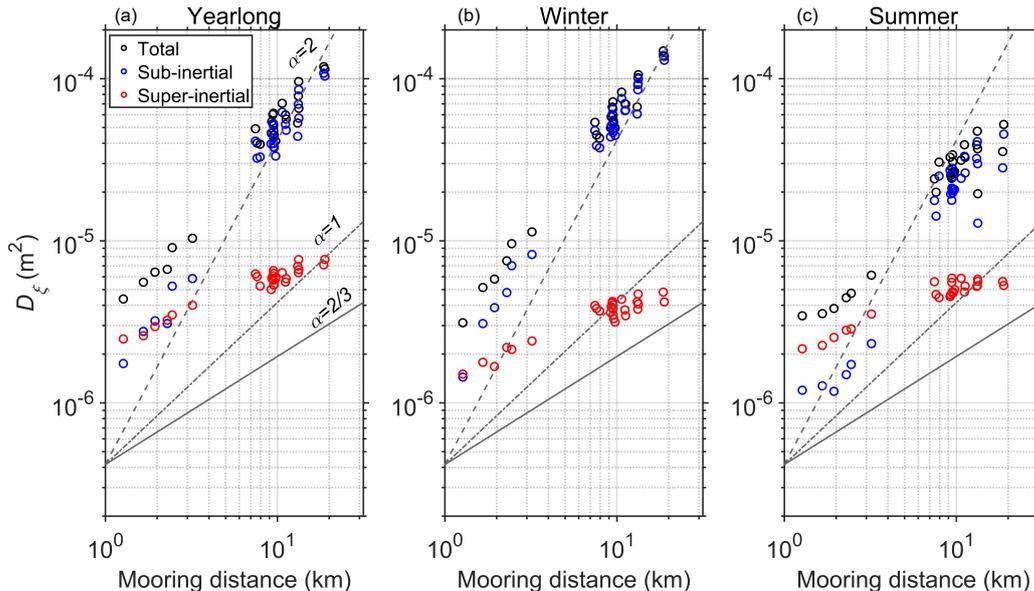
**Figure 2.** (a-b) Annual-mean and (c-d) seasonal-mean of the root-mean-square steric height at 80 db and for different frequency bands. Fraction of each component to the total steric height is marked. Black vertical bars illustrate the 95% confidence intervals estimated using a bootstrap approach.

### 247 3.2 Steric height variance across spatial scales

248 To assess the distribution of steric height variance across spatial scales, we compute  
 249 second-order structure functions of steric height at 80 db for yearlong, winter and summer,  
 250 respectively (Figure 3). Total steric height variance (black circles) exhibits a seasonality  
 251 across spatial scales of  $O(1-20 \text{ km})$ , elevated in winter and reduced in summer. The slope  
 252 for steric height structure functions in winter is comparable to 2, corresponding to a spectral  
 253 slope of  $k^{-3}$ , where  $k$  is the horizontal wavenumber. Note that the spectral slope of kinetic  
 254 energy estimated from the same observations data is between  $k^{-1.6}$  and  $k^{-2}$  (Erickson et al.,  
 255 2020), which indicates a sea surface height spectral slope of  $k^{-3.6}$  and  $k^{-4}$ , slightly steeper  
 256 than the slope  $k^{-3}$  found here. By contrast, the slope of steric height structure functions in  
 257 summer becomes flatter, indicating that steric height variance are reduced at larger scales  
 258 but enhanced at smaller scales.

259 The yearlong record also allows us to compare the sub-inertial (blue circles) and super-  
 260 inertial (red circles) steric height structure functions. In winter, the steric height variance is  
 261 predominately determined by its sub-inertial component, which also follows closely a spectral  
 262 slope of  $k^{-3}$ . The super-inertial component, however, contributes little to the total steric  
 263 height variance with a structure function slope near  $2/3$ , corresponding to a spectral slope

264 of  $k^{-5/3}$ . This slope is close to  $k^{-2}$  for sea level variance predicted by the GM spectrum  
 265 (Callies & Wu, 2019). Sub-inertial steric height structure functions at around 10 km are  
 266 one order of magnitude larger than the super-inertial ones, but their differences decrease  
 267 towards smaller scales. Notably, the sub-inertial and super-inertial steric height structure  
 268 functions are comparable in magnitude at smallest scales. However, a somewhat different  
 269 picture is seen in summer. While the sub-inertial steric height variance still dominates over  
 270 the super-inertial variance at spatial scales of order 10 km, the super-inertial component  
 271 is considerably larger than the sub-inertial component at spatial scales of O(1 km). This  
 272 suggests that high-frequency internal waves dominate over balanced (sub)mesoscale motions  
 273 at scales of O(1 km) in summer.

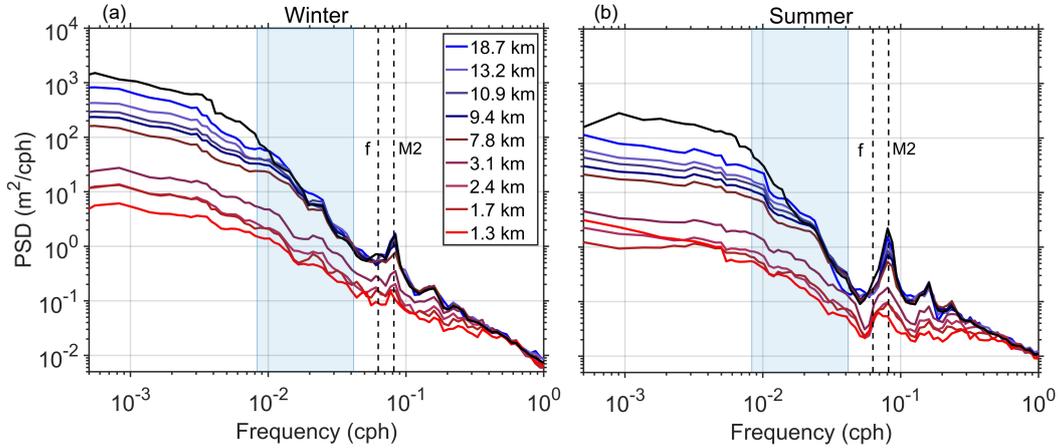


**Figure 3.** The second-order structure functions of steric height at 80 db during (a) the whole year, (b) winter and (c) summer. The black, blue and red circles denotes all, sub-inertial (>16 hours) and super-inertial (<16 hours) motions, respectively. Black lines provide a reference for 2, 1 and 2/3 power-law slopes.

274 The frequency-resolved structure functions of steric height provide useful insights into  
 275 the link between temporal and spatial scales resolved by the mooring array (Figure 4).  
 276 At lowest frequencies, the frequency spectrum (black curve) is considerably larger than  
 277 the frequency-resolved structure functions (colored curves) at all separation distances for  
 278 both winter and summer. This indicates that the steric height variance at low frequen-  
 279 cies is primarily controlled by spatial scales larger than the largest mooring sampled scale  
 280  $2r=37.4$  km. Within the submesoscale band, the frequency spectrum is notably larger in  
 281 winter than in summer, and the frequency-resolved structure functions for the larger sep-  
 282 arations ( $r=7.8-18.7$  km) converge to the frequency spectrum. At highest frequencies, the  
 283 frequency-resolved structure functions converges to the frequency spectrum at all separa-  
 284 tions, indicating that at these periods, the steric height variance becomes decorrelated across  
 285 all mooring pairs.

286 Comparing our work to Callies et al. (2020), who carried out frequency-resolved struc-  
 287 ture functions of horizontal velocity, there are two notable distinctions. Firstly, at the  
 288 near-inertial frequency, the frequency-resolved structure functions of steric height vari-  
 289 ance do not exhibit peaks, primarily due to the negligible signature of near-inertial motion on the

290 steric height variance. Secondly, at semidiurnal frequency, the frequency-resolved structure  
 291 functions of steric height at large separations (e.g.,  $r=7.8-18.7$  km) are found to more closely  
 292 align with the frequency spectrum, unlike the significant gap showed in velocity frequency-  
 293 resolved structure functions. We speculate this difference is partially attributed to the  
 294 presence of large-scale barotropic tide signals in the velocity measurements while absence  
 295 in steric height estimates. To confirm this hypothesis, we further compared the frequency-  
 296 resolved structure functions and frequency spectrum by incorporating the barotropic semi-  
 297 diurnal tide sea level component into the steric height time series, and found a discernible  
 298 increase of spectral peak at the semidiurnal frequency but not for structure functions (Figure  
 299 S2). Another reason may be that low-mode baroclinic tides, which are featured with larger  
 300 spatial scales compared to high-mode components, are not fully captured due to a shallow  
 301 reference level of 480 db. Sea level for low-mode semidiurnal tides predicted by High Resolu-  
 302 tion Empirical Tides (HRET; Zaron et al., 2022) yield a root-mean-square value of 0.28 cm  
 303 at the mooring site, nearly a factor of 2 larger than 0.15 cm estimated from the semi-diurnal  
 304 band of mooring-derived steric height time series. This suggests that a fair amount of steric  
 305 height variability at semi-tidal frequencies, most likely associated with large-scale low-mode  
 306 components, is being missed due to the limited vertical extent of mooring observations.



**Figure 4.** Frequency-resolved structure functions of steric height at 80 db for (a) winter and (b) summer. The black lines show the frequency spectrum and the colored lines show the frequency-resolved structure functions for the separations  $r$  given in the legend. The light blue area represents the submesoscale range, corresponding to 1-5 days. The two vertical lines denote the near-inertial frequency (1/16 cph) and semidiurnal frequency (1/12.42 cph), respectively.

## 307 4 Conclusion

308 In this study, we provide observational insights, for the first time, into seasonal to sub-  
 309 mesoscale steric height variability down to scales of 1 km at a typical mid-latitude ocean site  
 310 of northeast Atlantic. Throughout the year, large-scale atmospheric forcing and mesoscale  
 311 eddies are the two largest contributors to steric height, with their respective relative contribu-  
 312 tions of 32.8% and 34.1%, followed by submesoscale motions (15.2%), semidiurnal (8%),  
 313 supertidal (6.1%) and near-inertial (3.8%) waves. Another important finding is that contribu-  
 314 tions of balanced motions and unbalanced internal waves to steric height display a strong  
 315 seasonal cycle and scale dependence. Low-frequency balanced motions largely dominate  
 316 the upper-ocean steric height variance in winter and at scales of  $O(10)$  km. In contrast,  
 317 high-frequency unbalanced internal waves become increasingly important in summer, and  
 318 are able to dominate over low-frequency balanced motions at scales of  $O(1)$  km).

319 The study by Miao et al. (2021) found some similar results to our findings regarding  
320 the variations in steric height between seasons caused by multiscale dynamic processes. One  
321 notable difference is the relative contribution of semidiurnal tides on steric height was found  
322 to be significantly greater than that of submesoscale processes in the South China Sea,  
323 regardless of the season. By contrast, our analysis found that the relative contribution of  
324 semidiurnal tides to the steric height is lower than that of submesoscale motions and that  
325 reported in Miao et al. (2021). It should be emphasized that the South China Sea is well  
326 known for its strong internal tides radiated from the Luzon Strait (Alford et al., 2015),  
327 while the OSMOSIS observational area is representative of mid-latitude open ocean regions  
328 characterised by moderate eddy kinetic energy and weak mean flow and internal waves.

329 Our results add observational assessment on the steric height variance in the subme-  
330 soscale range, which is crucial to interpret future altimetric high-resolution sea surface height  
331 maps. Apart from steric height, the sea surface height also includes contributions from the  
332 bottom pressure anomaly (mainly caused by barotropic tides in the open ocean) and the  
333 atmospheric pressure loading (i.e., the inverted barometer effect). A key outcome of the  
334 satellite-derived sea surface height is to estimate surface velocity via geostrophy. To this end,  
335 one needs to filter out the signals associated with bottom pressure anomaly and atmospheric  
336 pressure loading, and then subtract unbalanced internal wave signals. The barotropic tides  
337 and inverted barometer effect are large-scale signals in the sea surface height, and can be  
338 largely removed by spatial filtering. Such corrections have been widely applied to present  
339 altimetry measurements. It has been recently suggested that sea surface height variance  
340 spectra associated with balanced flow drop off steeply with wavelength and internal tides  
341 are likely the main factor affecting the accuracy to infer submesoscale balanced flow at  
342 the smallest scales resolved by SWOT (Callies & Wu, 2019; de Marez et al., 2023). For  
343 the SWOT mission, spatial filtering is likely the practical approach to further mitigate the  
344 effects of fast unbalanced variability given its long repeat sampling cycle (21 days).

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### 352 **Data Availability Statement**

353 All OSMOSIS mooring data are freely available, and are archived at the British Oceanographic  
354 Data Centre. Moored observations can be obtained from [https://www.bodc.ac](https://www.bodc.ac.uk/projects/data_management/uk/osmosis/)  
355 [.uk/projects/data\\_management/uk/osmosis/](https://www.bodc.ac.uk/projects/data_management/uk/osmosis/). Altimeter data were obtained from [https://](https://data.marine.copernicus.eu/product/)  
356 [data.marine.copernicus.eu/product/](https://data.marine.copernicus.eu/product/).

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# Supporting Information for “Effects of balanced motions and unbalanced internal waves on steric height in the mid-latitude ocean”

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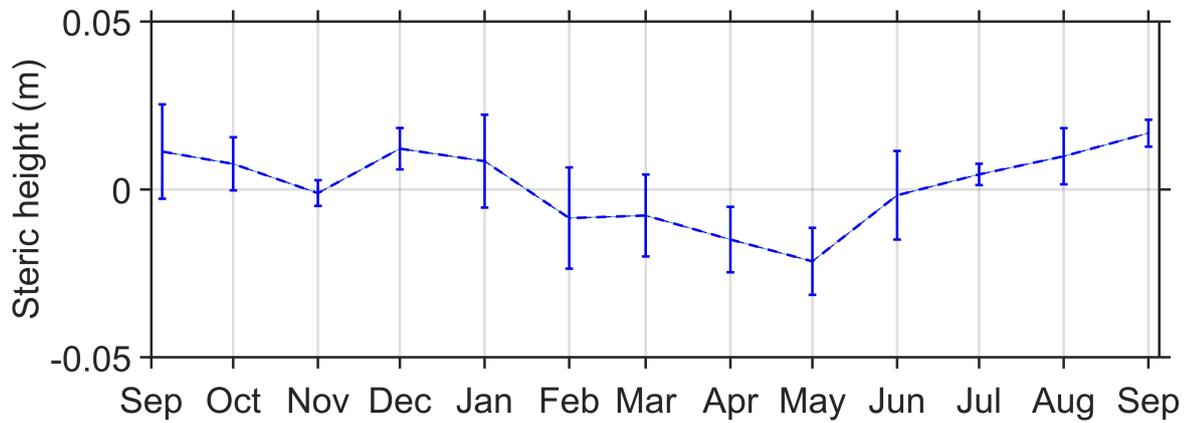
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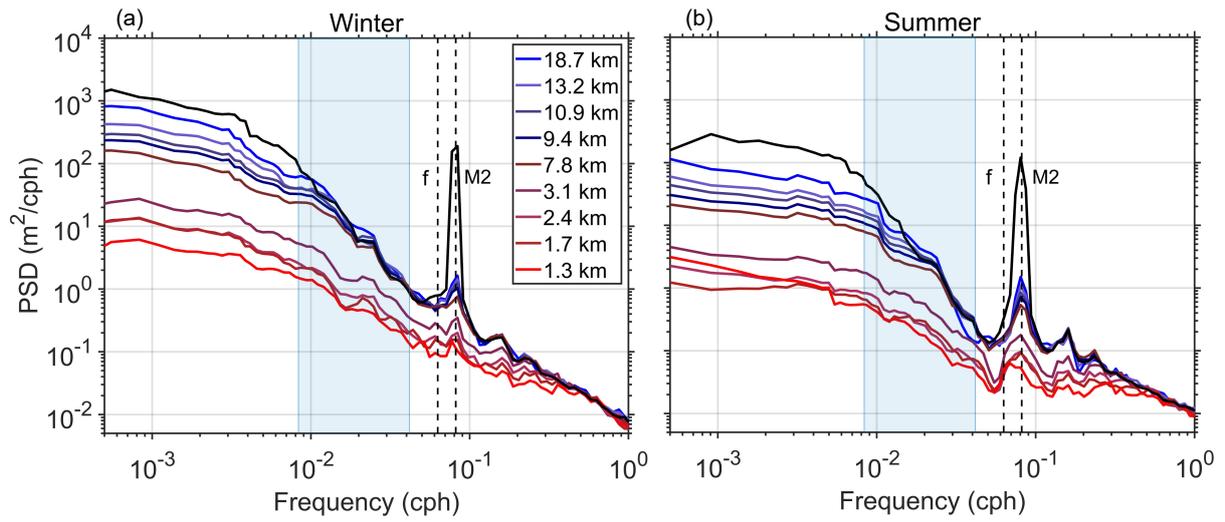
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1. Figures S1 to S2

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**Figure S1.** Monthly averaged values of steric height at 80 db, and the standard deviation computed for each month is illustrated by the vertical bars.



**Figure S2.** Same as Figure 4, but with the sea level associated with barotropic tides artificially added in the steric height time series.