Monthly climatologies of zonal-mean and tidal winds in the thermosphere as observed by ICON/MIGHTI during April 2020-March 2022

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Abstract

Version 5 (v05) of the thermospheric wind data from the Michelson Interferometer for Global High-resolution Thermospheric Imaging (MIGHTI) instrument on the Ionospheric CONnections (ICON) mission has been recently released, which largely avoids local-time dependent artificial baseline drifts that are found in previous versions of the ICON/MIGHTI wind data. This paper describes monthly climatologies of zonal-mean winds and tides based on the v05 ICON/MIGHTI data under geomagnetically quiet conditions (Hp30 < 30) during April 2020-March 2022. Green-line winds in the lower thermosphere (90-110 km) and red-line winds in the middle thermosphere (200-300 km) are analyzed, as these data cover both daytime and nighttime. The altitude and latitude structures of zonal-mean winds and tides are presented for each month, and the results are compared with the widely-used empirical model, Horizontal Wind Model 2014 (HWM14). The v05 wind retrieval algorithm does not involve HWM14. The ICON/MIGHTI and HWM14 results are in general agreement, providing a validation of the v05 ICON/MIGHTI data. The agreement is especially good for the zonal-mean winds. The tidal amplitudes in HWM14 are often too small compared with those from ICON/MIGHTI as well as previous studies. A more accurate description of tides in the thermosphere is key to the future improvement of HWM.

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12 Key Points:

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13	- Monthly climatologies of zonal-mean winds and tides at 90–110 km and 200–300 $$
14	km are determined using v05 ICON/MIGHTI observations.
15	• ICON/MIGHTI and HWM14 results are in general agreement, providing a val-
16	idation of the v05 ICON/MIGHTI data.
17	• HWM14 reproduces the zonal-mean winds well, but often underestimates tidal am-
18	plitude.

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19 Abstract

Version 5 (v05) of the thermospheric wind data from the Michelson Interferometer for 20 Global High-resolution Thermospheric Imaging (MIGHTI) instrument on the Ionospheric 21 Connection Explorer (ICON) mission has been recently released, which largely avoids 22 local-time dependent artificial baseline drifts that are found in previous versions of the 23 ICON/MIGHTI wind data. This paper describes monthly climatologies of zonal-mean 24 winds and tides based on the v05 ICON/MIGHTI data under geomagnetically quiet con-25 ditions (Hp30 < 30) during April 2020–March 2022. Green-line winds in the lower ther-26 mosphere (90–110 km) and red-line winds in the middle thermosphere (200–300 km) are 27 analyzed, as these data cover both daytime and nighttime. The altitude and latitude struc-28 tures of zonal-mean winds and tides are presented for each month, and the results are 29 compared with the widely-used empirical model, Horizontal Wind Model 2014 (HWM14). 30 The v05 wind retrieval algorithm does not involve HWM14. The ICON/MIGHTI and 31 HWM14 results are in general agreement, providing a validation of the v05 ICON/MIGHTI 32 data. The agreement is especially good for the zonal-mean winds. The tidal amplitudes 33 in HWM14 are often too small compared with those from ICON/MIGHTI as well as pre-34 vious studies. A more accurate description of tides in the thermosphere is key to the fu-35 ture improvement of HWM. 36

37 1 Introduction

The uppermost layer of the Earth's atmosphere, the thermosphere, extends from 38 \sim 90 km up to \sim 600 km (e.g., Richmond, 1983; Kato, 2007). Early studies evaluated den-30 sities of the thermosphere based on the measurement of orbital decay of artificial satel-40 lites. Jacchia (1965) developed a global empirical model of thermospheric densities un-41 der the assumption of diffusive equilibrium. A by-product of the model was an estimate 42 of the global distribution of air pressure. Theoretical studies found that the model pres-43 sure provides useful information for evaluating the global wind system in the thermo-44 sphere (Geisler, 1967; Kohl & King, 1967). It has been demonstrated that global mo-45 tion of the air above approximately 150 km is primarily driven by solar-induced pres-46 sure gradients. That is, horizontal winds blow from the higher-temperature (and higher-47 pressure) dayside to the lower-temperature (and lower-pressure) nightside. On the other 48 hand, the motion of the air in the lower thermosphere (<150 km) is often dominated by 49 waves from the lower layers of the atmosphere. In particular, atmospheric tides (e.g., Lindzen 50

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⁵¹ & Chapman, 1969) are known to play an important role for the meteorology of the meso-⁵² sphere and lower thermosphere.

Theoretical models of the thermosphere were developed and used to explain how 53 solar heating, as well as Joule heating in the polar region, drives the global circulation 54 of the thermosphere under different seasonal conditions (e.g., Dickinson et al., 1975, 1977; 55 Roble et al., 1977; Fuller-Rowell & Rees, 1980, 1981). These early modeling studies led 56 to the development of upper atmosphere models that self-consistently couple the ther-57 mosphere and ionosphere (e.g., Roble et al., 1988; Richmond et al., 1992; Fuller-Rowell 58 et al., 1994). Thermospheric winds can have a significant impact on ionospheric dynam-59 ics (e.g., Rishbeth, 1998) and electrodynamics (e.g., Heelis, 2004), and thus are impor-60 tant for the accurate description of space weather. 61

There are several ways to observe thermospheric winds. For instance, wind veloc-62 ities can be measured using an accelerometer onboard a low-Earth-orbit satellite. Past 63 satellite missions like Dynamic Explorer 2 (DE2) (Spencer et al., 1982), CHAllenging Min-64 isatellite Payload (CHAMP) (H. Liu et al., 2006; Sutton et al., 2007), and Gravity Field 65 and Steady State Ocean Circulation Explorer (GOCE) (Doornbos et al., 2010; H. Liu 66 et al., 2016) provided global in-situ observations of thermospheric winds. Wind veloc-67 ities can also be measured with a sounding rocket, which can reach the thermosphere. 68 For example, the chemical release technique (e.g., Larsen, 2002; Pfaff et al., 2020) uses 69 measurements of trails of a chemical tracer released by a rocket to derive thermospheric 70 wind velocities. Moreover, optical measurements of Doppler shifts in airglow emissions, 71 such as the 557.7 nm $O(^{1}S)$ green line and the 630.0 nm $O(^{1}D)$ red line, have also been 72 used to observe thermospheric wind velocities from ground stations (e.g., Shiokawa et 73 al., 1999; Meriwether, 2006; Makela et al., 2012) as well as from satellites such as DE2 74 (Hays et al., 1981), Upper Atmosphere Research Satellite (UARS) (Hays et al., 1993) 75 and Thermosphere-Ionosphere-Mesosphere Energetics and Dynamics (TIMED) (Killeen 76 et al., 2006). Ground-based meteor radars can be used to measure wind velocities in the 77 mesosphere and lower thermosphere around 80–100 km (e.g., Hocking et al., 2001; Chau 78 et al., 2019). Thermospheric wind velocities at E-region and F-region heights can also 79 be estimated using incoherent scatter radar measurements of ionospheric parameters (e.g., 80 Salah & Holt, 1974; Harper, 1977). 81

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82	Global empirical models of thermospheric winds have been developed based on the
83	measurements obtained through the techniques mentioned above and others. The most
84	commonly used empirical model is the Horizontal Wind Model (HWM) series (e.g., Hedin
85	et al., 1991, 1996; Drob et al., 2008, 2015). HWM is constructed by fitting analytical func-
86	tions to a large volume of historical data. It predicts the zonal and meridional compo-
87	nents of the neutral wind velocity at a given location (latitude, longitude and altitude)
88	and time (day of year and UT). The latest version is HWM14 (Drob et al., 2015), and
89	since its release, the model has been widely used in the space physics community. The
90	validation of HWM is a community effort. Thermospheric wind measurements are of-
91	ten compared against HWM for a validation of the observational data as well as for a
92	performance evaluation of HWM (e.g., Englert et al., 2012; Jiang et al., 2018; Li et al.,
93	2021; Tang et al., 2021; Okoh et al., 2022). The present study shows comparisons of HWM14 $$
94	with thermospheric wind observations from the Ionospheric Connection Explorer (ICON)
95	mission, which was launched on October 11, 2019 (Immel et al., 2018).

The Michelson Interferometer for Global High-Resolution Thermospheric Imaging 96 (MIGHTI) instrument onboard ICON measures the horizontal wind velocity by observ-97 ing Doppler shifts of the atomic oxygen airglow emissions (e.g., Englert et al., 2017; Hard-98 ing et al., 2017). The green-line wind measurements extend from an altitude of 90 km 99 to 300 km during the daytime but to only 110 km at night, as the strength of the green-100 line emission varies considerably from day to night. The red-line wind data cover the height 101 range approximately 160–300 km during day and 200–300 km at night. These wind data 102 are useful not only for studying the neutral dynamics of the thermosphere (e.g., He et 103 al., 2021; Cullens et al., 2020; Yiğit et al., 2022; Forbes et al., 2022; Englert et al., 2017; 104 Triplett et al., 2023) but also for investigating atmosphere-ionosphere coupling processes, 105 which can be realized by combining the ICON/MIGHTI wind data with ionospheric mea-106 surements made by ICON (e.g., England et al., 2021; Immel et al., 2021; Forbes et al., 107 2021; Park et al., 2021; Heelis et al., 2022; R. Zhang et al., 2022) or by other missions 108 (e.g., Gasperini et al., 2021, 2022; G. Liu et al., 2021; Yamazaki et al., 2021; Yamazaki, 109 Arras, et al., 2022; Aa et al., 2022; Le et al., 2022; Harding et al., 2022; Oberheide, 2022). 110

The studies mentioned above used version 4 (v04) or an earlier version of the ICON/MIGHTI

- wind data. The v04 wind data, especially during the early period of the mission, showed
- reasonable agreement with other independent observations (e.g., Harding et al., 2021;
- ¹¹⁴ Makela et al., 2021; Dhadly et al., 2021; Chen et al., 2022). However, later it became clear

that the baseline of the v04 data has slowly drifted over time, leading to errors of 50-115 100 m/s for some cases in 2021. This issue was described in detail by Englert et al. (2023). 116 The baseline drift was found to be dependent on the local time and height, which has 117 made the reliable assessment of zonal-mean winds and tides difficult. Version 5 (v05) of 118 the ICON/MIGHTI wind data has been recently (in November 2022) released. A new 119 calibration method for the so-called "zero wind" has been developed for v05, which uses 120 a long-term comparison of the ascending- and descending-orbit data to perform a self-121 calibration of the zero baseline, independent of external data or models (Englert et al., 122 2023). The present study evaluates, for the first time, zonal-mean winds and tides us-123 ing the v05 ICON/MIGHTI wind data for the height ranges 90–110 km and 200–300 km, 124 where wind measurements are made during both day and night. Monthly climatologies 125 derived from the ICON/MIGHTI observations during April 2020–March 2022 are com-126 pared with HWM14 predictions. 127

¹²⁸ 2 Method to Determine Zonal-mean Winds and Tides

The v05 ICON/MIGHTI wind data (level 2.2, cardinal vector winds) during the 129 24-month period from April 2020 to March 2022 are analyzed to determine zonal-mean 130 winds and tides. The estimated accuracy of the v05 wind data is generally 10-25 m/s 131 (Englert et al., 2023). Only the data that are flagged as "Good" (Wind-Quality = 1) are 132 used. This largely eliminates (1) the observations from the South Atlantic Anomaly where 133 the retrieval of wind velocities is difficult due to increased radiation, (2) the observations 134 with little airglow signal, and (3) the observations from the day-night terminators where 135 mode changes of the instrument take place. Cullens et al. (2020), using synthetic data 136 sampled along the ICON/MIGHTI measurement points, demonstrated that (3) does not 137 have a large impact on the estimation of tidal amplitude. 138

We use only the measurements made during geomagnetically quiet periods. Our 139 criterion for the geomagnetically quiet periods is Hp30 < 30, where Hp30 is the geomag-140 netic activity index described by Yamazaki, Matzka, et al. (2022). Briefly, Hp30 is a plan-141 etary geomagnetic activity index, similar to Kp (Matzka et al., 2021) but with a higher 142 temporal resolution of 30 minutes in contrast to the 3-hourly Kp index. The higher tem-143 poral resolution has an advantage in accurately selecting quiet-time data. Hp30 is pro-144 duced at the GeoForschungsZentrum (GFZ) Potsdam and distributed at their website: 145 https://kp.gfz-potsdam.de/en/hp30-hp60. 146

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The green-line winds are given at every ~ 3 km for the height range 91–112 km, while 147 the red-line winds are given at every ~ 10 km for 203–301 km. At each height, the data 148 were binned in hourly UT bins, in 5° latitude bins every 2.5° latitude from 10° S to 40° N, 149 and in 15° longitude bins every 15° longitude. This was done separately for each month 150 of the year (but without distinction of different years) and for the zonal and meridional 151 components of the wind. The mean value and standard deviation were computed for each 152 bin. The standard deviation is used, in a later step, to evaluate $1-\sigma$ uncertainties in zonal-153 mean winds and tides. The bin-mean values at given latitude and height were expressed 154 as a function of UT (t in hours), longitude (λ in degrees), and month (M=1, 2, ..., 12) 155 using the following analytical representation: 156

$$\sum_{n=0}^{4} \sum_{s=-4}^{4} \sum_{m=0}^{3} \left\{ a_{nsm} \cos\left(n\frac{t}{24} - s\frac{\lambda}{360} + m\frac{M}{12}\right) + b_{nsm} \sin\left(n\frac{t}{24} - s\frac{\lambda}{360} + m\frac{M}{12}\right) \right\}.$$
 (1)

Formula (1) takes into account zonal-mean winds, tides and stationary planetary waves, 158 and their seasonal variations. n represents the tidal frequency. That is, n=1, 2, 3, 4 cor-159 respond to the 24-h (or diurnal) tide, 12-h (or semidiurnal) tide, 8-h (or terdiurnal) tide 160 and 6-h tide, respectively. The higher order tides are generally not as important in the 161 thermosphere (e.g., Oberheide et al., 2011). |s| denotes the zonal wavenumber, and the 162 sign of s indicates the direction of the zonal propagation of tides. That is, s>0 and s<0163 correspond to eastward- and westward-propagating tides, respectively. The standard tidal 164 nomenclature is used throughout this paper, such as DE3 and SW2, where the first let-165 ter indicates the period (i.e., "D" for diurnal and "S" for semidiurnal), the second let-166 ter represents the propagation direction (i.e., "E" for eastward and "W" for westward), 167 and the last number is the zonal wavenumber |s|. Going back to formula (1), the zonal-168 mean winds are represented by the terms with n=0 and s=0, while stationary planetary 169 waves are represented by the terms with n=0 and |s|>0. The seasonal variations of the 170 zonal-mean winds, tides and stationary planetary waves are represented by m=1, 2, 3,171 corresponding to the annual, semiannual and terannual cycles. The coefficients a_{nsm} and 172 b_{nsm} were determined in such a way that the deviation of formula (1) from the binned 173 values of the ICON/MIGHTI data will be the smallest in a least-squares sense. 174

The goodness-of-fit was evaluated using two statistical metrics. One is the correlation coefficient between the observations (X) and fit (Y):

$$r = \frac{Cov(X,Y)}{\sqrt{Cov(X,X)Cov(Y,Y)}},$$
(2)

where Cov is the covariance. The other is the root-mean-square error:

$$RMS = \sqrt{\frac{\sum (X - Y)^2}{N}},$$
(3)

where N is the number of the observations. The latitude and height distributions of the 180 correlation coefficient and root-mean-square error are presented in Figure 1. The cor-181 relation coefficient is generally higher in the middle thermosphere (r=0.8-1.0, based on 182 red-line winds) than in the lower thermosphere (r=0.65-0.85, based on green-line winds). 183 This is mainly due to the fact that the lower thermosphere is more strongly influenced 184 by the waves that are not described by formula (1) such as acoustic waves, gravity waves, 185 lunar tides, Kelvin waves, and Rossby waves (e.g., Yiğit & Medvedev, 2015; H.-L. Liu, 186 2016). These waves are generated in the lower layers of the atmosphere and propagate 187 into the thermosphere. They can interact with tides and other waves to produce secondary 188 waves, which makes the spatial temporal variability of the lower thermosphere rather com-189 plex (e.g., Chang et al., 2011; H.-L. Liu, 2014; Nystrom et al., 2018). The waves from 190 the lower atmosphere get strongly dissipated before reaching the middle thermosphere. 191 The middle thermosphere is dominated by the diurnal tide that is locally generated by 192 solar heating (Hagan et al., 2001), which can be represented well by formula (1). RMS 193 is somewhat larger in the middle thermosphere (20-30 m/s) than in the lower thermo-194 sphere (15-25 m/s). This reflects generally larger wind velocities in the middle thermo-195 sphere. These RMS values are much smaller than those reported for HWM14 (40–80 m/s) 196 by Drob et al. (2015). This is not surprising given that HWM14 involves more diverse 197 sources of data from many different years with various degrees of accuracy. 198

199 200

179

For a given month M, formula (1) can be rewritten in the following form, which more explicitly represents the zonal-mean winds and waves:

$$\overline{A} + \sum_{n=1}^{4} \sum_{s=-4}^{4} A_{ns} \cos\left(n\frac{t}{24} - s\frac{\lambda}{360} + P_{ns}\right) + \sum_{s=1}^{4} A'_s \cos\left(s\frac{\lambda}{360} + P'_s\right).$$
(4)

Here, \overline{A} is the zonal-mean wind velocity (in m/s), A_{ns} and P_{ns} are the amplitude (in m/s) and phase (in rad) of a tide, respectively. A'_s and P'_s are the amplitude and phase of a stationary planetary wave, respectively. 1- σ uncertainties in the zonal-mean winds, tides and stationary planetary waves were evaluated using the standard deviation obtained during the binning procedure described earlier. A Monte Carlo method was used for this purpose. That is, random noise was generated for each bin based on the standard deviation, and the noise was superimposed on the corresponding mean value. Fitting of



Figure 1. Correlation coefficient and root-mean-square error (RMS), as measures of goodnessof-fit of formula (1) to the v05 ICON/MIGHTI green-line data (top four panels) and red-line data (bottom four panels). The left and right panels are for the zonal and meridional winds, respectively.

formula (1) was repeated for 250 Monte Carlo samples, and 1- σ uncertainties were computed for the zonal-mean wind velocity, and the amplitude and phase of tides and stationary planetary waves at each latitude and height. The derived 1- σ uncertainty in the zonal-mean wind velocity is typically 1.0–3.5 m/s for both green-line and red-line winds. The 1- σ uncertainties in the amplitude and phase of tides and stationary planetary waves are typically less than 4.5 m/s and 20°, respectively. These uncertainty values are appreciably smaller compared to the features discussed in this paper.

Zonal-mean winds, tides and stationary planetary waves were also evaluated using HWM14 for the purpose of comparison. Hourly values of the zonal and meridional wind velocities were derived from HWM14 for each month by running the model for the 15th day of the month without including disturbance winds (Emmert et al., 2008). At each latitude and height, the zonal-mean wind velocity, and the amplitude and phase of tides and stationary planetary waves were determined by least-squares fitting of formula (4), which can be directly compared with the ICON/MIGHTI results.

223 3 Results

First, we examine seasonal climatologies of zonal-mean winds. Figure 2 depicts the 224 zonal-mean zonal and meridional winds in the lower thermosphere (91–110 km) as de-225 rived from the ICON/MIGHTI green-line measurements. Below ~ 105 km, the zonal-mean 226 zonal wind in the equatorial region $(10^{\circ}S-10^{\circ}N)$ tends to be weakly westward through-227 out the year. An eastward jet can be seen at 30°N during the Northern Hemisphere (N.H.) 228 summer. The reversal of the zonal-mean zonal wind is often seen around 105 km, which 229 was also noted by Yiğit et al. (2022). The zonal-mean meridional wind is generally weak 230 with little seasonal variation. The corresponding results obtained from HWM14 are pre-231 sented in Figure 3. HWM14 captures the salient features of the observed zonal-mean zonal 232 and meridional winds well. 233

Figure 4 shows the zonal-mean zonal and meridional winds in the middle thermosphere (203–300 km) as derived from the ICON/MIGHTI red-line measurements. An annual variation of the zonal-mean zonal wind is evident. That is, the zonal wind in the N.H. is largely eastward and westward during the local winter and summer, respectively. The seasonal variation of the zonal-mean meridional wind is also dominated by an annual cycle. That is, the meridional wind is primarily northward during the N.H. winter

-9-



ICON/MIGHTI (green-line, zonal mean zonal wind, m/s)

Figure 2. Quiet-time monthly climatologies of the zonal-mean zonal wind (top 12 panels) and zonal-mean meridional wind (bottom 12 panels) in the lower thermopshere (91–110 km) as derived from the v05 ICON/MIGHTI green-line data during April 2020–March 2022.

-10

0 10 20 30 40

Latitude (°)

-10

0 10 20 30 40

Latitude (°)

-10 0 10 20 30 40

Latitude (°)

-10

0 10 20 30 40

Latitude (°)



Figure 3. Same as Figure 2 but from Horizontal Wind Model 2014 (HWM14).

and southward during the N.H. summer. The seasonal transitions occur in March and
September. The annual variations in the zonal-mean zonal and meridional winds are reproduced well by HWM14 as shown in Figure 5.

Next, we examine seasonal climatologies of tides and stationary planetary waves. 243 Different components of waves, as expressed by different combinations of (n, s), have vary-244 ing degrees of significance in the thermosphere (e.g., Truskowski et al., 2014; Forbes et 245 al., 2014). Figures 6 and 7 depict wave spectra for two representative altitudes. Figure 246 6 shows the amplitude of different wave components for the green-line winds over the 247 equator at an altitude of 106 km. At this particular latitude and height, the eastward-248 propagating diurnal tide with zonal wavenumber 3 (DE3; n=1, s=3) dominates the tide 249 in the zonal wind, especially during July–November, and the migrating semidiurnal tide 250 (SW2; n=2, s=-2) dominates the tide in the meridional wind, especially during April-251 September. Figure 7 is similar to Figure 6 but for the red-line winds at 30°N at an al-252 titude of 273 km. In the middle thermosphere, the migrating diurnal tide (DW1; n=1, 253 s=-1) is by far dominant. Since DW1, SW2 and DE3 are found to be dominant within 254 the latitudinal and altitudinal range of the ICON/MIGHTI wind measurement, we fur-255 ther analyze these specific tides. 256

Figure 8 shows the amplitude and phase of DW1 in the meridional wind in the lower 257 thermosphere as derived from the ICON/MIGHTI green-line measurements. The am-258 plitude is largest at $15-20^{\circ}$ N and 95-97 km, and it shows a semiannual variation with 259 equinoctial maxima of ~ 60 m/s. The phase of DW1 tends to decrease with increasing 260 height. This 'downward phase propagation' is a fundamental feature of upward-propagating 261 tides (e.g., Forbes, 1995). The results suggest that DW1 in the lower thermosphere orig-262 inate from lower layers of the atmosphere. The latitude-height pattern of the DW1 phase 263 does not vary much with the season. DW1 derived from HWM14 is presented in Figure 264 9. The DW1 amplitude in HWM14 is largest at $15-20^{\circ}$ N, which is in agreement with 265 the ICON/MIGHTI results. HWM14 also reproduces the semiannual variation in the 266 DW1 amplitude. However, the DW1 amplitude in HWM14 is generally too small, and 267 its height structure does not agree well with the observations. The latitude and height 268 structures of the DW1 phase are reproduced by HWM14 during equinoctial months. 269

Figure 10 presents the SW2 amplitude and phase in the meridional wind in the lower thermosphere as derived from the ICON/MIGHTI green-line measurements. The am-



ICON/MIGHTI (red-line, zonal mean zonal wind, m/s)

Figure 4. Quiet-time monthly climatologies of the zonal-mean zonal wind (top 12 panels) and zonal-mean meridional wind (bottom 12 panels) in the middle thermosphere (203-300 km) as derived from the v05 ICON/MIGHTI red-line data during April 2020–March 2022.

-10

0 10 20 30 40

Latitude (°)

-10

0 10 20 30 40

Latitude (°)

220

-10 0 10 20 30 40

Latitude (°)

-10 0 10 20 30 40

Latitude (°)



HWM14 (zonal mean zonal wind, m/s)

Figure 5. Same as Figure 4 but from Horizontal Wind Model 2014 (HWM14).



ICON/MIGHTI (green-line, tides, zonal wind amplitude, m/s, 106 km, equator)

ICON/MIGHTI (green-line, tides, meridional wind amplitude, m/s, 106 km, equator)



Figure 6. Amplitude of tides and stationary planetary waves in the zonal wind (top 12 panels) and meridional wind (bottom 12 panels) at 106 km at the equator as derived from the v05 ICON/MIGHTI green-line data. n represents tidal frequency. That is, n=1 for diurnal tides, n=2 for semidiurnal tides, and so on. n=0 for stationary planetary waves. s is the zonal wavenumber. s>0 for eastward-propagating waves, while s<0 for westward-propagating waves.



ICON/MIGHTI (red-line, tides, zonal wind amplitude, m/s, 273 km, 30°N)

ICON/MIGHTI (red-line, tides, meridional wind amplitude, m/s, 273 km, 30°N)



Figure 7. Same as Figure 6 but at 273 km at 30°N as derived from the v05 ICON/MIGHTI red-line data.



ICON/MIGHTI (green-line, DW1, meridional wind amplitude, m/s)





Figure 8. Amplitude (top 12 panels) and phase (bottom 12 panels) of the migrating diurnal tide (DW1) in the meridional wind in the lower thermosphere as derived from the v05 ICON/MIGHTI green-line data. The corresponding results for the zonal wind can be found in Figure S1 of Supporting Information.



HWM14 (DW1, meridional wind amplitude, m/s)

HWM14 (DW1, meridional wind phase, deg.)



Figure 9. Same as Figure 8 but from Horizontal Wind Model 2014 (HWM14).

plitude is relatively large over the equator $(10^{\circ}S-10^{\circ}N)$ and at N.H. middle latitudes $(>30^{\circ}N)$. 272 In the equatorial region, the amplitude grows with height, reaching 60 m/s at 110 km273 during August–September. The maximum amplitude probably occurs above 110 km. At 274 middle latitudes, the amplitude peaks at 105 km. The downward phase propagation is 275 seen at both equatorial and middle-latitude regions, indicating that the SW2 energy prop-276 agates upward at these heights. The corresponding results derived from HWM14 are shown 277 in Figure 11. Again, the amplitude in HWM14 is generally too small, and its height struc-278 ture does not agree well with the observations. Interestingly, there is remarkable agree-279 ment in the phase of SW2 in the lower thermosphere between the ICON/MIGHTI and 280 HWM14 results. 281

Figure 12 shows the amplitude and phase of DE3 in the zonal wind in the lower thermosphere as derived from the ICON/MIGHTI green-line measurements. DE3 is the largest non-migrating (i.e., non-sun-synchronous) tidal component found in the greenline data. The zonal-wind amplitude is largest over the equator at a height of 105–110 km. The maximum amplitude exceeds 30 m/s during July–October. The downward phase propagation is visible, indicating upward energy propagation of DE3. DE3 is nonexistent in HWM14, as the model does not take into account any non-migrating tide.

We now look at DW1 in the middle thermosphere. Figure 13 shows the amplitude 289 and phase of DW1 in the zonal wind in the middle thermosphere as derived from the ICON/MIGHTI 290 red-line observations. It is noted that the scale range for the amplitude is different from 291 those used for the green-line results (Figures 8, 10 and 12). The DW1 amplitude grows 292 with height from ${\sim}50$ m/s at 200 km to ${\sim}90$ m/s at 300 km. It exceeds 100 m/s in some 293 months. The phase does not vary with height, indicating that DW1 in the middle ther-294 mosphere is a vertically-trapped (evanescent) tidal mode that is locally generated, rather 295 than an upward-propagating mode from below. The corresponding results derived from 296 HWM14 are presented in Figure 14. HWM14 reproduces the latitude and height struc-297 tures of the amplitude and phase well. Figure 15 also shows the amplitude and phase 298 of DW1 from the ICON/MIGHTI red-line measurements, but for the meridional wind. 299 The amplitude is small over the equatorial region but can exceed 100 m/s at middle lat-300 itudes (>30°N) above 280 km. The phase depends strongly on latitude. The phase struc-301 ture is well captured by HWM14 (Figure 16), but the model severely underestimates the 302 DW1 amplitude at middle latitudes. 303



ICON/MIGHTI (green-line, SW2, meridional wind amplitude, m/s)





Figure 10. Amplitude (top 12 panels) and phase (bottom 12 panels) of the migrating semidiurnal tide (SW2) in the meridional wind in the lower thermosphere as derived from the v05 ICON/MIGHTI green-line data. The corresponding results for the zonal wind can be found in Figure S2 of Supporting Information.



HWM14 (SW2, meridional wind amplitude, m/s)

HWM14 (SW2, meridional wind phase, deg.)



Figure 11. Same as Figure 10 but from Horizontal Wind Model 2014 (HWM14).



ICON/MIGHTI (green-line, DE3, zonal wind amplitude, m/s)



-10 0 10 20 30 40

Latitude (°)

-10 0 10 20 30 40 Latitude (°)

-10 0 10 20 30 40

Latitude (°)

-10 0 10 20 30 40

Latitude (°)



ICON/MIGHTI (red-line, DW1, zonal wind amplitude, m/s)

Figure 13. Amplitude (top 12 panels) and phase (bottom 12 panels) of the migrating diurnal tide (DW1) in the zonal wind in the middle thermosphere as derived from the v05 ICON/MIGHTI red-line data. The corresponding results for the migrating semidiurnal tide (SW2) can be found in Figure S4 of Supporting Information.

Latitude (°)

Latitude (°)

Latitude (°)

Latitude (°)



HWM14 (DW1, zonal wind amplitude, m/s)

Figure 14. Same as Figure 13 but from Horizontal Wind Model 2014 (HWM14).



ICON/MIGHTI (red-line, DW1, meridional wind amplitude, m/s)

ICON/MIGHTI (red-line, DW1, meridional wind phase, deg.)



Figure 15. Same as Figure 13 but in the meridional wind. The corresponding results for the migrating semidiurnal tide (SW2) can be found in Figure S5 of Supporting Information.



HWM14 (DW1, meridional wind amplitude, m/s)

HWM14 (DW1, meridional wind phase, deg.)



Figure 16. Same as Figure 14 but in the meridional wind.

304 4 Discussion

We have presented seasonal climatologies of the zonal-mean winds and tides derived from the v05 ICON/MIGHTI data, and compared the results with those from HWM14. Here we compare the ICON/MIGHTI results with those presented in earlier work based on other observations and models. Also, we discuss physical mechanisms behind some of the features observed in the ICON/MIGHTI winds, referring to previous theoretical studies.

Wang et al. (1997) created an empirical model of lower thermospheric winds (90– 311 120 km) using the measurements from the wind imaging interferometer (WINDII; Shep-312 herd et al., 1993) onboard UARS. They presented the zonal-mean zonal and meridional 313 winds for different seasons, which can be compared with our ICON/MIGHTI results (Fig-314 ures 2 and 4). S. P. Zhang et al. (2007) later analyzed an updated version of UARS/WINDII 315 data and obtained similar results as Wang et al. (1997). The UARS/WINDII results showed 316 a westward jet of 10-30 m/s over the equator at ~ 100 km throughout the year. This is 317 also seen in the ICON/MIGHTI results (Figure 2), as well as in HWM14 (Figure 3). It 318 is noted that the UARS/WINDII data are already incorporated in HWM. The westward 319 jet over the equator is considered to result from the westward momentum deposition by 320 dissipating migrating (thus westward-propagating) tides (e.g., Miyahara, 1981, 1978; Lieber-321 man & Hays, 1994; Jones Jr et al., 2014). Wang et al. (1997) and S. P. Zhang et al. (2007) 322 noted that the equatorial westward jet is sandwiched by eastward jets centered around 323 $\pm 40^{\circ}$ latitudes. The eastward jets were reported to be stronger in the summer hemisphere, 324 with the magnitude of 30–40 m/s. The ICON/MIGHTI results (Figure 2) clearly cap-325 ture the N.H. part of the eastward jets. The mechanism for the middle-latitude eastward 326 jets is not well understood. The numerical work by Forbes et al. (1993) predicted that 327 the equatorial westward jet induced by tidal dissipation is accompanied by eastward jets 328 at higher latitudes. However, the eastward jets due to tidal dissipation are predicted to 329 be much weaker than the westward jet, which agrees with neither ICON/MIGHTI nor 330 UARS/WINDII observations. Besides tides, Miyoshi and Fujiwara (2006) numerically 331 demonstrated that the momentum deposition by eastward-propagating equatorial Kelvin 332 waves also plays a significant role for the zonal-mean zonal wind in the equatorial lower 333 thermosphere. More studies are required to determine the relative importance of differ-334 ent waves to explain the observed westward and eastward jets. 335

The zonal-mean meridional wind in the lower thermosphere as derived from the ICON/MIGHTI 336 data is generally weak (Figure 2), which is consistent with the UARS/WINDII results 337 presented by Wang et al. (1997) as well as HWM14 (Figure 3). S. P. Zhang et al. (2007) 338 noted that the zonal-mean meridional wind sometimes show a cell-like structure in the 339 low latitude region, which is characterized by poleward winds on both sides of the equa-340 tor at altitudes of 95–105 km and equatorward winds at 105–115 km. There is some hint 341 of such a cell-like structure in the ICON/MIGHTI zonal-mean meridional wind (see, e.g., 342 August–September), but it is not well resolved because of the small magnitude. The cell-343 like structure in the zonal-mean meridional wind in the lower thermosphere is sometimes 344 found in numerical models and is considered to be driven by tidal dissipation (e.g., Miya-345 hara et al., 1993; Forbes et al., 1993). 346

ICON/MIGHTI red-line measurements revealed seasonal climatologies of the zonal-347 mean zonal and meridional winds at 200–300 km (Figure 4). HWM14 reproduces the 348 ICON/MIGHTI observations well for both the zonal and meridional components (Fig-349 ure 5). HWM14 at this height range is well constrained by UARS/WINDII red-line winds 350 as well as observations by ground-based Fabry-Perot interferometers. The zonal-mean 351 meridional wind in the middle thermosphere is directed from the summer to the winter 352 hemisphere, which is not surprising given the higher temperature and pressure in the sum-353 mer hemisphere. The seasonal transition in the meridional circulation occurs in March 354 and September. Using a numerical model, Roble et al. (1977) showed that the seasonal 355 transition of the zonal-mean circulation takes place within a few weeks of equinox. Such 356 an abrupt seasonal transition is not fully resolved in our monthly analysis. The zonal-357 mean zonal wind in the middle thermosphere arises mainly from the correlation between 358 diurnal variations of pressure gradient and ion drag (Dickinson et al., 1975, 1977). That 359 is, the wind is weaker on the dayside than the nightside as the ion drag is larger on the 360 dayside due to higher plasma concentration. Since the wind in the middle thermosphere 361 undergoes a diurnal cycle due to day-night pressure differences, an unbalance between 362 the daytime and nighttime winds leads to the zonal-mean winds. 363

The three most dominant tidal components in the ICON/MIGHTI green-line winds are DW1, SW2 and DE3 (Figures 8, 10, 12). This is as expected from previous studies on tides in the lower thermosphere (e.g., Forbes et al., 2008; Oberheide et al., 2011). DW1 and SW2 are sun-synchronous, while DE3 is non-sun-synchronous. In the lower thermosphere, they consist mainly of upward-propagating modes, which can be seen from their

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downward phase propagation. They are driven by radiative heating through insolation 369 of H_2O in the troposphere and O_3 in the stratosphere (e.g., Forbes, 1982b, 1982a) as well 370 as by latent heating in the troposphere (Hagan & Forbes, 2002, 2003; X. Zhang et al., 371 2010a, 2010b). DW1 in the meridional wind, as derived from the ICON/MIGHTI ob-372 servations, shows an amplitude maximum at $15-20^{\circ}$ N at an altitude of 95-98 km (Fig-373 ure 8). The results are consistent with those from UARS/WINDII (McLandress et al., 374 1996; S. P. Zhang et al., 2007) and TIMED/TIDI (Wu et al., 2008a). The amplitude is 375 larger during the equinoxes than the solstices, which is well known from previous stud-376 ies (e.g., Burrage et al., 1995; Xu et al., 2009). McLandress (2002b, 2002a) examined the 377 mechanism for the semiannual variation of DW1 using a numerical model, and concluded 378 that the change in the latitudinal shear of the zonal-mean zonal wind plays a leading role 379 for the seasonal variation of DW1 in the lower thermosphere. HWM14 reproduces the 380 semiannual variation of DW1 (Figure 9) but the model underestimates the amplitude 381 in comparison not only with the ICON/MIGHTI results but also with the UARS/WINDII 382 and TIMED/TIDI results (S. P. Zhang et al., 2007; Wu et al., 2008a). Previous stud-383 ies reported that the amplitude of DW1 at low latitudes can change by a few tens of m/s 384 from one year to the next (e.g., Burrage et al., 1995; Hagan et al., 1999). Variation as-385 sociated with the quasi-biennial oscillation (QBO) of the equatorial atmosphere is an im-386 portant part of the interannual variation of DW1 in the lower thermosphere, account-387 ing for up to 10 m/s (e.g., Xu et al., 2009). The interannual variability of tides is not 388 taken into account in the present study. Resolving the QBO effect would require a larger 389 data set. 390

SW2 in the meridional wind, as derived from the ICON/MIGHTI observations, is 391 relatively strong during May–September (Figure 10), which is consistent with the UARS/WINDII 392 observations (S. P. Zhang et al., 2007). HWM14 reproduces the seasonal variation of SW2 393 but with somewhat smaller amplitude (Figure 11). The mechanism for the seasonal vari-394 ation of SW2 is not well established. DE3 in the lower thermosphere has characteristics 395 of a Kelvin wave (e.g., Forbes et al., 2003). In classical theory, a Kelvin wave travels east-396 ward, and its zonal wind component has a Gaussian-shaped latitudinal profile with max-397 imum amplitude over the equator (e.g., Forbes, 2000). The latitude and height struc-398 tures of DE3 and its seasonal variation in the ICON/MIGHTI green-line zonal wind (Fig-399 ure 12) are consistent with those from the UARS/WINDII (Forbes et al., 2003) and TIMED/TIDI 400 observations (Oberheide et al., 2006; Wu et al., 2008b). As the zonal wind amplitude of 401

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DE3 reaches its maximum in the equatorial dynamo region at 105–110 km, it has a significant impact on the equatorial zonal electric field and current (e.g., England et al., 2006;
Fejer et al., 2008) as well as on the F-region plasma concentration (e.g., Immel et al.,
2006; Lin et al., 2007). Despite the importance of DE3 in low-latitude ionosphere-thermosphere

coupling, it is not included in HWM14 like other non-migrating tides.

DW1 in the middle thermosphere (Figures 13 and 15) is predominantly a vertically-407 trapped tidal mode that is excited by in-situ solar heating (e.g., Forbes, 1982b; Hagan 408 et al., 2001). This contrasts with DW1 in the lower thermosphere (Figure 8), which is 409 primarily an upward-propagating mode. The latitude and height structures of DW1 in 410 the middle thermosphere are not well documented, particularly those based on obser-411 vations. The simulation results by Hagan et al. (2001) showed that (1) the amplitude 412 of DW1 at 200–300 km grows with height at all latitudes, (2) both zonal and meridional 413 wind amplitudes are largest at high latitudes, (3) the meridional wind amplitude is van-414 ishingly small over the equator but it increases with latitude, (4) the zonal wind ampli-415 tude does not depend strongly on latitude over the middle- and low-latitude regions, (5) 416 both zonal and meridional wind phases do not depend strongly on height, (6) the zonal 417 wind phase does not vary strongly with latitude, and (7) the meridional wind phase also 418 does not vary strongly with latitude except that the phase reversal occurs at the equa-419 tor. The ICON/MIGHTI results (Figures 13 and 15) are consistent with these numer-420 ical predictions. 421

Some previous studies have addressed a potential impact of the solar flux, mainly 422 at the wavelengths of extreme ultraviolet (EUV), on neutral winds in the middle and up-423 per thermosphere (e.g., Hedin et al., 1994). The ICON/MIGHTI observations examined 424 in this paper are obtained during the period April 2020–March 2022. The mean value 425 of the $F_{10.7}$ index (Tapping, 2013), which is often used as a proxy of the EUV flux, was 426 82.8 sfu (1 sfu = 10^{-22} W·m⁻²·Hz⁻¹), with the minimum and maximum monthly val-427 ues of 69.2 sft in May 2020 and 117.8 sfu in March 2022, respectively. We have ignored 428 possible variations in the wind velocities associated with the change in the solar flux, as 429 the ICON/MIGHTI data used in this study are not sufficient for evaluating the solar ac-430 tivity effect on the thermospheric winds. HWM14 also does not take into account the 431 dependence of wind velocities on solar activity. Hedin et al. (1994) reported that although 432 the solar flux significantly influences the temperature, the zonal-mean winds in the mid-433 dle thermosphere do not strongly depend on solar activity. Hagan et al. (2001) noted in 434

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their simulation results that the solar flux has a marked effect on the temperature am-435 plitude of DW1 in the middle thermosphere but not on the wind amplitudes. H. Liu et 436 al. (2006) examined the effect of the solar flux on thermospheric winds (\sim 400 km) us-437 ing the CHAMP accelerometer data. They found that the solar flux effect can be sig-438 nificant depending on the season and local time. The solar flux effect on tides is gener-439 ally small below about 130 km, where tidal waves are mainly of lower atmospheric ori-440 gin (Oberheide et al., 2009; Dhadly et al., 2018). More observational studies are required 441 to establish the solar activity dependence of zonal-mean winds and tides in the thermo-442 443 sphere.

444 5 Conclusions

Monthly climatologies of quiet-time zonal-mean winds and tides are derived using the recently-released v05 of the ICON/MIGHTI thermopsheric wind measurements during April 2020–March 2022 at the altitude ranges 90–110 km and 200–300 km. Earlier versions of the ICON/MIGHTI wind data suffered from artificial baseline drifts that depend on local time. Thus, it was previously difficult to obtain reliable climatological estimates of zonal-mean winds and tides. The v05 data avoids this issue by the use of a renewed baseline calibration technique (Englert et al., 2023).

The ICON/MIGHTI results are compared with those from the latest version of HWM 452 (i.e., HWM14) as well as previous studies. Salient features of zonal-mean winds and tides 453 in the lower and middle thermosphere are in general agreement between ICON/MIGHTI 454 and HWM14, including latitude and height structures and their seasonal variations. This 455 provides a validation of the v05 ICON/MIGHTI data. HWM14 reproduces the zonal-456 mean zonal and meridional winds well in both the lower and middle thermosphere. How-457 ever, HWM14 tends to underestimate tidal amplitude. Also, HWM14 does not include 458 non-migrating tides such as DE3, which is especially important in the equatorial lower 459 thermosphere. The latitude and height structures of DE3 and their seasonal variations 460 in the ICON/MIGHTI green-line zonal wind are found to be consistent with those from 461 the UARS/WINDII and TIMED/TIDI observations. The future improvement of HWM 462 can benefit from the inclusion of the ICON/MIGHTI winds for better description of tides. 463

464 Open Research Section

The ICON/MIGHTI Level 2.2 product Cardinal Vector Winds (Version 5) is accessible from the ICON website https://icon.ssl.berkeley.edu/Data. The Hpo indices including Hp30 used in this study are available at the GFZ website https://kp .gfz-potsdam.de/en/hp30-hp60/data; see also data publication Matzka et al. (2022). The monthly F10.7 index is available at the website of the Canadian Space Weather Forecast Centre https://spaceweather.gc.ca/forecast-prevision/solar-solaire/solarflux/ sx-5-mavg-en.php.

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Monthly climatologies of zonal-mean and tidal winds in the thermosphere as observed by ICON/MIGHTI during April 2020–March 2022

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12 Key Points:

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13	- Monthly climatologies of zonal-mean winds and tides at 90–110 km and 200–300 $$
14	km are determined using v05 ICON/MIGHTI observations.
15	• ICON/MIGHTI and HWM14 results are in general agreement, providing a val-
16	idation of the v05 ICON/MIGHTI data.
17	• HWM14 reproduces the zonal-mean winds well, but often underestimates tidal am-
18	plitude.

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19 Abstract

Version 5 (v05) of the thermospheric wind data from the Michelson Interferometer for 20 Global High-resolution Thermospheric Imaging (MIGHTI) instrument on the Ionospheric 21 Connection Explorer (ICON) mission has been recently released, which largely avoids 22 local-time dependent artificial baseline drifts that are found in previous versions of the 23 ICON/MIGHTI wind data. This paper describes monthly climatologies of zonal-mean 24 winds and tides based on the v05 ICON/MIGHTI data under geomagnetically quiet con-25 ditions (Hp30 < 30) during April 2020–March 2022. Green-line winds in the lower ther-26 mosphere (90–110 km) and red-line winds in the middle thermosphere (200–300 km) are 27 analyzed, as these data cover both daytime and nighttime. The altitude and latitude struc-28 tures of zonal-mean winds and tides are presented for each month, and the results are 29 compared with the widely-used empirical model, Horizontal Wind Model 2014 (HWM14). 30 The v05 wind retrieval algorithm does not involve HWM14. The ICON/MIGHTI and 31 HWM14 results are in general agreement, providing a validation of the v05 ICON/MIGHTI 32 data. The agreement is especially good for the zonal-mean winds. The tidal amplitudes 33 in HWM14 are often too small compared with those from ICON/MIGHTI as well as pre-34 vious studies. A more accurate description of tides in the thermosphere is key to the fu-35 ture improvement of HWM. 36

37 1 Introduction

The uppermost layer of the Earth's atmosphere, the thermosphere, extends from 38 \sim 90 km up to \sim 600 km (e.g., Richmond, 1983; Kato, 2007). Early studies evaluated den-30 sities of the thermosphere based on the measurement of orbital decay of artificial satel-40 lites. Jacchia (1965) developed a global empirical model of thermospheric densities un-41 der the assumption of diffusive equilibrium. A by-product of the model was an estimate 42 of the global distribution of air pressure. Theoretical studies found that the model pres-43 sure provides useful information for evaluating the global wind system in the thermo-44 sphere (Geisler, 1967; Kohl & King, 1967). It has been demonstrated that global mo-45 tion of the air above approximately 150 km is primarily driven by solar-induced pres-46 sure gradients. That is, horizontal winds blow from the higher-temperature (and higher-47 pressure) dayside to the lower-temperature (and lower-pressure) nightside. On the other 48 hand, the motion of the air in the lower thermosphere (<150 km) is often dominated by 49 waves from the lower layers of the atmosphere. In particular, atmospheric tides (e.g., Lindzen 50

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⁵¹ & Chapman, 1969) are known to play an important role for the meteorology of the meso-⁵² sphere and lower thermosphere.

Theoretical models of the thermosphere were developed and used to explain how 53 solar heating, as well as Joule heating in the polar region, drives the global circulation 54 of the thermosphere under different seasonal conditions (e.g., Dickinson et al., 1975, 1977; 55 Roble et al., 1977; Fuller-Rowell & Rees, 1980, 1981). These early modeling studies led 56 to the development of upper atmosphere models that self-consistently couple the ther-57 mosphere and ionosphere (e.g., Roble et al., 1988; Richmond et al., 1992; Fuller-Rowell 58 et al., 1994). Thermospheric winds can have a significant impact on ionospheric dynam-59 ics (e.g., Rishbeth, 1998) and electrodynamics (e.g., Heelis, 2004), and thus are impor-60 tant for the accurate description of space weather. 61

There are several ways to observe thermospheric winds. For instance, wind veloc-62 ities can be measured using an accelerometer onboard a low-Earth-orbit satellite. Past 63 satellite missions like Dynamic Explorer 2 (DE2) (Spencer et al., 1982), CHAllenging Min-64 isatellite Payload (CHAMP) (H. Liu et al., 2006; Sutton et al., 2007), and Gravity Field 65 and Steady State Ocean Circulation Explorer (GOCE) (Doornbos et al., 2010; H. Liu 66 et al., 2016) provided global in-situ observations of thermospheric winds. Wind veloc-67 ities can also be measured with a sounding rocket, which can reach the thermosphere. 68 For example, the chemical release technique (e.g., Larsen, 2002; Pfaff et al., 2020) uses 69 measurements of trails of a chemical tracer released by a rocket to derive thermospheric 70 wind velocities. Moreover, optical measurements of Doppler shifts in airglow emissions, 71 such as the 557.7 nm $O(^{1}S)$ green line and the 630.0 nm $O(^{1}D)$ red line, have also been 72 used to observe thermospheric wind velocities from ground stations (e.g., Shiokawa et 73 al., 1999; Meriwether, 2006; Makela et al., 2012) as well as from satellites such as DE2 74 (Hays et al., 1981), Upper Atmosphere Research Satellite (UARS) (Hays et al., 1993) 75 and Thermosphere-Ionosphere-Mesosphere Energetics and Dynamics (TIMED) (Killeen 76 et al., 2006). Ground-based meteor radars can be used to measure wind velocities in the 77 mesosphere and lower thermosphere around 80–100 km (e.g., Hocking et al., 2001; Chau 78 et al., 2019). Thermospheric wind velocities at E-region and F-region heights can also 79 be estimated using incoherent scatter radar measurements of ionospheric parameters (e.g., 80 Salah & Holt, 1974; Harper, 1977). 81

-3-

82	Global empirical models of thermospheric winds have been developed based on the
83	measurements obtained through the techniques mentioned above and others. The most
84	commonly used empirical model is the Horizontal Wind Model (HWM) series (e.g., Hedin
85	et al., 1991, 1996; Drob et al., 2008, 2015). HWM is constructed by fitting analytical func-
86	tions to a large volume of historical data. It predicts the zonal and meridional compo-
87	nents of the neutral wind velocity at a given location (latitude, longitude and altitude)
88	and time (day of year and UT). The latest version is HWM14 (Drob et al., 2015), and
89	since its release, the model has been widely used in the space physics community. The
90	validation of HWM is a community effort. Thermospheric wind measurements are of-
91	ten compared against HWM for a validation of the observational data as well as for a
92	performance evaluation of HWM (e.g., Englert et al., 2012; Jiang et al., 2018; Li et al.,
93	2021; Tang et al., 2021; Okoh et al., 2022). The present study shows comparisons of HWM14 $$
94	with thermospheric wind observations from the Ionospheric Connection Explorer (ICON)
95	mission, which was launched on October 11, 2019 (Immel et al., 2018).

The Michelson Interferometer for Global High-Resolution Thermospheric Imaging 96 (MIGHTI) instrument onboard ICON measures the horizontal wind velocity by observ-97 ing Doppler shifts of the atomic oxygen airglow emissions (e.g., Englert et al., 2017; Hard-98 ing et al., 2017). The green-line wind measurements extend from an altitude of 90 km 99 to 300 km during the daytime but to only 110 km at night, as the strength of the green-100 line emission varies considerably from day to night. The red-line wind data cover the height 101 range approximately 160–300 km during day and 200–300 km at night. These wind data 102 are useful not only for studying the neutral dynamics of the thermosphere (e.g., He et 103 al., 2021; Cullens et al., 2020; Yiğit et al., 2022; Forbes et al., 2022; Englert et al., 2017; 104 Triplett et al., 2023) but also for investigating atmosphere-ionosphere coupling processes, 105 which can be realized by combining the ICON/MIGHTI wind data with ionospheric mea-106 surements made by ICON (e.g., England et al., 2021; Immel et al., 2021; Forbes et al., 107 2021; Park et al., 2021; Heelis et al., 2022; R. Zhang et al., 2022) or by other missions 108 (e.g., Gasperini et al., 2021, 2022; G. Liu et al., 2021; Yamazaki et al., 2021; Yamazaki, 109 Arras, et al., 2022; Aa et al., 2022; Le et al., 2022; Harding et al., 2022; Oberheide, 2022). 110

The studies mentioned above used version 4 (v04) or an earlier version of the ICON/MIGHTI

- wind data. The v04 wind data, especially during the early period of the mission, showed
- reasonable agreement with other independent observations (e.g., Harding et al., 2021;
- ¹¹⁴ Makela et al., 2021; Dhadly et al., 2021; Chen et al., 2022). However, later it became clear

that the baseline of the v04 data has slowly drifted over time, leading to errors of 50-115 100 m/s for some cases in 2021. This issue was described in detail by Englert et al. (2023). 116 The baseline drift was found to be dependent on the local time and height, which has 117 made the reliable assessment of zonal-mean winds and tides difficult. Version 5 (v05) of 118 the ICON/MIGHTI wind data has been recently (in November 2022) released. A new 119 calibration method for the so-called "zero wind" has been developed for v05, which uses 120 a long-term comparison of the ascending- and descending-orbit data to perform a self-121 calibration of the zero baseline, independent of external data or models (Englert et al., 122 2023). The present study evaluates, for the first time, zonal-mean winds and tides us-123 ing the v05 ICON/MIGHTI wind data for the height ranges 90–110 km and 200–300 km, 124 where wind measurements are made during both day and night. Monthly climatologies 125 derived from the ICON/MIGHTI observations during April 2020–March 2022 are com-126 pared with HWM14 predictions. 127

¹²⁸ 2 Method to Determine Zonal-mean Winds and Tides

The v05 ICON/MIGHTI wind data (level 2.2, cardinal vector winds) during the 129 24-month period from April 2020 to March 2022 are analyzed to determine zonal-mean 130 winds and tides. The estimated accuracy of the v05 wind data is generally 10-25 m/s 131 (Englert et al., 2023). Only the data that are flagged as "Good" (Wind-Quality = 1) are 132 used. This largely eliminates (1) the observations from the South Atlantic Anomaly where 133 the retrieval of wind velocities is difficult due to increased radiation, (2) the observations 134 with little airglow signal, and (3) the observations from the day-night terminators where 135 mode changes of the instrument take place. Cullens et al. (2020), using synthetic data 136 sampled along the ICON/MIGHTI measurement points, demonstrated that (3) does not 137 have a large impact on the estimation of tidal amplitude. 138

We use only the measurements made during geomagnetically quiet periods. Our 139 criterion for the geomagnetically quiet periods is Hp30 < 30, where Hp30 is the geomag-140 netic activity index described by Yamazaki, Matzka, et al. (2022). Briefly, Hp30 is a plan-141 etary geomagnetic activity index, similar to Kp (Matzka et al., 2021) but with a higher 142 temporal resolution of 30 minutes in contrast to the 3-hourly Kp index. The higher tem-143 poral resolution has an advantage in accurately selecting quiet-time data. Hp30 is pro-144 duced at the GeoForschungsZentrum (GFZ) Potsdam and distributed at their website: 145 https://kp.gfz-potsdam.de/en/hp30-hp60. 146

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The green-line winds are given at every ~ 3 km for the height range 91–112 km, while 147 the red-line winds are given at every ~ 10 km for 203–301 km. At each height, the data 148 were binned in hourly UT bins, in 5° latitude bins every 2.5° latitude from 10° S to 40° N, 149 and in 15° longitude bins every 15° longitude. This was done separately for each month 150 of the year (but without distinction of different years) and for the zonal and meridional 151 components of the wind. The mean value and standard deviation were computed for each 152 bin. The standard deviation is used, in a later step, to evaluate $1-\sigma$ uncertainties in zonal-153 mean winds and tides. The bin-mean values at given latitude and height were expressed 154 as a function of UT (t in hours), longitude (λ in degrees), and month (M=1, 2, ..., 12) 155 using the following analytical representation: 156

$$\sum_{n=0}^{4} \sum_{s=-4}^{4} \sum_{m=0}^{3} \left\{ a_{nsm} \cos\left(n\frac{t}{24} - s\frac{\lambda}{360} + m\frac{M}{12}\right) + b_{nsm} \sin\left(n\frac{t}{24} - s\frac{\lambda}{360} + m\frac{M}{12}\right) \right\}.$$
 (1)

Formula (1) takes into account zonal-mean winds, tides and stationary planetary waves, 158 and their seasonal variations. n represents the tidal frequency. That is, n=1, 2, 3, 4 cor-159 respond to the 24-h (or diurnal) tide, 12-h (or semidiurnal) tide, 8-h (or terdiurnal) tide 160 and 6-h tide, respectively. The higher order tides are generally not as important in the 161 thermosphere (e.g., Oberheide et al., 2011). |s| denotes the zonal wavenumber, and the 162 sign of s indicates the direction of the zonal propagation of tides. That is, s>0 and s<0163 correspond to eastward- and westward-propagating tides, respectively. The standard tidal 164 nomenclature is used throughout this paper, such as DE3 and SW2, where the first let-165 ter indicates the period (i.e., "D" for diurnal and "S" for semidiurnal), the second let-166 ter represents the propagation direction (i.e., "E" for eastward and "W" for westward), 167 and the last number is the zonal wavenumber |s|. Going back to formula (1), the zonal-168 mean winds are represented by the terms with n=0 and s=0, while stationary planetary 169 waves are represented by the terms with n=0 and |s|>0. The seasonal variations of the 170 zonal-mean winds, tides and stationary planetary waves are represented by m=1, 2, 3,171 corresponding to the annual, semiannual and terannual cycles. The coefficients a_{nsm} and 172 b_{nsm} were determined in such a way that the deviation of formula (1) from the binned 173 values of the ICON/MIGHTI data will be the smallest in a least-squares sense. 174

The goodness-of-fit was evaluated using two statistical metrics. One is the correlation coefficient between the observations (X) and fit (Y):

$$r = \frac{Cov(X,Y)}{\sqrt{Cov(X,X)Cov(Y,Y)}},$$
(2)

where Cov is the covariance. The other is the root-mean-square error:

$$RMS = \sqrt{\frac{\sum (X - Y)^2}{N}},$$
(3)

where N is the number of the observations. The latitude and height distributions of the 180 correlation coefficient and root-mean-square error are presented in Figure 1. The cor-181 relation coefficient is generally higher in the middle thermosphere (r=0.8-1.0, based on 182 red-line winds) than in the lower thermosphere (r=0.65-0.85, based on green-line winds). 183 This is mainly due to the fact that the lower thermosphere is more strongly influenced 184 by the waves that are not described by formula (1) such as acoustic waves, gravity waves, 185 lunar tides, Kelvin waves, and Rossby waves (e.g., Yiğit & Medvedev, 2015; H.-L. Liu, 186 2016). These waves are generated in the lower layers of the atmosphere and propagate 187 into the thermosphere. They can interact with tides and other waves to produce secondary 188 waves, which makes the spatial temporal variability of the lower thermosphere rather com-189 plex (e.g., Chang et al., 2011; H.-L. Liu, 2014; Nystrom et al., 2018). The waves from 190 the lower atmosphere get strongly dissipated before reaching the middle thermosphere. 191 The middle thermosphere is dominated by the diurnal tide that is locally generated by 192 solar heating (Hagan et al., 2001), which can be represented well by formula (1). RMS 193 is somewhat larger in the middle thermosphere (20-30 m/s) than in the lower thermo-194 sphere (15-25 m/s). This reflects generally larger wind velocities in the middle thermo-195 sphere. These RMS values are much smaller than those reported for HWM14 (40–80 m/s) 196 by Drob et al. (2015). This is not surprising given that HWM14 involves more diverse 197 sources of data from many different years with various degrees of accuracy. 198

199 200

179

For a given month M, formula (1) can be rewritten in the following form, which more explicitly represents the zonal-mean winds and waves:

$$\overline{A} + \sum_{n=1}^{4} \sum_{s=-4}^{4} A_{ns} \cos\left(n\frac{t}{24} - s\frac{\lambda}{360} + P_{ns}\right) + \sum_{s=1}^{4} A'_s \cos\left(s\frac{\lambda}{360} + P'_s\right).$$
(4)

Here, \overline{A} is the zonal-mean wind velocity (in m/s), A_{ns} and P_{ns} are the amplitude (in m/s) and phase (in rad) of a tide, respectively. A'_s and P'_s are the amplitude and phase of a stationary planetary wave, respectively. 1- σ uncertainties in the zonal-mean winds, tides and stationary planetary waves were evaluated using the standard deviation obtained during the binning procedure described earlier. A Monte Carlo method was used for this purpose. That is, random noise was generated for each bin based on the standard deviation, and the noise was superimposed on the corresponding mean value. Fitting of



Figure 1. Correlation coefficient and root-mean-square error (RMS), as measures of goodnessof-fit of formula (1) to the v05 ICON/MIGHTI green-line data (top four panels) and red-line data (bottom four panels). The left and right panels are for the zonal and meridional winds, respectively.

formula (1) was repeated for 250 Monte Carlo samples, and 1- σ uncertainties were computed for the zonal-mean wind velocity, and the amplitude and phase of tides and stationary planetary waves at each latitude and height. The derived 1- σ uncertainty in the zonal-mean wind velocity is typically 1.0–3.5 m/s for both green-line and red-line winds. The 1- σ uncertainties in the amplitude and phase of tides and stationary planetary waves are typically less than 4.5 m/s and 20°, respectively. These uncertainty values are appreciably smaller compared to the features discussed in this paper.

Zonal-mean winds, tides and stationary planetary waves were also evaluated using HWM14 for the purpose of comparison. Hourly values of the zonal and meridional wind velocities were derived from HWM14 for each month by running the model for the 15th day of the month without including disturbance winds (Emmert et al., 2008). At each latitude and height, the zonal-mean wind velocity, and the amplitude and phase of tides and stationary planetary waves were determined by least-squares fitting of formula (4), which can be directly compared with the ICON/MIGHTI results.

223 3 Results

First, we examine seasonal climatologies of zonal-mean winds. Figure 2 depicts the 224 zonal-mean zonal and meridional winds in the lower thermosphere (91–110 km) as de-225 rived from the ICON/MIGHTI green-line measurements. Below ~ 105 km, the zonal-mean 226 zonal wind in the equatorial region $(10^{\circ}S-10^{\circ}N)$ tends to be weakly westward through-227 out the year. An eastward jet can be seen at 30°N during the Northern Hemisphere (N.H.) 228 summer. The reversal of the zonal-mean zonal wind is often seen around 105 km, which 229 was also noted by Yiğit et al. (2022). The zonal-mean meridional wind is generally weak 230 with little seasonal variation. The corresponding results obtained from HWM14 are pre-231 sented in Figure 3. HWM14 captures the salient features of the observed zonal-mean zonal 232 and meridional winds well. 233

Figure 4 shows the zonal-mean zonal and meridional winds in the middle thermosphere (203–300 km) as derived from the ICON/MIGHTI red-line measurements. An annual variation of the zonal-mean zonal wind is evident. That is, the zonal wind in the N.H. is largely eastward and westward during the local winter and summer, respectively. The seasonal variation of the zonal-mean meridional wind is also dominated by an annual cycle. That is, the meridional wind is primarily northward during the N.H. winter

-9-



ICON/MIGHTI (green-line, zonal mean zonal wind, m/s)

Figure 2. Quiet-time monthly climatologies of the zonal-mean zonal wind (top 12 panels) and zonal-mean meridional wind (bottom 12 panels) in the lower thermopshere (91–110 km) as derived from the v05 ICON/MIGHTI green-line data during April 2020–March 2022.

-10

0 10 20 30 40

Latitude (°)

-10

0 10 20 30 40

Latitude (°)

-10 0 10 20 30 40

Latitude (°)

-10

0 10 20 30 40

Latitude (°)



Figure 3. Same as Figure 2 but from Horizontal Wind Model 2014 (HWM14).

and southward during the N.H. summer. The seasonal transitions occur in March and
September. The annual variations in the zonal-mean zonal and meridional winds are reproduced well by HWM14 as shown in Figure 5.

Next, we examine seasonal climatologies of tides and stationary planetary waves. 243 Different components of waves, as expressed by different combinations of (n, s), have vary-244 ing degrees of significance in the thermosphere (e.g., Truskowski et al., 2014; Forbes et 245 al., 2014). Figures 6 and 7 depict wave spectra for two representative altitudes. Figure 246 6 shows the amplitude of different wave components for the green-line winds over the 247 equator at an altitude of 106 km. At this particular latitude and height, the eastward-248 propagating diurnal tide with zonal wavenumber 3 (DE3; n=1, s=3) dominates the tide 249 in the zonal wind, especially during July–November, and the migrating semidiurnal tide 250 (SW2; n=2, s=-2) dominates the tide in the meridional wind, especially during April-251 September. Figure 7 is similar to Figure 6 but for the red-line winds at 30°N at an al-252 titude of 273 km. In the middle thermosphere, the migrating diurnal tide (DW1; n=1, 253 s=-1) is by far dominant. Since DW1, SW2 and DE3 are found to be dominant within 254 the latitudinal and altitudinal range of the ICON/MIGHTI wind measurement, we fur-255 ther analyze these specific tides. 256

Figure 8 shows the amplitude and phase of DW1 in the meridional wind in the lower 257 thermosphere as derived from the ICON/MIGHTI green-line measurements. The am-258 plitude is largest at $15-20^{\circ}$ N and 95-97 km, and it shows a semiannual variation with 259 equinoctial maxima of ~ 60 m/s. The phase of DW1 tends to decrease with increasing 260 height. This 'downward phase propagation' is a fundamental feature of upward-propagating 261 tides (e.g., Forbes, 1995). The results suggest that DW1 in the lower thermosphere orig-262 inate from lower layers of the atmosphere. The latitude-height pattern of the DW1 phase 263 does not vary much with the season. DW1 derived from HWM14 is presented in Figure 264 9. The DW1 amplitude in HWM14 is largest at $15-20^{\circ}$ N, which is in agreement with 265 the ICON/MIGHTI results. HWM14 also reproduces the semiannual variation in the 266 DW1 amplitude. However, the DW1 amplitude in HWM14 is generally too small, and 267 its height structure does not agree well with the observations. The latitude and height 268 structures of the DW1 phase are reproduced by HWM14 during equinoctial months. 269

Figure 10 presents the SW2 amplitude and phase in the meridional wind in the lower thermosphere as derived from the ICON/MIGHTI green-line measurements. The am-



ICON/MIGHTI (red-line, zonal mean zonal wind, m/s)

Figure 4. Quiet-time monthly climatologies of the zonal-mean zonal wind (top 12 panels) and zonal-mean meridional wind (bottom 12 panels) in the middle thermosphere (203-300 km) as derived from the v05 ICON/MIGHTI red-line data during April 2020–March 2022.

220

-10

0 10 20 30 40

Latitude (°)

-10

0 10 20 30 40

Latitude (°)

220

-10 0 10 20 30 40

Latitude (°)

-10 0 10 20 30 40

Latitude (°)



HWM14 (zonal mean zonal wind, m/s)

Figure 5. Same as Figure 4 but from Horizontal Wind Model 2014 (HWM14).



ICON/MIGHTI (green-line, tides, zonal wind amplitude, m/s, 106 km, equator)

ICON/MIGHTI (green-line, tides, meridional wind amplitude, m/s, 106 km, equator)



Figure 6. Amplitude of tides and stationary planetary waves in the zonal wind (top 12 panels) and meridional wind (bottom 12 panels) at 106 km at the equator as derived from the v05 ICON/MIGHTI green-line data. n represents tidal frequency. That is, n=1 for diurnal tides, n=2 for semidiurnal tides, and so on. n=0 for stationary planetary waves. s is the zonal wavenumber. s>0 for eastward-propagating waves, while s<0 for westward-propagating waves.



ICON/MIGHTI (red-line, tides, zonal wind amplitude, m/s, 273 km, 30°N)

ICON/MIGHTI (red-line, tides, meridional wind amplitude, m/s, 273 km, 30°N)



Figure 7. Same as Figure 6 but at 273 km at 30°N as derived from the v05 ICON/MIGHTI red-line data.



ICON/MIGHTI (green-line, DW1, meridional wind amplitude, m/s)





Figure 8. Amplitude (top 12 panels) and phase (bottom 12 panels) of the migrating diurnal tide (DW1) in the meridional wind in the lower thermosphere as derived from the v05 ICON/MIGHTI green-line data. The corresponding results for the zonal wind can be found in Figure S1 of Supporting Information.



HWM14 (DW1, meridional wind amplitude, m/s)

HWM14 (DW1, meridional wind phase, deg.)



Figure 9. Same as Figure 8 but from Horizontal Wind Model 2014 (HWM14).

plitude is relatively large over the equator $(10^{\circ}S-10^{\circ}N)$ and at N.H. middle latitudes $(>30^{\circ}N)$. 272 In the equatorial region, the amplitude grows with height, reaching 60 m/s at 110 km273 during August–September. The maximum amplitude probably occurs above 110 km. At 274 middle latitudes, the amplitude peaks at 105 km. The downward phase propagation is 275 seen at both equatorial and middle-latitude regions, indicating that the SW2 energy prop-276 agates upward at these heights. The corresponding results derived from HWM14 are shown 277 in Figure 11. Again, the amplitude in HWM14 is generally too small, and its height struc-278 ture does not agree well with the observations. Interestingly, there is remarkable agree-279 ment in the phase of SW2 in the lower thermosphere between the ICON/MIGHTI and 280 HWM14 results. 281

Figure 12 shows the amplitude and phase of DE3 in the zonal wind in the lower thermosphere as derived from the ICON/MIGHTI green-line measurements. DE3 is the largest non-migrating (i.e., non-sun-synchronous) tidal component found in the greenline data. The zonal-wind amplitude is largest over the equator at a height of 105–110 km. The maximum amplitude exceeds 30 m/s during July–October. The downward phase propagation is visible, indicating upward energy propagation of DE3. DE3 is nonexistent in HWM14, as the model does not take into account any non-migrating tide.

We now look at DW1 in the middle thermosphere. Figure 13 shows the amplitude 289 and phase of DW1 in the zonal wind in the middle thermosphere as derived from the ICON/MIGHTI 290 red-line observations. It is noted that the scale range for the amplitude is different from 291 those used for the green-line results (Figures 8, 10 and 12). The DW1 amplitude grows 292 with height from ${\sim}50$ m/s at 200 km to ${\sim}90$ m/s at 300 km. It exceeds 100 m/s in some 293 months. The phase does not vary with height, indicating that DW1 in the middle ther-294 mosphere is a vertically-trapped (evanescent) tidal mode that is locally generated, rather 295 than an upward-propagating mode from below. The corresponding results derived from 296 HWM14 are presented in Figure 14. HWM14 reproduces the latitude and height struc-297 tures of the amplitude and phase well. Figure 15 also shows the amplitude and phase 298 of DW1 from the ICON/MIGHTI red-line measurements, but for the meridional wind. 299 The amplitude is small over the equatorial region but can exceed 100 m/s at middle lat-300 itudes (>30°N) above 280 km. The phase depends strongly on latitude. The phase struc-301 ture is well captured by HWM14 (Figure 16), but the model severely underestimates the 302 DW1 amplitude at middle latitudes. 303



ICON/MIGHTI (green-line, SW2, meridional wind amplitude, m/s)





Figure 10. Amplitude (top 12 panels) and phase (bottom 12 panels) of the migrating semidiurnal tide (SW2) in the meridional wind in the lower thermosphere as derived from the v05 ICON/MIGHTI green-line data. The corresponding results for the zonal wind can be found in Figure S2 of Supporting Information.



HWM14 (SW2, meridional wind amplitude, m/s)

HWM14 (SW2, meridional wind phase, deg.)



Figure 11. Same as Figure 10 but from Horizontal Wind Model 2014 (HWM14).



ICON/MIGHTI (green-line, DE3, zonal wind amplitude, m/s)



-10 0 10 20 30 40

Latitude (°)

-10 0 10 20 30 40 Latitude (°)

-10 0 10 20 30 40

Latitude (°)

-10 0 10 20 30 40

Latitude (°)



ICON/MIGHTI (red-line, DW1, zonal wind amplitude, m/s)

Figure 13. Amplitude (top 12 panels) and phase (bottom 12 panels) of the migrating diurnal tide (DW1) in the zonal wind in the middle thermosphere as derived from the v05 ICON/MIGHTI red-line data. The corresponding results for the migrating semidiurnal tide (SW2) can be found in Figure S4 of Supporting Information.

Latitude (°)

Latitude (°)

Latitude (°)

Latitude (°)



HWM14 (DW1, zonal wind amplitude, m/s)

Figure 14. Same as Figure 13 but from Horizontal Wind Model 2014 (HWM14).



ICON/MIGHTI (red-line, DW1, meridional wind amplitude, m/s)

ICON/MIGHTI (red-line, DW1, meridional wind phase, deg.)



Figure 15. Same as Figure 13 but in the meridional wind. The corresponding results for the migrating semidiurnal tide (SW2) can be found in Figure S5 of Supporting Information.



HWM14 (DW1, meridional wind amplitude, m/s)

HWM14 (DW1, meridional wind phase, deg.)



Figure 16. Same as Figure 14 but in the meridional wind.

304 4 Discussion

We have presented seasonal climatologies of the zonal-mean winds and tides derived from the v05 ICON/MIGHTI data, and compared the results with those from HWM14. Here we compare the ICON/MIGHTI results with those presented in earlier work based on other observations and models. Also, we discuss physical mechanisms behind some of the features observed in the ICON/MIGHTI winds, referring to previous theoretical studies.

Wang et al. (1997) created an empirical model of lower thermospheric winds (90– 311 120 km) using the measurements from the wind imaging interferometer (WINDII; Shep-312 herd et al., 1993) onboard UARS. They presented the zonal-mean zonal and meridional 313 winds for different seasons, which can be compared with our ICON/MIGHTI results (Fig-314 ures 2 and 4). S. P. Zhang et al. (2007) later analyzed an updated version of UARS/WINDII 315 data and obtained similar results as Wang et al. (1997). The UARS/WINDII results showed 316 a westward jet of 10-30 m/s over the equator at ~ 100 km throughout the year. This is 317 also seen in the ICON/MIGHTI results (Figure 2), as well as in HWM14 (Figure 3). It 318 is noted that the UARS/WINDII data are already incorporated in HWM. The westward 319 jet over the equator is considered to result from the westward momentum deposition by 320 dissipating migrating (thus westward-propagating) tides (e.g., Miyahara, 1981, 1978; Lieber-321 man & Hays, 1994; Jones Jr et al., 2014). Wang et al. (1997) and S. P. Zhang et al. (2007) 322 noted that the equatorial westward jet is sandwiched by eastward jets centered around 323 $\pm 40^{\circ}$ latitudes. The eastward jets were reported to be stronger in the summer hemisphere, 324 with the magnitude of 30–40 m/s. The ICON/MIGHTI results (Figure 2) clearly cap-325 ture the N.H. part of the eastward jets. The mechanism for the middle-latitude eastward 326 jets is not well understood. The numerical work by Forbes et al. (1993) predicted that 327 the equatorial westward jet induced by tidal dissipation is accompanied by eastward jets 328 at higher latitudes. However, the eastward jets due to tidal dissipation are predicted to 329 be much weaker than the westward jet, which agrees with neither ICON/MIGHTI nor 330 UARS/WINDII observations. Besides tides, Miyoshi and Fujiwara (2006) numerically 331 demonstrated that the momentum deposition by eastward-propagating equatorial Kelvin 332 waves also plays a significant role for the zonal-mean zonal wind in the equatorial lower 333 thermosphere. More studies are required to determine the relative importance of differ-334 ent waves to explain the observed westward and eastward jets. 335
The zonal-mean meridional wind in the lower thermosphere as derived from the ICON/MIGHTI 336 data is generally weak (Figure 2), which is consistent with the UARS/WINDII results 337 presented by Wang et al. (1997) as well as HWM14 (Figure 3). S. P. Zhang et al. (2007) 338 noted that the zonal-mean meridional wind sometimes show a cell-like structure in the 339 low latitude region, which is characterized by poleward winds on both sides of the equa-340 tor at altitudes of 95–105 km and equatorward winds at 105–115 km. There is some hint 341 of such a cell-like structure in the ICON/MIGHTI zonal-mean meridional wind (see, e.g., 342 August–September), but it is not well resolved because of the small magnitude. The cell-343 like structure in the zonal-mean meridional wind in the lower thermosphere is sometimes 344 found in numerical models and is considered to be driven by tidal dissipation (e.g., Miya-345 hara et al., 1993; Forbes et al., 1993). 346

ICON/MIGHTI red-line measurements revealed seasonal climatologies of the zonal-347 mean zonal and meridional winds at 200–300 km (Figure 4). HWM14 reproduces the 348 ICON/MIGHTI observations well for both the zonal and meridional components (Fig-349 ure 5). HWM14 at this height range is well constrained by UARS/WINDII red-line winds 350 as well as observations by ground-based Fabry-Perot interferometers. The zonal-mean 351 meridional wind in the middle thermosphere is directed from the summer to the winter 352 hemisphere, which is not surprising given the higher temperature and pressure in the sum-353 mer hemisphere. The seasonal transition in the meridional circulation occurs in March 354 and September. Using a numerical model, Roble et al. (1977) showed that the seasonal 355 transition of the zonal-mean circulation takes place within a few weeks of equinox. Such 356 an abrupt seasonal transition is not fully resolved in our monthly analysis. The zonal-357 mean zonal wind in the middle thermosphere arises mainly from the correlation between 358 diurnal variations of pressure gradient and ion drag (Dickinson et al., 1975, 1977). That 359 is, the wind is weaker on the dayside than the nightside as the ion drag is larger on the 360 dayside due to higher plasma concentration. Since the wind in the middle thermosphere 361 undergoes a diurnal cycle due to day-night pressure differences, an unbalance between 362 the daytime and nighttime winds leads to the zonal-mean winds. 363

The three most dominant tidal components in the ICON/MIGHTI green-line winds are DW1, SW2 and DE3 (Figures 8, 10, 12). This is as expected from previous studies on tides in the lower thermosphere (e.g., Forbes et al., 2008; Oberheide et al., 2011). DW1 and SW2 are sun-synchronous, while DE3 is non-sun-synchronous. In the lower thermosphere, they consist mainly of upward-propagating modes, which can be seen from their

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downward phase propagation. They are driven by radiative heating through insolation 369 of H_2O in the troposphere and O_3 in the stratosphere (e.g., Forbes, 1982b, 1982a) as well 370 as by latent heating in the troposphere (Hagan & Forbes, 2002, 2003; X. Zhang et al., 371 2010a, 2010b). DW1 in the meridional wind, as derived from the ICON/MIGHTI ob-372 servations, shows an amplitude maximum at $15-20^{\circ}$ N at an altitude of 95-98 km (Fig-373 ure 8). The results are consistent with those from UARS/WINDII (McLandress et al., 374 1996; S. P. Zhang et al., 2007) and TIMED/TIDI (Wu et al., 2008a). The amplitude is 375 larger during the equinoxes than the solstices, which is well known from previous stud-376 ies (e.g., Burrage et al., 1995; Xu et al., 2009). McLandress (2002b, 2002a) examined the 377 mechanism for the semiannual variation of DW1 using a numerical model, and concluded 378 that the change in the latitudinal shear of the zonal-mean zonal wind plays a leading role 379 for the seasonal variation of DW1 in the lower thermosphere. HWM14 reproduces the 380 semiannual variation of DW1 (Figure 9) but the model underestimates the amplitude 381 in comparison not only with the ICON/MIGHTI results but also with the UARS/WINDII 382 and TIMED/TIDI results (S. P. Zhang et al., 2007; Wu et al., 2008a). Previous stud-383 ies reported that the amplitude of DW1 at low latitudes can change by a few tens of m/s 384 from one year to the next (e.g., Burrage et al., 1995; Hagan et al., 1999). Variation as-385 sociated with the quasi-biennial oscillation (QBO) of the equatorial atmosphere is an im-386 portant part of the interannual variation of DW1 in the lower thermosphere, account-387 ing for up to 10 m/s (e.g., Xu et al., 2009). The interannual variability of tides is not 388 taken into account in the present study. Resolving the QBO effect would require a larger 389 data set. 390

SW2 in the meridional wind, as derived from the ICON/MIGHTI observations, is 391 relatively strong during May–September (Figure 10), which is consistent with the UARS/WINDII 392 observations (S. P. Zhang et al., 2007). HWM14 reproduces the seasonal variation of SW2 393 but with somewhat smaller amplitude (Figure 11). The mechanism for the seasonal vari-394 ation of SW2 is not well established. DE3 in the lower thermosphere has characteristics 395 of a Kelvin wave (e.g., Forbes et al., 2003). In classical theory, a Kelvin wave travels east-396 ward, and its zonal wind component has a Gaussian-shaped latitudinal profile with max-397 imum amplitude over the equator (e.g., Forbes, 2000). The latitude and height struc-398 tures of DE3 and its seasonal variation in the ICON/MIGHTI green-line zonal wind (Fig-399 ure 12) are consistent with those from the UARS/WINDII (Forbes et al., 2003) and TIMED/TIDI 400 observations (Oberheide et al., 2006; Wu et al., 2008b). As the zonal wind amplitude of 401

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DE3 reaches its maximum in the equatorial dynamo region at 105–110 km, it has a significant impact on the equatorial zonal electric field and current (e.g., England et al., 2006;
Fejer et al., 2008) as well as on the F-region plasma concentration (e.g., Immel et al.,
2006; Lin et al., 2007). Despite the importance of DE3 in low-latitude ionosphere-thermosphere

coupling, it is not included in HWM14 like other non-migrating tides.

DW1 in the middle thermosphere (Figures 13 and 15) is predominantly a vertically-407 trapped tidal mode that is excited by in-situ solar heating (e.g., Forbes, 1982b; Hagan 408 et al., 2001). This contrasts with DW1 in the lower thermosphere (Figure 8), which is 409 primarily an upward-propagating mode. The latitude and height structures of DW1 in 410 the middle thermosphere are not well documented, particularly those based on obser-411 vations. The simulation results by Hagan et al. (2001) showed that (1) the amplitude 412 of DW1 at 200–300 km grows with height at all latitudes, (2) both zonal and meridional 413 wind amplitudes are largest at high latitudes, (3) the meridional wind amplitude is van-414 ishingly small over the equator but it increases with latitude, (4) the zonal wind ampli-415 tude does not depend strongly on latitude over the middle- and low-latitude regions, (5) 416 both zonal and meridional wind phases do not depend strongly on height, (6) the zonal 417 wind phase does not vary strongly with latitude, and (7) the meridional wind phase also 418 does not vary strongly with latitude except that the phase reversal occurs at the equa-419 tor. The ICON/MIGHTI results (Figures 13 and 15) are consistent with these numer-420 ical predictions. 421

Some previous studies have addressed a potential impact of the solar flux, mainly 422 at the wavelengths of extreme ultraviolet (EUV), on neutral winds in the middle and up-423 per thermosphere (e.g., Hedin et al., 1994). The ICON/MIGHTI observations examined 424 in this paper are obtained during the period April 2020–March 2022. The mean value 425 of the $F_{10.7}$ index (Tapping, 2013), which is often used as a proxy of the EUV flux, was 426 82.8 sfu (1 sfu = 10^{-22} W·m⁻²·Hz⁻¹), with the minimum and maximum monthly val-427 ues of 69.2 sft in May 2020 and 117.8 sfu in March 2022, respectively. We have ignored 428 possible variations in the wind velocities associated with the change in the solar flux, as 429 the ICON/MIGHTI data used in this study are not sufficient for evaluating the solar ac-430 tivity effect on the thermospheric winds. HWM14 also does not take into account the 431 dependence of wind velocities on solar activity. Hedin et al. (1994) reported that although 432 the solar flux significantly influences the temperature, the zonal-mean winds in the mid-433 dle thermosphere do not strongly depend on solar activity. Hagan et al. (2001) noted in 434

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their simulation results that the solar flux has a marked effect on the temperature am-435 plitude of DW1 in the middle thermosphere but not on the wind amplitudes. H. Liu et 436 al. (2006) examined the effect of the solar flux on thermospheric winds (\sim 400 km) us-437 ing the CHAMP accelerometer data. They found that the solar flux effect can be sig-438 nificant depending on the season and local time. The solar flux effect on tides is gener-439 ally small below about 130 km, where tidal waves are mainly of lower atmospheric ori-440 gin (Oberheide et al., 2009; Dhadly et al., 2018). More observational studies are required 441 to establish the solar activity dependence of zonal-mean winds and tides in the thermo-442 443 sphere.

444 5 Conclusions

Monthly climatologies of quiet-time zonal-mean winds and tides are derived using the recently-released v05 of the ICON/MIGHTI thermopsheric wind measurements during April 2020–March 2022 at the altitude ranges 90–110 km and 200–300 km. Earlier versions of the ICON/MIGHTI wind data suffered from artificial baseline drifts that depend on local time. Thus, it was previously difficult to obtain reliable climatological estimates of zonal-mean winds and tides. The v05 data avoids this issue by the use of a renewed baseline calibration technique (Englert et al., 2023).

The ICON/MIGHTI results are compared with those from the latest version of HWM 452 (i.e., HWM14) as well as previous studies. Salient features of zonal-mean winds and tides 453 in the lower and middle thermosphere are in general agreement between ICON/MIGHTI 454 and HWM14, including latitude and height structures and their seasonal variations. This 455 provides a validation of the v05 ICON/MIGHTI data. HWM14 reproduces the zonal-456 mean zonal and meridional winds well in both the lower and middle thermosphere. How-457 ever, HWM14 tends to underestimate tidal amplitude. Also, HWM14 does not include 458 non-migrating tides such as DE3, which is especially important in the equatorial lower 459 thermosphere. The latitude and height structures of DE3 and their seasonal variations 460 in the ICON/MIGHTI green-line zonal wind are found to be consistent with those from 461 the UARS/WINDII and TIMED/TIDI observations. The future improvement of HWM 462 can benefit from the inclusion of the ICON/MIGHTI winds for better description of tides. 463

464 Open Research Section

The ICON/MIGHTI Level 2.2 product Cardinal Vector Winds (Version 5) is accessible from the ICON website https://icon.ssl.berkeley.edu/Data. The Hpo indices including Hp30 used in this study are available at the GFZ website https://kp .gfz-potsdam.de/en/hp30-hp60/data; see also data publication Matzka et al. (2022). The monthly F10.7 index is available at the website of the Canadian Space Weather Forecast Centre https://spaceweather.gc.ca/forecast-prevision/solar-solaire/solarflux/ sx-5-mavg-en.php.

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Earth and Space Science

Supporting Information for

"Monthly climatologies of zonal-mean and tidal winds in the thermosphere as observed by ICON/MIGHTI during April 2020–March 2022"

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Figures S1 to S5

Introduction

This document presents latitude and height structures of atmospheric tides as derived using version 5 of the thermospheric wind data from the Michelson Interferometer for Global High-resolution Thermospheric Imaging (MIGHTI) onboard NASA's Ionospheric Connection Explorer (ICON) mission. The data used are those observed during April 2020–March 2022. The procedures for retrieving tidal amplitude and phase are described in the main text.



ICON/MIGHTI (green-line, DW1, zonal wind amplitude, m/s)

Figure S1. Amplitude (top 12 panels) and phase (bottom 12 panels) of the migrating diurnal tide (DW1) in the zonal wind in the lower thermosphere as derived from the v05 ICON/MIGHTI greenline data.



ICON/MIGHTI (green-line, SW2, zonal wind amplitude, m/s)

Figure S2. Amplitude (top 12 panels) and phase (bottom 12 panels) of the migrating semidiurnal tide (SW2) in the zonal wind in the lower thermosphere as derived from the v05 ICON/MIGHTI green-line data.

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ICON/MIGHTI (green-line, DE3, meridional wind amplitude, m/s)

ICON/MIGHTI (green-line, DE3, meridional wind phase, deg.)



Figure S3. Amplitude (top 12 panels) and phase (bottom 12 panels) of the eastward-propagating diurnal tide with zonal wavenumber 3 (DE3) in the meridional wind in the lower thermosphere as derived from the v05 ICON/MIGHTI green-line data.



ICON/MIGHTI (red-line, SW2, zonal wind amplitude, m/s)

Figure S4. Amplitude (top 12 panels) and phase (bottom 12 panels) of the migrating semidiurnal tide (SW2) in the zonal wind in the middle thermosphere as derived from the v05 ICON/MIGHTI red-line data.



ICON/MIGHTI (red-line, SW2, meridional wind amplitude, m/s)

Figure S5. Amplitude (top 12 panels) and phase (bottom 12 panels) of the migrating semidiurnal tide (SW2) in the meridional wind in the middle thermosphere as derived from the v05 ICON/MIGHTI red-line data.