# Breaking the Ice: Exploring the Changing Dynamics of Winter Breakup Events in the Beaufort Sea

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## Abstract

The Beaufort Sea has experienced a significant decline in sea ice, with thinner first-year ice replacing thicker multi-year ice. This transition makes the ice cover weaker and more mobile, making it more vulnerable to breakup during winter. Using a coupled ocean-sea-ice model, we investigated the impact of these changes on sea-ice breakup events and lead formation from 2000 to 2018. The simulation shows an increasing trend in the Beaufort Sea lead area fraction during winter, with a pronounced transition around 2007. A high lead area fraction in winter promotes a significant growth of new, thin ice within the Beaufort region while also leading to enhanced sea ice transport out of the area. The export offsets ice growth, resulting in negative volume anomalies and preconditioning a thinner and weaker ice pack at the end of the cool season. Our results indicate that large breakup events may become more frequent as the sea-ice cover thins and that such events only became common after 2007. This result highlights the need to represent these processes in global-scale climate models to improve projections of the Arctic.

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11	Key Points:
12	- Modelled leads in the Beaufort Sea during winter time are increasing at $4\%$ per decade
13	over the period 2000-2018
14	• The shift to thinner and younger sea ice, particularly after 2007, makes the Beau-
15	fort Sea more vulnerable to large breakup events by winds
16	• Winter breakup increases ice export from the Beaufort Sea and leads to a thin-
17	ner and weaker ice cover at the end of the cool season

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#### Abstract 18

The Beaufort Sea has experienced a significant decline in sea ice, with thinner first-year 19 ice replacing thicker multi-year ice. This transition makes the ice cover weaker and more 20 mobile, making it more vulnerable to breakup during winter. Using a coupled ocean-sea-21 ice model, we investigated the impact of these changes on sea-ice breakup events and lead 22 formation from 2000 to 2018. The simulation shows an increasing trend in the Beaufort 23 Sea lead area fraction during winter, with a pronounced transition around 2007. A high 24 lead area fraction in winter promotes a significant growth of new, thin ice within the Beau-25 fort region while also leading to enhanced sea ice transport out of the area. The export 26 offsets ice growth, resulting in negative volume anomalies and preconditioning a thin-27 ner and weaker ice pack at the end of the cool season. Our results indicate that large 28 breakup events may become more frequent as the sea-ice cover thins and that such events 29 only became common after 2007. This result highlights the need to represent these pro-30 cesses in global-scale climate models to improve projections of the Arctic. 31

### 32

## **Plain Language Summary**

The sea ice cover in the Beaufort Sea has been changing - it is getting thinner and 33 weaker. This makes the ice more likely to break apart from strong winds. Using a com-34 puter model, we study how these changes may have affected the frequency of large sea-35 ice breakup events from 2000 to 2018. We find that the amount of open areas in the sea 36 ice, called leads, is increasing during winter. This allows new, thin ice to form, but also 37 causes more ice to move out of the region under the action of winds and currents. This 38 movement of ice cancels the growth of new ice, resulting in less ice overall at the end of 39 winter in this region. Interestingly, these events became more common after 2007 and 40 the results suggests that bigger breakup events might happen more often as the sea ice 41 continues to thin. This study highlights how important it is to include these changes in 42 large climate models to better predict what might happen in the Arctic in the future. 43

#### 1 Introduction 44

Recent decades have seen dramatic reductions in the extent, age, and thickness of 45 Arctic sea ice (e.g. Kwok, 2018). Those changes are particularly pronounced in the Beau-46 fort Sea, which has experienced a rapid decline in sea ice extent and sea ice thickness, 47 during both summer and winter. There has been a notable shift in the composition of 48

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sea ice in the early-2000s (Babb et al., 2022), where the Beaufort ice cover transitioned 49 from a state which was dominated by thick and old multi-year ice (MYI) to an increas-50 ingly thinner, more fragmented and mobile seasonal ice cover around 2007 (Moore et al., 51 2022; Wood et al., 2013). This regime shift towards younger, thinner sea ice is affect-52 ing the dynamical properties of the ice cover (Zhang et al., 2012), reducing the ice's me-53 chanical strength, thereby making it more vulnerable to atmospheric forcing (Petty et 54 al., 2016) and contributing to the observed increase in sea ice deformation and drift speeds 55 in the Arctic Ocean and the Beaufort Sea in particular (Rampal et al., 2009; Kwok & 56 Cunningham, 2010; Spreen et al., 2011). These changes in sea ice properties and ice dy-57 namics have consequences for the stability and persistence of the Beaufort Sea ice cover, 58 potentially resulting in more frequent sea ice breakup and lead formation (Maslanik et 59 al., 2007). This potentially has important implications for the overall mass balance of 60 sea ice, ice-ocean interactions, and the Arctic climate system. However, due to the lack 61 of long-term observations and the difficulties in modelling sea-ice breakup, our knowl-62 edge is currently limited when it comes to understanding the relationship between these 63 changing sea ice characteristics and the frequency and intensity of breakup events. 64

During winter months, lead formation exposes the ocean to the colder atmosphere 65 resulting in large air-sea heat, moisture and gas fluxes. The intense heat loss from the 66 ocean promotes new ice formation, contributing to the sea ice mass balance in the Arc-67 tic winter (accounting for between 10 and 20% of the total ice growth in the Arctic dur-68 ing winter Heil & Hibler, 2002; Kwok, 2006). Recent estimates from Boutin et al. (2023) 69 show that this number could be as high as 25–35%. Brine rejection from sea ice forma-70 tion increases the stability of the Arctic halocline (Shimada et al., 2005), which protects 71 sea ice from melting by suppressing the entrainment of subsurface heat into the surface 72 layer. Brine-driven eddies under sea ice leads can affect thermohaline structure of the 73 mixed layer by transporting heat and salt laterally under the sea ice (Matsumura & Ha-74 sumi, 2008; Peralta-Ferriz & Woodgate, 2015). Leads are also key regions for marine bi-75 ological productivity due to increased access to sunlight, which is otherwise very limited 76 due to the presence of thick, snow-covered sea ice. For example, recent observations show 77 that leads in Arctic pack ice can enable early phytoplanktonic blooms (Assmy et al., 2017) 78 impacting primary production and Arctic marine food webs. This could become more 79 frequent due to thinner and more dynamic sea ice that is more vulnerable to breakup 80 (Fadeev et al., 2021). 81

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In the Beaufort Sea, leads regularly form throughout the winter season in response 82 to divergent sea ice motion driven by atmospheric weather systems or ocean currents (Lewis 83 & Hutchings, 2019; Jewell & Hutchings, 2023). Meanwhile, several large breakup events 84 have been identified from satellite observations in recent decades (most noteworthy in 85 the winters of 2013 (Beitsch et al., 2014; Rheinlænder et al., 2022) and 2016 (Babb et 86 al., 2019)). Wintertime breakup events are characterized by extensive fracturing of the 87 ice cover associated with atmospheric synoptic conditions persisting from a few days to 88 several weeks (Jewell & Hutchings, 2023). Such events have been shown to significantly 89 impact sea ice conditions in the Beaufort Sea, with potential implications for the wider 90 Arctic sea ice mass balance. The large breakup events in winter 2013 and 2016 resulted 91 in anomalous sea ice drift and enhanced ice export out of the Beaufort Sea (e.g. Babb 92 et al., 2016; Rheinlænder et al., 2022). This led to an overall reduction in the Beaufort 93 ice volume in April and a thinner, less compact ice cover prior to the onset of the melt-94 ing season. This conditions the ice cover for rapid summer melt (e.g. Maslanik et al., 95 2007) and could contribute to the low regional September sea ice area seen in recent decades 96 (Williams et al., 2016; Babb et al., 2019; Moore et al., 2022). 97

Winter breakup events can also have important consequences for the MYI cover. 98 Enhanced ice export during winter may increase the flushing of MYI through the Beau-99 fort Sea which increases the amount of MYI being advected into the region from the cen-100 tral Arctic (as was seen in 2013, e.g. Richter-Menge & Farrell, 2013). This could mo-101 bilize the oldest and thickest sea ice residing north of Greenland, also known as the Last 102 Ice Area, which is subsequently advected into the Beaufort Sea. For example, during sum-103 mer 2020/21 Moore et al. (2022) found large concentrations of thick and old ice in the 104 Beaufort Sea, which could be traced back to enhanced winter transport from the Last 105 Ice Area. Less MYI now survives through the summer melt season, making the Beau-106 fort Sea a major contributor to MYI loss in the Arctic (Howell et al., 2016; Babb et al., 107 2022). 108

Despite their importance, sea ice breakup and lead formation are generally not adequately reproduced in large-scale sea-ice and climate models (e.g. Spreen et al., 2017). This is partly due to the difficulty in representing small-scale deformation features, like cracks and leads, for horizontal resolutions coarser than  $\sim$ 5 km (Hutter et al., 2022). And while higher resolution sea-ice models (4–5 km) have demonstrated a certain degree of proficiency in representing the large-scale distribution of sea-ice leads in the Arctic (e.g.

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<sup>115</sup> Wang et al., 2016), they are currently considered too costly for global-scale climate mod-<sup>116</sup> els.

In this study we present a newly developed coupled ocean-sea ice model based on 117 the neXtSIM sea-ice model which employs a brittle sea-ice rheology making it partic-118 ularly suitable for simulating small-scale ice deformation and linear kinematic features 119 like fractures and leads in sea ice at comparatively low resolution (about 12 km here) 120 (Rampal et al., 2019; Bouchat et al., 2022; Ólason et al., 2022). Rheinlænder et al. (2022) 121 recently demonstrated neXtSIM's ability to provide a realistic and accurate represen-122 tation of sea ice fracturing and lead propagation associated with the 2013 breakup event 123 in the Beaufort Sea. The study highlighted that such extreme breakup events could be-124 come more frequent as the sea ice thins, raising concerns about the vulnerability of the 125 Beaufort ice cover. Here, we seek to understand how changes in the Beaufort sea-ice regime 126 during the early 21st century have affected the stability of the ice cover and the occur-127 rence of extreme breakup events focusing on the winters of 2000–2018. By addressing 128 this question, this study aims to provide new insights into the ongoing transformations 129 of the Beaufort Sea ice cover and its implications for regional sea ice volume, MYI cov-130 erage, and sea-ice transport. 131

## 132 2 Methods

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## 2.1 Model setup

The model used in this study is the new coupled sea-ice-ocean model recently pre-134 sented in Boutin et al. (2023). In brief, the ocean component is the Océan PArallélisé 135 model (OPA), which is part of the NEMO3.6 modelling platform (Madec, 2008). We use 136 the regional CREG025 configuration (Talandier & Lique, 2021), which encompasses the 137 Arctic and parts of the North Atlantic down to 27°N, and has 75 vertical levels and a 138 nominal horizontal resolution of  $1/4^{\circ} (\simeq 12 \text{ km} \text{ in the Arctic basin})$ . The sea ice com-139 ponent is neXtSIM, a state-of-the-art, finite element, sea ice model using a moving La-140 grangian mesh (Bouillon & Rampal, 2015; Rampal et al., 2016). Sea ice dynamics rely 141 on the Brittle Bingham-Maxwell (BBM) rheology described in Olason et al. (2022), while 142 sea ice thermodynamics are simulated following the Winton (2000) model. We refer to 143 Boutin et al. (2023) for detailed information about the model setup. 144

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As noted by Hutter et al. (2022), sea ice models generally struggle to simulate sea 145 ice dynamics when run at resolutions coarser than about 5 km; in particular, features 146 like fractures, shear zones, and lead openings. However, the BBM rheology has demon-147 strated its capability to reproduce deformations consistent with observations when run-148 ning at a resolution of O(10 km) in the neXtSIM model (Ólason et al., 2022) and in the 149 SI3 model (L. Brodeau, personal communication). Specifically, these models exhibit ex-150 cellent capability in accurately capturing the divergence rates associated with the open-151 ing of leads when using the BBM rheology (Ólason et al., 2022; Rheinlænder et al., 2022). 152

The simulation presented in this study is the same as in Regan et al. (2023) and 153 Boutin et al. (2023). The simulation starts in 1995 and ends in 2018. The first five years 154 were considered a spin-up period and disregarded for analysis. Atmospheric forcing is 155 taken from the hourly ERA5 reanalysis at a 1/4-degree horizontal resolution. This sim-156 ulation has been thoroughly evaluated in two recent publications (Boutin et al., 2023; 157 Regan et al., 2023). Boutin et al. (2023) showed that simulating key sea-ice quantities 158 like volume, extent, large-scale drift, and sea ice deformations are consistent with satel-159 lite observations. Regan et al. (2023) demonstrated that the simulation successfully re-160 produces the spatial distribution and evolution of observed MYI extent. They also found 161 a good agreement with observed estimates of the regional dynamic and thermodynamic 162 components of the winter sea ice mass balance from Ricker et al. (2021). 163

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## 2.2 Lead definition

neXtSIM uses three ice categories; open water, young ice, and consolidated ice. Newly 165 formed ice, thinner than  $h_{\rm max}$  (here set to 18 cm), is assigned to the young ice category, 166 representing the formation and growth of frazil and young ice in open water. Ice in the 167 young ice category is transferred to the consolidated ice category as its thickness exceeds 168  $h_{\rm max}$  (see appendix A of Rampal et al., 2009). In winter, when the Beaufort Sea is fully 169 ice-covered, lead opening is the only way open water can be exposed to the atmosphere, 170 and young ice can be formed. Therefore, we assume (as in Rheinlænder et al., 2022; Boutin 171 et al., 2023) that open water and young ice formed in winter are a proxy for the pres-172 ence of leads in the model. We use this assumption to estimate the rapid growth of thin, 173 newly formed ice in open-water and thin ice regions. A grid cell is considered a lead when 174 the combined fraction of open water and young ice exceeds a critical threshold  $c_{lim}$ , thereby 175 excluding the thicker pack ice. We found that a value of  $c_{lim} = 5\%$  gives a reasonable 176

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- lead distribution. The sensitivity of the simulated lead fraction to the value of  $c_{lim}$  is included in the Supplementary Material. The total lead area fraction (LAF) can then be calculated by multiplying the lead fraction with the area of each grid cell. A snapshot of the simulated sea ice concentration and lead fraction on 25 March 2016 is shown in Fig. 1. Here, leads are clearly identified as areas of open water and newly formed ice, whereas the pack ice is associated with low lead fraction values. The LAF calculated over the Beaufort region for this instance is 21%, which means that 21% of the Beaufort area
- 184 is covered by leads.



Figure 1. Snapshots of sea ice concentration (%) and lead fraction (%) in the Beaufort Sea on March 25, 2016. The total lead area fraction (LAF) is calculated using a threshold value of 5% (see 2).

## 185 **3 Results**

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## 3.1 Simulated changes in the Beaufort ice cover

Since the 2000s, the MYI extent in the Arctic has declined considerably (Fig. 2a). 187 During the period 2000–2018, the model simulates a decline in the winter MYI area, which 188 is part of a long-term negative trend in the Arctic as seen from satellite observations (e.g. 189 Babb et al., 2022). In the Beaufort Sea region (outlined in Fig. 2a), extensive areas of 190 thicker and older MYI were present during January–March in the early 2000s (i.e. 2000– 191 2004). For the later part of the simulation (years 2014–2018), the MYI extent is signif-192 icantly reduced and is consistent with the observed trend towards reduced MYI concen-193 tration (Howell et al., 2016). Both the average sea ice thickness and MYI concentration 194 computed over the Beaufort region exhibit considerable year-to-year variations (Fig. 2b), 195 but overall there is a shift towards thinner and younger ice types. The average winter 196 sea-ice thickness decreased from 1.9 m in 2000–2004 to 1.6 m in 2014–2018. Meanwhile, 197 some old and thick sea ice still remains located north of Canada and the Last Ice Area, 198 which are important source regions for MYI import to the Beaufort Sea (Moore et al., 199 2022). 200

The ice drift speeds in the Beaufort Sea are also increasing (Fig. 2c) with a 12% increase over the 2000–2018 period (not shown). Previously, this increase in ice drift speeds has been linked to thinning of the sea ice cover and enhanced deformation rates leading to more fracturing and lead opening (Rampal et al., 2009). In the following, we examine the impact of transitioning to a more seasonal and thinner ice cover on the formation of leads and sea-ice breakup in the Beaufort Sea.

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## 3.2 Simulated changes in wintertime leads

We show the simulated wintertime LAF in the Beaufort Sea for the period 2000– 209 2018 in Fig. 3. The LAF shows a large day-to-day variability ranging from 5% to 40%, which reflects the intermittent nature of sea-ice fracturing. Winter-mean values (January– 211 March) generally fall between 10% and 25% with a climatological average of 20%. The lead fraction is generally higher in January and decreases during February and March 213 as the ice becomes thicker and more compact (Fig. 3b and c).



**Figure 2.** (a) Maps of simulated sea-ice thickness and MYI extent (yellow contour corresponding to the MYI fraction of 0.4) in the Arctic for January–March averaged from 2000–2004 and 2014–2018. (b) Time series of JFM-mean sea-ice thickness and MYI extent averaged in the Beaufort region. (c) Histogram of the winter sea-ice drift speed (cm s<sup>-1</sup>) distributions for the 2000–2004 and 2014–2018 period. The Beaufort region is bounded by three gates shown in (**a**); North gate (78°N), East gate (120°W), and West gate (160°W). The definition of the Beaufort region is the same as in Moore et al. (2022).

We find a statistically significant increase in wintertime lead occurrences (4.2% per decade) over the period 2000–2018 based on a simple linear regression analysis (Supplementary Fig. S1). However, the linear relationship becomes less significant when we consider individual months, likely due to the larger spread in the monthly data (Fig. 3b). Here, we note that the modelled LAF is affected by the cutoff value used in the lead definition (see section 2), but the choice of this cutoff value has no impact on the simulated variability and trends (Supplementary Fig. S2).

The LAF values simulated by neXtSIM are consistent with observations of sea-ice leads from MODIS at 1 km<sup>2</sup> spatial resolution with observed winter-mean (NovemberApril) lead fraction area ranging between 10–20% in the Beaufort Sea (Willmes et al., 2023). Willmes et al. (2023) also found a significant trend in leads over the 2002–2021 period, but only for April. It is worth noting that the MODIS observations have uncertainties due to contamination by clouds and will only see opening leads that are relatively large. This highlights the need for a dedicated intercomparison study to determine how to best use MODIS imagery to classify and evaluate lead formation in sea-ice models, but this is beyond the scope of this paper.

Based on the LAF time series, several larger breakup events can be identified, all 230 occurring after 2007. Most noteworthy are the years 2008, 2010, 2013, 2016, and 2018, 231 which have high average wintertime LAFs (green triangles in Fig. 3a). Many of these 232 events have also been identified in satellite observations (e.g. Jewell & Hutchings, 2023). 233 For example, large breakup events were observed in 2008, 2013 and 2016 and have been 234 described in earlier studies (Wang et al., 2016; Babb et al., 2019; Rheinlænder et al., 2022). 235 The breakup being simulated by the model in 2018 is not seen in observations and is likely 236 a result of too strong melting simulated by the model in the summer of 2016 (not shown), 237 leading to thinner sea ice that could break up more easily. During these events, the daily 238 LAF exceeds the 90th percentile (about 30%; Fig. 3c) for a period of more than 15 days 239 during winter (Fig. 3d). We therefore expect these events to have a significant impact 240 on the Beaufort ice cover. Meanwhile, smaller breakup events are also present in other 241 years. For example, 2005 and 2006 exhibit high LAFs (daily values exceeding 35%), but 242 these occurrences are relatively short-lived and result in low winter-mean values over-243 all. Consequently, they will likely have less impact. 244

Around 2007, we identify a shift in the interannual variability of the LAF based 245 on the monthly values in Fig. 3b. For the 2000–2007 period, the variability (shown by 246 the standard deviation for each month) ranges from 3.1–4.8% during winter. After 2007, 247 this increases to 3.9-7.4% for 2008–2018, and the average LAF increases during all win-248 ter months. This coincides with more extreme breakups during this period. March, in 249 particular, stands out, showing a 38% increase in the mean LAF relative to the 2000-250 2007 period while also exhibiting the highest variability (standard deviation of 7.4%). 251 We will investigate this in more detail in section 3.3 and identify some common char-252 acteristics of the simulated breakup events and their impacts (section 3.4). 253

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Figure 3. Simulated lead area fraction (LAF; %) in the Beaufort Sea from January through March. (a) Daily LAF (grey) from 2000–2018 with circles showing the winter-mean values along with the standard deviation. The dashed line shows the 2000–2018 winter climatology. (b) Monthly LAF climatologies for the 2000–2007 (blue) and 2008–2018 (orange) periods. Diamonds represent the monthly mean with the standard deviation in whiskers. (c) Histograms of monthly LAF distributions, with the 90th percentile (~30%) shown by the dashed line. (d) Stacked barplot of binned LAF from 20% to 45% where the height of the bars corresponds to the number of days. Numbers denote the total number of days where the daily average LAF exceeds the 90th percentile.

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## 3.3 Driving mechanisms of sea-ice breakup events

The sea ice movement driven by wind and ocean currents can create stresses within the ice pack, leading to fracturing and the formation of leads (Lewis & Hutchings, 2019). In addition, changes in the material properties of the ice, such as ice thickness, concen-

tration, and strength, can also influence the susceptibility to breakup. Here we look at 258 the drivers of the simulated changes in winter lead formation, focusing particularly on 259 the shift occurring around 2007. 260

## 3.3.1 Winds

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Jewell and Hutchings (2023) analysed the synoptic conditions during breakup events 262 in the Beaufort Sea during winters 1993–2013. They show a consistent connection be-263 tween wind forcing and lead formation, where a breakup is typically associated with high 264 sea level pressure and relatively strong anticyclonic winds over the Beaufort Sea. Sim-265 ilarly, Wang et al. (2016) concluded that a stronger Beaufort High results in stronger south-266 easterly winds in the Beaufort Sea, which pushes sea ice away from the coast and thus 267 promotes higher ice divergence and lead formation. 268

In general, high wintertime LAFs in the Beaufort Sea are linked with persistently 269 higher wind speeds (Figure 4b and c) in agreement with Jewell and Hutchings (2023). 270 In 2010, 2013, and 2016, the daily wind speed exceeds  $8.5 \text{ m s}^{-1}$  (one standard devia-271 tion above the mean) for more than 20 days (Fig. 4d). These conditions are typically 272 associated with a positive sea level pressure difference across the Beaufort Sea and east-273 erly winds (Supplementary Fig. S3), creating favourable conditions for off-shore ice drift 274 and enhanced breakup. It is worth noting that there is considerable variability between 275 the different months (Figure 4c and Supplementary Fig. S3), and we do not find a sim-276 ple relationship between wind speed and sea-ice breakup. Jewell and Hutchings (2023) 277 came to the same conclusion indicating that breakup may occur for a wide range of at-278 mospheric conditions. Both the duration of strong winds as well as the wind direction 279 appear to be important for initiating a breakup. For example, in 2005, conditions were 280 comparable to other breakup years (e.g. 2013), with relatively strong and persistent winds 281 during winter, however the LAF remained relatively low throughout the winter of 2005 282 (Supplementary Fig. S4). This could be due to the fact that the ice was thicker and stronger 283 (Fig. 2b) and thus less sensitive to winds. 284

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We find no trend in the ERA5 winds in the Beaufort region during winter over the period 2000 to 2018 (Fig. 4b). The year-to-year variability is also relatively similar for 286 the 2000–2007 and the 2008–2018 winter periods. The same is true if we consider indi-287 vidual months rather than the winter-mean values (Fig. 4c), showing no major differ-288

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ence in wind strength for January versus March. Thus, changes in the wind forcing cannot explain the fact that we are seeing the strongest increase in lead variability in March. Based on this, we speculate that the shift in the lead formation dynamics seen after 2007 is linked to the thinning of the Beaufort ice cover (Fig. 2), making it more vulnerable to atmospheric forcing during winter.

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## 3.3.2 Changes in ice conditions

Changes in the dynamic properties of sea ice are generally attributed to changes 295 in concentration and thickness, which in turn affect the strength of the ice (Zhang et al., 296 2012). In sea ice models, including neXtSIM, ice strength is parameterized as a function 297 of ice thickness and concentration (Hibler, 1979). Therefore, we expect the simulated de-298 cline in ice thickness (Fig. 2b) to weaken the ice pack and reduce the internal ice stress. 299 As a result, this leads to an overall increase in the simulated deformation rates (Fig. 5b) 300 and increased drift speeds (Fig. 5a) in the Beaufort Sea. The positive trend in sea ice 301 drift speeds is consistent with earlier modelling studies (e.g. Zhang et al., 2012) and ob-302 servations (e.g. Rampal et al., 2009; Spreen et al., 2011). 303

Comparing the simulated LAF in Fig. 4a to the time series of the mean ice speed 304 and deformation rates (Fig. 5) strongly points to a shift in the dynamic sea ice prop-305 erties and lead formation dynamics after 2007. Both sea ice drift and deformation rates 306 show a pronounced change in the variability, fluctuating between relatively low and high 307 values mirroring the changes in the LAF. High LAF is associated with high deformation 308 rates and increased ice speed in the Beaufort Sea (Fig. 5a and 5b), exceeding  $7 \text{ cm s}^{-1}$ 309 during breakup events. It is worth noting that both drift and deformation values remain 310 relatively low for winters without significant breakup, i.e. their baseline values do not 311 seem to change much over the 2000–2018 period. This suggests that individual extreme 312 events can substantially alter the overall trend seen in the data. 313

By plotting the ratio between ice drift to wind speed (Fig. 5c) we see a clear increase in the ice drift to wind speed ratio, especially during breakup events. This reflects an increased sensitivity of the Beaufort ice cover to wind forcing during the late 2000s, whereas in the earlier period (e.g. in 2005 when the ice was thicker), there is a larger disconnect between strong winds and ice drift speeds. Overall, this points to changes in the

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Figure 4. Time series of (a) normalized wintertime lead fraction (%) in Beaufort region, (b) ERA5 daily mean wind speed ( $m s^{-1}$ ) between 70–75°N along 150°. The January–March average wind speed is highlighted by blue circles. (c) Monthly mean wind speed for January, February and March. (d) The number of days where the daily wind speed exceeds 8.5 m s<sup>-1</sup> (one standard deviation above the mean) is shown in bars. Its mean and standard deviation is shown by solid and dashed lines. The solid black line shows the wintertime (January through March) mean wind speed. The time series in (a) is normalized by subtracting its mean and dividing by its standard deviation. The transect used for calculating the winds is indicated in Figure 4b.

- material properties of sea ice being a major factor in driving the shift we see in the sim-
- <sup>320</sup> ulated lead area fraction.



Figure 5. Time series of (a) mean ice speed (cm s<sup>-1</sup>), (b) total deformation rate (1/day) and (c) ratio (%) between mean sea-ice drift and mean wind speed. The thin black line shows the average wind speed. All time series are based on wintertime (January through March) means and averaged over the Beaufort region. The area used for averaging is shown in Fig. 2. Green triangles highlight winters with significant breakup identified in Figure 3a.

## 321

## 3.4 Impacts on Beaufort ice volume and MYI

322 323 In this section, we seek to understand how winter sea ice breakup impacts the ice volume in the Beaufort Sea. Changes in regional ice volume during winter can be sep-

arated into two terms: (i) thermodynamic ice growth and (ii) sea ice transport. Note that we are omitting the term associated with sea ice melting as this can be considered negligible during winter (Graham et al., 2019).

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## 3.4.1 Thermodynamic ice growth in leads

In winter, the opening of leads results in intense heat loss from the underlying ocean and promotes new ice formation. More sea ice breakup could therefore increase the local winter ice growth overall and modulate the composition of the Beaufort ice pack by increasing the fraction of thinner and younger ice types.

The thermodynamic ice growth from January through March is shown in Fig. 6 332 for leads and pack ice, respectively, where growth in leads is associated with the forma-333 tion of new and thin sea ice (Rheinlænder et al., 2022). Note that the ice growth in leads 334 is independent of the lead detection algorithm. Overall, the growth of new ice in leads 335 is increasing over the period 2000–2018 (Fig. 6b). We find a statistically significant lin-336 ear trend of 4% per decade for wintertime ice production in leads. This is consistent with 337 the results of Boutin et al. (2023), who used the same model to find a 4.3% per decade 338 trend on the pan-Arctic scale. Our result is also consistent with the simulated trend in 339 LAF (4.2%; Supplementary Fig. S1) and suggests that leads play an increasingly key role 340 in the local sea-ice volume budget as the ice cover becomes thinner and more fractured. 341

By comparing the growth estimates to the LAF time series in Figure 3a, we see that winters with more breakups also have larger ice production overall. New ice production in leads is consistently higher for these years (top 5) compared to the climatology, and the fraction relative to the total growth is above 40% (except for 2008). Overall, these results show that winter breakup can significantly increase the local ice volume by enhancing ice growth.

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## 3.4.2 Volume transports

The simulated winter-mean ice volume fluxes are shown in Fig. 7 across the Beaufort Sea's eastern, northern and western gates. The transport has been separated into FYI and MYI contributions. The western gate captures primarily the ice export from the central Beaufort to the Chukchi Sea, whereas the northern and eastern gate captures the import of thicker and older sea ice from the Canadian Archipelago and the LIA. This

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Figure 6. (a) Thermodynamic sea-ice growth (km<sup>3</sup>) in leads (blue) and pack ice (orange) in the Beaufort Sea during winter (January–March). The dashed line shows the winter climatological mean ice volume growth in leads. Numbers in white indicate the fraction (in %) of the respective growth relative to the total growth (shown by black numbers above the bars). (b) Fraction of ice volume growth in leads relative to the total thermodynamic growth. A linear regression model has been fitted to the data and is shown by the grey line. Breakup years are highlighted by green triangles in (a).



Figure 7. Time series of January–March ice volume fluxes (km<sup>3</sup>) into the Beaufort Sea across the eastern (160°W), northern (78°N), and western gates (120°W) and the total fluxes in the bottom panel. The transport is separated into contributions from FYI (light blue) and MYI (dark blue), with dashed lines showing the winter-mean climatologies. Positive values indicate sea ice transport into the Beaufort region.

circulation pattern is associated with the anti-cyclonic circulation of the Beaufort Gyre (Howell et al., 2016). Overall, the total fluxes from January through March show a net import of MYI and a net export of FYI. The majority of the transport occurs at the northern and western gates, while the transport through the eastern gate is generally small (only accounting for about 6% of the total import).

During winters with enhanced lead activity, the ice transport through the Beaufort region increases (bottom panel in Fig. 7) and is consistent with the higher mean ice drift speeds seen in Fig. 5a. At the western gate, there is a large export of primarily thinner and younger FYI during breakup events. A smaller fraction of MYI is also exported, especially in the early 2000s, but is reduced during the later part of the simulation in line with previous studies (e.g. Howell et al., 2016; Babb et al., 2022).

At the northern gate, the simulated volume transports are primarily dominated by 365 the import of MYI from the central Arctic. There is generally a higher MYI transport 366 into the Beaufort Sea during breakup events, for example, in 2013 which shows a net im-367 port of 85  $\mathrm{km}^3$  across the eastern and northern gates. In comparison, the average MYI 368 import is  $\sim 45 \text{ km}^3$  over the period 2000–2018. However, years with relatively low lead 369 fractions (in the early 2000s and in 2015) also show high MYI import (and export), while 370 the breakup events in 2016 and 2018 have lower MYI fluxes despite high lead fractions. 371 A possible explanation is that there is simply less MYI in the Arctic in total and, there-372 fore, less to be transported into the Beaufort region (Babb et al., 2022). This is likely 373 the case in this simulation, which underestimates MYI extent in 2017 and 2018 partly 374 due to unrealistically high melting in the summer of 2016. 375

In total, ice export is larger than import during winter breakup events, which suggests that sea ice breakup contributes to regional dynamic ice loss in the short term. This will likely also affect ice transport in the following months and impact the regional ice volume before the beginning of the next melt season. For example, winter export from the Beaufort region could lead to enhanced flushing of MYI through the Beaufort Sea (e.g. Babb et al., 2019) providing dynamical replenishment for the ice loss during winter.

To understand the cumulative effects of winter breakups, we look at the transport into the Beaufort Sea through the entire cool season from January through June in Figure 8. The 2000–2018 climatology shows a net ice export at the end of the cool season

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**Figure 8.** Cumulative sea ice volume fluxes from January to June for all years from 2000 to 2018. Years with high wintertime lead fractions (2008, 2010, 2013, 2016 and 2018) are shown in colours. The 2000–2018 climatology is shown in black with the ±1 standard deviation.

(about -200 km<sup>3</sup> in June). Years with higher lead activity in winter (2008, 2010, 2016, 386 2018) exhibit larger cumulative net ice export (more than one standard deviation be-387 low the mean). A notable exception is 2013, which shows as net ice import, despite a 388 large export in February–March (Babb et al., 2019; Rheinlænder et al., 2022). This was 389 caused by enhanced advection of thicker MYI through the northern boundary from mid-390 April (Fig. 9) offsetting ice export (primarily thin FYI) at the western boundary. Mean-391 while, the other breakup events show little evidence of MYI flushing associated with the 392 increased winter ice export from the Beaufort Sea. 393

Overall, these results suggest that winter breakup events may have a negative im-394 pact on the Beaufort ice mass balance by enhancing ice export, despite also promoting 395 significant new ice growth. But what is the combined effect of winter breakup on the Beau-396 fort ice volume? In Figure 10, we show the relationship between wintertime lead area 397 fraction, i.e. over the period from January through March, and the ice conditions in the 398 Beaufort Sea at the end of the cool season (from January 1 to June 1). The cumulative 399 cool season transport out of the Beaufort Sea is generally larger when the LAF is high 400 and we see a clear separation of the breakup years that also have enhanced export. The 401

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**Figure 9.** June 1 cumulative ice volume fluxes separated into FYI (blue) and MYI (orange) contributions from 2000 to 2018 in the Beaufort region. The net June 1 ice volume flux is shown by black circles. Positive (negative) values corresponds to a net import (export).

402	correlation between cool season transport and the Beaufort ice volume at the beginning
403	of June exhibits a similar grouping, with breakup years showing lower ice volume val-
404	ues on June 1. A similar relationship was also found based on satellite observations (e.g.
405	Babb et al., 2019; Moore et al., 2022), suggesting that high winter export from the Beau
406	fort Sea results in an anomalously thin ice cover and negative regional volume anoma-
407	lies. This could preconditioning the ice cover for increased summer melt and ultimately
408	result in record low regional September sea ice minima as shown by Babb et al. (2019).

## 409 4 Discussion

In an earlier modelling study, Wang et al. (2016) simulated the time evolution of 410 lead formation in the Beaufort Sea over the last three decades (1985–2014) using a high-411 resolution (4.5 km) sea-ice model (Finite Element Sea Ice-Ocean Model; FESOM). In 412 contrast to the neXtSIM simulation, they observed no increase in the number of large-413 scale breakup events in winter, which they related to the absence of wind stress trends 414 in the Beaufort region. However, one of the notable contrasts between these two mod-415 els is the difference in sea ice rheology; m-EVP for FESOM versus BBM in neXtSIM. 416 This could lead to significant differences in how the ice cover responds dynamically to 417 changes in the mechanical ice properties and the sensitivity to wind forcing. Another dif-418 ference is the definition of leads in Wang et al. (2016), which are defined as locations where 419 the sea ice is at least 20% thinner than at its surroundings (within a 25-km radius). Firstly, 420



Figure 10. Scatterplot of the cool season (January 1 - June 1) ice volume flux and (a) winter mean (January–March) lead area fraction and (b) June 1 ice volume. Breakup years (2008, 2010, 2013, 2016 and 2018) are highlighted in colours.

this excludes very wide leads, and secondly may fail to capture very localized divergence 421 events leading to drops in sea ice concentration as seen in the neXtSIM simulation. Mean-422 while, recent observational data based on MODIS imagery (Willmes et al., 2023) show 423 a significant positive trend of 2% per decade in lead frequencies in the Beaufort Sea over 424 the period from 2002 to 2021 (during April only). This is similar to Hoffman et al. (2022) 425 who observed a small, but significant increase in pan-Arctic leads from satellite data over 426 the same period, despite large uncertainty due to the increasing cloud cover in the Arc-427 tic. 428

We find that the change in lead formation dynamics simulated by the neXtSIM model, 429 notably the increased variability in lead formation after 2007, can be linked to a shift 430 in the Beaufort ice dynamics. Such transitions have been reported in observations. Long-431 term sea ice data from satellites dating back to the 1980s show evidence that the Beau-432 fort Sea transitioned to a thinner state in 1998 (Hutchings & Rigor, 2012 and Fig. 1 in 433 Babb et al., 2019). Another transition occurred around 2007 (e.g. Moore et al., 2022; 434 Babb et al., 2022), which reflected a shift from an old ice regime (1979–2007) when the 435 region was dominated by MYI to a young ice regime (2007-present). Similarly, Wood 436 et al. (2013) pointed to a "new normal" climate in the Beaufort Sea since 2007, char-437

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acterized by an increasingly mobile and thus more dynamic ice pack, which agrees with
the increase in ice drift in Figure 2c.

Our results indicate, that sea ice thinning and loss of MYI in the Beaufort region 440 makes the ice cover less resilient to wind forcing thus increasing the likelihood of large 441 breakups. This could lead to enhanced inter-annual variability in Beaufort Sea ice con-442 ditions and may increase the potential for rapid sea ice loss (Moore et al., 2022; Maslanik 443 et al., 2007). Similarly, Petty et al. (2016) found an amplified sensitivity of the Beau-444 fort sea ice circulation in winter to wind forcing during the late-2000s. This increase in 445 winter ice drift is commonly attributed to general sea ice thinning and reduction in me-446 chanical ice strength (Zhang et al., 2012; Rampal et al., 2009), which is also evident from 447 our results in Fig. 5. Meanwhile, Jewell and Hutchings (2023) noted that changes in ice 448 thickness is not the only factor controlling breakup. Atmospheric conditions such as wind 449 direction, storm propagation and duration of strong winds are also important factors that 450 contribute to sea-ice breakup. In fact, the LAF timeseries in Figure 3 show that breakup 451 events also occurred during the early-2000s (for example in 2005 and 2006) when the ice 452 was considerably thicker. This emphasizes the importance of atmospheric forcing in ini-453 tiating breakup. 454

While the atmosphere plays a dominant role in triggering sea-ice breakup on short 455 time scales (days to weeks), the ocean may also play a role in preconditioning sea-ice breakup 456 on seasonal time scales (Willmes et al., 2023). For example, enhanced ocean heat fluxes 457 during summer and autumn may predispose the ice cover to enhanced melt, resulting 458 in a thinner and weaker ice cover before the beginning of the cooling season (Herbaut 459 et al., 2022; Graham et al., 2019). Lead formation can also be expected to have signif-460 icant impacts on the ocean underneath, for example by enhancing mechanical energy in-461 put available for mixing and through brine formation thereby affecting mixed layer prop-462 erties and halocline stability (Matsumura & Hasumi, 2008; Peralta-Ferriz & Woodgate, 463 2015; Shimada et al., 2005). Mixing up warmer waters through lead opening could en-464 hance basal melting and limit new ice growth in the leads (e.g. Graham et al., 2019). 465 Such feedbacks could be important for ice-ocean interactions even on longer time scales 466 but they are not assessed explicitly in this study. 467

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## <sup>468</sup> 5 Summary and conclusions

This paper presents a multi-decadal simulation using the coupled ocean-sea-ice neXtSIM-OPA model and investigates the temporal changes in wintertime sea ice leads in the Beaufort Sea and their impacts. The simulation shows a small but significant increasing trend in the Beaufort lead area fraction (4% per decade) over the winter season (January through March) for the period 2000 to 2018. This is consistent with a general decrease in ice thickness and MYI cover as well as enhanced drift speeds during winter in the Beaufort region.

Around 2007 we find a notable increase in the simulated lead area fraction vari-476 ability associated with enhanced sea ice breakup, high deformation rates and an increase 477 in the mean ice velocities. These changes coincide with the observed regime shift that 478 occurred in the Beaufort Sea in 2007 (Wood et al., 2013; Moore et al., 2022), character-479 ized by a transition from a state dominated by thicker and older MYI towards more sea-480 sonal, thinner and younger sea ice. We find no significant trend in the surface winds dur-481 ing winter over the simulated time period. This suggests that the changes in lead for-482 mation dynamics can be attributed to changes in the sea ice conditions (i.e. thinning, 483 loss of ice strength and enhanced deformation) rather than changes in the atmospheric 484 forcing. Consequently, the ice cover becomes more sensitive to wind forcing, which may 485 lead to enhanced inter-annual variability in Beaufort Sea ice conditions and more extreme 486 breakup during winter. 487

Several large breakup events are identified which significantly impact the regional 488 thermodynamic ice production, with new ice growth in leads contributing up to 40% of 489 the total winter ice growth. This implies that sea ice leads play an important role in the 490 local ice mass balance in the Arctic (as Boutin et al., 2023, also found). Meanwhile, years 491 with high lead activity in winter consistently exhibit increased ice export, primarily FYI, 492 from the Beaufort Sea throughout the entire cool season (January 1 to June 1). While 493 some breakup events also show an enhanced import of MYI to the Beaufort from the cen-494 tral Arctic (e.g. in 2013), we find no consistent evidence that winter breakup leads to 495 the flushing of MYI through the Beaufort Sea. 496

<sup>497</sup> Overall, these results suggest that winter breakups have a negative impact on the <sup>498</sup> Beaufort ice volume, preconditioning a thinner and weaker ice pack at the end of the cool <sup>499</sup> season (see also Babb et al., 2016, 2019; Moore et al., 2022). This could lead to earlier

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<sup>500</sup> breakup in spring and enhanced summer melt, thereby contributing to accelerating sea

- <sup>501</sup> ice loss in the Beaufort Sea. This further highlights the need to include small-scale sea-
- <sup>502</sup> ice deformation and fracturing in global climate models to accurately simulate future Arc-

tic sea-ice mass balance, particularly the evolution of MYI in the Arctic.

504

## **Data Availability Statement**

The output from the neXtSIM-OPA simulation is available as NetCDF files at https:// doi.org/10.5281/Zenodo.7277523. Jupyter Notebooks for data analysis and plotting are located in a public GitHub repository at https://zenodo.org/badge/latestdoi/ 682991902. The ERA-5 data was downloaded from the Copernicus Climate Change Service Climate Data Store (C3S) https://cds.climate.copernicus.eu/cdsapp#!/dataset/ reanalysis-era5-single-levels?tab=overview.

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# Breaking the Ice: Exploring the Changing Dynamics of Winter Breakup Events in the Beaufort Sea

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11	Key Points:
12	- Modelled leads in the Beaufort Sea during winter time are increasing at $4\%$ per decade
13	over the period 2000-2018
14	• The shift to thinner and younger sea ice, particularly after 2007, makes the Beau-
15	fort Sea more vulnerable to large breakup events by winds
16	• Winter breakup increases ice export from the Beaufort Sea and leads to a thin-
17	ner and weaker ice cover at the end of the cool season

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#### Abstract 18

The Beaufort Sea has experienced a significant decline in sea ice, with thinner first-year 19 ice replacing thicker multi-year ice. This transition makes the ice cover weaker and more 20 mobile, making it more vulnerable to breakup during winter. Using a coupled ocean-sea-21 ice model, we investigated the impact of these changes on sea-ice breakup events and lead 22 formation from 2000 to 2018. The simulation shows an increasing trend in the Beaufort 23 Sea lead area fraction during winter, with a pronounced transition around 2007. A high 24 lead area fraction in winter promotes a significant growth of new, thin ice within the Beau-25 fort region while also leading to enhanced sea ice transport out of the area. The export 26 offsets ice growth, resulting in negative volume anomalies and preconditioning a thin-27 ner and weaker ice pack at the end of the cool season. Our results indicate that large 28 breakup events may become more frequent as the sea-ice cover thins and that such events 29 only became common after 2007. This result highlights the need to represent these pro-30 cesses in global-scale climate models to improve projections of the Arctic. 31

### 32

## **Plain Language Summary**

The sea ice cover in the Beaufort Sea has been changing - it is getting thinner and 33 weaker. This makes the ice more likely to break apart from strong winds. Using a com-34 puter model, we study how these changes may have affected the frequency of large sea-35 ice breakup events from 2000 to 2018. We find that the amount of open areas in the sea 36 ice, called leads, is increasing during winter. This allows new, thin ice to form, but also 37 causes more ice to move out of the region under the action of winds and currents. This 38 movement of ice cancels the growth of new ice, resulting in less ice overall at the end of 39 winter in this region. Interestingly, these events became more common after 2007 and 40 the results suggests that bigger breakup events might happen more often as the sea ice 41 continues to thin. This study highlights how important it is to include these changes in 42 large climate models to better predict what might happen in the Arctic in the future. 43

#### 1 Introduction 44

Recent decades have seen dramatic reductions in the extent, age, and thickness of 45 Arctic sea ice (e.g. Kwok, 2018). Those changes are particularly pronounced in the Beau-46 fort Sea, which has experienced a rapid decline in sea ice extent and sea ice thickness, 47 during both summer and winter. There has been a notable shift in the composition of 48

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sea ice in the early-2000s (Babb et al., 2022), where the Beaufort ice cover transitioned 49 from a state which was dominated by thick and old multi-year ice (MYI) to an increas-50 ingly thinner, more fragmented and mobile seasonal ice cover around 2007 (Moore et al., 51 2022; Wood et al., 2013). This regime shift towards younger, thinner sea ice is affect-52 ing the dynamical properties of the ice cover (Zhang et al., 2012), reducing the ice's me-53 chanical strength, thereby making it more vulnerable to atmospheric forcing (Petty et 54 al., 2016) and contributing to the observed increase in sea ice deformation and drift speeds 55 in the Arctic Ocean and the Beaufort Sea in particular (Rampal et al., 2009; Kwok & 56 Cunningham, 2010; Spreen et al., 2011). These changes in sea ice properties and ice dy-57 namics have consequences for the stability and persistence of the Beaufort Sea ice cover, 58 potentially resulting in more frequent sea ice breakup and lead formation (Maslanik et 59 al., 2007). This potentially has important implications for the overall mass balance of 60 sea ice, ice-ocean interactions, and the Arctic climate system. However, due to the lack 61 of long-term observations and the difficulties in modelling sea-ice breakup, our knowl-62 edge is currently limited when it comes to understanding the relationship between these 63 changing sea ice characteristics and the frequency and intensity of breakup events. 64

During winter months, lead formation exposes the ocean to the colder atmosphere 65 resulting in large air-sea heat, moisture and gas fluxes. The intense heat loss from the 66 ocean promotes new ice formation, contributing to the sea ice mass balance in the Arc-67 tic winter (accounting for between 10 and 20% of the total ice growth in the Arctic dur-68 ing winter Heil & Hibler, 2002; Kwok, 2006). Recent estimates from Boutin et al. (2023) 69 show that this number could be as high as 25–35%. Brine rejection from sea ice forma-70 tion increases the stability of the Arctic halocline (Shimada et al., 2005), which protects 71 sea ice from melting by suppressing the entrainment of subsurface heat into the surface 72 layer. Brine-driven eddies under sea ice leads can affect thermohaline structure of the 73 mixed layer by transporting heat and salt laterally under the sea ice (Matsumura & Ha-74 sumi, 2008; Peralta-Ferriz & Woodgate, 2015). Leads are also key regions for marine bi-75 ological productivity due to increased access to sunlight, which is otherwise very limited 76 due to the presence of thick, snow-covered sea ice. For example, recent observations show 77 that leads in Arctic pack ice can enable early phytoplanktonic blooms (Assmy et al., 2017) 78 impacting primary production and Arctic marine food webs. This could become more 79 frequent due to thinner and more dynamic sea ice that is more vulnerable to breakup 80 (Fadeev et al., 2021). 81

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In the Beaufort Sea, leads regularly form throughout the winter season in response 82 to divergent sea ice motion driven by atmospheric weather systems or ocean currents (Lewis 83 & Hutchings, 2019; Jewell & Hutchings, 2023). Meanwhile, several large breakup events 84 have been identified from satellite observations in recent decades (most noteworthy in 85 the winters of 2013 (Beitsch et al., 2014; Rheinlænder et al., 2022) and 2016 (Babb et 86 al., 2019)). Wintertime breakup events are characterized by extensive fracturing of the 87 ice cover associated with atmospheric synoptic conditions persisting from a few days to 88 several weeks (Jewell & Hutchings, 2023). Such events have been shown to significantly 89 impact sea ice conditions in the Beaufort Sea, with potential implications for the wider 90 Arctic sea ice mass balance. The large breakup events in winter 2013 and 2016 resulted 91 in anomalous sea ice drift and enhanced ice export out of the Beaufort Sea (e.g. Babb 92 et al., 2016; Rheinlænder et al., 2022). This led to an overall reduction in the Beaufort 93 ice volume in April and a thinner, less compact ice cover prior to the onset of the melt-94 ing season. This conditions the ice cover for rapid summer melt (e.g. Maslanik et al., 95 2007) and could contribute to the low regional September sea ice area seen in recent decades 96 (Williams et al., 2016; Babb et al., 2019; Moore et al., 2022). 97

Winter breakup events can also have important consequences for the MYI cover. 98 Enhanced ice export during winter may increase the flushing of MYI through the Beau-99 fort Sea which increases the amount of MYI being advected into the region from the cen-100 tral Arctic (as was seen in 2013, e.g. Richter-Menge & Farrell, 2013). This could mo-101 bilize the oldest and thickest sea ice residing north of Greenland, also known as the Last 102 Ice Area, which is subsequently advected into the Beaufort Sea. For example, during sum-103 mer 2020/21 Moore et al. (2022) found large concentrations of thick and old ice in the 104 Beaufort Sea, which could be traced back to enhanced winter transport from the Last 105 Ice Area. Less MYI now survives through the summer melt season, making the Beau-106 fort Sea a major contributor to MYI loss in the Arctic (Howell et al., 2016; Babb et al., 107 2022). 108

Despite their importance, sea ice breakup and lead formation are generally not adequately reproduced in large-scale sea-ice and climate models (e.g. Spreen et al., 2017). This is partly due to the difficulty in representing small-scale deformation features, like cracks and leads, for horizontal resolutions coarser than  $\sim$ 5 km (Hutter et al., 2022). And while higher resolution sea-ice models (4–5 km) have demonstrated a certain degree of proficiency in representing the large-scale distribution of sea-ice leads in the Arctic (e.g.

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<sup>115</sup> Wang et al., 2016), they are currently considered too costly for global-scale climate mod-<sup>116</sup> els.

In this study we present a newly developed coupled ocean-sea ice model based on 117 the neXtSIM sea-ice model which employs a brittle sea-ice rheology making it partic-118 ularly suitable for simulating small-scale ice deformation and linear kinematic features 119 like fractures and leads in sea ice at comparatively low resolution (about 12 km here) 120 (Rampal et al., 2019; Bouchat et al., 2022; Ólason et al., 2022). Rheinlænder et al. (2022) 121 recently demonstrated neXtSIM's ability to provide a realistic and accurate represen-122 tation of sea ice fracturing and lead propagation associated with the 2013 breakup event 123 in the Beaufort Sea. The study highlighted that such extreme breakup events could be-124 come more frequent as the sea ice thins, raising concerns about the vulnerability of the 125 Beaufort ice cover. Here, we seek to understand how changes in the Beaufort sea-ice regime 126 during the early 21st century have affected the stability of the ice cover and the occur-127 rence of extreme breakup events focusing on the winters of 2000–2018. By addressing 128 this question, this study aims to provide new insights into the ongoing transformations 129 of the Beaufort Sea ice cover and its implications for regional sea ice volume, MYI cov-130 erage, and sea-ice transport. 131

## 132 2 Methods

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## 2.1 Model setup

The model used in this study is the new coupled sea-ice-ocean model recently pre-134 sented in Boutin et al. (2023). In brief, the ocean component is the Océan PArallélisé 135 model (OPA), which is part of the NEMO3.6 modelling platform (Madec, 2008). We use 136 the regional CREG025 configuration (Talandier & Lique, 2021), which encompasses the 137 Arctic and parts of the North Atlantic down to 27°N, and has 75 vertical levels and a 138 nominal horizontal resolution of  $1/4^{\circ} (\simeq 12 \text{ km} \text{ in the Arctic basin})$ . The sea ice com-139 ponent is neXtSIM, a state-of-the-art, finite element, sea ice model using a moving La-140 grangian mesh (Bouillon & Rampal, 2015; Rampal et al., 2016). Sea ice dynamics rely 141 on the Brittle Bingham-Maxwell (BBM) rheology described in Olason et al. (2022), while 142 sea ice thermodynamics are simulated following the Winton (2000) model. We refer to 143 Boutin et al. (2023) for detailed information about the model setup. 144

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As noted by Hutter et al. (2022), sea ice models generally struggle to simulate sea 145 ice dynamics when run at resolutions coarser than about 5 km; in particular, features 146 like fractures, shear zones, and lead openings. However, the BBM rheology has demon-147 strated its capability to reproduce deformations consistent with observations when run-148 ning at a resolution of O(10 km) in the neXtSIM model (Ólason et al., 2022) and in the 149 SI3 model (L. Brodeau, personal communication). Specifically, these models exhibit ex-150 cellent capability in accurately capturing the divergence rates associated with the open-151 ing of leads when using the BBM rheology (Ólason et al., 2022; Rheinlænder et al., 2022). 152

The simulation presented in this study is the same as in Regan et al. (2023) and 153 Boutin et al. (2023). The simulation starts in 1995 and ends in 2018. The first five years 154 were considered a spin-up period and disregarded for analysis. Atmospheric forcing is 155 taken from the hourly ERA5 reanalysis at a 1/4-degree horizontal resolution. This sim-156 ulation has been thoroughly evaluated in two recent publications (Boutin et al., 2023; 157 Regan et al., 2023). Boutin et al. (2023) showed that simulating key sea-ice quantities 158 like volume, extent, large-scale drift, and sea ice deformations are consistent with satel-159 lite observations. Regan et al. (2023) demonstrated that the simulation successfully re-160 produces the spatial distribution and evolution of observed MYI extent. They also found 161 a good agreement with observed estimates of the regional dynamic and thermodynamic 162 components of the winter sea ice mass balance from Ricker et al. (2021). 163

164

## 2.2 Lead definition

neXtSIM uses three ice categories; open water, young ice, and consolidated ice. Newly 165 formed ice, thinner than  $h_{\rm max}$  (here set to 18 cm), is assigned to the young ice category, 166 representing the formation and growth of frazil and young ice in open water. Ice in the 167 young ice category is transferred to the consolidated ice category as its thickness exceeds 168  $h_{\rm max}$  (see appendix A of Rampal et al., 2009). In winter, when the Beaufort Sea is fully 169 ice-covered, lead opening is the only way open water can be exposed to the atmosphere, 170 and young ice can be formed. Therefore, we assume (as in Rheinlænder et al., 2022; Boutin 171 et al., 2023) that open water and young ice formed in winter are a proxy for the pres-172 ence of leads in the model. We use this assumption to estimate the rapid growth of thin, 173 newly formed ice in open-water and thin ice regions. A grid cell is considered a lead when 174 the combined fraction of open water and young ice exceeds a critical threshold  $c_{lim}$ , thereby 175 excluding the thicker pack ice. We found that a value of  $c_{lim} = 5\%$  gives a reasonable 176

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- lead distribution. The sensitivity of the simulated lead fraction to the value of  $c_{lim}$  is included in the Supplementary Material. The total lead area fraction (LAF) can then be calculated by multiplying the lead fraction with the area of each grid cell. A snapshot of the simulated sea ice concentration and lead fraction on 25 March 2016 is shown in Fig. 1. Here, leads are clearly identified as areas of open water and newly formed ice, whereas the pack ice is associated with low lead fraction values. The LAF calculated over the Beaufort region for this instance is 21%, which means that 21% of the Beaufort area
- 184 is covered by leads.



Figure 1. Snapshots of sea ice concentration (%) and lead fraction (%) in the Beaufort Sea on March 25, 2016. The total lead area fraction (LAF) is calculated using a threshold value of 5% (see 2).

## 185 **3 Results**

186

## 3.1 Simulated changes in the Beaufort ice cover

Since the 2000s, the MYI extent in the Arctic has declined considerably (Fig. 2a). 187 During the period 2000–2018, the model simulates a decline in the winter MYI area, which 188 is part of a long-term negative trend in the Arctic as seen from satellite observations (e.g. 189 Babb et al., 2022). In the Beaufort Sea region (outlined in Fig. 2a), extensive areas of 190 thicker and older MYI were present during January–March in the early 2000s (i.e. 2000– 191 2004). For the later part of the simulation (years 2014–2018), the MYI extent is signif-192 icantly reduced and is consistent with the observed trend towards reduced MYI concen-193 tration (Howell et al., 2016). Both the average sea ice thickness and MYI concentration 194 computed over the Beaufort region exhibit considerable year-to-year variations (Fig. 2b), 195 but overall there is a shift towards thinner and younger ice types. The average winter 196 sea-ice thickness decreased from 1.9 m in 2000–2004 to 1.6 m in 2014–2018. Meanwhile, 197 some old and thick sea ice still remains located north of Canada and the Last Ice Area, 198 which are important source regions for MYI import to the Beaufort Sea (Moore et al., 199 2022). 200

The ice drift speeds in the Beaufort Sea are also increasing (Fig. 2c) with a 12% increase over the 2000–2018 period (not shown). Previously, this increase in ice drift speeds has been linked to thinning of the sea ice cover and enhanced deformation rates leading to more fracturing and lead opening (Rampal et al., 2009). In the following, we examine the impact of transitioning to a more seasonal and thinner ice cover on the formation of leads and sea-ice breakup in the Beaufort Sea.

207

## 3.2 Simulated changes in wintertime leads

We show the simulated wintertime LAF in the Beaufort Sea for the period 2000– 209 2018 in Fig. 3. The LAF shows a large day-to-day variability ranging from 5% to 40%, which reflects the intermittent nature of sea-ice fracturing. Winter-mean values (January– 211 March) generally fall between 10% and 25% with a climatological average of 20%. The lead fraction is generally higher in January and decreases during February and March 213 as the ice becomes thicker and more compact (Fig. 3b and c).



**Figure 2.** (a) Maps of simulated sea-ice thickness and MYI extent (yellow contour corresponding to the MYI fraction of 0.4) in the Arctic for January–March averaged from 2000–2004 and 2014–2018. (b) Time series of JFM-mean sea-ice thickness and MYI extent averaged in the Beaufort region. (c) Histogram of the winter sea-ice drift speed (cm s<sup>-1</sup>) distributions for the 2000–2004 and 2014–2018 period. The Beaufort region is bounded by three gates shown in (**a**); North gate (78°N), East gate (120°W), and West gate (160°W). The definition of the Beaufort region is the same as in Moore et al. (2022).

We find a statistically significant increase in wintertime lead occurrences (4.2% per decade) over the period 2000–2018 based on a simple linear regression analysis (Supplementary Fig. S1). However, the linear relationship becomes less significant when we consider individual months, likely due to the larger spread in the monthly data (Fig. 3b). Here, we note that the modelled LAF is affected by the cutoff value used in the lead definition (see section 2), but the choice of this cutoff value has no impact on the simulated variability and trends (Supplementary Fig. S2).

The LAF values simulated by neXtSIM are consistent with observations of sea-ice leads from MODIS at 1 km<sup>2</sup> spatial resolution with observed winter-mean (NovemberApril) lead fraction area ranging between 10–20% in the Beaufort Sea (Willmes et al., 2023). Willmes et al. (2023) also found a significant trend in leads over the 2002–2021 period, but only for April. It is worth noting that the MODIS observations have uncertainties due to contamination by clouds and will only see opening leads that are relatively large. This highlights the need for a dedicated intercomparison study to determine how to best use MODIS imagery to classify and evaluate lead formation in sea-ice models, but this is beyond the scope of this paper.

Based on the LAF time series, several larger breakup events can be identified, all 230 occurring after 2007. Most noteworthy are the years 2008, 2010, 2013, 2016, and 2018, 231 which have high average wintertime LAFs (green triangles in Fig. 3a). Many of these 232 events have also been identified in satellite observations (e.g. Jewell & Hutchings, 2023). 233 For example, large breakup events were observed in 2008, 2013 and 2016 and have been 234 described in earlier studies (Wang et al., 2016; Babb et al., 2019; Rheinlænder et al., 2022). 235 The breakup being simulated by the model in 2018 is not seen in observations and is likely 236 a result of too strong melting simulated by the model in the summer of 2016 (not shown), 237 leading to thinner sea ice that could break up more easily. During these events, the daily 238 LAF exceeds the 90th percentile (about 30%; Fig. 3c) for a period of more than 15 days 239 during winter (Fig. 3d). We therefore expect these events to have a significant impact 240 on the Beaufort ice cover. Meanwhile, smaller breakup events are also present in other 241 years. For example, 2005 and 2006 exhibit high LAFs (daily values exceeding 35%), but 242 these occurrences are relatively short-lived and result in low winter-mean values over-243 all. Consequently, they will likely have less impact. 244

Around 2007, we identify a shift in the interannual variability of the LAF based 245 on the monthly values in Fig. 3b. For the 2000–2007 period, the variability (shown by 246 the standard deviation for each month) ranges from 3.1–4.8% during winter. After 2007, 247 this increases to 3.9-7.4% for 2008–2018, and the average LAF increases during all win-248 ter months. This coincides with more extreme breakups during this period. March, in 249 particular, stands out, showing a 38% increase in the mean LAF relative to the 2000-250 2007 period while also exhibiting the highest variability (standard deviation of 7.4%). 251 We will investigate this in more detail in section 3.3 and identify some common char-252 acteristics of the simulated breakup events and their impacts (section 3.4). 253

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Figure 3. Simulated lead area fraction (LAF; %) in the Beaufort Sea from January through March. (a) Daily LAF (grey) from 2000–2018 with circles showing the winter-mean values along with the standard deviation. The dashed line shows the 2000–2018 winter climatology. (b) Monthly LAF climatologies for the 2000–2007 (blue) and 2008–2018 (orange) periods. Diamonds represent the monthly mean with the standard deviation in whiskers. (c) Histograms of monthly LAF distributions, with the 90th percentile (~30%) shown by the dashed line. (d) Stacked barplot of binned LAF from 20% to 45% where the height of the bars corresponds to the number of days. Numbers denote the total number of days where the daily average LAF exceeds the 90th percentile.

## 254

## 3.3 Driving mechanisms of sea-ice breakup events

The sea ice movement driven by wind and ocean currents can create stresses within the ice pack, leading to fracturing and the formation of leads (Lewis & Hutchings, 2019). In addition, changes in the material properties of the ice, such as ice thickness, concen-

tration, and strength, can also influence the susceptibility to breakup. Here we look at 258 the drivers of the simulated changes in winter lead formation, focusing particularly on 259 the shift occurring around 2007. 260

## 3.3.1 Winds

261

Jewell and Hutchings (2023) analysed the synoptic conditions during breakup events 262 in the Beaufort Sea during winters 1993–2013. They show a consistent connection be-263 tween wind forcing and lead formation, where a breakup is typically associated with high 264 sea level pressure and relatively strong anticyclonic winds over the Beaufort Sea. Sim-265 ilarly, Wang et al. (2016) concluded that a stronger Beaufort High results in stronger south-266 easterly winds in the Beaufort Sea, which pushes sea ice away from the coast and thus 267 promotes higher ice divergence and lead formation. 268

In general, high wintertime LAFs in the Beaufort Sea are linked with persistently 269 higher wind speeds (Figure 4b and c) in agreement with Jewell and Hutchings (2023). 270 In 2010, 2013, and 2016, the daily wind speed exceeds  $8.5 \text{ m s}^{-1}$  (one standard devia-271 tion above the mean) for more than 20 days (Fig. 4d). These conditions are typically 272 associated with a positive sea level pressure difference across the Beaufort Sea and east-273 erly winds (Supplementary Fig. S3), creating favourable conditions for off-shore ice drift 274 and enhanced breakup. It is worth noting that there is considerable variability between 275 the different months (Figure 4c and Supplementary Fig. S3), and we do not find a sim-276 ple relationship between wind speed and sea-ice breakup. Jewell and Hutchings (2023) 277 came to the same conclusion indicating that breakup may occur for a wide range of at-278 mospheric conditions. Both the duration of strong winds as well as the wind direction 279 appear to be important for initiating a breakup. For example, in 2005, conditions were 280 comparable to other breakup years (e.g. 2013), with relatively strong and persistent winds 281 during winter, however the LAF remained relatively low throughout the winter of 2005 282 (Supplementary Fig. S4). This could be due to the fact that the ice was thicker and stronger 283 (Fig. 2b) and thus less sensitive to winds. 284

285

We find no trend in the ERA5 winds in the Beaufort region during winter over the period 2000 to 2018 (Fig. 4b). The year-to-year variability is also relatively similar for 286 the 2000–2007 and the 2008–2018 winter periods. The same is true if we consider indi-287 vidual months rather than the winter-mean values (Fig. 4c), showing no major differ-288

-12-

ence in wind strength for January versus March. Thus, changes in the wind forcing cannot explain the fact that we are seeing the strongest increase in lead variability in March. Based on this, we speculate that the shift in the lead formation dynamics seen after 2007 is linked to the thinning of the Beaufort ice cover (Fig. 2), making it more vulnerable to atmospheric forcing during winter.

294

## 3.3.2 Changes in ice conditions

Changes in the dynamic properties of sea ice are generally attributed to changes 295 in concentration and thickness, which in turn affect the strength of the ice (Zhang et al., 296 2012). In sea ice models, including neXtSIM, ice strength is parameterized as a function 297 of ice thickness and concentration (Hibler, 1979). Therefore, we expect the simulated de-298 cline in ice thickness (Fig. 2b) to weaken the ice pack and reduce the internal ice stress. 299 As a result, this leads to an overall increase in the simulated deformation rates (Fig. 5b) 300 and increased drift speeds (Fig. 5a) in the Beaufort Sea. The positive trend in sea ice 301 drift speeds is consistent with earlier modelling studies (e.g. Zhang et al., 2012) and ob-302 servations (e.g. Rampal et al., 2009; Spreen et al., 2011). 303

Comparing the simulated LAF in Fig. 4a to the time series of the mean ice speed 304 and deformation rates (Fig. 5) strongly points to a shift in the dynamic sea ice prop-305 erties and lead formation dynamics after 2007. Both sea ice drift and deformation rates 306 show a pronounced change in the variability, fluctuating between relatively low and high 307 values mirroring the changes in the LAF. High LAF is associated with high deformation 308 rates and increased ice speed in the Beaufort Sea (Fig. 5a and 5b), exceeding  $7 \text{ cm s}^{-1}$ 309 during breakup events. It is worth noting that both drift and deformation values remain 310 relatively low for winters without significant breakup, i.e. their baseline values do not 311 seem to change much over the 2000–2018 period. This suggests that individual extreme 312 events can substantially alter the overall trend seen in the data. 313

By plotting the ratio between ice drift to wind speed (Fig. 5c) we see a clear increase in the ice drift to wind speed ratio, especially during breakup events. This reflects an increased sensitivity of the Beaufort ice cover to wind forcing during the late 2000s, whereas in the earlier period (e.g. in 2005 when the ice was thicker), there is a larger disconnect between strong winds and ice drift speeds. Overall, this points to changes in the

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Figure 4. Time series of (a) normalized wintertime lead fraction (%) in Beaufort region, (b) ERA5 daily mean wind speed ( $m s^{-1}$ ) between 70–75°N along 150°. The January–March average wind speed is highlighted by blue circles. (c) Monthly mean wind speed for January, February and March. (d) The number of days where the daily wind speed exceeds 8.5 m s<sup>-1</sup> (one standard deviation above the mean) is shown in bars. Its mean and standard deviation is shown by solid and dashed lines. The solid black line shows the wintertime (January through March) mean wind speed. The time series in (a) is normalized by subtracting its mean and dividing by its standard deviation. The transect used for calculating the winds is indicated in Figure 4b.

- material properties of sea ice being a major factor in driving the shift we see in the sim-
- <sup>320</sup> ulated lead area fraction.



Figure 5. Time series of (a) mean ice speed (cm s<sup>-1</sup>), (b) total deformation rate (1/day) and (c) ratio (%) between mean sea-ice drift and mean wind speed. The thin black line shows the average wind speed. All time series are based on wintertime (January through March) means and averaged over the Beaufort region. The area used for averaging is shown in Fig. 2. Green triangles highlight winters with significant breakup identified in Figure 3a.

## 321

## 3.4 Impacts on Beaufort ice volume and MYI

322 323 In this section, we seek to understand how winter sea ice breakup impacts the ice volume in the Beaufort Sea. Changes in regional ice volume during winter can be sep-

arated into two terms: (i) thermodynamic ice growth and (ii) sea ice transport. Note that we are omitting the term associated with sea ice melting as this can be considered negligible during winter (Graham et al., 2019).

327

## 3.4.1 Thermodynamic ice growth in leads

In winter, the opening of leads results in intense heat loss from the underlying ocean and promotes new ice formation. More sea ice breakup could therefore increase the local winter ice growth overall and modulate the composition of the Beaufort ice pack by increasing the fraction of thinner and younger ice types.

The thermodynamic ice growth from January through March is shown in Fig. 6 332 for leads and pack ice, respectively, where growth in leads is associated with the forma-333 tion of new and thin sea ice (Rheinlænder et al., 2022). Note that the ice growth in leads 334 is independent of the lead detection algorithm. Overall, the growth of new ice in leads 335 is increasing over the period 2000–2018 (Fig. 6b). We find a statistically significant lin-336 ear trend of 4% per decade for wintertime ice production in leads. This is consistent with 337 the results of Boutin et al. (2023), who used the same model to find a 4.3% per decade 338 trend on the pan-Arctic scale. Our result is also consistent with the simulated trend in 339 LAF (4.2%; Supplementary Fig. S1) and suggests that leads play an increasingly key role 340 in the local sea-ice volume budget as the ice cover becomes thinner and more fractured. 341

By comparing the growth estimates to the LAF time series in Figure 3a, we see that winters with more breakups also have larger ice production overall. New ice production in leads is consistently higher for these years (top 5) compared to the climatology, and the fraction relative to the total growth is above 40% (except for 2008). Overall, these results show that winter breakup can significantly increase the local ice volume by enhancing ice growth.

348

## 3.4.2 Volume transports

The simulated winter-mean ice volume fluxes are shown in Fig. 7 across the Beaufort Sea's eastern, northern and western gates. The transport has been separated into FYI and MYI contributions. The western gate captures primarily the ice export from the central Beaufort to the Chukchi Sea, whereas the northern and eastern gate captures the import of thicker and older sea ice from the Canadian Archipelago and the LIA. This

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Figure 6. (a) Thermodynamic sea-ice growth (km<sup>3</sup>) in leads (blue) and pack ice (orange) in the Beaufort Sea during winter (January–March). The dashed line shows the winter climatological mean ice volume growth in leads. Numbers in white indicate the fraction (in %) of the respective growth relative to the total growth (shown by black numbers above the bars). (b) Fraction of ice volume growth in leads relative to the total thermodynamic growth. A linear regression model has been fitted to the data and is shown by the grey line. Breakup years are highlighted by green triangles in (a).



Figure 7. Time series of January–March ice volume fluxes (km<sup>3</sup>) into the Beaufort Sea across the eastern (160°W), northern (78°N), and western gates (120°W) and the total fluxes in the bottom panel. The transport is separated into contributions from FYI (light blue) and MYI (dark blue), with dashed lines showing the winter-mean climatologies. Positive values indicate sea ice transport into the Beaufort region.

circulation pattern is associated with the anti-cyclonic circulation of the Beaufort Gyre (Howell et al., 2016). Overall, the total fluxes from January through March show a net import of MYI and a net export of FYI. The majority of the transport occurs at the northern and western gates, while the transport through the eastern gate is generally small (only accounting for about 6% of the total import).

During winters with enhanced lead activity, the ice transport through the Beaufort region increases (bottom panel in Fig. 7) and is consistent with the higher mean ice drift speeds seen in Fig. 5a. At the western gate, there is a large export of primarily thinner and younger FYI during breakup events. A smaller fraction of MYI is also exported, especially in the early 2000s, but is reduced during the later part of the simulation in line with previous studies (e.g. Howell et al., 2016; Babb et al., 2022).

At the northern gate, the simulated volume transports are primarily dominated by 365 the import of MYI from the central Arctic. There is generally a higher MYI transport 366 into the Beaufort Sea during breakup events, for example, in 2013 which shows a net im-367 port of 85  $\mathrm{km}^3$  across the eastern and northern gates. In comparison, the average MYI 368 import is  $\sim 45 \text{ km}^3$  over the period 2000–2018. However, years with relatively low lead 369 fractions (in the early 2000s and in 2015) also show high MYI import (and export), while 370 the breakup events in 2016 and 2018 have lower MYI fluxes despite high lead fractions. 371 A possible explanation is that there is simply less MYI in the Arctic in total and, there-372 fore, less to be transported into the Beaufort region (Babb et al., 2022). This is likely 373 the case in this simulation, which underestimates MYI extent in 2017 and 2018 partly 374 due to unrealistically high melting in the summer of 2016. 375

In total, ice export is larger than import during winter breakup events, which suggests that sea ice breakup contributes to regional dynamic ice loss in the short term. This will likely also affect ice transport in the following months and impact the regional ice volume before the beginning of the next melt season. For example, winter export from the Beaufort region could lead to enhanced flushing of MYI through the Beaufort Sea (e.g. Babb et al., 2019) providing dynamical replenishment for the ice loss during winter.

To understand the cumulative effects of winter breakups, we look at the transport into the Beaufort Sea through the entire cool season from January through June in Figure 8. The 2000–2018 climatology shows a net ice export at the end of the cool season

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**Figure 8.** Cumulative sea ice volume fluxes from January to June for all years from 2000 to 2018. Years with high wintertime lead fractions (2008, 2010, 2013, 2016 and 2018) are shown in colours. The 2000–2018 climatology is shown in black with the ±1 standard deviation.

(about -200 km<sup>3</sup> in June). Years with higher lead activity in winter (2008, 2010, 2016, 386 2018) exhibit larger cumulative net ice export (more than one standard deviation be-387 low the mean). A notable exception is 2013, which shows as net ice import, despite a 388 large export in February–March (Babb et al., 2019; Rheinlænder et al., 2022). This was 389 caused by enhanced advection of thicker MYI through the northern boundary from mid-390 April (Fig. 9) offsetting ice export (primarily thin FYI) at the western boundary. Mean-391 while, the other breakup events show little evidence of MYI flushing associated with the 392 increased winter ice export from the Beaufort Sea. 393

Overall, these results suggest that winter breakup events may have a negative im-394 pact on the Beaufort ice mass balance by enhancing ice export, despite also promoting 395 significant new ice growth. But what is the combined effect of winter breakup on the Beau-396 fort ice volume? In Figure 10, we show the relationship between wintertime lead area 397 fraction, i.e. over the period from January through March, and the ice conditions in the 398 Beaufort Sea at the end of the cool season (from January 1 to June 1). The cumulative 399 cool season transport out of the Beaufort Sea is generally larger when the LAF is high 400 and we see a clear separation of the breakup years that also have enhanced export. The 401

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**Figure 9.** June 1 cumulative ice volume fluxes separated into FYI (blue) and MYI (orange) contributions from 2000 to 2018 in the Beaufort region. The net June 1 ice volume flux is shown by black circles. Positive (negative) values corresponds to a net import (export).

402	correlation between cool season transport and the Beaufort ice volume at the beginning
403	of June exhibits a similar grouping, with breakup years showing lower ice volume val-
404	ues on June 1. A similar relationship was also found based on satellite observations (e.g.
405	Babb et al., 2019; Moore et al., 2022), suggesting that high winter export from the Beau
406	fort Sea results in an anomalously thin ice cover and negative regional volume anoma-
407	lies. This could preconditioning the ice cover for increased summer melt and ultimately
408	result in record low regional September sea ice minima as shown by Babb et al. (2019).

## 409 4 Discussion

In an earlier modelling study, Wang et al. (2016) simulated the time evolution of 410 lead formation in the Beaufort Sea over the last three decades (1985–2014) using a high-411 resolution (4.5 km) sea-ice model (Finite Element Sea Ice-Ocean Model; FESOM). In 412 contrast to the neXtSIM simulation, they observed no increase in the number of large-413 scale breakup events in winter, which they related to the absence of wind stress trends 414 in the Beaufort region. However, one of the notable contrasts between these two mod-415 els is the difference in sea ice rheology; m-EVP for FESOM versus BBM in neXtSIM. 416 This could lead to significant differences in how the ice cover responds dynamically to 417 changes in the mechanical ice properties and the sensitivity to wind forcing. Another dif-418 ference is the definition of leads in Wang et al. (2016), which are defined as locations where 419 the sea ice is at least 20% thinner than at its surroundings (within a 25-km radius). Firstly, 420



Figure 10. Scatterplot of the cool season (January 1 - June 1) ice volume flux and (a) winter mean (January–March) lead area fraction and (b) June 1 ice volume. Breakup years (2008, 2010, 2013, 2016 and 2018) are highlighted in colours.

this excludes very wide leads, and secondly may fail to capture very localized divergence 421 events leading to drops in sea ice concentration as seen in the neXtSIM simulation. Mean-422 while, recent observational data based on MODIS imagery (Willmes et al., 2023) show 423 a significant positive trend of 2% per decade in lead frequencies in the Beaufort Sea over 424 the period from 2002 to 2021 (during April only). This is similar to Hoffman et al. (2022) 425 who observed a small, but significant increase in pan-Arctic leads from satellite data over 426 the same period, despite large uncertainty due to the increasing cloud cover in the Arc-427 tic. 428

We find that the change in lead formation dynamics simulated by the neXtSIM model, 429 notably the increased variability in lead formation after 2007, can be linked to a shift 430 in the Beaufort ice dynamics. Such transitions have been reported in observations. Long-431 term sea ice data from satellites dating back to the 1980s show evidence that the Beau-432 fort Sea transitioned to a thinner state in 1998 (Hutchings & Rigor, 2012 and Fig. 1 in 433 Babb et al., 2019). Another transition occurred around 2007 (e.g. Moore et al., 2022; 434 Babb et al., 2022), which reflected a shift from an old ice regime (1979–2007) when the 435 region was dominated by MYI to a young ice regime (2007-present). Similarly, Wood 436 et al. (2013) pointed to a "new normal" climate in the Beaufort Sea since 2007, char-437

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acterized by an increasingly mobile and thus more dynamic ice pack, which agrees with
the increase in ice drift in Figure 2c.

Our results indicate, that sea ice thinning and loss of MYI in the Beaufort region 440 makes the ice cover less resilient to wind forcing thus increasing the likelihood of large 441 breakups. This could lead to enhanced inter-annual variability in Beaufort Sea ice con-442 ditions and may increase the potential for rapid sea ice loss (Moore et al., 2022; Maslanik 443 et al., 2007). Similarly, Petty et al. (2016) found an amplified sensitivity of the Beau-444 fort sea ice circulation in winter to wind forcing during the late-2000s. This increase in 445 winter ice drift is commonly attributed to general sea ice thinning and reduction in me-446 chanical ice strength (Zhang et al., 2012; Rampal et al., 2009), which is also evident from 447 our results in Fig. 5. Meanwhile, Jewell and Hutchings (2023) noted that changes in ice 448 thickness is not the only factor controlling breakup. Atmospheric conditions such as wind 449 direction, storm propagation and duration of strong winds are also important factors that 450 contribute to sea-ice breakup. In fact, the LAF timeseries in Figure 3 show that breakup 451 events also occurred during the early-2000s (for example in 2005 and 2006) when the ice 452 was considerably thicker. This emphasizes the importance of atmospheric forcing in ini-453 tiating breakup. 454

While the atmosphere plays a dominant role in triggering sea-ice breakup on short 455 time scales (days to weeks), the ocean may also play a role in preconditioning sea-ice breakup 456 on seasonal time scales (Willmes et al., 2023). For example, enhanced ocean heat fluxes 457 during summer and autumn may predispose the ice cover to enhanced melt, resulting 458 in a thinner and weaker ice cover before the beginning of the cooling season (Herbaut 459 et al., 2022; Graham et al., 2019). Lead formation can also be expected to have signif-460 icant impacts on the ocean underneath, for example by enhancing mechanical energy in-461 put available for mixing and through brine formation thereby affecting mixed layer prop-462 erties and halocline stability (Matsumura & Hasumi, 2008; Peralta-Ferriz & Woodgate, 463 2015; Shimada et al., 2005). Mixing up warmer waters through lead opening could en-464 hance basal melting and limit new ice growth in the leads (e.g. Graham et al., 2019). 465 Such feedbacks could be important for ice-ocean interactions even on longer time scales 466 but they are not assessed explicitly in this study. 467

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## <sup>468</sup> 5 Summary and conclusions

This paper presents a multi-decadal simulation using the coupled ocean-sea-ice neXtSIM-OPA model and investigates the temporal changes in wintertime sea ice leads in the Beaufort Sea and their impacts. The simulation shows a small but significant increasing trend in the Beaufort lead area fraction (4% per decade) over the winter season (January through March) for the period 2000 to 2018. This is consistent with a general decrease in ice thickness and MYI cover as well as enhanced drift speeds during winter in the Beaufort region.

Around 2007 we find a notable increase in the simulated lead area fraction vari-476 ability associated with enhanced sea ice breakup, high deformation rates and an increase 477 in the mean ice velocities. These changes coincide with the observed regime shift that 478 occurred in the Beaufort Sea in 2007 (Wood et al., 2013; Moore et al., 2022), character-479 ized by a transition from a state dominated by thicker and older MYI towards more sea-480 sonal, thinner and younger sea ice. We find no significant trend in the surface winds dur-481 ing winter over the simulated time period. This suggests that the changes in lead for-482 mation dynamics can be attributed to changes in the sea ice conditions (i.e. thinning, 483 loss of ice strength and enhanced deformation) rather than changes in the atmospheric 484 forcing. Consequently, the ice cover becomes more sensitive to wind forcing, which may 485 lead to enhanced inter-annual variability in Beaufort Sea ice conditions and more extreme 486 breakup during winter. 487

Several large breakup events are identified which significantly impact the regional 488 thermodynamic ice production, with new ice growth in leads contributing up to 40% of 489 the total winter ice growth. This implies that sea ice leads play an important role in the 490 local ice mass balance in the Arctic (as Boutin et al., 2023, also found). Meanwhile, years 491 with high lead activity in winter consistently exhibit increased ice export, primarily FYI, 492 from the Beaufort Sea throughout the entire cool season (January 1 to June 1). While 493 some breakup events also show an enhanced import of MYI to the Beaufort from the cen-494 tral Arctic (e.g. in 2013), we find no consistent evidence that winter breakup leads to 495 the flushing of MYI through the Beaufort Sea. 496

<sup>497</sup> Overall, these results suggest that winter breakups have a negative impact on the <sup>498</sup> Beaufort ice volume, preconditioning a thinner and weaker ice pack at the end of the cool <sup>499</sup> season (see also Babb et al., 2016, 2019; Moore et al., 2022). This could lead to earlier

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<sup>500</sup> breakup in spring and enhanced summer melt, thereby contributing to accelerating sea

- <sup>501</sup> ice loss in the Beaufort Sea. This further highlights the need to include small-scale sea-
- <sup>502</sup> ice deformation and fracturing in global climate models to accurately simulate future Arc-

tic sea-ice mass balance, particularly the evolution of MYI in the Arctic.

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## **Data Availability Statement**

The output from the neXtSIM-OPA simulation is available as NetCDF files at https:// doi.org/10.5281/Zenodo.7277523. Jupyter Notebooks for data analysis and plotting are located in a public GitHub repository at https://zenodo.org/badge/latestdoi/ 682991902. The ERA-5 data was downloaded from the Copernicus Climate Change Service Climate Data Store (C3S) https://cds.climate.copernicus.eu/cdsapp#!/dataset/ reanalysis-era5-single-levels?tab=overview.

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# Supporting Information for "Breaking the Ice: Exploring the Changing Dynamics of Winter Breakup Events in the Beaufort Sea"

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1. Figures S1 to S4

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**Figure S1.** Simulated lead area fraction in the Beaufort Sea (%) for January, February, March and the JFM-mean. Linear trend is shown by the grey line.



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**Figure S2.** Sensitivity of lead area fraction to different cutoff values. The default used is 5%. Lead fractions are calculated as the fraction of open water and thin, young ice for January-March.



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Atmospheric conditions during breakup events (years 2008, 2010, 2013, 2016, and Figure S3. 2018). Monthly average ERA5 wind speed (shading), direction (vectors) and sea level pressure contours for January, February and Mugust Contours for January, February and Mugust Contours are Shown at 2 hPa intervals.



**Figure S4.** Atmospheric conditions in Beaufort Sea for the 2005 winter, showing ERA5 wind speed (shading), direction (vectors) and sea level pressure contours (shown at 2 hPa intervals) for January, February and March.

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