

# First order controls of avulsion in river deltas

Octria Adi Prasajo<sup>1</sup>, Trevor B. Hoey<sup>2</sup>, Amanda Owen<sup>1</sup>, and Richard David Williams<sup>1</sup>

<sup>1</sup>University of Glasgow

<sup>2</sup>Brunel University London

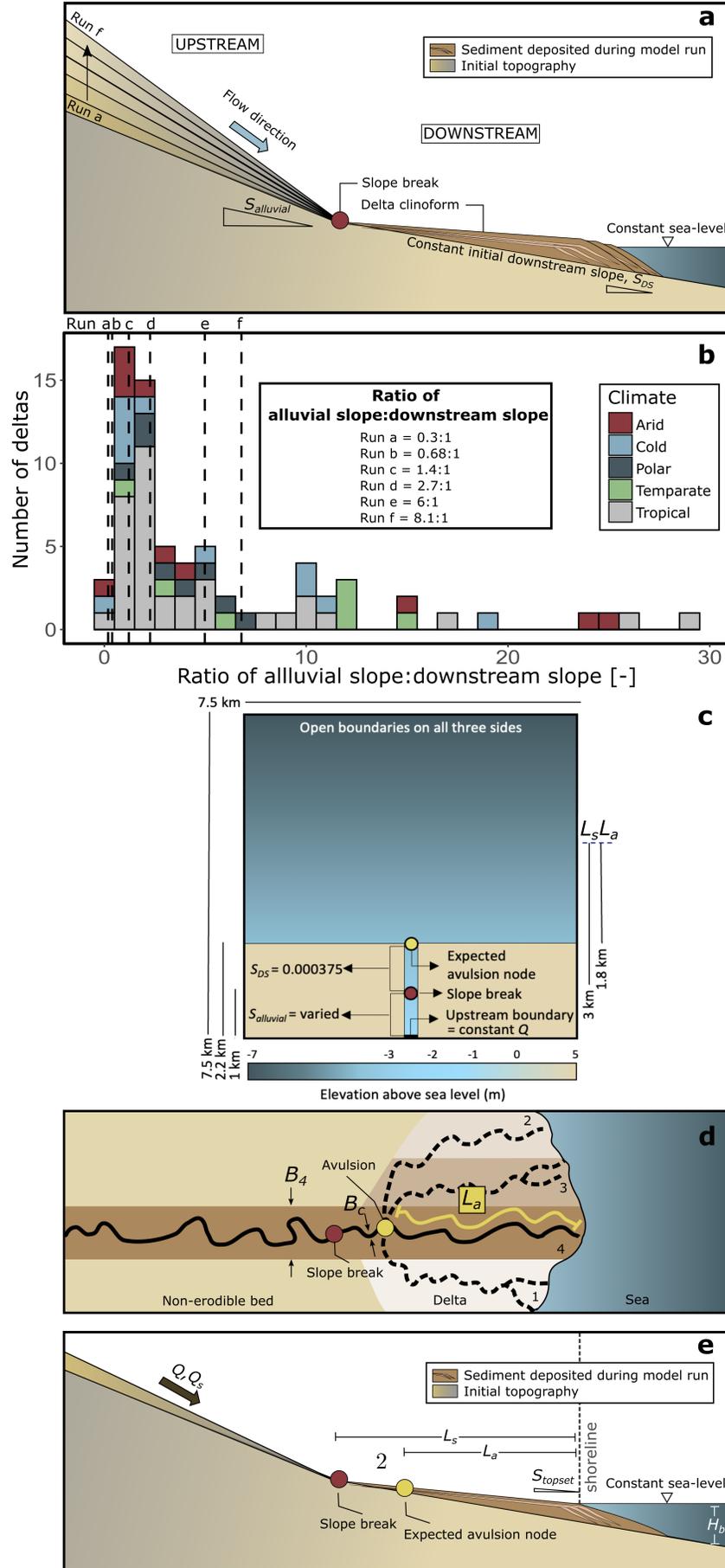
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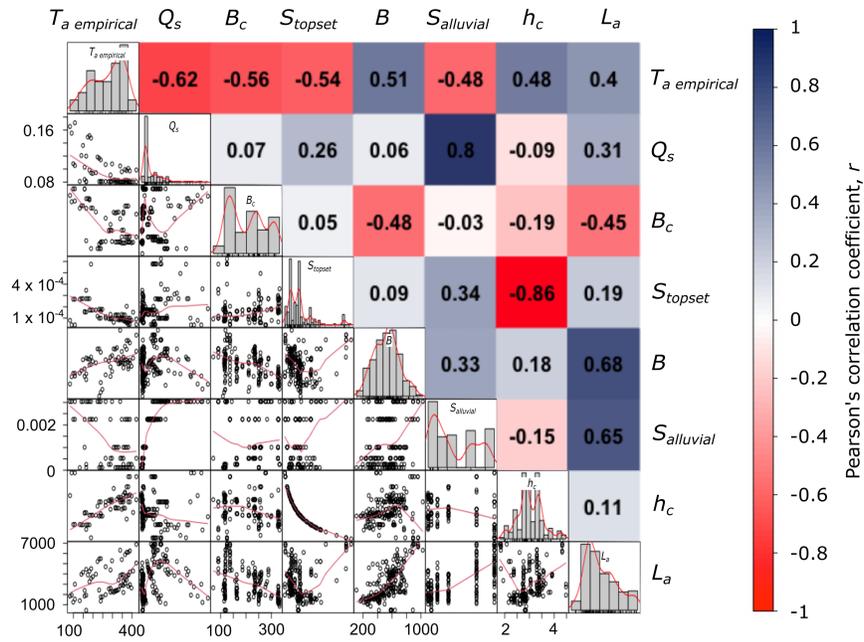
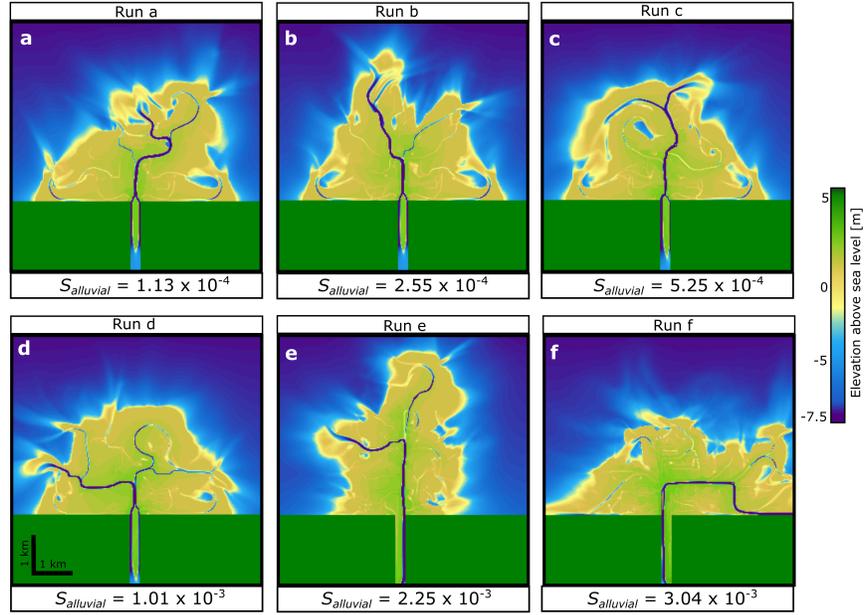
## Abstract

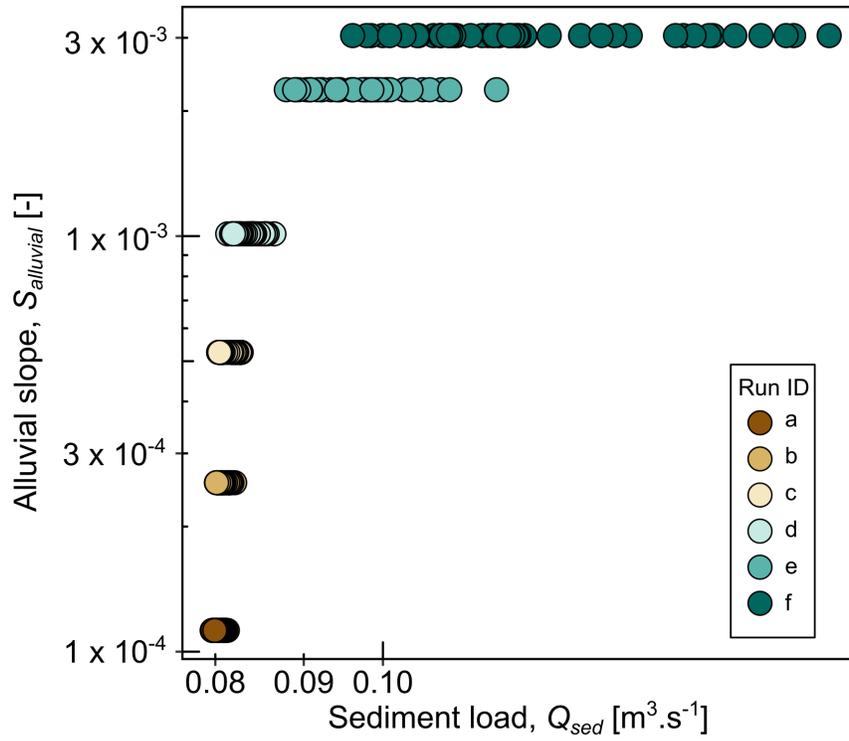
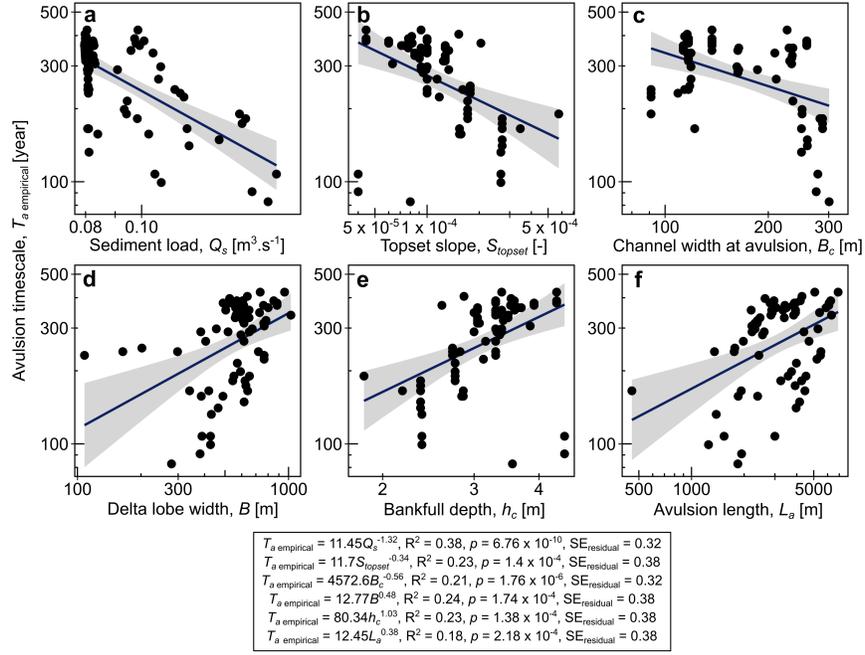
Changed hydrological regimes, sea-level rise, and accelerated subsidence are all putting river deltas at risk across the globe. Deltas may respond to these stressors through the mechanism of avulsion. Decades of delta avulsion studies have resulted in conflicting hypotheses that avulsion frequency and location are upstream (water and sediment discharge) or downstream (backwater and sea-level rise) controlled. In this study, we use Delft3D morphodynamic simulations to investigate the main controls over delta avulsion. Avulsion timing and location were recorded in six scenarios modelled over a 400-year period with varying alluvial slopes upstream of a delta slope break ( $1.13 \times 10^{-4}$  to  $3.04 \times 10^{-3}$ ) within a range representative global deltas. We measure several independent morphometric variables including avulsion length, delta lobe width, channel width at avulsion, delta topset slope and sediment load. Correlating these variables with the avulsion timescales observed in our model shows that avulsion timescale is mostly controlled by sediment load, which in turn is controlled by the alluvial slope upstream of a delta slope break. With higher stream power index in steeper alluvial slopes, more sediment can be carried within a channel, resulting in more frequent avulsions. Our results are consistent with the avulsion timescale derived from an analytical solution, 19 natural deltas and downscaled physical laboratory deltas. These results help mitigate delta avulsion risk by focusing management efforts on variables that primarily control avulsion in a river delta, but also induce further debate over whether sea-level rise may, or may not, trigger more avulsions in river deltas.

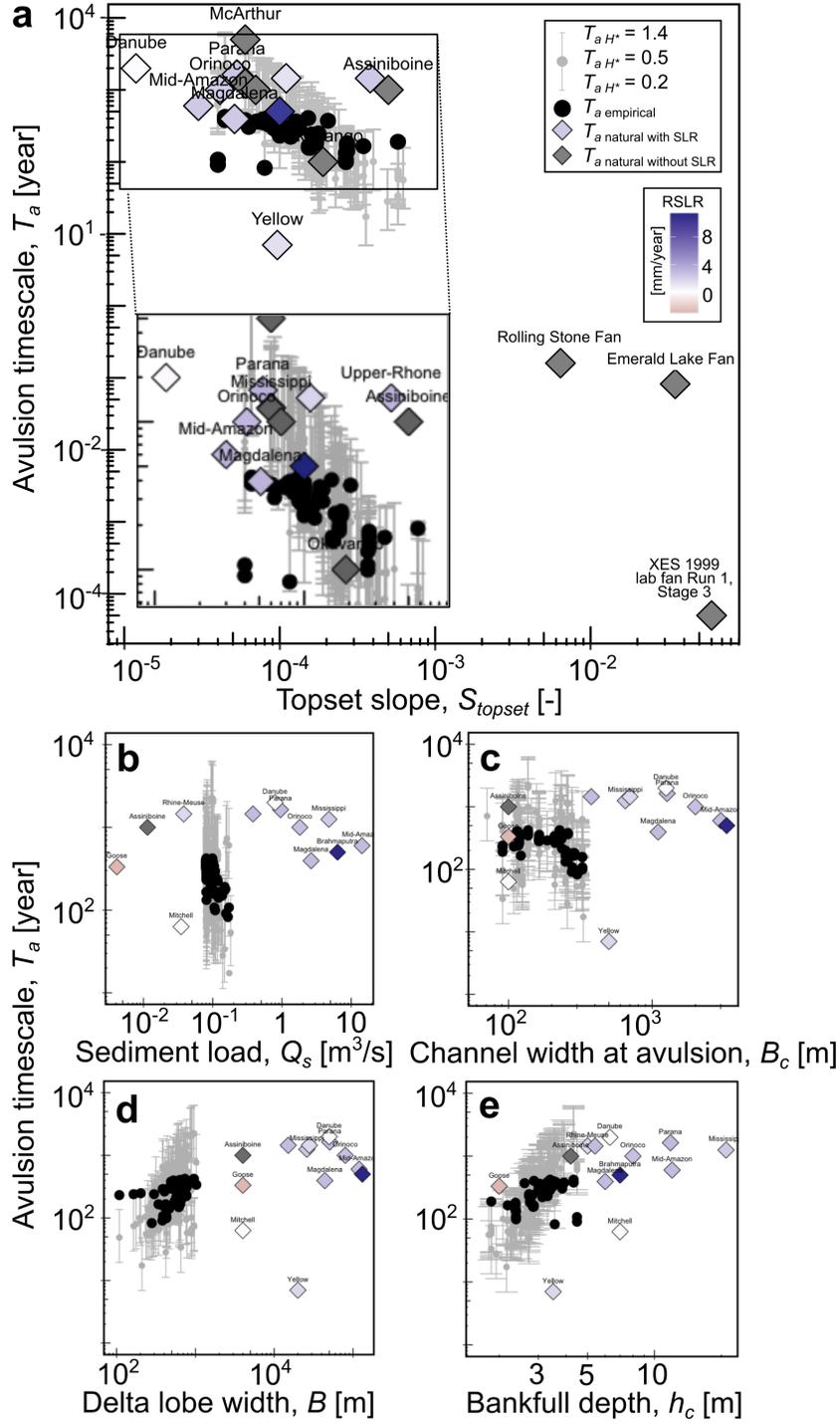
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# 1 First order controls of avulsion in river deltas

2 **O.A. Prasojo<sup>1\*</sup>, T. B. Hoey<sup>2</sup>, A. Owen<sup>1</sup> and R. D. Williams<sup>1</sup>**

3 *<sup>1</sup>School of Geographical and Earth Sciences, University of Glasgow, University Avenue,*  
4 *Glasgow, G12 8QQ, United Kingdom, Octria.prasojo@glasgow.ac.uk,*  
5 *Amanda.Owen@glasgow.ac.uk, Richard.Williams@glasgow.ac.uk*

6 *<sup>2</sup>Department of Civil and Environmental Engineering, Brunel University London, Uxbridge,*  
7 *UB8 3PH, United Kingdom, Trevor.Hoey@brunel.ac.uk*

## 8 **Highlights**

- 9     • The dominant variables controlling the avulsion timescale in river deltas are  
10     investigated through a Delft3D numerical model
- 11     • Results show sediment load, controlled by alluvial slope upstream of a delta plain,  
12     serves as the first-order control of delta avulsion frequency
- 13     • This supports the hypothesis of upstream forcing controlling delta avulsion timescale  
14     and location, rather than downstream, backwater or sea-level rise, controls
- 15     • Comparison with an analytical solution, natural systems and laboratory deltas all show  
16     consistency with our numerical results

## 17 **Abstract**

18 Changed hydrological regimes, sea-level rise, and accelerated subsidence are all putting river  
19 deltas at risk across the globe. Deltas may respond to these stressors through the mechanism of  
20 avulsion. Decades of delta avulsion studies have resulted in conflicting hypotheses that  
21 avulsion frequency and location are upstream (water and sediment discharge) or downstream  
22 (backwater and sea-level rise) controlled. In this study, we use Delft3D morphodynamic  
23 simulations to investigate the main controls over delta avulsion. Avulsion timing and location  
24 were recorded in six scenarios modelled over a 400-year period with varying alluvial slopes  
25 upstream of a delta slope break ( $1.13 \times 10^{-4}$  to  $3.04 \times 10^{-3}$ ) within a range representative global

26 deltas. We measure several independent morphometric variables including avulsion length,  
27 delta lobe width, channel width at avulsion, delta topset slope and sediment load. Correlating  
28 these variables with the avulsion timescales observed in our model shows that avulsion  
29 timescale is mostly controlled by sediment load, which in turn is controlled by the alluvial  
30 slope upstream of a delta slope break. With higher stream power index in steeper alluvial  
31 slopes, more sediment can be carried within a channel, resulting in more frequent avulsions.  
32 Our results are consistent with the avulsion timescale derived from an analytical solution, 19  
33 natural deltas and downscaled physical laboratory deltas. These results help mitigate delta  
34 avulsion risk by focusing management efforts on variables that primarily control avulsion in a  
35 river delta, but also induce further debate over whether sea-level rise may, or may not, trigger  
36 more avulsions in river deltas.

### 37 **Plain Language Summary**

38 River deltas grow by distributing sediment along their channel courses into a sea or a lake.  
39 During river delta growth, channels can abruptly change course, which can cause devastating  
40 floods to people, infrastructure and landscapes. The timing and the processes associated with  
41 this channel course switching are currently debated. Using a computer model, we create virtual  
42 river deltas to understand how their networks of channels develop and switch during delta  
43 growth over a 400-year period. We find that the steeper the topography upstream of a river  
44 delta, the faster this abrupt change of channel course occurs. This is because steeper channels  
45 will erode and transport more sediment than less steep channels. Our model predictions of  
46 channel course timings are consistent with those observed from 19 natural river deltas. We now  
47 better understand the timing and the main cause of abrupt channel changes on deltas, a finding  
48 that aids flood risk management in river delta environments.

## 49        **1. Introduction**

50            River deltas are home for ~339 million people worldwide, are hotspots for biodiversity,  
51 and crucial carbon sinks (Ericson et al., 2006; Hackney et al., 2020; Loucks, 2019; Shields et  
52 al., 2017; Syvitski & Saito, 2007). However, the geomorphic dynamism of river deltas has been  
53 altered by growing stressors such as change in hydrologic regimes, sea-level rise, and  
54 accelerated subsidence, putting human and other systems that rely on river deltas at  
55 considerable risk (Giosan et al., 2014; Syvitski et al., 2009; Tessler et al., 2015; Wallace et al.,  
56 2014). One mechanism by which deltas respond to these stressors is by avulsing, which creates  
57 additional, or relocated, flood risk. Many studies have proposed controls over avulsion  
58 frequency in river deltas (e.g. Aslan et al., 2005; Brooke et al., 2020; Edmonds et al., 2009;  
59 Kleinhans & Hardy, 2013; Nijhuis et al., 2015; Slingerland & Smith, 2004) and avulsion  
60 location correlates with backwater length, avulsion length and valley exit point location (Ganti  
61 et al., 2016a; Hartley et al., 2017; Prasojjo et al., 2022) but there is no consensus over which  
62 variable(s) is the most important factor contributing to delta avulsion frequency.

63            During avulsion, flow is abruptly diverted out of an established river channel into a new  
64 course on the adjacent floodplain or delta plain (Jones & Schumm, 2009; Slingerland & Smith,  
65 2004). When a delta channel avulses, the population and economic activities on the delta plain  
66 can be placed at risk. Avulsions may be considered rare, but this is partly due to anthropogenic  
67 impacts and natural systems often exhibit avulsion over decadal timescales, for example on  
68 average once every 12 years in the Yellow River Delta (Jerolmack, 2009) or 4 years in  
69 Sulengguole River, China (Li et al., 2022). Different scales of avulsion have been recognised:  
70 full, where a new course completely leaves its parent channel; and, partial in which only a  
71 portion of the flow is diverted (McEwan et al., 2023). There are also several styles of avulsion:  
72 annexation in which a pre-existing channel is reoccupied; incision, where a new channel is  
73 scoured into the floodplain surface as a direct result of the avulsion; and, progradation, where

74 extensive sediment deposition, such as a mouth bar, causes flow bifurcation and formation of  
75 a multi-channelled distributive network (Slingerland & Smith, 2004).

76 River deltas are initiated through repeated mouth bar deposition due to sudden  
77 expansion and deceleration of a sediment-laden jet of water entering relatively still water,  
78 usually a sea or lake (Bates, 1953; Edmonds et al., 2011; Kleinhans et al., 2013; Wright, 1977).  
79 Mouth bars grow in both upstream and downstream directions from the point of initiation,  
80 reach a height of c.0.4-0.6 of the initial flow depth, and stop growing once the sediment flux is  
81 advected around the mouth bar rather than accelerated over the bar (Edmonds & Slingerland,  
82 2007; Fagherazzi et al., 2015; Kleinhans et al., 2013). This is the point where avulsion by  
83 progradation starts in a river delta. Simultaneously, avulsion by incision takes place in the  
84 proximal parts of a delta plain when mouth-bar deposition and stagnation induce parent channel  
85 backfilling or in-channel aggradation, triggering an avulsion to create a smaller distributive  
86 channel network by breaching the channel levee (Ganti et al., 2016a; Prasajo et al., 2022). The  
87 most upstream point where a delta channel starts to avulse is correlated with the location of a  
88 break in bed slope (Prasajo et al., 2022; Ratliff et al., 2021), the limit of the backwater zone  
89 (Brooke et al., 2022; Ganti et al., 2016a), or the exit point from the river valley (Hartley et al.,  
90 2017).

91 A strong correlation has been identified from 105 global river deltas between the  
92 locations of breaks in delta slope and avulsion nodes (Prasajo et al., 2022). Consequently, it is  
93 hypothesised that the slope of the alluvial river upstream of a delta controls the frequency of  
94 avulsion in the proximal parts of deltas, with steeper alluvial slopes leading to more frequent  
95 avulsions. This control is due to greater sediment transport capacity on steeper slopes (Bagnold,  
96 1966). Hence, subject to sediment availability, steeper slopes transport more sediment per unit  
97 width into a delta plain. Assuming constant channel width, any reduction in energy slope across  
98 the delta plain leads to aggradation, the rate of which will be greater when upstream transport

99 capacity is higher, which in turn leads to increased avulsion frequency (Jerolmack & Mohrig,  
100 2007; Mohrig et al., 2000). Alternatively, lower alluvial slopes lead to lower sediment input  
101 rates and hence slower avulsion process.

102 To test this hypothesis, we use Delft3D morphodynamic simulation software to: 1)  
103 assess the effect of varying alluvial slopes upstream of a delta slope break on the avulsion  
104 timescale; and, 2) investigate what controls delta avulsion. Morphometric variables (delta lobe  
105 width, channel width at avulsion, avulsion length, topset slope, bankfull depth and sediment  
106 supply) were measured at every timestep during delta growth. These morphometric properties  
107 are considered to be independent variables that influence avulsion and so be correlated with  
108 avulsion timescales. This investigation aims to: (1) identify the first order controls of avulsion  
109 timescales from a suite of numerical model experiments with alluvial slope as the external  
110 forcing mechanism; (2) explain the mechanism of how the controlling variables control  
111 avulsion timescale; and, (3) compare the avulsion timescale from this numerical model to an  
112 analytical solution and natural river deltas. A robust understanding of these processes has  
113 practical implications due to their direct impact on coastal and inland flood risk on highly  
114 populated river deltas, as well as contributing to understanding of fundamental natural  
115 processes.

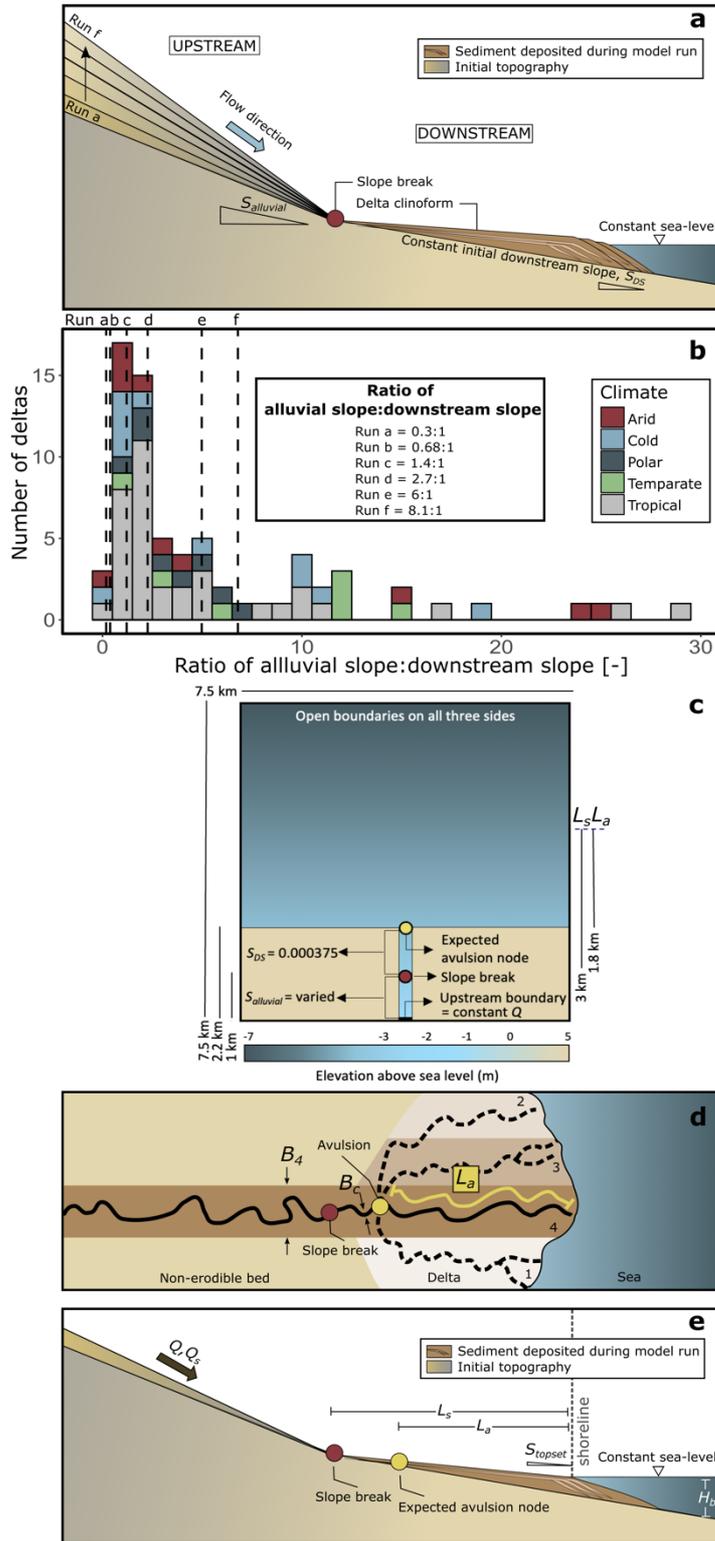
## 116 **2. Methods**

117 We designed a set of numerical experiments to model a natural scale river delta (7.5 x  
118 7.5 km, 300 by 300 computational cells, each 625m<sup>2</sup>) using Delft3D (v.4.04.02) software. We  
119 adopted physical parameters from the Delft3D river delta models from Edmonds & Slingerland  
120 (2010) and Caldwell & Edmonds (2014). Model bathymetry was designed to accommodate the  
121 six alluvial slopes defined below as our model scenarios.

122           **2.1.Scenario definition**

123           The model uses a range of alluvial slopes upstream of the delta's slope break ( $S_{alluvial}$ )  
124 (Fig. 1a), which are representative of natural river deltas (Fig. 1b). Ratios between  $S_{alluvial}$  and  
125 downstream delta slopes ( $S_{DS}$ ) measured from 105 global river deltas by Prasojo et al. (2022)  
126 were used to determine representative percentiles of this ratio (Fig. 1b; Table 1). Percentile 2.5  
127 to 75 from  $S_{alluvial}/S_{DS}$  ratio were used as our modelling scenario. These percentiles were then  
128 used to calculate model alluvial slopes using a constant initial downstream slope ( $S_{DS} =$   
129 0.000375) similar to that of the Atchafalaya Bay, Mississippi delta, Louisiana (Edmonds &  
130 Slingerland, 2010).

131



**Figure 1.** (a) Schematic diagram of the model design. The alluvial slope of each run was calculated from six percentiles from the alluvial slope-downstream slope ratios of modern river deltas shown in Fig. 1b. Initial downstream slope,  $S_{DS}$  is kept constant at 0.000375, the downstream slope of the modern Mississippi delta (Edmonds & Slingerland, 2010). (b) Distribution of the ratio between alluvial ( $S_{alluvial}$ ) and downstream ( $S_{DS}$ ) slopes from

105 modern river deltas distributed across five climate regions. Ratios used for numerical model runs are indicated by vertical dashed lines. (c) Plan view of the model design.  $L_s$  and  $L_a$  are slope break and avulsion lengths, respectively. The non-erodible bed at 5 m above sea level represents non-erodible bedrock. (d) Schematic diagram of a river delta showing avulsion location, inlet sediment supply ( $Q_s$ ), lobe width of each avulsion ( $B$ ), avulsion length ( $L_a$ ) and channel widths measured at avulsion ( $B_c$ ) modified from Chadwick et al. (2020). Numbers near the shoreline represent the number of delta lobes that were used to measure  $B$ ; e.g.  $B_4$  on (d) represents the width of the fourth lobe built. Schematic cross-section showing basin depth ( $H_b$ ) and topset slope ( $S_{topset}$ ). Parameters shown in Fig. 1d-e are measured at each timestep during delta growth.

## 132 2.2. Model setup

133 We use Delft3D software to model six scenarios. Delft3D is a physics-based model that  
 134 simulates hydrodynamics and morphodynamics (Edmonds & Slingerland, 2010; Caldwell &  
 135 Edmonds, 2014; Nienhuis et al., 2018a;b) and has been validated for a wide range of  
 136 environments, including self-formed river deltas (Edmonds & Slingerland, 2007, 2008;  
 137 Geleynse et al., 2011; Morgan et al., 2020; Nijhuis et al., 2015; Rossi et al., 2016; Williams et  
 138 al., 2016). Flow is computed using depth-averaged, nonlinear, shallow-water equations  
 139 obtained from three-dimensional Reynolds-averaged Navier-Stokes equations (Edmonds &  
 140 Slingerland, 2010). The modelled velocity distribution is then used to compute sediment  
 141 transport (only suspended load is applied in our model) and to update the bed elevation  
 142 according to divergence in sediment transport (Caldwell & Edmonds, 2014).

**Table 1.** Numerical modelling scenarios as defined in Fig. 1.

Run ID	Percentile from $S_{alluvial}$ to $S_{DS}$ ratio	Alluvial slope, $S_{alluvial}$	Initial downstream slope, $S_{DS}$	Ratio of alluvial slope to downstream slope
a	2.5	$1.13 \times 10^{-4}$	$3.75 \times 10^{-4}$	0.3
b	10	$2.55 \times 10^{-4}$	$3.75 \times 10^{-4}$	0.68
c	25	$5.25 \times 10^{-4}$	$3.75 \times 10^{-4}$	1.4
d	50	$1.01 \times 10^{-3}$	$3.75 \times 10^{-4}$	2.7
e	71	$2.25 \times 10^{-3}$	$3.75 \times 10^{-4}$	6.0
f	75	$3.04 \times 10^{-3}$	$3.75 \times 10^{-4}$	8.1

143

144 We adopted physical parameters from a previous synthetic self-formed river delta  
 145 numerical model ('scenario o') from Edmonds & Slingerland (2010) and Caldwell & Edmonds  
 146 (2014) (Fig. 1c). The model is rectangular with four boundaries, the incoming river discharge

147 being located at the ‘South’ boundary of the model and the other three boundaries set to 0 m  
148 elevation above sea level (Fig. 1c). The constant incoming river discharge, set at  $1050 \text{ m}^3 \cdot \text{s}^{-1}$ ,  
149 is uniformly distributed across the 250 m wide inlet channel, and inlet sediment discharge is in  
150 equilibrium with transport capacity. Various alluvial slopes are achieved by having various  
151 inlet channel bathymetry in each run while maintaining the receiving basin’s bathymetry. Our  
152 modelled deltas closely represent natural deltas because the discharge ratio and the differences  
153 in bed heights between bifurcating distributary channels follow similar ranges reported for  
154 natural deltas (Edmonds & Slingerland, 2010). Sea-level remains constant within the model,  
155 and no tide or wave effects are considered.

156         The model domain is 7.5 km x 7.5 km to avoid the delta plain extending across the  
157 model boundaries. We introduce a slope break 1 km from the inlet boundary to drive delta  
158 formation in the model’s initial bathymetry. Using the slope break-avulsion length scaling  
159 identified from global river deltas (Prasojo et al., 2022) the expected avulsion node location  
160 should emerge in each scenario at around 2.2 km from the inlet (Fig.1c). A constant sediment  
161 grain-size distribution is used throughout the model ( $D_{50} = 125\mu\text{m}$  with a normal Gaussian  
162 distribution, medium-grain silt  $D_{50} = 30\mu\text{m}$  is introduced as cohesive sediment), the critical bed  
163 shear stress for erosion =  $0.10 \text{ N} \cdot \text{m}^{-2}$ , and the model begins with 5 metres of fully mixed  
164 sediment. Other physical and numerical parameters were held constant across all scenarios  
165 (Table 2).

166         During an 18 days simulation, the model produces one output every 480 minutes. Hence  
167 at the end of simulation, the model stores 52 visualisation outputs (i.e. maps). Using a  
168 morphological scale factor (*morfac*) of 175, these 52 maps represent 3150 days (8.6 years) of  
169 prototype time with constant input discharge. Because bankfull discharge occurs for c.2% of  
170 time on average, 18 days of simulation thus represents around 430 years of ‘real’ time (i.e. 8.6  
171 years divided by 0.02).

**Table 2.** User-defined model parameters (adopted from Edmonds & Slingerland (2010); Caldwell & Edmonds (2014)).

Parameter	Value	Units
Grid size	300 x 300	cells
	7.5 x 7.5	km
Cell size	25 x 25	m
Run duration	18	days
Basin bed slope (downstream of slope break)	0.000375	(-)
Initial channel dimension (width x depth)	250 x 2.5	m
Upstream non-erodible bed elevation	5	m
Initial channel length upstream of slope break	1000	m
Initial avulsion length from the expected shoreline	1800	m
Water discharge	1050	m <sup>3</sup> .s <sup>-1</sup>
Constant water surface elevation at downstream open boundary	0	m
Initial sediment layer thickness at bed	5	m
Number of subsurface stratigraphy bed layers	1	(-)
Computational time step	0.2	min
Output interval	480	min
Morphological scale factor	175	(-)
Spin-up interval	720	min

### 172 2.3.Surface metric

173 When the model reaches an equilibrium (i.e. model's inlet shows constant sediment  
174 discharge and channel depth at timestep > 200 years), we begin the morphometric and avulsion  
175 timescale measurement. Avulsions throughout 18 days of simulation were empirically  
176 observed at the expected avulsion node location (Fig. 1c) or at the 'valley exit' in which an  
177 inlet channel meets the open water in the model. Avulsion in the model is defined when a  
178 distributive channel produced during delta formation changes its course and deposits a new  
179 delta lobe. Hence, we limit our avulsion observations to those caused by progradation or  
180 incision and not by annexation, which may occur on the more distal delta plain (Slingerland &  
181 Smith, 2004). Consequently, only the most upstream avulsions are observed in this study.  
182 Every time avulsion is observed in the model, computational timescale is noted and converted  
183 to a 'real' time as  $T_{a\text{ empirical}}$ .

184 Measured surface metrics are adopted from an analytical solution for avulsion timescale  
185 as a function of delta lobe width ( $B$ ), channel width at avulsion ( $B_c$ ), avulsion length ( $L_a$ ), basin  
186 depth ( $H_b$ ), magnitude of relative sea-level rise ( $z$ ), topset slope ( $S_{topset}$ ), bankfull depth ( $h_c$ ) and

187 sediment supply ( $Q_s$ ) (Eq. 4 from Chadwick et al., 2020). Avulsion length, delta lobe width,  
 188 channel width at avulsion and delta topset slope were measured over 52 maps. The delta lobe  
 189 width ( $B$ ), channel width at each avulsion node ( $B_c$ ) and avulsion length ( $L_a$ ) were measured in  
 190 QGIS from the georeferenced images produced by Delft3D (Fig. 1d, Table S1). Delta lobe  
 191 width ( $B$ ) is measured as the maximum width of each lobe, while avulsion length ( $L_a$ ) is  
 192 measured along the longest channel from the shoreline to the most upstream avulsion node  
 193 located at the ‘expected avulsion node’ mentioned in Fig. 1c. Topset slope ( $S_{topset}$ ) was  
 194 calculated by linear regression through topset elevations along a longitudinal cross-section  
 195 located through the centre of the model from the delta shoreline to the delta slope break  
 196 introduced in the model (i.e. located 1 km from the model’s South boundary) (Fig. 1e).  
 197 Sediment supply ( $Q_s$ ) at the channel inlet was obtained from a Delft3D visualisation software,  
 198 QUICKPLOT (v2.60.65942).

199 Bankfull depth ( $h_c$ ) was calculated using Eq. 1 (Parker, 2004).

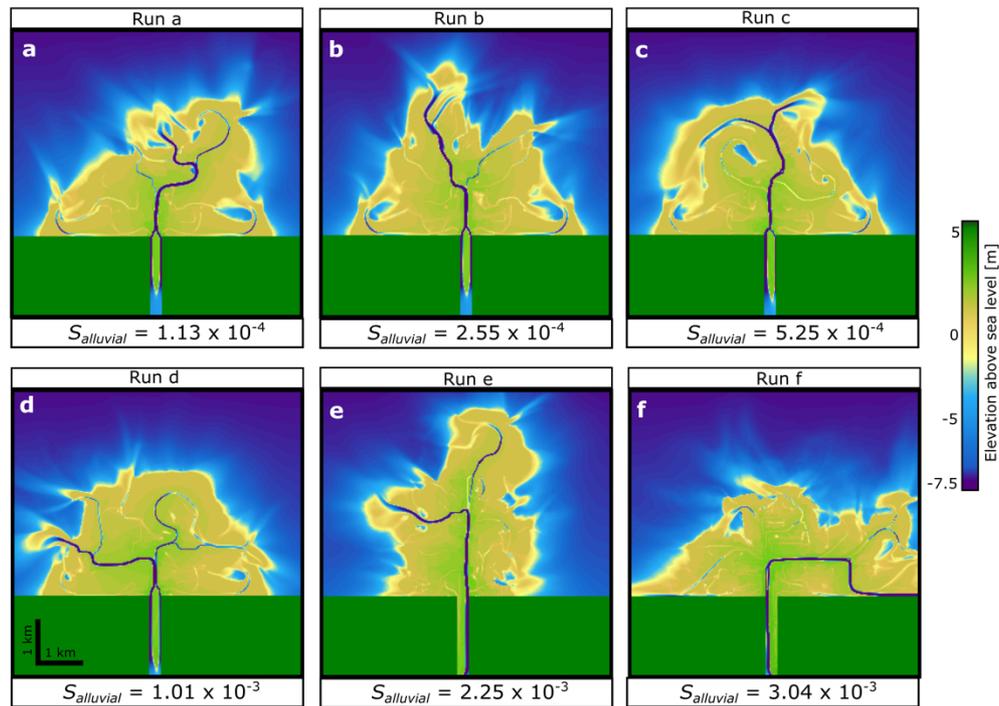
$$h_c = \left( \frac{C_f Q^2}{g B_c^2 S_{topset}} \right)^{\frac{1}{3}} \quad (1)$$

200  $C_f$  is defined as bed friction coefficient [-] = 0.002 for large lowland rivers (Parker et al., 2007),  
 201  $Q$  = bankfull discharge [ $\text{m}^3 \cdot \text{s}^{-1}$ ] = 1050  $\text{m}^3 \cdot \text{s}^{-1}$ ,  $g$  = gravitational acceleration [ $\text{m} \cdot \text{s}^{-2}$ ] = 9.81  $\text{m} \cdot \text{s}^{-2}$ ,  
 202  $B_c$  = channel width at avulsion node [m], and finally  $S_{topset}$  = topset slope [-].

203 The avulsion timescale empirically observed at each time an avulsion occurred ( $T_a$   
 204 empirical) was correlated with all the measured morphometric variables (e.g.  $Q_s$ ,  $L_a$ ,  $B_c$ ,  $B$ ,  $S_{topset}$ ,  
 205  $S_{alluvial}$ , and  $h_c$ ) from all 52 maps. Scatter plots and Pearson correlation coefficients ( $r$ ) were  
 206 used to assess the linearity of relationships and potential dependencies between all variables.

207 **3. Results**

208 Fig. 2 shows the morphology of the deltas in each scenario at the final timestep. Overall,  
 209 the different alluvial slopes produce delta plains that exhibit different shoreline configurations,  
 210 different numbers of active distributary channels and slightly different delta plain sizes. One  
 211 delta plain reached the model boundary (Run f) and this scenario was repeated with a larger  
 212 domain size (Fig. S1) and the avulsion timescales were observed from this larger domain.

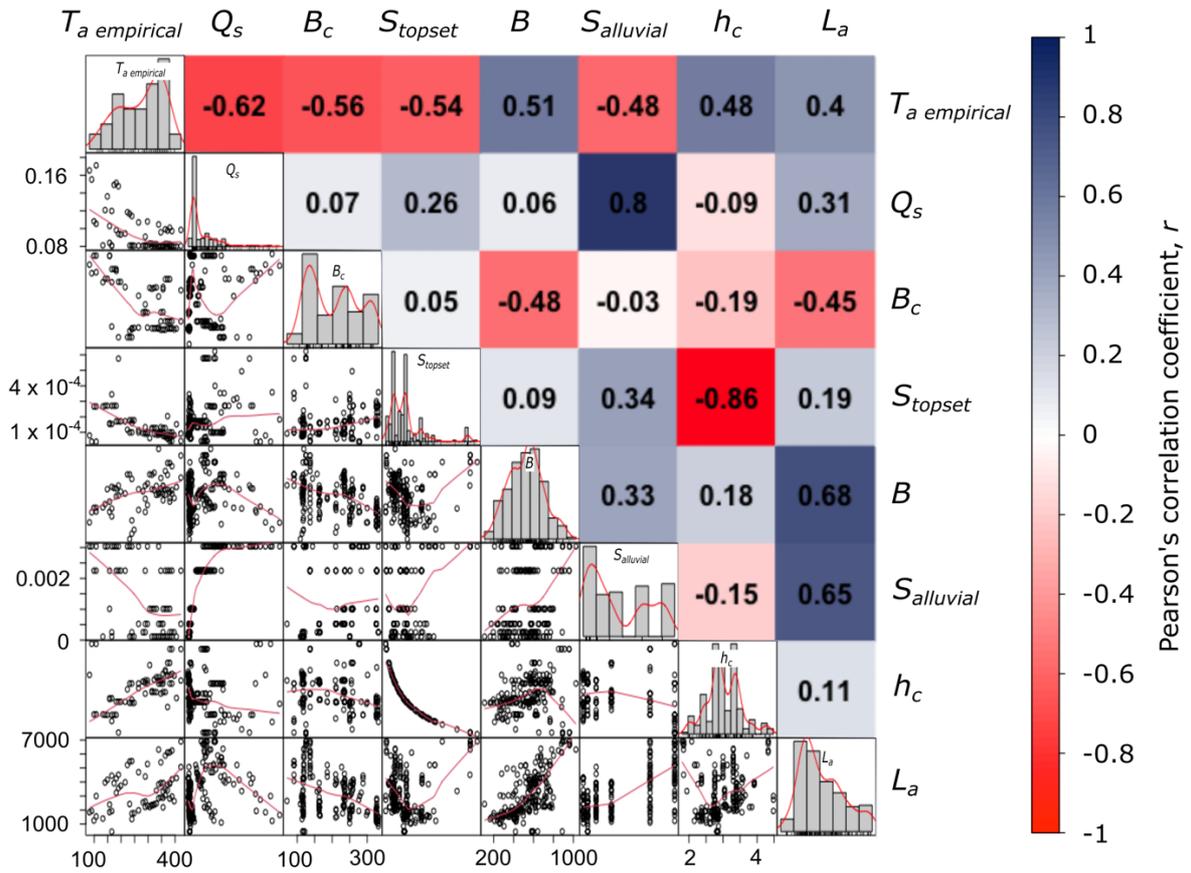


**Figure 2.** (a-f) River deltas for each run at the final simulation timestep. *Run f* was repeated with a larger (12.5 x 12.5 km) model size (Fig. S1) to avoid the delta plain reaching the model's boundary. Morphometric measurements for Run f were made on this larger model size.

213 Avulsion timescale observed in the model occurs as quickly as 82.7 years with the  
 214 longest of 428 years after the model starts. Skewed distribution of avulsion timescale can be  
 215 observed in most scenario runs (Fig. S2) with overall median value of 297.7 years and mean  
 216 value of 278.9 years. Run a and b do not show significant difference in avulsion timescale  
 217 range. However, later runs show a decrease of avulsion timescale as the median value of  
 218 avulsion timescale was found to be 340 years in run c, in contrast to 175 years in run f (Fig.  
 219 S2).

220 Fig. 3 shows correlations between observed avulsion timescales in the model ( $T_{a\text{ empirical}}$ )  
 221 and the independent morphometric variables measured in each timestep and relationships  
 222 between those independent variables.  $T_{a\text{ empirical}}$  has a high correlation with sediment supply,  $Q_s$   
 223 ( $r = -0.62$ ;  $p = 6.76 \times 10^{-10}$ ).  $T_{a\text{ empirical}}$  is also correlated with channel width at avulsion ( $B_c$ ;  $r$   
 224  $= -0.56$ ,  $p = 1.76 \times 10^{-6}$ ), topset slope ( $S_{\text{topset}}$ ;  $r = -0.54$ ;  $p = 1.4 \times 10^{-4}$ ), delta lobe width ( $B$ ;  $r =$   
 225  $0.51$ ,  $p = 1.74 \times 10^{-4}$ ), alluvial slope ( $S_{\text{alluvial}}$ ;  $r = -0.48$ ,  $p = 8.45 \times 10^{-4}$ ), bankfull depth ( $h_c$ ;  $r =$   
 226  $0.48$ ;  $p = 1.38 \times 10^{-4}$ ) and avulsion length ( $L_a$ ;  $r = 0.4$ ,  $p = 2.18 \times 10^{-4}$ ).

227 Sediment load ( $Q_s$ ) is driven by the alluvial slope ( $S_{\text{alluvial}}$ ) that is independently defined  
 228 in the model (Fig. 1a), leading to a correlation  $r = 0.8$  between these variables (Fig. 3). A  
 229 morphological consequence of sediment load is that bankfull depth ( $h_c$ ) is highly correlated ( $r$   
 230  $= -0.86$ ) with topset slope ( $S_{\text{topset}}$ ) as defined from Eq. 1. Moderate correlations are found  
 231 between other morphometric variables, such as  $B_c$ - $B$ ,  $B_c$ - $L_a$ ,  $B$ - $L_a$ ,  $S_{\text{alluvial}}$ - $L_a$ , because as a delta  
 232 grows the delta plain and its' constituent channels and islands enlarge in an allometric manner,  
 233 as observed in natural, physical laboratory and numerical deltas (Wolinsky et al., 2010).  
 234 Additionally, alluvial slope seems to control the avulsion length ( $r = 0.65$ ,  $p < 2.2 \times 10^{-16}$ ),  
 235 consistent with the findings from a survey of global river deltas (Prasojo et al., 2022).



**Figure 3.** Pearson correlations between avulsion timestep ( $T_a$  empirical) and independent morphometric variables with  $N = 233$  along with their distributions and correlations. Units on this figure are years for  $T_a$  empirical,  $m^3 \cdot s^{-1}$  for  $Q_s$ , meters for  $B_c$ ,  $h_c$  and  $L_a$ , consecutively.  $S_{topset}$  and  $S_{alluvial}$  are dimensionless. Note that  $h_c$  is autocorrelated with  $S_{topset}$  as shown in Eq. 1.

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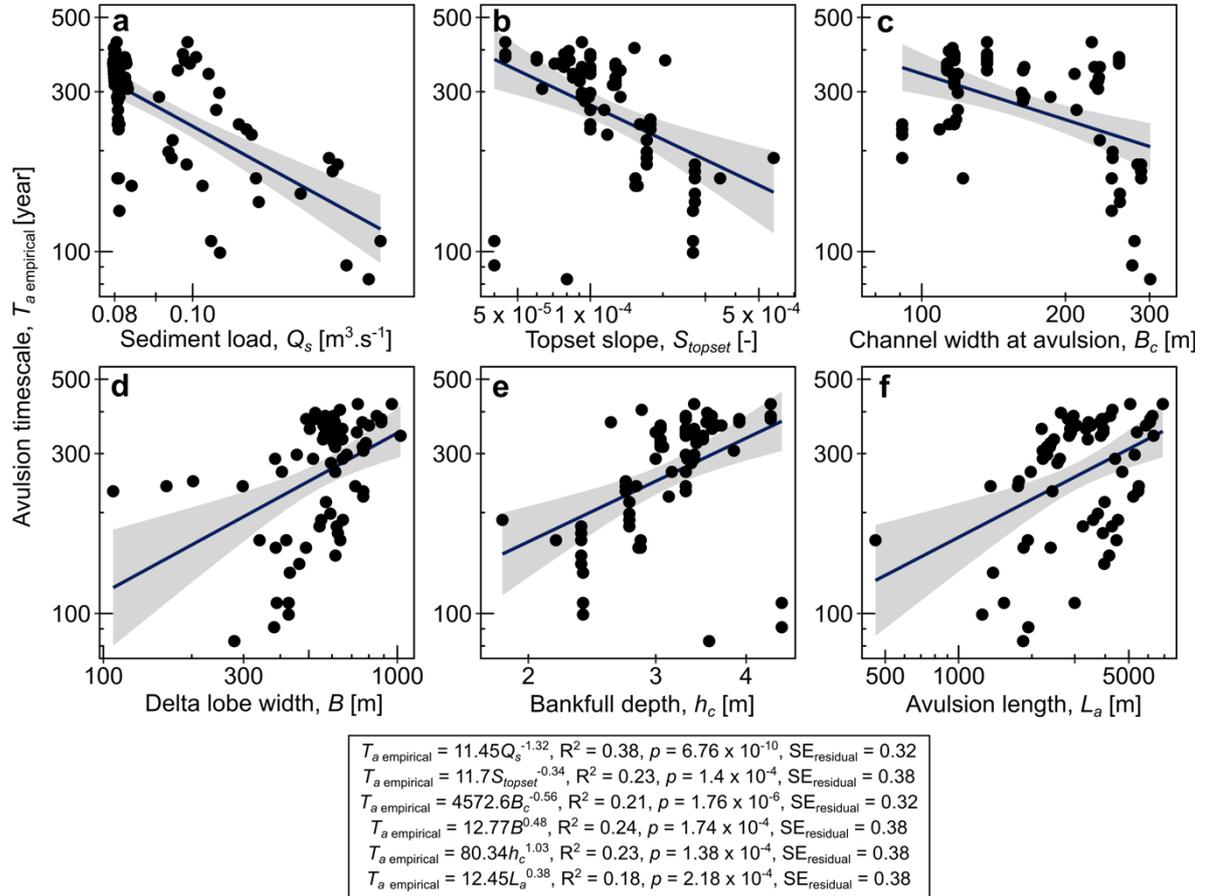
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Fig. 4 shows the data from the model and ordinary least square regressions for the highest correlations in Fig. 3. The regression relationships are statistically significant and have narrow confidence bands (grey shaded areas in Fig. 4), although the data exhibit scatter and some clustering. Avulsion timescale is inversely correlated with sediment load (Fig. 4a), which may reflect high aggradation rates when sediment loads are greater. This process link becomes more apparent when considering topset slope (Fig. 4b) which is a consequence of sediment load and aggradation. The correlation between channel width and slope that results from flow energy in turn produces a negative relationship between avulsion timescale and channel width at the avulsion location (Fig. 4c). Avulsion timescales increase with the size of the delta, represented by lobe width (Fig. 4d), bankfull depth (Fig. 4e) and avulsion length (Fig. 4f). Both

247 lobe width and avulsion length plots contain separate clusters of data points that may indicate  
 248 alternative patterns of morphological adjustment. The weak positive correlation between  
 249 avulsion length and timescale challenges the hypothesis that backwater length controls avulsion  
 250 timescale (Chadwick et al., 2020; Chatanantavet et al., 2012; Ganti et al., 2016a,b).



**Figure 4.** Regressions of independent morphometric variables against avulsion timescale ( $T_{a \text{ empirical}}$ ) observed in model runs with 95% confidence band in grey ( $N = 233$ ). Note that data exhibit scatter and some clustering (a,d) that may indicate alternative patterns of morphological adjustment.

#### 251 4. Discussion

252 The six scenarios used in this study provide insight into avulsion processes from the  
 253 inception of delta building. Since avulsion is infrequent, it is difficult to acquire large data sets  
 254 from field studies. By observing avulsions in a numerical river delta model, we can generate a  
 255 large data set from which to infer the process controls over avulsion timescales.

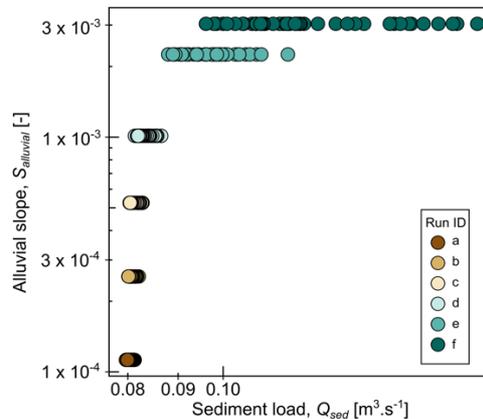
256 **4.1. Investigating the first-order controls of avulsion timescales**

257 Sediment load ( $Q_s$ ) has a high correlation with avulsion timescale observed in the model  
 258 ( $T_{a\text{ empirical}}$ ) (Figs. 3,4) and  $Q_s$  is controlled by the imposed alluvial slope ( $S_{\text{alluvial}}$ ) (Table 3, Fig.  
 259 5). Higher alluvial slopes and sediment loads ( $Q_s$ ) increase the vertical aggradation rate ( $v_a$ ) in  
 260 the proximal part of the delta leading to increased topset slopes (Chadwick et al., 2020). Since  
 261 avulsion timescale is proportional to the rate of vertical aggradation (Jerolmack & Mohrig,  
 262 2007; Mohrig et al., 2000), higher vertical aggradation rates lead to the more frequent  
 263 avulsions. Hence the avulsion timescale in our model is controlled by the alluvial slope that  
 264 was defined independently in our scenarios. This sediment mass-balance approach to  
 265 understanding avulsion timescales has been used in analytical solutions including a radially  
 266 averaged model (Muto, 2001; Muto & Steel, 1997), a channel-averaged model (Reitz et al.,  
 267 2010), and backwater-scaled models (Chadwick et al., 2019; Moodie et al., 2019).

**Table 3.** Cumulative sediment load and median avulsion timescale produced from each scenario. The median avulsion timescale is used to better represent the skewed distribution of avulsion timescale (Fig. 3).

Run ID	Alluvial slope, $S_{\text{alluvial}}$	Cumulative sediment load, $\Sigma Q_{\text{sed}}$ [ $\text{m}^3 \cdot \text{s}^{-1}$ ]	Median avulsion timescale, $T_{a\text{ median}}$ [year]
a	$1.13 \times 10^{-4}$	2.98	335.0
b	$2.55 \times 10^{-4}$	3.07	322.5
c	$5.25 \times 10^{-4}$	2.85	339.1
d	$1.01 \times 10^{-3}$	3.09	330.8
e	$2.25 \times 10^{-3}$	4.09	252.2
f	$3.04 \times 10^{-3}$	5.47	181.9

268



**Figure 5.** Relationship between sediment load ( $Q_{\text{sed}}$ ) and alluvial slopes ( $S_{\text{alluvial}}$ ) imposed to each scenario.

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270 In this study we show how, with a constant sea-level boundary condition, top-down  
 271 forcing of alluvial slope controls the likelihood for a channel to aggrade and then avulse most  
 272 of its water and sediment into a new channel or the surrounding delta plain. This finding is  
 273 consistent with a conceptual model of slope break-avulsion length scaling derived from a  
 274 database of global river deltas by Prasojo et al. (2022). They suggested that this scaling implies  
 275 that the slope break is the prevalent driver of avulsion rather than bottom-up control from  
 276 backwater length or sea-level rise (Chadwick et al., 2020; Chatanantavet et al., 2012; Ganti et  
 277 al., 2016b). In our model, varied in-channel aggradation due to the imposed alluvial slope acts  
 278 as the dominant trigger for avulsion. This reasoning is corroborated by a numerical model in  
 279 which the location of avulsion node consistently scales with the slope break due to linear  
 280 diffusion of aggradation and erosion of the riverbed, even under sea-level rise (Ratliff et al.,  
 281 2021).

#### 282 4.2. Comparison with analytical solution and natural deltas

283 Chadwick et al's (2020) mass-balance based analytical solution is used to calculate  
 284 expected avulsion timescales for our model conditions (Eqs. 3-6, Table S1). Measured  
 285 independent morphometric variables are used in Eqs. 3-6 to calculate avulsion frequency ( $f_a$ )  
 286 and timescale ( $T_a$ ).

$$f_a = \frac{1}{T_a} = \frac{1}{(1 - \lambda_p)} \frac{Q_s}{(L_a - D)BH + DB \left( H_b + z + \frac{DS_{topset}}{2} \right)} \text{ if } D \geq 0 \quad (3)$$

$$f_a = \frac{1}{T_a} = \frac{1}{(1 - \lambda_p)} \frac{Q_s}{L_a BH} \text{ if } D < 0 \quad (4)$$

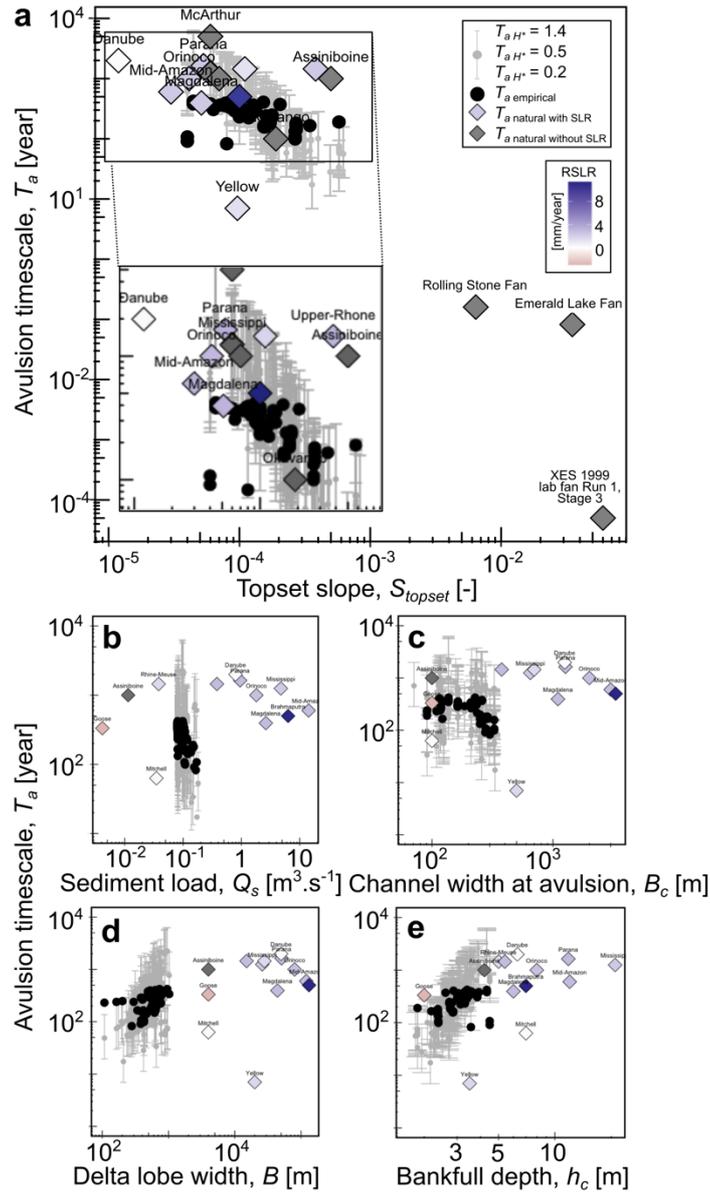
$$D = (H - z)/S_{topset} \quad (5)$$

$$H = H^* h_c \quad (6)$$

287 with  $f_a$  = avulsion frequency [ $\text{year}^{-1}$ ],  $Q_s$  = sediment load [ $\text{m}^3 \cdot \text{s}^{-1}$ ],  $\lambda_p$  = sediment porosity [-],  $L_a$   
 288 = avulsion length [m],  $D$  = delta lobe-progradation distance [km],  $B$  = delta lobe width of each  
 289 avulsion [m],  $H$  = aggradation thickness necessary for avulsion [m],  $H_b$  = basin depth [m],  $z$  =  
 290 magnitude of sea level rise [m],  $S_{topset}$  = topset slope [-],  $H^*$  = avulsion threshold [-], and  $h_c$  =  
 291 bankfull depth [m] calculated using Eq. 1.

292 In calculating these analytical avulsion timescales, sensitivity analyses were undertaken  
 293 using avulsion thresholds ( $H^*$ ) of 0.2, 0.5, and 1.4, which are realistic for lowland deltas (Ganti  
 294 et al., 2019), and  $D > 0$  since there is no allogenic forcing that would make the delta regress.  
 295 The analytical avulsion timescales for  $H^* = 0.2$ , 0.5, and 1.4 are  $T_{a H^* = 0.2}$ ,  $T_{a H^* = 0.5}$ , and  $T_{a H^*}$   
 296 = 1.4, respectively (Table S1). Since sea-level is constant in this study, sea level rise,  $z = 0$ .  
 297 Sediment porosity ( $\lambda_p$ ) is assumed to be 0.4 (Jerolmack, 2009; Paola et al., 2011), bed friction  
 298 coefficient ( $C_f$ ) = 0.002 for lowland rivers (Parker et al., 2007), and constant bankfull discharge  
 299 ( $Q$ ) =  $1050 \text{ m}^3 \cdot \text{s}^{-1}$ .

300 Analytical avulsion timescales were then compared to avulsion timescales observed  
 301 from 19 natural river deltas, two fan deltas and one downscaled physical laboratory fan delta  
 302 mentioned in Chadwick et al. (2020) and Jerolmack & Mohrig (2007), using topset slope values  
 303 from Prasojo et al. (2022) (Table S2). Fig. 6 shows how avulsion timescales observed in our  
 304 model fit both the pattern and the magnitude of both the analytical solution and natural delta  
 305 observations when correlated with topset slope ( $S_{topset}$ ) (Fig. 6a), sediment load ( $Q_s$ ) (Fig. 6b),  
 306 channel width at avulsion ( $B_c$ ) (Fig. 6c), delta lobe width ( $B$ ) (Fig. 6d) and bankfull depth ( $h_c$ )  
 307 (Fig. 6e).



**Figure 6.** Relationships between avulsion timescales and independent variables: (a) topset slope with inset showing a more detail plot; (b) total sediment load; (c) channel width at the avulsion location; (d) delta lobe width; and, (e) bankfull channel depth from model, analytical equations and natural deltas. The plots show model values ( $T_a$  empirical) and those calculated from analytical equations (eqs. 3-6). Solid black circles are empirical results from the model. Grey dots and bars are results from the analytical equations using three avulsion threshold  $H^*$  values ( $T_{aH^*} = 1.4$ ,  $T_{aH^*} = 0.5$ ,  $T_{aH^*} = 0.2$ ). Diamonds are results from the analytical equations applied to natural and laboratory deltas: grey diamonds have constant sea-level; purple diamonds are for deltas with relative sea-level rise ( $RSLR$ ;  $mm \cdot yr^{-1}$ ) colour-coded as shown. Data from natural deltas and the laboratory experiment are available in Table S2.

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In contrast, the avulsion timescales calculated for natural deltas ( $T_a$  natural) are not

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correlated with relative sea-level rise rate ( $RSLR$ ) (Fig. 6a and Fig. S3). This result supports the

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hypothesis that the location of avulsions is unaffected by sea-level rise, as also found in an

312 earlier numerical model study (Ratliff et al., 2021). Avulsion location is thus controlled by  
313 upstream forcing (alluvial slope) (Prasojo et al., 2022) rather than downstream forcing by sea-  
314 level rise or backwater influence (Chadwick et al., 2020; Chatanantavet et al., 2012; Ganti et  
315 al., 2016b). As in our previous field-based empirical study (Prasojo et al., 2022), these results  
316 demonstrate that total sediment load, controlled by alluvial slope upstream of a delta slope  
317 break, has a more dominant role in defining delta avulsion frequency than does sea-level.  
318 Alluvial channels may increase their sediment transport capacity through width adjustment,  
319 leading to the observed correlation between channel width and avulsion timescale (Fig. 6c).  
320 This consequently raises the question as to whether *RSLR* may lead to increased avulsion  
321 frequency in river deltas.

322 Previous literature on the relationship between the frequency of avulsion and sea-level  
323 rise is equivocal. A field study conducted in Mitchell River delta, Australia found that avulsion  
324 frequency increases with sea-level fall (Lane et al., 2017). Numerical model results suggest  
325 that avulsions on the Mississippi (faster) and Trinity (slower) Rivers showed different  
326 responses to Holocene sea-level rise even though they are geographically adjacent  
327 (Chatanantavet et al., 2012; Moran et al., 2017). An example during sea-level fall from the  
328 Goose River delta, Canada, shows that avulsion frequency remained constant during this base-  
329 level adjustment (Nijhuis et al., 2015). In contrast, avulsion frequency in the Rhine-Meuse  
330 delta, Netherlands, increased during the Holocene sea-level rise period (Törnqvist, 1994),  
331 possibly due to aggradation rate ( $Q_s$ ) being controlled by *RSLR*. However, we do not find  $Q_s$ -  
332 *RSLR* relationship in the 19 natural deltas reported in this study (Fig. S4).

333 Overall, these previous studies and our results (Fig. 6) imply that while the frequency  
334 of the most upstream avulsion on a delta is controlled by upstream sediment supply, the  
335 frequency of more distal avulsions has an unclear relationship with either upstream controls or

336 the rate of relative sea-level change. Further investigation from numerical models, analytical  
337 models and/or field data is needed to resolve this issue.

### 338 **4.3. Implications for delta management**

339 Our modelling results advance our understanding about how alluvial slope, which  
340 controls sediment load input, regulates the most upstream avulsion location and timescale. Note  
341 that the avulsion analysed in this modelling study is the most upstream avulsion node  
342 associated with the slope break and/or valley exit (Fig. 1c) and not avulsion nodes located  
343 further downstream the delta plain. The complex hydraulic and sediment transport processes  
344 that deliver these correlations are linked to aggradation rate and hence topset slope.  
345 Consequently, with the increase of anthropogenic forcing both directly within river deltas and  
346 throughout upstream catchment areas (Best, 2019; Darby et al., 2015; Dunn et al., 2019;  
347 Hackney et al., 2020), delta managers can use sediment load management to reduce the risk of  
348 avulsion driven flooding. Interventions that control sediment load may be more effective than  
349 those which address other less dominant factors such as flood variability, delta size, or channel  
350 morphology (Aslan et al., 2005; Brooke et al., 2020; Edmonds et al., 2009; Nienhuis et al.,  
351 2018; Slingerland & Smith, 2004; Valenza et al., 2020).

352 However, finding a perfect balance between maintaining sediment load to nourish delta  
353 environments and to hinder deltas' risk to coastal erosion is challenging. Current deforestation  
354 rate increases sediment supply, responsible for 25% of delta net land gain, but also hastening  
355 the next avulsion (Nienhuis et al., 2020). Alternatively, river damming is responsible for more  
356 than 50% reduction in sediment delivery, collectively leading to a loss of a delta land of 12  
357 km<sup>2</sup> annually (Nienhuis et al., 2020). This declining sediment load not only poses threats to the  
358 long-term sustainability of deltas but also renders them susceptible to adverse effects from  
359 rising sea levels, sand mining and ecological degradation due to sediment starvation (Jordan et

360 al., 2019). Therefore, gaining insights into the distribution patterns and quantities of sediments  
361 in the delta is imperative to ensure its continued sustainability.

#### 362 **4.4. Next steps**

363 An important extension of this modelling work is to vary water discharge ( $Q$ ) and  
364 sediment load ( $Q_s$ ) as variability in these may affect the geomorphic processes controlling  
365 avulsion timescale. Multi-temporal observation of well-studied natural river deltas, such as the  
366 Yellow (Moodie et al., 2019), Mississippi (Chamberlain et al., 2018) or Rhine-Meuse (Pierik  
367 et al., 2018; Stouthamer et al., 2015) deltas, could then be used to validate model results.  
368 Moreover, incorporating other variables such as grain size and sediment cohesion, forcing  
369 through sea-level rise and subsidence, and adding vegetation that controls crevassing and  
370 consequently increases avulsion timescale in future numerical modelling should be considered  
371 (Nienhuis et al., 2018; Pierik et al., 2023; Sanks et al., 2022). In particular, considering the  
372 importance of projected global sea-level changes and the variability of results reported in the  
373 literature, better understanding of sea-level rise impacts on delta avulsion is needed.

374 We have used a simplified modelling approach and have isolated one controlling  
375 variable while holding other factors constant. Observations of the processes and evolution in  
376 the numerical deltas shows the complexity of hydraulic and morphodynamic processes across  
377 delta plains. Future work will need to address this complexity, including: (a) How does the  
378 forcing studied here (alluvial slope) interact with a combination of other factors (e.g. sea-level,  
379 wave and tidal regimes, and anthropogenic effects)? (b) How do the other controls (e.g.  $Q_s$ ,  $Q$ ,  
380 riverbank material, vegetation) in river deltas influence avulsion timescales? And, (c) how  
381 might these avulsion signals be preserved or shredded in the rock record?

#### 382 **5. Conclusion**

383 We conducted a suite of numerical morphodynamic modelling experiments with  
384 variable river alluvial slopes (from  $1.13 \times 10^{-4}$  to  $3.04 \times 10^{-3}$ ) to understand the controls over

385 avulsion location and timescale in a river delta. Sediment load, which is directly driven by  
386 alluvial slope, is the dominant control of the timescale of avulsion. Mechanistically, this is due  
387 to greater sediment transport capacity over steeper alluvial slopes leading to increased sediment  
388 input to the delta plain, accelerated vertical aggradation and more frequent avulsion. The results  
389 support the hypothesis of upstream forcing controlling delta avulsion timescale and location,  
390 rather than downstream controls by backwater length or sea-level rise. A robust understanding  
391 of the main factors controlling avulsion in river deltas has significant implications due to their  
392 direct impacts on (i) coastal and inland hazards on highly populated river deltas and (ii) rock  
393 record interpretation.

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398 purpose of open access, the authors have applied a Creative Commons Attribution (CC-BY)  
399 licence to any Author Accepted Manuscript version arising from this submission.

#### 400 **Open Research**

401 The morphometric variables and avulsion timescales observed from our models are available  
402 in Table S1. The dataset from natural and laboratory river deltas used in this study (Table S2)  
403 and model scenarios (Run a-f) are available in the FigShare repository (Prasojo et al., 2023a,b).

#### 404 **References**

- 405 Aslan, A., Autin, W. J., & Blum, M. D. (2005). Causes of river avulsion: Insights from the  
406 late Holocene avulsion history of the Mississippi River, U.S.A. *Journal of Sedimentary*  
407 *Research*, 75(4), 650–664. <https://doi.org/10.2110/jsr.2005.053>  
408 Bagnold, R. A. (n.d.). *An Approach to the Sediment Transport Problem From General*  
409 *Physics*. Retrieved from doi:10.3133/pp422i  
410 Bates, C. C. (1953). Rational Theory of Delta Formation. *AAPG Bulletin*, 37(9), 2119–2162.  
411 <https://doi.org/10.1306/5ceadd76-16bb-11d7-8645000102c1865d>  
412 Best, J. (2019). Anthropogenic stresses on the world's big rivers. *Nature Geoscience*, 12(1),  
413 7–21. <https://doi.org/10.1038/s41561-018-0262-x>

- 414 Brooke, S., Chadwick, A. J., Silvestre, J., Lamb, M. P., Edmonds, D. A., & Ganti, V. (2022).  
 415 Where rivers jump course. *Science*, 376(6596), 2023.  
 416 [https://doi.org/10.1126/SCIENCE.ABM1215/SUPPL\\_FILE/SCIENCE.ABM1215\\_MO](https://doi.org/10.1126/SCIENCE.ABM1215/SUPPL_FILE/SCIENCE.ABM1215_MO)  
 417 [VIES\\_S1\\_AND\\_S2.ZIP](https://doi.org/10.1126/SCIENCE.ABM1215/SUPPL_FILE/SCIENCE.ABM1215_MO)
- 418 Brooke, Sam, Ganti, V., Chadwick, A. J., & Lamb, M. P. (2020). Flood Variability  
 419 Determines the Location of Lobe-Scale Avulsions on Deltas: Madagascar. *Geophysical*  
 420 *Research Letters*, 47(20), e2020GL088797. <https://doi.org/10.1029/2020GL088797>
- 421 Caldwell, R. L., & Edmonds, D. A. (2014). The effects of sediment properties on deltaic  
 422 processes and morphologies: A numerical modeling study. *Journal of Geophysical*  
 423 *Research: Earth Surface*, 119(5), 961–982. <https://doi.org/10.1002/2013JF002965>
- 424 Chadwick, A. J., Lamb, M. P., Moodie, A. J., Parker, G., & Nittrouer, J. A. (2019). Origin of  
 425 a Preferential Avulsion Node on Lowland River Deltas. *Geophysical Research Letters*,  
 426 46(8), 4267–4277. <https://doi.org/10.1029/2019GL082491>
- 427 Chadwick, A. J., Lamb, M. P., & Ganti, V. (2020). Accelerated river avulsion frequency on  
 428 lowland deltas due to sea-level rise. *Proceedings of the National Academy of Sciences of*  
 429 *the United States of America*, 117(30), 17584–17590.  
 430 <https://doi.org/10.1073/pnas.1912351117>
- 431 Chatanantavet, P., Lamb, M. P., & Nittrouer, J. A. (2012). Backwater controls of avulsion  
 432 location on deltas. *Geophysical Research Letters*, 39(1), 2–7.  
 433 <https://doi.org/10.1029/2011GL050197>
- 434 Darby, S. E., Dunn, F. E., Nicholls, R. J., Rahman, M., & Riddy, L. (2015). A first look at the  
 435 influence of anthropogenic climate change on the future delivery of fluvial sediment to  
 436 the Ganges–Brahmaputra–Meghna delta. *Environmental Science: Processes & Impacts*,  
 437 17(9), 1587–1600. <https://doi.org/10.1039/C5EM00252D>
- 438 Dunn, F. E., Darby, S. E., Nicholls, R. J., Cohen, S., Zarfl, C., & Fekete, B. M. (2019).  
 439 Projections of declining fluvial sediment delivery to major deltas worldwide in response  
 440 to climate change and anthropogenic stress. *Environmental Research Letters*, 14(8),  
 441 084034. <https://doi.org/10.1088/1748-9326/AB304E>
- 442 Edmonds, D. A., & Slingerland, R. L. (2007). Mechanics of river mouth bar formation:  
 443 Implications for the morphodynamics of delta distributary networks. *Journal of*  
 444 *Geophysical Research: Earth Surface*, 112(2). <https://doi.org/10.1029/2006JF000574>
- 445 Edmonds, D. A., & Slingerland, R. L. (2008). Stability of delta distributary networks and  
 446 their bifurcations. *Water Resources Research*, 44(9), 9426.  
 447 <https://doi.org/10.1029/2008WR006992>
- 448 Edmonds, D. A., & Slingerland, R. L. (2010). Significant effect of sediment cohesion on  
 449 deltamorphology. *Nature Geoscience*, 3(2), 105–109. <https://doi.org/10.1038/ngeo730>
- 450 Edmonds, D. A., Hoyal, D. C. J. D., Sheets, B. A., & Slingerland, R. L. (2009). Predicting  
 451 delta avulsions: Implications for coastal wetland restoration. *Geology*, 37(8), 759–762.  
 452 <https://doi.org/10.1130/G25743A.1>
- 453 Edmonds, D. A., Paola, C., Hoyal, D. C. J. D., & Sheets, B. A. (2011). Quantitative metrics  
 454 that describe river deltas and their channel networks. *Journal of Geophysical Research:*  
 455 *Earth Surface*, 116(4), 1–15. <https://doi.org/10.1029/2010JF001955>
- 456 Ericson, J. P., Vörösmarty, C. J., Dingman, S. L., Ward, L. G., & Meybeck, M. (2006).  
 457 Effective sea-level rise and deltas: Causes of change and human dimension implications.  
 458 *Global and Planetary Change*, 50(1–2), 63–82.  
 459 <https://doi.org/10.1016/j.gloplacha.2005.07.004>
- 460 Fagherazzi, S., Edmonds, D. A., Nardin, W., Leonardi, N., Canestrelli, A., Falcini, F., et al.  
 461 (2015). Dynamics of river mouth deposits. *Reviews of Geophysics*, 53(3), 642–672.  
 462 <https://doi.org/10.1002/2014RG000451>

- 463 Ganti, V., Chadwick, A. J., Hassenruck-Gudipati, H. J., & Lamb, M. P. (2016). Avulsion  
464 cycles and their stratigraphic signature on an experimental backwater-controlled delta.  
465 *Journal of Geophysical Research: Earth Surface*, 121(9), 1651–1675.  
466 <https://doi.org/10.1002/2016JF003915>
- 467 Ganti, V., Chadwick, A. J., Hassenruck-Gudipati, H. J., Fuller, B. M., & Lamb, M. P. (2016).  
468 Experimental river delta size set by multiple floods and backwater hydrodynamics.  
469 *Science Advances*, 2(5), e1501768. <https://doi.org/10.1126/sciadv.1501768>
- 470 Ganti, V., Lamb, M. P., & Chadwick, A. J. (2019). Autogenic Erosional Surfaces in Fluvio-  
471 deltaic Stratigraphy from Floods, Avulsions, and Backwater Hydrodynamics. *Journal of*  
472 *Sedimentary Research*, 89(8), 815–832. <https://doi.org/10.2110/jsr.2019.40>
- 473 Geleynse, N., Storms, J. E. A., Walstra, D. J. R., Jagers, H. R. A., Wang, Z. B., & Stive, M. J.  
474 F. (2011). Controls on river delta formation; insights from numerical modelling. *Earth*  
475 *and Planetary Science Letters*, 302(1–2), 217–226.  
476 <https://doi.org/10.1016/j.epsl.2010.12.013>
- 477 Giosan, L., Syvitski, J., Constantinescu, S., & Day, J. (2014). Climate change: Protect the  
478 world's deltas. *Nature*, 516(7529), 31–33. <https://doi.org/10.1038/516031a>
- 479 Hackney, C. R., Darby, S. E., Parsons, D. R., Leyland, J., Best, J. L., Aalto, R., et al. (2020).  
480 River bank instability from unsustainable sand mining in the lower Mekong River.  
481 *Nature Sustainability* 2020 3:3, 3(3), 217–225. [https://doi.org/10.1038/s41893-019-](https://doi.org/10.1038/s41893-019-0455-3)  
482 0455-3
- 483 Hartley, A. J., Weissmann, G. S., & Scuderi, L. (2017). Controls on the apex location of large  
484 deltas. *Journal of the Geological Society*, 174(1), 10–13.  
485 <https://doi.org/10.1144/jgs2015-154>
- 486 Jerolmack, D. J. (2009). Conceptual framework for assessing the response of delta channel  
487 networks to Holocene sea level rise. *Quaternary Science Reviews*, 28, 1786–1800.  
488 <https://doi.org/10.1016/j.quascirev.2009.02.015>
- 489 Jerolmack, D. J., & Mohrig, D. (2007). Conditions for branching in depositional rivers.  
490 *Geology*, 35(5), 463–466. <https://doi.org/10.1130/G23308A.1>
- 491 Jones, L. S., & Schumm, S. A. (2009). Causes of Avulsion: An Overview. *Fluvial*  
492 *Sedimentology VI*, 169–178. <https://doi.org/10.1002/9781444304213.CH13>
- 493 Jordan, C., Tiede, J., Lojek, O., Visscher, J., Apel, H., Nguyen, H. Q., et al. (2019). Sand  
494 mining in the Mekong Delta revisited - current scales of local sediment deficits.  
495 *Scientific Reports* 2019 9:1, 9(1), 1–14. <https://doi.org/10.1038/s41598-019-53804-z>
- 496 Kleinhans, M. G., & Hardy, R. J. (2013, March 15). River bifurcations and avulsion. *Earth*  
497 *Surface Processes and Landforms*. John Wiley & Sons, Ltd.  
498 <https://doi.org/10.1002/esp.3354>
- 499 Kleinhans, M. G., Ferguson, R. I., Lane, S. N., & Hardy, R. J. (2013). Splitting rivers at their  
500 seams: bifurcations and avulsion. *Earth Surface Processes and Landforms*, 38(1), 47–  
501 61. <https://doi.org/10.1002/esp.3268>
- 502 Lane, T. I., Nanson, R. A., Vakarelov, B. K., Ainsworth, R. B., & Dashtgard, S. E. (2017).  
503 Evolution and architectural styles of a forced-regressive Holocene delta and megafan,  
504 Mitchell River, Gulf of Carpentaria, Australia. *Geological Society Special Publication*,  
505 444(1), 305–334. <https://doi.org/10.1144/SP444.9>
- 506 Li, J., Ganti, V., Li, C., & Wei, H. (2022). Upstream migration of avulsion sites on lowland  
507 deltas with river-mouth retreat. *Earth and Planetary Science Letters*, 577, 117270.  
508 <https://doi.org/10.1016/J.EPSL.2021.117270>
- 509 Loucks, D. P. (2019). Developed river deltas: are they sustainable? *Environmental Research*  
510 *Letters*, 14(11), 113004. <https://doi.org/10.1088/1748-9326/AB4165>

- 511 McEwan, E., Stahl, T., Howell, A., Langridge, R., & Wilson, M. (2023). Coseismic river  
 512 avulsion on surface rupturing faults: Assessing earthquake-induced flood hazard.  
 513 *Science Advances*, 9(18). <https://doi.org/10.1126/SCIADV.ADD2932>
- 514 Mohrig, D., Heller, P. L., Paola, C., & Lyons, W. J. (2000). Interpreting avulsion process  
 515 from ancient alluvial sequences: Guadalope-Matarranya system (Northern Spain) and  
 516 Wasatch formation (Western Colorado). *Bulletin of the Geological Society of America*,  
 517 112(12), 1787–1803. [https://doi.org/10.1130/0016-  
 518 7606\(2000\)112<1787:IAPFAA>2.0.CO;2](https://doi.org/10.1130/0016-7606(2000)112<1787:IAPFAA>2.0.CO;2)
- 519 Moodie, A. J., Nittrouer, J. A., Ma, H., Carlson, B. N., Chadwick, A. J., Lamb, M. P., &  
 520 Parker, G. (2019). Modeling Deltaic Lobe-Building Cycles and Channel Avulsions for  
 521 the Yellow River Delta, China. *Journal of Geophysical Research: Earth Surface*,  
 522 124(11), 2438–2462. <https://doi.org/10.1029/2019JF005220>
- 523 Moran, K. E., Nittrouer, J. A., Perillo, M. M., Lorenzo-Trueba, J., & Anderson, J. B. (2017).  
 524 Morphodynamic modeling of fluvial channel fill and avulsion time scales during early  
 525 Holocene transgression, as substantiated by the incised valley stratigraphy of the Trinity  
 526 River, Texas. *Journal of Geophysical Research: Earth Surface*, 122(1), 215–234.  
 527 <https://doi.org/10.1002/2015JF003778>
- 528 Morgan, J. A., Kumar, N., Horner-Devine, A. R., Ahrendt, S., Istanbuloglu, E., &  
 529 Bandaragoda, C. (2020). The use of a morphological acceleration factor in the  
 530 simulation of large-scale fluvial morphodynamics. *Geomorphology*, 356, 107088.  
 531 <https://doi.org/10.1016/J.GEOMORPH.2020.107088>
- 532 Muto, T. (2001). Shoreline Autoretreat Substantiated in Flume Experiments. *Journal of*  
 533 *Sedimentary Research*, 71(2), 246–254. <https://doi.org/10.1306/091400710246>
- 534 Muto, Tetsuji, & Steel, R. J. (1997). Principles of regression and transgression; the nature of  
 535 the interplay between accommodation and sediment supply. *Journal of Sedimentary*  
 536 *Research*, 67(6), 994–1000. [https://doi.org/10.1306/D42686A8-2B26-11D7-  
 537 8648000102C1865D](https://doi.org/10.1306/D42686A8-2B26-11D7-8648000102C1865D)
- 538 Nienhuis, J. H., Ashton, A. D., Edmonds, D. A., Hoitink, A. J. F., Kettner, A. J., Rowland, J.  
 539 C., & Törnqvist, T. E. (2020). Global-scale human impact on delta morphology has led  
 540 to net land area gain. *Nature*, 577(7791), 514–518. [https://doi.org/10.1038/s41586-019-  
 541 1905-9](https://doi.org/10.1038/s41586-019-1905-9)
- 542 Nienhuis, Jaap H., Törnqvist, T. E., & Esposito, C. R. (2018). Crevasse Splays Versus  
 543 Avulsions: A Recipe for Land Building With Levee Breaches. *Geophysical Research*  
 544 *Letters*, 45(9), 4058–4067. <https://doi.org/10.1029/2018GL077933>
- 545 Nienhuis, Jaap H., Hoitink, A. J. F., & Törnqvist, T. E. (2018). Future Change to Tide-  
 546 Influenced Deltas. *Geophysical Research Letters*, 45, 3499–3507.  
 547 <https://doi.org/doi.org/10.1029/2018GL077638>
- 548 Nijhuis, A. G., Edmonds, D. A., Caldwell, R. L., Cederberg, J. A., Slingerland, R. L., Best, J.  
 549 L., et al. (2015). Fluvio-deltaic avulsions during relative sea-level fall. *Geology*, 43(8),  
 550 719–722. <https://doi.org/10.1130/G36788.1>
- 551 Paola, C., Twilley, R. R., Edmonds, D. A., Kim, W., Mohrig, D., Parker, G., et al. (2011).  
 552 Natural Processes in Delta Restoration: Application to the Mississippi Delta. *Annual*  
 553 *Review of Marine Science*, 3(1), 67–91. [https://doi.org/10.1146/annurev-marine-120709-  
 554 142856](https://doi.org/10.1146/annurev-marine-120709-142856)
- 555 Parker, G., Wilcock, P. R., Paola, C., Dietrich, W. E., & Pitlick, J. (2007). Physical basis for  
 556 quasi-universal relations describing bankfull hydraulic geometry of single-thread gravel  
 557 bed rivers. *Journal of Geophysical Research: Earth Surface*, 112(4).  
 558 <https://doi.org/10.1029/2006JF000549>
- 559 Pierik, H. J., Moree, J. I. M., van der Werf, K. M., Roelofs, L., Albernaz, M. B., Wilbers, A.,  
 560 et al. (2023). Vegetation and peat accumulation steer Holocene tidal–fluvial basin filling

- 561 and overbank sedimentation along the Old Rhine River, The Netherlands.  
 562 *Sedimentology*, 70(1), 179–213. <https://doi.org/10.1111/SED.13038>
- 563 Prasojo, O. A., Hoey, T. B., Owen, A., & Williams, R. D. (2022). Slope break and avulsion  
 564 locations scale consistently in global deltas. *Geophysical Research Letters*,  
 565 e2021GL093656. <https://doi.org/10.1029/2021GL093656>
- 566 Prasojo, O. A., Hoey, T. B., Owen, A., & Williams, R. D. (2023a). Supporting Information  
 567 Table S2: First order controls of avulsion in river deltas.  
 568 <https://doi.org/10.6084/m9.figshare.20654037.v2> [Dataset]
- 569 Prasojo, O. A., Hoey, T. B., Owen, A., & Williams, R. D. (2023b). Model runs: First order  
 570 controls of avulsion in river deltas. <https://doi.org/10.6084/m9.figshare.23912625.v1>.  
 571 [Dataset]
- 572 Ratliff, K. M., Hutton, E. W. H., & Murray, A. B. (2021). Modeling long-term delta  
 573 dynamics reveals persistent geometric river avulsion locations. *Earth and Planetary  
 574 Science Letters*, 559, 116786. <https://doi.org/10.1016/j.epsl.2021.116786>
- 575 Reitz, M. D., Jerolmack, D. J., & Swenson, J. B. (2010). Flooding and flow path selection on  
 576 alluvial fans and deltas. *Geophysical Research Letters*, 37(6), n/a-n/a.  
 577 <https://doi.org/10.1029/2009GL041985>
- 578 Rossi, V. M., Kim, W., López, J. L., Edmonds, D., Geleynse, N., Olariu, C., et al. (2016).  
 579 Impact of tidal currents on delta-channel deepening, stratigraphic architecture, and  
 580 sediment bypass beyond the shoreline. *Geology*, 44(11), 927–930.  
 581 <https://doi.org/10.1130/G38334.1>
- 582 Sanks, K. M., Zapp, S. M., Silvestre, J. R., Shaw, J. B., Dutt, R., & Straub, K. M. (2022).  
 583 Marsh Sedimentation Controls Delta Top Morphology, Slope, and Mass Balance.  
 584 *Geophysical Research Letters*, 49(12), e2022GL098513.  
 585 <https://doi.org/10.1029/2022GL098513>
- 586 Shields, M. R., Bianchi, T. S., Mohrig, D., Hutchings, J. A., Kenney, W. F., Kolker, A. S., &  
 587 Curtis, J. H. (2017). Carbon storage in the Mississippi River delta enhanced by  
 588 environmental engineering. *Nature Geoscience* 2017 10:11, 10(11), 846–851.  
 589 <https://doi.org/10.1038/ngeo3044>
- 590 Slingerland, R., & Smith, N. D. (2004). River Avulsions and Their Deposits. *Annual Review  
 591 of Earth and Planetary Sciences*, 32(1), 257–285.  
 592 <https://doi.org/10.1146/annurev.earth.32.101802.120201>
- 593 Syvitski, J. P. M., & Saito, Y. (2007). Morphodynamics of deltas under the influence of  
 594 humans. *Global and Planetary Change*, 57(3–4), 261–282.  
 595 <https://doi.org/10.1016/j.gloplacha.2006.12.001>
- 596 Syvitski, J. P. M., Kettner, A. J., Overeem, I., Hutton, E. W. H., Hannon, M. T., Brakenridge,  
 597 G. R., et al. (2009). Sinking deltas due to human activities. *Nature Geoscience*, 2(10),  
 598 681–686. <https://doi.org/10.1038/ngeo629>
- 599 Tessler, Z. D., Vorosmarty, C. J., Grossberg, M., Gladkova, I., Aizenman, H., Syvitski, J. P.  
 600 M., & Foufoula-Georgiou, E. (2015). Profiling risk and sustainability in coastal deltas of  
 601 the world. *Science*, 349, 638–643. <https://doi.org/10.1126/science.aab3574>
- 602 Törnqvist, T. E. (1994). Middle and late Holocene avulsion history of the River Rhine  
 603 (Rhine-Meuse delta, Netherlands). *Geology*, 22(August), 711–714. Retrieved from  
 604 [http://pubs.geoscienceworld.org/gsa/geology/article-pdf/22/8/711/3515257/i0091-7613-  
 605 22-8-711.pdf](http://pubs.geoscienceworld.org/gsa/geology/article-pdf/22/8/711/3515257/i0091-7613-22-8-711.pdf)
- 606 Valenza, J. M., Edmonds, D. A., Hwang, T., & Roy, S. (2020). Downstream changes in river  
 607 avulsion style are related to channel morphology. *Nature Communications*, 11(2116).  
 608 <https://doi.org/10.1038/s41467-020-15859-9>
- 609 Wallace, D. J., Storms, J. E. A., Wallinga, J., Dam, R. L. V. A. N., Blaauw, M., Derksen, M.  
 610 S., et al. (2014). Shrinking and Sinking Deltas : Major role of Dams in delta subsidence

- 611 and Effective Sea Level Rise. *Nature Geoscience*, 123(May), 1973–1984.  
612 <https://doi.org/10.1038/ngeo129>
- 613 Williams, R. D., Measures, R., Hicks, D. M., & Brasington, J. (2016). Assessment of a  
614 numerical model to reproduce event-scale erosion and deposition distributions in a  
615 braided river. *Water Resources Research*, 52(8), 6621–6642.  
616 <https://doi.org/10.1002/2015WR018491>
- 617 Wolinsky, M. A., Edmonds, D. A., Martin, J., & Paola, C. (2010). Delta allometry: Growth  
618 laws for river deltas. *Geophysical Research Letters*, 37(21).  
619 <https://doi.org/10.1029/2010GL044592>
- 620 Wright, L. D. (1977). Sediment transport and deposition at river mouths: A synthesis.  
621 *Bulletin of the Geological Society of America*, 88(6), 857–868.  
622 [https://doi.org/10.1130/0016-7606\(1977\)88<857:STADAR>2.0.CO;2](https://doi.org/10.1130/0016-7606(1977)88<857:STADAR>2.0.CO;2)  
623

Figure 1.

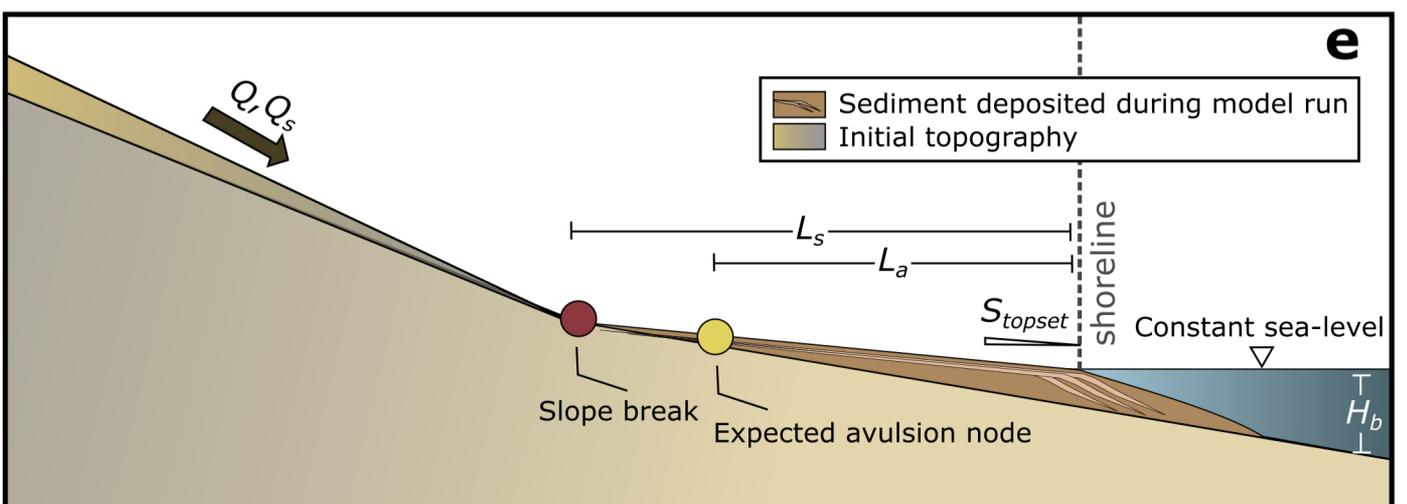
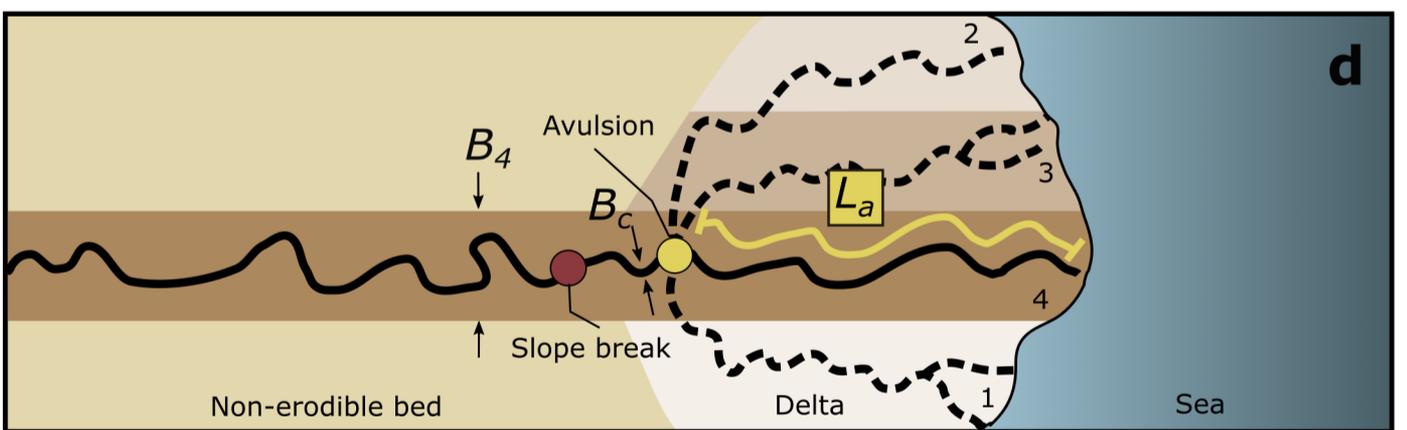
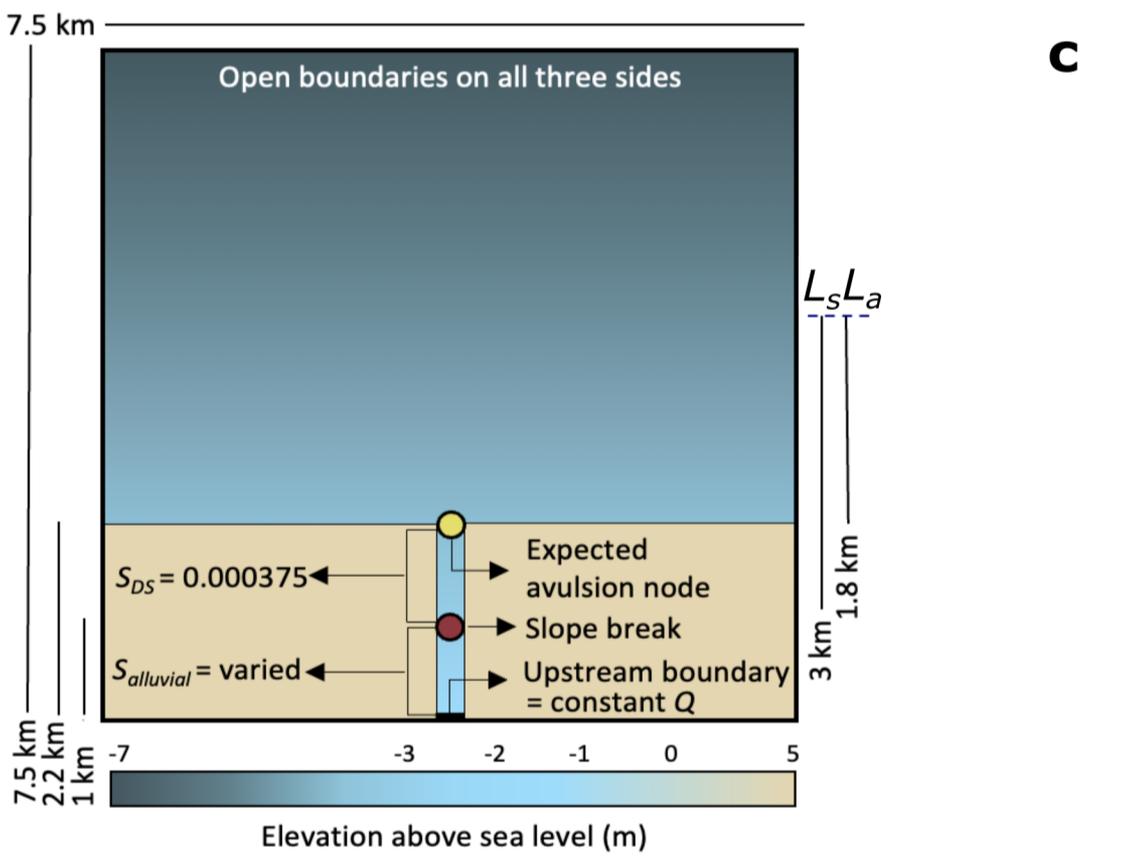
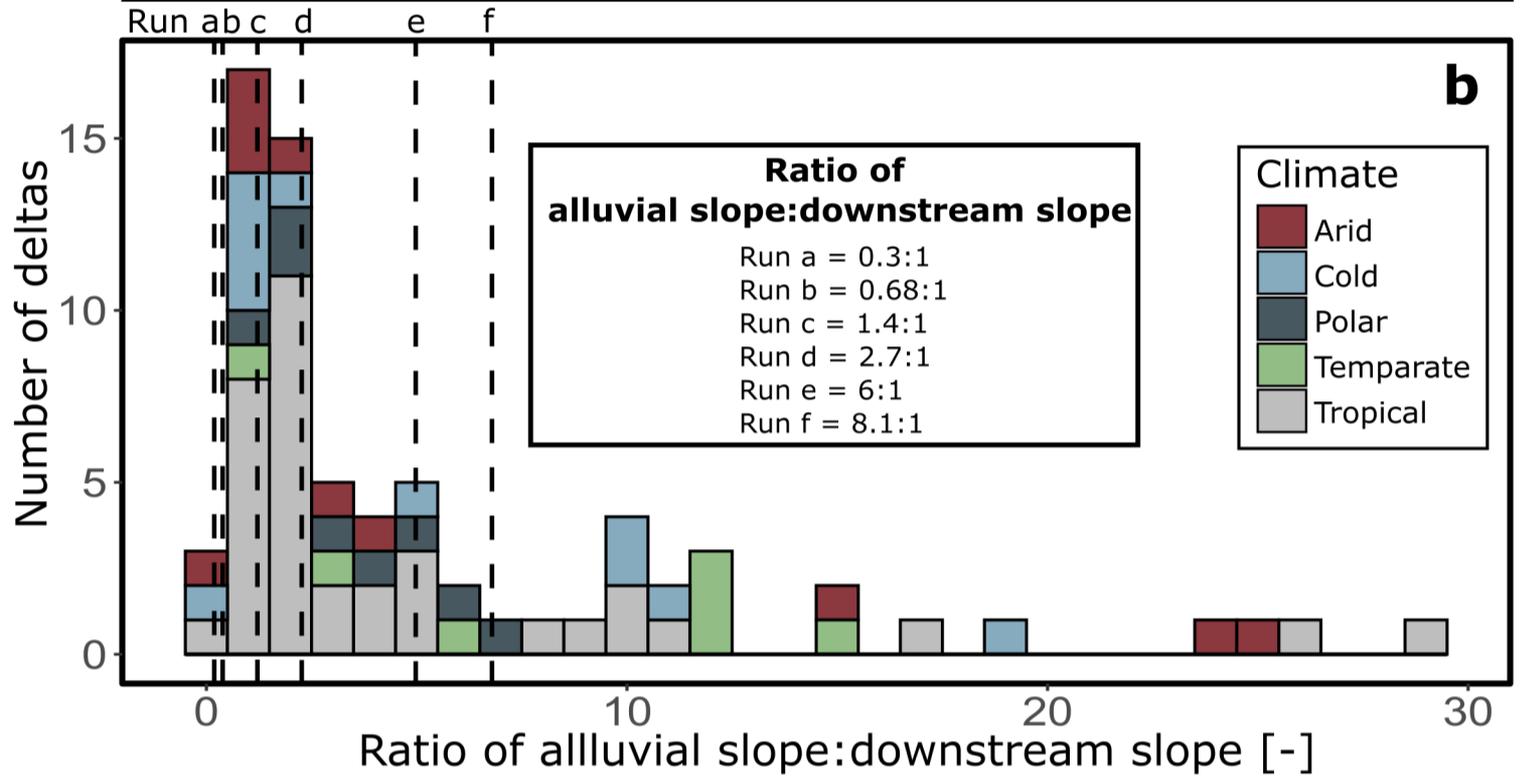
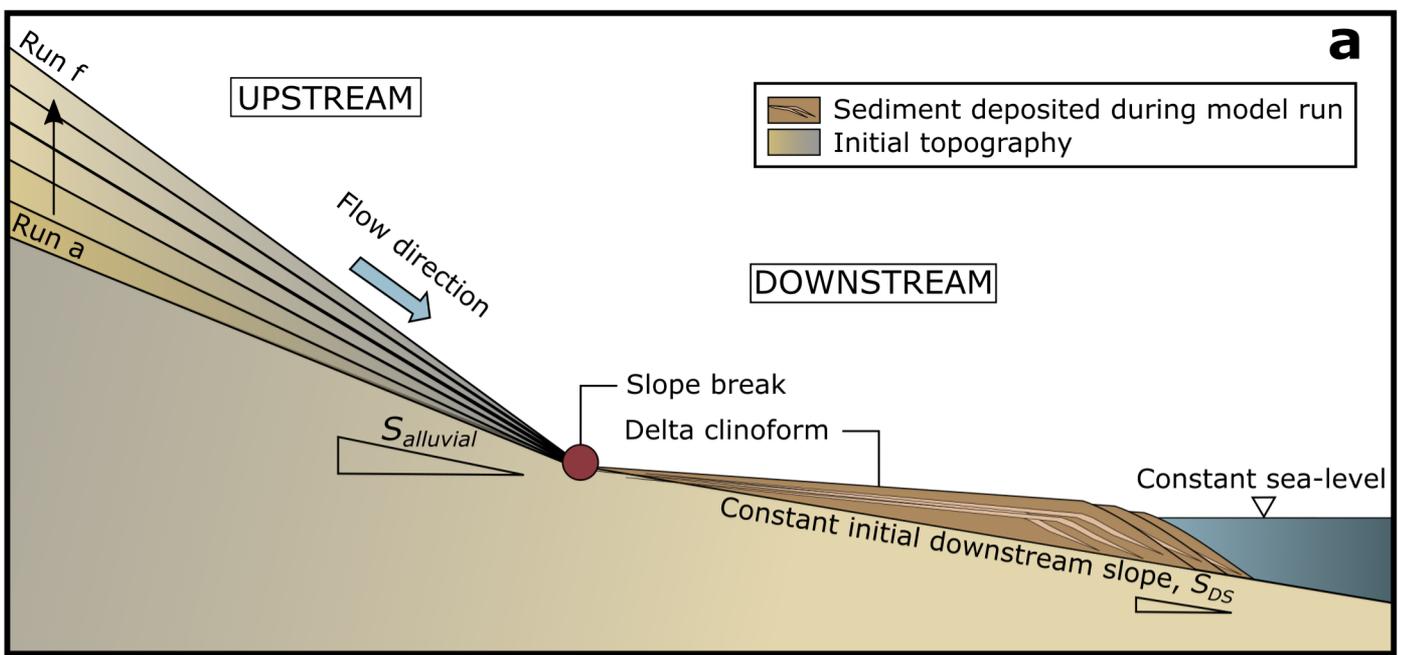
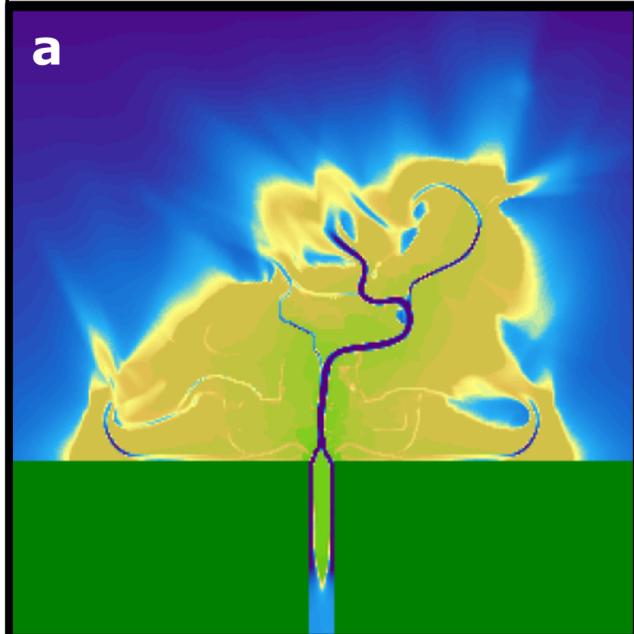


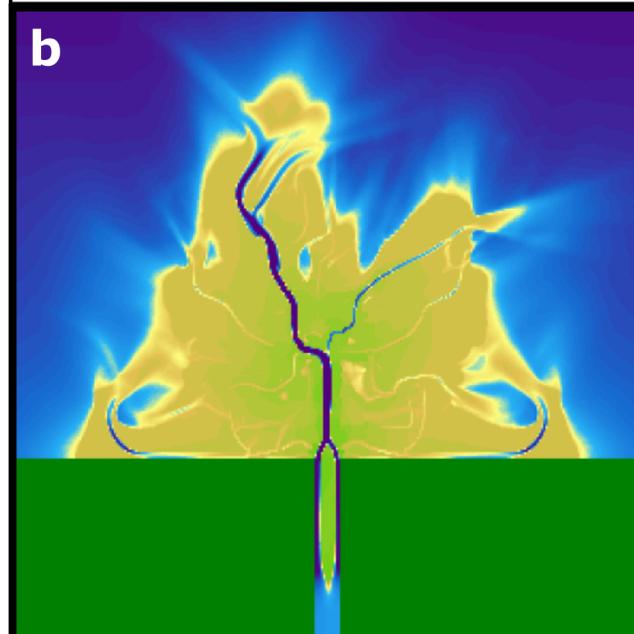
Figure 2.

Run a



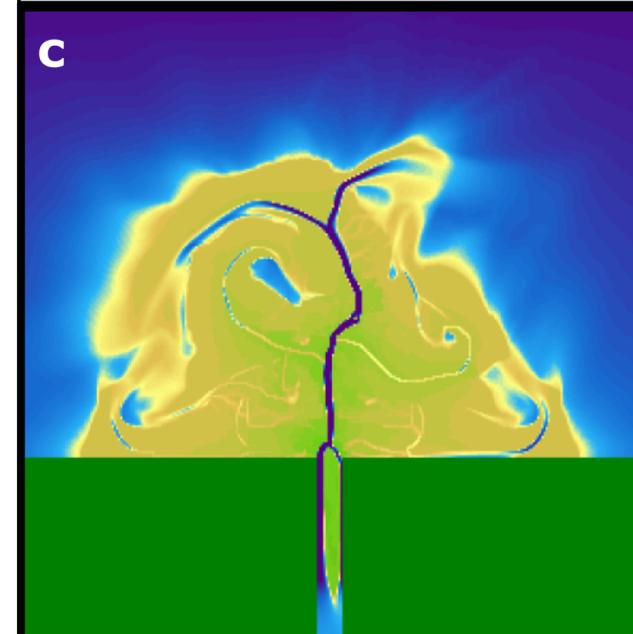
$$S_{alluvial} = 1.13 \times 10^{-4}$$

Run b



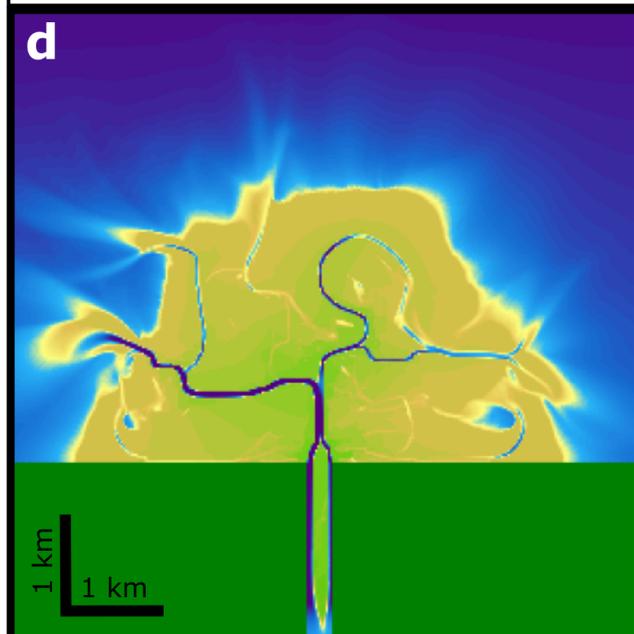
$$S_{alluvial} = 2.55 \times 10^{-4}$$

Run c



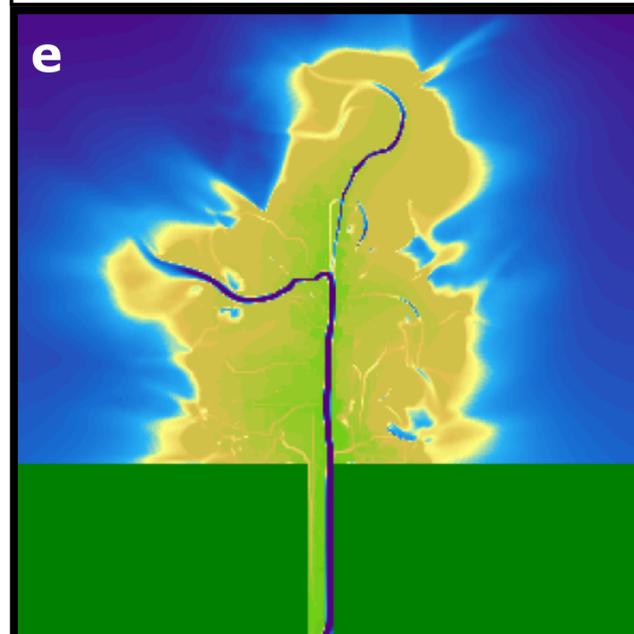
$$S_{alluvial} = 5.25 \times 10^{-4}$$

Run d



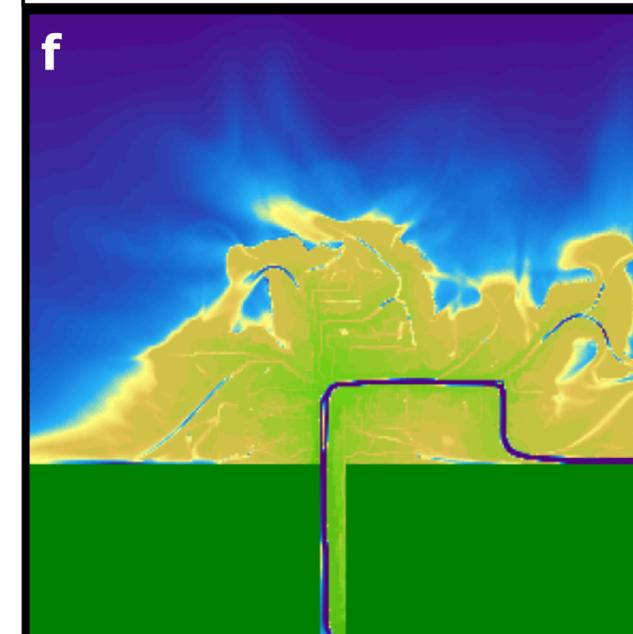
$$S_{alluvial} = 1.01 \times 10^{-3}$$

Run e



$$S_{alluvial} = 2.25 \times 10^{-3}$$

Run f



$$S_{alluvial} = 3.04 \times 10^{-3}$$

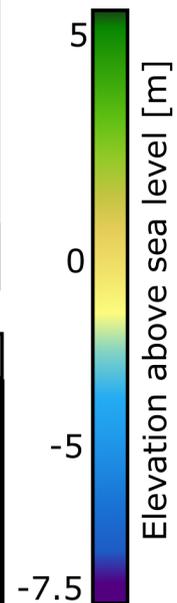


Figure 3.

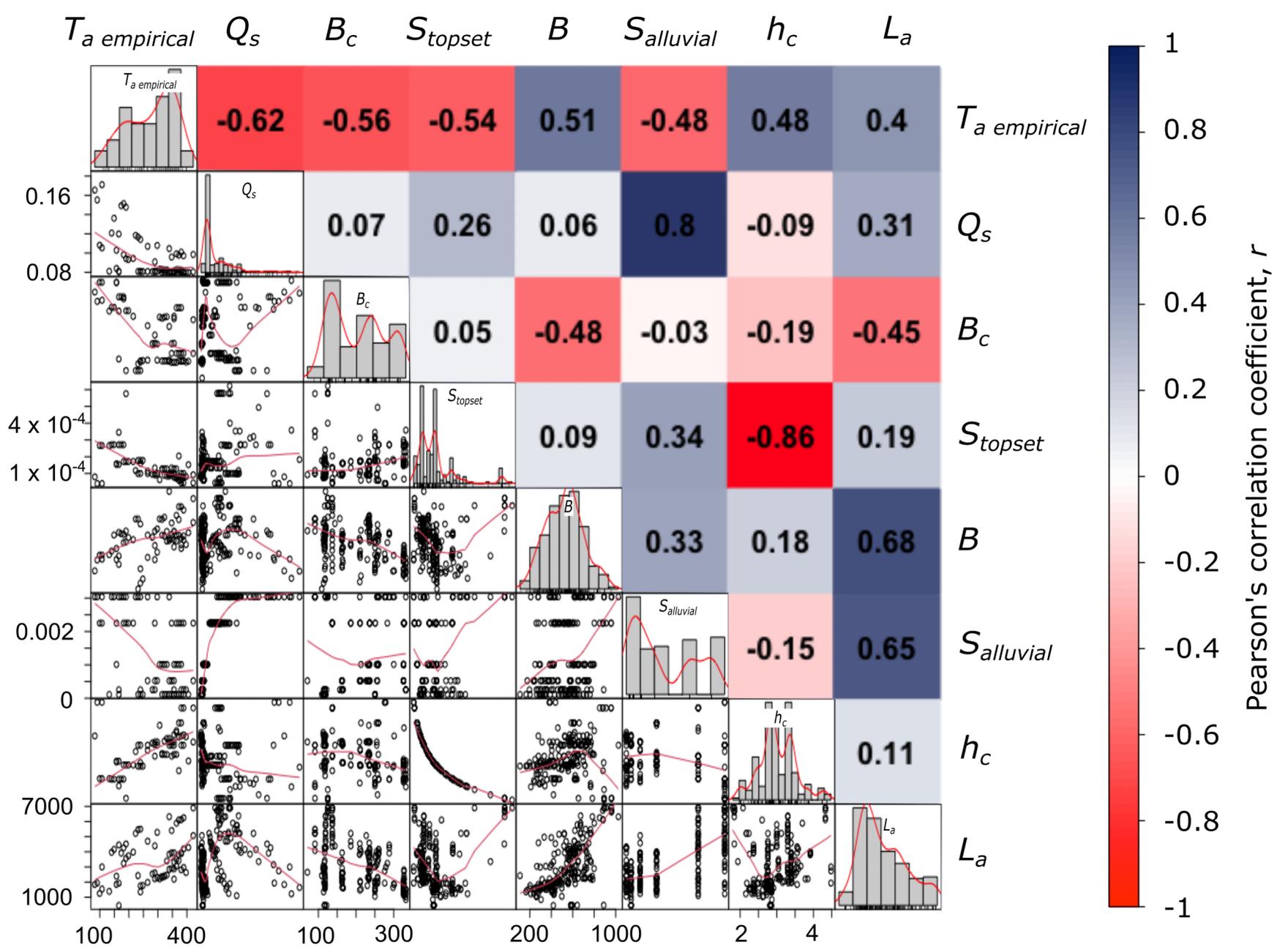
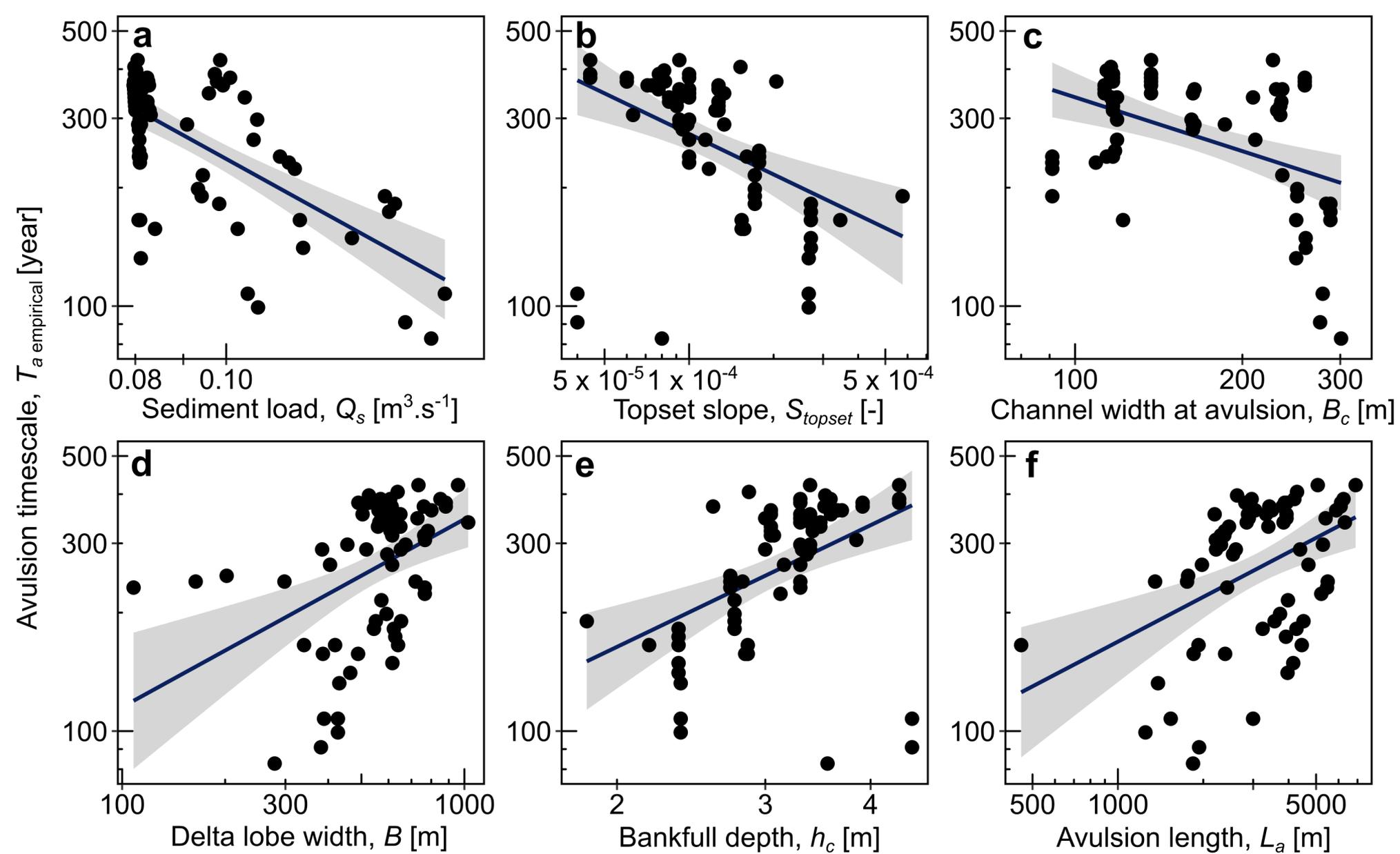


Figure 4.



$$T_{a \text{ empirical}} = 11.45Q_s^{-1.32}, R^2 = 0.38, p = 6.76 \times 10^{-10}, SE_{\text{residual}} = 0.32$$

$$T_{a \text{ empirical}} = 11.7S_{topset}^{-0.34}, R^2 = 0.23, p = 1.4 \times 10^{-4}, SE_{\text{residual}} = 0.38$$

$$T_{a \text{ empirical}} = 4572.6B_c^{-0.56}, R^2 = 0.21, p = 1.76 \times 10^{-6}, SE_{\text{residual}} = 0.32$$

$$T_{a \text{ empirical}} = 12.77B^{0.48}, R^2 = 0.24, p = 1.74 \times 10^{-4}, SE_{\text{residual}} = 0.38$$

$$T_{a \text{ empirical}} = 80.34h_c^{1.03}, R^2 = 0.23, p = 1.38 \times 10^{-4}, SE_{\text{residual}} = 0.38$$

$$T_{a \text{ empirical}} = 12.45L_a^{0.38}, R^2 = 0.18, p = 2.18 \times 10^{-4}, SE_{\text{residual}} = 0.38$$

Figure 5.

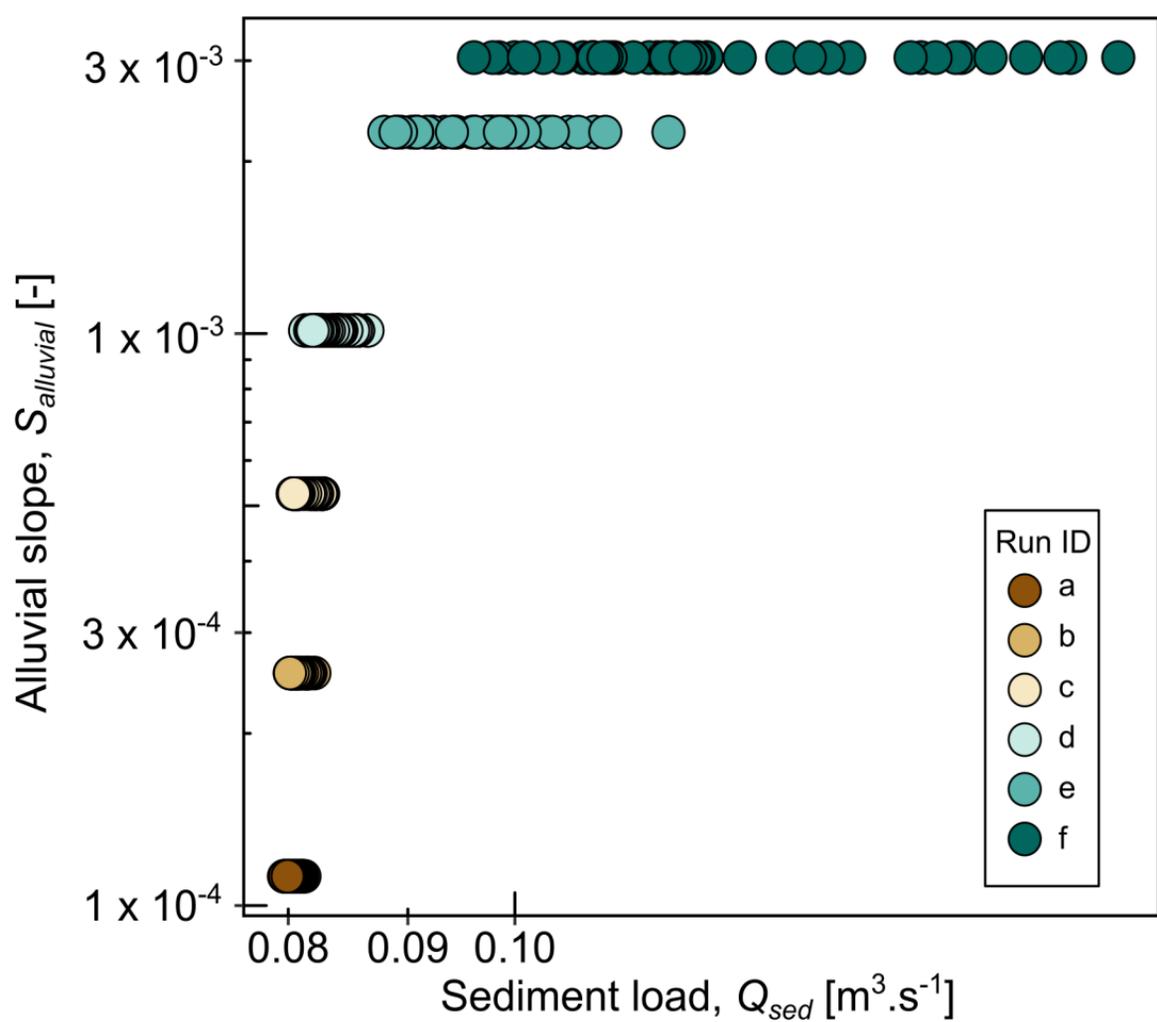


Figure 6.

