Modeling TEC Irregularities in the Northern Hemisphere Using Empirical Orthogonal Function Method

Yaqi Jin¹, Wojciech Jacek Miloch¹, Daria Kotova¹, Knut Stanley Jacobsen², -Dor⁻de Stevanovic³, Lasse Boy Novock Clausen¹, Nicholas Ssessanga¹, and Federico Da Dalt⁴

¹University of Oslo ²Norwegian Mapping Authority ³GMV Innovating Solutions ⁴Rhea System GmbH

April 16, 2023

Abstract

We develop a climatological model for the Northern Hemisphere based on a long-term dataset (2010-2021) of the rate of change of the total electron content (TEC) index (ROTI) maps from the International GNSS Service (IGS). The IGS ROTI maps are daily averaged in magnetic latitude and local time coordinates. To develop a climatological model, the ROTI maps are decomposed into a few base functions and coefficients using the empirical orthogonal function (EOF) method. The EOF method converges very quickly, and the first four EOFs reflect the majority (96%) of the total data variability. Furthermore, different EOF components can reflect different drivers of ionospheric irregularities. The first EOF reflects the averaged ROTI activity and the impact of the solar radiation and geomagnetic activity; the 2nd EOF reflects the impact of the interplanetary magnetic field (IMF) Bz and electric field; the 3rd and 4th EOFs reflect the dawn-dusk asymmetry around the auroral oval and polar cap, and they can be related to the IMF By. To build an empirical model, we fit the EOF coefficients using helio-geophysical indices from four different categories (solar activity; geomagnetic indices; IMF; the solar wind coupling function). The final EOF model is dependent on seven selected indices (F10.7P, Kp, Dst, Bt, By, Bz and Ekl). The statistical data-model comparisons show satisfactory results with a good correlation coefficient. However, the model cannot capture the significant expansion of the dayside ROTI activity during strong geomagnetic storms. Future effort is needed to provide corrections to the model for severe storms.

Hosted file

961109_0_art_file_10895187_rt3nv4.docx available at https://authorea.com/users/530554/ articles/635625-modeling-tec-irregularities-in-the-northern-hemisphere-using-empiricalorthogonal-function-method

1	Modeling TEC Irregularities in the Northern Hemisphere Using Empirical Orthogonal Function
2	Method
3	Yaqi Jin ^{1*} , Wojciech J. Miloch ¹ , Daria Kotova ¹ , Knut Stanley Jacobsen ² , Đorđe Stevanovic ³ , Lasse B. N.
4	Clausen ¹ , Nicholas Ssessanga ¹ , and Federico Da Dalt ⁴
5	¹ Department of Physics, University of Oslo, P.O. Box 1048 Blindern, 0316 Oslo, Norway
6	² Norwegian Mapping Authority, PO 600 Sentrum, 3507 Hønefoss, Norway
7	³ GMV Innovating Solutions, Warsaw, Poland
8	⁴ Rhea System GmbH, Darmstadt, Germany
9	
10	*Corresponding author: Yaqi Jin (<u>yaqi.jin@fys.uio.no</u>)
11	
12	Short title: Modeling of ROTI maps
13	
14	Key Points
15	1. An empirical model of ROTI maps in the Arctic is developed based on 12 years of ROTI data
16	using the EOF method
17	2. The statistical data-model comparisons show satisfactory results
18	3. The model fails to capture the significant expansion of the dayside ROTI activity during
19	severe geomagnetic storms

20 Abstract

21 We develop a climatological model for the Northern Hemisphere based on a long-term dataset 22 (2010-2021) of the rate of change of the total electron content (TEC) index (ROTI) maps from the 23 International GNSS Service (IGS). The IGS ROTI maps are daily averaged in magnetic latitude and local 24 time coordinates. To develop a climatological model, the ROTI maps are decomposed into a few base 25 functions and coefficients using the empirical orthogonal function (EOF) method. The EOF method 26 converges very quickly, and the first four EOFs reflect the majority (96%) of the total data variability. 27 Furthermore, different EOF components can reflect different drivers of ionospheric irregularities. The 28 first EOF reflects the averaged ROTI activity and the impact of the solar radiation and geomagnetic 29 activity; the 2nd EOF reflects the impact of the interplanetary magnetic field (IMF) Bz and electric 30 field; the 3rd and 4th EOFs reflect the dawn-dusk asymmetry around the auroral oval and polar cap, 31 and they can be related to the IMF By. To build an empirical model, we fit the EOF coefficients using 32 helio-geophysical indices from four different categories (solar activity; geomagnetic indices; IMF; the 33 solar wind coupling function). The final EOF model is dependent on seven selected indices (F10.7P, 34 Kp, Dst, Bt, By, Bz and Ek). The statistical data-model comparisons show satisfactory results with a 35 good correlation coefficient. However, the model cannot capture the significant expansion of the 36 dayside ROTI activity during strong geomagnetic storms. Future effort is needed to provide 37 corrections to the model during severe storms.

38 Plain Language Summary

39 The ionosphere is often highly structured and contains significant irregularities of plasma density, 40 which can impact the Global Navigation Satellite System (GNSS) services that rely on trans-41 ionospheric radio waves. Due to the practical use of GNSS services, there is a high demand for 42 modeling and forecasting of ionospheric irregularities. In this study, we develop a climatological 43 model for the Northern Hemisphere based on a long-term dataset from the International GNSS 44 Service (IGS). The final model is dependent on a few helio-geophysical indices that represent the 45 main drivers/sources of ionospheric irregularities. The statistical data-model comparisons show that 46 the model is able to reproduce all the climatological features of ionospheric irregularities. This model 47 can be useful in forecasting the space weather impact for GNSS users using real-time data as input.

48 1 Introduction

49 Ionospheric irregularities and scintillations have received more and more attention due to the

50 increasing number of Global Navigation Satellite System (GNSS) users. The GNSS service can be

51 heavily degraded due to space weather phenomena (Basu et al., 2002; Kintner et al., 2007; Jakowski

52 et al., 2012). To enable forecasting and mitigation of the adverse space weather effects, it is essential

to model ionospheric irregularities and scintillations (Wernik et al., 2007; Secan et al., 1997;

54 Priyadarshi, 2015; McGranaghan et al., 2018).

55 Ionospheric irregularities have been investigated using in-situ techniques (e.g., Basu et al., 1990;

56 Coley & Heelis, 1998; Moen et al., 2012; Jin et al., 2022b) and ground-based instruments (e.g., Rino

57 et al., 1983; Tsunoda, 1988). With the advent of GNSS, ground-based GNSS receivers have been

58 widely used to monitor and study ionospheric irregularities. Many researchers use amplitude and

59 phase scintillation indices based on the high-resolution measurements from specialized GNSS

scintillation receivers (Van Dierendonck et al., 1993; Mitchell et al., 2005; Spogli et al., 2009; Jin et al.,

61 2019a), while others focus on the rate of change of total electron content (TEC) index (ROTI) as a

62 measure of ionospheric irregularities (Pi et al., 1997; Basu et al., 1999). The advantage of ROTI is that

63 it can be obtained from any geodetic receivers (normally from dual-frequency measurements).

64 Therefore, the global extent and long-term variability of ionospheric irregularities can be investigated

using global networks of GNSS stations (e.g., Cherniak et al., 2014, 2018a; Jacobsen & Dahnn, 2014;

Jacobsen, 2014; Kotulak et al., 2020). More recently, daily ROTI maps in the northern mid- and high-

67 latitude regions have been developed by using GNSS data from about 700 selected permanent

68 stations (Cherniak et al., 2014; Cherniak et al., 2018b) The ROTI maps are provided to the

69 International GNSS Service (IGS) and are openly available for scientific purposes. The newly available

70 ROTI maps provide continuous information about the quiet-time background ionospheric

71 irregularities as well as intense ionospheric variability during storms (Cherniak et al., 2018b).

72 Therefore, the IGS ROTI maps provide a valuable dataset for developing a data-driven empirical

73 model of ionospheric irregularities.

74 In order to forecast ionospheric irregularities and scintillations, different approaches have been used.

75 The climatological models of ionospheric irregularities and scintillation have been widely used. For

reample, the WBMOD (WideBand MODel) ionospheric scintillation model is the most popular

climatological model that was first developed in the 1970s and upgraded later on (Secan et al., 1997).

78 Wernik et al. (2007) established a method to model scintillation effect using in-situ ionospheric

79 density data from Dynamics Explorer 2 satellite. A comprehensive review of existing climatological

80 models of ionospheric scintillation can be found in Priyadarshi (2015). More recently, Fabbro et al.

81 (2021) established a model of ROTI and GNSS positioning error by parameterization of the solar wind

measurements. Jin et al. (2022a) developed an empirical ROTI model based on GNSS data over two
solar cycles in Greenland using the Empirical Orthogonal Function (EOF) method. The latter model
only depends on F10.7P, Interplanetary Magnetic Field (IMF) Bt and Dst index. Despite its simplicity,
the model gives very satisfactory results with a high correlation coefficient. In the present paper, we
extend the 1-D ROTI model, and establish a 2-D ROTI model by using the IGS ROTI maps over one
solar cycle (2010-2021).

88 2 Dataset

89 2.1 ROTI maps

90 We make use of daily ROTI maps from the International GNSS Service (IGS) (Cherniak et al., 2014; 91 Cherniak et al., 2018b). ROTI is defined as the standard deviation of rate of change of TEC (ROT) in 5 92 minutes, where ROT is the time derivative of TEC (Pi et al., 1997). Both ROT and ROTI are in units of 93 TECU/min. To construct the ROTI maps, Cherniak et al. (2018b) collected Receiver Independent 94 Exchange (RINEX) format data with a 30-s sampling rate from about 700 GPS stations located at high 95 and middle latitudes in the Northern Hemisphere. The final product is presented as daily maps in 96 coordinates of magnetic latitude (MLAT) and magnetic local time (MLT). Each map starts from 50° 97 MLAT with a spatial resolution of 2° in MLAT and 2° in longitude (8 minutes in MLT). Although the 98 ROTI maps are from mid- to high latitudes, we focus on high latitudes in this present study. For more 99 detail about the methodology, please refer to Cherniak et al. (2014) and Cherniak et al. (2018b). 100 Figure 1 presents an example of ROTI maps during November 1 - 15, 2021. To assist eyes, the 101 Feldstein auroral oval (Q = 3) (Holzworth & Meng, 1975) is shown by magenta curves. The spatial 102 distribution of ROTI activity is generally guided by the auroral oval, where enhanced ROTI often 103 occurs around the auroral oval, and inside the polar cap (poleward of the auroral oval). Note that 104 there was a magnetic storm during November 3-5, 2021 that started at around 21 UT on November 3. 105 The daily ROTI map captures well the storm conditions, which are reflected as enhanced ROTI values 106 and expanded area of significant ROTI activity on November 4. The ROTI maps are available from 107 January 2010 to November 2022 through the ftp server of NASA's Crustal Dynamics Data Information 108 System (CDDIS) ftp://cddis.gsfc.nasa.gov/pub/gps/products/ionex (accessed in April 2023). In this 109 study, we use ROTI maps from 2010 to 2021 to develop the EOF model, while the data in 2022 are 110 used only for the purpose of validation.





117 2.2 Helio-geophysical proxies

118 The high-latitude ionosphere is directly driven by the solar activity, solar wind, and magnetospheric 119 phenomena. We use different solar and geophysical indices to represent these drivers. We use four 120 general classes of helio-geophysical proxies that characterize the space environment: 1) the solar 121 activity (F10.7, sunspot number etc.), 2) geomagnetic indices (Dst, Kp, ap etc.), 3) the IMF and solar 122 wind conditions (IMF Bt, Bx, By, Bz, the solar wind speed Vsw, solar wind density and pressure), 4) 123 interplanetary electric field (IEF) (IEF Ey, Ekl) and the Newell coupling function. 124 The F10.7 index measures the total emission of the solar disk at the wavelength of 10.7 cm, and it is 125 widely used to represent the solar activity (e.g., Tapping, 2013). Due to its simplicity, the F10.7 index 126 is often used to model the ionosphere (Bilitza et al., 2011). To represent the solar activity, there are 127 also variations of the F10.7 index. For example, F10.7 81 is an average of the F10.7 solar flux index in 128 a running window of 81 days (i.e., 3 solar rotations), while F10.7P is defined as (F10.7 + F10.7 81)/2 129 (Liu et al., 2011; Rentz & Lühr, 2008; Xiong et al., 2010). Liu et al. (2011) recommended using F10.7P 130 for the common use of proxy of the solar activity. We have compared F10.7, F10.7 81, F10.7P and

131 the sunspot number in our modeling work, and F10.7P turns out to be the best proxy for our purpose.

132 For the geomagnetic indices, we choose to use Kp, ap and Dst (Disturbance Storm Time) index. Kp

133 index is a three-hourly geomagnetic index in quasi-logarithm scale that is derived from the

- 134 standardized K index (Ks) of 13 worldwide observatories. It is designed to measure the solar particle
- 135 radiation, and it is considered as a proxy for the energy input from the solar wind to Earth. More
- details about the Kp index can be found on the official Kp index webpage (<u>https://www.gfz-</u>
- 137 <u>potsdam.de/en/kp-index/</u>). The 3-hourly ap (equivalent range) index can be derived from the Kp
- 138 index using a table from <u>https://www.ngdc.noaa.gov/stp/geomag/kp_ap.html</u>. The Dst index is an
- 139 hourly geomagnetic index that measures the intensity of the globally symmetrical equatorial
- 140 electrojet (the "ring current") and is derived from four magnetic observatories at low latitudes. The
- 141 variations of Dst provide quantitative information of the level of geomagnetic disturbances. The
- 142 original Dst index can be downloaded from the World Data Center for Geomagnetism, Kyoto

143 (http://wdc.kugi.kyoto-u.ac.jp/dstdir/). For the present study, we use the daily averaged Kp, ap and

- 144 Dst indices from the NASA CDAWeb (<u>https://cdaweb.gsfc.nasa.gov/index.html/</u>).
- 145 The upstream solar wind speed, solar wind dynamic pressure, IMF, and IEF are obtained from the
- daily averaged OMNI dataset (King & Papitashvili, 2005). In addition, we also use the Kan-Lee electric
- 147 field ($E_{KL} = V \cdot B_t \cdot \sin^2 \theta_c$, where V is the solar wind speed, B_t is the IMF strength, and θ_c is the
- 148 IMF clock angle) and the Newell coupling function $(d\Phi_{MP}/dt = V^{4/3}B_t^{2/3}\sin^{8/3}(\theta_c/2))$ to account for the
- solar wind-magnetosphere-ionosphere coupling (Kan & Lee, 1979; Newell et al., 2007). The clock
- angle (θ_c) is defined as the angle between the Z-axis and the projection of the IMF vector onto the Y-
- 151 Z plane of the geocentric solar magnetospheric (GSM) coordinate system.

152 3 Results

153 3.1 Climatology of ROTI maps

154 We first present the general climatology of ROTI maps at high latitudes. It has been shown that 155 ionospheric irregularities strongly depend on solar activity (Basu et al., 1988; Jin et al., 2018; De 156 Franceschi et al., 2019). Figure 2 shows the yearly averaged ROTI maps from 2010 to 2022. The yearly 157 averaged F10.7 index is also annotated on the top right of each panel. One can clearly observe solar 158 activity dependence of the global ROTI maps, i.e., the ROTI values increase from 2010 near solar 159 minimum to 2014 in solar maximum, then they decrease until 2019 followed by a slight increase of 160 ROTI in the new solar cycle 25 (from 2020). Besides the solar activity dependence, the distribution of 161 high ROTI values in each year is very similar. The averaged and standard deviation of ROTI maps are 162 shown on the last two panels in Figure 2. There are two regions of high ROTI, i.e., the dayside hotspot 163 around 75°-80° MLAT and 9-15 MLT, and the nightside hotspot around 68°-74° MLAT and 22-01 MLT. 164 The spatial distribution is similar to previous studies of ionospheric irregularities using regional ROTI

- data (Jacobsen, 2014), ionospheric scintillation data (Spogli et al., 2009; Jin et al., 2015), and in-situ
- 166 measurements (Jin et al., 2019b; Jin et al., 2021).
- 167 To present the seasonal variations, Figure 3 shows the averaged ROTI maps for two seasons using the
- data from 2010 to 2022. The summer months are averaged using data from May, June, and July,
- 169 while the winter months are averaged over November, December, and January. Clear seasonal
- 170 variations can be observed, i.e., the ROTI values are higher and more extended during winter.



Figure 2. The yearly averaged ROTI maps in 2010-2022. Note that only 11 months of ROTI data is
available in 2022. The last two panels show the mean and standard deviation (std) of ROTI over 20102022. The averaged F10.7 index for each time period is shown on the top right of each panel. The
Feldstein auroral oval (Q = 3) is shown by magenta curves to guide the analysis. Clear solar activity
dependence is observed.

177

178 Next, we present the temporal variations of ROTI activity in terms of season and solar cycle. Figure 4a 179 shows the F10.7P index (red) and sunspot number (black). The solar cycle 23 starts in December 2008 180 and ends in December 2019. However, the IGS ROTI maps start from 2010. Therefore, only ROTI data 181 from 2010 are presented in this paper. Clearly, the F10.7P index and sunspot number are high during 182 solar maximum near 2014 and low during 2019-2020 near solar minimum. For example, the sunspot 183 number reaches zero at the end of 2019. The solar activity greatly affects the climatological features 184 of ROTI as shown in Figures 4b-4c. Figure 4b shows the ROTI climatology using MLT and day of year. 185 To construct the statistics, the data are averaged along the MLAT from 50° to 90° at each fixed MLT.

186 This is similar to the study by Jin et al. (2022a), where only a single GNSS station was used, while the

- 187 present paper uses IGS ROTI maps that are based on global networks of about 700 GNSS stations.
- 188 The most remarkable variations of ROTI are associated with solar activity. During high solar activity
- 189 years (2014-2015), the ROTI values are significantly enhanced. Besides, the averaged ROTI also shows
- 190 clear seasonal variations, i.e., they are enhanced during equinox and winter seasons. In addition,
- 191 ROTI shows diurnal variations, i.e., it varies as a function of MLT. ROTI increases from MLT noon and
- remains enhanced until midnight. During high solar activity years (e.g., the winter of 2014-2015),
- 193 ROTI can remain enhanced all day.



Figure 3. The averaged ROTI map for summer (May, June and July) and winter (November, Decemberand January) months in 2010-2022.

- 197 Figure 4c shows the "keogram" format of ROTI data that is scanned along the noon-midnight
- 198 meridian. The data are averaged over fixed MLT (11-13 MLT on the dayside and 23-01 MLT on the
- 199 nightside). There are two bands of persistent ROTI enhancement both on the dayside and the
- 200 nightside. On the dayside, the enhancement locates around 75°-80° MLAT, while on the nightside, it
- 201 is located around 68°-74° MLAT. These two bands are associated with the cusp and nightside auroral
- 202 dynamics respectively, and they are consistent with the spatial distribution of ROTI as shown in
- 203 Figure 2. Though the dayside and nightside bands of ROTI enhancement are almost permanent, their
- intensity is dependent on the solar activity, such as they are more enhanced in 2012-2015. The
- 205 central polar cap is between the two bands of ROTI enhancements. The central polar cap can be filled
- 206 with enhanced ROTI during disturbed periods. For example, plasma structures can be associated with
- 207 polar cap patches or tongue of ionization (TOI) (Foster et al., 2005; Zhang et al., 2015; Jin et al., 2014;
- 208 van der Meeren et al., 2014). This phenomenon is dependent on the IMF orientations and
- 209 geomagnetic activity, and they are more prevalent during high solar activity (Spicher et al., 2017).

- 210 Another interesting feature is the periodic variations of ROTI along MLAT around noon, i.e., high ROTI
- 211 is located at the lowest MLAT in winter, and it drifts to the highest MLAT in summer. This feature is
- similar to the latitudinal variations of GPS phase scintillations (Jin et al., 2018). This phenomenon has
- 213 been explained by the seasonal motion of the cusp location and the impact of the solar terminator
- 214 (Jin et al., 2018; Newell & Meng, 1989).



^{Year}
Figure 4. (a) The F10.7P solar flux index (red) and sunspot number (black). (b) The averaged ROTI
over all MLAT (50°-90°). The data are presented in bins of 0.13 hour in MLT and 1 day. (c) The
averaged ROTI along MLAT around magnetic noon (11-13 MLT) and magnetic midnight (23-01 MLT).
The data are presented as a function of MLAT from the dayside (top) to the nightside (bottom). Note

that some data are missing at the end of 2020.

221 3.2 EOF decomposition

- 222 Since the IGS ROTI maps are daily maps in geomagnetic coordinates, the ionospheric condition at
- each day can be expressed as *ROTI(MLAT, MLT, D*). We use data during 2010-2021 (12 years) to
- account for the variations of time scales over one solar cycle. The EOF method is used to decompose
- 225 the 3-D matrix into a series of EOF base maps and coefficients as follows:

$$ROTI(MLAT, MLT, D) = \sum_{i=1}^{4} E_i(MLAT, MLT) \times A_i(D)$$
⁽¹⁾

- 226 Where E_i (MLAT, MLT) is the i-th EOF base function, and A_i is the i-th EOF coefficient. To keep
- 227 simplicity, we only use the first four EOFs to reconstruct the ROTI maps. This is feasible as the EOF

method converges very quickly and the first few EOFs dominate the total variability (Dvinskikh &
Naidenova, 1991; Zhang et al., 2009; A et al., 2011; A et al., 2012; Jin et al., 2022a). Table 1 presents
the percentage of total ROTI variability captured by the first four EOFs. The first four orders of EOFs
contribute to 96.32% of the total data variability. Therefore, it is adequate to use the first four orders
of EOFs to build the climatological model.

233 Figure 5 shows the first four decomposed EOF coefficients, while Figure 6 shows the corresponding 234 EOF base functions. The correlation of the EOF coefficients with various helio-geophysical indices is 235 presented in Table 2. The first EOF component represents the average behavior of the ROTI maps. To 236 justify this statement, we plot the daily averaged ROTI as a red line in Figure 5a. A good 237 correspondence between A_1 and daily mean ROTI is observed. In addition, this point can be 238 confirmed from the first base function in Figure 6a, where the distribution of E_1 resembles the 239 averaged ROTI maps in Figure 2. The second EOF component is mainly controlled by the geomagnetic 240 activity and IMF Bz. Figure 6b shows an enhancement of E₂ from pre-noon to pre-midnight around 80° 241 MLAT and a depletion equatorward of the statistical auroral oval. Table 2 shows that the correlation 242 with the IMF Bz is high and positive (R=0.61). When the IMF Bz is positive, there should be an 243 enhanced ROTI around 80° MLAT from noon to dusk. This corresponds to a contracted auroral oval 244 during northward IMF Bz. When the IMF Bz is negative, there should be a ring of enhanced ROTI 245 equatorward of the Feldstein auroral oval. This corresponds to the expanded auroral oval during negative IMF Bz due to dayside reconnection (Lockwood et al., 1990; Cowley & Lockwood, 1992). The 246 3rd and 4th EOF base functions show a clear dawn-dusk asymmetry. This indicates that the 3rd and 4th 247 248 EOF components represent the dawn-dusk asymmetry that is mainly controlled by the IMF By. A close look at Table 2 confirms that the 3rd and 4th EOF coefficients indeed correlate with the IMF By. 249



250

Figure 5. The first four decomposed EOF coefficients. The mean ROTI that is averaged over all MLAT and MLT in each day is displayed as red line in panel a).



254 Figure 6. The first four decomposed EOF base functions.

255 Table 1. The percentage of data variability captured by the first four Empirical Orthogonal Functions

²⁵⁶ *(EOFs)*

EOF	1	2	3	4
Percentage of Variability	87.30	4.43	2.42	2.16
Cumulative Percentage of Variability	87.30	91.73	94.16	96.32

257

Table 2. Correlation coefficients of the commonly used solar and geomagnetic indices with the first
 four EOF coefficients.

Coeff	A1	A2	A3	A4
F10.7P index	0.65	0.25	-0.55	-0.30
Sunspot Number	0.56	0.21	-0.45	-0.26
Kp index	0.61	-0.71	-0.04	-0.29
Dst index	-0.51	0.57	0.15	0.29
ap index	0.53	-0.68	-0.08	-0.33
IMF Bt	0.58	-0.30	-0.14	-0.34
IMF B _x	-0.00	0.09	0.31	-0.31
IMF B _y (GSM)	0.04	-0.14	-0.43	0.41
IMF B _z (GSM)	-0.19	0.61	-0.00	0.17
Solar Wind Speed	0.05	-0.09	0.07	-0.11
Solar Wind Pressure	0.33	-0.39	0.07	-0.09
IEF Ey	0.19	-0.62	-0.00	-0.19
Newell Coupling	0.45	-0.71	-0.03	-0.29
Function				
E _{kl}	0.53	-0.71	-0.06	-0.35

260 The next step is to fit the EOF coefficients by different classes of helio-geophysical indices. We use 4

261 general categories of helio-geophysical indices: 1) the solar activity (F10.7, sunspot number etc.), 2)

262 geomagnetic indices (Dst, Kp, ap etc.), 3) solar wind and IMF components (solar wind speed Vsw,

solar wind density, IMF Bt, Bx, By, Bz), 4) IEF and related indices (IEF, Ekl, the Newell coupling

- function). To ensure the simplicity and robustness of the model, we only select one, the most
- representative index in each category. This selection is mainly based on correlation analysis with the
- 266 EOF coefficients as shown in Table 2. The selected indices are shown in bold font.

267 We use the following equations to fit the EOF coefficients A_i :

$$A_{i}(d) = B_{i1}(d) + B_{i2}(d) + B_{i3}(d)$$
⁽²⁾

$$B_{i1}(d) = C_{i1} + D_{i1} \times P_1(d) + E_{i1} \times P_2(d) + F_{i1} \times P_3(d) + G_{i1}$$
(3)

$$\times P_4(d)$$
(3)

$$B_{i2}(d) = [C_{i2} + D_{i2} \times P_1(d) + E_{i2} \times P_2(d) + F_{i2} \times P_3(d) + G_{i2}$$
(4)
 $\times P_4(d)] \cos(2\pi d/365.25)$
 $+ [H_{i2} + I_{i2} \times P_1(d) + J_{i2} \times P_2(d) + K_{i2} \times P_3(d) + L_{i2} \times P_4(d)] \sin(2\pi d/365.25)$

$$B_{i3}(d) = [C_{i3} + D_{i3} \times P_1(d) + E_{i3} \times P_2(d) + F_{i3} \times P_3(d) + G_{i3}$$
(5)
 $\times P_4(d)] \cos(4\pi d/365.25)$
 $+ [H_{i3} + I_{i3} \times P_1(d) + J_{i3} \times P_2(d) + K_{i3} \times P_3(d) + L_{i3} \times P_4(d)] \sin(4\pi d/365.25)$

269 Where $A_i(d)$ is the i-th EOF coefficient, and it can be expressed by three components, B_{i1} , B_{i2} , B_{i3} , 270 which represent the solar cycle, annual and semi-annual variations in the EOF coefficients. For the 271 best fits, P1, P2, P3 and P4 are the selected helio-geophysical parameters from 4 classes of helio-272 geophysical indices. For A1, they are F10.7P index, Kp, Bt and Ek; for A2, they are F10.7P index, Kp, Bz and Eki; for A3, they are F10.7P index, Dst, By and Eki; for A4, they are F10.7P index, Kp, By and Eki. The 273 274 fitting parameters, namely C, D, E, F, G, H, I, J, K and L, can be calculated by linear regression. The 275 regression coefficients are presented in Table 3. Using these regression coefficients, the EOF 276 coefficients can be calculated from the selected helio-geophysical indices using equations 2)-5).

Regression	i=1	i=2	i=3	i=4
coefficients				
C _{i1}	-5.65	0.29	4.82	3.62
D _{i1}	0.12	0.03	-0.04	-0.02
E _{i1}	0.24	-0.18	0.01	-0.02
F _{i1}	0.40	0.76	-0.37	0.38
G _{i1}	0.83	-0.15	-0.08	-0.73
C _{i2}	-5.58	-1.63	1.72	1.69
D _{i2}	0.06	0.02	-0.01	-0.02
E _{i2}	0.04	-0.00	0.01	-0.02
F _{i2}	0.19	-0.05	-0.28	0.18
G _{i2}	-0.08	-0.19	-0.40	-0.11
H _{i2}	-0.14	0.15	-0.01	-0.13
l _{i2}	-0.00	-0.00	-0.00	0.00
J _{i2}	0.00	0.01	0.00	-0.00

277 Table 3. Regression coefficients used in Equations 2)-5).

K _{i2}	-0.04	-0.12	-0.06	-0.01
L _{i2}	0.12	0.02	0.01	0.06
C _{i3}	2.19	0.35	-1.34	0.25
D _{i3}	-0.04	-0.01	0.02	-0.00
E _{i3}	0.05	0.05	0.01	-0.00
F _{i3}	-0.01	-0.08	0.11	-0.12
G _{i3}	-0.04	-0.21	0.13	-0.03
H _{i3}	0.55	0.04	0.08	-0.19
l _{i3}	-0.00	-0.00	-0.00	0.00
J _{i3}	-0.01	-0.00	0.00	0.01
K _{i3}	-0.06	-0.09	0.03	0.04
L _{i3}	0.09	0.07	0.15	-0.05



279 Figure 7. The EOF coefficients and the fitted coefficients using selected helio-geophysical indices.

Figure 7 shows the fits of the four EOF coefficients. The original EOF coefficients are also plotted to
 compare with the corresponding fits. The fits of the first two EOF coefficients agree with the original
 EOF coefficients very well, with only small discrepancies during high solar activity years during 2014-

283 2015. However, the fits of the 3rd and 4th EOF coefficients are poorer, especially during high solar
 activity years. Note the 3rd and 4th EOF only contains 4.5% of the total variability. Therefore, the
 inperfect fits will not greatly affect the overall performance of the EOF model.

286 3.3 Data-Model Comparison

287 In section 3.2, we have developed an empirical EOF model of ROTI in the Northern Hemisphere. 288 Given the EOF base functions, the regression coefficients (C, D, E, F, G, H, I, J, K and L), and the 289 selected helio-geophysical indices (F10.7P, Kp, Dst, Bt, By, Bz and E_{kl}), we are able to calculate the 290 modeled ROTI map at any time. To perform a validation of the EOF model, we directly compare the 291 modeled ROTI maps with the experimental data. Figure 8 shows the comparison between 292 experimental ROTI data and EOF model from years 2010 to 2022. The ROTI data in years 2010-2021 293 are used to develop the EOF model. Thus, the comparisons in panels a)-l) stand for self-validation. 294 The EOF model can well represent the experimental data as indicated by the high values of 295 correlation coefficients. We also note slightly lower correlation coefficients during the high solar 296 activity years in 2014 and 2015 (R=0.90) as compared to the lower solar activity years in 2010 297 (R=0.94). This can be explained by a slightly worse fit of the EOF coefficients during the high solar 298 activity years (cf. Figure 7). This kind of behavior is opposite to the EOF modeling of other 299 ionospheric parameters such as the F2 peak height and TEC (Zhang et al., 2009; A et al., 2012). The 300 aggregated comparison during 2010-2021 is shown in Figure 8n, which shows a high correlation 301 coefficient of 0.94. Figure 80 shows the histogram of the bias (Data – Model). The mean bias is 0.001 302 TECU/min, which indicates a very low bias. The standard deviation is 0.031 TECU/min. In order to 303 present a cross-validation of the EOF model, we show in Figure 8m a comparison with independent data in 2022. The correlation between data and model is still high (R=0.84). This suggests that the 304 305 EOF model is able to represent the climatological feature of ROTI maps for independent dataset.





307 308 Figure 8. (a-n) Scatter plots to show comparisons between data and modeled results using the helio-309 geophysical indices as input. In each scatter plot, one point is the mean value of one ROTI map at one 310 day (for both data and model). The Pearson correlation coefficient is shown on the top left of each 311 panel. The linear fits of data and modeled results are presented as a black line in each panel. The 312 average value of F10.7 index during the corresponding time period is presented on the top right of 313 each panel. Note that the data in 2022 are not used in the development of the EOF model. (o) the 314 histogram of the bias (Data – Model) in a bin-step of 0.02 TECU/min. STD = standard deviation; UQ = 315 upper quartile; LQ = lower quartile.

316 4 Discussion

317 In this paper, we have developed an empirical model based on the daily ROTI maps in the northern 318 mid- and high latitudes over 12 years (2010-2021). By using the EOF method, the ROTI maps are 319 decomposed into a series of base functions and corresponding coefficients. Due the rapid 320 convergence of the method, the first four EOFs can capture 96% of the total data variability. We thus 321 use the first four EOFs to construct the EOF model. By correlation analysis of the coefficients with 322 respect to various helio-geophysical indices, we select four classes of helio-geophysical parameters 323 to fit the EOF coefficients, while the EOF base functions are fixed. Given the regression coefficients (C, 324 D, E, F, G, H, I, J, K and L in Table 3) and the selected helio-geophysical indices (F10.7P, Kp, Dst, Bt, By, 325 Bz and E_{kl}), the EOF model can predict the ROTI map at any given time. The data-model comparison 326 in section 3.3 gives satisfactory results and high correlation coefficient. Note that the good

327 performance of the EOF model is on the climatological scale. We also note that the correlation

328 between data and model is slightly lower during high solar activity (cf. Figure 8). This might be

329 expected, as the validation of the International Reference Ionosphere (IRI) also indicates that

empirical model tends to poorly represent the variability of high-latitude ionosphere during high

solar activity (Themens & Jayachandran, 2016; Bjoland et al., 2016). This has been attributed to the

high sensitivity to geomagnetic conditions and highly nonlinear dynamics (Themens et al., 2018).

333 The advantage of the EOF model is that it can decompose the data variability into different

334 categories. Different helio-geophysical indices represent different drivers/sources of ionospheric

335 plasma irregularities. For example, the first EOF represents the average behavior of the ROTI activity, 336 and it is positively correlated with F10.7P, the geomagnetic activity and the IMF strength. The second 337 EOF mainly represents the latitudinal variations as well as the expansion and contraction of the ROTI 338 activity. There is a high correlation between the second EOF coefficient versus Kp and IMF Bz. For 339 example, during negative IMF Bz, the auroral oval and ionospheric plasma irregularities can expand 340 to lower latitudes (Li et al., 2010; Alfonsi et al., 2011; Jin et al., 2020; Cherniak & Zakharenkova, 2017). The 3rd and 4th EOFs represent the dawn-dusk asymmetry. These two components are related to the 341 342 IMF By which modulates ionospheric convection pattern at high latitudes (e.g., Weimer, 2005). This

feature is also consistent with the IMF By dependence of the polar cap patches and ionospheric
plasma irregularities at high latitudes (Spicher et al., 2017; Jin et al., 2019b).

345 The EOF model only depends on seven helio-geophysical indices (F10.7P, Kp, Dst, Bt, By, Bz and E_{kl}). 346 These helio-geophysical indices reflect different helio-geophysical drivers/sources (i.e., solar activity, 347 geomagnetic storms, IMF and solar wind) of irregularities in the high-latitude ionosphere. A number 348 of different helio-geophysical indices have been used to develop space weather models in the 349 literature. For example, Fabbro et al. (2021) developed a ROTI model by using the Newell coupling 350 function as a driver to parameterize the lognormal distribution of ROTI. The model is able to forecast 351 the probability of ROTI that exceeds a given threshold. We also note that Fabbro et al. (2021) used 352 ROTI based on 1-s resolution data, whilst we use ROTI data based on 30-s resolution data. The 353 calculation of ROTI using different resolution may result in different level of activity (Jacobsen, 2014). 354 However, one major difference in the methodology from our model is that the model does not take 355 into account of the variations that are related to the general level of the solar EUV radiation. A 356 climatological model must take such variations into account. As shown in section 3.1, the global ROTI 357 activity is greatly affected by the solar activity that can be characterized by F10.7P. The solar EUV 358 radiation is the major source of the background ionospheric density even for the polar ionosphere 359 (Moen et al., 2008; Yang et al., 2016; Jin & Xiong, 2020). During the solar maximum, a small

360 percentage of perturbations in the high background electron density would produce significant ROTI

361 and scintillations (Aarons et al., 1981; Basu et al., 1988; Jin et al., 2018). In order to forecast ROTI 362 activity across different levels of solar activity, it is necessary to use indices that represent the solar 363 activity. In our model, we use F10.7P, and this allows us to forecast ROTI more accurately over different phases of a solar cycle. The high-latitude ionosphere is a highly dynamic system with energy 364 365 inputs from a variety of space sources. In order to resolve the complexity of the system, there are 366 also studies that use as many helio-geophysical indices as possible. For example, McGranaghan et al. 367 (2018) used 51 helio-geophysical parameters to predict the high-latitude GNSS phase scintillation. In 368 our study, we try to balance the simplicity and accuracy of the model by using information only from 369 the major sources/drivers to describe the system. We thus include four classes of major drivers and 370 the selection of indices is done according to the correlation analysis and physical meaning. The 371 simplicity of our model ensures the robustness and makes it easy to implement and use for general 372 users.

373 We are aware that our newly developed EOF model is a climatological model, which represents the 374 ROTI activity during relative quiet time. Though it is not designed for disturbed conditions, we next 375 show the worst scenario during very disturbed condition. The St. Patrick's Day storm in 2015 is the 376 strongest storm in solar cycle 24 (Kamide & Kusano, 2015; Jacobsen & Andalsvik, 2016; Cherniak et 377 al., 2015; Astafyeva et al., 2015). The minimum Dst reached -223 nT on March 17, 2015. Figures 9a-378 9c present the IGS ROTI maps before and during the magnetic storm, while Figures 9d-9f show the 379 modeled results for the same time. The EOF model gives a quite consistent prediction of the ROTI 380 observations before the storm on March 16. There was one enhancement from noon to dusk and 381 76°-82° MLAT. However, the ROTI observations and the modeled results are guite different 382 (especially on the dayside) during the main phase (March 17) of the storm. On March 17, the area of 383 high ROTI is greatly expanded as a result of the expanded auroral oval (Prikryl et al., 2016; Jin & 384 Oksavik, 2018). In addition, the enhancement of ROTI associated with storm-enhanced density (SED) 385 at mid-latitudes (15-18 MLT, 50°-60° MLAT) is not reflected by the EOF model. The transport of polar 386 cap patches and TOI across the polar cap is well captured by the EOF model. On March 18 (recovery 387 phase), the EOF model is more consistent with the IGS ROTI map except for an area of enhancement 388 around 13-14 MLT and 60°-70° MLAT in Figure 9c. The ROTI enhancement near the dayside and 389 duskside auroral oval is well reproduced by the EOF model.

The inadequate prediction during strong storms can be partly attributed to fixed EOF base functions.
In our model, only the EOF coefficients are adjusted by helio-geophysical parameters as inputs, while
the EOF base functions stay fixed. This means that the helio-geophysical indices can only modulate
the contribution (intensity) of each base function by changing the corresponding EOF coefficients.
For example, the first base function E₁ represents the average behavior of the high-latitude ROTI. The

395 enhancement on the dayside is around 10-15 MLT and 75°-80° MLAT, while it is located around 21-02 396 MLT and 68° -74° MLAT on the nightside (cf. Figure 6a). The second base function E₂ represents the 397 impact of IMF Bz and geomagnetic activity. There is a ring of depletion of E₂, i.e., 70°-76° MLAT 398 around noon and 60°-70° MLAT on the nightside (cf. Figure 6b). This component can reflect the 399 expansion of the auroral activity due to negative IMF Bz. However, the most significant expansion 400 during major geomagnetic storms can only be partly modeled by the second EOF (E_2 and A_2), i.e., the 401 equatorward expansion to 70° MLAT around noon and 60° MLAT on the nightside. However, during 402 the severe geomagnetic storm on March 17, 2015, enhanced ROTI expanded to as low as 62° MLAT 403 around noon (Figure 9b). In addition, enhanced ROTI associated with the SED is not captured by any 404 of the EOF base functions. This suggests that severe geomagnetic storms, such as the one on March 405 17, 2015, cannot be represented fully by the EOF model.





Figure 9. The comparison of IGS ROTI maps (top) and ROTI maps calculated using the EOF model
(bottom) before and during the St. Patrick Day storm in 2015. In each panel, noon is to the top, and
dawn is to the right.

410 One straightforward way to improve the model performance during severe geomagnetic storms is to

411 build a storm model based on data during storms. For example, it is possible to make a correction

412 model that can adjust the quiet time model to account for storm-time variability (Araujo-Pradere et

- 413 al., 2002; Themens et al., 2018). The most obvious difference between the model and observations is
- the significant expansion of the dayside ROTI region. To overcome this drawback, one can make the

- second EOF base function E₂ to be driven by the IMF Bz as well, in order to account for the significant
- 416 expansion of auroral oval during great storms. In such a way, the IMF Bz and geomagnetic activity not
- 417 only modulate the amplitude but also the location of ROTI. Another phenomenon that is not
- 418 captured in our model is the mid-latitude SED. It is necessary to build a new model for SED as an
- adjustment to the current model. These methods and possibilities will be explored in future studies.

420 5. Summary and conclusion

421 In this study, we have constructed an empirical model based on the daily ROTI maps in the Northern 422 Hemisphere over 12 years (2010-2021). We first presented the long-term climatology of TEC 423 irregularities from the IGS ROTI maps (section 3.1). The ROTI maps were derived using 30-s resolution 424 GNSS data from worldwide GNSS stations (about 700 stations). The global distribution of used GNSS 425 stations ensures that the ROTI data should have a very little longitudinal effect or universal time 426 dependence. This is a great advantage over previous models that often use regional data from the 427 European or American sectors. The 2-D ROTI maps show a variety of spatial and temporal variability. 428 The most obvious variation is the solar activity dependence, and this variability can be well 429 characterized by the F10.7P index. The seasonal variations are also clear, i.e., ROTI is more enhanced 430 during winter as compared to summer. The seasonal variations can be explained by the balance of 431 the energy input into the polar ionosphere and the solar illumination that creates E region 432 conductance and damps the amplitude of ionospheric irregularities. The diurnal variations (with 433 respect to MLT) are also observed such that enhanced ROTI is observed from noon to dusk on the 434 dayside, and around magnetic midnight. The latitudinal scan of ROTI along the noon-midnight 435 meridian can also reveal some key features of ionospheric irregularities. For example, the region of 436 enhanced ROTI around noon shows oscillations over a period of one year. This feature has been 437 explained by the seasonal shift of the cusp and the impact of the solar terminator. 438 By applying the EOF method, the ROTI maps are decomposed into a series of base functions and 439 corresponding coefficients. Due to the fast convergence, the first four EOFs capture 96% of the total 440 data variability. We use the first four EOFs to build the EOF model. The first EOF is the dominant

- 441 component that represents the averaged picture of ROTI activity. The first EOF is mostly determined
- 442 by the solar activity and geomagnetic activity. The second EOF reflects an important deviation (4.43%
- 443 of total data variability) from the averaged ROTI activity, and it represents the impacts of
- 444 geomagnetic activity and IMF. More specifically, it can reflect the equatorward expansion of ROTI
- 445 activity during moderate southward turn of IMF Bz. The third and fourth EOFs reflect the dawn-dusk
- 446 asymmetry that is modulated by IMF By.

447 By correlation analysis of the EOF coefficients with respect to various helio-geophysical indices, we 448 choose four classes of helio-geophysical parameters to fit the EOF coefficients, while the EOF base 449 functions remain fixed. The final EOF model is only dependent on seven selected helio-geophysical 450 indices (F10.7P, Kp, Dst, Bt, By, Bz and Ek). Given the seven selected helio-geophysical indices, the 451 EOF model is capable to make prediction of ROTI maps at any time. The data-model comparison gives 452 satisfactory results and high correlation coefficient (Pearson correlation coefficient of 0.94). The 453 validation with data in 2022 outside of the model training dataset also gives satisfactory result (R = 454 0.84).

Note that the good performance of the EOF model is related to the quiet time in a climatological sense. We have also presented a data-model comparison during the St. Patrick's Day storm in 2015. During this severe geomagnetic storm, enhanced ROTI associated with mid-latitude SED and great equatorward expansion of enhanced ROTI to as low as 62° MLAT around noon is not captured by the EOF model. Future effort is needed to provide supplemental adjustments/corrections to the climatological model to account for severe storms.

461 In this study, we have expanded our previous work about 1-D EOF model using a single station to a 2-462 D EOF model of ROTI irregularities over the entire Arctic area. Due to the long-term datasets and 463 global distribution of GNSS stations, this new model should be applicable to the whole Arctic. The 464 EOF model is only dependent on 7 commonly used helio-geophysical indices, and thus it is simple to 465 implement and use. It is also possible to develop models that are dependent on fewer helio-466 geophysical parameters, in order to accommodate the scenario when some indices are not available 467 at times. However, to ensure an adequate performance, the simplified models should at least include 468 F10.7, geomagnetic activity and IMF. This study demonstrates that the EOF method is very suitable 469 and efficient for developing models for ionospheric irregularities. The EOF method is very 470 straightforward and it can also be applied to other irregularity indices, such as the GNSS amplitude 471 and phase scintillation indices from ground-based GNSS scintillation receivers, as well as in-situ 472 irregularity indices obtained from Low Earth Orbiting satellites (Jin et al., 2022b; Kotova et al., 2022).

473 Acknowledgment

- 474 WJM, YJ and DK acknowledge funding from the European Research Council (ERC) under the European
- 475 Union's Horizon 2020 research and innovation programme (ERC Consolidator Grant agreement No.
- 476 866357, POLAR-4DSpace). This research is a part of the 4DSpace Strategic Research Initiative at the
- 477 University of Oslo and the European Space Agency's network of space weather services and service
- 478 development activities, supported under ESA contract number 4000139062/22/D/KS.

479 **Open Research**

- 480 The ROTI maps can be obtained from the ftp server of NASA's Crustal Dynamics Data Information
- 481 System (CDDIS) <u>https://cddis.nasa.gov/archive/gnss/products/ionex/</u>. The helio-geophysical indices
- 482 can be obtained through <u>https://omniweb.gsfc.nasa.gov/form/dx1.html</u>.

483 References

- 484 A, E., Zhang, D., Ridley, A. J., Xiao, Z., & Hao, Y. (2012). A global model: Empirical orthogonal function 485 analysis of total electron content 1999–2009 data. Journal of Geophysical Research: Space 486 *Physics*, 117(A3). doi:https://doi.org/10.1029/2011JA017238 487 A, E., Zhang, D. H., Xiao, Z., Hao, Y. Q., Ridley, A. J., & Moldwin, M. (2011). Modeling ionospheric foF2 488 by using empirical orthogonal function analysis. Ann. Geophys., 29(8), 1501-1515. 489 doi:10.5194/angeo-29-1501-2011 490 Alfonsi, L., Spogli, L., De Franceschi, G., Romano, V., Aquino, M., Dodson, A., & Mitchell, C. N. (2011). 491 Bipolar climatology of GPS ionospheric scintillation at solar minimum. Radio Science, 46. 492 doi:10.1029/2010rs004571 493 Araujo-Pradere, E. A., Fuller-Rowell, T. J., & Codrescu, M. V. (2002). STORM: An empirical storm-time
- 493 Araujo-Pradere, E. A., Fuller-Rowell, T. J., & Codrescu, M. V. (2002). STORM: An empirical storm-time
 494 ionospheric correction model 1. Model description. *Radio Science*, *37*(5).
 495 doi:10.1029/2001rs002467
- Astafyeva, E., Zakharenkova, I., & Forster, M. (2015). Ionospheric response to the 2015 St. Patrick's
 Day storm: A global multi-instrumental overview. *Journal of Geophysical Research-Space Physics, 120*(10), 9023-9037. doi:10.1002/2015ja021629
- Basu, S., Basu, S., Mackenzie, E., Coley, W. R., Sharber, J. R., & Hoegy, W. R. (1990). Plasma
 Structuring by the Gradient Drift Instability at High-Latitudes and Comparison With Velocity
 Shear Driven Processes. *Journal of Geophysical Research-Space Physics, 95*(A6), 7799-+.
 doi:10.1029/JA095iA06p07799
- Basu, S., Groves, K. M., Basu, S., & Sultan, P. J. (2002). Specification and forecasting of scintillations in
 communication/navigation links: current status and future plans. *Journal of Atmospheric and Solar-Terrestrial Physics, 64*(16), 1745-1754. doi:10.1016/s1364-6826(02)00124-4
- Basu, S., Groves, K. M., Quinn, J. M., & Doherty, P. (1999). A comparison of TEC fluctuations and
 scintillations at Ascension Island. *Journal of Atmospheric and Solar-Terrestrial Physics*, *61*(16),
 1219-1226. doi:Doi 10.1016/S1364-6826(99)00052-8
- Basu, S., Mackenzie, E., & Basu, S. (1988). Ionospheric Constraints on VHF UHF Communications Links
 during Solar Maximum and Minimum Periods. *Radio Science*, 23(3), 363-378. doi:DOI
 10.1029/RS023i003p00363
- 512 Bilitza, D., McKinnell, L.-A., Reinisch, B., & Fuller-Rowell, T. (2011). The international reference
 513 ionosphere today and in the future. *Journal of Geodesy, 85*(12), 909-920.
 514 doi:10.1007/s00190-010-0427-x
- Bjoland, L. M., Belyey, V., Løvhaug, U. P., & La Hoz, C. (2016). An evaluation of International
 Reference Ionosphere electron density in the polar cap and cusp using EISCAT Svalbard radar
 measurements. *Ann. Geophys.*, *34*(9), 751-758. doi:10.5194/angeo-34-751-2016
- 518 Cherniak, I., Krankowski, A., & Zakharenkova, I. (2014). Observation of the ionospheric irregularities
 519 over the Northern Hemisphere: Methodology and service. *Radio Science, 49*(8), 653-662.
 520 doi:https://doi.org/10.1002/2014RS005433
- 521 Cherniak, I., Krankowski, A., & Zakharenkova, I. (2018a). ROTI Maps: a new IGS ionospheric product
 522 characterizing the ionospheric irregularities occurrence. *Gps Solutions, 22*(3), 69.
 523 doi:10.1007/s10291-018-0730-1

524	Cherniak, I., Krankowski, A., & Zakharenkova, I. (2018b). ROTI Maps: a new IGS ionospheric product
525	characterizing the ionospheric irregularities occurrence. Gps Solutions, 22(3).
526	doi:10.1007/s10291-018-0730-1
527	Cherniak, I., & Zakharenkova, I. (2017). New advantages of the combined GPS and GLONASS
528	observations for high-latitude ionospheric irregularities monitoring: case study of June 2015
529	geomagnetic storm. Earth Planets and Space, 69. doi:10.1186/s40623-017-0652-0
530	Cherniak, I., Zakharenkova, I., & Redmon, R. J. (2015). Dynamics of the high-latitude ionospheric
531	irregularities during the 17 March 2015St. Patrick's Day storm: Ground-based GPS
532	measurements. Space Weather-the International Journal of Research and Applications, 13(9),
533	585-597. doi:10.1002/2015sw001237
534	Coley, W. R., & Heelis, R. A. (1998). Structure and occurrence of polar ionization patches. <i>Journal of</i>
535	Geophysical Research-Space Physics, 103(A2), 2201-2208. doi:10.1029/97ja03345
536	Cowley, S. W. H., & Lockwood, M. (1992). Excitation and Decay of Solar Wind-Driven Flows in the
537	Magnetosphere-Ionosphere System. Annales Geophysicae-Atmospheres Hydrospheres and
538	<i>Space Sciences, 10</i> (1-2), 103-115. Retrieved from <go isi="" to="">://WOS:A1992HL02700010</go>
539	De Franceschi, G., Spogli, L., Alfonsi, L., Romano, V., Cesaroni, C., & Hunstad, I. (2019). The
540	ionospheric irregularities climatology over Svalbard from solar cycle 23. Scientific Reports. 9.
541	doi:10.1038/s41598-019-44829-5
542	Dvinskikh, N. I., & Naidenova, N. I. (1991). An adaptable regional empirical ionospheric model.
543	Advances in Space Research, 11, 7, doi:10.1016/0273-1177(91)90312-8
544	Fabbro, V., Jacobsen, K. S., Andalsvik, Y. L., & Rougerie, S. (2021). GNSS positioning error forecasting
545	in the Arctic: ROTI and Precise Point Positioning error forecasting from solar wind
546	measurements. J. Space Weather Space Clim., 11, 43. Retrieved from
547	https://doi.org/10.1051/swsc/2021024
548	Foster, J. C., Coster, A. J., Erickson, P. J., Holt, J. M., Lind, F. D., Rideout, W., Rich, F. J. (2005).
549	Multiradar observations of the polar tongue of ionization. Journal of Geophysical Research-
550	<i>Space Physics, 110</i> (A9). doi:10.1029/2004ja010928
551	Holzworth, R. H., & Meng, C. I. (1975). Mathematical Representation of Auroral Oval. <i>Geophysical</i>
552	Research Letters, 2(9), 377-380, doi:DOI 10.1029/GL002i009p00377
553	Jacobsen, K. S. (2014). The impact of different sampling rates and calculation time intervals on ROTI
554	values. J. Space Weather Space Clim., 4, A33. Retrieved from
555	https://doi.org/10.1051/swsc/2014031
556	Jacobsen, K. S., & Andalsvik, Y. L. (2016). Overview of the 2015 St. Patrick's day storm and its
557	consequences for RTK and PPP positioning in Norway. J. Space Weather Space Clim., 6, A9.
558	Retrieved from https://doi.org/10.1051/swsc/2016004
559	Jacobsen, K. S., & Dahnn, M. (2014). Statistics of ionospheric disturbances and their correlation with
560	GNSS positioning errors at high latitudes. Journal of Space Weather and Space Climate,
561	4(A27). doi:10.1051/swsc/2014024
562	Jakowski, N., Beniguel, Y., De Franceschi, G., Pajares, M. H., Jacobsen, K. S., Stanislawska, I.,
563	Wautelet, G. (2012). Monitoring, tracking and forecasting ionospheric perturbations using
564	GNSS techniques. Journal of Space Weather and Space Climate, 2, 14.
565	doi:10.1051/swsc/2012022
566	Jin, Y., Clausen, L. B. N., Miloch, W. J., Høeg, P., Jarmołowski, W., Wielgosz, P., García-Rigo, A.
567	(2022a). Climatology and modeling of ionospheric irregularities over Greenland based on
568	empirical orthogonal function method. J. Space Weather Space Clim., 12, 23. Retrieved from
569	https://doi.org/10.1051/swsc/2022022
570	Jin. Y., Clausen, L. B. N., Spicher, A., Ivarsen, M. F., Zhang, Y., Miloch, W. J., & Moen, J. I. (2021).
571	(1, 1)
J/1	Statistical Distribution of Decameter Scale (50 m) ionospheric irregularities at High Latitudes.
572	Geophysical Research Letters, 48(19), e2021GL094794.
572 573	Geophysical Research Letters, 48(19), e2021GL094794. doi:https://doi.org/10.1029/2021GL094794
572 573 574	<i>Geophysical Research Letters, 48</i> (19), e2021GL094794. doi:h <u>ttps://doi.org/10.1029/2021GL094794</u> Jin, Y., Kotova, D., Xiong, C., Brask, S. M., Clausen, L. B. N., Kervalishvili, G., Miloch, W. J. (2022b).

576	Satellites. Journal of Geophysical Research: Space Physics, 127(4), e2021JA030183.
577	doi:h <u>ttps://doi.org/10.1029/2021JA030183</u>
578	Jin, Y., Moen, J. I., Spicher, A., Oksavik, K., Miloch, W. J., Clausen, L. B. N., Saito, Y. (2019a).
579	Simultaneous Rocket and Scintillation Observations of Plasma Irregularities Associated With a
580	Reversed Flow Event in the Cusp Ionosphere. Journal of Geophysical Research: Space Physics,
581	<i>124</i> (8), 7098-7111. doi:10.1029/2019ja026942
582	Jin, Y., & Xiong, C. (2020). Interhemispheric Asymmetry of Large-Scale Electron Density Gradients in
583	the Polar Cap Ionosphere: UT and Seasonal Variations. Journal of Geophysical Research:
584	Space Physics, 125(2), e2019JA027601. doi:10.1029/2019ja027601
585	Jin, Y., Xiong, C., Clausen, L., Spicher, A., Kotova, D., Brask, S., Miloch, W. (2020). Ionospheric
586	Plasma Irregularities Based on In Situ Measurements From the Swarm Satellites. Journal of
587	Geophysical Research: Space Physics, 125(7), e2020JA028103.
588	doi:h <u>ttps://doi.org/10.1029/2020JA028103</u>
589	Jin, Y. Q., Miloch, W. J., Moen, J. I., & Clausen, L. B. N. (2018). Solar cycle and seasonal variations of
590	the GPS phase scintillation at high latitudes. Journal of Space Weather and Space Climate, 8.
591	doi:10.1051/swsc/2018034
592	Jin, Y. Q., Moen, J. I., & Miloch, W. J. (2014). GPS scintillation effects associated with polar cap
593	patches and substorm auroral activity: direct comparison. Journal of Space Weather and
594	Space Climate, 4(A23). doi:10.1051/swsc/2014019
595	Jin, Y. Q., Moen, J. I., & Miloch, W. J. (2015). On the collocation of the cusp aurora and the GPS phase
596	scintillation: A statistical study. Journal of Geophysical Research-Space Physics, 120(10),
597	9176-9191. doi:10.1002/2015ja021449
598	Jin, Y. Q., & Oksavik, K. (2018). GPS Scintillations and Losses of Signal Lock at High Latitudes During
599	the 2015 St. Patrick's Day Storm. Journal of Geophysical Research-Space Physics, 123(9),
600	7943-7957. doi:10.1029/2018ja025933
601	Jin, Y. Q., Spicher, A., Xiong, C., Clausen, L. B. N., Kervalishvili, G., Stolle, C., & Miloch, W. J. (2019b).
602	Ionospheric Plasma Irregularities Characterized by the Swarm Satellites: Statistics at High
603	Latitudes. Journal of Geophysical Research-Space Physics, 124(2), 1262-1282.
604	doi:10.1029/2018ja026063
605	Kamide, Y., & Kusano, K. (2015). No Major Solar Flares but the Largest Geomagnetic Storm in the
606	Present Solar Cycle. Space Weather-the International Journal of Research and Applications,
607	13(6), 365-367. doi:https://doi.org/10.1002/2015SW001213
608	Kan, J., & Lee, L. (1979). Energy coupling function and solar wind-magnetosphere dynamo.
609	Geophysical Research Letters, 6(7), 577-580.
610	King, J. H., & Papitashvili, N. E. (2005). Solar wind spatial scales in and comparisons of hourly Wind
611	and ACE plasma and magnetic field data. Journal of Geophysical Research-Space Physics,
612	<i>110</i> (A2). doi:10.1029/2004ja010649
613	Kintner, P. M., Ledvina, B. M., & De Paula, E. R. (2007). GPS and ionospheric scintillations. Space
614	Weather-the International Journal of Research and Applications, 5(9).
615	doi:10.1029/2006sw000260
616	Kotova, D., Jin, Y., Spogli, L., Wood, A. G., Urbar, J., Rawlings, J. T., Miloch, W. J. (2022). Electron
617	density fluctuations from Swarm as a proxy for ground-based scintillation data: A statistical
618	perspective. Advances in Space Research. doi:h <u>ttps://doi.org/10.1016/j.asr.2022.11.042</u>
619	Kotulak, K., Zakharenkova, I., Krankowski, A., Cherniak, I., Wang, N., & Fron, A. (2020). Climatology
620	Characteristics of Ionospheric Irregularities Described with GNSS ROTI. Remote Sensing,
621	12(16), 2634. Retrieved from h <u>ttps://www.mdpi.com/2072-4292/12/16/2634</u>
622	Li, G. Z., Ning, B. Q., Ren, Z. P., & Hu, L. H. (2010). Statistics of GPS ionospheric scintillation and
623	irregularities over polar regions at solar minimum. <i>Gps Solutions</i> , 14(4), 331-341.
624	doi:10.1007/s10291-009-0156-x
625	Liu, L., Wan, W., Chen, Y., & Le, H. (2011). Solar activity effects of the ionosphere: A brief review.
626	Chinese Science Bulletin, 56(12), 1202-1211. doi:10.1007/s11434-010-4226-9

627 Lockwood, M., Cowley, S. W. H., & Freeman, M. P. (1990). The Excitation of Plasma Convection in the 628 High-Latitude Ionosphere. Journal of Geophysical Research-Space Physics, 95(A6), 7961-7972. 629 doi:DOI 10.1029/JA095iA06p07961 630 McGranaghan, R. M., Mannucci, A. J., Wilson, B., Mattmann, C. A., & Chadwick, R. (2018). New 631 Capabilities for Prediction of High-Latitude Ionospheric Scintillation: A Novel Approach With 632 Machine Learning. Space Weather-the International Journal of Research and Applications, 633 16(11), 1817-1846. doi:https://doi.org/10.1029/2018SW002018 634 Mitchell, C. N., Alfonsi, L., De Franceschi, G., Lester, M., Romano, V., & Wernik, A. W. (2005). GPS TEC 635 and scintillation measurements from the polar ionosphere during the October 2003 storm. 636 Geophysical Research Letters, 32(12). doi:10.1029/2004gl021644 637 Moen, J., Oksavik, K., Abe, T., Lester, M., Saito, Y., Bekkeng, T. A., & Jacobsen, K. S. (2012). First in-situ 638 measurements of HF radar echoing targets. Geophysical Research Letters, 39. 639 doi:10.1029/2012gl051407 640 Moen, J., Qiu, X. C., Carlson, H. C., Fujii, R., & McCrea, I. W. (2008). On the diurnal variability in F2-641 region plasma density above the EISCAT Svalbard radar. Annales Geophysicae, 26(8), 2427-642 2433. doi:DOI 10.5194/angeo-26-2427-2008 Newell, P. T., & Meng, C. I. (1989). Dipole Tilt Angle Effects on the Latitude of the Cusp and Cleft Low-643 644 Latitude Boundary-Layer. Journal of Geophysical Research-Space Physics, 94(A6), 6949-6953. 645 doi:DOI 10.1029/JA094iA06p06949 646 Newell, P. T., Sotirelis, T., Liou, K., Meng, C. I., & Rich, F. J. (2007). A nearly universal solar wind-647 magnetosphere coupling function inferred from 10 magnetospheric state variables. Journal 648 of Geophysical Research-Space Physics, 112(A1). doi:10.1029/2006ja012015 649 Pi, X., Mannucci, A. J., Lindqwister, U. J., & Ho, C. M. (1997). Monitoring of global ionospheric 650 irregularities using the worldwide GPS network. Geophysical Research Letters, 24(18), 2283-651 2286. doi:10.1029/97gl02273 652 Prikryl, P., Ghoddousi-Fard, R., Weygand, J. M., Viljanen, A., Connors, M., Danskin, D. W., . . . Sreeja, V. 653 (2016). GPS phase scintillation at high latitudes during the geomagnetic storm of 17–18 654 March 2015. Journal of Geophysical Research: Space Physics, 121(10), 10,448-410,465. 655 doi:10.1002/2016ja023171 656 Priyadarshi, S. (2015). A Review of Ionospheric Scintillation Models. Surveys in Geophysics, 36(2), 657 295-324. doi:10.1007/s10712-015-9319-1 658 Rentz, S., & Lühr, H. (2008). Climatology of the cusp-related thermospheric mass density anomaly, as 659 derived from CHAMP observations. Annales Geophysicae, pp. 2807-2823. 660 Rino, C. L., Livingston, R. C., Tsunoda, R. T., Robinson, R. M., Vickrey, J. F., Senior, C., . . . Klobuchar, J. 661 A. (1983). Recent Studies of the Structure and Morphology of Auroral-Zone F-Region 662 Irregularities. Radio Science, 18(6), 1167-1180. doi:DOI 10.1029/RS018i006p01167 663 Secan, J. A., Bussey, R. M., Fremouw, E. J., & Basu, S. (1997). High-latitude upgrade to the Wideband 664 ionospheric scintillation model. Radio Science, 32(4), 1567-1574. 665 doi:https://doi.org/10.1029/97RS00453 666 Spicher, A., Clausen, L. B. N., Miloch, W. J., Lofstad, V., Jin, Y., & Moen, J. I. (2017). Interhemispheric 667 study of polar cap patch occurrence based on Swarm in situ data. Journal of Geophysical 668 Research-Space Physics, 122(3), 3837-3851. doi:10.1002/2016ja023750 669 Spogli, L., Alfonsi, L., De Franceschi, G., Romano, V., Aquino, M. H. O., & Dodson, A. (2009). 670 Climatology of GPS ionospheric scintillations over high and mid-latitude European regions. 671 Annales Geophysicae, 27(9), 3429-3437. doi:DOI 10.5194/angeo-27-3429-2009 672 Tapping, K. F. (2013). The 10.7 cm solar radio flux (F10.7). Space Weather-the International Journal of 673 Research and Applications, 11(7), 394-406. doi:10.1002/swe.20064 674 Themens, D. R., & Jayachandran, P. T. (2016). Solar activity variability in the IRI at high latitudes: 675 Comparisons with GPS total electron content. Journal of Geophysical Research-Space Physics, 676 121(4), 3793-3807. doi:10.1002/2016ja022664 677 Themens, D. R., Jayachandran, P. T., Reid, B., & McCaffrey, A. M. (2018). The Limits of Empirical 678 Electron Density Modeling: Examining the Capacity of E-CHAIM and the IRI for Modeling

679 Intermediate (1- to 30-Day) Timescales at High Latitudes. Radio Science, 55(4), e2018RS006763. doi:https://doi.org/10.1029/2018RS006763 680 681 Tsunoda, R. T. (1988). High-Latitude F-Region Irregularities - a Review and Synthesis. Reviews of Geophysics, 26(4), 719-760. doi:DOI 10.1029/RG026i004p00719 682 van der Meeren, C., Oksavik, K., Lorentzen, D., Moen, J. I., & Romano, V. (2014). GPS scintillation and 683 684 irregularities at the front of an ionization tongue in the nightside polar ionosphere. Journal of 685 Geophysical Research: Space Physics, 119(10), 8624-8636. doi:10.1002/2014ja020114 686 Van Dierendonck, A. J., Klobuchar, J., & Hua, Q. (1993). Ionospheric Scintillation Monitoring Using 687 Commercial Single Frequency C/A Code Receivers. Proceedings of the 6th International 688 Technical Meeting of the Satellite Division of The Institute of Navigation (ION GPS 1993), Salt 689 Lake City, UT, September 1993, September 1993, pp. 1333 - 1342. 690 Weimer, D. R. (2005). Improved ionospheric electrodynamic models and application to calculating 691 Joule heating rates. Journal of Geophysical Research-Space Physics, 110(A5). 692 doi:10.1029/2004ja010884 693 Wernik, A. W., Alfonsi, L., & Materassi, M. (2007). Scintillation modeling using in situ data. Radio 694 Science, 42(1). doi:https://doi.org/10.1029/2006RS003512 695 Xiong, C., Park, J., Luehr, H., Stolle, C., & Ma, S. Y. (2010). Comparing plasma bubble occurrence rates 696 at CHAMP and GRACE altitudes during high and low solar activity. Annales Geophysicae, 697 28(9), 1647-1658. doi:10.5194/angeo-28-1647-2010 698 Yang, S. G., Zhang, B. C., Fang, H. X., Kamide, Y., Li, C. Y., Liu, J. M., . . . Hu, H. Q. (2016). New evidence 699 of dayside plasma transportation over the polar cap to the prevailing dawn sector in the 700 polar upper atmosphere for solar-maximum winter. Journal of Geophysical Research-Space 701 Physics, 121(6), 5626-5638. doi:10.1002/2015ja022171 702 Zhang, M. L., Liu, C., Wan, W., Liu, L., & Ning, B. (2009). A global model of the ionospheric F2 peak 703 height based on EOF analysis. Ann. Geophys., 27(8), 3203-3212. doi:10.5194/angeo-27-3203-704 2009 705 Zhang, Q. H., Lockwood, M., Foster, J. C., Zhang, S. R., Zhang, B. C., McCrea, I. W., . . . Ruohoniemi, J. 706 M. (2015). Direct observations of the full Dungey convection cycle in the polar ionosphere for 707 southward interplanetary magnetic field conditions. Journal of Geophysical Research-Space 708 Physics, 120(6), 4519-4530. doi:10.1002/2015ja021172 709 Aarons, J., Mullen, J. P., Whitney, H. E., Johnson, A. L., & Weber, E. J. (1981). Uhf Scintillation Activity 710 over Polar Latitudes. Geophysical Research Letters, 8(3), 277-280. doi:DOI 711 10.1029/GL008i003p00277

712