

# The ionospheric effects of the 2022 Hunga Tonga Volcano eruption and the associated impacts on GPS Precise Point Positioning across the Australian region

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## Abstract

The Hunga Tonga Volcano eruption launched a myriad of atmospheric waves that have been observed to travel around the world several times. These waves generated Traveling Ionospheric Disturbances (TIDs) in the ionosphere, which are known to adversely impact radio applications such as Global Navigation Satellite Systems (GNSS). One such GNSS application is Precise Point Positioning (PPP), which can achieve cm-level accuracy using a single receiver, following a typical convergence time of 30 mins to 1 hour. A network of ionosondes located throughout the Australian region were used in combination with GNSS receivers to explore the impacts of the Hunga-Tonga Volcano eruption on the ionosphere and what subsequent impacts they had on PPP. It is shown that PPP accuracy was not significantly impacted by the arrival of the TIDs and Spread-F, provided that PPP convergence had already been achieved. However, when the PPP algorithm was initiated from a cold start either shortly before or after the TID arrivals, the convergence times were significantly longer. GNSS stations in northeastern Australia experienced increases in convergence time of more than 5 hours. Further analysis reveals increased convergence times to be caused by a super equatorial plasma bubble (EPB), the largest observed over Australia to date. The EPB structure was found to be  $\sim 42$  TECU deep and  $\sim 300$  km across, traveling eastwards at 30 m/s. The Hunga Tonga Volcano eruption serves as an excellent example of how ionospheric variability can impact real-world applications and the challenges associated with modeling the ionosphere to support GNSS.



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## 41 Plain Language Summary

42 Global Navigation Satellite System (GNSS) applications permeate modern soci-  
 43 ety, with many industry sectors heavily relying on precision satellite positioning, nav-  
 44 igation and timing. Precise Point Positioning (PPP) is an advanced positioning technique  
 45 that can achieve cm-level accuracy without the need for nearby reference stations. How-  
 46 ever, the time that it takes for the PPP solution to ‘converge’ is typically in the range  
 47 of 10s of mins to hours, limiting the widespread uptake of PPP. There are numerous pre-  
 48 vious reports of waves and disturbances in the ionosphere, which are known to adversely  
 49 impact GNSS applications. In this study, the impact of the disturbances in the ionosphere  
 50 caused by the 2022 Hunga Tonga-Hunga Ha’apai Volcano eruption on PPP across the Aus-  
 51 tralian region is investigated. It is found that convergence times increased by more than  
 52 5 hours across northern Australia due to small-scale ionospheric turbulence. The source  
 53 of the turbulence was also found in this analysis to be due to a ‘super Equatorial Plasma  
 54 Bubble’ that persisted above northern Australia for several hours. This event serves as  
 55 an excellent example of how ionospheric disturbances can impact relied upon GNSS ap-  
 56 plications.

## 57 1 Introduction

58 At 04:14:45 UT on January 15, 2022, the Hunga Tonga-Hunga Ha’apai Volcano erupted  
 59 in what was one of the largest explosions on Earth in modern history (Matoza et al., 2022;  
 60 Wright et al., 2022). It has been estimated that the eruption released somewhere between  
 61  $3.7 \times 10^{16}$  to  $8.37 \times 10^{17}$  Joules (Wright et al., 2022; Astafyeva et al., 2022; Díaz & Rigby,  
 62 2022; Vergoz et al., 2022), making it comparable to the Krakatoa eruption in 1883 (Pyle,  
 63 2015). The explosion caused a tsunami that reached all sides of the Pacific Ocean with  
 64 an observed maximum wave-height of 3.4m on the Chilean shoreline (Carvajal et al., 2022).  
 65 The eruption was even audible as far away as Alaska, which is some 10,000 km away from  
 66 Tonga (Matoza et al., 2022). Not surprisingly, the eruption caused a myriad of waves  
 67 in the atmosphere and ionosphere (Themens et al., 2022; Wright et al., 2022; Aa et al.,  
 68 2022; Astafyeva et al., 2022; Ghent & Crowell, 2022; Maletckii & Astafyeva, 2022; Hong  
 69 et al., 2022) that were observed to encircle the Earth multiple times (S.-R. Zhang et al.,  
 70 2022; Matoza et al., 2022; Pradipta et al., 2023).

71 Waves in the ionosphere can be remotely detected using Global Navigation Satel-  
72 lite System (GNSS) signals in terms of the delay that the ionospheric plasma imparts  
73 upon the signals. The phase delay that is measured is related to the total electron con-  
74 tent (TEC) between the GNSS satellite and the receiver, with the majority of the TEC  
75 contribution coming from the ionosphere (Yizengaw et al., 2008). An expansive network  
76 of ground-based GNSS receivers therefore make it possible to geographically map iono-  
77 spheric TEC around the world (Mannucci et al., 1998) and track the propagation of any  
78 ionospheric waves (e.g., Otsuka et al., 2002; Kotake et al., 2006; Borries et al., 2009; Tsug-  
79 awa et al., 2011; Pradipta et al., 2016; H. Yang et al., 2017; Lay et al., 2018). To date,  
80 a range of ionospheric effects associated with the Hunga Tonga Volcano event have been  
81 published, and many of these studies have utilized the International GNSS Service net-  
82 work of receivers (Johnston et al., 2017).

83 Themens et al. (2022) and S.-R. Zhang et al. (2022) reported both large-scale and  
84 medium-scale traveling ionospheric disturbances (TIDs) propagating away from the vol-  
85 cano location. S.-R. Zhang et al. (2022) showed evidence of these TIDs continuing to prop-  
86 agate around the world for at least 4 days. Harding et al. (2022) and Le et al. (2022) ob-  
87 served changes to the equatorial electrojet that were caused by variations in the iono-  
88 spheric dynamo as a result of the eruption. Aa et al. (2022) reported a localized iono-  
89 spheric plasma depletion in the vicinity of the volcano and increased Equatorial Plasma  
90 Bubble (EPB) activity in the Asia-Oceania low-latitude region, which has been further  
91 supported by recent modeling efforts (Huba et al., 2022). Evidence of ‘super EPBs’ has  
92 been reported spanning across Chinese (Sun, Wenjie et al., 2022) and Japanese/Australian  
93 (Rajesh et al., 2022) longitude sectors following the eruption. Aa et al. (2022) also re-  
94 ported propagating ionospheric irregularities that exhibited a phase speed that matched  
95 the prevailing Lamb mode at  $\sim 315$  m/s.

96 GNSS Precise Point Positioning (PPP) is an advanced positioning technique that  
97 uses dual frequency observations made by a single receiver to achieve cm-level position-  
98 ing accuracy (Zumberge et al., 1997; Leick et al., 2015; Choy et al., 2017; Teunissen &  
99 Montenbruck, 2017). Over the previous few decades, the dominant method for achiev-  
100 ing cm-level accuracy has been relative positioning, in which a nearby accurately located  
101 reference station is used to determine the precise position of a ‘rover’ receiver (e.g., Odijk,  
102 2002; Hofmann-Wellenhof et al., 2007; Leick et al., 2015; Teunissen & Montenbruck, 2017).  
103 Utilizing a nearby reference station makes it possible to eliminate spatially correlated  
104 GNSS observations errors, such as the tropospheric and ionospheric delays imposed on  
105 the GNSS signals. Unsurprisingly, this precise positioning capability has now found its  
106 way into many applications across several major industries, including mining, agricul-  
107 ture and construction (e.g., Pérez-Ruiz et al., 2015; Choy et al., 2017; Woodgate et al.,  
108 2017; Rao et al., 2022). Over recent years, PPP is fast becoming the new global stan-  
109 dard for cm-level positioning applications due to its ability to model and account for GNSS  
110 observations errors without the need for a nearby reference station. However, the most  
111 significant drawback that is limiting the widespread uptake of PPP is the rather long  
112 convergence times, which are typically on the order of 10s of mins to hours (Bisnath &  
113 Gao, 2009; Van Bree et al., 2009; Choy et al., 2017). As a result, researchers have been  
114 investigating methods to significantly reduce PPP convergence times (e.g., Collins & Bis-  
115 nath, 2011; Collins et al., 2012; Geng & Bock, 2013; H. Zhang et al., 2013; Banville et  
116 al., 2014; Li et al., 2015; Duong, 2020). Given the significant influence of the ionosphere  
117 on GNSS signals, space weather poses a potential vulnerability to the PPP technique  
118 and its anticipated widespread up-take. As such, researchers have also been investigat-  
119 ing the impact of solar and geomagnetic activity on PPP (e.g., Luo et al., 2018; Poni-  
120 atowski & Nykiel, 2020; Z. Yang et al., 2020; Zha et al., 2021; Luo et al., 2022).

121 The Hunga Tonga Volcano eruption provides a unique opportunity to uncover new  
122 fundamental knowledge of the physics of atmosphere-ionosphere coupling. In addition,  
123 given the adverse influence of ionospheric variability on GNSS, this event can also pro-



Figure 1: The locations of the ionosondes used in this analysis. Tonga is indicated by the red cross.

124 vide insights into the impact of such ionospheric disturbances across broader society, par-  
 125 ticularly in industries that rely heavily on precise GNSS positioning. In this study, the  
 126 impact of the Tonga volcano eruption on Global Positioning System (GPS) PPP across  
 127 the Australian region is explored. Supporting this analysis are the data collected by ionoson-  
 128 des throughout the region that show both TID and Spread-F activity in the wake of the  
 129 volcano eruption. The PPP accuracy throughout the day of the eruption is explored, fol-  
 130 lowed by an analysis of the PPP convergence times from stations located throughout the  
 131 region. Finally, ionospheric observations using GPS receivers across the region are used  
 132 to investigate the physical mechanisms causing the disruptions identified in this study.

## 133 2 Data

134 In this analysis, the Australian Bureau of Meteorology Space Weather Services' ionosonde  
 135 data are used. Figure 1 shows the locations of the ionosonde stations (black triangles)  
 136 and Tonga volcano (red cross). The stations are Niue (NUE5F; 19.07°S, 190.07°E), Nor-  
 137 folk Island (NLK5D; 29.03°S, 167.97°E), Canberra (CBR5F; 35.32°S, 149.00°E), Townsville  
 138 (TVL5F; 19.63°S, 146.8°E), Darwin (DWN5D; 12.45°S, 130.95°E), Perth (PER5D; 31.94°S,  
 139 115.95°E), Learmonth (LEA5D; 22.25°S, 114.08°E) and Cocos Keeling Islands (CCK5F;  
 140 12.20°S, 96.80°E). Each ionosonde generates ionograms by sweeping through radio fre-  
 141 quencies between 2-22 MHz, transmitting and receiving ionospheric echoes that indicate  
 142 the electron density for a given virtual height determined by time-of-flight. In this anal-  
 143 ysis, the ionosonde data were used to indicate the presence of TIDs and spread-F traces.

144 The ionosonde data complement the primary dataset used in this study, namely  
 145 the GPS receiver data. GPS Continuously Operating Reference Station (CORS) data  
 146 from stations located across Australia, spanning into the Southeast Asian region and across  
 147 the South Pacific, were used in this analysis. While previous works have employed such  
 148 GPS CORS data to analyze ionospheric TEC fluctuations caused by the eruption, this  
 149 study focuses on the impact of this eruption on the GPS application of PPP.

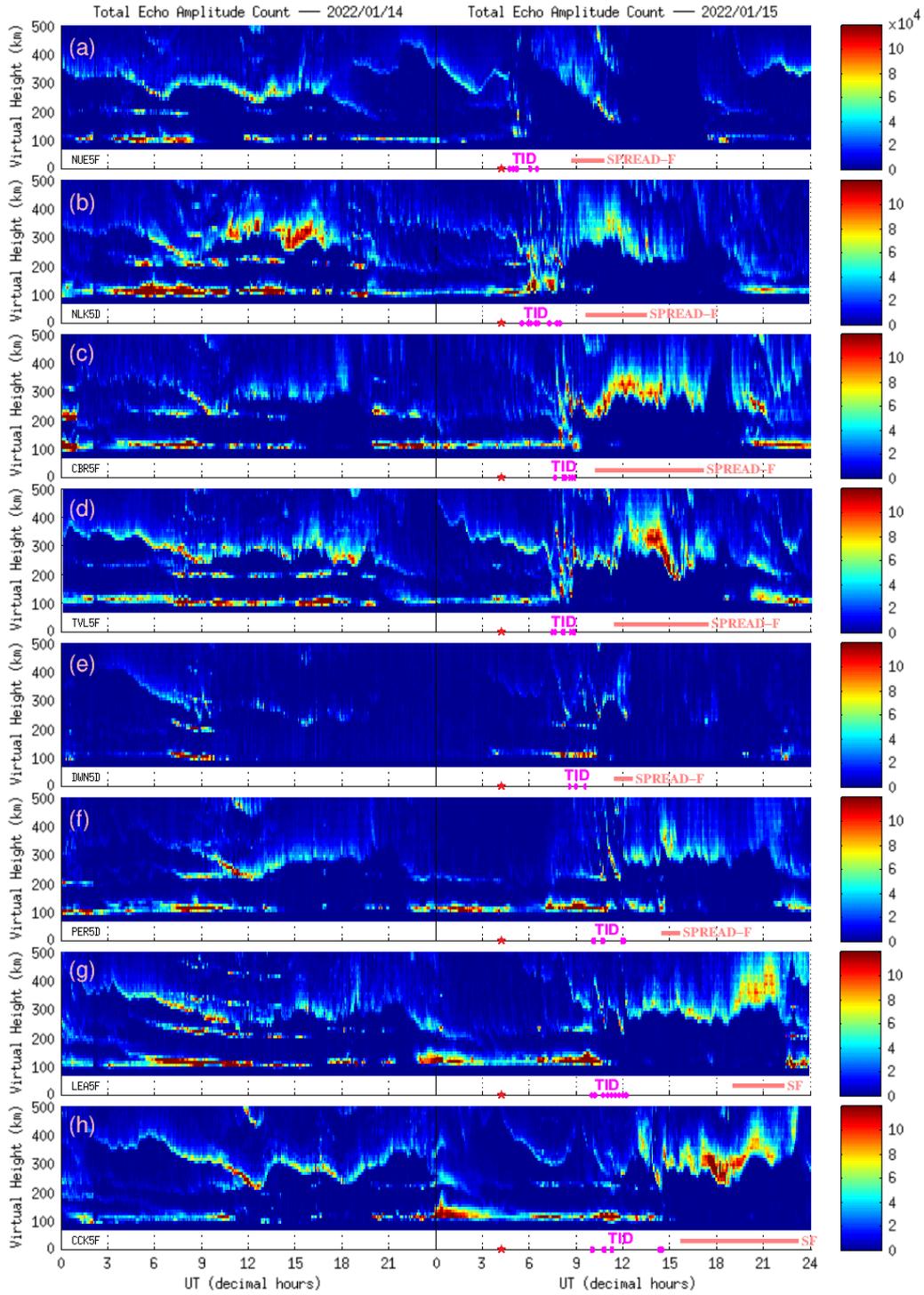


Figure 2: Range-time-intensity-style plots of ionospheric echoes received from ionosondes across the Australian region during January 14-15, 2022. Each panel corresponds to each station; (a) Niue (NUE5F), (b) Norfolk Island (NLK5D), (c) Canberra (CBR5F), (d) Townsville (TVL5F), (e) Darwin (DWN5D), (f) Perth (PER5D), (g) Learmonth (LEA5D) and (h) Cocos Keeling Islands (CCK5F). The color of each bin indicates the number of echoes received in that virtual height bin across all scanning frequencies, as indicated in the color bar on the right. The red stars indicate the time of the eruption; 04:14:45 UT. The pink dots (salmon bars) indicate the presences of TID (Spread-F) signatures.

150 The Geoscience Australia’s Ginan system was used for performing the PPP cal-  
 151 culations<sup>1</sup>. PPP is a high accuracy positioning method used to correct errors in GNSS  
 152 positioning based on the robust modelling and estimations of systematic errors in the  
 153 GNSS signals. The specific Ginan setup used in this analysis closely follows the Ginan  
 154 PPP example “Ex11”, in which the PPP is performed in ‘static mode’ using the ionosphere-  
 155 free combination<sup>2</sup>.

156 The carrier-phase measurements between the GPS satellites and receivers were used  
 157 to calculate the TEC along the signal path (i.e., the slant TEC, sTEC) following sim-  
 158 ilar methodology to Le Huy et al. (2016) and T. Dao et al. (2020). The possible jumps  
 159 in the sTEC values estimated from the carrier-phase measurements due to cycle slips were  
 160 eliminated by comparing them against the sTEC estimated from the pseudo-range mea-  
 161 surements that were smoothed by a fourth-degree polynomial approximation. A differ-  
 162 ence between the carrier-phase sTEC and the pseudo-range sTEC of more than 5 TECU  
 163 was taken to indicate an instrumental data jump/spike in the carrier-phase sTEC; in such  
 164 instances, the smoothed pseudo-range sTEC was used. The sTEC was then compared  
 165 with the Centre of Orbit Determination in Europe Global model to determine the total  
 166 delay of device (including biases), using elevation angles above 30° to remove mul-  
 167 tipath effects. Finally, the resulting sTEC was converted to the vertical total electron  
 168 content (VTEC) observed at the pierce point of the ionosphere by using a single-layer  
 169 model (Klobuchar, 1986) and an assumed altitude of 400 km. The presence of scintillation-  
 170 causing small-scale ionospheric irregularities can be detected using the rate of TEC in-  
 171 dex (ROTI), which is defined as the 5-min standard deviation in the rate of change in  
 172 the sTEC for each satellite-to-ground link (Pi et al., 1997).

### 173 3 Results

174 This analysis begins with an overview of the ionospheric conditions before and af-  
 175 ter the arrival of the disturbances caused by the volcano eruption using the ionosonde  
 176 network in Fig. 1. First, it is worth mentioning that a minor geomagnetic storm occurred  
 177 in the late hours of January 14, with Dst reaching -91 nT due to a small, short-lived re-  
 178 current solar wind stream that had a minimal effect on the equatorial electric field (Le  
 179 et al., 2022). A useful way to analyze temporal changes in ionograms is by representing  
 180 the data in a format similar to a ‘range-time-intensity’ plot (Pradipta et al., 2015; Carter  
 181 et al., 2018; Currie et al., 2021). Instead of using total power, Pradipta et al. (2015) in-  
 182 tegrated over the dBm amplitudes across all sounding frequencies, effectively creating  
 183 a sum of digitized echoes. This methodology was adopted for the present study. Figure  
 184 2 shows these range-time-intensity-style plots for the ionosondes during January 14-15,  
 185 2022; the panels are ordered by the station great circle distance to Tonga, closest to far-  
 186 thest. The colors indicate the number of echoes received across all frequencies in each  
 187 virtual height bin (10 kms). The presence of TIDs and spread-F traces are indicated at  
 188 the bottom of each plot by the pink dots and salmon bars, respectively; these have been  
 189 determined by visual inspection of the ionograms.

190 Figure 2 shows a range of ionospheric conditions across January 14 and 15; e.g.,  
 191 sporadic E features at virtual heights of 100 km are present for all stations at various  
 192 times and F region traces that display typical altitude changes with time. TID signa-  
 193 tures are also clear in the data on January 15 for all stations, beginning at approx. 4:30  
 194 UT at Niue (top panel) and at approx. 10 UT at Cocos Keeling Islands, with the other  
 195 stations showing the TID signatures at times in between. All stations, with the excep-  
 196 tion of Darwin which was suffering some intermittent hardware issues during this period,  
 197 show the presence of spread-F traces in the hours following the TIDs, and in all cases,

<sup>1</sup> <https://geoscienceaustralia.github.io/ginan/>

<sup>2</sup> [https://geoscienceaustralia.github.io/ginan/codeDocs/Pea\\_8Configuration\\_8Examples.html](https://geoscienceaustralia.github.io/ginan/codeDocs/Pea_8Configuration_8Examples.html)

198 the spread-F traces are persistent for a number of hours. In particular, the TID activ-  
 199 ity at Norfolk Island begins shortly before 6 UT, with Spread-F present from 09:15 UT  
 200 until 13:30 UT. Thus, the occurrence of spread-F traces measured by the Norfolk Island  
 201 ionosonde span local times (LT) 21:15-01:30. At Townsville, the TIDs are first observed  
 202 close to 7 UT (17 LT) and bring with them batches of Spread-F. The Spread-F then in-  
 203 tensifies close to 11:24 UT (21:24 LT) and remains strong until 15 UT (1:00 LT), before  
 204 finally ceasing at 17:30 UT (03:30 LT).

205 Next, the GPS data from receivers close to the Norfolk Island and Townsville ionosonde  
 206 stations are examined as initial examples of the features and trends present throughout  
 207 the GPS data. Figure 3 shows the observed vertical TEC (VTEC, upper panels) and the  
 208 rate of TEC index (ROTI, lower panels) observed by the Norfolk Island (NORF, left)  
 209 and Townsville (TOW2, right) GPS receivers on January 15, 2022. The VTEC data for  
 210 both stations show the presence of TIDs that disrupt the diurnal pattern starting at  $\sim$   
 211 6 and 7 UT for NORF and TOW2, respectively, in good agreement with the ionosonde  
 212 observations in Fig. 2. In the NORF data, the TIDs are observed in the individual satel-  
 213 lite traces as rather small-amplitude wave structures compared to the variations over the  
 214 24-hour period. Slight differences in the timings of these structures (or phase progres-  
 215 sion) for different satellite links are due to the geographical spread of the ionospheric pierce  
 216 points and the motion of the TIDs over the station. In the TOW2 data, the TIDs are  
 217 much clearer as strong changes in VTEC for specific satellite links at different times. The  
 218 phase progression of the VTEC structures overhead are much clearer in the TOW2 data.  
 219 The lower panels reveal the times in which increased ROTI values were observed. For  
 220 Norfolk Island (Fig. 3c) ROTI increases shortly after the TIDs arrive and then decreases  
 221 back to low levels by 9 UT. Townsville also observed increased ROTI values once the main  
 222 TIDs arrived and remained high until  $\sim$  15 UT once the strong Spread-F ceases, c.f. Fig.  
 223 2.

224 Next, we illustrate how the GINAN PPP software achieves convergence, as defined  
 225 here by a 3-D position error of less than 10 cm. Figure 4 shows six examples of PPP con-  
 226 vergence on January 15, 2022 using data collected by the NORF (top row) and TOW2  
 227 (bottom row) receivers. Each panel represents a ‘cold start’ of the PPP algorithm start-  
 228 ing at a different time; 02:30 UT, 04:30 UT and 06:30 UT, as indicated in the subfigure  
 229 titles. The positioning error was calculated as the difference between the PPP estima-  
 230 tions and the known geodetic station locations. The position errors in the X, Y, and Z  
 231 directions, as well as the full 3-D position, are shown in each panel. The 10 cm position  
 232 error threshold is indicated by the dash-dot line, and zero error is indicated by the dashed  
 233 line. Also displayed in each panel is the time it took for the 3-D position error to reach  
 234 below 10 cm, indicated by the vertical dotted line; i.e., the convergence time.

235 Given the TID and subsequent Spread-F activity shown in the ionosonde data in  
 236 Fig. 2, one might expect that PPP errors for nearby GPS receivers would noticeably in-  
 237 crease during these periods. However, Fig. 4a shows that despite the onset of TIDs at  
 238 approx. 6 UT for the NORF station, the position error remained below the 10-cm thresh-  
 239 old. The same is true in Fig. 4d for the TOW2 station, for which the onset of TIDs was  
 240 approx. 7 UT. While Fig. 4b shows that the NORF station position error increased fol-  
 241 lowing the commencement of Spread-F close to 9 UT, it remained well below 10 cm. In-  
 242 terestingly, despite the commencement of rather strong Spread-F at 11:24 UT at TOW2,  
 243 Fig. 4f shows no significant increase in the position error. This is likely the result of run-  
 244 ning the PPP algorithm in ‘static mode’. From these results, one can conclude that the  
 245 static-mode PPP algorithm was robust enough to remain at cm-level accuracy for these  
 246 stations throughout the turbulent ionospheric conditions caused by the Tonga eruption,  
 247 provided that convergence had already been achieved.

248 Figure 4 also illustrates how the convergence time is dependent on the start time  
 249 of the algorithm. For the NORF station, the convergence time ranges from less than 1  
 250 hour beginning at 04:30 UT (Fig. 4b) to almost 4 hours when started from 06:30 UT (Fig.

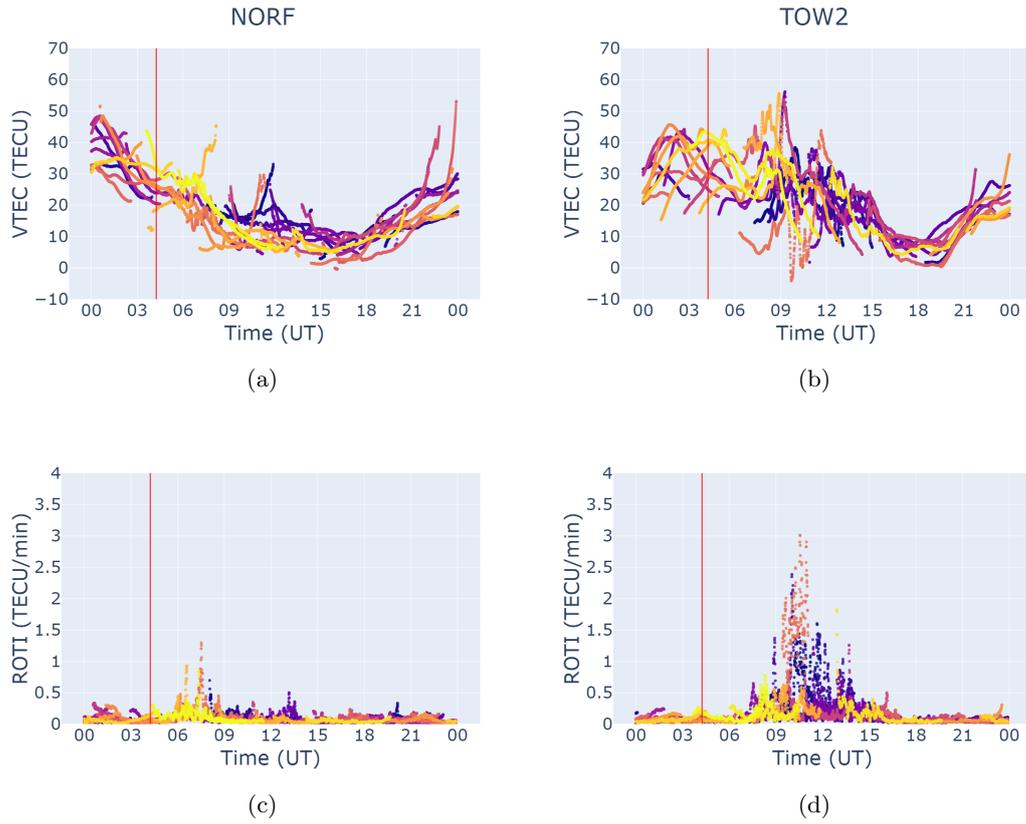


Figure 3: VTEC measured by (a) Norfolk Is. (NORF) and (b) Townsville (TOW2) GPS receivers with the different colors representing different satellite-to-ground links on January 15, 2022. The corresponding ROTI values for (c) NORF and (d) TOW2 stations. The red vertical lines indicate the time of the eruption.

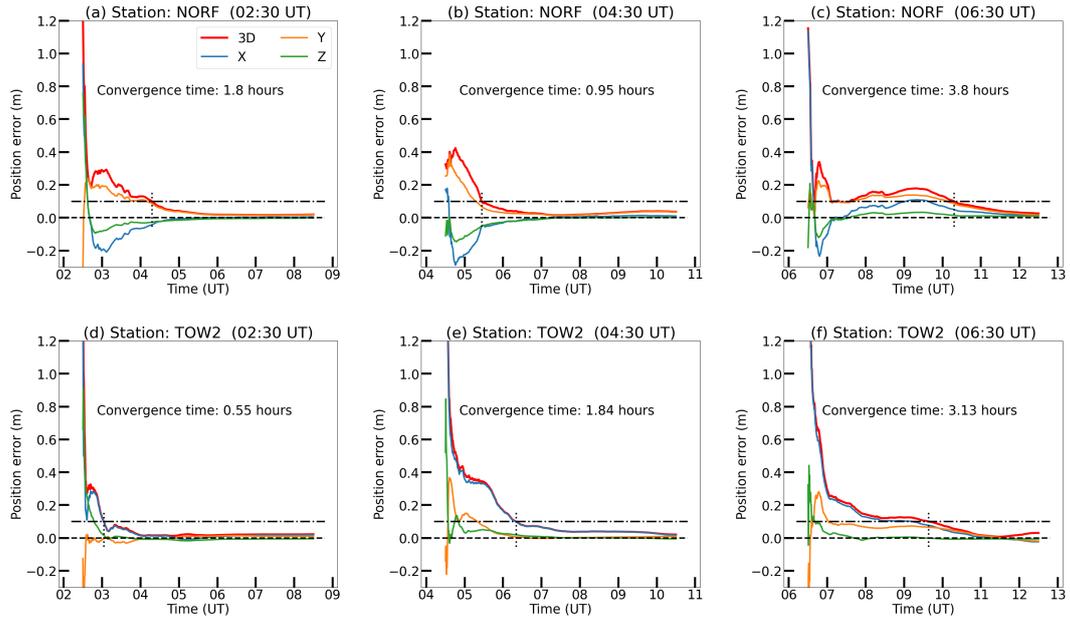


Figure 4: Examples of PPP convergence from a cold start using GPS data from the NORF (top row) and TOW2 (bottom row) commencing at 02:30 UT (left column), 04:30 UT (centre column) and 06:30 UT (right column) on January 15, 2022. Shown in each panel is the positioning error in X, Y and Z coordinates, in addition to the full 3-D position error. The dashed and dot-dashed horizon lines indicate position errors of 0 m and 10 cm, respectively. The moment when convergence was achieved and the associated convergence time are indicated in each panel by the vertical dotted lines and the text, respectively.

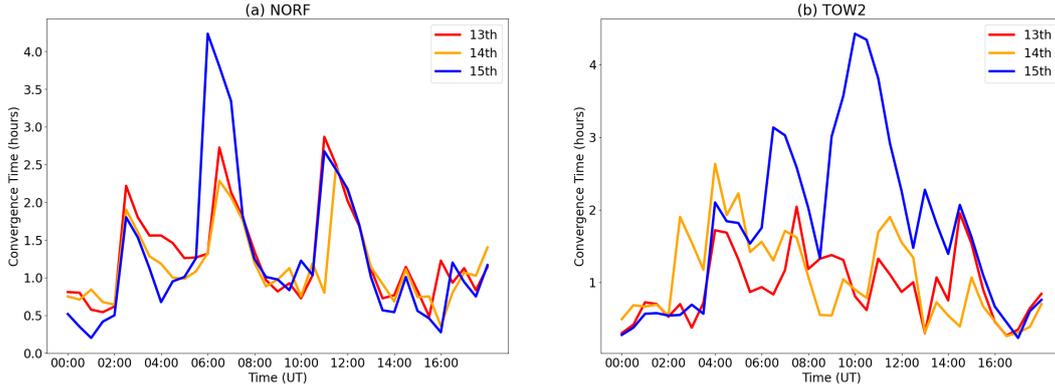


Figure 5: Diurnal variations in the PPP convergence time for the (a) NORF and (b) TOW2 GPS stations on the 13th, 14th and 15th of January, 2022.

251 4c). The TOW2 station shows a similar level of variation, ranging from 30 mins when  
 252 started from 02:30 UT (Fig. 4d) up to more than 3 hours when started from 06:30 UT  
 253 (Fig. 4c), which is  $\sim 30$  mins before the eruption effects reached Townsville. These re-  
 254 sults suggest a possible connection between the variable ionosphere in the wake of the  
 255 eruption and the PPP convergence times at Norfolk Island and Townsville.

256 The PPP convergence time is a complicated parameter that is dependent on the  
 257 number and geometry of available satellites, and the ability of the algorithm to model  
 258 and account for the signal errors including the atmospheric components caused by the  
 259 ionosphere and troposphere (e.g., Zumberge et al., 1997; Leick et al., 2015; Choy et al.,  
 260 2017; Kouba et al., 2017; Teunissen & Montenbruck, 2017). To further explore the possi-  
 261 ble impact of the disturbed ionosphere on the PPP convergence time, a numerical exper-  
 262 iment was performed using the GPS CORS across the Australian region. The GINAN  
 263 software was used to perform the PPP processing using data from staggered start times  
 264 throughout January 13-15, 2022, in 30-min steps. In other words, the experiment simu-  
 265 lated a series of ‘cold starts’ of the PPP algorithm every 30-mins during January 13-  
 266 15, 2022. At each time and for each station, the time it took for the PPP algorithm to  
 267 achieve convergence was recorded. Figure 5 shows how the convergence time for the NORF  
 268 (a) and TOW2 (b) stations varied throughout January 13-15. For January 13 and 14,  
 269 convergence times for both stations were typically on the order of 30 mins to 2.5 hours,  
 270 and appear to follow a similar diurnal variation for each station. On January 15, how-  
 271 ever, the convergence times are typically 2-4.5 hours, with some times showing signif-  
 272 icant differences compared to the January 13 and 14 values; particularly close to 6 UT  
 273 for NORF and between 6 and 12 UT for TOW2. For NORF, the maximum increase in  
 274 convergence time on January 15 (relative to January 14) was almost 3 hours at 6 UT.  
 275 For TOW2, the maximum increase in convergence time was 3.6 hours at 10:30 UT.

276 In order to investigate what impact, if any, the disturbed ionosphere had on GPS  
 277 PPP convergence time across the region, a reliable baseline was needed. At first glance  
 278 of Fig. 5, it appears that simply choosing the PPP convergence time from the day prior  
 279 as the baseline is a good option. However, some GPS stations showed significant differ-  
 280 ences between the diurnal variations in the convergence times between January 13 and  
 281 14; not shown here. In these cases, it is possible that the differences could be explained  
 282 by a difference in the handling of the tropospheric delays by the algorithm, as the geo-  
 283 magnetic activity level was quiet and the satellite geometry was very similar. There-  
 284 fore, in order to exclude times and GPS stations’ data for which the chosen baseline of  
 285 January 14 was not reliable, a simple selection criterion was used. Namely, if the con-

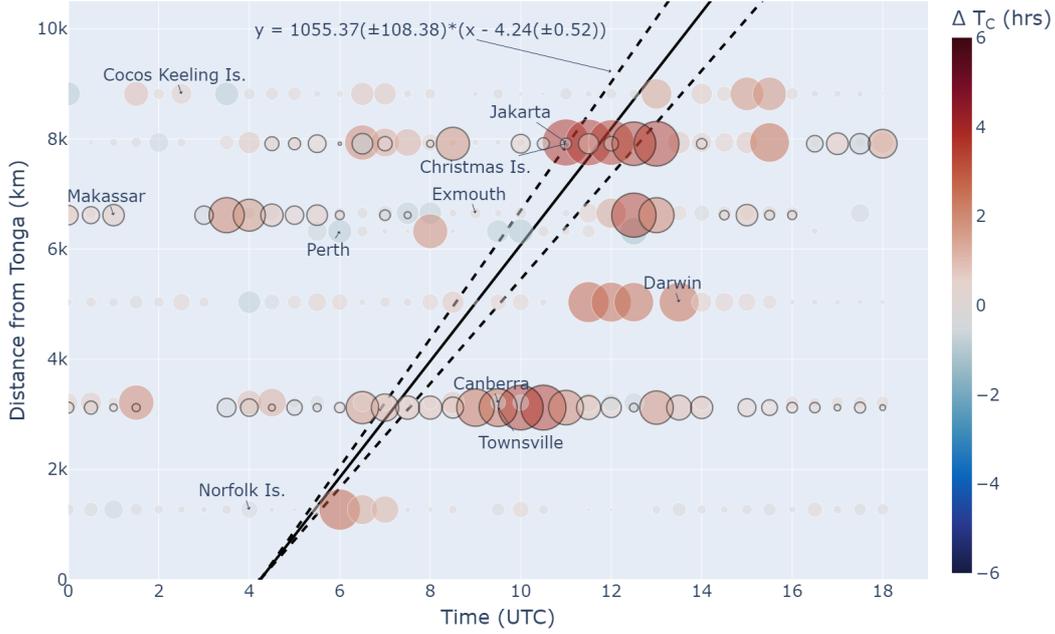


Figure 6: Distance from Tonga versus time colored according to the change in PPP convergence time,  $\Delta T_C$ , as determined using GPS stations located near the ionosonde stations in Figure 1, and two additional stations located at Makassar and Jakarta. The size of the circles also indicates the  $\Delta T_C$  value. For reference, the trend line representing the propagation of the ionospheric disturbances as identified in the ionosonde data in Fig. 2 is shown. Due to some stations' data overlapping in the plot, Townsville, Makassar and Christmas Island data points are shown with black outlines.

286 vergence time for a given moment in the day of January 14 was different compared to  
 287 January 13 by more than 1 hour, then that time was considered as a 'null' data point.  
 288 Otherwise, the difference between the convergence time on January 15 and January 14  
 289 for that time of day was taken as the change in convergence time due to the disturbed  
 290 ionosphere on the day of the Tonga volcano eruption,  $\Delta T_C$ .

291 Figure 6 shows how the change in convergence time varied throughout January 15  
 292 for GPS stations collocated with the ionosondes in Fig. 1, with two additional stations  
 293 at farther distances; Makassar and Jakarta in Indonesia. The size and color of the dots  
 294 indicates the change in convergence time,  $\Delta T_C$ . Also plotted is the line of best fit that  
 295 describes the propagation of the TIDs as detected by the ionosondes from Pradipta et  
 296 al. (2023).

297 Firstly, it is worth mentioning that prior to the arrival of the TIDs, some GPS stations  
 298 in Fig. 6 show some isolated increases in the convergence time on January 15 compared to  
 299 January 14. Makassar, Christmas Island, Canberra and Perth all show isolated  
 300 increases in the convergence time on the order of 2 hours before any ionospheric distur-  
 301 bances from the eruption arrive. However, most of the stations in Fig. 6 show signifi-  
 302 cant and lasting convergence time increases after the TIDs arrive at their respective loca-  
 303 tions. There are the exceptions of Canberra, Exmouth and Perth that do not observe  
 304 any clear convergence time increases. Interestingly, some GPS stations show convergence  
 305 time increases the moment the TIDs arrive, for example Norfolk Island and Jakarta, whereas  
 306 others observe their largest increases some hours later, for instance Darwin and Townsville.

307 In the Townsville ionosonde data shown in Fig. 2, it was noted that spread-F traces ac-  
 308 companied the arrival of TIDs at 7 UT (17 LT), with strong spread-F traces observed  
 309 during 11:24-15:00 UT (21:15-01:00 LT). It is interesting to note that the largest  $\Delta T_C$   
 310 increases for the TOW2/Townsville GPS station occurred for the algorithm start time  
 311 of 10:30 UT (20:30 LT), consistent with the beginning of the 3.5-hour period of strong  
 312 Spread-F activity detected by the Townsville ionosonde. Finally, it is also worth noting  
 313 that the magnitude of the  $\Delta T_C$  increases vary between the stations, from an increase of  
 314 2 hours observed at the Cocos Keeling Islands to an increase of 3.6 hours observed at  
 315 Townsville and 3.8 hours at Jakarta.

316 To further explore the impact of the disturbed ionosphere on the GPS PPP con-  
 317 vergence time across the region, Fig. 7 shows  $\Delta T_C$  for all of the GPS stations each hour  
 318 between 6 UT and 15 UT on January 15. At 6 and 7 UT, it can be seen that  $\Delta T_C$  val-  
 319 ues were mostly close to 0; some stations across South Eastern Australia show some el-  
 320 evated values at 6 UT, but most reduce to 0 by 7 UT. At 6 UT, the Norfolk Island sta-  
 321 tion – located to the southwest of Tonga approximately halfway to Australia – is already  
 322 showing elevated convergence times, as also shown in Fig. 6. By 8 UT, some stations  
 323 on Australia’s northeastern coast are showing some elevated convergence times, which  
 324 further increase at 9 UT. At this time, several stations show convergence time increases  
 325 of more than 5 hours. At 10 UT, the convergence time increases in the far-north Aus-  
 326 tralian region begin to decline as stations further south begin to show increases that reach  
 327 4 hours for some stations. The  $\Delta T_C$  values are still elevated at 11 UT in the east Aus-  
 328 tralian region, but some stations farther north and to the west are showing values close  
 329 to 5 hours. At 12 UT, the elevated  $\Delta T_C$  values across Australia’s north and across South-  
 330 east Asia remain at 3-4 hours as the eastern Australian stations approach 0. There is  
 331 a slight increase in  $\Delta T_C$  across Australia’s northeast once again at 13 UT to more than  
 332 3 hours, before almost all stations across the region approach  $\Delta T_C = 0$  by 15 UT. An  
 333 interactive map showing these data is included in the Supplementary Materials (S1).

334 In an effort to diagnose the physical phenomena that may be responsible for these  
 335 impacts on PPP convergence, ionospheric observations made by these GPS stations are  
 336 examined next. Figure 8 is the similar to Fig. 6, but the data points are colored accord-  
 337 ing to the maximum ROTI value detected for each GPS station in each 30-min inter-  
 338 val. Similar to the increased convergence times in Fig. 6, the ROTI values for most sta-  
 339 tions show a marked increase following the arrival of the TIDs from the volcano. Some  
 340 stations observed increased ROTI values some time after the primary TID arrivals, such  
 341 as Norfolk Island, Townsville and Darwin. The Perth and Exmouth stations generally  
 342 show very low ROTI values, as does the Canberra station; although at 08:30 UT, Can-  
 343 berra shows one elevated ROTI value. The agreement between the ROTI results in Fig.  
 344 8 and the convergence time results in Fig. 6 suggests a strong link between the two.

345 To explore the potential link between ROTI and PPP convergence time, Figure 9  
 346 is similar to Fig. 7 but shows the 30-min maximum ROTI value for each GPS receiver;  
 347 the full interactive map is available in the Supplementary Materials (S2). These maps  
 348 show a lot of similar trends to those in Fig. 7, particularly the increased ROTI values  
 349 over the northern Australian region compared to the south, and the increased values go-  
 350 ing into the Southeast Asian region. However, there is one notable difference between  
 351 Figs. 7 and 9; namely panel (d) that corresponds to 9 UT. While the change in conver-  
 352 gence times were low throughout South Eastern Australia, the ROTI values in Fig. 9  
 353 show that ROTI was actually quite high across this region. While the wave of increased  
 354 ROTI clearly sweeps across northern Australia, the increased ROTI across southern parts  
 355 of the country disappeared in place and did not make it to Western Australia.

356 Finally, VTEC observations made by the GPS stations are presented. Figure 10  
 357 is a series of maps of the 15-min averaged VTEC (i.e.,  $\langle VTEC \rangle$ ) for a select group of  
 358 GPS stations at latitudes between 24°S and 10°S and longitudes between 128°E and 152°E  
 359 (geographic). Each point represents the ionospheric pierce point (IPP) assuming an al-

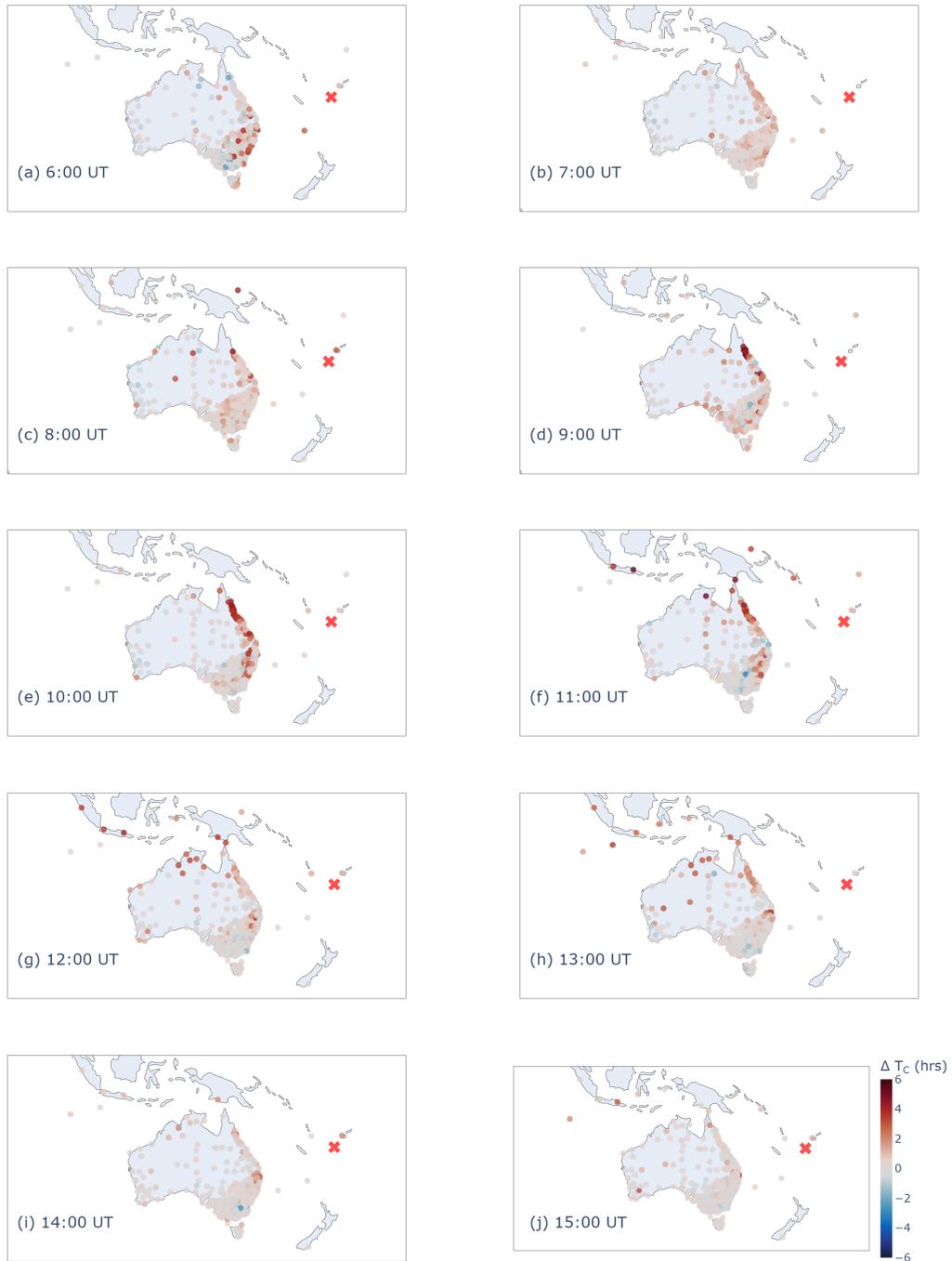


Figure 7: The change in PPP convergence time for GPS stations across the Australian/Southeast Asian region from (a) 6 UT until (j) 15 UT. The red cross indicates the location of Tonga. The full interactive map is included in the Supplementary Material S1.

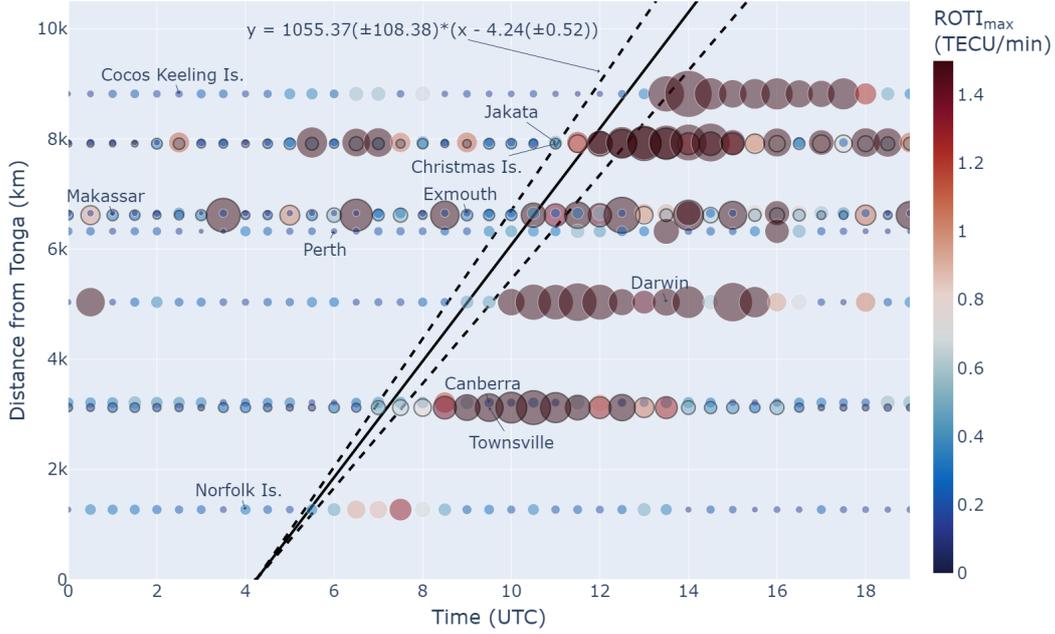


Figure 8: Same as Fig. 6, but colored according to the maximum ROTI observed for each station in every 30-min interval. Once again, Townsville, Makassar and Christmas Island data points are shown with black outlines.

360 titude of 400 km. A full interactive animation of this data sequence is included in the  
 361 Supplementary Materials (S3). Figure 10a shows the presence of a deep  $\langle VTEC \rangle$   
 362 valley over eastern of Papua New Guinea at 9 UT. By 09:30 UT (Fig. 10b) it appears as  
 363 though the  $\langle VTEC \rangle$  valley/depletion extends south to the Townsville station, and  
 364 remains almost in place for the remaining times plotted. The shape of the  $\langle VTEC \rangle$   
 365 depletion is quite pronounced in Fig. 10e (11 UT) as a thin dark blue feature with a north-  
 366 south alignment. Within this depletion, values as low as 6 TECU were observed, and were  
 367 as high as 48 TECU only 300 km to the west.

368 To track and measure the propagation of this  $\langle VTEC \rangle$  depletion feature, Figure.  
 369 11 shows the same data from Fig. 10, but it is restricted to spanning latitudes 17°S–14°S  
 370 and longitudes 140°E–155°E; i.e., the blue box in 10c. The depletion noted in Fig. 10  
 371 is quite clear in Fig. 11, along with some other  $\langle VTEC \rangle$  variations during this interval.  
 372 Fig. 11 also includes a manually plotted trendline that highlights the propagation of the  
 373 most pronounced depletion with longitude;  $y = 0.942x + 135.78$ . A noteworthy obser-  
 374 vation is that this ionospheric depletion is propagating to the east, albeit rather slowly.  
 375 As the trendline indicates, the depletion is propagating eastwards with a speed of  $\sim 0.9$   
 376 °/h, or equivalently  $\sim 30$  m/s (assuming 111 km/°).

#### 377 4 Discussion

378 In the results above, the ionospheric variability resulting from the Hunga Tonga  
 379 Volcano eruption across the Australian region was presented, followed by an investiga-  
 380 tion of the associated impact on PPP accuracy and convergence time. It was found that  
 381 positioning accuracy did not appear to be adversely impacted in the wake of the erup-  
 382 tion. Although, an experiment simulating ‘cold starts’ to the PPP algorithm showed that  
 383 convergence times were affected by the variable ionosphere in the wake of the eruption.

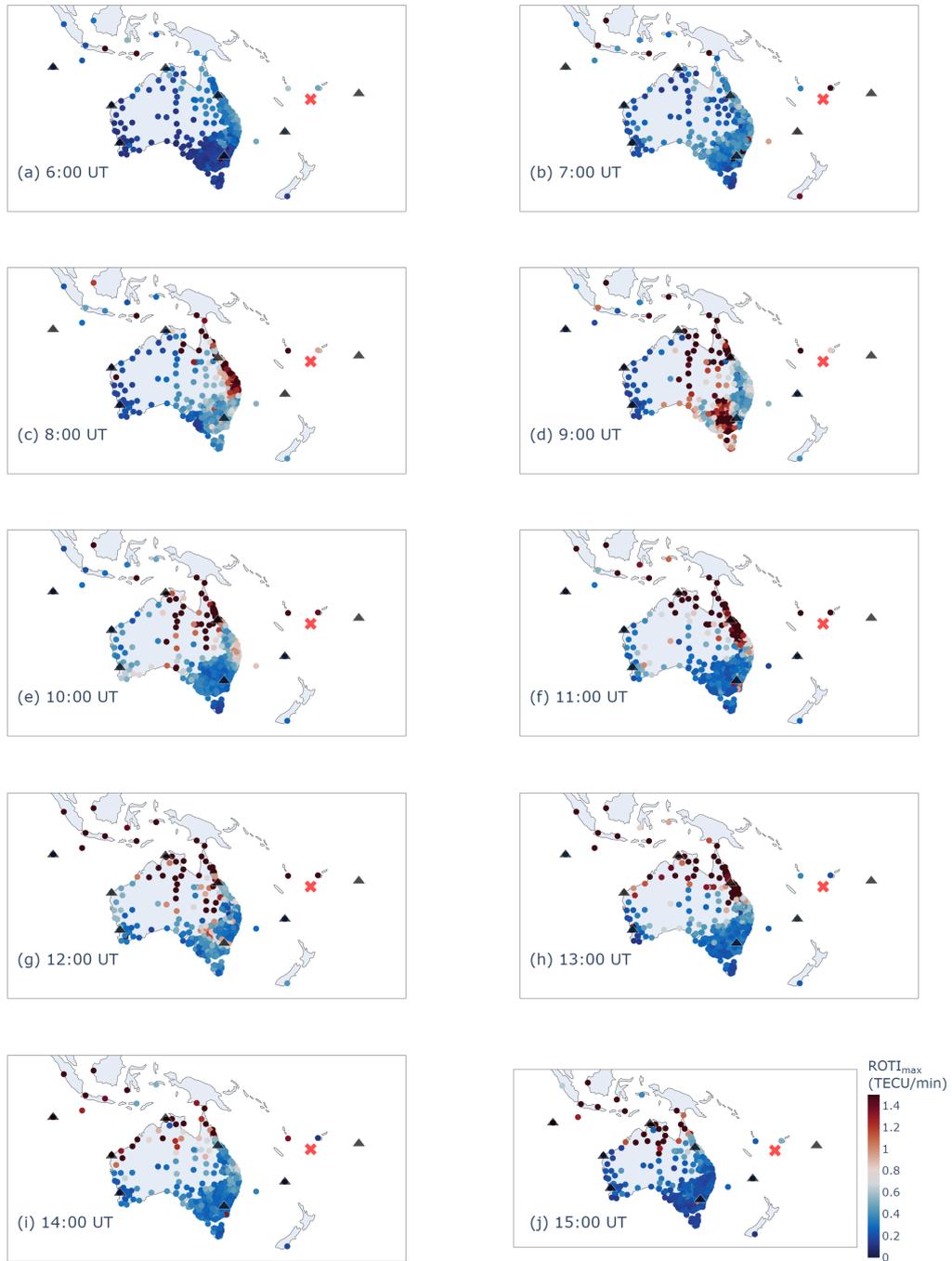


Figure 9: Same as Fig. 7, but showing the observed ROTI maximum in the 30-min following each time; i.e., panel a displays the maximum ROTI observed by each station between 6 and 6:30 UT. The full interactive map is included in the Supplementary Material S2.

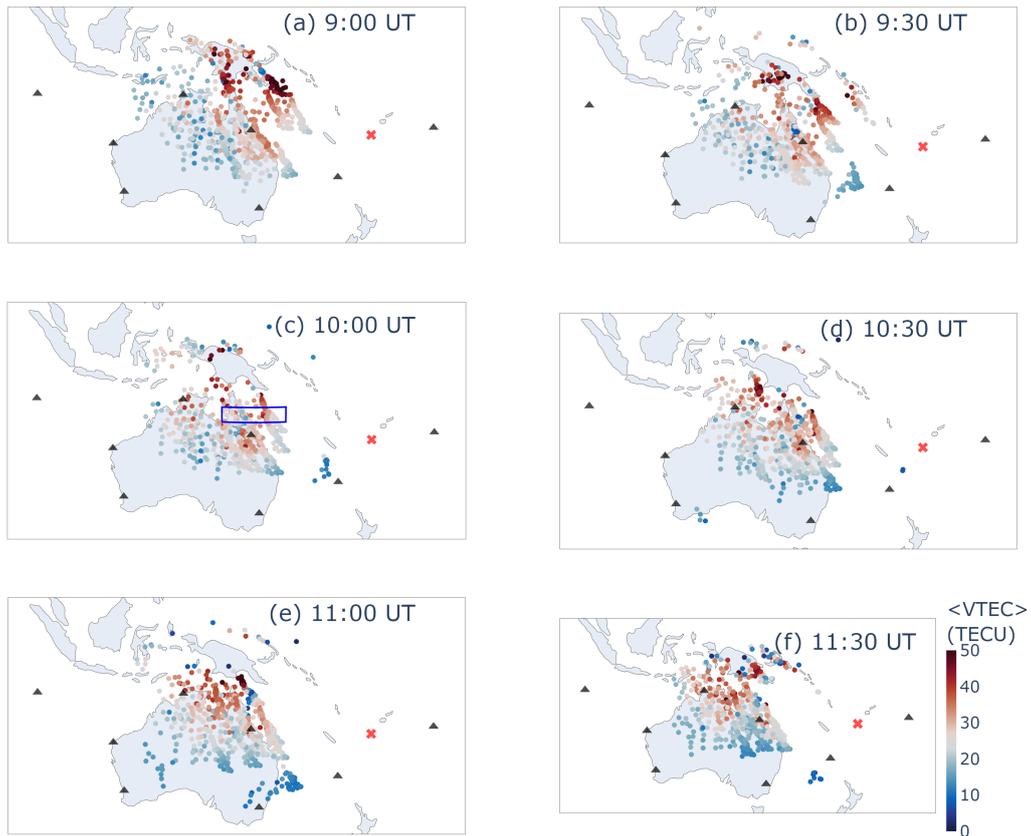


Figure 10: Maps showing the 15-min averaged VTEC (i.e.,  $\langle VTEC \rangle$ ) from 9 UT to 11:30 UT on January 15, 2022 for stations located between latitudes 24°S and 10°S and longitudes 128°E and 152°E (geographic). The blue box in panel (c) shows the range of latitudes and longitudes considered in the analysis to follow. The full interactive map is included in the Supplementary Material S3.

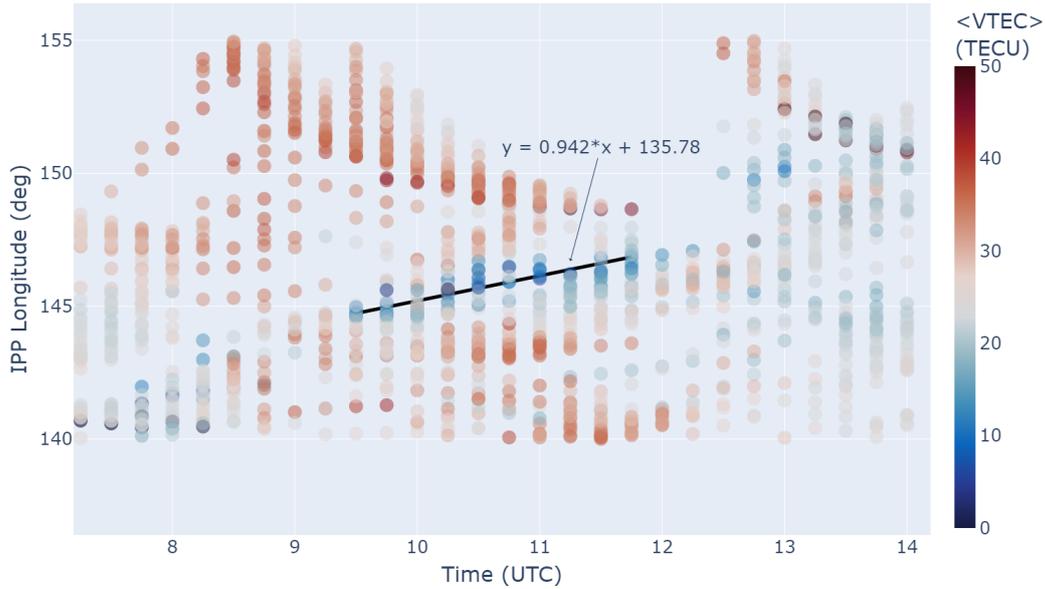


Figure 11: Ionospheric pierce point longitude versus time for data points for locations within the blue box in Fig. 10c – latitudes 17°S–14°S and longitudes 140°E–155°E – colored according to  $\langle VTEC \rangle$ .

384 Using the PPP convergence times from 13 and 14 January as a baseline, it was found  
 385 that stations across northern Australia and Southeast Asia experienced significant  
 386 convergence time increases on January 15, with some stations experiencing increases of more  
 387 than 5 hours. It was revealed that the stations experiencing the largest convergence time  
 388 increases measured extended periods of enhanced ROTI values. Examination of the 15-  
 389 min averaged VTEC data then revealed the presence of a significant depletion that  
 390 extended from Papua New Guinea into central eastern Queensland. A close examination  
 391 of this depletion revealed that it was propagating eastwards with a velocity of  $\sim 30$  m/s.

392 PPP accuracy is at the core of PPP as a usable application. As such, previous works  
 393 that have studied the impact of adverse ionospheric conditions on positioning focus on  
 394 the solution accuracy, and this has typically been done using ‘kinematic’ mode. For in-  
 395 stance, Poniatowski and Nykiel (2020) reported the degradation of PPP accuracy across  
 396 central Europe during the 2015 St. Patrick’s Day storm, with root mean square errors  
 397 of 0.58m, 0.37m and 0.26m in the vertical, north and east directions, respectively. Z. Yang  
 398 et al. (2020) also investigated the impact of the 2015 St. Patrick’s Day storm on PPP  
 399 accuracy, but their analysis included stations located around the world. They found that  
 400 intense auroral activity played a significant role in degrading the PPP solutions, with  
 401  $\sim 70\%$  of high-latitude stations suffering position errors of greater than 1 m. Importantly,  
 402 the most severe degradations in position accuracy were found to coincide with intense  
 403 ionospheric irregularities, as determined by the GPS ROTI parameter. Z. Yang et al.  
 404 (2020) also attributed some observed position accuracy degradation in low latitudes to  
 405 large-scale TIDs –  $\sim 1$  TECU in amplitude – that were generated by the storm. In the  
 406 present study, the PPP errors during the passage of the TIDs from the volcano eruption  
 407 – and the subsequent Spread-F activity – appeared to be smaller compared to that re-  
 408 ported by Z. Yang et al. (2020), despite the larger amplitude of the initial TIDs for this  
 409 event,  $\sim 3$  TECU (S.-R. Zhang et al., 2022). A similar analysis of the 2017 September  
 410 geomagnetic storms also revealed a strong relationship between enhanced ROTI and de-

411 graded PPP accuracy (Zakharenkova & Cherniak, 2021). In their analysis, it was found  
 412 that the 3-D errors rose to several meters due to the presence of small-scale ionospheric  
 413 irregularities. The lower errors in the present study can be understood by the fact that  
 414 ‘static mode’ was used in the positioning algorithm; i.e., the range of potential receiver  
 415 velocities passed to the filter via the ‘process noise’ parameter was set to zero. Interest-  
 416 ingly, the positioning errors in this study remained lower than 10 cm, provided that the  
 417 solution had already converged. One exception is shown in Fig. 5c, where the arrival of  
 418 TIDs at NORF delayed the convergence by  $\sim 2$  hours. While Figure 5 only shows a few  
 419 examples for two stations, their results largely reflect the results of the other stations con-  
 420 sidered in this analysis. Given the rather minimal impact of the volcano eruption on static-  
 421 mode PPP accuracy across the Australian region, the focus of this study shifted to the  
 422 potential impact on the PPP convergence time, which is a key limiting factor in the widespread  
 423 use of PPP (Choy et al., 2017).

424 Here, the investigation into the PPP convergence time made use of an experiment  
 425 in which ‘cold starts’ to the positioning algorithm were performed every 30-min during  
 426 January 15, 2022. Perhaps the most significant finding in this study is that the change  
 427 in convergence time – i.e., the difference between the 15th and 14th – reached 5 hours  
 428 or higher for some stations. Interestingly, these convergence time increases were consis-  
 429 tently observed over northern Australia, and were largely absent across central and south-  
 430 ern Australia, c.f. Figs. 6 and 7.

431 At first glance, the ionosonde data showed several inconsistencies with the PPP con-  
 432 vergence time results. All ionosonde stations observed clear TIDs associated with the  
 433 volcano eruption, followed by some period of Spread-F activity in the ionograms. As re-  
 434 ported by Pradipta et al. (2023), the propagation of the disturbances in these ionosonde  
 435 data agrees well with other studies that have investigated the TID propagation using GPS  
 436 TEC data (e.g., Themens et al., 2022; S.-R. Zhang et al., 2022). As observed by most  
 437 ionosonde stations, the TIDs immediately gave way to Spread-F activity, but in some  
 438 cases there was a significant delay between the arrival of the TIDs and the onset of Spread-  
 439 F. For example, the Townsville station TVL5F, Fig. 2d, showed Spread-F activity com-  
 440 mencing some 3 hrs after the initial TIDs passed, and the Learmonth station showed Spread-  
 441 F commencing more than 6 hours after the TID passage. The intensity of the Spread-  
 442 F also widely varied across the stations. Stations like Canberra, Townsville and Cocos  
 443 Keeling Islands observed quite strong and prolonged Spread-F, whereas one of the sta-  
 444 tions closest to the eruption, Norfolk Island, and another among the farthest away, Perth,  
 445 observed relatively weak Spread-F activity.

446 To better understand the differences in the detection of Spread-F across the Aus-  
 447 tralian region, it helps to be reminded about what the presence of TID signatures and  
 448 Spread-F in the ionograms indicates about the ionospheric plasma. The TID signatures  
 449 highlighted in Fig. 2 that are characterized by descending traces are clear signatures of  
 450 large-scale TIDs, on the order of 100s of kms (e.g., Cervera & Harris, 2014; Pederick et  
 451 al., 2017). While Spread-F observed at low latitudes is typically associated with EPBs  
 452 generated by the Generalized Rayleigh-Taylor instability during the post-sunset to lo-  
 453 cal midnight hours (Sultan, 1996; Burke et al., 2004; Kelley et al., 2006; Carter et al.,  
 454 2014, 2020, and references therein), mid-latitude Spread-F is generally considered to be  
 455 due to the Perkins instability (e.g., Perkins, 1973; Kelley & Fukao, 1991) or specular re-  
 456 flections from medium-scale TIDs with scales on the order of 10s of kms (e.g., G. Bow-  
 457 man, 1981; G. G. Bowman & Monro, 1988; G. G. Bowman, 1990). Interestingly, the re-  
 458 sults in Fig. 6 do not appear to show a clear and consistent relationship between the pres-  
 459 ence of large-scale TIDs or medium-scale TIDs/spread F and increased PPP convergence  
 460 times.

461 The GPS stations used in this study are not high-rate receivers that enable a de-  
 462 tailed analysis into the full spectrum of ionospheric waves (e.g., Cervera & Thomas, 2006;  
 463 van de Kamp & Cannon, 2009; Carrano et al., 2012). However, given the apparent mo-

464 tion of the GPS satellites above ground-based receivers and assuming typical ionospheric  
 465 wave phase speeds, the GPS ROTI parameter can provide insight into the presence of  
 466 ionospheric irregularities on the order of a few kms (e.g., Pi et al., 1997; Ma & Maruyama,  
 467 2006; Nishioka et al., 2008; Zou & Wang, 2009). Generally speaking, there is quite good  
 468 agreement between the impact on the convergence time and enhanced ROTI, by com-  
 469 paring Figs. 6 and 8 and comparing Figs. 7 and 9. Although, one noted difference was  
 470 the presence of enhanced ROTI that propagated over southeast Australia and then dis-  
 471 appeared. The propagation of this ROTI disturbance agrees well with the propagation  
 472 of the Lamb wave across the region reported by Aa et al. (2022) (see their Fig. 4). From  
 473 these observations, one can conclude that the large-scale TIDs caused by the eruption  
 474 had secondary small-scale irregularity generation all the way down to the scale of kms,  
 475 making them detectable using the ROTI parameter. Although, an open question that  
 476 arises is why did these secondary waves not continue to propagate across southern Aus-  
 477 tralia? In other words, what caused the small-scale irregularities over southeast Australia  
 478 at 9 UT in Fig. 9d to either no longer be generated – or be more heavily damped – by  
 479 10 UT in Fig. 9e, when the presence of the Lamb wave in the ROTI data was still clear  
 480 over northern Australia? The answer either lies in the differences in the effectiveness of  
 481 the energy cascade from large-scale to small-scale waves, or in the factors that control  
 482 the damping of small-scale structures between northern and southern Australia. In ei-  
 483 ther case, this is one interesting observation that requires further study and explanation.

484 In terms of the impact of ionospheric irregularities on the PPP convergence time,  
 485 these results suggest that increased ROTI – and therefore the presence of small-scale (i.e.,  
 486 km-scale) ionospheric irregularities – tends to prolong PPP convergence times only in  
 487 the event that the period of increased ROTI lasts for 30 mins or longer. This likely ex-  
 488 plains why stations in Australia’s northern region are clearly impacted compared to the  
 489 stations in southern and central Australia in Fig. 7. Therefore, an ongoing challenge for  
 490 the GPS PPP community is to effectively mitigate the adverse impacts of extended pe-  
 491 riods of small-scale ionospheric irregularity activity.

492 The next question is by what physical mechanism are these small-scale ionospheric  
 493 irregularities being generated in order to have such a detrimental impact on GPS PPP?  
 494 At first glance, the most heavily influenced region in northeastern Australia are too far  
 495 south for typical EPBs to be an obvious likely candidate. At 400 km above Townsville  
 496 (i.e., the IPP), the magnetic latitude is  $31.5^{\circ}\text{S}$ , and the magnetic field line from this lo-  
 497 cation maps to the equator at an altitude close to 2100 km; using a combination of the  
 498 International Geomagnetic Reference Field (Maus et al., 2005; IAGA, 2010) and the Al-  
 499 titude Adjustment Corrected Geomagnetic (AACGM) (Baker & Wing, 1989; Laundal  
 500 & Richmond, 2017) models. While some observations have been reported of such EPBs  
 501 in the past (Ma & Maruyama, 2006; Cherniak & Zakharenkova, 2016), disturbances this  
 502 large are particularly uncommon. Further, EPBs in the Southeast Asian/Australian lon-  
 503 gitude sector are not commonly observed during the months surrounding the December  
 504 solstice (e.g., Burke et al., 2004; Nishioka et al., 2008; E. Dao et al., 2011; Carter et al.,  
 505 2014, 2020). However, Aa et al. (2022) and Rajesh et al. (2022) recently reported ob-  
 506 servations of EPBs in the vicinity of Australia in the hours following the eruption. The  
 507 detection of strong spread-F traces by the Townsville ionosonde coincides with local times  
 508 that are typical of postsunset EPBs; i.e., 21-01 LT (see Fig. 2). Further, plots of VTEC  
 509 in Fig. 10 in this study agree well with the observations reported by Rajesh et al. (2022)  
 510 in revealing the presence of a deep depletion over northeastern Australia. Here, we mea-  
 511 sure that depletion to be 48 TECU on the ridges and 6 TECU in the valley. The prop-  
 512 agation of this depletion was determined in this study to be eastwards (Fig. 11), on the  
 513 order of 30 m/s. If this TEC depletion was a signature of TIDs propagating away from  
 514 the eruption, then the propagation would have been westwards. Moreover, if it were TIDs  
 515 from Tonga then the propagation would have been expected to much larger than 30 m/s.  
 516 Such low eastward propagation speeds are typical of the F-region dynamo in the dusk  
 517 sector, within which EPBs drift following their non-linear growth to the topside F re-

gion (e.g., Kelley, 2009; Chapagain et al., 2013, and references therein). Therefore, the local time of the Spread-F observed over Townsville, the VTEC depletion depth of 42 TECU and the associated eastward propagation speed in results presented here strongly indicate that the structure responsible for hours of increased ROTI, and subsequently increased PPP convergence times, in northeastern Australia is an EPB, in agreement with Rajesh et al. (2022).

It can be seen from the full animation of Fig. 11 (Supplementary Material (S3)) that data points associated with the depletion over northeastern Australia appeared as far south as 18°S, equivalent to 30°S magnetic latitude with an apex altitude of 1900 km over the magnetic equator. The authors are unaware of any other previously reported EPB disturbance that has been observed this far south over the Australian continent before, although similar-sized EPBs (referred to as ‘super bubbles’ due to their high altitude) have been reported using dense GPS stations across Japan (Ma & Maruyama, 2006). The EPBs that featured in that study would have also appeared over Australia due to magnetic conjugacy, but the relative sparsity of GPS stations across northern Australia did not allow clear EPB observations. However, the density of GPS stations across northern Australia has improved significantly during the solar cycle since Ma and Maruyama (2006)’s study. It is also worth highlighting that elements of Figs. 10 and 11 indicate the presence of TEC depletions farther to the south, but the sparsity of data at these locations makes a clear separation of TID disturbances from EPB-like disturbances difficult for this event. Further, the low density of GPS stations (and therefore, IPPs) across the rest of northern Australia makes a clear determination of whether depletions/EPBs were present difficult in this analysis. However, the impact on PPP convergence times across the rest of northern Australia and southeast Asia shown in Fig. 7 strongly suggests the ongoing presence of EPBs throughout the region during this event; a result in good agreement with the observations of Sun, Wenjie et al. (2022) and Aa et al. (2022) (in particular, the ICON plasma density depletions shown to the north of Australia in their Fig. 6). Some key insights into how TIDs from the eruption caused unseasonal EPB activity in the western Pacific sector can be gained from recent work by Huba et al. (2022). In Huba et al. (2022)’s modeling analysis, waves launched by the eruption caused significant perturbations in the zonal neutral wind in the equatorial plane, which effectively modified the equatorial upward plasma drift and gave rise to a huge EPB that spanned 30° in longitude (between 140-170°E). Huba et al. (2022) referred to this EPB as a ‘super EPB’ due to its large longitudinal extent. In addition, two very-high-altitude EPBs that reached 4000 km at ~ 155°E and 180° were also generated by their simulation (altitudes well above those that Ma and Maruyama (2006) referred to as ‘super bubbles’ in their study). While the 30°-wide EPB in Huba et al. (2022)’s simulation was not clearly observed in the present study (possibly due to its limited latitudinal extent), the observation of a longitudinally narrow, high-altitude and low-density structure over northeastern Australia is consistent with the very-high-altitude EPBs in Huba et al. (2022)’s study; noting the rather reasonable differences in observed versus modeled location. This structure over northeastern Australia thus fits the definition of a ‘super bubble’ according to Ma and Maruyama (2006) and is in good agreement with the results of Rajesh et al. (2022). The impacts on GPS reported in the present study across the rest of northern Australia indicate that further analysis of this event farther towards the west across southern and southeastern Asia is needed.

## 5 Conclusions

The Hunga Tonga Volcano eruption is a unique and complicated event that has provided an unprecedented opportunity to study how the ionosphere couples to the lower atmosphere. In this study, the impact of the eruption – via the ionospheric disturbances that it generated – on GPS PPP was investigated. While static-mode PPP accuracy itself did not appear to be heavily affected, the time taken for convergence to be achieved

570 was found to be significantly impacted. Across northern Australia the impact was particu-  
 571 larly clear, with some GPS stations near Townsville reporting increases in conver-  
 572 gence time of more than 5 hours. Long convergence times are limiting the widespread  
 573 use of PPP, so the Hunga Tonga Volcano eruption presents a good opportunity to re-  
 574 search the impacts and vulnerabilities of a variable ionosphere on PPP and to learn how  
 575 to mitigate/account for them.

576 In this study, it was found that large- and medium-scale TIDs from the eruption  
 577 were not the cause of convergence time increases, but it was the presence of small-scale  
 578 ionospheric structures on the order of a few kms, as determined using the ROTI param-  
 579 eter. Further, it was revealed that PPP was robust enough to endure some enhanced ROTI,  
 580 but not if the duration is longer than  $\sim 30$  mins. This observation effectively differen-  
 581 tiated the impact observed on PPP in northern and southern Australia; northern Aus-  
 582 tralia/southeast Asia experienced extended periods of increased ROTI and subsequently  
 583 increased convergence times, whereas southern Australia only experienced a short burst  
 584 of increased ROTI with marginal impact on convergence time.

585 The results of this study indicate that the ionospheric phenomenon responsible for  
 586 the presence of small-scale irregularities in southern and northern Australia was differ-  
 587 ent. In southern Australia, the small-scale irregularities were the result of secondary gen-  
 588 eration from the large- and medium-scale TIDs propagating away from the eruption. The  
 589 enhanced ROTI region propagated to the west in a manner consistent with the Lamb  
 590 waves previously reported. These small-scale structures were rather promptly damped  
 591 and did not propagate far beyond southeastern Australia. In northern Australia, the ex-  
 592 tended periods of enhanced ROTI were found to be due to the presence of at least one  
 593 ‘super bubble’ that was observed as far south as  $\sim 30^\circ\text{S}$  magnetic latitude, with an es-  
 594 timated apex altitude of  $\sim 1900$  km above the magnetic equator; this is the same ‘su-  
 595 per EPB’ recently reported by Rajesh et al. (2022). The VTEC data revealed that the  
 596 EPB was  $\sim 42$  TECU deep and  $\sim 300$  km across in longitude. Further, it was shown that  
 597 the EPB traveled eastwards at  $\sim 30$  m/s, consistent with the F-region dynamo speed  
 598 and direction. This super bubble is the largest/most southward-reaching EPB observed  
 599 over the Australian continent; only recently made possible due to GPS station deploy-  
 600 ments in the region.

601 The Hunga Tonga Volcano eruption stands as an excellent example of how iono-  
 602 spheric variability can adversely influence satellite-based precise positioning that is in-  
 603 creasingly heavily relied upon across many industries and sectors around the world.

## 604 **6 Open Research**

605 The Bureau of Meteorology Space Weather Services’ World Data Centre provides  
 606 ionosonde data via their website; [https://www.sws.bom.gov.au/World\\_Data\\_Centre](https://www.sws.bom.gov.au/World_Data_Centre). Geo-  
 607 science Australia provides GNSS data for all of the stations used in this analysis via their  
 608 Global Navigation Satellite System Data Centre; <https://gnss.ga.gov.au/>. Geoscience  
 609 Australia’s Ginan platform is also accessible via their github repository;  
 610 <https://geoscienceaustralia.github.io/ginan/>. The higher level analysis products, par-  
 611 ticularly the PPP convergence times that were calculated using Ginan, are available on  
 612 the Zenodo data repository (doi:10.5281/zenodo.7694409).

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1           **The ionospheric effects of the 2022 Hunga Tonga**  
2           **Volcano eruption and the associated impacts on GPS**  
3           **Precise Point Positioning across the Australian region**

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15           **Key Points:**

- 16           • PPP convergence times across northern Australia were significantly impacted in  
17           the hours after the eruption  
18           • Extended periods of enhanced ROTI were the cause of convergence time increases  
19           • A super equatorial plasma bubble, the largest observed over Australia, was respon-  
20           sible for PPP convergence time increases after the eruption

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## 21 Abstract

22 The Hunga Tonga Volcano eruption launched a myriad of atmospheric waves that  
 23 have been observed to travel around the world several times. These waves generated Trav-  
 24 eling Ionospheric Disturbances (TIDs) in the ionosphere, which are known to adversely  
 25 impact radio applications such as Global Navigation Satellite Systems (GNSS). One such  
 26 GNSS application is Precise Point Positioning (PPP), which can achieve cm-level accu-  
 27 racy using a single receiver, following a typical convergence time of 30 mins to 1 hour.  
 28 A network of ionosondes located throughout the Australian region were used in combi-  
 29 nation with GNSS receivers to explore the impacts of the Hunga-Tonga Volcano erup-  
 30 tion on the ionosphere and what subsequent impacts they had on PPP. It is shown that  
 31 PPP accuracy was not significantly impacted by the arrival of the TIDs and Spread-F,  
 32 provided that PPP convergence had already been achieved. However, when the PPP al-  
 33 gorithm was initiated from a cold start either shortly before or after the TID arrivals,  
 34 the convergence times were significantly longer. GNSS stations in northeastern Australia  
 35 experienced increases in convergence time of more than 5 hours. Further analysis reveals  
 36 increased convergence times to be caused by a super equatorial plasma bubble (EPB),  
 37 the largest observed over Australia to date. The EPB structure was found to be  $\sim 42$  TECU  
 38 deep and  $\sim 300$  km across, traveling eastwards at 30 m/s. The Hunga Tonga Volcano erup-  
 39 tion serves as an excellent example of how ionospheric variability can impact real-world  
 40 applications and the challenges associated with modeling the ionosphere to support GNSS.

## 41 Plain Language Summary

42 Global Navigation Satellite System (GNSS) applications permeate modern soci-  
 43 ety, with many industry sectors heavily relying on precision satellite positioning, nav-  
 44 igation and timing. Precise Point Positioning (PPP) is an advanced positioning technique  
 45 that can achieve cm-level accuracy without the need for nearby reference stations. How-  
 46 ever, the time that it takes for the PPP solution to ‘converge’ is typically in the range  
 47 of 10s of mins to hours, limiting the widespread uptake of PPP. There are numerous pre-  
 48 vious reports of waves and disturbances in the ionosphere, which are known to adversely  
 49 impact GNSS applications. In this study, the impact of the disturbances in the ionosphere  
 50 caused by the 2022 Hunga Tonga-Hunga Ha’apai Volcano eruption on PPP across the Aus-  
 51 tralian region is investigated. It is found that convergence times increased by more than  
 52 5 hours across northern Australia due to small-scale ionospheric turbulence. The source  
 53 of the turbulence was also found in this analysis to be due to a ‘super Equatorial Plasma  
 54 Bubble’ that persisted above northern Australia for several hours. This event serves as  
 55 an excellent example of how ionospheric disturbances can impact relied upon GNSS ap-  
 56 plications.

## 57 1 Introduction

58 At 04:14:45 UT on January 15, 2022, the Hunga Tonga-Hunga Ha’apai Volcano erupted  
 59 in what was one of the largest explosions on Earth in modern history (Matoza et al., 2022;  
 60 Wright et al., 2022). It has been estimated that the eruption released somewhere between  
 61  $3.7 \times 10^{16}$  to  $8.37 \times 10^{17}$  Joules (Wright et al., 2022; Astafyeva et al., 2022; Díaz & Rigby,  
 62 2022; Vergoz et al., 2022), making it comparable to the Krakatoa eruption in 1883 (Pyle,  
 63 2015). The explosion caused a tsunami that reached all sides of the Pacific Ocean with  
 64 an observed maximum wave-height of 3.4m on the Chilean shoreline (Carvajal et al., 2022).  
 65 The eruption was even audible as far away as Alaska, which is some 10,000 km away from  
 66 Tonga (Matoza et al., 2022). Not surprisingly, the eruption caused a myriad of waves  
 67 in the atmosphere and ionosphere (Themens et al., 2022; Wright et al., 2022; Aa et al.,  
 68 2022; Astafyeva et al., 2022; Ghent & Crowell, 2022; Maletckii & Astafyeva, 2022; Hong  
 69 et al., 2022) that were observed to encircle the Earth multiple times (S.-R. Zhang et al.,  
 70 2022; Matoza et al., 2022; Pradipta et al., 2023).

71 Waves in the ionosphere can be remotely detected using Global Navigation Satel-  
72 lite System (GNSS) signals in terms of the delay that the ionospheric plasma imparts  
73 upon the signals. The phase delay that is measured is related to the total electron con-  
74 tent (TEC) between the GNSS satellite and the receiver, with the majority of the TEC  
75 contribution coming from the ionosphere (Yizengaw et al., 2008). An expansive network  
76 of ground-based GNSS receivers therefore make it possible to geographically map iono-  
77 spheric TEC around the world (Mannucci et al., 1998) and track the propagation of any  
78 ionospheric waves (e.g., Otsuka et al., 2002; Kotake et al., 2006; Borries et al., 2009; Tsug-  
79 awa et al., 2011; Pradipta et al., 2016; H. Yang et al., 2017; Lay et al., 2018). To date,  
80 a range of ionospheric effects associated with the Hunga Tonga Volcano event have been  
81 published, and many of these studies have utilized the International GNSS Service net-  
82 work of receivers (Johnston et al., 2017).

83 Themens et al. (2022) and S.-R. Zhang et al. (2022) reported both large-scale and  
84 medium-scale traveling ionospheric disturbances (TIDs) propagating away from the vol-  
85 cano location. S.-R. Zhang et al. (2022) showed evidence of these TIDs continuing to prop-  
86 agate around the world for at least 4 days. Harding et al. (2022) and Le et al. (2022) ob-  
87 served changes to the equatorial electrojet that were caused by variations in the iono-  
88 spheric dynamo as a result of the eruption. Aa et al. (2022) reported a localized iono-  
89 spheric plasma depletion in the vicinity of the volcano and increased Equatorial Plasma  
90 Bubble (EPB) activity in the Asia-Oceania low-latitude region, which has been further  
91 supported by recent modeling efforts (Huba et al., 2022). Evidence of ‘super EPBs’ has  
92 been reported spanning across Chinese (Sun, Wenjie et al., 2022) and Japanese/Australian  
93 (Rajesh et al., 2022) longitude sectors following the eruption. Aa et al. (2022) also re-  
94 ported propagating ionospheric irregularities that exhibited a phase speed that matched  
95 the prevailing Lamb mode at  $\sim 315$  m/s.

96 GNSS Precise Point Positioning (PPP) is an advanced positioning technique that  
97 uses dual frequency observations made by a single receiver to achieve cm-level position-  
98 ing accuracy (Zumberge et al., 1997; Leick et al., 2015; Choy et al., 2017; Teunissen &  
99 Montenbruck, 2017). Over the previous few decades, the dominant method for achiev-  
100 ing cm-level accuracy has been relative positioning, in which a nearby accurately located  
101 reference station is used to determine the precise position of a ‘rover’ receiver (e.g., Odijk,  
102 2002; Hofmann-Wellenhof et al., 2007; Leick et al., 2015; Teunissen & Montenbruck, 2017).  
103 Utilizing a nearby reference station makes it possible to eliminate spatially correlated  
104 GNSS observations errors, such as the tropospheric and ionospheric delays imposed on  
105 the GNSS signals. Unsurprisingly, this precise positioning capability has now found its  
106 way into many applications across several major industries, including mining, agricul-  
107 ture and construction (e.g., Pérez-Ruiz et al., 2015; Choy et al., 2017; Woodgate et al.,  
108 2017; Rao et al., 2022). Over recent years, PPP is fast becoming the new global stan-  
109 dard for cm-level positioning applications due to its ability to model and account for GNSS  
110 observations errors without the need for a nearby reference station. However, the most  
111 significant drawback that is limiting the widespread uptake of PPP is the rather long  
112 convergence times, which are typically on the order of 10s of mins to hours (Bisnath &  
113 Gao, 2009; Van Bree et al., 2009; Choy et al., 2017). As a result, researchers have been  
114 investigating methods to significantly reduce PPP convergence times (e.g., Collins & Bis-  
115 nath, 2011; Collins et al., 2012; Geng & Bock, 2013; H. Zhang et al., 2013; Banville et  
116 al., 2014; Li et al., 2015; Duong, 2020). Given the significant influence of the ionosphere  
117 on GNSS signals, space weather poses a potential vulnerability to the PPP technique  
118 and its anticipated widespread up-take. As such, researchers have also been investigat-  
119 ing the impact of solar and geomagnetic activity on PPP (e.g., Luo et al., 2018; Poni-  
120 atowski & Nykiel, 2020; Z. Yang et al., 2020; Zha et al., 2021; Luo et al., 2022).

121 The Hunga Tonga Volcano eruption provides a unique opportunity to uncover new  
122 fundamental knowledge of the physics of atmosphere-ionosphere coupling. In addition,  
123 given the adverse influence of ionospheric variability on GNSS, this event can also pro-



Figure 1: The locations of the ionosondes used in this analysis. Tonga is indicated by the red cross.

124 vide insights into the impact of such ionospheric disturbances across broader society, par-  
 125 ticularly in industries that rely heavily on precise GNSS positioning. In this study, the  
 126 impact of the Tonga volcano eruption on Global Positioning System (GPS) PPP across  
 127 the Australian region is explored. Supporting this analysis are the data collected by ionoson-  
 128 des throughout the region that show both TID and Spread-F activity in the wake of the  
 129 volcano eruption. The PPP accuracy throughout the day of the eruption is explored, fol-  
 130 lowed by an analysis of the PPP convergence times from stations located throughout the  
 131 region. Finally, ionospheric observations using GPS receivers across the region are used  
 132 to investigate the physical mechanisms causing the disruptions identified in this study.

## 133 2 Data

134 In this analysis, the Australian Bureau of Meteorology Space Weather Services' ionosonde  
 135 data are used. Figure 1 shows the locations of the ionosonde stations (black triangles)  
 136 and Tonga volcano (red cross). The stations are Niue (NUE5F; 19.07°S, 190.07°E), Nor-  
 137 folk Island (NLK5D; 29.03°S, 167.97°E), Canberra (CBR5F; 35.32°S, 149.00°E), Townsville  
 138 (TVL5F; 19.63°S, 146.8°E), Darwin (DWN5D; 12.45°S, 130.95°E), Perth (PER5D; 31.94°S,  
 139 115.95°E), Learmonth (LEA5D; 22.25°S, 114.08°E) and Cocos Keeling Islands (CCK5F;  
 140 12.20°S, 96.80°E). Each ionosonde generates ionograms by sweeping through radio fre-  
 141 quencies between 2-22 MHz, transmitting and receiving ionospheric echoes that indicate  
 142 the electron density for a given virtual height determined by time-of-flight. In this anal-  
 143 ysis, the ionosonde data were used to indicate the presence of TIDs and spread-F traces.

144 The ionosonde data complement the primary dataset used in this study, namely  
 145 the GPS receiver data. GPS Continuously Operating Reference Station (CORS) data  
 146 from stations located across Australia, spanning into the Southeast Asian region and across  
 147 the South Pacific, were used in this analysis. While previous works have employed such  
 148 GPS CORS data to analyze ionospheric TEC fluctuations caused by the eruption, this  
 149 study focuses on the impact of this eruption on the GPS application of PPP.

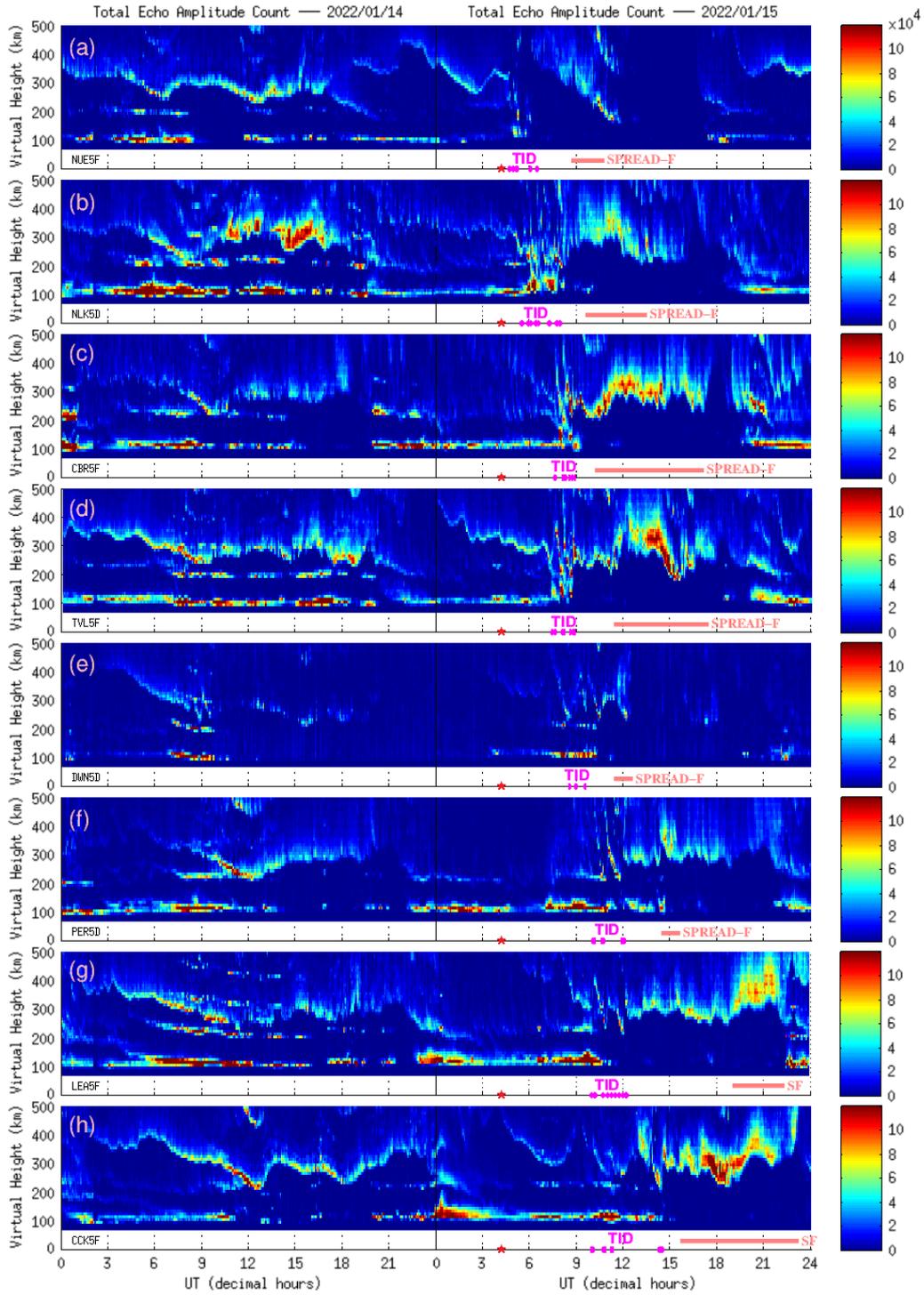


Figure 2: Range-time-intensity-style plots of ionospheric echoes received from ionosondes across the Australian region during January 14-15, 2022. Each panel corresponds to each station; (a) Niue (NUE5F), (b) Norfolk Island (NLK5D), (c) Canberra (CBR5F), (d) Townsville (TVL5F), (e) Darwin (DWN5D), (f) Perth (PER5D), (g) Learmonth (LEA5D) and (h) Cocos Keeling Islands (CCK5F). The color of each bin indicates the number of echoes received in that virtual height bin across all scanning frequencies, as indicated in the color bar on the right. The red stars indicate the time of the eruption; 04:14:45 UT. The pink dots (salmon bars) indicate the presences of TID (Spread-F) signatures.

150 The Geoscience Australia’s Ginan system was used for performing the PPP cal-  
 151 culations<sup>1</sup>. PPP is a high accuracy positioning method used to correct errors in GNSS  
 152 positioning based on the robust modelling and estimations of systematic errors in the  
 153 GNSS signals. The specific Ginan setup used in this analysis closely follows the Ginan  
 154 PPP example “Ex11”, in which the PPP is performed in ‘static mode’ using the ionosphere-  
 155 free combination<sup>2</sup>.

156 The carrier-phase measurements between the GPS satellites and receivers were used  
 157 to calculate the TEC along the signal path (i.e., the slant TEC, sTEC) following sim-  
 158 ilar methodology to Le Huy et al. (2016) and T. Dao et al. (2020). The possible jumps  
 159 in the sTEC values estimated from the carrier-phase measurements due to cycle slips were  
 160 eliminated by comparing them against the sTEC estimated from the pseudo-range mea-  
 161 surements that were smoothed by a fourth-degree polynomial approximation. A differ-  
 162 ence between the carrier-phase sTEC and the pseudo-range sTEC of more than 5 TECU  
 163 was taken to indicate an instrumental data jump/spike in the carrier-phase sTEC; in such  
 164 instances, the smoothed pseudo-range sTEC was used. The sTEC was then compared  
 165 with the Centre of Orbit Determination in Europe Global model to determine the tot-  
 166 al delay of device (including biases), using elevation angles above 30° to remove mul-  
 167 tipath effects. Finally, the resulting sTEC was converted to the vertical total electron  
 168 content (VTEC) observed at the pierce point of the ionosphere by using a single-layer  
 169 model (Klobuchar, 1986) and an assumed altitude of 400 km. The presence of scintillation-  
 170 causing small-scale ionospheric irregularities can be detected using the rate of TEC in-  
 171 dex (ROTI), which is defined as the 5-min standard deviation in the rate of change in  
 172 the sTEC for each satellite-to-ground link (Pi et al., 1997).

### 173 3 Results

174 This analysis begins with an overview of the ionospheric conditions before and af-  
 175 ter the arrival of the disturbances caused by the volcano eruption using the ionosonde  
 176 network in Fig. 1. First, it is worth mentioning that a minor geomagnetic storm occurred  
 177 in the late hours of January 14, with Dst reaching -91 nT due to a small, short-lived re-  
 178 current solar wind stream that had a minimal effect on the equatorial electric field (Le  
 179 et al., 2022). A useful way to analyze temporal changes in ionograms is by representing  
 180 the data in a format similar to a ‘range-time-intensity’ plot (Pradipta et al., 2015; Carter  
 181 et al., 2018; Currie et al., 2021). Instead of using total power, Pradipta et al. (2015) in-  
 182 tegrated over the dBm amplitudes across all sounding frequencies, effectively creating  
 183 a sum of digitized echoes. This methodology was adopted for the present study. Figure  
 184 2 shows these range-time-intensity-style plots for the ionosondes during January 14-15,  
 185 2022; the panels are ordered by the station great circle distance to Tonga, closest to far-  
 186 thest. The colors indicate the number of echoes received across all frequencies in each  
 187 virtual height bin (10 kms). The presence of TIDs and spread-F traces are indicated at  
 188 the bottom of each plot by the pink dots and salmon bars, respectively; these have been  
 189 determined by visual inspection of the ionograms.

190 Figure 2 shows a range of ionospheric conditions across January 14 and 15; e.g.,  
 191 sporadic E features at virtual heights of 100 km are present for all stations at various  
 192 times and F region traces that display typical altitude changes with time. TID signa-  
 193 tures are also clear in the data on January 15 for all stations, beginning at approx. 4:30  
 194 UT at Niue (top panel) and at approx. 10 UT at Cocos Keeling Islands, with the other  
 195 stations showing the TID signatures at times in between. All stations, with the excep-  
 196 tion of Darwin which was suffering some intermittent hardware issues during this period,  
 197 show the presence of spread-F traces in the hours following the TIDs, and in all cases,

<sup>1</sup> <https://geoscienceaustralia.github.io/ginan/>

<sup>2</sup> [https://geoscienceaustralia.github.io/ginan/codeDocs/Pea\\_8Configuration\\_8Examples.html](https://geoscienceaustralia.github.io/ginan/codeDocs/Pea_8Configuration_8Examples.html)

198 the spread-F traces are persistent for a number of hours. In particular, the TID activ-  
 199 ity at Norfolk Island begins shortly before 6 UT, with Spread-F present from 09:15 UT  
 200 until 13:30 UT. Thus, the occurrence of spread-F traces measured by the Norfolk Island  
 201 ionosonde span local times (LT) 21:15-01:30. At Townsville, the TIDs are first observed  
 202 close to 7 UT (17 LT) and bring with them batches of Spread-F. The Spread-F then in-  
 203 tensifies close to 11:24 UT (21:24 LT) and remains strong until 15 UT (1:00 LT), before  
 204 finally ceasing at 17:30 UT (03:30 LT).

205 Next, the GPS data from receivers close to the Norfolk Island and Townsville ionosonde  
 206 stations are examined as initial examples of the features and trends present throughout  
 207 the GPS data. Figure 3 shows the observed vertical TEC (VTEC, upper panels) and the  
 208 rate of TEC index (ROTI, lower panels) observed by the Norfolk Island (NORF, left)  
 209 and Townsville (TOW2, right) GPS receivers on January 15, 2022. The VTEC data for  
 210 both stations show the presence of TIDs that disrupt the diurnal pattern starting at  $\sim$   
 211 6 and 7 UT for NORF and TOW2, respectively, in good agreement with the ionosonde  
 212 observations in Fig. 2. In the NORF data, the TIDs are observed in the individual satel-  
 213 lite traces as rather small-amplitude wave structures compared to the variations over the  
 214 24-hour period. Slight differences in the timings of these structures (or phase progres-  
 215 sion) for different satellite links are due to the geographical spread of the ionospheric pierce  
 216 points and the motion of the TIDs over the station. In the TOW2 data, the TIDs are  
 217 much clearer as strong changes in VTEC for specific satellite links at different times. The  
 218 phase progression of the VTEC structures overhead are much clearer in the TOW2 data.  
 219 The lower panels reveal the times in which increased ROTI values were observed. For  
 220 Norfolk Island (Fig. 3c) ROTI increases shortly after the TIDs arrive and then decreases  
 221 back to low levels by 9 UT. Townsville also observed increased ROTI values once the main  
 222 TIDs arrived and remained high until  $\sim$  15 UT once the strong Spread-F ceases, c.f. Fig.  
 223 2.

224 Next, we illustrate how the GINAN PPP software achieves convergence, as defined  
 225 here by a 3-D position error of less than 10 cm. Figure 4 shows six examples of PPP con-  
 226 vergence on January 15, 2022 using data collected by the NORF (top row) and TOW2  
 227 (bottom row) receivers. Each panel represents a ‘cold start’ of the PPP algorithm start-  
 228 ing at a different time; 02:30 UT, 04:30 UT and 06:30 UT, as indicated in the subfigure  
 229 titles. The positioning error was calculated as the difference between the PPP estima-  
 230 tions and the known geodetic station locations. The position errors in the X, Y, and Z  
 231 directions, as well as the full 3-D position, are shown in each panel. The 10 cm position  
 232 error threshold is indicated by the dash-dot line, and zero error is indicated by the dashed  
 233 line. Also displayed in each panel is the time it took for the 3-D position error to reach  
 234 below 10 cm, indicated by the vertical dotted line; i.e., the convergence time.

235 Given the TID and subsequent Spread-F activity shown in the ionosonde data in  
 236 Fig. 2, one might expect that PPP errors for nearby GPS receivers would noticeably in-  
 237 crease during these periods. However, Fig. 4a shows that despite the onset of TIDs at  
 238 approx. 6 UT for the NORF station, the position error remained below the 10-cm thresh-  
 239 old. The same is true in Fig. 4d for the TOW2 station, for which the onset of TIDs was  
 240 approx. 7 UT. While Fig. 4b shows that the NORF station position error increased fol-  
 241 lowing the commencement of Spread-F close to 9 UT, it remained well below 10 cm. In-  
 242 terestingly, despite the commencement of rather strong Spread-F at 11:24 UT at TOW2,  
 243 Fig. 4f shows no significant increase in the position error. This is likely the result of run-  
 244 ning the PPP algorithm in ‘static mode’. From these results, one can conclude that the  
 245 static-mode PPP algorithm was robust enough to remain at cm-level accuracy for these  
 246 stations throughout the turbulent ionospheric conditions caused by the Tonga eruption,  
 247 provided that convergence had already been achieved.

248 Figure 4 also illustrates how the convergence time is dependent on the start time  
 249 of the algorithm. For the NORF station, the convergence time ranges from less than 1  
 250 hour beginning at 04:30 UT (Fig. 4b) to almost 4 hours when started from 06:30 UT (Fig.

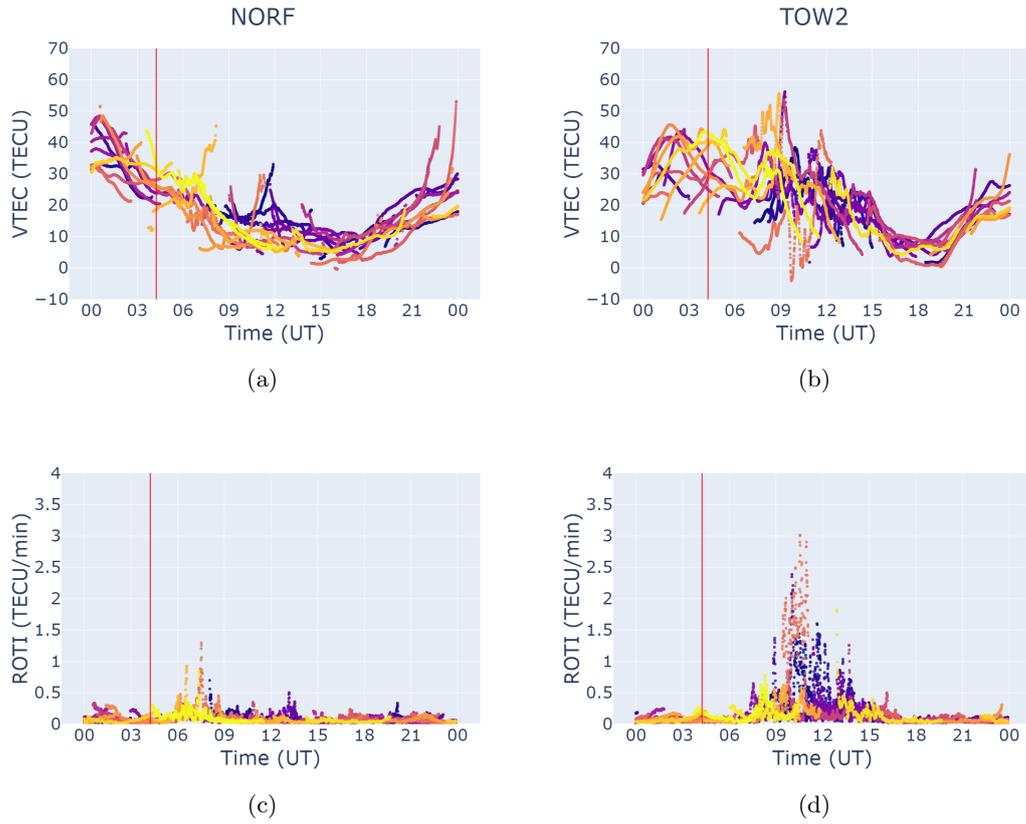


Figure 3: VTEC measured by (a) Norfolk Is. (NORF) and (b) Townsville (TOW2) GPS receivers with the different colors representing different satellite-to-ground links on January 15, 2022. The corresponding ROTI values for (c) NORF and (d) TOW2 stations. The red vertical lines indicate the time of the eruption.

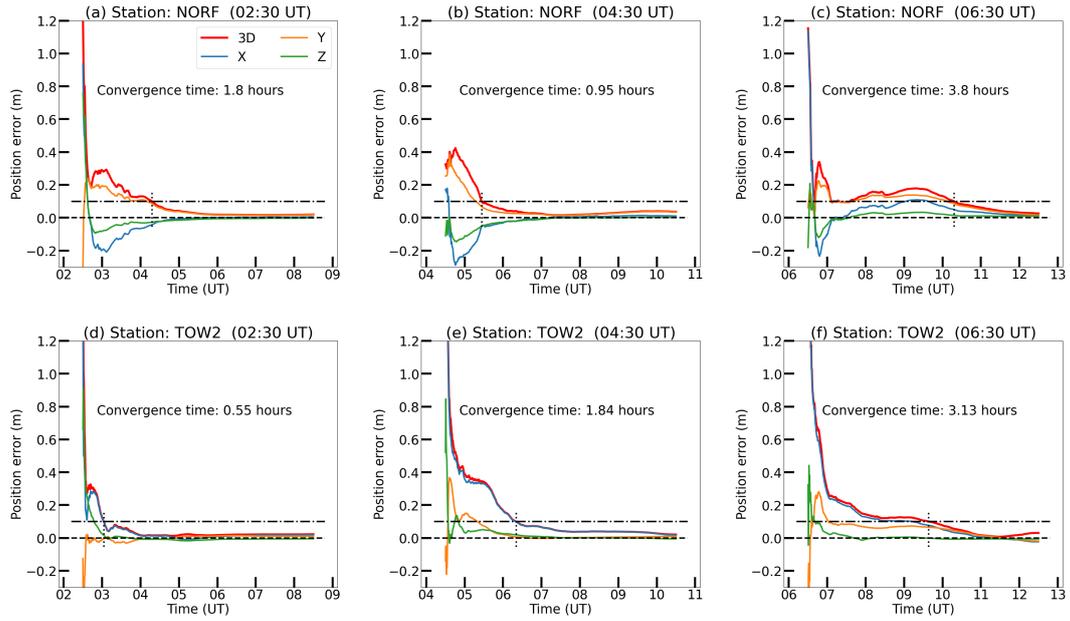


Figure 4: Examples of PPP convergence from a cold start using GPS data from the NORF (top row) and TOW2 (bottom row) commencing at 02:30 UT (left column), 04:30 UT (centre column) and 06:30 UT (right column) on January 15, 2022. Shown in each panel is the positioning error in X, Y and Z coordinates, in addition to the full 3-D position error. The dashed and dot-dashed horizontal lines indicate position errors of 0 m and 10 cm, respectively. The moment when convergence was achieved and the associated convergence time are indicated in each panel by the vertical dotted lines and the text, respectively.

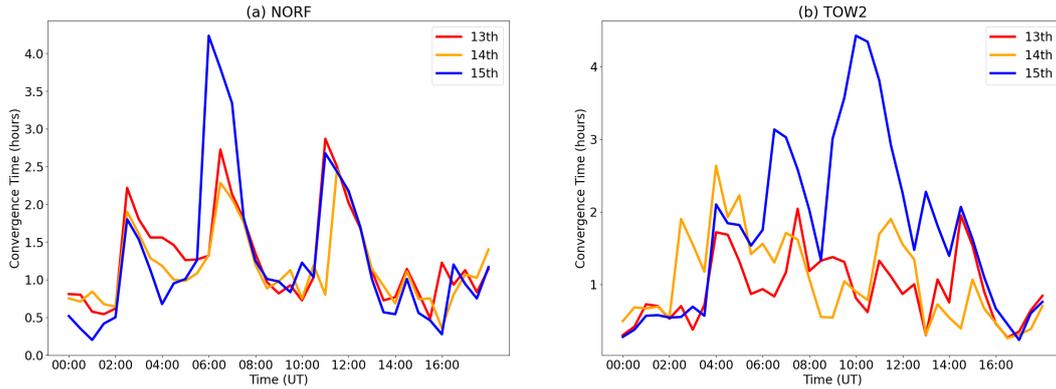


Figure 5: Diurnal variations in the PPP convergence time for the (a) NORF and (b) TOW2 GPS stations on the 13th, 14th and 15th of January, 2022.

251 4c). The TOW2 station shows a similar level of variation, ranging from 30 mins when  
 252 started from 02:30 UT (Fig. 4d) up to more than 3 hours when started from 06:30 UT  
 253 (Fig. 4c), which is  $\sim 30$  mins before the eruption effects reached Townsville. These re-  
 254 sults suggest a possible connection between the variable ionosphere in the wake of the  
 255 eruption and the PPP convergence times at Norfolk Island and Townsville.

256 The PPP convergence time is a complicated parameter that is dependent on the  
 257 number and geometry of available satellites, and the ability of the algorithm to model  
 258 and account for the signal errors including the atmospheric components caused by the  
 259 ionosphere and troposphere (e.g., Zumberge et al., 1997; Leick et al., 2015; Choy et al.,  
 260 2017; Kouba et al., 2017; Teunissen & Montenbruck, 2017). To further explore the possi-  
 261 ble impact of the disturbed ionosphere on the PPP convergence time, a numerical exper-  
 262 iment was performed using the GPS CORS across the Australian region. The GINAN  
 263 software was used to perform the PPP processing using data from staggered start times  
 264 throughout January 13-15, 2022, in 30-min steps. In other words, the experiment simu-  
 265 lated a series of ‘cold starts’ of the PPP algorithm every 30-mins during January 13-  
 266 15, 2022. At each time and for each station, the time it took for the PPP algorithm to  
 267 achieve convergence was recorded. Figure 5 shows how the convergence time for the NORF  
 268 (a) and TOW2 (b) stations varied throughout January 13-15. For January 13 and 14,  
 269 convergence times for both stations were typically on the order of 30 mins to 2.5 hours,  
 270 and appear to follow a similar diurnal variation for each station. On January 15, how-  
 271 ever, the convergence times are typically 2-4.5 hours, with some times showing signif-  
 272 icant differences compared to the January 13 and 14 values; particularly close to 6 UT  
 273 for NORF and between 6 and 12 UT for TOW2. For NORF, the maximum increase in  
 274 convergence time on January 15 (relative to January 14) was almost 3 hours at 6 UT.  
 275 For TOW2, the maximum increase in convergence time was 3.6 hours at 10:30 UT.

276 In order to investigate what impact, if any, the disturbed ionosphere had on GPS  
 277 PPP convergence time across the region, a reliable baseline was needed. At first glance  
 278 of Fig. 5, it appears that simply choosing the PPP convergence time from the day prior  
 279 as the baseline is a good option. However, some GPS stations showed significant differ-  
 280 ences between the diurnal variations in the convergence times between January 13 and  
 281 14; not shown here. In these cases, it is possible that the differences could be explained  
 282 by a difference in the handling of the tropospheric delays by the algorithm, as the geo-  
 283 magnetic activity level was quiet and the satellite geometry was very similar. There-  
 284 fore, in order to exclude times and GPS stations’ data for which the chosen baseline of  
 285 January 14 was not reliable, a simple selection criterion was used. Namely, if the con-

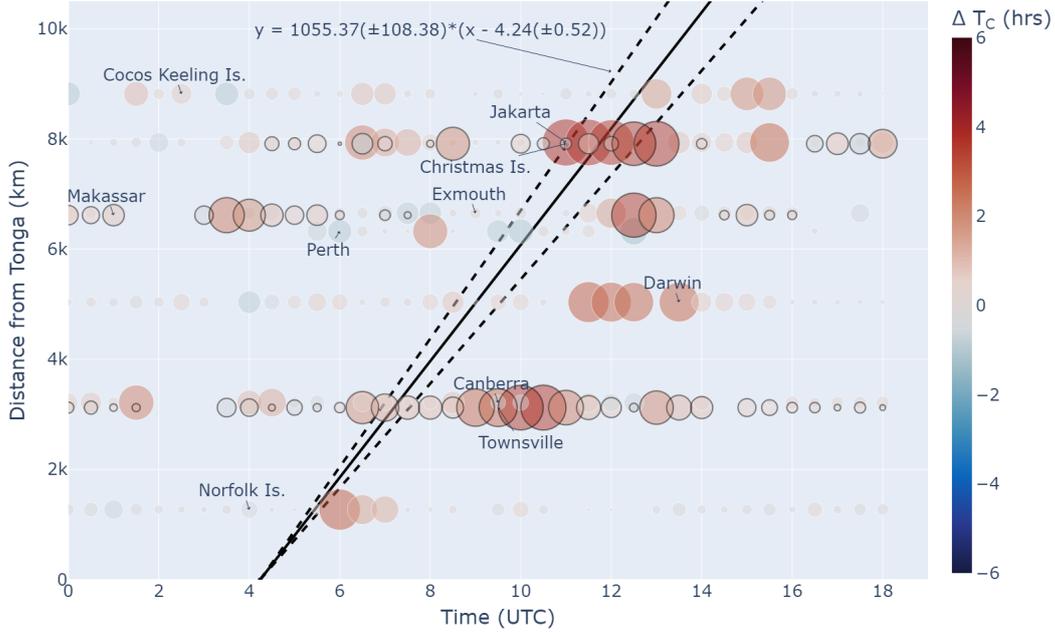


Figure 6: Distance from Tonga versus time colored according to the change in PPP convergence time,  $\Delta T_C$ , as determined using GPS stations located near the ionosonde stations in Figure 1, and two additional stations located at Makassar and Jakarta. The size of the circles also indicates the  $\Delta T_C$  value. For reference, the trend line representing the propagation of the ionospheric disturbances as identified in the ionosonde data in Fig. 2 is shown. Due to some stations' data overlapping in the plot, Townsville, Makassar and Christmas Island data points are shown with black outlines.

286 vergence time for a given moment in the day of January 14 was different compared to  
 287 January 13 by more than 1 hour, then that time was considered as a 'null' data point.  
 288 Otherwise, the difference between the convergence time on January 15 and January 14  
 289 for that time of day was taken as the change in convergence time due to the disturbed  
 290 ionosphere on the day of the Tonga volcano eruption,  $\Delta T_C$ .

291 Figure 6 shows how the change in convergence time varied throughout January 15  
 292 for GPS stations collocated with the ionosondes in Fig. 1, with two additional stations  
 293 at farther distances; Makassar and Jakarta in Indonesia. The size and color of the dots  
 294 indicates the change in convergence time,  $\Delta T_C$ . Also plotted is the line of best fit that  
 295 describes the propagation of the TIDs as detected by the ionosondes from Pradipta et  
 296 al. (2023).

297 Firstly, it is worth mentioning that prior to the arrival of the TIDs, some GPS stations  
 298 in Fig. 6 show some isolated increases in the convergence time on January 15 compared to  
 299 January 14. Makassar, Christmas Island, Canberra and Perth all show isolated  
 300 increases in the convergence time on the order of 2 hours before any ionospheric distur-  
 301 bances from the eruption arrive. However, most of the stations in Fig. 6 show signifi-  
 302 cant and lasting convergence time increases after the TIDs arrive at their respective loca-  
 303 tions. There are the exceptions of Canberra, Exmouth and Perth that do not observe  
 304 any clear convergence time increases. Interestingly, some GPS stations show convergence  
 305 time increases the moment the TIDs arrive, for example Norfolk Island and Jakarta, whereas  
 306 others observe their largest increases some hours later, for instance Darwin and Townsville.

307 In the Townsville ionosonde data shown in Fig. 2, it was noted that spread-F traces ac-  
 308 companied the arrival of TIDs at 7 UT (17 LT), with strong spread-F traces observed  
 309 during 11:24-15:00 UT (21:15-01:00 LT). It is interesting to note that the largest  $\Delta T_C$   
 310 increases for the TOW2/Townsville GPS station occurred for the algorithm start time  
 311 of 10:30 UT (20:30 LT), consistent with the beginning of the 3.5-hour period of strong  
 312 Spread-F activity detected by the Townsville ionosonde. Finally, it is also worth noting  
 313 that the magnitude of the  $\Delta T_C$  increases vary between the stations, from an increase of  
 314 2 hours observed at the Cocos Keeling Islands to an increase of 3.6 hours observed at  
 315 Townsville and 3.8 hours at Jakarta.

316 To further explore the impact of the disturbed ionosphere on the GPS PPP con-  
 317 vergence time across the region, Fig. 7 shows  $\Delta T_C$  for all of the GPS stations each hour  
 318 between 6 UT and 15 UT on January 15. At 6 and 7 UT, it can be seen that  $\Delta T_C$  val-  
 319 ues were mostly close to 0; some stations across South Eastern Australia show some el-  
 320 evated values at 6 UT, but most reduce to 0 by 7 UT. At 6 UT, the Norfolk Island sta-  
 321 tion – located to the southwest of Tonga approximately halfway to Australia – is already  
 322 showing elevated convergence times, as also shown in Fig. 6. By 8 UT, some stations  
 323 on Australia’s northeastern coast are showing some elevated convergence times, which  
 324 further increase at 9 UT. At this time, several stations show convergence time increases  
 325 of more than 5 hours. At 10 UT, the convergence time increases in the far-north Aus-  
 326 tralian region begin to decline as stations further south begin to show increases that reach  
 327 4 hours for some stations. The  $\Delta T_C$  values are still elevated at 11 UT in the east Aus-  
 328 tralian region, but some stations farther north and to the west are showing values close  
 329 to 5 hours. At 12 UT, the elevated  $\Delta T_C$  values across Australia’s north and across South-  
 330 east Asia remain at 3-4 hours as the eastern Australian stations approach 0. There is  
 331 a slight increase in  $\Delta T_C$  across Australia’s northeast once again at 13 UT to more than  
 332 3 hours, before almost all stations across the region approach  $\Delta T_C = 0$  by 15 UT. An  
 333 interactive map showing these data is included in the Supplementary Materials (S1).

334 In an effort to diagnose the physical phenomena that may be responsible for these  
 335 impacts on PPP convergence, ionospheric observations made by these GPS stations are  
 336 examined next. Figure 8 is the similar to Fig. 6, but the data points are colored accord-  
 337 ing to the maximum ROTI value detected for each GPS station in each 30-min inter-  
 338 val. Similar to the increased convergence times in Fig. 6, the ROTI values for most sta-  
 339 tions show a marked increase following the arrival of the TIDs from the volcano. Some  
 340 stations observed increased ROTI values some time after the primary TID arrivals, such  
 341 as Norfolk Island, Townsville and Darwin. The Perth and Exmouth stations generally  
 342 show very low ROTI values, as does the Canberra station; although at 08:30 UT, Can-  
 343 berra shows one elevated ROTI value. The agreement between the ROTI results in Fig.  
 344 8 and the convergence time results in Fig. 6 suggests a strong link between the two.

345 To explore the potential link between ROTI and PPP convergence time, Figure 9  
 346 is similar to Fig. 7 but shows the 30-min maximum ROTI value for each GPS receiver;  
 347 the full interactive map is available in the Supplementary Materials (S2). These maps  
 348 show a lot of similar trends to those in Fig. 7, particularly the increased ROTI values  
 349 over the northern Australian region compared to the south, and the increased values go-  
 350 ing into the Southeast Asian region. However, there is one notable difference between  
 351 Figs. 7 and 9; namely panel (d) that corresponds to 9 UT. While the change in conver-  
 352 gence times were low throughout South Eastern Australia, the ROTI values in Fig. 9  
 353 show that ROTI was actually quite high across this region. While the wave of increased  
 354 ROTI clearly sweeps across northern Australia, the increased ROTI across southern parts  
 355 of the country disappeared in place and did not make it to Western Australia.

356 Finally, VTEC observations made by the GPS stations are presented. Figure 10  
 357 is a series of maps of the 15-min averaged VTEC (i.e.,  $\langle VTEC \rangle$ ) for a select group of  
 358 GPS stations at latitudes between 24°S and 10°S and longitudes between 128°E and 152°E  
 359 (geographic). Each point represents the ionospheric pierce point (IPP) assuming an al-

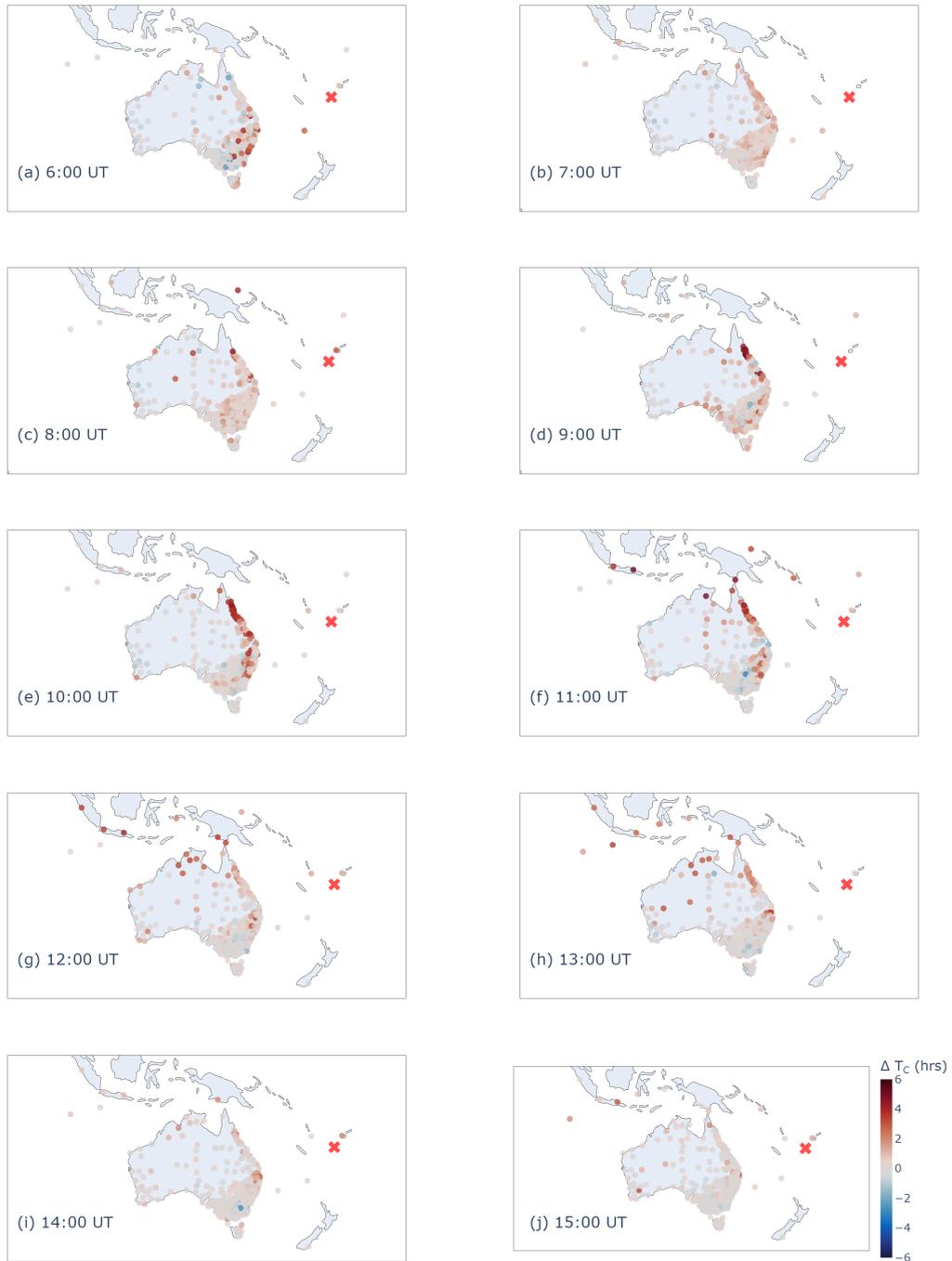


Figure 7: The change in PPP convergence time for GPS stations across the Australian/Southeast Asian region from (a) 6 UT until (j) 15 UT. The red cross indicates the location of Tonga. The full interactive map is included in the Supplementary Material S1.

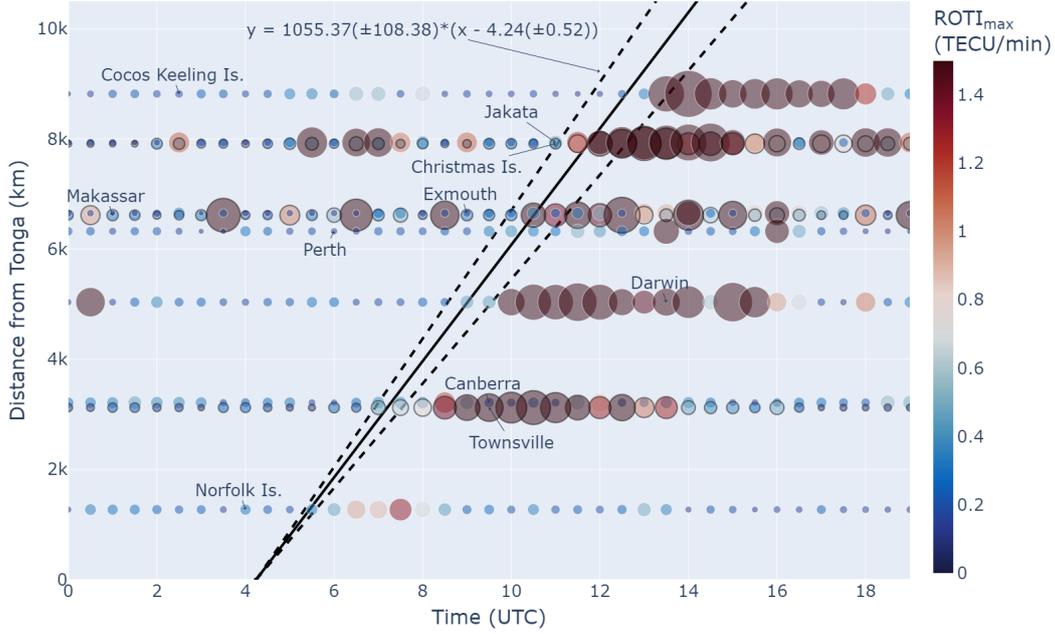


Figure 8: Same as Fig. 6, but colored according to the maximum ROTI observed for each station in every 30-min interval. Once again, Townsville, Makassar and Christmas Island data points are shown with black outlines.

360 titude of 400 km. A full interactive animation of this data sequence is included in the  
 361 Supplementary Materials (S3). Figure 10a shows the presence of a deep  $\langle VTEC \rangle$   
 362 valley over eastern of Papua New Guinea at 9 UT. By 09:30 UT (Fig. 10b) it appears as  
 363 though the  $\langle VTEC \rangle$  valley/depletion extends south to the Townsville station, and  
 364 remains almost in place for the remaining times plotted. The shape of the  $\langle VTEC \rangle$   
 365 depletion is quite pronounced in Fig. 10e (11 UT) as a thin dark blue feature with a north-  
 366 south alignment. Within this depletion, values as low as 6 TECU were observed, and were  
 367 as high as 48 TECU only 300 km to the west.

368 To track and measure the propagation of this  $\langle VTEC \rangle$  depletion feature, Figure.  
 369 11 shows the same data from Fig. 10, but it is restricted to spanning latitudes 17°S–14°S  
 370 and longitudes 140°E–155°E; i.e., the blue box in 10c. The depletion noted in Fig. 10  
 371 is quite clear in Fig. 11, along with some other  $\langle VTEC \rangle$  variations during this interval.  
 372 Fig. 11 also includes a manually plotted trendline that highlights the propagation of the  
 373 most pronounced depletion with longitude;  $y = 0.942x + 135.78$ . A noteworthy obser-  
 374 vation is that this ionospheric depletion is propagating to the east, albeit rather slowly.  
 375 As the trendline indicates, the depletion is propagating eastwards with a speed of  $\sim 0.9$   
 376 °/h, or equivalently  $\sim 30$  m/s (assuming 111 km/°).

#### 377 4 Discussion

378 In the results above, the ionospheric variability resulting from the Hunga Tonga  
 379 Volcano eruption across the Australian region was presented, followed by an investiga-  
 380 tion of the associated impact on PPP accuracy and convergence time. It was found that  
 381 positioning accuracy did not appear to be adversely impacted in the wake of the erup-  
 382 tion. Although, an experiment simulating ‘cold starts’ to the PPP algorithm showed that  
 383 convergence times were affected by the variable ionosphere in the wake of the eruption.

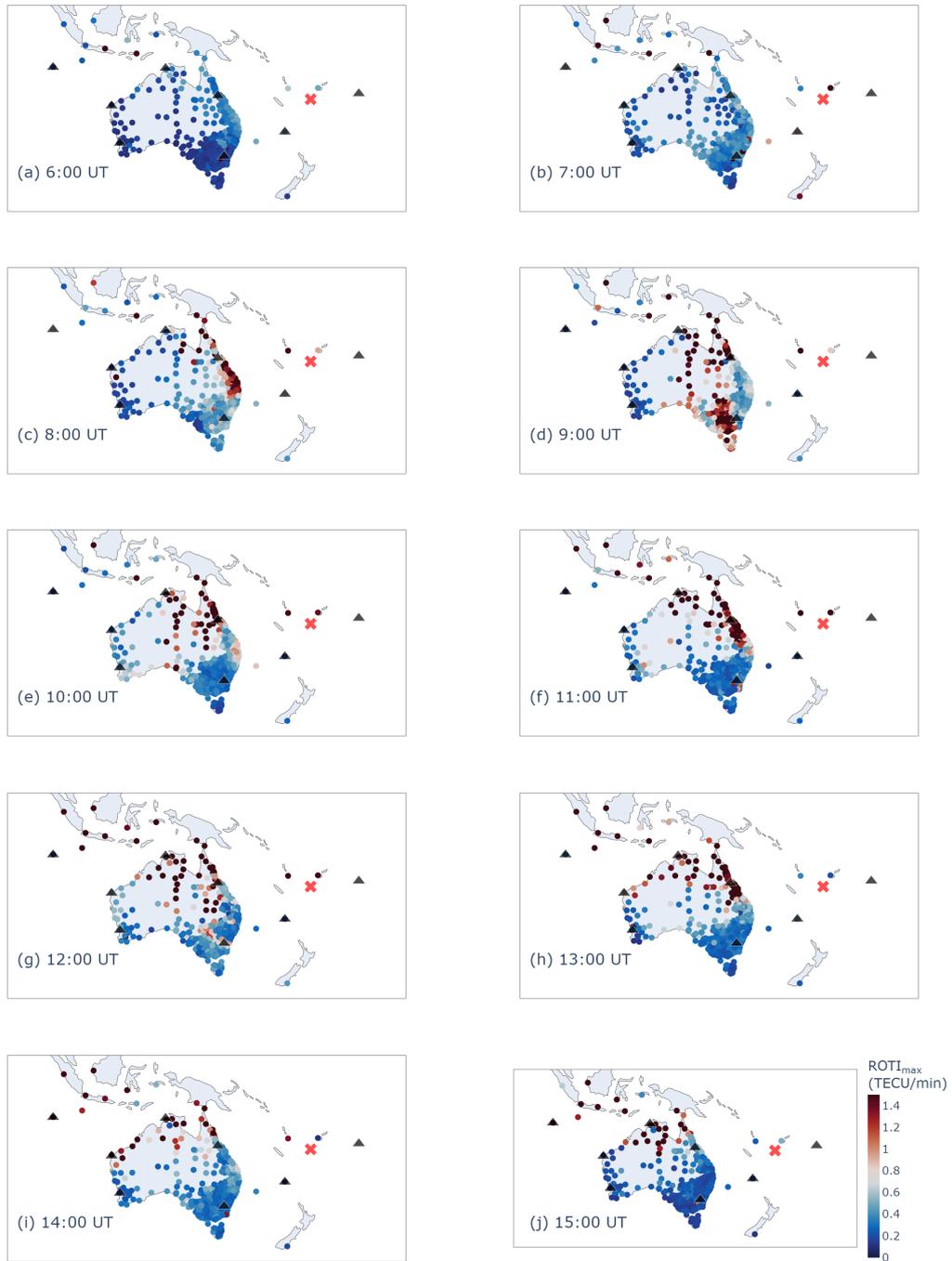


Figure 9: Same as Fig. 7, but showing the observed ROTI maximum in the 30-min following each time; i.e., panel a displays the maximum ROTI observed by each station between 6 and 6:30 UT. The full interactive map is included in the Supplementary Material S2.

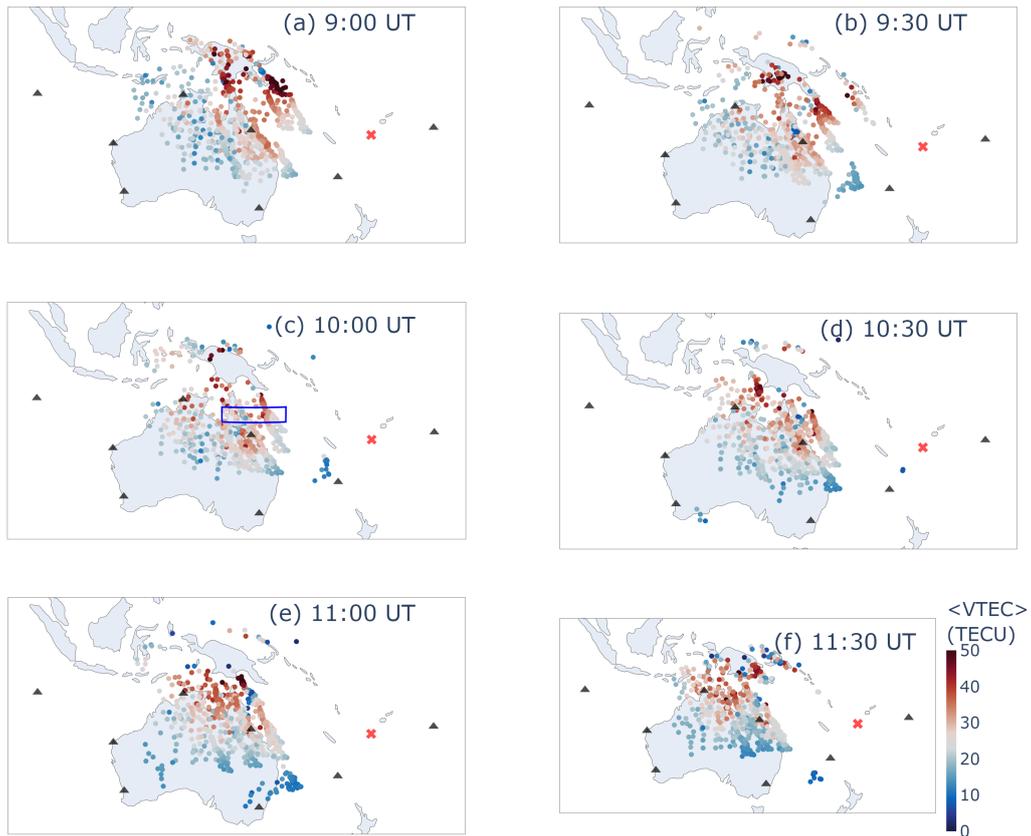


Figure 10: Maps showing the 15-min averaged VTEC (i.e.,  $\langle VTEC \rangle$ ) from 9 UT to 11:30 UT on January 15, 2022 for stations located between latitudes 24°S and 10°S and longitudes 128°E and 152°E (geographic). The blue box in panel (c) shows the range of latitudes and longitudes considered in the analysis to follow. The full interactive map is included in the Supplementary Material S3.

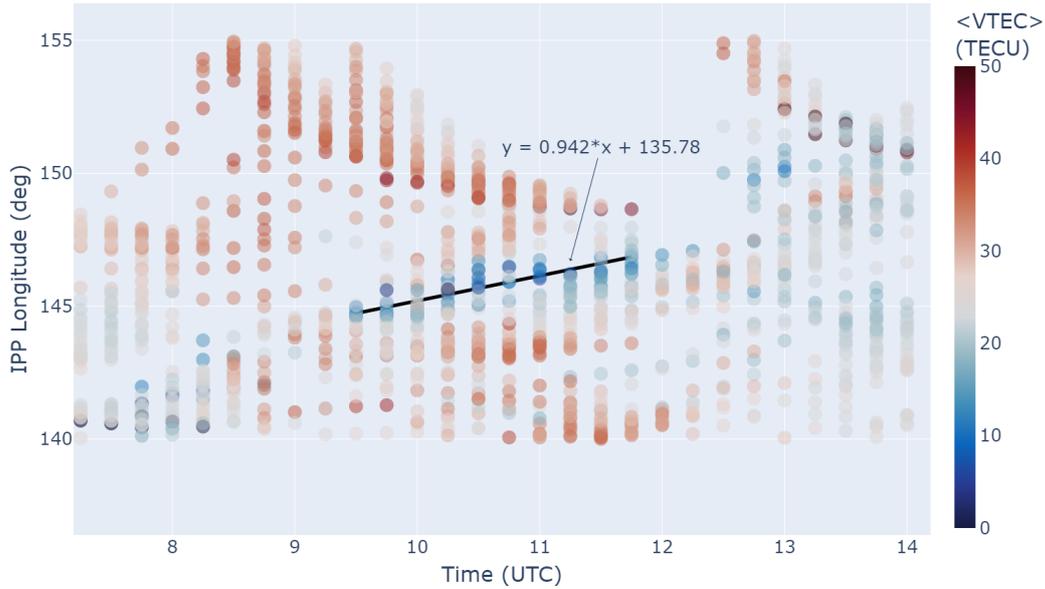


Figure 11: Ionospheric pierce point longitude versus time for data points for locations within the blue box in Fig. 10c – latitudes 17°S–14°S and longitudes 140°E–155°E – colored according to  $\langle VTEC \rangle$ .

384 Using the PPP convergence times from 13 and 14 January as a baseline, it was found  
 385 that stations across northern Australia and Southeast Asia experienced significant  
 386 convergence time increases on January 15, with some stations experiencing increases of more  
 387 than 5 hours. It was revealed that the stations experiencing the largest convergence time  
 388 increases measured extended periods of enhanced ROTI values. Examination of the 15-  
 389 min averaged VTEC data then revealed the presence of a significant depletion that  
 390 extended from Papua New Guinea into central eastern Queensland. A close examination  
 391 of this depletion revealed that it was propagating eastwards with a velocity of  $\sim 30$  m/s.

392 PPP accuracy is at the core of PPP as a usable application. As such, previous works  
 393 that have studied the impact of adverse ionospheric conditions on positioning focus on  
 394 the solution accuracy, and this has typically been done using ‘kinematic’ mode. For in-  
 395 stance, Poniatowski and Nykiel (2020) reported the degradation of PPP accuracy across  
 396 central Europe during the 2015 St. Patrick’s Day storm, with root mean square errors  
 397 of 0.58m, 0.37m and 0.26m in the vertical, north and east directions, respectively. Z. Yang  
 398 et al. (2020) also investigated the impact of the 2015 St. Patrick’s Day storm on PPP  
 399 accuracy, but their analysis included stations located around the world. They found that  
 400 intense auroral activity played a significant role in degrading the PPP solutions, with  
 401  $\sim 70\%$  of high-latitude stations suffering position errors of greater than 1 m. Importantly,  
 402 the most severe degradations in position accuracy were found to coincide with intense  
 403 ionospheric irregularities, as determined by the GPS ROTI parameter. Z. Yang et al.  
 404 (2020) also attributed some observed position accuracy degradation in low latitudes to  
 405 large-scale TIDs –  $\sim 1$  TECU in amplitude – that were generated by the storm. In the  
 406 present study, the PPP errors during the passage of the TIDs from the volcano eruption  
 407 – and the subsequent Spread-F activity – appeared to be smaller compared to that re-  
 408 ported by Z. Yang et al. (2020), despite the larger amplitude of the initial TIDs for this  
 409 event,  $\sim 3$  TECU (S.-R. Zhang et al., 2022). A similar analysis of the 2017 September  
 410 geomagnetic storms also revealed a strong relationship between enhanced ROTI and de-

411 graded PPP accuracy (Zakharenkova & Cherniak, 2021). In their analysis, it was found  
 412 that the 3-D errors rose to several meters due to the presence of small-scale ionospheric  
 413 irregularities. The lower errors in the present study can be understood by the fact that  
 414 ‘static mode’ was used in the positioning algorithm; i.e., the range of potential receiver  
 415 velocities passed to the filter via the ‘process noise’ parameter was set to zero. Interest-  
 416 ingly, the positioning errors in this study remained lower than 10 cm, provided that the  
 417 solution had already converged. One exception is shown in Fig. 5c, where the arrival of  
 418 TIDs at NORF delayed the convergence by  $\sim 2$  hours. While Figure 5 only shows a few  
 419 examples for two stations, their results largely reflect the results of the other stations con-  
 420 sidered in this analysis. Given the rather minimal impact of the volcano eruption on static-  
 421 mode PPP accuracy across the Australian region, the focus of this study shifted to the  
 422 potential impact on the PPP convergence time, which is a key limiting factor in the widespread  
 423 use of PPP (Choy et al., 2017).

424 Here, the investigation into the PPP convergence time made use of an experiment  
 425 in which ‘cold starts’ to the positioning algorithm were performed every 30-min during  
 426 January 15, 2022. Perhaps the most significant finding in this study is that the change  
 427 in convergence time – i.e., the difference between the 15th and 14th – reached 5 hours  
 428 or higher for some stations. Interestingly, these convergence time increases were consis-  
 429 tently observed over northern Australia, and were largely absent across central and south-  
 430 ern Australia, c.f. Figs. 6 and 7.

431 At first glance, the ionosonde data showed several inconsistencies with the PPP con-  
 432 vergence time results. All ionosonde stations observed clear TIDs associated with the  
 433 volcano eruption, followed by some period of Spread-F activity in the ionograms. As re-  
 434 ported by Pradipta et al. (2023), the propagation of the disturbances in these ionosonde  
 435 data agrees well with other studies that have investigated the TID propagation using GPS  
 436 TEC data (e.g., Themens et al., 2022; S.-R. Zhang et al., 2022). As observed by most  
 437 ionosonde stations, the TIDs immediately gave way to Spread-F activity, but in some  
 438 cases there was a significant delay between the arrival of the TIDs and the onset of Spread-  
 439 F. For example, the Townsville station TVL5F, Fig. 2d, showed Spread-F activity com-  
 440 mencing some 3 hrs after the initial TIDs passed, and the Learmonth station showed Spread-  
 441 F commencing more than 6 hours after the TID passage. The intensity of the Spread-  
 442 F also widely varied across the stations. Stations like Canberra, Townsville and Cocos  
 443 Keeling Islands observed quite strong and prolonged Spread-F, whereas one of the sta-  
 444 tions closest to the eruption, Norfolk Island, and another among the farthest away, Perth,  
 445 observed relatively weak Spread-F activity.

446 To better understand the differences in the detection of Spread-F across the Aus-  
 447 tralian region, it helps to be reminded about what the presence of TID signatures and  
 448 Spread-F in the ionograms indicates about the ionospheric plasma. The TID signatures  
 449 highlighted in Fig. 2 that are characterized by descending traces are clear signatures of  
 450 large-scale TIDs, on the order of 100s of kms (e.g., Cervera & Harris, 2014; Pederick et  
 451 al., 2017). While Spread-F observed at low latitudes is typically associated with EPBs  
 452 generated by the Generalized Rayleigh-Taylor instability during the post-sunset to lo-  
 453 cal midnight hours (Sultan, 1996; Burke et al., 2004; Kelley et al., 2006; Carter et al.,  
 454 2014, 2020, and references therein), mid-latitude Spread-F is generally considered to be  
 455 due to the Perkins instability (e.g., Perkins, 1973; Kelley & Fukao, 1991) or specular re-  
 456 flections from medium-scale TIDs with scales on the order of 10s of kms (e.g., G. Bow-  
 457 man, 1981; G. G. Bowman & Monro, 1988; G. G. Bowman, 1990). Interestingly, the re-  
 458 sults in Fig. 6 do not appear to show a clear and consistent relationship between the pres-  
 459 ence of large-scale TIDs or medium-scale TIDs/spread F and increased PPP convergence  
 460 times.

461 The GPS stations used in this study are not high-rate receivers that enable a de-  
 462 tailed analysis into the full spectrum of ionospheric waves (e.g., Cervera & Thomas, 2006;  
 463 van de Kamp & Cannon, 2009; Carrano et al., 2012). However, given the apparent mo-

464 tion of the GPS satellites above ground-based receivers and assuming typical ionospheric  
 465 wave phase speeds, the GPS ROTI parameter can provide insight into the presence of  
 466 ionospheric irregularities on the order of a few kms (e.g., Pi et al., 1997; Ma & Maruyama,  
 467 2006; Nishioka et al., 2008; Zou & Wang, 2009). Generally speaking, there is quite good  
 468 agreement between the impact on the convergence time and enhanced ROTI, by com-  
 469 paring Figs. 6 and 8 and comparing Figs. 7 and 9. Although, one noted difference was  
 470 the presence of enhanced ROTI that propagated over southeast Australia and then dis-  
 471 appeared. The propagation of this ROTI disturbance agrees well with the propagation  
 472 of the Lamb wave across the region reported by Aa et al. (2022) (see their Fig. 4). From  
 473 these observations, one can conclude that the large-scale TIDs caused by the eruption  
 474 had secondary small-scale irregularity generation all the way down to the scale of kms,  
 475 making them detectable using the ROTI parameter. Although, an open question that  
 476 arises is why did these secondary waves not continue to propagate across southern Aus-  
 477 tralia? In other words, what caused the small-scale irregularities over southeast Australia  
 478 at 9 UT in Fig. 9d to either no longer be generated – or be more heavily damped – by  
 479 10 UT in Fig. 9e, when the presence of the Lamb wave in the ROTI data was still clear  
 480 over northern Australia? The answer either lies in the differences in the effectiveness of  
 481 the energy cascade from large-scale to small-scale waves, or in the factors that control  
 482 the damping of small-scale structures between northern and southern Australia. In ei-  
 483 ther case, this is one interesting observation that requires further study and explanation.

484 In terms of the impact of ionospheric irregularities on the PPP convergence time,  
 485 these results suggest that increased ROTI – and therefore the presence of small-scale (i.e.,  
 486 km-scale) ionospheric irregularities – tends to prolong PPP convergence times only in  
 487 the event that the period of increased ROTI lasts for 30 mins or longer. This likely ex-  
 488 plains why stations in Australia’s northern region are clearly impacted compared to the  
 489 stations in southern and central Australia in Fig. 7. Therefore, an ongoing challenge for  
 490 the GPS PPP community is to effectively mitigate the adverse impacts of extended pe-  
 491 riods of small-scale ionospheric irregularity activity.

492 The next question is by what physical mechanism are these small-scale ionospheric  
 493 irregularities being generated in order to have such a detrimental impact on GPS PPP?  
 494 At first glance, the most heavily influenced region in northeastern Australia are too far  
 495 south for typical EPBs to be an obvious likely candidate. At 400 km above Townsville  
 496 (i.e., the IPP), the magnetic latitude is  $31.5^\circ\text{S}$ , and the magnetic field line from this lo-  
 497 cation maps to the equator at an altitude close to 2100 km; using a combination of the  
 498 International Geomagnetic Reference Field (Maus et al., 2005; IAGA, 2010) and the Al-  
 499 titude Adjustment Corrected Geomagnetic (AACGM) (Baker & Wing, 1989; Laundal  
 500 & Richmond, 2017) models. While some observations have been reported of such EPBs  
 501 in the past (Ma & Maruyama, 2006; Cherniak & Zakharenkova, 2016), disturbances this  
 502 large are particularly uncommon. Further, EPBs in the Southeast Asian/Australian lon-  
 503 gitude sector are not commonly observed during the months surrounding the December  
 504 solstice (e.g., Burke et al., 2004; Nishioka et al., 2008; E. Dao et al., 2011; Carter et al.,  
 505 2014, 2020). However, Aa et al. (2022) and Rajesh et al. (2022) recently reported ob-  
 506 servations of EPBs in the vicinity of Australia in the hours following the eruption. The  
 507 detection of strong spread-F traces by the Townsville ionosonde coincides with local times  
 508 that are typical of postsunset EPBs; i.e., 21-01 LT (see Fig. 2). Further, plots of VTEC  
 509 in Fig. 10 in this study agree well with the observations reported by Rajesh et al. (2022)  
 510 in revealing the presence of a deep depletion over northeastern Australia. Here, we mea-  
 511 sure that depletion to be 48 TECU on the ridges and 6 TECU in the valley. The prop-  
 512 agation of this depletion was determined in this study to be eastwards (Fig. 11), on the  
 513 order of 30 m/s. If this TEC depletion was a signature of TIDs propagating away from  
 514 the eruption, then the propagation would have been westwards. Moreover, if it were TIDs  
 515 from Tonga then the propagation would have been expected to much larger than 30 m/s.  
 516 Such low eastward propagation speeds are typical of the F-region dynamo in the dusk  
 517 sector, within which EPBs drift following their non-linear growth to the topside F re-

gion (e.g., Kelley, 2009; Chapagain et al., 2013, and references therein). Therefore, the local time of the Spread-F observed over Townsville, the VTEC depletion depth of 42 TECU and the associated eastward propagation speed in results presented here strongly indicate that the structure responsible for hours of increased ROTI, and subsequently increased PPP convergence times, in northeastern Australia is an EPB, in agreement with Rajesh et al. (2022).

It can be seen from the full animation of Fig. 11 (Supplementary Material (S3)) that data points associated with the depletion over northeastern Australia appeared as far south as 18°S, equivalent to 30°S magnetic latitude with an apex altitude of 1900 km over the magnetic equator. The authors are unaware of any other previously reported EPB disturbance that has been observed this far south over the Australian continent before, although similar-sized EPBs (referred to as ‘super bubbles’ due to their high altitude) have been reported using dense GPS stations across Japan (Ma & Maruyama, 2006). The EPBs that featured in that study would have also appeared over Australia due to magnetic conjugacy, but the relative sparsity of GPS stations across northern Australia did not allow clear EPB observations. However, the density of GPS stations across northern Australia has improved significantly during the solar cycle since Ma and Maruyama (2006)’s study. It is also worth highlighting that elements of Figs. 10 and 11 indicate the presence of TEC depletions farther to the south, but the sparsity of data at these locations makes a clear separation of TID disturbances from EPB-like disturbances difficult for this event. Further, the low density of GPS stations (and therefore, IPPs) across the rest of northern Australia makes a clear determination of whether depletions/EPBs were present difficult in this analysis. However, the impact on PPP convergence times across the rest of northern Australia and southeast Asia shown in Fig. 7 strongly suggests the ongoing presence of EPBs throughout the region during this event; a result in good agreement with the observations of Sun, Wenjie et al. (2022) and Aa et al. (2022) (in particular, the ICON plasma density depletions shown to the north of Australia in their Fig. 6). Some key insights into how TIDs from the eruption caused unseasonal EPB activity in the western Pacific sector can be gained from recent work by Huba et al. (2022). In Huba et al. (2022)’s modeling analysis, waves launched by the eruption caused significant perturbations in the zonal neutral wind in the equatorial plane, which effectively modified the equatorial upward plasma drift and gave rise to a huge EPB that spanned 30° in longitude (between 140-170°E). Huba et al. (2022) referred to this EPB as a ‘super EPB’ due to its large longitudinal extent. In addition, two very-high-altitude EPBs that reached 4000 km at ~ 155°E and 180° were also generated by their simulation (altitudes well above those that Ma and Maruyama (2006) referred to as ‘super bubbles’ in their study). While the 30°-wide EPB in Huba et al. (2022)’s simulation was not clearly observed in the present study (possibly due to its limited latitudinal extent), the observation of a longitudinally narrow, high-altitude and low-density structure over northeastern Australia is consistent with the very-high-altitude EPBs in Huba et al. (2022)’s study; noting the rather reasonable differences in observed versus modeled location. This structure over northeastern Australia thus fits the definition of a ‘super bubble’ according to Ma and Maruyama (2006) and is in good agreement with the results of Rajesh et al. (2022). The impacts on GPS reported in the present study across the rest of northern Australia indicate that further analysis of this event farther towards the west across southern and southeastern Asia is needed.

## 5 Conclusions

The Hunga Tonga Volcano eruption is a unique and complicated event that has provided an unprecedented opportunity to study how the ionosphere couples to the lower atmosphere. In this study, the impact of the eruption – via the ionospheric disturbances that it generated – on GPS PPP was investigated. While static-mode PPP accuracy itself did not appear to be heavily affected, the time taken for convergence to be achieved

570 was found to be significantly impacted. Across northern Australia the impact was particu-  
 571 larly clear, with some GPS stations near Townsville reporting increases in conver-  
 572 gence time of more than 5 hours. Long convergence times are limiting the widespread  
 573 use of PPP, so the Hunga Tonga Volcano eruption presents a good opportunity to re-  
 574 search the impacts and vulnerabilities of a variable ionosphere on PPP and to learn how  
 575 to mitigate/account for them.

576 In this study, it was found that large- and medium-scale TIDs from the eruption  
 577 were not the cause of convergence time increases, but it was the presence of small-scale  
 578 ionospheric structures on the order of a few kms, as determined using the ROTI param-  
 579 eter. Further, it was revealed that PPP was robust enough to endure some enhanced ROTI,  
 580 but not if the duration is longer than  $\sim 30$  mins. This observation effectively differen-  
 581 tiated the impact observed on PPP in northern and southern Australia; northern Aus-  
 582 tralia/southeast Asia experienced extended periods of increased ROTI and subsequently  
 583 increased convergence times, whereas southern Australia only experienced a short burst  
 584 of increased ROTI with marginal impact on convergence time.

585 The results of this study indicate that the ionospheric phenomenon responsible for  
 586 the presence of small-scale irregularities in southern and northern Australia was differ-  
 587 ent. In southern Australia, the small-scale irregularities were the result of secondary gen-  
 588 eration from the large- and medium-scale TIDs propagating away from the eruption. The  
 589 enhanced ROTI region propagated to the west in a manner consistent with the Lamb  
 590 waves previously reported. These small-scale structures were rather promptly damped  
 591 and did not propagate far beyond southeastern Australia. In northern Australia, the ex-  
 592 tended periods of enhanced ROTI were found to be due to the presence of at least one  
 593 ‘super bubble’ that was observed as far south as  $\sim 30^\circ\text{S}$  magnetic latitude, with an es-  
 594 timated apex altitude of  $\sim 1900$  km above the magnetic equator; this is the same ‘su-  
 595 per EPB’ recently reported by Rajesh et al. (2022). The VTEC data revealed that the  
 596 EPB was  $\sim 42$  TECU deep and  $\sim 300$  km across in longitude. Further, it was shown that  
 597 the EPB traveled eastwards at  $\sim 30$  m/s, consistent with the F-region dynamo speed  
 598 and direction. This super bubble is the largest/most southward-reaching EPB observed  
 599 over the Australian continent; only recently made possible due to GPS station deploy-  
 600 ments in the region.

601 The Hunga Tonga Volcano eruption stands as an excellent example of how iono-  
 602 spheric variability can adversely influence satellite-based precise positioning that is in-  
 603 creasingly heavily relied upon across many industries and sectors around the world.

## 604 **6 Open Research**

605 The Bureau of Meteorology Space Weather Services’ World Data Centre provides  
 606 ionosonde data via their website; [https://www.sws.bom.gov.au/World\\_Data\\_Centre](https://www.sws.bom.gov.au/World_Data_Centre). Geo-  
 607 science Australia provides GNSS data for all of the stations used in this analysis via their  
 608 Global Navigation Satellite System Data Centre; <https://gnss.ga.gov.au/>. Geoscience  
 609 Australia’s Ginan platform is also accessible via their github repository;  
 610 <https://geoscienceaustralia.github.io/ginan/>. The higher level analysis products, par-  
 611 ticularly the PPP convergence times that were calculated using Ginan, are available on  
 612 the Zenodo data repository (doi:10.5281/zenodo.7694409).

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# Supporting Information for “The ionospheric effects of the 2022 Hunga Tonga Volcano eruption and the associated impacts on GPS Precise Point Positioning across the Australian region”

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## Contents of this file

1. Description of interactive maps S1, S2 and S3

## Description of interactive maps S1, S2 and S3

Here we provide some additional interactive maps to supplement the figures presented in the main manuscript. The format is html, which is supported by internet browser

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programs. While Figures 7, 9 and 10 in the main manuscript provide snapshots of the map of Australia, with the change in convergence time  $\Delta T_C$ , ROTI and  $\langle VTEC \rangle$  data, S1, S2 and S3 here provide the reader with interactive maps to facilitate their own data exploration. Each map animates through time and is controllable with the slider at the bottom; S1 spans 0-18 UT, S2 spans 0-23:50 UT and S3 spans 07-16:45 UT on January 15, 2022. Further, the reader can zoom in and out of the maps plotted and click-and-drag to move around the map, and explore the plotted data points by hovering the mouse pointer over them. For reference, the location of Tonga is shown as the red cross, and the locations of the ionosonde stations are shown in S2 and S3 as black triangles.