

# Deepening of Southern Ocean gateway leads to abrupt onset of a deep-reaching meridional overturning circulation

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## Abstract

During the Eocene-Oligocene transition, the meridional overturning circulation underwent large changes, associated with the geological evolution of Southern Ocean gateways. These are crucial for the Cenozoic climate transition from Greenhouse to Icehouse, but their dynamics still remain elusive. We demonstrate, using an idealised eddy ocean model, that the opening of a gateway leads to an abrupt onset of a vigorous, deep-reaching, meridional overturning circulation. This meridional overturning circulation has a maximum transport for a shallow gateway, and decreases with further deepening of the gateway. This abrupt change in the meridional overturning circulation can be explained through the ability with which standing meanders – turbulent features located downstream of the gateway – can induce deep vertical heat transport at high latitudes where bottom waters are produced. Our results demonstrate the crucial role of turbulent processes, associated with tectonic evolution, in setting the strength of the global ocean’s deep-reaching meridional overturning circulation.

1 **Deepening of Southern Ocean gateway leads to abrupt**  
2 **onset of a deep-reaching meridional overturning**  
3 **circulation**

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8 **Key Points:**

- 9 • The shallow opening of an ocean gateway leads to an abrupt onset of a deep-reaching  
10 overturning circulation.
- 11 • The deep-reaching overturning circulation is a consequence of standing meanders  
12 allowing for full-depth vertical heat transport.
- 13 • Further deepening of the gateway leads to a weaker overturning due to a decrease  
14 in heat transport towards southern convection regions.

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## Abstract

During the Eocene-Oligocene transition, the meridional overturning circulation underwent large changes, associated with the geological evolution of Southern Ocean gateways. These are crucial for the Cenozoic climate transition from Greenhouse to Icehouse, but their dynamics still remain elusive. We demonstrate, using an idealised eddying ocean model, that the opening of a gateway leads to an abrupt onset of a vigorous, deep-reaching, meridional overturning circulation. This meridional overturning circulation has a maximum transport for a shallow gateway, and decreases with further deepening of the gateway. This abrupt change in the meridional overturning circulation can be explained through the ability with which standing meanders – turbulent features located downstream of the gateway – can induce deep vertical heat transport at high latitudes where bottom waters are produced. Our results demonstrate the crucial role of turbulent processes, associated with tectonic evolution, in setting the strength of the global ocean’s deep-reaching meridional overturning circulation.

## Plain Language Summary

Around 50-34 million years ago, the Southern Hemisphere witnessed a major re-organisation of continents. This led to the opening and deepening of two Southern Ocean gateways - the Tasmanian Gateway between Australia and Antarctica, and Drake Passage between Cape Horn and the Antarctica Peninsula. During this period Earth’s climate went through a major climate transition, from a hot Greenhouse world to a cold Icehouse world. One hypothesis to explain this dramatic climate transition is that the opening of these ocean gateways led to a major transition in the ocean’s overturning circulation (i.e. its vertical circulation) with important consequences for the ocean’s capability to store heat and carbon. In this study we use an ocean model to understand how the opening of an ocean gateway affects the ocean’s overturning circulation. We show that it is small-scale processes, and their ability to transport heat southward and downward, which lead to an abrupt onset of the ocean’s overturning circulation as soon as the ocean gateway opens. Further deepening of the ocean gateways then leads to a decrease in the overturning circulation. This study therefore highlights the crucial role of small-scale processes in changing Earth’s climate.

## 1 Introduction

The long-term investigation into the meridional overturning circulation (MOC) is primarily motivated by its prominent role in the global redistribution of heat and cycling of chemical elements (Kuhlbrodt et al., 2007; Talley, 2013; Cessi, 2019). In the widely accepted paradigm, the modern MOC is composed of an upper cell, associated with the formation of North Atlantic Deep Water in the Nordic Seas, and a lower overturning cell, associated with the formation of Antarctic Bottom Water around the Antarctic coast (G. C. Johnson, 2008; Marshall & Speer, 2012; H. L. Johnson et al., 2019). The initiation of the modern MOC can be traced to the Eocene-Oligocene transition (EOT, ~34 Ma) (Hohbein et al., 2012; Thomas et al., 2014; Boyle et al., 2017; Ferreira et al., 2018; Hutchinson et al., 2021). Before this time, throughout the Paleocene and early Eocene (~65–40 Ma), geological evidence points towards a bipolar mode of the MOC in the Pacific basin, with convection occurring both in the North Pacific and the Southern Ocean (Thomas, 2004; Hague et al., 2012; Thomas et al., 2014).

The transition of the Eocene MOC to a modern-like MOC is thought to have occurred around the EOT (Katz et al., 2011; Borrelli et al., 2014; Hutchinson et al., 2021), triggered by tectonic changes. These tectonic changes involved the opening and deepening of the ocean gateways, such as Greenland-Scotland Ridge and Southern Ocean gateways (Katz et al., 2011; Abelson & Erez, 2017; Borrelli et al., 2014; Hutchinson et al., 2021). The deepening of Drake Passage (DP), for example, is thought to have led to deep

65 water upwelling in the Southern Ocean driven by wind stress, closing the modern type  
 66 MOC (Toggweiler & Bjornsson, 2000; Sijp & England, 2004). While many studies fo-  
 67 cus on understanding changes in the MOC between fully open and fully closed gateways,  
 68 recent work has shown how a progressive deepening of the DP affects the global MOC  
 69 (Toumoulin et al., 2020). Before the opening of the DP, southern sinking occurs in the  
 70 Atlantic basin and is constrained to shallow depth (Toumoulin et al., 2020). After the  
 71 opening of the DP to 100 m depth, southern sinking partially shifts to the Pacific, and  
 72 the global MOC abruptly strengthens to almost twice modern values, but weakens for  
 73 a further gateway deepening (Toumoulin et al., 2020). These simulations, showing the  
 74 effect of the progressive deepening of the DP on the MOC, have been run with a com-  
 75 plex Earth System model with a coarse-resolution ocean, and hence turbulent processes,  
 76 such as mesoscale eddies, need to be parameterised. As such, it is difficult to pin down  
 77 the exact ocean dynamics leading to this radical change in the MOC.

78 The deepening of ocean gateways forms large local bathymetric features, such as  
 79 ridges and seamounts. These bathymetric features have profound effects on the dynam-  
 80 ics of the Antarctic Circumpolar Current (ACC), generating Rossby waves due to jets  
 81 in the ACC interacting with these bathymetric features. Rossby waves propagate west-  
 82 ward against the ACC mean flow (C. W. Hughes, 2005; Zhang et al., 2023). When their  
 83 propagation speed matches the speed of the eastward-flowing ACC, these Rossby waves  
 84 become standing Rossby waves, also known as standing meanders (C. W. Hughes, 2005).  
 85 Mathematically, standing meanders can be presented as time-mean deviations from the  
 86 zonal-mean component of the ACC (Ivchenko et al., 1996; Youngs et al., 2017). Stand-  
 87 ing meanders in the modern ACC influence heat transport and closure of the overturn-  
 88 ing circulation (Youngs et al., 2017). For example, hot spots of eddy heat flux along the  
 89 ACC occur around major bathymetric features where standing meanders occur. This in-  
 90 dicates that they may contribute to strong poleward heat transport across the ACC (Foppert  
 91 et al., 2017). However, the role of standing meanders during gateway deepening, such  
 92 as that during the EOT, is uncertain since coarse-resolution ocean models used in these  
 93 studies do not resolve these turbulent processes, and it is unclear how well parameter-  
 94 isations of these processes work.

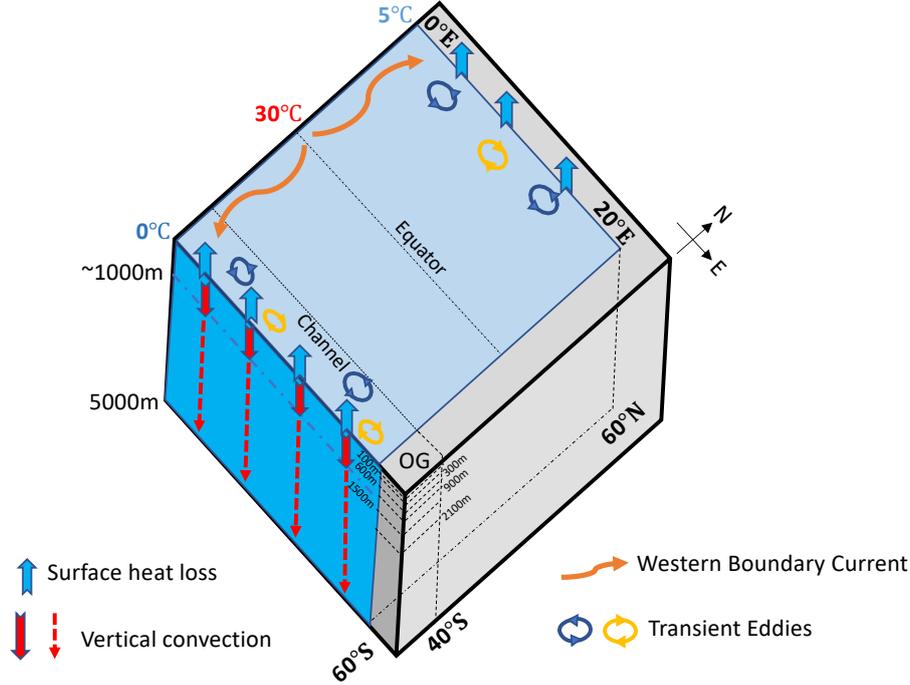
95 In this work, we use an idealised sector model to understand the role of gateway  
 96 deepening on the lower overturning cell. The model domain is that of a narrow sector  
 97 with a blocked ocean basin, such as the Atlantic basin, to the north, and a zonally re-  
 98 entrant channel, such as the Southern Ocean, to the south. Due to the limited size of  
 99 the domain, it is possible to use a horizontal resolution which allows for the represen-  
 100 tation of mesoscale eddies, whilst remaining computationally efficient. This allows for  
 101 the long spinup necessary to achieve statistical equilibrium. This work builds on results  
 102 of Klocker et al. (submitted), who used the same model configuration, but with only a  
 103 fully closed and fully open OG, to understand how buoyancy forcing alone can gener-  
 104 ate a deep-reaching lower overturning cell, similar to that observed. We use this config-  
 105 uration to analyse changes in ocean dynamics when we introduce different gateway depths  
 106 to the model geometry and seek to answer the following two questions based on Toumoulin  
 107 et al.’s simulation – why does the initial open of the DP generate a vigorous MOC, and  
 108 why does the simulated global MOC weaken with the deepening of the DP?

## 109 2 Methods

### 110 2.1 Model and simulations

111 The model configuration is based on an ocean-only sector model domain using the  
 112 MIT general circulation model (MITgcm) (Marshall, Hill, et al., 1997; Marshall, Adcroft,  
 113 et al., 1997). As shown in Figure 1, the model domain is composed of an ocean basin  
 114 extending from 60°S to 60°N, and 0°E to 20°E. An ocean gateway is located in the lat-  
 115 itudes of 60°S to 40°S. We refer to the ocean gateway as “OG” for the sake of conve-

116 nience. The model domain has  $1/6^\circ$  horizontal grid spacing, allowing for mesoscale ed-  
 117 dies. There are 42 unevenly spaced vertical levels with a 10 m thickness at the surface  
 118 level, stretching to 250 m at the bottom, for a total depth of 5000 m. At the surface, tem-  
 119 perature is restored to a fixed distribution, with a restoring time scale of 10 days. The  
 120 sea surface temperature at the equator is restored to  $30^\circ\text{C}$ , the southern end of the do-  
 121 main to  $0^\circ\text{C}$ , and the northern end of the domain to  $5^\circ\text{C}$ . A linear equation of state is  
 122 used for density with a constant salinity of 35 psu. There is no wind forcing or other me-  
 123 chanical forcing. The background vertical diffusivity is set to  $K_v = 10^{-6} \text{ m}^2\text{s}^{-1}$ . More  
 details of the model configuration can be found in Klocker et al.(submitted).



**Figure 1.** Schematic of model domain and ocean circulation. Orange curved arrows indicate western boundary current. In the channel ( $60^\circ\text{S} - 40^\circ\text{S}$ ), colored circles are transient eddies, blue vertical arrows are surface heat loss. Solid (dashed) red vertical arrows at the southern end indicate convection in the cases of *OG\_closed* (all experiments with an open OG).

124

125 An experiment with both a fully closed OG (*OG\_closed*), and a fully open OG (*OG\_open*),  
 126 were run for 3000 model years in Klocker et al.(submitted). Based on these two exist-  
 127 ing cases, we conduct six sensitivity experiments in which we change the OG (the  
 128 depth of the OG topography beneath sea surface) to 100 m (*OG\_100*), 300 m (*OG\_300*),  
 129 600 m (*OG\_600*), 900 m (*OG\_900*), 1500 m (*OG\_1500*) and 2100 m (*OG\_2100*). We run  
 130 all six sensitivity experiments for 2500 model years from the *OG\_closed* case, and use the  
 131 final 50 years for our analysis.

132

## 2.2 Calculation of the meridional overturning circulation

133

To accurately describe the MOC, composed of horizontal and vertical flows along  
 134 and across density surfaces, it is necessary to calculate the MOC in density coordinates

135 (Döös & Webb, 1994). We calculate the meridional transport ( $VH$ ) between every den-  
 136 sity layer, where  $V$  is meridional velocity and  $H$  is thickness of density layer. We take  
 137 the vertical integral of  $VH$ , then zonally integral and finally average over time to get the  
 138 total MOC in density coordinates (Ballarotta et al., 2013), which can be written as:

$$\Psi_{total}^{\sigma} = \frac{1}{t_1 - t_0} \int_{t_0}^{t_1} \int_{x_W}^{x_E} \sum_{z_{\sigma}^s(\sigma)}^{z_{\sigma}^b(\sigma)} VH dx dt, \quad (1)$$

139 where  $x_W$  and  $x_E$  are the longitudes of the western and eastern boundaries of the model  
 140 domain.  $z_{\sigma}^s$  and  $z_{\sigma}^b$  are surface and bottom density layer.  $t_0$  and  $t_1$  are the start and end  
 141 point of the time average, respectively. Noting that the vertical sum is vertically accu-  
 142 mulating density-binned  $VH$ .

143 We then use the thickness of density layers ( $z_{\sigma}(\sigma)$ ) and the time-average density  
 144 distribution in depth coordinates ( $\sigma(x, y, z)$ ) to linearly interpolate the density-coordinate  
 145 MOC into depth coordinates:

$$\Psi_{total}^{\sigma} \longrightarrow \Psi_{total}^z. \quad (2)$$

### 146 **2.3 Calculation of meridional heat transport and vertical heat transport**

147 The zonally integrated meridional heat transport ( $MHT$ ) is calculated on each depth  
 148 level as

$$MHT = \frac{\rho_0 C_{\rho}}{t_1 - t_0} \int_{t_0}^{t_1} \int_{x_W}^{x_E} VT dx, \quad (3)$$

149 where  $V$  is zonal velocity,  $T$  is potential temperature,  $\rho_0$  is the reference density, and  $C_{\rho}$   
 150 is the specific heat capacity. When zonally and vertically integrating the meridional heat  
 151 transport, we get

$$MHT_Z = \frac{\rho_0 C_{\rho}}{t_1 - t_0} \int_{t_0}^{t_1} \int_{x_W}^{x_E} \int_{z_b}^{z_s} VT dz dx dt, \quad (4)$$

152 where  $z^s$  and  $z^b$  are surface and bottom depth layer.

153 The vertical heat transport ( $VHT$ ) is calculated by horizontally integrating the ver-  
 154 tical advective flux of potential temperature at each depth level,

$$VHT = \frac{\rho_0 C_{\rho}}{t_1 - t_0} \int_{t_0}^{t_1} \int_A WT dA dt, \quad (5)$$

155 where  $W$  is the vertical velocity and  $A$  is horizontal area of the ocean domain.

### 156 **2.4 Decomposition into eddy, standing-meander, and mean components**

157 Using Reynolds averaging, an arbitrary field  $B$  can be decomposed into its time  
 158 mean,  $\overline{B}$ , and transient eddy,  $B'$ , components:

$$B = \overline{B} + B'. \quad (6)$$

159 Note that, by construction,  $\overline{B'} = 0$ .

160 The time mean component,  $\overline{B}$ , can be further separated into its zonal average part  
 161  $\langle \overline{B} \rangle$  and the excursion from this average  $\overline{B}^*$  (the so-called standing meanders compo-  
 162 nent):

$$\overline{B} = \langle \overline{B} \rangle + \overline{B}^*. \quad (7)$$

163 The total field,  $B$ , is then given by:

$$B = \langle \overline{B} \rangle + \overline{B}^* + B'. \quad (8)$$

164 We can use Reynolds averaging to decompose the total transport of any tracer, such  
 165 as temperature, into its mean, eddy, and standing meander components. The times and  
 166 zonal averages operate on the product of meridional velocity and temperature in the equa-  
 167 tion (3). This produces three terms, since the covariance of  $V'$  and  $T'$  is not zero by con-  
 168 struction. We can then decompose the total  $MHT_Z$  into three components:

$$MHT_Z = \rho_0 C_p \int_{x_W}^{x_E} \int_{z_b}^{z_s} \left[ \langle \bar{V} \rangle \langle \bar{T} \rangle + \overline{V^* T^*} + \overline{V' T'} \right] dz dx, \quad (9)$$

169 where the first term of right side is mean flow component ( $\langle \overline{MHT} \rangle_Z$ ), the second term  
 170 is the standing meander component ( $MHT_Z^*$ ), and the third term is transient eddy com-  
 171 ponent ( $MHT_Z'$ ). Similarly, we use this decomposition for the vertical heat transport,  
 172 giving  $\langle \overline{VHT} \rangle$ ,  $VHT^*$  and  $VHT'$ .

### 173 3 Results

#### 174 3.1 Overturning circulation with fully open and fully closed gateway

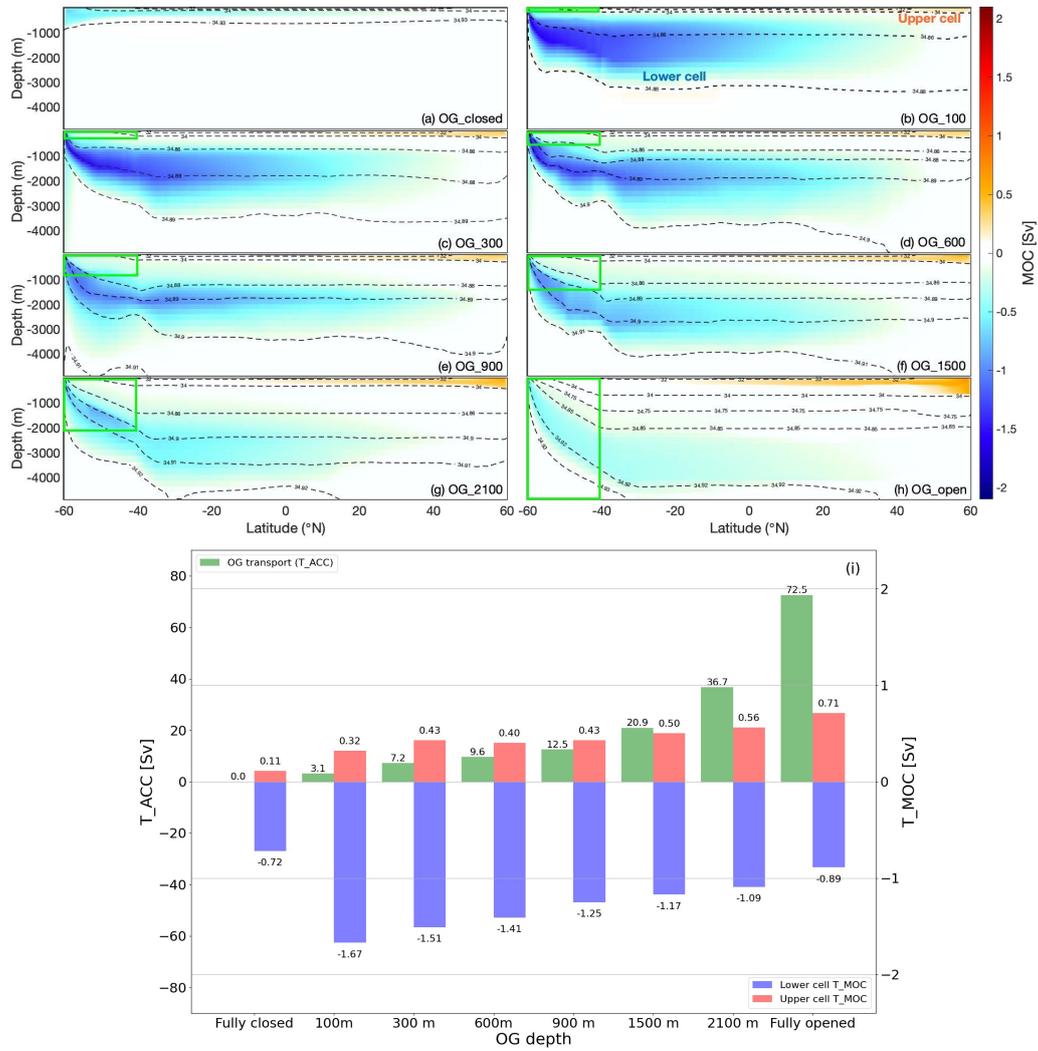
175 Due to the heating of the ocean surface at the equator and cooling at high latitudes,  
 176 ocean circulation adjusts to allow for a heat transport from the equator to high latitudes  
 177 through a process known as (rotating) horizontal convection (G. O. Hughes & Griffiths,  
 178 2008; Gayen & Griffiths, 2022). In a closed basin, such as in experiment *OG\_closed*, merid-  
 179 ional boundaries allow for the generation of an east-west pressure gradient, and hence  
 180 the poleward heat transport is due to western boundary currents (WBCs) associated with  
 181 ocean gyres (Figure S1b). These WBCs transport heat poleward very efficiently and lead  
 182 to a shallow MOC (Figure 2a), consistent with previous results (Klocker et al., submit-  
 183 ted).

184 In the absence of meridional boundaries, such as in experiment *OG\_open*, WBCs  
 185 cannot exist in the latitude range of the channel (Figure S1j). At these latitudes only  
 186 turbulent processes, such as mesoscale eddies, can transport heat across the channel (Fig-  
 187 ure 4i,j). These mesoscale eddies are generated by baroclinic instability due to steeply  
 188 sloping density surfaces across the channel. These surfaces also generate a circumpolar  
 189 current, resembling the Antarctic Circumpolar Current (Figure S1i), via thermal wind  
 190 balance. As shown by Klocker et al. (submitted), the presence of eddies allows for a deep-  
 191 reaching MOC, as opposed to the shallow MOC observed in the presence of meridional  
 192 boundaries. In experiment *OG\_open*, the surface buoyancy forcing generates a lower over-  
 193 turning cell with a strength ( $T_{MOC}$ ) of 0.89 Sv, which is equivalent to 16.02 Sv when ex-  
 194 trapolated from the width of the sector model to the full width of the ocean (Figure 2h,i).

195 For both *OG\_closed* and *OG\_open*, the MOC is closed by vertical plumes against  
 196 the northern and southern headwall of the domain (G. Hughes et al., 2007; Gayen & Grif-  
 197 fiths, 2022). In these regions, a destabilising buoyancy flux at the ocean surface leads  
 198 to the formation of deep and bottom waters (Gayen & Griffiths, 2022). The heat in both  
 199 cases can be transported by mean flows in the form of western boundary currents, or by  
 200 mesoscale eddies across the re-entrant channel to generate an overturning cell. Below we  
 201 will show that in the cases of partially open gateway there is a third process – standing  
 202 meanders – which can lead to lateral and vertical heat transport, allowing for the for-  
 203 mation of a deep-reaching MOC.

#### 204 3.2 Deep-reaching overturning circulation with gateway deepening

205 A shallow OG, as in experiment *OG\_100*, allows the formation of a vertically-sheared  
 206 circumpolar current (Figure 2i). This circumpolar current is the result of the combina-  
 207 tion of buoyancy loss at the southern boundary, which, together with the lack of merid-  
 208 ional boundaries above the OG depth, leads to the steepening of density surfaces and  
 209 hence a (weak) associated circumpolar current with a zonal transport ( $T_{ACC}$ ) of 3.1 Sv



**Figure 2.** (a)-(h) MOC [Sv] of all simulations; clockwise cell (red) is the upper overturning cell and counterclockwise cell (blue) is the lower overturning cell. (i) Transport through OG ( $T_{ACC}$  [Sv]; green), maximum MOC ( $T_{MOC}$  [Sv]) for the upper cell (red) and the lower cell (blue).

(Figure 2i). This circumpolar current is baroclinically unstable, resulting in the generation of transient eddies and standing meanders, with the latter being generated due to the circumpolar current interacting with the topography of the OG.

A striking change is the deep-reaching MOC in *OG\_100*; despite the very shallow OG, the lower overturning cell now extends to about 3500 m, and increases from 0.72 Sv in *OG\_closed* to 1.67 Sv (global scale: 30.06 Sv; Figure 2b,i). This formation of a strong deep-reaching lower overturning cell is consistent with results by Klocker et al. (submitted), who showed how mesoscale eddies allow for a full-depth overturning cell, and with results by Toumoulin et al. (2020) who used a coarse-resolution earth system model with a 100 m deep DP and Eocene boundary conditions to simulate changes in the MOC due to gateway deepening.

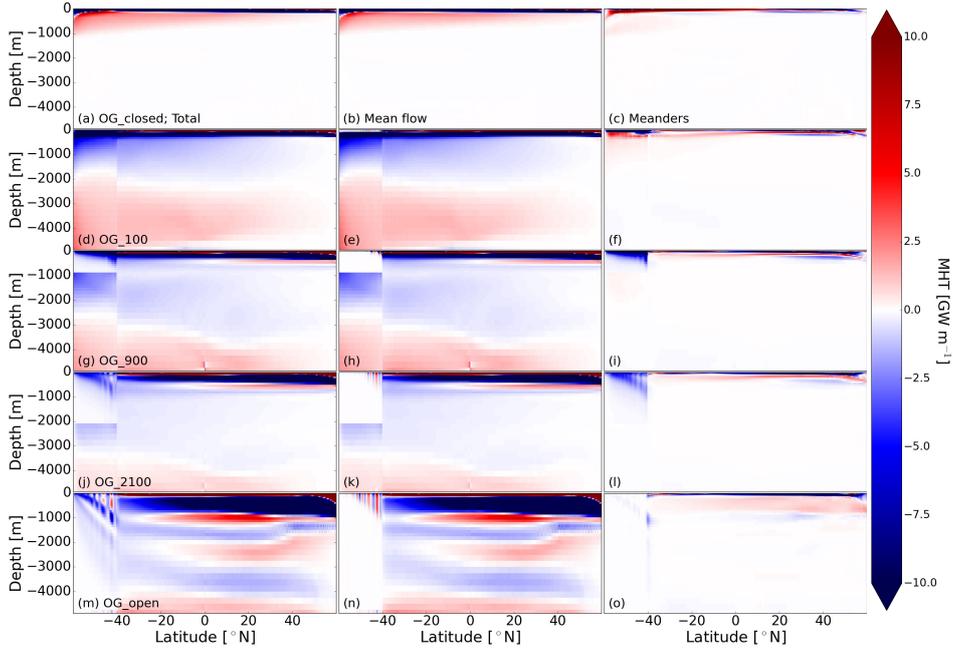
With the deepening of the OG from 100 m to 2100 m (experiments *OG\_100* to *OG\_2100*), the transport of the eastward circumpolar current ( $T_{ACC}$ ) gradually increases to 36.7 Sv (Figure 2i). In contrast with the increasing transport of the circumpolar current, the lower overturning cell continuously weakens to 1.09 Sv (global scale: 19.62 Sv) for the 2100 m OG experiment (Figure 2i). The shallowest OG therefore leads to the strongest lower overturning cell. This is consistent with recent results using a complex earth system model (Toumoulin et al., 2020).

### 3.3 The role of ocean heat transport

Buoyancy (heat) loss at the southern boundary of the channel leads to deep convection (Gayen & Griffiths, 2022) and hence the lower overturning cell. In an equilibrated ocean, maintaining the buoyancy loss and a deep-reaching overturning requires continuous lateral heat transport towards the southern ('Antarctic') boundary, and vertical heat transport to the deep ocean. The lower overturning cell is accordingly controlled by this heat transport. In this section, we will show how changes in OG depth affect the distribution of meridional and vertical heat transport by mean flows, standing meanders, and transient eddies.

For the experiment with a closed OG (*OG\_closed*), the zonally integrated meridional heat transport is dominated by the mean flow (Figure 3a,b), transporting heat southward (blue colours) at the surface and northward (red colours) below. This is consistent with the MOC being confined to the shallow ocean, with the deep ocean being largely at rest (Figure 2a and Figure 3a-c). In the southern hemisphere, standing meanders produce a southward MHT ( $MHT^*$ ) at the surface, and a northward MHT below. Transient eddies lead to a southward MHT across the channel, but contribute less to the heat transport than either the mean flow or standing meanders (Figure S4). The vertically integrated MHT ( $MHT_Z$ ) is strongly biased towards the southern high latitudes, consistent with experiments of thermally-driven flows in a domain lacking a re-entrant channel, in which the flow is biased towards the pole with the colder temperature forcing (Coman et al., 2010, Figure 4a). Consistent with the shallow confinement of the MHT and MOC, the VHT is confined to the shallow ocean (Figure 4b). In summary, in *OG\_closed* the entire heat transport is by baroclinic gyres which are confined to the ocean surface, leading to a shallow MOC. The deep ocean is at rest.

When the OG opens to a depth of 100 m (*OG\_100*), the MHT suddenly extends to the full depth of the basin rather than being confined to the shallow ocean, with a southward MHT from surface to mid-depth, and a northward MHT below (Figure 3e). This corresponds to a full-depth lower overturning cell, which increases in strength from 0.72Sv to 1.67Sv. Standing meanders lead to a southward MHT above the OG depth, and a northward MHT below the OG depth (Figure 3f). Transient eddies still provide weak southward MHT in the channel. The total vertically and zonally integrated southward MHT ( $MHT^*_Z$ ) in the southern half of the basin is reduced by about half (maximum about -10 TW) compared to *OG\_closed*, and is both due to the mean flow and standing me-

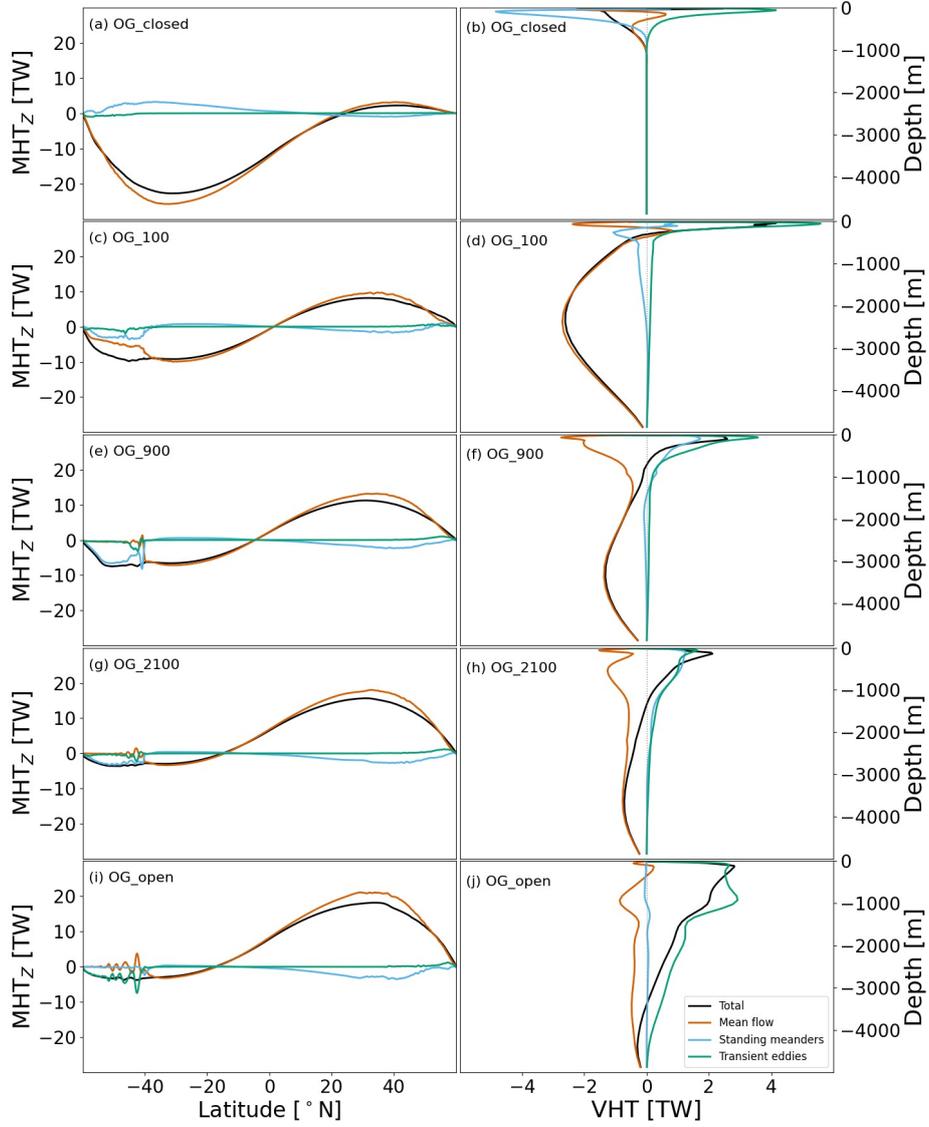


**Figure 3.** Hydrographic section of the zonally integrated meridional heat transport ( $MHT$ ) [ $\text{GW m}^{-1}$ ] for experiments  $OG\_closed$ ,  $OG\_100$ ,  $OG\_900$ ,  $OG\_2100$ , and  $OG\_open$ . Left panels show total  $MHT$ , middle panels show the  $MHT$  by the mean flow ( $\langle MHT \rangle$ ), and right panels show the  $MHT$  by standing meanders ( $MHT^*$ ).

261 anders (Figure 4c). On the other hand, the northward MHT in the northern half of the  
 262 basin increases by 8 TW compared to  $OG\_closed$ . Transient eddies again provide little  
 263 MHT ( $MHT'_z$ ) relative to the mean flow and standing meanders. In contrast to the large  
 264 decrease in MHT by the mean flow in  $OG\_100$ , the VHT by the mean flow shows a large  
 265 increase in strength and now extends to the full depth (Figure 4d).

266 For a further OG deepening (from  $OG\_100$  to  $OG\_2100$ ), the mean MHT in the  
 267 deep ocean decreases in strength compared with  $OG\_100$ , reflected by lighter blue and  
 268 red colors in the Figure 3h and k. This corresponds to a weakening of the deep-reaching  
 269 MOC. In the channel, standing meanders lead to a southward MHT above the OG depth.  
 270 The MHT by transient eddies increases, although it is still weaker than the other terms.  
 271 The gradual deepening of the OG strengthens the transport of the circumpolar current,  
 272 and leads to a state where the vertically integrated MHT is entirely due to standing me-  
 273 anders (Figure 4e and g) in the channel. The MHT by standing meanders becomes weaker  
 274 the deeper the OG, leading to a decrease in strength of the lower overturning cell for a  
 275 further OG deepening. This weakening of the southward MHT is accompanied by a strength-  
 276 ening of the northward MHT, and hence a stronger upper overturning cell.

277 When the OG is fully open ( $OG\_open$ ), standing meanders cannot exist since their  
 278 existence relies on the interaction with a ridge, and instead transient eddies provide south-  
 279 ward MHT. For this experiment the MOC is still full-depth, but weaker than in all cases  
 280 where standing meanders are present. This experiment therefore shows that eddies are  
 281 also capable of allowing for a full-depth lower overturning cell, but given that transient  
 282 eddies are less efficient at transporting heat towards the southern convection region, the  
 283 lower overturning cell is weaker than in all cases where standing meanders are present.



**Figure 4.** (a,c,e,g,i) Zonally and vertically integrated meridional heat transport ( $MHT_z$ ) [TW] for experiments  $OG\_closed$ ,  $OG\_100$ ,  $OG\_900$ ,  $OG\_2100$ , and  $OG\_open$ . (b,d,f,h,j) Horizontal integrated vertical heat transport ( $VHT$ ) [TW] in the channel for experiments  $OG\_closed$ ,  $OG\_100$ ,  $OG\_900$ ,  $OG\_2100$ , and  $OG\_open$ . In all panels, black lines indicate the total ( $MHT_z/VHT$ ), red lines the mean component ( $\langle \overline{MHT} \rangle_z / \langle \overline{VHT} \rangle$ ), green lines the standing meanders ( $MHT^*_z / VHT^*$ ), and blue lines the transient eddies ( $MHT'_z / VHT'$ ).

## 4 Discussion and Conclusion

In summary, the evolution of the deep-reaching lower overturning cell with the deepening of the ocean gateway can be divided into two parts: the abrupt onset of the deep-reaching lower overturning cell for a shallow ocean gateway, and the subsequent weakening for a further deepening of the ocean gateway. This behaviour of the lower overturning cell is closely linked to changes in southward and downward heat transport in the presence of standing meanders. First, in an equilibrated ocean, the strength of an overturning cell is limited by the amount of heat which can be supplied to the vertical plumes associated with strong surface buoyancy loss. This heat supply is most efficient if mean flows, such as western boundary currents, are present. This heat supply becomes less efficient if the heat transport is due to standing meanders, and even less efficient if heat supply is due to transient eddies. A deepening ocean gateway therefore leads to a reduced southward heat transport, which is compensated with an increase in northward heat transport. Second, a deep-reaching overturning cell is dependent on a process which can get this heat to depth. As shown by Klocker et al. (submitted), this can be achieved by transient eddies, and, as shown here, this can also be achieved by standing meanders. Nevertheless, heat transport by baroclinic gyres is always confined to the surface, and the presence of mesoscale turbulence, whether transient eddies or standing meanders, is necessary to get this heat to depth, and hence lead to a deep-reaching MOC.

The two-part evolution of the lower overturning cell for a deepening of an ocean gateway is therefore a combination of the two roles transient eddies and standing meanders play for ocean heat transport. These dynamics lead to a maximum lower overturning cell for the experiment with the shallowest OG, which is a combination of the domain where the southward heat transport is largest, while at the same time allowing for a deep-reaching heat transport by mesoscale turbulence. Further deepening of the OG leads to a weakening of the southward heat supply across the channel, and a strengthening of the northward heat transport, associated with a strengthening of the upper overturning cell. Once the OG is as deep as the rest of the domain, the southward heat transport is solely due to transient eddies, resulting in the weakest deep-reaching overturning cell. As opposed to studies which focus on just fully open and fully closed OGs, we highlight the importance of the dynamics of standing meanders in understanding the evolution of the MOC with gateway deepening, a turbulent process generally neglected in the paleoceanography community.

Previous studies have used complex, but coarse-resolution, earth system models to explain the evolution of the global MOC with the deepening of Southern Ocean gateways (Sijp & England, 2004; Toumoulin et al., 2020). However, in these studies results contradict each other, with (Sijp & England, 2004) showing the strongest MOC for a closed OG, while (Toumoulin et al., 2020), consistent with our study, showing the strongest MOC for a shallow OG. Nevertheless, both these studies show a weakening of the southward meridional heat transport with a deepening of the OG, and a strengthening of the northward meridional heat transport, consistent with our study. The different results for the lower overturning cell are possibly the consequence of the parameterisation of mesoscale turbulence in these models. It is the simple but high-resolution model configuration used in this study that helps to reveal the critical role of turbulence processes in generating the strong deep-reaching overturning cell with OG deepening. Future work is planned to extend this work on the evolution of the EOT MOC with a global ocean model and realistic paleo-bathymetry.

## 5 Data Availability Statement

All model output used in this study to perform the analysis and produce figures is available at <https://doi.org/10.5281/zenodo.7602996>.

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1 **Deepening of Southern Ocean gateway leads to abrupt**  
2 **onset of a deep-reaching meridional overturning**  
3 **circulation**

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8 **Key Points:**

- 9 • The shallow opening of an ocean gateway leads to an abrupt onset of a deep-reaching  
10 overturning circulation.
- 11 • The deep-reaching overturning circulation is a consequence of standing meanders  
12 allowing for full-depth vertical heat transport.
- 13 • Further deepening of the gateway leads to a weaker overturning due to a decrease  
14 in heat transport towards southern convection regions.

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**Abstract**

During the Eocene-Oligocene transition, the meridional overturning circulation underwent large changes, associated with the geological evolution of Southern Ocean gateways. These are crucial for the Cenozoic climate transition from Greenhouse to Icehouse, but their dynamics still remain elusive. We demonstrate, using an idealised eddying ocean model, that the opening of a gateway leads to an abrupt onset of a vigorous, deep-reaching, meridional overturning circulation. This meridional overturning circulation has a maximum transport for a shallow gateway, and decreases with further deepening of the gateway. This abrupt change in the meridional overturning circulation can be explained through the ability with which standing meanders – turbulent features located downstream of the gateway – can induce deep vertical heat transport at high latitudes where bottom waters are produced. Our results demonstrate the crucial role of turbulent processes, associated with tectonic evolution, in setting the strength of the global ocean’s deep-reaching meridional overturning circulation.

**Plain Language Summary**

Around 50-34 million years ago, the Southern Hemisphere witnessed a major re-organisation of continents. This led to the opening and deepening of two Southern Ocean gateways - the Tasmanian Gateway between Australia and Antarctica, and Drake Passage between Cape Horn and the Antarctica Peninsula. During this period Earth’s climate went through a major climate transition, from a hot Greenhouse world to a cold Icehouse world. One hypothesis to explain this dramatic climate transition is that the opening of these ocean gateways led to a major transition in the ocean’s overturning circulation (i.e. its vertical circulation) with important consequences for the ocean’s capability to store heat and carbon. In this study we use an ocean model to understand how the opening of an ocean gateway affects the ocean’s overturning circulation. We show that it is small-scale processes, and their ability to transport heat southward and downward, which lead to an abrupt onset of the ocean’s overturning circulation as soon as the ocean gateway opens. Further deepening of the ocean gateways then leads to a decrease in the overturning circulation. This study therefore highlights the crucial role of small-scale processes in changing Earth’s climate.

**1 Introduction**

The long-term investigation into the meridional overturning circulation (MOC) is primarily motivated by its prominent role in the global redistribution of heat and cycling of chemical elements (Kuhlbrodt et al., 2007; Talley, 2013; Cessi, 2019). In the widely accepted paradigm, the modern MOC is composed of an upper cell, associated with the formation of North Atlantic Deep Water in the Nordic Seas, and a lower overturning cell, associated with the formation of Antarctic Bottom Water around the Antarctic coast (G. C. Johnson, 2008; Marshall & Speer, 2012; H. L. Johnson et al., 2019). The initiation of the modern MOC can be traced to the Eocene-Oligocene transition (EOT, ~34 Ma) (Hohbein et al., 2012; Thomas et al., 2014; Boyle et al., 2017; Ferreira et al., 2018; Hutchinson et al., 2021). Before this time, throughout the Paleocene and early Eocene (~65–40 Ma), geological evidence points towards a bipolar mode of the MOC in the Pacific basin, with convection occurring both in the North Pacific and the Southern Ocean (Thomas, 2004; Hague et al., 2012; Thomas et al., 2014).

The transition of the Eocene MOC to a modern-like MOC is thought to have occurred around the EOT (Katz et al., 2011; Borrelli et al., 2014; Hutchinson et al., 2021), triggered by tectonic changes. These tectonic changes involved the opening and deepening of the ocean gateways, such as Greenland-Scotland Ridge and Southern Ocean gateways (Katz et al., 2011; Abelson & Erez, 2017; Borrelli et al., 2014; Hutchinson et al., 2021). The deepening of Drake Passage (DP), for example, is thought to have led to deep

65 water upwelling in the Southern Ocean driven by wind stress, closing the modern type  
 66 MOC (Toggweiler & Bjornsson, 2000; Sijp & England, 2004). While many studies fo-  
 67 cus on understanding changes in the MOC between fully open and fully closed gateways,  
 68 recent work has shown how a progressive deepening of the DP affects the global MOC  
 69 (Toumoulin et al., 2020). Before the opening of the DP, southern sinking occurs in the  
 70 Atlantic basin and is constrained to shallow depth (Toumoulin et al., 2020). After the  
 71 opening of the DP to 100 m depth, southern sinking partially shifts to the Pacific, and  
 72 the global MOC abruptly strengthens to almost twice modern values, but weakens for  
 73 a further gateway deepening (Toumoulin et al., 2020). These simulations, showing the  
 74 effect of the progressive deepening of the DP on the MOC, have been run with a com-  
 75 plex Earth System model with a coarse-resolution ocean, and hence turbulent processes,  
 76 such as mesoscale eddies, need to be parameterised. As such, it is difficult to pin down  
 77 the exact ocean dynamics leading to this radical change in the MOC.

78 The deepening of ocean gateways forms large local bathymetric features, such as  
 79 ridges and seamounts. These bathymetric features have profound effects on the dynam-  
 80 ics of the Antarctic Circumpolar Current (ACC), generating Rossby waves due to jets  
 81 in the ACC interacting with these bathymetric features. Rossby waves propagate west-  
 82 ward against the ACC mean flow (C. W. Hughes, 2005; Zhang et al., 2023). When their  
 83 propagation speed matches the speed of the eastward-flowing ACC, these Rossby waves  
 84 become standing Rossby waves, also known as standing meanders (C. W. Hughes, 2005).  
 85 Mathematically, standing meanders can be presented as time-mean deviations from the  
 86 zonal-mean component of the ACC (Ivchenko et al., 1996; Youngs et al., 2017). Stand-  
 87 ing meanders in the modern ACC influence heat transport and closure of the overturn-  
 88 ing circulation (Youngs et al., 2017). For example, hot spots of eddy heat flux along the  
 89 ACC occur around major bathymetric features where standing meanders occur. This in-  
 90 dicates that they may contribute to strong poleward heat transport across the ACC (Foppert  
 91 et al., 2017). However, the role of standing meanders during gateway deepening, such  
 92 as that during the EOT, is uncertain since coarse-resolution ocean models used in these  
 93 studies do not resolve these turbulent processes, and it is unclear how well paramet-  
 94 erisations of these processes work.

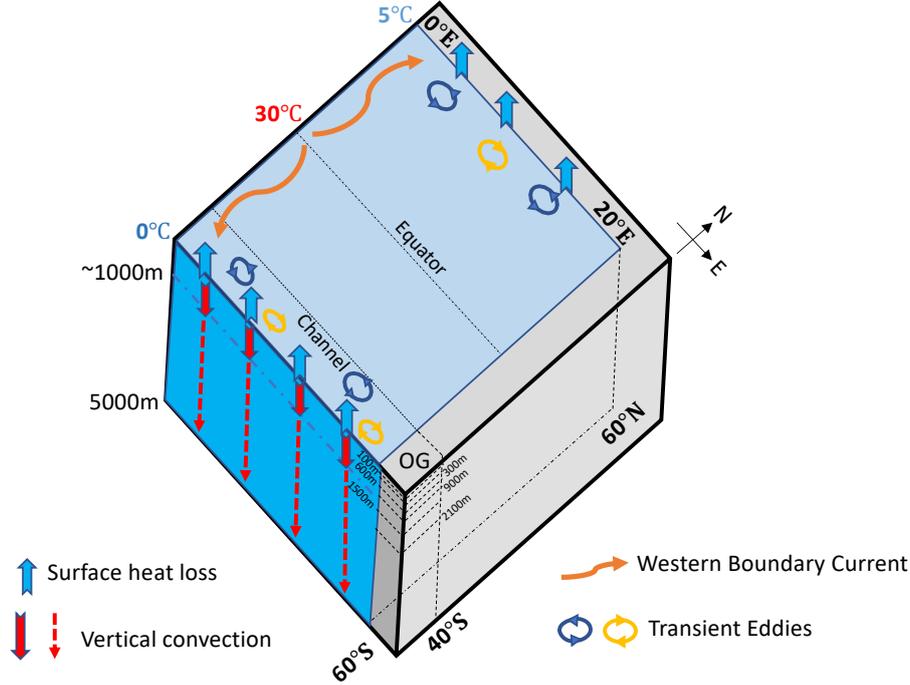
95 In this work, we use an idealised sector model to understand the role of gateway  
 96 deepening on the lower overturning cell. The model domain is that of a narrow sector  
 97 with a blocked ocean basin, such as the Atlantic basin, to the north, and a zonally re-  
 98 entrant channel, such as the Southern Ocean, to the south. Due to the limited size of  
 99 the domain, it is possible to use a horizontal resolution which allows for the represen-  
 100 tation of mesoscale eddies, whilst remaining computationally efficient. This allows for  
 101 the long spinup necessary to achieve statistical equilibrium. This work builds on results  
 102 of Klocker et al. (submitted), who used the same model configuration, but with only a  
 103 fully closed and fully open OG, to understand how buoyancy forcing alone can gener-  
 104 ate a deep-reaching lower overturning cell, similar to that observed. We use this config-  
 105 uration to analyse changes in ocean dynamics when we introduce different gateway depths  
 106 to the model geometry and seek to answer the following two questions based on Toumoulin  
 107 et al.’s simulation – why does the initial open of the DP generate a vigorous MOC, and  
 108 why does the simulated global MOC weaken with the deepening of the DP?

## 109 2 Methods

### 110 2.1 Model and simulations

111 The model configuration is based on an ocean-only sector model domain using the  
 112 MIT general circulation model (MITgcm) (Marshall, Hill, et al., 1997; Marshall, Adcroft,  
 113 et al., 1997). As shown in Figure 1, the model domain is composed of an ocean basin  
 114 extending from 60°S to 60°N, and 0°E to 20°E. An ocean gateway is located in the lat-  
 115 itudes of 60°S to 40°S. We refer to the ocean gateway as “OG” for the sake of conve-

116 nience. The model domain has  $1/6^\circ$  horizontal grid spacing, allowing for mesoscale ed-  
 117 dies. There are 42 unevenly spaced vertical levels with a 10 m thickness at the surface  
 118 level, stretching to 250 m at the bottom, for a total depth of 5000 m. At the surface, tem-  
 119 perature is restored to a fixed distribution, with a restoring time scale of 10 days. The  
 120 sea surface temperature at the equator is restored to  $30^\circ\text{C}$ , the southern end of the do-  
 121 main to  $0^\circ\text{C}$ , and the northern end of the domain to  $5^\circ\text{C}$ . A linear equation of state is  
 122 used for density with a constant salinity of 35 psu. There is no wind forcing or other me-  
 123 chanical forcing. The background vertical diffusivity is set to  $K_v = 10^{-6} \text{ m}^2\text{s}^{-1}$ . More  
 details of the model configuration can be found in Klocker et al.(submitted).



**Figure 1.** Schematic of model domain and ocean circulation. Orange curved arrows indicate western boundary current. In the channel ( $60^\circ\text{S} - 40^\circ\text{S}$ ), colored circles are transient eddies, blue vertical arrows are surface heat loss. Solid (dashed) red vertical arrows at the southern end indicate convection in the cases of *OG\_closed* (all experiments with an open OG).

124

125 An experiment with both a fully closed OG (*OG\_closed*), and a fully open OG (*OG\_open*),  
 126 were run for 3000 model years in Klocker et al.(submitted). Based on these two exist-  
 127 ing cases, we conduct six sensitivity experiments in which we change the OG (the  
 128 depth of the OG topography beneath sea surface) to 100 m (*OG\_100*), 300 m (*OG\_300*),  
 129 600 m (*OG\_600*), 900 m (*OG\_900*), 1500 m (*OG\_1500*) and 2100 m (*OG\_2100*). We run  
 130 all six sensitivity experiments for 2500 model years from the *OG\_closed* case, and use the  
 131 final 50 years for our analysis.

132

## 2.2 Calculation of the meridional overturning circulation

133

To accurately describe the MOC, composed of horizontal and vertical flows along  
 134 and across density surfaces, it is necessary to calculate the MOC in density coordinates

135 (Döös & Webb, 1994). We calculate the meridional transport ( $VH$ ) between every den-  
 136 sity layer, where  $V$  is meridional velocity and  $H$  is thickness of density layer. We take  
 137 the vertical integral of  $VH$ , then zonally integral and finally average over time to get the  
 138 total MOC in density coordinates (Ballarotta et al., 2013), which can be written as:

$$\Psi_{total}^{\sigma} = \frac{1}{t_1 - t_0} \int_{t_0}^{t_1} \int_{x_W}^{x_E} \sum_{z_{\sigma}^s(\sigma)}^{z_{\sigma}^b(\sigma)} VH dx dt, \quad (1)$$

139 where  $x_W$  and  $x_E$  are the longitudes of the western and eastern boundaries of the model  
 140 domain.  $z_{\sigma}^s$  and  $z_{\sigma}^b$  are surface and bottom density layer.  $t_0$  and  $t_1$  are the start and end  
 141 point of the time average, respectively. Noting that the vertical sum is vertically accu-  
 142 mulating density-binned  $VH$ .

143 We then use the thickness of density layers ( $z_{\sigma}(\sigma)$ ) and the time-average density  
 144 distribution in depth coordinates ( $\sigma(x, y, z)$ ) to linearly interpolate the density-coordinate  
 145 MOC into depth coordinates:

$$\Psi_{total}^{\sigma} \longrightarrow \Psi_{total}^z. \quad (2)$$

### 146 **2.3 Calculation of meridional heat transport and vertical heat transport**

147 The zonally integrated meridional heat transport ( $MHT$ ) is calculated on each depth  
 148 level as

$$MHT = \frac{\rho_0 C_{\rho}}{t_1 - t_0} \int_{t_0}^{t_1} \int_{x_W}^{x_E} VT dx, \quad (3)$$

149 where  $V$  is zonal velocity,  $T$  is potential temperature,  $\rho_0$  is the reference density, and  $C_{\rho}$   
 150 is the specific heat capacity. When zonally and vertically integrating the meridional heat  
 151 transport, we get

$$MHT_Z = \frac{\rho_0 C_{\rho}}{t_1 - t_0} \int_{t_0}^{t_1} \int_{x_W}^{x_E} \int_{z_b}^{z_s} VT dz dx dt, \quad (4)$$

152 where  $z^s$  and  $z^b$  are surface and bottom depth layer.

153 The vertical heat transport ( $VHT$ ) is calculated by horizontally integrating the ver-  
 154 tical advective flux of potential temperature at each depth level,

$$VHT = \frac{\rho_0 C_{\rho}}{t_1 - t_0} \int_{t_0}^{t_1} \int_A WT dA dt, \quad (5)$$

155 where  $W$  is the vertical velocity and  $A$  is horizontal area of the ocean domain.

### 156 **2.4 Decomposition into eddy, standing-meander, and mean components**

157 Using Reynolds averaging, an arbitrary field  $B$  can be decomposed into its time  
 158 mean,  $\overline{B}$ , and transient eddy,  $B'$ , components:

$$B = \overline{B} + B'. \quad (6)$$

159 Note that, by construction,  $\overline{B'} = 0$ .

160 The time mean component,  $\overline{B}$ , can be further separated into its zonal average part  
 161  $\langle \overline{B} \rangle$  and the excursion from this average  $\overline{B}^*$  (the so-called standing meanders compo-  
 162 nent):

$$\overline{B} = \langle \overline{B} \rangle + \overline{B}^*. \quad (7)$$

163 The total field,  $B$ , is then given by:

$$B = \langle \overline{B} \rangle + \overline{B}^* + B'. \quad (8)$$

164 We can use Reynolds averaging to decompose the total transport of any tracer, such  
 165 as temperature, into its mean, eddy, and standing meander components. The times and  
 166 zonal averages operate on the product of meridional velocity and temperature in the equa-  
 167 tion (3). This produces three terms, since the covariance of  $V'$  and  $T'$  is not zero by con-  
 168 struction. We can then decompose the total  $MHT_Z$  into three components:

$$MHT_Z = \rho_0 C_p \int_{x_W}^{x_E} \int_{z_b}^{z_s} \left[ \langle \bar{V} \rangle \langle \bar{T} \rangle + \overline{V^* T^*} + \overline{V' T'} \right] dz dx, \quad (9)$$

169 where the first term of right side is mean flow component ( $\langle \overline{MHT} \rangle_Z$ ), the second term  
 170 is the standing meander component ( $MHT_Z^*$ ), and the third term is transient eddy com-  
 171 ponent ( $MHT_Z'$ ). Similarly, we use this decomposition for the vertical heat transport,  
 172 giving  $\langle \overline{VHT} \rangle$ ,  $VHT^*$  and  $VHT'$ .

### 173 3 Results

#### 174 3.1 Overturning circulation with fully open and fully closed gateway

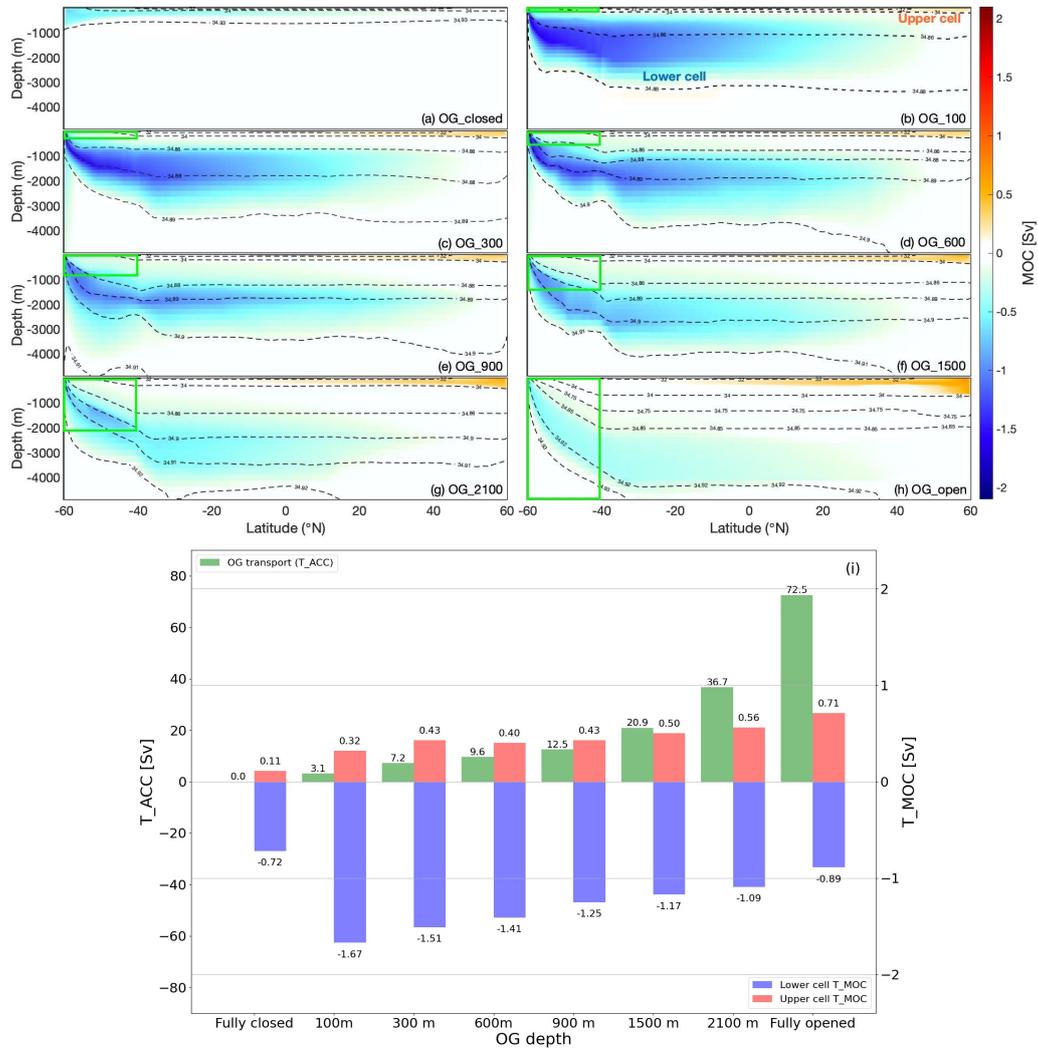
175 Due to the heating of the ocean surface at the equator and cooling at high latitudes,  
 176 ocean circulation adjusts to allow for a heat transport from the equator to high latitudes  
 177 through a process known as (rotating) horizontal convection (G. O. Hughes & Griffiths,  
 178 2008; Gayen & Griffiths, 2022). In a closed basin, such as in experiment *OG\_closed*, merid-  
 179 ional boundaries allow for the generation of an east-west pressure gradient, and hence  
 180 the poleward heat transport is due to western boundary currents (WBCs) associated with  
 181 ocean gyres (Figure S1b). These WBCs transport heat poleward very efficiently and lead  
 182 to a shallow MOC (Figure 2a), consistent with previous results (Klocker et al., submit-  
 183 ted).

184 In the absence of meridional boundaries, such as in experiment *OG\_open*, WBCs  
 185 cannot exist in the latitude range of the channel (Figure S1j). At these latitudes only  
 186 turbulent processes, such as mesoscale eddies, can transport heat across the channel (Fig-  
 187 ure 4i,j). These mesoscale eddies are generated by baroclinic instability due to steeply  
 188 sloping density surfaces across the channel. These surfaces also generate a circumpolar  
 189 current, resembling the Antarctic Circumpolar Current (Figure S1i), via thermal wind  
 190 balance. As shown by Klocker et al. (submitted), the presence of eddies allows for a deep-  
 191 reaching MOC, as opposed to the shallow MOC observed in the presence of meridional  
 192 boundaries. In experiment *OG\_open*, the surface buoyancy forcing generates a lower over-  
 193 turning cell with a strength ( $T_{MOC}$ ) of 0.89 Sv, which is equivalent to 16.02 Sv when ex-  
 194 trapolated from the width of the sector model to the full width of the ocean (Figure 2h,i).

195 For both *OG\_closed* and *OG\_open*, the MOC is closed by vertical plumes against  
 196 the northern and southern headwall of the domain (G. Hughes et al., 2007; Gayen & Grif-  
 197 fiths, 2022). In these regions, a destabilising buoyancy flux at the ocean surface leads  
 198 to the formation of deep and bottom waters (Gayen & Griffiths, 2022). The heat in both  
 199 cases can be transported by mean flows in the form of western boundary currents, or by  
 200 mesoscale eddies across the re-entrant channel to generate an overturning cell. Below we  
 201 will show that in the cases of partially open gateway there is a third process – standing  
 202 meanders – which can lead to lateral and vertical heat transport, allowing for the for-  
 203 mation of a deep-reaching MOC.

#### 204 3.2 Deep-reaching overturning circulation with gateway deepening

205 A shallow OG, as in experiment *OG\_100*, allows the formation of a vertically-sheared  
 206 circumpolar current (Figure 2i). This circumpolar current is the result of the combina-  
 207 tion of buoyancy loss at the southern boundary, which, together with the lack of merid-  
 208 ional boundaries above the OG depth, leads to the steepening of density surfaces and  
 209 hence a (weak) associated circumpolar current with a zonal transport ( $T_{ACC}$ ) of 3.1 Sv



**Figure 2.** (a)-(h) MOC [Sv] of all simulations; clockwise cell (red) is the upper overturning cell and counterclockwise cell (blue) is the lower overturning cell. (i) Transport through OG ( $T_{ACC}$  [Sv]; green), maximum MOC ( $T_{MOC}$  [Sv]) for the upper cell (red) and the lower cell (blue).

(Figure 2i). This circumpolar current is baroclinically unstable, resulting in the generation of transient eddies and standing meanders, with the latter being generated due to the circumpolar current interacting with the topography of the OG.

A striking change is the deep-reaching MOC in *OG\_100*; despite the very shallow OG, the lower overturning cell now extends to about 3500 m, and increases from 0.72 Sv in *OG\_closed* to 1.67 Sv (global scale: 30.06 Sv; Figure 2b,i). This formation of a strong deep-reaching lower overturning cell is consistent with results by Klocker et al. (submitted), who showed how mesoscale eddies allow for a full-depth overturning cell, and with results by Toumoulin et al. (2020) who used a coarse-resolution earth system model with a 100 m deep DP and Eocene boundary conditions to simulate changes in the MOC due to gateway deepening.

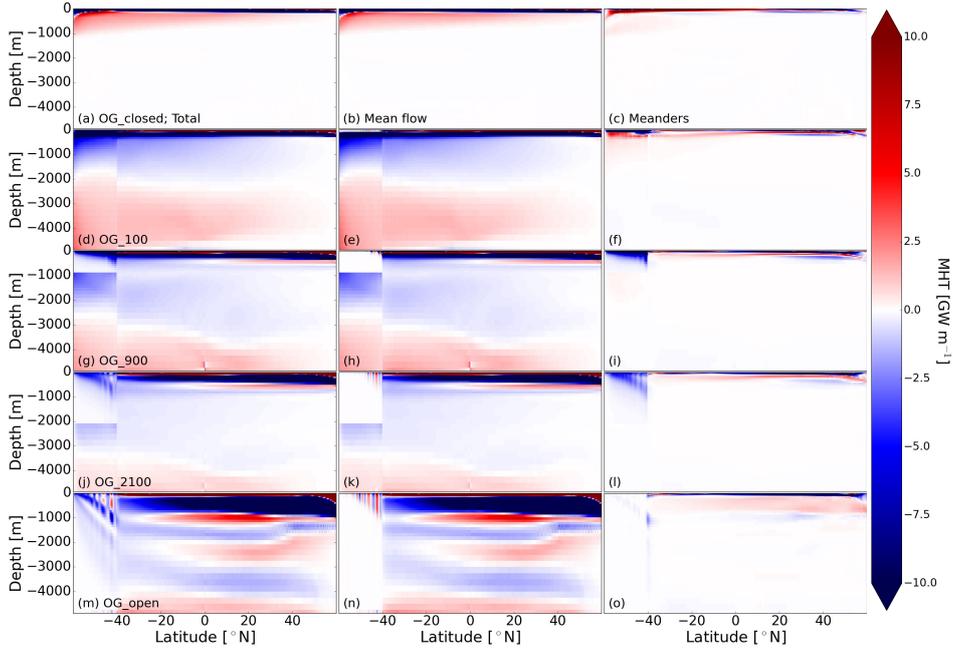
With the deepening of the OG from 100 m to 2100 m (experiments *OG\_100* to *OG\_2100*), the transport of the eastward circumpolar current ( $T_{ACC}$ ) gradually increases to 36.7 Sv (Figure 2i). In contrast with the increasing transport of the circumpolar current, the lower overturning cell continuously weakens to 1.09 Sv (global scale: 19.62 Sv) for the 2100 m OG experiment (Figure 2i). The shallowest OG therefore leads to the strongest lower overturning cell. This is consistent with recent results using a complex earth system model (Toumoulin et al., 2020).

### 3.3 The role of ocean heat transport

Buoyancy (heat) loss at the southern boundary of the channel leads to deep convection (Gayen & Griffiths, 2022) and hence the lower overturning cell. In an equilibrated ocean, maintaining the buoyancy loss and a deep-reaching overturning requires continuous lateral heat transport towards the southern ('Antarctic') boundary, and vertical heat transport to the deep ocean. The lower overturning cell is accordingly controlled by this heat transport. In this section, we will show how changes in OG depth affect the distribution of meridional and vertical heat transport by mean flows, standing meanders, and transient eddies.

For the experiment with a closed OG (*OG\_closed*), the zonally integrated meridional heat transport is dominated by the mean flow (Figure 3a,b), transporting heat southward (blue colours) at the surface and northward (red colours) below. This is consistent with the MOC being confined to the shallow ocean, with the deep ocean being largely at rest (Figure 2a and Figure 3a-c). In the southern hemisphere, standing meanders produce a southward MHT ( $MHT^*$ ) at the surface, and a northward MHT below. Transient eddies lead to a southward MHT across the channel, but contribute less to the heat transport than either the mean flow or standing meanders (Figure S4). The vertically integrated MHT ( $MHT_Z$ ) is strongly biased towards the southern high latitudes, consistent with experiments of thermally-driven flows in a domain lacking a re-entrant channel, in which the flow is biased towards the pole with the colder temperature forcing (Coman et al., 2010, Figure 4a). Consistent with the shallow confinement of the MHT and MOC, the VHT is confined to the shallow ocean (Figure 4b). In summary, in *OG\_closed* the entire heat transport is by baroclinic gyres which are confined to the ocean surface, leading to a shallow MOC. The deep ocean is at rest.

When the OG opens to a depth of 100 m (*OG\_100*), the MHT suddenly extends to the full depth of the basin rather than being confined to the shallow ocean, with a southward MHT from surface to mid-depth, and a northward MHT below (Figure 3e). This corresponds to a full-depth lower overturning cell, which increases in strength from 0.72Sv to 1.67Sv. Standing meanders lead to a southward MHT above the OG depth, and a northward MHT below the OG depth (Figure 3f). Transient eddies still provide weak southward MHT in the channel. The total vertically and zonally integrated southward MHT ( $MHT^*_Z$ ) in the southern half of the basin is reduced by about half (maximum about -10 TW) compared to *OG\_closed*, and is both due to the mean flow and standing me-

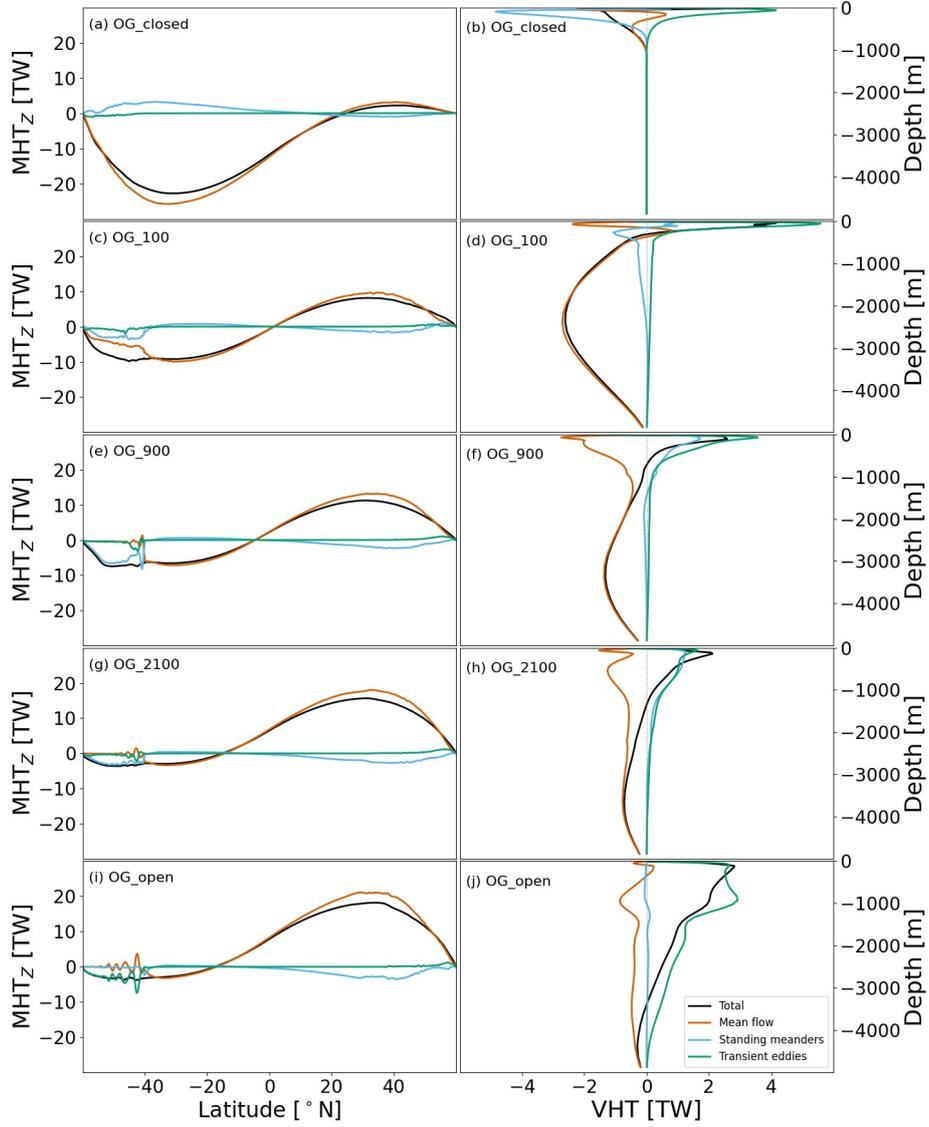


**Figure 3.** Hydrographic section of the zonally integrated meridional heat transport ( $MHT$ ) [ $\text{GW m}^{-1}$ ] for experiments  $OG\_closed$ ,  $OG\_100$ ,  $OG\_900$ ,  $OG\_2100$ , and  $OG\_open$ . Left panels show total  $MHT$ , middle panels show the  $MHT$  by the mean flow ( $\langle MHT \rangle$ ), and right panels show the  $MHT$  by standing meanders ( $MHT^*$ ).

261 anders (Figure 4c). On the other hand, the northward MHT in the northern half of the  
 262 basin increases by 8 TW compared to  $OG\_closed$ . Transient eddies again provide little  
 263 MHT ( $MHT'_z$ ) relative to the mean flow and standing meanders. In contrast to the large  
 264 decrease in MHT by the mean flow in  $OG\_100$ , the VHT by the mean flow shows a large  
 265 increase in strength and now extends to the full depth (Figure 4d).

266 For a further OG deepening (from  $OG\_100$  to  $OG\_2100$ ), the mean MHT in the  
 267 deep ocean decreases in strength compared with  $OG\_100$ , reflected by lighter blue and  
 268 red colors in the Figure 3h and k. This corresponds to a weakening of the deep-reaching  
 269 MOC. In the channel, standing meanders lead to a southward MHT above the OG depth.  
 270 The MHT by transient eddies increases, although it is still weaker than the other terms.  
 271 The gradual deepening of the OG strengthens the transport of the circumpolar current,  
 272 and leads to a state where the vertically integrated MHT is entirely due to standing me-  
 273 anders (Figure 4e and g) in the channel. The MHT by standing meanders becomes weaker  
 274 the deeper the OG, leading to a decrease in strength of the lower overturning cell for a  
 275 further OG deepening. This weakening of the southward MHT is accompanied by a strength-  
 276 ening of the northward MHT, and hence a stronger upper overturning cell.

277 When the OG is fully open ( $OG\_open$ ), standing meanders cannot exist since their  
 278 existence relies on the interaction with a ridge, and instead transient eddies provide south-  
 279 ward MHT. For this experiment the MOC is still full-depth, but weaker than in all cases  
 280 where standing meanders are present. This experiment therefore shows that eddies are  
 281 also capable of allowing for a full-depth lower overturning cell, but given that transient  
 282 eddies are less efficient at transporting heat towards the southern convection region, the  
 283 lower overturning cell is weaker than in all cases where standing meanders are present.



**Figure 4.** (a,c,e,g,i) Zonally and vertically integrated meridional heat transport ( $MHT_z$ ) [TW] for experiments  $OG\_closed$ ,  $OG\_100$ ,  $OG\_900$ ,  $OG\_2100$ , and  $OG\_open$ . (b,d,f,h,j) Horizontal integrated vertical heat transport ( $VHT$ ) [TW] in the channel for experiments  $OG\_closed$ ,  $OG\_100$ ,  $OG\_900$ ,  $OG\_2100$ , and  $OG\_open$ . In all panels, black lines indicate the total ( $MHT_z/VHT$ ), red lines the mean component ( $\langle \overline{MHT} \rangle_z / \langle \overline{VHT} \rangle$ ), green lines the standing meanders ( $MHT^*_z / VHT^*$ ), and blue lines the transient eddies ( $MHT'_z / VHT'$ ).

## 284 4 Discussion and Conclusion

285 In summary, the evolution of the deep-reaching lower overturning cell with the deep-  
286 ening of the ocean gateway can be divided into two parts: the abrupt onset of the deep-  
287 reaching lower overturning cell for a shallow ocean gateway, and the subsequent weak-  
288 ening for a further deepening of the ocean gateway. This behaviour of the lower over-  
289 turning cell is closely linked to changes in southward and downward heat transport in  
290 the presence of standing meanders. First, in an equilibrated ocean, the strength of an  
291 overturning cell is limited by the amount of heat which can be supplied to the vertical  
292 plumes associated with strong surface buoyancy loss. This heat supply is most efficient  
293 if mean flows, such as western boundary currents, are present. This heat supply becomes  
294 less efficient if the heat transport is due to standing meanders, and even less efficient if  
295 heat supply is due to transient eddies. A deepening ocean gateway therefore leads to a  
296 reduced southward heat transport, which is compensated with an increase in northward  
297 heat transport. Second, a deep-reaching overturning cell is dependent on a process which  
298 can get this heat to depth. As shown by Klocker et al. (submitted), this can be achieved  
299 by transient eddies, and, as shown here, this can also be achieved by standing meanders.  
300 Nevertheless, heat transport by baroclinic gyres is always confined to the surface, and  
301 the presence of mesoscale turbulence, whether transient eddies or standing meanders,  
302 is necessary to get this heat to depth, and hence lead to a deep-reaching MOC.

303 The two-part evolution of the lower overturning cell for a deepening of an ocean  
304 gateway is therefore a combination of the two roles transient eddies and standing me-  
305 anders play for ocean heat transport. These dynamics lead to a maximum lower over-  
306 turning cell for the experiment with the shallowest OG, which is a combination of the  
307 domain where the southward heat transport is largest, while at the same time allowing  
308 for a deep-reaching heat transport by mesoscale turbulence. Further deepening of the  
309 OG leads to a weakening of the southward heat supply across the channel, and a strength-  
310 ening of the northward heat transport, associated with a strengthening of the upper over-  
311 turning cell. Once the OG is as deep as the rest of the domain, the southward heat trans-  
312 port is solely due to transient eddies, resulting in the weakest deep-reaching overturn-  
313 ing cell. As opposed to studies which focus on just fully open and fully closed OGs, we  
314 highlight the importance of the dynamics of standing meanders in understanding the evo-  
315 lution of the MOC with gateway deepening, a turbulent process generally neglected in  
316 the paleoceanography community.

317 Previous studies have used complex, but coarse-resolution, earth system models to  
318 explain the evolution of the global MOC with the deepening of Southern Ocean gate-  
319 ways (Sijp & England, 2004; Toumoulin et al., 2020). However, in these studies results  
320 contradict each other, with (Sijp & England, 2004) showing the strongest MOC for a closed  
321 OG, while (Toumoulin et al., 2020), consistent with our study, showing the strongest MOC  
322 for a shallow OG. Nevertheless, both these studies show a weakening of the southward  
323 meridional heat transport with a deepening of the OG, and a strengthening of the north-  
324 ward meridional heat transport, consistent with our study. The different results for the  
325 lower overturning cell are possibly the consequence of the parameterisation of mesoscale  
326 turbulence in these models. It is the simple but high-resolution model configuration used  
327 in this study that helps to reveal the critical role of turbulence processes in generating  
328 the strong deep-reaching overturning cell with OG deepening. Future work is planned  
329 to extend this work on the evolution of the EOT MOC with a global ocean model and  
330 realistic paleo-bathymetry.

## 331 5 Data Availability Statement

332 All model output used in this study to perform the analysis and produce figures  
333 is available at <https://doi.org/10.5281/zenodo.7602996>.

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 340 model simulations.

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 438 10.1175/JPO-D-22-0086.1

# Supporting Information for Deepening of Southern Ocean gateway leads to abrupt onset of vigorous meridional overturning circulation

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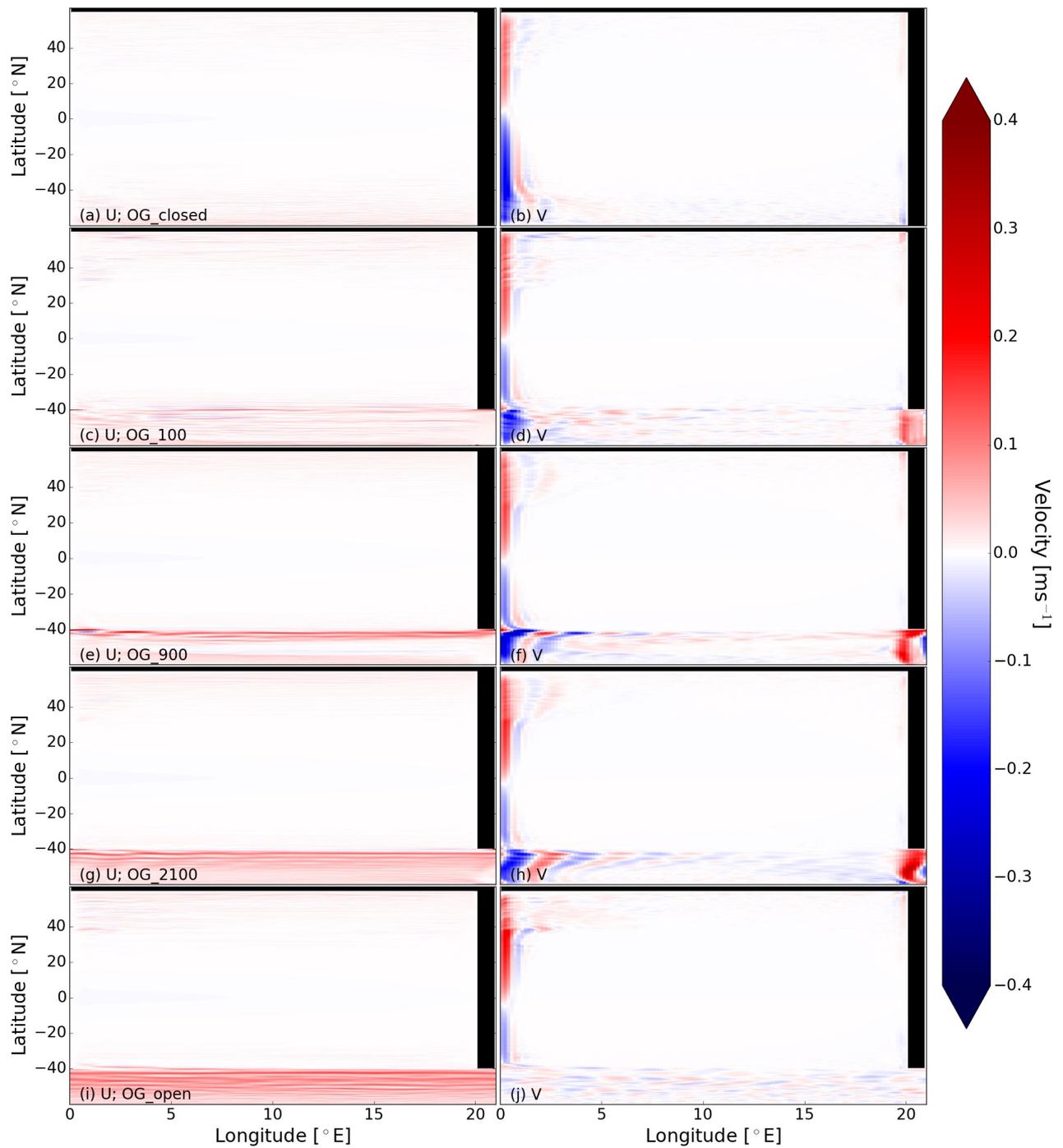
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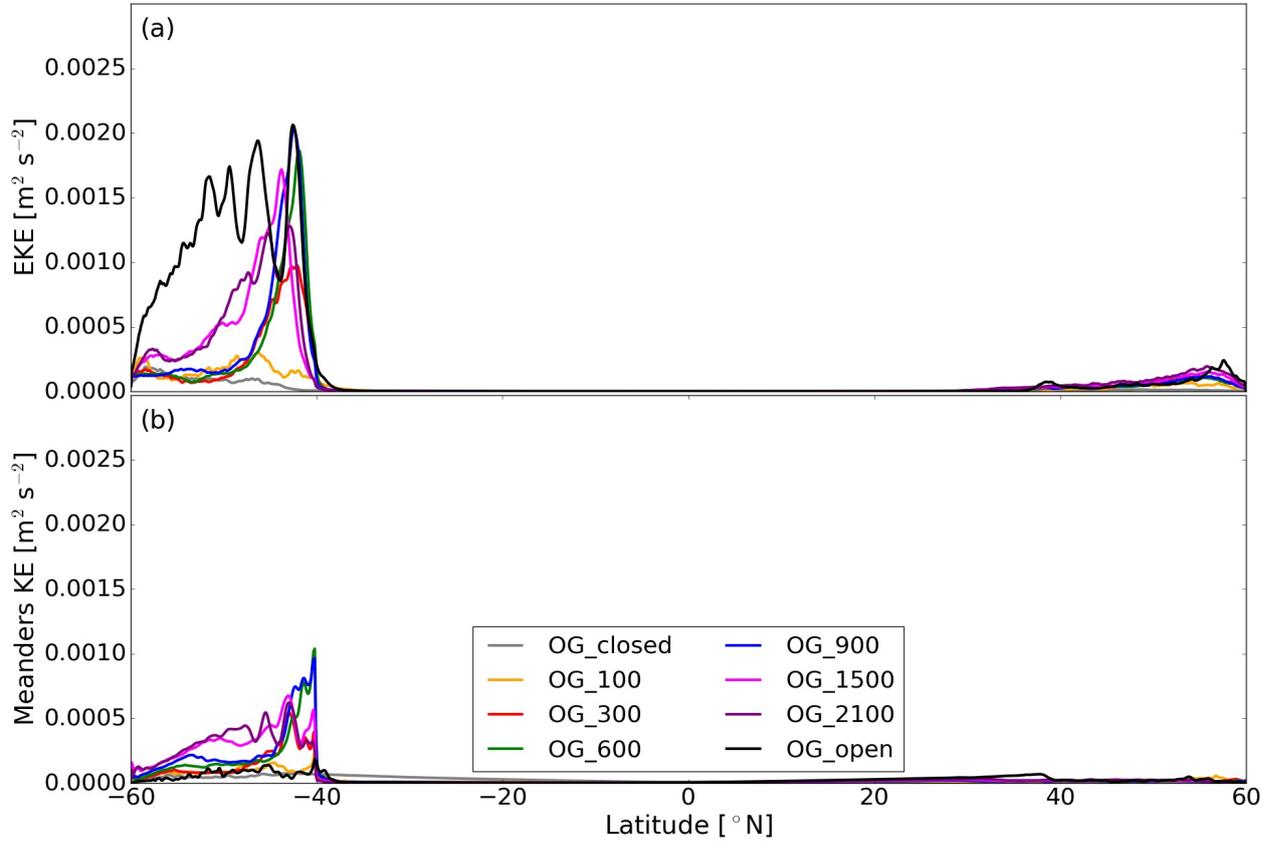
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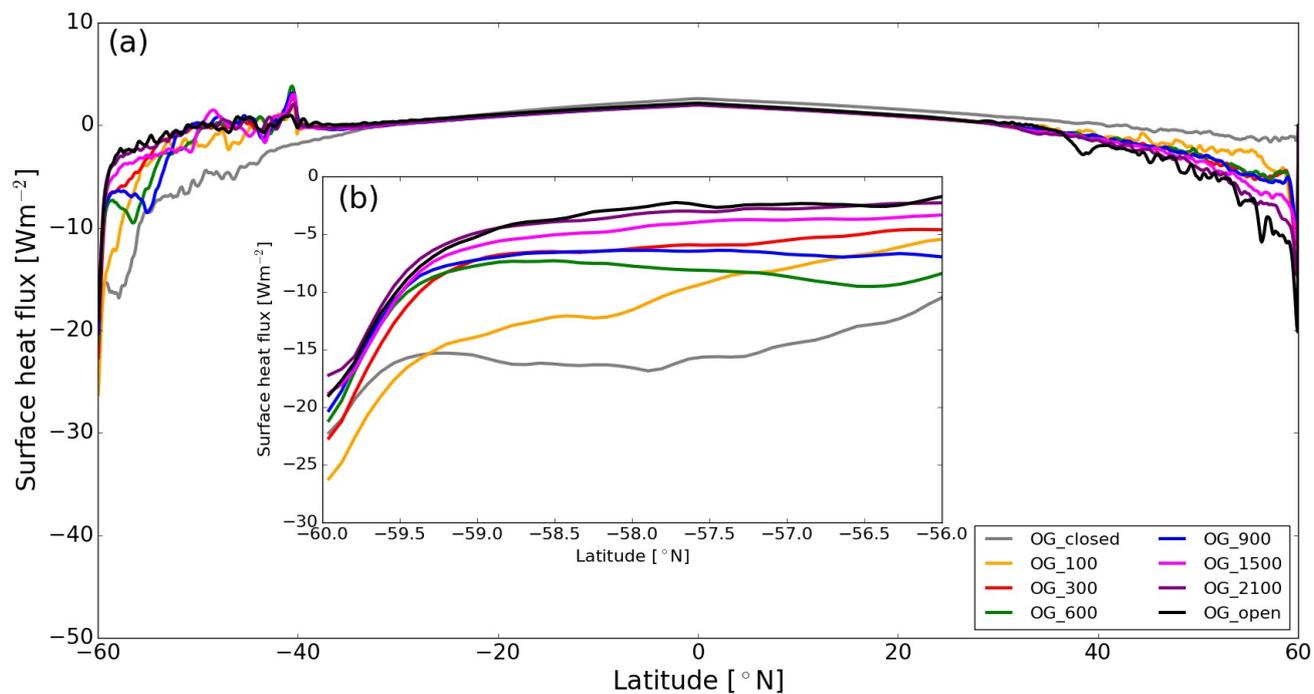
Figures S1 to S5



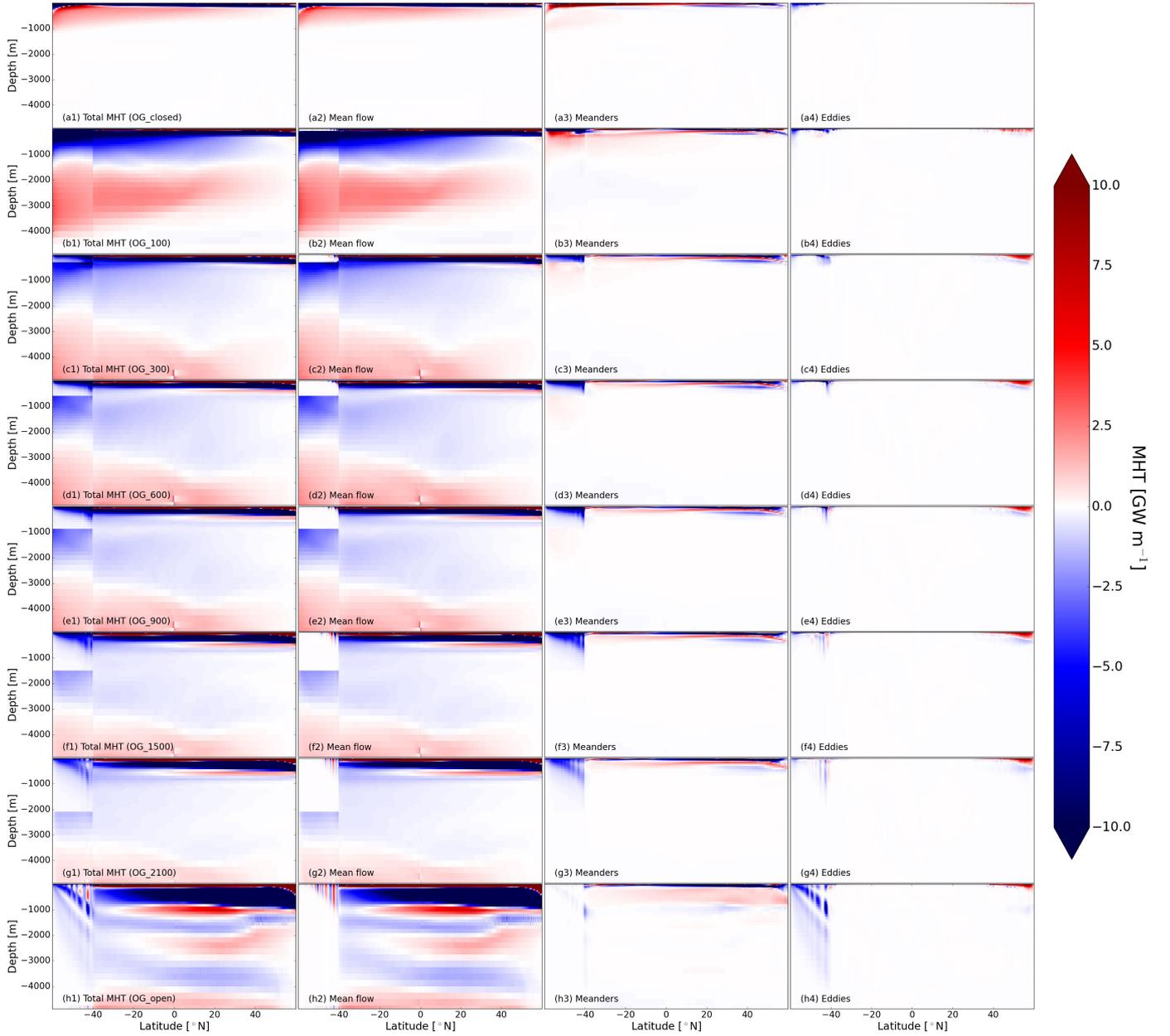
**Figure S1.** (a,c,e,g) Sea surface zonal velocity [ $\text{m s}^{-1}$ ], (b,d,f,h) Sea surface meridional velocity [ $\text{m s}^{-1}$ ] for *OG\_closed*, *OG\_100*, *OG\_900*, *OG\_2100*, and *OG\_open*.



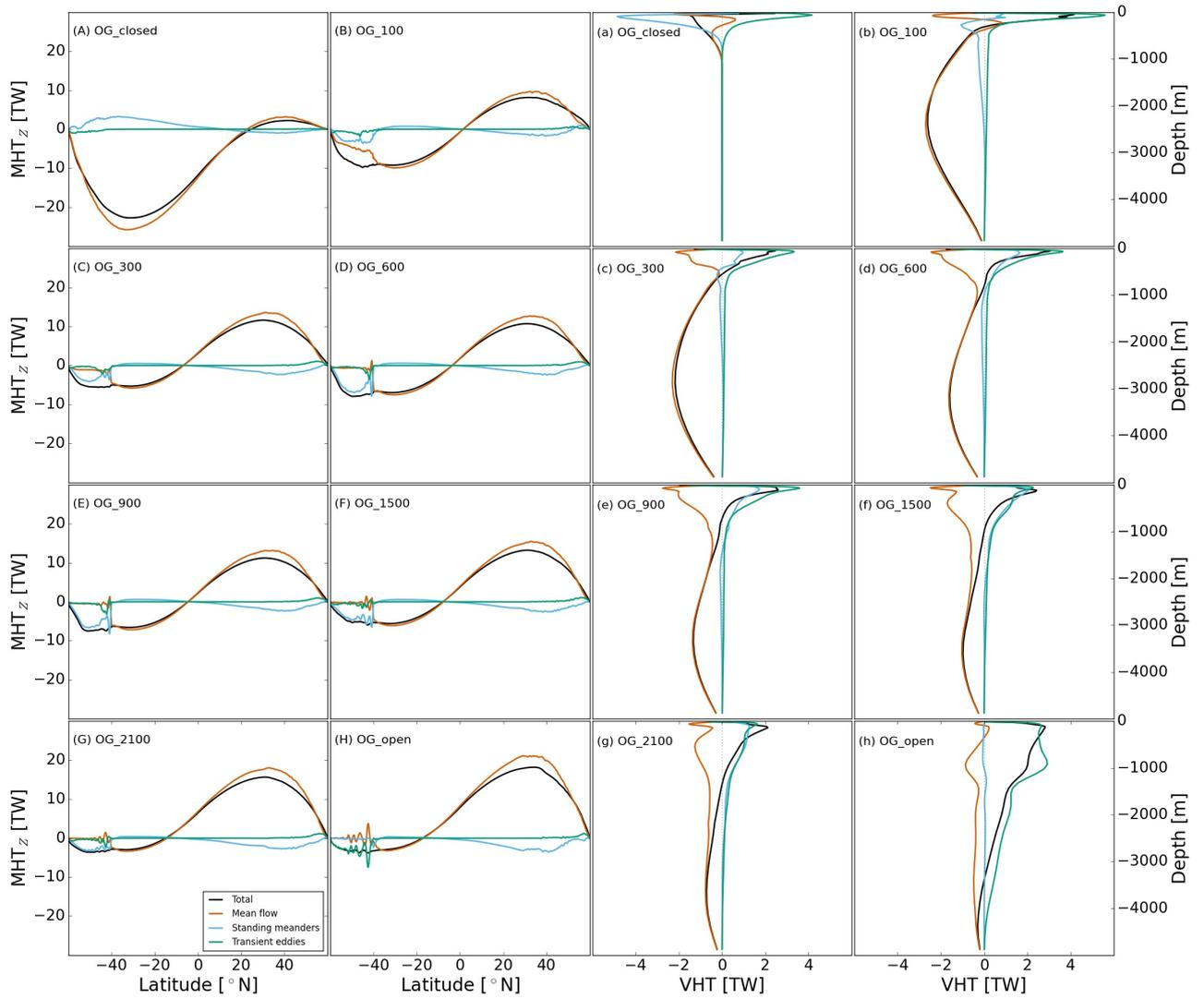
**Figure S2.** (a) Zonally and vertically mean kinetic energy of transient eddies; (b) Zonally and vertically mean kinetic energy of standing meanders. Colored lines show different simulations: *OG\_closed* (grey), *OG\_100* (orange), *OG\_300* (red), *OG\_600* (green), *OG\_900* (blue), *OG\_1500* (magenta), *OG\_2100* (purple), and *OG\_open* (black).



**Figure S3.** Zonally mean surface heat flux [ $\text{W m}^{-2}$ ] for (a) the full domain, and (b) zoomed in on the south of the domain. Colored lines show different simulations: *OG\_closed* (grey), *OG\_100* (orange), *OG\_300* (red), *OG\_600* (green), *OG\_900* (blue), *OG\_1500* (magenta), *OG\_2100* (purple), and *OG\_open* (black).



**Figure S4.** Hydrographic sections of zonally integrated meridional heat transport ( $MHT$ ) [ $10^{-2}$  TW  $m^{-1}$ ] for  $OG\_closed$ ,  $OG\_100$ ,  $OG\_300$ ,  $OG\_600$ ,  $OG\_900$ ,  $OG\_1500$ ,  $OG\_2100$ , and  $OG\_open$ . (a1, b1, c1, d1, e1, f1, g1, h1) total  $MHT$ , (a2, b2, c2, d2, e2, f2, g2, h2) mean flow  $\langle \overline{MHT} \rangle$ , (a3, b3, c3, d3, e3, f3, g3, h3) meanders  $MHT^*$ , (a4, b4, c4, d4, e4, f4, g4, h4) eddy  $MHT'$ .



**Figure S5.** (A-H) Zonally and vertically integrated meridional heat transport ( $MHT_z$ ) [TW]; (a-h) Horizontal mean vertical heat transport ( $VHT$ ) [TW] in the channel for  $OG\_closed$ ,  $OG\_100$ ,  $OG\_300$ ,  $OG\_600$ ,  $OG\_900$ ,  $OG\_1500$ ,  $OG\_2100$ , and  $OG\_open$ . Black lines show total  $MHT_z/VHT$ , red lines indicate the mean flow  $\langle \overline{MHT} \rangle_z / \langle \overline{VHT} \rangle$ , green lines describe standing meanders  $MHT^*_z / VHT^*$ , and blue lines show transient eddy  $MHT'_z / VHT'$ .