The Role of Climatological State on Driving US Heat Waves Through Rossby Waves Packets

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Abstract

While heat waves are local extreme weather events, a planetary Rossby wave pattern is statistically related to the occurrence of heat waves events in the U.S. However, whether such planetary wave patterns cause the enhanced statistics of local heat waves or as a coincidence is debatable. In this work, we hypothesize that the atmospheric climatological state dictates the slowly propagating wave pattern, which sets up a conducive large-scale environment for local US heat waves. We implement an idealized dry dynamic core model with an iterative approach to simulate the realistic North American summer climatological state. As the model can generate similar large-scale planetary wave patterns propagating throughout North America, significantly more heatwaves are generated, and the statistics of heat waves become consistent with that estimated in reanalysis products. The slowly propagating Rossby wave packets with a timescale of 20-30 days can serve as a new source of intraseasonal predictability.

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5	Key Points:
6	• The climatological state drives the stationary Rossby wave propagation provid-
7	ing critical dynamical conditions for heat waves in the US.
8	• A dry atmospheric model with a corrected climatological state generates heatwaves
9	that are statistically consistent with observations.
10	• The slowly propagating Rossby wave packets with a timescale of 20-30 days can
11	be a source of intraseasonal predictability.

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12 Abstract

While heat waves are local extreme weather events, a planetary Rossby wave pattern is 13 statistically related to the occurrence of heat waves events in the U.S. However, whether 14 such planetary wave patterns cause the enhanced statistics of local heat waves or as a 15 coincidence is debatable. In this work, we hypothesize that the atmospheric climatolog-16 ical state dictates the slowly propagating wave pattern, which sets up a conducive large-17 scale environment for local US heat waves. We implement an idealized dry dynamic core 18 model with an iterative approach to simulate the realistic North American summer cli-19 matological state. As the model can generate similar large-scale planetary wave patterns 20 propagating throughout North America, significantly more heatwaves are generated, and 21 the statistics of heat waves become consistent with that estimated in reanalysis prod-22 ucts. The slowly propagating Rossby wave packets with a timescale of 20-30 days can 23 serve as a new source of intraseasonal predictability. 24

²⁵ Plain Language Summary

Heatwaves are the leading weather-related killer in the United States and affect mainly 26 the most vulnerable communities. These extreme events are statistically related to a large-27 scale wave pattern of Rossby waves that features a zonal wavenumber five structure. To 28 find out what controls this wave pattern, we use a simple general circulation model that 29 only contains dry dynamics without complicated interactions with moisture or clouds. 30 31 We found that by modifying the climatological state of temperature and velocity fields based on the observed structure from the Northern Hemisphere summer, we can observe 32 the same wave number five pattern developing days before the heatwave events, resem-33 bling that from the observations. This result suggests that the climatological state of the 34 Northern Hemispheric summer provides a conducive environment for heatwaves in the 35 U.S. due to the slowed-down propagation speed of the Rossby wave packets. A deeper 36 understanding of its dynamics is crucial because, as this pattern develops up to 20 days 37 ahead of the extremes, the underlying physical process governing this pattern may serve 38 as a source of predictability on the Subseasonal-to-Seasonal (S2S) timescale, a current 39 gap of forecasts between weather and climate. 40

41 **1** Introduction

42 Over the coming century, climate change is expected to increase average summer 43 temperatures and the severity of extreme heat linked with heat waves events. In the United 44 States, the frequency, intensity, and duration of heat waves have been increasing rapidly 45 in recent decades, and this behavior is projected to continue in the next decades (Meehl 46 & Tebaldi, 2004). Currently, forecasters in the United States can only predict extreme 47 events up to 10 days in advance because, unlike in the tropics, circulation in the mid-48 latitudes is more chaotic as it is dominated by climatic noise (Feldstein, 2000).

However, it is intriguing to hypothesize that certain atmospheric circulation states 49 can be substantially more predictable than the average scenario, because these circula-50 tion regimes are associated with low-frequency patterns (Schubert et al., 2011). With 51 this in mind, and considering that different studies have suggested that propagating sta-52 tionary Rossby Waves play an important role in the mid-latitude atmospheric variabil-53 ity, the scientific community recently has increased interest in the connection between 54 extreme weather events and large-scale atmospheric patterns such as Rossby Wave Pack-55 ets (Chen & Newman, 1998; Schubert et al., 2011; Fragkoulidis & Wirth, 2020). 56

In many of the investigated cases, the extreme weather was linked to an upper-tropospheric
trough (i.e., a breaking Rossby Waves). For example, Chen and Newman (1998) suggests
that Rossby Waves originating in the west Pacific were the key in initiating intense anomalous anticyclones during the 1988 U.S. drought. Also, Schubert et al. (2011) relates this

large-scale pattern with monthly mean precipitation and surface temperature variabil-61 ity over many regions of the extratropical land areas, including the northern U.S., parts 62 of Canada, Europe, and Russia. Following the same approach, Ding and Wang (2005) 63 found an interannually varying Northern Hemisphere circumglobal pattern with a preferred wavenumber five structure. This particular structure is confined within the waveg-65 uide associated with the summer north jet stream in the stationary state and is linked 66 to significant surface air temperature and rainfall anomalies in western Europe, Euro-67 pean Russia, India, East Asia, and North America. In a more recent study Teng et al. 68 (2013), more concrete evidence has been shown of the relationship between the occur-69 rence of heatwayes in the U.S. and the same wavenumber five pattern developing as early 70 as 20 days ahead of the events. Despite the growing evidence for the concurrences of this 71 interesting wave number five pattern, the physical mechanism that drives the atmospheric 72 pattern and thus determines heatwaves has yet to be understood. 73

This study aims to understand better the fundamental role of atmospheric dynam-74 ics in the evolution of U.S. heat waves based on the hypothesis that the climatological 75 state for the Northern Hemisphere summer drives the stationary Rossby wave propaga-76 tion providing critical dynamical conditions for heat waves in the U.S. Inspired by the 77 motivation to look for the simplest possible model in the climate model hierarchies (Held, 78 2005), we adopt an idealized dry dynamical core that allows for a bias correction of the 79 mean flow structure without physical parameterizations. This approach allows us to iso-80 late the role of intrinsic planetary waves in the evolution of extreme events. 81

⁸² 2 Data and Methods

2.1 Reanalysis products

For the observational analysis, the NCEP/NCAR reanalysis of the National Oceanic and Atmospheric Administration (NOAA) was implemented. The surface air temperature (SAT), all levels of temperature and the meridional and zonal components of the wind at 300 hPa (V_{300} and U_{300}) were obtained from these databases with a spatial resolution of 2.5° and daily temporal resolution from 1948 to 2023.

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2.2 An Idealized GCM with realistic climatological basic state

Idealized general circulation models are commonly used for the study of atmospheric dynamics. A dry dynamical core model solves the primitive equations on the sphere by nudging the climatological temperature field toward a prescribed structure of radiative equilibrium temperature (T_{eq}) . Essentially, this process isolates the dry atmospheric dynamics from the complex physical parameterizations.

⁹⁵ We use the open-access updated version from (Wu & Reichler, 2018) of the spec-⁹⁶ tral dynamical core model proposed by (Held & Suarez, 1994) for the Geophysical Fluid ⁹⁷ Dynamics Laboratory (GFDL). The model has a horizontal resolution of T42 (64x128 ⁹⁸ grid) and 40 vertical σ levels between the surface and 0.01 hPa. To represent the boundary-⁹⁹ layer friction, Rayleigh drag is used to remove momentum in the lower troposphere be-¹⁰⁰ tween the surface and $\sigma = 0.7$. By default configurations of this model, the tempera-¹⁰¹ ture is forced by Newtonian relaxation toward a prescribed equilibrium temperature as ¹⁰² follows:

$$\frac{\partial T}{\partial t} = \frac{T - T_{eq}}{\tau} \tag{1}$$

where τ is the prescribed relaxation timescale.

Since the primary purpose of this study is to examine how the atmospheric basic state controls the large-scale Rossby Wave pattern with a zonal wavenumber five struc-

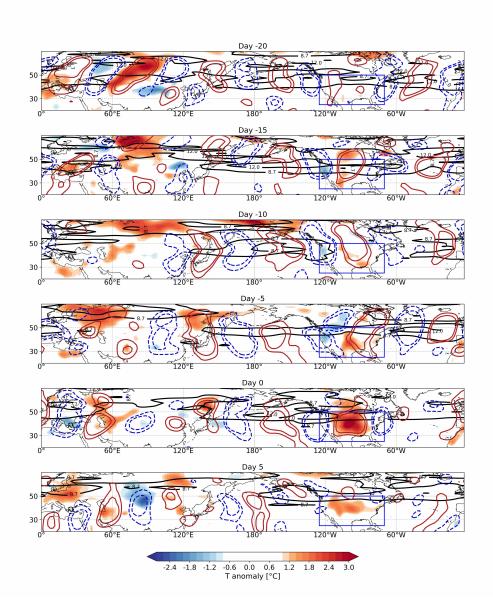


Figure 1. Evolution of heat waves 20 days ahead of the events: Red and dashed blue contours represent composites of 200hPa streamfunction anomalies at \pm 0.05, 0.08 (red as positives). Shading represents surface air temperature anomaly and black contours represent the Rossby Wave Packet envelope at levels 8.7, 12 m/s.

ture driving local heat waves in the US, the model includes the iterative procedure pro-106 posed by Chang (2006) to simulate the climatological basic state of the atmosphere in 107 the idealized model. This consists of iterating the radiative equilibrium temperature pro-108 file so that at the end of the iterations, the model climate closely resembles the desired 109 target climate (see Supporting Information for more details of the methodology). The 110 iteration uses a fixed equilibrium temperature for each N step (T_{eq}) in a run of 62 years, 111 for a total of 21630 days after eliminating the first 1000 days of simulation. Then, we 112 calculate the model simulated temperature climatology $T_{(N)}$ and correct it concerning 113 basic state from the NCEP reanalysis data T_R . The next iteration step N+1 is calcu-114 lated according to: 115

$$T_{eq(N+1)} = T_{eq(N)} - \frac{2}{3}(T_{(N)} - T_R), N = 1, 2, 3...$$
(2)

To assess the role of the climatological basic state, we carry out three experiments as follows:

- Held and Suarez 1994 configuration (CTR): the zonal-mean temperature 118 structure is relaxed toward a prescribed equinoctial radiative-equilibrium state. 119 It's the default configuration described in detail in (Held & Suarez, 1994) and leads 120 to a climatological state close to the observed annual mean climatological state. 121 **Zonally-symmetric boreal summer (ZOB)**: The ratiative forcing T_R described 122 in section 2.2 and in eq. 2 is set by the two-dimensional (i.e., latitude-pressure) 123 basic state from the NCEP reanalysis data. For this experiment, we use the zon-124 ally symmetrical distribution of temperature for the climatological boreal sum-125 mer (June, July, and August) between 1948 and 2021. 126 Realistic Northern Hemisphere summer (ROB): The full three-dimensional 127
- 128 structure of the radiative forcing T_R is obtained from the climatological boreal sum-129 mer from NCEP, especially including zonal asymmetries. This bias correction in-130 cludes zonal variations in T_{eq} , which can be related to zonal variations in diabatic 131 heating as discussed in Chang (2006) work.

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2.3 Heat waves identification

A heatwave event is set as at least five consecutive days with more than around 900000 km^2 over the US continental area (125W - 70W, 25N - 50N) with daily averaged SAT exceeding a threshold value, and the center of these warm points, considered as the point of maximum temperature, does not move faster than 5 latitude or longitude per day. The temperature threshold varies spatially as well as with the day of the year. It was established as the percentile 97.5 of the historical SAT. This approach follows Teng et al. (2013).

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2.4 Rossby Wave Packets evolution

As this study aims to understand the dynamics behind heat waves, it is necessary 141 to look at the evolution of the Rossby Waves Packets driving extreme events from ob-142 servational data and modeling outputs. Both the eddy phase speeds and their group ve-143 locity are important physical properties of Rossby Wave Packets' propagation. In par-144 ticular, the eddy phase speed indicates the propagation speed of individual troughs within 145 Rossby Waves Packets. It can thus be critical for the persistence of extreme surface weather 146 (Röthlisberger et al., 2019). On the other hand, the group velocity reflects how quickly 147 the whole packet propagates locally and represents the rate at which the Rossby waves 148 transfer energy horizontally (Pedlosky, 2003). Based on the Rossby Waves Packets' pro-149 gression in longitude and time, we compute the envelope of the meridional wind indi-150 cating their preferred regions of formation and decay. Specifically, we estimate the Rossby 151

Waves Packets' propagation by constructing Hovmöller diagrams using streamfunction
 anomalies at 300 hPa and the corresponding Rossby Waves Packets' envelope.

Subsequently, we calculate the local (in space and time) group velocity using the 154 Rossby Waves Packets' envelope, which reflects how the flow features of enhanced merid-155 ional wind amplitude propagate in the zonal and meridional directions (Fragkoulidis & 156 Wirth, 2020). We implemented the method proposed by Zimin et al. (2003) for the cal-157 culation of Rossby Waves Packets' envelope involving the Hilbert transform along cir-158 cles of constant latitude combined with a restriction of the zonal wavenumber to a spec-159 ified interval, which in this case corresponds to wavenumbers between 3 and 11. We im-160 plement this methodology using the meridional wind, which is usually implemented for 161 the diagnosis of meridional deviations from the zonal flow, and is particularly well suited 162 for the detection of Rossby waves. Specifically, we take the upper-tropospheric (300 hPa) 163 meridional wind which features strong Rossby Waves Packet's activities (Wirth et al., 164 2018). The meridional wind anomaly v(x) is considered on an equidistant grid along a 165 latitude circle, where $x = 2\pi/N$ with $0 < x \le 2\pi$, N is an even integer, and l = 1, 2, ..., N. 166 The Fourier transform of the real function v(x) is computed as: 167

$$\hat{v}_k = \frac{1}{N} \sum_{l=1}^N v\left(\frac{2\pi l}{N}\right) e^{-2\pi i k/N}, \left(K = -\frac{N}{2} + 1, \dots, \frac{N}{2}\right)$$
(3)

The inverse Fourier transform is applied to a selected band of the positive wavenumber half of the Fourier spectrum:

$$w\left(\frac{2\pi l}{N}\right) = 2\sum_{k=k_{min}}^{k_{max}} \hat{v}_k e^{2\pi i k/N} \tag{4}$$

Finally, the packet envelope is calculated as follows:

$$A(2\pi l/N) = |w(2\pi l/N)| \tag{5}$$

171 **3 Results**

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3.1 Stationary waves preceding heat waves in reanalysis

From the reanalysis data, we found 165 heat waves days from 26 events found in 173 the 75 summers from NCEP. We look at the temporal evolution of the planetary waves 174 preceding the events by constructing composites from day -20 to day 5 (Figure 1), where 175 we designate the first day of a heatwave as day 0. We used daily sub-seasonal anoma-176 lies at 300 hPa streamfunction calculated as the daily departures from the long-term mean 177 and from the seasonal (June-July-August) mean of the particular year. These anoma-178 lies were spatially filtered, considering only wavelengths between 2800 km and 10000 km 179 to consider only the wavenumbers of interest and avoid noise. Defined low and high-pressure 180 centers can be seen moving slowly westward, expected behavior for Rossby Wave prop-181 agation in the days leading up to the extreme event. Although several of the troughs and 182 ridges are clearly seen, it cannot be readily determined whether the wavenumber is 5 or 183 6. This difficulty may be explained by the scarcity of the events in the reanalysis, which 184 are rare by definition. On the other hand, the Rossby Waves amplitudes shown by the 185 envelope reflect the energy transport by the propagating waves. Notably, the large am-186 plitudes of the envelope are seen in the two ocean basins, consistent with the zones where 187 the Rossby Waves tend to reach large amplitudes and phase speeds. (Fragkoulidis & Wirth, 188 2020).189

3.2 Climatology on the idealized simulations

We first discuss the climatological circulation characteristics obtained for the 21630 191 analyzed days in each experiment (See Supporting Information Figures S1. and S2). As 192 we used a prescribed temperature from NCEP reanalysis data for the Newtonian relax-193 ation in the model, the bias-corrected experiments resemble the zonally symmetric and 194 asymmetric Northern Hemispheric summer in the real atmosphere by construction. It 195 means the temporal mean of the spatial distribution of the temperature, the barotropic 196 averaged zonal wind, and the vertical distribution of the zonal wind, in general, resem-197 ble the atmosphere. It should be clarified that although some features of the corrected 198 climatological state are not realistic, especially in the vertical distribution, they are re-199 alistic for the key latitudes of interest in this work (between 30°N and 60°N). Note that 200 the basic atmospheric state of the CTR differs spatially and in magnitude from the real 201 atmosphere. While in the real atmosphere, the peak jet intensity is about 10 m/s dur-202 ing the winter, the CTR exhibits a peak of more than 20 m/s. This is because the T_{eq} 203 in the initial configuration is set based on annual climatology values rather than sum-204 mer seasonality. 205

The expected phase speeds corresponding to the climatological zonal winds are cal-206 culated for different wavenumbers and compared to the results from reanalysis for the 207 boreal summer (Figure 2). This is to elucidate the expected stationarity for the wave 208 tracks preceding the events. The wavenumbers 5 and 6 have phase speeds close to zero 209 for both observations and the bias-corrected experiments ZOB and ROB. From this cal-210 culation, we expect that any pattern associated with any of these wavenumbers can be 211 developed before the events in the corrected experiments. However, the wavenumber five 212 tends to precede U.S. heat waves (Teng et al., 2013). Although in this result for NCEP, 213 it is wave number 6 that has phase velocity closest to 0, it should be noted that although 214 the Rossby wave dispersion relation delves into the large-scale dynamics of mid-latitudes, 215 it is based on a linearized approximation of the equations of motion, and may not ac-216 curately represent the behavior of atmospheric waves in situations where nonlinear in-217 teractions and other complicating factors are important. On the other hand, the overly 218 strong eddy phase speed obtained for the CTR run again reflect the unrealistic clima-219 to the boreal summer, which provides a non-favorable large-scale en-220 vironment for developing surface heatwaves. This explains the scarcity of heatwaves in 221 the CTR. 222

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3.3 Heat Waves driven by Intrinsic Planetary Wave pattern

Now that the model resembles the atmospheric basic state for the boreal summer, we can test our hypothesis that the basic state drives the stationary Rossby Waves propagation hence determining heat waves in the mid-latitudes. We must answer, firstly, if it resembles the actual statistics of heat waves, and secondly, if it resembles the atmospheric pattern preceding the events.

We compared the probability of heat waves days and the probability distribution 229 function (PDF) of the events between the experiments and NCEP (Figure 4). We found 230 111 events and 765 heat waves days in the bias-corrected symmetrical case and 65 and 231 464 heat waves days in the asymmetric one. Both bias-corrected experiments resemble 232 the PDF of the duration of the events in the reanalysis. However, only the one consid-233 ering asymmetries (ROB) has a probability of occurrence of heat waves days similar to 234 that found in the reanalysis. In contrast, the overly strong jetstream in the CTR run ex-235 periment can explain the extremely rare of similar heatwaves (only 5 events); the phase 236 speeds of troughs are too high to drive persistent weather conditions. Indeed, by chang-237 ing the detection methodology for this experiment, we can find slightly more events in 238 the CTR. For example, considering a lower velocity of the center of the warm points or 239 a minimum duration of the events shorter than five days, as shown in (Jiménez-Esteve 240

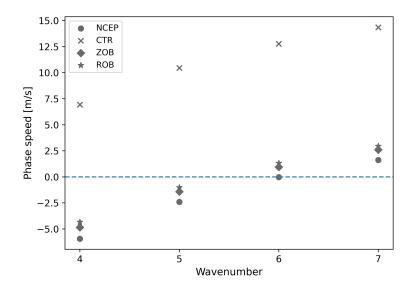


Figure 2. Phase speed for NCEP and the experiments in the Dry Core GCM. Calculated using the Rossby Wave dispersion relation (Eq. 7), considering the maximum barotropic averaged zonal wind in the northern mid-latitudes.

& Domeisen, 2022). However, the mere statistical difference of these five events found
with our methodology concerning the event statistics from NCEP and the bias-corrected
ones only confirms that an unrealistic basic state with respect to the boreal summer cannot provide a conducive environment for developing heatwave events that are statically
consistent with observations.

We construct Hovmöller diagrams for streamfunction anomalies and the Rossby 246 wave envelope for the bias-corrected experiments (Figure 3). In the ZOB and ROB, Rossby 247 wave propagation along the midlatitudes even 15 days ahead of the events can be ob-248 served. It can be seen clearly how the individual troughs and ridges repeatedly amplify 249 in almost the same longitudinal region, and the envelope shows the slow propagation of 250 the packets with an overall timescale of 20-30 days. In contrast, all five events in the CTR 251 are characterized by a much short timescale (3-7 days). Notice that when zonal varia-252 tions are included in the experiment ROB, this Hovmöller result has more realistic fea-253 tures, and most importantly, the pattern is characterized by a wavenumber 5, which co-254 incides with that associated with the occurrence of the events in a realistic atmosphere 255 (Teng et al., 2013).256

In addition, the latitudinal location of the jetstream changes for each experiment and it is known that the jetsetream serves as an efficient Rossby waveguide (Hoskins & Ambrizzi, 1993; Wirth et al., 2018), providing a preferred track for the Rossby Waves. The waveguide is shifted noticeably upward for the ZOB, while in the ROB, the latitudinal location closely resembles that observed in NCEP. As a consequence, the spatial location of the events is also more similar to that observed in the reanalysis.

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3.3.1 A back-of-the-envelop calculation of the phase speeds

According to the linear wave theory, features of any wave can be derived from its dispersion relation, which is obtained considering the equation of motion, the assumptions of two-dimensional barotropic flow and some other simplifications described in (Vallis,

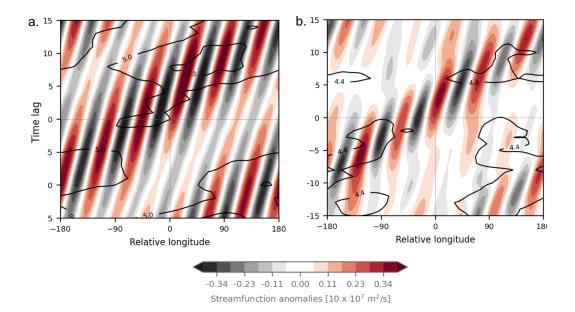


Figure 3. Hövmoller diagrams (longitude - time evolution) for streamfunction anomalies at 300 hPa between 30°N and 50°N with respect to all heatwaves detected (color shading) and Rossby Waves envelope (black contour) for the ZOB experiment (a.) and the ROB experiment (b.).

2017). Although these expression come from quite strong approximations, the solutions
turn out to be relevant to the real atmosphere, and provide insight into the large-scale
dynamics of mid-latitudes. The Rossby wave dispersion relation reads as follows:

$$\omega = Uk - \frac{\beta k}{k^2 + l^2} \tag{6}$$

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Then, the phase speed of Rossby waves in the x direction follows:

$$c_p^x = \frac{\omega}{k} = U - \frac{\beta}{k^2 + l^2} \tag{7}$$

Where ω is the constant wave angular frequency in units of radians/time, β the Coriolis parameter in β - plane approximation, $k^2 + l^2$ represent the square of the wavevector, where the wave numbers are given by $k = 2\pi/\lambda_x$, $l = 2\pi/\lambda_y$ with (λ_x, λ_y) the wavelengths in units of m. U is the density-weighted zonal-mean zonal-winds at the midlatitudes in m/s, for which the density of the air in each layer is calculated considering the proportionality described in 8, where z = -Hln(p/1000hPa).

$$\rho \propto e^{\frac{-z}{H}} \tag{8}$$

Since the phase speed relative to the background flow is always negative, these waves always propagate westward relative to the mean flow. When the phase lines are stationary relative to the ground, that is, $c_p^x = 0$ (Vallis, 2017), anomalous troughs lead to the occurrence and persistence of heat waves events.

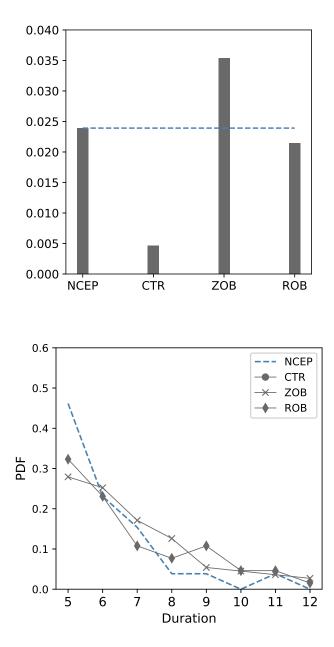


Figure 4. Comparison of statistics between NCEP and the experiments in the dry dynamical core GCM: Number of heatwaves waves over the total of analized days (upper) and the probability distribution function (PDF) of the duration of the events (bottom). For the CTR, the only 5 detected events correspond to a duration of 5 days.

²⁸¹ 4 Conclusions and discussions

The main objective of this study is to understand the role of the basic state on the 282 propagation of Rossby Wave packets driving heat waves in the U.S. We use an idealized 283 atmospheric model where physical parametrizations are substituted by a simple temper-284 ature relaxation. We implement the iterative procedure proposed by (Chang, 2006) to 285 simulate the climatological atmosphere for the boreal summer (June, July, August). At 286 the end, the model climate is forced to resemble the observed boreal summer mean flow 287 structure without explicitly considering complicated physical processes (e.g., moisture, 288 turbulence, clouds, etc.). This approach allows us to isolate the dry dynamics aspect of 289 atmospheric circulation on heatwaves evolution, and the conclusions from the results can 290 be summarized as follows: 291

292 293 • The climatological state drives the stationary Rossby wave propagation providing critical dynamical conditions for US heatwaves.

294 295 296 • A dry atmosphere with realistic boreal summer climatological state can produces a wavenumber-five Rossby Waves pattern, which often been seen in the observations preceding US heatwaves.

Previous studies show evidence of heat waves events preceded by a wavenumber 297 five structure contained in Rossby Waves, and the link between the amplification of this 298 pattern with its probability of heatwaves' occurrence. Our study helps to understand the 299 origin of this pattern. It is demonstrated that a dynamical model resembling the actual 300 zonal mean flow (ZOB) can clearly generate a similar Rossby Waves pattern, albeit char-301 acterized by zonal wavenumber six, preceding heat waves in the mid-latitudes. However, 302 including the zonal asymmetries in the mean flow (ROB) is crucial for developing the 303 specific wavenumber five structure, and consequently the statistics of the US heatwave 304 events that are consistent with those in reanalysis. 305

The idealized model enables us to isolate exclusively on the dry atmospheric dy-306 namics to understand the physical mechanism behind the interesting zonal wavenum-307 ber five pattern preceding frequent US heatwaves in the summer. While our idealized 308 model contains only dry dynamics, the specific structure of climatological state must arise 309 from the substantial contributions of various diabatic processes. On the one hand, lo-310 cal boundary conditions may play an important role on setting up a realistic climato-311 logical state. For example, (Donat et al., 2016; Lyon & Dole, 1995; McKinnon et al., 2016) 312 show that the combination of local anomalous sea surface temperature (SST) patterns 313 and atmospheric flow anomalies have contributed significantly to summer extremes. On 314 the other hand, other studies also show the contribution of soil moisture conditions to 315 heat waves by releasing surface diabatic heat (Seneviratne et al., 2010; Miralles et al., 316 2019) and even to the circumglobal circulation response of Rossby Waves (Douville & 317 Chauvin, 2000; Douville, 2002; Koster et al., 2016). During summer, local land and oceanic 318 conditions are important and may interact nonlinearly with atmospheric circulation states 319 preceding heatwave events. Our next step is to study the dynamics under which local 320 diabatic warming of soil moisture conditions impacts the amplitude of the wave num-321 ber five pattern and its associated impacts on the statistics of the air temperature ex-322 tremes in the US. 323

324 5 Open Research

The version of the Dry core GCM code used in this work is provided by the openaccess data from Wu and Reichler (2018). We construted the target climatology for our simulations from the NCEP-NCAR Reanalysis 1 data provided by the NOAA PSL, Boulder, Colorado, USA. All the NCEP datasets used for this work can be accessed from *https*: //downloads.psl.noaa.gov/Datasets/ncep.reanalysis/Dailies/. The temperature at all levels used to construct the inputs for the simulations, and also the velocity fields used for the comparison are in the folder called *pressure*/ in that link, and the temperature

at 2m in the folder *surface_gauss*/. The data also can be downloaded by following the

³³³ FTP link in the section Source & References in the NCEP-NCAR Reanalysis 1 webpage:

https: //psl.noaa.gov/data/gridded/data.ncep.reanalysis.html.

The codes for this work and the final datasets obtained from the three simulations are located in the open-access repository *https://github.com/castanev/Dynamics_Heat_Waves.git* linked to Zenodo *DOI*: 10.5281/zenodo.7844138

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342 References

343	Chang, E. K. M. (2006, July). An Idealized Nonlinear Model of the Northern Hemi-
344	sphere Winter Storm Tracks. Journal of the Atmospheric Sciences, $63(7)$,
345	1818-1839. Retrieved 2023-03-24, from https://journals.ametsoc.org/
346	view/journals/atsc/63/7/jas3726.1.xml (Publisher: American Me-
347	teorological Society Section: Journal of the Atmospheric Sciences) doi:
348	10.1175/JAS3726.1
349	Chen, P., & Newman, M. (1998, October). Rossby Wave Propagation and
350	the Rapid Development of Upper-Level Anomalous Anticyclones during
351	the 1988 U.S. Drought. Journal of Climate, $11(10)$, $2491-2504$. Re-
352	trieved 2023-04-07, from https://journals.ametsoc.org/view/journals/
353	clim/11/10/1520-0442_1998_011_2491_rwpatr_2.0.co_2.xml (Pub-
354	lisher: American Meteorological Society Section: Journal of Climate) doi:
355	10.1175/1520-0442(1998)011(2491:RWPATR)2.0.CO;2
356	Ding, Q., & Wang, B. (2005, September). Circumglobal Teleconnection in the
357	Northern Hemisphere Summer. Journal of Climate, 18(17), 3483–3505. Re-
358	trieved 2023-04-07, from https://journals.ametsoc.org/view/journals/
359	clim/18/17/jcli3473.1.xml (Publisher: American Meteorological Society
360	Section: Journal of Climate) doi: 10.1175/JCLI3473.1
361	Donat, M. G., King, A. D., Overpeck, J. T., Alexander, L. V., Durre, I., & Karoly,
362	D. J. (2016, January). Extraordinary heat during the 1930s US Dust Bowl
363	and associated large-scale conditions. Climate Dynamics, $46(1)$, $413-426$.
364	Retrieved 2023-04-07, from https://doi.org/10.1007/s00382-015-2590-5
365	doi: 10.1007/s00382-015-2590-5
366	Douville, H. (2002, April). Influence of Soil Moisture on the Asian and African
367	Monsoons. Part II: Interannual Variability. Journal of Climate, 15(7), 701–
368	720. Retrieved 2023-04-07, from https://journals.ametsoc.org/view/
369	journals/clim/15/7/1520-0442_2002_015_0701_iosmot_2.0.co_2.xml
370	(Publisher: American Meteorological Society Section: Journal of Climate)
371	doi: 10.1175/1520-0442(2002)015(0701:IOSMOT)2.0.CO;2
372	Douville, H., & Chauvin, F. (2000, October). Relevance of soil moisture for seasonal
373	climate predictions: a preliminary study. <i>Climate Dynamics</i> , 16(10), 719–736.
374	Retrieved 2023-04-07, from https://doi.org/10.1007/s003820000080 doi:
375	10.1007/s003820000080
376	Feldstein, S. B. (2000, December). The Timescale, Power Spectra, and Climate
377	Noise Properties of Teleconnection Patterns. Journal of Climate, 13(24),
378	4430-4440. Retrieved 2023-03-31, from https://journals.ametsoc.org/
379	view/journals/clim/13/24/1520-0442_2000_013_4430_ttpsac_2.0.co_2.xml
380	(Publisher: American Meteorological Society Section: Journal of Climate) doi: 10.1175/1520.0442(2000)013/4430:TTPSAC\2.0.CO:2
381	10.1175/1520-0442(2000)013(4430:TTPSAC)2.0.CO;2

382	Fragkoulidis, G., & Wirth, V. (2020, October). Local Rossby Wave Packet Ampli-
383	tude, Phase Speed, and Group Velocity: Seasonal Variability and Their Role
384	in Temperature Extremes. Journal of Climate, 33(20), 8767–8787. Retrieved
385	2023-04-07, from https://journals.ametsoc.org/view/journals/clim/33/
386	20/jcliD190377.xml (Publisher: American Meteorological Society Section:
387	Journal of Climate) doi: 10.1175/JCLI-D-19-0377.1
388	Held, I. M. (2005, November). The Gap between Simulation and Understanding in
389	Climate Modeling. Bulletin of the American Meteorological Society, $86(11)$,
390	1609–1614. Retrieved 2023-04-15, from https://journals.ametsoc.org/
391	view/journals/bams/86/11/bams-86-11-1609.xml (Publisher: Ameri-
392	can Meteorological Society Section: Bulletin of the American Meteorological
393	Society) doi: 10.1175/BAMS-86-11-1609
394	Held, I. M., & Suarez, M. J. (1994, October). A Proposal for the Intercomparison
395	of the Dynamical Cores of Atmospheric General Circulation Models. Bulletin
396	of the American Meteorological Society, 75(10), 1825–1830. Retrieved 2023-
397	03-24, from https://journals.ametsoc.org/view/journals/bams/75/10/
398	1520-0477_1994_075_1825_apftio_2_0_co_2.xml (Publisher: American Me- teorological Society Section: Publishin of the American Meteorological Society)
399 400	teorological Society Section: Bulletin of the American Meteorological Society) doi: 10.1175/1520-0477(1994)075(1825:APFTIO)2.0.CO;2
401	Hoskins, B. J., & Ambrizzi, T. (1993, June). Rossby Wave Propagation on a Realis-
402	tic Longitudinally Varying Flow. Journal of the Atmospheric Sciences, 50(12),
403	1661-1671. Retrieved 2023-04-07, from https://journals.ametsoc.org/
404	view/journals/atsc/50/12/1520-0469_1993_050_1661_rwpoar_2_0_co_2.xml
405	(Publisher: American Meteorological Society Section: Journal of the Atmo-
406	spheric Sciences) doi: $10.1175/1520-0469(1993)050(1661:RWPOAR)2.0.CO;2$
407	Jiménez-Esteve, B., & Domeisen, D. I. (2022, July). The role of atmospheric dynam-
408	ics and large-scale topography in driving heatwaves. Quarterly Journal of the
409	Royal Meteorological Society, 148(746), 2344–2367. Retrieved 2023-04-07, from
410	https://onlinelibrary.wiley.com/doi/10.1002/qj.4306 doi: $10.1002/qj$
411	.4306
412	Koster, R. D., Chang, Y., Wang, H., & Schubert, S. D. (2016, October). Impacts of
413	Local Soil Moisture Anomalies on the Atmospheric Circulation and on Remote
414	Surface Meteorological Fields during Boreal Summer: A Comprehensive Anal-
415	ysis over North America. Journal of Climate, 29(20), 7345–7364. Retrieved
416	2023-04-07, from https://journals.ametsoc.org/view/journals/clim/
417	29/20/jcli-d-16-0192.1.xml (Publisher: American Meteorological Society
418	Section: Journal of Climate) doi: 10.1175/JCLI-D-16-0192.1
419	Lyon, B., & Dole, R. M. (1995, June). A Diagnostic Comparison of the 1980 and
420	1988 U.S. Summer Heat Wave-Droughts. Journal of Climate, 8(6), 1658–
421	1675. Retrieved 2023-04-07, from https://journals.ametsoc.org/view/
422	journals/clim/8/6/1520-0442_1995_008_1658_adcota_2_0_co_2.xml (Pub- lisher: American Meteorological Society Section: Journal of Climate) doi:
423	10.1175/1520-0442(1995)008(1658:ADCOTA)2.0.CO;2
424	McKinnon, K. A., Rhines, A., Tingley, M. P., & Huybers, P. (2016, May). Long-
425 426	lead predictions of eastern United States hot days from Pacific sea surface
420	temperatures. Nature Geoscience, $9(5)$, 389–394. Retrieved 2023-04-07, from
428	https://www.nature.com/articles/ngeo2687 (Number: 5 Publisher: Na-
429	ture Publishing Group) doi: 10.1038/ngeo2687
430	Meehl, G. A., & Tebaldi, C. (2004, August). More Intense, More Frequent, and
431	Longer Lasting Heat Waves in the 21st Century. Science, 305 (5686), 994–
432	997. Retrieved 2023-04-07, from https://www.science.org/doi/10.1126/
433	science.1098704 (Publisher: American Association for the Advancement of
434	Science) doi: 10.1126/science.1098704
435	Miralles, D. G., Gentine, P., Seneviratne, S. I., & Teuling, A. J. (2019, January).
436	Land–atmospheric feedbacks during droughts and heatwaves: state of the sci-

437	ence and current challenges. Annals of the New York Academy of Sciences,
438	1436(1), 19-35. Retrieved 2023-04-07, from https://www.ncbi.nlm.nih.gov/
439	pmc/articles/PMC6378599/ doi: 10.1111/nyas.13912
440	Pedlosky, J. (2003). Waves in the Ocean and Atmosphere: Introduction to Wave Dy-
441	namics (2003rd edition ed.). Berlin; New York: Springer.
442	Röthlisberger, M., Frossard, L., Bosart, L. F., Keyser, D., & Martius, O. (2019,
443	June). Recurrent Synoptic-Scale Rossby Wave Patterns and Their Effect on
444	the Persistence of Cold and Hot Spells. Journal of Climate, 32(11), 3207–3226.
445	Retrieved 2023-04-07, from https://journals.ametsoc.org/view/journals/
446	clim/32/11/jcli-d-18-0664.1.xml doi: 10.1175/JCLI-D-18-0664.1
447	Schubert, S., Wang, H., & Suarez, M. (2011, September). Warm Season Subseasonal
448	Variability and Climate Extremes in the Northern Hemisphere: The Role of
449	Stationary Rossby Waves. Journal of Climate, 24(18), 4773–4792. Retrieved
450	2023-04-15, from https://journals.ametsoc.org/view/journals/clim/24/
451	18/jcli-d-10-05035.1.xml (Publisher: American Meteorological Society
452	Section: Journal of Climate) doi: 10.1175/JCLI-D-10-05035.1
453	Seneviratne, S. I., Corti, T., Davin, E. L., Hirschi, M., Jaeger, E. B., Lehner, I.,
454	Teuling, A. J. (2010, May). Investigating soil moisture-climate interactions
455	in a changing climate: A review. $Earth-Science Reviews, 99(3), 125-161.$
456	Retrieved 2023-04-15, from https://www.sciencedirect.com/science/
457	article/pii/S0012825210000139 doi: 10.1016/j.earscirev.2010.02.004
458	Teng, H., Branstator, G., Wang, H., Meehl, G. A., & Washington, W. M. (2013,
459	December). Probability of US heat waves affected by a subseasonal planetary
460	wave pattern. Nature Geoscience, $6(12)$, 1056–1061. Retrieved 2023-03-24,
461	from https://www.nature.com/articles/ngeo1988 (Number: 12 Publisher:
462	Nature Publishing Group) doi: 10.1038/ngeo1988
463	Vallis, G. K. (2017). Atmospheric and Oceanic Fluid Dynamics: Funda-
464	mentals and Large-Scale Circulation (2nd ed.). Cambridge: Cambridge
465	University Press. Retrieved 2023-03-01, from https://www.cambridge
466	.org/core/books/atmospheric-and-oceanic-fluid-dynamics/ 41379BDDC4257CBE11143C466F6428A4 doi: 10.1017/9781107588417
467	Wirth, V., Riemer, M., Chang, E. K. M., & Martius, O. (2018, July). Rossby Wave
468	Packets on the Midlatitude Waveguide—A Review. Monthly Weather Review,
469	146(7), 1965–2001. Retrieved 2023-04-05, from https://journals.ametsoc
470 471	.org/view/journals/mwre/146/7/mwr-d-16-0483.1.xml (Publisher: Amer-
471	ican Meteorological Society Section: Monthly Weather Review) doi: 10.1175/
473	MWR-D-16-0483.1
474	Wu, Z., & Reichler, T. (2018). Towards a more earth-like circulation in idealized
475	models. Retrieved from https://doi.org/ doi: 10.1029/2018MS001356
476	Zimin, A. V., Szunyogh, I., Patil, D. J., Hunt, B. R., & Ott, E. (2003, May). Ex-
477	tracting Envelopes of Rossby Wave Packets. Monthly Weather Review, 131(5),
478	1011-1017. Retrieved 2023-03-24, from https://journals.ametsoc.org/
479	view/journals/mwre/131/5/1520-0493_2003_131_1011_eeorwp_2.0.co_2.xml
480	(Publisher: American Meteorological Society Section: Monthly Weather Re-
481	view) doi: $10.1175/1520-0493(2003)131(1011:EEORWP)2.0.CO;2$

The Role of Climatological State on Driving US Heat Waves Through Rossby Waves Packets

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5	Key Points:
6	• The climatological state drives the stationary Rossby wave propagation provid-
7	ing critical dynamical conditions for heat waves in the US.
8	• A dry atmospheric model with a corrected climatological state generates heatwaves
9	that are statistically consistent with observations.
10	• The slowly propagating Rossby wave packets with a timescale of 20-30 days can
11	be a source of intraseasonal predictability.

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12 Abstract

While heat waves are local extreme weather events, a planetary Rossby wave pattern is 13 statistically related to the occurrence of heat waves events in the U.S. However, whether 14 such planetary wave patterns cause the enhanced statistics of local heat waves or as a 15 coincidence is debatable. In this work, we hypothesize that the atmospheric climatolog-16 ical state dictates the slowly propagating wave pattern, which sets up a conducive large-17 scale environment for local US heat waves. We implement an idealized dry dynamic core 18 model with an iterative approach to simulate the realistic North American summer cli-19 matological state. As the model can generate similar large-scale planetary wave patterns 20 propagating throughout North America, significantly more heatwaves are generated, and 21 the statistics of heat waves become consistent with that estimated in reanalysis prod-22 ucts. The slowly propagating Rossby wave packets with a timescale of 20-30 days can 23 serve as a new source of intraseasonal predictability. 24

²⁵ Plain Language Summary

Heatwaves are the leading weather-related killer in the United States and affect mainly 26 the most vulnerable communities. These extreme events are statistically related to a large-27 scale wave pattern of Rossby waves that features a zonal wavenumber five structure. To 28 find out what controls this wave pattern, we use a simple general circulation model that 29 only contains dry dynamics without complicated interactions with moisture or clouds. 30 31 We found that by modifying the climatological state of temperature and velocity fields based on the observed structure from the Northern Hemisphere summer, we can observe 32 the same wave number five pattern developing days before the heatwave events, resem-33 bling that from the observations. This result suggests that the climatological state of the 34 Northern Hemispheric summer provides a conducive environment for heatwaves in the 35 U.S. due to the slowed-down propagation speed of the Rossby wave packets. A deeper 36 understanding of its dynamics is crucial because, as this pattern develops up to 20 days 37 ahead of the extremes, the underlying physical process governing this pattern may serve 38 as a source of predictability on the Subseasonal-to-Seasonal (S2S) timescale, a current 39 gap of forecasts between weather and climate. 40

41 **1** Introduction

42 Over the coming century, climate change is expected to increase average summer 43 temperatures and the severity of extreme heat linked with heat waves events. In the United 44 States, the frequency, intensity, and duration of heat waves have been increasing rapidly 45 in recent decades, and this behavior is projected to continue in the next decades (Meehl 46 & Tebaldi, 2004). Currently, forecasters in the United States can only predict extreme 47 events up to 10 days in advance because, unlike in the tropics, circulation in the mid-48 latitudes is more chaotic as it is dominated by climatic noise (Feldstein, 2000).

However, it is intriguing to hypothesize that certain atmospheric circulation states 49 can be substantially more predictable than the average scenario, because these circula-50 tion regimes are associated with low-frequency patterns (Schubert et al., 2011). With 51 this in mind, and considering that different studies have suggested that propagating sta-52 tionary Rossby Waves play an important role in the mid-latitude atmospheric variabil-53 ity, the scientific community recently has increased interest in the connection between 54 extreme weather events and large-scale atmospheric patterns such as Rossby Wave Pack-55 ets (Chen & Newman, 1998; Schubert et al., 2011; Fragkoulidis & Wirth, 2020). 56

In many of the investigated cases, the extreme weather was linked to an upper-tropospheric
trough (i.e., a breaking Rossby Waves). For example, Chen and Newman (1998) suggests
that Rossby Waves originating in the west Pacific were the key in initiating intense anomalous anticyclones during the 1988 U.S. drought. Also, Schubert et al. (2011) relates this

large-scale pattern with monthly mean precipitation and surface temperature variabil-61 ity over many regions of the extratropical land areas, including the northern U.S., parts 62 of Canada, Europe, and Russia. Following the same approach, Ding and Wang (2005) 63 found an interannually varying Northern Hemisphere circumglobal pattern with a preferred wavenumber five structure. This particular structure is confined within the waveg-65 uide associated with the summer north jet stream in the stationary state and is linked 66 to significant surface air temperature and rainfall anomalies in western Europe, Euro-67 pean Russia, India, East Asia, and North America. In a more recent study Teng et al. 68 (2013), more concrete evidence has been shown of the relationship between the occur-69 rence of heatwayes in the U.S. and the same wavenumber five pattern developing as early 70 as 20 days ahead of the events. Despite the growing evidence for the concurrences of this 71 interesting wave number five pattern, the physical mechanism that drives the atmospheric 72 pattern and thus determines heatwaves has yet to be understood. 73

This study aims to understand better the fundamental role of atmospheric dynam-74 ics in the evolution of U.S. heat waves based on the hypothesis that the climatological 75 state for the Northern Hemisphere summer drives the stationary Rossby wave propaga-76 tion providing critical dynamical conditions for heat waves in the U.S. Inspired by the 77 motivation to look for the simplest possible model in the climate model hierarchies (Held, 78 2005), we adopt an idealized dry dynamical core that allows for a bias correction of the 79 mean flow structure without physical parameterizations. This approach allows us to iso-80 late the role of intrinsic planetary waves in the evolution of extreme events. 81

⁸² 2 Data and Methods

2.1 Reanalysis products

For the observational analysis, the NCEP/NCAR reanalysis of the National Oceanic and Atmospheric Administration (NOAA) was implemented. The surface air temperature (SAT), all levels of temperature and the meridional and zonal components of the wind at 300 hPa (V_{300} and U_{300}) were obtained from these databases with a spatial resolution of 2.5° and daily temporal resolution from 1948 to 2023.

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2.2 An Idealized GCM with realistic climatological basic state

Idealized general circulation models are commonly used for the study of atmospheric dynamics. A dry dynamical core model solves the primitive equations on the sphere by nudging the climatological temperature field toward a prescribed structure of radiative equilibrium temperature (T_{eq}) . Essentially, this process isolates the dry atmospheric dynamics from the complex physical parameterizations.

⁹⁵ We use the open-access updated version from (Wu & Reichler, 2018) of the spec-⁹⁶ tral dynamical core model proposed by (Held & Suarez, 1994) for the Geophysical Fluid ⁹⁷ Dynamics Laboratory (GFDL). The model has a horizontal resolution of T42 (64x128 ⁹⁸ grid) and 40 vertical σ levels between the surface and 0.01 hPa. To represent the boundary-⁹⁹ layer friction, Rayleigh drag is used to remove momentum in the lower troposphere be-¹⁰⁰ tween the surface and $\sigma = 0.7$. By default configurations of this model, the tempera-¹⁰¹ ture is forced by Newtonian relaxation toward a prescribed equilibrium temperature as ¹⁰² follows:

$$\frac{\partial T}{\partial t} = \frac{T - T_{eq}}{\tau} \tag{1}$$

where τ is the prescribed relaxation timescale.

Since the primary purpose of this study is to examine how the atmospheric basic state controls the large-scale Rossby Wave pattern with a zonal wavenumber five struc-

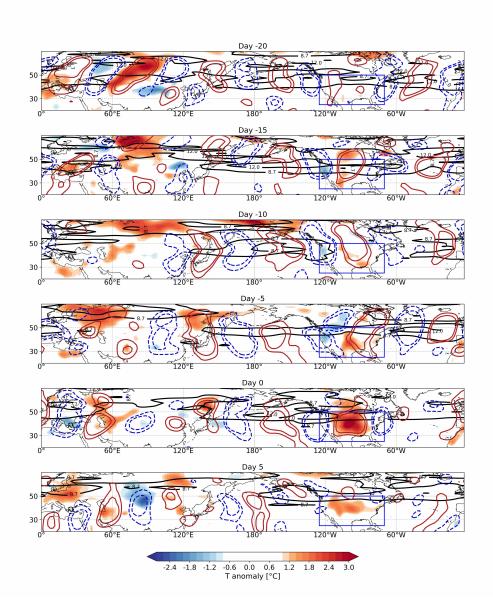


Figure 1. Evolution of heat waves 20 days ahead of the events: Red and dashed blue contours represent composites of 200hPa streamfunction anomalies at \pm 0.05, 0.08 (red as positives). Shading represents surface air temperature anomaly and black contours represent the Rossby Wave Packet envelope at levels 8.7, 12 m/s.

ture driving local heat waves in the US, the model includes the iterative procedure pro-106 posed by Chang (2006) to simulate the climatological basic state of the atmosphere in 107 the idealized model. This consists of iterating the radiative equilibrium temperature pro-108 file so that at the end of the iterations, the model climate closely resembles the desired 109 target climate (see Supporting Information for more details of the methodology). The 110 iteration uses a fixed equilibrium temperature for each N step (T_{eq}) in a run of 62 years, 111 for a total of 21630 days after eliminating the first 1000 days of simulation. Then, we 112 calculate the model simulated temperature climatology $T_{(N)}$ and correct it concerning 113 basic state from the NCEP reanalysis data T_R . The next iteration step N+1 is calcu-114 lated according to: 115

$$T_{eq(N+1)} = T_{eq(N)} - \frac{2}{3}(T_{(N)} - T_R), N = 1, 2, 3...$$
(2)

To assess the role of the climatological basic state, we carry out three experiments as follows:

- Held and Suarez 1994 configuration (CTR): the zonal-mean temperature 118 structure is relaxed toward a prescribed equinoctial radiative-equilibrium state. 119 It's the default configuration described in detail in (Held & Suarez, 1994) and leads 120 to a climatological state close to the observed annual mean climatological state. 121 **Zonally-symmetric boreal summer (ZOB)**: The ratiative forcing T_R described 122 in section 2.2 and in eq. 2 is set by the two-dimensional (i.e., latitude-pressure) 123 basic state from the NCEP reanalysis data. For this experiment, we use the zon-124 ally symmetrical distribution of temperature for the climatological boreal sum-125 mer (June, July, and August) between 1948 and 2021. 126 Realistic Northern Hemisphere summer (ROB): The full three-dimensional 127
- 128 structure of the radiative forcing T_R is obtained from the climatological boreal sum-129 mer from NCEP, especially including zonal asymmetries. This bias correction in-130 cludes zonal variations in T_{eq} , which can be related to zonal variations in diabatic 131 heating as discussed in Chang (2006) work.

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2.3 Heat waves identification

A heatwave event is set as at least five consecutive days with more than around 900000 km^2 over the US continental area (125W - 70W, 25N - 50N) with daily averaged SAT exceeding a threshold value, and the center of these warm points, considered as the point of maximum temperature, does not move faster than 5 latitude or longitude per day. The temperature threshold varies spatially as well as with the day of the year. It was established as the percentile 97.5 of the historical SAT. This approach follows Teng et al. (2013).

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2.4 Rossby Wave Packets evolution

As this study aims to understand the dynamics behind heat waves, it is necessary 141 to look at the evolution of the Rossby Waves Packets driving extreme events from ob-142 servational data and modeling outputs. Both the eddy phase speeds and their group ve-143 locity are important physical properties of Rossby Wave Packets' propagation. In par-144 ticular, the eddy phase speed indicates the propagation speed of individual troughs within 145 Rossby Waves Packets. It can thus be critical for the persistence of extreme surface weather 146 (Röthlisberger et al., 2019). On the other hand, the group velocity reflects how quickly 147 the whole packet propagates locally and represents the rate at which the Rossby waves 148 transfer energy horizontally (Pedlosky, 2003). Based on the Rossby Waves Packets' pro-149 gression in longitude and time, we compute the envelope of the meridional wind indi-150 cating their preferred regions of formation and decay. Specifically, we estimate the Rossby 151

Waves Packets' propagation by constructing Hovmöller diagrams using streamfunction
 anomalies at 300 hPa and the corresponding Rossby Waves Packets' envelope.

Subsequently, we calculate the local (in space and time) group velocity using the 154 Rossby Waves Packets' envelope, which reflects how the flow features of enhanced merid-155 ional wind amplitude propagate in the zonal and meridional directions (Fragkoulidis & 156 Wirth, 2020). We implemented the method proposed by Zimin et al. (2003) for the cal-157 culation of Rossby Waves Packets' envelope involving the Hilbert transform along cir-158 cles of constant latitude combined with a restriction of the zonal wavenumber to a spec-159 ified interval, which in this case corresponds to wavenumbers between 3 and 11. We im-160 plement this methodology using the meridional wind, which is usually implemented for 161 the diagnosis of meridional deviations from the zonal flow, and is particularly well suited 162 for the detection of Rossby waves. Specifically, we take the upper-tropospheric (300 hPa) 163 meridional wind which features strong Rossby Waves Packet's activities (Wirth et al., 164 2018). The meridional wind anomaly v(x) is considered on an equidistant grid along a 165 latitude circle, where $x = 2\pi/N$ with $0 < x \le 2\pi$, N is an even integer, and l = 1, 2, ..., N. 166 The Fourier transform of the real function v(x) is computed as: 167

$$\hat{v}_k = \frac{1}{N} \sum_{l=1}^N v\left(\frac{2\pi l}{N}\right) e^{-2\pi i k/N}, \left(K = -\frac{N}{2} + 1, \dots, \frac{N}{2}\right)$$
(3)

The inverse Fourier transform is applied to a selected band of the positive wavenumber half of the Fourier spectrum:

$$w\left(\frac{2\pi l}{N}\right) = 2\sum_{k=k_{min}}^{k_{max}} \hat{v}_k e^{2\pi i k/N} \tag{4}$$

Finally, the packet envelope is calculated as follows:

$$A(2\pi l/N) = |w(2\pi l/N)| \tag{5}$$

171 **3 Results**

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3.1 Stationary waves preceding heat waves in reanalysis

From the reanalysis data, we found 165 heat waves days from 26 events found in 173 the 75 summers from NCEP. We look at the temporal evolution of the planetary waves 174 preceding the events by constructing composites from day -20 to day 5 (Figure 1), where 175 we designate the first day of a heatwave as day 0. We used daily sub-seasonal anoma-176 lies at 300 hPa streamfunction calculated as the daily departures from the long-term mean 177 and from the seasonal (June-July-August) mean of the particular year. These anoma-178 lies were spatially filtered, considering only wavelengths between 2800 km and 10000 km 179 to consider only the wavenumbers of interest and avoid noise. Defined low and high-pressure 180 centers can be seen moving slowly westward, expected behavior for Rossby Wave prop-181 agation in the days leading up to the extreme event. Although several of the troughs and 182 ridges are clearly seen, it cannot be readily determined whether the wavenumber is 5 or 183 6. This difficulty may be explained by the scarcity of the events in the reanalysis, which 184 are rare by definition. On the other hand, the Rossby Waves amplitudes shown by the 185 envelope reflect the energy transport by the propagating waves. Notably, the large am-186 plitudes of the envelope are seen in the two ocean basins, consistent with the zones where 187 the Rossby Waves tend to reach large amplitudes and phase speeds. (Fragkoulidis & Wirth, 188 2020).189

3.2 Climatology on the idealized simulations

We first discuss the climatological circulation characteristics obtained for the 21630 191 analyzed days in each experiment (See Supporting Information Figures S1. and S2). As 192 we used a prescribed temperature from NCEP reanalysis data for the Newtonian relax-193 ation in the model, the bias-corrected experiments resemble the zonally symmetric and 194 asymmetric Northern Hemispheric summer in the real atmosphere by construction. It 195 means the temporal mean of the spatial distribution of the temperature, the barotropic 196 averaged zonal wind, and the vertical distribution of the zonal wind, in general, resem-197 ble the atmosphere. It should be clarified that although some features of the corrected 198 climatological state are not realistic, especially in the vertical distribution, they are re-199 alistic for the key latitudes of interest in this work (between 30°N and 60°N). Note that 200 the basic atmospheric state of the CTR differs spatially and in magnitude from the real 201 atmosphere. While in the real atmosphere, the peak jet intensity is about 10 m/s dur-202 ing the winter, the CTR exhibits a peak of more than 20 m/s. This is because the T_{eq} 203 in the initial configuration is set based on annual climatology values rather than sum-204 mer seasonality. 205

The expected phase speeds corresponding to the climatological zonal winds are cal-206 culated for different wavenumbers and compared to the results from reanalysis for the 207 boreal summer (Figure 2). This is to elucidate the expected stationarity for the wave 208 tracks preceding the events. The wavenumbers 5 and 6 have phase speeds close to zero 209 for both observations and the bias-corrected experiments ZOB and ROB. From this cal-210 culation, we expect that any pattern associated with any of these wavenumbers can be 211 developed before the events in the corrected experiments. However, the wavenumber five 212 tends to precede U.S. heat waves (Teng et al., 2013). Although in this result for NCEP, 213 it is wave number 6 that has phase velocity closest to 0, it should be noted that although 214 the Rossby wave dispersion relation delves into the large-scale dynamics of mid-latitudes, 215 it is based on a linearized approximation of the equations of motion, and may not ac-216 curately represent the behavior of atmospheric waves in situations where nonlinear in-217 teractions and other complicating factors are important. On the other hand, the overly 218 strong eddy phase speed obtained for the CTR run again reflect the unrealistic clima-219 to the boreal summer, which provides a non-favorable large-scale en-220 vironment for developing surface heatwaves. This explains the scarcity of heatwaves in 221 the CTR. 222

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3.3 Heat Waves driven by Intrinsic Planetary Wave pattern

Now that the model resembles the atmospheric basic state for the boreal summer, we can test our hypothesis that the basic state drives the stationary Rossby Waves propagation hence determining heat waves in the mid-latitudes. We must answer, firstly, if it resembles the actual statistics of heat waves, and secondly, if it resembles the atmospheric pattern preceding the events.

We compared the probability of heat waves days and the probability distribution 229 function (PDF) of the events between the experiments and NCEP (Figure 4). We found 230 111 events and 765 heat waves days in the bias-corrected symmetrical case and 65 and 231 464 heat waves days in the asymmetric one. Both bias-corrected experiments resemble 232 the PDF of the duration of the events in the reanalysis. However, only the one consid-233 ering asymmetries (ROB) has a probability of occurrence of heat waves days similar to 234 that found in the reanalysis. In contrast, the overly strong jetstream in the CTR run ex-235 periment can explain the extremely rare of similar heatwaves (only 5 events); the phase 236 speeds of troughs are too high to drive persistent weather conditions. Indeed, by chang-237 ing the detection methodology for this experiment, we can find slightly more events in 238 the CTR. For example, considering a lower velocity of the center of the warm points or 239 a minimum duration of the events shorter than five days, as shown in (Jiménez-Esteve 240

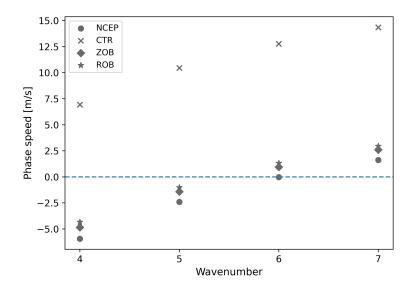


Figure 2. Phase speed for NCEP and the experiments in the Dry Core GCM. Calculated using the Rossby Wave dispersion relation (Eq. 7), considering the maximum barotropic averaged zonal wind in the northern mid-latitudes.

& Domeisen, 2022). However, the mere statistical difference of these five events found
with our methodology concerning the event statistics from NCEP and the bias-corrected
ones only confirms that an unrealistic basic state with respect to the boreal summer cannot provide a conducive environment for developing heatwave events that are statically
consistent with observations.

We construct Hovmöller diagrams for streamfunction anomalies and the Rossby 246 wave envelope for the bias-corrected experiments (Figure 3). In the ZOB and ROB, Rossby 247 wave propagation along the midlatitudes even 15 days ahead of the events can be ob-248 served. It can be seen clearly how the individual troughs and ridges repeatedly amplify 249 in almost the same longitudinal region, and the envelope shows the slow propagation of 250 the packets with an overall timescale of 20-30 days. In contrast, all five events in the CTR 251 are characterized by a much short timescale (3-7 days). Notice that when zonal varia-252 tions are included in the experiment ROB, this Hovmöller result has more realistic fea-253 tures, and most importantly, the pattern is characterized by a wavenumber 5, which co-254 incides with that associated with the occurrence of the events in a realistic atmosphere 255 (Teng et al., 2013).256

In addition, the latitudinal location of the jetstream changes for each experiment and it is known that the jetsetream serves as an efficient Rossby waveguide (Hoskins & Ambrizzi, 1993; Wirth et al., 2018), providing a preferred track for the Rossby Waves. The waveguide is shifted noticeably upward for the ZOB, while in the ROB, the latitudinal location closely resembles that observed in NCEP. As a consequence, the spatial location of the events is also more similar to that observed in the reanalysis.

263

3.3.1 A back-of-the-envelop calculation of the phase speeds

According to the linear wave theory, features of any wave can be derived from its dispersion relation, which is obtained considering the equation of motion, the assumptions of two-dimensional barotropic flow and some other simplifications described in (Vallis,

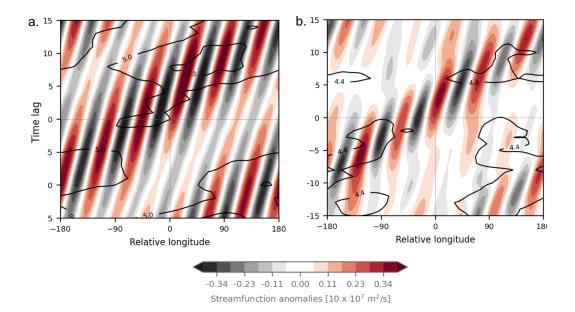


Figure 3. Hövmoller diagrams (longitude - time evolution) for streamfunction anomalies at 300 hPa between 30°N and 50°N with respect to all heatwaves detected (color shading) and Rossby Waves envelope (black contour) for the ZOB experiment (a.) and the ROB experiment (b.).

2017). Although these expression come from quite strong approximations, the solutions
turn out to be relevant to the real atmosphere, and provide insight into the large-scale
dynamics of mid-latitudes. The Rossby wave dispersion relation reads as follows:

$$\omega = Uk - \frac{\beta k}{k^2 + l^2} \tag{6}$$

270

Then, the phase speed of Rossby waves in the x direction follows:

$$c_p^x = \frac{\omega}{k} = U - \frac{\beta}{k^2 + l^2} \tag{7}$$

Where ω is the constant wave angular frequency in units of radians/time, β the Coriolis parameter in β - plane approximation, $k^2 + l^2$ represent the square of the wavevector, where the wave numbers are given by $k = 2\pi/\lambda_x$, $l = 2\pi/\lambda_y$ with (λ_x, λ_y) the wavelengths in units of m. U is the density-weighted zonal-mean zonal-winds at the midlatitudes in m/s, for which the density of the air in each layer is calculated considering the proportionality described in 8, where z = -Hln(p/1000hPa).

$$\rho \propto e^{\frac{-z}{H}} \tag{8}$$

Since the phase speed relative to the background flow is always negative, these waves always propagate westward relative to the mean flow. When the phase lines are stationary relative to the ground, that is, $c_p^x = 0$ (Vallis, 2017), anomalous troughs lead to the occurrence and persistence of heat waves events.

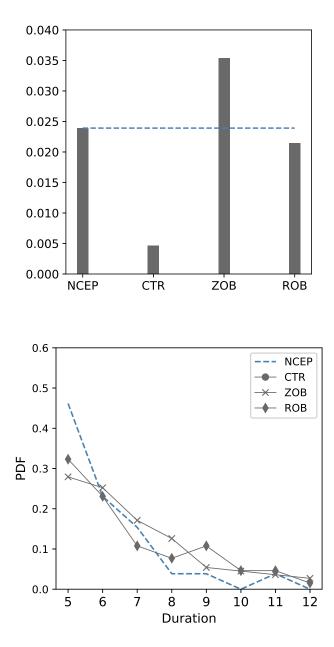


Figure 4. Comparison of statistics between NCEP and the experiments in the dry dynamical core GCM: Number of heatwaves waves over the total of analized days (upper) and the probability distribution function (PDF) of the duration of the events (bottom). For the CTR, the only 5 detected events correspond to a duration of 5 days.

²⁸¹ 4 Conclusions and discussions

The main objective of this study is to understand the role of the basic state on the 282 propagation of Rossby Wave packets driving heat waves in the U.S. We use an idealized 283 atmospheric model where physical parametrizations are substituted by a simple temper-284 ature relaxation. We implement the iterative procedure proposed by (Chang, 2006) to 285 simulate the climatological atmosphere for the boreal summer (June, July, August). At 286 the end, the model climate is forced to resemble the observed boreal summer mean flow 287 structure without explicitly considering complicated physical processes (e.g., moisture, 288 turbulence, clouds, etc.). This approach allows us to isolate the dry dynamics aspect of 289 atmospheric circulation on heatwaves evolution, and the conclusions from the results can 290 be summarized as follows: 291

292 293 • The climatological state drives the stationary Rossby wave propagation providing critical dynamical conditions for US heatwaves.

294 295 296 • A dry atmosphere with realistic boreal summer climatological state can produces a wavenumber-five Rossby Waves pattern, which often been seen in the observations preceding US heatwaves.

Previous studies show evidence of heat waves events preceded by a wavenumber 297 five structure contained in Rossby Waves, and the link between the amplification of this 298 pattern with its probability of heatwaves' occurrence. Our study helps to understand the 299 origin of this pattern. It is demonstrated that a dynamical model resembling the actual 300 zonal mean flow (ZOB) can clearly generate a similar Rossby Waves pattern, albeit char-301 acterized by zonal wavenumber six, preceding heat waves in the mid-latitudes. However, 302 including the zonal asymmetries in the mean flow (ROB) is crucial for developing the 303 specific wavenumber five structure, and consequently the statistics of the US heatwave 304 events that are consistent with those in reanalysis. 305

The idealized model enables us to isolate exclusively on the dry atmospheric dy-306 namics to understand the physical mechanism behind the interesting zonal wavenum-307 ber five pattern preceding frequent US heatwaves in the summer. While our idealized 308 model contains only dry dynamics, the specific structure of climatological state must arise 309 from the substantial contributions of various diabatic processes. On the one hand, lo-310 cal boundary conditions may play an important role on setting up a realistic climato-311 logical state. For example, (Donat et al., 2016; Lyon & Dole, 1995; McKinnon et al., 2016) 312 show that the combination of local anomalous sea surface temperature (SST) patterns 313 and atmospheric flow anomalies have contributed significantly to summer extremes. On 314 the other hand, other studies also show the contribution of soil moisture conditions to 315 heat waves by releasing surface diabatic heat (Seneviratne et al., 2010; Miralles et al., 316 2019) and even to the circumglobal circulation response of Rossby Waves (Douville & 317 Chauvin, 2000; Douville, 2002; Koster et al., 2016). During summer, local land and oceanic 318 conditions are important and may interact nonlinearly with atmospheric circulation states 319 preceding heatwave events. Our next step is to study the dynamics under which local 320 diabatic warming of soil moisture conditions impacts the amplitude of the wave num-321 ber five pattern and its associated impacts on the statistics of the air temperature ex-322 tremes in the US. 323

³²⁴ 5 Open Research

The version of the Dry core GCM code used in this work is provided by the openaccess data from Wu and Reichler (2018). We construted the target climatology for our simulations from the NCEP-NCAR Reanalysis 1 data provided by the NOAA PSL, Boulder, Colorado, USA. All the NCEP datasets used for this work can be accessed from *https*: //downloads.psl.noaa.gov/Datasets/ncep.reanalysis/Dailies/. The temperature at all levels used to construct the inputs for the simulations, and also the velocity fields used for the comparison are in the folder called *pressure*/ in that link, and the temperature

at 2m in the folder *surface_gauss*/. The data also can be downloaded by following the

³³³ FTP link in the section Source & References in the NCEP-NCAR Reanalysis 1 webpage:

https: //psl.noaa.gov/data/gridded/data.ncep.reanalysis.html.

The codes for this work and the final datasets obtained from the three simulations are located in the open-access repository *https://github.com/castanev/Dynamics_Heat_Waves.git* linked to Zenodo *DOI*: 10.5281/zenodo.7844138

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342 References

343	Chang, E. K. M. (2006, July). An Idealized Nonlinear Model of the Northern Hemi-
344	sphere Winter Storm Tracks. Journal of the Atmospheric Sciences, 63(7),
345	1818-1839. Retrieved 2023-03-24, from https://journals.ametsoc.org/
346	view/journals/atsc/63/7/jas3726.1.xml (Publisher: American Me-
347	teorological Society Section: Journal of the Atmospheric Sciences) doi:
348	10.1175/JAS3726.1
349	Chen, P., & Newman, M. (1998, October). Rossby Wave Propagation and
350	the Rapid Development of Upper-Level Anomalous Anticyclones during
351	the 1988 U.S. Drought. Journal of Climate, $11(10)$, $2491-2504$. Re-
352	trieved 2023-04-07, from https://journals.ametsoc.org/view/journals/
353	clim/11/10/1520-0442_1998_011_2491_rwpatr_2.0.co_2.xml (Pub-
354	lisher: American Meteorological Society Section: Journal of Climate) doi:
355	10.1175/1520-0442(1998)011(2491:RWPATR)2.0.CO;2
356	Ding, Q., & Wang, B. (2005, September). Circumglobal Teleconnection in the
357	Northern Hemisphere Summer. Journal of Climate, 18(17), 3483–3505. Re-
358	trieved 2023-04-07, from https://journals.ametsoc.org/view/journals/
359	clim/18/17/jcli3473.1.xml (Publisher: American Meteorological Society
360	Section: Journal of Climate) doi: 10.1175/JCLI3473.1
361	Donat, M. G., King, A. D., Overpeck, J. T., Alexander, L. V., Durre, I., & Karoly,
362	D. J. (2016, January). Extraordinary heat during the 1930s US Dust Bowl
363	and associated large-scale conditions. Climate Dynamics, $46(1)$, $413-426$.
364	Retrieved 2023-04-07, from https://doi.org/10.1007/s00382-015-2590-5
365	doi: 10.1007/s00382-015-2590-5
366	Douville, H. (2002, April). Influence of Soil Moisture on the Asian and African
367	Monsoons. Part II: Interannual Variability. Journal of Climate, 15(7), 701–
368	720. Retrieved 2023-04-07, from https://journals.ametsoc.org/view/
369	journals/clim/15/7/1520-0442_2002_015_0701_iosmot_2.0.co_2.xml
370	(Publisher: American Meteorological Society Section: Journal of Climate)
371	doi: $10.1175/1520-0442(2002)015(0701:IOSMOT)2.0.CO;2$
372	Douville, H., & Chauvin, F. (2000, October). Relevance of soil moisture for seasonal
373	climate predictions: a preliminary study. <i>Climate Dynamics</i> , 16(10), 719–736. Retrieved 2023-04-07, from https://doi.org/10.1007/s003820000080 doi:
374	10.1007/s00382000080
375	Feldstein, S. B. (2000, December). The Timescale, Power Spectra, and Climate
376	Noise Properties of Teleconnection Patterns. Journal of Climate, 13(24),
377	4430-4440. Retrieved 2023-03-31, from https://journals.ametsoc.org/
378 379	view/journals/clim/13/24/1520-0442_2000_013_4430_ttpsac_2.0.co_2.xml
379	(Publisher: American Meteorological Society Section: Journal of Climate) doi:
380	10.1175/1520-0442(2000)013(4430:TTPSAC)2.0.CO;2
501	10.11.0, 1020 0112(2000)010(1100.111) 010(2000)2

382	Fragkoulidis, G., & Wirth, V. (2020, October). Local Rossby Wave Packet Ampli-
383	tude, Phase Speed, and Group Velocity: Seasonal Variability and Their Role
384	in Temperature Extremes. Journal of Climate, 33(20), 8767–8787. Retrieved
385	2023-04-07, from https://journals.ametsoc.org/view/journals/clim/33/
386	20/jcliD190377.xml (Publisher: American Meteorological Society Section:
387	Journal of Climate) doi: 10.1175/JCLI-D-19-0377.1
388	Held, I. M. (2005, November). The Gap between Simulation and Understanding in
389	Climate Modeling. Bulletin of the American Meteorological Society, $86(11)$,
390	1609–1614. Retrieved 2023-04-15, from https://journals.ametsoc.org/
391	view/journals/bams/86/11/bams-86-11-1609.xml (Publisher: Ameri-
392	can Meteorological Society Section: Bulletin of the American Meteorological
393	Society) doi: 10.1175/BAMS-86-11-1609
394	Held, I. M., & Suarez, M. J. (1994, October). A Proposal for the Intercomparison
395	of the Dynamical Cores of Atmospheric General Circulation Models. Bulletin
396	of the American Meteorological Society, 75(10), 1825–1830. Retrieved 2023-
397	03-24, from https://journals.ametsoc.org/view/journals/bams/75/10/
398	1520-0477_1994_075_1825_apftio_2_0_co_2.xml (Publisher: American Me- taorological Society Section: Publishin of the American Metaorological Society)
399 400	teorological Society Section: Bulletin of the American Meteorological Society) doi: 10.1175/1520-0477(1994)075(1825:APFTIO)2.0.CO;2
401	Hoskins, B. J., & Ambrizzi, T. (1993, June). Rossby Wave Propagation on a Realis-
402	tic Longitudinally Varying Flow. Journal of the Atmospheric Sciences, 50(12),
403	1661-1671. Retrieved 2023-04-07, from https://journals.ametsoc.org/
404	view/journals/atsc/50/12/1520-0469_1993_050_1661_rwpoar_2_0_co_2.xml
405	(Publisher: American Meteorological Society Section: Journal of the Atmo-
406	spheric Sciences) doi: $10.1175/1520-0469(1993)050(1661:RWPOAR)2.0.CO;2$
407	Jiménez-Esteve, B., & Domeisen, D. I. (2022, July). The role of atmospheric dynam-
408	ics and large-scale topography in driving heatwaves. Quarterly Journal of the
409	Royal Meteorological Society, 148(746), 2344–2367. Retrieved 2023-04-07, from
410	https://onlinelibrary.wiley.com/doi/10.1002/qj.4306 doi: $10.1002/qj$
411	.4306
412	Koster, R. D., Chang, Y., Wang, H., & Schubert, S. D. (2016, October). Impacts of
413	Local Soil Moisture Anomalies on the Atmospheric Circulation and on Remote
414	Surface Meteorological Fields during Boreal Summer: A Comprehensive Anal-
415	ysis over North America. Journal of Climate, 29(20), 7345–7364. Retrieved
416	2023-04-07, from https://journals.ametsoc.org/view/journals/clim/
417	29/20/jcli-d-16-0192.1.xml (Publisher: American Meteorological Society
418	Section: Journal of Climate) doi: 10.1175/JCLI-D-16-0192.1
419	Lyon, B., & Dole, R. M. (1995, June). A Diagnostic Comparison of the 1980 and
420	1988 U.S. Summer Heat Wave-Droughts. Journal of Climate, 8(6), 1658–
421	1675. Retrieved 2023-04-07, from https://journals.ametsoc.org/view/ journals/clim/8/6/1520-0442_1995_008_1658_adcota_2_0_co_2.xml (Pub-
422	journals/clim/8/6/1520-0442_1995_008_1658_adcota_2_0_co_2.xml (Pub- lisher: American Meteorological Society Section: Journal of Climate) doi:
423	10.1175/1520-0442(1995)008(1658:ADCOTA)2.0.CO;2
424	McKinnon, K. A., Rhines, A., Tingley, M. P., & Huybers, P. (2016, May). Long-
425 426	lead predictions of eastern United States hot days from Pacific sea surface
420	temperatures. Nature Geoscience, $9(5)$, 389–394. Retrieved 2023-04-07, from
428	https://www.nature.com/articles/ngeo2687 (Number: 5 Publisher: Na-
429	ture Publishing Group) doi: 10.1038/ngeo2687
430	Meehl, G. A., & Tebaldi, C. (2004, August). More Intense, More Frequent, and
431	Longer Lasting Heat Waves in the 21st Century. Science, 305 (5686), 994–
432	997. Retrieved 2023-04-07, from https://www.science.org/doi/10.1126/
433	science.1098704 (Publisher: American Association for the Advancement of
434	Science) doi: 10.1126/science.1098704
435	Miralles, D. G., Gentine, P., Seneviratne, S. I., & Teuling, A. J. (2019, January).
436	Land–atmospheric feedbacks during droughts and heatwaves: state of the sci-

437	ence and current challenges. Annals of the New York Academy of Sciences,
438	1436(1), 19-35. Retrieved 2023-04-07, from https://www.ncbi.nlm.nih.gov/
439	pmc/articles/PMC6378599/ doi: 10.1111/nyas.13912
440	Pedlosky, J. (2003). Waves in the Ocean and Atmosphere: Introduction to Wave Dy-
441	namics (2003rd edition ed.). Berlin; New York: Springer.
442	Röthlisberger, M., Frossard, L., Bosart, L. F., Keyser, D., & Martius, O. (2019,
443	June). Recurrent Synoptic-Scale Rossby Wave Patterns and Their Effect on
444	the Persistence of Cold and Hot Spells. Journal of Climate, 32(11), 3207–3226.
445	Retrieved 2023-04-07, from https://journals.ametsoc.org/view/journals/
446	clim/32/11/jcli-d-18-0664.1.xml doi: 10.1175/JCLI-D-18-0664.1
447	Schubert, S., Wang, H., & Suarez, M. (2011, September). Warm Season Subseasonal
448	Variability and Climate Extremes in the Northern Hemisphere: The Role of
449	Stationary Rossby Waves. Journal of Climate, 24(18), 4773–4792. Retrieved
450	2023-04-15, from https://journals.ametsoc.org/view/journals/clim/24/
451	18/jcli-d-10-05035.1.xml (Publisher: American Meteorological Society
452	Section: Journal of Climate) doi: 10.1175/JCLI-D-10-05035.1
453	Seneviratne, S. I., Corti, T., Davin, E. L., Hirschi, M., Jaeger, E. B., Lehner, I.,
454	Teuling, A. J. (2010, May). Investigating soil moisture-climate interactions
455	in a changing climate: A review. $Earth-Science Reviews, 99(3), 125-161.$
456	Retrieved 2023-04-15, from https://www.sciencedirect.com/science/
457	article/pii/S0012825210000139 doi: 10.1016/j.earscirev.2010.02.004
458	Teng, H., Branstator, G., Wang, H., Meehl, G. A., & Washington, W. M. (2013,
459	December). Probability of US heat waves affected by a subseasonal planetary
460	wave pattern. Nature Geoscience, $6(12)$, 1056–1061. Retrieved 2023-03-24,
461	from https://www.nature.com/articles/ngeo1988 (Number: 12 Publisher:
462	Nature Publishing Group) doi: 10.1038/ngeo1988
463	Vallis, G. K. (2017). Atmospheric and Oceanic Fluid Dynamics: Funda-
464	mentals and Large-Scale Circulation (2nd ed.). Cambridge: Cambridge
465	University Press. Retrieved 2023-03-01, from https://www.cambridge
466	.org/core/books/atmospheric-and-oceanic-fluid-dynamics/ 41379BDDC4257CBE11143C466F6428A4 doi: 10.1017/9781107588417
467	Wirth, V., Riemer, M., Chang, E. K. M., & Martius, O. (2018, July). Rossby Wave
468	Packets on the Midlatitude Waveguide—A Review. Monthly Weather Review,
469	146(7), 1965–2001. Retrieved 2023-04-05, from https://journals.ametsoc
470 471	.org/view/journals/mwre/146/7/mwr-d-16-0483.1.xml (Publisher: Amer-
471	ican Meteorological Society Section: Monthly Weather Review) doi: 10.1175/
473	MWR-D-16-0483.1
474	Wu, Z., & Reichler, T. (2018). Towards a more earth-like circulation in idealized
475	models. Retrieved from https://doi.org/ doi: 10.1029/2018MS001356
476	Zimin, A. V., Szunyogh, I., Patil, D. J., Hunt, B. R., & Ott, E. (2003, May). Ex-
477	tracting Envelopes of Rossby Wave Packets. Monthly Weather Review, 131(5),
478	1011-1017. Retrieved 2023-03-24, from https://journals.ametsoc.org/
479	view/journals/mwre/131/5/1520-0493_2003_131_1011_eeorwp_2.0.co_2.xml
480	(Publisher: American Meteorological Society Section: Monthly Weather Re-
481	view) doi: $10.1175/1520-0493(2003)131(1011:EEORWP)2.0.CO;2$

Supporting Information for "The Role of Climatological State on Driving US Heat Waves Through Rossby Waves Packets"

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Contents of this file

- 1. Description of the Dry Core GCM
- 2. Figures S1 to S2

Introduction

This document includes a more detailed description of the model used for the three simulations in this work. In addition, two figures with a comparison of the climatological state between the experiments and the reanalysis data.

1. Description of the Dry Core GCM

Idealized general circulation models are commonly used for the study of atmospheric dynamics. This kind of models solve the primitive equations by nudging the temperature toward a prescribed equilibrium temperatures (T_{eq}) . Essentially, this process isolates the dynamics from the complex physical parametrizations. However, as the primitive

equations are solved on the sphere, there is some confidence that large-scale dynamics relate to the real atmosphere. In this work, we use the open-access updated version from (Wu & Reichler, 2018) of the spectral dynamical core model proposed by (Held & Suarez, 1994) for the Geophysical Fluid Dynamics Laboratory (GFDL).

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The model has horizontal resolution of T42 (64x128 grid) and 40 vertical σ levels between the surface and 0.01 hPa. Rayleigh drag with a prescribed rate, which decreases linearly with height from a surface value k_f to zero at ($\sigma = 0.7$) and higher levels, is used to remove momentum from the low levels, representing the boundary-layer friction. Mathematically, this can be written by:

$$\frac{\partial v}{\partial t} = \dots - k_v(\sigma)v \tag{1}$$

Where v represents the wind and σ the vertical sigma level. The damping rate (k_v) follows the form:

$$k_v = k_f max \left(0, \frac{\sigma - \sigma_b}{1 - \sigma_b}\right) \tag{2}$$

where $k_f = 1 day^{-1}$ and $\sigma_b = 0.7$ is the vertical level in the model where the top of the boundary layer is defined. The temperature is forced by Newtonian relaxation toward a prescribed equilibrium temperature as follows:

$$\frac{\partial T}{\partial t} = \frac{T - T_{eq}}{\tau} \tag{3}$$

where τ is the prescribed relaxation timescale following the distribution:

$$\tau^{-1} = k_a + (k_s - k_a)max\left(0, \frac{\sigma - \sigma_b}{1 - \sigma_b}\right)cos^4\phi \tag{4}$$

where $k_a = 1/40 \ day^{-1}$ and $k_s = 1/4 \ day^{-1}$ are parameters to set the distribution of the relaxation coefficient. In this model, the T_{eq} is zonally symmetric and is easily set through parameters. Since the primary purpose of this study is to examine how the atmospheric basic state controls the pattern contained in Rossby Waves driving heat waves, the model includes the iterative procedure proposed by (Chang, 2006) to simulate the climatological basic state of the atmosphere in the idealized model. This consists of iterating the radiative equilibrium temperature profile so that at the end of the iterations, the model climate closely resembles the desired target climate. The iteration uses a fixed equilibrium temperature for each N step (T_{eq}) in a run of Y years. Then, we calculate the model simulated temperature climatology $T_{(N)}$ and correct it concerning basic state from the NCEP reanalysis data T_R . The next iteration step N+1 is calculated according to:

:

$$T_{eq(N+1)} = T_{eq(N)} - \frac{2}{3}(T_{(N)} - T_R), N = 1, 2, 3...$$
(5)

References

Chang, E. K. M. (2006, July). An Idealized Nonlinear Model of the Northern Hemisphere Winter Storm Tracks. Journal of the Atmospheric Sciences, 63(7), 1818–1839. Retrieved 2023-03-24, from https://journals.ametsoc.org/view/journals/atsc/ 63/7/jas3726.1.xml (Publisher: American Meteorological Society Section: Journal of the Atmospheric Sciences) doi: 10.1175/JAS3726.1

Held, I. M., & Suarez, M. J. (1994, October). A Proposal for the Intercomparison of the Dynamical Cores of Atmospheric General Circulation Models. Bulletin of the American Meteorological Society, 75(10), 1825–1830. Retrieved 2023-03-24, from https://journals.ametsoc.org/view/journals/bams/75/10/1520-0477_1994_075_1825_apftio_2_0_co_2.xml (Publisher: American Meteorological Society Section: Bulletin of the American Meteorological Society) doi: 10.1175/1520-0477(1994)075(1825:APFTIO)2.0.CO;2

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Wu, Z., & Reichler, T. (2018). Towards a more earth-like circulation in idealized models. Retrieved from https://doi.org/ doi: 10.1029/2018MS001356

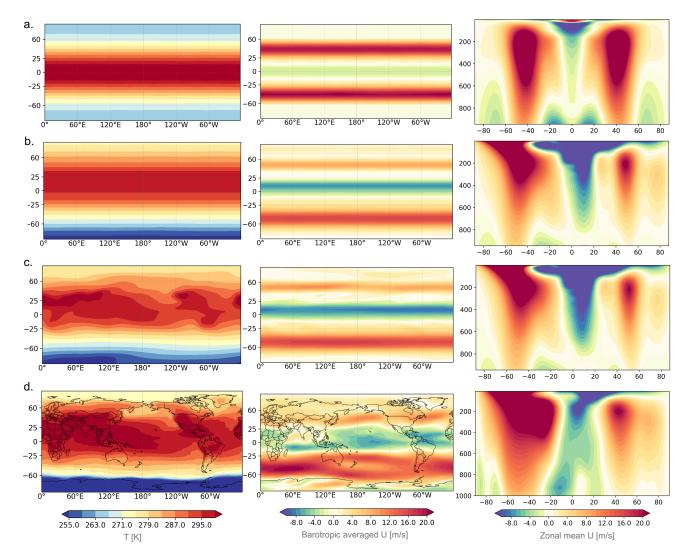


Figure S1. Basic state of the surface air temperature (SAT) (left column), the barotropic averaged zonal wind (middle column) and the vertical distribution of the zonal wind (right column) for the Held and Suarez 1994 (exp1_HS94) experiment (a.), the bias corrected experiment with the symmetrical SAT from NCEP (b.), the bias corrected experiment with the asymmetrical SAT from NCEP (d.).

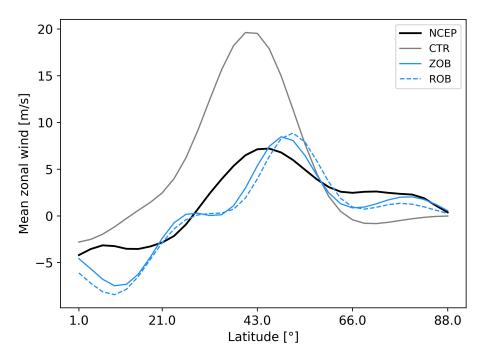


Figure S2. Comparison of the barotropic zonal mean zonal wind between NCEP, the Held and Suarez 1994 (exp1_HS94) experiment and the bias corrected experiments.

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