# Dynamical Heating in the Martian Thermosphere

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#### Abstract

Dynamical heating and cooling are prominent features of planetary atmospheres resulting in thermospheric structures on Venus, Earth and Mars. The purpose of this study is to determine the location and amplitude of localized heating regions in the Martian thermosphere, confirm that they occur in regions of wind convergence, and to compare the observed dynamical heating with that predicted by a global thermospheric model. This investigation uses several years of data from the NASA Mars Atmosphere and Volatile EvolutioN (MAVEN) mission including observations made by the Neutral Gas and Ion Mass Spectrometer (NGIMS) as well as the Extreme Ultraviolet Monitor (EUVM). Specifically, the analysis focuses on several years of horizontal wind, temperature, and composition data. EUVM measurements provide a solar forcing context for the neutral thermosphere datasets and aid in the statistical analysis. Statistical results are compared with two versions of the Mars Global Ionosphere Thermosphere (M-GITM) global circulation model; one that that includes gravity wave parametrization and a version without gravity wave effects. Data analysis indicates that heating features exists around 2-3 and 17-18 local solar time. These locations coincide with regions of converging winds and are in better agreement with M GITM when gravity wave parametrization is included in the model. A migrating oscillation in the observed wind field also results in convergence and a density enhancement near 15 local time. While a similar oscillation is reproduced by the model, the amplitude is much lower than observed and may be a result of modeled zonal winds that are too low.

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10	
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12	Key Points:
13 14	• Equinox analysis confirms the presence of dynamical heating post-midnight and near the dusk terminator
15 16	• Including gravity wave parametrization in the Mars Global Ionosphere Thermosphere Model improves specification of dynamical heating
17 18	• A global modulation of the thermospheric wind pattern leads to a migrating tide in dynamical effects with a 4 hour local time wavelength
19 20	

## 21 Abstract

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40

## 41 Plain Language Summary

The upper reaches of planetary atmospheres act as the interface between the planet and the 42 interplanetary space environment. On Mars this region is responsible for aerodynamic drag on 43 satellites orbiting the planet. It is also part of the pathway for molecules escaping to space and 44 45 contributing to the loss of most of the Martian atmosphere and water over the course of billions of years. The motion and structure of the upper atmospheric region is important to understanding and 46 predicting its behavior. We analyzed wind and temperature data collected in the upper atmosphere 47 of Mars by a NASA satellite called the Mars Atmosphere and Volatile EvolutioN spacecraft and 48 compared the results to a model of the Mars atmosphere. The results indicate that locations of 49 converging and diverging winds exist throughout the atmosphere and contribute to its temperature 50 and density structure. These locations cause hot spots in the upper atmospheres, particularly on the 51 night side. The wind pattern is also responsible the presence of large ripples or wave-like features 52 53 in the dayside atmosphere.

## 54 **1 Introduction**

The global structure of planetary thermospheres can shed light on their energy budgets and 55 56 transport processes, as well as on the ionospheres to which they couple. Thermospheric structure also affects aerodynamic drag experienced by spacecraft flying at sufficiently low altitudes. While 57 this structure is driven by photochemical, radiative, and dynamical processes, as well as by 58 59 coupling with the lower atmosphere through a variety of mechanisms including gravity waves, it is the dynamical effects that are more difficult to observe and not as well understood. On Earth, 60 dynamical thermospheric effects can cause cooling in the polar thermosphere (Crowley et al., 61 62 1995; Schoendorf et al., 1996) for example while at Venus, night-time heating is associated with adiabatic effects in a region of converging thermospheric winds (Brecht et al., 2011). In fact, this 63

dynamical heating effect is the reason that the nightside thermosphere can maintain temperatures
significantly above absolute zero in spite of the long duration of the Venutian night. On Mars,
dynamical effects lead to regions of significant thermospheric temperature and density
enhancements, especially near the dawn and dusk terminators (Forbes & Moudden, 2009; Pilinski
et al., 2018).

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Dynamical processes can lead to both heating and cooling with heating effects being associated with the presence of converging winds. Wind convergence can result from zonal or meridional wind direction reversals or by a rapid slowing of the horizontal winds. This in turn leads to a downwelling in the thermosphere and, often, adiabatic heating.

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The presence of dynamical heating regions or thermospheric "heat islands" on Mars was first theorized based on results of global circulation modeling (S. W. Bougher et al., 1990). Pilinski et al. (Pilinski et al., 2018) presented the first empirical evidence for the existence of heat islands using data from the Mars Atmosphere and Volatile Escape (MAVEN) mission (Jakosky et al., 2015).

80

In this paper, we explore the global structure of thermospheric temperature and winds at Mars in 81 order to map the location and amplitude of dynamical effects. These properties will also be 82 compared with two versions of a Mars global circulation model, one with and one without gravity 83 wave parametrization. The sections below describe the theory of dynamical heating and statistical 84 methods used to analyze MAVEN data. This is followed by a description of the global circulation 85 model. Datatypes used in the analysis are then reviewed. Results are presented in three parts: (a) a 86 broad multi-seasonal analysis of winds and temperatures, (b) a statistical analysis focusing on 87 equinox, and (c) an empirical reconstruction of equinox conditions at constant pressure and EUV 88 89 forcing. Model comparisons are made throughout the results section. Next we discuss the results in the context of global thermospheric structure and summarize the main conclusions. 90 91

# 92 2 Dynamical Heating Background

93 The contribution of dynamical heating to the overall energy budget can be examined through the 94 horizontal and vertical thermodynamic equations as described by Ridley et al. (Ridley et al., 2006) 95 and adopted here for the Martian thermosphere. The horizontal temperature equation is

96

$$0 = \frac{\partial \tau}{\partial t} + \frac{u_{\phi}}{r \cos \theta} \frac{\partial \tau}{\partial \phi} + u_{\theta} \frac{\partial \tau}{\partial \theta} + (\gamma - 1)\tau \left( \frac{1}{r} \frac{\partial u_{\theta}}{\partial \theta} + \frac{1}{r \cos \theta} \frac{\partial u_{\phi}}{\partial \phi} - \frac{u_{\theta} \tan \theta}{r} \right)$$
(1)

97

98 where the normalized temperature,  $\tau$ , is defined as 99

 $\tau = \frac{k_B}{\bar{m}}T\tag{2}$ 

100

- 101 where T is the neutral temperature,  $\overline{m}$  is the mean molecular mass,  $k_B$  is the Boltzmann constant, r
- 102 is the distance from the center of the planet,  $\theta$  is the latitude,  $\phi$  is the longitude,  $u_{\theta}$  and  $u_{\phi}$  are the
- 103 meridional and zonal wind components respectively, and  $\gamma$  is the ratio of specific heats, equal to

5/3. The normalized temperature is equivalent to the total pressure divided by the mass density(Ridley et al., 2006).

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107 The vertical temperature equation is

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$$\frac{\partial \tau}{\partial t} + u_r \frac{\partial \tau}{\partial r} + (\gamma - 1)\tau \left(\frac{2u_r}{r} + \frac{\partial u_r}{\partial r}\right) = \frac{k_B}{c_v \rho \bar{m}} Q_{tot}$$
(3)

109

where  $u_r$  is the vertical wind,  $c_v$  is the specific heat at constant volume,  $\rho$  is the mass density and  $Q_{tot}$  is the total energy source term. The specific heat is computed based on the weighted average of specific heats of individual species weighted by their number densities. Note that the  $2u_r/r$  term is due to divergence of velocities along the radial direction in a spherical coordinate system. The energy source term can be broken down into extreme UV (EUV), infrared (IR), 15 µm (CO<sub>2</sub> LTE and non-LTE terms), and conduction terms as follows

116

$$Q_{tot} = Q_{EUV} + Q_{IR} + Q_{15um} + \frac{\partial}{\partial r} \left( \frac{\partial T}{\partial r} \left( \kappa_c + \kappa_{eddy} \right) \right)$$
(4)

117

118 where  $\kappa_c$  and  $\kappa_{eddy}$  are the thermal and eddy diffusion heat conductivity respectively.

119

The net dynamical heating, shown below, is the combination of the horizontal hydrodynamic advection (subscript A) and horizontal divergence (subscript B) from equation 1, as well as vertical hydrodynamic advection (subscript C) and vertical divergence (subscript D) from equation 3. M-GITM solves equation 1, the left side of equation 3, and the source terms in the right side of equation 3 independently and the resulting dynamical heating terms can be isolated in postprocessing.

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127

$$Q_{TDYN} \stackrel{\text{def}}{=} \left[ \frac{u_{\phi}}{r \cos \theta} \frac{\partial \tau}{\partial \phi} + u_{\theta} \frac{\partial \tau}{\partial \theta} \right]_{A} + \left[ (\gamma - 1)\tau \left( \frac{1}{r} \frac{\partial u_{\theta}}{\partial \theta} + \frac{1}{r \cos \theta} \frac{\partial u_{\phi}}{\partial \phi} \right) \right]_{B} + \left[ u_{r} \frac{\partial \tau}{\partial r} \right]_{C}$$

$$+ \left[ (\gamma - 1)\tau \left( \frac{2u_{r}}{r} + \frac{\partial u_{r}}{\partial r} \right) \right]_{D}$$
(5)

128

Vertical neutral winds are not measured by MAVEN and therefore, the vertical terms C and D will be omitted when analyzing observed wind statistics using equation 5. The anticipated error is expected to be acceptable for the purposes of this analysis and is discussed in section 3.

#### 132 **3 Methods**

133 3.1 Statistical Analysis of Temperatures and Neutral Densities

134 Observations of Argon densities made by NGIMS are converted to neutral temperatures  $(T_n)$  using

the approach described by Snowden (Snowden et al., 2013) and used previously by (Pilinski et al.,

- 136 2018). These temperatures are combined with neutral observations of CO<sub>2</sub> and atomic oxygen (O)
- 137 number densities to compute a neutral pressure at each satellite location. Neutral temperatures are
- 138 analyzed for both descending and ascending portions of the MAVEN orbit due to the rapid nature

139 of Argon transmission through the NGIMS inlet. Neutral densities and pressures are only analyzed

- on the descending (inbound) segments of each orbit as the outbound CO<sub>2</sub> and O data is distorted
   by the buildup of these molecules on the instrument walls.
- 142

143 Several types of statistical analysis are performed using the resulting temperatures and pressures.

- 144 The first is a straightforward binning of existing data in season, local time, latitude, and altitude.
- 145 Unless otherwise noted, bins are 20 km wide in altitude, 20° wide in latitude, and 3 hours wide in 146 local solar time (LST)
- 146 local solar time (LST).
- 147

A variation of this binning is performed for the horizontal wind data in season ( $L_s$  angle), latitude, 148 and local time only. This is because horizontal winds are not expected to significantly vary with 149 altitude due to the high viscosity of the thermosphere in this region. During equinox, the modeled 150 wind pattern is approximately symmetric about the equator, particularly at low to mid latitudes. 151 As will be shown later, this is mostly true for the observed winds as well. Since wind data is 152 collected in intermittent campaigns lasting several days, taking advantage of this symmetry allows 153 for more complete coverage of the hemispheric wind pattern than would otherwise be possible. 154 The meridional winds in this case are converted into equatorward/anti-equatorward directions and, 155 along with the zonal component, are reflected around the equator. Since aggregating data for 156 equinox conditions includes a broad swath of  $L_s$  angles, it is possible that the line of symmetry will 157 depart slightly from the ideal of 0° latitude in IAU coordinates. Therefore, we perform data binning 158 Mars-sun-orbit (MSO) coordinates in order to at least partially compensate for this effect. The 159 "equator" in MSO coordinates follows the subsolar point which is the location of peak solar 160 heating and orients the diurnal wind pattern. We have confirmed this symmetry by comparing the 161 M-GITM winds in MSO coordinates at 200 km at various  $L_s$  angles. Once the equinoctial wind 162 statistics are compiled, a series of Fourier curve fits to the low-mid latitude meridional and zonal 163 164 wind components is used to compute an approximate low/mid-latitude wind-field divergence. The resulting average wind field is used to estimate divergence and to compute the dynamical heating 165 based on equation 5. Neglecting the vertical components is expected to cause an acceptable level 166 of error. To confirm this, a similar analysis is performed using model winds and temperatures from 167 M-GITM<sub>GW</sub>. Figure 1 shows the dynamical heating term at three pressure levels using all terms of 168 equation 5 (black) and using only the horizontal terms (red). 169

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Figure 1: Dynamical heating computed using M-GITM<sub>GW</sub> winds and temperatures at three

173 *different pressure levels in the thermosphere. The black lines represent dynamical heating using* 

- both horizontal and vertical advection and divergence terms. The red lines represent the
- 175 *dynamical heating computed using only the horizontal terms.*
- 176

Figure 1 demonstrates that omitting the vertical winds, while not inconsequential, still provides a reasonable representation of the dynamical heating (QD). The qualitative nature of the QD diurnal structure is also well represented.

179 180

The Mars thermosphere is highly variable at multiple temporal and spatial scales ranging from 181 gravity waves (GW) and tides (S. Bougher et al., 2015; England et al., 2017; Medvedev et al., 182 2011; Thaller et al., 2021) to EUV driven variability over the solar cycle (E.M.B. Thiemann et al., 183 2018) and due to the Mars orbital position transitioning from perihelion to aphelion. Furthermore 184 steep density gradients exist (Pilinski et al., 2018) across the terminator. Characterizing the spatial 185 variability of the thermosphere observed from a satellite in a processing, elliptical orbit is therefore 186 challenging due to the potential for aliasing the various sources of spatial and temporal variability. 187 To reduce the potential for aliasing, low-to-mid latitude neutral temperatures will be divided into 188 local time and pressure level bins. Each bin will include several days of data to reduce any GW 189 and nonmigrating tide contribution in the binned results. The neutral temperature response to EUV 190 heating within each bin can then be analyzed and controlled for. The Lyman alpha channel of the 191 EUVM instrument (F. G. Eparvier et al., 2015) serves as a proxy for EUV heating. For bins that 192 contain sufficiently diverse EUV sampling (spanning at least 0.005W/m<sup>2</sup>), a linear fit is used to 193 record the EUV- $T_n$  relationship. To evaluate changes in the potential energy (vertical movement 194 of pressure levels), the pressure level altitude  $(P_{height})$  is fit using both a linear and nonlinear fits. 195 The nonlinear fit represents the diminishing increase in  $P_{\text{height}}$  at higher EUV but requires that the 196 data in the bin includes high enough EUV values to be meaningful. Organizing the data by 197 pressure-level also allows any spatial differences in sampling within each bin to be reduced as 198 does the  $\pm 50^{\circ}$  degree MSO cutoff. In the last stage of our analysis, limiting the dataset to L<sub>s</sub> angles 199 of -45° to 45° and 135° to 225° (vernal and autumnal equinoxes combined) limits the seasonal 200 variability within the dataset and enables the assumption of approximate hemispheric symmetry. 201 The resulting set of linear and nonlinear fits for each pressure level and local time constitute a 202 simple empirical  $T_{\rm n}$  model that can represent the diurnal configuration of the low-latitude, equinox 203 thermosphere by interpolating to a fixed EUV irradiance. 204 205

- 206 3.2 M-GITM
- 207
- 208 Standard M-GITM Implementation

The Mars Global Ionosphere Thermosphere Model (M-GITM) is a model framework that 209 combines the terrestrial GITM source code (e.g., Ridley et al., 2006) with Mars fundamental 210 211 physical parameters, ion-neutral chemistry, and key radiative processes in order to capture the basic observed features of the thermal, compositional, and dynamical structure of the Mars 212 atmosphere from the ground to ~250 km e.g. (S. W. Bougher, Pawlowski, et al., 2015). The M-213 214 GITM model is typically setup with a 5x5° horizontal resolution and 2.5 km vertical resolution, and run using a ~2-second time step. Model physics and simulated features of the upper 215 atmosphere (above ~80 km) are investigated and characterized in most simulations of the M-GITM 216 217 model conducted to date. These upper atmosphere studies directly benefit from the lower to upper atmosphere coupling made possible with this whole atmosphere framework. 218

219

The modern M-GITM code currently simulates the thermal, dynamical and chemical drivers of the Martian atmosphere all the way to the surface (S. W. Bougher, Pawlowski, et al., 2015). For

the Mars lower atmosphere ( $\sim 0-80$  km), a state-of-the-art correlated-k radiation code was adapted 222 from the NASA Ames Mars General Circulation Model (MGCM) (Haberle et al., 1999) for 223 incorporation into M-GITM. This provides solar heating (long and short wavelength), seasonally 224 225 variable aerosol heating, and  $CO_2 15$ -µm cooling in the local thermal equilibrium (LTE) region of the Mars atmosphere (below ~80 km). In addition, dust opacity distributions are commonly 226 prescribed based upon empirical dust opacity maps obtained from several Martian years of 227 measurements (Smith, 2004, 2009). For the Mars upper atmosphere ( $\sim 80 - 250$  km), a modern 228 CO<sub>2</sub> NLTE 15-µm cooling scheme has recently been implemented (González-Galindo et al., 229 2013), enabling interactive CO<sub>2</sub> cooling to be simulated as atomic O abundances vary. M-GITM 230 utilizes MAVEN Extreme Ultraviolet Monitor (EUVM) L3 solar EUV and UV fluxes taken from 231 the Flare Irradiance Spectral Model for Mars (FISM-M) empirical model outputs on a daily 232 cadence (Thiemann et al., 2017). Subsequently, M-GITM thermospheric heating, dissociation and 233 ionization rates are simulated at each time step (e.g. Bougher et al., 2015). 234

235

236 Corresponding prognostic fields for neutral temperatures, neutral densities (CO<sub>2</sub>, CO, N<sub>2</sub>, O, O<sub>2</sub>,

- He, Ar, N<sup>4</sup>S) and photochemical ion densities  $(CO_2^+, O_2^+, O^+, NO^+, N_2^+)$ , plus 3-component neutral
- winds are computed on a three-dimensional regular grid. For application in this paper, FISM-M fluxes are used for driving new M-GITM simulations for the 3-cardinal seasons spanning MY33
- 239 fluxes are used for driving new M-G 240 (Ls =  $90^\circ$ ,  $180^\circ$  and  $270^\circ$ ).
  - 241

An evolving suite of M-GITM simulations have been compared with MAVEN measurements 242 obtained during its first three years of operations. This includes sampling during six Deep Dip 243 campaigns (Bougher et al., 2015b; Bougher et al., 2015c; Zurek et al., 2017) as well as dayside 244 science orbits (Bougher et al., 2017a; Bougher et al., 2017b). MAVEN NGIMS, Imaging 245 Ultraviolet Spectrograph (IUVS) and Accelerometer (ACC) measurements have been used to 246 validate the M-GITM code. For example, dayside (SZA < 60 degrees) NGIMS densities and 247 extracted temperatures were compared with corresponding M-GITM fields along Deep Dip 2 orbit 248 trajectories (Bougher et al., 2015b; Bougher et al., 2015c). These comparisons revealed that M-249 GITM dayside temperatures closely match NGIMS derived temperatures at low SZA. M-GITM 250 also produced significant dynamical heating and warm temperature bulges at the evening 251 terminator that were subsequently discovered in NGIMS datasets (Pilinski et al., 2018). By 252 contrast, the maintenance of the nightside density and temperature structure is still poorly 253 understood. For instance, Zurek et al. (Zurek et al., 2017) reported that M-GITM versus MAVEN 254 ACC mass density comparisons can be quite good from noon to the dusk terminator, but are poor 255 from midnight toward the dawn terminator. The latter is likely due to incomplete M-GITM physics 256 addressing gravity wave processes impacting global winds and the thermospheric circulation 257 patterns throughout all seasons. A version of M-GITM addressing gravity wave effects has been 258 developed by Roeten et al. (Roeten et al., 2022), and is also reviewed briefly at the end of this 259 260 section.

261

The significant evening terminator heat island features in the Martian thermosphere have been predicted by various three-dimensional global circulation models (GCMs) prior to the MAVEN

mission and NGIMS measurements (above). The coupled NASA Ames Mars Global Circulation

265 Model (MGCM)- NCAR Mars Thermospheric General Circulation Model (MTGCM) framework

266 initially predicted strong evening terminator heat island features, especially at Equinox and

267 Perihelion seasonal conditions at low to mid-latitudes (e.g. Bougher et al., 2008; Valeille et al.,

2009). In addition, a comparison study investigated temperature and wind features simulated by 268 the MGCM-MTGCM and the LMD-MGCM codes for common input conditions (González-269 Galindo et al., 2010). Both models produced strong dynamical heating at the evening terminator 270 271 giving rise to significant warming (in excess of ~30-50°K over what was otherwise expected). Weaker morning terminator warming was also revealed by each model. Two model dynamical 272 terms were compared to confirm the sources of this heating. Finally, strong evening terminator 273 dynamical heating was simulated for Aphelion, Equinox and Perihelion seasonal conditions by 274 initial M-GITM simulations prior to MAVEN (Bougher et al., 2015). All subsequent M-GITM 275 simulations have reproduced similar strong evening terminator heat island features. 276

277

Post-processing tools are typically used to visualize heat balance terms extracted from the M-278 GITM code thermal equation sources during model runtime. A separate history file is assembled 279 to contain radiative and dynamical terms taken from the code at 1-hour intervals throughout a 280 given simulation. These terms include: (a) EUV and UV heating rate, (b) NLTE CO<sub>2</sub> 15-µm 281 cooling rate, (c) molecular and eddy thermal conduction, (d) LTE and NLTE near-IR heating rate, 282 and (e) dynamical heating/cooling terms. For the latter, a combination of 4-dynamical terms gives 283 rise to the net dynamical heating rate: (a) horizontal hydrodynamic advection, (b) vertical 284 hydrodynamic advection, (c) horizontal divergence, and (d) vertical divergence. Two of these 285 terms (vertical advection and vertical divergence) are new to the M-GITM 5-moment equation 286 formulation, since the vertical momentum equation is solved explicitly in the GITM framework. 287 This means that the M-GITM solves the vertical and horizontal momentum equations separately 288 (splitting technique). Ultimately, vertical profiles of heating/cooling terms are combined into 289 single location (latitude/local time) plots for easy display of the steady-state balances of these 290 terms giving rise to the corresponding vertical temperature profile. 291

292 293

# 294 *M-GITM with a Gravity Wave Parametrization Scheme*

Recently, a modern whole atmosphere, nonlinear, non-orographic gravity wave (GW) parameterization scheme has been added to the M-GITM model in order to better account for the effects of GWs in the Martian thermosphere. Throughout this paper, we will refer to this version of the model as M-GITM<sub>GW</sub>. Results from new M-GITM<sub>GW</sub> simulations which include this GW scheme indicate the effects of GWs can be significant in the upper atmosphere, impacting wind speeds as well as the temperature structure at these altitudes (Roeten et al., 2022).

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These changes occur since the scheme allows M-GITM<sub>GW</sub> to account for both momentum 302 deposited by GWs as well as heating directly produced by wave dissipation and heating/cooling 303 corresponding to the divergence of wave-induced heat flux. Indirect changes to dynamical heating 304 can also result from changes in wind speed and the general circulation with the addition of the GW 305 drag in the forcing terms. It follows that an impact on the simulated dynamical heat island features 306 near the terminators might also be expected. To test this, new M-GITM<sub>GW</sub> simulations run over 307 the same time period were completed. Except for the addition of the GW scheme, all other key 308 parameters used in the M-GITM simulations were unchanged. 309

310

311 The GW scheme used in M-GITM<sub>GW</sub> was originally designed for Earth's upper atmosphere (Yiğit

et al., 2008) and has since been modified for use in Mars GCMS (Medvedev et al., 2015; Medvedev & Yiğit, 2012). It is specifically prescribed for the effects of non-orographic GWs,

includes nonlinear interactions between harmonics, and is appropriate for wave propagation and 314 315 dissipation in the Martian upper atmosphere. The parameterization in the model calculates the evolution of a spectrum of GWs as they propagate upward from a source level in the lower 316 317 atmosphere to the top of the model. The specified GW spectrum is an empirical Gaussian distribution that has been used previously in Earth and Mars GCMs, defined in terms of maximum 318 horizontal momentum fluxes at the source altitude as a function of horizontal phase speeds. For 319 the simulations run here, a source flux of  $0.0025 \text{ m}^2/\text{s}^2$  was used with a maximum phase speed of 320  $\pm 80$  m/s, and a source height of ~9 km. A horizontal wavelength of 300 km is used, which is 321 consistent with the findings of Creasey et al., 2006, Terada et al., 2017, and Siddle et al., 2019, 322 (Creasey et al., 2006; Siddle et al., 2019; Terada et al., 2017) among others, assuming the wave 323 structures they observed in the Martian thermosphere were due to horizontal variations. A more 324 detailed description of the GW scheme currently being used in M-GITM can be found in Roeten 325 et al., 2022 (Roeten et al., 2022), and the references therein. 326

327

## 328 **4 Data**

The MAVEN spacecraft was launched on November 18th, 2013 on a mission to explore the loss of 329 the Martian atmosphere to space. It arrived at Mars on September 21st, 2014 and began science 330 operations in October 2014. The MAVEN spacecraft is nominally in an elliptical, 150 km x 6000 331 km, orbit with a 4.5-hour period and an inclination of 75° with respect to the Martian equator. 332 Since August 2020, the MAVEN periapsis is kept closer to 200 km altitude. MAVEN carries a 333 number of neutral and plasma instruments to characterize the ionosphere, solar wind, neutral 334 atmosphere, magnetic field, and solar EUV forcing. An in-depth review of MAVEN and its 335 mission is provided by Jakosky et al. (Jakosky et al., 2015). In this work, we analyze MAVEN 336 Neutral Gas and Ion Spectrometer (NGIMS) data collected from October 2014 through August 337 2020. The Mars thermosphere is heated by solar extreme ultraviolet irradiance and we use data 338 from the Extreme Ultra-Violet Monitor (EUVM) on MAVEN to organize some of the analysis 339 described below. 340

- 341 4.1 Neutral Density and Temperature Data
- Neutral densities and composition are measured by the Neutral Gas and Ion Mass Spectrometer 342 (NGIMS) which is described in detail by (Mahaffy et al., 2015). Neutral density measurements 343 can be used to compute neutral temperature and pressure as described in a previous section. In this 344 paper, we use the closed source neutral mode of NGIMS as it has a higher signal to noise ratio. 345 Non-volatile species such as atomic oxygen and CO<sub>2</sub> tend to undergo reactions on the surfaces of 346 the accommodation chamber and transfer tube. The effects of these reactions are most prominent 347 during the ascending pass of each orbit where the increased "ram pressure" at periapsis leads to 348 large values of surface coverage for non-volatile species. For this reason, we only use closed source 349 measurements of O and CO<sub>2</sub> on the descending sides of each MAVEN orbit. 350
- 351

NGIMS measurements of Argon density are converted to neutral temperature. In calculating neutral temperature, this study assumed hydrostatic equilibrium along the argon density-altitude profiles and applied a method described in (Snowden et al., 2013). In the altitude region of interest,

argon has sufficient signal-to-noise ratio for this method and is expected to represent the  $T_n$ 

- adequately (Stone et al., 2018).
- 357

- Estimates of neutral temperature and the number densities of the primary neutral species (CO<sub>2</sub> and
- 359 O) are combined to compute neutral pressure.

$$\hat{P} = k_B (n_{44} + n_{16}) T_n \tag{6}$$

360

The density data corresponds to version 01 and revision 02 (v01\_r02) of the NGIMS L1b data

362 products found on the NASA PDS server.

363 4.2 Wind data

In addition to sampling neutral and ion species to determine the composition of the Martian upper atmosphere, the NGIMS instrument has another mode of operation that allows it to measure insitu thermospheric horizontal neutral wind velocities. During this mode, NGIMS' typical data collection is paused as the instrument platform sweeps back and forth by  $\pm 8^{\circ}$  across the spacecraft ram direction. Wind velocities are then extracted from the modulations observed in the neutral and ion fluxes as the pointing direction of the instrument varies. The technique by which NGIMS measures wind velocities is described in detail in (Benna et al., 2019).

371

Wind measurements take place along the spacecraft's track through the thermosphere at altitudes 372 373 ranging from ~140-220 km. This measurement technique assumes vertical winds are negligible, which for typical driving conditions is believed to be a reasonable assumption (Bougher et al., 374 2015). It is also assumed that winds do not vary significantly within the 30 seconds it takes for 375 the instrument to complete a full motion cycle. Uncertainties in the measurements are primarily 376 due to potential errors in the reconstructed ephemeris of the spacecraft trajectory and direction of 377 the NGIMS boresight, due to the energy resolution of the mass filter in the instrument, and in 378 379 counting statistics (Benna et al., 2019). Uncertainties typical of the along- and across- track wind magnitudes are 20 m/s and 6 m/s, respectively (Benna et al., 2019). 380

381

Wind observations most commonly take place within campaigns occurring over 2-3 days every month. Within a campaign, 5-10 consecutive orbits of neutral wind measurements are taken. Due to the nature of MAVEN's orbit, during an individual campaign, each orbit of wind observations tracks along nearly the same local times, latitudes, and altitudes, but different longitudes. Note that the maximum aerographic latitude MAVEN flies over is 75°.

387

This analysis relies on neutral wind data collected between March 2016 through April 2021. During this period, campaigns typically occurred at least once a month, with a few exceptions of

- During this period, campaigns typically occurred at least once a month, with a few exceptions of larger gaps between campaigns (i.e., gaps from October-December 2016 and February-April,
- 2019). The neutral winds dataset used is a NGIMS Level 3 data product (v03 r01).
- 392 4.3 EUV Irradiance Data

393 Irradiance observations are obtained from MAVEN Extreme UltraViolet Monitor (EUVM) (F. G.

Eparvier et al., 2015). The EUVM dataset nominally provides solar irradiance once per orbit in

three spectral bands including the Lyman-alpha and 17-22 nm bands. In this work, we use Lyman-

alpha measurements from version 13, revision 1 of the level 2b EUVM product (v13\_r01). This

irradiance product is provided once per orbit and is interpolated to the time of interest.

#### 398 **5 Results**

#### 399 5.1 Statistical Results Based on Season

Seasonal statistics for the median observed and M-GITM (without GW parametrization) 400 temperatures at 165 km are presented in Figure 2 for aphelion (northern summer), equinox, and 401 perihelion (southern summer) conditions in the top, middle, and bottom panels respectively. Note 402 that data from both autumnal and vernal equinoxes are combined in panel 2c. The left panels (a, c, 403 e) are based on a simple binning of the MAVEN data without controlling for EUV conditions. The 404 right hand panels (b, d, f) are M-GITM results at the equivalent altitude and season. Crosses 405 indicate the location of the terminator at 165 km altitude while the asterisk is the location of the 406 subsolar point. Starting with aphelion (panels a and b), we note that the warmest global 407 temperatures occur away from the subsolar point in the winter (southern) hemisphere. The 408 statistical results (2a) are qualitatively similar with the M-GITM temperature distribution in Figure 409 410 2b. As we will show later, the aphelion distribution of peak temperatures can be explained well by dynamical heating in regions of converging thermospheric winds. The equinox comparison 411 between data (2c) and model (2d) confirms warm temperatures near the subsolar point but the pre-412 413 dusk terminator feature that is so prominent in M-GITM results at 17-18 LST is difficult to see in the statistics. Nevertheless, the statistical results at  $\pm 30^{\circ}$  latitude do indicate a temperature increase 414 that occurs prior to 18 LST. The diminished amplitude of the observed feature may be a result of 415 the low spatial resolution of the bins but may also be due to other sources of variability introduced 416 into the statistics. For example, the equinox sampling incorporates a wider range of EUV forcing 417 conditions than the other seasons sampled in Figure 2. These EUV conditions are not always 418 represented equally in each bin. Meanwhile, each model result in panels (b, d, f) corresponds to a 419 single EUV irradiance value. Perihelion results appear in Figures 2e and 2f. The incomplete data 420 coverage during this season means that temperature statistics are not available in the entire dayside. 421 The available data does indicate temperature enhancements just prior to the dusk terminator, 422 especially near the equator. A temperature enhancement in the summer hemisphere (southern 423 latitudes) prior to the dusk terminator is also apparent in Figure 2e. These observed perihelion 424 features agree qualitatively with the M-GITM temperature structure near dusk in Figure 2f. 425

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Figure 2: Seasonal statistical analysis of observed (left) and MGITM temperatures at 165 km
altitude. The seasons are aphelion (a,b) Ls 220-320, equinox (c,d), and perihelion (e,f) or Ls 40140.

Next, we examine the observed statistical winds and compare them with M-GITM simulated 430 winds. The observed and modeled divergence of the horizontal wind field is computed using term 431 B in equation 5. Figure 3 displays the statistical (panels a, b, c) and M-GITM wind fields (panels 432 d, e, f) where the vector plots indicate the average wind direction and magnitude while the color 433 maps indicate divergence (blue for positive divergence associated with dynamical cooling and red 434 for negative divergence associated with dynamical heating). Three seasons are represented in 435 Figure 3, aphelion (a, d), equinox (b, e), and perihelion (c, f). As before, the season-averaged 436 terminator location is indicated by crosses while the subsolar point is indicated by an asterisk in 437 each panel. 438

439

426

At aphelion, the observed (3a) low latitude winds generally exhibit zonal divergence (blue) on the 440 dayside and convergence (red) just past the dusk terminator and also around midnight. 441 Convergence values are also high in the dayside summer hemisphere (north) especially near the 442 terminator. The M-GITM results for aphelion (3d) are qualitatively consistent with the statistical 443 winds in that convergence peaks around summer dayside hemisphere near the dusk terminator with 444 moderate convergence around the low latitude dusk terminator. The location of zonal wind reversal 445 and convergence in the nightside winter hemisphere occurs near midnight in the M-GITM result 446 but in the statistical results zonal convergence can be inferred to occur in the early morning (see 447

wind directions prior to dawn terminator southern hemisphere in panel a). The near equatorial dusk temperature enhancement in Figure 2a is located near the converging winds observed in the vicinity of the dusk terminator. Unfortunately, the statistically observed temperature enhancements in the southern hemisphere (Figure 2a) lack corresponding spatial coverage in terms of wind observations (Figure 3a). Also, the strong convergence around the summer hemisphere dusk terminator has no corresponding temperature enhancement in Figure 2a. This could again be due to the coarse resolution of the temperature analysis however.

455

456 A notable feature of the low latitude equinox winds is a region of convergence around 12-16 LST rather close to the subsolar point (Figure 3b). This region of convergence is caused primarily by 457 equatorward meridional winds on both North and South sides of the low-latitude region. 458 Examining the region south of the dayside equator, the equatorward winds near 15 LST are part 459 of a rotation pattern centered approximately at 30° south latitude and 12 LST. It seems reasonable 460 that this would be the case in the northern hemisphere as well although data coverage in this region 461 is incomplete. There is no equivalent dayside convergence feature in the equinox modeled 462 thermosphere near the subsolar point (Figure 3e) but there is a larger rotation pattern centered  $\pm 40^{\circ}$ 463 latitude and 15 LST which turns the modeled mid-latitude winds equatorward starting around 17 464 LST and culminating in the modeled peaks in convergence at 18-19 LST. The statistical dayside 465 equatorward winds originate from more than one wind campaign and cannot readily be ascribed 466 to an anomalous time period or atmospheric condition (EUV irradiance or dust storm for example). 467 As with aphelion, low latitude zonal convergence (zonal wind reversal) occurs post-midnight (3 468 LST) at equinox in the statistical results, and pre-midnight (18-19 LST) in the M-GITM outputs. 469 Where data is available, high convergence values are observed near the terminators, which is 470 consistent with model results. Another similarity with equinox M-GITM results is that the 471 observed winds are approximately symmetric about the equator. 472





Figure 3: Statistical (top) and MGITM (bottom) winds. Color maps represent the divergence of
the wind field in units of s<sup>-1</sup>. Results are shown for aphelion (a, d), equinox (b, e), and perihelion
(c, f). Crosses indicate the location of the terminator while the asterisk in each panel is the
location of the subsolar point.

Perihelion statistical winds have the sparsest coverage of the three seasons. Nevertheless, strong 478 479 convergence values are seen near the terminator and midnight in the summer hemisphere (south) in both the statistical and M-GITM results. As with aphelion and equinox, there is some evidence 480 that zonal wind reversal does not occur pre-midnight (in the summer hemisphere for example) 481 482 with eastward winds dominating across the terminator into later local times in the observed statistics for the southern hemisphere (3c) while the M-GITM winds indicate a reversal around the 483 same location. The sparsity of the perihelion wind statistics make it difficult to correlate 484 485 temperature enhancements in Figure 2e with statistical winds in Figure 3c.

486

The range of divergence values across the three seasons based on the statistical analysis is larger 487 than those estimated from model winds. The largest data-model difference in divergence 488 magnitudes occurs at aphelion. This is in spite of the wind magnitudes being comparable between 489 the statistical and model results. Statistical divergence magnitudes also seem to increase from 490 491 perihelion (3c) to aphelion (3a), a trend that is not present in the M-GITM results. When comparing divergence values in the high latitude dusk-side summer regions the divergence peaks at 492  $\sim 0.0005s^{-1}$  at aphelion (3a),  $0.0003s^{-1}$  at equinox (3b), and  $0.00015s^{-1}$  at perihelion (3c). This would 493 signify that as the solar EUV irradiance decreases at higher Mars-Sun distances, the dynamical 494 heating/cooling increases, reinforcing its relative contribution to the thermal state of the 495 thermosphere. 496 497

Another way in which the standard M-GITM runs disagree with observed wind statistics is that at 498 499 mid to high latitudes (beyond  $\sim 50^{\circ}$  north or south latitude), there is an absence of observed reversal in the zonal winds. This is most clearly seen in the northern aphelion hemisphere in panel a, 500 501 southern equinox hemisphere in panel b, and northern perihelion hemisphere in panel c where, for all seasons, the observed polar rotation is predominately eastward. In other words, observed winds 502 provide evidence of an eastward neutral wind jet above 50° latitude. Meanwhile, the M-GITM 503 zonal reversal can be seen in the model results around 12 LST along with a corresponding increase 504 in divergence at all seasons. 505

506

The results shown above indicate that regions of enhanced thermospheric temperatures often occur 507 away from the subsolar point at locations of converging winds (Figures 2a and 3a for example). 508 Such temperature enhancements are well represented by the M-GITM model. Furthermore, 509 observed winds confirm that significant convergence can exist near the equator, an important 510 ingredient for dynamical heating. Significant nighttime convergence is predicted by this M-GITM 511 version and occurs in the pre-midnight local times. In contrast, the observed nighttime convergence 512 occurs in the post-midnight region. Most surprising is a region of dayside equatorial convergence 513 observed at equinox. This feature is associated with converging meridional winds and does not 514 appear in the M-GITM results. Further analysis is needed to see if this feature is observationally 515

516 and statistically robust.

517 4.2 Equinox Empirical Reconstruction at Constant Solar Irradiance

To alleviate some of the challenges associated with data sparsity and binning resolution when 518 evaluating the wind field and convergence is taking advantage of the approximate hemispheric 519 symmetry in the equinox circulation. First, we use MSO coordinates to reduce any meridional 520 shifts in the circulation pattern within a limited set of solar longitudes  $(L_s)$ . The meridional wind 521 component is then converted into equatorward and anti-equatorward directions such that data from 522 the north and south hemisphere can be combined. High and very low EUV values are also excluded 523 from the analysis resulting in EUV conditions between 0.0025-0.0032 Wm<sup>-2</sup>. The results of this 524 analysis can be seen in Figure 4. The main equinox features described in the IAU wind statistics 525 526 (Figure 3) are also present in Figure 4a/d including the low latitude meridional convergence around 15 LST and the zonal convergence occurring post-midnight at approximately 2-3 LST. The M-527 GITM wind fields are analyzed in the same way and the results shown in Figure 4b for a version 528 of M-GITM without GW effects. Figure 4c shows the average low-latitude divergence from the 529 equator to 40° and 50° MSO latitude as indicated by two values shown at each local time. The 530 range of statistical convergence (asterisk values connected by a line) are compared to the 531 532 equivalent M-GITM analysis (diamonds). The modeled and observed horizontal divergence share the same approximate diurnal structure with predominately divergent flow on the dayside and 533 convergent flow on the nightside. Differences include a larger observed pre-midnight divergence 534 and larger observed post-midnight convergence as well as the aforementioned observed 535 convergence around 15 LST. Another feature of the observed convergence is a wave-like behavior 536 with a 4-hour LST wavelength (3,500 km) in some local time ranges, primarily 0-10 LST and 12-537 538 20 LST.

539

540 The right side of Figure 4 contains the same analysis of observed winds (4d) along with the 541 analogous model results from M-GITM<sub>GW</sub> in panel 4e and the average low-latitude divergence 542 based on data and M-GITM<sub>GW</sub>. The inclusion of GW wave effects removes the rotation pattern seen in the afternoon and centered on 16 LST in the standard M-GITM results (4b). Instead the afternoon winds in the M-GITM<sub>GW</sub> case continue in an eastward direction past midnight until 2-3 LST, resulting in enhanced convergence at these locations and better agreement with the statistical results. Overall wind magnitudes are reduced in M-GITM<sub>GW</sub> as are the model divergence values in panel 4f. Note that neither the standard M-GITM (4c) nor M-GITM<sub>GW</sub> (4f) manifest a dayside low/mid latitude convergence observed near 15 LST.







Figure 4: Equinox winds plotted in MSO coordinates assuming wind-field symmetry around the equator. The colormap indicates the value of divergence. Panels (a,d) are the observed equinox winds and are repeated at the top to aid in comparison with the plots below, panel (b) shows the

554 *M*-GITM results without gravity wave effects, panel (e) shows model winds from a version of M-

555  $GITM_{GW}$  that includes gravity waves, and panels (c,f) indicate the low to mid-latitude averaged

556 *divergence computed using modeled (diamonds) and measured (asterisk) winds. The range of* 

values indicated in panels (c,f) for both the model and measured divergence illustrates the

558 *difference between taking the average divergence between*  $0^{\circ}$ - $50^{\circ}$  *latitude and*  $0^{\circ}$ - $40^{\circ}$  *latitude.* 

559 Panel (e) M-GITM results do not include gravity wave effects and panel (f) M-GITM<sub>GW</sub> results

560 include gravity wave effects. Note that the divergence color scales are not the same in each of 561 the panels.

562

In order to further examine the low latitude divergence, we divide the low/mid-latitude wind data 563 into 1/3 LST bins and combine meridional wind data between 8° and 50° MSO latitude and zonal 564 data between 0° and 50° MSO latitude. The exclusion of lowest latitude meridional data is done 565 to avoid the near zero meridional components near the equator (see Figure 4a). This way, the 566 average meridional component (equatorward/anti-equatorward) is captured and can be used to 567 compute a horizontal divergence assuming that the meridional winds reach zero by 0° latitude. A 568 series of Fourier fits to the meridional and zonal components is then performed. The fits are based 569 on average, upper, and lower quartile winds within each bin. The results are shown in Figure 5a 570 and 5b for the meridional and zonal directions respectively. 571

572

Examining the peaks in equatorward (negative) winds in Figure 5a confirms some of the previous 573 results but with a higher resolution and indicates that equatorward winds lead to convergence post-574 midnight (2-3 LST), near 15 LST, and near the dusk terminator. There is also a region of 575 moderately equatorward winds at around 11 LST which was not apparent in the previous analysis 576 but which seems to form part of a wave pattern that extends from approximately 10 LST to 577 midnight. The associated wavelength is approximately 4 hours LST (3,500 km). Model meridional 578 winds are shown for M-GITM (dashed blue line), M-GITM<sub>GW</sub> (dashed black), and M-GITM<sub>GW</sub> 579 580 (dashed orange) where the meridional component is multiplied by a factor of three. Reasonable agreement between statistical and modeled meridional components is obtained when the model 581 meridional component is significantly increased and when gravity waves are included (dashed 582 orange). 583



Figure 5: Equinox mid-latitude meridional winds (a), zonal winds (b), and divergence (c) based
on the meridional and zonal fits. Dashed lines represent M-GITM model results and solid lines
represent fits to the data. The various fits shown represent the uncertainty in the fitting results.

The zonal winds in Figure 5b indicate convergence during east-to-west transitions (positive to negative) seen post-midnight, around 4 LST, 10-11 LST, and 14-15 LST. The zonal convergence locations are thus mostly consistent with the convergence of the meridional component. A wave pattern with an approximately 4 hour wavelength is also observed in the zonal measurements. The modeled zonal component is in general agreement with the statistical winds only when GW parametrization is included (black dashed line).

595

585

596 Combining the spatial derivatives from all six of the Fourier fits (three for each component) results 597 in nine divergence estimates plotted in Figure 5c using solid lines. The dashed lines in panel 5c

are the equinox M-GITM divergence values computed using the same sampling method as used

to analyze the observed winds. The three largest statistical convergence values occur post-midnight

600 (1-2 LST), near 15 LST, and near the dusk terminator (18-19 LST). The largest low-latitude M-

601 GITM convergence values are seen at 4 LST, and the dusk terminator (18 LST). A large oscillation

with a 4-hour wavelength appears in the statistical divergence estimates. The dayside M-GITM

divergence in panel c also exhibits a 3-4 hour (2,700 - 3,500 km) wavelength although the modeled
 amplitude is much lower than the observed signal.

605

The dashed black lines in Figure 5 representing M-GITM<sub>GW</sub> results are in better overall agreement 606 with the observed winds than the non-GW version of the model. A version of M-GITM<sub>GW</sub> 607 calculation that triples the magnitude of the meridional component improves data-model 608 agreement in convergence further (5a and 5c), especially in the night time. The inclusion of GW 609 parametrization significantly improves the zonal wind agreement between model and data as seen 610 in figure 5b. As noted before, this places the night-time zonal wind reversal in the morning sector. 611 Both the diurnal structure and magnitude of the M-GTIM<sub>GW</sub> zonal wind components agrees well 612 with the data. Horizontal divergence comparisons were shown in Figure 5c. Two versions of M-613 GITM<sub>GW</sub> divergence were included, those with the original modeled wind magnitudes (black 614 dashed line) and a version where the meridional wind magnitudes have been increased (orange 615 dashed line). Night-time convergence peaks appear at 1-2 LST, 4-5 LST, 18-19 LST in the 616 statistical analysis as well as M-GITM<sub>GW</sub>. A smaller convergence observed near 22-23 LST is not 617 seen in M-GITM<sub>GW</sub> although this feature may be a processing artifact resulting from the data gap 618 between 22 and 24 LST. In the dayside, the appearance of an oscillation with a ~4 hour wavelength 619 can be seen in both data and M-GITM<sub>GW</sub> and, compared to the standard M-GITM results in panel 620 5c, the phase of the oscillation matches the observed divergence more closely. As with the standard 621 M-GITM model, the amplitude of the M-GITM<sub>GW</sub> convergence oscillation is much smaller than 622 that based on the Fourier expansion fits. 623





#### 625

*Figure 6: Climatological thermospheric temperatures at three constant pressure levels (a, b, c, c, c)* 

- 627 *d)* and the corresponding pressure level heights (e, f, g, h) at equinox, low-to-mid latitudes ( $\pm 50^{\circ}$ 628 *MSO*). Dynamical heating terms are shown in panels (i, j, k, l). Panels (a, e, i) are an empirical
- 629 reconstruction based on a linear response to EUV irradiance, panels (b,f,j) are an empirical
- 630 reconstruction based on a nonlinear response for pressure level heights only. Panels (c, g, k) are
- 631 the M-GITM extracted at the equivalent conditions while panels (d, h, l) are equivalent M-GITM
- 632 results for a model version that includes gravity wave effects.

Figure 6 contains the constant-pressure empirical reconstruction of NGIMS derived temperatures 633 and pressure level heights that were described in section 3.1 (panels a, b, e, f). These results are 634 evaluated for a constant EUV Lyman-alpha irradiance of 0.0032 Wm<sup>-2</sup> and a range of latitudes 635 between  $\pm 50^{\circ}$  MSO. The corresponding M-GITM results appear in panels c, d, g, and h for a range 636 of latitudes between  $\pm 30^{\circ}$  MSO. The empirical reconstructions based on linear fits to EUV 637 irradiance are shown in panels a and e while those using non-linear fitting are in panels b and f. 638 Figure 6 also contains estimated dynamical heating based on horizontal terms A and B in equation 639 5 (panels i and j) and computed using the fits to zonal and meridional wind observations shown in 640 Figure 5. Note that panels i and j are identical and included for ease of comparison with the above 641 panels. The modeled dynamical heating (based on all terms in equation 5) is shown in red in panels 642 k and j along with the sum of the non-dynamical heating terms shown in black. Model results are 643 included for the standard M-GITM runs in panels c, g, k and for the M-GITM<sub>GW</sub> runs in panels d, 644 h, l. In all panels, results are shown for three pressure levels indicated by the dotted, solid, and 645 dashed lines. The pressure levels are separated from each other by a pressure factor of 646 approximately four. The mean altitude of these pressure levels is shown in panels e, f, g, and h for 647 the linear empirical reconstruction, non-linear reconstruction, standard M-GITM, and M-GITM<sub>GW</sub> 648 649 respectively.

650

First considering the empirical temperature reconstructions in panels a and b, we observe a dayside 651 peak temperature of ~250K near 12-15 LST and a nightside low temperature of 140-150K around 652 20-24 LST. As the thermosphere rotates further away from the dayside, we would expect radiative 653 cooling to result in ever decreased temperatures further from the dusk terminator. However, the 654 reconstructed post-midnight temperatures (0-4 LST) are actually warmer than those before it. 655 There are local temperature enhancements between 1-4 LST as well as 16-18 LST in both the 656 linear (a) and non-linear (b) empirical reconstructions. A smaller temperature peak is present at 657 14-15 LST in panel (a) and corresponds to a peak in pressure level altitudes at the same location 658 (panel e). These heating and pressure level uplift features correlate with the largest enhancements 659 in observed dynamical heating at 1-4 LST, 13-16 LST, and 17-19 LST. The morning pressure 660 enhancement, when interpreted at fixed altitude between 150-160 km, results in an almost 4x 661 increase in density relative to the surrounding nighttime thermosphere. In fact, dynamical heating 662 is present throughout the low-latitude morning region and is larger than the dynamical heating in 663 the pre-midnight sector (panel i/i). This may explain why the morning (post-midnight) nighttime 664 thermosphere in panels (a) and (b) is warmer than the pre-midnight thermosphere. A dip in 665 temperatures near 11-12 LST and 16-17 LST seen in panels a and b correlates with the dynamical 666 cooling (negative values) seen in panels i and j near the same locations. Taken together, these 667 observed temperature decreases and enhancements follow the dynamical heating estimated from 668 observed wind statistics and constitute an oscillation with an approximately 4 hour LST period. 669

670

671 Turning now to the model results, the empirical temperatures in panels a and b are best matched in terms of day-night ratio by the standard M-GITM run in panel c which has dayside temperatures 672 of ~250K at 12-15 LST and 140-150K around 20-24 LST. The location of the modeled dusk 673 dynamical heating feature at 17-18 LST in panel c also agrees with the empirical results but its 674 model amplitude is much higher. Meanwhile, the M-GITM<sub>GW</sub> temperatures in panel (d) have a 675 larger diurnal amplitude (day-night ratio) and lower overall values than those observed with the 676 677 dayside peak near ~240K and the nightside lows of ~120K. The dynamical heating temperature enhancement near dusk in M-GITM<sub>GW</sub> results occurs at a later local time (19 LST) but has a 678

smaller amplitude that better matches the empirical temperature enhancement (~10K increase). 679 680 Another feature of the M-GITM<sub>GW</sub> temperature results (panel d) is a ~10K peak near midnight that is absent in the standard M-GITM temperatures (panel c). This enhancement may correspond to 681 the post-midnight temperature increase seen in the data (panels a and b at 1-4LST). The modeled 682 morning temperature enhancement near 5 LST appears in both model versions and is accompanied 683 by an increase in modeled pressure-level height (panels g and h). A similar feature is difficult to 684 identify in the empirical reconstructions of both temperature and pressure-level height. 685 Qualitatively, the inclusion of GW parametrization improves the temperature structure associated 686 with dynamical heating in the morning and dusk regions by broadening and attenuating the 687 dynamical heating peaks (red lines in panel k compared to panel l). Both the dawn and dusk 688 enhancements in modeled temperature and pressure-level height correspond to the dynamical 689 heating enhancements in panels k and l. The dampening of meridional winds by a factor of ~3 in 690 the M-GITM<sub>GW</sub> relative to the wind statistics (Figure 5a) may be responsible for the reduction of 691 night-time dynamical heating leading to the cooler morning temperatures in this version of the 692 693 model.

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## 696 6 Discussion

697 The MAVEN data analyzed here suggests thermospheric enhancements in neutral temperature at a fixed pressure level that correspond to statistical wind convergence and dynamical heating. The 698 most robust signal associated with dynamical heating seems to be responsible for enhanced 699 temperatures in the deep night-side of the thermosphere. Additional evidence of dynamical heating 700 appears near the dusk terminator and in the afternoon region. Both data and model suggest that the 701 dynamical heating terms significantly modulate the overall thermal structure by modifying the 702 overall heat balance between EUV heating, IR heating, 15µm cooling, and conduction (responsible 703 primarily for cooling). The global variation seen in the model dynamical heating causes 704 705 corresponding temperature fluctuations in the model at fixed pressure. A similar dynamical heating oscillation is seen in the empirical reconstruction, especially in the dayside between 10-18 LST 706 707 but has a much larger amplitude than either of the model versions evaluated here.

708 709

710 An important question is how the large amplitude dayside oscillation in divergence and the resulting oscillation in dynamical heating is generated. One of the surprising results of this analysis 711 is that the observed dayside divergence pattern is responsible for dynamical heating near 15 LST, 712 713 very close to the subsolar point. The increase in dynamical heating at this location is associated with a small enhancement in temperature as well as a significant increase in pressure-level height. 714 It is possible that the presence of a smaller oscillation in the dayside model thermosphere (M-715 GITM and M-GITM<sub>GW</sub>) provides a clue. The phase and wavelength of the observed modulation 716 in dynamical heating/cooling signatures is better reproduced when GW parametrization is included 717 in the model. However, the inclusion of GW parametrization results in modeled meridional wind 718 719 components that are smaller than those observed at equinox. It is possible that the GW drag in the M-GITM model could be overestimated for this season, resulting in a meridional flow that is over-720 damped. This model-data comparison highlights the important coupling between GW effects, the 721 722 global circulation pattern, and the resulting dynamical effects seen throughout the Martian thermosphere. 723

### 724

One possible explanation for the global-scale temperature, pressure, and wind-field oscillation is a terminator wave which propagates through the density and wind field. Simulations by Forbes and Moudden (2009) have shown that such a wave is possible and that it results in a density oscillation with a wavelength of  $\sim$ 2-4 hours LST (1800-3600 km), similar to that observed. Furthermore, extrapolating the density profile modeled by Forbes and Moudden to  $\sim$ 200 km altitude does result in a density enhancement at 15 LST which is in qualitative agreement with the empirical results presented here as well as with the temperature structure in M-GITM<sub>GW</sub>.

732

733 An alternative perspective is the presence of two, ~2,500 km diameter, dayside circulation cells at low to mid latitudes originating from solar heating and rapidly turning eastward then equatorward, 734 leading to convergence around 15 LST. It is possible that these cells are mostly independent of 735 any wave originating from the dusk terminator. The modeled (with or without GW 736 parametrization) circulation patterns are much larger than this (>10,000 km diameter) with 737 circulation diverging from the dayside mid-high latitudes and converging on the nightside. This 738 difference can again be explained by excessive wind drag in the meridional direction represented 739 in the model, impeding the generation of smaller circulation cells. The over-dampening of 740 meridional wind magnitudes in the model is clearly visible in Figure 4a, especially in the night and 741 terminator regions. A difference in meridional and zonal wave drag may stem from anisotropy in 742 GW propagation at the source (~9km altitude) or from anisotropic filtering of GW's below the 743 thermosphere that is not represented in M-GITM<sub>GW</sub>. Any modeled errors in lower-altitude (10 -744 150 km) wind fields might also be manifested in the wind drag distribution. 745

746

### 747

## 748 7 Conclusions

Dynamical heating is critical to describing the large scale structure of the Mars thermosphere. The 749 presence of temperature and density features in the Martian thermosphere resulting from 750 dynamical heating has been suggested previously by global circulation models such as M-GITM. 751 752 The most prominent of these features are broad temperature enhancements in regions of converging winds occurring near the dawn and dusk terminators. In this study, we have compared 753 754 these modeled features with empirical representations based on several Mars years of data. In addition to evaluating the temperature at a fixed pressure, we have also looked at the observed 755 thermospheric wind divergence at similar locations and conditions. In order to take advantage of 756 the most complete empirical dataset, the study focuses on the equinox seasons. The results indicate 757 758 that dynamical heating predicted by models is indeed observed. The inclusion of GW parametrization in M-GITM generally improves the data-model agreement indicating that GW 759 drag is critical in modulating the observed global circulation and the resulting thermospheric 760 horizontal structure. The comparisons between M-GITM<sub>GW</sub> and the observations are not perfect 761 however and work remains to be done to investigate the differences in wind patterns between the 762 model and observations. 763

764

At equinox the most robust thermal features appear in the near-equatorial thermosphere in the postmidnight and near-dusk regions with the post-midnight enhancement having the larger relative

amplitude. The M-GITM<sub>GW</sub> model contains analogous features near the dusk terminator as well as

near midnight although the locations of these features are different by one or two hours of local

time. The appearance of a dayside convergence location and a corresponding pressure height and (subtle) temperature response was unexpected. The fact that this feature may be a part of a larger migrating tide across most of the Martian thermosphere is also a new result. M-GITM<sub>GW</sub> does contain a dynamical heating modulation with a peak at the same location (15 LST) and a corresponding broad feature in the dayside temperature. However, the modeled heating and the wind divergence pattern that causes it is much less pronounced. More observations are needed to ensure that this dayside equinox feature is not the result of sampling or aliasing of other effects.

- This is especially true of the surprising planetary scale oscillation found in the horizontal wind observations. Future analyses would benefit tremendously from continuous and dedicated wind observations to better understand this feature.
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# 781 8 Data Availability

- The MAVEN/NGIMS number densities used in this study are the Level 2, version 8, release 2
- data. The neutral winds were also collected by the NGIMS instrument and are the Level 3,
- version 3, release 1 data product. Both MAVEN/NGIMS data products are available on NASA's
- 785 Planetary Data System (Benna & Lyness, 2015).
- 786

787 Solar fluxes used in the analysis are from the L2b orbit-averaged product from the

- 788 MAVEN/EUVM instrument. The solar fluxes used to drive the M-GITM model are from the
- 789 MAVEN/EUVM FISM-M empirical model and are a Level 3, Version 14, Revision 3 data
- product. Both MAVEN/EUVM data products are hosted on the Planetary Data System (F.
- 791 Eparvier, 2017, 2022).
- 792
- M-GITM model runs used in this work can be retrieved at the University of Michigan Deep Blue
   repository (Steve Bougher & Pilinski, 2022).
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