

# Modeling the Day-to-Day Variability of Midnight Equatorial Plasma Bubbles with SAMI3/WACCM-X

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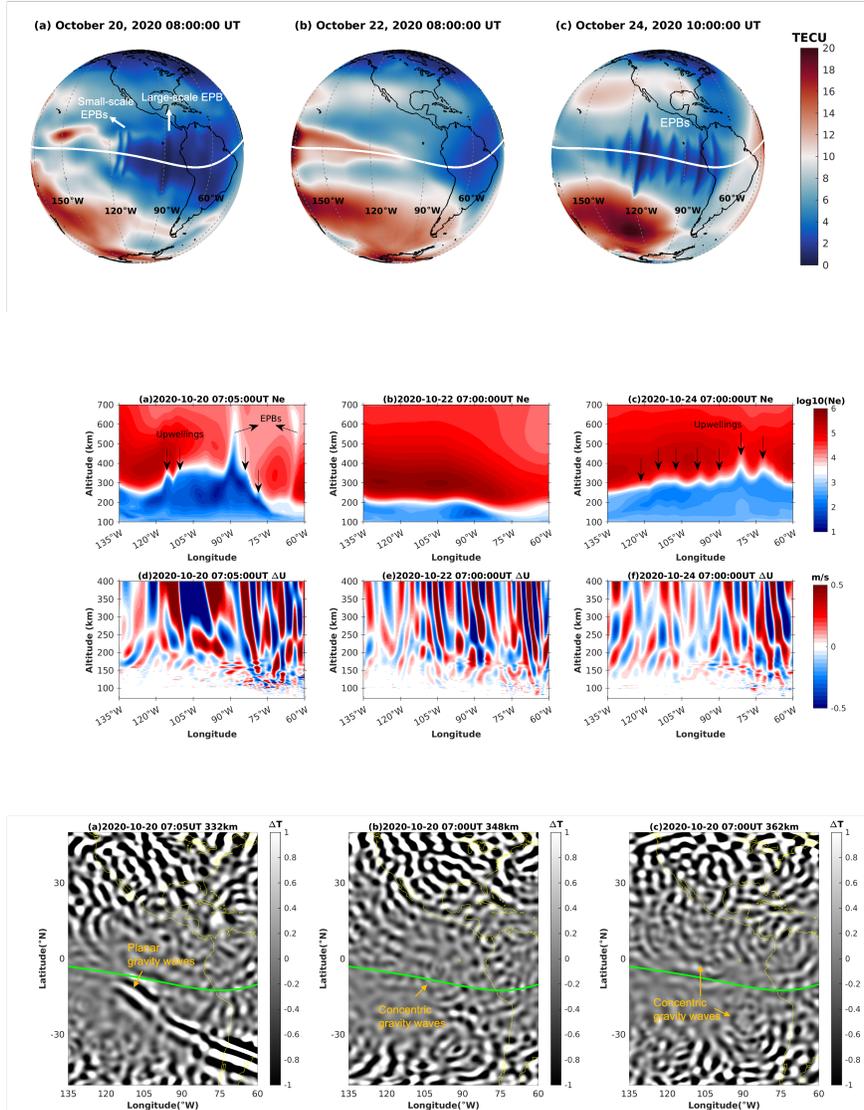
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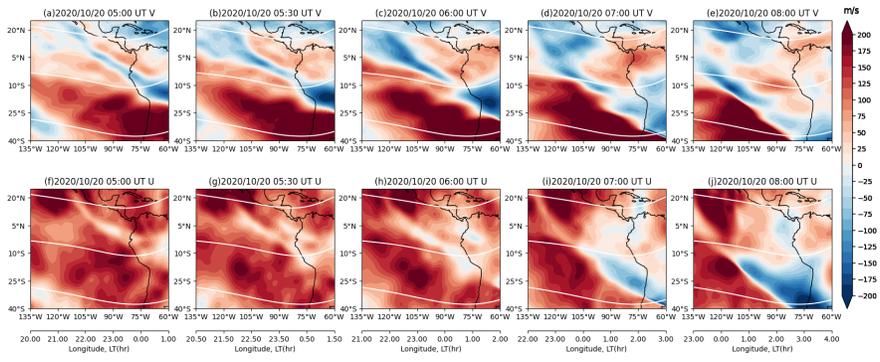
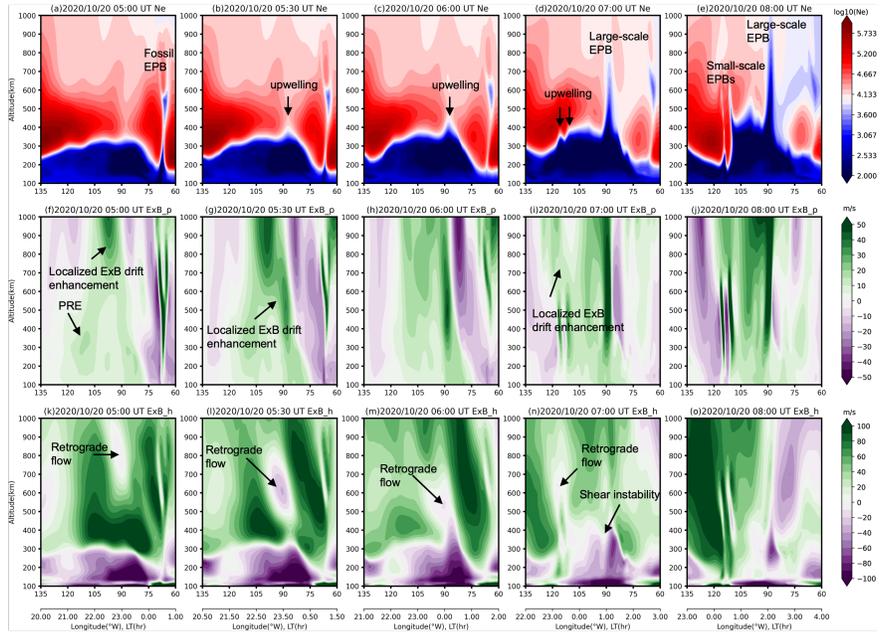
## Abstract

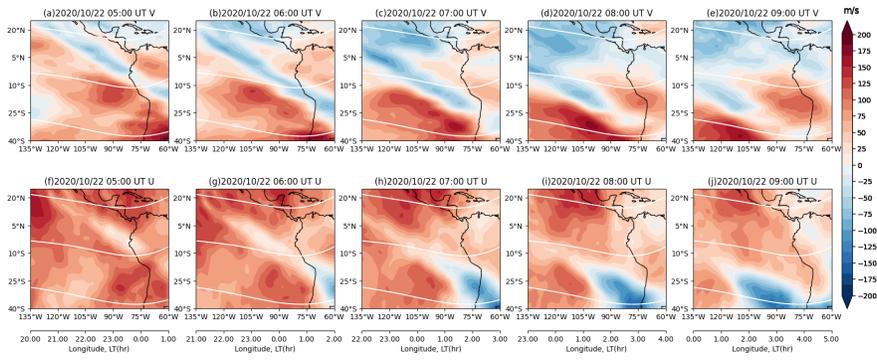
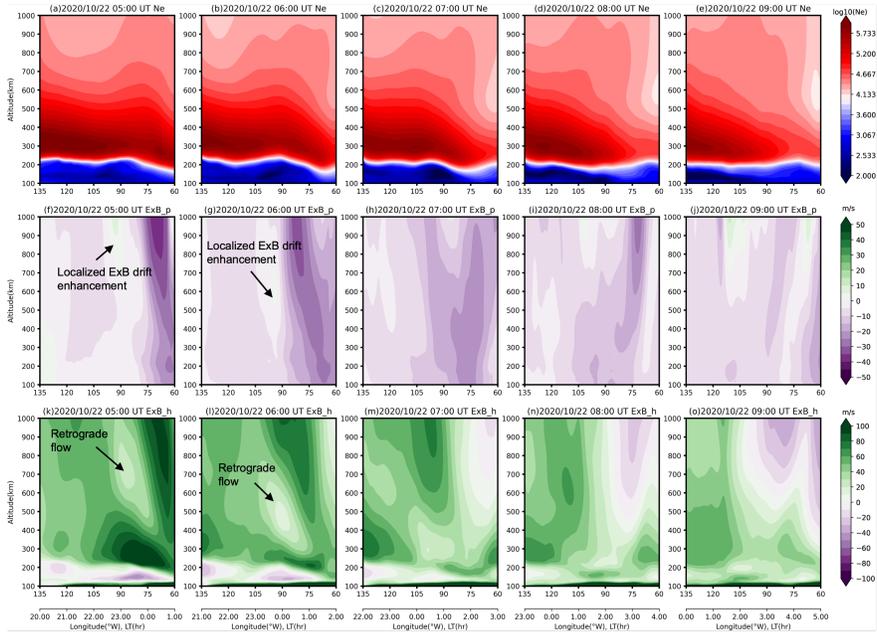
It is well-known that equatorial plasma bubbles (EPBs) are highly correlated to the post-sunset rise of the ionosphere on a climatological basis. However, when proceeding to the daily EPB development, what controls the day-to-day/longitudinal variability of EPBs remains a puzzle. In this study, we investigate the underlying physics responsible for the day-to-day/longitudinal variability of EPBs using the Sami3 is A Model of the Ionosphere (SAMI3) and the Whole Atmosphere Community Climate Model with thermosphere-ionosphere eXtension (WACCM-X). Simulation results on October 20, 22, and 24, 2020 were presented. SAMI3/WACCM-X self-consistently generated midnight EPBs on October 20 and 24, displaying irregular and regular spatial distributions, respectively. However, EPBs are absent on October 22. We investigate the role of gravity waves on upwelling growth and EPB development and discuss how gravity waves contribute to the distributions of EPBs. Of particular significance is that we found the westward wind associated with solar terminator waves and gravity waves causes midnight vertical drift enhancement and collisional shear instability, which provides conditions favorable for upwelling growth and EPB development. The converging and diverging winds associated with solar terminator waves and midnight temperature maximum also affect the longitudinal distribution of EPBs. The absence of EPBs on October 22 is related to the weak upward drift induced by weak westward wind associated with solar terminator waves.

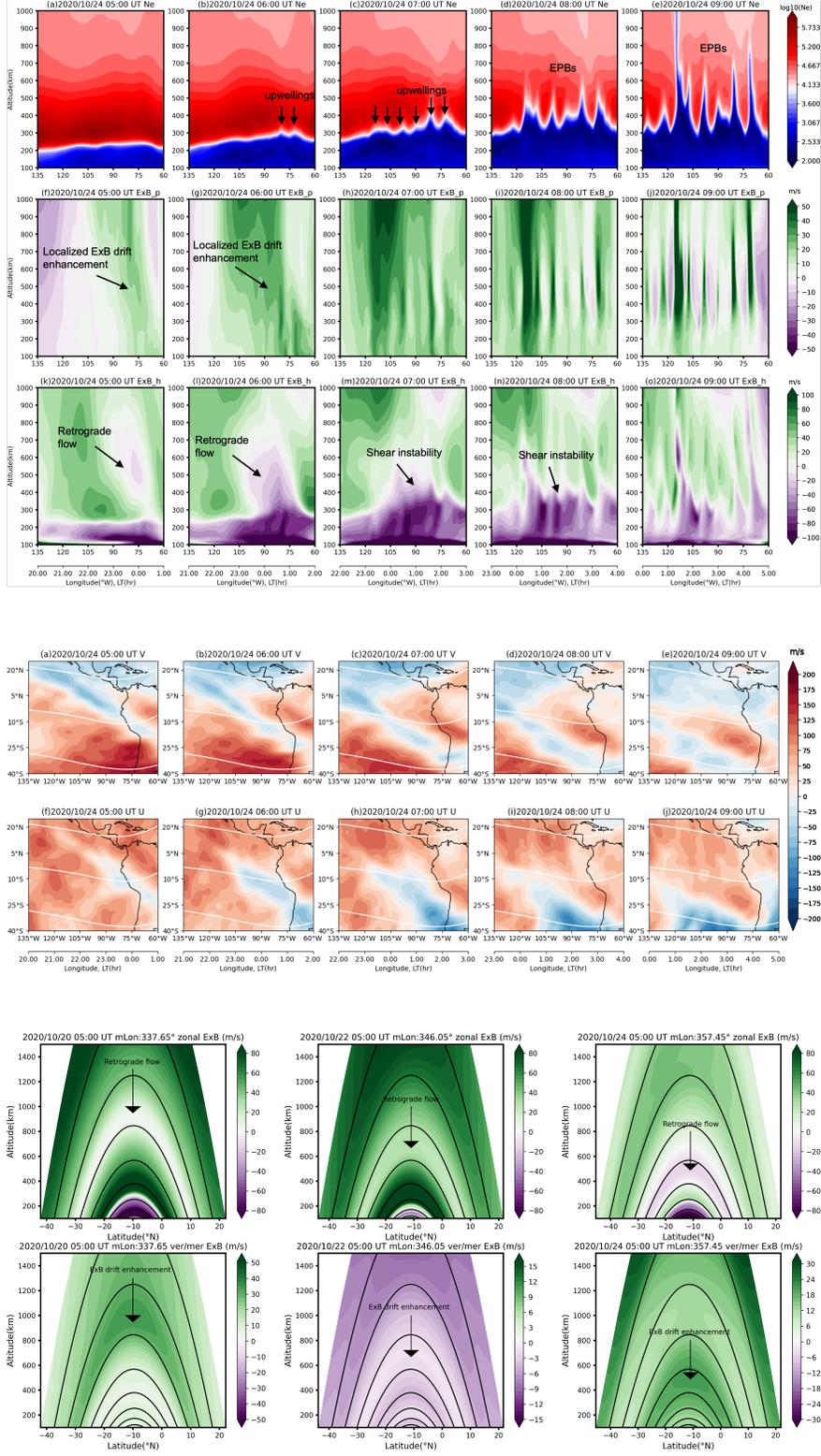
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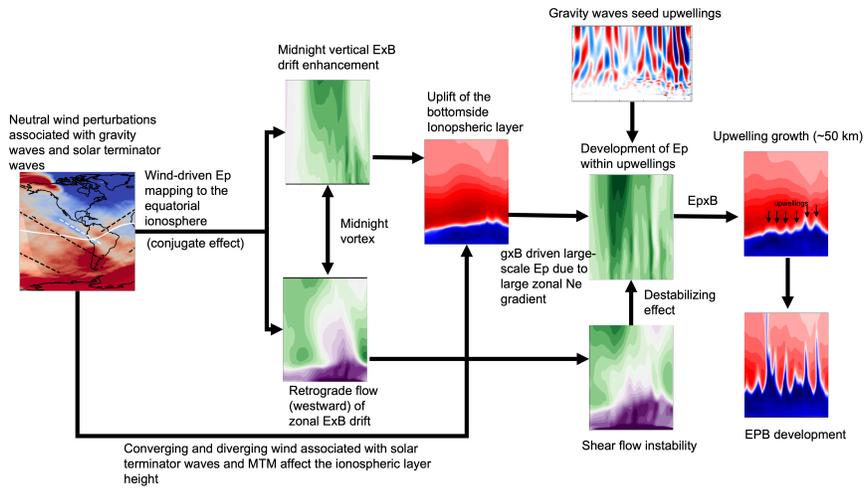
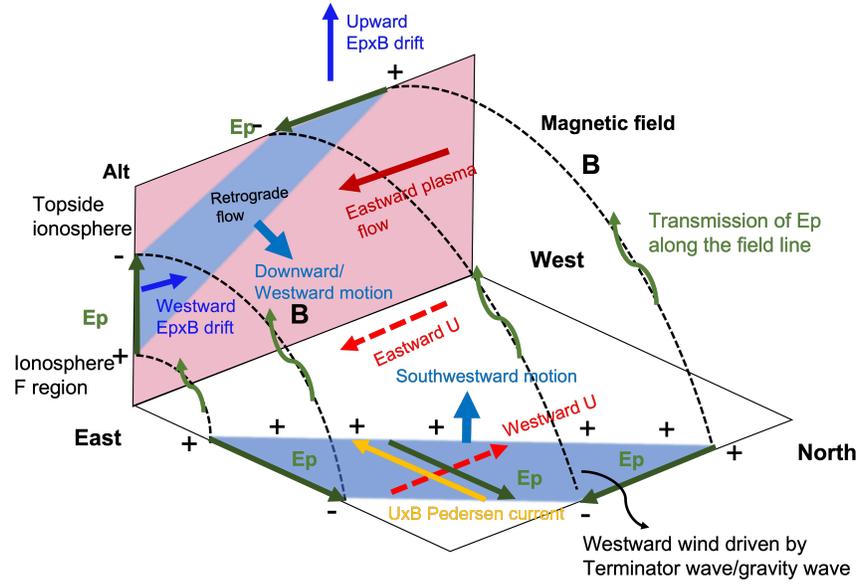
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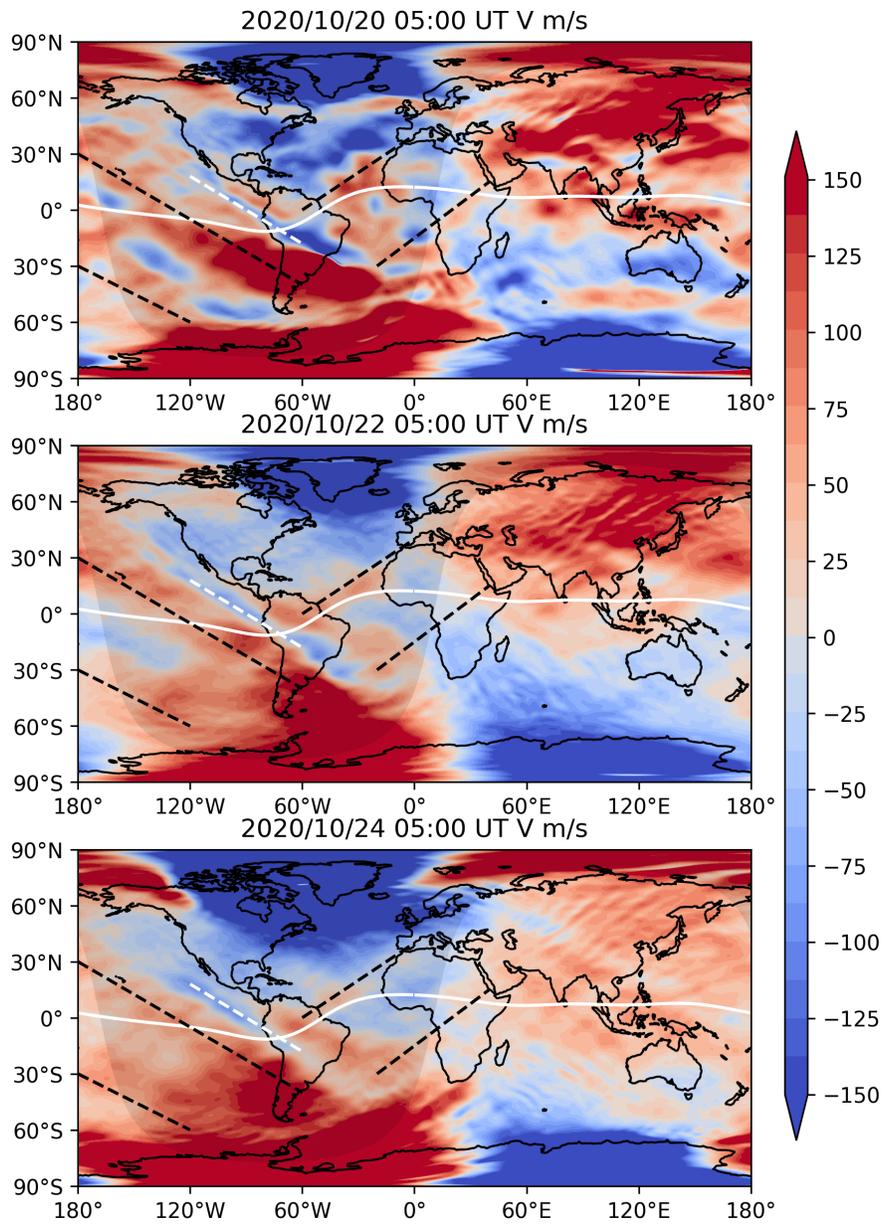












1 **Modeling the Day-to-Day Variability of Midnight Equatorial Plasma Bubbles with**  
2 **SAMI3/WACCM-X**

3

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5 Tate<sup>5</sup>, Nicholas Pedatella<sup>6,7</sup>, Cora E. Randall<sup>8,9</sup>, V. Lynn Harvey<sup>8,9</sup>

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22 **Key points:**

- 23 1. SAMI3/WACCM-X self-consistently generates EPBs at midnight.
- 24 2. Gravity waves and meridional winds affect the spatial and longitudinal distributions  
25 of EPBs.
- 26 3. Westward winds associated with solar terminator and gravity waves facilitate the  
27 midnight EPB development by generating midnight vortex.

28

29 **Abstract**

30 It is well-known that equatorial plasma bubbles (EPBs) are highly correlated to the post-  
31 sunset rise of the ionosphere on a climatological basis. However, when proceeding to the  
32 daily EPB development, what controls the day-to-day/longitudinal variability of EPBs  
33 remains a puzzle. In this study, we investigate the underlying physics responsible for the day-  
34 to-day/longitudinal variability of EPBs using the Sami3 is A Model of the Ionosphere  
35 (SAMI3) and the Whole Atmosphere Community Climate Model with thermosphere-  
36 ionosphere eXtension (WACCM-X). Simulation results on October 20, 22, and 24, 2020  
37 were presented. SAMI3/WACCM-X self-consistently generated midnight EPBs on October  
38 20 and 24, displaying irregular and regular spatial distributions, respectively. However, EPBs  
39 are absent on October 22. We investigate the role of gravity waves on upwelling growth and  
40 EPB development and discuss how gravity waves contribute to the distributions of EPBs. Of  
41 particular significance is that we found the westward wind associated with solar terminator  
42 waves and gravity waves causes midnight vertical drift enhancement and collisional shear  
43 instability, which provides conditions favorable for the upwelling growth and EPB  
44 development. The converging and diverging winds associated with solar terminator waves  
45 and midnight temperature maximum also affect the longitudinal distribution of EPBs. The  
46 absence of EPBs on October 22 is related to the weak upward drift induced by weak  
47 westward wind associated with solar terminator waves.

48

49

50 **Plain Language Summary**

51 Plasma bubbles are a particular space weather phenomenon that mainly occurs in the  
52 nighttime equatorial region. After sunset, the bottomside ionosphere (~100-200 km) becomes  
53 unstable due to the vertical motion of the ionosphere. Bubbles can develop from the  
54 bottomside ionosphere and stretch into the topside ionosphere (above 500 km), like wax  
55 bubbles in a lava lamp. Bubbles significantly reduce the plasma density in the ionosphere,  
56 displaying turbulent plume structures that can disrupt radio wave communications and GPS  
57 navigation. Understanding and predicting the development of plasma bubbles has baffled  
58 scientists for more than 80 years, especially in understanding the day-to-day variability. In  
59 this study, we aim to understand what controls the day-to-day variability of plasma bubbles  
60 by using the physics-based SAMI3/WACCM-X model. We found that gravity waves are  
61 ubiquitous and play a vital role in seeding and determining the spacing between plasma  
62 bubbles. The longitudinal distribution of plasma bubbles is affected by meridional wind. The  
63 most striking finding is that daily dusk solar terminator waves significantly impact neutral  
64 wind and electrodynamics, controlling the presence or absence of plasma bubbles at midnight.  
65 This study reveals that the day-to-day variability of plasma bubbles is considerably linked to  
66 the variations of the lower atmosphere.

## 67 1. Introduction

68 Equatorial spread F (ESF) and equatorial plasma bubbles (EPBs) are ionosphere  
69 irregularities that primarily occur in the nighttime equatorial ionosphere. Brook and Wells  
70 (1938) first observed spread echoes from the ionospheric F region, referred to as ESF, using  
71 ionosondes. Woodman and LaHoz (1976) proposed the concept of ionospheric “bubbles” to  
72 illustrate the nonlinear evolution of plasma depletions from the bottom to the topside  
73 ionosphere. EPBs are field-aligned structures in the form of meridionally-elongated wedges  
74 of plasma depletions in both hemispheres (e.g., Kil et al., 2009), which are characterized by  
75 bite-outs in ion density measurements (Heelis et al., 2010; Yokoyama et al., 2011), plume  
76 structures in radar observations (Kelley et al., 1981; Hysell et al., 2009), intensity depletions  
77 in airglow images (e.g., Kelley et al., 2003; Otsuka et al., 2002; Eastes et al., 2019; Chou et  
78 al., 2020a), and turbulent fluctuations in Global Navigation Satellite System Total Electron  
79 Content (TEC)(Nishioka et al., 2008; Cherniak and Zakharenkova, 2016). Understanding and  
80 forecasting the presence of EPBs is an essential topic since they can disrupt propagation of  
81 radio waves used in global communication and navigation systems (e.g., Kelley et al., 2014;  
82 Xiong et al., 2016) and cause scintillations in radio signals (e.g., Yeh and Liu, 1982; Kintner  
83 et al., 2007).

84 Tsunoda (1985) first proposed the longitudinal and seasonal distribution of EPBs is  
85 related to the angle between dusk solar terminator and geomagnetic field line at the magnetic  
86 equator. The pre-reversal enhancement (PRE) of upward  $E \times B$  drift occurs when the dusk  
87 solar terminator is aligned with the geomagnetic field line, resulting in the post-sunset rise  
88 (PSSR) of ionosphere. The PSSR destabilizes the ionosphere and allows EPBs to develop  
89 through the Rayleigh-Taylor (RT) instability (Sultan et al., 1996). The PSSR-to-EPB  
90 paradigm (Tsunoda et al., 2018) is supported by satellite observations that the PRE controls  
91 the EPB occurrence on a climatological basis (Burke et al., 2004; Gentile et al., 2006; Huang  
92 and Hairston, 2015).

93 However, PRE fails to explain the occurrence of EPBs on a day-to-day basis. EPB  
94 development during the post-midnight have been observed by the Formosa satellite-1  
95 (FORMOSAT-1 or ROCSAT-1), Communication/Navigation Outage Forecasting System  
96 (C/NOFS), and radar observations (e.g., Yizengaw et al., 2000; Yokoyama et al., 2011;  
97 Nishioka et al., 2012). Tsunoda (2015) further proposed an upwelling paradigm to describe  
98 the processes of EPB development from seeding, upwelling growth, and EPB formation.  
99 Upwelling (local uplift of the bottomside ionosphere) or large-scale wave structure (LSWS, a  
100 continuous distribution of upwellings) is the undulation of the bottomside ionosphere, mainly  
101 driven by an eastward polarization electric field (Ep). Tsunoda et al. (2018) suggested that  
102 the amplification of upwelling (i.e., upwelling growth) is comparable to the post-sunset rise  
103 (PSSR) of the ionosphere and can make an additive localized uplift to the PSSR by ~50 km  
104 (e.g., Chou et al., 2020a), leading to the conclusion that upwelling growth controls the EPB

105 development, instead of PRE. The source of upwelling remains a mystery; however, seed  
106 perturbations related to gravity waves are considered to be the most credible source of  
107 upwellings (e.g., Tulasi Ram et al., 2014; Chou et al., 2020a; Huba and Liu, 2020).

108 Understanding the complexities of the underlying physics responsible for day-to-day  
109 variability of EPBs remains a challenge. Various observation and modeling efforts have been  
110 conducted to investigate underlying physics responsible for the day-to-day variability of  
111 EPBs, such as seed perturbations (Singh et al., 1997; Abdu et al., 2009; Retterer et al., 2014;  
112 Krall et al., 2013), neutral winds (Maruyama and Matuura, 1984; Huba et al., 2009; Krall et  
113 al., 2009, 2021; Huba and Krall, 2013), vertical drifts (Retterer et al., 2005; Su et al., 2009),  
114 shear instability (Hysell and Kudeki, 2004; Yokoyama et al., 2015), Es-layer instability  
115 (Tsunoda, 2007; Huba et al., 2020), tidal forcing (Tsunoda et al., 2015; Chang et al., 2020;  
116 Chou et al., 2020b), upwelling growth (e.g., Tsunoda, 2015), and penetration electric fields  
117 due to geomagnetic storms (e.g., Cherniak and Zakharenkova, 2016; Rajesh et al., 2017).  
118 These studies primarily focus on a single driver that controls EPB development, and artificial  
119 seed perturbations are required in the initial conditions for EPB simulations (e.g., Yokoyama,  
120 2017). However, the onset of EPBs could concurrently involve multiple drivers and physical  
121 processes. Limited observational instruments and modeling capability prohibit a complete  
122 understanding of the complex physical processes of EPB onset. Therefore, comprehensive  
123 observations and coupled whole atmosphere/ionosphere models that consider more realistic  
124 background conditions and include all drivers (e.g., Huba and Liu, 2020; Hysell et al., 2022),  
125 are necessary to provide the whole picture for comprehending the morphology and day-to-  
126 day variability of EPBs.

127 Recent advances in satellite measurement techniques and modeling capabilities have  
128 enabled improved understanding of the complex processes that cause day-to-day variability  
129 of EPBs. The National Aeronautics and Space Administration (NASA) Global-scale  
130 Observations of the Limb and Disk (GOLD) mission has provided unprecedented daily  
131 observations of equatorial ionization anomaly (EIA) images from western Africa to South  
132 America. Eastes et al. (2019) reported that GOLD observed EPBs on most nights, displaying  
133 significant spatial and temporal variability that is unexpected during solar minimum  
134 conditions. Huba and Liu (2020) further conducted a high-resolution global simulation of  
135 EPBs using the Sami3 is A Model of the Ionosphere (SAMI3) and the high-resolution Whole  
136 Atmosphere Community Climate Model with thermosphere-ionosphere eXtension  
137 (WACCM-X). The coupled SAMI3/WACCM-X self-consistently generated EPBs for the  
138 first time, comparable to the GOLD observations (Eastes et al., 2019). They found that EPBs  
139 developed for a March case but not for a July case, which agrees well with the observations  
140 (e.g., Gentile et al., 2006). Huba and Liu (2020) suggested that gravity waves play an  
141 essential role in seeding EPBs because EPBs are absent when SAMI3 is coupled to empirical  
142 models, such as HWM and MSIS.

143        However, many questions remain unsolved with regard to EPBs: What is the linkage  
144 between gravity waves and upwellings? What is the most crucial factor that controls  
145 upwelling growth (e.g., Tsunoda, 2015)? What influences the spacing between EPBs and the  
146 longitudinal distribution of EPBs? Why do EPBs show isolated clusters separated by long  
147 distances on some nights but display a continuous distribution of EPB trains on other nights?  
148 Why do EPBs occur on some nights and not on others? What is the physics responsible for  
149 the EPBs occurred during midnight without PRE (e.g., Otsuka, 2018)? What is the  
150 underlying mechanism for generating large-scale EPBs (e.g., Eastes et al., 2019)?

151        In this study, the coupled SAMI3/WACCM-X model is utilized to investigate the day-  
152 to-day variability of EPBs. Simulation on October 20, 22, and 24 in 2020, during a solar  
153 minimum period, is presented. EPBs are generated on October 20 and 24 at midnight, but not  
154 on October 22. EPBs display irregular and regular spatial distributions on October 20 and 24,  
155 respectively. The underlying mechanisms and background conditions that cause the absence  
156 and presence of the midnight EPBs, as well as the spatial distribution are discussed. We  
157 outline the effects of gravity waves and neutral winds on the longitudinal distribution of  
158 EPBs and elucidate how solar terminator waves affect the ionospheric electrodynamics and  
159 facilitate the midnight EPB development. This study affords new insight into the day-to-day  
160 variability of EPBs during solar minimum.

161

## 162 **2. SAMI3/WACCM-X**

163        In this work we performed simulations using the SAMI3 model driven by WACCM-X  
164 (McDonald et al., 2015). SAMI3 is a global, three-dimensional, physics-based ionosphere  
165 model. It is based on the two-dimensional SAMI2 model (Huba et al., 2000). SAMI3 models  
166 the plasma and chemical evolution of seven ion species ( $H^+$ ,  $He^+$ ,  $N^+$ ,  $O^+$ ,  $N_2^+$ ,  $NO^+$ , and  $O_2^+$ )  
167 and solves the ion continuity and momentum equations for seven ion species. Ion inertia is  
168 included in the ion momentum equation for motion along the geomagnetic field. The electric  
169 fields driven by the neutral wind dynamo are self-consistently solved from the potential  
170 equation based on current conservation ( $\nabla \cdot J = 0$ ) and equipotential field lines (e.g., Huba et  
171 al., 2008). The model also solves the complete temperature equations for electrons and three  
172 ion species ( $H^+$ ,  $He^+$ , and  $O^+$ ). SAMI3 uses the solar EUV irradiance model for aeronomic  
173 calculations (EUVAC). The Richmond Apex model (Richmond, 1995) is used to specify the  
174 magnetic field (i.e., International Geomagnetic Reference Field, IGRF). The thermospheric  
175 inputs of neutral composition, temperature and winds can be specified in SAMI3 by  
176 analytical models, empirical models (e.g., HWM and MSIS), or physics-based models (e.g.,  
177 Huba et al., 2010, 2017).

178        In this study, the thermospheric variables (neutral densities, winds, and temperatures)  
179 from WACCM-X are inputs into SAMI3 (e.g., McDonald et al., 2015, 2018). A detailed  
180 description of WACCM-X is given in Liu et al. (2018). The WACCM-X resolution is

181 0.47°×0.625° in latitude and longitude. The upper boundary of WACCM-X is at  $4\times 10^{-10}$  hPa  
182 (approximately 450 km on average). SAMI3 uses a geomagnetic grid of dimension (nz, nf, nl)  
183 = (160,160,194), where nz is the number of grid points along the magnetic field line, nf is the  
184 number of field lines, and nl is the number of magnetic longitudes. SAMI3 used a non-  
185 uniform longitudinal grid in this study, including coarse- and high-resolution regions (e.g.,  
186 Huba et al., 2010). The longitudinal resolution is 0.6° from ~63.6°-136.5°W and 4° at the  
187 other longitudes. The latitudinal resolution is variable due to the nonlinear spacing of grid  
188 points along field lines. The resolution is approximately 0.2° near the magnetic equator and  
189 0.66° at 40° latitude at ~300 km altitude. Simulation on October 20, 22, and 24 in 2020 is  
190 performed using the following geophysical conditions: F10.7=74, 74.2, 71.3; F10.7A=78.2,  
191 79, 79.8; Ap = 4, 5, 17; Kp = 1, 1, 3. EPBs develop in the high-resolution region; thus, we  
192 focus on the region from ~63.6°-136.5°W.

193

### 194 **3. Results and Discussions**

#### 195 **3.1 Day-to-Day variability of EPBs**

196 Figure 1 shows the TEC simulated from SAMI3/WACCM-X at 08:00 UT, 08:00 UT,  
197 and 10:00 UT on 20, 22, and 24 October 2020, respectively. Note that different UT times are  
198 presented for each day due to the difference in EPB onset time. Distinct TEC depletions  
199 associated with EPBs are discernible in the equatorial ionosphere on October 20 and 24 but  
200 not on October 22. EPBs display irregular spatial distribution with two groups of EPBs on  
201 October 20. On October 20, the first group shows two isolated small-scale EPBs from  
202 105°W-120°W, and the other shows one large-scale EPB around 90°W over the Pacific  
203 Ocean. These EPBs developed around the local midnight. There are no EPBs on October 22,  
204 but a regular spatial distribution of successive post-midnight EPBs occurred on October 24.  
205 Approximately eight clusters of EPBs spanning ~75° in longitude can be discerned.

206 Of particular interest is the mechanism that causes regular and irregular spatial  
207 distributions of EPBs. Both irregular and regular spatial distributions of EPBs are commonly  
208 observed by satellite observations such as the C/NOFS and GOLD (e.g., Huang et al., 2013;  
209 Eastes et al., 2019). Makela et al. (2010) suggested that gravity waves in the bottomside  
210 ionosphere play a vital role in the quasi-periodically spaced EPBs (Figure 1c); however, the  
211 underlying mechanism responsible for the long-distance separation of the EPB groups  
212 (Figure 1a) remains unknown. Additionally, Figure 1a shows a large-scale EPB near the west  
213 coast of South America. Eastes et al. (2019) first identified the large-scale EPB with  
214 significant deviations in separation of the EIA crests compared to the adjacent longitudes.  
215 They suggested that penetration electric fields due to negative excursion in the interplanetary  
216 magnetic field Bz may be responsible for the abrupt shifts of EIA. Nevertheless, the exact  
217 mechanism responsible for the large-scale EPBs remains unknown.

218

219 **3.2 Gravity Wave Seeding and Upwelling growth**

220 Tsunoda et al. (2018) suggested that upwelling growth controls the EPB development  
221 and gravity waves appear to be the most credible source of upwellings (e.g., Tulasi Ram et al.,  
222 2014; Chou et al., 2020a). To investigate the linkage between gravity waves and upwellings,  
223 Figure 2 shows the electron density (top panels) and zonal wind perturbations (bottom panels)  
224 as a function of longitude and altitude on 20, 22, and 24 October. Wu et al. (2015) suggested  
225 that zonal and vertical wind perturbations associated with gravity waves were most effective  
226 in seeding EPBs because the zonal and vertical winds can effectively modify the electrostatic  
227 potential. Thus, we extract the zonal wind perturbations by applying a fifth-order high-pass  
228 filter with a cutoff period of 45 min, which covers typical period ranges for gravity waves  
229 from various sources in the upper atmosphere and ionosphere (e.g., Azeem et al., 2015; Chou  
230 et al., 2017; Sharon and Azeem, 2021; Heale et al., 2022; Yue et al., 2022).

231 Multiple instances of upwelling (indicated by black arrows) can be identified in the iso-  
232 density contours of  $\sim 10^{3.5} \text{ cm}^{-3}$  along the bottomside ionosphere before the EPB development  
233 at  $\sim 07:05$  UT and  $07:00$  UT on 20 and 24 October (Figures 2a and 2c). These upwelling  
234 structures are identical to the incoherent scatter radar observations (see Figure 1 of Tsunoda  
235 et al., 2018). The zonal scales of upwellings are also consistent with previous observations of  
236  $\sim 100$ - $1500$  km (Tsunoda, 2021). EPBs eventually developed from the crests of upwellings as  
237 shown in Figure 1. On the other hand, Figure 2b shows no evidence of upwellings; this is due  
238 to the lower bottomside ionospheric layer height of  $\sim 150$ - $200$  km (iso-density contour of  
239  $\sim 10^{3.5} \text{ cm}^{-3}$ , or peak density height (hmF2)  $\sim 250$ - $300$  km) on this night compared with  
240 Figures 2a and 2c (above  $\sim 300$  km in iso-density contour of  $\sim 10^{3.5} \text{ cm}^{-3}$ , or hmF2  $\sim 350$ - $400$   
241 km). EPBs tend to develop when hmF2 is around  $350$ - $400$  km, generally consistent with the  
242 FORMOSAT-3/COSMIC observations (e.g., Chou et al., 2020b). Note large-scale EPB and  
243 fossil EPB are presented at  $\sim 90^\circ\text{W}$  and  $\sim 65^\circ\text{W}$ , respectively, in Figure 2a. Two upwellings  
244 within the longitude range of  $75$ - $90^\circ\text{W}$  in Figure 2a do not lead to EPB development because  
245 the lower ionosphere height inhibits the upwelling growth.

246 The bottom panels of Figure 2 are the zonal wind perturbations extracted by a high-pass  
247 filter, which can be attributed to gravity waves in the WACCM-X model (Liu et al., 2014).  
248 Gravity wave seeding is important for EPB development (Huba and Liu, 2020). We found  
249 that the zonal scales of upwellings and zonal wind perturbations are generally comparable.  
250 This reveals that the zonal scale of gravity waves plays a vital role in determining the spacing  
251 between EPBs. However, gravity waves alone are insufficient for upwelling growth (Figure  
252 2b); a sufficiently high ionospheric layer is essential to facilitate the upwelling growth since  
253 the lower ionospheric layer height results in higher ion-neutral collision frequency and  
254 smaller growth rate (e.g., Saito and Maruyama, 2007). This also explains why upwellings  
255 tend to be amplified during the PSSR (Tsunoda, 2015). The physical mechanisms responsible  
256 for the ionospheric layer height variation will be discussed in the next section.

257 Note that the upwellings do not necessarily correspond to the specific phase front of  
258 gravity waves since the upwellings are stationary, but gravity waves are not. Upwellings are  
259 developed via  $E_p \times B$  drift (e.g., Tsunoda, 2015). The various zonal scales of gravity waves  
260 also partly explain why the EPBs occur in isolated regions on some nights (Figure 1a), but on  
261 other nights EPBs display a quasiperiodic wave-train, extending over thousands of kilometers  
262 in the zonal distance (Figure 1c).

263 There are two scenarios that could explain the interplay between gravity waves and  
264 upwellings. The first scenario is under ideal background conditions (e.g., solar maximum,  
265 equinoxes, strong upward drift, higher ionospheric layer height), when weak gravity wave  
266 perturbation is sufficient for the upwelling growth as shown in Figures 2a and 2c due to  
267 higher bottomside ionospheric layer height (i.e., large growth rate). The passage of gravity  
268 waves causes bottomside ionospheric undulations through ion-neutral coupling processes,  
269 leading to inhomogeneity of the Pedersen conductivity. A divergent charge would pile up on  
270 the edges of seed perturbations when eastward Pedersen current driven by gravity or  
271 equatorward neutral winds flow over this region, setting up polarization electric fields ( $E_p$ ) to  
272 satisfy ionospheric current-free conditions ( $\nabla \cdot J = 0$ ). Upwelling or LSWS eventually  
273 develop in the bottomside ionosphere via  $E_p \times B$  drifts.

274 The other scenario is when the background condition does not favor the upwelling  
275 growth (e.g., solar minimum, solstices, weak upward drift, lower ionospheric layer height),  
276 so strong gravity wave perturbations in the neutral wind become critical (e.g., Aa et al., 2022;  
277 Rajesh et al., 2022; Harding et al., 2022). Vertical oscillations of gravity wave-driven neutral  
278 winds can drive zonal divergent Pedersen currents ( $J_{\sim} = U \times B$ ) and  $E_p$  should be established to  
279 cancel the Pedersen currents, leading to upwelling growth (e.g., Eccles, 2004; Tsunoda, 2010;  
280 2021).

281 In Figure 3, we examine temperature perturbations from WACCM-X as a function of  
282 longitude and latitude at  $\sim 350$  km on 20, 22, and 24 October to confirm the presence and  
283 morphology of gravity waves. The morphology of wave patterns is quite complicated, likely  
284 due to the interference of gravity waves from different sources. We found that gravity waves  
285 are ubiquitous and could act as natural seeds for the formation of upwelling, albeit there are  
286 enhanced perturbations at mid-latitudes that may be related to mountain waves or  
287 convectively-generated gravity waves (cf. Ern et al., 2011). There are many sources that  
288 could generate gravity waves, such as deep convection (e.g., Yue et al., 2009), solar  
289 terminator waves (Bespalova et al., 2016), and oceanic waves (Zobotin et al., 2016).  
290 However, more careful studies of these gravity waves are out of scope for this paper. Future  
291 work will focus on analyzing the wave sources related to EPB development.

292 Of particular significance is that Figure 3a shows distinct southwestward propagating  
293 planar gravity waves at the magnetic equator from  $105^\circ$ - $120^\circ$ W. The large-scale zonal wind  
294 perturbations shown in Figure 2d are therefore related to the planar gravity waves. Tsunoda

295 (2010, 2013) and Krall et al. (2013) suggested that planar gravity waves cannot seed EPBs  
296 effectively because the coupling of planar gravity waves to the ionosphere tends to be weak  
297 when the wave phase fronts are not aligned with geomagnetic field lines. Thus, the  
298 alternating contributions of upward and downward winds to the electric potential cancel each  
299 other out along the same field line. In Figures 3b and 3c, multiple concentric waves can be  
300 identified near the magnetic equator. Tsunoda (2010) suggested that concentric gravity waves  
301 can seed EPBs effectively because the polarization response is more efficient when the  
302 wavefront is aligned with the geomagnetic field lines. The discrepancy of planar and  
303 concentric gravity waves could partly explain the longitudinal distribution of EPBs shown in  
304 Figures 1a and 1c. The zonal scale and wavefront orientation of gravity waves therefore  
305 control the spacing between EPBs.

306

### 307 **3.3 Electric Field and Neutral Wind Effects**

308 Gravity wave seeding is crucial for the formation of upwellings, but sufficient  
309 ionospheric layer height is necessary to facilitate upwelling growth and EPB development. In  
310 the nighttime topical ionosphere, the F region plasma dynamics are governed by a complex  
311 interplay between motions of electromagnetic forces, neutral winds, gravity, and pressure  
312 gradient (Kelley, 2009). The equilibrium of the ionosphere is primarily affected by neutral  
313 wind, gravitational and electromagnetic forces since the pressure gradient term produces  
314 negligible effects in the global electrodynamics (Perkins, 1973; Eccles, 2004; Maute et al.,  
315 2012). To understand the background conditions responsible for the day-to-day variability of  
316 EPBs, we examine the effects of  $E \times B$  drift and neutral wind on the ionospheric layer height  
317 variation. In this section, we will first discuss the background conditions related to the  
318 irregular spatial distribution of EPBs on October 20. Then, we will discuss the absence and  
319 regular spatial distribution of EPBs on October 22 and 24, respectively; both cases show  
320 similar initial background conditions at 05:00 UT.

321

#### 322 **3.3.1 Irregular Spatial Distribution of EPBs on October 20**

323 Figure 4 shows the time sequence of electron density (top panels), vertical  $E \times B$  drift  
324 (middle panels), and zonal  $E \times B$  drift (bottom panels) as a function of longitude (local time)  
325 and altitude on October 20. An EPB that occurred after sunset is discernible from  $60^\circ$ - $75^\circ$ W  
326 due to strong PRE vertical drifts after 00:00 UT. Here we focus on the EPBs that developed  
327 after 0500 UT. The PRE-related upward  $E \times B$  drift enhancement is visible around  $100^\circ$ -  
328  $135^\circ$ W below  $\sim 600$  km altitude (Figure 4f). EPBs do not develop following the PRE because  
329 of the weak upward  $E \times B$  drifts ( $\sim 20$  m/s) and lower bottomside ionospheric layer heights  
330 (below 300 km). However, significant localized upward  $E \times B$  drift enhancements of  $\sim 20$ - $50$   
331 m/s occurred around  $80$ - $120^\circ$ W in the topside ionosphere (700-1000 km) after 0500 UT. The  
332 localized upward  $E \times B$  drifts further moved downward and westward and made an additive

333 contribution to the PRE vertical drifts, raising the ionosphere to higher altitudes of ~350 km  
334 (Figure 4c) and contributing to the upwelling growth and large-scale EPB development at  
335 ~90°W at 0530 UT.

336 The localized upward  $E \times B$  drift enhancement causes significant undulations of the  
337 ionospheric layer height, resulting in large zonal and vertical plasma density gradients. Under  
338 such conditions, the large-scale gravity-driven Pedersen current becomes important in  
339 equatorial ionospheric electrodynamics (Eccles, 2004; Maus and Luhr, 2006; Burke et al.,  
340 2009). Eccles (2004) suggested that gravity-driven current is an essential source of large-  
341 scale  $E_p$  ( $\lambda > 1000$  km) during the nighttime ionosphere. As the eastward gravity-driven  
342 Pedersen current flows over the undulating bottomside ionosphere,  $E_p$  will develop and lead  
343 to more prominent ionospheric undulations through  $E_p \times B$  drifts. In our opinion, this explains  
344 the alternating large-scale upward and downward drifts in Figure 4 after 0500 UT. The  
345 presence of large-scale upward  $E \times B$  drifts further leads to the development of upwellings in  
346 the bottomside ionosphere near midnight with small-scale upward  $E \times B$  drifts of ~30-50 m/s  
347 (Figures 4d and 4i), which are superimposed on the large-scale upward  $E \times B$  drifts. We can  
348 identify two upwellings that developed around 110°-120°W at 07:00 UT and EPBs that  
349 developed from the crests of the upwellings after midnight. At 08:00 UT, more pronounced  
350 ionospheric undulations occur because of the contribution of gravity-driven eastward  
351 Pedersen current around 90-110°W (Figures 4e and 4j). Such large ionospheric undulations  
352 extending over ~200-300 km in altitude have been observed by Jicamarca radar (Kelley et al.,  
353 1981). The dynamic vertical  $E \times B$  drifts significantly affect the longitudinal variation of  
354 ionospheric layer height and the longitudinal distribution of EPBs. The distribution of large-  
355 scale upward  $E \times B$  drifts also explains why the EPBs are confined within ~85°-120°W.

356 An additional simulation excluding the gravity-driven current terms in the potential  
357 equation (Huba and Joyce, 2010) has been conducted. The gravity-driven electric current can  
358 contribute additional large-scale vertical  $E \times B$  drifts of ~10-20 m/s during the nighttime (not  
359 shown), consistent with the previous simulations and observations (e.g., Eccles, 2004;  
360 Stoneback et al., 2011). Such midnight upward drift enhancements have been observed by  
361 FORMOSAT-1 and C/NOFS during quiet time conditions (e.g., Yizengaw et al., 2009;  
362 Heelis et al., 2010; Stoneback et al., 2011).

363 The bottom panels of Figure 4 show that the corresponding zonal  $E \times B$  drifts display  
364 strong vertical shear flow with plasma moving eastward at up to ~150 m/s above 300 km and  
365 westward at up to ~150 m/s below 300 km, consistent with the NASA sounding rocket  
366 experiments during the postsunset equatorial ionosphere (Hysell et al., 2005). We noticed  
367 that clear localized retrograde flow (westward  $E \times B$  drifts) embedded within the F region  
368 eastward plasma flow emerged in the topside ionosphere at 05:00 UT (Figure 4k), and the  
369 movement of the retrograde flow is accompanied by the localized upward  $E \times B$  drift  
370 enhancement (Figures 4f and 4g). The retrograde flow keeps moving westward and

371 downward, eventually encountering the westward flow below 300 km at 06:00 UT, which in  
372 turn elevates the westward flow to a higher altitude of  $\sim 400$  km, destabilizing the bottomside  
373 ionosphere and resulting in shear instability (Figures 4m and 4n). The shear flow  
374 accompanies the upward and downward  $E \times B$  drifts between  $60^\circ W$  and  $105^\circ W$ , displaying a  
375 midnight equatorial vortex feature that is commonly observed in the post-sunset periods (e.g.,  
376 Kudeki and Bhattacharyya, 1999; Lee et al., 2014). These processes allow upwelling growth  
377 and EPB development, supporting the hypothesis proposed by Hysell and Kudeki (2004) that  
378 the shear instability could precondition the ionosphere for the RT instability. Considering that  
379 the nighttime eastward plasma flow in the F region is related to the eastward wind (cf. Heelis,  
380 2004), the retrograde flow could be related to the F region westward wind. To confirm this  
381 hypothesis, we further examine neutral winds since they affect the ionospheric layer height  
382 and electrodynamics (e.g., Heelis, 2004; Lin et al., 2007).

383 Figure 5 shows the time sequence of meridional (top) and zonal (bottom) winds at  $\sim 460$   
384 km on October 20. The cross-equatorial meridional winds are mainly northward around  
385  $\sim 75^\circ$ - $120^\circ W$  over the magnetic equator at 05:00 UT. However, we note that the meridional  
386 winds display distinct band structures with alternating wind directions. The wind patterns  
387 tend to move southwestward, extending from northwest to southeast. The meridional winds  
388 between  $\sim 95^\circ$ - $135^\circ W$  in the northern hemisphere are mainly northward at 05:00 UT and  
389 gradually turn southward, leading to a converging wind pattern between  $\sim 90^\circ$ - $135^\circ W$  over  
390 the magnetic equator at  $\sim 06:00$  UT. Converging winds can facilitate the RT instability  
391 because the converging winds can raise the ionosphere to higher altitude along the field line,  
392 leading to a decrease of integrated Pedersen conductivity and ion-neutral collision frequency  
393 (Huba and Krall, 2013). The downward component of converging winds is also an additive  
394 driver for the RT instability (Tsunoda, 2021) because the downward wind can drive eastward  
395 Pedersen current contributing to the RT instability, similar to the gravity-driven eastward  
396 electric current.

397 We note that the meridional winds gradually turn to a poleward direction between  $60^\circ$ -  
398  $90^\circ W$  after 05:00 UT (after 24:00 LT at  $75^\circ W$ ), exhibiting a typical midnight temperature  
399 maximum (MTM) wind pattern over the magnetic equator (c.f., Fang et al., 2016). The  
400 occurrence of MTM could result in localized reversal of the large-scale eastward and  
401 equatorial winds during the nighttime. This can be seen in Figures 5f-j, in that the eastward  
402 zonal wind over the MTM slows down and reverses to westward. The poleward winds further  
403 lower the ionosphere and weaken RT instability between  $\sim 60^\circ$ - $90^\circ W$  (cf., Huba and Krall,  
404 2013). The downward motion of the ionosphere from  $60^\circ$ - $90^\circ W$  is, therefore, due to the  
405 combined effects of the southward wind, downward  $E \times B$  drifts, and MTM winds. The  
406 distribution of meridional winds generally reflects the longitudinal variation of ionospheric  
407 layer height (top panel of Figure 4), demonstrating that the meridional wind is another  
408 important factor in determining the longitudinal distribution of EPBs.

409       Of particular significance is that the meridional winds display a blue narrow band  
410 structure (southward wind) accompanying with large-scale northward winds extending from  
411 northwest to southeast in the northern hemisphere, which appears to be related to the planar  
412 gravity waves (Figure 3a) and solar terminator waves. We will discuss the large-scale solar  
413 terminator waves in section 3.4. The narrow band wind structure also can be identified in the  
414 zonal wind, in that the presence of a westward wind causes cessation of the eastward wind.  
415 Compared with the vertical and zonal  $E \times B$  drifts (Figure 4), the retrograde flow and upward  
416  $E \times B$  drift enhancement in the equatorial F region can be attributed to this narrow band wind  
417 structure since the orientation of retrograde flow is consistent with the narrow band wind  
418 structure. Since the neutral wind perturbations can result in inhomogeneous electric  
419 conductivity distribution, the divergence and convergence of zonal wind driven dynamo  
420 currents cause accumulation of electric charges. The westward winds play a vital role on the  
421 development of retrograde flow and upward  $E \times B$  drift over the magnetic equator by  
422 generating  $E_p$  mapping to the magnetic equator, contrary to the post-sunset ionospheric  
423 conditions where the vertical drift is primarily driven by the eastward acceleration of zonal  
424 wind (Richmond et al., 2015).

425       Moreover, westward tilting eastward  $E \times B$  drift enhancements are also visible on either  
426 side of the retrograde flow, which could be modulated by the downward  $E_p$  generated by  
427 eastward winds at higher latitudes. Varney et al. (2009) observed the streak patterns in zonal  
428 and vertical  $E \times B$  drifts related to gravity waves in Jicamarca ion drift observations. This  
429 demonstrates that gravity waves can be the another source to drive midnight vertical drift and  
430 shear flow instability. Miller et al. (2009) reported the seeding of EPBs by the mid-latitude  
431 MSTIDs. They found that the  $E_p$  embedded within the MSTIDs can be mapped to the  
432 magnetic equator along the magnetic field lines (e.g., Chou et al., 2021) and lead to the post-  
433 midnight EPB development. Their MSTIDs displayed distinct westward-tilted band  
434 structures when the airglow images were projected to the Apex coordinate. Significant  
435 westward plasma flows embedded within the MSTIDs were also identified, consistent with  
436 our simulations.

437

### 438 **3.3.2 Absence of EPBs on October 22**

439       Figure 6 is the same as Figure 4, but for October 22. The PRE is small or absent near  
440 dusk due to weak eastward wind. The vertical  $E \times B$  drift generally shows typical downward  
441 drift throughout the night (e.g., Scherliess and Fejer, 1999), resulting in lower ionospheric  
442 layer heights. At 05:00-07:00 UT, there are gentle large-scale bottomside ionospheric  
443 undulations accompanied by nearly zero vertical drift above the crests of upwellings due to  
444 the cancellation of small-scale upward drifts and large-scale downward drifts (Figures 6f and  
445 6g). Significant downward drifts above the downwelling are also visible from 60°-90°W.  
446 Similar patterns of decreasing and increasing zonal  $E \times B$  drifts can also be identified from

447 75°-90°W (Figures 6k and 6i), most likely due to the zonal wind variations. The cessation of  
448 eastward  $E \times B$  drifts due to the westward-tilted retrograde flow in the F region can also be  
449 identified, and the corresponding upward  $E \times B$  drifts are weak as well. After 07:00 UT, there  
450 are significant westward  $E \times B$  drifts around 60°-105°W; however, the westward flow is  
451 accompanied by downward  $E \times B$  drifts.

452 Figure 7 shows the corresponding meridional and zonal wind variations on October 22.  
453 The meridional wind also displays distinct band structures, although wind velocities are  
454 weaker than in Figure 5. Weak converging winds are not able to push the ionosphere to  
455 sufficient altitude. A blue narrow band structure (southward wind) extends from northwest to  
456 southeast in the northern hemisphere (top panels). The narrow band wind structure can also  
457 be identified in the zonal wind around 75°-120°W in the northern hemisphere at 05:00-06:00  
458 UT (bottom panels), which is responsible for the west-tilted retrograde flow and upward  $E \times B$   
459 drifts in Figure 6. After 06:00 UT, the westward winds around 60°W in the southern  
460 hemisphere correspond to the large-scale westward plasma flow (bottom panels of Figure 6).  
461 Although the westward winds also display northwest to southeast alignment, the westward  
462 winds in the southern hemisphere can induce northward and downward Pedersen currents.  
463 Southward and westward  $E_p$  would be set up to drive westward and downward  $E_p \times B$  drifts,  
464 resulting in descending ionospheric layers. Together these processes explain the absence of  
465 EPBs on October 22.

466

### 467 3.3.3 Regular Spatial Distribution of EPBs on October 24

468 Figure 8 is the same as Figure 6, but for October 24. The PRE is also weak and no EPBs  
469 develop during the post-sunset period. At ~05:00 UT, Figure 8a shows that the ionospheric  
470 layer heights are comparable with the layer heights on October 22, the no-EPB case.  
471 However, large-scale upward drifts develop near midnight and extend from ~60°W to  
472 ~135°W (Figure 8f), uplifting the bottomside ionospheric layer by at least 50 km. We found  
473 that the large-scale upward  $E \times B$  drifts tilted westward at 05:00 UT, consistent with the  
474 morphology of westward retrograde flow in the F region (Figure 8k). At 06:00 UT, the  
475 retrograde flow in the F region merges with the westward  $E \times B$  drifts in the bottomside  
476 ionosphere, leading to shear instability and equatorial vortex features, consistent with the  
477 October 20 case. Two upwellings are further developed around 70°-80°W at 06:00 UT  
478 (Figure 8b). The large-scale upward  $E \times B$  drifts continue moving westward and lifting the  
479 ionosphere to ~300 km, leading to successive upwelling growth and EPB development.

480 Figure 9 shows that the meridional winds on October 24 have similar wind patterns and  
481 velocities compared with the meridional winds on October 22. However, zonal wind  
482 disparities exist, as discussed in more detail here. Much stronger westward winds can be  
483 identified over the magnetic equator during 05:00-07:00 UT on October 24 (Figures 9f-9h).  
484 Such westward winds are responsible for the retrograde flow and upward  $E \times B$  drift

485 enhancement in Figure 8, demonstrating that the zonal wind differences are the primary  
486 reason for the absence of EPBs on one day and the presence of EPBs on the other day.

487

### 488 **3.3.4 Mapping of Electric Fields Induced by Neutral Wind Perturbations**

489 The most striking discovery is that the neutral wind perturbations driven by solar  
490 terminator waves and gravity waves contribute to the midnight vertical drift enhancement and  
491 collisional shear instability, as discussed in section 3.3.1-3.3.3. To investigate the linkage  
492 between neutral wind perturbations, upward  $E \times B$  drift enhancement, and retrograde flow,  
493 Figure 10 shows the zonal (top panels) and vertical/meridional (bottom panels)  $E \times B$  drifts as  
494 a function of latitude and altitude at ~05:00 UT along the magnetic longitudes of 337.65°,  
495 346.05°, and 357.45° (~95.3°W, ~86.86°W and ~75.25°W in the geographic coordinate at the  
496 magnetic equator) on October 20, 22, and 24, respectively. Significant cessation of eastward  
497  $E \times B$  drift and westward  $E \times B$  drift related to retrograde flows (indicated by black arrows) can  
498 be identified along the field lines, accompanying with the upward  $E \times B$  drift enhancement, on  
499 October 20 and 24, respectively. However, the retrograde flow and upward  $E \times B$  drift are  
500 obscure on October 22. The retrograde flows extend from ~10°N and ~5°N to the southern  
501 hemisphere on October 20 and 24, consistent with the locations of blue narrow band  
502 structures of neutral winds discussed in previous sections. This reveals that the solar  
503 terminator waves and gravity waves are responsible for the retrograde flow and midnight  
504 vertical drift enhancement. Thus, the altitudinal variation in equatorial F region plasma  
505 motion is a direct mapping of the latitudinal variation of  $E_p$  generated by solar terminator  
506 wave and gravity wave induced neutral wind perturbations (e.g., Huba et al., 2015; Chou et  
507 al., 2022; Lin et al., 2022).

508 As opposed to the post-sunset ionosphere, the midnight upward  $E \times B$  drifts are attributed  
509 to the westward winds associated with the solar terminator waves and large-scale gravity-  
510 driven currents (e.g., Eccles., 2004). Figure 11 illustrates the mechanism of westward wind  
511 on the midnight upward  $E \times B$  drift and retrograde flow. The westward winds with northwest  
512 to southeast band structure at higher latitudes can drive southward Pedersen current ( $J_p \sim$   
513  $U \times B$ ), setting up northward  $E_p$ . Due to the specific wavefront alignment of the westward  
514 wind, the eastward component of  $E_p$  can also be established. Both northward and eastward  
515  $E_p$  will map along field lines to the magnetic equator due to high electrical conductivity of  
516 the field line, leading to westward and upward  $E_p \times B$  drifts in the topside equatorial  
517 ionosphere. Gravitational force further amplifies the  $E_p$  by generating eastward Pedersen  
518 current. The westward and upward  $E_p \times B$  drifts result in retrograde flow and midnight  
519 vertical drift enhancement. As the band structure of the westward wind moves southwestward,  
520 the retrograde flow moves downward and westward and merges with the westward plasma  
521 flow in the bottomside ionosphere, leading to collisional shear instability. Hysell and Kudeki  
522 (2004) and Hysell et al. (2005) suggested that the shear instability may destabilize the

523 bottomside ionosphere and generate precursor seed waves responsible for the EPB  
524 development. Our simulations suggest that the shear instability could be influenced by the  
525 retrograde flow associated with solar terminator wave or gravity waves. Additionally, Coley  
526 et al. (2014) showed westward plasma flow in the topside ionosphere after midnight during  
527 solar minimum. Forbes et al. (2008) also suggested that solar terminator waves are more  
528 prominent during solar minimum, implying that the westward plasma flow observed by  
529 Coley et al. (2014) may be related to solar terminator waves or gravity waves.

530 However, this scenario is only valid in the northern hemisphere. As the band structure of  
531 the westward wind moves to the southern hemisphere, the westward wind would induce  
532 downward  $E \times B$  drift, as shown in Figures 6 and 7, due to the direction of magnetic field.  
533 Under such condition, the eastward wind with the same wavefront orientation in the southern  
534 hemisphere could induce downward and eastward  $E_p$  mapping to the magnetic equator. This  
535 demonstrates that zero or weak zonal  $E_p \times B$  drift may occur due to the cancellation of vertical  
536  $E_p$ ; however, the upward  $E_p \times B$  drift should persist.

537 Figure 12 summarizes the coupled physical processes contributing to the midnight EPB  
538 development. Neutral wind perturbations associated with gravity waves and solar terminator  
539 waves destabilize the ionosphere by generating the midnight vertical drift enhancement and  
540 shear flow instability, resulting in the midnight vortex. The midnight vortex is therefore  
541 related to gravity waves or solar terminator waves, which is different from the post-sunset  
542 vortex associated with the PRE (e.g., Tsunoda et al., 1981). The midnight vertical drift  
543 enhancement and converging winds associated with solar terminator waves further uplift the  
544 ionosphere, and gravity waves seed upwellings. The zonal scale of gravity waves determines  
545 the spacing between upwellings (or EPBs).  $E_p$  developed within upwellings further leads to  
546 upwelling growth via  $E_p \times B$  drift and EPB development. The study reveals that gravity waves  
547 can not only contribute to seeding but create ionospheric conditions resembling the post-  
548 sunset ionosphere to facilitate EPB development.

549

### 550 **3.4 Influences of Solar Terminator Waves on the Midnight EPB Development**

551 In section 3.3, we demonstrated that the nighttime neutral wind displays distinct band  
552 structures with alternating wind directions. The converging winds and wind dynamo effect  
553 contribute to the midnight EPB development by lifting the ionosphere to higher altitude,  
554 providing conditions favorable for upwelling growth. Since the dayside solar heating and  
555 pressure bulge cannot explain the alternating wind patterns on the nightside, we propose that  
556 solar terminator waves could be the primary mechanism to explain the alternating band  
557 structures in neutral winds.

558 Figure 13 shows the global distribution of meridional winds at ~460 km altitude on 20,  
559 22, and 24 October. The nighttime meridional winds (shaded area) display large-scale wind  
560 perturbations with northwest to southeast alignment (indicated by dashed lines) after the dusk

561 solar terminator and southwest to northeast alignment before the dawn solar terminator,  
562 consistent with the solstice solar terminator waves in thermosphere winds and densities  
563 observed by the CHAMP satellite (Forbes et al., 2008; Liu et al., 2009). These large-scale  
564 wind patterns move westward with the solar terminator and can be identified daily, despite  
565 the morphology and amplitude being slightly different. Medium-scale meridional wind  
566 perturbations following the dusk terminator can also be identified over the continent of  
567 Eurasia, which could cause post-sunrise ionospheric perturbations (e.g., Zhang et al., 2021).  
568 It should be mentioned that the blue narrow band structures (southward wind) indicated by  
569 white dashed lines have smaller horizontal wavelength ( $\sim 10^\circ$  in latitude) on October 20  
570 compared with other days ( $\sim 20^\circ$  in latitude), which could be due to the modulation of planar  
571 gravity waves shown in Figure 3a.

572 Huang et al. (2014) showed significant longitudinal asymmetry of midnight EPBs around  
573  $60^\circ$ - $150^\circ$ W in October-January. This could be because the wavefront orientation of solar  
574 terminator waves is aligned with the inclination angle of the magnetic equator in this region.  
575 The resulting converging winds and zonal wind dynamo can lift the ionosphere to a higher  
576 altitude, providing conditions favorable for EPB development. We could expect that the solar  
577 terminator waves with southwest to northeast wavefront alignment near June solstice could  
578 also contribute to the high occurrence of midnight EPBs around  $20^\circ$ - $70^\circ$ W near June solstice  
579 (Gentile et al., 2011; Yizengaw et al., 2013).

580 On the other hand, the solar terminator waves on October 20 show more dynamic  
581 features. Miyoshi et al. (2009) suggested that the solar terminator wave is mainly generated  
582 by the superposition of the upward propagating migrating tides, which could contribute to the  
583 generation of MTM as shown in Figure 5. The MTM winds considerably impact EPB  
584 development (Krall et al., 2021) and its longitudinal distribution. McDonald et al. (2015)  
585 indicated that the nonmigrating tides play an important role in the nighttime ion upward drift.  
586 Tidal forcing contributes to the longitudinal distribution of EPBs (Dao et al., 2011; Chang et  
587 al., 2021; Chou et al., 2020b). It appears that the behavior of the diurnal tides in the neutral  
588 winds significantly affects the day-to-day variability of EPBs. Future investigation of the  
589 tidal forcing should advance our understanding of how atmospheric tides control the day-to-  
590 day variability of EPBs.

591

#### 592 **4. Conclusions**

593 We have investigated the day-to-day variability of EPBs using the coupled  
594 SAMI3/WACCM-X model. Simulations reveal that EPBs developed on October 20 and 24  
595 but not on October 22. We found that EPBs developed at midnight. Atmospheric gravity  
596 waves and solar terminator waves are critical to midnight EPB development. They  
597 significantly affect the neutral winds and electrodynamic, which could be responsible for the

598 day-to-day variability of midnight EPBs. The main findings of the present work are  
599 summarized as follows:

- 600 1. We found that gravity waves appear ubiquitous and could act as a natural seed for  
601 EPB development. The spacing between bottomside upwellings is consistent with the  
602 zonal scale of gravity wave perturbations in the zonal winds, suggesting that  
603 upwellings are related to gravity waves. However, upwelling growth requires  
604 sufficient ionospheric layer height, and the longitudinal variation of neutral wind  
605 becomes important.
- 606 2. Gravity waves do not necessarily lead to a quasiperiodic distribution of EPBs (e.g.,  
607 Makela et al., 2010), depending on their zonal scale and wavefront alignment. The  
608 equatorward propagation of planar gravity waves causes irregular or isolated spatial  
609 distributions of EPBs, which partly explain why EPBs show isolated clusters  
610 separated by long distances on some nights but display a continuous distribution of  
611 EPB trains on other nights. The other reason is the longitudinal variation of  
612 ionospheric layer height due to meridional winds.
- 613 3. The longitudinal variation of meridional winds can affect the longitudinal variation  
614 of ionospheric layer height, which in turn controls the occurrence and longitudinal  
615 distribution of EPBs. We found that the converging winds associated with solar  
616 terminator waves along the magnetic equator can lead to a continuous distribution of  
617 EPBs spanning a large zonal distance. On the contrary, diverging winds due to MTM  
618 over the magnetic equator could lower the ionosphere, inhibit upwelling growth and  
619 result in irregular spatial distribution of EPBs.
- 620 4. As opposed to the post-sunset upward vertical drift enhancement due to the eastward  
621 wind, the westward winds associated with the dusk solar terminator waves or gravity  
622 waves play a vital role in the midnight upward drift enhancement and retrograde  
623 plasma flow. Both upward  $E \times B$  drift and retrograde flow result in midnight vortex  
624 features, providing conditions favorable for upwelling growth and EPB development.
- 625 5. We found that solar terminator waves and/or gravity waves can be responsible not  
626 only for the seeding mechanism (e.g., Kelley et al., 1981), but also for collisional  
627 shear instability (Hysell and Kudeki, 2004). The dusk solar terminator waves (or  
628 gravity waves) generate a localized retrograde flow, merging with westward plasma  
629 flow in the bottomside ionosphere and leading to shear instability. The solar  
630 terminator waves and gravity waves also contribute to the large-scale EPB  
631 development (Figure 1a).
- 632 6. We found that the presence or absence of midnight EPBs connects to the westward  
633 winds driven by dusk solar terminator waves. Weak (or cessation of) westward winds  
634 prevented formation of midnight EPBs on October 20 because of weak upward  $E \times B$

635 drifts and retrograde flow, which in turn lead to lower ionospheric layer height and  
636 small growth rate.

637 This study provides a new perspective that the day-to-day variability of EPBs is  
638 significantly affected by the neutral wind perturbations driven by atmospheric waves.  
639 Therefore, ion drift, neutral wind, and atmospheric waves measurements in both hemispheres  
640 are helpful to the nowcasting of EPBs. Building a data assimilation system by incorporating  
641 state-of-the-art thermosphere data, such as ICON, into WACCM-X should improve the  
642 global neutral wind specification, advancing the capability of EPB nowcasting and  
643 forecasting (e.g., Hsu et al., 2021).

644

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#### 651 **Open Research**

652 Data Availability Statement

653 The SAMI3/WACCM-X files are available at <https://doi.org/10.5281/zenodo.7425959>.

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955 **Figure Captions:**

956

957 Figure 1. TEC maps from the SAMI3/WACCM-X simulation at 08:00, 08:00, and 10:00 UT  
958 on October 20, 22, and 24, respectively, in 2020. Clear dark band structures related to EPBs  
959 can be identified on October 20 and 24 with irregular and regular spatial distribution,  
960 respectively. EPBs are confined within 63.6°-136.5°W. The white line indicates the magnetic  
961 equator. The dashed lines indicate the longitudes from 30° -180°W with 30° interval.

962

963 Figure 2. The electron density (top panels, log scale) and zonal wind perturbations extracted  
964 by high-pass filter (bottom panels) as a function of longitude and altitude along the magnetic  
965 equator at 07:05, 07:00, and 07:00 UT on October 20, 22, and 24 in 2020. The black arrows  
966 denote the locations of upwellings and EPBs.

967

968 Figure 3. The neutral temperature perturbations from the WACCM-X simulation on October  
969 20 (left panel), 22 (middle panel), and 24 (right panel). The green line indicates the magnetic  
970 equator.

971

972 Figure 4. The electron density (top panels, log scale), vertical  $E \times B$  drift (middle panels) and  
973 zonal  $E \times B$  drift (bottom panels) from the SAMI3/WACCM-X simulation as a function of  
974 longitude (local time) and altitude along the magnetic equator at 05:00, 05:30, 06:00, 07:00  
975 and 08:00 UT on October 20, 2020.

976

977 Figure 5. Meridional (top panels) and zonal (bottom panels) winds from the WACCM-X  
978 simulation as a function of longitude (local time) and latitude at ~460 km at 05:00, 05:30,  
979 06:00, 07:00 and 08:00 UT on October 20, 2020. The white lines denote the magnetic  
980 latitudes at 0° and  $\pm 25^\circ$ .

981

982 Figure 6. The electron density (top panels, log scale), vertical  $E \times B$  drift (middle panels) and  
983 zonal  $E \times B$  drift (bottom panels) from the SAMI3/WACCM-X simulations as a function of  
984 longitude (local time) and altitude along the magnetic equator at 05:00, 06:00, 07:00, 08:00  
985 and 09:00 UT on October 22, 2020.

986

987 Figure 7. Meridional (top panels) and zonal (bottom panels) winds from the WACCM-X  
988 simulations as a function of longitude (local time) and latitude at ~460 km at 05:00, 06:00,  
989 07:00, 08:00 and 09:00 UT on October 22, 2020. The white lines denote the magnetic  
990 latitudes at 0° and  $\pm 25^\circ$ .

991

992 Figure 8. Same as Figure 6, but for October 24, 2020.

993 Figure 9. Same as Figure 7, but for October 24, 2020.

994

995 Figure 10. Zonal (top panels) and vertical/meridional  $E \times B$  drifts (bottom panels) as a  
996 function of latitude and altitude along the magnetic longitudes of  $337.65^\circ$ ,  $346.05^\circ$ , and  
997  $357.45^\circ$  ( $\sim 95.3^\circ W$ ,  $86.86^\circ W$ , and  $75.25^\circ W$  at the magnetic equator) at 05:00 UT on October  
998 20, 22, and 24, respectively.

999

1000 Figure 11. Schematic of upward and westward (retrograde flow)  $E \times B$  drifts generated by  
1001 westward wind associated with solar terminator wave/gravity waves in the northern  
1002 hemisphere.

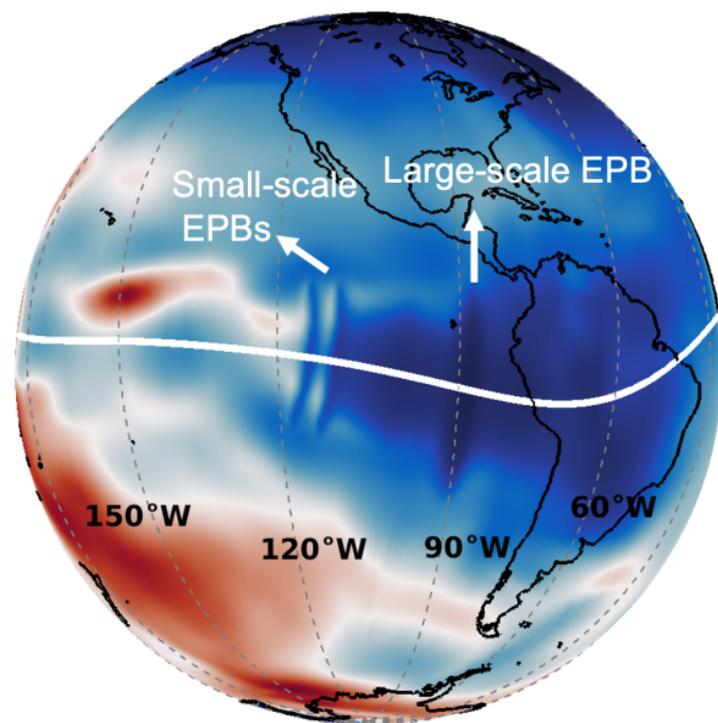
1003

1004 Figure 12. A schematic representation of the coupled processes controlling the midnight EPB  
1005 development.

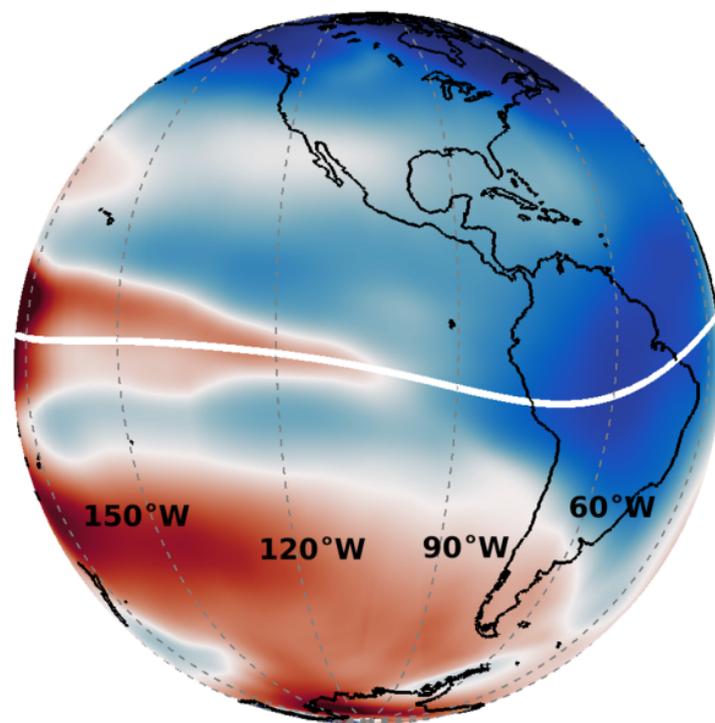
1006

1007 Figure 13. Meridional winds at  $\sim 460$  km altitude from the WACCM-X simulation at 05:00  
1008 UT on October 20 (top panel), 22 (middle panel), and 24 (bottom panel), 2020. The white  
1009 line indicates the magnetic equator. The shaded area represents dawn-dusk solar terminators.  
1010 The dashed lines indicate the wavefronts of solar terminator waves.

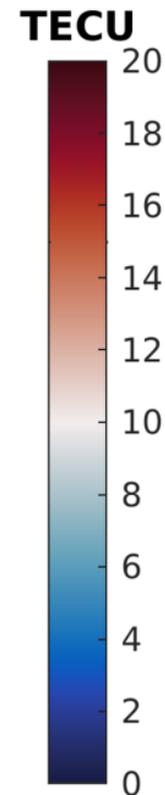
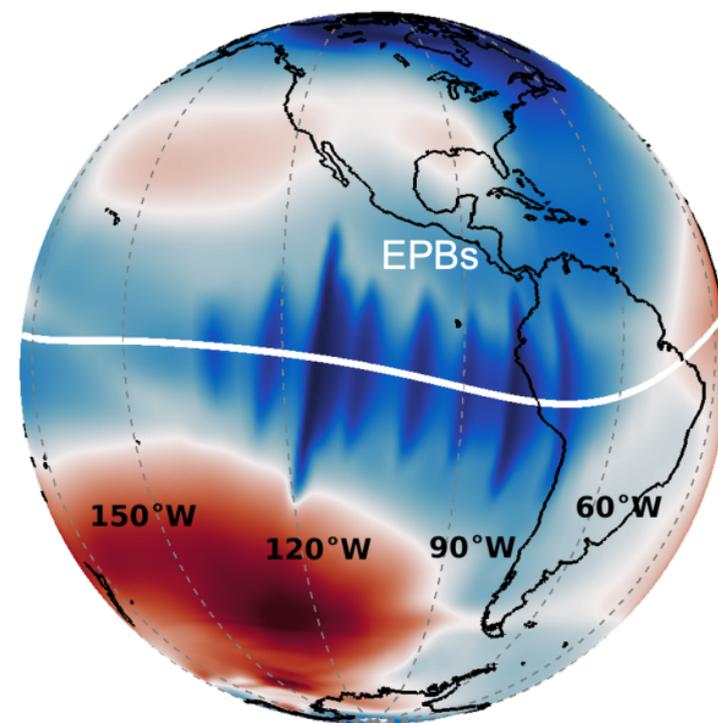
(a) October 20, 2020 08:00:00 UT

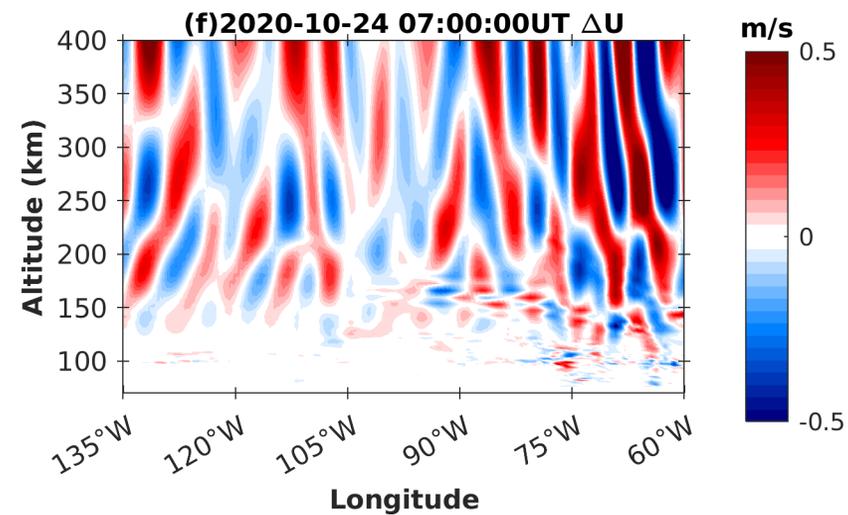
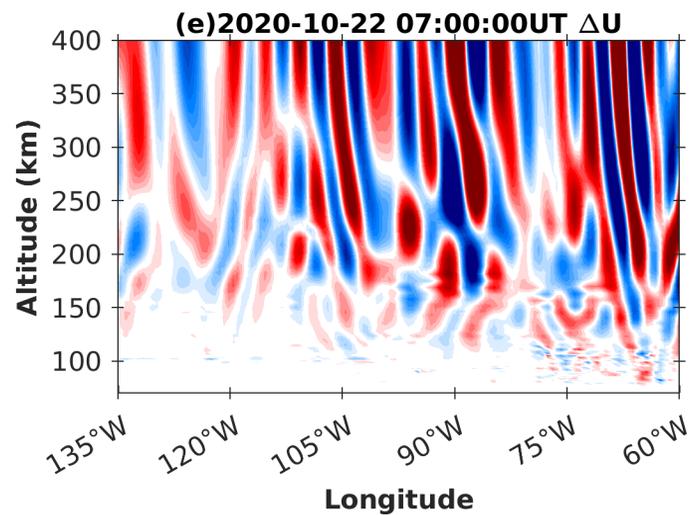
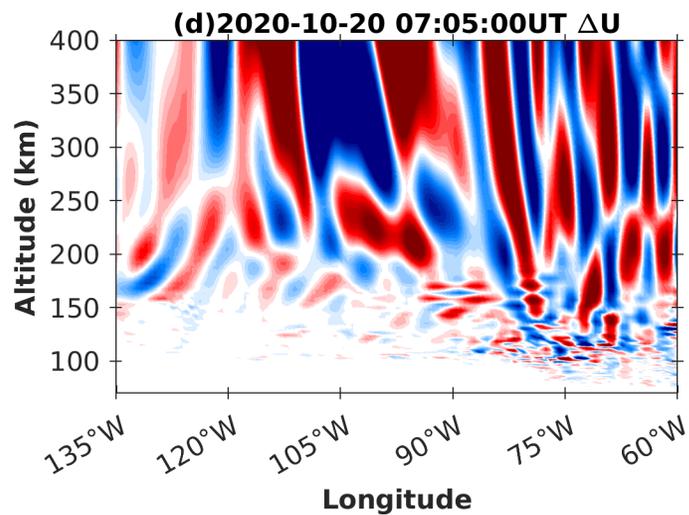
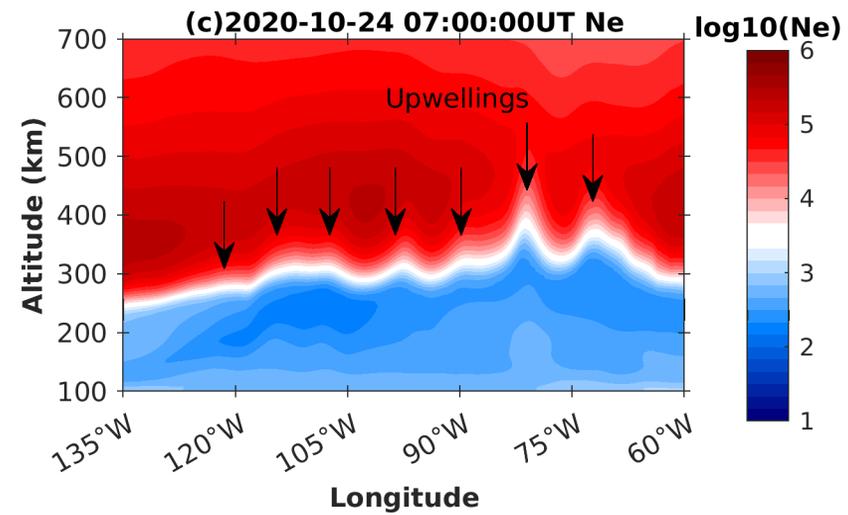
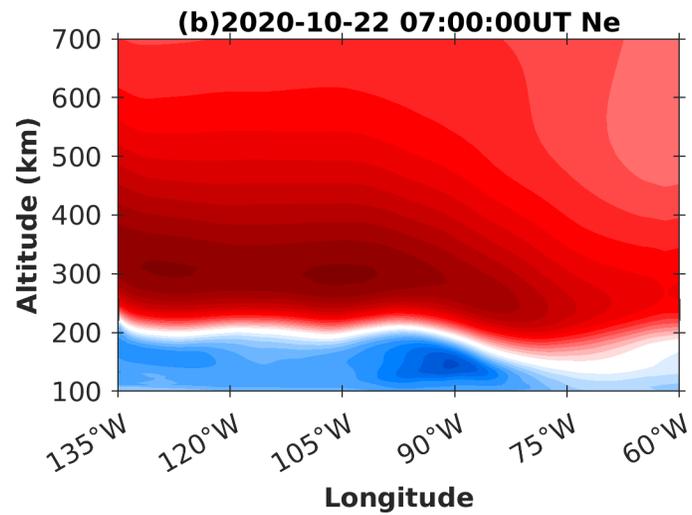
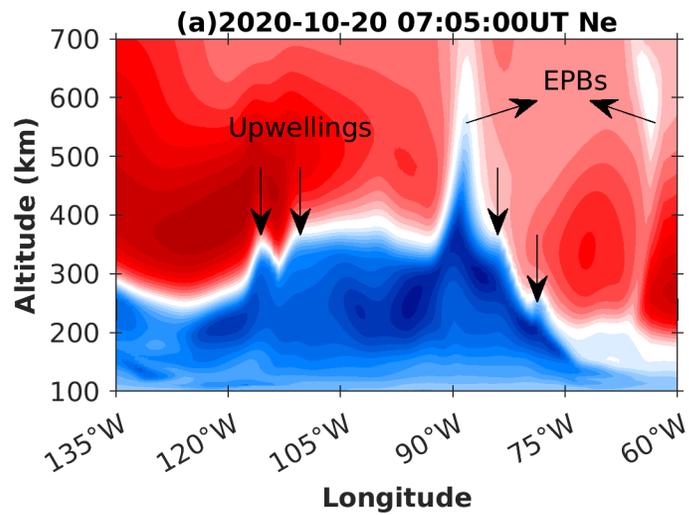


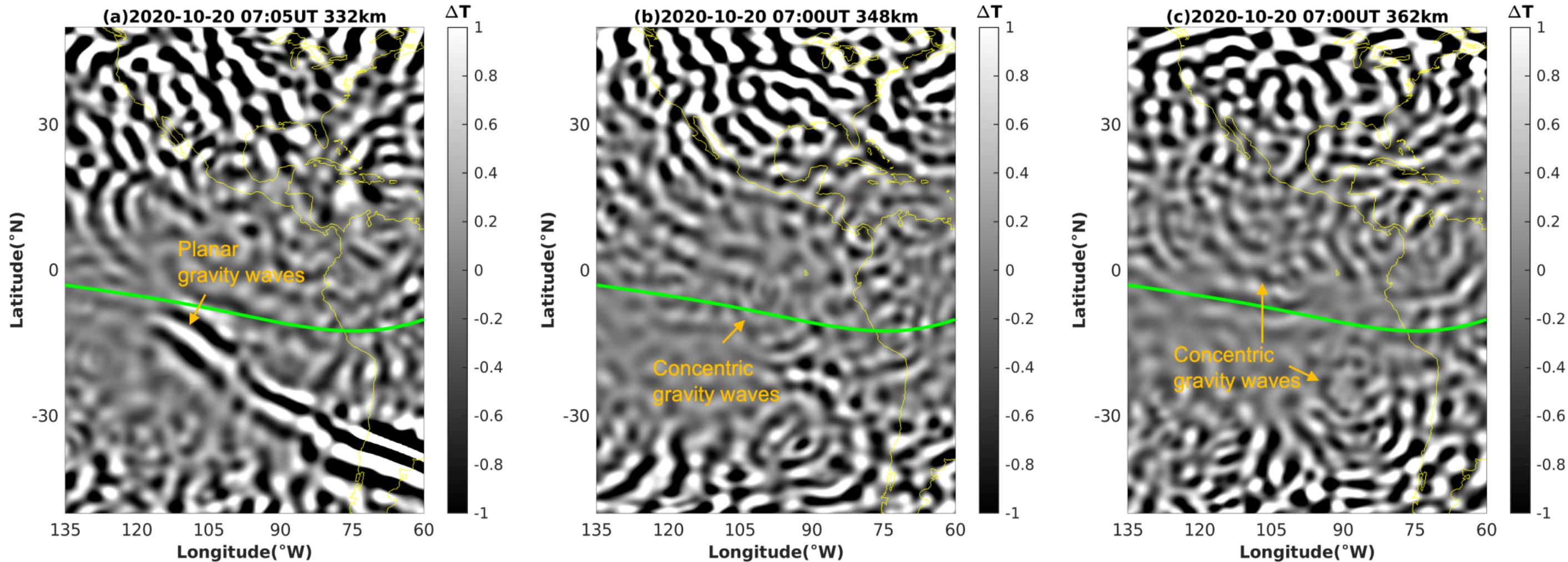
(b) October 22, 2020 08:00:00 UT

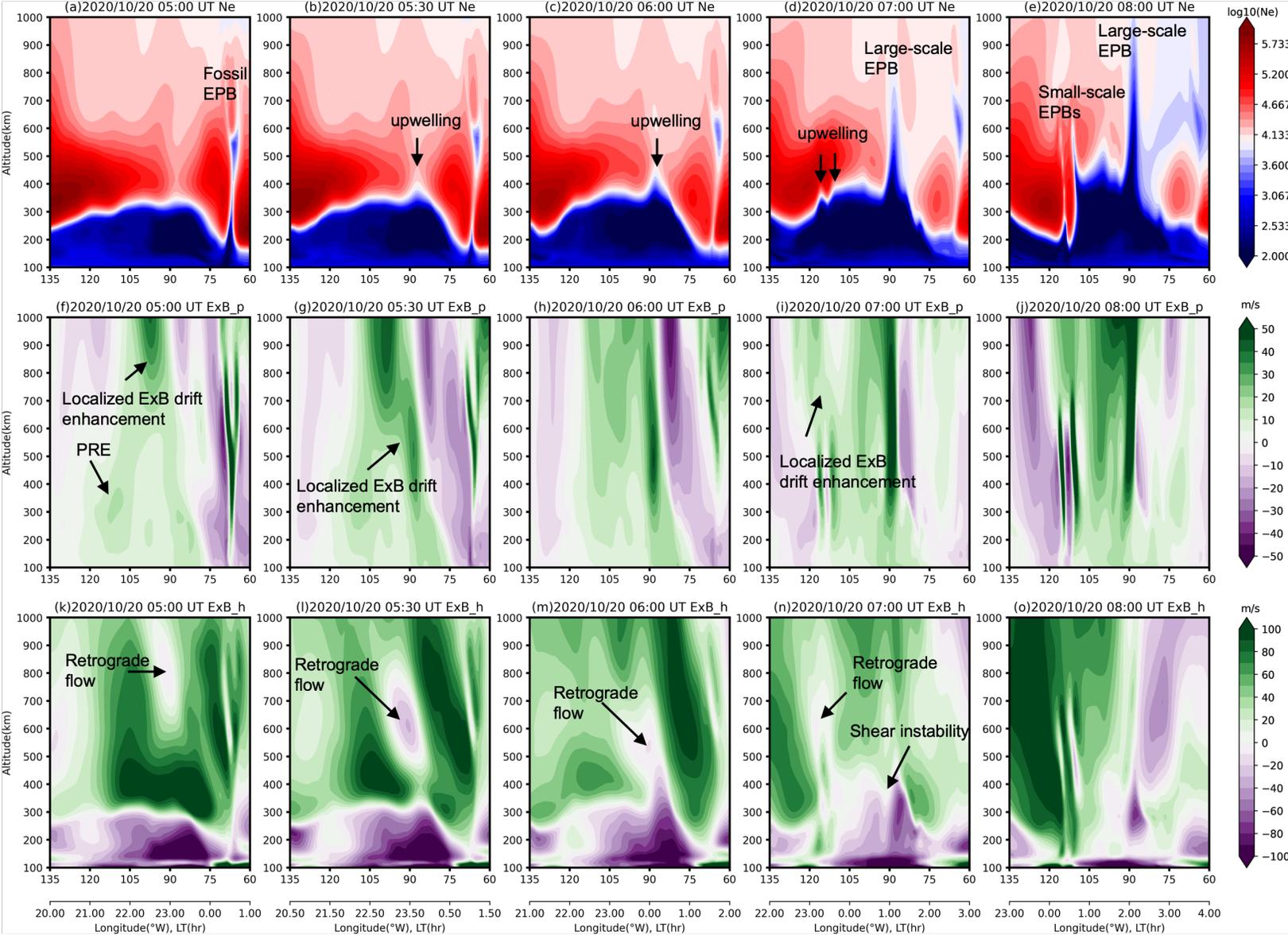


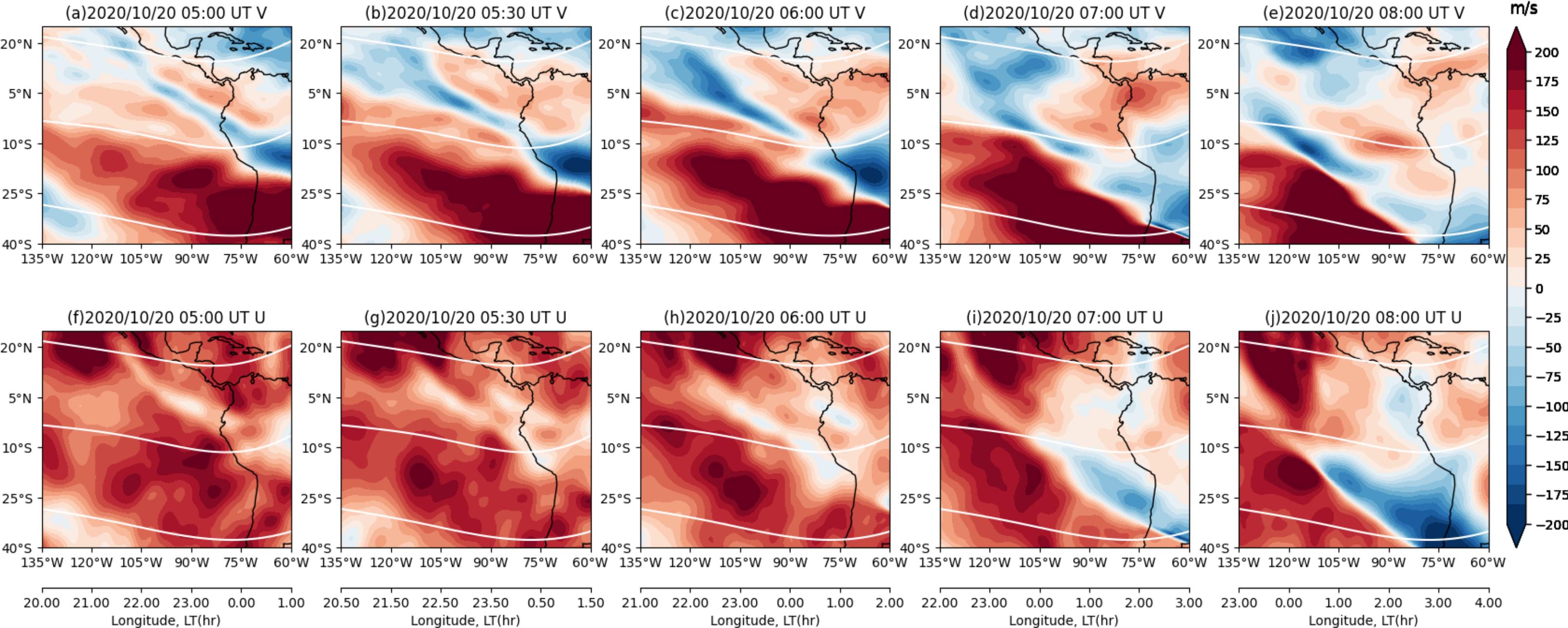
(c) October 24, 2020 10:00:00 UT

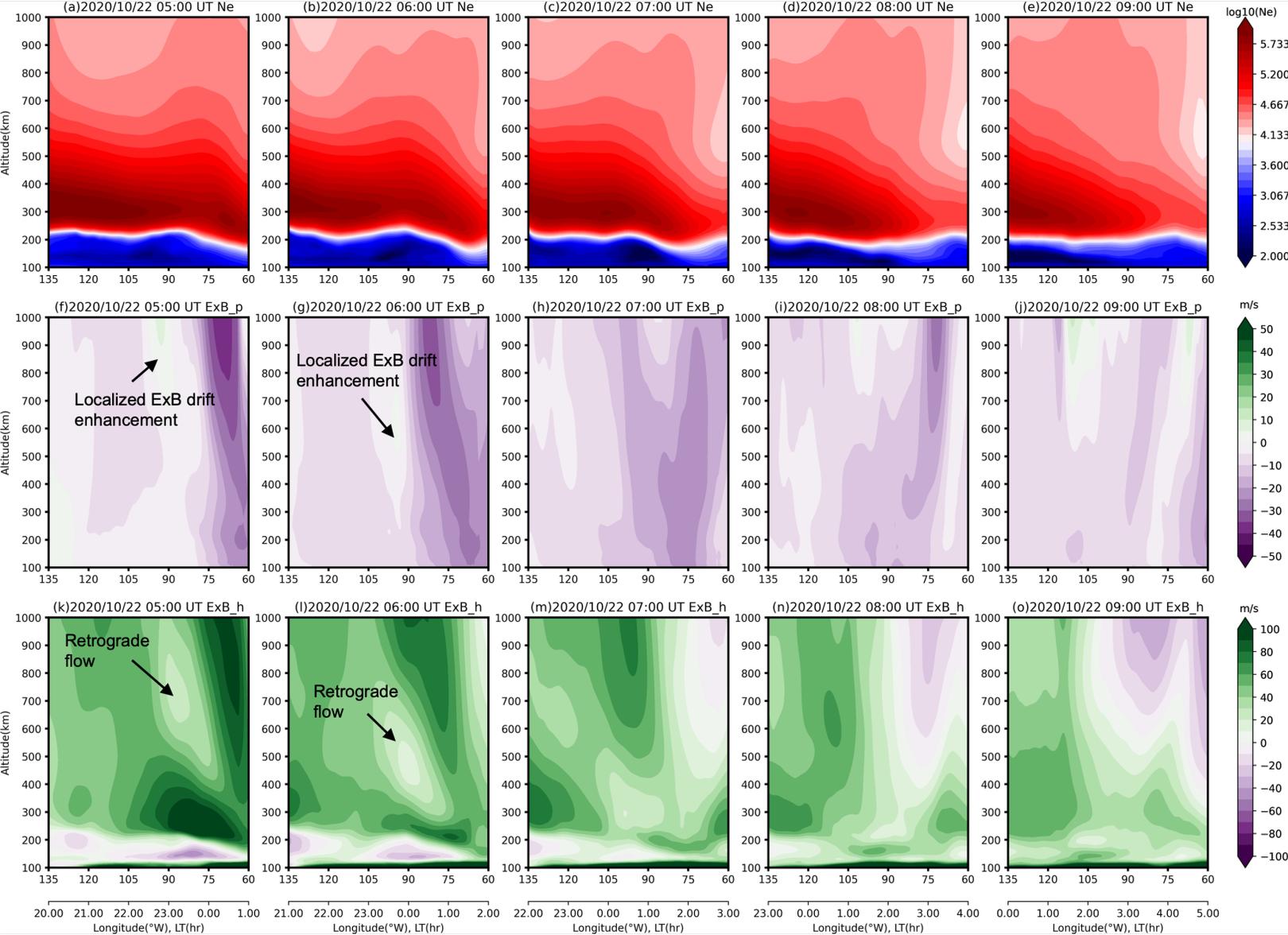


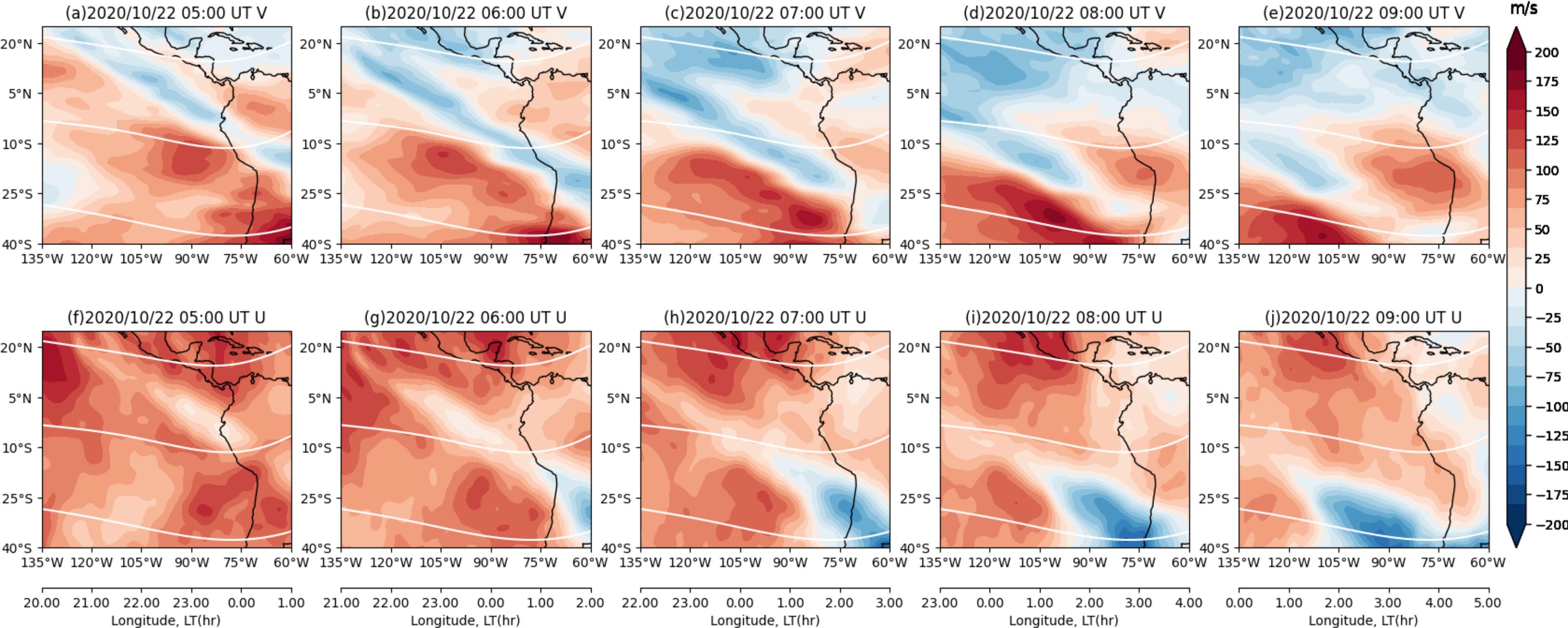


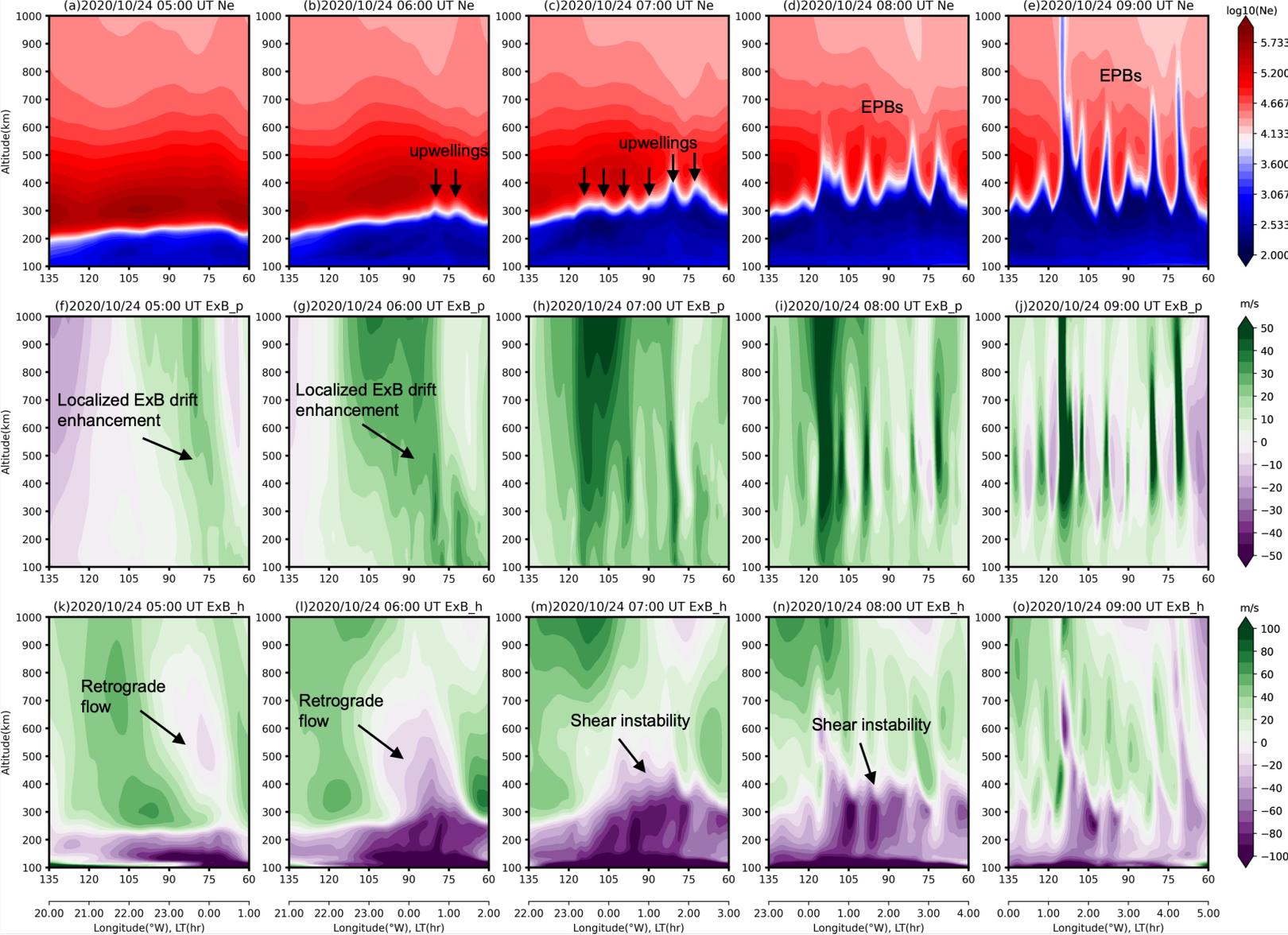


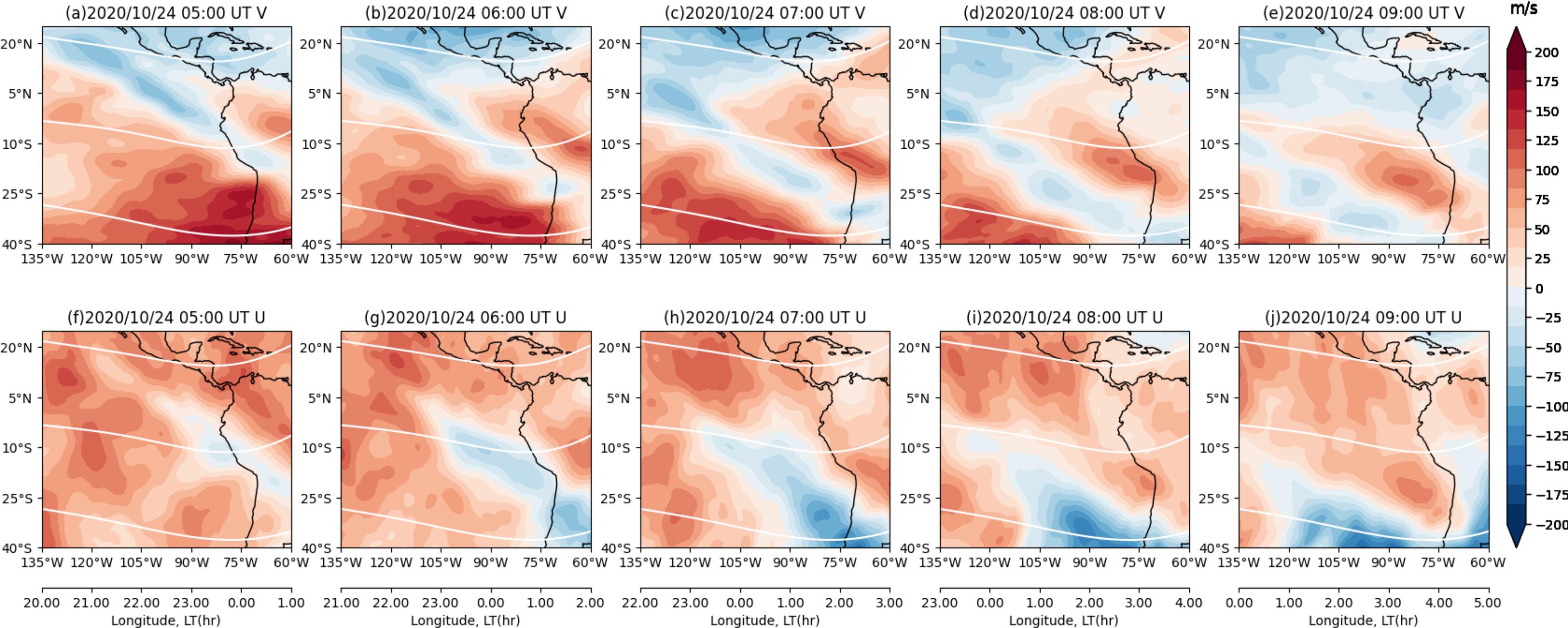




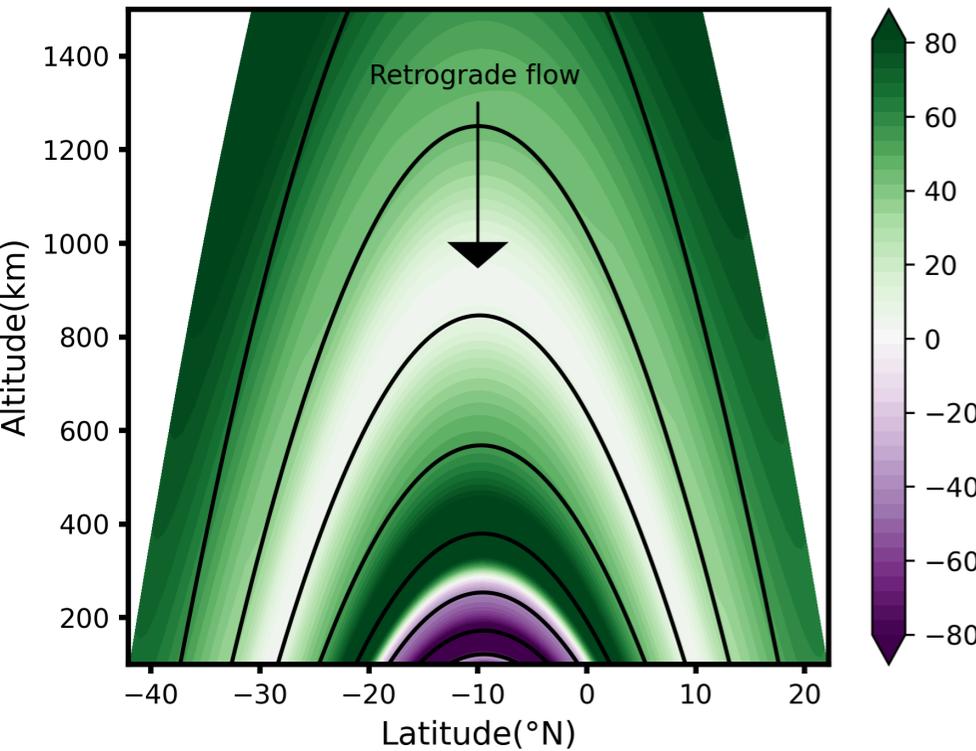




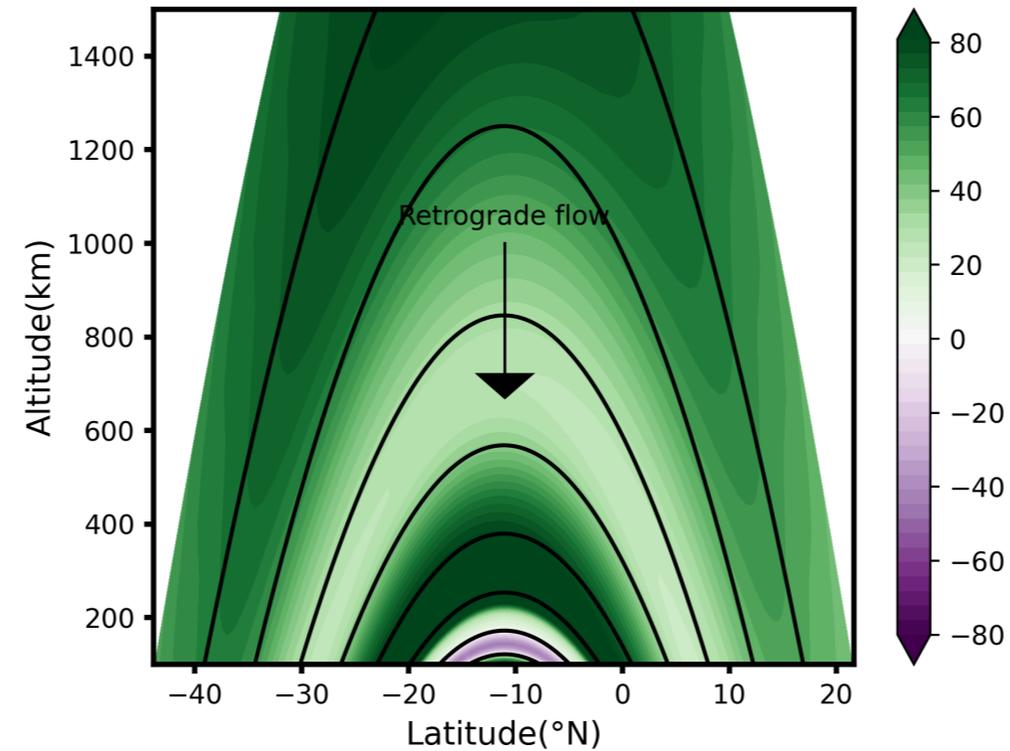




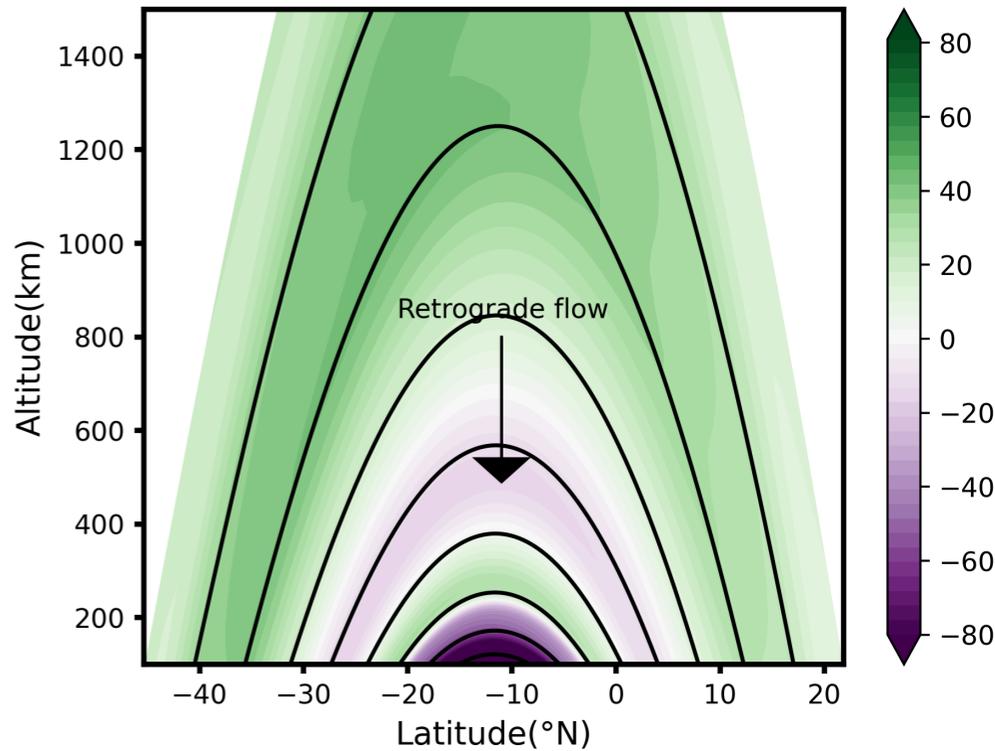
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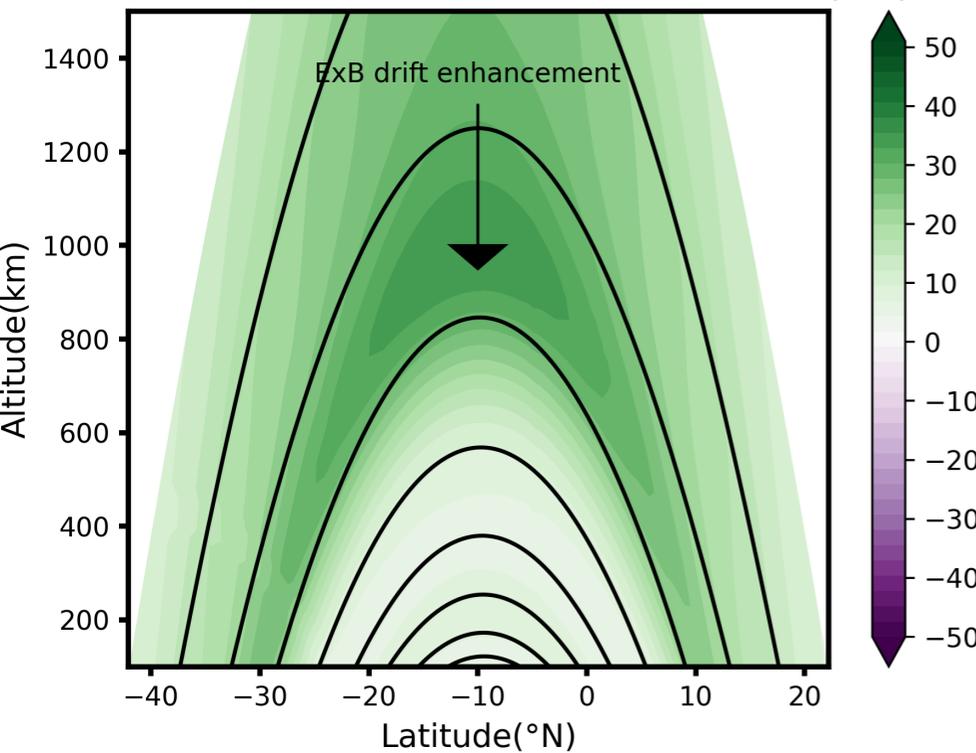
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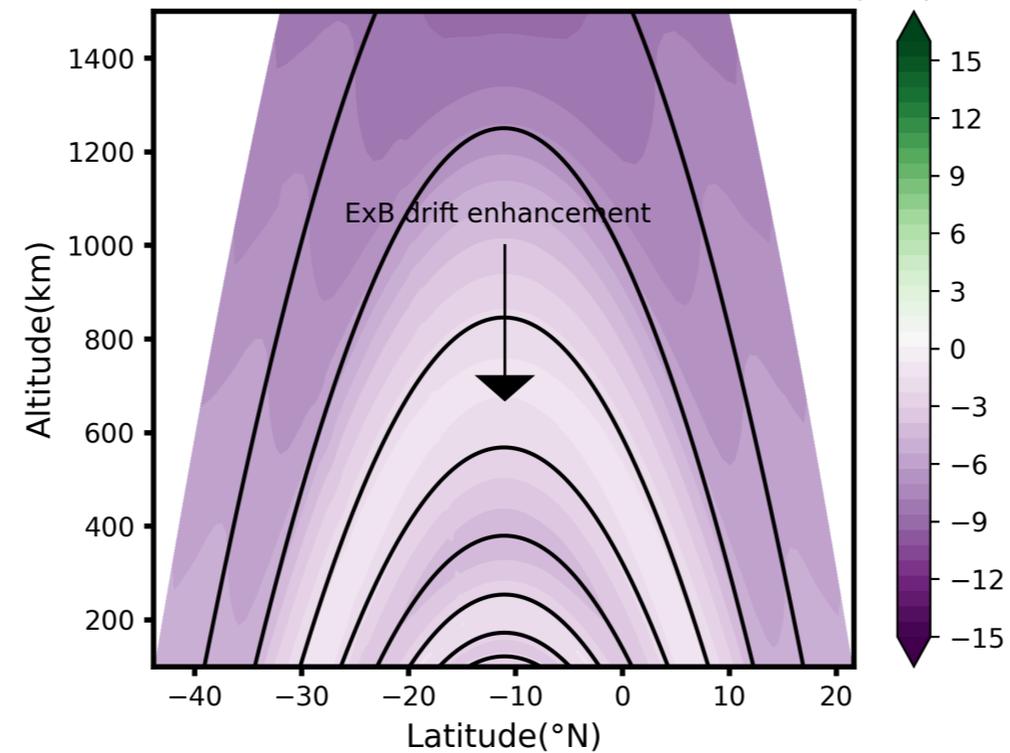
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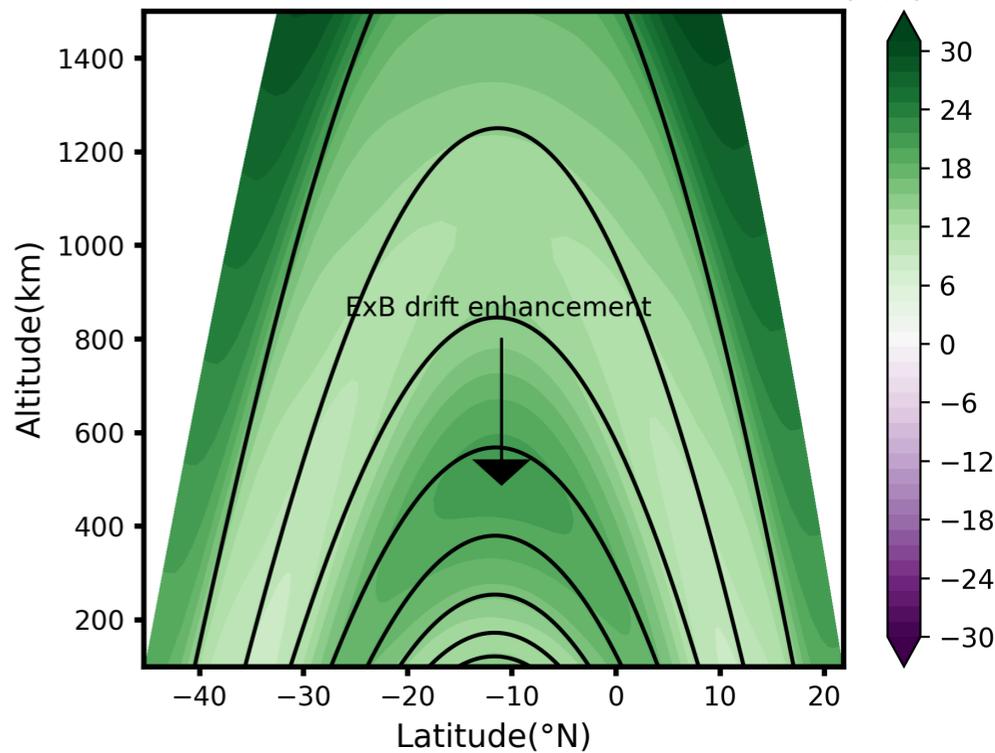
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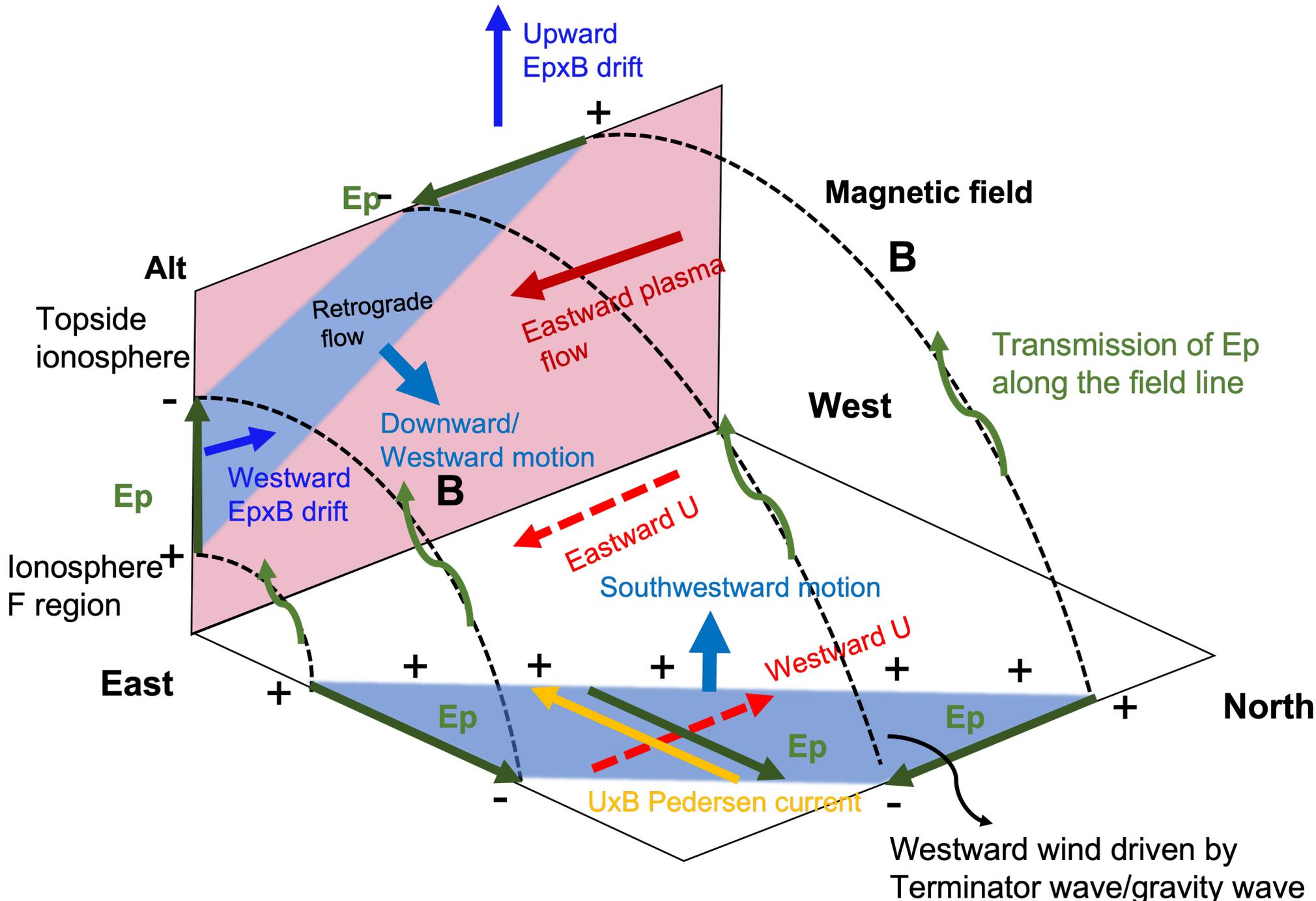


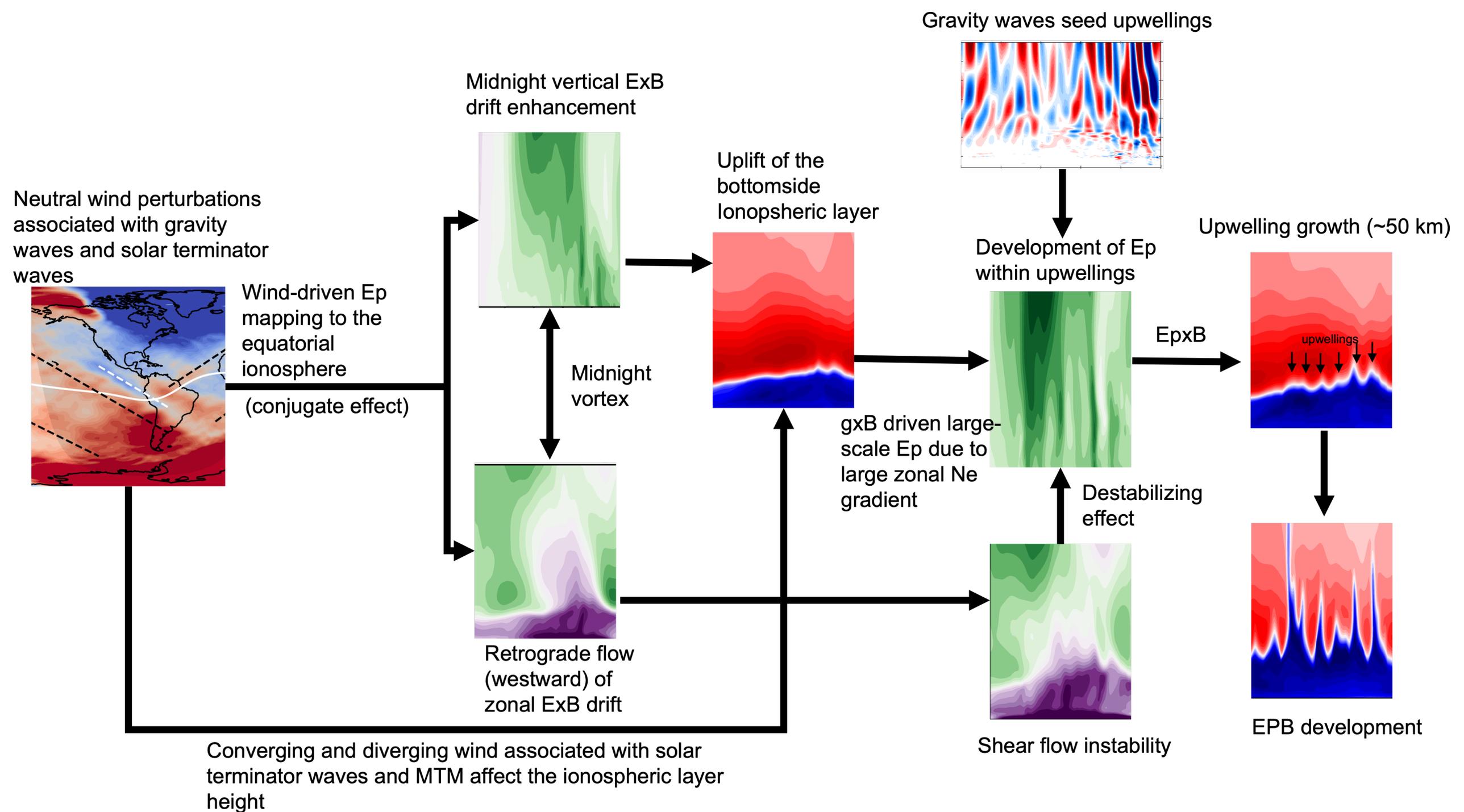
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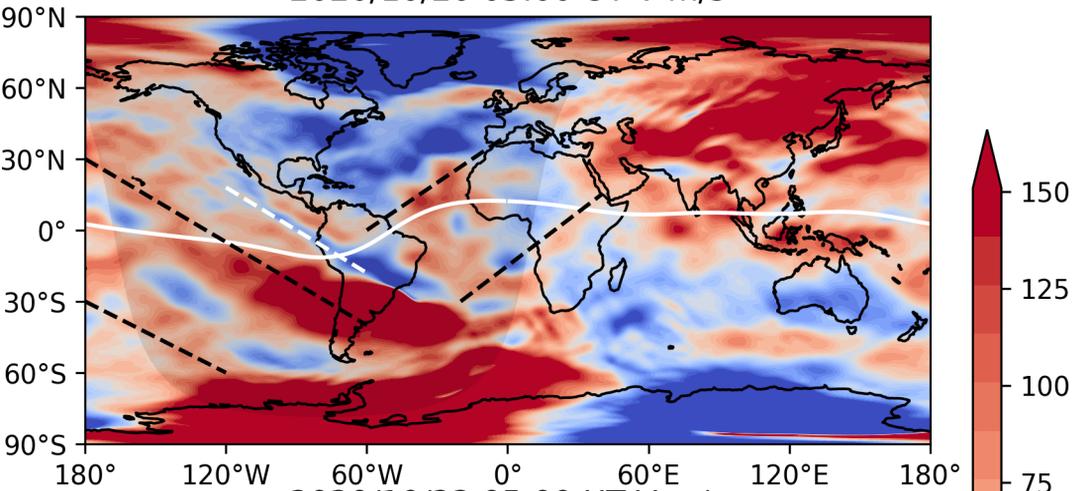
2020/10/24 05:00 UT mLon:357.45 ver/mer ExB (m/s)



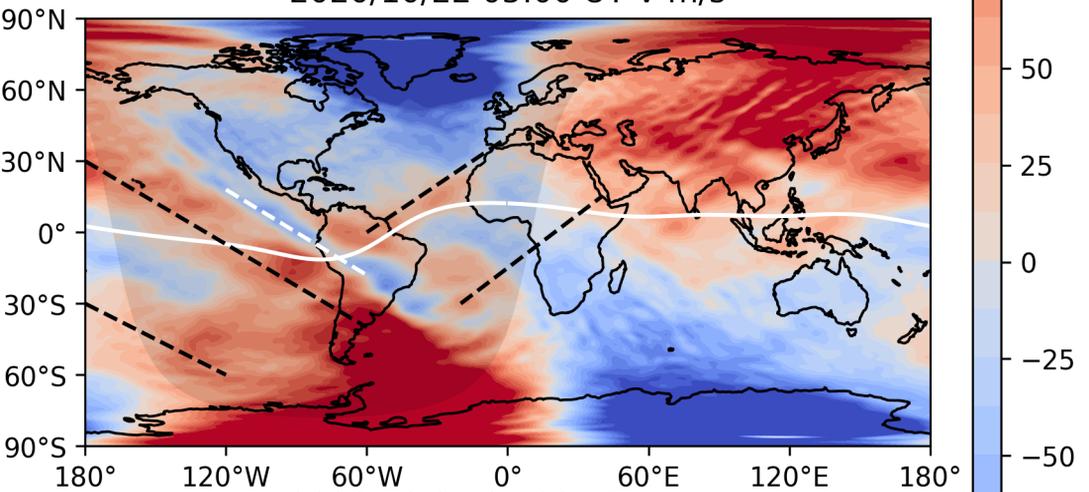




2020/10/20 05:00 UT V m/s



2020/10/22 05:00 UT V m/s



2020/10/24 05:00 UT V m/s

