A regionally refined and mass-consistent atmospheric and hydrological de-aliasing product for GRACE, GRACE-FO and future gravity missions

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Key Points:

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9	•	We provide an atmosphere-hydrology de-aliasing product with regional mass-consistent
10		refinement over Europe.
11	•	Using non-hydrostatic as opposed to hydrostatic numerical weather prediction model
12		output significantly impacts the de-aliasing product.
13	•	We found that for extreme events additional moisture fields unaccounted in present
14		AOD models can reach magnitudes relevant for de-aliasing.

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15 Abstract

De-aliasing products are used in the estimation process of satellite-based gravity field 16 computation to reduce errors from high-frequency mass variations that alias into monthly 17 gravity fields. The latest official product is AOD1B RL07 and describes non-tidal atmo-18 sphere and oceanic mass variations at 3-hourly resolution. However, the model-based de-19 aliasing products are inevitably incomplete and prone to temporally and spatially cor-20 related errors that substantially contribute to errors in the estimated gravity fields. Here, 21 we investigate possible enhancement of current de-aliasing products by nesting a regional 22 high-resolution atmospheric reanalysis over Europe into a global reanalysis. As further 23 novelty we include almost mass consistent terrestrial water storage variability from a re-24 gional hydrological model nested into a global model as additional component of the de-25 aliasing product. While we find in agreement with earlier studies only minor contribu-26 tions from increasing the temporal resolution beyond 3-hourly data, our investigations 27 suggest that contributions from continental hydrology and from regional non-hydrostatic 28 atmospheric modeling to sub-monthly mass variations could be relevant already for grav-29 ity fields estimated from current gravity missions. Moreover, in the context of extreme 30 events, we find regionally contributions from additional moisture fields, such as cloud liq-31 uid water, in the order of a few mm over Europe. We suggest this needs to be taken into 32 account when preparing data analysis schemes for future space gravimetric missions. 33

³⁴ Plain Language Summary

Observing temporal variations in the Earth's gravity field with satellite gravime-35 try plays an essential role for monitoring mass transports on and underneath the Earth's 36 surface. This is crucial for understanding the evolution of floods and droughts, the role 37 of groundwater pumping, and to quantify the melting of ice sheets and glaciers and the 38 resulting sea level rise. In order to isolate the target variable (e.g., terrestrial water stor-39 age changes) unwanted signals (e.g., fast mass variations in the atmosphere) need to be 40 removed in the gravity field estimation process using background models, so-called de-41 aliasing models. This paper focuses on improving background models by incorporating 42 regional high-resolution models, which more specifically resolve certain processes in the 43 atmosphere. Our hypothesis is that this will lead to better gravity fields with increased 44 spatial resolution and less noise. Moreover, we find that considering fast hydrological vari-45 ations as additional background model could improve gravity fields from the current satel-46 lite mission GRACE-FO. For the first time, we quantify contributions from so far ne-47 glected atmospheric moisture fields, such as cloud liquid water, to enhance background 48 models in the context of extreme events – which, however, will likely be in particular rel-49 evant for more sensitive gravity missions in the future. 50

51 **1** Introduction

The GRACE and GRACE Follow-On (GRACE-FO) satellite missions have pro-52 vided unique data sets for studying the mass loss of the world's ice sheets and glaciers 53 (Tapley et al., 2019; Wouters et al., 2019), variability and extremes in terrestrial water 54 storage (TWS) (Kusche et al., 2016; Zhao et al., 2017; Gerdener et al., 2020; Han et al., 55 2021) and its long-term response to anthropogenic forcing (Rodell et al., 2018), and the 56 mass-related component of sea level change (WCRP Global Sea Level Budget Group, 57 2018). However, the spatial resolution of the standard monthly GRACE/-FO data prod-58 ucts is limited to about 300-400 km, which is insufficient for many potential applications 59 (Pail et al., 2015; Wiese et al., 2022). 60

The biggest obstacle to improving the resolution of gravity fields with a single-pair satellite mission is, next to instrument noise, the temporal aliasing of high frequency mass variations due to the poor sampling geometry (Flechtner et al., 2016; Behzadpour et al., 2019). In the current GRACE/-FO processing (see, e.g. Dahle et al., 2019), the effect

of tidal and sub-monthly non-tidal mass variations in ocean and atmosphere is removed 65 from level-1 data on the basis of model forecast simulations, so-called Atmosphere and 66 Ocean De-aliasing (AOD) data (Dobslaw, Bergmann-Wolf, Dill, Poropat, Thomas, et al., 67 2017). Corresponding monthly mean AOD fields are provided to users in level-2 data (spher-68 ical harmonic coefficients) for optionally restoring them depending on the application 69 (Uebbing et al., 2019). However, these AOD products are inevitably imperfect, and er-70 rors map into the GRACE/-FO level-2 data affecting science results even at longer timescales 71 (Velicogna et al., 2001; Han et al., 2004; Boy & Chao, 2005; Wahr et al., 2006; Seo et 72 al., 2008; Bonin & Chambers, 2011; Kvas, Behzadpour, et al., 2019b; Kvas & Mayer-Gürr, 73 2019; Zhou et al., 2023). 74 Candidate future gravity missions either as a successor of GRACE-FO, such as GRACE-75 C (envisaged from about 2027; Flechtner (2020)) and the Mass-Change and Geosciences 76 International Constellation (MAGIC, from about 2031; according to Massotti et al. (2021)), 77

or in the long run, concepts based on quantum technology (Rossi et al., 2023), seek to
 overcome the temporal aliasing problem partly by flying multiple satellite pairs, possibly some in inclined orbits. However, due to the still limited sampling even of multi-pair
 missions (Elsaka et al., 2014), contamination with unmodeled high frequency mass variations will likely become the single most dominant error source, after instrument technology improves. In addition, we speculate that mission targets will be set higher in temporal and spatial resolution, with higher demands in AOD modeling.

The biggest contribution to non-tidal sub-monthly time variability of the Earth's gravity field is actually due to atmospheric mass variations (Thompson et al., 2004). It should be mentioned here that in non-ocean applications of GRACE/-FO data, the atmospheric de-aliasing fields removed during gravity field processing are usually not restored to the estimated monthly gravity fields since one is interested in groundwater, snow or ice mass change, and atmospheric mass variations have never been part of the mission objectives.

The Earth's lower atmosphere mainly consists of dry air constituents, and of wa-92 ter in different states; i.e. water vapor, cloud water, rain water, cloud ice and snow. Mod-93 eling the atmosphere contribution to AOD entails that multi-level fields from numeri-94 cal weather prediction (NWP) model runs are converted to 4D mass density, and then further to time series of spherical harmonic potential coefficients. This involves a 3D in-96 tegration approach including various approximations (Swenson & Wahr, 2002; Boy & 97 Chao, 2005). Many efforts have been made by previous studies to identify the effect of 98 lateral and vertical discretization, quadrature and interpolation approaches, orography 99 representation (Dobslaw, 2016), and the geometrical and physical approximations be-100 ing applied in the georeferencing of pressure-level fields (Forootan et al., 2013; Yang et 101 al., 2021) during the vertical integration (Table B2). 102

The latest release of the official AOD data, AOD1B (RL07; Shihora et al., 2022), 103 makes use of European Centre for Medium-Range Weather Forecasts (ECMWF) ERA5 104 reanalysis data concatenated with ECMWF forecast (hereafter IFS) simulations. AOD 105 coefficients are computed at 3-hourly resolution, and completed up to degree and order 106 180 using an ellipsoidal reference surface. The vertical integration procedures follows Dobslaw, 107 Bergmann-Wolf, Dill, Poropat, and Flechtner (2017) and implements findings by Swenson 108 and Wahr (2002) and Boy and Chao (2005), who demonstrated the importance of con-109 sidering the vertical structure of the atmosphere in the de-aliasing approach (i.e., 3D vs 110 2D surface pressure method). Here, as to the best of our knowledge in all previous stud-111 ies, only the dry air and the water vapor (i.e., specific humidity) fields are being used. 112 while other (additional) moisture fields, i.e., rain water content, cloud water content, cloud 113 ice content and snow content, are disregarded. 114

A number of studies (Forootan et al., 2013; Yang et al., 2021; Shihora et al., 2022) have compared the use of atmospheric fields from the IFS and the reanalyses ERA-Interim (ERA-I hereafter) and ERA5 for de-aliasing. The spatial resolution of ECMWF-IFS has

changed over time and is about 9 km since 2016, while ERA-I and ERA5 have a grid res-118 olution of 79 km and 31 km, respectively. Dobslaw (2016) showed that mapping surface 119 pressure grids from both the operational ECMWF model and ERA-I to a common ref-120 erence orography allows to reduce relative biases and residual variability by about one 121 order of magnitude, and achieves consistency at a level of about 1 hPa. Moreover, Dobslaw 122 et al. (2016) derived pairwise RMS differences of band-passed surface pressure for pe-123 riods of 10–30 days; for related ECMWF models they found differences of $\sim 0.2 \text{ hPa}$ (equiv-124 alent to 2 mm change in water height), whereas comparisons to the Climate Forecast Sys-125 tem Reanalysis (CFSR; Saha et al., 2010) and NASA's Modern Era Retrospective Anal-126 ysis for Research and Applications (MERRA; Rienecker et al., 2011) reanalyses led to 127 RMS differences of up to 0.7 hPa in Antarctica. One may conclude at this point that global 128 atmospheric modeling has already, or will in the near-future, converge to an extent suf-129 ficient for AOD modeling. 130

Independent of the definition of the equipotential surface used for referencing the 131 height coordinate, all current NWP models simulate the atmosphere with a spherical geopo-132 tential assumption (White et al., 2005). The shallow-atmosphere assumption is a geo-133 metric approximation that satisfies Lagrangian symmetries and thus conservation equa-134 tions; its main advantage is that it allows a separation of the variables in the vertical and 135 the horizontal direction (White et al., 2005). The hydrostatic approximation commonly 136 employed in global atmospheric models further assumes that the vertical acceleration $D/Dtz \approx 0$ 137 is negligible. As a consequence, vertical changes in density are determined by temper-138 ature only, since pressure is constraint by the vertical velocity equation, which is reduced 139 to $dp = -\rho g dz$. All models considered for AOD products so far have in common that 140 their physics implementation makes use of the hydrostatic and shallow-atmsophere as-141 sumptions. 142

It is largely unknown how big the errors in temperature, surface pressure, and spe-143 cific humidity and in the resulting AOD coefficients are, and to what extent they can be 144 estimated from comparing model forecasts and reanalyses. It is further unclear to what 145 extent such errors propagate through the GRACE/-FO level-2 processing and whether 146 they can be partly mitigated by restoring monthly mean atmospheric mass fields, whether 147 reliable error estimates can potentially be used in least-squares weighting procedures (Zenner 148 et al., 2010), or in the assessment of future mission concepts (Dobslaw et al., 2015, 2016). 149 Abrykosov et al. (2021) demonstrated additional benefit for gravity retrieval performance 150 151 from incorporating ocean tide error covariance matrices. However, the additional benefit was found to be limited by the performance of the non-tidal atmospheric and oceanic 152 components of the de-aliasing product, which stresses again the importance of improv-153 ing these products. 154

Over the global ocean and at a wide range of frequencies, the sea surface reacts to 155 pressure forcing in a way that is generally described as "inverse barometric effect" (Wunsch 156 & Stammer, 1997). This means that the mass of an atmosphere-ocean column is con-157 served to first order. In the atmosphere-ocean de-aliasing procedure, an ocean model is 158 forced by pressure and surface winds and thus provides a more realistic picture. How-159 ever, it is clear that errors in atmospheric surface pressure are still to first order absorbed 160 in the ocean response that is forced with the same atmosphere model. Since the focus 161 of this contribution is not on providing an operational AOD product, we chose not to 162 force an ocean model here and the evaluations in this study will, for the time being, re-163 fer to the atmosphere only. 164

Since atmospheric water mass redistribution occurs between water vapor, cloud liquid water and other constituents via phase changes, total mass can only be conserved if all mass-related model fields are consistently considered in the AOD computation, which is currently not the case. While the total mass fraction of moisture fields other than water vapor, i.e., the above defined "additional moisture fields", is thought to be smaller than 1 mm in equivalent water height (EWH) in the global average, we will show that the sum of these fields may reach in localized extremes, e.g., in case of a thunderstorm,
values larger than 5 mm EWH. For tropical storms we expect even larger values. This
means, in the case of extremes the additional moisture fields can reach about the same
magnitude as water vapor variability in regular conditions and, thus, become relevant
not only for de-aliasing products but also when analyzing extremes based on TWS maps.

A somewhat related consistency problem then occurs when rain or snow hits the 176 land surface and evapotranspiration adds to atmospheric water vapor. The mass flux across 177 the land-atmosphere interface would need to be derived from coupled model components, 178 if one were to confront GRACE/-FO or future gravimetric mission data consistently with 179 a simulation of mass change in the full atmosphere and hydrological column. Such fully 180 coupled models do not yet exist globally with a resolution sufficient for de-aliasing. In-181 stead, in the present situation, TWS change is mostly derived from off-line hydrologi-182 cal models – while this is a best-practice approach if, e.g., precipitation data introduced 183 as forcing data into these models is inconsistent with the atmospheric de-aliasing prod-184 uct, it should at least be noted that derived TWS change is not mass-consistent with the 185 AOD procedures and also modeled evapotranspiration does not feedback into the atmo-186 sphere. We speculate that this again will lead to likely small but essentially unknown 187 errors. 188

In order to explain the conundrum, let us follow the path of water in a thought ex-189 periment: If we imagine an area of strong local water vapor convergence, possibly driven 190 by evapotranspiration over land, at some point cloud ice or liquid water may rapidly form 191 and then lead to convective rainfall events. Such extreme weather events then often lead 192 to catastrophic flooding. While a space gravimetric observable such as intersatellite track-193 ing will track the entire, uninterrupted history of the resulting mass convergence in the 194 atmosphere and later at the land surface, the conventional AOD procedure will remove 195 a modeled large-scale water vapor field from these. As a result the de-aliased observa-196 tions still contain a residual water vapor mass signal, the full cloud ice and liquid wa-197 ter signal, and of course the signature of the resulting rainfall in terms of land water stor-198 age change. 199

High-frequency mass redistribution of TWS, albeit occurring at longer timescales 200 as compared to atmospheric mass variability, may cause additional temporal aliasing arte-201 facts that are not taken into account when using the standard de-aliasing product AOD1B 202 for estimating monthly GRACE fields. However, some studies (Zenner et al., 2014; Eicker 203 & Springer, 2016; Ghobadi-Far et al., 2022) have suggested that removing sub-monthly 204 hydrological mass redistribution prior to level-1 GRACE analysis may help in mitigating aliasing errors. This is also confirmed by, e.g., Kvas, Behzadpour, et al. (2019a) who 206 considered terrestrial hydrology based on the Land and Surface Discharge Model (LSDM; 207 Dill, 2008) as additional de-aliasing product and by Schindelegger et al. (2021) who used 208 experimental (un-reduced, full-signal) daily GRACE solutions for de-aliasing in lieu of 209 model simulations. In this paper we will revisit the question whether AOD and hydro-210 logical de-aliasing (called AOHD in the following) may be an asset for future de-aliasing 211 products, while at the same time suggesting a way to maintain (as far as possible) mass 212 consistency across the land-atmosphere interface. 213

With expanding HPC opportunities, global models move to even higher spatial res-214 olution and non-hydrostatic computations become a necessity. Since it is unclear what 215 this will mean for AOD computations, both in terms of required mathematical proce-216 dures and in terms of expected differences in mass variability compared to the conven-217 tional analyses, we decided to investigate the use of the non-hydrostatic, regional COSMO-218 REA6 atmospheric data set (Bollmeyer et al., 2015) for de-aliasing. COSMO-REA6 has 219 been successfully validated in several studies, including its land-atmosphere water fluxes 220 compared against GRACE and runoff data and global atmospheric reanalyses (Springer 221 et al., 2017). 222

We used COSMO-REA6 atmospheric forcings to run the Community Land Model 223 (CLM) version 3.5. This allows us to complement the COSMO-REA6 atmospheric mass 224 variability with nearly mass-consistent hydrological storage variations, again, of course, 225 limited to the European Coordinated Regional Climate Downscaling Experiment (EURO-226 CORDEX) domain. A caveat may be that the evaporative fluxes in COSMO-REA6 were 227 not consistently derived from our own CLM run; yet we believe this provides a very good 228 approximation to mass-consistent atmosphere (A) and hydrology (H) fields. For consis-229 tency, we use ERA-I for the remaining part of the atmosphere, and output from the Wa-230 terGAP Global Hydrology Model (WGHM; Müller Schmied et al., 2021) for the global 231 hydrological component. We developed mathematical procedures for the re-gridding and 232 nesting of global and regional data sets prior to spherical harmonics computation. The 233 de-aliasing product AHD-UB generated in this study is available from Mielke et al. (2023). 234

To the best of our knowledge, this is the first investigation on the use of regional, 235 non-hydrostatic atmospheric model data for mass de-aliasing, and it is also the first study 236 that looks at mass-consistent hydrological de-aliasing. The methods that we describe here 237 238 could be applied to other continental regions where high-resolution non-hydrostatic reanalyses exist already now (e.g., the North American Regional Reanalysis, NARR; Hunter 239 et al., 2020). The same AOHD data could then be used for removing atmospheric and 240 hydrological mass variability from other geodetic data where higher spatial resolution 241 than for space gravimetry is required, e.g., in the case of terrestrial gravimetry or GNSS. 242 To our knowledge, this is also the first study which explicitly separates the contribution of different atmospheric constituents, such as dry air and water vapor also including rain 244 water, cloud water, cloud ice, and snow to total mass variability. 245

This paper is organized as follows: in the main part we provide results, discussion and possible future investigations. We focus on December 2007 for the results in the main part and we make additional figures available for February, June, and December 2007 in the supplementary material (SM). In the Appendix, we (i) introduce the applied models and data sets, (ii) provide a comprehensive summary of the vertical integration method applied in this study and in previous publications (Table B2) and (iii) explain the nesting of regional into global models.

253 2 Results

We investigate sub-monthly vertically integrated atmospheric and hydrological mass variations and discuss the contribution from: (i) increasing the temporal resolution of numerical models, (ii) applying regional high-resolution (and non-hydrostatic) models, (iii) considering additional atmospheric moisture fields with particular regard to extreme events, and (iv) taking into account continental hydrology.

For the atmospheric component we included two global and two regional NWP reanalyses in our study, the global ECMWF models ERA-Interim (ERA-I) and ERA5 that are part of the official AOD models (Dobslaw, Bergmann-Wolf, Dill, Poropat, Thomas, et al., 2017; Shihora et al., 2022), the Copernicus regional reanalysis for Europe (CERRA), and the Consortium for Small-scale Modeling reanalysis at 6 km resolution (COSMO-REA6), with the latter being the only non-hydrostatic model (Annex A1).

For the hydrological component we make use of the WaterGAP Global Hydrology Model (WGHM; Müller Schmied et al., 2021) and the regional high-resolution Community Land Model (CLM; Oleson et al., 2008) in version 3.5 (Annex A2). We integrated output from the NWP models vertically based on the ITG-3D approach (Forootan et al., 2013) as described in Annex B2 and nested regional into global models following Annex B4.

The first order mass consistent global atmosphere and hydrology de-aliasing data set with regional refinement over the EURO-CORDEX domain, the so-called AHD-UB data set, is available from Mielke et al. (2023) for the year 2007. In this manuscript we show examples for December 2007; results for March, June, and September are provided in the supplementary material (SM).

In the following, we will (i) describe how atmospheric vertically integrated mass 276 (VIM) changes when using different NWP models at different temporal and spatial res-277 olution, (ii) discuss the magnitude of the individual wet atmosphere components, and 278 (ii) assess contributions of continental hydrology to sub-monthly mass variability. Tem-279 poral variability of a specific data set, e.g., within one month, is provided as standard 280 deviation computed with respect to the monthly mean. Differences of the temporal vari-281 ability represented by the individual data sets are computed as root mean square devi-282 ation (RMSD). 283

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2.1 Contribution of high-resolution atmospheric models

High-resolution regional models may provide improved VIM variations compared to global models due to their increased spatial and temporal resolution and by explicitly implementing non-hydrostatic dynamics. Both aspects lead to better resolved smallscale processes that potentially provide an asset for obtaining more realistic VIM estimates in particular in the case of localized extreme events.

As expected, we find that sub-monthly VIM varies within a range of 870 mm (SM Figure S4). The standard deviation of VIM with respect to the monthly mean reaches values between 10 and 250 mm in December 2007 (Figure 1). The mid-latitudes are characterized by strong baroclinic instabilities leading to mass fluctuations due to cyclone activity and by large-scale pressure variations such as the North Atlantic Oscillation.



Figure 1. Standard deviation of ERA5-based sub-monthly vertically integrated atmospheric mass in December 2007.

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The geometric height at the surface level, referred to as orography here, varies among 295 individual NWP models and contributes to differences in VIM. The orography of each 296 of the models is compared in Figure 2 over Europe. For ERA5 and ERA-I the orogra-297 phy was derived from the geopotential following the OFCM (Office of the Federal Co-298 ordinator for Meteorological Services and Supporting Research) conventions according 299 to Forootan et al. (2013); see also Equation B7. For CERRA the geopotential heights 300 were converted to orography, accordingly. As COSMO-REA6 is based on terrain-following 301 coordinates (Doms, 2011), the geometric height is provided directly as surface orogra-302 phy. 303

The largest differences in orography are found in mountainous areas, where higher spatial resolution allows a more detailed representation of elevation differences. Comparing the global models, where orography is averaged over 31 km (ERA5) or 79 km (ERA-I), and the regional models with ~6 km grid width (COSMO-REA6, CERRA), differences of several hundred meters are found in regions such as the Alps (Figure 2). For the two high-resolution models, we generally find smaller differences of tens of meters, although differences of several hundred meters are possible locally due to different grid conventions.



Figure 2. Differences in orography in meters over the EURO-CORDEX domain for selected models plotted on a 0.25° grid for the global models in (a) and on a linearly interpolated 0.1° grid in (b) and (c).

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In 2017, with the RL06 product, the temporal resolution of the AOD1B product 312 was increased from 6-hourly to 3-hourly sampling (Dobslaw, Bergmann-Wolf, Dill, Poropat, 313 Thomas, et al., 2017). We computed the impact of the increased temporal resolution by 314 first sampling to the coarser time intervals and then linearly interpolating to compare 315 to the data set with the higher sampling frequency. After subtracting the respective monthly 316 mean values, we calculated the root mean square deviation (RMSD) between the two 317 data sets. Over Europe, Figure 3b shows differences in VIM variability of about 1 to 5%318 (RMSD differences of 2 to 6 mm) when moving from 6-hourly to 3-hourly fields. Mov-319 ing further from 3-hourly de-aliasing fields to 1-hourly fields captures additionally about 320 0 to 3% (RMSD differences of 1 to 4 mm) of total sub-monthly variability (Figure 3c and 321 SM Figure S5-S7). Generally, the change in VIM variability when moving from 6-hourly 322 fields to 3-hourly fields is about twice as large as the effect when moving from 3-hourly 323 to 1-hourly fields. This means, in terms of VIM variability, that the impact of switch-324 ing from 3-hourly fields to 1-hourly fields is comparable to switching the version of the 325 de-aliasing model from ERA-I to ERA5 (Figure 4). Over mountainous areas, differences 326 in the overall variability are smaller due to smaller absolute values of VIM. 327

The change from ERA-I to ERA5 results in only small differences in larger-scale 328 VIM patterns of about 5 mm (2 to 4% of total sub-monthly variability, Figure 4a). Dif-329 ferences in VIM between the regional CERRA reanalysis and the global ERA5 reanal-330 ysis show small-scale differences of similar magnitude with larger differences possible in 331 mountainous regions (Figure 4b) due to differences in orography. Interestingly, the dif-332 ferences in sub-monthly VIM variability between COSMO-REA6 and ERA5 models can 333 be twice as large as those between CERRA and ERA5, or between ERA-I and ERA5. 334 The differences between COSMO-REA6 and CERRA reach about the same magnitude 335 as between COSMO-REA6 and ERA5. In this case, however, small-scale differences do 336



Figure 3. RMSD of sub-monthly atmospheric mass in December 2007 over the EURO-CORDEX region for (a) 1-hourly ERA5, (b) 6-hourly (linearly interpolated) ERA5 minus 3-hourly ERA5, and (c) 3-hourly (linearly interpolated) ERA5 minus 1-hourly ERA5. In order to compute differences between time series with different temporal resolution, the lower-resolution time series is linearly interpolated to the higher-resolution time steps.

not occur because the models have approximately the same spatial resolution. One reason for the differences in the VIM variability of COSMO-REA6 compared to the other
 models may be that COSMO-REA6 is the only non-hydrostatic model.

To take a closer look at the temporal evolution of the VIM of each model, we examined VIM averaged over five regions. Each of them covers about 280,000 km² (Figure 5), which roughly corresponds to the spatial footprint of the GRACE mission. The regions are over the North Sea on the British coast (A), over France (B), over the Alps (C), over the Mediterranean Sea on the east coast of Sicily (D), and over Turkey (E).

In general, sub-monthly VIM varies in December 2007 in a range of 600 mm EWH 346 for footprint A (North Sea) with a significant increase after one third of the month (Fig-347 ure 6). Indeed, a series of low pressure systems passed over Northern Germany and the 348 North Sea in the first half of December, whereas the second half of December was char-349 acterized by stable high-pressure conditions. Differences between CERRA vs. ERA5 tend 350 to be slightly smaller than differences between ERA-I vs. ERA5. Generally, differences 351 between these models are below 10 mm EWH and they resemble random noise (see also 352 SM Figure S11-S19). In contrast, the temporal evolution of VIM from COSMO-REA6 353 deviates more significantly and in a non-random, systematic way from the other mod-354 els. This could be related to the fact that the local dynamics are better captured in COSMO-355 REA6 due to the higher model resolution and the non-hydrostatic formulation. In fact, 356 over the North Sea, VIM from COSMO-REA6 differs by up to 20 mm from ERA5 (Fig-357 ure 6b and SM Figure S11). 358

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2.2 Contribution of wet atmosphere components

Water vapor content should not be neglected when calculating the VIM, as it can account for about 10% of the total variability (Figure 7, second column). However, the contribution of moisture in the form of rain, cloud water, cloud ice, and snow to VIM variability has not been studied so far. To estimate their contribution, we compare VIM from COSMO-REA6 of (i) dry air components only VIM_{dry} , (ii) dry air including wa-



Figure 4. RMSD of vertically integrated atmospheric mass in December 2007 over the EURO-CORDEX domain between 3-hourly data from (a) ERA5 and ERA-I (interpolated from 6-hourly fields), (b) ERA5 and CERRA, (c) ERA5 and COSMO-REA6, and (d) CERRA and COSMO-REA6.



Figure 5. Orography of the COSMO-REA6 model and location of investigated GRACE footprints (A-E) over Europe with a radius of 300 km.



Figure 6. (a) Area-averaged sub-monthly vertically integrated atmospheric mass of all models within footprint A (North Sea at the British coast) and (b) the respective differences in comparison to ERA5 in December 2007.

ter vapor VIM_{hum} , and (iii) moist air including rainwater, cloud water, cloud ice, and snow VIM_{wet} additionally to VIM_{hum} .

As expected, the contribution of the additional moisture fields (<1 mm with respect)367 to the RMSD) is small compared to that of water vapor (Figure 7a-c). However, dur-368 ing extreme events, the VIM variability due to the additional moisture fields can locally 369 reach values of several mm EWH, accounting for up to 2% of the variability of the to-370 tal VIM_{wet} (Figure 7d-f). As an example we show daily RMSD computations of VIM 371 on 5 December 2007 when heavy rainfall lead to flooding over Cyprus and Turkey (Fig-372 ure 7d-f). The signature of the additional moisture fields is clearly visible for the spe-373 cific day with a magnitude of about 3 mm in terms of RMSD. 374

Other examples are a heavy precipitation event that occurred in the Mediterranean Sea near Sicily on 4 June (SM Figure S20f); and an extreme event in the East of Great Britain on 25 June (SM Figure S20l). Particularly interesting is also the situation of the period between 19-23 December 2007, when heavy thunderstorms with supercells and hail struck Germany (Hechler & Bissolli, 2011). Figure S20i in SM shows contributions of additional moisture fields to VIM variations over large parts of Germany with a magnitude of up to 5.9 mm.

We show the impact of additional moisture fields in more detail for individual footprints over Europe. The contribution from water vapor to VIM variability reaches values of about 7 mm over Turkey (Figure 7g), and exceeds 15 mm in the first half of December 2007 over France (Figure S23b). In the case of the extreme event over Turkey the contribution of the additional moisture fields to VIM, which is not taken into account in the official AOD products, reaches 2 mm, i.e., 30 % of the contribution from water vapor.

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2.3 Contribution of continental hydrology

Unlike water vapor variability, short-term water storage changes at the land surface and the sub-surface are not considered as part of the standard AOD de-aliasing. Figure 8a shows sub-monthly variability of TWS globally based on output of the global hy-



Figure 7. VIM based on (i) only dry air density variations (VIM_{dry}) , (ii) dry air density and water vapor (VIM_{hum}) , and (iii) dry air, water vapor including rain, cloud water, cloud ice and snow (VIM_{wet}) . First row shows (a) standard deviation of VIM_{wet} , and the RMSD for (b) VIM_{hum} versus VIM_{dry} , and (c) VIM_{wet} versus VIM_{hum} for the entire December 2007. The second row shows standard deviation and RMSD for 5 December 2007. Exemplarily, (g) shows VIM differences for footprint E over Turkey (red dashed circle) in (e) and (f). Differences between VIM_{hum} and VIM_{dry} are in orange, and between VIM_{wet} and VIM_{hum} in blue.

drological WGHM model for December 2007. The standard deviation of submonthly vari-393 ability of TWS reaches values larger than 50 mm, and it is particularly high in equato-394 rial regions. Most significant variations are observed over the Amazon, the Ganges Delta, 395 and the Saint Elias Mountains (North America). Over Europe, sub-monthly TWS vari-396 ability varies regionally and the model suggests a magnitude between 5 and 30 mm. Dur-397 ing the month of June particularly high variations arose over the Scandinavian Moun-398 tains with a standard deviation of around 100 mm EWH likely due to snow variability 399 (SM Figure S25). In general, the sub-monthly TWS variability from the regional high-400 resolution model CLM (Figure 8b) shows more pronounced local features and larger RMSD 401 values than WGHM. Over Europe, differences between sub-monthly TWS variability from 402 the global model WGHM and the regional model CLM can be as large as total sub-monthly 403 TWS variability (Figure 8c). We suspect that the way how WGHM disaggregates me-404 teorological forcing data within the month is partly responsible for this. 405

Sub-monthly variability of TWS is characterized by linear evolution over several 406 days over Europe (Figure 8d). In December 2007 strong rainfalls occurred in large parts 407 of Europe leading to an increase in TWS. The storm event of 2 December 2007 is clearly 408 visible for the footprints over France and over the Alps (SM Figure S30). The two mod-100 els, WGHM and CLM, show generally the same patterns with a slightly larger range of 410 variability for CLM. In the Alpine region small ($\sim 1 \,\mathrm{mm}$) sub-daily variations of TWS 411 are visible in the CLM model (SM Figure S30). For extreme events, terrestrial water stor-412 age anomalies (TWSA) simulated by CLM show immediate responses with steep TWSA 413 curves, whereas WGHM simulates a smoother change in TWSA. 414

415 **3** Discussion

Future satellite gravity missions are expected to provide gravity fields and mass 416 change maps with higher resolution as compared to GRACE and GRACE-FO. However, 417 we will achieve substantial progress only if de-aliasing models can be improved. We sug-418 gest that it is time to reconsider, besides the tidal and non-tidal ocean components, also 419 the atmosphere contribution. Recent improvements on the atmospheric part of the AOD 420 product included (i) homogenizing surface pressure time-series from different NWP mod-421 els to a common reference orography, (ii) moving from ERA-I to ERA5, (iii) moving from 422 6-hourly resolution to 3-hourly resolution or even 1-hourly resolution, and (iv) updat-423 ing geometric, physical, and numerical approximations. Forootan et al. (2013) and Yang 424 et al. (2021) showed that a more realistic geometrical and physical shape of the Earth, 425 together with better approximation of latitude- and altitude-dependent gravity accel-426 eration in the vertical integration process, leads to differences in VIM of few 0.01 mm 427 in terms of geoid height without degree 0 and 1 coefficients, i.e. less than 1 mm EWH. 428 With AOD-RL06 the temporal resolution of the de-aliasing product was increased from 429 6-hourly data to 3-hourly data, which lead to differences in the geoid of about 0.1 mm 430 averaged over one year (Yang et al., 2021), i.e. few mm EWH. Further increasing the 431 temporal resolution to 1-hourly fields was assessed by Shihora et al. (2022) and Yang et 432 al. (2021). Yang et al. (2021) quantified the impact on the good to about 0.04 mm over 433 Europe and up to 0.08 mm in equatorial regions again averaged over one year. In line 434 with this, Shihora et al. (2022) suggested that the differences between hourly data and 435 hourly fields interpolated from 3-hourly fields are negligible for current applications. The 436 impact of changing the NWP model from ERA-I to ERA5 was found to alter geoid height 437 variability by about 0.1 to 0.3 mm, which corresponds to 1 to 5 mm EWH (Yang et al., 438 2021). For long-term consistency of the official de-aliasing product, Dobslaw (2016) mapped 439 the global reanalysis and analysis fields to a common reference orography, thus achiev-440 ing consistency of residual variability of surface pressure at the level of 10 mm. The im-441 plication is that the atmospheric component of the official AOD-RL07 product is well 442 tailored for current gravity missions, while for future missions with significantly more 443 demanding requirements, steps towards an even higher temporal resolution and better 444 long-term consistency will become necessary. 445



Figure 8. Standard deviation of sub-monthly variability of (a) WGHM- and (b) CLM-based water mass in December 2007. (c) RMSD between both models. (d) Sub-monthly terrestrial water mass variations of WGHM (blue) and CLM (orange) models for footprint E in December 2007.

All previous studies relied on global NWP models which are currently only avail-446 able at limited spatial resolution (>30 km), and they are all based on the hydrostatic as-447 sumption. In this paper, we investigated the use of regional high-resolution NWP model 448 output for future de-aliasing products, including also the non-hydrostatic model COSMO-449 REA6 for the European CORDEX region. Moreover, we analyzed contributions from 450 currently neglected additional atmospheric moisture fields and the terrestrial hydrology 451 to high-frequency mass variability. We interpret our results also in the context of future 452 gravity missions. 453

Over Europe, moving from 6-hourly to 3-hourly adds variability in the order of 3 454 to 10 mm EWH to VIM, whereas the effect of moving from 3-hourly to 1-hourly fields 455 is half as large and usually smaller than 5 mm (Figure 3). More specifically, averaged over 456 the study domain, VIM variability computed from ERA5 data increased by 4 mm EWH 457 for 6-hourly vs 3-hourly data, and by 1.6 mm EWH for 3-hourly versus 1-hourly data. 458 For COSMO-REA6 increased temporal resolution has the same effect like for ERA5 in 459 terms of magnitude of surface pressure variability, which means that better resolving small-460 scale processes within a regional model does not necessarily capture more variability at 461 shorter time scales. Yet, for ground-based applications like GNSS or terrestrial gravime-462 try a product with higher spatial and temporal resolution is an interesting option for ver-463 tical loading corrections. It should be kept in mind, that atmospheric tides were not re-464 moved here. However, their impact can be assumed to be not significant for our anal-465 ysis in the mid-latitudes (Yang et al., 2021). 466

Shihora et al. (2022) found, after careful removal of tidal signals, contribution of 467 hourly temporal resolution versus 3-hourly resolution of about 5 mm and decided based 468 on that analysis to keep the official AOD product at 3-hourly resolution. Likewise, Yang 469 et al. (2021) computed the impact from higher temporal resolution to reach values of about 470 0.2 mm geoid height (which corresponds to about 5 mm EWH) and suggested to test them 471 as AOD product in the gravity field processing chain. Indeed, the impact from increas-472 ing the temporal resolution from 3-hourly fields to 1-hourly fields on VIM is as large as the impact from updating the NWP model version from ERA-I to ERA5, and about 10 474 times as large as the impact from the refined vertical integration approach suggested by 475 Yang et al. (2021). 476

Recently, the MAGIC mission was decided to be launched as a successor of GRACEFO. Monthly fields are expected to possibly have an accuracy of 10 mm EWH (GRACEFO 20 mm) at 300 km spatial scale (Heller-Kaikov et al., 2023; Wiese et al., 2022). Future quantum gravimetry missions aim at providing monthly fields with an accuracy smaller
than 10 mm EWH and daily fields with an accuracy of few centimeters. Given these numbers, for future gravity remote sensing application we suggest to transition to 1-hourly
AOD products.

Will it be possible to improve computations of vertically integrated atmospheric 484 mass through the use of high-resolution models, which are currently available only as re-485 gional (limited-area) configurations? We find that differences between VIM from ERA-486 I, ERA5 and CERRA are all in the same range with mean RMSD values averaged over 487 the EURO-CORDEX area between 3 to 7 mm for the individual months. These differ-188 ences may arise from (i) differences in the forecast model applied by the NWP models, 489 (ii) assimilation of differently processed or new observation data sets, and (iii) different 490 spatial resolution and, thus, distinct parameterization schemes of smaller scale processes. 491 In contrast, the VIM variability from COSMO-REA6 differs noticeably from the other 492 models, with RMSD values that can be more than twice as large compared to those from 493 the other model comparisons. Our sensitivity studies show that this is not due to higher spatial resolution or the consideration of the additional moisture fields. Indeed, we as-495 sume that the larger differences can be explained by non-hydrostatic contributions from 496 the COSMO-REA6 model. Of course, minor differences may also arise due to different 497 choice of the orography. Bollmeyer et al. (2015) validated integrated water vapor from 498

different models against GPS observations over Germany, and found smaller biases for 499 COSMO-REA6 in comparison to ERA-I. At this stage, however, we can only conclude 500 that a de-aliasing product applying COSMO-REA6 over Europe will differ significantly 501 from current products, at a scale comparable to the transition from AOD1B-RL05 to AOD1B-502 RL06. It is an open question whether this will lead to better estimation of monthly grav-503 ity fields already for GRACE and GRACE-FO missions and needs to be assessed in fu-504 ture. To address this issue, Yang et al. (2018) evaluated different de-aliasing products 505 via computing level-1B pre-fit range-rate residuals. To be more precise, they investigated 506 how well measured range-rates and range-rates modeled based on different AOD prod-507 ucts fit together, and they found significant improvement of the atmospheric part of RL06 508 in comparison to RL05. As our product includes a refinement only over Europe, we ex-509 pect only minor impact on global gravity field estimation. 510

The impact of extreme events on GRACE observations has been investigated solely 511 looking at the contribution of hydrology in terms of TWS changes, either by analyzing 512 estimated sub-monthly gravity fields (Save et al., 2016; Nie et al., 2023; Zhang et al., 2023) 513 514 or by studying range-rate observations (Li et al., 2023; Han et al., 2021; Ghobadi-Far et al., 2022). In this paper, we moved a step further and investigated water mass anoma-515 lies that occur during extreme events before, in the form of precipitation, atmospheric 516 water flux increases hydrological storage. In other words, we seek to track the fate of wa-517 ter during its entire travel through the atmosphere-land continuum. In more detail, we 518 analyzed the contribution of atmospheric moisture fields that are currently disregarded 519 in AOD modeling (rain water, cloud water, cloud ice and snow) to VIM in the context 520 of extreme events. We found that locally these additional fields can reach values of few 521 mm EWH over Europe, which amounts to about 2% of total sub-monthly variability and 522 up to 30% of the contribution from water vapor. We speculate that for the catastrophic 523 July 2021 event that led to widespread flooding over Western Germany, these values were 524 exceeded by far. On small spatial scales the impact from additional moisture fields is com-525 parable to the impact of switching, e.g., from ERA-I to ERA5 over continental areas. 526 As we restricted our analysis to the EURO-CORDEX area, we cannot address the ques-527 tion here which magnitude the impact of this kind of events can reach over the oceans 528 during hurricane season, and over land in tropical cyclone regions. Future studies are 529 required to understand whether such localized strong events indeed map into monthly 530 gravity field estimates. 531

532 Only few studies investigated the impact of sub-monthly hydrological variability on GRACE range-rate observations (Zenner et al., 2014; Eicker & Springer, 2016; Ghobadi-533 Far et al., 2022) or on monthly GRACE fields (Thompson et al., 2004). Indeed, Eicker 534 and Springer (2016) showed that taking sub-monthly hydrological variability into account 535 reduces K-band range rate residuals, i.e. the fit of modeled range-rates to measured range-536 rates, in particular in equatorial regions where substantial water mass changes at short 537 time scales occur. At the time of the study conducted by Thompson et al. (2004) hy-538 drological models did not provide reasonable short-term variability and, thus, did not 539 improve monthly gravity field estimation. Ten years later, Luthcke et al. (2013) found 540 that GRACE Mascon solutions improved when taking into account continental hydrol-541 ogy based on output from the GLDAS-NOAH model (Ek et al., 2003; Rodell et al., 2004). 542 Kvas, Behzadpour, et al. (2019a) applied submonthly TWSA from the global Land Sur-543 face Discharge Model (LSDM; Dill, 2008) as additional de-aliasing product. However, 544 global hydrological models usually focus on climate change and water resource assess-545 ments but not on flood forecasting, and as a consequence processes that can lead to fast-546 moving water transports are not well represented; also these models sometimes lack a 547 suitable temporal resolution of the meteorological forcing data. 548

In our study, we compared high-frequency hydrological mass variations based on the output of two different hydrological models, the global model WGHM and the regional high-resolution model CLM3.5, which was forced with COSMO-REA6 data. As expected, sub-monthly TWS variability over Europe was found twice as large for CLM

 $(\approx 10 \text{ mm})$ compared to WGHM ($\approx 5 \text{ mm}$ over Europe), although large-scale patterns 553 of TWS evolution agree for the two models. To some extent the higher variability of CLM3.5 554 can be explained by the higher temporal and spatial resolution and more hydrological 555 processes being represented physically – some of differences will be averaged out at satel-556 lite height. Dobslaw et al. (2015) found high frequency variability of hydrological sig-557 nals of about 5 mm in Europe and of about 15 mm in regions near to the equator, and 558 isolated single events with fast water mass changes of few centimeters based on the LSDM 559 model. 560

Although the contribution from continental hydrology to sub-monthly mass vari-561 ations is about one magnitude smaller as compared to the atmospheric mass variability, we find that including hydrological variability into total mass variability would lead 563 to similar changes like increasing the temporal resolution from 6-hourly data to 3-hourly 564 data for the atmospheric model, or changing the model version from ERA-I to ERA5 565 (but obviously only over the continents). In this light, the potential for improving the 566 standard de-aliasing product by including sub-monthly hydrology should be assessed care-567 fully in particular for future missions. However, two important obstacles will arise here: 568 (i) current global hydrological models lack resolution and exhibit large uncertainties and 569 we showed in this paper that differences can be almost as large as the signal while few, 570 if any, high-resolution data sets exist at global scale. Alternatively, daily GRACE solu-571 tions could be assessed as additional de-aliasing product that accounts for hydrological 572 mass variability; however, such solutions inevitably lack the spatial resolution required 573 to resolve local strong events of significant size. (ii) We will have to avoid double book-574 keeping of water masses that transition from the atmosphere to the land surface and to 575 some part back into the atmosphere, involving phase changes. This can be achieved cor-576 rectly only by using a fully coupled atmosphere-hydrology model. In this paper, we en-577 sure consistency to first order over Europe by using output from the applied atmospheric 578 model as forcings for the hydrological model. 579

In summary, from the experiments covered in this study, we can conclude that the 580 impact on sub-monthly mass variability over Europe from (i) wet atmosphere contribu-581 tions, (ii) updating hydrostatic model versions, (iii) increasing the temporal resolution 582 of the NWP model from 6-hourly fields to 3-hourly fields, and (iv) adding TWS varia-583 tions modeled by a global hydrological model are of about the same size. We find that the impact from using a non-hydrostatic NWP model and a regional hydrological model 585 are considerably larger. Impacts from increasing the temporal resolution of the NWP 586 model from 3-hourly fields to 1-hourly are smaller, and only very limited impact arises 587 from considering additional moisture fields over Europe – which might be different for 588 equatorial regions. One major asset of our regional refinement is a first order mass con-589 sistency across atmosphere and land surface. Due to the altitude of the satellites, fur-590 ther increasing the spatial resolution of NWP and hydrological models may likely not 591 directly improve the de-aliasing process, but we suggest that it will indirectly lead to a 592 better representation of local processes. In contrast to previous studies, which assessed 593 anomalies with respect to a long-term mean, we assessed sub-monthly variability here. 594 Thus, when comparing to previous results, one should keep in mind that the RMSD values we compute here are most likely slightly smaller. 596

For current satellite gravimetry missions, where de-aliasing includes already wet 597 atmosphere contributions, a recent NWP model, and 3-hourly fields, at least the impact of hydrology on estimated gravity fields should be assessed carefully. For future more 599 sensitive satellite missions also 1-hourly fields might become important. Contributions 600 from moisture fields other than water vapor, i.e., rain water, cloud water, cloud ice and 601 snow, might not only map into estimated monthly gravity fields in the case of extreme 602 events, but are relevant also for other applications, such as loading computations for GNSS 603 studies, in particular if one is interested in cyclones (Zhan et al., 2021). We suggest that 604 these contributions will also be relevant if one aims at studying the evolution of extreme 605

precipitation, surface runoff and resulting flooding in integrative studies including space
 gravimetry data.

608 4 Outlook

We suggest that high-resolution de-aliasing data sets from refined atmospheric mod-609 els, such as provided in this study, should be assessed at the level of K-band and laser 610 ranging interferometry observations, both in view of improving the accuracy of monthly 611 gravity field model estimates and the "direct" evaluation of level-1 data for mass trans-612 port monitoring. Particularly interesting will be a systematic impact assessment of the 613 different contributions to mass variability shown in this contribution in the context of gravity field estimation. In future, we plan to include newly available reanalysis prod-615 ucts, e.g., the new version of COSMO-REA6, and analysis with the Icosahedral Nonhy-616 drostatic (ICON) model. A consistent global atmospheric data set will then be comple-617 mented by ERA5 outside of Europe. As long as no high-resolution non-hydrostatic global 618 NWP reanalysis exists, another option will be to nest multiple regional NWP products 619 simultaneously into one global model. Currently available regional high-resolution re-620 analyses focus on the Northern hemisphere and include the North American Regional 621 Reanalysis (NARR; Hunter et al., 2020) and the East Asia Reanalysis System (EARS; 622 Yin et al., 2023). Yet, we first need to understand if high-resolution non-hydrostatic mod-623 els indeed represent more realistic mass variability. 624

Regarding future enhanced de-aliasing products, we recommend that the contri-625 bution of additional atmospheric moisture fields, i.e., cloud water, rain water, cloud ice 626 and snow, should be investigated globally and possibly be included in the product. With 627 respect to sub-monthly TWS variability we need to further investigate how output from 628 different hydrological products influence the estimated gravity fields and which model 629 is favorable and guarantees near real-time availability. Furthermore, we argue that in 630 order to maintain mass consistency, improvements in Earth System Models (ESMs), which 631 provide fully coupled representation of water fluxes, will be of major interest for a con-632 sistent dealasing product. For instance, for ERA7 enhanced coupling with the land data 633 assimilation system is expected, which will provide new perspectives for de-aliasing and/or 634 signal separation. 635

We envision two different strategies evolving: On the one hand, given that we now 636 know very well that fast hydrological signals should be removed during de-aliasing, one 637 could argue that signals in all compartments that can be modeled with reasonable con-638 fidence should be removed. This could lead to a situation where the gravity missions would 639 "only" provide corrections to mass changes derived from compartmental models of some 640 advanced ESM, with the challenge of separating observations into hydrology, ocean, and 641 further contributions. On the other hand one could argue that model simulations and 642 data assimilation are better at separating the data; in this situation one might assim-643 ilate space gravimetric observations directly into a coupled ESM - the challenge here might 644 be that probably few coupled assimilation approaches could handle gravimetric data now. 645

Our study was constrained to the EURO-CORDEX domain. However, in partic-646 ular continental hydrology and contributions from additional moisture fields are much 647 larger in other regions. This means, a detailed analysis of these contributions, i.e. rain 648 water, cloud water, cloud ice and snow, on a global scale would be interesting. In future, 649 the sensitivity of gravity missions might also become interesting for improving NWP mod-650 els, e.g., via data assimilation. Another future implication is – whether a future grav-651 ity mission, e.g., with standard daily gravity field products and spatial resolution as to-652 day or better, would consider either dry air mass or water vapor as target variables -653 this entails that we would need to derive model-free estimates of these from the space-654 gravimetric data; this appears challenging but may be facilitated via separating the dom-655 inating time-scales. 656

Other geodetic techniques that require the removal or consideration of time vari-657 able gravity effects due to atmospheric mass variations are terrestrial gravimetry (Neumeyer 658 et al., 2004) and precise orbit determination (Cerri et al., 2010). Moreover, for GNSS 659 and satellite altimetry it is required to derive vertical loading corrections due to atmo-660 spheric pressure (Mémin et al., 2020; König et al., 2021) and to model the dynamic re-661 sponse of the sea surface (Andersen et al., 2018). While in particular for station observ-662 ables the requirements on spatial resolution may be much higher in the "near-zone", it 663 is clear that consistency across all geodetic techniques would be a welcome goal. 664

665 Appendix A Data

666

A1 Atmospheric reanalyses

NWP models represent various physical processes in the atmosphere and at the ocean and land surface, on complex topography, and propagate their impact on the temporal evolution of pressure, temperature, wind, water vapor fields and clouds and precipitation via some simplified form of the Navier-Stokes equations. Many processes like cloud formation cannot be explicitly resolved and their impact is approximated via parameterization schemes.

The global reanalysis ERA-I and its successor ERA5 were developed by the Eu-673 ropean Centre for Medium-Range Weather Forecasts (ECMWF) and provide global at-674 mospheric variables at multi-level fields. Compared to ERA-I, ERA5 has a higher hor-675 izontal grid resolution of 31 km (ERA-I: 79 km) and considers 137 pressure levels (ERA-676 I: 60 levels). The output frequency increased from 6-hourly data to 1-hourly data. Fur-677 thermore, the land surface model, i.e. the Tiled ECMWF Scheme for Surface Exchanges 678 over Land (TESSEL), was updated to HTESSEL, which incorporates an improved for-679 mulation of soil hydraulic properties and more realistic surface runoff generation. Ma-680 jor improvements of ERA5 with respect to ERA-I are related to the number and han-681 dling of observations; ERA5 incorporated more comprehensive satellites observations and 682 applied a new data-assimilation system. Additionally, ERA5 provides a 3-hourly uncer-683 tainty estimate for each parameter from a low-resolution 10-members ensemble, which 684 was not available for ERA-I. The improved resolution, wider vertical coverage, and en-685 hanced sampling of ERA5 enable a better representation of atmospheric patterns com-686 pared with previous global reanalyses, and show a considerable improvement especially in the troposphere with respect to ERA-I (Hoffmann et al., 2019; Hersbach et al., 2020) 688

Two high-resolution regional atmospheric reanalysis data sets are used in this study, 689 namely the COSMO-REA6 and CERRA reanalyses. COSMO-REA6 (Bollmeyer et al., 690 2015) is based on the non-hydrostatic Consortium for Small-scale Modeling (COSMO) 691 model (Doms, 2011). COSMO-REA6 applies a continuous nudging scheme (Schraff & 692 Hess, 2003) to assimilate meteorological observations from radiosondes, aircraft and weather 693 stations; the lateral boundary is constrained by ERA-I. The model grid of COSMO-REA6 is composed by an approximately 6 km staggered Arakawa-C horizontal grid and the terrain-695 following Gal-Chen hybrid height-based vertical coordinates with layer thickness increas-696 ing with altitude up to 22700 m corresponding to 40 hPa (Gal-Chen & Somerville, 1975; 697 Arakawa & Lamb, 1981; Bollmeyer et al., 2015). The terrain-following coordinate sys-698 tem reduces the complexity of formulating lower boundary conditions when surface ter-699 rain is involved, where the lowest surface of constant vertical coordinate becomes con-700 formal to the orography (Doms, 2011). Topography set-up for COSMO-REA6 is taken 701 from Global Land One-kilometer Base Elevation (GLOBE; Hastings & Dunbar, 1999; 702 Asensio et al., 2020) topography. 703

The Copernicus European Regional ReAnalysis (CERRA) (Schimanke et al., 2021) is a collaborative effort led by the Swedish Meteorological and Hydrological Institute in cooperation with the Norwegian Meteorological Institute and Météo-France. In contrast to the non-hydrostatic, fully compressible hydro-thermodynamic of COSMO-REA6, CERRA

adheres to the hydrostatic assumption in line with global climate models. CERRA is based 708 on the HARMONIE-ALADIN data assimilation system (Bengtsson et al., 2017; Termo-709 nia et al., 2018), includes the 3DVar deterministic CERRA reanalysis (CERRA-EDA) 710 with 10 ensemble members for the upper air, and a surface reanalysis CERRA-LAND, 711 which combines CERRA forcast fields and additional surface observation using an op-712 timal interpolation algorithm. Compared to COSMO-REA6, more observations are as-713 similated into CERRA such as satellite radiance observations and other non-conventional 714 observations, and the up-to-date ERA5 is applied for boundary conditions. The hori-715 zontal resolution of CERRA is 5.5 km (11 km for CERRA-EDA), the spatial domain cov-716 ers entire Europe and surrounding seas, and vertically extends up to 1 hPa height with 717 106 levels (El-Said et al., 2021; Z. Q. Wang & Randriamampianina, 2021). For this study, 718 we collected surface pressure, orography, air temperature, and specific humidity profiles 719 from both regional reanalysis. Additional moist air contents including specific rain wa-720 ter content and specific cloud water/ice content from COSMO-REA6 were extracted. CERRA 721 provides comparable variables but were not included in this study, because these vari-722 ables were only available for forecast runs. 723

724 A2 Hydrological models

The global water resources and use model WaterGAP v2.2d (Müller Schmied et 725 al., 2021) simulates daily water flows between all continental water storage compartments 726 except glaciers at a 0.5° grid. The model is calibrated against mean annual river discharge 727 of more than thousand gauging stations, located at the outlet of the respective drainage 728 catchments. Human water abstraction from groundwater aquifers and surface water bod-729 ies and return flows are modeled by a linked groundwater and surface water use model 730 (Döll et al., 2012, 2014). Climate forcings at daily resolution are obtained from WFDEI 731 (WATCH Forcing Data methodology applied to ERA-Interim data; Weedon et al., 2014). 732 In this study, we use daily TWS estimates that include canopy, snow, soil, groundwa-733 ter, lakes, man-made reservoirs, wetlands and rivers. 734

Over Europe, we use TWS outputs from the Community Land Model version 3.5 735 (CLM3.5; Oleson et al., 2008) to increase the temporal and spatial resolution of our hy-736 drological de-aliasing model in this region. CLM3.5 is a land-surface model that consists 737 of multiple modules representing biogeophysical and biogeochemical processes, dynamic 738 vegetation composition and structure, plant phenology, and the hydrological cycle. We 739 run this model over the European COordinated Regional Downscaling Experiment (CORDEX) 740 area at 0.11° (~12.5 km) resolution with hourly time steps. The model is forced with COSMO-741 REA6 6-hourly meteorological data (see Springer et al. (2019) for details on the set up). 742 TWS estimates are aggregated over soil water and soil ice at different levels, snow wa-743 ter, canopy water, and water in the unconfined aquifer with hourly output intervals. 744

745 Appendix B Methods

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B1 Water mass consistency across COSMO-REA6 and CLM

To obtain consistent atmospheric-hydrologic mass variations over Europe for the de-aliasing product, we use COSMO-REA6 to drive the CLM3.5 land surface model. Therefore, water masses in the hydrological model are consistent with atmospheric precipitation. However, the feedback from the land surface to the atmosphere is not yet consistent because evaporation in COSMO-REA6 differs from evaporation in our CLM3.5 run. We believe this is a second order effect, as evaporation fluxes generally evolve much more regularly than precipitation.

A substantial part of the sub-daily variability in atmospheric mass variability is due to solar forcing and thus corresponds to S_1 , S_2 tides (solar or thermal tides). It is wellknown that 6h fields do not permit an adequate resolution of the S_2 tide (Ray & Ponte, ⁷⁵⁷ 2003), and therefore it is generally recommended to remove S_1 and S_2 signals from the ⁷⁵⁸ atmospheric mass grids. Recent studies even consider minor tides (Yang et al., 2021; Shi-⁷⁵⁹ hora et al., 2022). However, the focus of this paper is on the influence of atmospheric ⁷⁶⁰ processes and their representation in models, and so we refrain from computing and re-⁷⁶¹ moving a dedicated solar tide model from each of the different cases. Instead, we pro-⁷⁶² pose that, for example, the tide model in Shihora et al. (2022) is simply removed con-⁷⁶³ sistently across different cases.

Below, we explain (i) how output from hydrostatic and non-hydrostatic atmospheric
models is integrated to compute VIM, (ii) how VIM was computed in previous studies
including the approach applied for the official atmospheric de-aliasing product RL07 and
(iii) how regional and global atmospheric and hydrological models are nested together
using a tapering approach.

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B2 Vertical integration

⁷⁷⁰ Spatiotemporal mass variations for de-aliasing products are represented by sets of ⁷⁷¹ spherical harmonic coefficients $\Delta C_n^m(t)$ and $\Delta S_n^m(t)$ (Stokes coefficients) up to degree ⁷⁷² *n* and order *m* (Hofmann-Wellenhof & Moritz, 2006; Boy & Chao, 2005):

$$\frac{\Delta C_n^m(t)}{\Delta S_n^m(t)} = \frac{(1+k'_n)a^2}{(2n+1)M_e} \int_0^{2\pi} \int_0^{\pi} \Delta I_n(\theta,\lambda,t) \tilde{P}_{nm}(\cos\theta) \left\{ \begin{array}{c} \cos\left(m\lambda\right)\\ \sin(m\lambda) \end{array} \right\} \sin\theta d\theta d\lambda,$$
(B1)

where (θ, λ) denote co-latitude and longitude of a given point in spherical coordinates, 774 \dot{P}_{nm} are normalized associated Legendre Functions, M_e is the average mass of the Earth, 775 a is a scaling factor that is typically chosen as the semi-major axis of the reference el-776 lipsoid (here the WGS84), and k'_n are the load Love numbers (Farrell, 1972; Dong et al., 777 1996). The factor $1+k'_n$ takes into account indirect elastic loading effects, additionally 778 to direct gravity effects due to mass changes. The vertical degree-dependent so-called 779 "inner integral" ΔI_n represents the changes over all masses in the Earth system in ra-780 dial direction 781

$$\Delta I_n(\theta, \lambda, t) = \left(\int_0^\infty \left(\frac{r}{a}\right)^{(n+2)} \rho(\theta, \lambda, r, t) dr \right) - \overline{I_n}(\theta, \lambda)$$

= $I_n(\theta, \lambda, t) - \overline{I_n}(\theta, \lambda),$ (B2)

where r is the geometric distance to the Earth's center of mass, $\rho(\theta, \lambda, r, t)$ are time variable density anomalies at a certain radius r, and $\overline{I_n}(\theta, \lambda)$ represents the monthly mean of VIM. As water storage variations occur within a thin layer (10-15 km) at the Earth's surface, the inner integral is reduced to representing changes in surface density, which then simplifies Equation B1 to a simple surface integral when we apply it to hydrological model output (Wahr et al., 1998).

In the case of the atmosphere, Boy and Chao (2005) showed that the impact from 789 3-D mass redistribution is non-negligible. Thus, for current AOD products the full con-790 tribution of atmospheric mass distribution is computed, including surface pressure ef-791 fects and contributions of density anomalies above ground. For global atmospheric mod-792 els and the CERRA reanalysis the hydrostatic assumption allows to introduce the re-793 lationship $\rho dz = -dp/q$ into Equation B2. For the official AOD products of RL06 and 794 RL07, atmospheric mass distribution is then computed at each time step (time index-795 ing dropped in the following) according to Dobslaw, Bergmann-Wolf, Dill, Poropat, and 796 Flechtner (2017) as 797

$$I_n(\theta,\lambda) = \int_0^\infty \left(\frac{r_e(\theta) + z(\theta,\lambda)}{a_{45}}\right)^{n+2} \frac{dp(\theta,\lambda,z)}{g_{45}},\tag{B3}$$

where a_{45} is the mean radius of the Earth computed at 45° latitude from semi-major axis

 g_{45} of the reference ellipsoid and g_{45} is the standard gravity defined by the World Meteo-

rological Organization (WMO). The ellipsoidal latitude-dependent radius $r_e(\theta) = a\sqrt{(1 - e^2 \sin(\theta)^2)}$

is computed from the eccentricity e of the reference ellipsoid. The orthometric heights $z(\theta, \lambda)$ were derived from the reference orography by a transformation of the geopotential heights that the model levels refer to (Dobslaw, Bergmann-Wolf, Dill, Poropat, & Flechtner, 2017). In discrete format, the vertical integral in Equation B3 is approximated by a summation over N vertical layers from the bottom level k = N to the top of the atmosphere (top-most layer) k = 1:

$$I_n(\theta, \lambda) = \sum_{k=1}^N \left(\frac{r_e(\theta) + z_k}{a_{45}}\right)^{n+2} \frac{\Delta p_k}{g_{45}},$$
(B4)

Here, the pressure difference $\Delta p_k = \Delta p_k(\theta, \lambda)$ related to model level k at height $z_k = z_k(\theta, \lambda)$ is computed between the two adjacent interfaces k-1/2 and k+1/2 using air pressure derived at the model boundaries according to

$$p_{k+1/2} = a_{k+1/2} + b_{k+1/2} p_S,\tag{B5}$$

with p_S being surface pressure, and $a_{k+1/2}$ and $b_{k+1/2}$ being model constants which define the vertical model grid at hybrid sigma levels.

For our study, we now have to distinguish between two cases:

(i) For global atmospheric models and the CERRA reanalysis, where the hydro static assumption applies, we follow Forootan et al. (2013) and compute the – in comparison to Equation B4 – refined inner integral of atmospheric contributions

$$I_n(\theta,\lambda) = \sum_{k=1}^N \left(\frac{r_e(\theta) + \xi + z_k}{a}\right)^{n+2} \frac{\Delta p_k}{g(\theta, z_k)},\tag{B6}$$

where *a* is the semi-major axis of the reference ellipsoid. Here, in contrast to the vertical integral computed for the official AOD product (Equation B4), the time-invariant height contribution from the geoid $\xi = \xi(\theta, \lambda)$, which is given as height in meter above an ellipsoid of revolution, is taken into account and latitude- and altitude-dependent gravity acceleration $g(\theta, z_k)$ instead of standard gravity is applied. Geometric height z_k is computed at each model level from geopotential height Φ_k^g following the conventions of the Office of the Federal Coordinator for Meteorology (OFCM, 1997),

$$z_k = \frac{r_e(\theta)\Phi_k^g}{\frac{g(\theta)r_e(\theta)}{g_{45}} - \Phi_k^g},\tag{B7}$$

and refers to the good ξ .

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(ii) For non-hydrostatic models, such as the regional COSMO-REA6 model applied in this study, we write the inner integral as

$$I_n(\theta,\lambda) = \int_0^\infty \left(\frac{r_e(\theta) + \xi + h + z}{a}\right)^{n+2} \rho(\theta,\lambda,z) dz.$$
(B8)

Here, the model height level z is referred to the orography $h = h(\theta, \lambda)$, and not to the geoid like in case (i), since COSMO-REA6 makes use of terrain-following coordinates (see Section A1). In discrete format, we sum again over N vertical layers according to

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$$I_n(\theta,\lambda) = \left(\frac{r_e(\theta) + \xi + h + z}{a}\right)^{n+2} \left[\frac{\rho_1}{g(\theta)} + \sum_{k=2}^N \rho_k(z_{k-1/2} - z_{k+1/2})\right].$$
 (B9)

The interface heights $z_{k+1/2}$ represent the height between model levels and $z_{N+1/2}$ equals to the surface height. The density ρ_k at each model level (k = 2, ..., N) is multiplied with the vertical distance of the two adjacent interfaces $z_{k-1/2}-z_{k+1/2}$. Additionally, the mass between z_1 and $z_{1+1/2}$ and the mass above the uppermost level need to be added (described by the first summand in Equation B9). In order to estimate this term, we again
 assume a hydrostatic equilibrium state in the top atmosphere:

$$\int_{z_1}^{\infty} \rho dz = \int_0^{p_1} \frac{dp}{g(\theta)} = \frac{p_1}{g(\theta)}$$
(B10)

Here, p_1 represents the mass above the center of the upper most level and is prescribed by ERA-Interim top level pressure. In order to avoid computing the vertical integral repeatedly for each spherical harmonic degree, we can approximate Equation B8 by a Taylor expansion described in the supplementary material (not applied for the results shown here).

For the dry atmosphere we use $\rho = \rho_{dry}$ of dry air, whereas the wet atmosphere density ρ_{wet} includes density variations from specific humidity q, specific rain water content q_r , specific cloud water content q_c , specific cloud ice content q_i and specific snow content q_s as

$$\rho_{wet} = \rho_{dry} \frac{1}{1 - (q + q_c + q_r + q_i + q_s)}.$$
(B11)

The computation of VIM makes use of a number of physical constants and parameters. These are summarized in Table B2.

Constant	Value / Unit	Description	Source / Ref.
a	$6378137.0{ m m}$	Semi major axis	WGS84 (NIMA, 2000)
1/f	298.257223563	Reciprocal of flatten- ing	WGS84 (NIMA, 2000)
ω	$7.292115 \cdot 10^{-5} \mathrm{rad/s}$	Angular rotation rate	WGS84 (NIMA, 2000)
GM_e	$3.986004418 \cdot 10^{14} \mathrm{m}^3/\mathrm{s}^2$	Geocentric gravita- tional constant	WGS84 (NIMA, 2000)
G_e	$6.6743 \cdot 10^{-11} \mathrm{m^2/(kg \ s^2)}$	Gravitational constant	Tiesinga et al. (2021)
g_{45}	$9.80665 \mathrm{m/s^2}$	Standard gravity	WMO (2021)
M_e	kg	Mass of Earth	derived from GM_e and G_e
k'_n	dimensionless	Load Love Number	H. Wang et al. (2012)
ξ	m	Geoid height	GOCO06s (Kvas, Mayer-Gürr, et al., 2019)
h	m	Mean orography	COSMO-REA6 (Doms, 2011)
$\Phi^g_{k+1/2}$	m	Altitude-dependent geopotential height	ERA-I / ERA5 (Hersbach et al., 2020) & Forootan et al. (2013, Eq. 10)
$g(\theta)$	m/s^2	Latitude-dependent gravity	Forootan et al. (2013, Eq. 17)
$g(\theta, z_{k+1/2})$	m/s^2	Latitude- and altitude-dependent gravity	Forootan et al. (2013, Eq. 18)

Table B1. Summary of physical constants and models used here

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B3 Overview on de-aliasing products

The computation of atmospheric de-aliasing products has been refined and extended during the last few years regarding the applied NWP model and its temporal and spatial resolution, the treatment of tides, the consideration of the inverse barometer effect, the temporal coverage, and the vertical integration procedure. Table B2 provides an overview on AOD products of the last decade. The refined vertical integration method, which is also applied in this paper, differs from the classical one applied in the official AOD products by considering the geoid undulation as component of the radial distance and by taking into account latitude-dependency of gravity acceleration (Yang et al., 2021).

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B4 Nesting of regional models into global models

Regional refinement of global atmospheric or hydrological models involves (i) regridding or coarse-graining to obtain homogeneous grids for the computation of spherical harmonic coefficients and (ii) tapering of the boundary between the global and regional model to avoid jumps which would translate to the so-called "Gibbs" phenomenon in spectral space characterized by ring patterns.

Before nesting, the temporal resolution of regional and global data sets needs to be unified. Common de-aliasing products used in the gravity field estimation process consist of 3-hourly fields of mass variations. Consequently, we temporally averaged hourly fields and applied temporal interpolation to the 6-hourly fields. Nesting of the regional model into the global model is performed for the hydrological models using TWS output of each model. For the atmospheric models, the inner integrals (Equation B2) are nested for each degree n separately. The monthly mean is removed for each data set before nesting.

Then, in the first step, the model grids are homogenized by re-gridding them to a 879 regular geographical grid of similar resolution like the rotated grid of the regional (high 880 resolution) model, i.e., 0.1° . In the second step, the transition area (Figure B1) is de-881 fined where regional and global models are smoothed to avoid data jumps. A width of 882 1.5° for the transition area and a filter half width increasing from 0° at the borders of 883 the transition area to 0.5° at the middle of the transition area for Gaussian filter Kernel leads to a satisfying smoothing effect at rather small differences between original and 885 filtered data sets. In a final step, global, regional, and filtered data sets are put together 886 (Figure B1) and converted into SH coefficients. As the nested gridded data sets are given



Figure B1. Nesting Procedure – Global model (ERA-I) without European CORDEX area (left), regional model (COSMO-REA6) given within the European CORDEX (right), both models and spatially filtered data in transition area (center) shown exemplarily for one snapshot of sub-monthly atmospheric mass.

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on equally spaced geographical grids they satisfy the Driscoll and Healy sampling theorem (Driscoll & Healy, 1994). Thus, they can be converted into spherical harmonic co-

orem (Driscoll & Healy, 1994). Thus, they can be converted in

		I		I		
Data set	AOD-RL05	AOD-RL06	AOD-RL07	ITG3D	HUST-ERA5	AHD-UB
Reference	Dobslaw et al. (2013)	Dobslaw, Bergmann-Wolf, Dill, Poropat, Thomas, et al. (2017)	Shihora et al. (2022)	Forootan et al. (2013)	Yang et al. (2021)	
NWP model	IFS	ERA-40/ERA- I/IFS	$\mathrm{ERA5}/\mathrm{IFS}$	ERA-I	ERA5	ERA-I/COSMO- REA6
lat. res.	$\sim 0.5^{\circ}$	$\sim 0.5^{\circ}$	${\sim}0.31^{\circ}$	$\sim 0.5^{\circ}$	${\sim}0.31^{\circ}$	$\sim 0.5^\circ \ / \ \sim 0.055^\circ$
degree/order	100	180	180	100	100	$180 \ / \ 360$
temp. res.	6-hourly	3-hourly	3-hourly	6-hourly	1-hourly	3-hourly
coverage	1989 - 2017	1976 - today	1975 - today	2002 - 2020	2002 - today	2007
tides	included	removed (12 fre- quencies)	removed (16 fre- quencies)	included	removed (12 fre- quencies)	included
temp. mean	2001 - 2002	2003 - 2014	2003 - 2014	2001 - 2002, monthly	2007 - 2014	monthly
IB effect	model	model	model	not corr.	model	not corr.
Vertical Integra- tion Method	classic	classic	classic	refined	refined	refined

 Table B2.
 Comparison of key attributes of available de-aliasing data sets.

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efficients using the Python toolbox pyshtools 4.1 (Wieczorek & Meschede, 2018), which implements the efficient Legendre transform algorithm described by Driscoll and Healy

⁸⁹² (1994).

⁸⁹³ Open Research Section

The de-aliasing data sets evaluated within this study are published as spherical harmonic coefficients at the German Research Centre For Geosciences (GFZ) data services repository via Mielke et al. (2023) with CC BY 4.0 licence. Please note that this paper refers to the updated data set version 2.0.

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Figure 1.



Figure 2.





Figure 3.



RMSD [mm EWH]

(a) ERA5(1h)

Std. [mm EWH]

(b) ERA5(6h) - ERA5(3h)

(c) ERA5(3h) - ERA5(1h)

Figure 4.

(a) ERA5 - ERA-I



(b) ERA5 - CERRA



(c) ERA5 - COSMO-REA6

(d) CERRA - COSMO-REA6



RMSD [mm EWH]

Figure 5.





Figure 6.





Figure 7.

- (a) VIM_{wet} , Dec.
- min:16.9, mean:102.6, max:212.6
- (b) VIM_{hum} VIM_{dry} , Dec.
- (c) $VIM_{wet} VIM_{hum}$, Dec.





(d) VIM_{wet} , 5 Dec.

- (e) VIM_{hum} VIM_{dry} , 5 Dec.
- (f) VIM_{wet} VIM_{hum} , 5 Dec.





Figure 8.





(d) Differences of hydrology for footprint E (Turkey)



Figure B1.



Supporting Information for "A regionally refined and mass-consistent atmospheric and hydrological de-aliasing product for GRACE, GRACE-FO and future gravity missions"

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Text S1: Resource-saving approximate computation of vertically integrated mass for non-hydrostatic models

For non-hydrostatic models, such as the regional COSMO-REA6 model applied in this study, the inner integral described by Eq. B8 of the main document can be rewritten as

$$I_n(\theta,\lambda) = \int_0^\infty \left(\frac{r_e(\theta) + \xi + h + z}{a}\right)^{n+2} \rho(\theta,\lambda,z) dz$$

$$= \left(\frac{r_e(\theta) + \xi + h}{a}\right)^{n+2} \int_0^\infty \left(1 + \frac{z}{r_e(\theta) + \xi + h}\right)^{n+2} \rho(\theta,\lambda,z) dz.$$
(1)

Here, the model height level z is referred to the orography $h = h(\theta, \lambda)$, and not to the geoid like in the case of the hydrostatic models, since COSMO-REA6 makes use of terrain-following coordinates (see Section A1 of the Appendix in the main publication). The integrated term involving z to the power of (n+2) is approximated to the first order using Taylor expansion, and can be dissected to the integration of ρdz and $\rho z dz$:

$$I_n(\theta,\lambda) \approx \left(\frac{r_e(\theta) + \xi + h}{a}\right)^{n+2} \left(\int_0^\infty \rho(\theta,\lambda,z) dz + \frac{(n+2)}{(r_e(\theta) + \xi + h)} \int_0^\infty \rho(\theta,\lambda,z) z dz\right).$$
(2)

As the degree n of the spherical harmonics is pulled out of the integral, the vertical integration now has to be performed only once and not for each degree separately. This means, all computations involving pressure and density fields can be prepared independent of the terms containing surface height and degree information. The first integral of Equation (2) is approximated in discrete form as

$$\int_0^\infty \rho dz \approx \frac{p_1}{g} + \rho_1(z_1 - z_{1+1/2}) + \sum_{k=2}^N \rho_k(z_{k-1/2} - z_{k+1/2}).$$
(3)

The vertical integral is approximated by a summation over N vertical layers from top to the surface. The density at each model level (k = 2, ..., N) is multiplied with the vertical distance of two adjacent interface $z_{k-1/2} - z_{k+1/2}$. The second integral is approximated

accordingly as

$$\int_0^\infty \rho z dz = \int_0^\infty \rho d\frac{z^2}{2} \approx \frac{p_1 z_1}{g} + \frac{p_1^2}{\rho_1 g^2} + \rho_1 \frac{z_1^2 - z_{1+1/2}^2}{2} + \sum_{k=2}^N \rho_k \frac{z_{k-1/2}^2 - z_{k+1/2}^2}{2}.$$
 (4)

Finally, the inner integral is computed according to

$$I_{n}(\theta,\lambda) \approx \left(\frac{r_{e}(\theta) + \xi + h}{a}\right)^{n+2} \left[\left(\frac{p_{1}}{g} + \rho_{1}(z_{1} - z_{1+1/2}) + \sum_{k=2}^{N} \rho_{k}(z_{k-1/2} - z_{k+1/2})\right) + \frac{(n+2)}{(r_{e}(\theta) + \xi + h)} \left(\frac{p_{1}z_{1}}{g} + \frac{p_{1}^{2}}{\rho_{1}g^{2}} + \rho_{1}\frac{z_{1}^{2} - z_{1+1/2}^{2}}{2} + \sum_{k=2}^{N} \rho_{k}\frac{z_{k-1/2}^{2} - z_{k+1/2}^{2}}{2}\right) \right].$$
(5)

To estimate the mass and inner product above the uppermost layer of COSMO-REA6 (first term on the right side of Equation 3 and first two terms on the right side of Equation 4), we again assume a hydrostatic equilibrium state in the top atmosphere:

$$\int_{z_1}^{\infty} \rho dz = \int_0^{p_1} \frac{dp}{g} = \frac{p_1}{g}$$
(6)
$$\int_{z_1}^{\infty} \rho z dz = \int_0^{p_1} z \frac{dp}{g}$$
$$= \frac{1}{g} \left(p_1 z_1 - \int_{z_1}^{\infty} p dz \right)$$
$$= \frac{p_1 z_1}{g} + \frac{R_L}{g^2} \int_0^{p_1} T dp$$
$$\approx \frac{p_1 z_1}{g} + \frac{R_L T_1 p_1}{g^2} = \frac{p_1 z_1}{g} + \frac{p_1^2}{\rho_1 g^2}.$$
(7)



Figure S1. Standard deviation of ERA5-based sub-monthly vertically integrated atmospheric mass in March 2007.



Figure S2. Standard deviation of ERA5-based sub-monthly vertically integrated atmospheric mass in June 2007.



Figure S3. Standard deviation of ERA5-based sub-monthly vertically integrated atmospheric mass in September 2007.



Figure S4. Range between minimum and maximum of ERA5-based sub-monthly vertically integrated atmospheric mass in December 2007.



Figure S5. Root mean square deviation (RMSD) of sub-monthly atmospheric mass in March 2007 over the EURO-CORDEX region for (a) 1-hourly ERA5, (b) 6-hourly (linearly interpolated) ERA5 minus 3-hourly ERA5, and (c) 3-hourly (linearly interpolated) ERA5 minus 1-hourly ERA5. In order to compute differences between time series with different temporal resolution, the lower-resolution time series is linearly interpolated to the higher-resolution time steps.



Figure S6. Root mean square deviation (RMSD) of sub-monthly atmospheric mass in June 2007 over the EURO-CORDEX region for (a) 1-hourly ERA5, (b) 6-hourly (linearly interpolated) ERA5 minus 3-hourly ERA5, and (c) 3-hourly (linearly interpolated) ERA5 minus 1-hourly ERA5. In order to compute differences between time series with different temporal resolution, the lower-resolution time series is linearly interpolated to the higher-resolution time steps.



Figure S7. Root mean square deviation (RMSD) of sub-monthly atmospheric mass in September 2007 over the EURO-CORDEX region for (a) 1-hourly ERA5, (b) 6-hourly (linearly interpolated) ERA5 minus 3-hourly ERA5, and (c) 3-hourly (linearly interpolated) ERA5 minus 1-hourly ERA5. In order to compute differences between time series with different temporal resolution, the lower-resolution time series is linearly interpolated to the higher-resolution time steps.



Figure S8. Root mean square deviation (RMSD) of vertically integrated atmospheric mass in March 2007 over the EURO-CORDEX domain between (a) 3-hourly ERA5 and 3-hourly ERA-I data (interpolated from 6-hourly fields), (b) 3-hourly ERA5 and 3-hourly CERRA data, (c) 3-hourly ERA5 and 3-hourly COSMO-REA6 data, and (d) 3-hourly CERRA and 3-hourly COSMO-REA6 data.



Figure S9. Root mean square deviation (RMSD) of vertically integrated atmospheric mass in June 2007 over the EURO-CORDEX domain between (a) 3-hourly ERA5 and 3-hourly ERA-I data (interpolated from 6-hourly fields), (b) 3-hourly ERA5 and 3-hourly CERRA data, (c) 3-hourly ERA5 and 3-hourly COSMO-REA6 data, and (d) 3-hourly CERRA and 3-hourly COSMO-REA6 data.



Figure S10. Root mean square deviation (RMSD) of vertically integrated atmospheric mass in September 2007 over the EURO-CORDEX domain between (a) 3-hourly ERA5 and 3-hourly ERA-I data (interpolated from 6-hourly fields), (b) 3-hourly ERA5 and 3-hourly CERRA data, (c) 3-hourly ERA5 and 3-hourly COSMO-REA6 data, and (d) 3-hourly CERRA and 3-hourly COSMO-REA6 data.



Figure S11. (a) Area-averaged sub-monthly vertically integrated atmospheric mass of all models within the investigated footprint A (North Sea at the British coast) and (b) the respective differences in comparison to ERA5 in June 2007.





Figure S12. (a) Area-averaged sub-monthly vertically integrated atmospheric mass of all models within the investigated footprint B (France) and (b) the respective differences in comparison to ERA5 in June 2007.



Figure S13. (a) Area-averaged sub-monthly vertically integrated atmospheric mass of all models within the investigated footprint C (Alps) and (b) the respective differences in comparison to ERA5 in June 2007.





Figure S14. (a) Area-averaged sub-monthly vertically integrated atmospheric mass of all models within the investigated footprint D (Mediterranean Sea) and (b) the respective differences in comparison to ERA5 in June 2007.



Figure S15. (a) Area-averaged sub-monthly vertically integrated atmospheric mass of all models within the investigated footprint E (Turkey) and (b) the respective differences in comparison to ERA5 in June 2007.





Figure S16. (a) Area-averaged sub-monthly vertically integrated atmospheric mass of all models within the investigated footprint B (France) and (b) the respective differences in comparison to ERA5 in December 2007.



Figure S17. (a) Area-averaged sub-monthly vertically integrated atmospheric mass of all models within the investigated footprint C (Alps) and (b) the respective differences in comparison to ERA5 in December 2007.




Figure S18. (a) Area-averaged sub-monthly vertically integrated atmospheric mass of all models within the investigated footprint D (Mediterranean Sea) and (b) the respective differences in comparison to ERA5 in December 2007.



Figure S19. (a) Area-averaged sub-monthly vertically integrated atmospheric mass of all models within the investigated footprint E (Turkey) and (b) the respective differences in comparison to ERA5 in December 2007.



Figure S20. Vertically integrated mass (VIM) was computed for three cases based on (i) only dry air density variations (VIM_{dry}), (ii) dry air density and water vapor contributions (VIM_{hum}), and (iii) dry air, water vapor, and additional moisture fields including rain water content, cloud water content, cloud ice content and snow content (VIM_{wet}). The first row shows (a) standard deviation of VIM_{wet}, and the root mean square deviation (RMSD) for (b) VIM_{hum} versus VIM_{dry}, and (c) VIM_{wet} versus VIM_{hum} for the entire June 2007. The other rows show the corresponding fields for 4, 21, and 25 June. Note the different scaling of the colorbars in each column. March 8, 2024, 12:47pm



Figure S21. Vertically integrated mass (VIM) was computed for three cases based on (i) only dry air density variations (VIM_{dry}), (ii) dry air density and water vapor contributions (VIM_{hum}), and (iii) dry air, water vapor, and additional moisture fields including rain water content, cloud water content, cloud ice content and snow content (VIM_{wet}). Shown is the (a) standard deviation of VIM_{wet}, and the root mean square deviation (RMSD) for (b) VIM_{hum} versus VIM_{dry}, and (c) VIM_{wet} versus VIM_{hum} for the 20 December 2007. Note the different scaling of the colorbars in each column.





Figure S22. Contribution of water vapor and the liquid moisture fields for selected footprints over Europe in June 2007. We show differences between VIM_{hum} and VIM_{dry} in orange, and differences between VIM_{wet} and VIM_{hum} in blue.

Days (June, 2007) -5 -10



Contribution of water vapor and the liquid moisture fields for selected footprints Figure S23. over Europe in December 2007. We show differences between VIM_{hum} and VIM_{dry} in orange, and differences between VIM_{wet} and VIM_{hum} in blue.

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Figure S24. Standard deviation of sub-monthly variability of (a) WGHM- and (b) CLM-based water mass in March 2007. (c) Root mean square deviation (RMSD) between both models.

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Figure S25. Standard deviation of sub-monthly variability of (a) WGHM- and (b) CLM-based water mass in June 2007. (c) Root mean square deviation (RMSD) between both models.



Figure S26. Standard deviation of sub-monthly variability of (a) WGHM- and (b) CLM-based water mass in September 2007. (c) Root mean square deviation (RMSD) between both models.

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Figure S27. Sub-monthly terrestrial water mass variations of WGHM (blue) and CLM (orange) models for selected footprints over Europe in March 2007.





Figure S28. Sub-monthly terrestrial water mass variations of WGHM (blue) and CLM (orange) models for selected footprints over Europe in June 2007.



Figure S29. Sub-monthly terrestrial water mass variations of WGHM (blue) and CLM (orange) models for selected footprints over Europe in September 2007.





Figure S30. Sub-monthly terrestrial water mass variations of WGHM (blue) and CLM (orange) models for selected footprints over Europe in December 2007.