# Cosmology constant and quantum mechanics equation based on the rotational gravitational field

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#### Abstract

In this work, the gravitational field is investigated in detailed and the quantum mechanics equation under the gravitational field has been derived. Then, the Schrodinger and Dirac equations are accordantly solved under the gravitational field condition by separating variables. As a result, the Rydberg formula is deduced in such conditions, which proves that the change of the external gravitational field intensity will cause the overall spectral movement. Obviously, the partial redshift of quasar spectrum should assign to this effect. Furthermore, by applying this gravitational field together with the energy and mass concepts into the symmetry, gravity theory and gauge theory, it is deduced that the interaction of "gravity" between matter and anti-matter is repulsive force, which is the originator of the accelerated expansion phenomenon for dark energy in the universe. It is found that the calculated cosmological constant is a small variable related to the radial and angular direction of the universe, and the "spontaneous breaking of vacuum symmetry" is caused by this gravitational field. Further, the gravitational field lead to the non-conservation of weak action parity. The equal number of baryon and antibaryon as well as the energy conservation in the universe are confirmed. In this work, the gravitational field is introduced into quantum theory, which will promote the integrality of the quantum mechanics, and explain the dark energy phenomenon constitutionally. This study will push the astrophysical theory and the gauge theory of particle physics for the further study of energy level, basic particle structure, and quantum gravity theory.

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8	Key Points:
9 10	• The change of the external gravitational field intensity will cause the overall spectral movement
11	• The interaction of "gravity" between matter and anti-matter is repulsive force
12 13 14	• The equal number of baryon and antibaryon as well as the energy conservation in the universe are confirmed

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#### 34 **1 Introduction**

35 Dark energy phenomena [1-4], dark matter [5,6], neutrino mass, asymmetry of material and antimatter are the known experimental or observable facts, which cannot be explained or does not 36 be included by either the standard particle physics model or the standard cosmic model. As we 37 know, the electromagnetic interaction, the strong interaction, and the weak interaction have been 38 unified as gauge theory interactions. However, it is the hot issue in the past 50 years or even nearly 39 100 years to explore a simple universal physical principle to unify all kinds of interaction forces. 40 In order to explain the new phenomena such as those originated from the experiments and 41 observations and also those cannot be explained by the above standard models, and as well in order 42 to intergrade gravity theory with quantum theory, it is necessary to quantify the gravitational field. 43 By such an approach, a quantum gravity theory is constructed which could response the 44 completeness of quantum theory. Actually, it has encountered many insurmountable difficulties: 45 1) The quantum theory is different from the general relativity in the concept of time. For instance, 46 the time corresponded to the instantaneous collapse of the quantum state is absolute whereas it is 47 relative for the time of general relativity. 2) It is much more complicated for the quantal 48 gravitational field than the quantization of electromagnetic field. The metric is a second order 49 tensor, that is, it includes the gravitational information as well as the time and space geometry. 50 51 This makes the change of gravitational field in space and the evolution of gravitational field with time being ambiguous. Further, it is ambiguous for the motion and evolution of gravitational field 52 after the quantization. The perturbation method of quantization gravity has also been tried, but the 53 problem of gravity non-renormalization is not perfectly solved. Later the attempts were tried to 54 find solutions to gravitational field quantization from different views, such as the non-perturbed 55 method is applied to the circle quantum gravitation (Loop Quantum Gravity) [7,8] of the 56 gravitational quantization, the string theory, and superstring theory [9-16]. Horava-Lifshity theory 57 have been proposed by modifying Einstein general relativity. All these theories formed by either 58 replacing the metric tensor with nonlocal field operators, or considering particles as spatiotemporal 59

nonlocal representations of strings and interactions in high dimensional supersymmetric space-60 time, or correction theory of Lorentz symmetry breaking at high energy. Therefore, although these 61 theories could find solutions for several problems, most of them are locally questionable or 62 imitative. For examples, the circle quantum gravitation theory has dynamic problems while string 63 theory (superstring theory) has encountered serious difficulties in compacting to the real 4-64 dimensional space-time, and Horava-Lifshit [17] theory could not response the Lorentz symmetry 65 at low energy. Recently, several quantum effective theories were proposed, such as, the double 66 special relativity (DSR) model [18,19] which was constructed by modifying the energy momentum 67 relationship of the relativistic quantum, the general uncertainty principle (GUP) [20-23] which was 68 derived from the uncertain relationship in quantum mechanics, and the Lorentz symmetry breaking 69 70 (SME)[24,25] which was based on the standard model extending. These theories have gained some valuable conclusions in the investigation of the high energy phase quantum gravity effect in Planck 71 scale or black hole thermodynamics. Whereas, further research and exploration are significantly 72 needed due to the limitative (or defective) parameters described the quantum gravity effect. There 73 are also investigations on the scattering effects [26-28] and the Higgs push [29,30] of the cosmic 74 inflation originated from the non-minimal coupling term RH<sup>+</sup>H (H is the Higgs field and R is the 75 standard curvature of space time scalar) between Higgs field and the gravity in the framework of 76 effective field theory. The spin value of the graviton and high-energy phases have been assumed 77

to be 2 in this investigation. A colored Higgs field  $H^{c}$  was introduced to the SU(5) grand unity theory (GUT)[31-35] whereas the proton decay is synchronously caused. Consequently, the progress has been made in high energy and black hole processing by fitting the gravity into the standard model of the gauge theory. Nevertheless, the gravitational fields used are all the Einstein gravitational field with the plane wave solution, in which spin value is 2. It is not the ideal candidates for the interaction theory between Higgs and gravitational forth. Here, a complete theory of quantum gravity is still not established until now.

The Higgs mechanism of the gauge theory and its particle physics standard model is to 85 assume that there naturally is a scalar field  $\Phi(\chi)$ . In the vacuum state of this scalar field, the 86 interaction between the scalar field and the gauge field or the fermion field makes the gauge 87 particles and fermions obtain excess mass. However, this is only a scalar field and the 88 corresponding spontaneous breaking vacuum state assumed at the requirements of gauge theory. 89 In one hand, it is not clear whether such a real scalar field and the required spontaneous breaking 90 vacuum state exist. On the other hand, our understanding of the physical real vacuum is relatively 91 vague. If the gravitational field is quantized and the physical vacuum is described on the basis of 92 such quantum gravitational field, the physical vacuum contained the quantum gravitational field 93 and its vacuum state will be more suitable for the gauge theory. As a result, not only the real 94 physical field of the scalar field required by the gauge theory is determined, but also the physical 95 reality of the real vacuum is correspondingly determined. Therefore, the theoretical vacuum and 96 the real vacuum is equated, and the gravitational field could be introduced into the gauge theory, 97 which will promote to specify the standard model of the particle physics and the gauge theory. 98

In the section 2 of this manuscript, we have defined the chiral energy, the chiral mass and the chiral gravitational field (the only one assumption in this manuscript), and then the CPT theorem is proved to be held in this gravitational field. After that, in the section 3, the 1/2 spinor gravity field is naturally introduced under the concept of chiral gravity field. After fully understanding the different expressions of microgravity field and macrogravity field and the original establishment of quantum mechanics, the quantum mechanics equation under the

condition of gravitational field is established through Poisson's equation. The energy solution of 105 this equation is the Rydberg formula under the condition of gravitational field. Moreover, the 106 overall shift effect of the atomic characteristic spectrum can be derived. It is proposed the dense 107 matter will appear if the external gravitational field tends to infinity but the mass matter is unstable 108 and tends to decompose when the gravitational field become infinitesimal. In section 4, the 109 repulsive gravitational interaction between positive and negative mass matter is established. Next, 110 the scalar field of the gauge theory is namely identified as the scalar of the spinor gravitational 111 field via the mass generation of the Higgs mechanism. Hence, there existed the transverse energy 112 and the longitudinal gravitational field conditions for the generation of the mass. Because the 113 different contributions from both the positive-anti particle spin in the term coupled the scalar field 114 in Lagrange quantity with the fermion field and the different positive and negative masses, the 115 positive and antiparticles will present the various lifetimes. It makes a positive matter stable and 116 anti-matter unstable in the right-handed gravitational field. And the stronger the field strength, the 117 greater the difference of the positive particle and anti-particle in the Lagrange quantity value. That 118 is, the right-handed gravitational field is a positive matter stable field and the unstable field of the 119 antimatter. Combining the CPT theorem with the limit case of the Rydberg formula in the 120 gravitational field condition, it derives a conclusion that a microscopic chiral symmetric particle 121 world is existed and also a macroscopic chiral-symmetric universe. The universe consists of the 122 positive and antimatter sky, in which the positive and antimatter cannot mix together to form the 123 124 universe. Based on this derivation, the cosmological constant is calculated under the positive matter sky and antimatter sky model, which is a quasi-constant associated with the cosmological 125 radial and angular direction. 126

127 The cosmic model presented in this manuscript is obtained naturally based on the theory 128 of particle physics and is an extension of the Standard Universe Model ( $\Lambda$  CDM), which is 129 different from the Dirac-Milne universe [36]. The  $\Lambda$  solved in this manuscript is also simple in 130 principle, and it will be better if it derived from the observations and n-body simulations [37-40] 131 (our observations are limited to the positive sky, but the long-range effect of the anti-sky is still 132 valid). However, the redshift caused by the change of matter structure in the gravitational field will 133 have a big impact on cosmic observations and cosmic theory.

Nowadays, it is still inharmonic for the physics gauge theory and gravity field theory. Therefore, it is good choice for us to distinguish the representation and function of the microscopic gravitational field from the macroscopic classical gravitational field. Moreover, the microscopic gravitational field, the quantum mechanics, and the quantum field theory will be combined. It is truth that this combination of a series of interrelated conclusion can be identified by experiment, such as the spectral redshift and the existence of the dense matter (The white dwarf and the neutron star) have actually proved this theory.

## 2 Definition of chiral energy, mass and gravitational field, and the field quantization and the certification for CPT's Theorem

143 2.1 Definition of right-hand energy and left-hand energy

144 The propagation behavior of electromagnetic wave in medium is determined by the 145 permittivity  $\varepsilon$  and permeability  $\mu$  of medium. The relationship between wave vector K and 146 electromagnetic vector E, and H can be deduced by Maxwell's equations:

147  $\mathbf{K} \times \mathbf{E} = \mathbf{w} \mu_0 \mu_r H \tag{2.1}$ 

148  $\mathbf{K} \times \mathbf{H} = -\mathbf{w}\varepsilon_0\varepsilon_r E$  (2.2)

149 Where, w is the angular frequency,  $\mu_r$  is the dielectric permeability,  $\varepsilon_{\mu}$  is the dielectric 150 permittivity. By equations (2.1) and (2.2), the following equations are obtained:

151  $K^2 = \left(\frac{\omega}{c}\right)^2 n^2$  (2.3) 152 Where the refractive index of medium  $n^2 = \varepsilon_r \mu_r$ . Then,

153 
$$K = \pm \frac{w}{a}n$$

i) When  $\varepsilon > 0$  and  $\mu > 0$ , n is large than 0 and  $K_R = \frac{\omega}{c}n$ . In addition, E, H, K satisfy the right-hand relationship. The energy is thus defined as the right-hand energy  $\epsilon_R$  or the nominal positive energy  $\epsilon_R = \epsilon$ .

(2.4)

157 ii) When  $\varepsilon < 0$  and  $\mu < 0$ , n is large than 0 and  $K_L = -\frac{\omega}{c}$ n (negative value). In addition, E, H, 158 K satisfy the left-hand relationship. Meanwhile, the direction of light propagation (e.g., direction 159 of the K and it is also the phase velocity direction) is opposite to the energy propagation direction 160 (S direction of the Poynting vector). The energy is thus defined as the left-hand energy  $\epsilon_L$  or the

161 negative energy:  $\epsilon_L = -\epsilon$ ,  $K_R = -K_L$ .

162 2.2 Definition of chiral mass and chiral gravity field

163 Under the location-time four vectors  $x^{\mu}$  ( $\mu=0,1,2,3$ ) :  $x^{0} = ct$ , x' = x,  $x^{2} = y$ ,  $x^{3} = 164$  z, the energy-momentum four-vectors is  $P^{\mu} = (\frac{E}{c}, P_{x}, P_{y}, P_{z})$ .

According to the above definition of chiral energy and 4-momentum, the chiral mass and the 165 chiral gravitational field are defined as the right-handed energy  $E_R$  (in order to distinguish the 166 electric field vector energy E, the energy in section (2.1) is written as  $\epsilon$ , and here E is also used to 167 present the energy) and left-hand energy E<sub>L</sub>. Then,  $P_R^0 = \frac{E_R}{c} = \frac{E}{c}$ ,  $P_L^0 = \frac{E_L}{c} = \frac{-E}{c}$ . Accordingly, the 168 right-hand mass (intrinsic property) is defined as  $m_R=m$ , that is, the mass of normal matter while 169 the left-handed mass is defined as mL=-m, a antimatter mass. Furthermore, the corresponding 170 gravitational potentials (take the single particle Newton potential as an example to illustrate the 171 two kinds of the gravitational potentials) are defined as: 172

173  $\varphi_R = -\frac{Km_R}{v} = -\frac{Km}{v}$ 

174 
$$\varphi_L = -\frac{V_L}{\gamma} = \frac{Km_L}{\gamma}$$

The gravitational potential carried by the right-handed mass is right-handed whereas the lefthanded mass is the left-handed gravitational potential.

177 2.3 CPT theorem under chiral energy and mass [41]

The CPT invariant representation of various spin fields is written by quantizing each quantity field under the definition of left-handed energy, right-handed energy, and chiral 4-momentum. Then, the CPT theorem  $\oint \mathcal{L}(x) \oint^{-1} = \mathcal{L}^+(-x)$  ( $\oint = CPT$ ) has been derived, and the detail is listed in Appendix 1. Thus, the following conclusions could be followed with the CPT theorem:

- 182 1) The existence of right-handed energy  $E_R$  and the chiral symmetrical left-handed energy  $E_L$ . 183 It is set  $E_R$ =E during the calculation which is known as the positive energy while  $E_L$ =-E known as 184 the negative energy.
- 185 2) There is right-handed mass matter (also called positive matter  $m_R=m$ ) and left-handed mass 186 matter (also called antimatter  $m_L=-m$ ). This mass chirality is the intrinsic property of matter and 187 does not change due to the selection of reference system.

3) The mass matter carries the gravitational field and the mass matter exhibits chiral
 symmetry. Then, the gravitational field accordingly presents chiral symmetry. That is, the positive
 matter carries a right-handed gravitational field and antimatter carries a left-handed gravitational
 field.

4) There is a microscopic world of chiral symmetric particles, and also there is a macroscopic 192 chiral symmetry universe. Such universe consists of the positive matter sky and the antimatter sky 193 (the matter and antimatter in the universe form their own positive and antimatter sky, respectively. 194 But it cannot be a mixture of matter and anti-matter to form the universe which will be discussed 195 in the Appendix 4 particle physics parts). In a physical vacuum in the sky of positive matter, it 196 is  $\varepsilon_0 > 0$ ,  $\mu_0 > 0$ . While, in a vacuum for an antimatter sky, it is  $\varepsilon_0 < 0$ ,  $\mu_0 < 0$ . It means there 197 are two gravitational field vacuums in the universe. The defined "positive sky" is actually what 198 people now call the universe. In fact, the cosmic model in this manuscript doubles the original 199 universe (a pair of positive and anti-sky). 200

201

# 3 Establishment of the quantum mechanics equation under gravitational field and its physical meaning

204 3.1 Establishment of the quantum mechanics equation under gravity field

Today, the superstring theory (or not yet the so-called M theory) is the hottest research by 205 quantizing the gravity field or combining the gauge theory with the gravity theory. It is a more 206 complicated mathematical process by introducing boundary conditions under the light cone 207 specification, or repeating the classical string theory, gravity theory and quantum theory in string 208 length  $(10^{-30} \text{ cm})$  regions. Whether or not is it correct (because we still cannot verify the 209 authenticity of the string theory under Planck energy), it is truth that it can only study gravity in 210 Planck length ( $10^{-33}$  cm) and explore the universe in Planck time ( $10^{-44}$  sec). For relatively low-211 energy quantum mechanics or quantum field theory, we are hardly able to consider the effects of 212 gravity. 213

The expression and action of gravitational field in microscopic gauge theory and macroscopic gravity theory should be different, which is also a manifestation of the inconsistency of microscopic and macroscopic theories. The Eeinstain equation of the macroscopic gravitational field is suitable for the large-scale physics. For microscopic gravitational fields, we need to seek out from Schrodinger and the initial quantum mechanics of Dirac.

According to the definition of section 2, the natural introduction of spin 1/2 of the microscopic gravitational field, leads to its quantization, microscopic gravitational field quantum state, microscopic equation of motion, Lagrangian Hamiltonian are all 1/2 field.

The gravity is very weak than the other three interactions, so the quantum mechanics that describes the interaction of microscopic particles does not consider the gravitational field at all. However, the physical behavior is all occurred in the gravitational field, which forces us to consider the influence of the gravitational field. Thus, a complete quantum mechanical system has to include gravity in it.

Here, on the basis of fully understanding the concept of primitive quantization, the quantization mechanics equation under gravitational field is obtained by using the classical 3D gravitational field Poisson equation.

First of all, it is clear that the following properties of the gravitational field are included: a) The gravitational field is a form of energy (the other form of energy is electromagnetic wave), a spin field, and a helical field of spin 1/2, with a component of 1/2 or -1/2. b) As a form of energy

existence, the gravitational field can exchange energy momentum directly with mass matter, which 233 does not require the intermediate particles. Unlike the electromagnetic energy transfer, 234 gravitational fields do not need to be exchanged with gravitons in the concept of gravitational 235 waves (radiation). For example, neutrinos produced in particle decay are the direct exchange of 236 energy momentum and angular momentum between mass matter and gravitational field. That is to 237 say, the energy and momentum transfer between the gravitational field and the mass matter is not 238 carried out in the form of radiation, but by the gravitational field itself, which is a function of the 239 gravitational field and different from the electromagnetic field. It is also the most essential 240 difference between the present quantization of gravitational field and the previous quantized 241 gravitational field. c) All the executable physical experiments and observations are going in the 242 gravitational field of the positive sky space, and the vacuum state is the lowest energy state of the 243 gravitational field. 244

The second quantization of spin 1/2 field is listed detailed in section 3 of Appendix 1. Here only the quantization process is simply written. The description of 0 mass spin 1/2 field can be described by a single spin  $\psi$ , which satisfies the Dirac equation:

248

$$(\hat{\mathbf{p}}_{0} + \hat{\mathbf{p}} \cdot \boldsymbol{\sigma}) \boldsymbol{\psi} = 0$$
 (1)

249 where  $\hat{p}_{\mu} = i\partial\mu$  is a four-momentum operator,  $\sigma$  is a Pauli matrix. Between the energy 250 and momentum of the zero mass fields, it is

1  $_{int}$   $_{inx}$ 

equation (2) and (3), the state of a particle with a certain momentum is necessarily nerval. Based on the above concepts of gravitational field properties and quantization, a classical three-dimensional gravitational field Poisson equation ( $\Delta \varphi = 4\pi G p$ ) is used to establish a quantization mechanics equation under the gravitational field. As it is known, the classical Poisson equation is about mass flux. In view of microcosmic quantum mechanics, the unit volume mass substance ( $\rho$ ) converted into energy substance (E=mc<sup>2</sup>) can be divided into two parts: one is in

the form of transverse polarization energy of the electromagnetic radiation, and the other part is 266 the existence of energy in the form of longitudinal polarization gravitational field (the Higgs 267 mechanism in the following sections will show the existence of energy during analyzing the 268 formation of mass). Or in other words, the quantum form of mass matter per unit volume is 269 composed of two parts: one is corresponding to the transverse polarization energy part and the 270 other is the longitudinal polarization energy. The quantum understanding of a classical  $\Phi$ 271 gravitational potential is a scalar quantity of the spinning gravitational field where the quantum 272 form is  $(\bar{\psi}\psi)$ ,  $\varphi \sim \bar{\psi}\psi$ , that is the gravitational potential varies according to the the scalar of the 273 spinning gravitational field. It can be defined as: 274

 $\phi = k \overline{\psi} \psi$ 275

Here  $\kappa$  is the quantity associated with the strength of the macroscopic gravitational field and 276 can be considered as C constant in the equation. 277

(5)

(4)

Mass density per unit volume: 278

$$\rho = \frac{E}{C^2}$$

 $E \to i\hbar \frac{\partial}{\partial t}$ The energy E in quantum mechanics can be replaced by 280

The quantum substitutions of equation (4), (5) and energy E were placed into the classical 281 Poisson equation. Meanwhile, the longitudinal polarization energy part is added into the equation. 282 The longitudinal energy is expressed as  $\gamma \cdot \hat{p}$  with the zero mass Dirac field (space 3D). The 283 corresponding Dirac matrix is introduced and applied to the wave function to obtain the equation: 284

$$(\gamma^0 \Delta k \overline{\psi} \psi + k \gamma \cdot \hat{p}) \psi(r, t, \sigma) = \frac{4\pi G}{C^2} i \hbar \gamma^0 \frac{\partial}{\partial t} \psi(r, t, \sigma)$$

285

$$\mathbf{k}(\gamma^{0}\nabla^{2}\bar{\psi}\psi + \gamma \cdot \hat{p}) \ \psi(r, t, \sigma) = \frac{4\pi G}{C^{2}} i\hbar\gamma^{0} \frac{\partial}{\partial t}\psi(r, t, \sigma)$$

Or represented as: 286

An infinite small volume element is taken, and in which the scalar of the curl field  $\bar{\psi}\psi$  can 287 be regarded as a constant. Moreover, only the wave function  $\psi(\mathbf{r},\mathbf{t},\sigma)$  changes with the quadratic 288 derivative. Compared with the Schrodinger equation of the fundamental wave equation of quantum 289

 $\bar{\psi}\psi = -\frac{\hbar^2}{2m}$ , and finally a mechanics, the "normalized" transverse energy is partially obtained as 290 quantization mechanics equation under the quantization gravitational field is obtained as following: 291

$$\mathbf{k}\left(-\frac{\hbar^{2}}{2m}\gamma^{0}\nabla^{2}+\gamma\cdot\hat{p}\right)\psi(r,t,\sigma)=\frac{4\pi G}{C^{2}}i\hbar\gamma^{0}\frac{\partial}{\partial t}\psi(r,t,\sigma)$$
(6)

292 293

Several conclusions are derived by analyzing the physical meaning of equation (6):

1) when all unit mass matter is converted into the transverse energy without forming 294 longitudinal gravitational field energy, or when the unit mass matter exists in the form of the mass 295 matter particles and the transverse energy, but not the longitudinal energy, there is no rotation field 296 in equation (6). In this case, the equation (6) becomes the Schrodinger of the free particles in the 297

$$(-\frac{k\hbar^2}{2m}\nabla^2)\psi(\mathbf{r},t) = \frac{4\pi G}{C^2}i\hbar\frac{\partial}{\partial t}\psi(\mathbf{r},t)$$

298 context of gravitational field:

Clearer after adjusting: 299

$$\frac{\mathrm{k}\mathrm{C}^2}{4\pi G} \left(-\frac{\hbar^2}{2m}\,\nabla^2\right) \psi(\mathbf{r},t) = i\hbar \,\frac{\partial}{\partial t}\,\psi(\mathbf{r},t) \tag{7}$$

300

$$C_{M} = \frac{kC^{2}}{4\pi G}, \text{ now, the equation (3.7) is transformed into:}$$

301

$$C_{M}(-\frac{\hbar^{2}}{2m}\nabla^{2})\psi(r,t) = i\hbar\frac{\partial}{\partial t}\psi(r,t)$$
(8)

General form of the Schrodinger equation:

304

$$H' = C_{M}H = C_{M}(-\frac{\hbar^{2}}{2m}\nabla^{2})$$

 $i\hbar \frac{\partial}{\partial t} \psi(\mathbf{r}, t) = H' \psi$ 

305

The coefficient  $\kappa$  is a constant variable related to the strength of the macroscopic gravitational field, and the corresponding C<sub>M</sub> is also a quantity related to the strength of the gravitational field. Therefore, it is dominated as the quantum conditional coefficient of the gravitational field. The stronger the gravitational field is, the smaller the  $\kappa$  and C<sub>M</sub> are. Therefore, the different quantum mechanical equations possess the different gravitational field strengths and the earth region is the

(9)

$$C_{M} = \frac{kC^{2}}{4\pi G} = 1$$

311 quantum conditional region of

2) The equation (6) will become the equation of 0 mass field Dirac equation when the unit mass is completely converted to the longitudinal gravitational field energy and there is no transverse electromagnetic field energy.

$$k\gamma \cdot \hat{p}\psi(r,t,\sigma) = \frac{4\pi G}{C^2} i\hbar\gamma^0 \frac{\partial}{\partial t}\psi(r,t,\sigma)$$
(10)

315

Equation (2.10) is the 0-mass spin field Dirac equation under the gravitational field conditions.
 This equation could be deformed as follows:

$$C_{M}\gamma \cdot \hat{p}\psi(r,t,\sigma) = i\hbar\gamma^{0}\frac{\partial}{\partial t}\psi(r,t,\sigma)$$

318 *CT* 319 Or a similar form of Schrodinger equation:

$$i\hbar \frac{\partial}{\partial t} \psi = H'\psi$$

320

321 Here, the Hamiltonian is:

$$H' = C_M H = C_M \gamma \cdot p = C_M (-i\hbar\gamma \cdot \nabla), \hat{p} = -i\hbar\nabla$$

Actually, the Schrodinger equation and the Dirac equation can be obtained by separating the variables from equation (6) under the condition of gravitational field.

325 3) the quantum of the microcosmic gravitational field satisfies the longitudinal energy 326 polarization condition. Here, the neutrino is regarded as the quantum of the gravitational field, and 327 three kinds of neutrinos have been discovered so far.

The gravitational field quantum defined in this work is the Neutrinos. There are three kind of the Neutrinos which are all classified into energy levels of a 0-mass particles. It contradicts with the neutrino mass required in the Neutrino oscillation theory, in which the Neutrinos are not equal and cannot be all zero mass. But it does not contradict with the oscillation phenomenon itself. The

(11)

reason is that the Neutrino mass in the weak electric unified standard model of  $SU(2) \times U(1)$  is 332 strictly equal to zero. Whereas, the right-handed Neutrino component of the SU(2) singlet state 333 and the coupling of the Yukawa are introduced to form the Dirac mass term in the extended 334 standard model. Then, Neutrinos is sure to have the theoretical mass. In addition, in the mixing 335 theory of Neutrino quantum mechanics, the Dirac mass shall be of the same order as the other 336 lepton e,  $\mu$ ,  $\tau$  mass. But the Neutrino mass is very small, and the "seesaw mechanism" was 337 introduced to give a small mass to the Neutrinos. It shows that the Neutrino theory is debatable. 338 Furthermore, the cosmological observations of the Neutrino mass are  $m_e + m_\mu + m_\tau < 0.28 \text{ eV}$  (2010) 339 year). Generally, the mass of Neutrino is not exceeding 1 eV and its controversial mass is indeed 340 very small. As it is reported the mass of Neutrino is zero or very small in several theories. Finally, 341 if we confirm the Neutrino to be a gravitational field quantum, then it is certain to be affected by 342 the gravitational field during the transmission process. Consequently, the specially change appears, 343 that is, Neutrino oscillation, which is the nature of the gravitational field, to be further studied. 344 Based on the above discussions, there is no specific explanation between the 0-mass definition in 345 the gravity field quantum and the neutrino oscillation theory. 346

347 3.2 The solution of the quantization mechanics equation under the gravitation field and its348 physical significance

Energy solutions of the Schrodinger and Dirac equations in the gravitational field background and the energy solutions of hydrogen atoms are described in detail in Appendix 2. As an example, the hydrogen atoms are discussed concretely below.

### 352 The energy level difference of hydrogen atoms is: $\Delta E = C_M \left[ \frac{\mu e^4}{2\hbar^2} \left( \frac{1}{m^2} - \frac{1}{n^2} \right) \right]$ (12) 353 Take the wave number as: $\bar{\nu} = C_M R \left( \frac{1}{m^2} - \frac{1}{n^2} \right)$ (13)

354 Where R = 
$$\frac{2\pi^2 \mu e^4}{h^3 c}$$
 is the constant of Rydberg

In the Earth's quantum condition region, the  $C_M=1$ , and the formula (13) becomes Rydberg formul  $\bar{\nu} = R\left(\frac{1}{m^2} - \frac{1}{n^2}\right)$ .

In terms of the spectral formula (13), when  $C_M$  takes different values, that is, the differently external gravitational field strength, the overall shifts of the hydrogen characteristic spectrum occur. Unlike the energy level splitting of the Stark effect in the external electric field and the Zeeman effect in the external magnetic field, the atomic characteristic spectrum in the external gravitational field has the overall shift effects with the various strength of the gravitational field, such as the quasar spectral redshift. The results here is the greater the field strength changes, the greater the spectral moves.

Using the data and conclusions in the literature [42], the formula (13) can be solved: firstly, defining  $C_M = \frac{1}{Z+1}$ , for instance, for QSO Q1442+231 the red shift of the emission line  $\alpha$  in the Lyman spectrum is Z=3.625. At this point,  $C_M$ =0.2162, then the solution of (13) is  $\bar{\nu}$ =1/5622 Å, that is,  $\lambda$ =5622 Å.

368 On the one hand, it shows that the quantization mechanics equation under the gravitational 369 field is suitable for different gravitational field under the strong backgrounds. The energy solution 370 of the spectral formula is the observed wavelength (the intrinsic wavelength of the hydrogen 371 Lyman  $\alpha$  spectrum or the Earth zone wavelength that is 1216 Å). On the other hand, it is confirmed 372 that the external strong gravitational field changes the atomic structure, making the electron 373 "orbital" energy reduce, the energy level difference decrease, and the spectral shift to red, which 374 explains the spectral red shift from the perspective of the matter structure.

The number of quasars reaches a peak when the redshift value Z=0.3, 0.6, 0.96, 1.41, 1.96, 375 corresponding to the peak value of the quasars when C<sub>M</sub> value is 0.77, 0.63, 0.51, 0.41, 0.33. That 376 is, a more stable quantum conditional region is formed when the difference of C<sub>M</sub> value is about 377 0.1. It indicates that the gravitational field condition is quantized and discontinuous. Although 378 there is only one set of data, it can be concluded that there are several stable quantum conditional 379 regions in the universe (A large amount of astronomical observation data should be able to 380 determine the stable quantum conditions. The Doppler redshift is the rest of the redshift that 381 excludes the  $C_M$  effect ones in the large redshift of the quasar spectrum). The quantum mechanics 382 in Earth region ( $C_M$ = 1) is the quantum mechanics that we have not considered the background 383 conditions of the gravitational field in the past hundred years, or that it does not walk out of the 384 Earth. The detailed data and analysis of the guasar spectral redshift are listed in our previous 385 investigation [42]. 386

According to the formula (12), when  $C_M \rightarrow 0$ , the strength of the gravitational field tends to 387 388 infinity. At this point,  $\Delta E \rightarrow 0$ , and the energy difference of "atomic orbit" is zero. The density of matter is enormous and then to form the dense material that does not interact with the 389 electromagnetic fields. This dense matter is so called dark matter. Therefore, in terms of the 390 material structure, the dark matter is a very dense ordinary matter. As the C<sub>M</sub> changes from 0 to 1, 391 there are atoms with different structures, and the material structures varies with the different 392 densities. From the spectral red shift of the atomic structure, the atomic structure of the white dwarf 393 394 or neutron star is gradually formed. If  $C_M > 1$ , the spectral blue shift. When  $C_M \rightarrow \infty$  and  $\Delta E \rightarrow \infty$ , the energy levels tend to be infinite. Consistently, the atomic structure is instable and it is 395 disintegrated. When  $C_M \rightarrow \infty$  and the strength of the gravitational field tends to zero, the material 396 is instable. It indicates in view of the energy solution quantization mechanics equation that the 397 gravitational field is a condition for the existence of the mass matter and it is an indispensable 398 condition. There is no the mass material if without the external gravitational field. 399

Based on the above discussion, it infers Einstein's equivalence principle is not suitable for the material structure or the quantum mechanics. Since the effect of the gravitational mass  $m_g$  on the material structure in the field is not available for the inertial mass  $m_i$ , it is difficult to combine the general relativity under space-time background with the quantum mechanics related to space-time background. Therefore, the quantum of the gravitational field in this investigation does not adopt the metric field, and the superposition principle is applied to the vacuum gravitational potential.

#### 406 **4 The application of the chiral and quantum gravitational fields**

407 4.1 The application in gravitation theory

408 4.1.1 The interactions between mass matter derived by Newton's law

Taking the Newtonian potential of single particle as an example to show the gravity interaction between the mass matter and anti-mass matter.

411 Newtonian potential: 
$$\varphi_{\alpha} = -\frac{\kappa m_{\alpha}}{\gamma}$$
 ( $\alpha = \begin{cases} R & right - hand mass \\ L & left - hand mass \end{cases}$ )

412 Force F that acts on another particle  $m_{\beta}$  in the  $m_{\alpha}$  field is::

413 
$$\mathbf{F} = -m_{\beta} \frac{\partial \varphi}{\partial \gamma} = -\frac{\kappa m_{\alpha} m_{\beta}}{\gamma^{2}}$$
414 
$$\mathbf{If} \ \alpha = \beta, \quad (\beta = \begin{cases} R \\ L \end{cases}) \quad (m_{R} = m, m_{L} = -m)$$

$$F = -\mathfrak{m}_{\beta} \frac{\partial \varphi}{\partial \gamma} = -\frac{k\mathfrak{m}_{\alpha}\mathfrak{m}_{\beta}}{\gamma^{2}} = \begin{cases} -\frac{k\mathfrak{m}}{\gamma^{2}}(\mathfrak{m}_{\alpha} = \mathfrak{m}_{\beta} = \mathfrak{m}_{R} = \mathfrak{m}) \\ -\frac{k(-\mathfrak{m})(-\mathfrak{m})}{\gamma^{2}} = -\frac{k\mathfrak{m}}{\gamma^{2}}(\mathfrak{m}_{\alpha} = \mathfrak{m}_{\beta} = \mathfrak{m}_{L} = -\mathfrak{m}) \\ \mathbf{F} < \mathbf{0} , \end{cases}$$

415

that is, the "gravity" between matter and matter, or antimatter and antimatter is mutually attractive,a gravity force.

418 If  $\alpha \neq \beta$ ,

$$F = -\frac{km_{\alpha}m_{\beta}}{\gamma^{2}} = \begin{cases} -\frac{km(-m)}{\gamma^{2}} = \frac{kmm}{\gamma^{2}} \begin{pmatrix} m_{\alpha} = m_{R} = m\\ m_{\beta} = m_{L} = -m \end{pmatrix} \\ -\frac{k(-m)(m)}{\gamma^{2}} = \frac{kmm}{\gamma^{2}} \begin{pmatrix} m_{\alpha} = m_{L} = -m\\ m_{\beta} = m_{R} = m \end{pmatrix} \end{cases}$$

419

420 F > 0, the "gravity" between matter and antimatter is mutually exclusive, a repulsion force. 421 4.1.2 The interaction between mass materials in view of the symmetry

The gravitation between the normal matters is mutually attractive, that is, it is mutually attractive between the right-handed mass or the gravitational fields. In view of the symmetry, the gravity between the two left-handed gravitational fields should also be mutually attractive. It is namely mutually attractive between the antimatters. Then, it had to be mutually exclusive between the right-hand and the left-hand gravitational field. That is, the "gravity" between matter and antimatter is the repulsion as well.

Summarily, by deriving and analyzing of the parts 4.1.1 and 4.1.2, it can be proved that the "gravity" between the matter and antimatter is mutually exclusive. As expected, the repulsion of the antimatter to the matter causes the accelerated expansion of cosmic galaxies, which is the real originating of the dark energy phenomenon.

432 4.1.3 Einstein cosmological constant derived by the "gravitation" repulsion between the 433 matter and anti-matter

According to CPT theory, there exists a positive-antimatter sky in the universe. It is the cooperation between the repulsion action of the anti-sky to the positive sky galaxies and the gravity effect of the positive sky on its internal galaxies that can accelerate the expansion of galaxies in the positive sky in the universe, manifested as a dark energy phenomenon.

Under the cosmological assumption, the positive sky (M) and the anti-sky (-M) are applied 438 to the physical laws of FRW (Friedmann-Robertson-Walker), respectively. However, the 439 440 geometry influence cannot be clearly determined due to the large scale, the selection of coordinates and direction of action. Moreover, the positive sky metric tensor under anti-sky action cannot be 441 determined. Therefore, the 4-dimensional space-time tensor equation cannot be applied to the 442 connection between the positive sky and the anti-sky. On one hand, this paper focuses on clarifying 443 the physical principle and does not pay special attention to the calculative details. On the other 444 hand, in considering the role of the mass center in the anti-sky on galaxies in the positive sky, it 445 can be regarded as the role of the Newtonian potential at t moments in a very large-scale space. It 446 is applicable the positive sky evolution under the FRW metric at 1-dimensional time and 3-447 dimensional Euclidean space with the same time horizon. The detailed calculation can be seen in 448 449 Appendix 3 and here the evolution equation of the positive sky under the anti-sky directly describes as: 450

(16)

451 
$$\ddot{a} = \frac{4\pi G}{3} \left( \rho + \frac{3P}{c^2} \right) a + \frac{GM}{a^2 (4 + \gamma^2 - 4\gamma \cos\theta)} (\hat{e}_-)$$
 (13)

452 
$$a\ddot{a} + 2\dot{a}^2 = 4\pi G \left(\rho - \frac{P}{c^2}\right) a^2 - 2kc^2 + \frac{3GM}{a^2(4+\gamma^2 - 4\gamma \cos\theta)}(\hat{e}_-)$$
 (14)

453 
$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8}{3}\pi G\rho + \frac{GM}{a^3(4+\gamma^2 - 4\gamma\cos\theta)}(\hat{e}_-) - \frac{kc^2}{a^2}$$
 (15)

454 
$$\frac{d\rho}{d\tau} + 3\left(\frac{a}{a}\right)\left(\rho + \frac{p}{c^2}\right) = 0$$

- 455 The cosmological constant  $\Lambda = \frac{3GM}{a^3c^2(4+\gamma^2-4\gamma\cos\theta)} = \frac{3GM}{a^3c^2} \frac{1}{4\sin^2\theta + (\gamma-2\cos\theta)^2}$  (17)
- Here, M is the total mass in anti-sky, and it is a definite and invariant value when acts on m.
  m: the mass of galaxy in the positive sky.
- 458  $\theta$ : It is the approximation clip angle between the anti-sky center of mass action on the 459 positive sky galaxy m and the positive sky center of mass action on the m. It is also the spherical 460 coordinate θ angle,  $\theta \in [0,2\pi]$ .
- $\hat{e}_{-}$ : The direction of the anti-sky on galaxy m and it cannot be determined. It is only showed the difference from the metric.
- 463 *a*: It is the radius of the positive sky, which is also a radius in anti-sky.
- 464 G: The Newton's gravitational constant.
- 465 P: The fluid pressure in the positive sky.
- 466  $\rho$ : The energy density of the positive sky.
- 467 K: The curvature of the three-dimensional cosmic space.
- 468  $\gamma a$ : Total motion distance from m to the center of mass in the positive sky.  $\gamma$  is the proportional 469 constant,  $0 < \gamma < 1$ .
- 470 Any two combination of the equation (13), (14), (15), (16) can be called as the Lemaitre 471 equation under the antimatter sky. It is the instantaneous Newtonian potential introduced the anti-472 sky at a very large scale which is obtained by Lemaitre's  $\Lambda$  equation. Here, the cosmic constant 473  $\Lambda = \frac{3GM}{a^3c^2} \frac{1}{4sin^2\theta + (\gamma - 2cos\theta)^2}$ , and  $\Lambda$  is a positive value.

At the fixed moment with a constant a, the value of  $\Lambda$  is only related to the common 474 coordinate  $\gamma$  and the direction angle of the galaxy. When  $\theta$  takes a fixed value (certain spatial 475 orientation),  $\Delta \gamma$  change of  $\gamma$  has little influence on the change of  $\Lambda$  value. When  $\gamma$  takes fixed value 476 (certain spatial radius), the  $\Delta\theta$  change of  $\theta$  also has very small influence on  $\Lambda$ . That is, the radial 477 478 and angular changes have a small impact on the  $\Lambda$  value, manifested as "quasi-constant". It is physically manifested as constant negative energy acting on the cosmic positive sky galaxies. 479 However, the radial and angular variations of the  $\Lambda$  make the cosmic space inhomogeneous, 480 producing a small anisotropy. 481

482 Considering the cosmological evolution of the cosmic age greater than the moment  $T_0$  (the 483 entry of the non-relativistic time), e.g., T>T<sub>0</sub>, the equations (13) and (16) are approximately written 484 as:

485 
$$\ddot{a} = -\frac{4\pi G}{3}(\rho a) + \frac{GM}{a^2(4+\gamma^2-4\gamma \cos\theta)}(\hat{e}_-)$$
 (18)  
486  $\dot{\rho} = -3\left(\frac{\dot{a}}{2}\right)\rho$  (19)

487 The Solution of the equation (19) is  $\rho = \frac{\rho_0}{a^3}$ , and  $\rho_0$  is the density of a = 1. Then substitute 488  $\rho = \frac{\rho_0}{a}$  into the equation (18)

489 It is: 
$$\ddot{a} = -\frac{4\pi G}{3a^2}\rho_0 + \frac{GM}{a^2(4+\gamma^2 - 4\gamma \cos\theta)}(\hat{e}_-)$$
 (20)

490 Let 
$$\ddot{a} = 0$$
, we obtained  $4 + \gamma^2 - 4\gamma \cos\theta = \frac{3M}{4\pi\rho_0}$  (the fixed value) (21)

491 That is, when the co-dynamic coordinates  $\gamma$  and  $\theta$  satisfy the equation (21), the FRW sky 492 transforms from the deceleration expansion to the accelerated expansion. Equation (21) is the 493 transformation point of the FRW accelerated expansion, which is not only radial but also angular. 494 Also, the solution can be get  $\cos\theta = \frac{1}{\gamma} \left(1 - \frac{3M}{16\pi\rho}\right) + \frac{\gamma}{4}$ .

Therefore, the following conclusions can be inferred for the positive sky FRW cosmology under the anti-sky (-FRW) repulsion:

497 1) The physical reason for the accelerated expansion of the universe is due to 498 the chiral symmetry positive and anti-sky. It is the repulsion of the anti-sky that accelerates 499 the expansion of galaxies of the positive sky, manifested as a dark energy phenomenon. 500 The change in the  $\Lambda$  value caused the inflation disturbance is only originated form the 501 term of  $(\frac{1}{4sin^2\theta + (\gamma - 2cos\theta)^2})$ , which is small in radial, angular changes and appears as 502 constant.

503 2) A specific critical radius of cosmic accelerated or deceleration expansion is no longer 504 existed. There are different acceleration or deceleration expansion turning points at different points 505 in space. The universe is non-spherically symmetric with the small spatial anisotropy.

3) When 
$$t \to \infty$$
, then  $a \to \infty$ , thus  $\rho \to 0$ . The equation (15) then transformed to

507 
$$\dot{a}^2 = \frac{GM}{a(4+\gamma^2-4\gamma\cos\theta)}\widehat{e_-} - Kc^2$$

508 When K=0, 
$$a^{\frac{3}{2}}(t) = \frac{3}{2} \left( \frac{GM}{a(4+\gamma^2 - 4\gamma \cos\theta)} \right)^{\frac{1}{2}} t + \text{constant}$$

509 When K=1, 
$$a^{\frac{3}{2}}(t) = \frac{3}{2} \left( \frac{GM}{a(4+\gamma^2 - 4\gamma \cos\theta)} - c^2 \right)^{\frac{1}{2}} t + \text{constant}$$

510 When K=-1, 
$$a^{\frac{3}{2}}(t) = \frac{3}{2} \left( \frac{GM}{a(4+\gamma^2 - 4\gamma \cos\theta)} + c^2 \right)^{\frac{1}{2}} t + \text{constant}$$

511 Whether FRW positive sky is flat, open, or closed, the radius*a*varies is substantially same 512 with the time t.

4) The total kinetic energy of the positive sky (FRW) in the universe.

514 
$$\mathbf{E} = \frac{1}{2}M\dot{a}^2 = \frac{4M\pi G\rho_0}{3a} + \frac{GM^2}{2a(4+\gamma^2 - 4\gamma \cos\theta)} - \frac{1}{2}\mathbf{MK}c$$

The total kinetic energy of the FRW space with the antimatter sky model equation is one more than that of the Friedmann total kinetic energy,  $\frac{GM^2}{2a(4+\gamma^2-4\gamma cos\theta)}$ , which enables the FRW space to accelerated expansion if it matches with  $\cos \theta > \frac{1}{\gamma} \left(1 - \frac{1}{16\pi\rho_0}\right) + \frac{\gamma}{4}$ . Otherwise, it will be a decreased expansion. Since the action of the positive and anti-sky is mutual, the energy in the total universe remains conservation, namely, the law of energy conservation is also held on large scales. 4.2 Application in symmetry and Higgs Mechanism

Now, due to the the limitation of space only the significant conclusions are summarized here (others like the discussion of vacuum symmetry breaking caused by gravitational field and the detailed Higgs mechanism have listed in Appendix 4). Based on the analysis of the Higgs mechanism and the energy solution of the quantization equation of section 3.2, it can be concluded that the mass matter generated by energy matter must have transverse electromagnetic field conditions and longitudinal gravitational field conditions. The stronger region of the positive sky gravitational field is, the matter will be more stable in this region. Meanwhile, the antimatter is 528 obviously more unstable. Similarly, in the strong region of the anti-sky gravitational field, the 529 antimatter is more stable while the matter is more unstable. So, the matter and antimatter cannot 530 be intermingled, and they can only form a chiral symmetric positive-antimatter sky to construct 531 the universe.

#### 532 **5 Discussion**

The boundaries of the positive and antimatter sky will not be areas of extremely intense 533 534 activity. The mass particles cannot reach the boundary because the positive and anti-sky is far away and the boundary area is broad. Thus, the mass particles in the positive material sky side 535 cannot have enough kinetic energy to overcome the continuous force towards the center of the sky 536 to reach the boundary. It is similarly for the mass particles in the anti-material sky. So, the extreme 537 violent phenomenon of the positive and anti-particle annihilation in the positive and anti-sky 538 boundary will not occur. In addition, one kind of the 0 mass energy particles, such as 539 540 electromagnetic wave, can converse the right-hand light to left hand light if they enter the gravitational field from the right-hand side into the left-hand area. The boundary of the energy 541 conversion from the positive sky to anti-sky is equal to the conversion from anti-sky to positive 542 sky due to CPT conservation. There will be no violent energy fluctuations but only the weak energy 543 exchange. Another energy particle, such as neutrinos, have the nature of the gravitational field, let 544 alone a violent energy conversion. 545

This work cannot completely solve the understanding problem of the gravitational field. In different theories (such as Newtonian gravity theory, the general relativity, the symmetry, the gravitational field under quantum mechanics, the gauge theory of the "scalar field", etc.), there are different performances and mathematical expression in the gravitational field. It results in no unified and ineffective integration of gauge theory and gravitational theory.

The cosmological observations are related to the location of the universe we are in, and a large amount of data should be able to determine the general direction of the line of the centroid between the positive sky and anti-sky\_which can be understood as the "magic axis" of the microwave background of K. Land and T. Magueijo[43]. If so, the credibility of the cosmic model of the positive-antisky is also increased.

After all, the gravity is relatively weak, and its effect is not very obvious. However, at least 556 three experiments can be used to confirm or falsify the present theory: (1) anti-hydrogen atomic 557 experiments. Synthetic anti-hydrogen atoms should drift upward in the earth's gravitational field. 558 But the antihydrogen atoms trace is a problem. (2) the positive-antiparticle lifetime experiments 559 or half-life experiments. The positive-antiparticle life is different in the gravity field. The stronger 560 the gravitational field, the greater the positive-antiparticle life difference is. The similar 561 experiments are the half-life of the radioactive matter changes with the strength of the gravitational 562 field. The stronger the gravitational field, the longer the half-life of the radioactive matter in it is. 563 However, it is unlike the difference in positive and antiparticle life, the strength of the gravitational 564 field that could affect the half-life needs to change greatly with the earth's gravitational field due 565

to the more inert Half-life. (3) the parity breaking experiment. Like  $C_0^{60}$ , the gravitational field condition changes while the degree of breaking changes. We do not know whether the experiment can be tested for the weak change of the parity breaking caused by the change of the very small gravitational field strength such as the tide-induced force.

Finally, it should be pointed out that we only define dark matter as ordinary dense matter when the quantum condition factor  $C_M \rightarrow 0$  (gravitational field strength is infinite) in term of mater 572 structure, which infers that the dark matter only exists in the center of the galaxy. But this work 573 does not consider the relative movement of the dark matter found in the literatures [5,6].

#### 574 **5 Conclusions**

The following a set of the interrelated conclusions are derived based on the forementioned analysis:

577 1) We have quantized the gravitational field, and obtained the quantization equation of the 578 gravitational field. Furthermore, the Schrodinger equation and Dirac equation under a variable-579 gravitational field condition have been obtained by variables separation method.

2) The following conclusions have been derived from the quantization equation of the 580 gravitational field: a) The strong change of an external gravitational field will introduce the change 581 of the atomic structure and then result in the spectral shift effect. The redshift of quasar spectrum 582 is just this effect. b) the dark matter is an ordinary dense matter when the gravitational field 583 quantum condition coefficient  $C_M \rightarrow 0$ , that is, the gravitational field tends to be infinite. It is 584 recognized to be the essence of the dark matter in view of the matter structure. c) Due to the 585 different gravitational field strength (e.g., the different C<sub>M</sub> values), there are several discrete stable 586 quantum condition regions in the universe, and the Earth region is only that of the gravitational 587 field quantum condition factor C<sub>M</sub>=1. 588

3) The universe is composed of the chiral symmetric positive matter sky and antimatter sky. The "gravitational" interaction between the positive and anti-sky is exclusive. It is the repulsion that accelerates the expansion of galaxies in the positive sky and appears as a dark energy phenomenon, which is the physical essence of dark energy.

4) The cosmological constant is obtained by this calculation, which has small changes in both the radial and angular directions, producing a small anisotropy of the cosmic space.

595 5) The positive-anti baryon number in the universe is equal and the total energy of the 596 positive-anti sky is conserving.

6) Not only a vacuum symmetry breaking but also a parity breaking could be caused by the gravitational field. The stronger the gravitational field will lead to the greater breaking degree.

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655

## Appendix 1. Evidence of the CPT theorems under the definition of manual performance energy and mass

- This appendix defines various fields in the presence of CPT invariance, and then proves the CPT theorem:
- 649 1.1 The quantization of the free electromagnetic field (spin is 1 and m=0)
- From the invariance of CPT, there are two kinds of vacuum: one is the positive matter vacuum
- $(ε_0 > 0, μ_0 > 0)$ , and the right-hand energy is in it,  $K_R = K$ . The other one is antimatter vacuum (ε<sub>0</sub>  $< 0, μ_0 < 0$ ), and there is the left-hand energy in it,  $K_L = -K$ .

Let A (t, r) to be the vector potential of the free electromagnetic field, and satisfy the "transverse condition" divA = 0, then scalar potential  $\phi = 0$ . Whereas, field E and H are for:

$$E = -\dot{A}$$
,  $H = rot A$ 

- 656 Maxwell's equations can be transformed into wave equations of A:  $\Delta A \frac{\partial^2 A}{\partial t^2} = 0$
- In classical electrodynamics, the localized fields can be expanded into the plane waves, and
   their potentials can be represented as a series (Fourier expansion):

659 
$$A = \sum_{K} (a_{K}e^{iK\cdot\gamma} + a_{K}e^{-iK\cdot r})$$
  
660 Where  $a_{K}$  is a function of time,

- 660 Where  $a_K$  is a function of time, 661  $a_K \sim e^{-i\omega K}, \ \omega = |K|$
- (3)

(2)

(1)

Define the regular variable for the field:  $Q_K = \frac{1}{\sqrt{4\pi}}(a_K + a_K^*), P_K = \frac{i\omega}{\sqrt{4\pi}}(a_K - a_K^*) = \dot{Q}_K$ 662 Each vector of  $P_K$  and  $Q_K$  is perpendicular to the wave vector K, that is, there is two 663 independent components. The direction of the vector determines the polarization direction of the 664 corresponding wave, which the two components are represented by using  $Q_{K\alpha}$  and  $P_{K\alpha}$  ( $\alpha = 1, 2$ ). 665 Now, quantizing the free field, and then the classical description of the above field could be 666 transited to the quantum theory. Meanwhile, the regular variables could be treated as operators that 667 satisfy the commutation relations:  $\hat{P}_{K\alpha}\hat{Q}_{K\alpha} - \hat{Q}_{K\alpha}\hat{P}_{K\alpha} = -i$ . All operators with the different K and 668  $\alpha$  can be commutative with each other. The potential A and the fields E and H (by formula (1)) 669 670 also form the Hermitian operators. Define the operators  $\hat{C}_{K\alpha} = \frac{1}{\sqrt{Z\omega}} (\omega \hat{Q}_{K\alpha} + i \hat{P}_{K\alpha}), \quad C^+_{K\alpha} = \frac{1}{\sqrt{Z\omega}} (\omega \hat{Q}_{K\alpha} - i \hat{P}_{K\alpha})$ 671  $\hat{C}_{Ka}$  and  $\hat{C}_{Ka}^+$  meet the commutation relationship  $\hat{C}_{Ka}\hat{C}_{Ka}^+ - C_{K\alpha}^+\hat{C}_{Ka} = 1$ 672 The operator of the potential A is (see formula (2)) 673  $\hat{A} = \sum_{K} (\hat{C}_{K\alpha} A_{K\alpha} + C^{+}_{K\alpha} A^{*}_{K\alpha})$ (4)674 where 675  $A_{K\alpha} = \sqrt{4\pi} \frac{e^{(\alpha)}}{\sqrt{2\omega}} e^{iK \cdot r}$ (5)676 The symbol  $e^{(\alpha)}$  is the unit vector of the oscillator polarization direction, and it is 677 perpendicular to the wave vector K. Also, it has two independent polarization directions for each 678 K, similarly to write the operator 679  $\hat{E} = \sum_{K,\alpha} (\hat{C}_{K\alpha} E_{K\alpha} + \hat{C}_{K\alpha}^{+} E_{K\alpha}^{*})$ 680  $\widehat{H} = \sum_{K\alpha} (\widehat{C}_{K\alpha} H_{K\alpha} + \widehat{C}_{K\alpha}^+ H_{K\alpha}^*)$ (6)681 Where  $E_{K\alpha} = i\omega A_{K\alpha}$ ,  $H_{K\alpha} = n \times E_{K\alpha}$ 682  $E^*_{K\alpha} = -i\omega A^*_{K\alpha}, \ H^*_{K\alpha} = n \times E^*_{K\alpha}$ 683 (7)

In the vacuum of  $\varepsilon_0 < 0$ ,  $\mu_0 < 0$ , the left-hand energy is running. The three vectors of electric magnetic field strength vector E, the magnetic field strength vector H, and the wave vector K form the left-hand system, thus  $K_L = -K_R = -K$ . Then, the  $a_K$ ,  $Q_K$ ,  $P_K$ ,  $\hat{P}_{Ka}$ ,  $\hat{C}_{Ka}$ ,  $\hat{E}$ ,  $\hat{H}$ ,  $A_{K\alpha}$  and their corresponding conjugators are defined as the right hand energy. Then, define each quantity in the left-hand energy as a complex vector (comparison with the formula (3))  $a_{-K} \sim e^{i\omega K} = a_K^*$ Thus, let

690 
$$a_{-K} = a_K^*, \ a_{-K}^* = a_K$$
 (8)

692 
$$P_{-K} = \frac{\iota\omega}{\sqrt{4\pi}} (a_{-K} - a_{-K}^*) = \frac{1}{\sqrt{4\pi}} (a_K^* - a_K) = -P_K$$

After the operation of 
$$Q_{-K} = Q_K$$
,  $P_{-K} = -P_K$ , it can be defined similarly:

694 
$$\hat{C}_{-Ka} = \frac{1}{\sqrt{2\omega}} \left( \omega \hat{Q}_{-Ka} + i \hat{P}_{-Ka} \right) = \frac{1}{\sqrt{2\omega}} \left( \omega \hat{Q}_{-Ka} - i \hat{P}_{-Ka} \right) = C_{Ka}^+$$

695 
$$\hat{C}_{Ka}^{+} = \frac{1}{\sqrt{2\omega}} \left( \omega \hat{Q}_{-Ka} - i \hat{P}_{-Ka} \right) = \frac{1}{\sqrt{2\omega}} \left( \omega \hat{Q}_{Ka} + i \hat{P}_{Ka} \right) = \hat{C}_{Ka}$$
 (9)  
697  $\hat{A}_{L} = \sum \left( \hat{C}_{-Ka} A_{-Ka} + \hat{C}_{-Ka}^{+} A_{-Ka}^{+} \right)$ 

$$\hat{A}_L = \sum_{-Ka} \left( \hat{C}_{-Ka} A_{-K\alpha} + \hat{C}_{-Ka}^+ A_{-K\alpha}^+ \right)$$

696 
$$A_{-K\alpha} = -\sqrt{4\pi} \frac{e^{(\alpha)}}{\sqrt{2\omega}} e^{-iK \cdot r} = -A_{K\alpha}^*,$$

With the left-hand energy, the vector A takes negative values, 698

699 
$$A^*_{-K\alpha} = -\sqrt{4\pi} \frac{e^{(\alpha)}}{\sqrt{2\omega}} e^{iK \cdot r} = -A_{K\alpha}$$

$$E_{-K\alpha} = i\omega A_{-K\alpha} = -i\omega A_{K\alpha}^* = E_{K\alpha}^*$$

701 
$$E_{-K\alpha}^* = -i\omega A_{-K\alpha}^* = -i\omega(-A_{K\alpha}) = i\omega A_{K\alpha} = E_{K\alpha}$$

$$H_{-K\alpha} = n_L \times E_{-K\alpha} = -n \times E_{K\alpha} = -H_{K\alpha}$$

$$H_{-K\alpha}^* = n_L \times E_{-K\alpha}^* = (-n) \times E_{K\alpha} = -H_{K\alpha}$$
(10)

704 Then it can be derived:

705 
$$\hat{E}_{L} = \sum_{-\kappa a} (\hat{C}_{-\kappa a} E_{-\kappa a} + \hat{C}_{-\kappa a}^{+} E_{-\kappa a}^{*}) = \sum_{\kappa a} (\hat{C}_{\kappa a}^{+} E_{\kappa a}^{*} + \hat{C}_{\kappa a} E_{\kappa a}) = \hat{E} = \hat{E}_{R}$$

706 
$$\widehat{H}_{L} = \sum_{-\kappa a} \left( \widehat{C}_{-\kappa a} H_{-\kappa a} + \widehat{C}_{-\kappa a}^{+} H_{-\kappa a}^{*} \right) = \sum_{\kappa a} \left[ \widehat{C}_{\kappa a}^{+} (-H_{\kappa a}^{*}) + \widehat{C}_{\kappa a} (-H_{\kappa a}) \right]$$

$$= -\sum_{Ka} \left( C_{K\alpha}^{+} H_{K\alpha}^{*} + \hat{C}_{Ka} H_{K\alpha} \right) = -\hat{H} = -\hat{H}_{R}$$

708 
$$A_{L} = \sum_{-Ka} (\hat{C}_{-Ka}A_{-K\alpha} + \hat{C}_{-Ka}A_{-K\alpha}^{*}) = \sum_{Ka} [\hat{C}_{Ka}(-A_{K\alpha}^{*}) + \hat{C}_{Ka}(-A_{K\alpha})]$$
709 
$$= -\sum_{Ka} (\hat{C}_{Ka}A_{K\alpha}^{*} + \hat{C}_{Ka}A_{K\alpha}) = -\hat{A} = -\hat{A}_{R}$$
(11)

$$= -\sum_{Ka} \left( \hat{C}_{Ka}^{+} A_{K\alpha}^{*} + \hat{C}_{Ka} A_{K\alpha} \right) = -\hat{A} = -\hat{A}_{R}$$
  
Since the above quantization of the electromagnetic field

Since the above quantization of the electromagnetic field start from the classical definition, 710 the production and annihilation operators of the electromagnetic field are expressed by  $\hat{C}_{Ka}$  to 711 avoid confusion with the classical plane wave expansion. Then, when the spin is defined as 0 or 712 1/2 field, both *a* or  $\alpha$  are used to represent the generation and annihilation operators. 713

The plane wave is defined as spin is 1 and mass m of particles does not require to be 0, 714 similarly with the electromagnetic field with 0 mass and spin =1. Corresponding to the right-hand 715 plane wave  $A_R = \frac{1}{\sqrt{2\omega\Omega}} e^{iK_R \cdot r}$  is the left-hand plane wave. By referencing to 0 mass definition, it 716 has  $A_L = \frac{1}{\sqrt{2+0}} e^{iK_L \cdot r}$ , here,  $K_R = K = -K_L$ , and the Fourier expansion is: 717

718 
$$A_{R}(t,r) = \sum_{KR} \frac{1}{\sqrt{2\omega\Omega}} \left[ a_{K_{R}}(t)e^{iK_{R}\cdot r} + a_{K_{R}}^{+}(t)e^{-iK_{R}\cdot r} \right]$$

720 
$$= \sum_{K} \frac{1}{\sqrt{2\omega\Omega}} \left[ a_{K}(t)e^{iK\cdot r} + a_{K}^{+}(t)e^{-iK\cdot r} \right]$$

719 
$$(12)$$

721 
$$A_{L}(t,r) = \sum_{K_{L}} \frac{1}{\sqrt{2\omega\Omega}} \left[ a_{K_{L}}(t)e^{iK_{L}\cdot r} + a_{K_{L}}^{+}(t)e^{-iK_{L}\cdot r} \right]$$

$$= -\sum_{-K} \frac{1}{\sqrt{2\omega\Omega}} \left[ a_{-K}(t) e^{-iK \cdot r} + a_{-K}^{+}(t) e^{iK \cdot r} \right]$$

$$= -\sum_{K} \frac{1}{\sqrt{2\omega\Omega}} \left[ a_{K}^{+}(t)e^{-iK\cdot r} + a_{K}(t)e^{iK\cdot r} \right] = -A_{R}(t,r)$$
(13)

723

Such as the transformation of  $C_{K\alpha}$  in the electromagnetic field with 0 mass and spin 1, here 725  $a_{K}^{+} = a_{-K}, \ a_{K} = a_{-K}^{+}$ 726

1.2 Secondary quantization of the field where spin is 0 727

Following the quadratic quantization methods, the arbitrary wave functions unfold by an 728 eigen function of a complete set of possible states, such as the plane waves could be represent as: 729

731 
$$\phi = \sum_{P} a_{P} \phi_{P}, \quad \phi^{*} = \sum_{P} a_{P}^{*} \phi_{P}^{*}$$

 $\phi_P = \frac{1}{\sqrt{2\omega\Omega}}e^{-iPx}$ ,  $\phi^*$  could be viewed as the plane wave of -P:  $\phi_{-P}^* = \frac{1}{\sqrt{2\omega\Omega}}e^{iPx}$ . 730 As defined by the four-momentum momentum,  $\phi_{P_R} = \frac{1}{\sqrt{2\omega\Omega}} e^{-iP_R x} = \frac{1}{\sqrt{2\omega\Omega}} e^{-iPx}$ 732

 $\phi_{P_L} = \frac{1}{\sqrt{2\omega\Omega}} e^{-iP_L x} = \frac{1}{\sqrt{2\omega\Omega}} e^{iPx}$ ,  $\phi_R$  and  $\phi_L$  conjugate each other 733

Corresponding to the annihilation operator  $a_P$  of the particle, the antiparticle production 734 operator  $b^+$  constitutes a complete set. 735

736  

$$\phi(t,r) = \sum_{P} \frac{1}{\sqrt{2\omega\Omega}} \Big[ a_{P}(t)e^{-iP_{R}\cdot r} + b_{P}^{+}(t)e^{iP_{L}\cdot r} \Big]$$
738  

$$= \sum_{P} \frac{1}{\sqrt{2\omega\Omega}} \Big[ a_{P}(t)e^{-iP\cdot r} + b_{P}^{+}(t)e^{-iP\cdot r} \Big]$$

738
$$=\sum_{P} \frac{1}{\sqrt{2\omega\Omega}} \left[a_{P}(t)e^{-tP\cdot t} + b_{P}^{+}(t)e^{-t}\right]$$
737
$$(14)$$

737 739

conjugation:

739 The Hermit  
740 
$$\phi_{(t,r)}^{+} = \sum_{P} \left[ a_{P}^{+}(t)e^{iP\cdot r} + b_{P}(t)e^{-iP\cdot r} \right]$$

1.3 The secondary quantization of the spin-1/2 field 741 In the four-vectors,  $P^{\mu} = \begin{pmatrix} P_R^{\mu} \\ P_{\cdot}^{\mu} \end{pmatrix}$ , the right-hand plane wave  $\psi_R(x)$  and the left-hand plane 742

743 wave could be introduced:

744 
$$\psi_R(x) = \frac{1}{\sqrt{\Omega}} u_{P_R,S} e^{-iP_R x} = \frac{1}{\sqrt{\Omega}} u_{P,S} e^{-iP x}$$

745 
$$\psi_{L}(x) = \frac{1}{\sqrt{\Omega}} v_{P_{L},S} e^{-iP_{L}x} = \frac{1}{\sqrt{\Omega}} v_{-P,S} e^{iPx}$$

Here the spinor transformation follows the agreement of the phase factor: 746

(i)  $v_{P,S} = \gamma_2 u_{P,S}^*$ ,  $u_{-P,S} = \gamma_2 v_{P,S}^*$ 747

(ii)  $\gamma_o u_{P,S} = u_{-P,-S}, \ \gamma_o v_{P,S} = -v_{-P,-S}$ 748

749 (iii) 
$$\sigma_2 u_{PS}^* = e^{i\theta_{PS}} u_{-PS}, \ \sigma_2 v_{PS}^* = e^{-i\theta_{-PS}} v_{-PS}, \ \text{ $\pm e^{i\theta_{PS}} = -e^{i\theta_{-PS}} = -e^{i\theta_{-PS}} - e^{i\theta_{-PS}} = -e^{i\theta_{-PS}} - e^{i\theta_{-PS}} - e^{i\theta_{-PS}} = -e^{i\theta_{-PS}} - e^{i\theta_{-PS}} - e^{i\theta_{-PS}}$$

Here,  $\gamma$  is the Dirac matrix. 750

At any given moment t, like the original-spin 1/2 field, the spin 1/2 field operator under the 751 left-handed four-momentum is introduced and  $\psi(x) = \psi(t, r)$  could be extended with the Fourier 752 753 series:

755 
$$\psi(t,r) = \sum_{P} S_{P}(t) \frac{e^{iP \cdot r}}{\sqrt{\Omega}}$$

754 (15)

Where,  $S_P(t)$  is 4×1 matrix which is not relative with  $\gamma$ , and  $\Omega$  is the volume taken. It also 756 takes *C*-number base vetors, like  $u_{P_R,S} = u_{P,S}$ ,  $v_{P_L,S} = v_{-P,S}$  in the spinor space and all meet with 757 the Dirac equation  $(\gamma^{\mu}P_{\mu} + m)u = 0$  that is,  $(\alpha \cdot P_i + m)u = E_P u$ . 758

759 Here m = 
$$\begin{cases} m_R \\ m_L \end{cases} = \begin{cases} m \\ -m \end{cases}$$
,  $E_P = \begin{cases} E_R \\ E_L \end{cases} = \begin{cases} E_P \\ -E_P \end{cases}$ ,  $u = \begin{cases} u_{PR,S} \\ v_{PL,S} \end{cases} = \begin{cases} u_{P,S} \\ v_{-P,S} \end{cases}$ 

760

761

 $E_L$ ,  $P_L$ ,  $m_L$ :

In above formula, substitute a group  $E_R$ ,  $P_R$  (including three vectors) simultaneously,  $m_R$  or

 $(\sigma \cdot P + m)u_{PS} = E_P u_{PS}$ 762  $(\sigma \cdot P - m)v_{-P,S} = -E_P v_{-P,S}$ 763 To normalizing the vector, let  $u_{P,S}^+ u_{P,S} = v_{-P,S}^+ v_{-P,S} = 1$ 764  $u_{P,S}$  and  $v_{-P,S}$  form a complete set of the orthogonal base vectors in spinor space,  $S_P(t)$  could 765 be deployable by this set of base vectors: 766  $S_P(t) = \sum_{s=\pm\frac{1}{2}} \left[ a_{P,S}(t) u_{P,S} + b^+_{-P,S}(t) v_{-P,S} \right]$ (16)767 In the formula,  $a_{P,S}(t)$  and  $b^+_{-P,S}(t)$  is the Operator in the Hilbert space, and combining (15) and 768 (16), we have 769  $\psi(x) = \psi(t,r) = \frac{1}{\sqrt{O}} \sum_{P,S} \left[ a_{P,S}(t) u_{P,S} e^{iP \cdot r} + b_{P,S}^+(t) v_{P,S} e^{-iP \cdot r} \right]$ (17)770  $\psi^{+}(x) = \psi^{+}(t,r) = \frac{1}{\sqrt{\Omega}} \sum_{P,S} \left[ a_{P,S}^{+}(t) u_{P,S}^{+} e^{-iP \cdot r} + b_{P,S}(t) v_{P,S}^{+} e^{iP \cdot r} \right]$ 771 (18)772 Below is a proof of CPT's theorem under the definition of a chiral field. 773 Considering the localized field theory, where there is  $N_i$  the field of the Spin j and it is 774 represented as: 775 Spin 0:  $\phi_1(x)$ ,  $\phi_2(x)$  ...  $\phi_{N_0}(x)$ 776 Spin  $\frac{1}{2}$ :  $\psi_1(x)$ ,  $\psi_2(x) \dots \psi_{N_{\frac{1}{2}}}(x)$ 777 Spin 1:  $[A_1(x)]_{\mu}$ ,  $[A_2(x)]_{\mu} \dots [A_{N_1}(x)]_{\mu}$ (19)778 Under the Lorentz group or the C.P.T transformation, in terms of transformation properties, 779 Field is available a symmetric tensor representation with the spin j as an integer:  $T_{\mu_1...\mu_i}(x)$ 780 (20)781 A field with a spin j as a half-integer is represented as  $S_{\mu_1...\mu_{j-\frac{1}{2}}}(x) \sim T_{\mu_1...\mu_{j-\frac{1}{2}}}\psi_{\alpha}(x)$ 782 (21)It can be seen as transforming as the direct product like the order symmetric tensor of 783  $T_{\mu_1 \dots \mu_{j-\frac{1}{2}}}$  with a spin j as a half-integer and Dirac spinor. 784 Assuming the Lagrangian density 785  $\mathcal{L}(x) = \left(\frac{\partial}{\partial x_{a}}, \phi_{a}, \phi_{a}^{+}, \psi_{b}, \psi_{b}^{+}, (A_{c})_{\mu}, (A_{c}^{+})_{\mu} \dots\right)$  sum of the formal products. 786 All of the fields are taken at the same space time point x,  $x_{\mu} = (it, r)$ 787 Operator  $\oint \equiv CPT$  (or the other arrangement and combination forms of the CPT) 788 (22)789 The theorems can be Introduced: Any the Lorentz-invariant  $\mathcal{L}(x)$  meets with  $\oint \mathcal{L}(x) \oint^{-1} =$ 790  $\mathcal{L}^+(-x)$ 791 If we make the following choices 792 For all  $a = 1, 2, ..., N_0$ ,  $\oint \phi_a(x) \oint^{-1} = \phi_a^+(-x)$ 793 For all  $b = 1, 2, \dots N_{\frac{1}{2}}, \quad \oint \left(\psi_b(x)\right)_{\alpha} \oint^{-1} = i(\gamma_5)_{\alpha\beta} \left(\psi_b^+(-x)\right)_{\beta}$ (24)794 For all  $c = 1, 2, ..., N_1$ ,  $\oint (A_c(x))_{\mu} \oint^{-1} = -(A_c^+(-x))_{\mu}$  (25) 795 For all integer j field (20),  $\oint T_{\mu_1...\mu_i}(x) \oint^{-1} = (-1)^j T^+_{\mu_1...\mu_i}(-x)$ 796 (26)For all semi-integer j fields (21), 797

798 
$$\oint S_{\mu_1 \dots \mu_{j-\frac{1}{2}},\alpha}(x) \oint^{-1} = (-1)^j (i\gamma_5)_{\alpha\beta} S^+_{\mu_1 \dots \mu_{j-\frac{1}{2}},\beta}(-x)$$
(27)

Proof: 799

(1) Consider the field with the spin 1/2800

under the introduction of the defined spin 1/2 field of the left-hand energy, the final Fourier 801 expansion is formula 17. 802

803 
$$\psi(x) = \frac{1}{\sqrt{\Omega}} \sum_{P,S} \left[ a_{P,S}(t) u_{P,S} e^{iP \cdot r} + b_{P,S}^+(t) v_{P,S} e^{-iP \cdot r} \right]$$

(1) if  $Ca_{P,S}C^+ = \eta_c b_{P,S}$ 804  $Cb_{P,S}^+C^+ = \eta_c a_{P,S}^+$ 805

the conjugated  $Ca_{P,S}^+C^+ = \eta_c^+ b_{P,S}^+$ 

and under the spinor transformation convention of the formula (19), it has  $C\psi(x)C^+ =$ 807  $\eta_c \varphi^c(x)$ (28)808 ② if  $Pb_{P,S}^+P^+ = -\eta_P b_{-P,-S}^+$ 

809

806

810 
$$Pa_{P,S}P^+ = \eta_P a_{-P,S}$$

811 the conjugated 
$$Pa_{P,S}^+P^+ = \eta_P^*a_{-P,-S}^+$$

And formula (19), it has  $P\psi(t,r)P^+ = \eta_P \gamma_0 \psi(t,-r)$ (29)812

813 ③ if 
$$Ta_{P,S}T^{-1} = e^{-i\theta_{-P,S}}a_{-P,S}$$

814 
$$Tb_{P,S}^{+}T^{-1} = e^{i\theta_{-P,S}}b_{-P,S}$$

815 the conjugated 
$$Ta_{P,S}^{-1}T^{-1} = e^{i\theta_{-P,S}}a_{-P,S}^{-1}$$
  
816  $Tb_{P,S}T^{-1} = e^{-i\theta_{-P,S}}b_{-P,S}$ 

817 and formula (19), it has 
$$T\psi(t,r)T^{-1} = \eta_t \sigma_2 \psi(-t,r)$$
 (30)

by formula (28), (29) and (30) 818

819 
$$\oint \psi(x) \oint^{-1} = CPT\psi(t,r)T^{-1}P^{-1}C^{-1} = CP\eta_t \sigma_2 \psi(-t,r)P^{-1}C^{-1}$$
  
820 
$$= C\eta_t \eta_P \sigma_2 \gamma_0 \psi(-t,-r)C^{-1} = \eta_t \eta_P \eta_c \sigma_2 \gamma_0 \psi^c(-x) = \eta(ir_5)\psi^+(-x)$$

Here, 
$$\psi^c = \gamma_2 \psi^+$$
,  $\sigma_2 \gamma_0 \gamma_2 = \sigma_2 \rho_2 \rho_3 \sigma_2 = -i\rho = i\gamma_5$ , let  $\eta_t \eta_P \eta_c = 1$ 

It matches with the conditions of CPT invariance (24) 822

(2) Consider a field with a spin=0823

The final Fourier expansion is formula (14) by defining the spin as 0 field with left-hand 824 825 energy

826 
$$\phi(t,r) = \sum_{P} \frac{1}{\sqrt{2\omega\Omega}} \left[ a_P(t) e^{-iP \cdot r} + b_P^+(t) e^{-iP \cdot r} \right]$$

827 Under the transform of 
$$Ca_PC^+ = \eta_c b_P$$
,  $Pa_PP^+ = \eta_P a_{-P}$ ,  $Ta_PT^{-1} = \eta_T a_{-P}$   
828  $Ca_P^+C^+ = \eta_c^* b_P^+$ ,  $Pa_P^+P^+ = \eta_P^* a_{-P}^+$ ,  $Ta_P^+T^{-1} = \eta_T^* a_{-P}^+$ 

$$\mathcal{L}u_{P}\mathcal{L}^{*} = \eta_{c}u_{P}, \quad \mathcal{P}u_{P}\mathcal{P}^{*} = \eta_{P}u_{-P}, \quad \mathcal{I}u$$

829 It has 
$$C\phi(x)C^+ = \eta_c \phi^+(x)$$

830 
$$P\phi(t,r)P^{+} = \eta_{P}\phi(t,-r)$$

831  $T\phi(t,r)T^{-1} = \eta_t\phi(-t,-r)$ 

That is,  $\oint \phi(x) \oint^{-1} = \phi^+(-x)$ , and it matches with the conditions of CPT invariance (23) 832 (3) Consider a field with a spin=1833

The final Fourier expansion is formula (12) and (13) by defining the spin as 1 field 834 with left-hand energy. 835

$$A(t,r) = \sum_{K} \frac{1}{\sqrt{2\omega\Omega}} \left[ a_{K}(t)e^{iK\cdot r} + a_{K}^{+}(t)e^{-iK\cdot r} \right]$$

under the transform of 
$$Pa_{K}^{+}P^{+} = -a_{K}^{+}$$
,  $Ca_{K}^{+}C^{+} = -a_{K}^{+}$ ,  $Ta_{K}^{+}T^{-1} = -a_{-K}^{+}$   
 $Pa_{K}P^{+} = -a_{-K}$ ,  $Ca_{K}C^{+} = -a_{K}$ ,  $Ta_{K}T^{-1} = -a_{-K}^{+}$  it has

$$Pa_{K}P^{+} = -a_{-K}, \quad Ca_{K}C^{+} = -a_{K}, \quad Ta_{K}T^{+} = -a_{-K}^{-} \text{ if n}$$

$$PA(x)P^{+} = PA(t,r)P^{+} = -A(t,-r), \quad CA(x)C^{+} = -A(x),$$

839 
$$PA(x)P^+ = PA(t,r)P^+ = -A(t,-r), \quad CA(x)C^+ = -A$$

 $TA(t,r)T^{-1} = -A(-t,r)$ 840

841 That is, 
$$\oint (A_C(x))_{\mu} \oint^{-1} = -(A_C^+(-x))_{\mu}$$

Here  $A_{\mu} = (iA_0, A)$ , and  $A = A^+$  is Hermitian. Whereas,  $A_0 = -A^+$  is anti-Hermitian. 842 However, i becomes -i under the time inversion. So  $A_{\mu} = A_{\mu}^{+}$  will meet with the conditions for 843 CPT invariance of the formula (25). 844

The spin 0.1 and 1/2 fields based on the definition of the above right hand or the left-hand 845 energy meet with the conditions of CPT invariance under CPT joint operation, where the proof of 846 the CPT theorem  $\oint \mathcal{L}(x) \oint^{-1} = \mathcal{L}^+(-x)$  is consistent with any textbook on particle physics, and 847 thus here it is no longer repeated discussed. 848

From the CPT's theorem, the following conclusions are drawn: 849

1) The existence of right-handed energy  $E_R$  and the chiral symmetrical left-handed energy 850  $E_{L}$ . It is set  $E_{R}=E$  during the calculation which is known as the positive energy while  $E_{L}=-E$  known 851 as the negative energy. 852

2) There is right-handed mass matter (also called normal matter  $m_R=m$ ) and left-handed mass 853 matter (also called antimatter m<sub>L</sub>=-m). This mass chirality is the intrinsic property of matter and 854 does not change due to the selection of reference system. 855

3) The mass matter carries the gravitational field and the mass matter exhibits chiral 856 symmetry. Then, the gravitational field accordingly presents chiral symmetry. That is, the normal 857 matter carries a right-handed gravitational field and antimatter carries a left-handed gravitational 858 field. 859

4) There is a microscopic world of chiral symmetric particles, and also there is a macroscopic 860 chiral symmetry universe. Such universe consists of the positive matter sky and the antimatter sky 861 862 with the chiral symmetry (the matter and antimatter in the universe form their own positive and antimatter sky, respectively. But it cannot be a mixture of matter and anti-matter to form the 863 universe which will be discussed in the Appendix4 particle physics parts). In a physical vacuum 864 in the sky of matter, it is  $\varepsilon_0 > 0$ ,  $\mu_0 > 0$ . While, in a vacuum for an antimatter sky, it is  $\varepsilon_0 < 0$ , 865  $\mu_0 < 0$ . It means there are two gravitational field vacuums in the universe. The defined "positive 866 sky" is actually what people now call the universe. In fact, the cosmic model in this manuscript 867 doubles the original universe (a pair of positive and anti-sky). 868

869

836

#### Appendix 2. Analysis of the quantum mechanics equation under the gravitational field 870

Under the energy scale nucleon mass m=1GeV condition, the strengths of the four 871 interactions are divided into the strong 1, electromagnetic 10<sup>-2</sup>, weak 10<sup>-5</sup>, gravitational 10<sup>-38</sup>, 872 respectively. The gravitational interaction is extremely weak than the other three interactions, so 873 the quantum mechanics described the microscopic particle interactions does not consider the 874 gravitational field at all. However, all the physical behavior occurs in the gravitational field and 875 thus it is inseparable from the gravitational field, which forces us to consider the influence of the 876 gravitational field. Meanwhile, the gravitational field is also a longitudinal condition for the 877 878 generation of mass matter, so a complete quantum mechanical system must include gravity.

3.1 Quantum mechanics equation of the gravitation field (3.6)

880 
$$k(-\frac{\hbar^2}{2m}\gamma^0\nabla^2 + \gamma\hat{p})\psi(\mathbf{r}, t, \sigma) = \frac{4\pi G}{c^2}i\hbar\gamma^0\frac{\partial}{\partial t}\psi(\mathbf{r}, t, \sigma) \quad (1)$$

When the mass matter exists at a mass particle and the transverse energy without the longitudinal energy, the equation (3.6) transforms into the Schrodinger equation:

883 
$$C_{M}\left(-\frac{\hbar^{2}}{2m}\nabla^{2}\right)\psi(\mathbf{r},t) = i\hbar\frac{\partial}{\partial t}\psi(\mathbf{r},t) \qquad (2)$$

Equation (2) is the Schrodinger equation on the background of the gravitational field.

When all the mass material is converted into the longitudinal gravitational field energy, there is no transverse electromagnetic field energy, and thus the equation (3.6) becomes a massless spinor field Dirac equation:

888 
$$C_{M}\gamma \cdot \hat{p}\psi(r,t,\sigma) = i\hbar\gamma^{0} \frac{\partial}{\partial t}\psi(r,t,\sigma)$$
(3)

Equation (3) is the Dirac equation in the context of a gravitational field. Equations (2) and (3) can also be seen as the equation obtained by separating variables from equations (1).

The energy solutions to the Schrodinger equation and Dirac equation in the context of the gravitational field are explained as follows:

893 (1) One-dimensional fixed state

894 The Schrodinger equation is 
$$H'\psi(x) = E\psi(x)$$

895 
$$H' = C_M H = C_M \left[ -\frac{\hbar^2}{2m} \nabla^2 + \mathsf{V}(x) \right]$$

896 One-dimensional square potential trap  $\lor (x) = \begin{cases} 0, & 0 < x < a \\ \infty, & x < 0 \text{ or } x > a \end{cases}$ 

897 The Schrodinger equation in the one-dimensional fixed situation well is:  $-C_M \frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi = E\psi$ 

898 Let 
$$K = \sqrt{\frac{2mE}{C_M \hbar^2}}$$
, we obtain:

899 
$$\psi(x) = A\sin(Kx + \sigma), \sin Ka = 0, Ka = n\pi, n = 1,2,3...$$

900 
$$\frac{n\pi}{a} = \sqrt{\frac{2mE}{C_M\hbar^2}}, \quad E = \frac{C_M\hbar^2\pi^2n^2}{2ma^2}$$

901 The level distance between the adjacent energy levels:

902 
$$\Delta E_n = \frac{C_M \hbar^2 \pi^2}{2ma^2} [(n+1)^2 - n^2] = \frac{C_M \hbar^2 \pi^2}{ma^2} n$$
(4)

803 Known from the formula (4), if the gravitational field strength coefficient  $C_M$  changes, the 804 energy level  $\Delta E_n$  also changes. When the gravitational field increases,  $C_M$  decreases, and  $\Delta E_n$ 805 decreases. That is, the the energy level difference decreases with the external gravitational field 806 strength increases. Therefore, the change of the field strength of the external gravitational field (or 807 background gravitational field) will alter the atomic structure of the material.

908 (2) Hydrogen atoms

910

909 The energy eigen equation is 
$$H'\psi = E\psi$$

$$H' = C_M H = C_M \left[ -\frac{\hbar^2}{2\mu} \nabla^2 + \mathsf{V}(\gamma) \right]$$

911  $\mu$ : electronic quality,  $\forall (\gamma) = -\frac{e^2}{\gamma}$ , The Coulomb action energy.

912 Substitute the H' to the eigen equation, it gets

913 
$$C_M \left[ -\frac{\hbar^2}{2\mu} \nabla^2 + \vee (\gamma) \right] \psi = E \psi$$
 (5)

914 Take the ball coordinates, now

915 
$$\nabla^2 = \frac{1}{\gamma^2} \frac{\partial}{\partial \gamma} \gamma^2 \frac{\partial}{\partial \gamma} - \frac{l^2}{\hbar^2 \gamma^2} = \frac{1}{\gamma} \frac{\partial^2}{\partial \gamma^2} \gamma - \frac{l^2}{\hbar^2 \gamma^2}$$
(6)

916 The formula (6) is replaced into (5) and the equation (5) is:

917 
$$\left[-\frac{\hbar^2}{2\mu}\frac{1}{\gamma}\frac{\partial^2}{\partial\gamma^2}\gamma + \frac{l^2}{2\mu\gamma^2} + V(\gamma)\right]\psi = \frac{E}{C_M}\psi$$
(7)

From the separation of the variables  $\psi(\gamma, \theta, \psi) = R_l Y_{lm}(\theta, \psi)$ , the radial equations are available

920 
$$\left[\frac{1}{\gamma}\frac{d^2}{d\gamma^2}\gamma + \frac{2\mu}{\hbar^2}\left(\frac{E}{C_M} - \nabla\left(\gamma\right) - \frac{l(l+1)}{\gamma^2}\right)\right]R_l = 0 \tag{8}$$

Divided the singularity  $\gamma = 0$ , let  $R_l(\gamma) = \frac{X_l(\gamma)}{\gamma}$ , and then substitute it into equation (8):

922 
$$X_l'' + \left[\frac{2\mu}{\hbar^2} \left(\frac{E}{C_M} - \nabla(\gamma)\right) - \frac{l(l+1)}{\gamma^2}\right] X_l = 0$$
(9)

923 Coulomb energy is substituted into equation (9)

924 
$$X_l'' + \left[\frac{2\mu}{\hbar^2}\left(\frac{E}{C_M} + \frac{e^2}{\gamma}\right) - \frac{l(l+1)}{\gamma^2}\right]X_l = 0$$

Considering  $\gamma \to 0, \gamma \to \infty$  and the boundary conditions, the physically allowed solutions can be get by the application of the confluence super geometric equation:

927 
$$E = E_n = -\frac{c_M \mu e^4}{2\hbar^2} \frac{1}{n^2} = -\frac{c_M e^2}{2a} \frac{1}{n^2} \quad n = 1, 2, 3, ... \quad (10)$$
928 
$$a = \frac{\hbar^2}{\mu e^2}$$

929 The energy difference 
$$\Delta E = h\nu = h\frac{\overline{\nu}}{c} = E_n - E_m = C_M \left[\frac{\mu e^4}{2\hbar^2} \left(\frac{1}{m^2} - \frac{1}{n^2}\right)\right], \quad (n > m)$$

930 
$$\bar{\nu} = C_M \left[ \frac{2\pi^2 \mu e^4}{\hbar^3 c} \left( \frac{1}{m^2} - \frac{1}{n^2} \right) \right] = C_M R \left( \frac{1}{m^2} - \frac{1}{n^2} \right)$$
(11)

931 
$$R = \frac{2\pi^2 \mu e^4}{\hbar^3 c}$$
为 Rydberg constant。

In the gravitational field quantum condition in the Earth region  $C_M=1$ ,  $\bar{\nu} = R\left(\frac{1}{m^2} - \frac{1}{n^2}\right)$  is converted into the Rydberg formula.

According to the spectral formula (11), when C<sub>M</sub> take different values, e.g., the various 934 external gravitational field strength, the hydrogen characteristic spectrum moves as an whole. 935 Unlike the energy level splitting of the Stark effectin in the external electric field and the Zeeman 936 937 effect in the external magnetic field, the atomic characteristic spectra in the external gravitational field have the effect to move with the strength of the gravitational field overall. Such as, the quasar 938 spectral redshift is this effect (after all, the big red shift of the mass matter almost the speed of light 939 is difficult to be widely accepted). The greater the field strength will lead to the greater the spectral 940 movement. 941

In the strong gravitational field region,  $\bar{\nu} = C_M R \left(\frac{1}{m^2} - \frac{1}{n^2}\right)$ ,  $C_M < 1$ , a red shift is produced. Therefore, it could explain the quasar big red shift, which is a spectral redshift caused by the external gravitational field (or the gravitational field in the background) changing the atomic structure. See the literature [18] for more details.

Referring to the literature [18], the presence of n stable quantum condition regions in the universe is defined by the redshift distribution of the number of quasars. It is only one of n stable quantum condition regions for Earth. The number of quasars reaches its peak at the redshift values of Z = 0.3, 0.6, 0.96, 1.41, 1.96,respectively.

If we define  $C_M = \frac{1}{Z+1}$ , the peak values are corresponding to  $C_M = 0.77, 0.63, 0.51, 0.41, 0.33$ .

That is, when the  $C_M$  difference is about 0.1, there is a stable quantum condition region, and the gravitational field quantum condition is also a gradient, non-continuous.

955 (3) The Dirac equation

956 The Dirac equation for free electrons

$$i\hbar\frac{\partial}{\partial t}\psi = H'\psi \qquad (12)$$

958 
$$H' = C_M H = C_M (-i\hbar c\alpha \cdot \nabla + mc^2 \beta) = C_M (c\alpha \cdot P + mc^2 \beta)$$
  
959 
$$\psi_{P,E}(r,t) = U(P) exp \left[ \frac{i(P \cdot r - Et)}{\hbar} \right]$$

Pipe  $n_1$  Replacing the upper 2 formula into equation (12), it gets:

$$C_M(c\alpha P + mc^2\beta)u = Eu$$
(13)

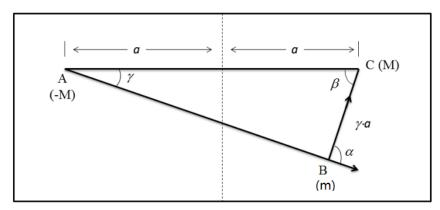
Then, it is obtained:  $E = \pm C_M \sqrt{m^2 c^4 + c^2 \hbar^2 K^2}$ , here  $P = \hbar K$ ,  $\alpha$  and  $\beta$  are the matrix formula.

It is learned that the energy "content" is still related to the background gravitational field strength. When the gravitational field increases the C<sub>M</sub> decreases and the particle energy decreases. **Appendix 3. The Einstein cosmological constant derived from the ''gravitational'' repulsion** between the positive-sky and anti-sky

According to CPT invariance, there exists a chiral symmetrically positive matter sky and anti-968 matter sky in the universe (it will explain in the section of the particle physics part that the matter 969 and anti-matter will form the positive and anti-matter sky, respectively, but not the matter and anti-970 mater mixed together to form the universe). It is the cooperation between the repulsion action of 971 the anti-sky to the positive sky galaxies and the gravity effect of the positive sky on its internal 972 galaxies that can accelerate the expansion of galaxies in the positive sky in the universe, manifested 973 as a dark energy phenomenon. The following calculation is the evolution of the positive sky under 974 975 anti-sky repulsion.

Under the cosmological assumption, the positive sky (M) and the anti-sky (-M) are applied 976 to the physical laws of FRW (Friedmann-Robertson-Walker), respectively. However, the 977 978 geometry influence cannot be clearly determined due to the large scale, the selection of coordinates and direction of action. Moreover, the positive sky metric tensor under anti-sky action cannot be 979 determined. Therefore, the 4-dimensional space-time tensor equation cannot be applied to the 980 connection between the positive sky and the anti-sky. In considering the role of the mass center in 981 the anti-sky on galaxies in the positive sky, it can be regarded as the role of the Newtonian potential 982 at t moments in a very large-scale space. It is applicable the positive sky evolution under the FRW 983 984 metric at 1-dimensional time and 3-dimensional Euclidean space with the same time horizon (FRW represents the positive matter sky, and -FRW represents the anti-matter sky in the followed 985 descriptions). 986

The universe consists of the chiral symmetric positive sky (FRW) and anti-sky(-FRW) and begins to expand in the two directions after explosion, forming the today's universe (see Fig. 3.1).



**Figure 3.1** a cosmic model formed by the explosion of the positive and anti-sky. A: anti-center of sky, M is total mass. B: the positive sky galaxy m position, *a* is the radius of the positive and antisky, *ra* is the co-moving distance of m from the positive center of mass,  $\alpha$  is the angle between the positive and anti-sky on m force. C: the positive center of mass, and M is total mass.

993

996

994 In Fig. 3.1, AC = 2a, BC = ra

995  $(AB)^2 = (2a)^2 + (ra)^2 - 2(2a)(ra)\cos\beta$ 

$$= 4a^2 + (ra)^2 - 4a^2r\cos\mu$$

<sup>997</sup> Here, we do not know the centroid of mass position of-FRW and FRW and angle  $(\alpha - \gamma)$  cannot <sup>998</sup> be measured.

When the galaxy is selected as the coordinate origin, and the spherical coordinates is taken, 999 the angle  $\beta = \alpha - \gamma$  is  $0 \sim 2\pi$  if  $\theta$  is  $0 \sim 2\pi$  and  $\varphi$  is fixed. Accordingly in this way, we can always let 1000 1001 the angle of  $\theta$  and  $(\alpha - \gamma)$  is roughly the same by selecting the coordinate system. If the angle difference is  $\theta_1$ ,  $\theta = (\alpha - \gamma) \cdot \theta_1$ ,  $\alpha \in [0, 2\pi]$ ,  $\theta \in [0, 2\pi]$ . Then, we can replace the  $(\alpha \cdot \gamma)$  value 1002 with  $\theta$ . Thus, when  $\theta$  value takes  $0 \sim 2\pi$ , ( $\alpha - \gamma$ ) also takes the corresponding values. Consequently, 1003 the cos values for  $\theta$  and  $(\alpha - \gamma)$  are always in the same region, only the order of the taken value 1004 1005 is different. Therefore, the instead  $(\alpha \cdot \gamma)$  with  $\theta$  does not affect the calculation of the cos value. Furthermore, the impact is even less for our astronomical observations due to the  $\Delta\theta$  values from 1006 1007 the two galaxies.

1008 Therefore, 
$$(AB)^2 \approx 4a^2 + (\gamma a)^2 - 4a^2\gamma\cos\theta$$

We can only use the Einstein tensor equations of four-dimensional space time in observing and computing cosmological evolution. But at a very large scales, the role of the-FRW sky on galaxies in the FRW sky can only use the Newtonian potential and the interaction on m is in case at an invariant mass (-M) center of mass:

1013 
$$\mathbf{F} = -\frac{G(-M)m}{(AB)^2} = \frac{GMm}{a^2(4+r^2-4r\cos\theta)}(\hat{e}_-)$$

- The acceleration of the Newtonian potential due to the interaction of anti-sky on galaxy m in
   FRW (Euclidean space)
- 1016  $\frac{d^2(AB)}{dt^2} = \frac{GM}{a^2(4+r^2-4r\cos\theta)}(\hat{e}_-)$

1017  $\hat{e}_{-}$ : it is the interaction direction of the anti- sky on the galaxy, which cannot be determined 1018 now, and it is only shows the difference from the later metric.

1019 In the positive sky, the spherical coordinate system and the FRW metric are taken, and the m 1020 point is defined as the coordinate origin:

1021 
$$ds^{2} = -c^{2}dt^{2} + a^{2}\left[\frac{dr^{2}}{1-Kr^{2}} + r^{2}(d\theta^{2} + \sin^{2}\theta \,d\varphi^{2})\right]$$
(1)

1022 Let 
$$(x^0, x^1, x^2, x^3)$$
 represents  $(ct, r, \theta, \varphi)$ , and thus this metric tensor component is:

1023 
$$g_{00} = -1, \ g_{11} = \frac{a^2}{1-Kr^2}, \ g_{22} = a^2r^2, \ g_{33} = a^2r^2\sin^2\theta, \ g_{\mu\nu} = 0, \ \forall \mu \neq \nu$$
  
1024 (2)

1024 (2)
 1025 Under the symmetry conditions, the FRW medium energy momentum tensor in the universe
 1026 takes the following form:

1027 
$$T_{\mu\nu} = \begin{pmatrix} \rho C^2 & 0 & 0 & 0 \\ 0 & g_{11}P & 0 & 0 \\ 0 & 0 & g_{22}P & 0 \\ 0 & 0 & 0 & g_{33}P \end{pmatrix}$$
(3)

1028 The connection formula by diagonalization  $\Gamma_{ij}^{K} = 0$ , when *i*, *j*, K is not equal to each other

1029 
$$\Gamma_{ii}^{K} = -\frac{1}{2} \frac{1}{g_{KK}} \frac{\partial g_{ii}}{\partial x^{K}}, \quad i \neq K, \quad 1 \le i$$

1030 
$$\Gamma_{Ki}^{K} = \Gamma_{iK}^{K} = \frac{1}{2} \frac{1}{g_{kk}} \frac{\partial g_{K\mu}}{\partial x^{i}}, \quad 1 \le i$$

1031 For (1), All 
$$\Gamma^{\alpha}_{\mu\nu}$$
 (not equal to zero) are listed as below:

1032 
$$\Gamma_{\rm KK}^0 = \frac{1}{2c} \frac{\partial g_{kk}}{\partial t}, \quad {\rm K=1}, 2, 3$$

1033 
$$\Gamma_{11}^1 = \frac{Kr}{1-Kr^2}, \ \Gamma_{22}^1 = -r(1-Kr^2)$$

1034 
$$\Gamma_{33}^{1} = -r \sin^{2} \theta \left(1 - Kr^{2}\right), \ \Gamma_{10}^{1} = \frac{1}{ca} \frac{\partial a}{\partial t}$$
 (4)

1035 
$$\Gamma_{20}^2 = \frac{1}{ca} \frac{\partial a}{\partial t}, \quad \Gamma_{21}^2 = \frac{1}{r}, \quad \Gamma_{33}^2 = -\sin\theta\cos\theta$$
  
1036 
$$\Gamma_{32}^3 = \frac{1}{ca} \frac{\partial a}{\partial t}, \quad \Gamma_{33}^3 = \frac{1}{ca}, \quad \Gamma_{35}^3 = \frac{\cos\theta}{ca}$$

1036 
$$\Gamma_{30} = \frac{1}{ca} \frac{1}{\partial t}, \ \Gamma_{31} = \frac{1}{r}, \ \Gamma_{32} = \frac{1}{\sin \theta}$$
  
1037 According to the formula of the Ricci tensor  $R_{\mu\nu}$ :

1038 
$$R_{\mu\nu} = \frac{\partial \Gamma^{\lambda}_{\mu\lambda}}{\partial x^{\nu}} - \frac{\partial \Gamma^{\lambda}_{\mu\nu}}{\partial x^{\lambda}} + \Gamma^{\alpha}_{\mu\lambda}\Gamma^{\lambda}_{\nu\alpha} - \Gamma^{\lambda}_{\mu\nu}\Gamma^{\alpha}_{\lambda\alpha}$$

1040 
$$R_{00} = \frac{1}{c} \frac{\partial}{\partial t} \Gamma_{K0}^{K} + \Gamma_{0\lambda}^{\alpha} \Gamma_{0\alpha}^{\lambda} = \frac{3}{c^{2}} \frac{1}{a} \frac{\partial^{2} a}{\partial t^{2}}$$

1041 
$$R_{11} = \frac{\partial \Gamma_{\lambda_1}^{\alpha}}{\partial r} - \frac{\partial \Gamma_{1_1}^{\alpha}}{\partial r} - \frac{\partial \Gamma_{1_1}^{\alpha}}{\partial t} - \Gamma_{\lambda_1}^{\alpha} \Gamma_{\alpha_1}^{\lambda} - \Gamma_{1_1}^{1} \Gamma_{\alpha_1}^{\alpha} - \Gamma_{1_1}^{0} \Gamma_{\alpha_0}^{\alpha}$$
1042 
$$= -\frac{1}{2\pi} \left[ a \frac{\partial^2 a}{\partial t} + 2 \left( \frac{\partial a}{\partial t} \right)^2 + 2Kc^2 \right]$$

$$= -\frac{1}{c^2(1-Kr^2)} \left[ a \frac{\partial t^2}{\partial t^2} + 2 \left( \frac{\partial t}{\partial t} \right)^2 + 2Kc^2 \right]$$

1043 
$$R_{22} = \frac{\partial \Gamma_{\lambda 2}}{\partial \theta} - \frac{\partial \Gamma_{22}}{\partial t} - \frac{\partial \Gamma_{22}}{\partial r} + \Gamma_{2\lambda}^{\alpha} \Gamma_{2\alpha}^{\lambda} - \Gamma_{22}^{\lambda} \Gamma_{\lambda\alpha}^{\alpha}$$
  
1044 
$$= -\frac{1}{c^2} \left[ a \frac{\partial^2 a}{\partial t^2} + 2 \left( \frac{\partial a}{\partial t} \right)^2 + 2Kc^2 \right]$$
(5)

1045 
$$R_{33} = -\frac{\partial \Gamma_{33}^{\lambda}}{\partial x^{\lambda}} + \Gamma_{3\lambda}^{\alpha} \Gamma_{3\alpha}^{\lambda} - \Gamma_{33}^{\lambda} \Gamma_{\lambda\alpha}^{\alpha}$$

1046 
$$= -\frac{r^{2}}{c^{2}}\sin^{2}\theta \left[a\frac{\partial^{2}a}{\partial t^{2}} + 2\left(\frac{\partial a}{\partial t}\right)^{2} + 2Kc^{2}\right]$$
  
1047 
$$R_{\mu\nu} = 0, \quad \forall \mu \neq \nu$$

1048 For 
$$T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T$$
, it has

1049 
$$T = g_{\mu\nu}T_{\mu\nu} = -\rho c^2 + 3P$$

1050 
$$T_{00} - \frac{1}{2}g_{00}T = \frac{c^2}{2}\left(\rho + \frac{3P}{c^2}\right)$$

1051 
$$T_{11} - \frac{1}{2}g_{11}T = \frac{1}{2}\frac{a^2c^2}{1-Kc^2}\left(\rho - \frac{P}{c^2}\right)$$

(6)

(8)

1052 
$$T_{22} - \frac{1}{2}g_{22}T = \frac{1}{2}r^2a^2c^2\left(\rho - \frac{P}{c^2}\right)$$

1053

1054

$$T_{33} - \frac{1}{2}g_{22}T = \frac{1}{2}\gamma^2 a^2 \sin^2 \theta \ c^2 \left(\rho - \frac{P}{c^2}\right)$$
$$T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T = 0. \ \forall \ \mu \neq \nu$$

1055 Here, the  $\Lambda$  equation for Lemaitre of the FRW metric is applied

1056 
$$R_{\mu\nu} = -\frac{8\pi G}{c^4} \left( T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T \right) + \Lambda g_{ij}$$
(7)

1057 Now,  $g_{ij} = -g_{\mu\nu}$ 

1058 Replace formula (5) and (6) into formula (7), then by equation  $R_{00}$ , it can derive:

1059 
$$\ddot{a} = -\frac{4\pi G}{3} \left( \rho + \frac{3P}{c^2} \right) a + \frac{\Lambda c^2}{3} a$$

1060 The acceleration of the Newtonian potential from the action of the cosmic anti-sky on the m 1061 galaxy in the positive sky

1062 
$$\frac{d_{AB}^2}{dt^2} = \frac{GM}{a^2(4+r^2-4r\cos\theta)}(\hat{e}_-)$$

1063 Comparing to formula (8) and the Friedmann universe model, there is an extra term  $\frac{\Lambda c^2}{3}a$ . It 1064 is the physical effect of anti-sky action on the positive sky galaxies. Thus, at the determined 1065 moment, the galaxies moving in FRW can be treated as the instantaneous cooperation of -FRW 1066 and FRW. Replacing the item containing  $\Lambda$  in formula (8) with  $\frac{d^2AB}{dt^2}$ , it will derive:

1067 
$$\ddot{a} = -\frac{4\pi G}{3} \left( \rho + \frac{3P}{c^2} \right) a + \frac{GM}{a^2 (4 + r^2 - 4r \cos \theta)} (\hat{e}_-)$$
(9)  
1068 Comparing formula (8) and (9)  
1069 
$$\Lambda = \frac{3GM}{a^3 c^2 (4 + r^2 - 4r \cos \theta)} = \frac{3GM}{a^3 c^2} \frac{1}{[4 \sin^2 \theta + (r - 2 \cos \theta)^2]}$$

1070 It obtains by equation 
$$R_{11}$$
:

1071 
$$a\ddot{a} + 2\dot{a}^2 = 4\pi G \left(\rho - \frac{P}{c^2}\right) a^2 - 2Kc^2 + \Lambda c^2 a^2$$

1072 Substituting the A value into the above formula:

1073 
$$a\ddot{a} + 2\dot{a}^2 = 4\pi G \left(\rho - \frac{P}{c^2}\right) a^2 - 2Kc^2 + \frac{3GM}{a[4\sin^2\theta + (r-2\cos\theta)^2]}$$
 (10)

1074 By equation (9) and (10), to eliminate the  $\ddot{a}$ :

1075 
$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8}{3}\pi G\rho + \frac{GM}{a^3(4+r^2-4r\cos\theta)}(\hat{e}_-) - \frac{Kc^2}{a^2}$$
 (11)

1076 By the divergence equation  $D^{\mu}T_{\mu\nu} = 0$ , it derives:

1077 
$$\frac{d\rho}{dt} + 3\left(\frac{\dot{a}}{a}\right)\left(\rho + \frac{P}{c^2}\right) = 0$$
(12)

Any two combination of the equation (9), (10), (11), and (12) could be called the improved Lemaitre equation, which is a cosmic model composed of the positive and anti-sky. By applying the Newtonian potential and Euclidean space on a very large scale, the evolution equation of the positive sky in Universe could be derived. This time, the FRW metric and the Lemaitre fourdimensional tensor equation in the positive matter sky should be adopted.

1083

### 1084 Appendix 4. Application in symmetry and Higgs Mechanism

1085 1. Vacuum symmetry breaking

1086 The universe is composed of the positive and anti-sky. Taking any point  $x (x_{\mu} = (x_0, x))$  in the 1087 positive sky, where the effective gravitational field at this space time point is the summing field of 1088 the positive and anti-sky gravitational field (the general relativity curved space time is not 1089 considered here, and it is stated in the part 2.2 that the equivalence principles are not suitable for 1090 quantum mechanics). It is traditionally represented as:

1091 
$$\varphi = \varphi_R + \varphi_L = -\frac{Km_R}{\gamma_R} - \frac{KM_L}{\gamma_L}$$
(1)

 $m_R$  is the positive sky effective material mass that produces the gravitational field at the fixed 1092 point of the positive sky while  $M_L$  is the total mass of the anti-sky. The difference is small for them. 1093 However, because the taken point is in the positive sky, thus,  $\gamma_R \ll \gamma_L$ , and  $\varphi_R \gg \varphi_L$ . As a result, 1094 the physical vacuum in the positive sky appears as a right-hand gravitational field. In order to adopt 1095 the quantum theory, the classical gravitational field can be phenomenally defined as a spinor field 1096  $\psi$  with a spin value of 1/2 according to the definition of some chiral gravitational field in sections 1097 1098 1 and 2 of this paper. In the quantum mechanics equation under gravitational field (2.6), the 1099 Newton scalar potential of the macroscopic gravitational field can be understood as a quantum hierarchical spinor field scalar  $\overline{\psi}\psi$ . The physical process observed and experiment test is always 1100 1101 occurred in the "vacuum" of the positive sky, which always exists a gravitational field, e.g., the summing field of the positive and anti-sky. 1102

1103

$$\varphi = k \overline{\psi}_R \psi_R + k \overline{\psi}_L \psi_L \tag{2}$$

1104 As mentioned above, the right-hand field is far larger than the left-hand field. Thus, the 1105 physical effect is manifested as the right-hand gravitational field vacuum, and this vacuum is invariant under any the parity transformation. It always represents as the right-hand gravitational 1106 field vacuum. In the (2) formula, if  $k\bar{\psi}_R\psi_R = -k\bar{\psi}_I\psi_I$  (here  $\varphi=0$ ), it is the symmetry point of the 1107 cosmic space in the positive and anti-sky, at which the symmetry of the vacuum remains constant. 1108 1109 It is  $k\bar{\psi}_R\psi_R \gg k\bar{\psi}_L\psi_L$  and  $\varphi = k\bar{\psi}_R\psi_R + k\bar{\psi}_L\psi_L \neq 0$  at an any point in the positive sky vacuum. Therefore, the vacuum gravitational field involved in the interaction is actually symmetry breaking, 1110 and this breaking is caused by the asymmetry of the positive and anti-sky gravity field at any point 1111 in the positive sky space, which is the physical interpretation of the "vacuum symmetry 1112 spontaneous breaking" of the real physical vacuum. The breaking increases if the gravitational 1113 field is enhanced at x point, that is, the strong difference between the right-hand gravitational field 1114 and the left-hand gravitational field increases (The variety of the left-hand gravitational field in 1115 1116 the positive sky is very small). Further, the parity transformation cannot change gravitational field 1117 of our experimental environment in the positive sky (our experiment or observation phenomenon is always in the positive sky gravitational field, and the positive sky gravitational field can be 1118 1119 mathematically transformed into the anti-sky gravitational field, which is so-called mathematical parity conservation). The parity breaking caused in this case is assigned to the intensity of the 1120 gravitational field. The difference of the parity breaking degree is found originated from the 1121 gravitation field different from the earth region. A more detailed study is the observable parity 1122 1123 breaking arising from the action difference between the right-handed gravitational field in the positive sky vacuum and the positive-antiparticle or the differences in spatial orientation, such as 1124 the  $C_0^{60}$  experiment. It is the topic of replacing the scalar field with the spinor field in gauge theory. 1125 2. Analysis of the quality generation using Higgs mechanism 1126

In the Higgs mechanism of the standard model in the particle physics, it is assumed there is a scalar field  $\Phi(\chi)$  in nature. It interacts with the gauge and fermionic fields by maintaining the localized gauge symmetry, and the scalar field strength is not the lowest energy state of the system by assuming a suitable scalar potential. However, the vacuum state of the system in the quantum theory is the lowest energy state of the system, causing the spontaneous breaking of vacuum symmetry. The vacuum state of the scalar field interacted with the gauge and fermionic fields gives

a mass to the gauge particles and fermions. However, it is only a guessed scalar field and the 1133 corresponding spontaneous breaking vacuum state under the requirements of gauge theory. On one 1134 hand, we are still not clear whether such a physical scalar field and its required spontaneous 1135 1136 breaking vacuum state exists. On the other hand, our understanding of the physical real vacuum is still vague, and we can contact only one physical real vacuum state if regardless of the varies of 1137 the vacuum gravitational field strength. Whereas, it had to be the different vacuum or vacuum 1138 states, which is required by different physical theories. In fact, we can create countless vacuum 1139 states mathematically. It is known that there must be a gravitational field in the physical vacuum. 1140 All physical experiments or phenomena are performed in a vacuum contained a gravitational field. 1141 The gravitational field in the vacuum, especially at a relatively low energy, has still not been 1142 introduced into the gauge theory. If the gravitational field is quantized and the physical vacuum is 1143 described on the basis of the quantum gravity field, the physical vacuum and its vacuum state 1144 contained the quantum gravitational field are very suitable for the gauge theory. Therefore, it 1145 determines the real physical field required by the scalar field and determines the physical reality 1146 of the real vacuum (the vacuum is a gravitational field medium). Then, the theoretical vacuum and 1147 the real field have been equated and finally the gravity field is introduced into the gauge theory, 1148 1149 promoting to build the gauge theory and the standard model of particle physics.

1150 As for the analysis in the section 1, the gravitational field may be represented as:

1151 
$$\varphi = \begin{pmatrix} \overline{\psi}_L \psi_L \\ \overline{\psi}_R \psi_R \end{pmatrix}$$
(3)

The formula (3) is the scalar field required by the standard model gauge theory, so that the 1152 gravitational field existed in the positive sky space (the gravitational summary field caused in 1153 positive and antisky) corresponds to the scalar field required by the gauge theory. Then the 1154 gravitational field vacuum state in the positive sky corresponds to the scalar field vacuum state, 1155 and the scalar field and vacuum state in the gauge theory present the physical reality. In considering 1156 of the different requirements, such as for the convenient calculation and analysis for the Unitary 1157 gauge and the renormalizable non-unitary gauge contained Goldstone particles, the Scalar fields 1158 of unitary or other specifications will be used. For example, to match with the Weinberg-Salam-1159 Glashow weak electro-uniform gauge theory, the scalar field is: 1160

1161 
$$\Phi(x) = \begin{pmatrix} \phi^+(x) \\ \phi^0(x) \end{pmatrix} or \Phi(x) = \begin{pmatrix} \phi^+(x) \\ [v + H(x) + i\chi(x)] / \sqrt{2} \end{pmatrix}$$
(4)

1162 The potential energy term is

1163

 $V(\Phi) = -\mu^2 \Phi^+ \Phi + \lambda (\Phi^+ \Phi)^2$ (5)

The breaking of vacuum symmetry will spontaneously appear and it will lead the gauge 1164 particles to have mass and then Higgs particles will presence. All this will not be repeated here. 1165 This work will focus on two inseparable aspects of the same topic: one is to explain the dark energy 1166 phenomenon and calculate the cosmological constant with the interaction between macroscopic 1167 mass matter. On the other hand, the microscopic gravitational field is quantized and then the real 1168 1169 vacuum and vacuum states of physics are defined based on this gravitational quantization and mass generation, to derive parity breaking and confirm the scalar field vacuum in the Higgs mechanism 1170 of the standard model without further Feynman rules and cross-section calculations. As we know, 1171 the interaction between macroscopic mass matter needs to recognize and define gravitational fields. 1172 1173 The gravitational field quantization, the generation and the stability of mass matter all need to analyze the microscopic particle mass generation and interaction information. In order to obtain 1174

the particle mass generation and interaction information, the gravitational field vacuum physics reality is not considering here. Then, the non-abelian group Su(2) gauge symmetry spontaneous breaking is used as an example and the U(1) gauge transformation of QED is directly analyzed.

1178 The real scalar field  $\varphi$  is transformed in 3 dimensions of the Su (2) group. The matrix 1179 generated from Su (2) denoted as  $L_j(j = 1,2,3)$ . Let  $A_\mu \equiv A_{j\mu}L_j$ , the covariant differential 1180 quotient  $[(\partial_\mu + igA_\mu)\varphi]_i = \partial_\mu \varphi_i - gA_{j\mu}\varepsilon_{ijk}\varphi_k$ , and the partial of  $\varphi$  was included in the 1181 Lagrangian density.

1182 
$$\mathcal{L} = \frac{1}{2} \Big( \partial_{\mu} \varphi_{i} - g A_{j\mu} \varepsilon_{ijk} \varphi_{k} \Big) \Big( \partial^{\mu} \varphi_{i} - g A_{l}^{\mu} \varepsilon_{ilm} \varphi_{m} \Big) - V(\varphi^{T} \cdot \varphi) \quad (6)$$

- 1183 Let SU (2) has symmetry spontaneous breaking, that is,  $V(\varphi^T \cdot \varphi)$  has non-zero minimal 1184 points, thus it is always  $\langle \varphi \rangle_0 = \begin{pmatrix} 0 \\ 0 \\ u \end{pmatrix}$  by rotating SU (2). Thus, the nonzero vacuum expectation
- 1185 appears only in the third component of  $\langle \varphi \rangle_0$ .

1186 The 3 D representation matrix of SU (2) is 
$$L_1 = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{bmatrix}, L_2 = \begin{bmatrix} 0 & 0 & i \\ 0 & 0 & 0 \\ -i & 0 & 0 \end{bmatrix}, \text{ and } L_3 = \begin{bmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \text{ respectively. Therefore, there is } L_1 \langle \varphi \rangle_0 = \begin{bmatrix} 0 \\ -i\nu \\ 0 \end{bmatrix} \neq 0, L_2 \langle \varphi \rangle_0 = \begin{bmatrix} i\nu \\ 0 \\ 0 \end{bmatrix} \neq 0,$$

1188  $L_3\langle \varphi \rangle_0 = 0.$ 1189 It suggests that the vacuum is no longer invariant under  $L_1$  and  $L_2$ , whereas it is still invariant

1190 under  $L_3$ . By parameterizing  $\varphi$ , it could obtain  $\varphi = \exp\left[-\frac{i}{2}(\zeta_1 L_1 + \zeta_2 L_2)\right] \begin{bmatrix} 0\\ 0\\ \nu + \eta \end{bmatrix} = \langle \varphi \rangle_0 + \int \zeta_2 1$ 

- 1191  $\begin{bmatrix} \zeta_2 \\ -\zeta_1 \\ \eta \end{bmatrix}$  + high-secondary term of the field.
- 1192  $\zeta_1$  and  $\zeta_2$  are two Goldstone bosons associated with the breaking generation  $L_1$  and  $L_2$ . 1193 Make the domain Su(2) specification transformation:  $\varphi(x) \rightarrow \varphi'(x) = U(x)\varphi(x) =$

1194 
$$exp\left\{\frac{i}{\nu}[\zeta_{1}(x)L_{1}+\zeta_{2}(x)L_{2}]\right\}\varphi(x) = \begin{pmatrix} 0\\0\\\nu+\eta \end{pmatrix}$$
  
1195  $A_{\mu i}(x)L_{i} \to A_{\mu i}^{'}(x)L_{i} = U(x)\left[A_{i i}(x)L_{i}-\frac{1}{2}\partial_{\mu}\right]U^{-1}$ 

1195 
$$A_{\mu j}(x)L_j \to A_{\mu j}(x)L_j = U(x) \left[A_{ij}(x)L_j - \frac{1}{g}\partial_{\mu}\right] U^+(x)$$

1196 Substituting the upper binary formula into the equation (6), we obtained:

1197 
$$\mathcal{L}(\varphi, A_{\mu}) = \mathcal{L}(\varphi', A_{\mu}') = \frac{1}{2}(\partial_{\mu}\eta)(\partial^{\mu}\eta) + \frac{1}{2}g^{2}\nu^{2}(A_{1\mu}A_{1}^{'\mu} + A_{2\mu}A_{2}^{'\mu}) + \cdots$$

1198 It is noteworthy that in  $\mathcal{L}$  the scalar field  $\zeta_1 \,\, \zeta_2$  corresponding  $L_1$  and  $L_2$  vanishes. 1199 However, the gauge fields  $A'_{1\mu}$  and  $A'_{2\mu}$  corresponding  $L_1$ ,  $L_2$  gain the mass. Known by the 1200 infinitesimal transform  $U(x) = 1 + \frac{i}{2}(\zeta_1 L_1 + \zeta_2 L_2)$ ,

1201 
$$A'_{j\mu}(x) = A_{j\mu}(x) - \frac{1}{g\nu} \left( \partial_{\mu} \zeta_j(x) + o(\zeta^2, \zeta A) \right), \quad j=1,2$$
 (7)

1202 Thus, the vanishing scalar fields  $\zeta_1$  and  $\zeta_2$  are absorbed into the longitudinal components of 1203 the gauge fields  $A'_{1\mu}$  and  $A'_{2\mu}$ , respectively, making them massive fields. For the gauge transformation of u(1) in QED, in the invariant L-quantity after gauge transformation,  $A'_{\mu}(x) = A_{\mu}(x) - \frac{1}{ev} \partial_{\mu} \zeta(x)_{\circ}$  The vanishing scalar field  $\zeta$  enters the longitudinal component of the  $A'_{\mu}$ , making the massless vector field  $A_{\mu}$  a massive vector field.

In terms of the mass production of the Higgs mechanism of the above Su(2), u(1) gauge fields, 1207 the scalar field is necessarily required both in the electromagnetic gauge theory or the standard 1208 model gauge theory (containing U(1) gauge). Moreover, the gauge particles could obtain mass 1209 only when the vanishing scalar fields such as  $\zeta_1$  and  $\zeta_2$  in SU(2), become longitudinal components 1210 of the gauge field. In the previous discussion, it is clear that the scalar field required by the standard 1211 1212 model gauge theory of the particle physics is the scalar of the gravitational spinor field in the real physical vacuum. The mathematical formula of the scalar field introduced to the theory is 1213 represented as formula (4) and its potential energy terms is showed as formula (5), which gives 1214 the gauge particle mass during the symmetry breaking. Therefore, the generation of mass can be 1215 1216 understood as: although the gravitational field is weak, the longitudinal polarized gravitational field is one of the conditions for the generation of mass matter (that is "mass matter"). The other 1217 1218 is the transverse condition and it is the transverse polarized electromagnetic field for electrons. The two conditions are indispensable, neither deficiency can constitute a stable mass matter 1219 (particles). Furthermore, it is concluded that the mass material generated by the energy levels must 1220 have transverse electromagnetic field conditions and longitudinal gravitational field conditions. 1221 Consequently, the stronger the gravitational field in the positive sky is, the more stable the particle 1222 is, and the more unstable the anti-particle is. Obviously, this particle stability conclusion is related 1223 to the Yukawa coupling of the fermions and scalar field ( $\mathcal{L}_{Yukawa} = -g_i^L \bar{l}_{iR} \bar{\Phi}^+ l_L^i + g_{ij}^d \bar{d}'_{ik} \Phi^+ q_L^j - g_i^u \bar{u}_{iR} \tilde{\Phi}^+ q_L^i + h.c$ ).Due to the positive and antiparticle spin in the coupling term 1224 1225 between the scalar field (gravitational field) in the Lagrange quantity and the fermion field, the 1226 1227 different positive and negative masses make different contributions to the Lagrange quantity values. Thus, the positive and antiparticle show the different lifetimes. Consequently, a positive 1228 1229 matter is stable and antimatter is unstable in the right-handed gravitational field. And the stronger 1230 the field strength, the greater the difference between the Lagrange quantity value of the positive 1231 and antiparticles is. That is, the stronger the right-hand field strength, the more stable the positive matter is, and the more unstable the antimatter is. Therefore, the right-handed gravitational field is 1232 a stable field for the positive matter and the unstable field for the antimatter. It derives there existed 1233 a microscopic chiral symmetric particle world and also a macroscopic chiral-symmetric universe 1234 1235 via combining the CPT theorem and the limit case of the Rydberg formula in the gravitational field condition. The universe consists of the positive and antimatter sky, in which the positive and 1236 1237 antimatter cannot mix together to form the universe. We do analysis by above approaches but not to calculate particle lifetime through the Green's function, because it is a real new physics by 1238 introducing the gravitational field scalar into gauge theory. It is involving into the gravitational 1239 field (scalar field) and gauge field, fermion field, Faddeev-Popov ghost field. The issues to select 1240 the scalar gravitational field or spinor is still existence, and it is the problem we will continue to 1241 1242 solve in future.

Based on the analysis of the Higgs mechanism and the energy solution of the quantization equation of section 3.2, we can conclude that the mass matter generated by energy matter must have transverse electromagnetic field conditions and longitudinal gravitational field conditions. The stronger region of the positive sky gravitational field is, the positive matter will be more stable in this region. Meanwhile, the antimatter is more unstable. Similarly, in the strong region of the anti-sky gravitational field, the antimatter is more stable while the positive matter is more unstable. 1249 So, the positive and antimatter cannot be intermingled, and they can only form a chiral symmetric 1250 positive-antimaterter sky.

In terms of the micromatter structure, the synthetic particles should have two conditions: one 1251 1252 is the transverse energy and the other is the longitudinal gravitational field. Under a fixed gravity field condition of  $C_M=1$  in the earth area, we cannot synthesize the big massive particles even if 1253 we promote the energy conditions. If you need to synthesize the big massive particles, it is 1254 necessary to reduce C<sub>M</sub>, that is, to improve the gravitational field strength. Therefore, to synthesize 1255 the big massive particles, you have to find a larger field strength zone (by the way, I therefore don't 1256 advocate building a large high-energy collider on earth, it is firmed not get any experimental 1257 results). With the enhancement of the gravitational field, it performs in macroscopic as a 1258 spectroscopy red-shifted of the atomic structure, the atomic structure of the white dwarf, and the 1259 dense material structure of the neutron stars. In fact, these macroscopic phenomena themselves are 1260 proof that the external gravitational field will change the matter structure. 1261