# Potential role of volcanic glass-smectite mixtures in slow earthquakes in shallow subduction zones: Insights from low- to high-velocity friction experiments

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#### Abstract

Volcanic glass and its mixture with smectite are commonly observed in shallow parts of subduction zones. As volcanic glass layers often act as a glide plane to induce mass transportation such as submarine landslides, and because its alteration product, smectite, is one of the frictionally weakest geological materials, the frictional characteristics of volcanic glass-smectite mixtures are important for fault slip behavior in shallow parts of subduction zones. We performed a series of friction experiments on volcanic glass-smectite mixtures with different smectite contents at various velocity conditions from 10  $\mu$ m/s to 1 m/s under an effective normal stress of 5 MPa and pore pressure of 10 MPa. In general, friction coefficients negatively depend on the smectite content at any velocity tested. We found that samples with smectite contents of 15-30 % showed a drastic slip-weakening behavior at intermediate velocities of 1-3 mm/s with a characteristic slip displacement of ~0.1 m. Finite element method modeling shows that thermal pressurization does not contribute to the observed weakening. The slip-weakening behavior at intermediate velocities enlarges a critical nucleation length for frictional instability to 1-30 km, or prevent acceleration to seismic slip velocities. Therefore, gouges with minor amount of clay, such as subducting volcanic ash layers, may contribute to the occurrence of the at shallow depths in subduction zones.

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1	Potential role of volcanic glass-smectite mixtures in slow earthquakes
2	in shallow subduction zones: Insights from low- to high-velocity
3	friction experiments
4	
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#### 14 Abstract

15 Volcanic glass and its mixture with smectite are commonly observed in shallow 16 parts of subduction zones. As volcanic glass layers often act as a glide plane to induce 17 mass transportation such as submarine landslides, and because its alteration product, 18 smectite, is one of the frictionally weakest geological materials, the frictional 19 characteristics of volcanic glass-smectite mixtures are important for fault slip behavior 20 in shallow parts of subduction zones. We performed a series of friction experiments on 21 volcanic glass-smectite mixtures with different smectite contents at various velocity 22 conditions from 10  $\mu$ m/s to 1 m/s under an effective normal stress of 5 MPa and pore 23 pressure of 10 MPa. In general, friction coefficients negatively depend on the smectite 24 content at any velocity tested. We found that samples with smectite contents of 15-30 % 25 showed a drastic slip-weakening behavior at intermediate velocities of 1-3 mm/s with a 26 characteristic slip displacement of ~0.1 m. Finite element method modeling shows that 27 thermal pressurization does not contribute to the observed weakening behavior. We 28 propose that gouge fluidization or compaction-induced pore pressure increase may be 29 the cause of the weakening. The slip-weakening behavior at intermediate velocities 30 enlarges a critical nucleation length for frictional instability to 1-30 km, or prevent 31 acceleration to seismic slip velocities. Therefore, gouges with minor amount of clay, 32 such as subducting volcanic ash layers, may contribute to the occurrence of the slow 33 earthquakes at shallow depths in subduction zones.

## 35 Plain language summary

36	Materials erupted from volcanoes deposit on the seafloor and subduct at the trench. One
37	such material, volcanic glass, easily turns into mechanically weak clay minerals such as
38	smectite that can cause an enormous slip during an earthquake in subduction zones. In
39	this study, we experimentally examined the frictional properties of mixtures of volcanic
40	glass and smectite to elucidate fault slip behavior in shallow depths of subduction zones.
41	Experiments with varying smectite content showed that a drastic reduction in fault
42	frictional strength was induced when a small amount of smectite was present at
43	moderately high velocity conditions. This behavior may induce slow earthquakes that
44	could be related to huge earthquakes in subduction zones.
45	
46	Key points
47	• Small amount of clay mineral induces a drastic slip-weakening at ~1 mm/s.
48	• Thermal processes do not contribute to the weakening behavior.
49	• Slip-weakening enlarges critical nucleation length leading to slow earthquakes.
50	
51	1. Introduction
52	In subduction zones, the downgoing slab provides fluids to the upper plate
53	leading to melt production and the formation of arc volcanoes. Erupted materials from
54	these arc volcanoes, which mainly consist of volcanic glass, deposit on the incoming
55	plate and subduct at the trench as ash layers or dispersed ash in the incoming sediment
56	(Screaton et al., 2009). At the Nankai Trough, for example, approximately 35% of the
57	incoming sediment is of volcanic origin (Scudder et al., 2018). As volcanic glass reacts

58	progressively with water to form smectite through diagenetic and weathering processes
59	(Compton, 1991), and because smectite is frictionally weak (Morrow et al., 2017), the
60	frictional properties of volcanic glass-smectite mixtures are important for understanding
61	natural hazards associated with subduction zones and volcanic regions, such as
62	earthquakes and landslides. For example, hemipelagic smectite may have been
63	responsible for the large fault slip that occurred near the trench during the 2011 Tohoku-
64	Oki earthquake at the Japan Trench (Chester et al., 2013; Ujiie et al., 2013), with the
65	smectite originating from altered volcanic glass that was deposited on the northwestern
66	part of the Pacific Plate (Kameda et al., 2015). Seismic fault behavior can also occur in
67	these volcanic materials during caldera collapse or magma ascent (Han et al., 2019;
68	Lavallée et al., 2014), and volcanic ash layers may also act as glide planes for
69	subaqueous and terrestrial landslides (Wiemer & Kopf, 2017). Such examples of ash
70	layers at the base of the mass-transport deposits have been reported at the Nankai
71	Trough (Laberg et al., 2017; Strasser et al., 2012), the Middle America Trench (Harders
72	et al., 2010) and Lake Villarrica in Chile (Wiemer et al., 2015), as well as for terrestrial
73	landslides in New Zealand and Japan (Kluger et al., 2020; Moon, 2016; S. Nakamura et
74	al., 2014; Shimizu & Ono, 2016).
75	Despite its potentially important role in natural hazards, our understanding of
76	the frictional behavior of volcanic glass-smectite mixtures is limited. Previous
77	experimental studies have shown that volcanic glass has friction coefficients of 0.6-0.8
78	at low velocity conditions less than 1 mm/s (Lavallée et al., 2014; Wiemer & Kopf,
79	2015), consistent with "Byerlee's rule" for rock friction (Byerlee, 1978). When it is
80	mixed with smectite, the frictional strength of the glass-smectite mixture shows an

81	inverse relationship with smectite content (Wiemer & Kopf, 2015), similar to other
82	previous results where weak clay minerals have been mixed with non-clay minerals
83	(Bedford et al., 2022; Ikari et al., 2018; Logan & Rauenzahn, 1987; Okuda et al., 2021;
84	Takahashi et al., 2007; Tembe et al., 2010). At a high slip-velocity up to 1.3 m/s, the
85	frictional strength of volcanic glass is dynamically weakened (Lavallée et al., 2014). As
86	for smectite, a low friction coefficient of 0.1-0.2 under wet conditions has been reported
87	at low sliding velocities (Morrow et al., 2017). Friction is also low at high sliding
88	velocities, but thermal pressurization additionally plays a significant role in the
89	weakness of wet smectite during high velocity slip (Faulkner et al., 2011; Ujiie &
90	Tsutsumi, 2010). Although we have some constraints on the individual frictional
91	behaviors of smectite and volcanic glass from the previous studies mentioned above,
92	there is a lack of data on their frictional behavior when they are mixed together and
93	sheared over a range of velocity conditions, particularly intermediate velocities of an
94	order of 1 mm/s.
95	To understand better the frictional behavior of volcanic glass-smectite mixtures,
96	we performed a series of friction experiments from low to high velocity under
97	controlled pore pressure conditions. Based on the obtained frictional properties of glass-
98	smectite mixtures, we discuss the implications for fault slip behavior at shallow
99	subduction zones where volcanic glass is abundant (e.g., the Nankai Trough, SW Japan).
100	
101	2. Method
102	2.1. Materials
103	Volcanic glass samples erupted from the Baekdusan volcano in the North Korea,

104	collected at the southern part of Hokkaido, Japan (Y. Nakamura, 2016; Nakanishi et al.,
105	2020), were used for experiments. This volcanic glass has been used as a chronological
106	marker of the eruption at Baekdusan volcano in $10^{th}$ century, and few minerals other
107	than glass are present after being transported over 1000 km (Y. Nakamura, 2016).
108	Samples were sieved to obtain particle sizes between 63-250 $\mu$ m. The broad halo
109	without significant peaks in the XRD profile indicates the sample has little impurity and
110	is predominantly comprised of amorphous glass (Figure 1a). Scanning electron
111	microscope (SEM) images shows that the glass consists of angular shards (Figure 1a).
112	For the smectite sample, we used commercially available purified smectite ("Kunipia-F"
113	provided by Kunimine Industries Co. Ltd., Japan), comprised of montmorillonite with
114	less than 2% quartz as an impurity. Mean grain size is reported to be 2.5 $\mu$ m, but
115	particle aggregates can be several hundred micrometers in diameter (Figure 1b). No
116	disaggregation procedure was employed. We prepared 10 g testing samples for each
117	experiment by mixing volcanic glass and smectite with different gravimetric smectite
118	content under room dry condition (Table 1).
119	





- 122 XRD profiles and SEM micrographs for (a) volcanic glass and (b) smectite. (c-h)
- 123 Unsheared microstructures (SEM backscattered electron images) for each tested clay
- 124 mineral content after applying the effective normal stress of 5 MPa. Bright, angular
- 125 grains are volcanic glass particles, dark areas are smectite particles, and black areas are
- 126 epoxy resin.
- 127
- 128 **2.2. Experimental procedure** 
  - 7



$$V_{eq} = \frac{4\pi R(r_o^2 + r_o r_i + r_i^2)}{3(r_o + r_i)}, \#(1)$$

147 where *R* is the revolutions rate,  $r_o$  and  $r_i$  are the inner and outer radii of the hollow 148 cylindrical specimen, which are 15 and 30 mm, respectively. In most tests,  $V_{eq}$  was 149 initially set to be 100 µm/s for the first 300 mm of shear displacement (nominally "pre-

150	slip"), and then reduced to 10 $\mu$ m/s followed by stepwise increases by a factor of 3 up to
151	300 mm/s ("multi-V test"). Acceleration and deceleration rates were set to be $0.5 \text{ m/s}^2$ .
152	Before stepping from 3 mm/s to 10 mm/s, rotation was stopped for a short time to shift
153	from the slow gear to the fast gear. During rotation with the slow gear, the acceleration
154	rate during the velocity steps is quick enough to be regarded as an abrupt change in $V_{eq}$ ,
155	up to the step from 1 mm/s to 3 mm/s where the acceleration phase finished in less than
156	0.004 s. Shear displacements were set to be 15 mm for $V_{eq} = 10-300 \mu\text{m/s}$ , 150 mm for
157	$V_{eq} = 1$ and 3 mm/s, 1.5 m for $V_{eq} = 10$ and 30 mm/s, and 4.5 m for $V_{eq} = 100$ and 300
158	mm/s (Figure 2b). For samples with smectite contents of 0, 15, and 70%, experiments
159	with low $V_{eq}$ from 10 µm/s to 3 mm/s were also conducted (PHV596, 592, and 595;
160	"low-V test"). For the 15%-smectite sample, $V_{eq}$ was reduced to 10 µm/s after it reached
161	to 3 mm/s, and then increased again to 3 mm/s (PHV592). Additional tests for 100%-
162	smectite and 100%-glass samples were conducted with $V_{eq}$ of 1 m/s for 12.5 m after the
163	pre-slip of 300 mm with $V_{eq} = 1$ mm/s (PHV593 and 594; "single-velocity test"). Test
164	conditions are summarized in Table 1. Friction coefficient $\mu$ is calculated by dividing $\tau$
165	with $\sigma_{eff}$ . The $\tau$ value between Teflon liner and O-ring was less than 0.15 MPa under $P_f$
166	= 10 MPa conditions at $V_{eq} < 3$ mm/s; no correction for the seal friction was applied.
167	Axial displacement, volume of pore water expelled from the gouge, and $T$ in the vicinity
168	of the gouge were measured during experiments to evaluate the gouge thickness and
169	fluid expulsion. All experimental data were collected at a sampling rate of 100 Hz.
170	





173 (a) Schematic view of the experimental configuration. (b and c) Velocity sequences

- 174 employed in this study. (d) Location for the microstructural observation and the
- 175 schematic deformation texture.

#### **Table 1.**

Run	Glass /	$V_{eq}$ [m/s]	Displacement	Comment
	smectite		[m]	
	[%]			

170	T	•	•	.1 *	. 1
1.1.8	List of e	vnorimonte	111	thic	ofudy
1/0			ш	uns	stuuv.
		1			5

	smectite		[m]	
	[%]			
PHV591	100 / 0	$10^{-5} - 10^{-0.5}$	13	Multi-V test
PHV586	85 / 15	$10^{-5} - 10^{-0.5}$	13	Multi-V test
PHV587	70 / 30	$10^{-5} - 10^{-0.5}$	13	Multi-V test
PHV588	50 / 50	$10^{-5} - 10^{-0.5}$	13	Multi-V test
PHV589	30 / 70	$10^{-5} - 10^{-0.5}$	13	Multi-V test
PHV590	0 / 100	$10^{-5} - 10^{-0.5}$	13	Multi-V test
PHV593	100 / 0	$10^{0}$	13	Single-V test
PHV594	0 / 100	$10^{0}$	13	Single-V test
PHV592	85 / 15	$10^{-5} - 10^{-2.5}$	1.2	Low- $V$ test with additional $V$
				steps
PHV595	30 / 70	$10^{-5} - 10^{-2.5}$	0.7	Low-V test
PHV596	100 / 0	$10^{-5} - 10^{-2.5}$	0.7	Low-V test

#### **2.3. Microstructural observations**

181 After the experiments, gouge samples were retrieved from the apparatus and

182 impregnated with epoxy resin so they could be cut and polished to obtain cross-

183 sectional microstructural images normal to the gouge layer and subparallel to the shear

184 direction (Figure 2d). The microstructural images were acquired at the Atmosphere and

185 Ocean Research Institute (University of Tokyo, Japan) using a JEOL JXA-8900

- 186 scanning electron microscope (SEM) with an accelerating voltage of 15 kV. As it was
- 187 difficult to recover the entire sample due to the unconsolidated nature of the gouge, we
- 188 focused our observations on the structures in the vicinity of the shear band that forms

189 near the edge of the gouge layer (SB, Figure 2d), where the shear deformation is

- 190 localized compared to surrounding undeformed zone (UDZ, Figure 2d).
- 191

#### 192 2.4. Finite Element Method (FEM) modeling

193 To constrain the T and  $P_f$  conditions inside the gouge, we used COMSOL

194 Multiphysics® software and modeled time-dependent heat transfer and fluid diffusion

195 which follow the equations below (Rice, 2006):

$$\frac{\partial T}{\partial t} = \frac{1}{\rho_g C_p} \nabla (K \nabla T) + \frac{A(r)}{\rho_g C_p}, \#(2)$$

$$\frac{\partial P_f}{\partial t} = \frac{1}{\rho_f \phi(\beta_f + \beta_\phi)} \nabla \left(\frac{\rho_f k}{\eta_f} \nabla P_f\right) + \frac{(\gamma_f - \gamma_\phi)}{(\beta_f + \beta_\phi)} \frac{\partial T}{\partial t}, \#(3)$$

196 where *t* is time,  $\rho_f$  is water density,  $\rho_g$  is density of a material,  $C_p$  is heat capacity of a 197 material, *K* is thermal conductivity of a material,  $\gamma_f$  is water thermal expansivity,  $\gamma_{\phi}$  is 198 pore space thermal expansivity,  $\beta_f$  is water compressibility,  $\beta_{\phi}$  is pore space 199 compressibility,  $\phi$  is porosity,  $\eta$  is water viscosity, *k* is permeability (see Table 2 for 200 details). We assume  $k = 10^{-20}$  m<sup>2</sup> to simulate the low permeability nature of clay-mixed 201 gouge (Oohashi et al., 2015; Takahashi et al., 2007). *A*(*r*) is heat generation per unit 202 volume at the distance from the rotation axis *r* calculated as follows:

$$A(r) = \frac{1}{w_{hs}} V(r)\tau = \frac{1}{w_{hs}} \left(\frac{3}{4} (2r) \frac{r_o + r_i}{r_o^2 + r_o r_i + r_i^2} V_{eq}\right)\tau, \#(4)$$

203 where  $w_{hs}$  is thickness of the heat source where shear localized within the gouge and 204 V(r) is velocity at r. We assumed that  $\tau$  is independent of r within the limited velocity 205 range from  $V(r_i) = 0.64V_{eq}$  to  $V(r_o) = 1.29V_{eq}$ . Dilation/compaction of the gouge and 206 thermoelastic response of the apparatus are not considered in the FEM modeling. 207 A schematic representation of FEM configuration is shown in Figure 3. The 208 model was made using the axial symmetry, thus only one side of the gouge layer shown. 209 Fluid diffusion (equation 3) is only considered inside the gouge zone with thickness of 2 210 mm. We set the initial  $P_f$  conditions throughout the entire gouge layer to be 10 MPa, 211 and the  $P_f$  was fixed at this value at the top and bottom boundaries of the gouge layer 212 (Dirichlet boundary condition). The inner and outer boundaries have no flow of fluid 213 crossing the boundaries. Based on microstructural observations, the heat source inside 214 the gouge was assumed to be concentrated in the shear localized zone (SB) with  $w_{hs}$  = 215 0.3 mm at 0.1 mm below the top of gouge. We extract T data at r = 22.5 mm at a 216 distance of 1 mm above the gouge layer and compare it with the measured T during 217 experiments. For simplicity, shear forces between O-rings and Teflon liners and the 218 frictional heat generated from the outer and inner contacts are ignored in the FEM 219 model. Based on the computed  $P_f$  inside the heat source where shear localized, we 220 define  $P_{feq}$  which represents the "equivalent" influence of  $P_f$  on frictional behavior as 221 follows (see Appendix A1 for details):

$$P_{feq} = \frac{3}{r_o^3 - r_i^3} \int_{r_i}^{r_o} \int_{z_{hs}^{bot}}^{z_{hs}^{top}} \frac{P_f(r, z)r^2}{w_{hs}} dz dr, \#(5)$$

where  $P_f(r, z)$  is the calculated  $P_f$  at r and z (direction vertical to the gouge) inside the 13

shear localized zone,  $z_{hs}^{top}$  and  $z_{hs}^{bot}$  are z for top and bottom ends of the shear localized zone, respectively, thus  $w_{hs} = z_{hs}^{top} - z_{hs}^{bot}$ . Equation 5 leads to:

$$\mu_{obs}(\sigma_n - P_{fo}) = \mu_{cor}(\sigma_n - P_{feq}), \#(6)$$

where  $\mu_{cor}$  is the corrected  $\mu$  value without the influence of pore pressure variation due to frictional heat, and  $P_{fo}$  is the imposed  $P_f$  (= 10 MPa in this study).

227



228

229 **Figure 3**.

230 (a) Geometry and mesh structure for the FEM modeling. (b) Close up view of the gouge

231 with locations of heat source and the thermocouple.

232

#### **Table 2.**

- 234 Mechanical properties of components in the FEM model at around room temperature.
- 235 Dependence on temperature was employed for the parameter with an asterisk.

Gouge Pore space SUS304 Ti-6Al-4V Teflon O-ring Water

$\rho  [\text{kg/m}^3]$	2400		*7900	*4420	2200	1800	*1004
$C_p$	1000		*460	*516	1050	1500	*1716
[J/kg/K]	1000		*462	*340	1030	1300	*4210
K	15		*116	*7	0.25	0.2	*0 56
[W/m/K]	1.3		14.0	• /	0.23	0.5	0.30
γ [1/K]		$2.4 \times 10^{-5}$					$5.7 \times 10^{-4}$
β[1/Pa]		2.49×10 <sup>-9</sup>					0.49×10 <sup>-9</sup>
$\phi$	0.25						
η [Pa.s]							*1.8×10 <sup>-3</sup>
$k [m^2]$	$10^{-20}$						

#### 237 **3. Results**

#### 238 **3.1. Steady state friction coefficient**

239 Results on frictional strength at all  $V_{eq}$  conditions are summarized in Figures 4

and S1. Friction coefficients ( $\mu$ ) of 100%-glass sample ranged from 0.6 to 0.7 with a V-

strengthening behavior from  $V_{eq} = 10 \,\mu\text{m/s}$  to 30 mm/s. At  $V_{eq} = 100 \,\text{mm/s}, \mu$ 

dynamically weakened from a peak value of 0.8 to a steady state value of 0.2. Such a

low  $\mu$  is also observed at a  $V_{eq}$  of 300 mm/s in the multi-V test (Figure 4a) and 1 m/s in

the single-*V* test (Figure 4b).

As smectite content was increased, frictional strength decreased at every V

- values were about 0.5 and 0.3, respectively, with V-neutral or slightly V-weakening
- 248 trends. At  $V_{eq} = 1-3$  mm/s (intermediate V range),  $\mu$  decreased to 0.1-0.2 with a

<sup>246</sup> condition. For samples with 15 and 30% smectite, at  $V_{eq} \leq 300 \,\mu\text{m/s}$ , steady state  $\mu$ 

249 characteristic weakening time or distance (see section 3.2) and remained low at higher  $V_{eq}$  conditions (Figures 4a and 4c). For samples with more than 50% smectite,  $\mu$  was as 250 low as 0.1 at every  $V_{eq}$  condition, with nearly V-neutral trends. 251 252 According to the FEM modeling (section 3.4), at high-V and low- $\mu$  conditions (V  $> 10^{-2}$  m/s and  $\mu < 0.3$ ), the measured  $\mu$  may be significantly overestimated by friction 253 254 between O-ring and Teflon liners. However, it is difficult to estimate exact effect of the 255 seal friction because the amount of gouge materials that extrude from the layer and fill 256 the space between upper and lower pistons and inner and outer Teflon liners may vary 257 with experimental conditions, such as velocity, displacement, and mineral compositions.







260 (a) Mechanical behavior of multi-V tests ( $V_{eq} = 10 \mu m/s$  to 300 mm/s) with different

smectite contents. (b) Results of single-V tests ( $V_{eq} = 1 \text{ m/s}$ ) for 100%-glass and 100%-

smectite cases. Light colors represent the results of multi-V tests and dark colors

263 represent those of single-V tests. (c) Close view of the low- $V(V_{eq} < 3 \text{ mm/s})$  part of

- 264 multi-V tests shown in (a). (d) Relation between  $\mu$  and  $V_{eq}$  for multi-V and single-V tests
- 265  $(V_{eq} = 1 \text{ m/s})$  after the pre-slip stage. Single-V tests were only conducted for the smectite 17

266 contents of 0 and 100%. Gray area indicates a possible range where friction coefficients267 are overestimated due to the seal friction.

268

#### 269 **3.2.** Frictional behavior for a sample with 15% smectite at low to intermediate

270 velocity range

To see if the observed weakening behavior for 15%- and 30%-smectite samples at the intermediate V range depend on  $V_{eq}$  or on shear strain, we performed an additional experiment (PHV592, 15% smectite) that included both velocity upsteps and downsteps. We found that  $\mu$  recovered to 0.5 during the V downsteps which is consistent with  $\mu$  in the low V range before the V upsteps (Figure 5).  $\mu$  decreased again in the subsequent V upsteps, which implies that the steady-state  $\mu$  in the intermediate V range is independent of shear strain and only dependent on V.

278 Just after the V step,  $\mu$  showed a positive spike followed by a gentle increase 279 and then a decrease to its steady state  $\mu$  at each V condition (Figures 5b-e). The spikes 280 appear to be comprised of a direct effect and an evolutionary effect described by the 281 rate- and state-dependent friction (RSF) law (Figures 5f-i) (Dieterich, 1979; Ruina, 282 1983). The evolutionary effect has a characteristic slip distance of  $\sim 100 \,\mu\text{m}$ . The extent 283 of the increase in  $\mu$  after the RSF-like behavior seems to be positively related to the 284 imposed  $V_{eq}$  conditions, except for the V step from  $V_{eq} = 1$  mm/s to 3 mm/s, although 285 the V steps may be affected by the  $P_f$  condition at these slip rates (gray markers in 286 Figures 5b-e; see section 3.4). Little variation in gouge thickness during the V steps was 287 observed (Figures 5b-i). As the V step was imposed within less than 0.004 seconds, and 288 because the gentle increase and decrease in  $\mu$  occurred after the direct and evolutionary

effects, the observed weakening behavior at the intermediate V range to  $\mu = 0.1-0.2$  is not a response to V step (like the RSF law) but a slip-weakening behavior toward the steady state  $\mu$  value for a given V condition. During this behavior, gouge thickness did not vary, whereas the thermocouple captured gentle increases in T especially at  $V_{eq} = 1$ and 3 mm/s (section 3.4).







297 (a) Experimental result of PHV592 (15% smectite, low-velocity test with additional

298	velocity steps). (b-i) Close-up views of variation in friction coefficient at the velocity
299	steps for PHV592. Red lines are experimental data and gray lines in (b-e) are modeled
300	friction coefficient by FEM after correcting pore pressure effect, and black lines are
301	gouge thickness.

#### 303 **3.3. Microstructure**

304 Before imposing shear deformation, volcanic glass grains kept their original, 305 angular shapes regardless of the smectite content (Figure 1). As the smectite content 306 was increased, volcanic glass grains distributed in the smectite matrix and the numbers 307 of contacts between glass grains reduced. After the shear deformation, samples often 308 showed a localized shear band (SB) in the upper part of the gouge layer, characterized 309 by grain size reduction and an increase in the number of rounded grains compared to the 310 surrounding undeformed zone (UDZ). 311 The thickness of the SB for the 100%-glass sample sheared at low-V (PHV596, 312 Figures 6a-c) was about 300 µm, with the SB containing more rounded grains than the UDZs. The mean grain size in the SB was less than  $\sim$ 50 µm, with the minimum grain 313 314 size being  $<1 \mu m$  (Figure 6c). 315 For the sample with  $V_{eq}$  up to 300 mm/s (PHV591, Figures 6d-f), the SB had a 316 similar thickness as the low-V sample (PHV596), however a localized shear band (LSB) 317 with a thickness of 10-20 µm had also formed in the middle of the main SB (Figure 6f). 318 The largest grain size in the LSB was  $<1 \mu m$ , which was significantly smaller than the 319 surrounding SB. The mean grain size in the main SB was several tens of microns, which 320 was a little smaller than that in the SB of the low-V sample (PHV596). In some portions

321 of the SB and UDZ, larger fragments (FR; Figures 6d and 6e) containing features that 322 look like LSBs were observed, possibly indicating reworking of localized bands that 323 formed during previous V-steps in the multi-V test. 324 For low-smectite samples (e.g., 15%-smectite, PHV592, low-V test), a SB was 325 also observed, however it was a little wider ( $\sim$ 500 µm) than 100%-glass case (Figure 326 6g). The grain size of glass particles was reduced to less than 50  $\mu$ m (Figure 6h). The 327 grain shape appeared to be a little more angular than 100%-glass case (Figure 6i), and 328 the crushed glass particles were dispersed and mixed with smectite (Figure 6i). 329 For high-smectite samples, no SB was observed (Figure 6j). No obvious 330 localized feature was observed for the 70%-smectite sample (PHV589) and one shear 331 plane parallel to the deformation was likely observed for 100%-smectite sample 332 (PHV590), although it is potentially an unloading feature. Glass grains in the 70%-333 smectite sample were not in contact with each other and preserved their original angular 334 shapes and sizes (Figures 1, 6h, and j). As the size of smectite particles was less than 2 335  $\mu$ m, we could not easily resolve the deformation within the smectite matrix, although 336 the deformation may be primarily accommodated by smectite matrix. 337



#### 339 Figure 6.

- 340 Microstructures of deformed samples for (a-c) 100%-glass with low-V test of PHV596,
- 341 (d-f) that with multi-V test of PHV591, (g-i) 15%-smectite with low-V test of PHV592,
- and (j-l) 70%- or 100%-smectite with multi-V tests of PHV589 or 590. The left side of
- (j) is the microstructure for PHV589 (70% smectite) and the right one is that for22

344 PHV590 (100% smectite). Locations of figures in the center and right columns are

345 shown in the left and center columns, respectively. SB: shear band; UDZ: undeformed

346 zone; LSB: localized shear band; FR: fragment.

347

#### 348 **3.4.** Temperature and pore pressure conditions inside the gouge

349 In the FEM modeling, we obtained fairly consistent result of T with the 350 measured value for the low-V tests (Figure 7a). The modeled maximum T increase 351 within the gouge is 12°C for 100%-glass, 6-7°C for 85%-glass, and less than 3°C for 352 <70%-glass samples (Figure 7b). For the 100%-glass sample at high-V conditions, the 353 modeled T at the thermocouple position showed consistent result with measured T, at 354 least for shear displacements of <5 m, which is before the friction dynamically 355 weakened (Figure 7c, see also Figure 4a). The maximum T inside the gouge reached 356 about 340°C when the shear displacement was 5 m, which was the onset of dynamic 357 weakening (Figure 7d). After the dynamic weakening behavior, the modeled T was 358 significantly overestimated, which implies a significant shear resistance between the O-359 rings and Teflon liners at high-V conditions. Such a discrepancy between modeled and 360 measured T conditions was also observed for high-V conditions with other samples. 361 Hence, the  $\mu$  of the gouge may be nearly zero at high-V conditions, but as we cannot 362 know the true  $\mu$  as a result of the contribution from O-ring and Teflon friction, we do 363 not discuss the friction values in detail at the V conditions of the shaded region in Figure 364 4d.

The distribution of  $P_f$  within the gouge was simultaneously modeled in the FEM modeling. In the case of PHV592 (15% smectite), the maximum *T* increase was only 6-

367	7°C (Figure 7b) because of slow V and relatively low $\mu$ , meaning T within the gouge is
368	far lower than the boiling T of water at $P_f = 10$ MPa (Figure 8c). However, the low
369	permeability nature of the clay-mixed gouge ( $k = 10^{-20} \text{ m}^2$ ) caused an increase in $P_f$
370	inside the gouge of up to 0.3 MPa (Figure 8e), with $P_{feq}$ being about 10.2 MPa just after
371	the V-step from $V_{eq} = 1$ to 3 mm/s (Figure 8a). This $P_f$ build-up caused a reduction in $\sigma_{eff}$
372	(Figure 8a) and thus the measured $\mu$ was apparently lowered. We corrected the $\mu$ in
373	PHV592 by using equations 5 and 6 and found that the measured $\mu$ value just after the
374	<i>V</i> -steps from $V_{eq} = 1$ to 3 mm/s was underestimated by about 0.02 (Figure 5e), however
375	this is insufficient to explain the slip-weakening behavior we observe (Figure 8b).
376	The $P_f$ conditions inside the gouge are also influenced by gouge compaction and
377	dilation, which are not implemented in the FEM modeling. As shear displacement
378	increased during the test, gouge thickness reduced as measured by the shortening in the
379	axial displacement data. The majority of axial shortening occurred during the pre-slip
380	stage, but gouge thickness continuously decreased afterward (Figure 8c). However, the
381	volume of expelled water was consistent with the amount of axial shortening (ideally,
382	the volume of expelled water [cc] = 0.001 [cc/mm <sup>3</sup> ] × (30 [mm]) <sup>2</sup> × $\pi$ × axial shortening
383	[mm]), suggesting that the $P_f$ build up due to the gouge compaction would be minimal.
384	Furthermore, gouge compaction occurred simultaneously with the slip-hardening in the
385	V downsteps. Therefore, the influence of gouge compaction on $P_f$ is likely to be
386	insignificant and not a dominant cause of the observed slip-weakening behavior and low
387	steady state $\mu$ values at $V_{eq} = 1$ and 3 mm/s.
388	







391 (a) Measured T (light color) and modeled T (solid line) at the thermocouple position for

392 low-V conditions. (b) Modeled maximum T increase within the gouge for low-V

393 conditions. (c) Measured T (gray), modeled T at the thermocouple position (blue) and

- 394 the maximum T inside the gouge (red) for high-V conditions of PHV591 (glass 100%).
- 395 (d) Distribution of T around the gouge at the shear displacement of  $\sim$ 5 m indicated in (c).





Equivalent pore pressure and corrected effective normal stress conditions modeled by FEM calculations for PHV592. Permeability of  $10^{-20}$  m<sup>2</sup> was assumed. (b) Apparent (gray line) and corrected friction coefficients (orange line) by the corrected effective normal stress condition. (c) Comparison between the volume of expelled water (gray line) and axial shortening (black line). Results of FEM simulation on (d) temperature and (e) pore pressure conditions at the shear displacement of ~0.5 m, indicated in (a).

## **4. Discussion**

# **4.1. Dynamic weakening for 100% glass caused by thermal processes**

408	The $\mu$ for 100%-glass sample at $V_{eq} = 10$ mm/s was higher than that typically
409	observed for quartz-rich rocks, which often shows a reduction in frictional strength at
410	similar intermediate slip velocities due to silica-gel lubrication or the formation of
411	amorphous wear materials (Goldsby & Tullis, 2002; Hayashi & Tsutsumi, 2010; Rowe
412	et al., 2019; Di Toro et al., 2004). In addition, because the dynamic weakening in the
413	100% glass sample occurred just after the maximum $T$ within the gouge reached the
414	boiling $T$ of 311°C for water with 10 MPa (Figure 7c), the observed dynamic weakening
415	in the multi- $V$ test may be caused by vaporization of the pore water (Acosta et al., 2018;
416	Chen et al., 2017; Hunfeld et al., 2021). This result supports the idea that thermal
417	processes are important for dynamic weakening and that the amorphous grains
418	themselves do not induce the weakening (Kanagawa et al., 2020). Although the FEM
419	modeling of the single- $V$ test could not estimate the $T$ condition inside the gouge,
420	because the measured $T$ was lower than in the multi- $V$ test, the dynamic weakening in
421	the single- $V$ test may instead be caused by flash heating or thermal pressurization, rather
422	than vaporization. The apparent difference in the slip distance for the dynamic
423	weakening between multi- $V$ and single- $V$ tests might indicate a difference in the
424	weakening mechanisms for the same materials under similar pressure conditions but
425	under different ambient $T$ conditions (higher $T$ for multi- $V$ tests because of the frictional
426	heat generating during preceding slips); further experimental studies will be needed to
427	validate this.

# 429 4.2. Athermal slip weakening for low-smectite gouges at intermediate velocity 430 range

431 The result of PHV592 which included both velocity upsteps and downsteps (section 3.1; Figure 4) indicated that the slip weakening behavior at the intermediate  $V_{eq}$ 432 433 range (1-3 mm/s) for the 15%-smectite sample is not a result of high shear strain but 434 purely a V dependent phenomenon. Also, the FEM modeling showed that the  $P_f$  build-435 up due to frictional heat had minor influence on the reduction in apparent  $\mu$  values for 436 such a low  $V_{eq}$  conditions (Figure 8). Previous studies reported similar weakening at ~1 437 mm/s for low-clay samples and discussed the possibility of flash heating followed by 438 local thermal pressurization (Oohashi et al., 2013, 2015). In the case of this study, the 439 flash heating *T* can be calculated by the following equation (Archard, 1959):

$$\Delta T_{flash} = \frac{1}{8} \frac{\pi \mu_a H l V_a}{K}, \#(7)$$

440 where  $V_a$  is the velocity at an asperity contact, and l is the size of an asperity contact,  $\mu_a$ 441 is the friction coefficient at an asperity contact, and H is the compressive yield strength 442 of 7 GPa (Ben Abdelounis et al., 2009; Yonekura, 2015). We assume that K is 1.5 443 W/m/K for glass (Romine et al., 2012) and  $\mu_a$  is identical to the  $\mu$  value for 100% glass, 444 i.e.,  $\mu_a = 0.6$ , because the asperity contact is the contact between glass grains. As the 445 grain size in the SB was less than 30 µm (Figure 5i), we used an l value of 1-30 µm. The 446 estimated  $\Delta T_{flash}$  for  $V_a = 1$  mm/s is 1-33°C, which is far lower than the melting T of 447 typical rocks (~1,000°C). The  $\Delta T_{flash}$  of 1-33°C may not be high enough to induce local 448 reduction in shear strength by thermal pressurization near asperity contacts; therefore, 449 other mechanisms would be needed to reduce frictional strength. 450 One possible process that could reduce  $\mu$  is a fluid-like behavior of the gouge
(fluidization). For the 15%-smectite sample at intermediate 
$$V$$
 (PHV592), deformation  
likely occurred within the shear band (SB) with thickness of ~500 µm (Figures 6g-i)  
rather than a discrete shear surface like observed in the 100% glass sample with  $V_{eq}$  up  
to 300 mm/s (Figures 6d-f). Such distributed shear deformation is one of the  
characteristics of fluidization. However, further experimental and microstructural  
studies would be needed to ensure the occurrence of fluid-like behavior because we did  
not observe grain segregations, which is recognized as a form of evidence for gouge  
fluidization (Demurtas et al., 2021; Ujiie & Tsutsumi, 2010). In addition, because the  
fluidized gouge should be insensitive to  $\sigma_{eff}$ , we need to test the same material under  
different  $\sigma_{eff}$  conditions (Ujiie & Tsutsumi, 2010), which is planned in future work.  
Another possibility is compaction-induced pore fluid pressurization. During the  
experiments, we observed a continuous reduction in gouge thickness which appears to  
depend on displacement (Figure 8c). Compaction occurred in a short time when  $V_{eq}$  was  
in a higher range. If we consider a simple form of  $P_f$  increase without diffusion:

. .

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1

$$\frac{\Delta P_f}{\Delta t} = -\frac{1}{\beta_f \phi w_g} \frac{\Delta w_g}{\Delta t}, \#(8)$$

465 where  $w_g$  is the thickness of gouge (2 mm), and if we assume the rate of reduction in 466 gouge thickness  $\Delta w_g/\Delta t$  of 0.0025 mm/s,  $P_f$  increases by 10 MPa in a second. Therefore, 467 the observed weakening in friction can be explained by compaction-induced 468 pressurization, however there is a large uncertainty regarding the value of  $\phi$  because of 469 gouge leakage to the narrow spaces between the O-rings and Teflon liner, thus we 470 cannot separate effects of compaction and gouge leakage on Pf. Note that compaction-471 induced pressurization has been reported at a high velocity conditions of  $\sim 1$  m/s for the

(0. 11)

. •

472 sediment gouge from the Hikurangi subduction zone (Aretusini et al., 2021), and at a 473 low velocity conditions of  $\sim 1 \mu m/s$  for gouges with low permeability (Faulkner et al., 474 2018).

475

#### 476 4.3. Low friction coefficient at all velocity conditions for high-smectite samples 477 As the clay mineral content increases, the permeability of the fault gouge continuously decreases down to $\sim 10^{-21}$ m<sup>2</sup> for 100% smectite (Takahashi et al., 2007). 478 479 This low permeability contributes to the fluid overpressure within the gouge. 480 Consequently, the frictional strength becomes quite low even at low-V conditions, such 481 that the samples with more than 50% smectite showed low $\mu$ close to or less than 0.1 482 (Figure 3). Due to the intrinsic weakness of wet smectite (Morrow et al., 2017) and its 483 impermeable nature, frictional heat generation was low and thus thermal pressurization 484 likely did not play a significant role in the observed weakness of the clay. In addition, 485 the chances of grain contacts among volcanic glass grains are less than low-clay 486 samples because the volcanic glass grains are dispersed in the pervasive clay matrix and 487 not comminuted (Figures 6k and 1). This observation is consistent with the similar 488 friction coefficients for 50%-, 70%-, and 100%-smectite samples suggesting that the 489 smectite deformation dominantly controls the frictional behavior of the entire gouge 490 when the clay mineral content exceeds 50% (Oohashi et al., 2013). If higher slip 491 velocity and/or longer shear displacement were imposed, increased frictional heat would 492 be generated, and thermal pressurization of pore fluid or water dehydrated from clay 493 could potentially further decrease the apparent frictional strength (Hirose & Bystricky, 494 2007).

# 496 5. Possible impacts of slip weakening at intermediate velocity range on slow and 497 fast earthquakes in the shallow subduction zone

The weakening with V and slip displacement can cause acceleration of fault slip and result in an earthquake. We observed V-neutral or slightly V-weakening trends up to  $V_{eq} = 300 \,\mu\text{m/s}$  for low-clay samples (see section 3.1). Considering a spring-slider model, unstable, self-accelerating slip occurs when the system stiffness K, that is the stiffness of the country rock surrounding a fault, is lower than the critical stiffness  $K_c$ , defined by frictional properties of a fault. The  $K_c$  for the (quasi-static) RSF law is described as follows:

$$K_{c,RSF} = -\frac{\sigma_{eff}(a-b)}{d_c}, \#(9)$$

505 where  $d_c$  is a characteristic slip distance at the V-step, and a-b is defined as follows:

$$a-b=\frac{\Delta\mu_{ss}}{\Delta\ln V}$$
, #(10)

506 where  $\Delta \mu_{ss}$  and  $\Delta \ln V$  are the variation in steady-state friction coefficient and the

507 logarithm of velocity, respectively (Dieterich, 1979; Ruina, 1983). Similarly, the  $K_c$  for

508 a slip-weakening behavior can be defined as follows:

$$K_{c,sw} = \frac{\Delta \tau}{\delta_c}, \#(11)$$

509 where  $\Delta \tau$  is the drop in shear stress during a characteristic slip distance  $\delta_c$  (Ikari et al.,

510 2013). The  $K_c$  value can be converted to a critical nucleation length of slip area  $L_c$  to

511 initiate self-acceleration of slip as follows:

$$L_c = \frac{\xi G}{K_c}, \#(12)$$

512	where $\xi$ is a shape-dependent parameter and G is the shear modulus of surrounding rock.
513	When the size of the slip length $L$ becomes larger than $L_c$ , then the dynamic rupture
514	initiates (Dieterich, 1992; McLaskey, 2019). For simplicity, we adopt a $\xi$ value of $7\pi/12$
515	considering a circular crack with a diameter of $L_c$ (Dieterich, 1986; Eshelby, 1957), and
516	a $G$ value of 10 GPa considering partially consolidated, underthrust or accreted
517	sediment at a shallow depth in a subduction zone (e.g., S-wave velocities of 1.5-2.5
518	km/s and density of 2.0 g/cm <sup>3</sup> lead to 4.5-12.5 GPa) (Akuhara et al., 2020).
519	The obtained $a-b$ , $d_c$ , and calculated $L_c$ for low V conditions, and $\Delta \tau$ , $\delta_c$ , and $L_c$
520	for intermediate $V$ conditions are summarized in Table 3. We used the most negative
521	a-b value for each smectite content to estimate the possible unstable bound of slip. For
522	the 0%-smectite sample, V-strengthening behavior was observed for $V_{eq} = 10 \ \mu\text{m/s}$ to
523	300 mm/s; therefore, earthquake nucleation would not be expected to occur under either
524	low or intermediate $V$ conditions. Low-smectite samples (15-30% smectite) showed
525	negative $a-b$ values of $-0.012$ to $-0.003$ for 10 $\mu$ m/s to 300 $\mu$ m/s, with $d_c$ values of
526	~0.1 mm. At intermediate velocities, the drastic drops in shear stress $\Delta \tau$ were 1.5-0.9
527	MPa with $\delta_c \sim 0.1$ m. For high-smectite samples with more than 50% smectite, the $a-b$
528	values at the low-V conditions are almost velocity neutral (0 to 0.003). The $\Delta \tau$ values at
529	the intermediate-V condition were 0.15-0.05 MPa with $\delta_c \sim 0.1$ m.
530	The estimated $L_c$ values for low V conditions ( $V_{eq} < 300 \ \mu m/s$ ) are 30-122 m,
531	whereas those for intermediate V conditions ( $V_{eq} > 1 \text{ mm/s}$ ) are 1.2-36.7 km (Figure 9a,
532	Table 3). As self-accelerating slip starts when <i>L</i> exceeds $L_c$ , the dependence of $L_c$ on <i>V</i>
533	leads to the following two possible cases of earthquake nucleation and rupture,
534	depending on the slip length $L$ and the slip velocity.

535	(case i) If L exceeds ~10-100 m ( $L_c$ for low V conditions) before the slip velocity
536	reaches 300 $\mu$ m/s, a rapid acceleration of slip initiates resulting in dynamic rupture
537	propagation. This situation leads to the emission of seismic waves as a result of the
538	stress drop and subsequent slip, which means a regular earthquake. It would be also
539	possible that the slip velocity after the rupture propagation may be suppressed around
540	intermediate $V$ conditions when sufficient strain energy is accumulated on a fault
541	because the shear stress of the gouge becomes minimum at the intermediate $V$ range and
542	fault strength cannot decrease further. Therefore, in some cases, the slip may result in a
543	tremor that are seismologically detected slow earthquakes (Figure 9b).
544	(case ii) If L does not reach ~10-100 m when the slip velocity reaches 300 $\mu$ m/s,
545	a coseismic rupture does not initiate and the nucleation phase continues, because $L$ is
546	much smaller than $L_c$ which becomes ~1-30 km ( $L_c$ for intermediate V conditions) and it
547	will take a longer duration for nucleating an earthquake resulting in a slow slip event
548	(SSE) like slip on the fault. If L exceeds ~1-30 km, then coseismic rupture will
549	eventually occur. In the case of the Nankai subduction zone, 30 km corresponds to the
550	entire length of the plate boundary fault beneath the outer wedge. Therefore, the
551	coseismic rupture initiation at $L > \sim 1-30$ km would generate a megathrust earthquake
552	that ruptures all the shallow part of the plate boundary fault.
553	The above two cases can be interpreted to be caused by a difference in rupture
554	velocity during the quasi-static rupture propagation: the slower the rupture velocity is,
555	the higher the probability for a long nucleation phase becomes. This means that the
556	large $L_c$ for intermediate V conditions is not necessarily required for an SSE like slip.
557	However, a larger $L_c$ at intermediate V conditions could contribute to suppress

558	acceleration of slip and elongate the duration of nucleation phase. Therefore, the
559	weakening behavior at the intermediate $V$ range could be an important factor that
560	controls a complex spectrum of slip behavior on a fault, although the physical
561	weakening mechanisms and constitutive law must be constrained in detail to be
562	implemented in numerical models. Since the transition from volcanic glass to smectite
563	occurs in the shallow part of subduction zones, and as the smectite content increases
564	with depth, the plate boundary fault would gradually be more likely to host slow
565	earthquakes, as observed in the Nankai Trough where slow earthquakes have been
566	observed beneath the accretionary prism (Figure 9c). At deeper depths where smectite
567	dehydrates into illite at $T \sim 150^{\circ}$ C or $\sigma_{eff}$ will become high, frictional properties may
568	change, or $L_c$ will decrease, which will make the plate boundary fault more prone to
569	regular earthquakes.

570







573 (a) The  $L_c$  values for different smectite contents with low velocity range (open square) 574 and intermediate velocity range (filled square). The  $d_c$  values are assumed to be 1 mm 575 for low velocity range and 0.1 m for intermediate velocity range. (b) Schematic 576 illustration for the evolutions of slip velocity that can generate slow earthquakes (see text for details). Orange and gray areas represent the critical nucleation length  $L_c$  for low 577 578 and intermediate velocity conditions, respectively. (c) A sketch of the cross section of 579 the Nankai Trough, SW Japan (Moore et al., 2009) overlaid by locations of slow 580 earthquakes and the seismogenic zone. 581

582 **Table 3.** 

583 Summary of velocity dependence at low and intermediate velocity conditions. All

584 parameters are the case for the effective normal stress of 5 MPa. "X" means that  $L_c$ 

<sup>585</sup> cannot be defined because a-b is positive.

	Low V condi	tions (RSF)		Intermediate V conditions (slip				
				weakening)				
	$(a-b)\sigma_{eff}$	$d_c$	$L_c$	Δτ	$\delta_c$	$L_c$		
Glass/Smectite	0.075 MPa	~0.1 mm	Х	0.25 MPa	~0.1 mm	Х		
= 100/0								
Glass/Smectite	-0.06 MPa	~0.1 mm	30	-1.5 MPa	~0.1 m	1.2 km		
= 85/15			m					
Glass/Smectite	-0.015 MPa	~0.1 mm	122	-0.9 MPa	~0.1 m	2.0 km		
= 70/30			m					
Glass/Smectite	0.015 MPa	~0.1 mm	Х	-0.15 MPa	~0.1 m	12.2 km		
= 50/50								
Glass/Smectite	0 MPa	~0.1 mm	Х	-0.05 MPa	~0.1 m	36.7 km		
= 30/70								
Glass/Smectite	0 MPa	~0.1 mm	Х	-0.05 MPa	~0.1 m	36.7 km		
= 0/100								

### 586

# 587 **6.** Conclusions

We performed a series of friction experiments on a range of volcanic glass smectite mixtures at different velocity conditions. For low-velocity conditions (<1 mm/s), friction coefficients decreased with increased smectite content. Samples with 15

591	and 30% smectite characteristically showed a negative velocity dependence of friction
592	coefficients. At intermediate-velocity conditions (1-3 mm/s), samples with 15 and 30%
593	smectite showed marked slip-weakening behavior with a characteristic slip distance of
594	$\sim$ 0.1 m. At high-velocity conditions (>100 mm/s), the volcanic glass exhibited dynamic
595	weakening behavior caused by thermal processes such as vaporization. Silica-gel
596	lubrication likely does not explain the observed dynamic weakening behavior. Although
597	the underlying physical processes are still undetermined, the marked weakening
598	behavior at intermediate velocity conditions will produce large critical nucleation
599	lengths on the order of 1-30 km. Such weakening behavior at intermediate velocity
600	conditions potentially induces slow earthquakes, as observed in shallow subduction
601	zones such as the Nankai Trough.

#### 603 Appendix A: Derivation of equivalent pore pressure

604 To define the equivalent pore pressure  $P_{feq}$  (equation 5), we consider the total

605 work for shear deformation similar to the derivation of equivalent velocity  $V_{eq}$ 

606 (Shimamoto & Tsutsumi, 1994). Assuming  $\mu$  does not depend on r, the shear stress at r

607 can be written as:

$$\tau(r) = \mu_{obs} \left( \sigma_n - P_{fo} \right) = \mu_{cor} \left( \sigma_n - P_f(r) \right). \#(A1)$$

608 The work for shear deformation at *r* can be described as:

$$dw = V(r)\tau(r)(2\pi r)dr = \frac{3}{4}V_{eq}(2r)\frac{r_o + r_i}{r_o^2 + r_o r_i + r_i^2}\tau(r)(2\pi r)dr, \#(A2)$$

and thus the total work on the entire shear plane is:

$$W = \int_{r_i}^{r_o} dw = \int_{r_i}^{r_o} C r^2 \tau(r) dr, \#(A3)$$

610 where

$$C = 3\pi V_{eq} \frac{r_o + r_i}{r_o^2 + r_o r_i + r_i^2} . \# (A4)$$

611 Since equation A3 can be converted as follows:

$$W = \int_{r_i}^{r_o} C r^2 \mu_{obs} (\sigma_n - P_{fo}) dr = \int_{r_i}^{r_o} C r^2 \mu_{cor} (\sigma_n - P_f(r)) dr, \#(A5)$$

612 we obtain the following equation:

$$\mu_{obs} \left(\sigma_n - P_{fo}\right) \frac{r_o^3 - r_i^3}{3} = \mu_{cor} \left(\sigma_n \frac{r_o^3 - r_i^3}{3} - \int_{r_i}^{r_o} P_f(r) r^2 dr\right) . \# (A6)$$

613 Considering the average  $P_f$  within the shear localized zone with the thickness  $w_{hs}$ , we

614 can describe  $P_f(r)$  as follows:

$$P_f(r) = \int_{z_{hs}^{bot}}^{z_{hs}^{top}} \frac{P_f(r, z)}{w_{hs}} dz. \#(A7)$$

615 If we define  $P_{feq}$  as follows:

$$P_{feq} = \frac{3}{r_o^3 - r_i^3} \int_{r_i}^{r_o} P_f(r) r^2 dr = \frac{3}{r_o^3 - r_i^3} \int_{r_i}^{r_o} \int_{z_{hs}^{bot}}^{z_{hs}^{top}} \frac{P_f(r, z) r^2}{w_{hs}} dz dr, \#(A8)$$

616 equation A6 can be converted to:

$$\mu_{obs}(\sigma_n - P_{fo}) = \mu_{cor}(\sigma_n - P_{feq}). \# (A9)$$

617 Therefore, *P<sub>feq</sub>* represents the "equivalent" pore pressure for radially heterogenous pore

618 pressure condition, which can be considered as the representative  $P_f$  value in the shear

619 localized zone.

620

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#### 631 **Open research**

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#### 634 References

- 635 Ben Abdelounis, H., Elleuch, K., Vargiolu, R., Zahouani, H., & Le Bot, A. (2009). On
- 636 the behaviour of obsidian under scratch test. *Wear*, *266*(7–8), 621–626.
- 637 https://doi.org/10.1016/j.wear.2008.07.007
- Acosta, M., Passelegue, F., Schubnel, A., & Violay, M. (2018). Dynamic weakening
- 639 during earthquakes controlled by fluid thermodynamics. *Nature Communications*,
- 640 9(1), 3074. https://doi.org/10.1038/s41467-018-05603-9
- 641 Akuhara, T., Tsuji, T., & Tonegawa, T. (2020). Overpressured Underthrust Sediment in
- 642 the Nankai Trough Forearc Inferred From Transdimensional Inversion of High-
- 643 Frequency Teleseismic Waveforms. *Geophysical Research Letters*, 47(15).
- 644 https://doi.org/10.1029/2020GL088280
- Archard, J. F. (1959). The temperature of rubbing surfaces. *Wear*, 2(6), 438–455.

646 https://doi.org/10.1016/0043-1648(59)90159-0

- 647 Aretusini, S., Meneghini, F., Spagnuolo, E., Harbord, C. W. A., & Di Toro, G. (2021).
- 648 Fluid pressurisation and earthquake propagation in the Hikurangi subduction zone.
- 649 *Nature Communications*, *12*(1), 2481. https://doi.org/10.1038/s41467-021-22805-
- 650

w

- 651 Bedford, J. D., Faulkner, D. R., & Lapusta, N. (2022). Fault rock heterogeneity can
- 652 produce fault weakness and reduce fault stability. *Nature Communications*, 13(1),
- 653 326. https://doi.org/10.1038/s41467-022-27998-2
- 654 Byerlee, J. D. (1978). Friction of rocks. Pure and Applied Geophysics PAGEOPH,
- 655 *116*(4–5), 615–626. https://doi.org/10.1007/BF00876528
- 656 Chen, J., Niemeijer, A., Yao, L., & Ma, S. (2017). Water vaporization promotes
- coseismic fluid pressurization and buffers temperature rise. *Geophysical Research Letters*, 44(5), 2177–2185. https://doi.org/10.1002/2016GL071932
- 659 Chester, F. M., Rowe, C. D., Ujiie, K., Kirkpatrick, J. D., Regalla, C., Remitti, F., et al.
- 660 (2013). Structure and Composition of the Plate-Boundary Slip Zone for the 2011
- 661 Tohoku-Oki Earthquake. *Science*, *342*(6163), 1208–1211.
- 662 https://doi.org/10.1126/science.1243719
- 663 Compton, J. S. (1991). Origin and Diagenesis of Clay Minerals in the Monterey
- 664 Formation, Santa Maria Basin Area, California. *Clays and Clay Minerals*, 39(5),
- 665 449–466. https://doi.org/10.1346/CCMN.1991.0390501
- 666 Dieterich, J. H. (1979). Modeling of rock friction: 1. Experimental results and
- 667 constitutive equations. *Journal of Geophysical Research*, 84(B5), 2161.
- 668 https://doi.org/10.1029/JB084iB05p02161
  - 40

- 669 Faulkner, D. R., Mitchell, T. M., Behnsen, J., Hirose, T., & Shimamoto, T. (2011).
- 670 Stuck in the mud? Earthquake nucleation and propagation through accretionary
- 671 forearcs. *Geophysical Research Letters*, 38(18), n/a-n/a.
- 672 https://doi.org/10.1029/2011GL048552
- 673 Faulkner, D. R., Sánchez-Roa, C., Boulton, C., den Hartog, S. A. M., Sanchez-Roa, C.,
- Boulton, C., & den Hartog, S. A. M. (2018). Pore Fluid Pressure Development in
- 675 Compacting Fault Gouge in Theory, Experiments, and Nature. *Journal of*
- 676 *Geophysical Research: Solid Earth*, *123*(1), 226–241.
- 677 https://doi.org/10.1002/2017JB015130
- 678 Goldsby, D. L., & Tullis, T. E. (2002). Low frictional strength of quartz rocks at
- 679 subseismic slip rates. *Geophysical Research Letters*, 29(17), 1844.
- 680 https://doi.org/10.1029/2002GL015240
- 681 Han, R., Kim, J.-S., Kim, C.-M., Hirose, T., Jeong, J. O., & Jeong, G. Y. (2019).
- 682 Dynamic weakening of ring faults and catastrophic caldera collapses. *Geology*,
- 683 47(2), 107–110. https://doi.org/10.1130/G45687.1
- Harders, R., Kutterolf, S., Hensen, C., Moerz, T., & Brueckmann, W. (2010). Tephra
- 685 layers: A controlling factor on submarine translational sliding? *Geochemistry*,
- 686 *Geophysics, Geosystems, 11*(5), n/a-n/a. https://doi.org/10.1029/2009GC002844
- 687 Hayashi, N., & Tsutsumi, A. (2010). Deformation textures and mechanical behavior of
- a hydrated amorphous silica formed along an experimentally produced fault in
- 689 chert. *Geophysical Research Letters*, 37(12), n/a-n/a.
- 690 https://doi.org/10.1029/2010GL042943
- Hirose, T., & Bystricky, M. (2007). Extreme dynamic weakening of faults during
  - 41

- dehydration by coseismic shear heating. *Geophysical Research Letters*, 34(14),
- 693 L14311. https://doi.org/10.1029/2007GL030049
- Hunfeld, L. B., Chen, J., Niemeijer, A. R., Ma, S., & Spiers, C. J. (2021). Seismic Slip-
- 695 Pulse Experiments Simulate Induced Earthquake Rupture in the Groningen Gas
- 696 Field. *Geophysical Research Letters*, 48(11).
- 697 https://doi.org/10.1029/2021GL092417
- 698 Ikari, M. J., Marone, C., Saffer, D. M., & Kopf, A. J. (2013). Slip weakening as a
- 699 mechanism for slow earthquakes. *Nature Geoscience*, *6*(6), 468–472.
- 700 https://doi.org/10.1038/ngeo1818
- 701 Ikari, M. J., Kopf, A. J., Hüpers, A., & Vogt, C. (2018). Lithologic control of frictional
- strength variations in subduction zone sediment inputs. *Geosphere*, 14(2), 604–625.
- 703 https://doi.org/10.1130/GES01546.1
- Kameda, J., Shimizu, M., Ujiie, K., Hirose, T., Ikari, M. J., Mori, J. J., et al. (2015).
- 705 Pelagic smectite as an important factor in tsunamigenic slip along the Japan Trench.
- 706 *Geology*, 43(2), 155–158. https://doi.org/10.1130/G35948.1
- 707 Kanagawa, K., Murayama, H., Sugita, A., Takahashi, M., Sawai, M., Furukawa, N., &
- 708 Hirose, T. (2020). Weakening of quartz rocks at subseismic slip rates due to
- frictional heating, but not to lubrication by wear materials of hydrated amorphous
- 710 silica or silica gel. *Tectonophysics*, 784(April 2019), 228429.
- 711 https://doi.org/10.1016/j.tecto.2020.228429
- 712 Kluger, M. O., Jorat, M. E., Moon, V. G., Kreiter, S., de Lange, W. P., Mörz, T., et al.
- 713 (2020). Rainfall threshold for initiating effective stress decrease and failure in
- 714 weathered tephra slopes. *Landslides*, *17*(2), 267–281.
  - 42

- 715 https://doi.org/10.1007/s10346-019-01289-2
- 716 Laberg, J. S., Strasser, M., Alves, T. M., Gao, S., Kawamura, K., Kopf, A. J., & Moore,
- 717 G. F. (2017). Internal deformation of a muddy gravity flow and its interaction with
- 718 the seafloor (site C0018 of IODP Expedition 333, Nankai Trough, SE Japan).
- 719 *Landslides*, 14(3), 849–860. https://doi.org/10.1007/s10346-016-0766-7
- 720 Lavallée, Y., Hirose, T., Kendrick, J. E., De Angelis, S., Petrakova, L., Hornby, A. J., &
- 721 Dingwell, D. B. (2014). A frictional law for volcanic ash gouge. *Earth and*
- 722 Planetary Science Letters, 400, 177–183.
- 723 https://doi.org/10.1016/j.epsl.2014.05.023
- Logan, J. M., & Rauenzahn, K. A. (1987). Frictional dependence of gouge mixtures of
- 725 quartz and montmorillonite on velocity, composition and fabric. *Tectonophysics*,

726 *144*(1–3), 87–108. https://doi.org/10.1016/0040-1951(87)90010-2

- 727 Moon, V. (2016). Halloysite behaving badly: geomechanics and slope behaviour of
- halloysite-rich soils. *Clay Minerals*, *51*(3), 517–528.
- 729 https://doi.org/10.1180/claymin.2016.051.3.09
- 730 Morrow, C. A., Moore, D. E., & Lockner, D. A. (2017). Frictional strength of wet and
- 731 dry montmorillonite. Journal of Geophysical Research: Solid Earth, 122(5), 3392–
- 732 3409. https://doi.org/10.1002/2016JB013658
- 733 Nakamura, S., Wakai, A., Umemura, J., Sugimoto, H., & Takeshi, T. (2014).
- 734 Earthquake-induced landslides: Distribution, motion and mechanisms. *Soils and*
- 735 *Foundations*, 54(4), 544–559. https://doi.org/10.1016/j.sandf.2014.06.001
- 736 Nakamura, Y. (2016). Stratigraphy, distribution, and petrographic properties of
- 737 Holocene tephras in Hokkaido, northern Japan. Quaternary International, 397, 52–
  - 43

738 62. https://doi.org/10.1016/j.quaint.2015.07.056

- 739 Nakanishi, R., Ashi, J., & Okamura, S. (2020). A dataset for distribution and
- 740 characteristics of Holocene pyroclastic fall deposits along the Pacific coasts in
- 741 western Hokkaido, Japan. *Data in Brief*, 33, 106565.
- 742 https://doi.org/10.1016/j.dib.2020.106565
- 743 Okuda, H., Ikari, M. J., Roesner, A., Stanislowski, K., Hüpers, A., Yamaguchi, A., &
- 744 Kopf, A. J. (2021). Spatial Patterns in Frictional Behavior of Sediments Along the
- 745 Kumano Transect in the Nankai Trough. Journal of Geophysical Research: Solid

746 *Earth*, *126*(11). https://doi.org/10.1029/2021JB022546

- 747 Oohashi, K., Hirose, T., & Shimamoto, T. (2013). Graphite as a lubricating agent in
- fault zones: An insight from low- to high-velocity friction experiments on a mixed
- 749 graphite-quartz gouge. Journal of Geophysical Research: Solid Earth, 118(5),
- 750 2067–2084. https://doi.org/10.1002/jgrb.50175
- 751 Oohashi, K., Hirose, T., Takahashi, M., & Tanikawa, W. (2015). Dynamic weakening
- 752 of smectite-bearing faults at intermediate velocities: Implications for subduction
- zone earthquakes. Journal of Geophysical Research: Solid Earth, 120(3), 1572–
- 754 1586. https://doi.org/10.1002/2015JB011881
- 755 Rempe, M., Smith, S. A. F., Mitchell, T. M., Hirose, T., & Di Toro, G. (2017). The
- effect of water on strain localization in calcite fault gouge sheared at seismic slip
- rates. *Journal of Structural Geology*, 97, 104–117.
- 758 https://doi.org/10.1016/j.jsg.2017.02.007
- 759 Rice, J. R. (2006). Heating and weakening of faults during earthquake slip. *Journal of*
- 760 *Geophysical Research: Solid Earth*, 111(B5), n/a-n/a.
  - 44

761 https://doi.org/10.1029/2005JB004006

- 762 Romine, W. L., Whittington, A. G., Nabelek, P. I., & Hofmeister, A. M. (2012).
- 763 Thermal diffusivity of rhyolitic glasses and melts: effects of temperature, crystals
- and dissolved water. *Bulletin of Volcanology*, 74(10), 2273–2287.
- 765 https://doi.org/10.1007/s00445-012-0661-6
- 766 Rowe, C. D., Lamothe, K., Rempe, M., Andrews, M., Mitchell, T. M., Di Toro, G., et al.
- 767 (2019). Earthquake lubrication and healing explained by amorphous nanosilica.
- 768 *Nature Communications*, 10(1), 1–11. https://doi.org/10.1038/s41467-018-08238-y
- Ruina, A. L. (1983). Slip instability and state variable friction laws. *Journal of*
- 770 *Geophysical Research: Solid Earth*, 88(B12), 10359–10370.
- 771 https://doi.org/10.1029/JB088iB12p10359
- 772 Screaton, E. J., Kimura, G., Curewitz, D., Moore, G. F., Chester, F. M., Fabbri, O., et al.
- 773 (2009). Interactions between deformation and fluids in the frontal thrust region of
- the NanTroSEIZE transect offshore the Kii Peninsula, Japan: Results from IODP
- Expedition 316 Sites C0006 and C0007. *Geochemistry, Geophysics, Geosystems*,

776 *10*(12), n/a-n/a. https://doi.org/10.1029/2009GC002713

- 777 Scudder, R. P., Murray, R. W., Kutterolf, S., Schindlbeck, J. C., Underwood, M. B., &
- 778 Wang, K.-L. (2018). Sedimentary inputs to the Nankai subduction zone: The
- importance of dispersed ash. *Geosphere*, *14*(4), 1451–1467.
- 780 https://doi.org/10.1130/GES01558.1
- 781 Shimamoto, T., & Tsutsumi, A. (1994). A new rotary-shear high-speed frictional testing
- 782 machine: its basic design and scope of research. *Journal of the Tectonic Research*
- 783 *Group, Japan.* Retrieved from https://ci.nii.ac.jp/naid/80008741644/
  - 45

784	Shimizu, O., & Ono, M. (2016). Relationship of tephra stratigraphy and hydraulic
785	conductivity with slide depth in rainfall-induced shallow landslides in Aso
786	Volcano, Japan. Landslides, 13(3), 577-582. https://doi.org/10.1007/s10346-015-
787	0666-2
788	Strasser, M., Henry, P., Kanamatsu, T., Thu, M. K., & Moore, G. F. (2012). Scientific
789	Drilling of Mass-Transport Deposits in the Nankai Accretionary Wedge: First
790	Results from IODP Expedition 333. In Submarine Mass Movements and Their
791	Consequences (pp. 671-681). Dordrecht: Springer Netherlands.
792	https://doi.org/10.1007/978-94-007-2162-3_60
793	Takahashi, M., Mizoguchi, K., Kitamura, K., & Masuda, K. (2007). Effects of clay
794	content on the frictional strength and fluid transport property of faults. Journal of
795	Geophysical Research, 112(B8), B08206. https://doi.org/10.1029/2006JB004678
796	Tanikawa, W., Mukoyoshi, H., Tadai, O., Hirose, T., Tsutsumi, A., & Lin, W. (2012).
797	Velocity dependence of shear-induced permeability associated with frictional
798	behavior in fault zones of the Nankai subduction zone. Journal of Geophysical
799	Research: Solid Earth, 117(5), 1-16. https://doi.org/10.1029/2011JB008956
800	Tanikawa, W., Ishikawa, T., Honda, G., Hirono, T., & Tadai, O. (2015). Trace element
801	anomaly in fault rock induced by coseismic hydrothermal reactions reproduced in
802	laboratory friction experiments. Geophysical Research Letters, 42(9), 3210-3217.
803	https://doi.org/10.1002/2015GL063195
804	Tembe, S., Lockner, D. A., & Wong, TF. (2010). Effect of clay content and
805	mineralogy on frictional sliding behavior of simulated gouges: Binary and ternary
806	mixtures of quartz, illite, and montmorillonite. Journal of Geophysical Research,

- 807 *115*(B3), B03416. https://doi.org/10.1029/2009JB006383
- 808 Di Toro, G., Goldsby, D. L., & Tullis, T. E. (2004). Friction falls towards zero in quartz
- rock as slip velocity approaches seismic rates. *Nature*, *427*(6973), 436–439.
- 810 https://doi.org/10.1038/nature02249
- 811 Ujiie, K., & Tsutsumi, A. (2010). High-velocity frictional properties of clay-rich fault
- gouge in a megasplay fault zone, Nankai subduction zone. *Geophysical Research*
- 813 *Letters*, *37*(24), L24310. https://doi.org/10.1029/2010GL046002
- Ujiie, K., Tanaka, H., Saito, T., Tsutsumi, A., Mori, J. J., Kameda, J., et al. (2013). Low
- 815 Coseismic Shear Stress on the Tohoku-Oki Megathrust Determined from
- 816 Laboratory Experiments. *Science*, *342*(6163), 1211–1214.
- 817 https://doi.org/10.1126/science.1243485
- 818 Wiemer, G., & Kopf, A. J. (2015). Altered marine tephra deposits as potential slope
- failure planes? *Geo-Marine Letters*, *35*(4), 305–314.
- 820 https://doi.org/10.1007/s00367-015-0408-4
- 821 Wiemer, G., & Kopf, A. J. (2017). On the role of volcanic ash deposits as preferential
- submarine slope failure planes. *Landslides*, *14*(1), 223–232.
- 823 https://doi.org/10.1007/s10346-016-0706-6
- Wiemer, G., Moernaut, J., Stark, N., Kempf, P., De Batist, M., Pino, M., et al. (2015).
- 825 The role of sediment composition and behavior under dynamic loading conditions
- 826 on slope failure initiation: a study of a subaqueous landslide in earthquake-prone
- 827 South-Central Chile. *International Journal of Earth Sciences*, *104*(5), 1439–1457.
- 828 https://doi.org/10.1007/s00531-015-1144-8
- 829 Yonekura, K. (2015). Rock properties and material selection for blade manufacture in
  - 47

- 830 upper paleolithic Japan. *Lithic Technology*, *40*(2), 85–93.
- 831 https://doi.org/10.1179/2051618515Y.0000000001

Figure 1.



Figure 2.



Figure 3.



Figure 4.



Figure 5.



Figure 6.

## PHV596, glass = 100%, 10 µm/s – 3 mm/s



#### PHV592, smectite = 15%, 10 µm/s - 3 mm/s



PHV589/590, smectite = 70/100%, 10  $\mu$ m/s – 300 mm/s



Figure 7.



Figure 8.






Figure 9.

